

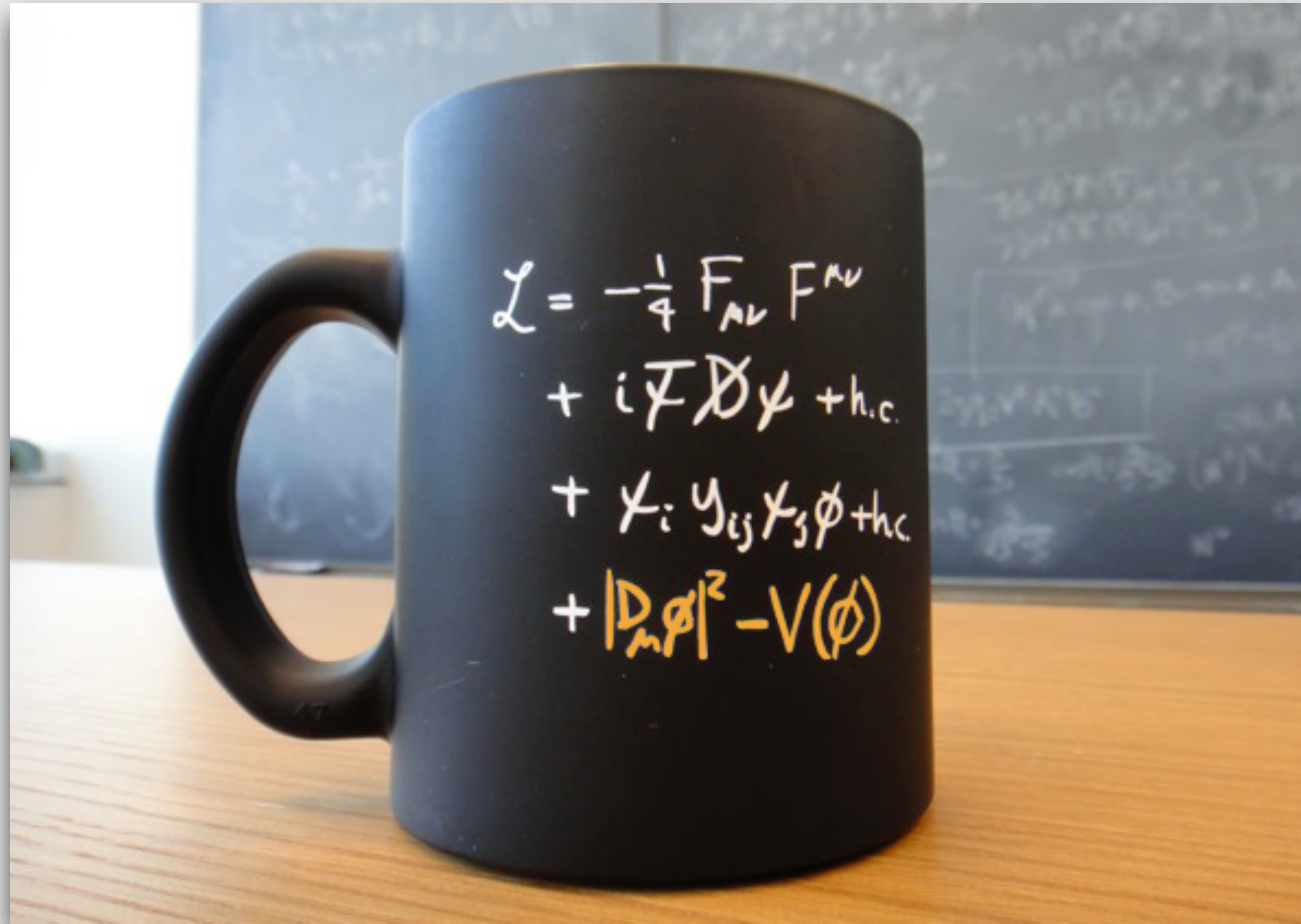
Motivation for high energy lepton colliders

APR. 29 2020

ROBERTO FRANCESCHINI (ROMA 3 UNIVERSITY)

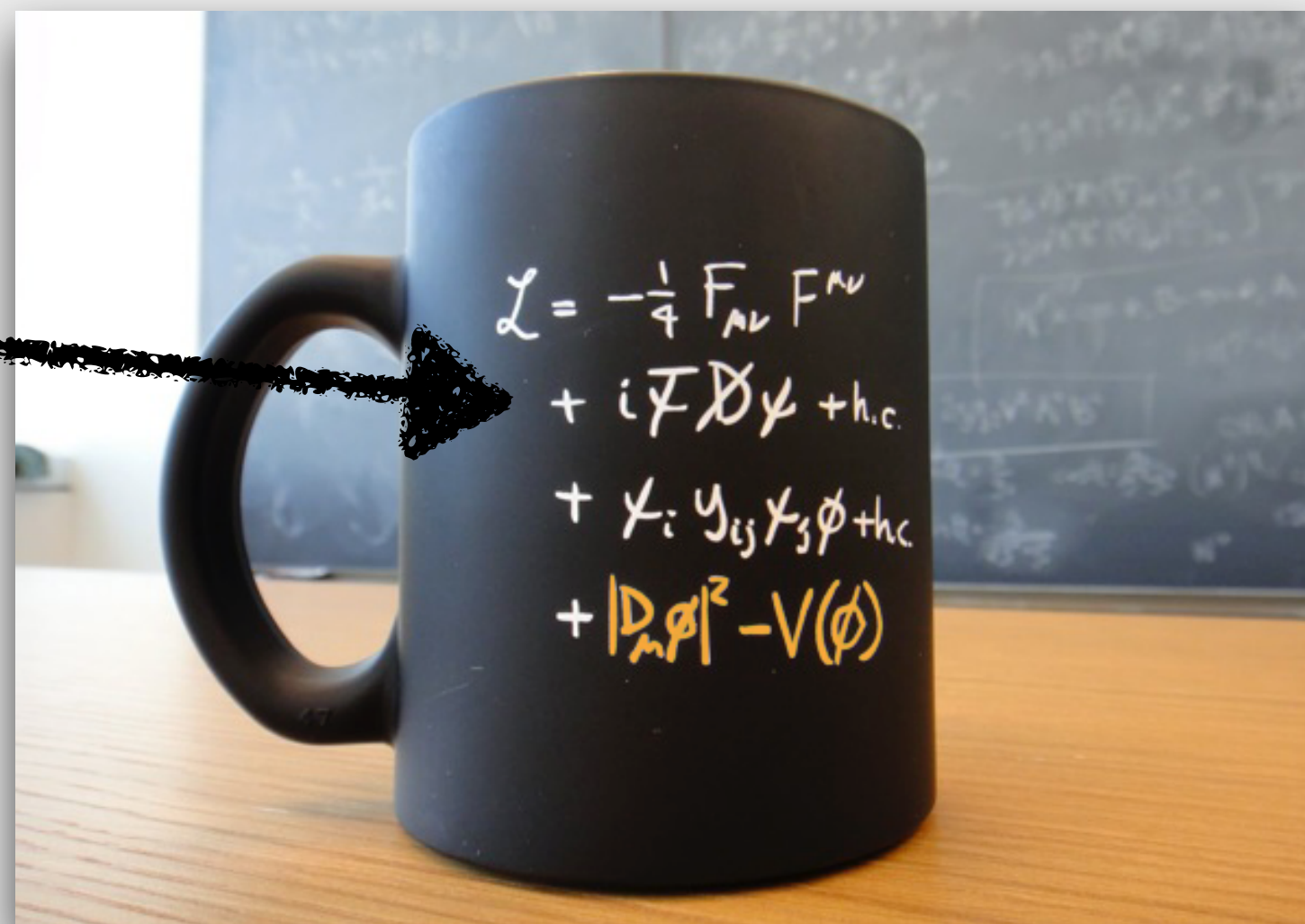
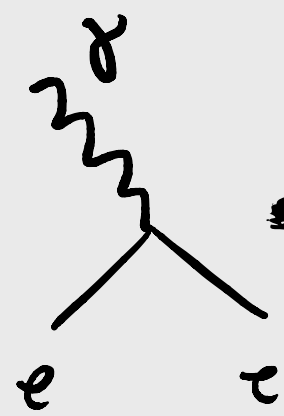


Where do we stand?



**We have got “the” formula
... and it is surprisingly short!**

Where do we stand?



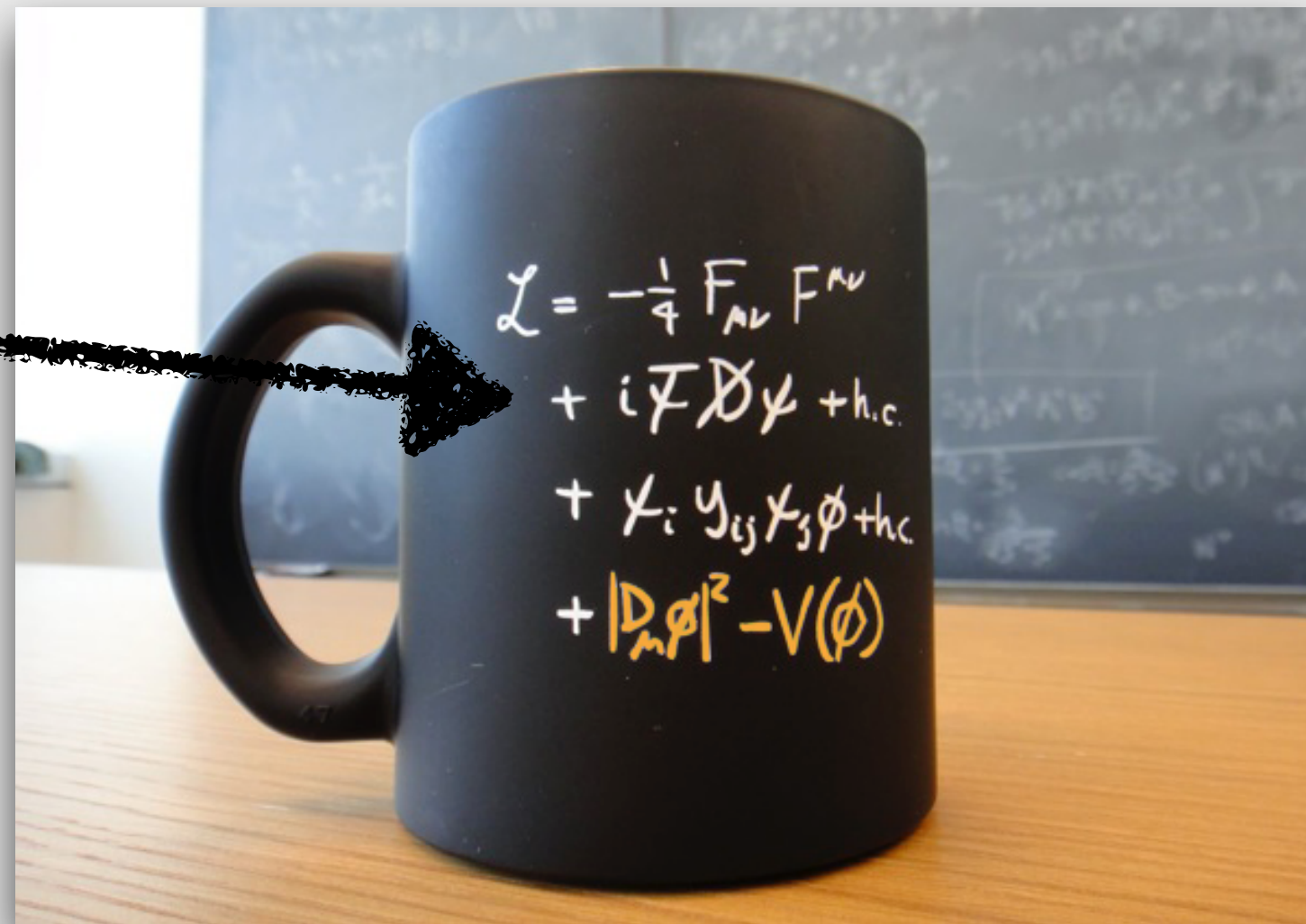
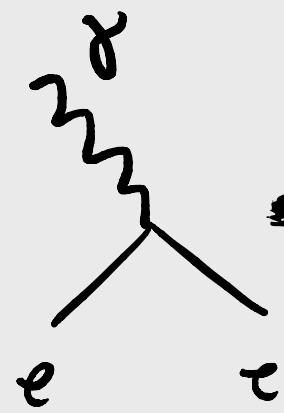
electro-weak interactions

strong interactions

Where do we stand?

SYMMETRY

AS A FUNDAMENTAL CHARACTER OF NATURE



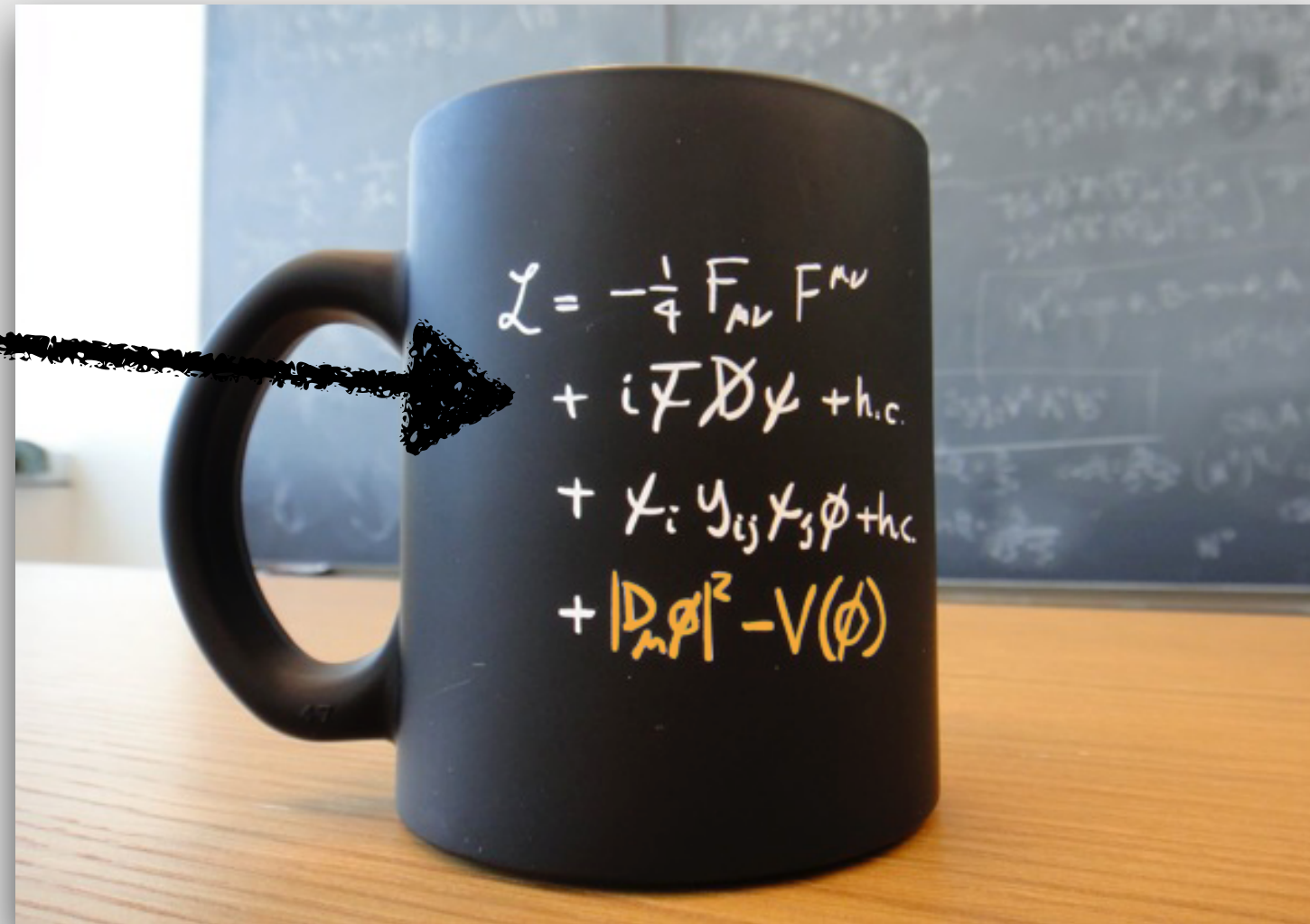
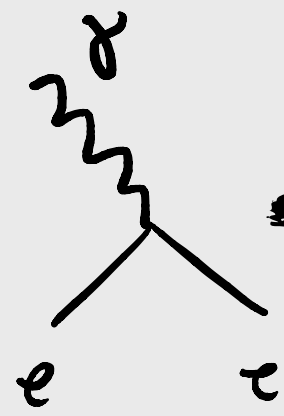
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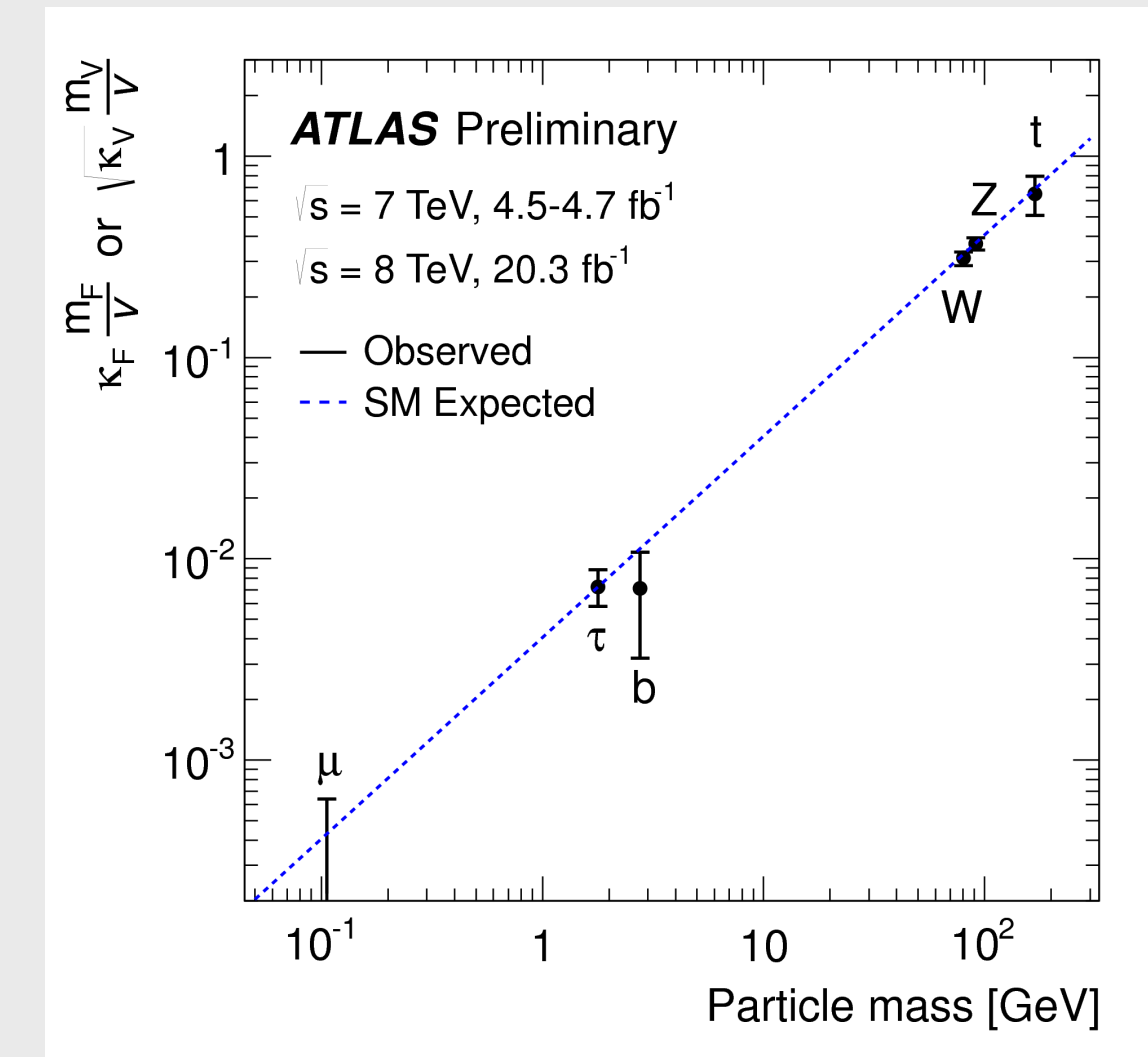
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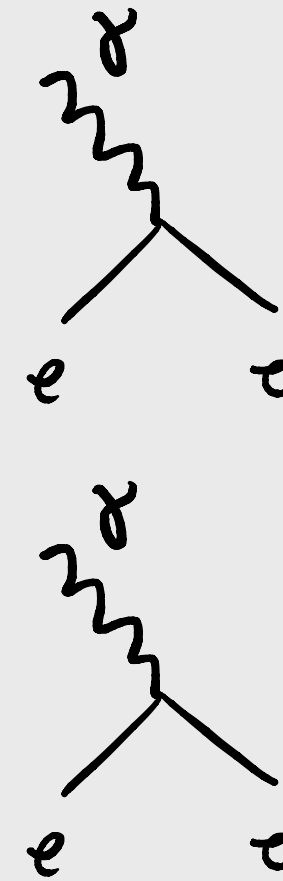
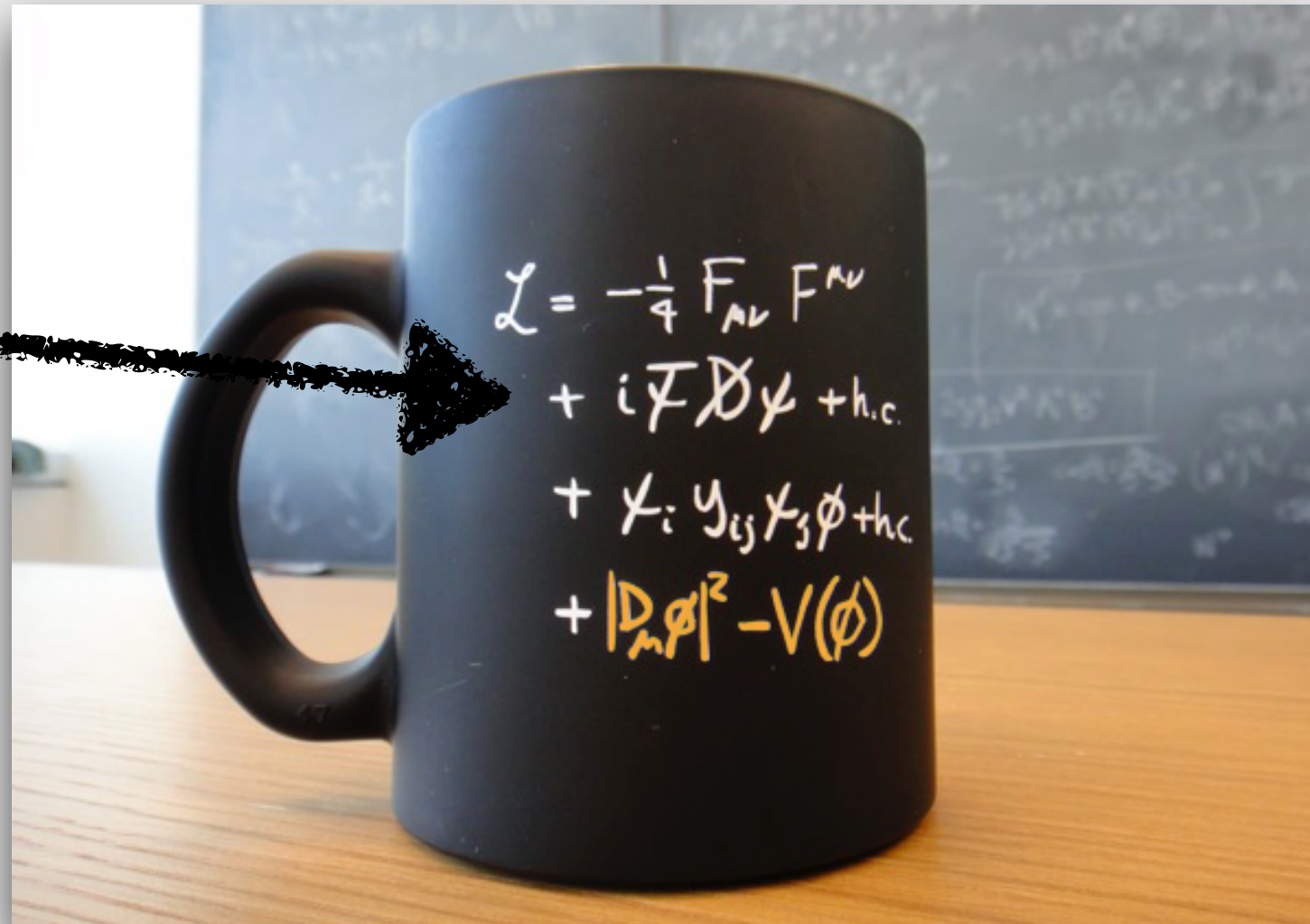
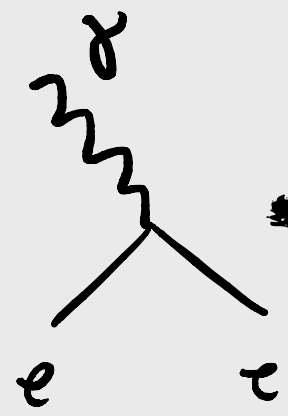


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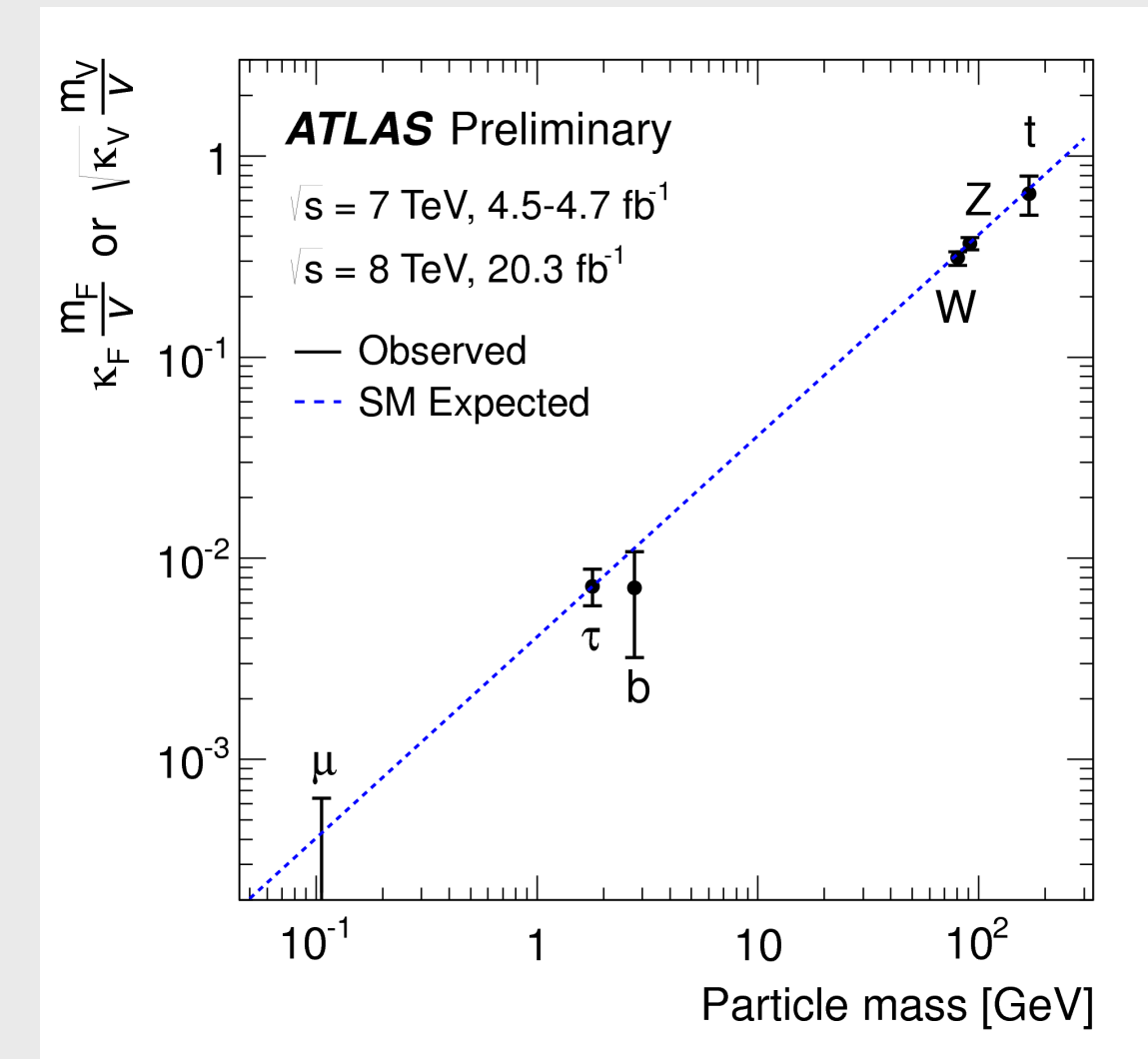
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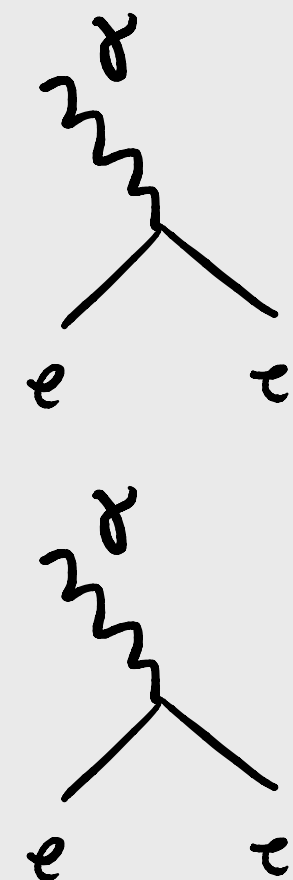
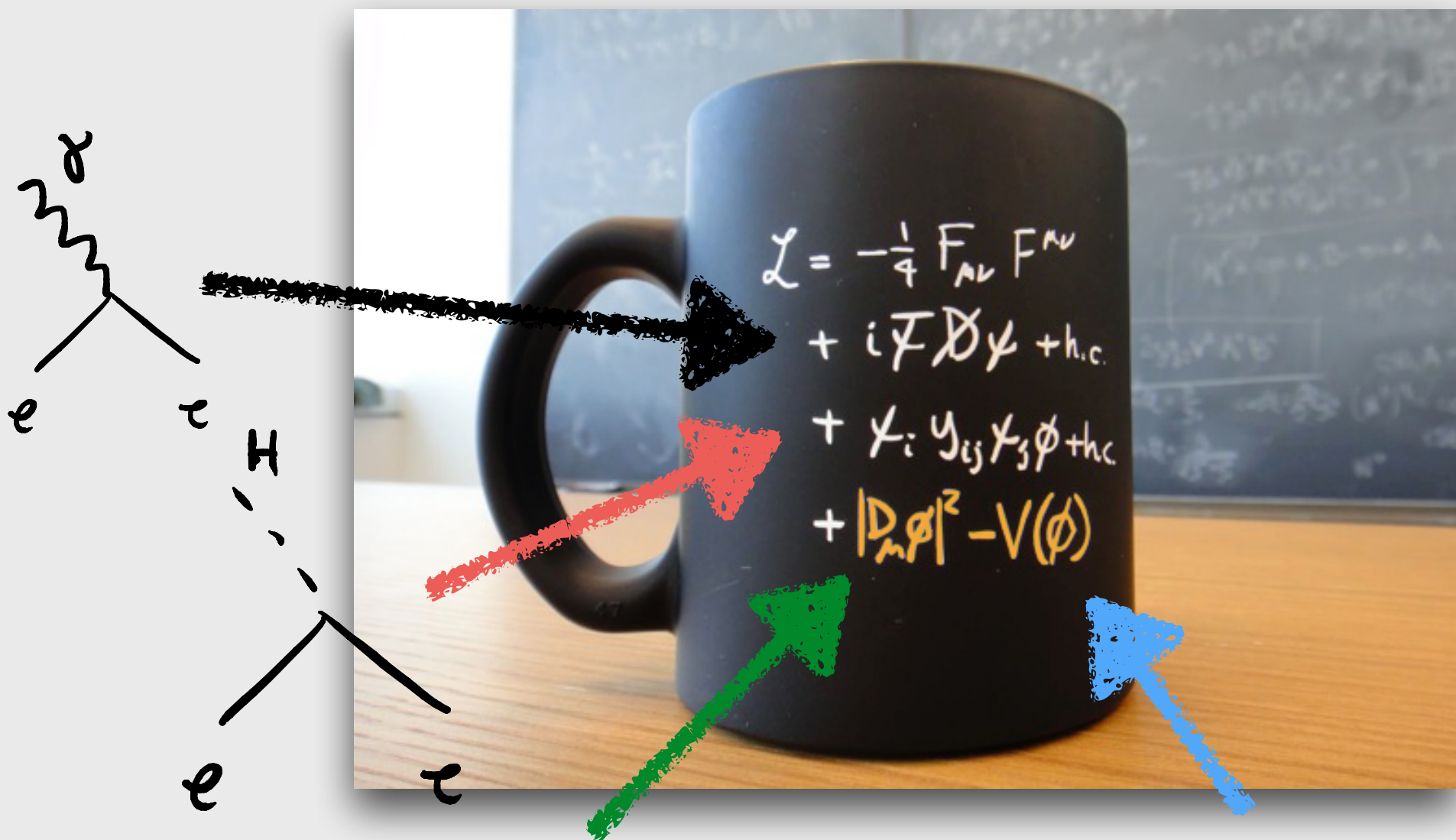


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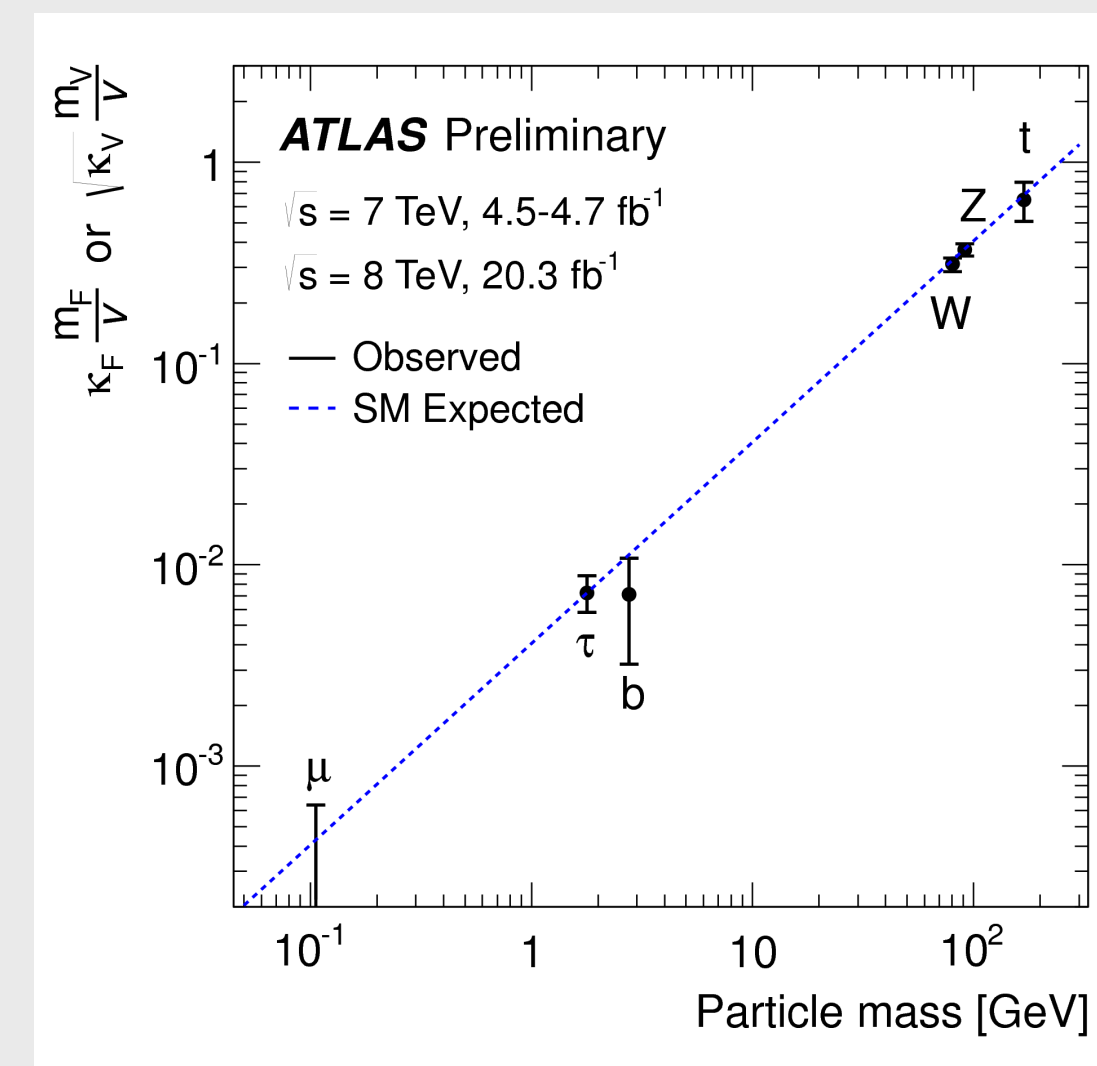
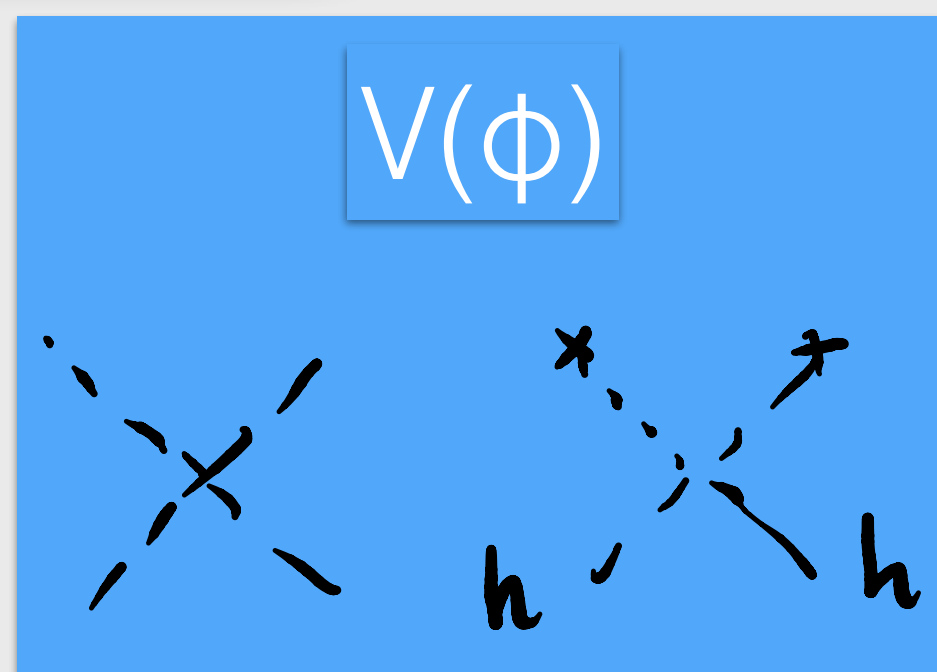
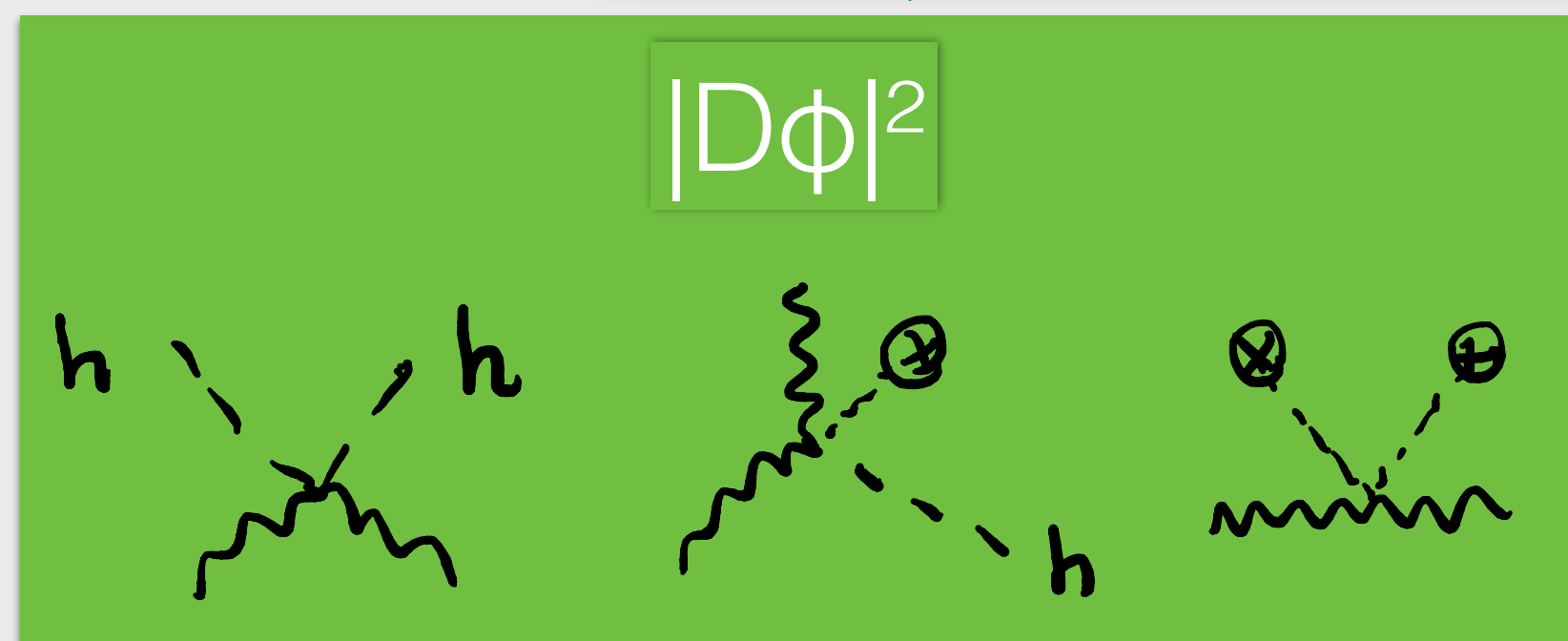
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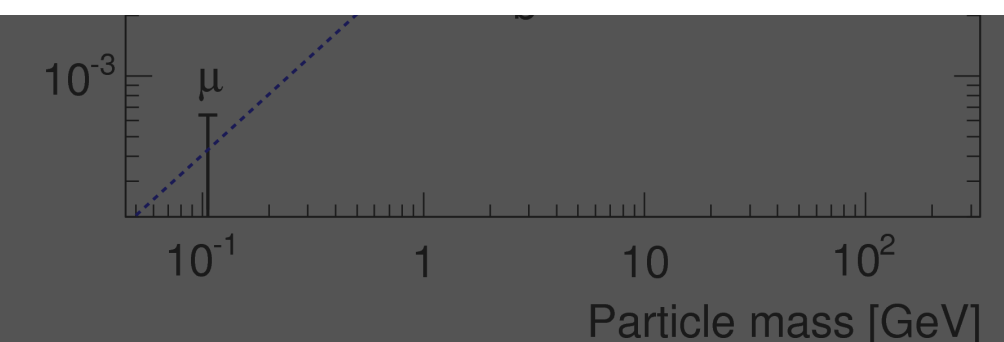
SYMMETRY

AS A FUNDAMENTAL CHARACTER OF NATURE

?????

electro-weak interactions

- We established the principles behind electroweak and strong interaction very well
- We measured the Higgs boson only very “broad brush”
- The Higgs boson may be a whole new thing compared to strong and electroweak interactions



And there is more than “just” the Higgs boson

The Standard Model is:

- Observationally “unfit” (misses Gravity, Dark Matter, ...)
- Symmetry, the very idea at the basis of “the” formula, is challenged by a number of phenomena, which may, at best, be described in this language.

Open Questions on the “big picture” on fundamental physics circa 2020

?

• what is the dark matter in the Universe?



• why QCD does not violate CP?



• how have baryons originated in the early Universe?



• what originates flavor mixing and fermions masses?



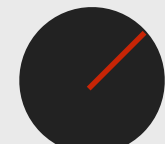
• what gives mass to neutrinos?

EFT



• why gravity and weak interactions are so different?

EFT



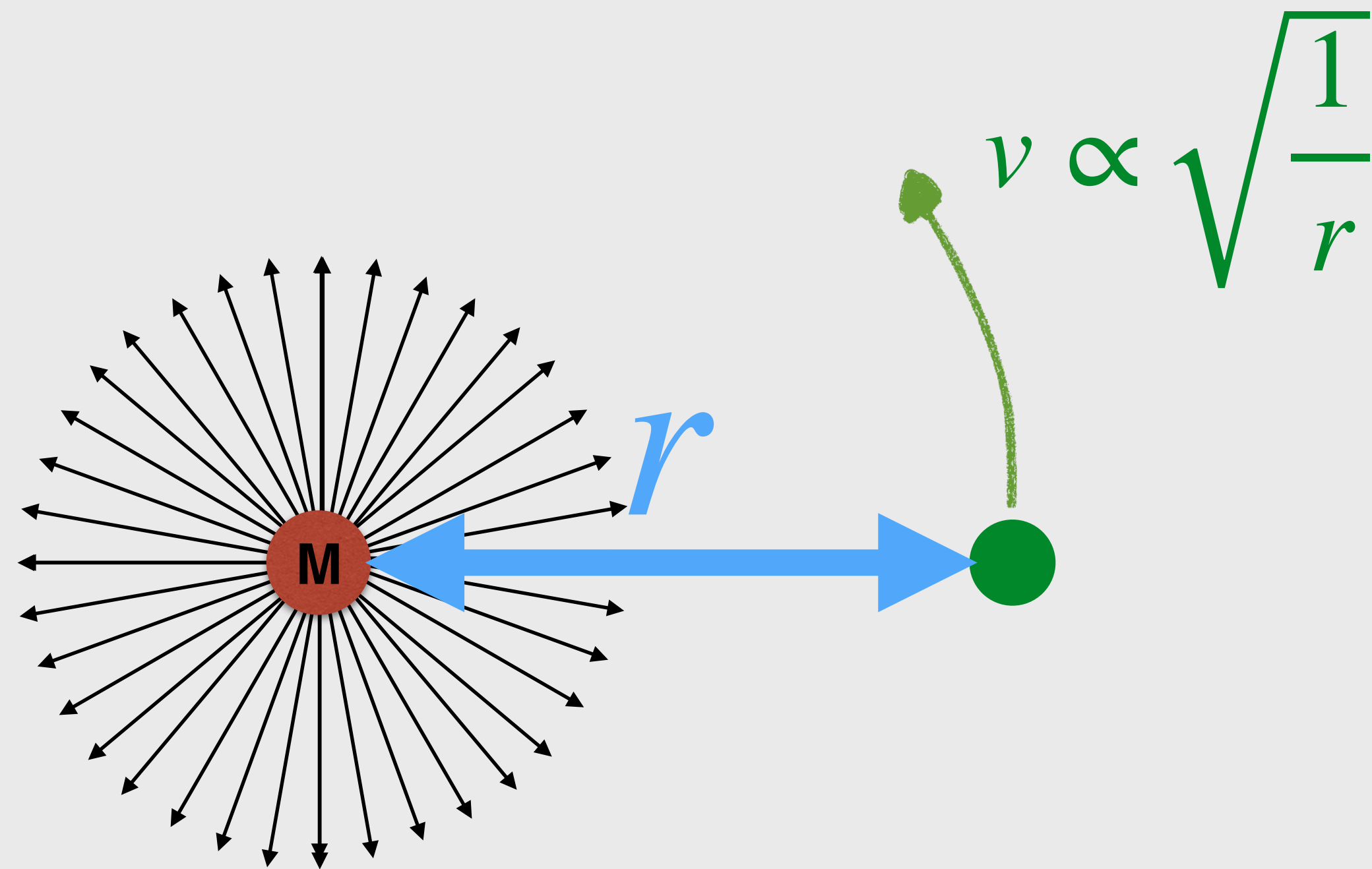
• what fixes the cosmological constant?

EACH of these issues one day will teach us a lesson

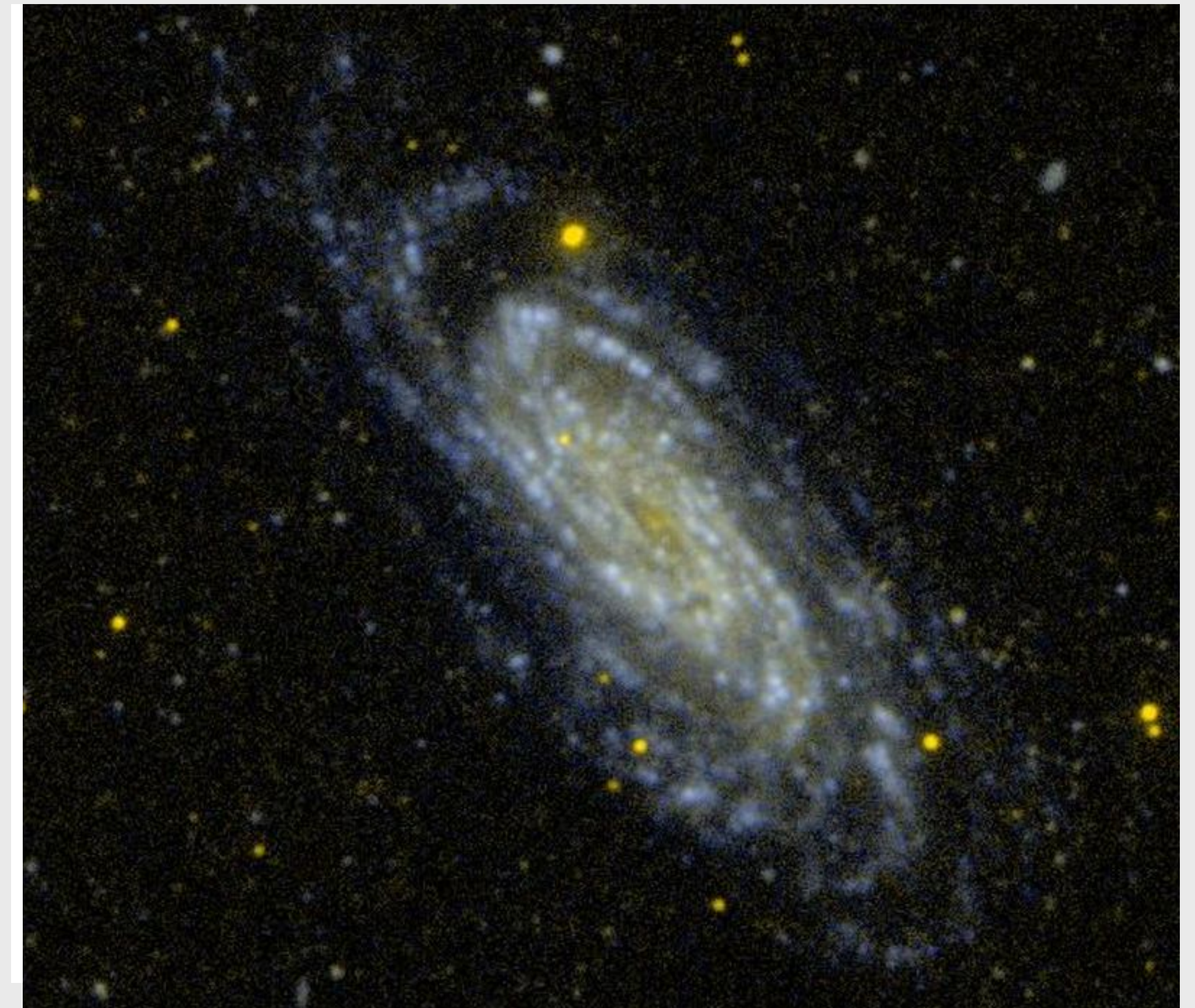
A puzzle we have no idea how to solve

NEWTONIAN

MECHANICS FAILS?



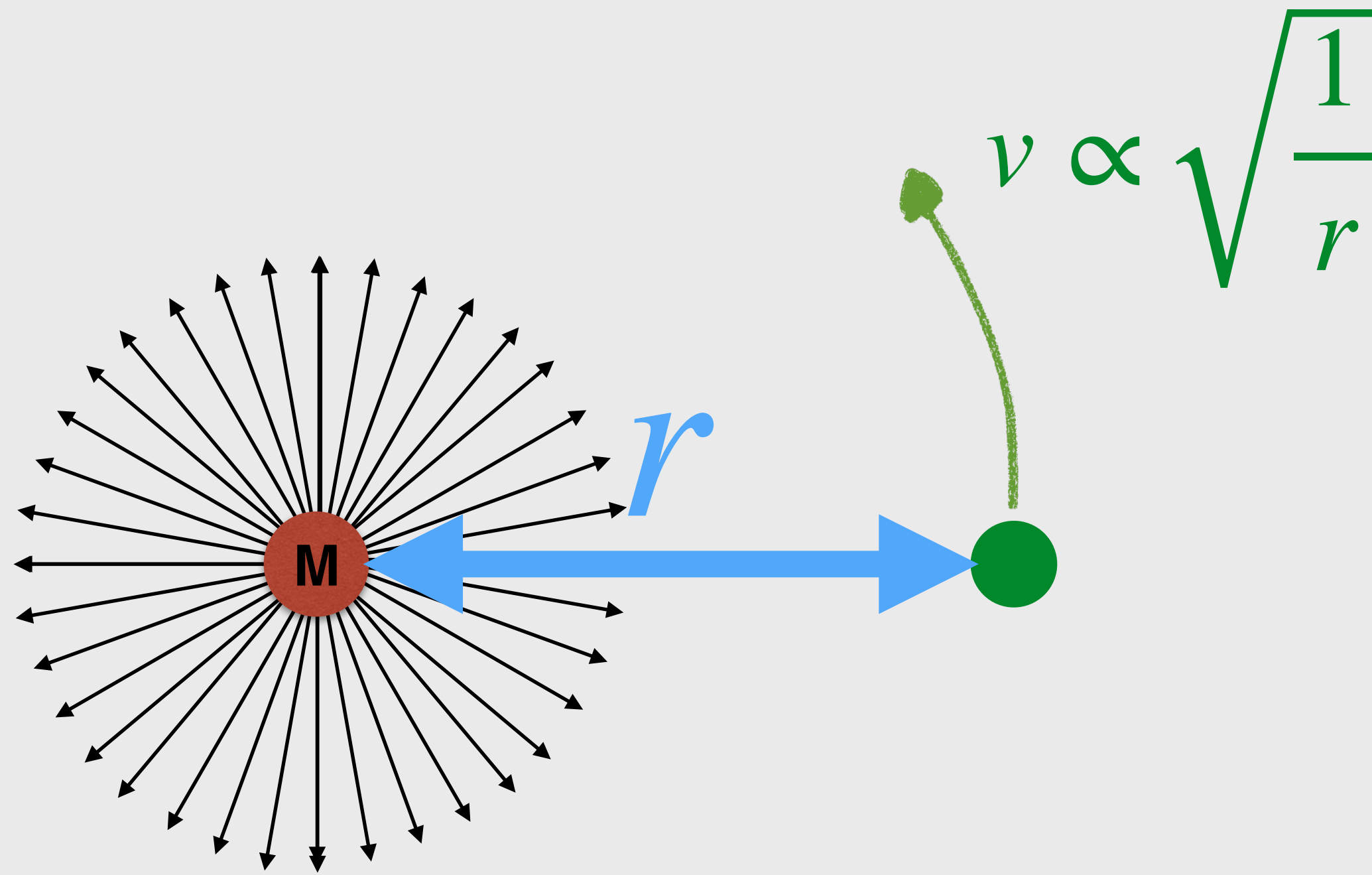
Perfect in our “neighborhood”



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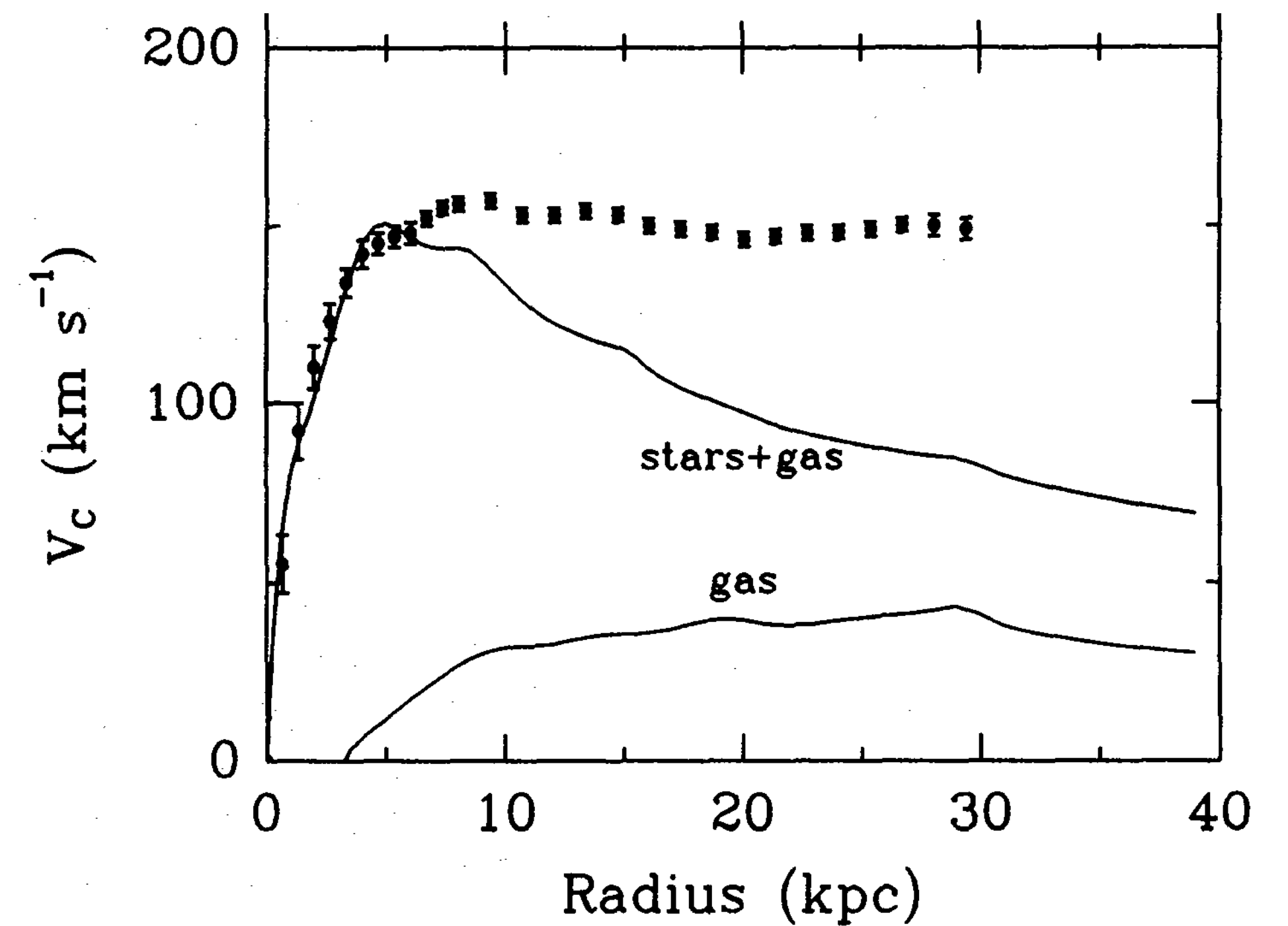
MECHANICS FAILS?



