

# Unsupervised ML & Bayesian Inference @ four-tops

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**Measured cross-sections** 

ATLAS: 
$$25^{+7}_{-6}$$
 fb  
CMS:  $12.6^{+5.8}_{-5.2}$  fb

(Details in previous talk by Kong)



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(Details in previous talk by Kong)





Four tops

p t t









• Flavour rescaling

**Before** 

• Sequential kinematic reweighting



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In multilepton and/or multijets/b, some results have tuning

- $\rightarrow$  Quite clever techniques
- $\rightarrow$  In principle do not modify results
- $\rightarrow$  Some times after having seen the data

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- → Quite clever techniques
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But....

What if we are missing something ?



The algorithm recognizes similarities and differences and clusters the data



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?



?





We've started studying

- Latent Dirichlet Allocation (LDA)
- Autoencoders (AE)
- Variational Autoencoders (VAE)

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Dataset 2LSS++

Signal = tttt Background = ttW+ 10 fb<sup>-1</sup> / 600 fb<sup>-1</sup>  $\rightarrow \sim 1/1$  events

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Solve Bayesian inference numerically using Gibbs Sampling (also EMCEE and others)

#### **Bayesian Inference Results**



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Each parameter approaches the true values with the posterior!







Select a signal region sample Bayesian Inference to get Signal and Background features Tune Montecarlo using background features in signal region!

#### Conclusions

- Unsupervised ML & Bayes provide a new way of using data
- Four-tops is a very suitable physics case
- Presented algorithm: As simple as providing all (N<sub>i</sub>,N<sub>b</sub>) pairs!
- Tune MC with Background in Signal region
- New ways to test for SM & NP @ four-tops
- Subtleties and details in upcoming arXiv:2106.XXXX. Also more to explore & understand.

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More discussion after this session:

zoom.us/j/95915707476?pwd=NEdsekM5d05KVS9VZWdKOHEwY3l5QT09

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#### 7.2 Sequential kinematic reweighting

Following the flavour rescaling, a sequential reweighting is used to mitigate the kinematic mismodelling observed in  $t\bar{t}$ +jets MC. The reweighting corrects for the distributions of  $N_{jets}$ , the number of large-R jets  $(N_{LR-jets})$ , the scalar sum of all jet and lepton  $p_T$  in the event  $(H_T^{all})$ , and the average  $\Delta R$  between any two jets  $(\Delta R_{avg.}^{ij})$ . These variables are related to the overall jet activities in the events and are observed to be mismodelled, especially the  $N_{jets}$  and  $H_T^{all}$  spectra. These variables capture the most representative global kinematics of the events, as well as kinematic properties of the individual jets such as  $p_T$  and their angular distributions.

The  $t\bar{t}$ +jets events in  $\geq 3b$  regions are reweighted according to the discrepancy between data and MC in the 2b regions. The reweighting factors are derived such that the overall MC prediction matches the data in the 2b regions. This is done based on the assumption that the deficiency of the radiation modelling in the parton shower is independent of the flavour of the radiated jets. Systematic variations on the  $t\bar{t}$ +jets modelling cover possible deviations from such assumption.

#### Four-tops ATLAS-CONF-2021-013

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#### 7.1 $t\bar{t}$ +jets flavour rescaling

The  $t\bar{t}$ +jets flavour rescaling corrects for the overall yields of  $t\bar{t}$ +light,  $t\bar{t}$ + $\geq 1c$  and  $t\bar{t}$ + $\geq 1b$  categories. The rescaling factors are derived from a dedicated profile likelihood fit to data using the event yields in the regions defined by various *b*-tagging requirements. Events with  $\geq 8j$  in the 1L channel and  $\geq 6j$  in the 2LOS channel are split into 2b, 3bL, 3bH and  $\geq 4b$  regions, using the same criteria as defined in Table 1. The fit exploits the different  $t\bar{t}$ +jets flavour fractions in the eight fitted regions. The largest signal to background ratio in these regions is 2.5%, estimated using MC prior to the fit. Systematic uncertainties due to the tagging efficiency of *b*-jets and the mis-tag rate of *c*- and light-jets are considered as nuisance parameters in the profile likelihood fit. The measured rescaling factors for  $t\bar{t}$ +light,  $t\bar{t}$ + $\geq 1c$  and  $t\bar{t}$ + $\geq 1b$  are  $1.0 \pm 0.1$ ,  $1.6 \pm 0.2$  and  $1.3 \pm 0.1$ , respectively, where the quoted uncertainties are from the statistical uncertainty on data and from uncertainties on the *b*-tagging calibration.



