Light Pseudoscalar and Axial Meson Spectroscopy via an AdS/QCD Modified Soft Wall Model

Santiago Cortés, Universidad Nacional de Colombia

In Collaboration with

Miguel Ángel Martín Contreras and Alfredo Vega, Universidad de Valparaiso, Chile



#### 1 Motivation

- 2 Holographic Setup
- O Pseudoscalar Mesons

#### Axial Mesons

#### 5 Results



æ

#### **QCD** perturbative regime: E > 1 GeV.

- In Non-Perturbative QCD ( $E \leq 1$  GeV):
  - $\alpha_{QCD} \geq 1.$
  - Strongly bounded quarks.
    - Hadron states with light flavors (u, d and s).
      - Mass generation and chiral symmetry breaking (χ58).

#### Possible solutions:

- Effective Field Theories (χPT) and Effective Models (NJL, -N-LSM).
   Nonconformal AdS/QCD holographic approaches.
  - AdS/CET-like formulation without scale invariance.
  - ♦ Nonperturbative FT → Perturbative gravitational approach.

イロト 不得 トイヨト イヨト

- **QCD** perturbative regime: E > 1 GeV.
- **W** Non-Perturbative QCD ( $E \leq 1$  GeV):
  - $\alpha_{QCD} \geq 1$ .
  - Strongly bounded quarks.
    - Hadron states with light flavors (*u*, *d* and *s*).
    - Mass generation and chiral symmetry breaking (χSB).
- D Possible solutions:
  - Effective Field Theories (χPT) and Effective Models (NJL, -N-LSM).
     Nonconformal AdS/QCD holographic approaches.
    - AdS/CET-like formulation without scale invariance.
    - $\sim$  Nonperturbative FT  $\rightarrow$  Perturbative gravitational approach.

- **QCD** perturbative regime: E > 1 GeV.
- **W** Non-Perturbative QCD ( $E \leq 1$  GeV):
  - $\alpha_{QCD} \geq 1.$
  - Strongly bounded quarks.
    - Hadron states with light flavors (*u*, *d* and *s*).
    - Mass generation and chiral symmetry breaking ( $\chi$ SB).
- Possible solutions:
  - Effective Field Theories (χPT) and Effective Models (NJL, -N-LSM).
     Nonconformal AdS/QCD holographic approaches.
    - AdS/CET-like formulation without scale invariance.
    - $\mathbb{P} := \mathbb{P} = \mathbb{P}$  Perturbative gravitational approach.

- **QCD** perturbative regime: E > 1 GeV.
- **W** Non-Perturbative QCD ( $E \leq 1$  GeV):
  - $\alpha_{\text{QCD}} \geq 1.$
  - Strongly bounded quarks.
    - Hadron states with light flavors (*u*, *d* and *s*).
    - Mass generation and chiral symmetry breaking ( $\chi$ SB).
- Possible solutions:
  - Effective Field Theories (χPT) and Effective Models (NJL, -N-LSM).
     Nonconformal AdS/QCD holographic approaches.
    - AdS/CET-like formulation without scale invariance.
    - Nonperturbative FT ightarrow Perturbative gravitational approach.

- **QCD** perturbative regime: E > 1 GeV.
- **W** Non-Perturbative QCD ( $E \leq 1$  GeV):
  - $\alpha_{QCD} \geq 1.$
  - Strongly bounded quarks.
    - Hadron states with light flavors (u, d and s).
    - Mass generation and chiral symmetry breaking ( $\chi$ SB).
- Possible solutions:
  - Effective Field Theories (χPT) and Effective Models (NJL, -N-LSM).
     Nonconformal AdS/QCD holographic approaches.
    - AdS/CET-like formulation without scale invariance.
    - $\mathbb{P} : \mathbb{P} \to \mathbb{P}$ erturbative gravitational approach.

- **QCD** perturbative regime: E > 1 GeV.
- **W** Non-Perturbative QCD ( $E \leq 1$  GeV):
  - $\alpha_{QCD} \geq 1.$
  - Strongly bounded quarks.
    - Hadron states with light flavors (u, d and s).
    - Mass generation and chiral symmetry breaking ( $\chi$ SB).

#### Possible solutions:

- Effective Field Theories (χPT) and Effective Models (NJL, -N-LSM).
   Nonconformal AdS/QCD holographic approaches
  - AdS/CFT-like formulation without scale invariance.
  - Nonperturbative  $FT \rightarrow Perturbative gravitational approach.structure <math>FT \rightarrow Perturbative gravitational$

- **QCD** perturbative regime: E > 1 GeV.
- **W** Non-Perturbative QCD ( $E \leq 1$  GeV):
  - $\alpha_{QCD} \geq 1.$
  - Strongly bounded quarks.
    - Hadron states with light flavors (*u*, *d* and *s*).
    - Mass generation and chiral symmetry breaking ( $\chi$ SB).

