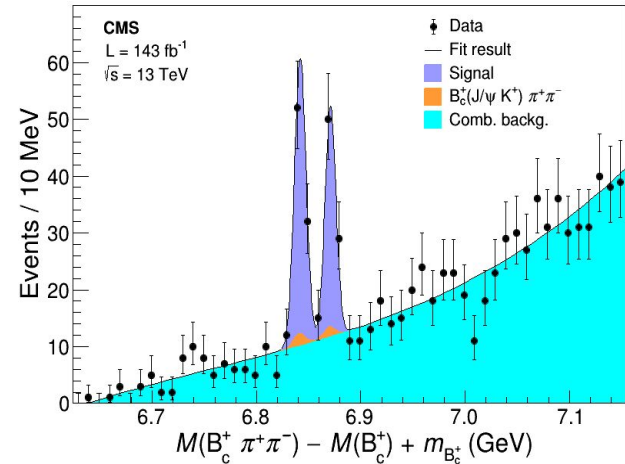


Measurement of $B_c(2S)^+$ and $B_c^*(2S)^+$ cross section ratios in proton-proton collisions

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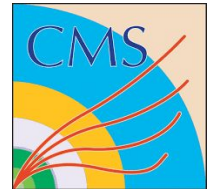


Grupo de Fenomenología e Interacciones Fundamentales (GFIF)

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Introduction

Particle	Predicted M(MeV)
B_c	6247-6286
B_c^*	6308-6341
$B_c(2S)$	6835-6882
$B_c(2S)^*$	6881-6914

$B_c(2S)^* \rightarrow B_c^* \pi^+ \pi^-$ followed by $B_c^* \rightarrow B_c \gamma_{lost}$

Since the photon is not detected, we end up seeing

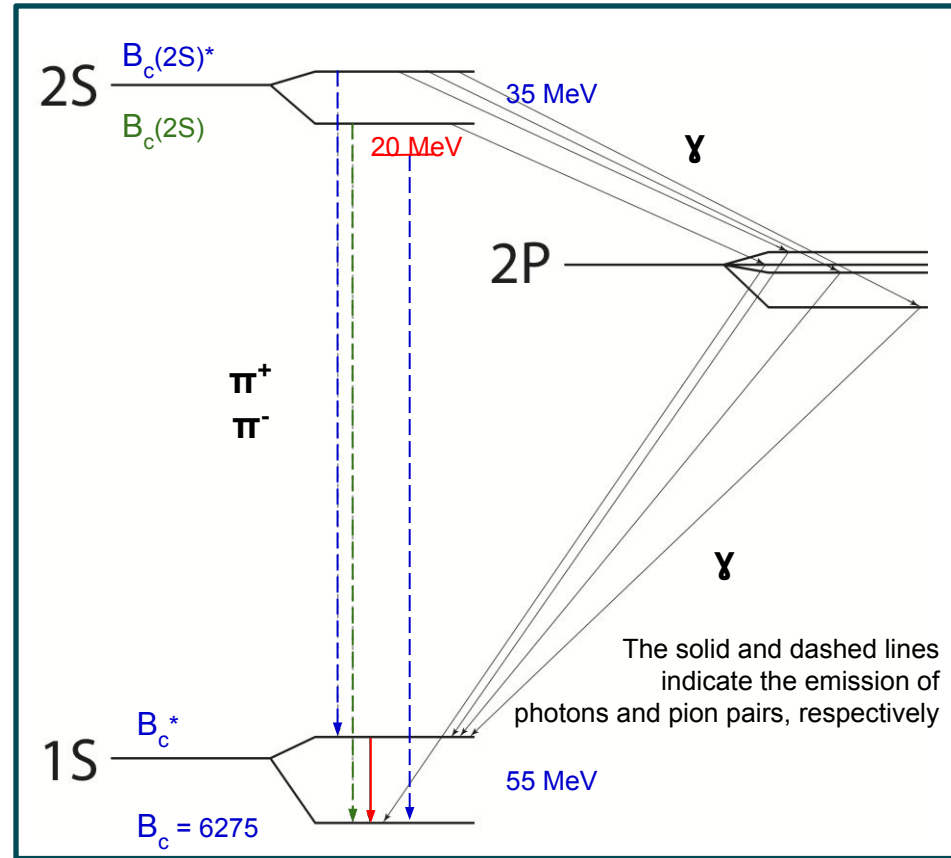
$B_c(2S)^* \rightarrow B_c \pi^+ \pi^-$ plus “missing energy”

