

# Measurements of the strong coupling constant at the LHC and perspectives

Fernando Barreiro  
On behalf of the ATLAS and CMS Collaborations

Universidad Autónoma de Madrid

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- The strong coupling, one of the main free parameters in QCD, is known with least precision than any other SM coupling,  
 $\alpha_s(m_Z) = 0.1181 \pm 0.0011$
- Uncertainties in  $\alpha_s$  are important in theoretical calculations of
  - Higgs boson production cross sections and decay widths,
  - top quark mass
  - EW precision observables
  - physics at Planck/GUT scales
- At hadron colliders the only observable known at NNLO accuracy used so far for determination of the strong coupling is  $\sigma(t\bar{t})$ ,

New measurements considered this year:

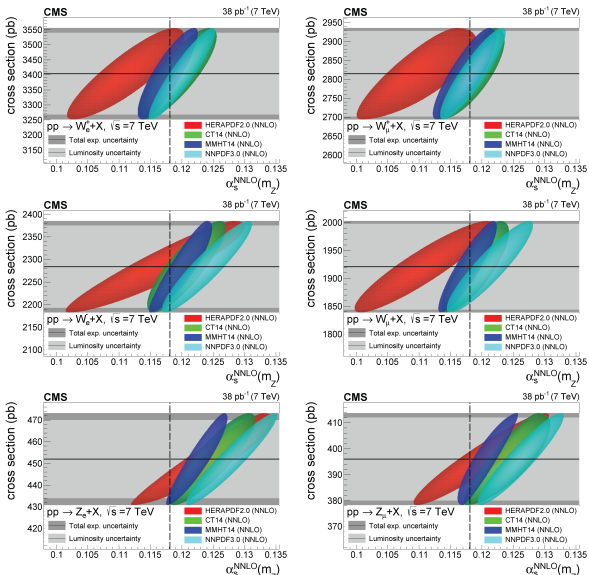
- CMS: Determination of  $\alpha_s(M_Z)$  from inclusive  $W^\pm$  and  $Z$  boson production in  $p - p$  at 7 and 8 TeV ([JHEP 06 \(2020\) 018](#) )
- ATLAS : Determination of  $\alpha_s(M_Z)$  and test of asymptotic freedom from Transverse Energy-Energy Correlations (TEEC) at 13 TeV ([ATLAS-CONF-2020-025](#) )

The table below shows a summary of the twelve  $W^\pm$  and  $Z$  boson fiducial cross sections in the  $e, \mu$  decay channels at 7 TeV ( $38.0 \text{ pb}^{-1}$ ) and 8 TeV ( $18.2 \text{ pb}^{-1}$ ) together with their uncertainties : (JHEP 06 (2020) 018 )