#### Possible solutions:

- Effective Field Theories ( $\chi$ PT) and Effective Models (NJL, -N-LSM).
- Nonconformal AdS/QCD holographic approaches.
  - AdS/CFT-like formulation without scale invariance.
  - Nonperturbative  $FT \rightarrow Perturbative gravitational approach.$

イロト 不得 トイヨト イヨト

- **QCD** perturbative regime: E > 1 GeV.
- **W** Non-Perturbative QCD ( $E \leq 1$  GeV):
  - $\alpha_{QCD} \geq 1.$
  - Strongly bounded quarks.
    - Hadron states with light flavors (*u*, *d* and *s*).
    - Mass generation and chiral symmetry breaking ( $\chi$ SB).
- Possible solutions:
  - Effective Field Theories ( $\chi$ PT) and Effective Models (NJL, -N-LSM).
  - Nonconformal AdS/QCD holographic approaches.
    - AdS/CFT-like formulation without scale invariance.
    - Nonperturbative FT → Perturbative gravitational approach.

イロト 不得 トイヨト イヨト

- **QCD** perturbative regime: E > 1 GeV.
- **W** Non-Perturbative QCD ( $E \leq 1$  GeV):
  - $\alpha_{QCD} \geq 1.$
  - Strongly bounded quarks.
    - Hadron states with light flavors (*u*, *d* and *s*).
    - Mass generation and chiral symmetry breaking ( $\chi$ SB).
- Possible solutions:
  - Effective Field Theories ( $\chi$ PT) and Effective Models (NJL, -N-LSM).
  - Nonconformal AdS/QCD holographic approaches.
    - AdS/CFT-like formulation without scale invariance.
    - Nonperturbative  $FT \rightarrow Perturbative gravitational approach.$

(日)

- **QCD** perturbative regime: E > 1 GeV.
- **W** Non-Perturbative QCD ( $E \leq 1$  GeV):
  - $\alpha_{QCD} \geq 1.$
  - Strongly bounded quarks.
    - Hadron states with light flavors (*u*, *d* and *s*).
    - Mass generation and chiral symmetry breaking ( $\chi$ SB).
- Possible solutions:
  - Effective Field Theories ( $\chi$ PT) and Effective Models (NJL, -N-LSM).
  - Nonconformal AdS/QCD holographic approaches.
    - AdS/CFT-like formulation without scale invariance.
    - Nonperturbative  $FT \rightarrow Perturbative gravitational approach.$

イロト 不得 トイヨト イヨト

- **QCD** perturbative regime: E > 1 GeV.
- **W** Non-Perturbative QCD ( $E \leq 1$  GeV):
  - $\alpha_{QCD} \geq 1.$
  - Strongly bounded quarks.
    - Hadron states with light flavors (*u*, *d* and *s*).
    - Mass generation and chiral symmetry breaking ( $\chi$ SB).
- Possible solutions:
  - Effective Field Theories ( $\chi$ PT) and Effective Models (NJL, -N-LSM).
  - Nonconformal AdS/QCD holographic approaches.
    - AdS/CFT-like formulation without scale invariance.
    - $\bullet~$  Nonperturbative FT  $\rightarrow~$  Perturbative gravitational approach.

イロト 不得 トイヨト イヨト

## Modified Soft-Wall Model

Model AdS Poincare patch with an extra UV cutoff:

$$dS^{2} = g_{MN} dx^{M} dx^{N} = \frac{R^{2}}{z^{2}} \left[ dz^{2} + \eta_{\mu\nu} dx^{\mu} dx^{\nu} \right] \Theta \left( z - z_{0} \right).$$
(1)

Associated action:

$$I = I_{\text{Scalar}} + I_{\text{Vector}},\tag{2}$$

#### ) wher

$$I_{\text{Scalar}} = -\frac{1}{2g_{S}^{2}} \int d^{5} x \sqrt{-g} e^{-\Phi(z)} \left[ g^{MN} \partial_{M} S \partial_{N} S + M_{5}^{2} S^{2} \right], \quad (3)$$

$$I_{\text{Vector}} = -\frac{1}{2g_{V}^{2}} \int d^{5} x \sqrt{-g} e^{-\Phi(z)} \left[ \frac{1}{2} F_{MN} F^{MN} + \tilde{M}_{5}^{2} g^{MN} A_{M} A_{N} \right], \quad (4)$$

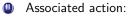
$$\Phi(z) = \kappa^{2} z^{2}. \quad (5)$$

Ref: Braga, Nelson R. F. et al, Phys. Lett. B763 (2016), 203-207.

