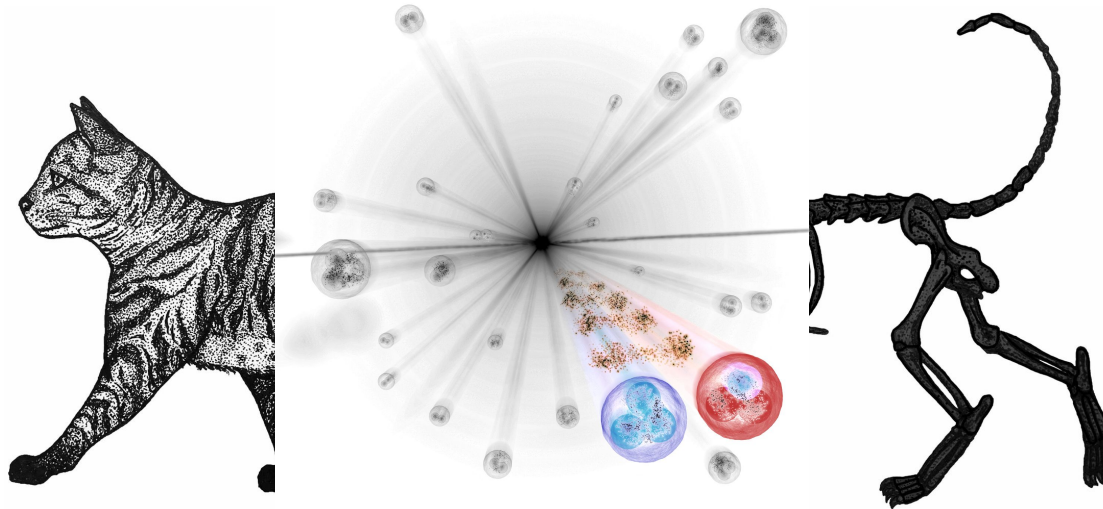


Hyperon-nucleon and hyperon-nucleon-nucleon interaction studies via femtoscopy



Dimitar Mihaylov

29nd June 2022, Prague, Czech Republic





Strange

p

Λ



Strange

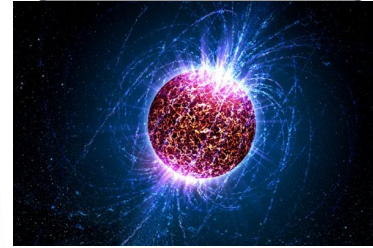
p

Λ

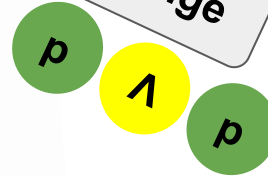
p



Neutron stars



Strange





Charming

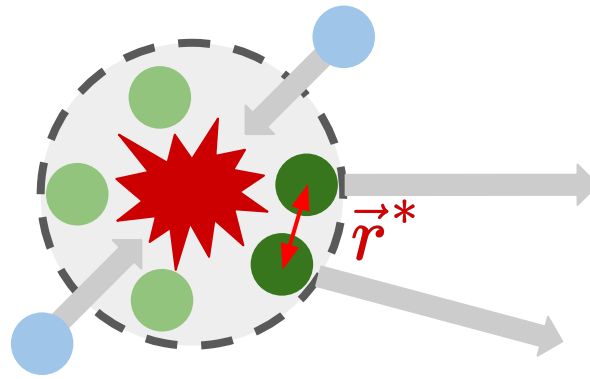
D

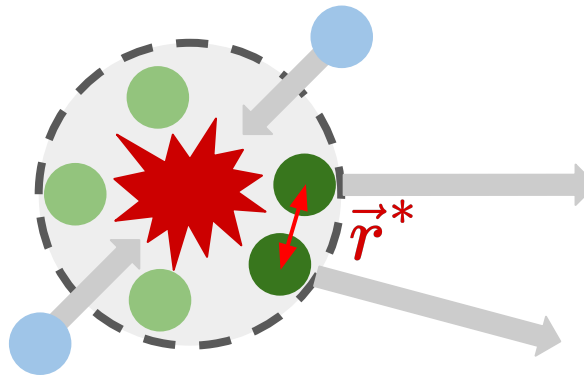
π

D

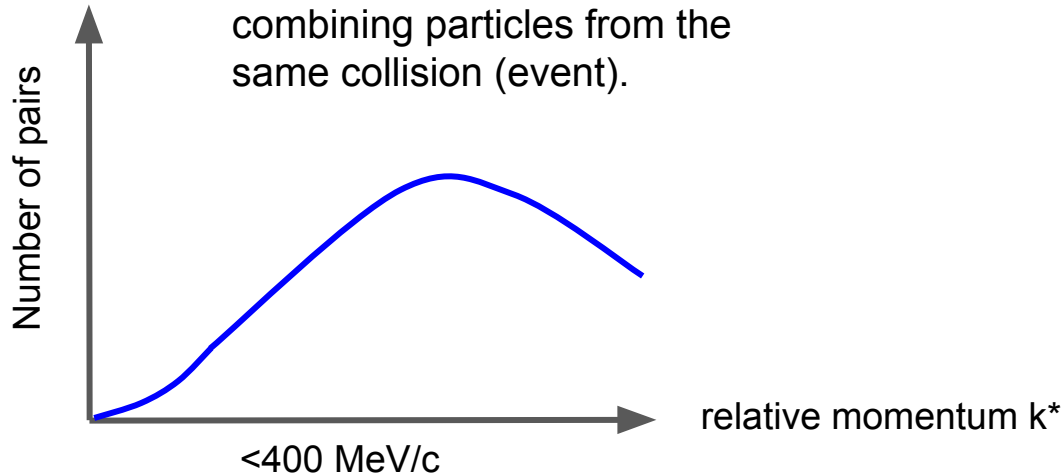
p

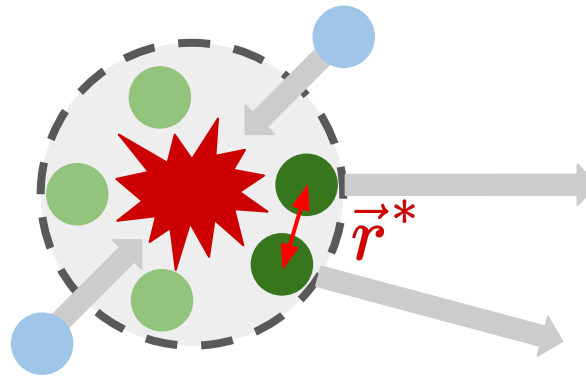
TUM Femtoscopy
Overview



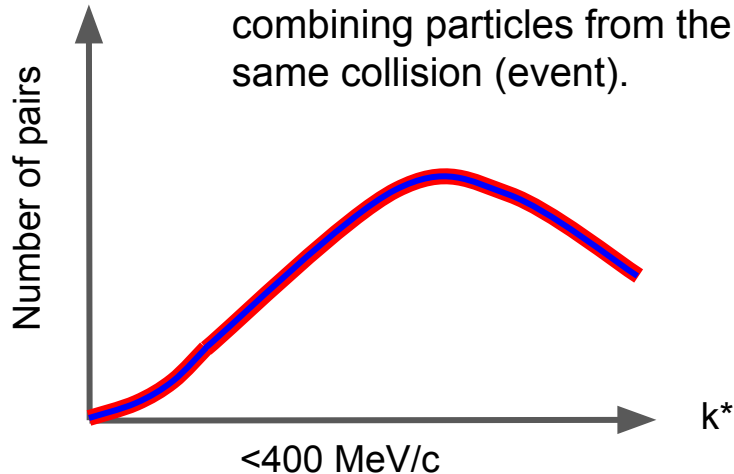


- **Same event sample (SE):**
Correlated pairs, obtained by combining particles from the same collision (event).

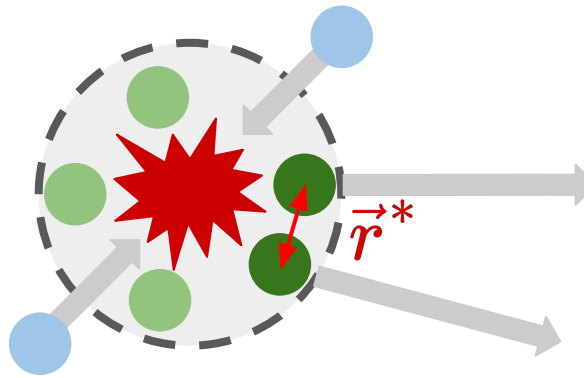




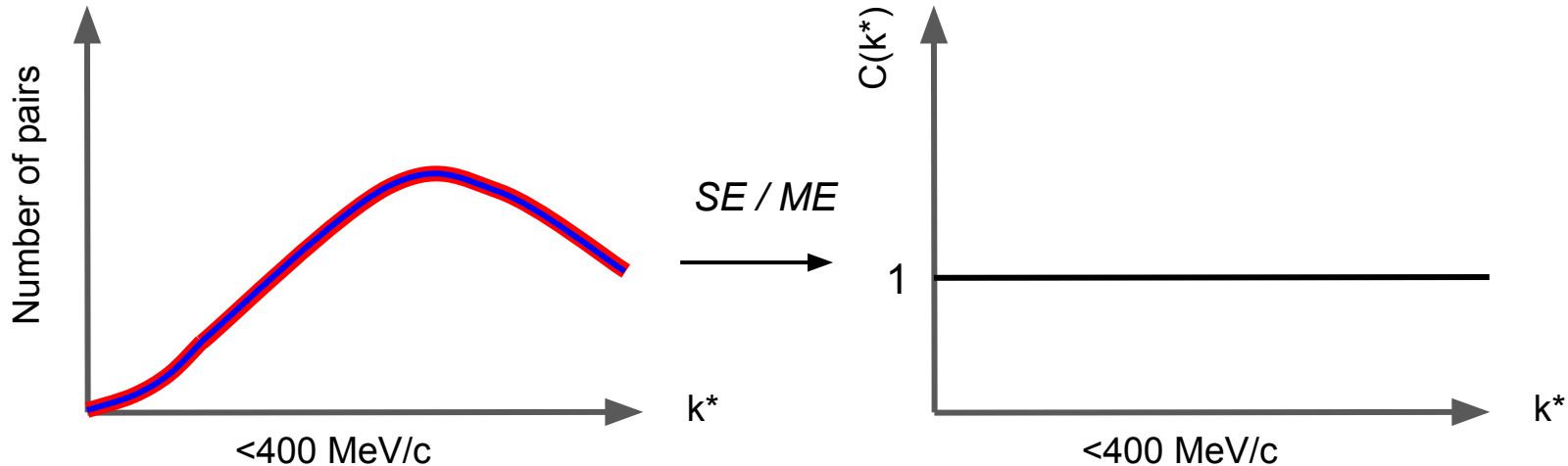
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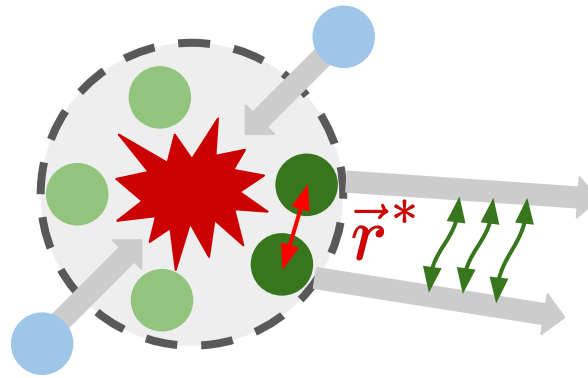


- **Mixed event sample (ME):**
Uncorrelated pairs, obtained by combining particles from two different collisions (events).

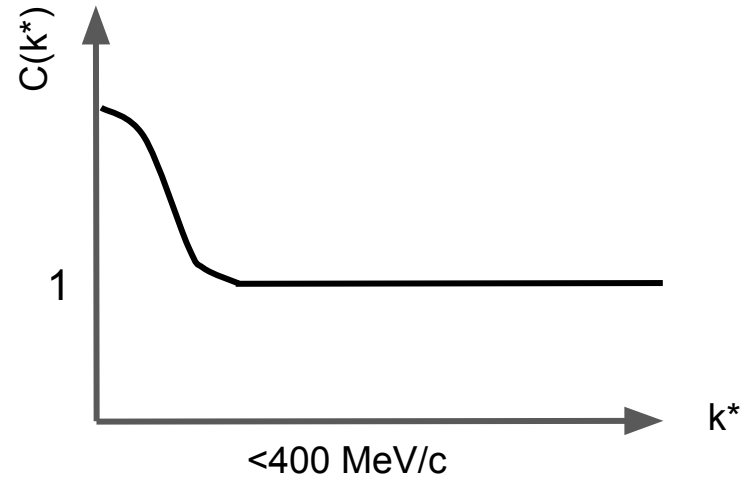
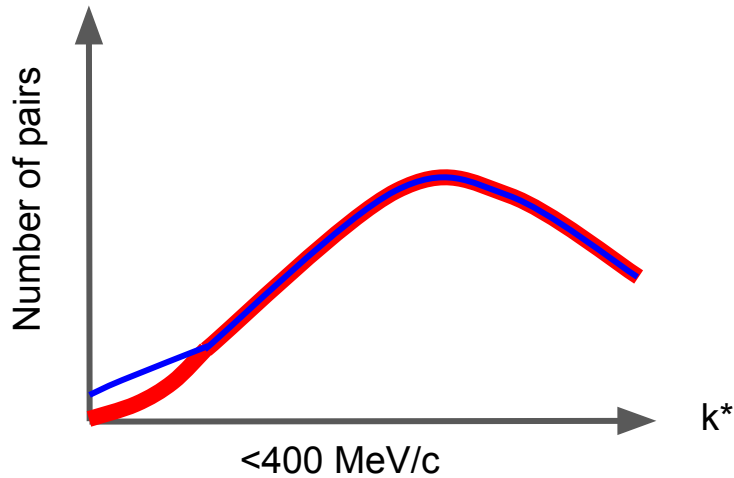


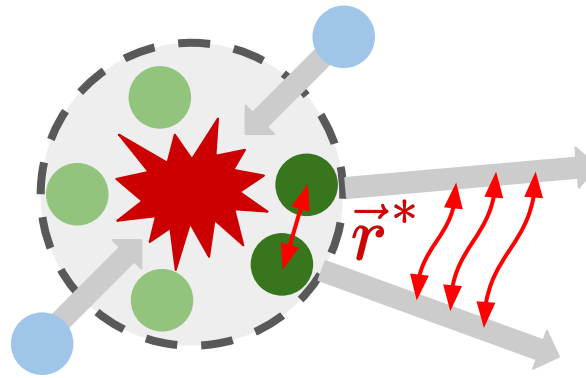
- The **correlation function** $C(k^*) = SE / ME$, ideally equal to unity in the absence of any correlations.