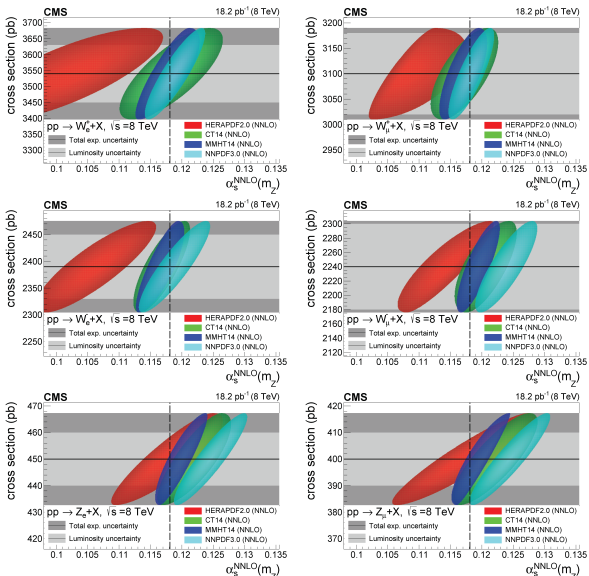
Measurement	Fiducial cross section
pp at $\sqrt{s} = 7 \text{ TeV}$ [? ]	
$W_e^+, p_T^e > 25 \text{ GeV},  \eta^e  < 2.5$	$3404 \pm 12 \text{ (stat)} \pm 67 \text{ (syst)} \pm 136 \text{ (lumi)} \text{ pb} = 3404 \pm 152 \text{ pb}$
$W_e^-, p_T^e > 25 \text{ GeV},  \eta^e  < 2.5$	$2284 \pm 10 \text{ (stat)} \pm 43 \text{ (syst)} \pm 91 \text{ (lumi)} \text{ pb} = 2284 \pm 101 \text{ pb}$
$Z_e, p_T^e > 25 \text{ GeV},  \eta^e  < 2.5, 60 < m_Z < 120 \text{ GeV}$	$452 \pm 5 \text{ (stat)} \pm 10 \text{ (syst)} \pm 18 \text{ (lumi)} \text{ pb} = 452 \pm 21 \text{ pb}$
$W_\mu^+, p_T^\mu > 25 \text{ GeV},  \eta^\mu  < 2.1$	$2815 \pm 9 \text{ (stat)} \pm 42 \text{ (syst)} \pm 113 \text{ (lumi)} \text{ pb} = 2815 \pm 121 \text{ pb}$
$W_\mu^-, p_T^\mu > 25 \text{ GeV},  \eta^\mu  < 2.1$	$1921 \pm 8 \text{ (stat)} \pm 27 \text{ (syst)} \pm 77 \text{ (lumi)} \text{ pb} = 1921 \pm 82 \text{ pb}$
$Z_\mu, p_T^\mu > 20 \text{ GeV},  \eta^\mu  < 2.1, 60 < m_Z < 120 \text{ GeV}$	$396 \pm 3 \text{ (stat)} \pm 7 \text{ (syst)} \pm 16 \text{ (lumi)} \text{ pb} = 396 \pm 18 \text{ pb}$
pp at $\sqrt{s} = 8 \text{ TeV}$ [? ]	
$W_e^+, p_T^e > 25 \text{ GeV},  \eta^e  < 1.44, 1.57 <  \eta^e  < 2.5$	$3540 \pm 20 \text{ (stat)} \pm 110 \text{ (syst)} \pm 90 \text{ (lumi)} \text{ pb} = 3540 \pm 140 \text{ pb}$
$W_e^-, p_T^e > 25 \text{ GeV},  \eta^e  < 1.44, 1.57 <  \eta^e  < 2.5$	$2390 \pm 10 \text{ (stat)} \pm 60 \text{ (syst)} \pm 60 \text{ (lumi)} \text{ pb} = 2390 \pm 90 \text{ pb}$
$Z_e, p_T^e > 25 \text{ GeV},  \eta^e  < 1.44, 1.57 <  \eta^e  < 2.5, 60 < m_Z < 120 \text{ GeV}$	$450 \pm 10 \text{ (stat)} \pm 10 \text{ (syst)} \pm 10 \text{ (lumi)} \text{ pb} = 450 \pm 20 \text{ pb}$
$W_\mu^+, p_T^\mu > 25 \text{ GeV},  \eta^\mu  < 2.1$	$3100 \pm 10 \text{ (stat)} \pm 40 \text{ (syst)} \pm 80 \text{ (lumi)} \text{ pb} = 3100 \pm 90 \text{ pb}$
$W_\mu^-, p_T^\mu > 25 \text{ GeV},  \eta^\mu  < 2.1$	$2240 \pm 10 \text{ (stat)} \pm 20 \text{ (syst)} \pm 60 \text{ (lumi)} \text{ pb} = 2240 \pm 60 \text{ pb}$
$Z_\mu, p_T^\mu > 25 \text{ GeV},  \eta^\mu  < 2.1, 60 < m_Z < 120 \text{ GeV}$	$400 \pm 10 \text{ (stat)} \pm 10 \text{ (syst)} \pm 10 \text{ (lumi)} \text{ pb} = 400 \pm 20 \text{ pb}$

- Calculations at NNLO accuracy with MCFM (R. Boughezal et al., EPJ C 77 (2017) 7 ) interfaced with LHAPDF (CT14, MMHT14, NNPDF3.0 and HERAPDF2.0) and  $\mu_R = \mu_F = m_W, m_Z$
- EW corrections are taken into account with the code MCSANC (S.G. Bondarenko and A.A. Sapronov, Comp. Phys. Comm. 184 (2013) 2343 )

Experimental fiducial x-sections on the Y-axis are shown at 7 TeV together with joint probability densities for 4 PDF sets as a function of  $\alpha_s(M_Z)$  on the X-axis