## Modified Soft-Wall Model

Model AdS Poincare patch with an extra UV cutoff:

$$dS^{2} = g_{MN} \, dx^{M} \, dx^{N} = \frac{R^{2}}{z^{2}} \left[ dz^{2} + \eta_{\mu\nu} \, dx^{\mu} \, dx^{\nu} \right] \, \Theta \left( z - z_{0} \right). \tag{1}$$



$$I = I_{\text{Scalar}} + I_{\text{Vector}},\tag{2}$$

#### ) where

$$I_{\text{Scalar}} = -\frac{1}{2g_{S}^{2}} \int d^{5} x \sqrt{-g} e^{-\Phi(z)} \left[ g^{MN} \partial_{M} S \partial_{N} S + M_{5}^{2} S^{2} \right], \quad (3)$$

$$I_{\text{Vector}} = -\frac{1}{2g_{V}^{2}} \int d^{5} x \sqrt{-g} e^{-\Phi(z)} \left[ \frac{1}{2} F_{MN} F^{MN} + \tilde{M}_{5}^{2} g^{MN} A_{M} A_{N} \right], \quad (4)$$

$$\Phi(z) = \kappa^{2} z^{2}. \quad (5)$$

Ref: Braga, Nelson R. F. et al, Phys. Lett. B763 (2016), 203-207.

## Modified Soft-Wall Model

Model AdS Poincare patch with an extra UV cutoff:

$$dS^{2} = g_{MN} dx^{M} dx^{N} = \frac{R^{2}}{z^{2}} \left[ dz^{2} + \eta_{\mu\nu} dx^{\mu} dx^{\nu} \right] \Theta \left( z - z_{0} \right).$$
(1)

#### Associated action:

$$I = I_{\text{Scalar}} + I_{\text{Vector}},\tag{2}$$

#### where

$$I_{\text{Scalar}} = -\frac{1}{2g_{S}^{2}} \int d^{5} x \sqrt{-g} e^{-\Phi(z)} \left[ g^{MN} \partial_{M} S \partial_{N} S + M_{5}^{2} S^{2} \right], \quad (3)$$

$$I_{\text{Vector}} = -\frac{1}{2g_{V}^{2}} \int d^{5} x \sqrt{-g} e^{-\Phi(z)} \left[ \frac{1}{2} F_{MN} F^{MN} + \tilde{M}_{5}^{2} g^{MN} A_{M} A_{N} \right], \quad (4)$$

$$\Phi(z) = \kappa^{2} z^{2}. \quad (5)$$

Ref: Braga, Nelson R. F. et al, Phys. Lett. B763 (2016), 203-207.

Holographic Setup

#### Modified Soft-Wall Model - Hadronic Identity

**(**) Hadronic dimension  $\Delta$  and Bulk Mass:

$$M_5^2 R^2 = \Delta (\Delta - 4) - s (s - 4).$$
(6)

イロト イボト イヨト イヨト

Modified Bulk Mass:

$$M_5^2 R^2 = \left(\Delta_{\text{Phys}} + \Delta_P\right) \left(\Delta_{\text{Phys}} + \Delta_P - 4\right) - s\left(s - 4\right). \tag{7}$$

#### Table:

Meson Identity	$\Delta_P$	$M_5^2 R^2$
Scalar meson		-3
Vector meson		
Pseudoscalar meson	-1	-4
Axial vector meson	-1	-1

Ref: He, Song et al, Eur. Phys. J. C66 (2010), 187-196; Vega, Alfredo et al, Phys. Rev. D79 (2009), 055003.

э

Holographic Setup

#### Modified Soft-Wall Model - Hadronic Identity

**(**) Hadronic dimension  $\Delta$  and Bulk Mass:

$$M_5^2 R^2 = \Delta (\Delta - 4) - s (s - 4).$$
(6)

イロト イボト イヨト イヨト

Modified Bulk Mass:

$$M_5^2 R^2 = (\Delta_{\mathsf{Phys}} + \Delta_P) \left( \Delta_{\mathsf{Phys}} + \Delta_P - 4 \right) - s \left( s - 4 \right). \tag{7}$$

Table

Meson Identity	$\Delta_P$	$M_5^2 R^2$
Scalar meson		-3
Vector meson		
Pseudoscalar meson	-1	-4
Axial vector meson	-1	-1

Ref: He, Song et al, Eur. Phys. J. C66 (2010), 187-196; Vega, Alfredo et al, Phys. Rev. D79 (2009), 055003.

