

# Divers topics on EW virtual corrections

**E. Richter-Was (IF UJ, Krakow)**

- **Predictions for  $\Delta A_4$  (EW) with experimental binning. TauSpinner + Dizet 6.45**
- **Factorizing QCD from EW corrections.**
- **Draft v05 status**

# Status of the YR draft (v04)

 Recently updated

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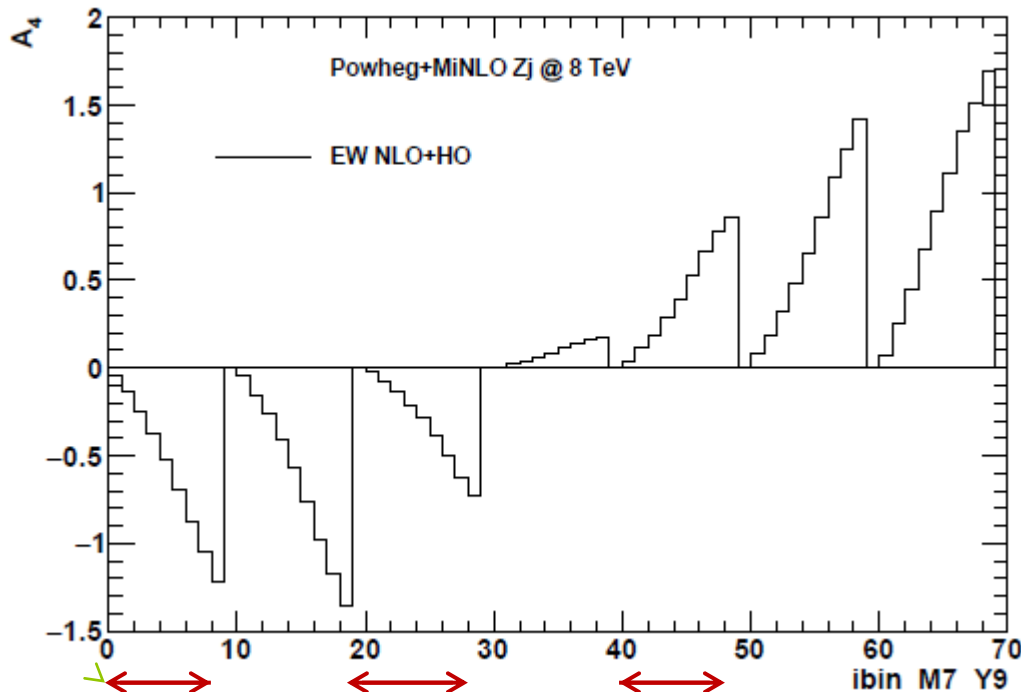
# The $\Delta A_4(\text{EW})$ in the full phase-space

- Experimental binning

Table 29: The binning in  $M_{bin}, Y_{bin}$  for tabulating  $A_4$  sensitivity to  $\sin^2 \theta_W^{eff}$ .

Observable	Bin thresholds
$M_{bin}$	[60, 66, 76, 86, 96, 106, 116, 150] GeV
$Y_{bin}$	[0.0, 0.4, 0.8, 1.2, 1.6, 2.0, 2.4, 2.8, 3.2, 3.6]

Should start at 50 GeV, but I have no events below 60 GeV



$m_{ll} = 60-66 \text{ GeV}$   
E. Richter-Was, IF JU

76-86 GeV

96-106 GeV

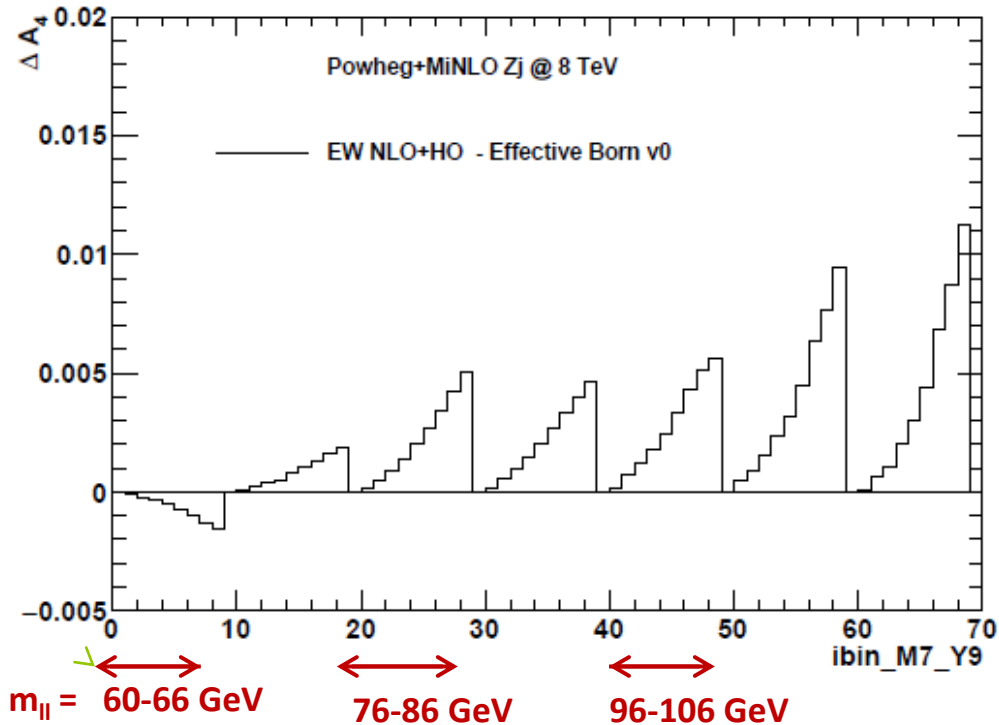
LHC EWWG meeting, 27.03.2020

Dizet v6.45

Parameter	$(\alpha(0), G_\mu, M_Z)$ $\alpha(0) v_0$
$M_Z$ (GeV)	91.1876
$1/\alpha(M_Z)$	0.0077549256
$\alpha(M_Z)$	128.9503020
$G_\mu$ (GeV <sup>-2</sup> )	$1.1663787 \cdot 10^{-5}$
$M_W$ (GeV)	80.358935
$s_W^2$	0.223401084
$\sin^2 \theta_{eff}^l$	0.231499
$\sin^2 \theta_{eff}^u$	0.231392
$\sin^2 \theta_{eff}^d$	0.231265
$\sin^2 \theta_{eff}^b$	0.232733

# The $\Delta A_4(\text{EW})$ in the full phase-space

- EW NLO+HO vs Effective Born v0



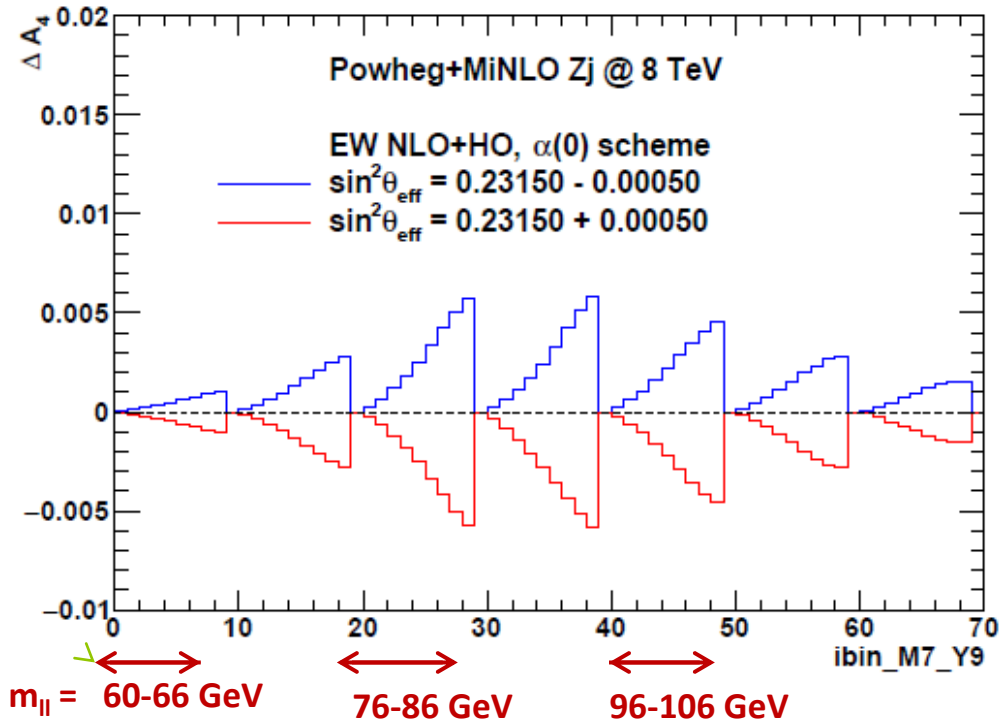
Effective Born v0
$\alpha = 1/128.950302$ $s_W^2 = 0.231499$
$\rho_{eff} = 1.0$

Dizet v6.45

Parameter	$(\alpha(0), G_\mu, M_Z)$ $\alpha(0) \text{ v0}$
$M_Z$ (GeV)	91.1876
$1/\alpha(M_Z)$	0.0077549256
$\alpha(M_Z)$	128.9503020
$G_\mu$ (GeV <sup>-2</sup> )	$1.1663787 \cdot 10^{-5}$
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$\sin^2 \theta_{eff}^b$	0.232733

# The $\Delta A_4(\text{EW})$ in the full phase-space

- EW NLO+HO vs Effective Born v0



Effective Born v0
$\alpha = 1/128.9503022$ $s_W^2 = 0.231499$
$s_W^2 = s_W^2 + \delta_V$

EW NLO+HO,  $\alpha(0)$  scheme

$$v_\ell = (2 \cdot T_3^\ell - 4 \cdot q_\ell \cdot (s_W^2 \cdot \mathcal{K}_\ell(s,t) + \delta_V)) / \Delta,$$

$$v_f = (2 \cdot T_3^f - 4 \cdot q_f \cdot (s_W^2 \cdot \mathcal{K}_f(s,t) + \delta_V)) / \Delta,$$

$$vv_{\ell f} = \frac{1}{v_\ell \cdot v_f} [(2 \cdot T_3^\ell)(2 \cdot T_3^f) - 4 \cdot q_\ell \cdot (s_W^2 + \mathcal{K}_f(s,t) + \delta_V)(2 \cdot T_3^\ell) - 4 \cdot q_f \cdot (s_W^2 \cdot \mathcal{K}_\ell(s,t) + \delta_V)(2 \cdot T_3^f) + (4 \cdot q_\ell \cdot s_W^2)(4 \cdot q_f \cdot s_W^2) \mathcal{K}_{\ell f}(s,t) + 2 \cdot (4 \cdot q_\ell)(4 \cdot q_f) \cdot s_W^2 \cdot \mathcal{K}_{\ell f}(s,t) \delta_V] \frac{1}{\Delta^2}.$$

$$\delta_V = \pm 0.00050$$

# The $A_4$ in the full phase-space

## Powheg+MiNLO events with EW NLO+HO corrections (TauSpinner + Dizet 6.45)

Table 30: The  $A_4$  calculated including full EW corrections, in experimental bins  $M_{bin}, Y_{bin}$ . Updated with Dizet 6.45 form factors.

$M_{bin}, Y_{bin}$	$\sin^2 \theta_W^{eff} = 0.23100$	$\sin^2 \theta_W^{eff} = 0.23150$	$\sin^2 \theta_W^{eff} = 0.23200$	$\Delta_P$	$\Delta_M$
1, 1	-0.04712	-0.04716	-0.04720	0.00004	-0.00004
1, 2	-0.13503	-0.13515	-0.13528	0.00012	-0.00012
1, 3	-0.24463	-0.24485	-0.24506	0.00021	-0.00021
1, 4	-0.37830	-0.37862	-0.37893	0.00032	-0.00032
1, 5	-0.52417	-0.52461	-0.52505	0.00044	-0.00045
1, 6	-0.69617	-0.69676	-0.69735	0.00059	-0.00059
1, 7	-0.87865	-0.87940	-0.88013	0.00074	-0.00075
1, 8	-1.04540	-1.04630	-1.04719	0.00089	-0.00090
1, 9	-1.22346	-1.22447	-1.22548	0.00100	-0.00101
2, 1	-0.04599	-0.04610	-0.04621	0.00011	-0.00011
2, 2	-0.15099	-0.15133	-0.15168	0.00035	-0.00035
2, 3	-0.25526	-0.25587	-0.25648	0.00061	-0.00061
2, 4	-0.41167	-0.41258	-0.41349	0.00091	-0.00091
2, 5	-0.56949	-0.57077	-0.57204	0.00127	-0.00128
2, 6	-0.76368	-0.76537	-0.76705	0.00168	-0.00169
2, 7	-0.97291	-0.97501	-0.97710	0.00209	-0.00210
2, 8	-1.17256	-1.17506	-1.17755	0.00249	-0.00250
2, 9	-1.35532	-1.35813	-1.36092	0.00279	-0.00281
3, 1	-0.00223	-0.02245	-0.02267	0.00022	-0.00022
3, 2	-0.07747	-0.07815	-0.07884	0.00068	-0.00068
3, 3	-0.13806	-0.13926	-0.14046	0.00120	-0.00120
3, 4	-0.20723	-0.20904	-0.21085	0.00181	-0.00181
3, 5	-0.28220	-0.28473	-0.28726	0.00253	-0.00253
3, 6	-0.38354	-0.38690	-0.39024	0.00335	-0.00335
3, 7	-0.49968	-0.50390	-0.50810	0.00421	-0.00421
3, 8	-0.61852	-0.62357	-0.62861	0.00504	-0.00505
3, 9	-0.72616	-0.73188	-0.73758	0.00571	-0.00572
4, 1	0.00742	0.00721	0.00691	0.00030	-0.00021
4, 2	0.02350	0.02285	0.02204	0.00081	-0.00065
4, 3	0.04065	0.03950	0.03811	0.00139	-0.00114
4, 4	0.06283	0.06111	0.05904	0.00207	-0.00172
4, 5	0.08847	0.08604	0.08324	0.00280	-0.00243
4, 6	0.11597	0.11267	0.10908	0.00359	-0.00330
4, 7	0.14134	0.13709	0.13271	0.00438	-0.00424
4, 8	0.16457	0.15941	0.15424	0.00517	-0.00516
4, 9	0.18159	0.17575	0.16990	0.00585	-0.00584
5, 1	0.03252	0.03232	0.03213	0.00020	-0.00020
5, 2	0.11502	0.11441	0.11380	0.00061	-0.00061
5, 3	0.19154	0.19047	0.18940	0.00107	-0.00107
5, 4	0.28932	0.28773	0.28612	0.00160	-0.00160
5, 5	0.39524	0.39303	0.39083	0.00220	-0.00220
5, 6	0.52786	0.52499	0.52213	0.00287	-0.00286
5, 7	0.66747	0.66395	0.66042	0.00353	-0.00352
5, 8	0.77857	0.77447	0.77035	0.00411	-0.00411
5, 9	0.86671	0.86218	0.85764	0.00454	-0.00453

Table 31: The  $A_4$  calculated including full EW corrections, in experimental bins  $M_{bin}, Y_{bin}$ . Continuation of Table 30: Updated with Dizet 6.45 form factors.

