

# Nanga Parbat 2019 (Drell-Yan) fit

## Transverse-momentum-dependent parton distributions up to $N^3$ LL from Drell-Yan data

Alessandro Bacchetta, Valerio Bertone, Chiara Bissoletti (INFN, Pavia & Pavia U.), Giuseppe Bozzi (Florence U. & INFN, Florence), Filippo Delcarro (Pavia U.), Fulvio Piacenza (INFN, Pavia & Pavia U.), Marco Radici (INFN, Pavia)

Dec 16, 2019 - 43 pages

JLAB-THY-19-3121

e-Print: [arXiv:1912.07550](https://arxiv.org/abs/1912.07550) [hep-ph] | [PDF](#)

### Abstract (arXiv)

We present an extraction of unpolarised Transverse-Momentum-Dependent Parton Distribution Functions based on Drell-Yan production data from different experiments, including those at the LHC, and spanning a wide kinematic range. We deal with experimental uncertainties by properly taking into account correlations. We include resummation of logarithms of the transverse momentum of the vector boson up to  $N^3$  LL order, and we include non-perturbative contributions. These ingredients allow us to obtain a remarkable agreement with the data.

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$$\frac{d\sigma}{dQ dy dq_T} = \frac{8\pi\alpha^2 q_T \mathcal{P}}{9Q^3} H(Q, \mu) \times \sum_q c_q(Q) \int_0^\infty db_T b_T J_0(b_T q_T) x_1 \hat{f}_1^q(x_1, b_T; \mu, \zeta_1) x_2 \hat{f}_1^{\bar{q}}(x_2, b_T; \mu, \zeta_2)$$

Diagram illustrating the components of the Drell-Yan cross-section formula:

- Phase-space factor (lepton cuts)** points to  $\mathcal{P}$ .
- Hard factor** points to  $H(Q, \mu)$ .
- electroweak charges** points to  $c_q(Q)$ .
- TMD PDFs** points to  $\hat{f}_1^q(x_1, b_T; \mu, \zeta_1)$  and  $\hat{f}_1^{\bar{q}}(x_2, b_T; \mu, \zeta_2)$ .

# Regularisation

**b\* prescription**

$$b_*(b_T) = b_{\max} \left( \frac{1 - \exp\left(-\frac{b_T^4}{b_{\max}^4}\right)}{1 - \exp\left(-\frac{b_T^4}{b_{\min}^4}\right)} \right)^{\frac{1}{4}}$$

$$\mathbf{b_{\max}} \quad \alpha_s \left( \frac{2e^{-\gamma E}}{b_{\max}} \right) \ll 1$$

$$\mathbf{b_{\min}} \quad b_{\min} = \frac{2e^{-\gamma E}}{Q}$$

(not strictly necessary)

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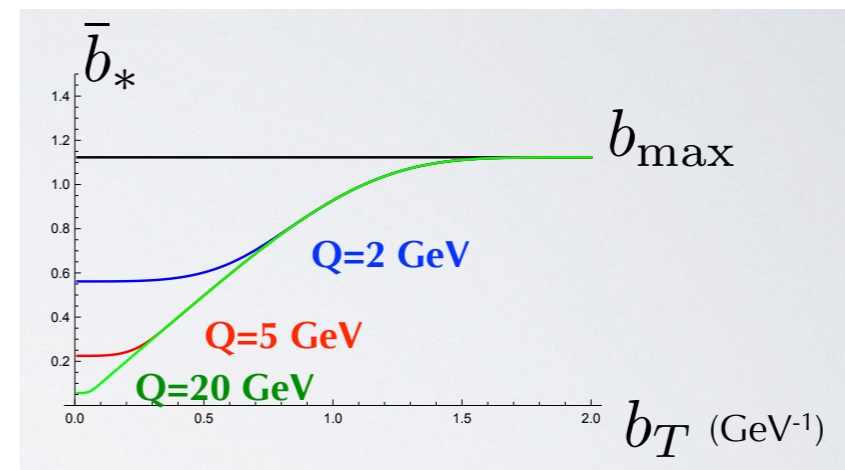
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## limiting values

$$b_*(b_T) \rightarrow b_{\min} \quad \text{for } b_T \rightarrow 0 ,$$

$$b_*(b_T) \rightarrow b_{\max} \quad \text{for } b_T \rightarrow \infty .$$



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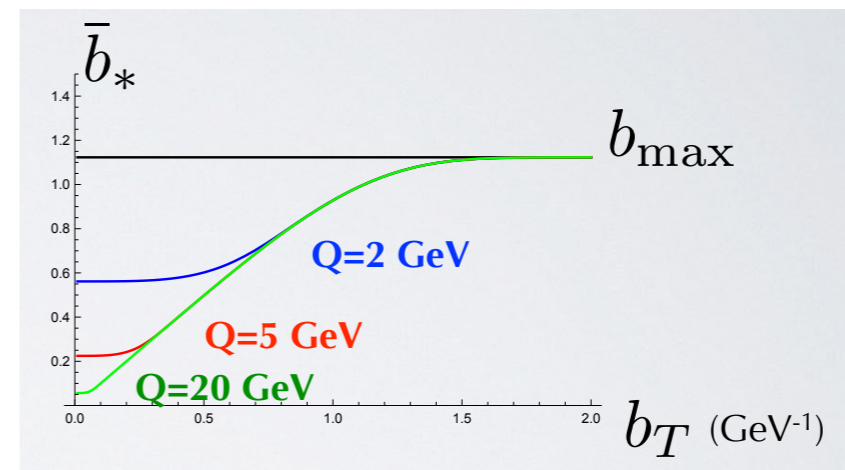
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## Definition of $f_{\text{NP}}$

$$\hat{f}_1(x, b_T; \mu, \zeta) = \left[ \frac{\hat{f}_1(x, b_T; \mu, \zeta)}{\hat{f}_1(x, b_*(b_T); \mu, \zeta)} \right] \hat{f}_1(x, b_*(b_T); \mu, \zeta)$$

$$\equiv f_{\text{NP}}(x, b_T, \zeta) \hat{f}_1(x, b_*(b_T); \mu, \zeta).$$

# Parameterisation

$$f_{\text{NP}}(x, b_T, \zeta) = \left[ \frac{1 - \lambda}{1 + g_1(x) \frac{b_T^2}{4}} + \lambda \exp \left( -g_{1B}(x) \frac{b_T^2}{4} \right) \right] \\ \times \exp \left[ - (g_2 + g_{2B} b_T^2) \ln \left( \frac{\zeta}{Q_0^2} \right) \frac{b_T^2}{4} \right],$$

$$g_1(x) = \frac{N_1}{x\sigma} \exp \left[ -\frac{1}{2\sigma^2} \ln^2 \left( \frac{x}{\alpha} \right) \right],$$

$$g_{1B}(x) = \frac{N_{1B}}{x\sigma_B} \exp \left[ -\frac{1}{2\sigma_B^2} \ln^2 \left( \frac{x}{\alpha_B} \right) \right].$$

