

QCD@LHC

Aug 31–Sept 3 · 2020



QCD impact in BSM searches at LHC

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On behalf of the ATLAS and CMS Collaboration

QCD@LHC-X, the 2020 edition of the joint theoretical-experimental conference on applications of QCD to collider physics, will be held online via Zoom.

QCD@LHC-X , 31st August– 3rd September 2020,
Virtually in Bologna, Italy

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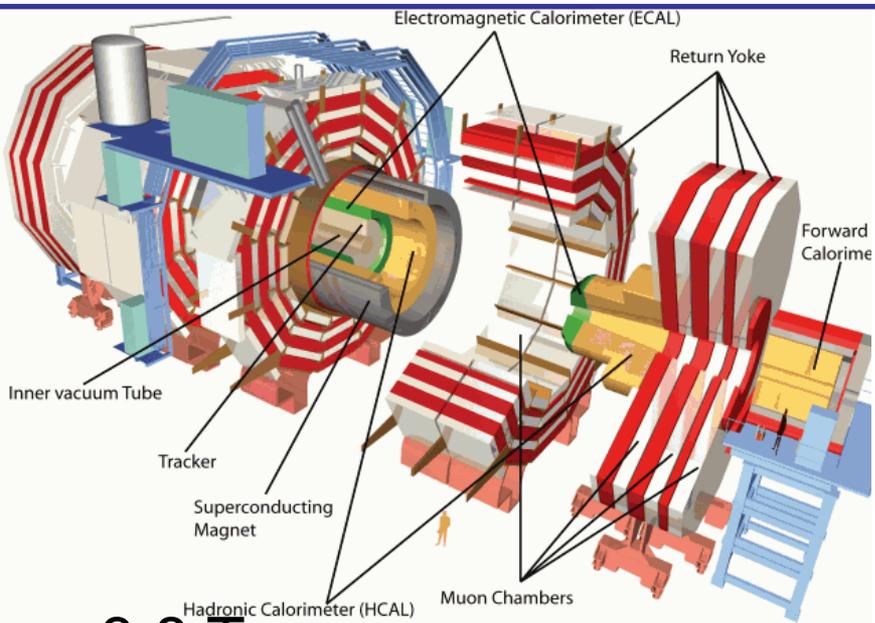
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01/09/2020



- **Introduction**
- **QCD uncertainties in *selected powerful and general* BSM searches**
 - Dijet and Monojet searches (M_{jj} , χ , $E_{T\text{miss}}$)
 - Dilepton resonant and non-resonant searches (M_{ll})
- **QCD uncertainties when *looking inside jets***
 - Quark-gluon separation
 - Jet taggers, jet mass
- **Summary and Outlook**





3.8 T

Pixels

$$\sigma/pT \sim 1.5 \cdot 10^{-4} pT(\text{GeV}) \oplus 0.005$$

Electromagnetic Calorimeter

$$\sigma E/E \approx 2.9\%/\sqrt{E(\text{GeV})} \oplus 0.5\% \oplus 0.13 \text{GeV}/E$$

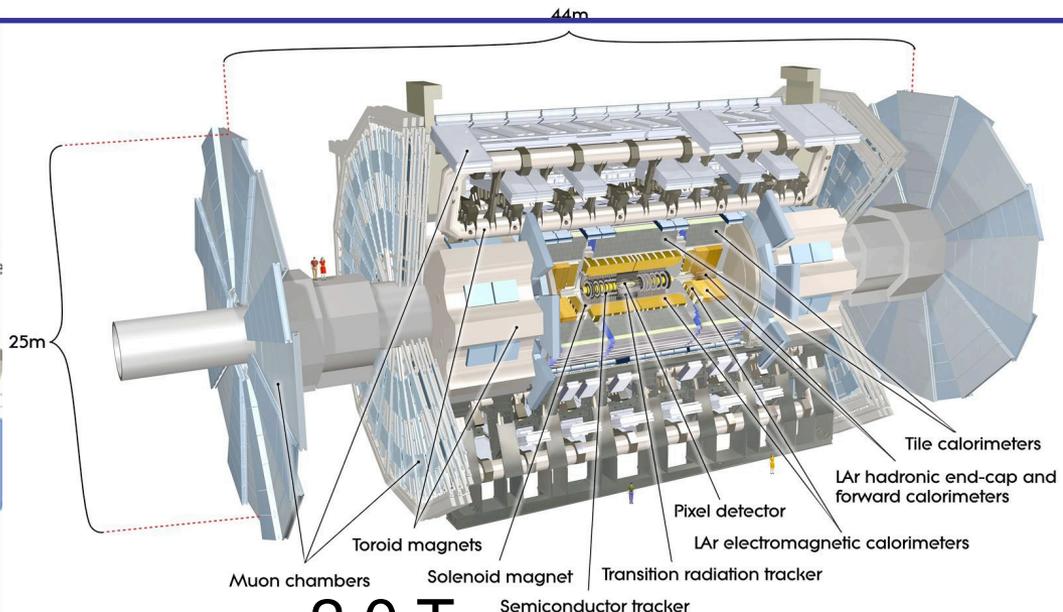
Hadronic Calorimeter

$$\sigma E/E \approx 120\%/\sqrt{E(\text{GeV})} \oplus 6.9\%$$

Muon Spectrometer

$$\sigma pT/pT \approx 1\% \text{ for low } pT \text{ muons}$$

$$\sigma pT/pT \approx 5\% \text{ for } 1 \text{ TeV muons}$$



2.0 T

Pixels, Si strips & Straw tubes

$$\sigma/pT \sim 3.8 \cdot 10^{-4} pT(\text{GeV}) \oplus 0.015$$

Electromagnetic Calorimeter

$$\sigma E/E \approx 10\%/\sqrt{E(\text{GeV})} \oplus 0.7\% \oplus 0.2 \text{GeV}/E$$

Hadronic Calorimeter

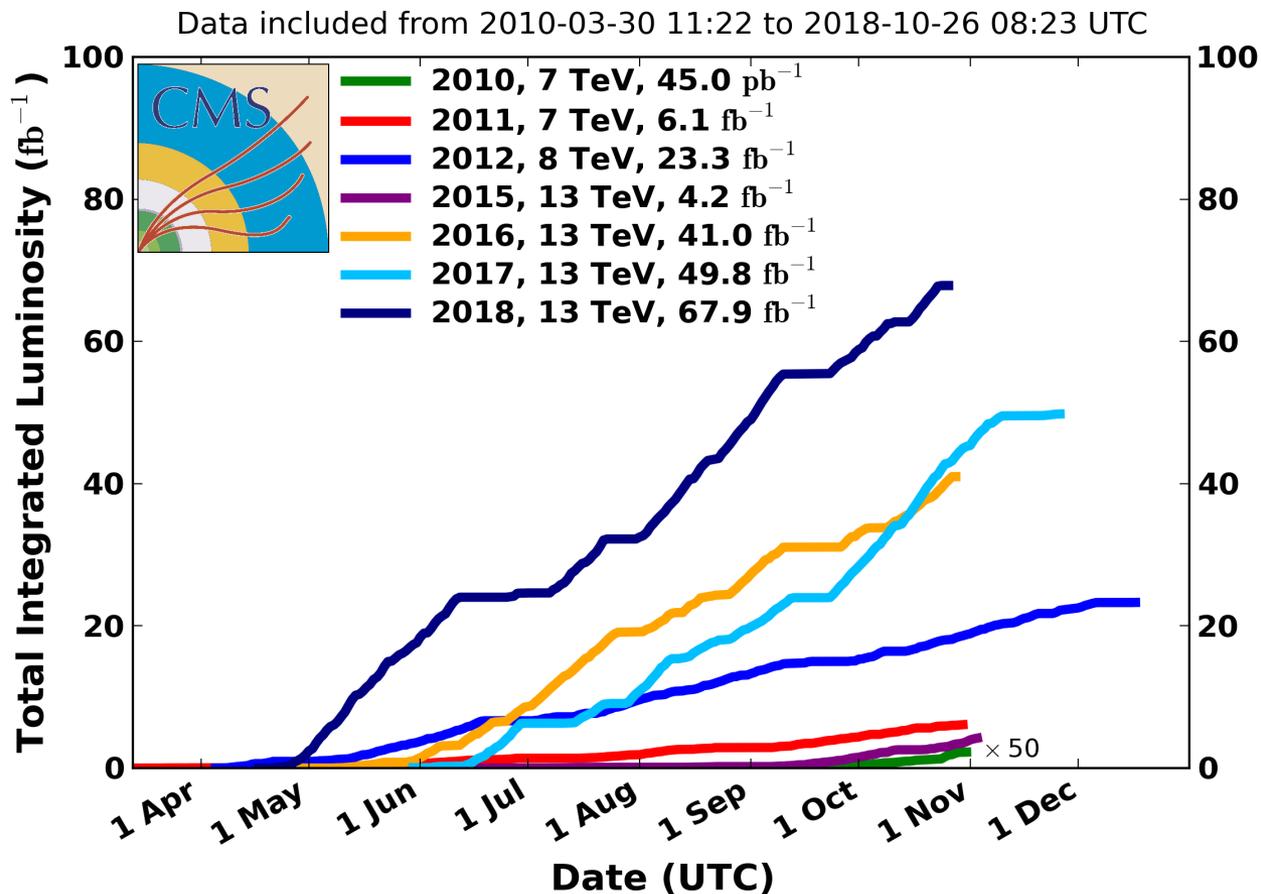
$$\sigma E/E \approx 60-100\%/\sqrt{E(\text{GeV})} \oplus 3\%$$

Muon Spectrometer

$$\sigma pT/pT < 10\% \text{ up to } 1 \text{ TeV muons}$$



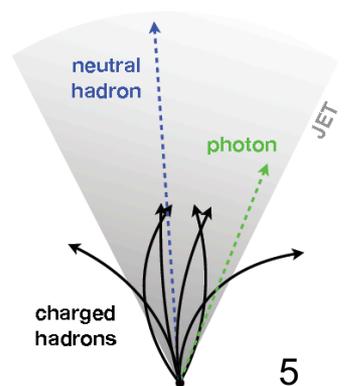
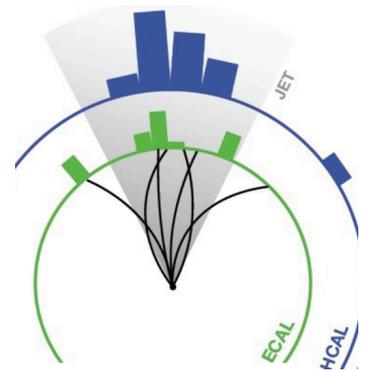
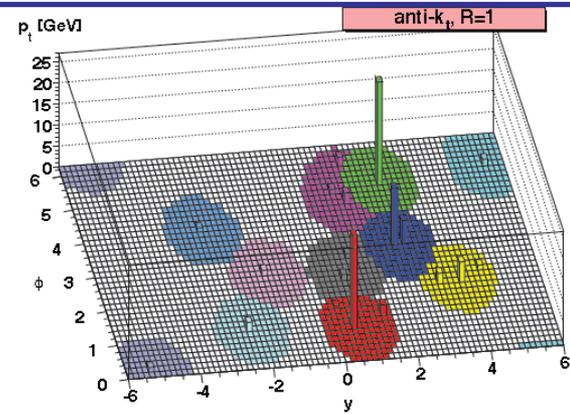
CMS Integrated Luminosity Delivered, pp



- LHC Accelerator had so far a superb performance.

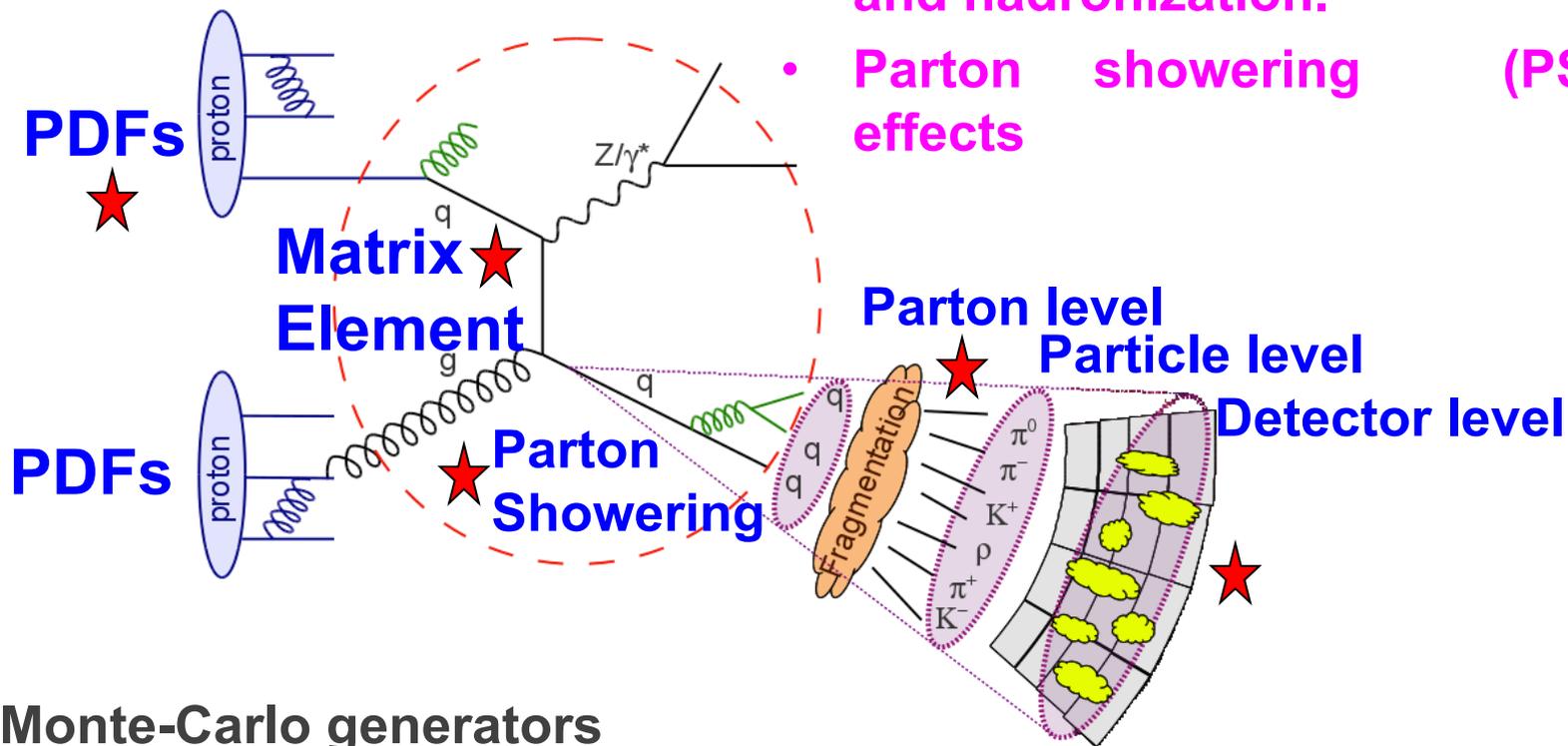


- **Most analyses use the Anti-kt clustering algorithm** : with a cone $R = 0.4$ and 0.8 for CMS and 0.4 and 1.0 for ATLAS, which is infrared and collinear safe, geometrically well defined, and tends to cluster around the hard energy deposits.
- **Calorimeter Jets** : Clustering of Calorimeter Towers composed of ECAL and HCAL energy deposits.
- **Particle Flow Jets**: Clustering of Particle Flow candidates constructed combining information from all sub-detector systems.
- **Jet taggers and jet - substructure** : several jet taggers and jet substructure techniques developed.



Sources of uncertainties★

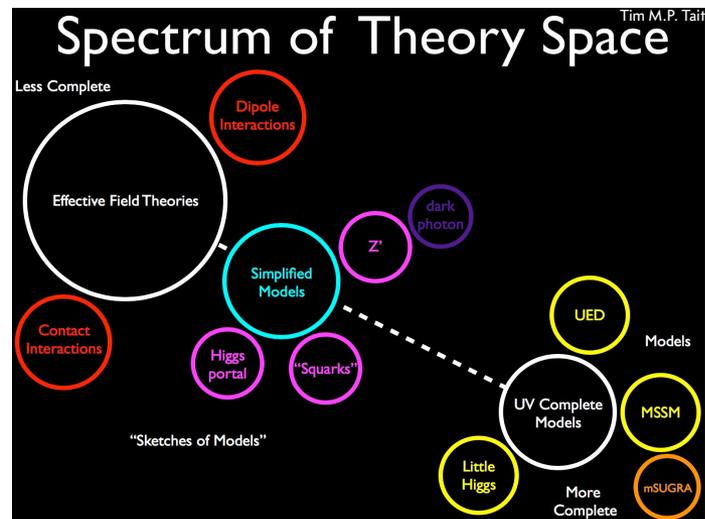
- **Perturbative QCD calculations @ NLO**
- **Non-perturbative corrections for multi-parton interaction and hadronization.**



- **LO QCD Monte-Carlo generators**



Dark Matter Searches at the LHC : A key part of the physics search program

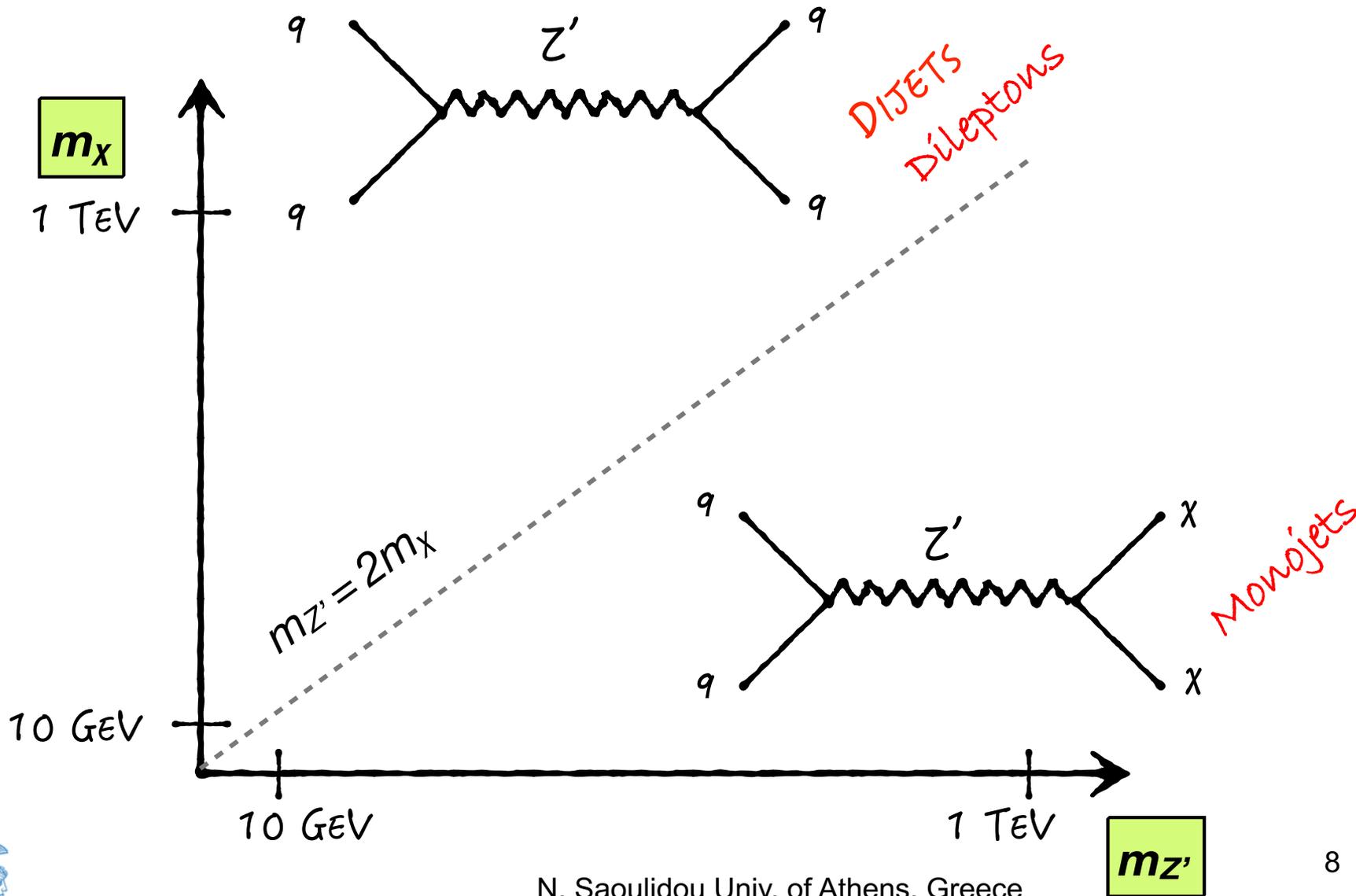


- There is plenty of evidence for the existence of **Dark Matter** which we have only **seen** so far **gravitationally**.
- **Direct Dark Matter searches:** Detect interactions of DM particle (or particles) with terrestrial detectors
- **Indirect Dark Matter searches:** Detect DM-DM interactions in the cosmos, ie DM-DM interactions at the centre of the galaxy
- **Collider Searches : Produce DM and DM mediators in the Lab**

