

Searching for the d^* in photoreactions on transversely polarized deuterons

Dan Watts

Mikhail Baskanov, Michael Ostrick, ...



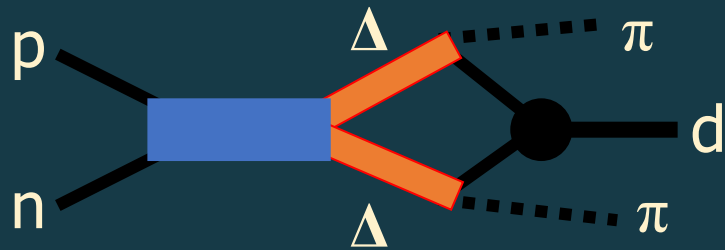
Overview

- Evidence from pn scattering
- EM properties – recent MAMI results
- d^* theory
- Potential added impact: Neutron stars, DM candidate
- Opportunities with nucleon polarisation observables

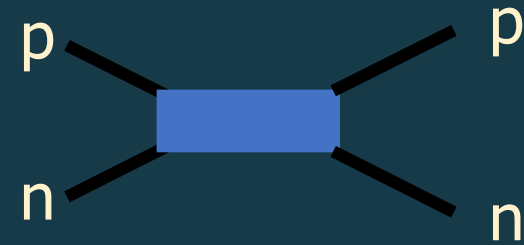
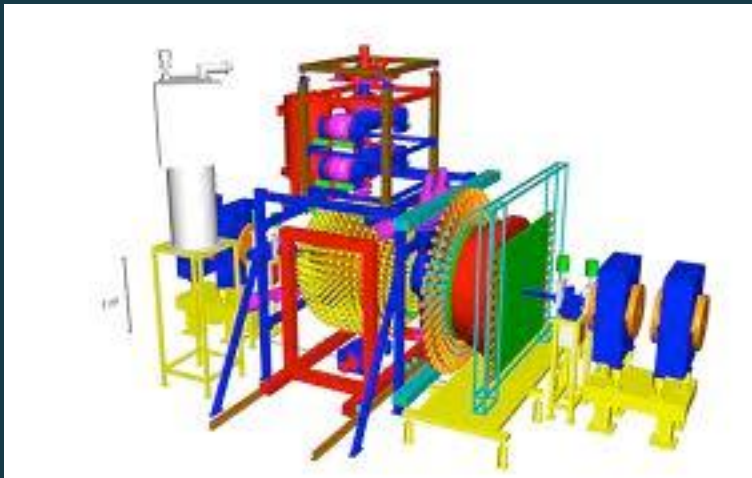
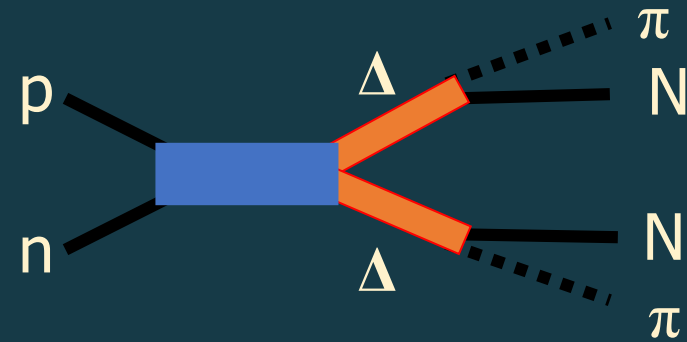
The $d^*(2380)$ in pn scattering

Nucleon scattering with large acceptance

- $pn \rightarrow d^* \rightarrow \Delta\Delta \rightarrow d\pi\pi$

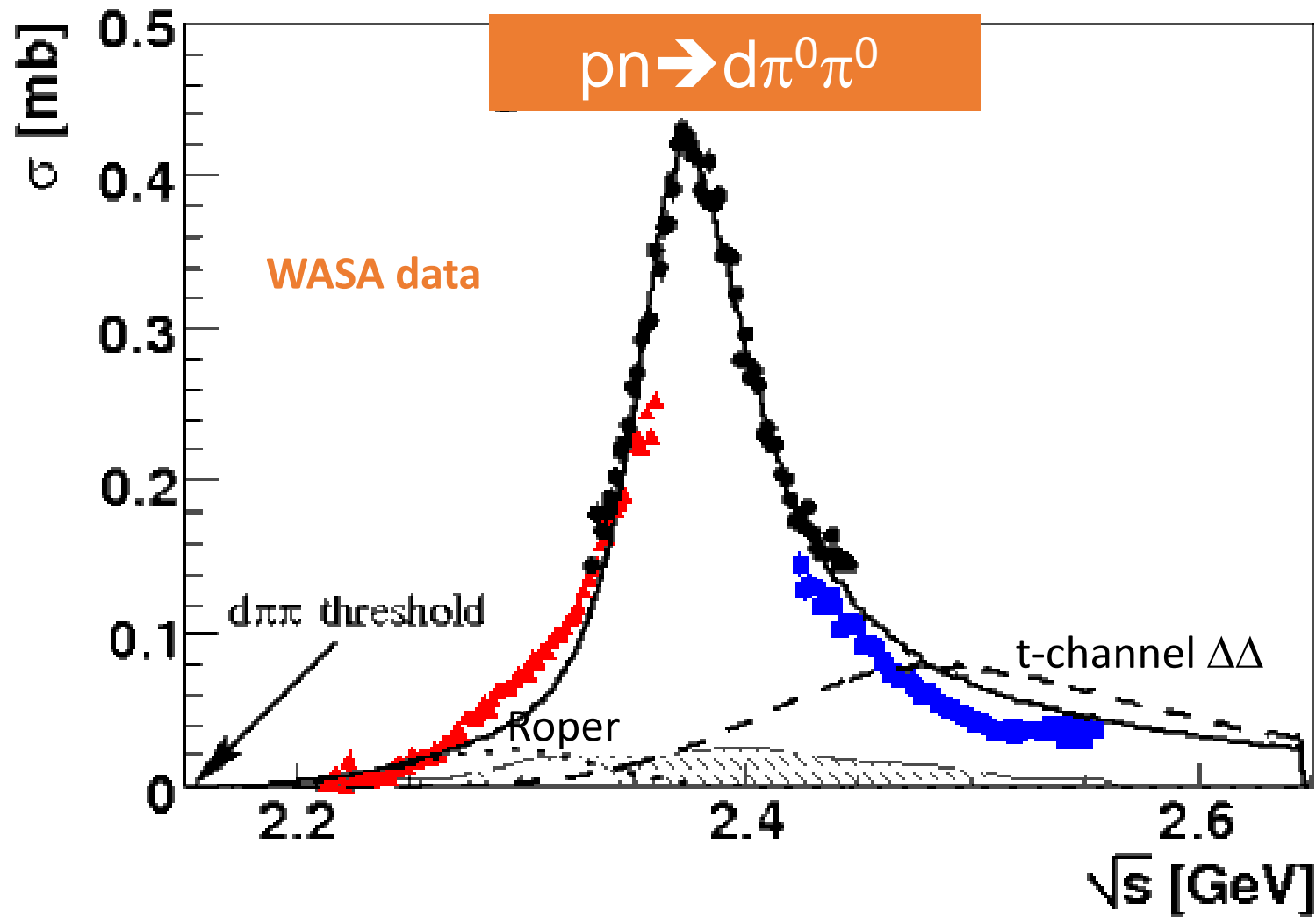


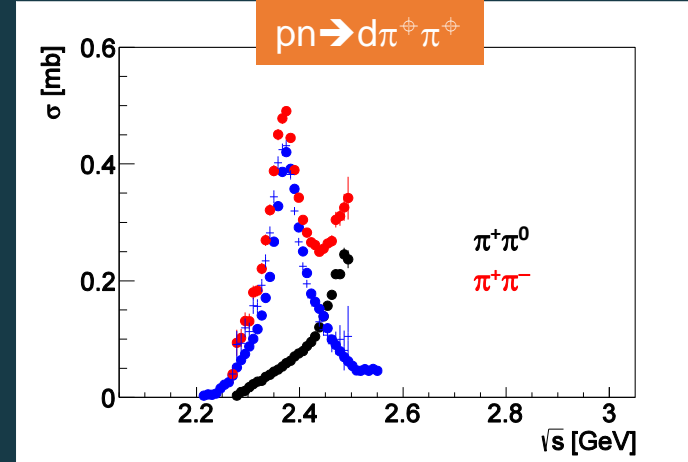
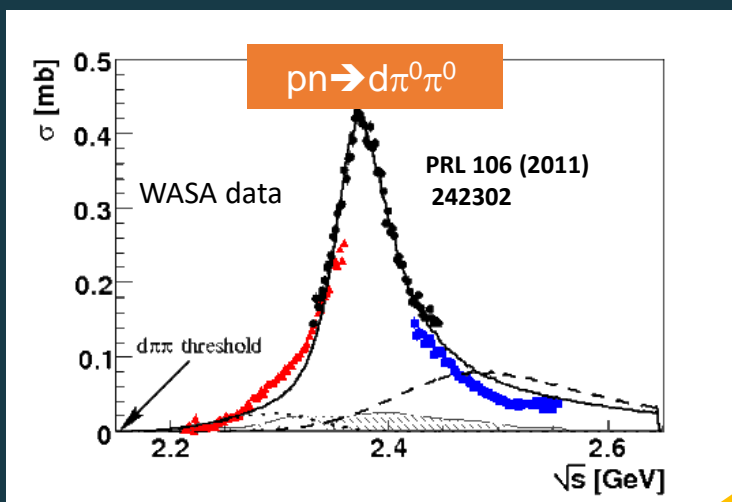
- $pn \rightarrow d^* \rightarrow \Delta\Delta \rightarrow NN\pi\pi$



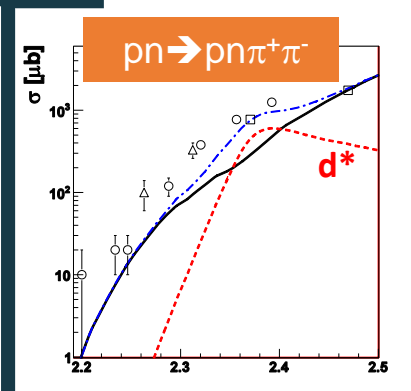
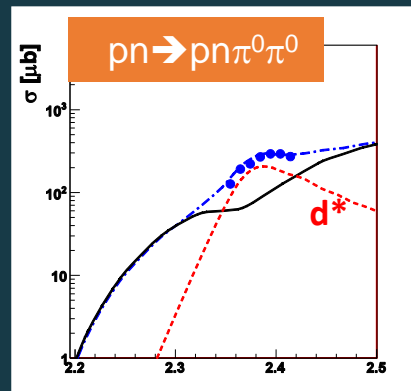
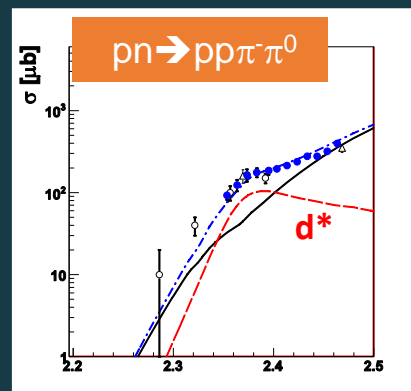
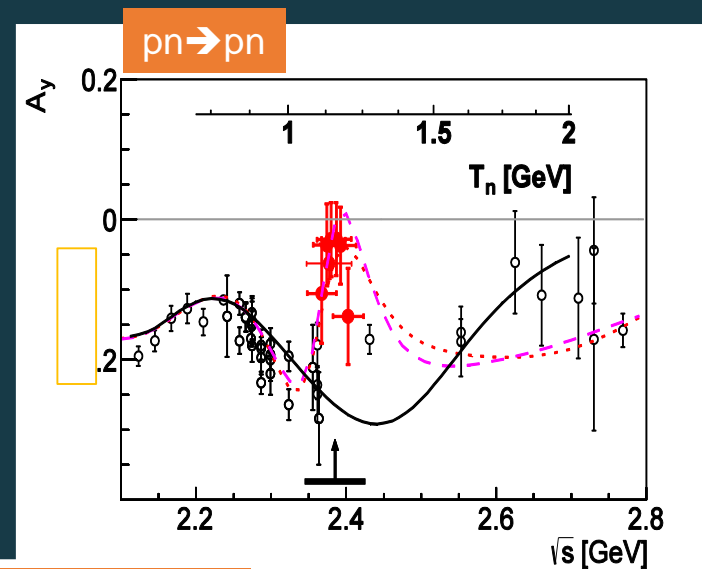
$d^*(2380)$ signals

PRL 106 (2011) 242302





$p + n \rightarrow$ **$d^*(2380)$**
 $I(J^P) = 0(3^+)$



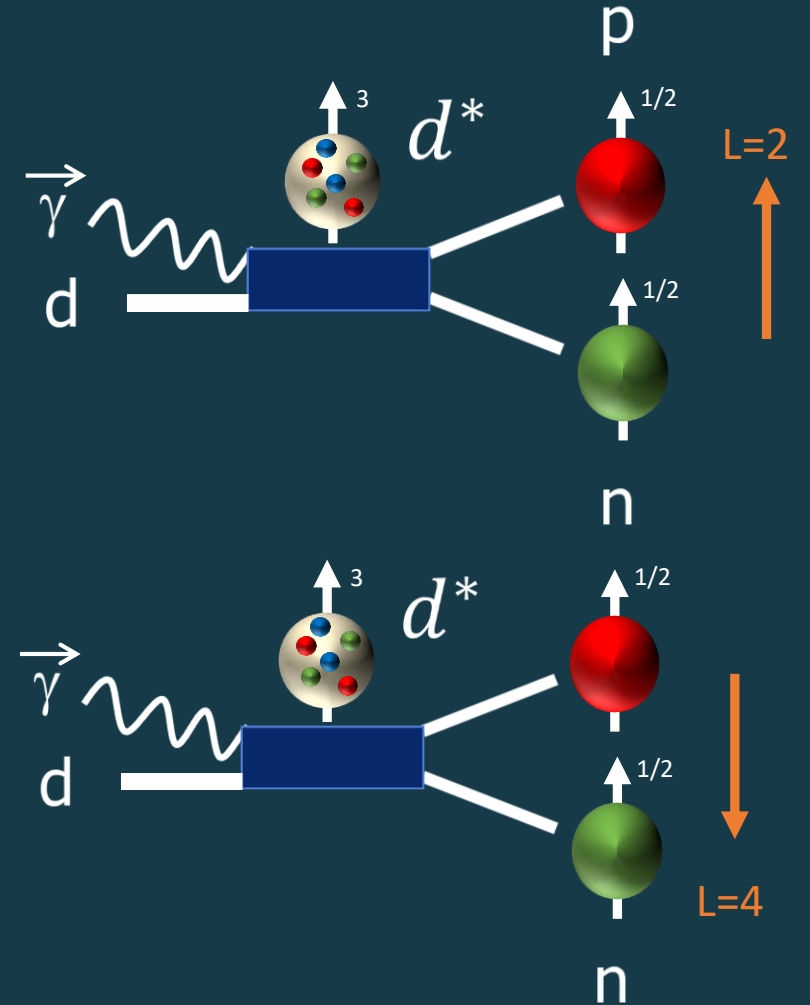
$\Delta\Delta$ decay $\sim 90\%$
 pn decay $\sim 10\%$

d^* decays to pn

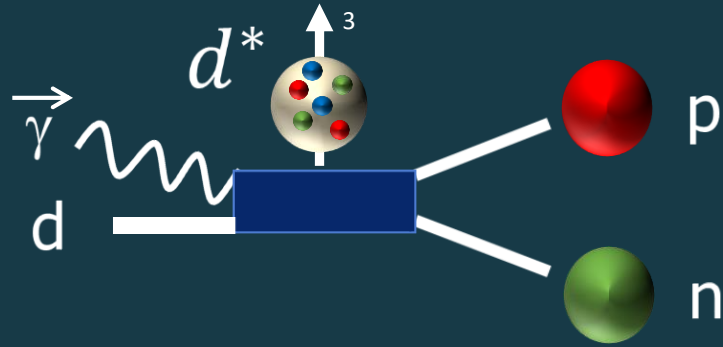
- High partial waves in decay to a proton-neutron final state ($J=3$)

90% of cases this is via the 3D_3 partial wave
($L = 2$, nucleon spins and L all aligned)

10% of cases via the 3G_3 partial wave
($L = 4$, nucleon spins aligned, spin and L anti-aligned).



