

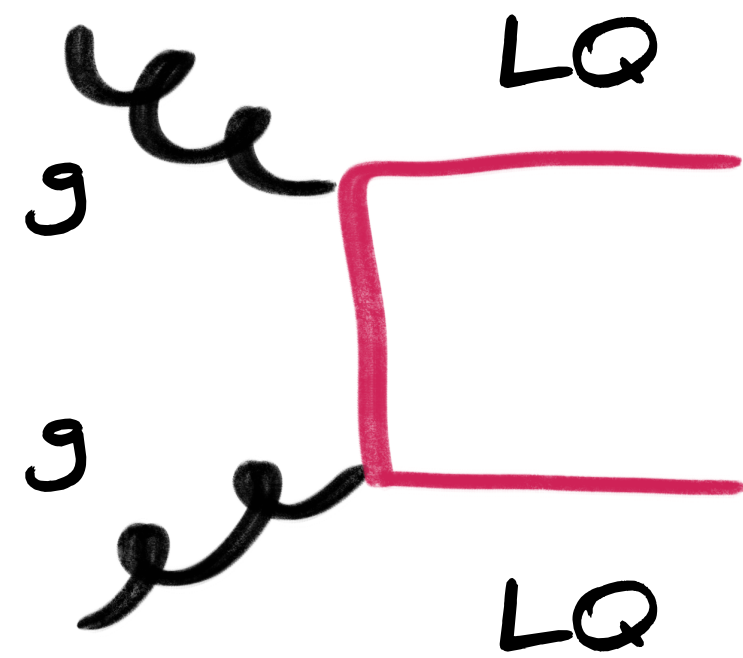
Resonant leptoquark production at hadron colliders

SM/BSM seminar, CERN, 12. June 2020

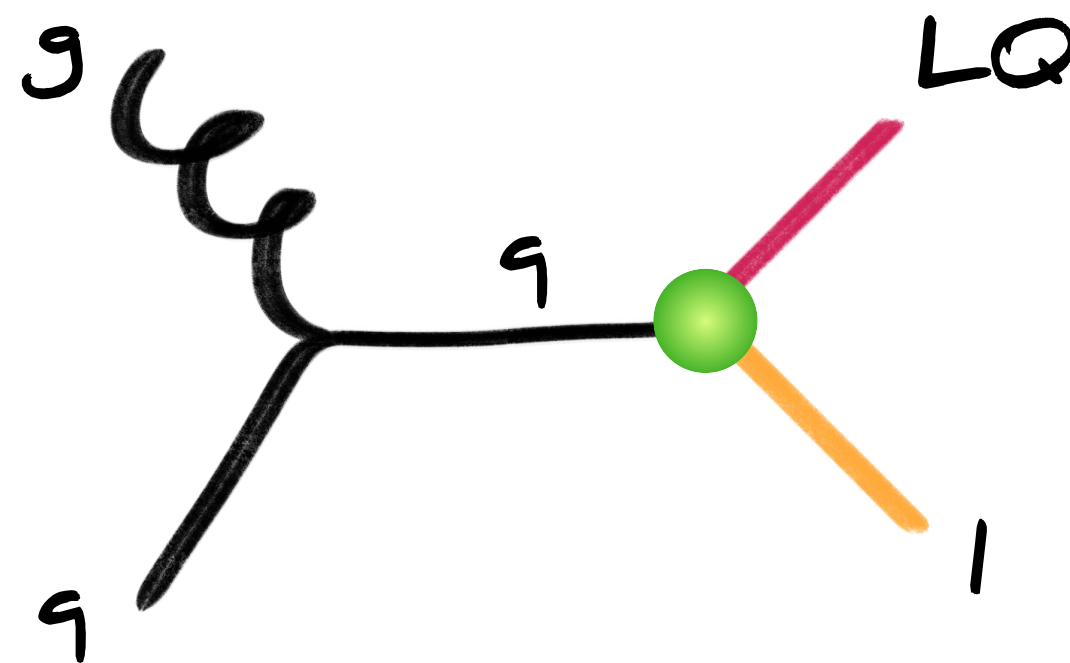
Luca Buonocore, Ulrich Haisch, Paolo Nason, Francesco Tramontano & Giulia Zanderighi, 2005.06475

Leptoquarks (LQs) in a nutshell

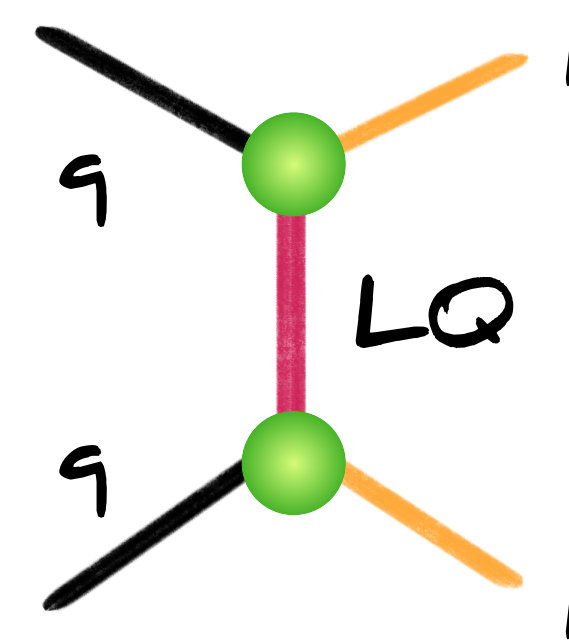
$$\mathcal{L} \supset |D_\mu LQ|^2 + \lambda_{ql} LQ (ql)$$



pair production (PP)



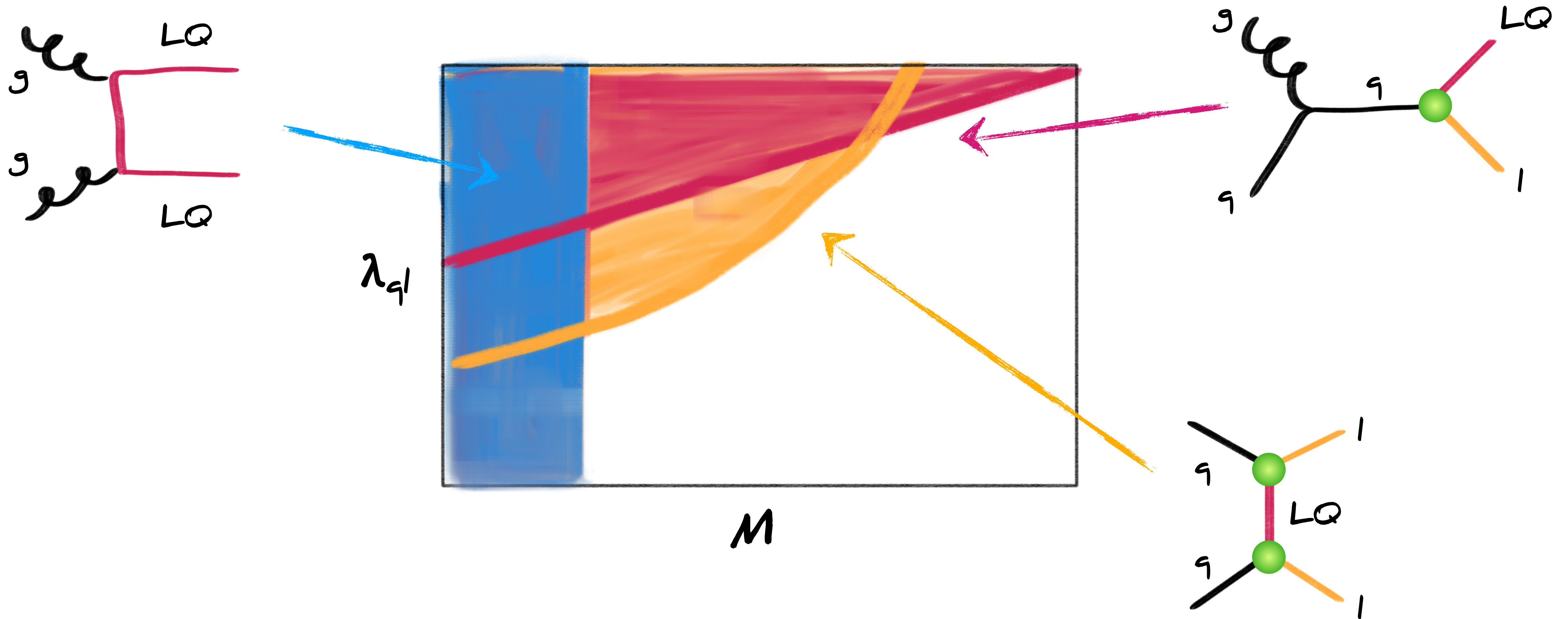
single production (SP)



Drell-Yan (DY) t-channel

for a recent LQ review see Dorsner et al., 1603.04993

Complementarity of LQ searches



sketch adopted from Dorsner & Greljo, 1801.07641

May 2020 summary of LHC limits

LQ	Scalar LQ 1 st gen	LQ mass	1.4 TeV	1902.00377
	Scalar LQ 2 nd gen	LQ mass	1.56 TeV	1902.00377
	Scalar LQ 3 rd gen	LQ_3^u mass	1.03 TeV	1902.08103
	Scalar LQ 3 rd gen	LQ_3^d mass	970 GeV	1902.08103

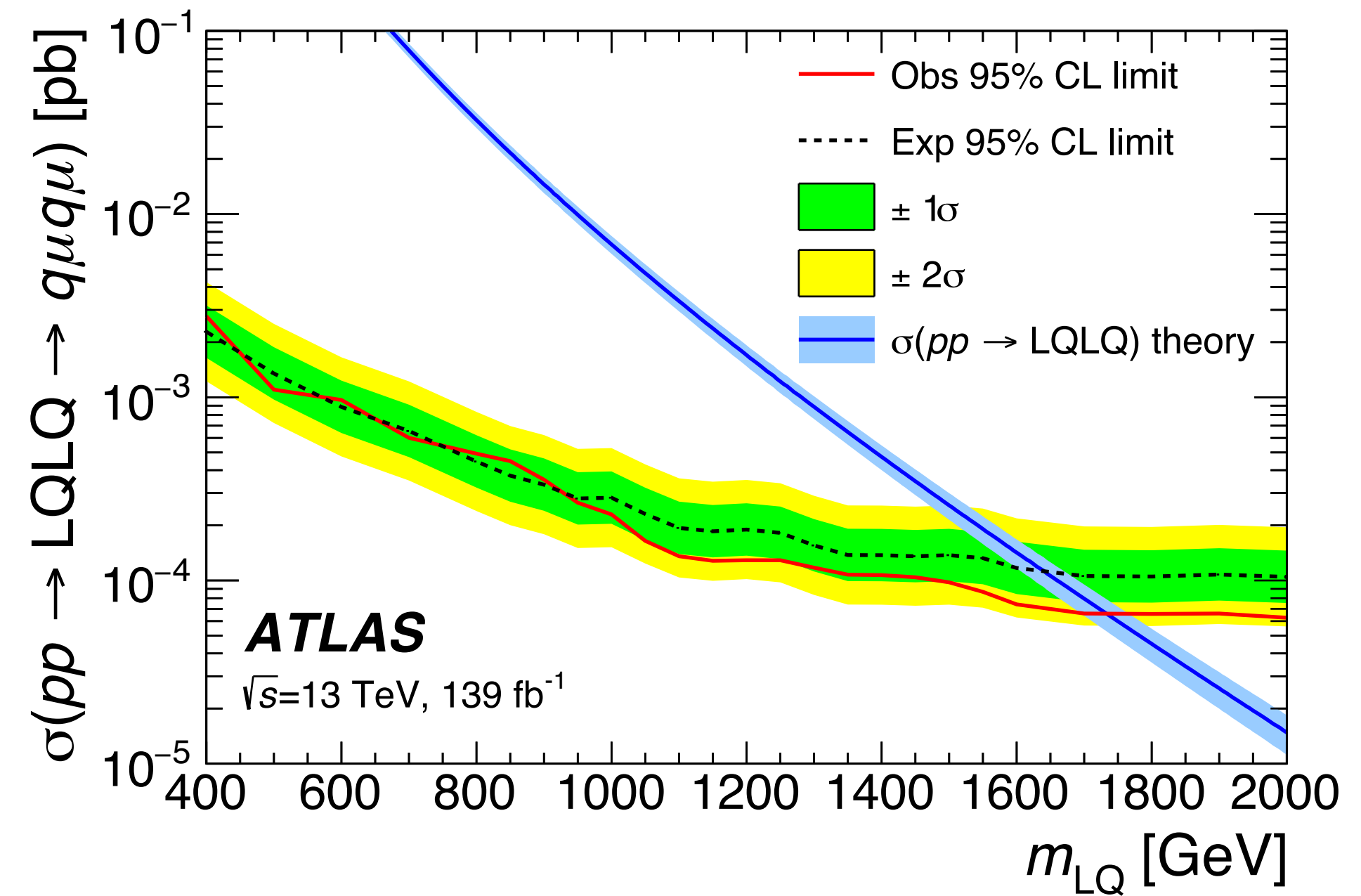
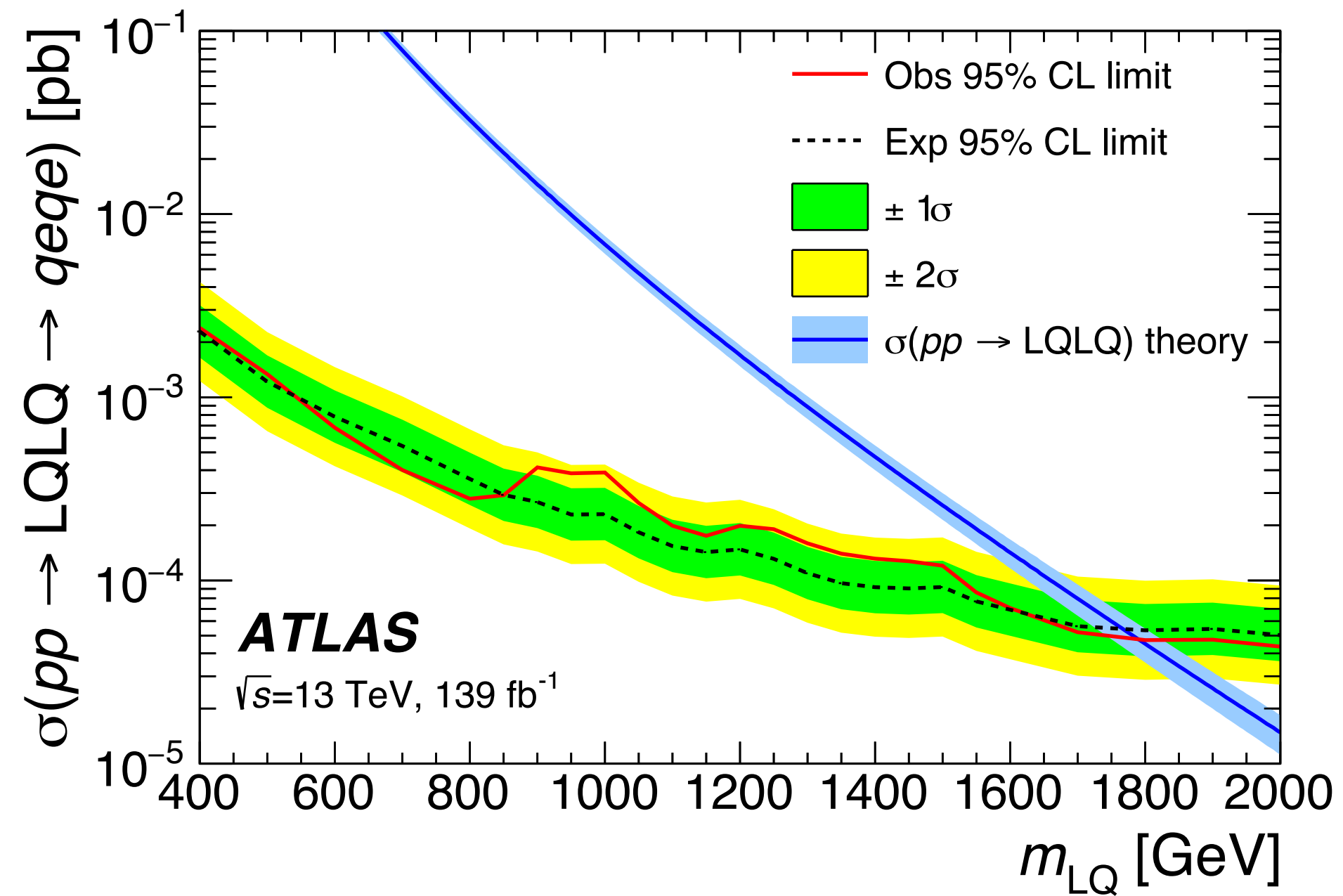
ATLAS Exotics Searches

Leptoquarks	scalar LQ (pair prod.), coupling to 1 st gen. fermions, $\beta = 1$	<1.44	1811.01197 (2e + 2j)	36 fb ⁻¹
	scalar LQ (pair prod.), coupling to 1 st gen. fermions, $\beta = 0.5$	<1.27	1811.01197 (2e + 2j; e + 2j + E _T ^{miss})	36 fb ⁻¹
	scalar LQ (pair prod.), coupling to 2 nd gen. fermions, $\beta = 1$	<1.53	1808.05082 (2μ + 2j)	36 fb ⁻¹
	scalar LQ (pair prod.), coupling to 2 nd gen. fermions, $\beta = 1$	0.8–1.5	1811.10151 (1μ + 1j + E _T ^{miss})	77 fb ⁻¹
	scalar LQ (pair prod.), coupling to 2 nd gen. fermions, $\beta = 0.5$	<1.29	1808.05082 (2μ + 2j; μ + 2j + E _T ^{miss})	36 fb ⁻¹
	scalar LQ (pair prod.), coupling to 3 rd gen. fermions, $\beta = 1$	<1.02	1811.00806 (2τ + 2j)	36 fb ⁻¹
	scalar LQ (single prod.), coup. to 3 rd gen. ferm., $\beta = 1, \lambda = 1$	<0.74	1806.03472 (2τ + b)	36 fb ⁻¹

CMS EXO results

Most existing searches have targeted LQ PP. Scalar LQ masses below around 1.5 TeV (1.0 TeV) are excluded for 1st & 2nd (3rd) generation LQs

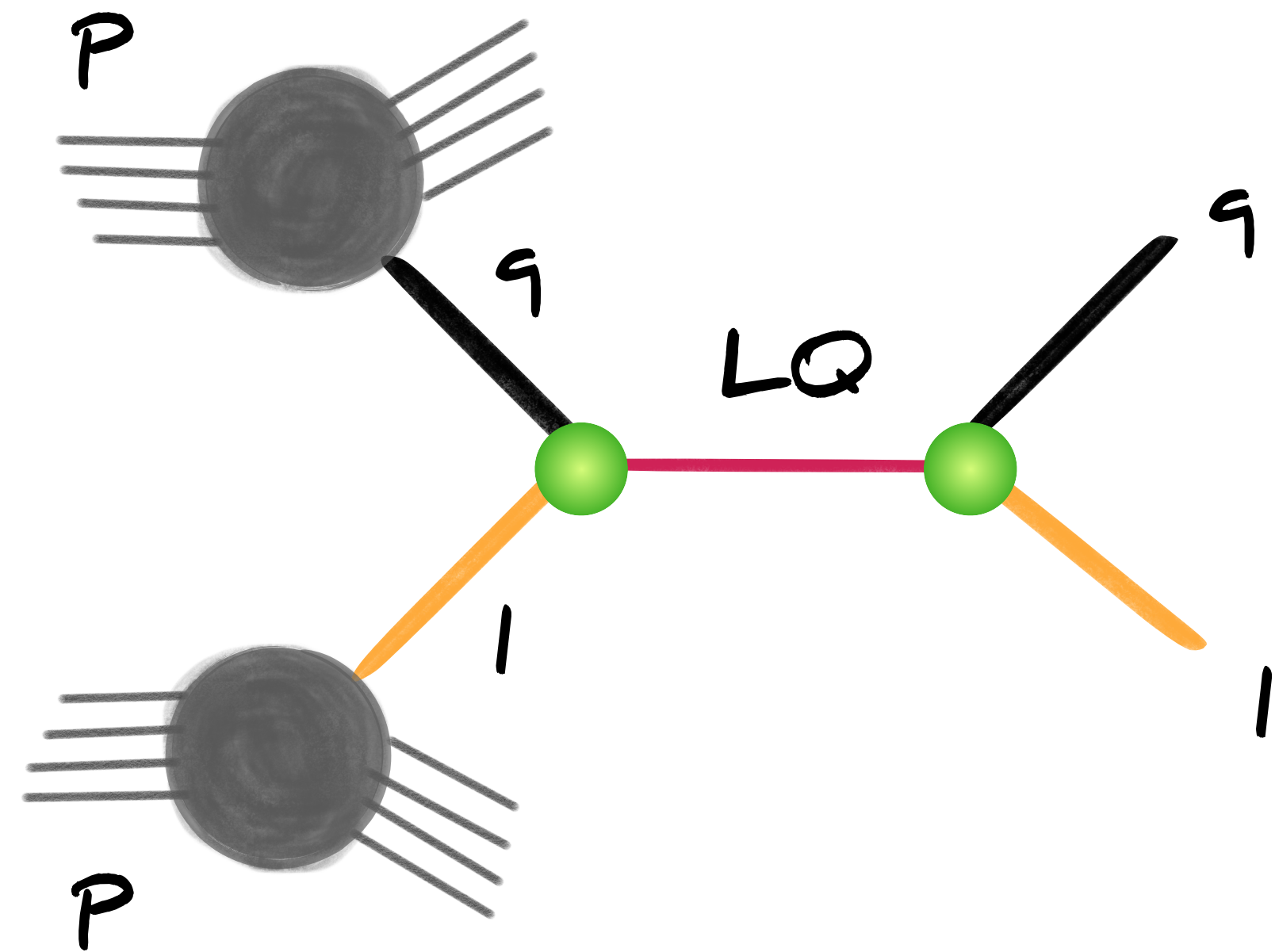
First full LHC Run II results of ATLAS



LQs masses below 1.8 TeV (1.7 TeV) are excluded in the electron (muon) channel, assuming a branching ratio into a charged lepton & a quark of 100%

Main idea of 2005.06475

But as Luca explained before, "a proton is a quark is a photon is a lepton", so LQs can also be produced resonantly in proton-proton collisions. Are the corresponding rates large enough to be measurable at the LHC?



