

JHU generator framework: New features for Higgs boson studies

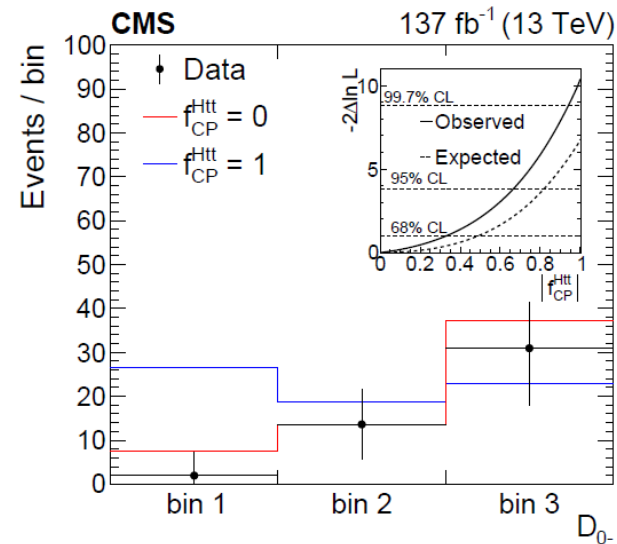
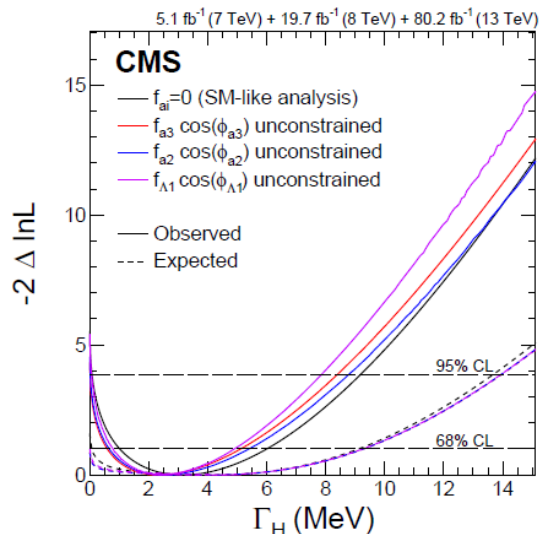
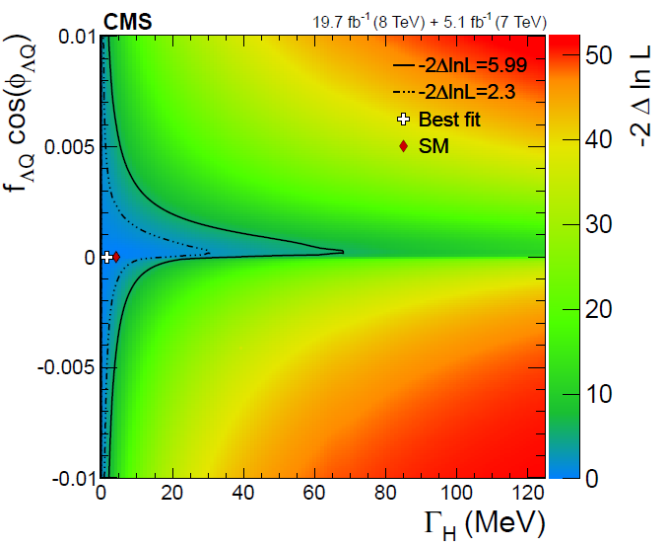
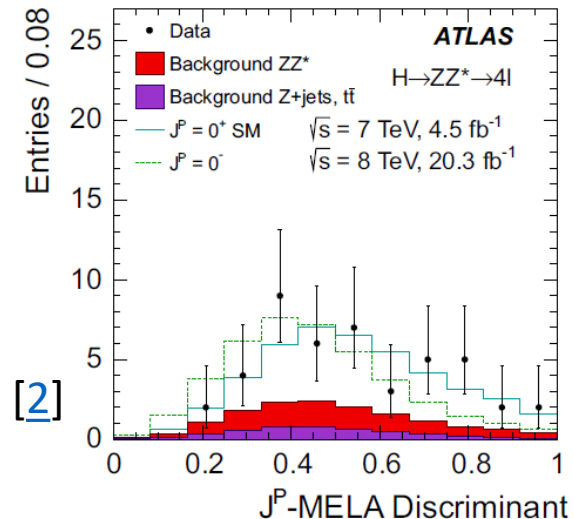
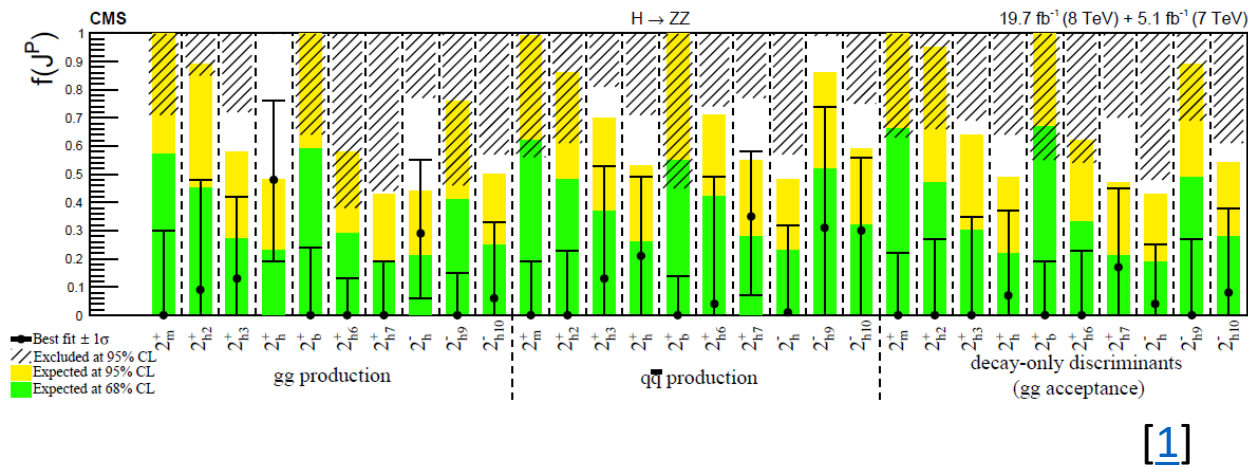
U. Sarica for

I. Anderson, S. Bolognesi, F. Caola, J. Davis, Y. Gao, A. Gritsan,
Z. Guo, C. Martin, K. Melnikov, R. Rontsch, H. Roskes, U. Sarica,
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Available packages

- JHUGenerator: Event generator for a spin-0, 1 or 2 resonance:
 - A key simulation tool in experiments to study resonances with BSM effects since Run 1
 - Has been a framework to study anomalous couplings/EFT effects in Run 2 as well
 - Focus on the anomalous/EFT couplings of the Higgs boson today
 - Interface to MCFM matrix elements for off-shell Higgs processes
- JHUGenMELA: Matrix element library with an object-oriented C++/Python interface and related numerical convenience tools for analysis, usable for
 - Reweighting of existing simulation via ratios of $|M|^2$.
 - Constructing discriminants for the analysis of couplings and background suppression
- JHUGenLexicon: Translation tool for between different EFT and Warsaw bases, and the JHUGen amplitude convention of anomalous couplings
- More details on tool availability and versions: <https://spin.pha.jhu.edu>
- References:
 - [Phys.Rev. D81 \(2010\) 075022](#), [arXiv:1001.3396](#)
 - [Phys.Rev. D86 \(2012\) 095031](#), [arXiv:1208.4018](#)
 - [Phys.Rev. D89 \(2014\) 035007](#), [arXiv:1309.4819](#)
 - [Phys.Rev. D94 \(2016\) 055023](#), [arXiv:1606.03107](#)
 - [Phys.Rev. D102 \(2020\) 056022](#), [arXiv:2002.09888](#)