э

## Modified Soft-Wall Model - Hadronic Identity

**(**) Hadronic dimension  $\Delta$  and Bulk Mass:

$$M_5^2 R^2 = \Delta (\Delta - 4) - s (s - 4).$$
 (6)

Modified Bulk Mass:

$$M_5^2 R^2 = (\Delta_{\mathsf{Phys}} + \Delta_P) \left( \Delta_{\mathsf{Phys}} + \Delta_P - 4 \right) - s \left( s - 4 \right). \tag{7}$$



Meson Identity	$\Delta_P$	$M_5^2 R^2$
Scalar meson	0	-3
Vector meson	0	0
Pseudoscalar meson	-1	-4
Axial vector meson	-1	-1

Ref: He, Song et al, Eur. Phys. J. C66 (2010), 187-196; Vega, Alfredo et al, Phys. Rev. D79 (2009), 055003.

э

#### 5D scalar/vector action.

- Small variations in the fields involved.
- Euler-Lagrange Equations of Motion.
  - Boundary/On-shell action term evaluated at  $z = z_0$ .
  - Bulk-to-boundary propagator (BtBP).
- BtBP into On-shell action with source terms.
- Two-point scalar/vector function (2PF).
  - 2PF poles = mass spectra.
  - (Non)linear Regge trajectories in (M<sup>2</sup>, n) space.

- 5D scalar/vector action.
- Small variations in the fields involved.
- Euler-Lagrange Equations of Motion.
  - Boundary/On-shell action term evaluated at  $z = z_0$ .
  - Bulk-to-boundary propagator (BtBP).
- BtBP into On-shell action with source terms.
- Two-point scalar/vector function (2PF).
  - 2PF poles = mass spectra.
  - (Non)linear Regge trajectories in (M<sup>2</sup>, n) space.

- 5D scalar/vector action.
- Small variations in the fields involved.
- Euler-Lagrange Equations of Motion.
  - Boundary/On-shell action term evaluated at  $z = z_0$ .
  - Bulk-to-boundary propagator (BtBP).
- BtBP into On-shell action with source terms.
- Two-point scalar/vector function (2PF).
  - 2PF poles = mass spectra.
  - (Non)linear Regge trajectories in (M<sup>2</sup>, n) space.

- 5D scalar/vector action.
- Small variations in the fields involved.
- Euler-Lagrange Equations of Motion.
  - Boundary/On-shell action term evaluated at  $z = z_0$ .
  - Bulk-to-boundary propagator (BtBP).
- BtBP into On-shell action with source terms.
- Two-point scalar/vector function (2PF).
  - 2PF poles = mass spectra.
  - (Non)linear Regge trajectories in (M<sup>2</sup>, n) space.

- 5D scalar/vector action.
- Small variations in the fields involved.
- Euler-Lagrange Equations of Motion.
  - Boundary/On-shell action term evaluated at  $z = z_0$ .
  - Bulk-to-boundary propagator (BtBP).
- BtBP into On-shell action with source terms.
- Two-point scalar/vector function (2PF).
  - 2PF poles = mass spectra.
  - (Non)linear Regge trajectories in (M<sup>2</sup>, n) space.

- 5D scalar/vector action.
- Small variations in the fields involved.
- Euler-Lagrange Equations of Motion.
  - Boundary/On-shell action term evaluated at  $z = z_0$ .
  - Bulk-to-boundary propagator (BtBP).
- BtBP into On-shell action with source terms.
- Two-point scalar/vector function (2PF).
  - 2PF poles = mass spectra.
  - (Non)linear Regge trajectories in (M<sup>2</sup>, n) space.

- 5D scalar/vector action.
- Small variations in the fields involved.
- Euler-Lagrange Equations of Motion.
  - Boundary/On-shell action term evaluated at  $z = z_0$ .
  - Bulk-to-boundary propagator (BtBP).
- BtBP into On-shell action with source terms.
- Solution (2007) Two-point scalar/vector function (2007).
  - 2PF poles = mass spectra.
  - (Non)linear Regge trajectories in  $(M^2, n)$  space.

イロト 不得下 イヨト イヨト

- 5D scalar/vector action.
- Small variations in the fields involved.
- Euler-Lagrange Equations of Motion.
  - Boundary/On-shell action term evaluated at  $z = z_0$ .
  - Bulk-to-boundary propagator (BtBP).
- BtBP into On-shell action with source terms.
- Solution (2007) Two-point scalar/vector function (2007).
  - 2PF poles = mass spectra.
  - (Non)linear Regge trajectories in  $(M^2, n)$  space.