$M_{bin}, Y_{bin}$	$\sin^2 \theta_W^{eff} = 0.23100$	$\sin^2 \theta_W^{eff} = 0.23150$	$\sin^2 \theta_W^{eff} = 0.23200$	$\Delta_P$	$\Delta_M$
6, 1	0.08041	0.08027	0.08013	0.00014	-0.00014
6, 2	0.18749	0.18705	0.18660	0.00044	-0.00044
6, 3	0.32530	0.32454	0.32378	0.00076	-0.00076
6, 4	0.48828	0.48715	0.48602	0.00113	-0.00113
6, 5	0.65712	0.65557	0.65402	0.00155	-0.00155
6, 6	0.85586	0.85389	0.85191	0.00198	-0.00197
6, 7	1.08547	1.08313	1.08078	0.00235	-0.00234
6, 8	1.24701	1.24433	1.24165	0.00268	-0.00267
6, 9	1.42075	1.41796	1.41516	0.00280	-0.00279
7, 1	0.07507	0.07498	0.07489	0.00009	-0.00009
7, 2	0.25610	0.25582	0.25553	0.00028	-0.00028
7, 3	0.44532	0.44484	0.44435	0.00049	-0.00049
7, 4	0.67565	0.67494	0.67423	0.00071	-0.00071
7, 5	0.89229	0.89135	0.89041	0.00095	-0.00094
7, 6	1.11025	1.10907	1.10788	0.00119	-0.00119
7, 7	1.35502	1.35364	1.35226	0.00139	-0.00138
7, 8	1.51281	1.51129	1.50975	0.00153	-0.00153
7, 9	1.69631	1.69481	1.69330	0.00151	-0.00150

$\Delta_P, \Delta_M$  correspond to  $\Delta A_4$  due to  $\delta_V = \pm 0.00050$

$m_{II} = 86-96$  GeV

# The $A_4$ in the full phase-space

## Powheg+MiNLO events with Effective EW Born corrections (TauSpinner)

Table 32: The  $A_4$  calculated with Effective Born v0, in experimental bins  $M_{bin}, Y_{bin}$ .

$M_{bin}, Y_{bin}$	$\sin^2 \theta_W^{eff} = 0.23100$	$\sin^2 \theta_W^{eff} = 0.23150$	$\sin^2 \theta_W^{eff} = 0.23200$	$\Delta_P$	$\Delta_M$
1, 1	-0.04709	-0.04713	-0.04716	0.00004	-0.00004
1, 2	-0.13492	-0.13503	-0.13514	0.00011	-0.00011
1, 3	-0.24443	-0.24462	-0.24482	0.00019	-0.00020
1, 4	-0.37801	-0.37831	-0.37860	0.00029	-0.00029
1, 5	-0.52369	-0.52409	-0.52450	0.00040	-0.00041
1, 6	-0.69549	-0.69604	-0.69657	0.00054	-0.00054
1, 7	-0.87772	-0.87840	-0.87908	0.00068	-0.00068
1, 8	-1.04412	-1.04494	-1.04576	0.00082	-0.00082
1, 9	-1.22195	-1.22288	-1.22379	0.00091	-0.00093
2, 1	-0.04604	-0.04615	-0.04625	0.00010	-0.00010
2, 2	-0.15122	-0.15155	-0.15188	0.00033	-0.00033
2, 3	-0.25571	-0.25629	-0.25686	0.00057	-0.00058
2, 4	-0.41223	-0.41309	-0.41395	0.00086	-0.00086
2, 5	-0.57033	-0.57154	-0.57274	0.00120	-0.00121
2, 6	-0.76484	-0.76644	-0.76802	0.00158	-0.00160
2, 7	-0.97431	-0.97630	-0.97826	0.00197	-0.00198
2, 8	-1.17431	-1.17667	-1.17901	0.00234	-0.00236
2, 9	-1.35734	-1.35998	-1.36260	0.00262	-0.00264
3, 1	-0.02243	-0.02264	-0.02285	0.00021	-0.00021
3, 2	-0.07801	-0.07866	-0.07932	0.00065	-0.00065
3, 3	-0.13899	-0.14015	-0.14129	0.00115	-0.00115
3, 4	-0.20869	-0.21043	-0.21216	0.00173	-0.00174
3, 5	-0.28436	-0.28678	-0.28920	0.00242	-0.00242
3, 6	-0.38640	-0.38961	-0.39281	0.00320	-0.00321
3, 7	-0.50327	-0.50731	-0.51133	0.00402	-0.00404
3, 8	-0.62293	-0.62776	-0.63257	0.00481	-0.00483
3, 9	-0.73141	-0.73688	-0.74233	0.00545	-0.00547
4, 1	0.00725	0.00703	0.00681	0.00022	-0.00022
4, 2	0.02298	0.02229	0.02155	0.00074	-0.00069
4, 3	0.03972	0.03851	0.03721	0.00130	-0.00121
4, 4	0.06141	0.05962	0.05769	0.00193	-0.00179
4, 5	0.08642	0.08398	0.08135	0.00263	-0.00245
4, 6	0.11317	0.10997	0.10656	0.00341	-0.00320
4, 7	0.13780	0.13372	0.12951	0.00422	-0.00407
4, 8	0.16030	0.15539	0.15042	0.00497	-0.00492
4, 9	0.17675	0.17115	0.16560	0.00555	-0.00559
5, 1	0.03233	0.03214	0.03195	0.00019	-0.00019
5, 2	0.11426	0.11368	0.11310	0.00058	-0.00058
5, 3	0.19031	0.18930	0.18829	0.00101	-0.00101
5, 4	0.28747	0.28596	0.28444	0.00151	-0.00151
5, 5	0.39271	0.39063	0.38854	0.00208	-0.00208
5, 6	0.52438	0.52167	0.51897	0.00271	-0.00271
5, 7	0.66295	0.65962	0.65630	0.00332	-0.00333
5, 8	0.77326	0.76938	0.76551	0.00387	-0.00388
5, 9	0.86079	0.85652	0.85226	0.00426	-0.00427

Table 33: The  $A_4$  calculated Effective Born v0, in experimental bins  $M_{bin}, Y_{bin}$ . Continuation of Table 32:

$M_{bin}, Y_{bin}$	$\sin^2 \theta_W^{eff} = 0.23100$	$\sin^2 \theta_W^{eff} = 0.23150$	$\sin^2 \theta_W^{eff} = 0.23200$	$\Delta_P$	$\Delta_M$
6, 1	0.07990	0.07977	0.07963	0.00013	-0.00013
6, 2	0.18659	0.18618	0.18577	0.00041	-0.00041
6, 3	0.32375	0.32304	0.32233	0.00071	-0.00071
6, 4	0.48585	0.48479	0.48374	0.00105	-0.00105
6, 5	0.65382	0.65238	0.65094	0.00144	-0.00144
6, 6	0.85126	0.84943	0.84761	0.00183	-0.00183
6, 7	1.07892	1.07675	1.07458	0.00217	-0.00217
6, 8	1.23912	1.23665	1.23418	0.00247	-0.00247
6, 9	1.41110	1.40854	1.40599	0.00256	-0.00256
7, 1	0.07496	0.07488	0.07479	0.00008	-0.00008
7, 2	0.25547	0.25521	0.25496	0.00025	-0.00025
7, 3	0.44422	0.44378	0.44334	0.00044	-0.00044
7, 4	0.67355	0.67291	0.67228	0.00064	-0.00064
7, 5	0.88920	0.88836	0.88751	0.00085	-0.00085
7, 6	1.10575	1.10469	1.10362	0.00106	-0.00106
7, 7	1.34808	1.34684	1.34561	0.00123	-0.00123
7, 8	1.50391	1.50256	1.50120	0.00136	-0.00136
7, 9	1.68489	1.68358	1.68226	0.00131	-0.00132

$\Delta_P, \Delta_M$  correspond to  $\Delta A_4$  due to  $\delta_V = \pm 0.00050$

# The $A_4$ in the full phase-space

## EW NLO+HO - Effective Born (TauSpinner + Dizet 6.45)

Table 34: The  $\Delta A_4$  for  $\Delta \sin^2 \theta_W = +0.00050$ , estimated with full EW corrections and Effective Born v0, in experimental bins  $M_{bin}, Y_{bin}$ . Updated with Dizet 6.45 form factors.

$M_{bin}, Y_{bin}$	$\Delta A_4$ (full EW)	$\Delta A_4$ (Effective v0)
1, 1	0.00004	0.00004
1, 2	0.00012	0.00011
1, 3	0.00021	0.00019
1, 4	0.00032	0.00029
1, 5	0.00044	0.00040
1, 6	0.00059	0.00054
1, 7	0.00074	0.00068
1, 8	0.00089	0.00082
1, 9	0.00100	0.00091
2, 1	0.00011	0.00010
2, 2	0.00035	0.00033
2, 3	0.00061	0.00057
2, 4	0.00091	0.00086
2, 5	0.00127	0.00120
2, 6	0.00168	0.00158
2, 7	0.00209	0.00197
2, 8	0.00249	0.00234
2, 9	0.00279	0.00262
3, 1	0.00022	0.00021
3, 2	0.00068	0.00065
3, 3	0.00120	0.00115
3, 4	0.00181	0.00173
3, 5	0.00253	0.00242
3, 6	0.00335	0.00320
3, 7	0.00421	0.00402
3, 8	0.00504	0.00481
3, 9	0.00571	0.00545
4, 1	0.00030	0.00022
4, 2	0.00081	0.00074
4, 3	0.00139	0.00130
4, 4	0.00207	0.00193
4, 5	0.00280	0.00263
4, 6	0.00359	0.00341
4, 7	0.00438	0.00422
4, 8	0.00517	0.00497
4, 9	0.00585	0.00555
5, 1	0.00020	0.00019
5, 2	0.00061	0.00058
5, 3	0.00107	0.00101
5, 4	0.00160	0.00151
5, 5	0.00220	0.00208
5, 6	0.00287	0.00271
5, 7	0.00353	0.00332
5, 8	0.00411	0.00387
5, 9	0.00454	0.00426

Table 35: The  $\Delta A_4$  for for  $\Delta \sin^2 \theta_W = +0.00050$ , estimated with full EW corrections and Effective Born v0, in experimental bins  $M_{bin}, Y_{bin}$ . Updated with Dizet 6.45 form factors.

$M_{bin}, Y_{bin}$	$\Delta A_4$ (full EW)	$\Delta A_4$ (Effective v0)
6, 1	0.00014	0.00013
6, 2	0.00044	0.00041
6, 3	0.00076	0.00071
6, 4	0.00113	0.00105
6, 5	0.00155	0.00144
6, 6	0.00198	0.00183
6, 7	0.00235	0.00217
6, 8	0.00268	0.00247
6, 9	0.00280	0.00256
7, 1	0.00009	0.00008
7, 2	0.00028	0.00025
7, 3	0.00049	0.00044
7, 4	0.00071	0.00064
7, 5	0.00095	0.00085
7, 6	0.00119	0.00106
7, 7	0.00139	0.00123
7, 8	0.00153	0.00136
7, 9	0.00151	0.00131

$\Delta A_4$  due to  $\delta_V = +0.00050$



# Factorizing QCD: Reweighting technique

Reweighting possible because of Drell-Yan factorisation properties,  
Mirkes et al. arXiv:9406381.

Method follows technique developed for **TauSpinner** program (for LHC!),  
arXiv:1201.0117; 1802.05459

Define per event electroweak weight  $wt^{EW} = \sigma_{Born}^{new} / \sigma_{Born}^{old}$

$$wt^{EW} = \frac{d\sigma_{Born+EW}(x_1, x_2, \hat{s}, \cos\theta^*, s_W^2)}{d\sigma_{Born}(x_1, x_2, \hat{s}, \cos\theta^*, s_W^2)}$$

$$d\sigma_{Born}(x_1, x_2, \hat{s}, \cos\theta^*, s_W^2) = \sum_{q_f, \bar{q}_f} [f^{q_f}(x_1, \dots) f^{\bar{q}_f}(x_2, \dots) d\sigma_{Born}^{q_f \bar{q}_f}(\hat{s}, \cos\theta^*, s_W^2) \\ + f^{q_f}(x_2, \dots) f^{\bar{q}_f}(x_1, \dots) d\sigma_{Born}^{\bar{q}_f q_f}(\hat{s}, -\cos\theta^*, s_W^2)]$$

$x_1, x_2, \cos\theta^*$  (symmetrised)  
calculated using 4-momenta  
of outgoing leptons;  
asymmetry in sign of  $\cos\theta^*$   
from average over PDFs

Allows to reweight MC event generated with any EW LO scheme to

- **Improved Born Approximation** including:
  - EW loop corrections to propagators
  - EW loop corrections to vector couplings without boxes
  - EW Loop corrections with boxes
- **Effective Born** with LEP with improved norm. parametrisation

# Different rest frames (reminder)

- Polar and azimuthal angle are defined in the rest-frame of outgoing leptons. But how?
  - 1) **Collins-Soper frame**: used for Ais 8 TeV measurement. **All Ais non-zero at high pT.**
  - 2) **Mustraal frame**: proven in 80's (F.A.Berends et al. Comp. Phys. Com 29(1983) 185) that for  $q\bar{q} \rightarrow Z \rightarrow l\bar{l}$  and single spin-1 emission in initial state (gluon or photon) matrix element can be presented as weighted sum of Borns. Adapted (ERW&ZW, Eur. Phys. J C76 (2016) 473) to pp case and added definition of azimuthal angle. **Only A4 non-zero at high pT.**
  - 3)  **$\cos\theta^*$  frame**: used in precision calculations for LEP to minimise some higher order corrections to optimal observables (eg. asymmetries).
- Any of those frames can be used for Ais definition or calculating  $wt_{sw2}$ ,  $wt_{EWloop}$ . For now, for measurement we use (1) and for calculating weights we use (3).

