NLO+PS matching for loopinduced processes in SHERPA

Presented by Simon Luca Villani

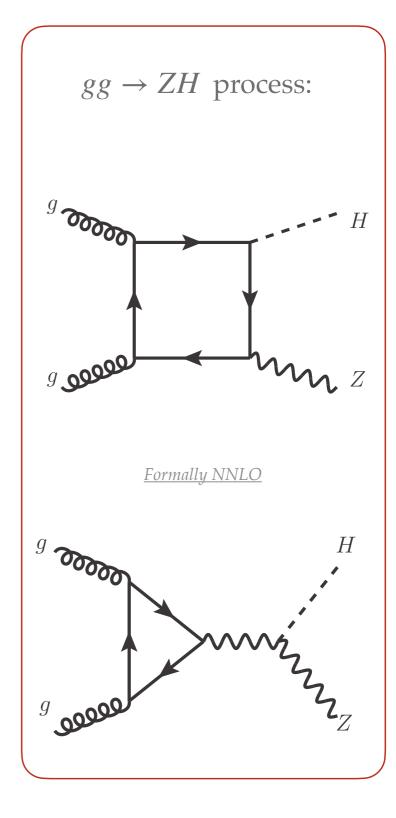




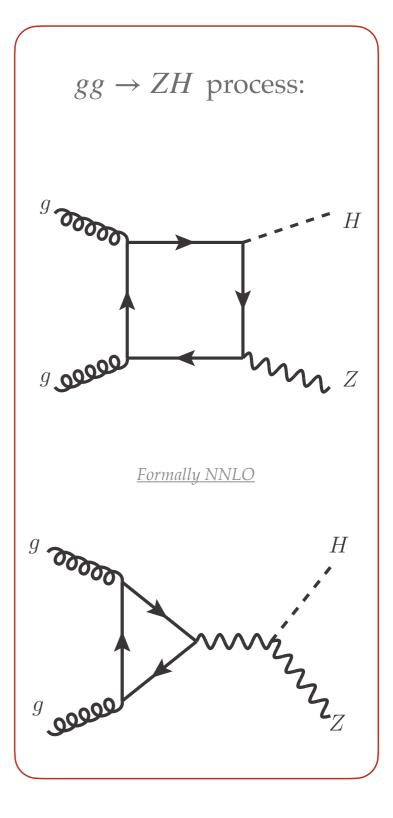


30.04.2020 Internet MCnet Meeting

• Loop-induced processes feature first non-trivial contribution at loop-level only



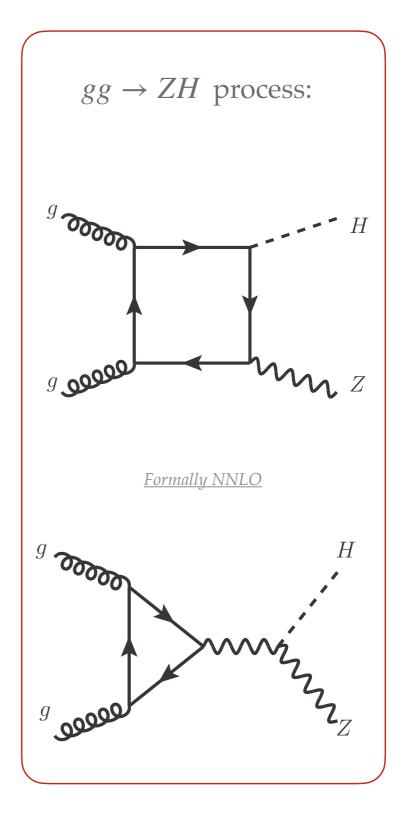
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Peculiarities:

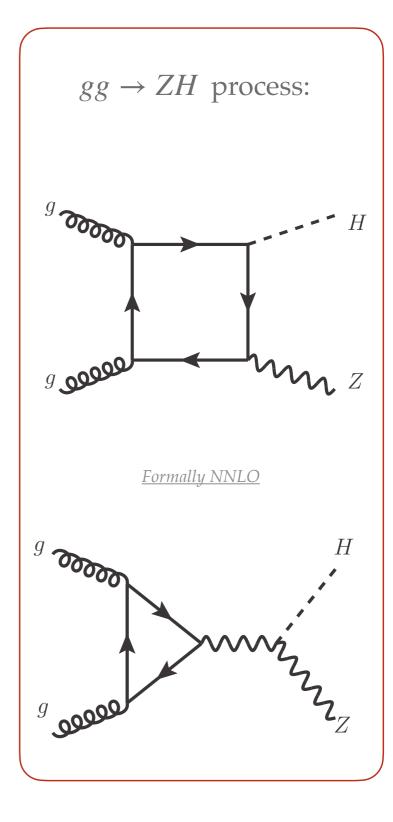
- Ren/fac uncertainty rather big
- NLO very challenging due to the presence of massive multi-scale double box integrals
- Sensitive to theory parameters variation



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Peculiarities:

- Ren/fac uncertainty rather big \rightarrow ~30% in ggHZ [de Florian, D. et al.: 1610.07922]
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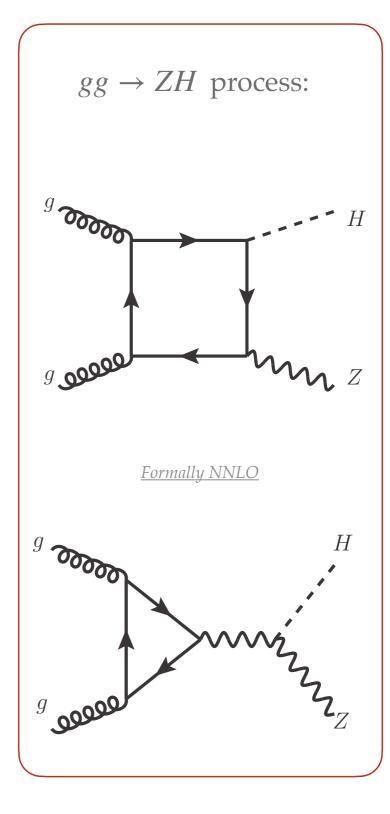


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Particularly suitable for BSM studies and SM precision studies



$gg \rightarrow ZZ$ case:

I	Grazzini	M. et	al:	HEP030	(2019)	0701
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\sqrt{s}	8 TeV	$13\mathrm{TeV}$	8 TeV	13 TeV
	σ [1	fb]	$\sigma/\sigma_{ m NJ}$	LO-1
LO	$8.1881(8)^{+2.4\%}_{-3.2\%}$	$13.933(1)_{-6.4\%}^{+5.5\%}$	-27.5%	-29.8%
NLO	$11.2958(4)_{-2.0\%}^{+2.5\%}$	$19.8454(7)_{-2.1\%}^{+2.5\%}$	0%	0%
$qar{q}$ NNLO	$12.09(2)_{-1.1\%}^{+1.1\%}$	$21.54(2)_{-1.2\%}^{+1.1\%}$	+7.0%	+8.6%
	σ [1	$\sigma/\sigma_{ m ggLO}-1$		
ggLO	$0.79355(6)_{-20.9\%}^{+28.2\%}$	$2.0052(1)_{-17.9\%}^{+23.5\%}$	0%	0%
$ggNLO_{gg}$	$1.4787(4)_{-13.1\%}^{+15.9\%}$	$3.626(1)_{-12.7\%}^{+15.2\%}$	+86.3%	+80.8%
ggNLO	$1.3892(4)_{-13.6\%}^{+15.4\%}$	$3.425(1)_{-12.0\%}^{+13.9\%}$	+75.1%	+70.8%
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$gg \rightarrow HZ$ case:

[Astill, W. et al.: 1804.08141]

Fiducial cross section	HZJ-MiNLO	MCFM-8.0	HZ-NNLOPS(LHEF)	HZNNLOPS
no $gg \rightarrow HZ$	$6.59^{+7.2\%}_{-6.2\%}$ fb	$7.14_{-0.9\%}^{+0.5\%}$ fb	$7.14^{+0.3\%}_{-0.4\%}$ fb	$6.49^{+0.8\%}_{-0.6\%}$ fb
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Ī	Grazzini,	M.	et al:	THEP03	(2019)	0701