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“intrinsic” NP contribution  
(x- and b<sub>T</sub>-dependent)



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- “intrinsic” contribution: q-Gaussian (Tsallis) + simple Gaussian
  - q-Gaussian has larger tail → bigger contribution at small q<sub>T</sub>

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- NP correction: quadratic and quartic terms
  - quartic term allows to reproduce energy evolution



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x-dep. width of TMDs

- “intrinsic” contribution: q-Gaussian (Tsallis) + simple Gaussian
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“intrinsic” NP contribution  
( $x$ - and  $b_T$ -dependent)

NP correction to pert. evolution  
( $b_T$ -dependent)

$$g_1(x) = \frac{N_1}{x\sigma} \exp \left[ -\frac{1}{2\sigma^2} \ln^2 \left( \frac{x}{\alpha} \right) \right],$$

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$x$ -dep. width of TMDs

- “intrinsic” contribution:  $q$ -Gaussian (Tsallis) + simple Gaussian
  - $q$ -Gaussian has larger tail  $\rightarrow$  bigger contribution at small  $q_T$
- NP correction: quadratic and quartic terms
  - quartic term allows to reproduce energy evolution

**9 parameters**  $(\lambda, g_2, g_{2B}, N_1, \sigma, \alpha, N_{1B}, \sigma_B, \alpha_B)$

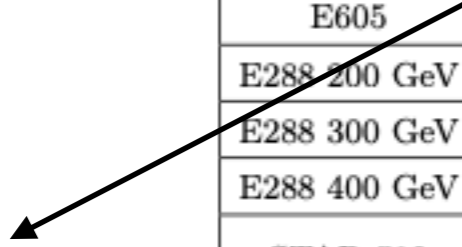
# Experimental data

Experiment	$N_{\text{dat}}$	Observable	$\sqrt{s}$ [GeV]	$Q$ [GeV]	$y$ or $x_F$	Lepton cuts
E605	50	$Ed^3\sigma/d^3q$	38.8	7 - 18	$x_F = 0.1$	-
E288 200 GeV	30	$Ed^3\sigma/d^3q$	19.4	4 - 9	$y = 0.40$	-
E288 300 GeV	39	$Ed^3\sigma/d^3q$	23.8	4 - 12	$y = 0.21$	-
E288 400 GeV	61	$Ed^3\sigma/d^3q$	27.4	5 - 14	$y = 0.03$	-
STAR 510	7	$d\sigma/dq_T$	510	73 - 114	$ y  < 1$	$p_{T\ell} > 25$ GeV $ \eta_\ell  < 1$
CDF Run I	25	$d\sigma/dq_T$	1800	66 - 116	Inclusive	-
CDF Run II	26	$d\sigma/dq_T$	1960	66 - 116	Inclusive	-
D0 Run I	12	$d\sigma/dq_T$	1800	75 - 105	Inclusive	-
D0 Run II	5	$(1/\sigma)d\sigma/dq_T$	1960	70 - 110	Inclusive	-
D0 Run II ( $\mu$ )	3	$(1/\sigma)d\sigma/dq_T$	1960	65 - 115	$ y  < 1.7$	$p_{T\ell} > 15$ GeV $ \eta_\ell  < 1.7$
LHCb 7 TeV	7	$d\sigma/dq_T$	7000	60 - 120	$2 < y < 4.5$	$p_{T\ell} > 20$ GeV $2 < \eta_\ell < 4.5$
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ATLAS 8 TeV off-peak	4 8	$(1/\sigma)d\sigma/dq_T$	8000	46 - 66 116 - 150	$ y  < 2.4$	$p_{T\ell} > 20$ GeV $ \eta_\ell  < 2.4$
Total	353	-	-	-	-	-

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Only data with  
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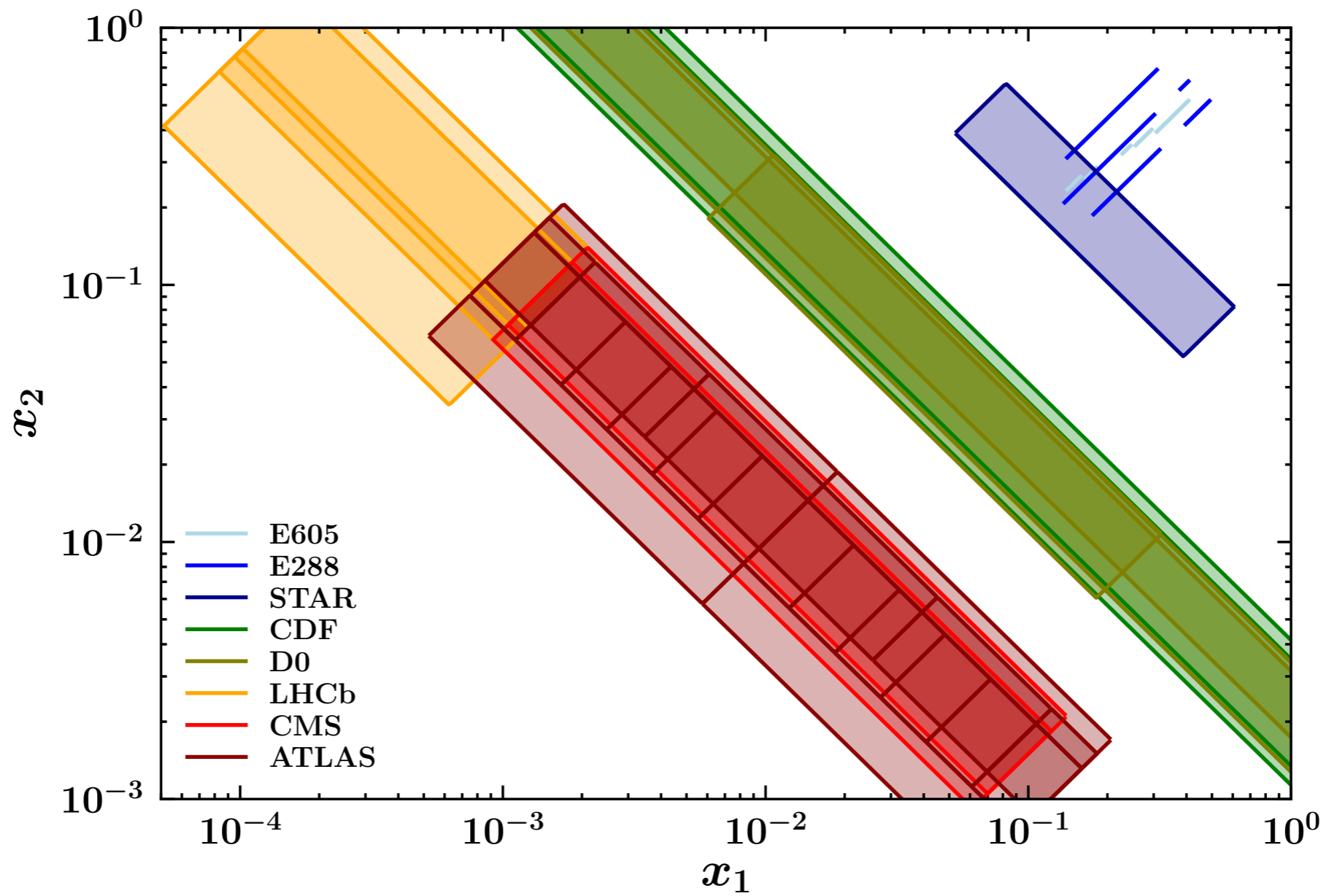
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Total	353	-	-	-	-	-