#### ATLAS CONF Note

ATLAS-CONF-2019-045

16th October 2019 Minor revision: 24th August 2020



#### Analysis of $t\bar{t}H$ and $t\bar{t}W$ production in multilepton final states with the ATLAS detector

A search for the associated production of a top-quark pair with the Higgs boson ( $t\bar{t}H$ ) in multilepton final states is presented. The search is based on a dataset of proton–proton collisions at  $\sqrt{s} = 13$  TeV recorded with the ATLAS detector at the CERN Large Hadron Collider and corresponding to an integrated luminosity of 80 fb<sup>-1</sup>. Six final states, defined by the number and flavour of charged-lepton candidates, and 25 event categories are defined to simultaneously search for the  $t\bar{t}H$  signal and constrain several leading backgrounds. The  $t\bar{t}W$  background normalisation is left unconstrained in the statistical analysis and the resulting  $t\bar{t}W$  normalisation is found to be higher than the theoretical prediction. An excess of events consistent with  $t\bar{t}H$  production, over the expected background from Standard Model processes, is found with an observed significance of 1.8 standard deviations, compared to an expectation of 3.1 standard deviations. Assuming Standard Model branching fractions, the best-fit value of the  $t\bar{t}H$  production cross section is  $\sigma_{t\bar{t}H} = 294^{+182}_{-162}$  fb, which is consistent with the Standard Model prediction. The impact on the  $t\bar{t}H$  cross section measurement of the assumptions made on the  $t\bar{t}W$  background modelling is discussed.

#### Four-tops ATLAS-CONF-2019-045

Eur. Phys. J. C (2016) 76:11 DOI 10.1140/epjc/s10052-015-3852-4 THE EUROPEAN PHYSICAL JOURNAL C

Regular Article - Experimental Physics

Measurements of fiducial cross-sections for  $t\bar{t}$  production with one or two additional *b*-jets in *pp* collisions at  $\sqrt{s} = 8$  TeV using the ATLAS detector

**ATLAS Collaboration\*** 

#### (see discussion in 1701.04427)

malised to the NNLO+NNLL result [32–37]. PYTHIA 8 offers several options for modelling  $g \rightarrow b\bar{b}$  splittings in the final-state parton showers, which may be accessed by varying the TIMESHOWER:WEIGHTGLUONTOQUARK (wgtq) parameter [75]. Differences between the models arise by neglecting (wgtq5) or retaining (wgtq3, wgtq6) the massdependent terms in the  $g \rightarrow b\bar{b}$  splitting kernels. Differences also arise with respect to the treatment of the high $m_{b\bar{b}}$  region, with specific models giving an enhanced or suppressed  $g \rightarrow b\bar{b}$  rate. The model corresponding to wgtq3 was chosen to maximise this rate. Finally, some of the models (wgtq5, wgtq6) offer the possibility to choose sgtq $\cdot m_{b\bar{b}}$  instead of the transverse momentum as the argument of  $\alpha_{\rm S}$  in the  $g \rightarrow b\bar{b}$  vertices. Here sgtq refers to the TIMESHOWER:SCALEGLUONTOQUARK parameter, and is allowed to vary in the range  $0.25 \leq \text{sgtq} \leq 1$ , with larger values giving a smaller  $g \rightarrow b\bar{b}$  rate and vice versa. For the model wgtq5, sgtq was set to 1, a combination that minimises the  $g \rightarrow b\bar{b}$  rate, while for wgtq6, sgtq was set to 0.25.

Eur. Phys. J. C (2016) 76:379 DOI 10.1140/epjc/s10052-016-4105-x THE EUROPEAN PHYSICAL JOURNAL C

Regular Article - Experimental Physics

Measurement of  $t\bar{t}$  production with additional jet activity, including b quark jets, in the dilepton decay channel using pp collisions at  $\sqrt{s} = 8$  TeV

CMS Collaboration\*

11. CMS Collaboration, Measurement of the cross section ratio  $\sigma_{t\bar{t}b\bar{b}}/\sigma_{t\bar{t}jj}$  in pp collisions at  $\sqrt{s} = 8$  TeV. Phys. Lett. B **746**, 132 (2015). doi:10.1016/j.physletb.2015.04.060. arXiv:1411.5621

#### (see discussion in 1701.04427)

PYTHIA6 and HERWIG6. The normalization factors applied to the MADGRAPH and POWHEG predictions are found to be about 1.3 for results related to the leading additional b jet. The predictions from both generators underestimate the ttbb cross sections by a factor 1.8, in agreement with the results from Ref. [11]. The normalization factors applied to MC@NLO are approximately 2 and 4 for the leading and subleading additional b jet quantities, respectively, reflecting the observation that the generator does not simulate sufficiently large jet multiplicities. All the predictions have slightly harder  $p_{\rm T}$ spectra for the leading additional b jet than the data, while they describe the behaviour of the  $|\eta|$  and  $m_{bb}$  distributions within the current precision. The predictions favour smaller  $\Delta R_{\rm bb}$  values than the measurement, although the differences are in general within two standard deviations of the total uncertainty.