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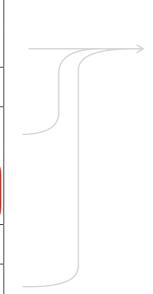
[Astill, W. et al.: 1804.08141]

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[Grazzini,	M.	et al:	IHEP03	2019	070

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	8TeV	13 TeV
ggLO =	~ 6 %	~ 8.5 %

- It even doubles the total uncertainty
- ggLO largely underestimates the NLO

$gg \rightarrow HZ$ case:

[Astill, W. et al.: 1804.08141]

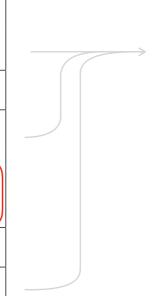
NNLO

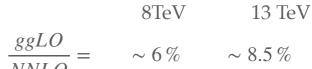
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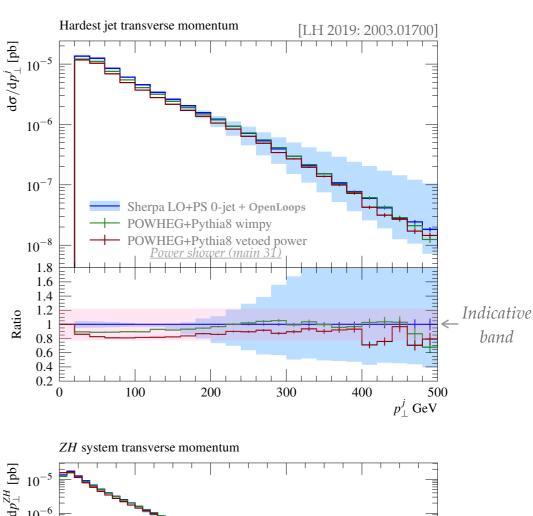


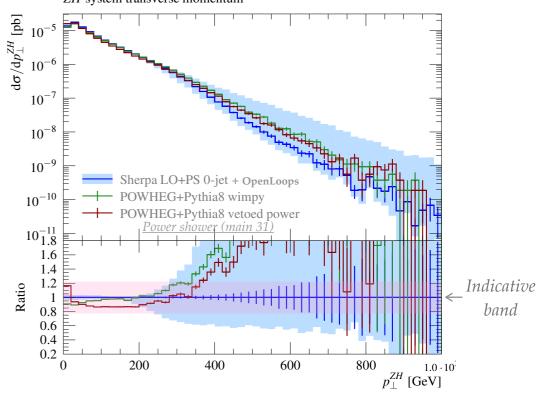
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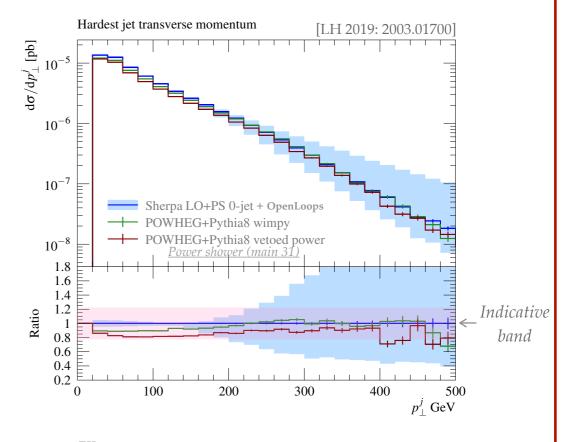
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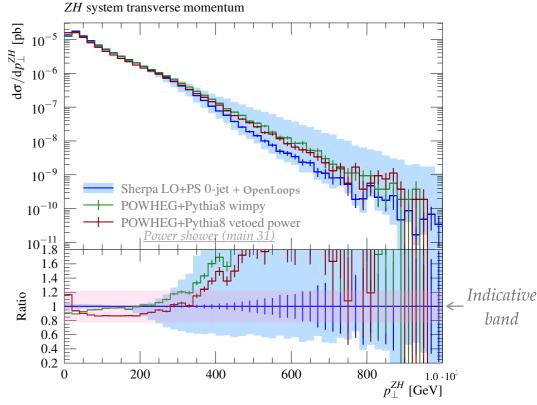