Same final state as

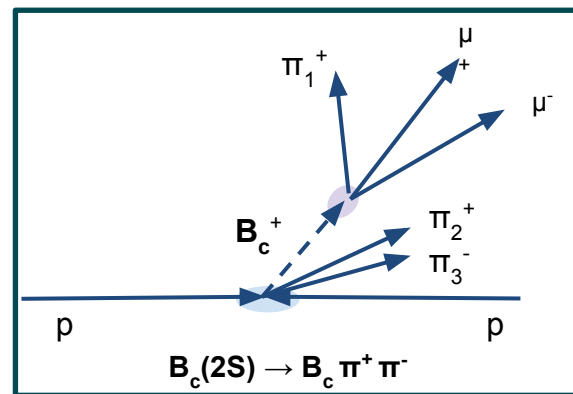
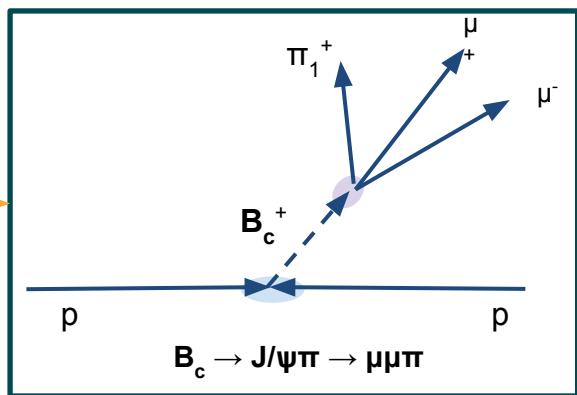
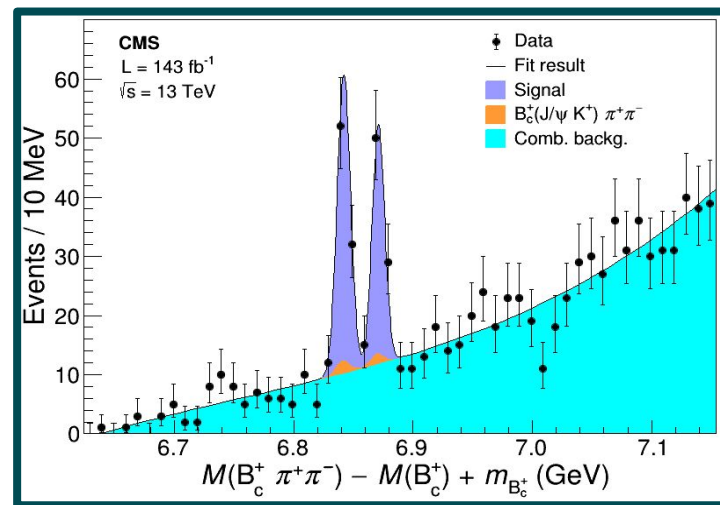
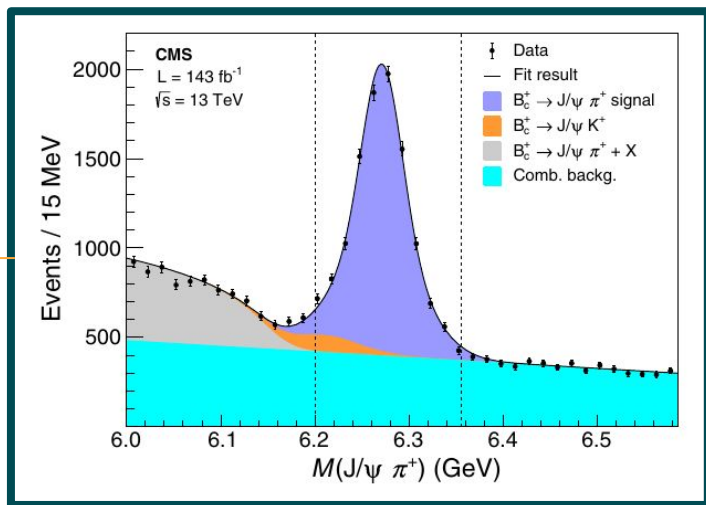
$B_c(2S) \rightarrow B_c \pi^+ \pi^-$