Examples of past and recent uses in analyses



Available processes for Higgs simulation

Single resonance
production modes

$$gg \rightarrow H$$

QCD H production with 1 or 2 jets

VBF

$$q\bar{q} \rightarrow Z/W/\gamma H$$

(LO or NLO in QCD)

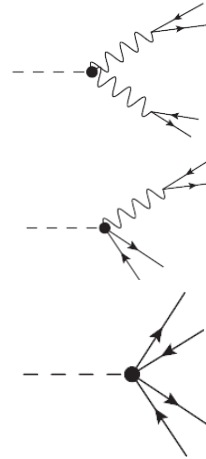
$$gg \rightarrow ZH$$

$$t\bar{t}H$$

(LO or NLO* in QCD)

$$tH$$

$$b\bar{b}H$$



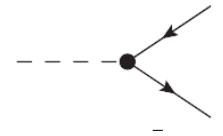
Single resonance decay modes
with production coming from
elsewhere (e.g. POWHEG, JHUGen)

$$H \rightarrow ZZ/Z\gamma^*/\gamma^*\gamma^*/WW \rightarrow 4f$$

$$H \rightarrow \gamma Z/\gamma^* \rightarrow \gamma 2f$$

$$H \rightarrow \gamma\gamma$$

$$H \rightarrow \tau\tau$$



Off-shell mixed modes

$ZZ/Z\gamma^*/\gamma^*\gamma^*/WW \rightarrow 4f$ production
via gluon fusion, and

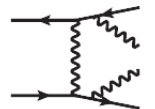
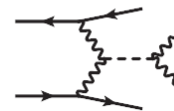
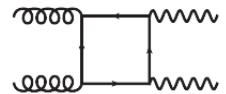
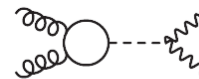
$ZZ/Z\gamma^*/\gamma^*\gamma^*/WW \rightarrow 4f + 2$ jets or V
electroweak production, both featuring

Higgs, VV continuum and

Higgs+continuum with their

interference, or an additional spin-0

resonance



Resonance couplings

HV_1V_2 : $A(HV_1V_2) = \frac{1}{v} \left\{ M_{V_1}^2 \left[g_1^{VV} + \frac{\kappa_1^{VV} q_1^2 + \kappa_2^{VV} q_2^2}{(\Lambda_1^{VV})^2} + \frac{\kappa_3^{VV} (q_1 + q_2)^2}{(\Lambda_0^{VV})^2} + \frac{2q_1 \cdot q_2}{M_{V_1}^2} g_2^{VV} \right] (\varepsilon_1 \cdot \varepsilon_2) - 2g_2^{VV} (\varepsilon_1 \cdot q_2)(\varepsilon_2 \cdot q_1) - 2g_4^{VV} \varepsilon_{\varepsilon_1 \varepsilon_2 q_1 q_2} \right\}$
 =2 for $V_1V_2 = ZZ, WW$ in SM at tree level additional HVV anomalous couplings

Vff : $\varepsilon_i^\mu(q_i) \rightarrow j_i^\mu = e \frac{\bar{\psi}_{f'} \gamma^\mu (g_L^{Vf'f} \omega_L + g_R^{Vf'f} \omega_R) \psi_f}{q_i^2 - M_V^2 + iM_V \Gamma_V}$ or $\frac{\bar{\psi}_{f'} \gamma^\mu (e_L^{V'f'f} \omega_L + e_R^{V'f'f} \omega_R) \psi_f}{q_i^2 - M_{V'}^2 + iM_{V'} \Gamma_{V'}}$
 SM Vff couplings additional $V'ff$ or contact interaction under $M_{V'} \rightarrow \infty$
 with a separate set of HV(')V(') couplings

Hff : $A(Hf\bar{f}) = -\frac{m_f}{v} \bar{\psi}_{f'} (\kappa_f + i\tilde{\kappa}_f \gamma_5) \psi_f$
 =1 in SM pseudoscalar anomalous coupling

The parametrization of the amplitude as above can easily be interpreted in other conventions, e.g. the EFT Higgs basis [1]:

$$\begin{aligned} \mathcal{L}_{\text{hvv}} = & \frac{h}{v} \left[(1 + \delta c_z) \frac{(g^2 + g'^2)v^2}{4} Z_\mu Z_\mu + c_{zz} \frac{g^2 + g'^2}{4} Z_{\mu\nu} Z_{\mu\nu} + c_{z\Box} g^2 Z_\mu \partial_\nu Z_{\mu\nu} + \tilde{c}_{zz} \frac{g^2 + g'^2}{4} Z_{\mu\nu} \tilde{Z}_{\mu\nu} \right. \\ & + (1 + \delta c_w) \frac{g^2 v^2}{2} W_\mu^+ W_\mu^- + c_{ww} \frac{g^2}{2} W_{\mu\nu}^+ W_{\mu\nu}^- + c_{w\Box} g^2 (W_\mu^- \partial_\nu W_{\mu\nu}^+ + \text{H.c.}) + \tilde{c}_{ww} \frac{g^2}{2} W_{\mu\nu}^+ \tilde{W}_{\mu\nu}^- \\ & + c_{z\gamma} \frac{e\sqrt{g^2 + g'^2}}{2} Z_{\mu\nu} A_{\mu\nu} + \tilde{c}_{z\gamma} \frac{e\sqrt{g^2 + g'^2}}{2} Z_{\mu\nu} \tilde{A}_{\mu\nu} + c_{\gamma\Box} g g' Z_\mu \partial_\nu A_{\mu\nu} \\ & \left. + c_{\gamma\gamma} \frac{e^2}{4} A_{\mu\nu} A_{\mu\nu} + \tilde{c}_{\gamma\gamma} \frac{e^2}{4} A_{\mu\nu} \tilde{A}_{\mu\nu} + c_{gg} \frac{g_s^2}{4} G_{\mu\nu}^a G_{\mu\nu}^a + \tilde{c}_{gg} \frac{g_s^2}{4} G_{\mu\nu}^a \tilde{G}_{\mu\nu}^a \right], \end{aligned}$$