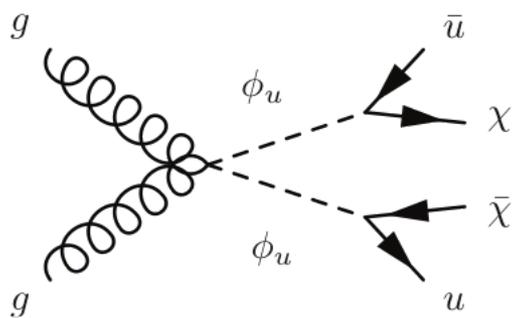


Dark Matter Searches @ the LHC :

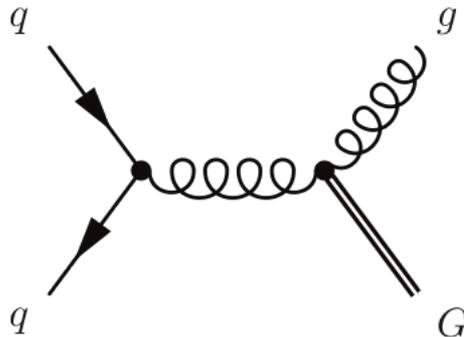
Complementarity between difference searches



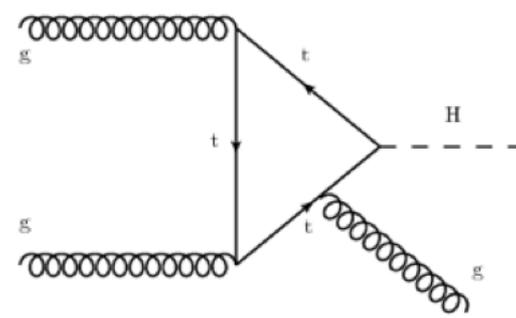
Access a broad range of new physics hypothesis



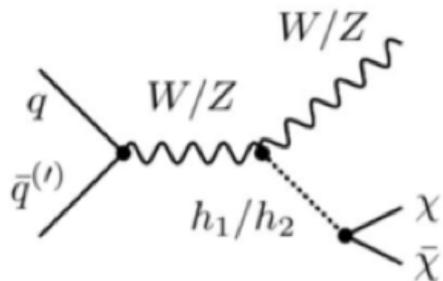
Dark Matter



Gravitons (ADD model)



Higgs -> invisible



Dark Energy

- **Main Experimental Signature :**
 - **A jet or a Vector Boson decaying hadronically and Missing transverse energy**

- **Main irreducible backgrounds:**
 - **Z → vv+jets, W → lv+jets**



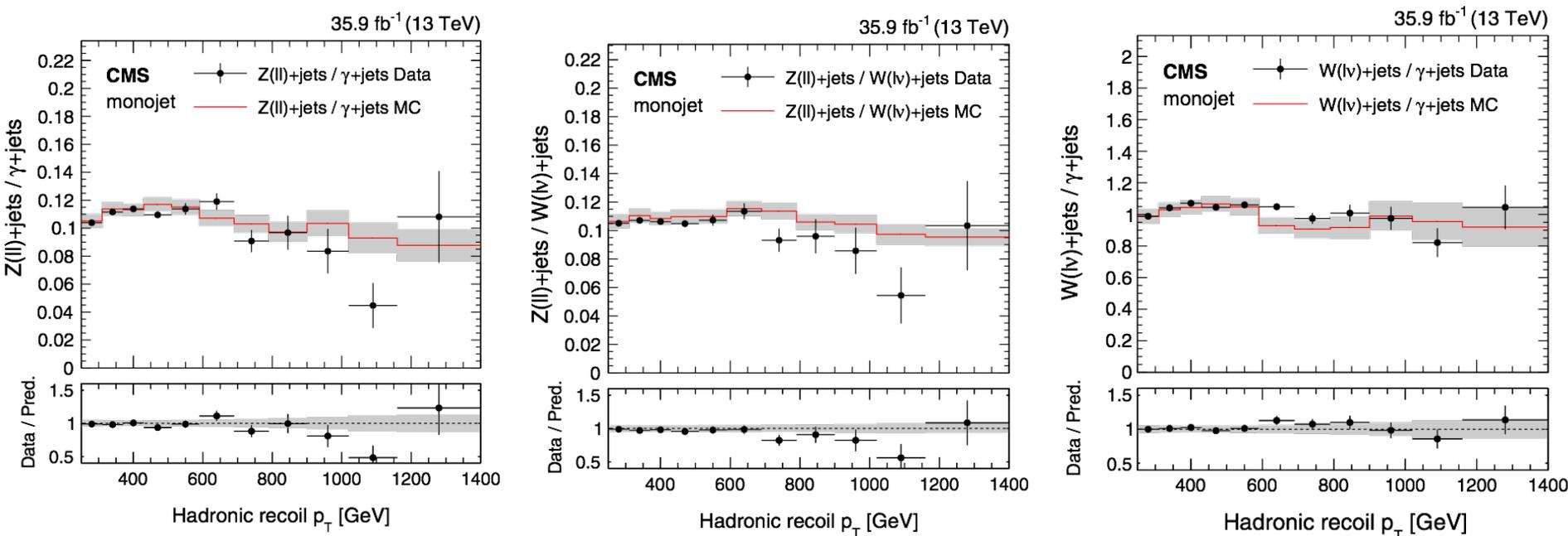
Background estimation

- **Background estimation** : Data-driven, in order to minimize dependence on simulation, and hence **significantly reduce theoretical systematic uncertainties**.
- Five **control regions** are used, $Z \rightarrow ee$, $Z \rightarrow \mu\mu$, $W \rightarrow ev$, $W \rightarrow \mu\nu$, **top (ATLAS) or γ +jets (CMS)**, to measure the E_{Tmiss} spectrum excluding leptons and photons. The main assumption, for the CMS approach, is that Z+jets and γ +jets behave similarly.
- **Transfer factors, from simulation** are used to predict the backgrounds in the signal region. These, especially at the high end of the spectrum, exhibit theoretical uncertainties due to:
 - QCD renormalization and factorization scales ★
 - QCD effects in γ +jets production being different than those in W+jets and Z+jets ones. ★
 - Mixed QCD-EW corrections ★



Transfer Factors from simulation

Physical Review D 97, 092005 (2018)

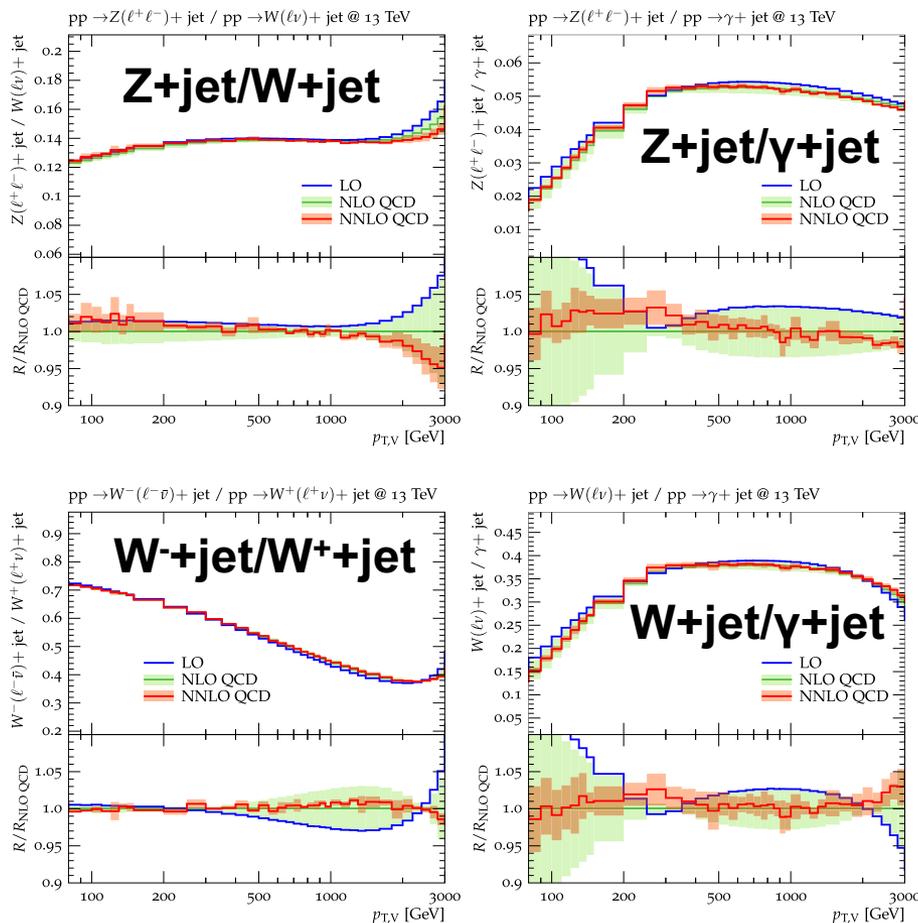


- The irreducible $\mathbf{Z} \rightarrow \mathbf{vv+jets}$ background could be measured from the $\mathbf{Z} \rightarrow \mathbf{ll+jets}$ events.
- Due to low statistics, gamma+jets and $\mathbf{W} \rightarrow \mathbf{lv+jets}$ events are used instead and the Z/gamma and W/gamma ratios from simulation.



Guidance from Theory

Eur. Phys. J. C (2017) 77:829



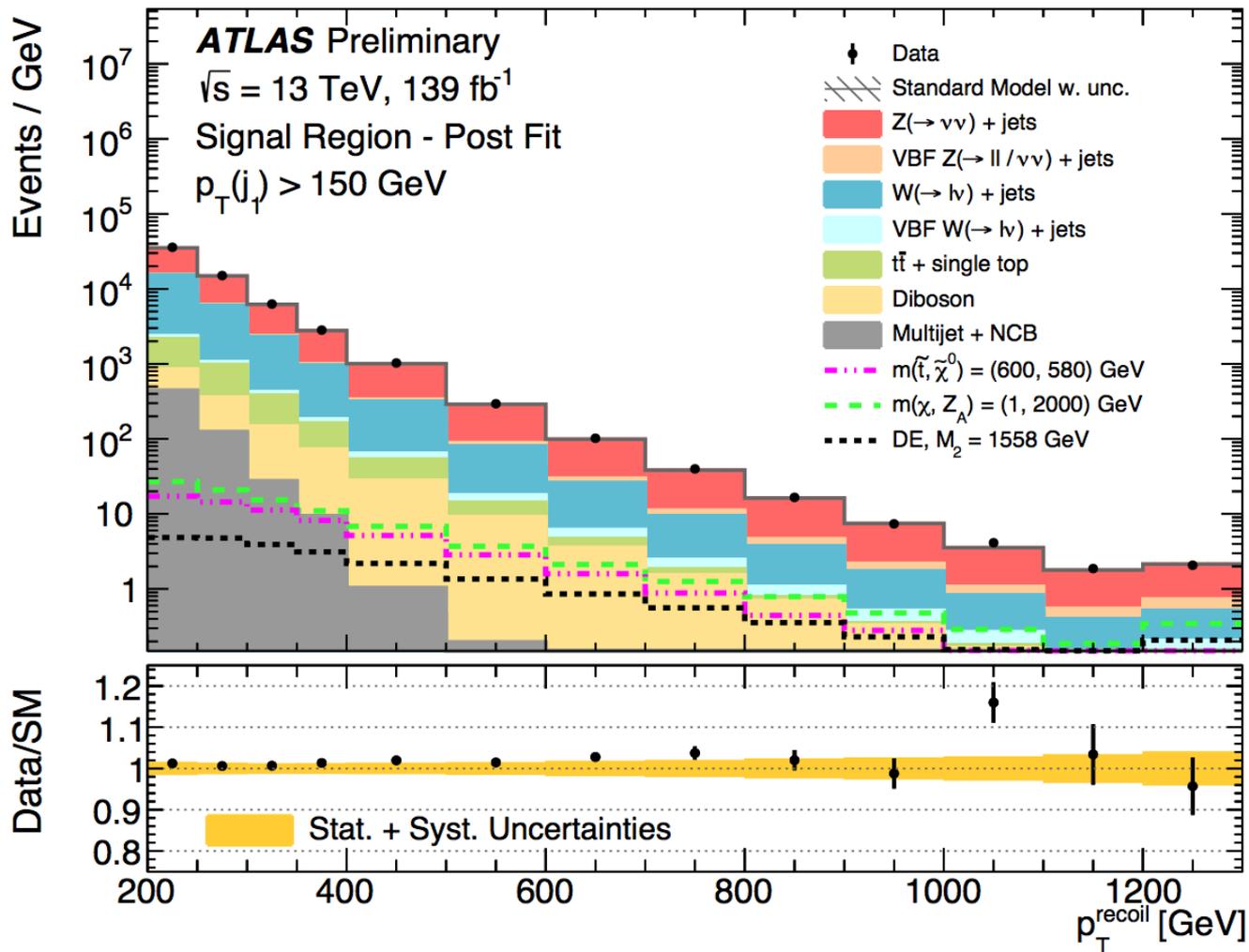
Simulated Transfer factors

Provided accurate theoretical predictions and detailed prescriptions on how theoretical systematic uncertainties should be taken into account in the experimental analyses.

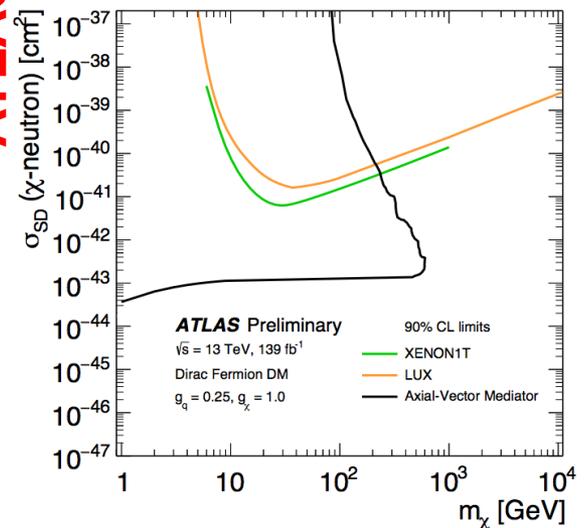
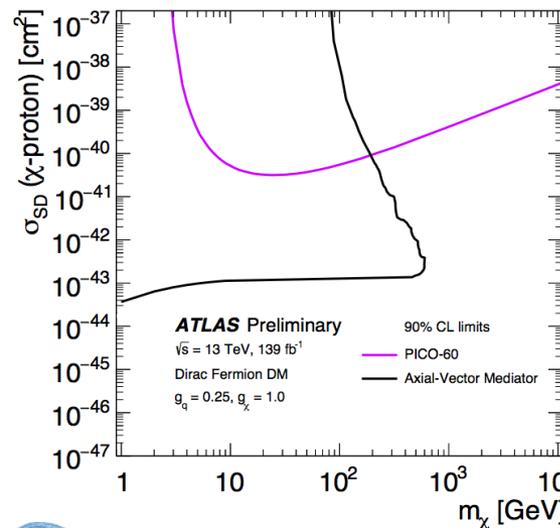
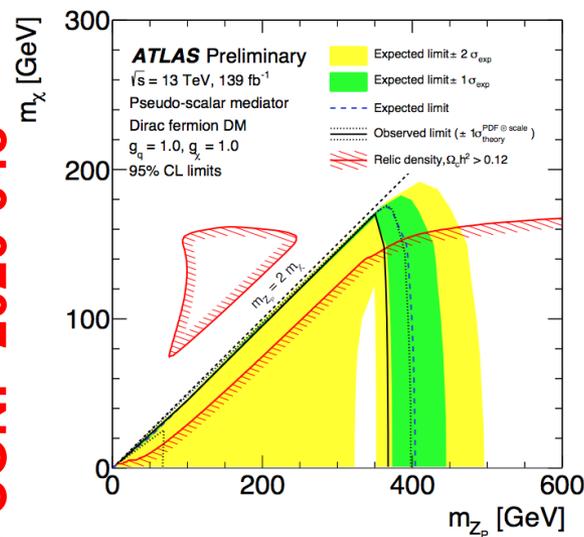
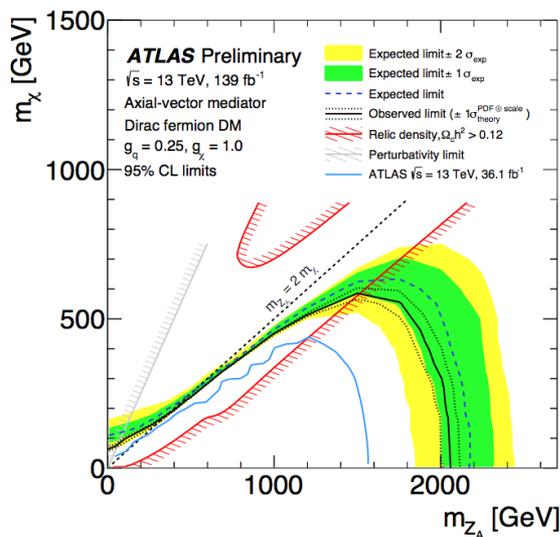


Theoretical uncertainties (QCD+EW) account for \sim half of the 5% uncertainty at the high end of the missing E_T spectrum.

ATLAS-CONF-2020-048



Interpretations

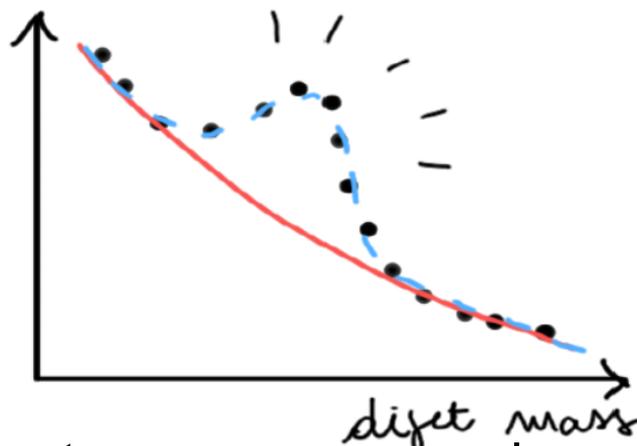
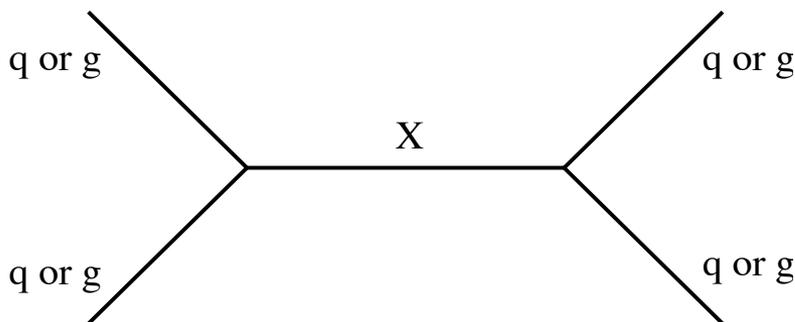


ATLAS-CONF-2020-048

This search is one of the most “direct” probes for dark matter at the LHC! Expecting further improvements in theoretical (and experimental) uncertainties in the future?



