

Dark Matter from Exponential Growth

Originally proposed as Pandemic Dark Matter

Based on Phys. Rev. Lett. 127 (2021) 191802 and 2206.10630

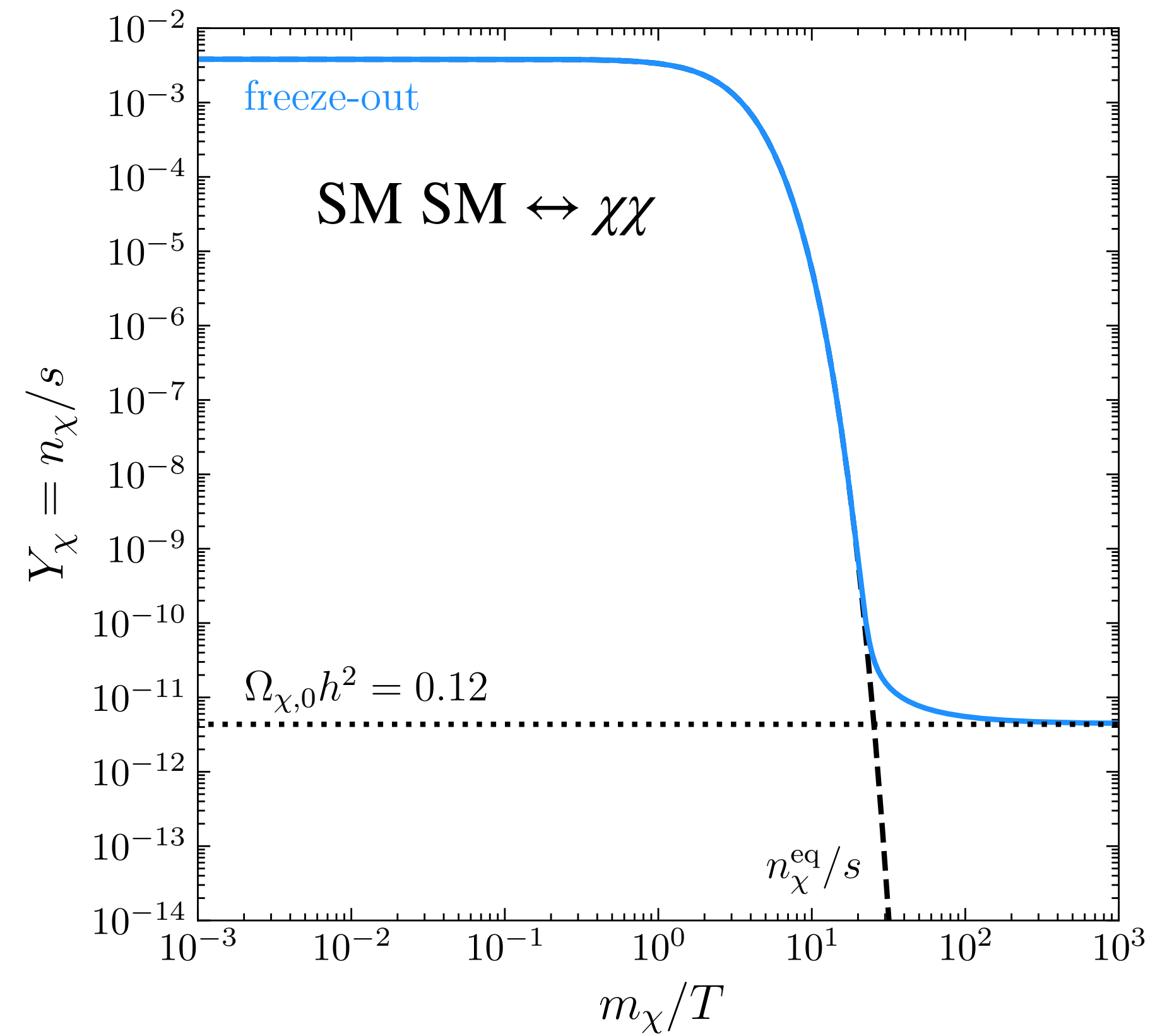
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In collaboration with T. Bringmann, M. Hufnagel, J. Kersten, J. T. Ruderman, and
K. Schmidt-Hoberg