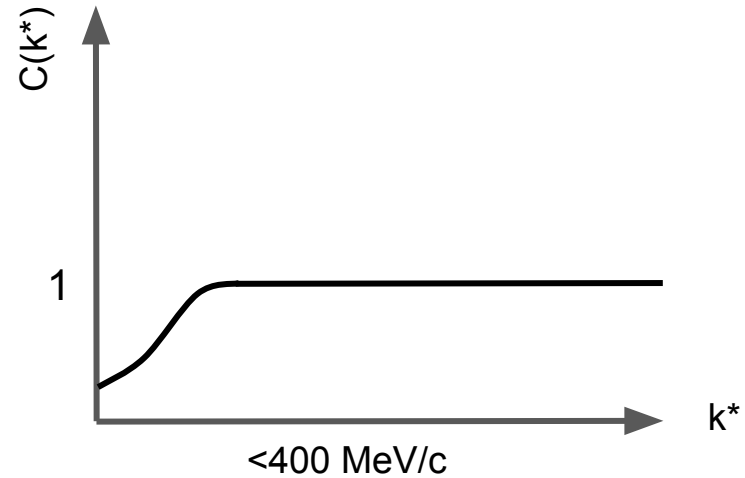
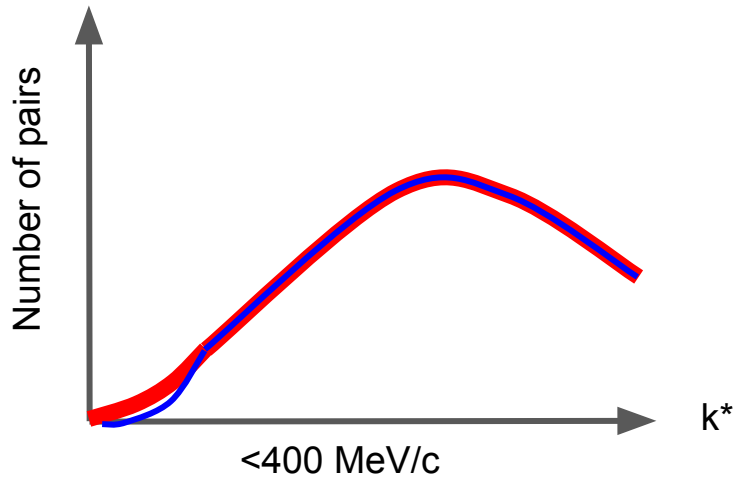


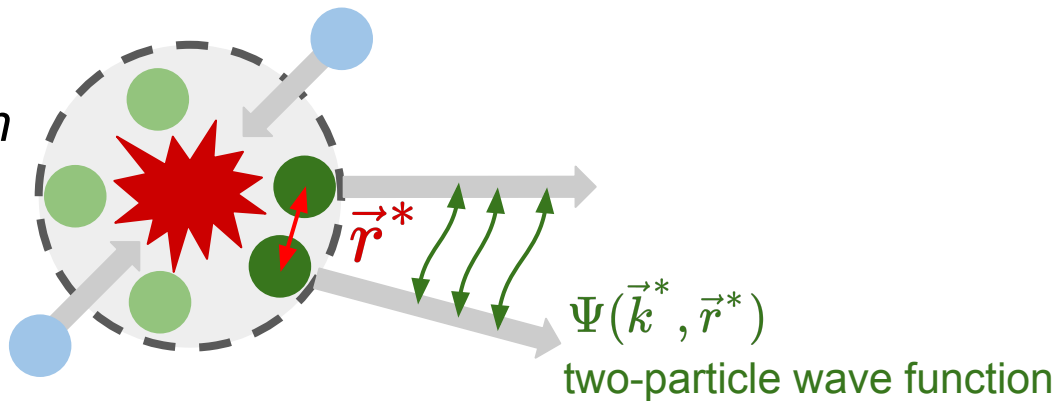
- **Attractive** final state interaction (FSI).





- **Repulsive** final state interaction (FSI).



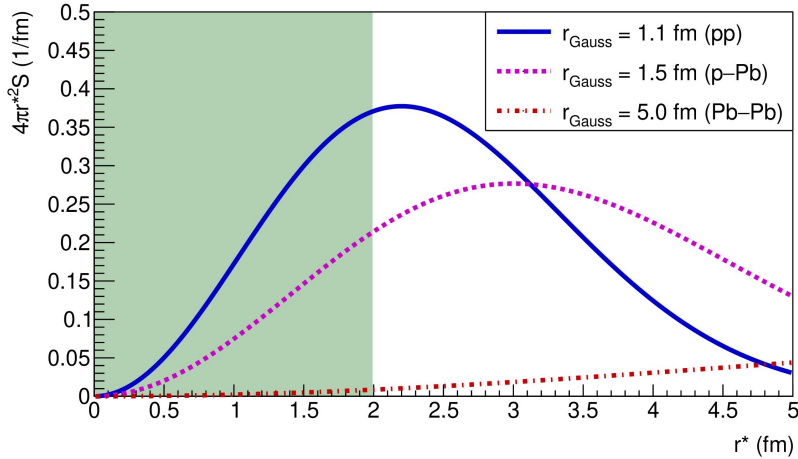


$$C(k^*) = \frac{N_{\text{SE}}(k^*)}{N_{\text{ME}}(k^*)} = \int S(r^*) \left| \Psi(\vec{k}^*, \vec{r}^*) \right|^2 d^3 r^* \xrightarrow{k^* \rightarrow \infty} 1$$

[Ann.Rev.Nucl.Part.Sci.55:357-402, 2005](#)

Relative distance and $\frac{1}{2}$ relative momentum
evaluated in the pair rest frame

- Measure $C(k^*)$, fix $S(r^*)$, study the interaction.
- Extension to coupled channels: [Wed-I: Yuki Kamiya](#)



$$C(k^*) = \int S(r^*) \left| \Psi(\vec{k}^*, \vec{r}^*) \right|^2 d^3 r^*$$

Measure

Fix $S(r^*)$

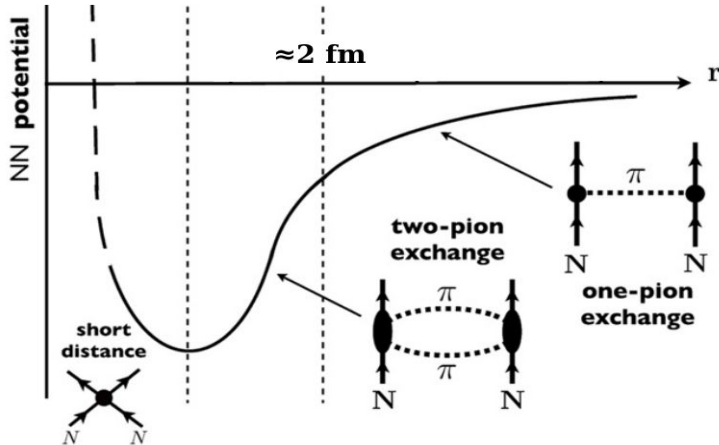
Study

- Enhanced sensitivity in **small collision systems (pp)**.
- Scan over the range of the interaction by using different collision systems.

[Wed-I: Yuki Kamiya](#), [Thu-I: Ramona Lea](#)

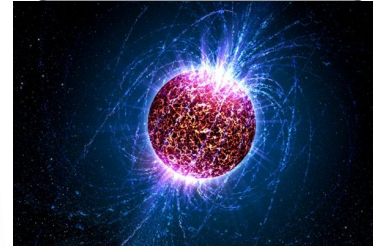
- **Common emission** of primordial particles in pp collisions, allowing to fix $S(r^*)$ for any baryon-baryon pair using the pp correlation function.

[Phys. Lett. B 811 \(2020\) 135849](#)





Neutron stars



Strange

p

Λ

TUM Neutron stars (NS)

Very compact, very dense

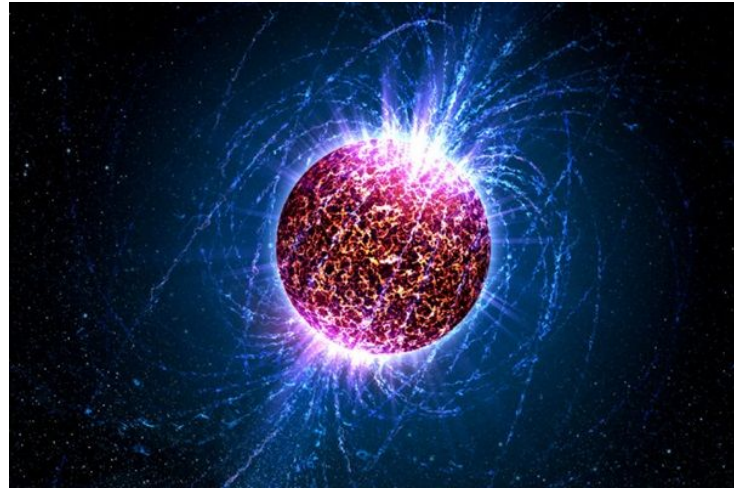
The sun packed in Manhattan

Radius: $R \sim 10$ km

Mass: $M \sim 2$ solar masses (M_{\odot})

Density: few times ρ_0^*)

*) $\rho_0 = 0.16 \text{ fm}^{-3}$ is the density of the nucleus of an atom, called “nuclear saturation density”.



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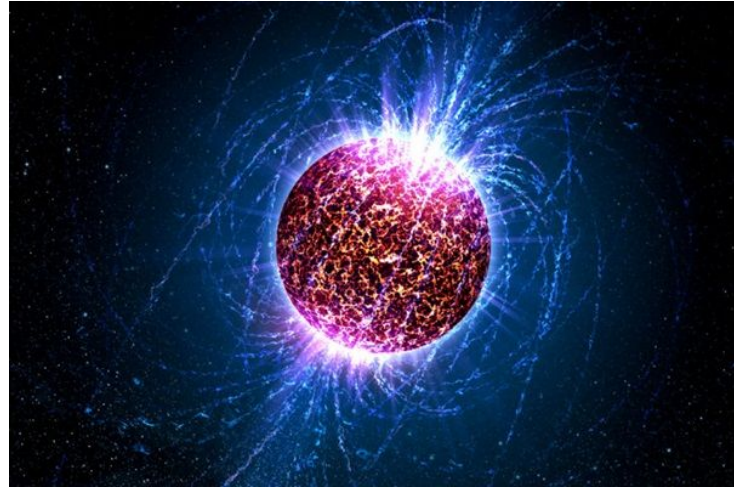
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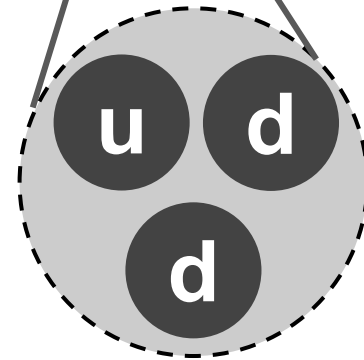
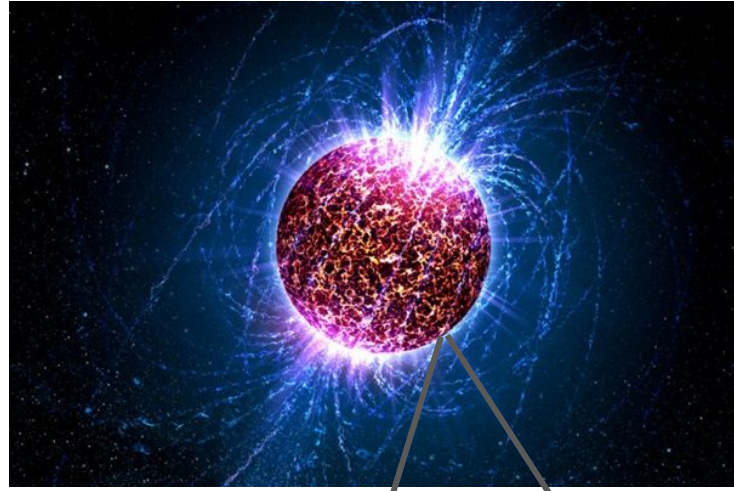


What is inside?