Access a broad range of new physics hypothesis



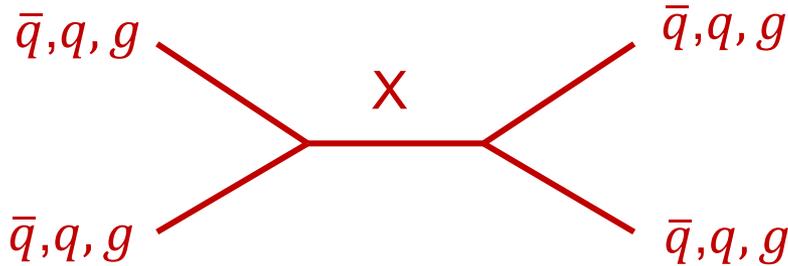
- **Powerful:** LHC is a dijet resonance factory at a new energy scale
- **Broad:** search for many sources of new physics in a single simple search
 - String resonances** from string theory
 - Excited quarks** from theories of quark compositeness
 - W', Z' and scalar diquarks** from grand unified theory
 - Gravitons** from the Randall-Sundrum model of extra dimensions
 - Axiguons, Colorons and Color Octet Scalars** from other new models
 - DM mediators**
- **Model Independent:** Search results are applicable to any model of narrow qq, qg, or gg resonances.



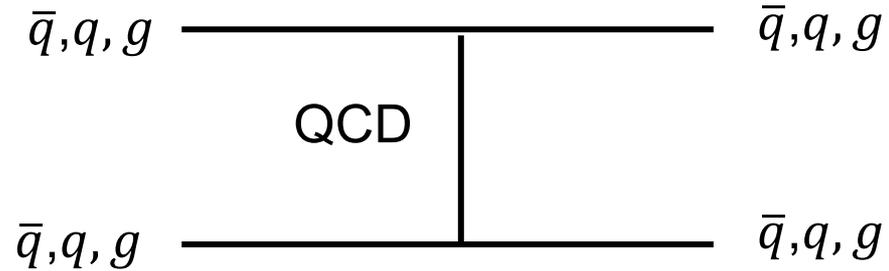
Dijet searche in a nutshell

Access a broad range of new physics hypothesis

Resonance Signal



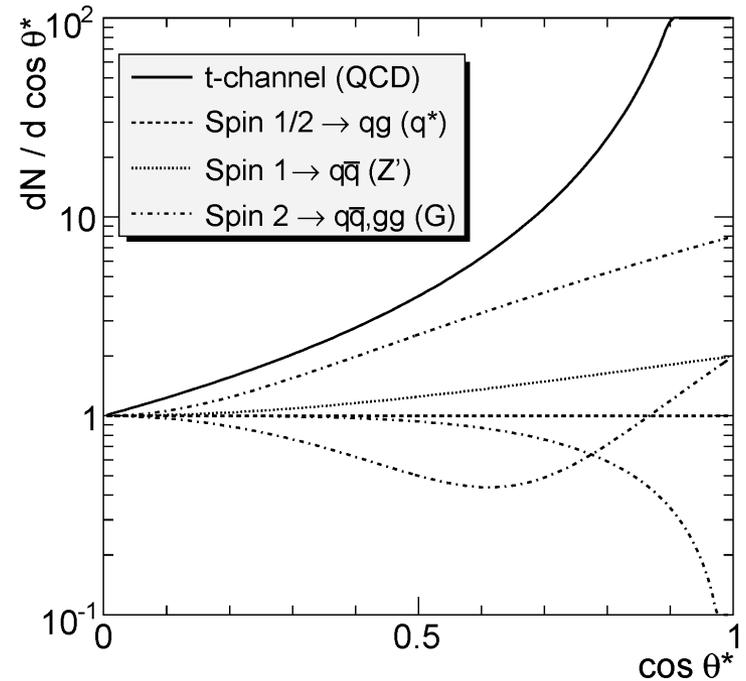
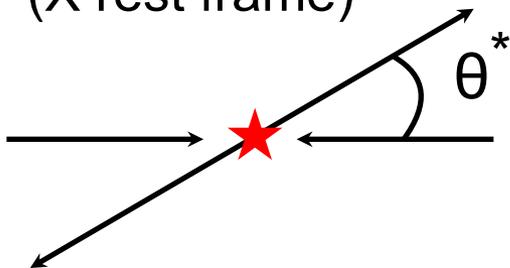
SM Background



Physics observables
 $M(jj) \rightarrow$ Resonance Mass
 $\Delta\eta(jj) \rightarrow$ Resonance Spin

$$\Delta\eta_{12} = |\eta_{jet1} - \eta_{jet2}| = \ln \frac{1 + |\cos\theta^*|}{1 - |\cos\theta^*|}$$

(X rest frame)

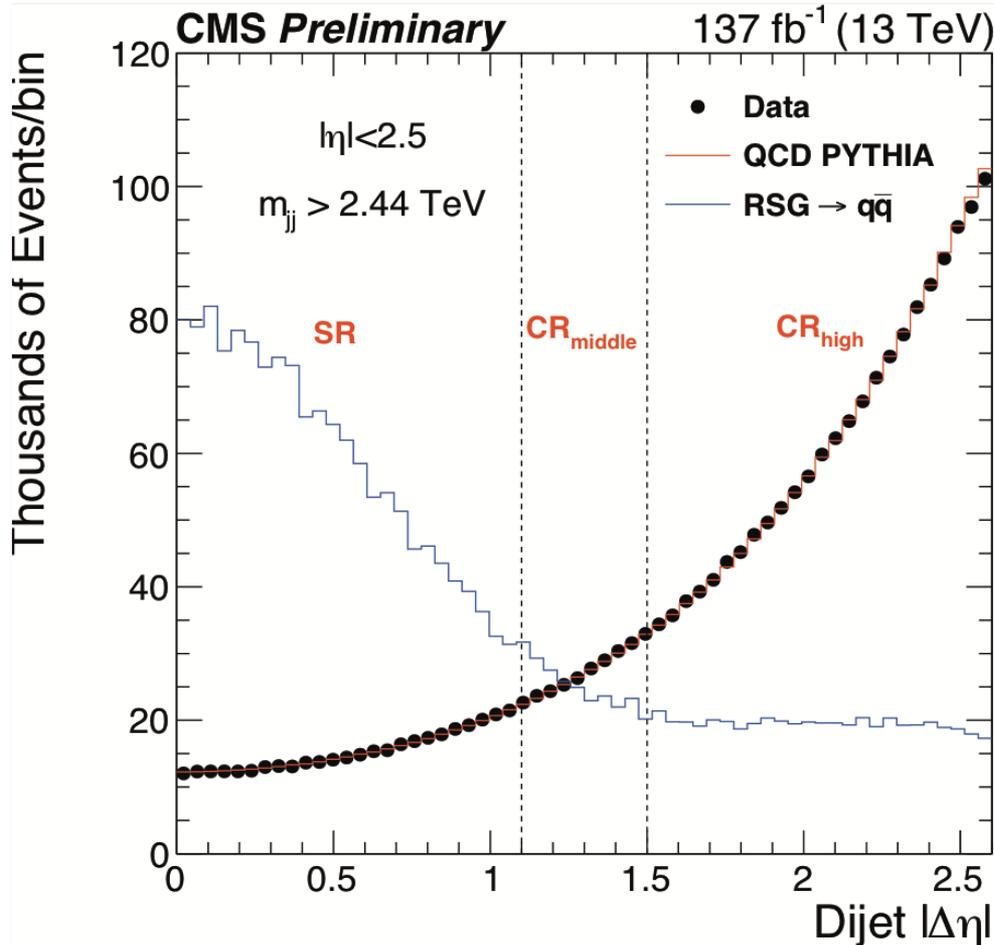


JHEP 05 (2020) 033

- **Traditional Background Modeling (used in the past 25 years)** : A fit with **empirical parametrizations** is performed to the data, with **its parameters treated as unconstrained (in CMS) nuisance parameters** in the final limit setting and significance estimation procedure.
- **Advantages** : Well tested, well understood, significant expertise gained in many years.
- **Disadvantages** : Less robust and less performant in searches for wide resonances, or ones that might appear at the high mass tail of the distribution where not enough data exist to constrain the parametric fit.



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- The signal region (SR) is defined with $|\Delta\eta| < 1.1$.
- Two control regions are defined
 - **CR_{middle} with $1.1 < |\Delta\eta| < 1.5$** used to estimate the corrections to the simulated transfer factor, constraining systematic uncertainties
 - **CR_{high} with $1.5 < |\Delta\eta| < 2.6$** used to predict the dijet mass distribution of the QCD background



QCD Prediction for the invariant mass of the dijet system

$$N_{SR}^{\text{Prediction}} = R \times N_{CR_{\text{high}}}^{\text{Data}}$$

Transfer factor

$$R = C \times N_{SR}^{\text{Simulation}} / N_{CR_{\text{high}}}^{\text{Simulation}}$$

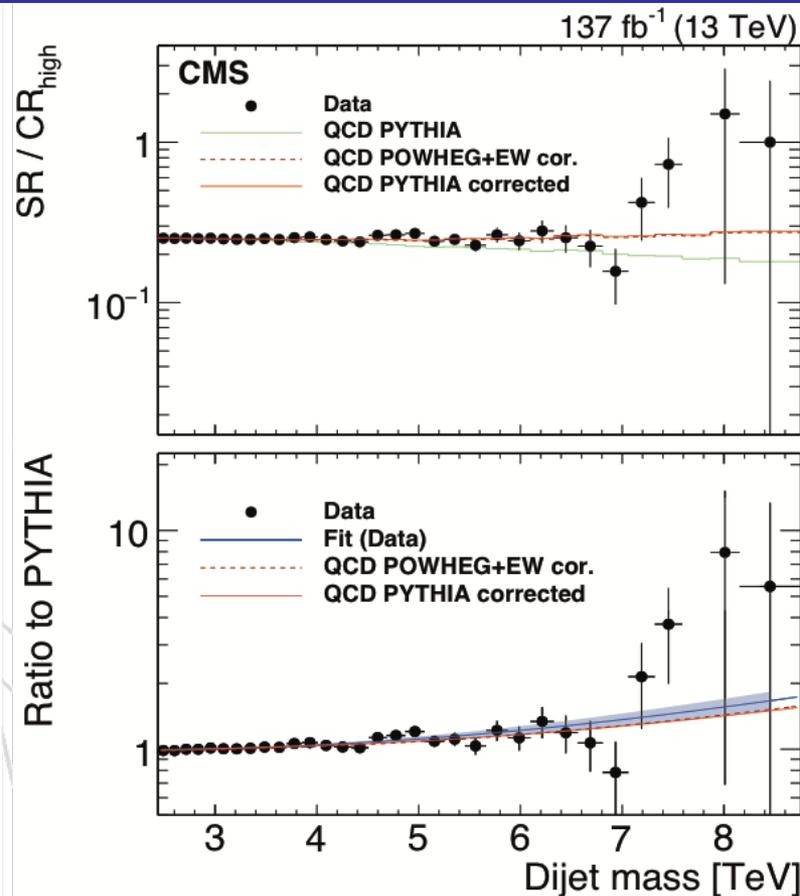
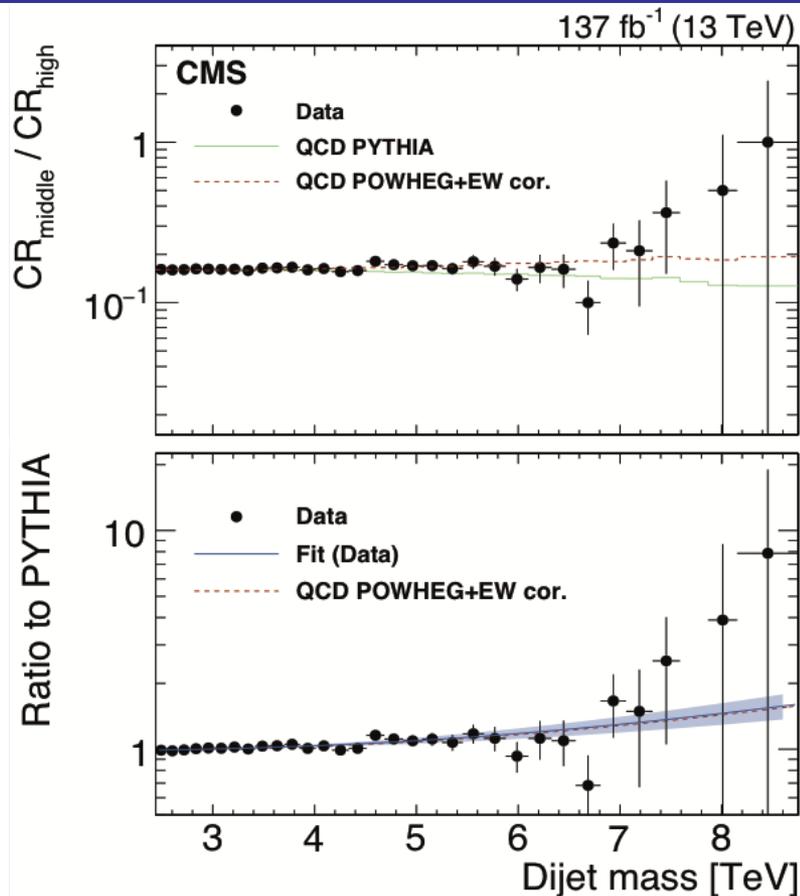
Transfer factor correction

$$R_{\text{aux.}} = N_{CR_{\text{middle}}} / N_{CR_{\text{high}}}$$

$$C = \frac{R_{\text{aux.}}^{\text{Data}}}{R_{\text{aux.}}^{\text{Simulation}}} = p_0 + p_1 \times (m_{jj} / \sqrt{s})^3$$

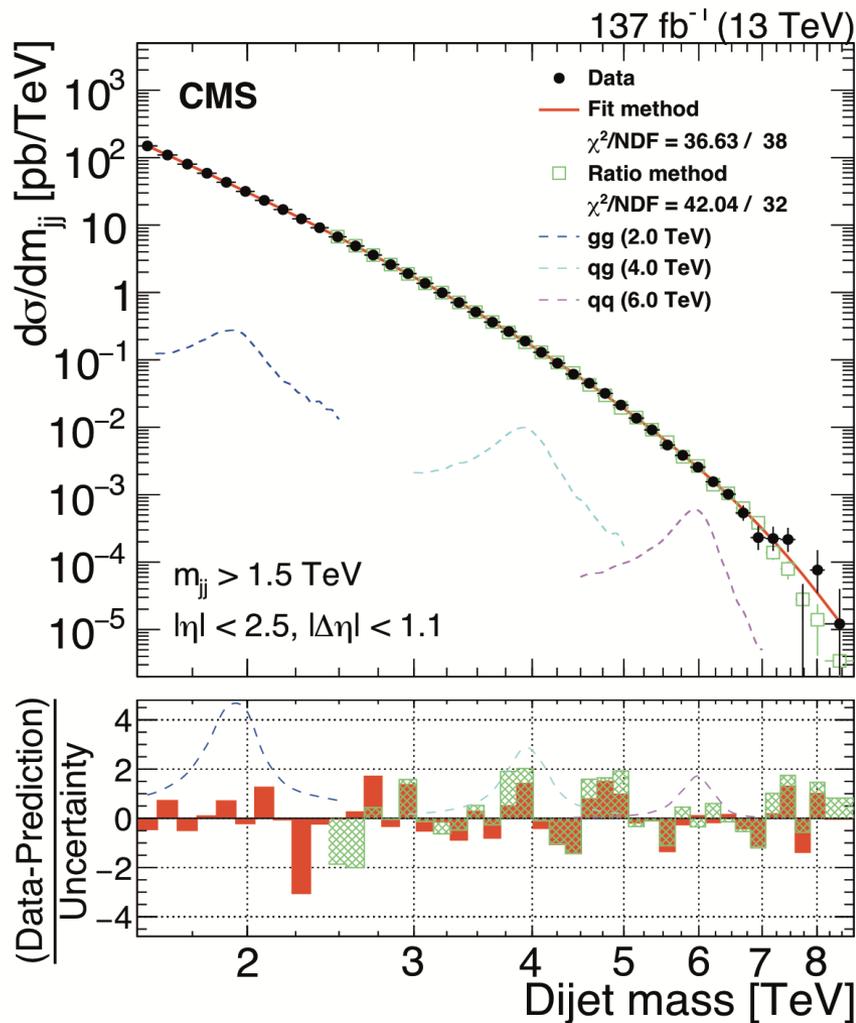
- N is the number of events in each bin of the **data dijet mass distribution**, and **R and R_{aux} the transfer factors from simulation.**
- We **constraint theoretical and experimental systematic effects**, responsible for the **differences between the data and simulated transfer factors**, in a **data-driven way using CR_{middle}**
- For the final result a **simultaneous fit** is performed in **SR, CR_{middle} and CR_{high}** ,





- **Transfer factors** are **flat** as a **function of dijet mass**, and **very similar** between **data and simulation** ; both the main and the auxiliary ones.
- **Missing higher QCD orders** from the simulation and **electroweak effects** are **well described by the functional form used**, and can account for the differences seen between data and simulation. ★





- **Fitting** : Modified frequentist CL_s used for limit setting, performing a binned fit with a background and signal template.

- **Systematic uncertainties:**

- **Related to signal modeling:**

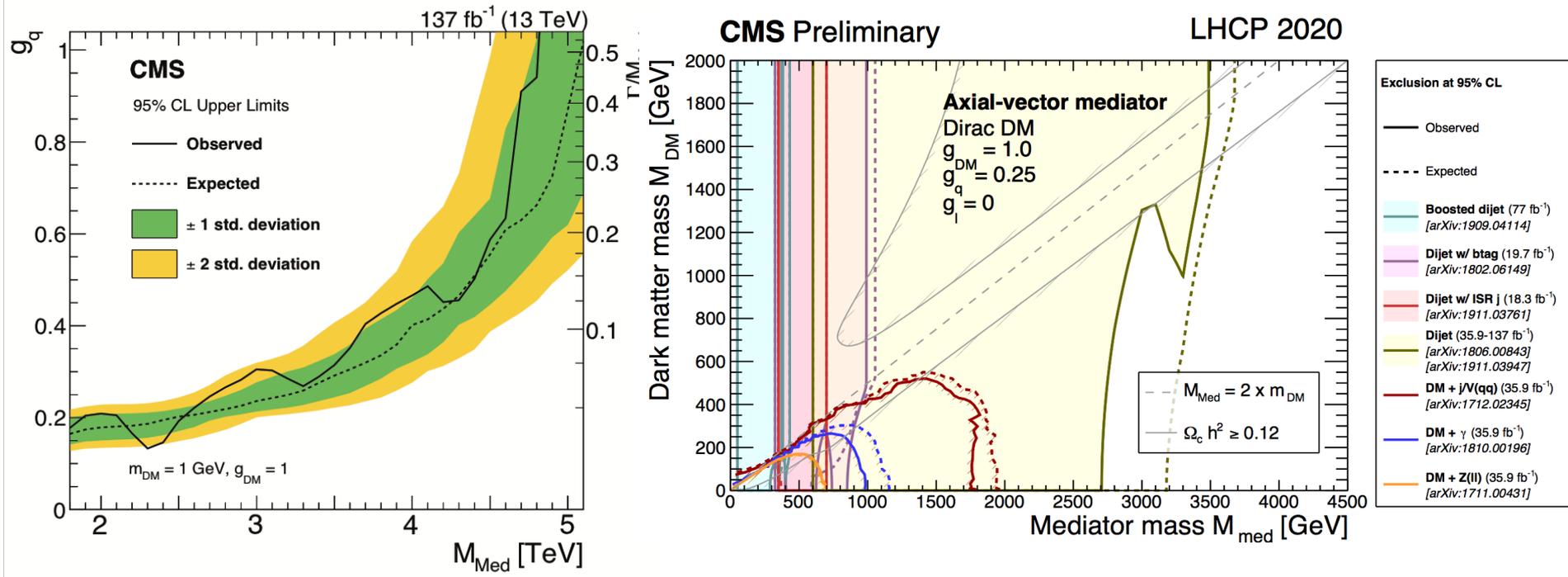
luminosity (log-normal constraints), jet energy scale and jet resolution (gaussian constraints).