Perfect in our “neighborhood”

Begeman, K. 1989, A&A, 223, 47

NGC 3198



A puzzle we have no idea how to solve

A number of observations (including CMB from early Universe) suggest

a new form of matter must exist

It may well be not of the kind we are used to:

- It may have only weak interactions (even possible it feels only gravity)
- There are candidates “particles” with Compton length $1/M$ ranging from the size of a Galaxy down to High Energy Physics scales (GeV-TeV) and even beyond

It is not necessarily material for particle physics and accelerators

A puzzle we have no idea how to solve

A number of observations (including CMB from early Universe) suggest

- We know the scope of the search for Dark Matter is huge
- In principle, it can be very elusive (to all experiments)
- The simplest history of the early Universe suggests the “TeV” mass range
- Accelerators are the only way to go see it and study it in detail

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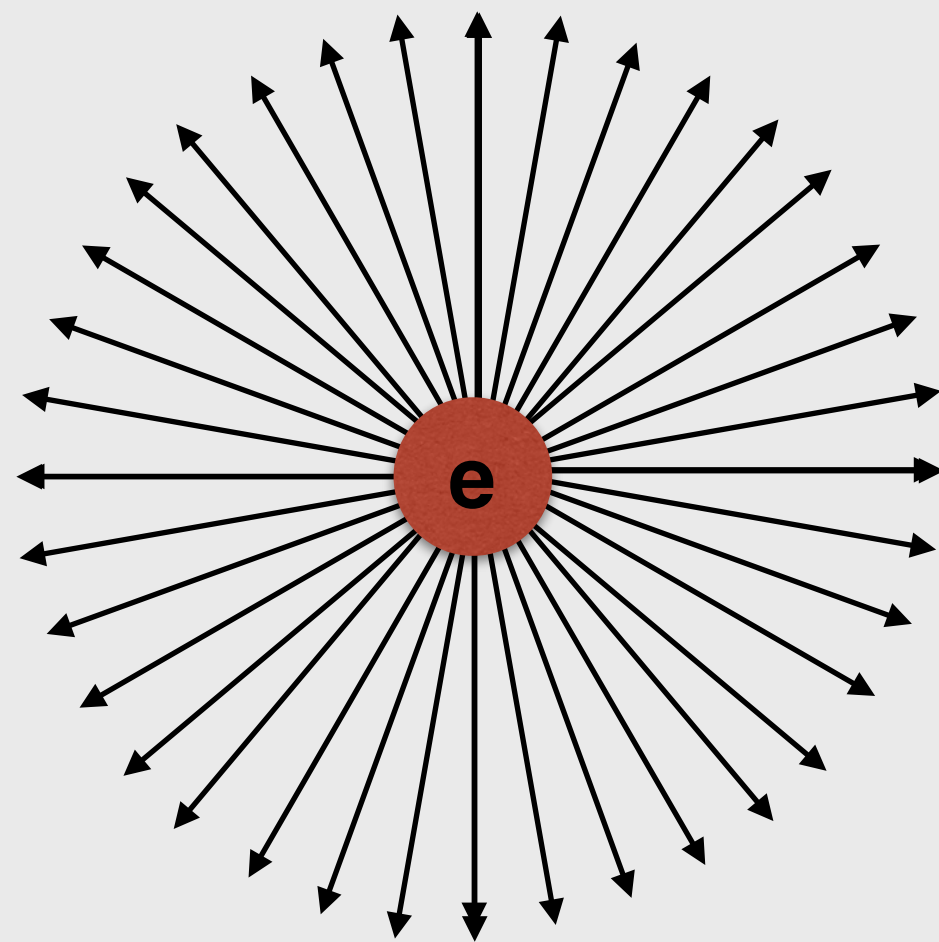
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A puzzle (today) we know how to solve

AFTER

RELATIVITY



$$m_e = m_e^{(0)} + \int_{r_e}^{\infty} \mathcal{E}$$

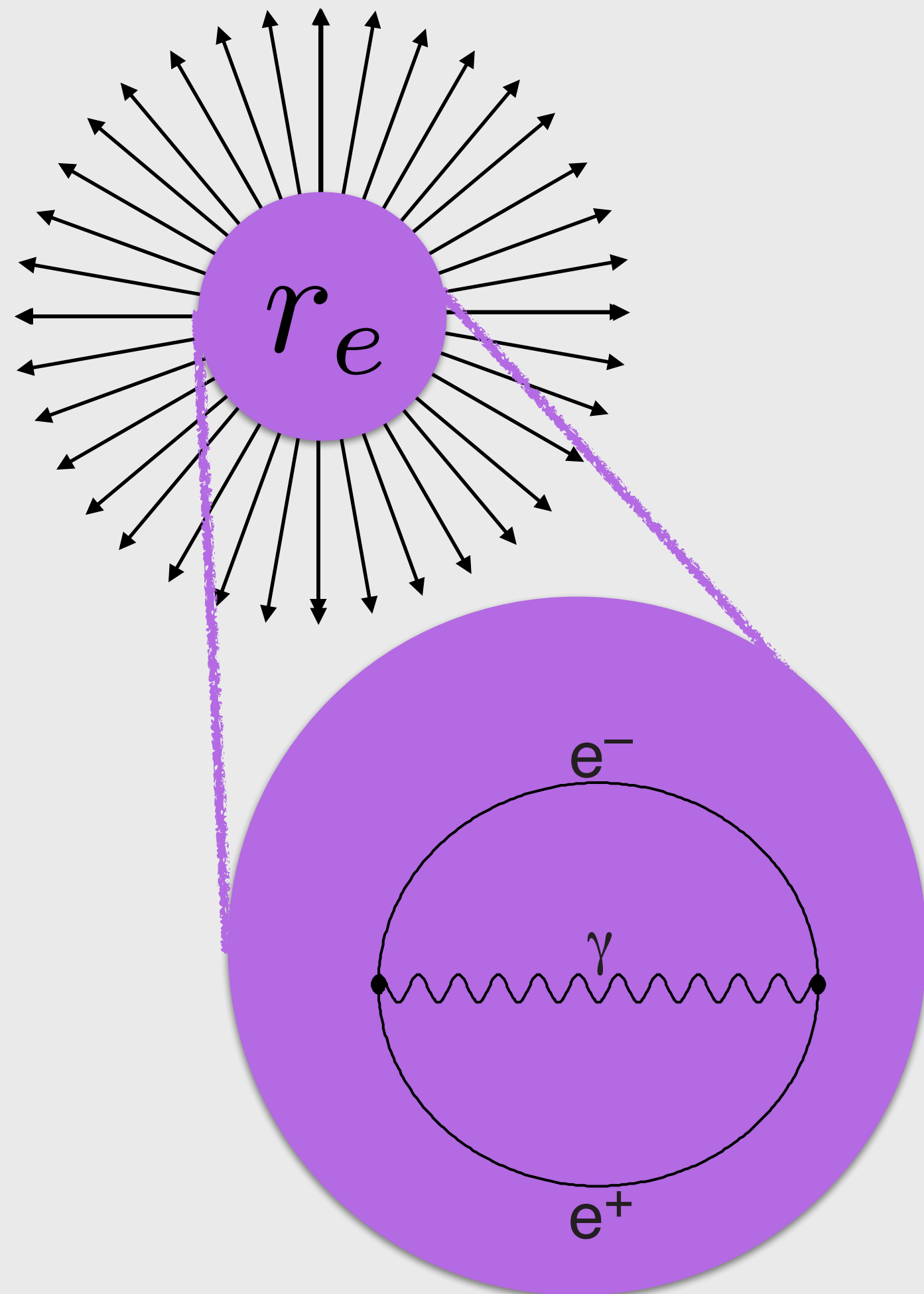
$$\int_{r_e}^{\infty} \mathcal{E} = \frac{1}{4\pi\epsilon_0} \frac{e^2}{r_e}$$

$$\delta m_e \simeq \frac{\alpha_{em}}{r_e} \xrightarrow{r_e \rightarrow 0} \infty$$

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RELATIVITY & QUANTUM MECHANICS



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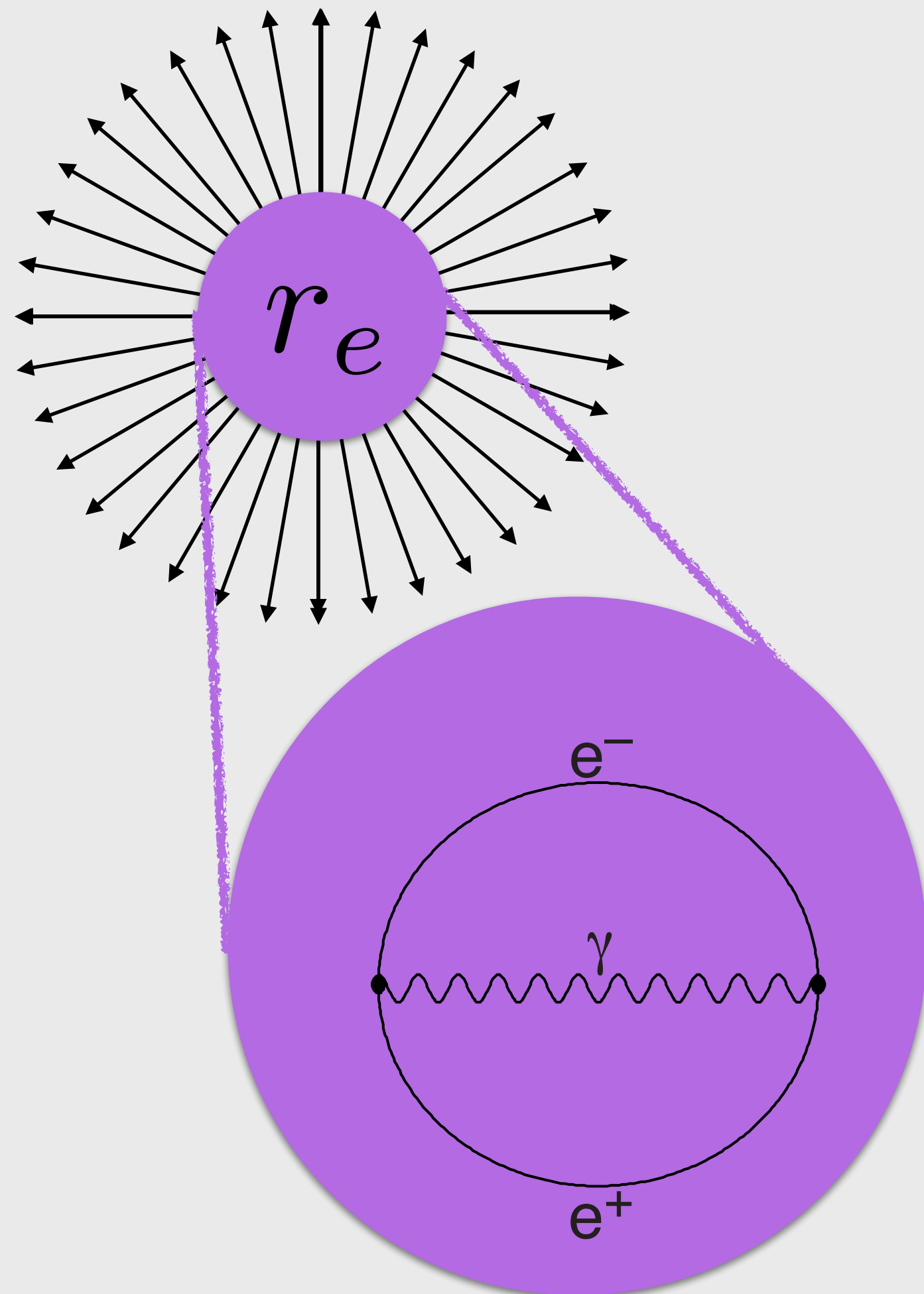
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RELATIVITY & QUANTUM MECHANICS



New symmetry (particle-antiparticle) which brought a new particle: the positron

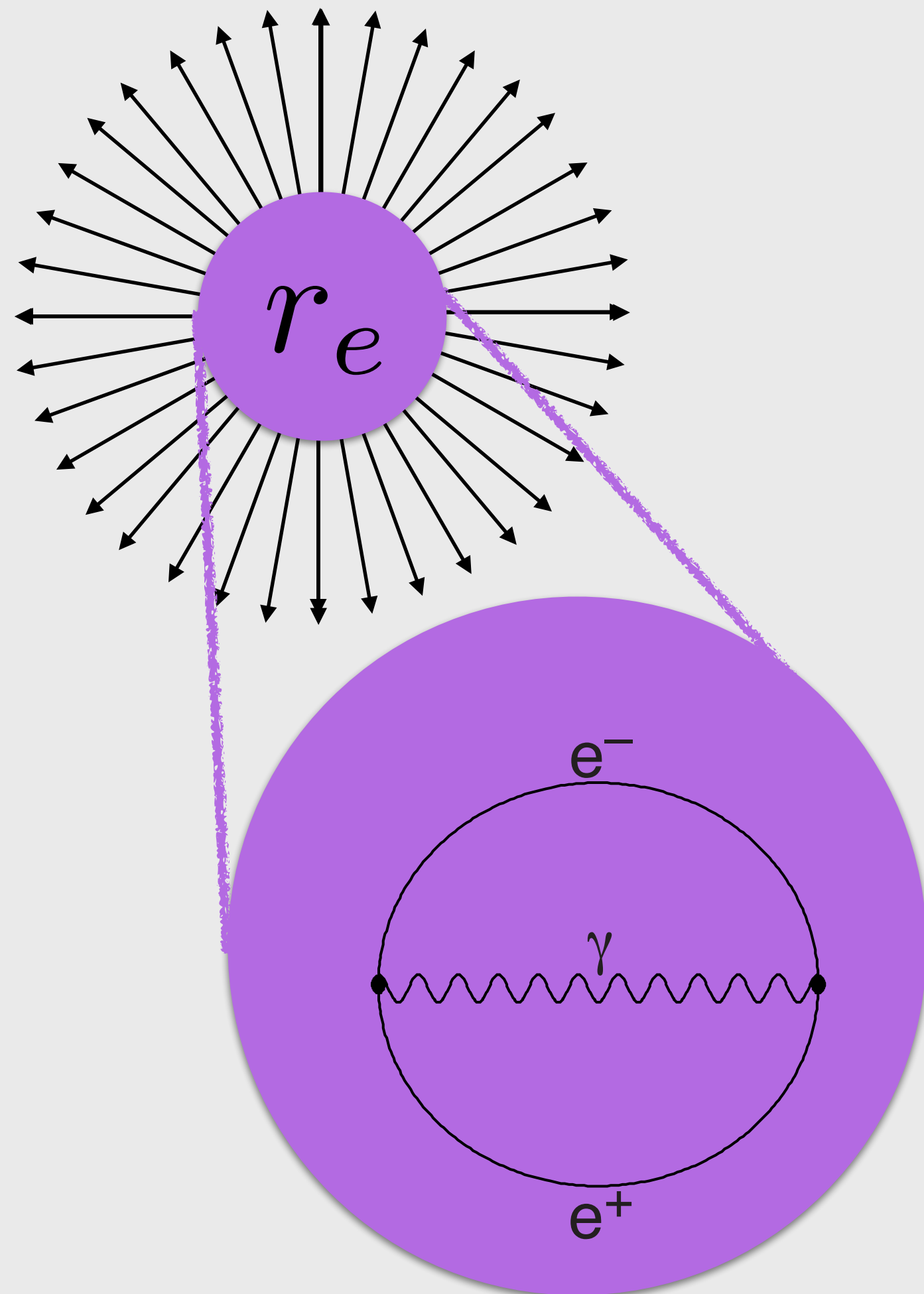
We learned a lesson on physics **at the same mass scale** as where the puzzle arises:

$$m_{\text{positron}} = m_{\text{electron}} \ll m_{\text{electron}} / \alpha_{em}$$

A puzzle (today) we know how to solve

AFTER

RELATIVITY & QUANTUM MECHANICS



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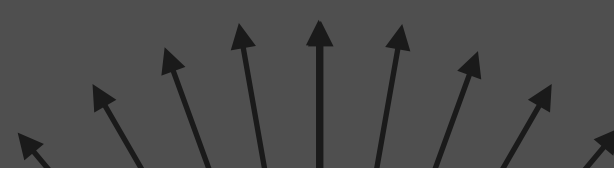
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A puzzle (today) we know how to solve

AFTER

RELATIVITY & QUANTUM MECHANICS

- 
- Similar arguments would require a contribution of the electric field to the mass of the charged pion
 - In that case the solution is not an antiparticle, but a “heavy photon”, the ρ meson, somewhat heavier than the pion
 - In the grand picture, both the positron and the ρ meson appear at the same scale where the problem arises.



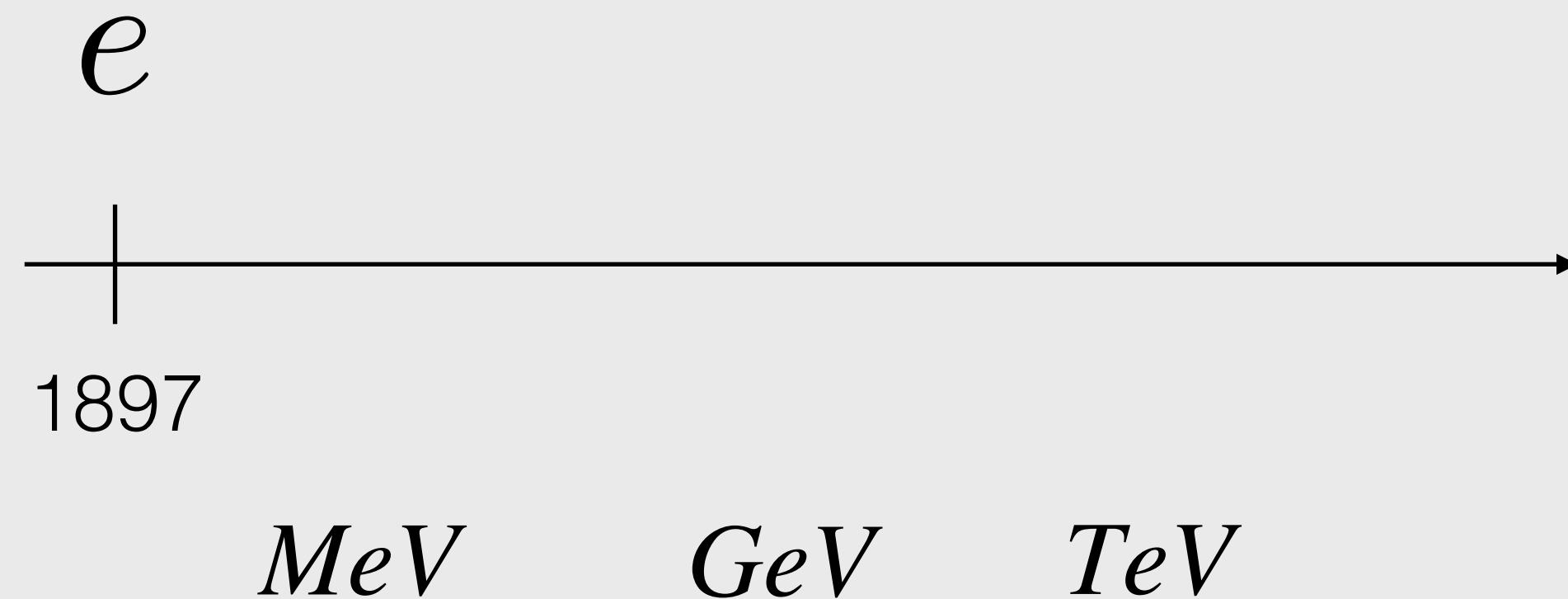
e^+

How we got there

SYMMETRY

AS A FUNDAMENTAL CHARACTER OF NATURE

Symmetries and particles

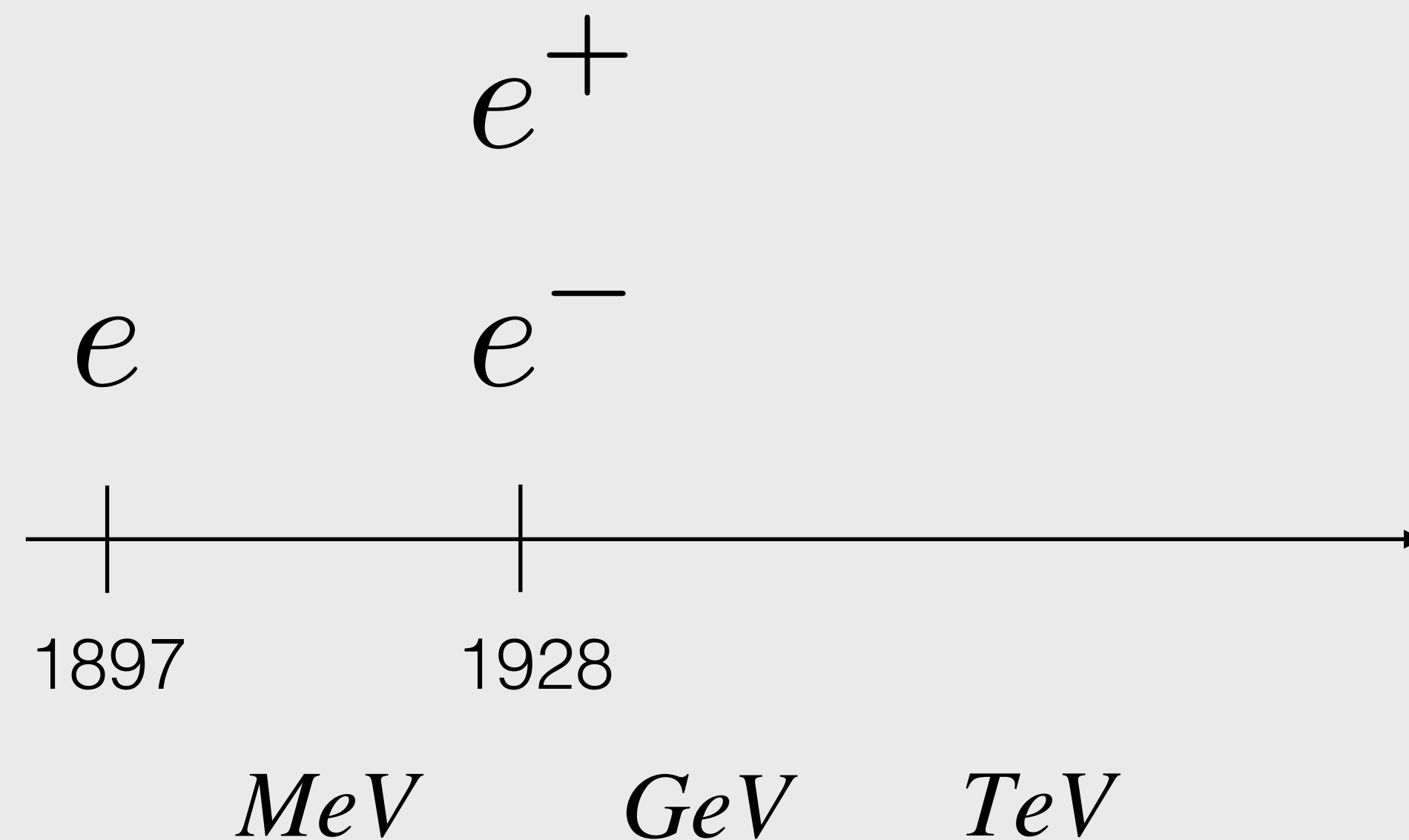


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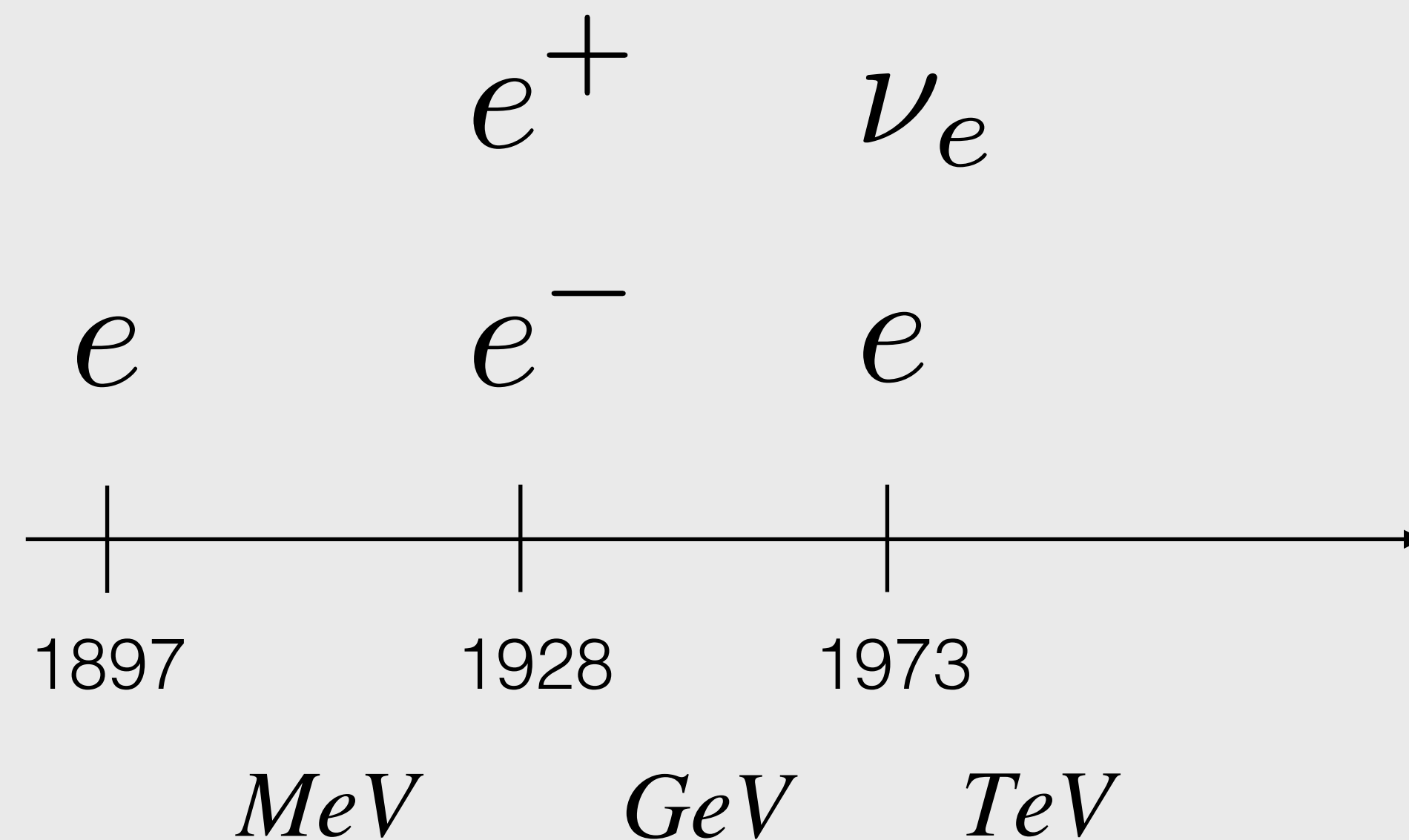


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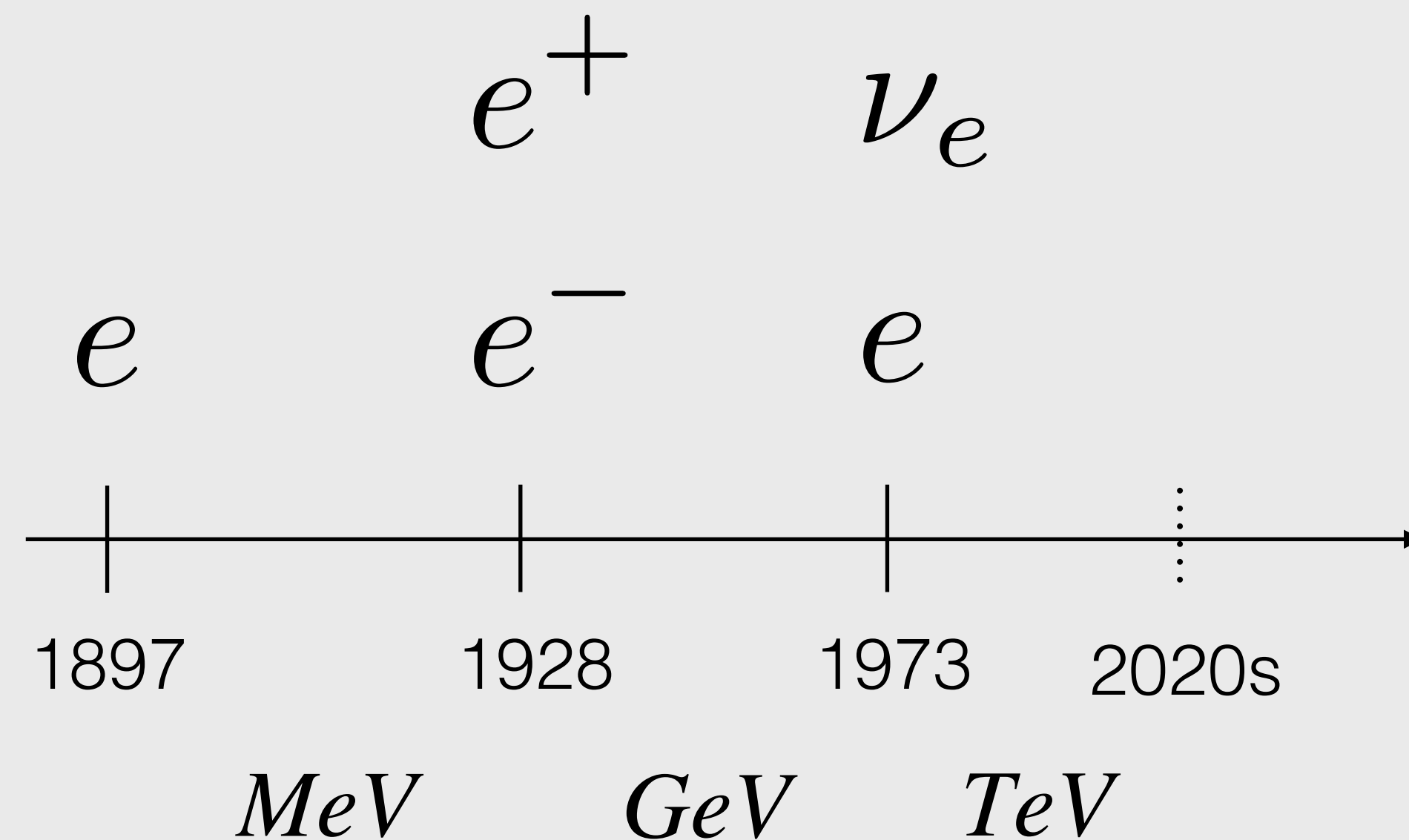


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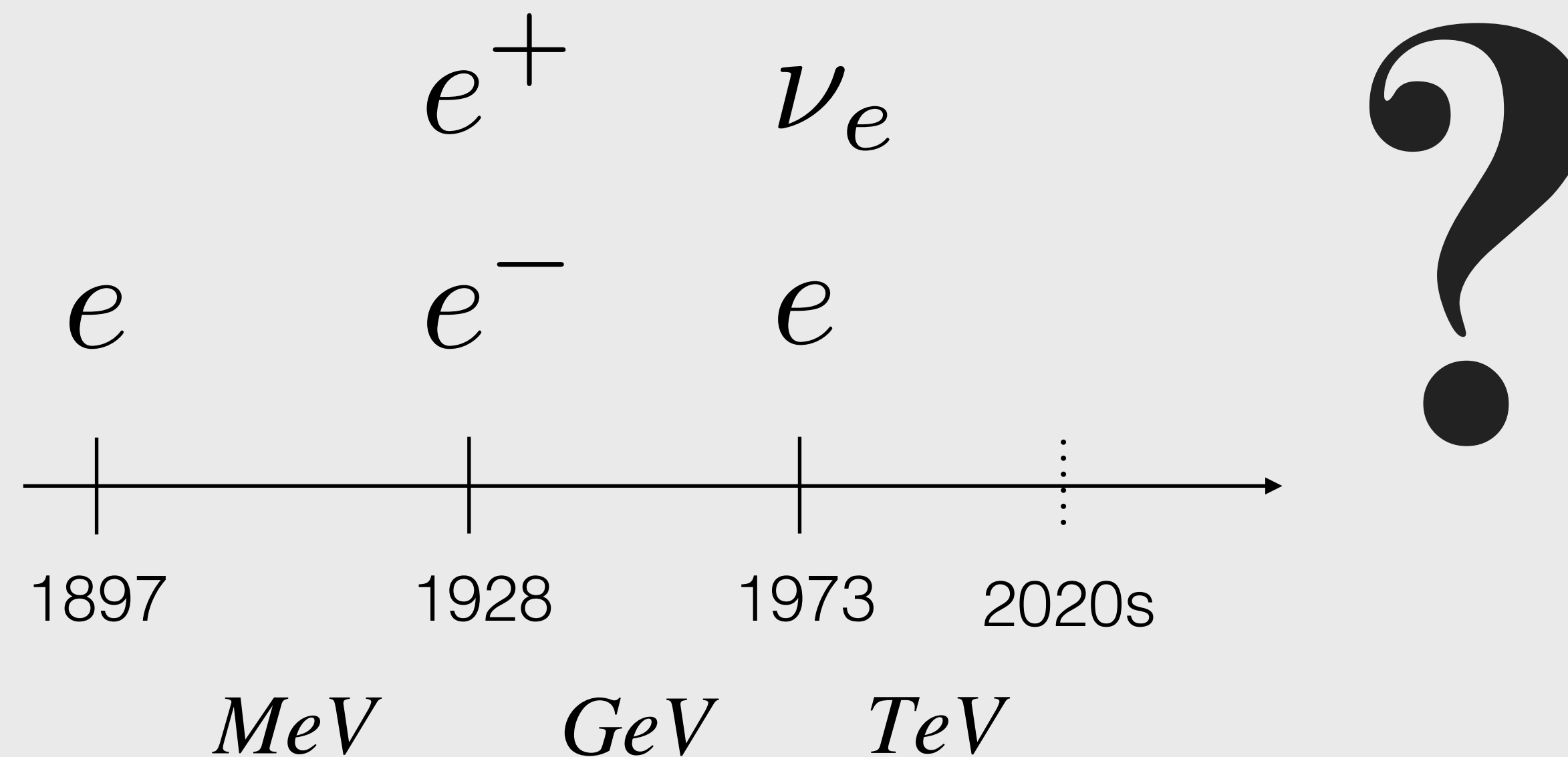


How we got there

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Symmetries and particles



Open Questions on the “big picture” on fundamental physics circa 2020

- what is the dark matter in the Universe?
- why QCD does not violate CP?
- how have baryons originated in the early Universe?
- what originates flavor mixing and fermions masses?
- what gives mass to neutrinos?
- why gravity and weak interactions are so different?
- what fixes the cosmological constant?

EFT

EFT

EACH of these issues one day will teach us a lesson

Open Questions on the “big picture” on fundamental physics circa 2020

WEAK INTERACTIONS

STRONG INTERACTIONS

?

- what is the dark matter in the Universe?



- why QCD does not violate CP?



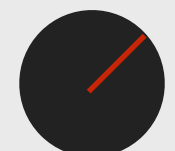
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- what fixes the cosmological constant?

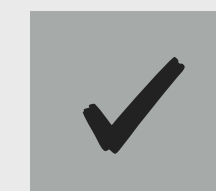
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Open Questions on the “big picture” on fundamental physics circa 2020

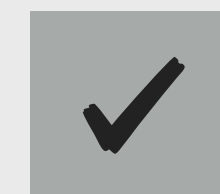
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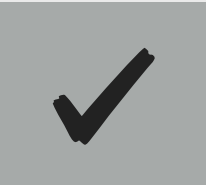
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We need to explore weak interactions

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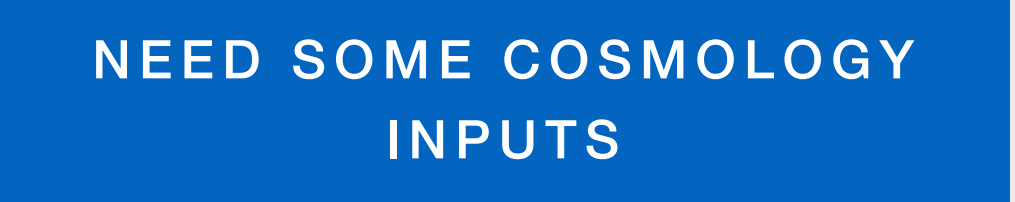










WEAK INTERACTIONS

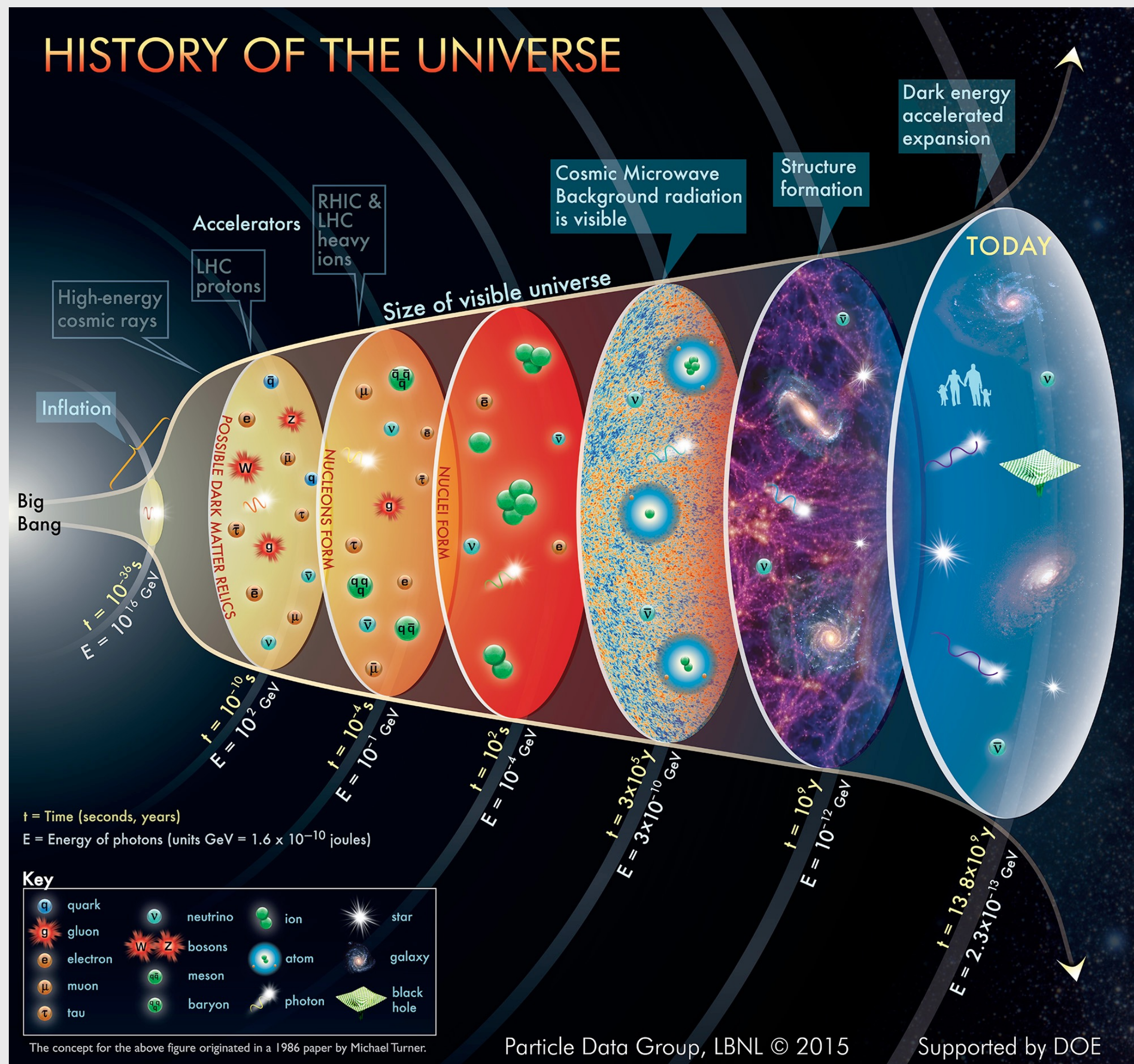
STRONG INTERACTIONS

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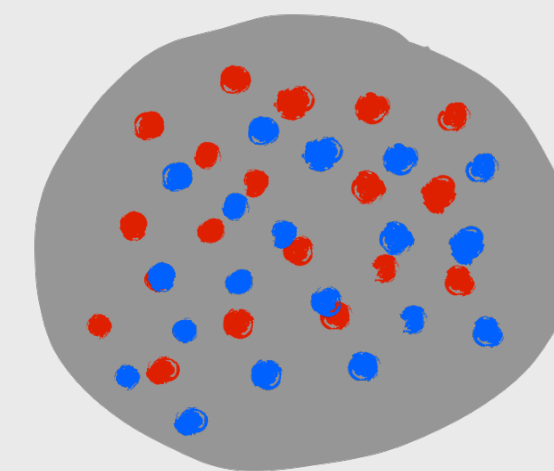
Open Questions on the “big picture” on fundamental physics circa 2020



Nothing we have measured in high energy physics makes so much of a distinction between particles and anti-particles.

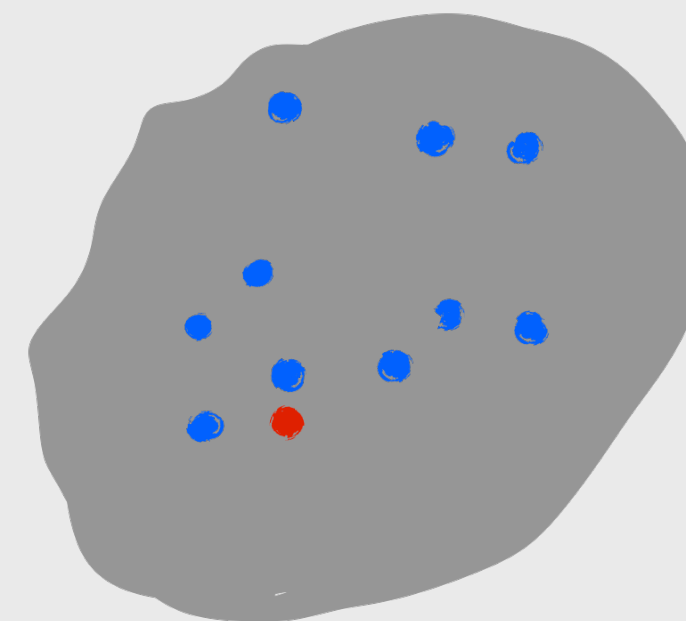
The observable Universe is made of matter, no antimatter

We need to go from this



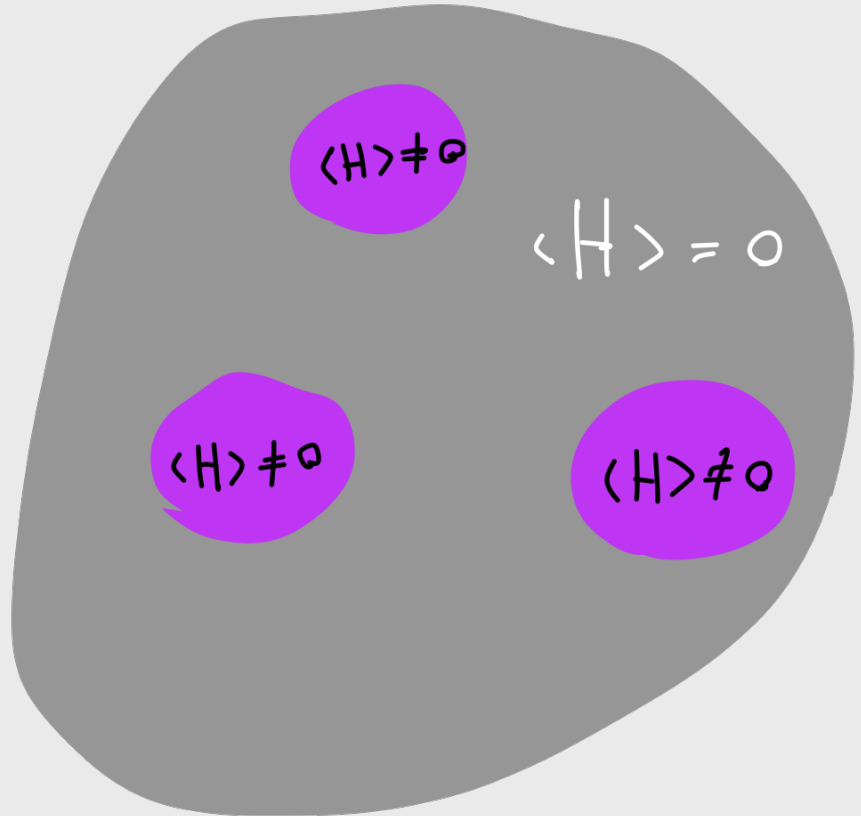
particles
antiparticles

to this

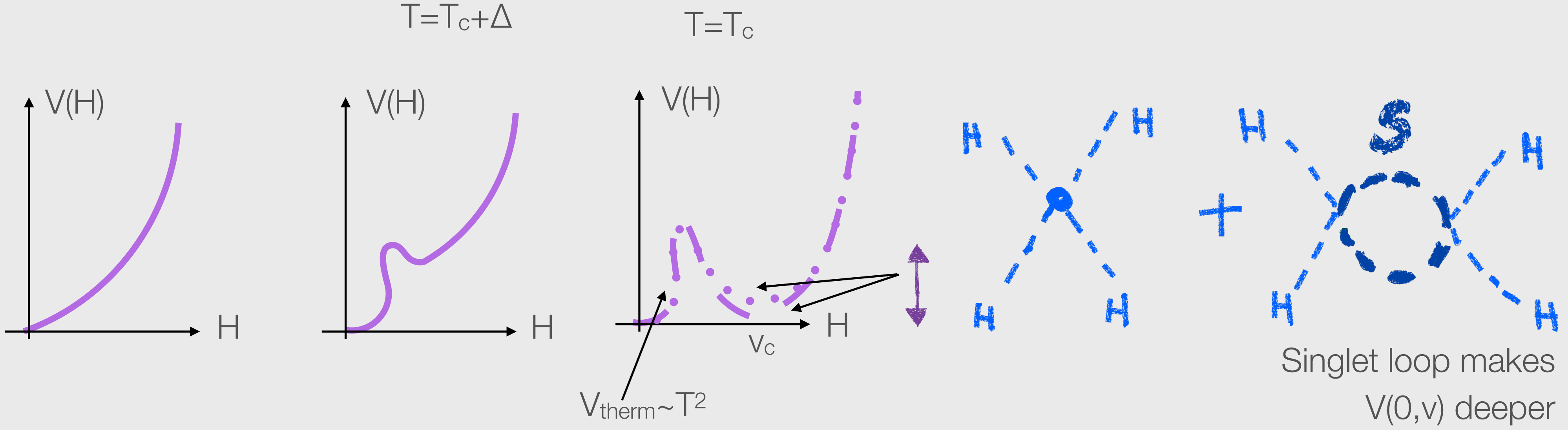


out-of-equilibrium processes are necessary

Electroweak phase transition



- Modifications of the Higgs potential \Rightarrow Out of Equilibrium transition from one vacuum to a new energetically favorable one



Singlet loop makes $V(0, v)$ deeper

Electroweak phase transition

- We need to study all possible new states that induce a change in the Higgs boson potential.
- For these new state to have sizable effects in the early Universe they must be light, around 1 TeV at most.
- All searches for new Higgs bosons (or general electroweak particles) probe such fundamental issue of the origin of matter in the early Universe!

$$V_{\text{therm}} \sim T^2$$

$V(0,v)$ deeper

flashing concrete results for

EW phase transition

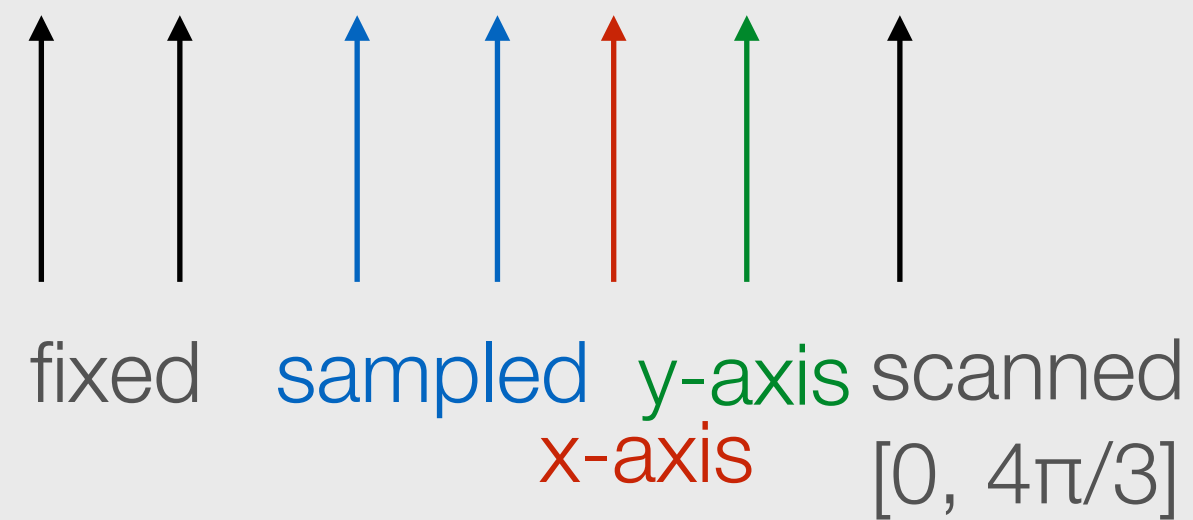
Mixed Singlet for EW phase transition

EW PHASE TRANSITION

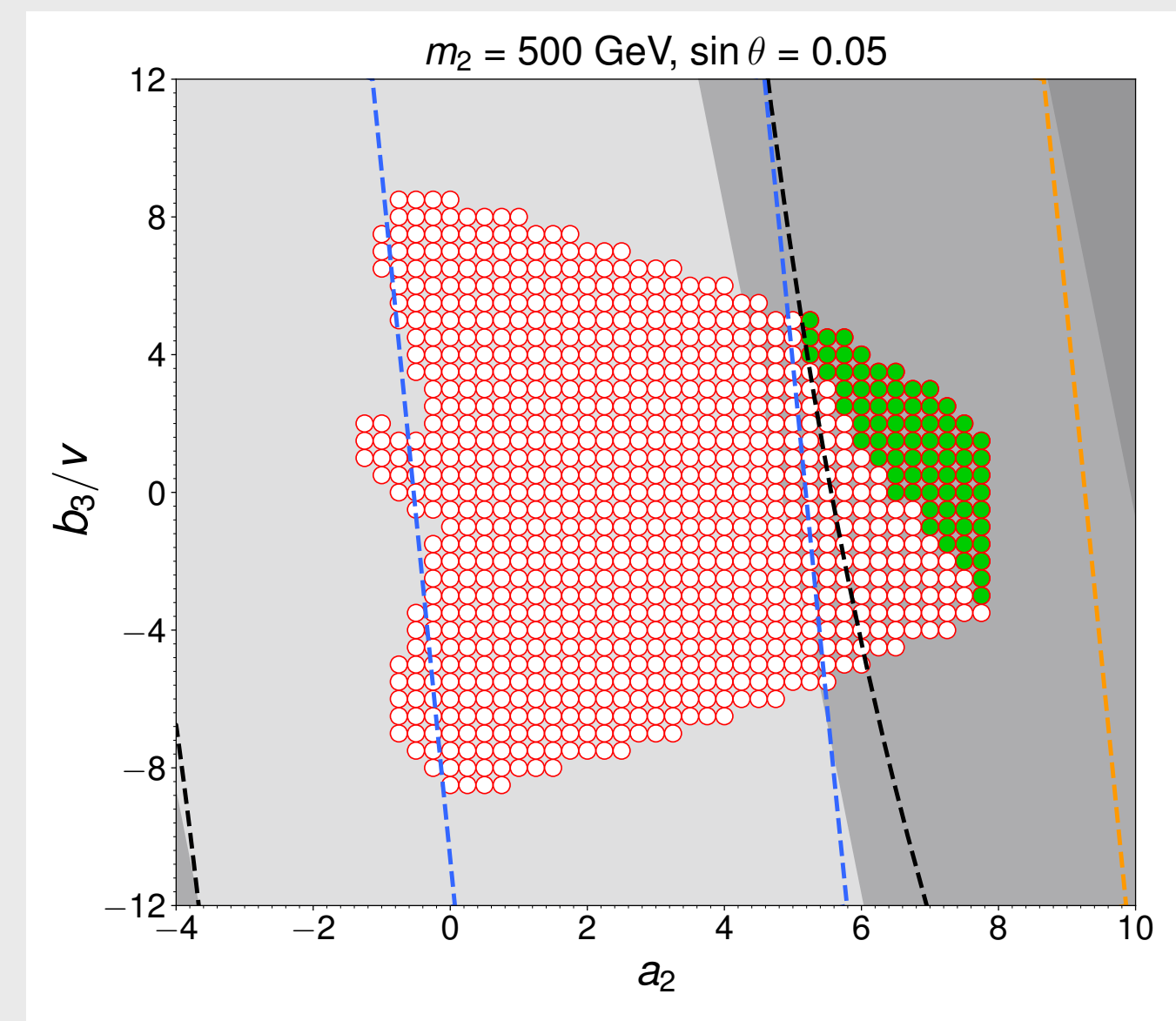
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$$V(\Phi, S) = -\mu^2 (\Phi^\dagger \Phi) + \lambda (\Phi^\dagger \Phi)^2 + \frac{a_1}{2} (\Phi^\dagger \Phi) S + \frac{a_2}{2} (\Phi^\dagger \Phi) S^2 + b_1 S + \frac{b_2}{2} S^2 + \frac{b_3}{3} S^3 + \frac{b_4}{4} S^4.$$

independent parameters
 $\{v, m_1, m_2, \theta, a_2, b_3, b_4\}$.



- “healty” potential (no runaway, minimum $v=246$ GeV, perturbative)
- 1st order phase transition
- HL-LHC sensitivity (from $pp \rightarrow S \rightarrow ZZ$)
- CLIC380/3TeV Single Higgs couplings
- CLIC 1.4 TeV 3 TeV WBF $S \rightarrow hh \rightarrow 4b$
- CLIC hhh 20% @ 95% CL coupling measurement



parameters space of 1st order phase transition accessible by several probes

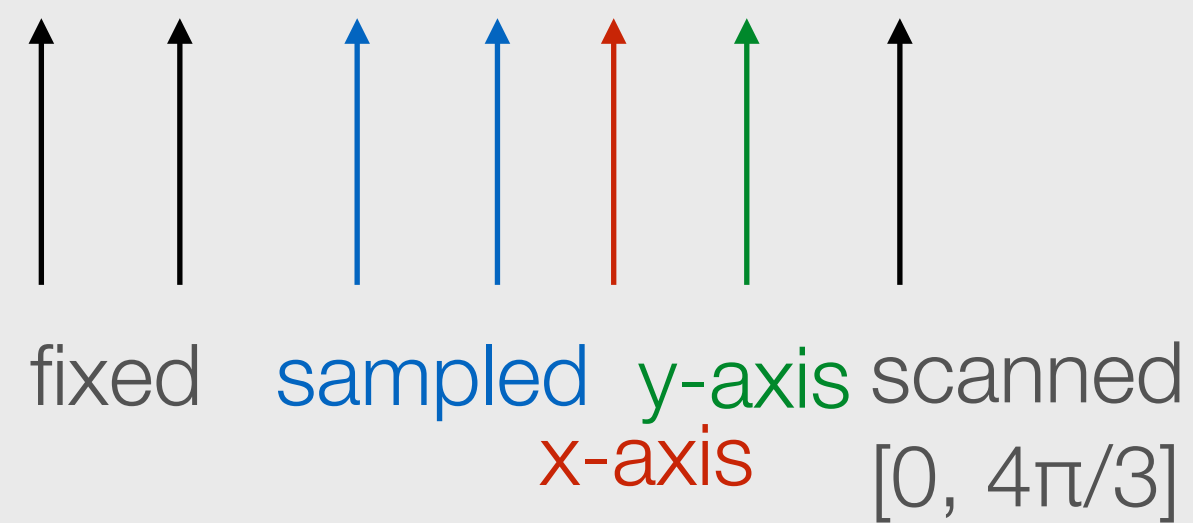
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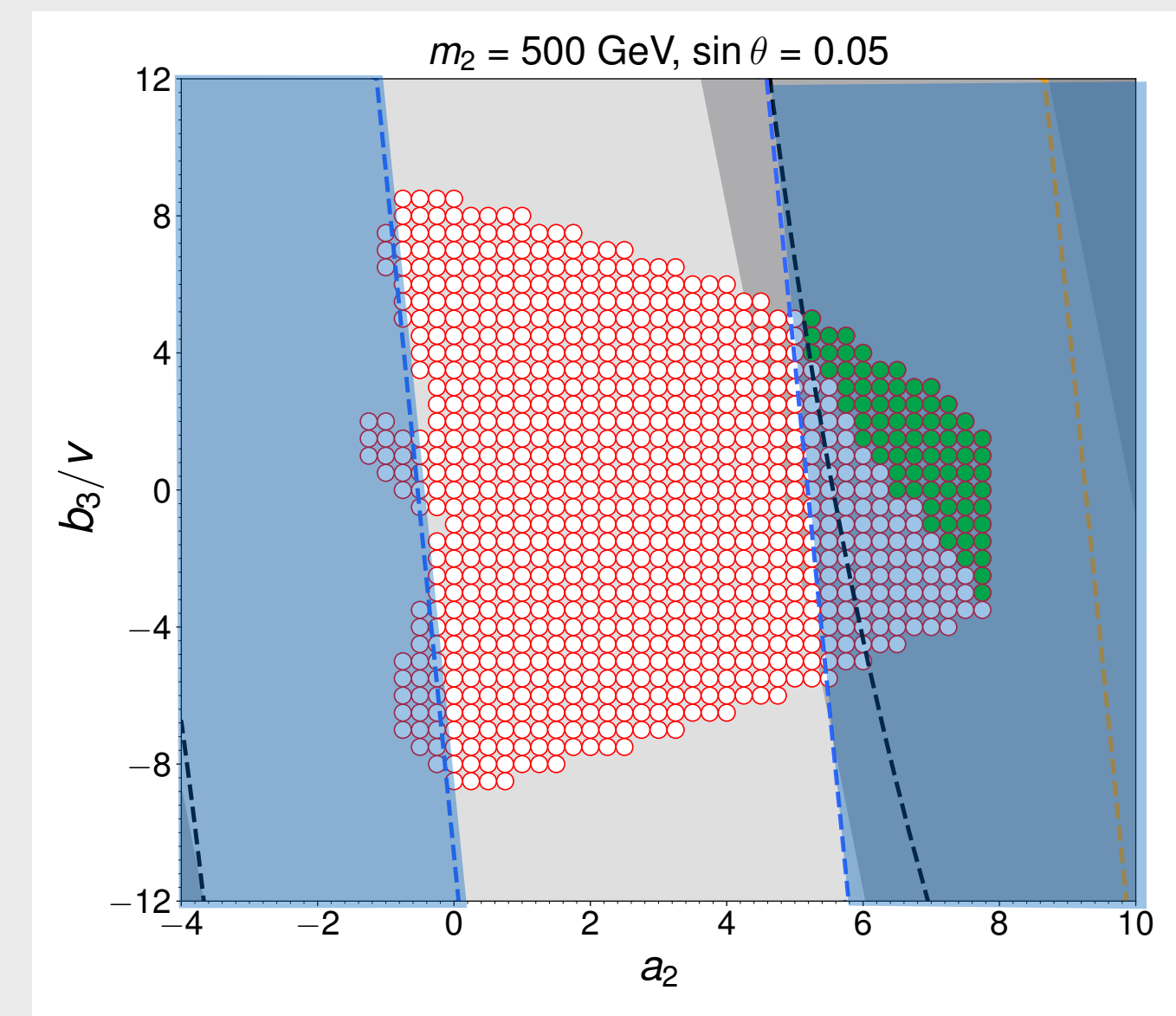
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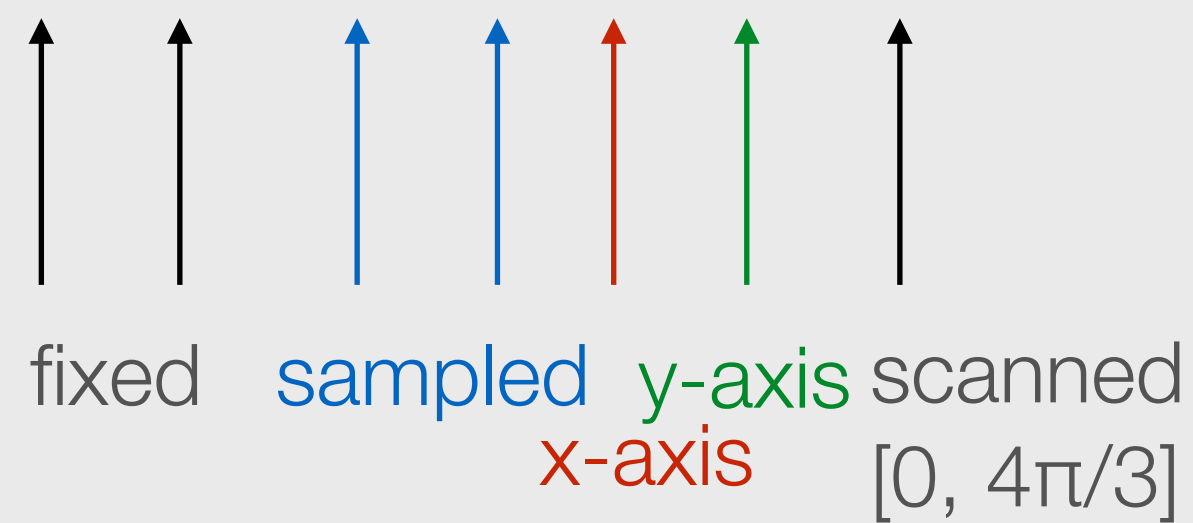
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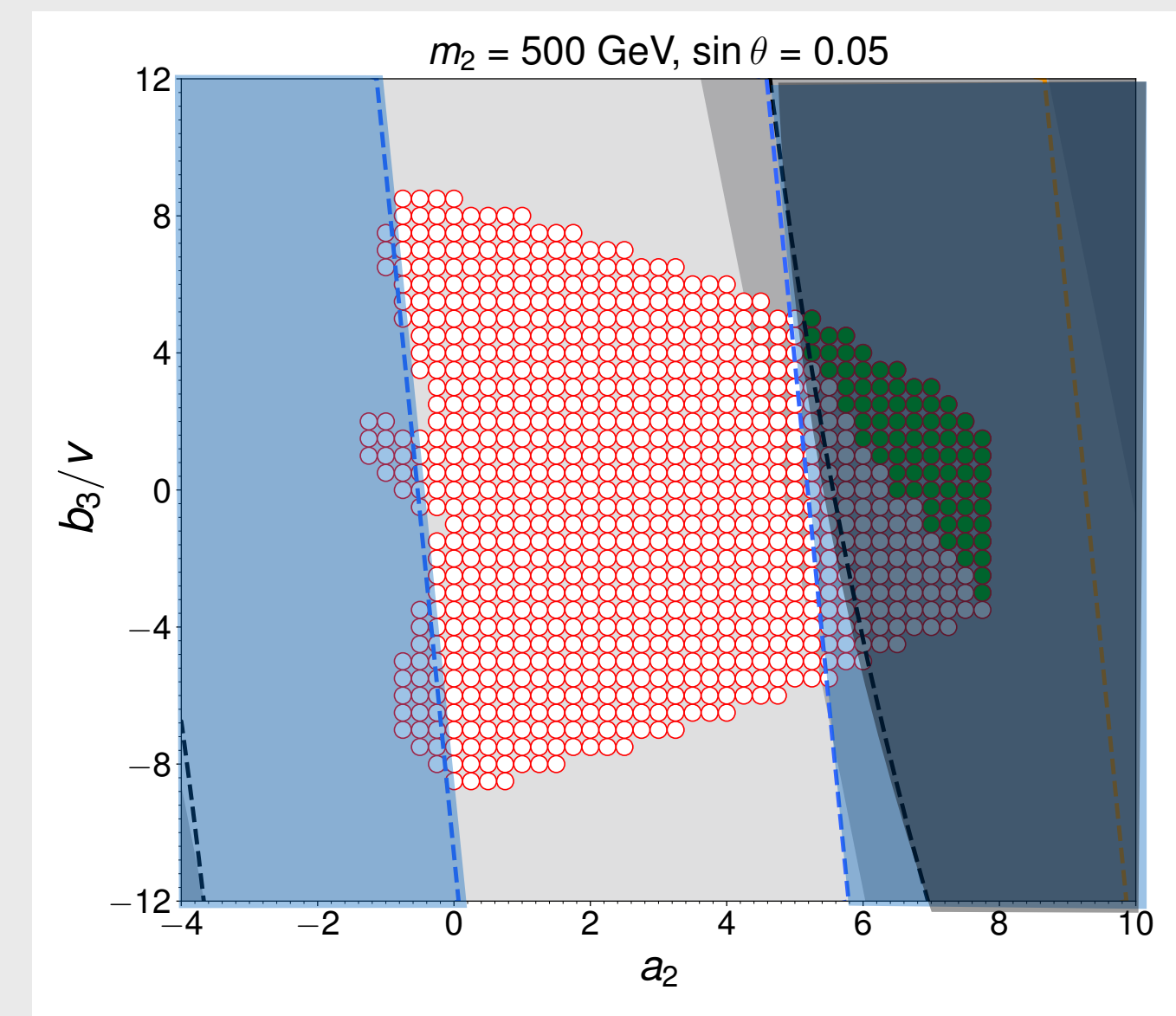
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parameters space of 1st order phase transition accessible by several probes

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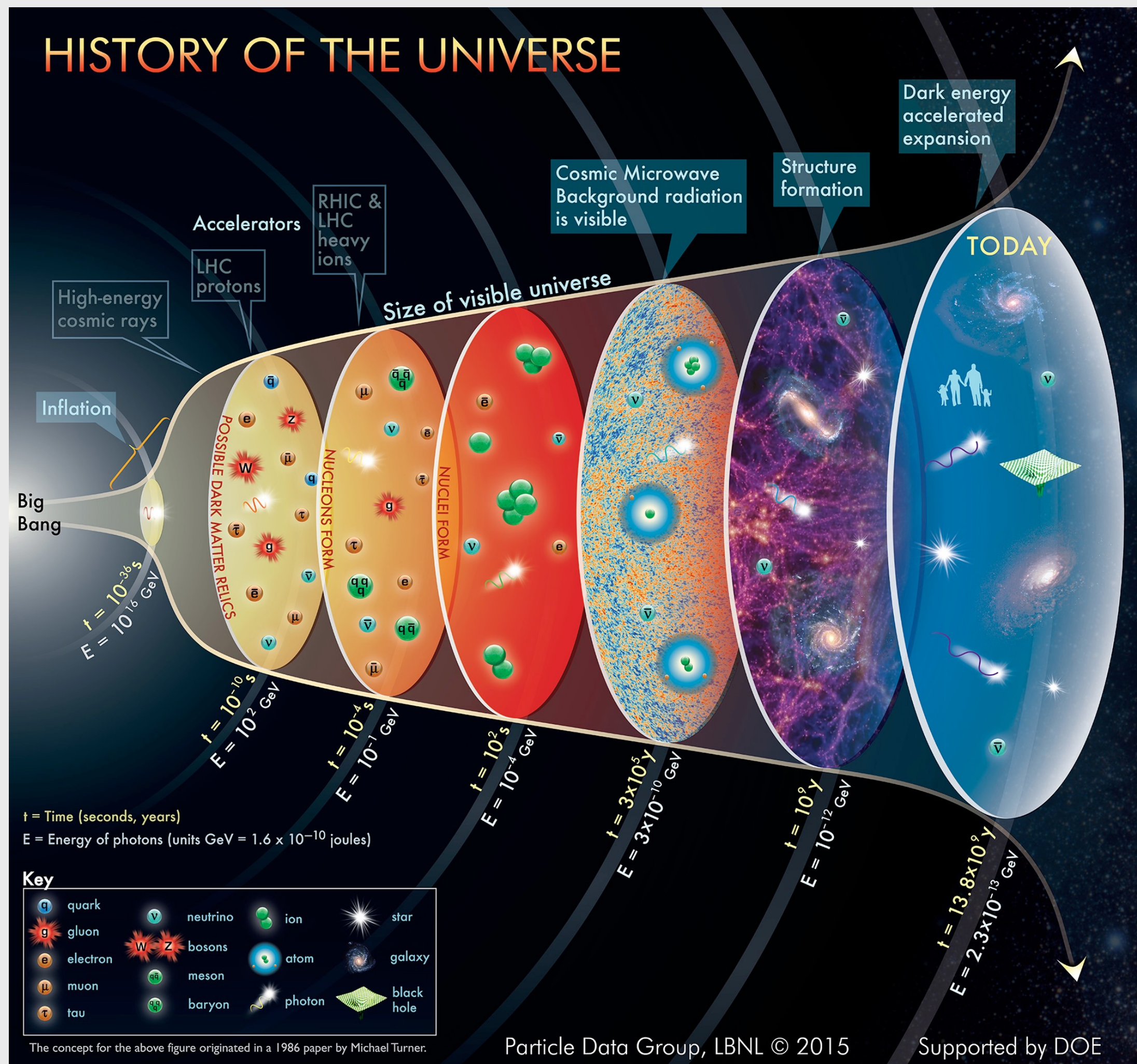
● 1st order phase transition

- Lepton colliders can observe directly these states
- Clean events allow to search for the most elusive particles
- The mass range is clearly well above the ZH threshold, we need higher energies!

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 a_2

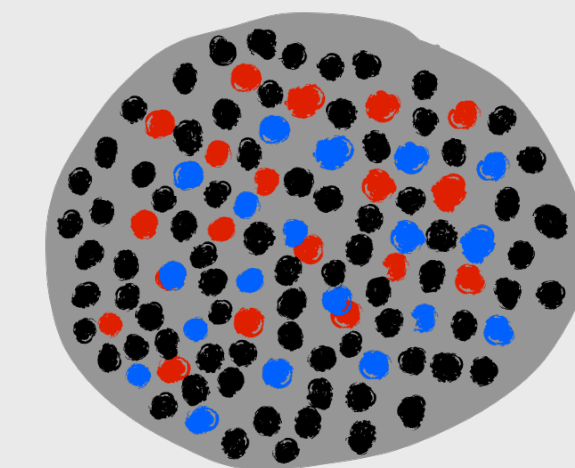
parameters space of 1st order phase transition accessible by several probes

Open Questions on the “big picture” on fundamental physics circa 2020



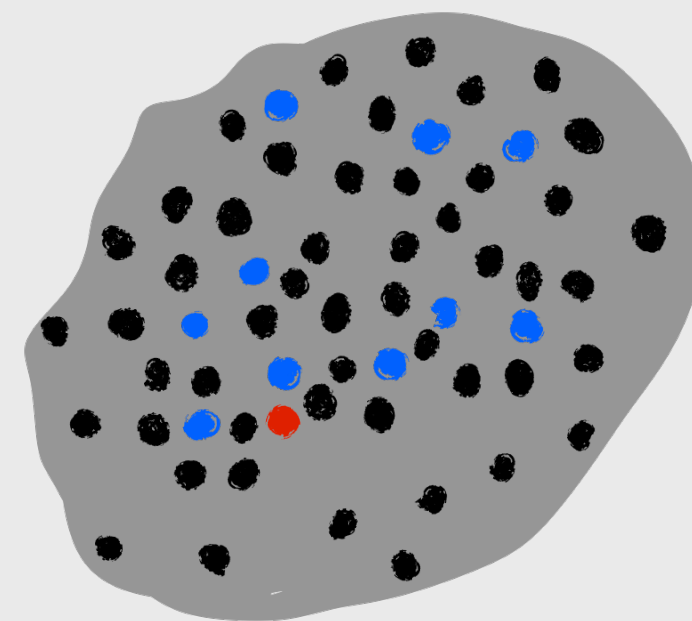
The observable Universe is made of matter, plus about 5 times as much dark matter

We need to go from this



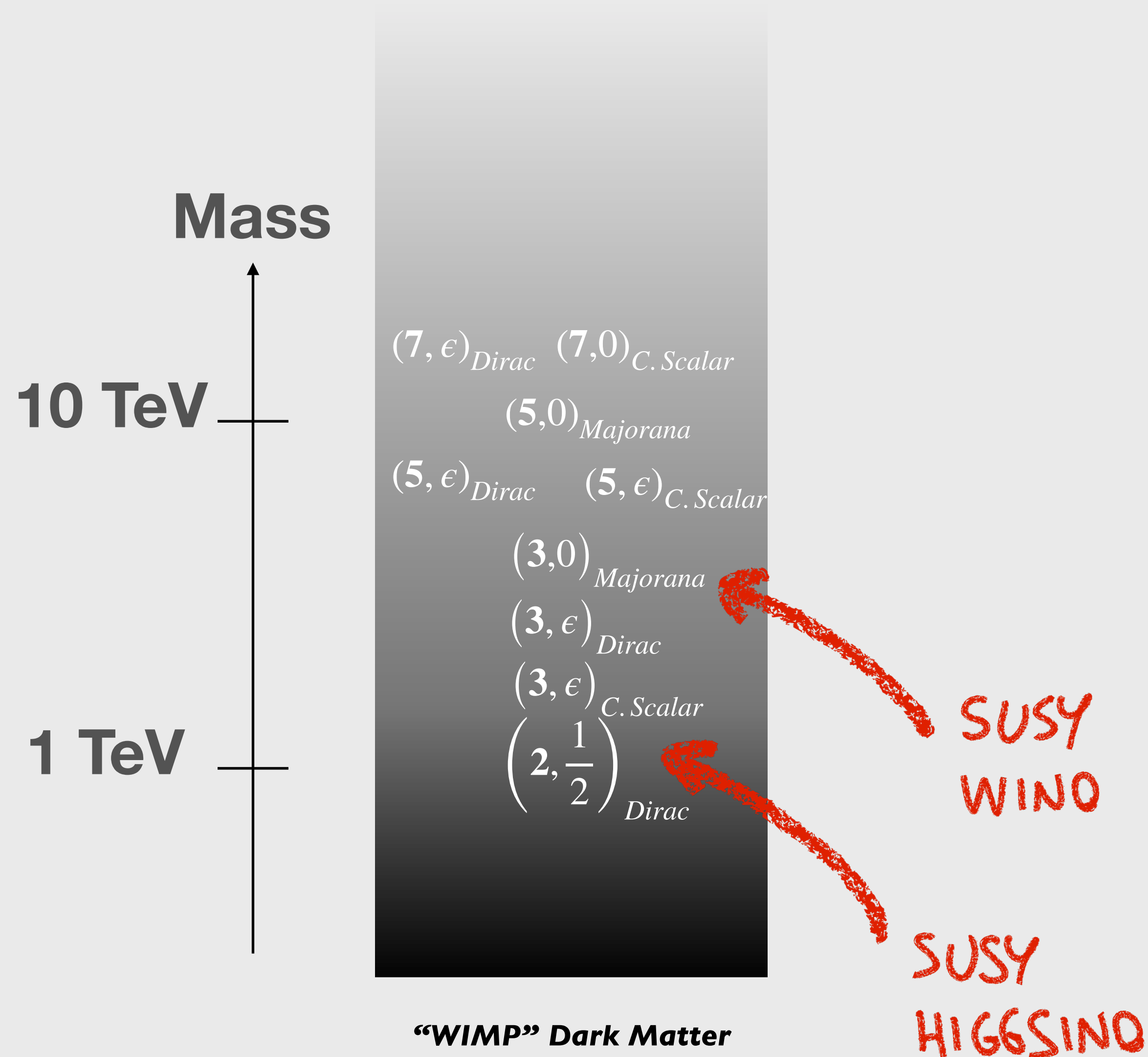
normal particles
dark matter
antiparticles

to this



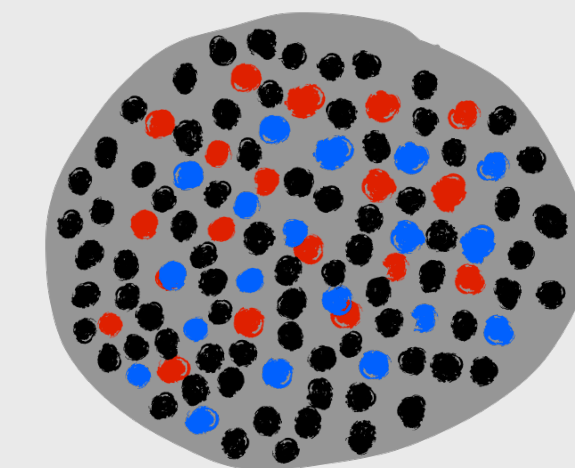
interactions rate from $\sigma = \left(\frac{g_{weak}}{M_{weak}} \right)^2$ are just about right!

