#### The Momentum Amplituhedron Boundaries



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> Zoomplitudes 2020 13.05.2020

Based on: arXiv:2003.13704 with T. Lukowski, R. Moerman



### Outline

- \* Introduction and Motivation
- \* Reminder: Momentum Amplituhedron
- \* Singularities of Amplitudes from its Boundaries
- \* Conclusions

Scattering amplitudes in (planar) N=4 sYM: encoded in "Amplituhedra"

More general paradigm: **positive geometries** -> regions with boundaries equipped with rational differential forms -> the forms are logarithmic on any boundary



(picture by A. Gilmore)

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Scattering amplitudes in (planar) N=4 sYM: encoded in "Amplituhedra"

> More general paradigm: positive geometries

- \* amplituhedron (N. Arkani-Hamed, J. Trnka)
- \* Kinematic associatedron (N. Arkani-Hamed, Y. Bai, S. He, G. Yan)
- \* cosmological polytope (N. Arkani-Hamed, P. Benincasa, A. Postnikov)
- \* CFTS (B. Eden, P. Heslop, L.Mason; N. Arkani-Hamed, Y.-T. Huang, S.-H. Shao)
- \* ..
- \* momentum amplituhedron (D. Damgaard, LF, T. Lukowski, M. Parisi)

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Conjecture:

amplituhedra encode physical singularities of amplitudes in the structure of their boundaries -> Facets (= codim 1): straightforward -> What happens deeper in geometry?

In this talk:

Boundaries of Mom Ampl<->Singularities of Tree Ampls Boundary Stratification + Full Classification of Sings

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Color-ordered amplitudes: pure Yang-Mills

\* Multi-particle poles when sum of adjacent momenta goes on-shell:

$$P_{i,j}^2 = (p_i + p_{i+1} + \dots + p_j)^2 \to 0$$

$$A_{n}^{\text{tree}}(1,...,n) \xrightarrow{P_{i,j}^{2} \to 0} \sum_{h=\pm 1} A_{j-i+2}^{\text{tree}}(i,...,j,P_{i,j}^{h}) \frac{1}{P_{i,j}^{2}} A_{n-j+i}^{\text{tree}}(P_{i,j}^{-h},j+1,...,i-1)$$

\* Collinear and soft singularities: universal factorization properties

$$A_{n}^{\text{tree}}(\dots, i^{h_{i}}, (i+1)^{h_{i+1}}, \dots) \xrightarrow{p_{i}||p_{i+1}} \sum_{h=\pm 1} A_{n-1}^{\text{tree}}(\dots, P_{i,i+1}^{h}, \dots) \operatorname{Split}_{-h}^{\text{tree}}(i^{h_{i}}, (i+1)^{h_{i+1}})$$

$$A_n^{\text{tree}}(\ldots, i, s^h, j, \ldots) \xrightarrow{p_s \to 0} A_{n-1}^{\text{tree}}(\ldots, i, j, \ldots) \text{ Soft}^{\text{tree}}(i, s^h, j)$$

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#### Color-ordered amplitudes: N=4 super Yang-Mills

\* Multi-particle poles when sum of adjacent momenta goes on-shell:

$$A_{n,k}(1,...,n) \xrightarrow{P_{i,j}^2 \to 0} \sum_{k'=1}^k \int d^4 \eta_{P_{i,j}} A_{j-i+2,k'}(i,...,j,P_{i,j}) \frac{1}{P_{i,j}^2} A_{n-j+i,k-k'+1}(P_{i,j},j+1,...,i-1)$$

\* Collinear and soft singularities:

