

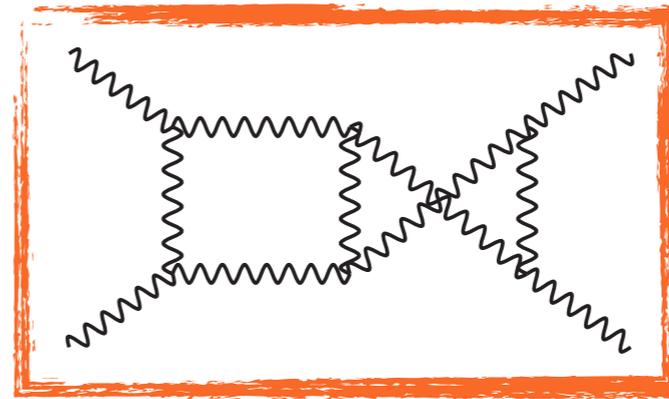
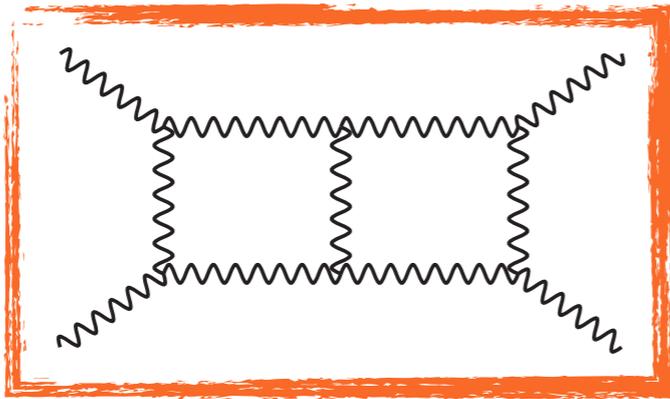
THE TWO-LOOP FOUR-GRAVITON SCATTERING AMPLITUDES

ZOOMPLITUDES 2020

Samuel Abreu — UCLouvain

arXiv:2002.12374, to appear in PRL

with: F. Febres Cordero, H. Ita, M. Jaquier, B. Page, M. Ruf,
V. Sotnikov



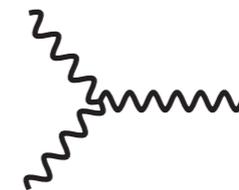
- ◆ Compute an important amplitude in pure GR:
 - ✓ All-plus known, compute remaining helicities [Bern et al.], [Dunbar et al.]
 - ✓ UV well studied, compute finite pieces [Weinberg], ['t Hooft and Veltman], [Goroff and Sagnotti], [Donoghue], [Bern et al]...
- ◆ Insights into relation between multi-loop amplitudes/classical gravitational dynamics
- ◆ Develop tools to compute GR amplitudes
- ◆ Explore two-loop numerical unitarity framework beyond planar QCD

$$\mathcal{L} = \mathcal{L}_{\text{EH}} + \mathcal{L}_{\text{GB}} + \mathcal{L}_{\text{R}^3}$$

Emil's talk

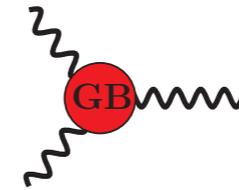
[Weinberg], ['t Hooft and Veltman], [Goroff and Sagnotti], [Donoghue], ...

$$\mathcal{L}_{\text{EH}} = -\frac{2}{\kappa^2} \sqrt{|g|} R$$



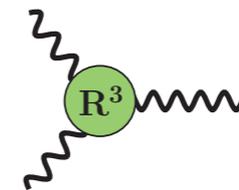
$$\mathcal{O}(\kappa)$$

$$\mathcal{L}_{\text{GB}} = \frac{\mathcal{C}_{\text{GB}}}{(4\pi)^2} \sqrt{|g|} (R^2 - 4R_{\mu\nu}R^{\mu\nu} + R_{\mu\nu\rho\sigma}R^{\mu\nu\rho\sigma})$$



$$\mathcal{O}(\kappa^3)$$

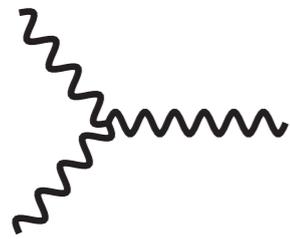
$$\mathcal{L}_{\text{R}^3} = \frac{\mathcal{C}_{\text{R}^3}}{(4\pi)^4} \left(\frac{\kappa}{2}\right)^2 \sqrt{|g|} R_{\alpha\beta}{}^{\mu\nu} R_{\mu\nu}{}^{\rho\sigma} R_{\rho\sigma}{}^{\alpha\beta}$$



