

Flavours - Classical spectroscopy

V.Mariani for the
ALICE, ATLAS, CMS, LHCb Collaborations

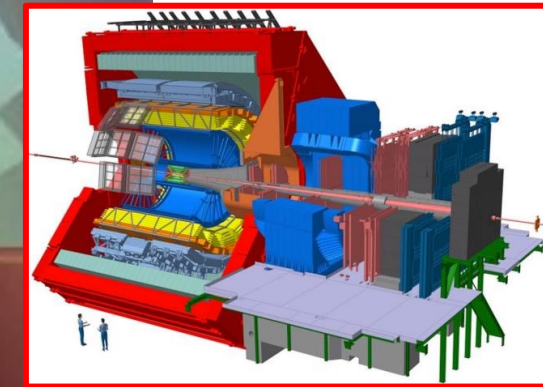
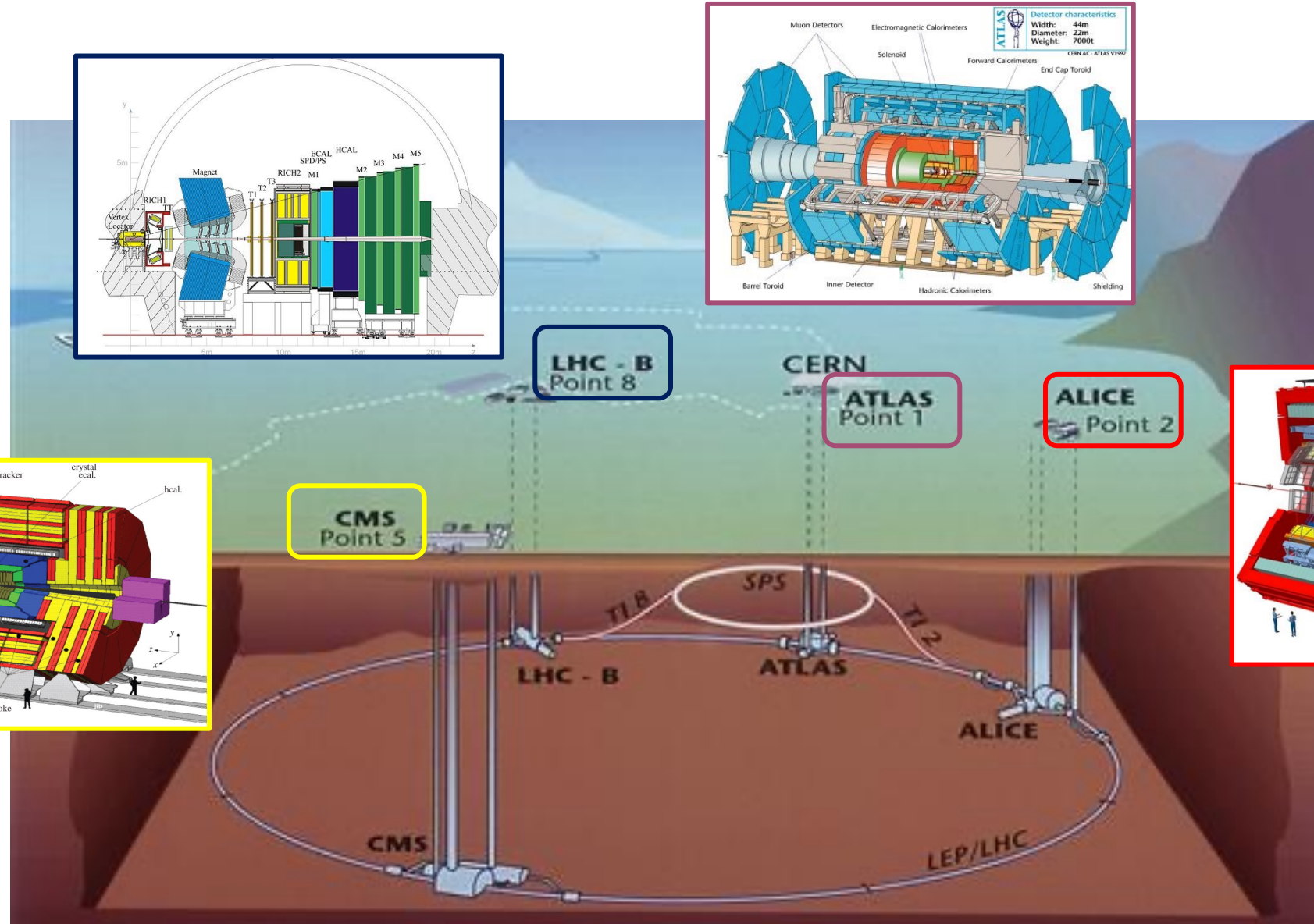
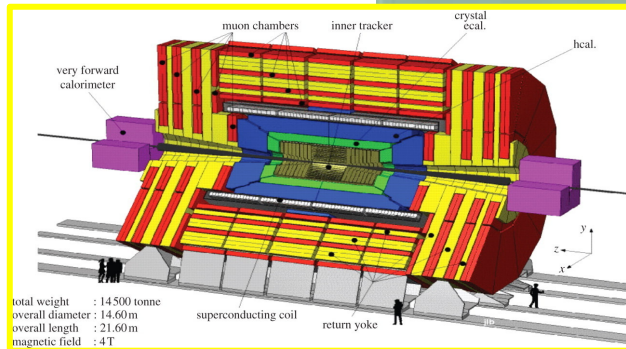
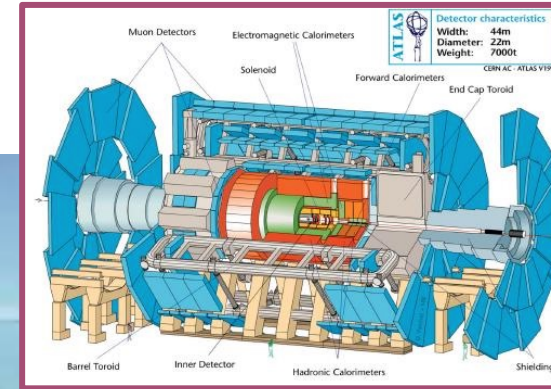
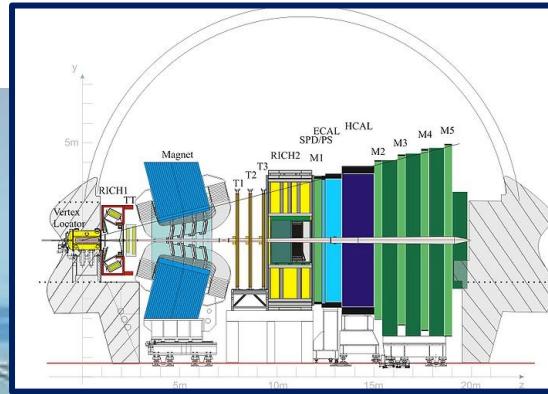


Heavy Flavour physics

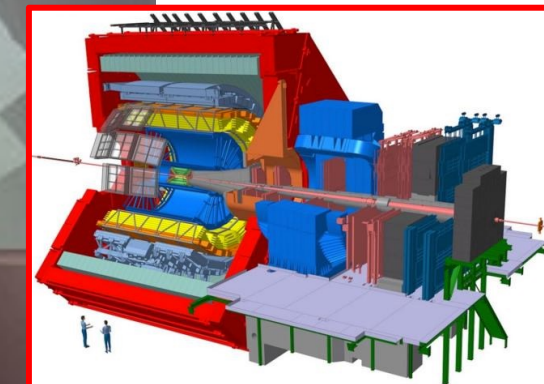
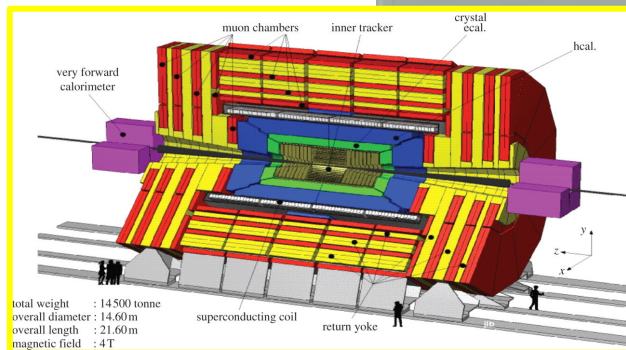
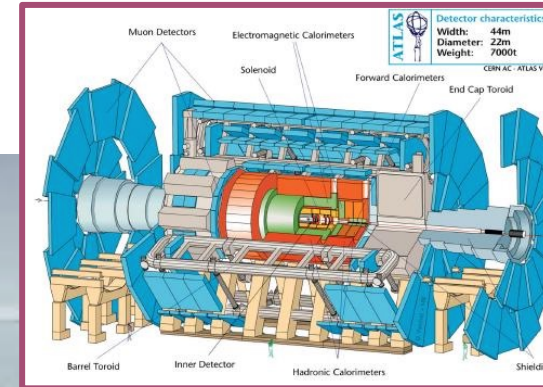
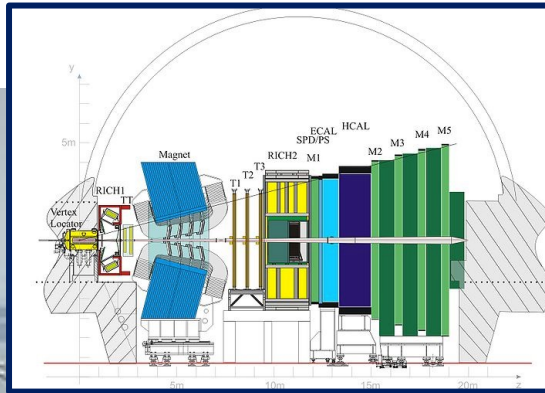
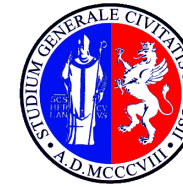


- Measurements and observations of heavy flavour production provide **important tests of QCD** and give insight into particle production at colliders
- **Hadronization challenging** to understand -> measurements needed
- Form **baseline or background** for other physics studies at the LHC
- LHC provides access to **wide kinematic range** with a **very high production cross section** if compared to e^+e^- colliders.