$$p_{1} \| p_{2} A_{n,k}(1,2,3,...,n) \xrightarrow{P_{12}^{2} \rightarrow 0} \int d^{4} \eta_{P_{12}} \left[ \begin{array}{c} \text{Split}_{0}^{\text{tree}} A_{n-1,k}(P_{12},3,...,n) + \begin{array}{c} \text{Split}_{-1}^{\text{tree}} A_{n-1,k-1}(P_{12},3,...,n) \right] \\ \text{helicity-preserving} \\ \text{super-splitting function} \\ \text{Split}_{0}^{\text{tree}}(z;\eta_{1},\eta_{2},\eta_{P_{12}}) \equiv \frac{1}{\sqrt{z(1-z)}} \frac{1}{\langle 12 \rangle} \prod_{A=1}^{4} (\eta_{P_{12}A} - \sqrt{z}\eta_{1A} - \sqrt{1-z}\eta_{2A}) \\ \text{split}_{-1}^{\text{tree}}(z;\eta_{1},\eta_{2},\eta_{P_{12}}) \equiv \frac{1}{\sqrt{z(1-z)}} \frac{1}{\langle 12 \rangle} \prod_{A=1}^{4} (\eta_{P_{12}A} - \sqrt{z}\eta_{1A} - \sqrt{1-z}\eta_{2A}) \\ \text{split}_{-1}^{\text{tree}}(z;\eta_{1},\eta_{2},\eta_{P_{12}}) \equiv \frac{1}{\sqrt{z(1-z)}} \frac{1}{|12|} \prod_{A=1}^{4} (\eta_{1A}\eta_{2A} + \sqrt{1-z}\eta_{1A}\eta_{P_{12A}} - \sqrt{z}\eta_{2A}\eta_{P_{12}A} - \sqrt{z}\eta_{P_{12}A} - \sqrt{z}$$

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#### Color-ordered amplitudes: N=4 super Yang-Mills

\* Multi-particle poles when sum of adjacent momenta goes on-shell:

$$A_{n,k}(1,...,n) \xrightarrow{P_{i,j}^2 \to 0} \sum_{k'=1}^k \int d^4 \eta_{P_{i,j}} A_{j-i+2,k'}(i,...,j,P_{i,j}) \frac{1}{P_{i,j}^2} A_{n-j+i,k-k'+1}(P_{i,j},j+1,...,i-1)$$

\* Collinear and soft singularities:

$$p_{1} || p_{2} : A_{n,k}(1,2,3,...,n) \xrightarrow{P_{12}^{2} \rightarrow 0} \int d^{4} \eta_{P_{12}} \left[ \begin{array}{c} \text{Split}_{0}^{\text{tree}} A_{n-1,k}(P_{12},3,...,n) + \text{Split}_{-1}^{\text{tree}} A_{n-1,k-1}(P_{12},3,...,n) \right] \\ p_{\text{arametrised by}} \\ (12) \rightarrow 0 \quad (23) \rightarrow 0 \end{array} \right] p_{12} \rightarrow 0 \quad (23) \rightarrow 0$$

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#### Color-ordered amplitudes: N=4 super Yang-Mills



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#### Color-ordered amplitudes: N=4 super Yang-Mills



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#### Scattering amplitudes in N=4 sYM

on-shell superspace  $(\lambda_i^{lpha}, \tilde{\lambda}_i^{\dot{lpha}}, \eta_i^A)$ 

$$(\lambda_i^a, \eta_i^r | \tilde{\lambda}_i^{\dot{a}}, \tilde{\eta}_i^{\dot{r}}), a, \dot{a}, r, \dot{r} = 1, 2$$

$$\widetilde{q}^{\dot{a}r} = \sum_{i=1}^{n} \widetilde{\lambda}_{i}^{\dot{a}} \eta_{i}^{r}$$
$$q^{ar} = \sum_{i=1}^{n} \lambda_{i}^{a} \widetilde{\eta}_{i}^{\dot{r}}$$

Associate a, à with SU(2)xSU(2) R-symmetry indices:

$$\eta^a 
ightarrow d\lambda^a \,, \quad ilde{\eta}^{\dot{a}} 
ightarrow d ilde{\lambda}^{\dot{a}}$$

n-point super-amplitude in non-chiral space <-> 2n form

(S. He, C. Zhang)

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 $(\lambda_i^a, \eta_i^r | \tilde{\lambda}_i^{\dot{a}}, \tilde{\eta}_i^{\dot{r}}), a, \dot{a}, r, \dot{r} = 1,2$ 

n-point super-amplitude in non-chiral space <-> 2n form

$$\mathcal{A}_{n,k} := (dq)^{4} \wedge \Omega_{n,k}$$
Example: 4-point MHV
$$\Omega_{4,2} = \frac{(d\tilde{q})^{4}}{st} = \operatorname{dlog} \frac{\langle 12 \rangle}{\langle 13 \rangle} \wedge \operatorname{dlog} \frac{\langle 23 \rangle}{\langle 13 \rangle} \wedge \operatorname{dlog} \frac{\langle 34 \rangle}{\langle 13 \rangle} \wedge \operatorname{dlog} \frac{\langle 41 \rangle}{\langle 13 \rangle}$$
Geometry whose canonical form gives the amplitude form?

