

Calorimetry

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Virtual Warwick week 2020



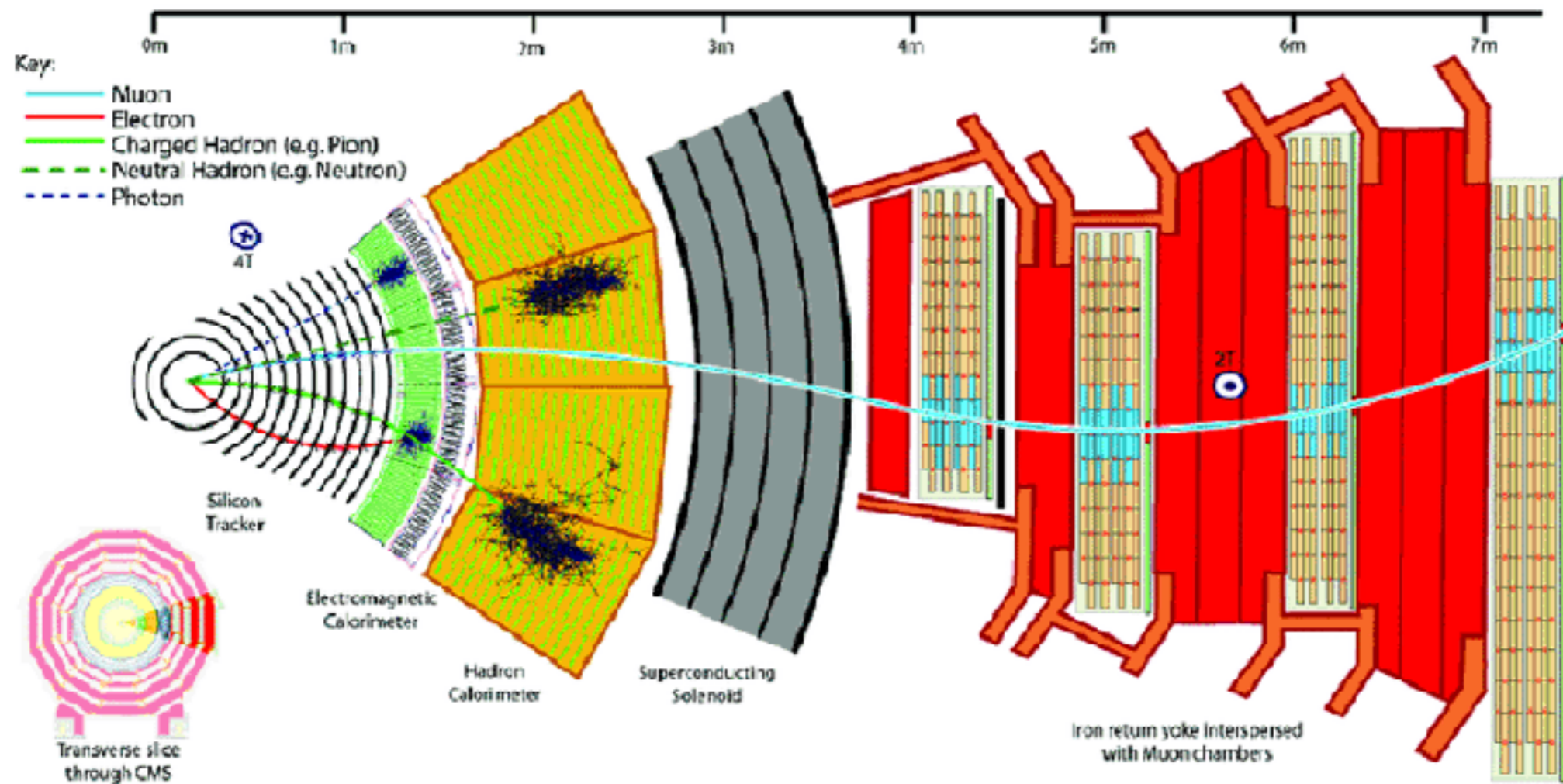
Science and
Technology
Facilities Council

1. Short overview
2. EM showers
3. EM calorimeters
4. Hadronic showers/calorimeters

Calorimetry

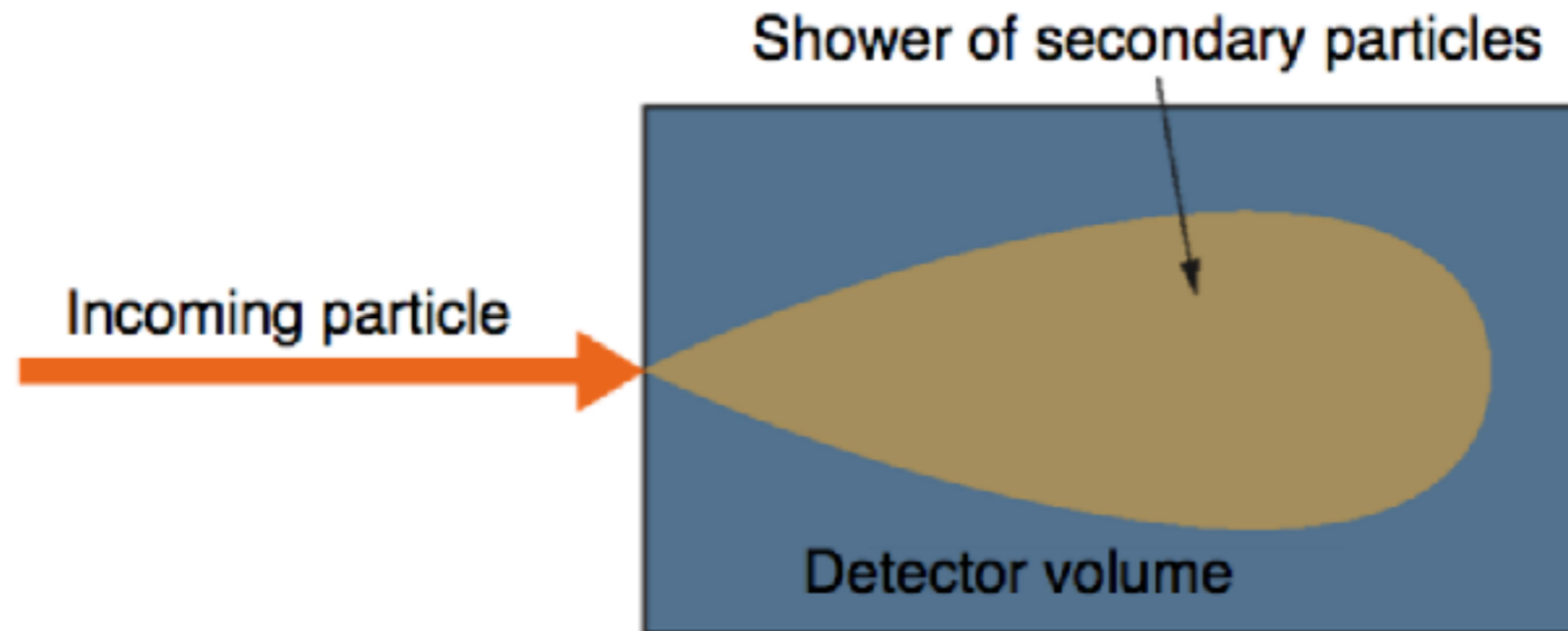
- Wide usage in particle physics, e.g.
 - 4 π (or LHCb-like) collider experiments
 - Instrumented targets
 - Shower counters
- Various detection mechanisms
 - Scintillation
 - Ionisation
 - Cerenkov
 - Cryogenics

Typical collider detector schematic



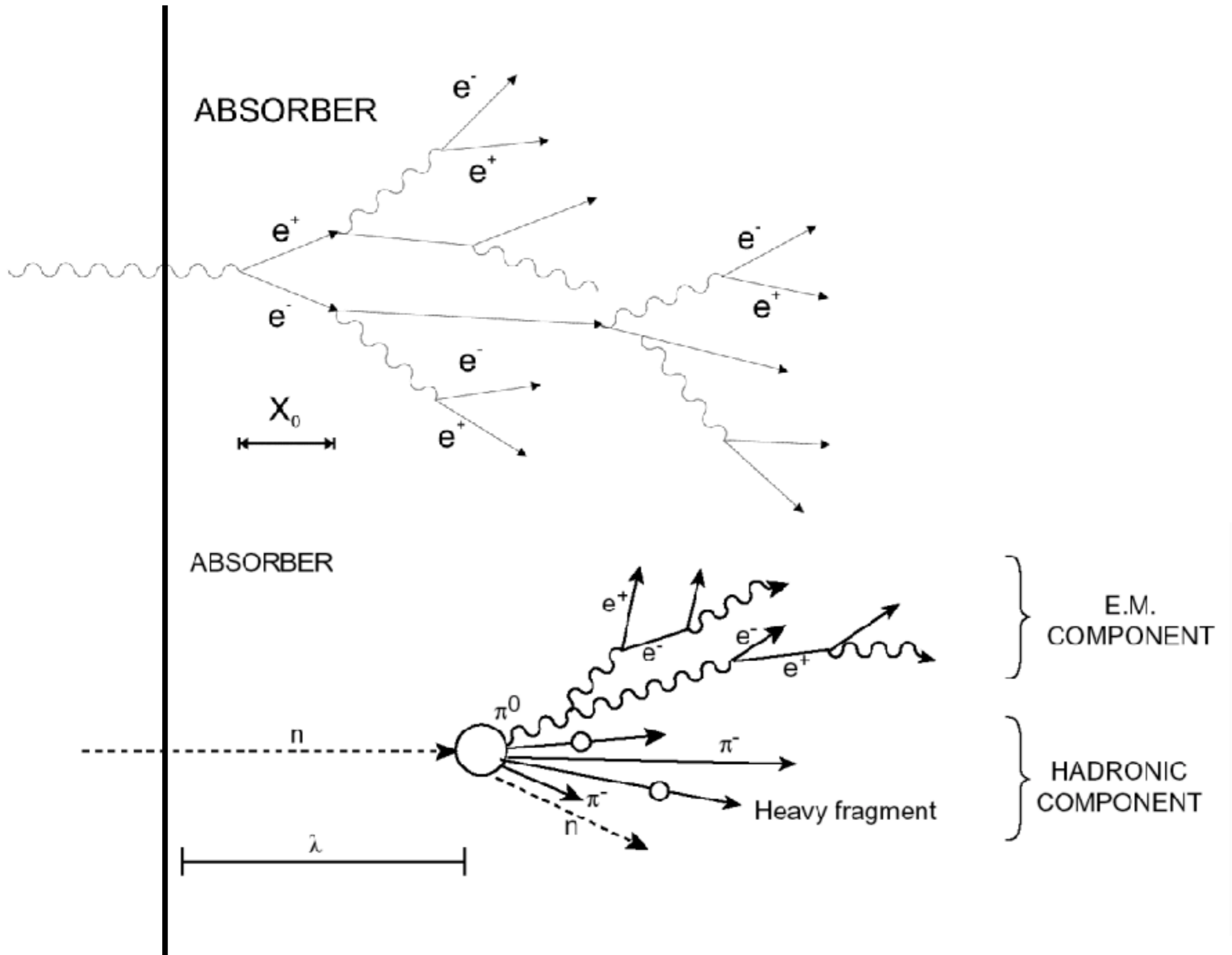
- Tracking system is ideally massless
- Calorimeter is *massive* and should totally absorb the energy of a particle [or jet] in 1 GeV to 1 TeV range.
- Electromagnetic and hadronic calorimeters.

The basic idea



- Stop/contain particle/jet by **shower** and absorption processes.
- Convert energy to signal with ionisation, scintillation etc..
- **Linearity** and good **resolution** desirable.
- Direction measurement for neutral particles.
- Missing transverse energy in $\sim 4\pi$ detectors.
- Intrinsically fast \rightarrow triggering.

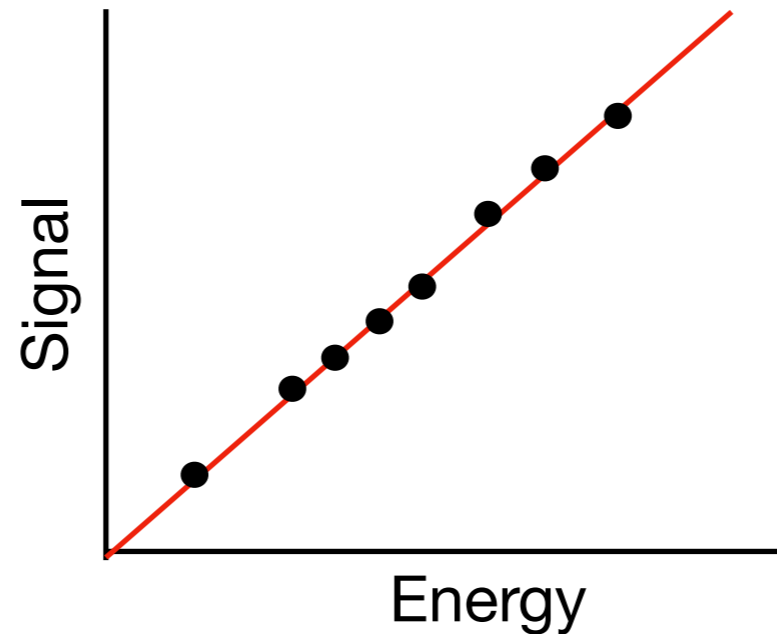
Electromagnetic and hadronic shower processes



Interplay with *stable visible* particles

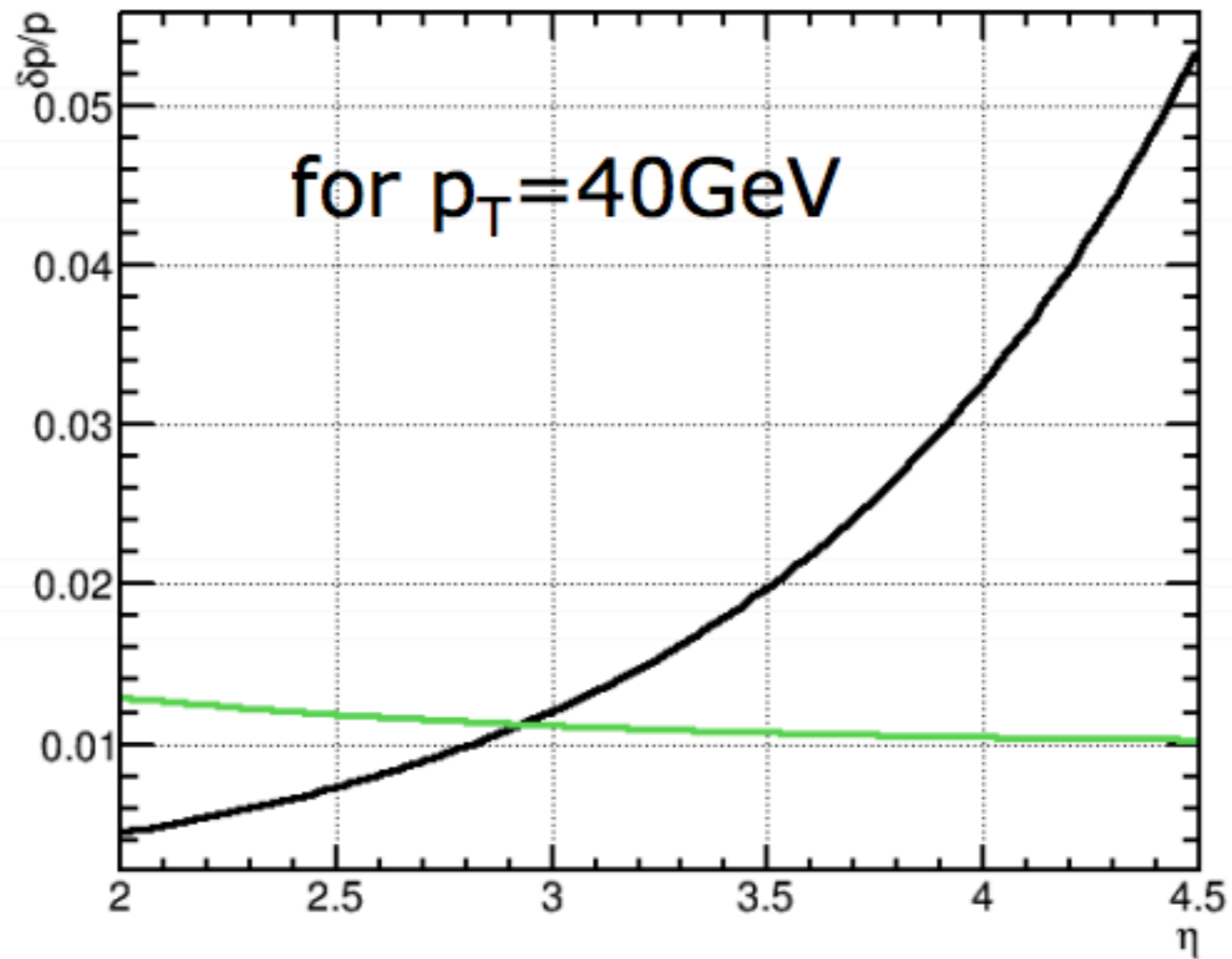
- Charged hadrons (π , K , p)
Hadronic showers
- Electrons and photons
Electromagnetic showers
- Neutral hadrons (n , K_L)
Hadronic showers
- Muons
Minimum ionising (track in calorimeter)

Linearity



- Readily achieved in EM calorimeters
 - Non-linearity can still be caused by, e.g., shower leakage, variation of response with depth, saturation of electronics etc...
- Hadronic calorimeters are intrinsically non-linear...

Complementarity with tracking

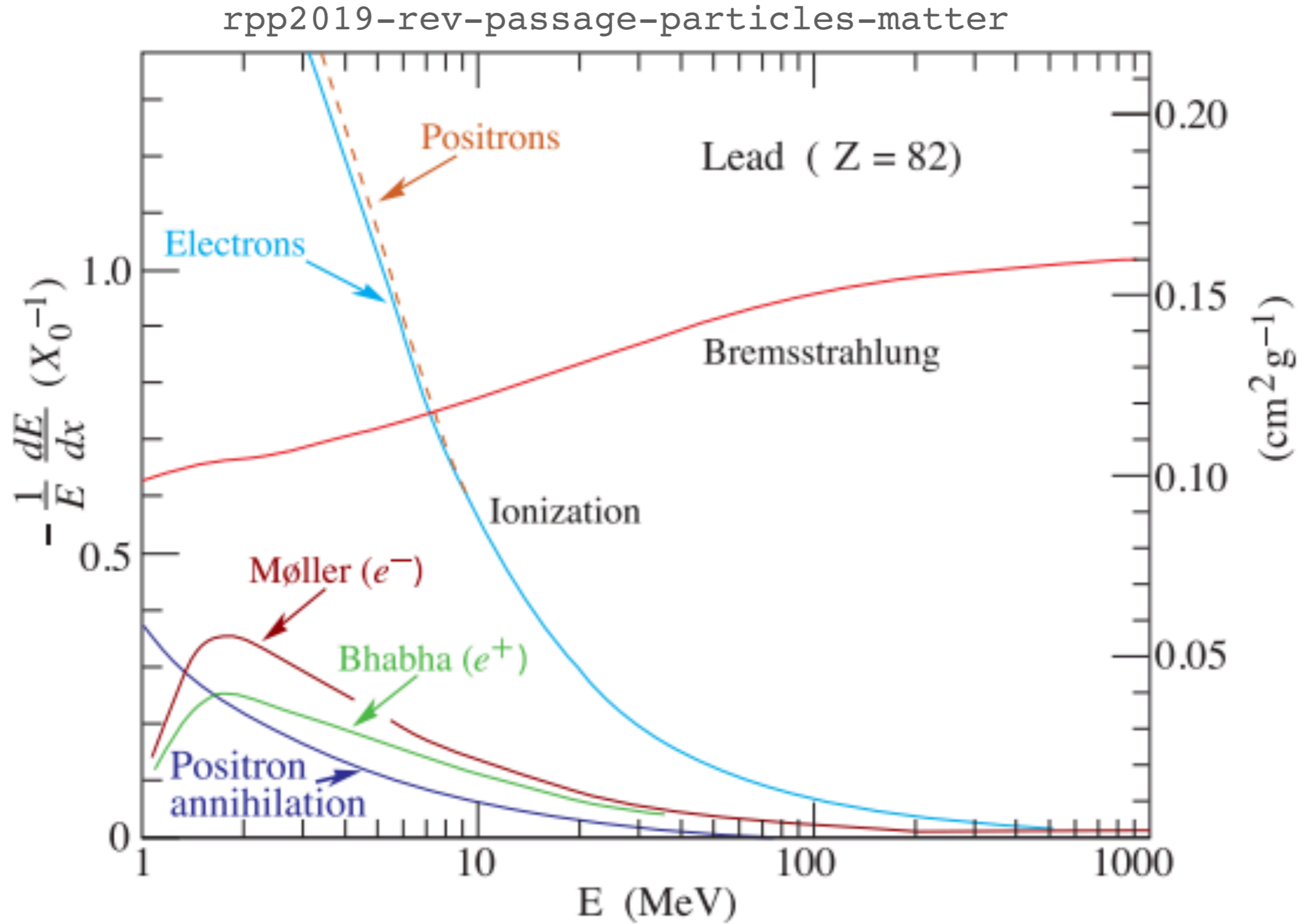


$$\frac{\sigma_p}{p} \sim p$$

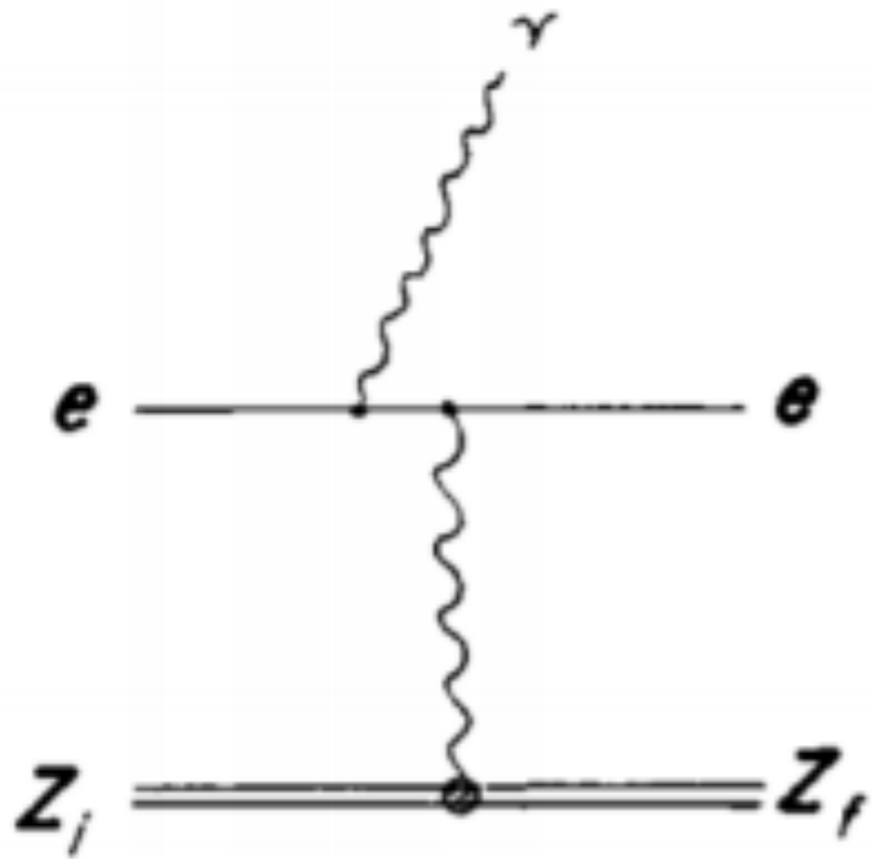
$$\frac{\sigma_E}{E} \sim 1/\sqrt{E}$$

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Electron interactions with matter