Experiments at the LHC



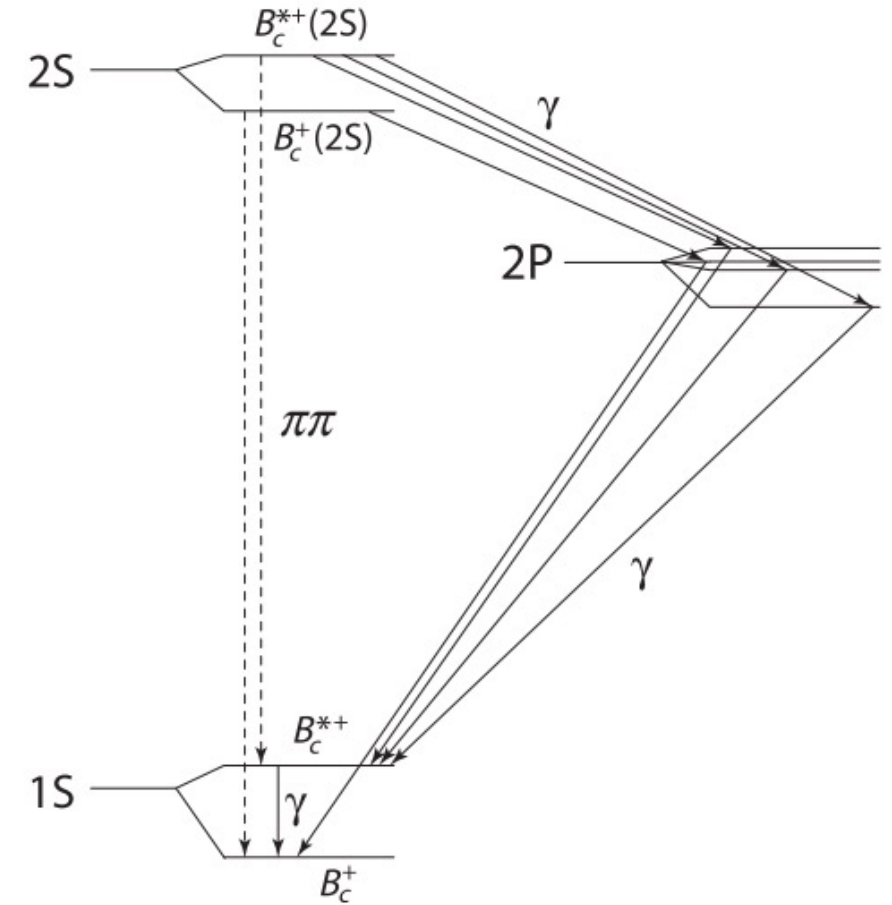
Experiments at the LHC



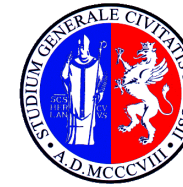
I'll show the main / latest results from the four Collaborations on the Heavy Flavour spectroscopy

$B_c^{*}(2S)$

- The B_c mesons ($b\bar{c}$) family is predicted to be **very populated** but the spectroscopy and property studies are still scarce.
- Less explored compared to quarkonia because of the **small production rate**: dominant mechanism requires the production of both $c\bar{c}$ and $b\bar{b}$.
- The $b\bar{c}$ excited states decay to the ground states via the cascade emission of γ and π pairs -> total width of $O(100\text{KeV})$
-> **hard to detect**

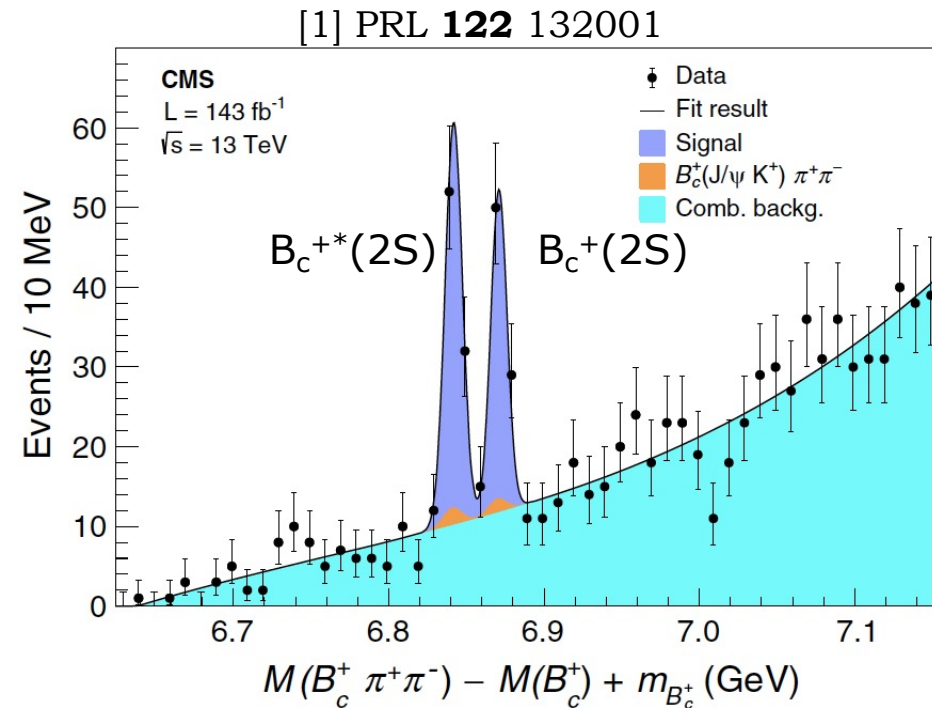
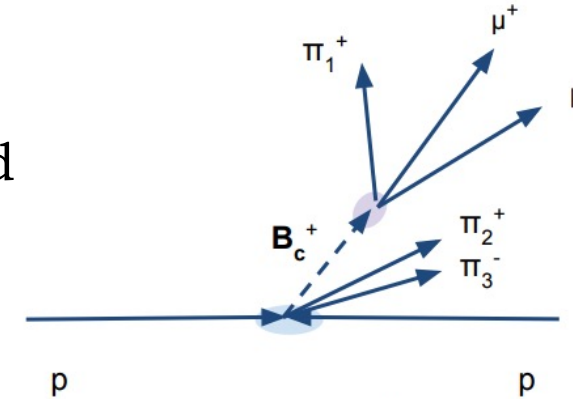


$B_c^{*}(2S)$

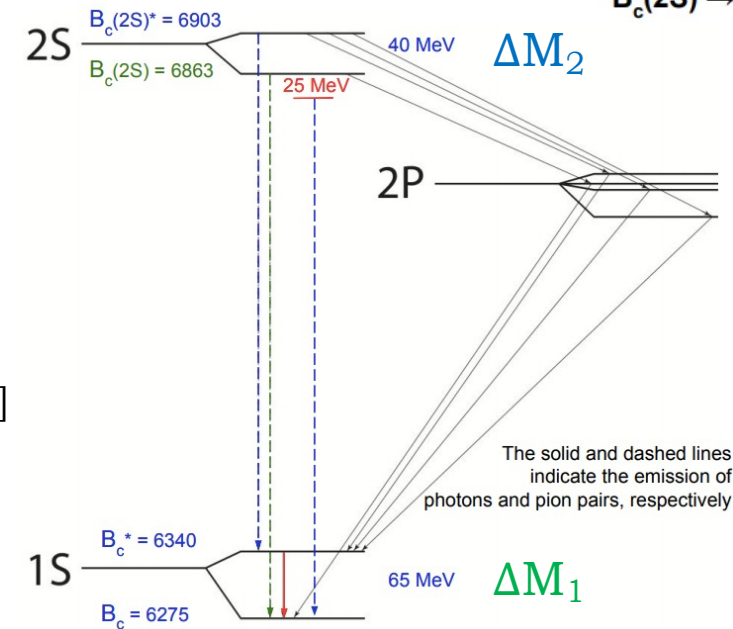


The existence of two separate states, $B_c^+(2S)$ and $B_c^{*+}(2S)$, was announced by CMS in [1], using data from Run2.

Significance of 2 peaks instead of 1 is 6.5σ
 \Rightarrow **First observation of two excited states**



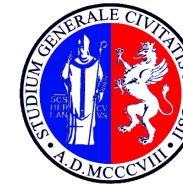
Consistent with ATLAS observation of one unresolved $B^*(2S)$ state [PRL 113 (2014) 212004] -> backup



$$M(B_c^+(2s)) = 6871.0 \pm 1.2 \text{ (stat)} \pm 0.8 \text{ (syst)} \pm 0.8 \text{ (} B_c^+ \text{) MeV}$$

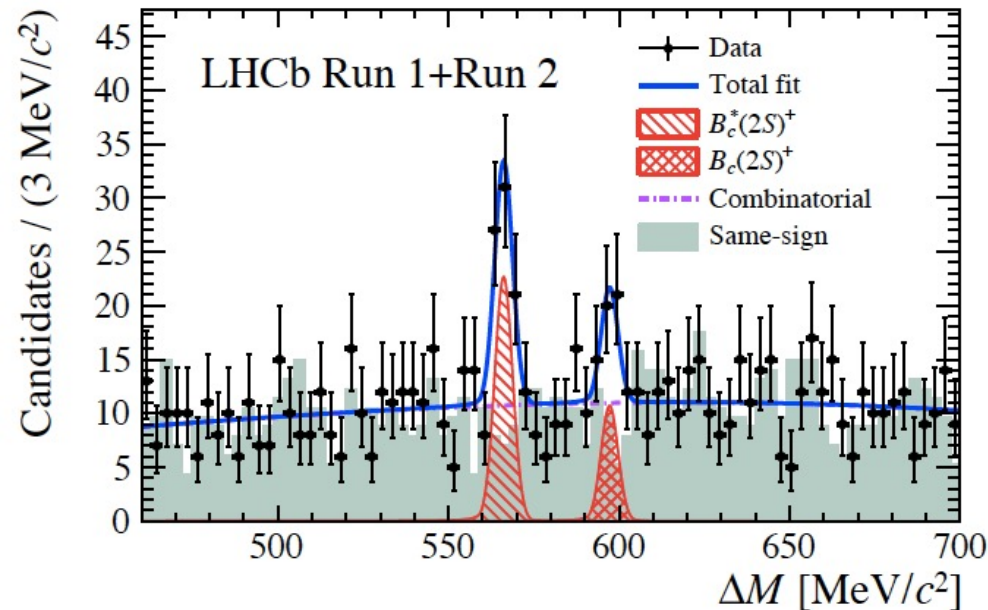
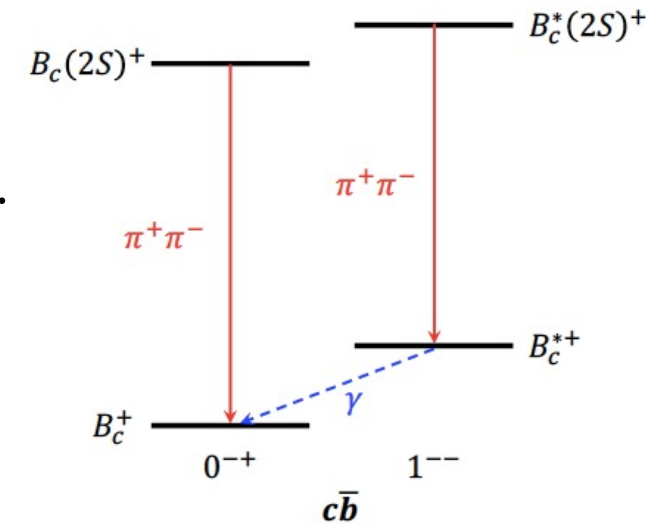
$$\Delta M = \Delta M_1 - \Delta M_2 = 29.1 \pm 1.5 \text{ (stat)} \pm 0.7 \text{ (syst) MeV}$$

$B_c^{*}(2S)$



LHCb could confirm the result using the statistics of Run1+Run2 (8.5 fb^{-1}) [2]

The second state is seen with a global (local) significance of 2.2σ (3.2σ).



$$M(B_c^+(2s)) = 6872.1 \pm 1.3 \text{ (stat)} \pm 0.1 \text{ (syst)} \pm 0.8 \text{ (} B_c^+ \text{) MeV}$$

$$\Delta M = \Delta M_1 - \Delta M_2 = 31.1 \pm 1.4 \text{ (stat) MeV}$$

Consistent with CMS and with the theoretical predictions

Barions with heavy-flavor quark



- Spectroscopy of baryons that contain a heavy-flavor quark can test predictions of heavy-quark effective theory.
 - Several theoretical calculations exist for various excitations of the ground state baryons containing a b quark.
- The widths and the production cross sections of various excited states are generally unknown.

