



# Technology of superfluid helium

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Helium Week, WUT & CERN  
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# Contents



- Introduction to superfluid helium
- Superfluid helium as a technical coolant
- Practical cooling schemes
- Refrigeration below 2 K
- Specific technology for He II systems
- Conclusion and outlook

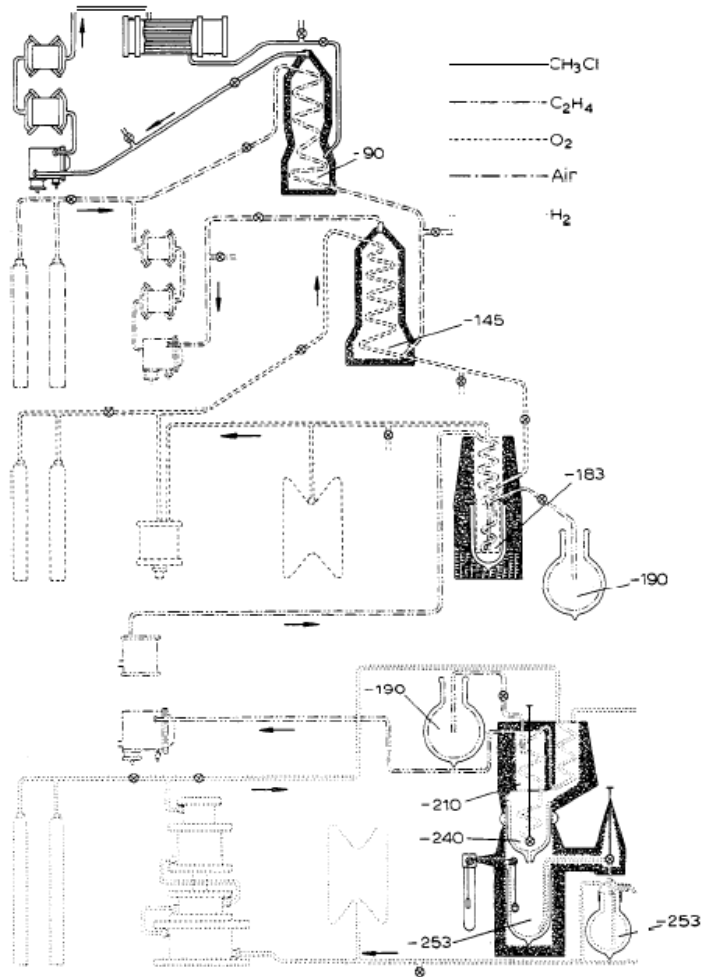


# Contents

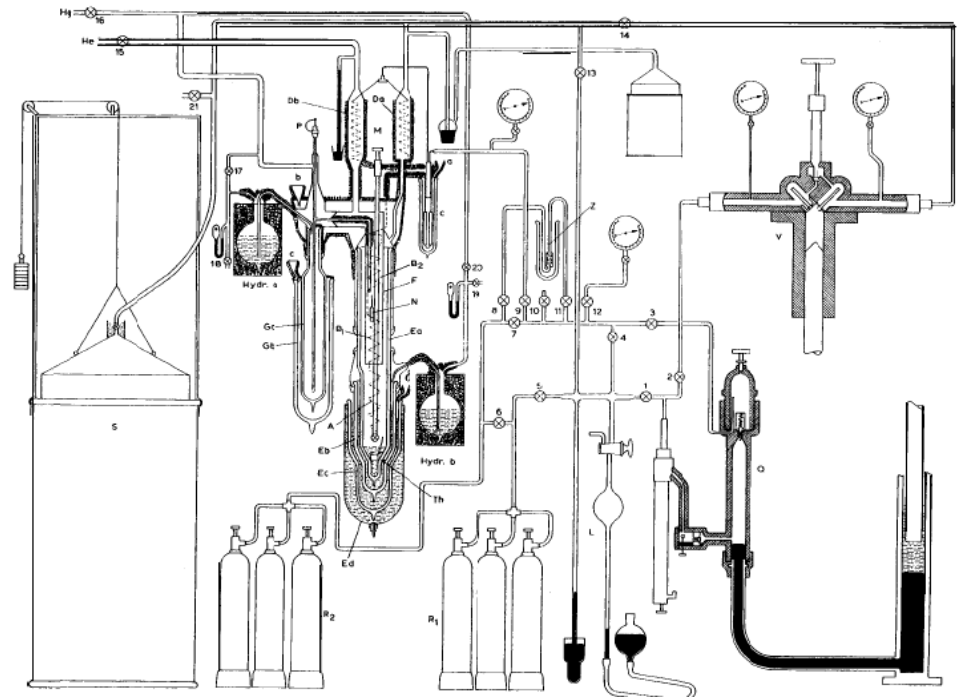


- **Introduction to superfluid helium**
  - Superfluid helium as a technical coolant
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# First liquefaction of helium (1908)



Leiden « cascade » to produce liquid hydrogen



Helium liquefaction stage



HEIKE KAMERLINGH ONNES

Investigations into the properties of substances at low temperatures, which have led, amongst other things, to the preparation of liquid helium

*Nobel Lecture, December 11, 1913*



# Unsuccessful attempt to solidify helium

Naturally the question arose as to whether helium can also be converted into the solid state. An experiment aimed at lowering the temperature of helium sufficiently by evaporating it without supply of heat was not successful, and only served to reach the lowest temperature recorded up to that time.

The evaporation of even a very small quantity, when the pressure of the vapour is small, demands the continuous carrying away of colossal volumes of vapour. With vacuum pumps of very large capacity we succeeded in lowering the pressure to 0.2 millimetre. The temperature then reached was 1.15°K according to the law of vapour pressure found. (Of course we can only make an estimate here. The working out of the thermometry of these low temperatures with, amongst other things, the aid of the Knudsen hot wire manometer is still in its initial stages.) Since it would have needed new equipment, I deferred the question as to whether helium can be made to freeze in favour of other, more urgent problems, which could be tackled with the equipment available.



## Hint of a quantum effect...?

It is very noticeable that the experiments indicate that the density of the helium, which at first quickly drops with the temperature, reaches a maximum at 2.2°K approximately, and if one goes down further even drops again. Such an extreme could possibly be connected with the quantum theory.

# Zero-point energy prevents helium from solidifying

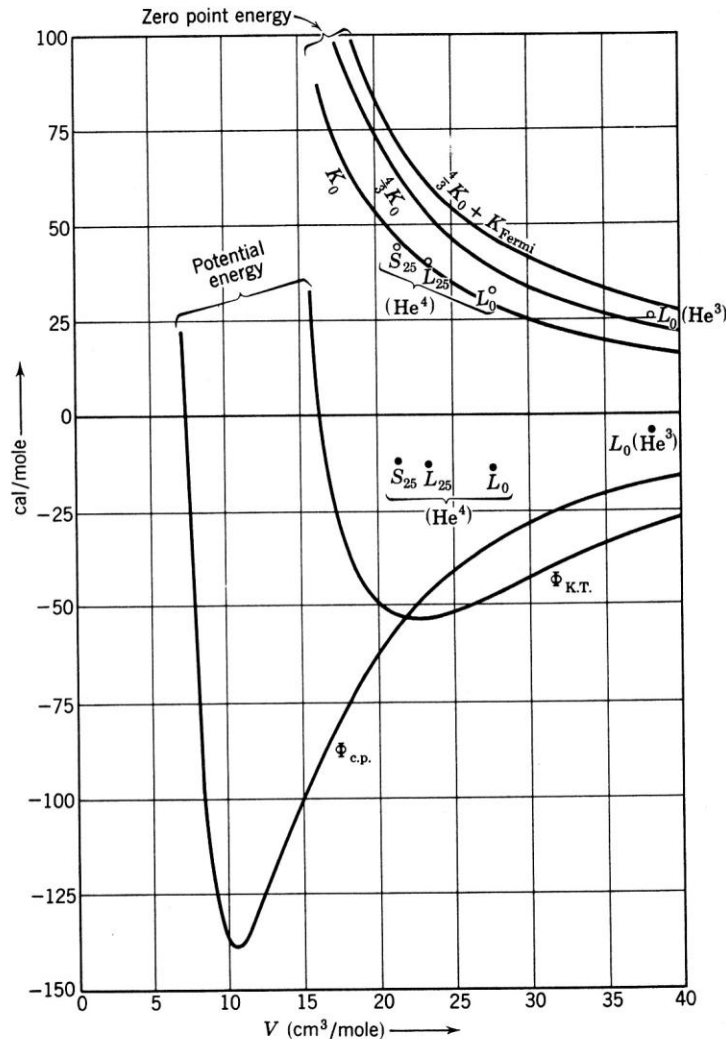


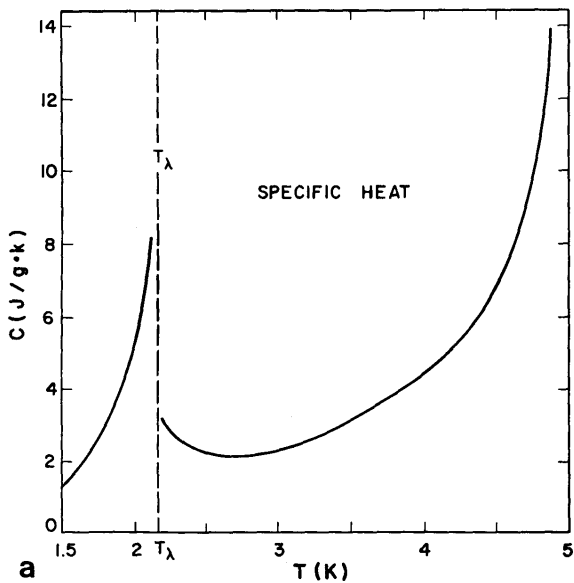
Fig. 15. The potential energy of the close-packed structure,  $\Phi_{c.p.}$ , and of the  $T_d^2$  configuration suggested by Keesom and Taconis,  $\Phi_{K.T.}$ . The curve  $K_0$  gives the zero point energy of eq. (6), §5. The solid circles refer to the experimental energy content of condensed  $\text{He}^4$  and  $\text{He}^3$  at  $0^\circ\text{K}$ . The open circles refer to the "experimental zero point energies," defined as the difference between the experimental total energies and the lowest potential energy.



F. London



# Discovery of He II phase transition (1928) Helium phase diagram (1933)



W.H. Keesom

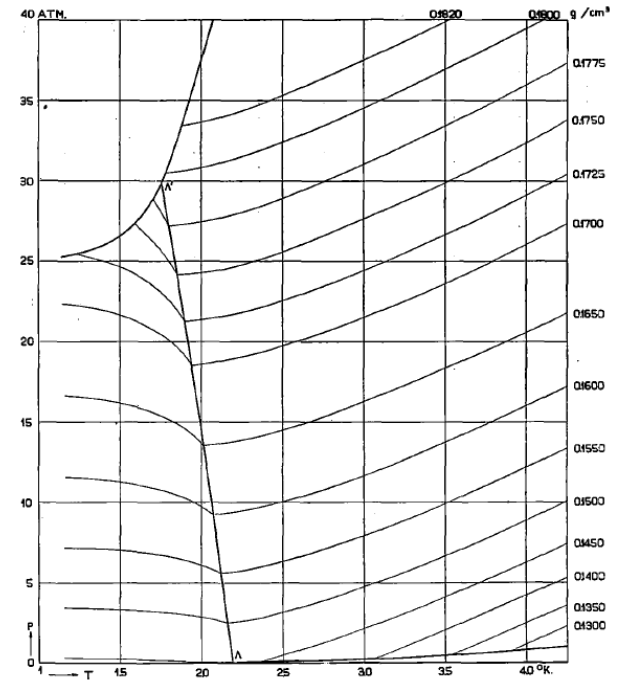


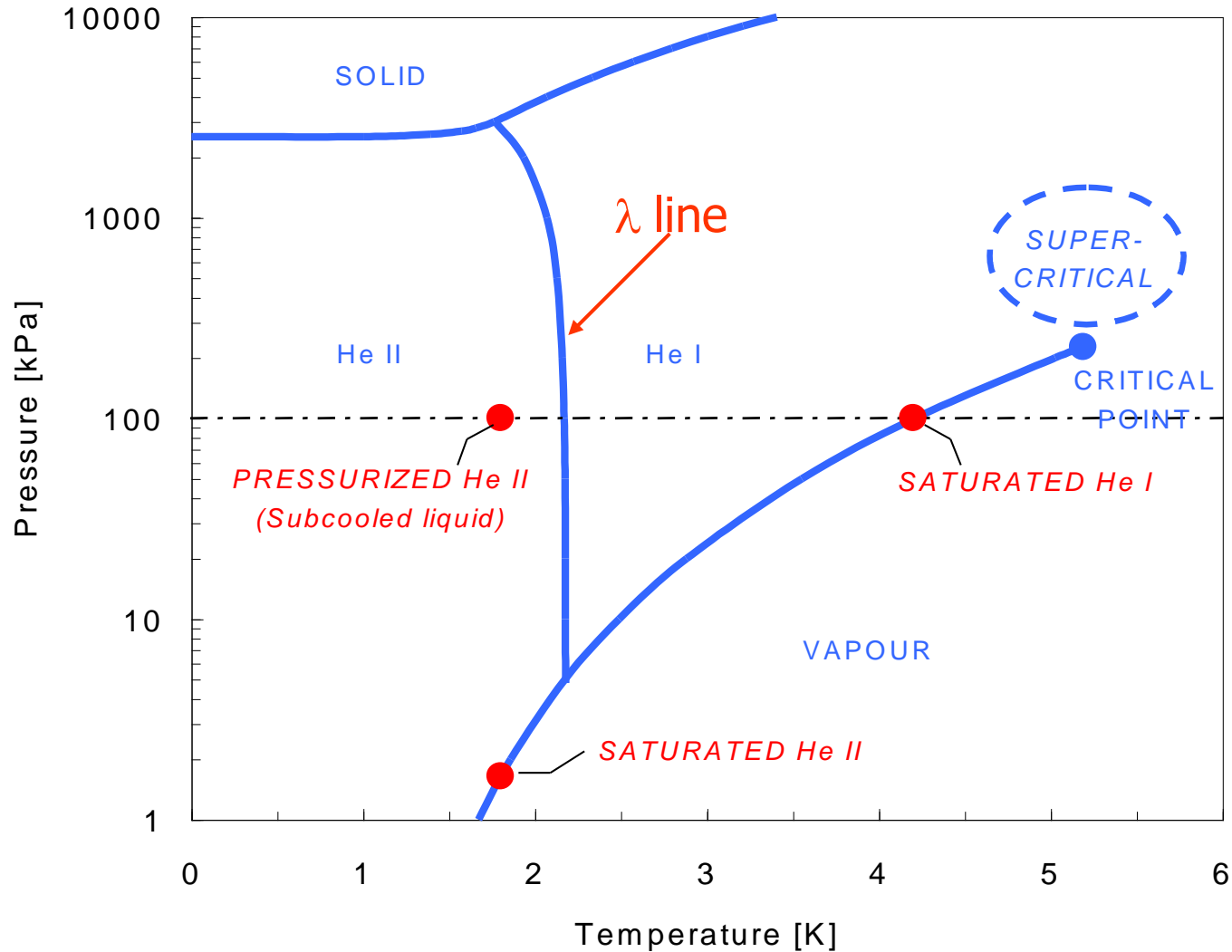
Fig. 1.

## THERMODYNAMIC DIAGRAMS OF LIQUID HELIUM

by W. H. KEESOM and Miss A. P. KEESOM

Supplement No. 76b to the Communications from the Kamerlingh  
Onnes Laboratory at Leiden

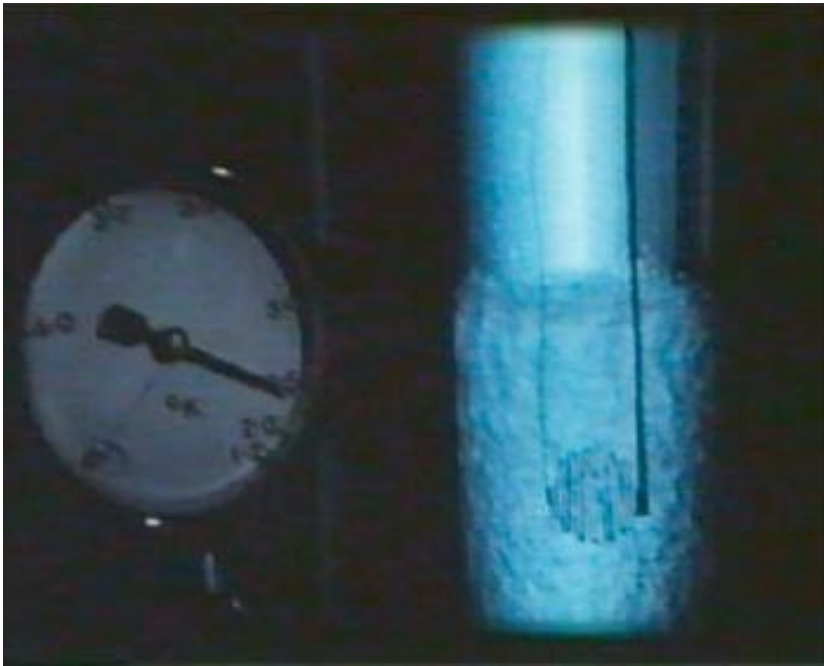
# Phase diagram of helium



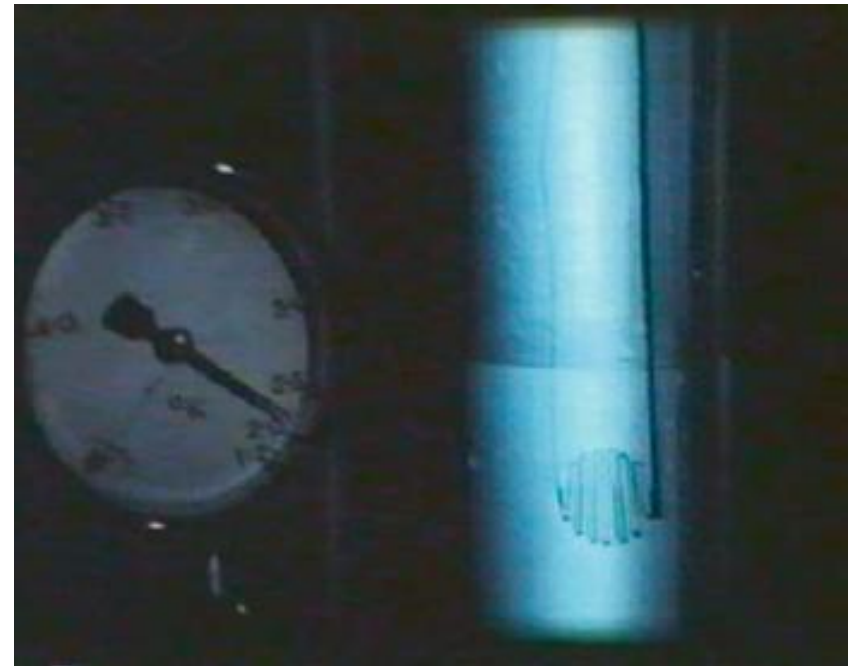
# Discovery of superfluidity in He II (1938)

J.F. Allen & A.D. Misener (Cambridge)  
P.L. Kapitsa (Moscow)

*Vaporization of liquid helium under applied heat load*



He I (T=2.4 K)



He II (T=2.1 K)

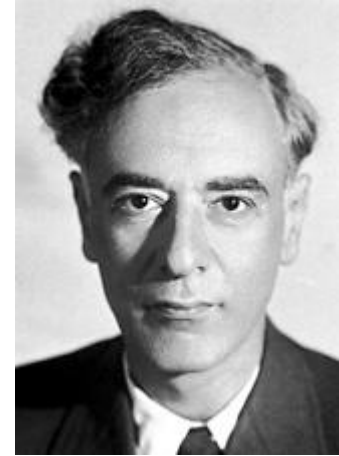
# Theoretical approaches to superfluid helium



Fritz London



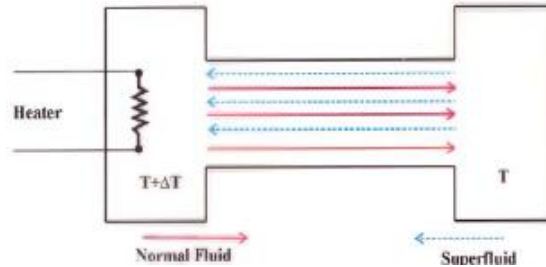
Laszlo Tisza



Lev Davidovich Landau

$$T_{\text{BEC}} = \left( \frac{2\pi\hbar^2}{1.897mk_B} \right) n^{2/3}$$

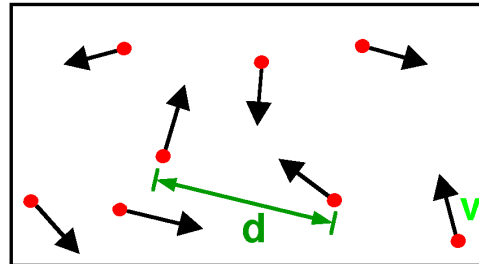
Bose-Einstein condensation



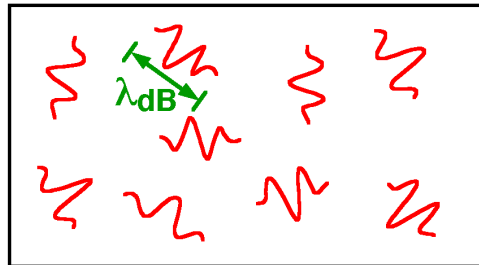
Two-fluid model

$$\epsilon = \hbar\omega = \Delta + \frac{(p - p_0)^2}{2\mu}$$

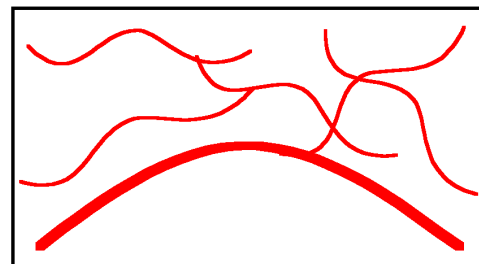
Quasi-particle description



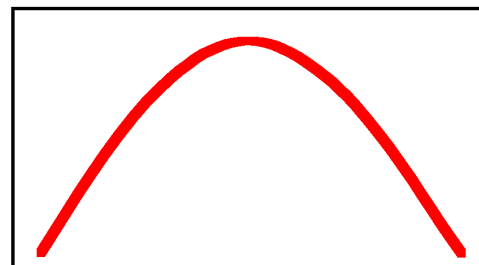
**High Temperature T:**  
 thermal velocity  $v$   
 density  $d^{-3}$   
 "Billiard balls"



**Low Temperature T:**  
 De Broglie wavelength  
 $\lambda_{dB} = h/mv \propto T^{-1/2}$   
 "Wave packets"



**T = T<sub>crit</sub>:**  
 Bose-Einstein  
 Condensation  
 $\lambda_{dB} \approx d$   
 "Matter wave overlap"



**T = 0:**  
 Pure Bose  
 condensate  
 "Giant matter wave"

## Phenomenological model

Two interpenetrating fluids

$$\rho = \rho_s + \rho_n$$

Normal & superfluid fractions varying with T

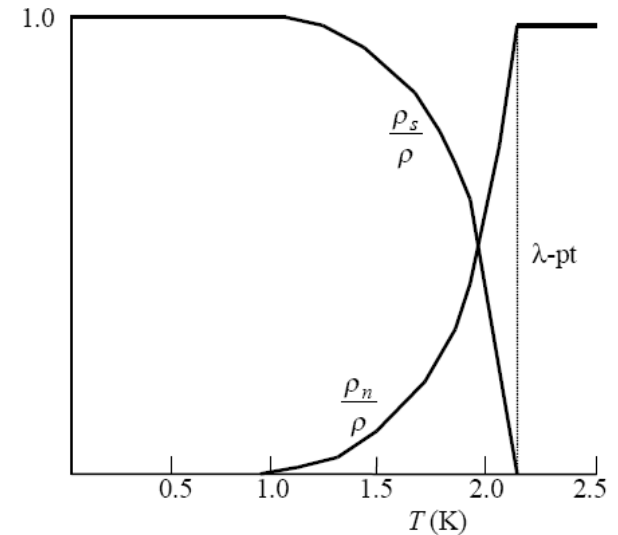
$$\rho \mathbf{v} = \rho_s \mathbf{v}_s + \rho_n \mathbf{v}_n$$

$$\rho s = \rho_n s_n \text{ since } s_s = 0$$

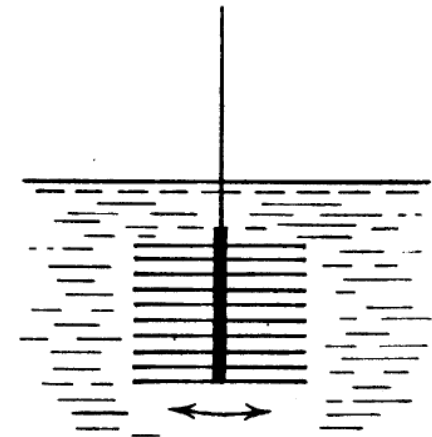
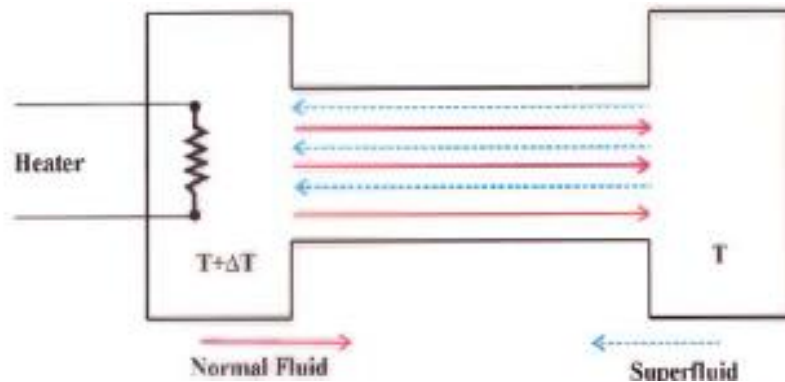
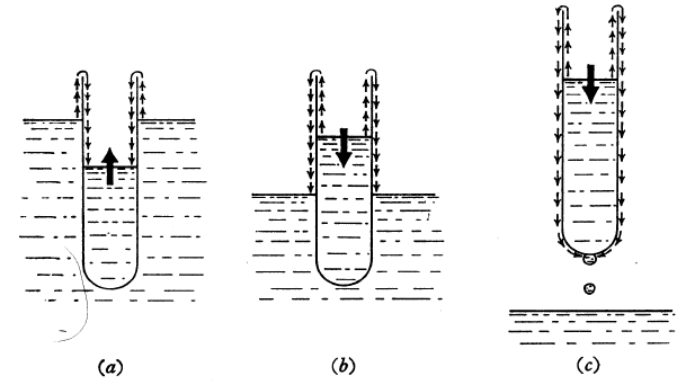
All entropy carried by normal component