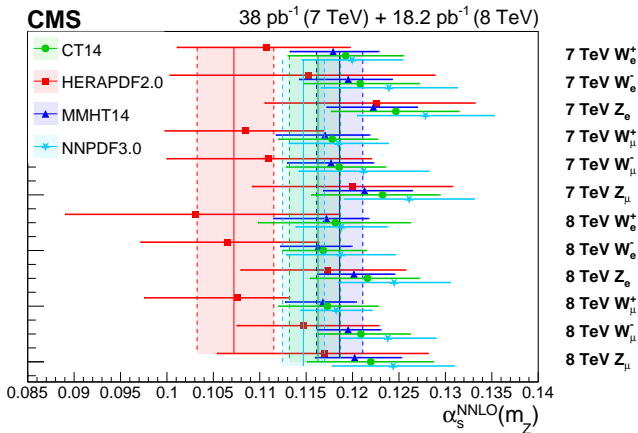
Experimental fiducial x-sections on the Y-axis are shown at 8 TeV together with joint probability densities for 4 PDF sets as a function of  $\alpha_s(M_Z)$  on the X-axis



# $W^\pm, Z$ production at 7 and 8 TeV with CMS at the LHC

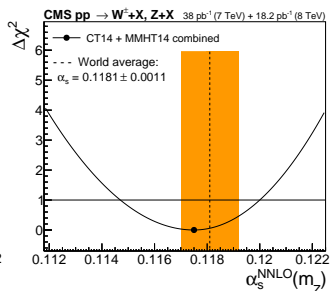
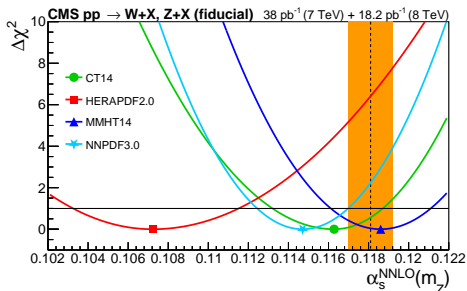
- From previous two figures: predictions using the HERAPDF2.0 set, which is obtained from fits to DIS data only, tend to be above the data
- Those using NNPDF3.0 tend to be systematically below
- While MMHT14 and CT14 overlap

These features translate into these values of  $\alpha_s(m_Z)$  extracted for each channel and for each PDF separately



## Results for $\alpha_s(m_Z)$ extracted from a global fit

PDF	$\alpha_s(m_Z)$	$\delta(\text{stat})$	$\delta(\text{lumi})$	$\delta(\text{syst})$	$\delta(\text{PDF})$	$\delta(\text{scale})$	$\delta(\text{num})$	$\chi^2/\text{ndf}$
CT14	$0.1163^{+0.0024}_{-0.0031}$	0.0007	0.0013	0.0010	$+0.0016$ $-0.0022$	0.0009	0.0006	13.3/11
HERAPDF2.0	$0.1072^{+0.0043}_{-0.0040}$	0.0012	0.0027	0.0012	$+0.0027$ $-0.0020$	0.0012	0.0009	14.2/11
MMHT14	$0.1186 \pm 0.0025$	0.0003	0.0018	0.0009	0.0013	0.0007	0.0002	10.2/11
NNPDF3.0	$0.1147 \pm 0.0023$	0.0009	0.0008	0.0007	0.0014	0.0006	0.0010	29.2/11



- Event shapes have been measured in  $e^+e^-$ ,  $e-p$  and hadron colliders
- One particularly interesting example of event shapes : TEEC and ATEEC
- Measured by ATLAS at 7 TeV ( $158 \text{ pb}^{-1}$ ) and 8 TeV ( $20.2 \text{ fb}^{-1}$ )

Extension at 13 TeV desirable because

- The interaction scale is extended from below 1 TeV to 4 TeV
- Improved statistical accuracy from 1.5% to **0.5%**
- Reduction of systematic uncertainties from 5% to **2%**
- Reduction in JES and JER uncertainties due to combined use of tracking and calorimetric information ([ATLAS Coll.](#), [CERN-EP-2020-083](#) )
- Aim: leave as a legacy a measurement which supersedes those from LEP

For a sample of  $N$  multijet events, the TEEC function is defined as :

$$\frac{1}{\sigma'} \frac{d\Sigma'}{d\cos\phi} = \frac{1}{N} \sum_{A=1}^N \frac{1}{\Delta\cos\phi} \sum_{\substack{a,b=1 \\ \text{pairs in } \Delta\cos\phi}}^{N_{\text{jets}}} \frac{2E_{T_a}^A E_{T_b}^A}{(E_T^A)^2} \quad (1)$$

