



Beam-beam induced crabbing

X. Buffat

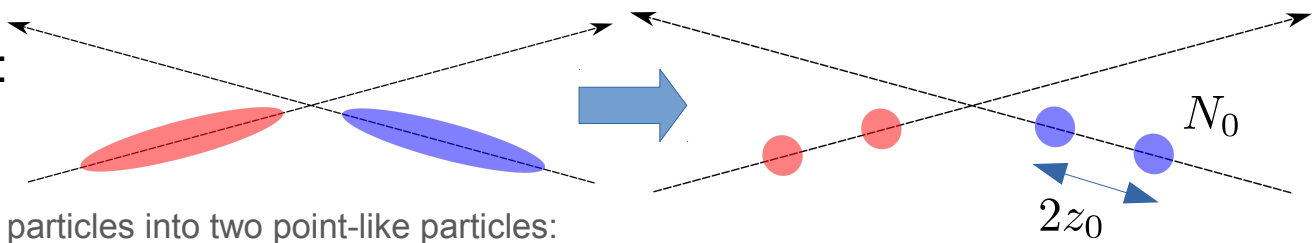


HL-LHC WP2 meeting 21.04.2020

Content

- A pessimistic model based on a simplified incoherent 6D beam-beam kick using the two particle model
- Coherent kick based on Hirata's 6D approach
- Results for the HL-LHC at start and end of collision

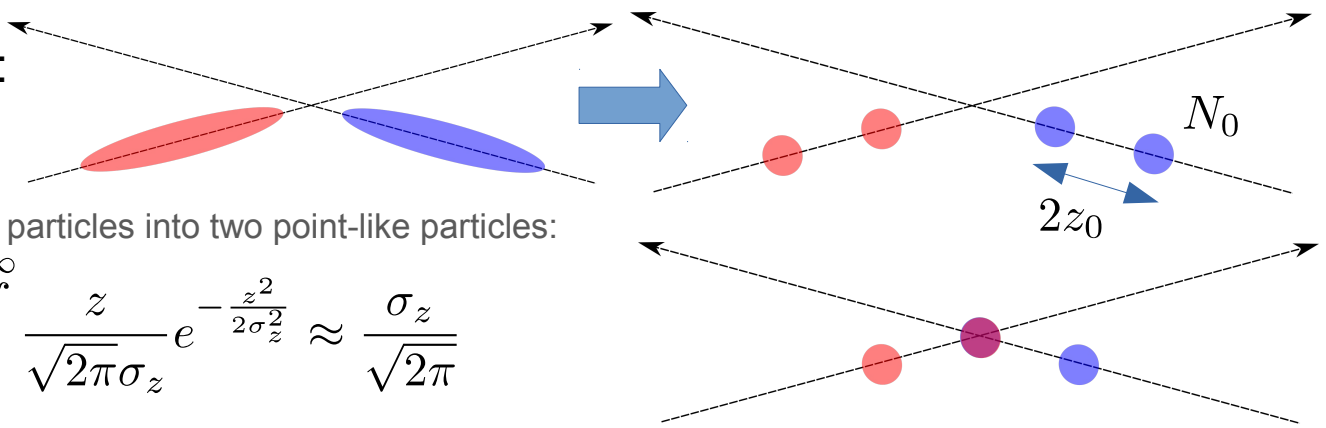
Two particle model : (Incoherent)



Reducing the distribution of N particles into two point-like particles:

$$N_0 = \frac{N}{2}, \quad z_0 = \int_0^{\infty} \frac{z}{\sqrt{2\pi}\sigma_z} e^{-\frac{z^2}{2\sigma_z^2}} \approx \frac{\sigma_z}{\sqrt{2\pi}}$$