Bremsstrahlung: dominant for electrons at high energy



$$\sigma \sim \frac{1}{m^4}$$

$$\frac{dE}{dx} \propto E$$

Radiation length* (X_0)

$$\frac{1}{E} \left(\frac{dE}{\rho dx} \right) = -\frac{1}{X_0} \quad E = E_0 e^{-x/X_0}$$

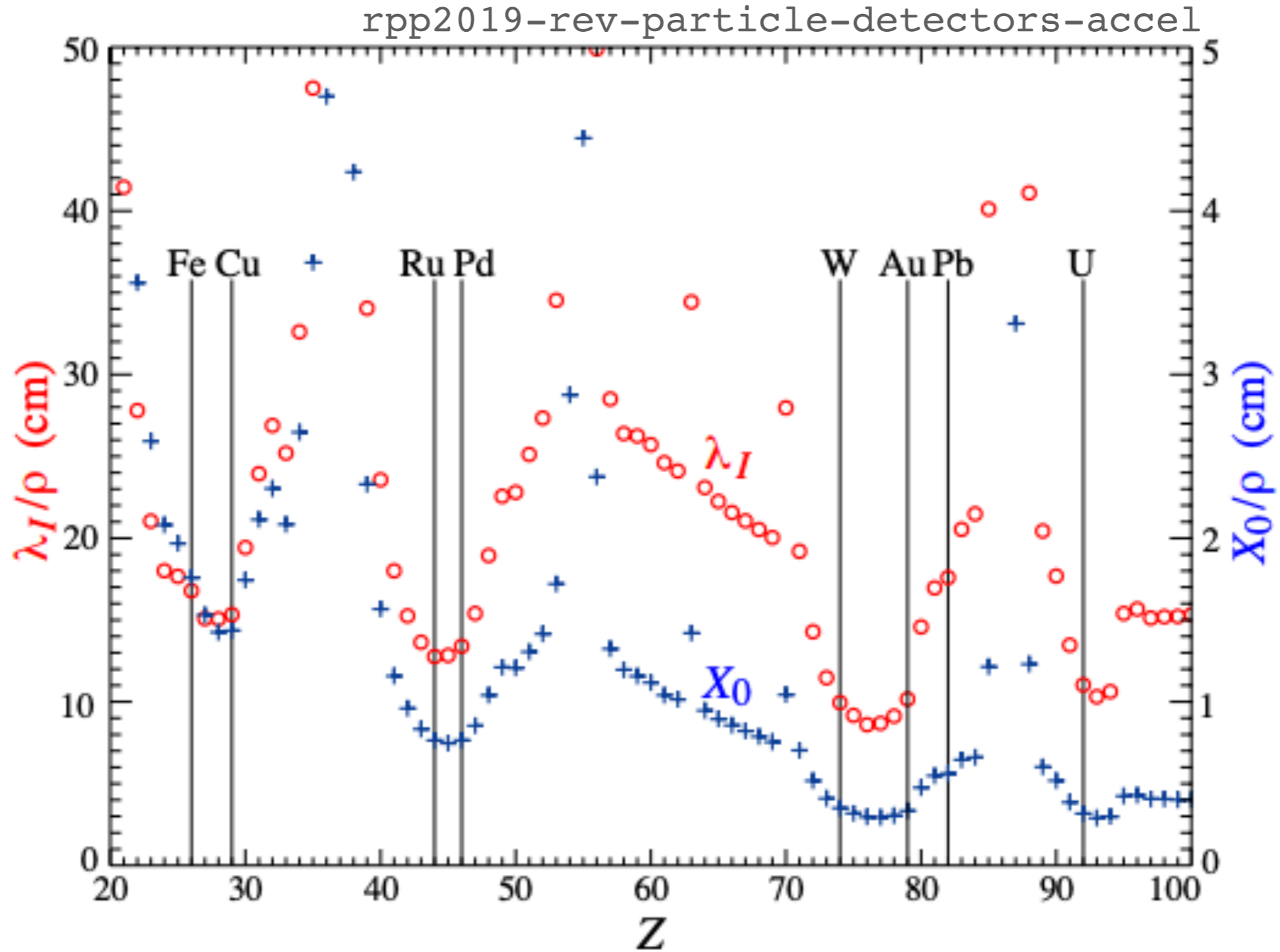
~Density of scattering centres

Approximation:

$$X_0 \sim \left[180 \frac{\text{g}}{\text{cm}^2} \right] \frac{A}{Z^2} \quad \sigma_{\text{brem}} \propto Z^2$$

👍 if we express material thickness in X_0 then the radiation loss is independent of material.

Material dependence



👍 X_0/ρ is a convenient quantity [with length units].

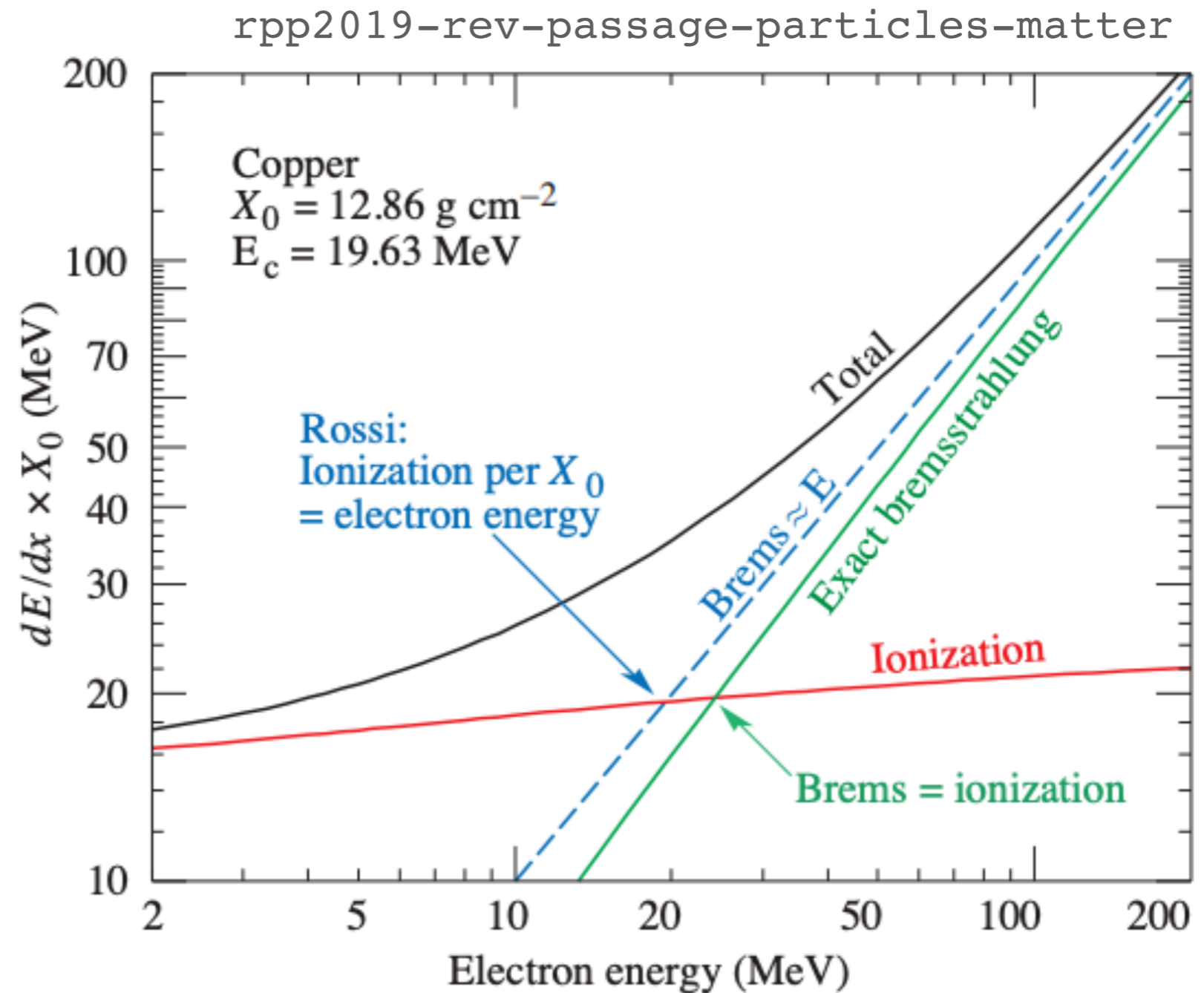
The critical energy (E_c),

at which Brem. and ionisation losses are equal.

$$\text{Ionisation} \sim \frac{Z}{A} \beta^{-2}$$

$$\text{Brem} \sim Z^2$$

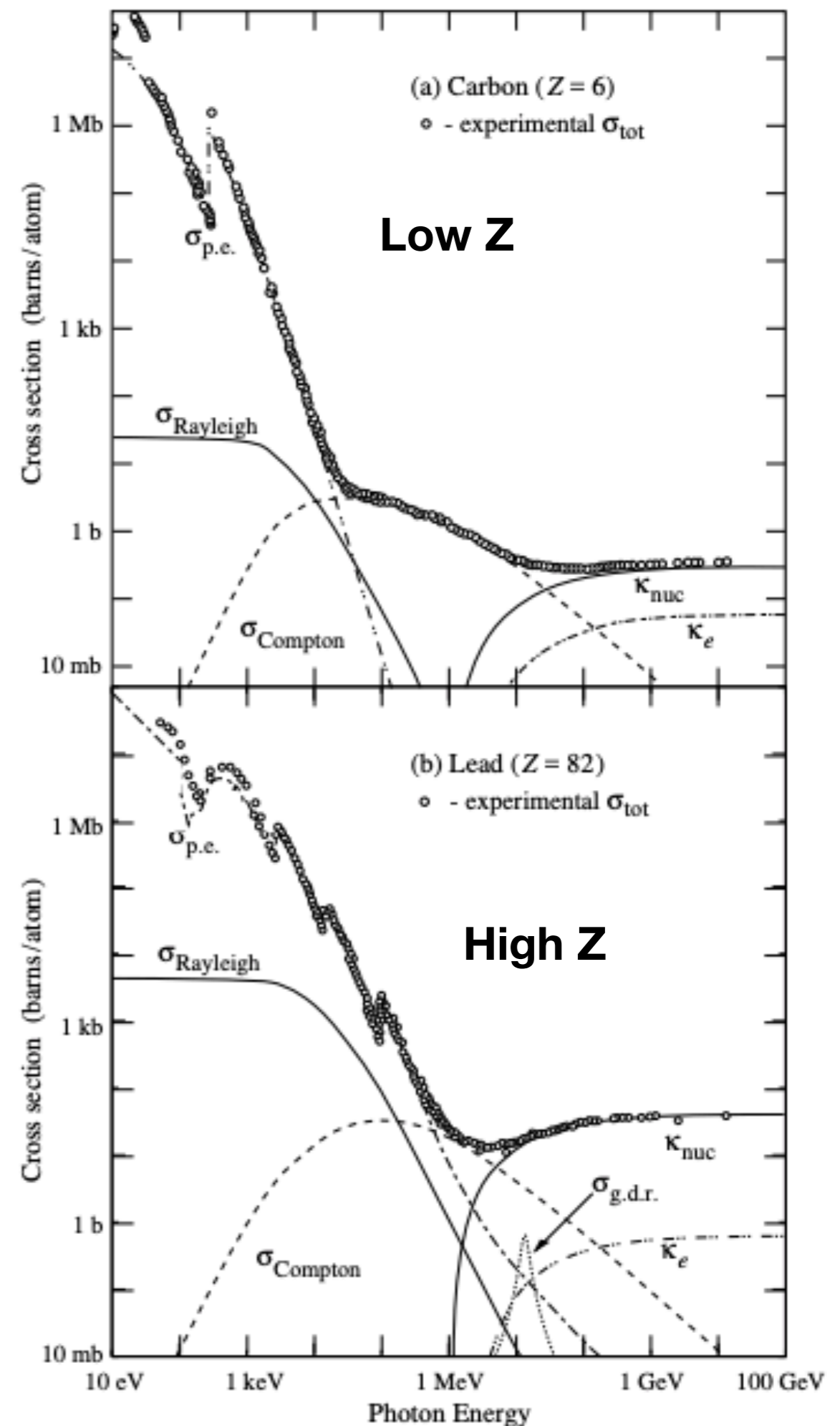
$$E_c \approx \frac{600 \text{ MeV}}{Z}$$



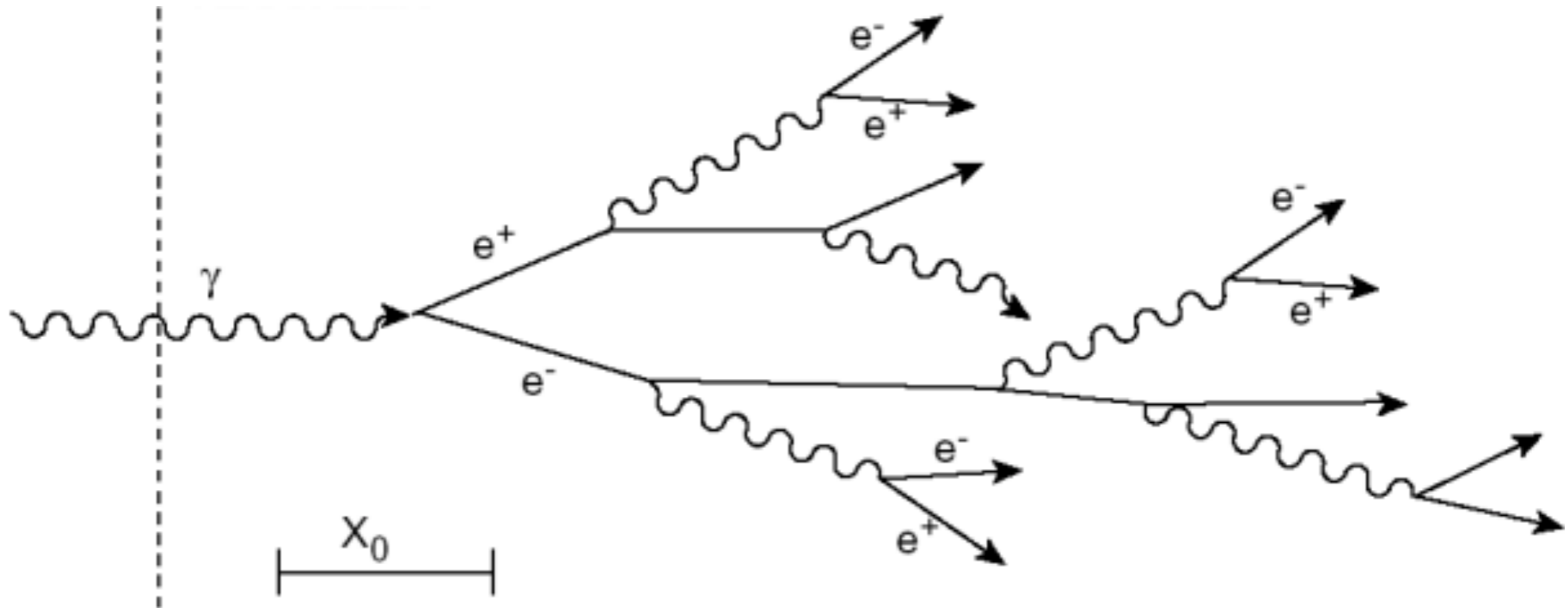
Photon interactions

1. Pair production at high energy
2. Compton scattering at lower energy
3. PE effect at even lower energy

$$\lambda_\gamma \approx \frac{9}{7} X_0$$



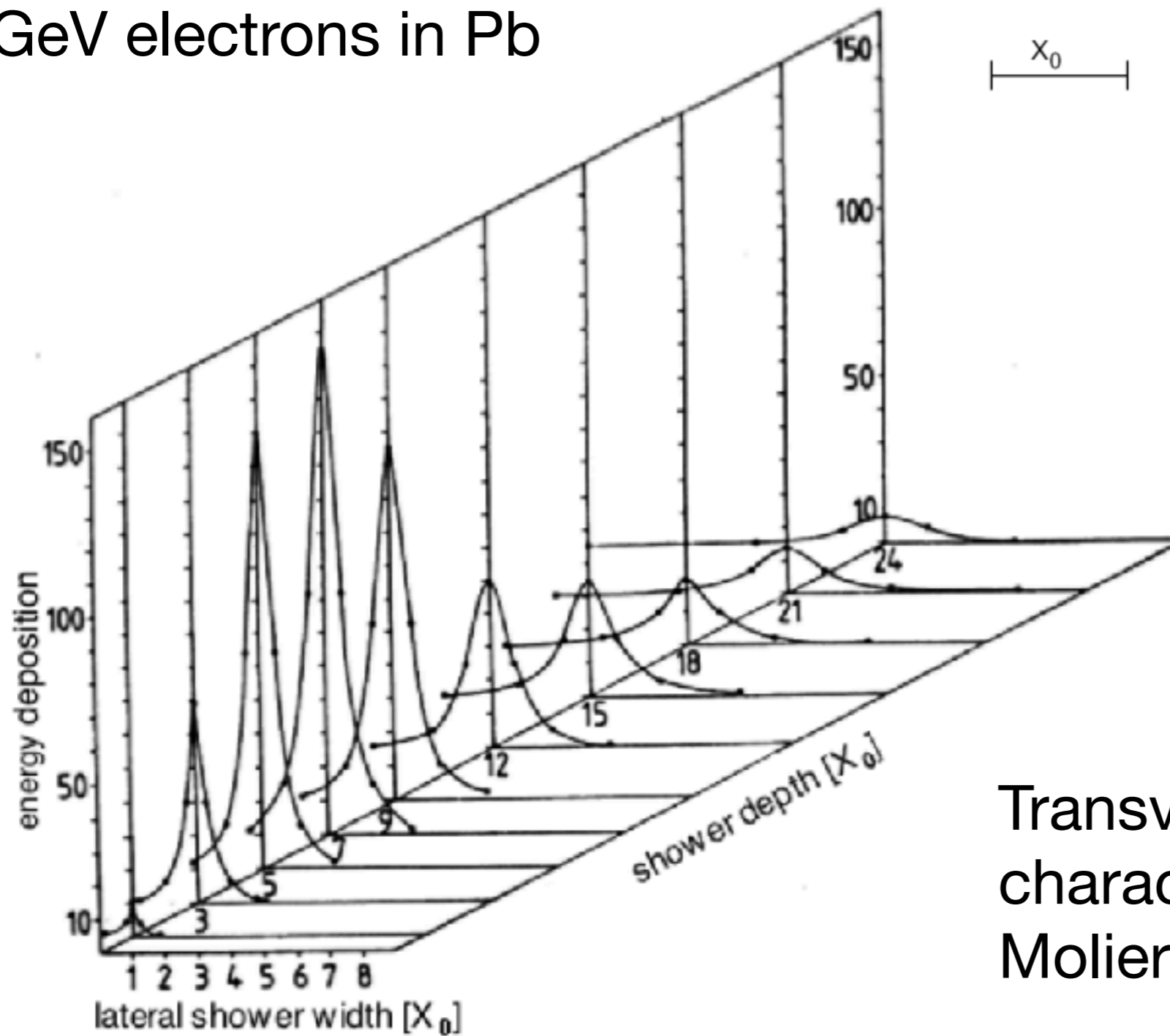
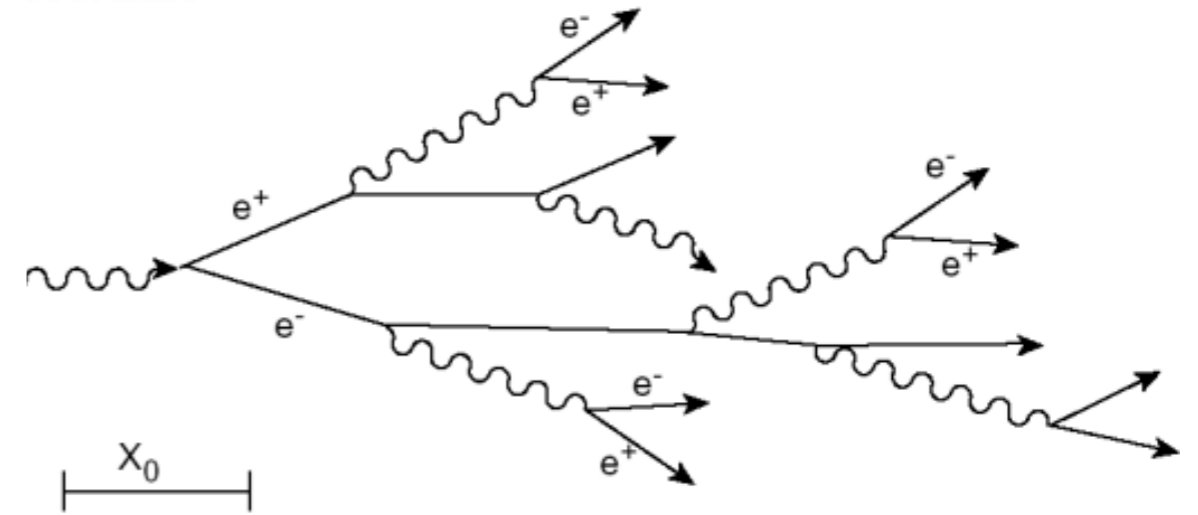
Electromagnetic shower



- Secondary electrons/photons from pair production and bremsstrahlung.
- Number *increases* but mean energy *decreases*.
- Ionisation and excitation take over when mean energy falls below E_C .

