



Extension of Glauber-like model for Proton-Proton collisions using anisotropic and inhomogeneous density profile

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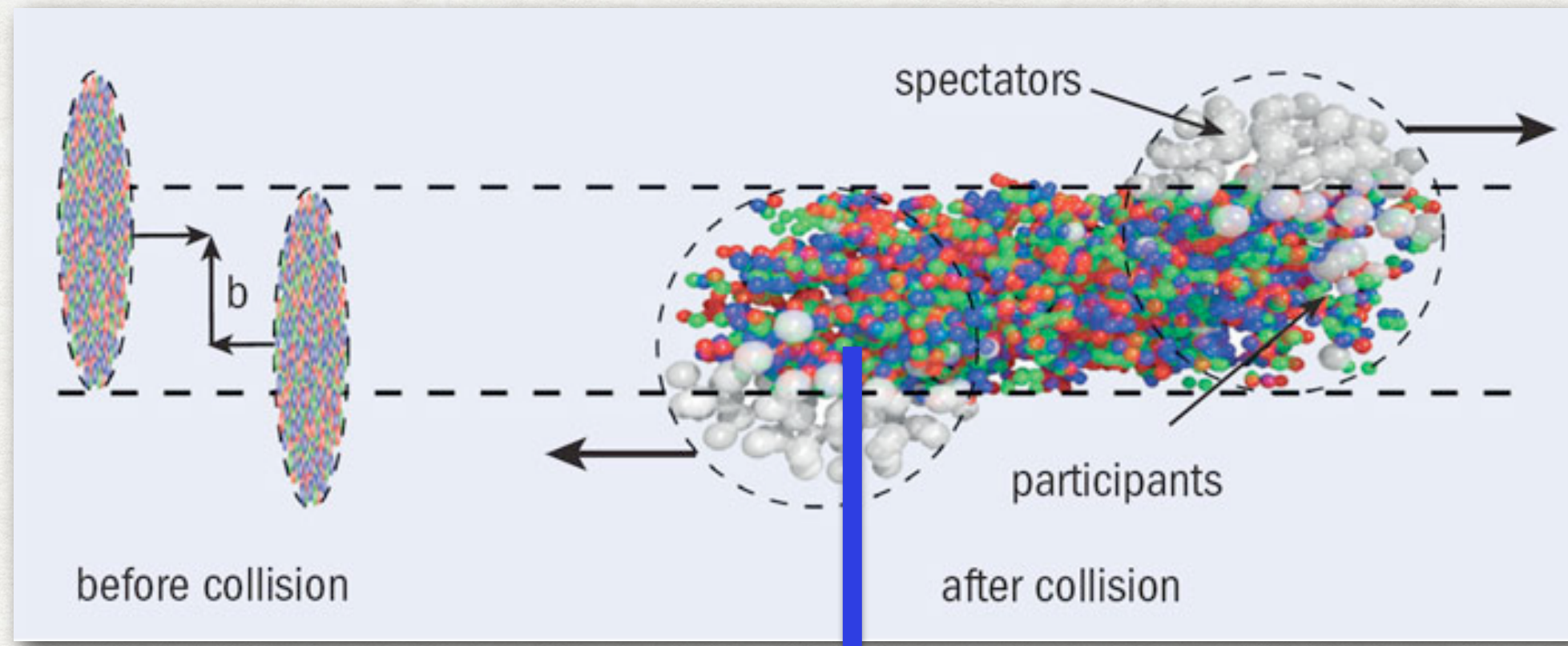
1. Introduction

- Results of relativistic proton proton (pp) collisions are traditionally used as base line for interpreting results of heavy ion collisions at relativistic energies to understand the Physics of Quark-Gluon Plasma (QGP)
- For instance, enhancement of strange hadrons and suppression in the number of J/Ψ in heavy ion collisions with respect to pp are taken as signatures of QGP in such collisions
- For such interpretations, it is assumed that, in pp collisions, no partonic medium is formed
- However recent results of pp collisions like enhancement of Multi-strange Particles and Multi-particle Ridge-like Correlations shows such assumptions may not be correct
- Understanding of pp collisions is crucial for characterization of the QCD medium formed in a heavy ion collision