- **Related to background modelling:**

Missing higher QCD orders and EW effects, treated as freely floating automatically evaluated via profiling. ★

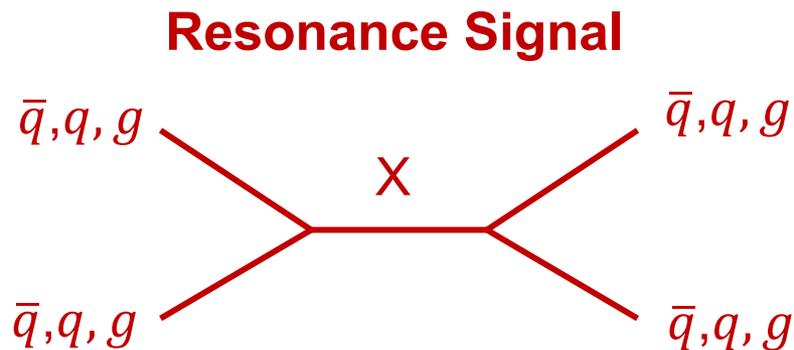


Narrow and Wide resonances : DM mediator limits

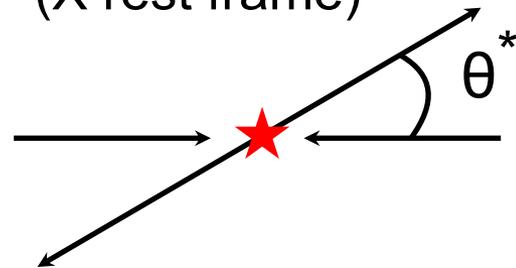


- Can further improve with more data and a better, more precise, QCD NLO simulation including also EW effects.★





(χ rest frame)



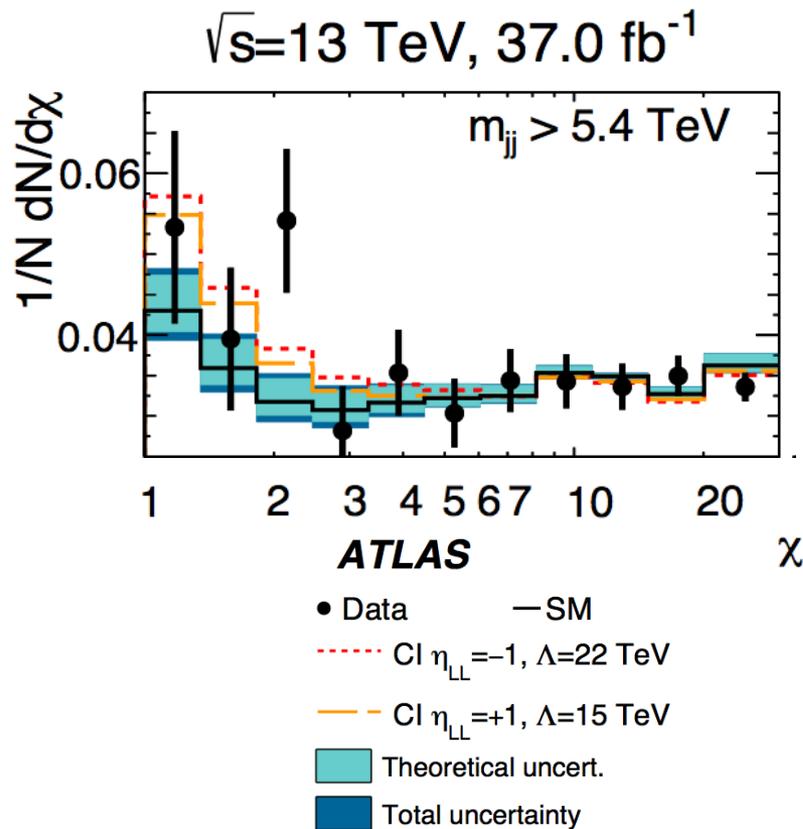
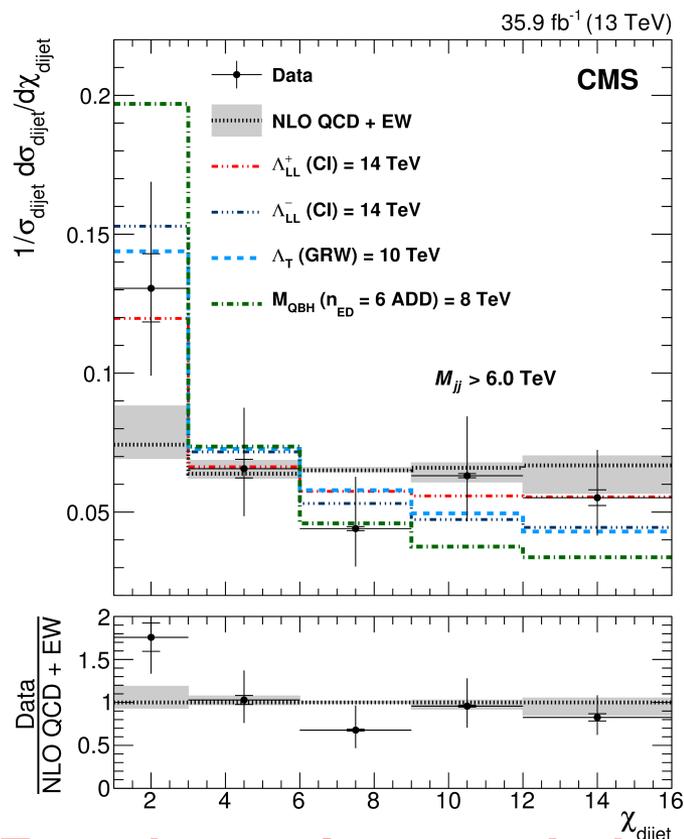
$$\chi = e^{|y_1 - y_2|} \approx \frac{1 + |\cos \theta^*|}{1 - |\cos \theta^*|}$$

- χ chosen since **main QCD background is flat.**
- **Unique analysis with coupling limits extending to higher DM mediator masses and larger DM mediator widths.**
- Additional **limits** are placed to a **variety of new physics models like quark contact interactions, extra spatial dimensions, quantum black holes, and Dark Matter models.**



Eur. Phys. J. C 78 (2018) 789

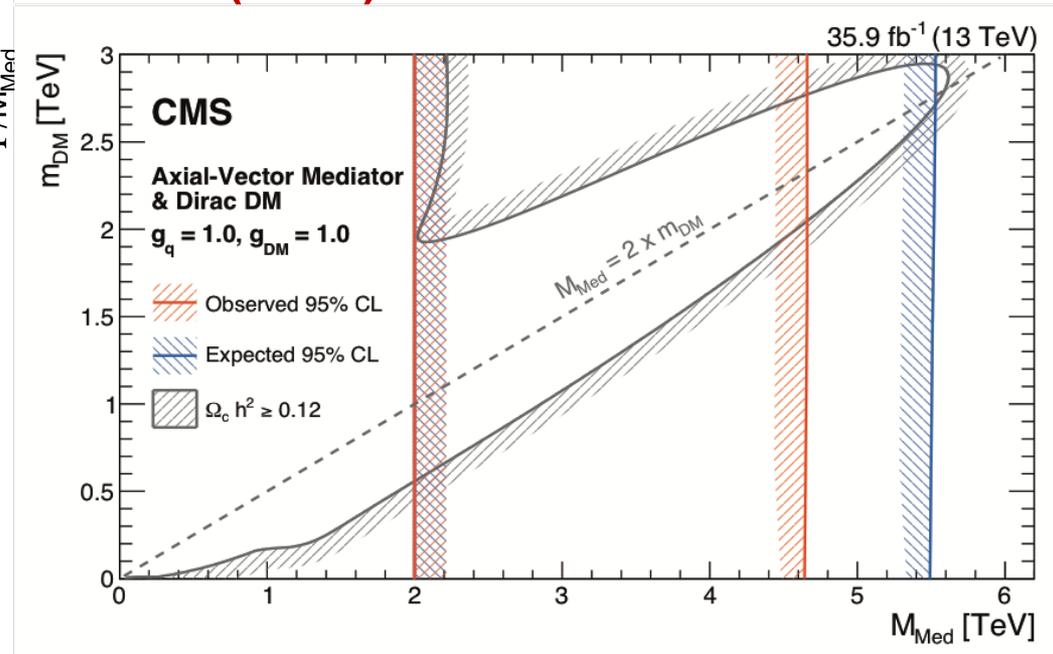
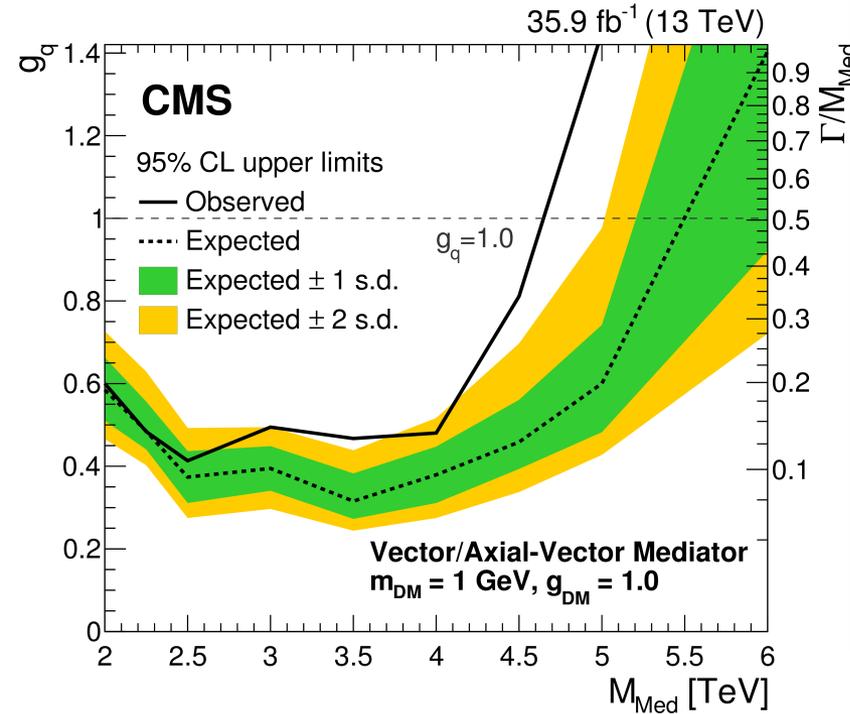
Phys. Rev. D 96 (2017) 052004



- **Experimental uncertainties** dominated by **jet resolution and relative JES** (absolute cancels)
- **Theoretical uncertainties dominated by non perturbative corrections and renormalization and factorization scales.** ★



Eur. Phys. J. C 78 (2018) 789



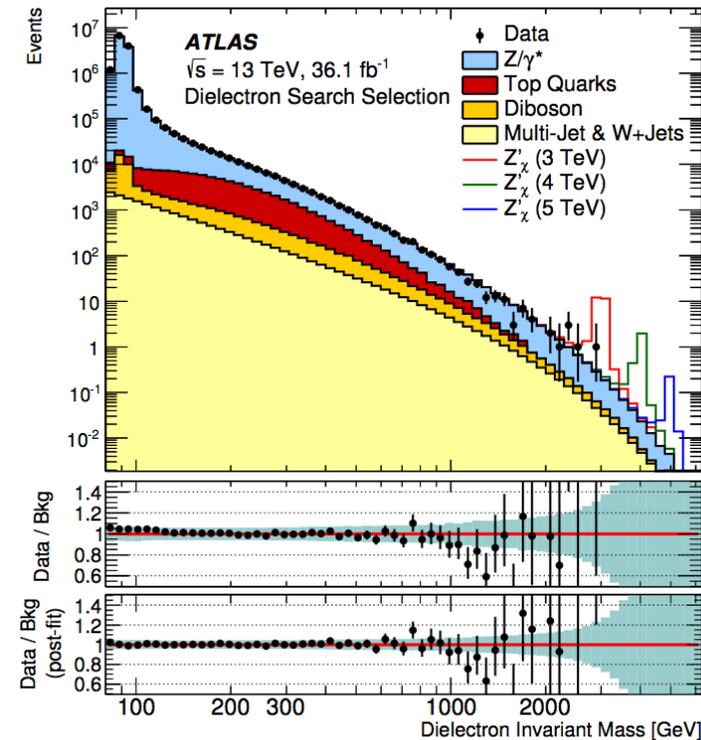
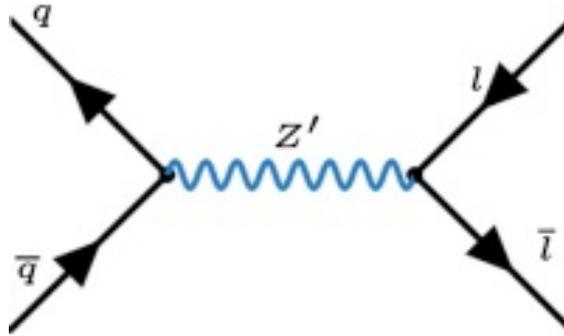
- This analysis is the **only one** being able to access so large resonance widths.
- For full Run II result changed background estimate from NLO to NNLO QCD prediction, which significantly reduces uncertainties★



Phys. Lett. B 796 (2019) 68

ATLAS-EXOT-2019-16

JHEP 10 (2017) 182



- Search for **resonant and non-resonant signals** in final states **with a pair of leptons**, coming from **a variety of new physics models**:
 - Grand unified theories (GUTs)
 - Superstrings
 - Left-right-symmetric models (LR)
 - Dark Matter mediators
 - Contact Interactions



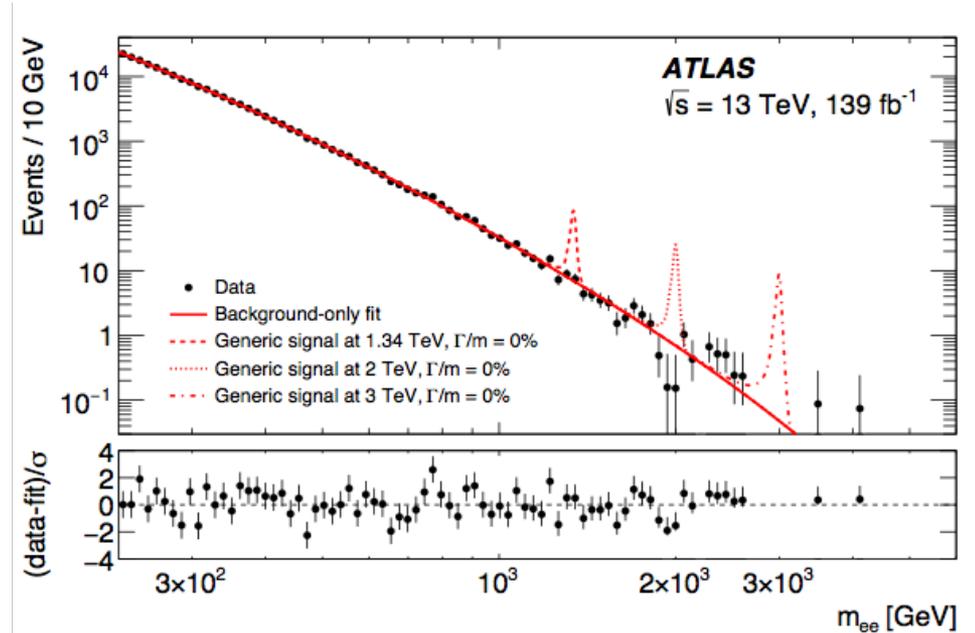
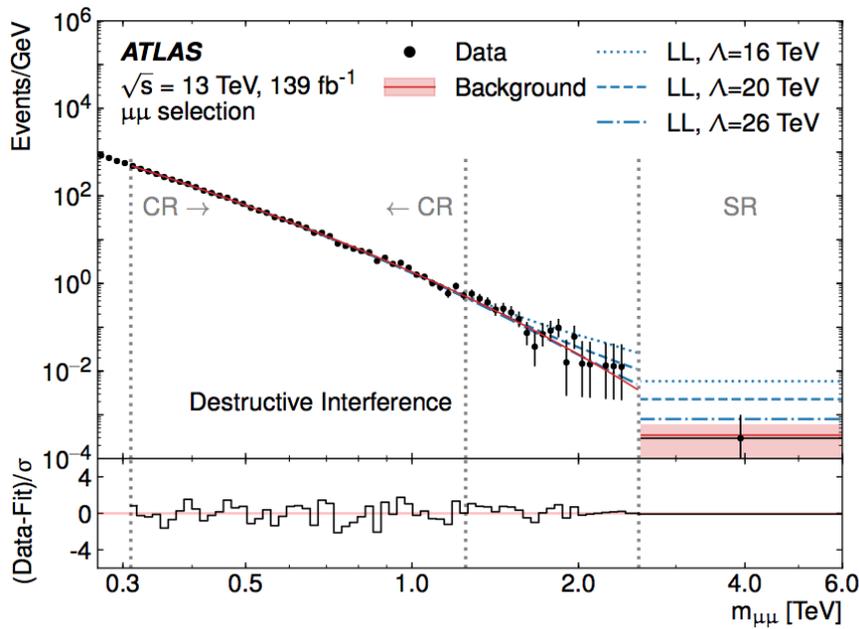
Phys. Lett. B 796 (2019) 68
ATLAS-EXOT-2019-16

JHEP04(2019)114
CMS-PAS-EXO-19-019

- **Resonant analysis strategy :**
 - To estimate the background **fit the simulated dilepton spectrum with an empirical function**, and use the **uncertainties on the arbitrary function parameters as constraints in the data fit.**
- **Non resonant analysis strategy :**
 - To estimate the background **fit the low mass end of the dilepton spectrum and extrapolate to the high end (ATLAS)**, or use **Bayesian inference using the simulation for the background template (CMS)**



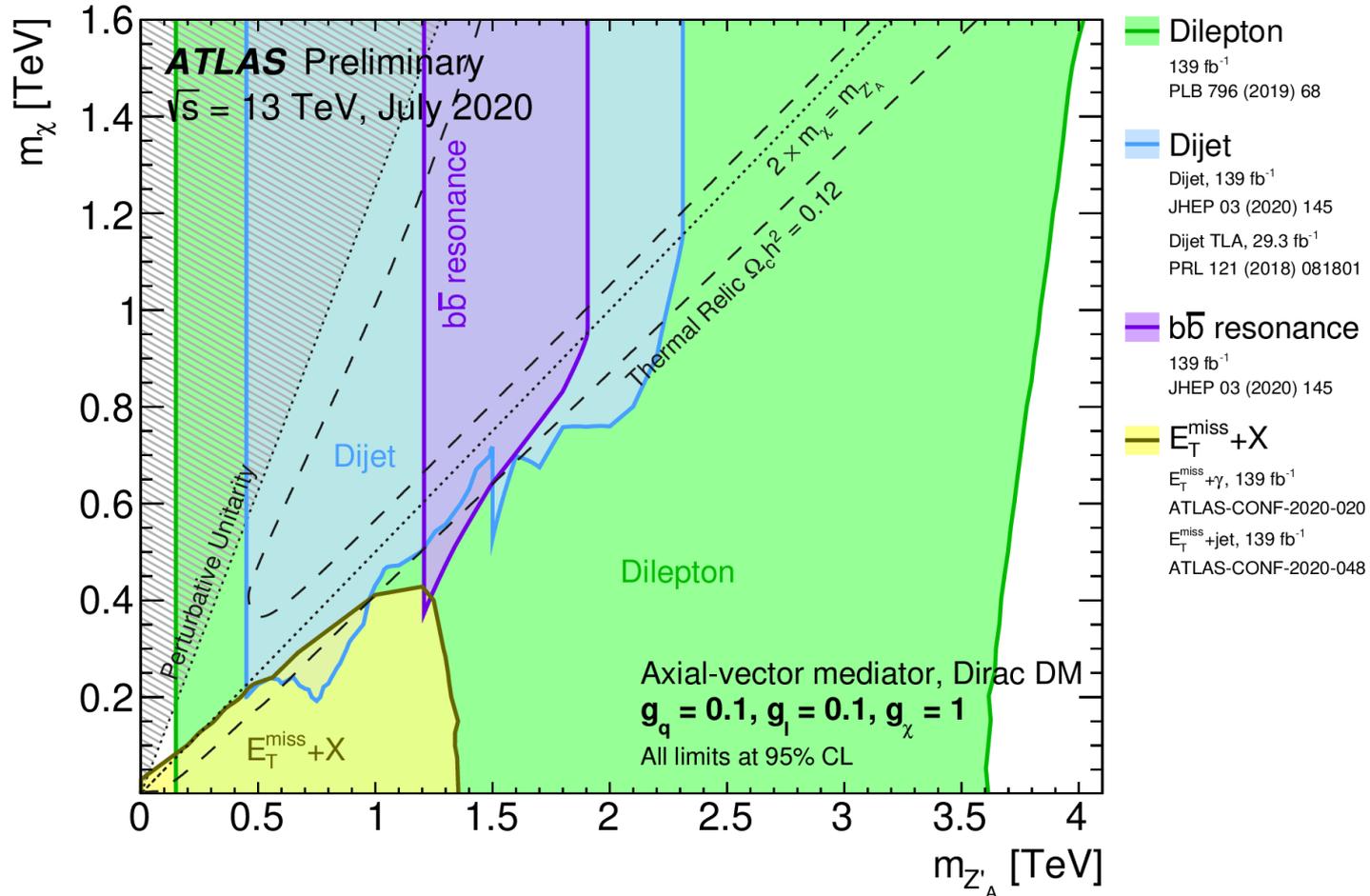
Phys. Lett. B 796 (2019) 68
ATLAS-EXOT-2019-16



