

Status of EPPS nPDF global analysis and connection with the LHC HI data combination work

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**XUNTA
DE GALICIA**

- Define nPDFs in terms of

$$f_i^{p/A}(x, Q^2) = R_i^A(x, Q^2) f_i^p(x, Q^2)$$

nuclear modification
bound-proton PDF free-proton PDF

- Parametrize the x and A dependence of

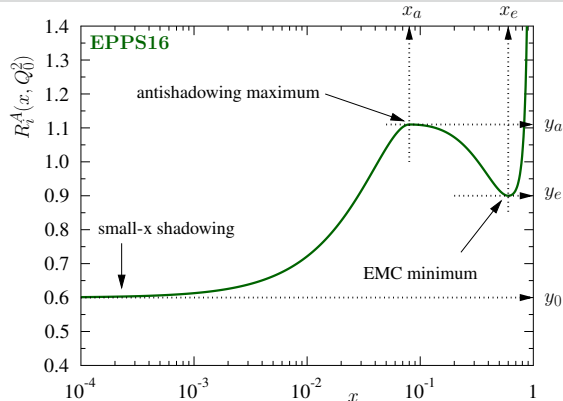
$$R_i^A(x, Q_0^2) \text{ at } Q_0^2 = m_{\text{charm}}^2$$

- PDFs of the full nucleus are then constructed with

$$f_i^A(x, Q^2) = Z f_i^{p/A}(x, Q^2) + N f_i^{n/A}(x, Q^2),$$

where the neutron content is obtained via isospin symmetry

- Allow **full flavour separation** and include heavy-quark mass effects with a general-mass variable flavour number scheme (GM-VFNS)
- Most **extensive data set** to date, with νA DIS, πA DY, LHC pPb dijets and EW bosons



LHC constraints in EPPS16

- No 5.02 TeV pp baseline was available at this time, used forward-to-backward ratios

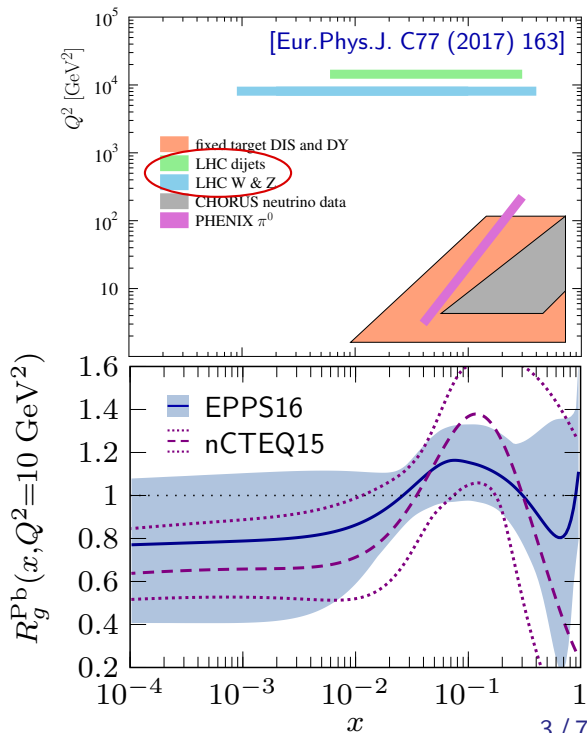
$$R_{\text{FB}} = \frac{d\sigma(\eta > 0)}{d\sigma(\eta < 0)}$$