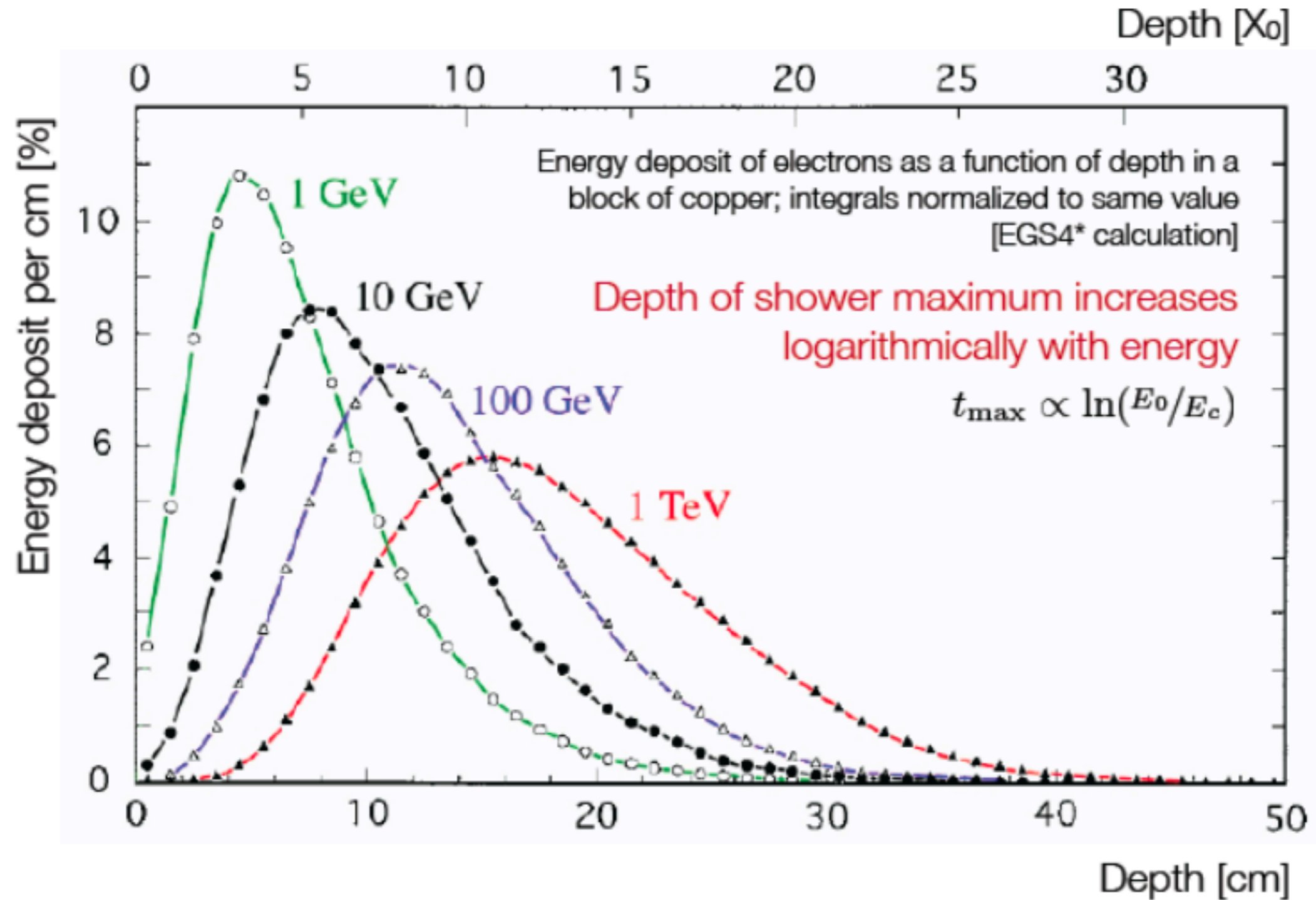
Shower development

6 GeV electrons in Pb



Transverse shower extent characterised by the Moliere radius (R_M).

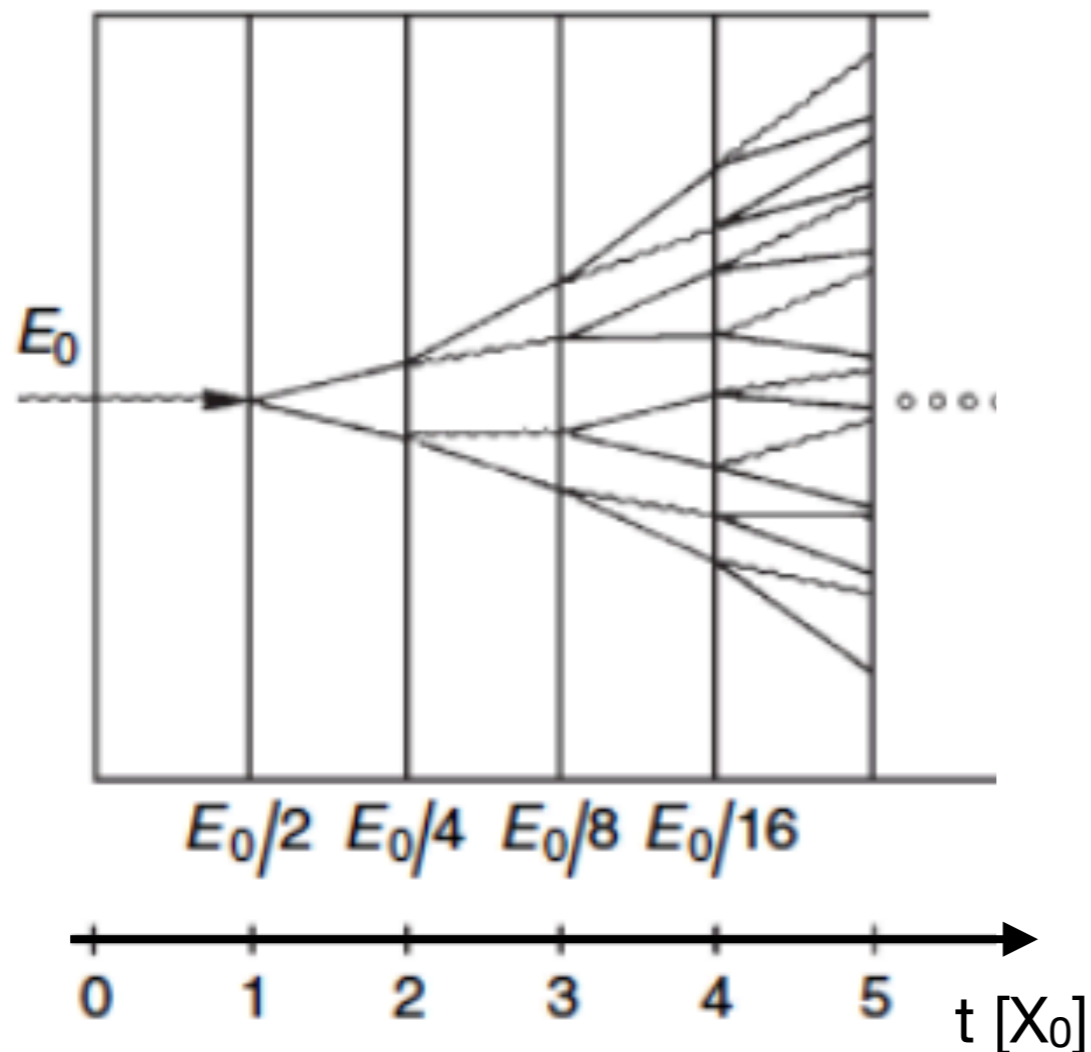
Depth of shower max



Depth only has $\log(E)$ scaling \rightarrow can build compact calorimeters!

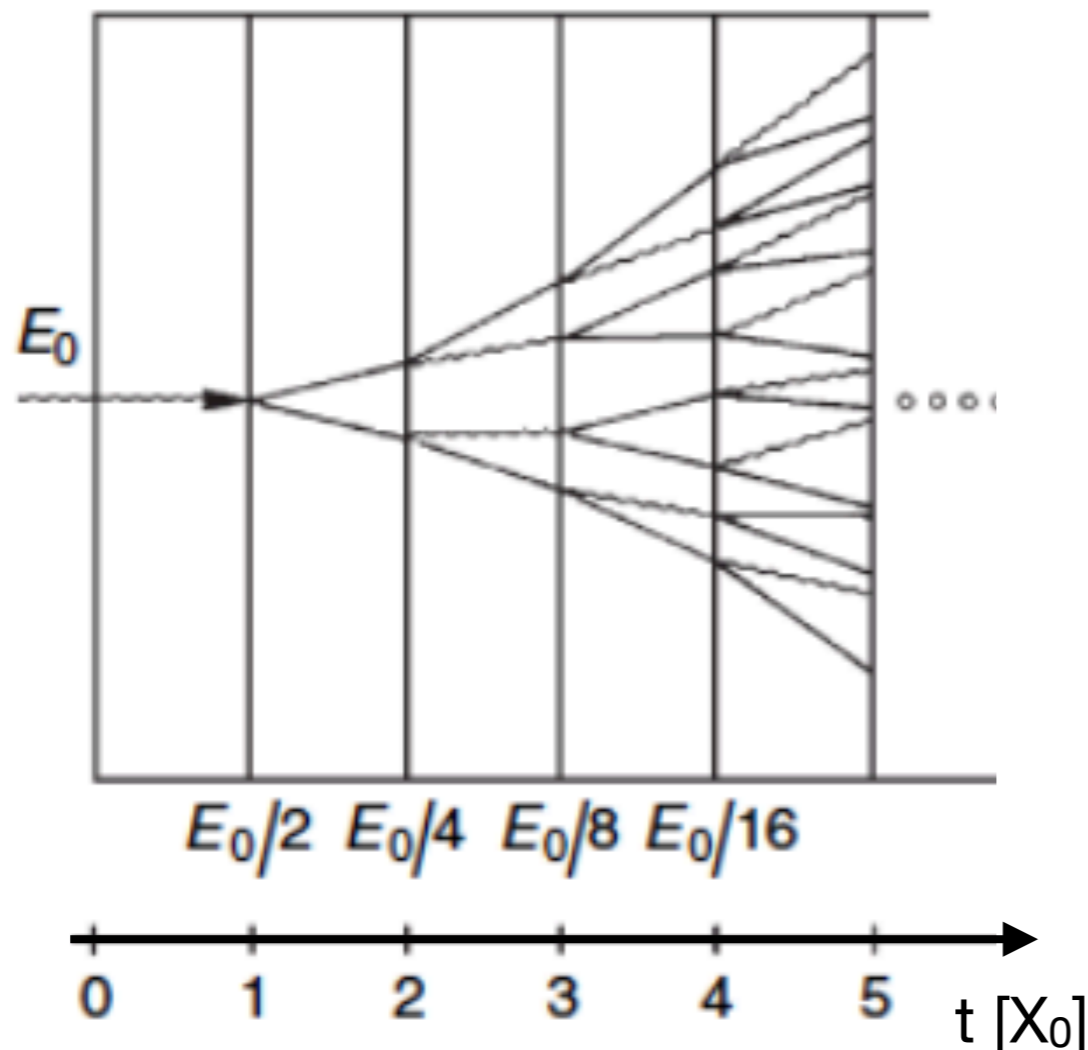
Simple shower model

1. In 1 X_0 Electron loses $[1-1/e]$ of energy to Brem
2. Mean free path of photon is $9/7 X_0$ (pair production)
3. No(only) ionisation/excitation above(below) E_C .



Simple shower model

1. In 1 X_0 Electron loses $[1-1/e]$ of energy to Brem
2. Mean free path of photon is $9/7 X_0$ (pair production)
3. No(only) ionisation/excitation above(below) E_c .



After $t [X_0]$ we have 2^t particles with energy $E/2^t$

Shower stops when $E < E_c$, with $N=E/E_c$ particles

Shower max at $t_{\max} \sim \ln(E_0/E_c)$

Lateral shower development

- Bremsstrahlung and pair production at small angles because m_e is small.

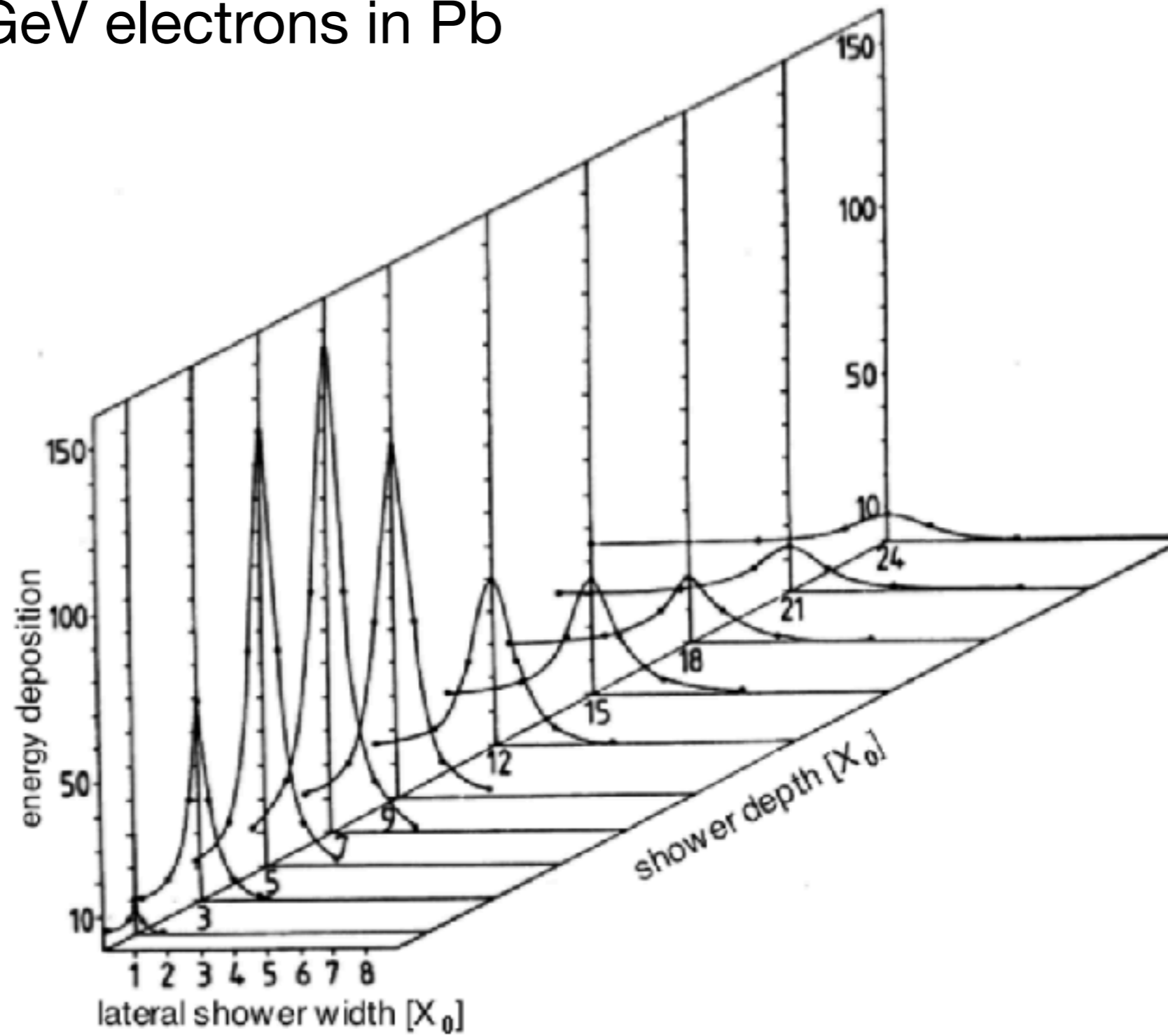
$$\langle \theta^2 \rangle \sim (m/E)^2$$

- Multiple coulomb scattering [Mollier theory] of low energy electrons dominates lateral spread.
- Characteristic Mollier radius

$$R_M \approx \left[7 \frac{\text{g}}{\text{cm}^2} \right] \frac{A}{Z}$$

R_M is a crucial consideration when specifying the segmentation (calorimeter cell size).

6 GeV electrons in Pb



Back of the envelope EM shower characteristics

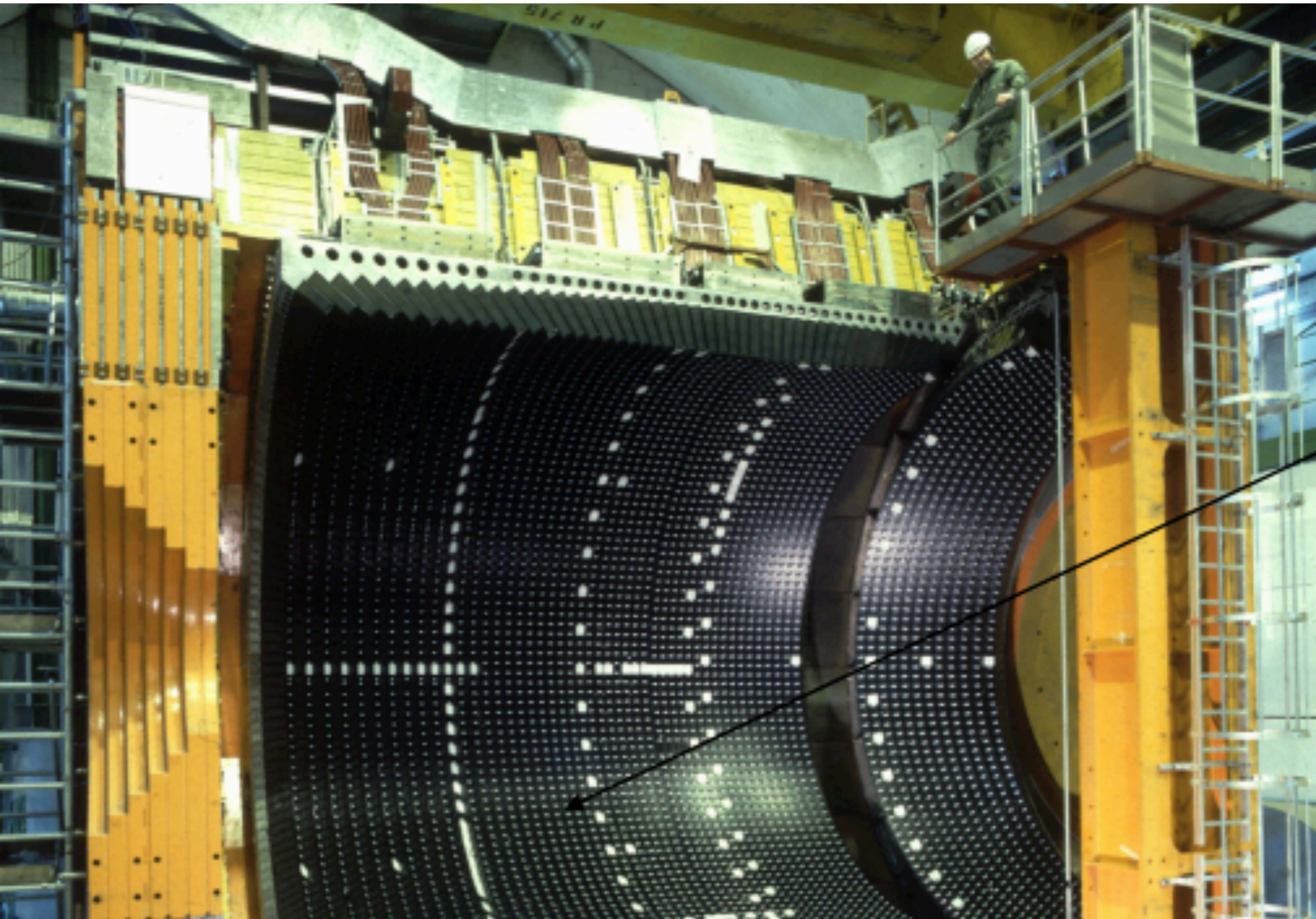
Radiation length $X_0 \sim \left[180 \frac{\text{g}}{\text{cm}^2} \right] \frac{A}{Z^2}$

Critical energy $E_c \approx \frac{600 \text{ MeV}}{Z}$

Shower max $t_{\text{max}} \approx \ln \frac{E}{E_c}$

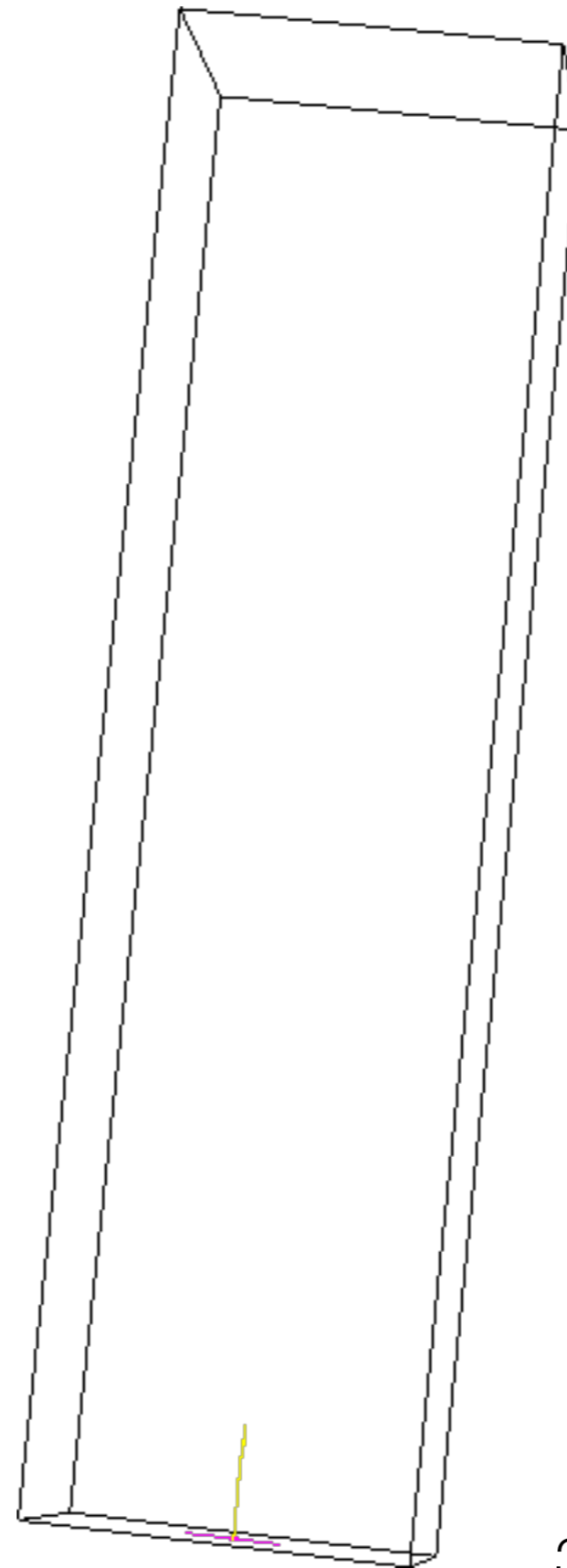
Lateral $R_M \approx \left[7 \frac{\text{g}}{\text{cm}^2} \right] \frac{A}{Z}$