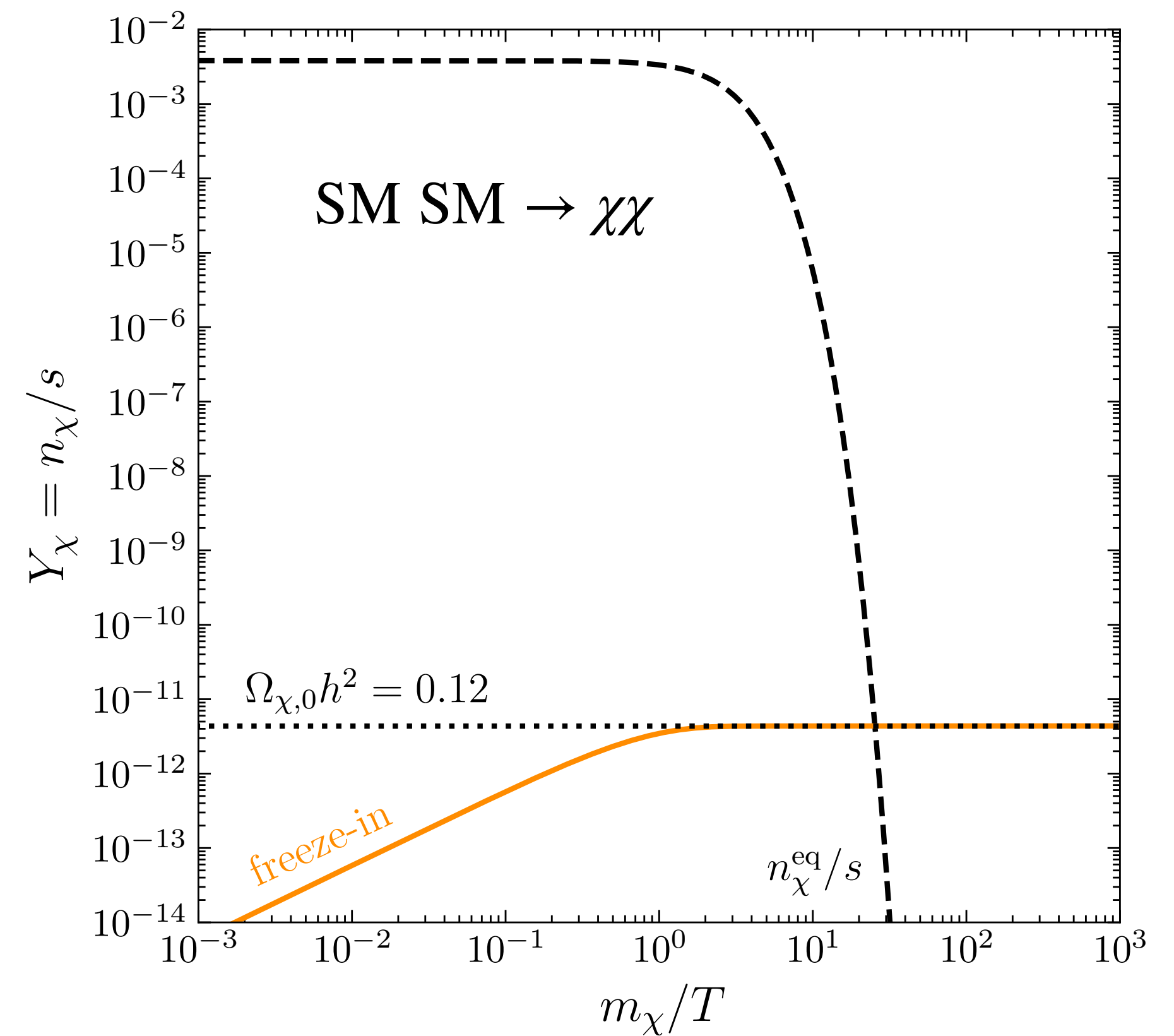
IDM 2022
18 July 2022



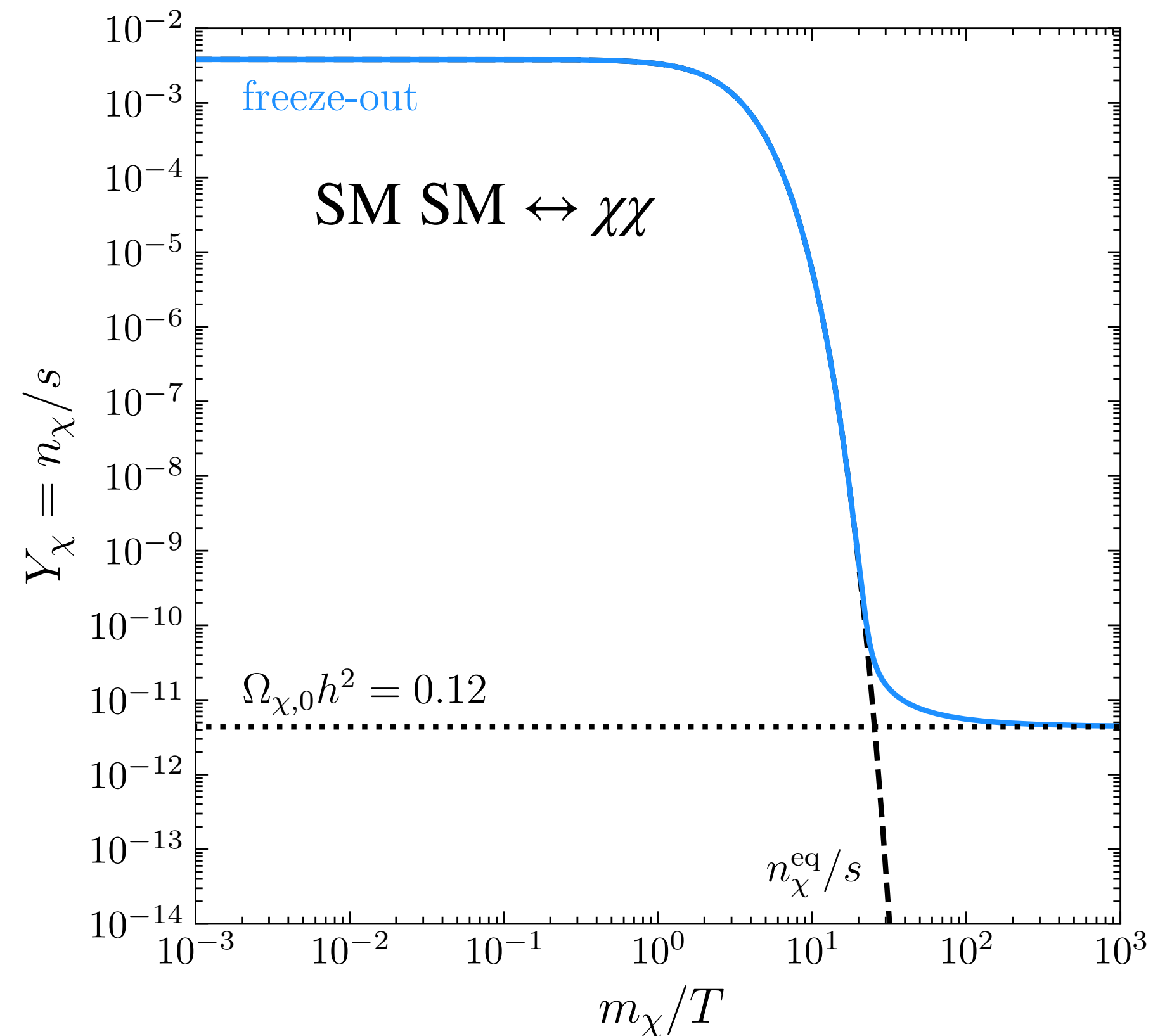