Momentum amplituhedron M!

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Bosonized spinor helicity variables:

$$\tilde{\Lambda}_{i}^{\dot{A}} = \begin{pmatrix} \tilde{\lambda}_{i}^{\dot{a}} \\ \tilde{\phi}_{\dot{a}}^{\dot{\alpha}} \cdot \tilde{\eta}_{i}^{\dot{a}} \end{pmatrix}, \quad \dot{A} = (\dot{a}, \dot{\alpha}) = 1, \dots, k+2 \qquad \Lambda_{i}^{A} = \begin{pmatrix} \lambda_{i}^{a} \\ \phi_{a}^{\alpha} \cdot \eta_{i}^{a} \end{pmatrix}, \quad A = (a, \alpha) = 1, \dots, n-k+2$$

$$\left\{ \text{matrix } \tilde{\Lambda} \text{ positive and matrix } \Lambda^{\perp} \text{ positive} \right\}$$

Positive region

Momentum amplituhedron: Image of the positive Grassmannian  $G_{+}(k,n)$  through the map

$$\Phi_{(\Lambda,\tilde{\Lambda})}: G_+(k,n) \to G(k,k+2) \times G(n-k,n-k+2)$$

defined as:

$$\tilde{Y}_{\dot{\alpha}}^{\dot{A}} = \sum_{i=1}^{n} c_{\dot{\alpha}i} \tilde{\Lambda}_{i}^{\dot{A}} \qquad Y_{\alpha}^{A} = \sum_{i=1}^{n} c_{\alpha i}^{\perp} \Lambda_{i}^{A}$$

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$$\tilde{Y}_{\dot{\alpha}}^{\dot{A}} = \sum_{i=1}^{n} c_{\dot{\alpha}i} \tilde{\Lambda}_{i}^{\dot{A}} \qquad Y_{\alpha}^{A} = \sum_{i=1}^{n} c_{\alpha i}^{\perp} \Lambda_{i}^{A}$$

Facets: codimension-1 boundaries

 $\langle Yii+1 \rangle = 0, \quad [\tilde{Y}ii+1] = 0$  Collinear limits  $S_{i,i+1...,i+p} = 0, \quad p = 2,...,n-4$  Factorizations Uplift of planar Mandelstam variables  $S_{i,i+1...,i+p} = \sum_{i \leq j_1 < j_2 \leq i+p} \langle Yj_1j_2 \rangle [\tilde{Y}j_1j_2]$ 

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$$\tilde{Y}_{\dot{\alpha}}^{\dot{A}} = \sum_{i=1}^{n} c_{\dot{\alpha}i} \tilde{\Lambda}_{i}^{\dot{A}} \qquad Y_{\alpha}^{A} = \sum_{i=1}^{n} c_{\alpha i}^{\perp} \Lambda_{i}^{A}$$

Facets: codimension-1 boundaries

 $\langle Yii+1 \rangle = 0$ ,  $[\tilde{Y}ii+1] = 0$  Collinear limits  $S_{i,i+1...,i+p} = 0$ , p = 2,...,n-4 Factorizations Volume form

Differential form with log singularities on all boundaries = Sum over cells of push-forwards of canonical diff-form

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$$\tilde{Y}_{\dot{\alpha}}^{\dot{A}} = \sum_{i=1}^{n} c_{\dot{\alpha}i} \tilde{\Lambda}_{i}^{\dot{A}} \qquad Y_{\alpha}^{A} = \sum_{i=1}^{n} c_{\alpha i}^{\perp} \Lambda_{i}^{A}$$

Facets: codimension-1 boundaries

$$\langle Yii+1 \rangle = 0$$
,  $[\tilde{Y}ii+1] = 0$  Collinear limits  
 $S_{i,i+1...,i+p} = 0$ ,  $p = 2,...,n-4$  Factorizations