- **Main (or among largest) uncertainties in the non-resonant (resonant) search : PDF uncertainties!** ★

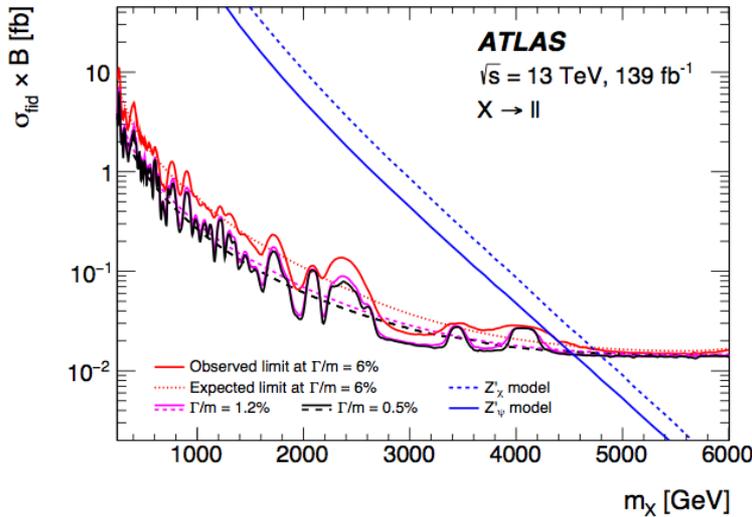


If DM mediator is leptophilic dilepton searches play a key role

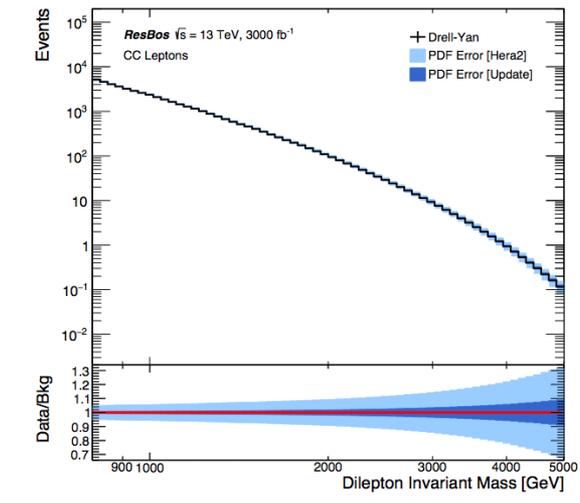
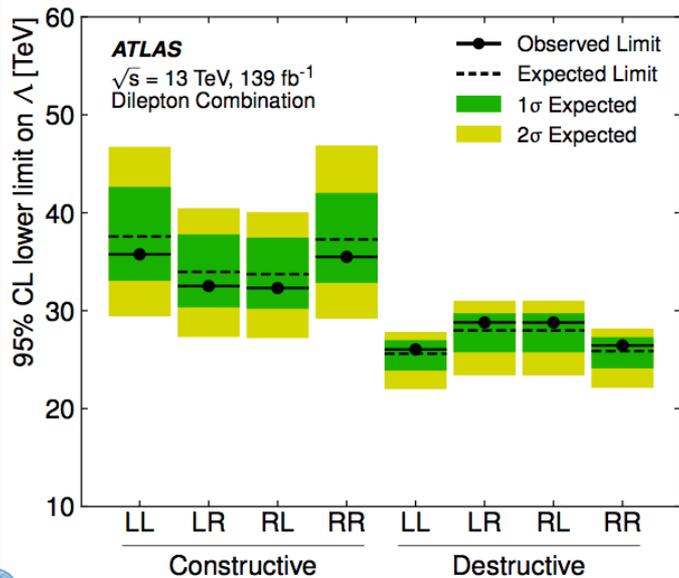


Phys. Lett. B 796 (2019) 68

ATLAS-EXOT-2019-16



- Reduction of PDF uncertainties will increase the sensitivity of these searches! ★
- A nice suggestion already : *Phys.Rev.D 99 (2019) "A new method for reducing PDF uncertainties in the high-mass Drell Yan spectrum"*.



Looking Inside Jets \Leftrightarrow

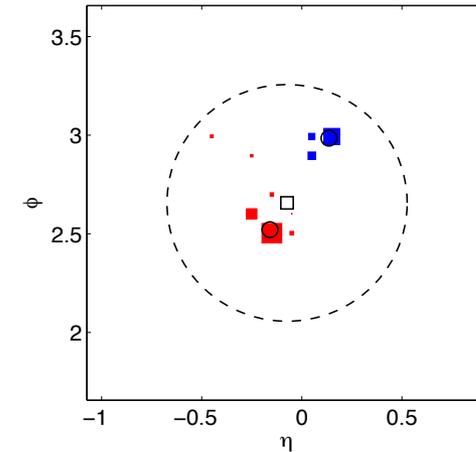
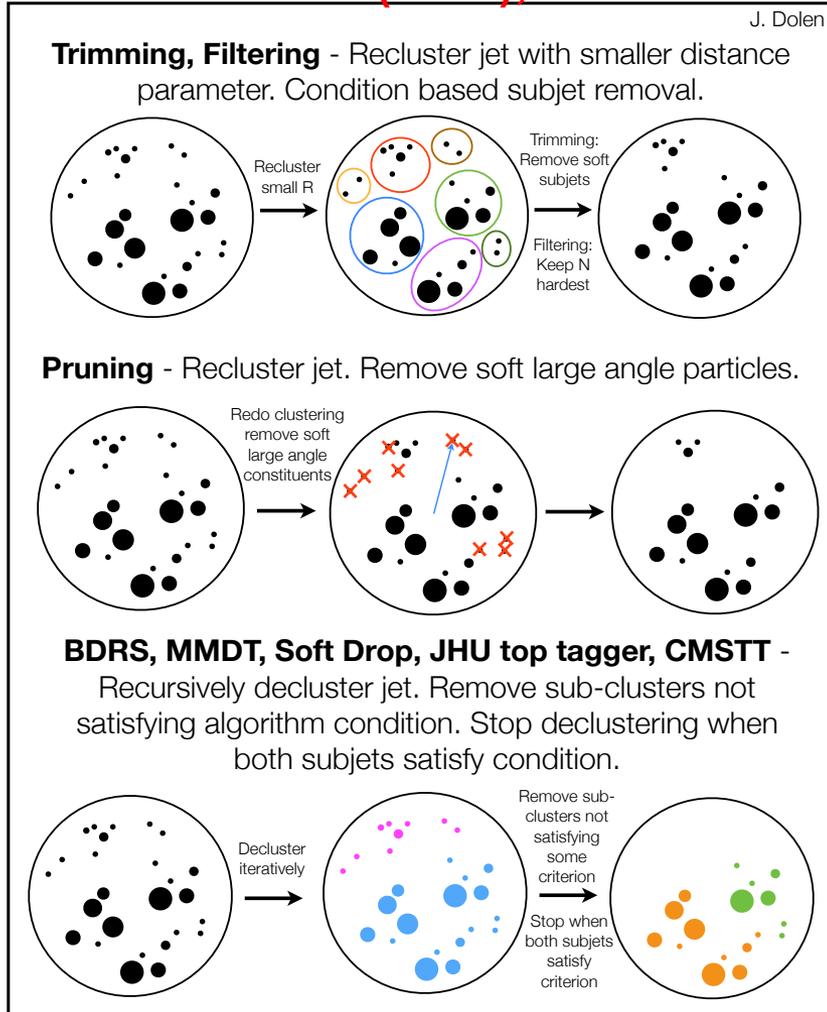
Sources of uncertainties

Phys.Rev.D80,05150 (2009)

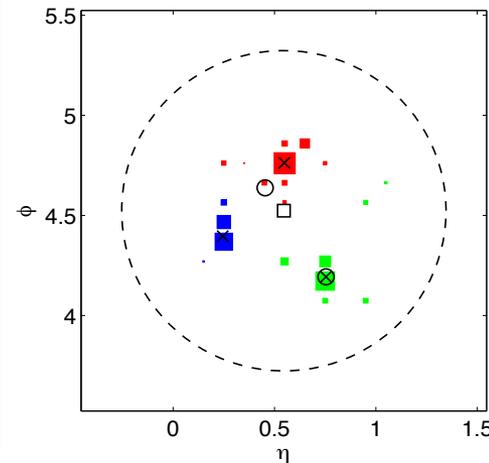
JHEP05(2014),146

JHEP 1103:015,2011

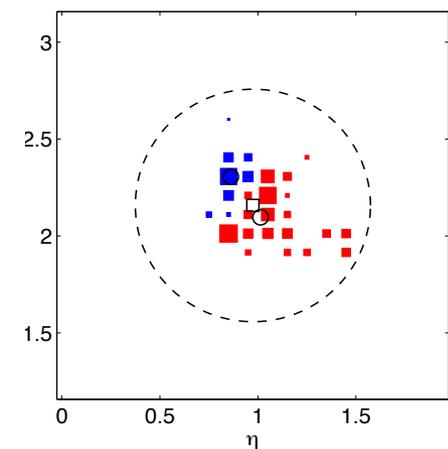
W jet, $\tau_2/\tau_1=0.15$



Top jet, $\tau_3/\tau_2=0.21$



QCD jet, $\tau_2/\tau_1=0.73$

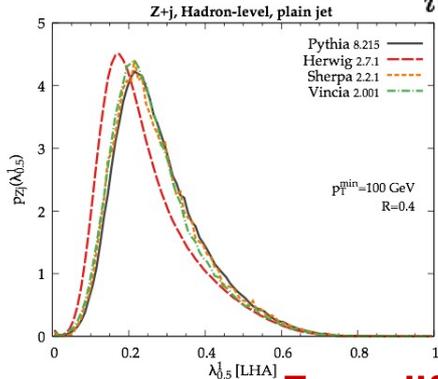


Quark – Gluon Separation

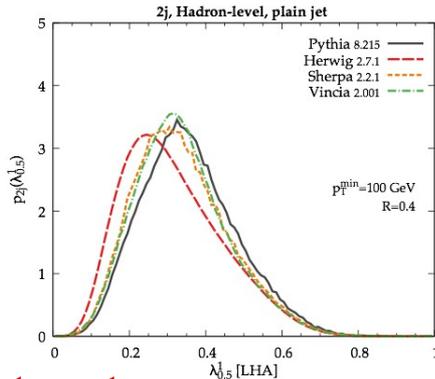
Quark-gluon discriminant (angularity)

$$\lambda_{\beta}^{\kappa} = \sum_{i \in \text{jet}} z_i^{\kappa} \theta_i^{\beta},$$

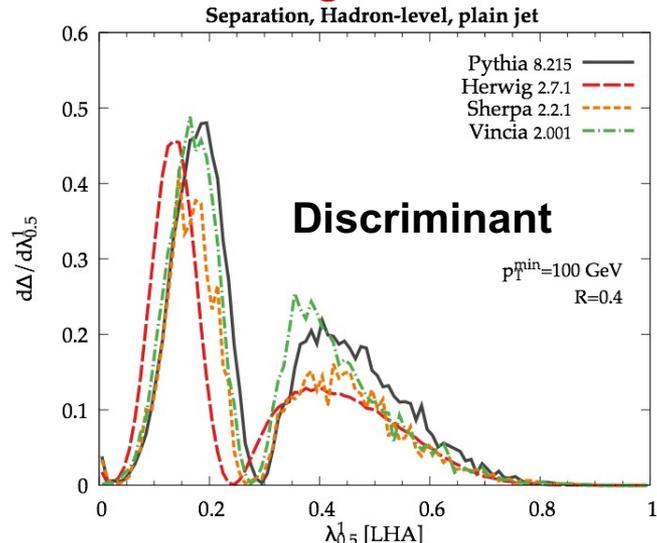
Quark enriched



Gluon enriched



Four different parton shower generators



(c)

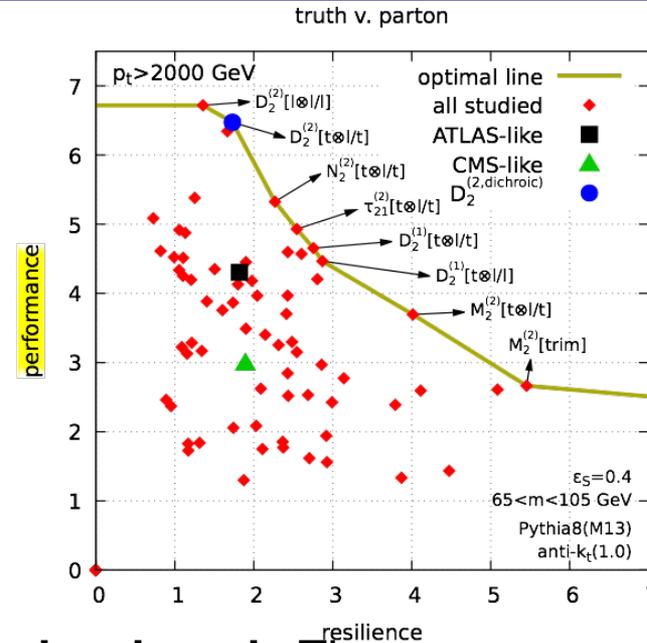
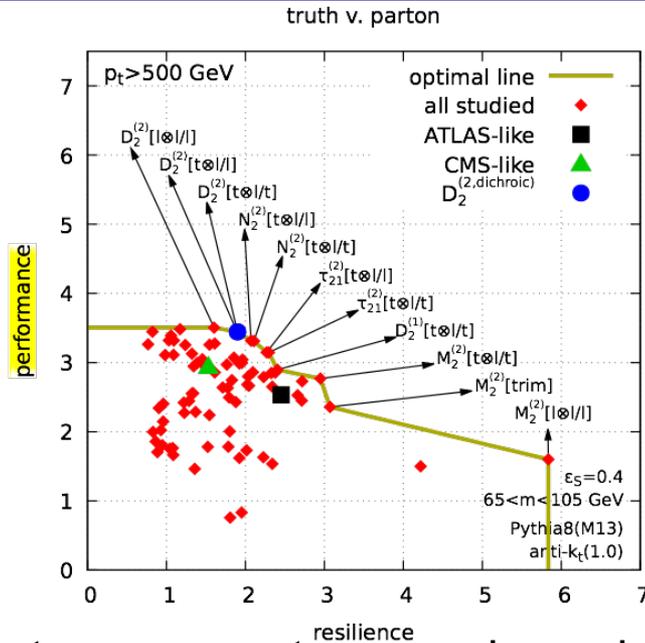
- Identification of quark initiated vs gluon initiated jets is key for many **BSM** searches : often **signals** are **dominated** by **quarks** whereas **backgrounds** by **gluons**

Radiation patterns of quark and gluon jets are not well understood, and their modelling is dependent on parton-shower generators.

- “LHC measurements are the best near-term strategy to constrain quark/gluon radiation patterns and enable quark/gluon discrimination to become a robust experimental tool.”

JHEP07(2017)091





Many two-prong taggers have been developed. These are very important tools as we move into probing lower resonance masses, and lower couplings.

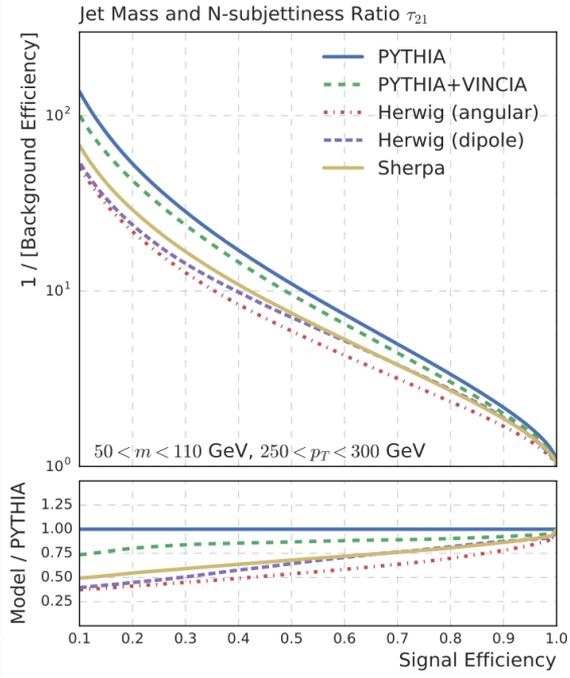
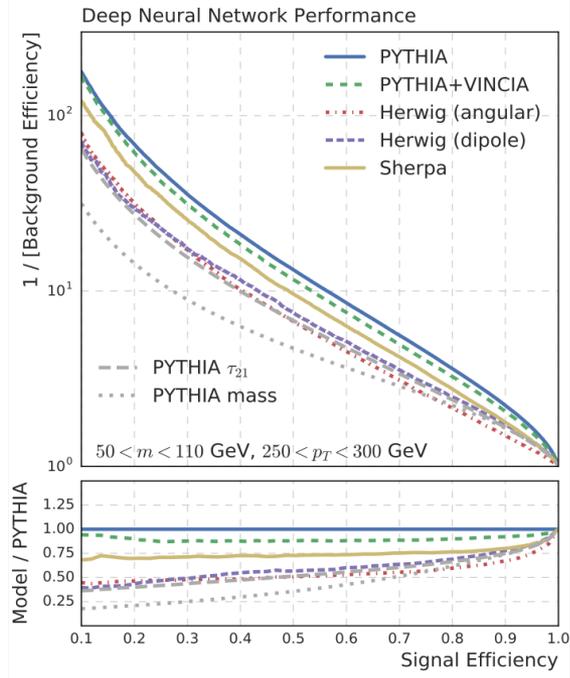
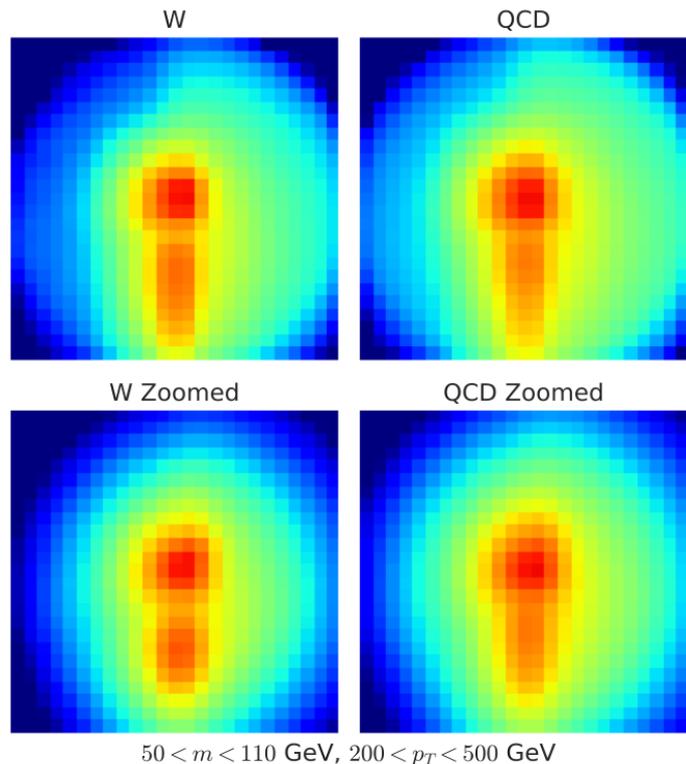
- **Their performance depends on the non-perturbative effects, and some are more robust (resilient) than others.** ★
- *“Keeping more radiation in the jet or putting tighter constraints on soft radiation at larger angles typically leads to more efficient taggers but at the same time yields more sensitivity to the regions where hadronisation and the Underlying Event have a larger impact, hence reducing resilience.”*