2. Motivation

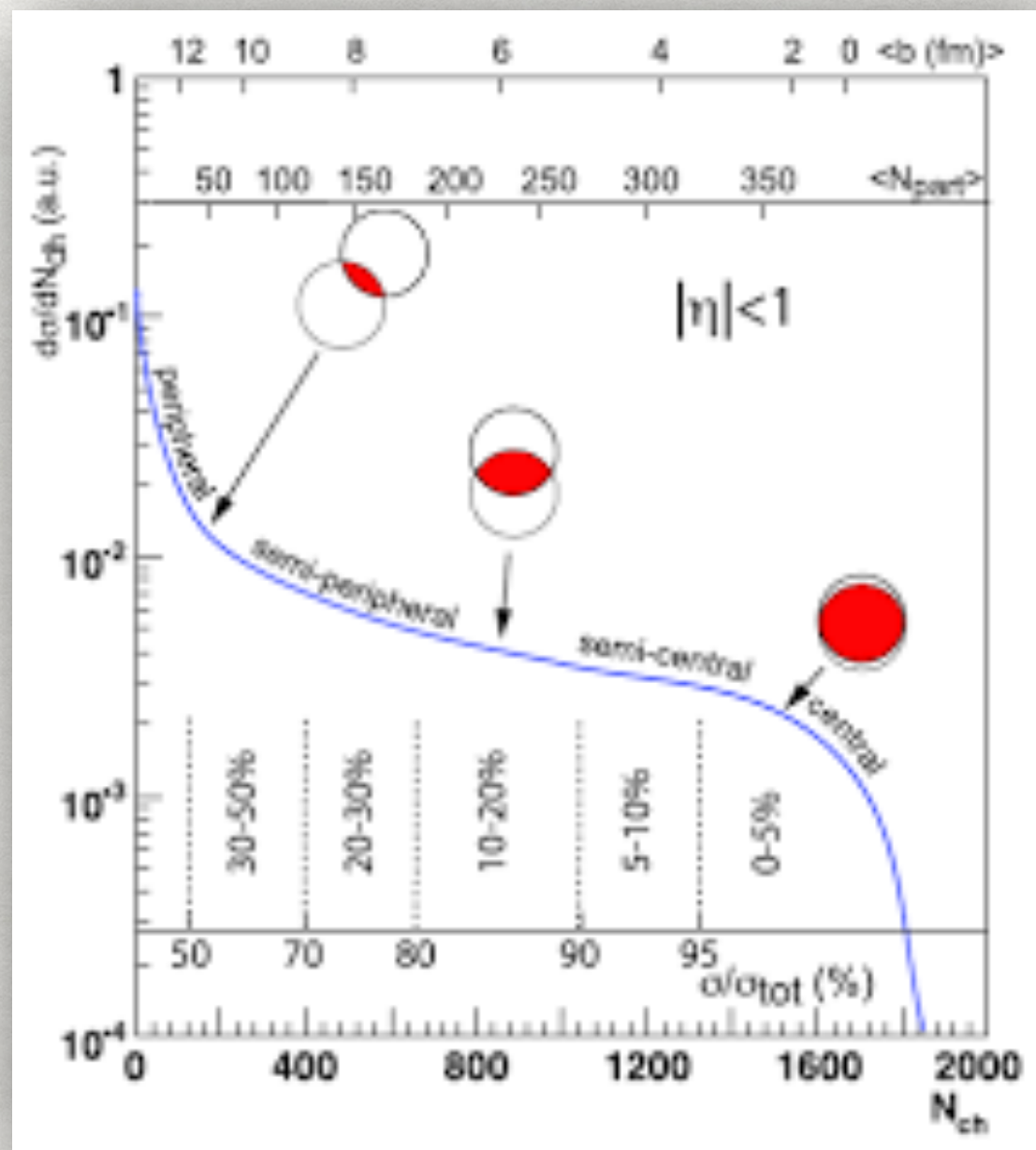
- The interpretation of results in the heavy-ion collisions, relies on models like Glauber model based on initial matter distribution resulting from the overlap of the two colliding nuclei at a given impact parameter (b)
- Thus, it is imperative to understand the initial condition of the medium formed in pp collisions in approximately similar way

3. Formalism: Glauber model for heavy-ion collision- main idea



Two-component model

$$\frac{dN_{ch}}{d\eta} \approx \left[(1-f) \frac{N_{part}}{2} + f N_{coll} \right]$$



4. Glauber Formalism from pp system

- Azimuthally asymmetric and inhomogeneous density distribution of a proton

$$\rho_{G-f}(\mathbf{r}; \mathbf{r}_1, \mathbf{r}_2, \mathbf{r}_3) = N_g \frac{1-\kappa}{3} \sum_{k=1}^3 \rho_q(\mathbf{r} - \mathbf{r}_k; r_q) + N_g \frac{\kappa}{3} \sum_{k=1}^3 \rho_g[\mathcal{R}^{-1}[\theta_k, \phi_k](\mathbf{r} - \frac{\mathbf{r}_k}{2}; r_q, \frac{r_k}{2}]$$

where, $\rho_q(\mathbf{r}; r_q) = \frac{1}{(2\pi)^{3/2} r_q^3} e^{-\frac{r^2}{2r_q^2}}$, $\rho_g(\mathbf{r}; r_s, r_l) = \frac{1}{(2\pi)^{3/2} r_s^2 r_l} e^{-\frac{x^2+y^2}{2r_s^2} - \frac{z^2}{2r_l^2}}$