Only data with  
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DYNNLO  
for total  $\sigma$

# Experimental data



# Uncertainties

**Correlation matrix**

$$V_{ij} = (\sigma_{i,\text{stat}}^2 + \sigma_{i,\text{unc}}^2) \delta_{ij} + \sum_{l=1}^k \sigma_{i,\text{corr}}^{(l)} \sigma_{j,\text{corr}}^{(l)}$$

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**Chi-square**

$$\begin{aligned} \chi^2 &= \sum_{i,j=1}^n (m_i - t_i) V_{ij}^{-1} (m_j - t_j) \\ &= \sum_{i=1}^n \frac{1}{s_i^2} \left( m_i - t_i - \sum_{\alpha=1}^k \lambda_{\alpha} \sigma_{i,\text{corr}}^{(\alpha)} \right)^2 + \sum_{\alpha=1}^k \lambda_{\alpha}^2 \end{aligned}$$



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nuisance  
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nuisance parameters

$s_i^2 = \sigma_{i,\text{stat}}^2 + \sigma_{i,\text{unc}}^2$

- Minimising  $\chi^2$  with respect to  $\lambda_{\alpha}$   $\rightarrow$  optimal values for  $\lambda_{\alpha}$

# Uncertainties

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$$V_{ij} = (\sigma_{i,\text{stat}}^2 + \sigma_{i,\text{unc}}^2) \delta_{ij} + \sum_{l=1}^k \sigma_{i,\text{corr}}^{(l)} \sigma_{j,\text{corr}}^{(l)}$$

**Chi-square**

$$\begin{aligned} \chi^2 &= \sum_{i,j=1}^n (m_i - t_i) V_{ij}^{-1} (m_j - t_j) \\ &= \sum_{i=1}^n \frac{1}{s_i^2} \left( m_i - t_i - \sum_{\alpha=1}^k \lambda_{\alpha} \sigma_{i,\text{corr}}^{(\alpha)} \right)^2 + \sum_{\alpha=1}^k \lambda_{\alpha}^2 \end{aligned}$$

nuisance parameters

$$s_i^2 = \sigma_{i,\text{stat}}^2 + \sigma_{i,\text{unc}}^2$$

- Minimising  $\chi^2$  with respect to  $\lambda_{\alpha}$   $\rightarrow$  optimal values for  $\lambda_{\alpha}$
- $\sum_{\alpha=1}^k \lambda_{\alpha} \sigma_{i,\text{corr}}^{\alpha}$  = shift induced by systematic correlated uncertainties