< ロ > < 同 > < 回 > < 回 > < 回 > <

- 5D scalar/vector action.
- Small variations in the fields involved.
- Euler-Lagrange Equations of Motion.
  - Boundary/On-shell action term evaluated at  $z = z_0$ .
  - Bulk-to-boundary propagator (BtBP).
- BtBP into On-shell action with source terms.
- Solution (2007) Two-point scalar/vector function (2007).
  - 2PF poles = mass spectra.
  - (Non)linear Regge trajectories in  $(M^2, n)$  space.

< □ > < □ > < □ > < □ > < □ > < □ >

## Holographic Description of Pseudoscalar Mesons

Equation of Motion (after taking 
$$S(z,q) = S_0(q)\mathcal{V}(z,q)$$
):

$$\partial_z \left[ \frac{e^{-\kappa^2 z^2}}{z^3} \partial_z \mathcal{V} \right] + \frac{e^{-\kappa^2 z^2}}{z^3} q^2 \mathcal{V} + \frac{4 e^{-\kappa^2 z^2}}{z^5} \mathcal{V} = 0.$$
(8)

Normalized pseudoscalar propagator:

$$\mathcal{V}_{\eta}\left(z,q\right) = \frac{z^{2} {}_{1}F_{1}\left(1 - \frac{q^{2}}{4\kappa^{2}}, 1, \kappa^{2} z^{2}\right)}{z_{0}^{2} {}_{1}F_{1}\left(1 - \frac{q^{2}}{4\kappa^{2}}, 1, \kappa^{2} z_{0}^{2}\right)}.$$
(9)

 $\eta$ -meson two-point function:

$$\Pi_{\eta}\left(q^{2}\right) = -\frac{R^{3}}{g_{5}^{2}} \frac{e^{-\kappa^{2} z_{0}^{2}}}{z_{0}^{3}} \left[\frac{2}{z_{0}} + \left(1 - \frac{q^{2}}{4\kappa^{2}}\right) \frac{2\kappa^{2} z_{0\,1}F_{1}\left(2 - \frac{q^{2}}{4\kappa^{2}}, 2, \kappa^{2} z_{0}^{2}\right)}{{}_{1}F_{1}\left(1 - \frac{q^{2}}{4\kappa^{2}}, 1, \kappa^{2} z_{0}^{2}\right)}\right].$$
(10)

Ref: Witten, Edward, Adv. Theor. Math. Phys. 2 (1998) 253-291; Braga, Nelson R. F. et al, Phys. Lett. B763 (2016), 203-207.

#### Holographic Description of Pseudoscalar Mesons

**(**) Equation of Motion (after taking  $S(z,q) = S_0(q)\mathcal{V}(z,q)$ ):

$$\partial_{z} \left[ \frac{e^{-\kappa^{2} z^{2}}}{z^{3}} \partial_{z} \mathcal{V} \right] + \frac{e^{-\kappa^{2} z^{2}}}{z^{3}} q^{2} \mathcal{V} + \frac{4 e^{-\kappa^{2} z^{2}}}{z^{5}} \mathcal{V} = 0.$$
(8)

Normalized pseudoscalar propagator:

$$\mathcal{V}_{\eta}(z,q) = \frac{z^2 {}_{1}F_{1}\left(1 - \frac{q^2}{4\kappa^2}, 1, \kappa^2 z^2\right)}{z_{0\ 1}^2 F_{1}\left(1 - \frac{q^2}{4\kappa^2}, 1, \kappa^2 z_{0}^2\right)}.$$
(9)

 $\eta$ -meson two-point function:

$$\Pi_{\eta}\left(q^{2}\right) = -\frac{R^{3}}{g_{S}^{2}} \frac{e^{-\kappa^{2} z_{0}^{2}}}{z_{0}^{3}} \left[\frac{2}{z_{0}} + \left(1 - \frac{q^{2}}{4\kappa^{2}}\right) \frac{2\kappa^{2} z_{0\,1}F_{1}\left(2 - \frac{q^{2}}{4\kappa^{2}}, 2, \kappa^{2} z_{0}^{2}\right)}{{}_{1}F_{1}\left(1 - \frac{q^{2}}{4\kappa^{2}}, 1, \kappa^{2} z_{0}^{2}\right)}\right].$$
(10)

Ref: Witten, Edward, Adv. Theor. Math. Phys. 2 (1998) 253-291; Braga, Nelson R. F. et al, Phys. Lett. B763 (2016), 203-207.