Jet taggers :

A machine learning example

PhysRevD.95.014018

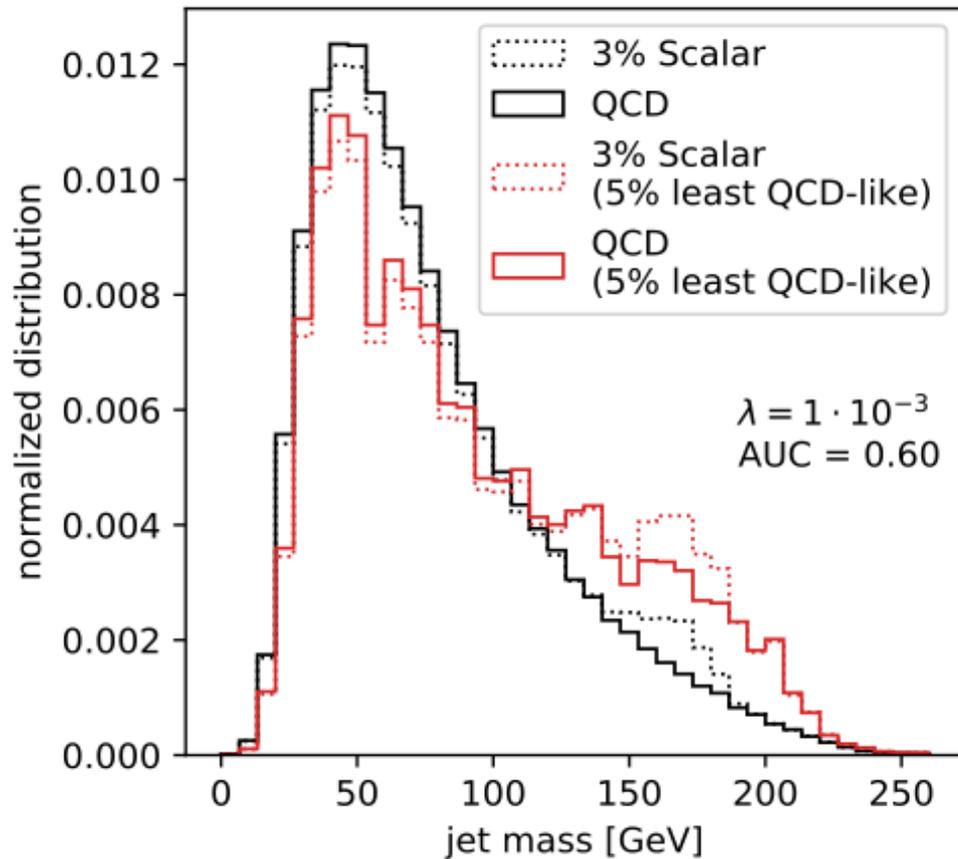


- Parton shower uncertainties, and soft radiation related ones (hadronization) are large, both for the ML cases, and for the “traditional” substructure ones (n-subjettiness) ★
- “...cautious and detailed studies of uncertainties...” are needed in order for these methodologies to be successfully applied to data analyses.



Jet taggers looking into the future : Unsupervised learning

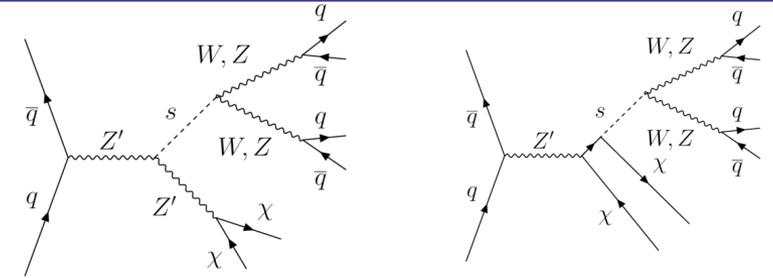
SciPost Phys. 6 (2019) 3, 030



- Another way one can **overcome large parton showering and hadronization uncertainties in ML jet taggers** is to use unsupervised learning and use real data to “teach” the tagger QCD jets.
- An example of a **dark shower model**, with **dark mesons identified** through a **peak** in the **jet mass distributions**.

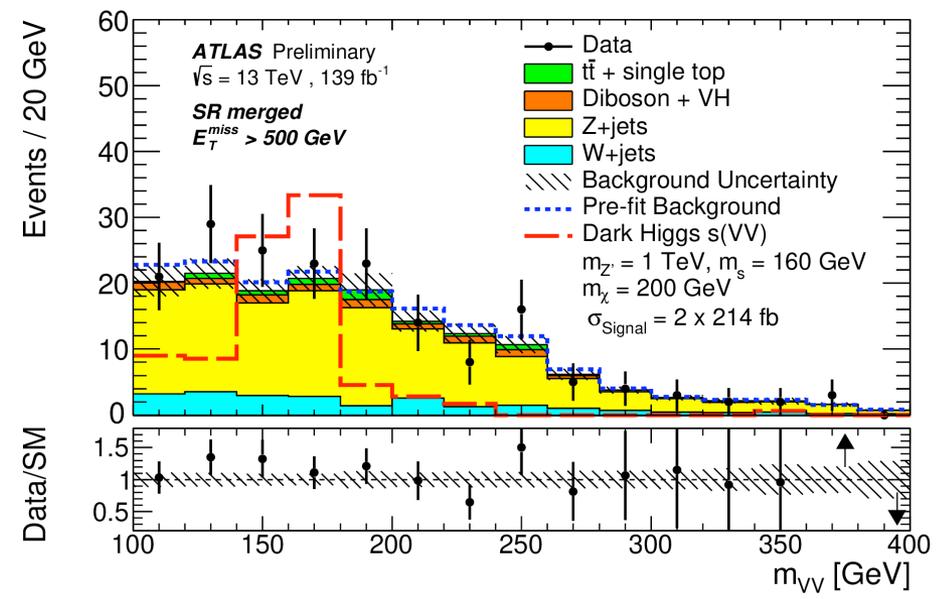
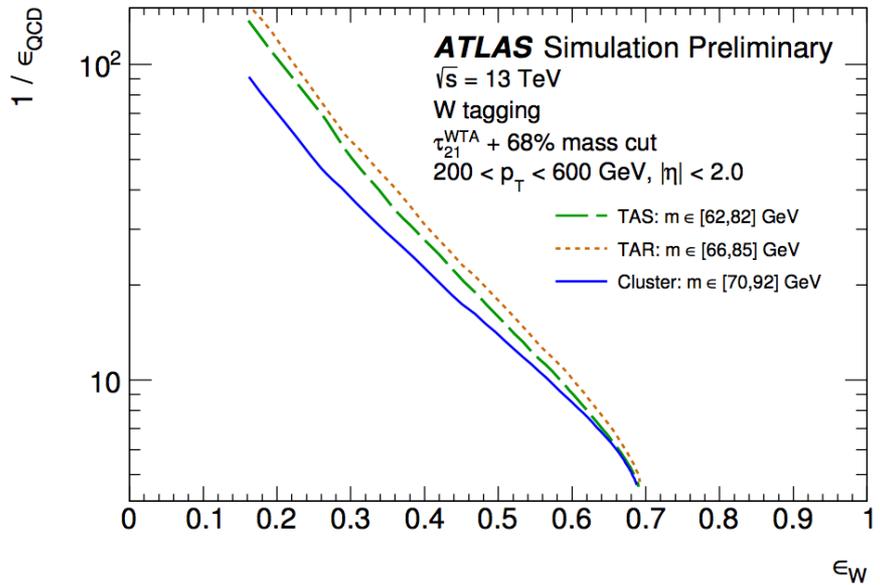


New algorithms developed that combine tracking and calorimeter information to calculate large-R jet substructure observables



ATLAS-CONF-2020-036/

<https://cds.cern.ch/record/2630864>

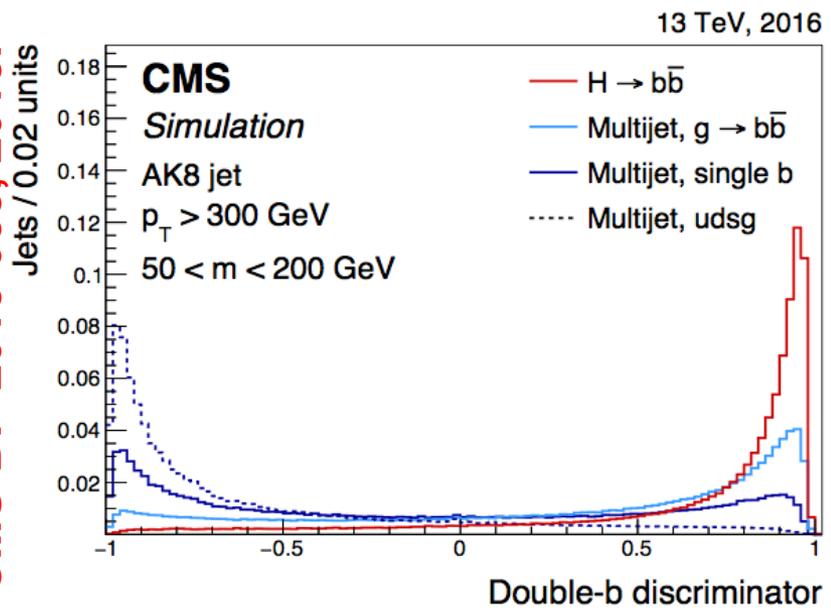


- **Hunting for Dark Higgs with novel jet taggers** in ATLAS.

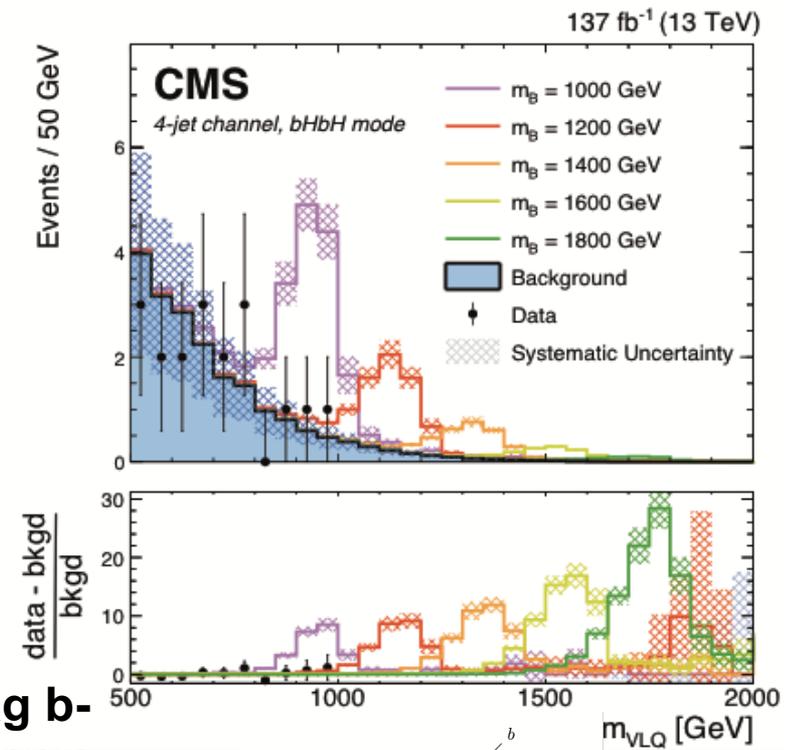
- **Theoretical systematic uncertainties due to PDF, renormalization and factorization scales, among dominant ones**



JINST 13 (2018) P05011
CMS-DP-2018-058, 2018.



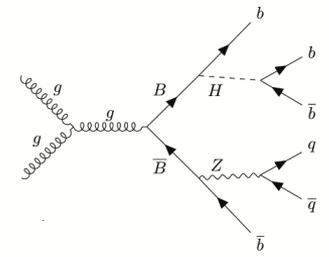
CERN-EP-2020-154 CMS-
B2G-19-005



DeepJet : jet-tagger using **DNNs** is utilized to tag **b-jets**, together with a **dedicated double b-jet tagger**

- **Hunting for vector-like quarks in CMS**

- **Main QCD background estimated in data-driven way. All background uncertainties related to the lack of knowledge of the QCD background (freely floating parameters of empirical function used), and the jet tagger efficiency and background rejection.**



- **QCD related uncertainties** concerning **PDFs, parton-showering, non-perturbative effects (hadronization and multi-parton interactions), higher orders, enter and play a key role in many powerful BSM searches.**
- **These uncertainties play a key role in developing sophisticated jet taggers, that can significantly increase the LHC discovery potential.**
- Many **efforts** from the **theory** side in **collaboration** with **experiments** aim to **improve on these, and significantly enhance the LHC discovery potential.**



Thank you!

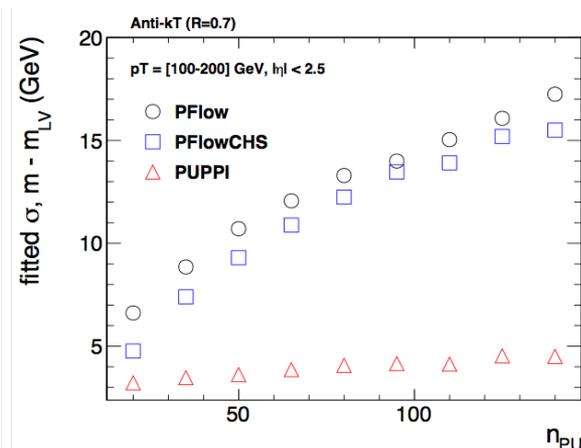
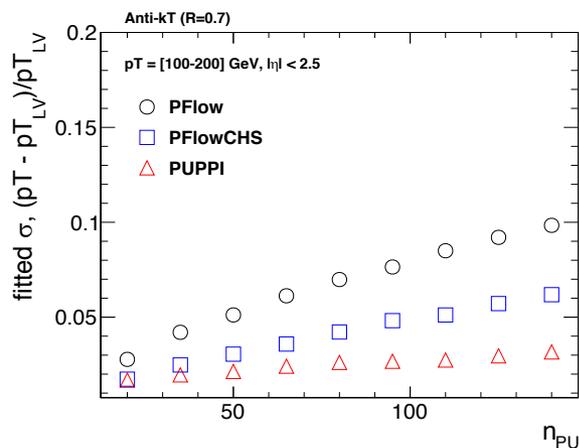
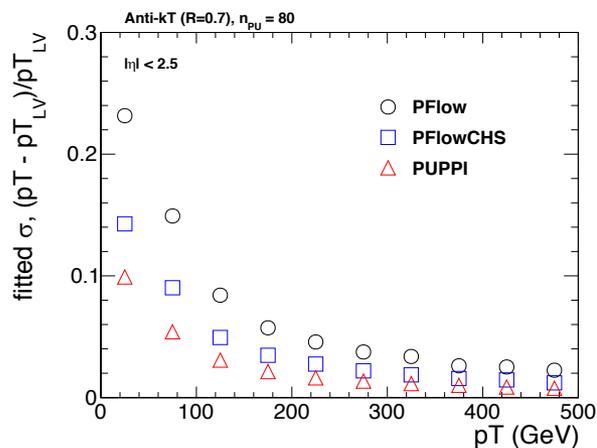
Backup slides follow



Jet Pileup Per Particle : PUPPI



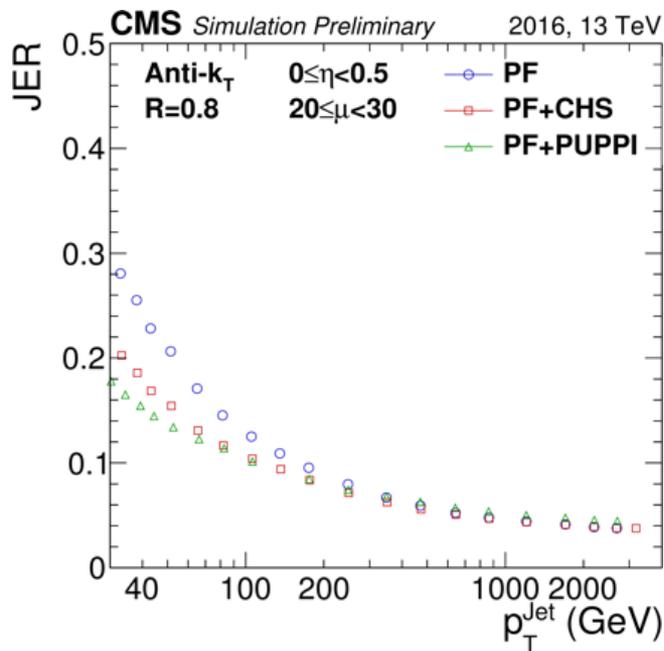
JHEP 1410 (2014) 059



$$\alpha_i = \log \sum_{j \in \text{event}} \xi_{ij} \times \Theta(R_{\min} \leq \Delta R_{ij} \leq R_0),$$

$$\text{where } \xi_{ij} = \frac{p_{Tj}}{\Delta R_{ij}}.$$

- PUPPI assigns a **weight to particles of unknown origin to be coming from pileup interactions**, using it to rescale the particle's four-momentum.
- **By applying corrections at the particle level, before jet clustering one can simultaneously perform pileup jet mitigation, and jet four-vector and shape corrections.**





PILEUP PER PARTICLE ID, PUPPI

Bertolini, Harris, Low, NT, arXiv:1407.6013

asymptotic behavior

vertexing information

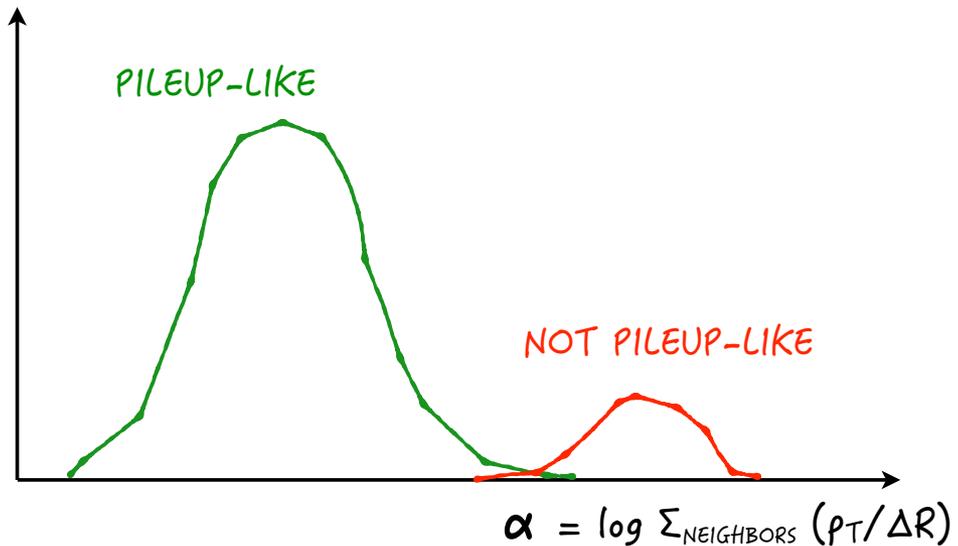
collinear QCD radiation

precision timing



A general **framework**, that assigns on a **per particle** basis a weight for **how likely** a particle is to be from pileup

key insight: using the QCD ansatz to infer neutral particles as pileup



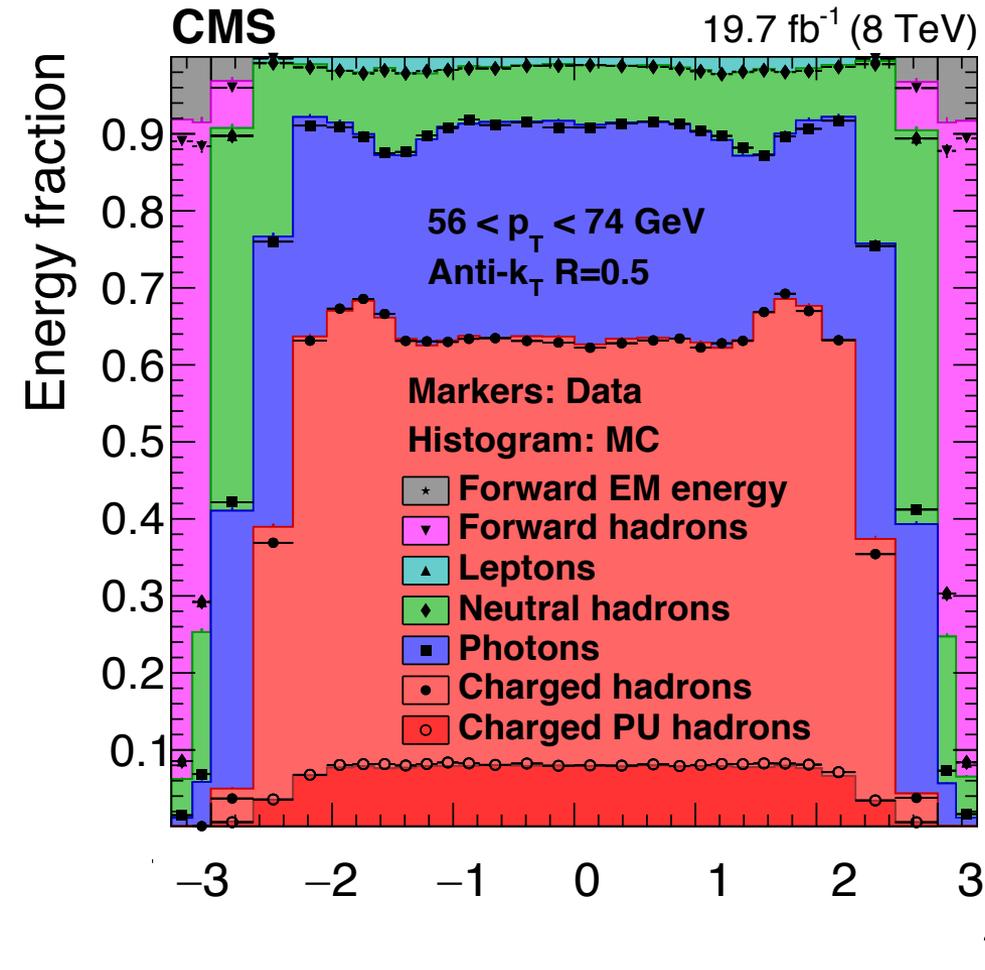
PUPPI DEMONSTRATES LARGE GAINS, EVEN FOR CURRENT 2016 DATASET



Jet Quality



CMS-PRF-14-001



Particle flow jets, described by:

- Energy fractions
- Neutral and charged particle multiplicities
- Pileup weights per particle

provide several handles on noise, pileup, and mis-reconstruction rejection.