## Physical basis of the two-fluid model

- *Collective excitations constitute the normal component (Landau)*
- *B-E condensate in liquid (Penrose & Onsager)*



- Frictionless flow through small channels
  - Film flow
  - Andronikashvili experiment
- Thermal transport by counterflow
  - Laminar
  - Turbulent
- Thermomechanical effect
- Second sound





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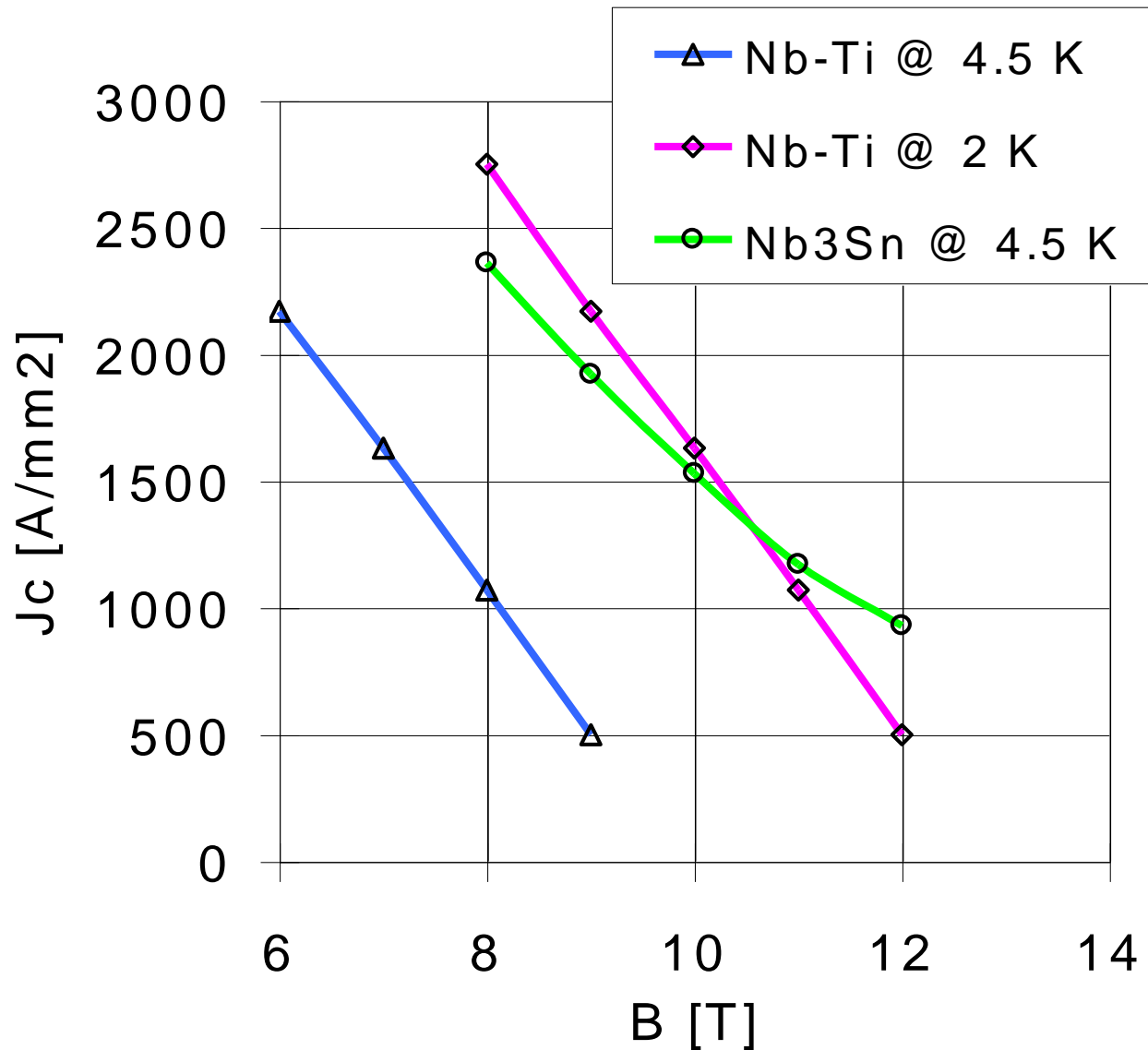


# Benefits of He II cooling



- Lower the operating temperature
  - Achieve higher magnetic field through increase of critical current density of superconductor
  - Minimize overall energy dissipation in RF cavities
- Enhance heat transfer
  - At solid-liquid interface  $\Rightarrow$  conductor cooling
  - In the bulk liquid
    - $\Rightarrow$  device/system cooling scheme
    - $\Rightarrow$  calorimetry in isothermal bath

# Critical current density of superconductors



# Optimization of operating temperature for superconducting RF cavity

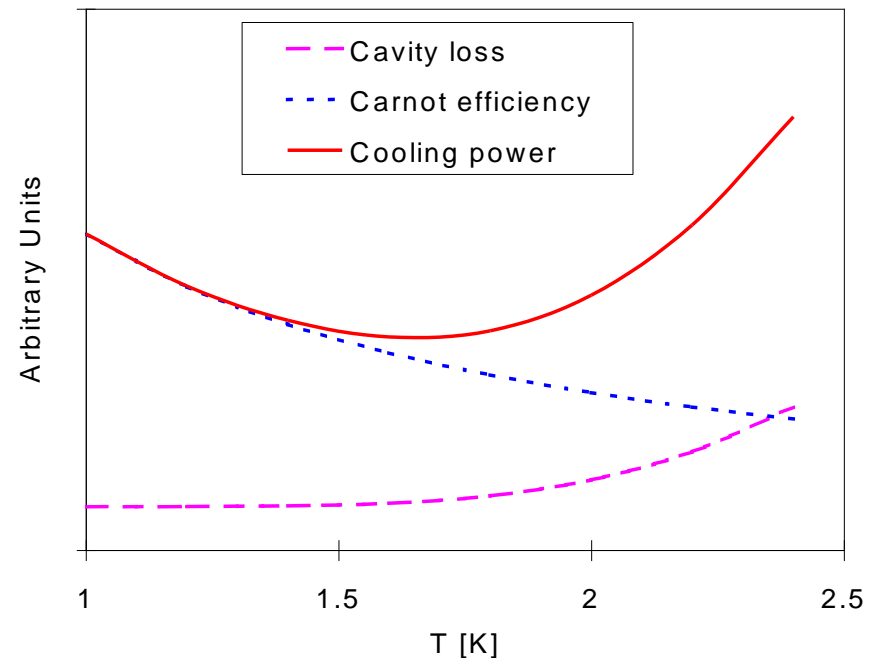
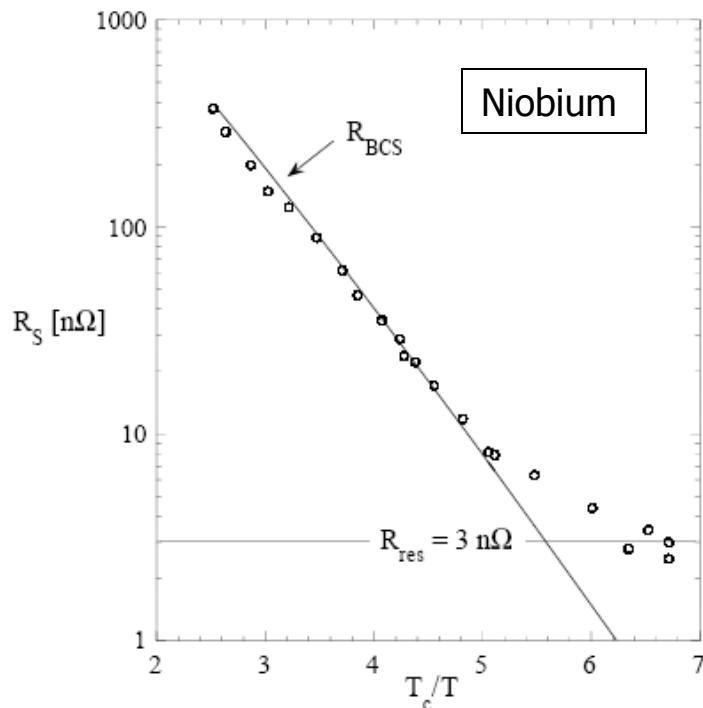
- BCS theory
- For practical materials
- Refrigeration (Carnot)

$$R_{\text{BCS}} = (A \omega^2 / T) \exp(-B T_c / T)$$

$$R_S = R_{\text{BCS}} + R_0$$

$$P_a = P (T_a / T - 1)$$

⇒ *depending upon  $\omega$  and  $R_0$ , optimum operating temperature for superconducting cavities*





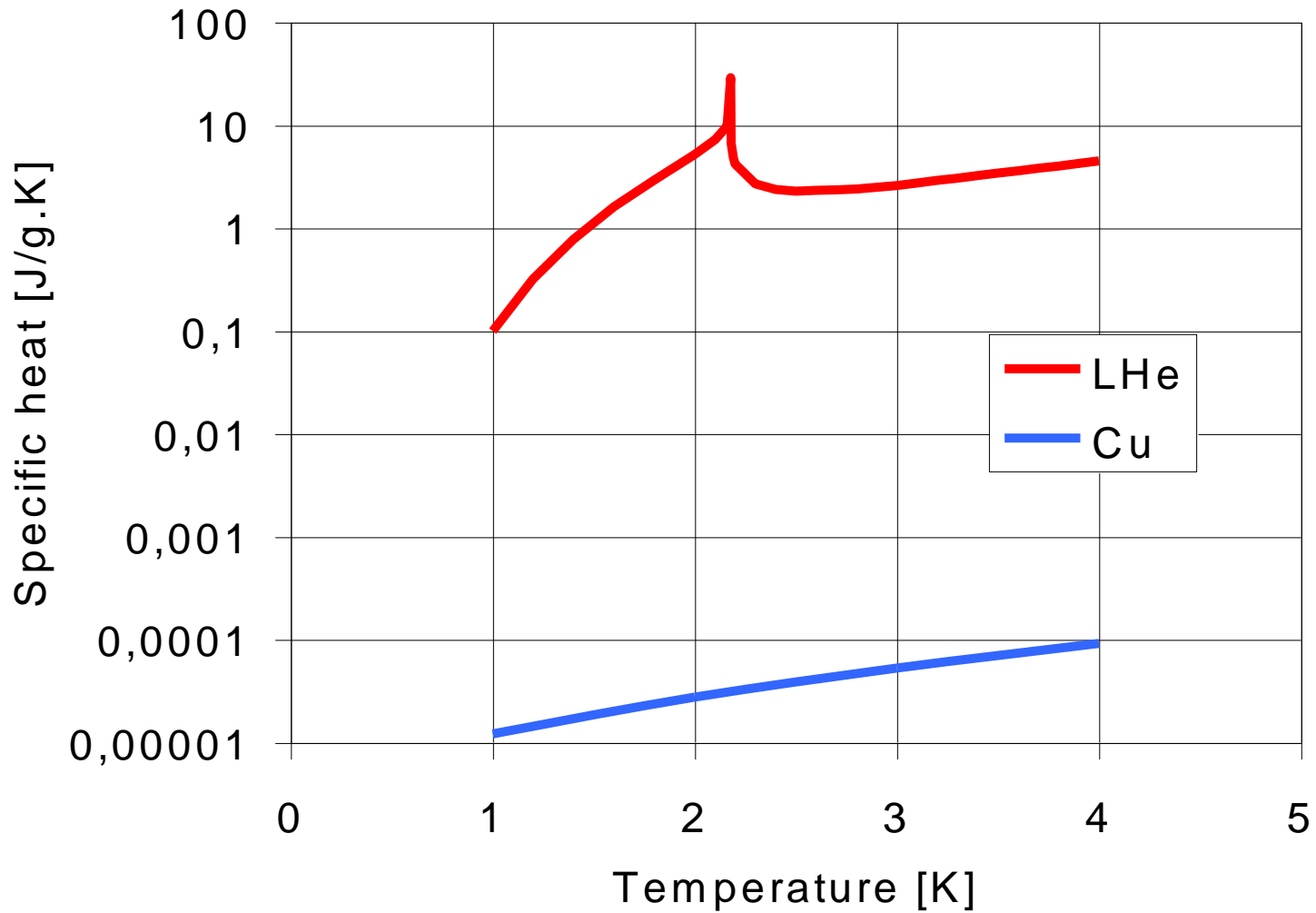
# Enhancement of heat transfer

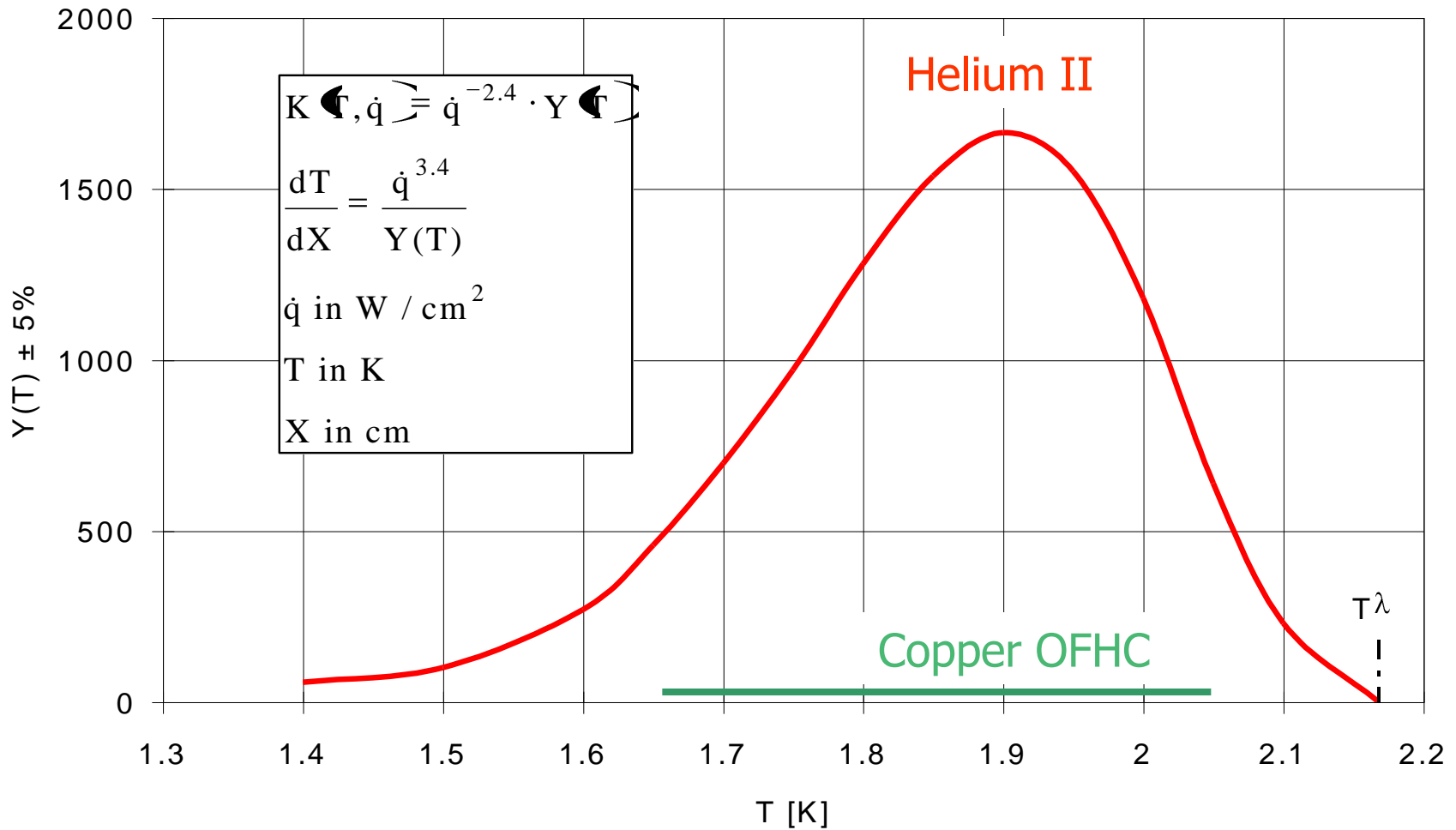


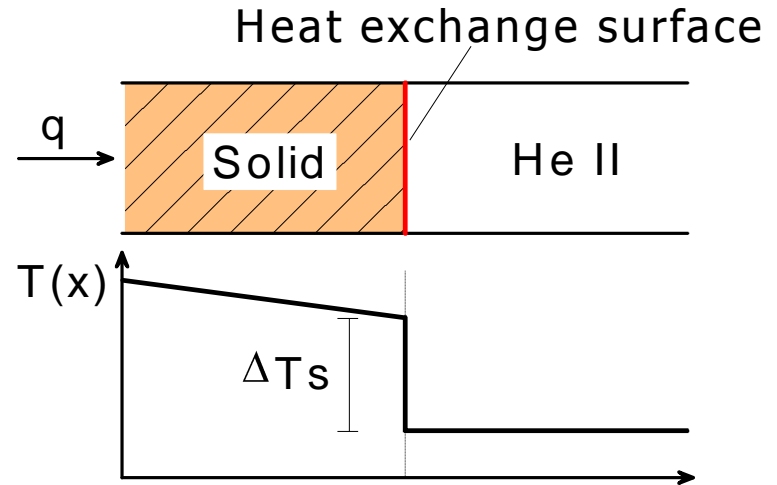
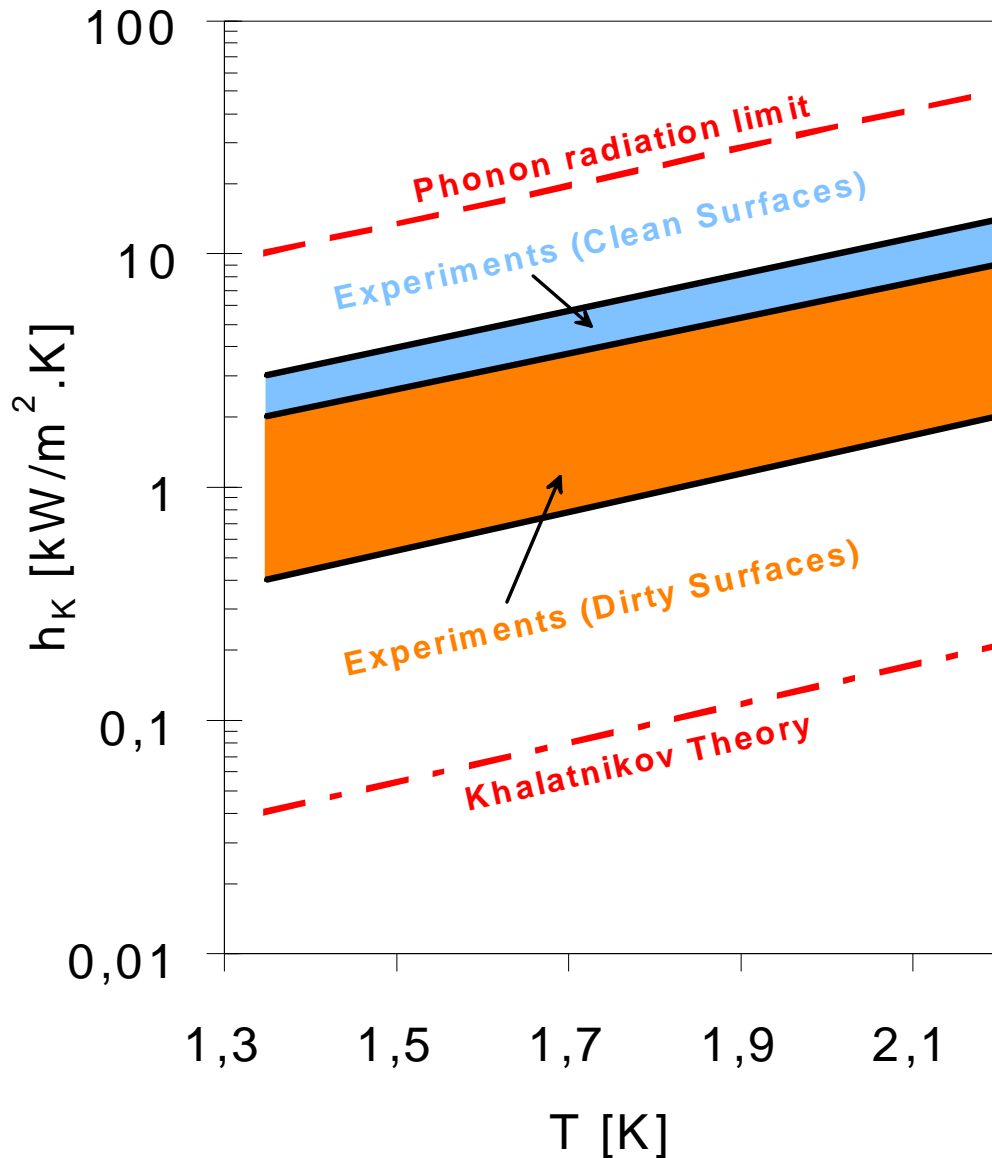
- Low viscosity  $\Rightarrow$  *permeation*
- Very high specific heat  $\Rightarrow$  *stabilization*
  - $10^5$  times that of the conductor per unit mass
  - $2 \times 10^3$  times that of the conductor per unit volume
- Very high thermal conductivity  $\Rightarrow$  *heat transport*
  - $10^3$  times that of cryogenic-grade OFHC copper
  - peaking at 1.9 K

*Full benefit of these transport properties can only be reaped by appropriate design providing good wetting of the superconductors and percolation paths in the insulation, often in conflict with other technical requirements*

# Specific heat of liquid helium and copper







Experimental data for Copper  
(S. Van Sciver, "Helium  
Cryogenics")

$$h_K \sim T^3$$

Valid for small heat flux  
(when  $\Delta T \ll T$ )