## Holographic Description of Pseudoscalar Mesons

**Q** Equation of Motion (after taking  $S(z,q) = S_0(q)\mathcal{V}(z,q)$ ):

$$\partial_{z} \left[ \frac{e^{-\kappa^{2} z^{2}}}{z^{3}} \partial_{z} \mathcal{V} \right] + \frac{e^{-\kappa^{2} z^{2}}}{z^{3}} q^{2} \mathcal{V} + \frac{4 e^{-\kappa^{2} z^{2}}}{z^{5}} \mathcal{V} = 0.$$
(8)

Normalized pseudoscalar propagator:

$$\mathcal{V}_{\eta}(z,q) = \frac{z^2 {}_{1}F_{1}\left(1 - \frac{q^2}{4\kappa^2}, 1, \kappa^2 z^2\right)}{z_{0}^2 {}_{1}F_{1}\left(1 - \frac{q^2}{4\kappa^2}, 1, \kappa^2 z_{0}^2\right)}.$$
(9)

 $\eta$ -meson two-point function:

$$\Pi_{\eta}\left(q^{2}\right) = -\frac{R^{3}}{g_{5}^{2}} \frac{e^{-\kappa^{2} z_{0}^{2}}}{z_{0}^{3}} \left[\frac{2}{z_{0}} + \left(1 - \frac{q^{2}}{4\kappa^{2}}\right) \frac{2\kappa^{2} z_{0\,1}F_{1}\left(2 - \frac{q^{2}}{4\kappa^{2}}, 2, \kappa^{2} z_{0}^{2}\right)}{{}_{1}F_{1}\left(1 - \frac{q^{2}}{4\kappa^{2}}, 1, \kappa^{2} z_{0}^{2}\right)}\right].$$
(10)

Ref: Witten, Edward, Adv. Theor. Math. Phys. 2 (1998) 253-291; Braga, Nelson R. F. et al, Phys. Lett. B763 (2016), 203-207.

#### Axial Mesons

#### Holographic Description of Axial Mesons



Equation of Motion (after taking  $A_m(z,q) = A_m^0(q) \mathcal{V}(z,q)$ ):

$$\partial_z \left[ \frac{e^{-\kappa^2 z^2}}{z} \, \partial_z \, \mathcal{V} \right] + q^2 \, \frac{e^{-\kappa^2 z^2}}{z} \, \mathcal{V} + \frac{e^{-\kappa^2 z^2}}{z^3} \, \mathcal{V} = 0, \tag{11}$$

Normalized axial propagator:

$$\mathcal{V}_{a_1}(z,q) = \frac{z_1 F_1\left(\frac{1}{2} - \frac{q^2}{4\kappa^2}, 1, \kappa^2 z^2\right)}{z_{0\,1} F_1\left(\frac{1}{2} - \frac{q^2}{4\kappa^2}, 1, \kappa^2 z_0^2\right)}.$$
(12)

*a*<sub>1</sub>-meson two-point function:

$$\Pi_{a_{1}}^{\mu\nu}\left(q^{2}\right) = -\frac{R}{g_{V}^{2}} \frac{e^{-\kappa^{2} z_{0}^{2}}}{z_{0}} \left[\frac{1}{z_{0}} + \left(\frac{1}{2} - \frac{q^{2}}{4\kappa^{2}}\right) \frac{2\kappa^{2} z_{0} {}_{1}F_{1}\left(\frac{3}{2} - \frac{q^{2}}{4\kappa^{2}}, 2, \kappa^{2} z_{0}^{2}\right)}{{}_{1}F_{1}\left(\frac{1}{2} - \frac{q^{2}}{4\kappa^{2}}, 1, \kappa^{2} z_{0}^{2}\right)}\right] \left(\eta^{\mu\nu} - \frac{q^{\mu} q^{\nu}}{q^{2}}\right).$$
(13)

### Holographic Description of Axial Mesons

**(**) Equation of Motion (after taking  $A_m(z,q) = A_m^0(q) \mathcal{V}(z,q)$ ):

$$\partial_z \left[ \frac{e^{-\kappa^2 z^2}}{z} \partial_z \mathcal{V} \right] + q^2 \frac{e^{-\kappa^2 z^2}}{z} \mathcal{V} + \frac{e^{-\kappa^2 z^2}}{z^3} \mathcal{V} = 0, \qquad (11)$$

Normalized axial propagator:

$$\mathcal{V}_{a_1}(z,q) = \frac{z_1 F_1\left(\frac{1}{2} - \frac{q^2}{4\kappa^2}, 1, \kappa^2 z^2\right)}{z_{0\,1} F_1\left(\frac{1}{2} - \frac{q^2}{4\kappa^2}, 1, \kappa^2 z_0^2\right)}.$$
(12)