Deuterium photodisintegration (Σ)



- First detailed measurement in region of d^*
- Good agreement with sparse existing data
- Almost complete acceptance – amenable to (truncated) PWA

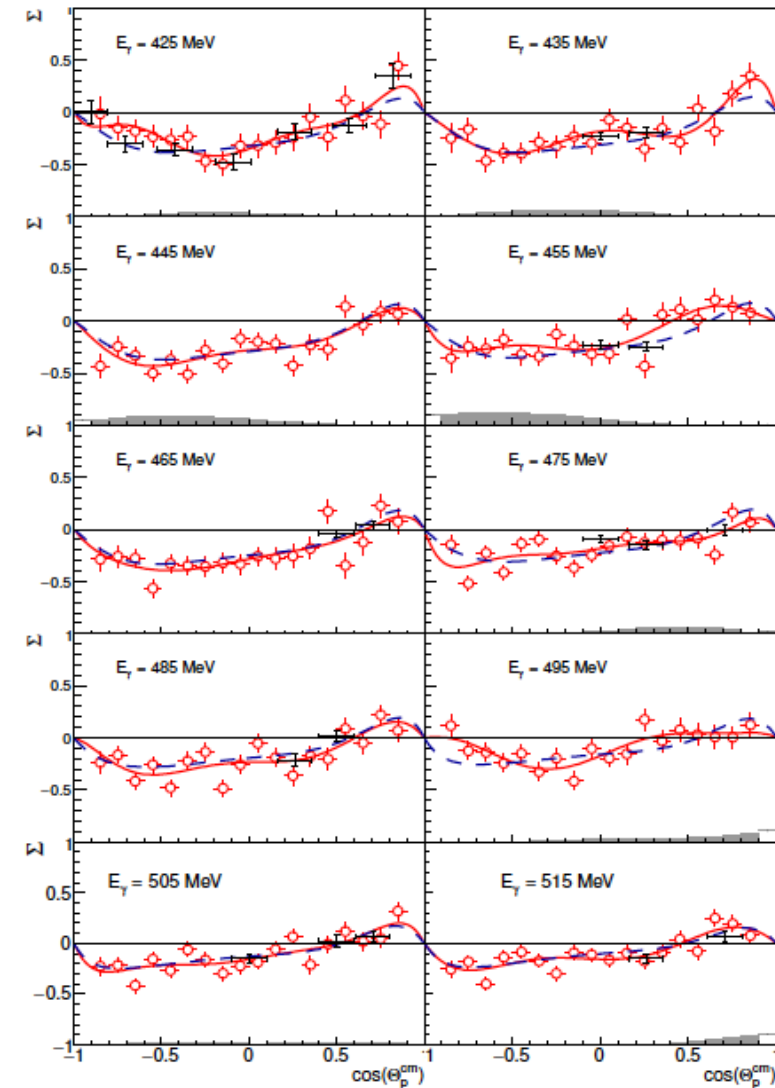


Figure 2: (Color online) Beam-spin asymmetry (Σ) results from this experiment (red open circles) in comparison with previous results (black crosses) [30, 31, 32]. The corresponding systematic uncertainties are depicted as shaded bars on the bottom. Energy independent (energy dependent) $a_1 P_1^2$ fits are shown as solid red (dashed blue) lines (see text).

Σ decomposition



$$\sigma_1 = \Sigma(\Theta, E_\gamma) \cdot \sigma(\Theta, E_\gamma)$$

$$\frac{\sigma_1}{\sigma_{\text{tot}}} = \sum_{l=2}^7 a_l P_l^2.$$

Discrete gaussian sampling method

3 gaussians centroids 420,520,620 width 100 MeV



2 methods: Single energy & energy dependent (smooth functions for a_l)

Additional term: BW, mass 2380 MeV, width=70 MeV, arbitrary strength

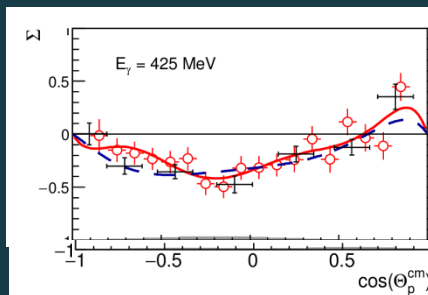
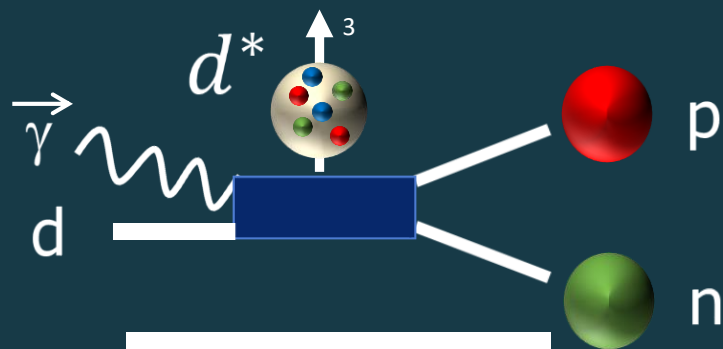
Pure M3 should manifest in a_6 but not a_4

Pure E2 should manifest in a_4 but not a_6

(E2 may contribute in a_6 only with interference with higher multipoles E4 or higher)

$$\begin{aligned} a_6 &\sim d_1 |{}^3G_3(M3)|^2 + \dots \\ &+ d_i |{}^3D_3(M3)| |{}^3G_3(M3)| \cos\delta_i + \dots \\ &+ d_j |{}^3D_3(E2)| |{}^{2S+1}L \geq 4_J(E4)| \cos\delta_j + \dots \end{aligned}$$

Deuterium photodisintegration (Σ)

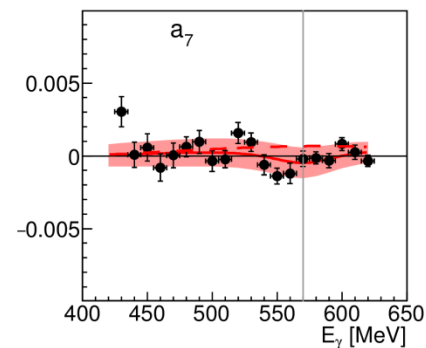
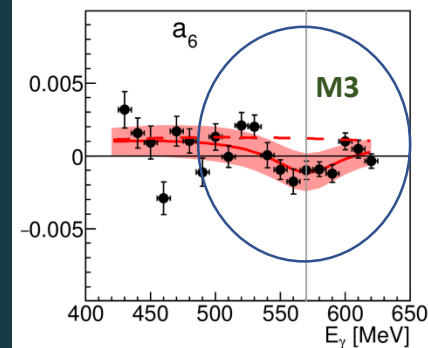
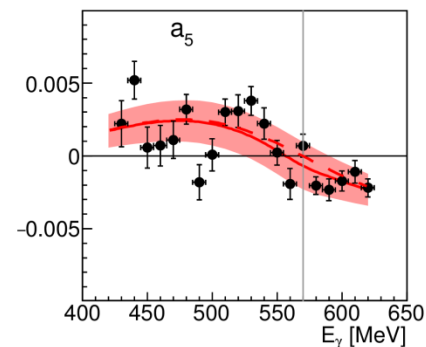
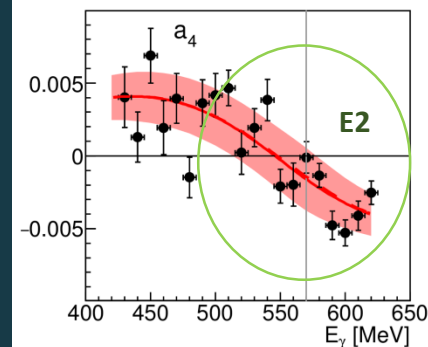
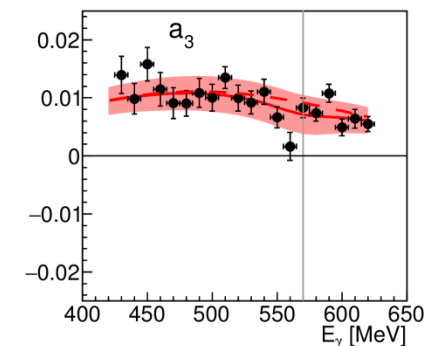
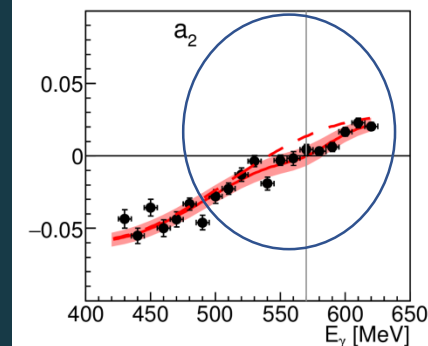


$$\Sigma \sim \sum_{l=2} a_l P_l^2(\cos\Theta)$$

H.Ikeda *et al.*, NPB172,509,(1980)

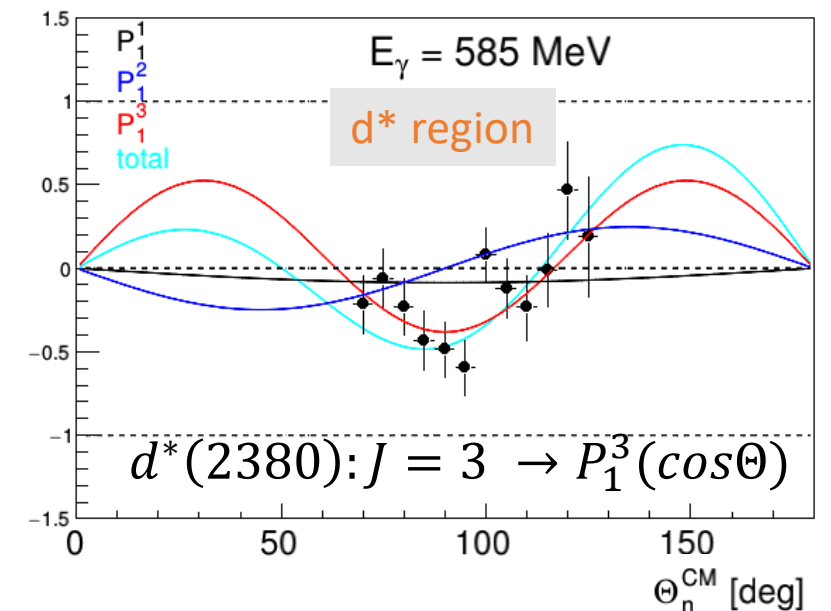
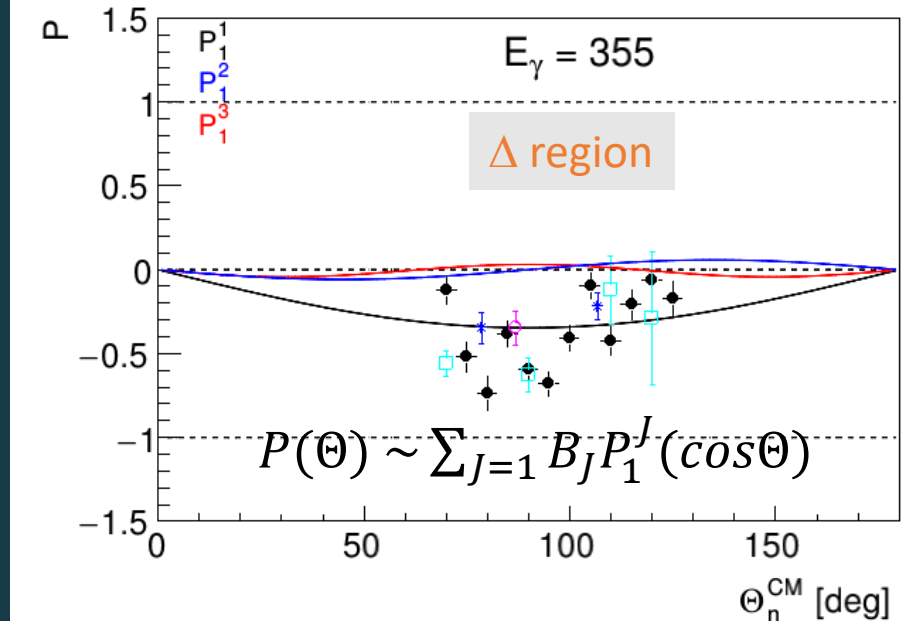
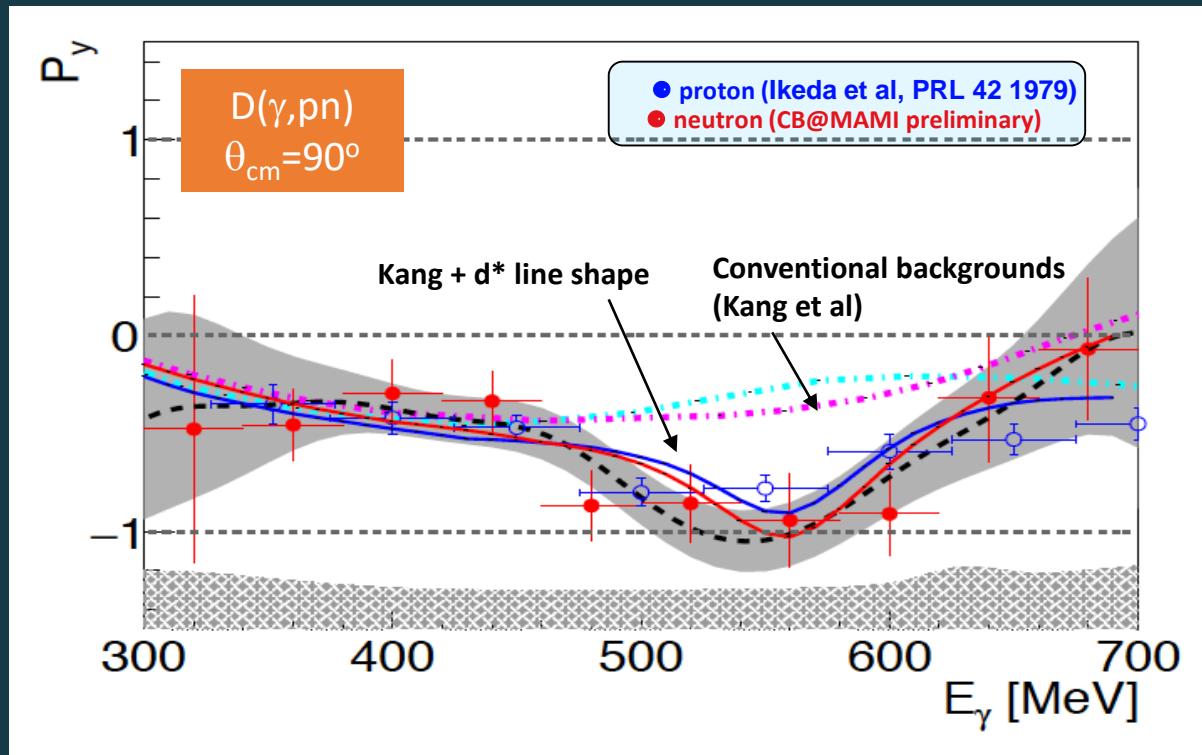
E2 transition \Rightarrow small
M3 transition \Rightarrow dominant

Consistent with $d^*(2380)$
as a compact object



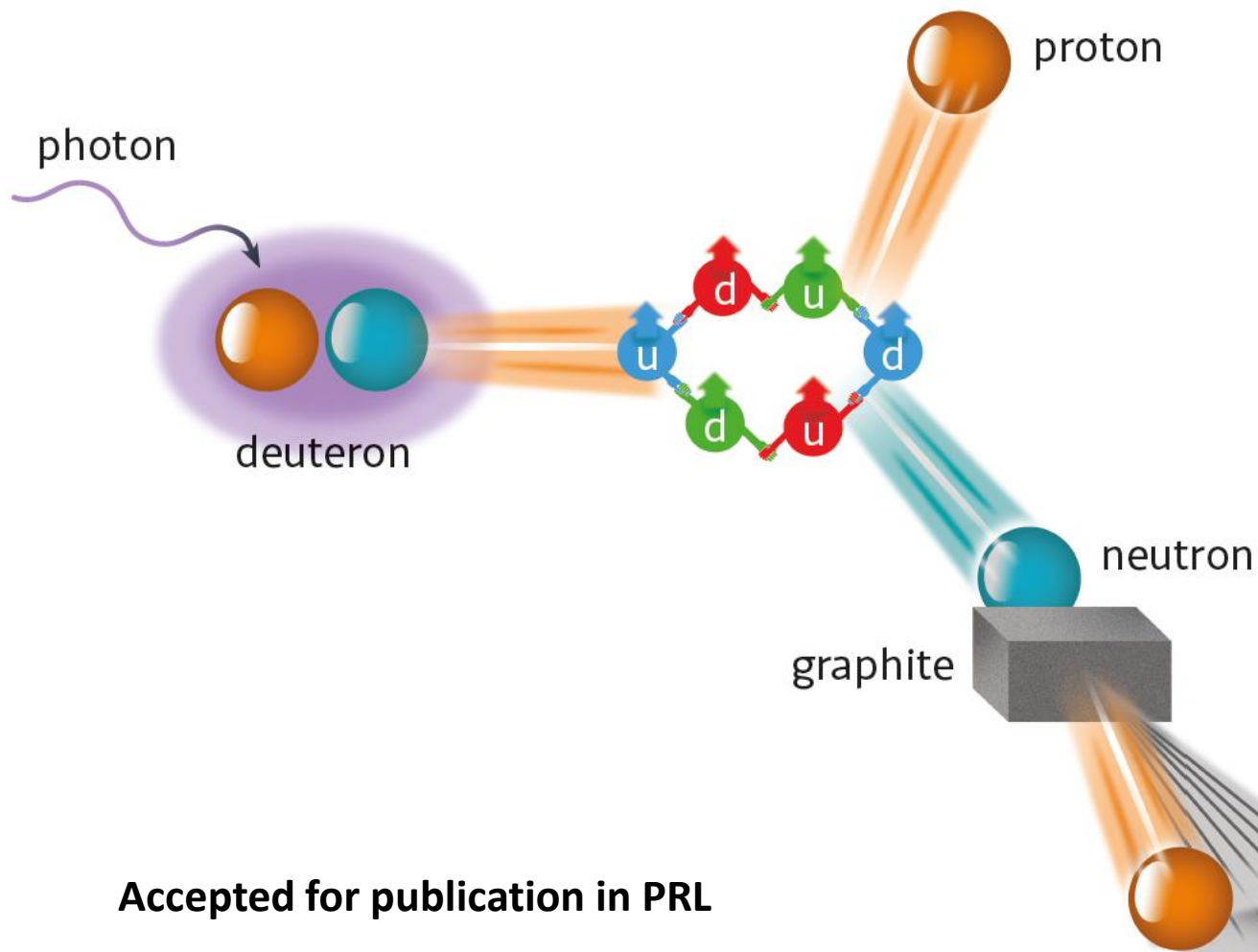
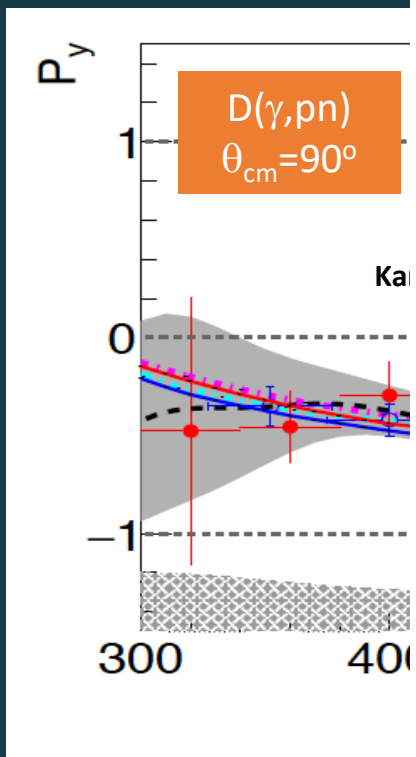
Deuterium photodisintegration (P_γ)