This situation makes **experimental searches** for excited heavy-quark baryons both **challenging** and **important** for testing various theoretical models.

Λ_b^0 excited states



LHCb already measured two narrow states [3] denoted as $\Lambda_b(5912)^0$ and $\Lambda_b(5920)^0$

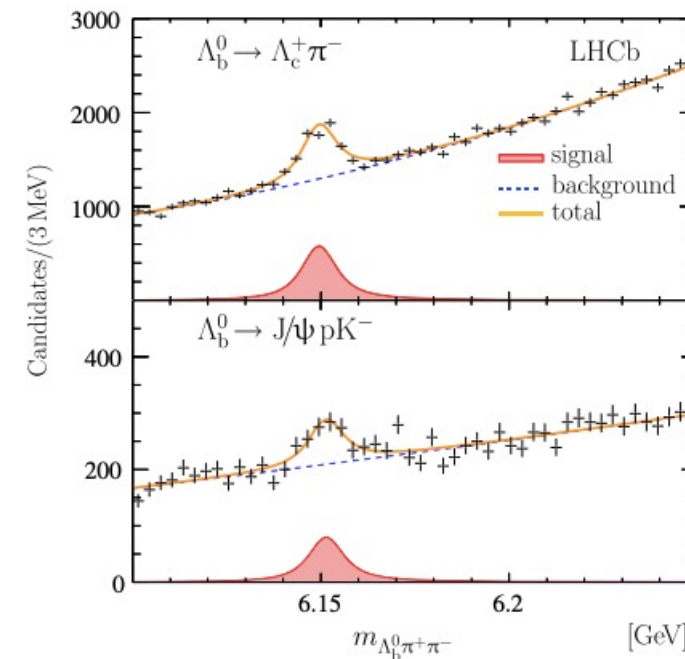
- Using data collected at 7, 8 and 13 TeV ($L=9/\text{fb}$) a search in the $\Lambda_b^0 \pi^+ \pi^-$ spectrum in the range 6.10- 6.25 GeV has been performed [4].

Two Λ_b^0 decays: $\Lambda_b^0 \rightarrow \Lambda_c^+ \pi^-$ and $\Lambda_b^0 \rightarrow J/\psi p K^-$

A clear excess of $\Lambda_b^0 \pi^+ \pi^-$ candidates around 6.15 GeV for both Λ_b^0 decay modes

Mass above the $\Sigma_b^{*\pm} \pi^\pm$ threshold => three regions investigated:

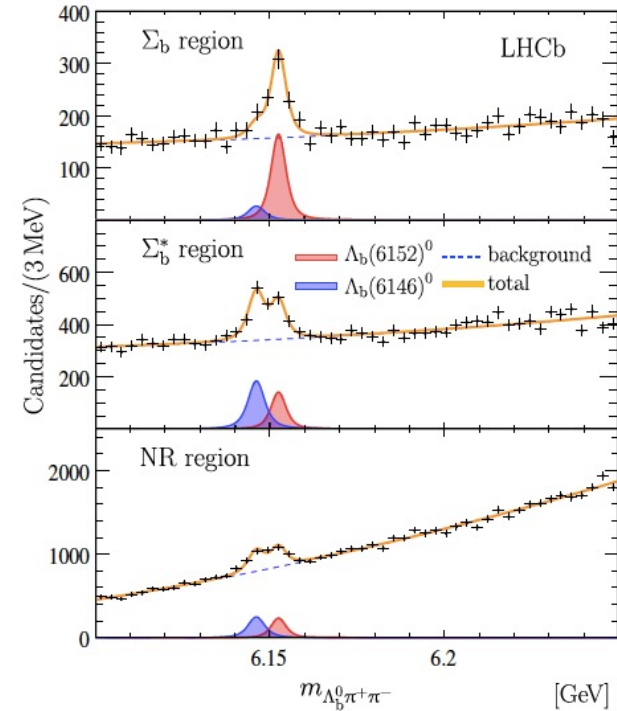
- $\Lambda_b^0 \pi^\pm$ mass within the width of Σ_b^\pm
- $\Lambda_b^0 \pi^\pm$ mass within the width of $\Sigma_b^{*\pm}$
- Nonresonant region



[3] PRL **109**, 172003

[4] PRL **123**, 152001

Λ_b^0 excited states



The shapes suggest two different narrow peaks \Rightarrow states are called $\Lambda_b(6146)^0$ and $\Lambda_b(6152)^0$ and masses and widths are measured:

$$m_{\Lambda_b(6146)^0} = 6146.17 \pm 0.33 \pm 0.22 \pm 0.16 \text{ MeV},$$

$$m_{\Lambda_b(6152)^0} = 6152.51 \pm 0.26 \pm 0.22 \pm 0.16 \text{ MeV},$$

$$\Gamma_{\Lambda_b(6146)^0} = 2.9 \pm 1.3 \pm 0.3 \text{ MeV},$$

$$\Gamma_{\Lambda_b(6152)^0} = 2.1 \pm 0.8 \pm 0.3 \text{ MeV},$$

The mass difference between the two states is $6.34 \pm 0.32 \pm 0.02 \text{ MeV}$.

States measured are **consistent with the predictions for the doublet of $\Lambda_b(1D)^0$ states** with quantum numbers $J^P = 3/2^+$ and $5/2^+$

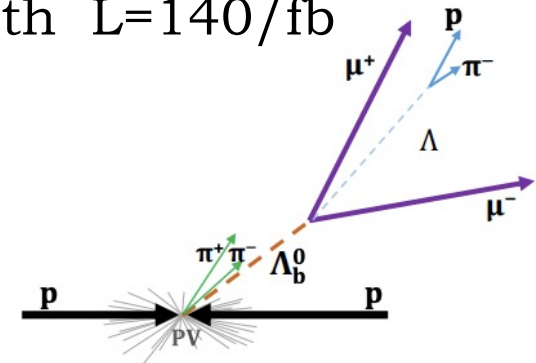
However, the **interpretation of these states as excited Σ_b^0 states cannot be excluded.**

Λ_b^0 excited states



CMS recently measured [5] excited states of the Λ_b^0 using Run2 data 16-18 with $L=140/\text{fb}$

- Two Λ_b^0 decays: $\Lambda_b^0 \rightarrow \psi(2S)\Lambda \rightarrow \psi(2S)p\pi^-$ and $\Lambda_b^0 \rightarrow J/\psi\Lambda \rightarrow J/\psi p\pi^-$
- $\Lambda_b(5912)^0$ and $\Lambda_b(5920)^0$ states have been confirmed (backup)

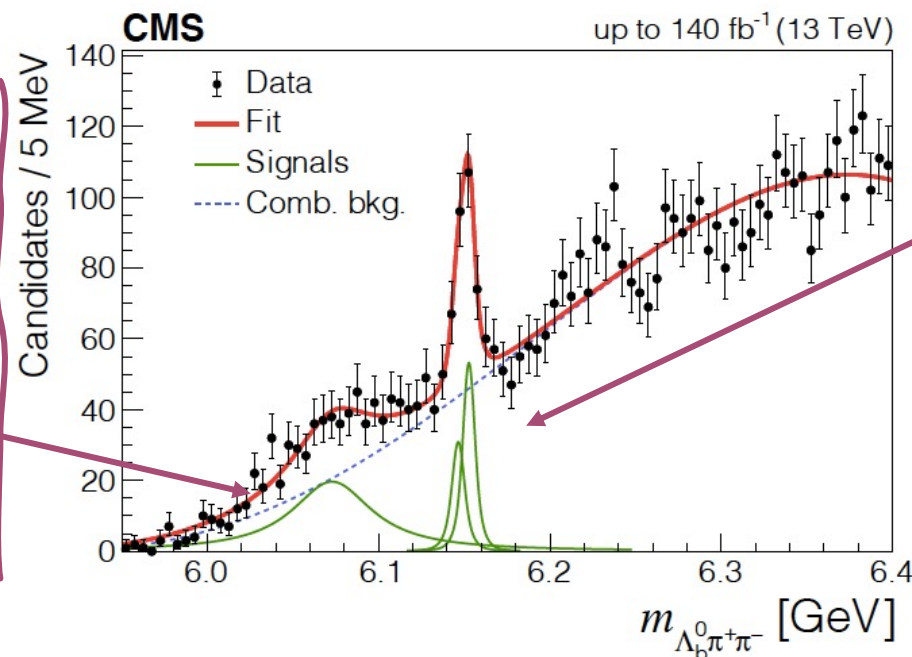


In the higher mass the $\Lambda_b^0\pi^+\pi^-$ spectrum show some structures

Broad excess (4σ significance) maybe related to intermediate states of $\Sigma_b^{*\pm}$ and Σ_b^\pm .
 $m=6073\pm 5$ (stat) MeV
 width= 55 ± 11 (stat) MeV.

More data to better elucidate

=> LHCb also confirmed a similar structure (backup)

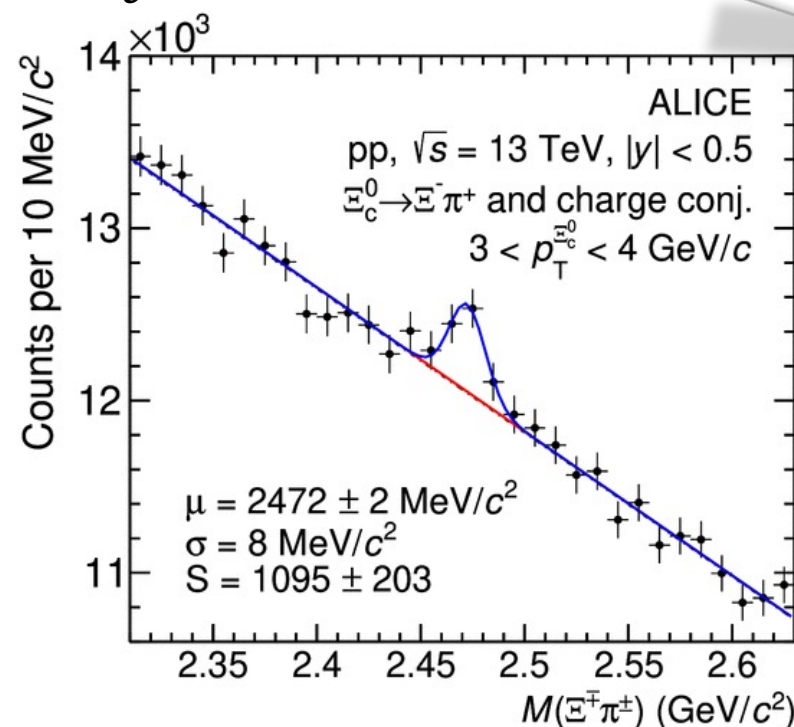


Narrow peak around 6150 MeV (more than 5σ significance) consistent with $\Lambda_b(6146)^0$ and $\Lambda_b(6152)^0$, but also with a single peak hypothesis.
 $m_{\Lambda_b(6146)^0}=6146.5\pm 1.9\pm 0.8\pm 0.2$ MeV
 $m_{\Lambda_b(6152)^0}=6152.7\pm 1.1\pm 0.4\pm 0.2$ MeV