The ATEEC is defined as its forward-backward asymmetry

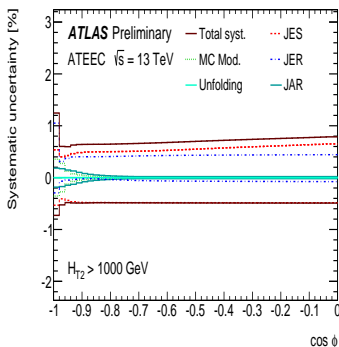
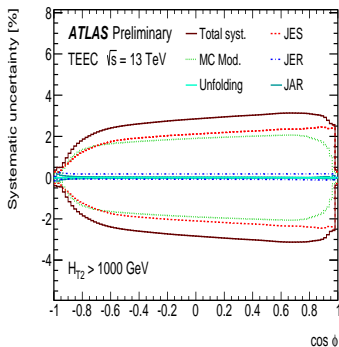


- Full Run 2 data set at  $\sqrt{s} = 13 \text{ TeV}$  with an integrated luminosity of  $139 \text{ fb}^{-1}$
- *HLT\_j460* unprescaled single jet trigger
- Jet clustering: particle flow objects
- Jet reconstruction: anti- $k_T$  algorithm with  $R = 0.4$
- Event and jet cleaning algorithms applied
- Jet selection:  $p_T > 60 \text{ GeV}$  and  $|\eta| < 2.4$  for each jet
- Event selection :  $H_{T2} = p_{T1} + p_{T2} > 1 \text{ TeV}$
- Ten bins in  $H_{T2}$  considered: From  $1.0 < H_{T2} < 1.2 \text{ TeV}$  up to  $H_{T2} > 3.5 \text{ TeV}$

# ATLAS TEEC at 13 TeV: Unfolding and experimental systematic uncertainties

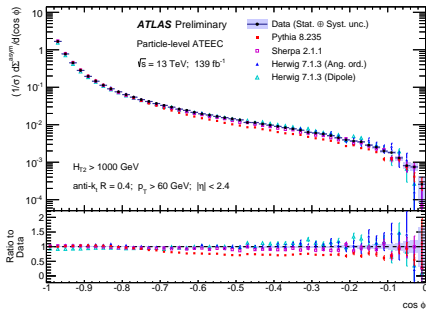
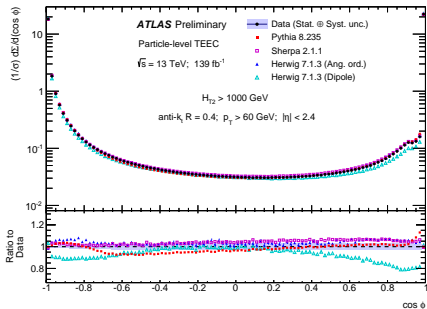
Data unfolded with an iterative Bayesian procedure

The plots below show the systematic uncertainties for the inclusive TEEC (left) and ATEEC (right)



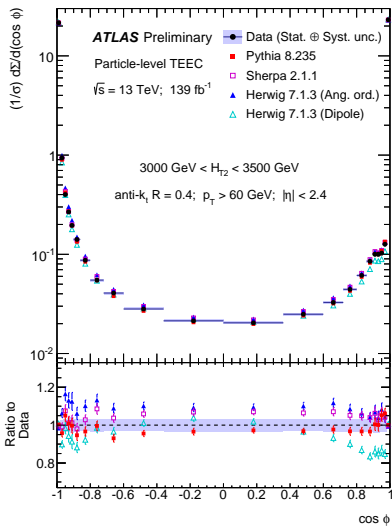
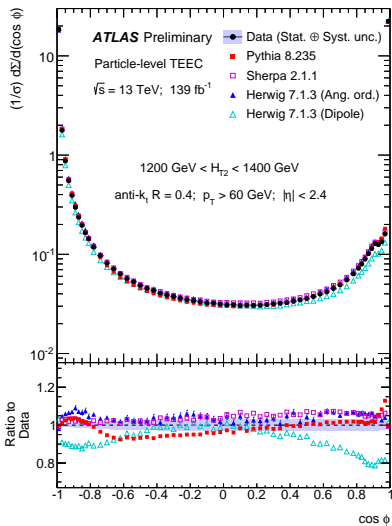
- Dominant systematic sources are the JES and MC modelling
- They amount to **1 – 2%** in the central plateau for any  $H_{T2}$  bin