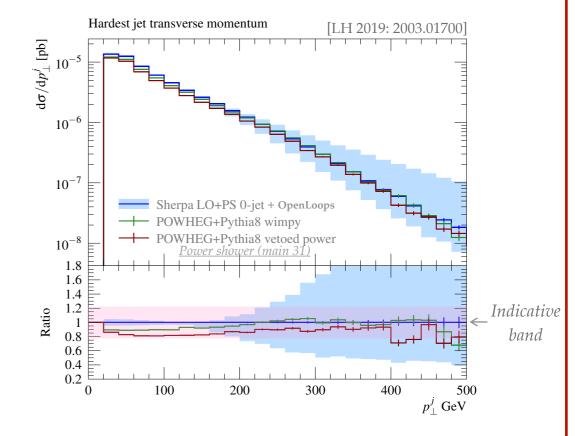
Again $gg \rightarrow HZ$ case:

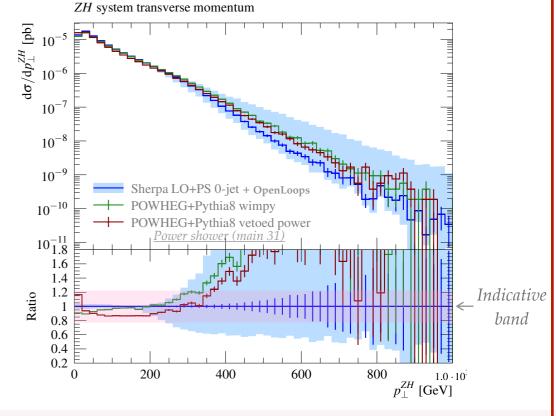
• SHERPA shower starting scale has been varied by a factor $\sqrt{2}$ around its central value of m_{ZH}





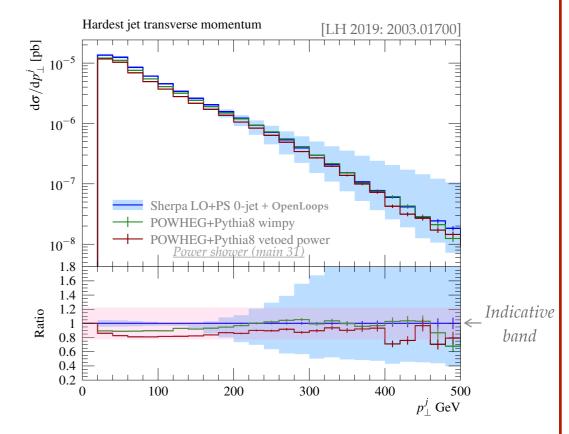
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- For $p_{\perp} > 250$ **GeV** parton shower bands largely exceed the (μ_R, μ_F) variation