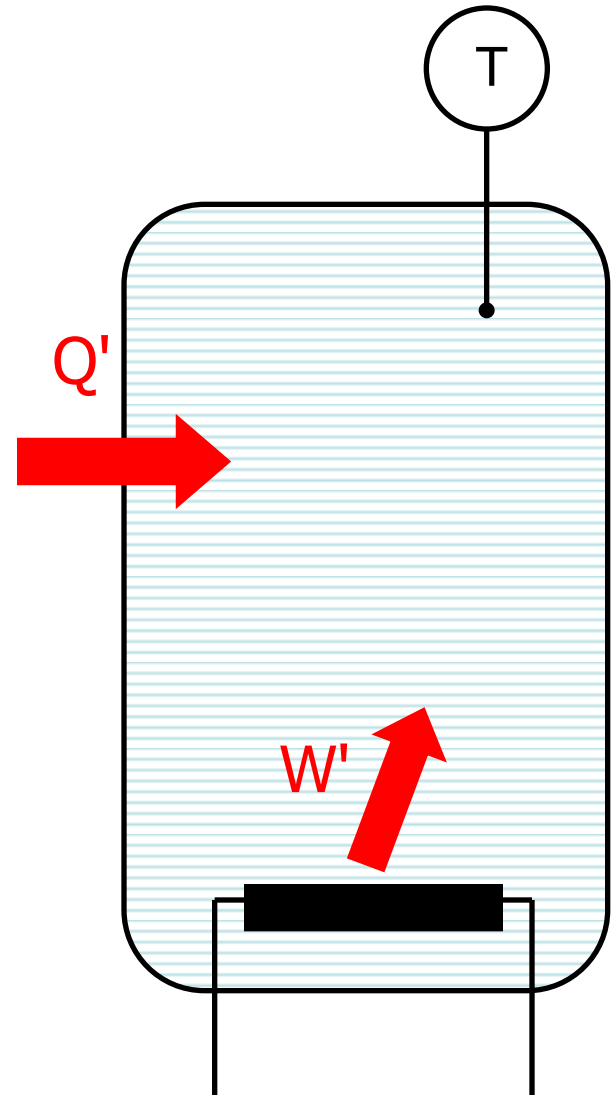
# Calorimetry in isothermal He II bath

- For slow thermal transients, the He II bath is quasi-isothermal: a single temperature measurement allows to estimate heat deposition/generation  $Q'$

$$Q' = M_{\text{bath}} \left. \frac{dH}{dt} \right|_1$$

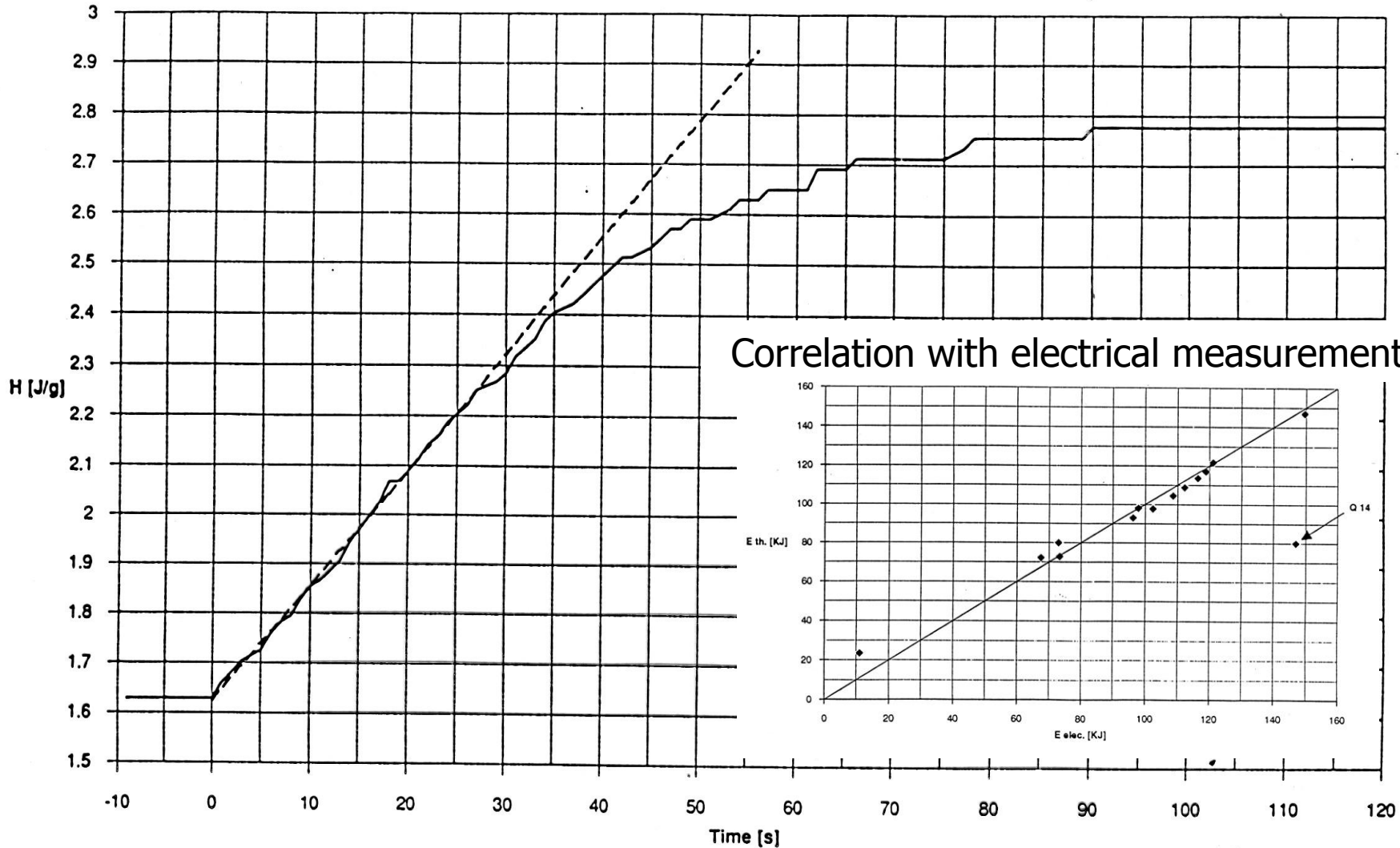
- $M_{\text{bath}}$  can be estimated by *in situ* calibration, using applied heating power  $W'$

$$W' = M_{\text{bath}} \left. \frac{dH}{dt} \right|_2$$

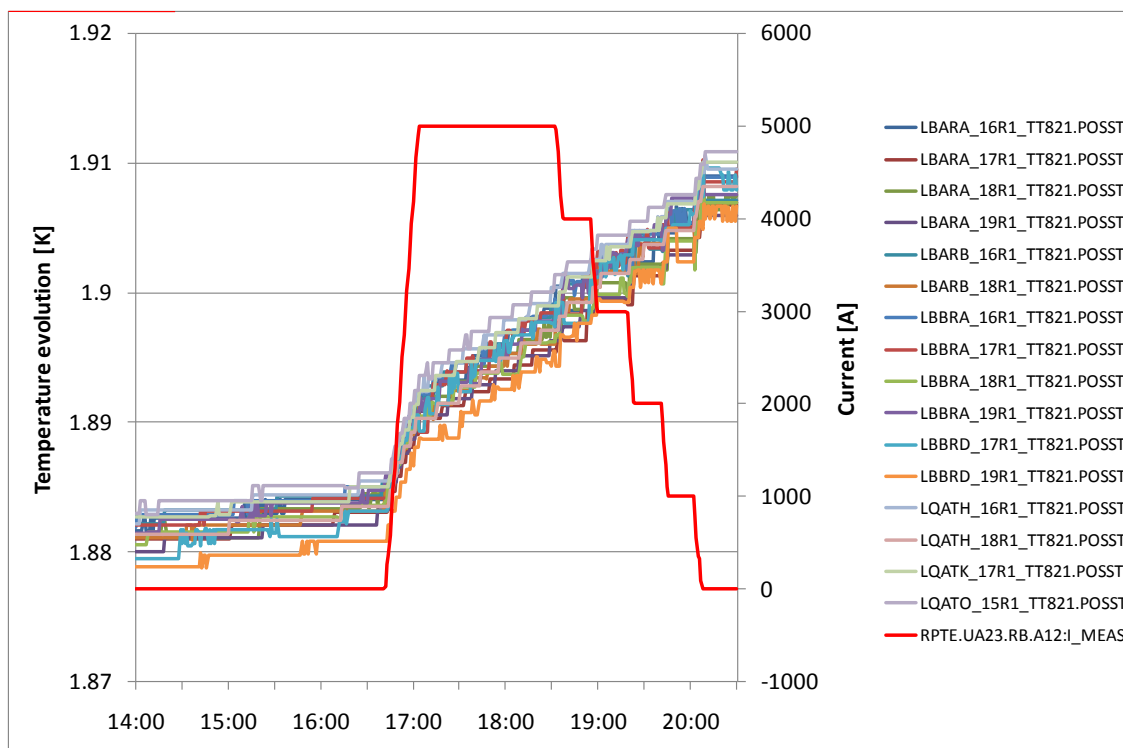




Time evolution of bath temperature



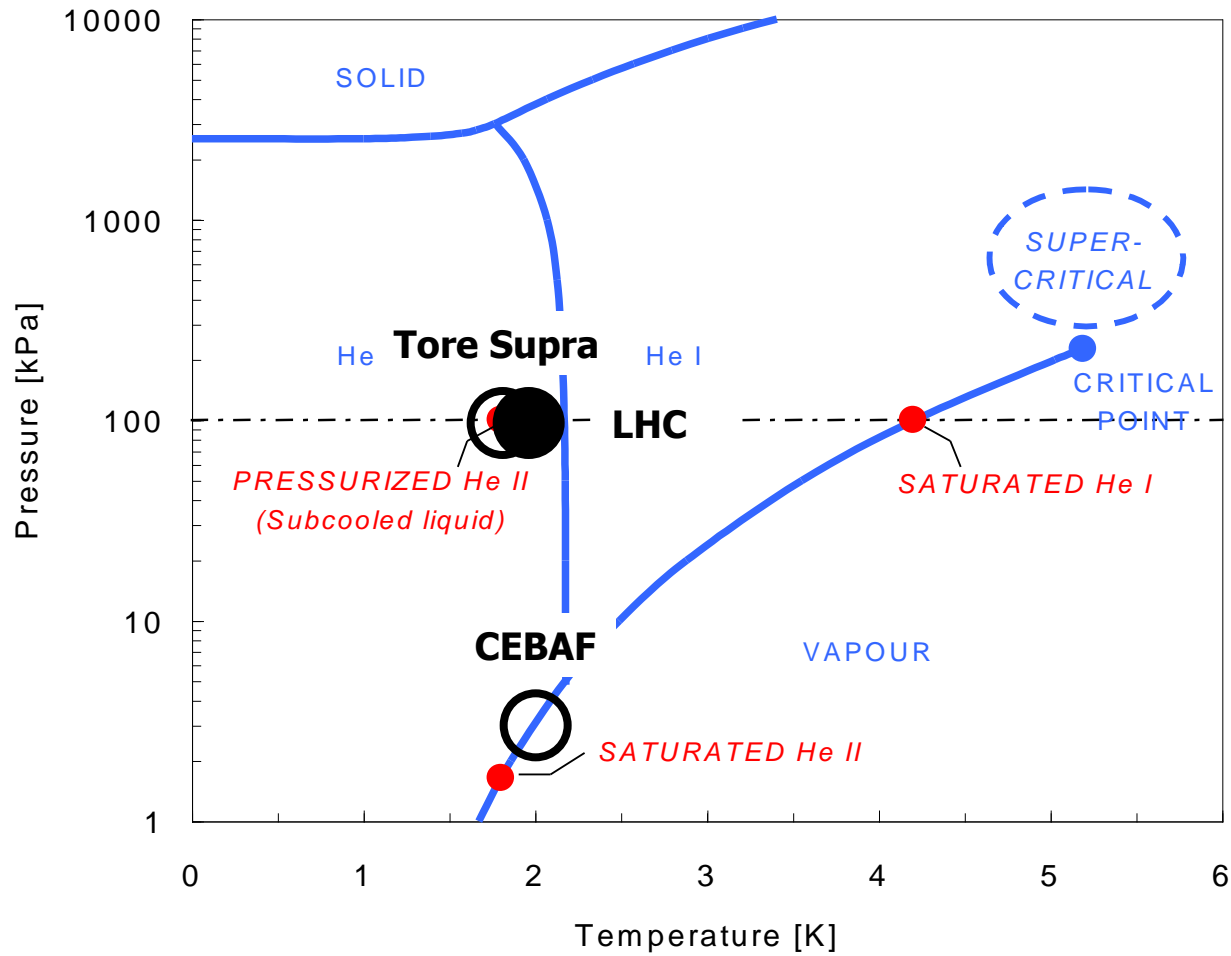
# Precision thermometry allows calorimetric detection of faulty joints in LHC at safe powering level



Current	Total (measured)		Nominal Splices*	Add. local dissipation	Uncertainty
[A]	[mW/m]	[W]	[W]	[W]	[W]
3000	4.4	1.0	0.4	0.6	0.6
5000	14.9	3.2	1.1	2.1	0.6
7000	32.2	6.9	2.1	4.8	0.6

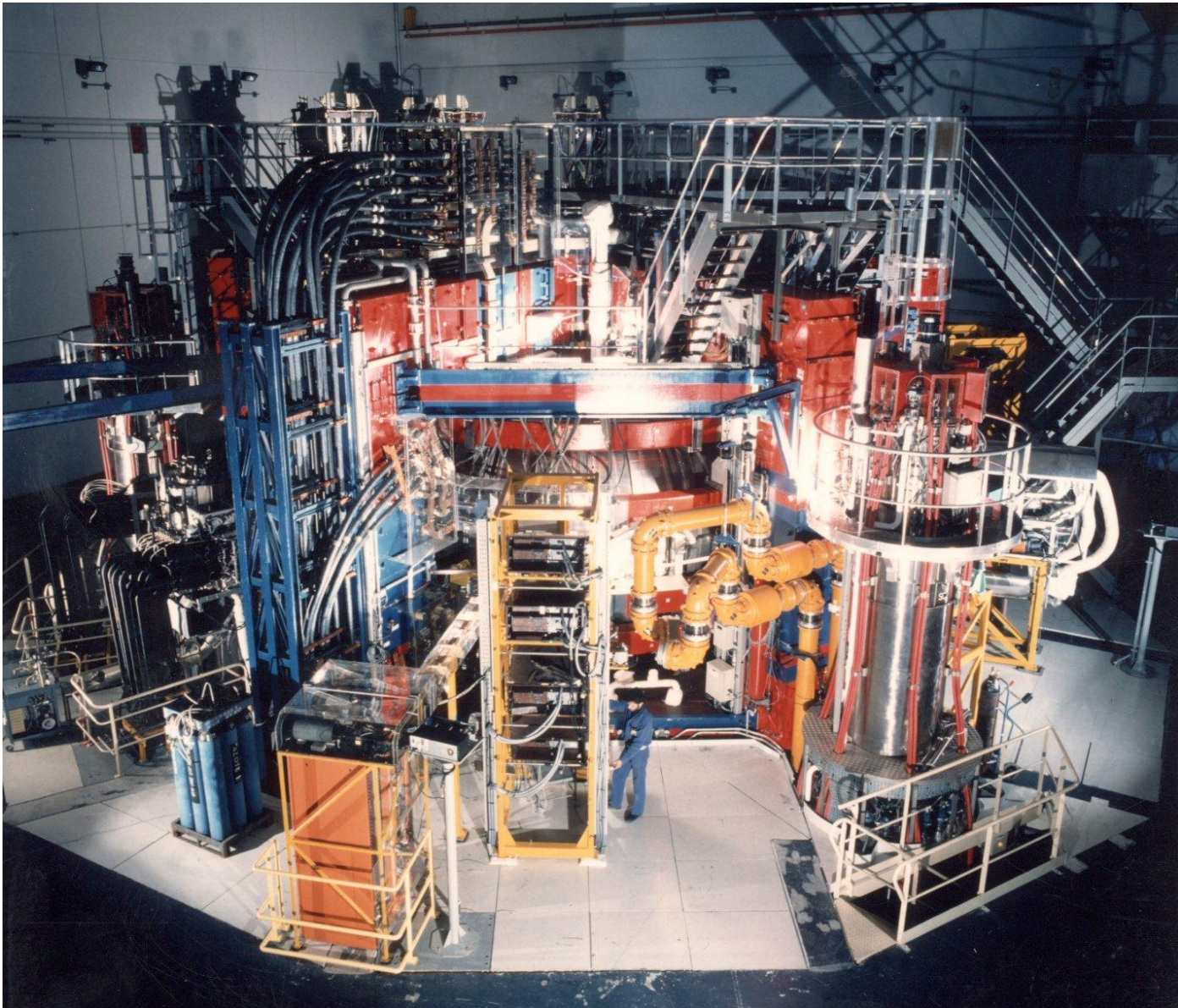
→ Local resistance:  $\sim 90 \text{ n}\Omega$ , confirmed by electrical measurement

# He II cooling of large systems





# Tore Supra tokamak at CEA Cadarache, France

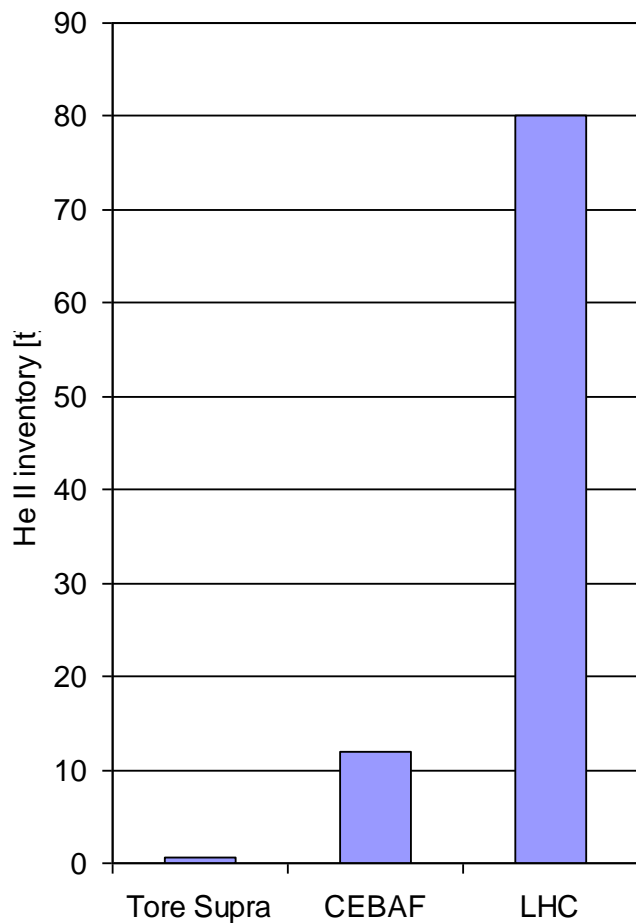


# CEBAF at Jefferson Laboratory, USA

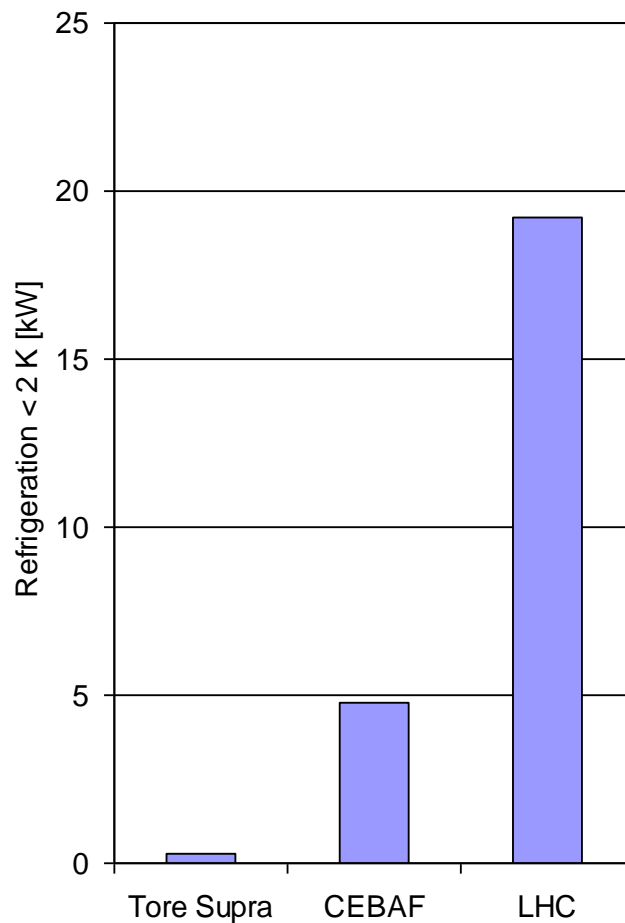


# Characteristics of large-scale He II systems

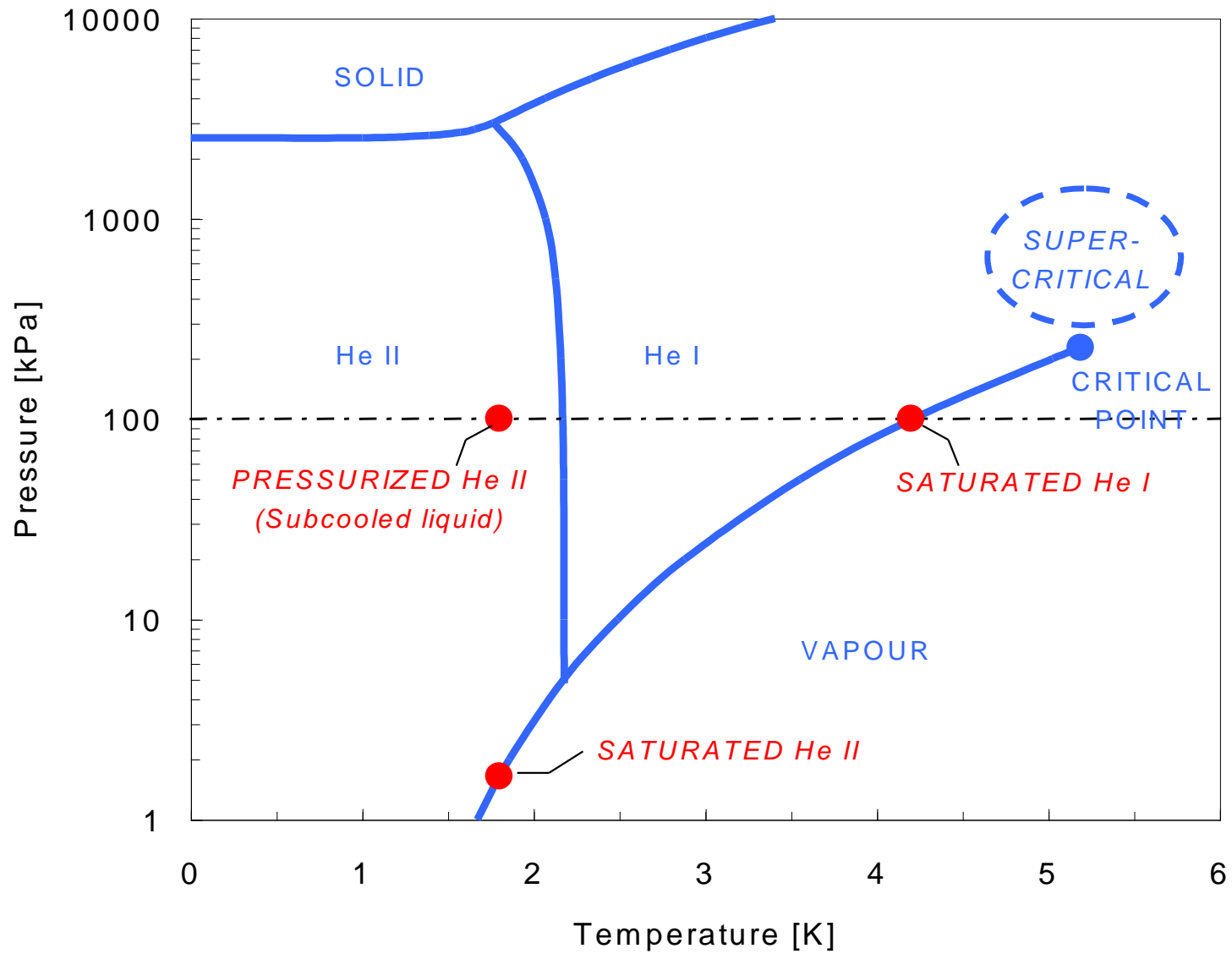
### He II inventory



### Refrigeration power < 2 K



# Pressurized vs saturated He II





# Properties of fluids at saturation



	He I	He II	Water
Psat [bar]	1	0.016	1
Tsat [K]	4.2	1.8	373
L/V Density ratio	7.5	320	1600
Latent heat [J/g]	20.8	23.4	2260
Viscosity [ $\mu$ Pa.s]	3.2	1.3	282



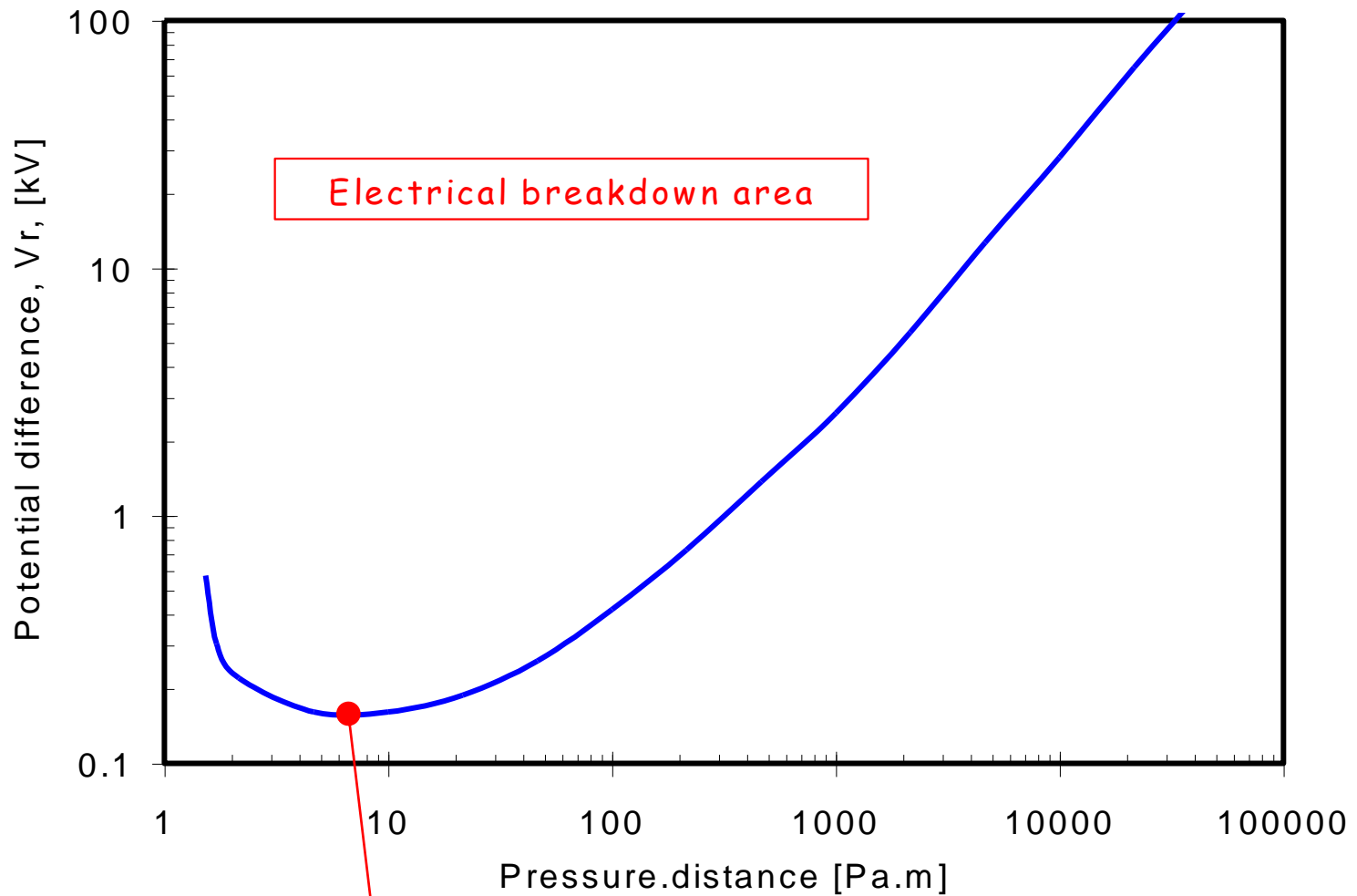


# Advantages & drawbacks of He II p



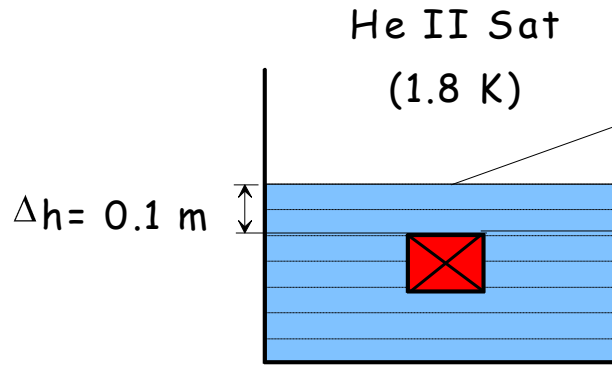
- Advantages
  - limits the risk of air inleaks and contamination in large and complex cryogenic systems
  - for electrical devices, limits the risk of electrical breakdown at fairly low voltage due to the bad dielectric characteristics of helium vapour (Paschen curve)
  - better stabilizer for heat buffering
- Drawbacks
  - one more level of heat transfer
  - additional process equipment (pressurized-to-saturated helium II heat exchanger)