OPAL experiment lead glass



OPAL experiment lead glass

<https://www.mpp.mpg.de/~menke/elss/>

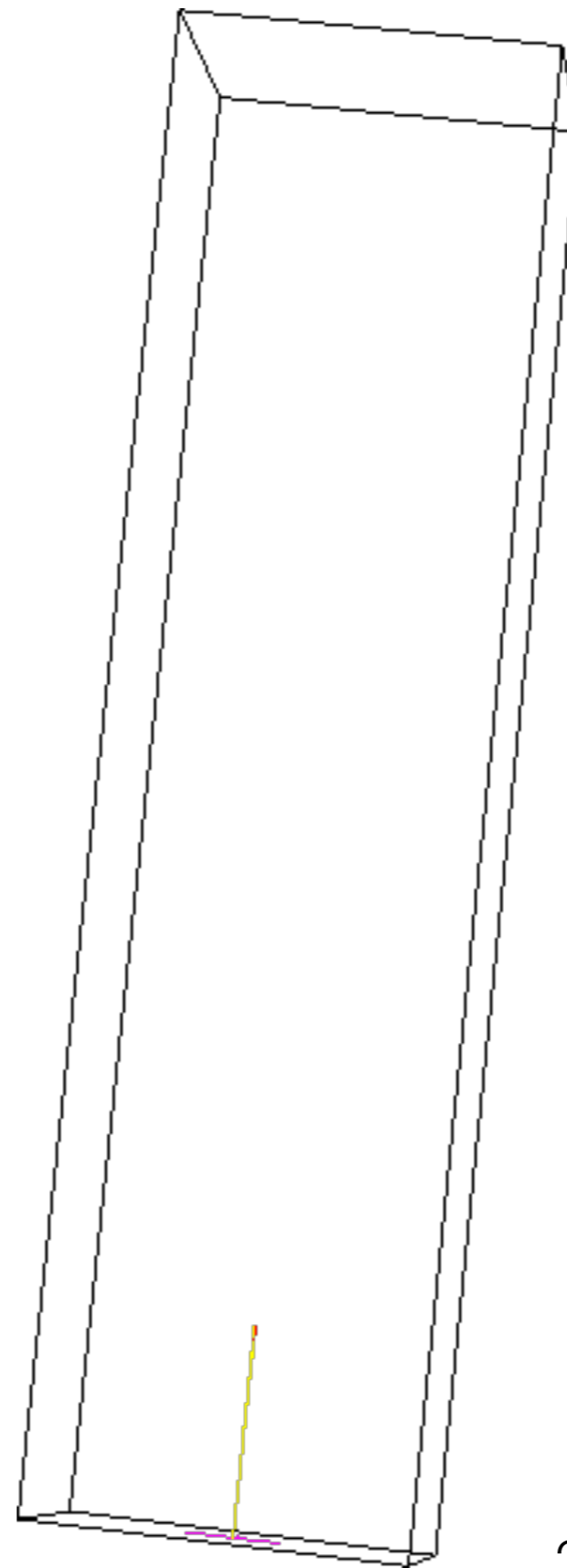


1 GeV e^-

37 x 10 x 10 cm

OPAL experiment lead glass

<https://www.mpp.mpg.de/~menke/elss/>

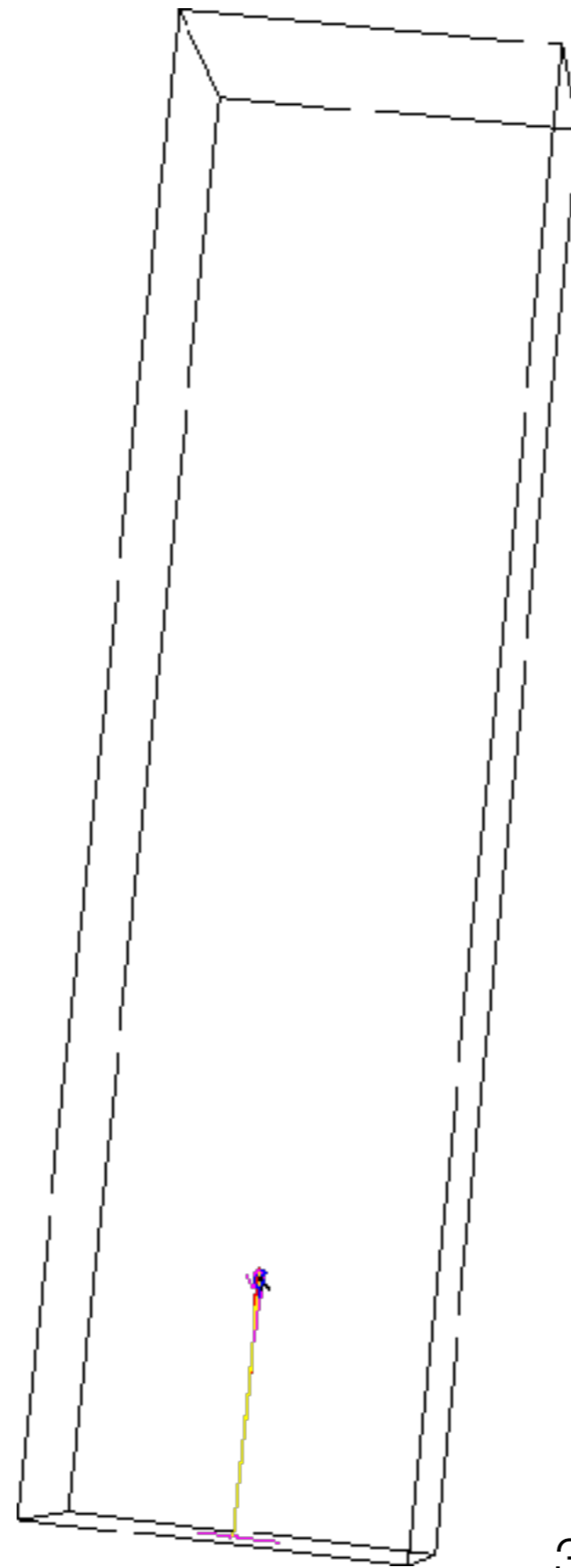


10 GeV e^-

37 x 10 x 10 cm

OPAL experiment lead glass

<https://www.mpp.mpg.de/~menke/elss/>

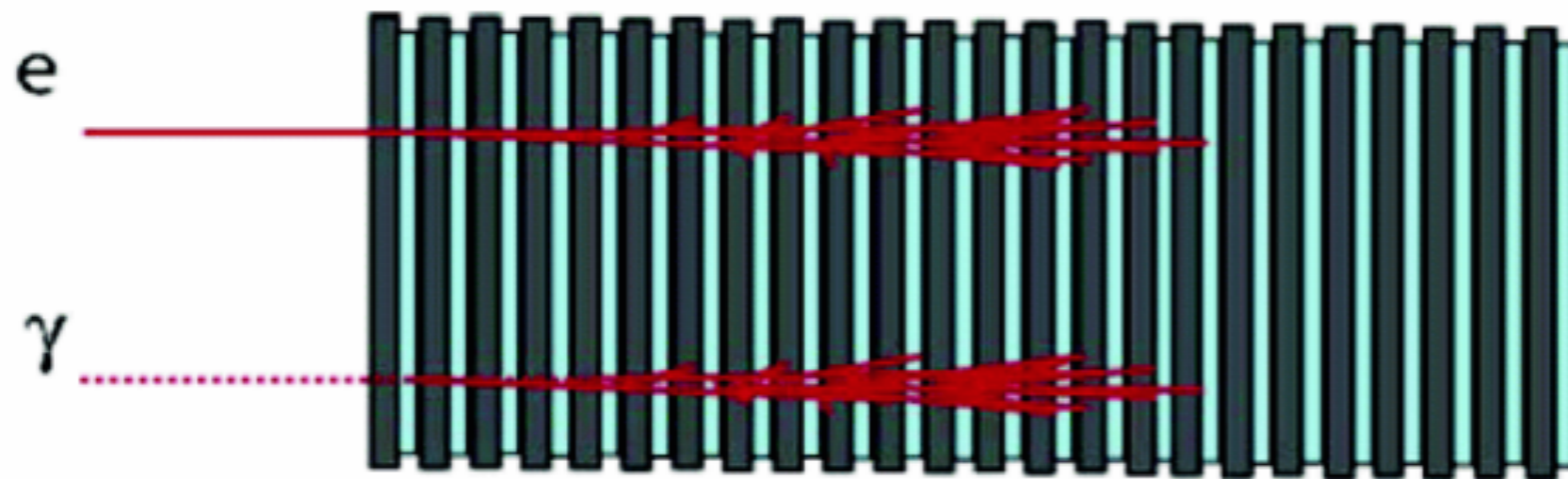


80 GeV e^-

37 x 10 x 10 cm

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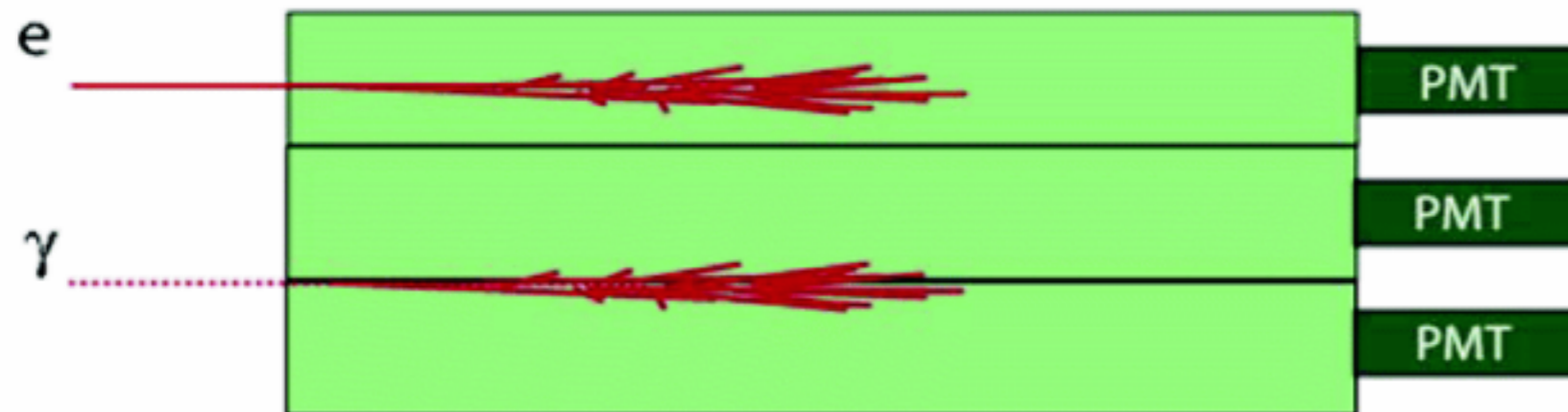
Sampling



Typically:

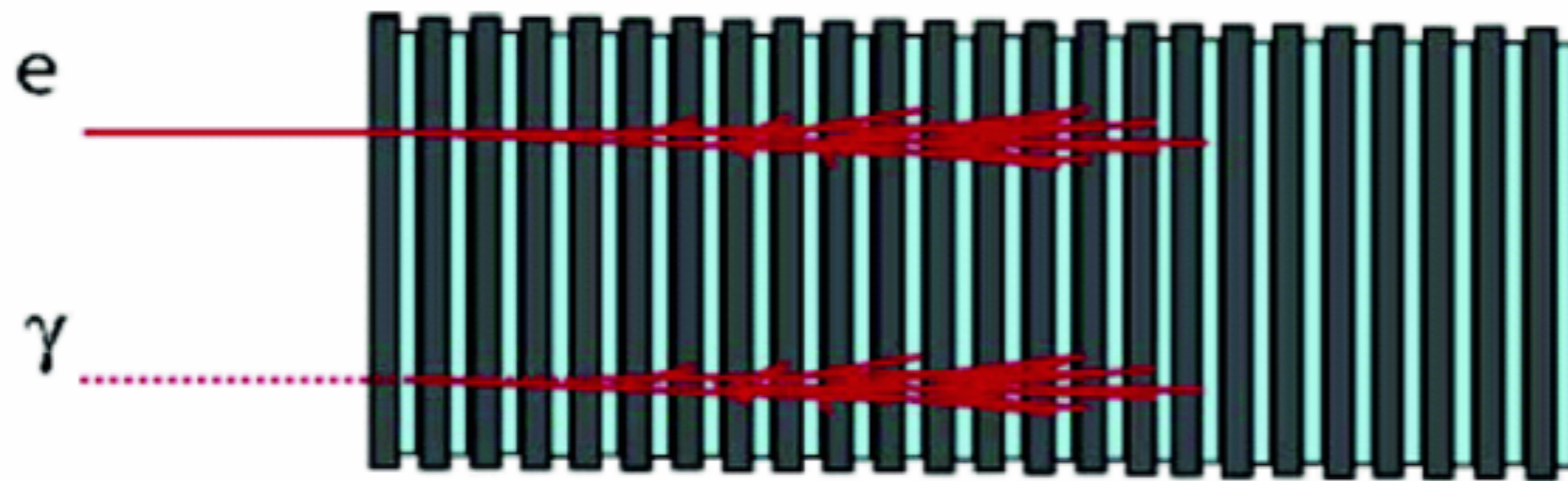
$$\sigma_E/E \sim 10\% / \sqrt{E}$$

Homogenous



$$\sigma_E/E \sim 1\% / \sqrt{E}$$

Sampling

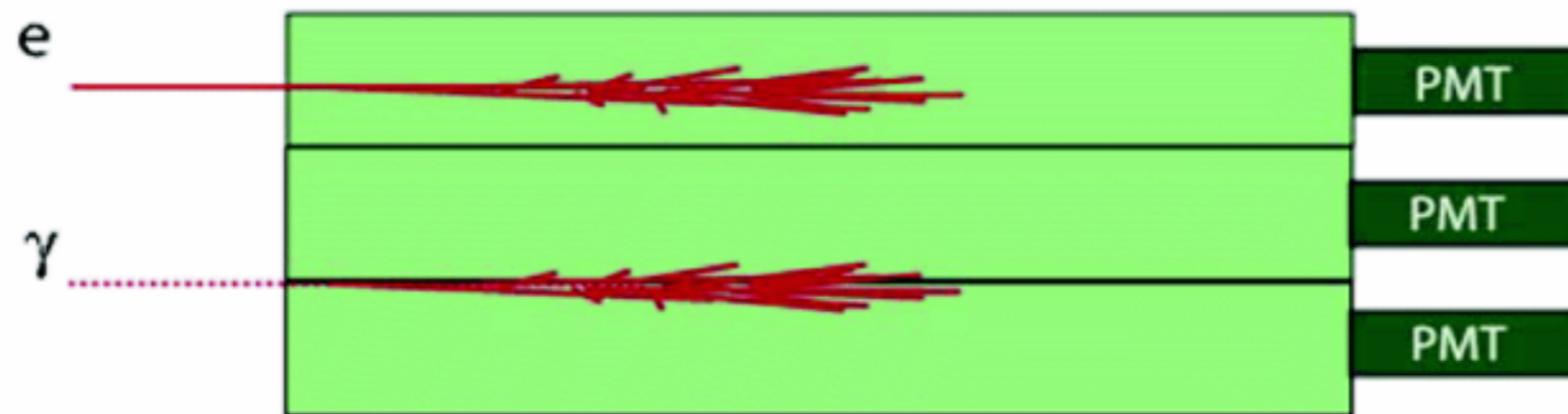


Typically:

$$\sigma_E/E \sim 10\% / \sqrt{E}$$

- 👍 Freedom to independently choose optimal absorber and active detector material
- 👍 Dense absorber → compact calorimeters
- 👍 Can be cost effective (cheap absorber)
- 👎 Only the sampling fraction of energy is measured
→ compromised resolution

Homogenous



Typically:

$$\sigma_E/E \sim 1\% / \sqrt{E}$$

- 👍 Good resolution because all shower particles seen
- 👍 Uniform response → linearity
- 👎 Expensive and limited segmentation

Special use cases e.g.

1. “medium energy” ECAL-only B-factory experiments,
2. CMS and ultimate $H \rightarrow \gamma\gamma$ mass resolution

EM energy resolution

$$\frac{\sigma_E}{E} = \frac{S}{\sqrt{E}} \oplus \frac{N}{E} \oplus C$$

S: sampling or stochastic term

Fluctuations in the signal generating process

The ideal calorimeter has $E \sim N$, $\sigma \sim \sqrt{N} \sim \sqrt{E}$

N: noise term

E.g., readout electronics

C: constant term

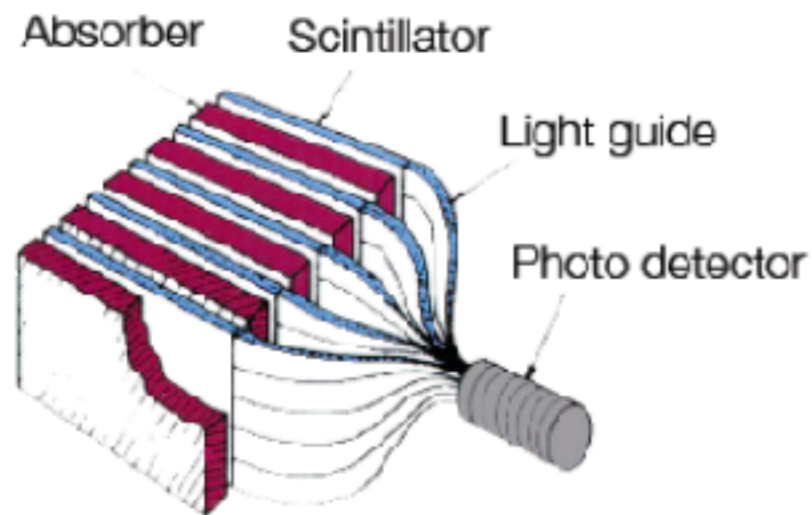
E.g., Non uniformity, calibration etc...

Table 34.8: Resolution of typical electromagnetic calorimeters. E is in GeV.

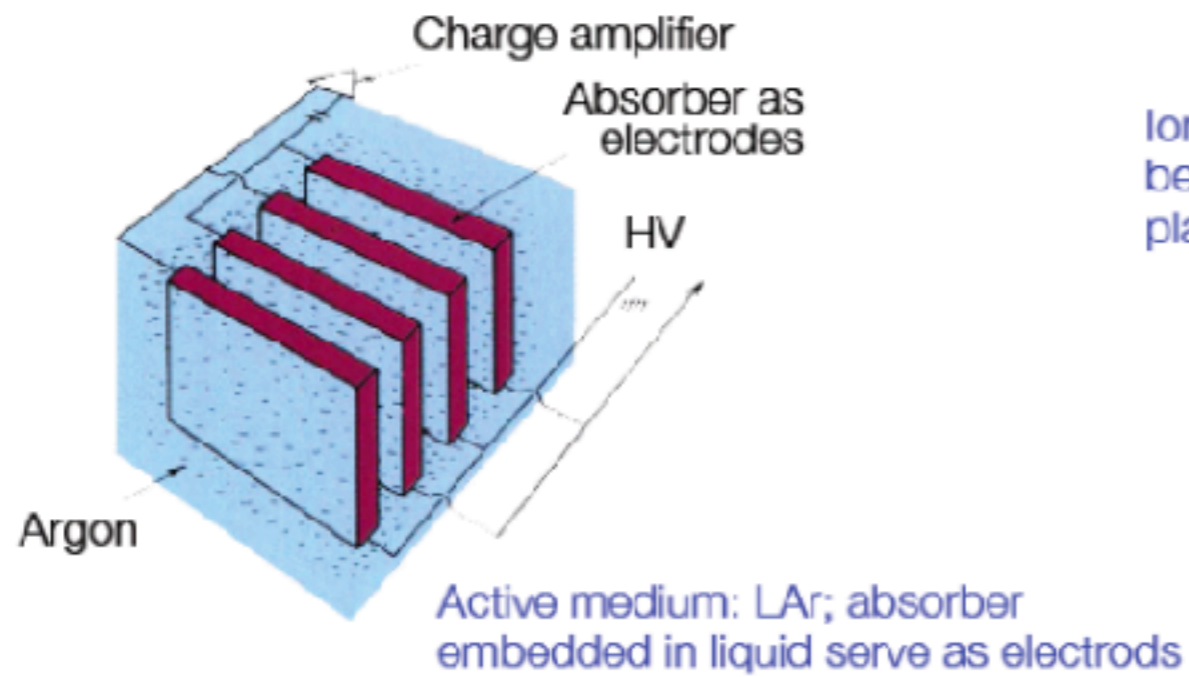
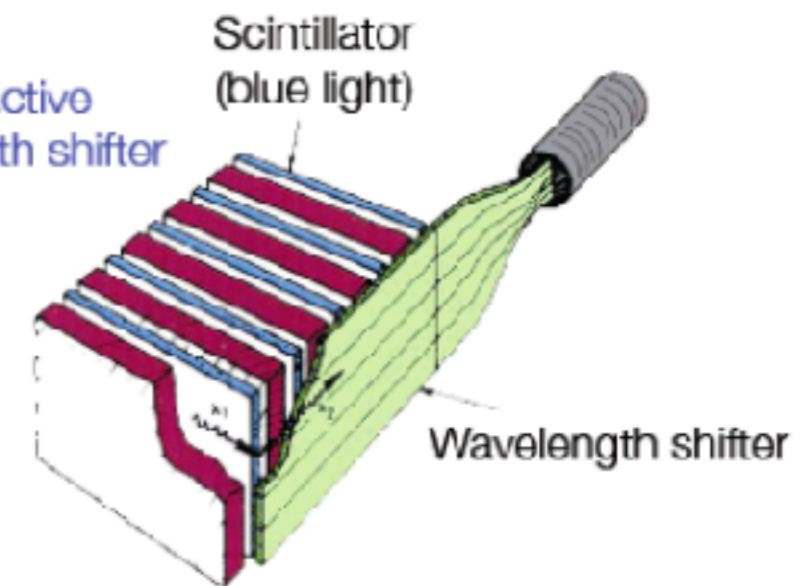
	Technology (Experiment)	Depth	Energy resolution	Date
Homogenous	NaI(Tl) (Crystal Ball)	$20X_0$	$2.7\%/E^{1/4}$	1983
	Bi ₄ Ge ₃ O ₁₂ (BGO) (L3)	$22X_0$	$2\%/\sqrt{E} \oplus 0.7\%$	1993
	CsI (KTeV)	$27X_0$	$2\%/\sqrt{E} \oplus 0.45\%$	1996
	CsI(Tl) (BaBar)	$16\text{--}18X_0$	$2.3\%/E^{1/4} \oplus 1.4\%$	1999
	CsI(Tl) (BELLE)	$16X_0$	1.7% for $E_\gamma > 3.5$ GeV	1998
	CsI(Tl) (BES III)	$15X_0$	2.5% for $E_\gamma = 1$ GeV	2010
	PbWO ₄ (CMS)	$25X_0$	$3\%/\sqrt{E} \oplus 0.5\% \oplus 0.2/E$	1997
	PbWO ₄ (ALICE)	$19X_0$	$3.6\%/\sqrt{E} \oplus 1.2\%$	2008
	Lead glass (OPAL)	$20.5X_0$	$5\%/\sqrt{E}$	1990
	Liquid Kr (NA48)	$27X_0$	$3.2\%/\sqrt{E} \oplus 0.42\% \oplus 0.09/E$	1998
Sampling	Scintillator/depleted U (ZEUS)	$20\text{--}30X_0$	$18\%/\sqrt{E}$	1988
	Scintillator/Pb (CDF)	$18X_0$	$13.5\%/\sqrt{E}$	1988
	Scintillator fiber/Pb spaghetti (KLOE)	$15X_0$	$5.7\%/\sqrt{E} \oplus 0.6\%$	1995
	Liquid Ar/Pb (NA31)	$27X_0$	$7.5\%/\sqrt{E} \oplus 0.5\% \oplus 0.1/E$	1988
	Liquid Ar/Pb (SLD)	$21X_0$	$8\%/\sqrt{E}$	1993
	Liquid Ar/Pb (H1)	$20\text{--}30X_0$	$12\%/\sqrt{E} \oplus 1\%$	1998
	Liquid Ar/depl. U (DØ)	$20.5X_0$	$16\%/\sqrt{E} \oplus 0.3\% \oplus 0.3/E$	1993
	Liquid Ar/Pb accordion (ATLAS)	$25X_0$	$10\%/\sqrt{E} \oplus 0.4\% \oplus 0.3/E$	1996