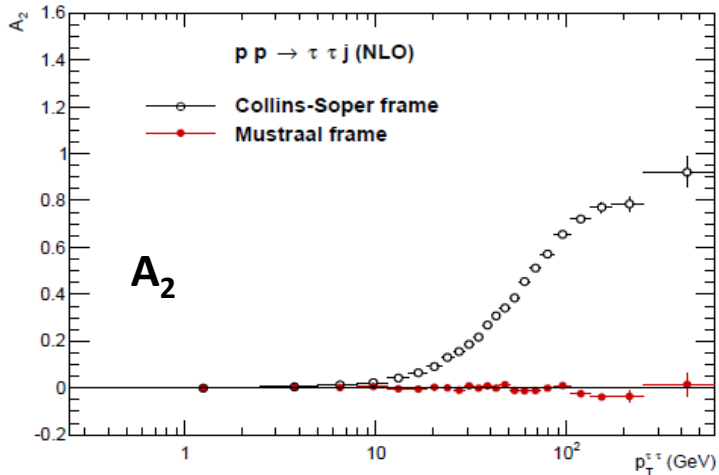
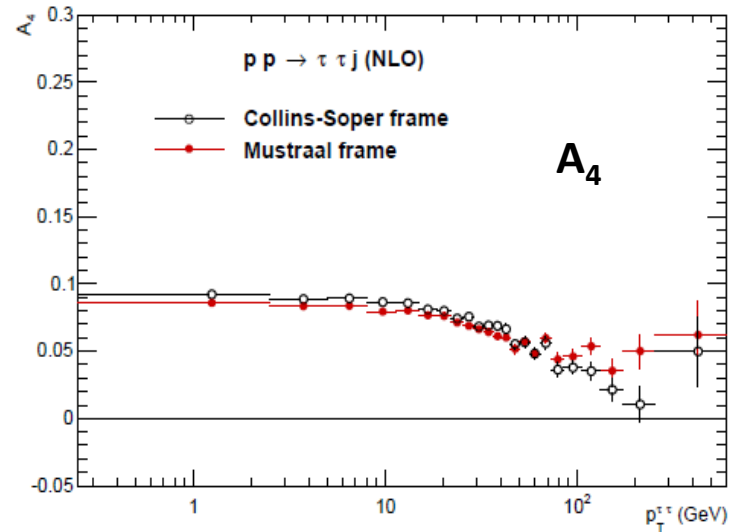
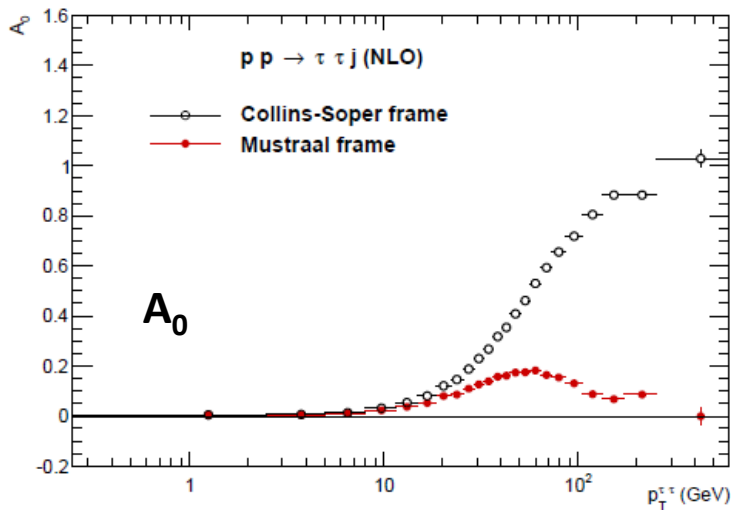
# $A_i$ 's in different rest frames

Black: Collins-Soper frame

Red: Mustraal frame

ERW & Z. Was

arXiv:1605.05450



MC events generated with Powheg+MiNLO Zj (NLO)

In **Mustraal frame**, NLO QCD event has „Born-like” angular correlations of leptons from Z decay. Preferred frame to **factorise EW and QCD corrections**. Remaining  $A_i$ 's close to zero.

# How sensitive is $A_{fb}$ to rest frame used for calculating $w_t^{EW}$

Table 24: The difference in forward-backward asymmetry,  $\Delta A_{FB}$  around Z-pole,  $m_{ee} = 89 - 93$  GeV. The difference is calculated using  $\cos\theta^{CS}$  to define forward and backward hemisphere. The EW weight is calculated with  $\cos\theta^*$ ,  $\cos\theta^{Mstraal}$  or  $\cos\theta^{CS}$ .

Updated with Dizet 6.45 form factors.

Corrections to $A_{FB}$ ( $89 < m_{ee} < 93$ GeV)	$w_t^{EW}(\cos\theta^*)$	$w_t^{EW}(\cos\theta^{ML})$	$w_t^{EW}(\cos\theta^{CS})$
$A_{FB}(\text{EW/QCD corr. to } m_W) - A_{FB}(\text{EW LO } \alpha(0))$	-0.02076	-0.02091	-0.02080
$A_{FB}(\text{EW/QCD corr. to } \chi(Z), \chi(\gamma)) - A_{FB}(\text{EW LO } \alpha(0))$	-0.02047	-0.02062	-0.02051
$A_{FB}(\text{EW/QCD FF no boxes}) - A_{FB}(\text{EW LO } \alpha(0))$	-0.03491	-0.03517	-0.03497
$A_{FB}(\text{EW/QCD FF with boxes}) - A_{FB}(\text{EW LO } \alpha(0))$	-0.03489	-0.03516	-0.03496
$A_{FB}(\text{Eff. } v_0) - A_{FB}(\text{EW/QCD FF with boxes})$	0.00039	0.00037	0.00039
$A_{FB}(\text{Eff. } v_1) - A_{FB}(\text{EW/QCD FF with boxes})$	0.00042	0.00038	0.00042
$A_{FB}(\text{Eff. } v_2) - A_{FB}(\text{EW/QCD FF with boxes})$	0.00022	0.00024	0.00022

$$\delta A_{fb}^{\text{frame}} = 20 \cdot 10^{-5}$$

Shift in predicted  $A_{fb}(M_Z)$  by  $20 \cdot 10^{-5} \rightarrow \Delta \sin^2\theta_w^{\text{eff}} = \sim 10 \cdot 10^{-5}$

We can take it as proxy for systematics uncertainties from QCD corrections. Should fall into same category as PDFs uncertainties.

# Summary

- **Started preparing Tables with binned  $\Delta A_4$  (EW)**
  - Quantify difference between EW NLO+HO corrections and Effective Born
  - Next step will be to compare with those presented by Aleko.  
Not straightforward, because I don't have 13 TeV samples generated consistently with setup of CMS pseudodata.
- **To which extend we can factorise EW and QCD corrections**
- **Draft v05 status**
  - **No comments nor contributions received!**

# SPARES slides

# Z-boson propagator

For now we took pragmatic approach: use defaults of each code:

- Powheg\_ew and MCSANC: pole-mass and fixed width propagator
  - Not clear to me if implementation includes  $N_Z$  modification to couplings (?)
  - If yes, how it is share between  $g_a$ ,  $g_f$  couplings, does it affect  $\sin^2\theta_{\text{eff}}$  interpretation (?)
  - If not included, does it count for „missing HO corrections“ (?)
- $\text{wt}^{\text{EW}}$  : calculated with on-shell masses and running width propagator, as it is standard used by Zfitter+Dizet

We should keep it in mind, that ones we reach precision of the comparisons which might be sensitive to the effect of  $\chi(s)$  implementation.

It should be discussed as component of theoretical uncertainties of the predictions.

# Z-boson propagator

Running width, on-shell  $M_Z, \Gamma_Z$

$$\chi_Z(s) = \frac{1}{s - M_Z^2 + i \cdot \Gamma_Z \cdot s / M_Z}$$

Fixed width, on-shell  $M_Z, \Gamma_Z$

$$\chi'_Z(s) = \frac{1}{s - M_Z^2 + i \cdot \Gamma_Z \cdot M_Z}$$

Equivalent to  $\chi'_Z(s)$

shifted  $M_Z, \Gamma_Z$ , scaled by  $N_Z$

$$\chi_Z(s) = N'_Z \frac{1}{s - M'^2_Z + i \Gamma'_Z M'_Z}$$

$$M'_Z = \frac{M_Z}{\sqrt{1 + \Gamma_Z^2 / M_Z^2}}$$

$$\Gamma'_Z = \frac{\Gamma_Z}{\sqrt{1 + \Gamma_Z^2 / M_Z^2}}$$

$$N'_Z = \frac{(1 - i \cdot \Gamma_Z / M_Z)}{(1 + \Gamma_Z^2 / M_Z^2)} = \frac{(1 - i \cdot \Gamma'_Z / M'_Z)}{(1 + \Gamma'^2_Z / M'^2_Z)}$$

Shifted  $M_Z, \Gamma_Z$ , no scaling  $N_Z$

$$\chi_Z(s) = \frac{1}{s - M'^2_Z + i \Gamma'_Z M'_Z}$$

$$M'_Z = \frac{M_Z}{\sqrt{1 + \Gamma_Z^2 / M_Z^2}}$$

$$\Gamma'_Z = \frac{\Gamma_Z}{\sqrt{1 + \Gamma_Z^2 / M_Z^2}}$$



# Z-boson propagator

Reference: LEP convention: running width propagator, nominal  $M_Z, \Gamma_Z$

Table 15: Ratio of the cross-section  $\sigma$  calculated with different form Z-boson propagator, integrated over specified mass windows. Shown in case of EW LO and EW NLO+HO predictions with  $O(\alpha^0)$  v0 EW scheme.

$\sigma$ (Fixed/Running)	$90.5 < m_{ee} < 91.5$ GeV	$89 < m_{ee} < 93$ GeV	$60 < m_{ee} < 81$ GeV	$81 < m_{ee} < 101$ GeV	$101 < m_{ee} < 150$ GeV
EW LO					
with $M_Z, \Gamma_Z$ shift, no scaling	1.00087	1.00087	1.00062	1.00086	1.00071
no $M_Z, \Gamma_Z$ shift, no scaling	0.99620	1.00074	0.99716	0.99977	1.00392
EW NLO+HO					
with $M_Z, \Gamma_Z$ shift, no scaling	1.00113	1.00085	1.00043	1.00083	1.00075
no $M_Z, \Gamma_Z$ shift, no scaling	0.99746	1.00122	0.99719	1.00013	1.00392

Table 16: Difference in  $A_{fb}$  calculated with different form of Z-boson propagator, integrated over specified mass windows. Shown in case of EW LO and EW NLO+HO predictions with  $O(\alpha^0)$  v0 EW scheme.

$\Delta A_{fb}$ (Running - Fixed)	$90.5 < m_{ee} < 91.5$ GeV	$89 < m_{ee} < 93$ GeV	$60 < m_{ee} < 81$ GeV	$81 < m_{ee} < 101$ GeV	$101 < m_{ee} < 150$ GeV
EW LO					
with $M_Z, \Gamma_Z$ shift, no scaling	-0.00048	-0.00047	-0.00047	-0.00047	-0.00030
no $M_Z, \Gamma_Z$ shift, no scaling	-0.00006	-0.00026	-0.00012	-0.00040	-0.00005
EW NLO+HO					
with $M_Z, \Gamma_Z$ shift, no scaling	-0.00053	-0.00053	-0.00052	-0.00053	-0.00024
no $M_Z, \Gamma_Z$ shift, no scaling	-0.00007	-0.00030	-0.00026	-0.00048	-0.00004



used by LHC experiments



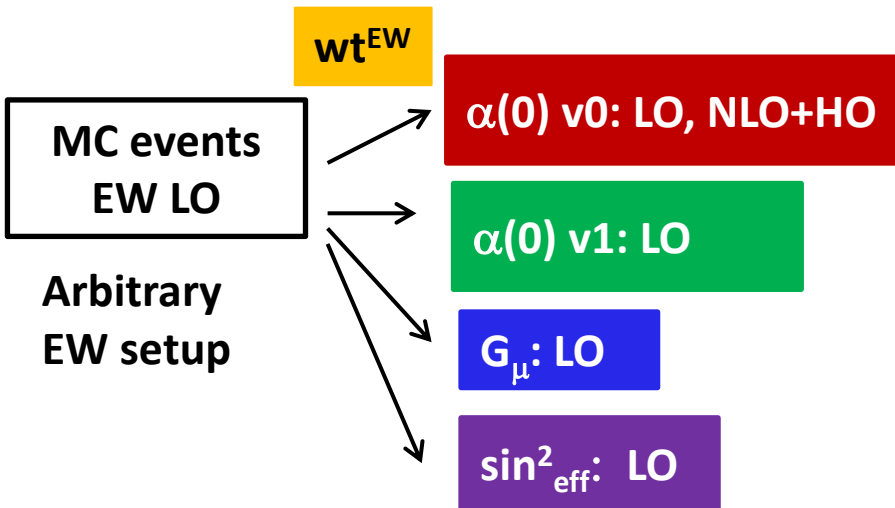
for  $A_{fb}$  and around Z-pole  
very close to LEP convention



used by Powheg\_ew

# Virtual EW corrections

$wt^{EW}$  : TauSpinner + Dizet 6.45



Powheg\_ew: QCD LO

$\alpha(0) v0$ : LO

$\alpha(0) v1$ : LO, NLO, NLO+HO

$G_\mu$ : LO, NLO, NLO+HO

$\sin^2_{eff}$ : LO, NLO, NLO+HO

MCSANC: QCD LO

$\alpha(0) v1$ : LO, NLO, NLO+HO

$G_\mu$ : LO, NLO, NLO+HO

# EW schemes: benchmark input parameters

SM relation used to calculate **EW LO parameters** for different schemes. On-shell mass.

	LEP-legacy		LHC-paradigm		New scheme
Parameter	$(\alpha(0), G_\mu, M_Z)$ $\alpha(0) \text{ v0}$	$(\alpha(0), M_W, M_Z)$ $\alpha(0) \text{ v1}$	$(G_\mu, M_Z, M_W)$ $G_\mu$	$(\alpha(0), s_W^2, M_Z)$ $\sin_{eff}^2 \text{ v1}$	$(G_\mu, s_W^2, M_Z)$ $\sin_{eff}^2 \text{ v2}$
$M_Z$ (GeV)	91.1876	91.1876	91.1876	91.1876	91.1876
$\Gamma_Z$ (GeV)	2.4952	2.4952	2.4952	2.4952	2.4952
$\Gamma_W$ (GeV)	2.085	2.085	2.085	2.085	2.085
$1/\alpha$	137.035999139	137.035999139	132.23323	137.035999139	128.744939484
$\alpha$	0.007297353	0.007297353	0.007562396	0.007297353	0.007767296
$G_\mu$ (GeV <sup>-2</sup> )	$1.1663787 \cdot 10^{-5}$	$1.1254734 \cdot 10^{-5}$	$1.1663787 \cdot 10^{-5}$	$1.09580954 \cdot 10^{-5}$	$1.1663787 \cdot 10^{-5}$
$M_W$ (GeV)	80.93886	80.385	80.385	79.93886984	79.93886984
$s_W^2$	0.2121517	0.2228972	0.2228972	0.231499	0.231499
$\frac{G_\mu M_Z^2 \cdot 16c_W^2 s_W^2}{\sqrt{2} \cdot 8\pi \cdot \alpha} = 1.0$	$\rightarrow s_W^2, M_W$	$\rightarrow G_\mu, s_W^2$	$\rightarrow \alpha, s_W^2$	$\rightarrow G_\mu, m_W$	$\rightarrow \alpha, m_W$
$s_W^2 = 1 - m_W^2/m_Z^2$					
$\alpha_s(M_Z)$	0.120178900000	0.120178900000	0.120178900000	0.120178900000	0.120178900000

$$s_W^2 = 1 - m_W^2/m_Z^2$$

$$G_\mu = \frac{\pi\alpha}{\sqrt{2}M_W^2 s_W^2}$$

# Pseudo-observables at Z-pole: benchmarks

„Best predictions” in each EW scheme, i.e. EW NLO+HO

Parameter	Dizet v6.45	Powheg_ew, MCSANC		Powgeh_ew	
	$(\alpha(0), G_\mu, M_Z)$ $\alpha(0)$ v0	$(\alpha(0), M_W, M_Z)$ $\alpha(0)$ v1	$(G_\mu, M_Z, M_W)$ $G_\mu$	$(\alpha(0), s_W^2, M_Z)$ $\sin_{eff}^2$ v1	$(G_\mu, s_W^2, M_Z)$ $\sin_{eff}^2$ v2
$M_Z$ (GeV)	91.1876	91.1876	91.1876	91.1876	91.1876
$1/\alpha(M_Z)$	0.0077549256		???		???
$\alpha(M_Z)$	128.9503020				
$G_\mu$ (GeV <sup>-2</sup> )	$1.1663787 \cdot 10^{-5}$		$1.1663787 \cdot 10^{-5}$		$1.1663787 \cdot 10^{-5}$
$M_W$ (GeV)	80.358935	80.385	80.385		???
$s_W^2$	0.223401084	0.22289722	0.22289722		???
$\sin^2 \theta_{eff}^l$	0.231499			0.231499	0.231499
$\sin^2 \theta_{eff}^u$	0.231392		???		???
$\sin^2 \theta_{eff}^d$	0.231265				
$\sin^2 \theta_{eff}^b$	0.232733				

Experiments measure observables: cross-sections, asymmetries, distributions.

We need predictions to interpret these measurements.