# Paschen curve for gaseous He at 20 °C



e.g., 1.6 kPa pressure and 6 mm gap

# Thermal stabilisation in He II s



**At the liquid surface:**

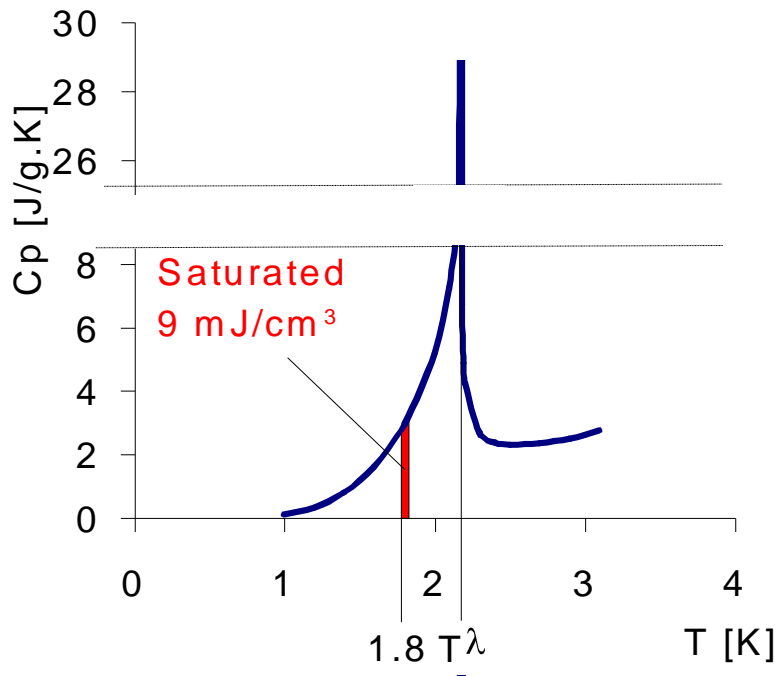
$$P = P_{\text{sat}} = 16.4 \text{ mbar}$$

**At the device surface:**

(adding hydrostatic head)

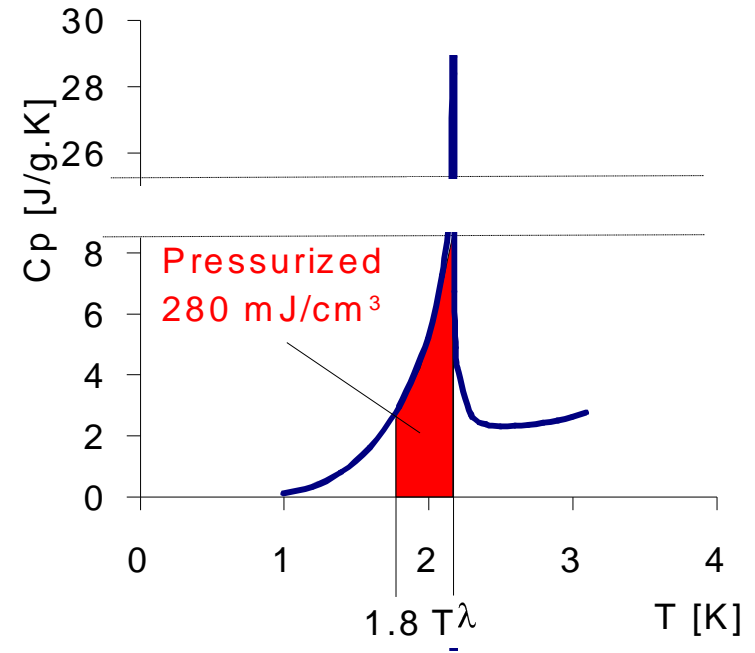
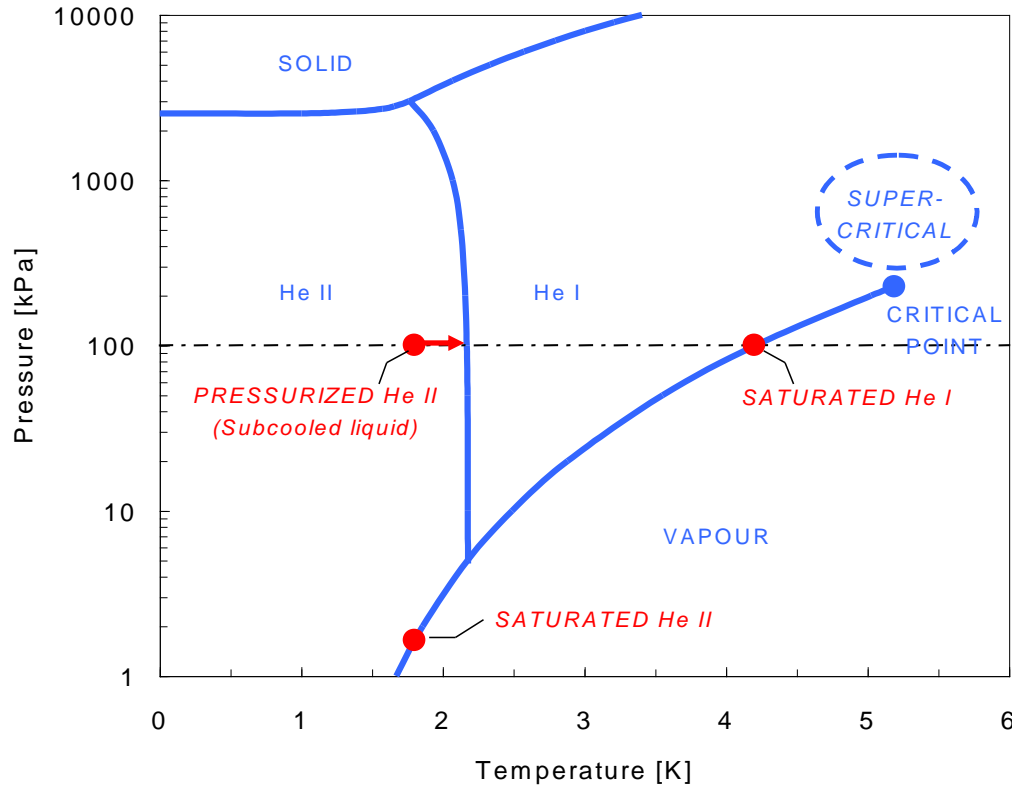
$$P^* = P + \rho \cdot g \cdot \Delta h = 17.8 \text{ mbar}$$

for which corresponds a temperature of saturation of  $T^* = 1.82 \text{ K}$



The corresponding  $\Delta H$  is  $9 \text{ mJ/cm}^3$   
(vs.  $1 \text{ mJ/cm}^3$  in He I at 4.2 K)

# Thermal stabilisation in He II p



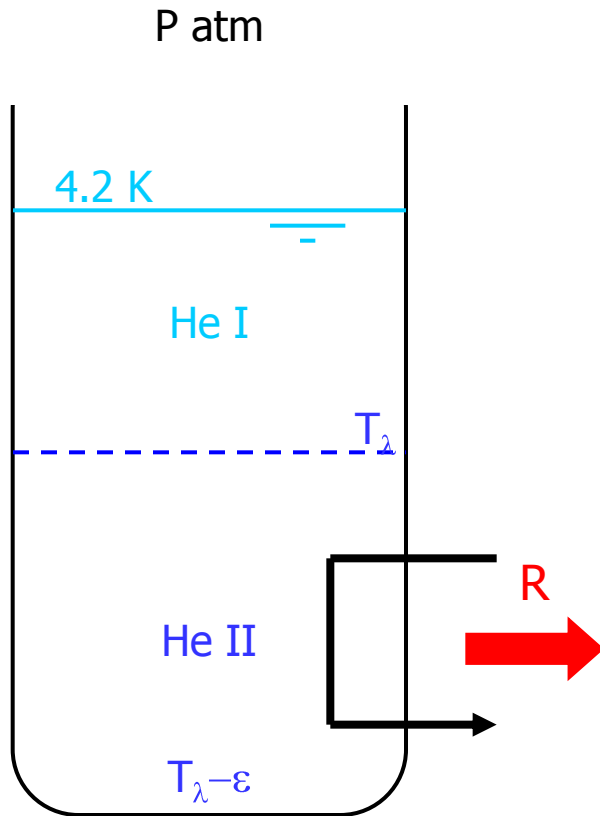
The corresponding  $\Delta H$  is 280 mJ/cm<sup>3</sup>  
(30 times higher than in He II s)



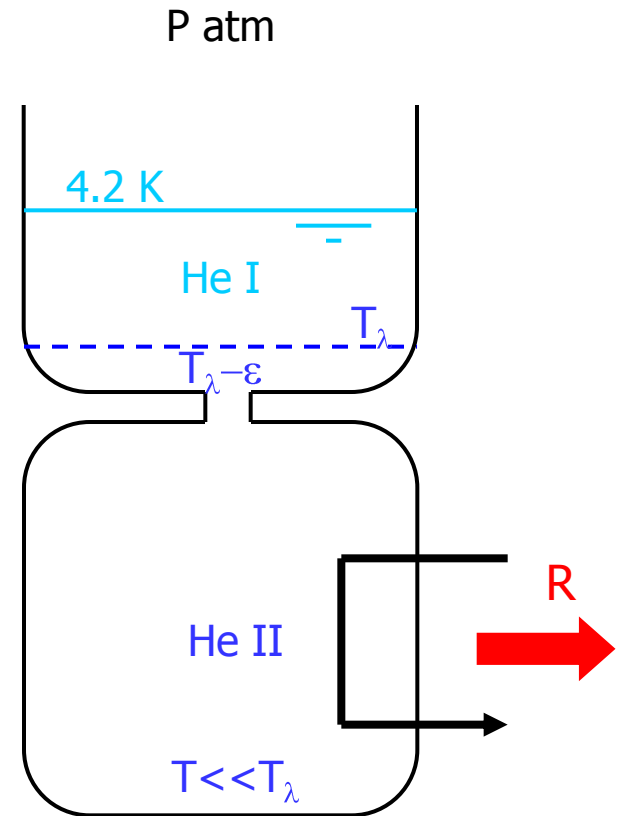
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# Working with superfluid helium at atmospheric pressure



« *Roubeau bath* »: He II conduction prevents from lowering the bath temperature well below  $T_\lambda$



« *Claudet bath* »: restriction in cryostat allows subcooling He II bath to temperatures well below  $T_\lambda$

# A practical Claudet bath

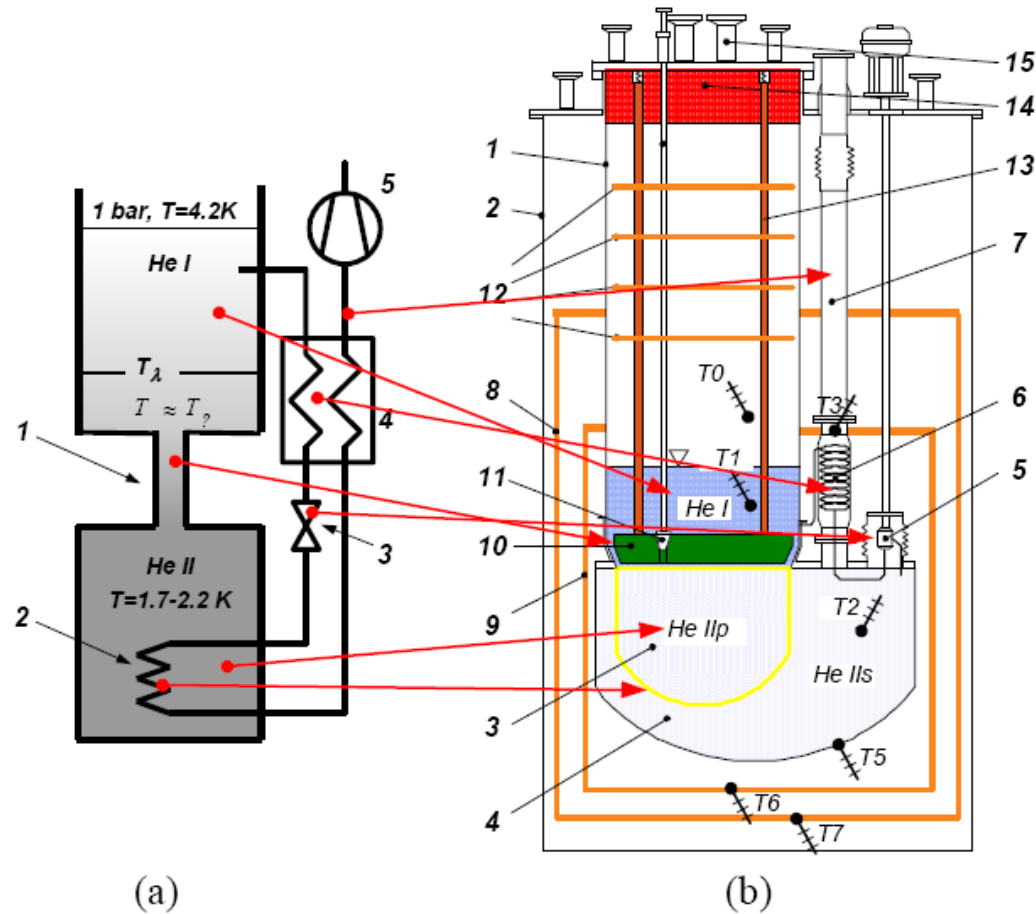
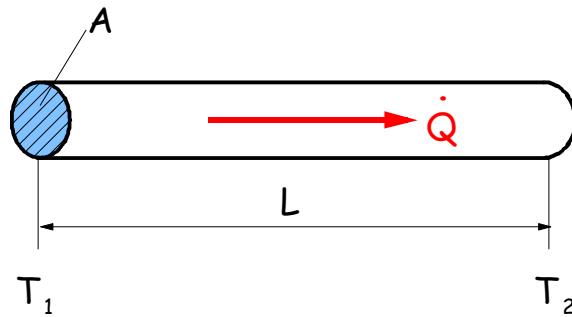


Figure 1. a) Claudet bath principle: 1 – constriction, 2 – He IIs/He IIp heat exchanger, 3 – J-T valve, 4 – recuperative heat exchanger, 5 – vacuum pump; b) NED cryostat scheme: 1 – He I vessel, 2 – vacuum container, 3 – He IIp vessel, 4 – He II s vessel, 5 – J-T valve, 6 – recuperative heat exchanger, 7 – heat exchanger pipe, 8/9 – external/internal radiation shield, 10 –  $\lambda$ -plate, 11 –  $\lambda$ -valve, 12 – insert radiation shields, 13 –  $\lambda$ -plate supports, 14 – foam insulation, 15 – instrumentation ports, T0 – T7 temperature measurement points

# Conduction cooling in static He II p



Cross-section:  $A$

Length:  $L$

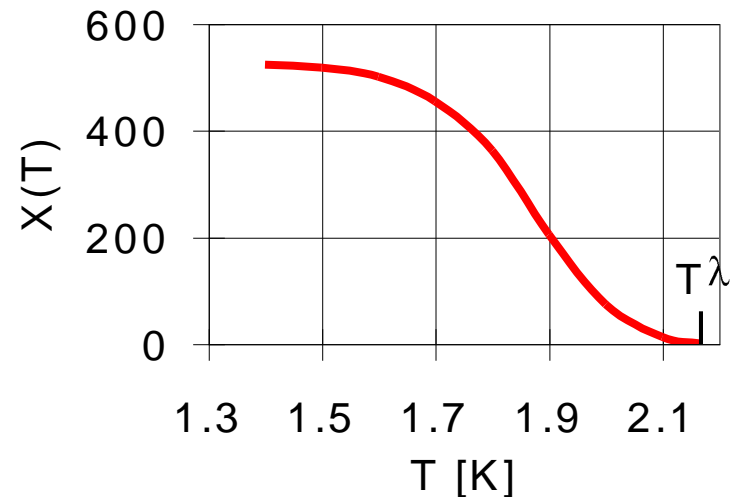
Power:  $\dot{Q}$

Heat Flux:  $\dot{q} = \frac{\dot{Q}}{A}$

• Gorter-Mellink law:  $\dot{q}^m \cdot L = X \left( \frac{T}{T_\lambda} \right)^{-m} - X \left( \frac{T}{T_\lambda} \right)^{-m}$

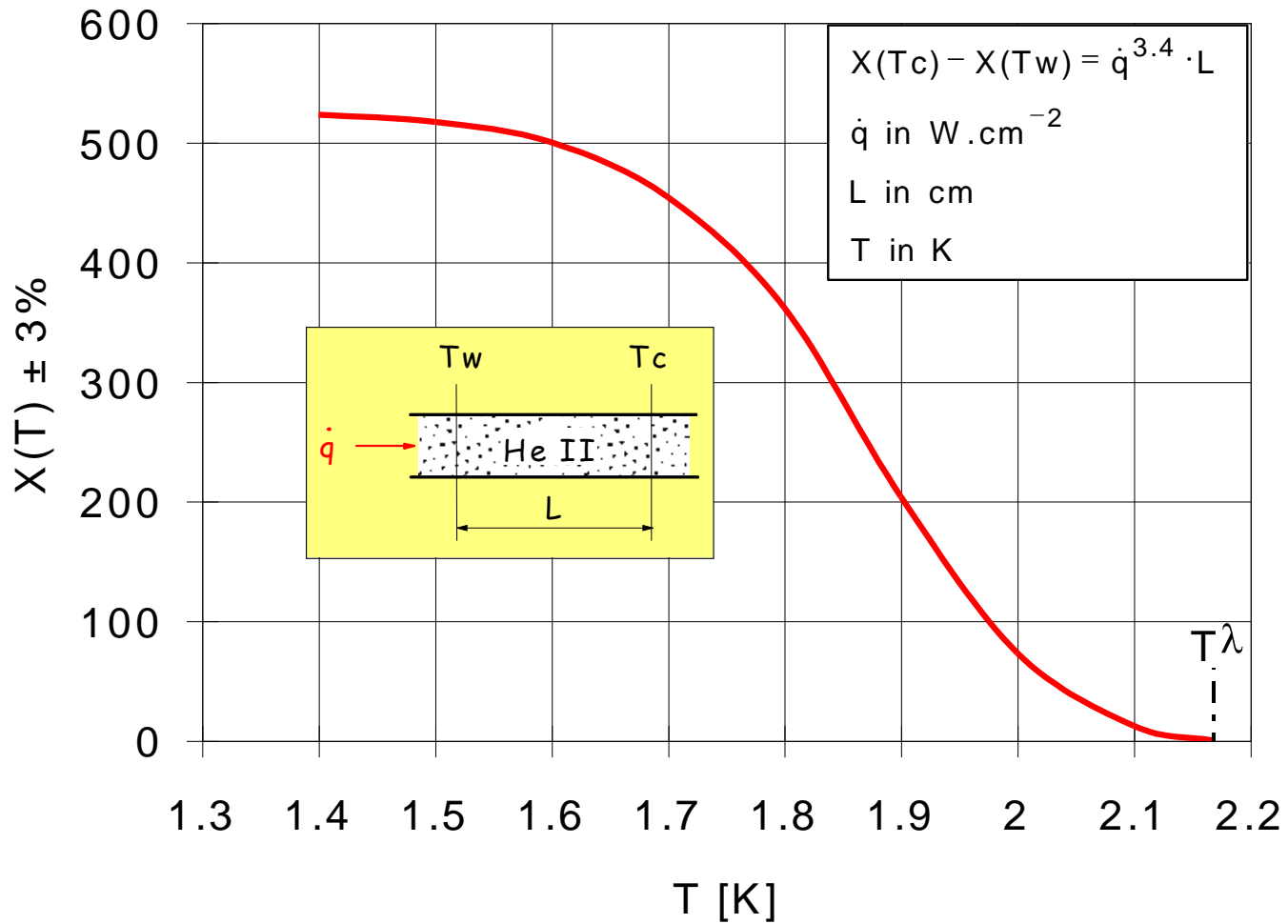
• Experimental work of Bon Mardion, Claudet & Seyfert:

- $m = 3.4$
- tabulation of  $X(T)$

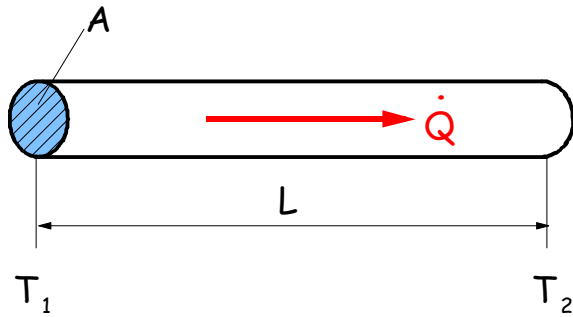




# Thermal conduction integral function of He II p



# Conduction cooling: a numerical example



$$A = 1 \text{ cm}^2$$

$$L = 1 \text{ m} = 100 \text{ cm} \quad (\text{UNITS !})$$

$$T_1 = 1.9 \text{ K} \quad \rightarrow \quad X(T_1) = 200$$

$$T_2 = 1.8 \text{ K} \quad \rightarrow \quad X(T_2) = 360$$

$$\text{Then: } \dot{q}^{3.4} \cdot L = 360 - 200 \quad \rightarrow \quad \dot{q}^{3.4} = 1.6 \quad \rightarrow \quad \dot{q} = 1.15 \text{ W/cm}^2$$

Comparison with "good solid conductor", e.g. **Copper**

$$\dot{q} = k \cdot \frac{\Delta T}{L} \quad \text{With } k \text{ thermal conductivity at } 1.8 \text{ K}$$

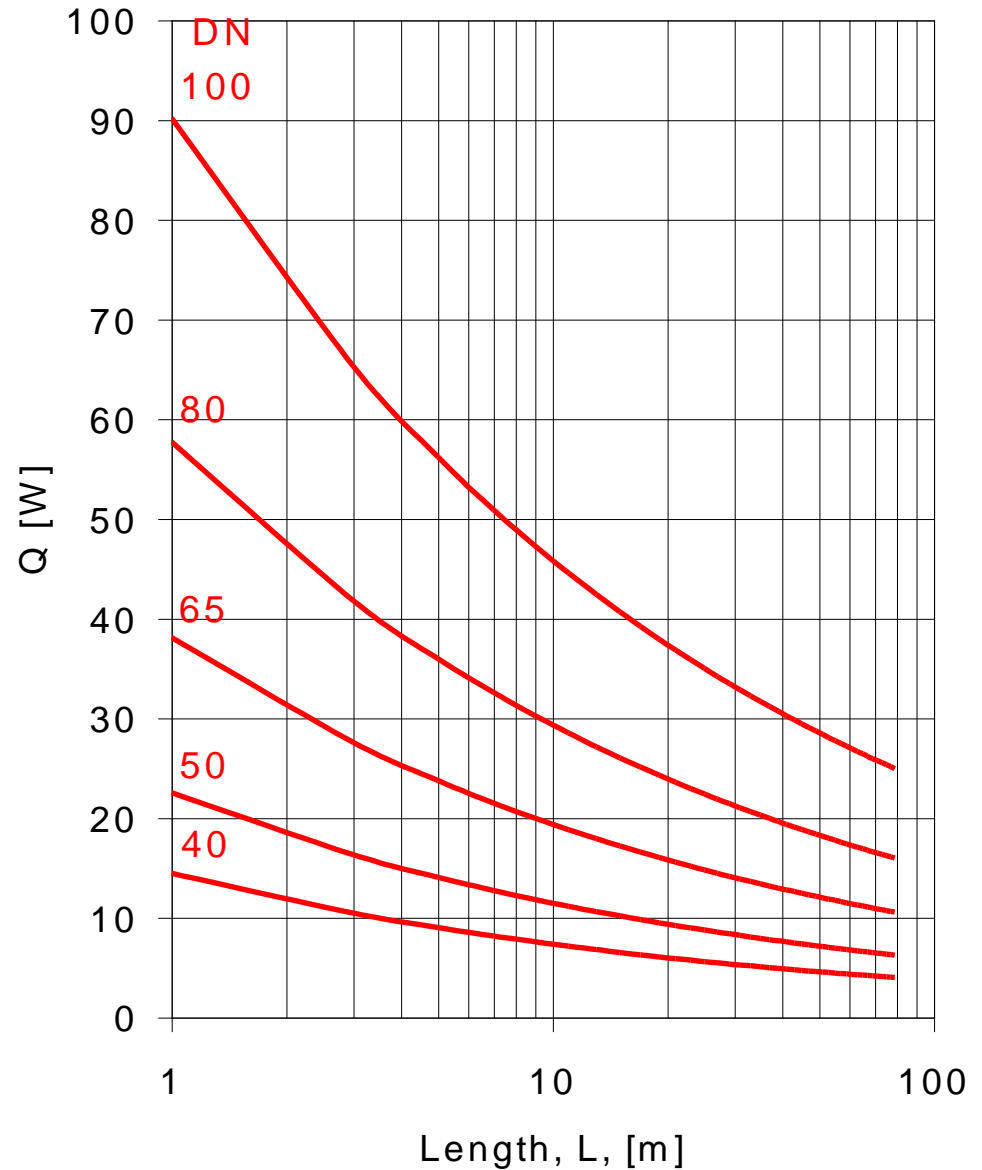
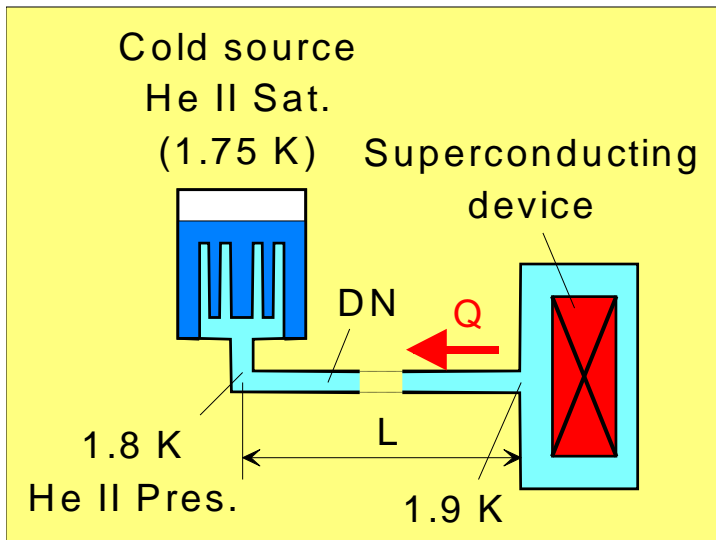
Cu type	k [W/m.K]	$\Delta T$ [K]	L [m]	q [mW/cm <sup>2</sup> ]
<b>OFHC</b>	120	0.1	1	<b>1.2</b>
<b>DHP</b>	3	0.1	1	<b>0.03</b>