see also Ohnemus et al., hep-ph/9406235

Analysis strategy

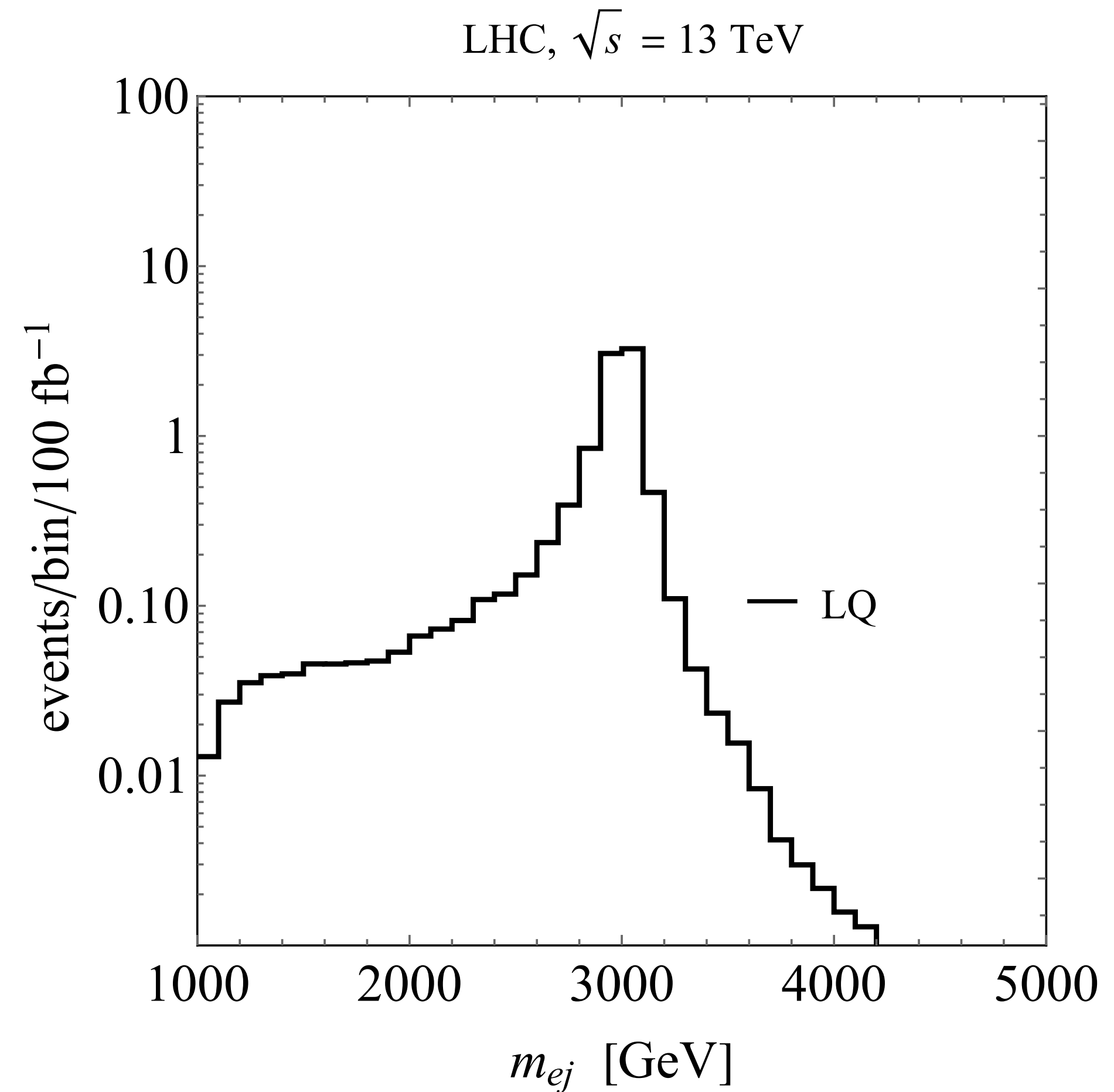
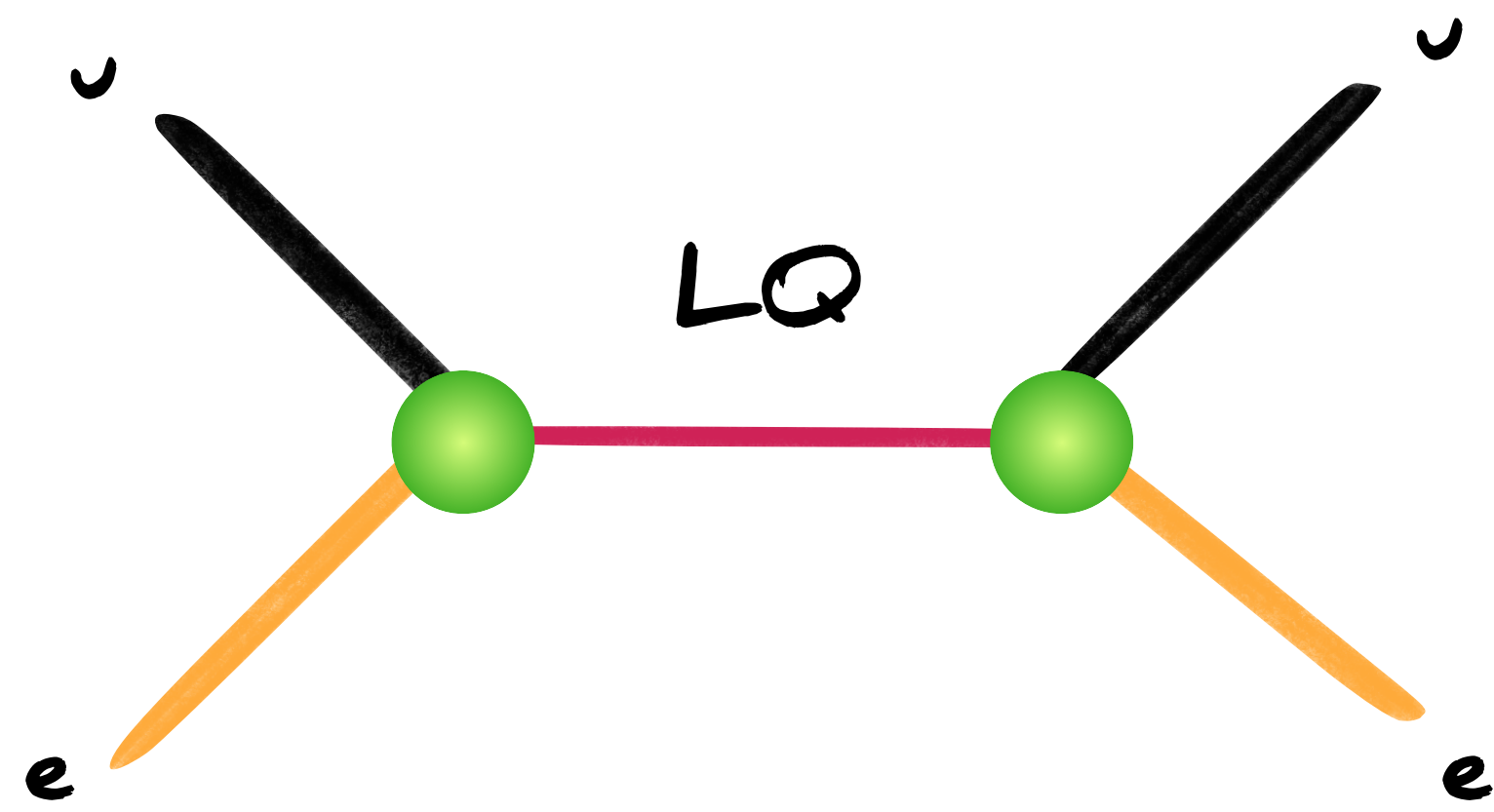
- Target final states with a high transverse momentum (p_T) lepton & a high- p_T jet:

$$|\eta_{l_1}| < 2.5, \quad p_{T,l_1} > 500 \text{ GeV}, \quad |\eta_{j_1}| < 2.5, \quad p_{T,j_1} > 500 \text{ GeV}$$

- To suppress backgrounds, cut on missing energy ($E_{T,\text{miss}}$), veto extra leptons & jets:

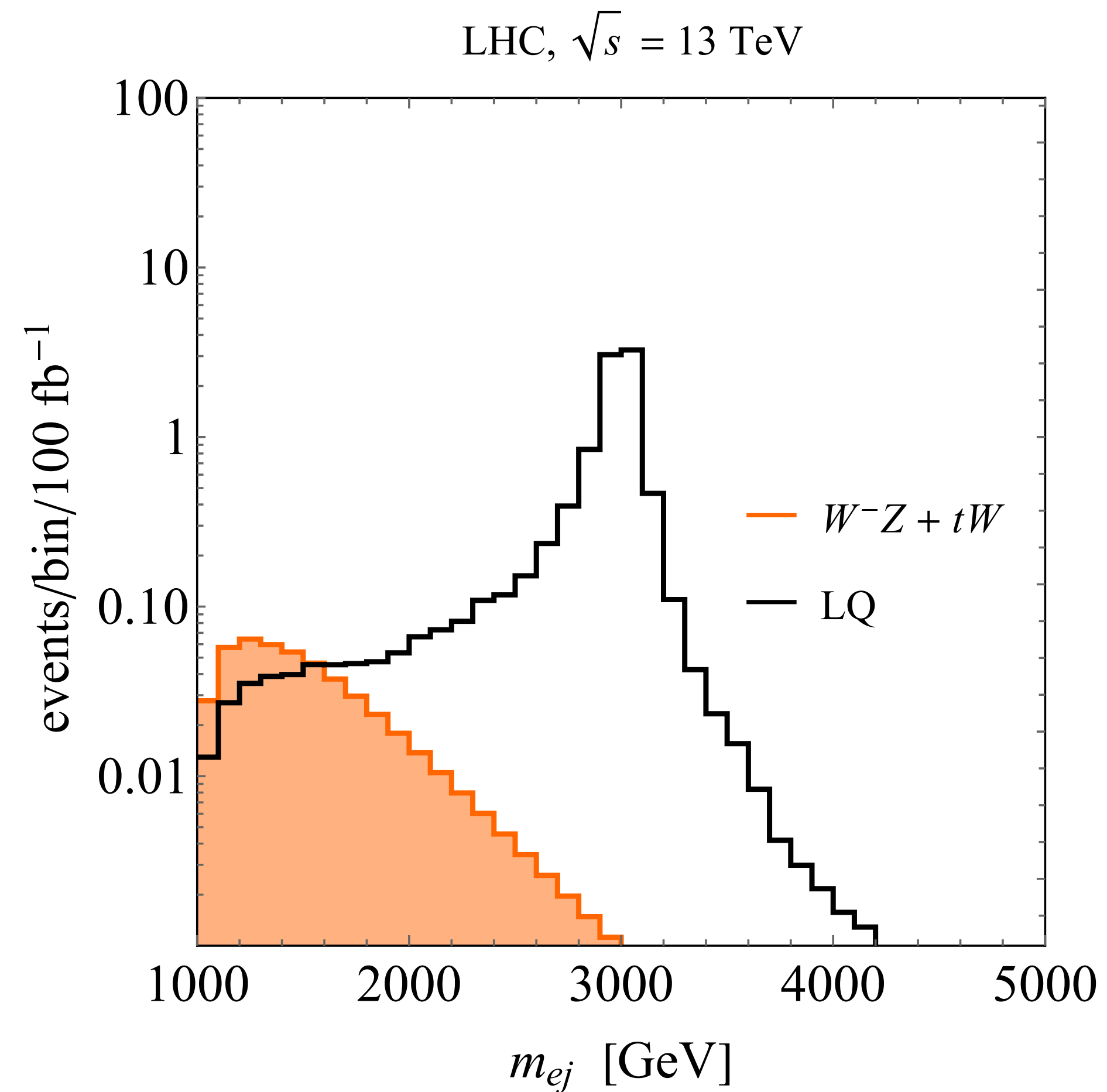
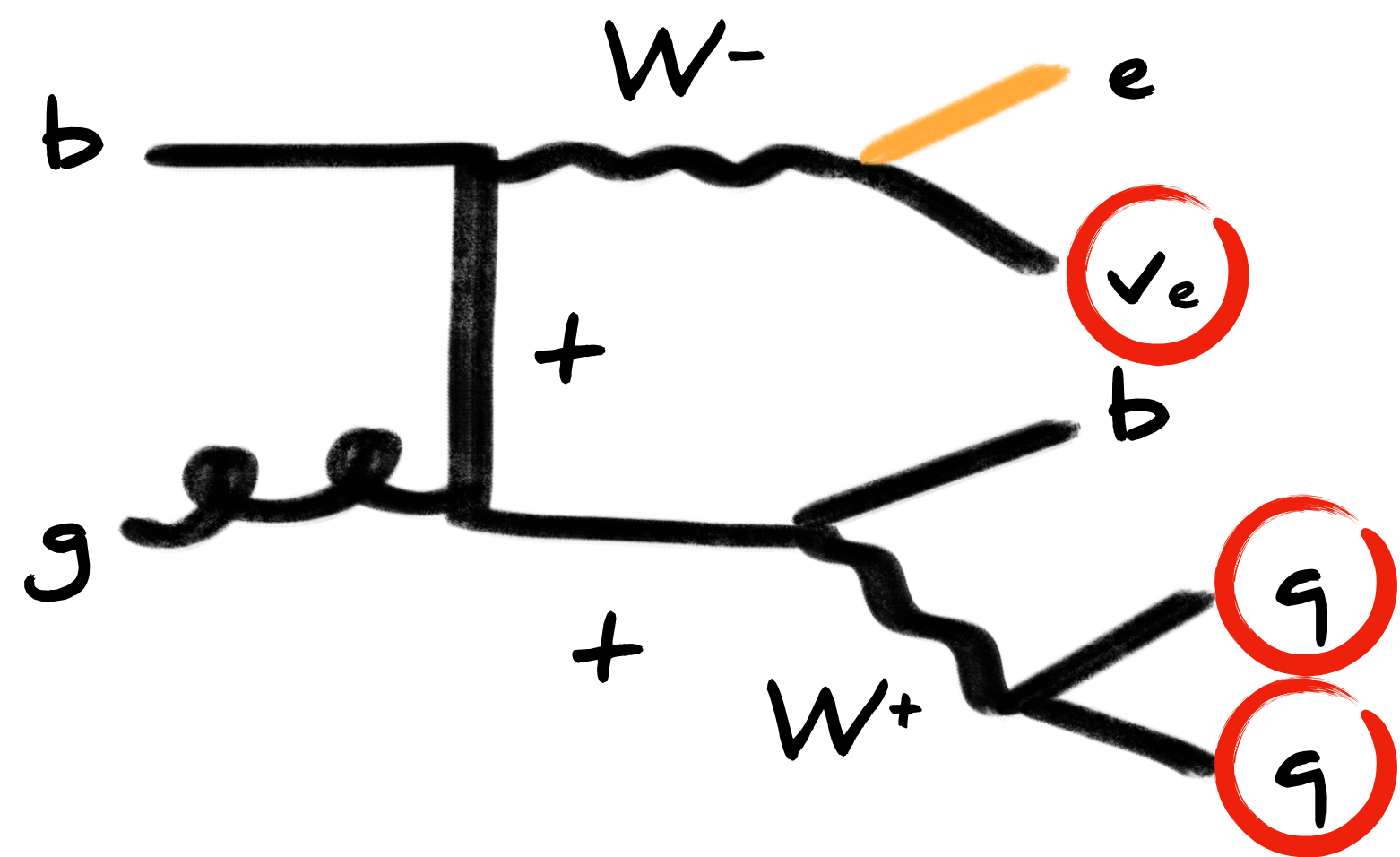
$$E_{T,\text{miss}} < 50 \text{ GeV}, \quad |\eta_{l_2}| < 2.5, \quad p_{T,l_2} > 7 \text{ GeV}, \quad |\eta_{j_2}| < 2.5, \quad p_{T,j_2} > 30 \text{ GeV}$$

Benchmark signal



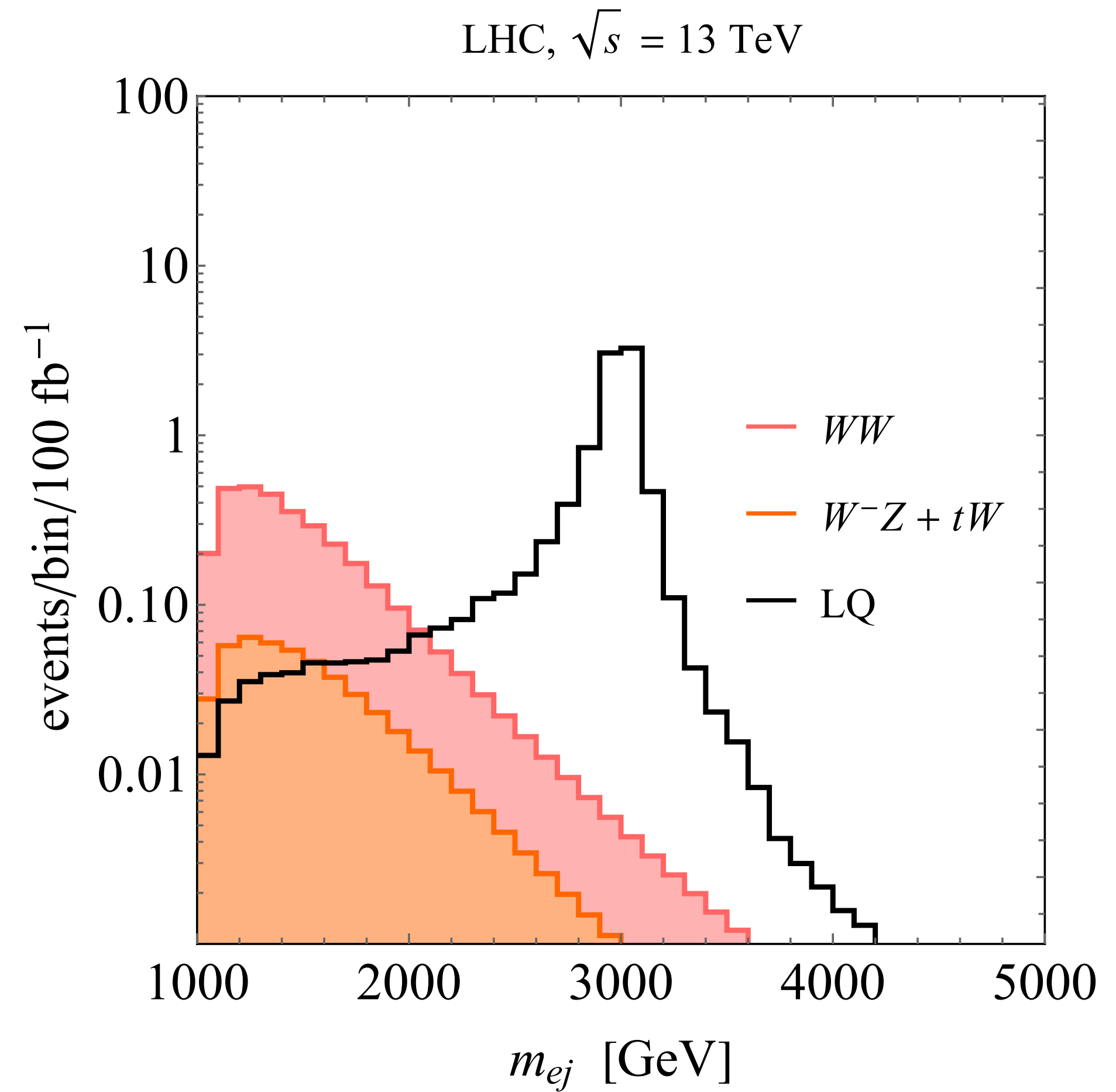
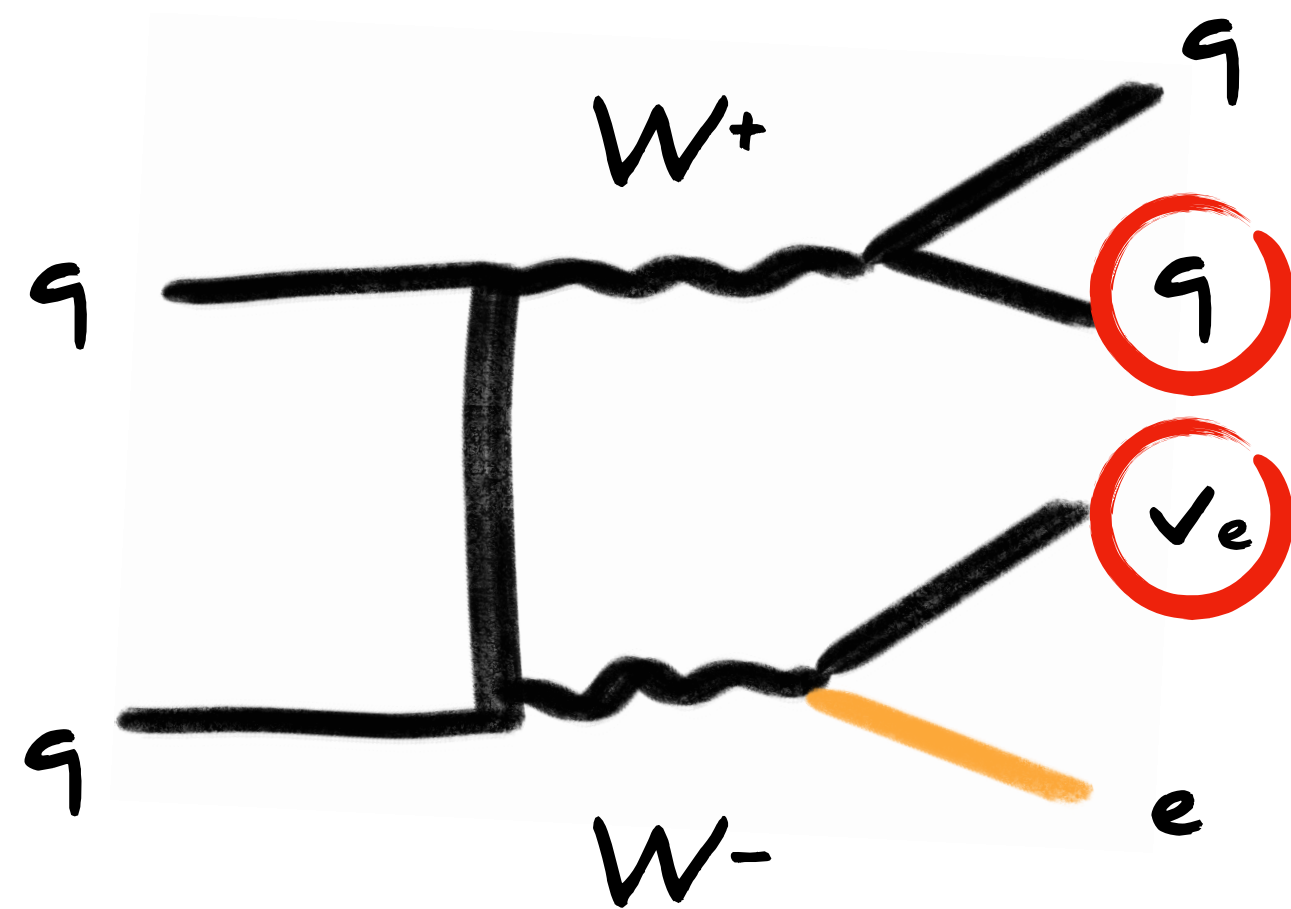
At 13 TeV LHC, 9 events per 100 fb⁻¹ for minimal scalar LQ of $M = 3$ TeV & $g_{eu} = 1$