# Uncertainties

**Correlation matrix**

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nuisance parameters

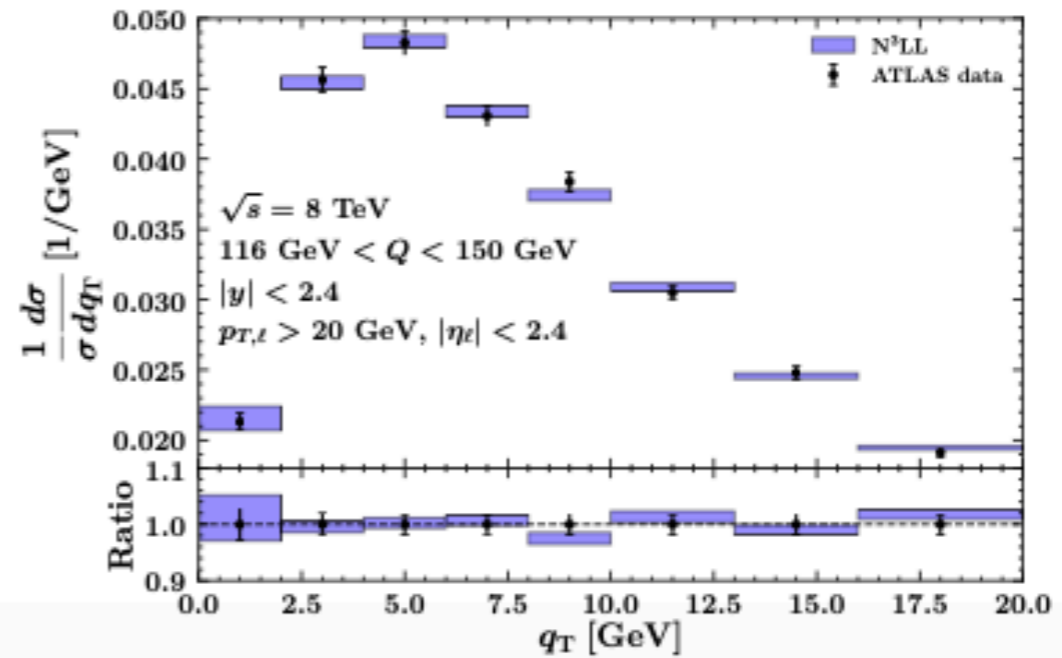
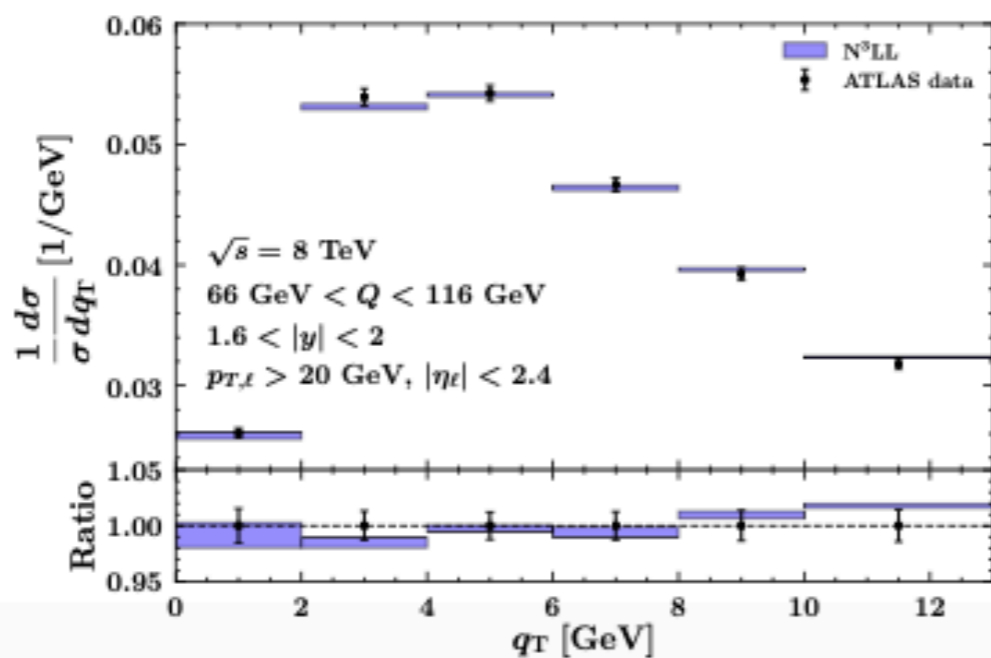
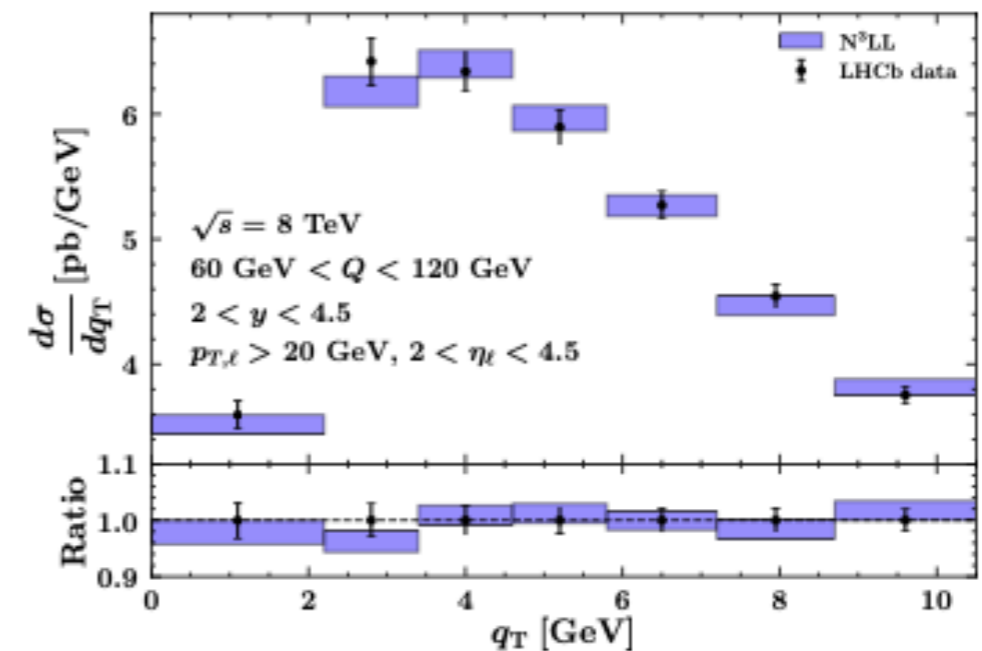
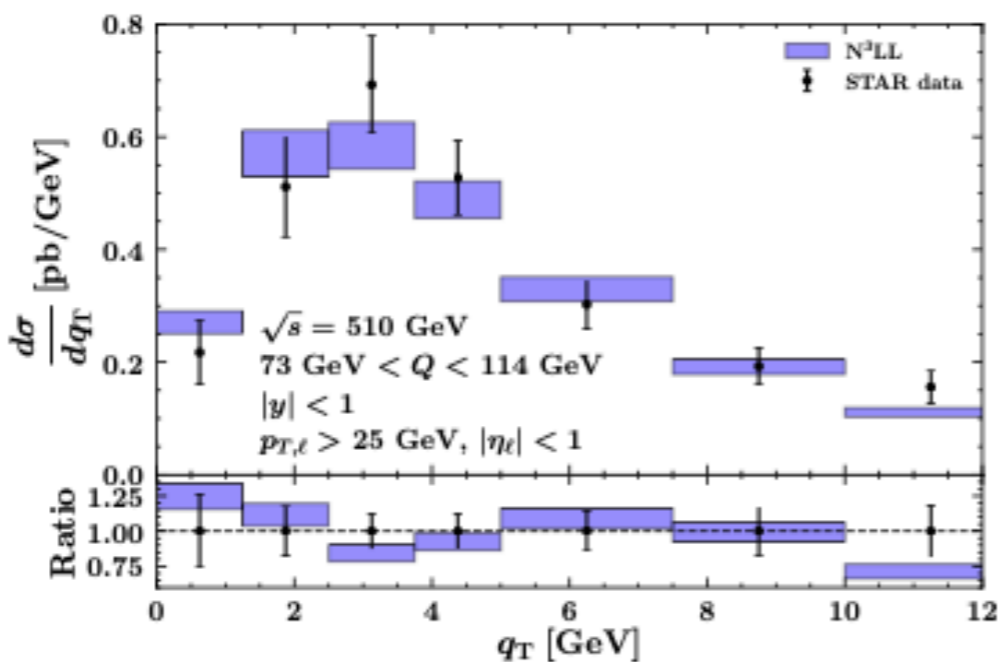
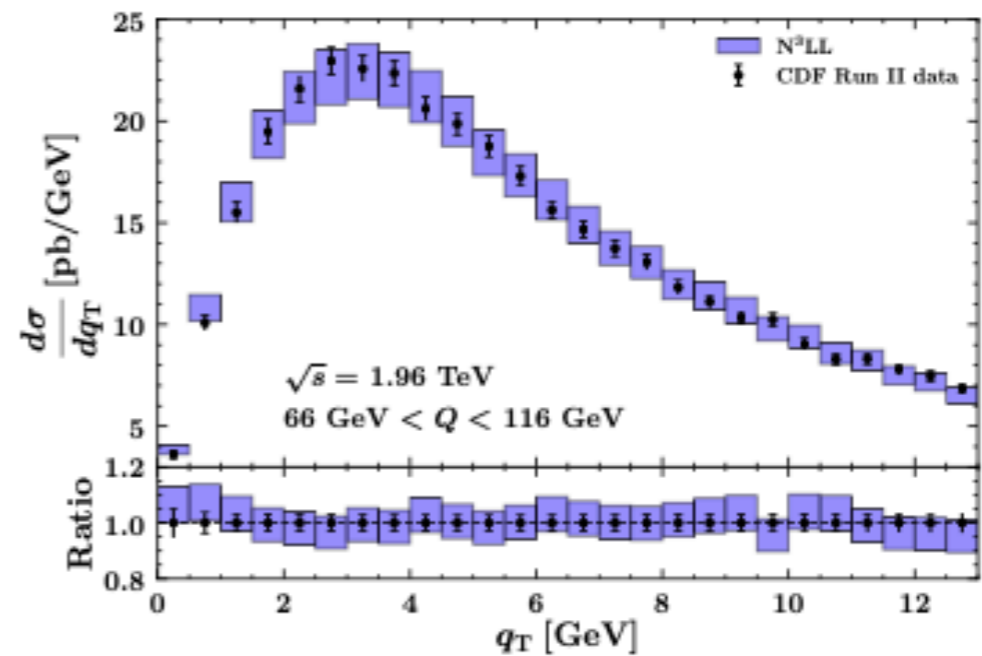
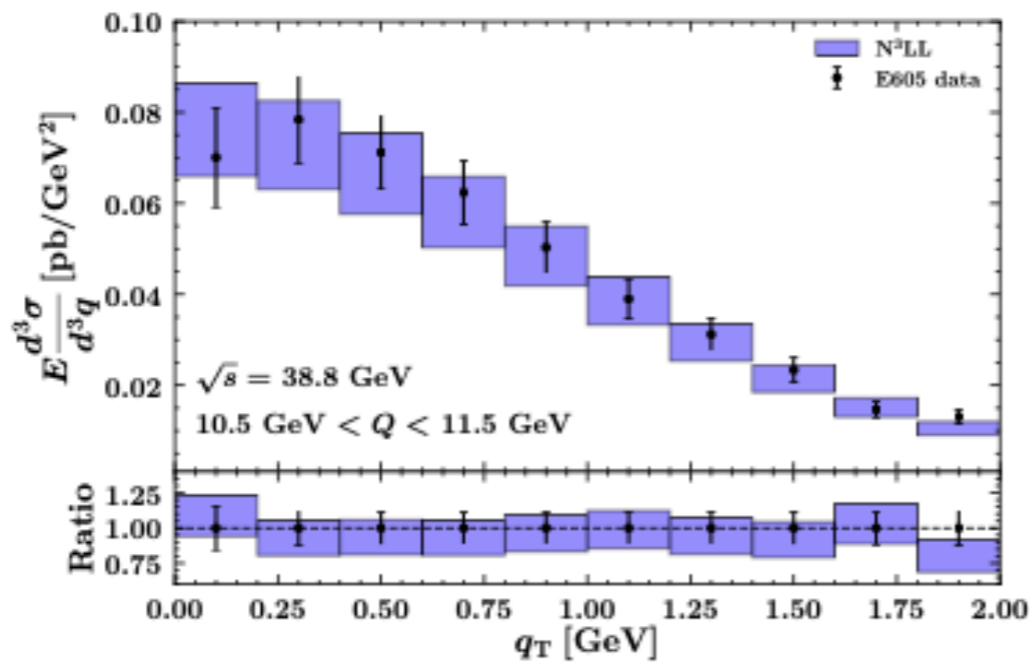
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- Minimising  $\chi^2$  with respect to  $\lambda_{\alpha}$   $\rightarrow$  optimal values for  $\lambda_{\alpha}$
- $\sum_{\alpha=1}^k \lambda_{\alpha} \sigma_{i,\text{corr}}^{\alpha}$  = shift induced by systematic correlated uncertainties
- By shifting theoretical predictions,  $\chi^2 = \chi_D^2 + \chi_{\lambda}^2$  (diagonal + penalty term)