Open Questions on the “big picture” on fundamental physics circa 2020



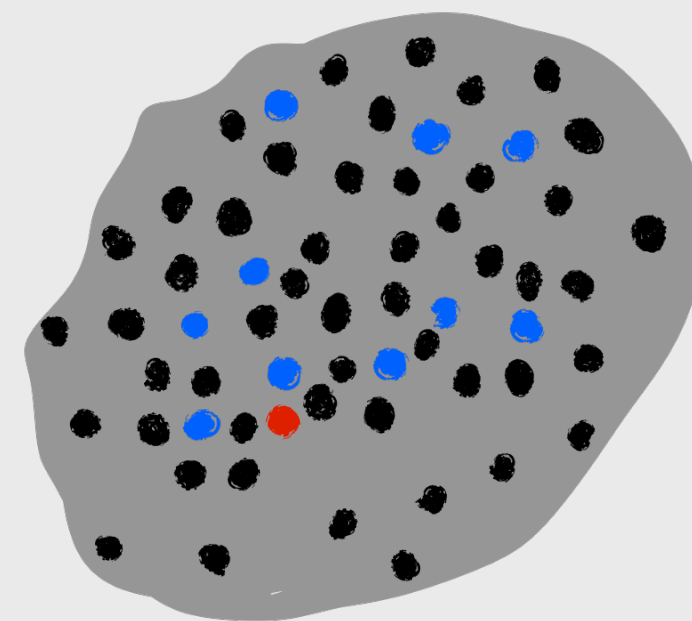
The observable Universe is made of matter, plus about 5 times as much dark matter

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normal particles
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to this



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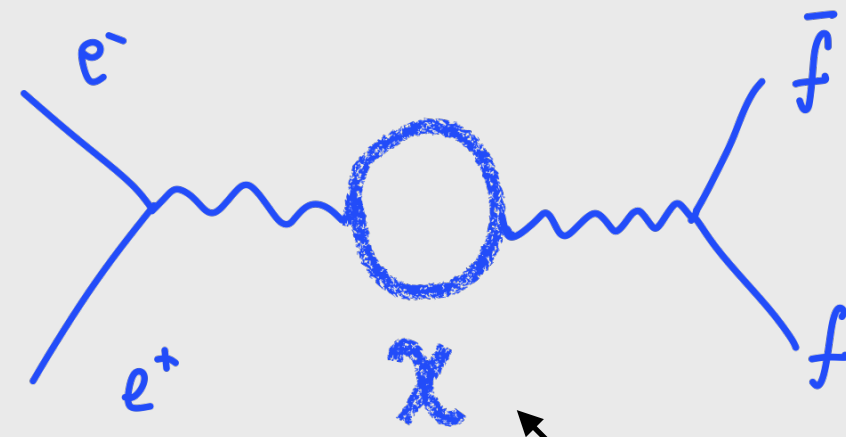
flashing concrete results for

Dark Matter at the weak scale

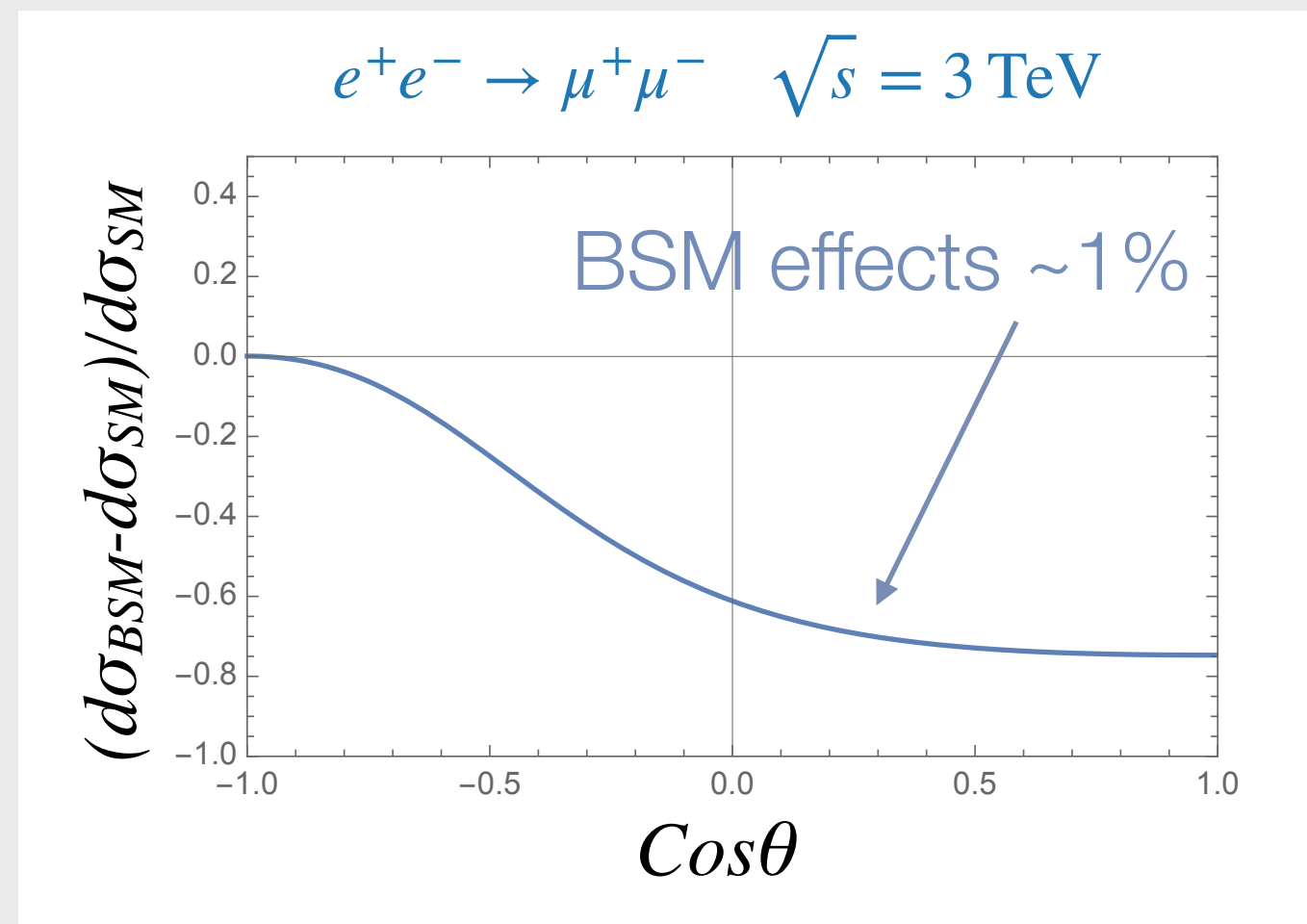
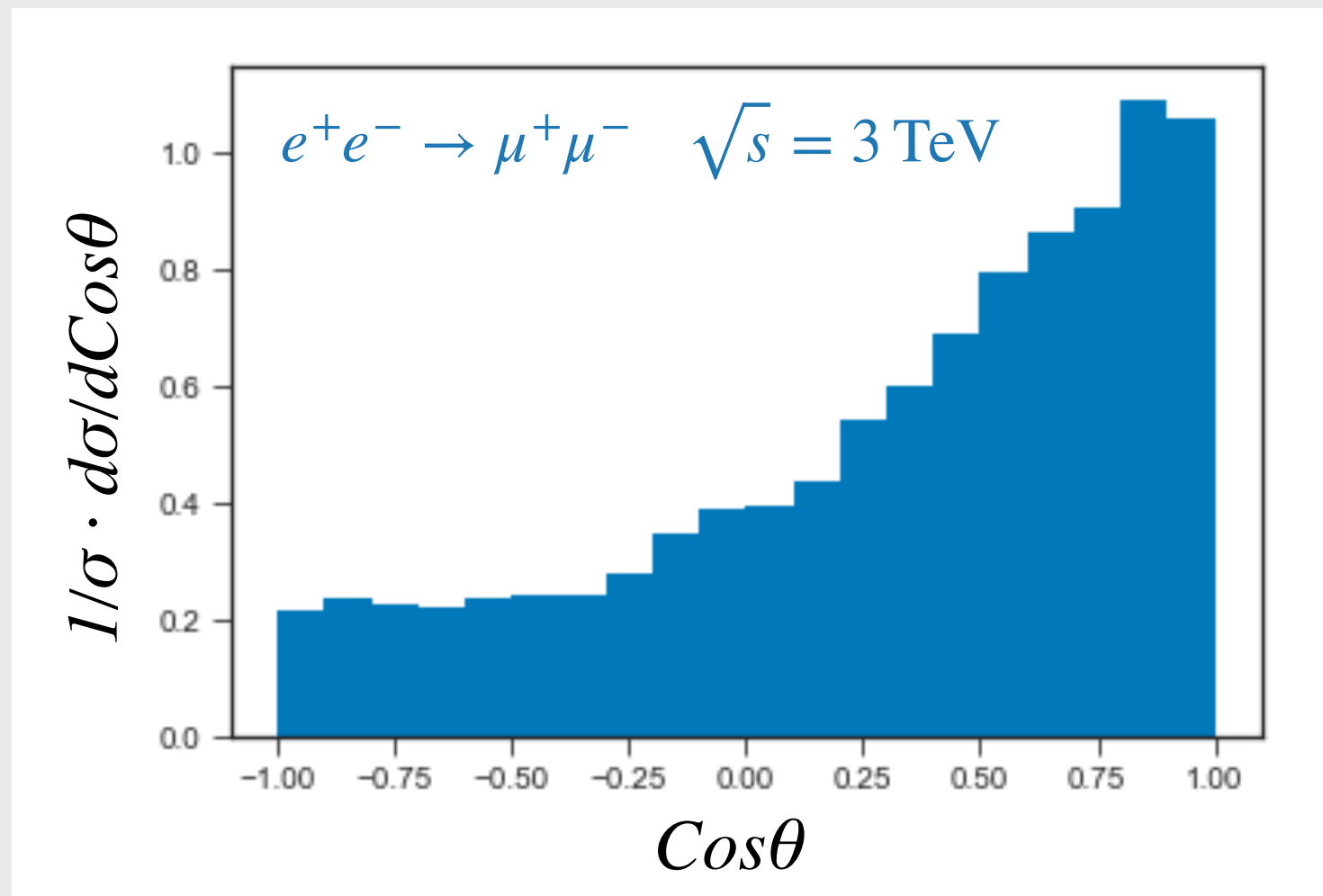
$$e^+ e^- \rightarrow f \bar{f}$$

PRECISION

ANGULAR DISTRIBUTION



χ is heavy/light new physics

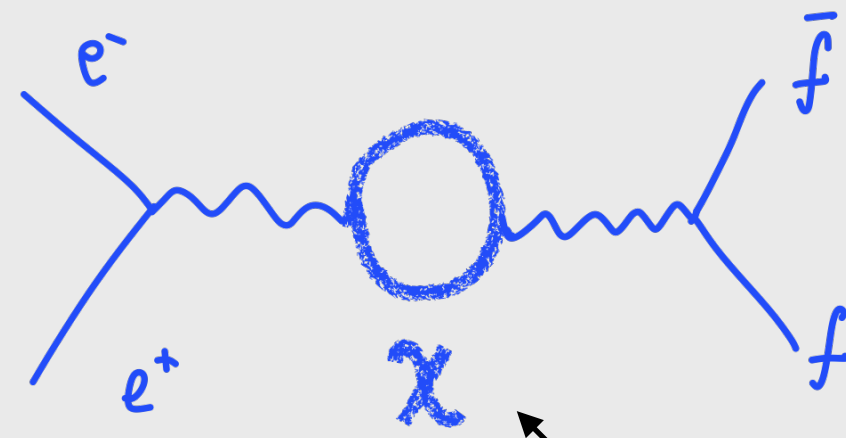


beams polarization is beneficial to increase NP effects

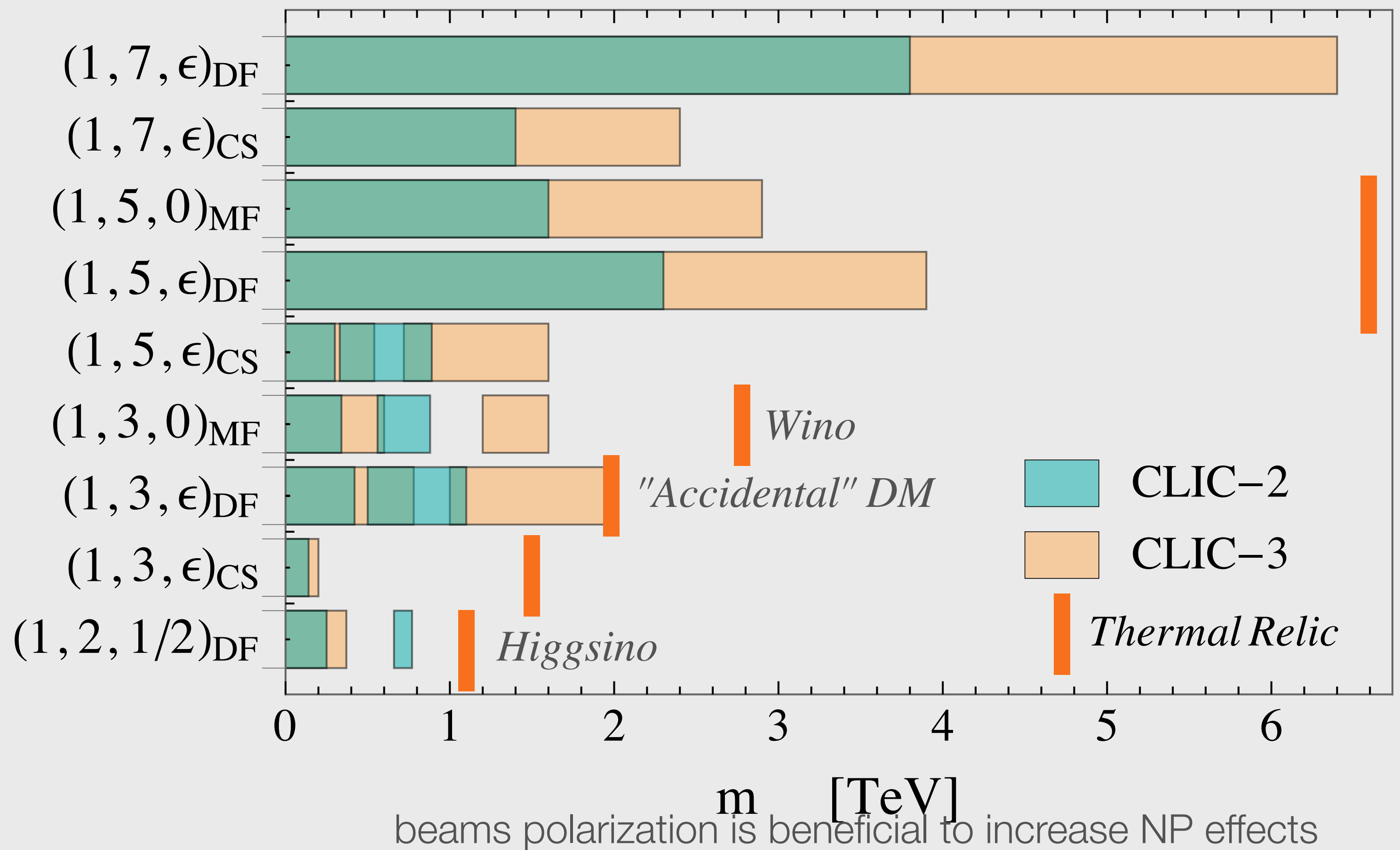
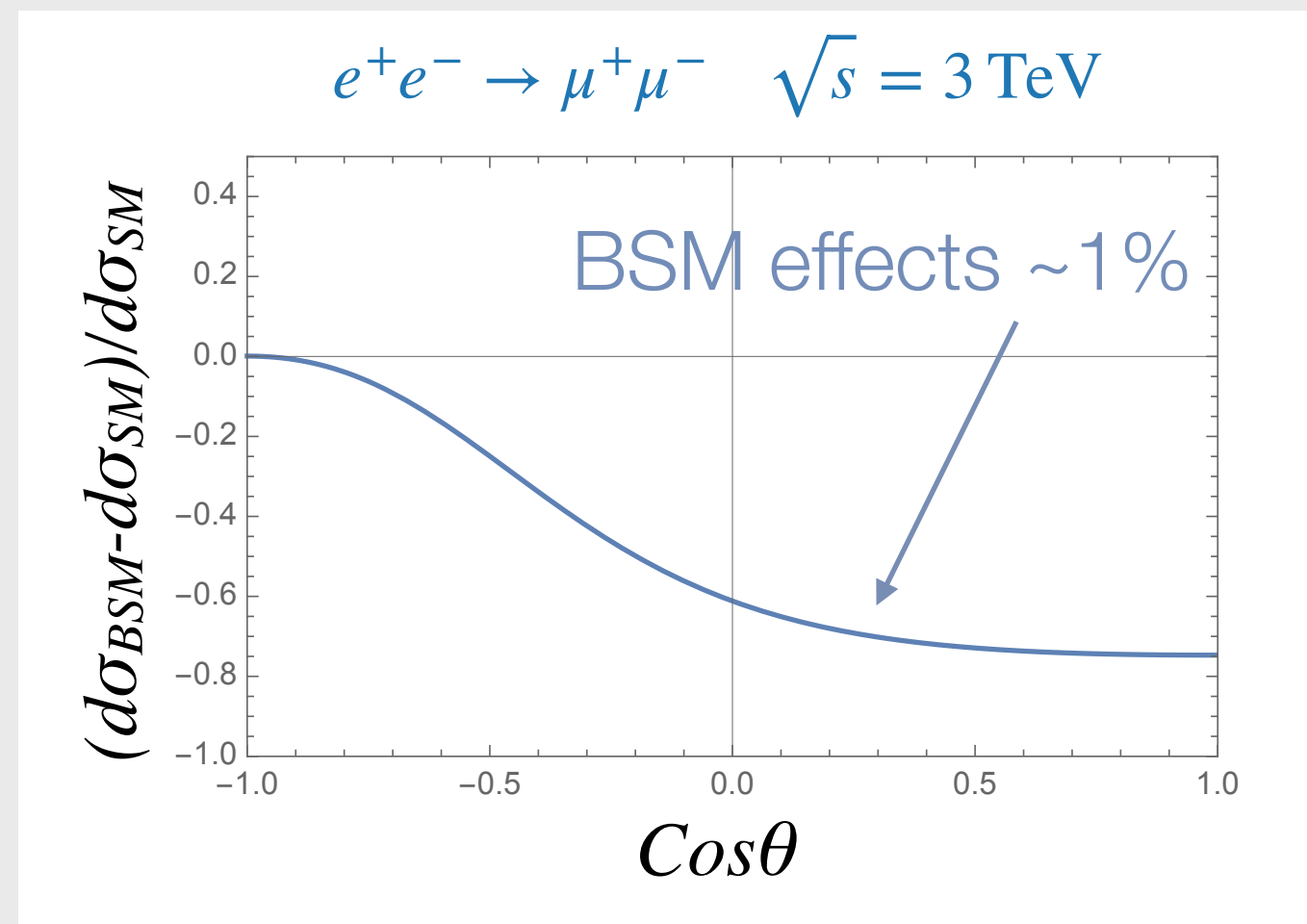
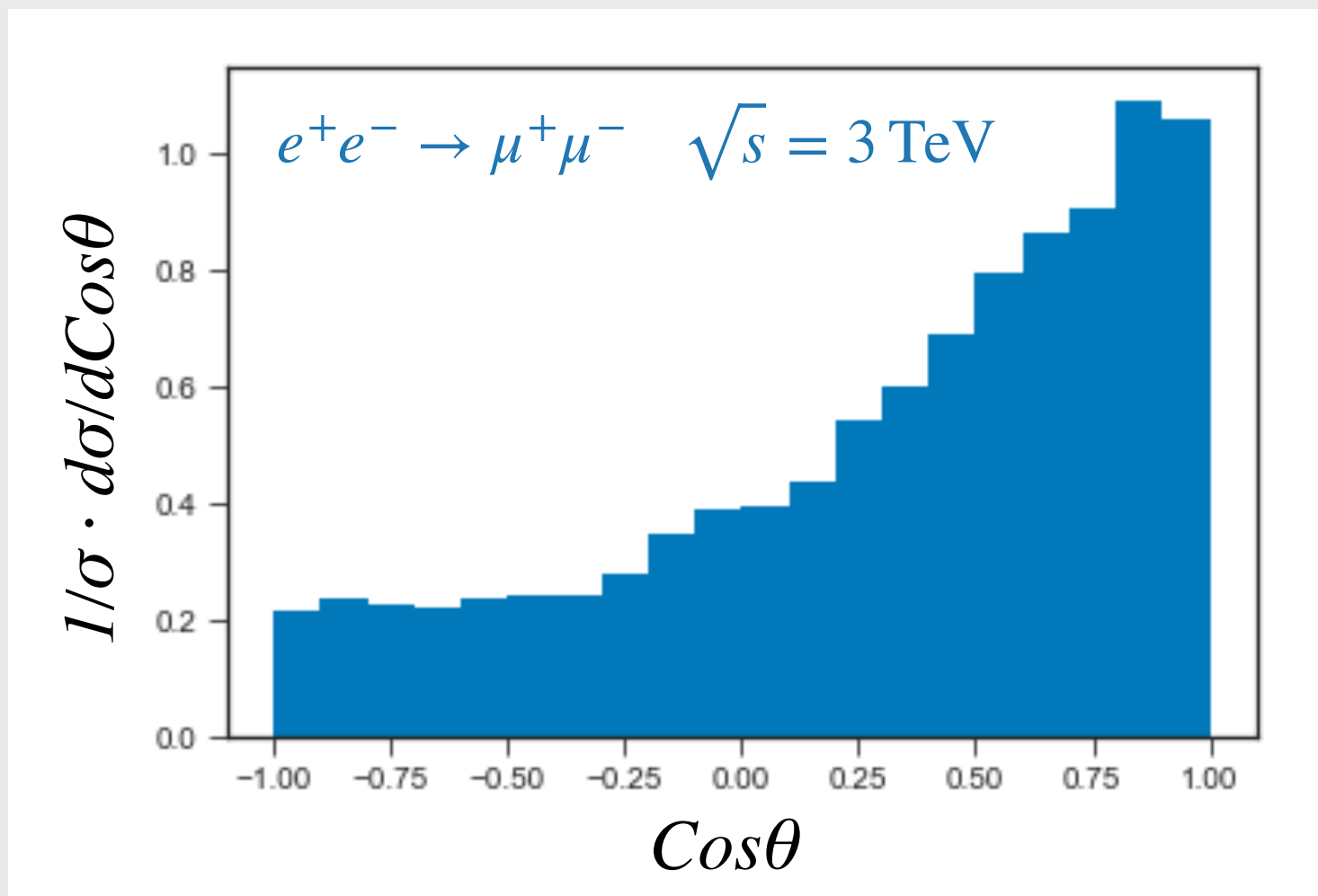
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PRECISION

ANGULAR DISTRIBUTION



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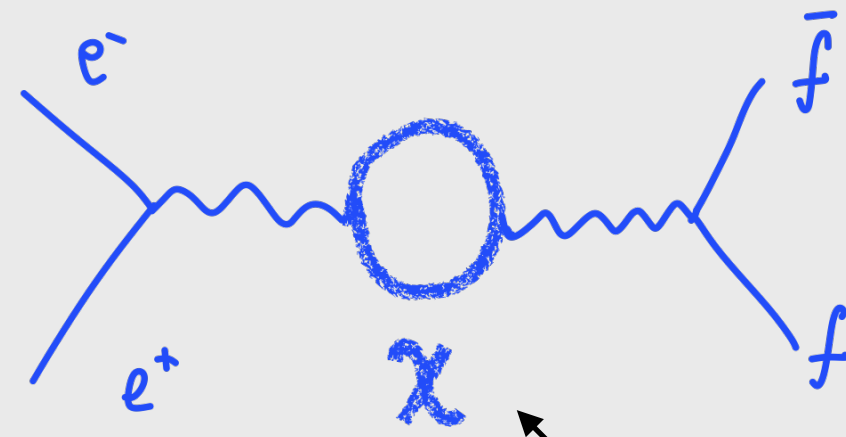


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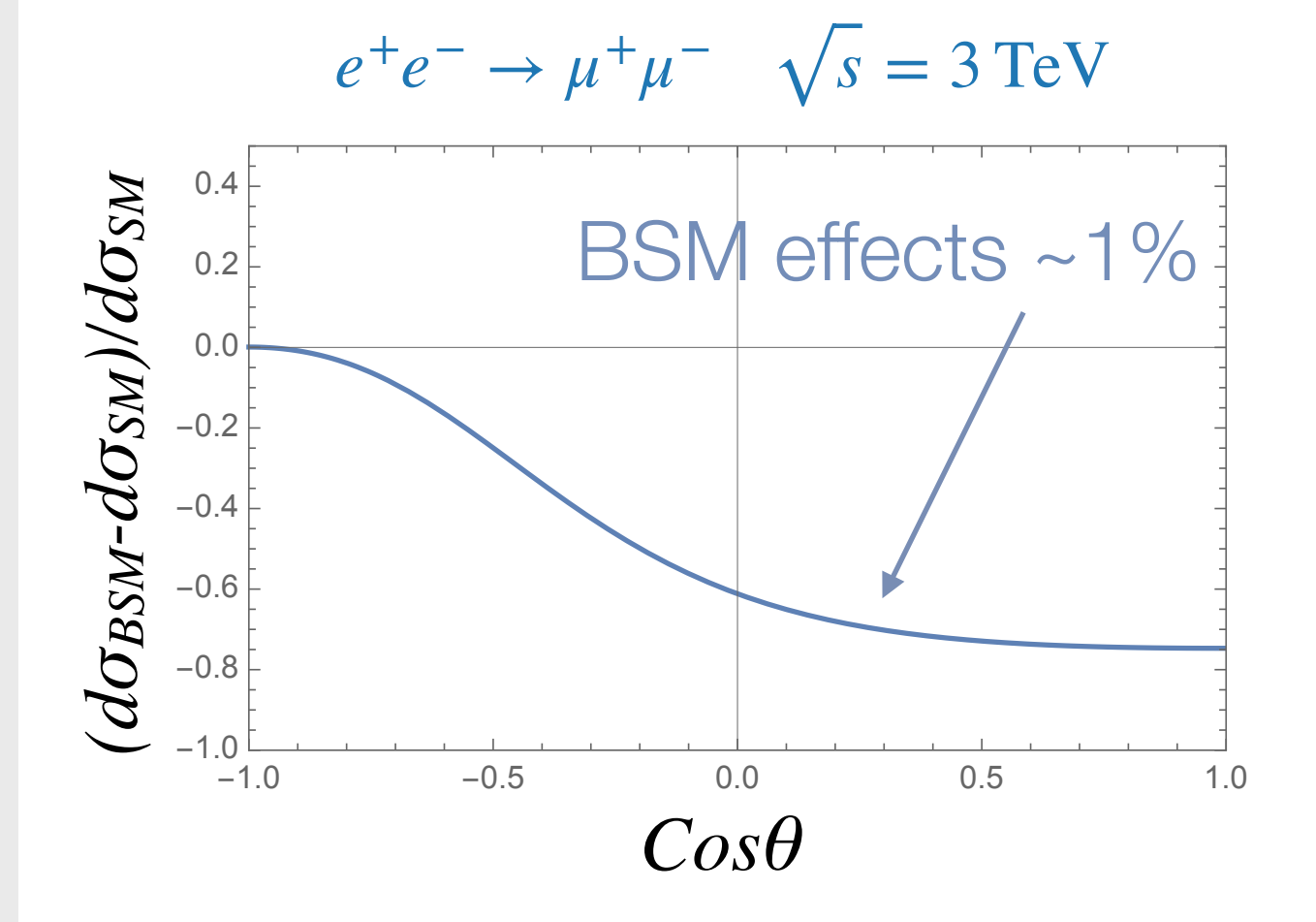
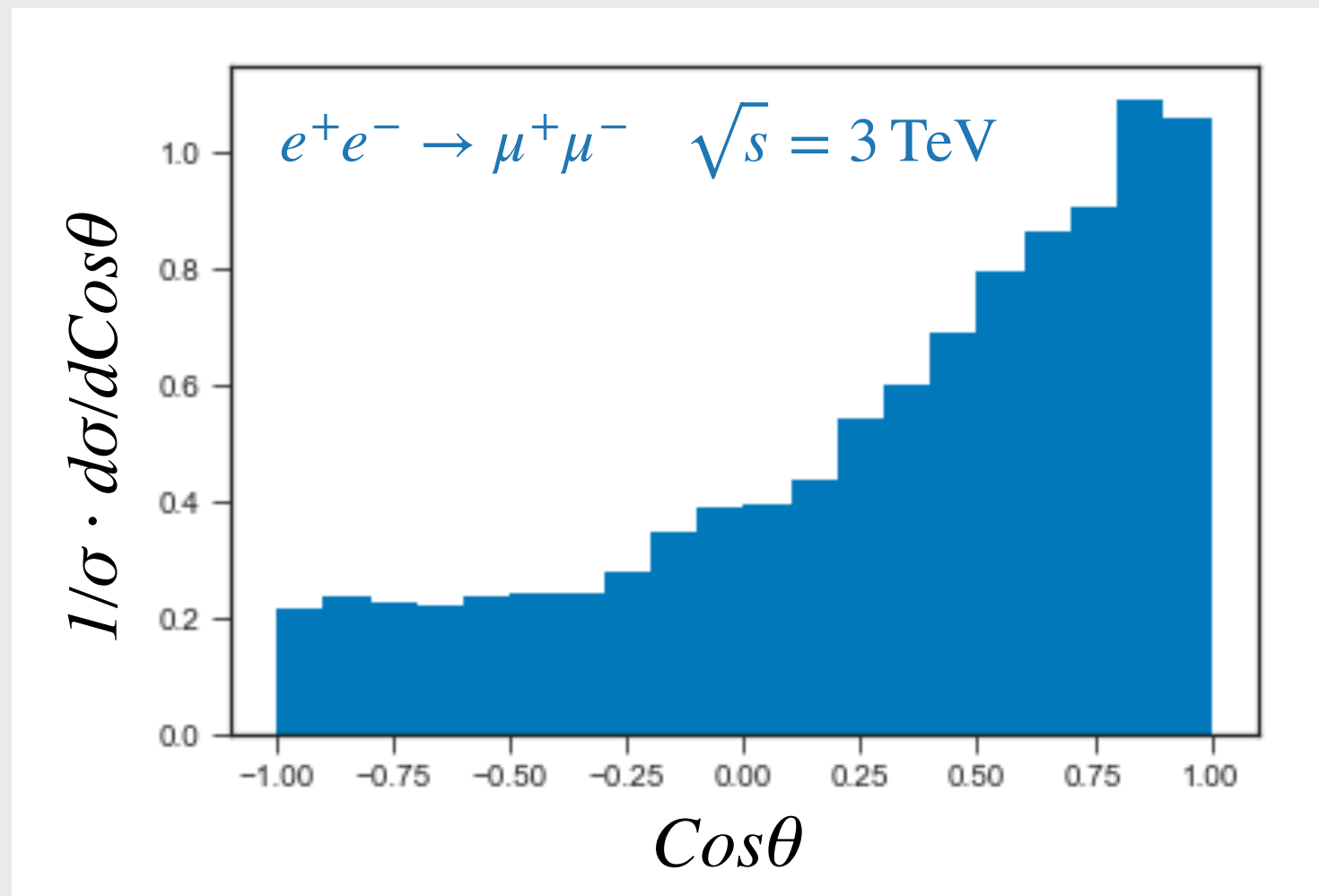
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PRECISION

ANGULAR DISTRIBUTION



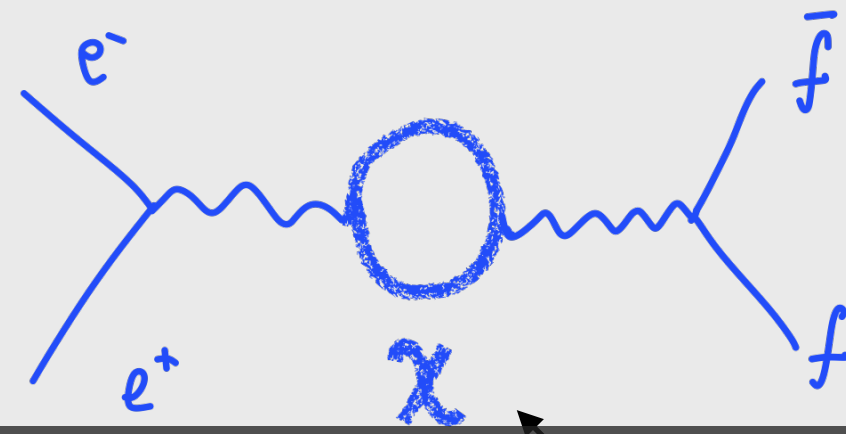
χ is heavy/light new physics



χ / m_χ [TeV]	DM	HL-LHC	HE-LHC	FCC-100	CLIC-3	Muon-14
$(1, 2, 1/2)_{DF}$	1.1	—	—	—	0.4	0.6
$(1, 3, \epsilon)_{CS}$	1.6	—	—	—	0.2	0.2
$(1, 3, \epsilon)_{DF}$	2.0	—	0.6	1.5	0.8 & [1.0, 2.0]	2.2 & [6.3, 7.1]
$(1, 3, 0)_{MF}$	2.8	—	—	0.4	0.6 & [1.2, 1.6]	1.0
$(1, 5, \epsilon)_{CS}$	6.6	0.2	0.4	1.0	0.5 & [0.7, 1.6]	1.6
$(1, 5, \epsilon)_{DF}$	6.6	1.5	2.8	7.1	3.9	11
$(1, 5, 0)_{MF}$	14	0.9	1.8	4.4	2.9	3.5 & [5.1, 8.7]
$(1, 7, \epsilon)_{CS}$	16	0.6	1.3	3.2	2.4	2.5 & [3.5, 7.4]
$(1, 7, \epsilon)_{DF}$	16	2.1	4.0	11	6.4	18

- Comprehensive tool to explore new electroweak particles
- Can probe valid dark matter candidates!

$$e^+ e^- \rightarrow f \bar{f}$$



χ is heavy/light new physics

PRECISION

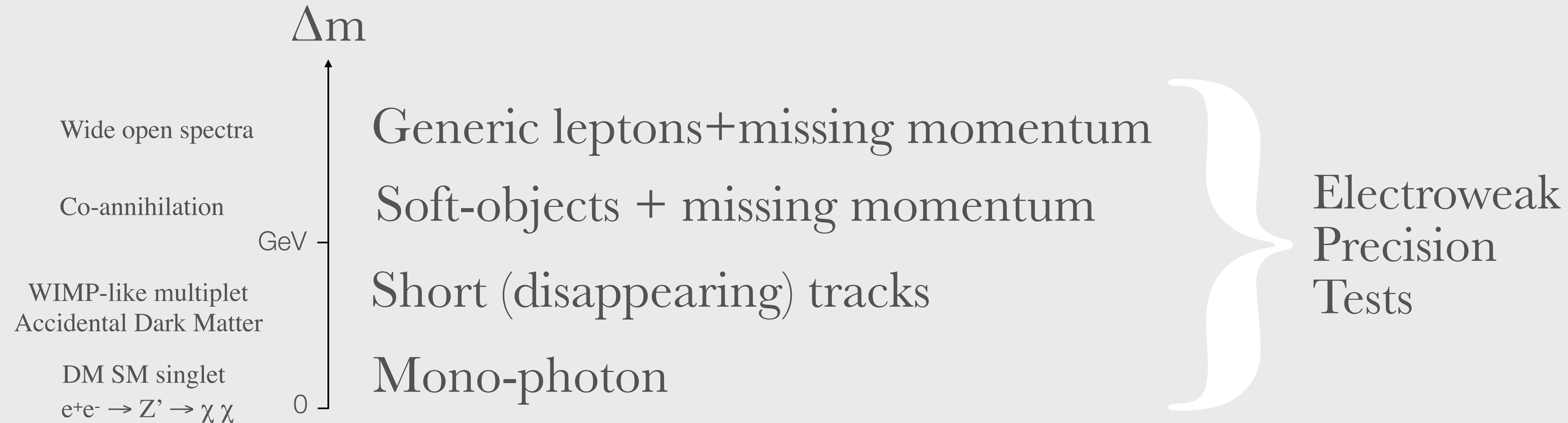
ANGULAR DISTRIBUTION

- Lepton colliders can observe these states directly and indirectly
- Clean events allow to search for the most elusive particles
- The mass range is clearly (very) well above the ZH threshold, we need higher energies!

-1.0 -0.5 0.0 0.5 1.0

$\cos\theta$

Electroweak Dark Matter: LSP (+NLSP)



flashing concrete results for

The size of the Higgs boson

Effects of the size of the Higgs boson

$h \sim \pi$

STRONGLY INTERACTING LIGHT HIGGS

$$\begin{aligned}
 \mathcal{L}_{universal}^{d=6} = & c_H \frac{g_*^2}{m_*^2} \mathcal{O}_H + c_T \frac{N_c \epsilon_q^4 g_*^4}{(4\pi)^2 m_*^2} \mathcal{O}_T + c_6 \lambda \frac{g_*^2}{m_*^2} \mathcal{O}_6 + \frac{1}{m_*^2} [c_W \mathcal{O}_W + c_B \mathcal{O}_B] \\
 & + \frac{g_*^2}{(4\pi)^2 m_*^2} [c_{HW} \mathcal{O}_{HW} + c_{HB} \mathcal{O}_{HB}] + \frac{y_t^2}{(4\pi)^2 m_*^2} [c_{BB} \mathcal{O}_{BB} + c_{GG} \mathcal{O}_{GG}] \\
 & + \frac{1}{g_*^2 m_*^2} [c_{2W} g^2 \mathcal{O}_{2W} + c_{2B} g'^2 \mathcal{O}_{2B}] + c_{3W} \frac{3! g^2}{(4\pi)^2 m_*^2} \mathcal{O}_{3W} \\
 & + c_{y_t} \frac{g_*^2}{m_*^2} \mathcal{O}_{y_t} + c_{y_b} \frac{g_*^2}{m_*^2} \mathcal{O}_{y_b}
 \end{aligned}$$

$$1/f \sim g_*/m_*$$

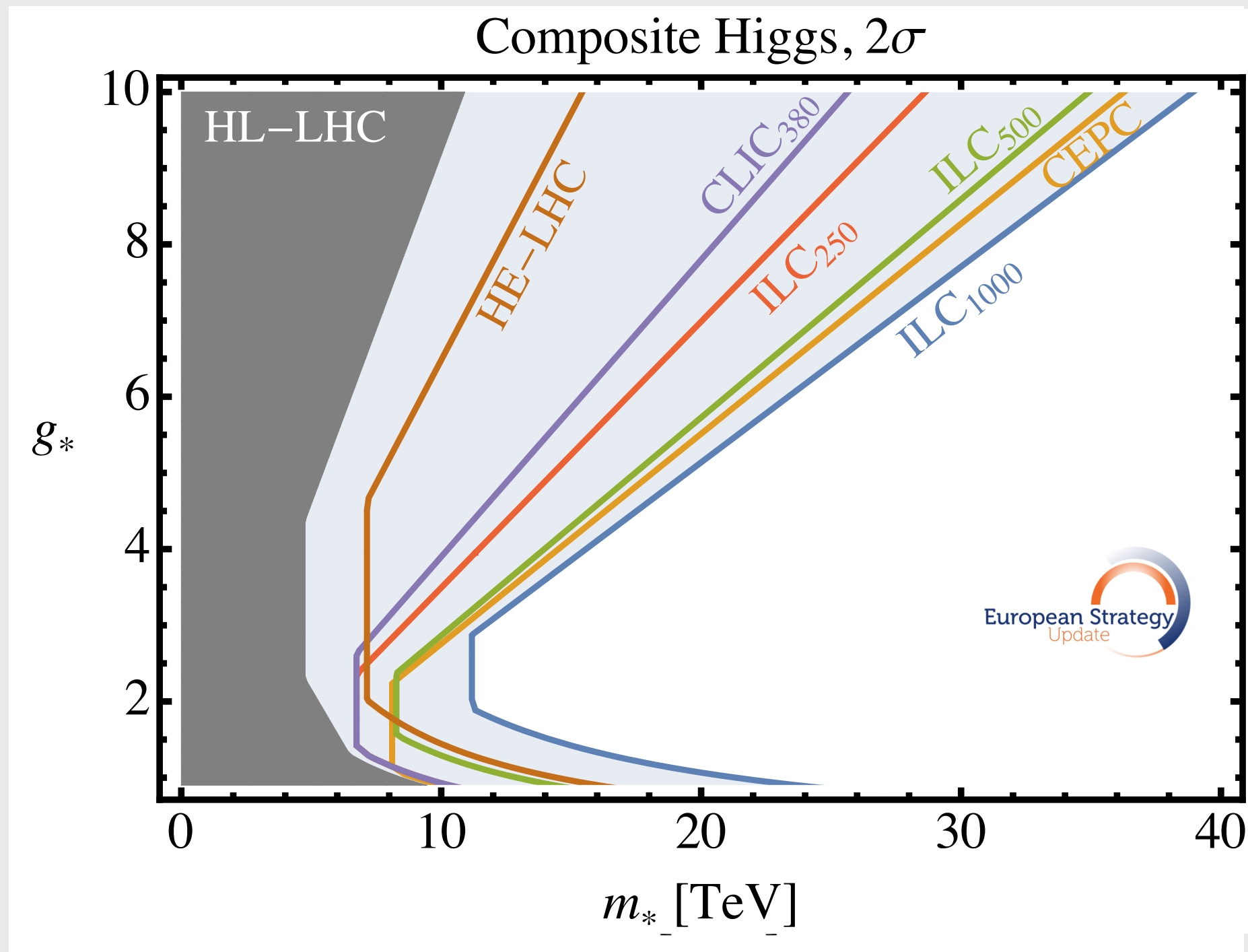
$$1/(g_* f) \sim 1/m_*$$

$$g_{SM}/(g_* f) \sim g_{SM}/m_*$$

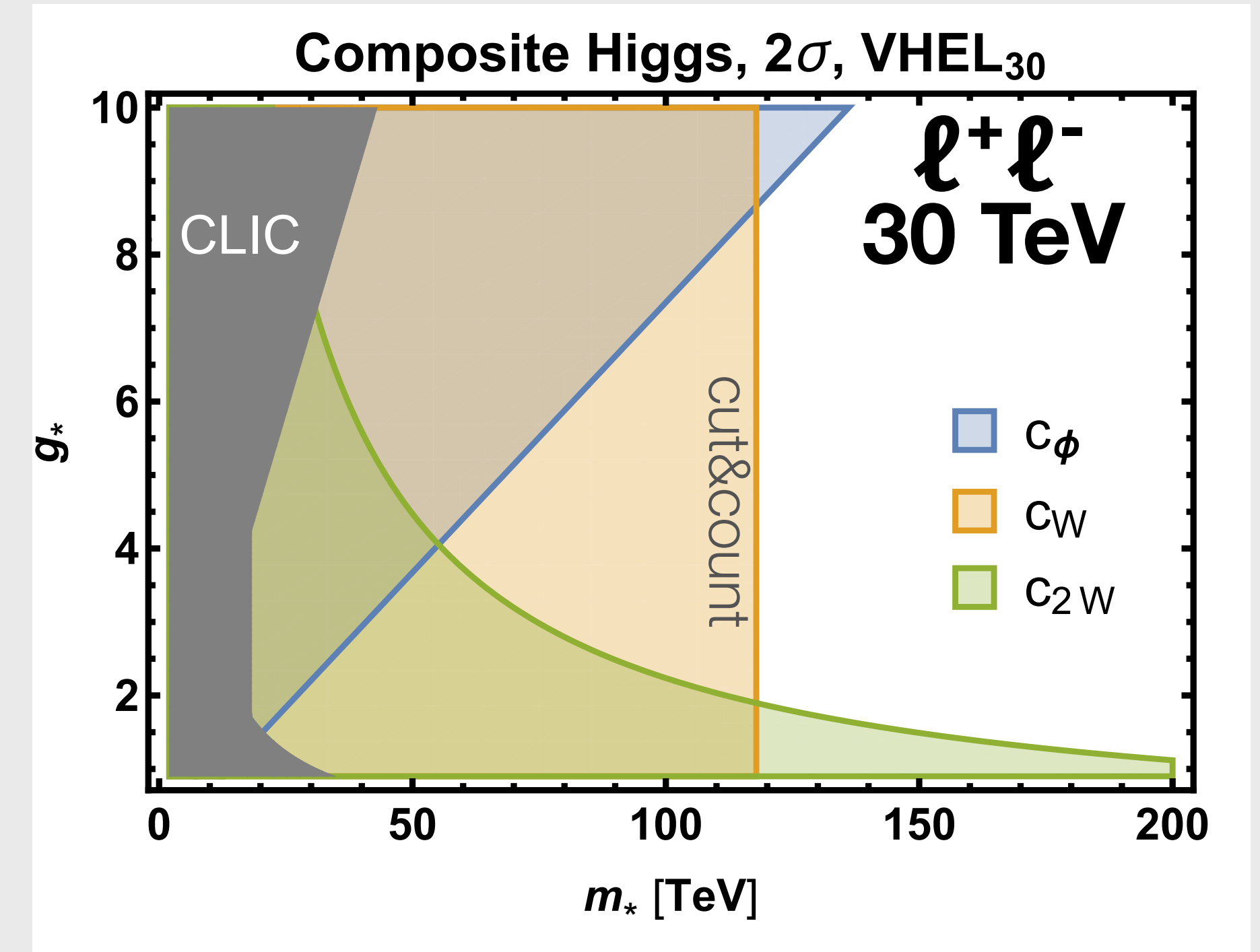
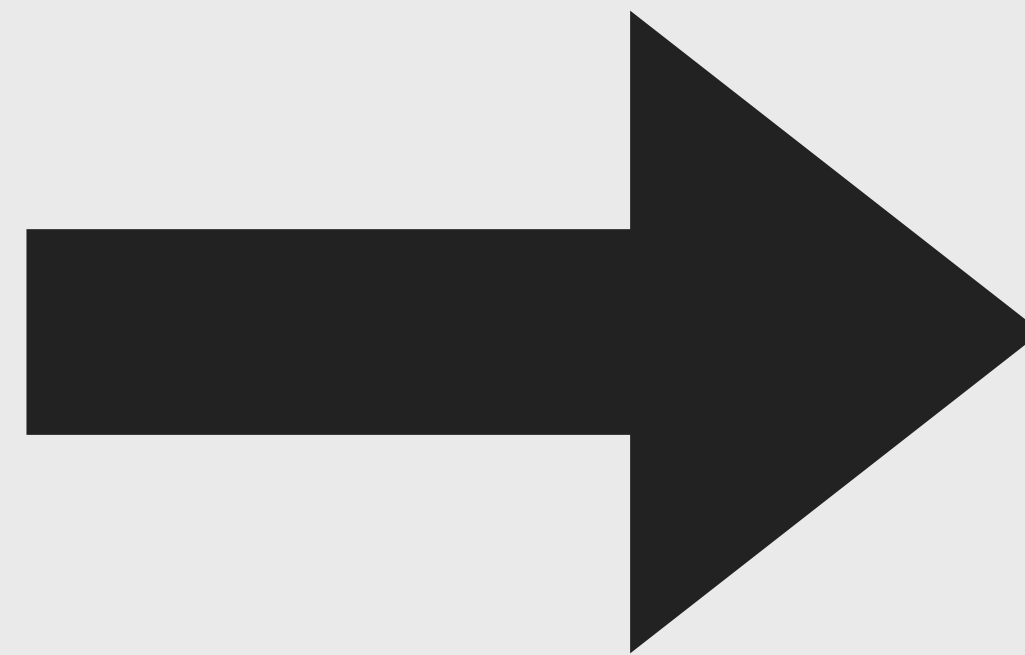
$$\ell_{Higgs} \sim 1/m_*$$



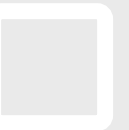
Looking ahead



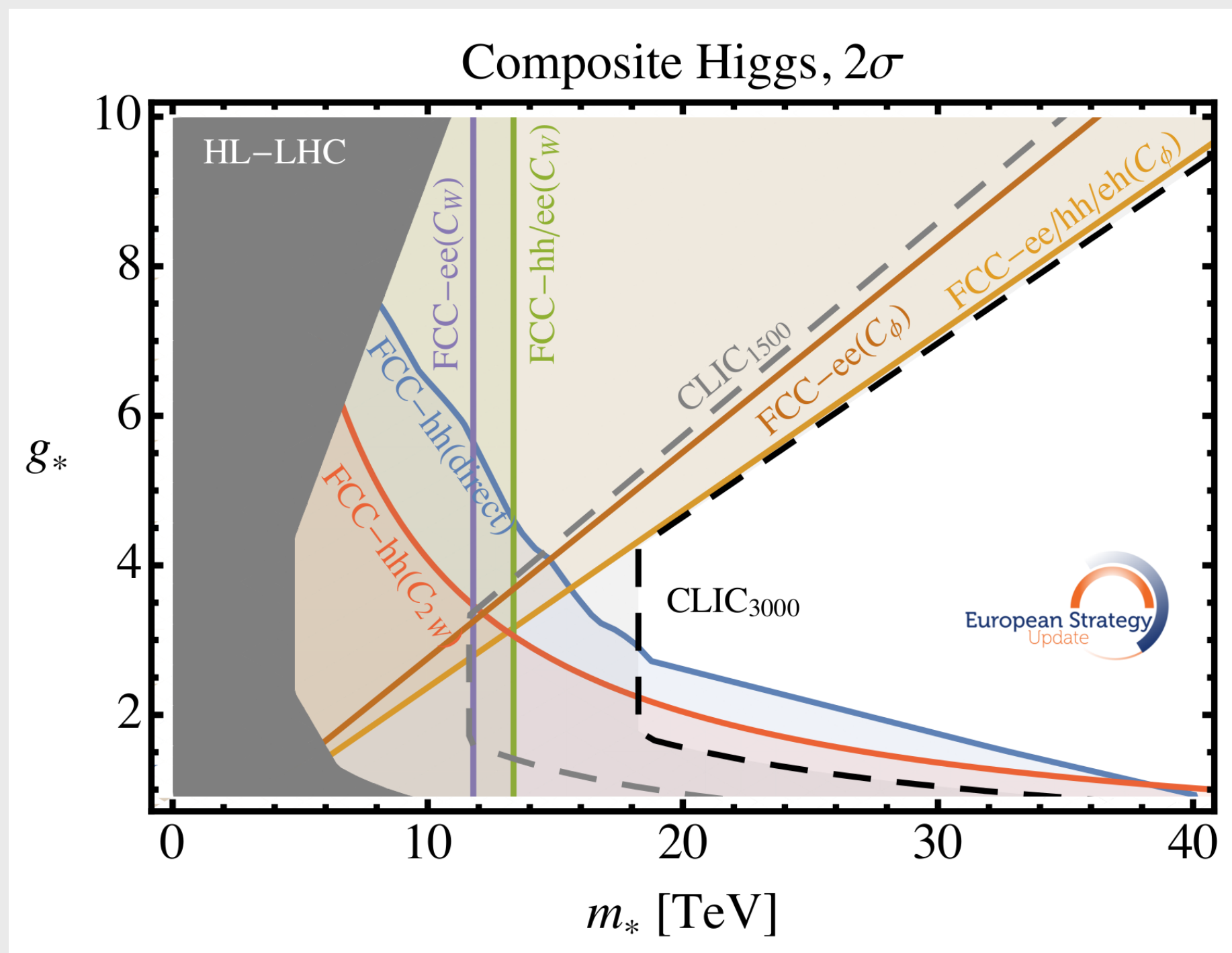
**compositeness at
10 TeV-20 TeV**



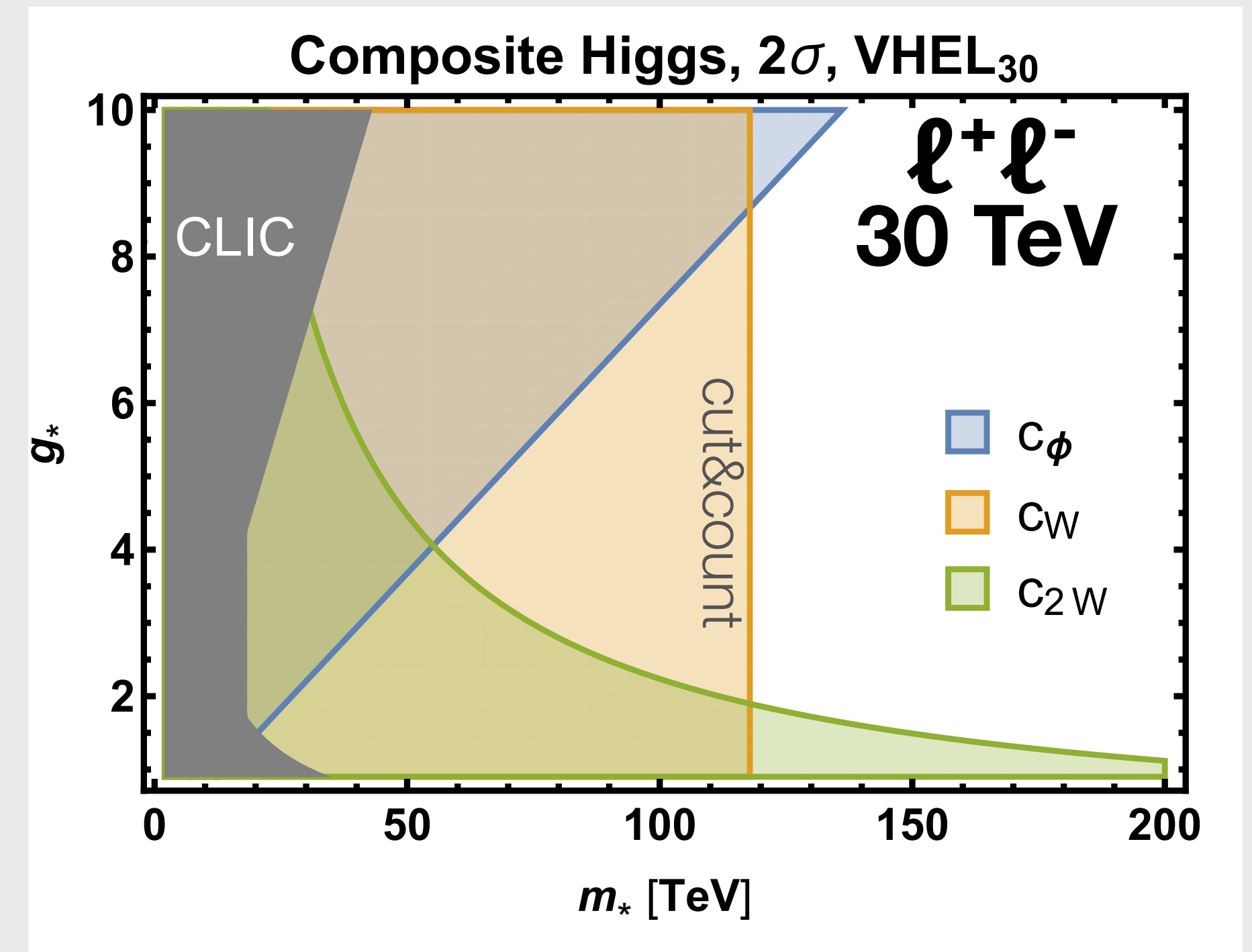
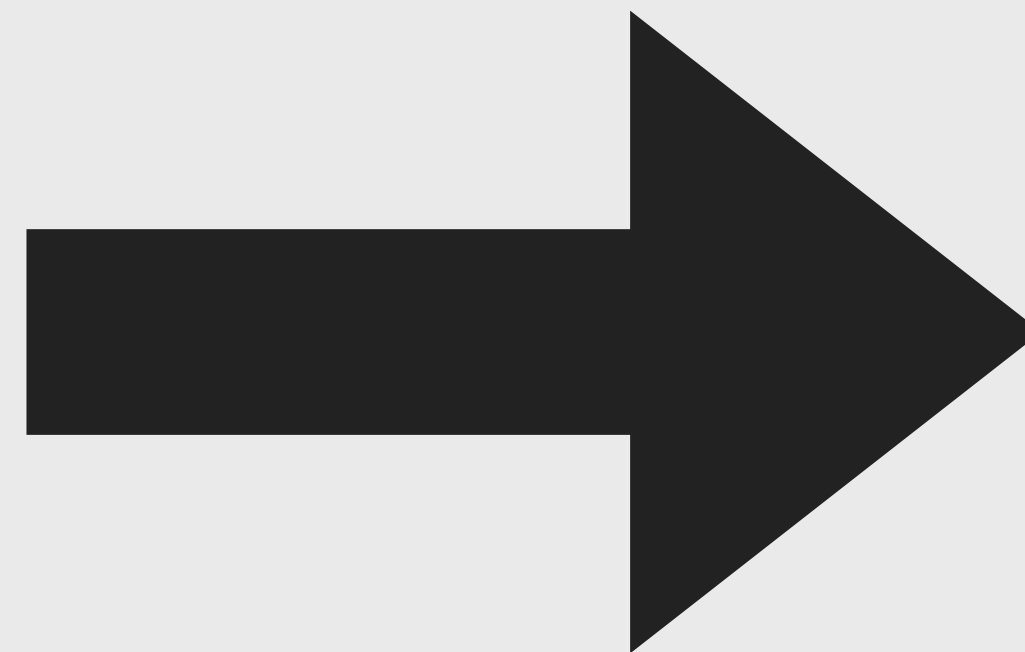
**compositeness at
100 TeV-500 TeV**



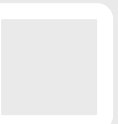
Looking ahead



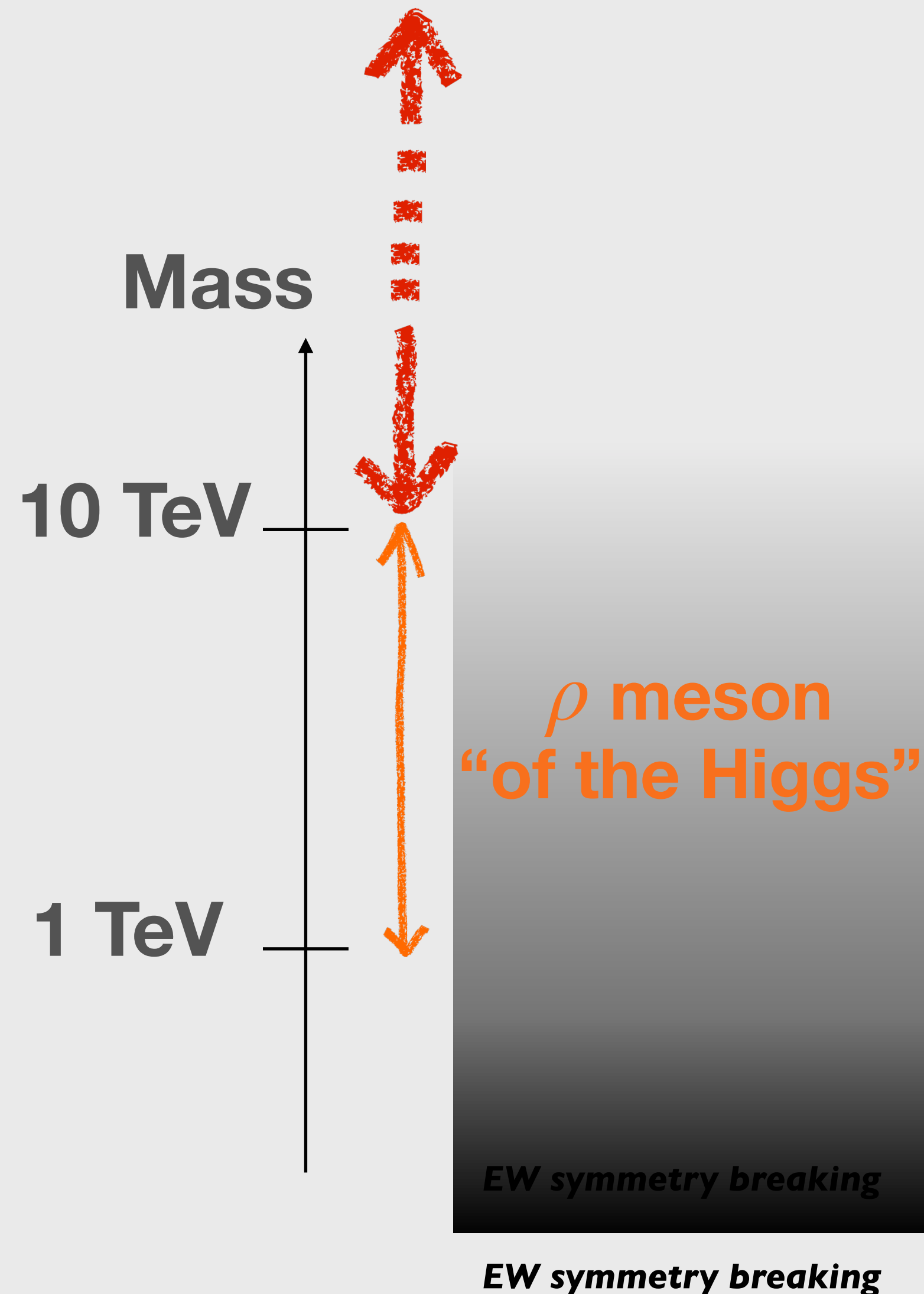
**compositeness at
10 TeV-20 TeV**



**compositeness at
100 TeV-500 TeV**



Open Questions on the “big picture” on fundamental physics circa 2020



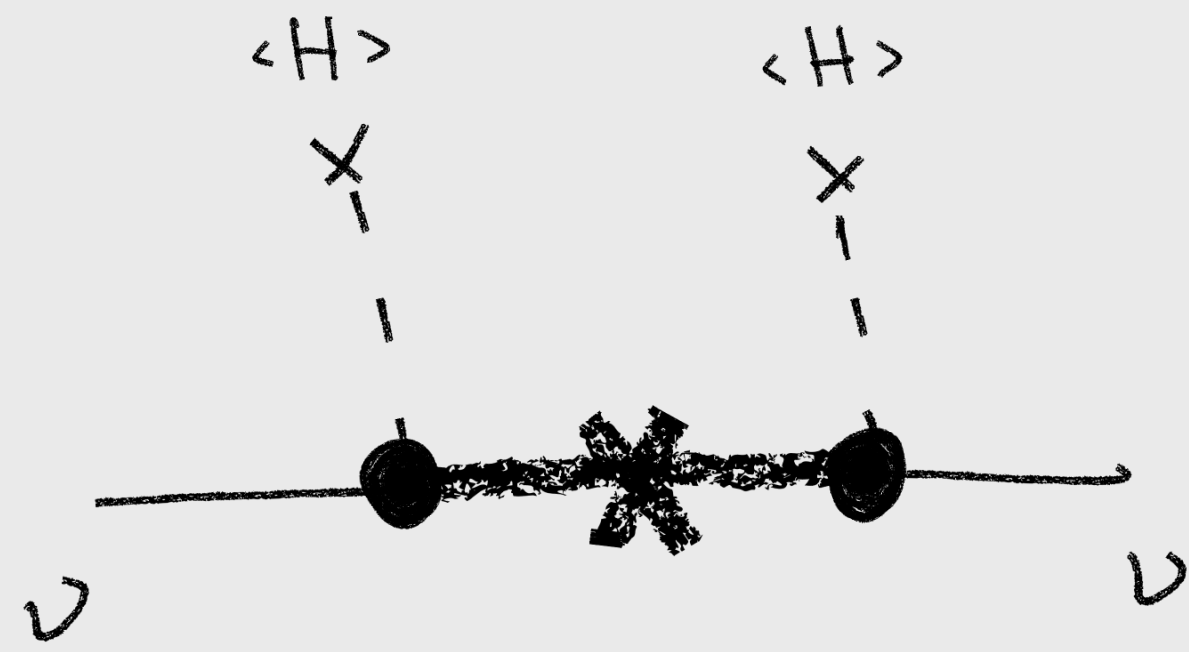
We might be in a situation like QCD, where the ρ meson is only somewhat heavier than the pion, or in a situations where it is much heavier.

