

# CASCADE vrs pp data at LHC

H. Jung (CERN, DESY & Antwerp)

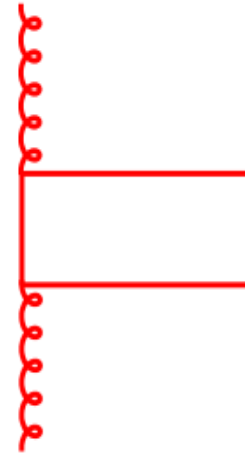
- basic ideas of the CASCADE MC generator
- first comparison with pp data
  - min bias
  - heavy flavor (open & hidden)
- prospects for forward jets

# CASCADE basic idea

- CASCADE elements are:

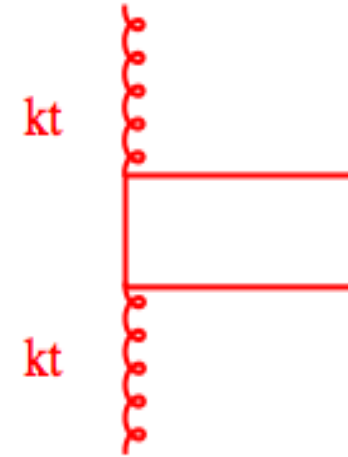
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- CASCADE elements are:
  - Matrix Elements:



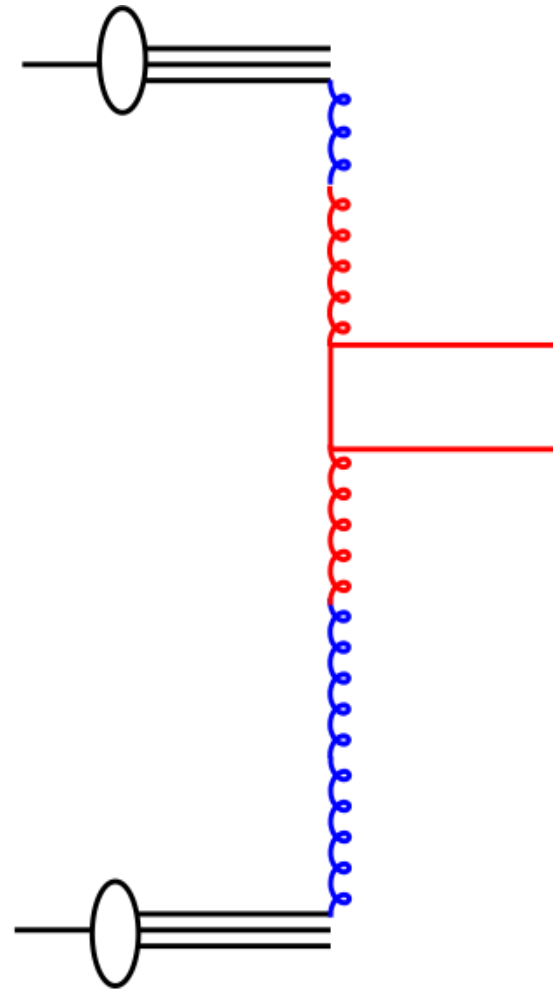
# CASCADE basic idea

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    - on shell/off shell



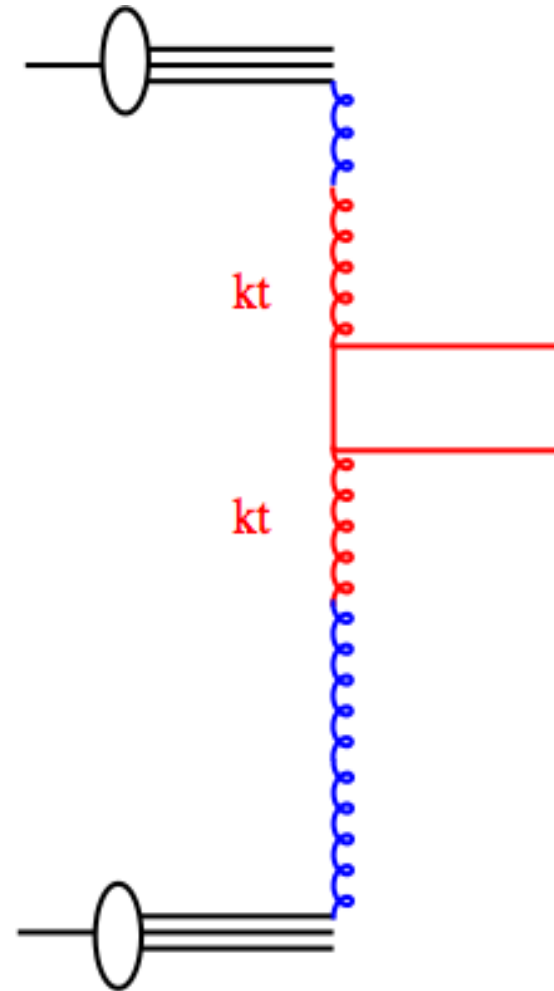
# CASCADE basic idea

- CASCADE elements are:
  - Matrix Elements:
    - on shell/off shell
  - PDFs



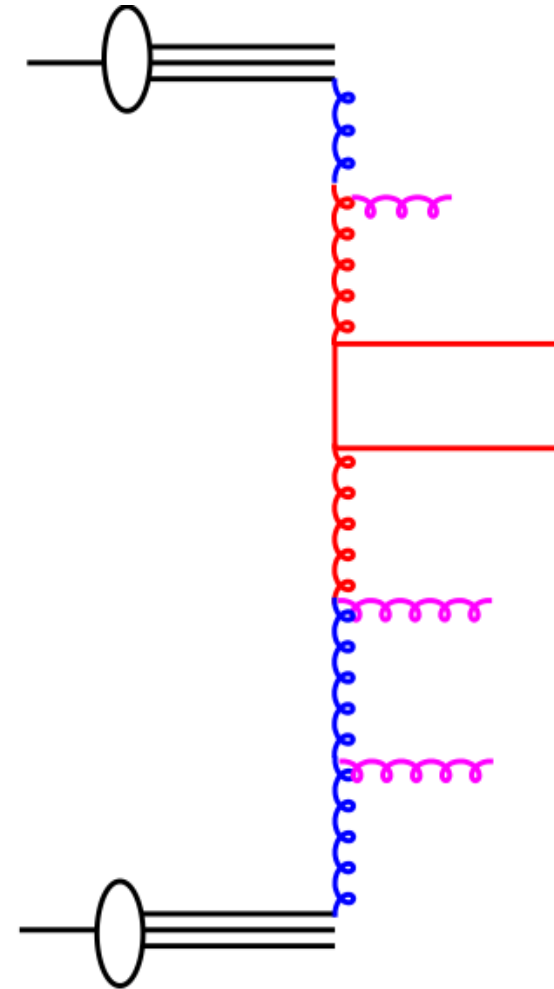
# CASCADE basic idea

- CASCADE elements are:
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  - PDFs
    - unintegrated PDFs



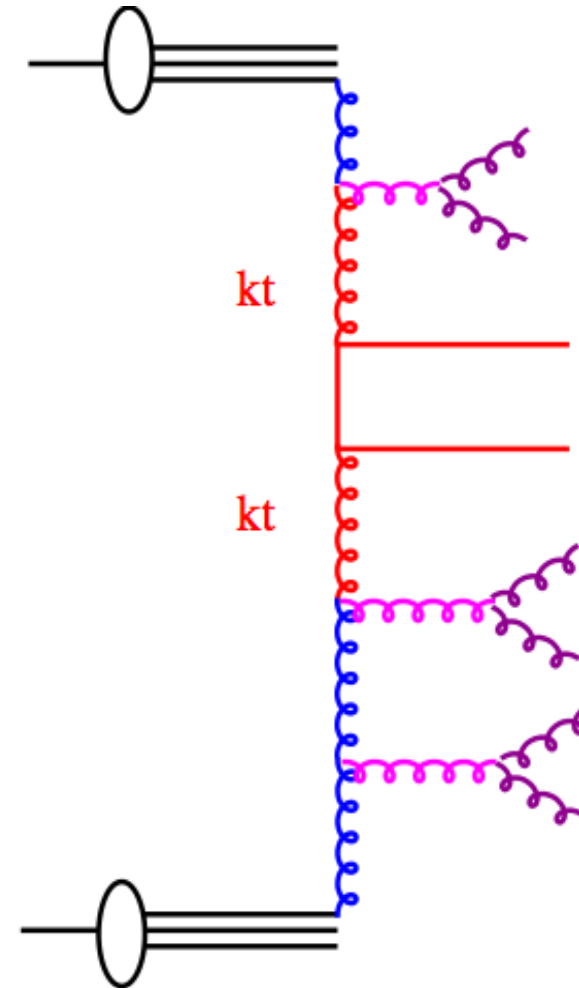
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    - unintegrated PDFs
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# CASCADE basic idea

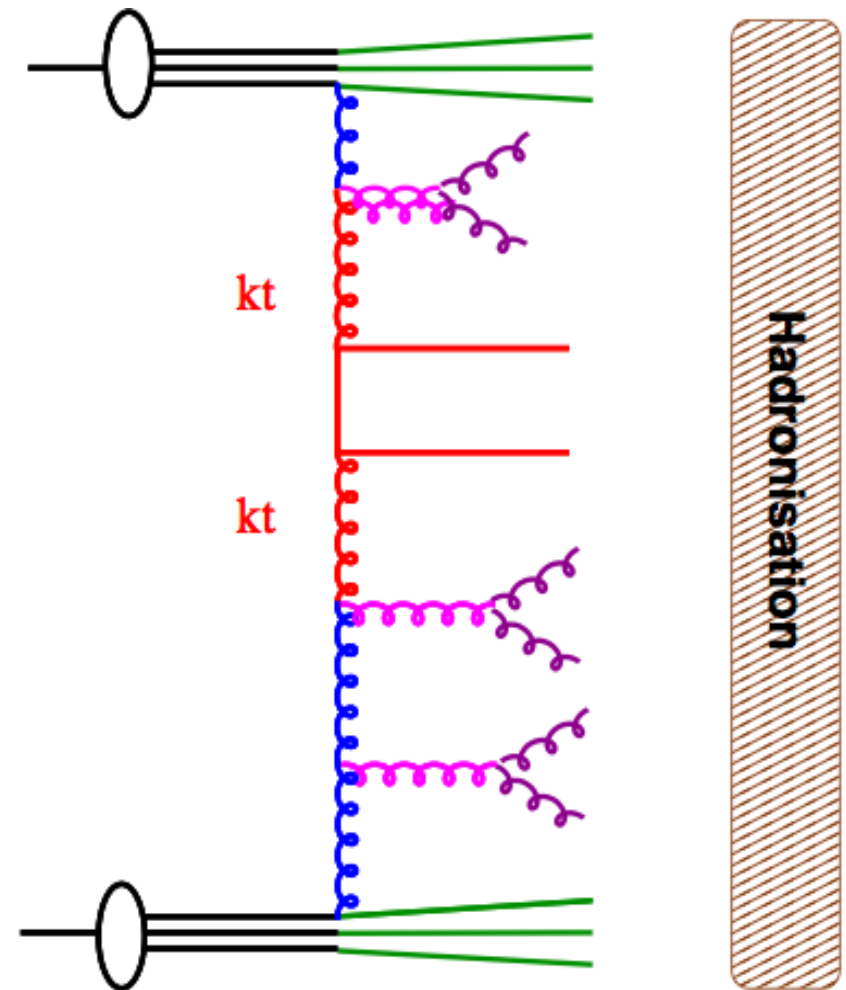
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    - unintegrated PDFs
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    - angular ordering (CCFM)