- 1st measurement of **final state neutron polarisation**
 -> **Both** p and n highly polarised in region of d^*
- θ_{CM} dependence – consistent with spin 3

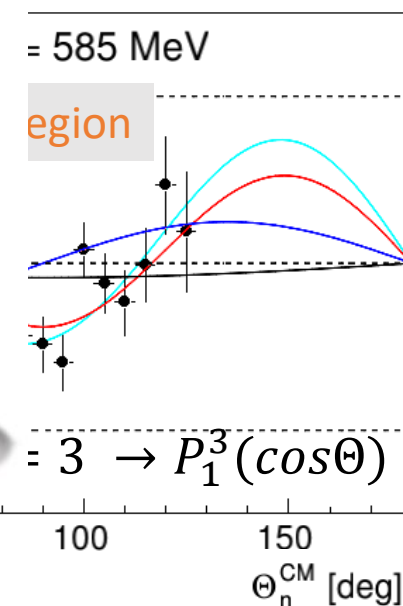
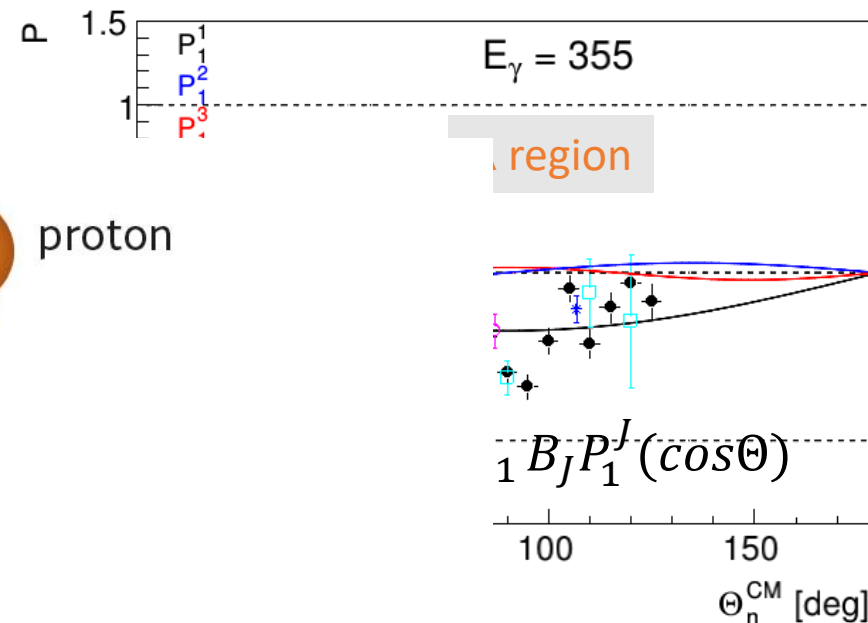


Deuterium photodisintegration (P_γ)

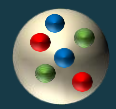
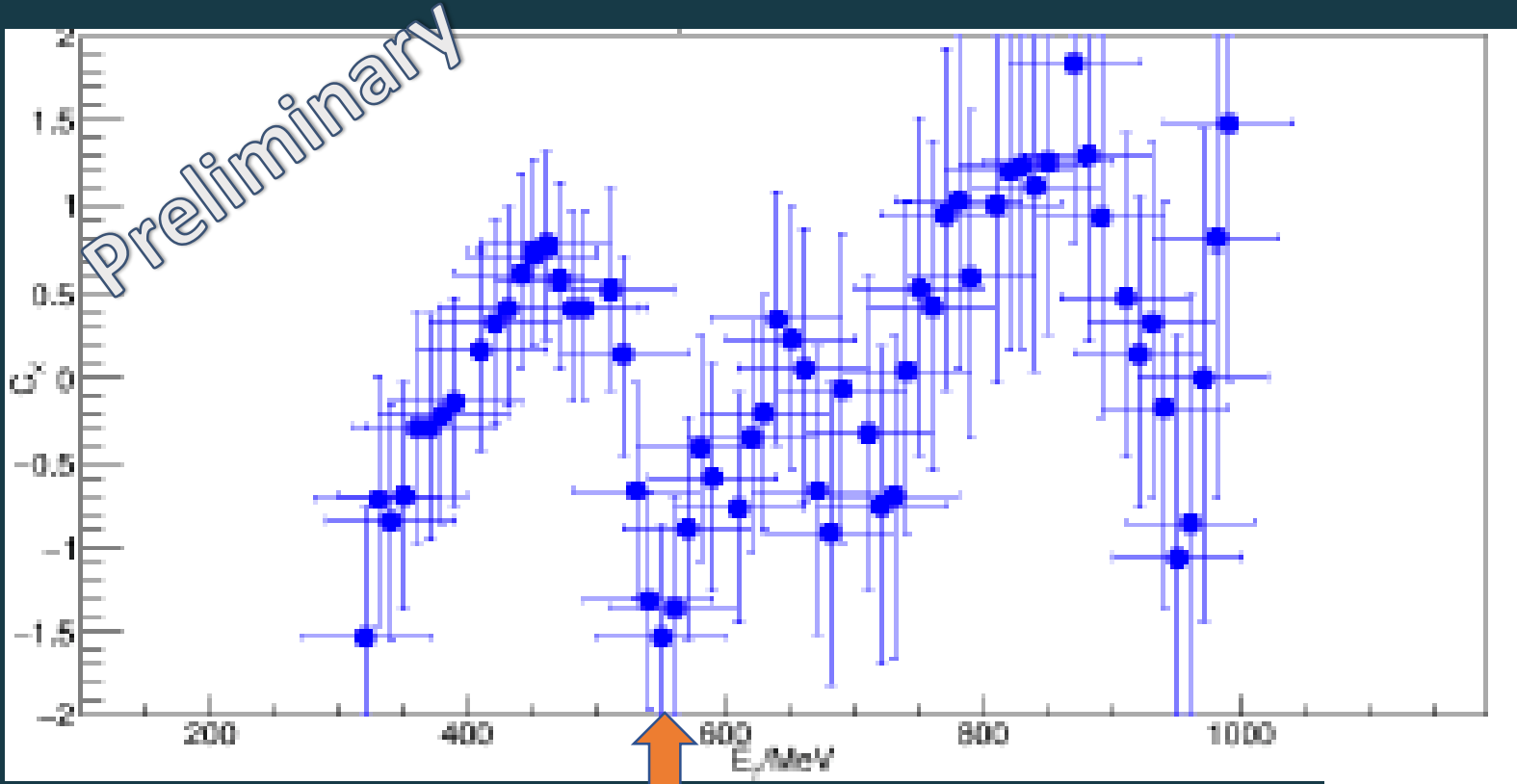
- 1st measurement of P_γ dependence of $D(\gamma, pn)$
- > **Both** p and n
- θ_{CM} dependence



Accepted for publication in PRL



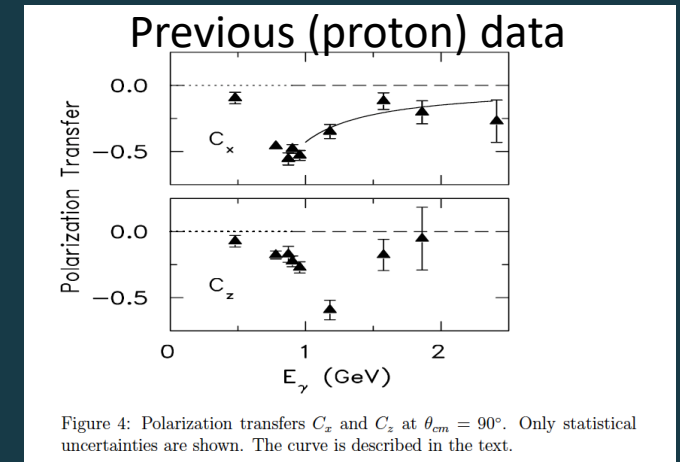
C_x^* for neutron final state



Combination of C_x^* and P

-> real and imaginary parts of same amplitude combination

- Indications of strong polarisation *transfer* in region of d^*
- Caveat: approx. (n,p) analyzing power used in analysis

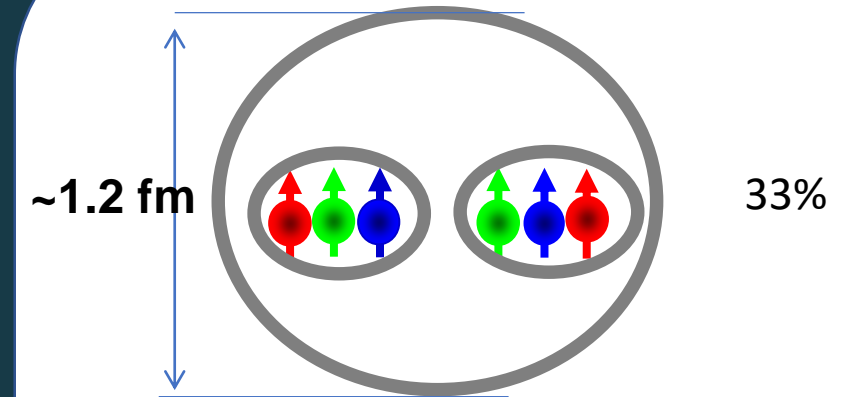
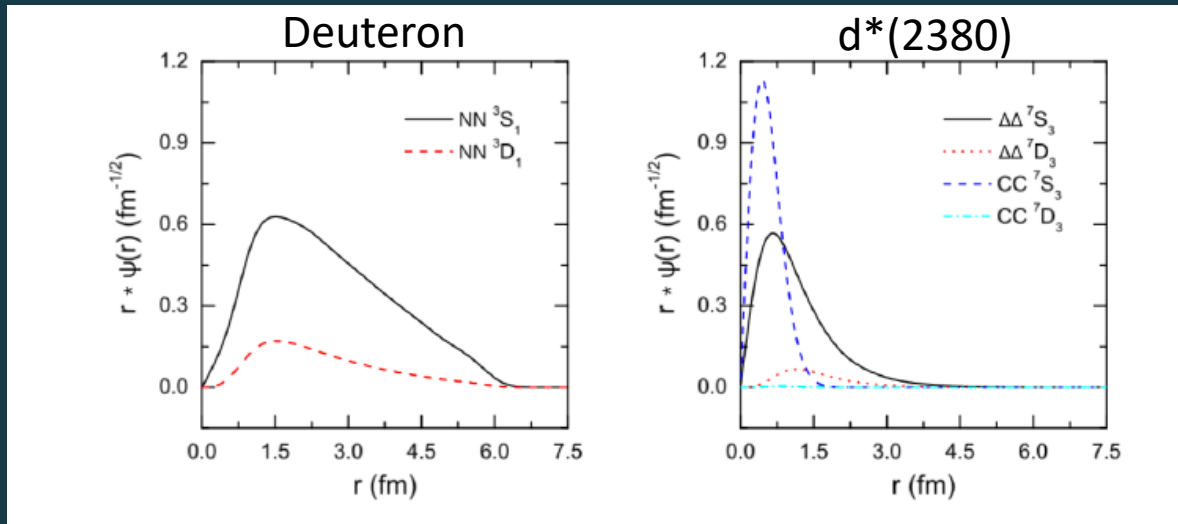


Observable	Structure function	Helicity amplitude combination
p_y	$R_T(y)$	$2\text{Im} \sum_{i=1}^3 [F_{i+}^* F_{(i+3)-} + F_{i-} F_{(i+3)+}^*]$
T	$R_T(\text{Im}T_{11})$	$2\text{Im} \sum_{i=1}^2 \sum_{j=0}^1 [F_{(i+3j)+} + F_{(i+3j+1)+}^* + F_{(i+3j)-} - F_{(i+3j+1)-}^*]$
Σ	R_{TT}	$2\text{Re} \sum_{i=1}^3 (-)^i [-F_{i+} F_{(4-i)-}^* + F_{(3+i)+} + F_{(7-i)-}^*]$
T_1	$R_{TT}(y)$	$2\text{Im} \sum_{i=1}^3 (-)^i [-F_{i+} F_{(7-i)+}^* + F_{i-} F_{(7-i)-}^*]$
$C_{x'}$	$R_T(x')$	$2\text{Re} \sum_{i=1}^3 [F_{i+}^* F_{(i+3)-} + F_{i-} F_{(i+3)+}^*]$
$C_{z'}$	$R_T(z')$	$\sum_{i=1}^6 \{ F_{i+} ^2 - F_{i-} ^2\}$
$O_{x'}$	$R_{TT}(x')$	$2\text{Im} \sum_{i=1}^3 (-)^{i+1} [F_{i+} F_{(7-i)+}^* + F_{i-} F_{(7-i)-}^*]$
$O_{z'}$	$R_{TT}(z')$	$2\text{Im} \sum_{i=1}^3 (-)^{i+1} [F_{i+} F_{(4-i)-}^* + F_{(3+i)+} + F_{(7-i)-}^*]$

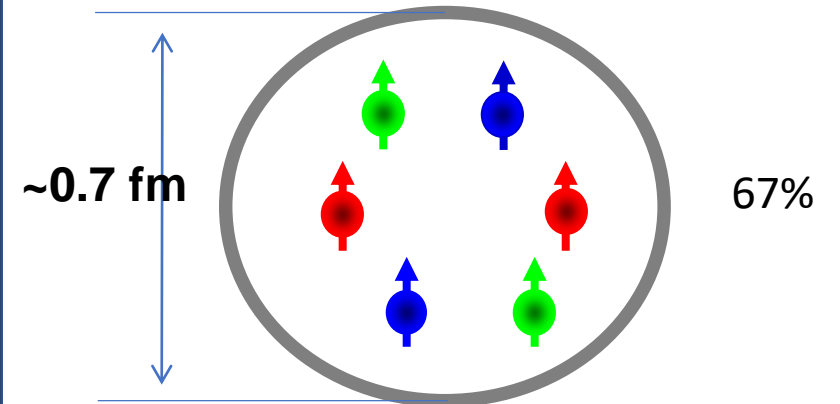
The $d^*(2380)$ in models

Any quark model with confinement and one gluon exchange *inevitably* predicts a 6-quark object with $(I)J^P=(0)3^+$

T Goldman et. al. Phys. Rev. C 39, 1889 (1989)



33%



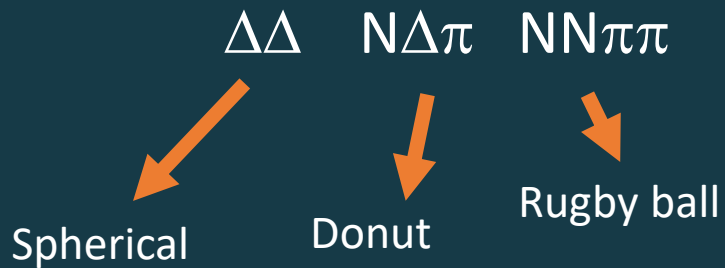
67%

Recent microscopic chiral quark models
 $\Delta\Delta$ + **hidden colour**

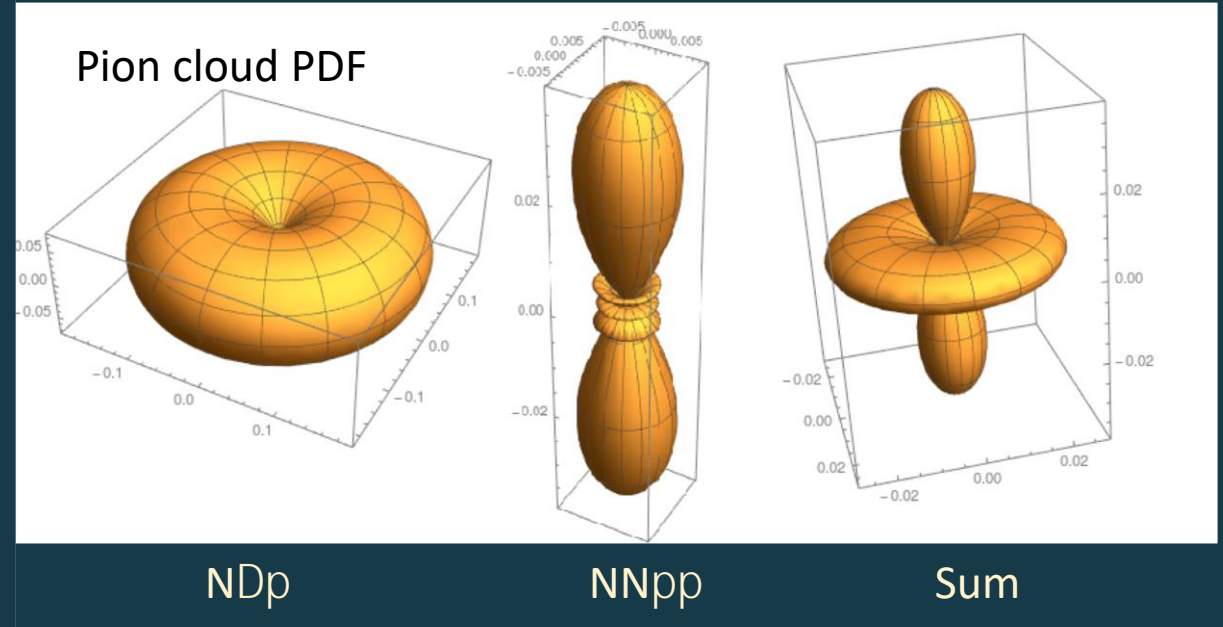
F. Huang et al, Chin.Phys. C39 (2015) 7, 071001

d^* in pion cloud model

- EM properties determined in simple overlapping Δ pion cloud model:



- EQM \rightarrow $\Delta\Delta$ no contribution
 $N\Delta\pi$, $NN\pi\pi$ \sim equal contribution but opposite sign \rightarrow cancellation
- MOM \rightarrow $\Delta\Delta$ no contribution
 $N\Delta\pi$, $NN\pi\pi$ \sim equal contribution but **same** sign

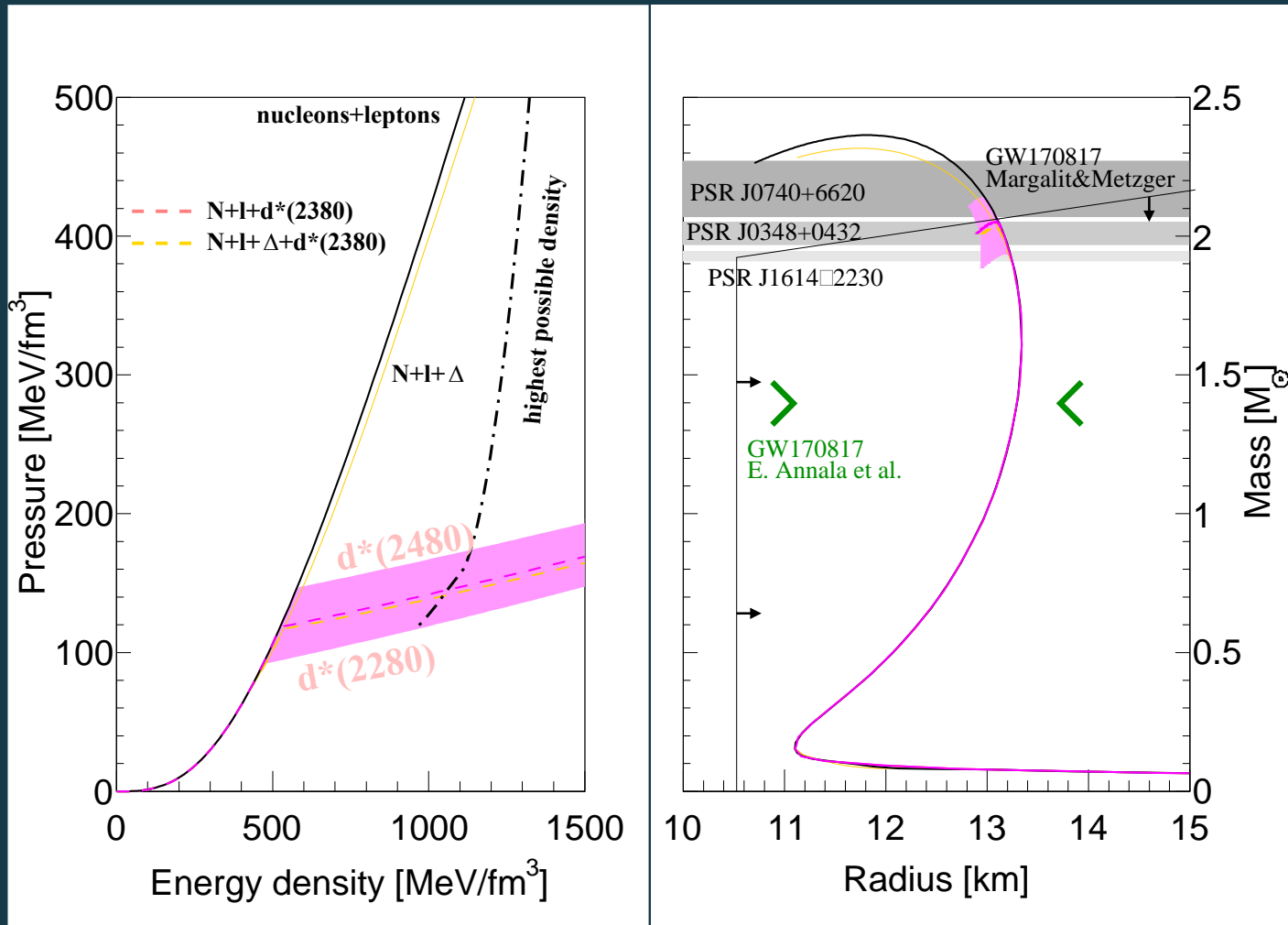


The $d^*(2380)$ in neutron stars

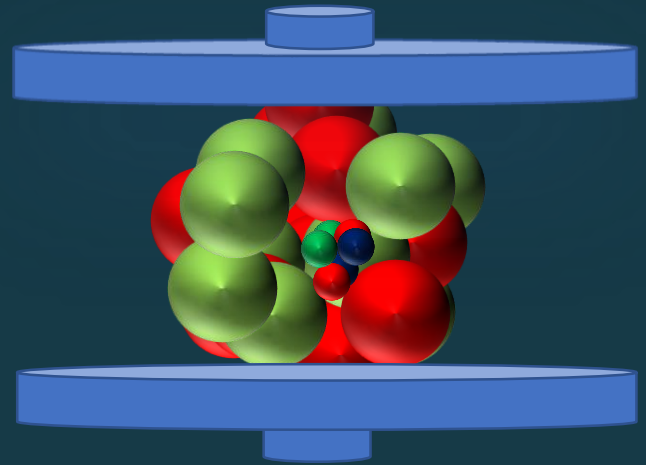
The $d^*(2380)$ in Neutron Stars – A New Degree of Freedom?

I. Vidaña^a, M. Bashkanov^{b,*}, D.P. Watts^b, A. Pastore^c

^a INFN Sezione di Catania, Dipartimento di Fisica, Università di Catania, Via Santa Sofia 64, 95123 Catania, Italy
^b School of Physics and Astronomy, University of Edinburgh, James Clerk Maxwell Building, Peter Guthrie Tait Road, Edinburgh EH9 3FD, UK
^c Department of Physics, University of York, Heslington York, YO10 5DD, UK



- d^* -> forms copiously above $2.5\rho_0$
 -> ~20% d^* at centre of heavy stars



- Star mass limit - around $2.1M_0$
- d^* in medium an important topic

More $d^*(2380)$ in neutron stars

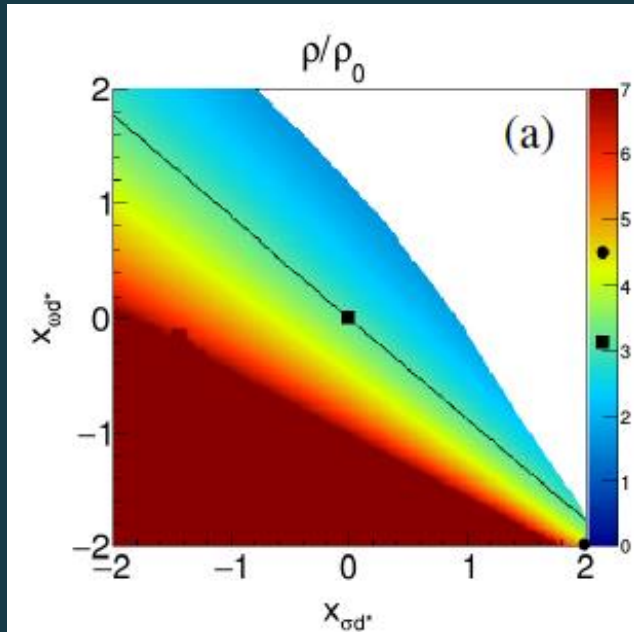
Neutron star matter equation of state including d^* -hexaquark degrees of freedom

A. Mantziris^{1,2}, A. Pastore¹, I. Vidaña³, D. P. Watts¹, M. Bashkanov¹, and A. M. Romero¹

¹ Department of Physics, University of York, Heslington, York, YO10 5DD, United Kingdom

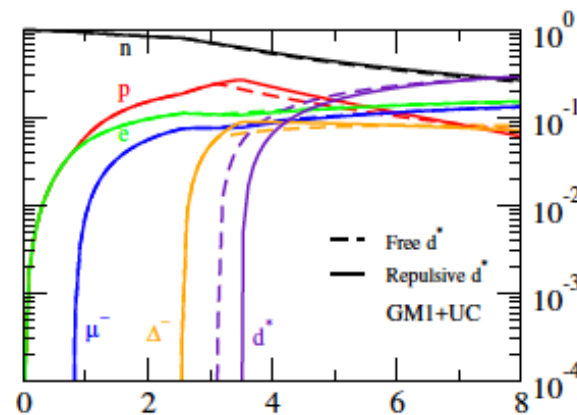
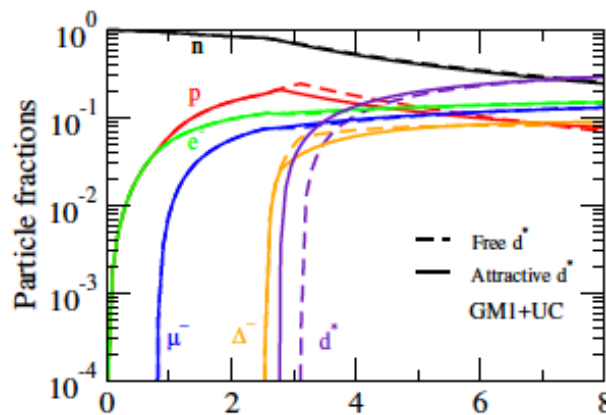
² Department of Physics, Imperial College London, London SW7 2AZ, United Kingdom

³ INFN Sezione di Catania, Dipartimento di Fisica "Ettore Majorana", Università di Catania, Via Santa Sofia 64, I-95123 Catania, Italy



- Non-linear Walecka model to describe d^* - N interaction via σ, ω
 -> vary the couplings

d^* in medium (nuclei) crucial -> Mihai's thesis

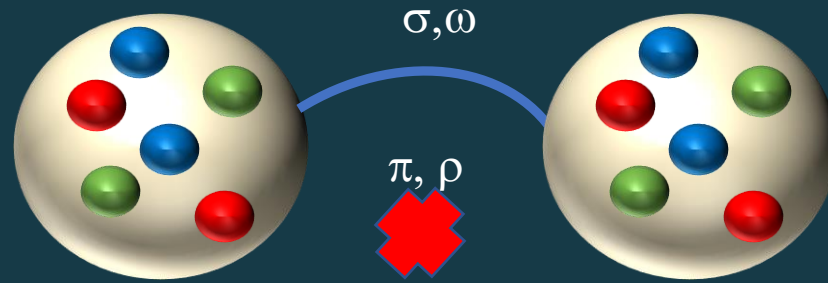


A&A, under review
 arxiv:2002.06571

The $d^*(2380)$ hexaquark and dark matter

Many body $d^*(2380)$ systems

- d^* hexaquark - a new compact, **bosonic** and **isoscalar** form of light quark matter

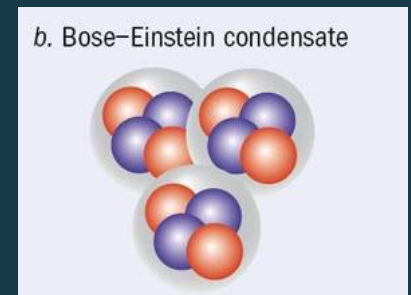


- Its a boson - can the d^* form a Bose-Einstein condensate ?

The Hoyle state – a condensate of alpha

- The Hoyle state of ^{12}C (7.65 MeV) - crucial in nucleosynthesis
- α is also an isoscalar, bosonic object
-> analogous to the d^* (but much larger)
- The Hoyle state properties reproduced remarkably well assuming it is a dilute Bose-Einstein **condensate** of 3 alpha particles.

70% condensed state with 30% of time scattered out from interactions



Alpha-Particle Condensate Structure of the Hoyle State: where do we stand?

To cite this article: P Schuck et al 2017 J. Phys.: Conf. Ser. 863 012005

The d^* condensate ?

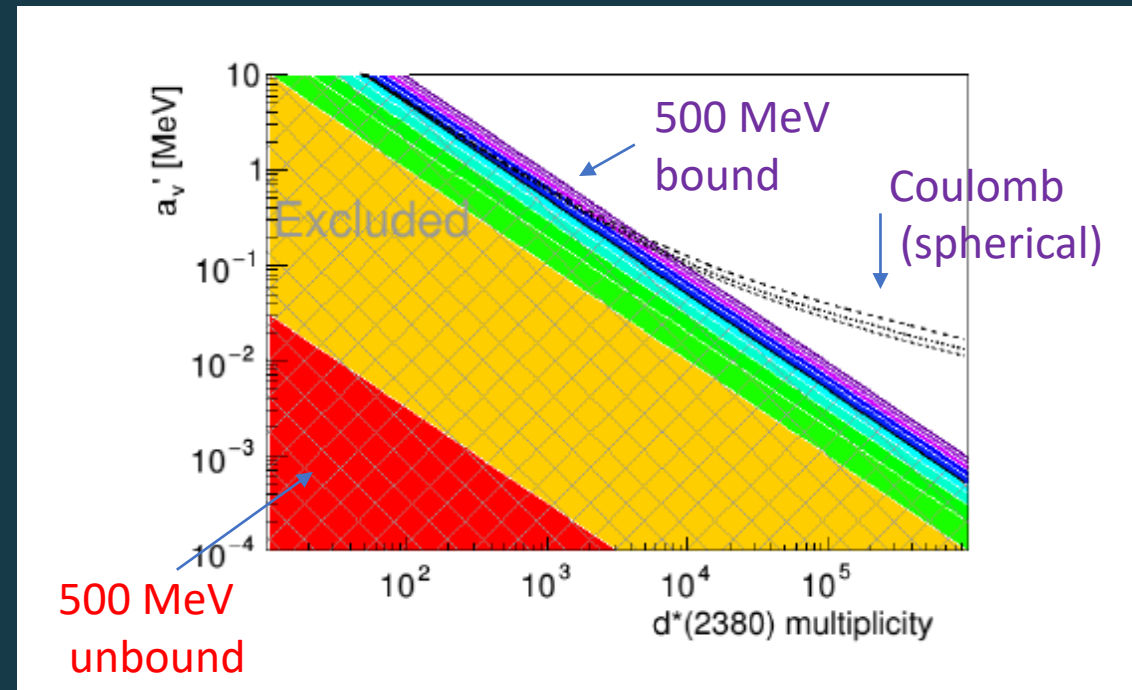
- The Binding Energy of a condensate can be calculated as a function of d^* multiplicity
- Account for different scaling of binding and Coulomb terms (common wavefunction for all particles).
Surface term not relevant for condensate – linear chain configuration
-> Vary volume coefficient (reflecting the d^* - d^* interaction potential) over physical ranges

Standard nuclei

$$B_{LD}/A = a_v - a_s \cdot \frac{1}{A^{1/3}} - a_c \cdot \frac{Z(Z-1)}{A^{4/3}}$$

d^* condensate

$$B/D = a'_v \cdot (D-1) - a'_c \cdot \frac{(D-1)}{D}; \quad a'_c [MeV] = 0.064 \cdot \frac{\rho}{\rho_0}$$



A new possibility for light-quark Dark matter

- The Independent, Daily Mail, Daily Express, Newsweek... 40+ newspapers and journals all over the globe.
- 12k downloads a day, 15k+ downloads total.

The screenshot shows a news aggregator page for the article "A new possibility for light-quark Dark Matter" published in the Journal of Physics: G Nuclear & Particle Physics in January 2020. The page features a grid of news snippets from various international media outlets, including CNN, BBC, Reuters, and many others. Each snippet includes a small logo, a headline, and a brief summary of the article's content. The headlines are in multiple languages, including English, Spanish, French, and Russian. The page also includes a sidebar on the left with a "422" icon and a "SHARE THIS ARTICLE" section with social media icons for Facebook, Twitter, and LinkedIn. The top of the page displays the IOP Publishing logo and the article title.