In agreement with the LHCb measurement

Ξ_c and Ξ_{cc} states

Results from ALICE [6] in the $\Xi_c^0 \rightarrow \Xi^\mp \pi^\pm$ at 13 TeV

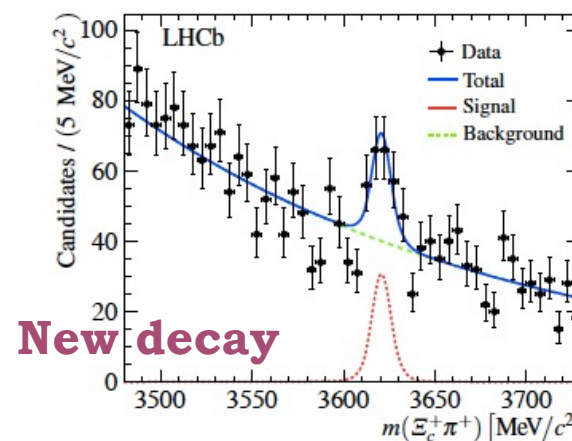


ALI-PUB-488829

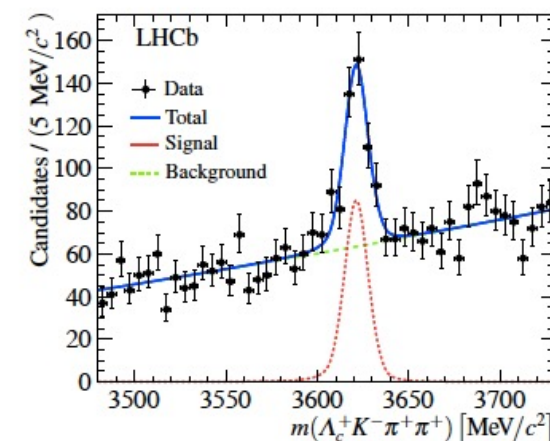
$\Xi_c^0 \rightarrow \Xi^\mp \pi^\pm$ and charge conjugation in $3 < p_T < 4$ GeV in pp collisions



First measurement of $\Xi_{cc}^{++} \rightarrow \Xi_c^+ \pi^+$ critical towards understanding the dynamics of weak decays of doubly heavy baryons [7]



$\Xi_{cc}^{++} \rightarrow \Xi_c^+ \pi^+$ reconstruction with $m(\Xi_{cc}^{++}) = 3620.6 \pm 1.5 \pm 0.4 \pm 0.3$ MeV



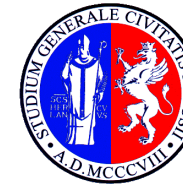
$\Xi_{cc}^{++} \rightarrow \Lambda_c^+ K^- \pi^+ \pi^+$ reconstruction with $m(\Xi_{cc}^{++}) = 3621.40 \pm 0.72 \pm 0.27 \pm 0.14$ MeV

Average of the Ξ_{cc}^{++} mass: $3621.24 \pm 0.65 \pm 0.31$ MeV

[6] [arXiv:2105.05187](https://arxiv.org/abs/2105.05187)

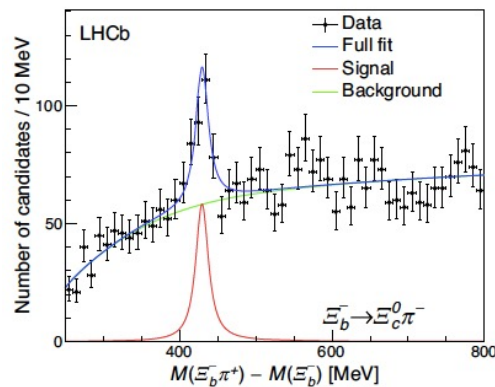
[7] PRL **121**, 162002

New $\Xi_b^{0,-}$ states

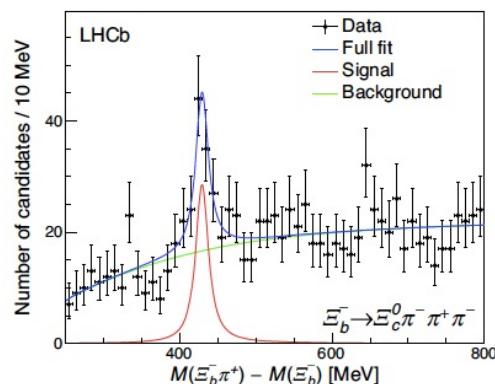


- New excited Ξ_b^0 state $\Xi_b(6227)^0$ observed in the $\Xi_b^-\pi^+$ using Run1 and Run2 data (L=8.5/fb) [8]
- Two Ξ_b^- decays: $\Xi_b^- \rightarrow \Xi_c^0\pi^-$ and $\Xi_b^- \rightarrow \Xi_c^0\pi^-\pi^+\pi^-$

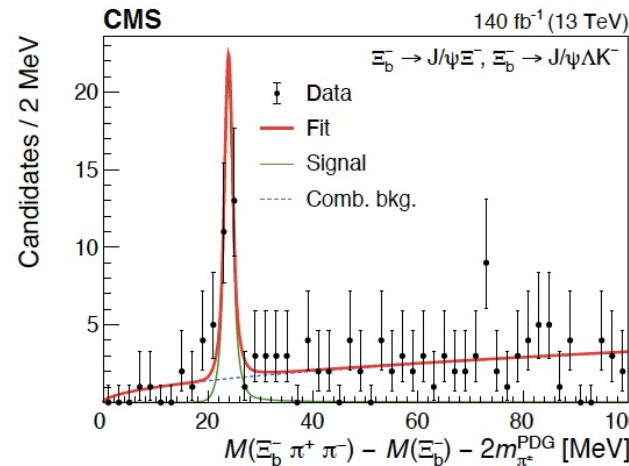
- New excited Ξ_b^- state $\Xi_b(6100)^-$ observed in the $\Xi_b^-\pi^+\pi^-$ using Run2 data (L=140/fb) [9]
- Two Ξ_b^- decays: $\Xi_b^- \rightarrow \Xi^- J/\psi$ and $\Xi_b^- \rightarrow J/\psi \Lambda K^-$



$\delta M_\pi = M(\Xi_b^-\pi^+) - M(\Xi_b^-)$
 \Rightarrow Clear peak shown in the two decays with 10σ significance



Identified as $\Xi_b(6227)^0$ with:
 $\delta m_\pi^{peak} = 429.8_{-1.5}^{+1.4} \text{ MeV}$
 $m(\Xi_b(6227)^0) = 6227.1_{-1.5}^{+1.4} \text{ MeV}$
 $\Gamma(\Xi_b(6227)^0) = 18.6_{-4.1}^{+5.0} \text{ MeV}$



$\Delta M = M(\Xi_b^-\pi^+\pi^-) - M(\Xi_b^-) + 2m_\pi^{PDG}$
 \Rightarrow Clear peak shown in the two decays with $> 6\sigma$ significance

Identified as $\Xi_b(6100)^-$ with:
 $\Delta M = 24.14 \pm 0.22 \pm 0.05 \text{ MeV}$
 $m(\Xi_b(6100)^-) = 6100.3 \pm 0.2 \pm 0.1 \pm 0.6 \text{ MeV}$
 Natural width consistent with 0
 Consistent with $J^P=3/2^-$

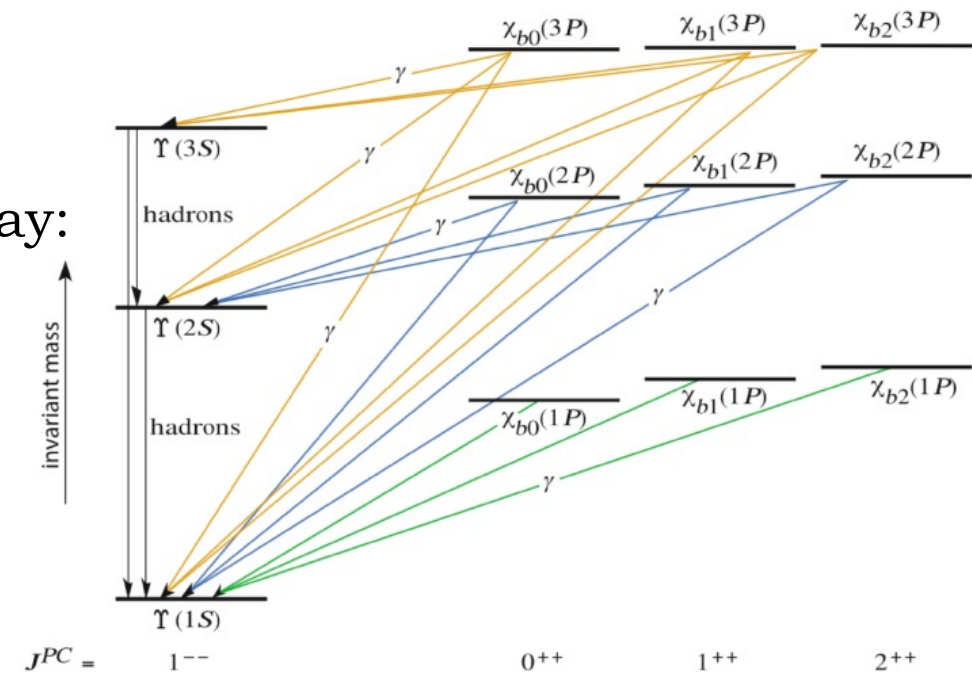
[8] PHYSICAL REVIEW D 103, 012004

[9] CMS-BPH-20-004 accepted by PRL

Bottomonium



- The bottomonium family ($b\bar{b}$) plays a special role in understanding how the QCD binds quarks into hadrons
- The $\chi_b(3P)$ state is especially interesting since its properties could be affected by the proximity of the open-beauty threshold
- The $\chi_b(3P)$ state is reconstructed through the decay:
 $\chi_b \rightarrow \Upsilon(nS)\gamma \rightarrow 2\mu\gamma$ where $n = 1, 2, 3$