Cancel part of the free-proton uncertainty, but lose also some information

- Main LHC constraints come from the CMS dijet R_{FB} data
 - ▶ Better control over the gluon antishadowing & EMC effect
- Too low statistics for W & Z R_{FB} to make strong impact

$$R_i^A(x, Q^2) = f_i^{p/A}(x, Q^2) / f_i^p(x, Q^2)$$

nuclear modification bound-proton PDF free-proton PDF

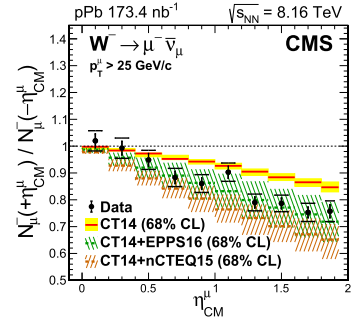
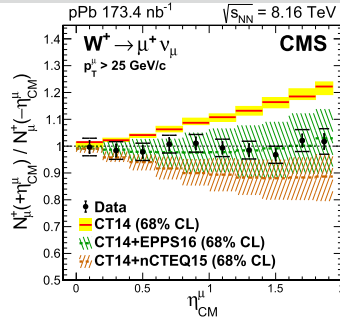


Since then...

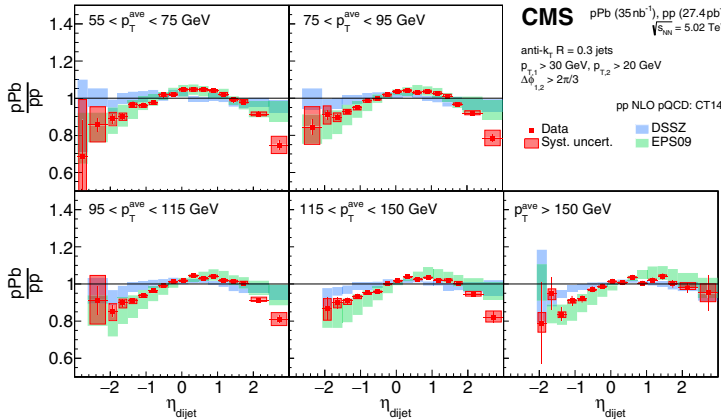
Plentiful of new measurements

- 5.02 TeV R_{pA}
 - ▶ More direct probes of the R_i^A than R_{FB}
- 8.16 TeV R_{FB}

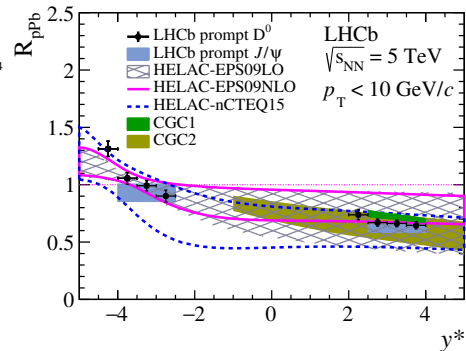
[Phys.Lett.B 800 (2020) 135048]



[Phys.Rev.Lett. 121 (2018) 062002]



[JHEP 1710 (2017) 090]



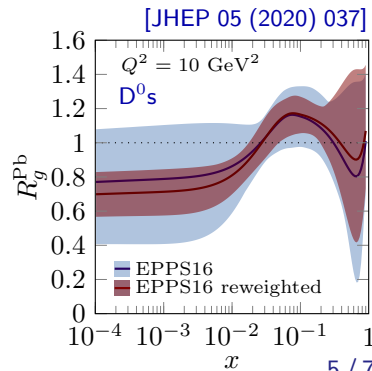
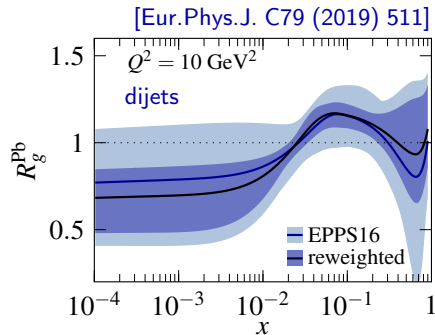
Progress in EPPS analysis

We have performed Hessian PDF reweighting studies to see the impact of dijets and D^0 s

- Large reduction of EPPS16 gluon uncertainties
- Support for mid- x antishadowing and rather deep small- x gluon shadowing
- D-meson R_{pPb} sensitive to nPDFs down to $x \sim 10^{-5}$
- Constraints from dijet and D-meson data mutually consistent!