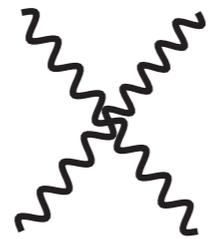
$$\mathcal{O}(\kappa^5)$$

$$\mathcal{A}^{(2)} = \text{[tree-level box diagram]} + \text{[R}^3 \text{ tree-level vertex]} + \text{[GB tree-level vertex]} + \text{[GB tree-level vertex]} \quad \mathcal{O}(\kappa^6)$$

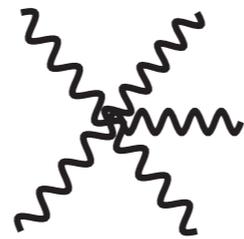
- ◆ EH Feynman rules are complicated



Terms: $\mathcal{O}(100)$



$\mathcal{O}(1000)$



$\mathcal{O}(10000)$

Feynman-diagram based calculation out of question
⇒ **(Generalised) Unitarity**

- ◆ EH interactions have high power-counting (QCD²)



Complicated integrand
⇒ **Analytics from numerics**

Issues addressed by
two-loop numerical unitarity



Caravel

[Abreu, Febres Cordero, Ita, Jaquier, Page, Zeng, 17,18,19]

- ◆ Computing the amplitude, $D = 4 - 2\epsilon$

$$\mathcal{A} = \sum_{\Gamma} \sum_{i \in M_{\Gamma}} c_{\Gamma,i}(D) I_{\Gamma,i}(D)$$

- ◆ All integrals (4-point 2-loop massless) are known

[Anastasiou, Smirnov, Tausk, Tejada-Yeomans, Veretin, 99,00]

- ◆ Decomposition of integrand

$$\mathcal{A}(\ell_l) = \sum_{\Gamma} \sum_{i \in M_{\Gamma} \cup S_{\Gamma}} c_{\Gamma,i}(D) \frac{m_{\Gamma,i}(\ell_l, D)}{\prod_{j \in P_{\Gamma}} \rho_j}$$

- ✓ Dimension $|M_{\Gamma} \cup S_{\Gamma}|$ simple to determine
- ✓ Construct surface S_{Γ} terms that integrate to zero
- ✓ Compute coefficients $c_{\Gamma,i}$ from generalised unitarity