# ATLAS Inclusive TEEC (lhs) and ATEEC (rhs) at particle level and comparison with Monte Carlo models

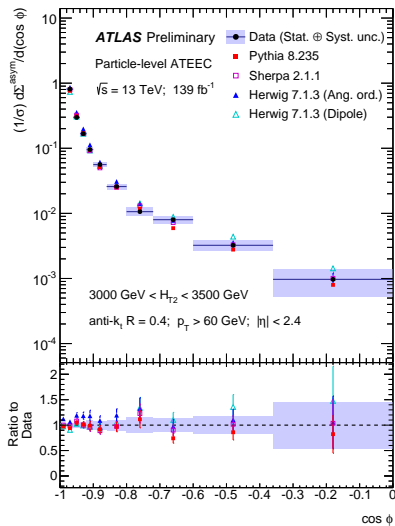
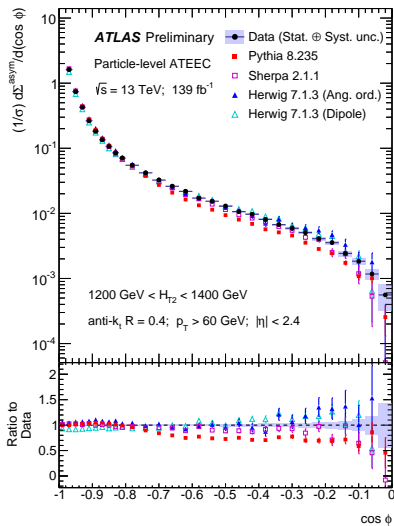


- LO SHERPA ( $2 \rightarrow 2$  and  $2 \rightarrow 3$ ) and NLO HERWIG 7 (angular ordered PS) give similar overall good description of the data
- LO PYTHIA 8 ( $2 \rightarrow 2$ ) underestimates the asymmetry at large angles
- NLO HERWIG 7 (dipole PS) needs a better tune