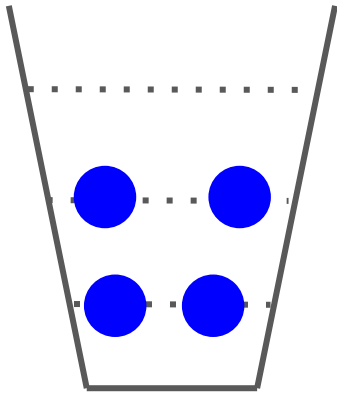
[Tue-I @ HYP](#)

(I. Vidana, L. Tolos, A. Ohnishi)

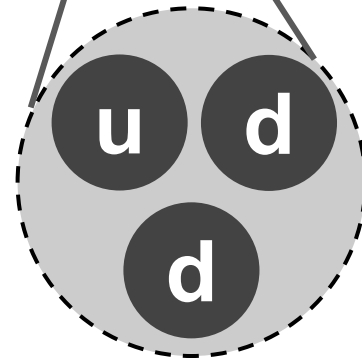
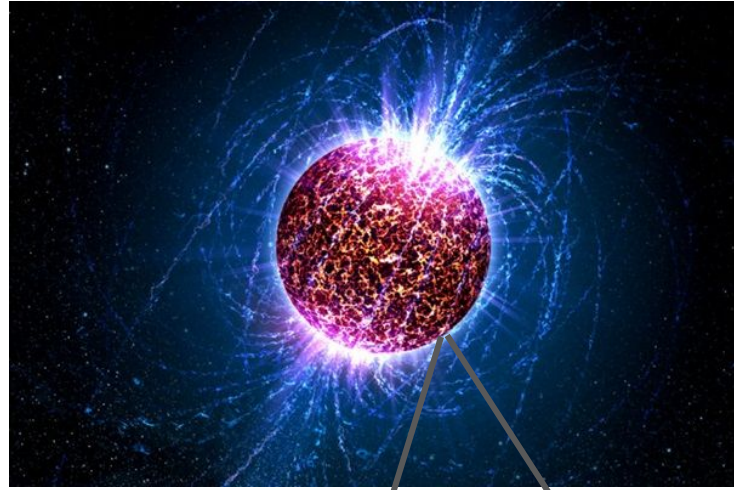
neutrons



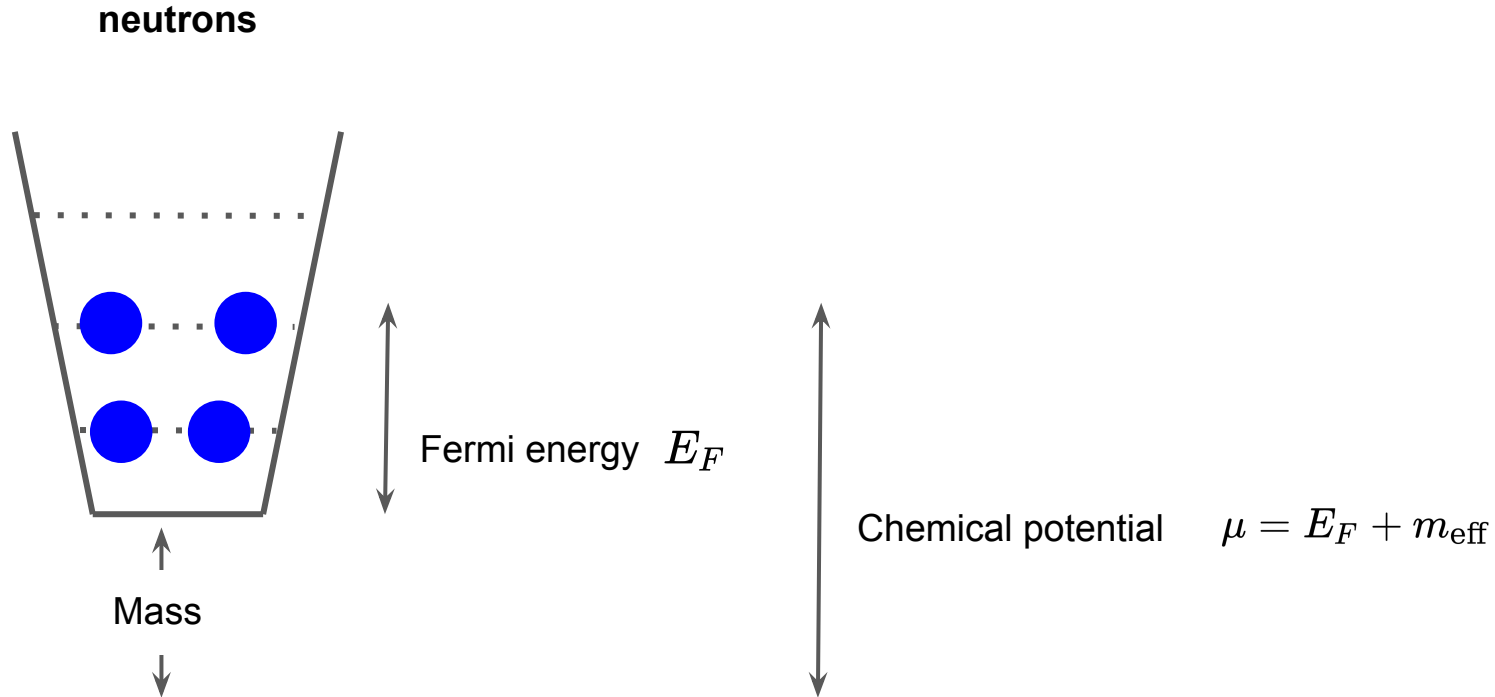
neutrons



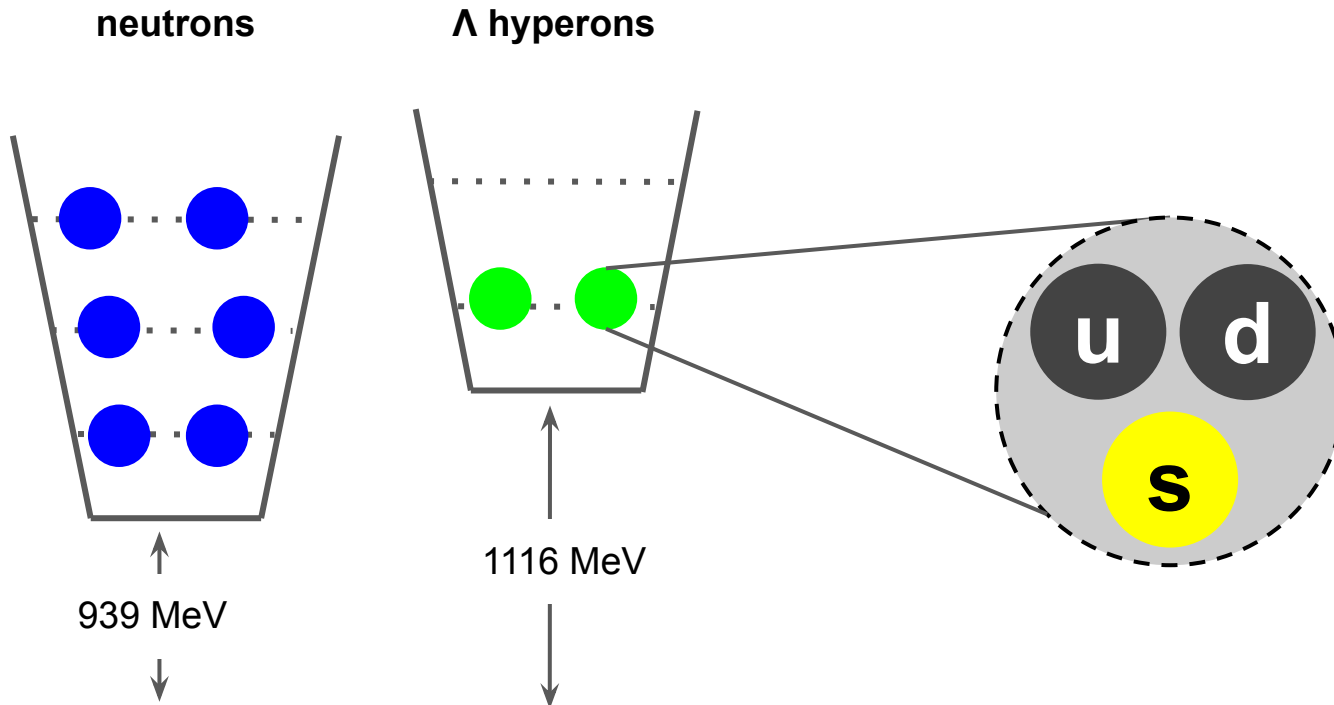
↑
Mass
↓



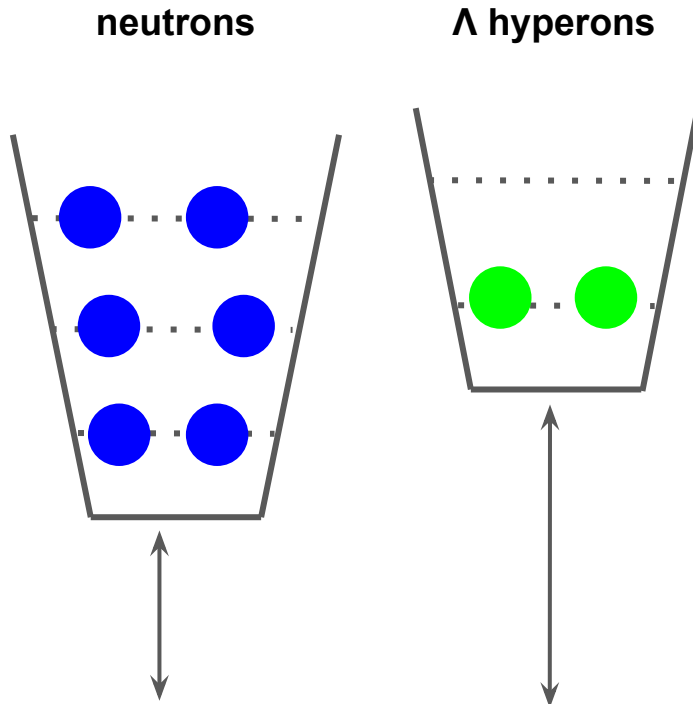
- As density increases, so does the Fermi energy and chemical potential.

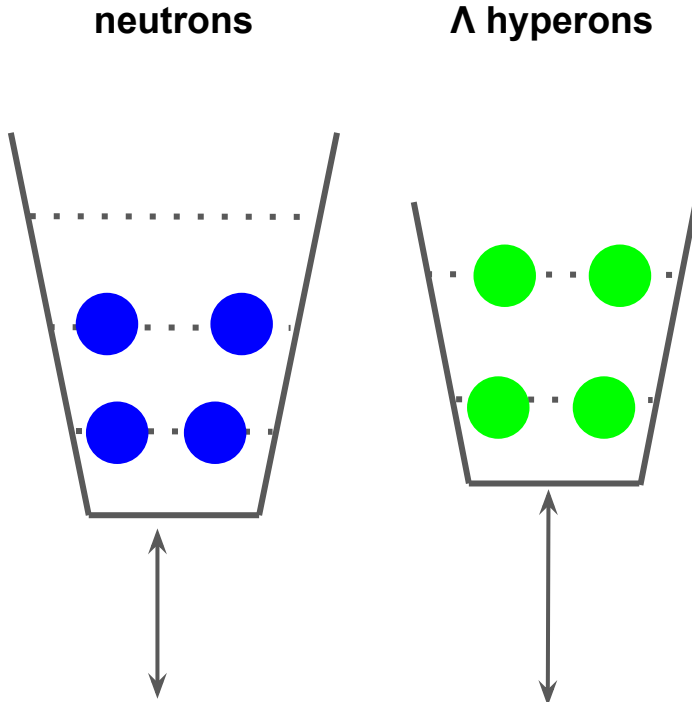


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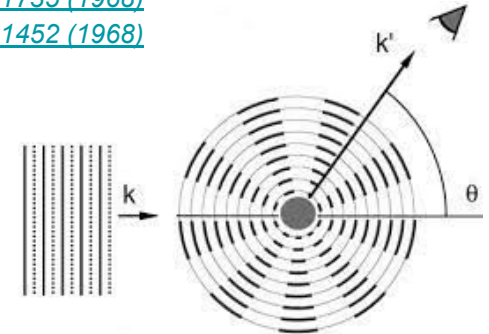