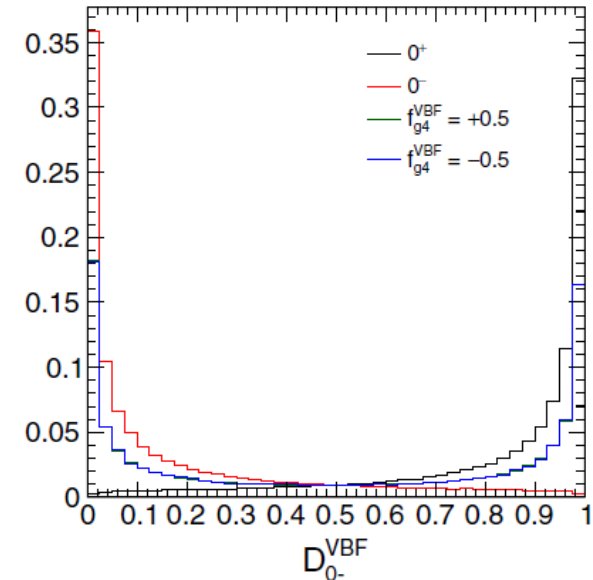
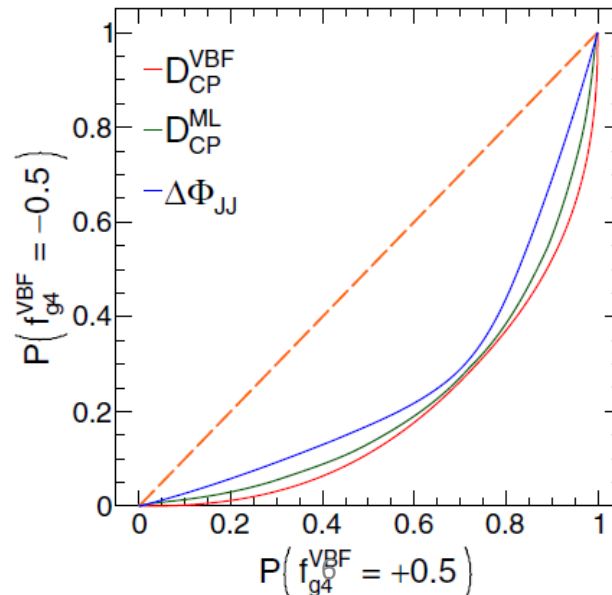
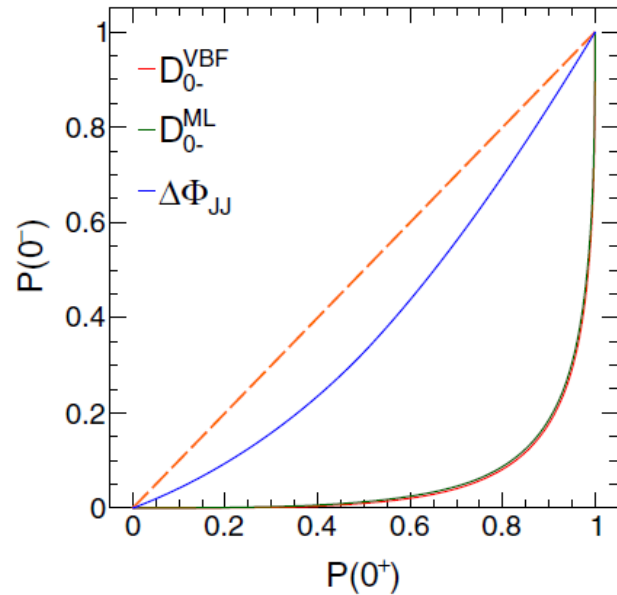
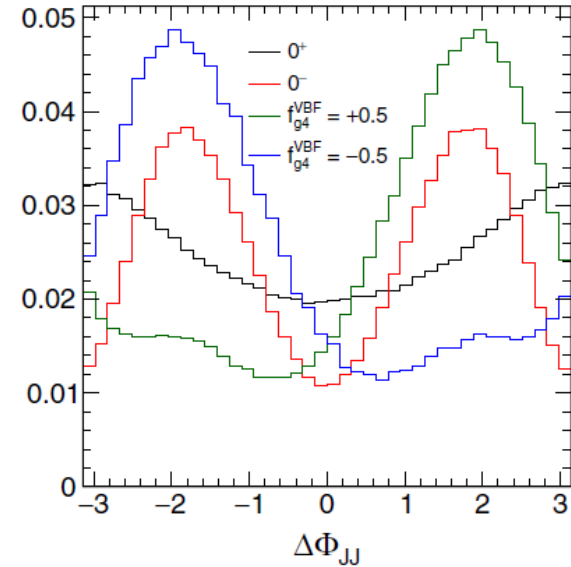
Case study: VBF CP analysis (HVV)

- For a VBF CP analysis, one could examine $\Delta\phi_{JJ}$ or a multivariate discriminant using more kinematic information
- For optimal analysis, construct MELA discriminants of form

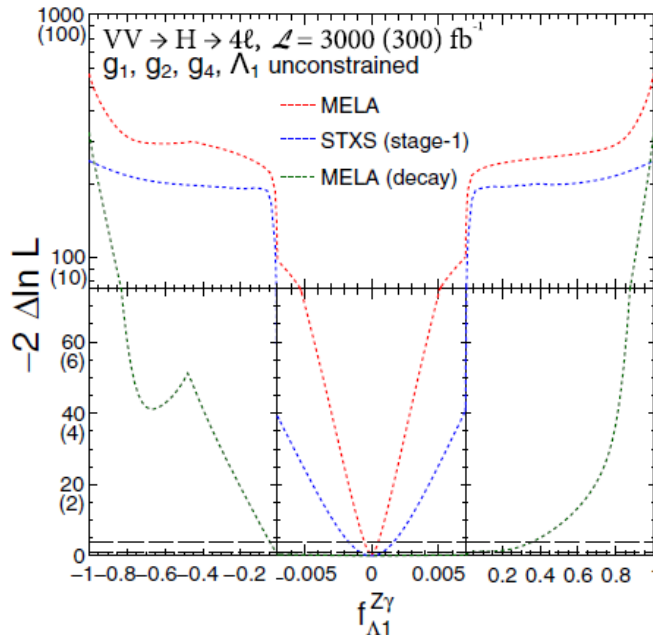
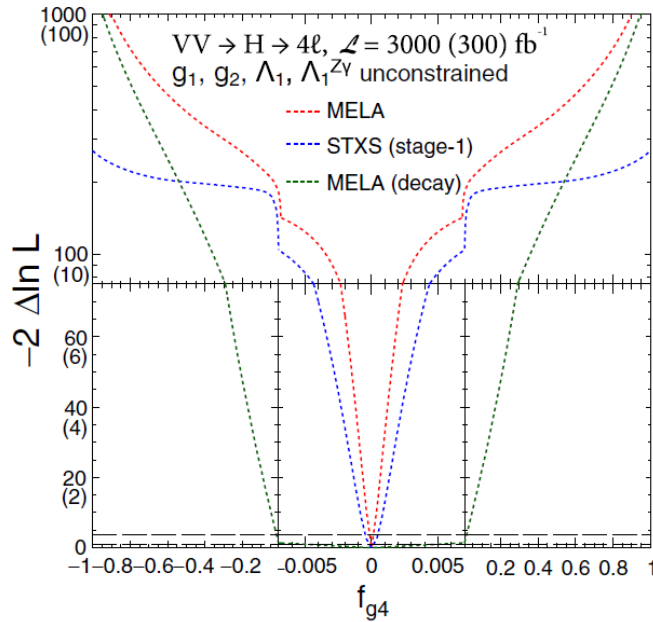
$$\mathcal{D}_{\text{alt}}(\Omega) = \frac{\mathcal{P}_{\text{sig}}(\Omega)}{\mathcal{P}_{\text{sig}}(\Omega) + \mathcal{P}_{\text{alt}}(\Omega)} \quad \mathcal{D}_{\text{int}}(\Omega) = \frac{\mathcal{P}_{\text{int}}(\Omega)}{2\sqrt{\mathcal{P}_{\text{sig}}(\Omega) \times \mathcal{P}_{\text{alt}}(\Omega)}}$$

(here, sig= 0^- , alt= 0^+ [SM], int [CP]=interference extracted using $f_{g4}^{\text{VBF}} = +0.5$)

- MELA discriminants contain full kinematic information and can outperform single-variable or other multivariate techniques