# Fit quality

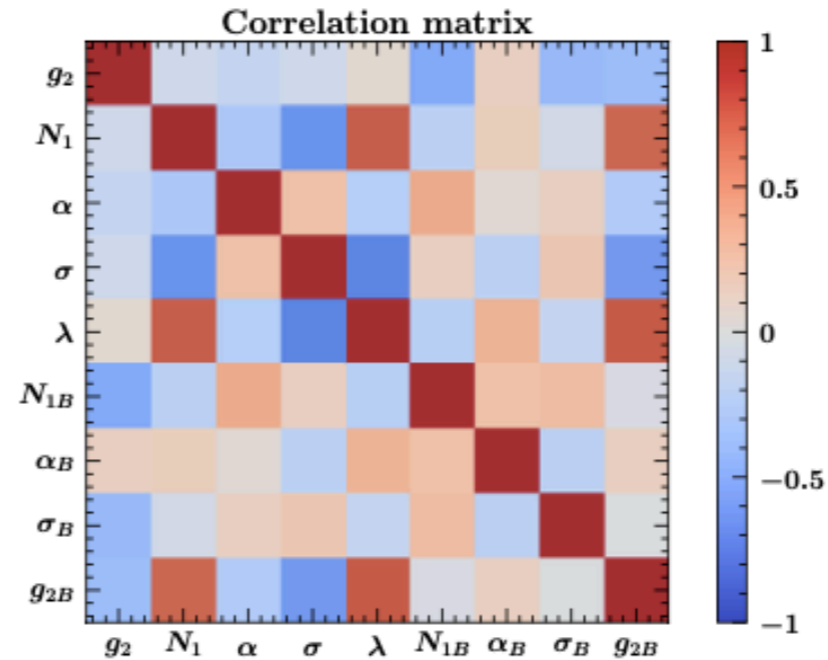
Experiment		$\chi_D^2/N_{\text{dat}}$	$\chi_\lambda^2/N_{\text{dat}}$	$\chi^2/N_{\text{dat}}$
E605	7 GeV < Q < 8 GeV	0.419	0.068	0.487
	8 GeV < Q < 9 GeV	0.995	0.034	1.029
	10.5 GeV < Q < 11.5 GeV	0.191	0.137	0.328
	11.5 GeV < Q < 13.5 GeV	0.491	0.284	0.775
	13.5 GeV < Q < 18 GeV	0.491	0.385	0.877
E288 200 GeV	4 GeV < Q < 5 GeV	0.213	0.649	0.862
	5 GeV < Q < 6 GeV	0.673	0.292	0.965
	6 GeV < Q < 7 GeV	0.133	0.141	0.275
	7 GeV < Q < 8 GeV	0.254	0.014	0.268
	8 GeV < Q < 9 GeV	0.652	0.024	0.676
E288 300 GeV	4 GeV < Q < 5 GeV	0.231	0.555	0.785
	5 GeV < Q < 6 GeV	0.502	0.204	0.706
	6 GeV < Q < 7 GeV	0.315	0.063	0.378
	7 GeV < Q < 8 GeV	0.056	0.030	0.086
	8 GeV < Q < 9 GeV	0.530	0.017	0.547
	11 GeV < Q < 12 GeV	1.047	0.167	1.215
E288 400 GeV	5 GeV < Q < 6 GeV	0.312	0.065	0.377
	6 GeV < Q < 7 GeV	0.100	0.005	0.105
	7 GeV < Q < 8 GeV	0.018	0.011	0.029
	8 GeV < Q < 9 GeV	0.437	0.039	0.477
	11 GeV < Q < 12 GeV	0.637	0.036	0.673
	12 GeV < Q < 13 GeV	0.788	0.028	0.816
	13 GeV < Q < 14 GeV	1.064	0.044	1.107
STAR		0.782	0.054	0.836

CDF Run I		0.480	0.058	0.538
CDF Run II		0.959	0.001	0.959
D0 Run I		0.711	0.043	0.753
D0 Run II		1.325	0.612	1.937
D0 Run II ( $\mu$ )		3.196	0.023	3.218
LHCb 7 TeV		1.069	0.194	1.263
LHCb 8 TeV		0.460	0.075	0.535
LHCb 13 TeV		0.735	0.020	0.755
CMS 7 TeV		2.131	0.000	2.131
CMS 8 TeV		1.405	0.007	1.412
ATLAS 7 TeV	0 <  y  < 1	2.581	0.028	2.609
	1 <  y  < 2	4.333	1.032	5.365
	2 <  y  < 2.4	3.561	0.378	3.939
ATLAS 8 TeV on-peak	0 <  y  < 0.4	1.924	0.337	2.262
	0.4 <  y  < 0.8	2.342	0.247	2.590
	0.8 <  y  < 1.2	0.917	0.061	0.978
	1.2 <  y  < 1.6	0.912	0.095	1.006
	1.6 <  y  < 2	0.721	0.092	0.814
	2 <  y  < 2.4	0.932	0.348	1.280
ATLAS 8 TeV off-peak	46 GeV < Q < 66 GeV	2.138	0.745	2.883
	116 GeV < Q < 150 GeV	0.501	0.003	0.504
<b>Global</b>		<b>0.88</b>	<b>0.14</b>	<b>1.02</b>



# TMD distributions

Parameter	Value
$g_2$	$0.036 \pm 0.009$
$N_1$	$0.625 \pm 0.282$
$\alpha$	$0.205 \pm 0.010$
$\sigma$	$0.370 \pm 0.063$
$\lambda$	$0.580 \pm 0.092$
$N_{1B}$	$0.044 \pm 0.012$
$\alpha_B$	$0.069 \pm 0.009$
$\sigma_B$	$0.356 \pm 0.075$
$g_{2B}$	$0.012 \pm 0.003$

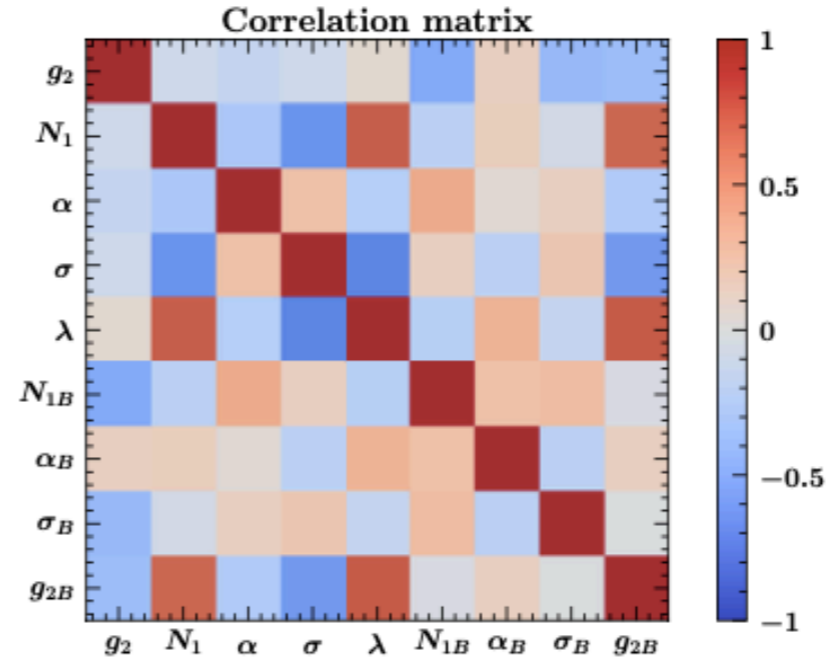


- $\lambda \sim 0.5$ : gaussian and q-gaussian equally important
- $g_{2B}$  (NP quartic term) small but significantly different from zero
- off-diagonal elements (correlations) not very large except for  $\lambda$  and  $\sigma$