Work in progress: Include these and the 8.16 TeV CMS W bosons into a global analysis

- Studies in more relaxed parametrization ongoing
- “Final call” for new observables

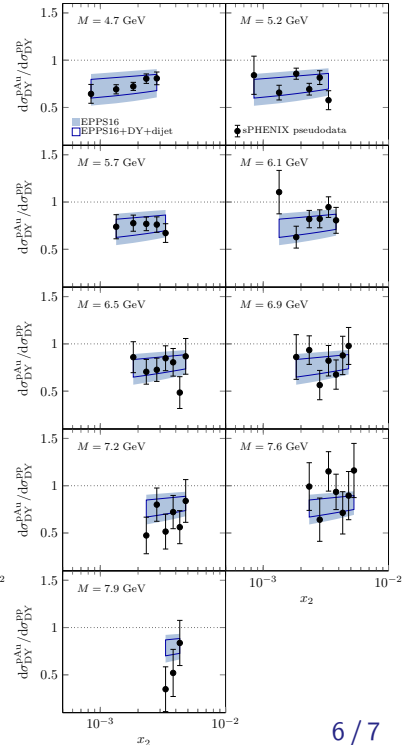
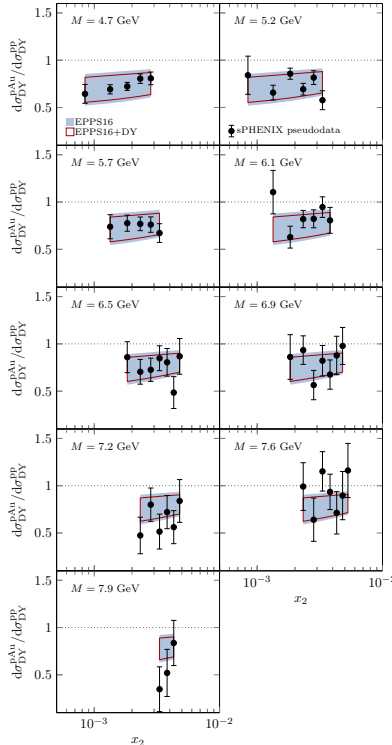


Correlations are important!

[Phys.Rev.D 100 (2019) 014004]

We have studied the prospects of using future sPHENIX Drell–Yan data for nPDF constraints

- If we include the Drell–Yan data alone, a large luminosity uncertainty prevents getting any meaningful constraints
- Since this uncertainty is correlated across observables, we can use dijet measurement to fix this problem, but we need to know the correlations to do so!



- For $k \in \{1, \dots, N\}$ separate data sets, we define the global χ^2 figure of merit as

$$\chi_{\text{global}}^2 = \sum_k \chi_k^2$$

where (simplifying)

$$\chi_k^2 = [D_k - T_k]^T C_k^{-1} [D_k - T_k]$$

..... please publish your correlations!
↓

- Now, if a combined data set $(D_{\text{comb}}, C_{\text{comb}})$ contains the same information as separate data sets (D_1, C_1) and (D_2, C_2) , i.e.

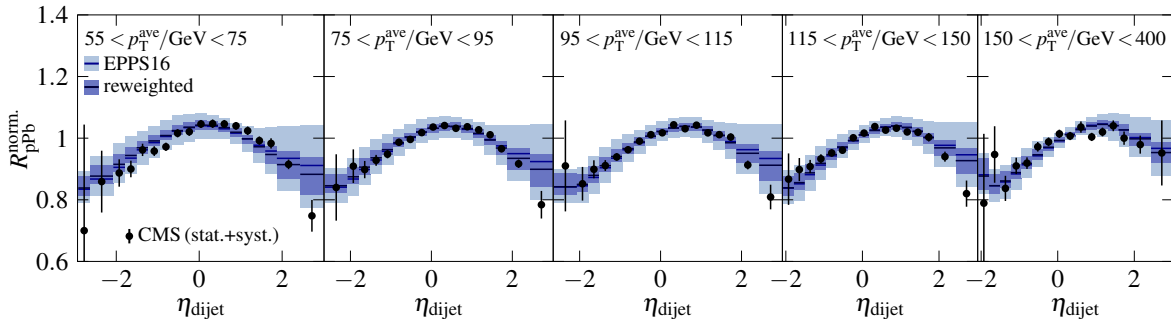
$$\chi_{\text{comb}}^2 = \chi_1^2 + \chi_2^2$$