Two-loop (numerical) unitarity — surface terms I

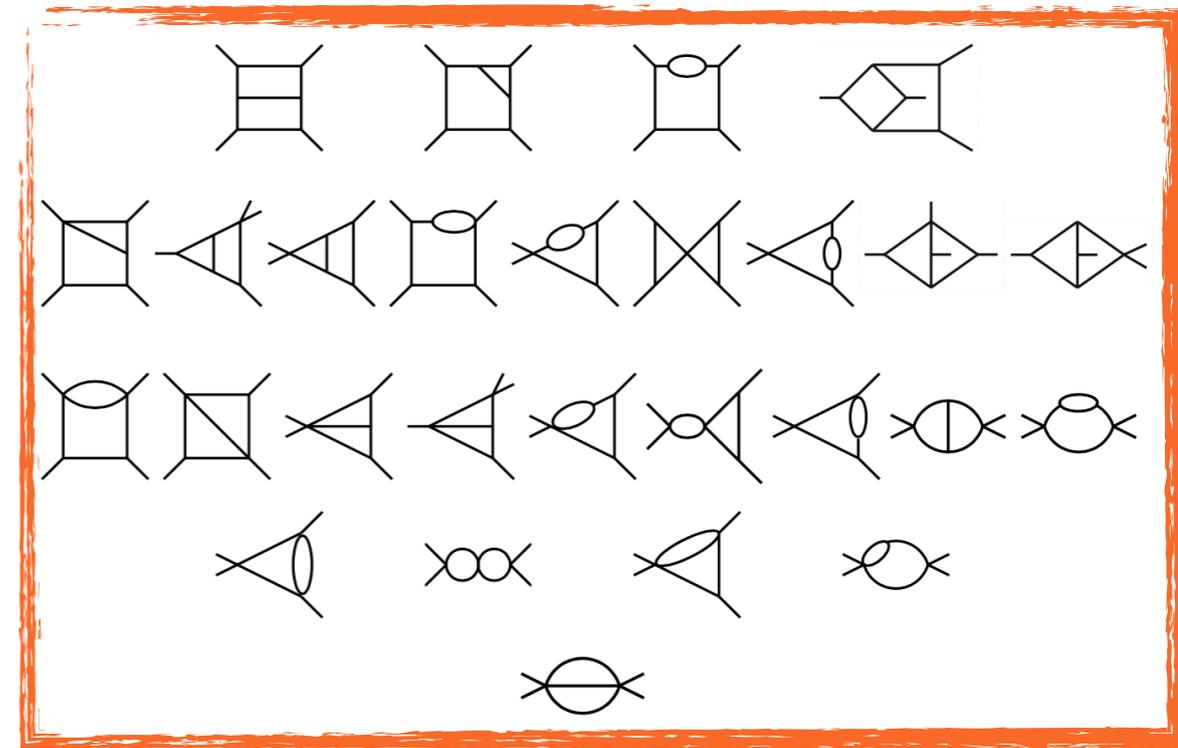
[Ita 15; Abreu, Febres Cordero, Ita, Jaquier, Page, Zeng 17]

- ◆ For each Γ , construct numerators such that

$$\int [d^D \ell_l] \frac{m_{\Gamma,i}(\ell_l, D)}{\prod_{j \in P_\Gamma} \rho_j} = 0$$

- ◆ “Integration-by-parts” IBP relations ...

$$\int [d^D \ell_l] \frac{\partial}{\partial \ell_i^\nu} \left[\frac{u_i^\nu}{\prod_{k \in P_\Gamma} \rho_k} \right] = 0$$



- ◆ ... that don't increase propagator power

$$u_i^\nu \frac{\partial}{\partial \ell_i^\nu} \rho_j = f_j \rho_j$$

[Gluza, Kajda, Kosower 10; Schabinger 11]

- ✓ Compute the u_i^ν by solving syzygy equations

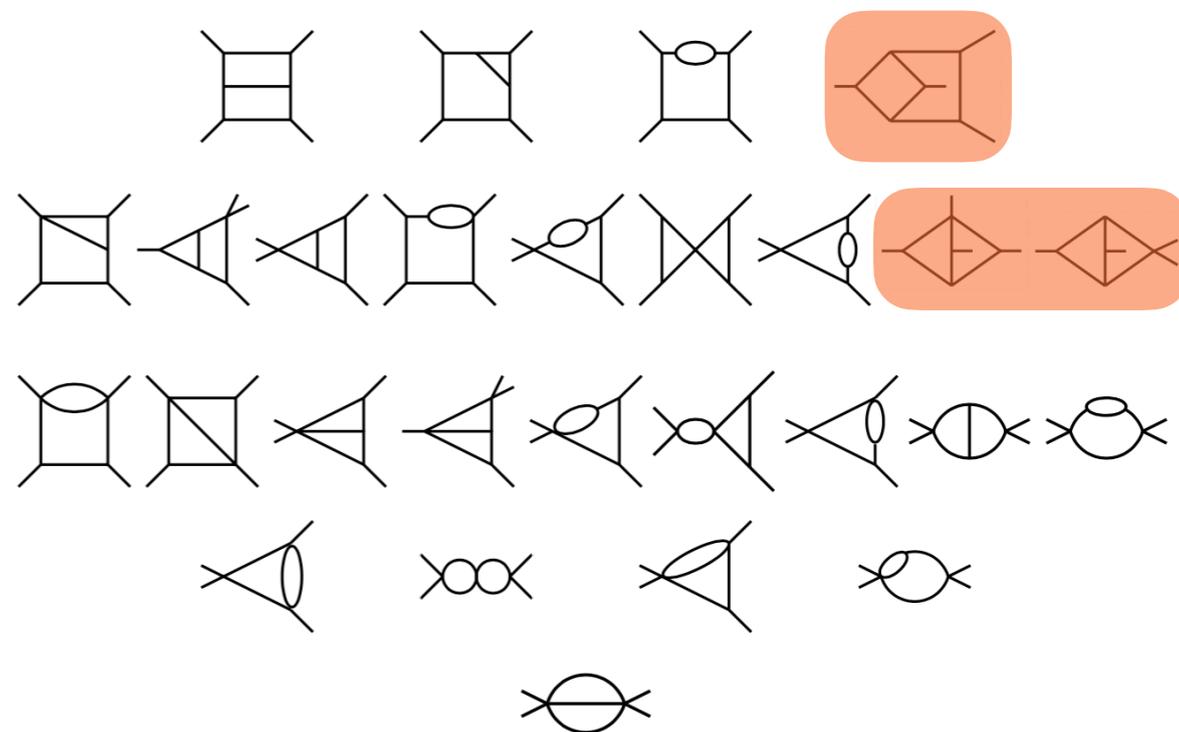
see also [Boehm, Georgoudis, Larsen, Schulze, Zhang, ... 16 - 19] and [Agarwal, von Manteuffel 19]