(日)

a<sub>1</sub>-meson two-point function:

$$\Pi_{a_{1}}^{\mu\nu}\left(q^{2}\right) = -\frac{R}{g_{V}^{2}} \frac{e^{-\kappa^{2} z_{0}^{2}}}{z_{0}} \left[\frac{1}{z_{0}} + \left(\frac{1}{2} - \frac{q^{2}}{4\kappa^{2}}\right) \frac{2\kappa^{2} z_{0} {}_{1}F_{1}\left(\frac{3}{2} - \frac{q^{2}}{4\kappa^{2}}, 2, \kappa^{2} z_{0}^{2}\right)}{{}_{1}F_{1}\left(\frac{1}{2} - \frac{q^{2}}{4\kappa^{2}}, 1, \kappa^{2} z_{0}^{2}\right)}\right] \left(\eta^{\mu\nu} - \frac{q^{\mu} q^{\nu}}{q^{2}}\right).$$
(13)

### Holographic Description of Axial Mesons

**(**) Equation of Motion (after taking  $A_m(z,q) = A_m^0(q) \mathcal{V}(z,q)$ ):

$$\partial_z \left[ \frac{e^{-\kappa^2 z^2}}{z} \partial_z \mathcal{V} \right] + q^2 \frac{e^{-\kappa^2 z^2}}{z} \mathcal{V} + \frac{e^{-\kappa^2 z^2}}{z^3} \mathcal{V} = 0, \qquad (11)$$

Normalized axial propagator:

$$\mathcal{V}_{a_1}(z,q) = \frac{z_1 F_1\left(\frac{1}{2} - \frac{q^2}{4\kappa^2}, 1, \kappa^2 z^2\right)}{z_{0\,1} F_1\left(\frac{1}{2} - \frac{q^2}{4\kappa^2}, 1, \kappa^2 z_0^2\right)}.$$
(12)

 $a_1$ -meson two-point function:

$$\Pi_{a_{1}}^{\mu\nu}\left(q^{2}\right) = -\frac{R}{g_{V}^{2}} \frac{e^{-\kappa^{2} z_{0}^{2}}}{z_{0}} \left[\frac{1}{z_{0}} + \left(\frac{1}{2} - \frac{q^{2}}{4\kappa^{2}}\right) \frac{2\kappa^{2} z_{0} {}_{1}F_{1}\left(\frac{3}{2} - \frac{q^{2}}{4\kappa^{2}}, 2, \kappa^{2} z_{0}^{2}\right)}{{}_{1}F_{1}\left(\frac{1}{2} - \frac{q^{2}}{4\kappa^{2}}, 1, \kappa^{2} z_{0}^{2}\right)}\right] \left(\eta^{\mu\nu} - \frac{q^{\mu} q^{\nu}}{q^{2}}\right).$$
(13)

イロト 不得 トイラト イラト 一日

#### $\eta$ Mesons

Table: Mass spectrum for  $\eta$  pseudoscalar mesons with  $\kappa = 0.45$  GeV and  $z_0 = 5.0$  GeV<sup>-1</sup>.

$\eta$ trajectory with $\Delta_{P}=-1$				
n	State	$M_{Exp}$ (MeV)	$M_{\rm Th}~({\rm MeV})$	%M
1	$\eta$ (550)	$547.86\pm0.017$	975.25	43.8
2	$\eta(1295)$	$1294\pm4$	1233.6	4.9
3	$\eta$ (1405)	$1408.8\pm1.8$	1455.3	3.2
4	$\eta(1475)$	$1476\pm4$	1652.9	10.7
5	$\eta(1760)$	$1760\pm11$	1829.2	3.8
6	$\eta$ (2225)	$2216 \pm 21$	1992.7	11.3

*Ref:* Martín Contreras, Miguel Ángel *et al*, Chin. J. Phys. 66 (2020) 715-723; Tanabashi, M. *et al*, Review of Particle Physics Phys. Rev. D98 (2018) 030001; Wang, Li-Ming *et al*, Phys. Rev. D96 (2017) 034013.

< □ > < □ > < □ > < □ > < □ > < □ >

#### a<sub>1</sub> Mesons

Table: Mass spectrum for  $a_1$  axial mesons with  $\kappa = 0.45$  GeV and  $z_0 = 5.0$  GeV<sup>-1</sup>

$a_1$ trajectory with $\Delta_P=-1$				
n	State	$M_{Exp}$ (MeV)	$M_{\rm Th}$ (MeV)	%M
1	$a_1(1260)$	$1230\pm40$	808.1	52.2
2	$a_1(1420)$	$1414^{+15}_{-13}$	1114.7	26.8
3	$a_1(1640)$	$1654\pm19$	1351.3	22.4
4	$a_1(1930)$	$1930^{+19}_{-70}$	1558.7	23.8
5	$a_1(2095)$	$2096\pm17\pm121$	1744.3	20.1
6	<i>a</i> <sub>1</sub> (2270)	$2270^{+55}_{-40}$	1913.4	18.6

*Ref:* Martín Contreras, Miguel Ángel *et al*, Chin. J. Phys. 66 (2020) 715-723; Tanabashi, M. *et al*, Review of Particle Physics Phys. Rev. D98 (2018) 030001; Adolph, C. *et al*, Phys. Rev. Lett. 115 (2015) 082001.