Both cases have profound consequences for telling **what the Higgs boson really is.**

Neutrino mass mechanisms

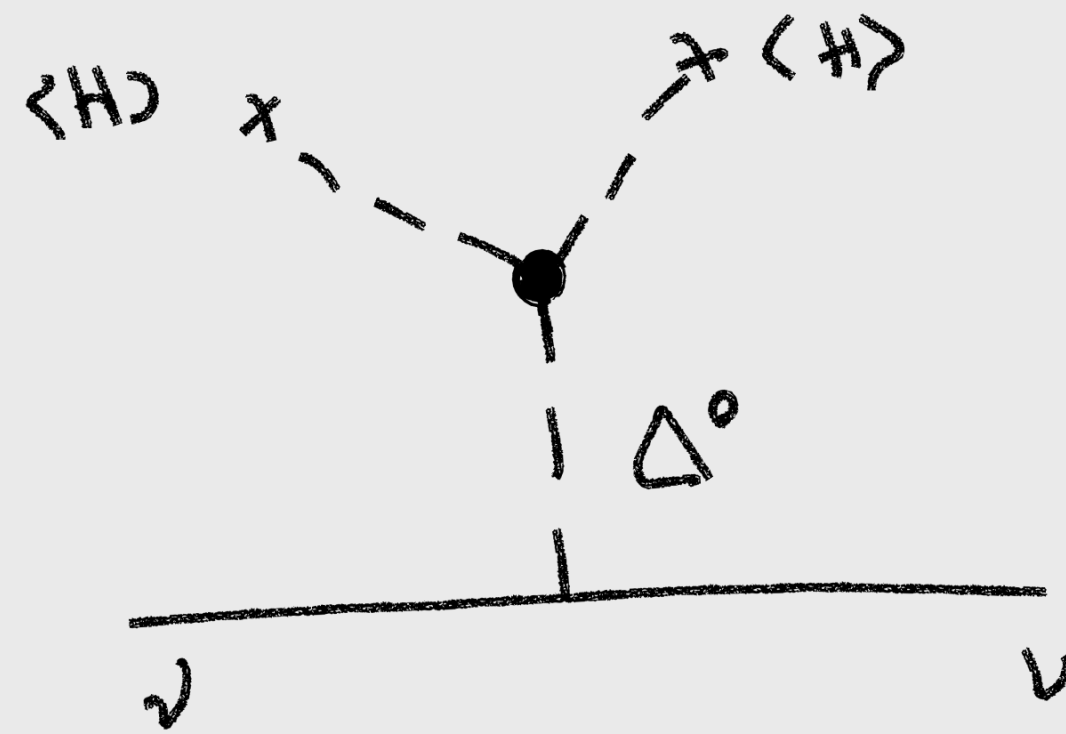
LEPTON

NUMBER BREAKING



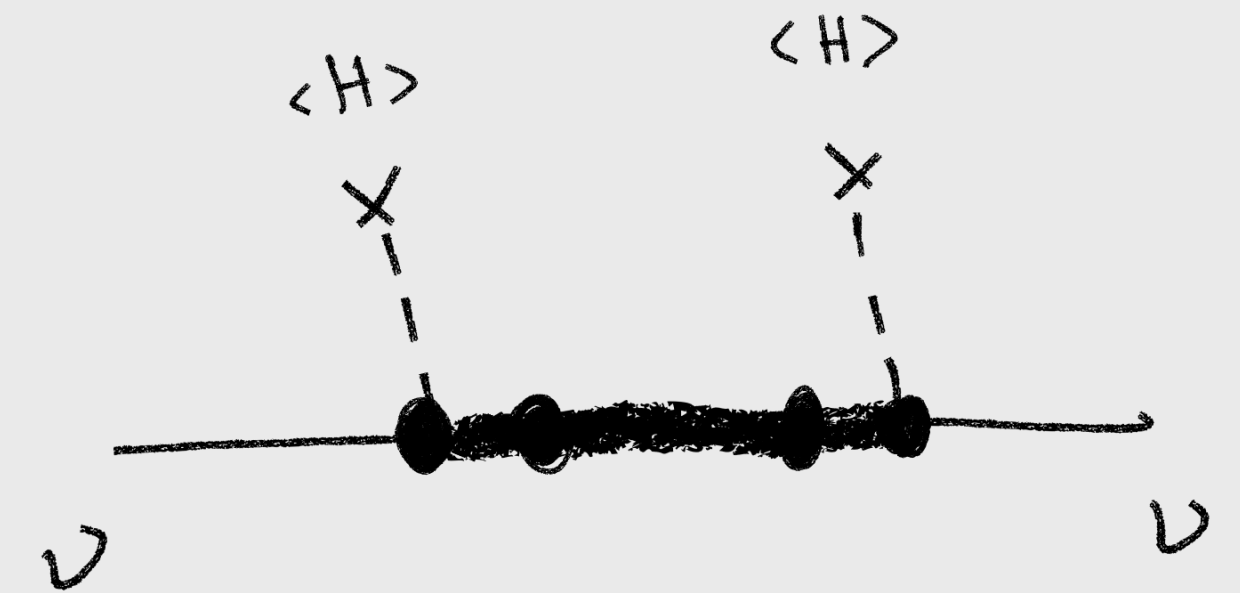
$$m_\nu = \frac{(\text{coupling})^2 \langle H \rangle^2}{M_{\text{heavy}}} \rightarrow \text{SMALL}$$

$M_{\text{heavy}} \rightarrow \text{LARGE}$



$$m_\nu = \frac{(\text{coupling})^2 \langle H \rangle^2}{M_{\text{heavy}}} \rightarrow \text{SMALL}$$

coupling $\rightarrow \text{SMALL}$



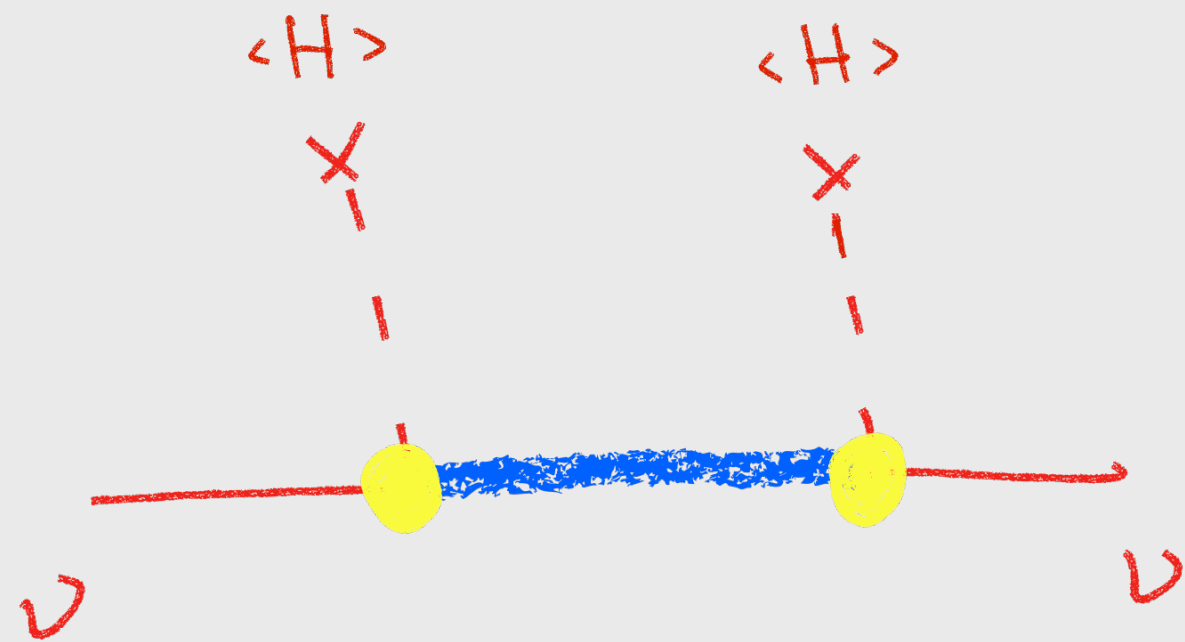
$$m_\nu = \mu \cdot \frac{(\text{coupling})^2 \langle H \rangle^2}{M_{\text{heavy}}^2} \rightarrow \text{SMALL}$$

$\mu \rightarrow \text{SMALL}$

Neutrino mass mechanisms

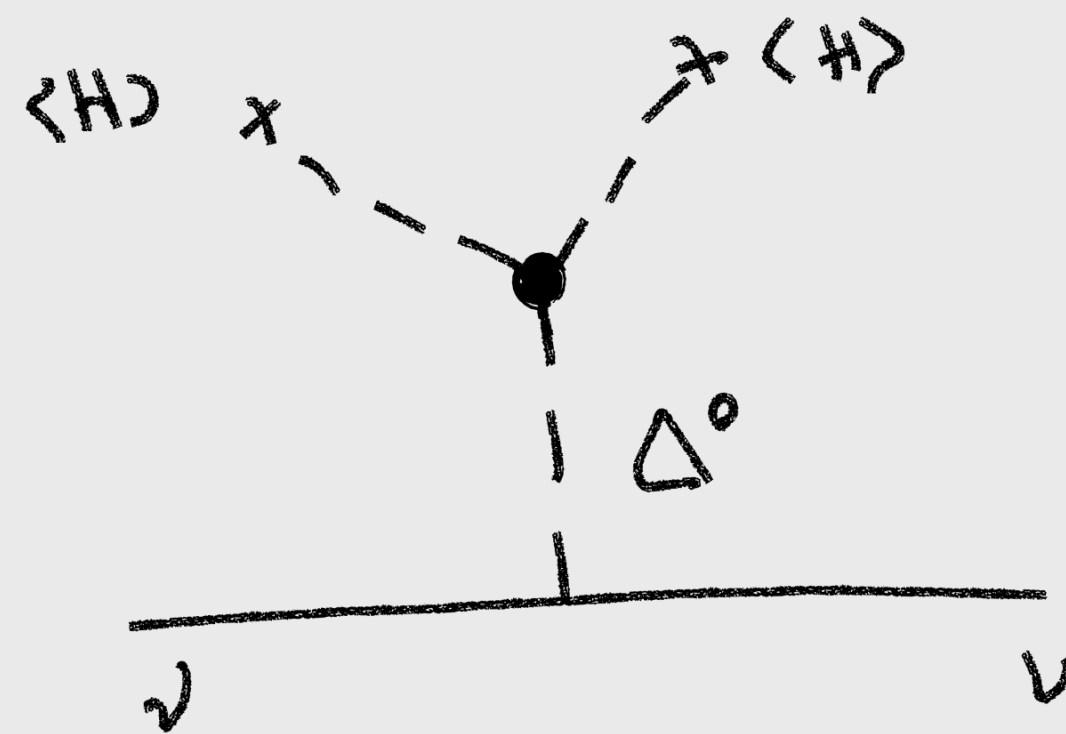
LEPTON

NUMBER BREAKING



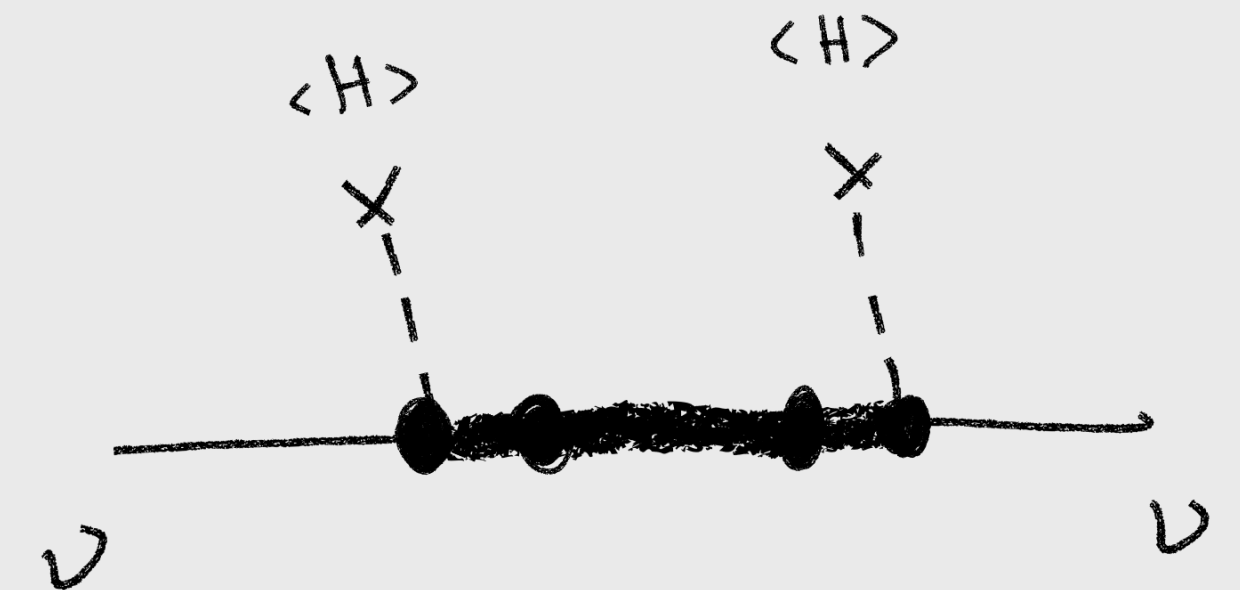
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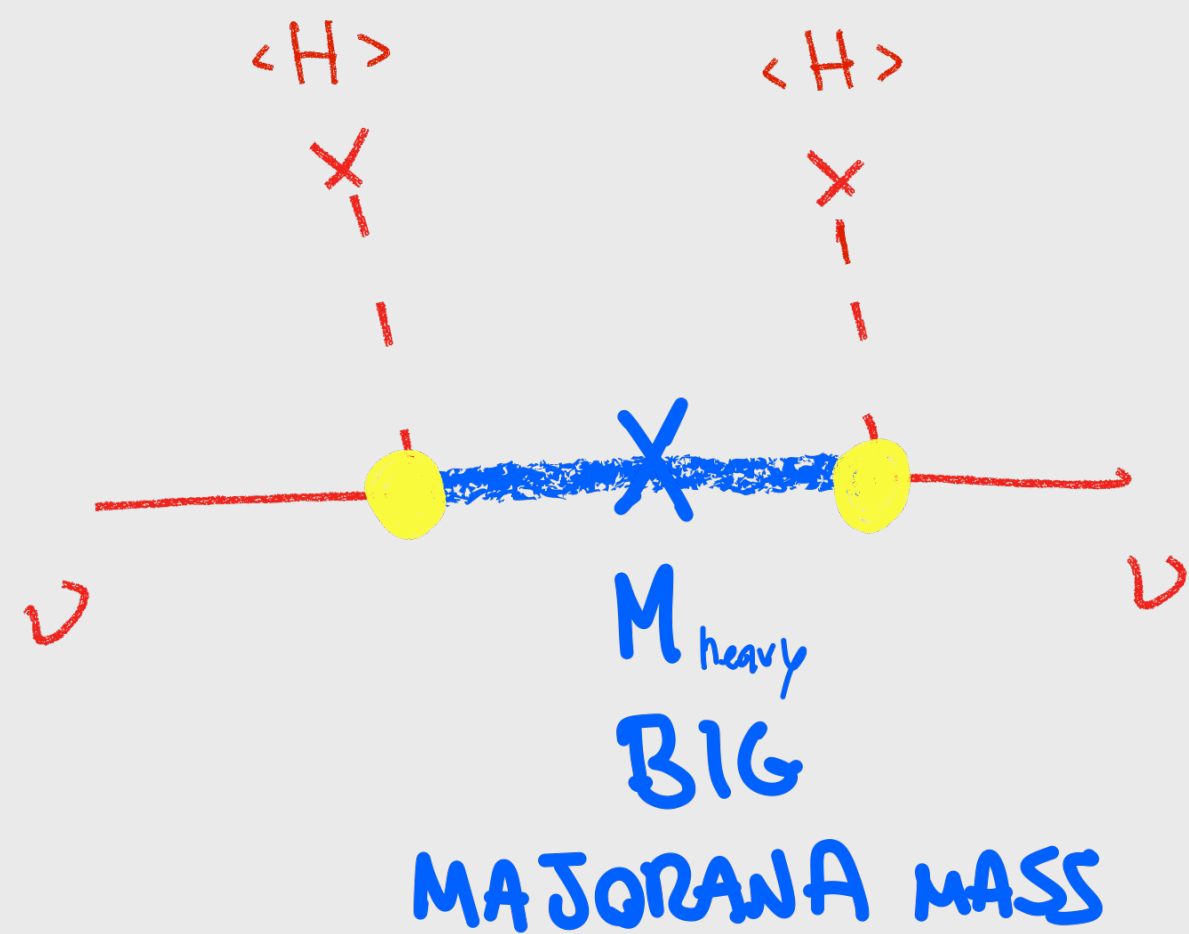
$$m_\nu = \mu \cdot \frac{(\text{coupling})^2 \langle H \rangle^2}{M_{\text{heavy}}^2} \rightarrow \text{SMALL}$$

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Neutrino mass mechanisms

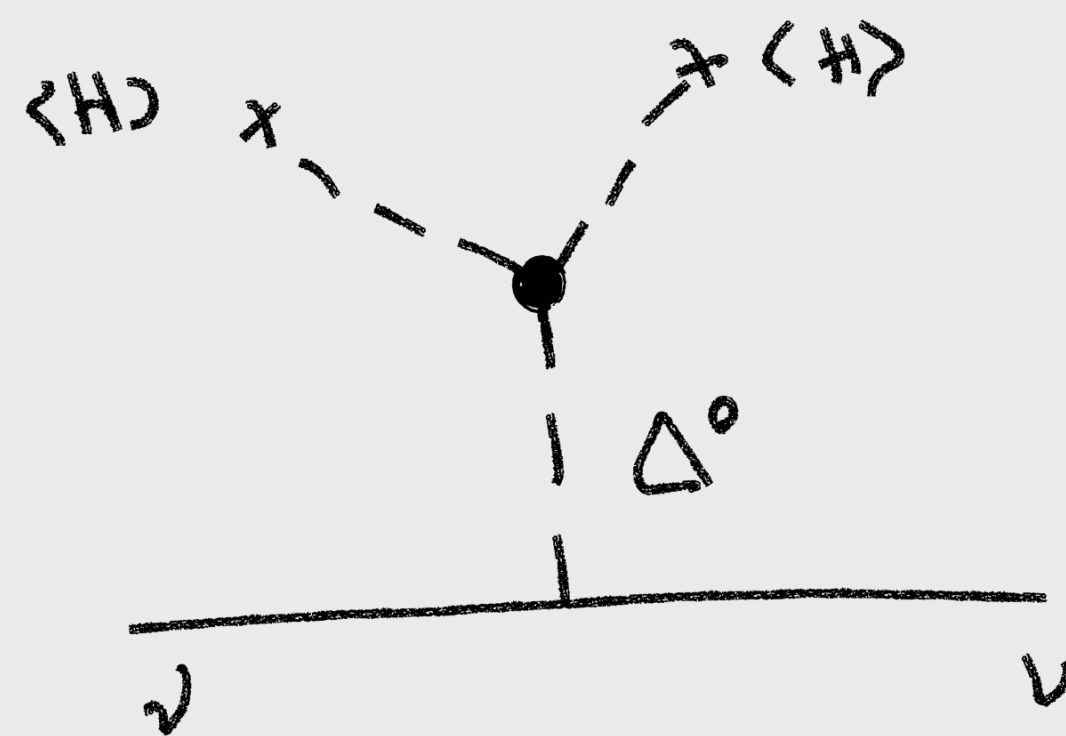
LEPTON

NUMBER BREAKING



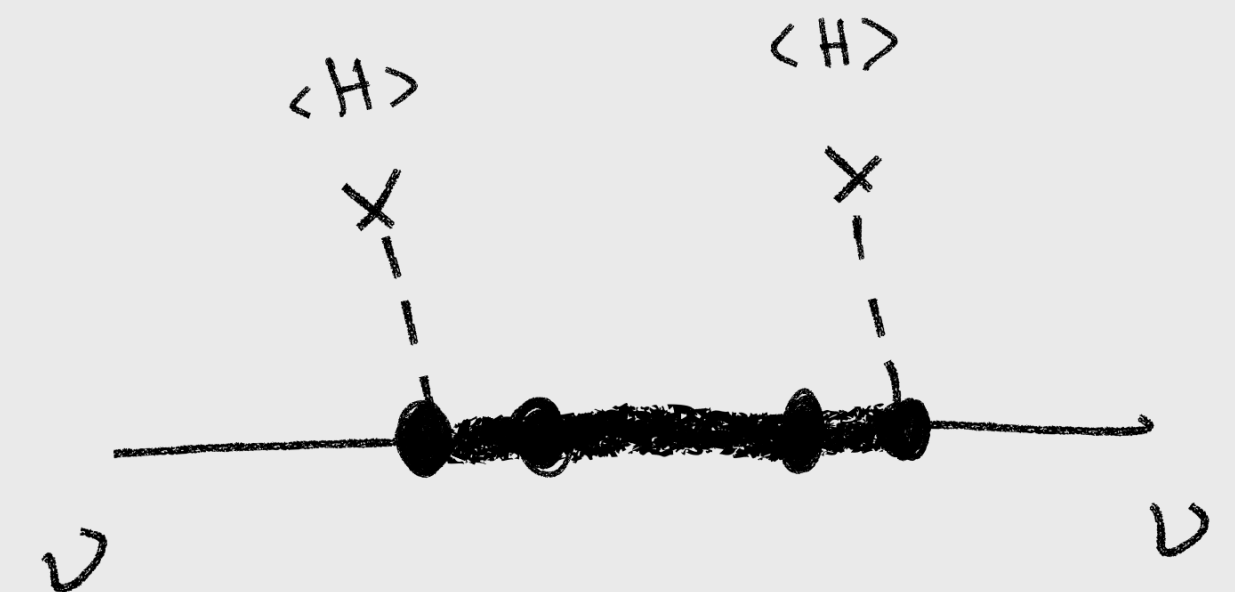
$$m_\nu = \frac{(coupling)^2 \langle H \rangle^2}{M_{heavy}} \rightarrow \text{SMALL}$$

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$$m_\nu = \frac{(coupling)^2 \langle H \rangle^2}{M_{heavy}} \rightarrow \text{SMALL}$$

coupling \rightarrow SMALL



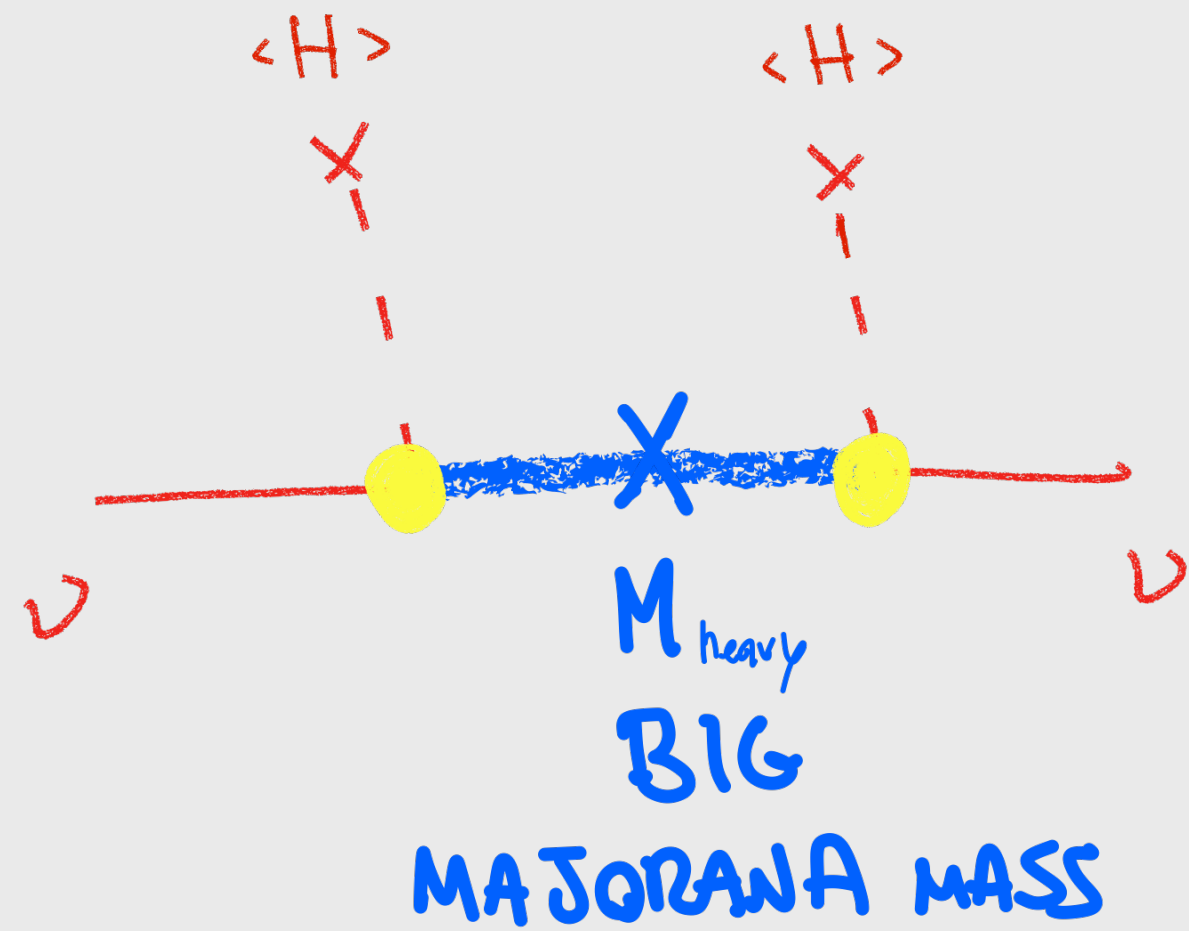
$$m_\nu = \mu \cdot \frac{(coupling)^2 \langle H \rangle^2}{M_{heavy}^2} \rightarrow \text{SMALL}$$

$\mu \rightarrow \text{SMALL}$

Neutrino mass mechanisms

LEPTON

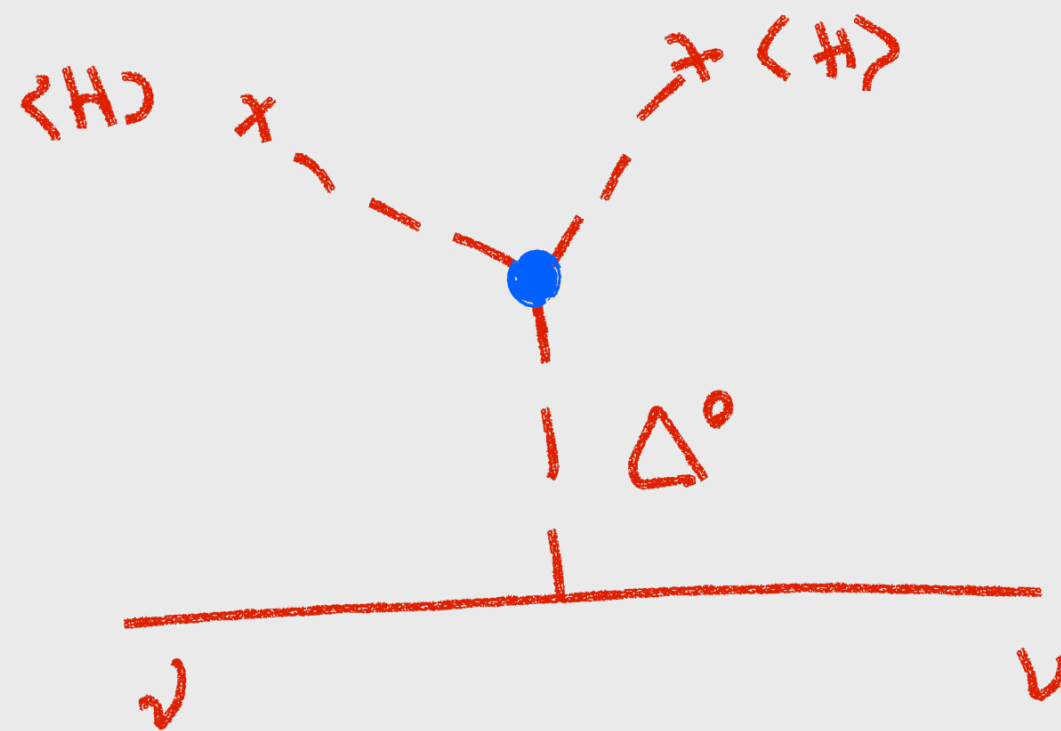
NUMBER BREAKING



$$m_\nu = \frac{(coupling)^2 \langle H \rangle^2}{M_{heavy}} \rightarrow \text{SMALL}$$

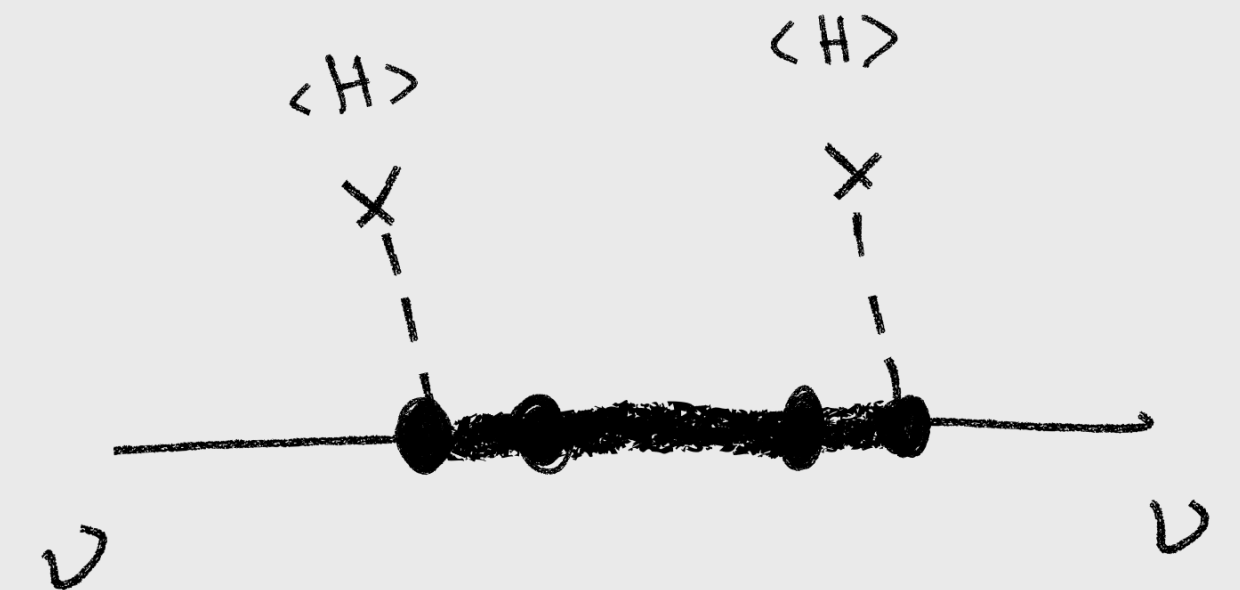
$M_{heavy} \rightarrow \text{LARGE}$

SMALL COUPLING



$$m_\nu = \frac{(coupling)^2 \langle H \rangle^2}{M_{heavy}} \rightarrow \text{SMALL}$$

coupling $\rightarrow \text{SMALL}$



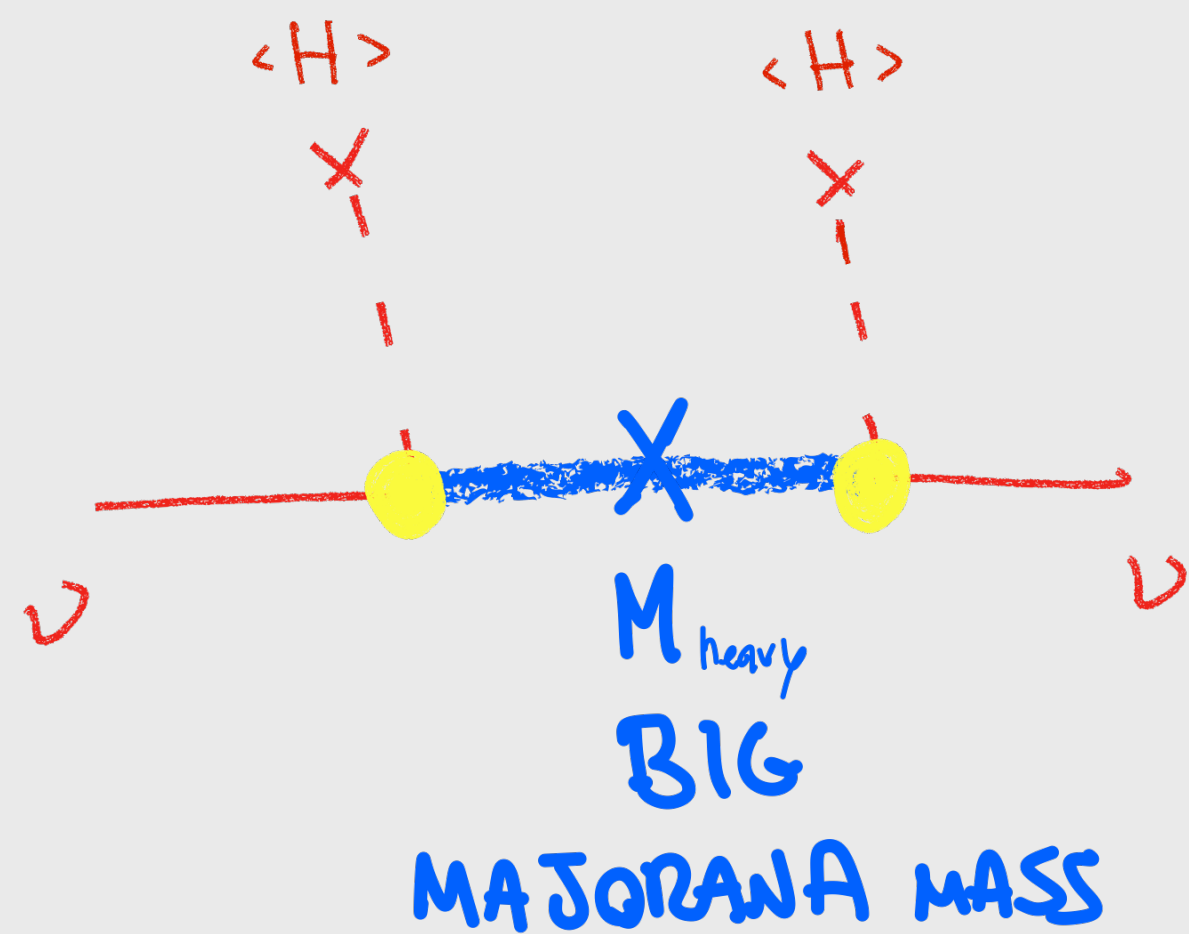
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$\mu \rightarrow \text{SMALL}$

Neutrino mass mechanisms

LEPTON

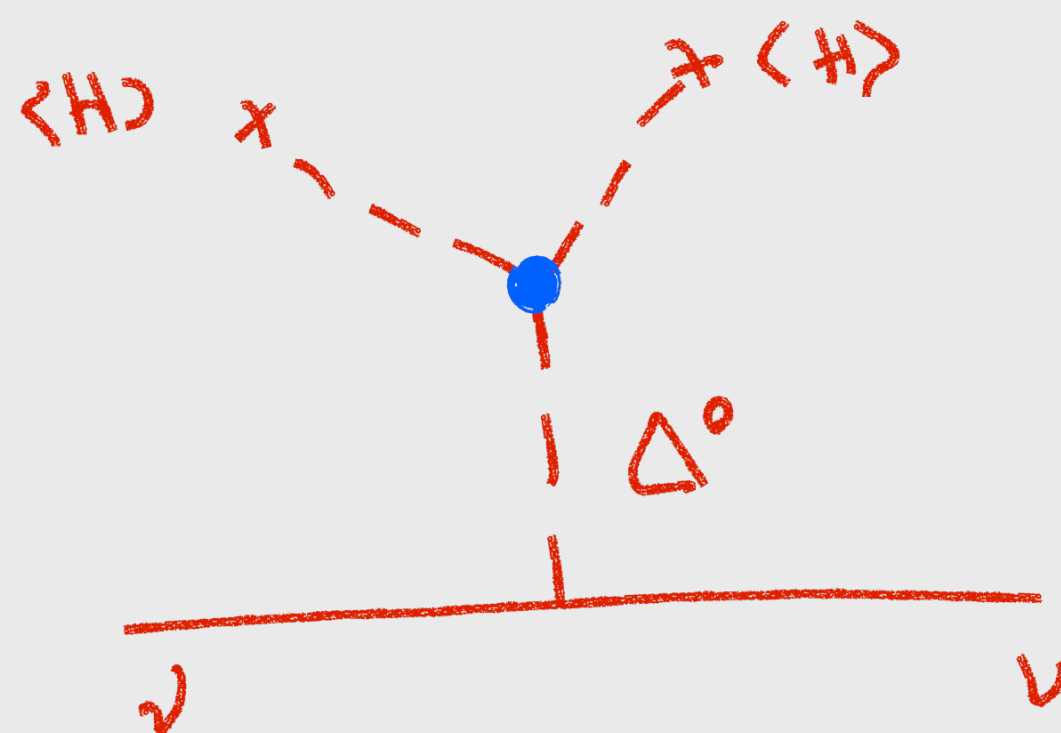
NUMBER BREAKING



$$m_\nu = \frac{(coupling)^2 \langle H \rangle^2}{M_{heavy}} \rightarrow \text{SMALL}$$

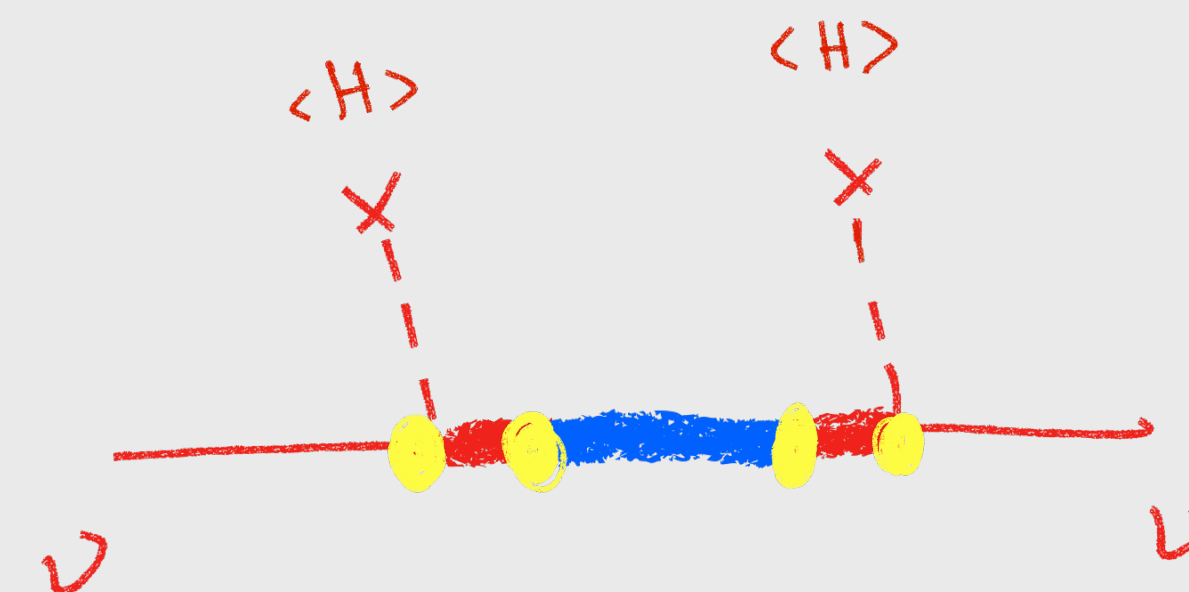
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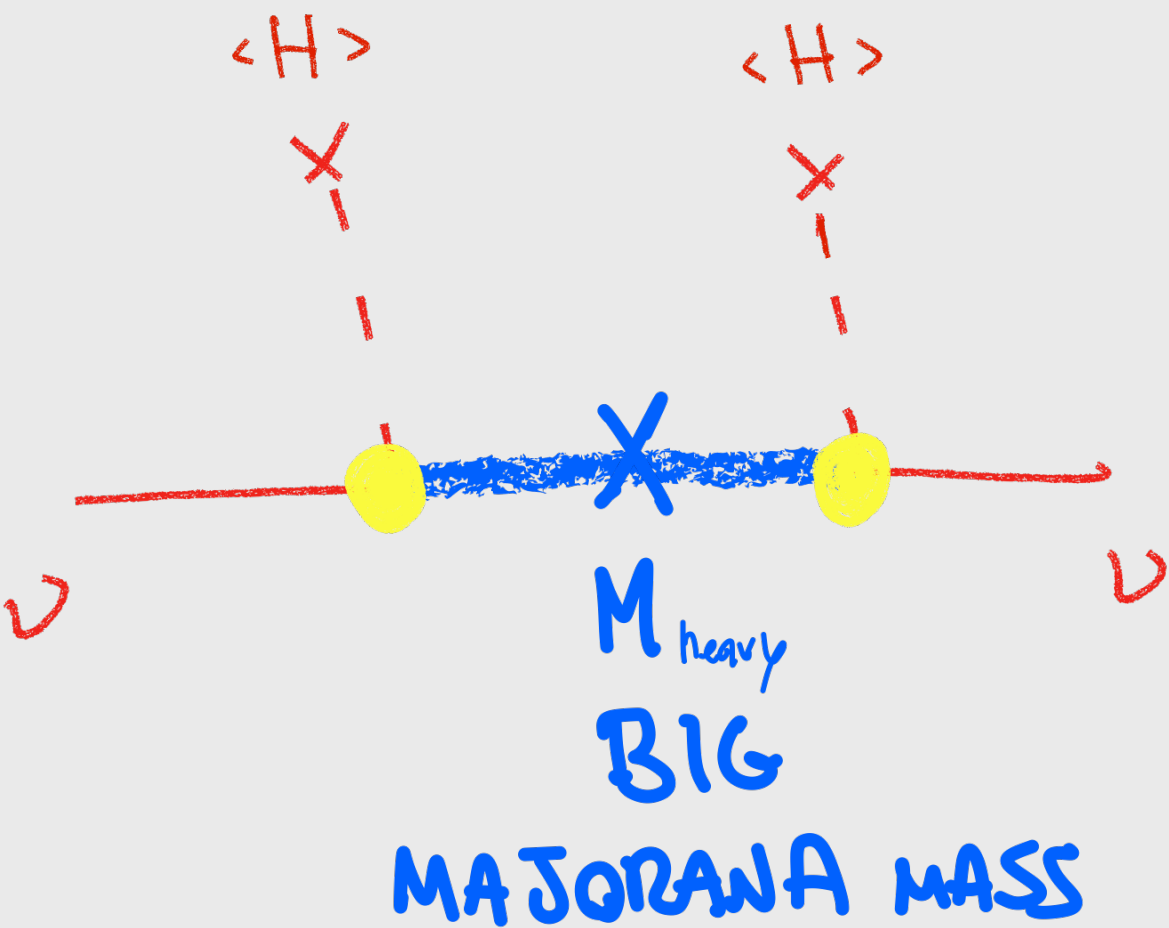
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Neutrino mass mechanisms

LEPTON

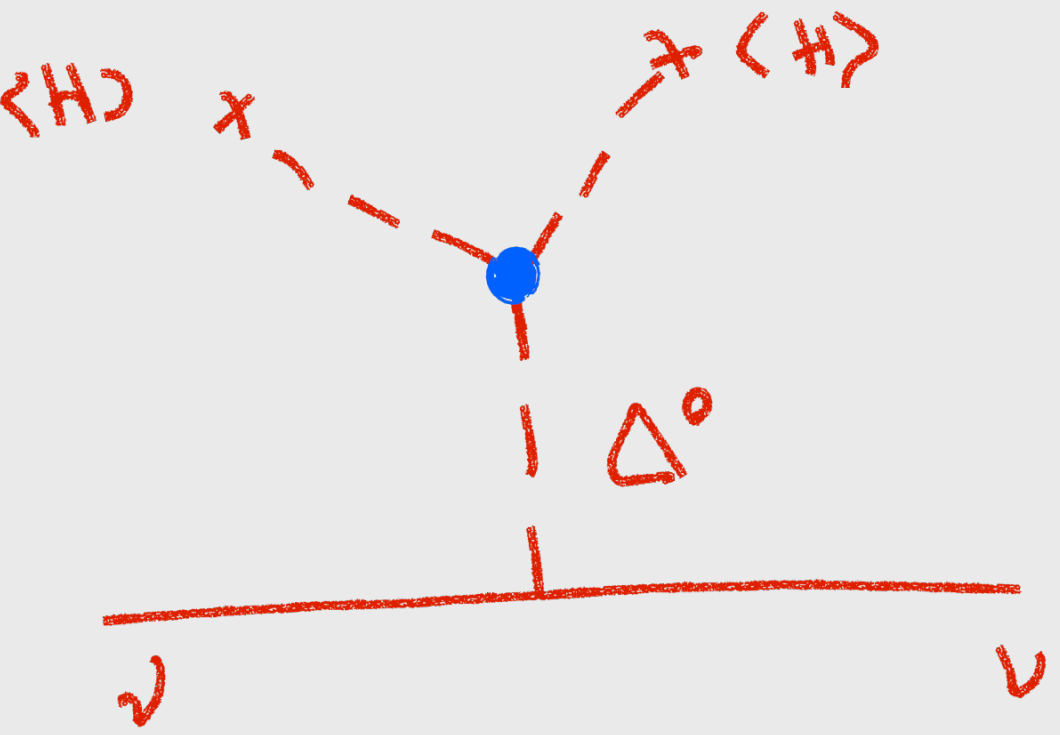
NUMBER BREAKING



$$m_\nu = \frac{(\text{coupling})^2 \langle H \rangle^2}{M_{\text{heavy}}} \rightarrow \text{SMALL}$$

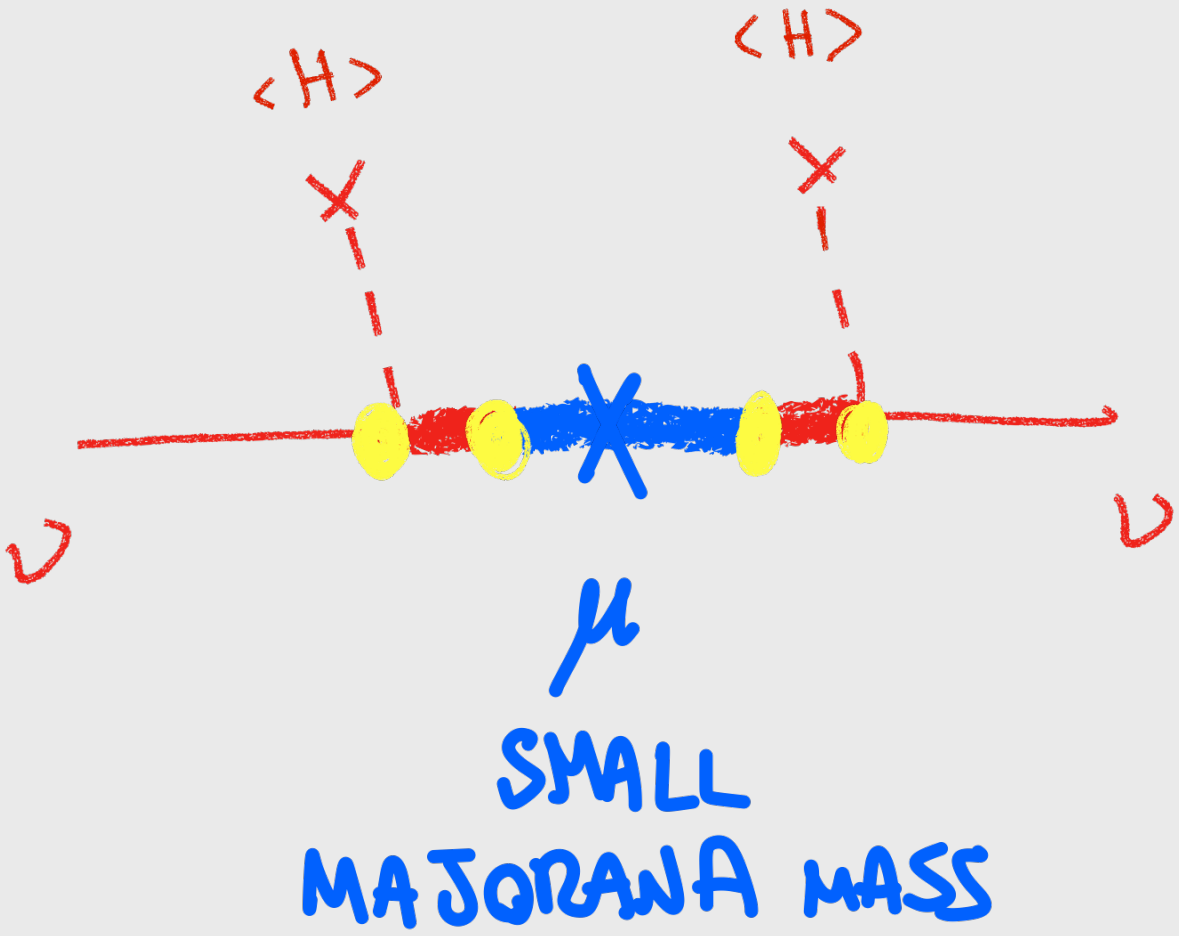
$M_{\text{heavy}} \rightarrow \text{LARGE}$

SMALL COUPLING



$$m_\nu = \frac{(\text{coupling})^2 \langle H \rangle^2}{M_{\text{heavy}}} \rightarrow \text{SMALL}$$

coupling $\rightarrow \text{SMALL}$



$$m_\nu = \mu \cdot \frac{(\text{coupling})^2 \langle H \rangle^2}{M_{\text{heavy}}^2} \rightarrow \text{SMALL}$$

$\mu \rightarrow \text{SMALL}$

Neutrino mass

LEPTON

NUMBER BREAKING

Mass

10 TeV

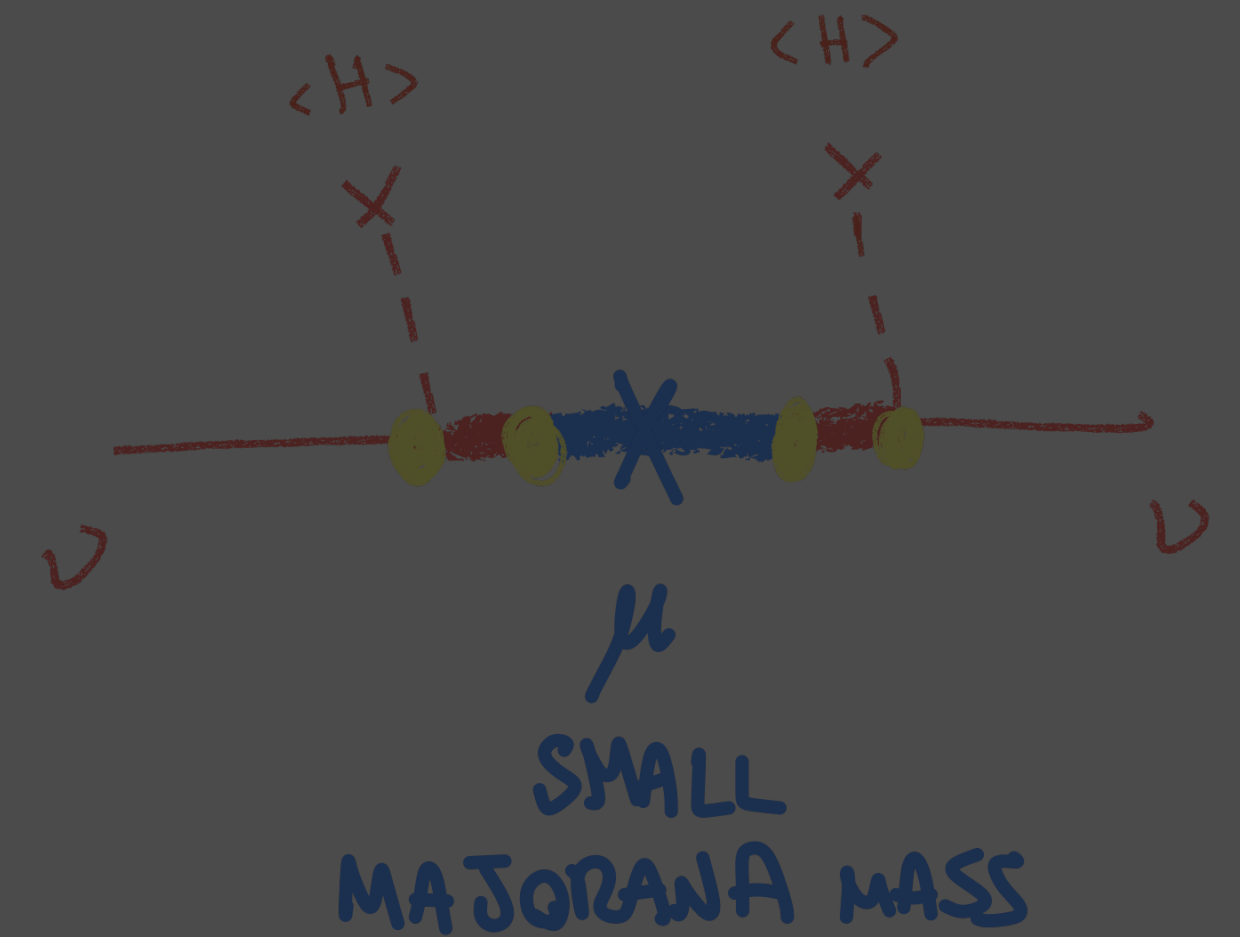
1 TeV

LEPTON
NUMBER
VIOLATION



$$m_\nu = \frac{(\text{coupling})^2 \langle H \rangle^2}{M_{\text{heavy}}} \rightarrow \text{SMALL}$$

$M_{\text{heavy}} \rightarrow \text{LARGE}$



$$m_\nu = \mu \cdot \frac{(\text{coupling})^2 \langle H \rangle^2}{M_{\text{heavy}}^2} \rightarrow \text{SMALL}$$

$\mu \rightarrow \text{SMALL}$

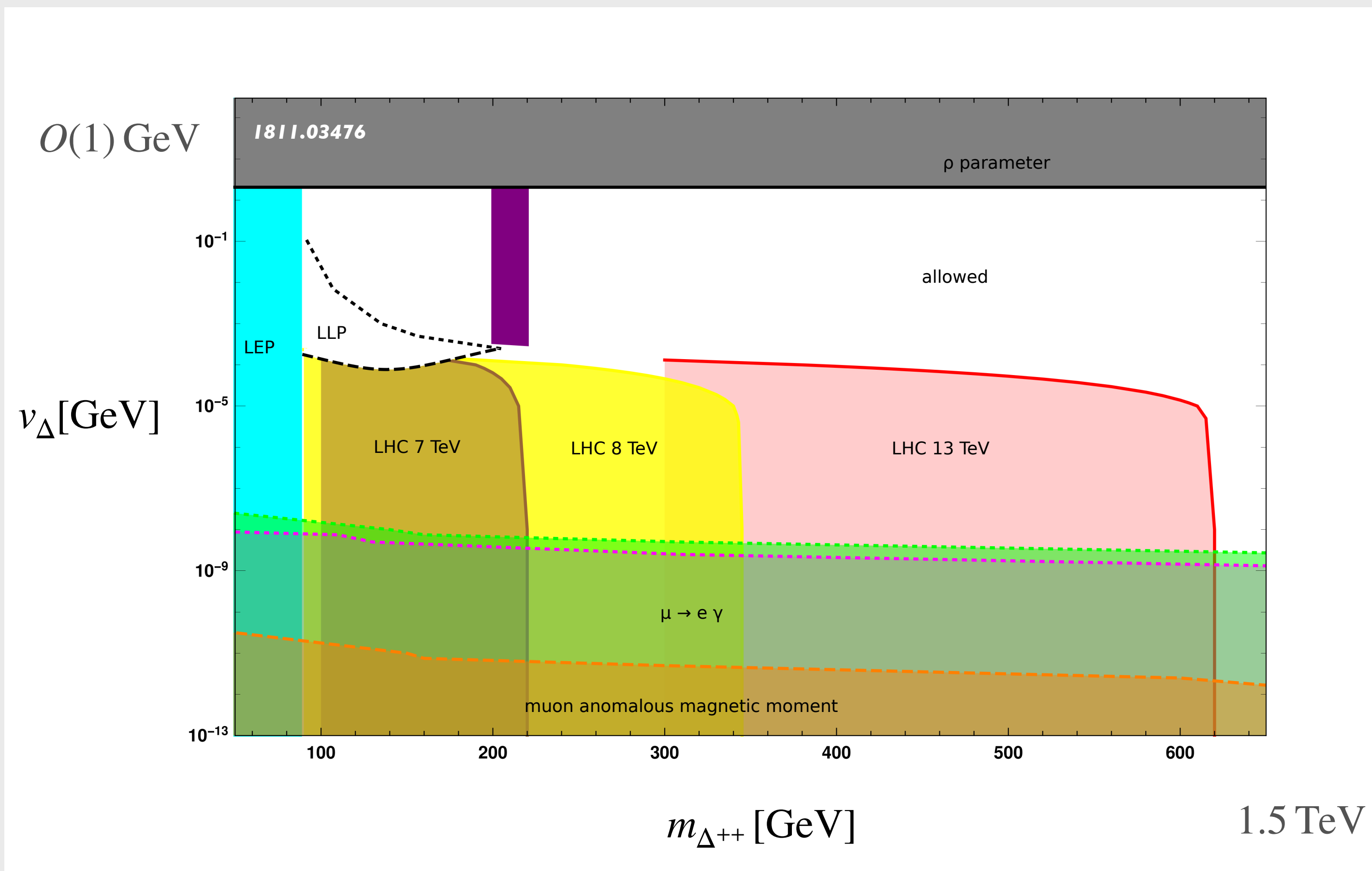
mechanisms

flashing concrete results for

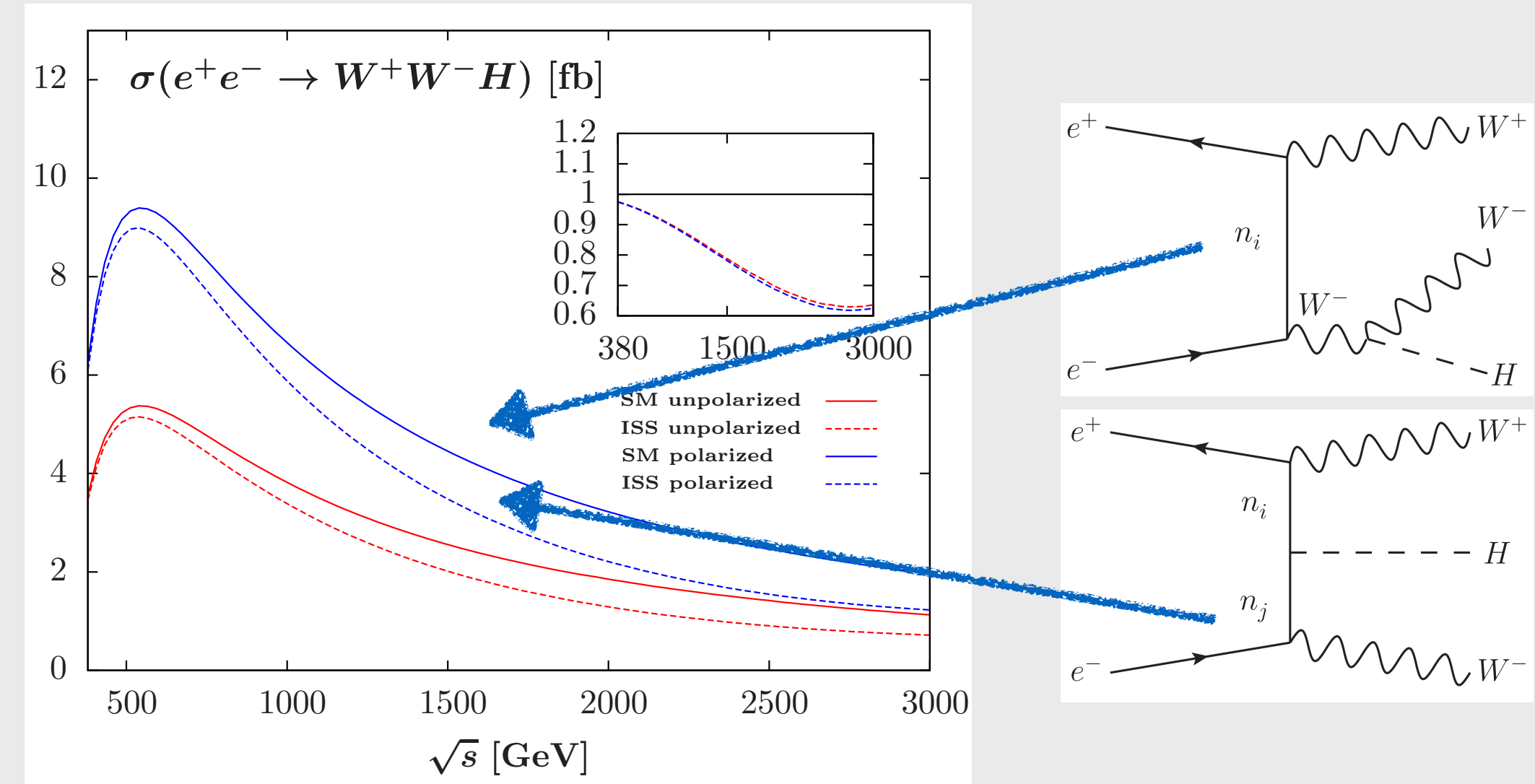
The origin of neutrino masses

Plenty of neutrino mass models in reach

Type-2 See-Saw 1803.00677 - Agrawal, Mitra, Niyogi, Shil, Spannowsky

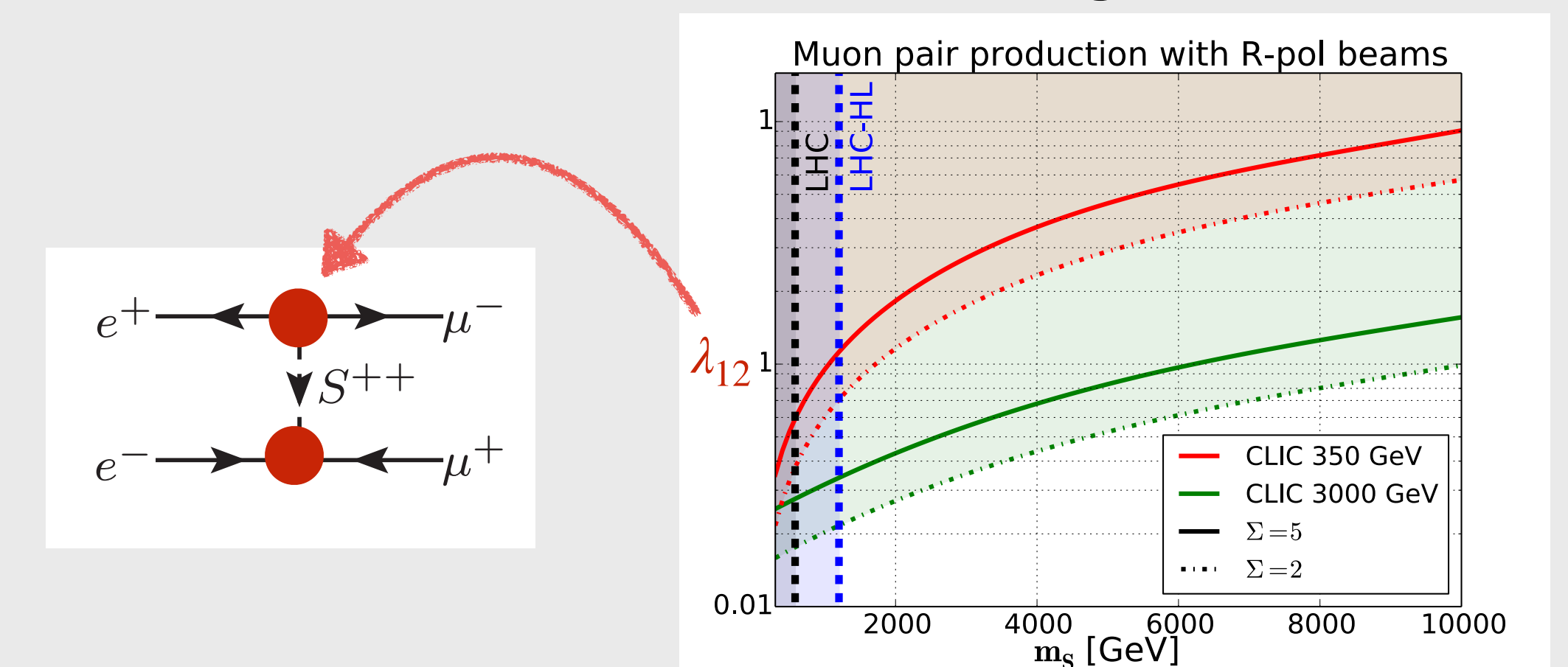


Inverse See-Saw 1712.07621 - Baglio, Pascoli, Weiland



Exclude ISS RH Neutrino up to 10 TeV for Yukawa ~ 1

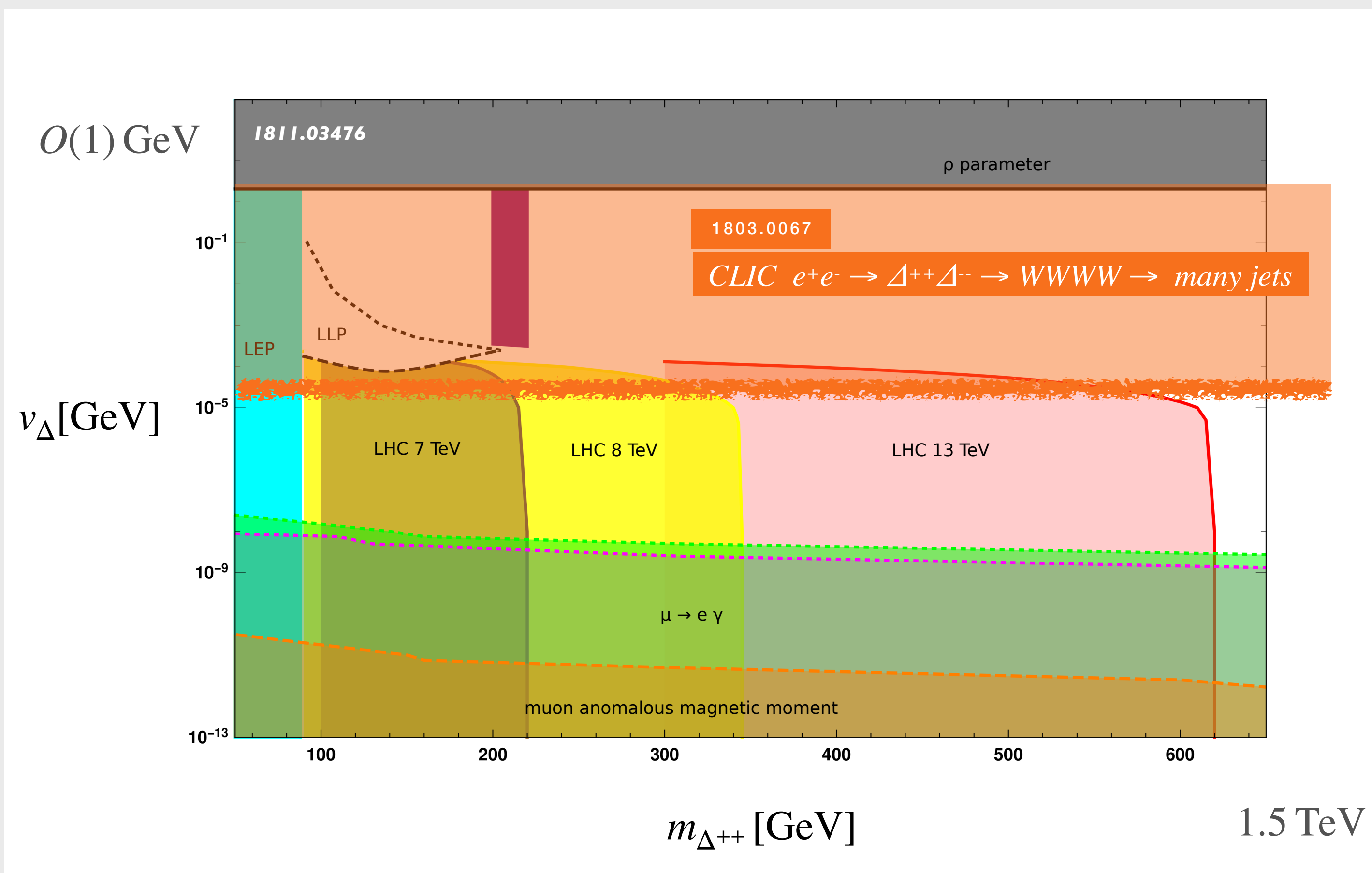
1807.10224 - Crivellin, Ghezzi, Panizzi, Pruna, Signer