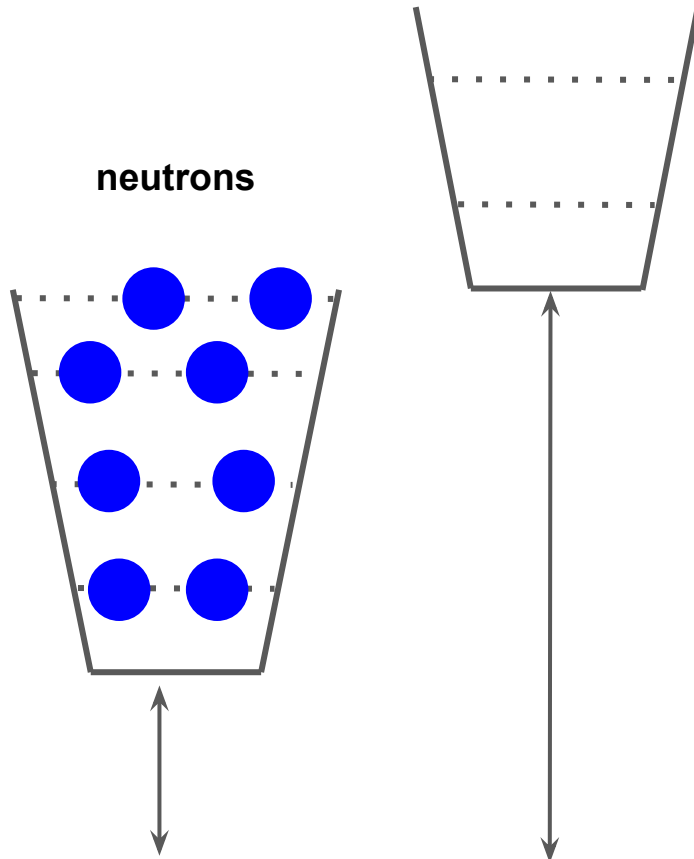


- As density increases, so does the Fermi energy and chemical potential.
- It might become **possible to form hyperons**.
- Their effective **in-medium mass depends on the interaction** with the surrounding particles.
- For the **attractive $N\Lambda$ interaction**, more Λ formation.
Known from e.g. scattering experiments

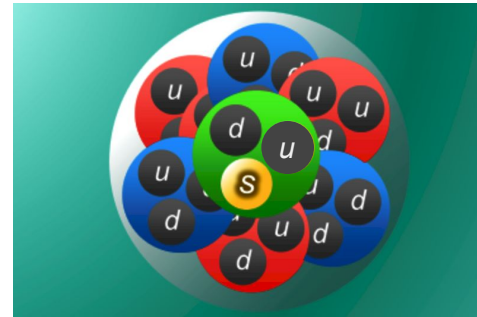
[Phys. Rev. 175, 1735 \(1968\)](#)

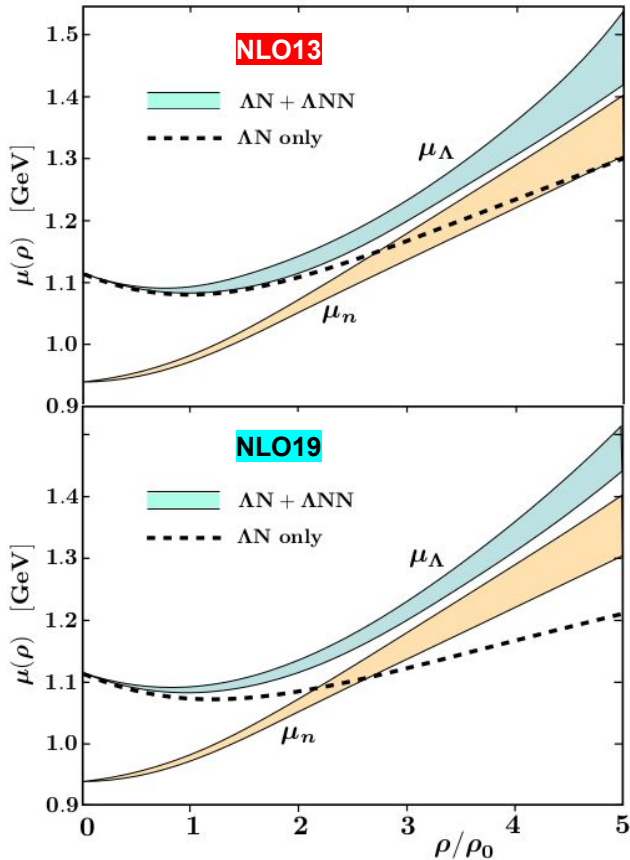
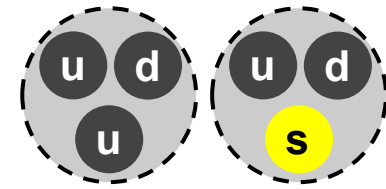
[Phys. Rev. 173, 1452 \(1968\)](#)





- As density increases, so does the Fermi energy and chemical potential.
- It might become **possible to form hyperons**.
- Their effective **in-medium mass depends on the interaction** with the surrounding particles.
- Any **repulsion**, e.g. due to many- (three-) body forces, may prohibit their formation.
In agreement with hypernuclei experiments

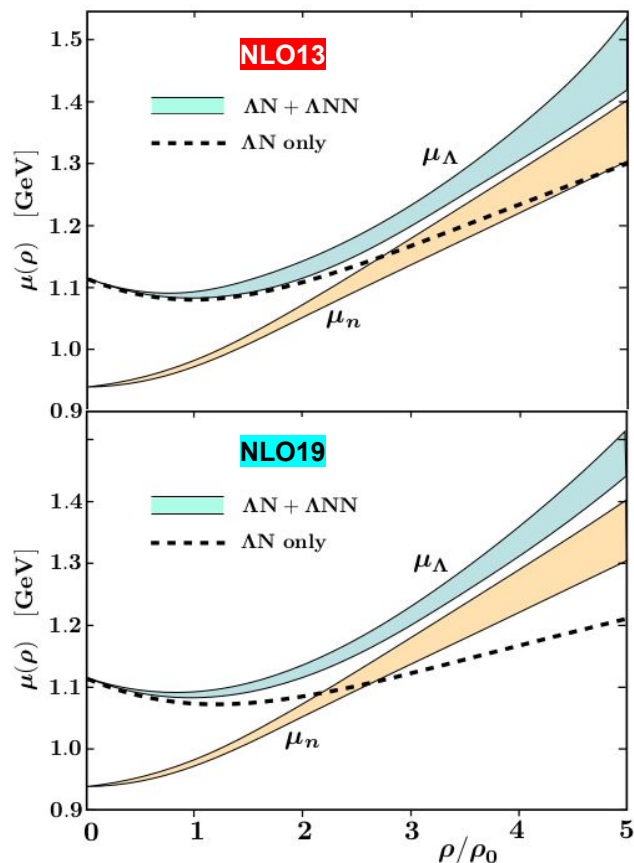
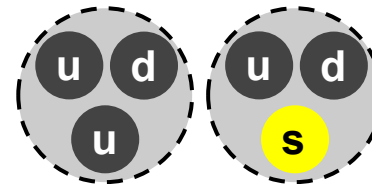




[Wed-II: Wolfram Weise](#)

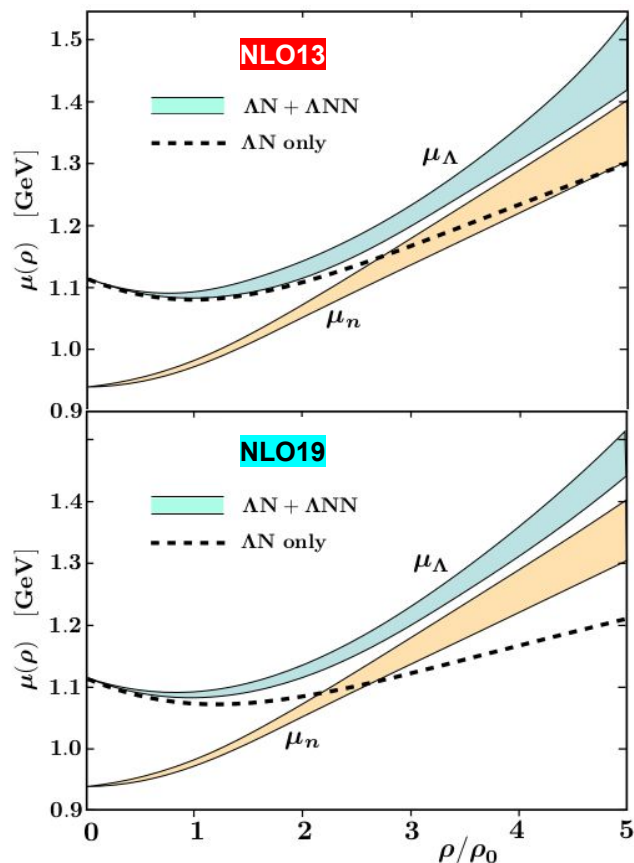
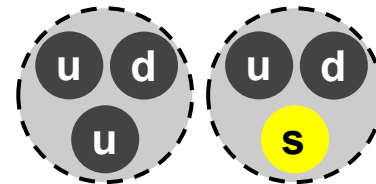
- The Next-to-Leading Order (NLO) calculation can be fine tuned to reproduce existing data using different parameters.

[Wed-I: Johann Haidenbauer](#)



[Wed-II: Wolfram Weise](#)

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- **NLO13** has slightly stronger 2-body attraction in vacuum. **NLO19** leads to stronger 3-body repulsion in-medium.
- Within this model, Λ s cannot form inside neutron stars!
This will explain the existence of measured massive neutron stars ($M > 2 M_{\odot}$).



[Wed-II: Wolfram Weise](#)

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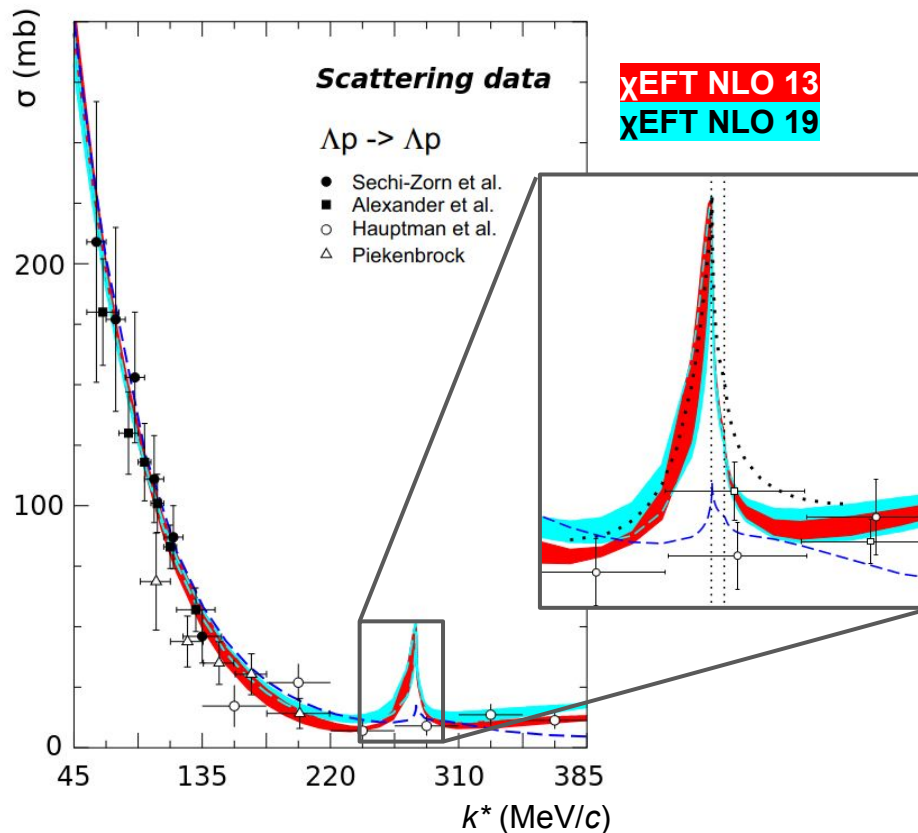
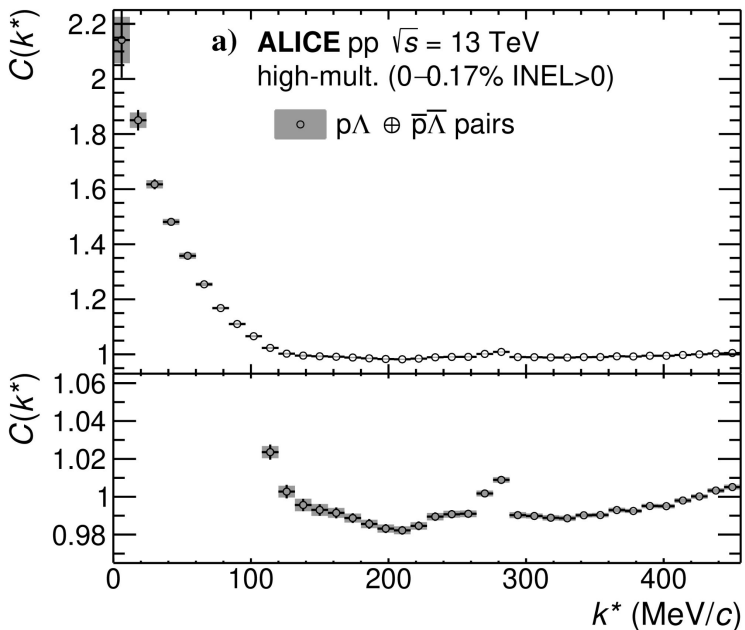
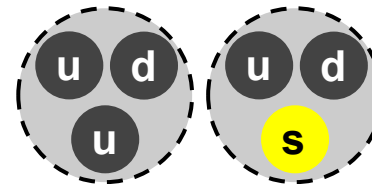
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This will explain the existence of measured massive neutron stars ($M > 2 M_\odot$).

- **Experimental data is needed for both the 2-body and 3-body interaction** to obtain any quantitative conclusions.