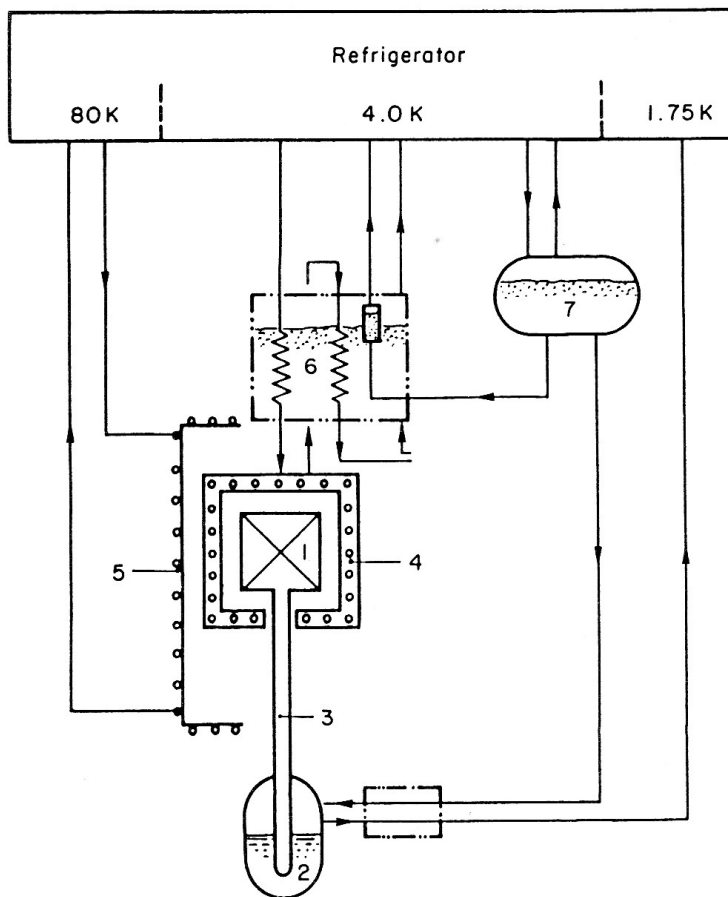
**He II conducts heat 1000 times better than OFHC Cu**

# Steady-state conduction duct in He II

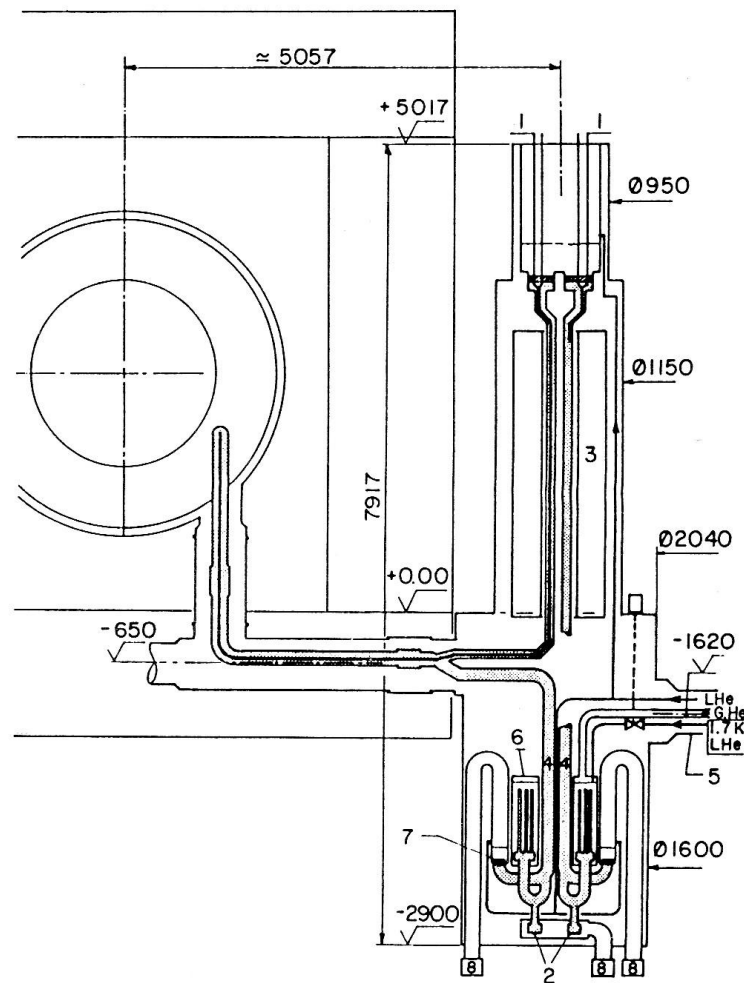


From 1.80 K to 1.90 K

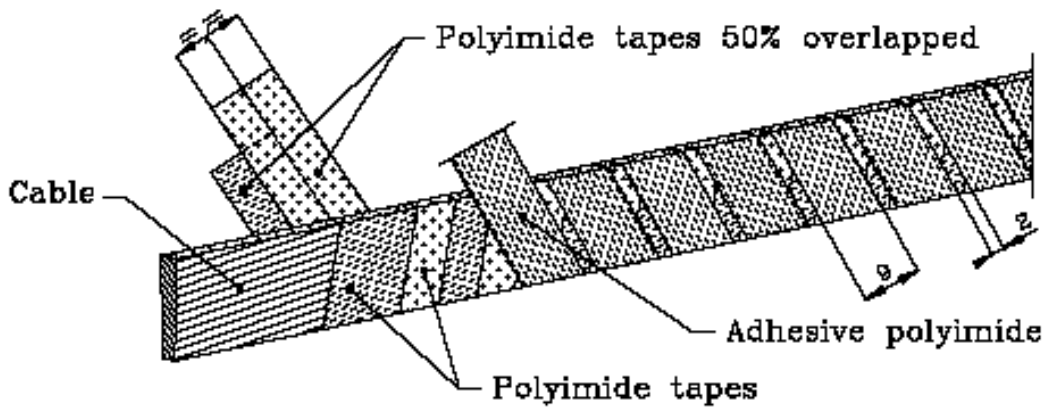




**Figure 2** Schematic design of the cryogenic system. 1, 1.8 K coil; 2, 1.75 K cold box; 3, static pressurized superfluid helium; 4, thick casing total weight = 120 tonnes; 5, 80 K shield total weight = 22 tonnes; 6, thermal ballast; 7, 20 000 dm<sup>3</sup> He tank

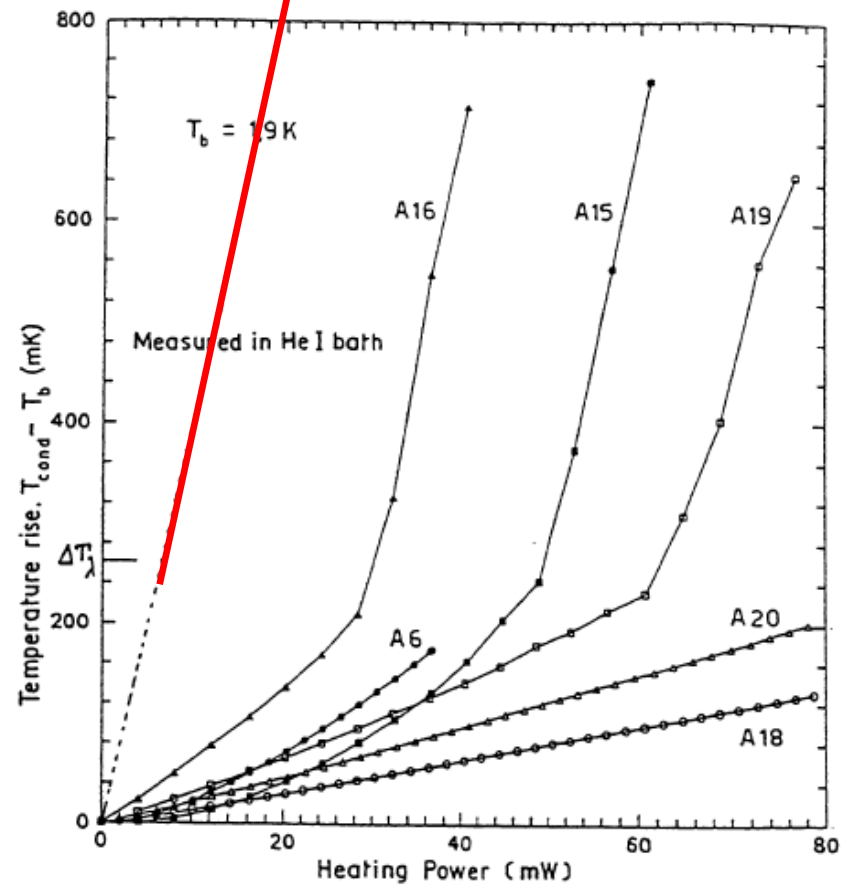


**Figure 3** 1 atm He II circuit. 1, 12 Current leads; 2, cold burst discs; 3, thermal ballast = 1500 dm<sup>3</sup>; 4, six He II pipes; 5, cryogenic line; 6, 1.7 K cold source; 7, cold valves; 8, safety valves



Conduction in polyimide

# Heat transfer across electrical insulation of superconducting cable

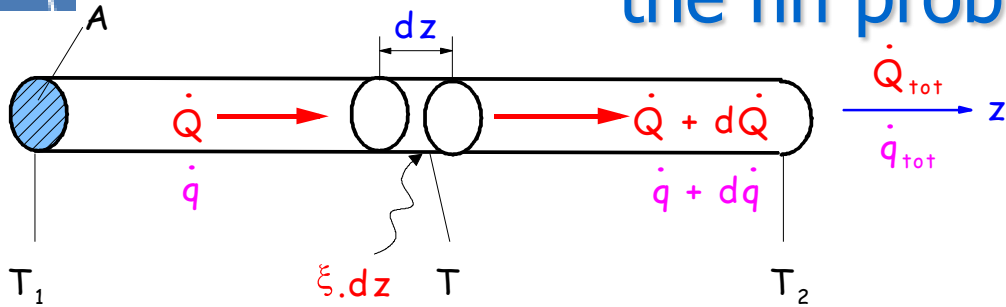


The LHC: 1716 superconducting magnets  
cooled at 1.9 K in 3 km long strings:  
how to transport the heat over such distances?





# Conduction cooling of accelerator string: the fin problem



Distance : \$z\$  
Linear heat load: \$\xi\$

G-M law applied on \$dz\$:  $\dot{q}^{3.4} \cdot dz = dX$

Energy conservation:  $d\dot{q} = \frac{\xi}{A} \cdot dz \Rightarrow dz = \frac{A}{\xi} \cdot d\dot{q}$

then,  $\dot{q}^{3.4} \cdot d\dot{q} = \frac{\xi}{A} \cdot dX \Rightarrow \dot{q}_{tot}^{4.4} = 4.4 \cdot \frac{\xi}{A} \cdot X$

Calling \$Q\_{tot}\$ the total heat load on the string of length \$L\$:

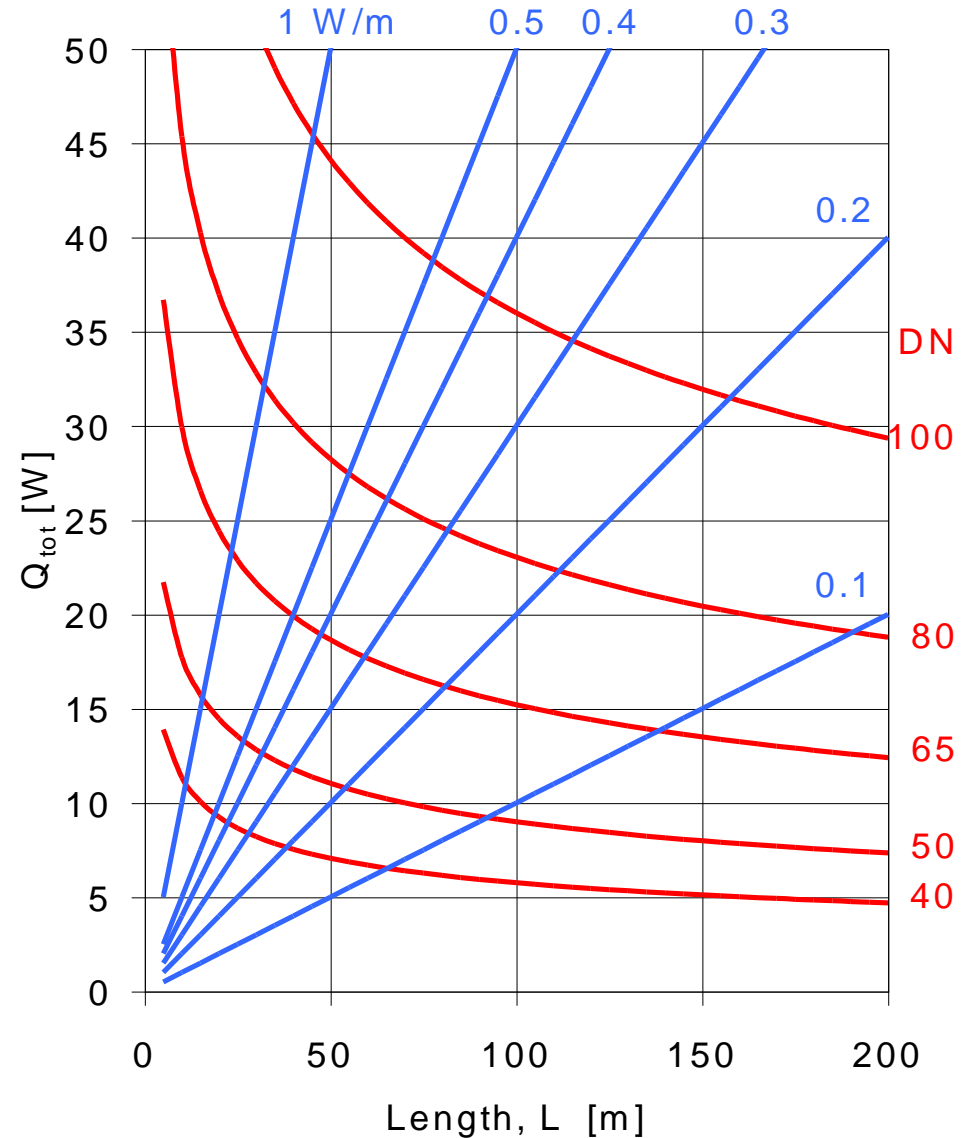
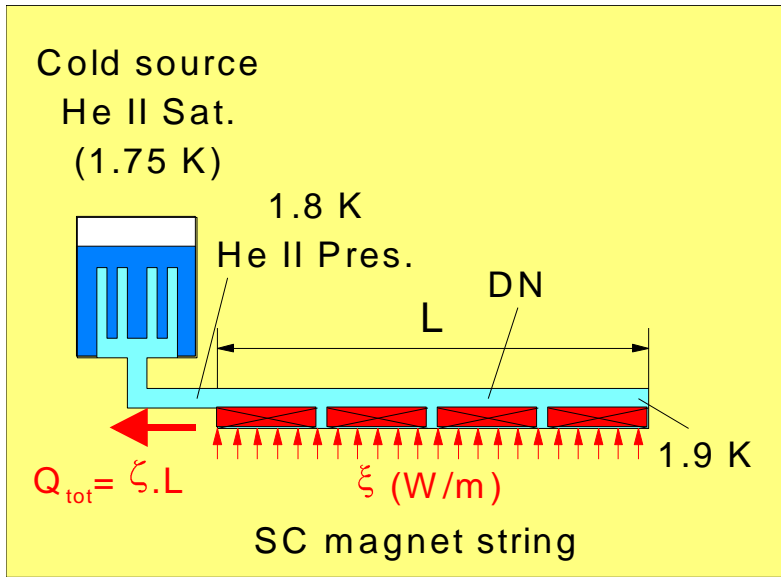
$$\dot{Q}_{tot} = L \cdot \xi = \dot{q}_{tot} \cdot A \Rightarrow \frac{\xi}{A} = \frac{\dot{q}_{tot}}{L}$$

Hence,  $\dot{q}_{tot}^{3.4} \cdot L = 4.4 \cdot X$

# Conduction in He II with linear heat load

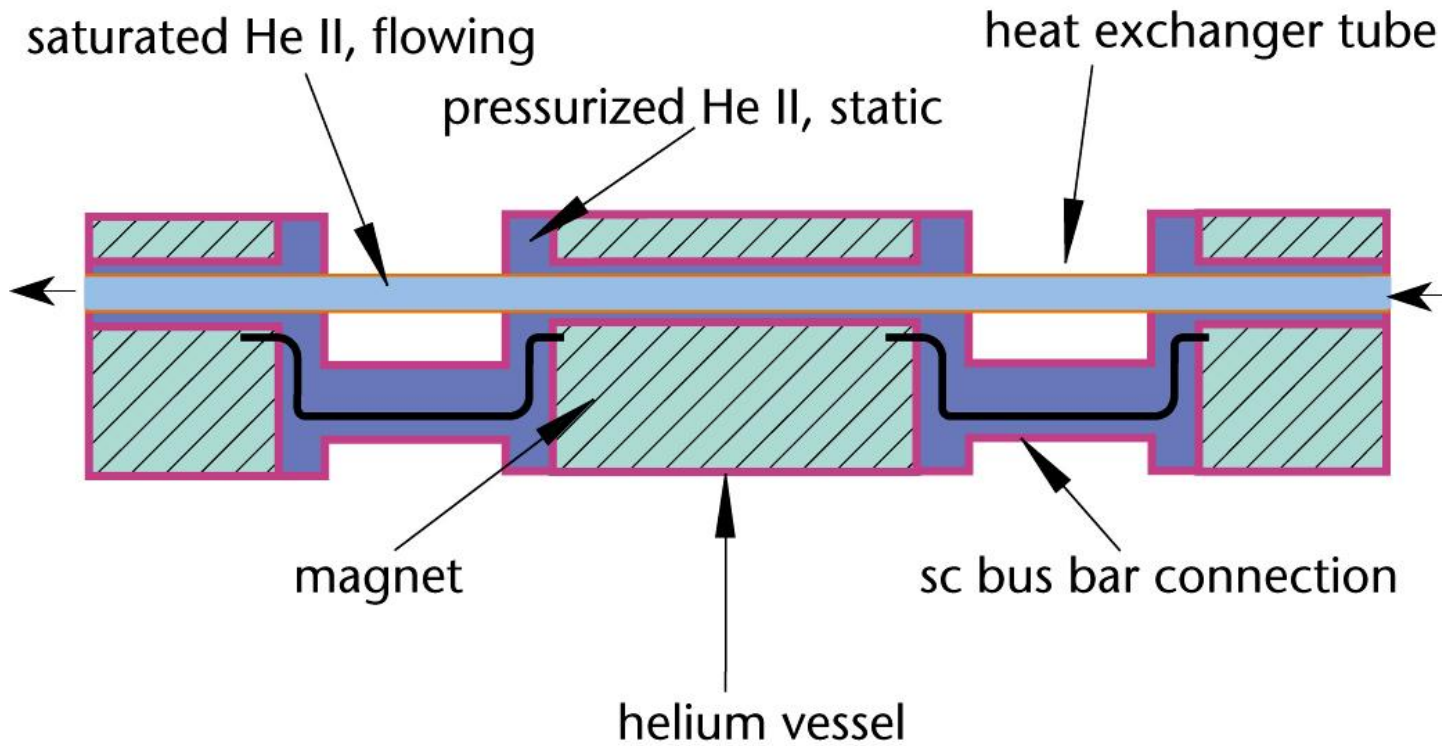


From 1.80 K to 1.90 K





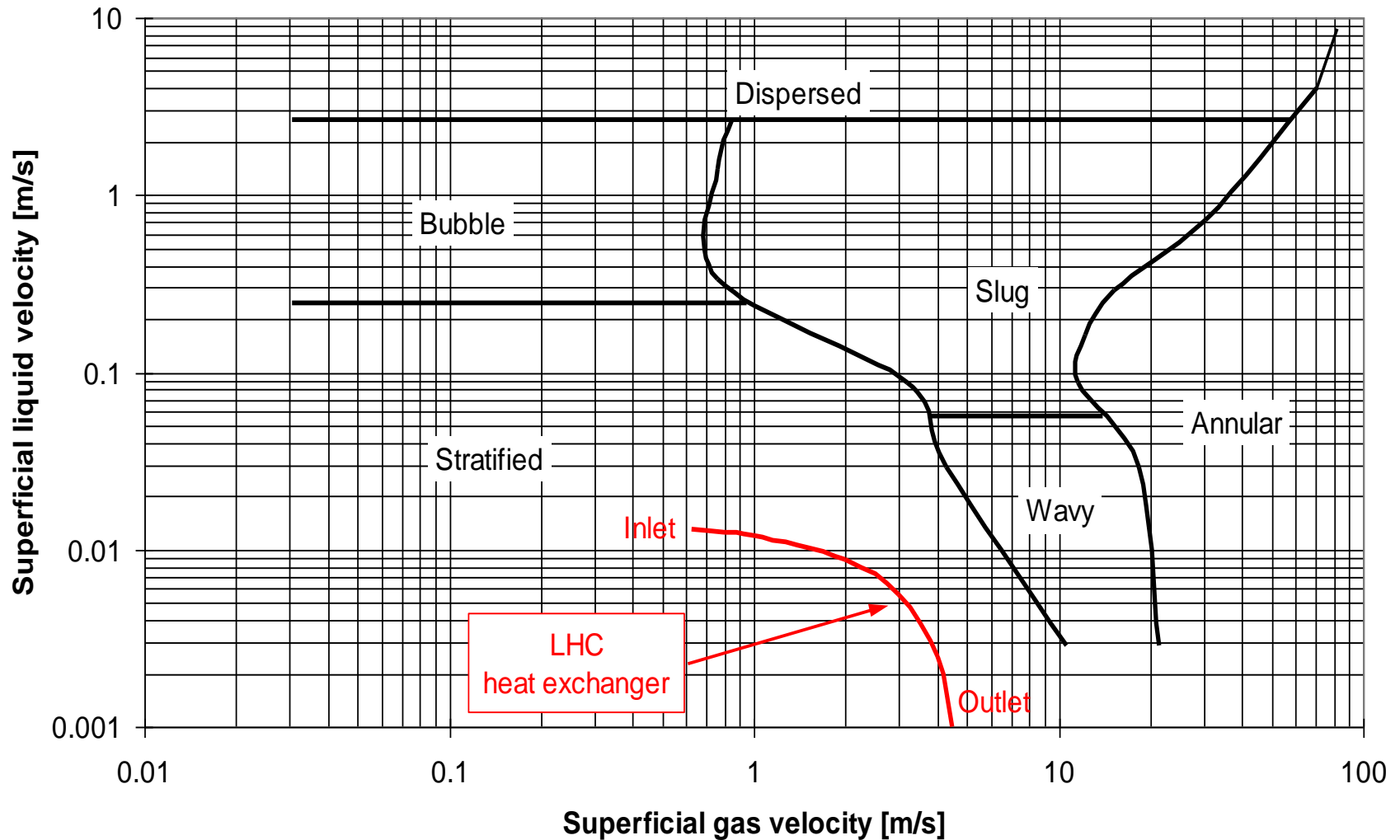
# Cooling by two-phase flow of He II s of the LHC magnet strings



- Heat exchanger tube in copper with a diameter DN50
- Overall thermal conductance:  $\sim 100 \text{ W/m.K}$   
(i.e., for  $1 \text{ W/m}$ , a temperature difference of  $10 \text{ mK}$ )

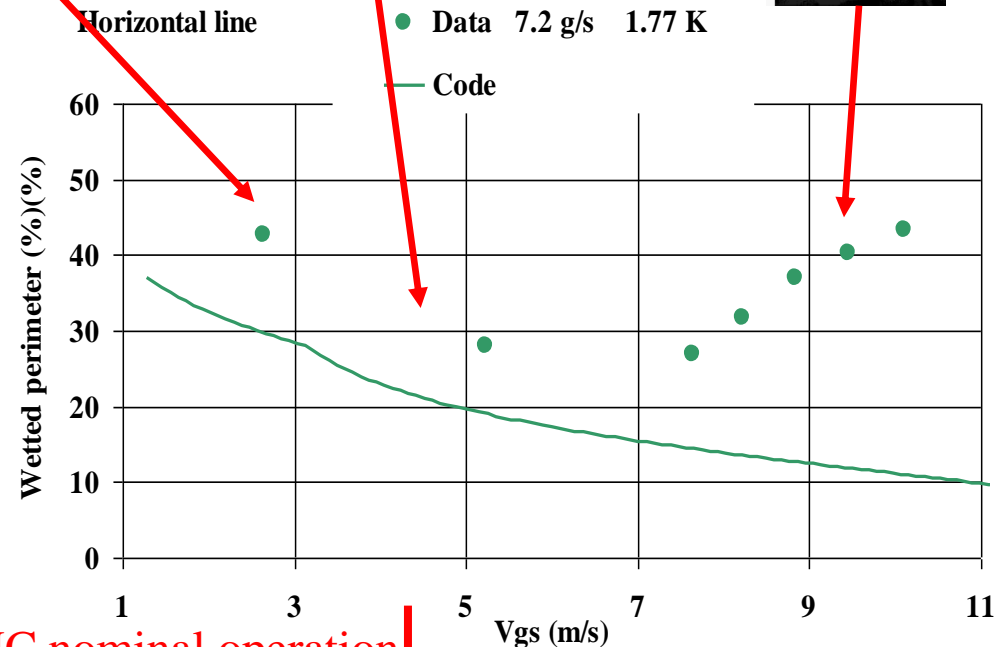
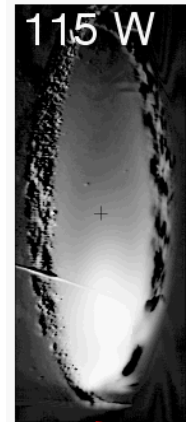
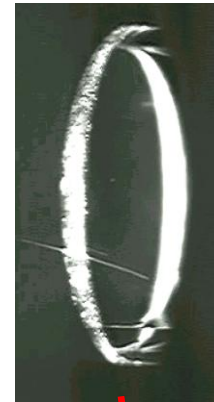
# Two-phase Flow of Saturated He II