Full case study: HVV coupling constraints from 4ℓ

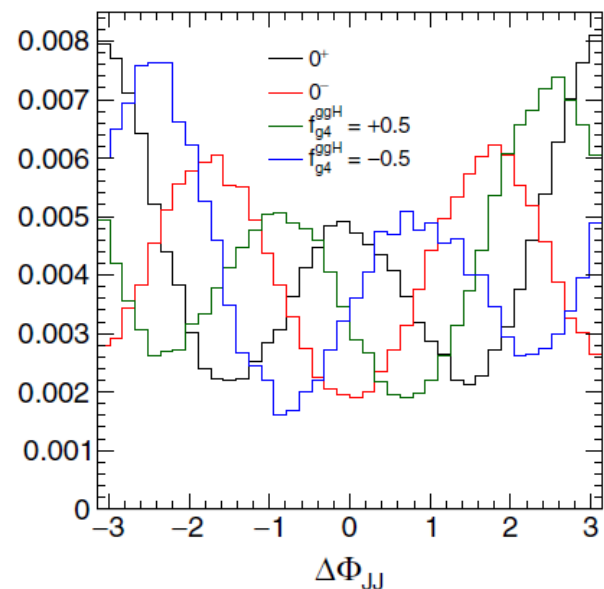
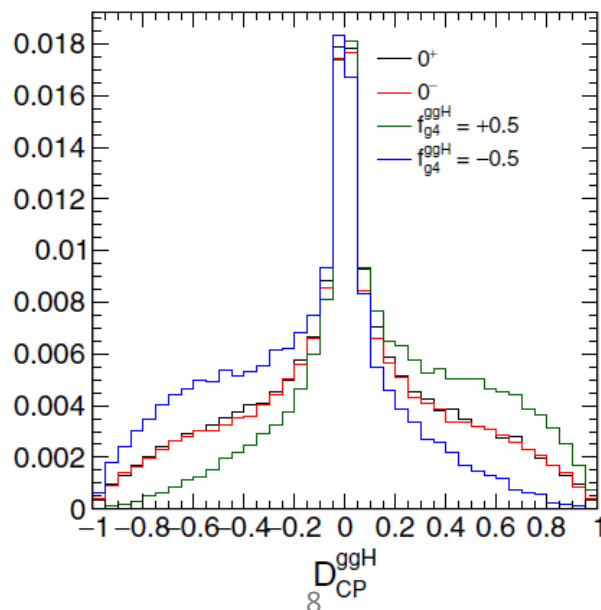
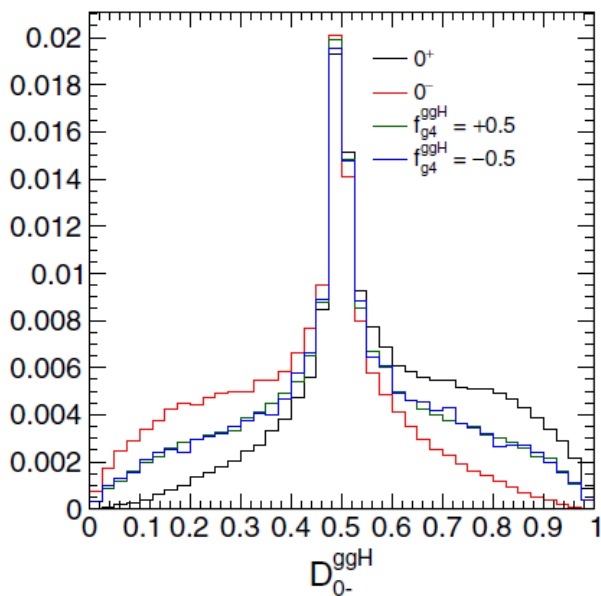
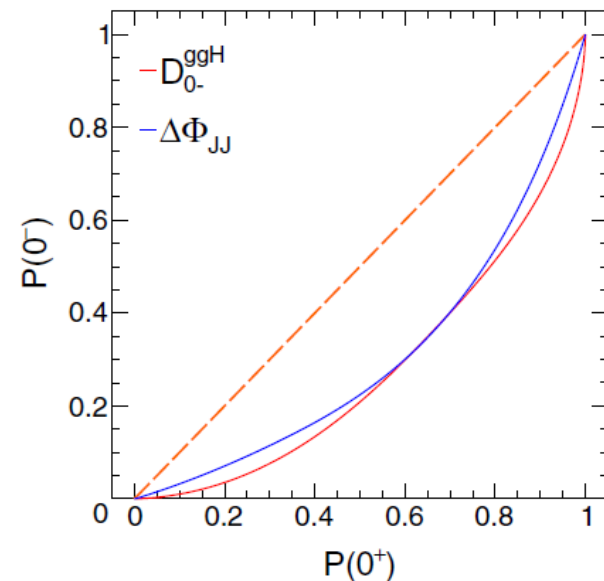


- Case study of constraints on the HVV couplings using VBF $H \rightarrow 4\ell$ events
- Different interference components can be obtained via dedicated simulation or reweighting existing simulated events
- Comparison of constraints at 300 and 3000 fb^{-1} made using
 - a) full set of MELA discriminants D_{alt} and D_{int} with production and decay information
 - b) set of MELA discriminants D_{alt} and D_{int} with decay information only
 - c) Binning of observables following the STXS 1.1 binning
- Analysis with a MELA-based binning in 4ℓ with production and decay information combined can be more sensitive to couplings than a decay-only analysis or STXS 1.1 binning

Case study: ggH+2 jets CP analysis (Hgg)

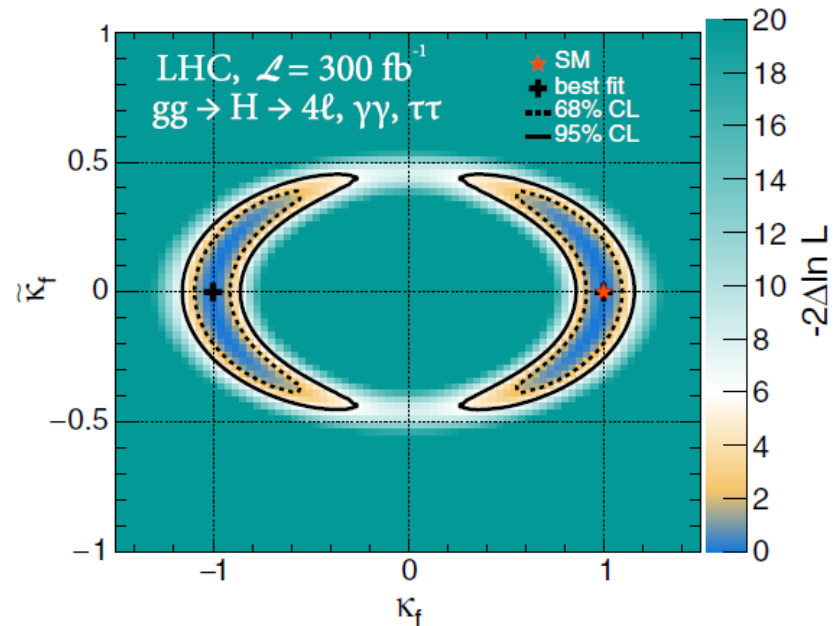
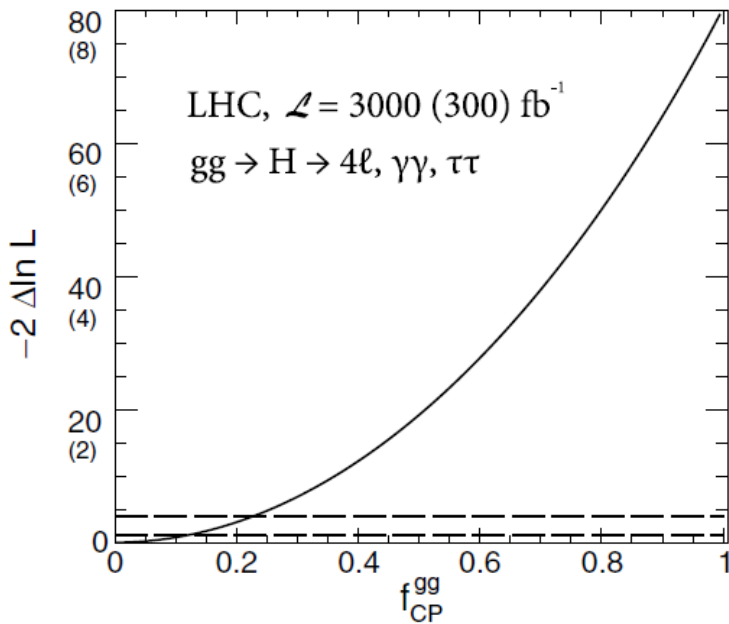
→ Similar to a VBF CP analysis, one could examine $\Delta\phi_{JJ}$ or construct MELA discriminants
 → Only quark-initiated MEs are sensitive, so discriminants are constructed only out of those.