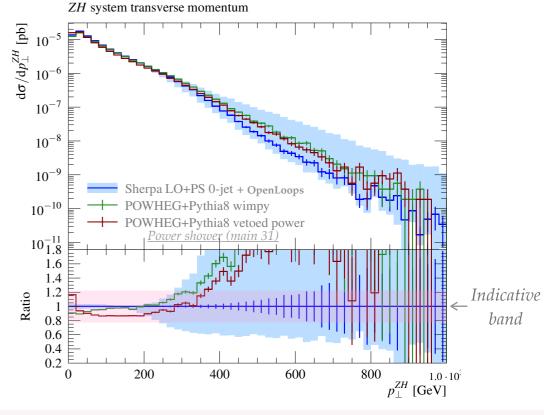




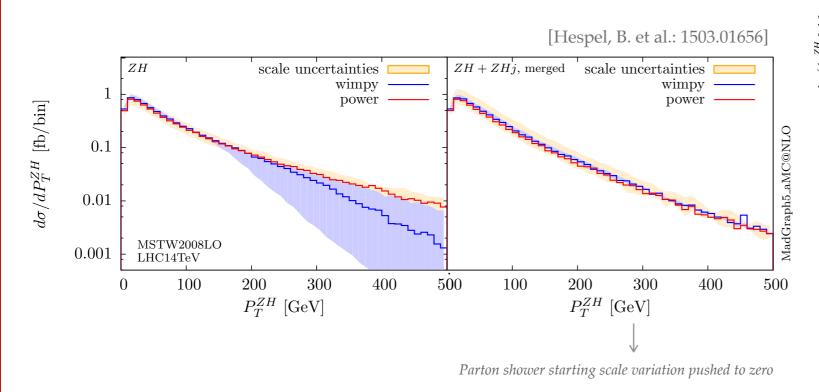


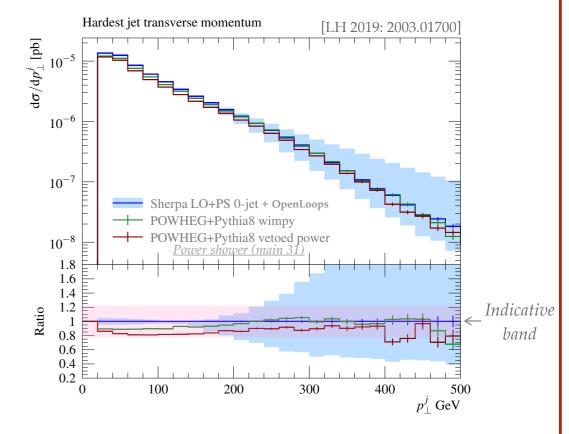
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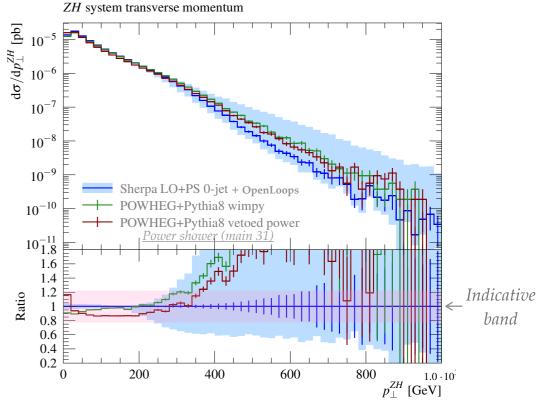




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In Sherpa 3.0.0 $gg \rightarrow HH$ is available at full Standard Model only for on-shell final states and can thus be used as test case

Run card configuration example:

```
ME_GENERATORS:
- External
- OpenLoops

BEAMS: 2212
BEAM_ENERGIES: 6500

PROCESSES:
- "21 21 -> 25 25":
    Order: {QCD: 2, EW: 2}
    NLO_Mode: MC@NLO
    NLO_Order: {QCD: 1, EW: 0}
    Loop_Generator: DiHiggsNLO
    Integrator: Rambo
```