Two-loop (numerical) unitarity — surface terms II

[Ita 15; Abreu, Febres Cordero, Ita, Jaquier, Page, Zeng 17]

- ◆ Surface terms with u_i^ν and $t(\ell_l)$: $m_{\Gamma,k} = u_i^\nu \frac{\partial t(\ell_l)}{\partial \ell_i^\nu} + t(\ell_l) \left(\frac{\partial u_i^\nu}{\partial \ell_i^\nu} - \sum_{k \in P_\Gamma} f_k \right)$
- ◆ Linear algebra: independent surface terms, determine $|S_\Gamma|$, infer $|M_\Gamma|$
- ◆ u_i^ν are power-counting/theory independent

Analytic decomposition of integrand into master integrands and surface terms

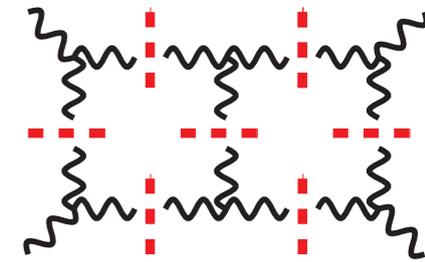


$$\mathcal{A}(\ell_l) = \sum_{\Gamma} \sum_{i \in M_{\Gamma} \cup S_{\Gamma}} c_{\Gamma,i}(D) \frac{m_{\Gamma,i}(\ell_l, D)}{\prod_{j \in P_{\Gamma}} \rho_j}$$

[Bern, Dixon, Dunbar, Kosower 94,95]

- ◆ In on-shell configuration of ℓ_l , integrand factorises

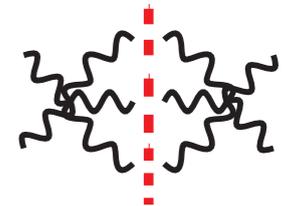
$$\sum_{\text{states } k \in T_{\Gamma}} \prod \mathcal{A}_k^{\text{tree}}(\ell_l^{\Gamma}) = \sum_{\Gamma' \geq \Gamma} \frac{c_{\Gamma'} m_{\Gamma'}(\ell_l^{\Gamma})}{\prod_{j \in (P_{\Gamma'} \setminus P_{\Gamma})} \rho_j(\ell_l^{\Gamma})}$$



- ◆ Need efficient computation of tree amplitudes (numerical)

- ✓ Berends-Giele recursion [Berends, Giele, 88]

- ✓ D_s -dimensional state sum, $D_s = 6, \dots, 10$



- ◆ Solve cut equations numerically: never construct analytic integrand

Compute amplitude at a numerical phase-space point

$$\mathcal{A} = \sum_{\Gamma} \sum_{i \in M_{\Gamma}} c_{\Gamma,i}(\epsilon, D_s, s, x = t/s) I_{\Gamma,i}$$

[Peraro, 16]

[Abreu, Dormans, Febres Cordero, Ita, Page, Sotnikov 18, 19]

