

# Non-linear optimization of the BDS using MADX-PTC

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# Motivation & contents

Compensate the residual high order aberrations of the FFS to increase luminosity

- Analytical computation of beam sizes at the IP
- The beam sigmas order by order
- Optimization of beam sizes
- What about the usable luminosity?
- Outlook

# Computation of beam sizes at the IP

Given the transfer map between one location of the accelerator and the IP in the form:

$$x_{IP} = \sum X_{jklmn} x^j p_x^k y^l p_y^m \delta^n$$

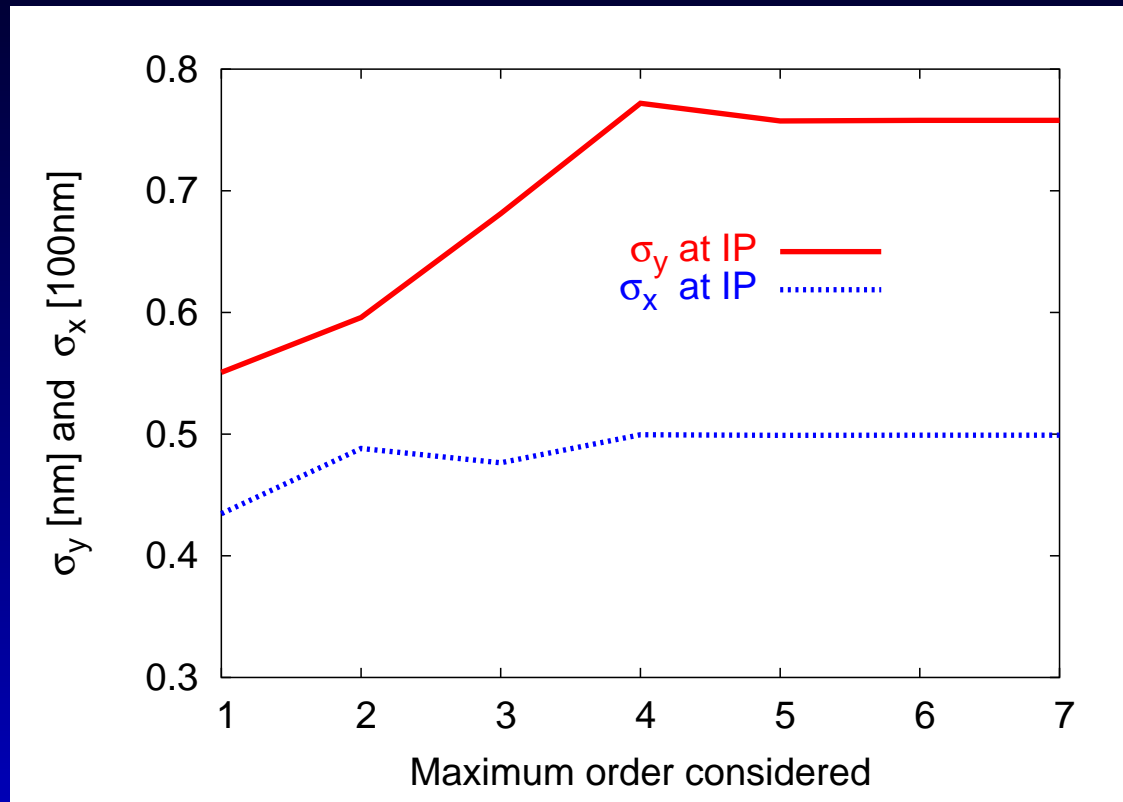
and given the particle density at the initial location, the rms beam size at the IP is given by:

$$\sigma_{IP}^2 = \sum X_{jklmn} X_{j'k'l'm'n'} \int x^{j+j'} p_x^{k+k'} y^{l+l'} p_y^{m+m'} \delta^{n+n'} \rho dv$$

The integral is performed depending on  $\rho$  and the  $X_{jklmn}$  are obtained to arbitrary order from MADX-PTC (thanks to E. Forest, F. Schmidt, et al).