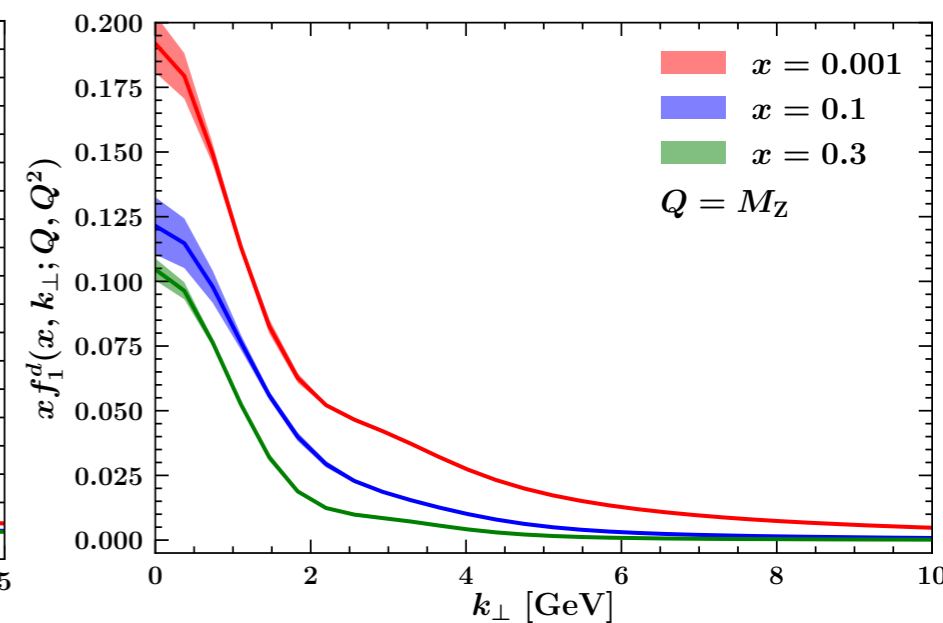
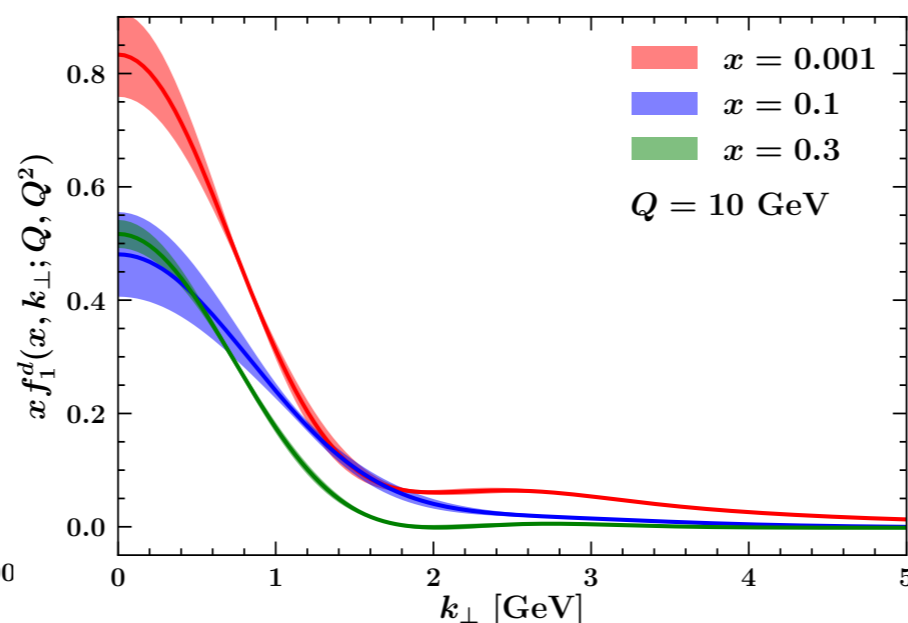
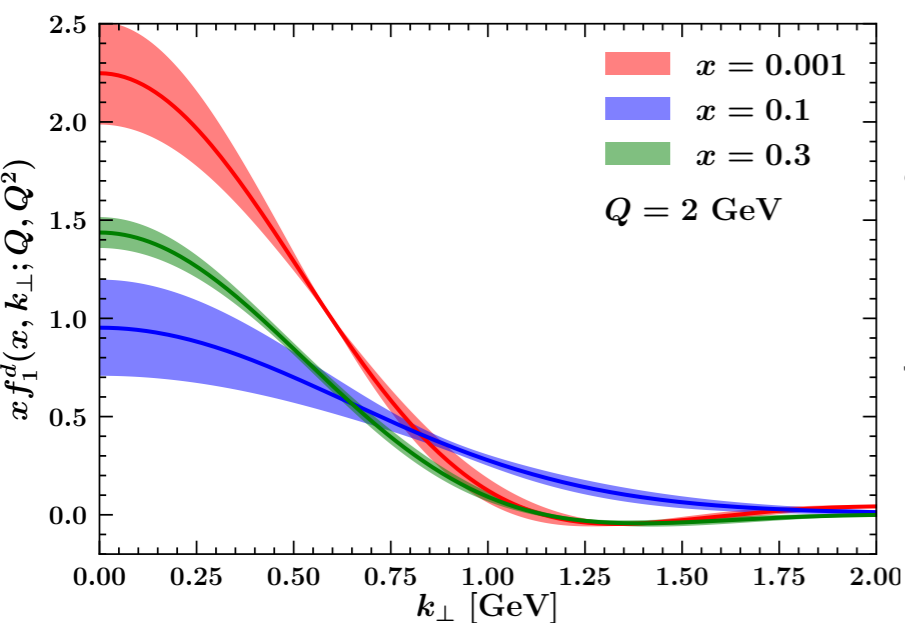