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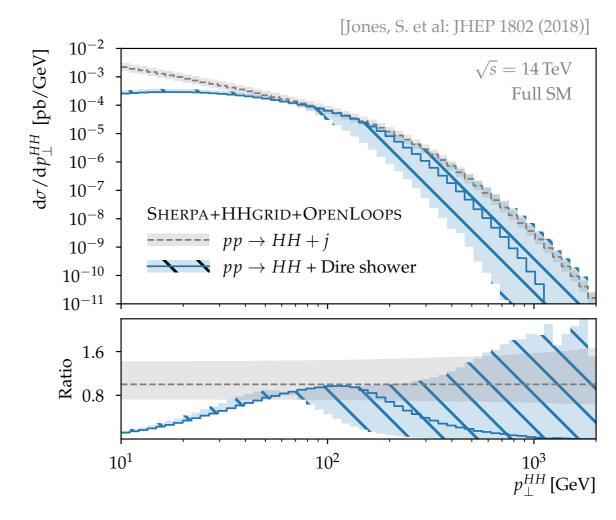
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                                                    From interpolation
     Loop_Generator: DiHiggsNLO
     Integrator: Rambo
                                           [Heinrich, G. et al: 1703.09252]
       {
m d} \sigma / {
m d} m_{
m hh} \, {
m [pb/GeV]}^{-4}
                                                        POWHEG
                                                       Original —
                        Full SM NLO
                        LHC 14 \text{ TeV}
                        PDF4LHC15 NLO
                        \mu = m_{\rm hh}/2
            1.2
             1.1
             1.0
            0.9
             0.8
                                 500
                                               700
                                                      800
                                                             900
                          400
                                        600
                                                                    1000
                                      m_{\rm hh} \; [{\rm GeV}]
```

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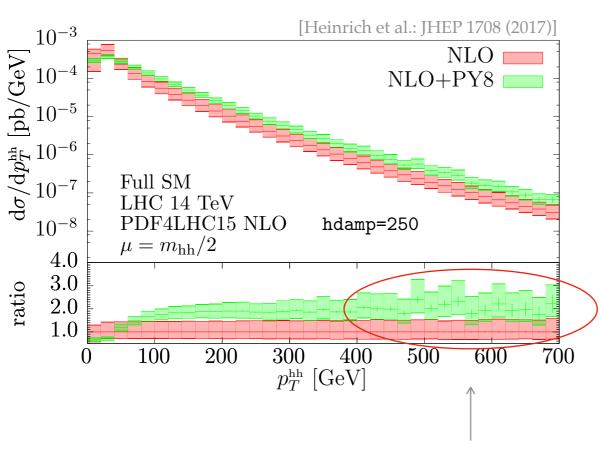
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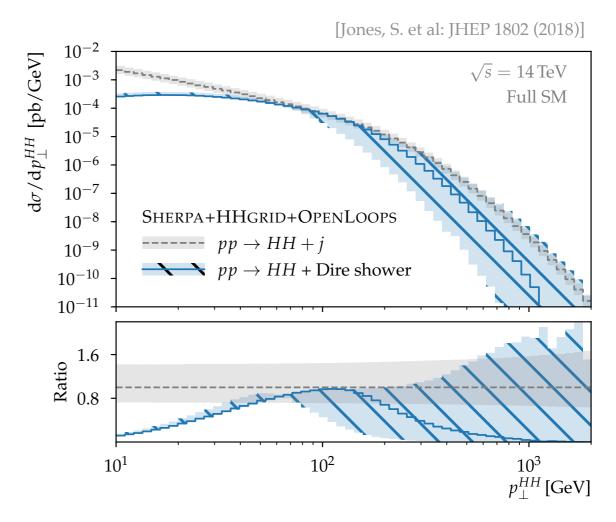


LO suffers by the same large PS starting scale uncertainty shown before for other processes

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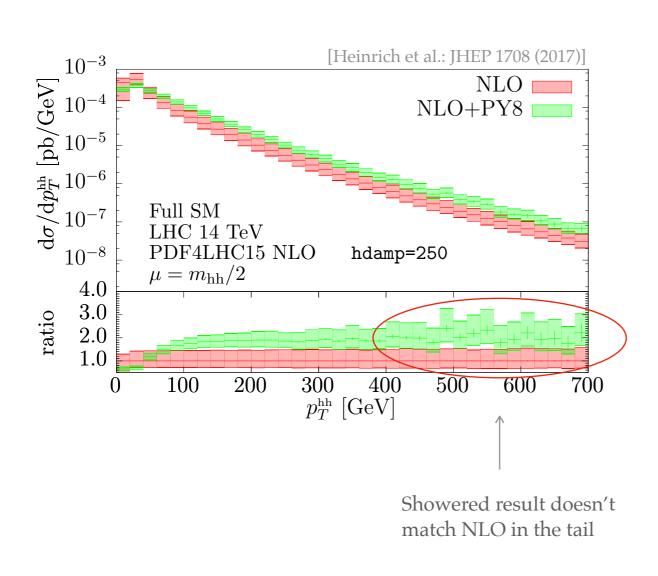


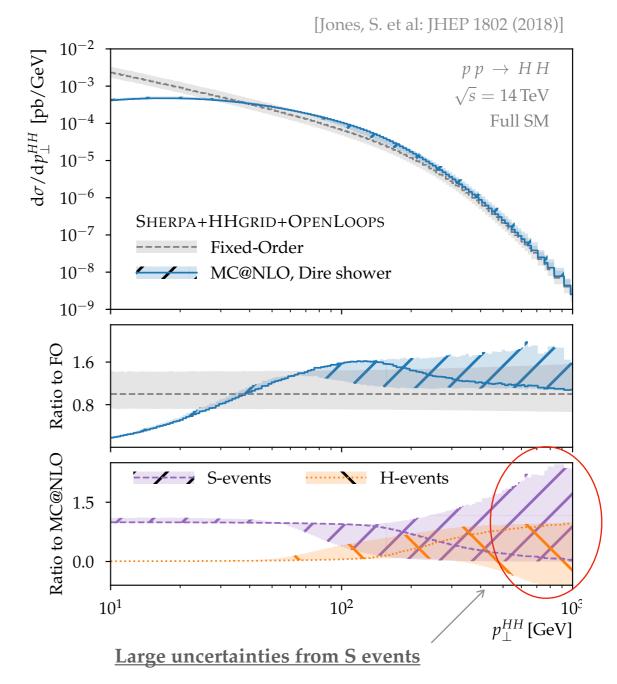
Showered result doesn't match NLO in the tail



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What causes this mismatch?

[Jones, S. et al: JHEP 1802 (2018)]

MC@NLO general expression:

$$\langle O \rangle = \int \! d\phi_B \left(B(\phi_B) + V(\phi_B) + I(\phi_B) \right) O(\phi_B) \times$$

$$\times \left[\Delta(t_0, \mu_{PS}^2) + \int \! d\phi_1 \Delta(t, \mu_{PS}^2) \frac{D(\phi_B, \phi_1)}{B(\phi_B)} \Theta(\mu_{PS}^2 - t) \Theta(t - t_0) \right] + \qquad \text{S events}$$

$$+ \int \! d\phi_R H(\phi_R) O(\phi_R) \qquad \qquad \text{H events}$$