# ATLAS TEEC at particle level for two representative $H_{T2}$ bins

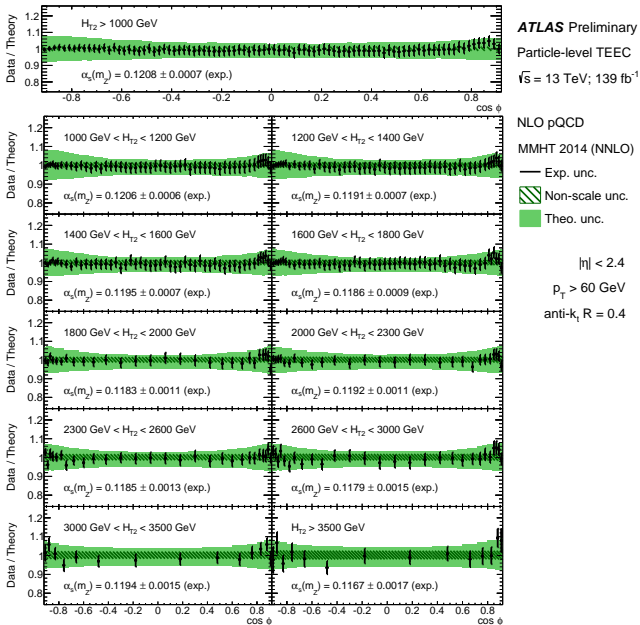


# ATLAS ATEEC at particle level for two representative $H_{T2}$ bins

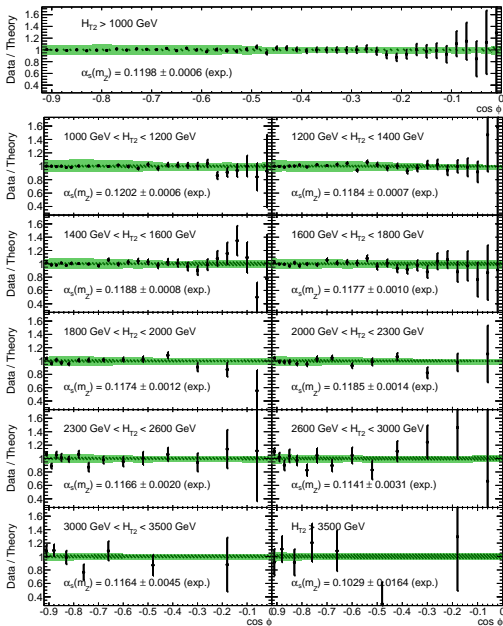


- Analogue of the EEC in  $e^+e^-$  annihilation is the transverse EEC: (A. Ali, E. Pietarinen and W.J. Stirling, PL B141 (1984) 447 )
- NLO corrections are found moderate : (A. Ali, F.B., J. Llorente and W. Wang., Phys. Rev. D86 (2012) 114017 )
- Theoretical predictions computed at parton level using NLOJET++ (Z. Nagy, Phys. Rev. D 68 (2003) 094002 )
- Collinear emissions are controlled with the cut  $|\cos\phi| < 0.92$
- Renormalization scale choice  $\mu_R = \hat{H}_T = \sum_i p_{Ti}$
- Cross sections are convoluted with NNLO PDF sets available from LHAPDF namely MMHT 2014, NNPDF 3.0 and CT14
- Factorization scale  $\mu_F = \mu_R/2$
- Jets are reconstructed with FASTJET
- Non perturbative corrections are computed with PYTHIA 8.240 with Tune A14, compatible with unity within 1%
- Scale uncertainty is the dominant theoretical uncertainty source  $O(6\%)$
- PDF uncertainties obtained from variations of the PDF group eigenvectors are subdominant  $O(1\%)$

# $\alpha_S(M_Z)$ determined from fits to ATLAS TEEC particle level data



# $\alpha_s(M_Z)$ determined from fits to ATLAS ATEEC particle level data



ATLAS Preliminary

Particle-level ATEEC

$\sqrt{s} = 13$  TeV;  $139 \text{ fb}^{-1}$

NLO pQCD

MMHT 2014 (NNLO)

— Exp. unc.

▨ Non-scale unc.

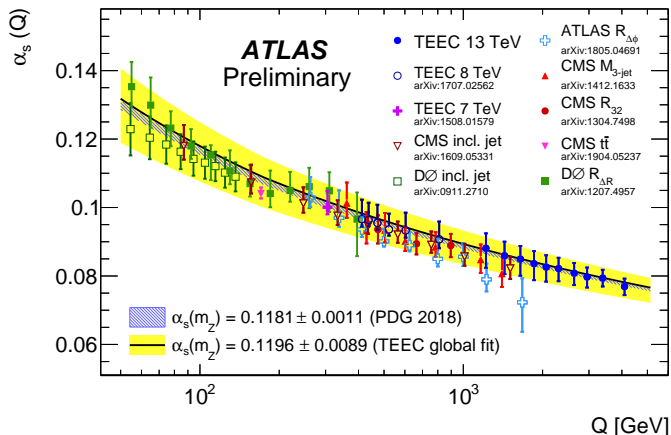
■ Theo. unc.

$|\eta| < 2.4$

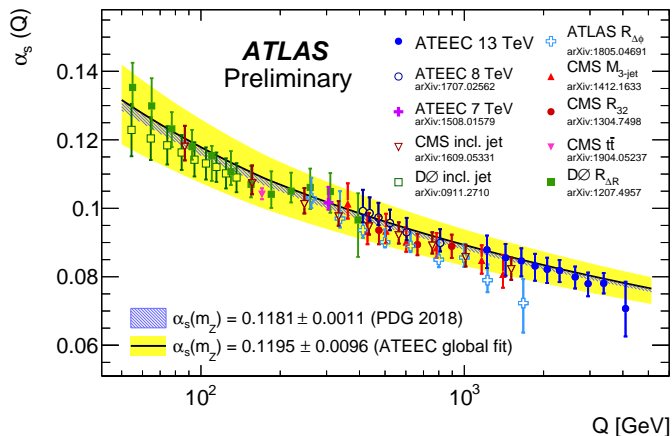
$p_T > 60$  GeV

anti- $k_t$  R = 0.4





These data could be used to improve limits on masses for new coloured fermions obtained by [J. Llorente and B. Nachman, Nucl. Phys. B936 \(2018\) 106](#) using previous measurements at 8 TeV.



- CMS has measured  $W^\pm$  and  $Z$  boson production x-sections in the  $e, \mu$  channels at 7 and 8 TeV
- Comparisons with MCFM calculations at NNLO accuracy interfaced with LHAPDF  $\Rightarrow$  extracted  $\alpha_s(m_Z) = 0.1186 \pm 0.0025$  with MMHT 14
  
- ATLAS measured TEEC and ATTEC at 13 TeV
- Comparison with NLOJET++ predictions  $\Rightarrow$  extracted  $\alpha_s(Q)$  with  $Q$  in the range 1 – 4 TeV
- From a global fit to the TEEC data  $\Rightarrow$   
 $\alpha_s(m_Z) = 0.1196 \pm 0.0004(\text{exp}) \pm 0.0089(\text{theo})$
- **NNLO corrections** needed to bring down the scale systematic uncertainties