Thermal



Non-Thermal



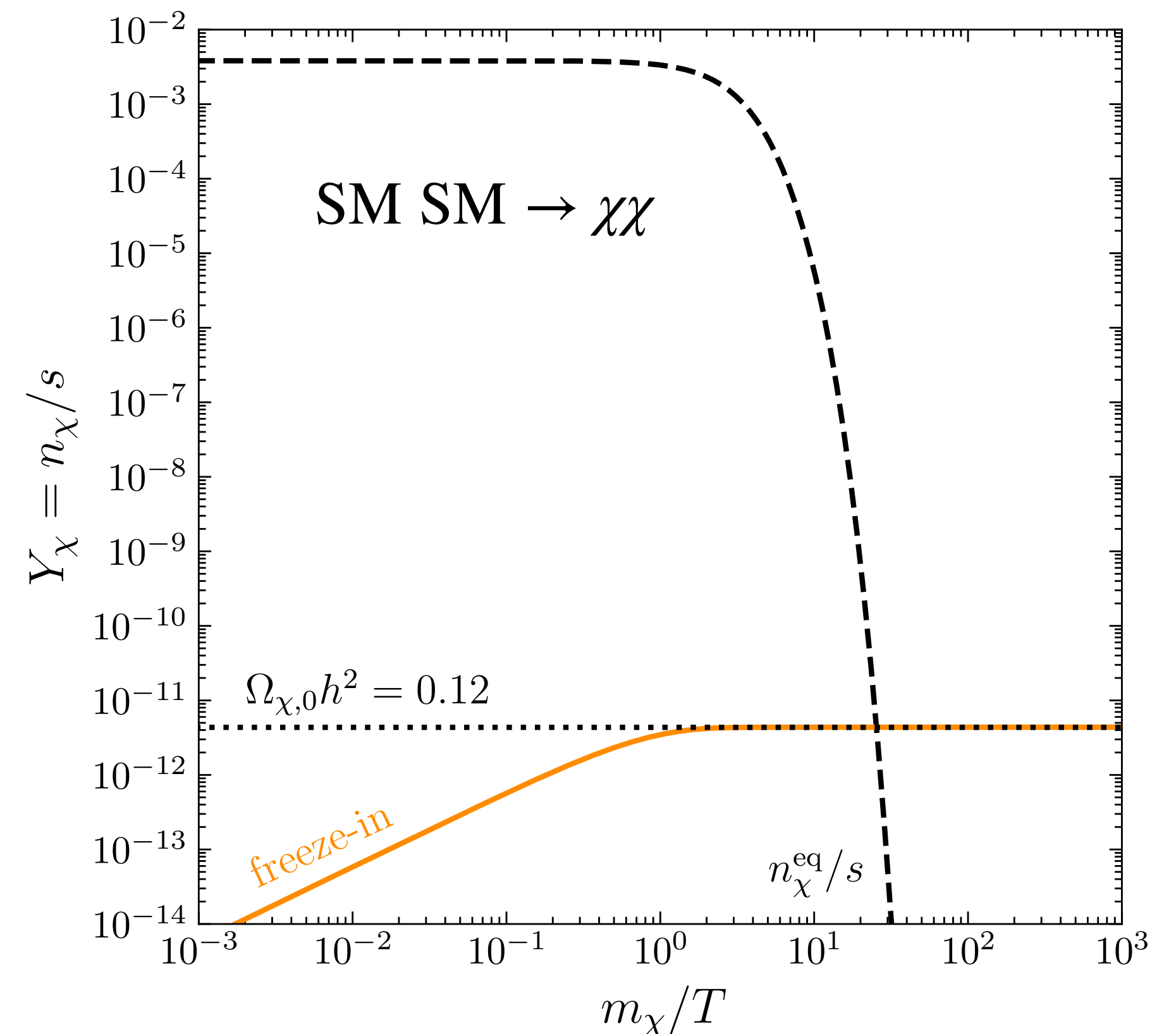
Thermal



Many variants of freeze-out:

- Semi-annihilations
- Hidden sector
- Cannibal DM
- Forbidden DM
- ...