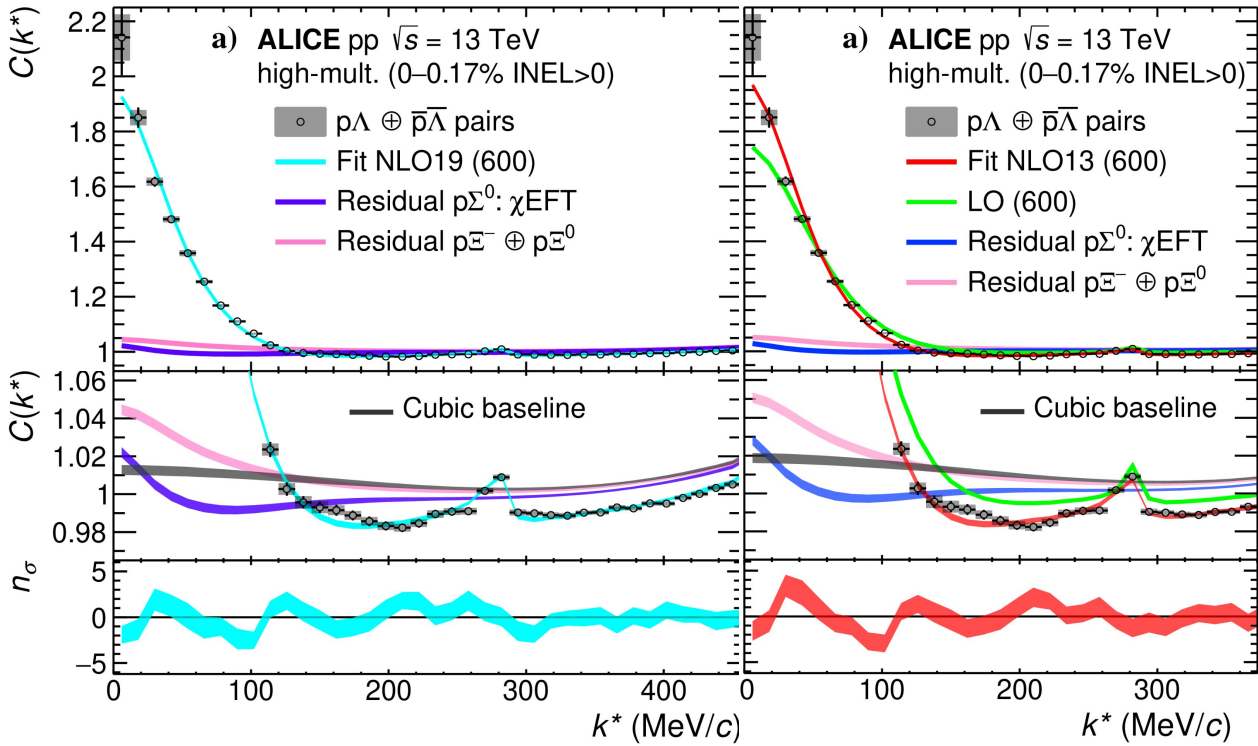
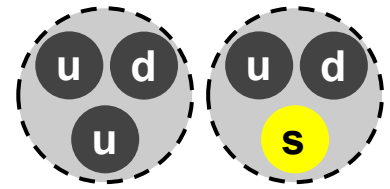
[Tue-II: Koji Miwa](#)

[Tue-II: Nick Zachariou](#)



TUM $p\Lambda$ interaction

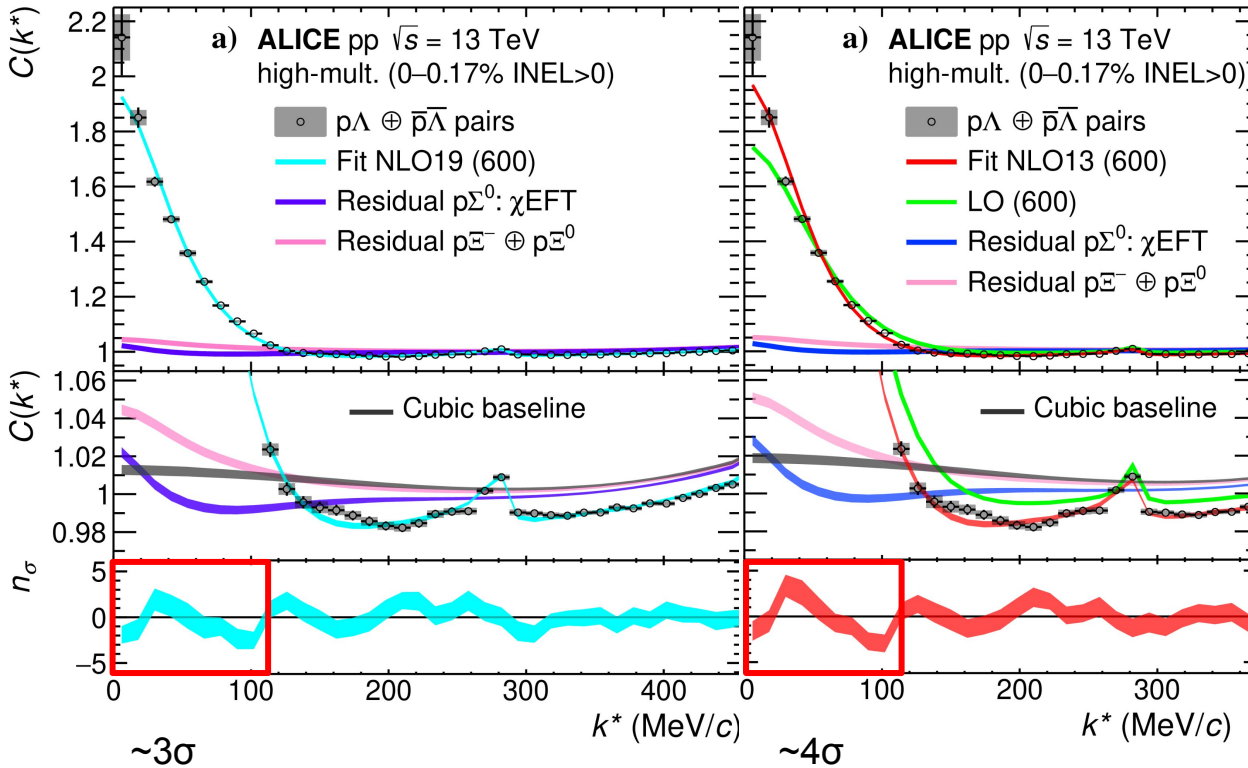
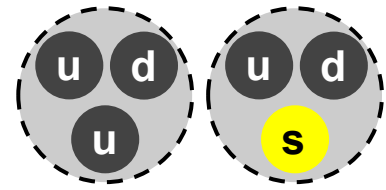
Results



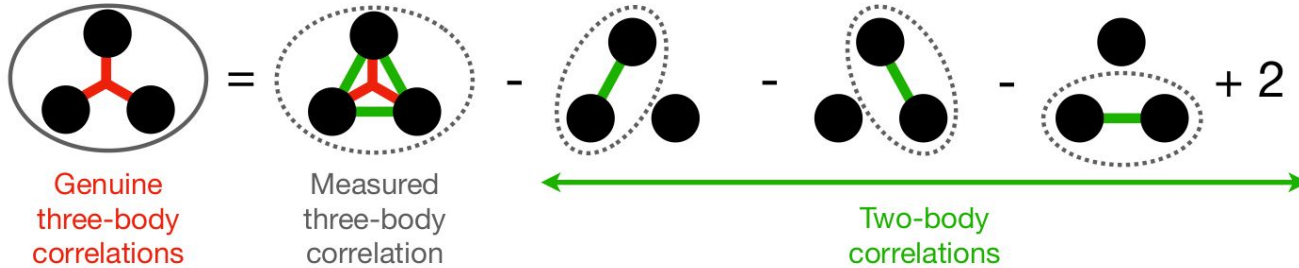
- Observation of the $N\Lambda \leftrightarrow N\Sigma$ cusp.
- Superior precision at low momenta over existing data.

TUM $p\Lambda$ interaction

Results



- Observation of the $N\Lambda \leftrightarrow N\Sigma$ cusp.
- Superior **precision at low momenta** over existing data.
- Preference towards the NLO19. *Differences in the coupling to $N\Sigma$, and in the interplay between two- and three-body forces. Important for the equation of state.*
- NLO19 deviates by $\sim 3\sigma$ at low k^* . ***Further improvement of the model is possible!***



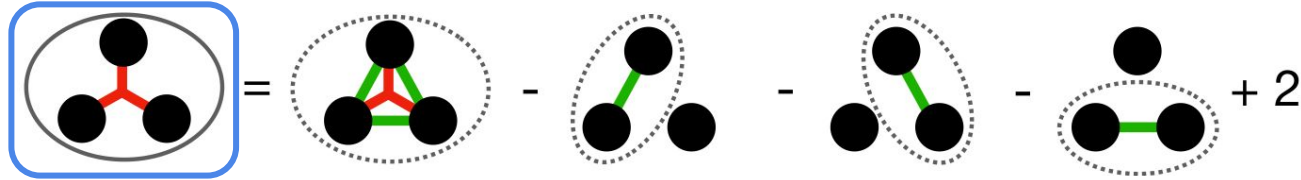
Kubo's **cumulant** expansion method to extract the **genuine 3-body** interaction

[J. Phys. Soc. Jpn. 17, pp. 1100-1120 \(1962\)](#)

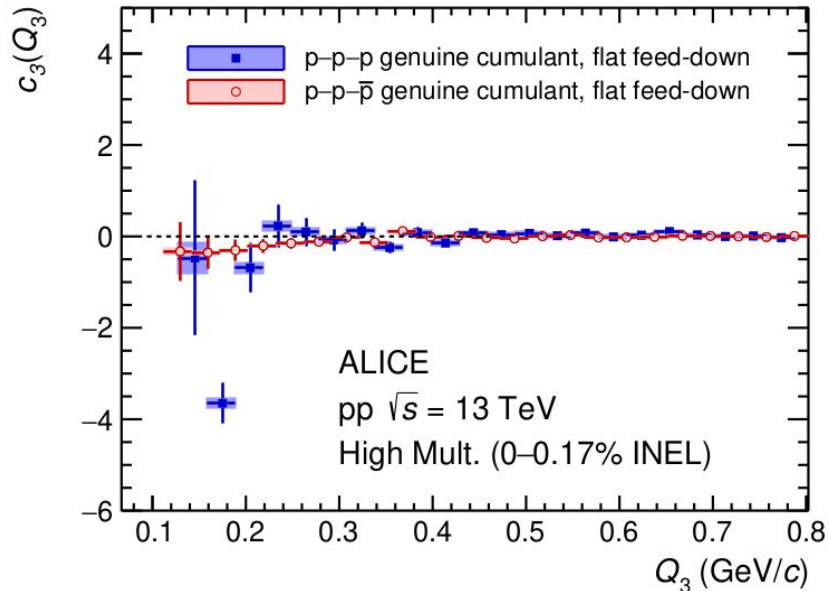
[Wed-II: Laura Šerkšnytė](#)

TUM pp Λ three-body interaction

Going into the future

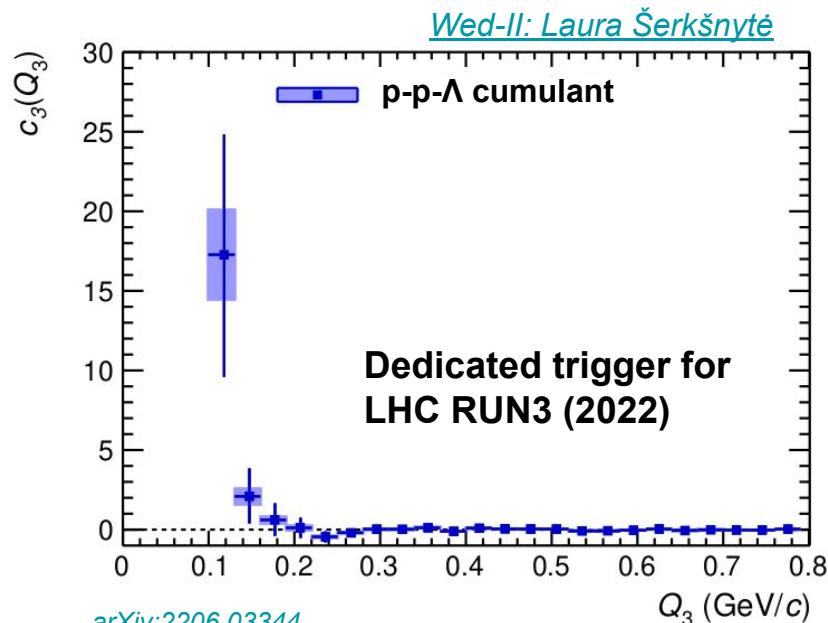
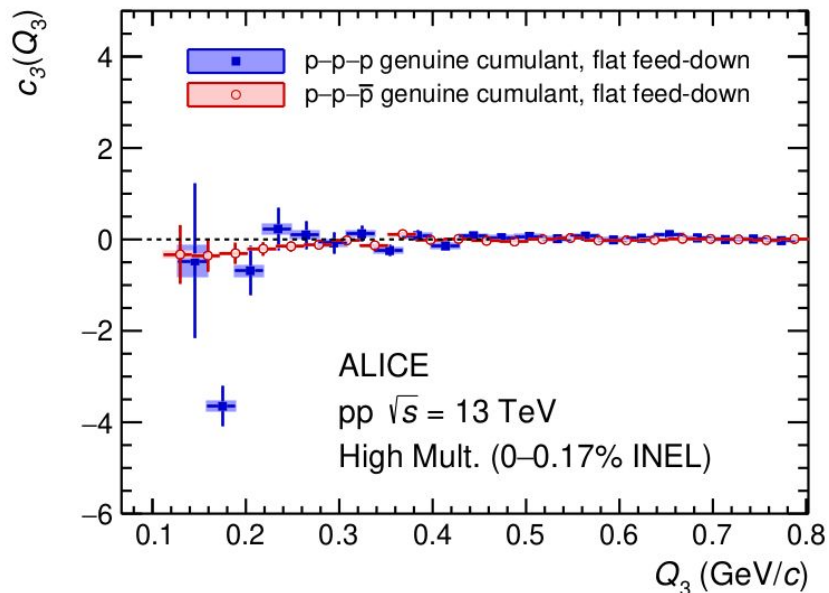