$\mathcal{R}[\theta, \phi]$ transform vector (0,0,1) to $(\cos\phi \sin\theta, \sin\phi \cos\theta, \cos\theta)$

- Thickness function:

$$T(x, y) = \int \rho(x, y, z) dz$$

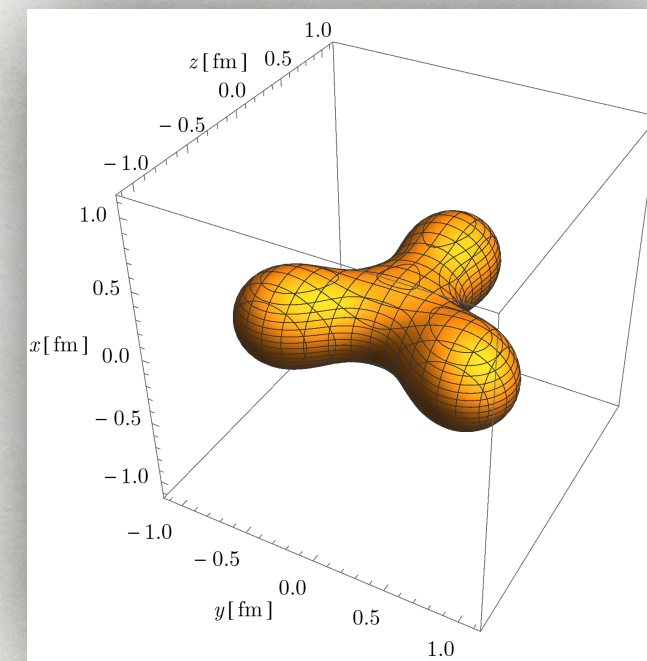
- Overlap function:

$$T_{pp}(b) = \iint T_A(x - \frac{b}{2}, y) T_B(x + \frac{b}{2}, y) dx dy$$

- Different orientation of three quarks within proton are considered

$$\begin{aligned} \mathbf{r}_1 &= \left(\frac{d}{2} \cos\left(\frac{\pi}{3} + \psi\right), \frac{d}{2} \sin\left(\frac{\pi}{3} + \psi\right) \cos \alpha, -\frac{d}{2} \sin\left(\frac{\pi}{3} + \psi\right) \sin \alpha \right), \\ \mathbf{r}_2 &= \left(\frac{d}{2} \cos\left(\frac{5\pi}{3} + \psi\right), \frac{d}{2} \sin\left(\frac{5\pi}{3} + \psi\right) \cos \alpha, -\frac{d}{2} \left(\sin\left(\frac{\pi}{3} + \psi\right)\right) \sin \alpha \right), \\ \text{and } \mathbf{r}_3 &= \left(\frac{d}{2} \cos(\psi), \frac{d}{2} \sin \psi \cos \alpha, -\frac{d}{2} \sin \psi \sin \alpha \right), \end{aligned}$$

Generalised Configuration considering the tilt Ψ along the x-axis and rotation by the angle α



5. Testing model

- The model here provides number of participant (N_{part}) and Number of binary collisions (N_{coll}) for an event with a given impact parameter (b)

$$N_{coll}(b) = \sigma_{gg} T_{pp}(b) , \quad N_{part}(b) \propto N_{coll}^{3/4}(b) \quad \text{and} \quad \sigma_{gg} = 4.3 \pm 0.6 \text{ mb} \quad [1]$$

- In heavy ion collisions, the concept of number of “ancestors” (independently emitting sources of particles) is parametrised in terms of N_{part} and N_{coll} as follows