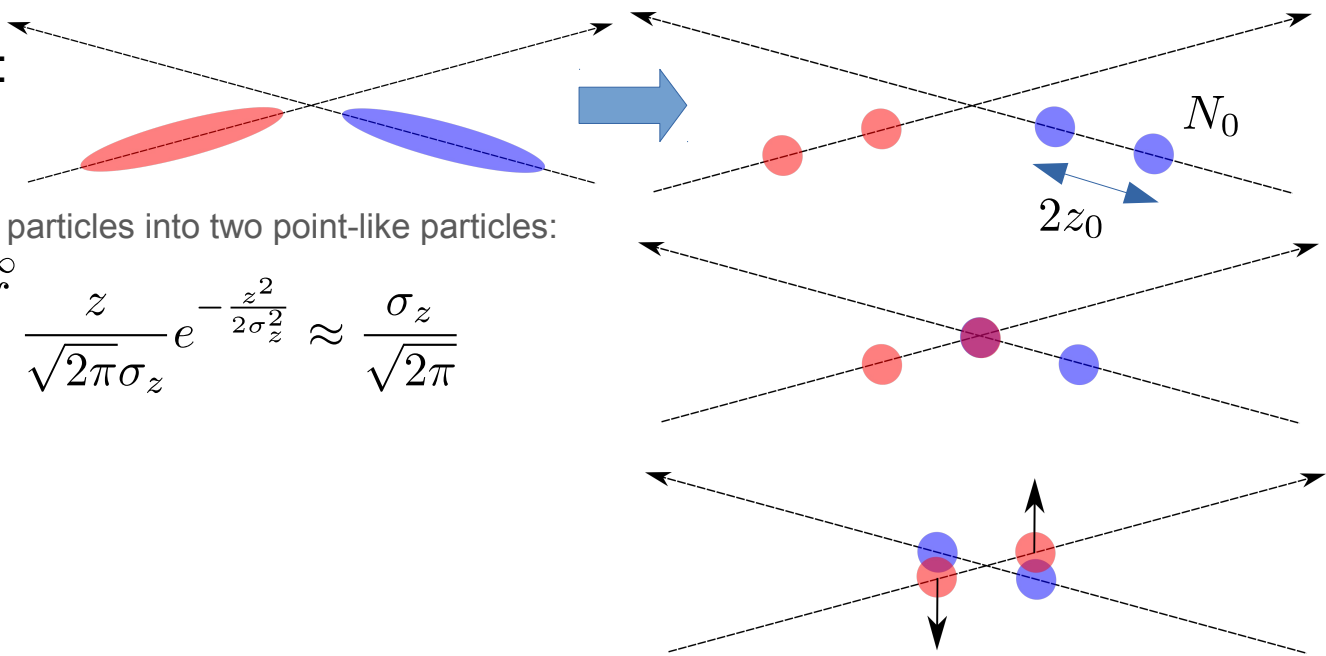
Two particle model : (Incoherent)



Reducing the distribution of N particles into two point-like particles:

$$N_0 = \frac{N}{2}, \quad z_0 = \int_0^{\infty} \frac{z}{\sqrt{2\pi}\sigma_z} e^{-\frac{z^2}{2\sigma_z^2}} \approx \frac{\sigma_z}{\sqrt{2\pi}}$$