(Mandhane, Gregory & Aziz flow map)



# Investigation of two-phase He II s flow (CEA Grenoble, France)

- Phenomenology
- Stability
- Pressure drop
- Heat transfer at wall

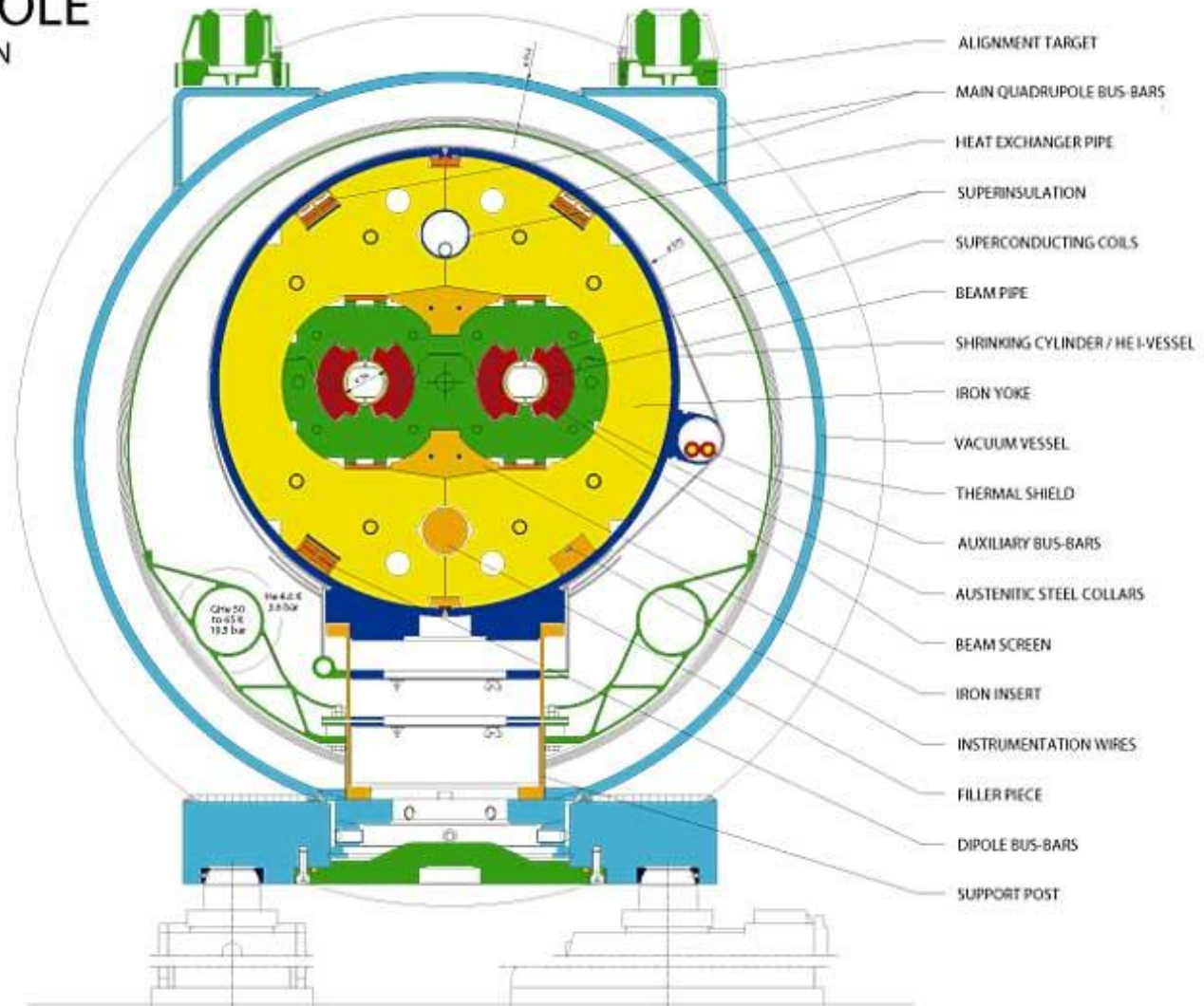


LHC nominal operation

# LHC magnet cross-section

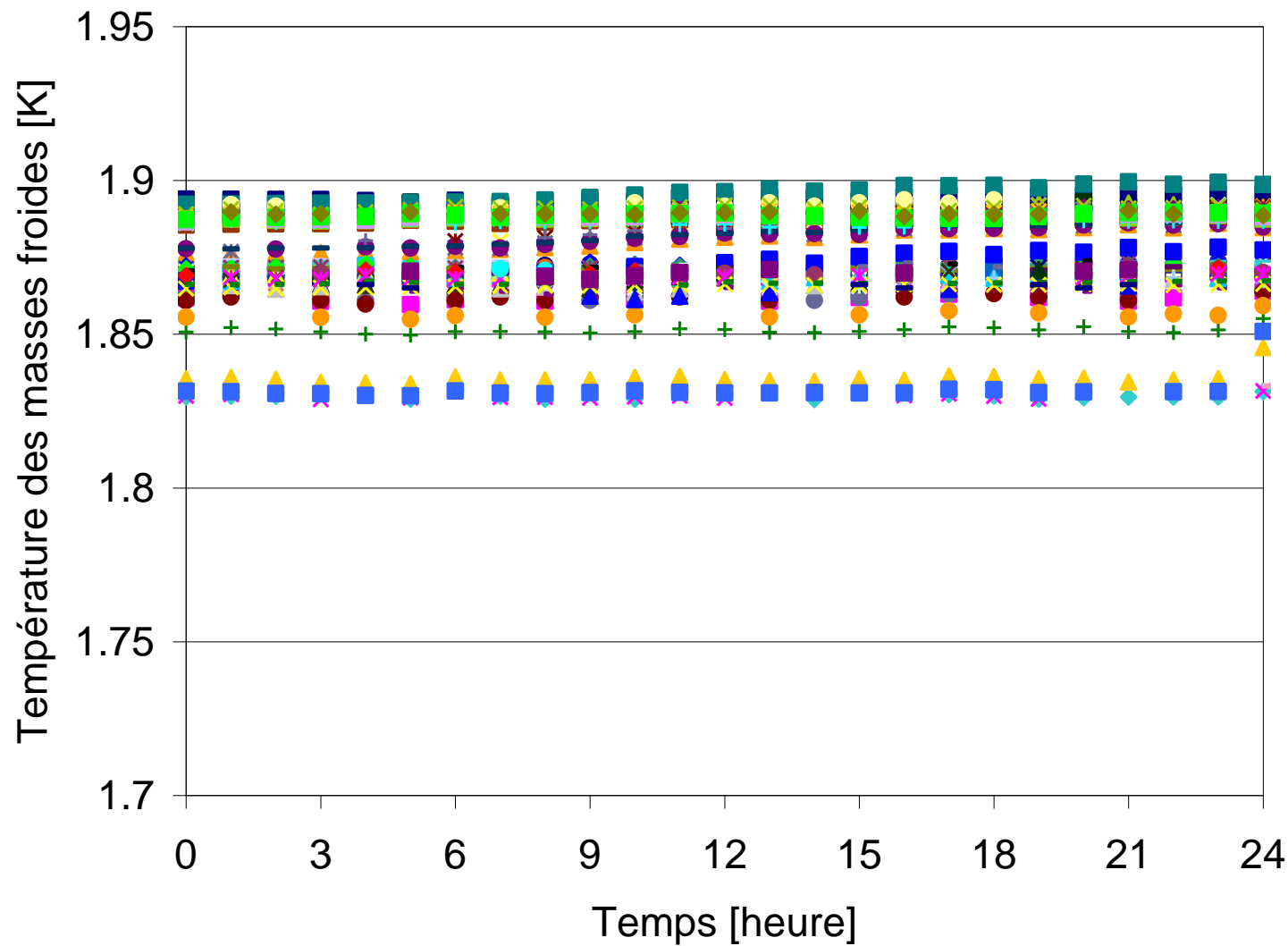


## LHC DIPOLE CROSS SECTION





# Temperature stability along LHC sector

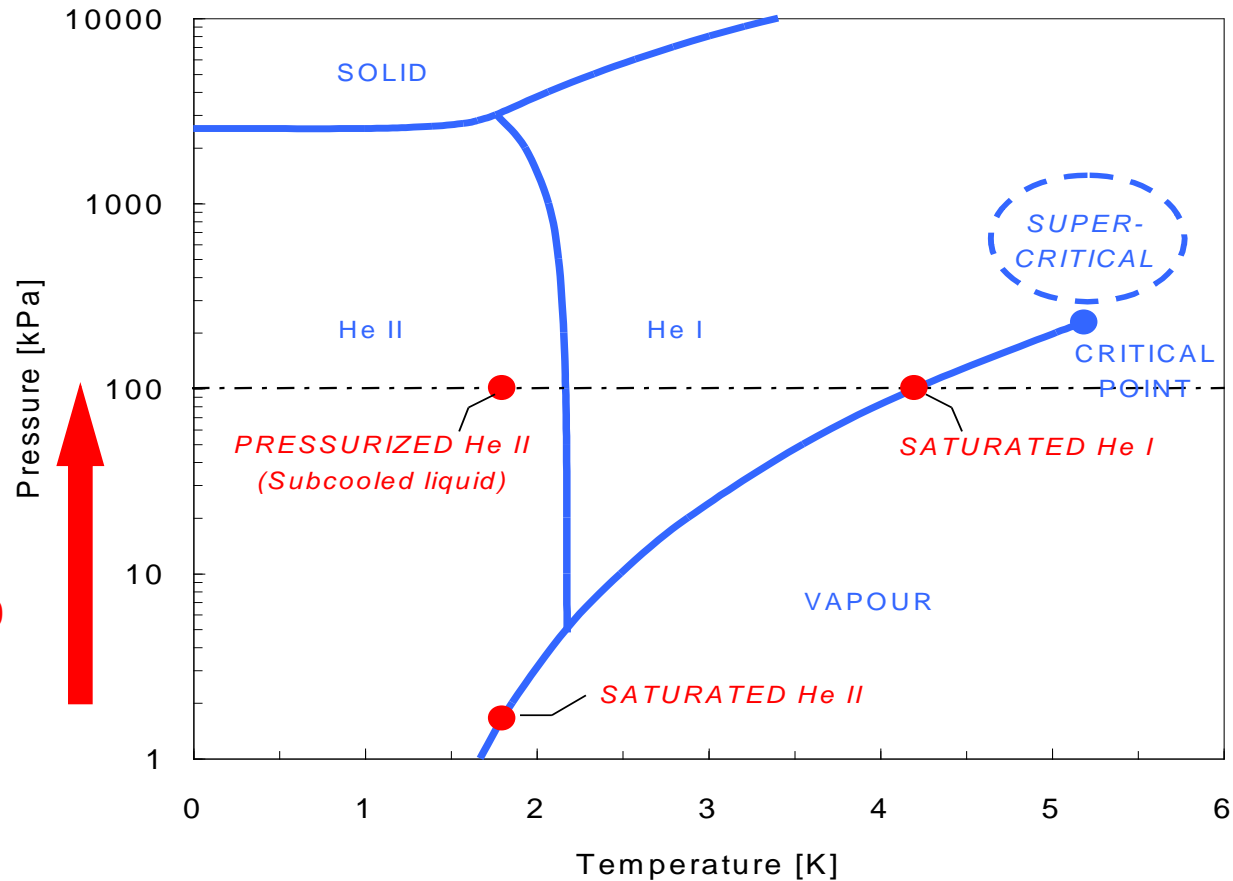




# Contents

- Introduction to superfluid helium
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# Challenges of power refrigeration at 1.8 K



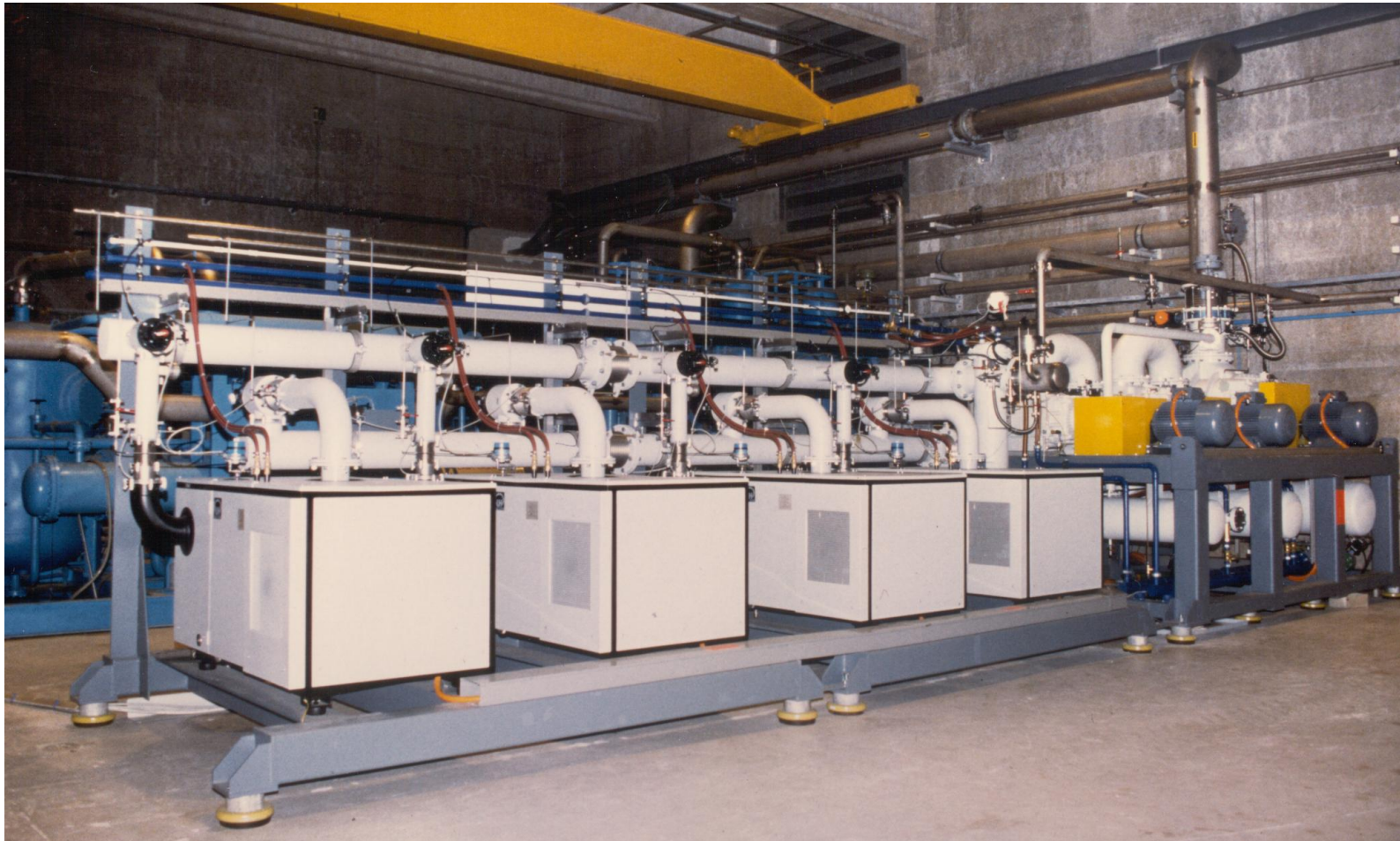
- Compression of large mass flow-rate of He vapor across high pressure ratio  $\Rightarrow$  intake He at maximum density, i.e. cold
- Need contact-less, vane-less machine  $\Rightarrow$  hydrodynamic compressor
- Compression heat rejected at low temperature  $\Rightarrow$  thermodynamic efficiency



# Warm pumping unit for LHC magnet tests

3 stages of Roots + 1 stage rotary-vane pumps

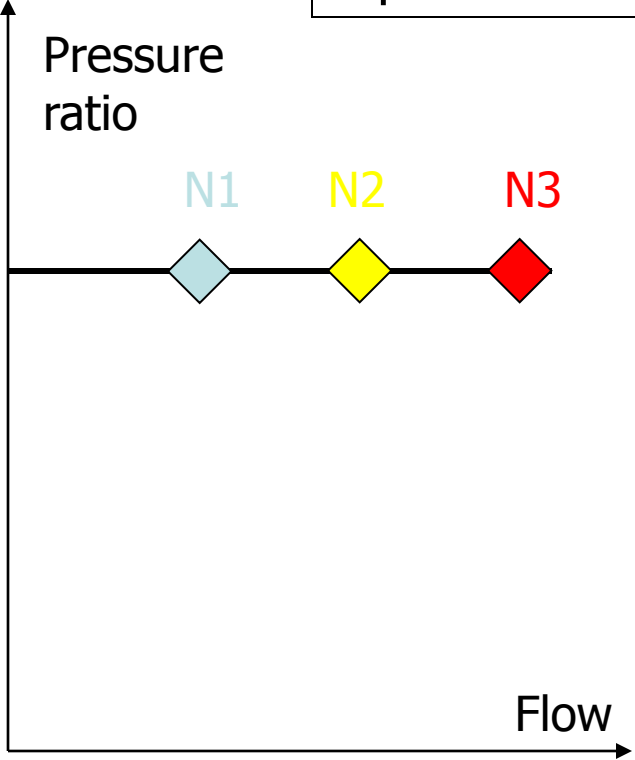
6 g/s @ 10 mbar (*1/160 of LHC!*)



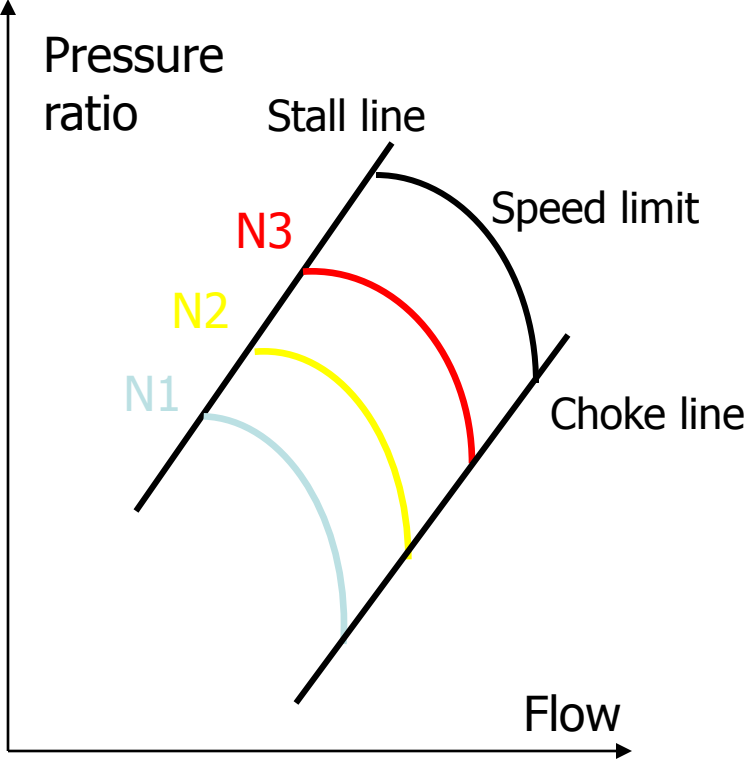


# Operating ranges of volumetric & hydrodynamic compressors

Speed  $N1 < N2 < N3$

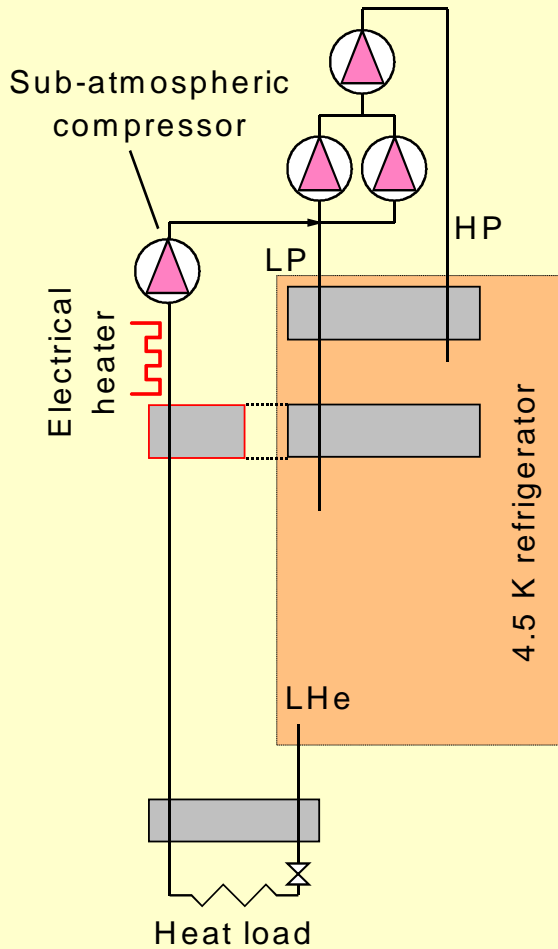


Volumetric (ideal)