Sampling calorimeter designs

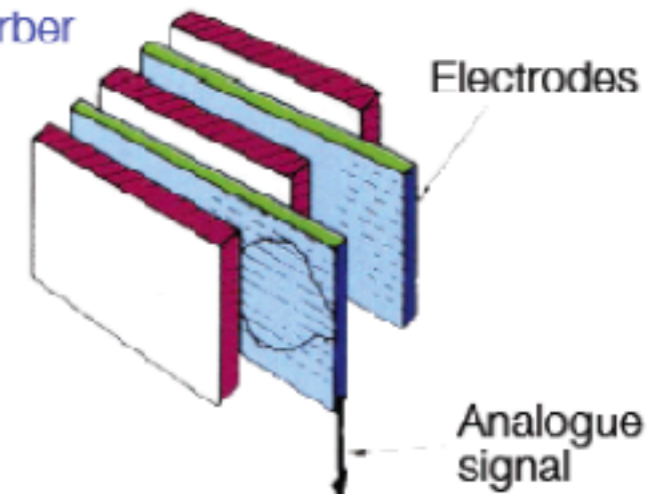
Scintillators as active layer;
signal readout via photo multipliers



Scintillators as active layer; wave length shifter to convert light



Ionization chambers between absorber plates



LHCb ECAL

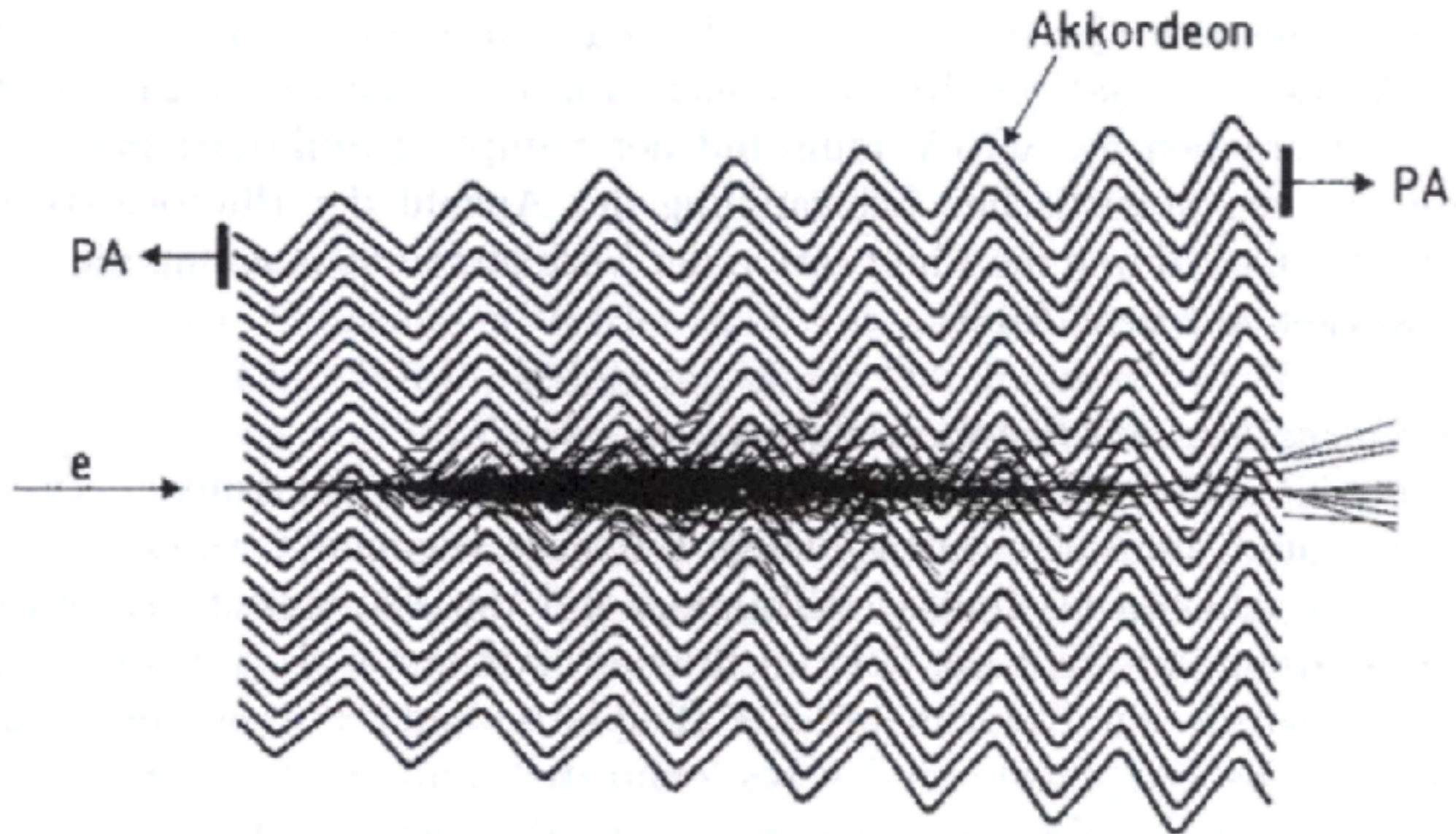


“Shashlik” design.

Alternating Pb absorber and scintillator.

Energy resolution sampling term of approximately $10\%/\sqrt{E/\text{GeV}}$

ATLAS LAr ECAL

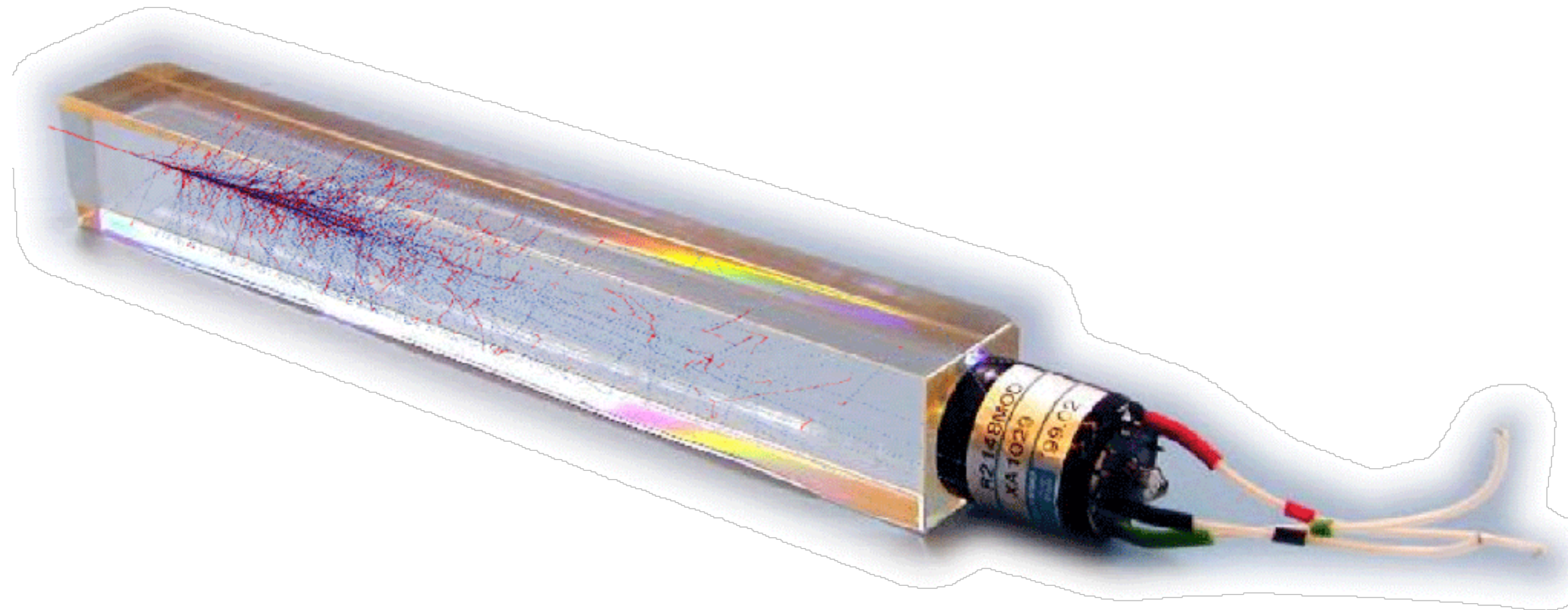


Homogenous calorimeters

Scintillating crystals (sodium iodide NaI, bismuth germanate BGO, caesium iodide CsI, lead tungstate PbWO_4 , etc.)

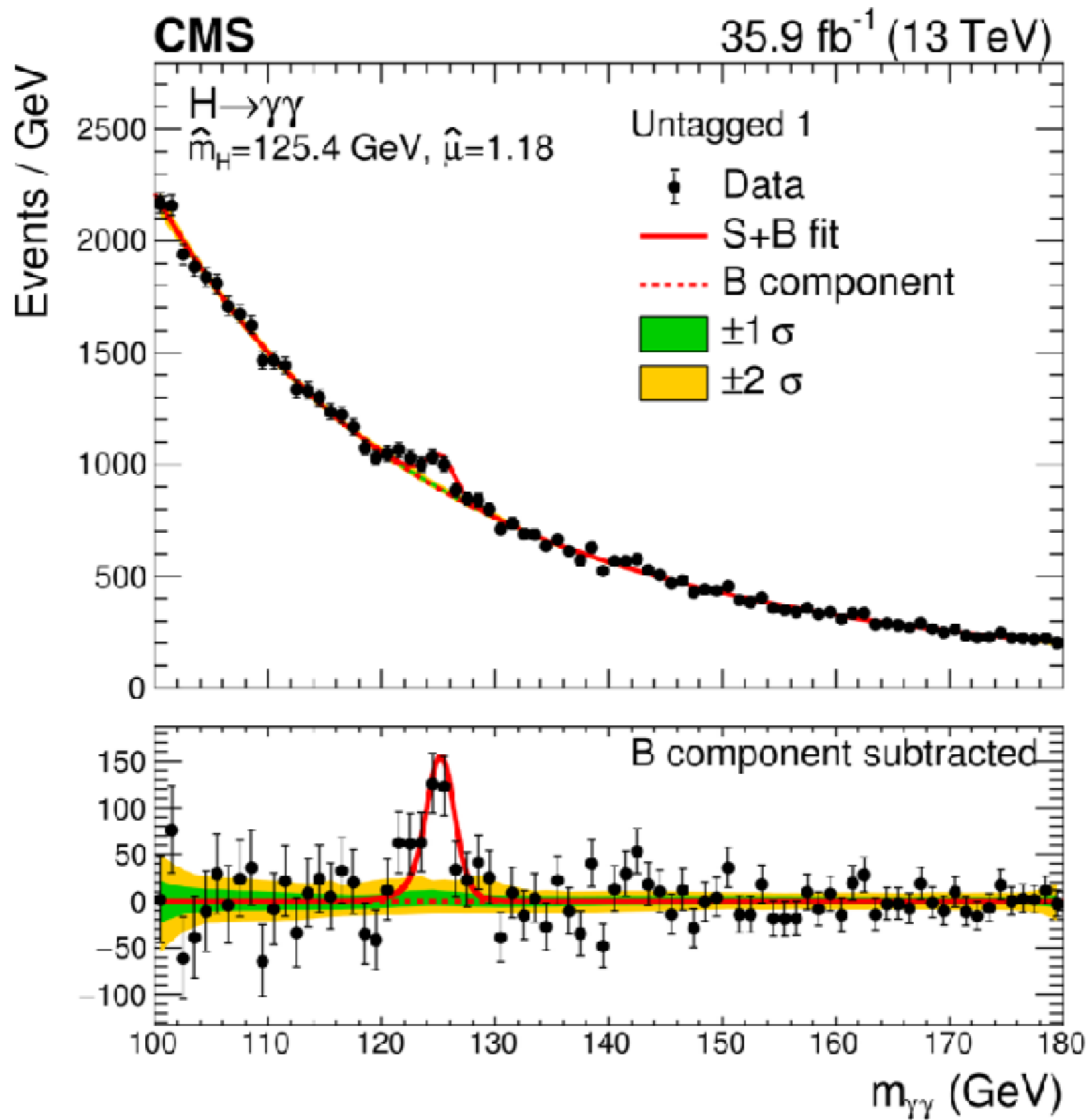
	NaI(Tl)	BGO	CsI(Tl)	PbWO_4
density (g/cm^3)	3.67	7.13	4.53	8.28
X_0 (cm)	2.59	1.12	1.85	0.89
R_M (cm)	4.5	2.4	3.8	2.2
dE/dx_{mip} (MeV/cm)	4.8	9.2	5.6	13.0
light yield (photons/MeV)	$4 \cdot 10^4$	$8 \cdot 10^3$	$5 \cdot 10^4$	$3 \cdot 10^2$
energy resolution σ_E/E	$1\%/\sqrt{E}$	$1\%/\sqrt{E}$	$1.3\%/\sqrt{E}$	$2.5\%/\sqrt{E}$

CMS ECAL

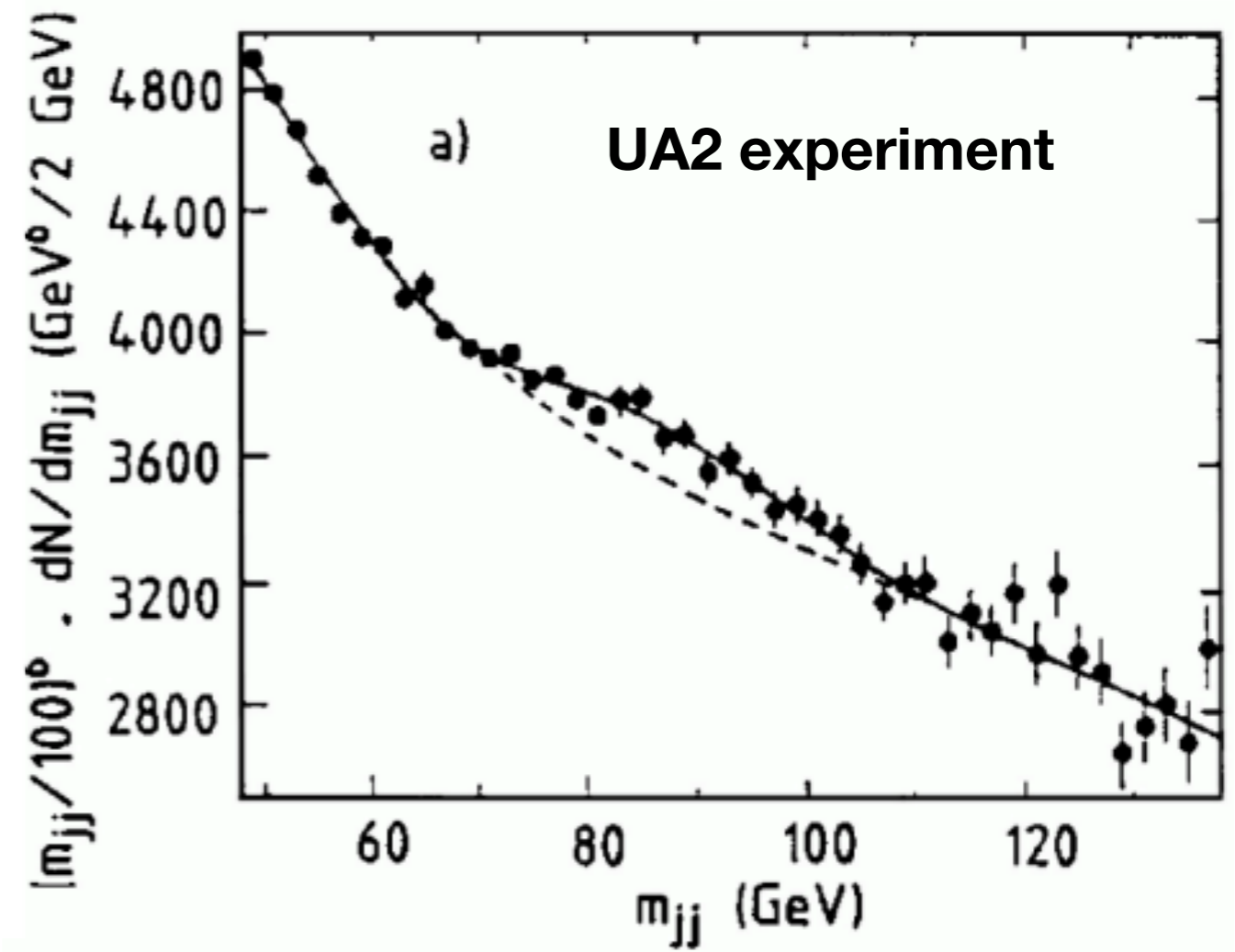
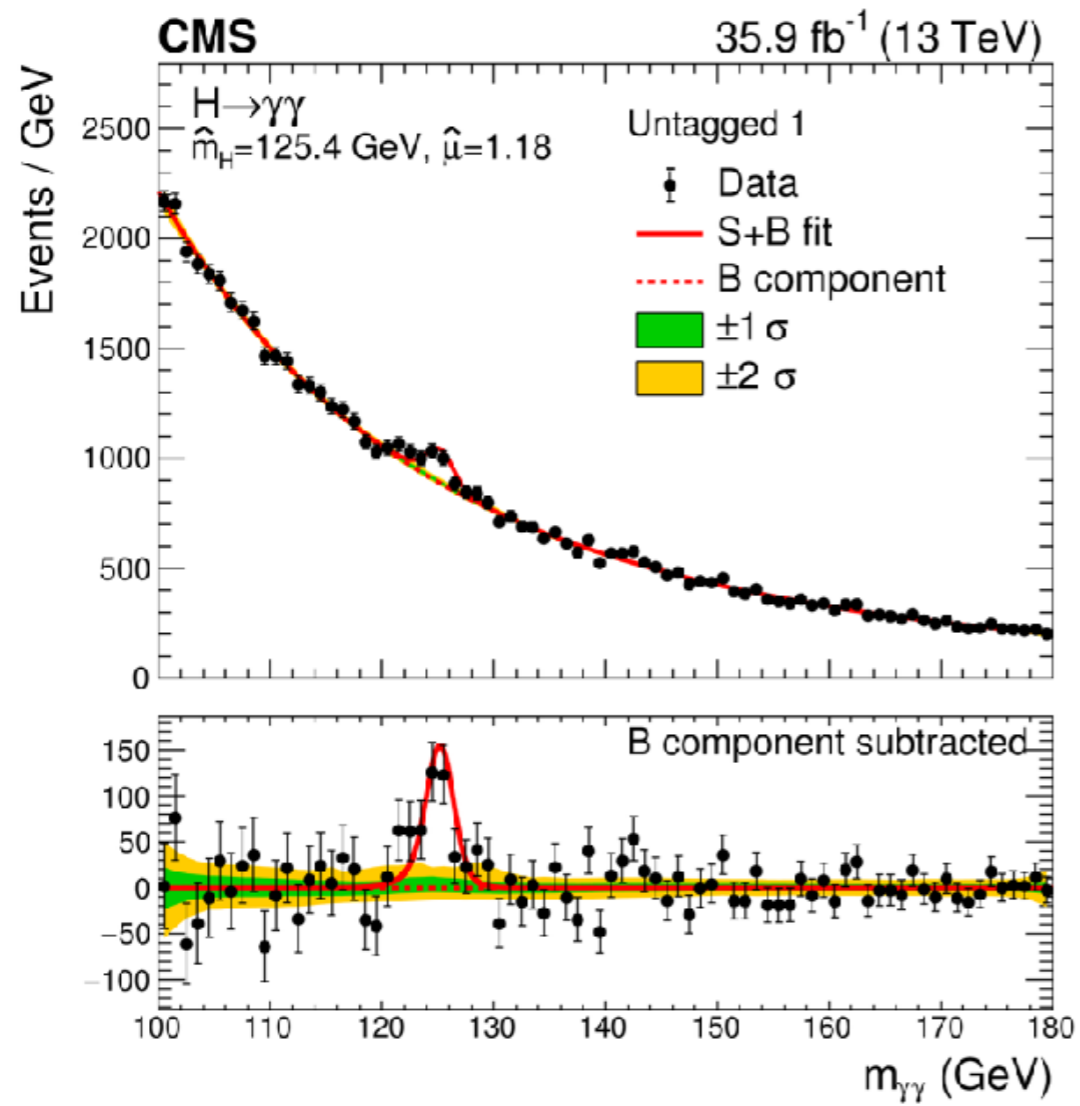


$$\frac{\sigma}{E} = \frac{2.8\%}{\sqrt{E(\text{GeV})}} \oplus \frac{125}{E(\text{MeV})} \oplus 0.3\%$$