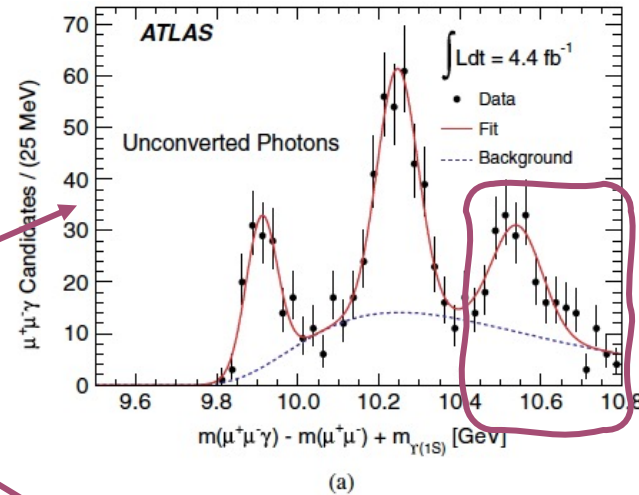
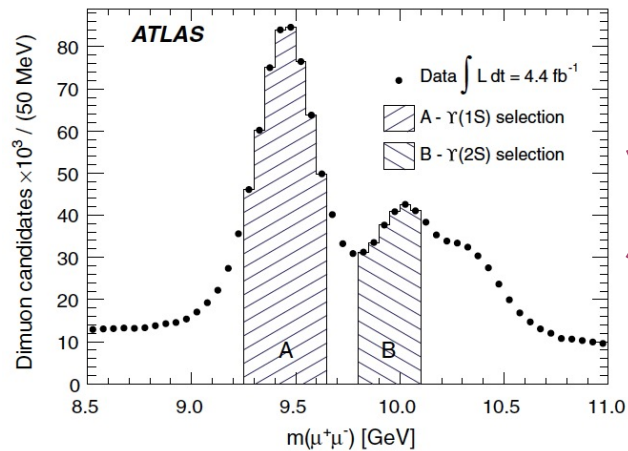


Picture from : V. Knünz, Measurement of Quarkonium Polarization to Probe QCD - DOI 10.1007/978-3-319-49935-2_2

$\chi_b(3P)$



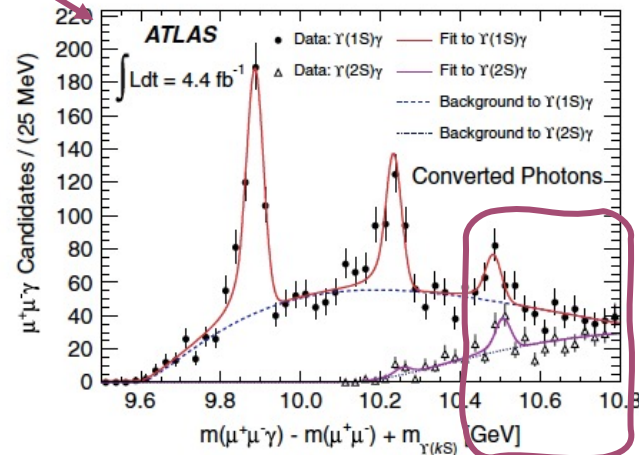
- **Firstly observed** by ATLAS in [10] at 7 TeV through the radiative decays $\chi_b(nP) \rightarrow \Upsilon(1S)\gamma$ and $\chi_b(nP) \rightarrow \Upsilon(2S)\gamma$



$\chi_b(3P)$ identified as third peak around 10.5 GeV

Significance of about 6σ

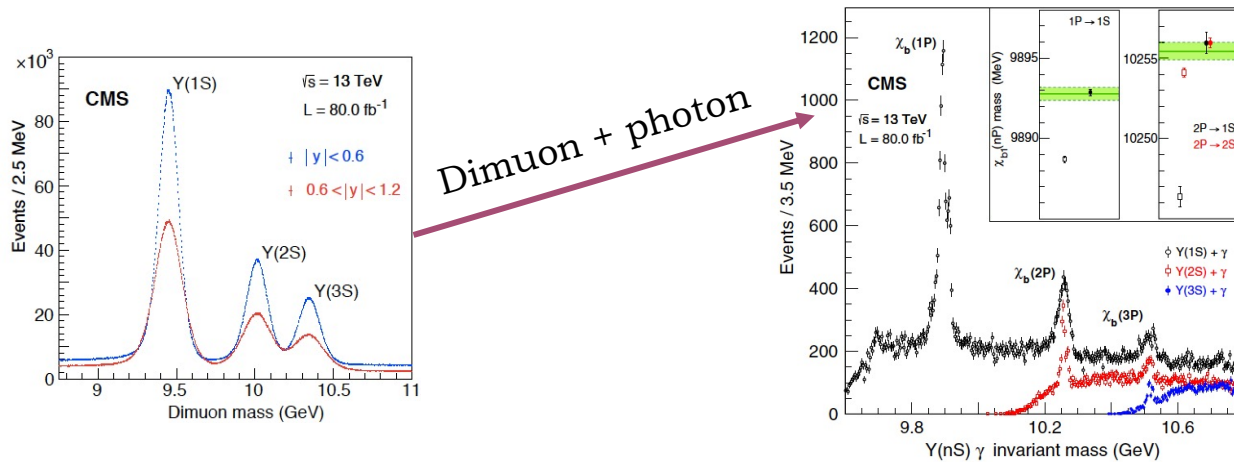
Mass of the $\chi_b(3P)$ baricenter
 $m = 10.530 \pm 0.005$ (stat) ± 0.009 (syst)
 MeV



$\chi_b(3P)$

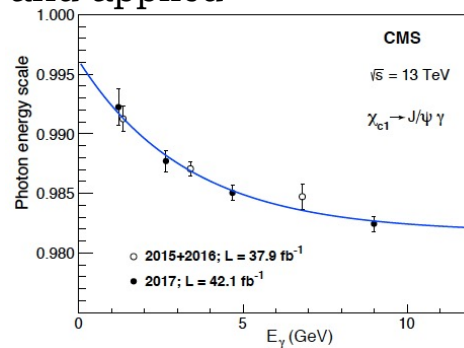


The **first observation of the $\chi_{b1}(3P)$ and $\chi_{b2}(3P)$ resolved states** performed with 13 TeV data from 2015 to 2017 ($L=80/\text{fb}$) [11]



Dimuon + photon

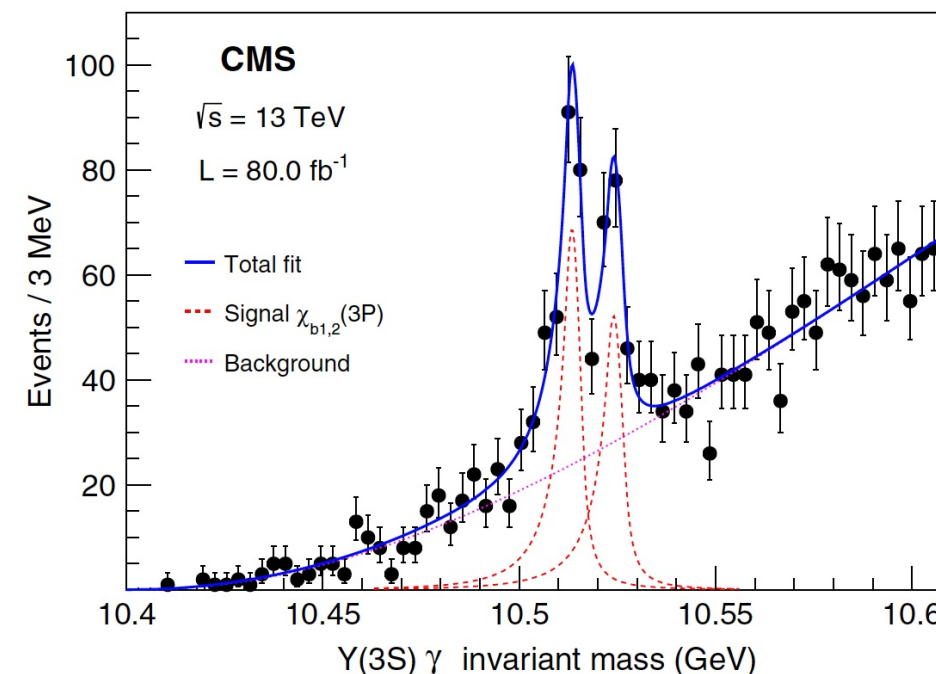
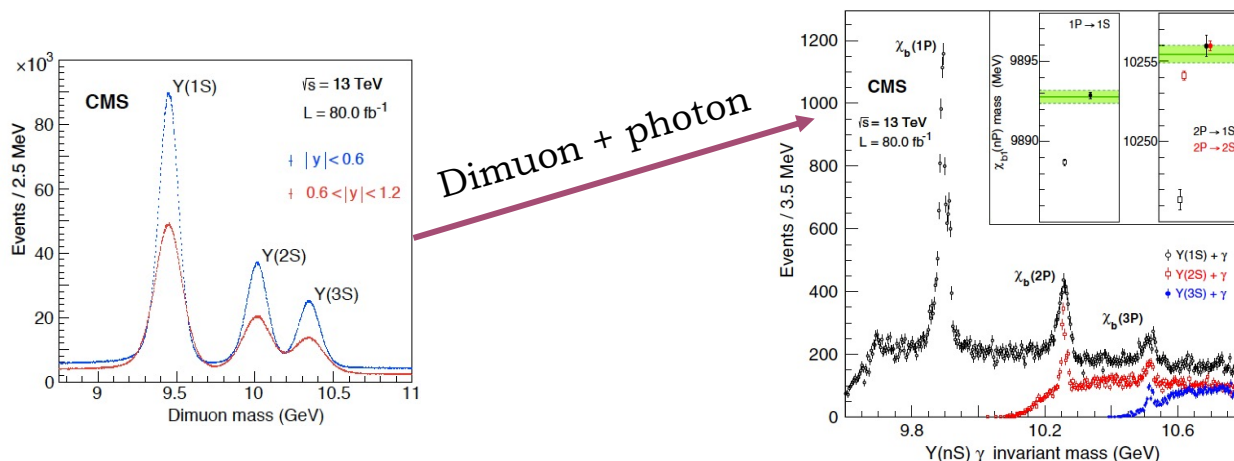
Photons energy scale calibrated in a control sample and applied event by event



$\chi_b(3P)$



The **first observation of the $\chi_{b1}(3P)$ and $\chi_{b2}(3P)$ resolved states** performed with 13 TeV data from 2015 to 2017 ($L=80/\text{fb}$) [11]

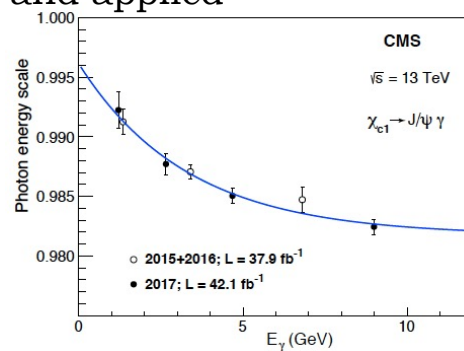


$$M(\chi_{b1}(3P)) = 10513.42 \pm 0.41 \text{ (stat)} \pm 0.18 \text{ (syst)} \pm 0.5 \text{ (Y(3S)) MeV}$$

$$M(\chi_{b2}(3P)) = 10524.02 \pm 0.57 \text{ (stat)} \pm 0.18 \text{ (syst)} \pm 0.5 \text{ (Y(3S)) MeV}$$

$$\Delta M = 10.60 \pm 0.64 \text{ (stat)} \pm 0.17 \text{ (syst) MeV}$$

compatible with the theoretical expectations

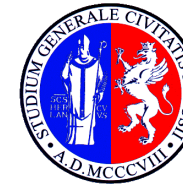


Conclusion



- I showed the latest results on heavy flavor classical spectroscopy from the four LHC collaborations: ALICE, ATLAS, CMS and LHCb
- Hadron spectroscopy and help to shed light on the mechanisms responsible for dynamics of quarks and baryon formation.
- Very important results have been produced, mainly from Run2 data, with an impressive precision and sensitivity
- Heavy flavor spectroscopy is a very active field: much more results has to come!