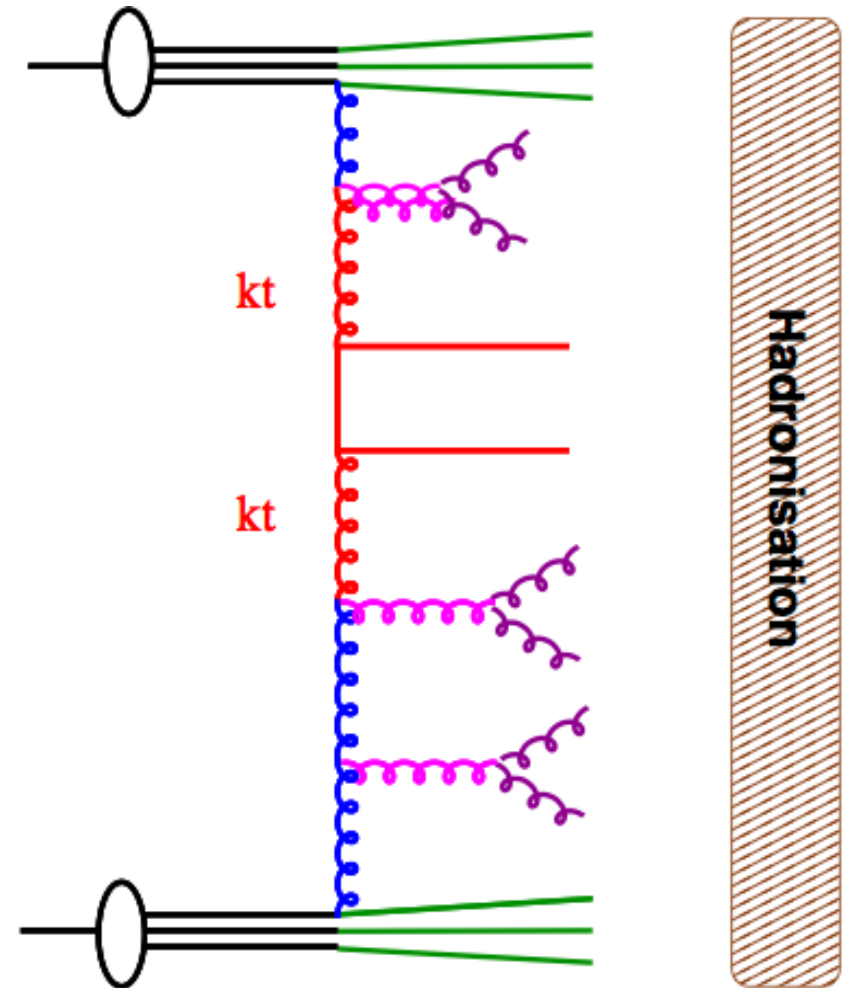
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$$\sigma(pp \rightarrow q\bar{q} + X) = \int \frac{dx_{g1}}{x_{g1}} \frac{dx_{g2}}{x_{g2}} \int d^2 k_{t1} d^2 k_{t2} \hat{\sigma}(\hat{s}, k_t, \bar{q}) \times x_{g1} \mathcal{A}(x_{g1}, k_{t1}, \bar{q}) x_{g2} \mathcal{A}(x_{g2}, k_{t2}, \bar{q})$$

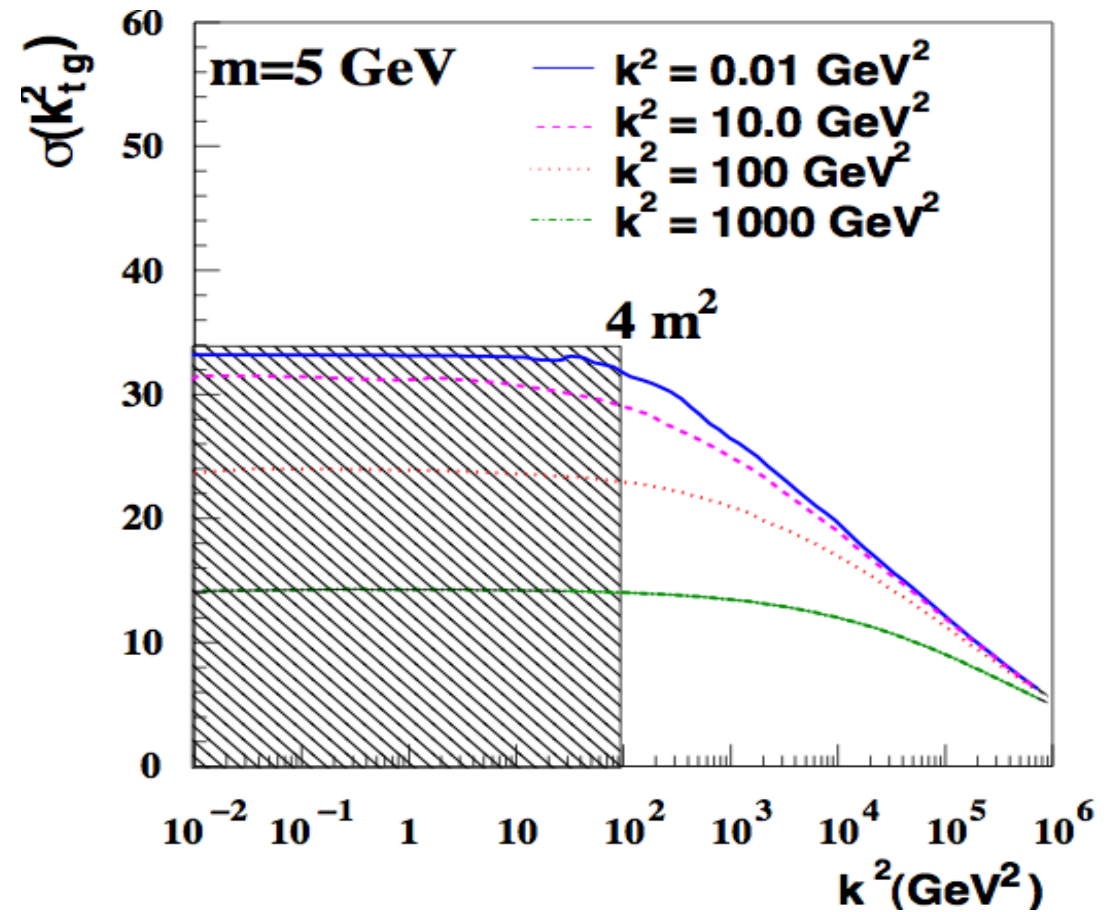
# Why off-shell matrix elements ?

- Example:  $g^* g^* \rightarrow Q \bar{Q}$ 
  - ME is finite for  $k_{\perp} \rightarrow 0$
  - ME has tail to large  $k_{\perp}$
- collinear factorization:

- integration over  $k_{\perp}$

$$\int_0^{\mu^2} dk_{\perp} \hat{\sigma}(k_{\perp}, \dots)$$

up to  $\mu^2 \sim 4m^2$



# Which off-shell matrix elements ?

- heavy quarks

**NEW**

including excited states:

- 1S, 2S, 3S
- 1P, 2P

$$g^* g^* \rightarrow Q\bar{Q}$$

$$g^* g^* \rightarrow J/\psi(\Upsilon)g$$

$$g^* g^* \rightarrow \chi_c(\chi_b)$$

- Gauge boson & Higgs

**NEW**

$$g^* g^* \rightarrow h$$

$$g^* g^* \rightarrow Z + Q\bar{Q}$$

$$g^* g^* \rightarrow W + q_i q_j$$

$$qg^* \rightarrow Zq$$

- QCD processes – forward jets

$$g^* g^* \rightarrow q\bar{q}$$

$$qg^* \rightarrow qg$$

$$gg^* \rightarrow gg$$

# Which uPDFs ?

➡ take derivative of integrated PDF:

$$f(x, k_{\perp}^2) = \frac{dg(x, k_{\perp}^2)}{dk_{\perp}^2} = \left[ \frac{\alpha_s}{2\pi} \int_x^{1-\delta} P(z) g\left(\frac{x}{z}, k_{\perp}^2\right) dz \right]$$

➡ KMR approach:

$$f(x, k_{\perp}^2, \mu^2) = \exp\left(-\int_{k_{\perp}^2}^{\mu^2} \frac{\alpha_s}{2\pi} d \log k_{\perp}^2 \sum_i \int_0^1 P(z') dz'\right) \times \frac{dg(x, \mu^2)}{d\mu^2}$$

using integrated PDF, only last emission generates transverse momentum via sudakov form factor. ...

➡ appropriate for DGLAP with strong ordering....

➡ this is what is done in all standard parton shower MCs

# Which uPDFs ? CCFM approach

- Color coherence requires angular ordering instead of  $p_t$  ordering ...

$$q_i > z_{i-1} q_{i-1} \quad \text{with} \quad q_i = \frac{p_{ti}}{1 - z_i}$$

- recover DGLAP with  $q$  ordering

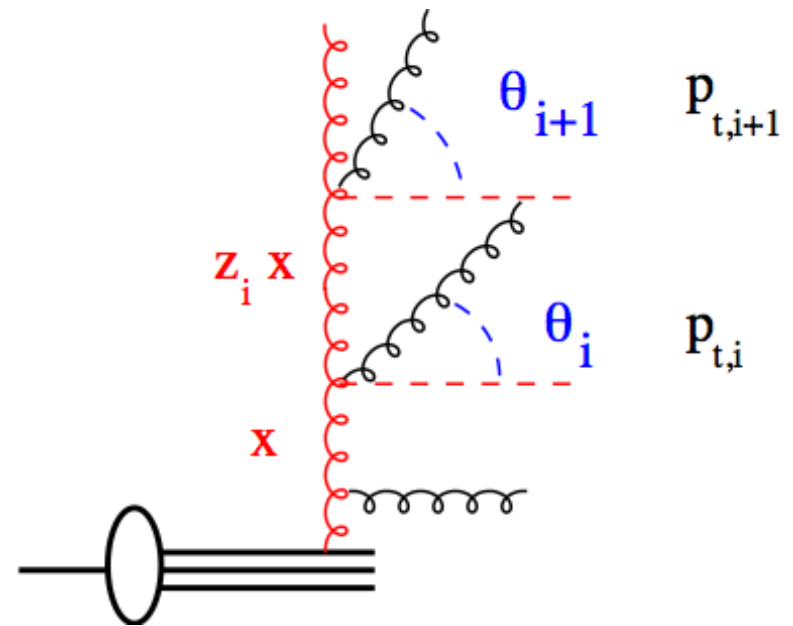
at medium and large  $x$

- HERWIG uses:

$$q_i > q_{i-1}$$

- at small  $x$ , no restriction on  $q$

$p_{ti}$  can perform a random walk



- Catani-Ciafaloni-Fiorani-Marchesini evolution forms a bridge between DGLAP and BFKL evolution

- important for comparison with collinear NLO calculations ...

# uPDF fit to $F_2$ : x-dependence

$$\chi^2 = \sum_i \left( \frac{(T - D)^2}{\sigma_i^2 \text{stat} + \sigma_i^2 \text{uncor}} \right)$$

