Charming CP Violation

Yuval Grossman

Cornell

The birth of charm CPV

$$\Delta A_{CP} \equiv a_{K^+K^-} - a_{\pi^+\pi^-} = (-1.54 \pm 0.29) \times 10^{-3}$$

Charm CPV is still a baby

- Can the SM explain LHCb?
- What next for charm CPV?

The effective 2-gen SM

The effective 2-generation SM

- Kaon and charm physics: only the first two generation are on-shell
- In many cases we can forget about the 3rd generation
- In some cases, like for CPV, we cannot do it
- The effective 2-generation model: We work with an EFT with two generation that is valid below m_b
- There are two main effects
  - The 2×2 CKM is not unitary (NU)
  - There are NR terms, like four Fermi operators (box diagram)

In charm we only care about the NU of the 2×2 CKM

The effective 2×2 CKM

- Consider the non-unitary 2×2 block of the CKM

$$V \approx \begin{pmatrix} \cos\theta_C & \sin\theta_C \\ -\sin\theta_C + \cos\theta_C \Delta e^{i\gamma} & \cos\theta_C + \sin\theta_C \Delta e^{i\gamma} \end{pmatrix}$$

- CPV effects are proportional to the Non Unitarity (NU) parameter

$$\Delta = |V_{cb} V_{ub}| \sim \lambda^5 \qquad \lambda = \sin\theta_C \approx 0.2$$

- We also define

$$\lambda_i \equiv V_{ci} V_{ui}^* \qquad \varepsilon_{\rm NU} \equiv \frac{\Delta}{2\sin\theta_C} \approx 6 \times 10^{-4}$$

The small parameters for charm

We can map all the parameters in charm based on the following small parameters

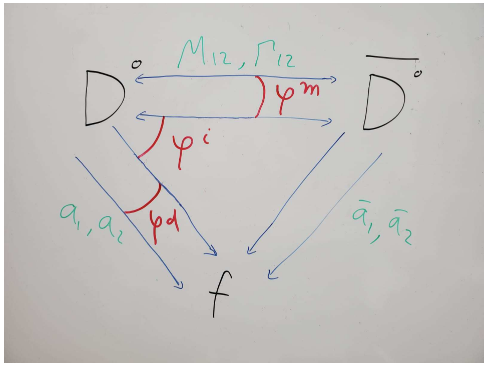
- Non-unitarity of the 2×2 CKM: $\varepsilon_{\rm NU} \sim 10^{-3}$
- SU(3)/U–spin breaking: $\varepsilon_{\rm SU(3)} \sim 0.2$
- The Wolfenstein parameter of the CKM: $\lambda \sim 0.2$
- For example [$x = \Delta M/\Gamma$, $y = \Delta\Gamma/(2\Gamma)$]
  - $x_{\rm th} \sim y_{\rm th} \sim \lambda^2 \varepsilon_{\rm SU(3)}^2 \sim 0.2\%$
  - $x_{\rm ex} \sim y_{\rm ex} \sim 0.5\%$

CP asymmetry

The time integrated CP asymmetry to leading order in x, y

$$a_f \equiv \frac{\Gamma(D \to f) - \Gamma(\overline{D} \to \overline{f})}{\Gamma(D \to f) + \Gamma(\overline{D} \to \overline{f})} \approx a_f^d + a^m + a_f^i$$

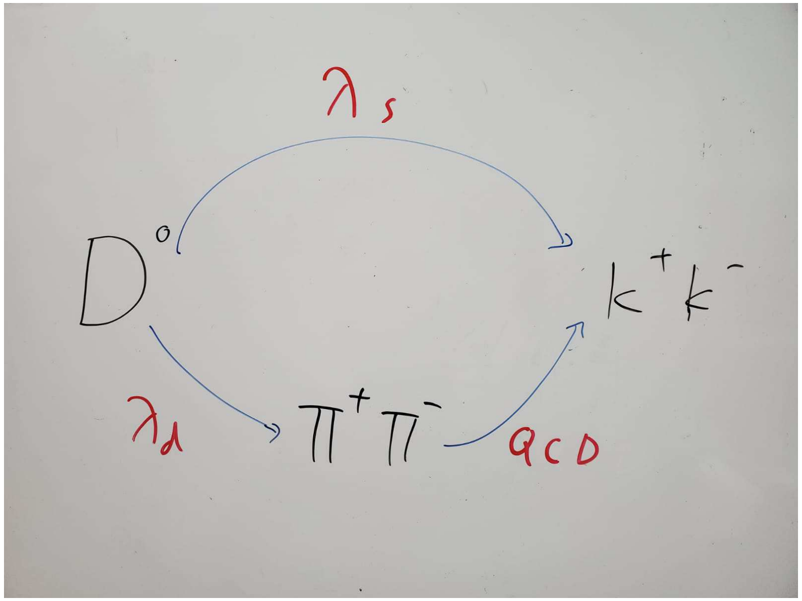
1. $a_f^d \sim r_f \sin\varphi_f^d$
2. $a^m \sim y \sin\varphi^m$
3. $a_f^i \sim x \sin\varphi_f^i$