 $\Omega_{n,k}$ : volume form with logarithmic singularities on all boundaries of the space

$$\Omega_{n,k} = \sum_{\sigma} \operatorname{dlog} \alpha_1^{\sigma} \wedge \operatorname{dlog} \alpha_2^{\sigma} \wedge \ldots \wedge \operatorname{dlog} \alpha_{2n-4}^{\sigma}$$
Logarithmic differential form

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$$\tilde{Y}_{\dot{\alpha}}^{\dot{A}} = \sum_{i=1}^{n} c_{\dot{\alpha}i} \tilde{\Lambda}_{i}^{\dot{A}} \qquad Y_{\alpha}^{A} = \sum_{i=1}^{n} c_{\alpha i}^{\perp} \Lambda_{i}^{A}$$

→ Volume form

$$\mathbf{\Omega}_{n,k} = \sum_{\sigma} \operatorname{dlog} \alpha_1^{\sigma} \wedge \operatorname{dlog} \alpha_2^{\sigma} \wedge \ldots \wedge \operatorname{dlog} \alpha_{2n-4}^{\sigma}$$

→ Volume function

$$\Omega_{n,k} \wedge d^4 P \,\delta^4(P) = \prod_{\alpha=1}^{n-k} \langle Y_1 \dots Y_{n-k} d^2 Y_\alpha \rangle \prod_{\dot{\alpha}=1}^k [\tilde{Y}_1 \dots \tilde{Y}_k d^2 \tilde{Y}_{\dot{\alpha}}] \,\delta^4(P) \,\Omega_{n,k}$$

- Amplitude  $\mathcal{A}_{n,k}^{\text{tree}} = \delta^4(p) \int d\phi_a^1 \dots d\phi_a^{n-k} \int d\tilde{\phi}_{\dot{a}}^1 \dots d\tilde{\phi}_{\dot{a}}^k \Omega_{n,k}(Y^*, \tilde{Y}^*, \Lambda, \tilde{\Lambda})$ 

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Example

\* MHV4 amplitude:

$$\tilde{Y}_{\dot{\alpha}}^{\dot{A}} = \sum_{i=1}^{n} c_{\dot{\alpha}i} \tilde{\Lambda}_{i}^{\dot{A}} \qquad Y_{\alpha}^{A} = \sum_{i=1}^{n} c_{\alpha i}^{\perp} \Lambda_{i}^{A}$$

$$\alpha_1 = \frac{\langle Y12 \rangle}{\langle Y13 \rangle}, \alpha_2 = \frac{\langle Y23 \rangle}{\langle Y13 \rangle}, \alpha_3 = \frac{\langle Y34 \rangle}{\langle Y13 \rangle}, \alpha_4 = \frac{\langle Y14 \rangle}{\langle Y13 \rangle}$$

$$\Omega_{4,2} = \bigwedge_{j=1}^{4} \operatorname{dlog} \alpha_{j} = \operatorname{dlog} \frac{\langle Y12 \rangle}{\langle Y13 \rangle} \wedge \operatorname{dlog} \frac{\langle Y23 \rangle}{\langle Y13 \rangle} \wedge \operatorname{dlog} \frac{\langle Y34 \rangle}{\langle Y13 \rangle} \wedge \operatorname{dlog} \frac{\langle Y14 \rangle}{\langle Y13 \rangle}$$

$$= \frac{\langle 1234 \rangle^{2}}{\langle Y12 \rangle \langle Y23 \rangle \langle Y34 \rangle \langle Y41 \rangle} \langle Yd^{2}Y_{1} \rangle \langle Yd^{2}Y_{2} \rangle$$
Divergences on the 4 facets of the momentum amplituhedron:  

$$\langle Yii + 1 \rangle = 0, i = 1, ..., 4$$

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Examples

\* NMHV6 amplitudes:

$$\tilde{Y}_{\dot{\alpha}}^{\dot{A}} = \sum_{i=1}^{n} c_{\dot{\alpha}i} \tilde{\Lambda}_{i}^{\dot{A}} \qquad Y_{\alpha}^{A} = \sum_{i=1}^{n} c_{\alpha i}^{\perp} \Lambda_{i}^{A}$$

$$\Omega_{6,3} = \Omega_{6,3}^{(612)} + \Omega_{6,3}^{(234)} + \Omega_{6,3}^{(456)} = \Omega_{6,3}^{(123)} + \Omega_{6,3}^{(345)} + \Omega_{6,3}^{(561)}$$