- fit parameters of starting distributic

$$x\mathcal{A}_0(x, \mu_0) = Nx^{-B_g} \cdot (1-x)^4$$

- using  $F_2$  data H1

(H1 Eur. Phys. J. C21 (2001) 33-61, DESY 00-181)

$$x < 0.05 \quad Q^2 > 5 \text{ GeV}^2$$

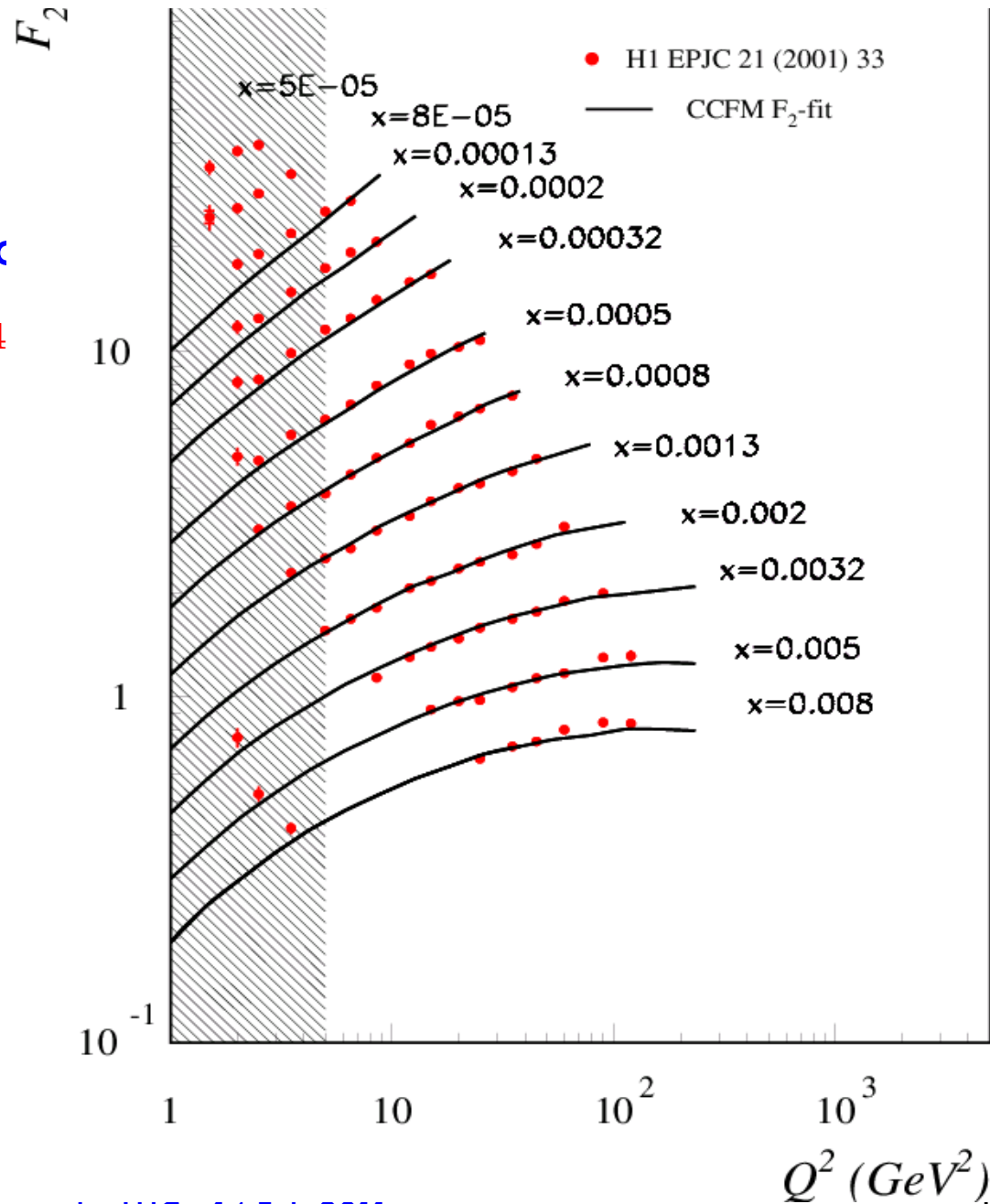
- parameters:  $\mu_r^2 = p_t^2 + m_{q,Q}^2$

$$m_q = 250 \text{ MeV}, m_c = 1.5 \text{ GeV}$$

- Fit (only stat+uncorr):

$$\frac{\chi^2}{\text{ndf}} = \frac{111.8}{61} = 1.83$$

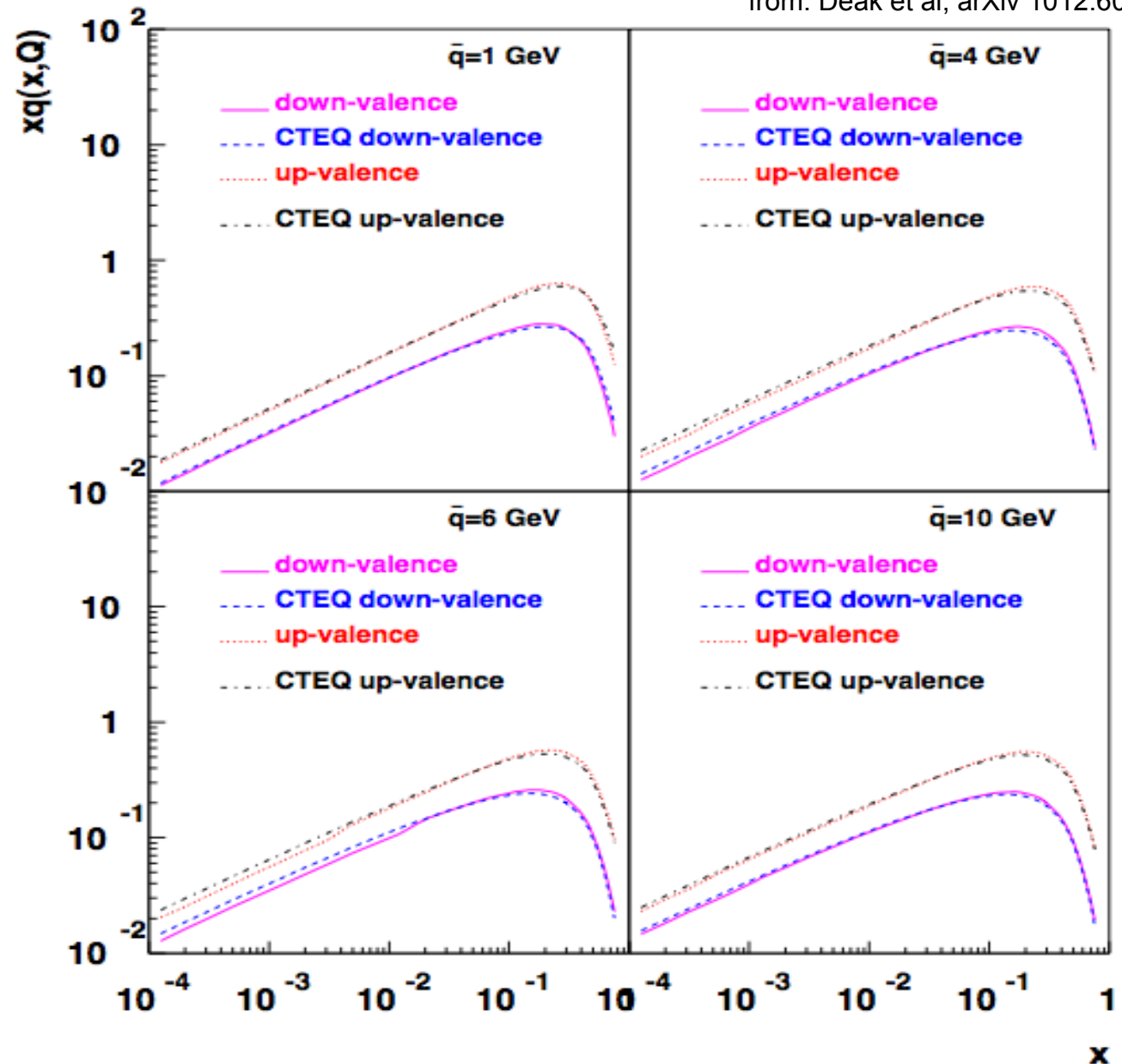
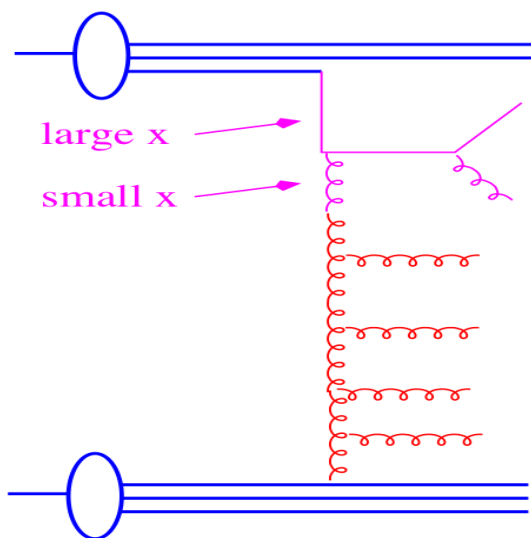
$$B_g = 0.028 \pm 0.003$$



# unintegrated valence quarks

from: Deak et al, arXiv 1012.6037

- **unintegrated valence quarks:**
- use CTEQ61 as initial condition
- evolve with "CCFM-type" splitting function
- needed for:

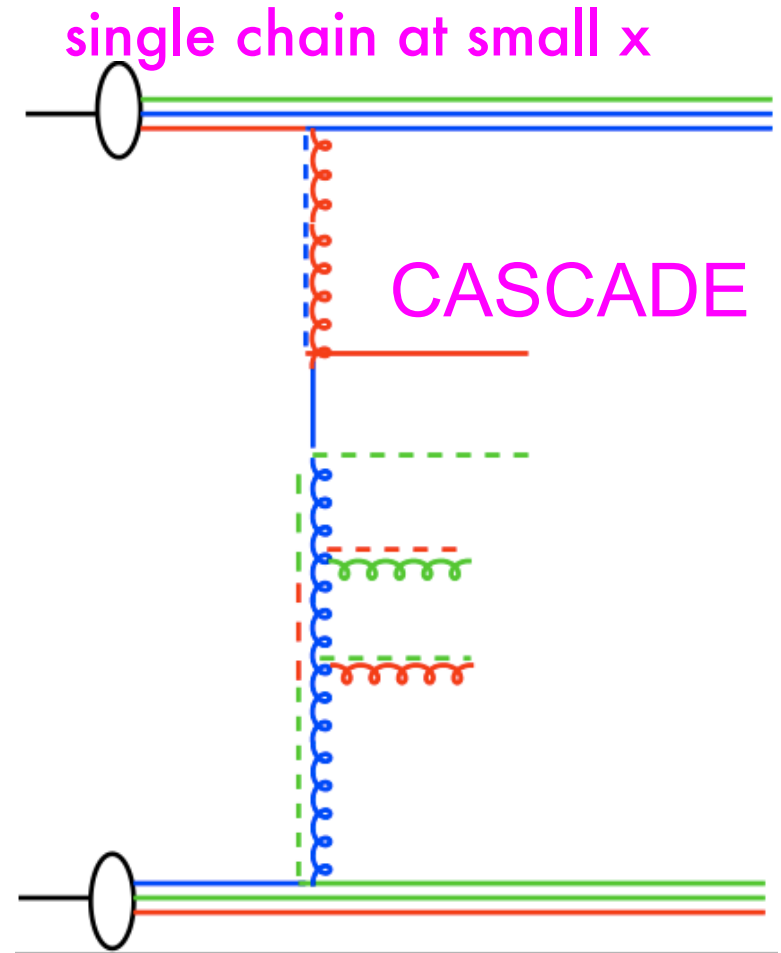
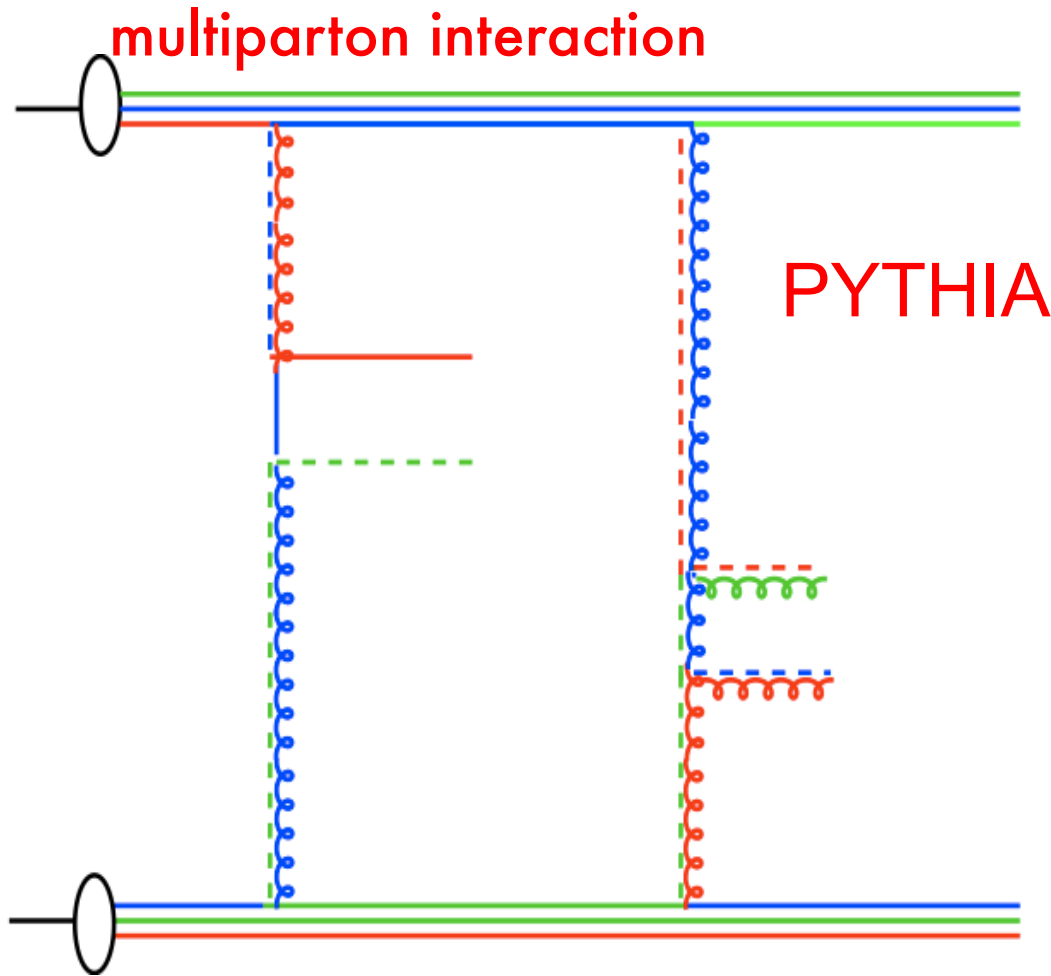