Backgrounds



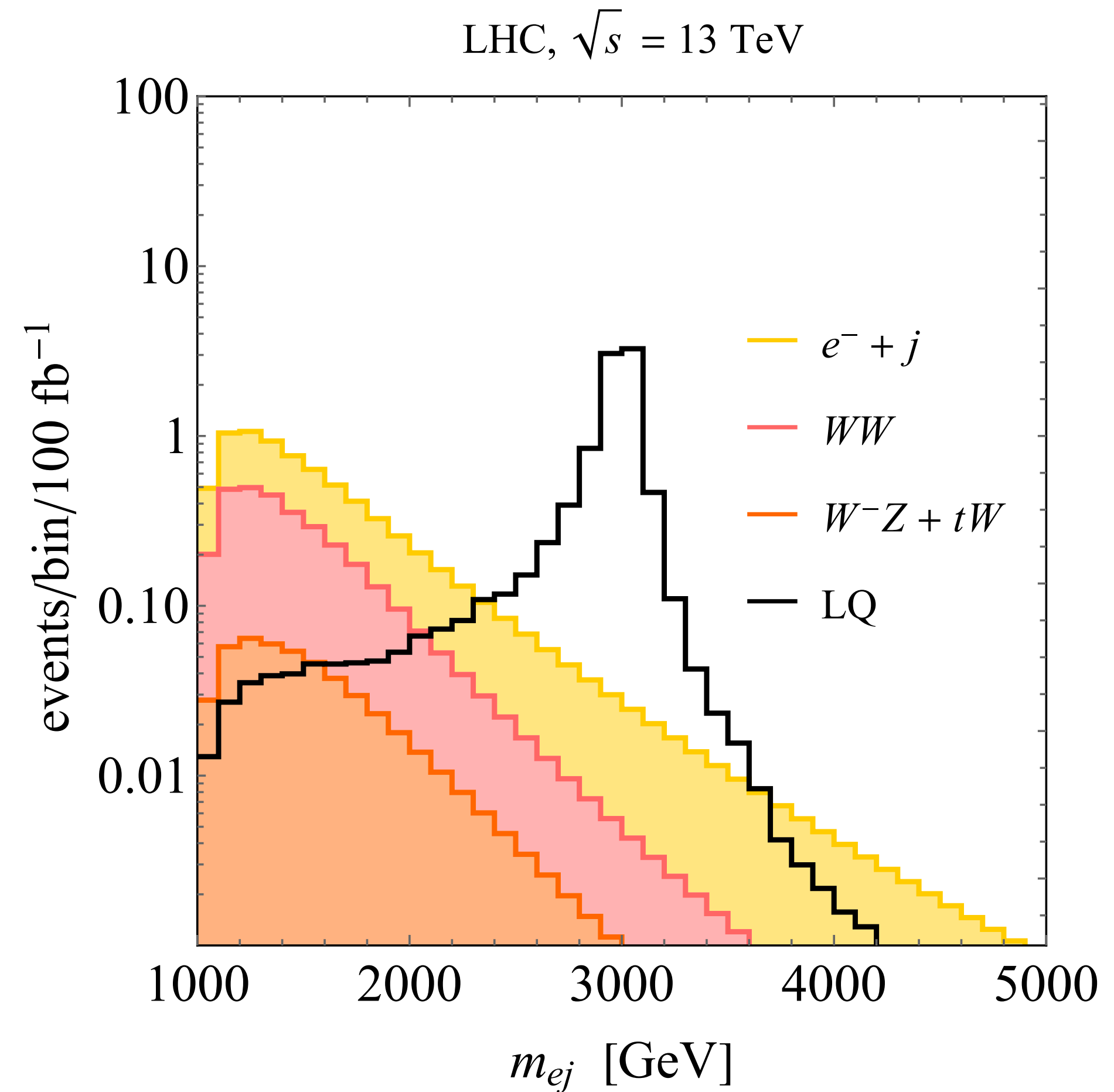
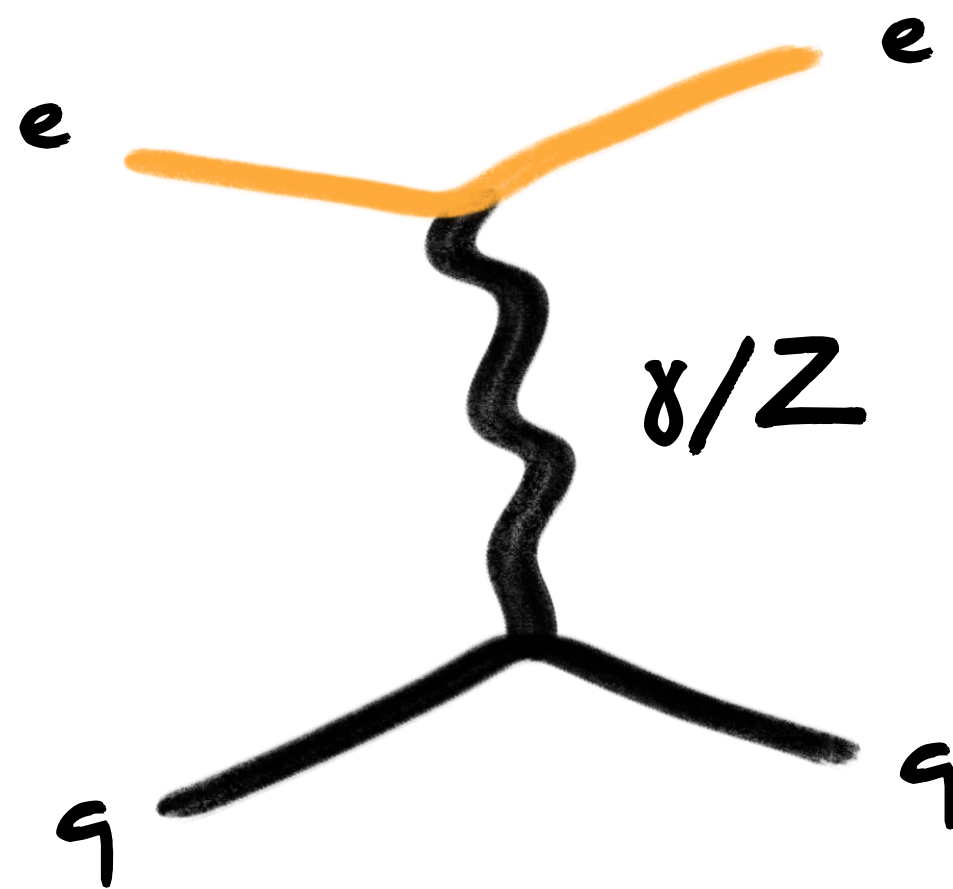
Suppressed by $E_{T,miss}$ requirement & jet veto

Backgrounds



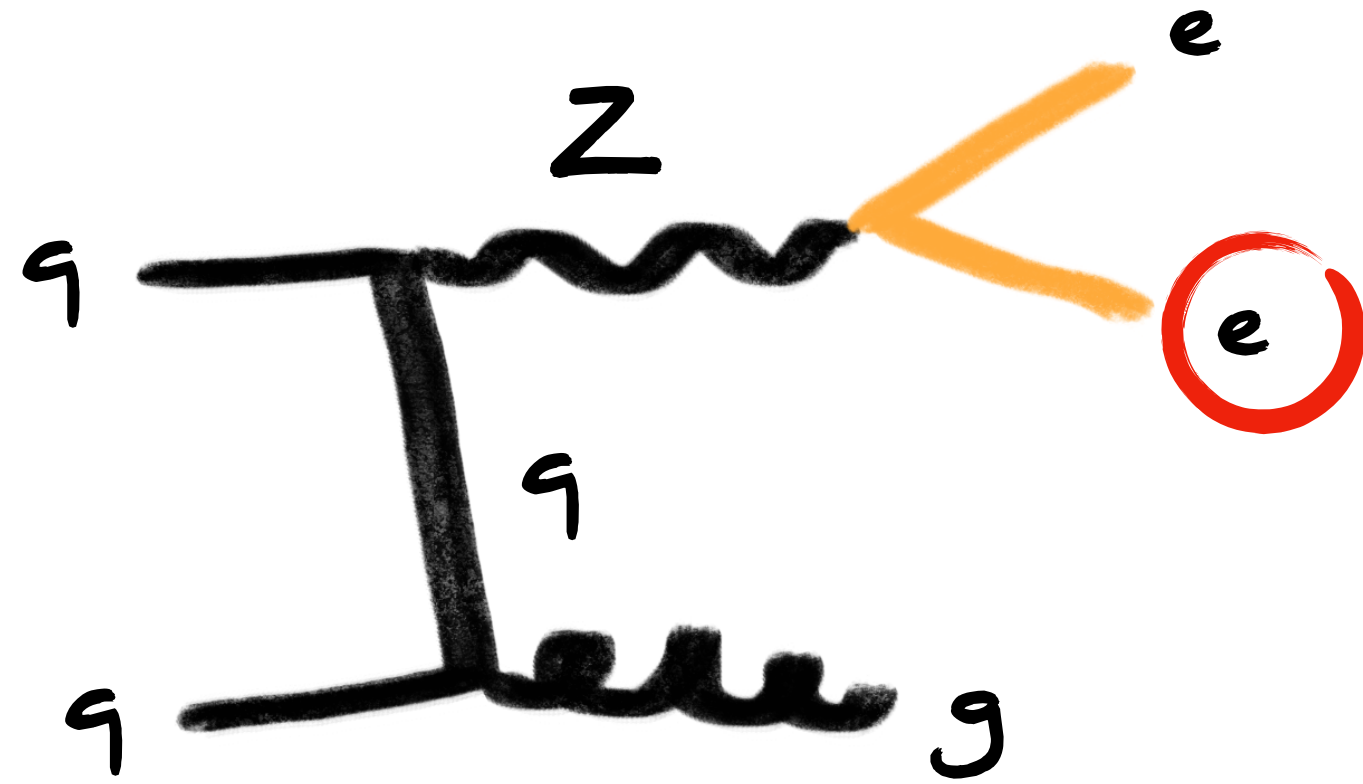
Suppressed by $E_{T,miss}$ requirement & jet veto