For now, only TauSpinner + Dizet provides predictions for LEP-style pseudo-observables.

# $\sin^2 \theta_W^{\text{eff}}$ at NLO in $G_\mu$ scheme

## H Results from analytical programs

From S. Dittmaier

Table 34: The  $\sin^2 \theta_W^{\text{eff}}$  predictions in EW  $G_\mu$  scheme.

$\sin^2 \theta_W^{\text{eff}}$	EW LO	EW NLO	EW NLO+HO	Comments
lepton	0.2228972225239183	0.2323557983674498		
neutrino	0.2228972225239183	0.2320009933224815		
up-quark	0.2228972225239183	0.2322559935838819		
down-quark	0.2228972225239183	0.2321377252355592		
bottom-quark	0.2228972225239183	0.2337274233845253		

Dizet v6.45

Parameter	$(\alpha(0), G_\mu, M_Z)$ $\alpha(0) \text{ v0}$
$\sin^2 \theta_{eff}^l$	0.231499
$\sin^2 \theta_{eff}^u$	0.231392
$\sin^2 \theta_{eff}^d$	0.231265
$\sin^2 \theta_{eff}^b$	0.232733

# Effective Born

## Can we parametrise Effective Born to bring it closer to EW LO+HO Improved Born Approximation?

Table 4: The EW parameters used for: (i) the EW LO  $\alpha(0)$  v0 scheme, (ii) effective Born spin amplitude around the Z-pole. The  $G_\mu = 1.1663887 \cdot 10^{-5} \text{ GeV}^{-2}$ ,  $M_Z = 91.1876 \text{ GeV}$  and  $\mathcal{K}_f, \mathcal{K}_e, \mathcal{K}_{\ell f} = 1$ .

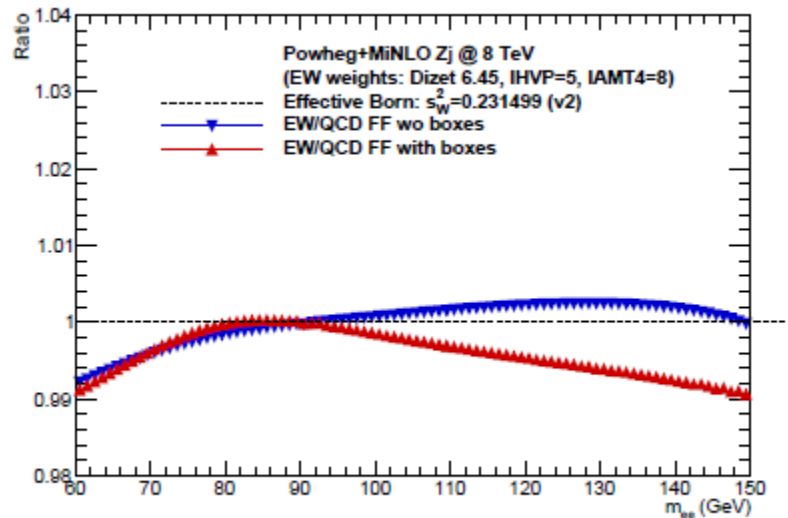
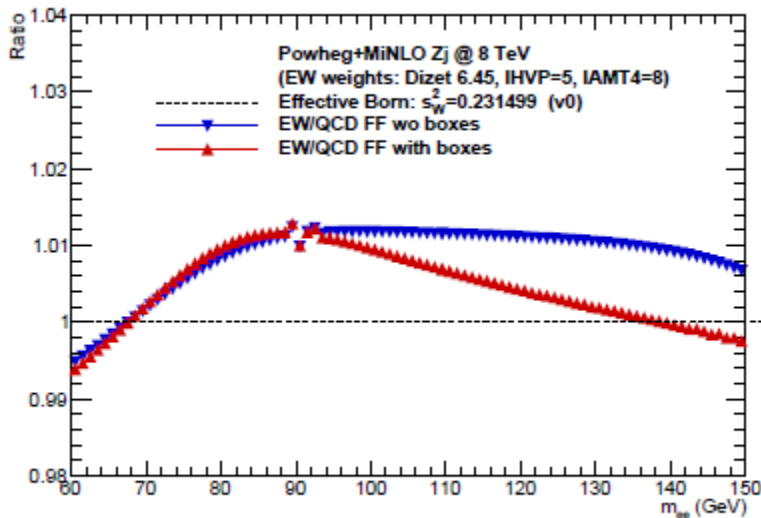
EW LO $\alpha(0)$ scheme	Effective Born v0	Effective Born v1	Effective Born v2
$\alpha = 1/137.03599$	$\alpha = 1/128.9503022$	$\alpha = 1/128.9503022$	$\alpha = 1/128.9503022$
$s_W^2 = 0.21215$	$s_W^2 = 0.231499$	$s_W^2 = 0.231499$	$s_W^{2\ell} = 0.231499$ $s_W^{2up} = 0.231392$ $s_W^{2down} = 0.231265$
$\rho_{\ell f} = 1.0$	$\rho_{\ell f} = 1.0$	$\rho_{\ell f} = 1.005$	$\rho_{\ell up} = 1.005403$ $\rho_{\ell down} = 1.005889$

Dizet v6.45

Parameter	$(\alpha(0), G_\mu, M_Z)$ $\alpha(0)$ v0
$M_Z$ (GeV)	91.1876
$1/\alpha(M_Z)$	0.0077549256
$\alpha(M_Z)$	128.9503020
$G_\mu$ ( $\text{GeV}^{-2}$ )	$1.1663787 \cdot 10^{-5}$
$M_W$ (GeV)	80.358935
$s_W^2$	0.223401084
$\sin^2 \theta_{eff}^\ell$	0.231499
$\sin^2 \theta_{eff}^u$	0.231392
$\sin^2 \theta_{eff}^d$	0.231265
$\sin^2 \theta_{eff}^b$	0.232733

# Effective Born

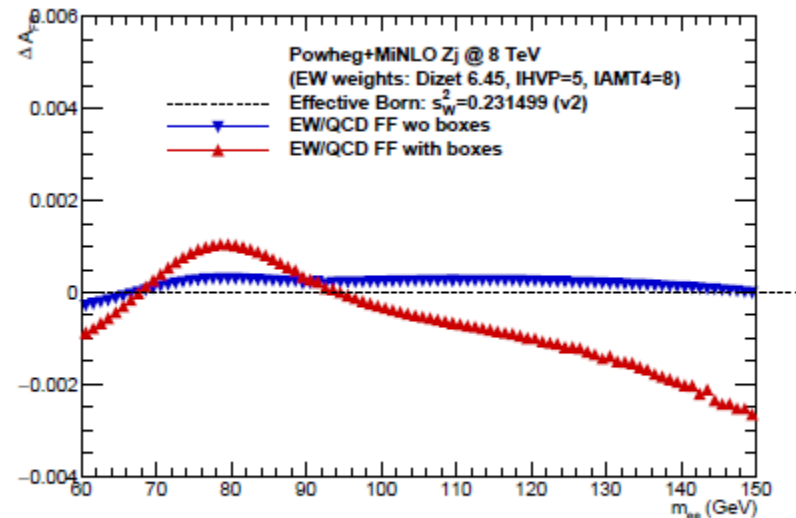
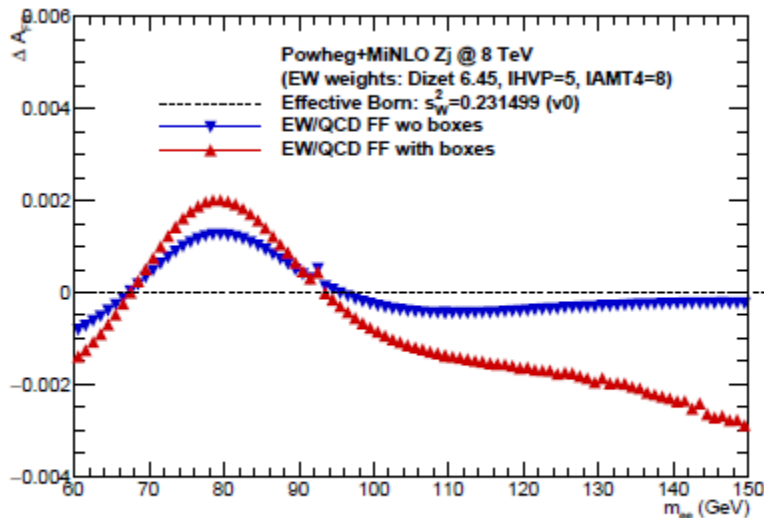
Can we parametrise Effective Born to bring it closer to EW LO+HO Improved Born Approximation?



$\sigma^{e\bar{e}}(NLO+HO)/\sigma(effect.)$	$89 < m_{ee} < 93$ GeV	$60 < m_{ee} < 81$ GeV	$81 < m_{ee} < 101$ GeV	$101 < m_{ee} < 150$ GeV
Effective v0	1.01142	1.00411	1.01135	1.00627
Effective v1	1.00130	0.99780	1.00132	0.99800
Effective v2	0.99989	0.99701	0.99987	0.99654

# Effective Born

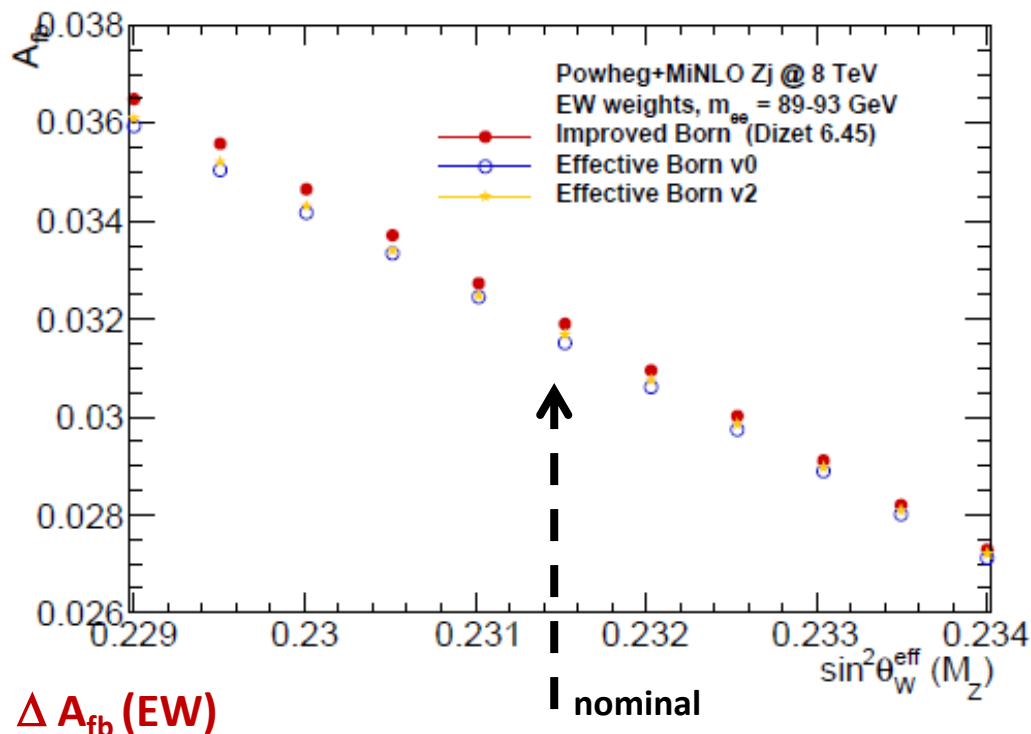
Can we parametrise Effective Born to bring it closer to EW LO+HO Improved Born Approximation?



$A_{fb}^{ref}(NLO+HO) - A_{fb}(effect.)$	$89 < m_{ee} < 93$ GeV	$60 < m_{ee} < 81$ GeV	$81 < m_{ee} < 101$ GeV	$101 < m_{ee} < 150$ GeV
Effective v0	0.00039	0.00104	0.00042	-0.00153
Effective v1	0.00042	0.00068	0.00042	-0.00094
Effective v2	0.00022	0.00052	0.00024	-0.00087



# $\sin^2\theta_{\text{eff}}$ scan for $A_{\text{FB}}$ and $A_4$



Effective Born v0	Effective Born v1	Effective Born v2
$\alpha = 1/128.9503022$	$\alpha = 1/128.9503022$	$\alpha = 1/128.9503022$
$s_W^2 = 0.231499$	$s_W^2 = 0.231499$	$s_W^{\ell} = 0.231499$
		$s_W^{up} = 0.231392$
		$s_W^{down} = 0.231265$
$\rho_{\ell f} = 1.0$	$\rho_{\ell f} = 1.005$	$\rho_{\ell up} = 1.005403$
		$\rho_{\ell down} = 1.005889$

## Dizet v6.45

Parameter	$(\alpha(0), G_\mu, M_Z)$ $\alpha(0)$ v0
$M_Z$ (GeV)	91.1876
$1/\alpha(M_Z)$	0.0077549256
$\alpha(M_Z)$	128.9503020
$G_\mu$ ( $\text{GeV}^{-2}$ )	$1.1663787 \cdot 10^{-5}$
$M_W$ (GeV)	80.358935
$s_W^2$	0.223401084
$\sin^2\theta_{\text{eff}}^{\ell}$	0.231499
$\sin^2\theta_{\text{eff}}^u$	0.231392
$\sin^2\theta_{\text{eff}}^d$	0.231265
$\sin^2\theta_{\text{eff}}^b$	0.232733

$A_{\text{fb}}^{\text{ref}}(NLO+HO) - A_{\text{fb}}(\text{effect.})$	$89 < m_{ee} < 93$ GeV	$60 < m_{ee} < 81$ GeV	$81 < m_{ee} < 101$ GeV	$101 < m_{ee} < 150$ GeV
Effective v0	0.00039	0.00104	0.00042	-0.00153
Effective v1	0.00042	0.00068	0.00042	-0.00094
Effective v2	0.00022	0.00052	0.00024	-0.00087

## EW NLO+HO corrections vs „Effective Born v0”

- Shift in predicted  $A_{\text{fb}}(M_Z)$  by  $40 \cdot 10^{-5} \rightarrow \Delta \sin^2\theta_w^{\text{eff}} = \sim 20 \cdot 10^{-5}$
- Different slope for  $A_{\text{fb}}(\sin^2\theta_w^{\text{eff}}(M_Z))$

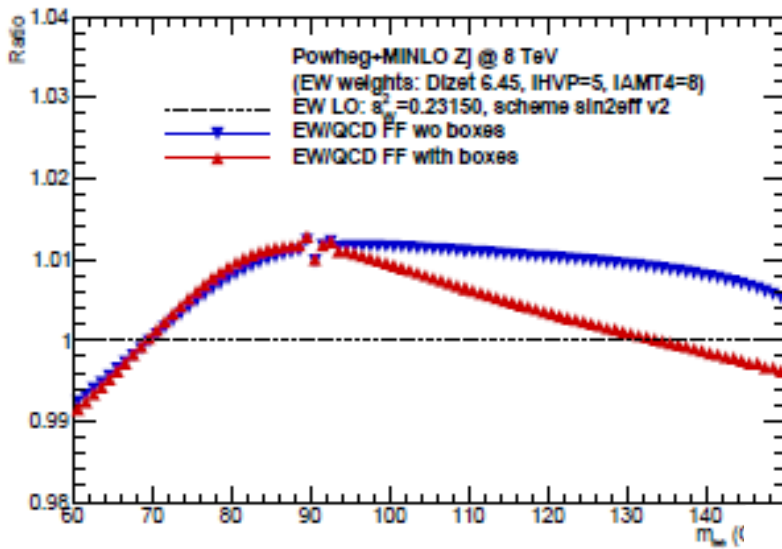
# $\sigma(\text{EW})/\sigma(\text{LO})$ „expected” for different EW schemes

Ref: NLO+HO with  $\alpha(0)$  scheme, Dizet 6.45 form-factors,  $w_t^{\text{EW}}$  from TauSpinner

EW LO

$\sigma^{\text{ref}}(NLO+HO)/\sigma(LO)$	$89 < m_{ee} < 93 \text{ GeV}$	$60 < m_{ee} < 81 \text{ GeV}$	$81 < m_{ee} < 101 \text{ GeV}$	$101 < m_{ee} < 150 \text{ GeV}$
EW scheme $\alpha(0)$ v0	0.96510	1.04695	0.96632	0.96508
EW scheme $\alpha(0)$ v1	1.06558	1.09892	1.06613	1.06202
EW scheme $G_\mu$	0.99211	1.02321	0.99264	0.98884
EW scheme $\sin^2\theta_{\text{eff}} v2$	1.01141	1.00293	1.01132	1.00572

Using Dizet 6.45 form-factors and  $w_t^{\text{EW}}$ , table above presents expected size of EW corrections in different EW schemes.



