# Controlling uncertainties of the IAM extension of EFTs Alexandre Salas and Felipe Llanes-Estrada (UCM)



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 High energy scattering: V << T, Feynman diagrams, Madgraph, etc.



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• Low energy respect to new physics (strongly interacting?  $V \sim T$  requires resummation)

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# Expand partial wave amplitudes

$$T_I(s,t,u) = 16\eta\pi \sum_{J=0}^{\infty} (2J+1)t_{IJ}(s)P_J(\cos\theta_s)$$

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$$t_{IJ}(s) \simeq \underbrace{t_{0}}_{O(s)} + \underbrace{t_{1}}_{O(s^{2})} + \dots$$
$$(\text{typical HEFT expansion})$$

# Inverse Amplitude Method

$$rac{1}{t}\simeq rac{1}{t_0+t_1}\simeq rac{1}{t_0}-rac{t_1}{t_0^2} \implies \left| t^{IAM}\simeq rac{t_0^2}{t_0-t_1} 
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Advantage: for  $s > s_{th}$ ,

$$\operatorname{Im} rac{1}{t_{IJ}(s)} = -\sigma(s) \simeq -1$$

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Exact in IAM

Only order by order in EFT

$$\operatorname{Im} t_{\mathbf{1}}(s) = \sigma(s)|t_{\mathbf{0}}(s)|^2$$

Why would anyone care?

▶ EFT reliable only near threshold



#### Prediction of resonances from HEFT



#### Much used in hadron physics to obtain resonances



(This is an IAM prediction from threshold data, not a fit)

#### Use its dispersive derivation: 2010.13709



# Master formula is a dispersion relation for $G(s) \equiv rac{t_0^2(s)}{t(s)}$

$$G(s) = G(0) + G'(0)s + \frac{1}{2}G''(0)s^{2} + PC(G) + \frac{s^{3}}{\pi} \int_{RC} ds' \frac{\operatorname{Im} G(s')}{s'^{3}(s'-s)} + \frac{s^{3}}{\pi} \int_{LC} ds' \frac{\operatorname{Im} G(s')}{s'^{3}(s'-s)}$$

#### Dispersion relation: approximations



#### Sources of uncertainty

Neglected pole contributions of t<sup>-1</sup>: subthreshold Adler zeroes and CDD zeroes of t.

▶ Inelasticities due to *KK* (*hh* in HEFT),  $4\pi$ , etc.

•  $\mathcal{O}(p^4)$  truncation of subtraction constants.

• Left cut approximation  $Im \ G \simeq -Im \ t_1$ .

# Adler zeroes of t near threshold



 $t_0 + t_1 = a + bs + cs^2$  vanishes near s = -a/b

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#### 0712.2763



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$$t_{\mathrm{IAM}} = \frac{t_0^2}{t_0 - t_1} \rightarrow \frac{t_0^2}{t_0 - t_1 + \frac{s}{s - s_c} \mathrm{Re}(t_1)}$$

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• Hadrons:  $\pi\pi \to \pi\pi$  couples to *KK* 

$$\operatorname{Im} \frac{1}{t_{\pi\pi}} \to -\sigma_{\pi\pi} \left( 1 + \frac{\sigma_{K\bar{K}}}{\sigma_{\pi\pi}} \frac{|t_{\pi\pi \to K\bar{K}}|^2}{|t_{\pi\pi \to \pi\pi}|^2} \right)$$

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suppressed by phase-space  $rac{\sigma_{Kar{K}}}{\sigma_{\pi\pi}}$  and low inelasticity in  $t_{\pi\pi o Kar{K}}$ 

- ▶ In HEFT only inelasticity in  $\omega\omega hh$  (actually zero in SM)
- We can use the coupled channel IAM directly or to estimate uncertainty in elastic IAM

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Inelastic 2-body	$(\sqrt{s}/(4\pi f_{\pi}))^4$	10 <sup>-3</sup>	Yes



 Difference with SMEFT: here, in ChPT or HEFT, additional particles \*not\* suppressed by the chiral counting. But phase space helps.



# In hadron physics, (with elastic and 4- $\pi$ inelastic amplitudes taken as similar )

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# $O(p^4)$ truncation



Estimate based on size of NNLO counterterms ( $\implies$  subtraction constants) from Resonance Effective Field Theory

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# Approximate left cut



Need to check  $\int_{LC} ds' \frac{\operatorname{Im} G + \operatorname{Im} t_1}{s'^3(s' - s)} .$  *i.e.*, failure of IAM's  $Im \ G = -Im \ t_1$ over the left cut

#### Approximate left cut

Split interval in 3:

- ► Low-|s| (ChPT/HEFT  $\checkmark$ )  $|s|^{\frac{1}{2}} < 470 \text{MeV}.$
- Intermediate-|s|: Match to ChPT + natural-size counterterm. Currently studying LC parameterizations from GKPY eqns.
- High -|s|: Sugawara-Kanazawa relates it to right cut: Regge asymptotics there. Far from RC anyway.

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$O(p^4)$ truncation	$(\sqrt{s}/(4\pi f_\pi))^4$	$10^{-2}$	Yes
Left Cut	$(\sqrt{s}/(4\pi f_\pi))^4$	0.17	Perhaps

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- To make it more useful for BSM searches

# Thank You!

This project has received funding from the European Union's Horizon 2020 research and innovation programme under grant agreement No 824093; grants MINECO:FPA2016-75654-C2-1-P, FPA2016-77313-P MICINN: PID2019-108655GB-I00, PID2019-106080GB-C21, PID2019-106080GB-C22 (Spain); UCM research group 910309 and the IPARCOS institute; and the VBSCAN COST action CA16108.