# The beam sigmas order by order

## Nominal FFS design and beam parameters



Aberrations exist up to the fifth order!

# Biggest contributors to vertical size

Contribution to $\sigma_y^2$	j	k	l	m	n	j'	k'	l'	m'	n'
6.109200e-19	0	0	1	0	0	0	0	1	0	0
-1.43715e-19	0	0	1	2	0	0	0	1	0	0
1.221256e-19	1	1	0	1	1	1	1	0	1	1
1.166679e-19	0	2	0	1	1	0	2	0	1	1
1.035384e-19	0	0	0	1	0	0	0	0	3	0
-1.02905e-19	0	2	0	1	1	0	0	0	1	1
9.762807e-20	0	0	1	0	0	0	0	1	0	2
-8.07838e-20	2	0	0	1	1	0	0	0	1	1
7.476994e-20	0	0	0	1	0	0	0	0	1	0
7.189904e-20	2	0	0	1	1	2	0	0	1	1
6.807446e-20	0	0	0	1	1	0	0	0	1	1
-6.58350e-20	0	0	0	1	2	0	0	0	1	0
...										

Third and Fourth order terms are the heaviest after linear

# Optimization of the beam sizes

MAPCLASS is a Python code that:

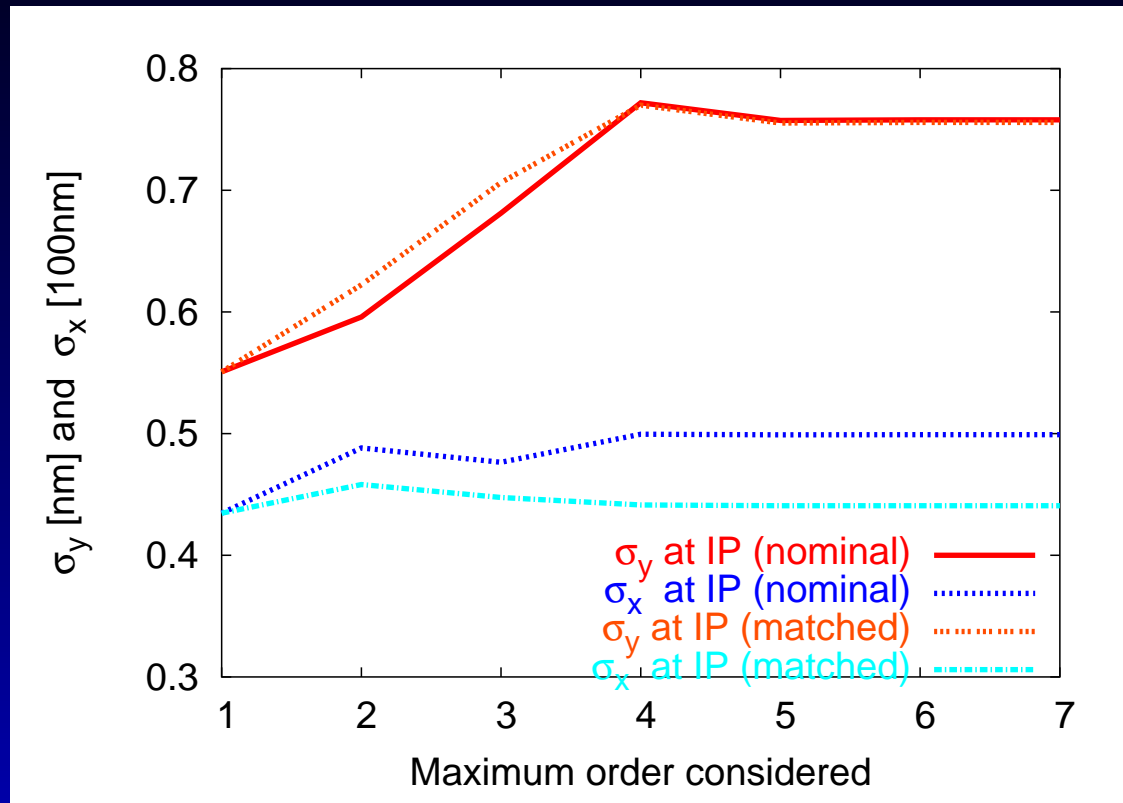
- Generates/reads MADX input/output
- Computes  $\sigma_x$  and  $\sigma_y$  from map coefficients
- Uses either SVD or the Simplex algorithms to optimize the sigmas<sup>†</sup>

Four examples follow on the optimization of the FFS by varying strengths.