Conclusion



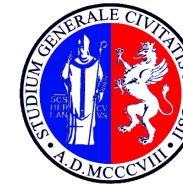
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- Hadron spectroscopy and help to shed light on the mechanisms responsible for dynamics of quarks and baryon formation.
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STAY TUNED

backup



$B_c^{*+}(2S)$

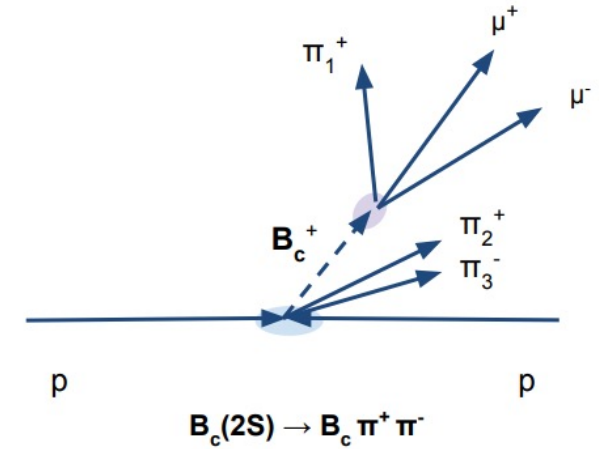


Recently also the relative cross sections of the $B_c^+(2S)$ and $B_c^{*+}(2S)$ states with respect to the B_c^+ have been measured

$$R^+ \equiv \frac{\sigma(B_c^+(2S))}{\sigma(B_c^+)} \cdot \mathcal{B}(B_c^+(2S) \rightarrow B_c^+ \pi^+ \pi^-) = \frac{N(B_c^+(2S))}{N(B_c^+)} \frac{\epsilon(B_c^+)}{\epsilon(B_c^+(2S))},$$

$$R^{*+} \equiv \frac{\sigma(B_c^{*+}(2S))}{\sigma(B_c^+)} \cdot \mathcal{B}(B_c^{*+}(2S) \rightarrow B_c^+ \pi^+ \pi^-) = \frac{N(B_c^{*+}(2S))}{N(B_c^+)} \frac{\epsilon(B_c^+)}{\epsilon(B_c^{*+}(2S))},$$

$$R^{*+}/R^+ = \frac{\sigma(B_c^{*+}(2S))}{\sigma(B_c^+(2S))} \cdot \frac{\mathcal{B}(B_c^{*+}(2S) \rightarrow B_c^+ \pi^+ \pi^-)}{\mathcal{B}(B_c^+(2S) \rightarrow B_c^+ \pi^+ \pi^-)} = \frac{N(B_c^{*+}(2S))}{N(B_c^+(2S))} \frac{\epsilon(B_c^+(2S))}{\epsilon(B_c^{*+}(2S))}.$$

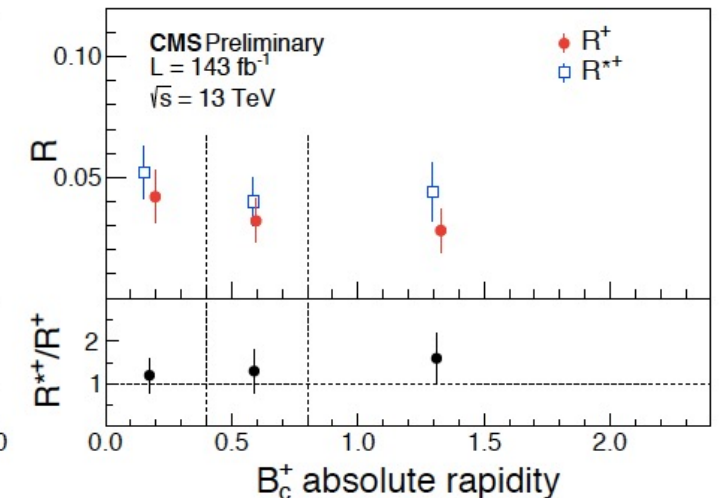
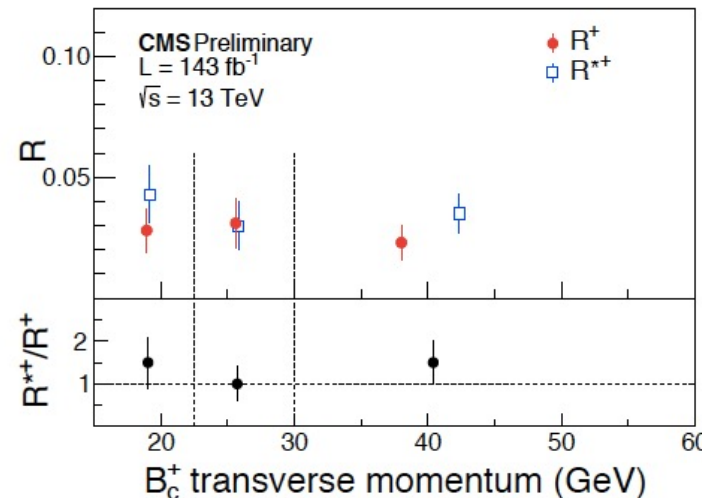


$$R^+ = 3.57 \pm 0.69 \text{ (stat)} \pm 0.32 \text{ (syst)} \%,$$

$$R^{*+} = 4.91 \pm 0.69 \text{ (stat)} \pm 0.57 \text{ (syst)} \%,$$

$$R^{*+}/R^+ = 1.39 \pm 0.35 \text{ (stat)} \pm 0.09 \text{ (syst)},$$

[2] CMS PAS BPH-19-001



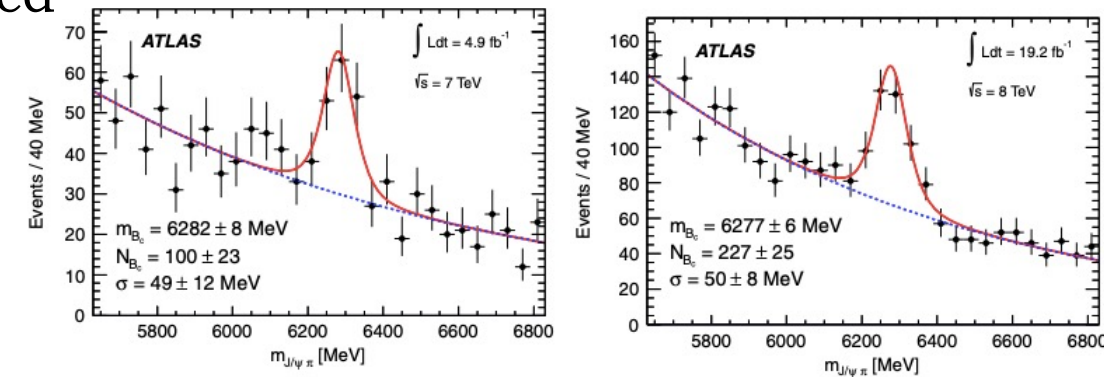
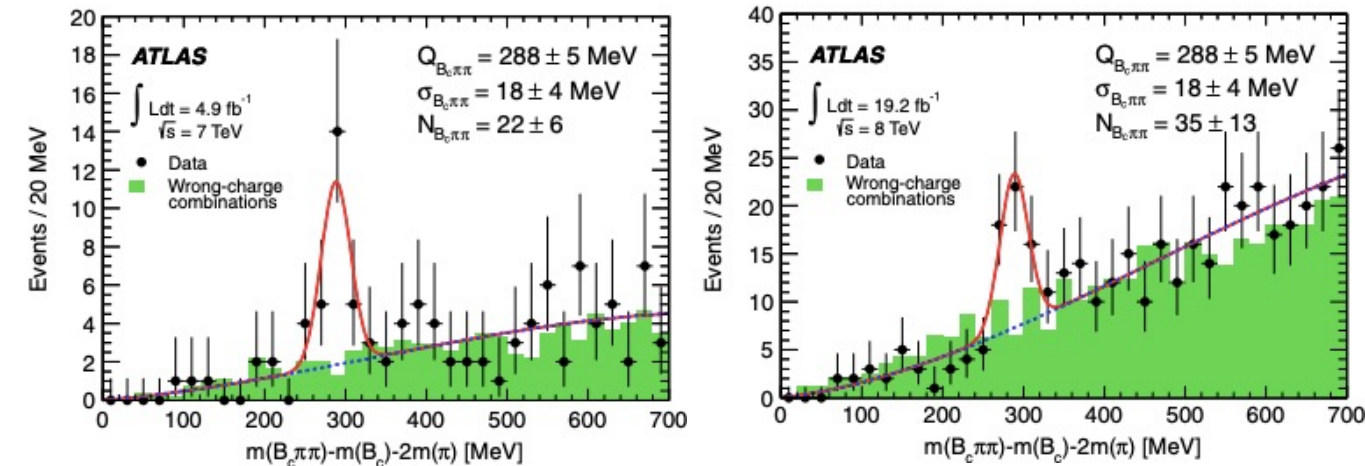
$B_c^*(2S)$



Observed for the first time by ATLAS in 2014 using Run1 data (7 + 8 TeV) [*].

The B_c^\pm selection criteria for the events are optimized separately for 7 and 8 TeV

[*] Phys. Rev. Lett. 113, 212004 (2014)



mass difference distribution:

$$m(B_c^\pm \pi \pi) - m(B_c^\pm) - 2m(\pi)$$

Significance of the observation: 3.7σ @7TeV
and 4.5σ @8 TeV data.

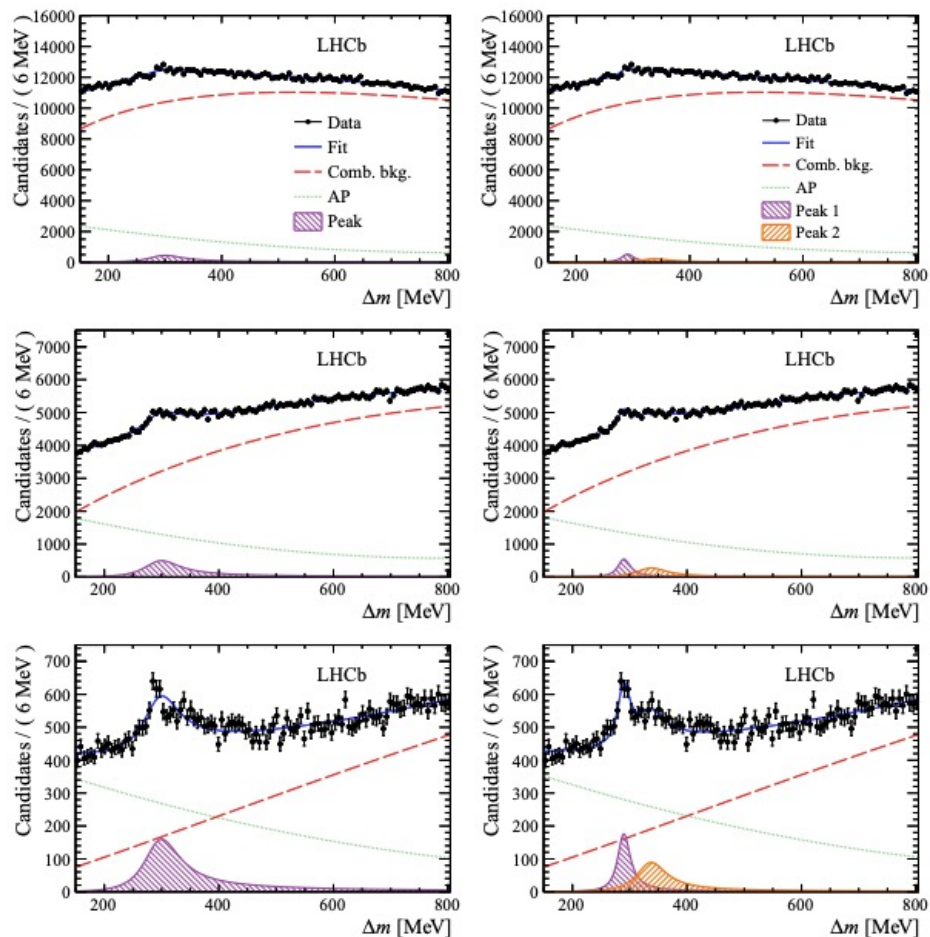
Combined 7 + 8 TeV significance 5.2σ

Consistent with the predicted mass of the $B_c^\pm(2S)$ state.