# Comparison to pp data: Minimum Bias

# How well do we know soft parton radiation?

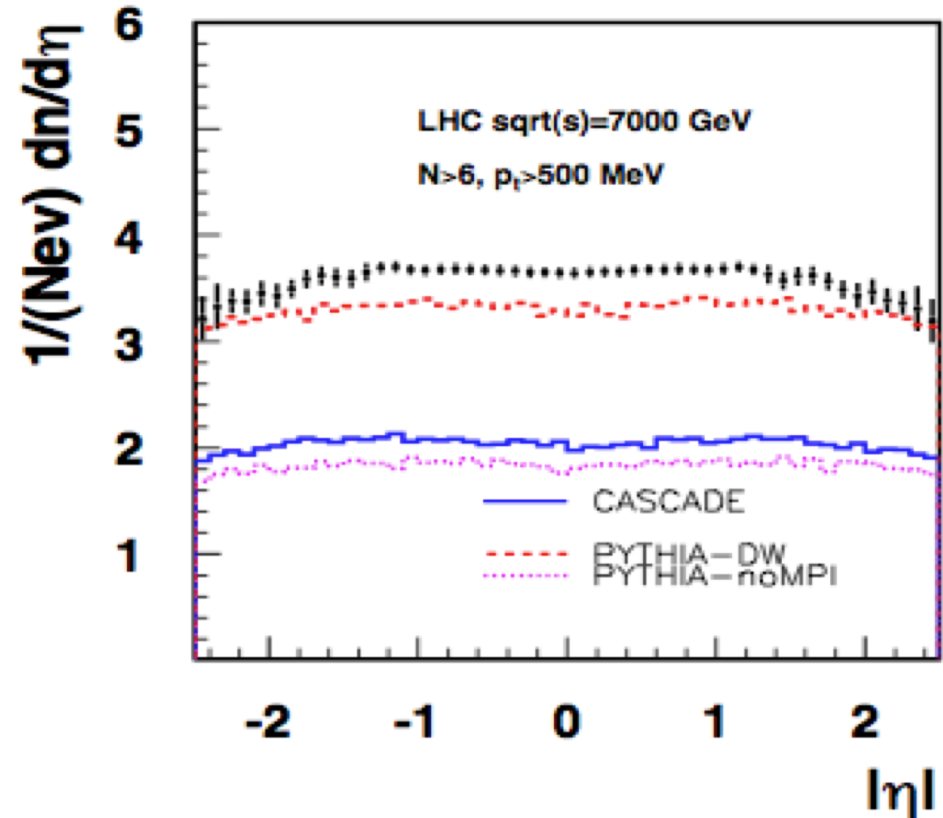
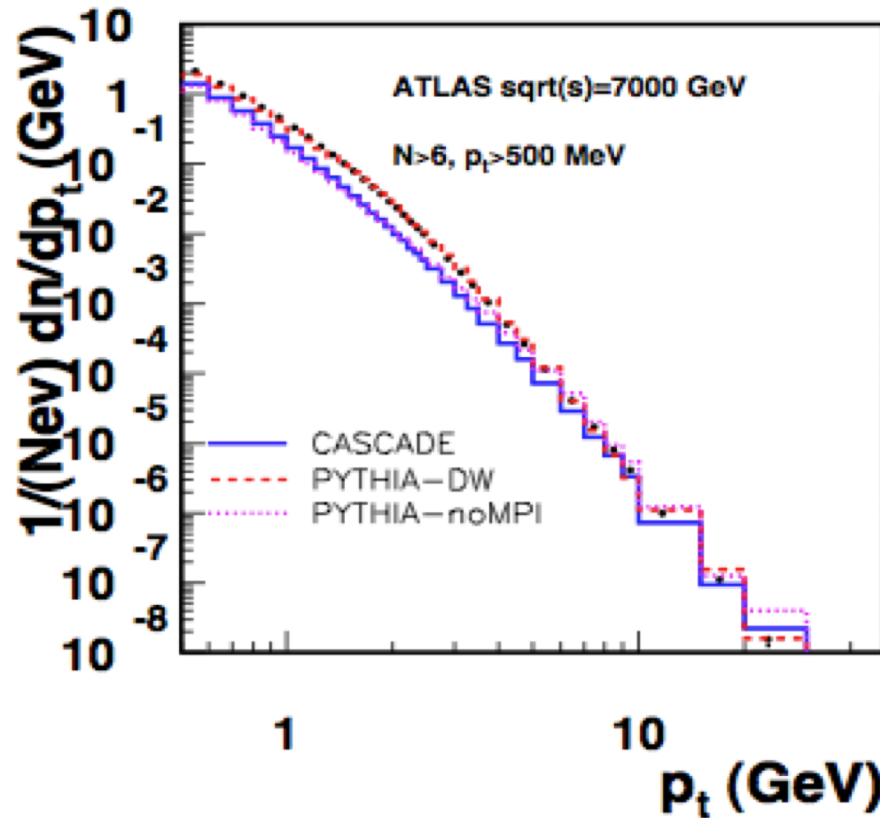
- soft parton radiation:



- which of the two is correct or are they both describing the same ... ???

# Charged particle spectra in Minbias

$$g^*g^* \rightarrow qq\bar{q}, g^*g \rightarrow gg, g^*q \rightarrow qq$$



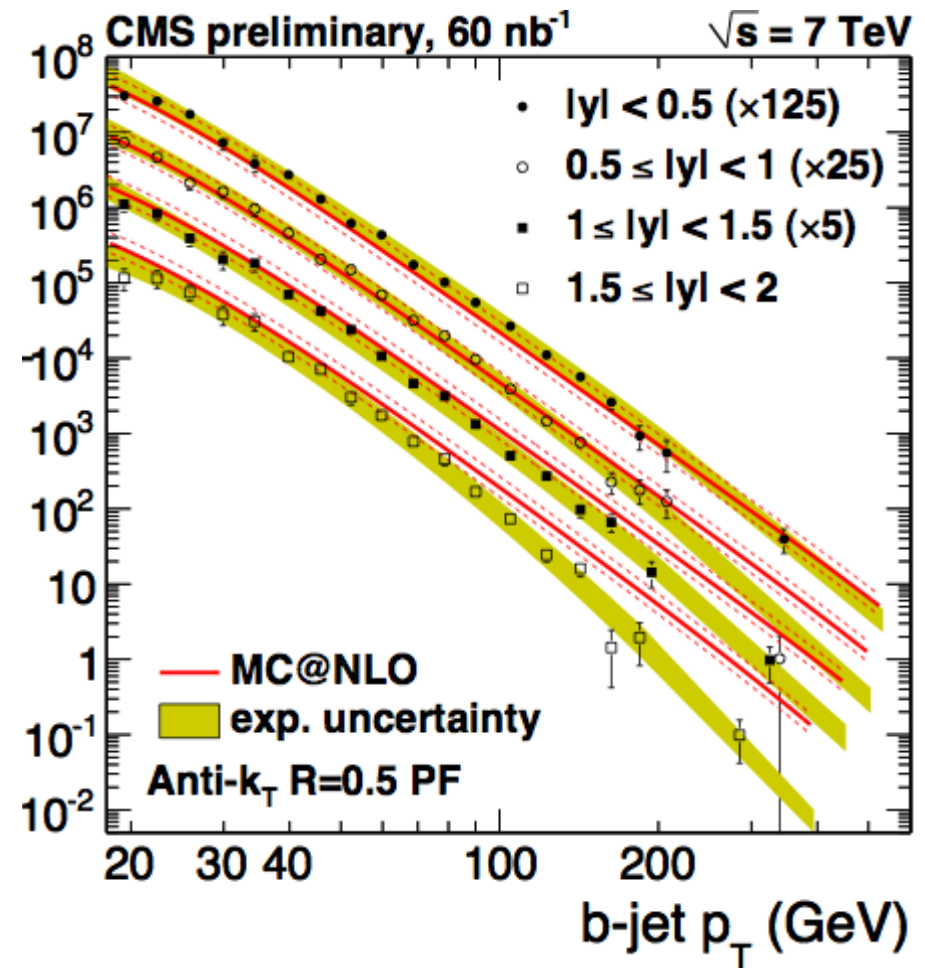
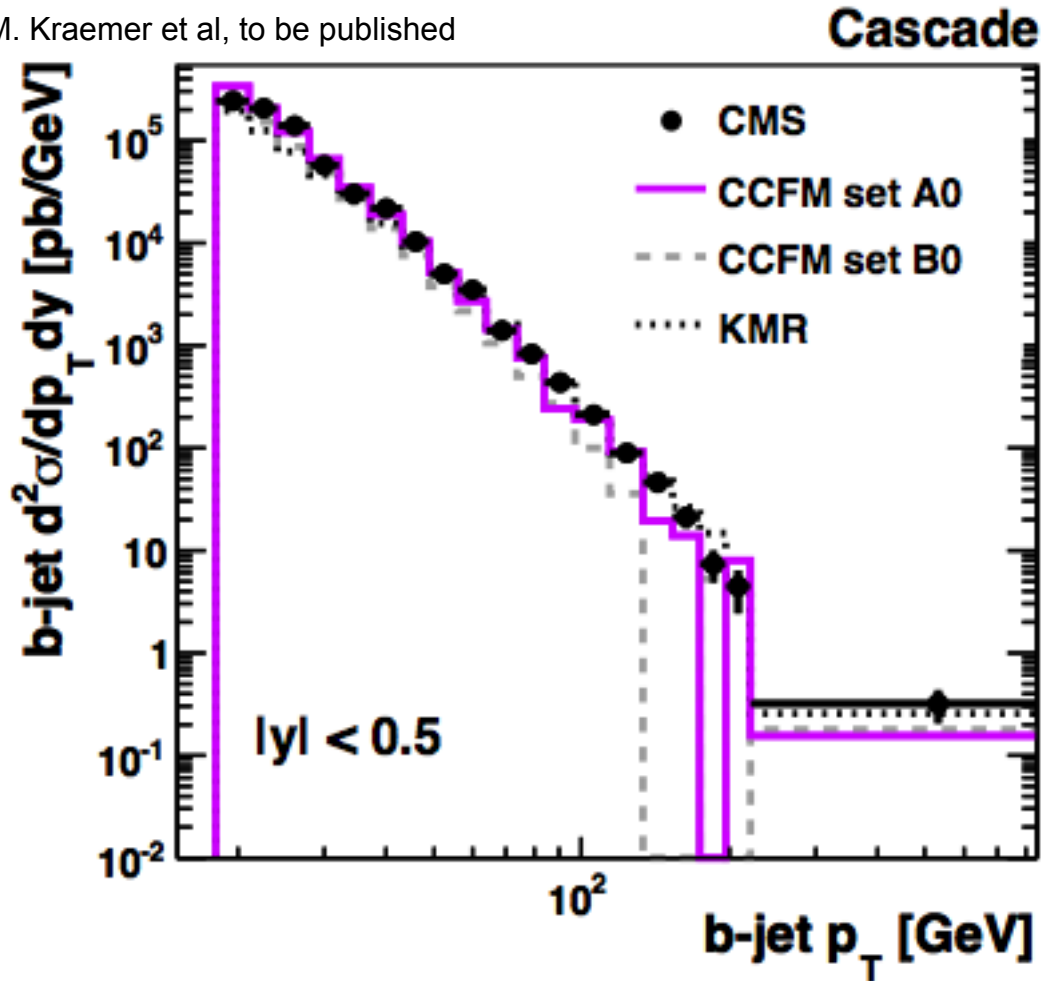
- clear deficit of particle production in the soft region
  - region where multiparton interactions play dominant role
  - CASCADE larger than PYTHIA w/o MPI