$$N_{ancestors} = f N_{part} + (1 - f) N_{coll}$$

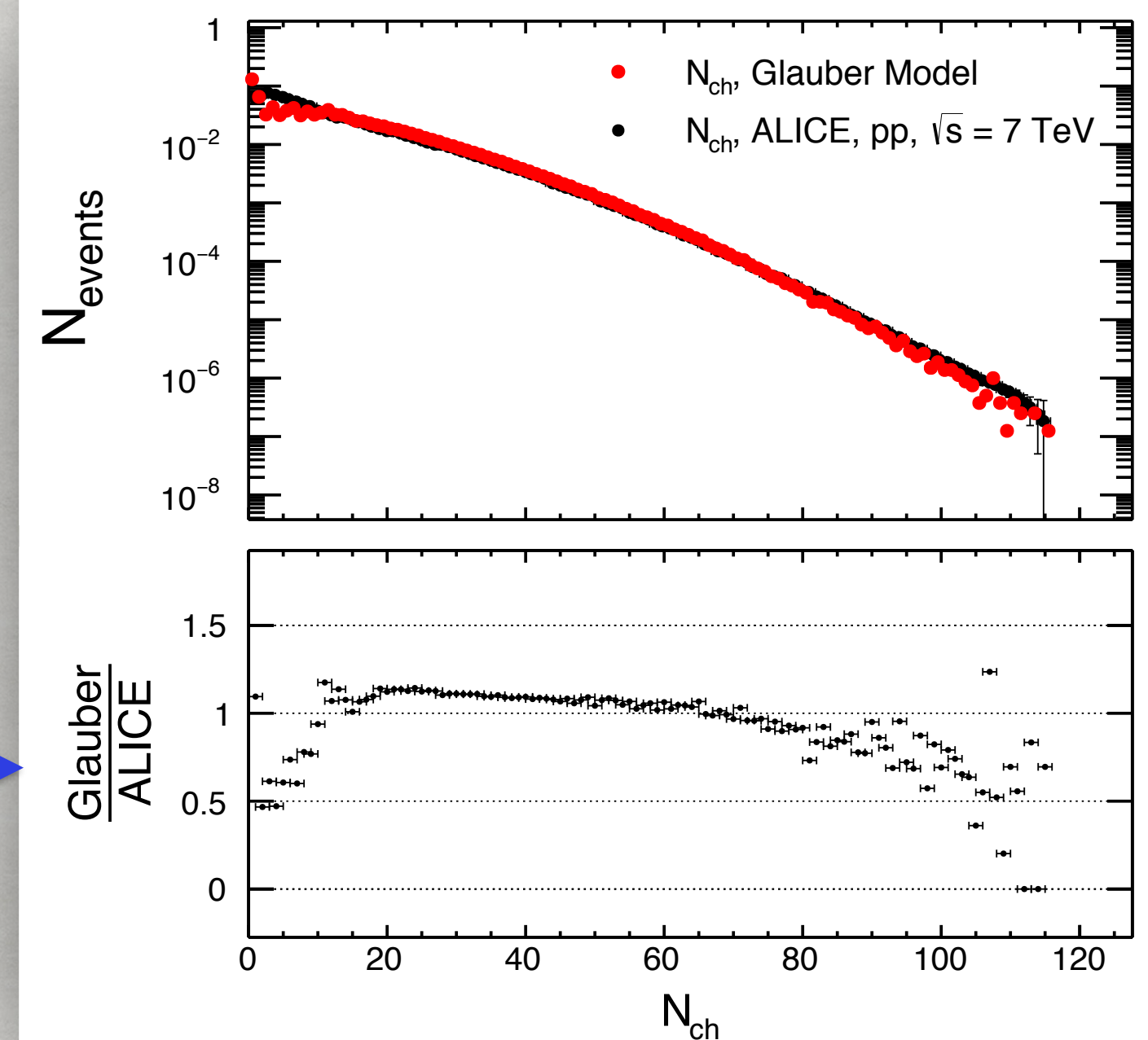
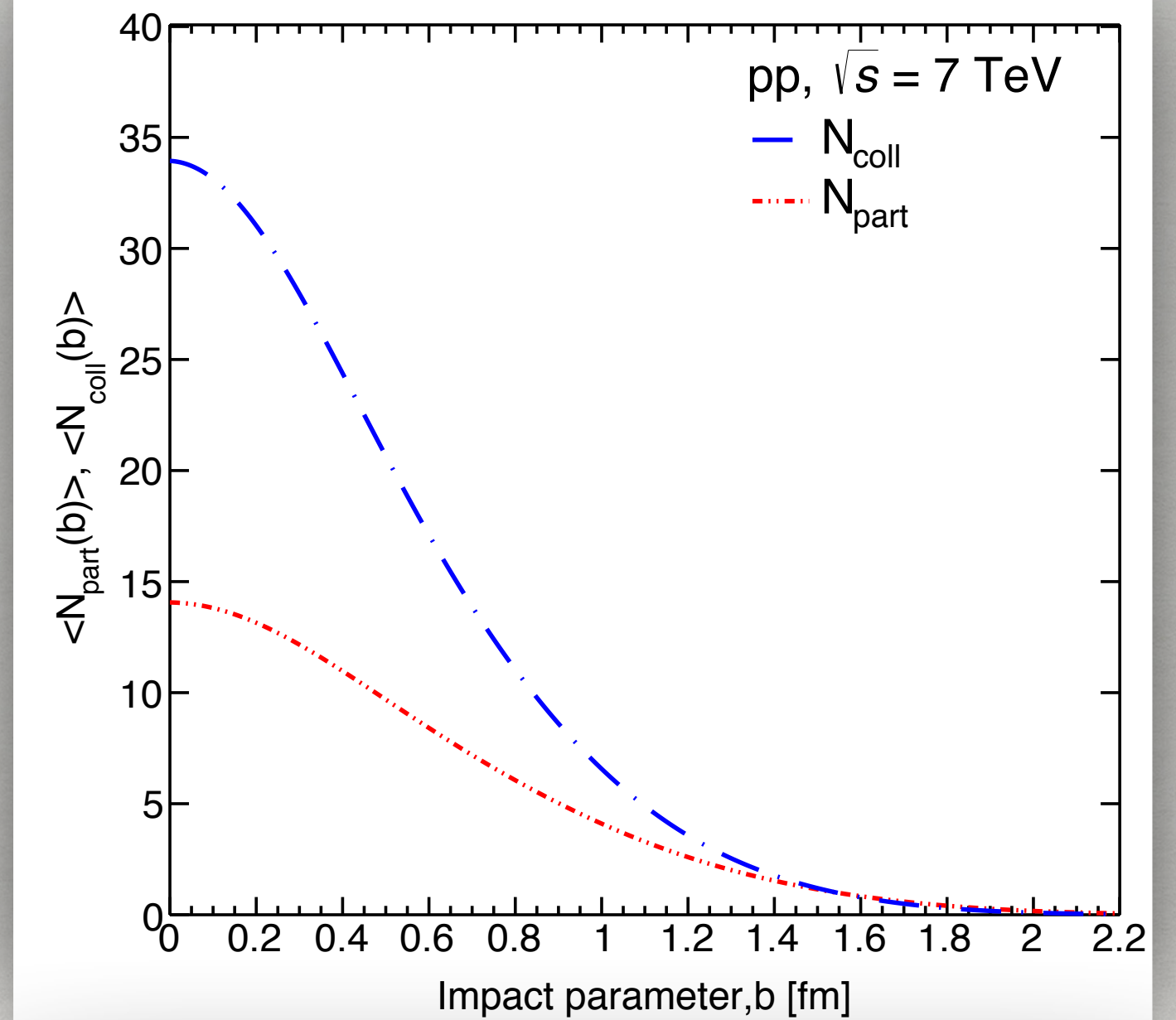
- As the negative binomial distribution (NBD) is able to well reproduce the charged-particle distribution in pp system, we use the two-parameter NBD to calculate the probability of producing n particles per ancestor:

$$P(n; \bar{n}, k) = \frac{\Gamma(n + k)}{\Gamma(k)\Gamma(n + 1)} \left[\frac{\bar{n}}{k + \bar{n}} \right]^n \left[\frac{k}{k + \bar{n}} \right]^k$$

where \bar{n} is the average multiplicity and k characterizes the width of the distribution.

- By using different combinations of f [2], \bar{n} and k [3], the process of obtaining the multiplicity distribution for a large sample of events repeated, until our model simulates the experimental multiplicity distribution.

1. H. J. Drescher et.al., [Phys. Rev. C 76, 024905 \(2007\)](#)
2. C. Loizides, [Phys. Rev. C 94, 024914 \(2016\)](#)
3. D.d'Enterria et.al., [Eur. Phys. J. C 66, 173 \(2010\)](#)

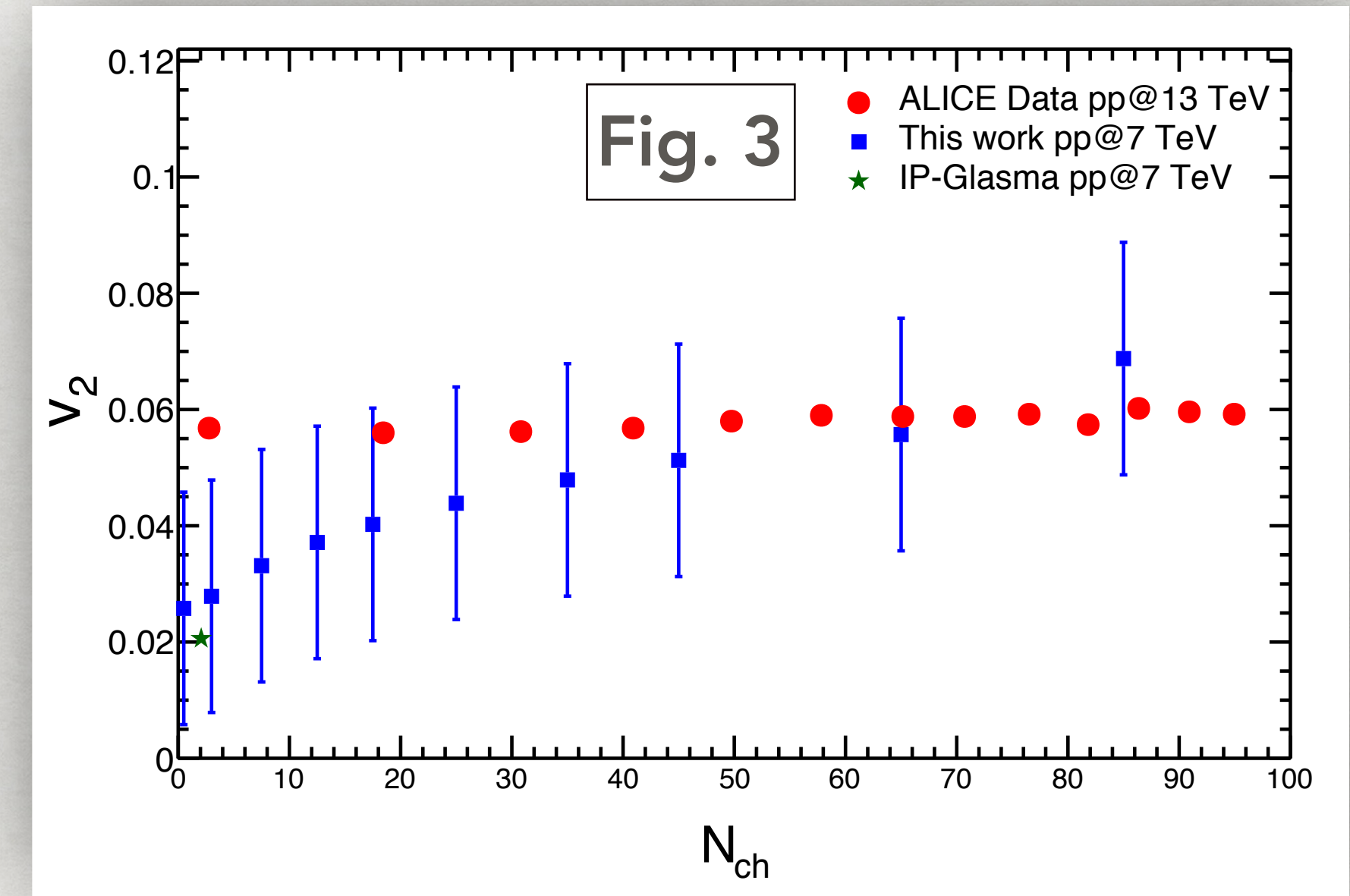
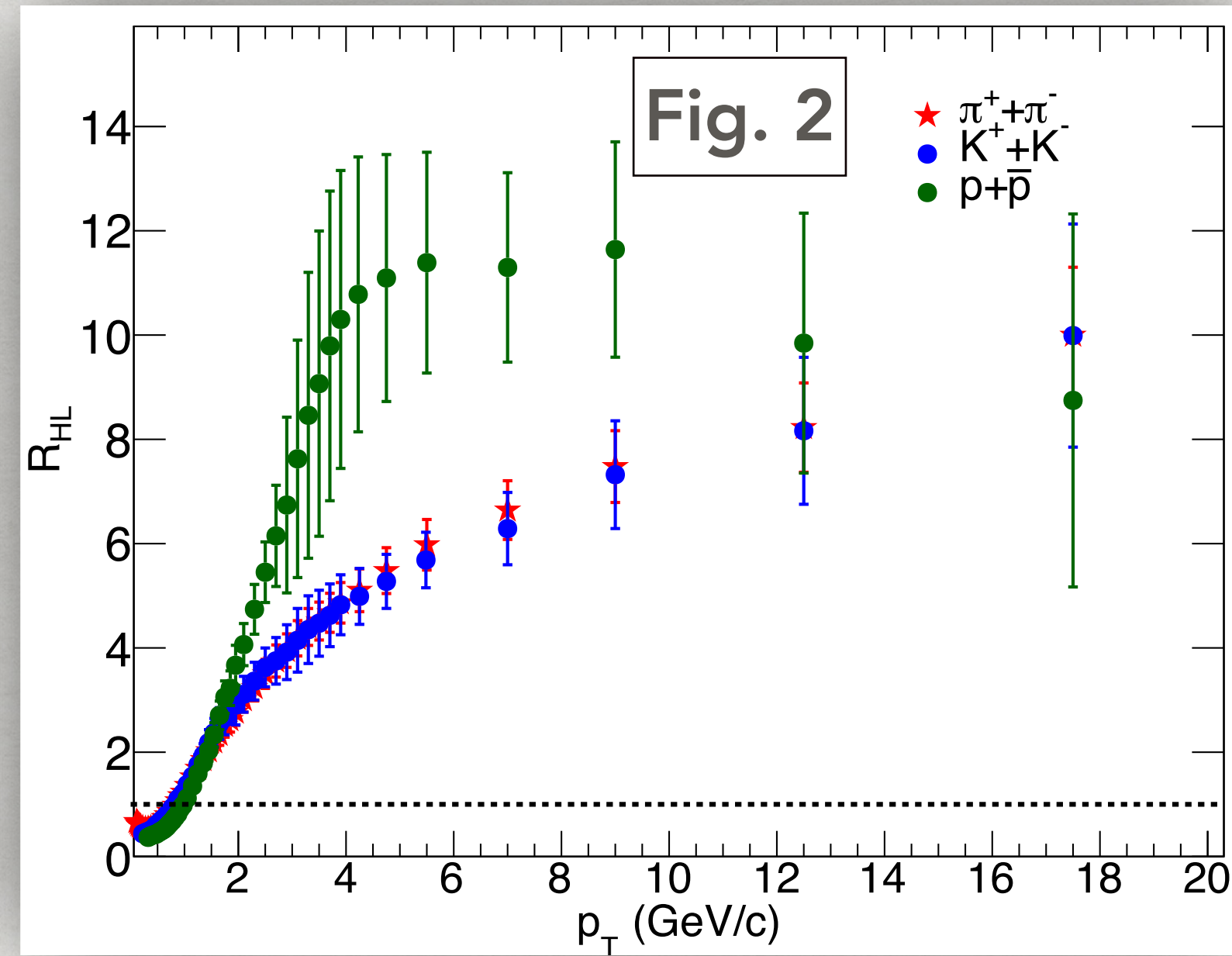
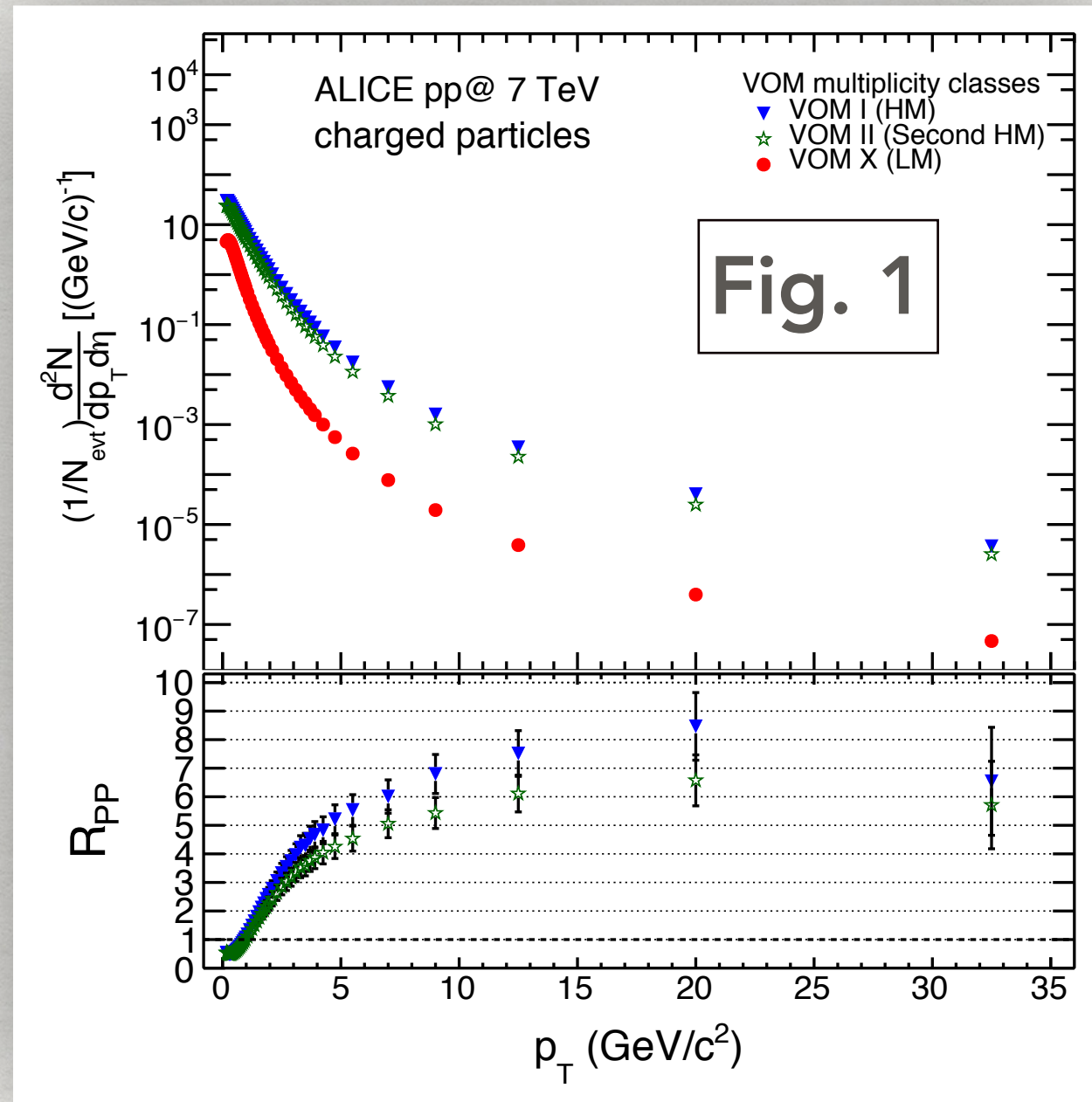