The proposed programme with polarised deuteron targets

Possible beam-target $D(\gamma, pn)$ observables

- Longitudinal target polarisation

$$T_{10}^c(-E) \frac{d\sigma_0}{d\Omega_p} = \frac{1}{2} \sum_{sm_s} (|t_{sm_s 11}|^2 - |t_{sm_s 1-1}|^2)$$

Circular γ -polarisation

$$T_{10}^l(G) \frac{d\sigma_0}{d\Omega_p} = \Im m \sum_{sm_s} (t_{sm_s 11}^* t_{sm_s -11})$$

Linear γ -polarisation

- Transverse target polarisation

$$T_{11}^0(-T) \frac{d\sigma_0}{d\Omega_p} = \Im m \sum_{sm_s} (t_{sm_s 1-1}^* t_{sm_s 10} + t_{sm_s 10}^* t_{sm_s 11})$$

Unpolarised

$$T_{11}^c(-F) \frac{d\sigma_0}{d\Omega_p} = -\Re e \sum_{cm} (t_{sm_s 1-1}^* t_{sm_s 10} + t_{sm_s 10}^* t_{sm_s 11})$$

Circular γ -polarisation

$$T_{11}^l \frac{d\sigma_0}{d\Omega_p} = \Im m \sum_{sm_s} (t_{sm_s 1-1}^* t_{sm_s -10})$$

Linear γ -polarisation

$$T_{1-1}^l \frac{d\sigma_0}{d\Omega_p} = -\Im m \sum_{sm_s} (t_{sm_s 11}^* t_{sm_s -10})$$

Linear γ -polarisation

**Photodisintegration of polarized deuterons –
measurement of angular distributions at $E_\gamma = 450, 550$ and 650 MeV**

K.H. Althoff, G. Anton, B. Bock¹, D. Bour, R. Dostert², P. Erbs, T. Jahnen, O. Kaul³, E. Kohlgarth, B. Lücking⁴, D. Menze, W. Meyer, E.P. Schilling, W.J. Schwillie, D. Sundermann⁵, W. Thiel, D. Thiesmeyer⁶
Physikalisches Institut der Universität Bonn, Nussallee 12, D-5300 Bonn 1, Federal Republic of Germany

Received 14 February 1989

$$T = \frac{3}{2} \frac{\sigma_+ - \sigma_-}{\sigma_+ + \sigma_0 + \sigma_-}$$

$$Z = (-2) \frac{\sigma_+ + \sigma_- - 2\sigma_0}{\sigma_+ + \sigma_0 + \sigma_-}$$

Significant improvement
In data quality possible
at MAMI

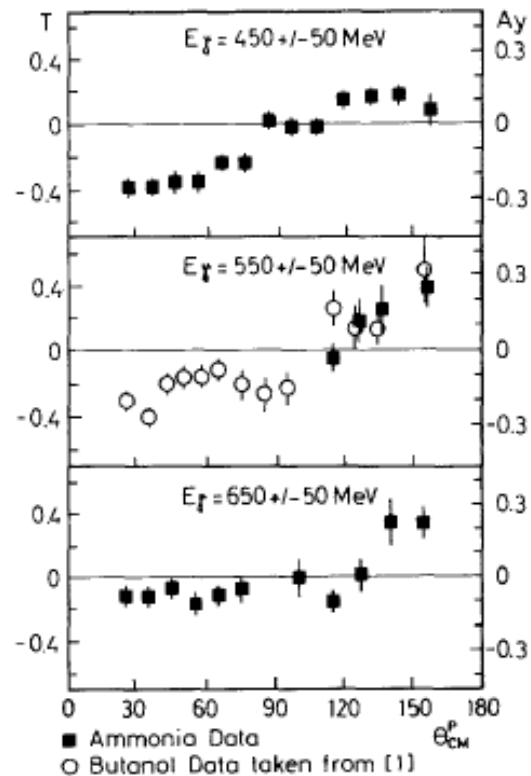


Fig. 4. Target asymmetry vs. proton angle θ_{CM}^p

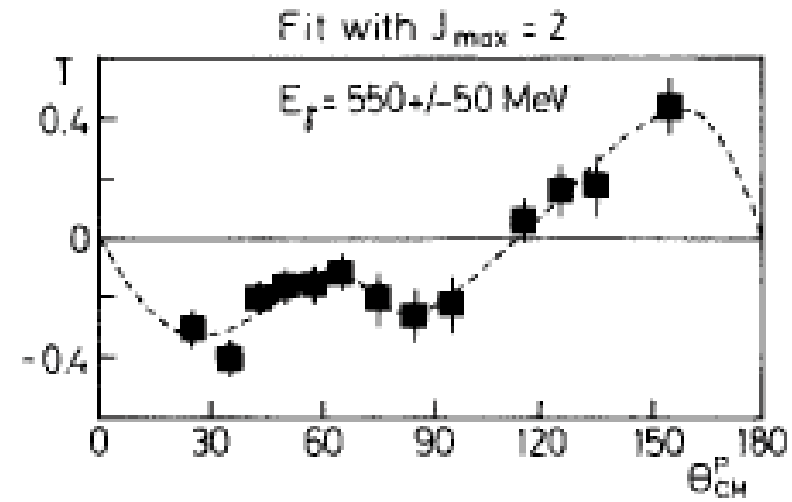
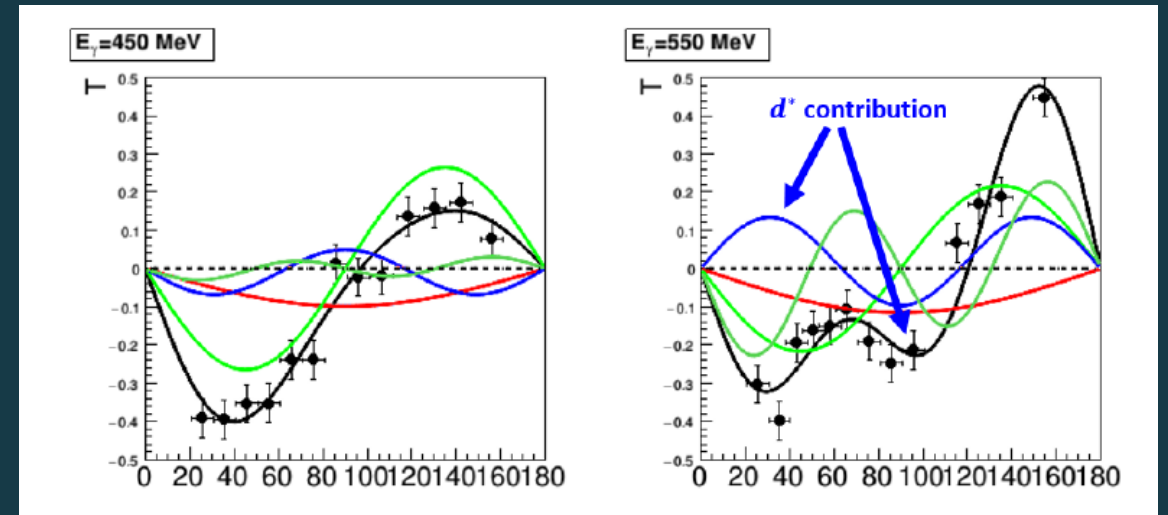


Fig. 6. Angular distribution at 550 MeV fitted with angular momentum $J_{max} = 2$

Taken as contribution of higher multipoles!



Expected sensitivity

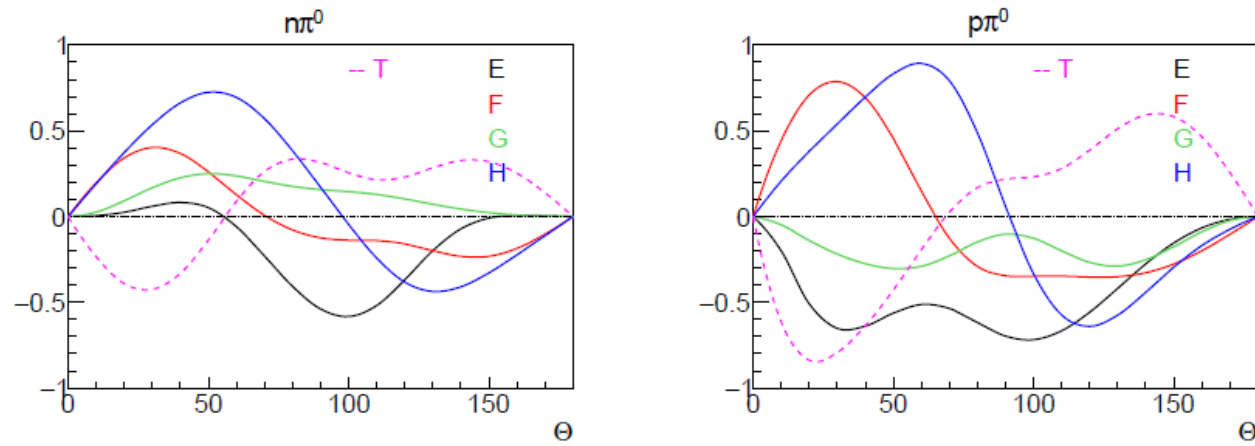
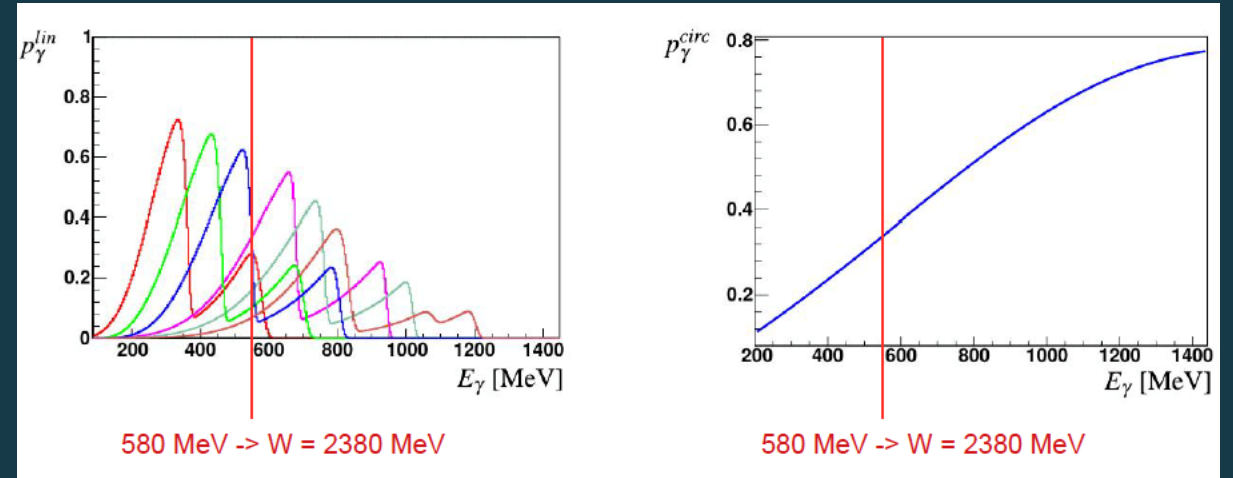


Figure 6: MAID predictions for photoexcitation of the $N^*(1680)$ (along with background from Born terms only). The predictions of various single- and double-polarisation observables for the $n\pi^0$ reaction (left) and $p\pi^0$ reaction (right) are shown. Photon energy = 1035 MeV ($W=1680$ MeV).

Mainz MAID: $N^*(1680) \frac{5^+}{2}$ - same multipoles (E2,M3) as d^*

Count rate estimate

- Measure all 4 transverse target asymmetries



- Stat accuracy $\delta A \sim 0.01$ for ± 15 MeV bin in d^* region, $P_T \sim 0.7$, $P_\gamma \sim 0.4$

$$N_{tot} = \frac{1}{P_T^2 \cdot P_\gamma^2 \cdot \delta A^2}$$

-> requires $\sim 100k$ events bin^{-1} (good angular binning required)

$$\begin{aligned} \text{Event rate} &= \text{Beam intensity} \times \text{tag eff} \times \text{Number nuclei} \times \text{efficiency} \times \text{cross section} \\ &= 1 \text{ s}^{-1} \end{aligned}$$

Beam time ~ 7 weeks

Required beamtime

Table 3: Expected beamtime length

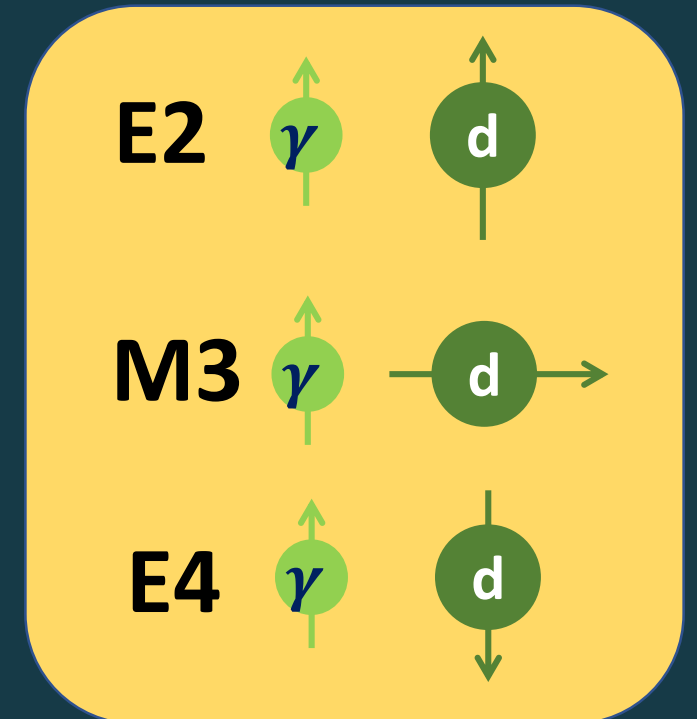
Target	Radiator	Beam polarisation	Target polarisation	days
D-butanol	Diamond	\parallel	\uparrow	7
D-butanol	Diamond	\parallel	\downarrow	7
D-butanol	Diamond	\perp	\uparrow	7
D-butanol	Diamond	\perp	\downarrow	7
D-butanol	Moeller	\odot/\otimes	\uparrow	3
D-butanol	Moeller	\odot/\otimes	\downarrow	3
Carbon	Diamond	\parallel	—	3
Carbon	Diamond	\perp	—	3
				6 weeks

+ 1 week to setup/maintain the target

What will we learn?

- Clean separation of the $d^*(2380)$ from background channels
 - Precise knowledge of the d^* photo coupling
- d^* can be excited in E2, M3 or E4 transitions
 - E2 \leftrightarrow Electric Quadrupole moment
 - M3 \leftrightarrow Magnetic Octupole moment
 - E4 \leftrightarrow Electric Hexadecapole moment

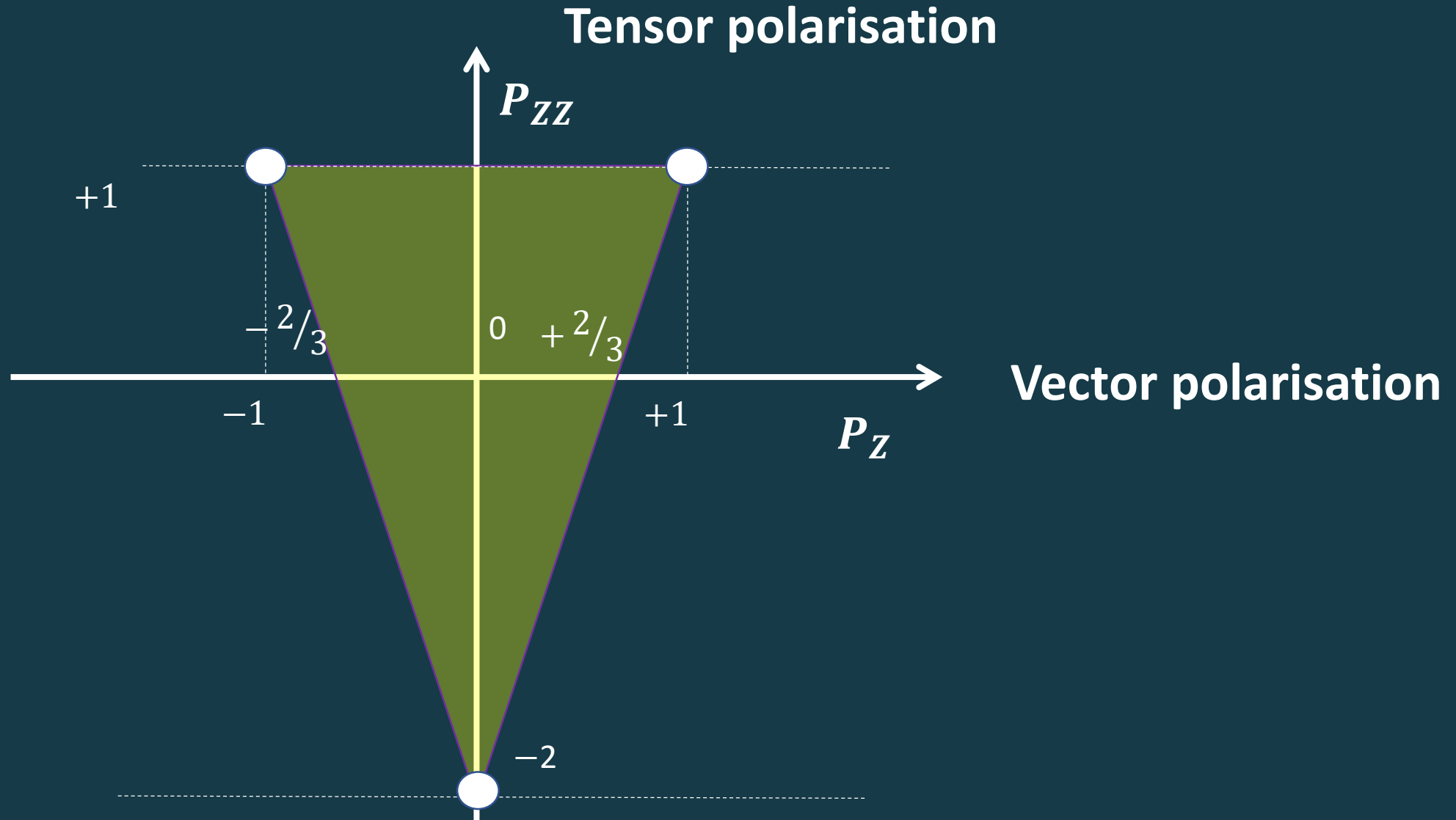
Size and structure of the d^*



Summary

- MAMI can lead the way in providing key data to elucidate the $d^*(2380)$ with EM probes
- Important ramifications for hadron physics - potentially for astrophysics
- Campaign with polarized deuteron target would provide T and a trio of beam-target observables

Deuteron polarisation



Extracted (approximate!) d^* cross section from σ and Σ

Coefficients: assuming M3 dominance !

$$a_2 = 0.375|M3(^3D_3)|^2 + 0.108|M3(^3D_3)||M3(^3G_3)| \\ \cdot \cos(\delta_{3D_3} - \delta_{3G_3}) + 0.391|M3(^3G_3)|^2$$

$$a_4 = 0.014|M3(^3D_3)|^2 + 0.021|M3(^3D_3)||M3(^3G_3)| \\ \cdot \cos(\delta_{3D_3} - \delta_{3G_3}) + 0.017|M3(^3G_3)|^2$$

$$a_6 = 0.172|M3(^3D_3)||M3(^3G_3)| \\ \cdot \cos(\delta_{3D_3} - \delta_{3G_3}) + 0.025|M3(^3G_3)|^2$$

-> Use extracted even coefficients to estimate d^* contribution

-> 4.6 +-1%

Similar order as suggested by structure in $\gamma + D \rightarrow d\pi^0\pi^0$

(d^* coefficients only contribute to odd coefficients via Interference with other partial waves -> therefore assumed small)

Primordial production at the quark hadron phase transition

- DM/matter ratio calculated in general terms assuming Boltzmann distribution in QGP