Backgrounds

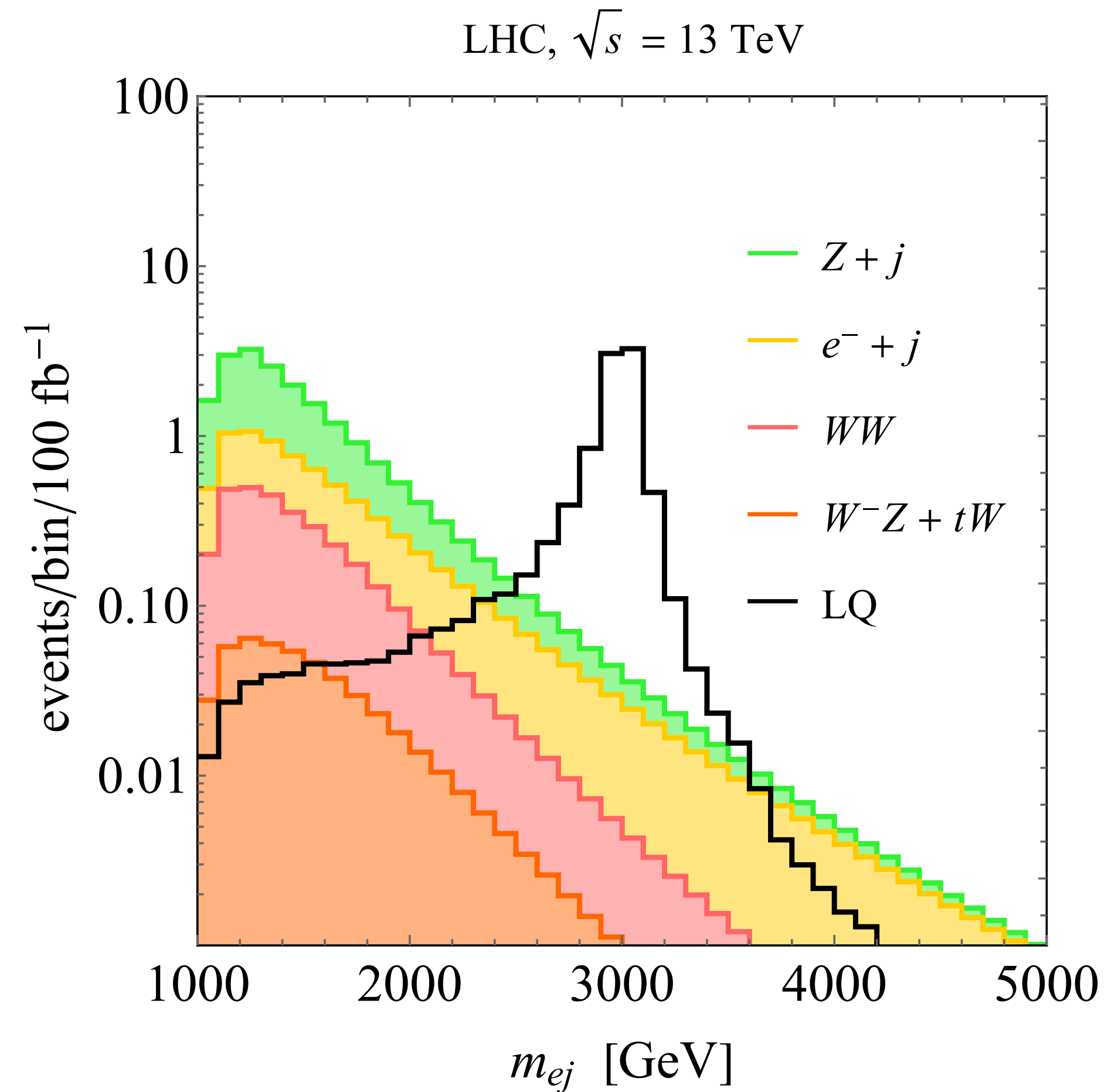


Irreducible background particularly relevant at high invariant lepton-jet mass

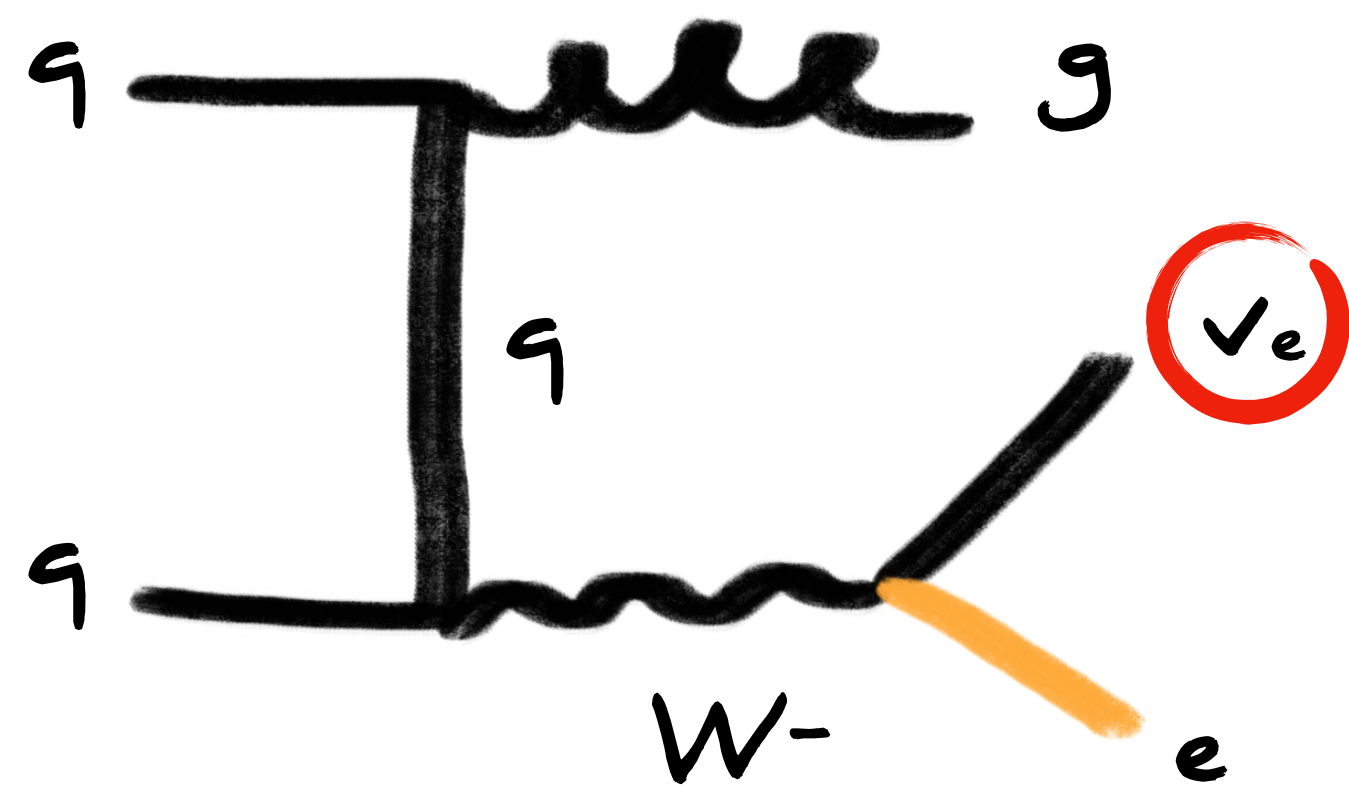
Backgrounds



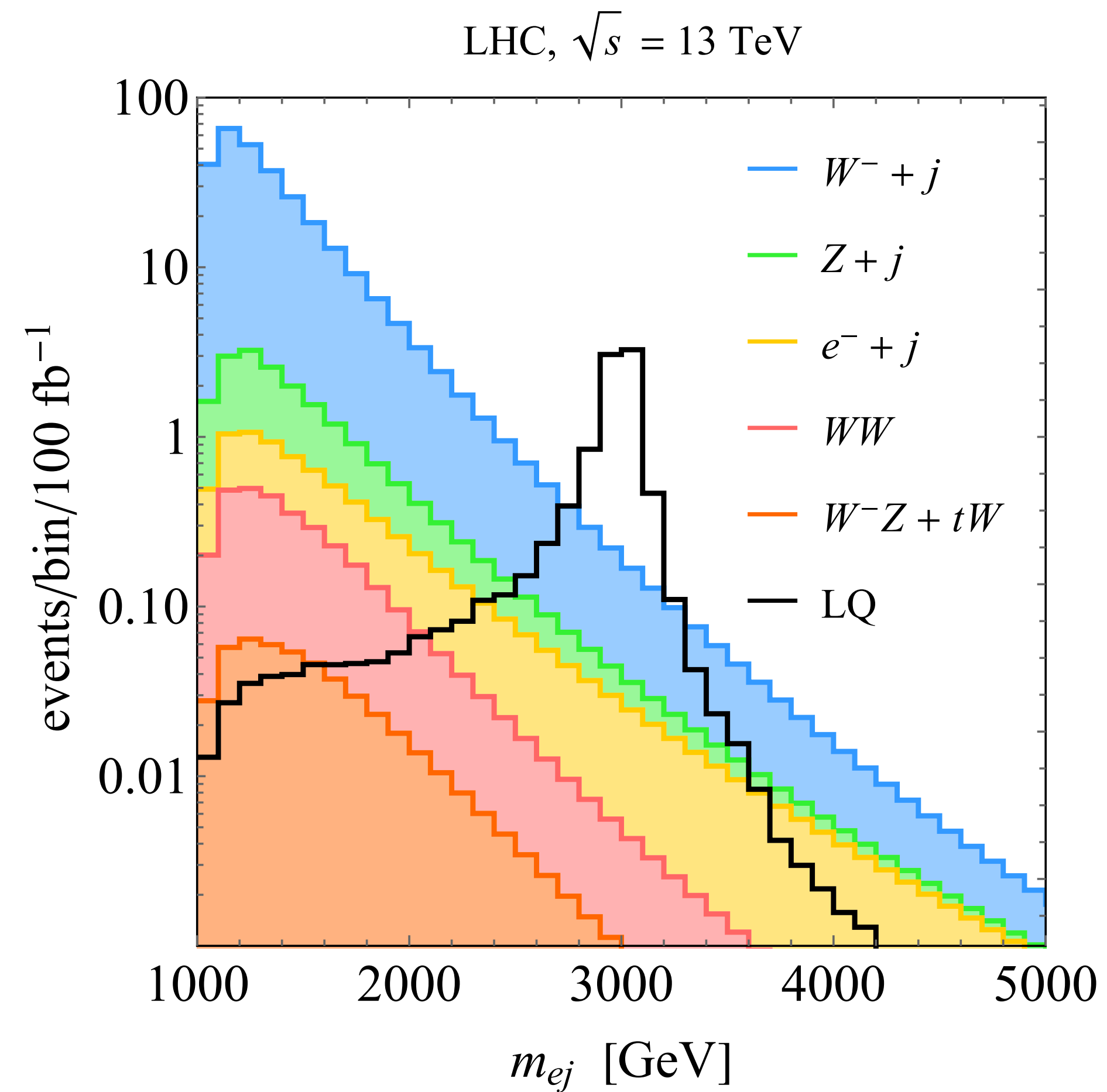
Suppressed by lepton veto



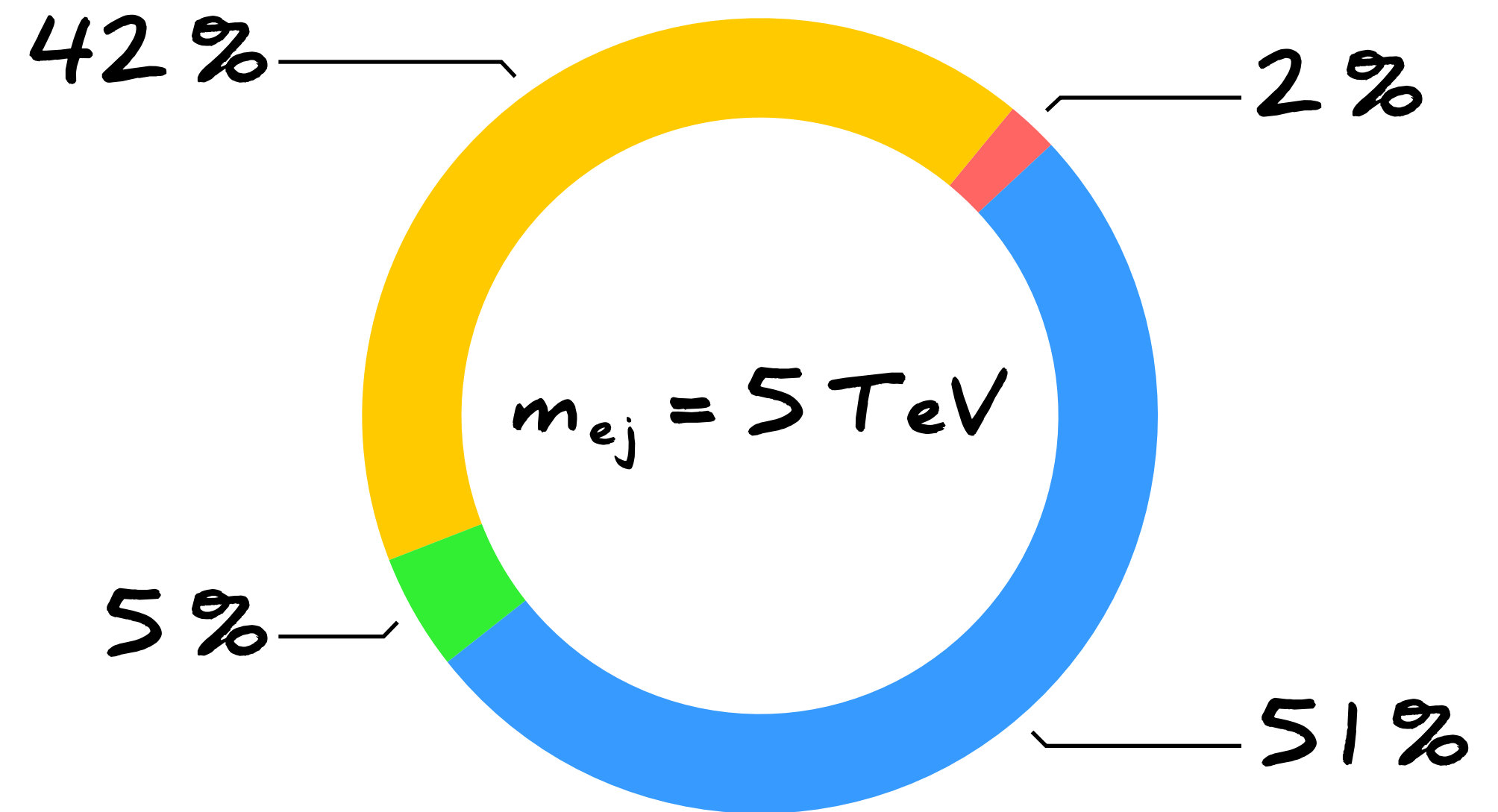
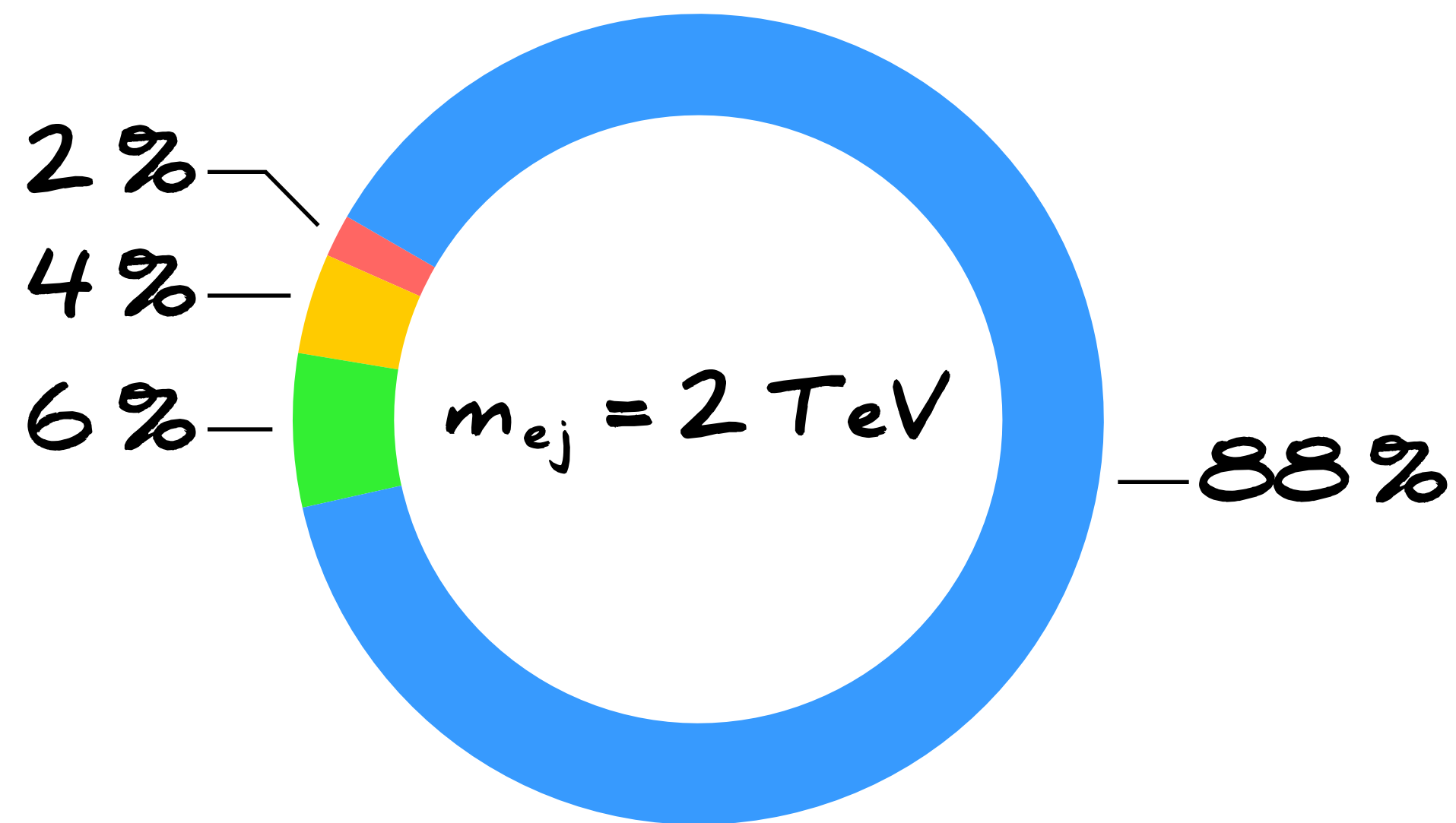
Backgrounds



Suppressed by $E_{T,miss}$ requirement

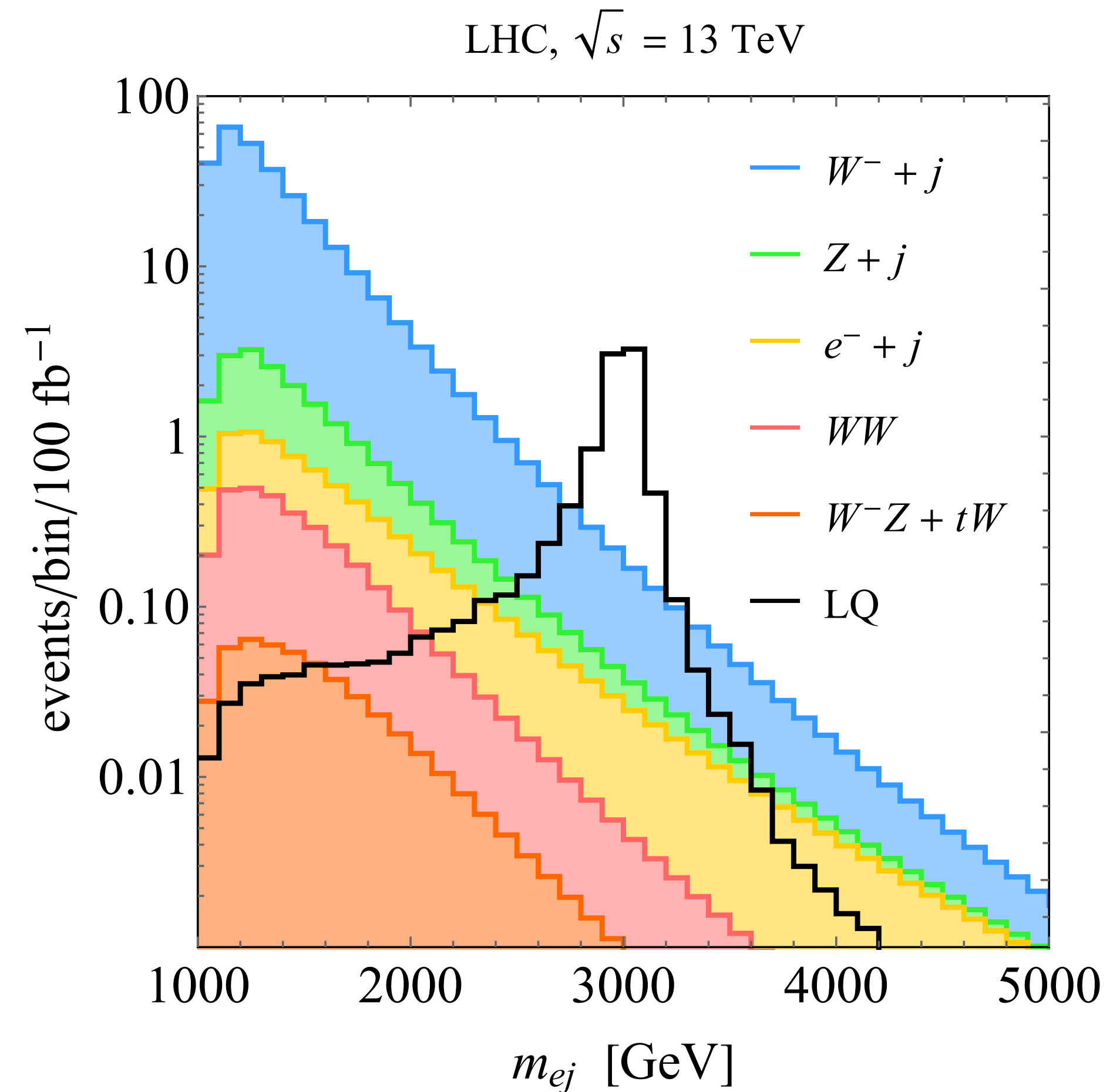
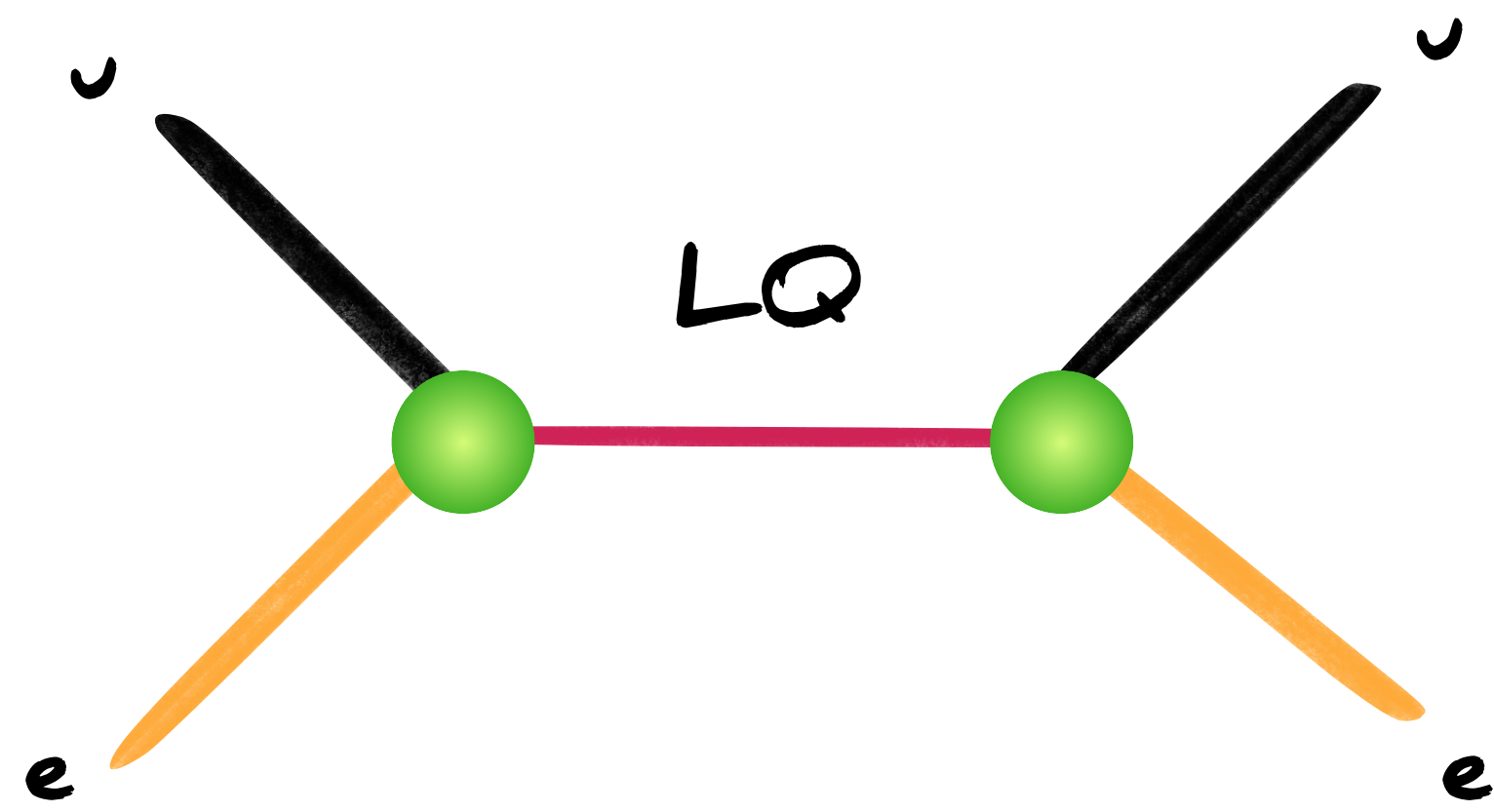


Background decomposition



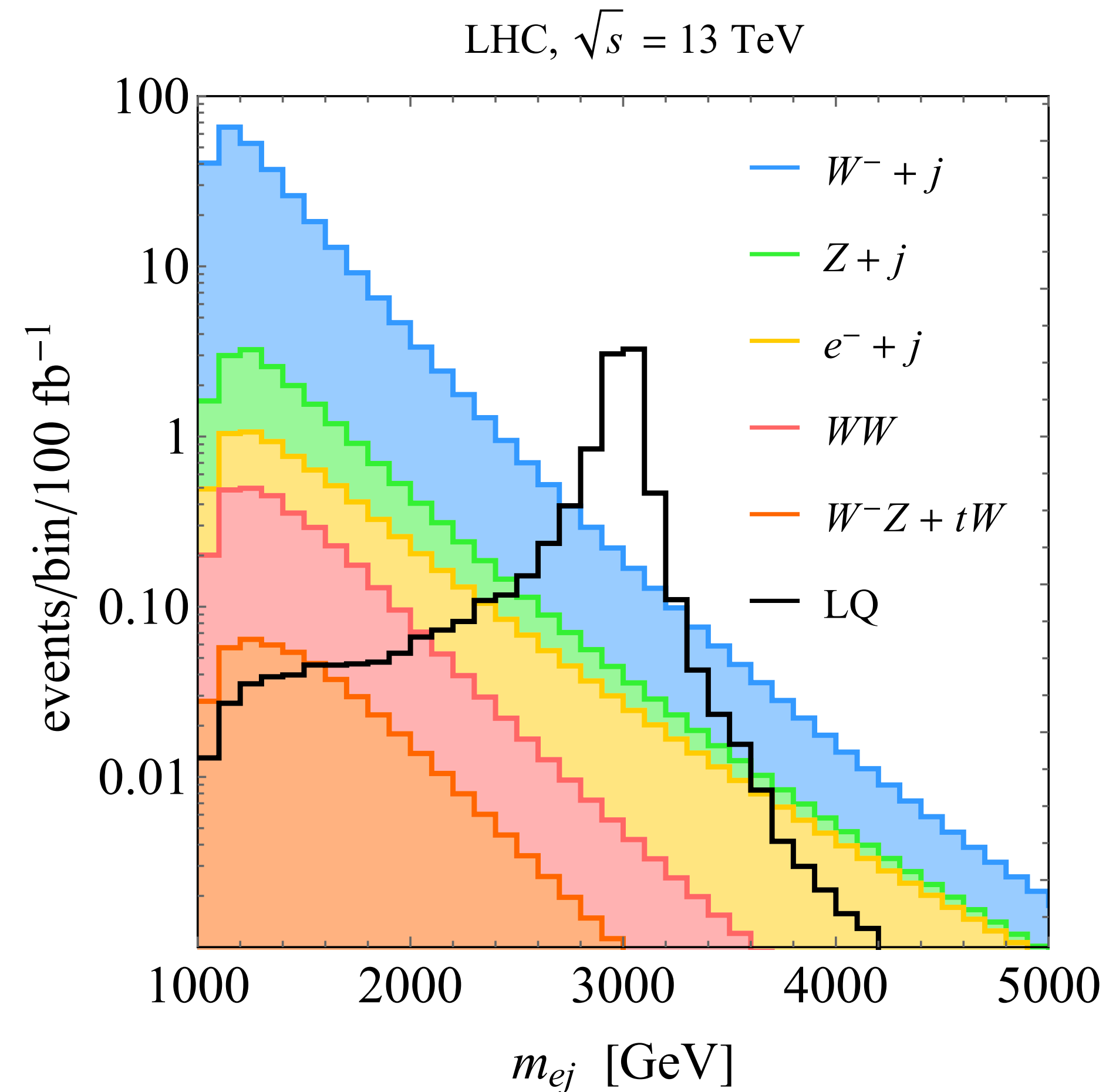
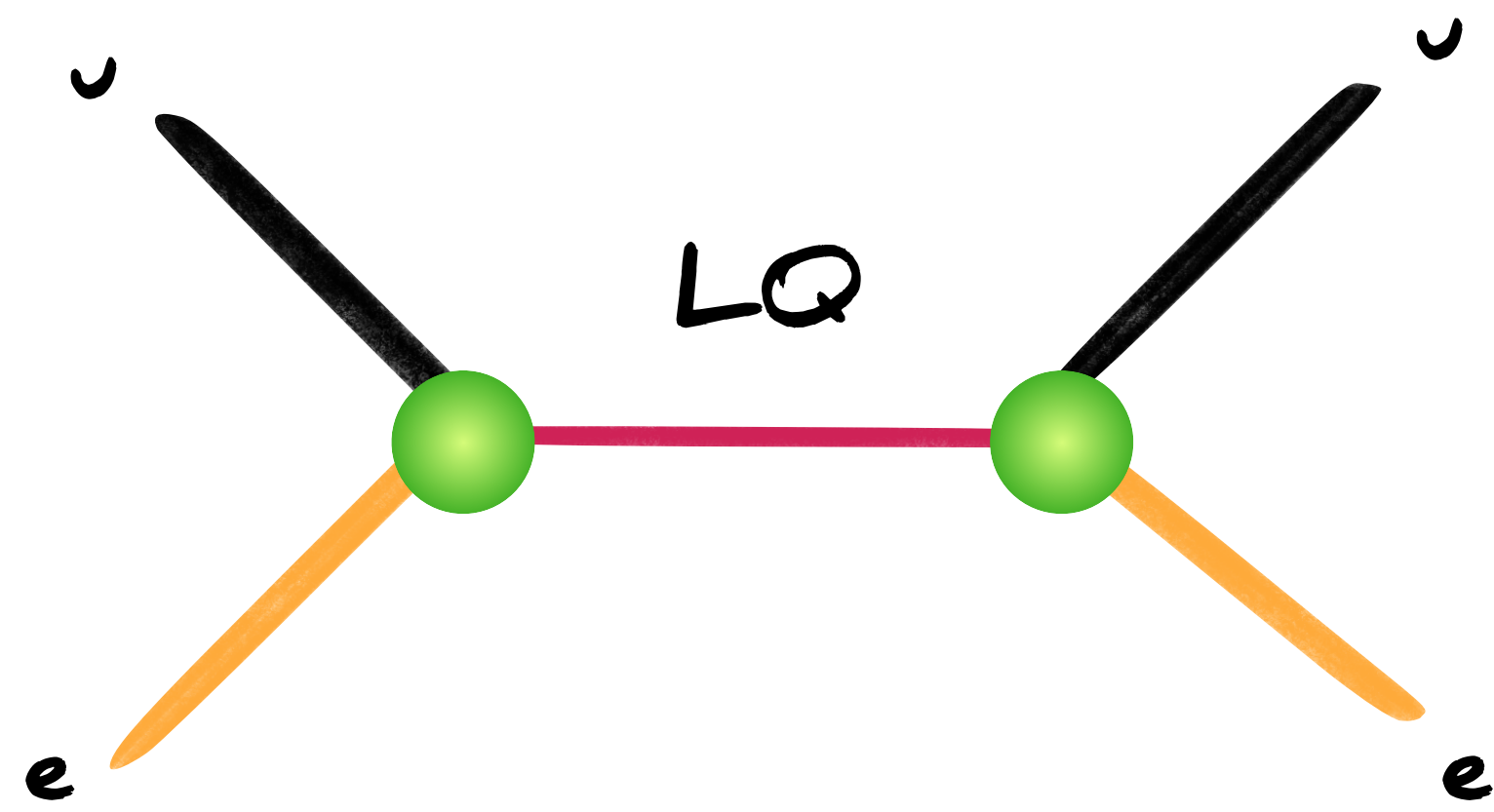
● $W^- + j$ ● $Z + j$ ● $e^- + j$ ● WW

Signal vs. backgrounds



Sum over backgrounds is a steeply falling distribution, while signal exhibits a narrow peak