B_s^0



Structure observed in the B^+K^- mass spectrum using data from Run1+Run2 (9 fb^{-1}) [*]



$K \text{ pT} \in [0.5;1] \text{ GeV}$

Significance of single peak wrt null hypothesis > 20 sigma, double peak wrt single 7.7 sigma

$K \text{ pT} \in [1;2] \text{ GeV}$

$$\begin{aligned} m_1 &= 6063.5 \pm 1.2 \text{ (stat)} \pm 0.8 \text{ (syst)} \text{ MeV,} \\ \Gamma_1 &= 26 \pm 4 \text{ (stat)} \pm 4 \text{ (syst)} \text{ MeV,} \\ m_2 &= 6114 \pm 3 \text{ (stat)} \pm 5 \text{ (syst)} \text{ MeV,} \\ \Gamma_2 &= 66 \pm 18 \text{ (stat)} \pm 21 \text{ (syst)} \text{ MeV.} \end{aligned}$$

$K \text{ pT} > 2 \text{ GeV}$

$$\Delta m = m_{BK} - m_B - m_K$$

[*] LHCb-PAPER-2020-026

$B^*(2S) - CMS$



Selection:

- Two muons: $|\eta| < 2.5$, $p_T > 4$ GeV, distance of closest approach between the two muons < 0.5 cm, dimuon vertex fit $\chi^2 > 10\%$, dimuon invariant mass in the range 2.9–3.3 GeV, distance between the dimuon vertex and the beam axis larger than three times its uncertainty, dimuon p_T aligned with the transverse displacement vector $\cos \theta > 0.98$
- third track produced at the dimuon vertex with: $|\eta| < 2.4$, $p_T > 3.5$ GeV, ≥ 1 hit in the pixel layers, ≥ 5 hits in the tracker, transverse IP $>$ two times its uncertainty
- B_c^\pm candidate obtained from a kinematic fit between the 2 muons and the track: $p_T > 15$ GeV, $|y| < 2.4$, $l > 100 \mu\text{m}$, and a kinematic fit χ^2 probability $> 10\%$
- $B_c^\pm(2S)$ and $B_c^{\pm*}(2S)$ candidates reconstructed from kinematic fit between B_c^\pm candidate with two opposite-sign tracks with a common vertex.
Only B_c^\pm candidates with invariant mass in the range 6.2–6.355 GeV are selected.
One of the pion candidates must have $p_T > 0.8$ GeV and the other $p_T > 0.6$ GeV.
- The $B_c^\pm \pi^+ \pi^-$ candidates must have $|y| < 2.4$ and a vertex χ^2 probability larger than 10%

$B^{*}(2S) - LHCb$



Selection:

- Two muons with $p_T > 550$ MeV and good track-fit quality. Required to form a common decay vertex with an invariant mass in the range 3040-3140 MeV
- Pion with $p_T > 1000$ MeV, good track-fit quality, and be inconsistent with originating from any PV
- The J/ψ candidate is combined with a charged pion to form the B_c^\pm candidate.
- The B_c^\pm candidate is required to have a good-quality vertex, a trajectory consistent with coming from PV, and a decay time larger than 0.2 ps.
- A BDT is used to further suppress bkg (pt / eta / decay length / decay time / χ^2)
- B_c^\pm candidates with M in 6200--6320 MeV are combined with a pair of oppositely charged particles identified as pions. These pion candidates are required to originate from the PV, and each have $p_T > 300$ MeV, $p > 1500$ MeV, and a good track-fit quality.
- The $B_c^{\pm*}(2S)$ candidate is required to have a good vertex-fit quality

Λ_b^0 - LHCb



- Samples of Λ_b^0 candidates are formed from $\Lambda_c^+\pi^-$ combinations, where the Λ_c^+ baryon is reconstructed in the $pK^-\pi^+$ final state.
- The reconstructed Λ_c^+ vertex is required to have a good fit quality and to be significantly displaced from all PVs and to have a mass of $\pm 25\text{MeV}$ wrt known value
- BDT used to reduce bkg (inputs: kinematic variables of the Λ_c^+ and Λ_b^0 , the lifetime of the Λ_b^0 , kinematic variables and quality of particle identification for pions, kaons, and protons, and variables describing the consistency of the selected candidates)
- Selected Λ_b^0 candidates are combined with pairs of pions originating from the same PV as the Λ_b^0 . Only pion pairs with $p_T^{\pi^+\pi^-} > 500\text{ MeV}$ are used, to suppress the otherwise large combinatorial

Λ_b^0 - CMS

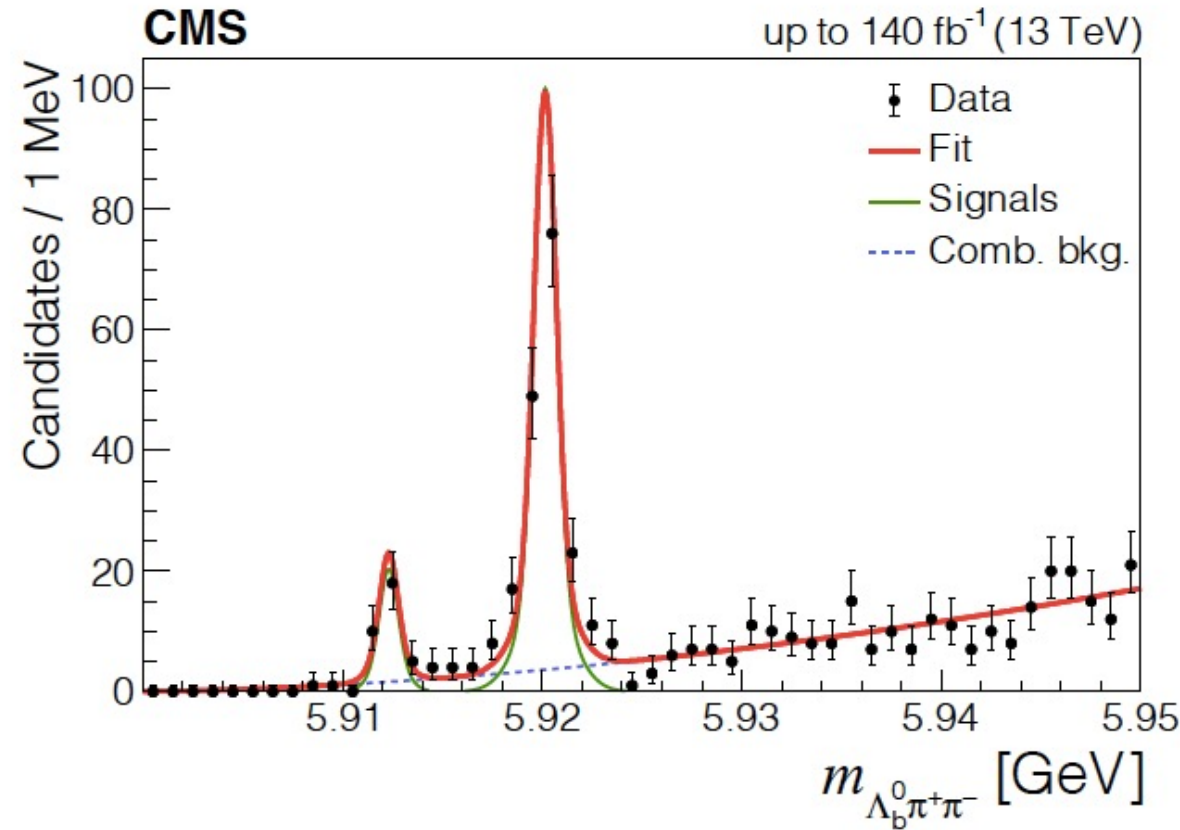


$\Lambda_b^0 \rightarrow \psi \Lambda$ ($\mu^+ \mu^- + \Lambda \rightarrow p \pi^-$ from V0 collection)
 Λ vtx fit probability $> 1\%$, $p_T(\Lambda) > 1$ GeV,
 $\cos(3D \Lambda \text{ pointing angle to } \Lambda_b^0 \text{ vertex}) > 0.99$,
 $L_{xy}/\sigma_{Lxy}(\Lambda \text{ vtx} \rightarrow \Lambda_b^0 \text{ vtx}) > 5$
 Λ mass within 10 MeV from PDG

Λ_b^0 vtx fit probability $> 1\%$
 $p_T(\Lambda_b^0) > 5$ GeV,
 $\cos(2D \Lambda_b^0 \text{ pointing angle to PV}) > 0.99$,
 $L_{xy}/\sigma_{Lxy}(\Lambda_b^0 \text{ vtx} \rightarrow \text{PV}) > 3$

Additional channel $\Lambda_b^0 \rightarrow \psi(2S) \Lambda$, $\psi(2S) \rightarrow J/\psi \pi^+ \pi^-$, $J/\psi \rightarrow \mu^+ \mu^-$: as above, additional 2 pions – high-purity tracks,
 $p_T > 0.35$ GeV; Λ_b^0 obtained by vertexing $\mu^+ \mu^- \pi^+ \pi^- \Lambda$ with J/ψ and Λ mass constraints; $3672 < M(J/\psi \pi^+ \pi^-) < 3700$ MeV

Λ_b^0 - CMS



$$M(\Lambda_b(5912)^0) = 5912.32 \pm 0.12 \pm 0.01 \pm 0.17 \text{ MeV},$$
$$M(\Lambda_b(5920)^0) = 5920.16 \pm 0.07 \pm 0.01 \pm 0.17 \text{ MeV},$$

Λ_b^0 - LHCb



Table 1: Quark-model predictions for the masses of the lightest Λ_b and $\Sigma_b^{(*)}$ states (in MeV).