EW LO  $\sin^2\theta_{\text{eff}} v2$

EW NLO+HO  $\alpha(0)$  v0

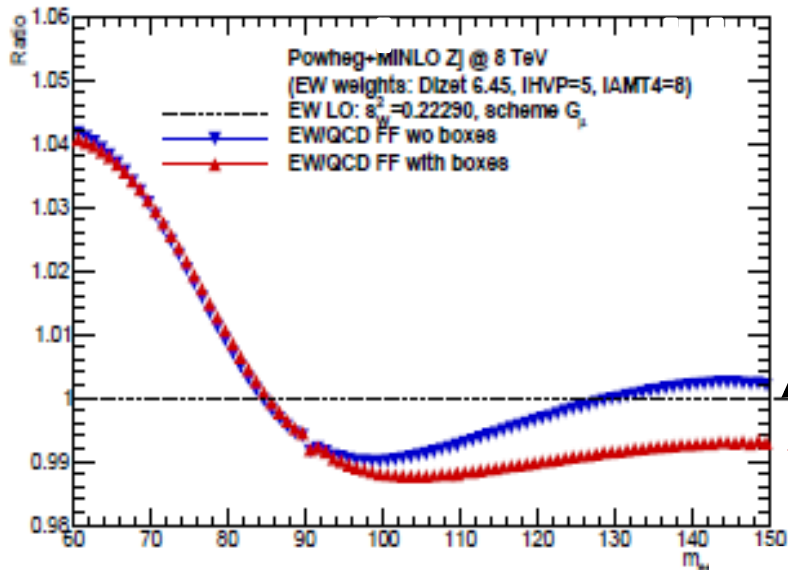
# $\sigma(\text{EW})/\sigma(\text{LO})$ „expected” for different EW schemes

Ref: NLO+HO with  $\alpha(0)$  scheme, Dizet 6.45 form-factors,  $w_t^{\text{EW}}$  from TauSpinner

EW LO

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Using Dizet 6.45 form-factors and  $w_t^{\text{EW}}$ , table above presents expected size of EW corrections in different EW schemes.



EW LO  $G_\mu$

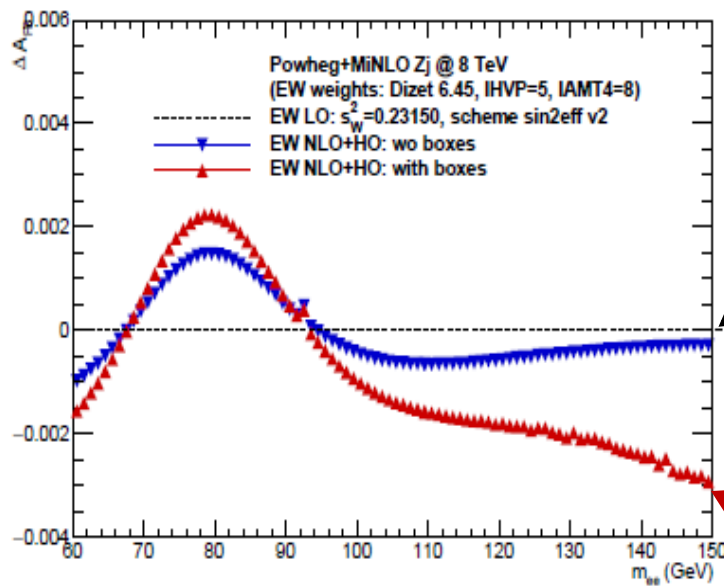
EW NLO+HO  $\alpha(0)$  v0

# $\Delta A_{fb}(\text{EW})$ „expected” for different EW schemes

Ref: NLO+HO with  $\alpha(0)$  scheme, Dizet 6.45 form-factors,  $wt^{\text{EW}}$  from TauSpinner

EW LO

$A_{fb}^{\text{ref}}(\text{NLO}+\text{HO}) - A_{fb}(\text{LO})$	$89 < m_{ee} < 93 \text{ GeV}$	$60 < m_{ee} < 81 \text{ GeV}$	$81 < m_{ee} < 101 \text{ GeV}$	$101 < m_{ee} < 150 \text{ GeV}$
EW scheme $\alpha(0)$ v0	-0.03489	-0.02880	-0.03514	-0.01334
EW scheme $\alpha(0)$ v1	-0.01508	-0.01104	-0.01515	-0.00684
EW scheme $G_\mu$	-0.01507	-0.01104	-0.01514	0.00684
EW scheme $\sin^2 \theta_{\text{eff}} \text{ v2}$	-0.00039	0.00115	-0.00046	-0.00171



Using Dizet 6.45 form-factors and  $wt^{\text{EW}}$ , table above presents expected size of EW corrections in different EW schemes.

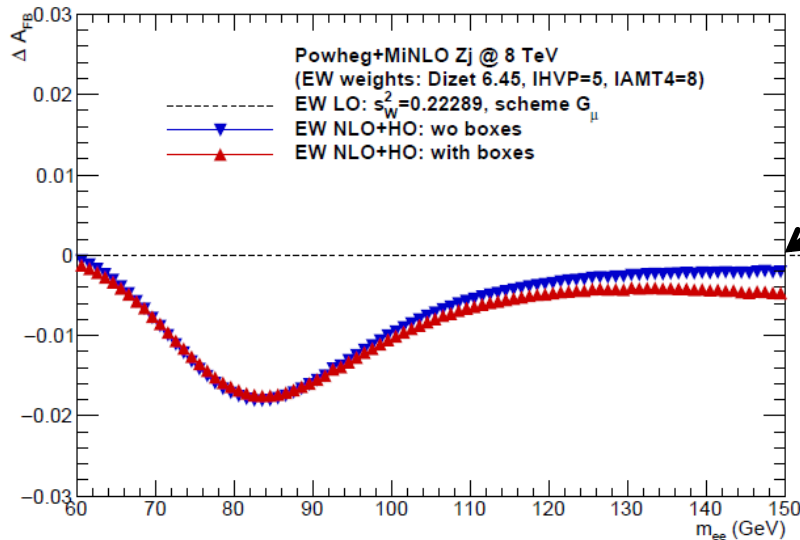
EW LO  $\sin^2 \theta_{\text{eff}} \text{ v2}$

EW NLO+HO  $\alpha(0)$  v0

# $\Delta A_{fb}(EW)$ „expected” for different EW schemes

Ref: NLO+HO with  $\alpha(0)$  scheme, Dizet 6.45 form-factors,  $wt^{EW}$  from TauSpinner

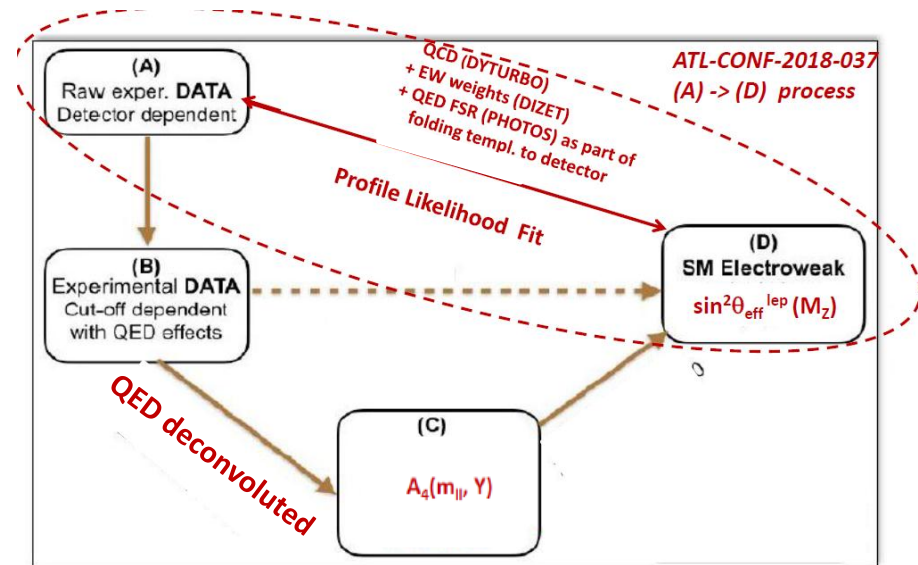
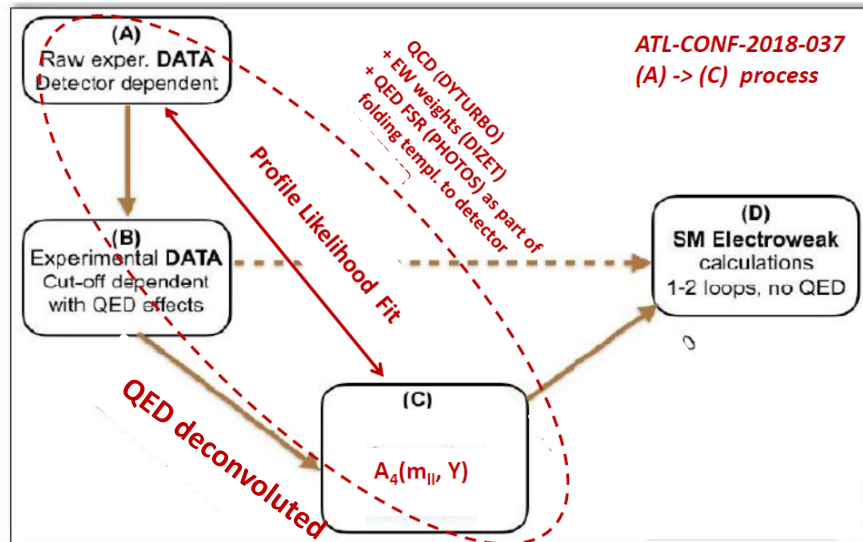
$A_{fb}^{ref}(NLO+HO) - A_{fb}(LO)$	$89 < m_{ee} < 93$ GeV	$60 < m_{ee} < 81$ GeV	$81 < m_{ee} < 101$ GeV	$101 < m_{ee} < 150$ GeV
EW scheme $\alpha(0)$ v0	-0.03489	-0.02880	-0.03514	-0.01334
EW scheme $\alpha(0)$ vL	-0.01508	-0.01104	-0.01515	-0.00684
EW scheme $G_\mu$	-0.01507	-0.01104	-0.01514	0.00684
EW scheme $\sin^2 \theta_{eff}$ v2	-0.00039	0.00115	-0.00046	-0.00171



Using Dizet 6.45 form-factors and  $wt^{EW}$ ,  
table above presents expected size of  
EW corrections in different EW schemes.

EW LO  $G_\mu$   
EW NLO+HO  $\alpha(0)$  v0

# Electroweak Pseudo-Observables at LHC: the meeting point between data and theory



$$N_{\text{exp}}^n(A, \sigma, \theta) = \left\{ \sum_{j=0}^{N_{\text{bins}}} \sigma_j \times L \times \left[ t_{8j}^n(\beta) + \sum_{i=0}^7 A_{ij} \times t_{ij}^n(\beta) \right] \right\} \times \gamma^n + \sum_B^{\text{bkgs}} T_B^n(\beta),$$

$$A_{4,j}(\sin^2 \theta_{\text{eff}}^{\ell}, \theta) = a_j(\theta) \times \sin^2 \theta_{\text{eff}}^{\ell} + b_j(\theta)$$

# Electroweak Pseudo-Observables at LHC: the meeting point between data and theory

ATL-CONF-2018-037

$ y^{\ell\ell} $	$70 < m^{\ell\ell} < 80$ GeV			$80 < m^{\ell\ell} < 100$ GeV				$100 < m^{\ell\ell} < 125$ GeV		
	0 – 0.8	0.8 – 1.6	1.6 – 2.5	0 – 0.8	0.8 – 1.6	1.6 – 2.5	2.5 – 3.6	0 – 0.8	0.8 – 1.6	1.6 – 2.5
Central value (NNLO QCD)	-0.0870	-0.2907	-0.5970	0.0144	0.0471	0.0928	0.1464	0.1045	0.3444	0.6807
$\Delta A_4$ (NNLO - NLO QCD)	0.0003	0.0010	0.0021	-0.0001	-0.0005	-0.0009	-0.0015	-0.0007	-0.0022	-0.0041
$\Delta A_4$ (EW)	0.0008	0.0028	0.0056	0.0002	0.0007	0.0015	0.0026	-0.0008	-0.0026	-0.0048
$\Delta \sin^2 \theta_{\text{eff}}^{\ell}$ (EW)	0.00129	0.00130	0.00133	0.00024	0.00024	0.00025	0.00026	-0.00120	-0.00123	-0.00119
	Uncertainties			Uncertainties				Uncertainties		
Total	0.0035	0.0094	0.0137	0.0007	0.0017	0.0021	0.0021	0.0040	0.0102	0.0140
PDF	0.0034	0.0092	0.0127	0.0007	0.0016	0.0020	0.0019	0.0039	0.0100	0.0131
QCD scales	0.0006	0.0019	0.0052	0.0003	0.0003	0.0004	0.0008	0.0005	0.0022	0.0049

## Observables:

cross-sections and asymmetries ( $A_{\text{FB}}$ ,  $A_4$ ), unfolded to truth level, in different  $m_{\ell\ell}$  and  $y$  bins.