- a^m is universal but a_f^d and a_f^i depend on f
- All the weak phases depend on $\varepsilon_{\rm NU}$
- In principle, each is a separate observable

CPV in decay

Interference via Rescattering



- We need two amplitude to interfere
- $\pi\pi$ represents many similar states like $\pi\rho$, $\rho\rho$
- Interference of trees with λ_s and λ_d
- We do not talk about penguins

Tree rescattering in Nature

The factors

$$\frac{\mathcal{A}(D \to \text{“}\pi\pi\text{''} \to KK)}{\mathcal{A}(D \to KK)} = \left(r_{\rm QCD} e^{i\delta}\right) \left(r_{\rm CKM} e^{i\varphi}\right)$$

$$a^d = 2(r_{\rm QCD} \sin\delta)(r_{\rm CKM} \sin\varphi)$$

- $r_{\rm QCD}$: ratio of rescattering amplitudes
- $\sin\delta = O(1)$: strong phase
- $r_{\rm CKM} = 1$: ratio of CKM factors, $|\lambda_d/\lambda_s|$
- $\sin\varphi \sim \varepsilon_{\rm NU} \sim 10^{-3}$: deviation from 2×2 unitarity

$$a^d \sim \varepsilon_{\rm NU} \times r_{\rm QCD} \sim 10^{-3} \times r_{\rm QCD}$$

The ratios

$$a^d \sim 10^{-3} \times r_{\rm QCD} \qquad r_{\rm QCD} \sim \left| \frac{\mathcal{A}(D \to \pi\pi \to KK)}{\mathcal{A}(D \to KK)} \right|$$

What is $r_{\rm QCD}$?

- Light Cone Sum Rules (LCSR)

$$r_{\rm QCD} \sim O\left(\frac{\alpha_s}{\pi}\right) \sim 10^{-1}$$

- Low energy QCD, rescattering is $O(1)$

$$r_{\rm QCD} \sim O(1)$$

What we learn from direct CPV

Within the SM the data implies $r_{\rm QCD} \sim 1$

- Theory: $a^d \sim 10^{-3} \times r_{\rm QCD}$
- Data: $a^d \sim 10^{-3}$

We conclude

- The assumption of large rescattering agrees with the data
- It is hard to argue that the LHCb result requires BSM
- Yet, BSM can still be present

Charm CPV agrees with the SM with large rescattering

CPV involving mixing (future)