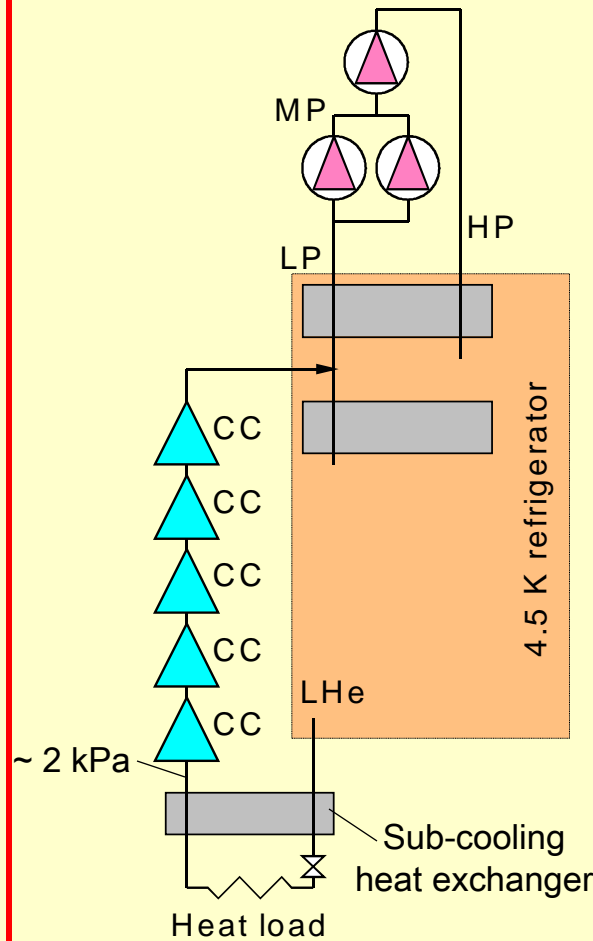


Hydrodynamic

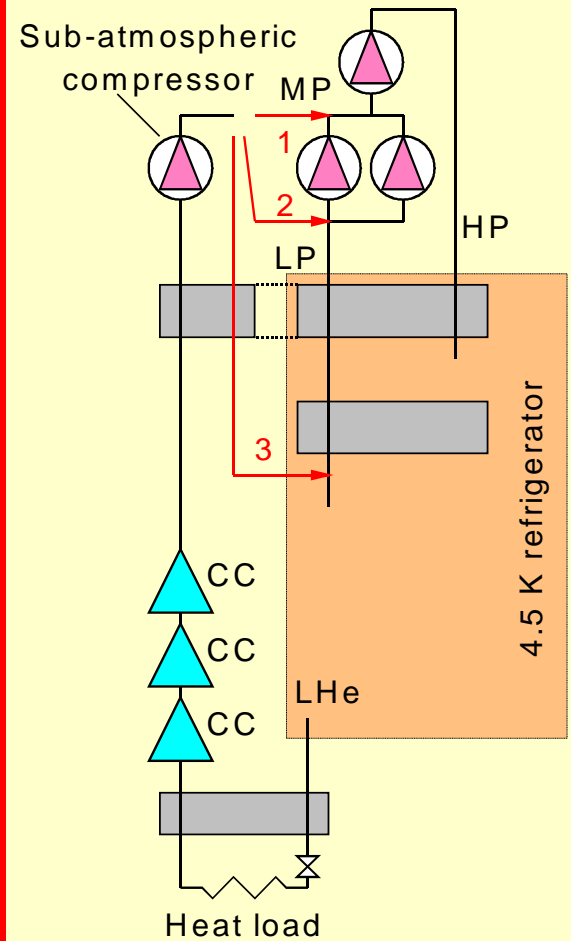
# Cycles for refrigeration below 2 K



**“Warm”  
Compression**

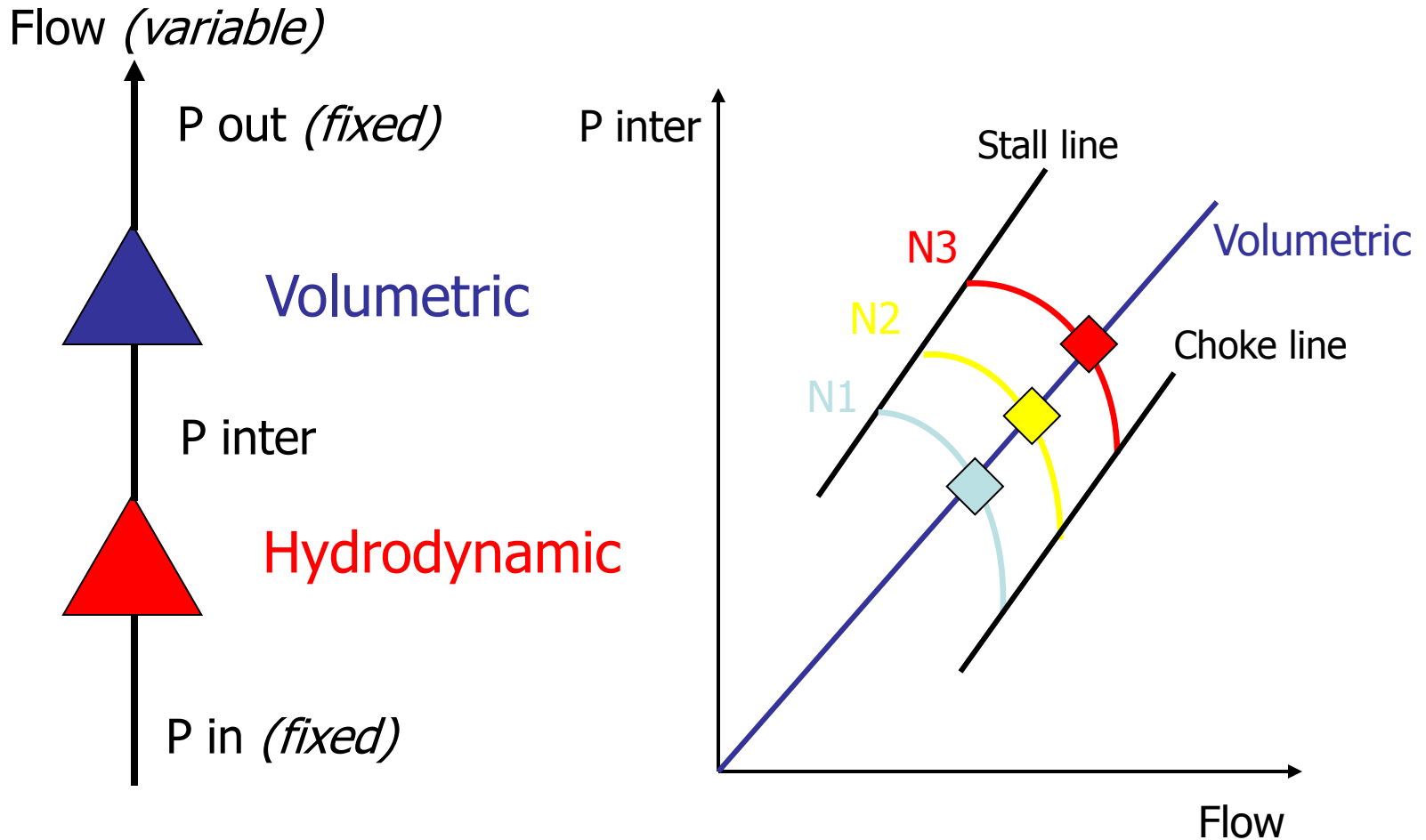


**“Integral Cold”  
Compression**



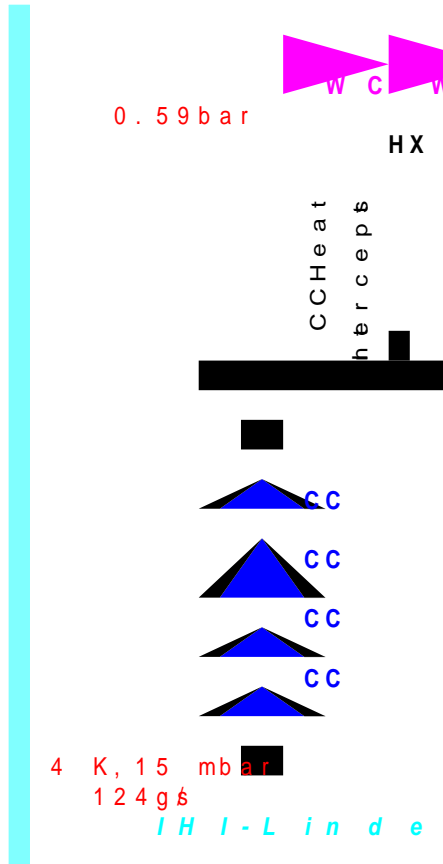
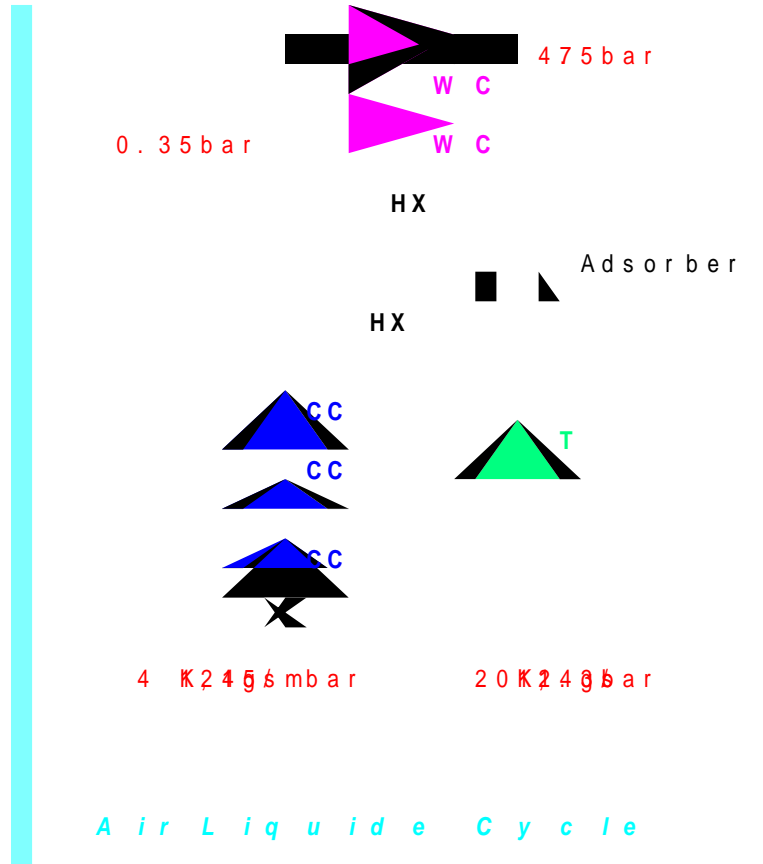
**“Mixed”  
Compression**

# Flow compliance of "mixed" compression



For fixed overall inlet & outlet conditions, coupling of the two machines *via*  $P_{inter}$  maintains the operating point in the allowed range

# Simplified flow-schemes of the 1.8 K refrigeration units of LHC



## The four-stage LHC cold compressors



1<sup>st</sup> stage

IHI





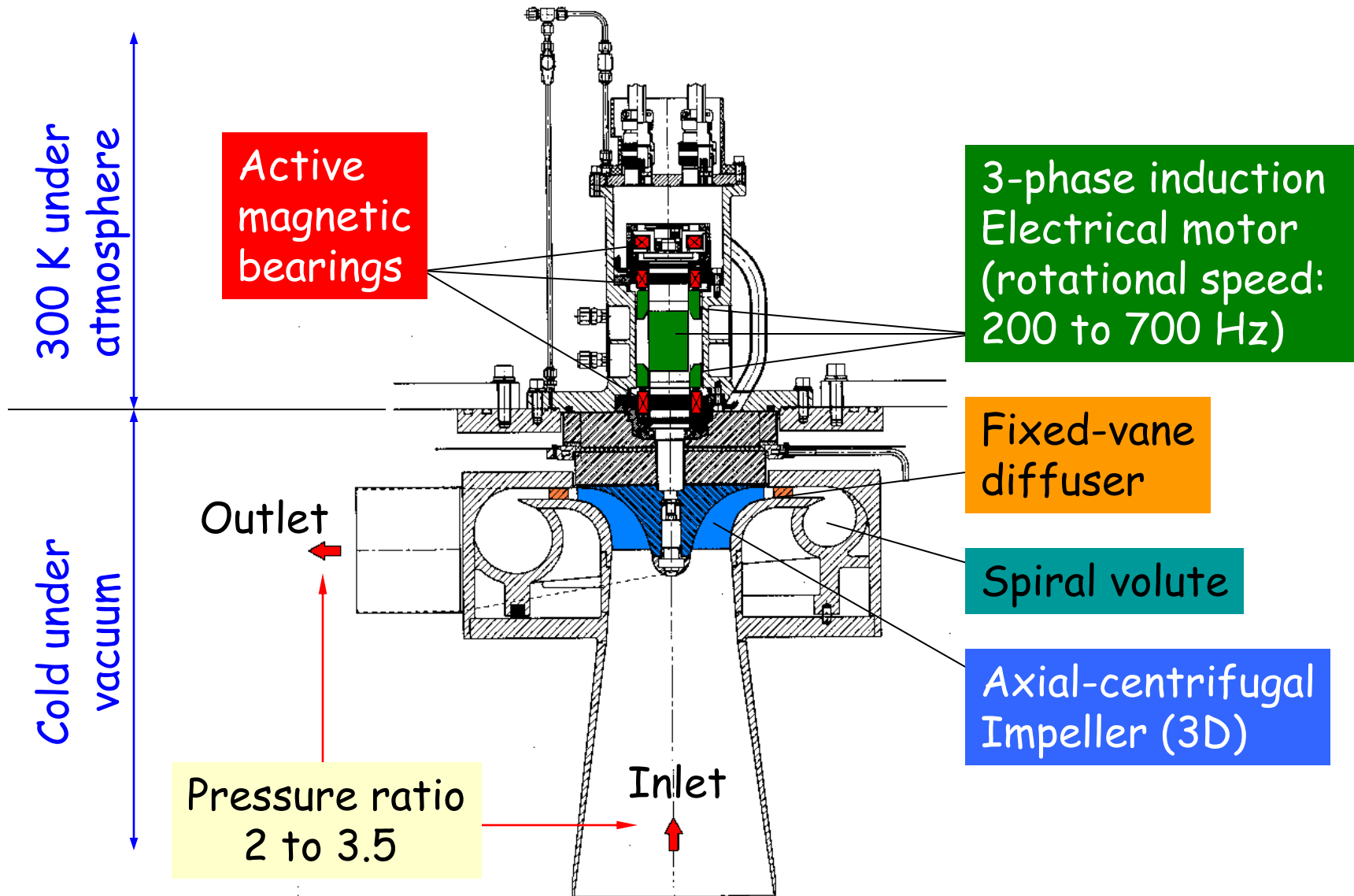
1<sup>st</sup> stage



3<sup>rd</sup> stage

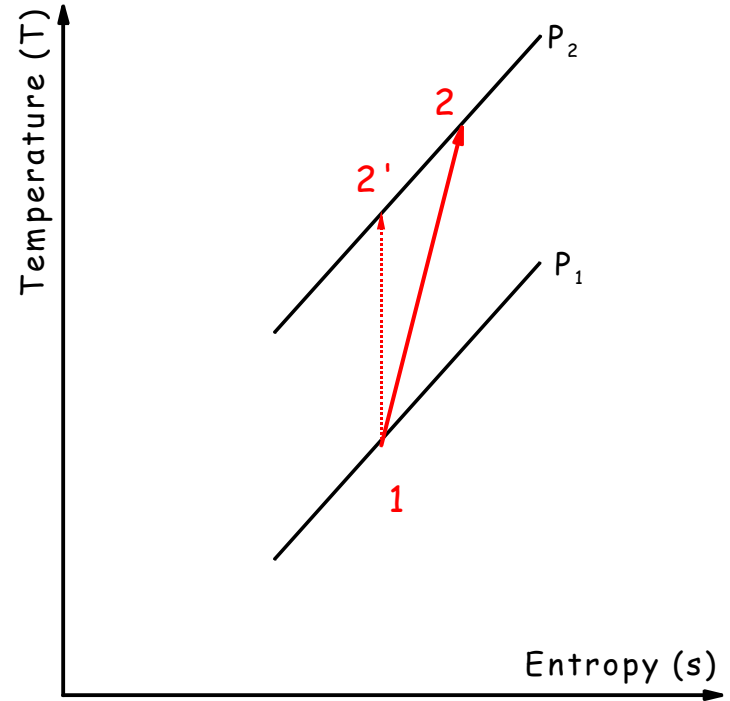
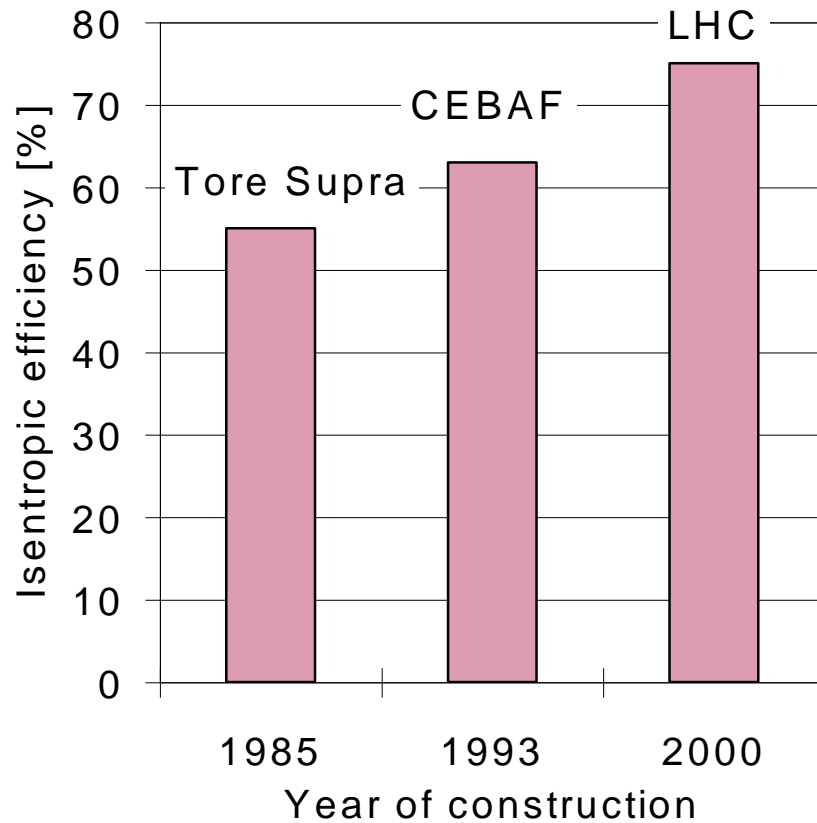
Air Liquide

# Specific features of LHC cold compressors





# Performance of LHC cold compressors



$$\eta_{is} = \frac{H_{2'} - H_1}{H_2 - H_1}$$



Compound two-stage screw compressor

Mycom

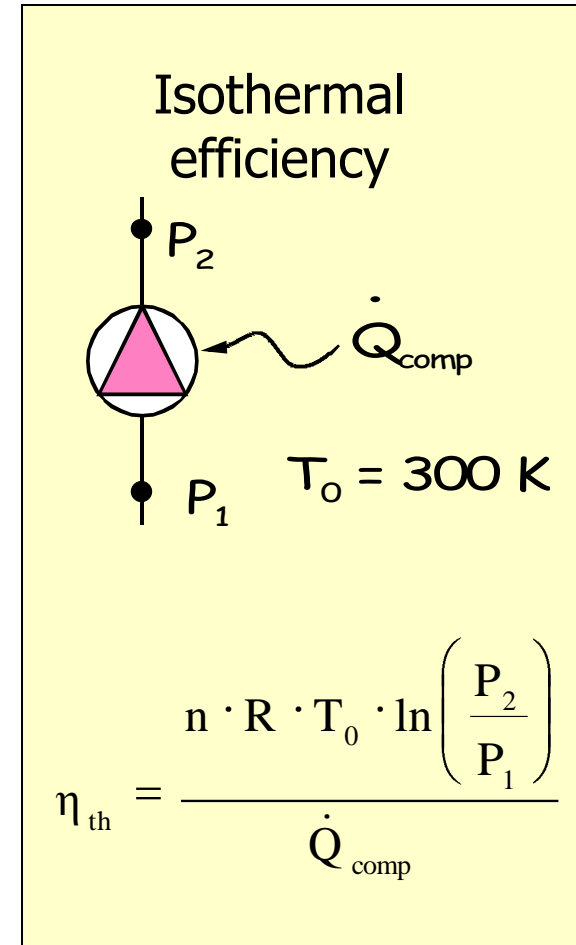
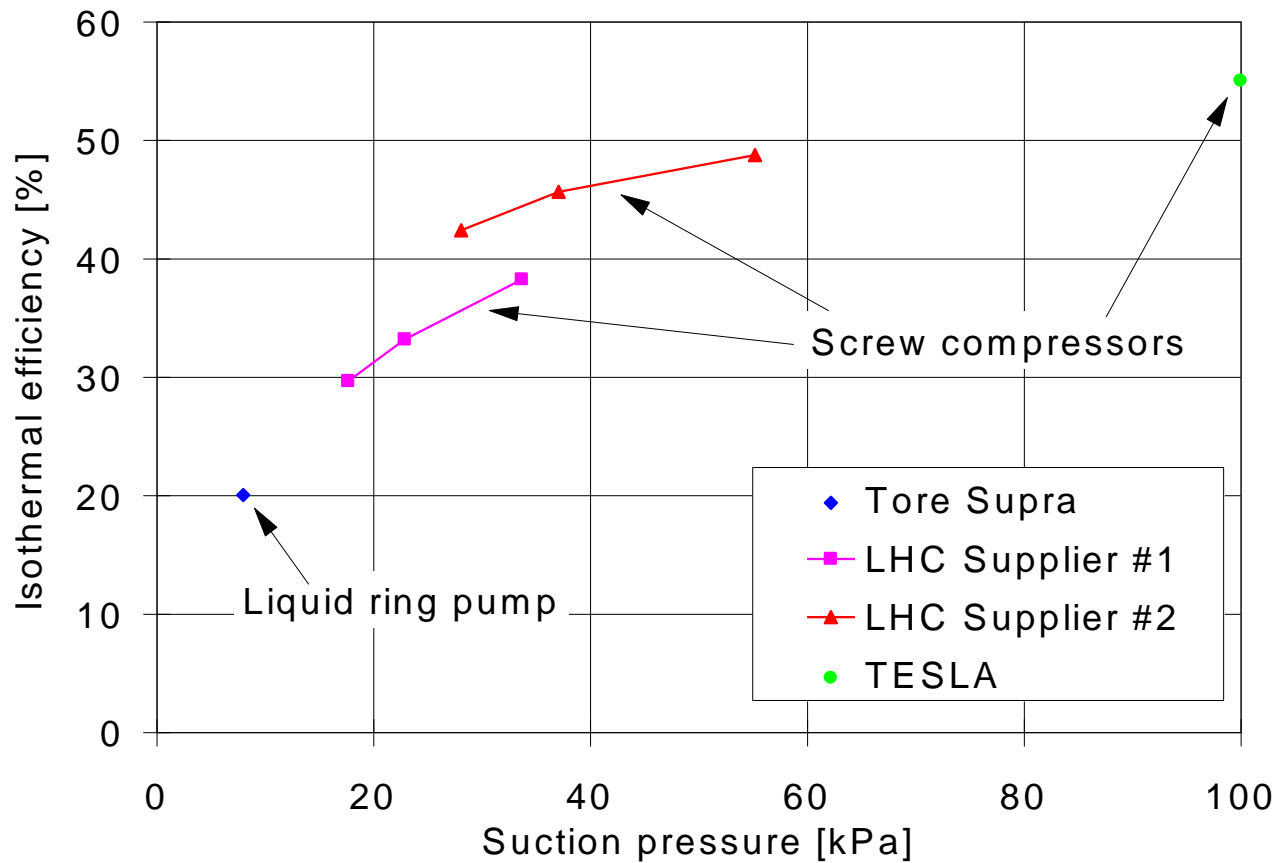
WCS at CERN:  
125 g/s @ 0.6 bar  
or  
4600 m<sup>3</sup>/h @ 15 °C





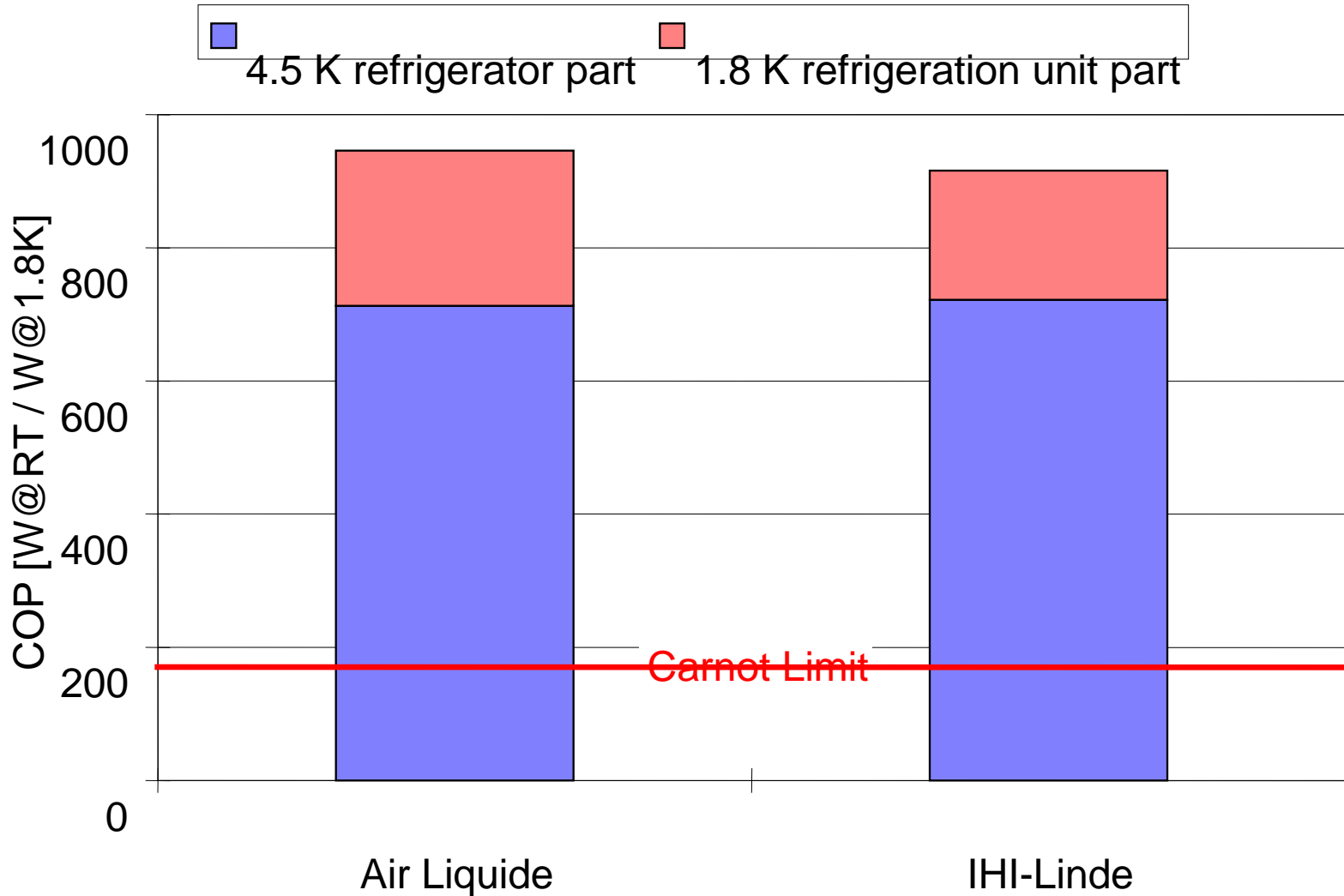
WCS at CERN:  
125 g/s @ 0.35 bar  
or  
2 x 3900 m<sup>3</sup>/h @ 15 °C