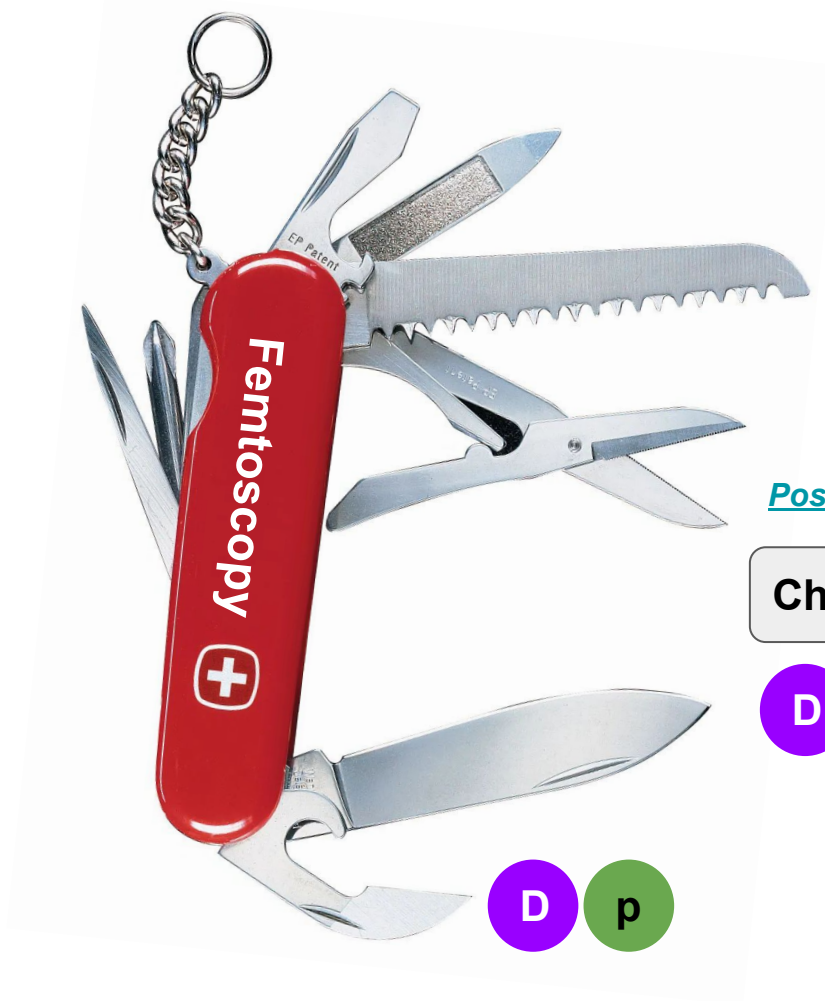
Wed-II: Laura Šerkšnytė



- Clear three body effect in ppp, in contrast to pp \bar{p} .

- No significant deviation from zero.





Poster: Daniel Battistini

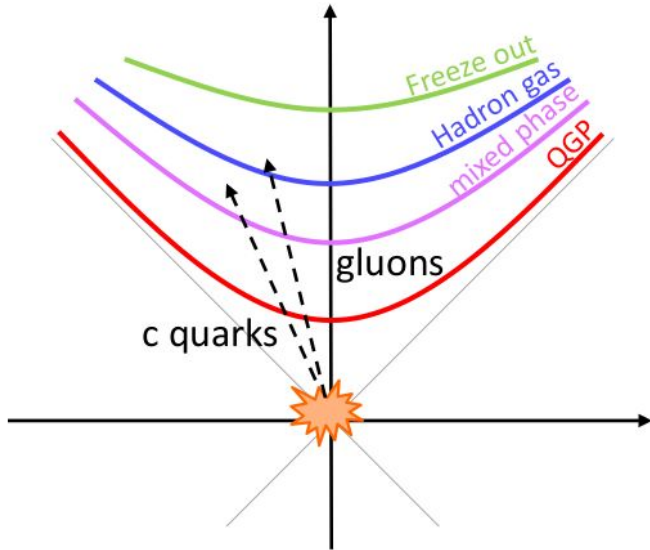
Charming

D

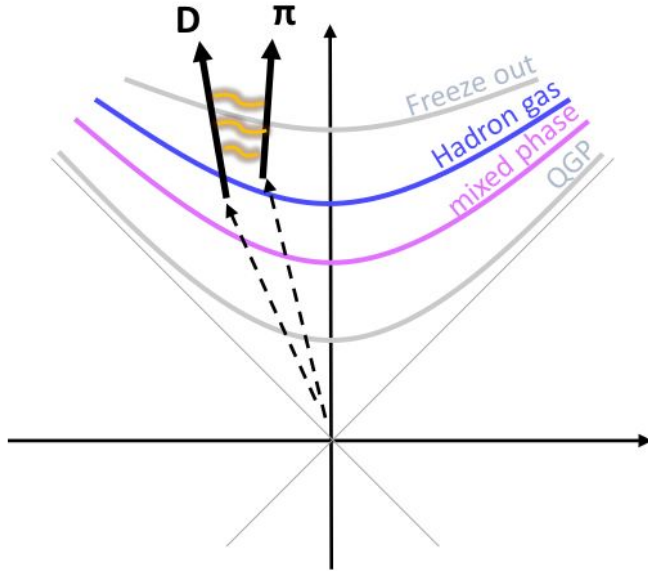
π

D

p



- Charm quarks: in heavy ion collisions, produced before the formation of the quark-gluon plasma (QGP).
⇒ **Ideal probes of the QGP**

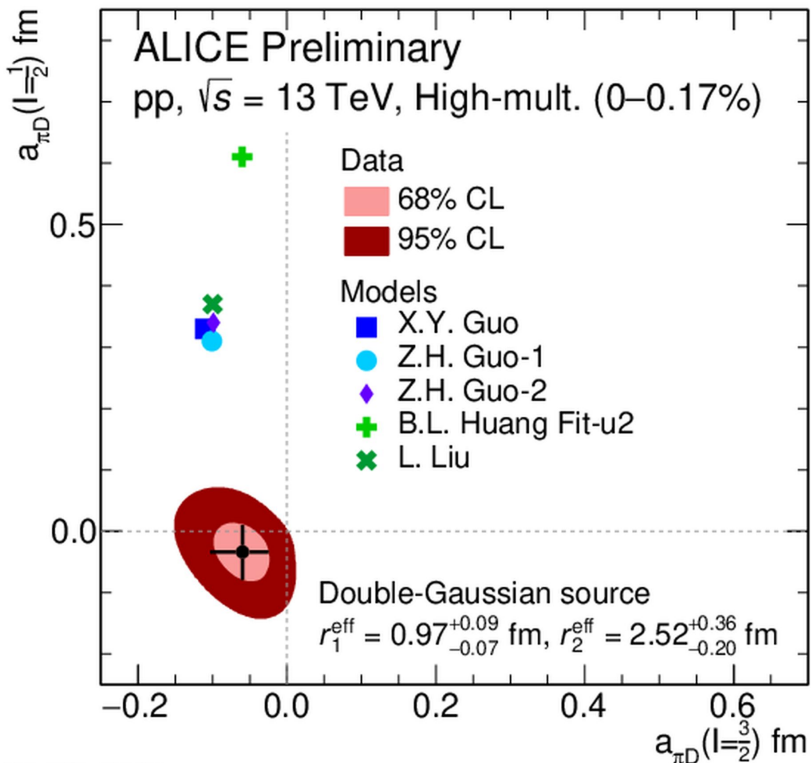


- Charm quarks: in heavy ion collisions, produced before the formation of the quark-gluon plasma (QGP).
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- After hadronization, charm hadrons still interact with the light hadrons, which has to be accounted for.
 ⇒ **No experimental constraints**



- Charm quarks: in heavy ion collisions, produced before the formation of the quark-gluon plasma (QGP).
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- After hadronization, charm hadrons still interact with the light hadrons, which has to be accounted for.
⇒ **No experimental constraints**
- **Lednický model:**
Relate the correlation function to the scattering parameters.
[Sov. J. Nucl. Phys., 35:770, 1982](#)

Poster: Daniel Battistini



- Simultaneous analysis of π^+D^+ ($l=1/2$) and π^+D^- (67% $l=1/2$ and 33% $l=3/2$).
- Scattering length for $l=3/2$ in agreement with lattice calculations.
- Scattering length for $l=1/2$ significantly smaller than lattice calculations.

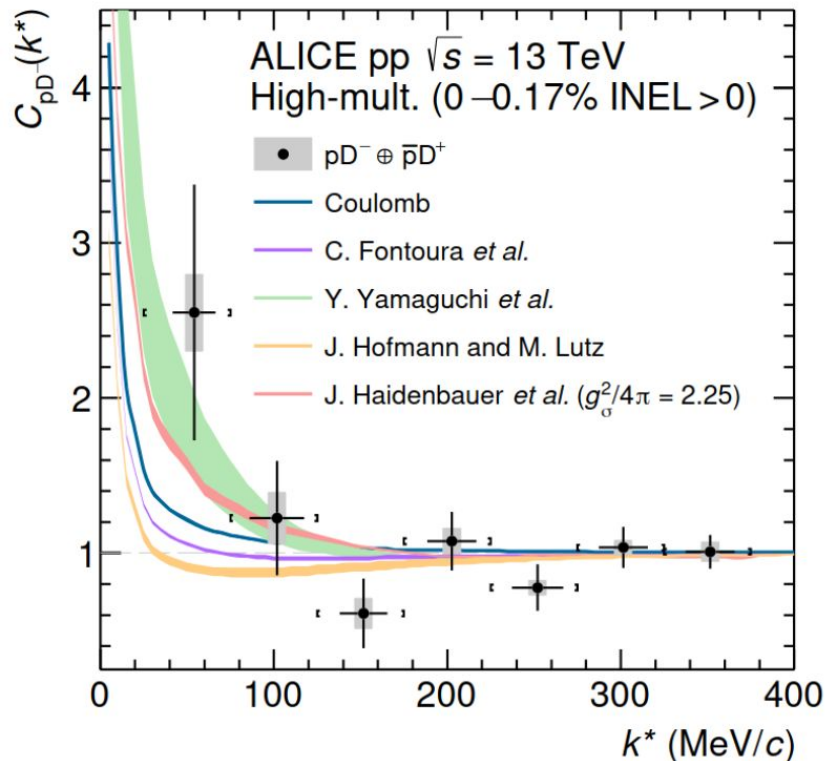
[L. Liu et al, Phys. Rev. D87 \(2013\) 014508](#)

[X.-Y. Guo et al, Phys. Rev. D 98 \(2018\) 014510](#)

[B.-L. Huang et al, Phys. Rev. D 105 \(2022\) 036016](#)

[Z.-H. Guo et al Eur. Phys. J. C \(2019\) 79:13](#)

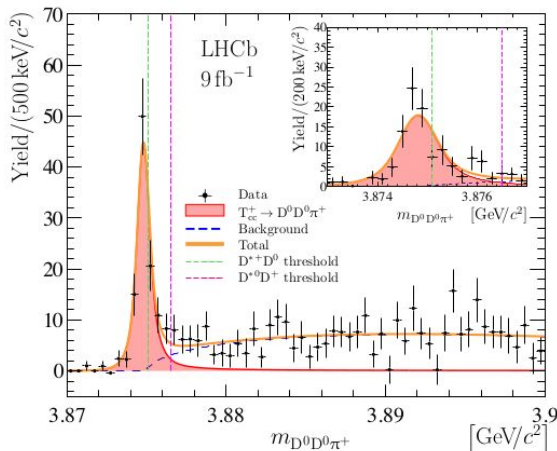
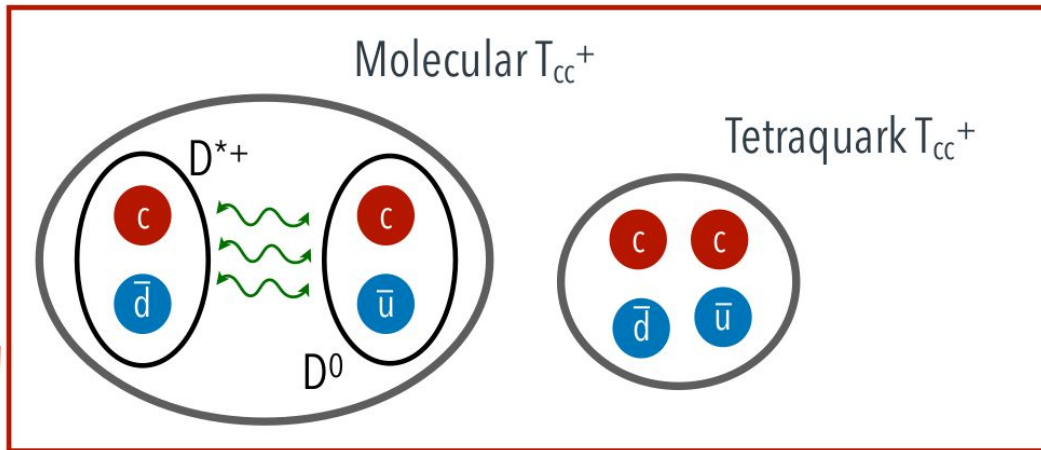
Poster: Daniel Battistini, [arXiv: 2201.05352](https://arxiv.org/abs/2201.05352)



- Compatible with Coulomb only.
- Preference towards models with attractive, perhaps binding, strong interaction.
- To be clarified with LHC RUN3 data!