- ◆ Rational dependence on all variables: compute in finite field

- ✓ No loss of precision in numerical manipulations

[von Manteuffel, Schabinger, 15], [Peraro, 16]

- ◆ Analytic ϵ dependence

- ✓ Compute at enough values to determine rational function

[Thiele's formula]

- ◆ Analytic D_s dependence

- ✓ EH: Quartic polynomial

- ✓ GB: Cubic polynomial / $(D_s - 2)$, vanishes at $D_s = 4$

Analytic ϵ and D_s dependence at numerical phase-space point amplitude from numerics

[Abreu, Dormans, Febres Cordero, Ita, Page, Sotnikov 18, 19]

$$\mathcal{A} = \sum_{\Gamma} \sum_{i \in M_{\Gamma}} c_{\Gamma,i}(\epsilon, D_s, s, x = t/s) I_{\Gamma,i}$$

◆ Analytic dependence in x is rational but complicated

✓ Target genuine two-loop contribution

◆ Reconstruct analytic form of finite remainder ($D_s = 4 - 2\epsilon$)

✓ Remove soft $\mathcal{R}^{(2)} = \mathcal{A}^{(2)} - \mathcal{S}\mathcal{A}^{(1)} + \frac{\mathcal{S}^2}{2}\mathcal{A}^{(0)} \quad \mathcal{S} = \sum_{i < j} \frac{1}{\epsilon^2} (-(p_i + p_j)^2)^{1-\epsilon}$

✓ UV finite by construction

$$\mathcal{R}_{\vec{h}}^{(2)} = \sum_i d_i(x) h_i(x)$$

Compute analytic amplitude from numerics

- ◆ Full framework automated in Caravel

[Abreu, Dormans, Febres Cordero, Ita, Kraus, Page, Pascual, Ruf, Sotnikov 20]

- ◆ Well tested in QCD calculations

- ✓ Analytic planar 5-parton amplitudes

[Abreu, Dormans, Febres Cordero, Ita, Page, Sotnikov 18, 19]

- ◆ New developments for four-graviton amplitudes

- ✓ Non-planar topologies

- ✓ GB, R^3 Feynman rules

- ✓ For EH: cubic-interaction reformulation [Cheung, Remmen 17]

- ◆ 20 numerical evaluations for each helicity

[Abreu, Febres Cordero, Ita, Jaquier, Page, Ruf, Sotnikov 20]

Analytic 4-graviton 2-loop amplitudes:

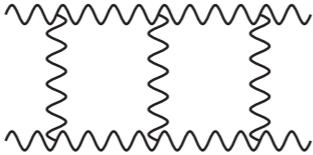
$$\mathcal{A}(++++), \mathcal{A}(-+++), \mathcal{A}(- - ++)$$

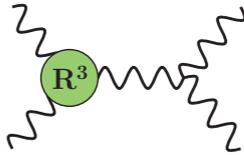


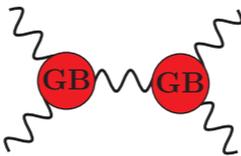
Caravel

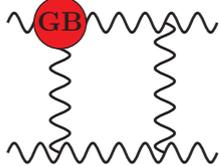
Soon to be released!

◆ Computed the three independent helicities for

✓ EH gravity 

✓ Tree-level R^3 

✓ Tree-level GB-GB 

✓ One-loop GB 

◆ Checks

✓ Remainders are finite, correct symmetry, no spurious poles

✓ 1-loop amplitudes [Dunbar, Norridge 95], [Bern, Cheung, Chi, Davies, Dixon, Nohle, unpub.]

✓ GB tree and 1-loop: + + + + and - - + + [Bern, Cheung, Chi, Davies, Dixon, Nohle, 15, unpub.]

✓ R^3 tree [Bern, Cheung, Chi, Davies, Dixon, Nohle, 15], [Dunbar, Jehu, Perkins, 17]

✓ 2-loop: + + + + [Bern, Cheung, Chi, Davies, Dixon, Nohle, unpub.], [Dunbar, Jehu, Perkins, 17]

◆ Remainders

- ✓ + + + +: up to weight 1, symmetric on s, t, u
- ✓ - + + +: up to weight 2, only logs, symmetric on s, t, u
- ✓ - - + +: up to weight 4, classical polylogs

$$\mathcal{R}_{++++}^{(2)} = f(s, t, u) + f(t, u, s) + f(u, s, t) + \frac{117617}{21600} s t u + 30 s t u (c_{\text{GB}}(\mu) - 2c_{\text{R}^3}(\mu))$$

$$f(s, t, u) = \frac{s}{60} (14s^2 - 9t u) \log(-s)$$

[Bern, Cheung, Chi, Davies, Dixon, Nohle, unpub.]

