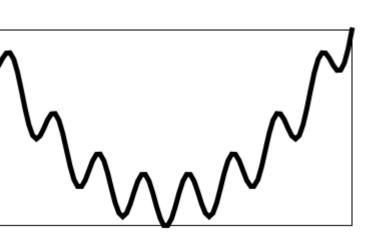
Gravity waves from nonlinear axion-like particle dynamics Aleksandr Chatrchyan¹, Joerg Jaeckel² CLUSTER OF EXCELLENCE ¹ DESY Hamburg, ² ITP Heidelberg QUANTUM UNIVERSE Based on: arxiv:2004.07844, in preparation

Axion-like particles (ALPs) and nonlinear dynamics

Axions and ALPs play an important role in cosmology, including dark matter and inflationary model building.

Wiggly potentials: A characteristic feature of ALPs is the cosine-type potential, which conserves the discrete shift symmetry. In the more general case, nonperiodic, wiggly potentials are possible (e.g. via monodromy in string theory [1]).



Homogeneity: Inflation sets almost homogeneous initial conditions for the field. At the same time, the post-inflationary evolution over the nonquadratic regions of the potential leads to the amplification of fluctuations, via parametric/tachyonic instabilities.

This impacts the dynamics, by possibly leading to the **fragmentation** of the field [2,3], **bubble nucleation** and transitions between local minima, and by sourcing gravity waves.

Production of gravity waves: the misalignment mechanism

The rapid conversion of energy from the homogeneous field into fluctuations is accompanied by the production of a stochastic GW background, peaked at similar wavelengths, determined by the mass of the scalar field (or, in the case of a phase transition, by the typical bubble distance).

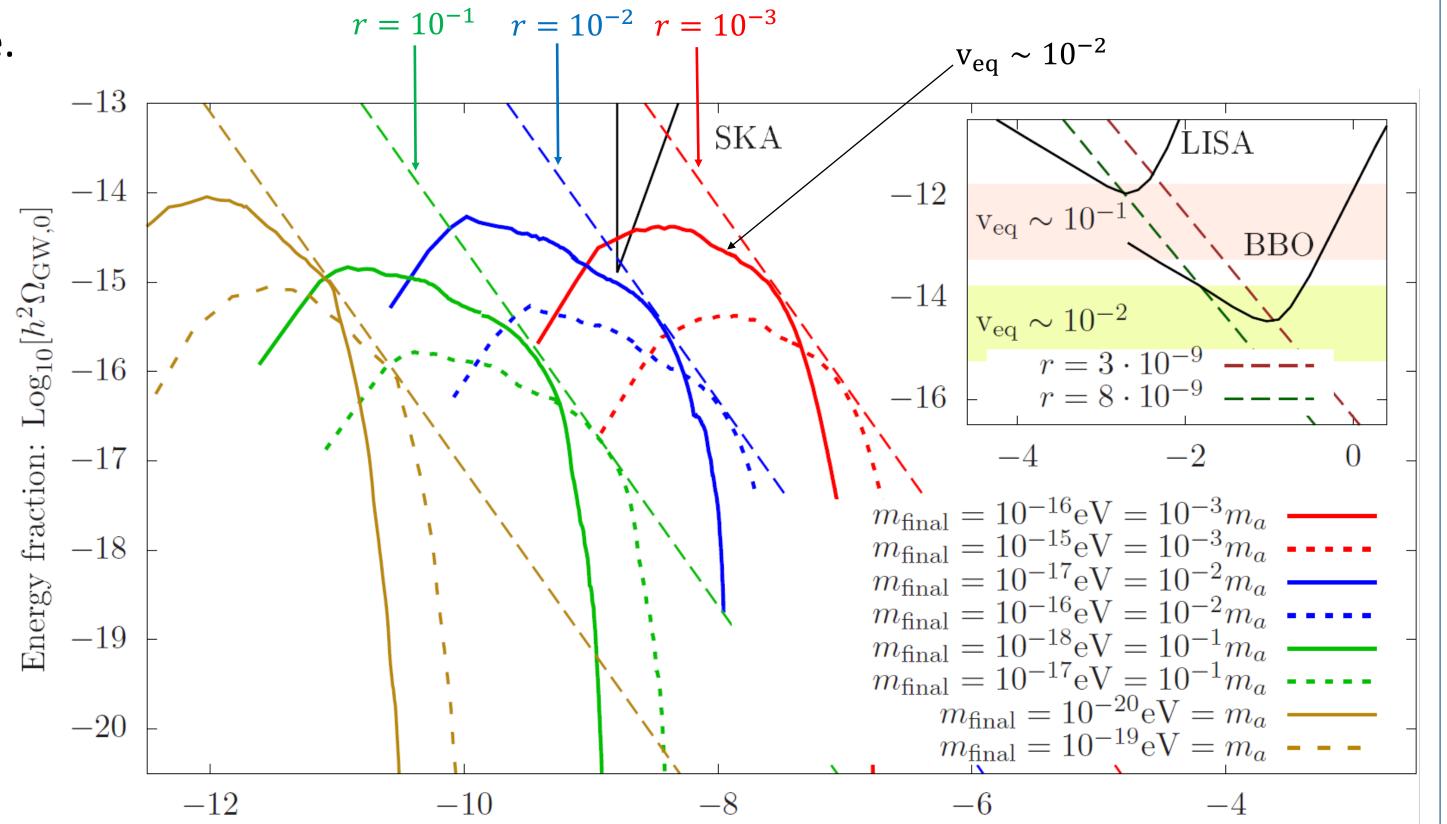
ALP DM from the misalignment mechanism: Light ALPs with masses from a wide range, $(10^{-21} - 10^3)$ eV, can behave as dark matter. The resulting GW background was investigated by means of numerical lattice simulations, using a modified version of the "Hlattice" code [4], for the ALP potential $U(\varphi) = \frac{1}{2}m^2\varphi^2 + \Lambda^4\left(1 - \cos\frac{\varphi}{f}\right)$, with a large misalignment angle $\varphi_1/f \gg 1$.