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$$+ \int d\phi_R H(\phi_R) O(\phi_R)$$

$$\text{H events}$$

For observables insensitive to Born kinematical configuration

$$\begin{split} \langle O \rangle &= \int \! d\phi_B \left(B(\phi_B) + V(\phi_B) + I(\phi_B) \right) O(\phi_B) \times \\ &\times \left[\Delta(t_0, \mu_{PS}^2) + \int \! d\phi_1 \Delta(t, \mu_{PS}^2) \frac{D(\phi_B, \phi_1)}{B(\phi_B)} \Theta(\mu_{PS}^2 - t) \Theta(t - t_0) \right] + \\ &+ \int \! d\phi_R H(\phi_R) O(\phi_R) \end{split}$$

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[Jones, S. et al: JHEP 1802 (2018)]

MC@NLO general expression:

$$\begin{split} \langle O \rangle &= \int \! d\phi_B \left(B(\phi_B) + V(\phi_B) + I(\phi_B) \right) O(\phi_B) \, \times \\ &\times \left[\Delta(t_0, \mu_{PS}^2) + \int \! d\phi_1 \Delta(t, \mu_{PS}^2) \frac{D(\phi_B, \phi_1)}{B(\phi_B)} \Theta(\mu_{PS}^2 - t) \Theta(t - t_0) \right] + \qquad \text{S events} \\ &+ \int \! d\phi_R H(\phi_R) O(\phi_R) \end{split}$$

For observables insensitive to Born kinematical configuration and focussing on the high energy tail

$$\langle O \rangle = \int d\phi_B \left(B(\phi_B) + V(\phi_B) + I(\phi_B) \right) O(\phi_B) \times$$

$$\times \left[\Delta(t_0, \mu_{PS}^2) + \int d\phi_1 \Delta(t, \mu_{PS}^2) \frac{D(\phi_B, \phi_1)}{B(\phi_B)} \Theta(\mu_{PS}^2 - t) \Theta(t - t_0) \right] +$$

$$+ \int d\phi_R H(\phi_R) O(\phi_R) \qquad \text{It goes to 1}$$

What causes this mismatch?

[Jones, S. et al: JHEP 1802 (2018)]

MC@NLO general expression:

$$\langle O \rangle = \int \! d\phi_B \left(B(\phi_B) + V(\phi_B) + I(\phi_B) \right) O(\phi_B) \times$$

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$$+ \int \! d\phi_R H(\phi_R) O(\phi_R) \qquad \qquad \text{H events}$$

For observables insensitive to Born kinematical configuration and focussing on the high energy tail

$$\langle O \rangle = \int d\phi_B \left(B(\phi_B) + V(\phi_B) + I(\phi_B) \right) O(\phi_B) \times$$

$$\times \int d\phi_1 \frac{D(\phi_B, \phi_1)}{B(\phi_B)} \Theta(\mu_{PS}^2 - t) +$$

$$+ \int d\phi_R H(\phi_R) O(\phi_R)$$

$$\langle O \rangle = \int d\phi_B \left(V(\phi_B) + I(\phi_B) \right) O(\phi_B) \times$$

$$\times \int d\phi_1 \frac{D(\phi_B, \phi_1)}{B(\phi_B)} \Theta(\mu_{PS}^2 - t) +$$