So, we see a two-peak structure in the $B_c \pi^+ \pi^-$ mass distribution, with the $B_c(2S)^*$ peak at a mass shifted by

$$\Delta M = [M(B_c^*) - M(B_c)] - [M(B_c(2S)^*) - M(B_c(2S))]$$



Reconstruction of the mass of B_c



Cross section ratios as functions of p_T and $|y|$

Efficiency

$$\epsilon_{B_c^*(2S)^+} = \frac{N_{B_c^*(2S)^+ \rightarrow B_c^{*+} \pi^+ \pi^-}^{rec}}{N_{B_c^*(2S)^+ \rightarrow B_c^{*+} \pi^+ \pi^-}^{gen}}$$

$$\epsilon_{B_c(2S)^+} = \frac{N_{B_c(2S)^+ \rightarrow B_c^+ \pi^+ \pi^-}^{rec}}{N_{B_c(2S)^+ \rightarrow B_c^+ \pi^+ \pi^-}^{gen}}$$

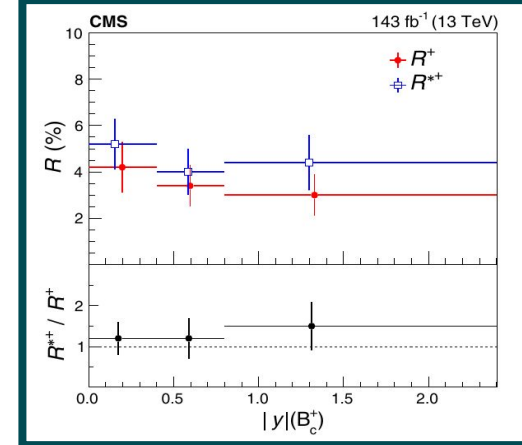
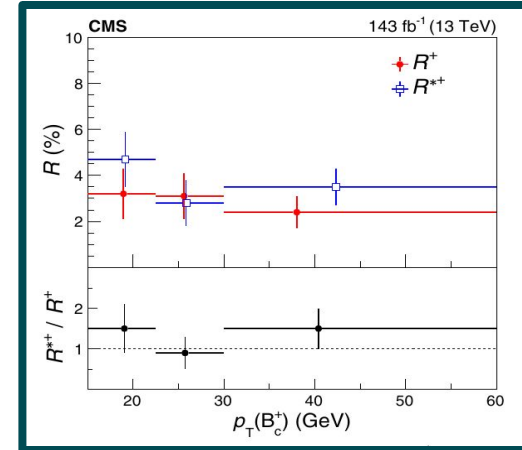
where $N_{B_c(2S)^+ \rightarrow B_c^+ \pi^+ \pi^-}^{rec}$ is the number of reconstructed $B_c(2S)^+ \rightarrow B_c^+ \pi^+ \pi^-$ events after the full selection and $N_{B_c(2S)^+ \rightarrow B_c^+ \pi^+ \pi^-}^{gen}$ is the number of generated $B_c(2S)^+ \rightarrow B_c^+ \pi^+ \pi^-$ decays in the fiducial region of the analysis specified by the B_c kinematic window $p_T(B_c) > 15$ GeV and $|y(B_c)| < 2.4$. Analogous efficiency definitions apply for B_c and $B_c(2S)$

Cross section ratios

$$R^+ \equiv \frac{\sigma(B_c(2S)^+)}{\sigma(B_c^+)} \mathcal{B}(B_c(2S)^+ \rightarrow B_c^+ \pi^+ \pi^-) = \frac{N(B_c(2S)^+)}{N(B_c^+)} \frac{\epsilon(B_c^+)}{\epsilon(B_c(2S)^+)},$$

$$R^{*+} \equiv \frac{\sigma(B_c^*(2S)^+)}{\sigma(B_c^+)} \mathcal{B}(B_c^*(2S)^+ \rightarrow B_c^{*+} \pi^+ \pi^-) = \frac{N(B_c^*(2S)^+)}{N(B_c^+)} \frac{\epsilon(B_c^+)}{\epsilon(B_c^*(2S)^+)},$$

$$R^{*+}/R^+ = \frac{\sigma(B_c^*(2S)^+) \mathcal{B}(B_c^*(2S)^+ \rightarrow B_c^{*+} \pi^+ \pi^-)}{\sigma(B_c(2S)^+) \mathcal{B}(B_c(2S)^+ \rightarrow B_c^+ \pi^+ \pi^-)} = \frac{N(B_c^*(2S)^+) \epsilon(B_c(2S)^+)}{N(B_c(2S)^+) \epsilon(B_c^*(2S)^+)}.$$