then it does not matter if we include the separate or combined data in our analysis

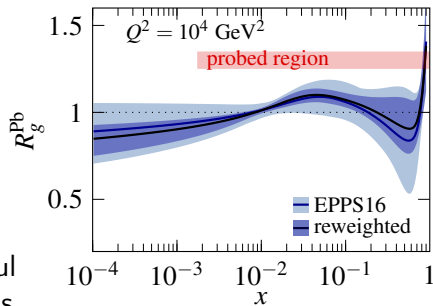
- For a data combination to be useful to us, some uncertainty reduction **beyond** a simple quadratic sum would be needed (requires understanding the inter-detector correlations)
 - ▶ At HERA, they were able to “cross calibrate” the detectors [JHEP 01 (2010) 109]

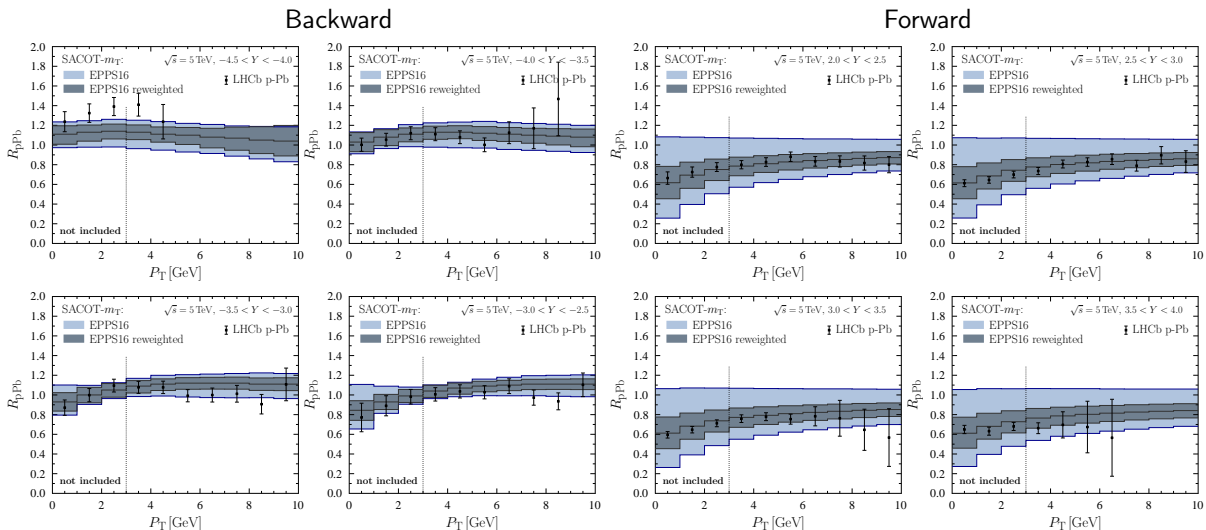
Whether this can be done at the LHC is an experimental problem (**I am all ears!**)

Backup



- Drastic reduction in EPPS16 uncertainties!
- Downward pull in the forward region
- The most forward data points lie systematically below the reweighted uncertainty band – could be due to
 - ▶ inflexibility in EPPS16 parametrization at small x
 - ▶ systematics of the measurement – would be helpful to have correlations of uncertainties available to us