# TMD distributions

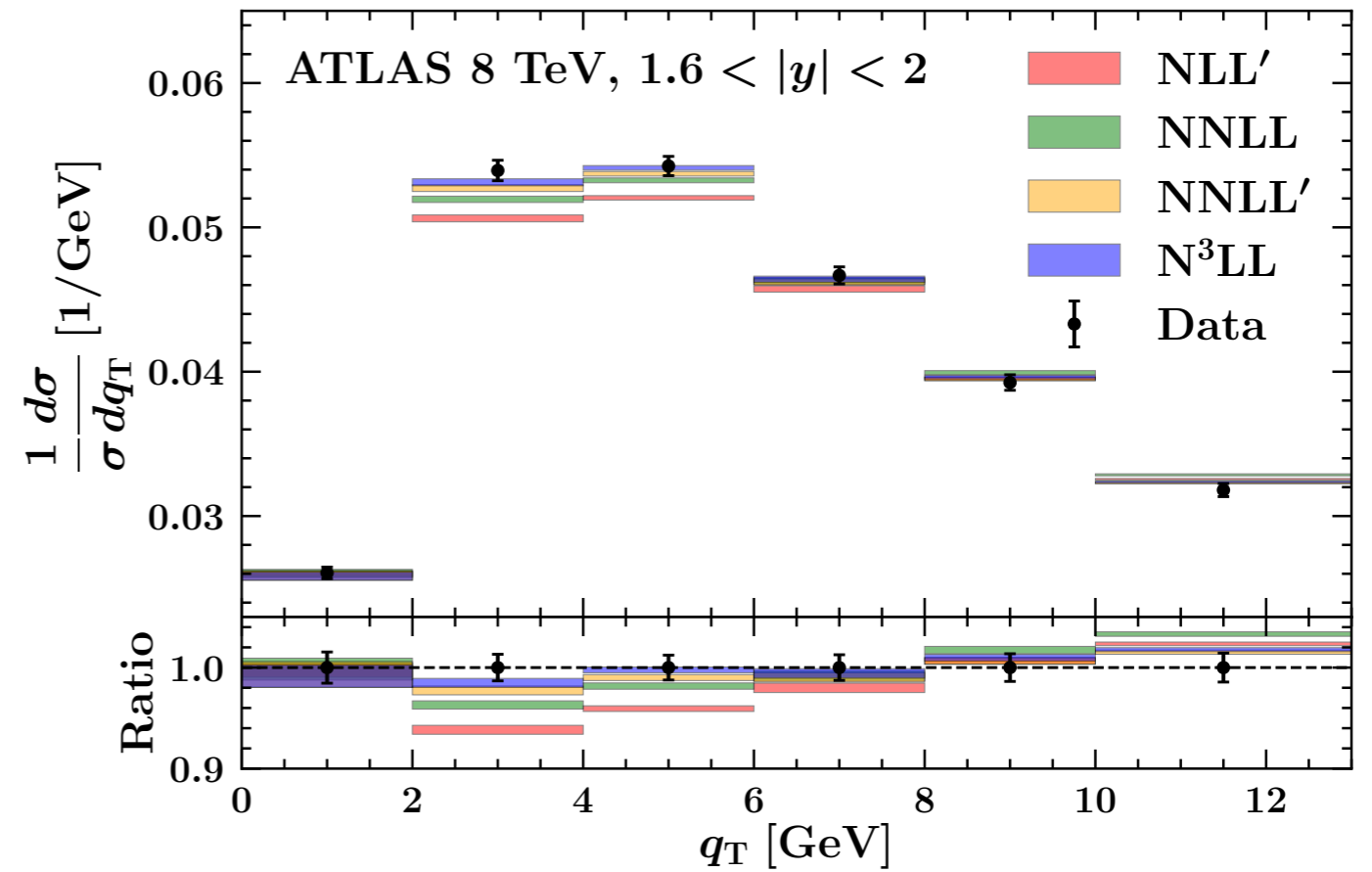
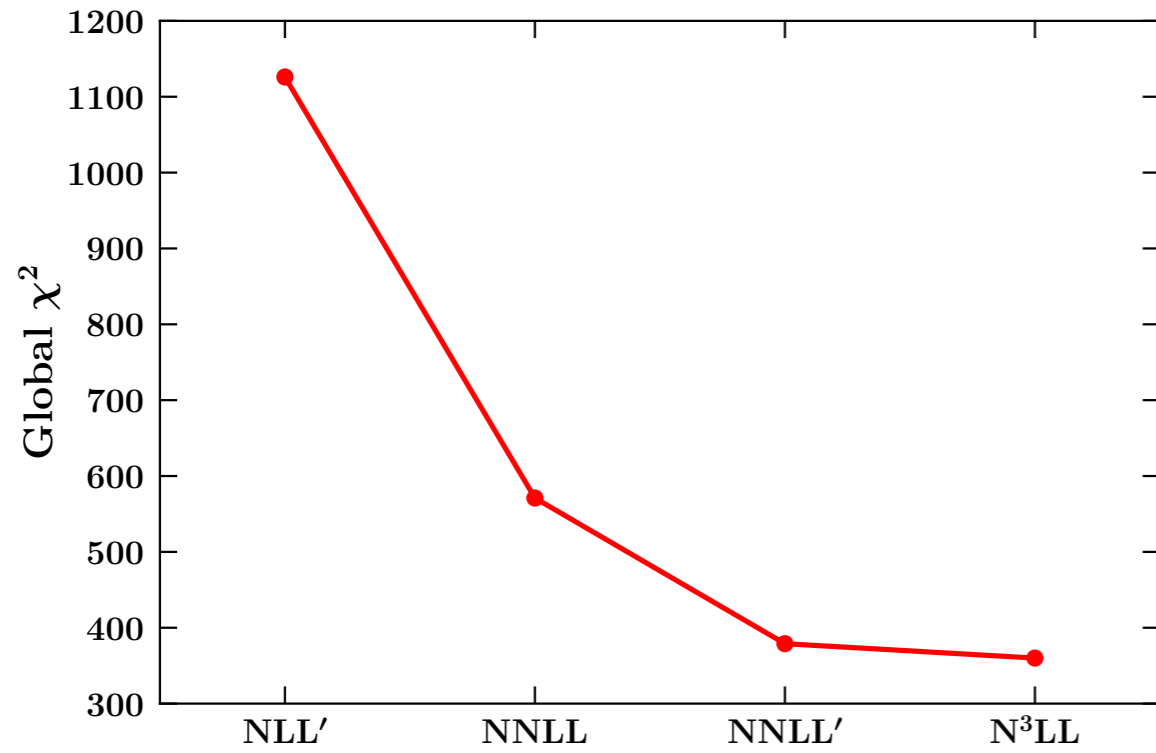
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# Impact of different logarithmic accuracies



# Test of $x$ -dependence

**Tried a fit at N3LL with Davies, Webber, Stirling (1985) NP parameterisation:**

$$f_{\text{NP}}^{\text{DWS}}(b_T, \zeta) = \exp \left[ -\frac{1}{2} \left( g_1 + g_2 \ln \left( \frac{\zeta}{2Q_0^2} \right) \right) b_T^2 \right]$$

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with and without ATLAS data

	Full dataset	No $y$ -differential data
Global $\chi^2/N_{\text{dat}}$	1.339	0.895
$g_1$	0.304	0.207
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—>  $x$ -dependence of TMDs at N3LL driven by ATLAS data