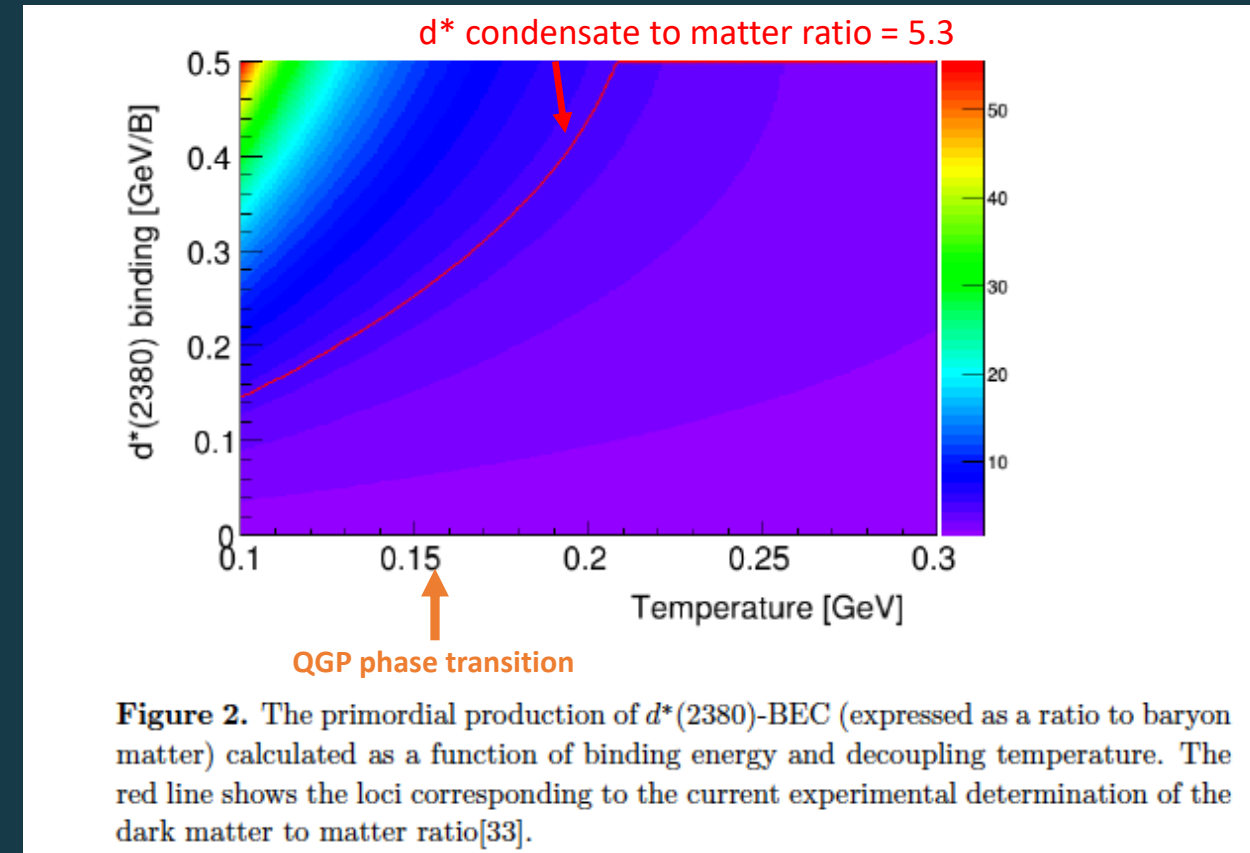
arxiv.org/abs/1805.03723

$$\Omega_{DM}/\Omega_{baryon} = \frac{7}{4} M_N \left(1 - \frac{y_D}{M_N}\right)^{3/2} e^{\frac{M_N y_D}{T}}$$

y_D - d^* binding energy (units of nucleon mass)
T - temperature

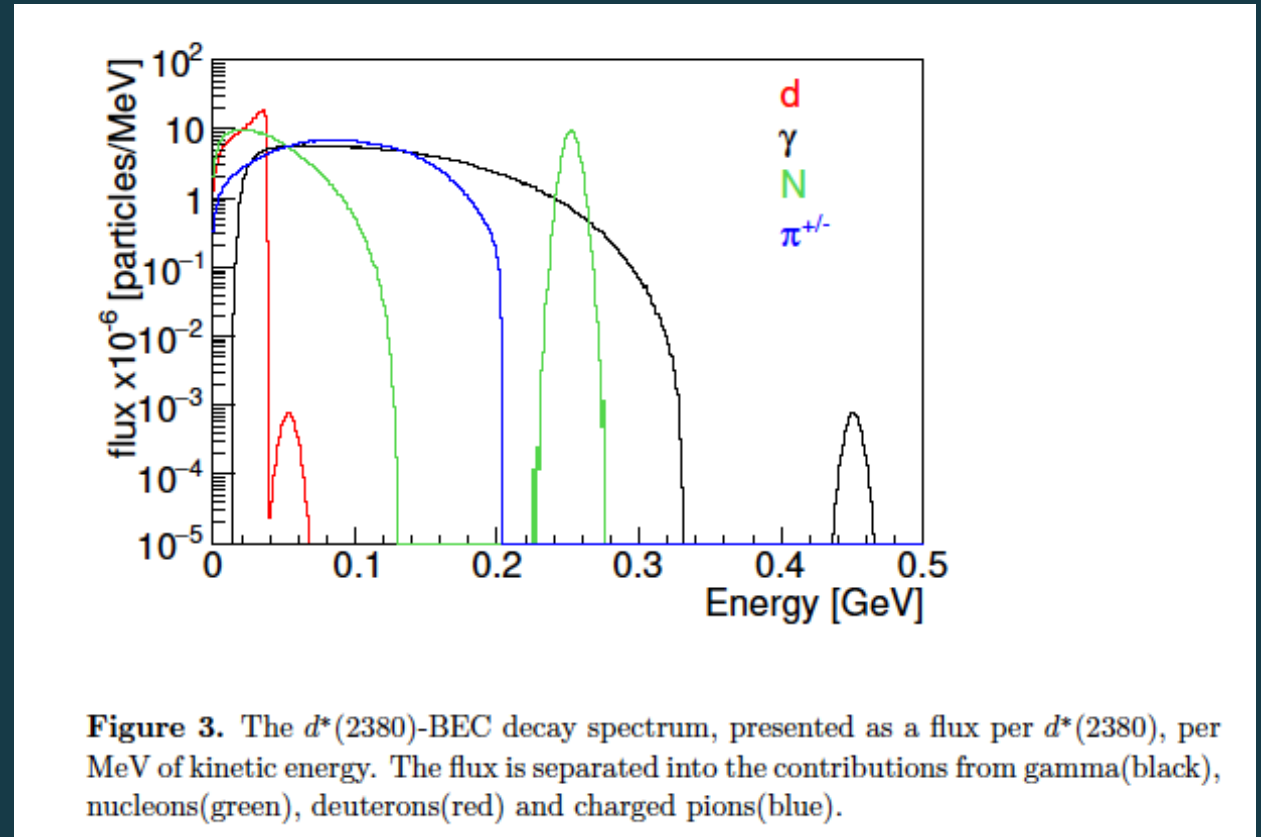
At QGP phase transition (0.16 GeV) a BE~300 MeV would produce measured DM ratio.

Achievable with d^* condensate multiplicity $>\sim 1000$



Signatures of d^* condensates

- The (charged) chain condensates would attract electrons to become neutral objects (DFT calculations for linear chain configuration in progress)
- Breakup of a condensate “nucleus” would produce a characteristic decay spectrum



Measurements of the P_x component of neutron polarisation in the reaction $\gamma d \rightarrow pn$ by linearly polarised photons in the energy range 0.3-0.5 GeV

To cite this article: F V Adamian *et al* 1988 *J. Phys. G: Nucl. Phys.* 14 831

- The target asymmetry T :

$$T = \frac{3}{2} \cdot \frac{\sigma_+ - \sigma_-}{\sigma_+ + \sigma_0 + \sigma_-} \quad \text{with} \quad -3/2 \leq T \leq 3/2. \quad (1)$$

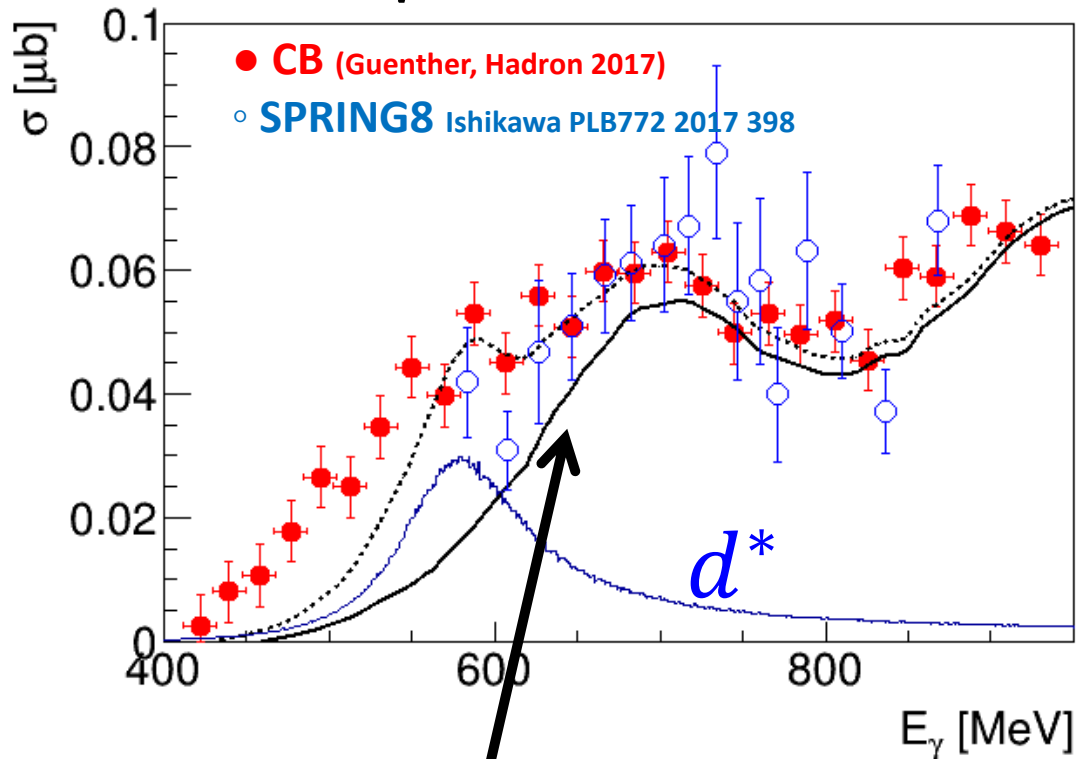
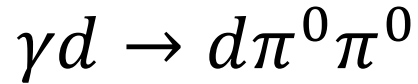
- The tensor asymmetry Z :

$$Z = (-2) \cdot \frac{\sigma_+ + \sigma_- - 2\sigma_0}{\sigma_+ + \sigma_0 + \sigma_-} \quad \text{with} \quad -2 \leq Z \leq 4 \quad (2)$$

σ_+ , σ_0 and σ_- are the differential cross sections of the process with pure deuteron states. σ_+ for example is the cross section with all deuterons parallel to the polarization axis, in our case the y -axis. The latter is given by the $\mathbf{k} \times \mathbf{q}$ -direction (\mathbf{k} photon- and \mathbf{q} pro-

E_γ (GeV)

Photoproduction of $d\pi\pi$ final state



- Photoproduction kinematics challenging \rightarrow deuteron only detected at forward angles
- Active deuteron target prototype under construction

Conventional Background

M. Egorov, A. Fix, Nucl.Phys. A933 (2015) 104-113

See talk of B Krusche



Polarization observables - disintegration

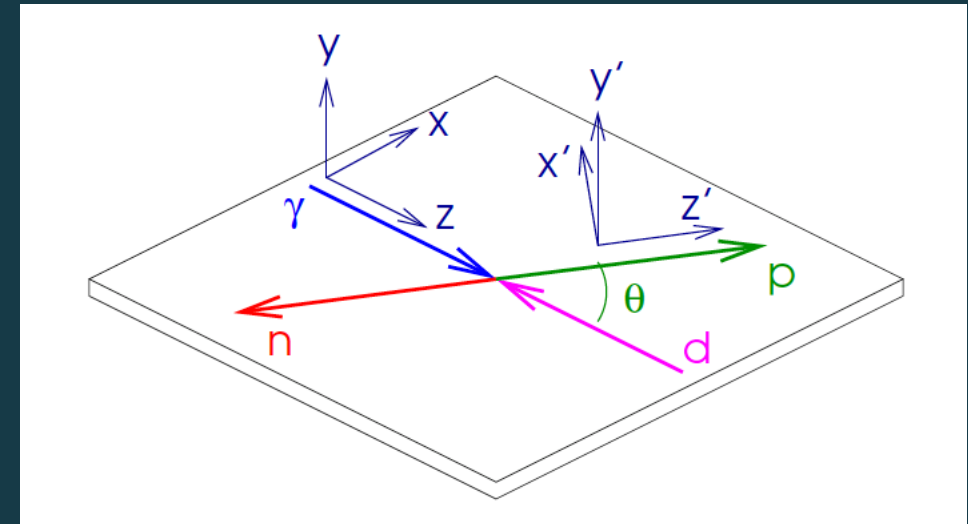
12 amplitudes – 23 independent observables

For the differential cross section in the cm frame one obtains

$$\begin{aligned} \frac{d\sigma}{d\Omega} = & \frac{d\sigma_0}{d\Omega} \left[1 + P_l^\gamma \Sigma^l(\Theta) \cos 2\Phi \right. \\ & + \sum_{I=1,2} P_I^\alpha \left\{ \sum_{M \geq 0} (T_{IM}(\Theta) \cos(M(\Phi_d - \Phi) - \delta_{I1}\pi/2) \right. \\ & \quad \left. + P_z^\gamma T_{IM}^c(\Theta) \sin(M(\Phi_d - \Phi) + \delta_{I1}\pi/2)) d_{M0}^I(\Theta_d) \right. \\ & \quad \left. + P_l^\gamma \sum_{M=-I}^I T_{IM}^l(\Theta) \cos(\Psi_M - \delta_{I1}\pi/2) d_{M0}^I(\Theta_d) \right\} \Big], \end{aligned}$$

where

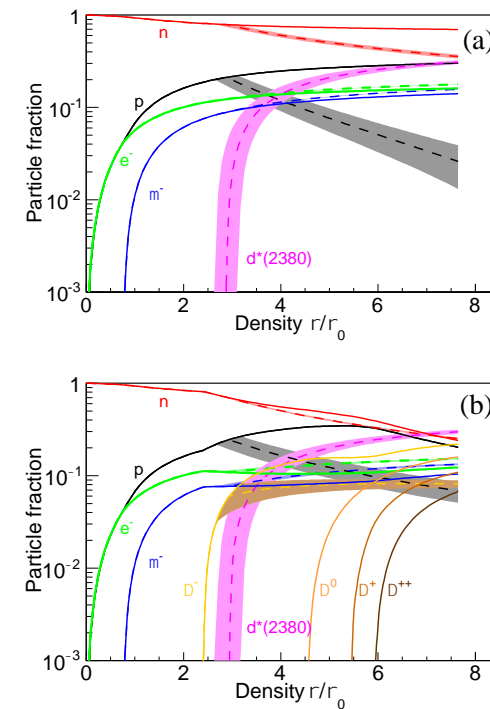
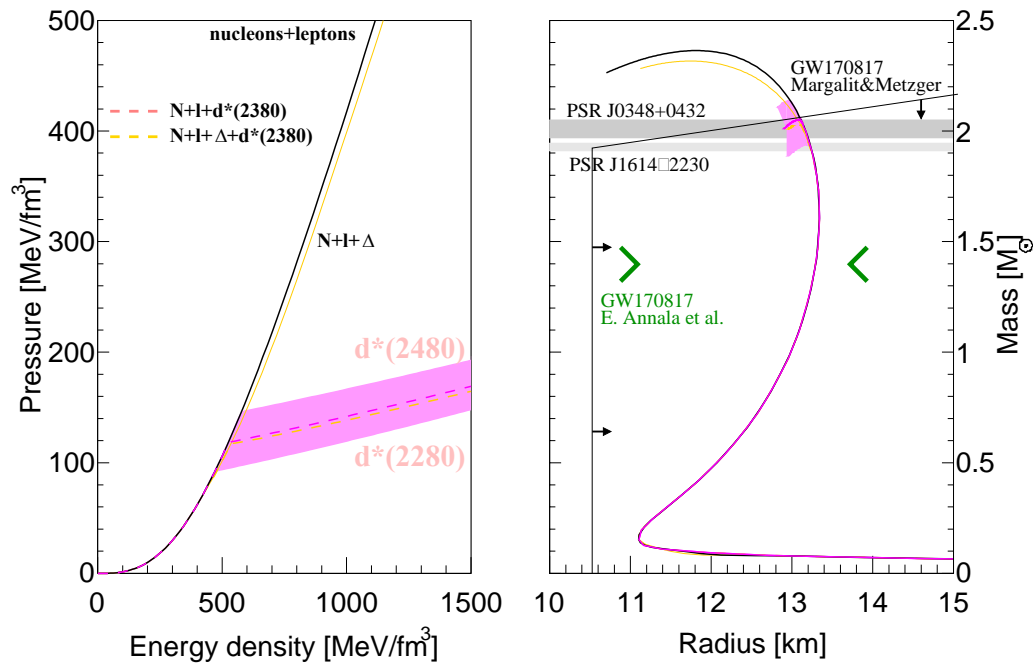
$$\Psi_M = M(\Phi_d - \Phi) + 2\Phi.$$