[Dunbar, Jehu, Perkins, 17]

s -channel Regge limit ($s \gg -t > 0$):

$$\begin{aligned} \mathcal{R}_{\{-, -, +, +\}}^{(2)} = & s^3 \left\{ 2 \frac{s}{t} \pi^2 \left(\frac{i\pi}{2} - L \right)^2 - 3\pi^2 L^2 \right. \\ & + \frac{107}{10} \pi^2 L + \frac{14191}{1350} \pi^2 - \frac{158}{45} \pi^4 - \frac{13049}{2160} \\ & + i\pi \left[-\frac{14}{3} L^3 + \frac{87}{10} L^2 - \left(8\pi^2 - \frac{17749}{450} \right) L \right. \\ & \left. \left. - 20\zeta_3 + \frac{2621}{210} \pi^2 - \frac{11221}{375} \right] + \mathcal{O}(-t/s) \right\} \end{aligned}$$

$$L = \log(-s/t)$$

Julio's
talk

$$\mathcal{C}_{\text{GB}} = \left(\frac{53}{90} \frac{1}{\epsilon} + c_{\text{GB}}(\mu) \right) \mu^{-2\epsilon} \quad \mathcal{C}_{\text{R}^3} = \left(\frac{209}{1440} \frac{1}{\epsilon} + c_{\text{R}^3}(\mu) \right) \mu^{-4\epsilon}$$

◆ Remainders depend on a single coupling

$$c(\mu) = c_{\text{GB}}(\mu) - 2 c_{\text{R}^3}(\mu)$$

$$\begin{aligned} \mathcal{R}_{\{+,+,+,+\}}^{(2)} &= 30 \, st u c(\mu) + \dots & \mathcal{R}_{\{-,+,+,+\}}^{(2)} &= 3 \, st u c(\mu) + \dots \\ \mathcal{R}_{\{-,-,+,+\}}^{(2)} &= 0 \, c(\mu) + \dots \end{aligned}$$

✓ same combination in 5-point 2-loop all-plus

[Dunbar, Jehu, Perkins, 17]

◆ μ -dependence of remainder

$$\mu \frac{\partial}{\partial \mu} \mathcal{R}_{\vec{h}}^{(2)} = \left[\frac{1}{120} + \mu \frac{\partial}{\partial \mu} c(\mu) \right] \frac{\mathcal{A}_{\vec{h}}^{\text{tree}, \text{R}^3}}{\mathcal{C}_{\text{R}^3}}$$

$$\mathcal{A}_{+,+,+,+}^{\text{tree}, \text{R}^3} = -\mathcal{C}_{\text{R}^3} \frac{209}{24\epsilon} st u \quad \mathcal{A}_{-,+,+,+}^{\text{tree}, \text{R}^3} = -\mathcal{C}_{\text{R}^3} \frac{209}{240\epsilon} st u \quad \mathcal{A}_{-,-,+,+}^{\text{tree}, \text{R}^3} = 0$$

✓ GB counter-term is evanescent

[Bern, Cheung, Chi, Davies, Dixon, Nohle, 15]

[Bern, Chi, Dixon, Edison, 17]

[Dunbar, Godwin, Jehu, Perkins, 17]

- ◆ Two-loop numerical unitarity is a mature framework
 - ✓ Efficient approach to obtain analytic results
 - ✓ Generic: state of the art results in QCD & GR
- ◆ Computed the full set of four-graviton helicity amplitudes
 - ✓ Confirm known $++++$
 - ✓ New results for $-+++$ and $--++$
 - ✓ Led to new insights into relation between multi-loop amplitudes and classical gravitational dynamics
[Bern, Ita, Parra-Martinez, Ruf, 20]
 - ✓ Interesting coupling and μ dependence in remainder

THANK YOU!