Strong bounds on the signal (gold line) imposed by the DM abundance.

Weaker bounds are possible, if the field remains relativistic for a longer period after fragmentation, so that the field would dilute its energy more efficiently. This can be achieved if the mass near the minimum is small.

Modified potential, $U \to U + \xi \Lambda^4 \left(1 - \cos\left(\frac{2\varphi}{f}\right)\right)$, can lead to a smaller final mass, $m_{final} = rm_a$, $r \ll 1$ (green, blue, red lines), which extends the parameter space for the signal.

Constraints on ALP velocities: typical velocities at matter-radiation equality v_{eq} slightly higher than for typical warm dark matter.



Nonequilibrium bubble nucleation

The field can get trapped in one of the local minima, and perform a transition to a lower minimum by means of **bubble nucleation**.

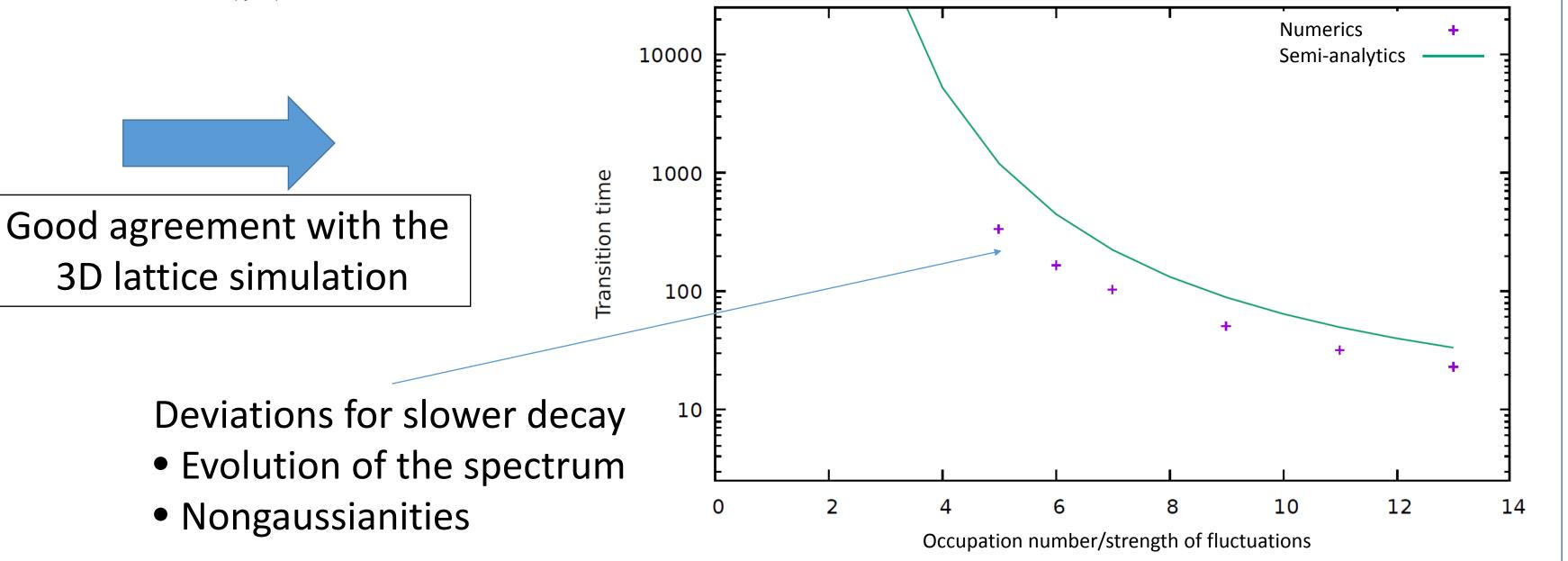
Instabilities produce strong fluctuations at small momenta. This can strongly modify the transition from the usual thermal case.