$\Delta \sin^2_{\text{eff}}^{\text{lep}}$  (scan)  $\rightarrow \Delta A_4$ (EW, QCD) predicted  $\leftrightarrow A_4$  (measured)  $\rightarrow$  fitted  $\sin^2_{\text{eff}}^{\text{lep}}$  (best)

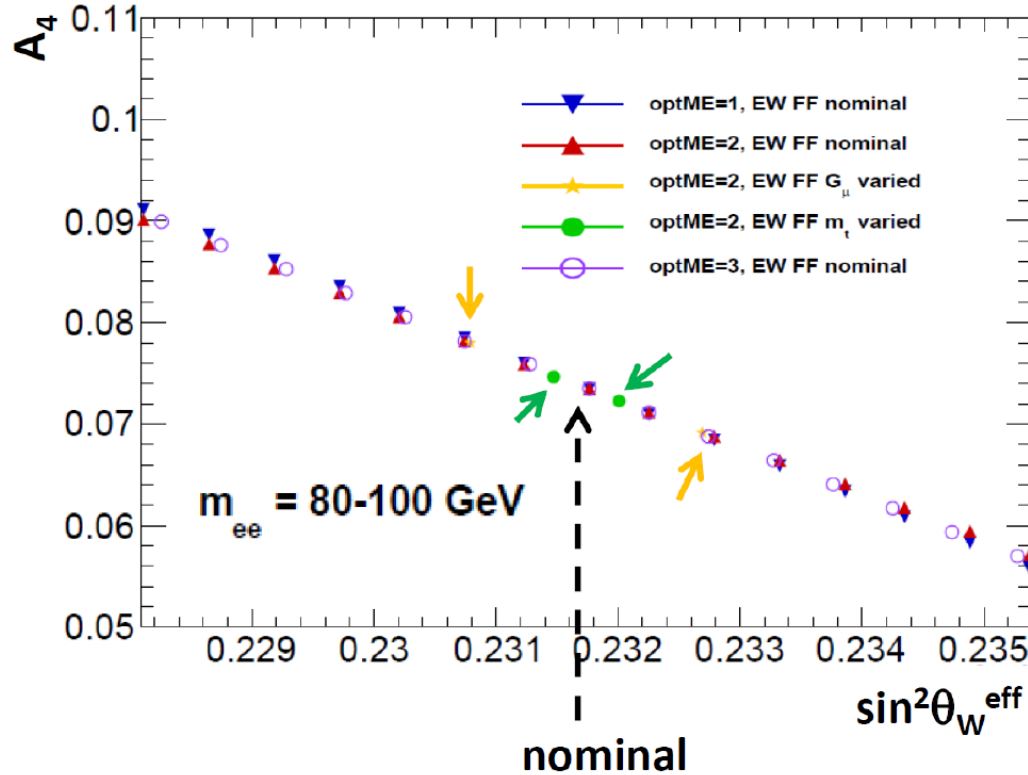
# $\sin^2\theta_{\text{eff}}$ scan for $A_{\text{FB}}$ and $A_4$

- This is not a „global fit”, but scan targeted to sensitivity to  $\sin^2\theta_W^{\text{eff}}$ 
  - One should therefore not impose EW LO relations for the scan, namely only  $\sin^2\theta_W^{\text{eff}}$  should be varied and nothing else. Is this agreed by everyone? See arguments below.
- What are options?
  - Change input parameter:  $m_t$ ,  $m_W$ ,  $G_\mu$ , recalculated EW corrections, find the „best matching”.
    - Cons: input parameters for scan outside measured values; it is indirect fit of that parameter not of the  $\sin^2\theta_W^{\text{eff}}$
  - Change  $\sin^2\theta_W^{\text{eff}}$  by adding  $\delta_v$  term, propagate to matrix element, find „best matching”
    - Cons: going beyond SM in arbitrary manner.
- Compare scans between EW LO and EW NLO+HO.
  - Each calculation should precisely determine the difference in behaviour of asymmetry versus  $\sin^2\theta_W^{\text{eff}}$ , since  $A_{\text{FB}}/A_4(\sin^2\theta_W^{\text{eff}})$  is sensitive to EW corrections.



# $\sin^2\theta_{\text{eff}}$ scan for $A_{\text{FB}}$ and $A_4$

EW corr. with  $\alpha(0)$  scheme



- Varied  $m_\tau$ , recalculated form-factors
- ★ Varied  $G_\mu$ , recalculated form-factors
- Varied  $\delta_V$
- ▼ Varied  $\delta_{S2W}$

**optME=1** ▼

$$v_\ell = (2 \cdot T_3^\ell - 4 \cdot q_\ell \cdot (s_W^2 + \delta_{S2W}) \cdot \mathcal{K}_\ell(s, t)) / \Delta,$$

$$v_f = (2 \cdot T_3^f - 4 \cdot q_f \cdot (s_W^2 + \delta_{S2W}) \cdot \mathcal{K}_f(s, t)) / \Delta,$$

$$vv_{\ell f} = \frac{1}{v_\ell \cdot v_f} [(2 \cdot T_3^\ell)(2 \cdot T_3^f) - 4 \cdot q_\ell \cdot (s_W^2 + \delta_{S2W}) \cdot \mathcal{K}_f(s, t)(2 \cdot T_3^\ell) - 4 \cdot q_f \cdot (s_W^2 + \delta_{S2W}) \cdot \mathcal{K}_\ell(s, t)(2 \cdot T_3^f) + (4 \cdot q_\ell \cdot s_W^2)(4 \cdot q_f \cdot s_W^2) \mathcal{K}_{\ell f}(s, t) + 2 \cdot (4 \cdot q_\ell)(4 \cdot q_f) \cdot s_W^2 \cdot \delta_{S2W} \mathcal{K}_{\ell f}(s, t)] \frac{1}{\Delta^2}$$

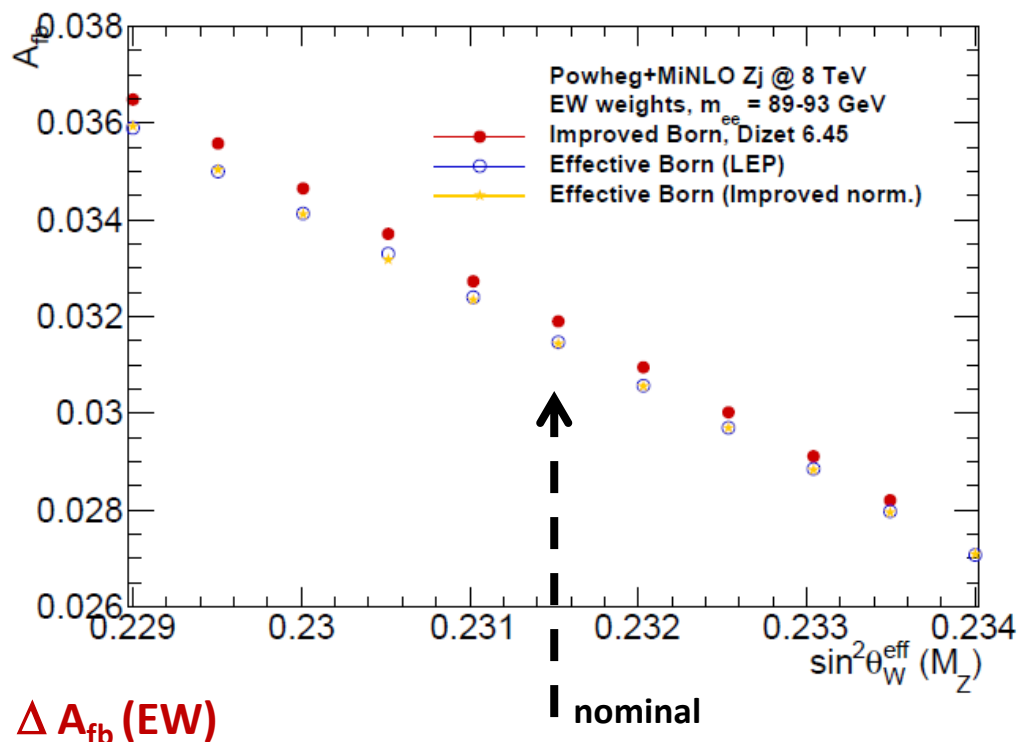
**optME=3** ○

$$v_\ell = (2 \cdot T_3^\ell - 4 \cdot q_\ell \cdot (s_W^2 \cdot \mathcal{K}_\ell(s, t) + \delta_V)) / \Delta,$$

$$v_f = (2 \cdot T_3^f - 4 \cdot q_f \cdot (s_W^2 \cdot \mathcal{K}_f(s, t) + \delta_V)) / \Delta,$$

$$vv_{\ell f} = \frac{1}{v_\ell \cdot v_f} [(2 \cdot T_3^\ell)(2 \cdot T_3^f) - 4 \cdot q_\ell \cdot (s_W^2 \cdot \mathcal{K}_f(s, t) + \delta_V)(2 \cdot T_3^\ell) - 4 \cdot q_f \cdot (s_W^2 \cdot \mathcal{K}_\ell(s, t) + \delta_V)(2 \cdot T_3^f) + (4 \cdot q_\ell \cdot s_W^2)(4 \cdot q_f \cdot s_W^2) \mathcal{K}_{\ell f}(s, t) + 2 \cdot (4 \cdot q_\ell)(4 \cdot q_f) \cdot s_W^2 \cdot \mathcal{K}_{\ell f}(s, t) \cdot \delta_V] \frac{1}{\Delta^2}$$

# $\sin^2\theta_{\text{eff}}$ scan for $A_{\text{FB}}$ and $A_4$



Effective Born <i>LEP</i>	Effective Born <i>LEP with improved norm.</i>
$\alpha = 1/128.8667$	$\alpha = 1/128.8667$
$s_W^2 = 0.23152$	$s_W^2 = 0.23152$
$\rho_{\ell f} = 1.0$	$\rho_{\ell f} = 1.005$

## Dizet v6.45

Parameter	$(\alpha(0), G_\mu, M_Z)$ $\alpha(0)$ v0
$M_Z$ (GeV)	91.1876
$1/\alpha(M_Z)$	0.0077549256
$\alpha(M_Z)$	128.9503020
$G_\mu$ ( $\text{GeV}^{-2}$ )	$1.1663787 \cdot 10^{-5}$
$M_W$ (GeV)	80.358935
$s_W^2$	0.223401084
$\sin^2\theta_{\text{eff}}^\ell$	0.231499
$\sin^2\theta_{\text{eff}}^u$	0.231392
$\sin^2\theta_{\text{eff}}^d$	0.231265
$\sin^2\theta_{\text{eff}}^b$	0.232733

$A_{\text{fb}}^{\text{ref}}(NLO+HO) - A_{\text{fb}}(\text{effect.})$	$89 < m_{ee} < 93$ GeV	$60 < m_{ee} < 81$ GeV	$81 < m_{ee} < 101$ GeV	$101 < m_{ee} < 150$ GeV
LEP	-0.00043	0.00111	0.00046	-0.00158
LEP with improv. norm.	-0.00046	0.00076	0.00047	-0.00099

## EW NLO+HO corrections vs „Effective LEP”

- Shift in predicted  $A_{\text{fb}}(M_Z)$  by  $45 \cdot 10^{-5} \rightarrow \Delta \sin^2\theta_w^{\text{eff}} = \sim 20 \cdot 10^{-5}$
- Different slope for  $A_{\text{fb}}(\sin^2\theta_w^{\text{eff}}(M_Z))$

# Status of the YR draft (v03)

 Recently updated

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# Status of the YR draft (v03)

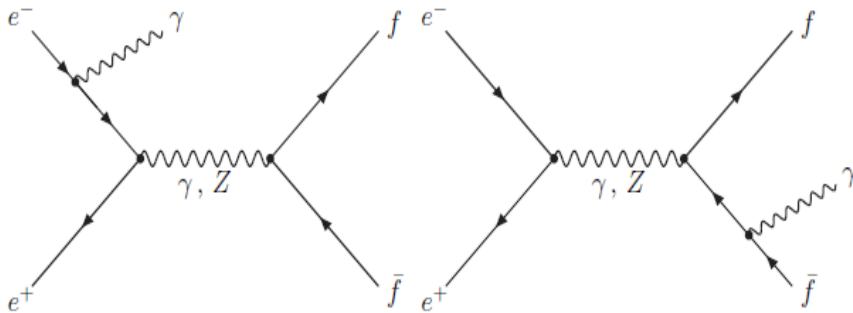
- **New since December 2019**
  - Appendix from KKMC\_hh on QED ISR/IFI
  - Updated TauSpinner + Dizet 6.45 tables/plots
  - Final proposal for mass binning of calculations and how results could be tabulated
- **Expected soon:**
  - Updates from MCSANC, Powheg\_ew:
    - tables, plots, inputs
    - write-up for Appendices
  - Feedback on the content/layout of the draft
- **Not settled yet:**
  - What is the level of agreement in predictions from different codes and EW schemes.
  - Theory and parametric uncertainties of different EW schemes.
  - Discussion on how to scan  $\sin^2\theta_w^{\text{eff}}$

# LEP legacy: QED (radiative) corrections

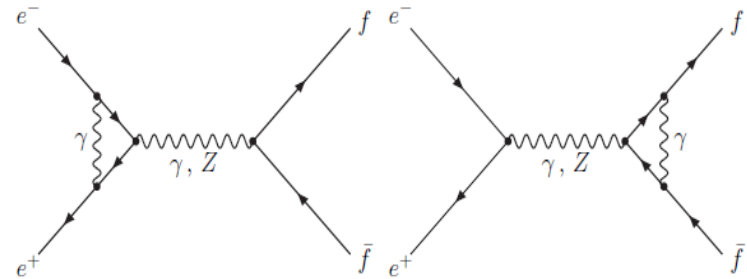
**NOT discussed here.**

QED FSR can be simulated by PHOTOS (exponentiated multi-photon emission) implemented as after-burner step on already generated event.

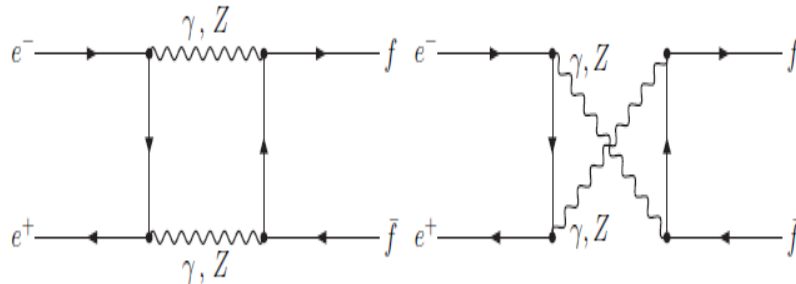
**Real emission + pairs creation**



**Vertex corrections**



**$\gamma\gamma$  and  $\gamma Z$  box diagrams**



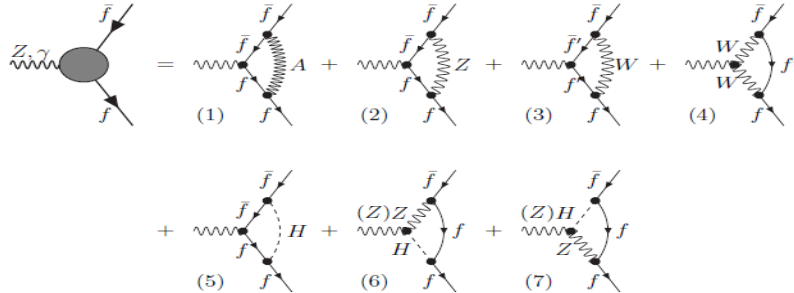
It is **QED gauge-invariant set of diagrams** (D. Bardin, hep-ph/9908433) which can be factorised out and/or convoluted with QCD corrections.