5. Result and discussions

Below table shows geometric properties ($\langle b \rangle$, $\langle N_{ch} \rangle$, $\langle N_{part} \rangle$, $\langle N_{Coll} \rangle$) of pp collisions for different multiplicity (mult) classes using Glauber Monte Carlo calculation along with a Negative binomial distribution fit to charged particle multiplicity distribution at $\sqrt{s} = 7$ TeV for the ALICE experiment at the LHC

Multiplicity [%]	b range [fm]	$\langle dN_{ch}/d\eta \rangle^{\text{Glauber}}$	$\langle dN_{ch}/d\eta \rangle^{\text{expt}}$	$\langle N_{part} \rangle$	$\langle N_{coll} \rangle$
0–1	0–0.25534	19.69	$28.82^{+0.86}_{-0.84}$	13.142	31.156
1–5	0.25535–0.46909	16.24	$20.34^{+0.58}_{-0.57}$	11.164	24.815
5–10	0.46909–0.58484	13.37	$15.80^{+0.34}_{-0.32}$	9.244	19.478
10–15	0.58484–0.66430	11.61	$13.07^{+0.24}_{-0.22}$	8.037	16.153
15–20	0.66431–0.72766	10.28	$11.25^{+0.19}_{-0.18}$	7.131	13.818
20–30	0.72767–0.83026	8.94	$9.21^{+0.15}_{-0.14}$	6.116	11.326
30–40	0.83027–0.91819	7.48	$7.13^{+0.12}_{-0.11}$	5.268	9.215
40–50	0.91820–1.00117	6.49	$5.65^{+0.11}_{-0.09}$	4.418	7.340
50–70	1.00118–1.17163	5.12	$3.81^{+0.07}_{-0.06}$	3.395	5.208
70–100	1.17164–2.54998	3.66	$1.66^{+0.05}_{-0.04}$	1.968	2.591

- We observe the calculated $\langle dN_{ch}/d\eta \rangle$ is consistent with the experimental value expect at high and low mult regions
- This is because of artifact of the model where $\sigma_{gg} = 4.3 \pm 0.6$ mb contains 14% uncertainties which is reflected in the calculation of $\langle dN_{ch}/d\eta \rangle$
- Example: Value of $\langle dN_{ch}/d\eta \rangle$ for minimum bias (0-100%) collisions : Our model = 7.47 and Experimental data = $6.01 \pm 0.01^{+0.20}_{-0.12}$ [5]



- To understand the possibility of a formation of a medium, Nuclear modification factor like observables for charged particle and identified hadrons are plotted in Fig. 1 and Fig 2 respectively using following relation

$$R_{pp}(p_T) = \frac{d^2 N / dp_T d\eta|^{HM}}{d^2 N / dp_T d\eta|^{LM}} \times \frac{\langle N_{coll}^{LM} \rangle}{\langle N_{coll}^{HM} \rangle}$$

- Elliptic flow (v_2) is plotted as a function of charged pseudo rapidity density in Fig. 3 and compared with experimental data and IP- Glasma model. Elliptic flow is obtained here by considering linear scaling as

$$v_2 = \Omega \epsilon \quad \text{where, } \Omega = 0.3 \pm 0.02 \text{ [4]}$$