Observable	Structure function	Helicity amplitude combination
p_y	$R_T(y)$	$2\text{Im} \sum_{i=1}^3 [F_{i+}^* F_{(i+3)-} + F_{i-} F_{(i+3)+}^*]$
T	$R_T(\text{Im}T_{11})$	$2\text{Im} \sum_{i=1}^2 \sum_{j=0}^1 [F_{(i+3j)+} F_{(i+3j+1)+}^* + F_{(i+3j)-} F_{(i+3j+1)-}^*]$
Σ	R_{TT}	$2\text{Re} \sum_{i=1}^3 (-)^i [-F_{i+} F_{(4-i)-}^* + F_{(3+i)+} F_{(7-i)-}^*]$
T_1	$R_{TT}(y)$	$2\text{Im} \sum_{i=1}^3 (-)^i [-F_{i+} F_{(7-i)+}^* + F_{i-} F_{(7-i)-}^*]$
$C_{x'}$	$R_T(x')$	$2\text{Re} \sum_{i=1}^3 [F_{i+}^* F_{(i+3)-} + F_{i-} F_{(i+3)+}^*]$
$C_{z'}$	$R_T(z')$	$\sum_{i=1}^6 \{ F_{i+} ^2 - F_{i-} ^2\}$
$O_{x'}$	$R_{TT}(x')$	$2\text{Im} \sum_{i=1}^3 (-)^{i+1} [F_{i+} F_{(7-i)+}^* + F_{i-} F_{(7-i)-}^*]$
$O_{z'}$	$R_{TT}(z')$	$2\text{Im} \sum_{i=1}^3 (-)^{i+1} [F_{i+} F_{(4-i)-}^* + F_{(3+i)+} F_{(7-i)-}^*]$

The density frontier – heavy stars and mergers

- $d^*(2380)$ may have very significant effect on heavy neutron stars and the physics of mergers. (Vidal, Bashkanov, DPW, Pastore)



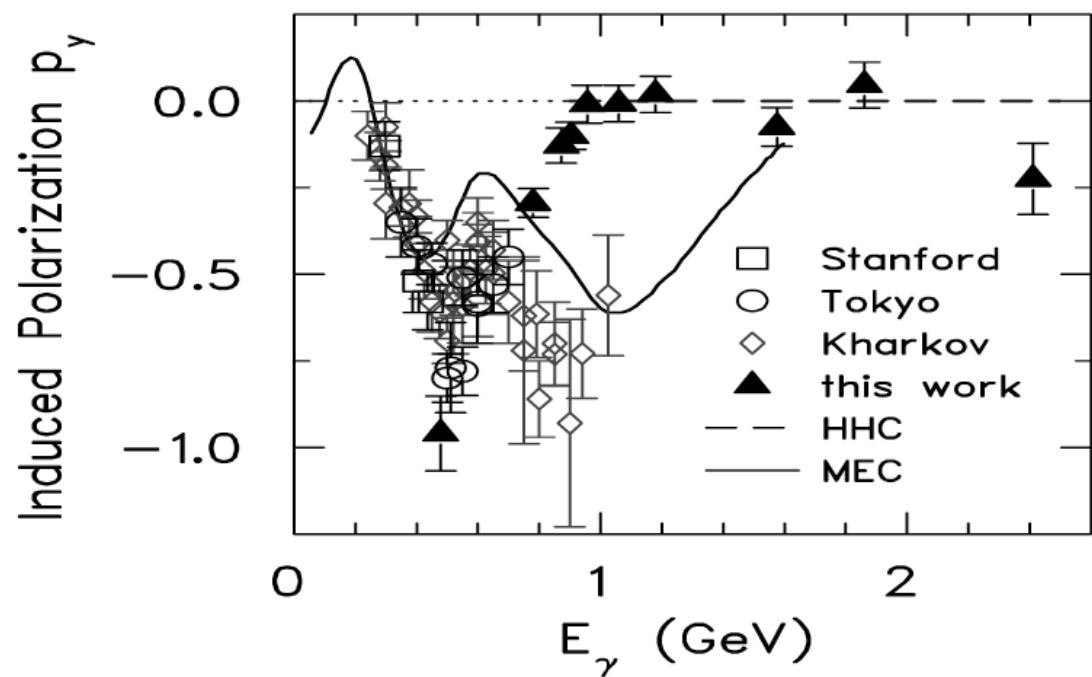


Figure 3: Induced polarization p_y in deuteron photodisintegration at $\theta_{cm} = 90^\circ$. Only statistical uncertainties are shown. The curves are described in the text.

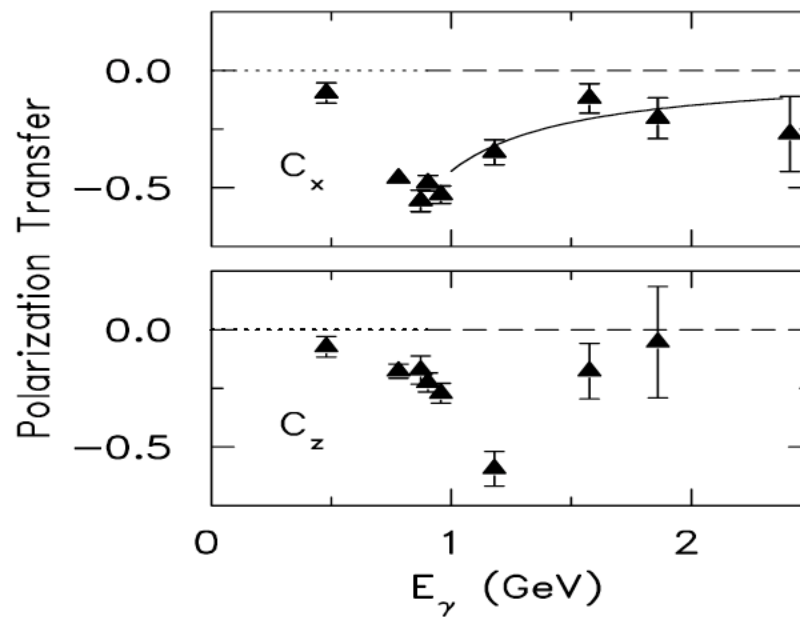


Figure 4: Polarization transfers C_x and C_z at $\theta_{cm} = 90^\circ$. Only statistical uncertainties are shown. The curve is described in the text.

$d^*(2380)$ multipole expansion

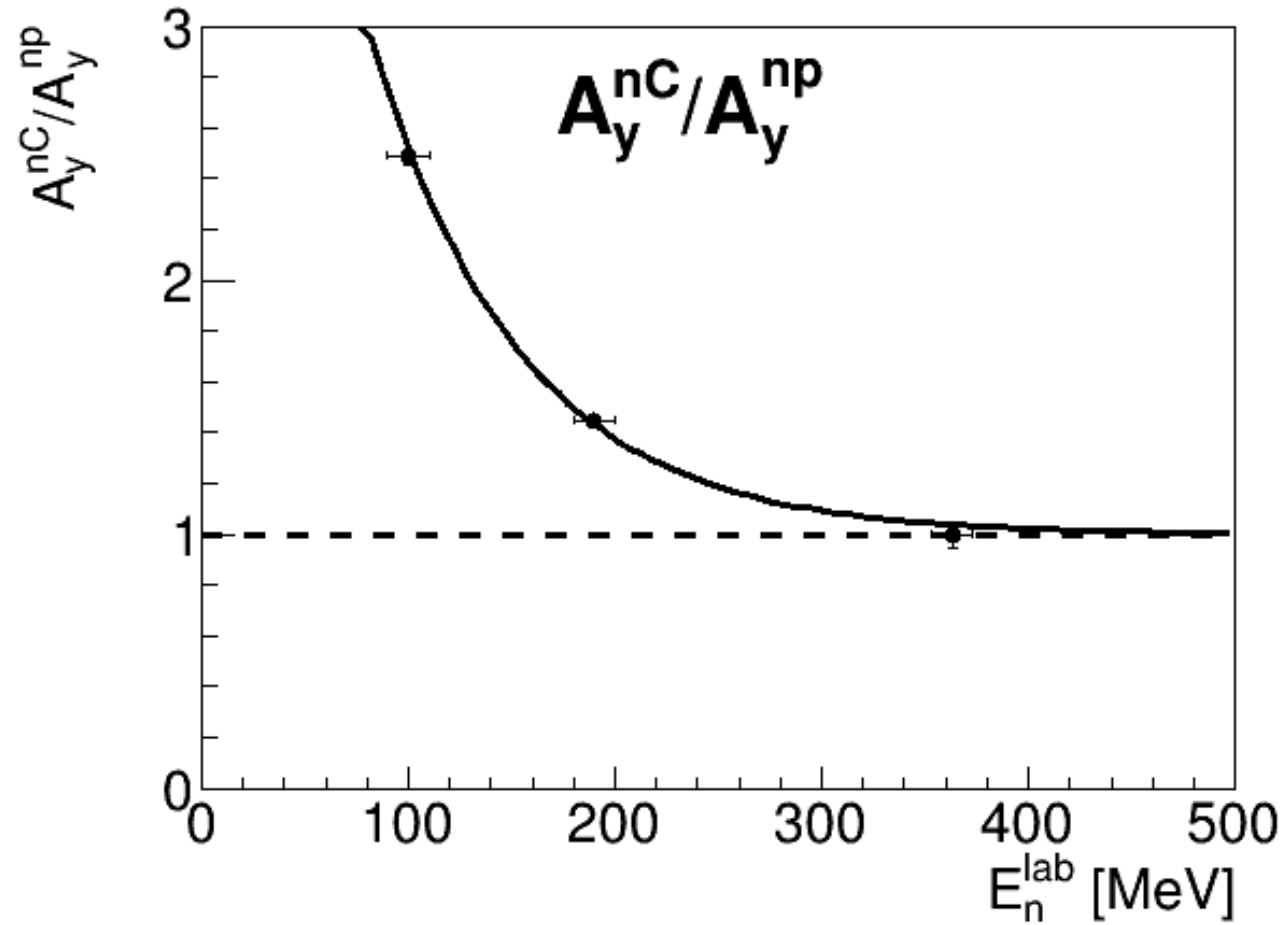
$$\frac{M3}{E2} \sim \frac{\langle \Omega \rangle_{d \rightarrow d^*}}{\langle Q \rangle_{d \rightarrow d^*}} \sim \frac{\langle Q \rangle_{d \rightarrow d^*} \mu_{d \rightarrow d^*}}{\langle Q \rangle_{d \rightarrow d^*}} \sim M_d \omega_\gamma^2 \mu_{d \rightarrow d^*}$$

If $\mu_{d \rightarrow d^*}$ is large the $d \rightarrow d^*$ M3 transition might be dominant

$$\mu_{d^*} \sim 7.6 \mu_N$$

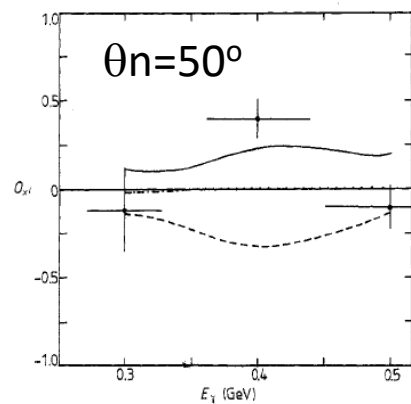
Yubing Dong, Pengnian Shen, Zongye Zhang Phys.Rev. D 97, (2018), no.11, 114002

Neutron A_y on Carbon



Measurements of the P_x component of neutron polarisation in the reaction $\gamma d \rightarrow pn$ by linearly polarised photons in the energy range 0.3-0.5 GeV

To cite this article: F V Adamian *et al* 1988 *J. Phys. G: Nucl. Phys.* 14 831



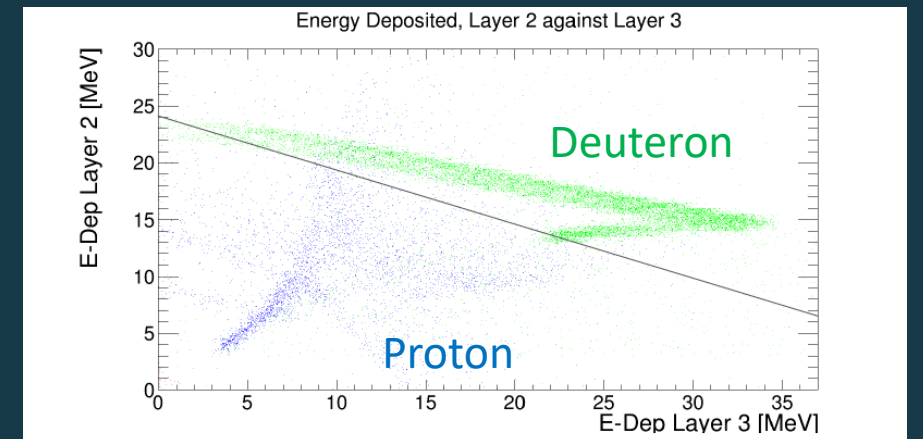
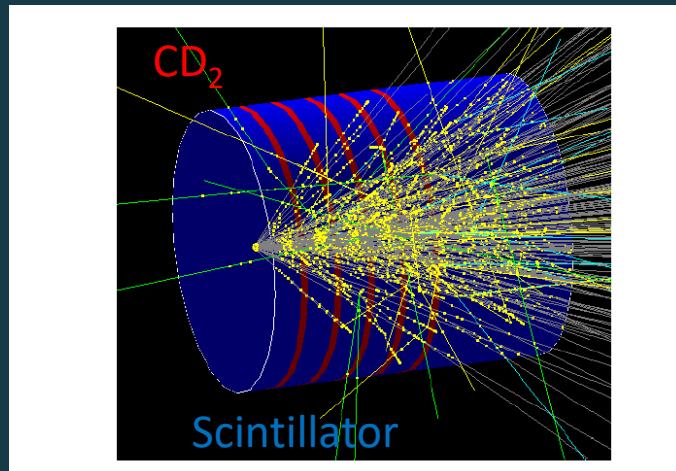
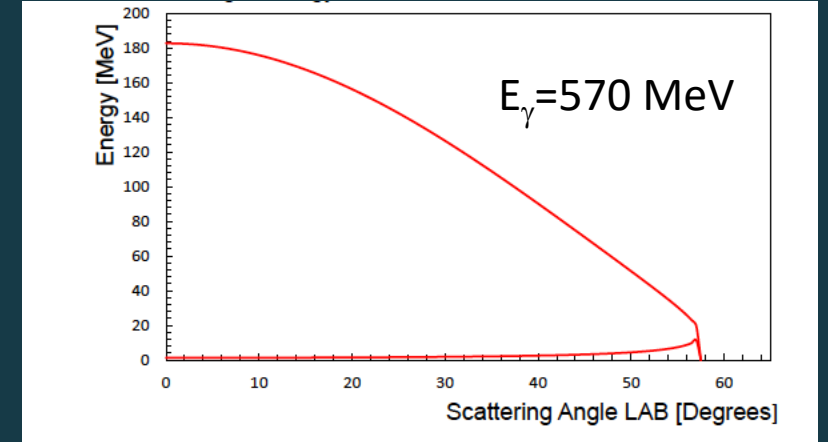
Active deuteron target?

Detecting recoil deuteron in photoproduction of $d\pi\pi$ challenging
-> e.g stops in cryogenic target for much of phase space

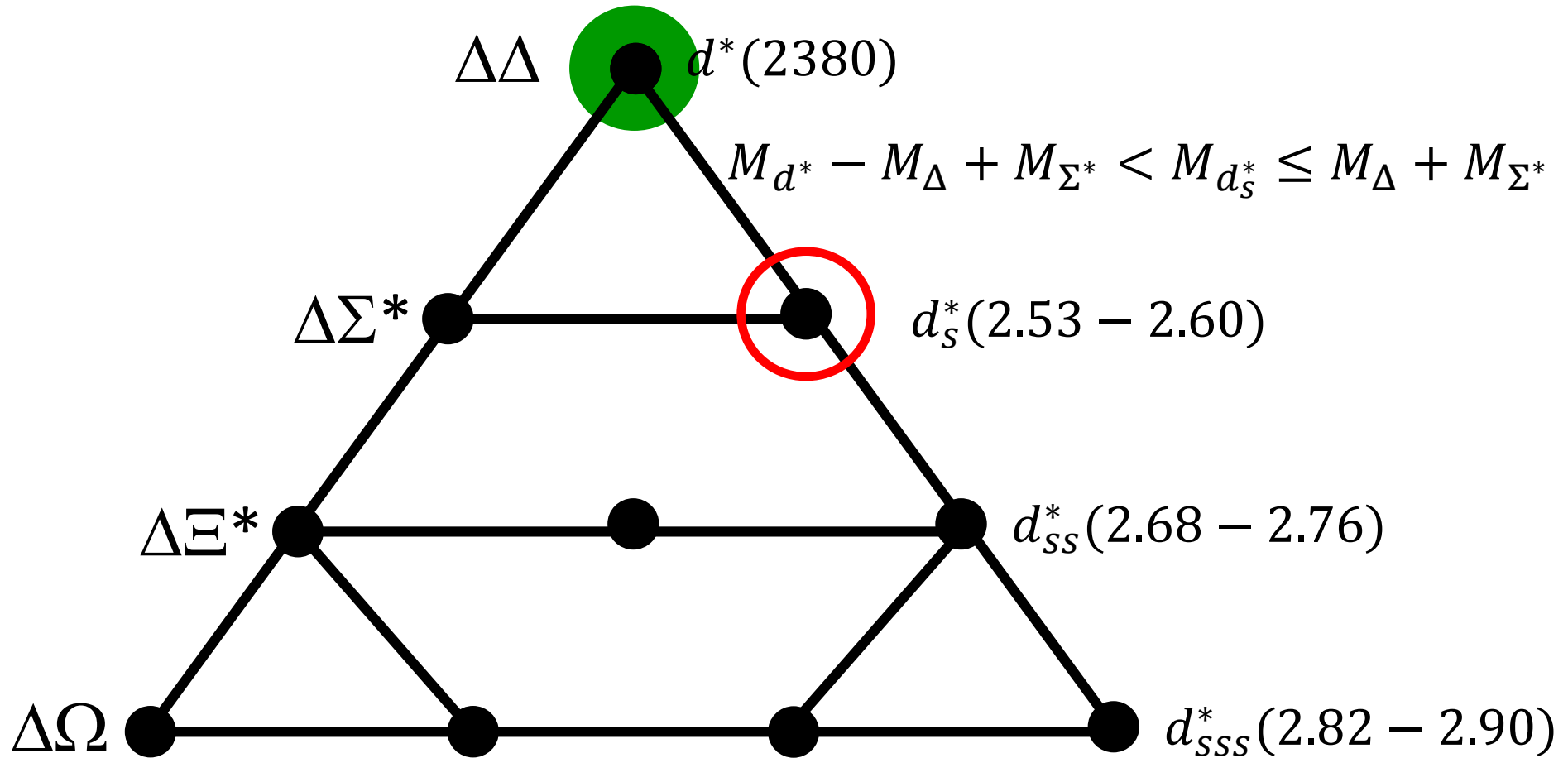
Active target prototype developed at UoY