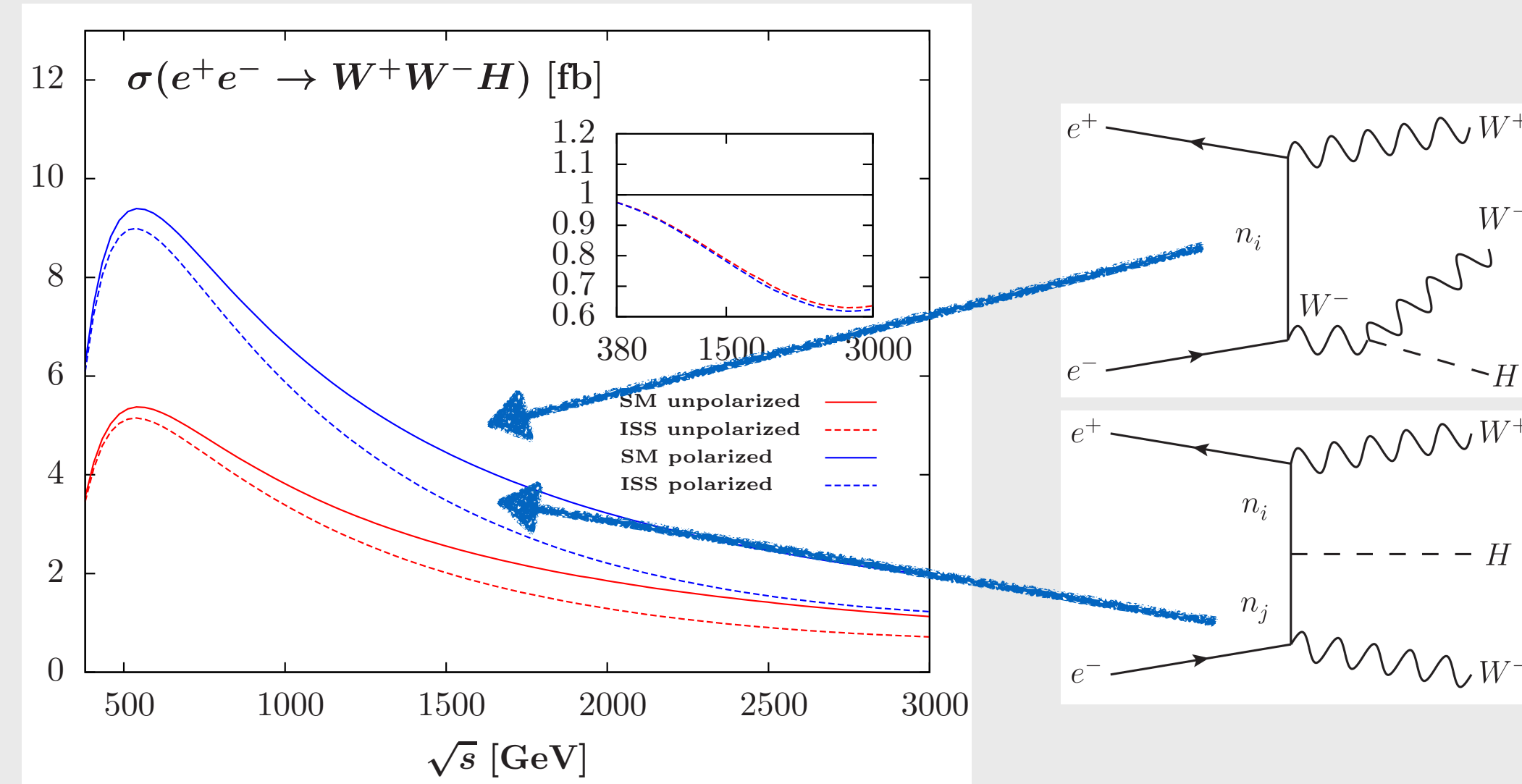
Exclude S^{++} up to 10 TeV for triplet Yukawa ~ 0.1

Plenty of neutrino mass models in reach

Type-2 See-Saw 1803.00677 - Agrawal, Mitra, Niyogi, Shil, Spannowsky

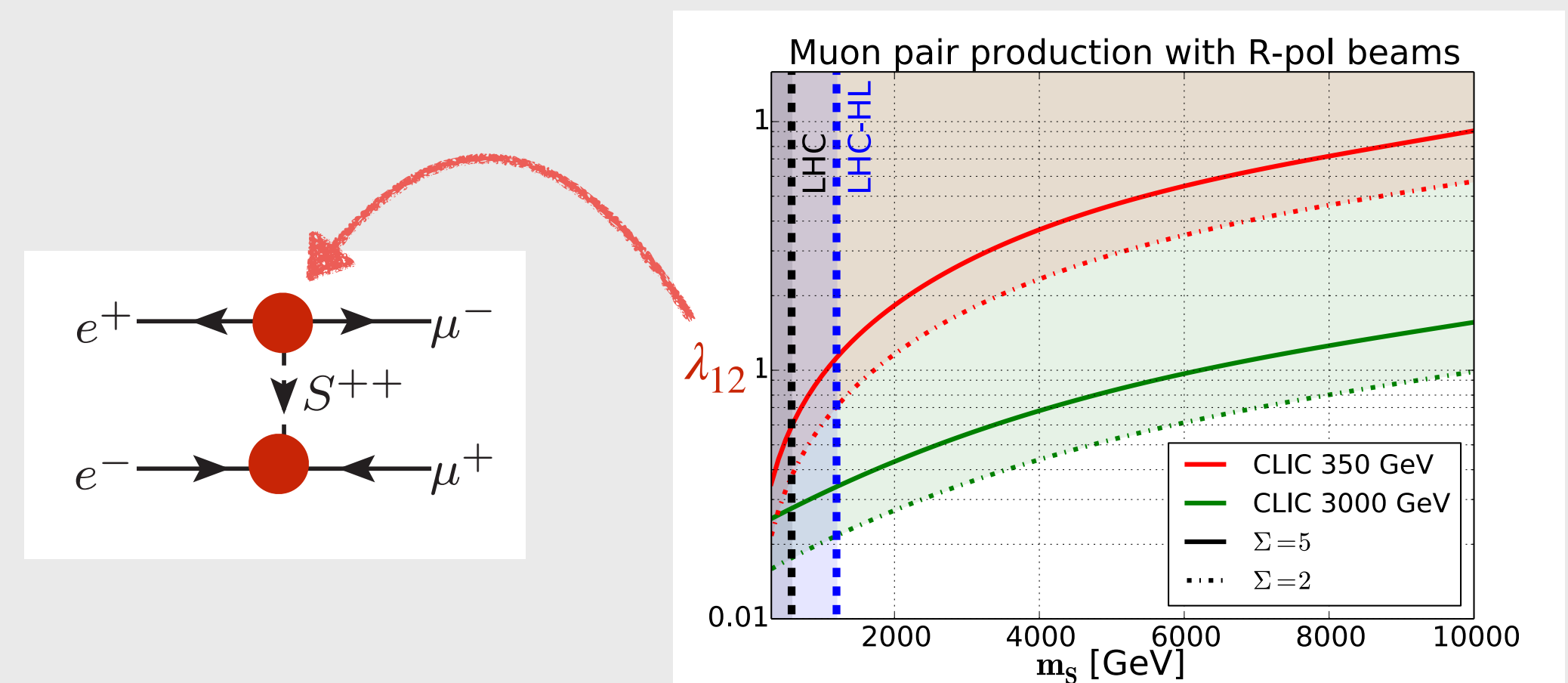


Inverse See-Saw 1712.07621 - Baglio, Pascoli, Weiland



Exclude ISS RH Neutrino up to 10 TeV for Yukawa ~ 1

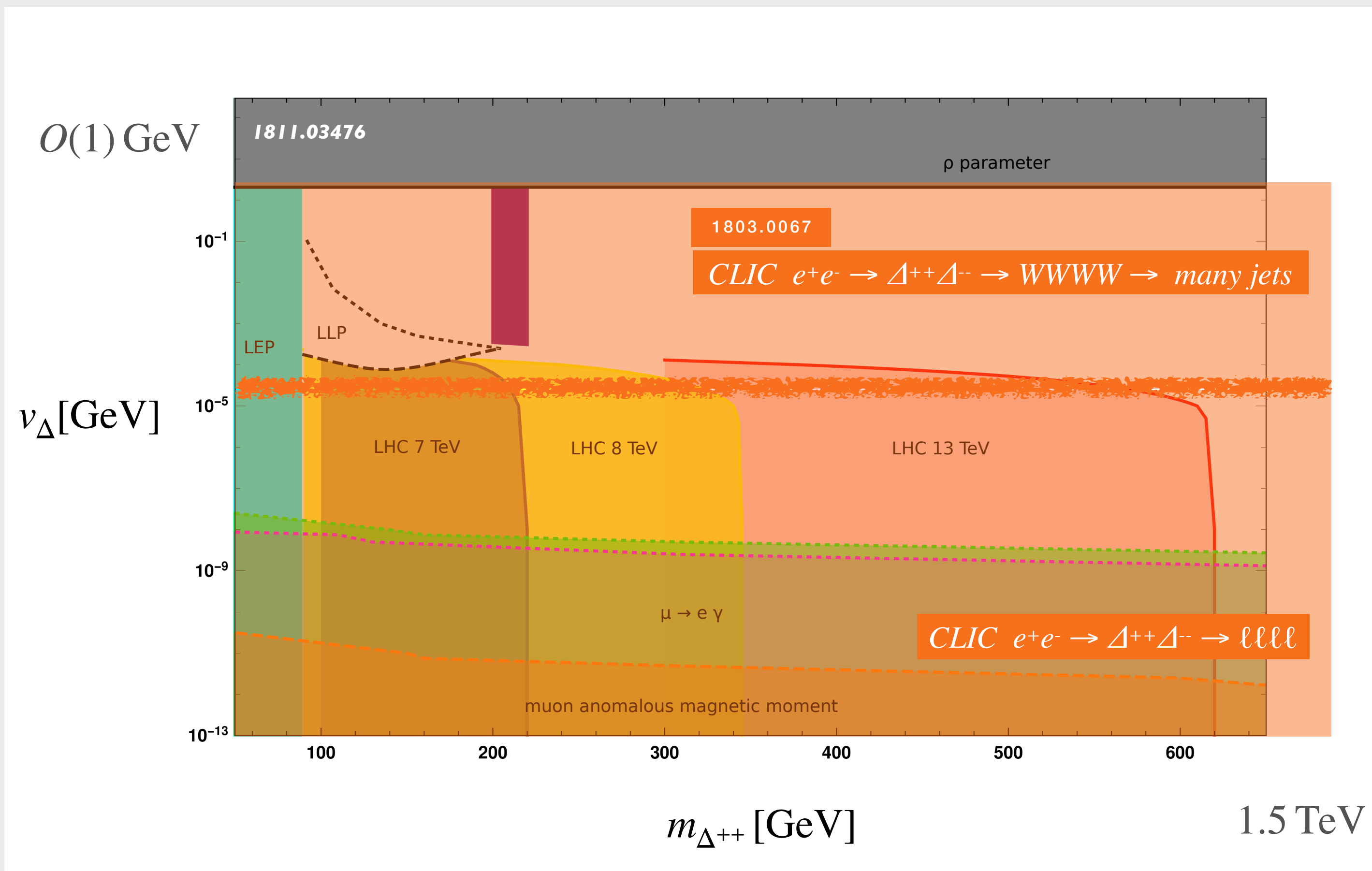
1807.10224 - Crivellin, Ghezzi, Panizzi, Pruna, Signer



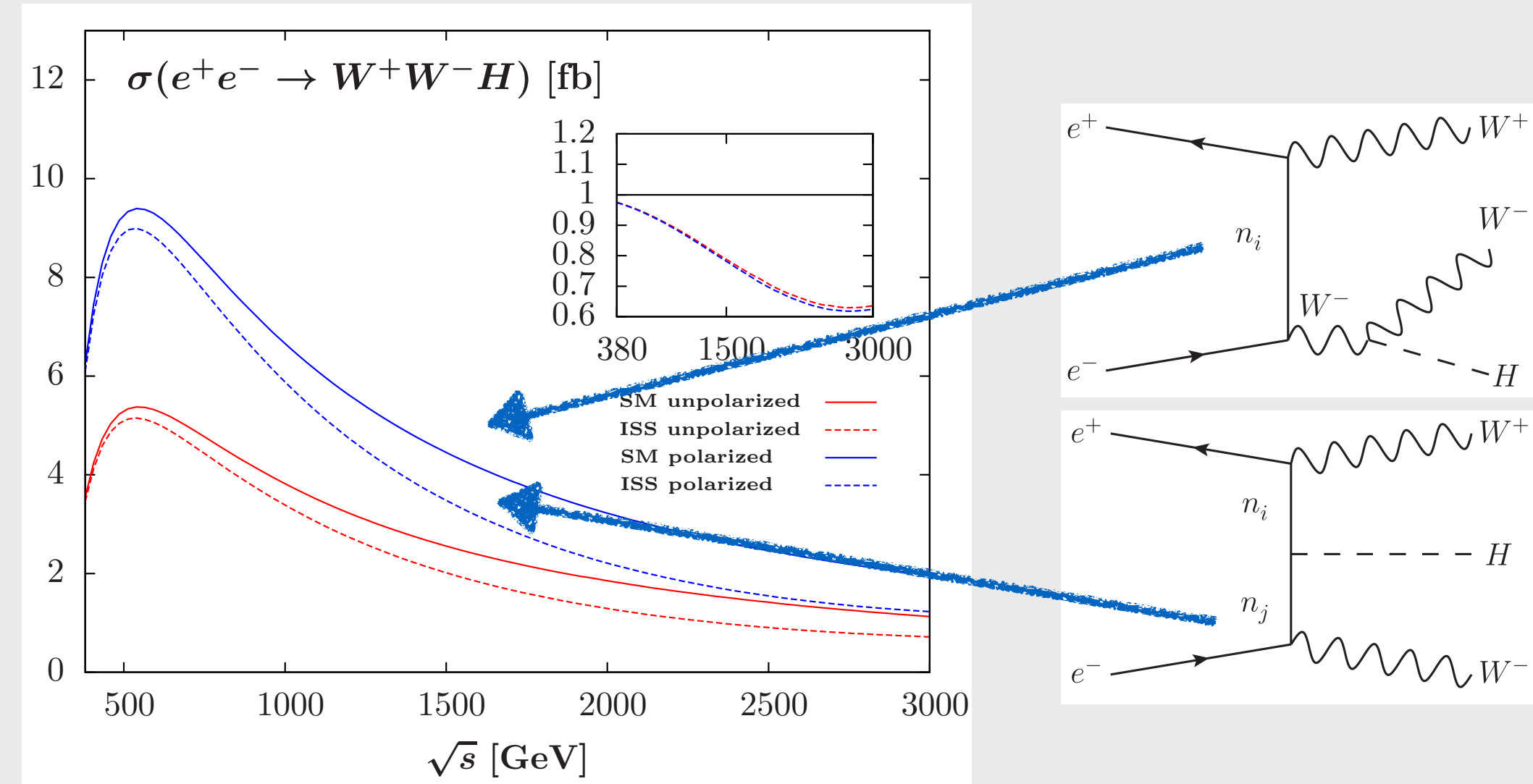
Exclude S^{++} up to 10 TeV for triplet Yukawa ~ 0.1

Plenty of neutrino mass models in reach

Type-2 See-Saw 1803.00677 - Agrawal, Mitra, Niyogi, Shil, Spannowsky

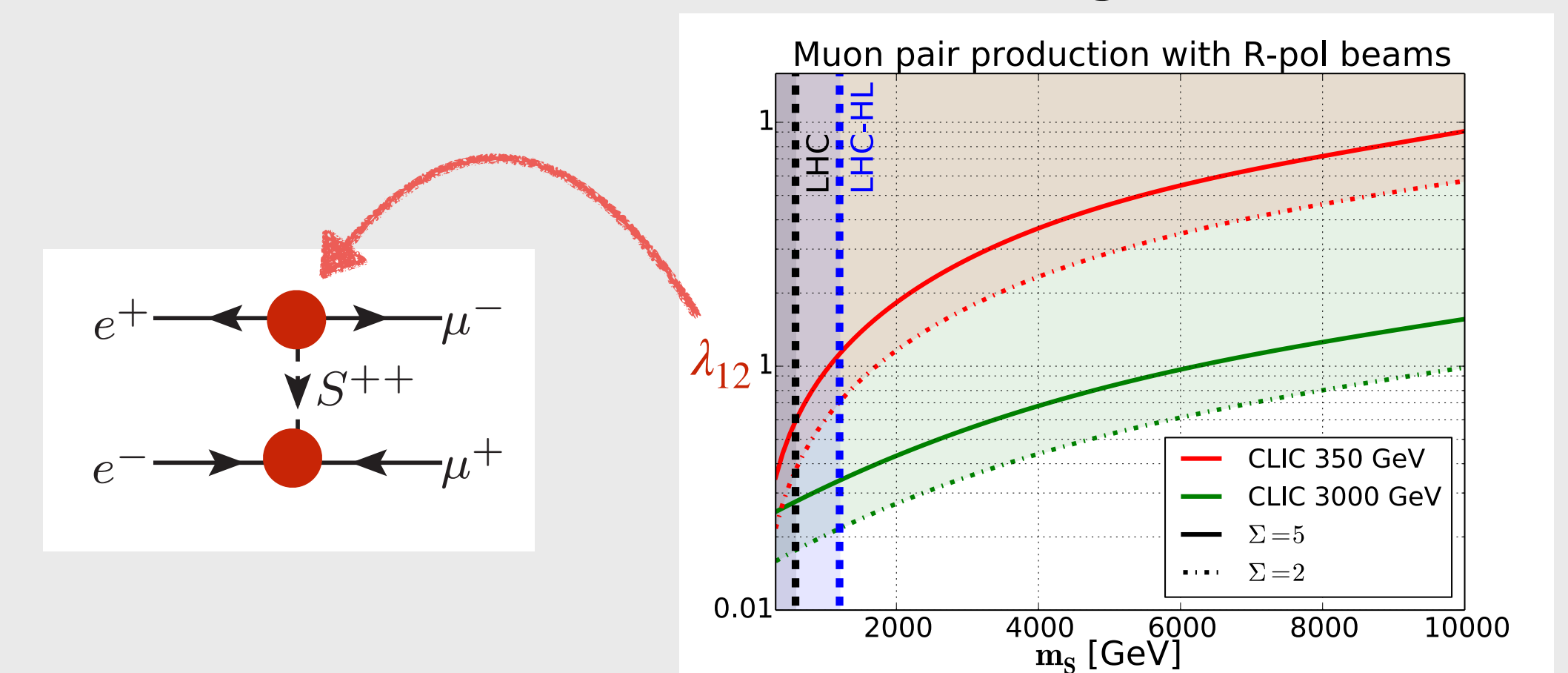


Inverse See-Saw 1712.07621 - Baglio, Pascoli, Weiland



Exclude ISS RH Neutrino up to 10 TeV for Yukawa ~ 1

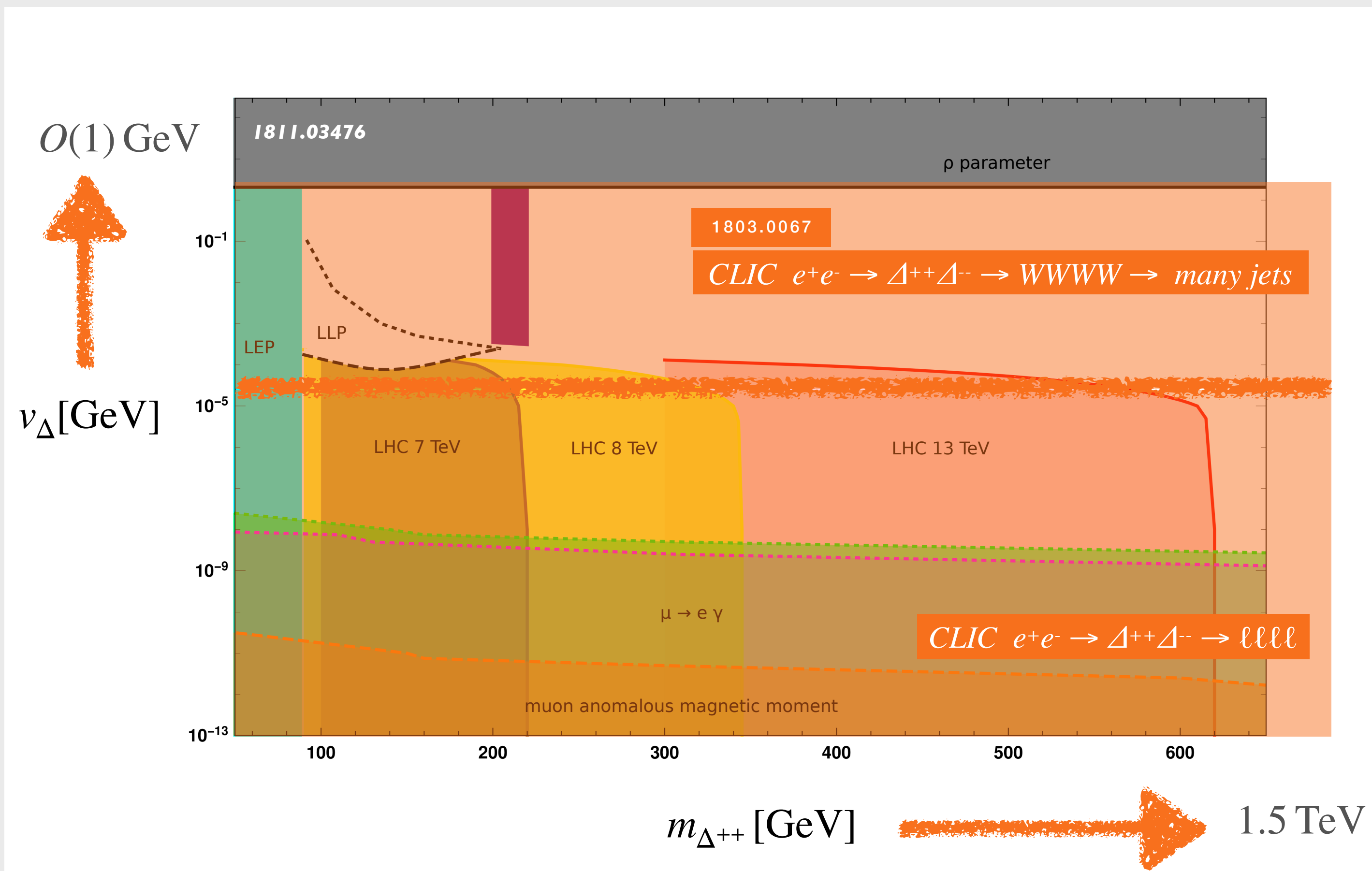
1807.10224 - Crivellin, Ghezzi, Panizzi, Pruna, Signer



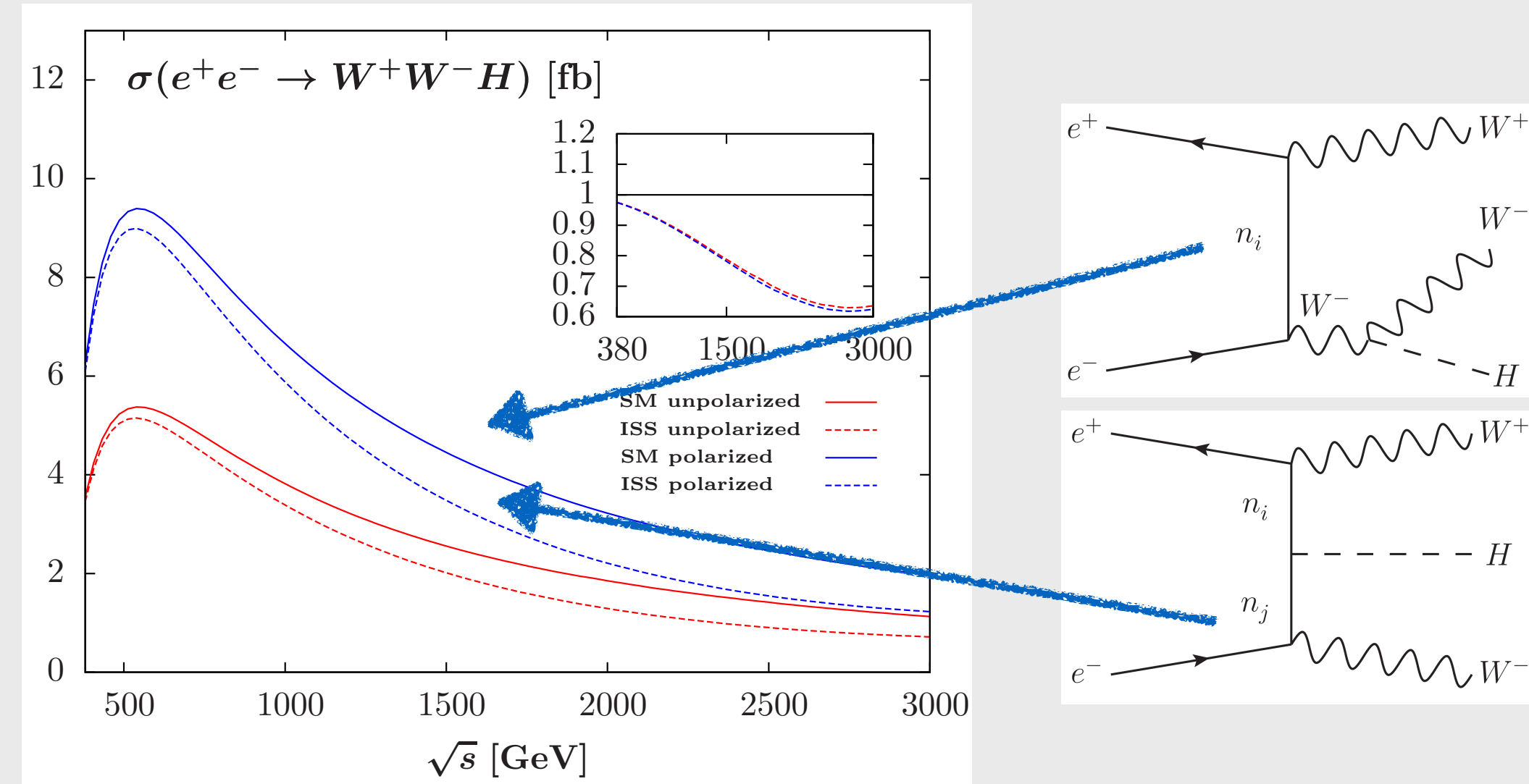
Exclude S^{++} up to 10 TeV for triplet Yukawa ~ 0.1

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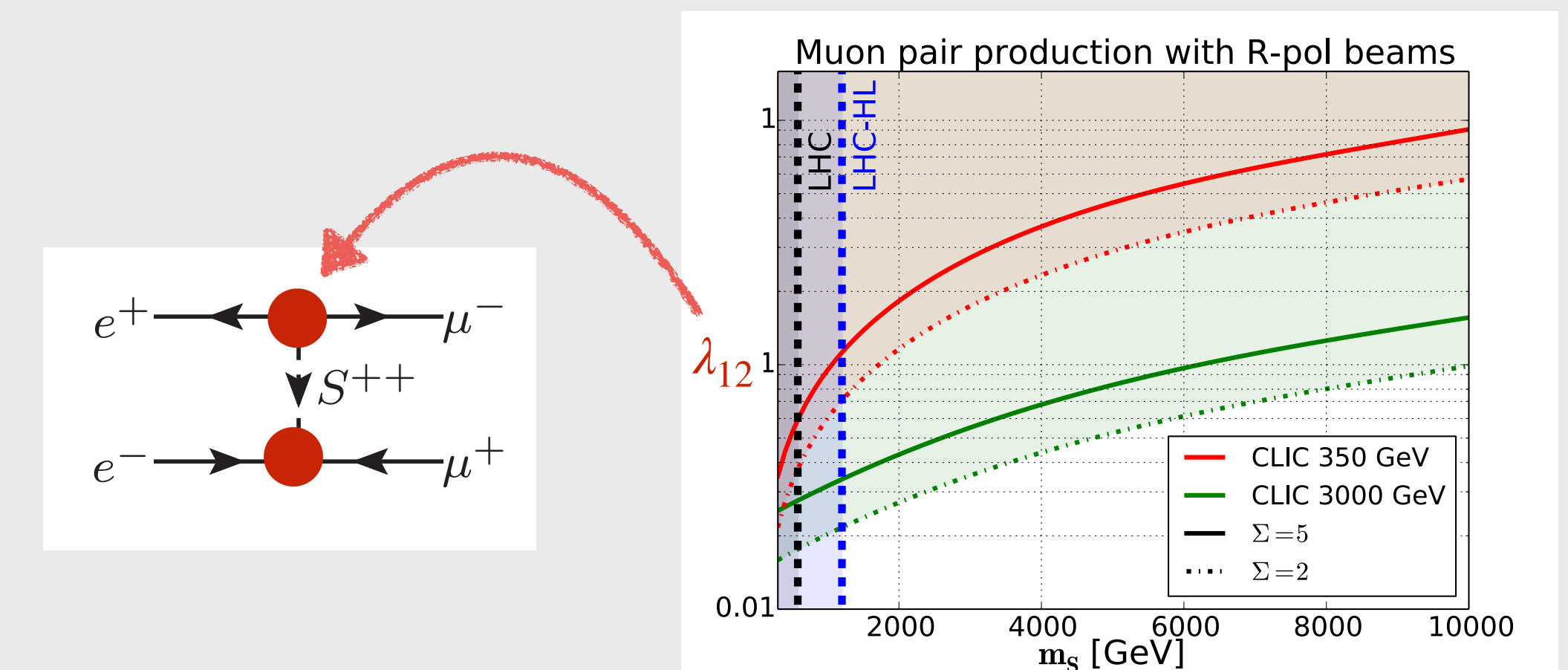


Inverse See-Saw 1712.07621 - Baglio, Pascoli, Weiland



Exclude ISS RH Neutrino up to 10 TeV for Yukawa ~ 1

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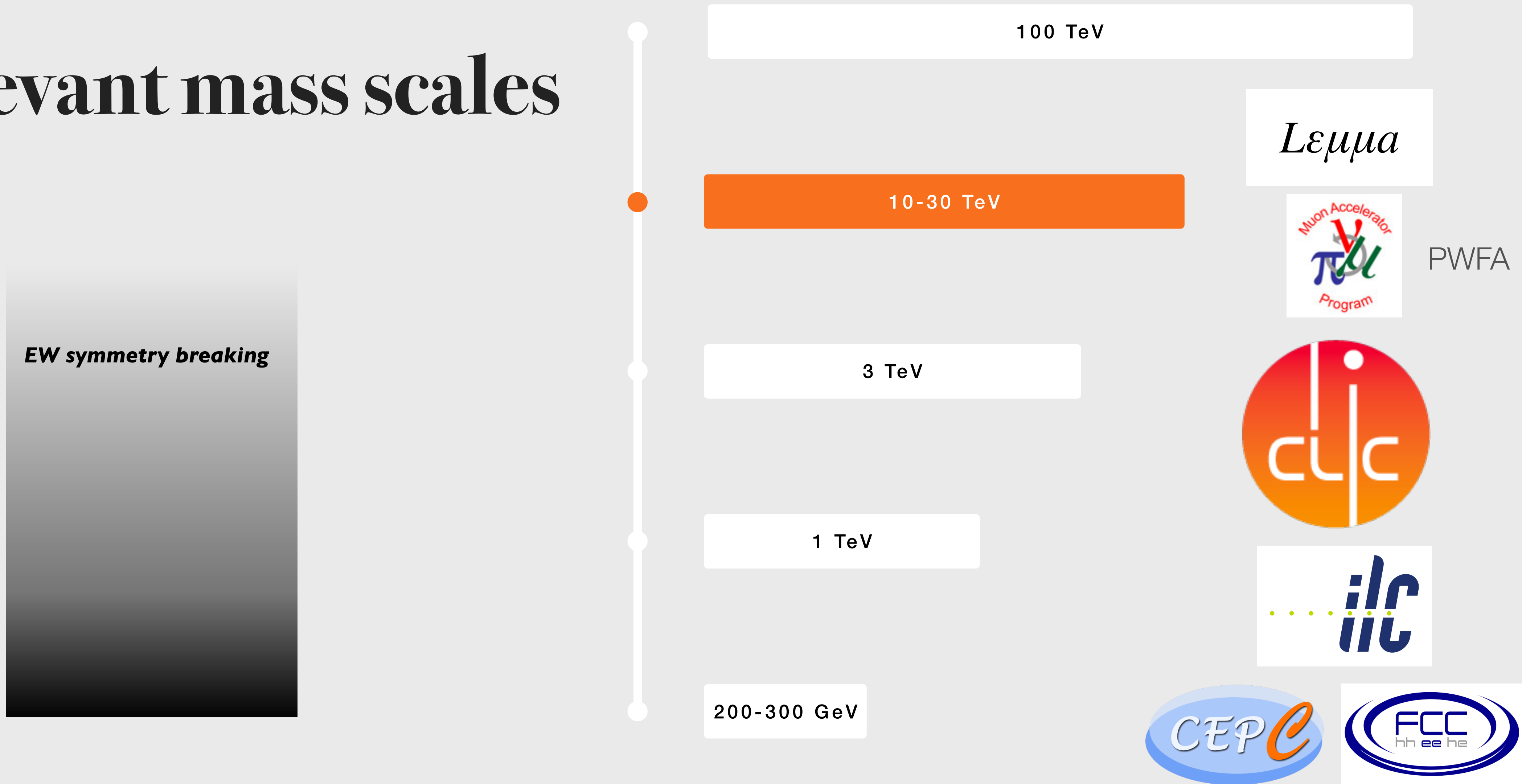


Exclude S^{++} up to 10 TeV for triplet Yukawa ~ 0.1

Relevant mass scales



Relevant mass scales



EW symmetry breaking

100 TeV

Lemna

10-30 TeV



PWFA

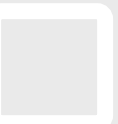
3 TeV



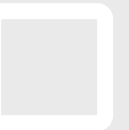
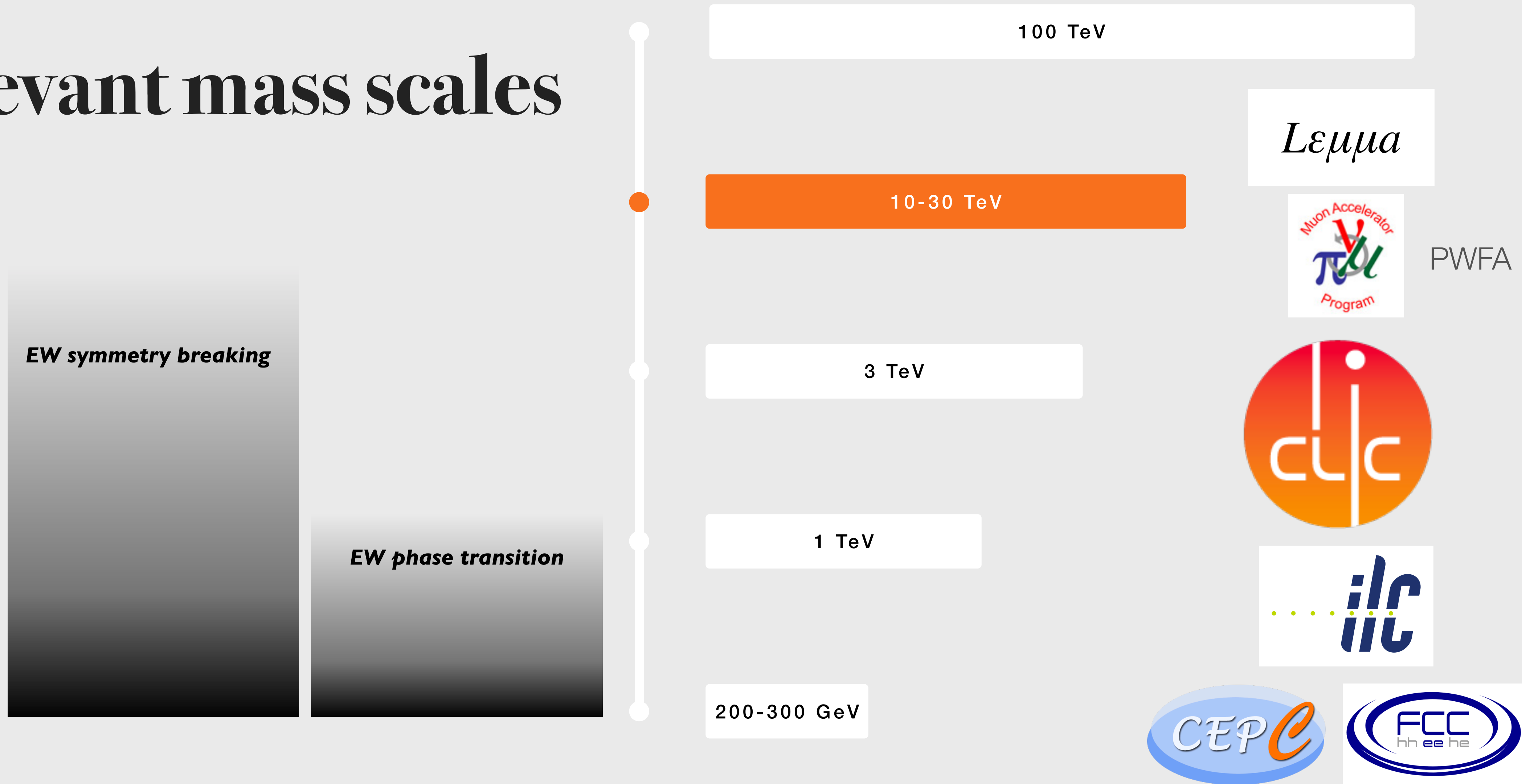
1 TeV



200-300 GeV



Relevant mass scales

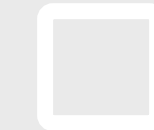


Relevant mass scales

WIMP Dark Matter

EW symmetry breaking

EW phase transition



Relevant mass scales



Relevant mass scales

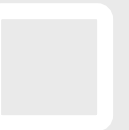
Fermions masses
and mixings



Relevant mass scales

Fermions masses
and mixings

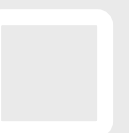
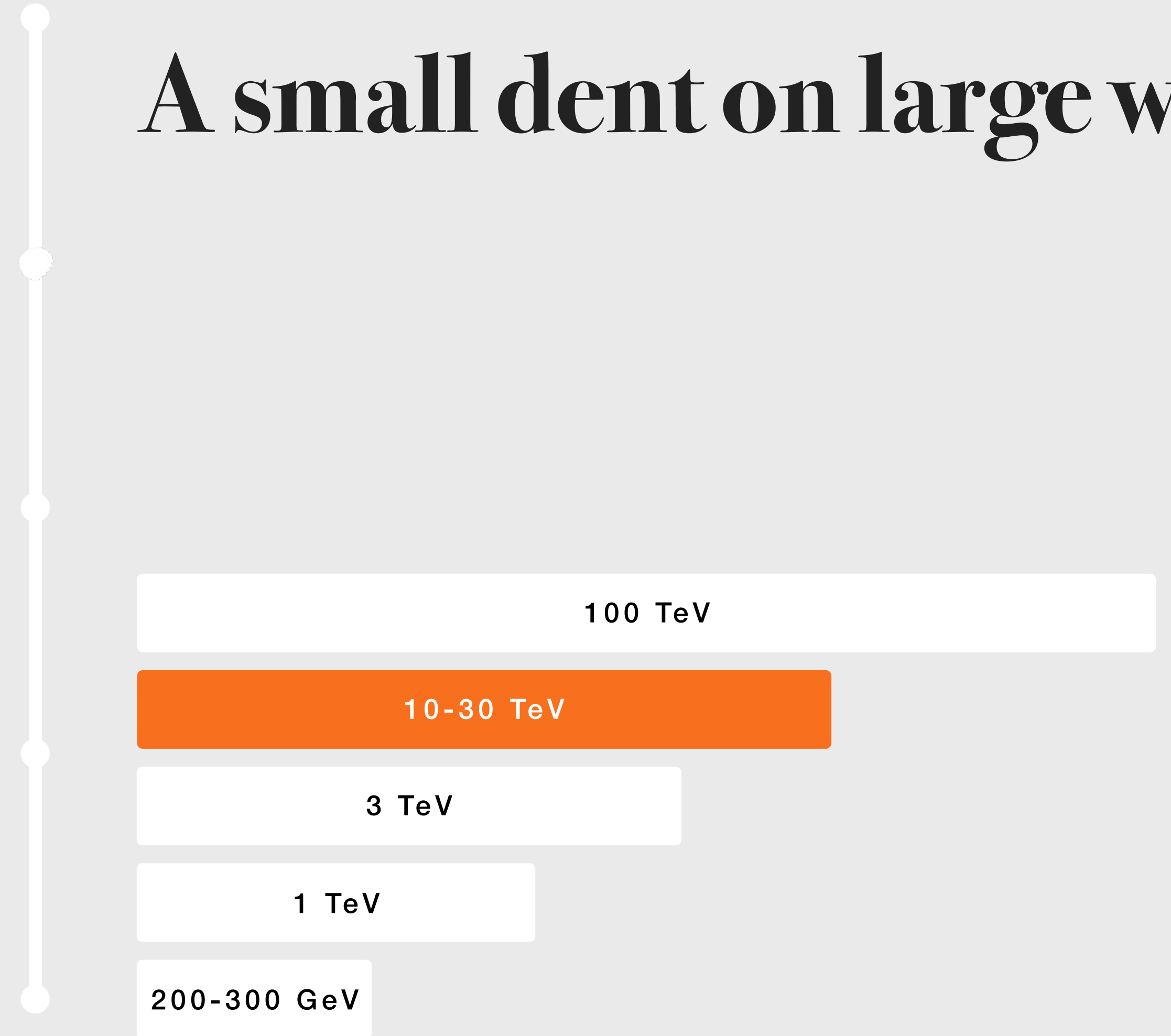
BSM source of CPV










*Fermions masses
and mixings*

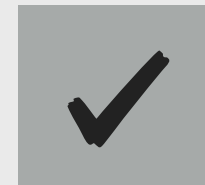
BSM source of CPV

A small dent on large wall



Open Questions on the “big picture” on fundamental physics circa 2020

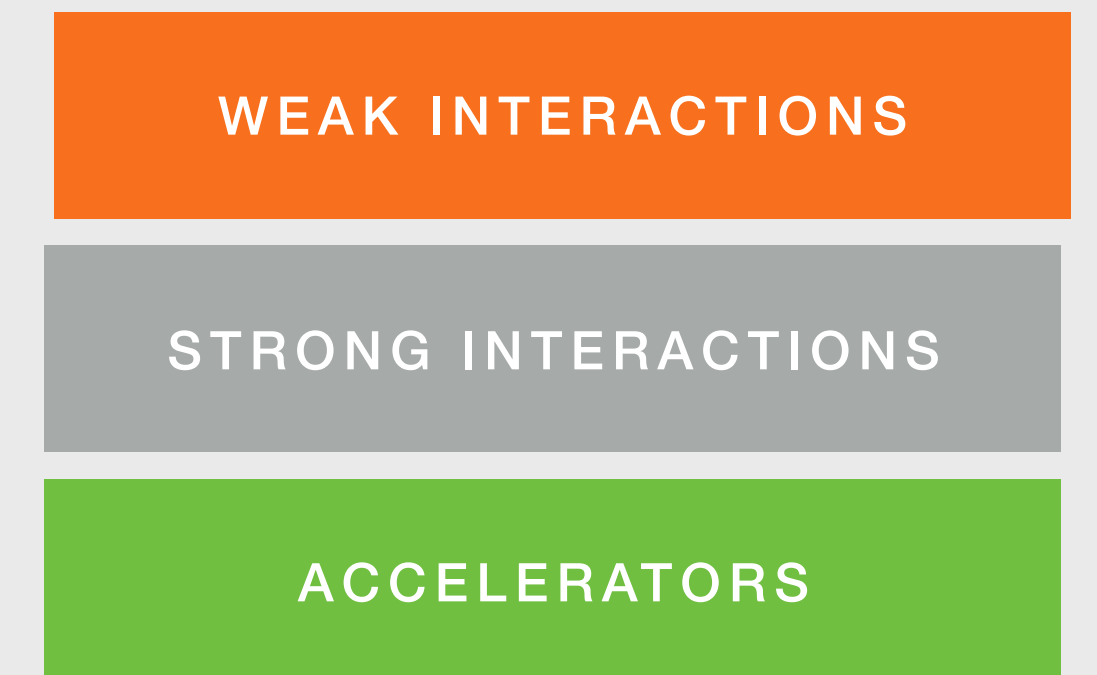
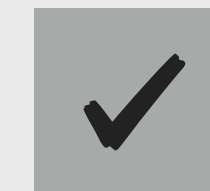
-  • what is the dark matter in the Universe?
-  • why QCD does not violate CP?
-  • how have baryons originated in the early Universe?
-  • what originates flavor mixing and fermions masses?
-  • what gives mass to neutrinos?
- EFT*  • why gravity and weak interactions are so different?
- EFT*  • what fixes the cosmological constant?



We need to explore weak interactions

Open Questions on the “big picture” on fundamental physics circa 2020

- ? • what is the dark matter in the Universe?
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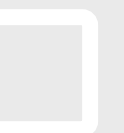


We need to explore weak interactions

Motivations?

Future Lepton Colliders can deliver if they have
mass reach

far enough to see states beyond the LHC exclusions
(unless LHC leaves “light blindspots” as a legacy),
AND have convincing indirect sensitivity to NP through
precision



$(7, \epsilon)_{Dirac}$ $(7, 0)_{C. Scalar}$
 $(5, 0)_{Majorana}$
 $(5, \epsilon)_{Dirac}$ $(5, \epsilon)_{C. Scalar}$
 $(3, 0)_{Majorana}$
 $(3, \epsilon)_{Dirac}$
 $(3, \epsilon)_{C. Scalar}$
 $\left(2, \frac{1}{2}\right)_{Dirac}$

“WIMP” Dark Matter



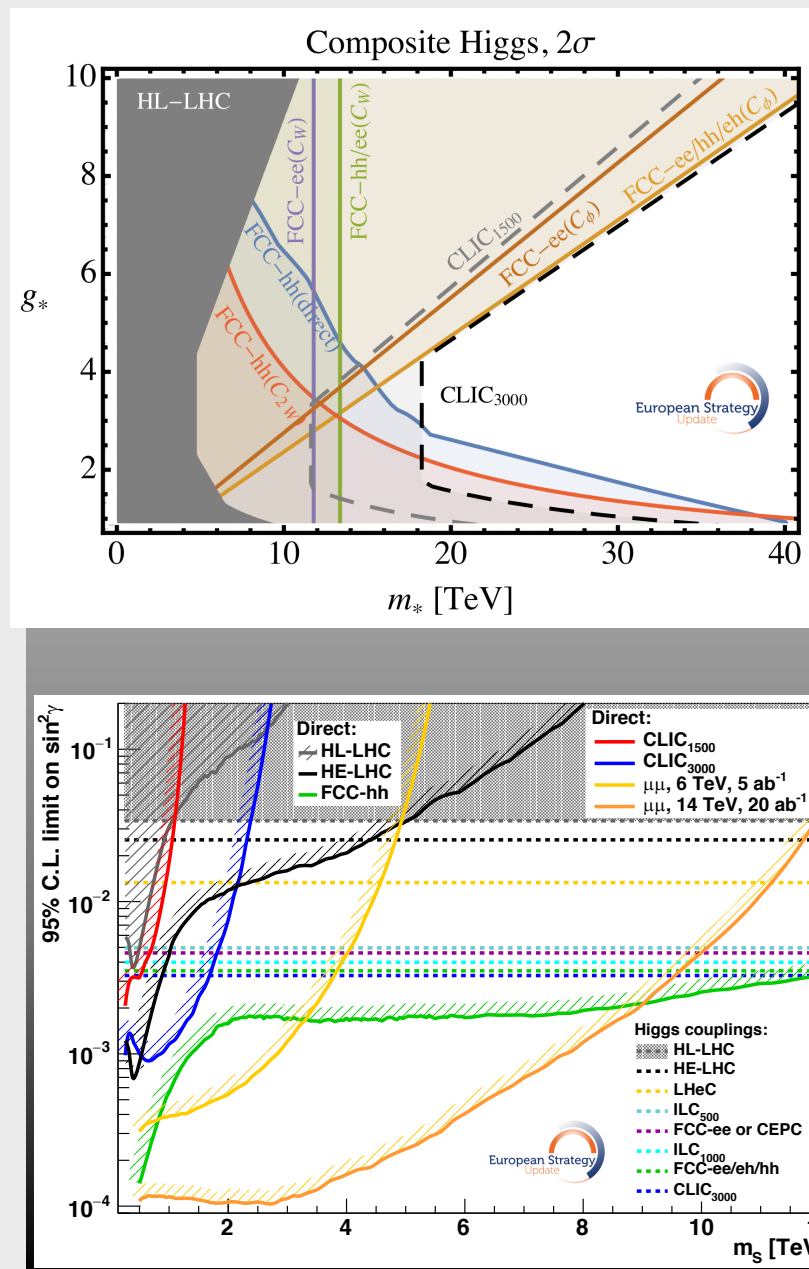
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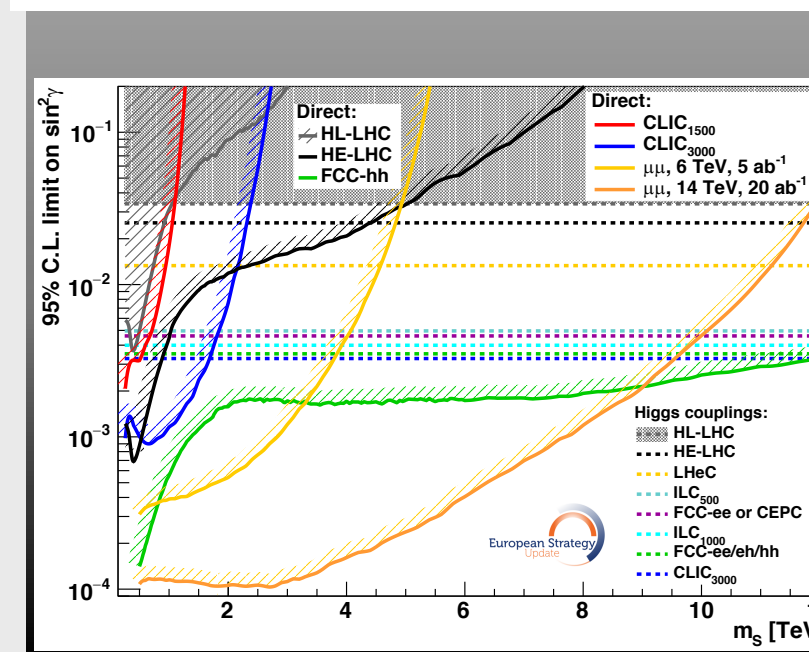
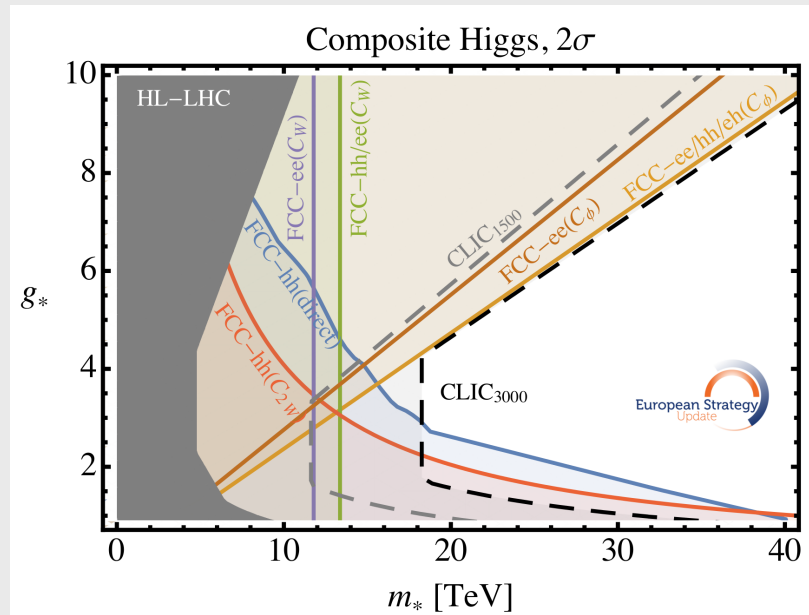


EW symmetry breaking

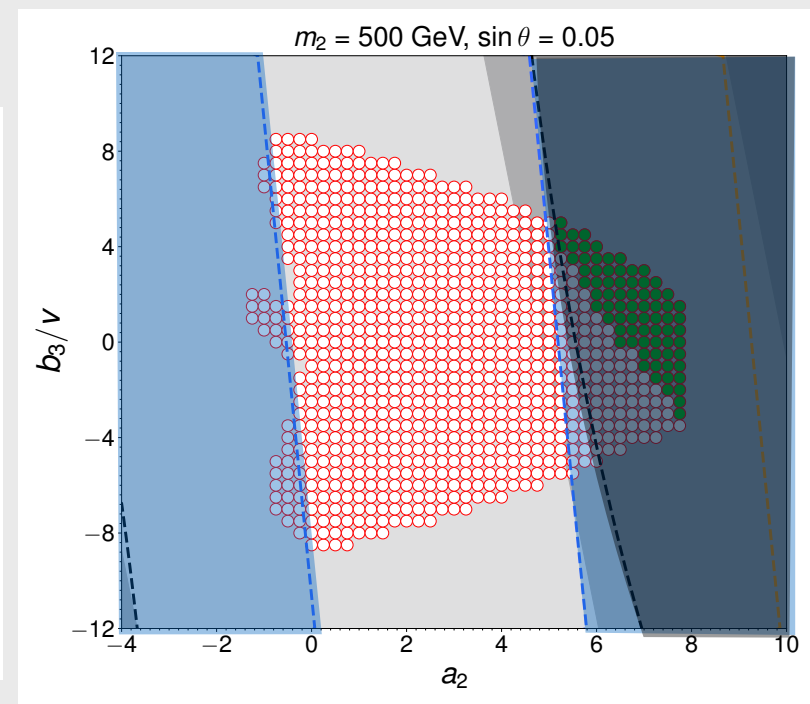


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“WIMP” Dark Matter



EW symmetry breaking



EW phase transition

100 TeV

10-30 TEV

3 TEV

Several important milestones: **full exploration of TeV EW states, EW phase transition, TeV Dark Matter**

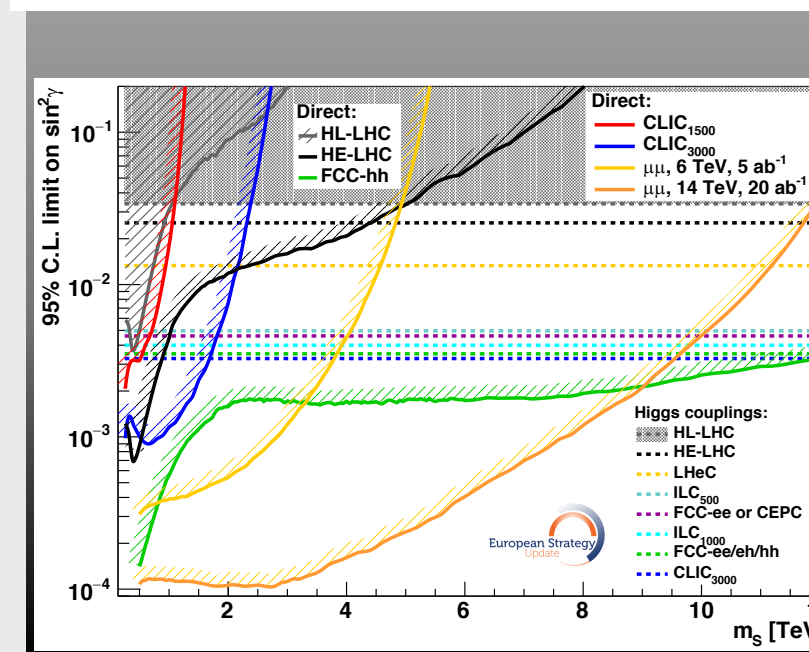
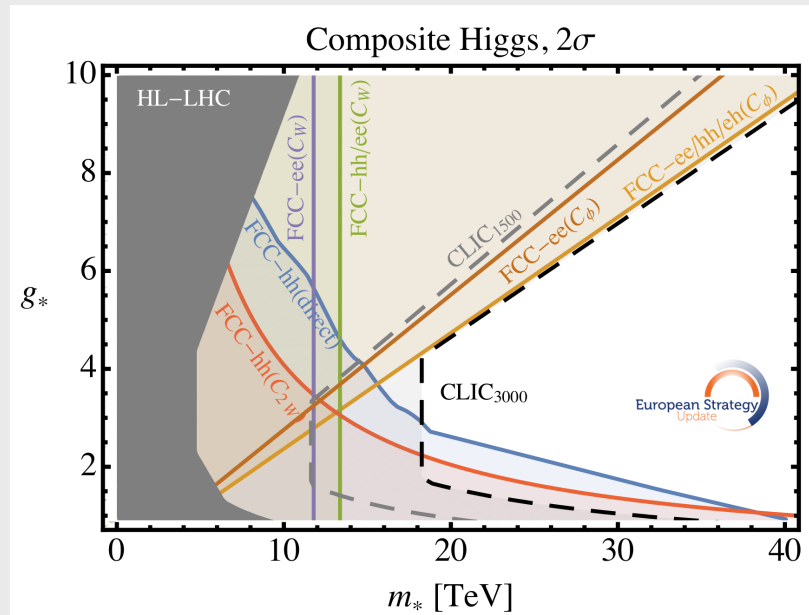
1 TeV

200-300 GeV

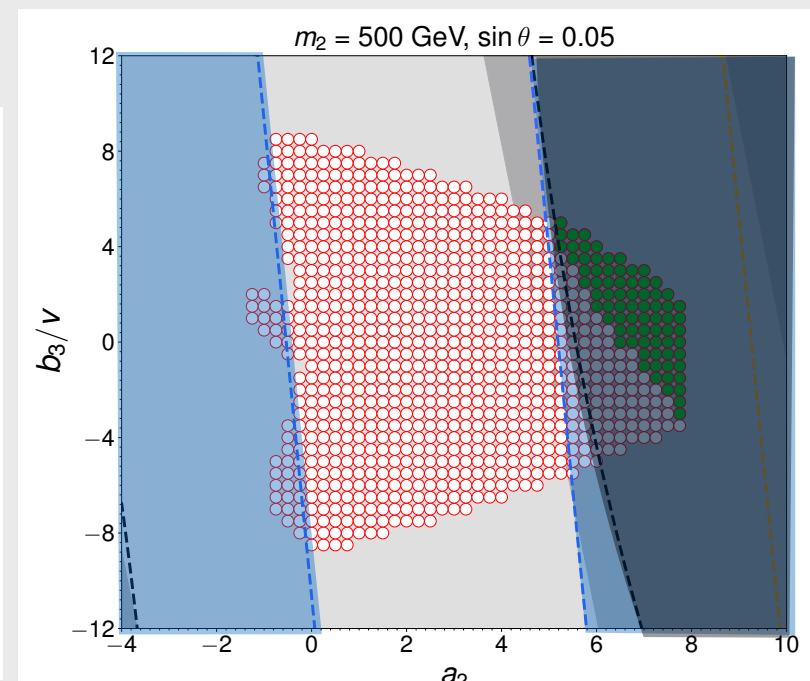


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Λεμμα

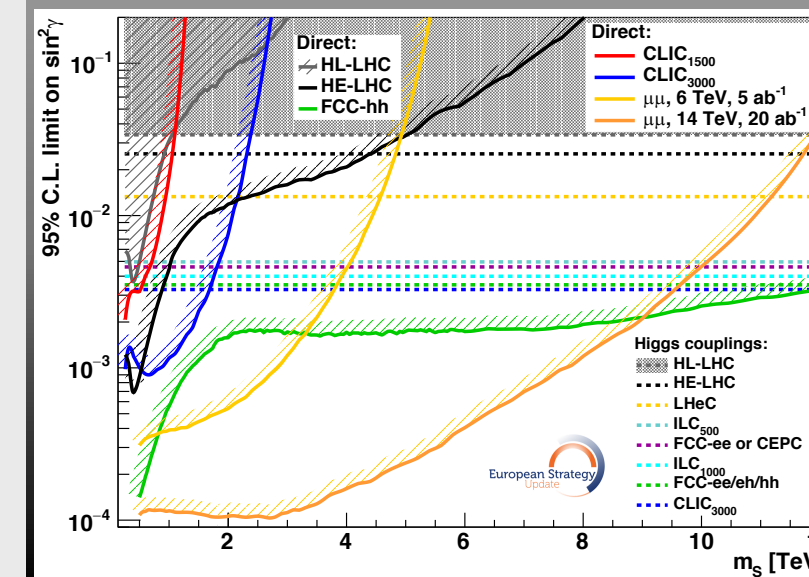
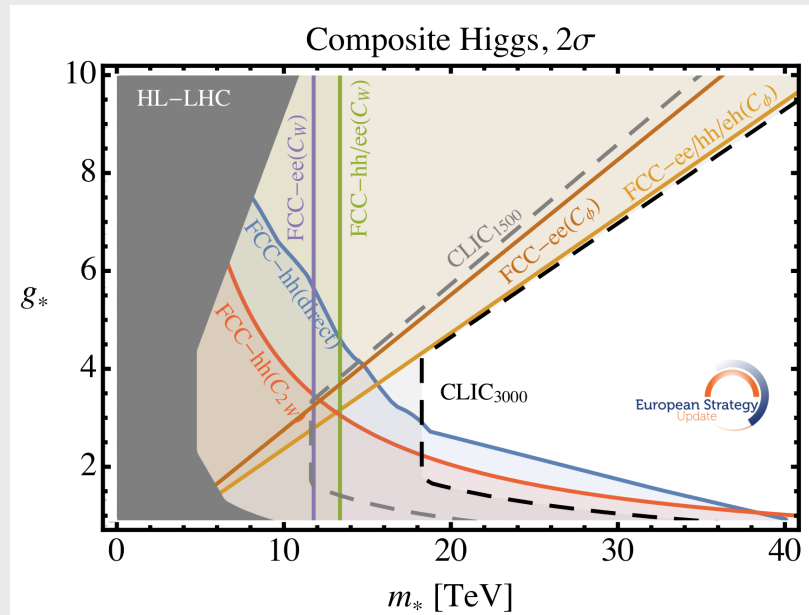


PWFA

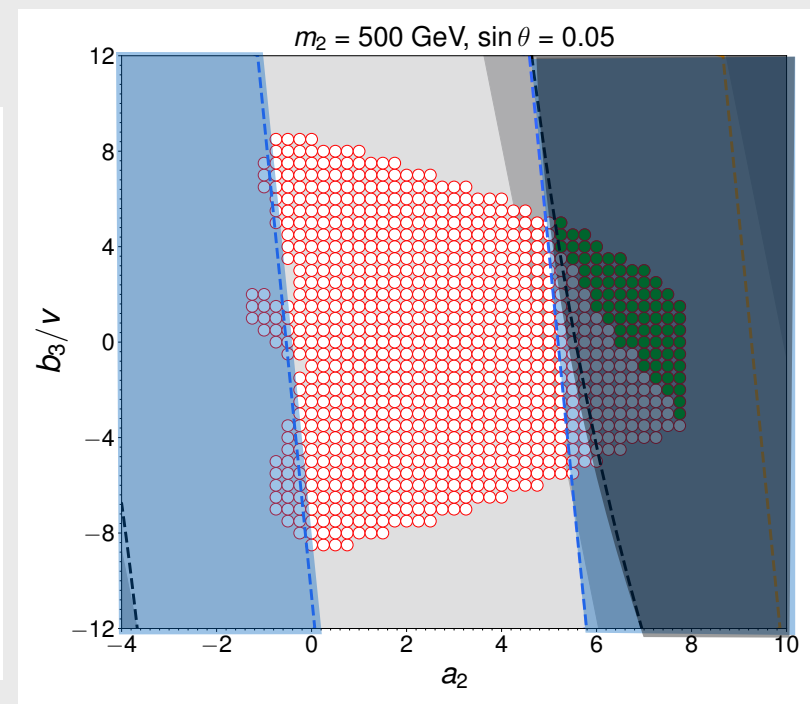


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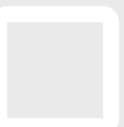
1 TeV