- Charm molecules?

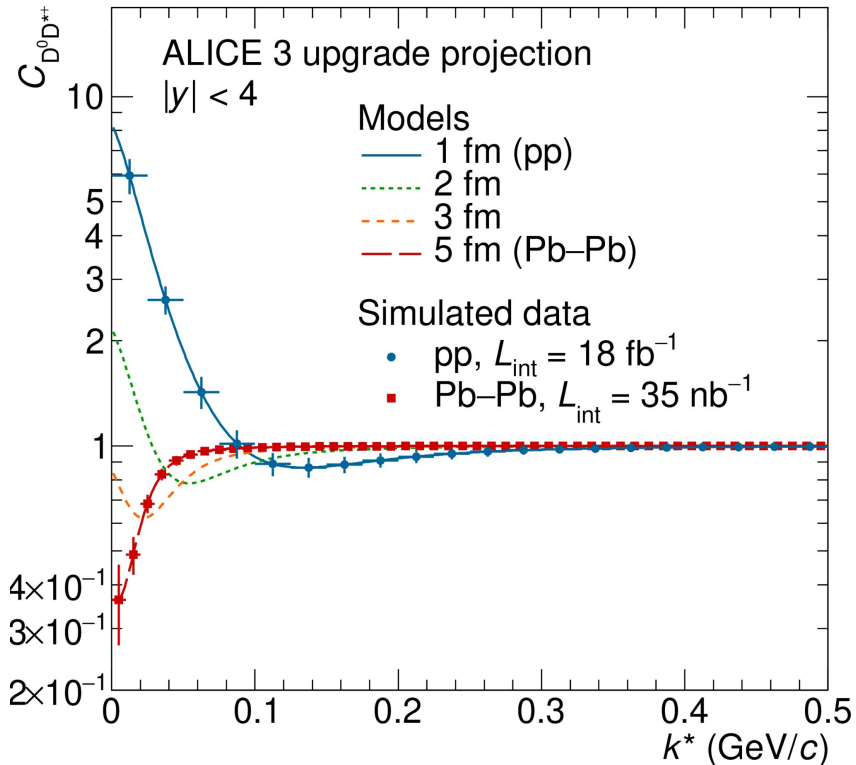
System	$I(J^{P(C)})$	Candidate
np	$0(1^+)$	deuteron
ND	$0(1/2^-)$	$\Lambda_c(2765)$
ND*	$0(3/2^-)$	$\Lambda_c(2940)$
ND	$0(1/2^-)$	$\Sigma_c(2800)$
$D^*\bar{D}$	$0(1^{++})$	$X(3872)$
D^*D	$0(1^+)$	T_{cc}
$D_1\bar{D}$	$0(1^{--})$	$Y(4260)$
$D_1\bar{D}^*$	$0(1^{--})$	$Y(4360)$
$\Sigma\bar{D}$	$1/2(1/2^-)$	$P_c(4312)$
$\Sigma\bar{D}^*$	$1/2(1/2^-)$	$P_c(4457)$
$\Sigma\bar{D}^*$	$1/2(3/2^-)$	$P_c(4440)$



- Just below the D^*D threshold
Ideal candidate to be a molecular state

- Charm molecules?

[ALICE3 LOI: CERN-LHCC-2022-009](#)



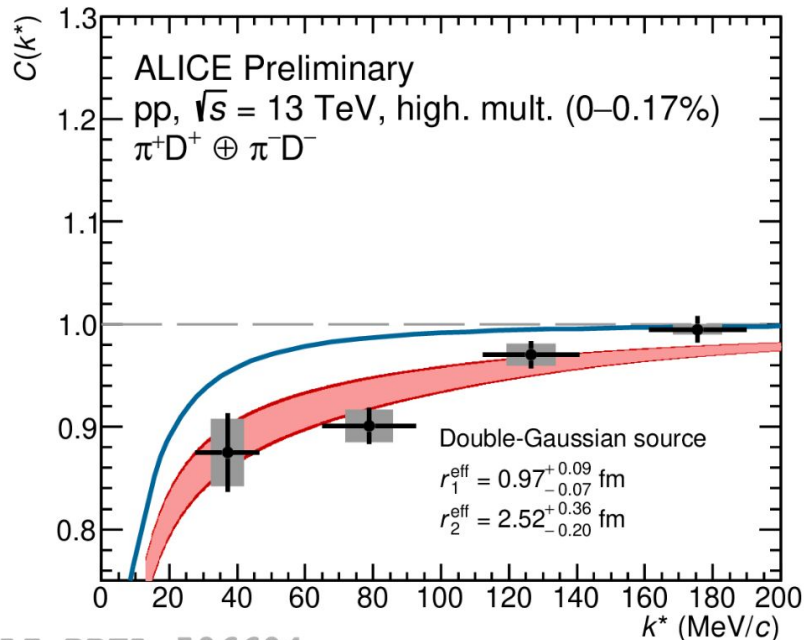
- ALICE 3: large acceptance, high luminosity, excellent spatial resolution.
- Run 5: ideal laboratory for the measurement of charm-hadron momentum correlations in different colliding systems.
- Source size dependent modification of the correlation function in presence of a bound state.

[Wed-I: Yuki Kamiya](#)

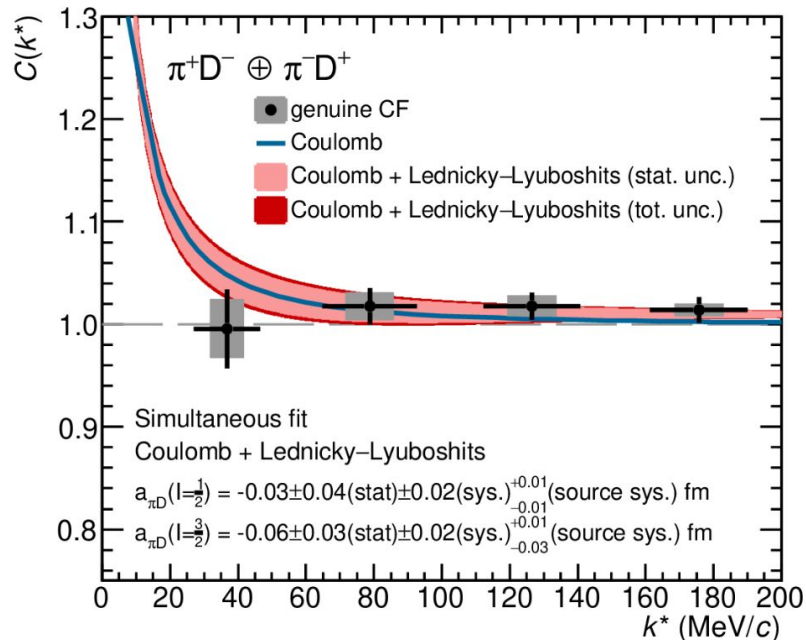
[Yuki Kamiya et al. arXiv:2203.13814](#)

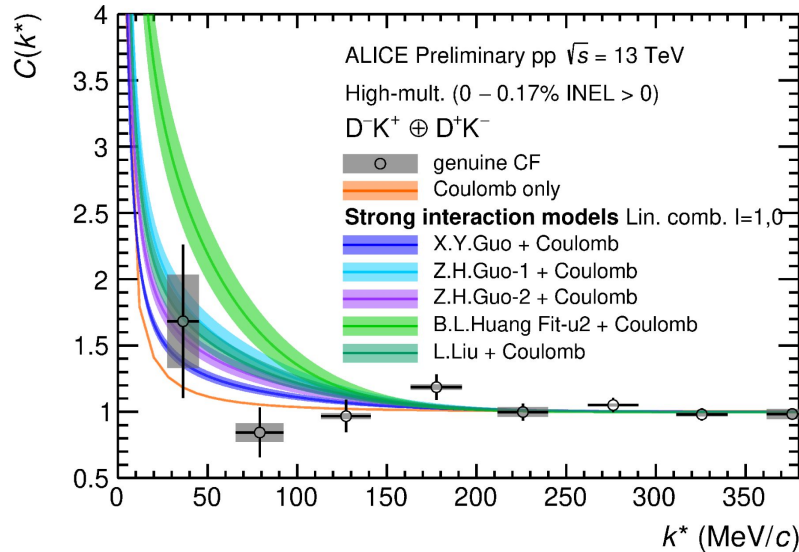
- Particle correlations as a tool to study two- and three-body interactions among any hadrons measured at accelerators.
[Wed-I: Yuki Kamiya](#), [Wed-II: Laura Šerkšnytė](#)
- **p Λ and pp Λ measurements** allowing to constrain existing χ EFT models, helping to describe the composition of neutron stars.
- Further applications:
 - KN interaction and kaonic bound states
[Wed-II: Laura Šerkšnytė](#), [Thu-I: Ramona Lea](#), [Fri-I: Oton V. Doce](#)
 - pd correlations
[Wed-III: Bhawani Singh](#), [Poster: Wioleta Rzesza](#)
 - Studying multi-strangeness systems
[Wed-IVa: Georgios Mantzaridis](#)
 - **Testing the nature of charm**
[Poster: Daniel Battistini](#)
 - **Great prospects at LHC RUN3 and beyond**



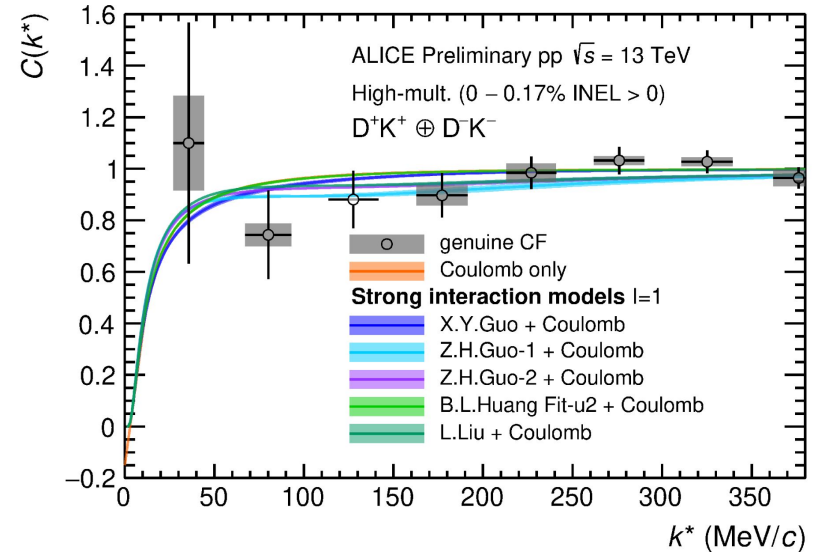


ALI-PREL-506604





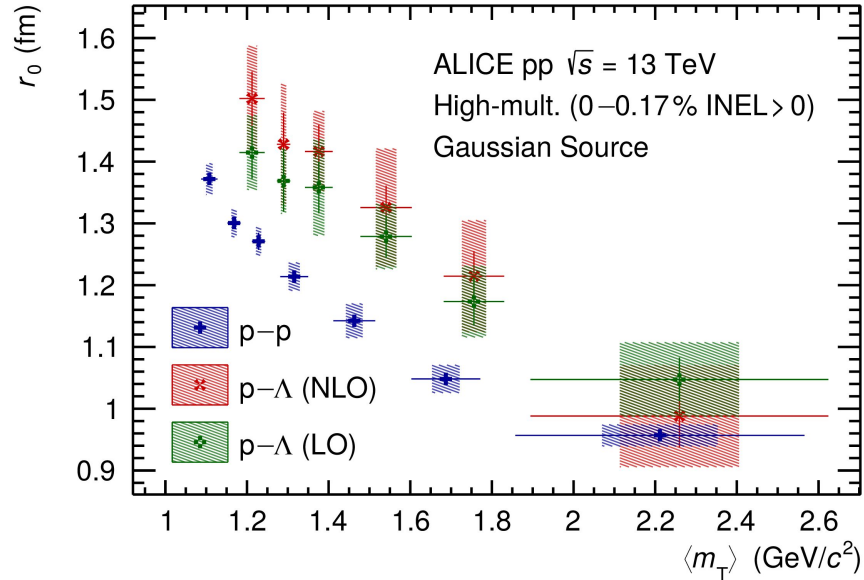
ALI-PREL-506581



ALI-PREL-506586

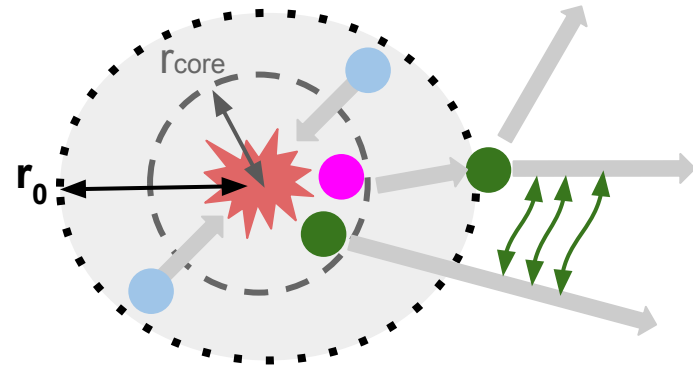
- Statistics is yet insufficient to impose constraints
To be improved during LHC RUN3