CMS ECAL

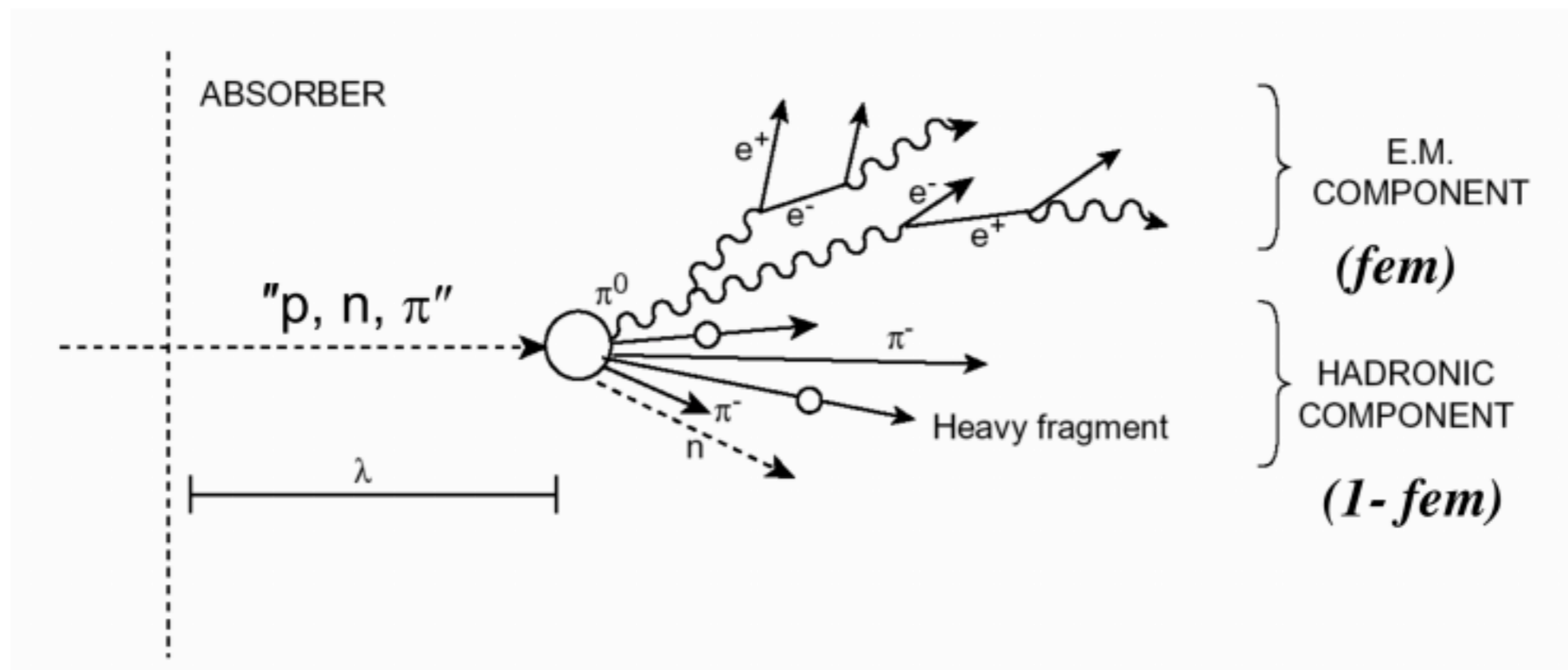


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3. EM calorimeters
4. Hadronic showers/calorimeters



Hadronic showers

- High energy hadrons interact with nuclei to produce secondary hadrons.
- Number of secondary hadrons $\sim \ln(E)$.
- Characteristic interaction length λ_I .
- Multiplication until $\langle E \rangle$ below [few $\times m_\pi$].
- Two distinct components



Detector response to EM and had components is different

Hadronic showers

EM component from π^0, η^0

Hadronic component

~20% Charged hadrons

~25% Nuclear fragments

~15% neutrons + soft photons

~40% nuclear breakup (invisible)

Nuclear interaction length and containment

- Nuclear interaction length

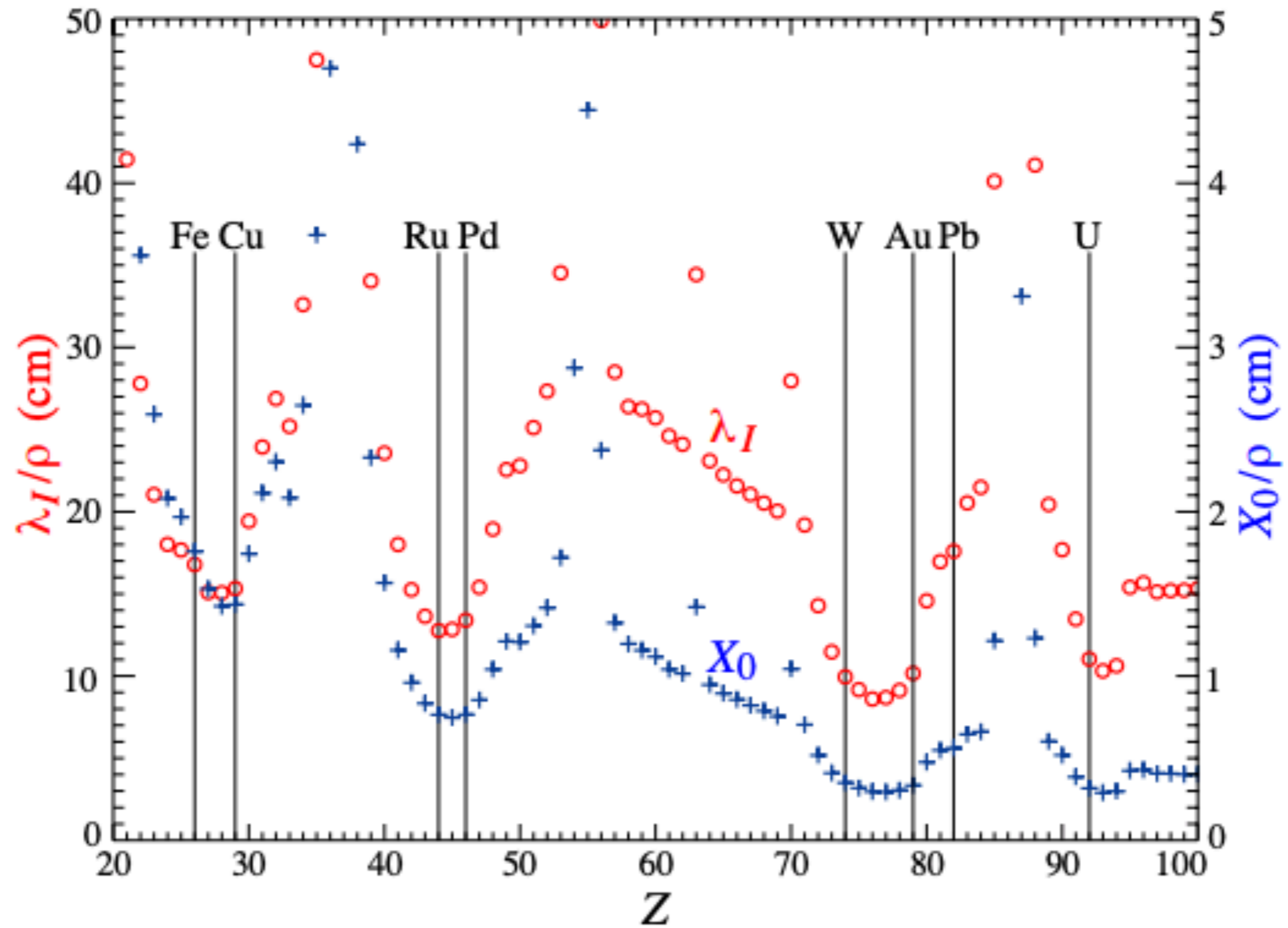
$$\lambda_i \approx \frac{1}{\sigma_{\text{tot}}} \frac{A}{N_A} \sim \left[35 \frac{\text{g}}{\text{cm}^2} \right] A^{1/3}$$

- Typically order of magnitude larger than X_0
- Typically require about 10λ for containment
- Hadronic calorimeters are always of the sampling type

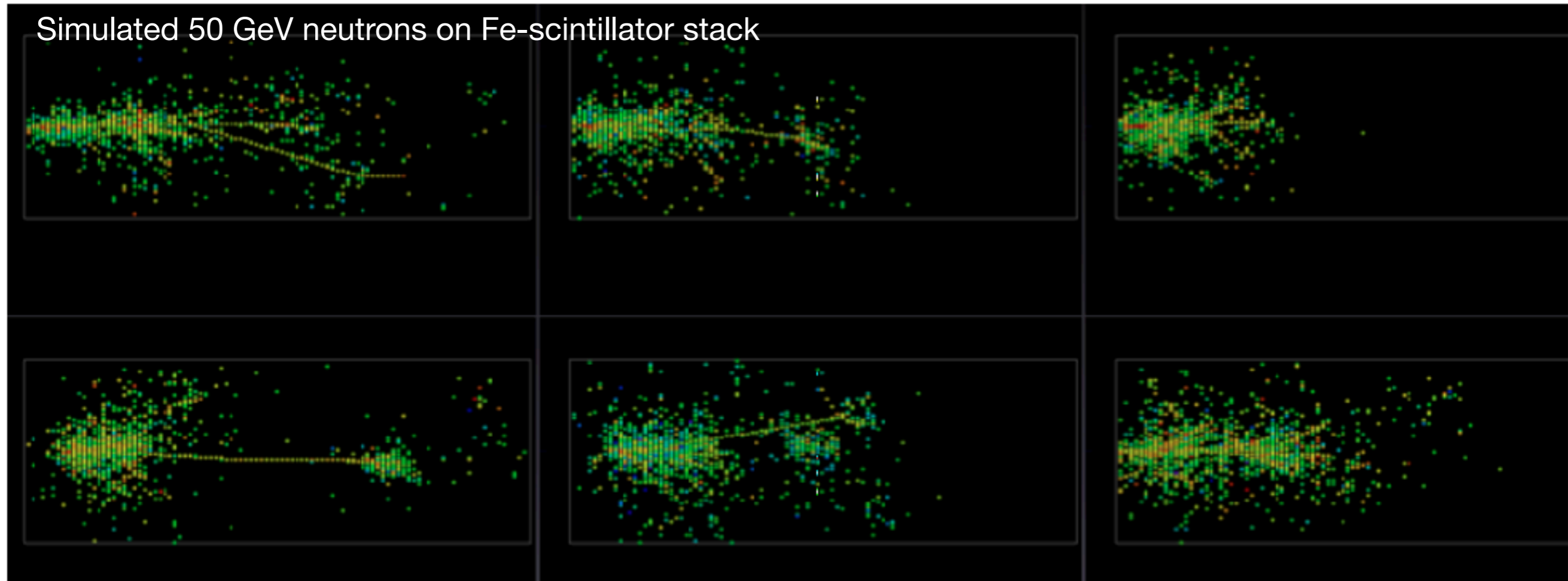
$$X_0 \sim \left[180 \frac{\text{g}}{\text{cm}^2} \right] \frac{A}{Z^2}$$

Material dependence

rpp2019-rev-particle-detectors-accel



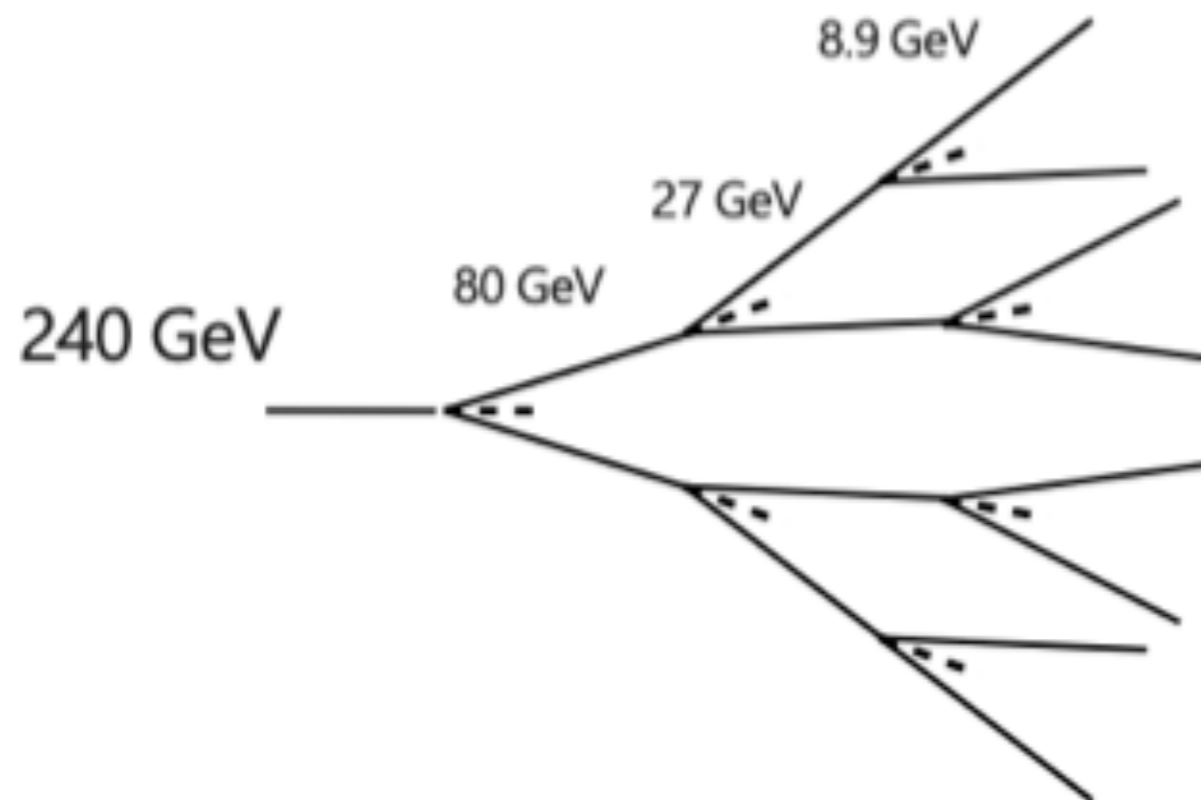
Hadron shower characteristics



- Large fluctuations in shape/profile!
- Charge hadrons propagate shower over large scale (λ)
- Local EM showers from π^0 , η^0

Simple hadronic shower model

- Shower is series of interactions producing, on average, $1/3 \pi^0$ and $2/3 \pi^\pm$
- Shower stops when $\langle E \rangle < 3m_\pi$



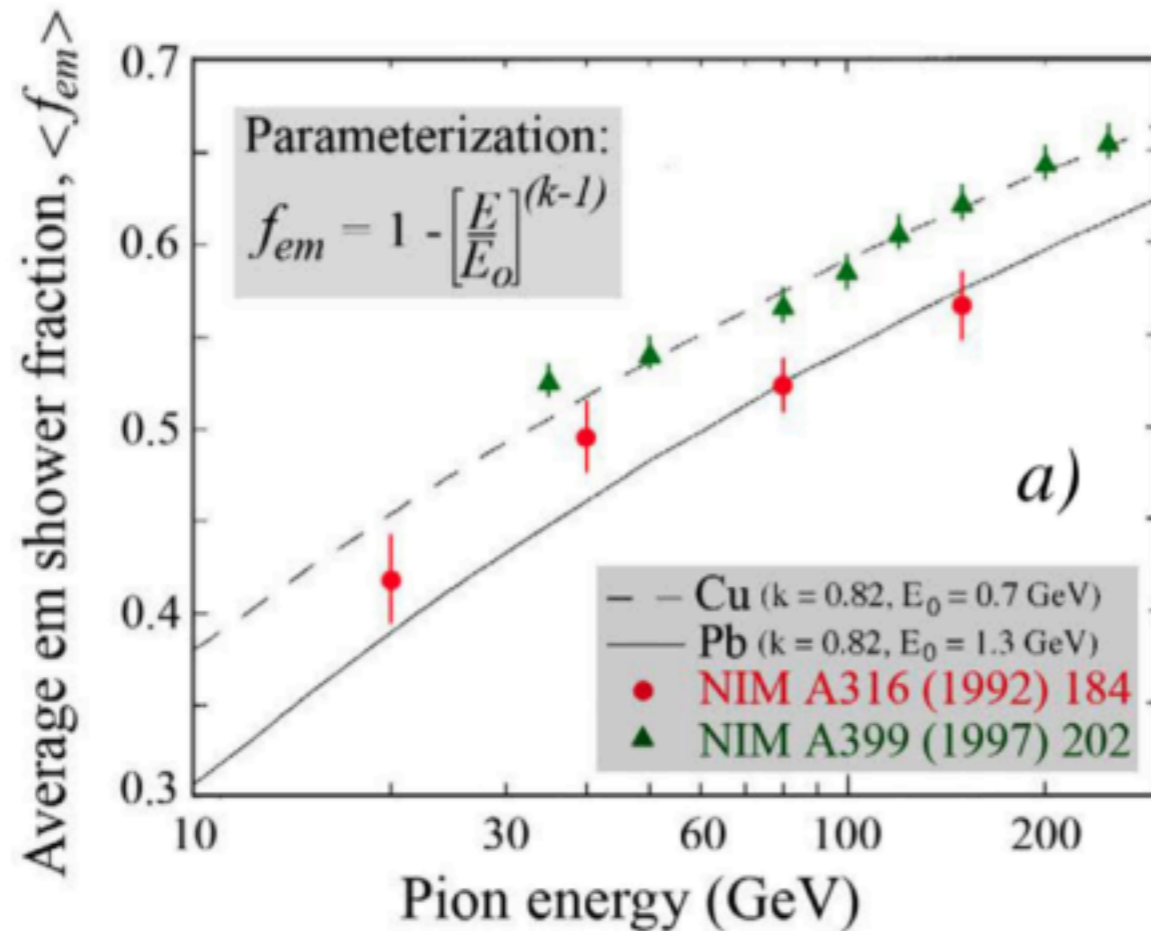
$$f_{\text{em}} = 1 - \left(\frac{2}{3}\right)^n$$

... f_{EM} increases with energy

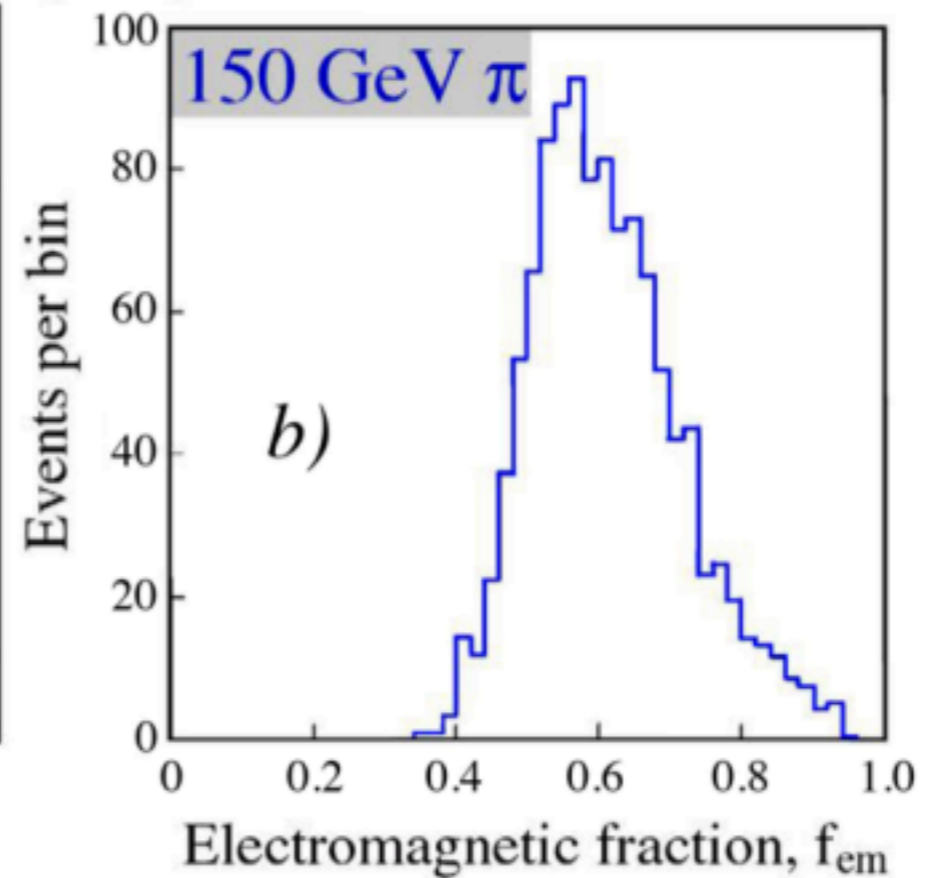
$$f_{\text{em}} = 1 - \left(\frac{E}{E_{\text{th}}}\right)^{k-1}$$

Challenge of hadron calorimeters

Energy dependence of f_{EM}



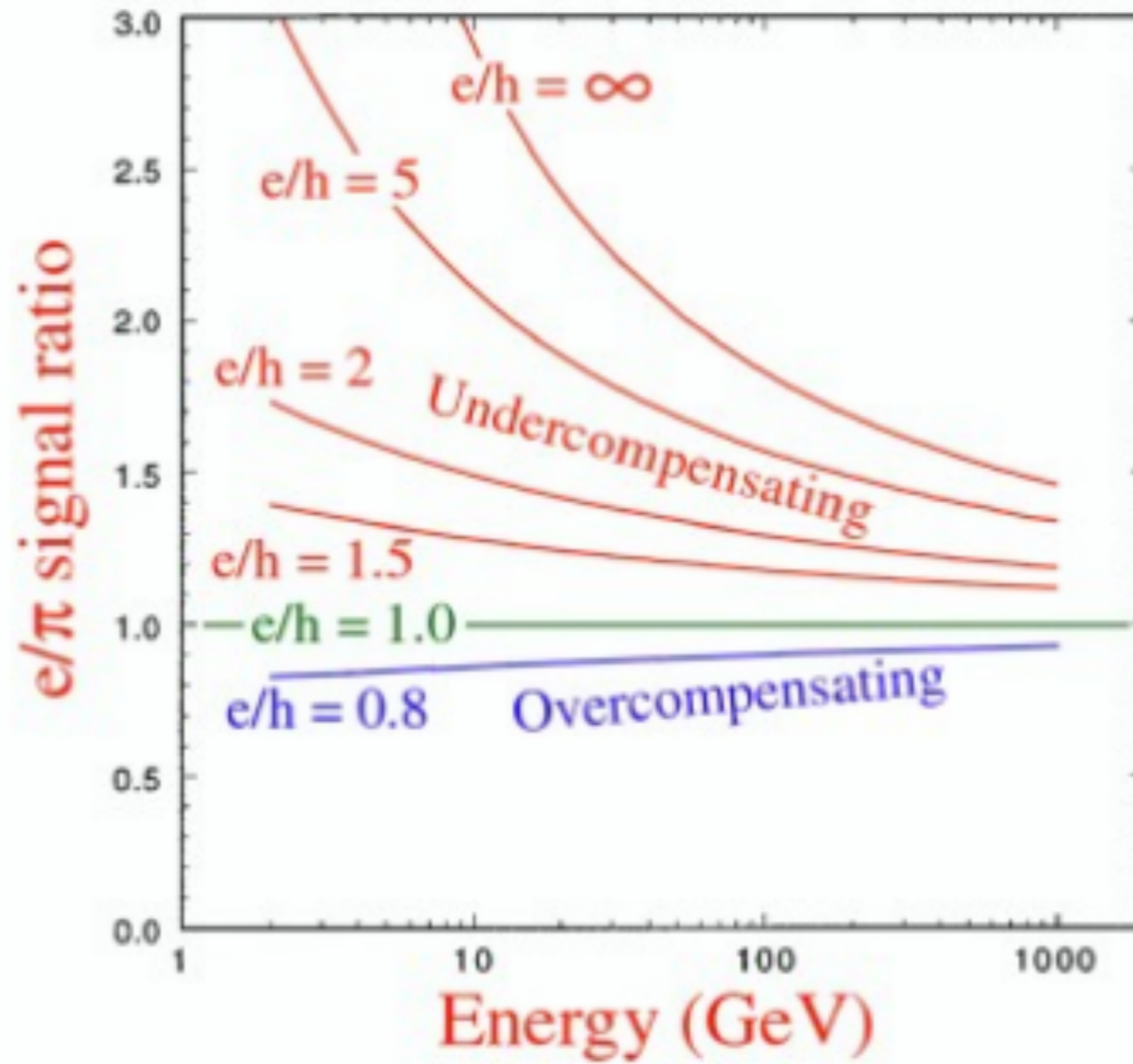
Fluctuations in f_{EM}



And the calorimeter response to the hadronic component tends to be smaller than to the electromagnetic component.