200-300 GeV

Λεμμα

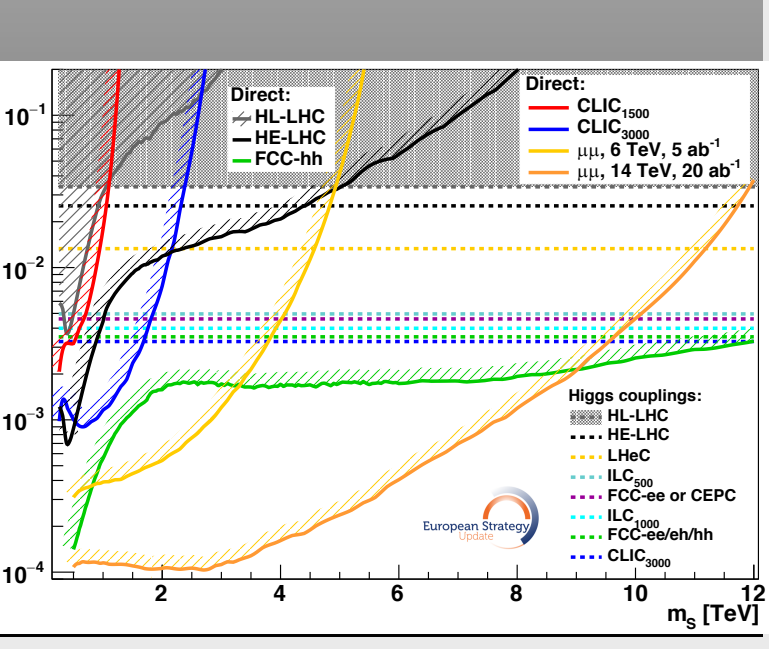
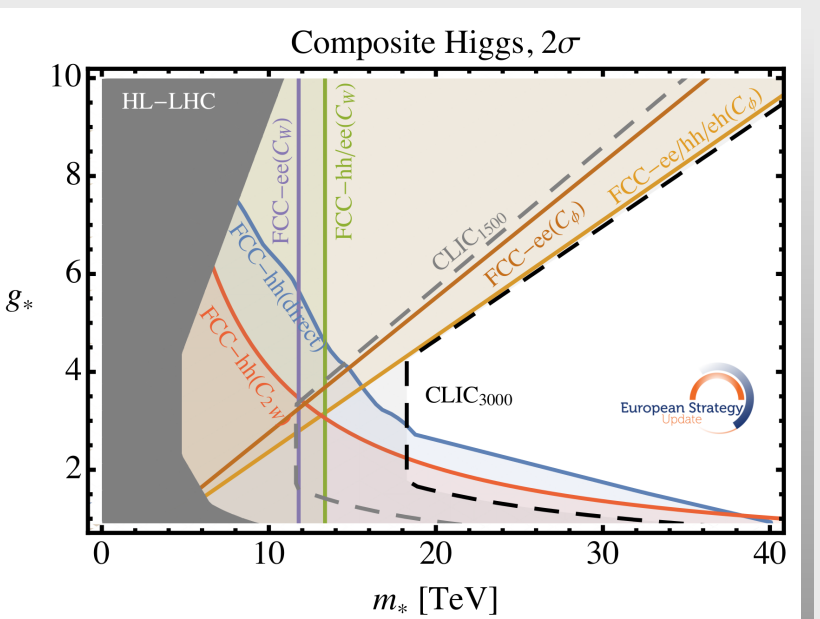


PWFA

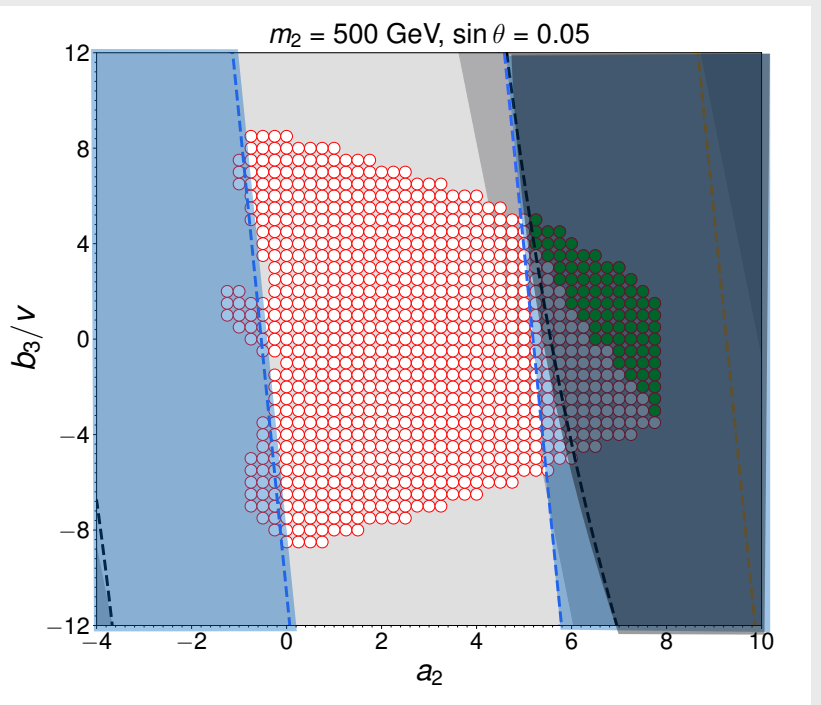


~~(7, ϵ) Dirac~~ ~~(7, 0) C. Scalar~~
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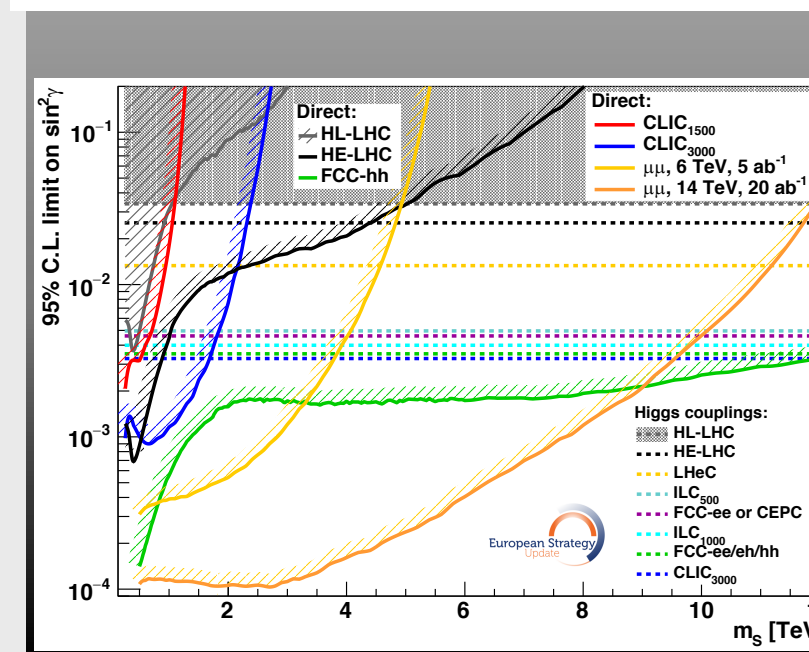
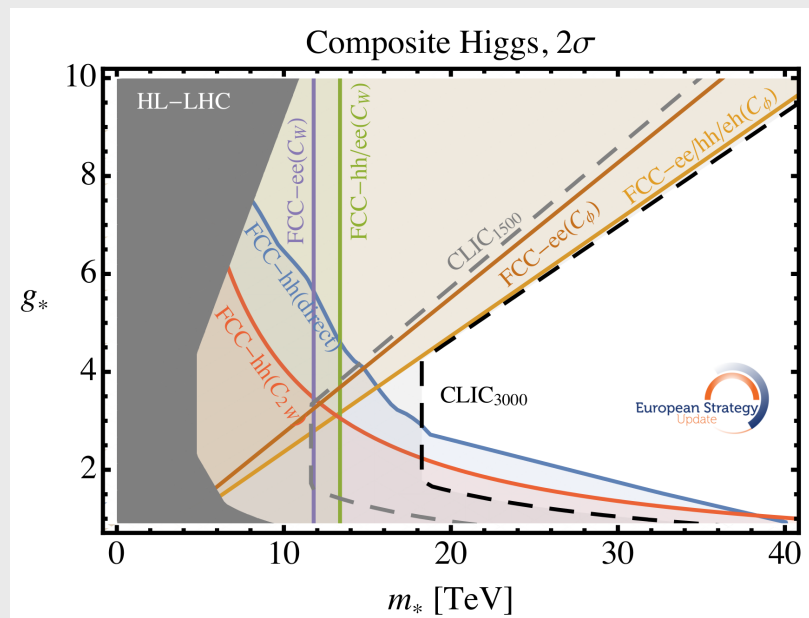


PWFA

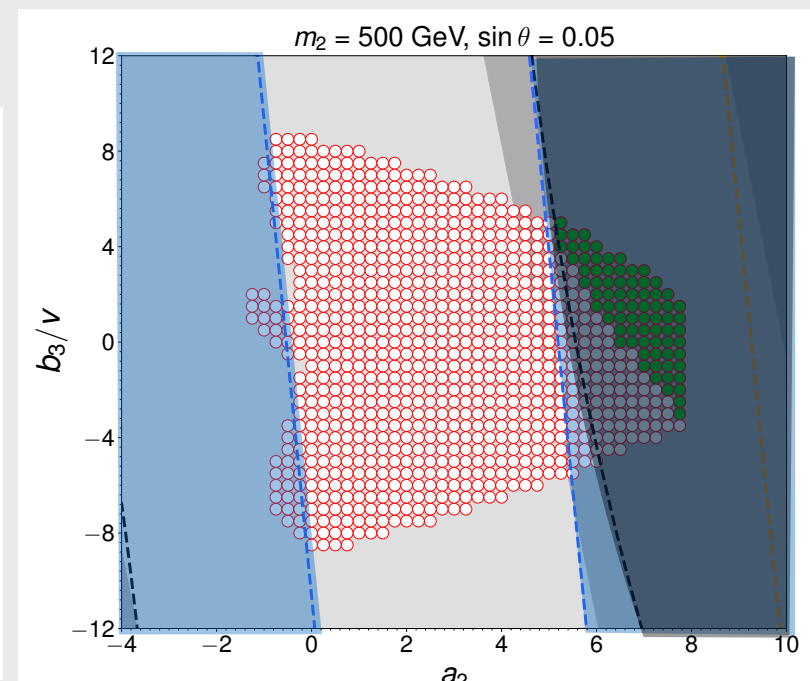


~~(1, 0) Dirac~~ ~~(7, 0) C. Scalar~~
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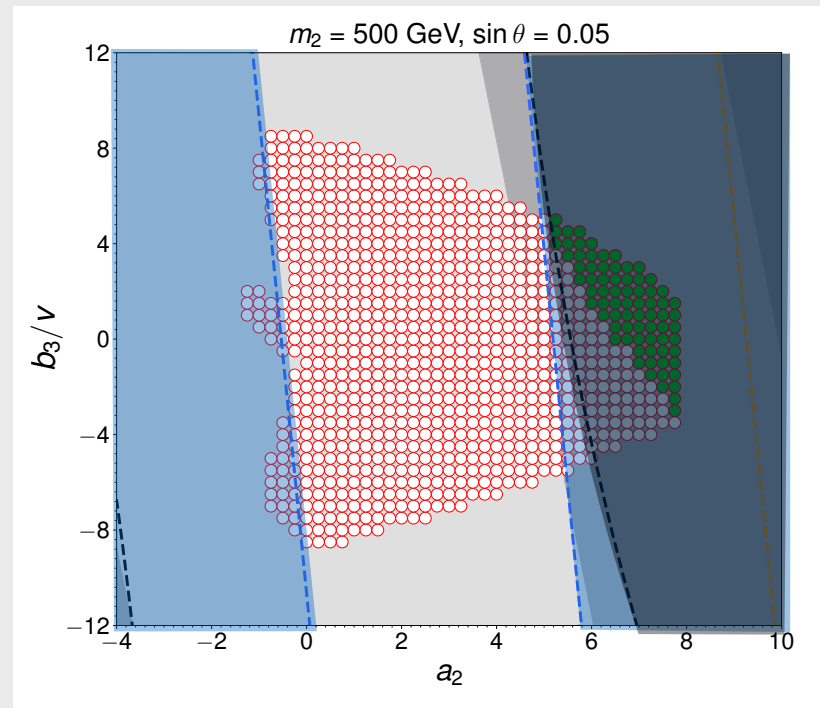
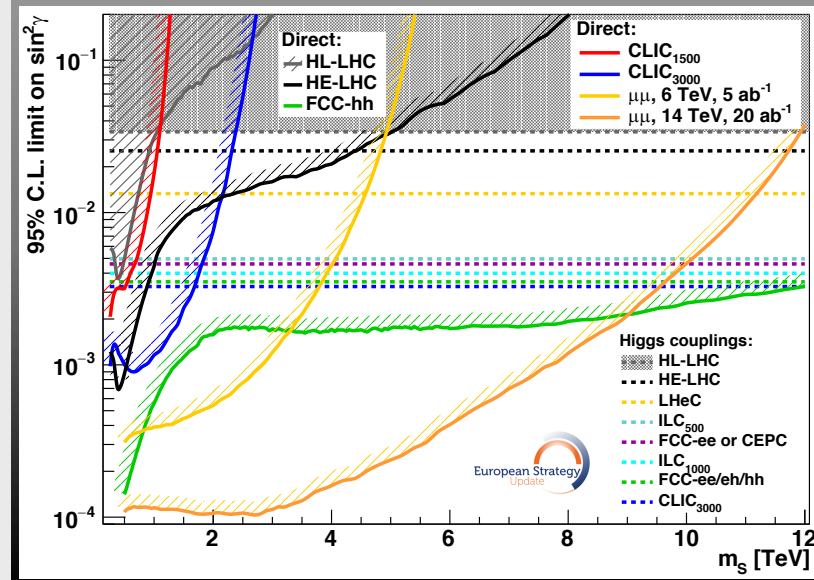
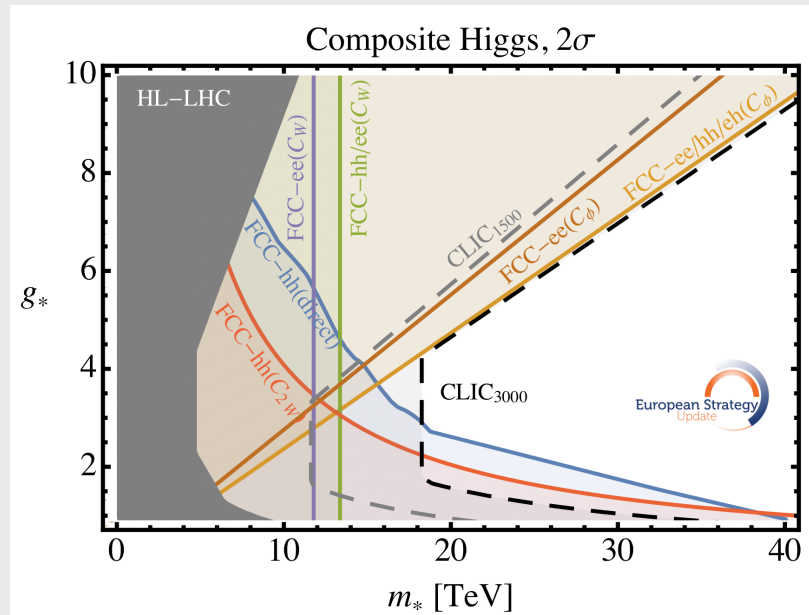
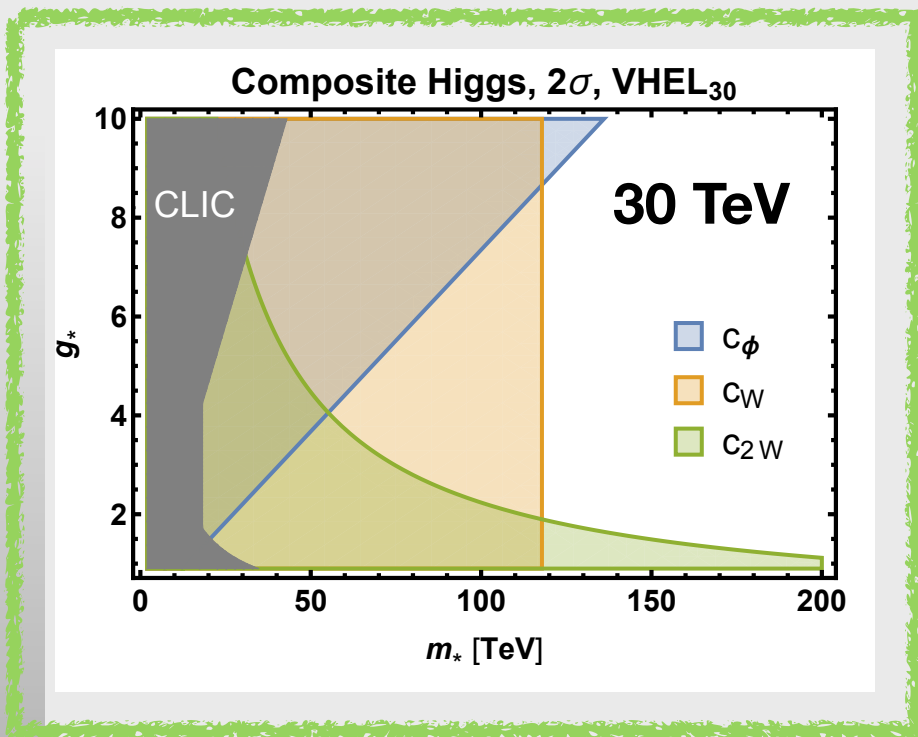
200-300 GeV

Λεμμα



PWFA





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PWFA



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“WIMP” Dark Matter

EW symmetry breaking

EW phase transition



Conclusion

- Leptons beam structure enables *qualitatively new investigations* of the electroweak/Higgs sector
- CLIC definitively shows that lepton colliders can have *both precision and mass reach* to probe new physics well beyond TeV
- If novel acceleration technologies can deliver even larger *energy* and keep the *luminosity* on track with $\mathcal{L} \propto E_{com}^2$ we can start probing fundamental interactions in novel and deeper ways.

Thank You!

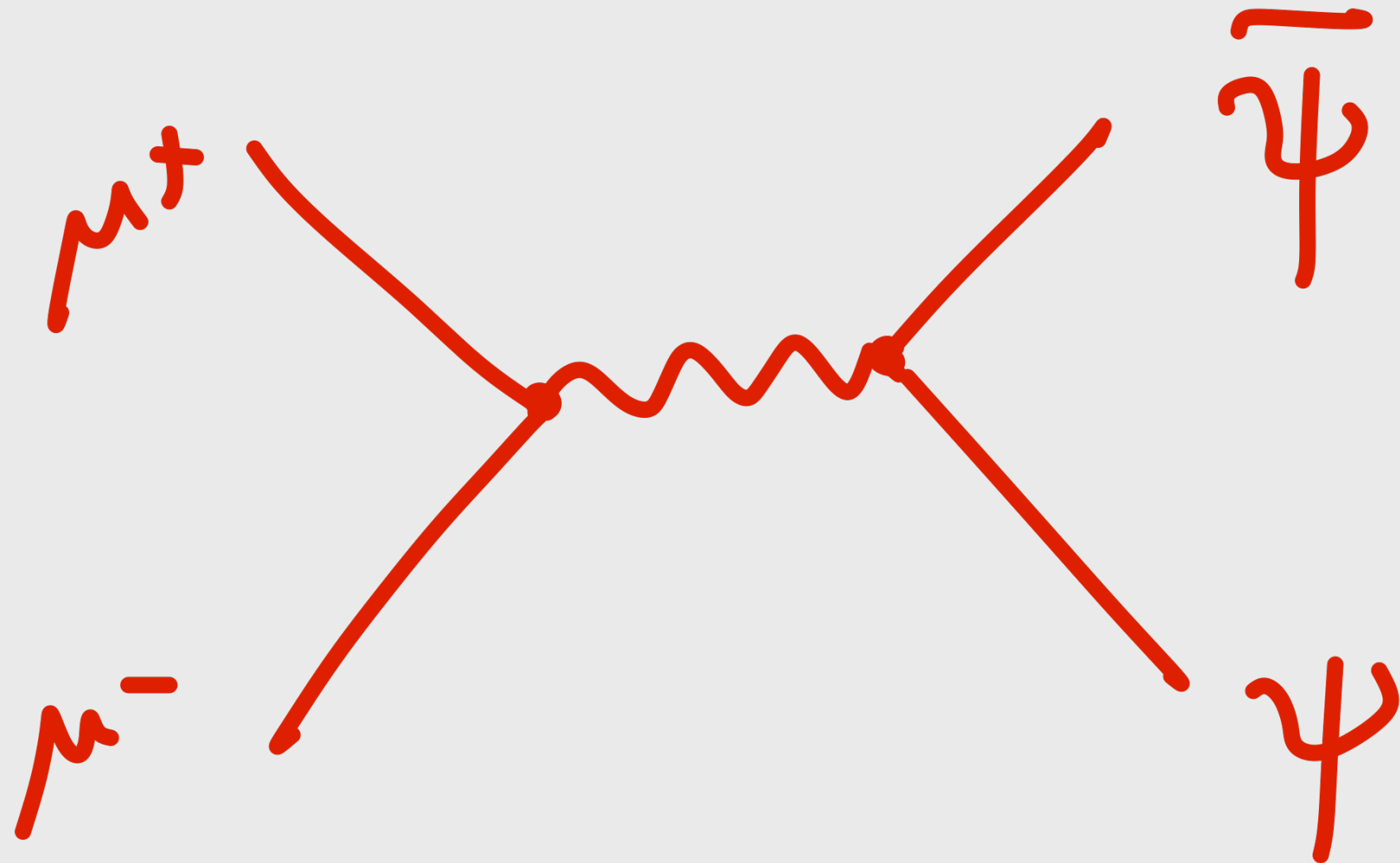
“Valence” Leptons

$\ell^+ \ell^- \rightarrow$ new physics

VALENCE

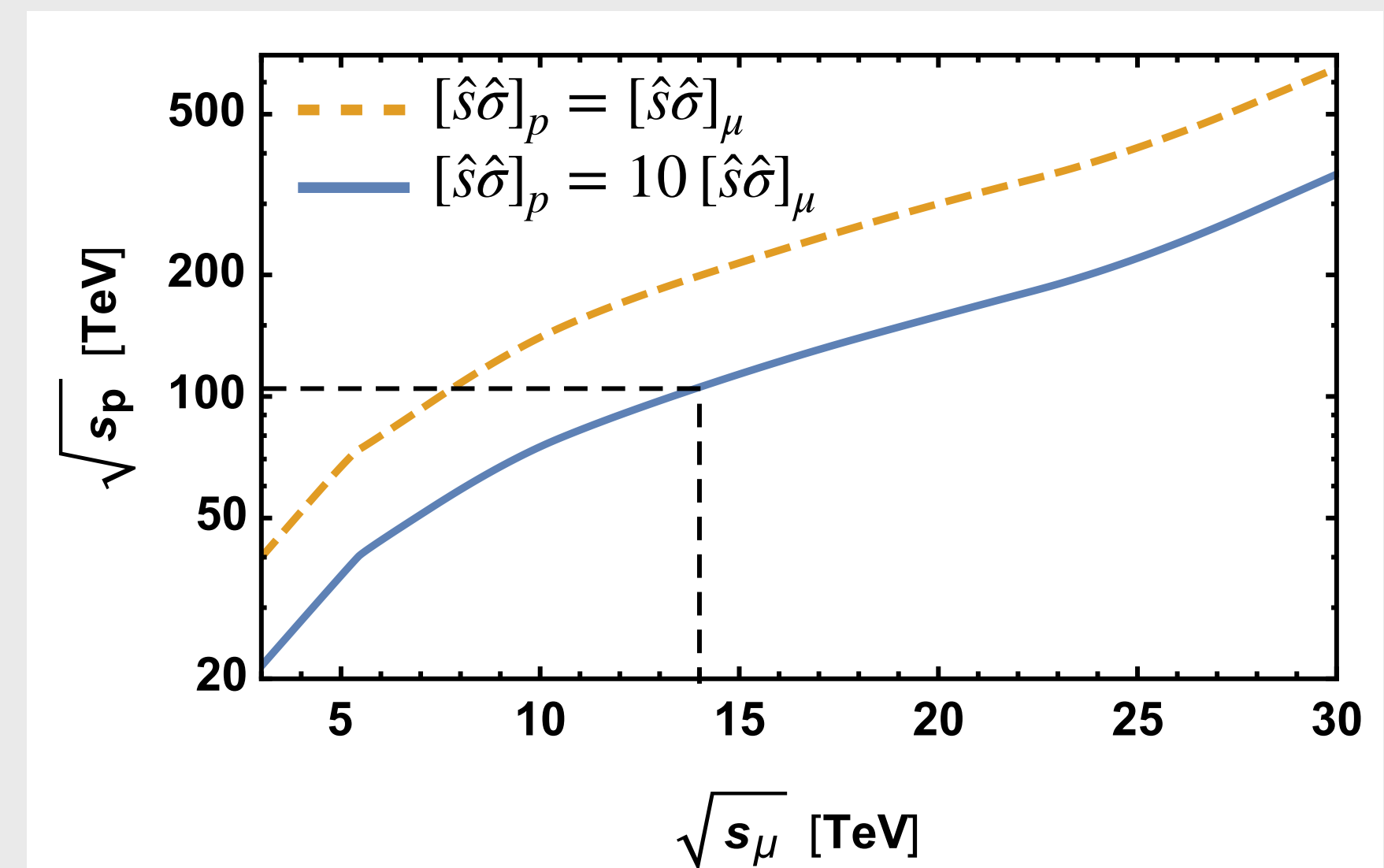
LEPTONS

Can produce heavy new physics (colored or not)



in principle can probe directly new states at O(10) TeV scale!

Compares pretty well with a pp collider

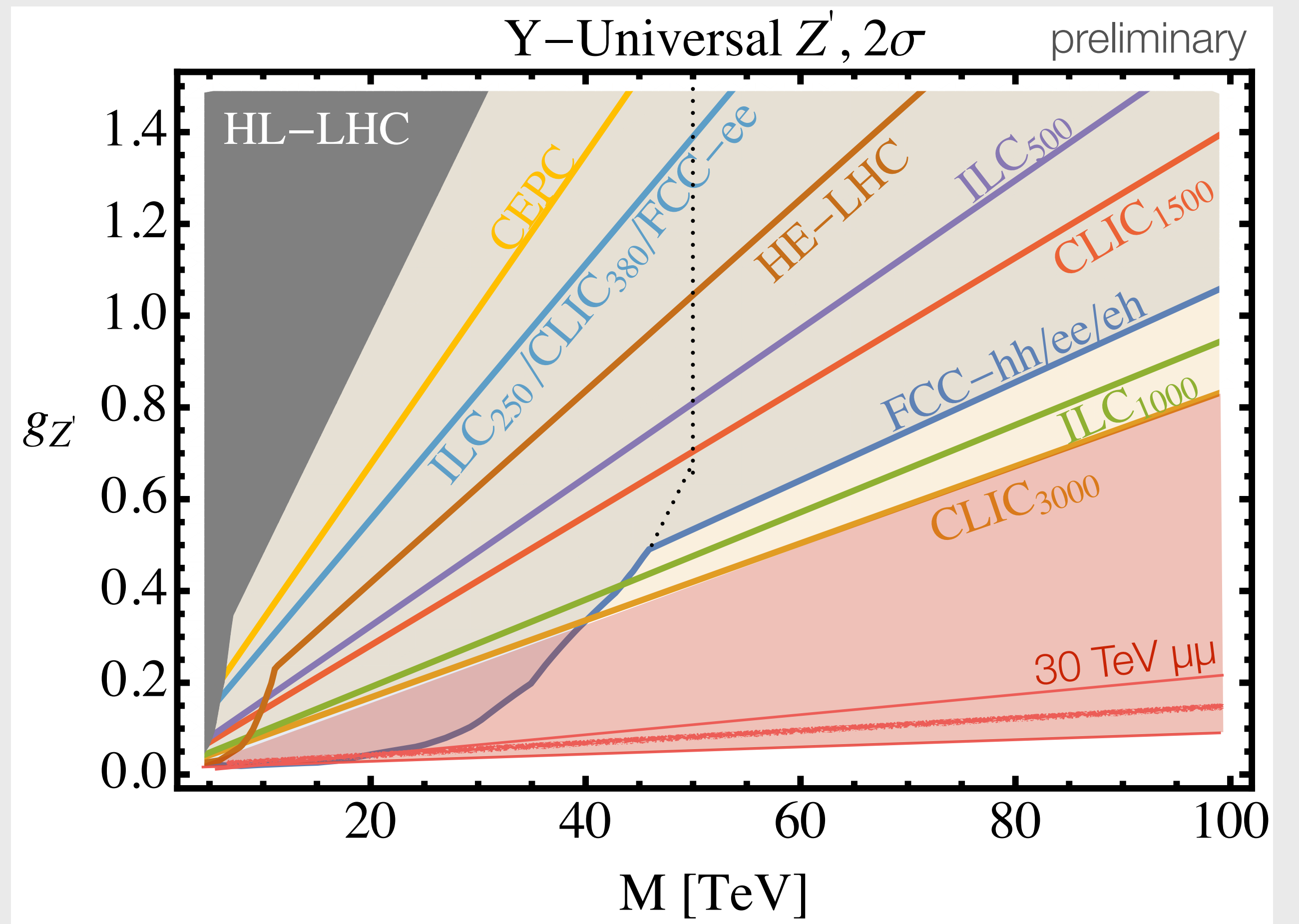
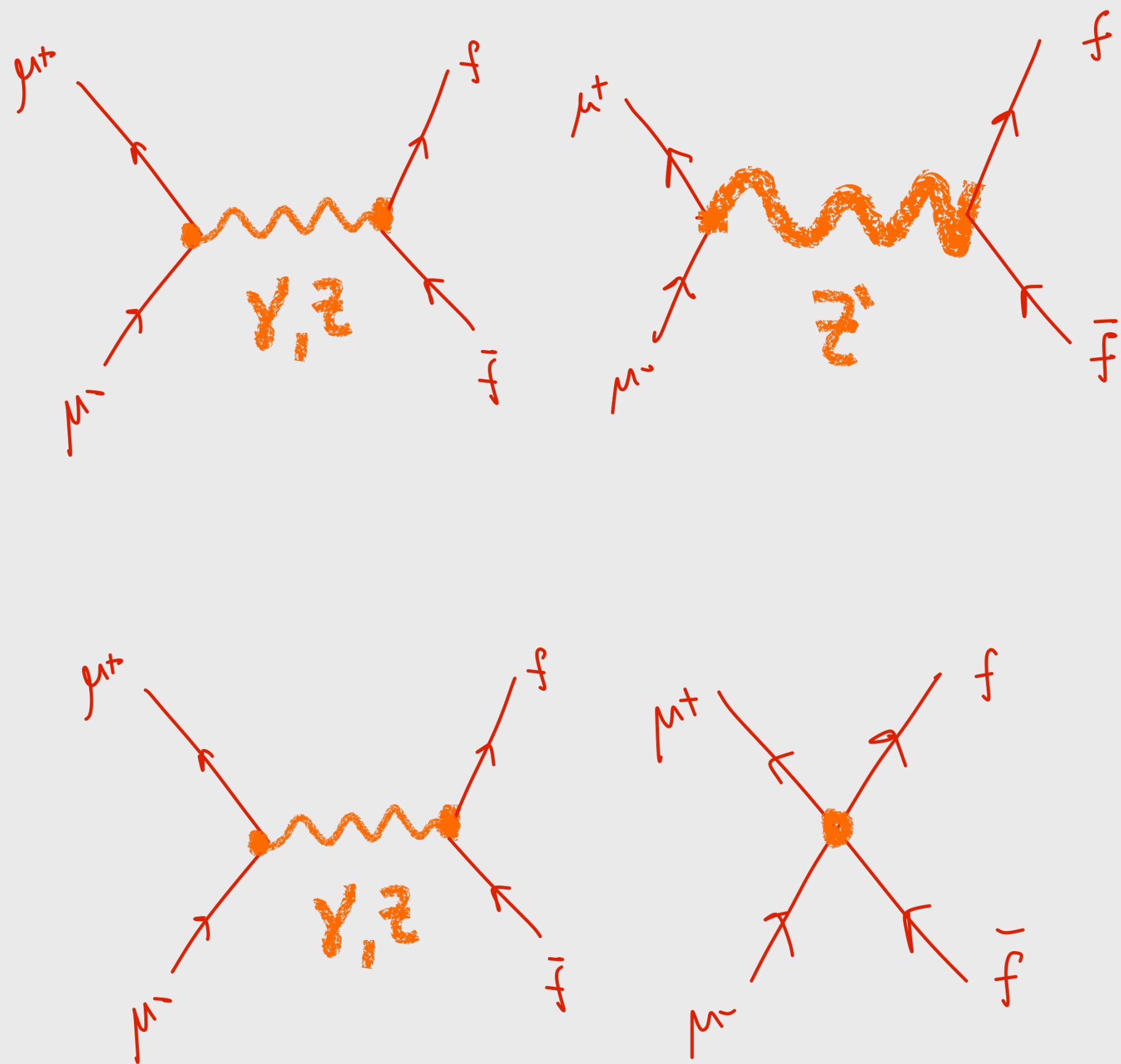


14 TeV $\mu\mu$ roughly equivalent to 100 TeV pp

A heavy Z'

DRELL-YAN

RATES AND ANGULAR DISTRIBUTIONS



UnMixed Singlet for EW phase transition

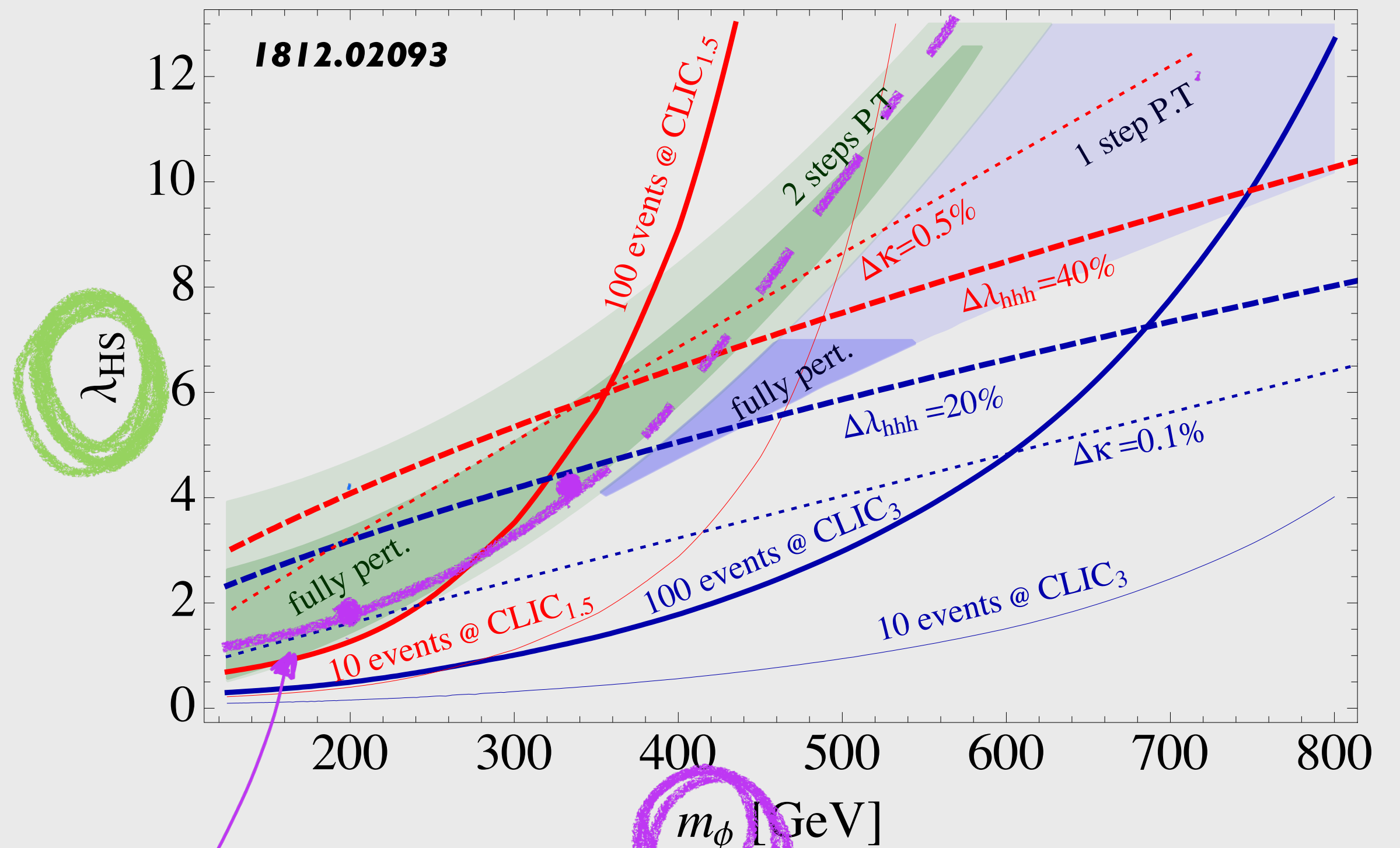
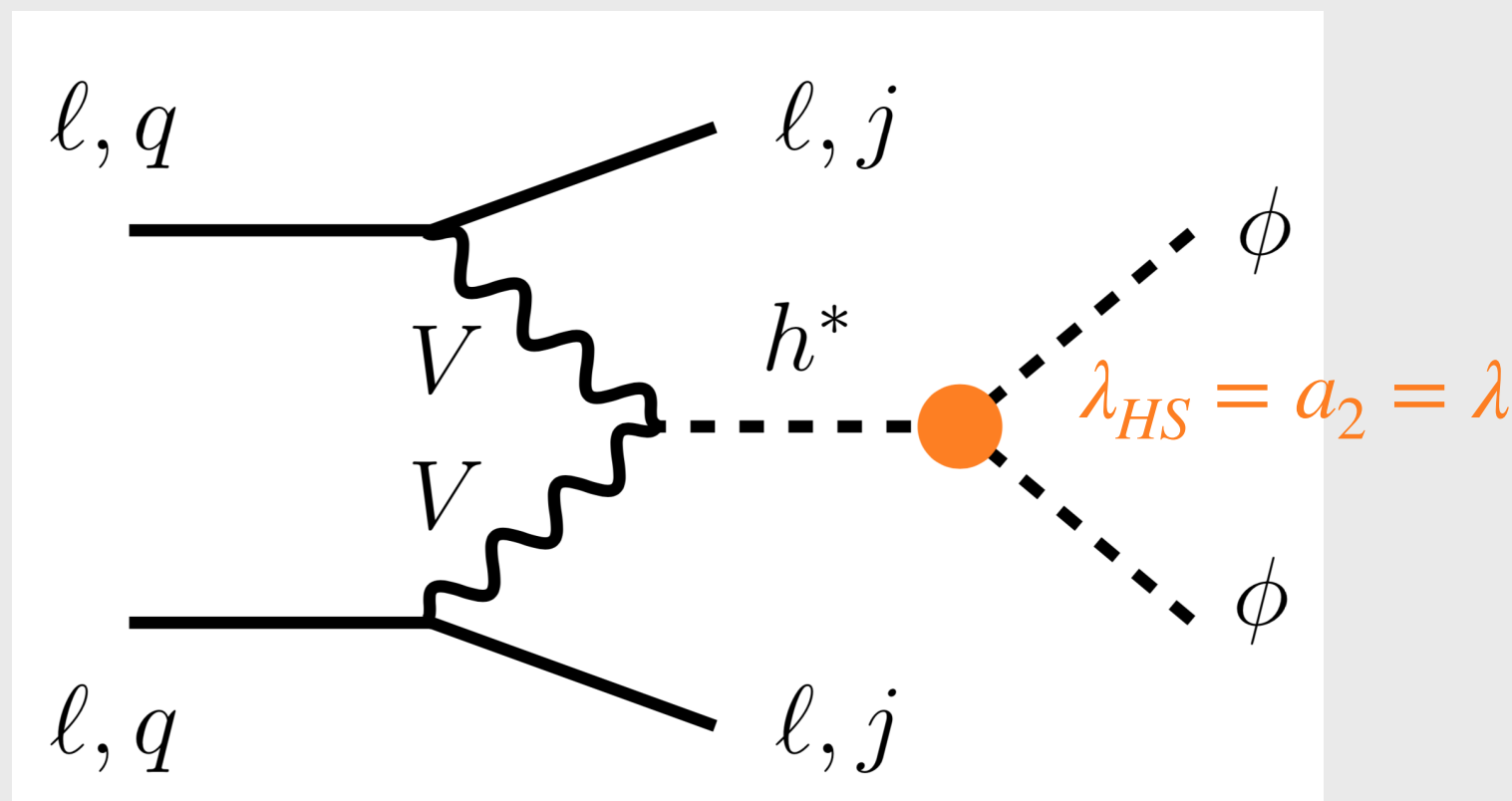
EW PHASE TRANSITION

IS IT FIRST ORDER?

$$V(\Phi, S) = -\mu^2 (\Phi^\dagger \Phi) + \lambda (\Phi^\dagger \Phi)^2 + \frac{a_1}{2} (\Phi^\dagger \Phi) S + \frac{a_2}{2} (\Phi^\dagger \Phi) S^2 + b_1 S + \frac{b_2}{2} S^2 + \frac{b_3}{3} S^3 + \frac{b_4}{4} S^4.$$

independent parameters

$\{v, m_1, m_2, \lambda, a_2, b_3, b_4\}$.



zero-mixing: Higgs couplings as good or better than VBF
 tiny mixing: displaced decays events

parameters space of 1st order phase transition accessible by several probes

UnMixed Singlet for EW phase transition

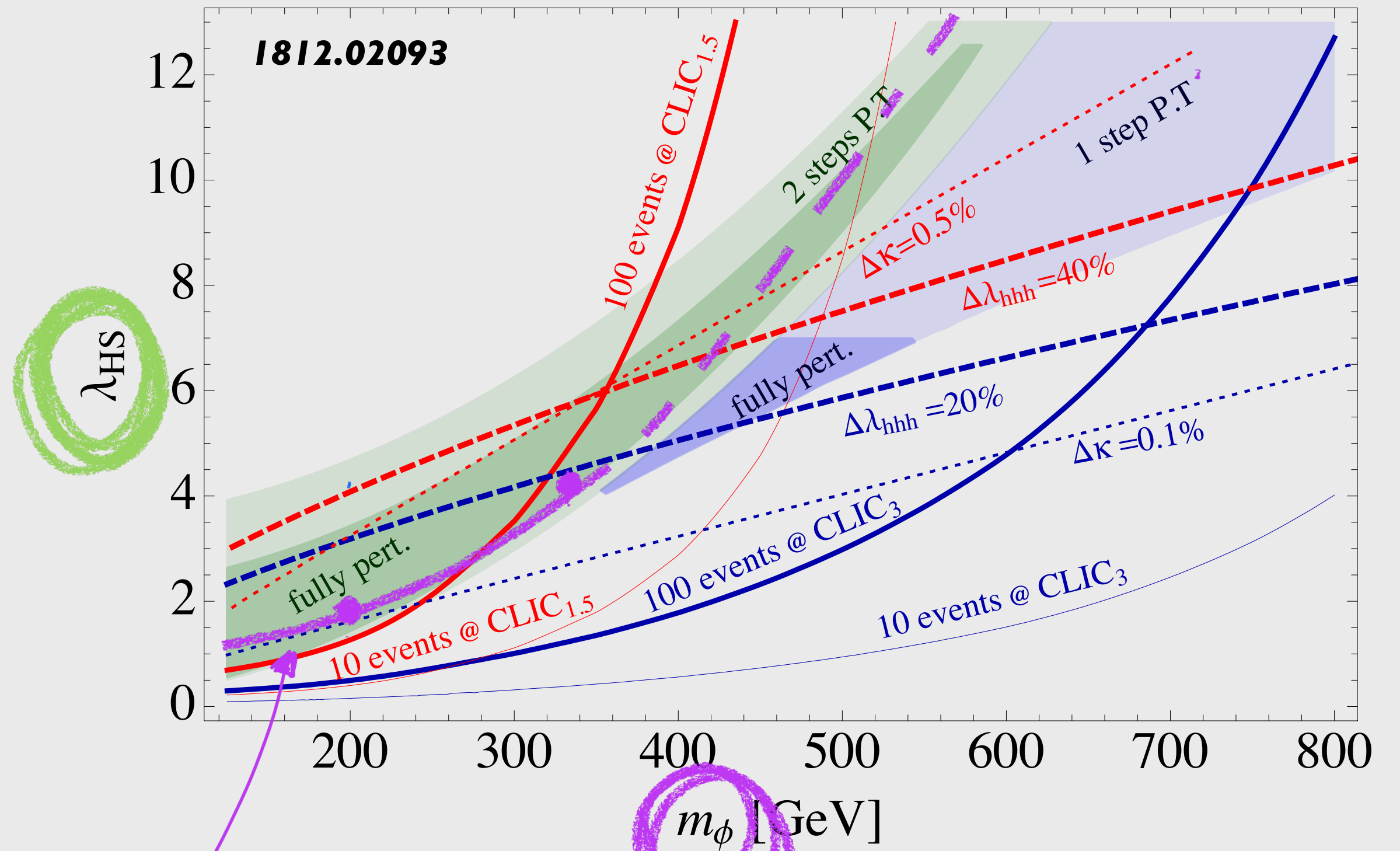
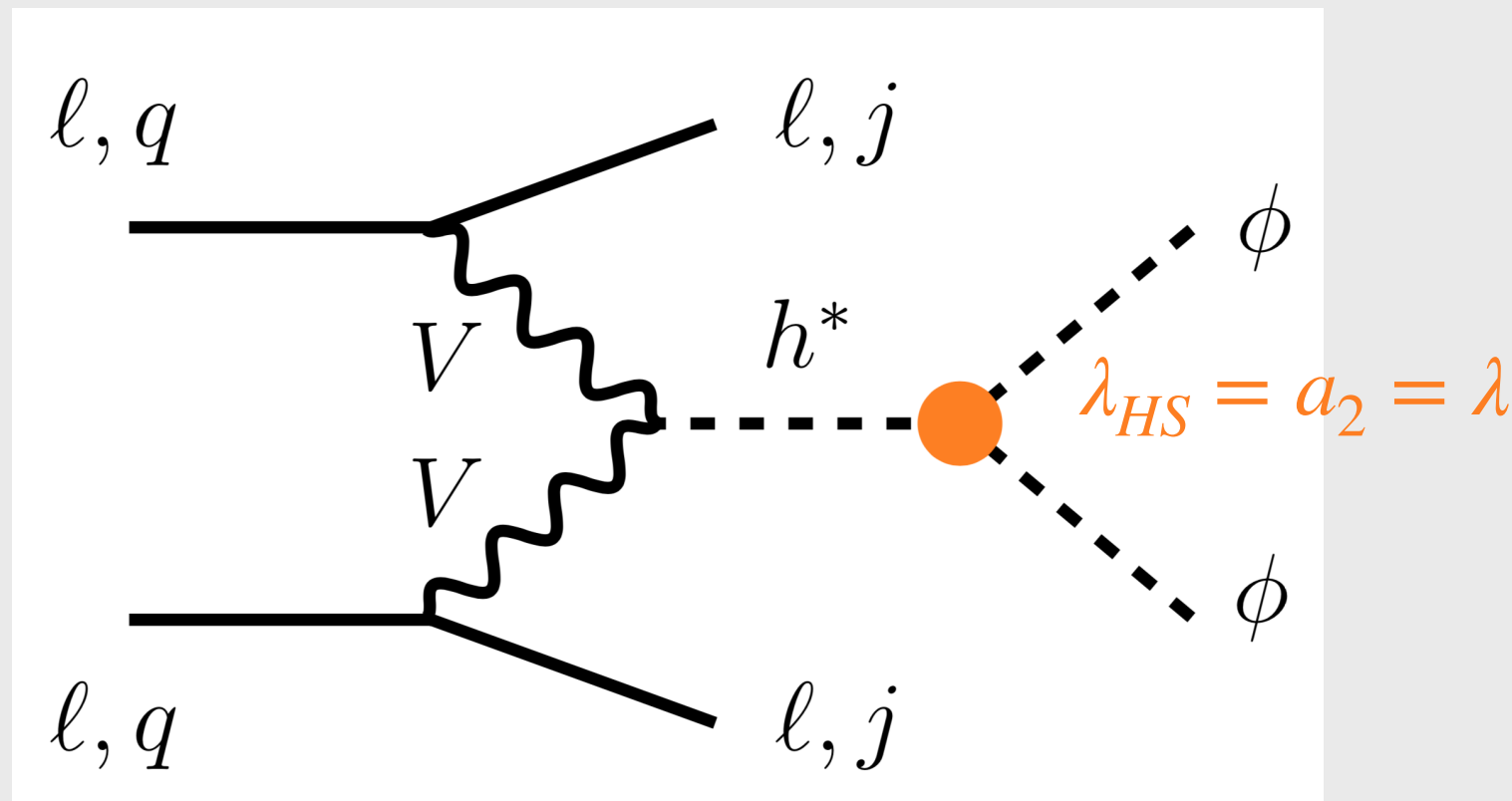
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1910.04170 VBF

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EW symmetry breaking

Electroweak symmetry breaking

Big picture questions:

- Higgs compositeness
- Extended Higgs Sector

Electroweak symmetry breaking

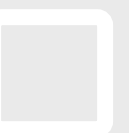
Big picture questions:

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“The size of the Higgs boson”

it matters because being “point-like” is the source of all the theoretical questions on the Higgs boson and weak scale

... and if it is not ... well, that is physics beyond the Standard Model!



Effects of the size of the Higgs boson

$h \sim \pi$

STRONGLY INTERACTING LIGHT HIGGS

$$\begin{aligned}
 \mathcal{L}_{universal}^{d=6} = & c_H \frac{g_*^2}{m_*^2} \mathcal{O}_H + c_T \frac{N_c \epsilon_q^4 g_*^4}{(4\pi)^2 m_*^2} \mathcal{O}_T + c_6 \lambda \frac{g_*^2}{m_*^2} \mathcal{O}_6 + \frac{1}{m_*^2} [c_W \mathcal{O}_W + c_B \mathcal{O}_B] \\
 & + \frac{g_*^2}{(4\pi)^2 m_*^2} [c_{HW} \mathcal{O}_{HW} + c_{HB} \mathcal{O}_{HB}] + \frac{y_t^2}{(4\pi)^2 m_*^2} [c_{BB} \mathcal{O}_{BB} + c_{GG} \mathcal{O}_{GG}] \\
 & + \frac{1}{g_*^2 m_*^2} [c_{2W} g^2 \mathcal{O}_{2W} + c_{2B} g'^2 \mathcal{O}_{2B}] + c_{3W} \frac{3! g^2}{(4\pi)^2 m_*^2} \mathcal{O}_{3W} \\
 & + c_{y_t} \frac{g_*^2}{m_*^2} \mathcal{O}_{y_t} + c_{y_b} \frac{g_*^2}{m_*^2} \mathcal{O}_{y_b}
 \end{aligned}$$

$$1/f \sim g_*/m_*$$

$$1/(g_* f) \sim 1/m_*$$

$$g_{SM}/(g_* f) \sim g_{SM}/m_*$$



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$$\ell_{Higgs} \sim 1/m_*$$

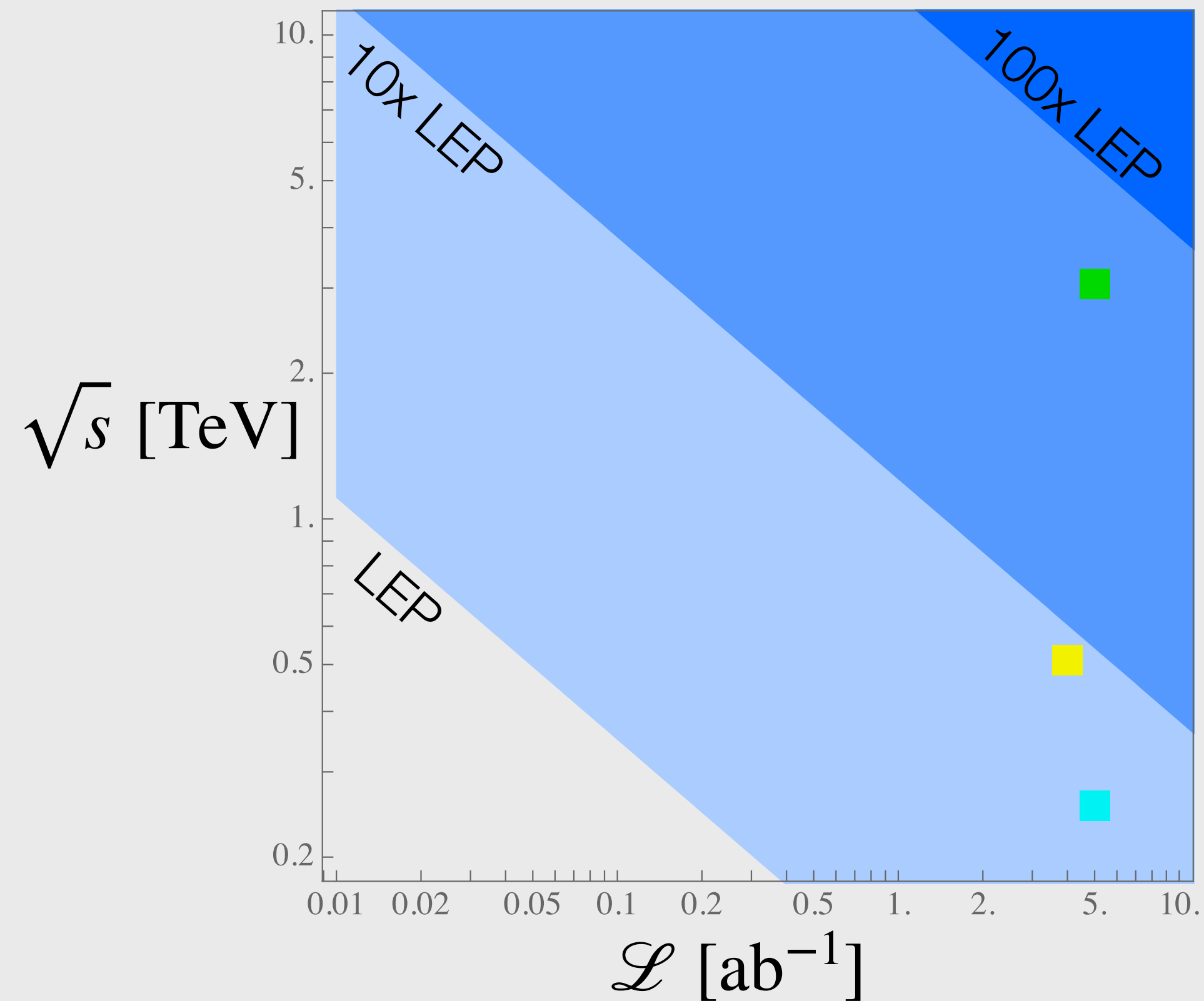




Ever higher energy colliders can exploit “precise” measurements at the 10% level

TOTAL **RATE** $\left| A_{SM}^{(00)} \right|^2 + A_{SM}^{00} \cdot A_{BSM}^{00} + \dots$

$$\hat{S}_{95\%} \lesssim 1.2 \cdot 10^{-4} \frac{1}{E_{beam}/\text{TeV}} \cdot \frac{1}{\sqrt{\mathcal{L}/\text{ab}^{-1}}}$$





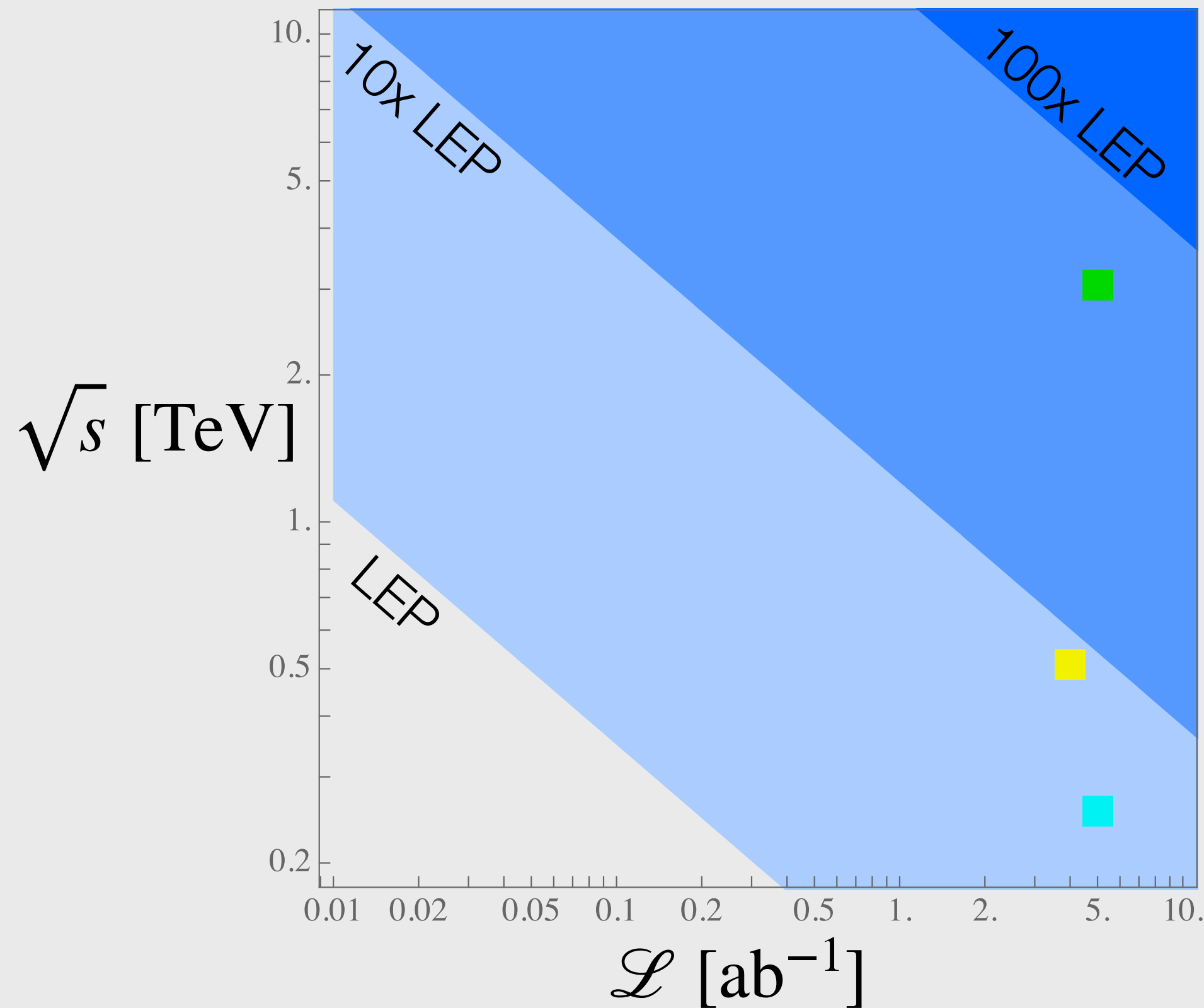
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$$c_W = \hat{S}/m_W^2$$

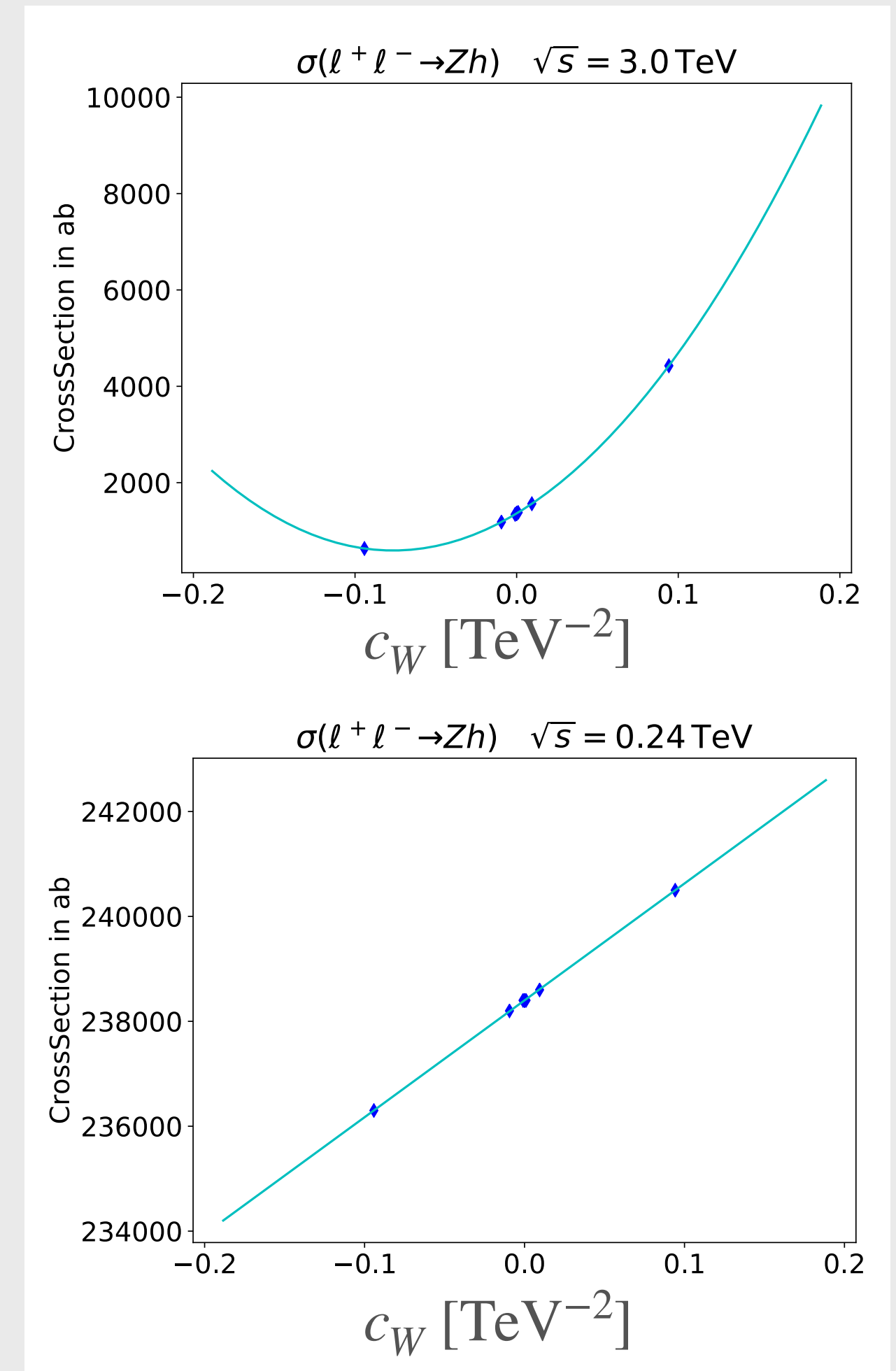
$$c_W \lesssim 0.02 \text{ TeV}^{-2} \frac{1}{E_{beam}/\text{TeV}} \cdot \frac{1}{\sqrt{\mathcal{L}/\text{ab}^{-1}}}$$

$$\hat{S}_{95\%} \lesssim 1.2 \cdot 10^{-4} \frac{1}{E_{beam}/\text{TeV}} \cdot \frac{1}{\sqrt{\mathcal{L}/\text{ab}^{-1}}}$$



$$\hat{S} < 3 \cdot 10^{-5} \text{ (95 \% CL)} \quad \leftarrow \mathcal{L} = 5 \text{ ab}^{-1}$$

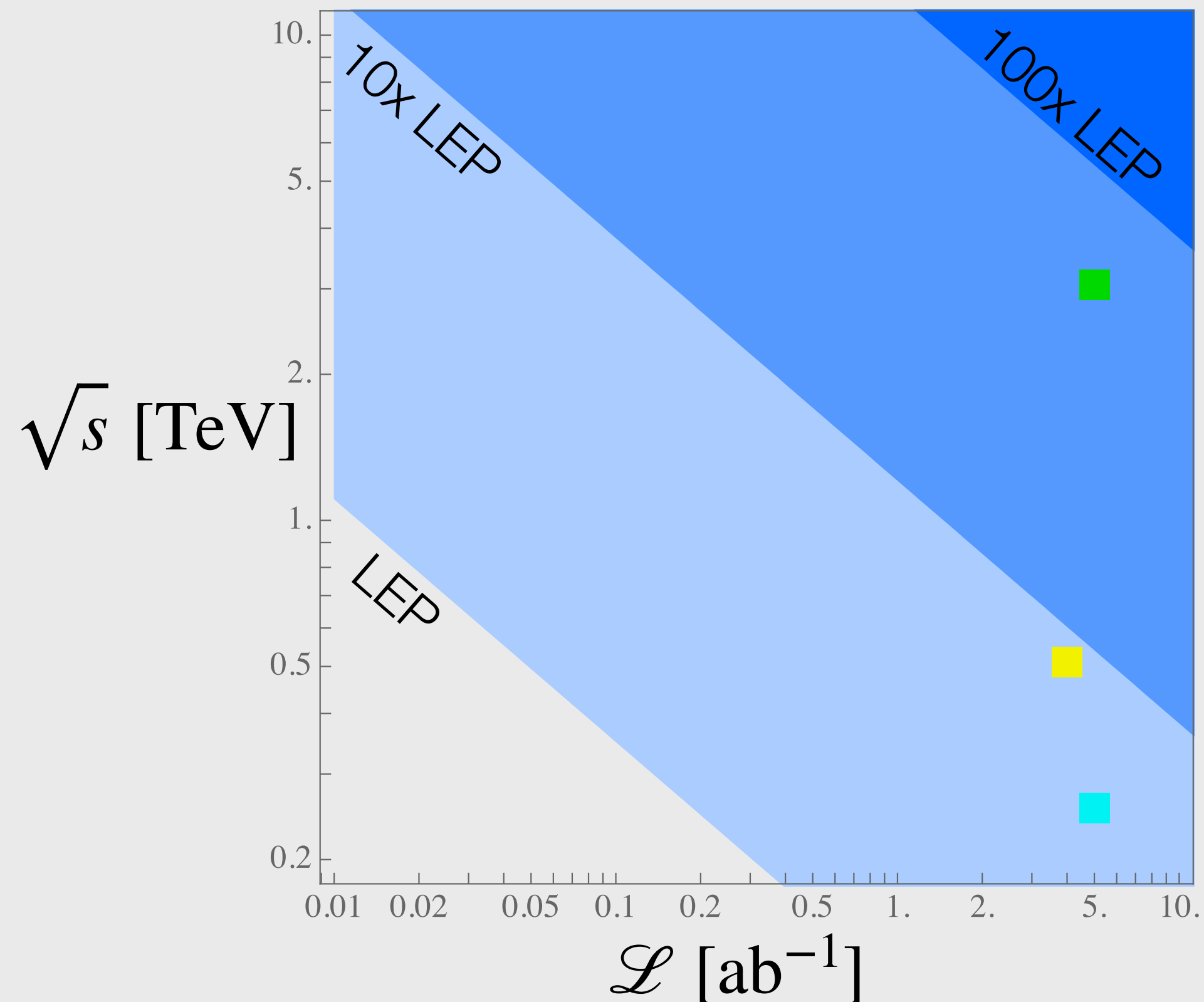
$$\hat{S} < 2 \cdot 10^{-4} \text{ (95 \% CL)} \quad \leftarrow \mathcal{L} = 5 \text{ ab}^{-1}$$