Gaussian source



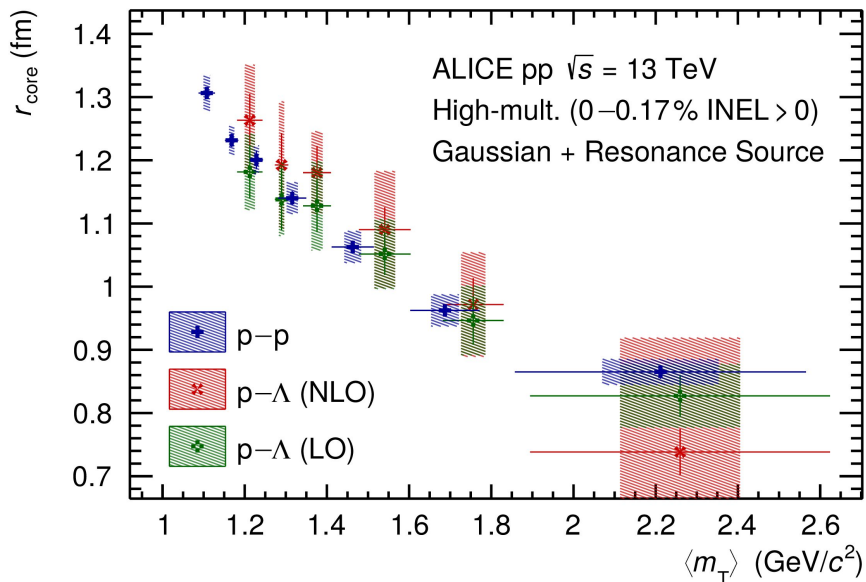
- Different source size for p-p and p- Λ pairs

- The Statistical Hadronization Model tells us: c.a. $\frac{2}{3}$ of protons and Λ s stem from resonances. The average lifetimes (τ) are:
1.6 fm for $X \rightarrow \text{proton}$
4.7 fm for $X \rightarrow \Lambda$
- Production through short-lived resonances



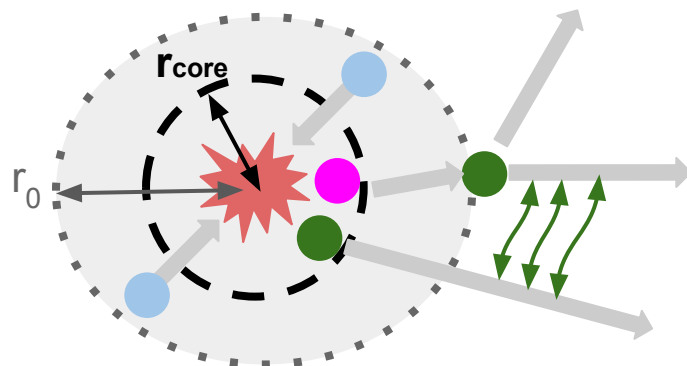
[Phys. Lett. B 811 \(2020\) 135849](#)

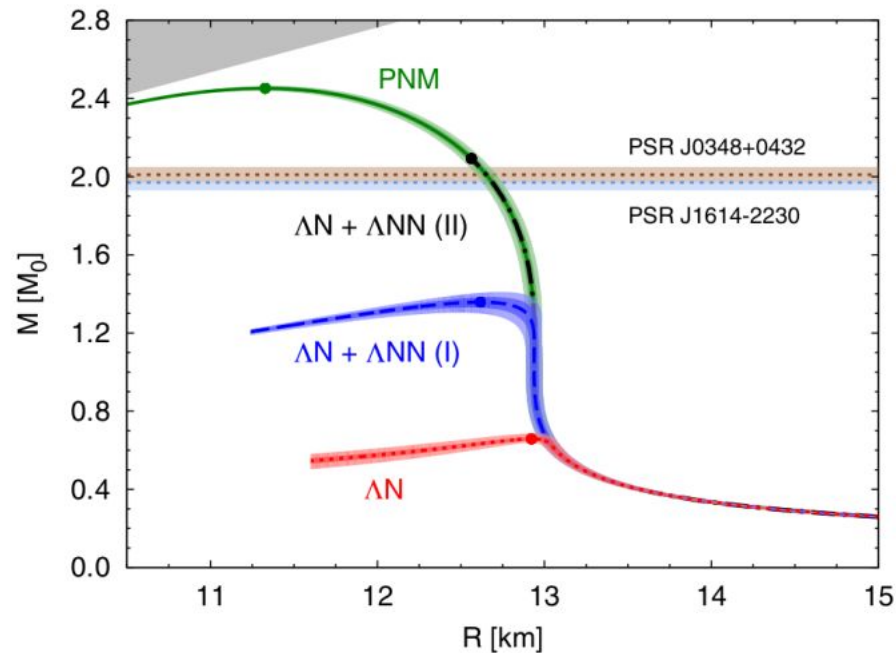
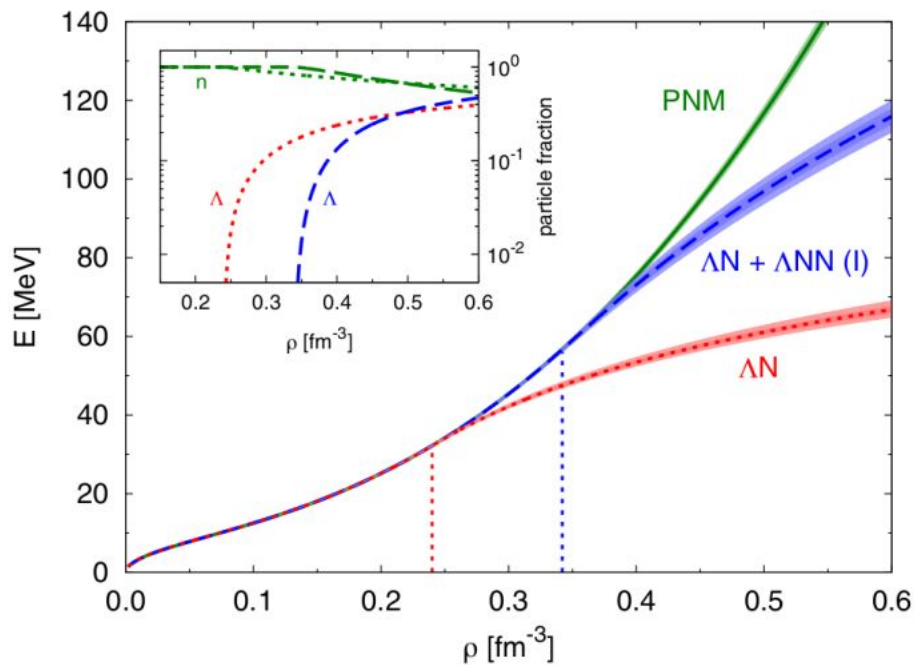
Gaussian core + resonances



- **Common source** for p-p and p-Λ pairs. Ones measured (from p-p), it is fixed for ANY baryon-baryon pair!

- The Statistical Hadronization Model tells us: c.a. $\frac{2}{3}$ of protons and Λs stem from resonances. The average lifetimes (τ) are:
1.6 fm for $X \rightarrow \text{proton}$
4.7 fm for $X \rightarrow \Lambda$
- **Production through short-lived resonances**

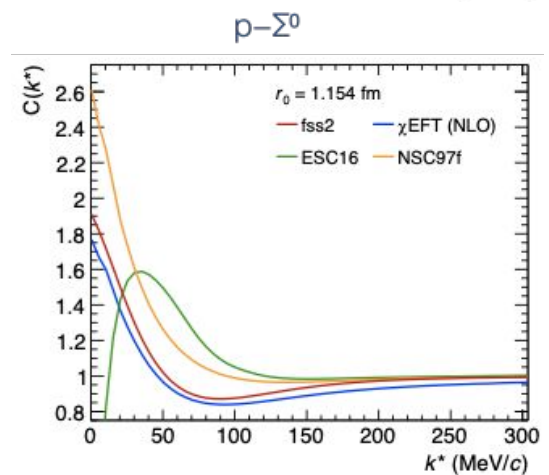
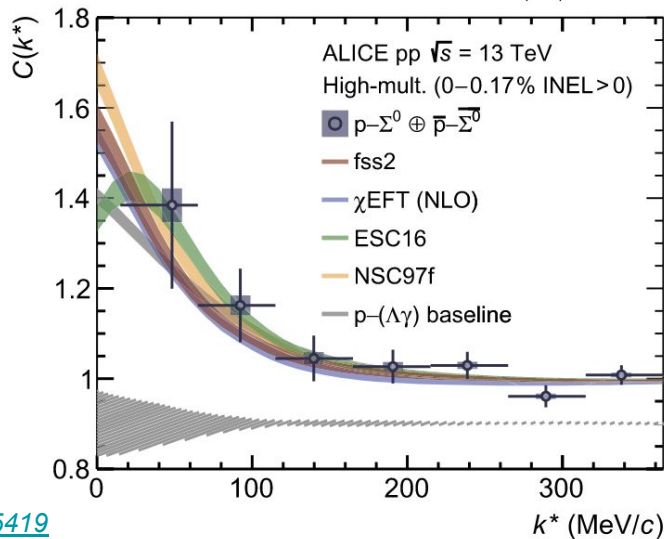
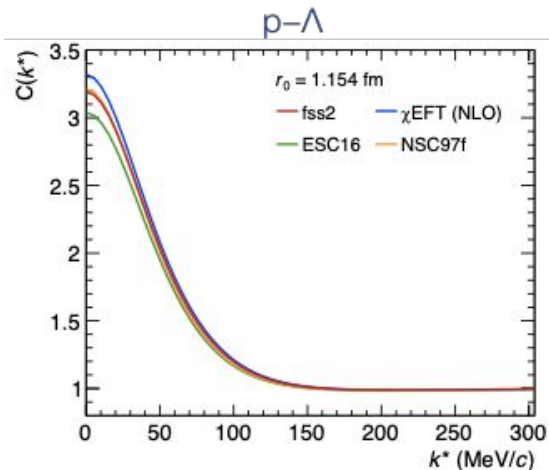
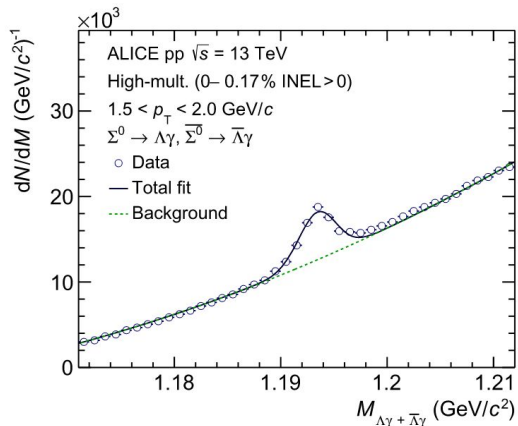




[PRL 114 \(2015\) 9, 092301](#)

A soft EoS, allowing hyperons within NS, leads to an underprediction of measured NS masses. This is known as the **hyperon puzzle**.

$p\Sigma^0$

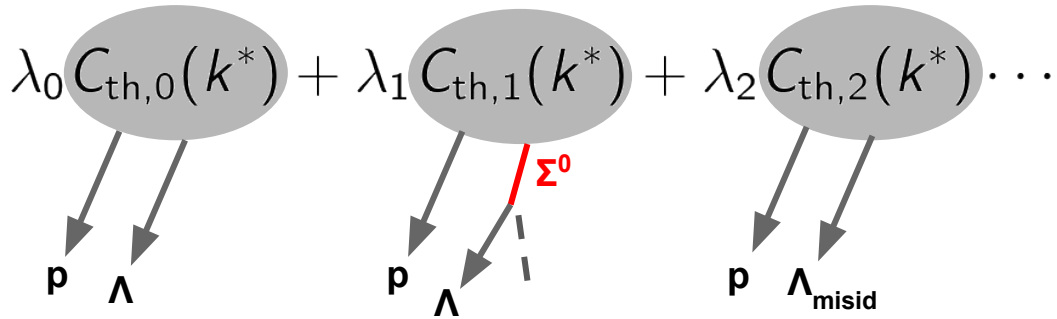


TUM $p\Lambda$ correlation function

Overview

$$C_{\text{exp}}(k^*) = b(k^*) C_{\text{th}}(k^*) = b(k^*) \sum_i \lambda_i C_{\text{th},i}(k^*)$$

Non-femto baseline



Modeled using χEFT

corrected for using a "sideband" analysis