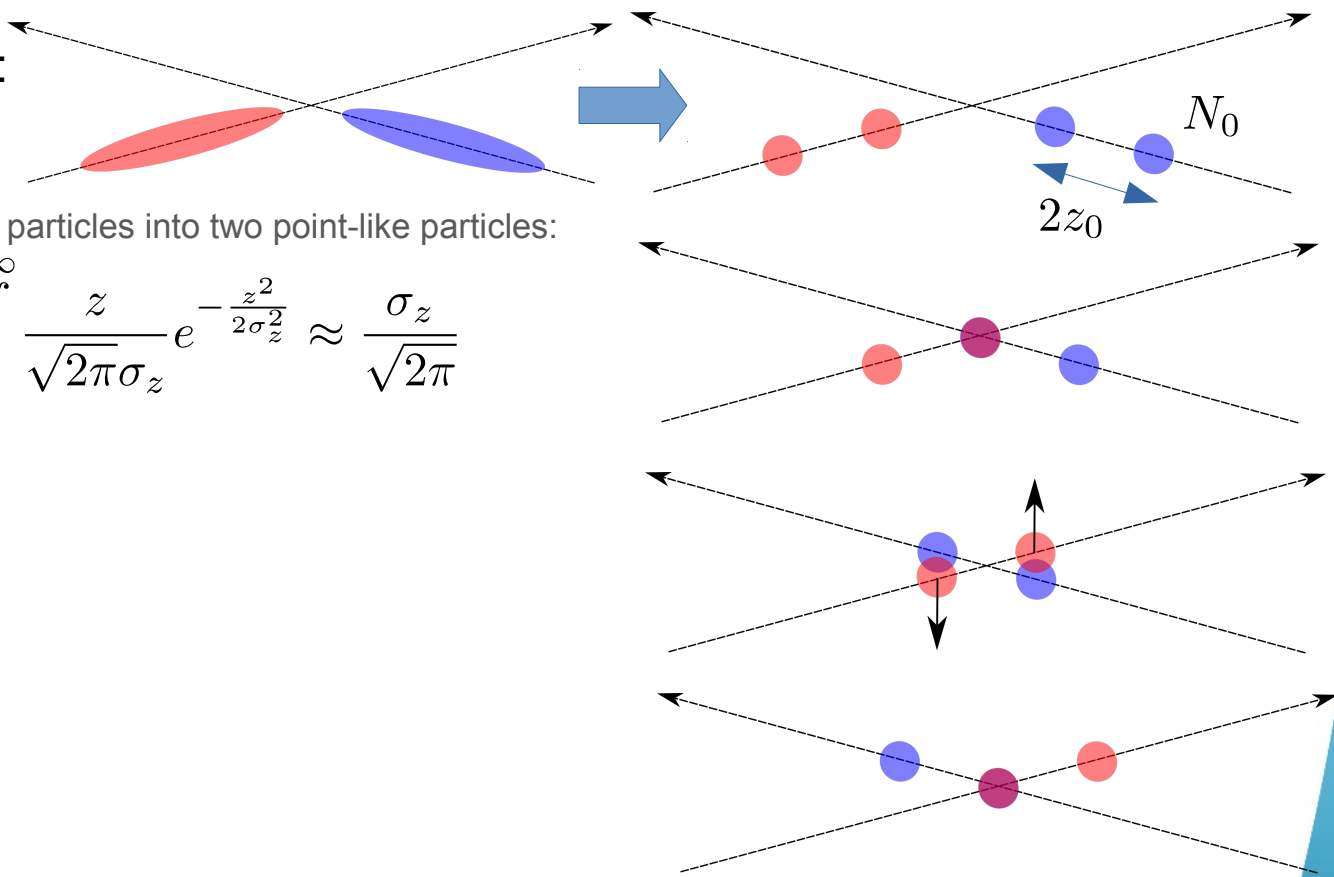
Two particle model : (Incoherent)



Reducing the distribution of N particles into two point-like particles:

$$N_0 = \frac{N}{2}, \quad z_0 = \int_0^{\infty} \frac{z}{\sqrt{2\pi}\sigma_z} e^{-\frac{z^2}{2\sigma_z^2}} \approx \frac{\sigma_z}{\sqrt{2\pi}}$$