Baryon	State	J^P	Ref. [15]	Ref. [16]	Ref. [17]	Ref. [18]
Λ_b^0	1S	$\frac{1}{2}^+$	5585	5612	5620	5619
	1P	$\frac{1}{2}^-$	5912	5939	5930	5911
		$\frac{3}{2}^-$	5920	5941	5942	5920
	2S	$\frac{1}{2}^+$	6045	6107	6089	
	1D	$\frac{3}{2}^+$	6145	6181	6190	6147
		$\frac{5}{2}^+$	6165	6183	6196	6153
$\Sigma_b^{(*)0}$	1S	$\frac{1}{2}^+$	5795	5833	5800	
		$\frac{3}{2}^+$	5805	5858	5834	
	1P	$\frac{1}{2}^-$	6070	6099	6101	
		$\frac{3}{2}^-$	6070	6101	6096	
		$\frac{5}{2}^-$	6090	6172	6084	
	2S	$\frac{1}{2}^+$	6200	6294	6213	
		$\frac{3}{2}^+$	6250	6308	6226	

Λ_b^0 - LHCb



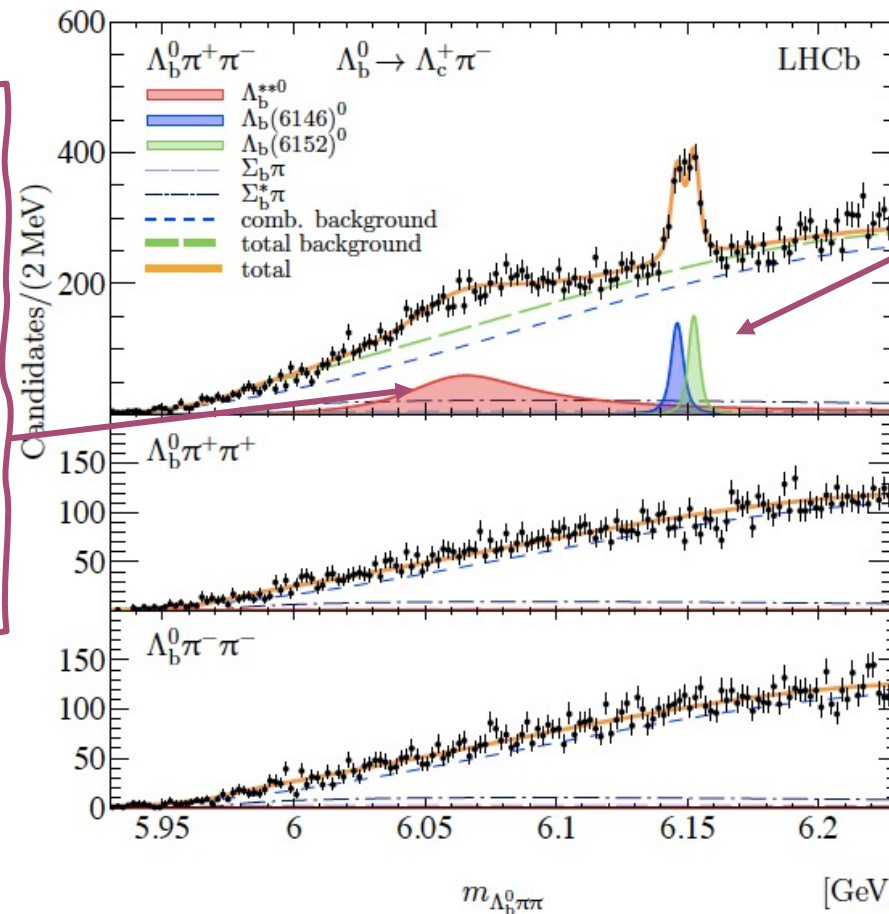
- Data collected at 7, 8 and 13 TeV ($L=9/\text{fb}$) \Rightarrow new state observed in the $\Lambda_b^0 \pi^+ \pi^-$ spectrum [*]
- Two Λ_b^0 decays: $\Lambda_b^0 \rightarrow \Lambda_c^+ \pi^-$ and $\Lambda_b^0 \rightarrow J/\psi p K^-$

Broad excess referred as Λ_b^{**0}

Mass difference wrt Λ_b^0
 $\Delta m = 452.7 \pm 2.9 \pm 0.5 \text{ MeV}$
 Width = $72 \pm 11 \pm 2 \text{ MeV}$

Λ_b^{**0} mass = $6072.3 \pm 2.9 \pm 0.6 \pm 0.2 \text{ MeV}$

Consistent with the CMS measurement



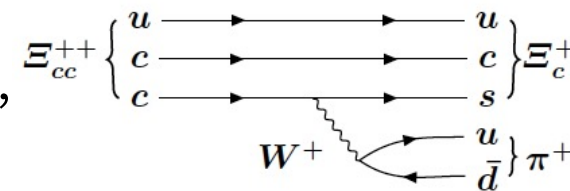
Narrow peaks already measured as $\Lambda_b(6146)^0$ and $\Lambda_b(6152)^0$

Can't exclude that the broad structure corresponds to a superposition of more than one narrow states, but the interpretation of these states as excited Σ_b resonances is disfavoured.

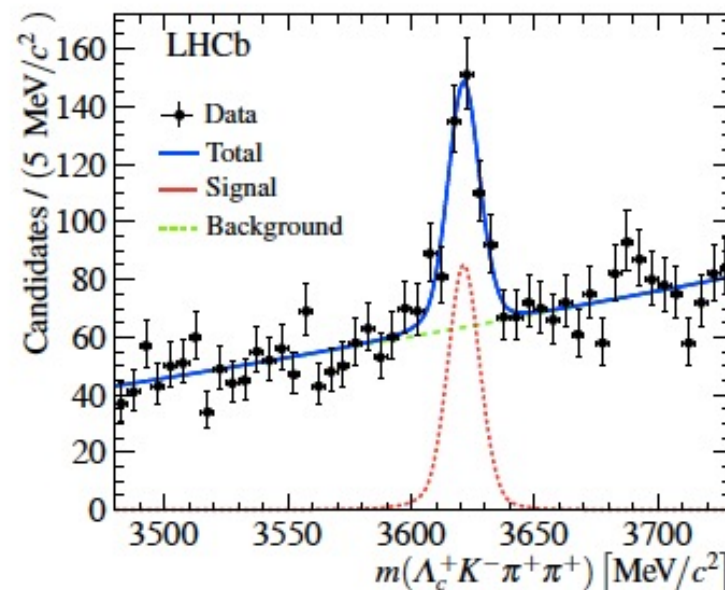
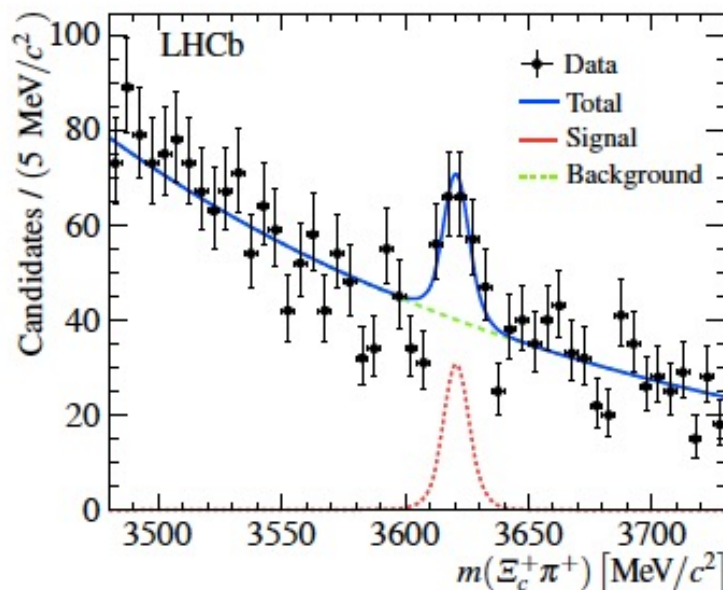
Ξ_{cc}^{++}



- Measured for the first time by LHCb in [10] in the $\Xi_{cc}^{++} \rightarrow \Lambda_c^+ K^- \pi^+ \pi^+$ decay with $m = 3621.40 \pm 0.72 \pm 0.27 \pm 0.14$ MeV and its lifetime $0.256^{+0.024}_{-0.022}$ ps \Rightarrow compatible with a weak interacting particle
- The first measurement of the decay $\Xi_{cc}^{++} \rightarrow \Xi_c^+ \pi^+$ is presented here [11], critical towards understanding the dynamics of weak decays of doubly heavy baryons



$\Xi_{cc}^{++} \rightarrow \Xi_c^+ \pi^+$
reconstruction with
 $m(\Xi_{cc}^{++}) = 3620.6 \pm 1.5 \pm 0.4 \pm 0.3$ MeV



$\Xi_{cc}^{++} \rightarrow \Lambda_c^+ K^- \pi^+ \pi^+$
reconstruction with
 $m(\Xi_{cc}^{++}) = 3621.40 \pm 0.72 \pm 0.27 \pm 0.14$ MeV

Average of the Ξ_{cc}^{++} mass: $3621.24 \pm 0.65 \pm 0.31$ MeV

[10] Phys. Rev. Lett. **119**, 112001

[11] Phys. Rev. Lett. **121**, 162002

Ξ_{cc}^{++}



Selection of $\Xi_{cc}^{++} \rightarrow \Xi_c^+ (\rightarrow p K^- \pi^+) \pi^+$ decay:

- three charged tracks with $p_T > 0.5 \text{ GeV}$ fitted to a common vertex
- Ξ_c^+ vertex displayed with lifetime $> 0.15 \text{ ps}$ and invariant mass between 2450-2488 MeV
- Additional positive track with $p_T > 0.2 \text{ GeV}$ combined to the Ξ_c^+
- Ξ_{cc}^{++} candidate with $p_T > 2 \text{ GeV}$ and originating from PV

BDT used to minimise combinatorial bkg with 3 sets of variables (particle kinematics, vertex fitting, lifetime)

Also the ratio between the two decays is measured as
with $\mathcal{R}(\mathcal{B}) = 0.035 \pm 0.009 \text{ (stat)} \pm 0.003 \text{ (syst)}$
in agreement with theory

$$\mathcal{R}(\mathcal{B}) \equiv \frac{\mathcal{B}(\Xi_{cc}^{++} \rightarrow \Xi_c^+ \pi^+) \times \mathcal{B}(\Xi_c^+ \rightarrow p K^- \pi^+)}{\mathcal{B}(\Xi_{cc}^{++} \rightarrow \Lambda_c^+ K^- \pi^+ \pi^+) \times \mathcal{B}(\Lambda_c^+ \rightarrow p K^- \pi^+)}$$

$\chi_b(3P)$



Selection:

- Two OS muons with $p_T > 4$ GeV and $|\eta| < 2.3$ fitted to a common vertex
- Dimuon $p_T > 12$ GeV and $|y| < 2$
- $9.25 < m_{\mu\mu} < 9.65$ GeV for the $\Upsilon(1S)$ and $9.80 < m_{\mu\mu} < 10.10$ GeV for the $\Upsilon(2S)$
- Photons are reconstructed as converted (in e^+e^-) or unconverted

Figure of merit $\tilde{m}_k = \Delta m + m_{\Upsilon(kS)}$