Luminosity is computed for the different cases with PLACET without SR.

<sup>†</sup> in general any quantity can be optimized

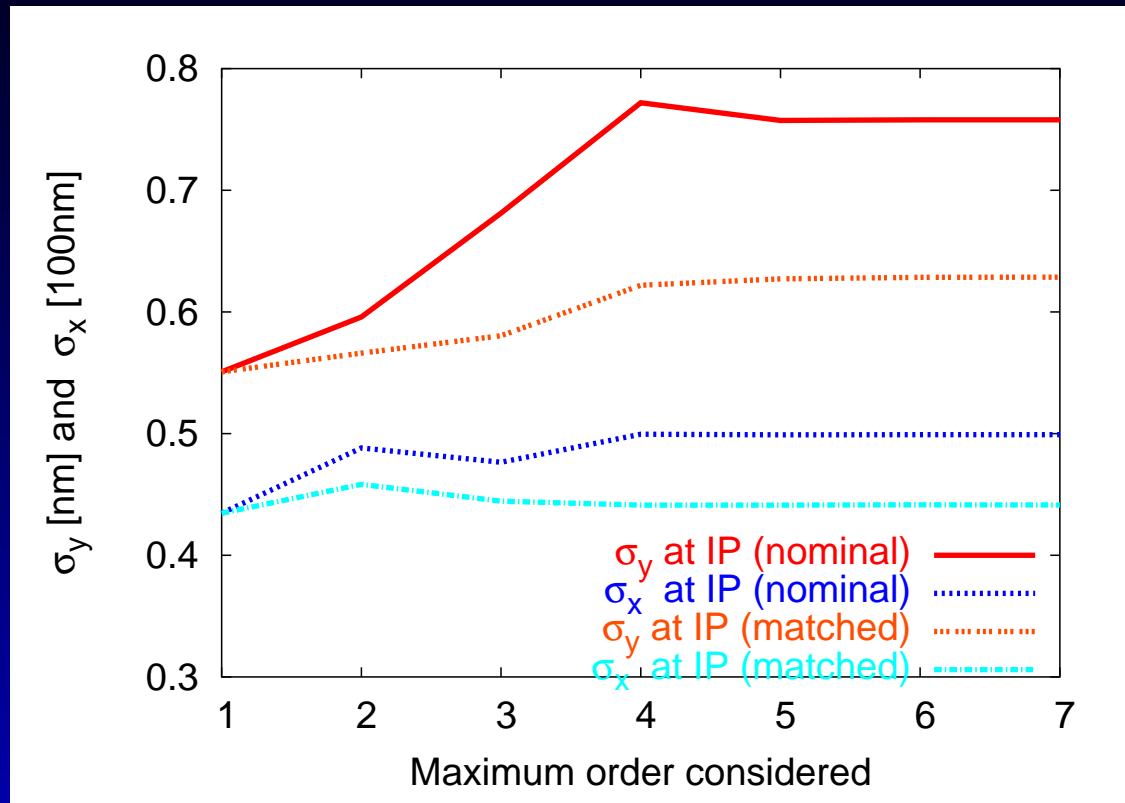
# Using only the sextupoles



A small adjustment in the strengths of the 5 sextupoles gives a 12% reduction in  $\sigma_x$  leaving  $\sigma_y$  unchanged.