Divergences on the 15 facets of the momentum amplituhedron:

 $\langle Yii+1 \rangle = 0, i = 1,...,6, [\tilde{Y}ii+1] = 0, i = 1,...,6, S_{i,i+1,i+2} = 0, i = 1,2,3$ 

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Boundary stratification from Mathematica package amplituhedronBoundaries.m up to n=9 and all k (extended to n=11 via dual graphs) (T. Lukowski, R. Moerman) based on positroids.m

(J. Bourjaily)

Boundary stratification from Mathematica package amplituhedronBoundaries.m (T. Lukowski, R. Moerman) up to n=9 and all k (extended to n=11 via dual graphs) based on positroids.m (J. Bourjaily)

-> labeled by plabic diagrams of cells in G+(k,n)

Codim 1 boundaries



 $\langle Yii + 1 \rangle = 0$ 



 $[\tilde{Y}ii + 1] = 0$ 



factorization

 $S_{i,i+1,\ldots,i} = 0$ 

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-> labeled by plabic diagrams of cells in G+(K,n)



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Codim 1 boundaries

Codim 2 boundaries



further factorizations

and collinear limits





soft



i  $A_{n-1,k}$ 



 $\langle Yii + 1 \rangle = 0 = \langle Yi - 1i \rangle \qquad [\tilde{Y}ii + 1] = 0 = [\tilde{Y}i - 1i]$ 

Intersection of two consecutive codim 1 collinear boundaries

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-> labeled by plabic diagrams of cells in G+(k,n)

Codim 1 boundaries







Codim 2 boundaries

further factorizations and collinear limits





(J. Bourjaily)

+ generalization deeper in the geometry

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- \* black lollipop helicity-preserving soft limit
- \* white lollipop helicity-reducing soft limit
- \* single line forward-limit
- \* top cell, collinear limit or factorization channel for an amplitude A<sub>n'k'</sub> with n'<n and k'≤k. Can be any boundary of M as long as it is a connected diagram

#### Examples

 $\therefore$  MHVn <-> for k=2 it agrees with boundary stratification of Gr+(2,n)

Representatives:

Codim 1 boundaries



Codim 2 boundaries

Codim 3 boundaries

Codim 4 boundaries



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#### \* Number of boundaries for a given dimension:

(n,k)ackslash d	0	1	2	3	4	5	6	7	8	9	10	11	12
(4,2)	6	12	10	4	1								
(5,2)	10	30	40	30	15	5	1						
(6,2)	15	60	110	120	90	50	21	6	1				
(6,3)	20	90	180	215	180	114	54	15	1				
(7,2)	21	105	245	350	350	266	161	77	28	7	1		
(7,3)	35	210	560	910	1050	938	665	350	119	21	1		
(8,2)	28	168	476	840	1050	1008	784	504	266	112	36	8	1
(8,3)	56	420	1400	2870	4200	4788	4424	3262	1820	720	188	28	1
(8, 4)	70	560	1960	4200	6426	7672	7420	5696	3264	1280	300	32	1

Generating fc  $F_{n,k}(x) = \sum_{\sigma_{n,k}} (-x)^{\dim \sigma}$  $F_{n,k}(1) = 1$ 

+ parity conjugates (n,k) -> (n,n-k)

 $\bigstar$  Euler characteristic = 1 -> indicates that M is (2n-4)-diml ball

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Dual graphs: enumeration enabling to go up to n=11 any k -> Each boundary labelled by a partial triangulation of a regular n-gon + additional decorations



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Conclusions

positive geometry for tree-level amplitudes in N=4 sYM in spinor helicity space

Classified all physical singularities of tree-level ampls in N=4 sYM by studying its boundaries
 Each singularity comes from a subsequent multi-particle factorization and collinear limit
 Singularities translated to geometry as appropriate intersection of facets

Proof that momentum amplituhedron is a ball?
 Loop amplitudes?

Thank you!