Calculated with fixed value of  $\alpha_{\text{QED}}$   
 **$\alpha_{\text{QED}} = 1./137.0359895$**

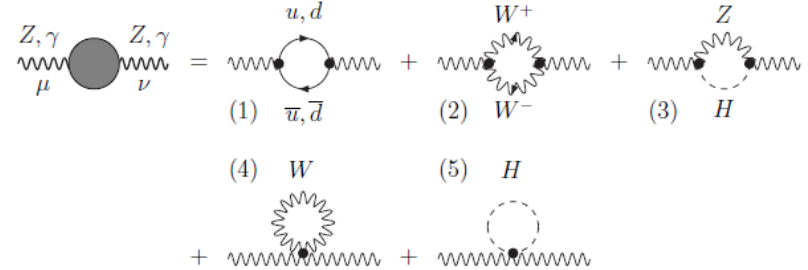
# LEP legacy: Genuine EW and lineshape corrections

Also gauge-invariant set of diagrams. Calculated as form-factor corrections to couplings, propagators and masses.  
 Eg. running  $\alpha_{\text{QED}}(s)$ ,  $\alpha_{\text{QED}}(M_Z) = 1./128.86674175$

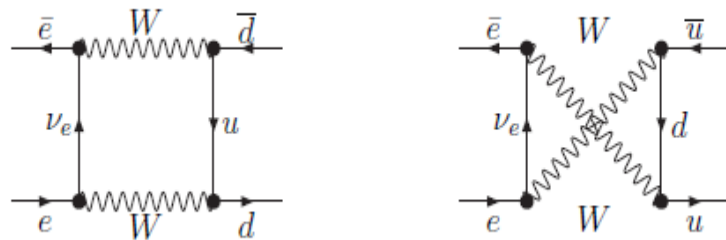
## Zff and $\gamma$ ff vertices



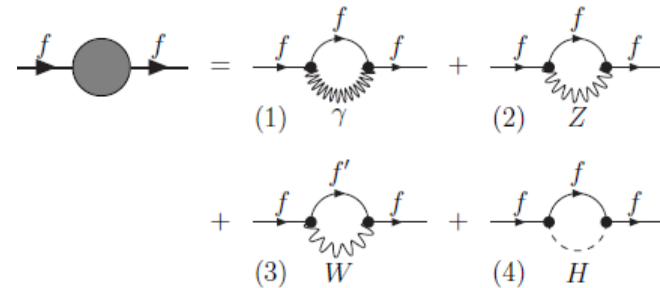
## Bosonic self-energies



## WW, ZZ boxes (shown only WW diagrams)



## Fermionic self-energies



# From Zfitter/Dizet documentation

D. Bardin et al.  
arXiv:9908433

**Zfitter** is a **semi-analytical program** for calculating total cross-sections and pseudo-observables (eg.  $A_{fb}$ ,  $\sin^2\theta_W^{\text{eff}}$ ), used by LEP1, and to a lesser degree by LEP2.

**DIZET** is a library for calculating form-factors and some other corrections. Provides complete EW  $O(\alpha)$  weak-loop corrections supplemented with selected higher order terms (eg. vacuum polarisation,  $\alpha_{\text{QED}}(Q^2)$ ).

For analyses at LEP1, LEP2 used always in parallel with **MC generators (KoralZ, KoralW)** eg. to evaluate systematics of simplified cuts used in analysis integration.

$$\begin{aligned}
 A_Z^{OLA}(s, t) = & i\sqrt{2}G_\mu I_e^{(3)} I_f^{(3)} M_Z^2 \chi_Z(s) \rho_{ef}(s, t) \left\{ \gamma_\mu(1 + \gamma_5) \otimes \gamma_\mu(1 + \gamma_5) \right. \\
 & - 4|Q_e|s_W^2 \kappa_e(s, t) \gamma_\mu \otimes \gamma_\mu(1 + \gamma_5) - 4|Q_f|s_W^2 \kappa_f(s, t) \gamma_\mu(1 + \gamma_5) \otimes \gamma_\mu \\
 & \left. + 16|Q_e Q_f|s_W^4 \kappa_{e,f}(s, t) \gamma_\mu \otimes \gamma_\mu \right\}. \tag{A.4.75}
 \end{aligned}$$

one loop amplitude

$$A_\gamma^{OLA} = i\chi_\gamma(s) \alpha(s) \gamma_\mu \otimes \gamma_\mu. \tag{2.2.36}$$

Dyson summation leads to the change of  $\alpha$  into  $\alpha(s)$ :

$$\alpha(s) = \frac{\alpha(0)}{1 - \Delta\alpha^{\text{fer}}(s)} = \frac{\alpha(0)}{1 - \Delta\alpha^{(5)}(s) - \Delta\alpha^t(s) - \Delta\alpha^{\alpha\alpha_s}(s)}. \tag{2.2.37}$$

Vacuum polarisation corrections

# LEP legacy: from Zfitter/Dizet documentation

After some trivial algebra one derives the final expressions:

$$\boxed{\rho_{ef}} = 1 + \frac{g^2}{16\pi^2} \left\{ -\Delta\rho_z^F + \mathcal{D}_z^F(s) + \frac{5}{3}B_0^F(-s; M_W, M_W) - \frac{9}{4} \frac{c_w^2}{s_w^2} \ln c_w^2 - 6 \right. \\ \left. + \frac{5}{8}c_w^2(1+c_w^2) + \frac{1}{4c_w^2}(3v_e^2+a_e^2+3v_f^2+a_f^2)\mathcal{F}_z(s) + \hat{\mathcal{F}}_w^0(s) + \hat{\mathcal{F}}_w(s) \right. \\ \left. - \frac{\tau_t}{4}[B_0^F(-s; M_W, M_W) + 1] - c_w^2(R_z - 1)s\hat{\mathcal{B}}_{ww}^d(s, t) \right\}, \quad (\text{A.4.80})$$

$$\boxed{\kappa_e} = 1 + \frac{g^2}{16\pi^2} \left\{ -\frac{c_w^2}{s_w^2}\Delta\rho^F - \Pi_{z\gamma}^F(s) - \frac{1}{6}B_0^F(-s; M_W, M_W) - \frac{1}{9} - \frac{v_e\sigma_\epsilon}{2c_w^2}\mathcal{F}_z(s) \right. \\ \left. - \hat{\mathcal{F}}_w^0(s) + (R_z - 1) \left[ \frac{|Q_f|}{2}(1 - 4|Q_f|s_w^2)\mathcal{F}_z(s) + c_w^2[\hat{\mathcal{F}}_{wn}(s) \right. \right. \\ \left. \left. - |Q_f|\mathcal{F}_{wa}(s) + s\hat{\mathcal{B}}_{ww}^d(s, t) \right] \right\}, \quad (\text{A.4.81})$$

$$\boxed{\kappa_f} = 1 + \frac{g^2}{16\pi^2} \left\{ -\frac{c_w^2}{s_w^2}\Delta\rho^F - \Pi_{z\gamma}^F(s) - \frac{1}{6}B_0^F(-s; M_W, M_W) - \frac{1}{9} - \frac{v_f\sigma_f}{2c_w^2}\mathcal{F}_z(s) \right. \\ \left. - \hat{\mathcal{F}}_w(s) + (R_z - 1) \left[ \frac{|Q_e|}{2}(1 - 4|Q_e|s_w^2)\mathcal{F}_z(s) + c_w^2[\hat{\mathcal{F}}_{wn}^0(s) \right. \right. \\ \left. \left. - |Q_e|\mathcal{F}_{wa}(s) + s\hat{\mathcal{B}}_{ww}^d(s, t) \right] - \frac{\tau_t}{4}[B_0^F(-s; M_W, M_W) + 1] \right\}, \quad (\text{A.4.82})$$

interference

$$\boxed{\kappa_{ef}} = 1 + \frac{g^2}{16\pi^2} \left\{ -2\frac{c_w^2}{s_w^2}\Delta\rho^F - 2\Pi_{z\gamma}^F(s) - \frac{1}{3}B_0^F(-s; M_W, M_W) - \frac{2}{9} \right. \\ \left. - \frac{1}{4c_w^2} \left[ \frac{\delta_e^2 + \delta_f^2}{s_w^2}(R_w - 1) + 3v_e^2 + a_e^2 + 3v_f^2 + a_f^2 \right] \mathcal{F}_z(s) \right. \\ \left. - \hat{\mathcal{F}}_w^0(s) - \hat{\mathcal{F}}_w(s) - \frac{\tau_t}{4}[B_0^F(-s; M_W, M_W) + 1] \right. \\ \left. + c_w^2(R_z - 1) \left[ \frac{2}{3} - \hat{\Pi}_{\gamma\gamma}^{\text{bos},F}(s) + s\hat{\mathcal{B}}_{ww}^d(s, t) \right] \right\}. \quad (\text{A.4.83})$$

Fermionic loops in  $\gamma$  propagator

BOX

LHC EWWG meeting, 27.03.2020

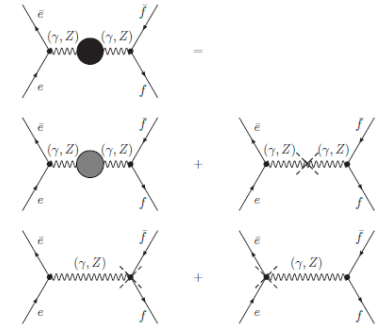


Figure A.11. Bosonic self-energies and bosonic counter-terms for  $e\bar{e} \rightarrow (Z, \gamma) \rightarrow f\bar{f}$

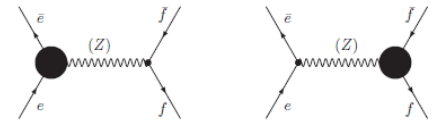


Figure A.10. Electron (a) and final fermion (b) vertices in  $e\bar{e} \rightarrow (Z) \rightarrow f\bar{f}$

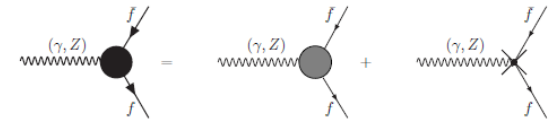


Figure A.6. Off-shell  $Zff$  and  $\gamma ff$  vertices



Figure A.7. The  $WW$  boxes

etc. etc.



# Constructing $wt^{EW}$ : EW Improved Born (IBA)

ERW and Z.Was,  
arXiv: 1808.08616

$$\mathcal{A}^{Born+EW} = \frac{\alpha}{s} \left\{ \begin{aligned} & [\bar{u}\gamma^\mu v g_{\mu\nu} \bar{v}\gamma^\nu u] \cdot (q_\ell \cdot q_f) \Gamma_{V\Pi} \chi_\gamma(s) \\ & + [\bar{u}\gamma^\mu v g_{\mu\nu} \bar{v}\gamma^\nu u \cdot (v_\ell \cdot v_f \cdot vv_{\ell f}) + \bar{u}\gamma^\mu v g_{\mu\nu} \bar{v}\gamma^\nu \gamma^5 u \cdot (v_\ell \cdot a_f) \\ & + \bar{u}\gamma^\mu \gamma^5 v g_{\mu\nu} \bar{v}\gamma^\nu u \cdot (a_\ell \cdot v_f) + \bar{u}\gamma^\mu \gamma^5 v g_{\mu\nu} \bar{v}\gamma^\nu \gamma^5 u \cdot (a_\ell \cdot a_f)] \cdot Z_{V\Pi} \chi_Z(s) \end{aligned} \right\}$$

$$\chi_\gamma(s) = 1$$

$$\chi_Z(s) = \frac{G_\mu \cdot M_Z^2 \cdot \Delta^2}{\sqrt{2} \cdot 8\pi \cdot \alpha} \cdot \frac{s}{s - M_Z^2 + i \cdot \Gamma_Z \cdot M_Z}$$

$$v_\ell = (2 \cdot T_3^\ell - 4 \cdot q_\ell \cdot s_W^2) \cdot K_\ell(s, t) / \Delta$$

$$v_f = (2 \cdot T_3^f - 4 \cdot q_f \cdot s_W^2) \cdot K_f(s, t) / \Delta$$

$$a_\ell = (2 \cdot T_3^\ell) / \Delta$$

$$a_f = (2 \cdot T_3^f) / \Delta$$

$$\Delta = \sqrt{16 \cdot s_W^2 \cdot (1 - s_W^2)}$$

$$Z_{V\Pi} = \rho_{e,f}(s, t)$$

**EW form-factors, functions of  $(s, t) = (m_{\Pi}, \cos\theta)$   
Calculated with Dizeit 6.21 library.**

$$\Gamma_{V\Pi} = \frac{1}{2 - (1 + \Pi_{\gamma\gamma}(s))}$$

Vacuum polarisation corrections, used low-energy experiment input.

Warning: problem for analytic continuation.

$$vv_{\ell f} = \frac{1}{v_\ell \cdot v_f} \left[ (2 \cdot T_3^\ell)(2 \cdot T_3^f) - 4 \cdot q_\ell \cdot s_W^2 \cdot K_f(s, t) (2 \cdot T_3^\ell) - 4 \cdot q_f \cdot s_W^2 \cdot K_\ell(s, t) (2 \cdot T_3^f) + (4 \cdot q_\ell \cdot s_W^2)(4 \cdot q_f \cdot s_W^2) K_{\ell f}(s, t) \right] \frac{1}{\Delta^2}$$

# EW schemes: details

## EW schemes: come with „on-shell” or „pole” definitions!

Table 44: The EW parameters used at tree-level EW, with on-mass-shell definition (LEP convention).

Parameter	$\alpha(0) \text{ v0}$	$\alpha(0) \text{ v1}$	$G_\mu$
$M_Z$	91.1876 GeV	91.1876 GeV	91.1876 GeV
$\Gamma_Z$	2.4952 GeV	2.4952 GeV	2.4952 GeV
$\Gamma_W$	2.085 GeV	2.085 GeV	2.085 GeV
$\alpha$	1/137.03599	1/137.03599	1/132.23323
$G_\mu$	$1.1663787 \cdot 10^{-5} \text{ GeV}^{-2}$	$1.1254734 \cdot 10^{-5} \text{ GeV}^{-2}$	$1.1663787 \cdot 10^{-5} \text{ GeV}^{-2}$
$M_W$	80.93886 GeV	80.385 GeV	80.385 GeV
$s_W^2$	0.2121517	0.2228972	0.2228972
$\frac{G_\mu \cdot M_Z^2 \cdot \Delta^2}{\sqrt{2} \cdot 8\pi \cdot \alpha}$	1.0	1.0	1.0

Running  $\Gamma_Z$  in  
Z-propagator

Shift:

- -30 MeV for  $M_Z$
- change on  $\Gamma_Z$
- -0.00006 for  $s_W^2$

Scaling

- 0.99906 for  $\alpha$

Fixed  $\Gamma_Z$  in  
Z-propagator

Table 45: The EW parameters used at tree-level EW, with pole definition of the Z, W masses.

Parameter	$\alpha(0) \text{ v0}$	$\alpha(0) \text{ v1}$	$G_\mu$
$M_Z$	91.15348 GeV	91.15348 GeV	91.15348 GeV
$\Gamma_Z$	2.494266 GeV	2.494266	2.494266 GeV
$\Gamma_W$	2.085 GeV	2.085 GeV	2.085 GeV
$\alpha$	1/137.03599	1/137.03599	1/132.3572336357709
$G_\mu$	$1.1663787 \cdot 10^{-5} \text{ GeV}^{-2}$	$1.126555497 \cdot 10^{-5} \text{ GeV}^{-2}$	$1.1663787 \cdot 10^{-5} \text{ GeV}^{-2}$
$M_W$	80.91191 GeV	80.35797 GeV	80.35797 GeV
$s_W^2$	0.21208680	0.22283820939	0.22283820939
$\frac{G_\mu \cdot M_Z^2 \cdot \Delta^2}{\sqrt{2} \cdot 8\pi \cdot \alpha}$	1.0	1.0	1.0

# Form of the Z-boson propagator

- **Discussed since fall last year, problem in nutshell**
  - **LEP1 legacy (Dizet+Zfitter, experiments):**
    - use running width in the Born propagator
    - form-factors calculated with pole-mass/fixed width (internally converted), applied to Born with on-shell mass/running width
    - see references: hep-ex/0509008, hep-ph/9908433
  - **LEP2, LHC standard**
    - use complex-mass scheme, pole masses, fixed width propagator
  - **Zfitter+Dizet v6.42, v6.45, FCCee standard**
    - stayed with LEP1 convention

**Is that a concern for  $\sin^2\theta_{\text{eff}}$  measurement at LHC ?**

# Z-boson propagator

Topic discussed in Fulvio's talks at EW meetings on 13.03, 7.05 and 1.07

How to model „resonance”

Is the Breit-Wigner form good enough?