Signal vs. backgrounds



For 36 fb⁻¹ of 13 TeV LHC data, benchmark disfavoured with a significance of 2σ

Comments on simulation of signal

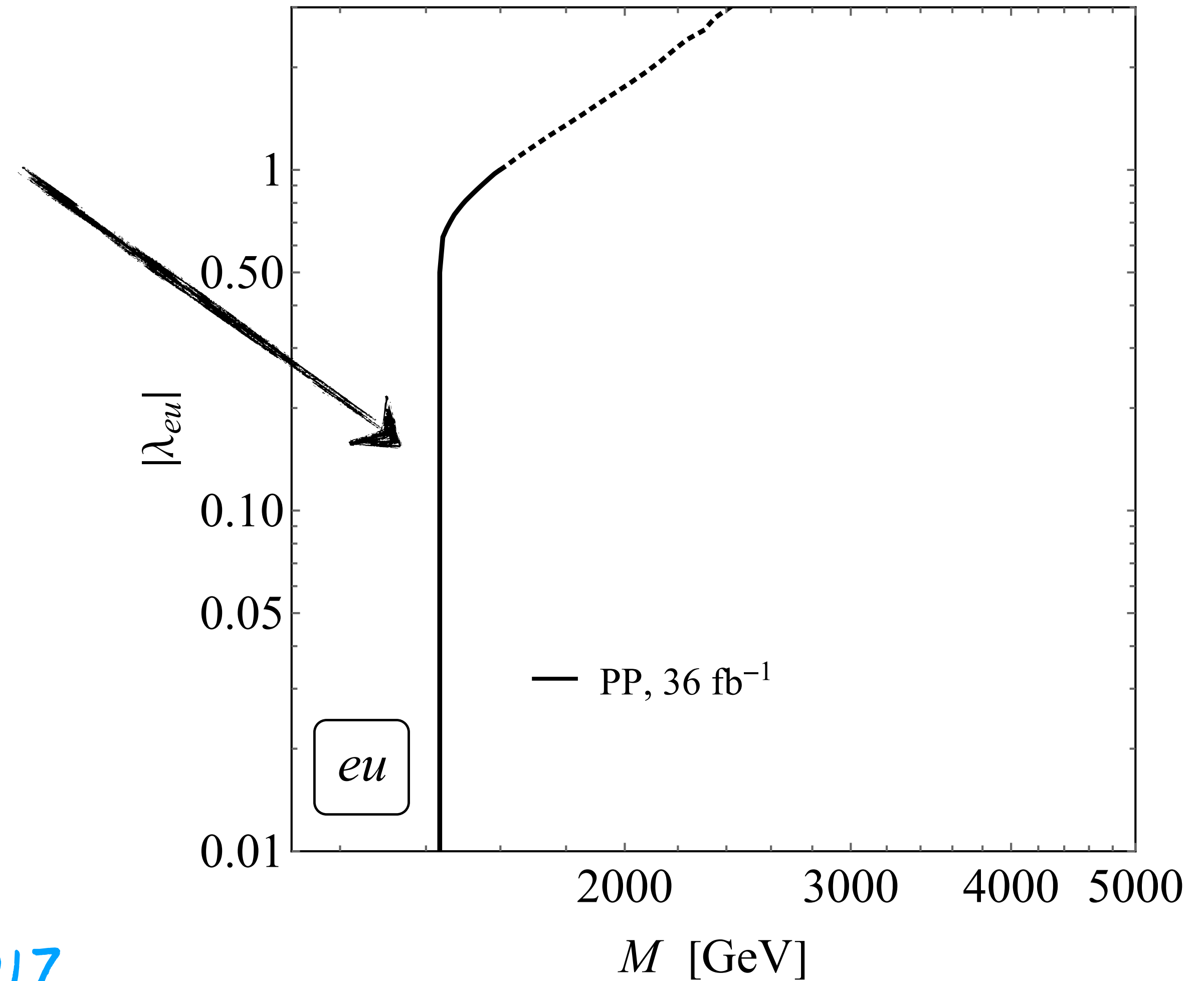
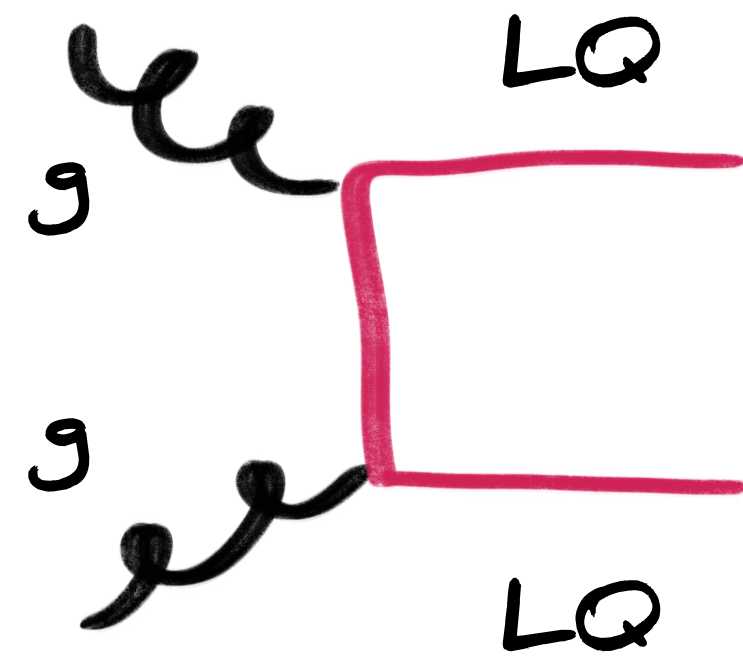
- Since PYTHIA currently[†] cannot handle incoming leptonic partons, initial-state leptons have been replaced by photons to shower events. Our simulation does thus not include leptons but quarks from photon splitting in the parton (PS) shower backward evolution
- In consequence, the jet- & lepton-veto induce a mismodelling of the signal strength. By studying the process $q\bar{q} \rightarrow LQ \rightarrow qll$, we estimate this effect to be of $O(10\%)$ & therefore to only mildly affect the derived LQ limits
- The above PS issue needs to be resolved before next-to-leading order corrections for the signal can be correctly included

[†]In the meantime, Peter Richardson implemented initial-state leptons into HERWIG. Now undergoing testing...

Comments on backgrounds & analysis

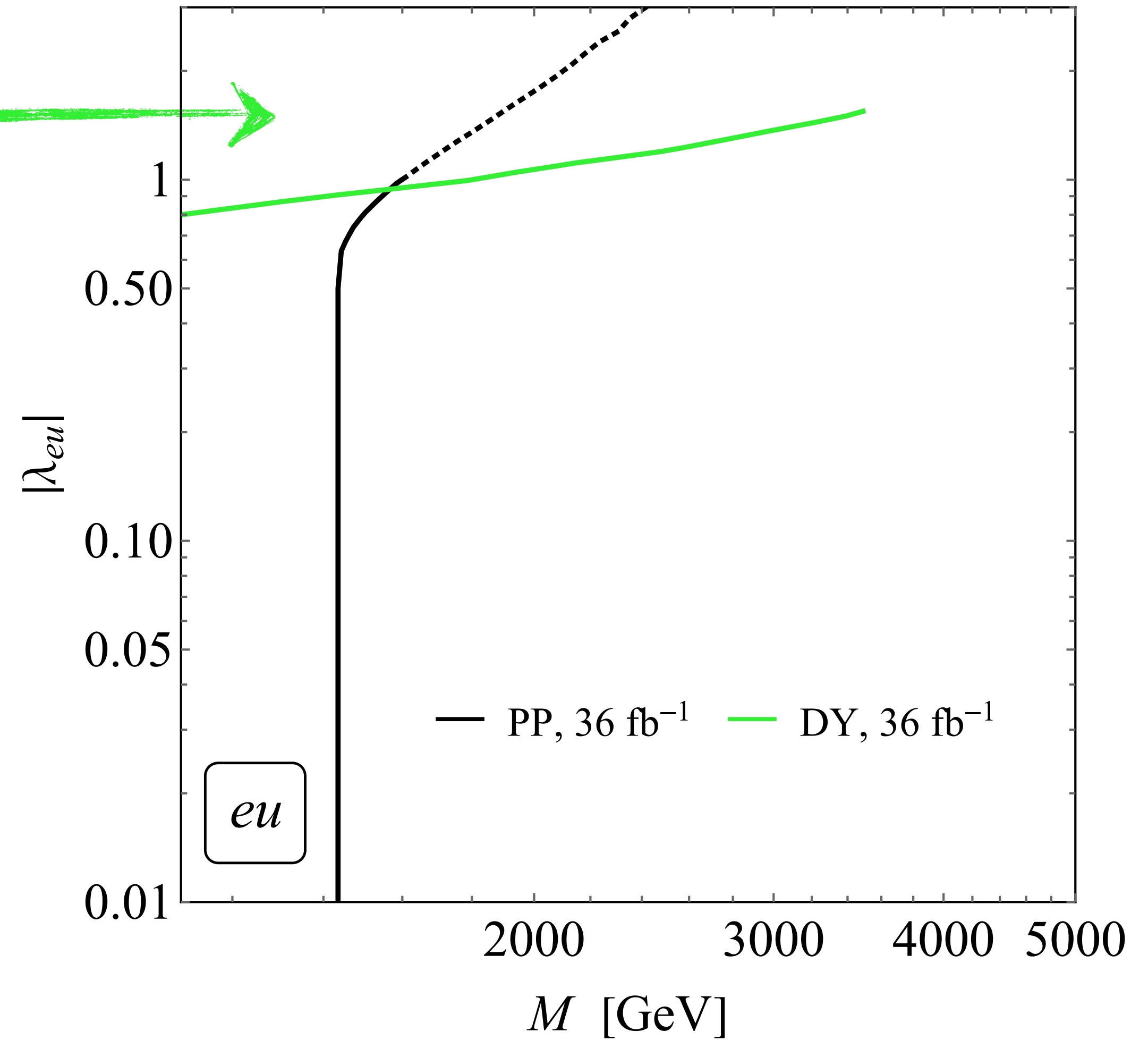
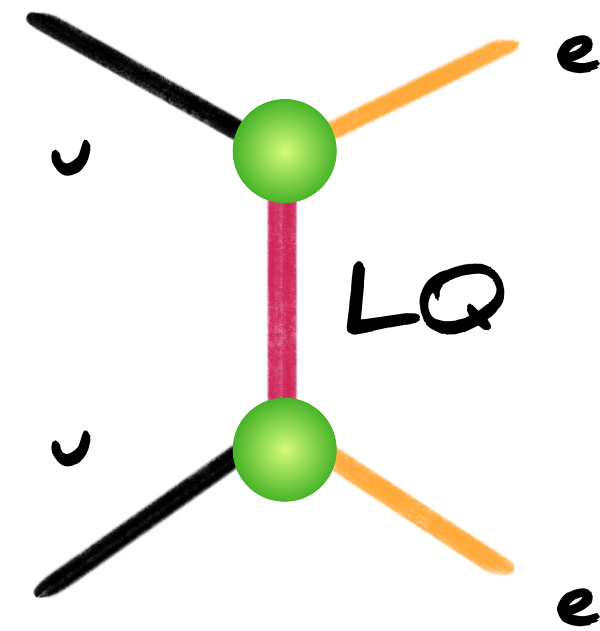
- The impact of multijet backgrounds is incorporated in our analysis by employing the post-fit systematic uncertainties of the ATLAS $l + E_{T,miss}$ search [1906.05609](#). Doubling the quoted errors we find uncertainties of 3%, 30% & 250% (3%, 24% & 50%) for m_{ej} ($m_{\mu j}$) values of 1 TeV, 3 TeV & 5 TeV
- The resolution of the invariant lepton-jet mass is estimated by combining the information on the dilepton & dijet mass resolutions given in [1903.06248](#) & [1910.08447](#). We find that the mass resolution amounts to 2.2% (4.3%) at 1 TeV & 1.5% (11%) at 5 TeV in the electron (muon) case
- Derived limits $IMHO$ conservative, as ATLAS & CMS could probably improve analysis

Summary of 95% CL limits: eu case



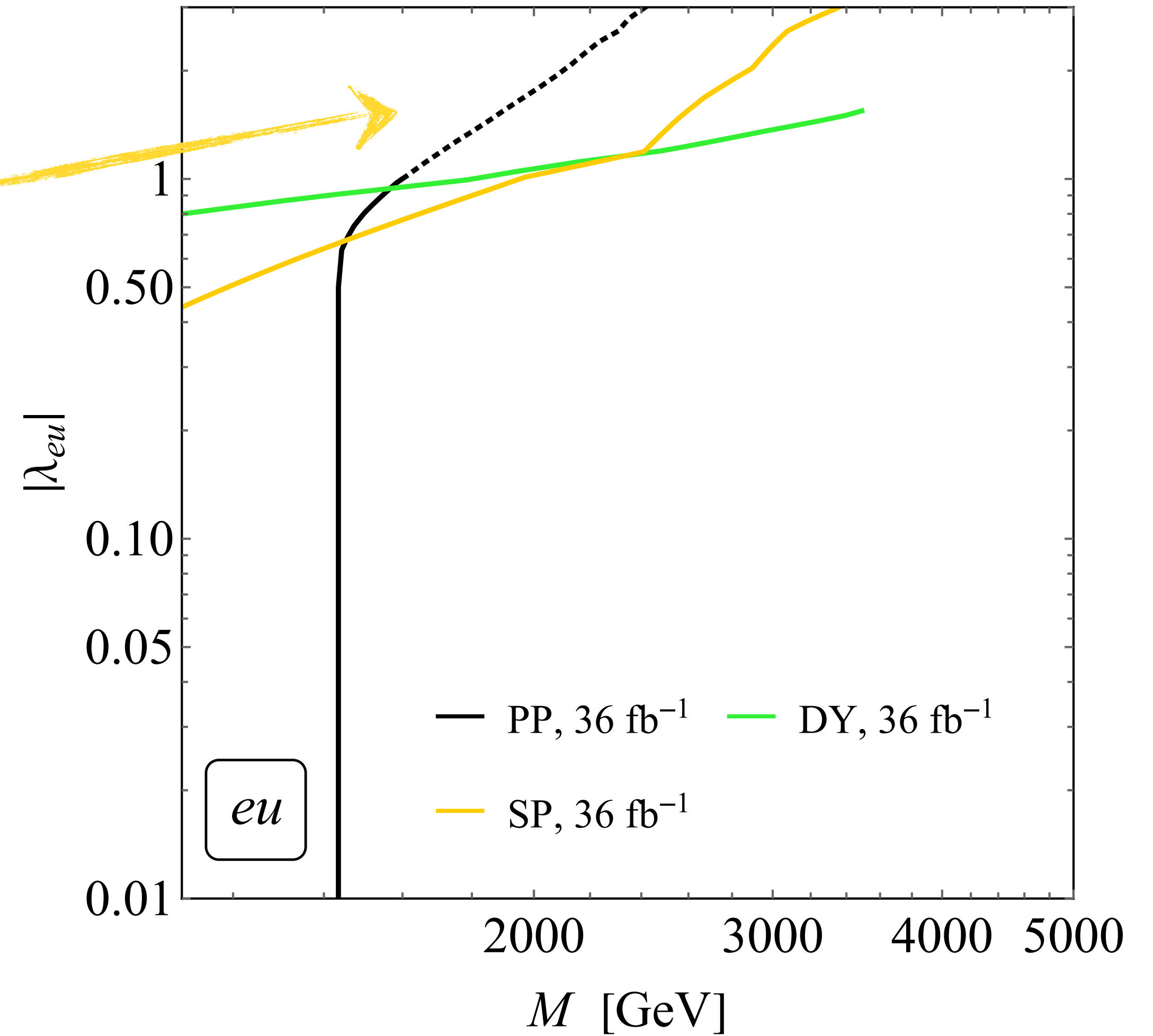
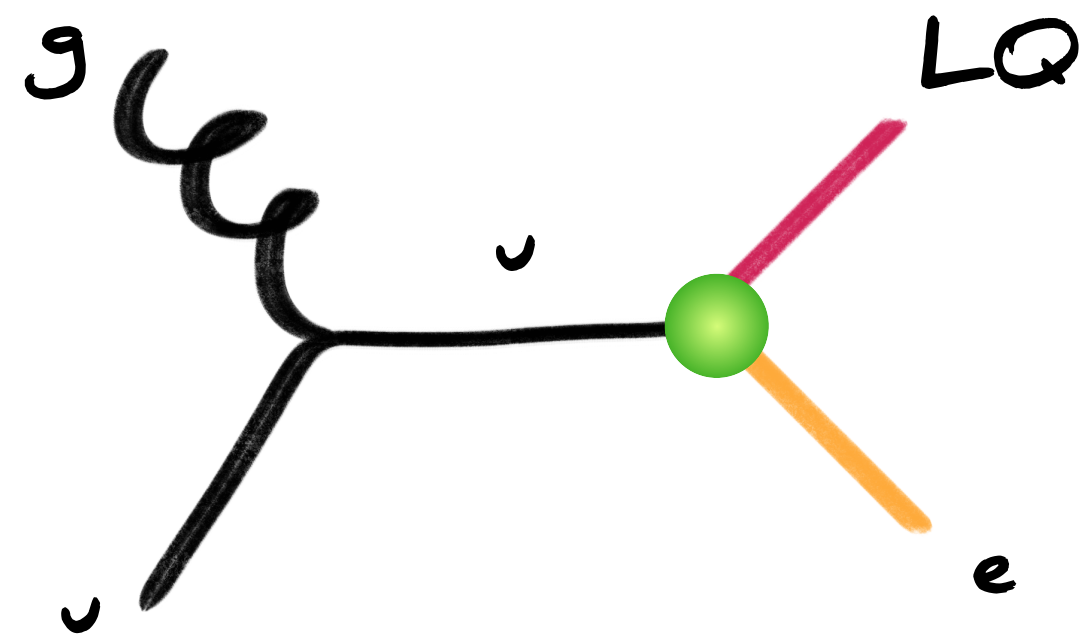
PP limit from Schmaltz & Zhong, 1810.10017

Summary of 95% CL limits: eu case



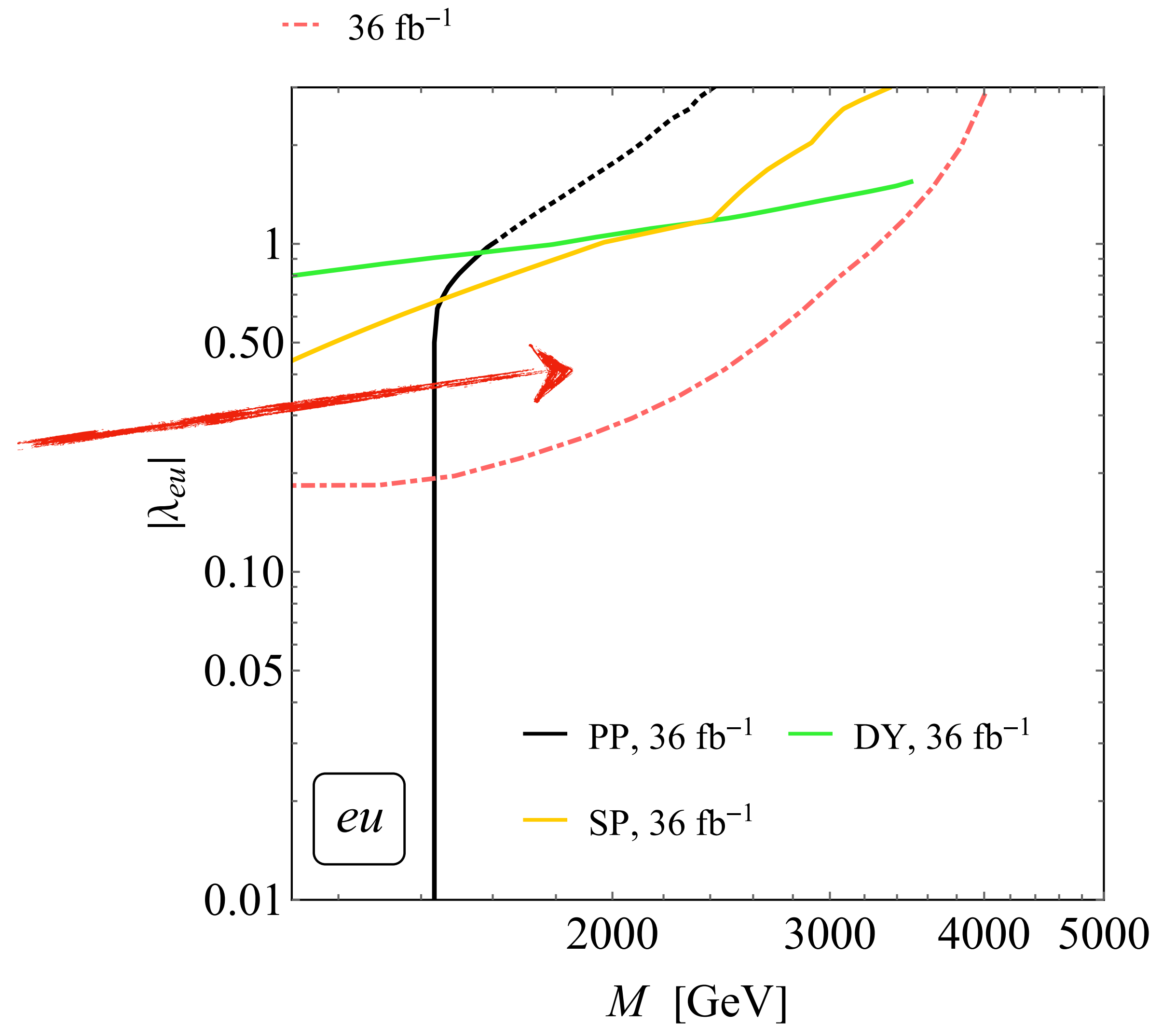
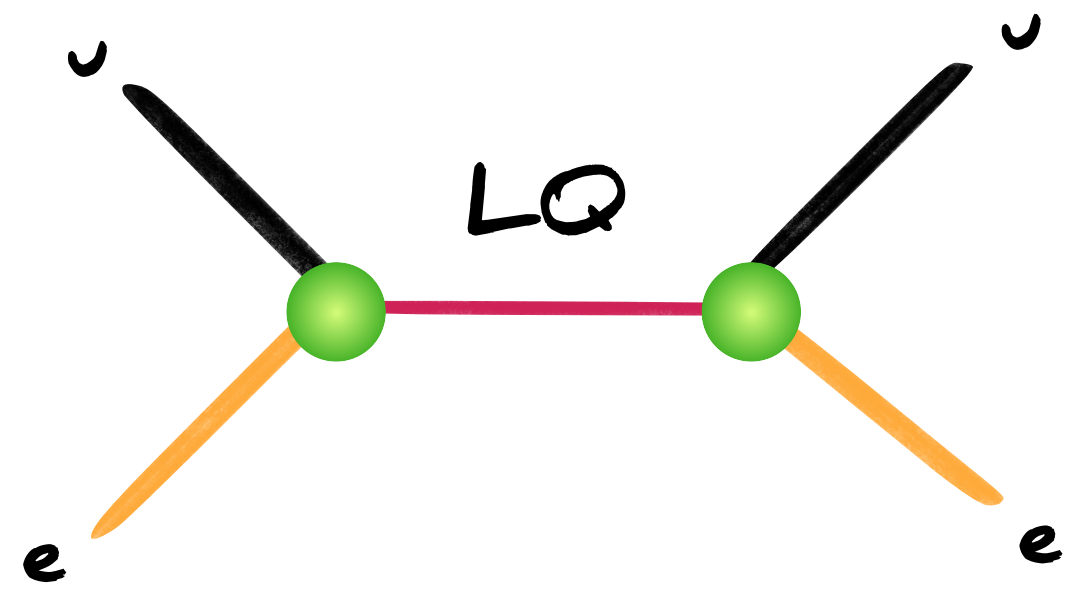
DY limit from Schmaltz & Zhong, 1810.10017