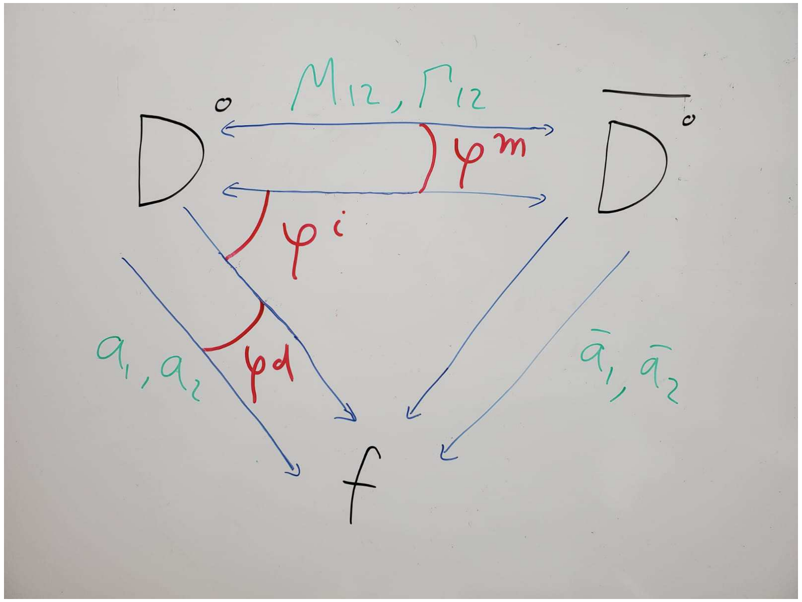
What are the phases in mixing

Kagan, Silvestrini, arXiv:2001.07207

YG et al. in preparation

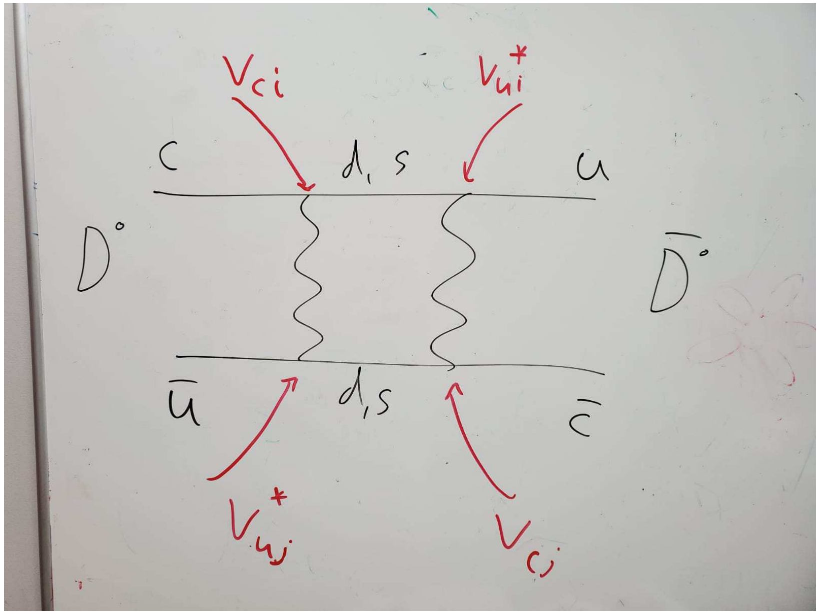
$$a_f \approx a_f^d + a^m + a_f^i$$

We care about a^m and a^i



The mixing amplitude

$$M_{12}, \Gamma_{12} \propto \lambda_s^2 f_{ss} + 2\lambda_s \lambda_d f_{sd} + \lambda_d^2 f_{dd}$$



- We cannot calculate f_{ij} reliably
- We can use SU(3)

Evaluation of the mixing amplitude

$$M_{12}, \Gamma_{12} \propto \lambda_s^2 f_{ss} + 2\lambda_s \lambda_d f_{sd} + \lambda_d^2 f_{dd}$$

- Vanishes in combined SU(3) $+$ two generation limit:
  - In the SU(3) limit: $f_{ss} = f_{dd} = f_{sd}$
  - In the 2-gen SM: $\lambda_s + \lambda_d = 0$
- It is non-zero only at second order of the breaking

$$M_{12}, \Gamma_{12} \sim \lambda^2 \left[\varepsilon_{\mathrm{SU(3)}}^2 + 2\varepsilon_{\mathrm{SU(3)}}\varepsilon_{\mathrm{NU}} + \varepsilon_{\mathrm{NU}}^2\right]$$

- Recall:

$$\varepsilon_{\mathrm{SU(3)}} \sim 0.2 \qquad \varepsilon_{\mathrm{NU}} \sim 10^{-3}$$

The phases of the mixing

$$M_{12}, \Gamma_{12} \sim \lambda^2 \left[\varepsilon_{\rm SU(3)}^2 + 2\varepsilon_{\rm SU(3)}\varepsilon_{\rm NU} + \varepsilon_{\rm NU}^2\right]$$

- The CPV phase enters with $\varepsilon_{\rm NU}$
- We can neglect the $\varepsilon_{\rm NU}^2$ term
- The mixing phases are

$$\arg(M_{12}) \sim \arg(\Gamma_{12}) \sim \frac{\varepsilon_{\rm NU}}{\varepsilon_{\rm SU(3)}}$$

- The phases of the decays are $O(\varepsilon_{\rm NU})$

The universal phase is enhanced

The prediction

- The relevant phases are

$$\phi^m \sim \frac{\varepsilon_{\rm NU}}{\varepsilon_{\rm SU(3)}} \qquad \phi_f^i \sim \frac{\varepsilon_{\rm NU}}{\varepsilon_{\rm SU(3)}} + (\varepsilon_{\rm NU})_f \qquad \phi_f^d \sim (\varepsilon_{\rm NU})_f$$

- To leading order in SU(3) breaking the time dependent asymmetries are universal
- Numerically, it is only a rough prediction
- It can be tested, hopefully soon
- We will learn something
  - If it fail, we found BSM
  - If it is confirmed, we will understand QCD better

Conclusion

A charming baby

- While we cannot do precision with charm, we can still learn a lot
- I think that we will hear about that baby in the future