#### Scalar Mesons

Table: Mass spectrum for  $f_0(500)$  mesons with  $\kappa = 0.45$  GeV and  $z_0 = 5$  GeV<sup>-1</sup>.

f <sub>0</sub>	$M_{\rm th}~({\rm MeV})$	$M_{\rm exp}$ (MeV)	%M
$f_0(980)$	1.070	0.99	7.46
$f_0(1370)$	1.284	1.370	5.11
$f_0(1500)$	1.487	1.504	1.13
$f_0(1710)$	1.674	1.723	2.93
$f_0(2020)$	1.846	1.992	7.94
$f_0(2100)$	2.153	2.101	2.39
$f_0(2200)$	2.292	2.189	4.49
$f_0(2330)$	2.424	2.314	4.52

*Ref:* Cortés, Santiago *et al*, Phys. Rev. D (2017) 106002; C. Patrignani *et al.*, [Particle Data Group] Chin. Phys. C **40**, no. 10, 100001 (2016).

< □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > <

## Vector Mesons

Table: Mass spectrum for  $\rho$  and mesons with  $\kappa = 0.45$  GeV and  $z_0 = 5$  GeV<sup>-1</sup>.

ρ	$M_{\rm th}~({\rm GeV})$	$M_{\rm exp}~({\rm GeV})$	%M
$\rho(775)$	0.975	0.775	20.53
$\rho$ (1450)	1.455	1.465	0.66
$\rho(1570)$	1.652	1.570	4.96
$\rho(1700)$	1.829	1.720	5.97
$\rho(1900)$	1.992	1.909	4.15
ρ(2150)	2.142	2.153	0.50

*Ref:* Cortés, Santiago *et al*, Phys. Rev. D (2017) 106002; C. Patrignani *et al.*, [Particle Data Group] Chin. Phys. C **40**, no. 10, 100001 (2016).

イロト 不得 トイヨト イヨト

## Conclusions

- We have fitted the radial trajectory for both the  $\eta$  and the  $a_1$  mesons. The first excited states, i. e.,  $\eta(550)$  and  $a_1(1260)$ , were not well predicted by the model. However, we could fit 26 states (among scalar, pseudoscalar, vector and axial particles) by just taking three parameters within a  $\delta_{RMS}$  error close to 21.1%.
- We infer that the small amount of theoretical parameters taken in our approach must have had an important influence in our results for the masses described in this work. Something similar happens with the radial states in non-conformal models of scalar and vector mesons, as shown previously with the first vector state.
- We could confirm that chiral symmetry is restored at highly orbital excited states; this was checked after confirming the rising of some degeneracy in their respective masses.

## Conclusions

- We have fitted the radial trajectory for both the  $\eta$  and the  $a_1$  mesons. The first excited states, i. e.,  $\eta(550)$  and  $a_1(1260)$ , were not well predicted by the model. However, we could fit 26 states (among scalar, pseudoscalar, vector and axial particles) by just taking three parameters within a  $\delta_{RMS}$  error close to 21.1%.
- We infer that the small amount of theoretical parameters taken in our approach must have had an important influence in our results for the masses described in this work. Something similar happens with the radial states in non-conformal models of scalar and vector mesons, as shown previously with the first vector state.
- We could confirm that chiral symmetry is restored at highly orbital excited states; this was checked after confirming the rising of some degeneracy in their respective masses.

## Conclusions

- We have fitted the radial trajectory for both the  $\eta$  and the  $a_1$  mesons. The first excited states, i. e.,  $\eta(550)$  and  $a_1(1260)$ , were not well predicted by the model. However, we could fit 26 states (among scalar, pseudoscalar, vector and axial particles) by just taking three parameters within a  $\delta_{RMS}$  error close to 21.1%.
- We infer that the small amount of theoretical parameters taken in our approach must have had an important influence in our results for the masses described in this work. Something similar happens with the radial states in non-conformal models of scalar and vector mesons, as shown previously with the first vector state.
- We could confirm that chiral symmetry is restored at highly orbital excited states; this was checked after confirming the rising of some degeneracy in their respective masses.

(日)