# Isothermal efficiency of warm subatmospheric compressors



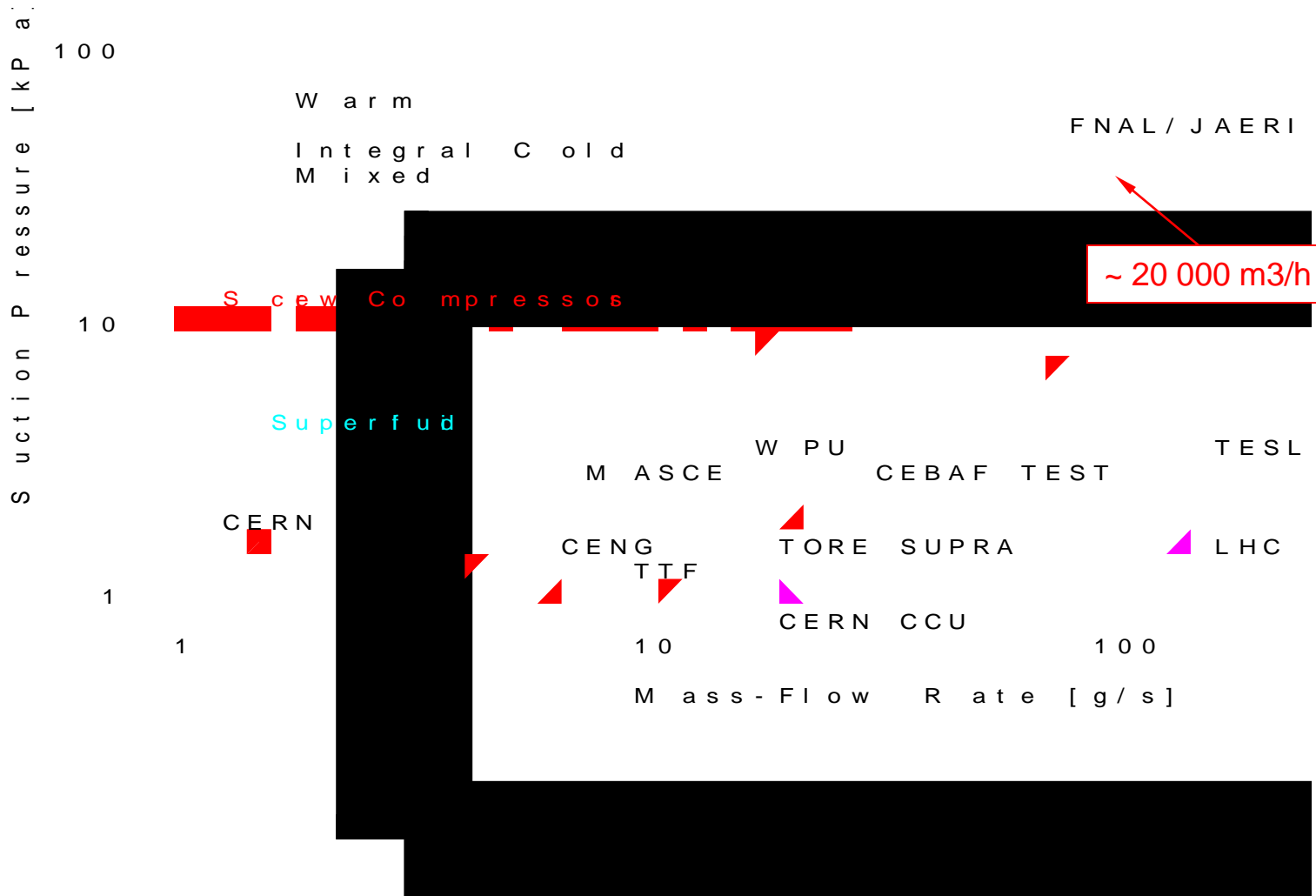


# C.O.P. of LHC 1.8 K units

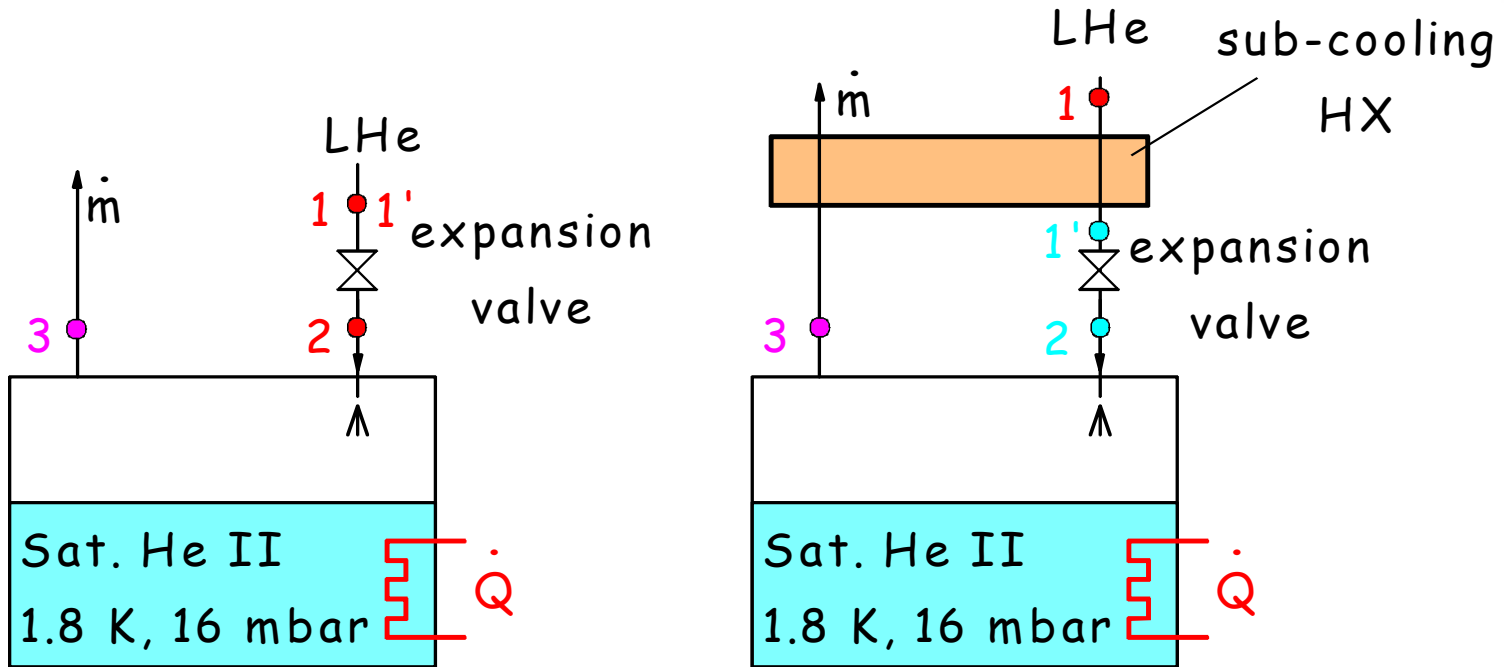




# Application range of low-pressure He compressors



# Efficiency of Joule-Thomson expansion



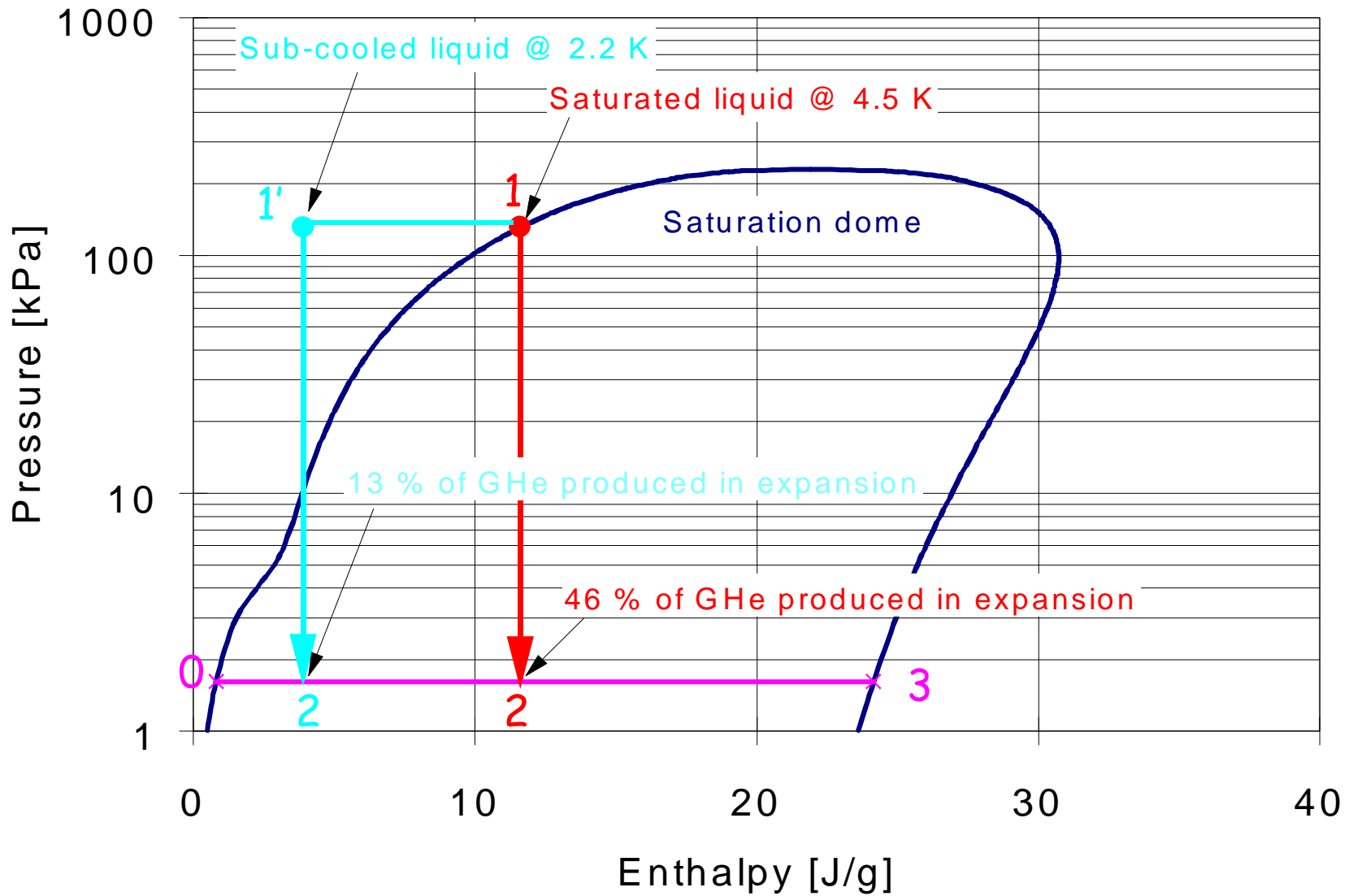
Sub-cooling efficiency :

$$\eta_{sc} = \frac{H_3 - H_2}{H_3 - H_0}$$

Sub-cooling	$T_{1'}$ [K]	$H_3 - H_2$ [J/g]	$H_3 - H_0$ [J/g]	$\eta_{sc}$ [%]
without	4.5	12.6	23.4	54
with	2.2	20.4	23.4	87

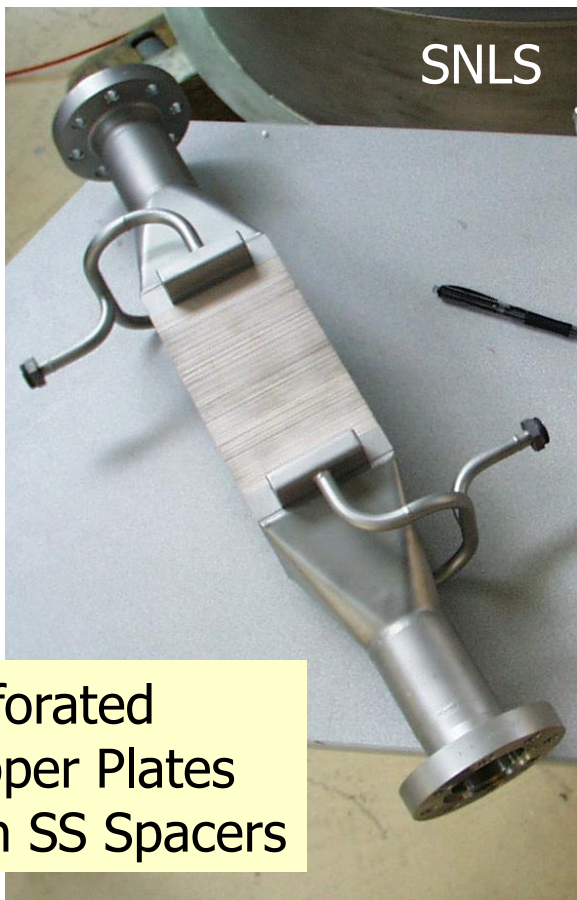


# Subcooling before J-T expansion



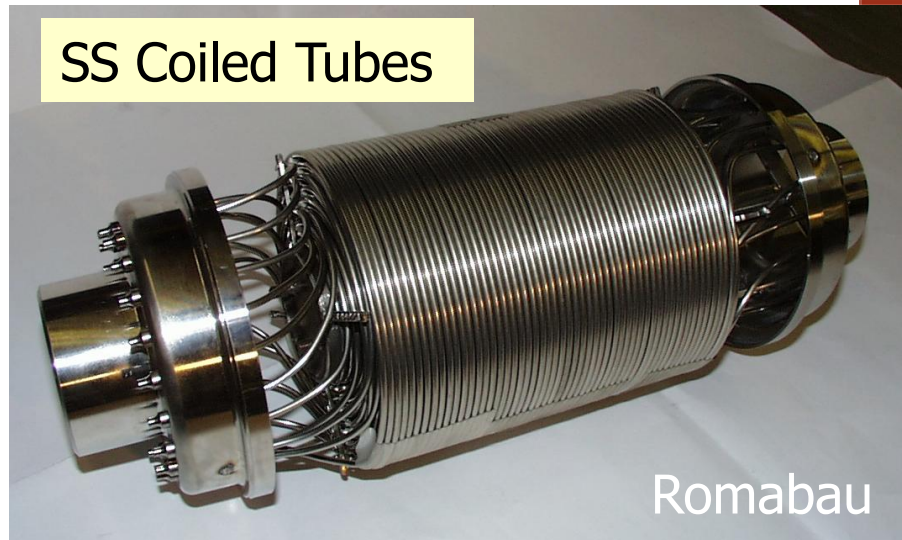
# Prototypes of subcooling HX for LHC

Mass-flow: 4.5 g/s  
 $\Delta P$  VLP stream: < 1 mbar  
Sub-cooling T: < 2.2 K



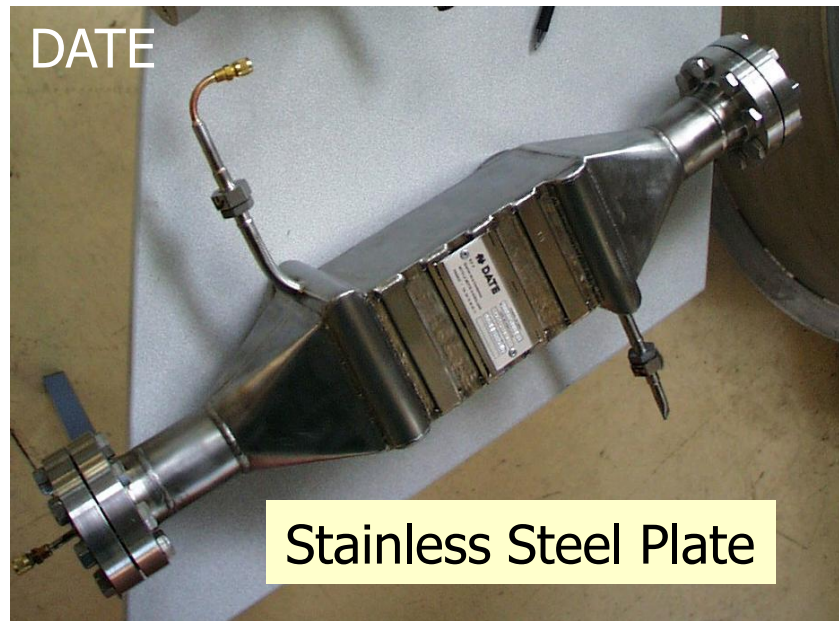
SNLS

Perforated  
Copper Plates  
with SS Spacers



SS Coiled Tubes

Romabau



DATE

Stainless Steel Plate



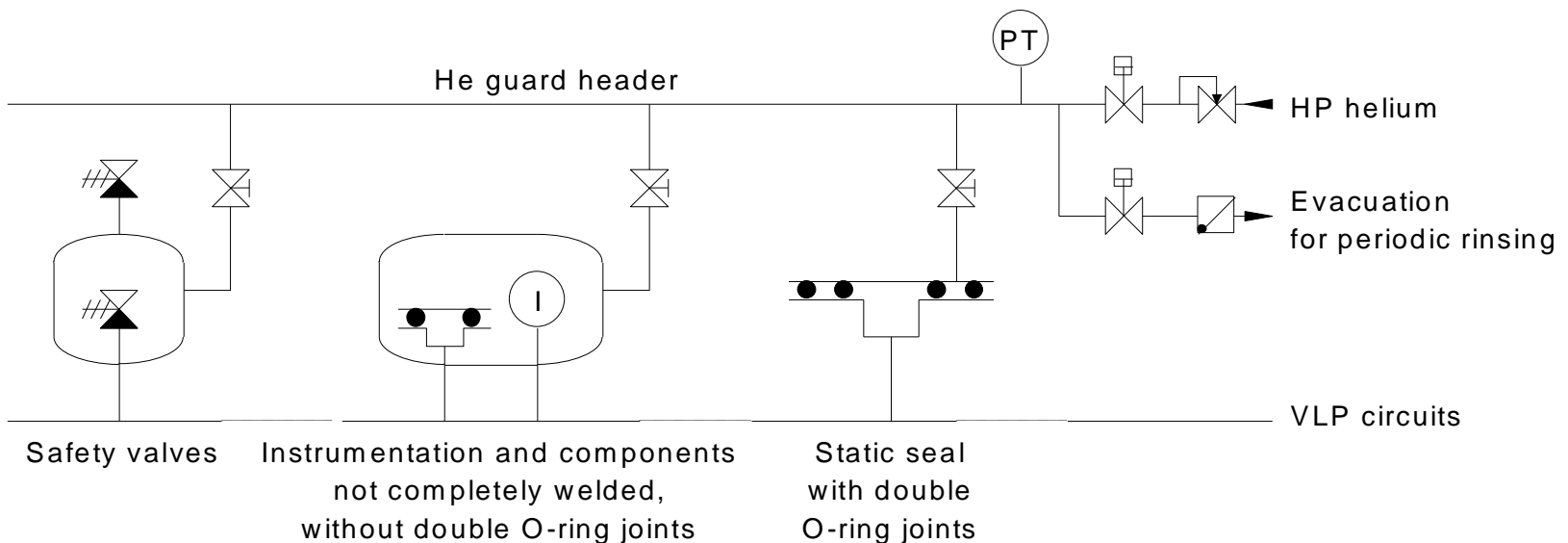
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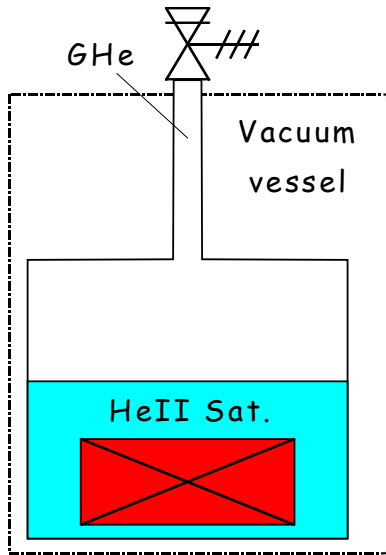
# Protection against air inleaks



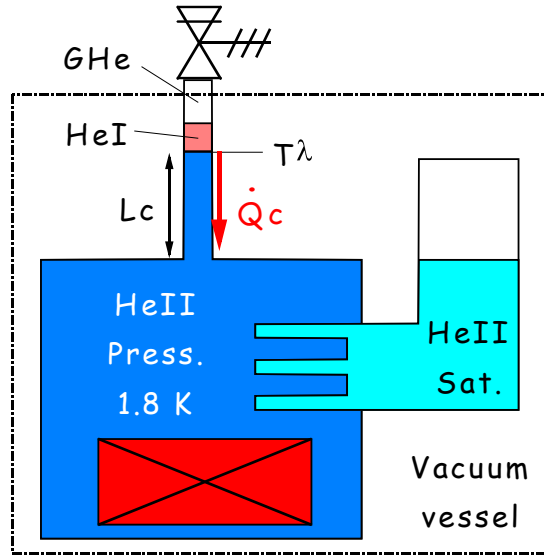
- Motor shaft of warm sub-atmospheric compressors placed at the discharge side to work above atmospheric pressure.
- For sub-atmospheric circuits which are not under vacuum or not completely welded, helium guard protection on dynamic seal of valves, instrumentation, safety relief valves and critical static seal



# Protection of He II p enclosure



"Standard" Protection of saturated HeII bath with "warm" safety valve

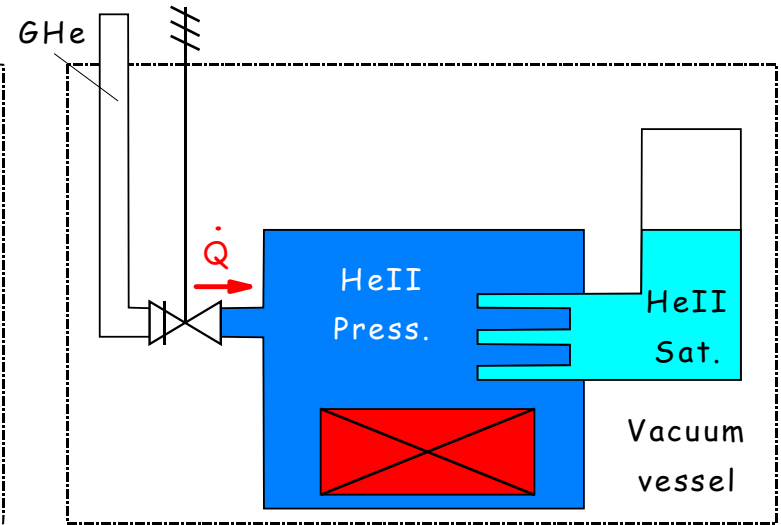


"Standard" Protection of pressurized HeII bath → Critical heat flux in safety valve piping:

$$L_c = 1 \text{ m} \rightarrow \dot{q}_c = 1.5 \text{ W/cm}^2$$

$$\text{DN40 Piping} \rightarrow A = 12 \text{ cm}^2$$

$$\dot{Q}_c = 18 \text{ W} !$$



Protection of pressurized HeII with a cold safety relief valve:

$$\dot{Q} \approx 0.35 \text{ W}$$

(350 valves in LHC, i.e. 123 W instead of 6300 W with "standard" protection)

# Specifications for LHC He II Safety Valves

## Valve Characteristics

Number of valves:	<b>362</b>
Valve type:	<b>direct loaded proportional safety valve</b>
Valve sizing:	<b>DN40</b>
Discharge capacity:	<b>Kv &gt; 30 m<sup>3</sup>/h</b>

## Pressure relief

Set pressure:	<b>17 bar</b>
Fully open pressure:	<b>20 bar</b>
Time to full open:	<b>&lt; 350 ms @ 10 bar/s</b>

## Thermal & leak

	<i>Specified</i>	<i>Measured</i>
Thermal anchoring heat load @ 100K:	<b>&lt; 3 W</b>	<b>3.5 W</b>
Conduction heat load to 1.9K:	<b>&lt; 0.35 W</b>	<b>0.38 W</b>
Mass leakage at 1.9K and $\Delta P = 100$ mbar :	<b>&lt; 0.01 g/s</b>	<b>0.0007 g/s</b>



Series safety relief valve



## Conclusion and outlook

- From a laboratory curiosity and a hot research topic in condensed-matter physics, superfluid helium has become a state-of-the-art cryogen for cooling large superconducting devices such as high-energy accelerators, tokamaks and research magnets
- Projects such as TORE SUPRA, CEBAF and LHC have triggered vigorous development programmes in laboratories and industry concerning flow and heat transfer, refrigeration techniques, instrumentation and engineering
- Superfluid helium remains an enabling technology for NMR magnets and future large projects using high-field superconducting devices, e.g. the European X-FEL, ILC. The unique hydrodynamic properties of the fluid can also be used *per se*, e.g. in turbulence research

# High-field magnets fo high-frequency NMR





# The International Linear Collider (ILC)



$e^+ e^-$  linear collider

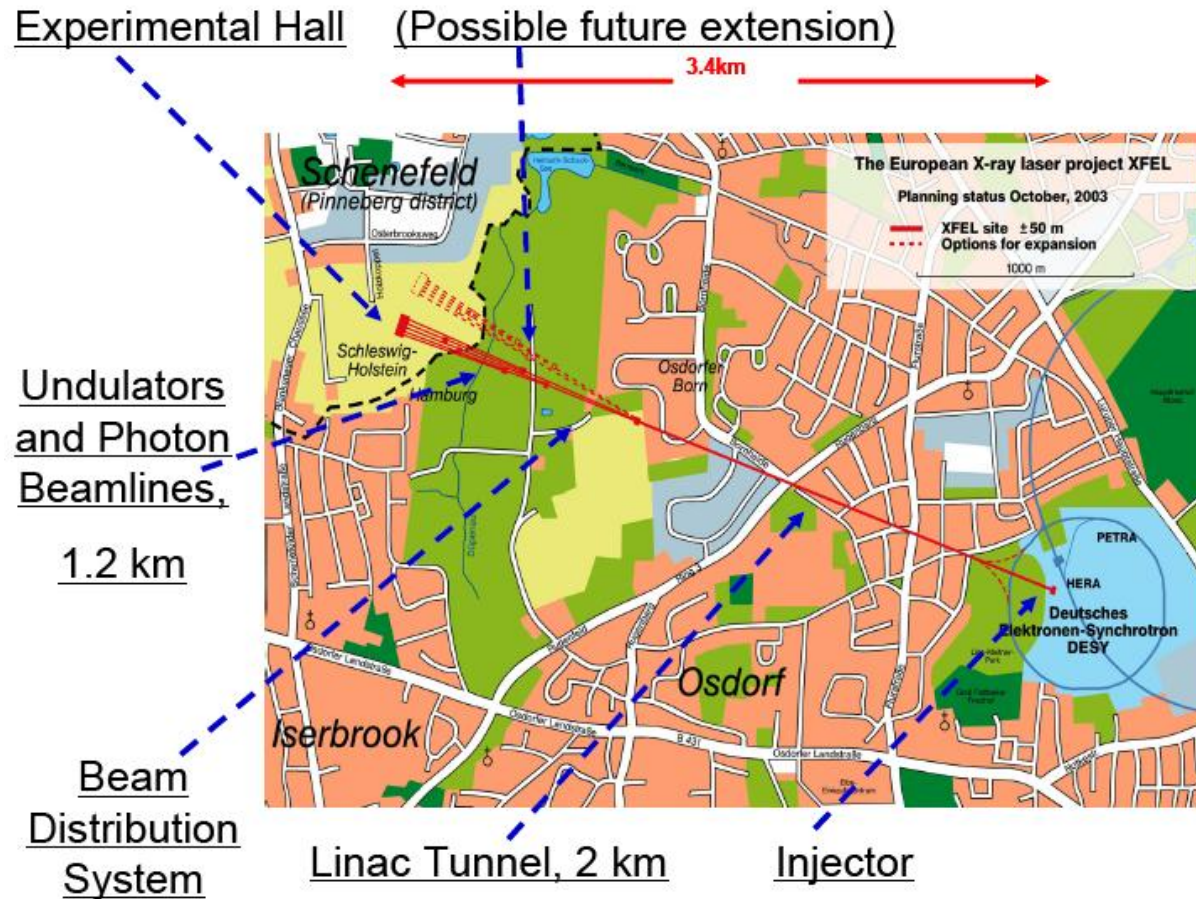
Collision energy 500 GeV c.m.  
initially, later upgrade to 1 TeV  
c.m.

Overall length 47 km, of which  
22 km linacs

16'000 superconducting RF  
cavities operated at 2 K

Refrigeration 45 kW @ 2 K

# The European X-ray FEL, DESY, Hamburg (Germany)



Very brilliant, ultra-short (100 fs) pulses of X-rays down to 0.1 nm

Based on s.c. e linac

Beam energy 17.5 GeV

Beam power 600 kW

Linac length 1.7 km

928 superconducting RF cavities operated at 2 K

Refrigeration 2.5 kW @ 2 K



## Reference

- Ph. Lebrun & L. Tavian, *The technology of superfluid helium* in *Cryogenics for particle accelerators & detectors*, CERN-LHC-2002-011 (2002) pp. 71-92

This write-up contains a comprehensive list of references and bibliography on superfluid helium technology