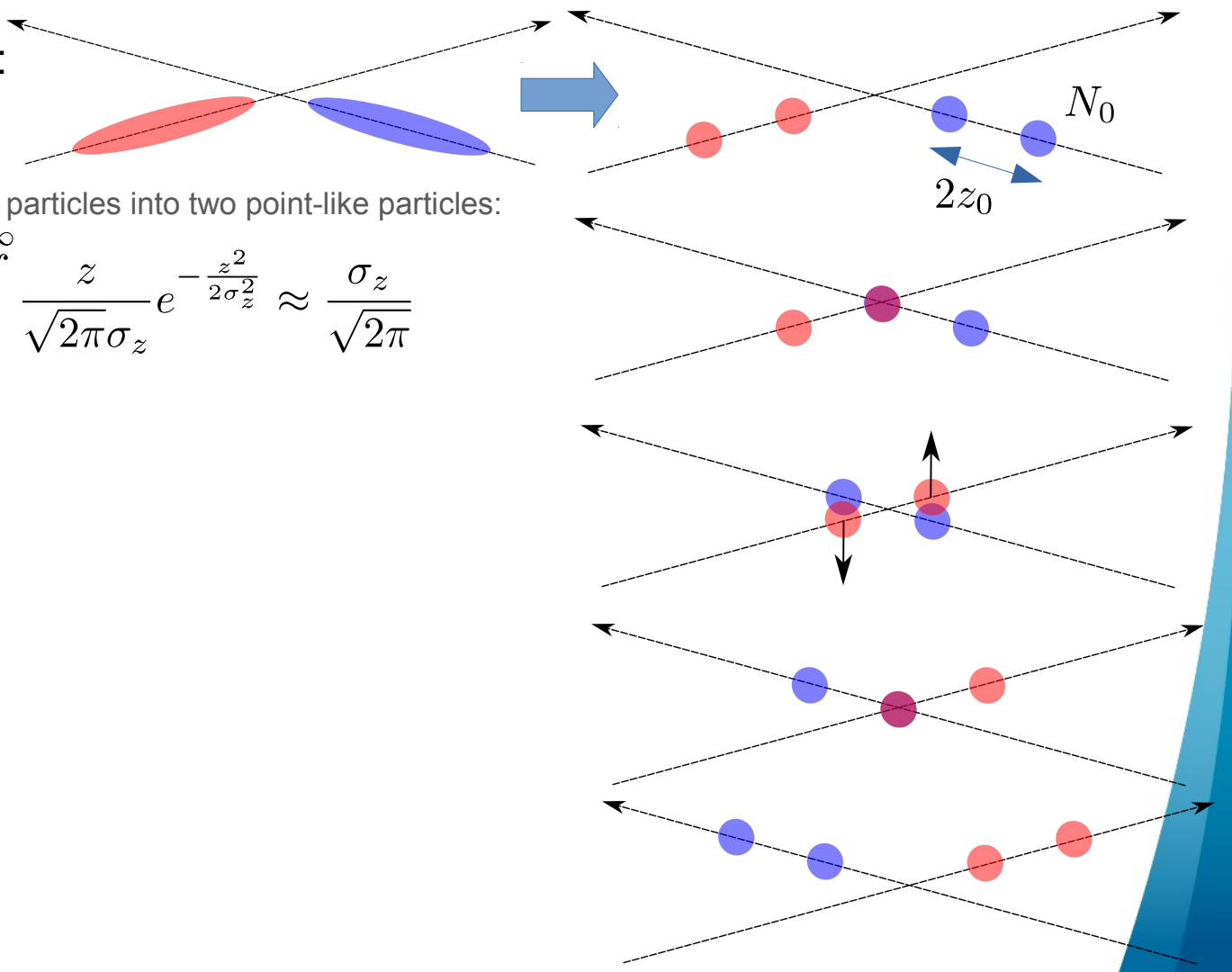
Two particle model : (Incoherent)



Reducing the distribution of N particles into two point-like particles:

$$N_0 = \frac{N}{2}, \quad z_0 = \int_0^{\infty} \frac{z}{\sqrt{2\pi}\sigma_z} e^{-\frac{z^2}{2\sigma_z^2}} \approx \frac{\sigma_z}{\sqrt{2\pi}}$$

Two particle model : (Incoherent)



Reducing the distribution of N particles into two point-like particles:

$$N_0 = \frac{N}{2}, \quad z_0 = \int_0^{\infty} \frac{z}{\sqrt{2\pi}\sigma_z} e^{-\frac{z^2}{2\sigma_z^2}} \approx \frac{\sigma_z}{\sqrt{2\pi}}$$