The response to hadrons is energy dependent and fluctuates a lot.

Challenge of hadron calorimeters



Compensation methods [for $e/h \approx 1$]

1. Software based

Pattern recognition and reweighing.

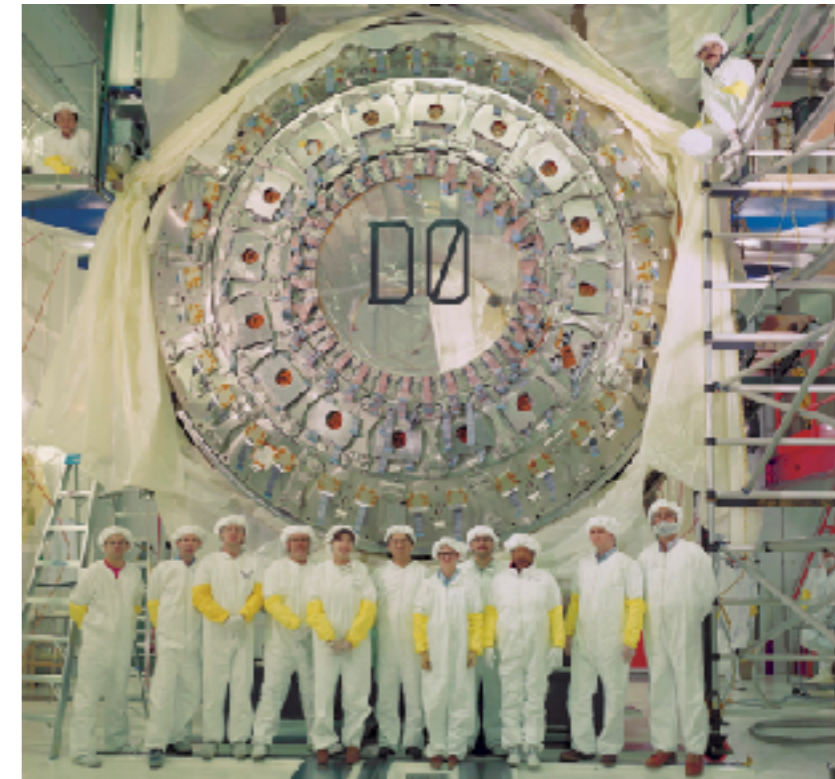
2. Reduce EM component

High Z material to filter out photo-electrons.

3. Boost the hadronic response

Organic (hydrogen rich) materials with high neutron cross section.

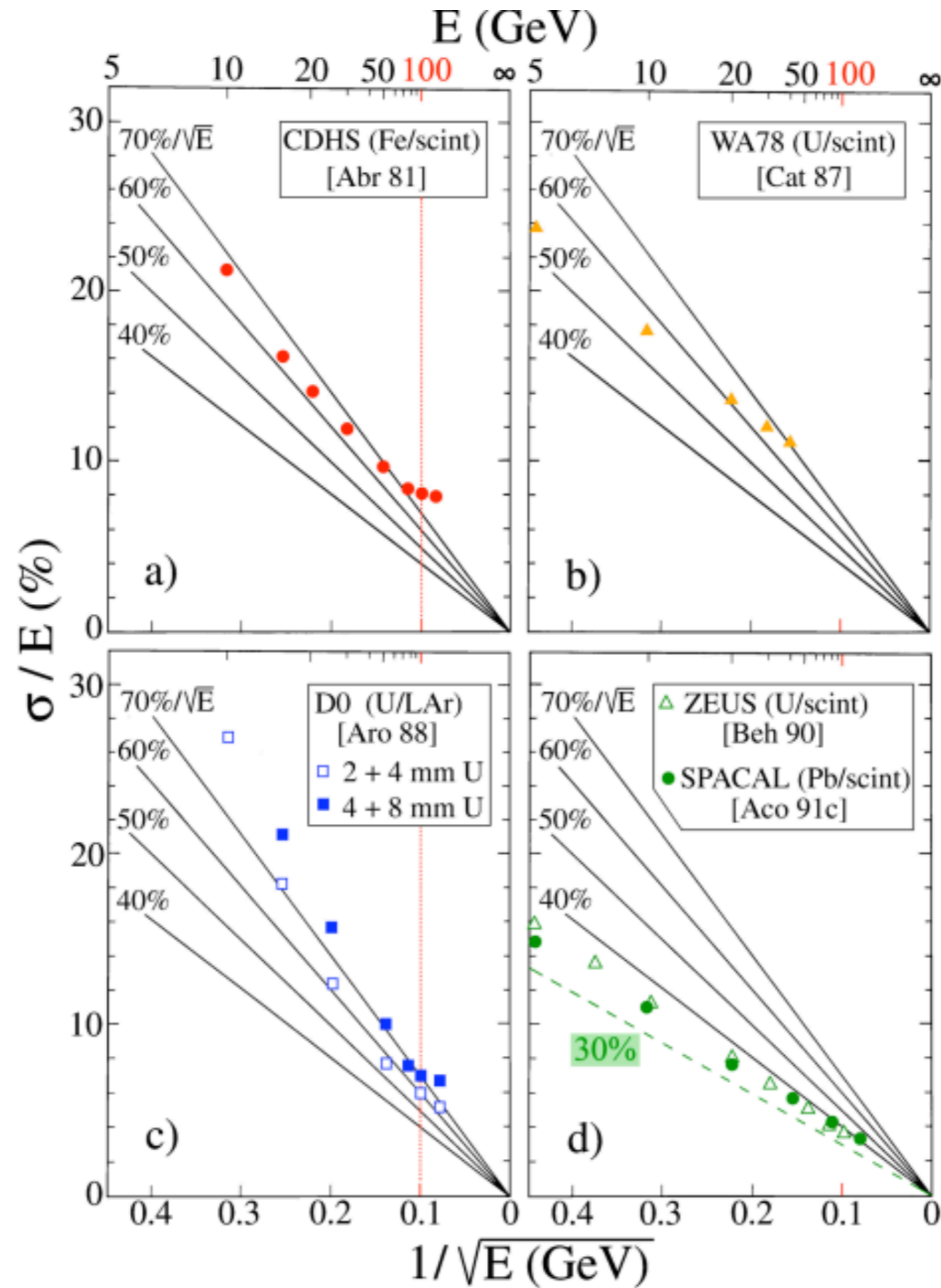
Uranium (nuclear fission triggered by neutrons).



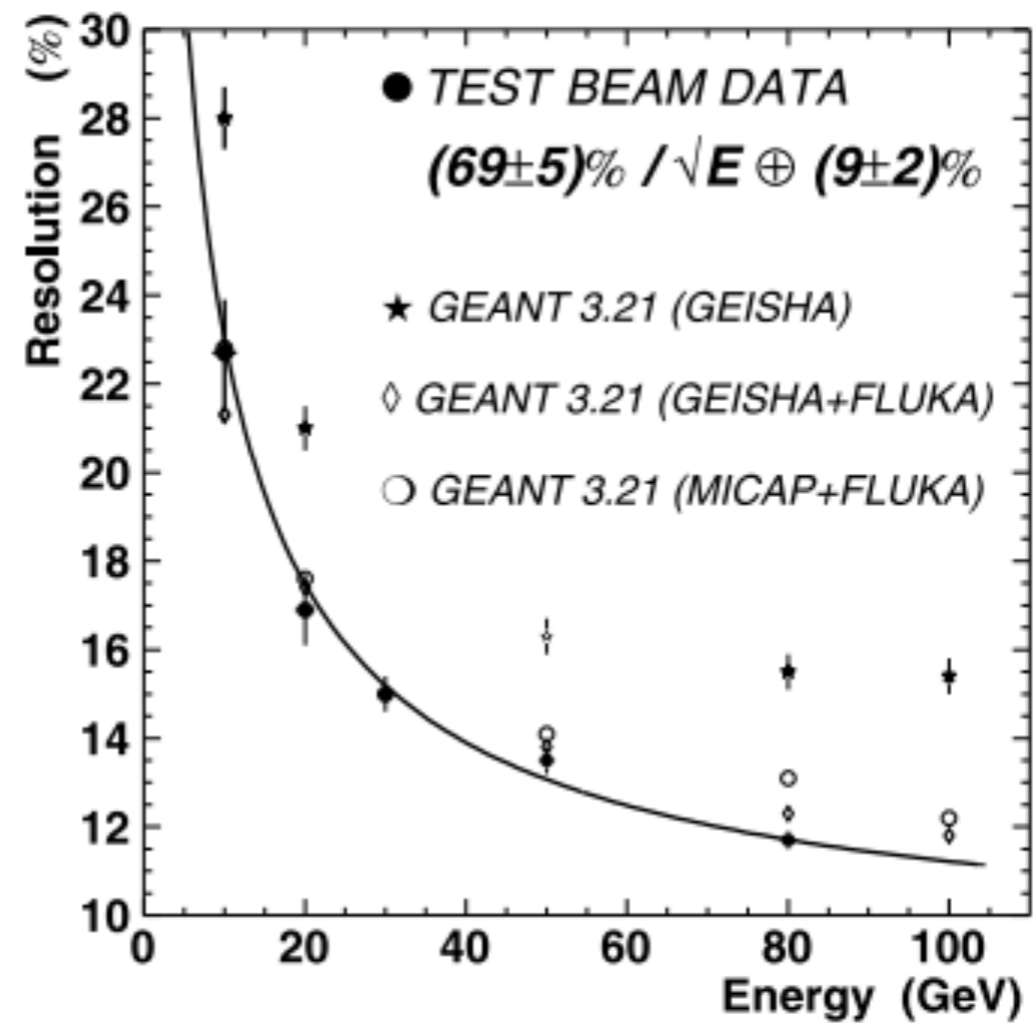
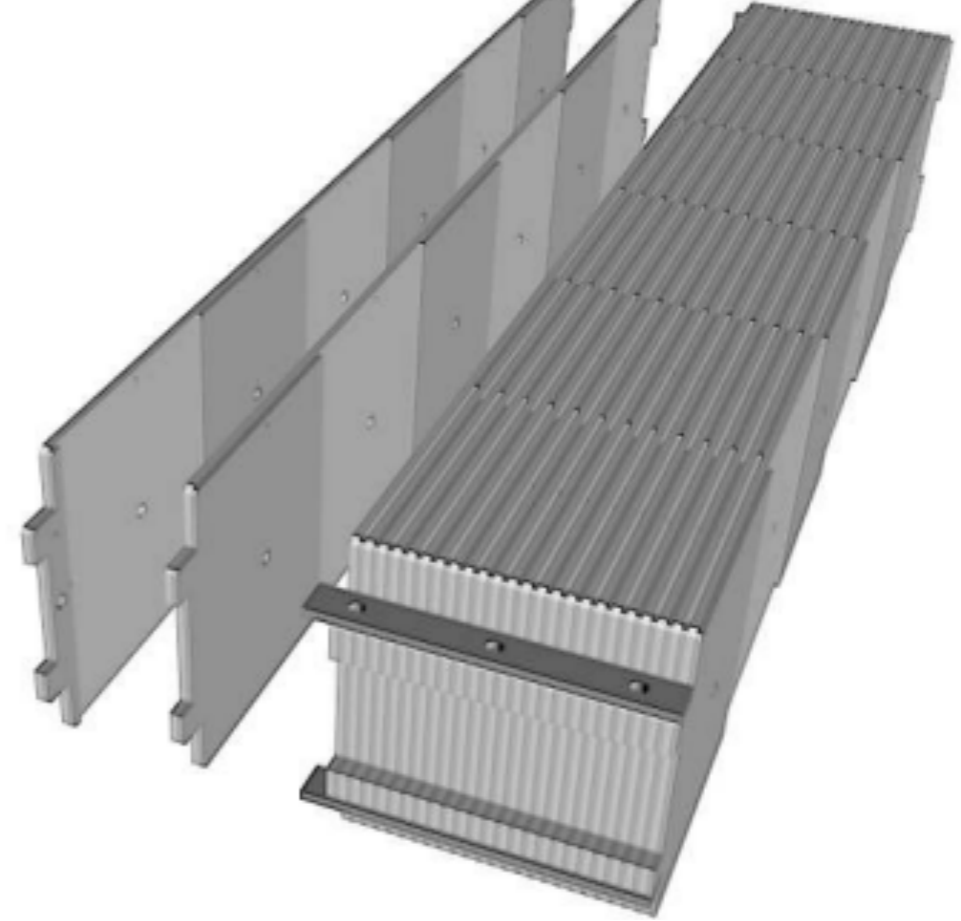
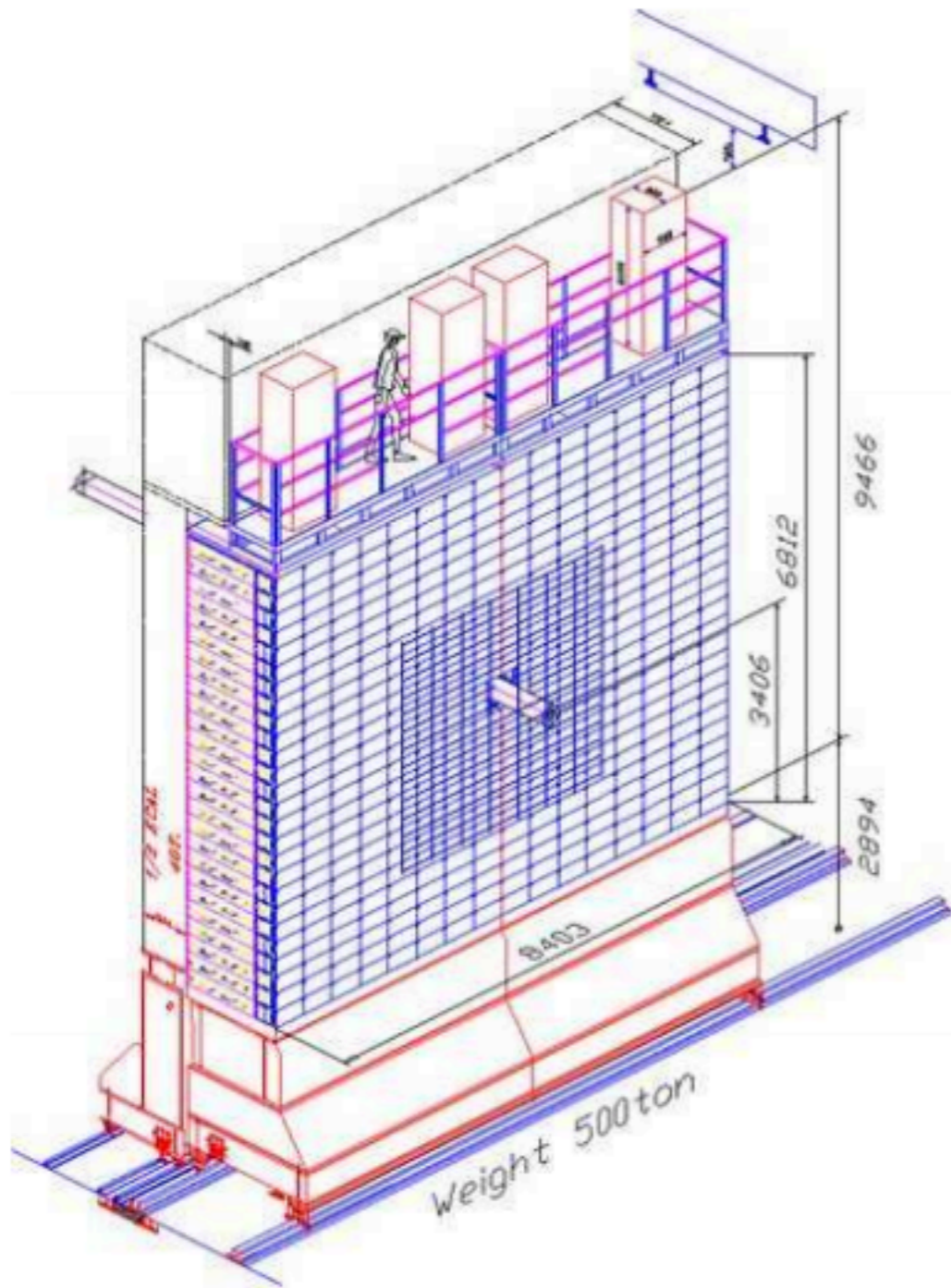
D0 HCAL with U absorber

Not compatible with good EM resolution!

Some example performances

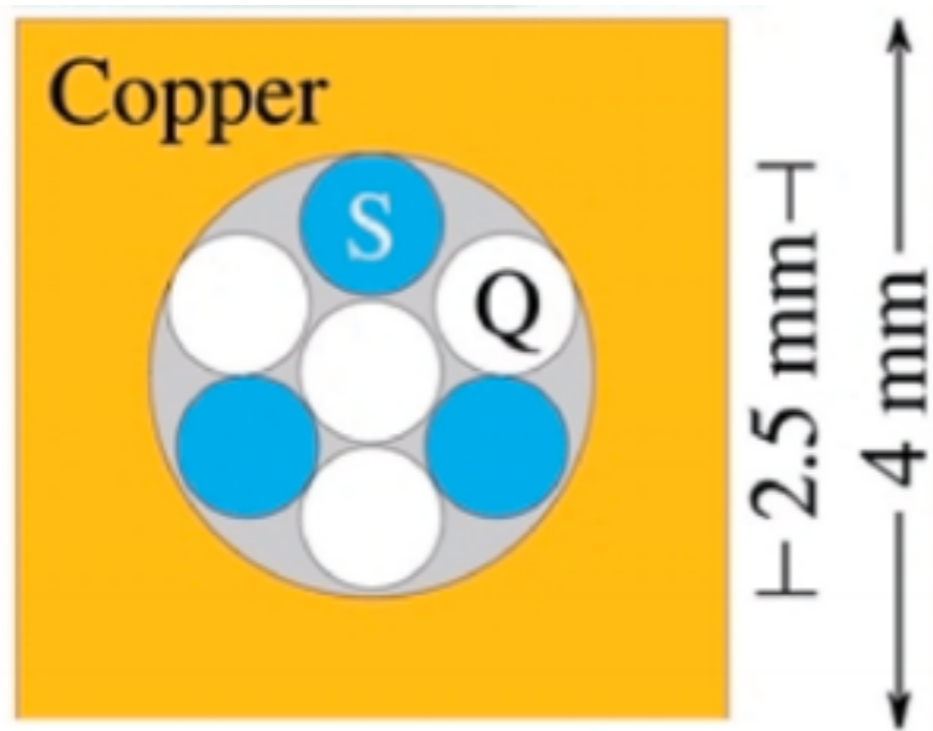


LHCb HCAL

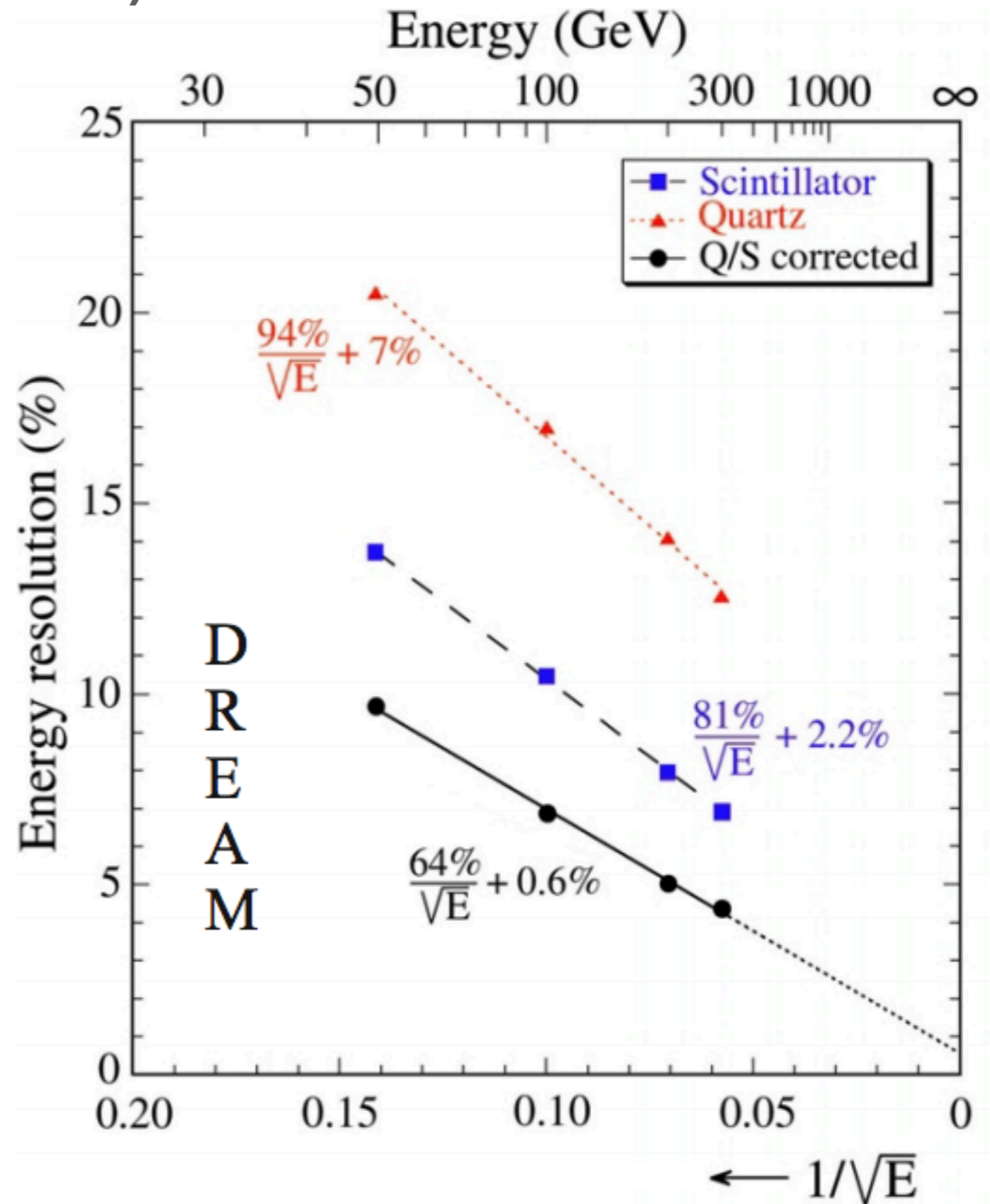


Dual readout R&D (compensation)