Non-Thermal



Less variants for freeze-in

Dark Matter from Exponential Growth

Bringmann, PFD et al. 2103.16572

Hryczuk, Laletin 2104.05684

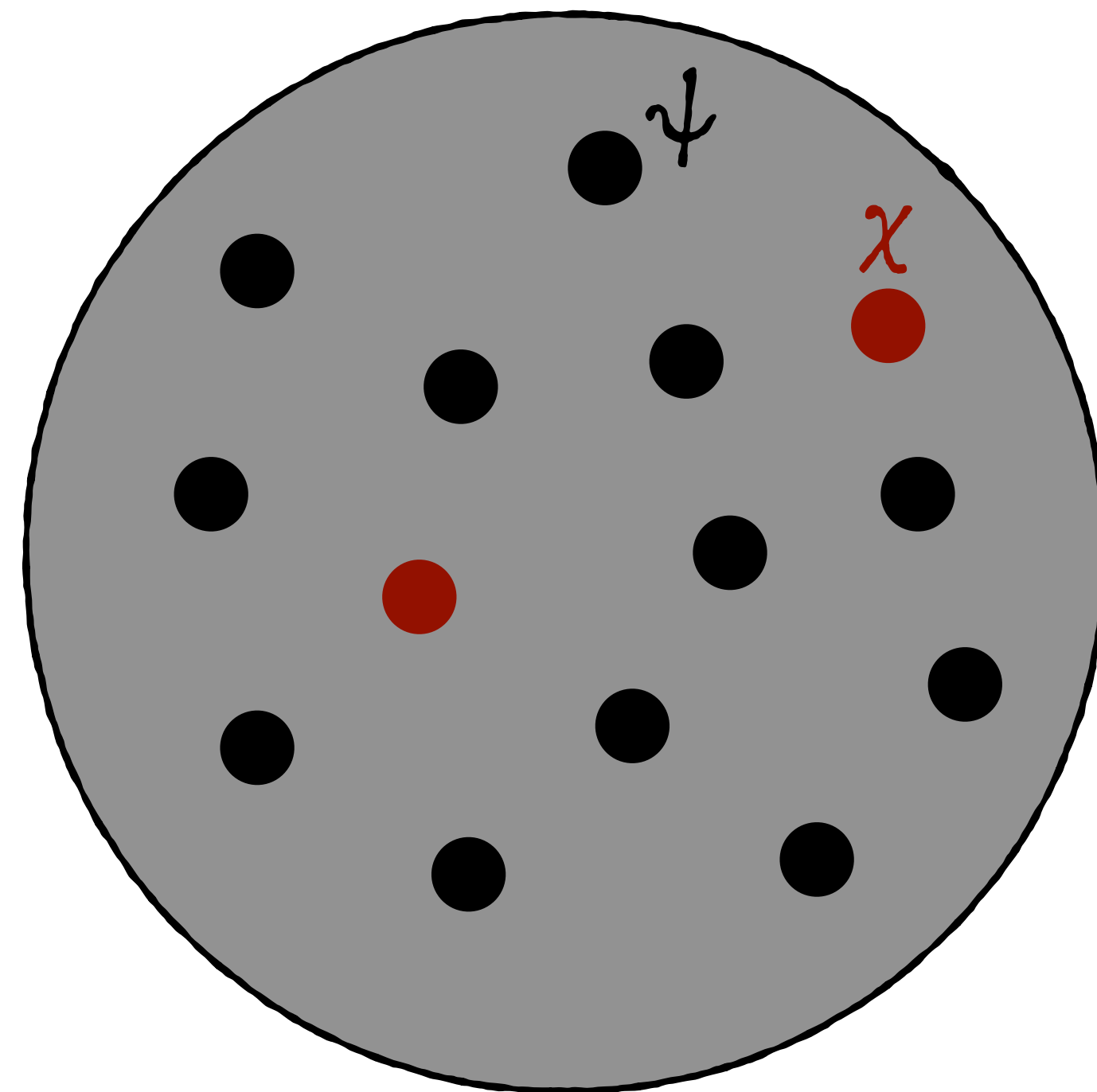
Very interesting application to sterile neutrinos