Summary of 95% CL limits: eu case

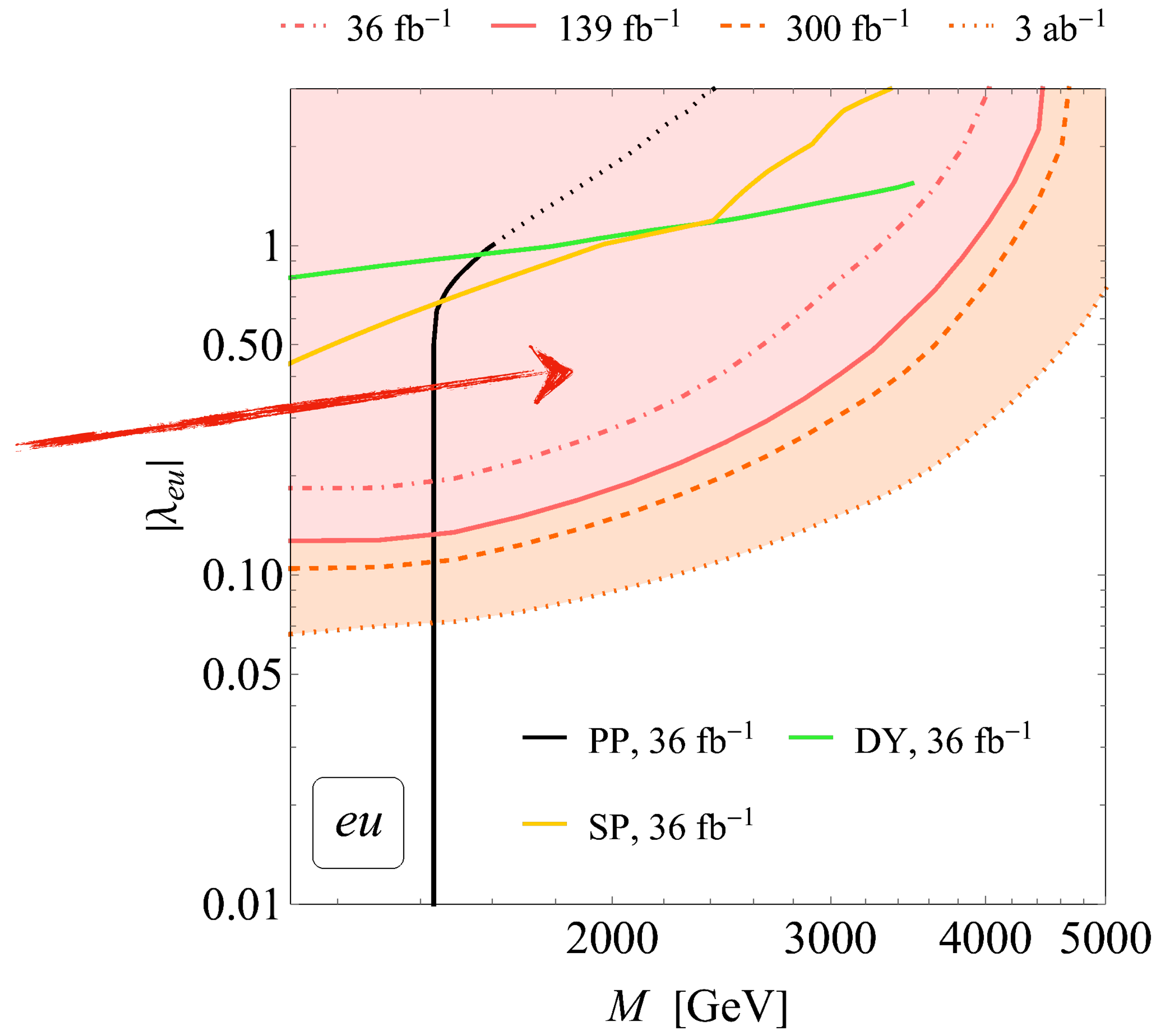
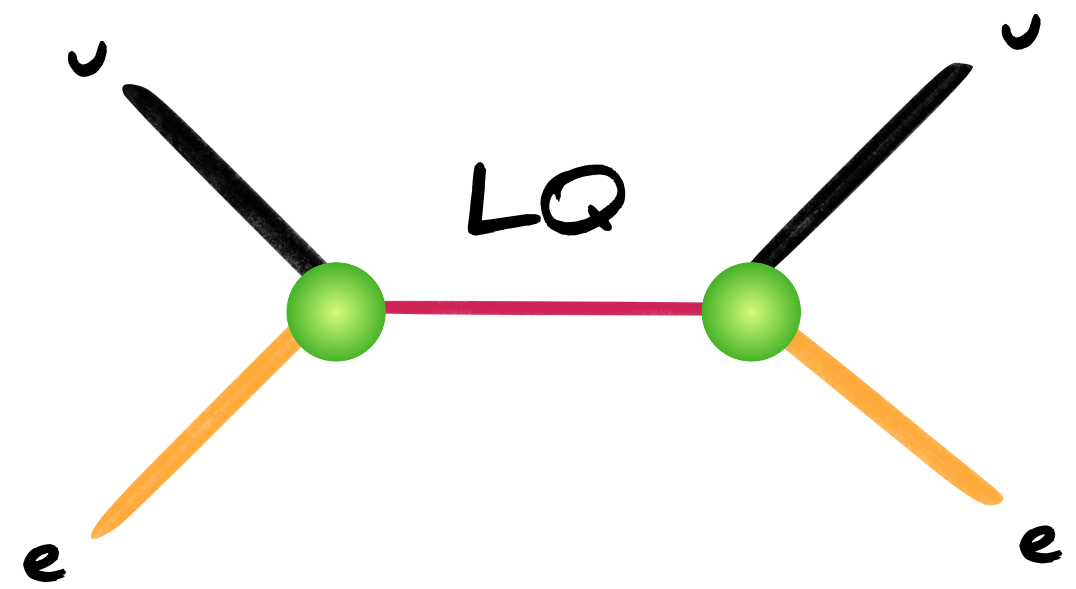


SP limit from Schmaltz & Zhong, 1810.10017

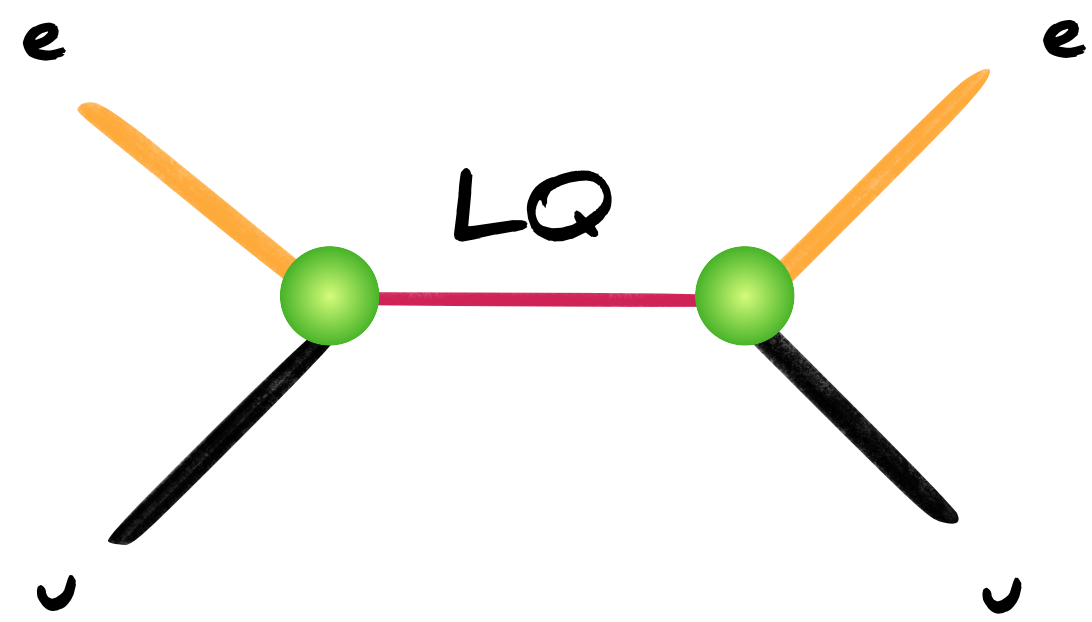
Summary of 95% CL limits: eu case



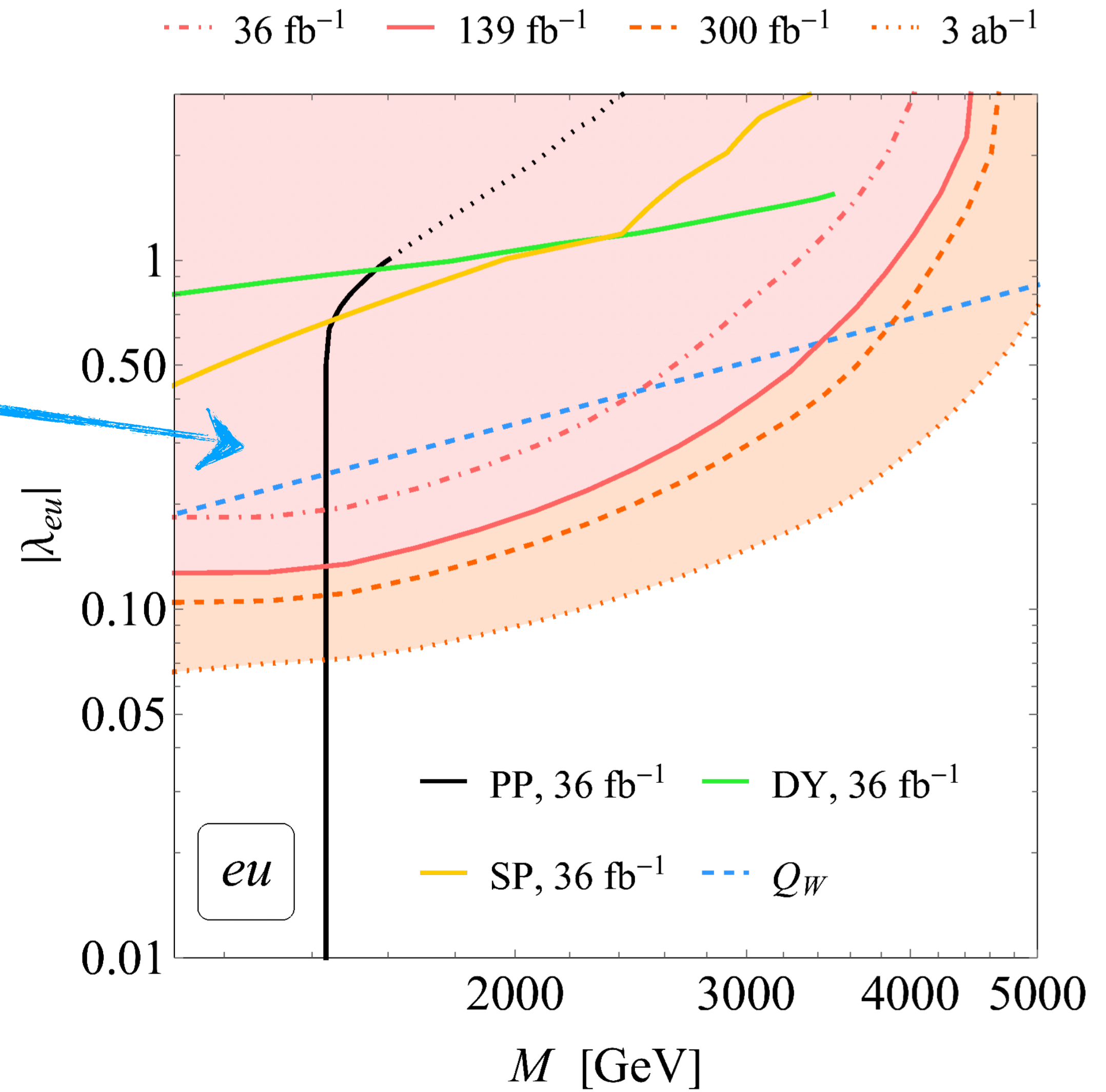
Summary of 95% CL limits: eu case



Summary of 95% CL limits: eu case



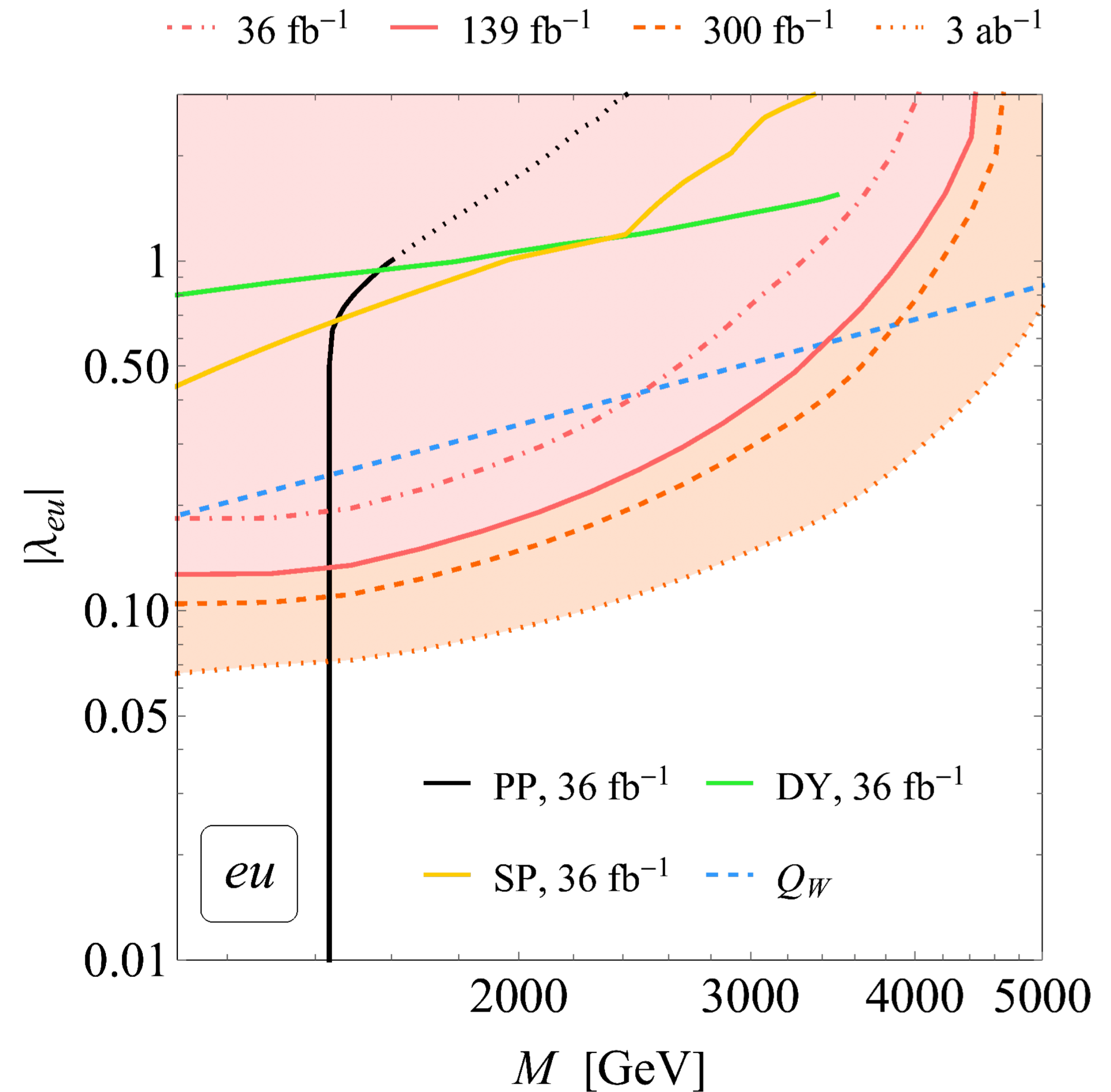
weak charge measurements (Q_W)
 from parity-violating electron
 scattering & atomic parity violation



Q_W limit from Schmaltz & Zhong, 1810.10017

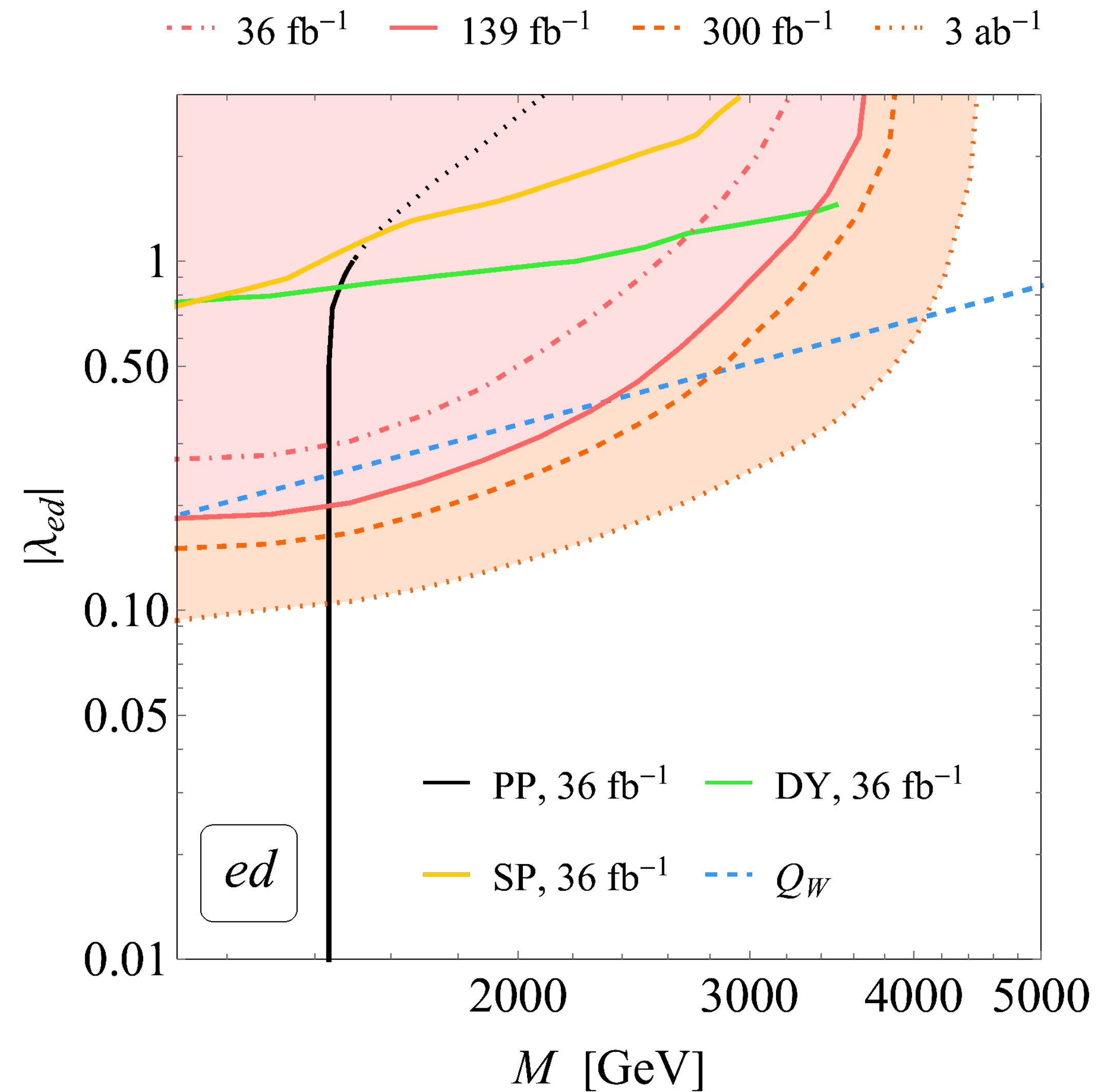
Summary of 95% CL limits: $e\nu$ case

In $e\nu$ case our hypothetical 139 fb^{-1} bounds are more stringent than the constraints from Q_W measurements for LQ masses below around 3.4 TeV . At high-luminosity LHC the corresponding limits can be improved to 5.2 TeV