Anti-Kt Jet Algorithm

- For each input object (Topological Clusters), d_{ij} & d_{iB} are defined as:

$$d_{ij} = \min(p_{Ti}^{-2}, p_{Tj}^{-2}) \frac{\Delta R_{ij}^2}{R^2}$$

$$d_{iB} = p_{Ti}^{-2}$$

$$\Delta R_{ij}^2 = (y_i - y_j)^2 + (\phi_i - \phi_j)^2$$

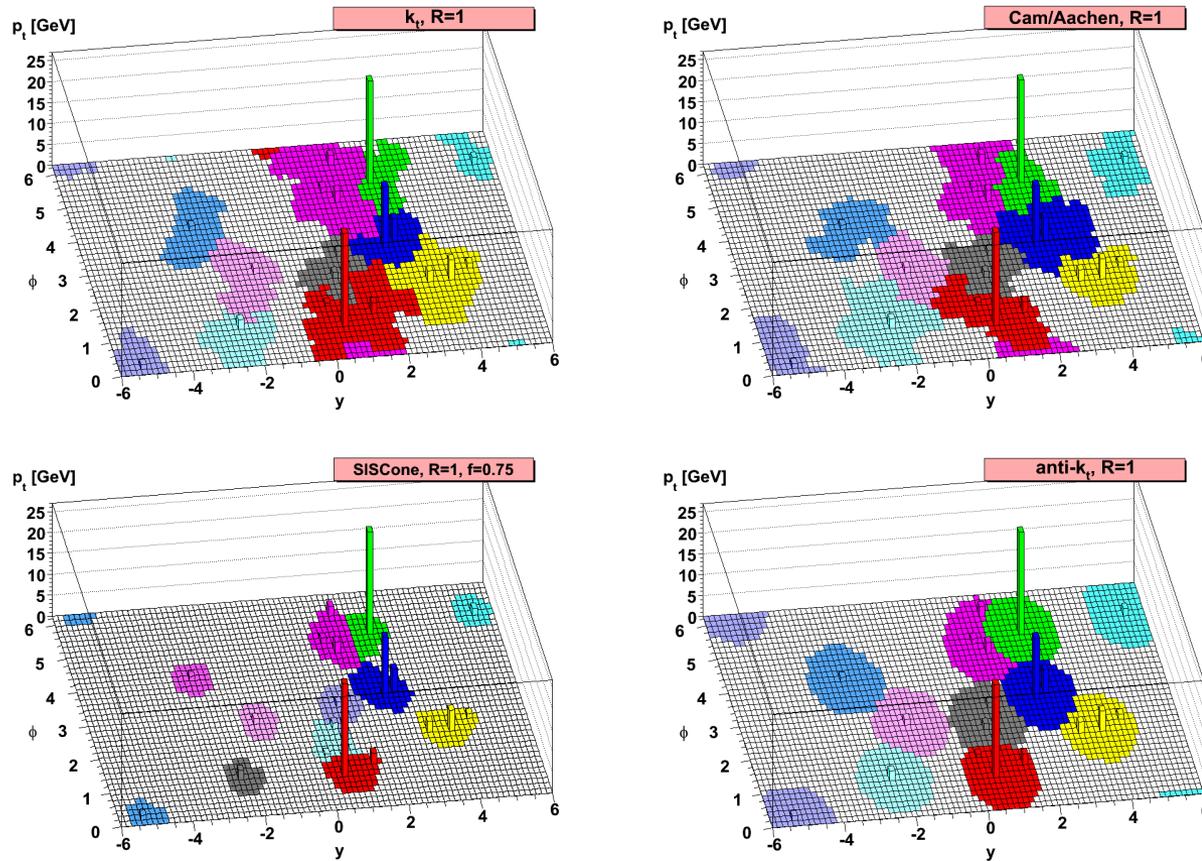
- A list of d_{ij} & d_{iB} are formed;
 - If d_{ij} is the smallest entry; objects i & j are combined & the list is remade
 - If d_{iB} is smallest, it is a jet by itself
- Anti-Kt algorithm can be implemented in NLO QCD calculations
- The algorithm also produces geometrically well-defined (cone-like) jets.



Anti-kT Algorithm



arXiv:0802.1189



$p=1$ kT
 $p=0$ CA
 $p=-$ Anti-kT

$$d_{ij} = \min(k_{ti}^{2p}, k_{tj}^{2p}) \frac{\Delta_{ij}^2}{R^2},$$

$$d_{iB} = k_{ti}^{2p},$$

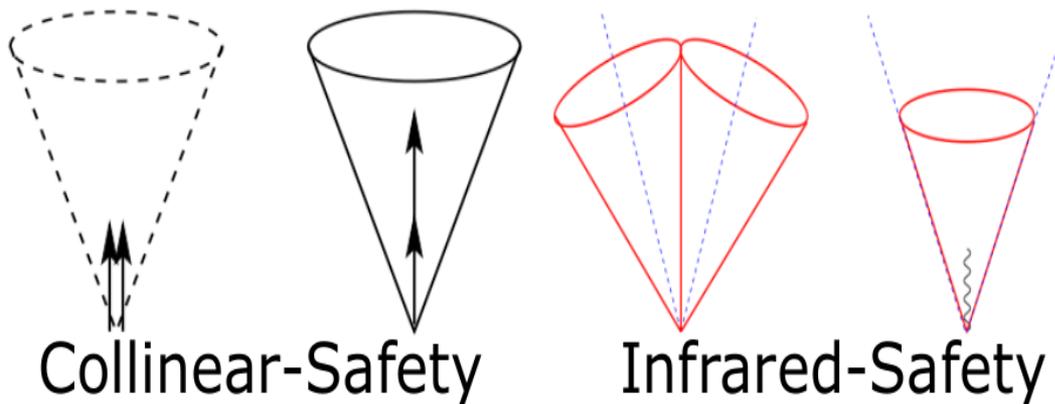
Figure 1: A sample parton-level event (generated with Herwig [8]), together with many random soft “ghosts”, clustered with four different jets algorithms, illustrating the “active” catchment areas of the resulting hard jets. For k_t and Cam/Aachen the detailed shapes are in part determined by the specific set of ghosts used, and change when the ghosts are modified.



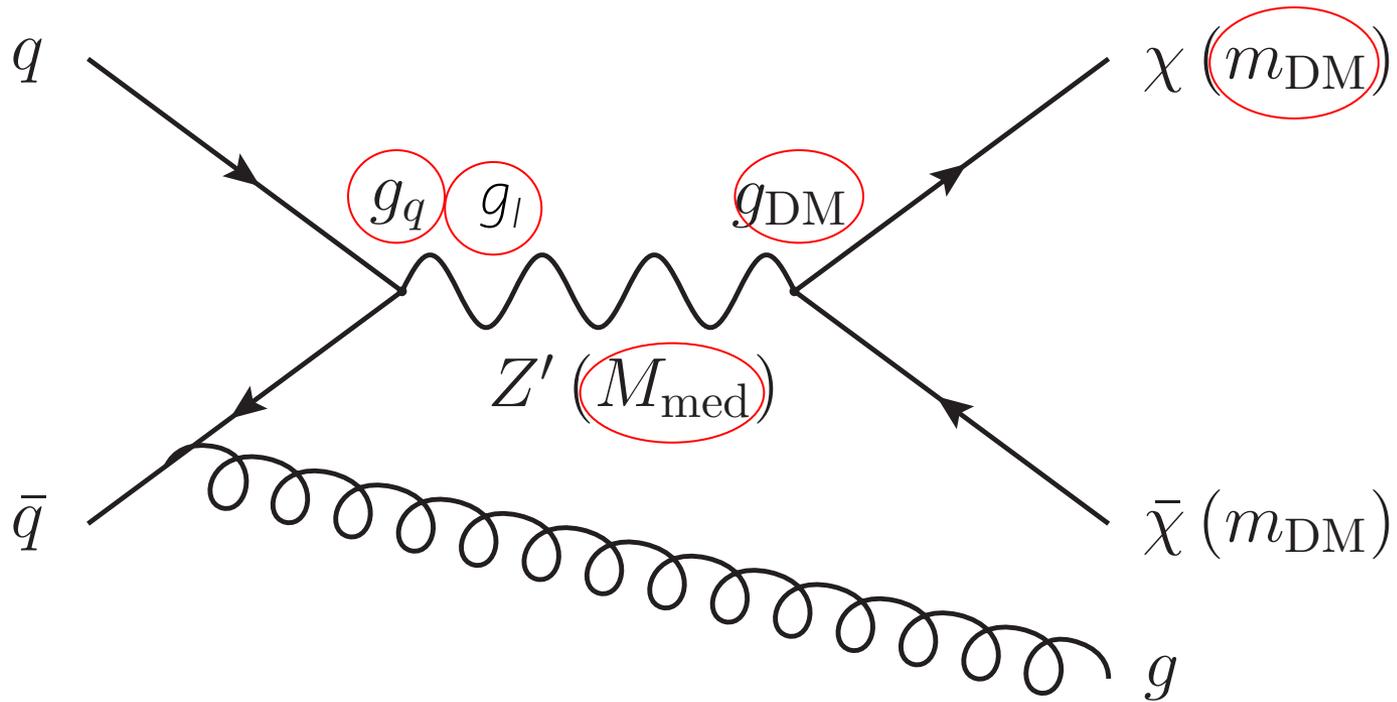
Jet clustering requirements



- **Collinear safety** : collinear splitting should not change jets
- **Infrared safety** : soft emissions should not change jets

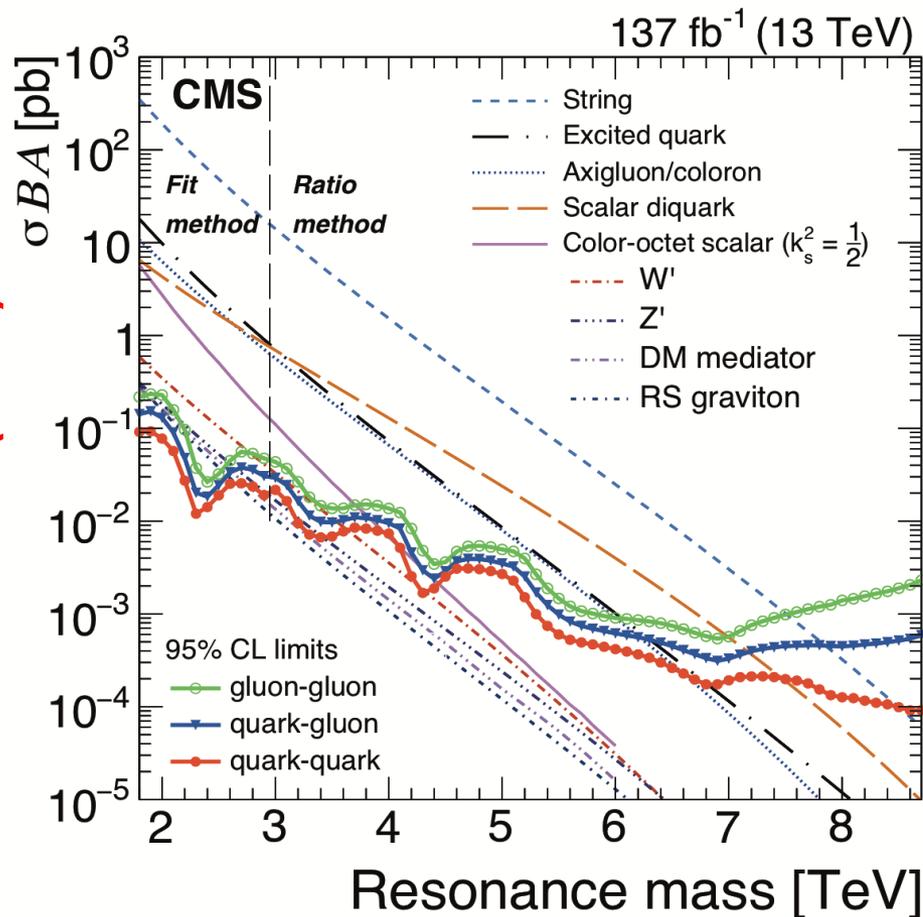


Dark Matter Searches@ the LHC: Simplified Models



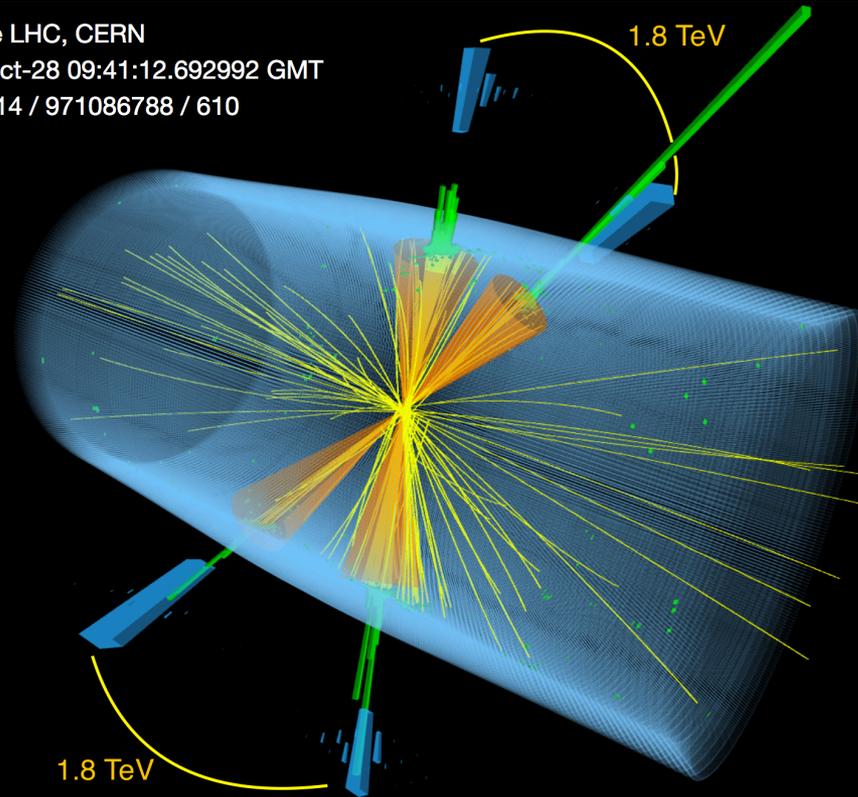
<https://arxiv.org/pdf/1703.05703.pdf>





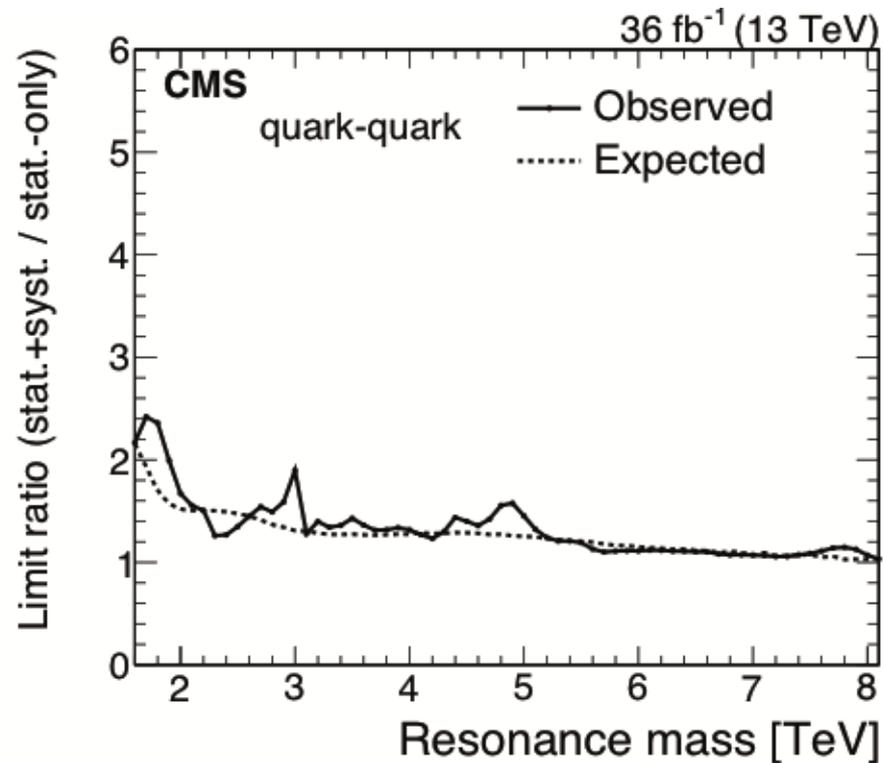
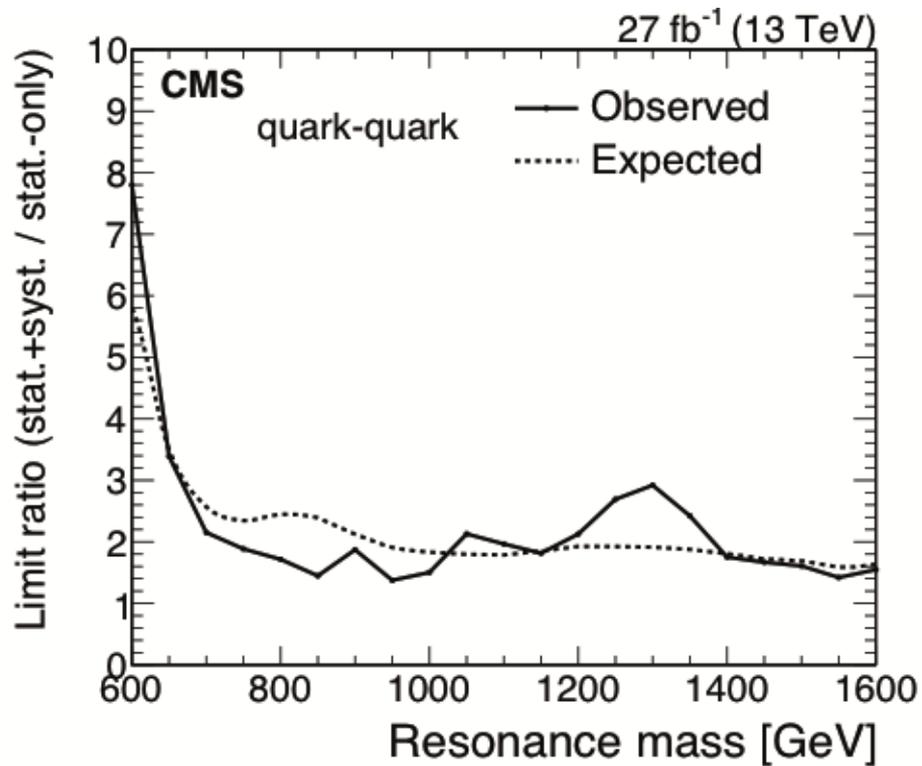