Summary

We have look at 2015, 2016, 2017 and 2018 data (143 fb⁻¹), looking for the excited states $B_c(2S)^{*\pm}$

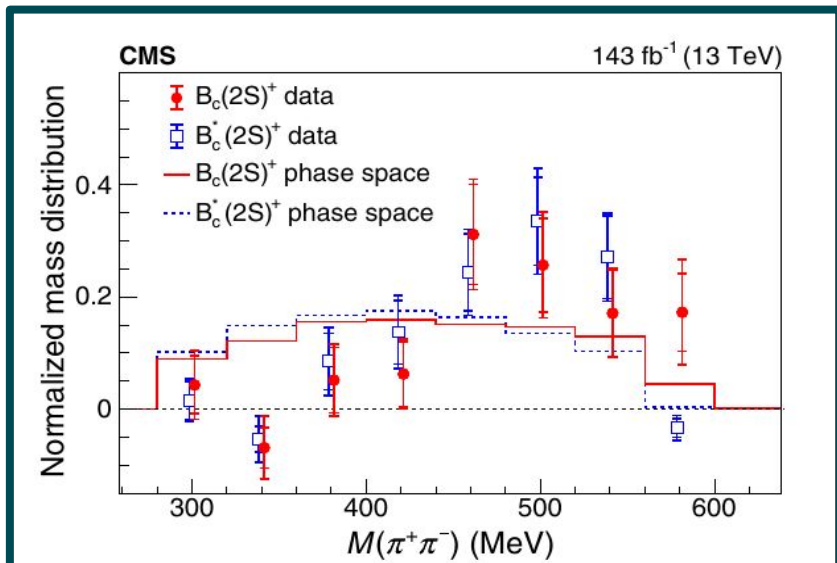
$B_c(2S)^\pm \rightarrow B_c^\pm \pi^+ \pi^-$ followed by $B_c^\pm \rightarrow J/\psi \pi^\pm$ and $J/\psi \rightarrow \mu^+ \mu^-$

$B_c(2S)^{*\pm} \rightarrow B_c^{*\pm} \pi^+ \pi^- \rightarrow (B_c^\pm \gamma) \pi^+ \pi^-$ followed by $B_c^\pm \rightarrow J/\psi \pi^\pm$ and $J/\psi \rightarrow \mu^+ \mu^-$

$$R(B_c(2S)^+) = 0.034 \pm 0.007 \text{ (stat)} \pm 0.006 \text{ (syst)} \quad \text{for } p_T(B_c) > 15 \text{ GeV and } |y(B_c)| < 2.4$$

$$R(B_c(2S)^{*+}) = 0.047 \pm 0.007 \text{ (stat)} \pm 0.004 \text{ (syst)} \quad \text{for } p_T(B_c) > 15 \text{ GeV and } |y(B_c)| < 2.4$$

$$R(B_c(2S)^{*+}/B_c(2S)^+) = 1.394 \pm 0.349 \text{ (stat)} \pm 0.132 \text{ (syst)} \quad \text{for } p_T(B_c) > 15 \text{ GeV and } |y(B_c)| < 2.4$$



Systematic uncertainties

	R^+	R^{*+}	R^{*+}/R^+
$J/\psi \pi^+$ fit model	5.5	5.5	...
$B_c^+ \pi^+ \pi^-$ fit model	5.9	2.9	2.9
Efficiencies: statistical uncertainty	1.1	1.0	1.4
Efficiencies: spread among years	1.8	1.6	0.9
Efficiencies: pion tracking	4.2	4.2	...
Decay kinematics	1.5	6.9	4.2
Helicity angle	1.0	6.0	3.5
Total	9.5	12.0	6.4