Bringmann, PFD et al. 2206.10630

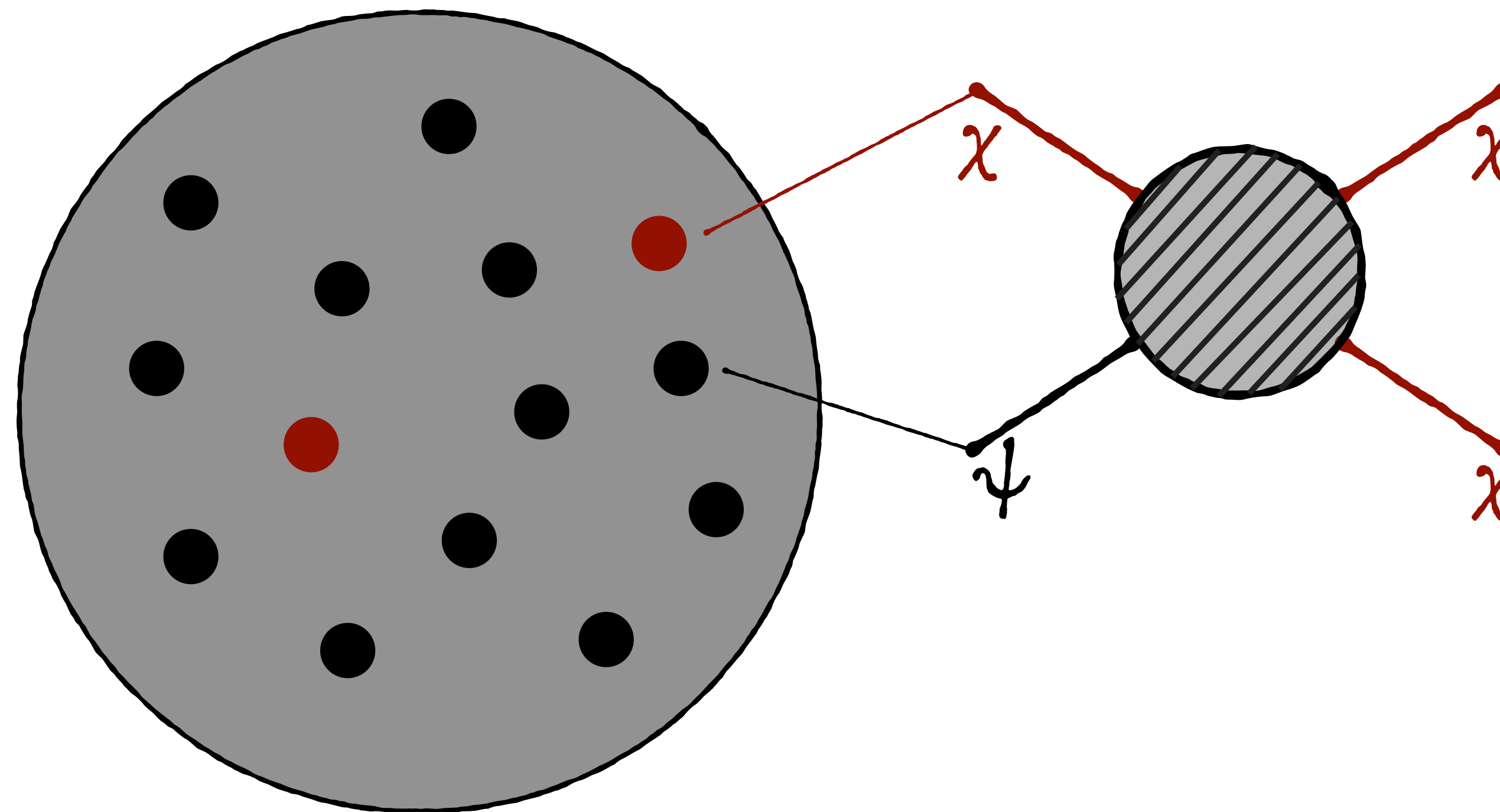
Outline

- Production by transformation
- Evolution of the DM abundance
- Phase diagram
- A new life for sterile neutrino DM after the pandemic (exponential growth)