TOTAL **RATE** $\left| A_{SM}^{(00)} \right|^2 + A_{SM}^{00} \cdot A_{BSM}^{00} + \dots$

$$\hat{S}_{95\%} \lesssim 1.2 \cdot 10^{-4} \frac{1}{E_{beam}/\text{TeV}} \cdot \frac{1}{\sqrt{\mathcal{L}/\text{ab}^{-1}}}$$



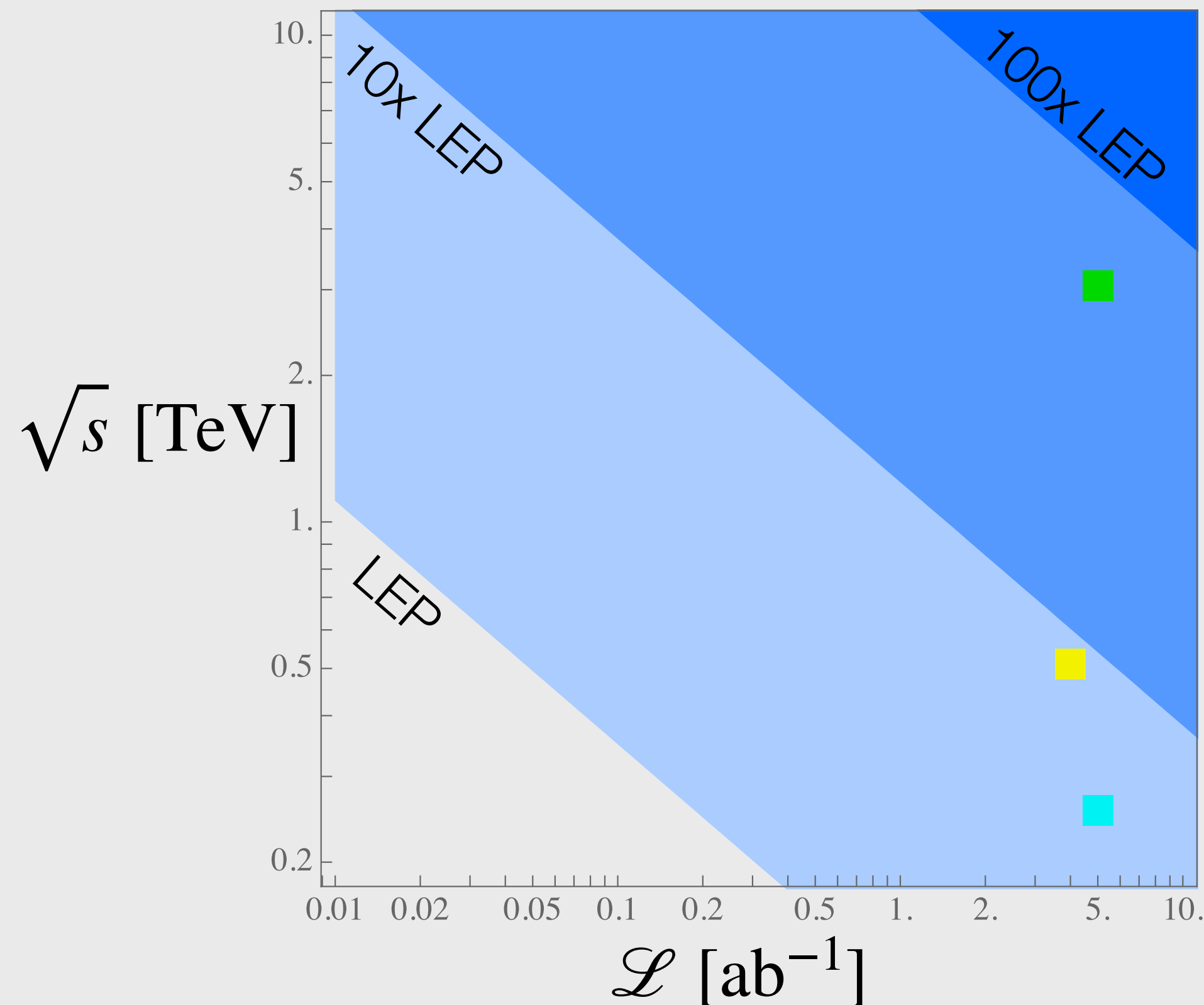
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Ever higher energy colliders can exploit “precise” measurements at the 10% level

$$\hat{S} \equiv c_W / m_W^2 \simeq \frac{\delta O}{O} \text{ at Z pole}$$

GOING TO HIGHER ENERGY WE CAN EXPLOIT “PRECISE” MEASUREMENTS AT THE 10% LEVEL, AVOIDING THE BOTTLENECK OF SYSTEMATIC UNCERTAINTIES



TOTAL RATE

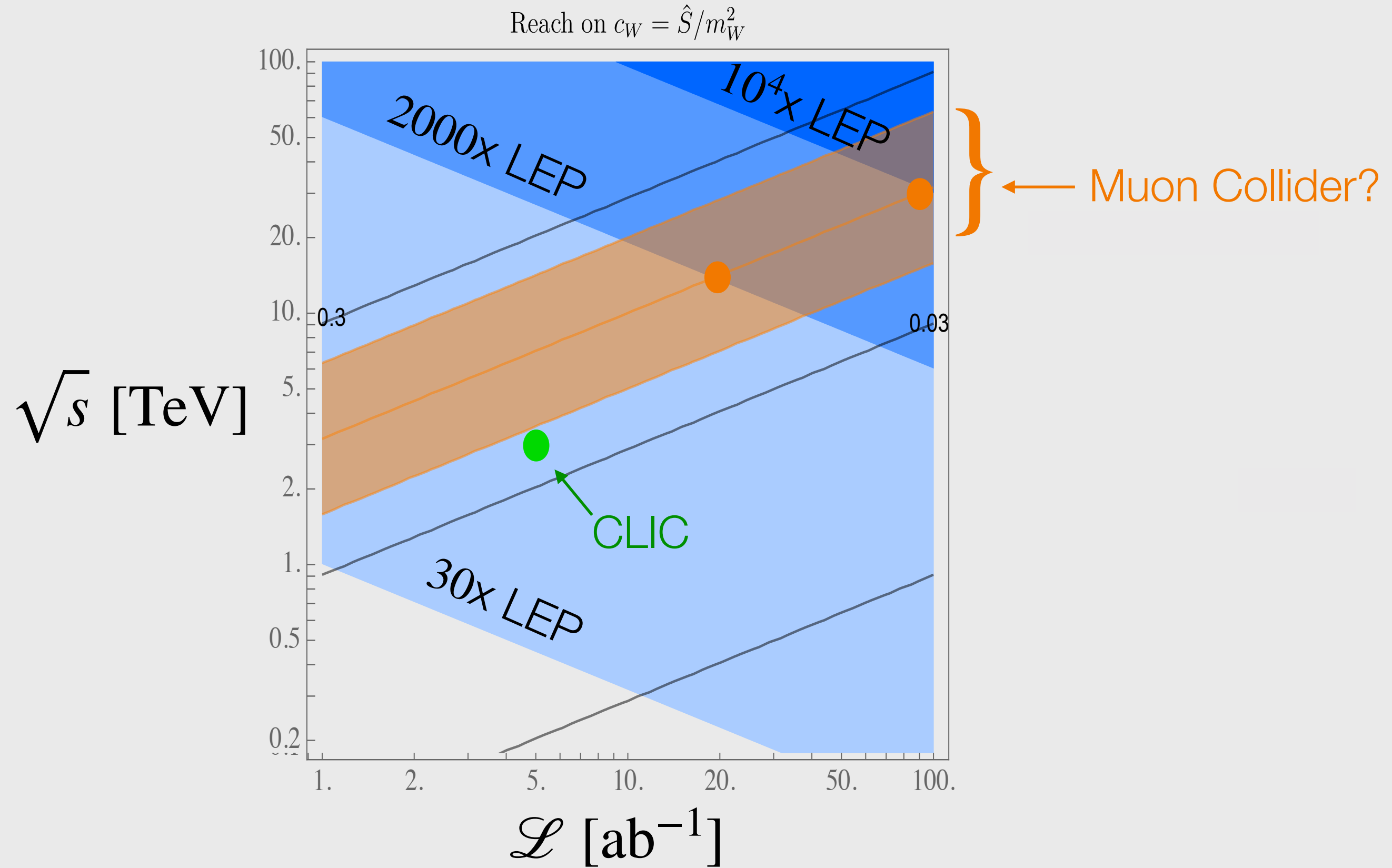
$$\left| A_{SM}^{(00)} \right|^2 + A_{SM}^{00} \cdot A_{BSM}^{00} + \dots$$

ZH	Cross-Section	m* 95% CL
3 TeV	1362 ab	12 TeV
14 TeV	62 ab	57 TeV
30 TeV	13 ab	120 TeV



Very important to keep up the lumi $\mathcal{L} \propto E^2$

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TOTAL

RATE

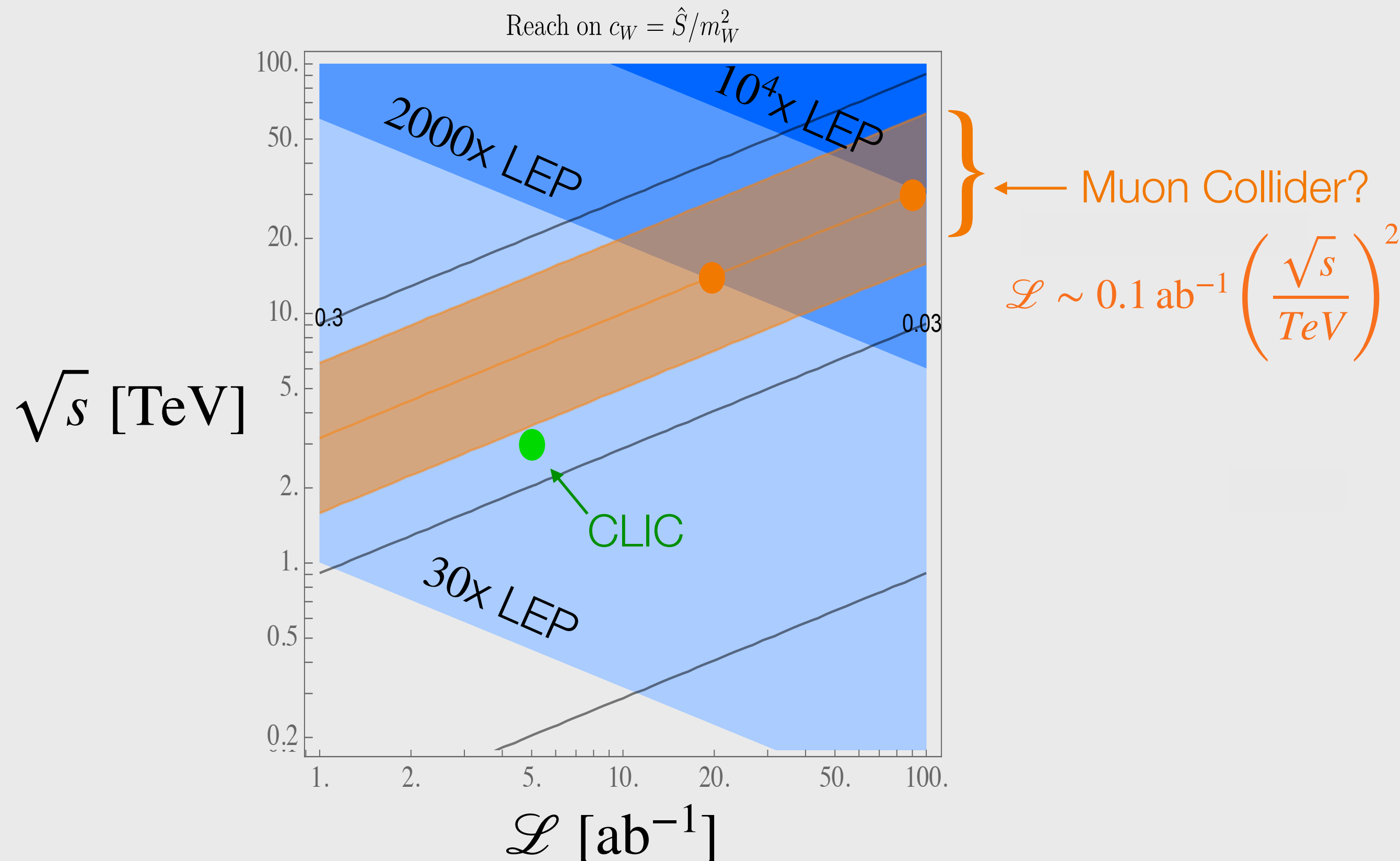
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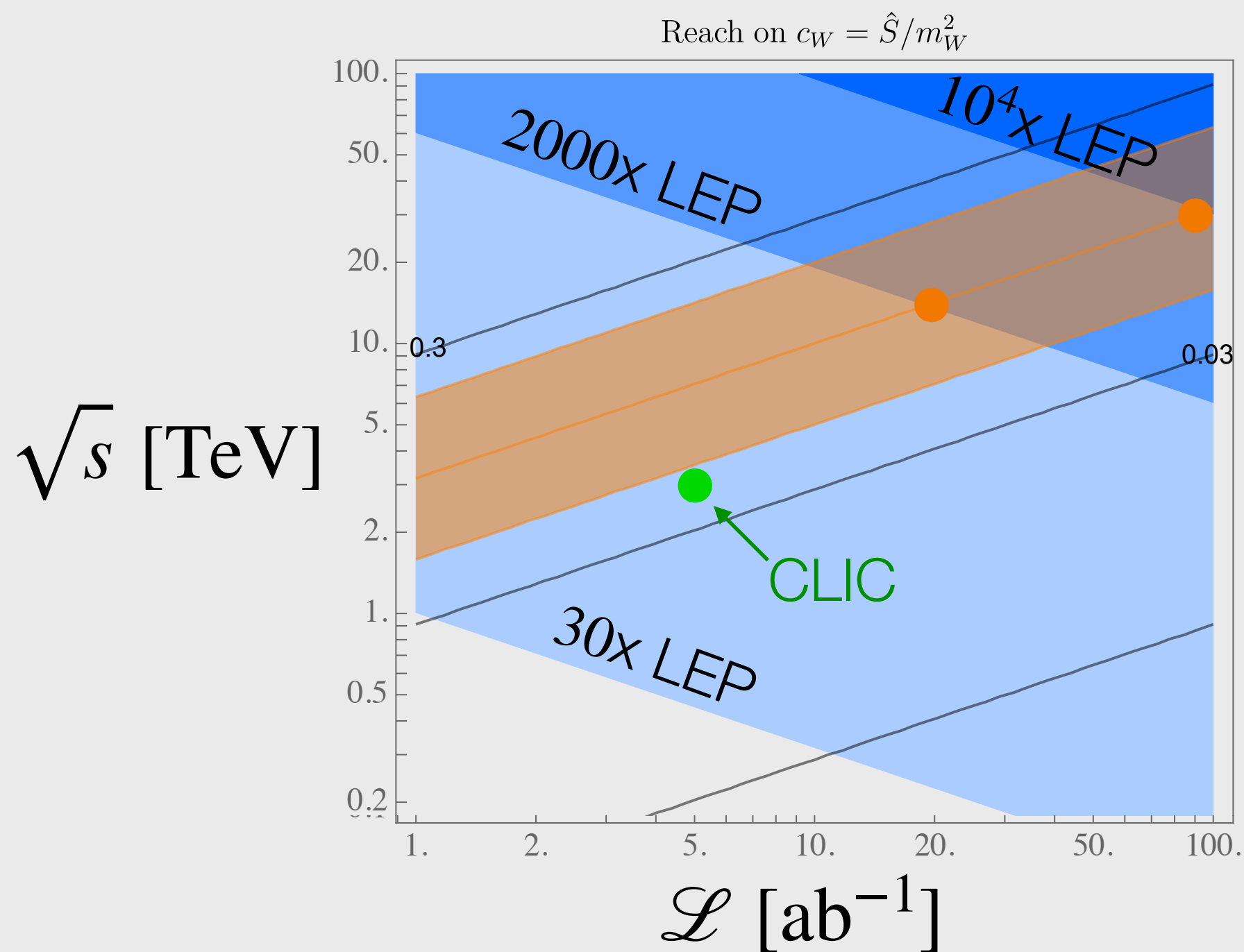


GOING TO HIGHER ENERGY WE CAN EXPLOIT “PRECISE” MEASUREMENTS AT THE 10% LEVEL, AVOIDING THE BOTTLENECK OF SYSTEMATIC UNCERTAINTIES



EFT EPOCH

LESSON FROM LHC

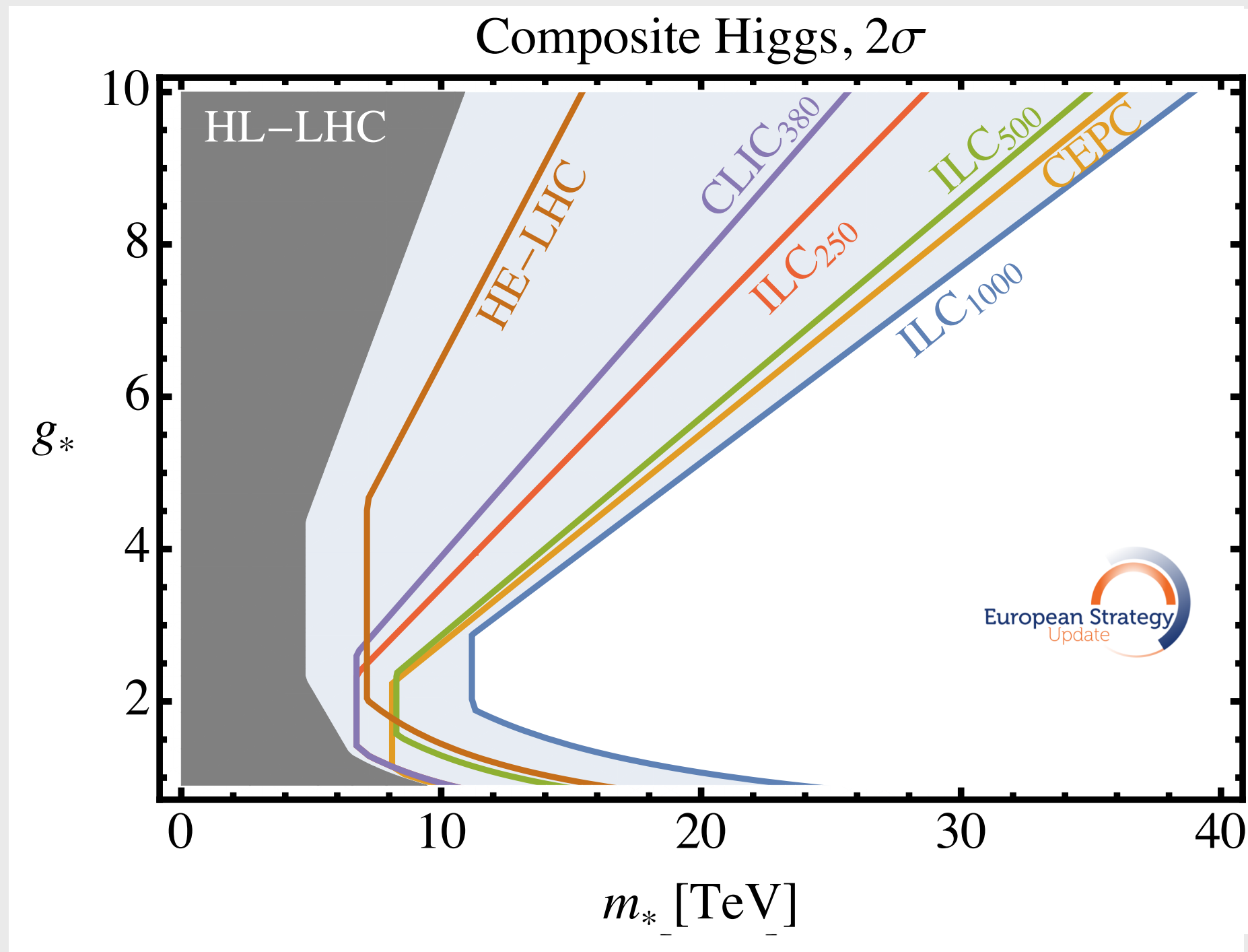


Orders of magnitude of improvement of the bounds on the mass scale of new physics

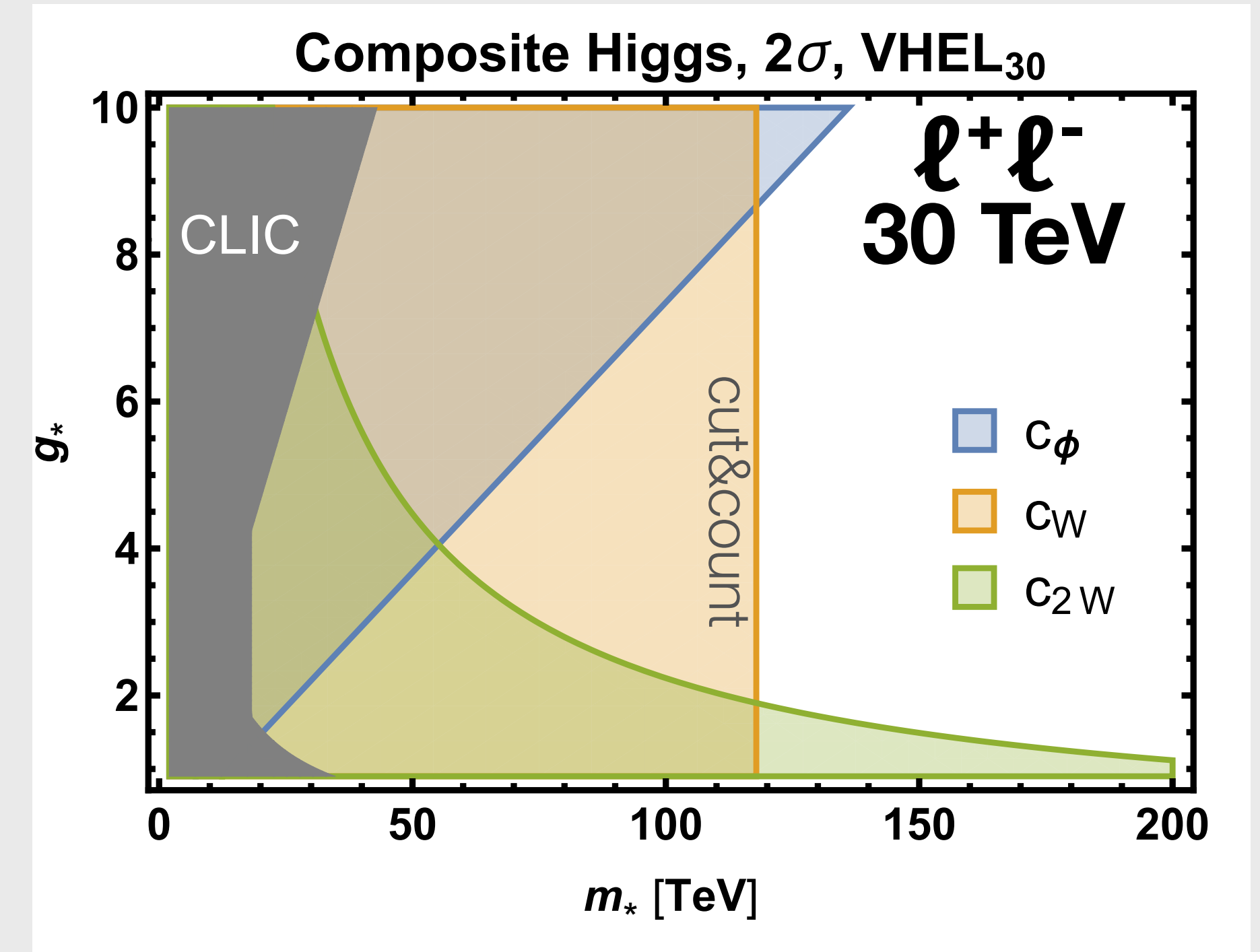
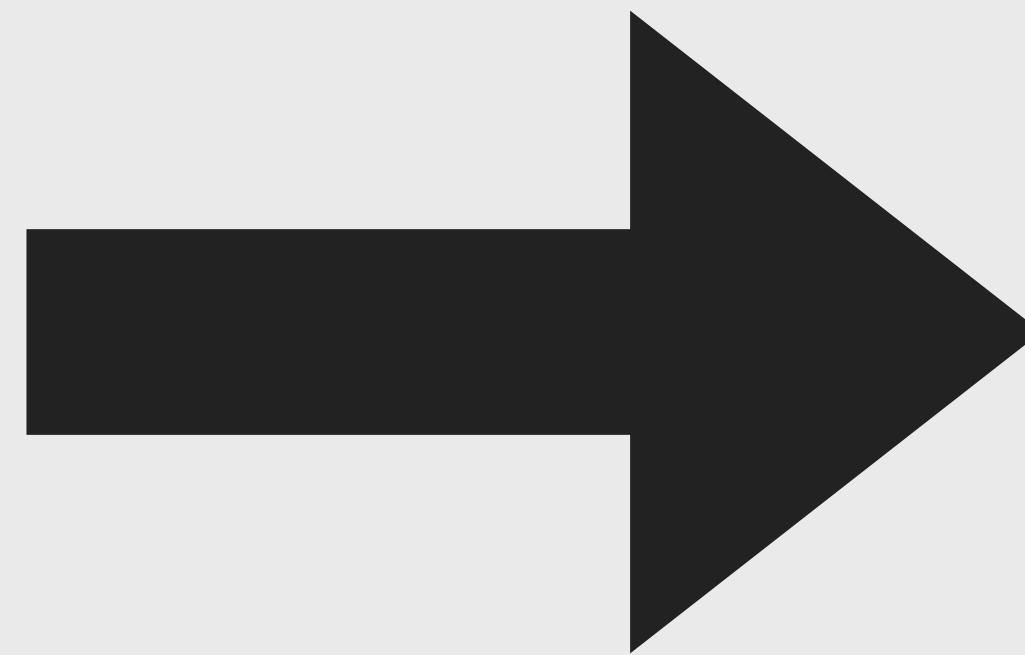
	Cross-Section @ 30 TeV	m^* 95% CL	
HH	non trivial c&c analysis	60 TeV ($g^*/4$)	\mathcal{O}_H
WW	pTW > 7.5 TeV: 180 ab	84 \oplus 76 TeV \approx 113 TeV	$a_q^{(3)}$
ZH	inclusive: 13 ab	120 TeV	$a_q^{(1)}$
ff	angular analysis	120 TeV ($4/g^*$)	W, Y

All-round progress up to $m^* \sim 10^3 m_{\text{Higgs}}$

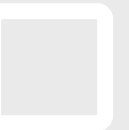
Looking ahead



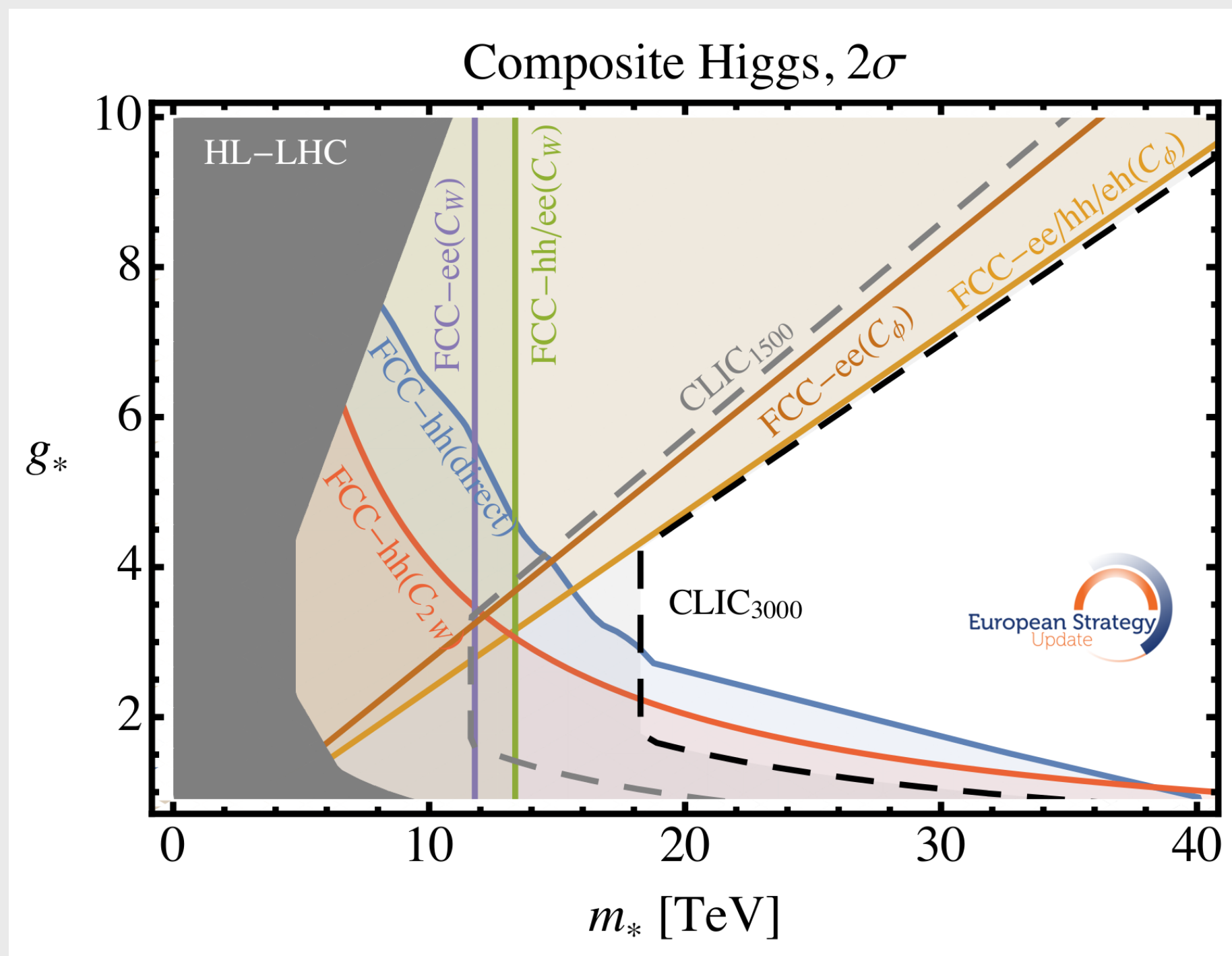
compositeness at
10 TeV-20 TeV



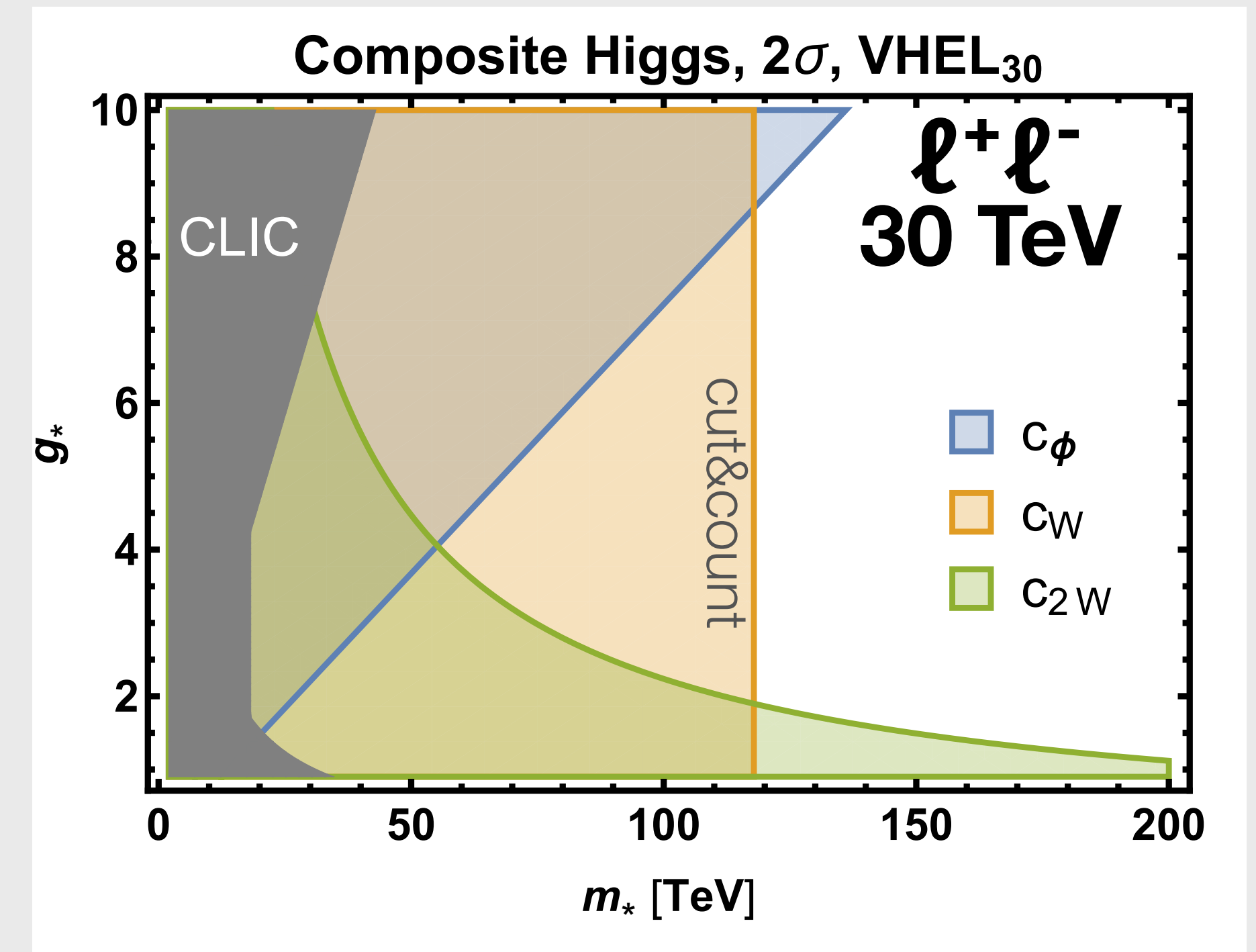
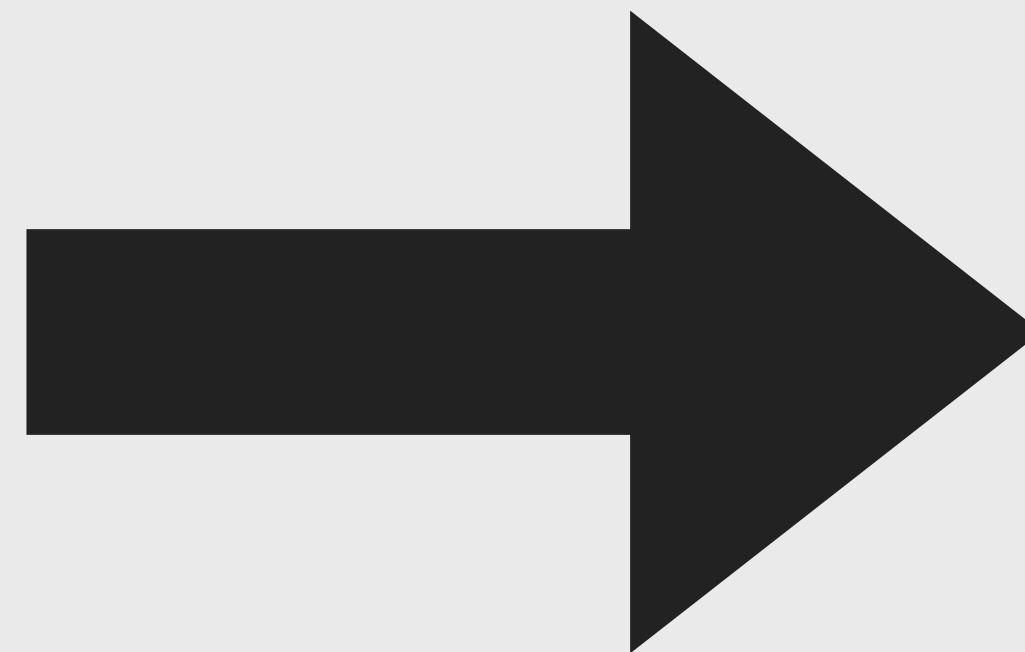
compositeness at
100 TeV-500 TeV



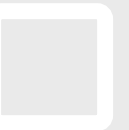
Looking ahead



**compositeness at
10 TeV-20 TeV**

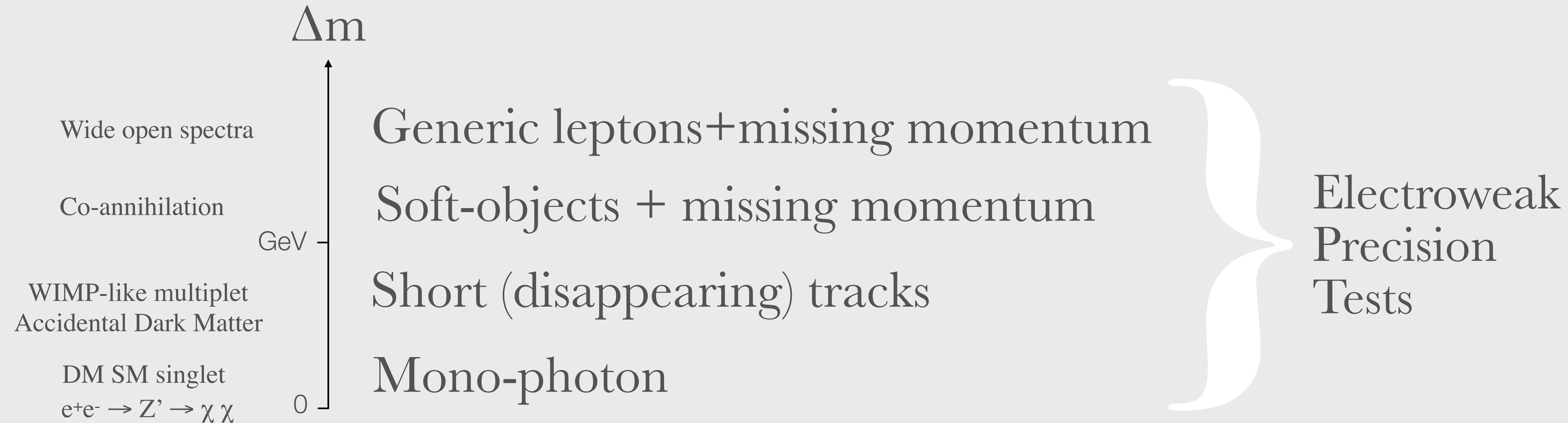


**compositeness at
100 TeV-500 TeV**

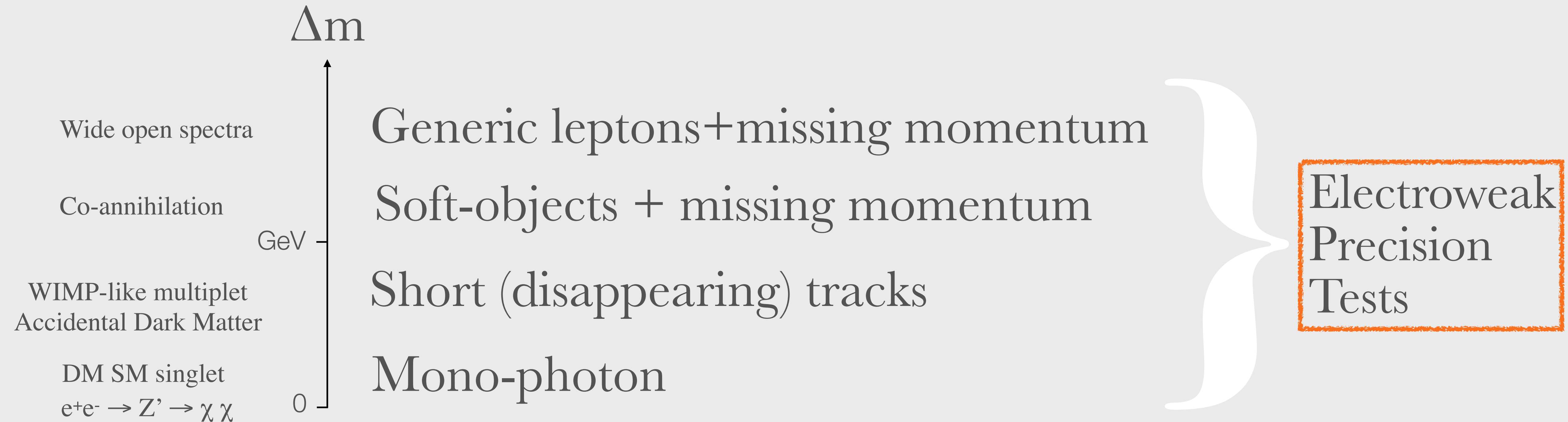


Wimp Dark Matter

Electroweak Dark Matter: LSP (+NLSP)



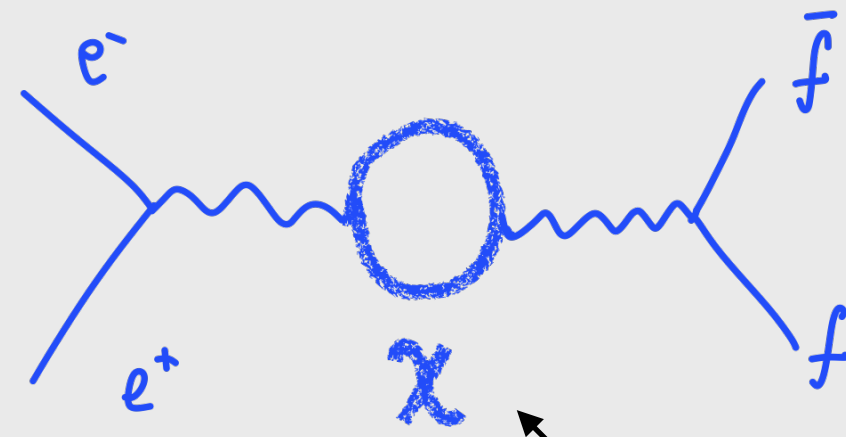
Electroweak Dark Matter: LSP (+NLSP)



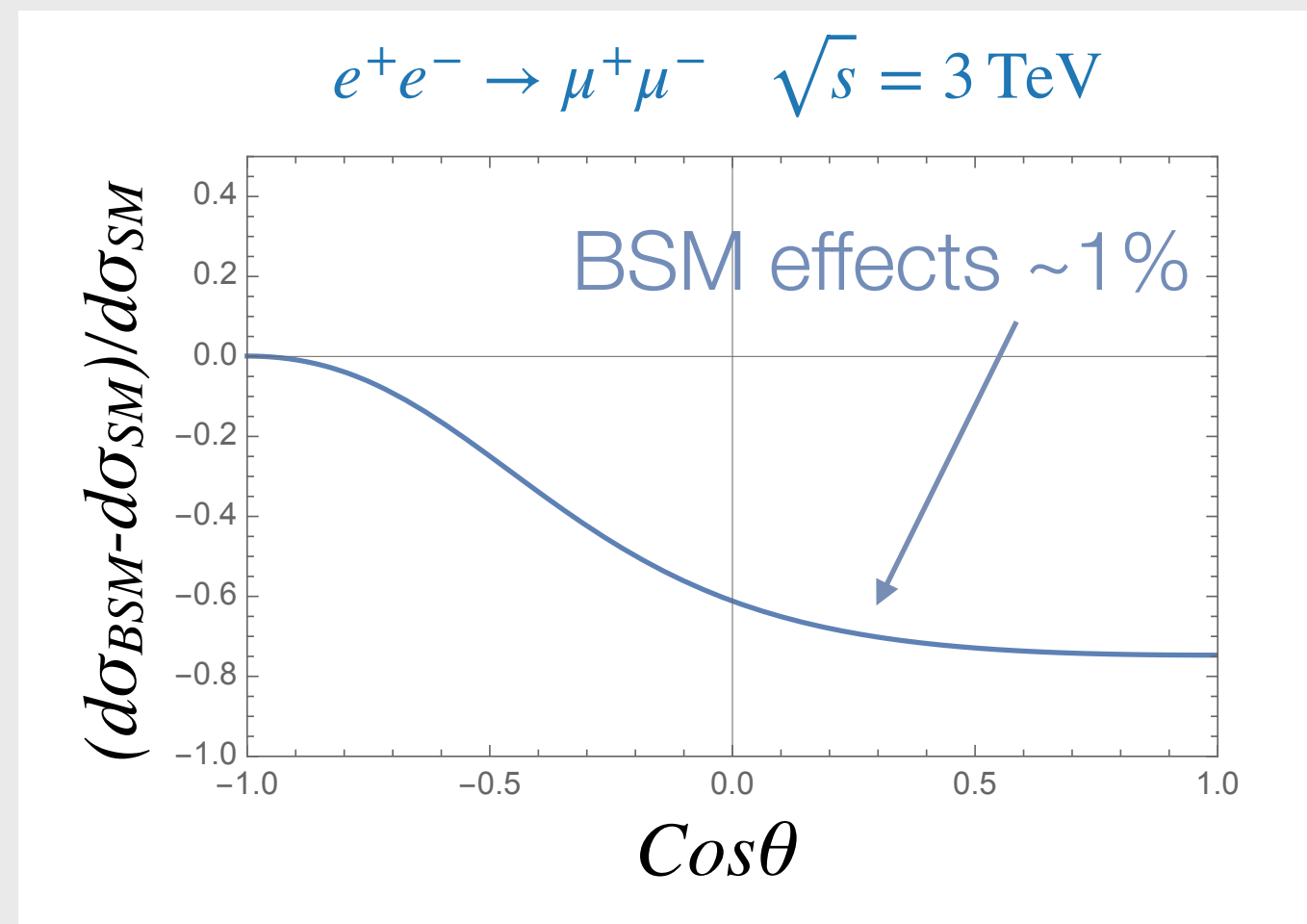
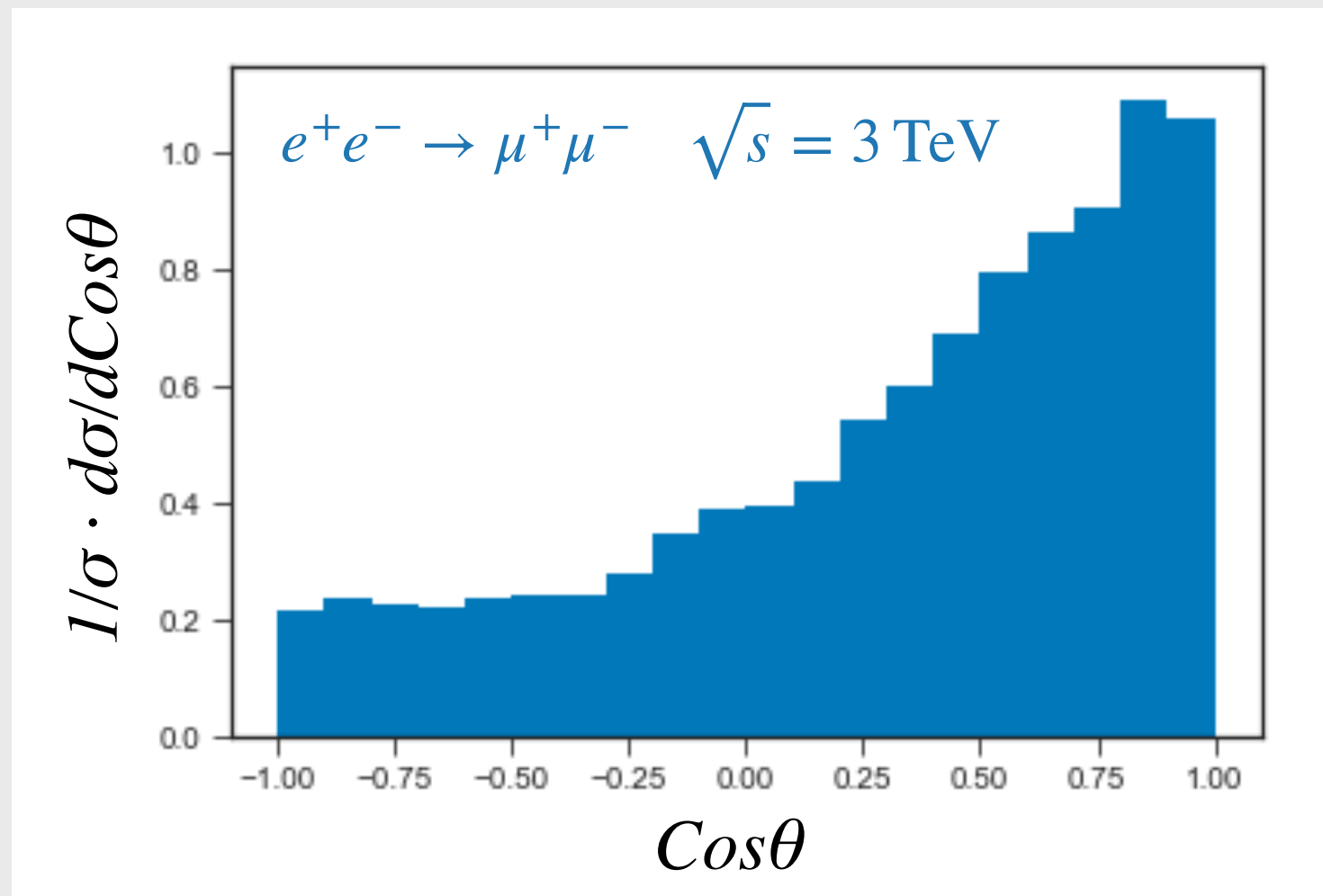
$$e^+ e^- \rightarrow f \bar{f}$$

PRECISION

ANGULAR DISTRIBUTION



χ is heavy/light new physics

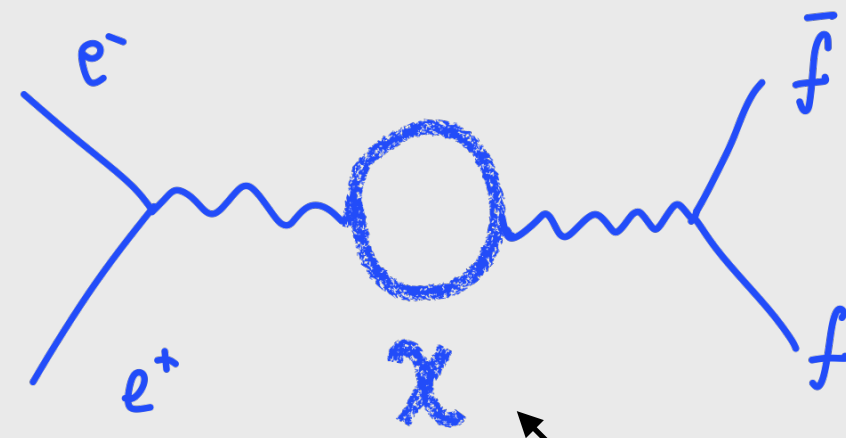


beams polarization is beneficial to increase NP effects

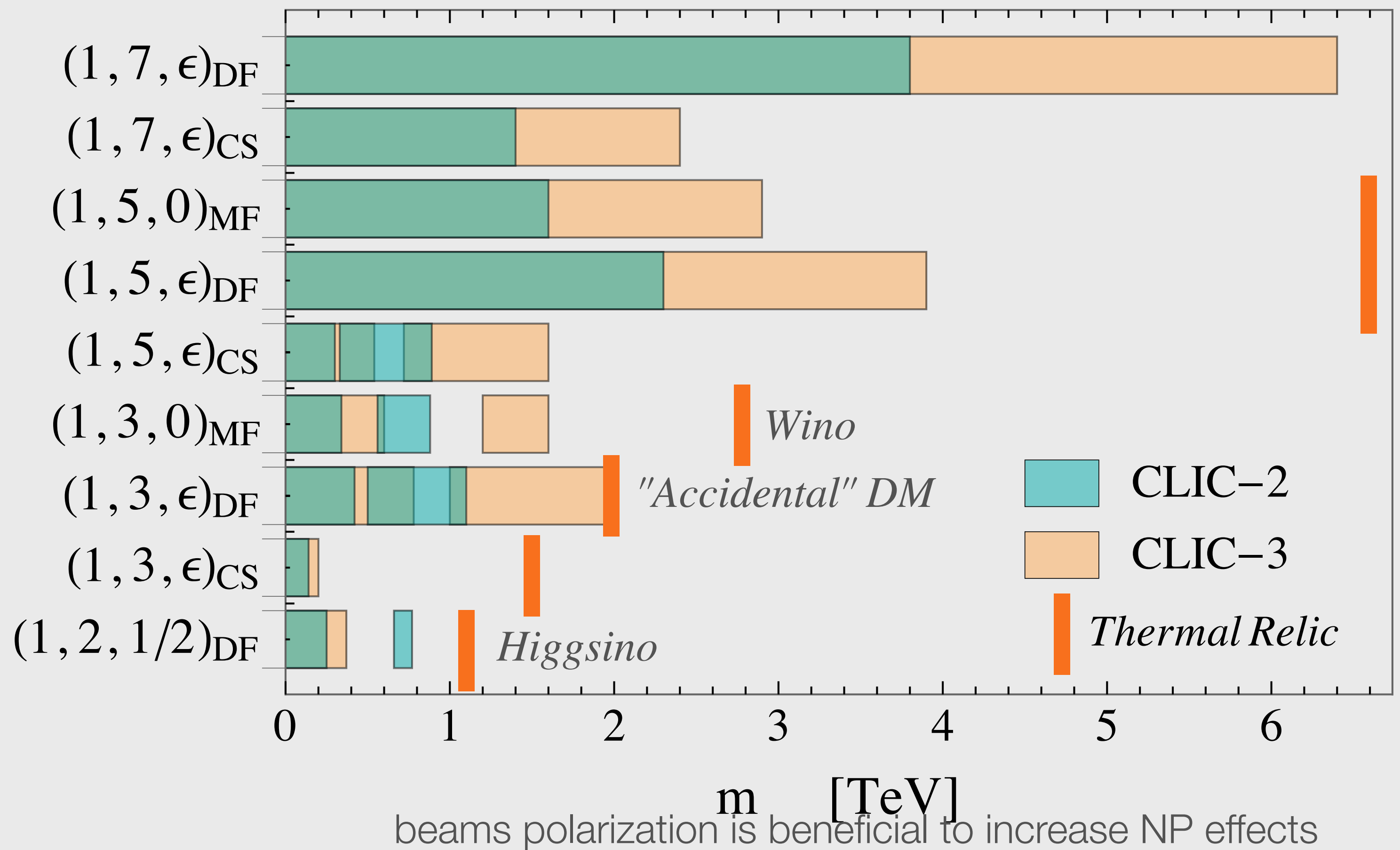
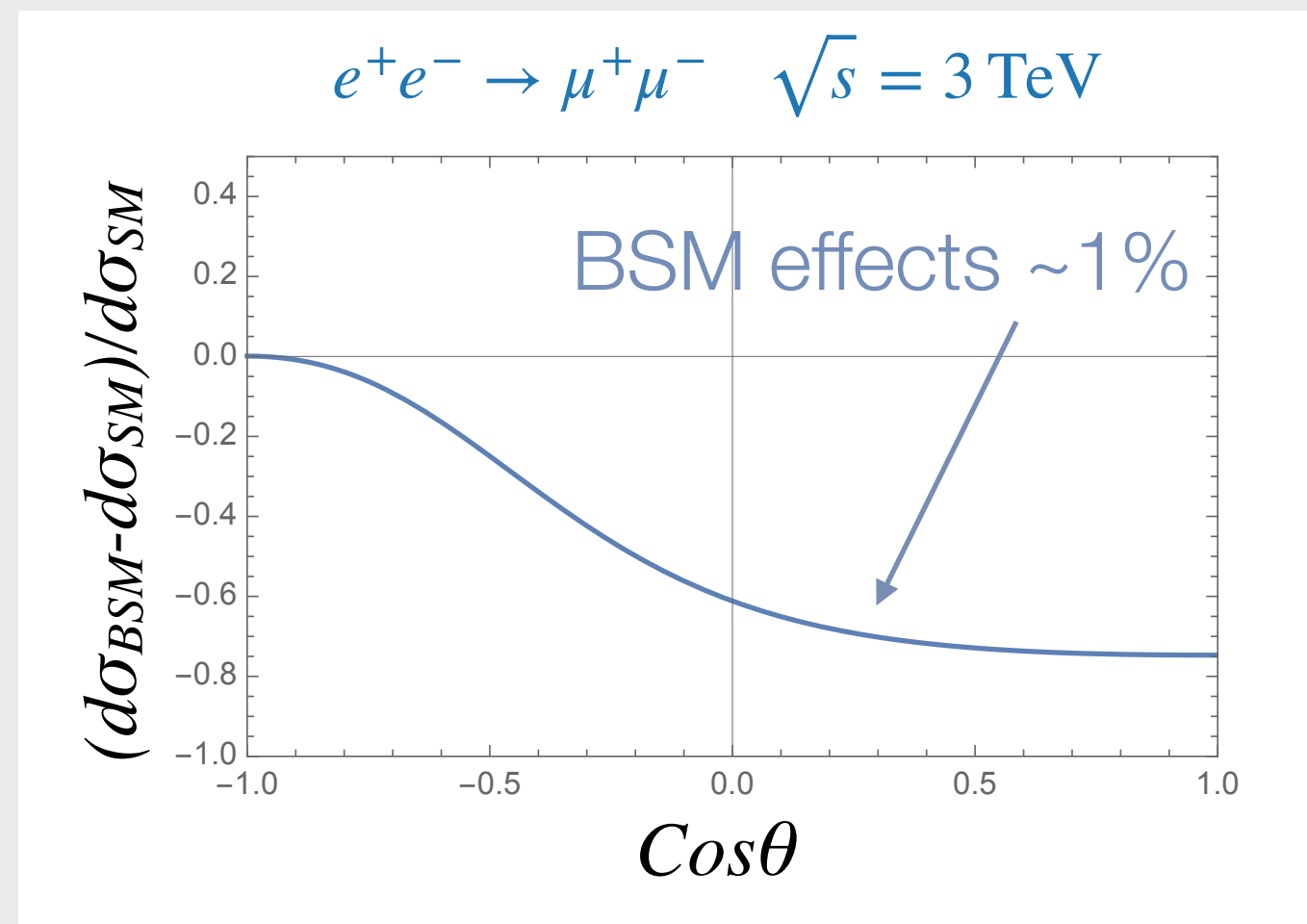
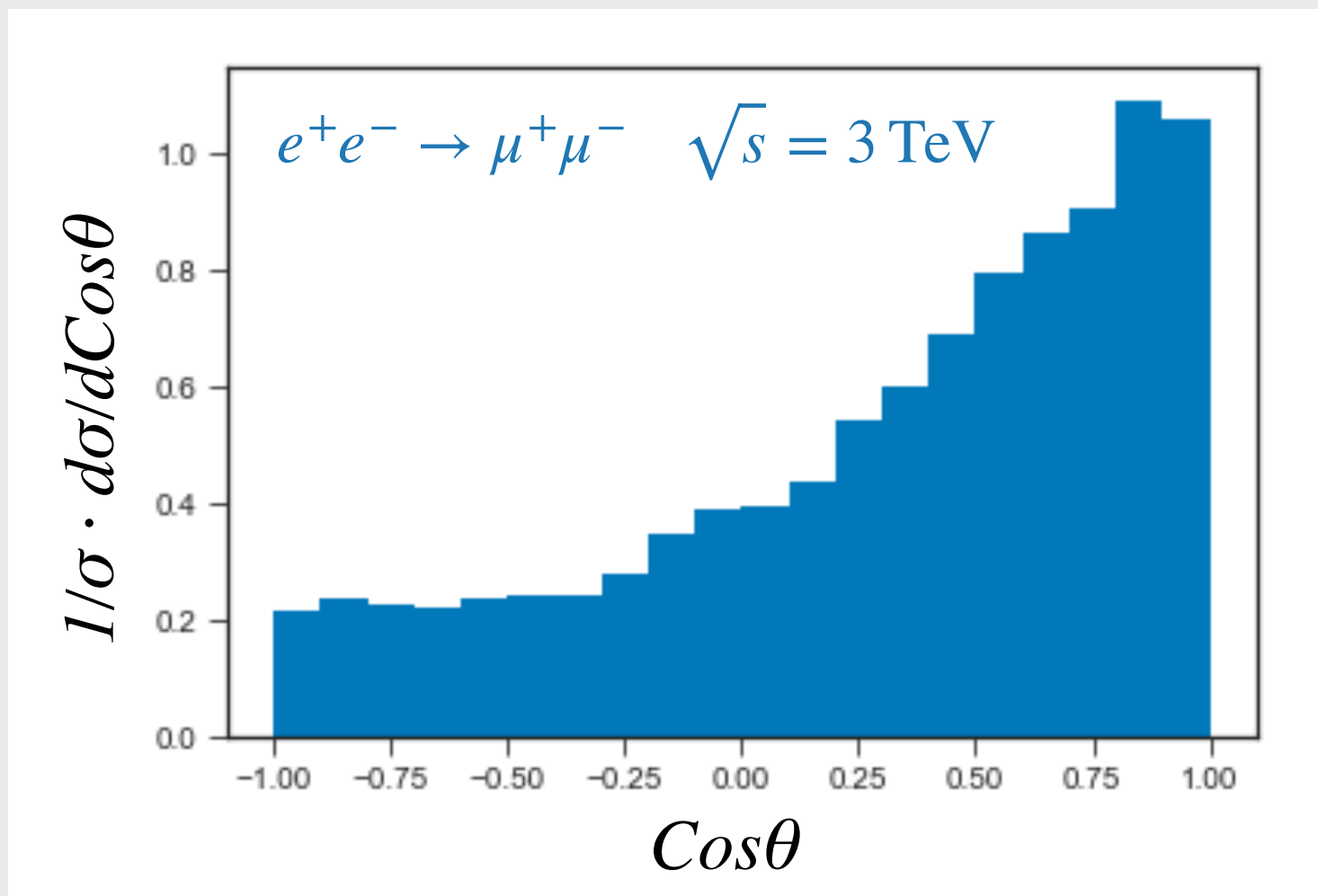
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PRECISION

ANGULAR DISTRIBUTION



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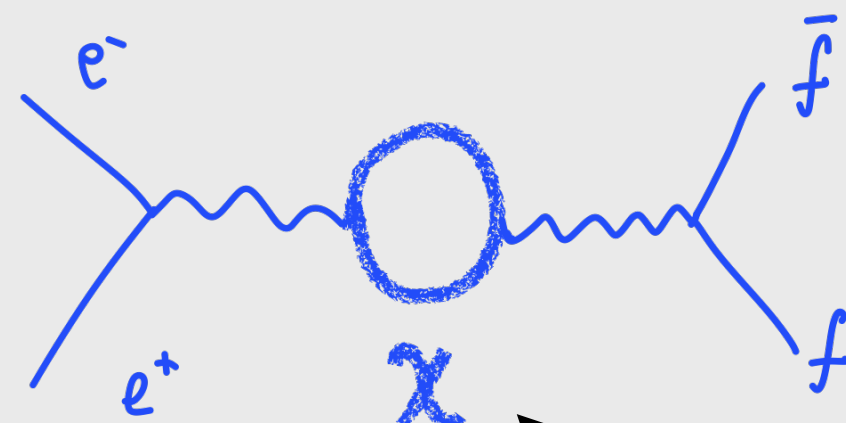


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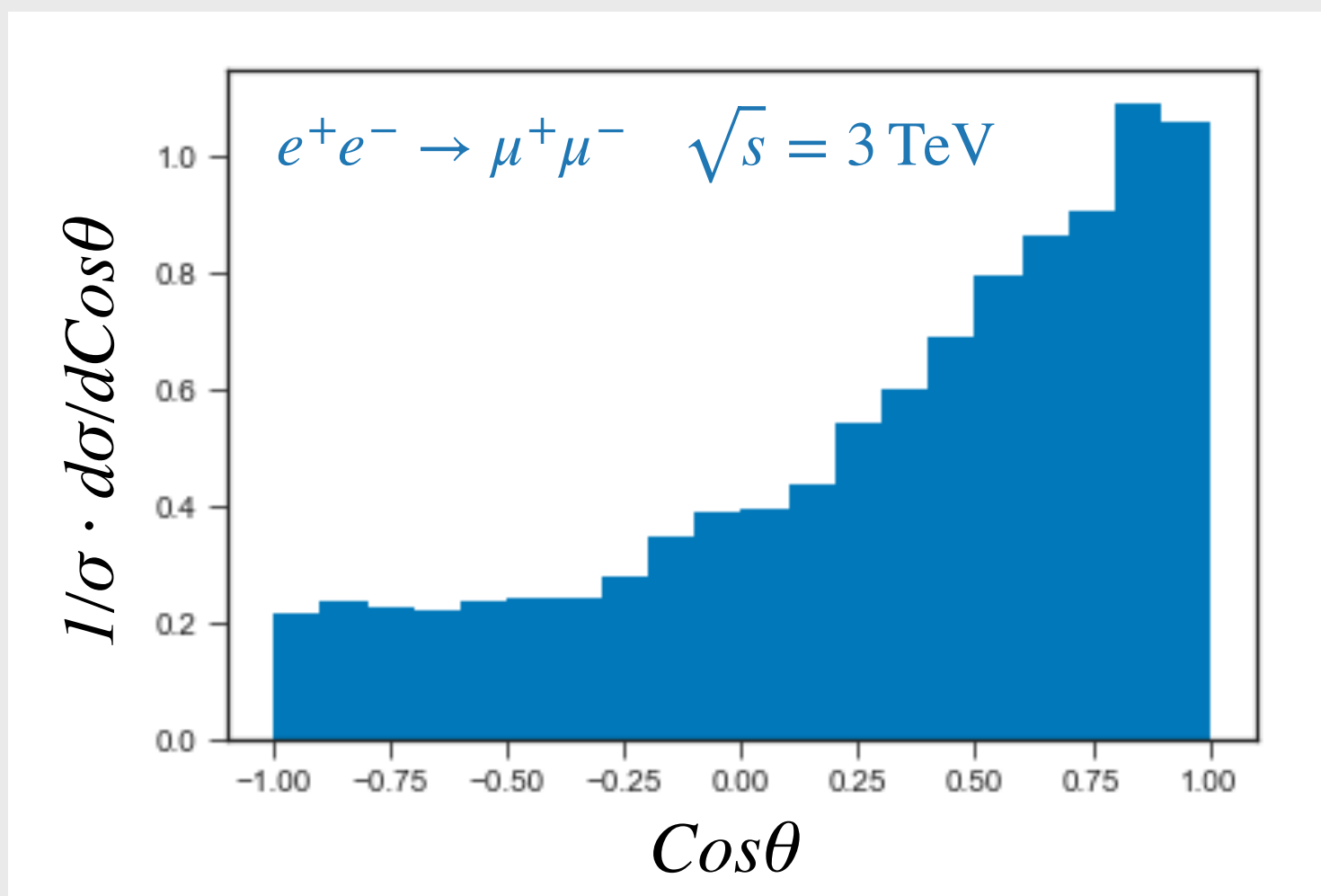
$$e^+ e^- \rightarrow f \bar{f}$$