→ MELA discriminants could outperform $\Delta\phi_{JJ}$



Full case study: Hgg coupling constraints

- Shown constraints are obtained using 4ℓ events, but expected yields are scaled to include $\gamma\gamma$ and $\tau\tau$ channels to illustrate the effect of combination
- No statistical sensitivity on the modification of the loop to Higgs p_T , ignored at the moment in the projections.



Translation of couplings

→ Conversion between the couplings in the JHUGen amplitude conventions and the EFT Higgs basis, or the Warsaw basis, is possible using the JHUGenLexicon:

$$\begin{aligned}
 \delta c_z &= \frac{1}{2} g_1^{ZZ} - 1, & c_{zz} &= -\frac{2s_w^2 c_w^2}{e^2} g_2^{ZZ}, & c_{z\Box} &= \frac{M_Z^2 s_w^2}{e^2} \frac{\kappa_1^{ZZ}}{(\Lambda_1^{ZZ})^2}, & \tilde{c}_{zz} &= -\frac{2s_w^2 c_w^2}{e^2} g_4^{ZZ}, \\
 \delta c_w &= \frac{1}{2} g_1^{WW} - 1, & c_{ww} &= -\frac{2s_w^2}{e^2} g_2^{WW}, & c_{w\Box} &= \frac{M_W^2 s_w^2}{e^2} \frac{\kappa_1^{WW}}{(\Lambda_1^{WW})^2}, & \tilde{c}_{ww} &= -\frac{2s_w^2}{e^2} g_4^{WW}, \\
 c_{Z\gamma} &= -\frac{2s_w c_w}{e^2} g_2^{Z\gamma}, & \tilde{c}_{Z\gamma} &= -\frac{2s_w c_w}{e^2} g_4^{Z\gamma}, & c_{\gamma\Box} &= \frac{s_w c_w}{e^2} \frac{M_Z^2}{(\Lambda_1^{Z\gamma})^2} \kappa_2^{Z\gamma}, \\
 c_{\gamma\gamma} &= -\frac{2}{e^2} g_2^{\gamma\gamma}, & \tilde{c}_{\gamma\gamma} &= -\frac{2}{e^2} g_4^{\gamma\gamma}, & c_{gg} &= -\frac{2}{g_s^2} g_2^{gg}, & \tilde{c}_{gg} &= -\frac{2}{g_s^2} g_4^{gg}.
 \end{aligned}$$

Higgs basis coefficients, see [1]

$$d_1^{\gamma} = 1 + (g_2^{\gamma\gamma} - g_2^{ZZ}) c_w^2 + g_2^{Z\gamma} \left(\frac{c_w}{s_w} - 2s_w c_w \right),$$

$$d_4^{\gamma} = (g_4^{\gamma\gamma} - g_4^{ZZ}) c_w^2 + g_4^{Z\gamma} \left(\frac{c_w}{s_w} - 2s_w c_w \right),$$

$$d_1^Z = 1 - 2 \frac{s_w^2 c_w^2}{c_w^2 - s_w^2} (g_2^{\gamma\gamma} - g_2^{ZZ}) - 2s_w c_w g_2^{Z\gamma} - \frac{M_Z^2}{2(c_w^2 - s_w^2)} \frac{\kappa_1^{ZZ}}{(\Lambda_1^{ZZ})^2},$$

$$d_2^Z = d_3^Z = 1 - \frac{s_w^2}{c_w^2 - s_w^2} (g_2^{\gamma\gamma} - g_2^{ZZ}) - \frac{s_w}{c_w} g_2^{Z\gamma} - \frac{M_Z^2}{2(c_w^2 - s_w^2)} \frac{\kappa_1^{ZZ}}{(\Lambda_1^{ZZ})^2},$$

$$d_4^Z = -\frac{s_w^2}{c_w^2} d_4^{\gamma},$$

$$d^{ZZWW} = \frac{c_w^2}{s_w^2} (2d_2^Z - 1), \quad d^{Z\gamma WW} = \frac{c_w}{s_w} d_2^Z.$$

aT/QGC contributions, see [2]

Warsaw basis coefficients, see [3]

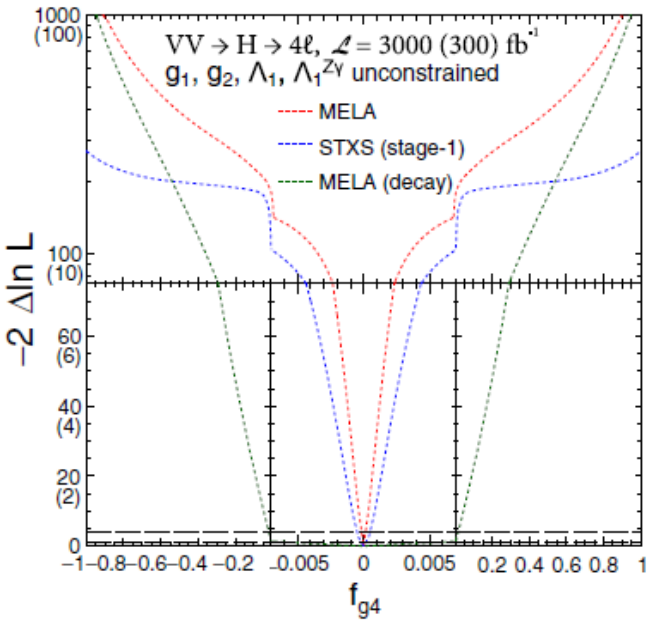
$$g_4^{ZZ} = -2 \frac{v^2}{\Lambda^2} \left(s_w^2 w_{\phi\tilde{B}} + c_w^2 w_{\phi\tilde{W}} + s_w c_w w_{\phi B\tilde{W}} \right),$$

$$g_4^{\gamma\gamma} = -2 \frac{v^2}{\Lambda^2} \left(c_w^2 w_{\phi\tilde{B}} + s_w^2 w_{\phi\tilde{W}} - s_w c_w w_{\phi B\tilde{W}} \right),$$

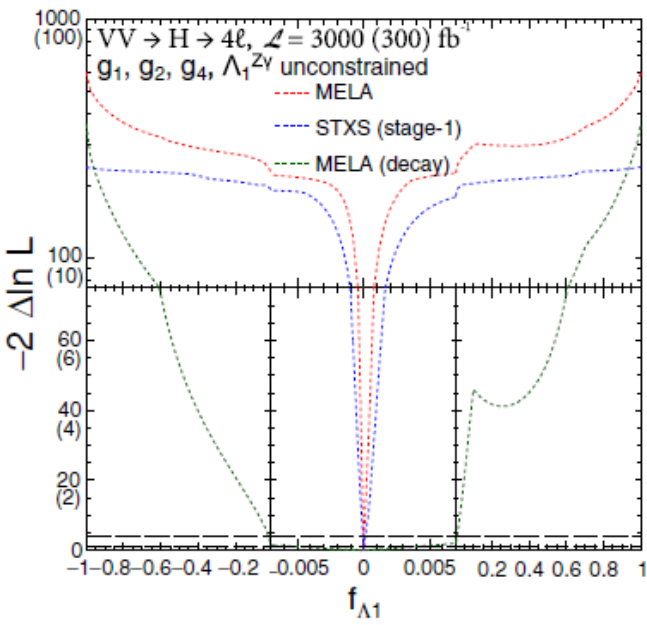
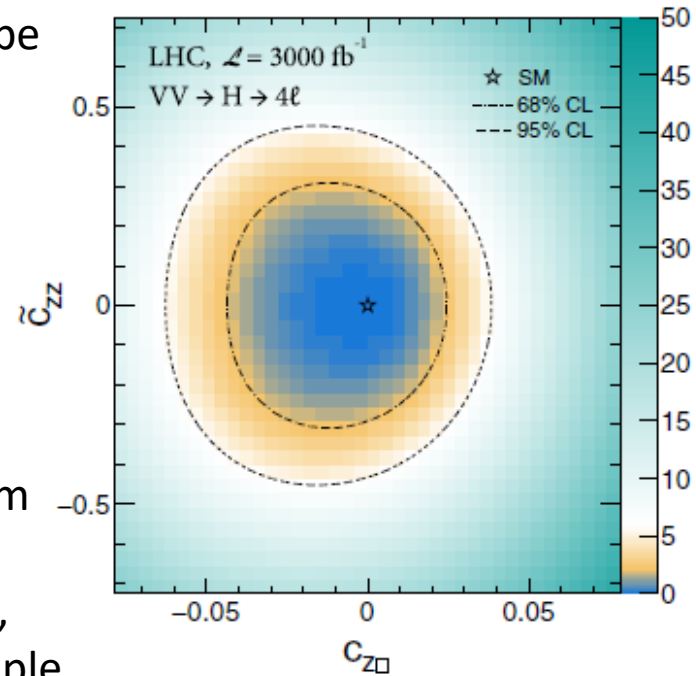
$$g_4^{Z\gamma} = -2 \frac{v^2}{\Lambda^2} \left(s_w c_w (w_{\phi\tilde{W}} - w_{\phi\tilde{B}}) + \frac{1}{2} (s_w^2 - c_w^2) w_{\phi B\tilde{W}} \right)$$