Production by transformation

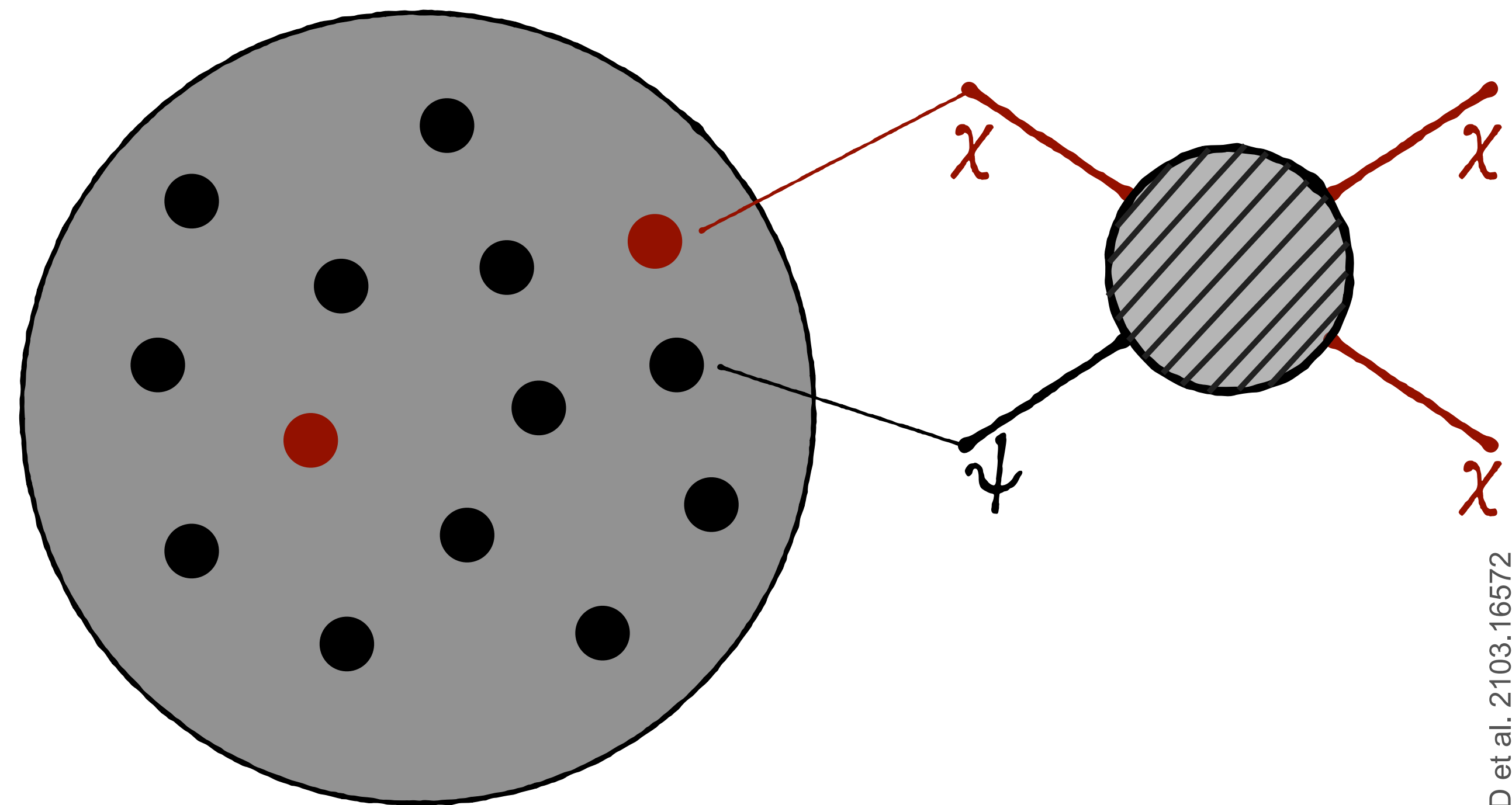


Production by transformation



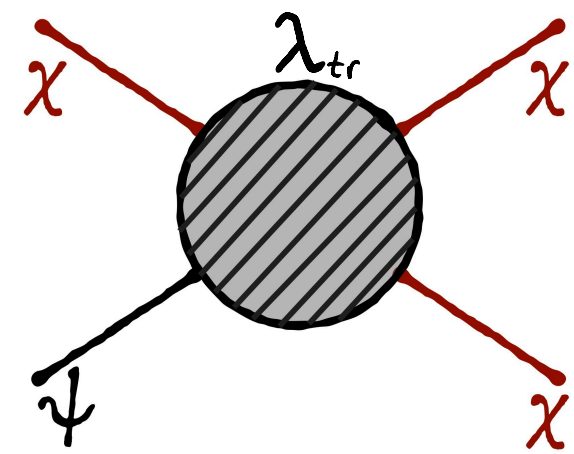
Production by transformation

- $\dot{n}_\chi + 3Hn_\chi = \langle \sigma v \rangle_{\text{tr}} n_\psi^{\text{eq}} n_\chi$
- $Y_\chi(x_\psi) \equiv n_\chi/s \simeq Y_\chi^0 \exp\left(3 \int_{x_\psi^0}^{x_\psi} \frac{dx}{x} R(x)\right)$
- $R(x) = \frac{n_\psi^{\text{eq}} \langle \sigma v \rangle_{\text{tr}}}{3H}$: # of transformations of DM particle per Hubble time
- \rightarrow Phase of exponential production
- Shutoff by kinematical or Boltzmann suppression
- First: constant matrix element for simplicity