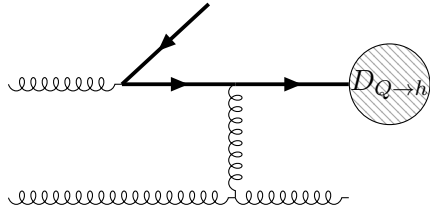
- Data well reproduced with the reweighted results
- Significant reduction in EPPS16 uncertainties especially in forward bins
- Good agreement with data below cut – no physics beyond collinear factorization needed

Heavy-flavour production mass schemes

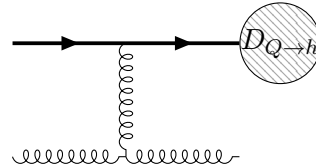
FFNS

In *fixed flavour number scheme*, valid at small p_T , heavy quarks are produced only at the matrix element level

Contains $\log(p_T/m)$ and m/p_T terms



– subtraction term +



ZM-VFNS

In *zero-mass variable flavour number scheme*, valid at large p_T , heavy quarks are treated as massless particles produced also in ISR/FSR

Resums $\log(p_T/m)$ but ignores m/p_T terms

GM-VFNS

A *general-mass variable flavour number scheme* combines the two by supplementing subtraction terms to prevent double counting of the resummed splittings, valid at all p_T

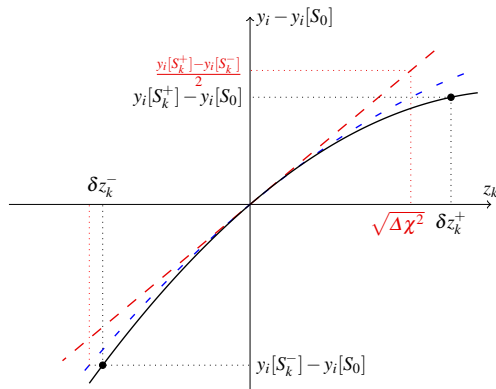
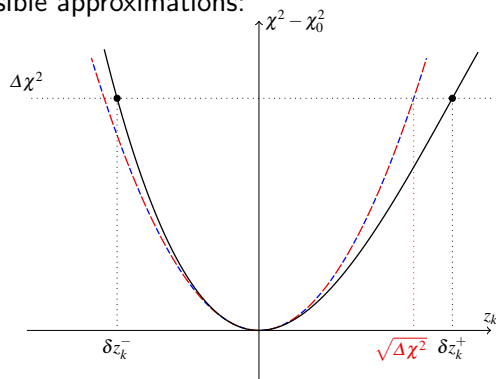
Resums $\log(p_T/m)$ and includes m/p_T terms in the FFNS matrix elements

Important: includes also **gluon-to-HF fragmentation** – large contribution to the cross section!

The Hessian reweighting is a method to study the impact of a new set of data on the PDFs without performing a full global fit

$$\chi_{\text{new}}^2(\mathbf{z}) = \chi_{\text{old}}^2(\mathbf{z}) + \sum_{ij} (y_i(\mathbf{z}) - y_i^{\text{data}}) C_{ij}^{-1} (y_j(\mathbf{z}) - y_j^{\text{data}})$$

Possible approximations:



quadratic-linear: $\chi_{\text{old}}^2 \approx \chi_0^2 + \sum_k z_k^2$,

quadratic-quadratic: $\chi_{\text{old}}^2 \approx \chi_0^2 + \sum_k z_k^2$,

cubic-quadratic: $\chi_{\text{old}}^2 \approx \chi_0^2 + \sum_k (a_k z_k^2 + b_k z_k^3)$,

$y_i \approx y_i[S_0] + \sum_k d_{ik} z_k$

$y_i \approx y_i[S_0] + \sum_k (d_{ik} z_k + e_{ik} z_k^2)$

$y_i \approx y_i[S_0] + \sum_k (d_{ik} z_k + e_{ik} z_k^2)$