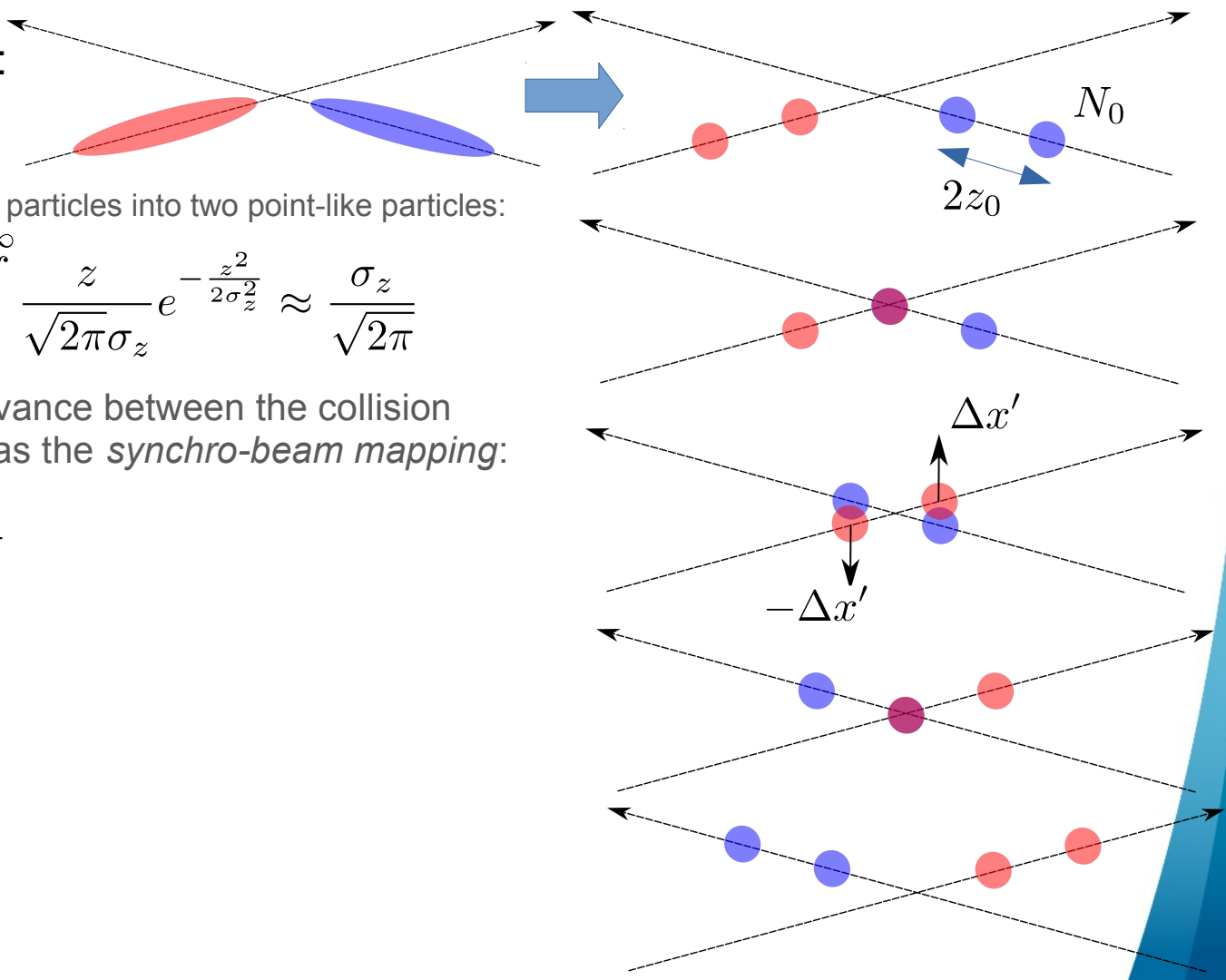
Two particle model : (Incoherent)

- Reducing the distribution of N particles into two point-like particles:

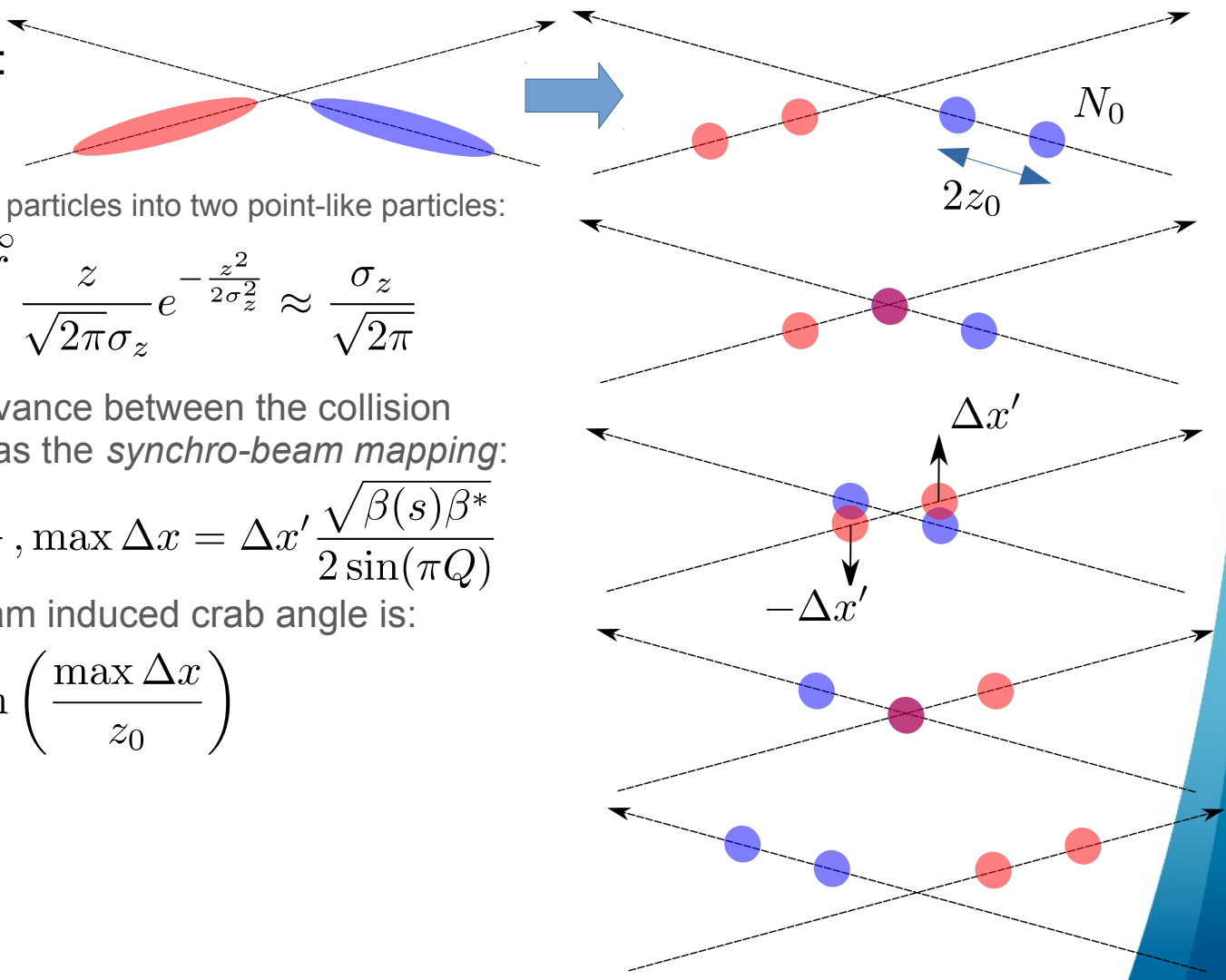
$$N_0 = \frac{N}{2}, \quad z_0 = \int_0^\infty \frac{z}{\sqrt{2\pi}\sigma_z} e^{-\frac{z^2}{2\sigma_z^2}} \approx \frac{\sigma_z}{\sqrt{2\pi}}$$

- Neglecting the phase advance between the collision point and the IP, as well as the *synchro-beam mapping*:

$$\max \Delta x' \approx 0.9 \frac{N_0 r_p}{\gamma \sigma}$$



Two particle model : (Incoherent)



- Reducing the distribution of N particles into two point-like particles:

$$N_0 = \frac{N}{2}, \quad z_0 = \int_0^{\infty} \frac{z}{\sqrt{2\pi}\sigma_z} e^{-\frac{z^2}{2\sigma_z^2}} \approx \frac{\sigma_z}{\sqrt{2\pi}}$$