$$g_4^{gg} = -2 \frac{v^2}{\Lambda^2} w_{\phi\tilde{G}}.$$

Translation of couplings: HVV example



→ Constraints shown in red for either of the couplings can be interpreted in a joint multi-parameter form.



→ The multi-parameter form can be translated to other bases using JHUGenLexicon, here Higgs basis as an example.

Off-shell (mixed) processes

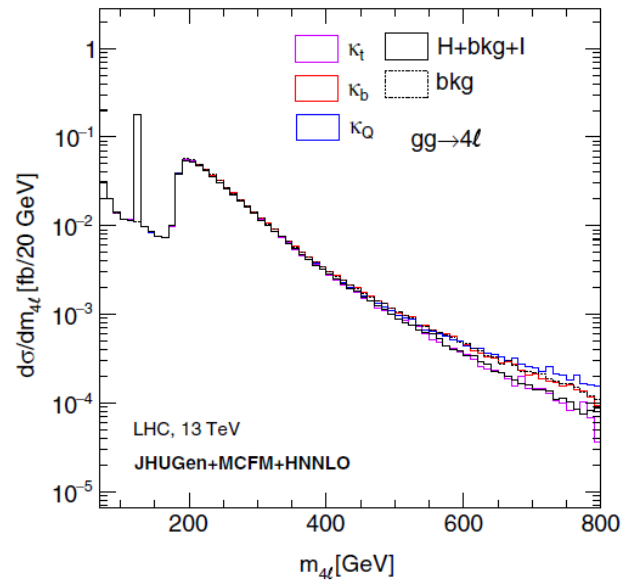
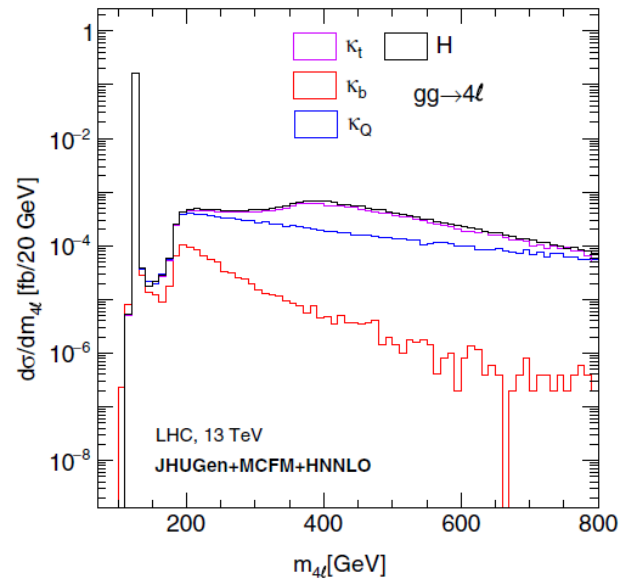
→ We utilize the MCFM matrix elements for the gg -initiated 0-jet VV production (since 2015), or the EW VV production with 2 jets (since 2017):

→ Added additional anomalous couplings of the Higgs (HVV, Hgg, or $Hf\bar{f}$), an additional t'/b' loop in gluon fusion case, also triple and quartic V coupling contributions in EW VV production with 2 jets continuum amplitudes for a full EFT treatment.

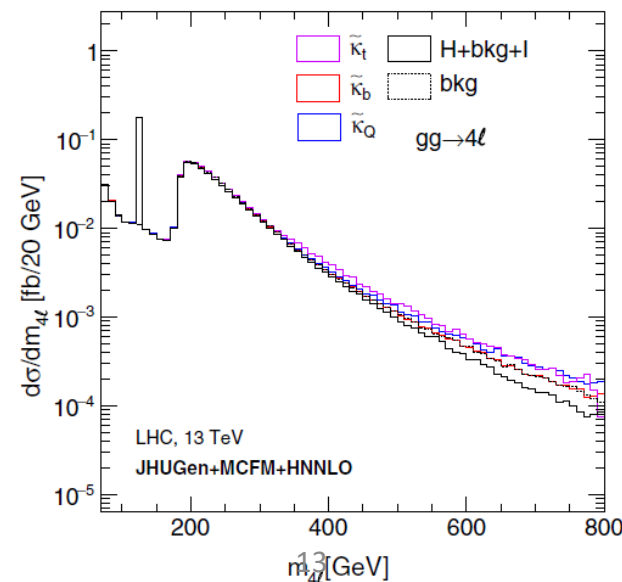
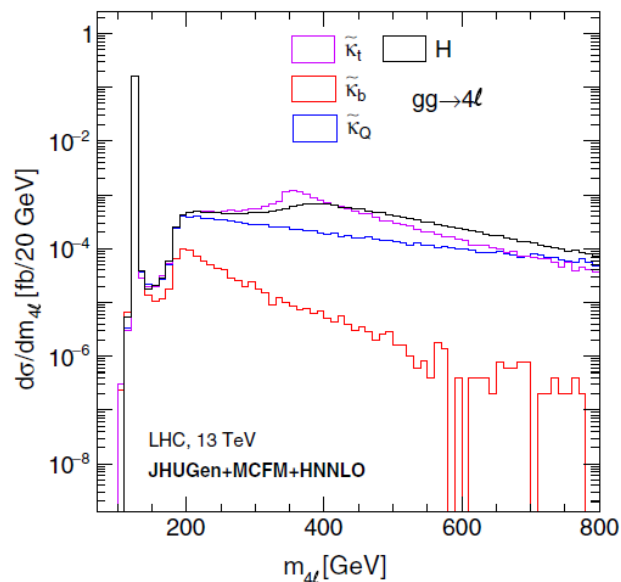
→ Also added a second resonance to the amplitudes with equivalent but separate sets of couplings

→ Expanded the EW MEs to allow separate generation for each initial-final state configuration, and the triple-V production MEs to also feature leptonic decays of the associated V for ME reweighting purposes.