Scintillator pixels interleaved with CD₂ target material

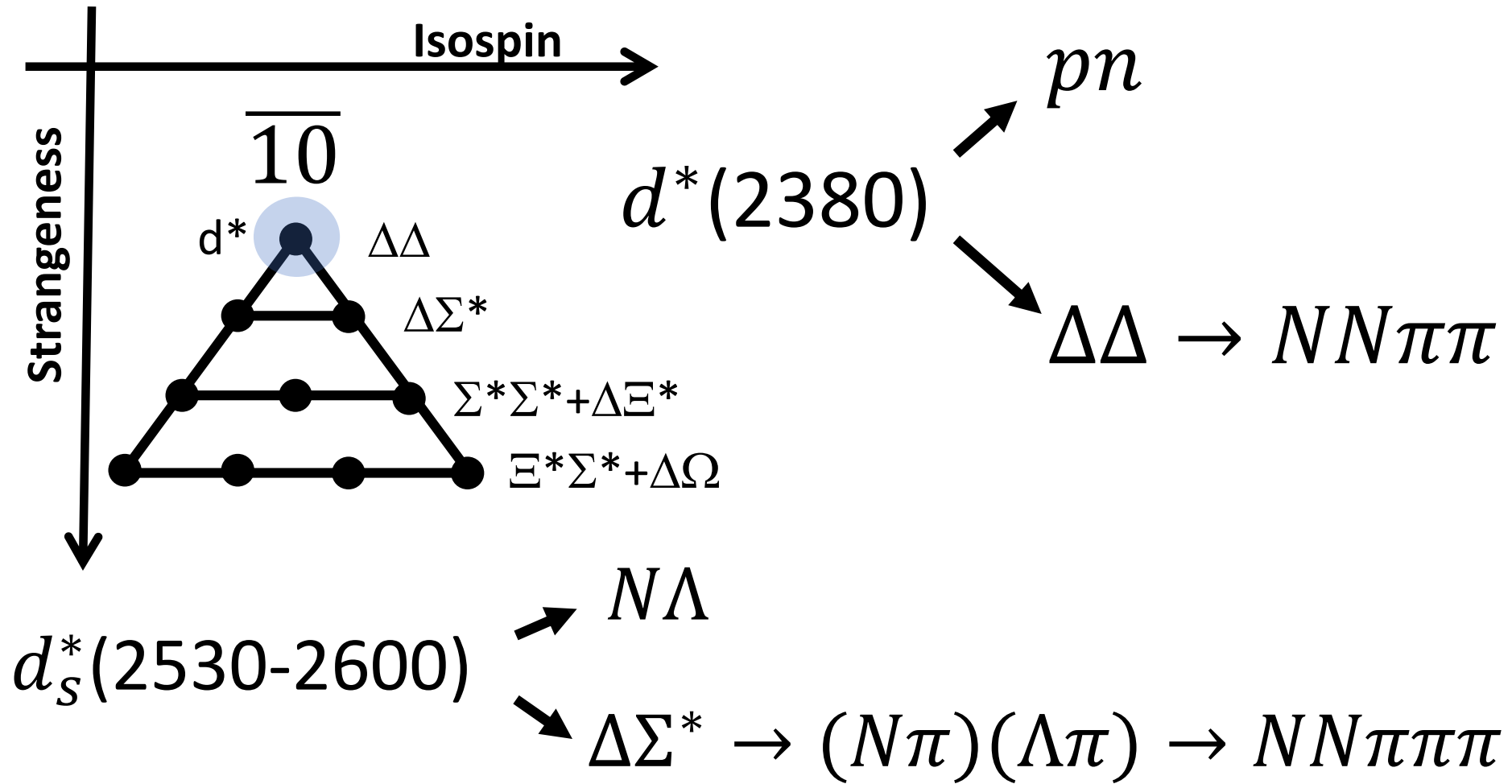
Tagging and separation of low energy Deuteron/spectator proton feasible
-> G4 simulation



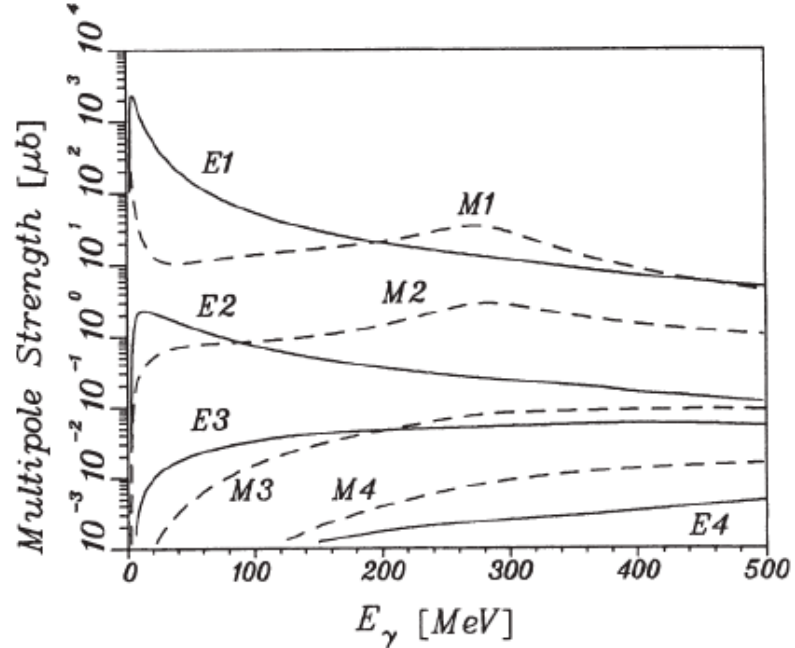
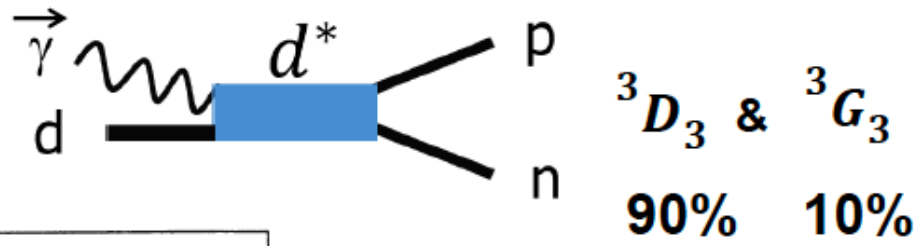
$d^*(2380)$ SU(3) multiplet



Strange Dibaryon decays



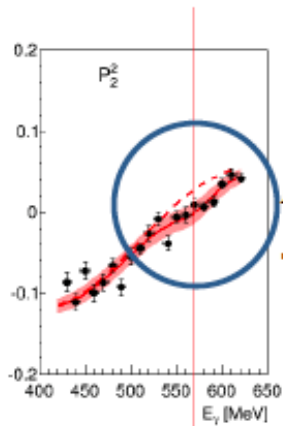
E2 transition (2⁺)
M3 transition (3⁺)
E4 transition (4⁺)



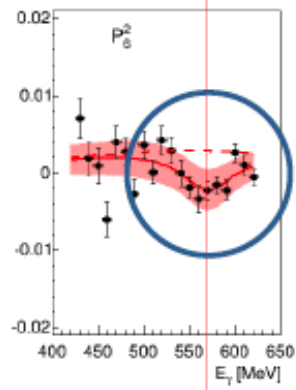
$$\frac{\sigma_{\perp} - \sigma_{\parallel}}{\sigma_{\perp} + \sigma_{\parallel}} = P_{\gamma} \Sigma \cos 2\phi$$

Fig. 7.1.3: Multipole strengths up to $L = 4$ contributing to the total cross section with inclusion of MEC, IC and RC for the Bonn r-space potential.

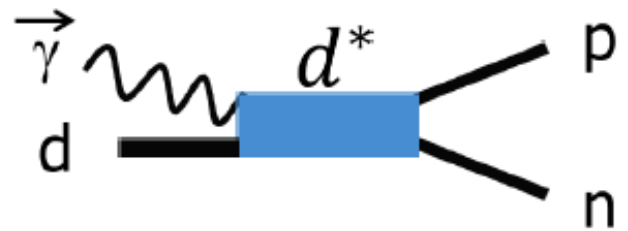
$$\frac{\Sigma(\theta)\sigma(\theta)}{\sigma_0} \sim \sum_{J=2} B_J P_J^2(\cos\theta)$$



$$B_2 \sim c_1 |^3D_3(E2)|^2 + c_2 |^3D_3(M3)|^2 + c_3 |^3G_3(E2)|^2 + c_4 |^3G_3(M3)|^2 + \dots$$



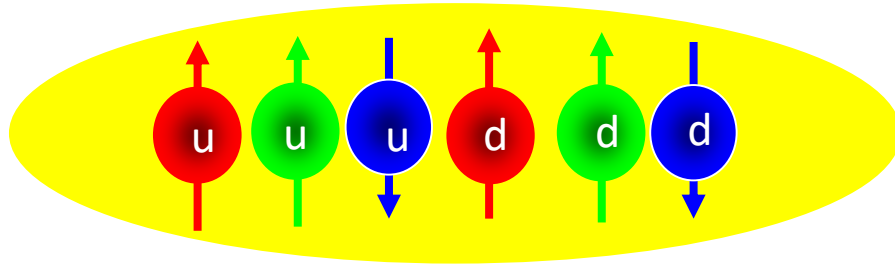
$$B_6 \sim d_1 |^3G_3(M3)|^2 + \dots + d_i |^3D_3(M3)| |^3G_3(M3)| \cos\delta_i + \dots + d_j |^3D_3(E2)| |^{2S+1}L \ge 4_J(E4)| \cos\delta_j + \dots$$



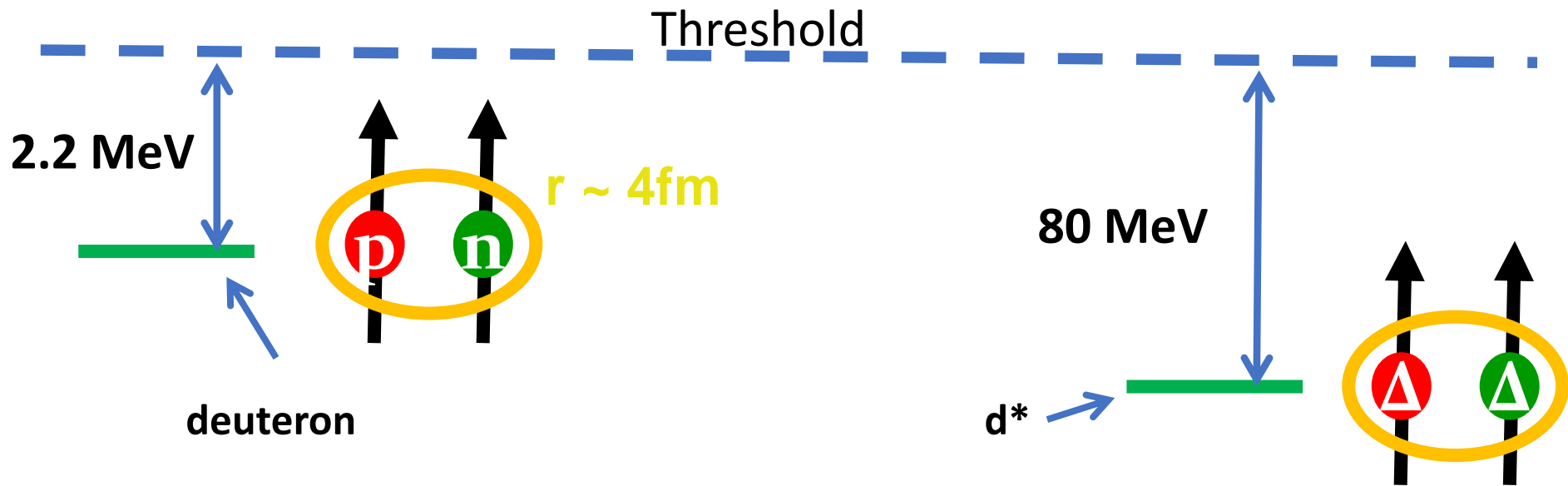
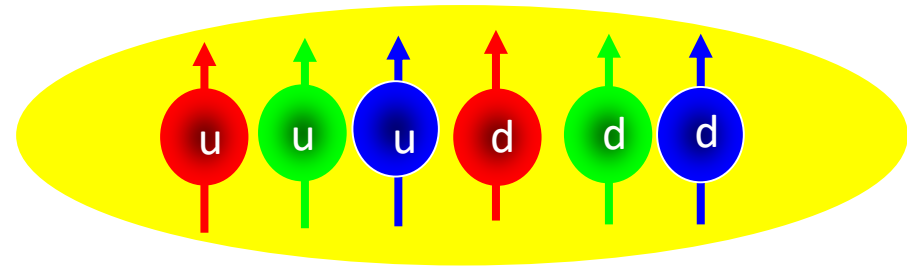
Magnetic M3 transition might be sizable

What is the d^* - Deltaron hypothesis

$$I(J^P) = 0(1^+)$$



$$I(J^P) = 0(3^+)$$



But width of d^* is 70 MeV – width($\Delta\Delta$) \sim 240 MeV?
& strong decay open ($NN\pi\pi$, $N\Delta\pi$) ?

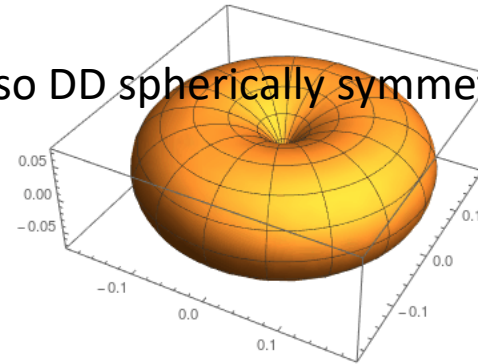
$r \sim 1.2\text{fm}$

d* shape

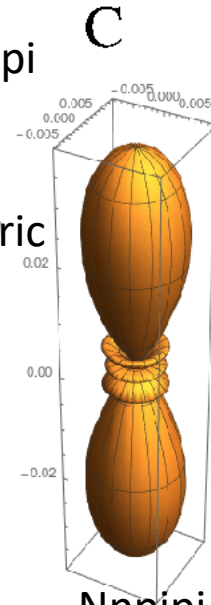
- Estimates made using simple pion cloud model.
- d* made from 2D spend part of time as nucleon plus pion – pure quark DD component, NNpp component and NDpi components
- Enables estimate M3/E2 ratios and ind. strengths
- Indicates d* is spherical (electrical quadropole deformation small)
- However magnetic octupole deformation may be large (don't cancel)
- Octopole moment is quadropole x magnetic dipole
- For 2 cases magnetic dipole moment of opposite sign (D0npi+ mag moment neut -2mu_b p = 3mu_b)
- Nnpi - Dpipi like (but more compact deuteron) NN is like deuteron therefore magnetic moment +1mu_b. The plot is a Y30 shape octupole.
- More polarization observables (long polarised target with transeverse beam or transeverse target circular beam) , Ox, may help to experimentally determine these ratios M3/E2. This ratio is proportional to mag dipole moment of d* (cancellation)
- M3 prop octupole d* E2 quadropole d*
- E2 photin spin up target L=1 in same direction 1- + 1+. 1- intrinsic spin 1- from angular gives 2+
- 2+ from E2 plus 1 from deut must be aligned t give 3+
- E4 1-L=1 in same direction gives 4+ E4 transition. 4+E4 and deuteron spin now in OPPOSITE direction to get 3+
- Even if d* is not compact the E2 is small because of cancellations. But in this model it is small. Radii of above similar in size to single delta
- Label arbitrary units probability density. Prefactors for use in calculation. Product of wavefunction with isospin factors.
- Quadropole deformation of delat from pion cloud only, but on lattic shown comes from quarks and has similar size starting rom just quarks.
- Doorway d* excited from d-wave of deuteron. Predictions give huge quadropole transition. Real cross section lot smaller. Therefore maybe not d-wave, maybe from 6q compinsnt. 3oom. If 6q get 0.15% so this might be it. 6q still deformed has cloud like nucleon and delta.

Pion cloud probablility
That would form around nDpi
nD spheric symmetric

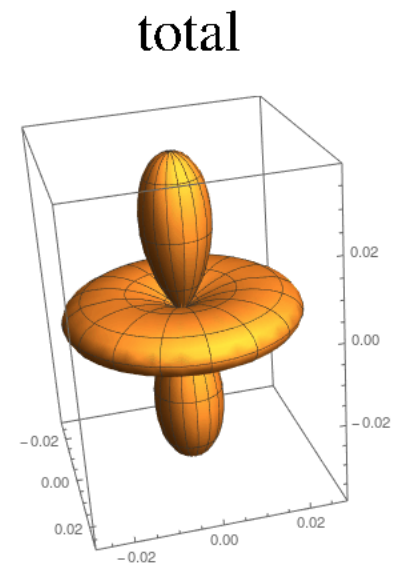
Also DD spherically symmetric



Ndpi
+0.73QD+



Nnpipi
-0.52QD+



Sum