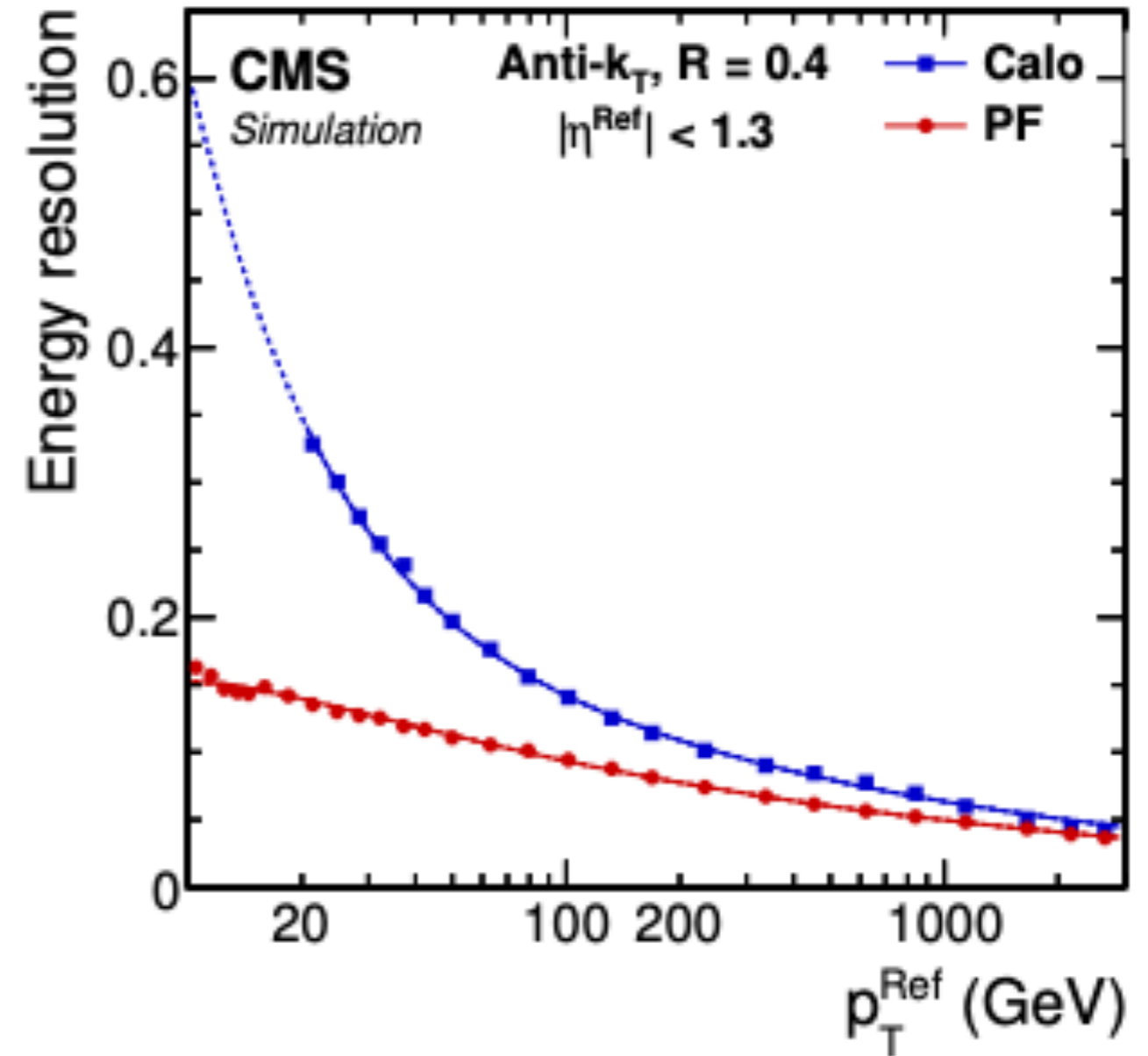
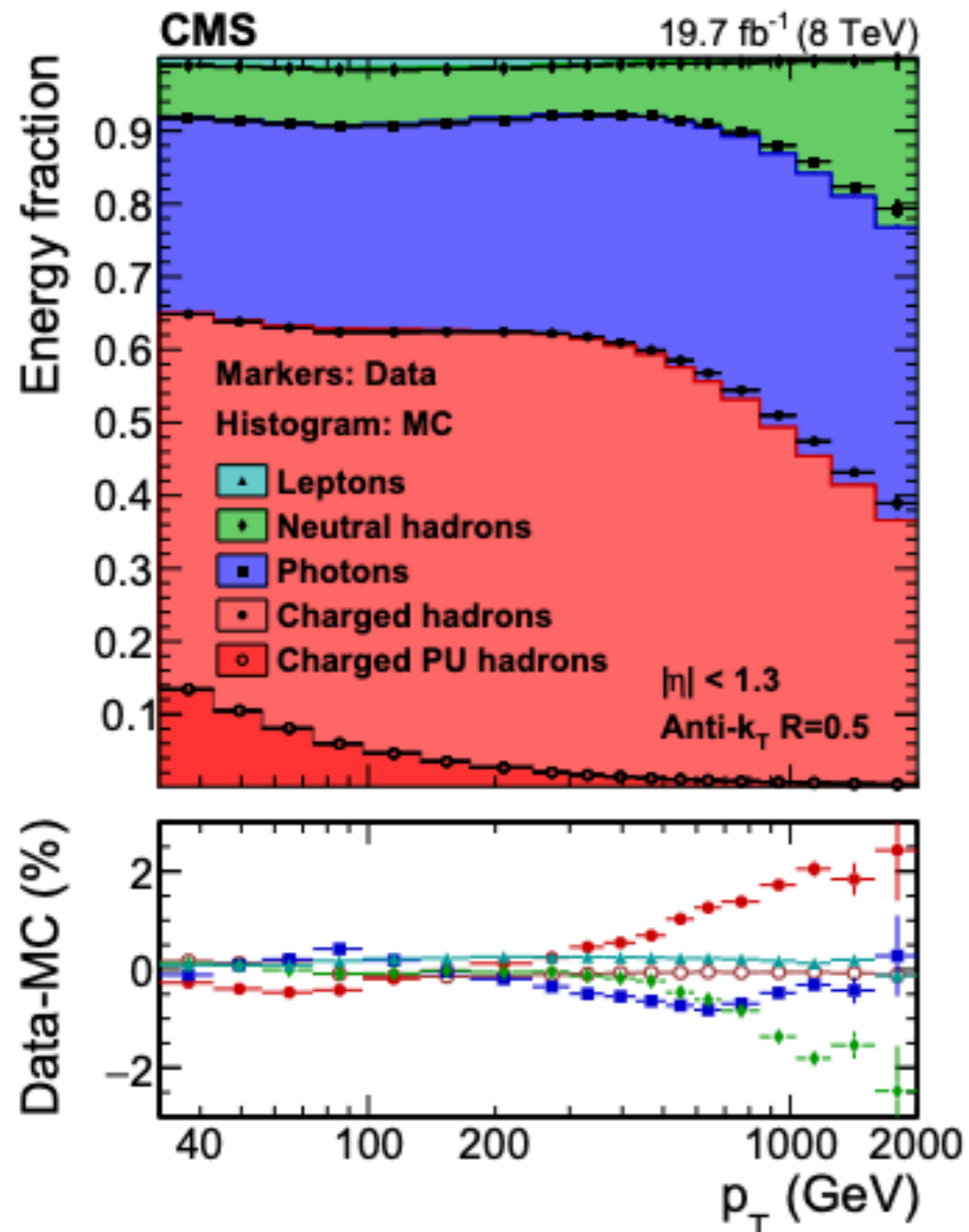
E.g. the original DREAM prototype with scintillating fibres and quartz fibres that have different e/h



Area of ongoing R&D activity



Particle flow



Useful references

Book

Calorimetry: Energy Measurement in Particle Physics, Richard Wigmans

PDG reviews

<http://pdg.lbl.gov/2019/reviews/rpp2019-rev-passage-particles-matter.pdf>

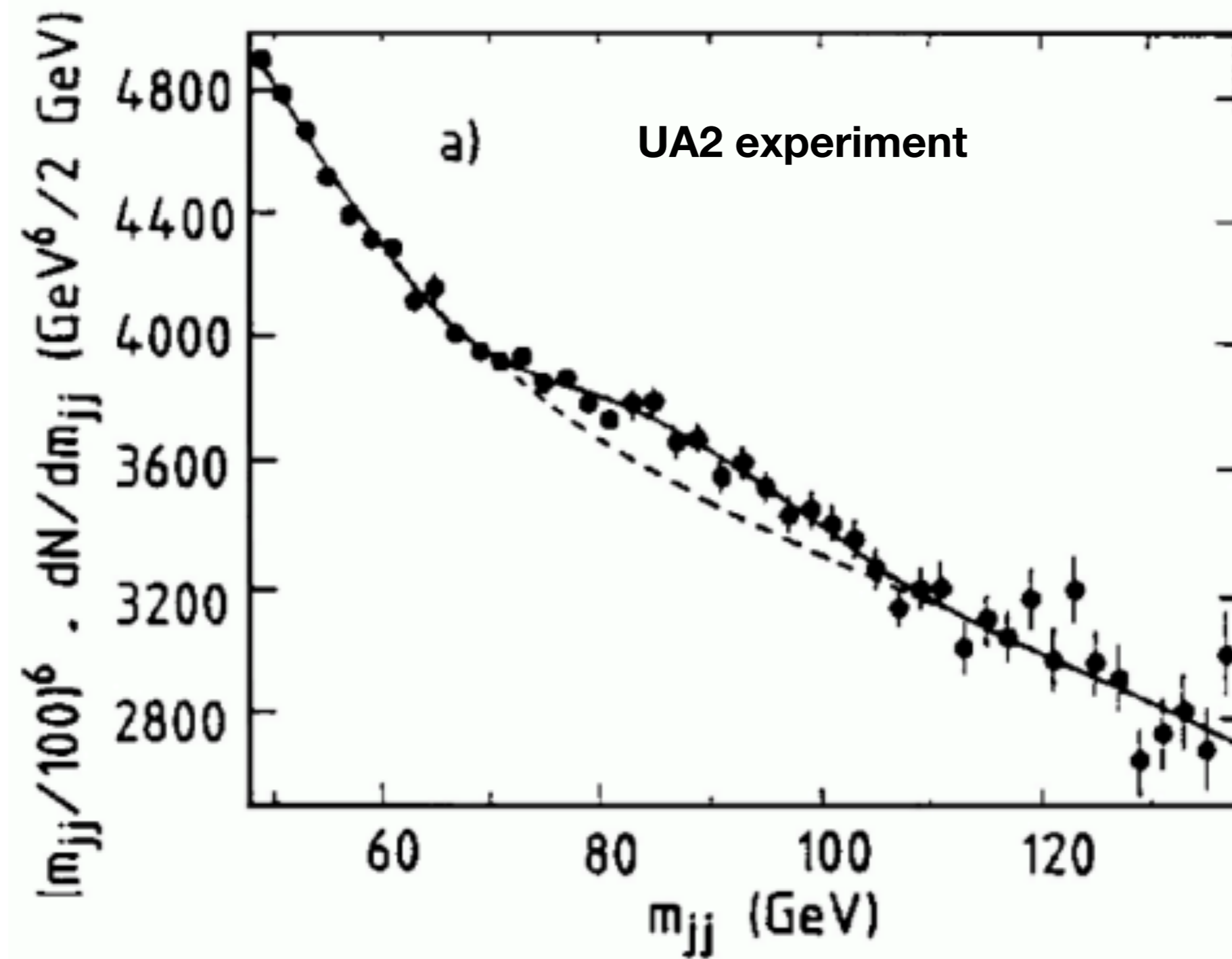
<https://pdg.lbl.gov/2019/reviews/rpp2019-rev-particle-detectors-accel.pdf>

Animated gifs of shower simulations

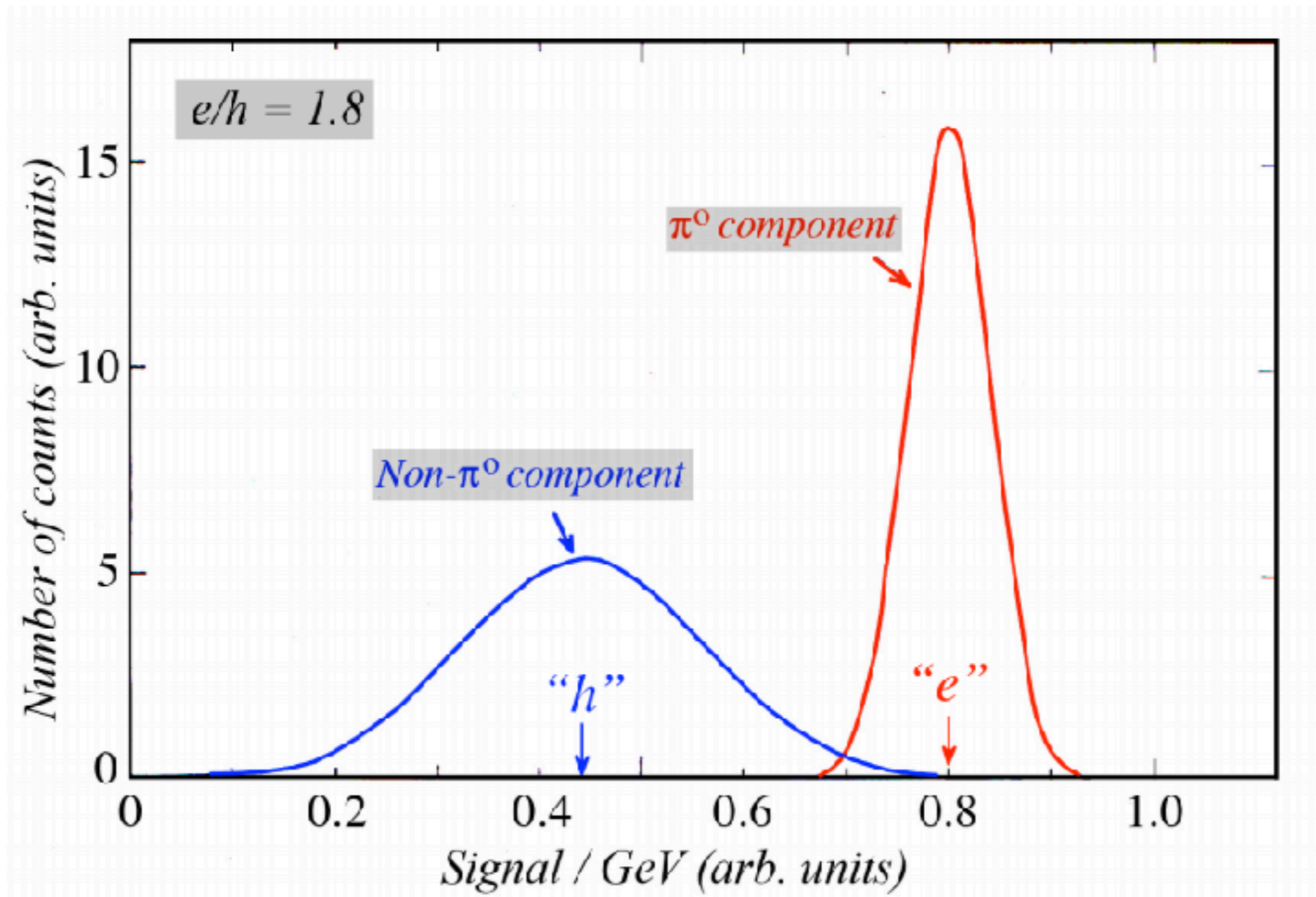
<https://www.mpp.mpg.de/~menke/elss/home.shtml>

Backup slides

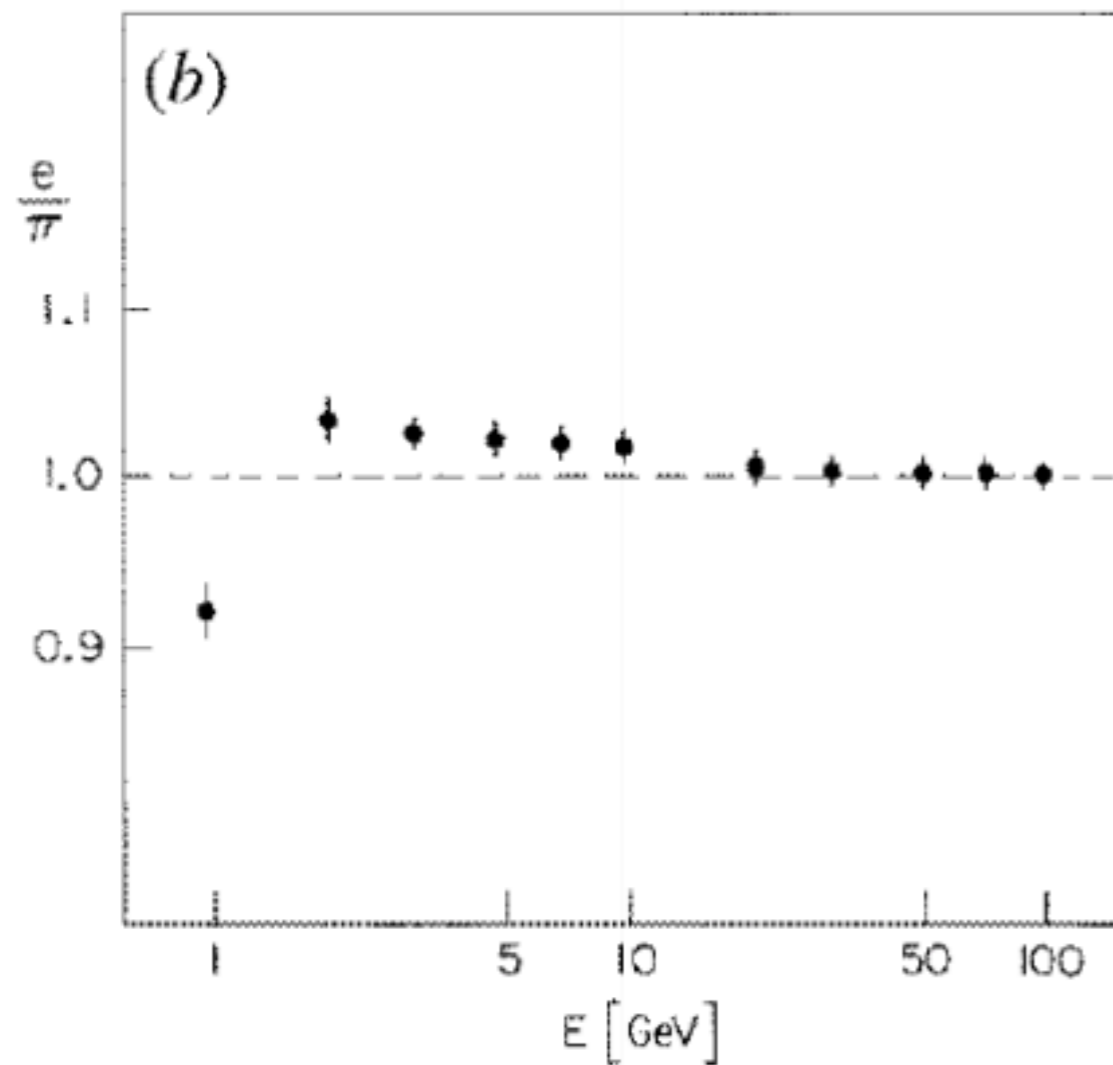
Hadronic showers



e/h example



ZEUS HCAL



ZEUS

- U/scintillator or Pb/scintillator?
- If U, then 1:1 absorber/scintillator ratio is compensating, if Pb then a 4:1 ratio is required
- The *intrinsic* fluctuations in a Pb sampling calorimeter are smaller than those for a U calorimeter
 - Pb: 13% vs U: 20% for hadrons
 - Pb: 0.3% vs U: 2.2% for EM
- However the much poorer sampling ratio for Pb resulted in the choice of Uranium.

Lateral containment

- Secondary hadrons have p_T of a few hundred MeV
- Comparable to energy lost in 1λ
- Characteristic lateral extent of 1λ
- High energy showers have pronounced core with exponential halo.

Moliere radius

An infinite cylinder of radius R_M contains 90% of the energy.