Off-shell processes: Hgg couplings

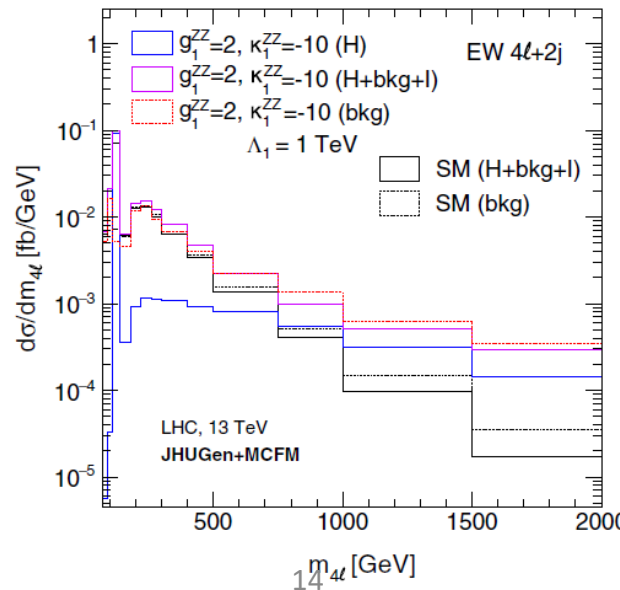
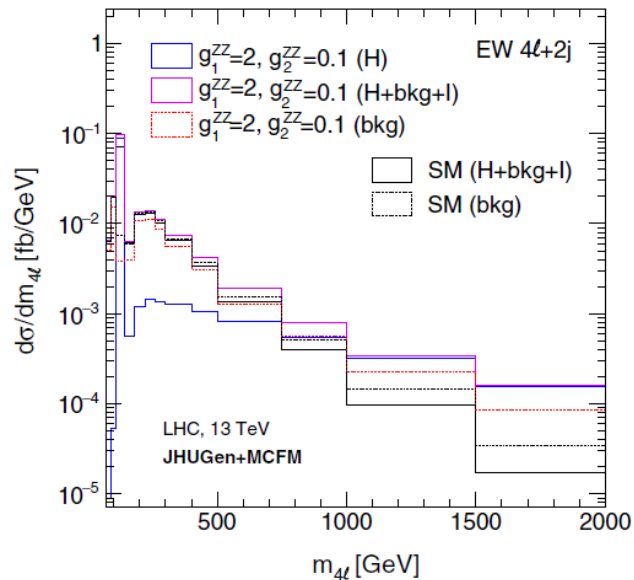
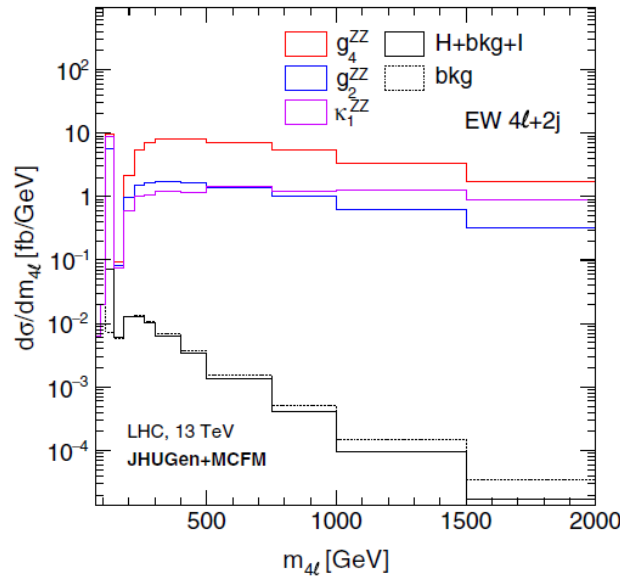
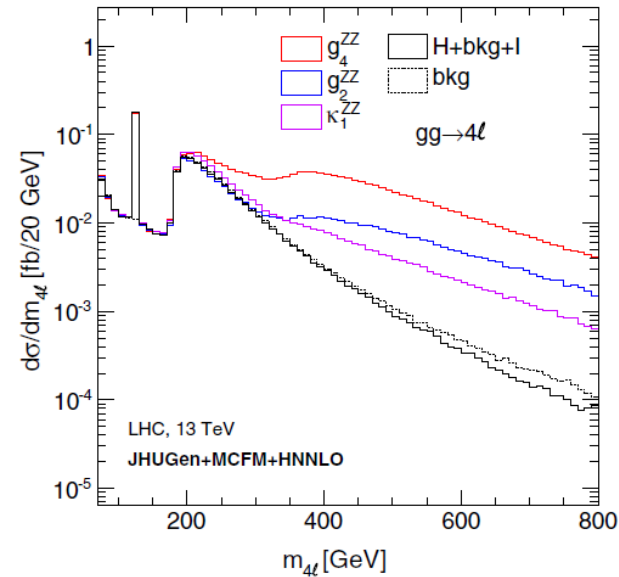


→ Individual top (κ_t) and bottom (κ_b) contributions for scalar Higgs scenario (+ full continuum bkg), as well as their sum, or a point-like interaction ($\kappa_Q \sim g_2^{gg}$)



→ Individual top ($\tilde{\kappa}_t$) and bottom ($\tilde{\kappa}_b$) contributions for pseudoscalar Higgs scenario (+ full continuum bkg), as well as their sum, or a point-like interaction ($\tilde{\kappa}_Q \sim g_4^{gg}$)

Off-shell processes: HVV couplings

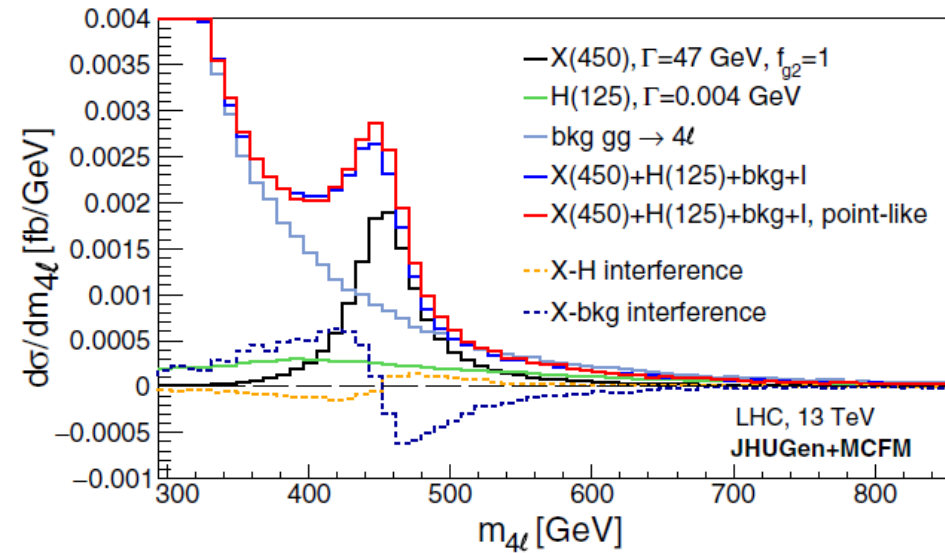


→ Changes in the VV invariant mass shape due to HVV anomalous couplings in $gg \rightarrow 4\ell$ or $4\ell + 2j$ EW production (assuming no changes in the continuum amplitudes)

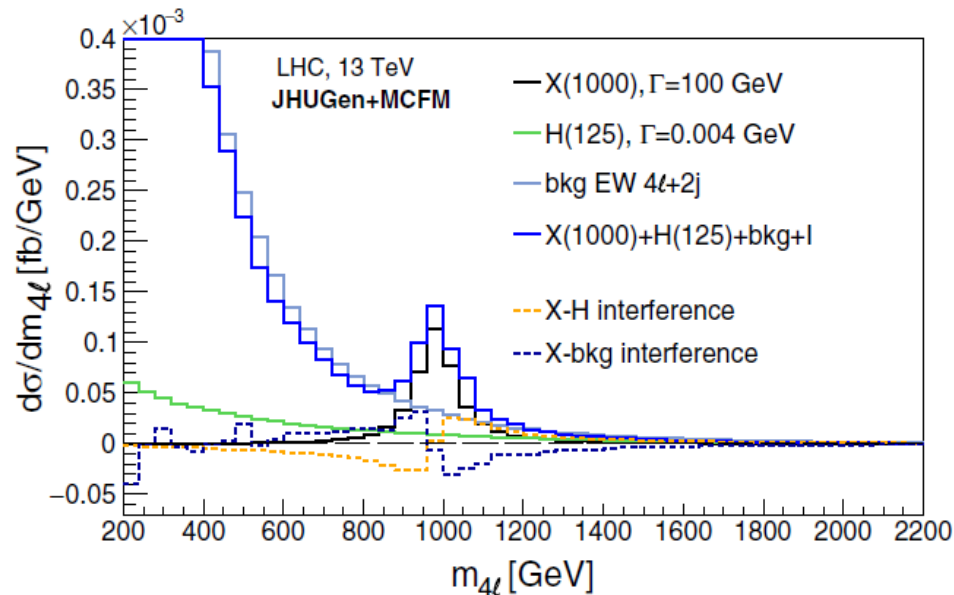
→ m_{VV} shape with two different exemplary HVV anomalous couplings in $4\ell + 2j$ EW production, assuming triple and quartic gauge couplings are affected by the changes in Higgs couplings (see [1] for correspondence details)

→ Different components can be obtained using dedicated simulation or reweighting

Off-shell processes: Second resonance



→ Illustrating different spin-0 X masses and couplings in the $gg \rightarrow 4\ell$ (top) or EW $4\ell + 2j$ (bottom) processes.