Luminosity increase of 4%

# Plus 1 sextupole, 1 octupole & 2 decapoles

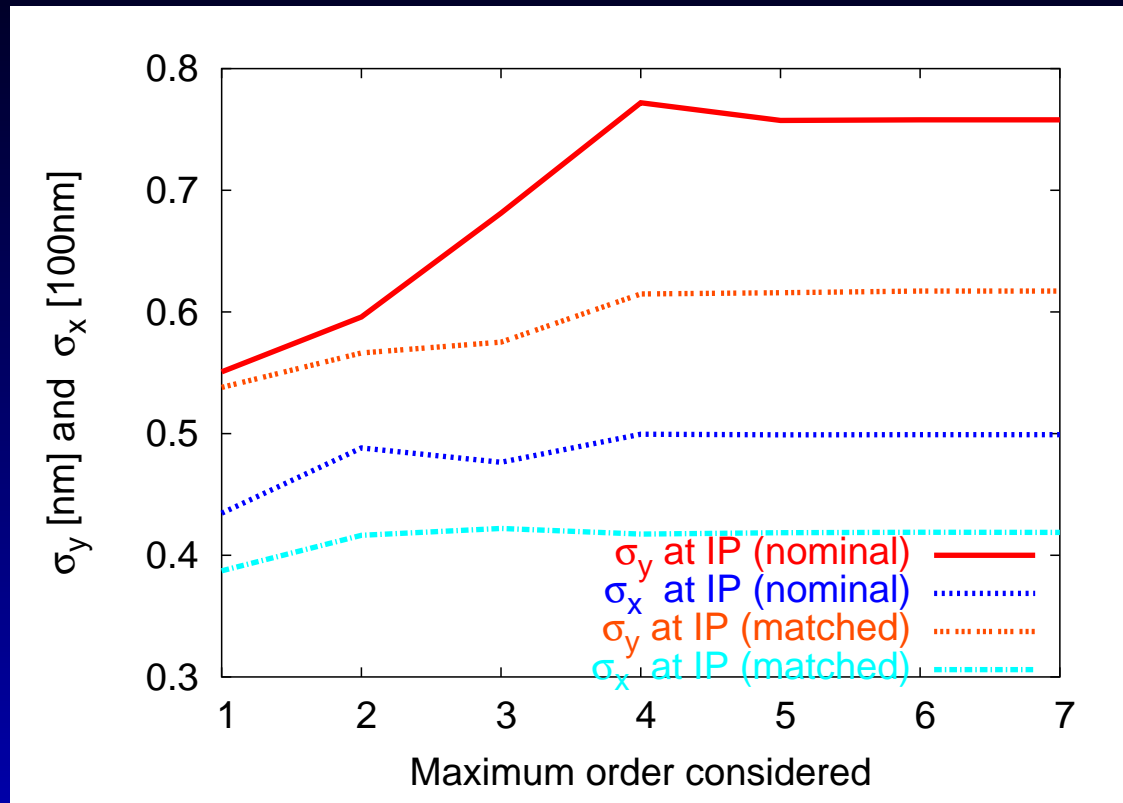


With a total of 9 parameters  $\sigma_x$  and  $\sigma_y$  are reduced by 12% and 17% respectively.