**Comparison to pp data:  
hard processes - heavy flavors**

# Open Heavy Flavor production

$$g^* g^* \rightarrow Q \bar{Q}$$

M. Kraemer et al, to be published

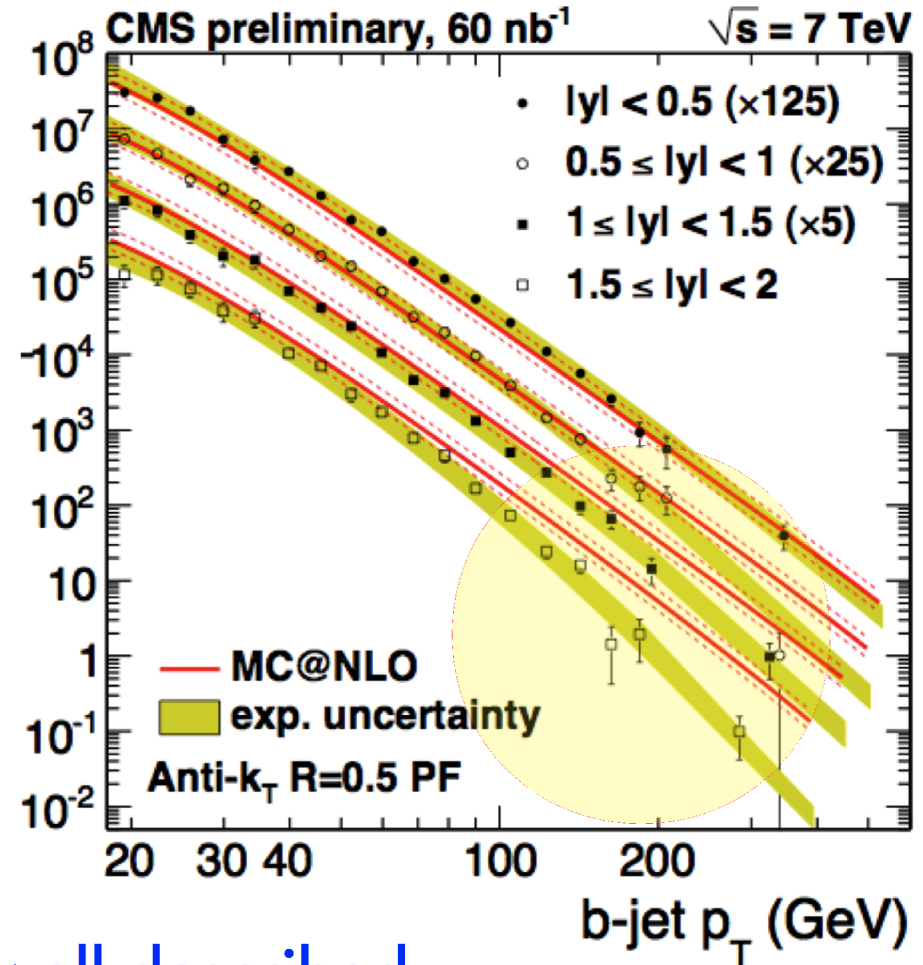
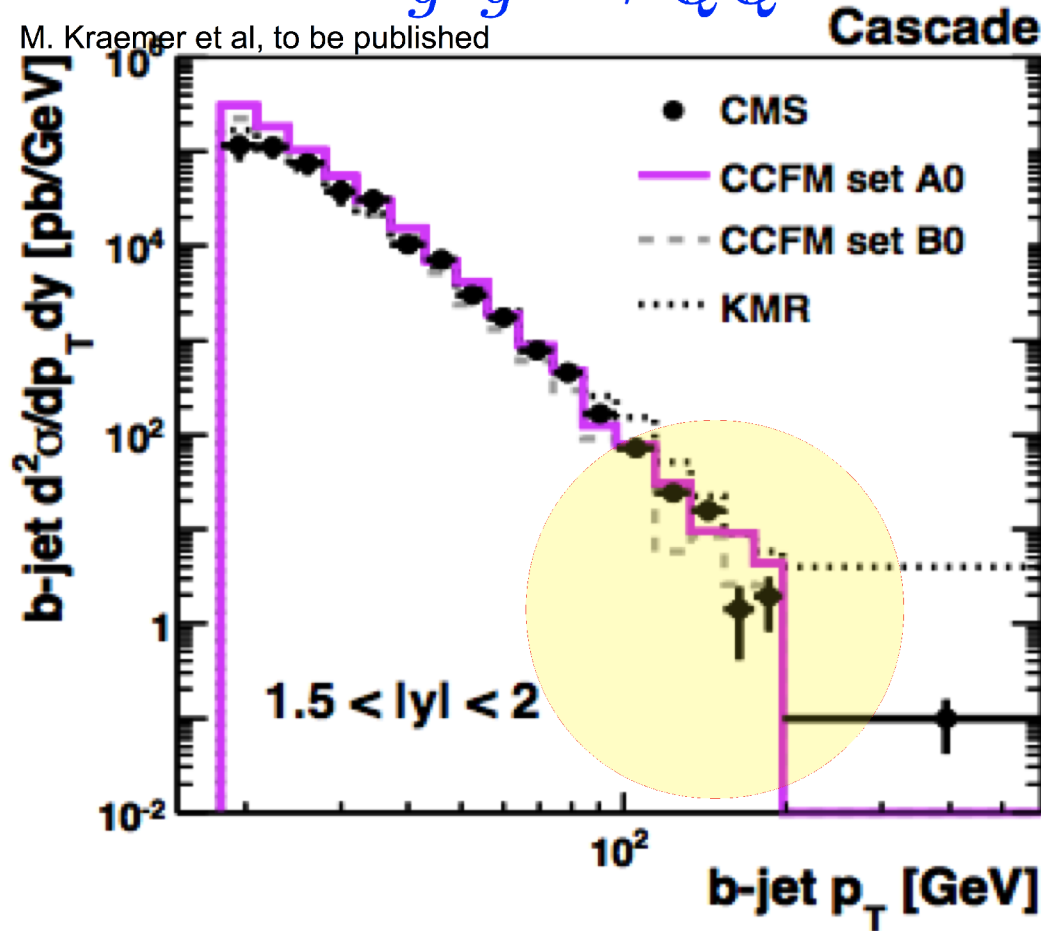


- at small  $|y|$  xsection well described (similar to MC@NLO)

# Open Heavy Flavor production

$$g^* g^* \rightarrow Q Q \bar{Q}$$

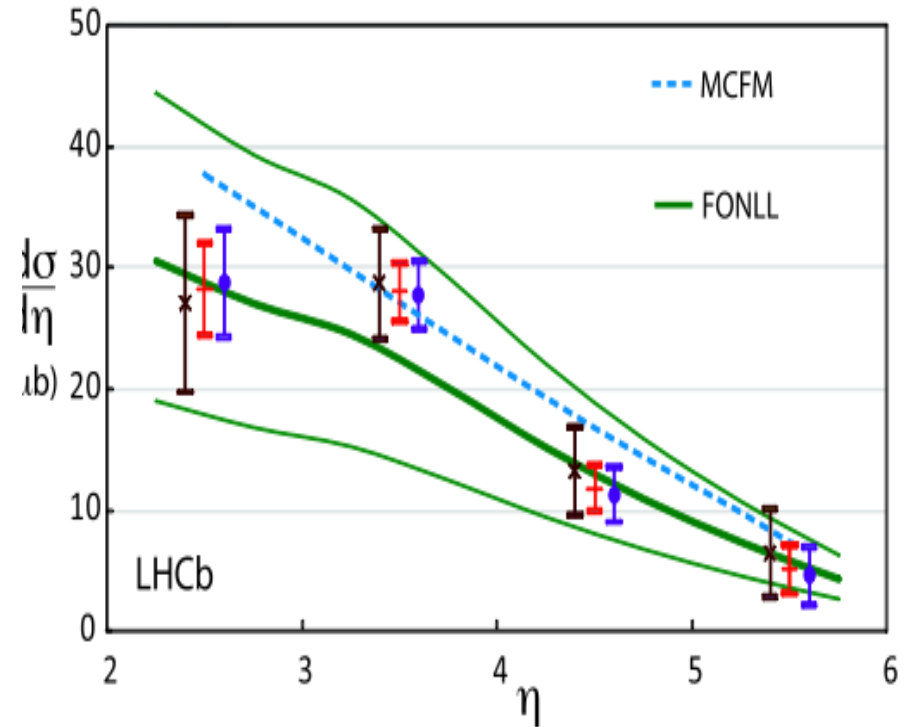
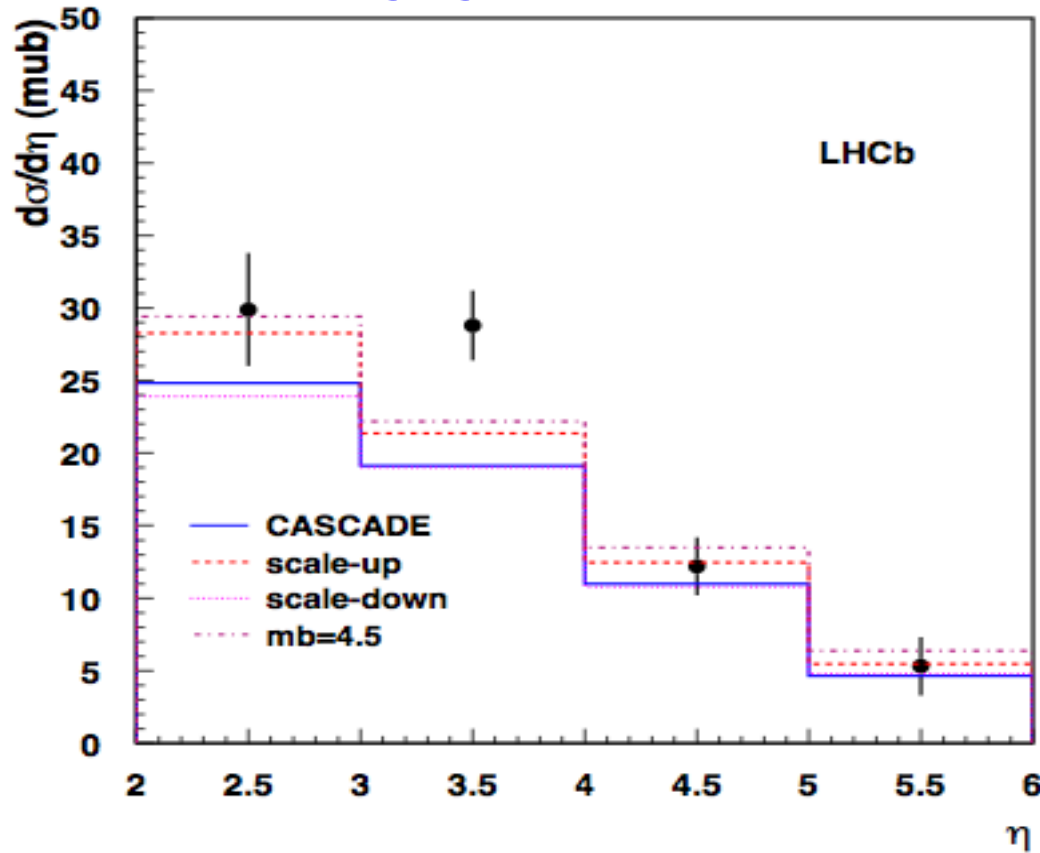
M. Kraemer et al, to be published



- at large  $1.5 < |y| < 2.0$  x-section well described
  - effect of suppression in off-shell ME ?
  - note: MC@NLO is too large at large  $p_T$  ...