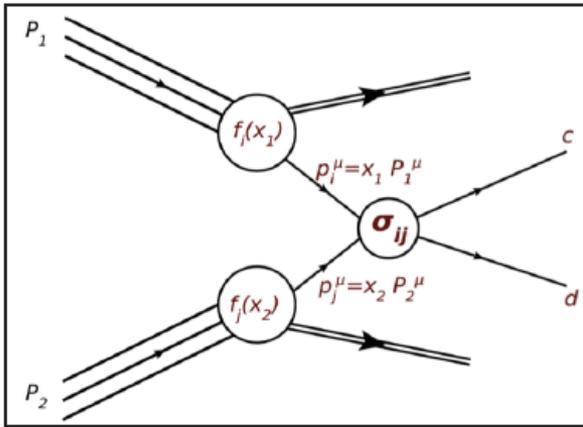
CMS Experiment at the LHC, CERN
Data recorded: 2017-Oct-28 09:41:12.692992 GMT
Run / Event / LS: 305814 / 971086788 / 610





Systematics vs Statistical Uncertainties





- Events with two energetic partons (quarks or gluons) arise in proton-proton collisions from parton-parton scattering.
- We observe this outgoing parton pair as two hadronic jets (dijets) in the detector.
- The dijet mass spectrum predicted by Quantum Chromodynamics (QCD) falls smoothly and steeply with increasing dijet mass.

$$d\sigma(P_1 P_2 \rightarrow cd) = \int_0^1 dx_1 dx_2 \sum_{q_i, q_j} f_i(x_1, \mu_F^2) f_j(x_2, \mu_F^2) d\hat{\sigma}_{q_i q_j \rightarrow cd}(\alpha_s(\mu_F^2), Q^2/\mu_F^2)$$

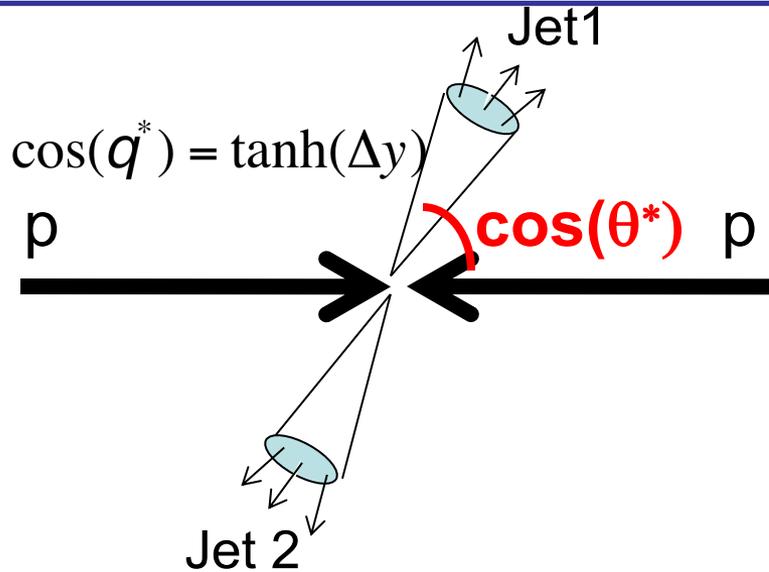
$$M_{jj}^2 = 2 p_{T1} p_{T2} [\cosh(\Delta\eta) - \cos\Delta\phi] \approx 2 p_T^2 [\cosh(\Delta\eta) + 1]$$

Higher η means increased $\Delta\eta$, which makes the Dijet mass higher

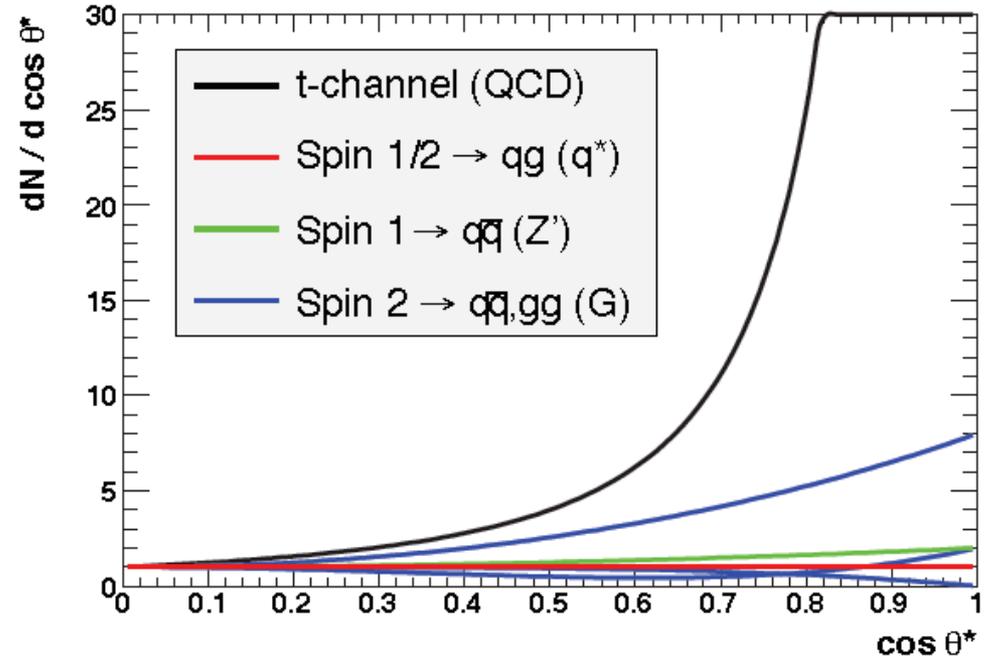
$$y = 1/2 \ln [(E + p_L)/(E - p_L)], |y| = \max(|y_1|, |y_2|)$$



Dijet Angular Distributions



Dijet Angular Distributions from Resonances and QCD



- Parton-parton scattering in QCD is t-channel dominated.
- **Stringent test of pQCD with no dependence on PDFs.**
- **New physics would show deviations from expectation at large scattering angles.**



Dilepton searches : uncertainties



JHEP04(2019)114
CMS-PAS-EXO-19-019

Source	Dielectron channel [%]		Dimuon channel [%]	
	Signal	Background	Signal	Background
Luminosity	3.2 (3.2)	3.2 (3.2)	3.2 (3.2)	3.2 (3.2)
MC statistical	<1.0 (<1.0)	<1.0 (<1.0)	<1.0 (<1.0)	<1.0 (<1.0)
Beam energy	2.0 (4.1)	2.0 (4.1)	1.9 (3.1)	1.9 (3.1)
Pile-up effects	<1.0 (<1.0)	<1.0 (<1.0)	<1.0 (<1.0)	<1.0 (<1.0)
DY PDF choice	-	<1.0 (8.4)	-	<1.0 (1.9)
DY PDF variation	-	8.7 (19)	-	7.7 (13)
DY PDF scales	-	1.0 (2.0)	-	<1.0 (1.5)
DY α_s	-	1.6 (2.7)	-	1.4 (2.2)
DY EW corrections	-	2.4 (5.5)	-	2.1 (3.9)
DY γ -induced corrections	-	3.4 (7.6)	-	3.0 (5.4)
Top quarks theoretical	-	<1.0 (<1.0)	-	<1.0 (<1.0)
Dibosons theoretical	-	<1.0 (<1.0)	-	<1.0 (<1.0)
Reconstruction efficiency	<1.0 (<1.0)	<1.0 (<1.0)	10 (17)	10 (17)
Isolation efficiency	9.1 (9.7)	9.1 (9.7)	1.8 (2.0)	1.8 (2.0)
Trigger efficiency	<1.0 (<1.0)	<1.0 (<1.0)	<1.0 (<1.0)	<1.0 (<1.0)
Identification efficiency	2.6 (2.4)	2.6 (2.4)	-	-
Lepton energy scale	<1.0 (<1.0)	4.1 (6.1)	<1.0 (<1.0)	<1.0 (<1.0)
Lepton energy resolution	<1.0 (<1.0)	<1.0 (<1.0)	2.7 (2.7)	<1.0 (6.7)
Multi-jet & W +jets	-	10 (129)	-	-
Total	10 (11)	18 (132)	11 (18)	14 (24)

- **Resonant CMS search.**
- **Among largest uncertainties are the PDF related ones.**



Dilepton searches : uncertainties



Phys. Lett. B 796 (2019) 68
ATLAS-EXOT-2019-16

Uncertainty source for m_X [GeV]	Dielectron			Dimuon		
	300	2000	5000	300	2000	5000
Spurious signal	± 12.5 (12.0)	± 4.6 (10.8)	± 0.1 (1.0)	± 11.7 (11.0)	± 3.8 (3.5)	± 2.1 (2.2)
Lepton identification	± 1.6 (1.6)	± 5.6 (5.6)	± 5.6 (5.6)	± 1.8 (1.8)	$^{+12}_{-10}$ $\left(\begin{smallmatrix} +12 \\ -10 \end{smallmatrix} \right)$	$^{+25}_{-20}$ $\left(\begin{smallmatrix} +25 \\ -20 \end{smallmatrix} \right)$
Isolation	± 0.3 (0.3)	± 1.1 (1.2)	± 1.1 (1.1)	± 0.4 (0.4)	± 0.4 (0.4)	± 0.4 (0.5)
Luminosity	± 1.7 (1.7)	± 1.7 (1.7)	± 1.7 (1.7)	± 1.7 (1.7)	± 1.7 (1.7)	± 1.7 (1.7)
Electron energy scale	$^{-1.7}_{-4.0}$ $\left(\begin{smallmatrix} +1.0 \\ -1.8 \end{smallmatrix} \right)$	$^{-1.9}_{-6.0}$ $\left(\begin{smallmatrix} +1.7 \\ -2.9 \end{smallmatrix} \right)$	$^{+0.1}_{-0.4}$ (± 0.8)	-	-	-
Electron energy resolution	$^{+7.9}_{-8.3}$ $\left(\begin{smallmatrix} +1.1 \\ -0.9 \end{smallmatrix} \right)$	$^{+9.0}_{-11.8}$ $\left(\begin{smallmatrix} +0.7 \\ -0.5 \end{smallmatrix} \right)$	$^{+0.4}_{-0.9}$ (± 0.1)	-	-	-
Muon ID resolution	-	-	-	$^{+0.8}_{-2.3}$ $\left(\begin{smallmatrix} +0.3 \\ -0.8 \end{smallmatrix} \right)$	$^{+0.9}_{-1.3}$ $\left(\begin{smallmatrix} +0.7 \\ -1.1 \end{smallmatrix} \right)$	$^{+0.6}_{-0.4}$ $\left(\begin{smallmatrix} +0.5 \\ -0.3 \end{smallmatrix} \right)$
Muon MS resolution	-	-	-	$^{+2.8}_{-3.8}$ $\left(\begin{smallmatrix} +1.0 \\ -1.3 \end{smallmatrix} \right)$	$^{+3.2}_{-3.0}$ $\left(\begin{smallmatrix} +2.6 \\ -2.4 \end{smallmatrix} \right)$	± 2.4 (2.1)
‘Good muon’ requirement	-	-	-	± 0.6 (0.6)	$^{+9.0}_{-8.2}$ $\left(\begin{smallmatrix} +9.0 \\ -8.2 \end{smallmatrix} \right)$	$^{+55}_{-35}$ $\left(\begin{smallmatrix} +55 \\ -35 \end{smallmatrix} \right)$

- **Resonant ATLAS search.**
- **Among largest uncertainties are the background modeling related ones depicted in the “spurious” signal category.**



Jet Substructure : N-subjetiness

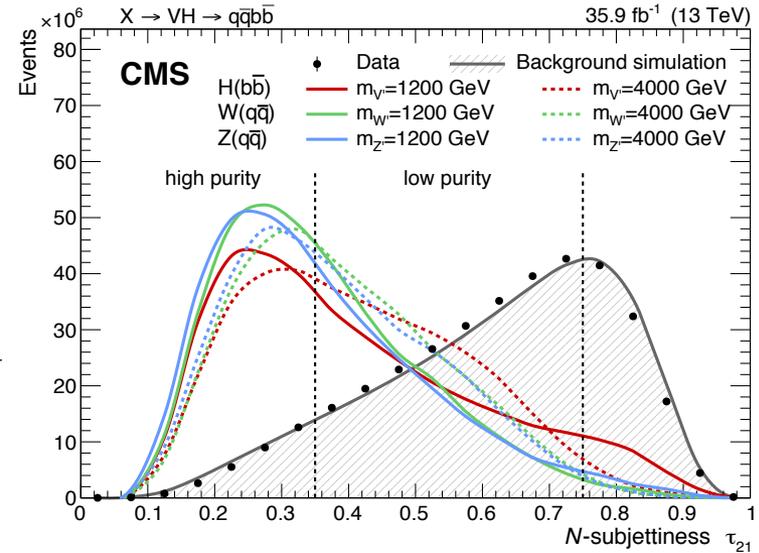


JHEP 1103:015,2011

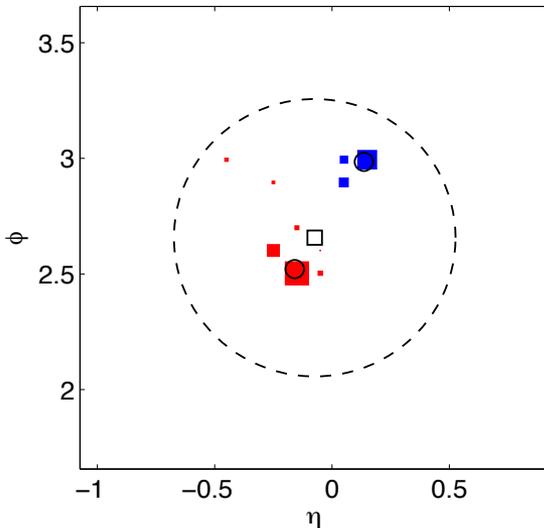
- Main discriminating variables between a "uniform" jet and one experiencing inner "sub-structure" are the N-subjetiness ones

$$\tau_N = \frac{1}{d_0} \sum_k p_{T,k} \min \{ \Delta R_{1,k}, \Delta R_{2,k}, \dots, \Delta R_{N,k} \}$$

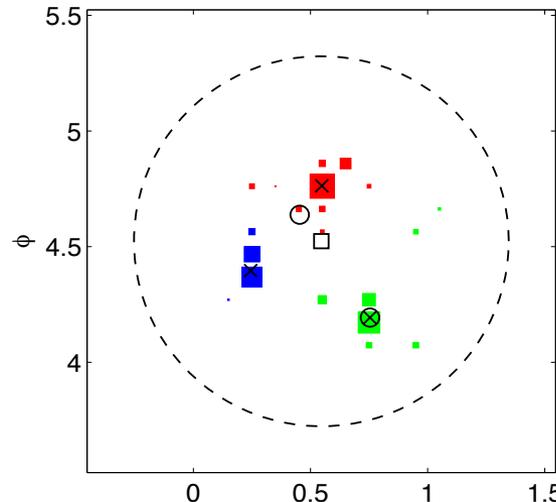
$$d_0 = \sum_k p_{T,k} R_0,$$



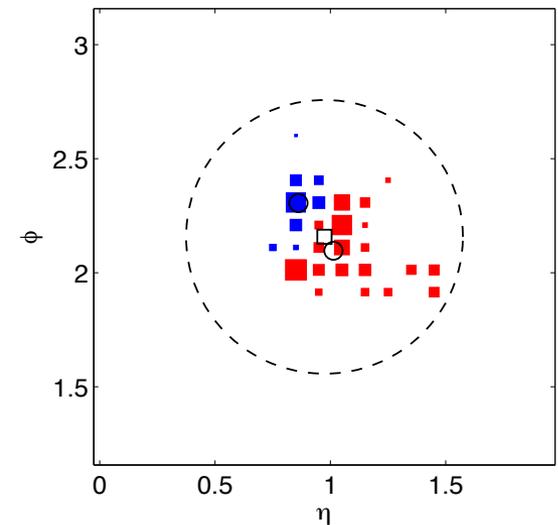
W jet, $\tau_2/\tau_1=0.15$



Top jet, $\tau_3/\tau_2=0.21$



QCD jet, $\tau_2/\tau_1=0.73$

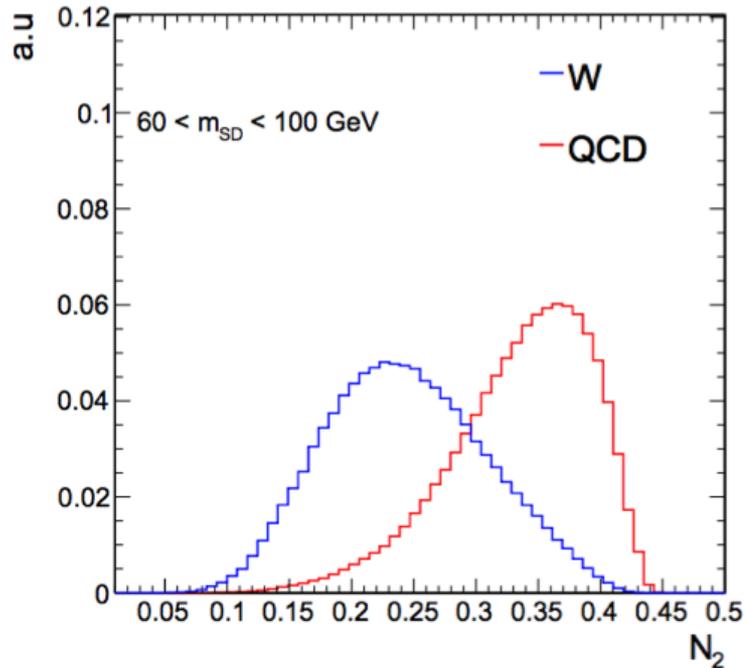




Jet Substructure : Energy Flow



JHEP 1306 (2013) 108



$$\text{ECF}(N, \beta) = \sum_{i_1 < i_2 < \dots < i_N \in J} \left(\prod_{a=1}^N E_{i_a} \right) \left(\prod_{b=1}^{N-1} \prod_{c=b+1}^N \theta_{i_b i_c} \right)^\beta \cdot$$

Infrared, collinear safe, for all $\beta > 0$.

- Main discriminating variables between a "uniform" jet and one experiencing inner "sub-structure" using **generalized energy correlation functions**.
- They use **information about the energies and pair-wise angles of particles within a jet**.
- They can identify **N-prong jet substructure without requiring a subjet finding procedure**, and hence can **optimize discriminating power depending on the problem**.



Jet Mass



Two main methods for jet grooming: **Pruning** [Phys.Rev.D80,05150 (2009)] (recluster) and **Softdrop** [JHEP05(2014),146] (decluster).

Eur.Phys.J. C77 (2017) no.9, 636