Luminosity increase of 5%  $\Rightarrow$  Not corresponding



# Plus using the strength of the quadrupoles

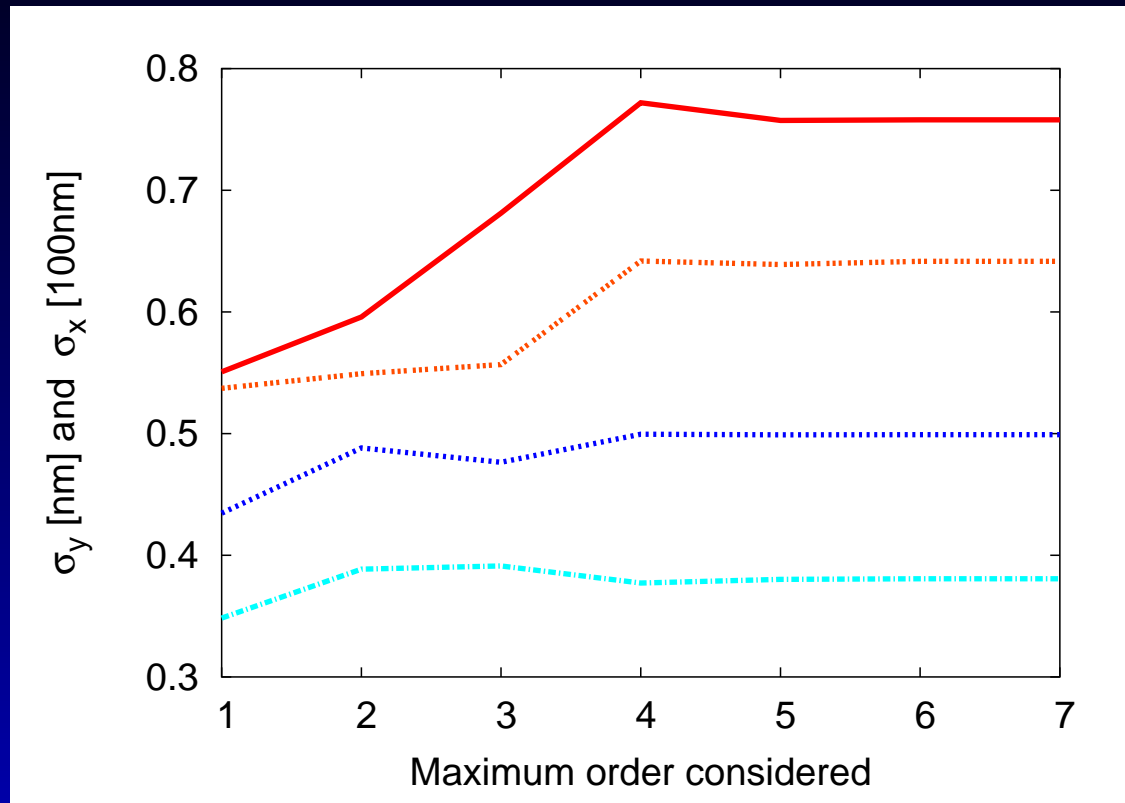


With a total of 21 parameters  $\sigma_x$  and  $\sigma_y$  are reduced by 16% and 18% respectively.

Luminosity increase of 17%

( $\beta_x^* = 6.3\text{mm} \approx -10\%$ ,  $\beta_y^* = 86.8\mu\text{m} \approx -6\%$ )

# Pushing the envelope



With a total of 21 parameters  $\sigma_x$  and  $\sigma_y$  are reduced by 24% and 15% respectively.

Luminosity increase of 35%

( $\beta_x^* = 5.2\text{mm} \approx -26\%$ ,  $\beta_y^* = 86.5\mu\text{m} \approx -5\%$ )

# What about the usable luminosity?

$L_{1\%}$  is the luminosity of the particles with an energy larger than  $0.99 \times E_{peak}$ , the *usable* luminosity.

<i>Case</i>	$-\frac{\Delta\sigma_x}{\sigma_x}$	$-\frac{\Delta\sigma_y}{\sigma_y}$	$\frac{\Delta L_{tot}}{L_{tot}}$	$\frac{\Delta L_{1\%}}{L_{1\%}}$	$\frac{L_{1\%}}{L_{tot}}$
Nominal	0	0	0	0	35
Only sexts	12	0	4	5	35
Sexts, oct & decs	12	17	5	13	38
Full I	16	18	17	14	34
Full II	24	15	35	15	30

(All numbers are percent)

$\Rightarrow$  Looks like  $L_{1\%}$  is mostly affected by  $\sigma_y$

# Outlook

- Use more parameters and evaluate different optics
- Consider synchrotron radiation
- Optimize the full BDS at the same time (the non-linear collimation already optimized independently, see Javier's talk)
- Improve the model with  $\langle xy \rangle$ ,  $\langle x^3 \rangle$ ,  $\langle y^4 \rangle$ , ... ?