# Forward B – production (LHCb)

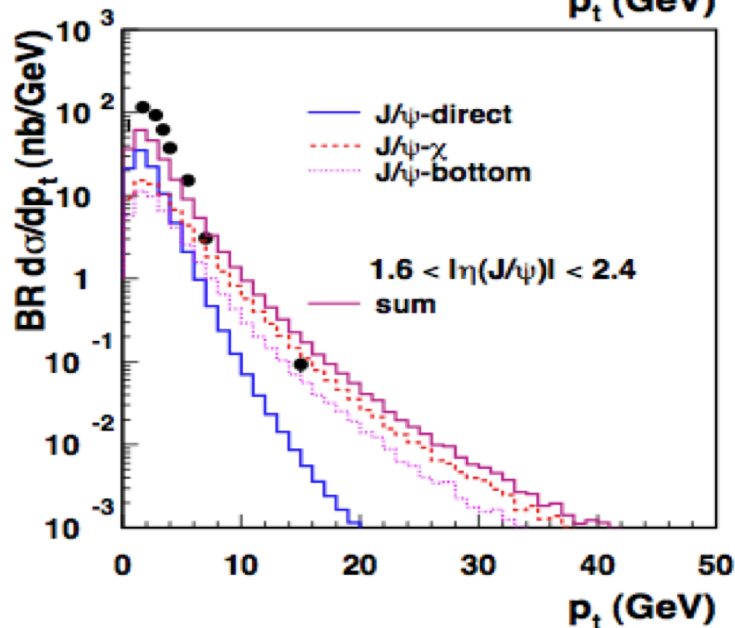
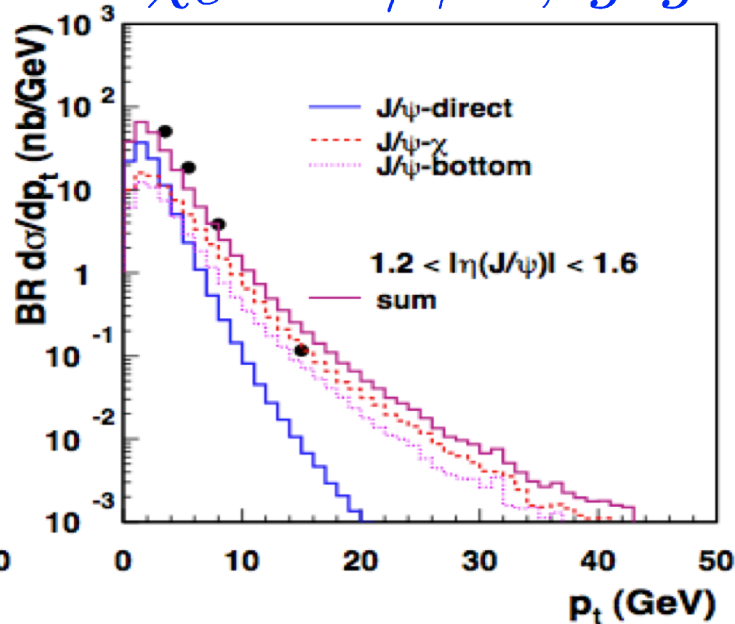
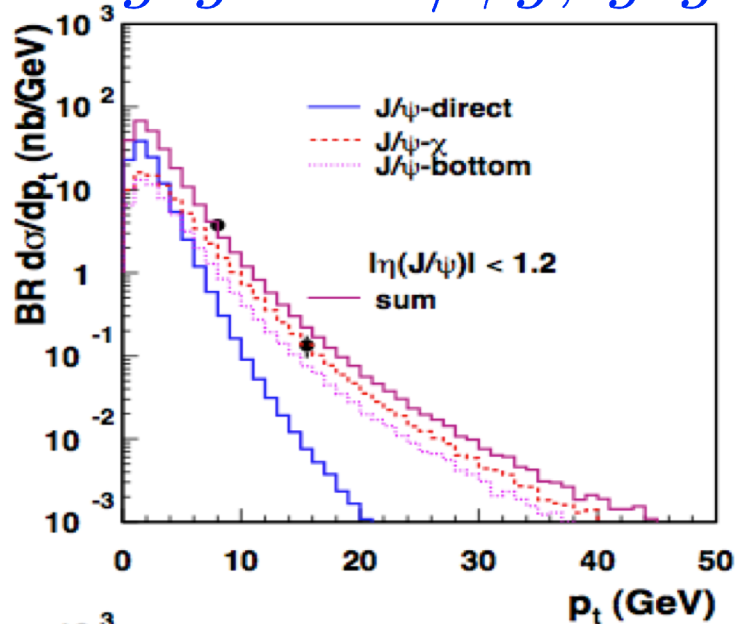
$$g^* g^* \rightarrow Q Q \bar{Q}$$



- Forward B-hadron production described well
  - “small” uncertainties from  $u$ PDF and  $b$ -mass

# Inelastic $J/\psi$ production

$$g^* g^* \rightarrow J/\psi g, g^* g^* \rightarrow \chi_c \rightarrow J/\psi X, g^* g^* \rightarrow b\bar{b} \rightarrow J/\psi X$$

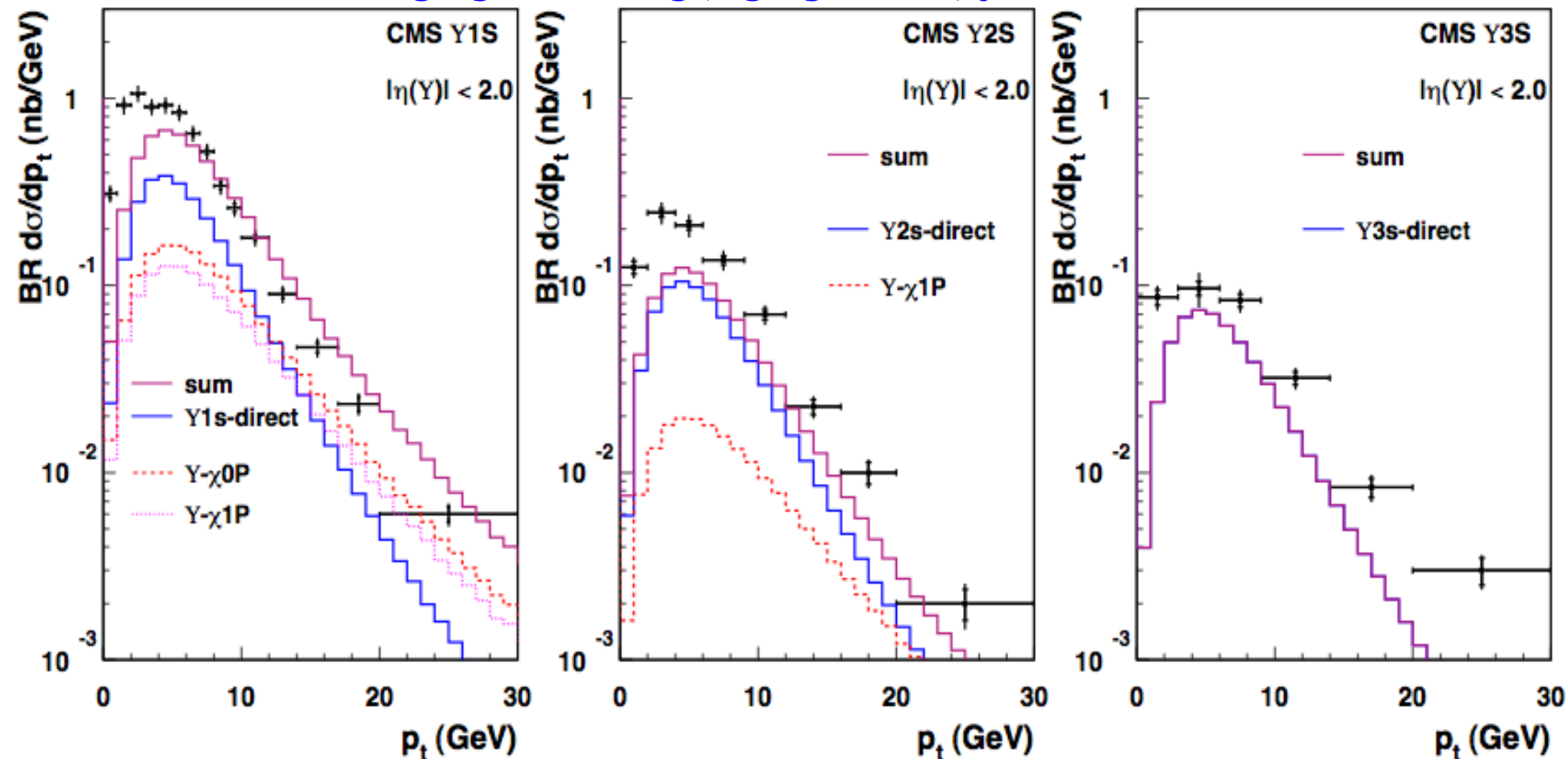


- only CSM contributions included
- cross section depends on wave function
- no free tunable parameters
- possible contribution from color octet processes expected to be very small



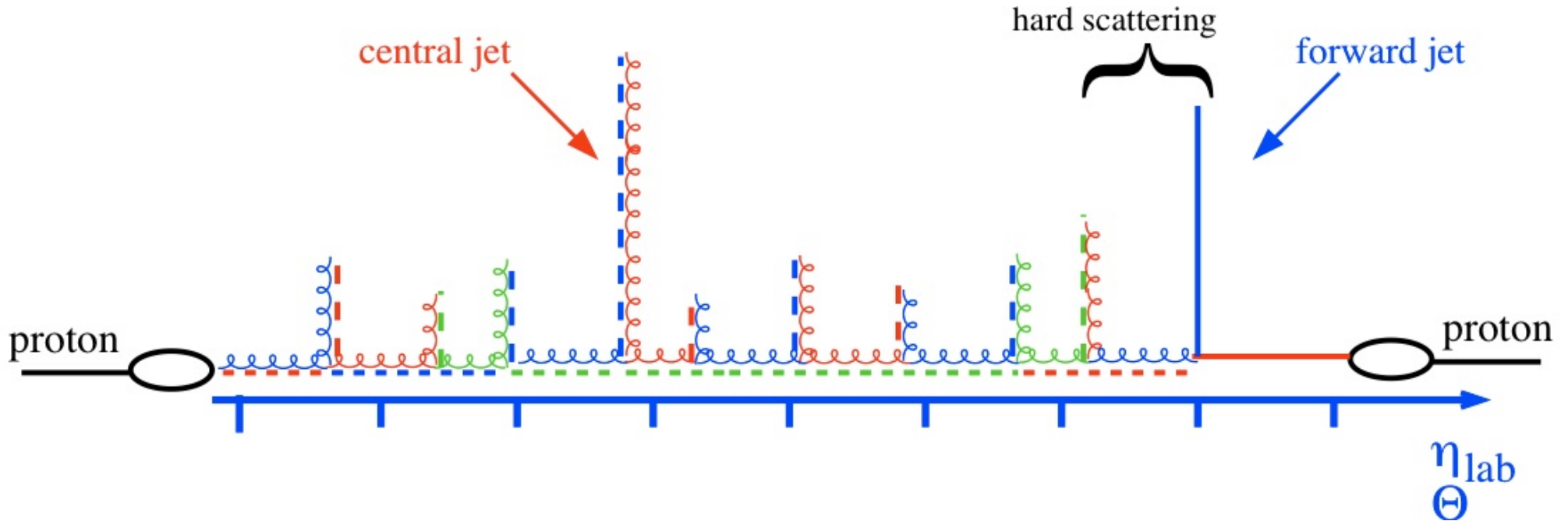
# Inelastic Upsilon production

$$g^*g^* \rightarrow \Upsilon g, g^*g^* \rightarrow \chi_b \rightarrow \Upsilon + X$$



- Only color singlet model contributions included
- excited states calculated explicitly including proper wavefunctions
- NO tunable parameters, except uPDF, scales in  $\alpha_s$