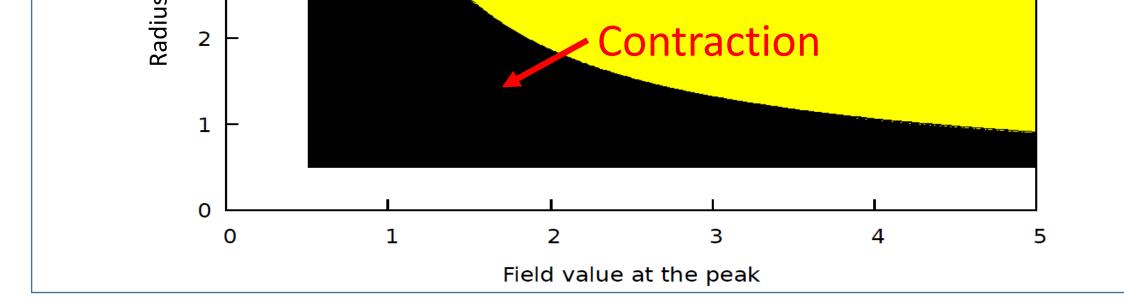
This "nonequilibrium" regime of bubble nucleation was explored using a real-time approach, based on the classical-statistical approximation [5], for the simple case of an asymmetric double well potential.

A semi-analytical method to compute the transition time/average bubble distance:

Estimate the number density of peaks in the field (depends on the spectrum via the variance and the first few moments)

e.g. for a gaussian random field $n_{peak}(\nu) \approx \frac{1}{8\pi^2} \left(\frac{\langle |\nabla \varphi|^2 \rangle}{3 \langle \varphi^2 \rangle} \right)^{\frac{3}{2}} \nu^2 e^{-\frac{\nu^2}{2}}, \qquad \nu^2 = \frac{\varphi_{peak}^2}{\langle \varphi^2 \rangle}$ Determine which peaks turn into expanding bubbles, 5 **Expanding bubbles** of the peak





Summary

- Axions and ALPs can exhibit very rich dynamics in the early universe, accompanied by the production of a stochastic GW background, potentially within reach of future experiments.
- Nonlinear effects can be important for the dynamics and should be taken into account.

References

[1] E. Silverstein, A. Westphal, Phys. Rev. D78, 106003 (2008). [2] J. Berges, A. Chatrchyan, J. Jaeckel, JCAP 08 (2019) 020, 1903.03116 [3] N. Fonseca, E. Morgante, R. Sato, G. Servant, JHEP 04 (2020) 010, 1911.08472 [4] Z. Huang, Phys. Rev. D83, 123509 (2011).

[5] A. Linde, Nucl.Phys.B 372 (1992) 421, 9110037