PRECISION

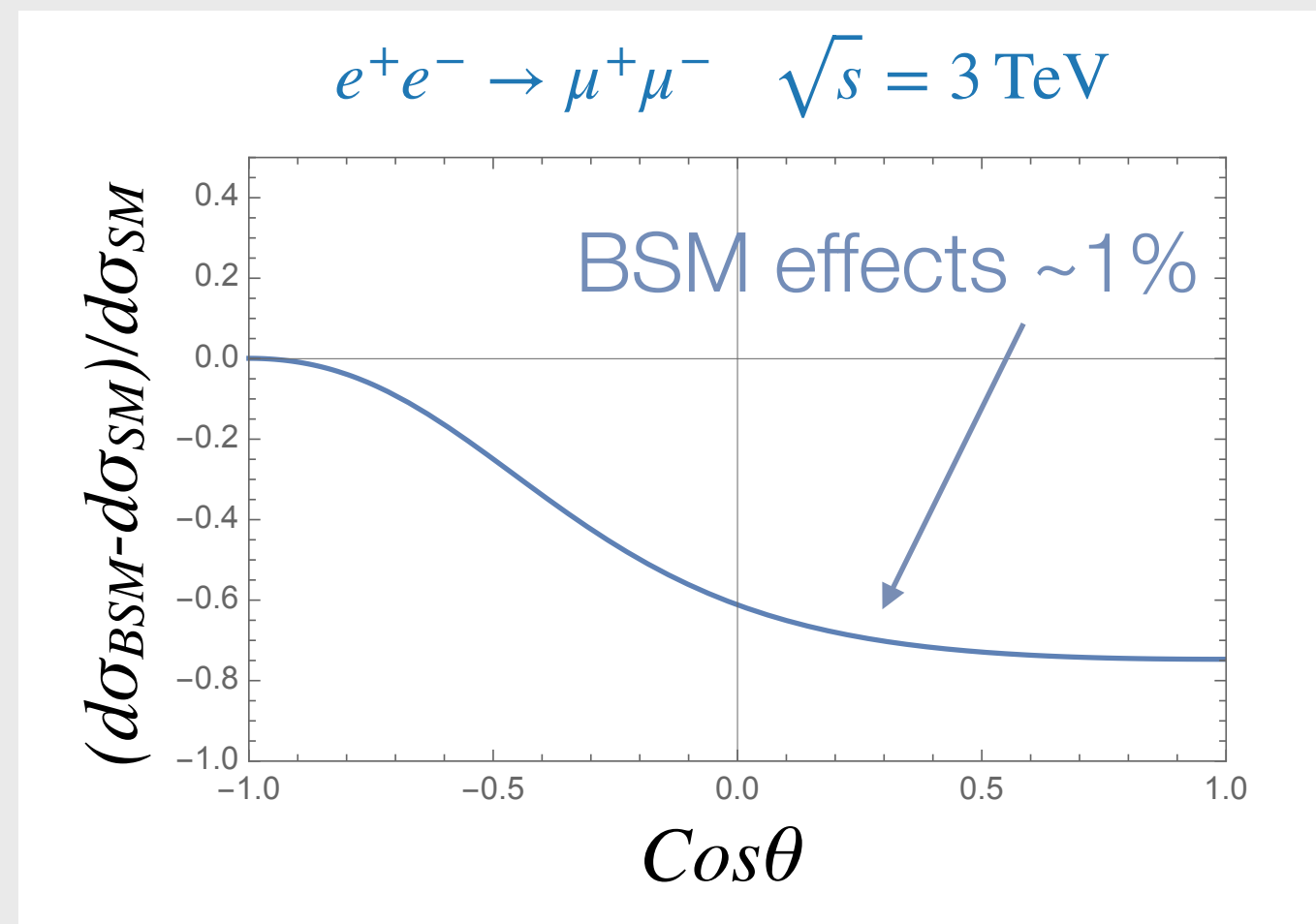
ANGULAR DISTRIBUTION



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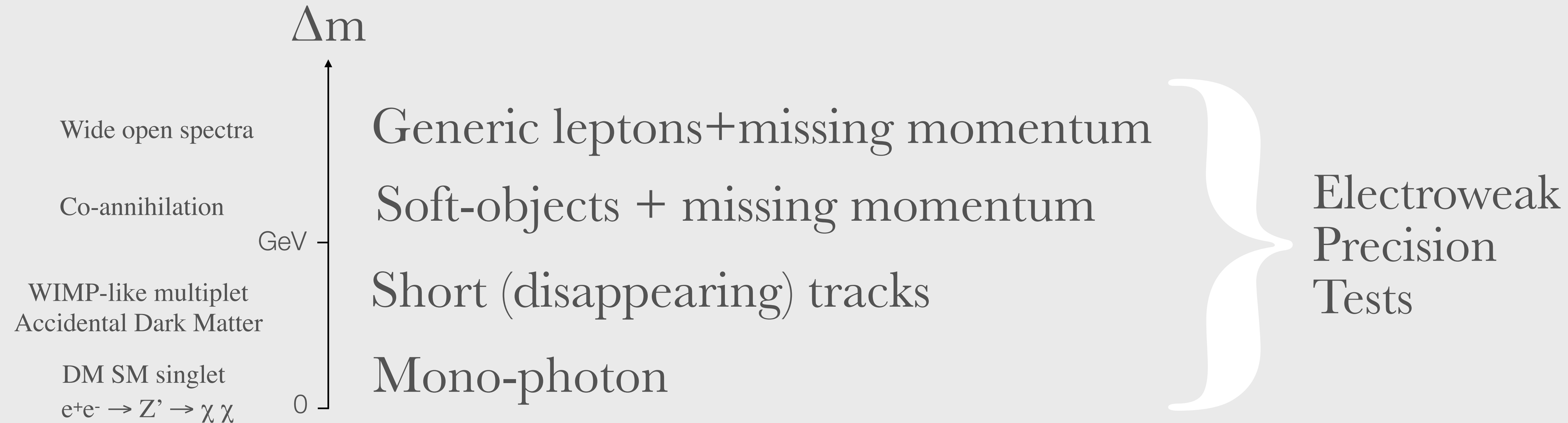
χ / m_χ [TeV]	DM	HL-LHC	HE-LHC	FCC-100	CLIC-3	Muon-14
$(1, 2, 1/2)_{DF}$	1.1	—	—	—	0.4	0.6
$(1, 3, \epsilon)_{CS}$	1.6	—	—	—	0.2	0.2
$(1, 3, \epsilon)_{DF}$	2.0	—	0.6	1.5	0.8 & [1.0, 2.0]	2.2 & [6.3, 7.1]
$(1, 3, 0)_{MF}$	2.8	—	—	0.4	0.6 & [1.2, 1.6]	1.0
$(1, 5, \epsilon)_{CS}$	6.6	0.2	0.4	1.0	0.5 & [0.7, 1.6]	1.6
$(1, 5, \epsilon)_{DF}$	6.6	1.5	2.8	7.1	3.9	11
$(1, 5, 0)_{MF}$	14	0.9	1.8	4.4	2.9	3.5 & [5.1, 8.7]
$(1, 7, \epsilon)_{CS}$	16	0.6	1.3	3.2	2.4	2.5 & [3.5, 7.4]
$(1, 7, \epsilon)_{DF}$	16	2.1	4.0	11	6.4	18



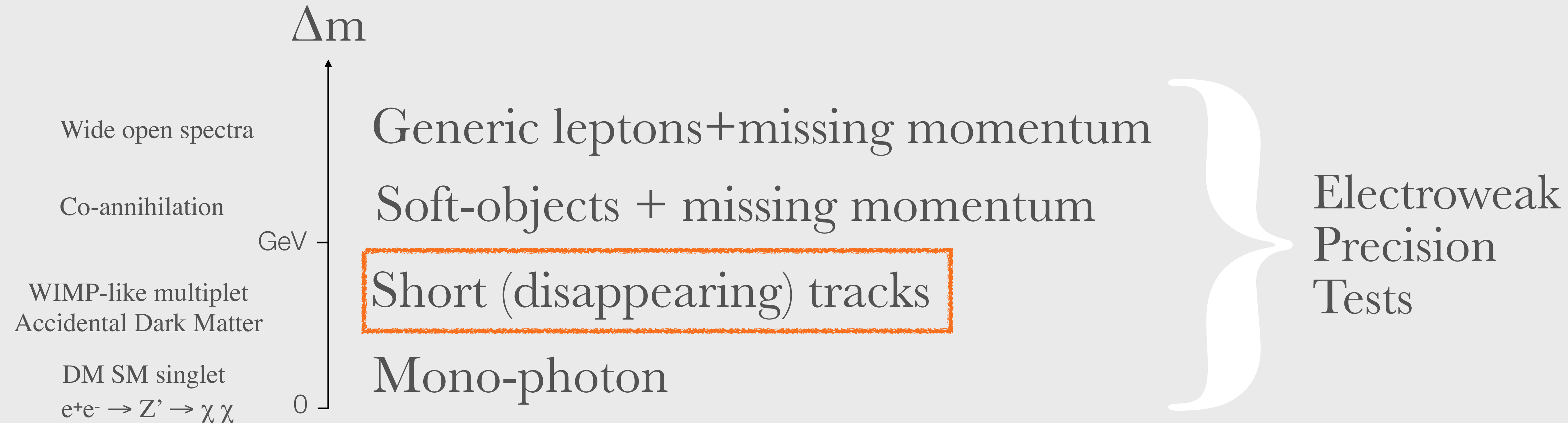
- Comprehensive tool to explore new electroweak particles
- Can probe valid dark matter candidates!

beams polarization is beneficial to increase NP effects

Electroweak Dark Matter: LSP (+NLSP)



Electroweak Dark Matter: LSP (+NLSP)

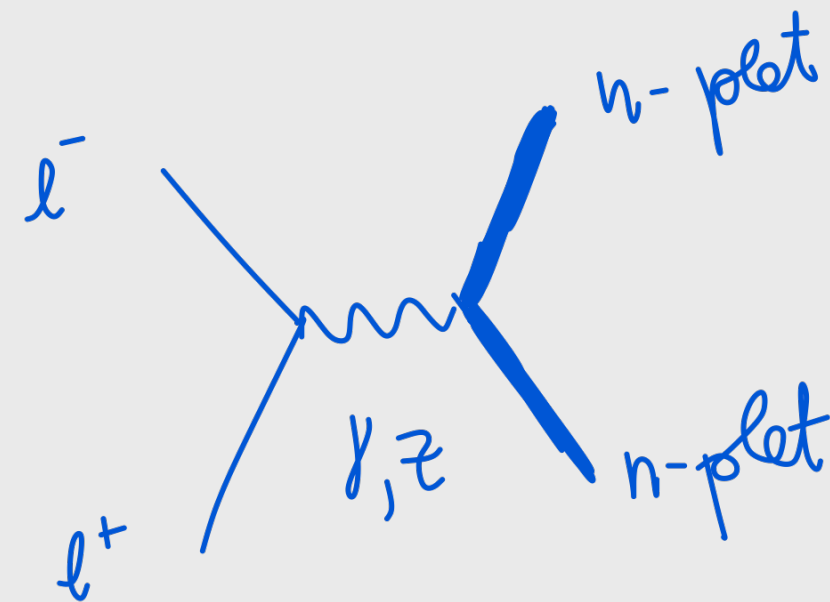


Degenerate EW multiplets

STUB-TRACKS

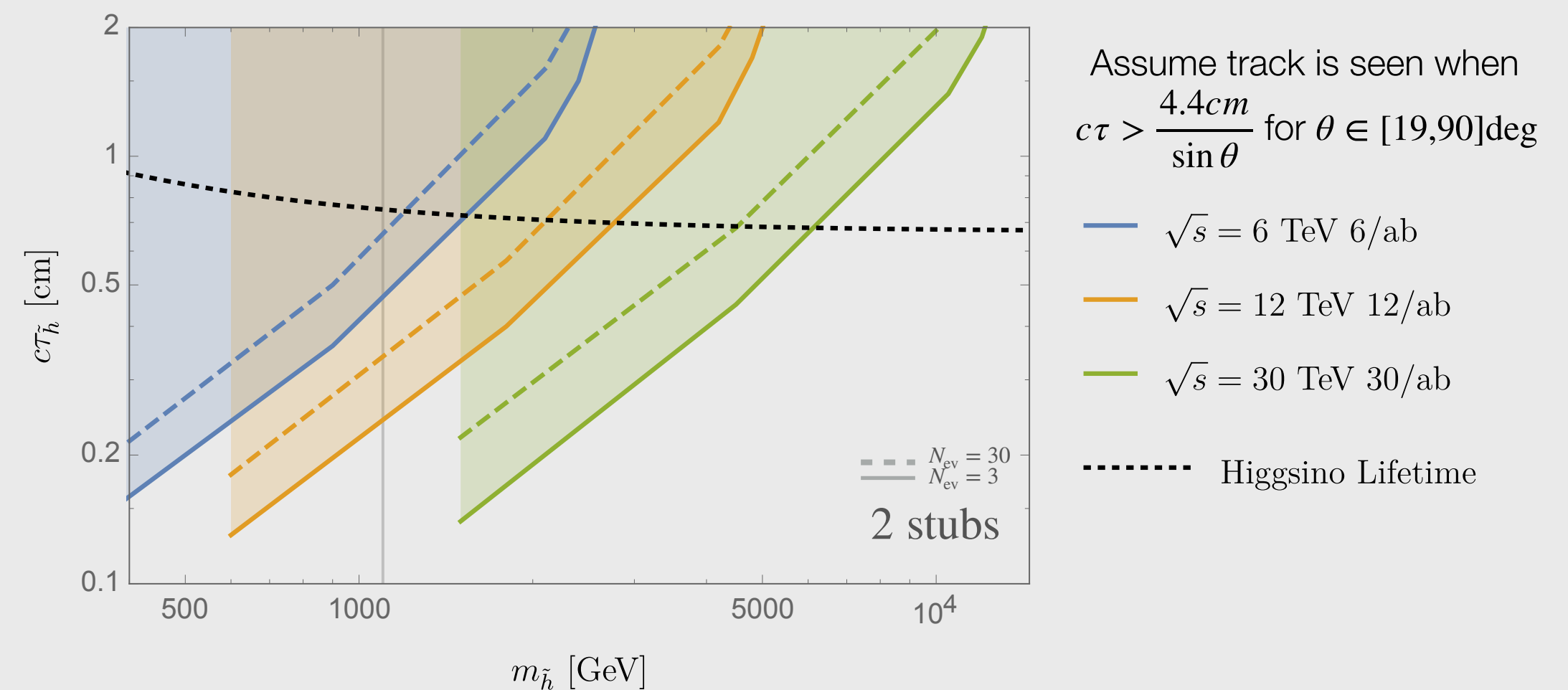
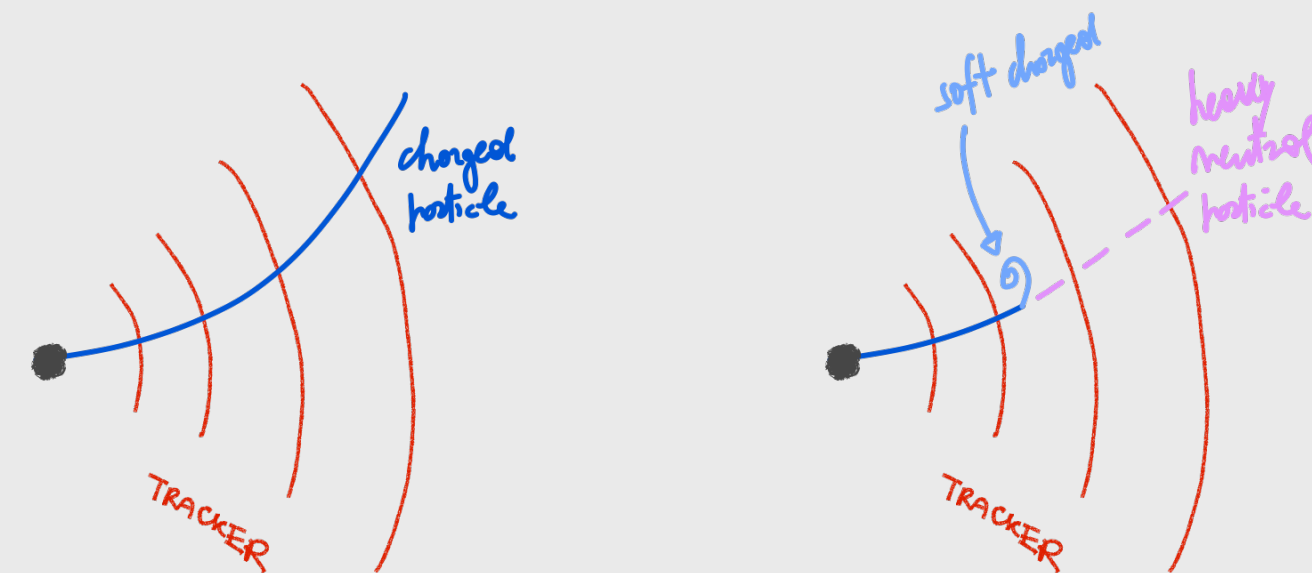
EXTRAPOLATION FROM CLIC

- Heavy n-plet of SU(2)
- Mass splitting $\sim \alpha_w m_W \sim 0.1 \text{ GeV} - \text{GeV}$





LARGE RATES, BUT NEEDS TO LIGHT UP THE DETECTOR IN A DISCERNIBLE WAY

- Heavily subject to detector design issues
- Even in CLIC needs full detector simulation



Full detector simulation study



Long-lived particles at CLIC

Ulrike Schnoor (CERN)
Erica Brondolin, Cecilia Ferrari, Emilia Leogrande
on behalf of the CLICdp collaboration

LCWS 2019

Ulrike Schnoor (CERN) Long-lived particles at CLIC - LCWS 2019 1 / 12



Long-lived particle reconstruction at CLIC

30 Oct 2019, 15:08

Parallel : Physics/Detect...

20m

Room1 (Conference building)

Speaker

Ulrike Schnoor

Primary author

Ulrike Schnoor

Presentation Materials

UlrikeSchnoor_chargedLLP_CLIC3TeV.pdf



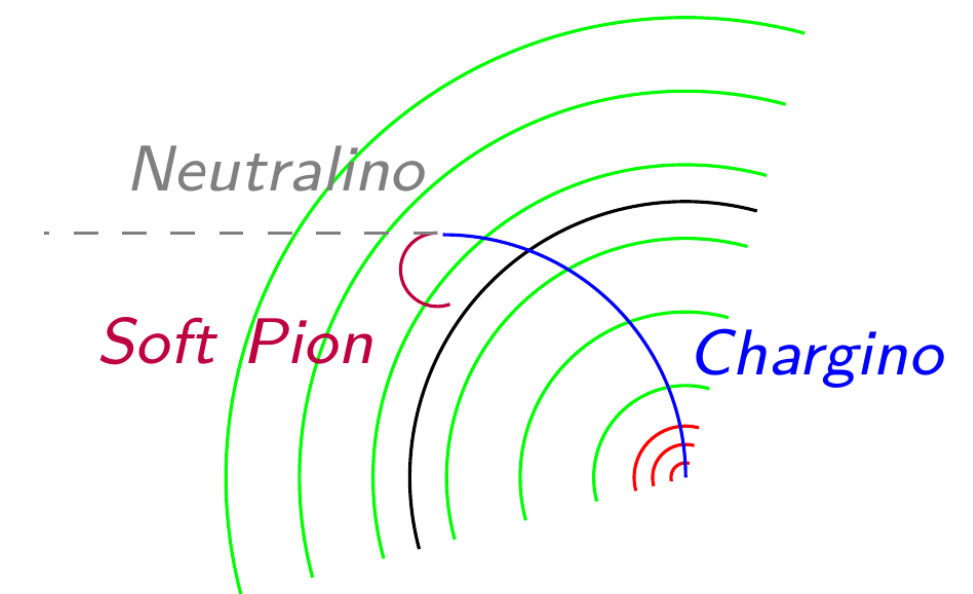
Analysis strategy

Stub track analysis at 3 TeV with CLICdet



Signal selection

- ▶ Stub track candidate definition:
 - ▶ at least four hits in the tracking system
 - ▶ disappearing within the tracking system volume
 - ▶ no energy deposition in the calorimeter
 - ▶ prompt, isolated track
 - ▶ minimum transverse momentum
 - ▶ dE/dx requirement
- ▶ At least one stub candidate per event
- ▶ Additional: Requirements on soft displaced pion(s)
- ▶ Additional: Requirements on additional photons




Backgrounds:

- ▶ Beam-induced $\gamma\gamma \rightarrow$ hadrons:
 - ▶ algorithmic
 - ▶ split tracks
 - ▶ conversion
- ▶ final states with low multiplicity of isolated leptons

Full detector simulation study





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Ulrike Schnoor (CERN) Long-lived particles at CLIC - LCWS 2019 1 / 12

International Workshop on Future Linear Colliders
LCWS2019 Sendai
October 28 - November 1

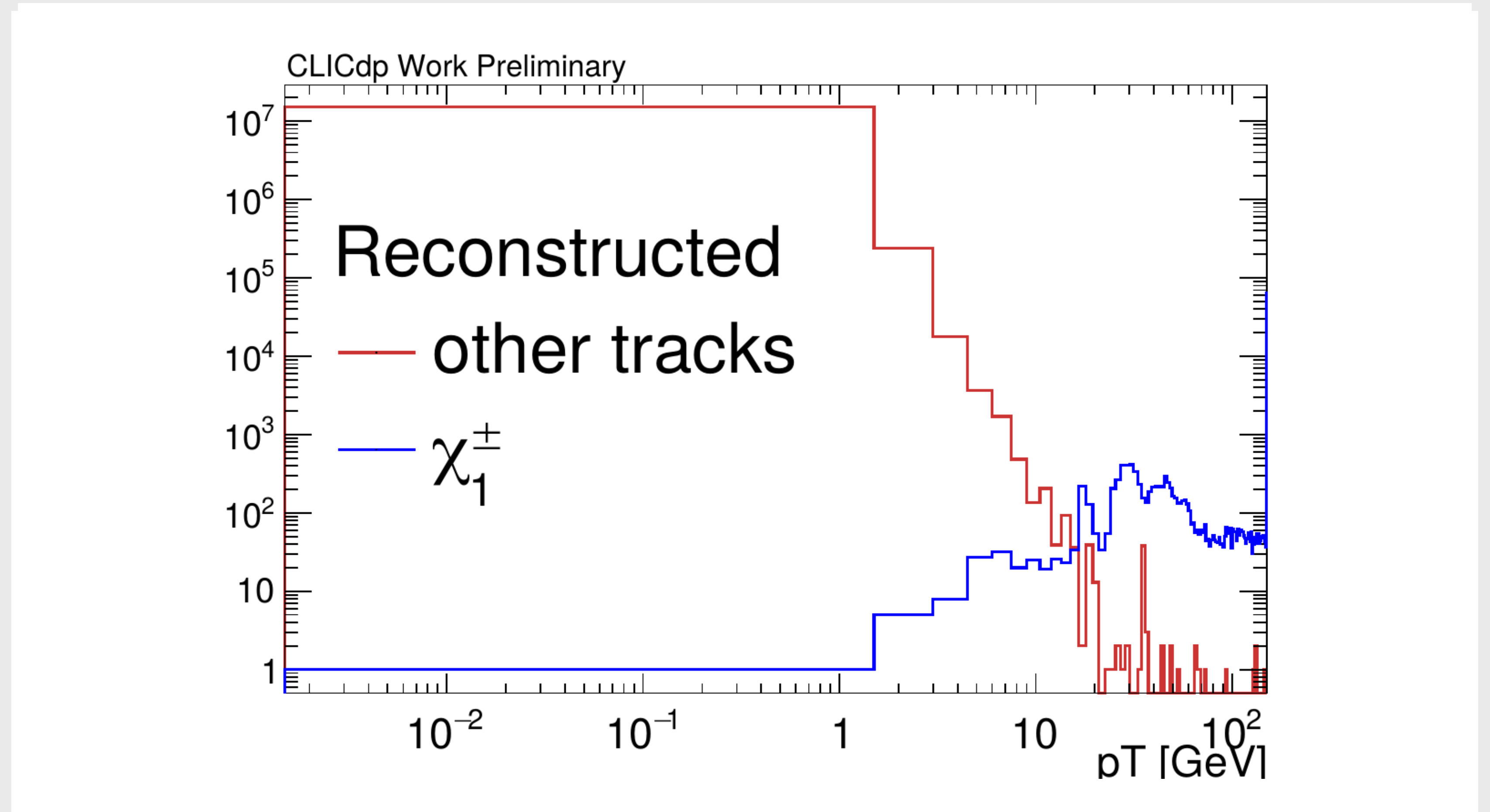
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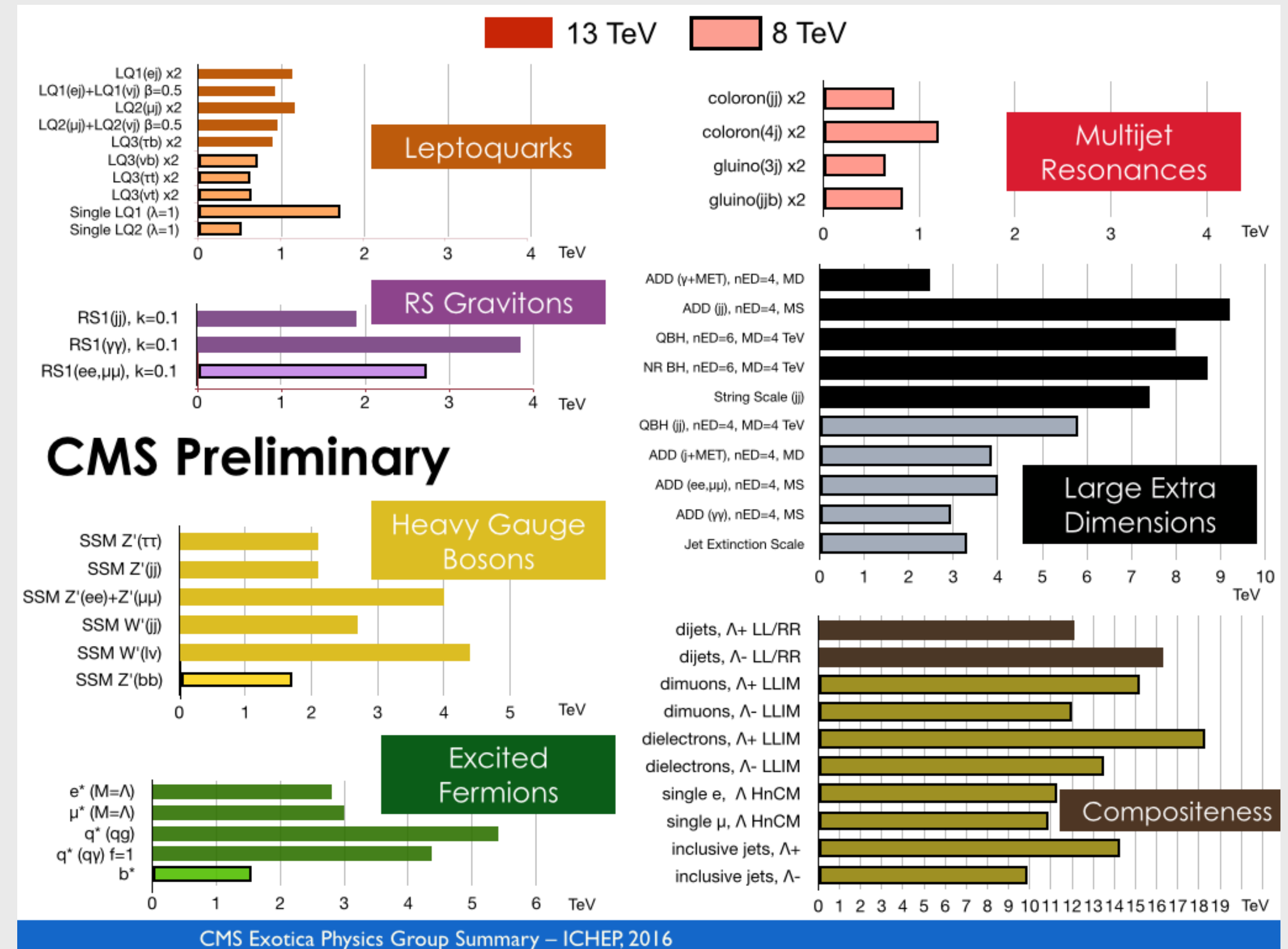
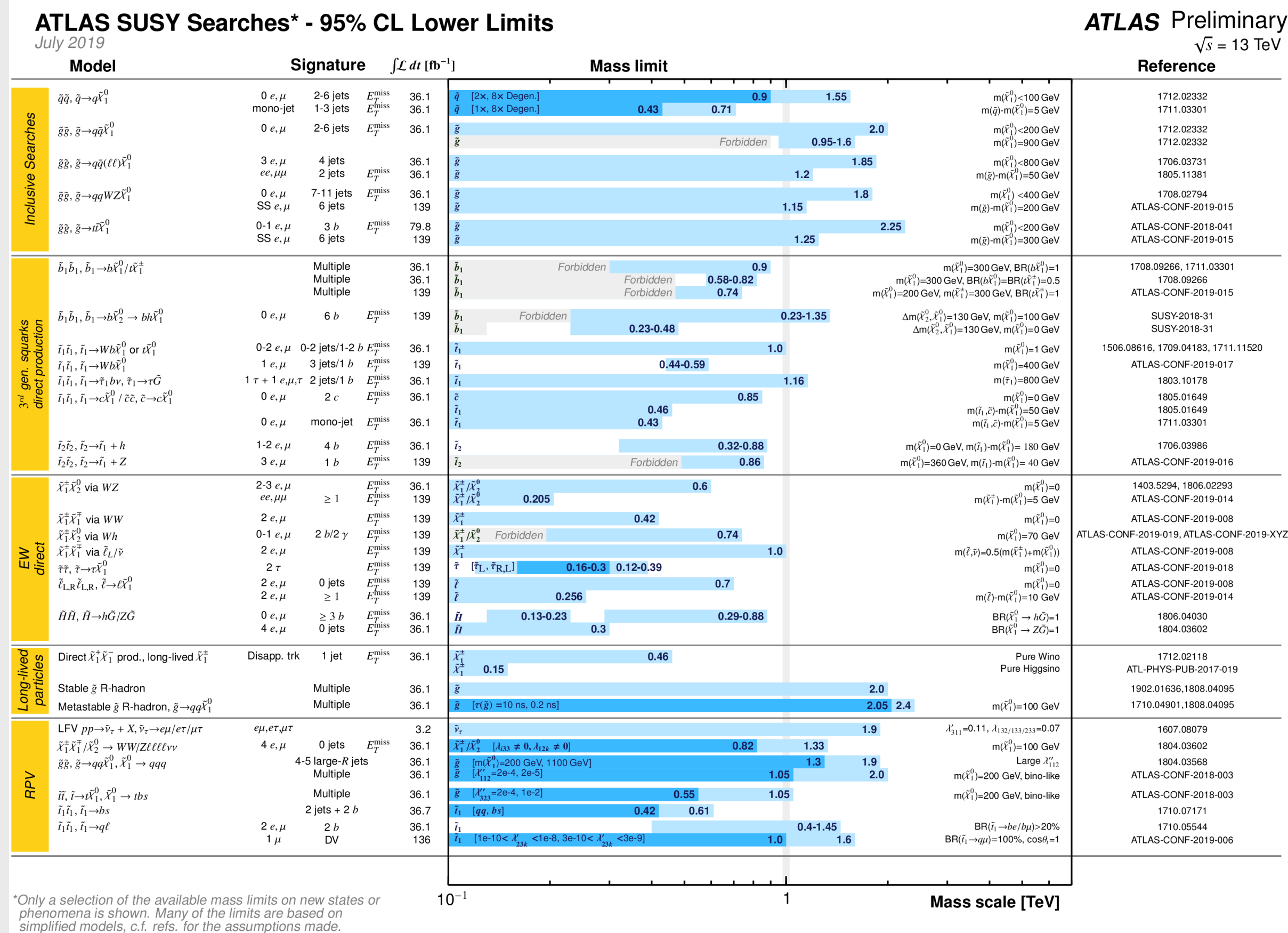


LHC ruled out new
physics at N TeV ...

LHC ruled out new physics at the TeV ...

SUMMARY

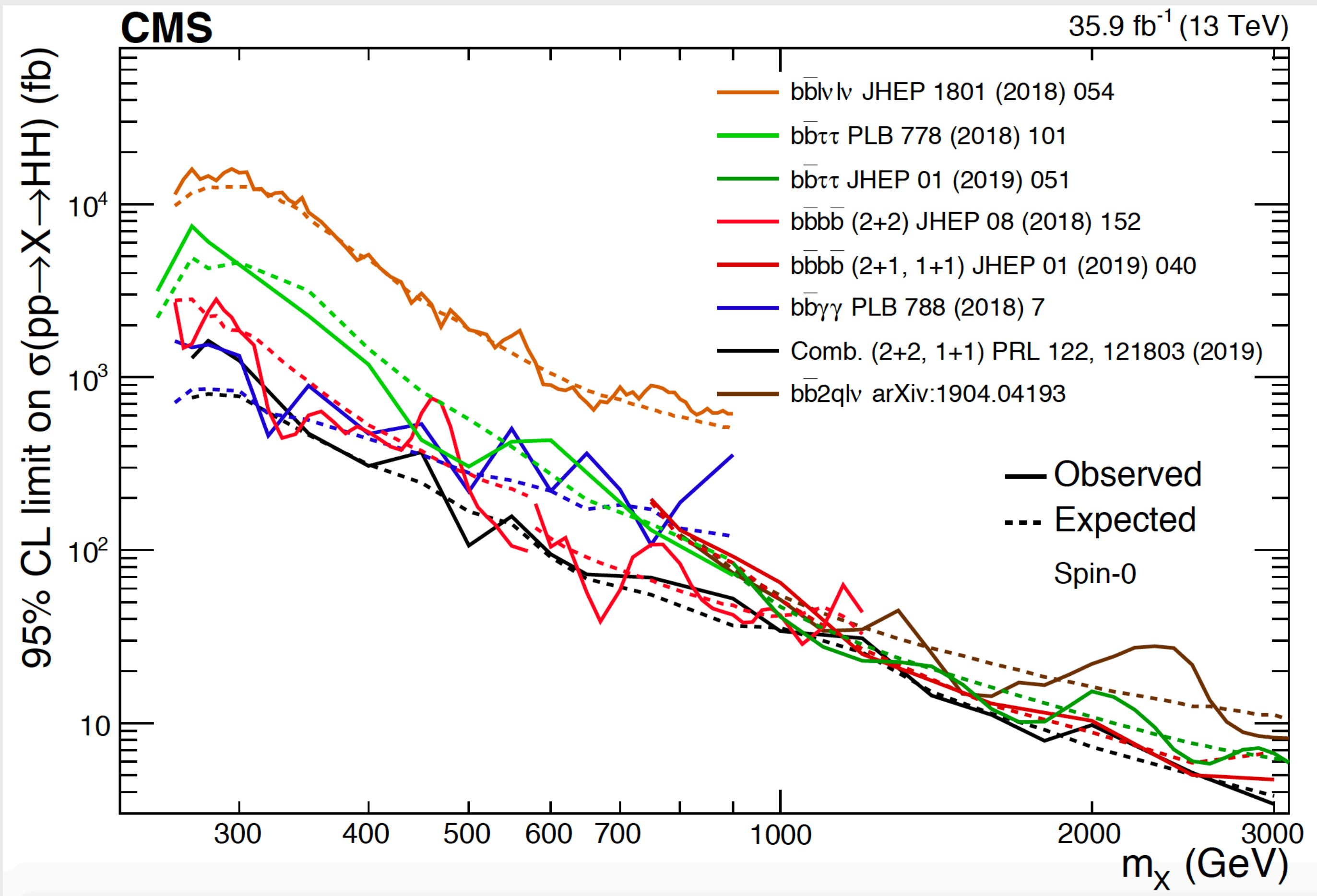
OF THE SUMMARIES



What about electroweak scalars?

SINGLET

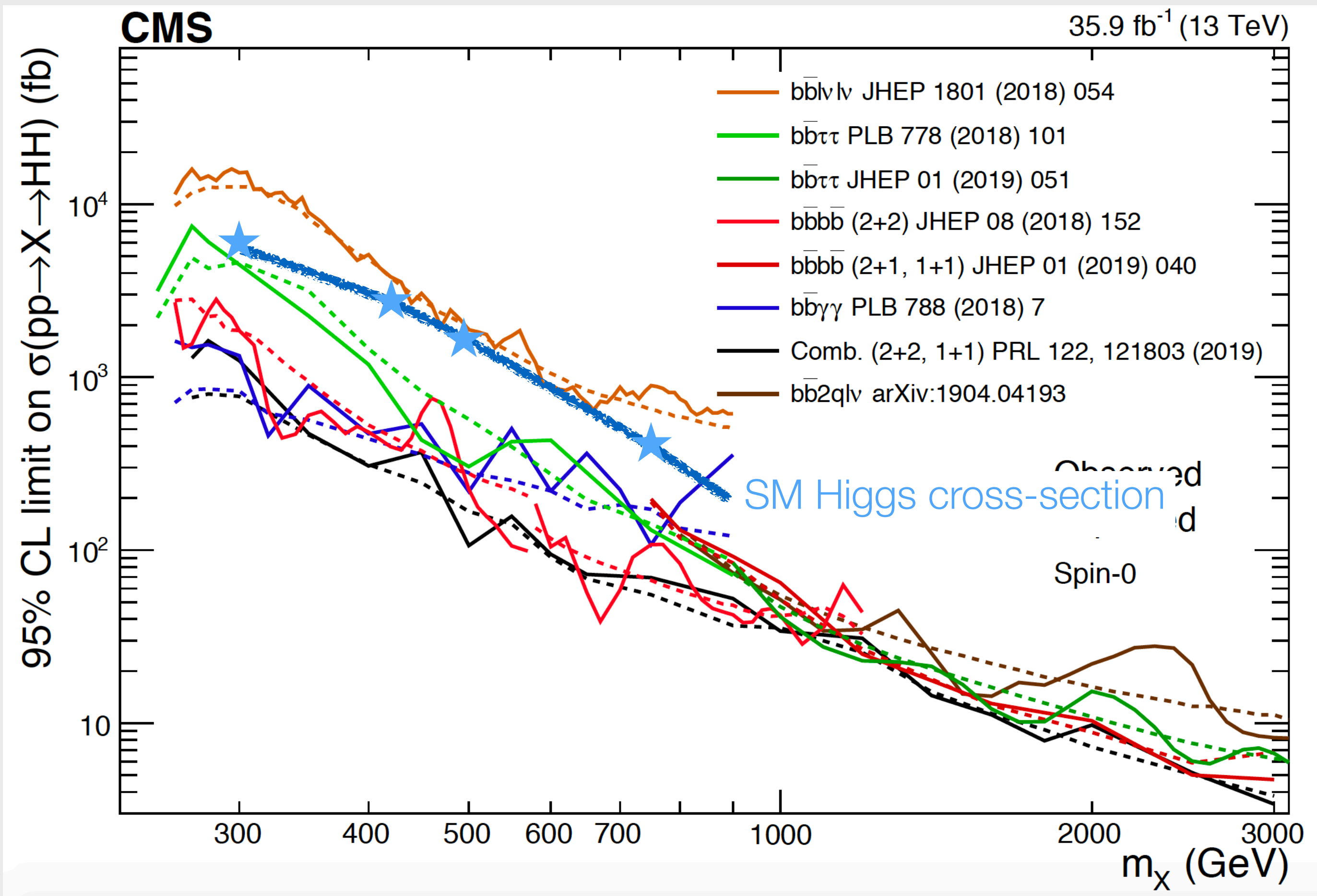
ARE ELUSIVE



What about electroweak scalars?

SINGLETs

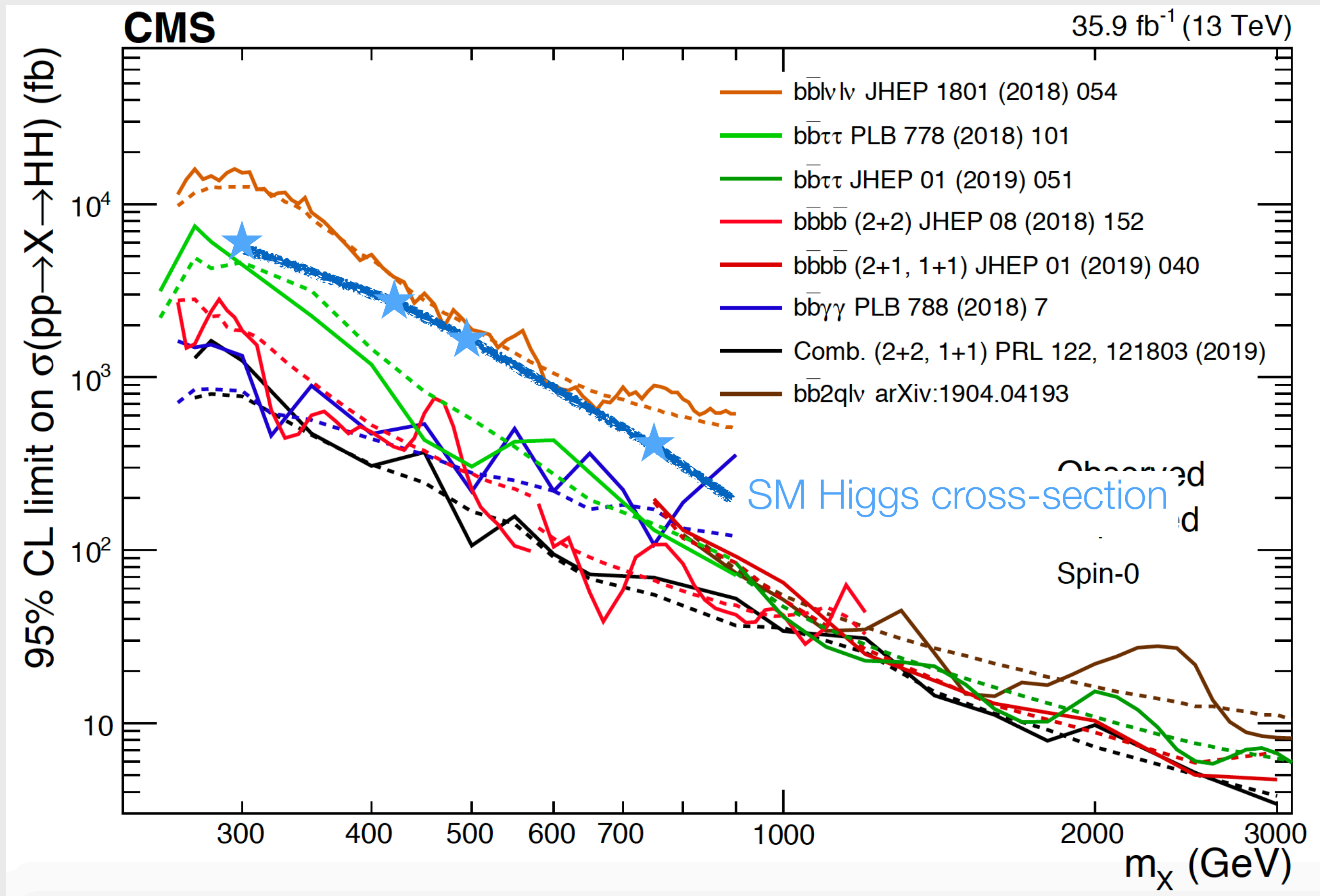
ARE ELUSIVE



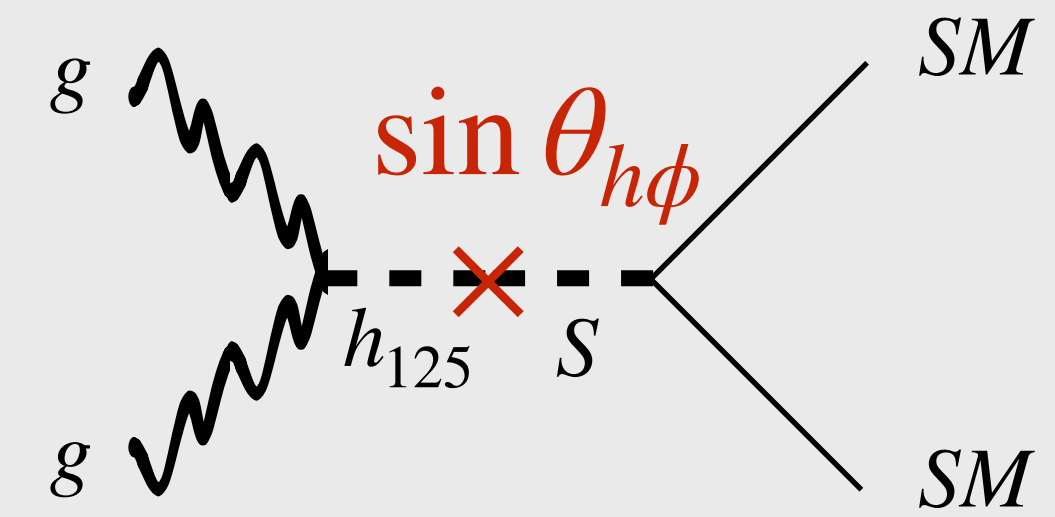
What about electroweak scalars?

SINGLETs

ARE ELUSIVE



$$\sigma(\phi) \sim \sin^2 \theta_{h\phi} \cdot \sigma(h_{SM} \text{ with } m_\phi)$$



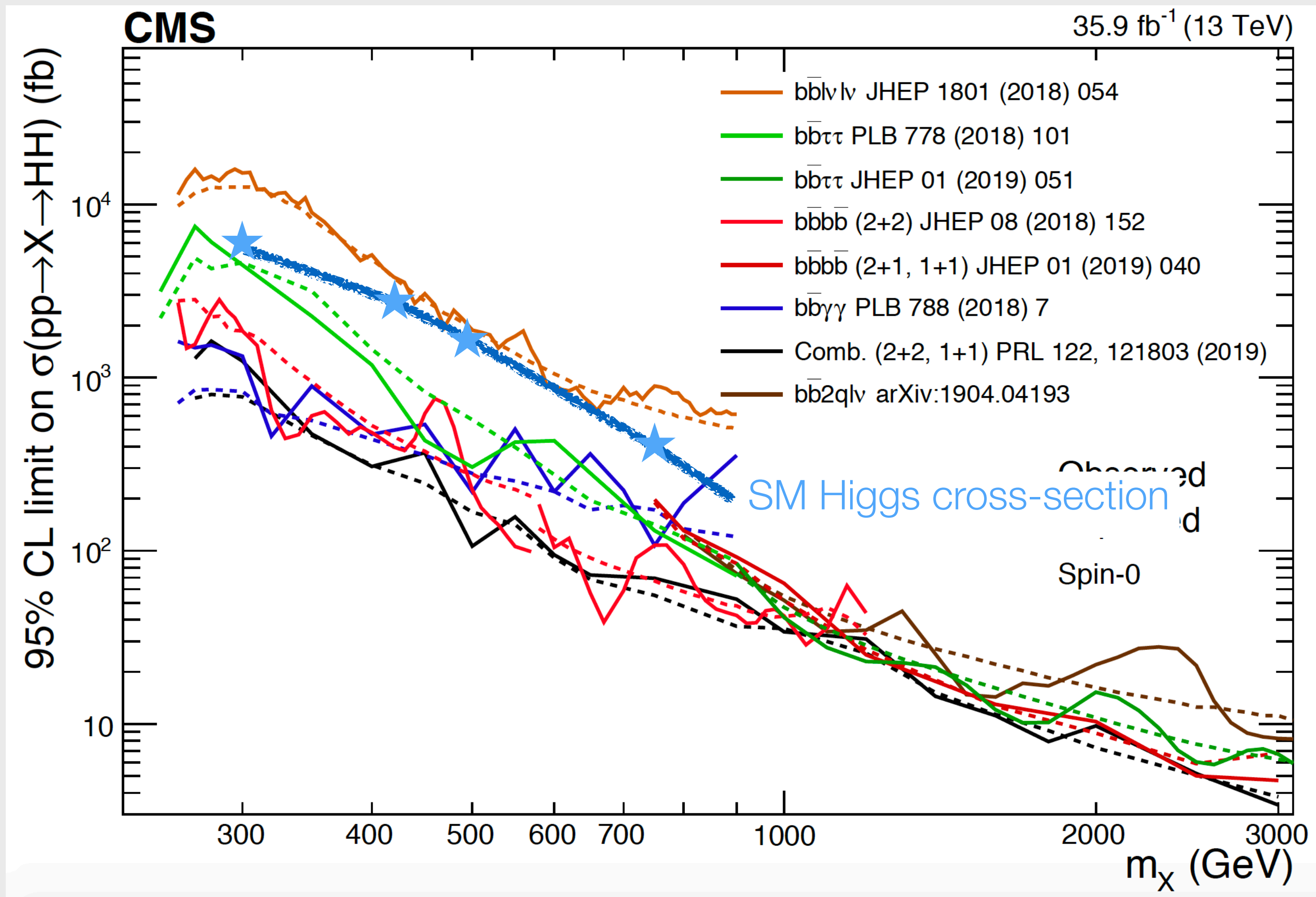
$$\Rightarrow \sin \theta_{h\phi} \lesssim 0.3$$

$$\sin \theta_{h\phi} \simeq \left(\frac{m_h}{m_\phi} \right)^\alpha \Rightarrow m_\phi \simeq 2 \div 3 \cdot m_h$$

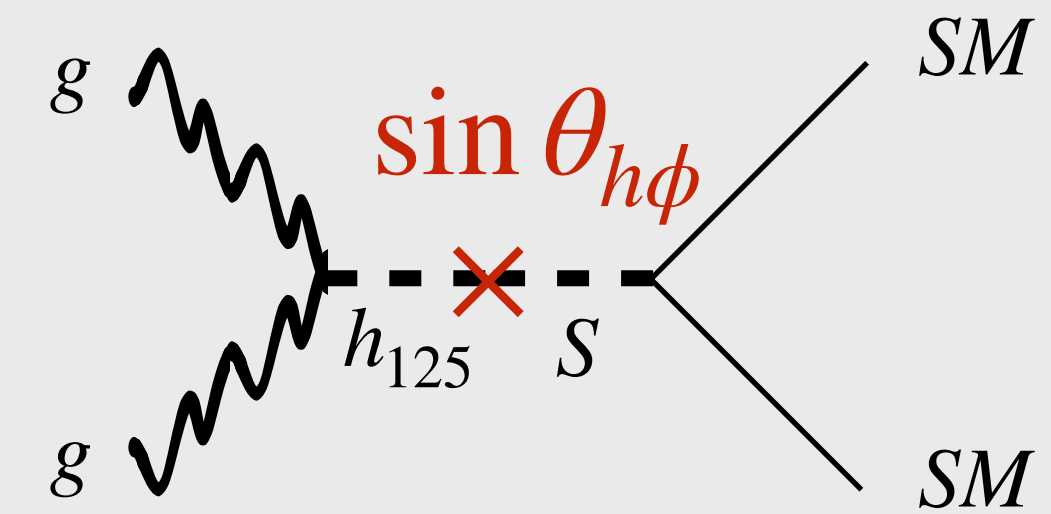
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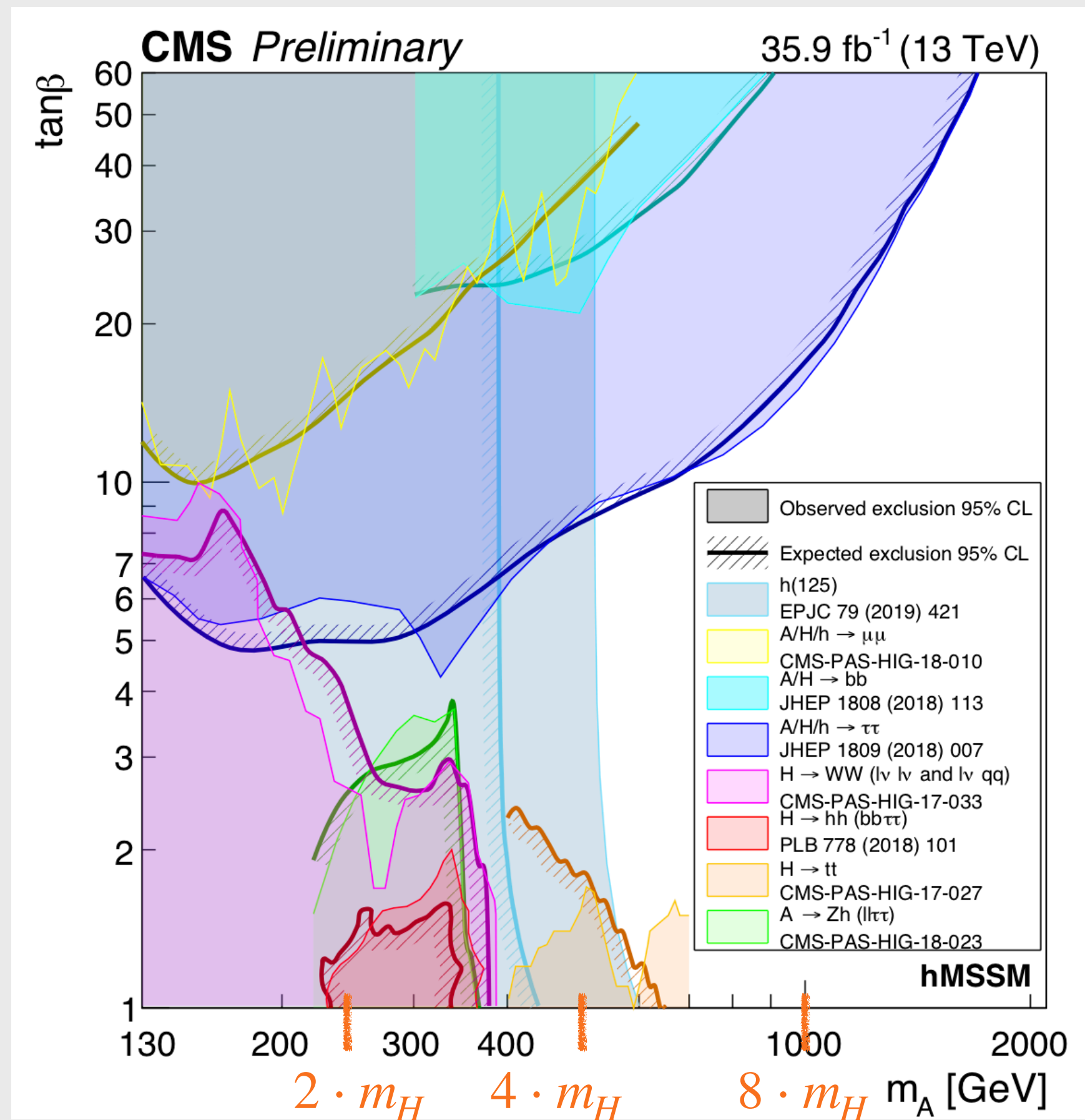
$$\sin \theta_{h\phi} \simeq \left(\frac{m_h}{m_\phi} \right)^\alpha \Rightarrow m_\phi \simeq 2 \div 3 \cdot m_h$$

$1 \leq \alpha \leq 2$

What about electroweak scalars?

DOUBLETS

ARE ABOUT AS TOUGH TO CATCH

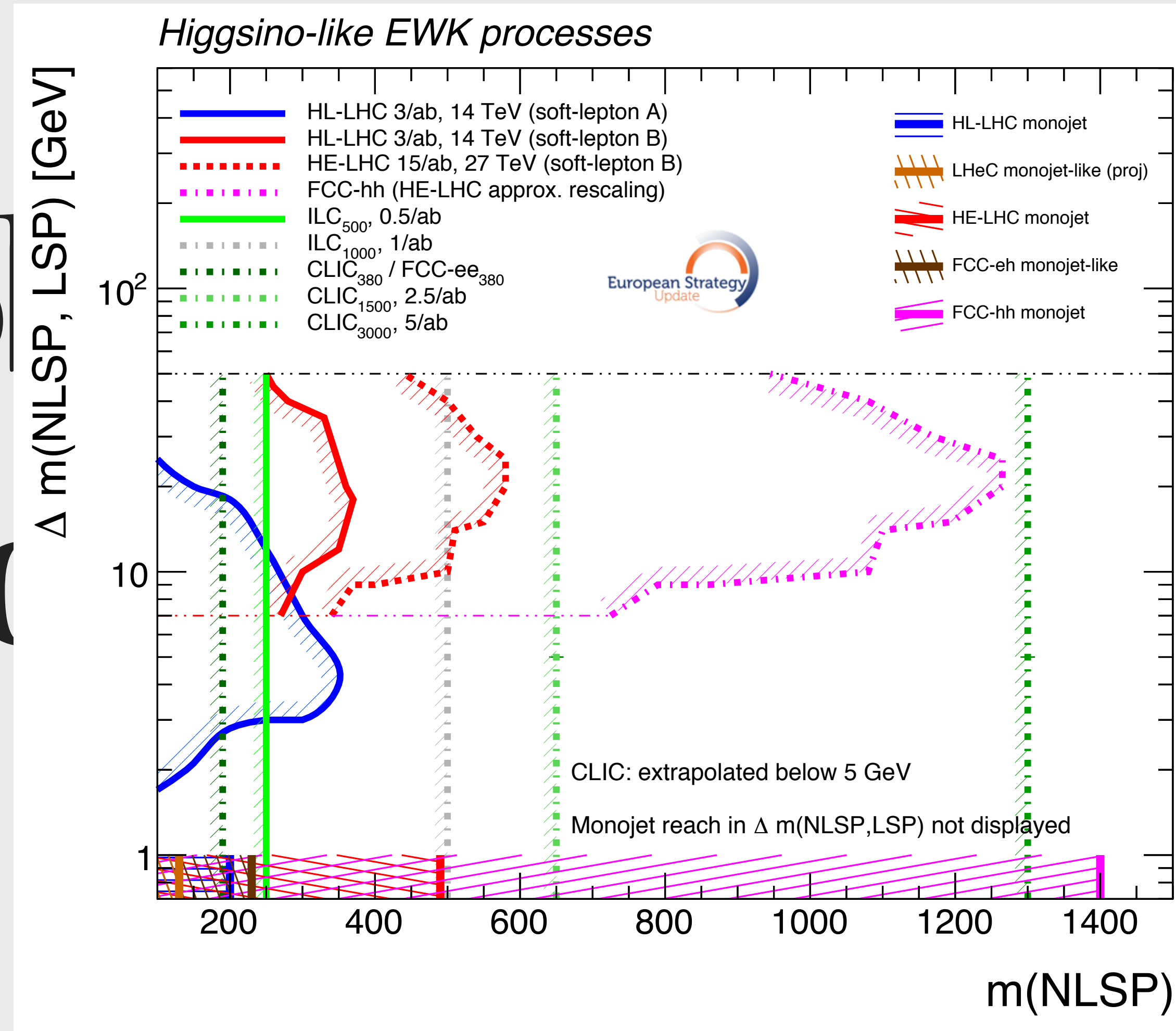


There is in general a weak sensitivity to new scalars, because of:

- “small” cross-sections
- large backgrounds

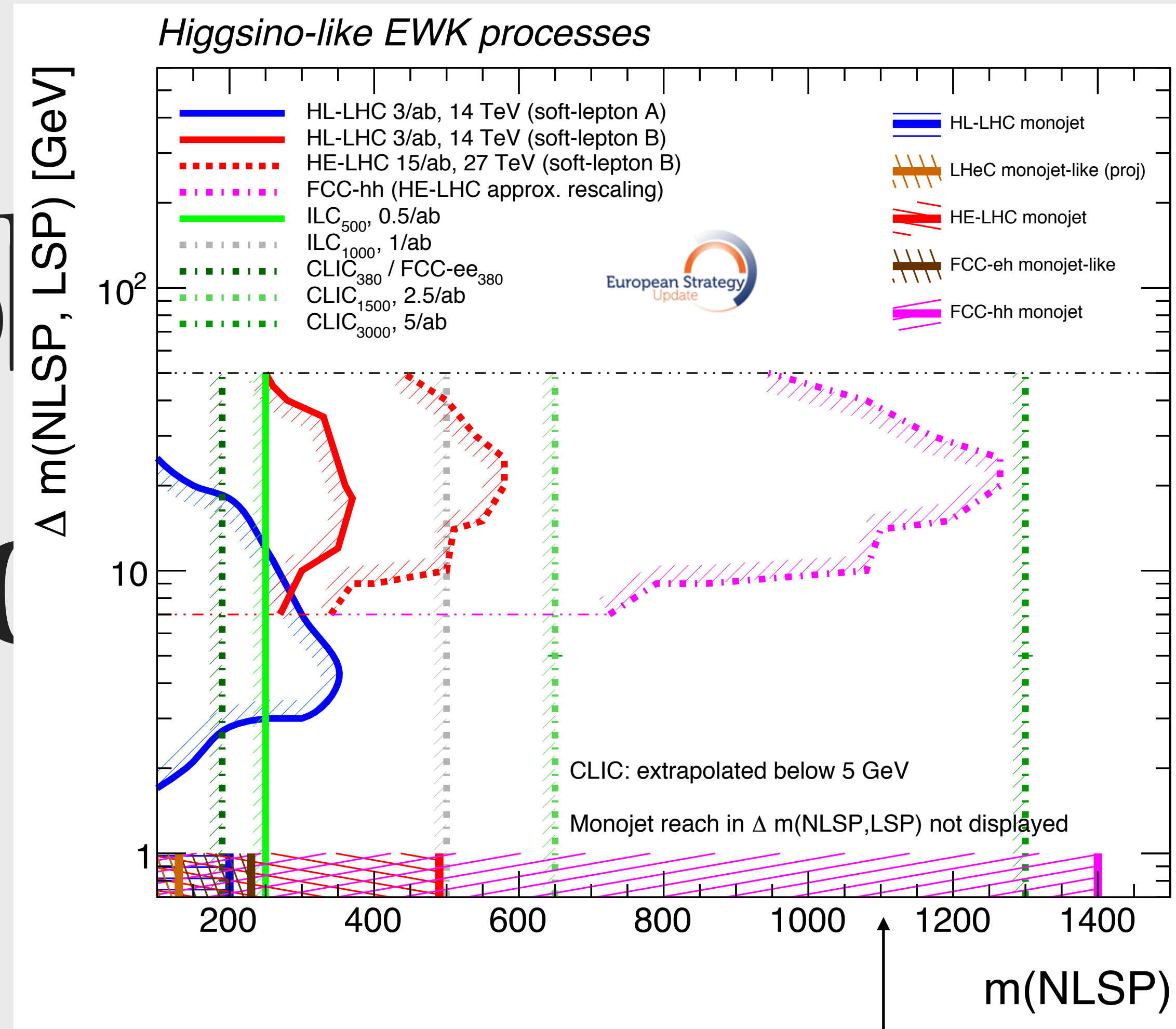
**this problem is common to lots
of electroweak new physics states**

this pro
of electro



on to lots
sics states

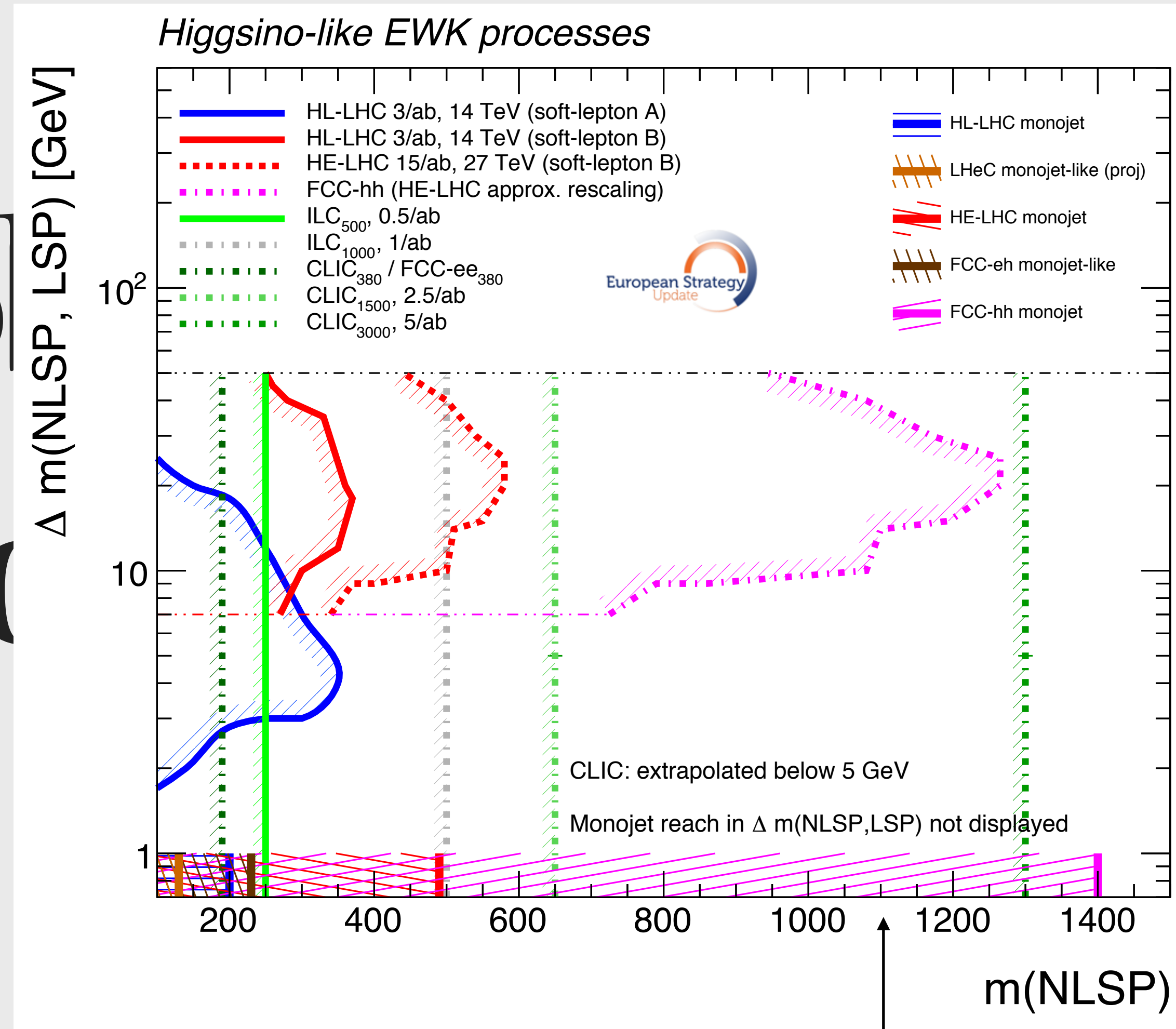
this pro
of electro



on to lots
sics states

This could be Dark Matter

this pro
of electro



on to lots
sics states

This could be Dark Matter
(inaccessible in Direct Detection!)

Yes, after HL-LHC there is going to be a uncharted territory as low as

- **Fermionic pure Doublet: 200 GeV; 400 GeV if you are really pessi/opti-misitic**
- **Scalar Doublet: 1 TeV**
- **Scalar Singlet: 500-900 GeV (depending on the UV origin of the singlet) ***

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- **Some searches have “loopholes”**