# Associated forward jet production



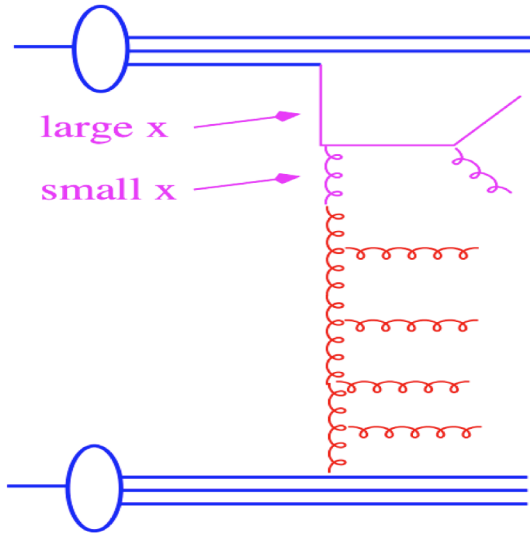
- **forward jet**  $E_t > 10(30) \text{ GeV}, 3 < |\eta| < 5$
- **central jet**  $E_t > 10(30) \text{ GeV}, |\eta| < 2$

# Hard processes: forward jets

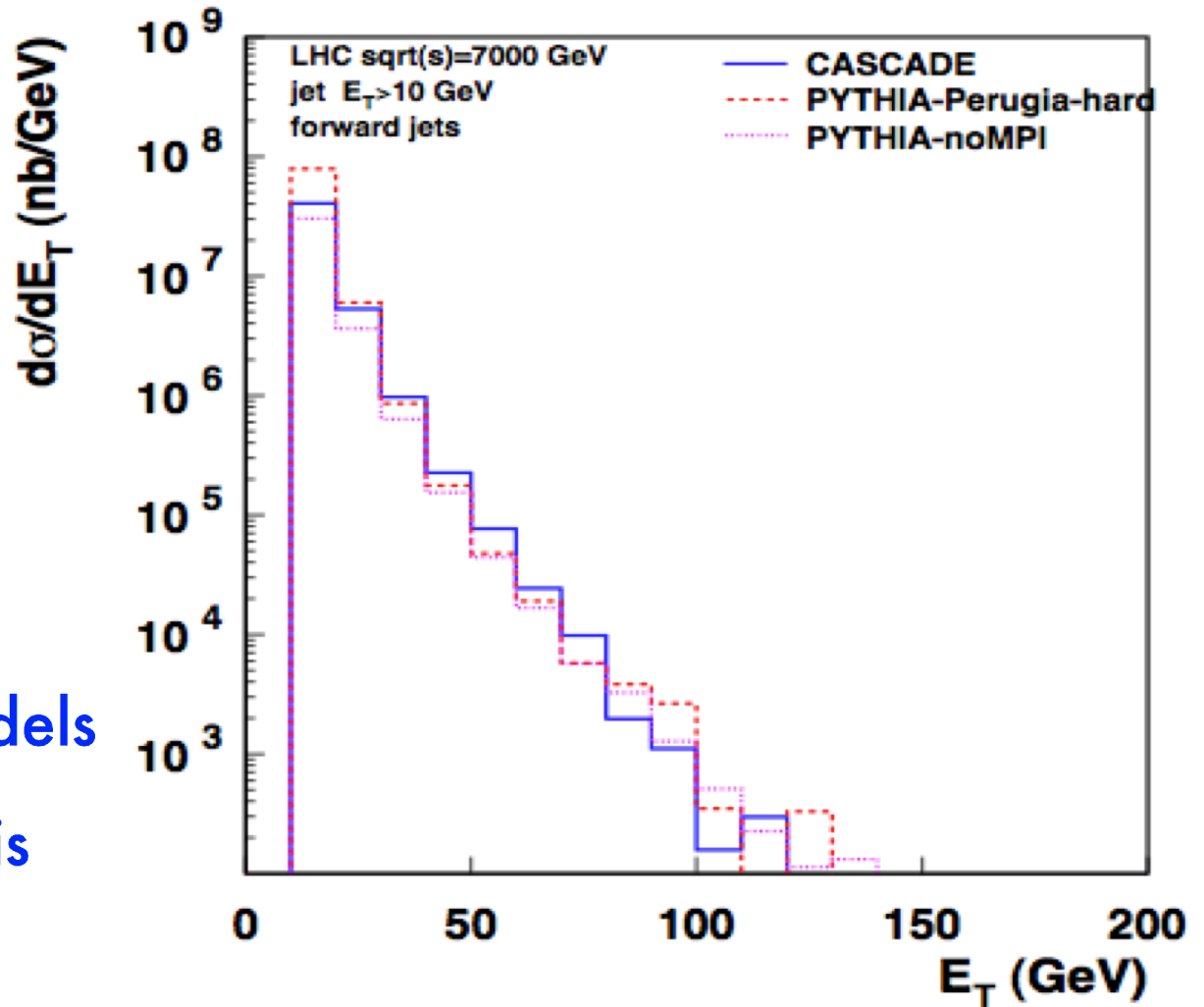
# Forward – central jets

$$g^* g^* \rightarrow q\bar{q}, g^* g \rightarrow gg, g^* q \rightarrow qg$$

from: Deak et al, arXiv 1012.6037



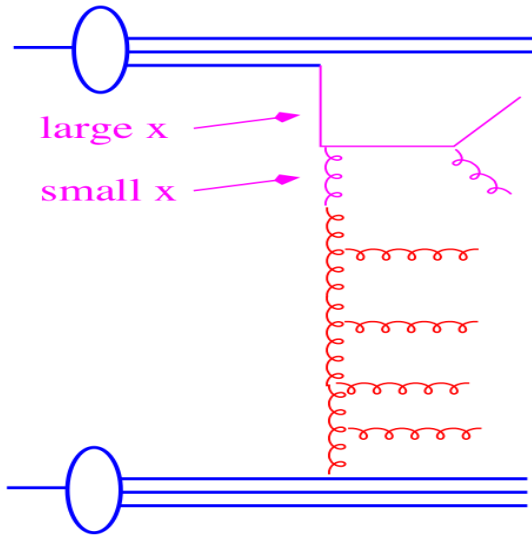
- pt spectra are similar in CASCADE and MPI models
- consistent with what is observed in data



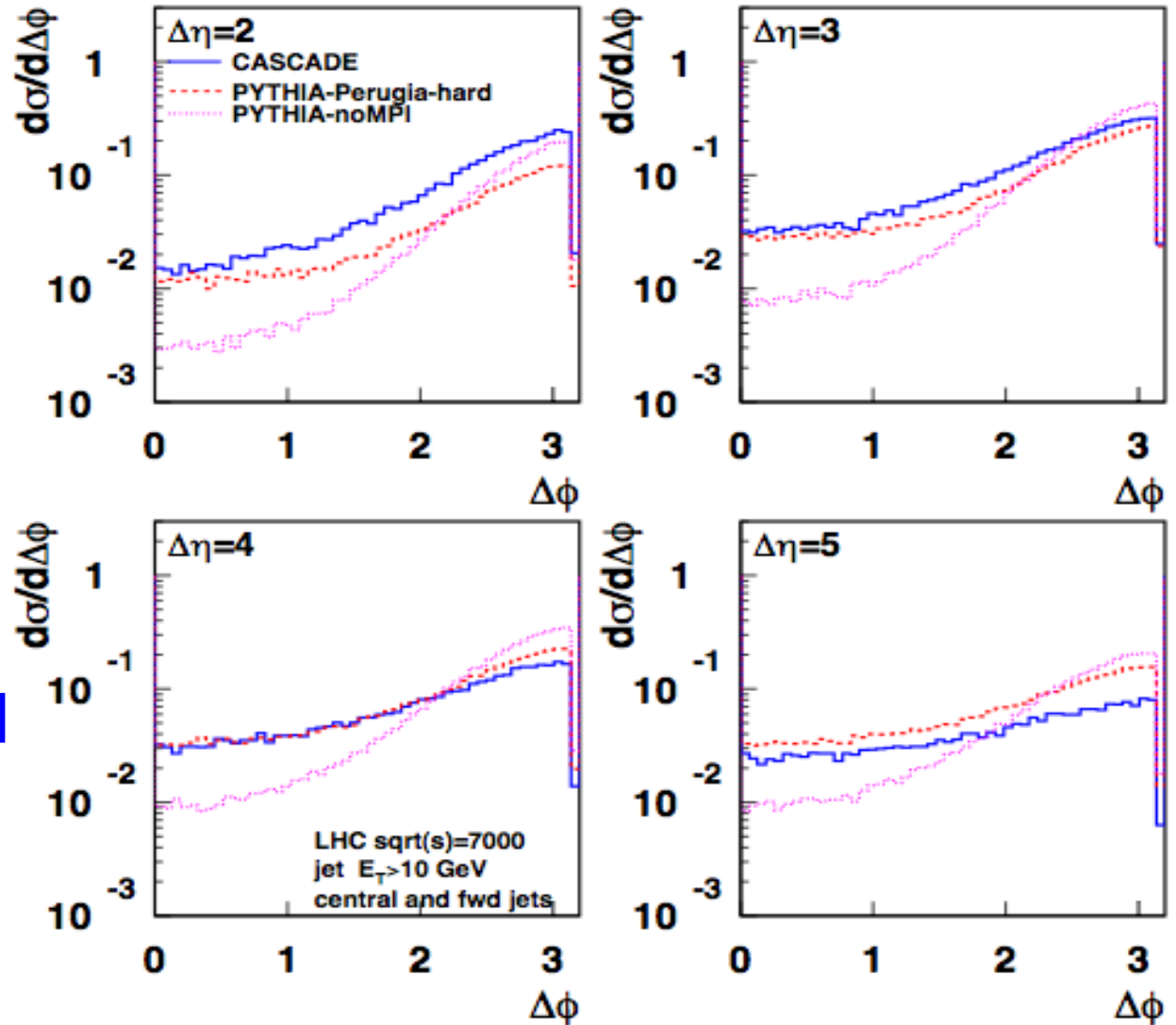
# Azimuthal correlations in fwd-cent jets

$$g^* g^* \rightarrow q\bar{q}, g^* g \rightarrow gg, g^* q \rightarrow qg$$

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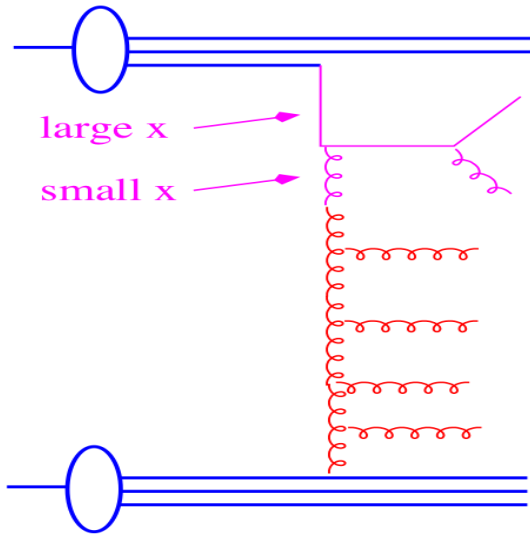
- significant differences between CASCADE and MPI models