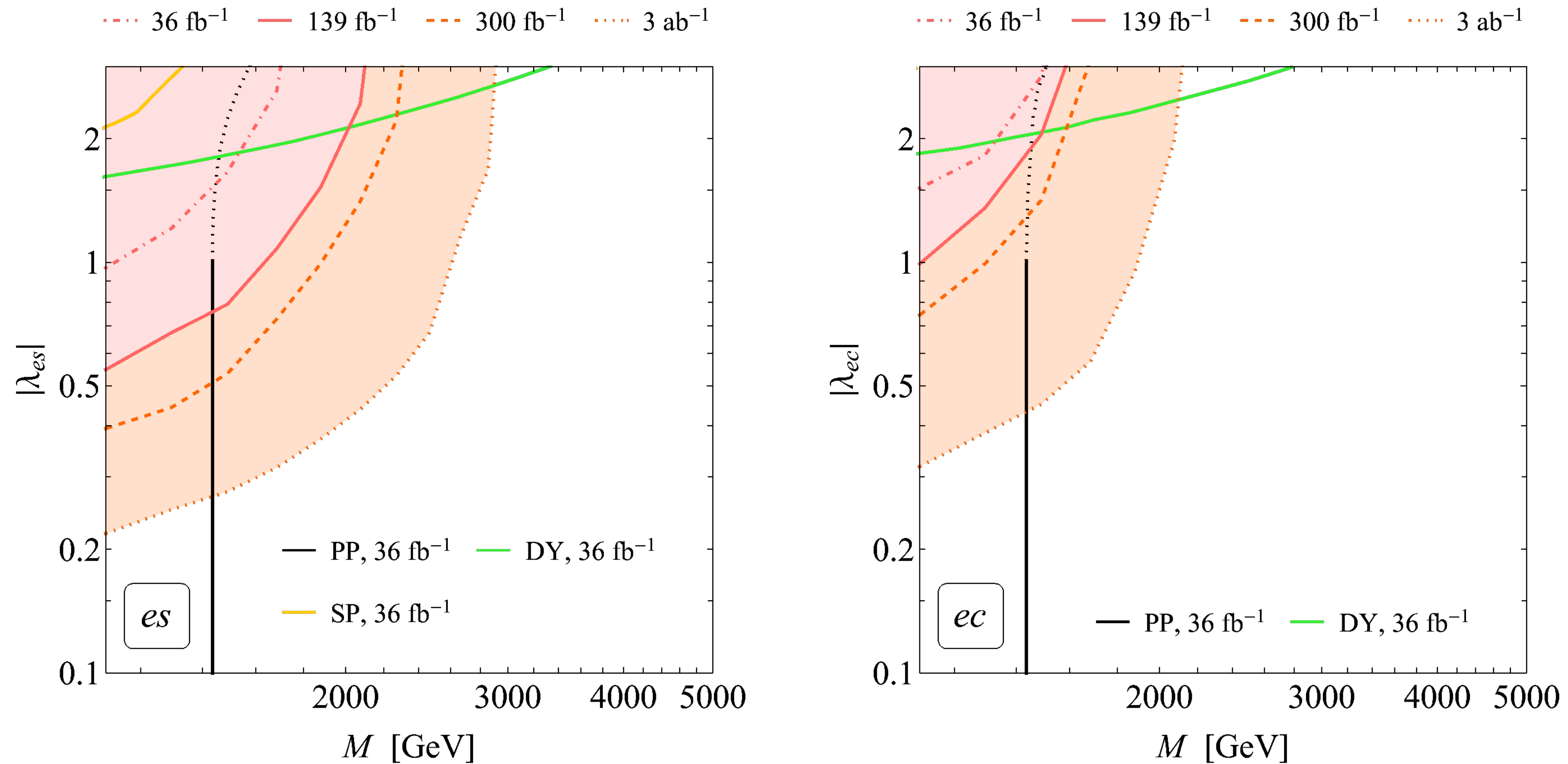


Summary of 95% CL limits: ed case

Slightly weaker limits in ed case due to smaller d luminosity. Nevertheless our bounds are still superior to Q_W measurements for scalar LQ masses below roughly 2.3 TeV (4.1 TeV) assuming 139 fb^{-1} (3 ab^{-1}) of data

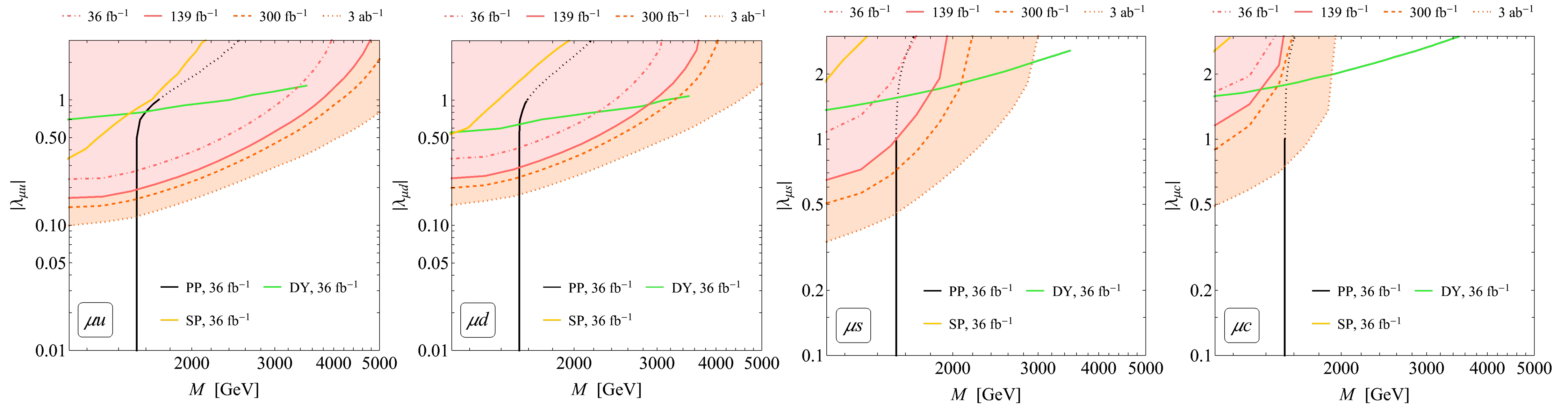


Summary of 95% CL limits: other e cases



Notable weaker constraints in s & c case due to smaller parton distribution functions (PDFs)

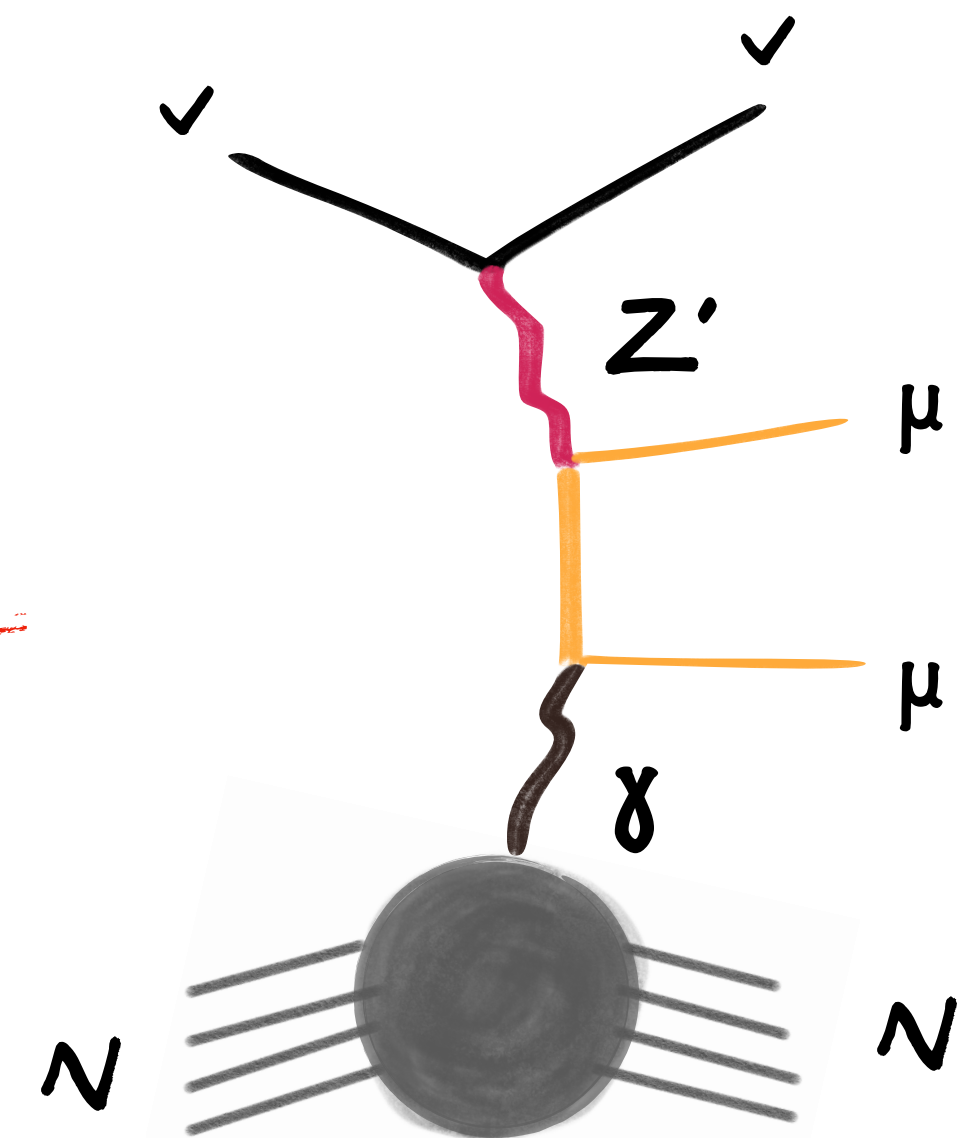
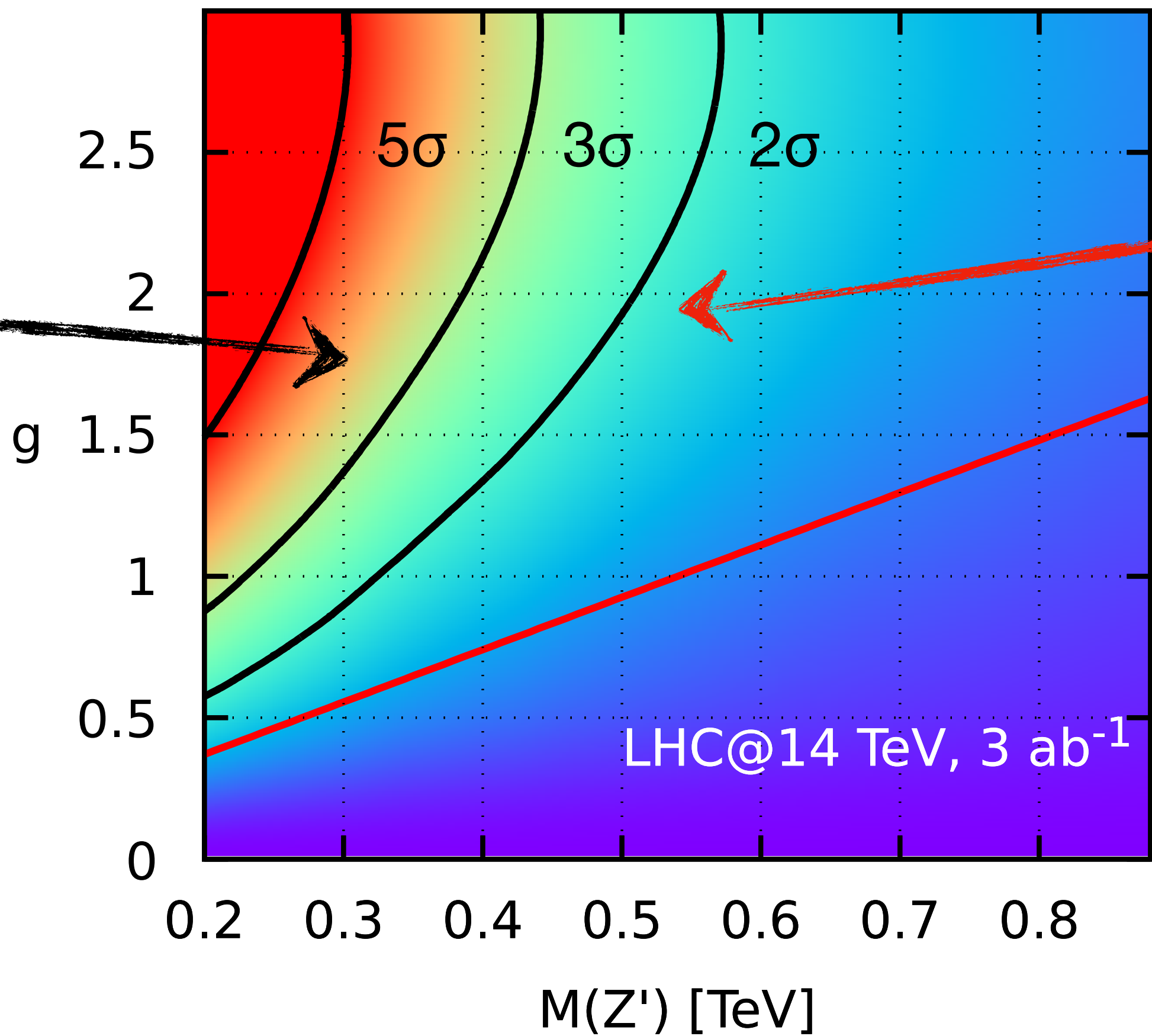
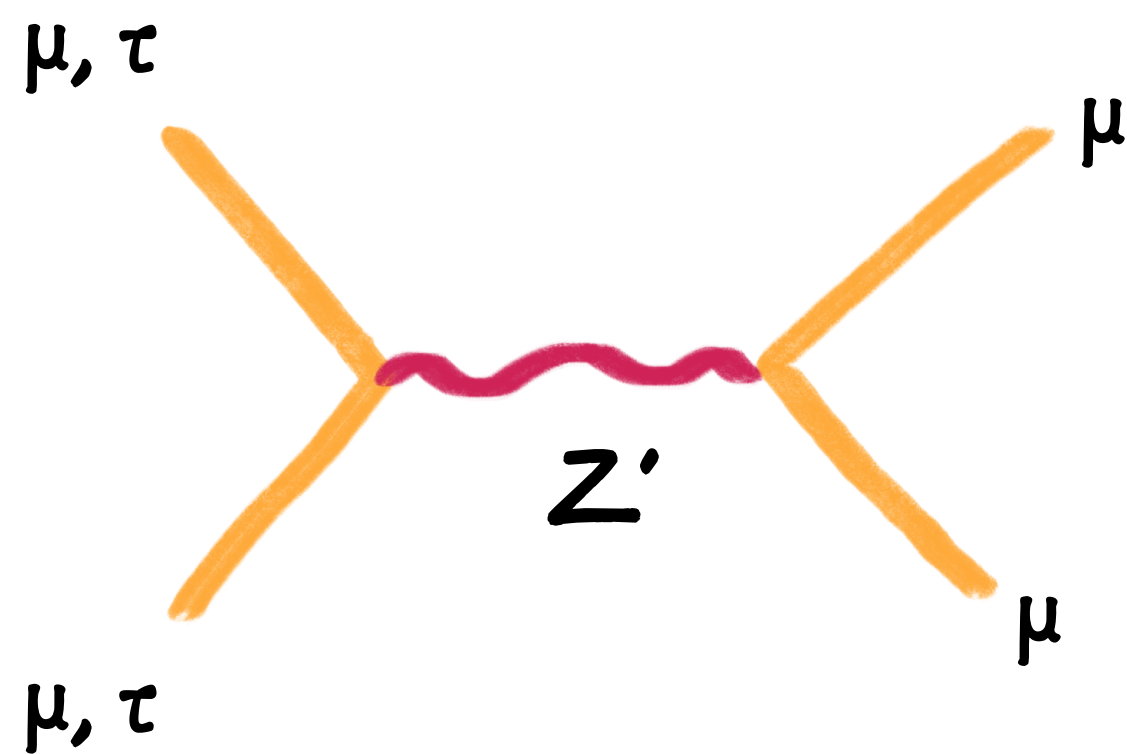
Summary of 95% CL limits: μ cases



Constraints on μq & $e q$ couplings similar. Weaker μq limits at low LQ mass due to lower muon detection efficiencies, but stronger bounds at high LQ mass given lower multijet background

$L_\mu-L_\tau$ Z' bosons

Buonocore et al., 2005.06477

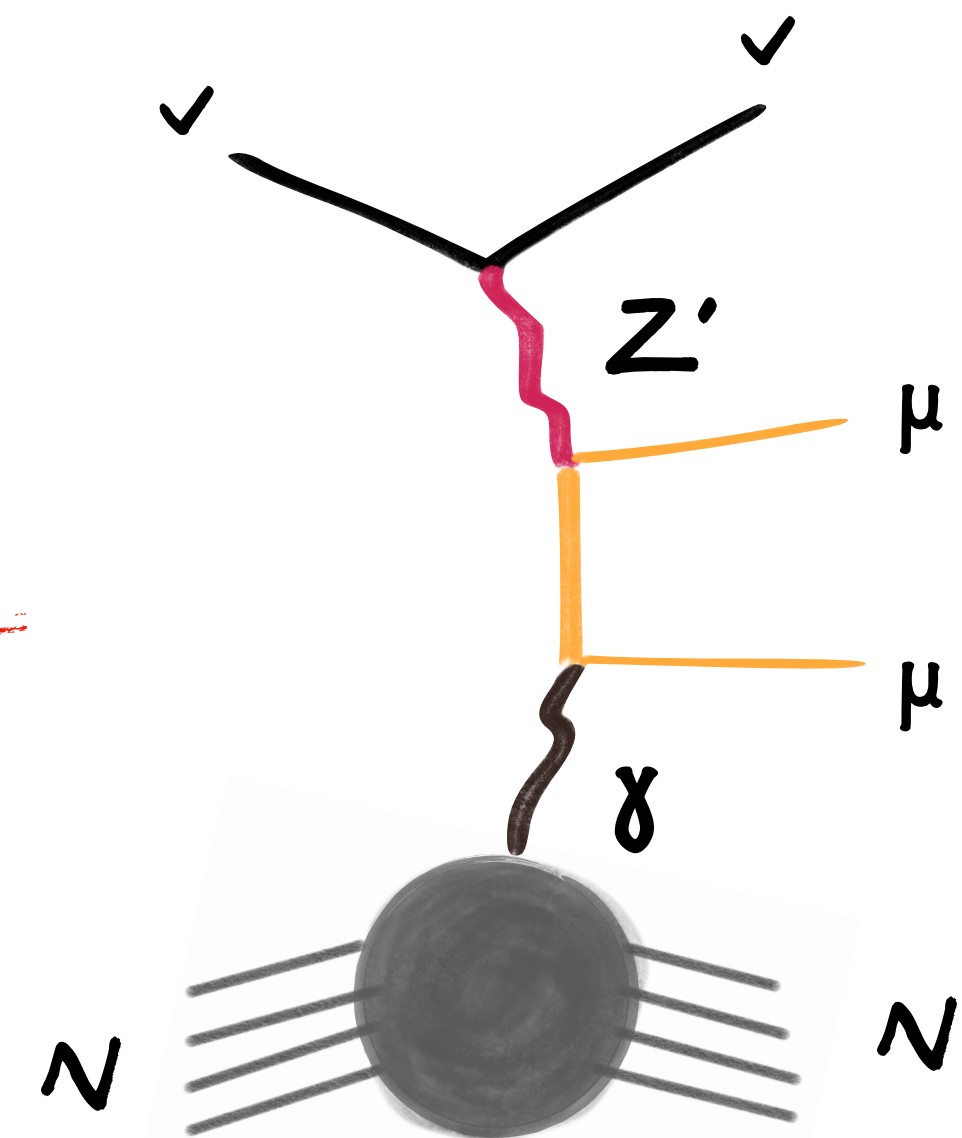
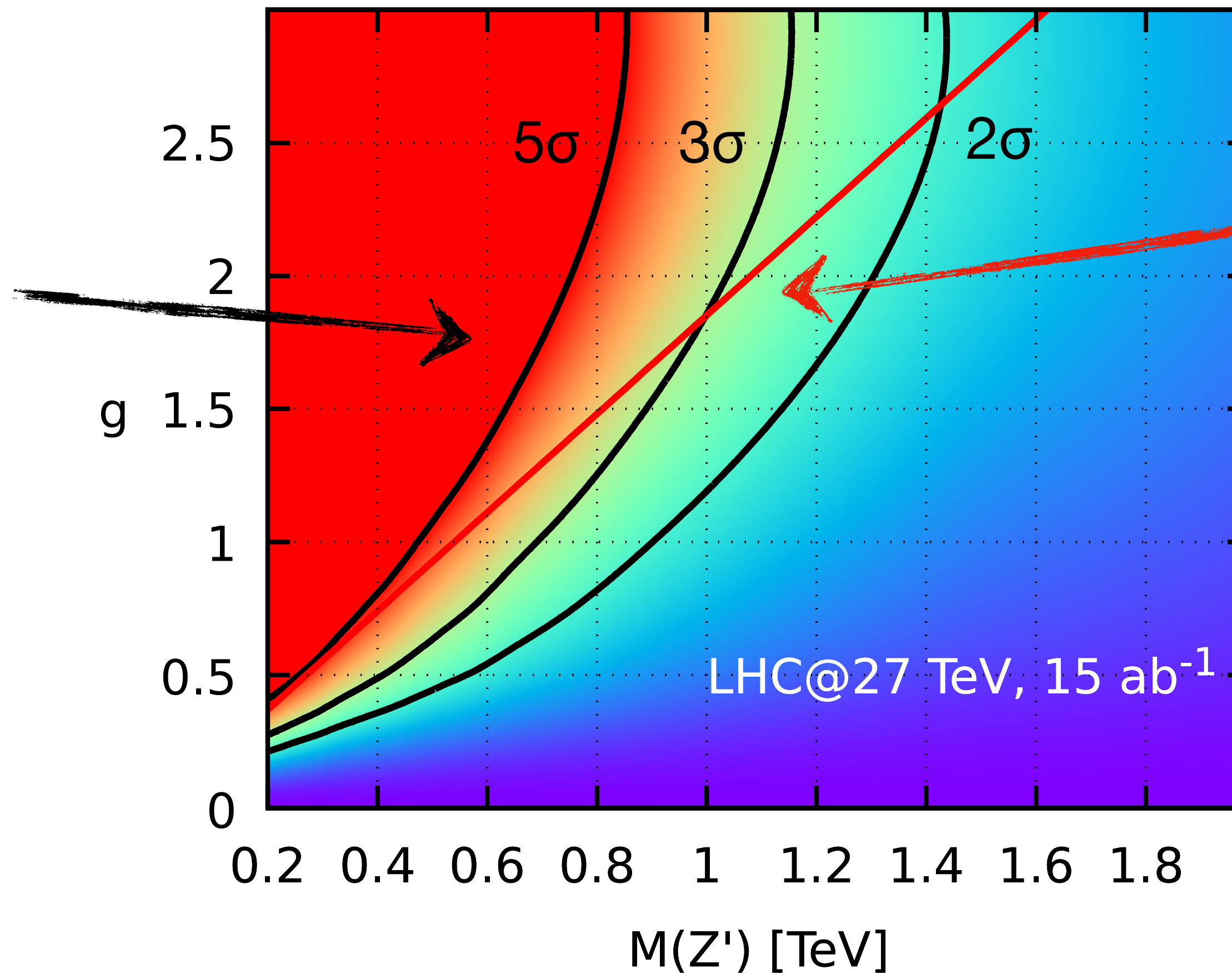
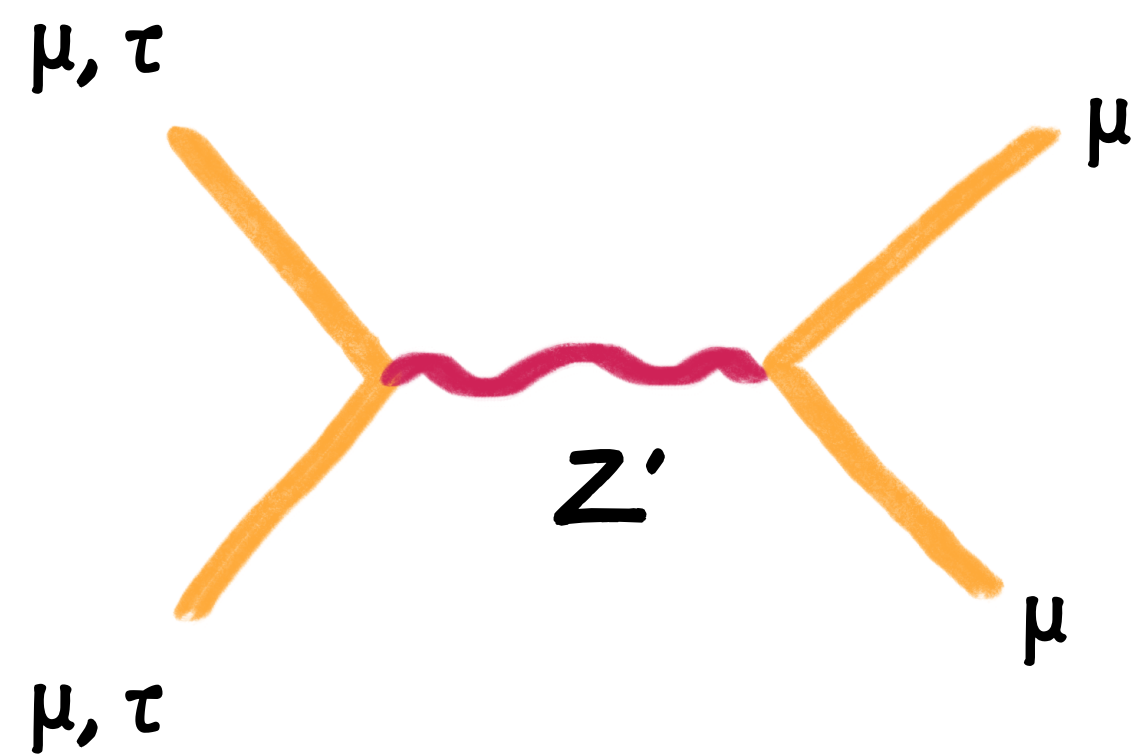


neutrino trident production

Direct LHC limits weaker than indirect constraints from neutrino trident production