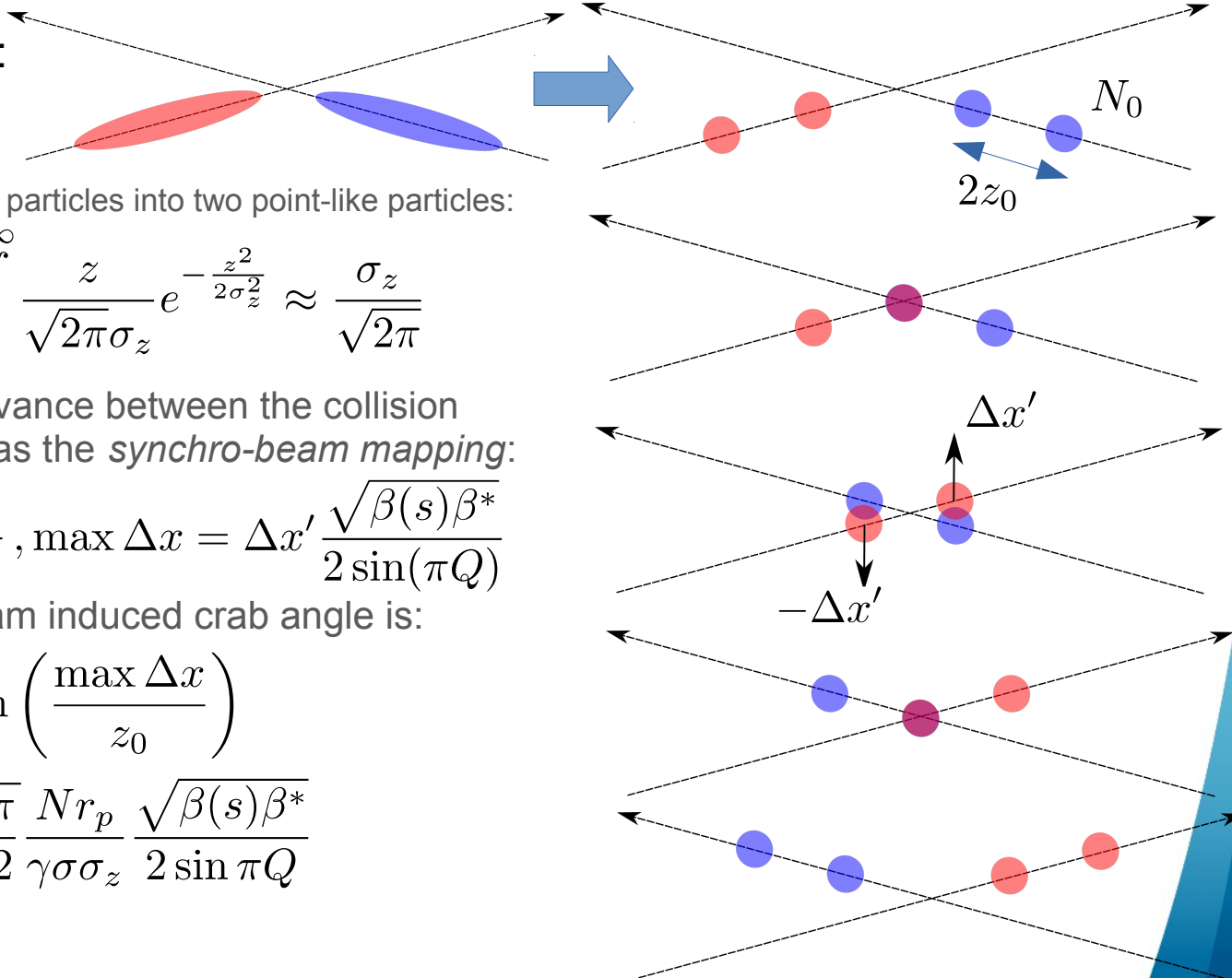
- Neglecting the phase advance between the collision point and the IP, as well as the *synchro-beam mapping*:

$$\max \Delta x' \approx 0.9 \frac{N_0 r_p}{\gamma \sigma}, \quad \max \Delta x = \Delta x' \frac{\sqrt{\beta(s)\beta^*}}{2 \sin(\pi Q)}$$

- The maximum beam-beam induced crab angle is:

$$\max \phi = \arctan \left(\frac{\max \Delta x}{z_0} \right)$$

Two particle model : (Incoherent)



- Reducing the distribution of N particles into two point-like particles:

$$N_0 = \frac{N}{2}, \quad z_0 = \int_0^\infty \frac{z}{\sqrt{2\pi}\sigma_z} e^{-\frac{z^2}{2\sigma_z^2}} \approx \frac{\sigma_z}{\sqrt{2\pi}}$$

- Neglecting the phase advance between the collision point and the IP, as well as the *synchro-beam mapping*:

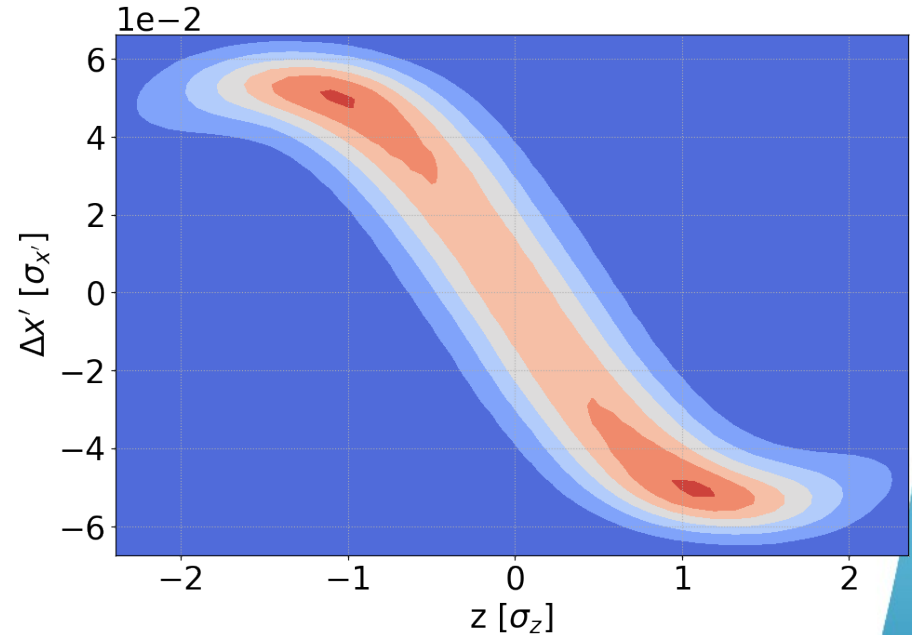
$$\max \Delta x' \approx 0.9 \frac{N_0 r_p}{\gamma \sigma}, \quad \max \Delta x = \Delta x' \frac{\sqrt{\beta(s)\beta^*}}{2 \sin(\pi Q)}$$

- The maximum beam-beam induced crab angle is:

$$\begin{aligned} \max \phi &= \arctan \left(\frac{\max \Delta x}{z_0} \right) \\ &\approx 0.9 \sqrt{\frac{\pi}{2}} \frac{N r_p}{\gamma \sigma \sigma_z} \frac{\sqrt{\beta(s)\beta^*}}{2 \sin \pi Q} \end{aligned}$$

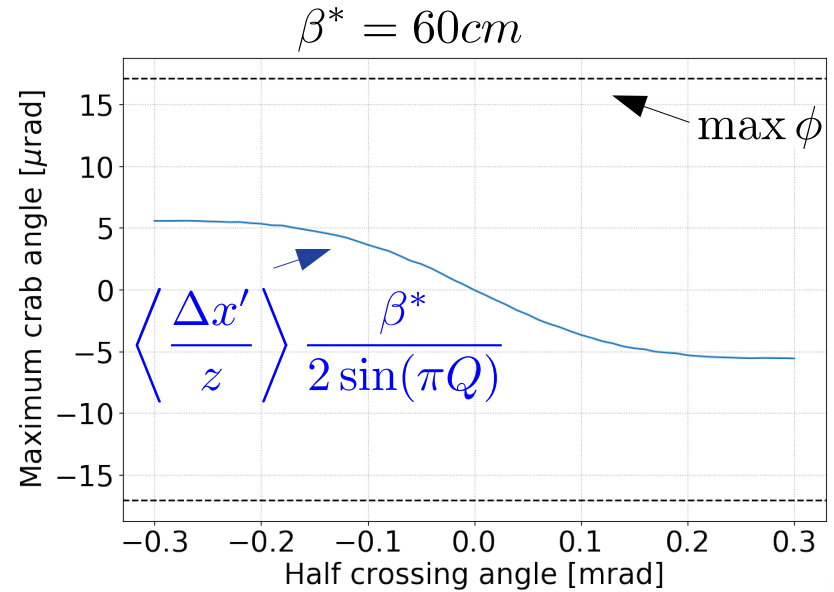
Using Hirata's model

- Due to the non-linearity of the beam-beam force (in all d.o.f.) the transverse kick is not linear with the longitudinal position (~ CC RF curvature)
 - We use an average of the beam-beam kick over a 6D distribution of particle (i.e. coherent kick, but a not self-consistent treatment)*



Dependence on the crossing angle

- The maximum of crab angle is expressed with the β^* such that it can be compared to the CC induced angle
- The analytical estimation is quite pessimistic as expected



➤ A crab angle of maximum 5 μrad can be expected due to the crossing angle

- The contribution from beam-beam is reduced by the CC induced crabbing
- If needed, the contribution from beam-beam and the CC non-closure can be disentangled outside of the IR given that they are out of phase (given the proper hardware, i.e. well placed head-tail monitors)

Dependence on the β^*

- Similar conclusions may be drawn for lower β^* , in spite of the hourglass effect