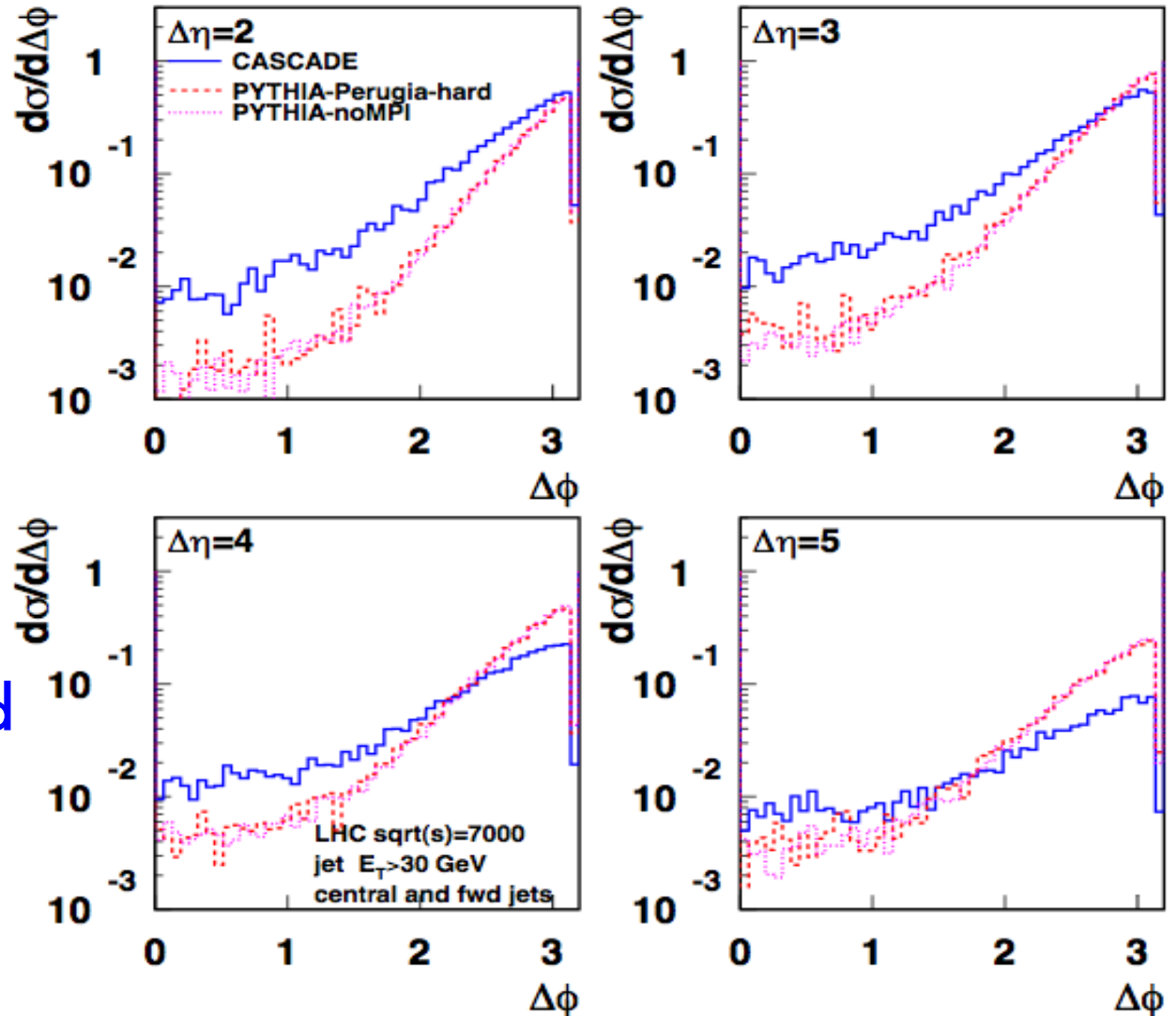
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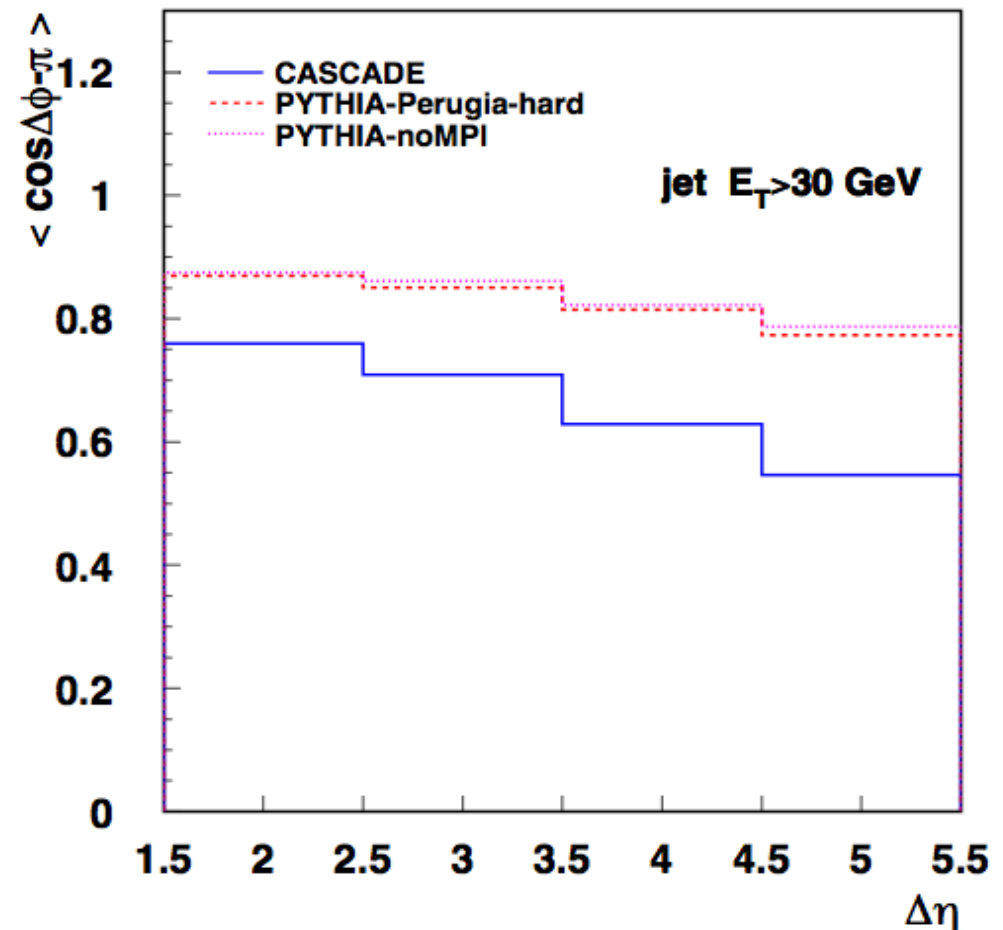
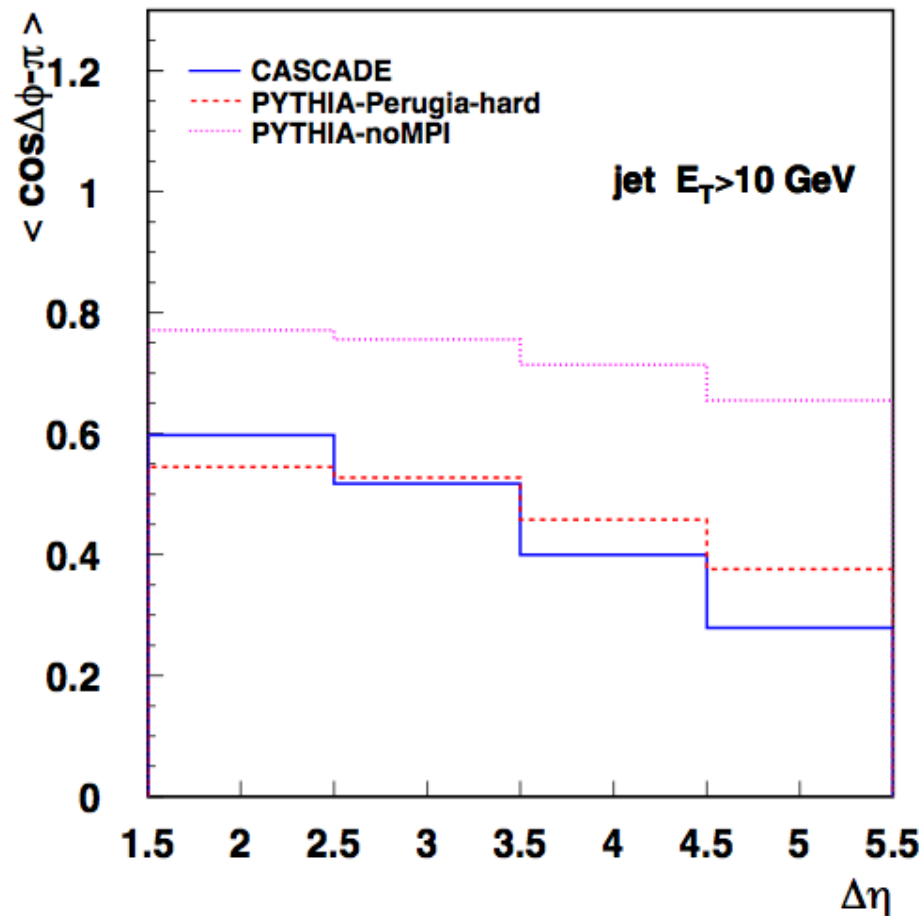
- significant differences between CASCADE and MPI models
- differences also at  $E_t > 30$  GeV



# Azimuthal correlations in fwd-cent jets

$$g^*g^* \rightarrow q\bar{q}, g^*g \rightarrow gg, g^*q \rightarrow qg$$

from: Deak et al, arXiv 1012.6037



- significant de-correlation effects observable
- **BUT**, differential distributions has **discriminate better !**

# Conclusions

- **CASCADE** works promising well in pp
  - specialities are (up to now):
    - heavy flavors (open and hidden)
    - forward (& central) jets
  - even does not too bad for min – bias
- **CASCADE** has essentially **NO** free parameters
  - all parameters are fixed by uPDF
  - gives real predictions



# Conclusions

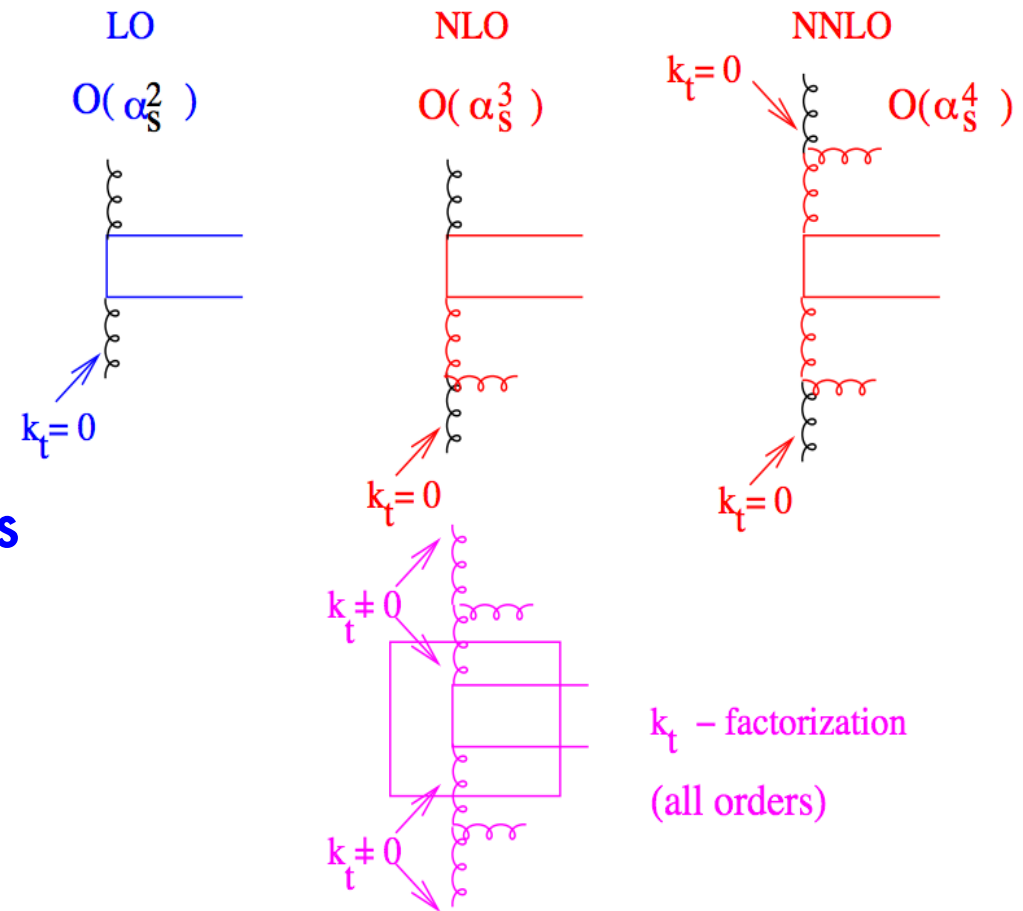
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**You are most welcome  
to use **CASCADE**  
for **your** analysis !**

# Backup slides

# CASCADE and coll. NLO calculations

- fit of uPDF to inclusive structure functions /x-sections used to determine normalization
- includes "all-orders" !!!!
- off-shell matrix element simulates part of real NLO corrections



uPDFs are important....

# CASCADE and NLO: pp

- compare CASCADE with MC@NLO for  $t\bar{t}$  production at LHC

- sudakov suppression at small  $p_t(t\bar{t}\text{-pair})$
- even larger  $p_t$  tail, coming from 2 off-shell gluons