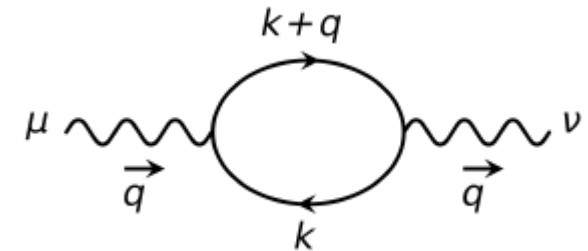
idea behind running width

$$\sigma_{\text{ff}}^Z = \sigma_{\text{ff}}^{\text{peak}} \frac{s\Gamma_Z^2}{(s - M_Z^2)^2 + s^2\Gamma_Z^2/M_Z^2} = \sigma_{\text{ff}}^{\text{peak}} \frac{s\Gamma_Z^2}{(s - M_Z^2)^2 + \Gamma(s)^2 M_Z^2}$$

from my slides at our 13 March 2019 meeting

- $M_Z$  and  $\Gamma_Z$  above are “OS” quantities,  $\Gamma(s) = \Gamma \frac{s}{M^2}$
- let's express  $\sigma_{\text{ff}}^Z$  in terms of “pole” quantities, with  $\gamma \equiv \frac{\Gamma_{\text{pole}}}{M_{\text{pole}}}$

$$\begin{aligned} \sigma_{\text{ff}}^Z &= \sigma_{\text{ff}}^{\text{peak}} \frac{s\Gamma_{\text{pole}}^2(1 + \gamma^2)}{(s - M_{\text{pole}}^2(1 + \gamma^2))^2 + s^2\gamma^2} \\ &= \sigma_{\text{ff}}^{\text{peak}} \frac{s\Gamma_{\text{pole}}^2(1 + \gamma^2)}{s^2 + M_{\text{pole}}^4(1 + \gamma^2)^2 - 2sM_{\text{pole}}^2(1 + \gamma^2) + s^2\gamma^2} \\ &= \sigma_{\text{ff}}^{\text{peak}} \frac{s\Gamma_{\text{pole}}^2(1 + \gamma^2)}{s^2(1 + \gamma^2) + M_{\text{pole}}^4(1 + \gamma^2)^2 - 2sM_{\text{pole}}^2(1 + \gamma^2)} \\ &= \sigma_{\text{ff}}^{\text{peak}} \frac{s\Gamma_{\text{pole}}^2}{(s - M_{\text{pole}}^2)^2 + \Gamma_{\text{pole}}^2 M_{\text{pole}}^2} \end{aligned}$$



$$\begin{aligned} M_{OS}^2 &= M_{\text{pole}}^2 \left( 1 + \frac{\Gamma_{\text{pole}}^2}{M_{\text{pole}}^2} \right) \\ \Gamma_{OS}^2 &= \Gamma_{\text{pole}}^2 \left( 1 + \frac{\Gamma_{\text{pole}}^2}{M_{\text{pole}}^2} \right) \end{aligned}$$

But the propagator in ME is of the form

$$\chi_Z(s) = \frac{1}{s - M_Z^2 + i \cdot \Gamma_Z \cdot s / M_Z}$$

# Z-boson propagator

Topic discussed in Fuvio talks at EW meetings on 13.03, 7.05 and 1.07

including photon exchange

$$\frac{d\sigma_0^\gamma}{d\Omega} = \frac{\alpha^2 Q_f^2 N_c}{4s} (1 + \cos^2 \vartheta)$$

$$\frac{d\sigma_0^{\gamma Z}}{d\Omega} = -\frac{\alpha^2 Q_f N_c}{4\sqrt{2}s_\theta^2 c_\theta^2 s} \text{Re}(\chi(s)) [g_V^e g_V^f (1 + \cos^2 \vartheta) + 2g_A^e g_A^f \cos \vartheta]$$

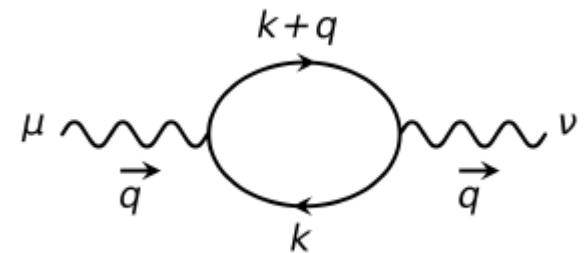
$$\frac{d\sigma_0^Z}{d\Omega} = -\frac{\pi\alpha^2 N_c}{32s_\theta^4 c_\theta^4 s} |\chi(s)|^2 [f(g_V^{e,f}, g_A^{e,f})(1 + \cos^2 \vartheta) + g(g_V^{e,f}, g_A^{e,f}) \cos \vartheta]$$

$$\chi(s) = \frac{s}{(s - M_Z^2) + i\Gamma_Z M_Z}$$

$$\chi(s)_{\text{running}} = \frac{1}{(1 + i\gamma)} \chi(s)_{\text{pole}} \quad \gamma \simeq 0.0274$$

- the couplings, in schemes where  $\sin^2 \theta$  is connected to  $M_Z$  and  $M_W$ , get modified when changing from running- to fixed-width scheme
- the relative weights of channels can get modified

idea behind running width



$$M_{OS}^2 = M_{\text{pole}}^2 \left( 1 + \frac{\Gamma_{\text{pole}}^2}{M_{\text{pole}}^2} \right)$$

$$\Gamma_{OS}^2 = \Gamma_{\text{pole}}^2 \left( 1 + \frac{\Gamma_{\text{pole}}^2}{M_{\text{pole}}^2} \right)$$

!!!

# Z-boson propagator

Mathematically formulas for  $\chi(s)$  are equivalent, ones  $M_Z, \Gamma_Z, N_Z$  are properly implemented.

At the Z-pole both formulas should lead to same calculated cross-section.

Running width

$$\chi_Z(s) = \frac{1}{s - M_Z^2 + i \cdot \Gamma_Z \cdot s / M_Z}$$



$$\begin{aligned} \chi_Z(s) &= \frac{1}{s(1 + i \cdot \Gamma_Z / M_Z) - M_Z^2} \\ &= \frac{(1 - i \cdot \Gamma_Z / M_Z)}{s(1 + \Gamma_Z^2 / M_Z^2) - M_Z^2(1 - i \cdot \Gamma_Z / M_Z)} \\ &= \frac{(1 - i \cdot \Gamma_Z / M_Z)}{(1 + \Gamma_Z^2 / M_Z^2)} \frac{1}{s - \frac{M_Z^2}{1 + \Gamma_Z^2 / M_Z^2} + i \cdot \frac{\Gamma_Z M_Z}{1 + \Gamma_Z^2 / M_Z^2}} \end{aligned}$$

Fixed width

$$\chi'_Z(s) = \frac{1}{s - M'_Z{}^2 + i \cdot \Gamma'_Z \cdot M'_Z}$$

$$\begin{aligned} &= N_Z \frac{1}{s - M'_Z{}^2 + i \Gamma'_Z M'_Z} && \text{Equivalent to} \\ &M'_Z = \frac{M_Z}{\sqrt{1 + \Gamma_Z^2 / M_Z^2}} && \text{running width} \\ &\Gamma'_Z = \frac{\Gamma_Z}{\sqrt{1 + \Gamma_Z^2 / M_Z^2}} \\ &N_Z = \frac{(1 - i \cdot \Gamma_Z / M_Z)}{(1 + \Gamma_Z^2 / M_Z^2)} = \frac{(1 - i \cdot \Gamma'_Z / M'_Z)}{(1 + \Gamma'_Z{}^2 / M'_Z{}^2)} \end{aligned}$$

# Z-boson propagator

Running width, on-shell  $M_Z, \Gamma_Z$

$$\chi_Z(s) = \frac{1}{s - M_Z^2 + i \cdot \Gamma_Z \cdot s / M_Z}$$

Fixed width, on-shell  $M_Z, \Gamma_Z$

$$\chi'_Z(s) = \frac{1}{s - M_Z^2 + i \cdot \Gamma_Z \cdot M_Z}$$

Equivalent to  $\chi'_Z(s)$

shifted  $M_Z, \Gamma_Z$ , scaled by  $N_Z$

$$\chi_Z(s) = N'_Z \frac{1}{s - M'^2_Z + i \Gamma'_Z M'_Z}$$

$$M'_Z = \frac{M_Z}{\sqrt{1 + \Gamma_Z^2 / M_Z^2}}$$

$$\Gamma'_Z = \frac{\Gamma_Z}{\sqrt{1 + \Gamma_Z^2 / M_Z^2}}$$

$$N'_Z = \frac{(1 - i \cdot \Gamma_Z / M_Z)}{(1 + \Gamma_Z^2 / M_Z^2)} = \frac{(1 - i \cdot \Gamma'_Z / M'_Z)}{(1 + \Gamma'^2_Z / M'^2_Z)}$$

Shifted  $M_Z, \Gamma_Z$ , no scaling  $N_Z$

$$\chi_Z(s) = \frac{1}{s - M'^2_Z + i \Gamma'_Z M'_Z}$$

$$M'_Z = \frac{M_Z}{\sqrt{1 + \Gamma_Z^2 / M_Z^2}}$$

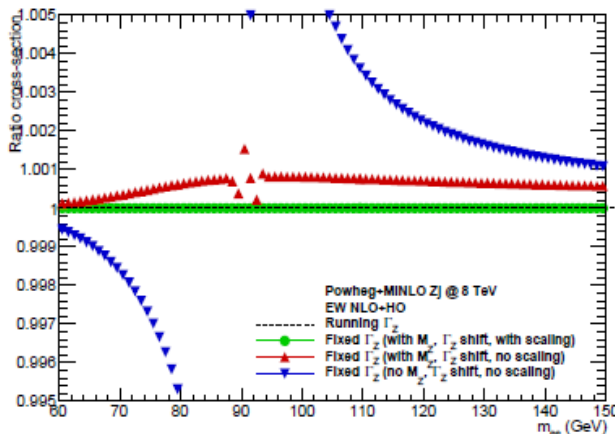
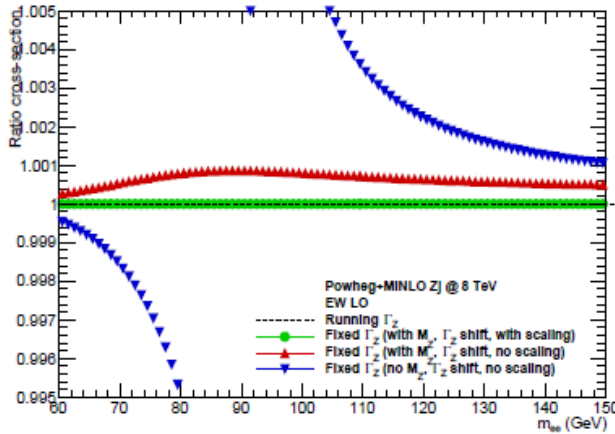
$$\Gamma'_Z = \frac{\Gamma_Z}{\sqrt{1 + \Gamma_Z^2 / M_Z^2}}$$

# Z-boson propagator

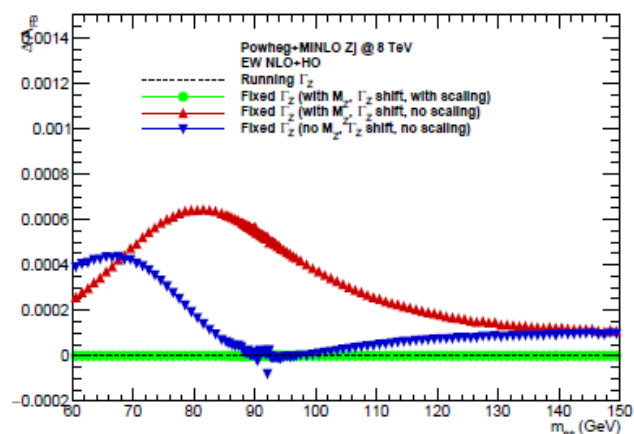
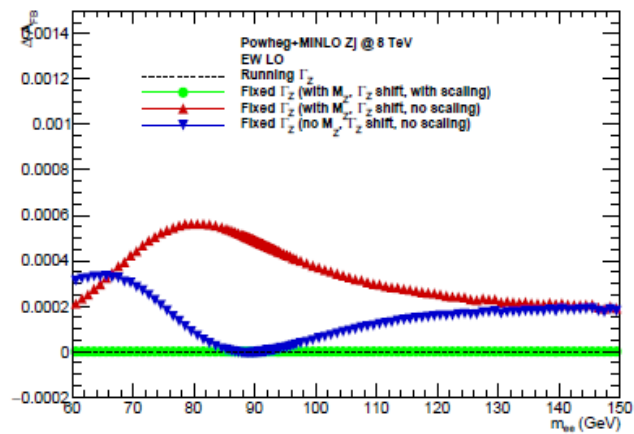
EW LO

EW NLO+HO

Cross-section ratios



$\Delta A_{fb}$



- Powheg+MiNLO Zj @ 8 TeV
- EW LO
  - Running  $\Gamma_Z$
  - Fixed  $\Gamma_Z$  (with  $M_Z, \Gamma_Z$  shift, with scaling)
  - Fixed  $\Gamma_Z$  (with  $M_Z, \Gamma_Z$  shift, no scaling)
  - Fixed  $\Gamma_Z$  (no  $M_Z, \Gamma_Z$  shift, no scaling)



# Z-boson propagator

Table 15: Ratio of the cross-section  $\sigma$  calculated with different form Z-boson propagator, integrated over specified mass windows. Shown in case of EW LO and EW NLO+HO predictions with  $O(\alpha(0))$  v0 EW scheme.

$\sigma$ (Fixed/Running)	$89 < m_{ee} < 93$ GeV	$60 < m_{ee} < 81$ GeV	$81 < m_{ee} < 101$ GeV	$101 < m_{ee} < 150$ GeV
EW LO				
with $M_Z, \Gamma_Z$ shift, no scaling	1.00087	1.00062	1.00086	1.00071
no $M_Z, \Gamma_Z$ shift, no scaling	1.00074	0.99716	0.99977	1.00392
EW NLO+HO				
with $M_Z, \Gamma_Z$ shift, no scaling	1.00085	1.00043	1.00083	1.00075
no $M_Z, \Gamma_Z$ shift, no scaling	1.00122	0.99719	1.00013	1.00392

Table 16: Difference in  $A_{fb}$  calculated with different form of Z-boson propagator, integrated over specified mass windows. Shown in case of EW LO and EW NLO+HO predictions with  $O(\alpha(0))$  v0 EW scheme.

$\Delta A_{fb}$ (Running - Fixed)	$89 < m_{ee} < 93$ GeV	$60 < m_{ee} < 81$ GeV	$81 < m_{ee} < 101$ GeV	$101 < m_{ee} < 150$ GeV
EW LO				
with $M_Z, \Gamma_Z$ shift, no scaling	-0.00047	-0.00047	-0.00047	-0.00030
no $M_Z, \Gamma_Z$ shift, no scaling	-0.00026	-0.00012	-0.00040	-0.00005
EW NLO+HO				
with $M_Z, \Gamma_Z$ shift, no scaling	-0.00053	-0.00052	-0.00053	-0.00024
no $M_Z, \Gamma_Z$ shift, no scaling	-0.00030	-0.00026	-0.00048	-0.00004