$$+ \int d\phi_R R(\phi_R) O(\phi_R)$$

What causes this mismatch?

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MC@NLO general expression:

$$\begin{split} \langle O \rangle &= \int \! d\phi_B \left(B(\phi_B) + V(\phi_B) + I(\phi_B) \right) O(\phi_B) \times \\ &\times \left[\Delta(t_0, \mu_{PS}^2) + \int \! d\phi_1 \Delta(t, \mu_{PS}^2) \frac{D(\phi_B, \phi_1)}{B(\phi_B)} \Theta(\mu_{PS}^2 - t) \Theta(t - t_0) \right] + \qquad \text{S events} \\ &+ \left[d\phi_R H(\phi_R) O(\phi_R) \right] \end{split}$$

For observables insensitive to Born kinematical configuration and focussing on the high energy tail

$$\langle O \rangle = \int d\phi_B \left(V(\phi_B) + I(\phi_B) \right) O(\phi_B) \times$$

$$\times \int d\phi_1 \frac{D(\phi_B, \phi_1)}{B(\phi_B)} \Theta(\mu_{PS}^2 - t) +$$

$$+ \int d\phi_R R(\phi_R) O(\phi_R)$$

To recover the real emission result the first term in the r.h.s. must be negligible. This requirement is spoiled if the following conditions are met:

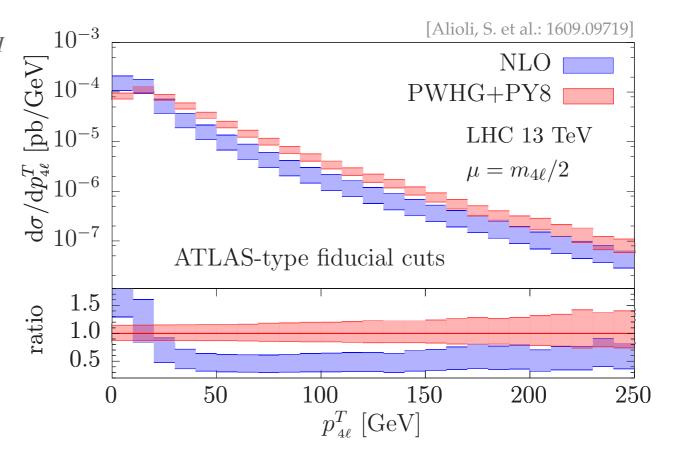
- Large **K factor**
- Non-negligible <u>splitting function</u> in that energy region
- Energy <u>region accessible</u> to the parton shower

What comes next?

• Study parton shower matching uncertainty for other processes, e.g. $gg \rightarrow VV$:

[von Manteuffel, A. et al.: 1503.08835] Using <u>VVamp</u> c++ code

 $gg \rightarrow ZZ$ shows a similar discrepancy in the tail as $gg \rightarrow HH$



What comes next?

• Study parton shower matching uncertainty for other processes, e.g. $gg \rightarrow VV$:

• Including Top quark effect in the loop using high and low energy approximation [Davies, J. et al.: 2002.05558]

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• Resummation effects and relative uncertainties using dedicated Sherpa module

Conclusions

- Moving towards LHC@HL makes the good modeling of these processes an important step for future high precision studies and BSM analyses
- These processes suffer from theoretical uncertainties more than others. In particular for what concern the parton shower matching uncertainties
- A more detailed study of these uncertainties is needed to have a reliable MC@NLO and solutions to improve the showered sample are required

Thanks for the attention