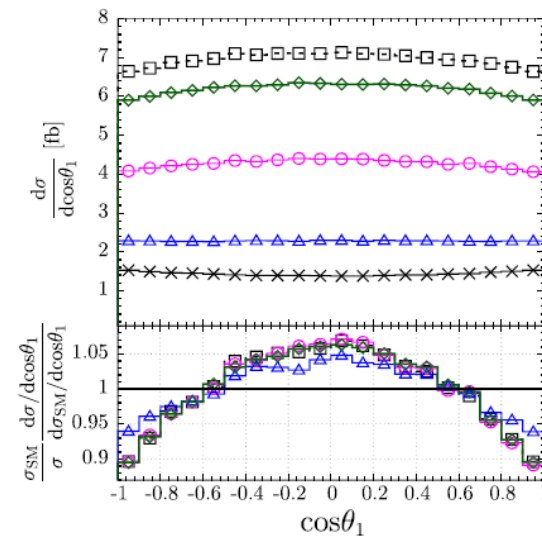
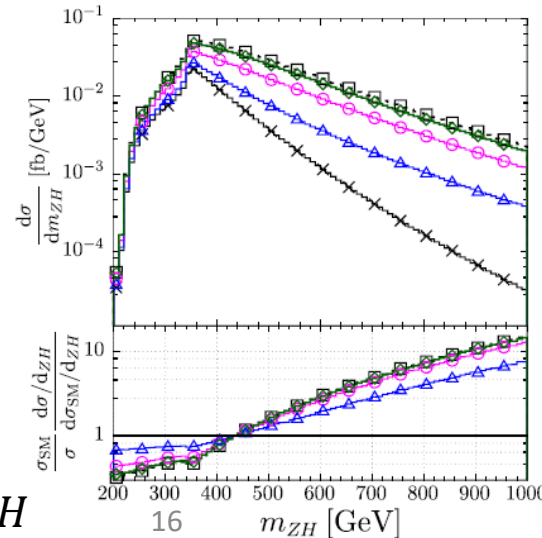
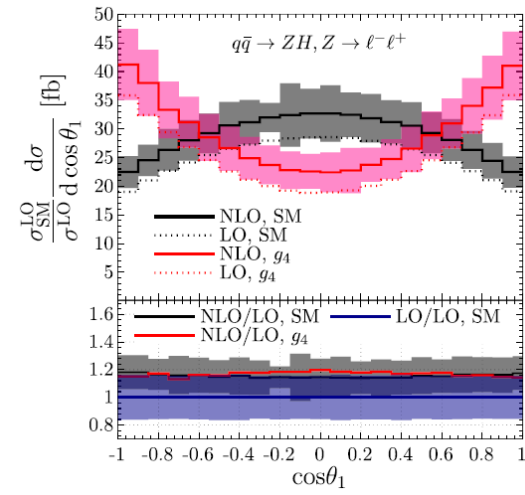
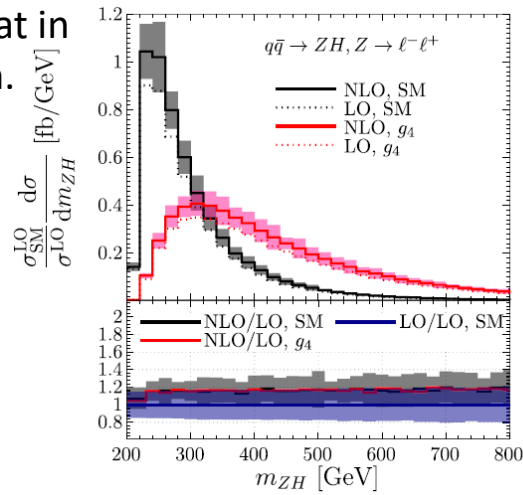
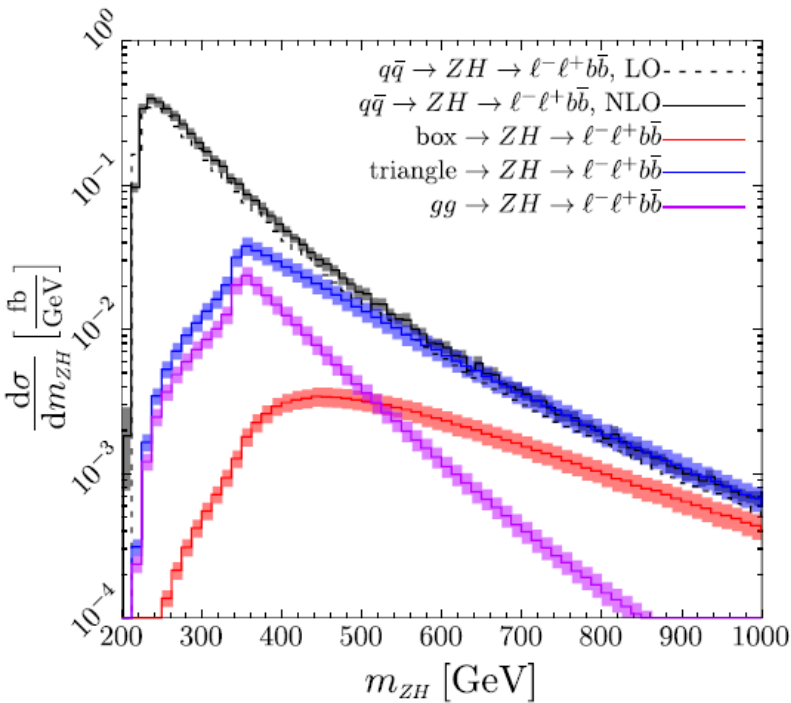


→ Full interference with the first resonance (H) and the continuum amplitude are implemented, allowing expansion of additional resonance searches in ZZ and WW final states.

New features in VH production

→ VH production at LO in QCD had been available in our package for several years. We now add NLO production capability, and the $gg \rightarrow ZH$ contribution.

LO → NLO corrections are relatively flat in relevant kinematics within a 5% margin.



Shown are different $\kappa, \tilde{\kappa}$ combinations in $gg \rightarrow ZH$

Summary

- Presented a coherent and self-contained framework to simulate, measure and analyze Higgs boson properties
 - Provide an amplitude implementation to analyze and translate the couplings to any preferred basis for analysis – anomalous couplings, EFT, pseudo-observables...
 - Exemplify ME discriminant-based analyses for optimal observables in the analysis of the couplings
- Interface to a modified version of MCFM matrix elements to model changes in kinematics due to Higgs and continuum amplitudes under different anomalous couplings/EFT treatment
 - Also adding the capability to simulate a second resonance for additional resonance searches
- Expand the VH production to include NLO treatment with anomalous couplings
 - NLO in $q\bar{q} \rightarrow VH$, and adding $gg \rightarrow ZH$
- See our latest [paper](#), or <https://spin.pha.jhu.edu> for more details

Thank you!