$L_\mu-L_\tau$ Z' bosons

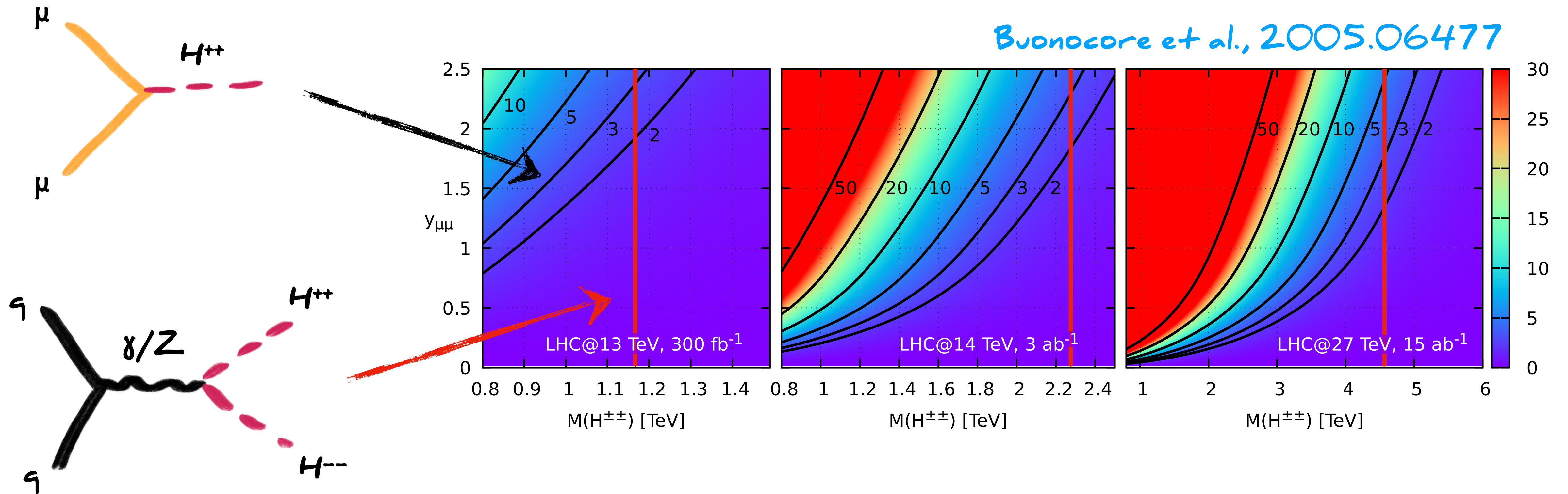
Buonocore et al., 2005.06477



neutrino trident production

Need LHC energy upgrade to make direct & indirect bounds comparable in strength

Doubly-charged Higgses



For sufficiently large Yukawa couplings $y_{\mu\mu}$ s-channel production of a doubly-charged Higgs may have a mass reach comparable to analyses relying upon PP

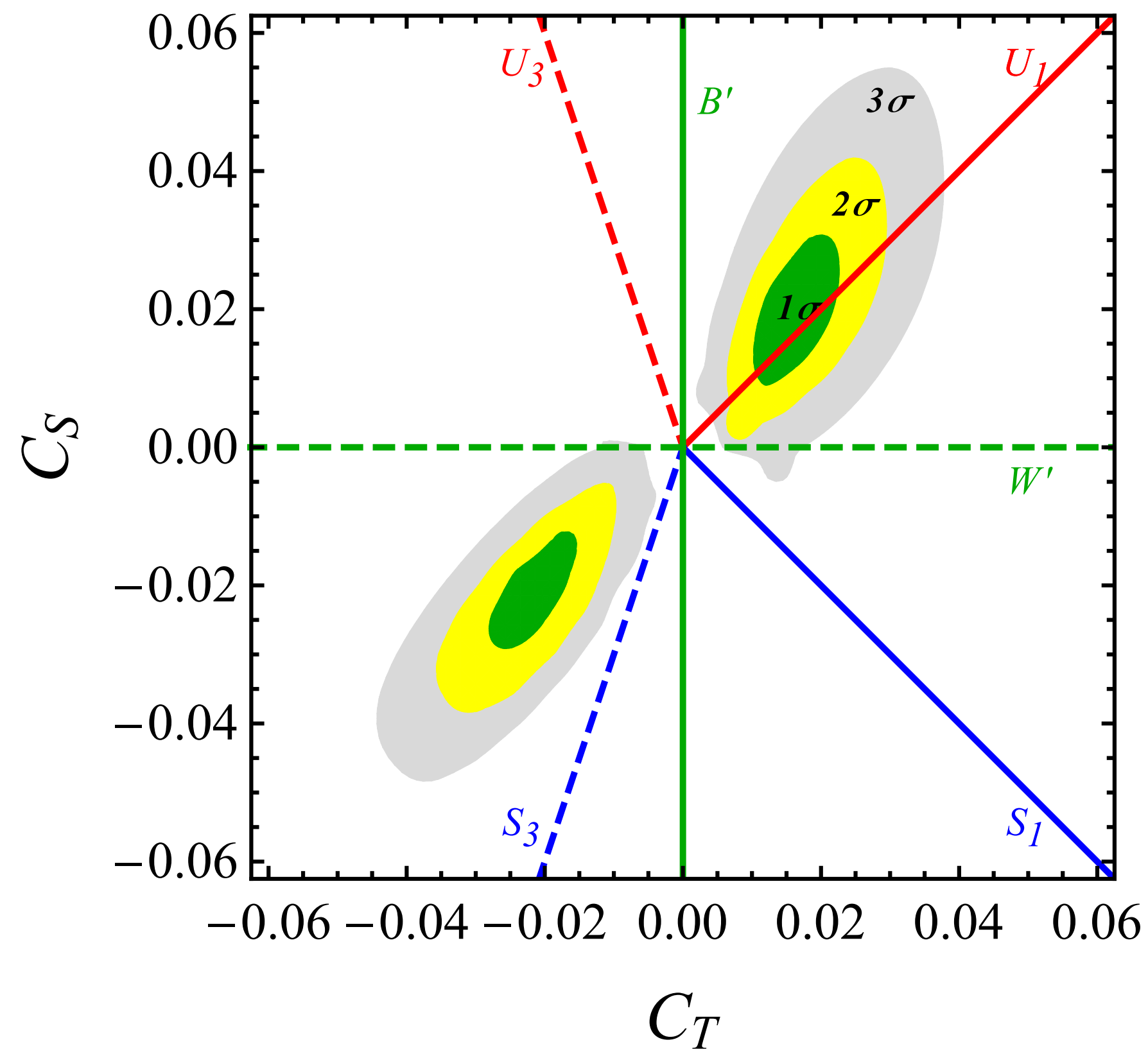
Conclusions & outlook

- Resonant LQ production induced by lepton PDFs provides sensitive direct probes of 1st- & 2nd-generation scalar LQs at the LHC
- In view of the simplicity of the proposed LQ signature & its discovery reach, ATLAS & CMS should perform dedicated resonance searches in lepton-jet final states at future LHC runs
- After some modifications our general search strategy can also be applied to other new-physics searches: leptophilic Z' bosons, doubly-charged Higgses, vector LQs explaining B anomalies, ...



Simplified models for B anomalies

$$\lambda_{ij}^q \lambda_{\alpha\beta}^l \left(C_T (\bar{Q}_L^i \gamma_\mu \sigma^a Q_L^j) (\bar{L}_L^\alpha \gamma^\mu \sigma^a L_L^\beta) + C_S (\bar{Q}_L^i \gamma_\mu Q_L^j) (\bar{L}_L^\alpha \gamma^\mu L_L^\beta) \right)$$



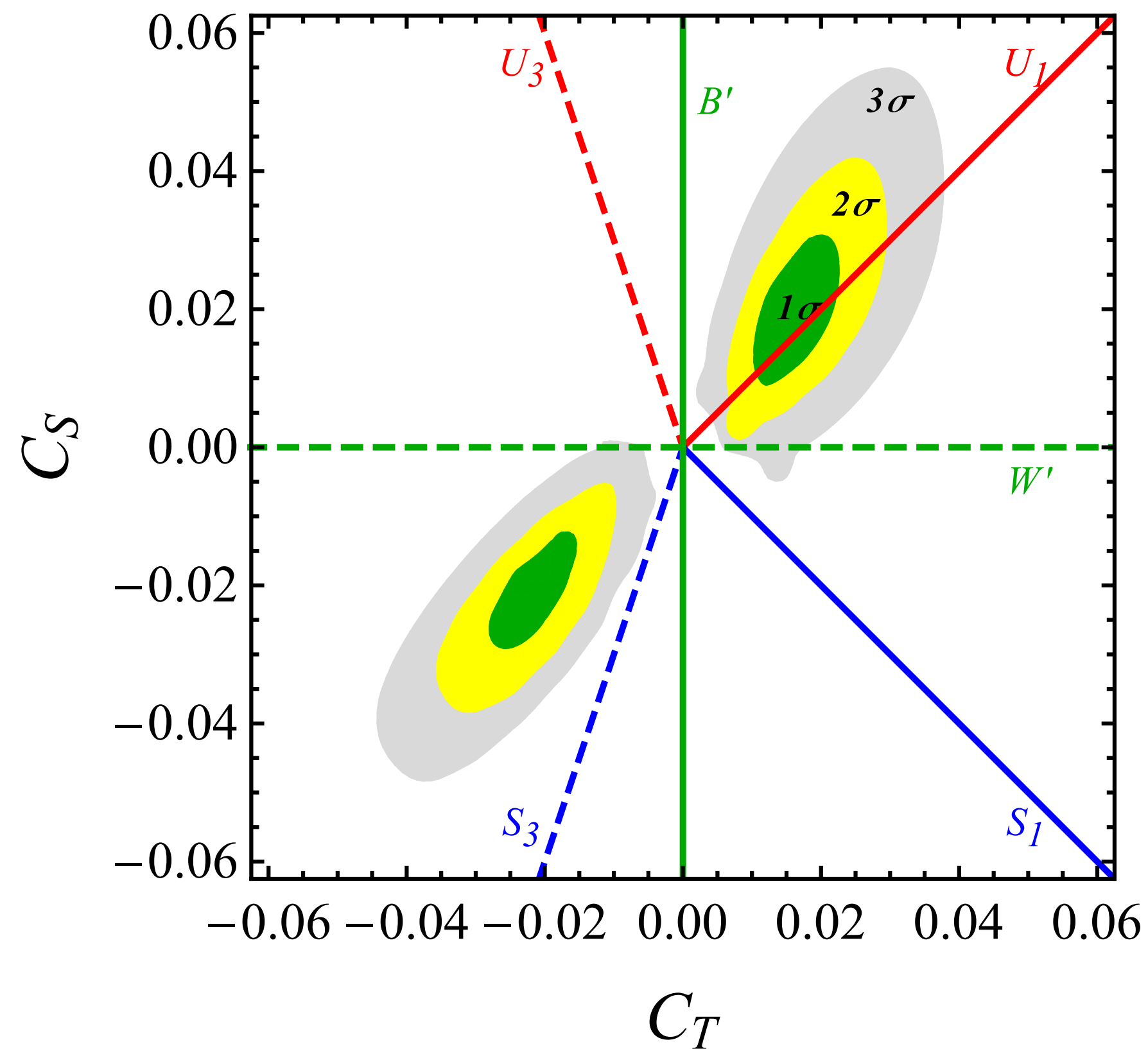
Model	Mediator	$b \rightarrow s$	$b \rightarrow c$
Colorless vectors	$B' = (1, 1, 0)$	✓	✗
	$W' = (1, 3, 0)$	✗	✓
Scalar leptoquarks	$S_1 = (\bar{3}, 1, 1/3)$	✗	✓
	$S_3 = (\bar{3}, 3, 1/3)$	✓	✗
Vector leptoquarks	$U_1 = (3, 1, 2/3)$	✓	✓
	$U_3 = (3, 3, 2/3)$	✓	✗

$b \rightarrow s$ ($b \rightarrow c$) anomalies alone can be explained by several simple single-mediator models

see for instance Buttazzo et al., 1706.07808

Simplified models for B anomalies

$$\lambda_{ij}^q \lambda_{\alpha\beta}^l \left(C_T (\bar{Q}_L^i \gamma_\mu \sigma^a Q_L^j) (\bar{L}_L^\alpha \gamma^\mu \sigma^a L_L^\beta) + C_S (\bar{Q}_L^i \gamma_\mu Q_L^j) (\bar{L}_L^\alpha \gamma^\mu L_L^\beta) \right)$$



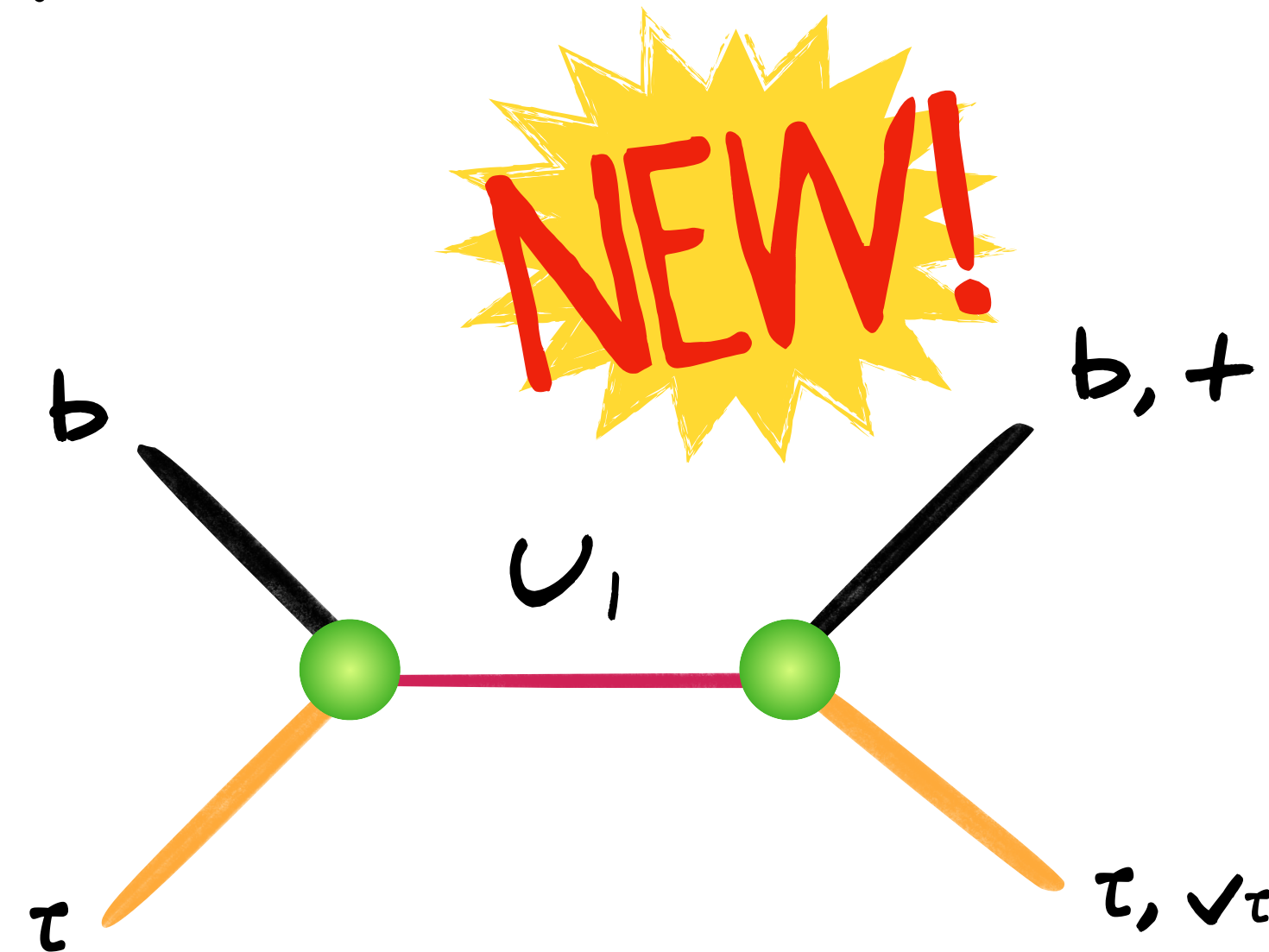
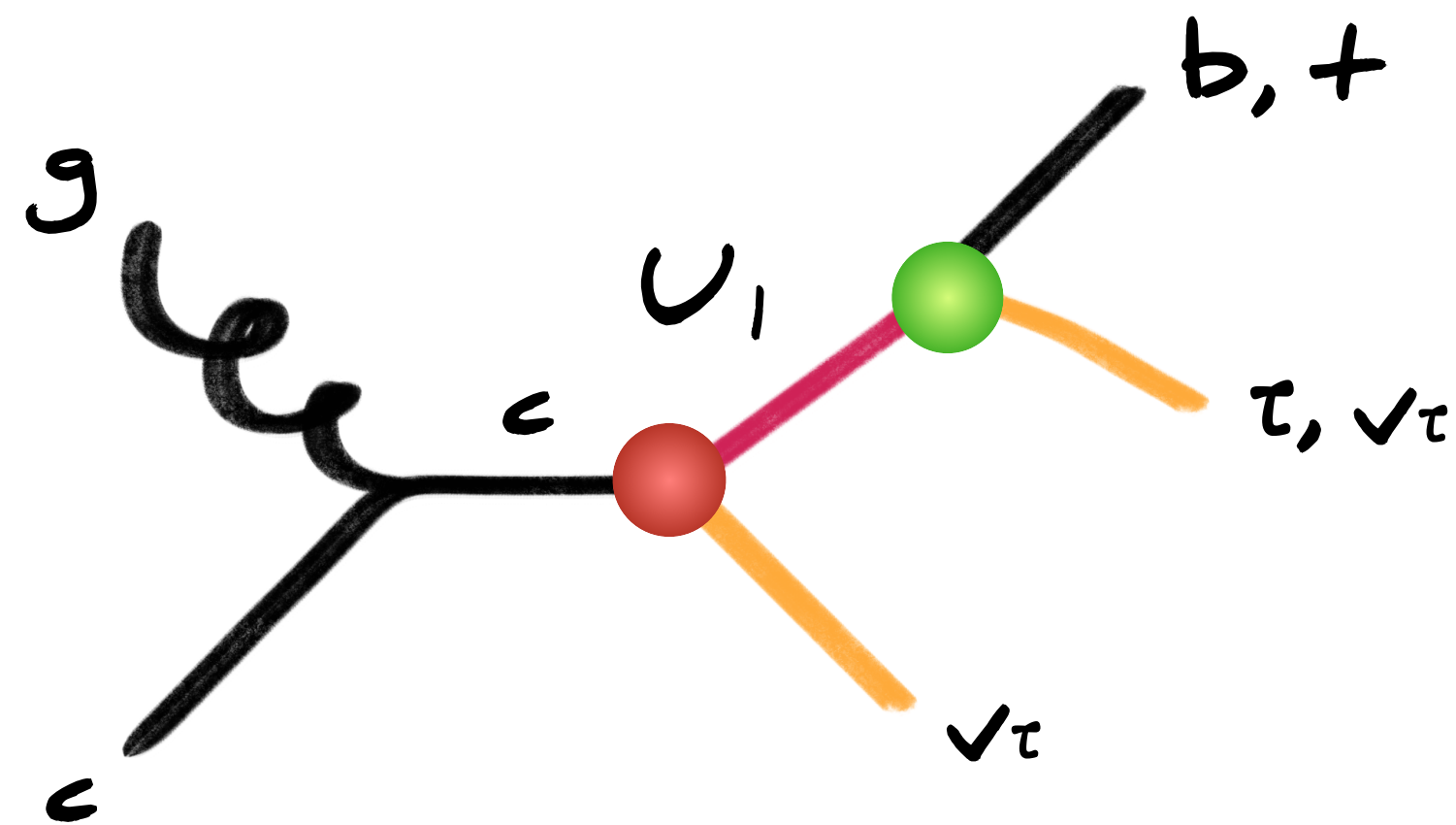
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	$W' = (1, 3, 0)$	✗	✓
Scalar leptoquarks	$S_1 = (\bar{3}, 1, 1/3)$	✗	✓
	$S_3 = (\bar{3}, 3, 1/3)$	✓	✗
Vector leptoquarks	$U_1 = (3, 1, 2/3)$	✓	✓
	$U_3 = (3, 3, 2/3)$	✓	✗

U_1 vector LQ is the only single-mediator model that can explain both sets of anomalies

see for instance Buttazzo et al., 1706.07808

New resonant U_1 search channels

$$\mathcal{L} \supset (g_{33} \bar{q}_L^3 \gamma_\mu l_L^3 + g_{23} \bar{q}_L^2 \gamma_\mu l_L^3) U_1^\mu + \dots$$



Lepton-initiated scatterings allow for new resonant U_1 contributions to $b\tau$, mono-top & mono-jet production. This feature has not been explored yet