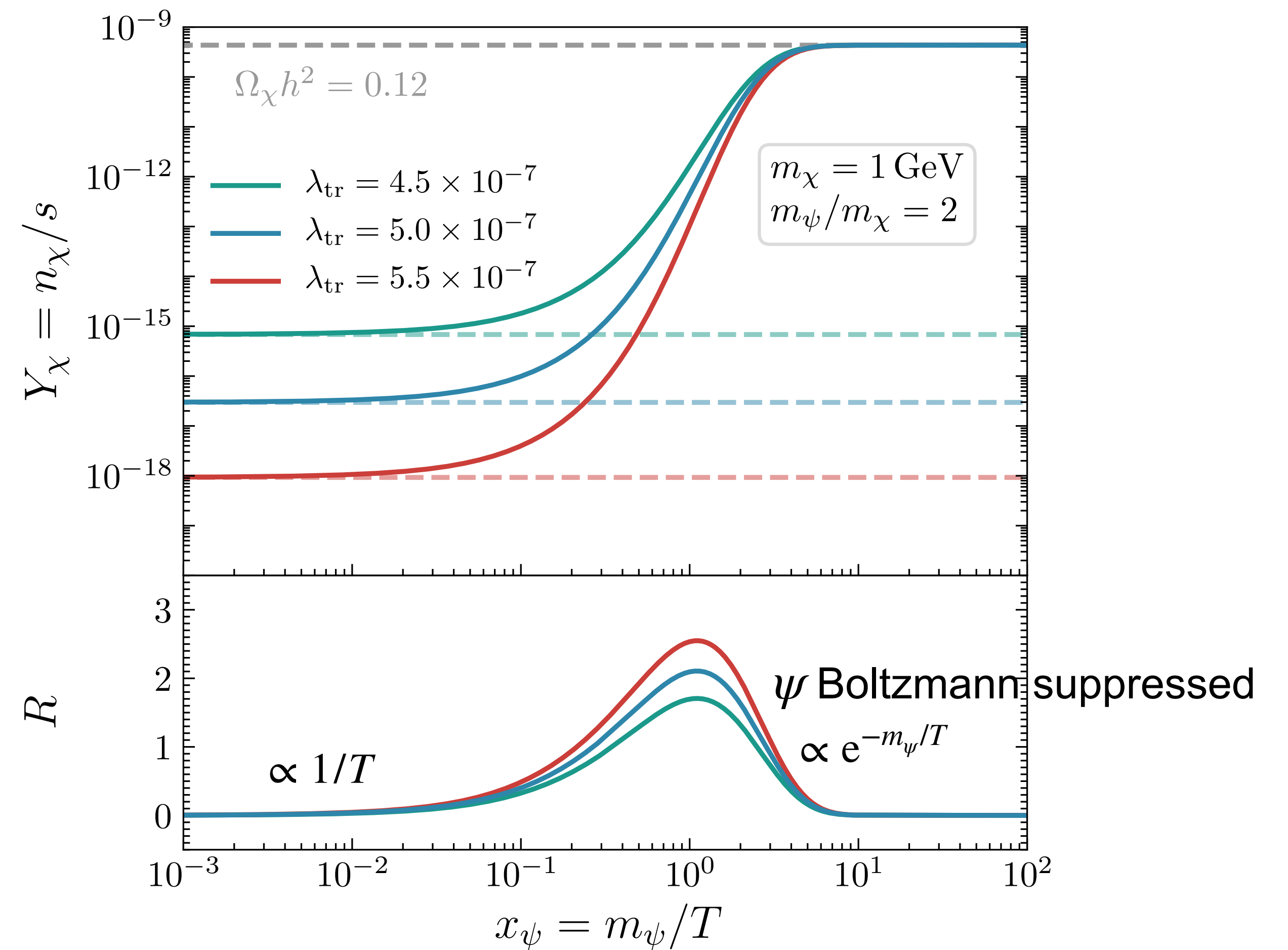


Evolution of DM abundance

Fixed initial abundance

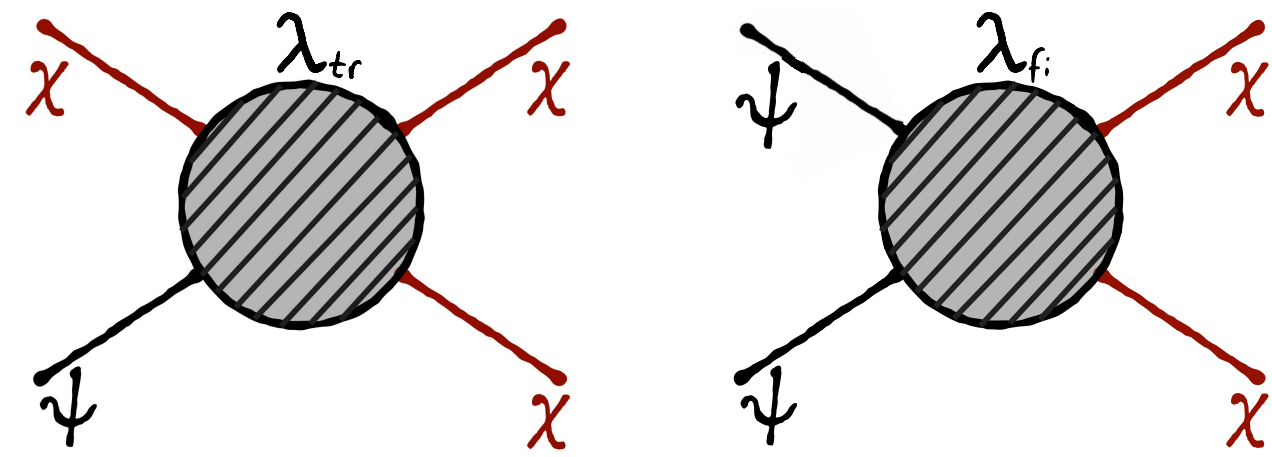


$$\dot{n}_\chi + 3Hn_\chi = \langle \sigma v \rangle_{\text{tr}} n_\psi^{\text{eq}} n_\chi$$

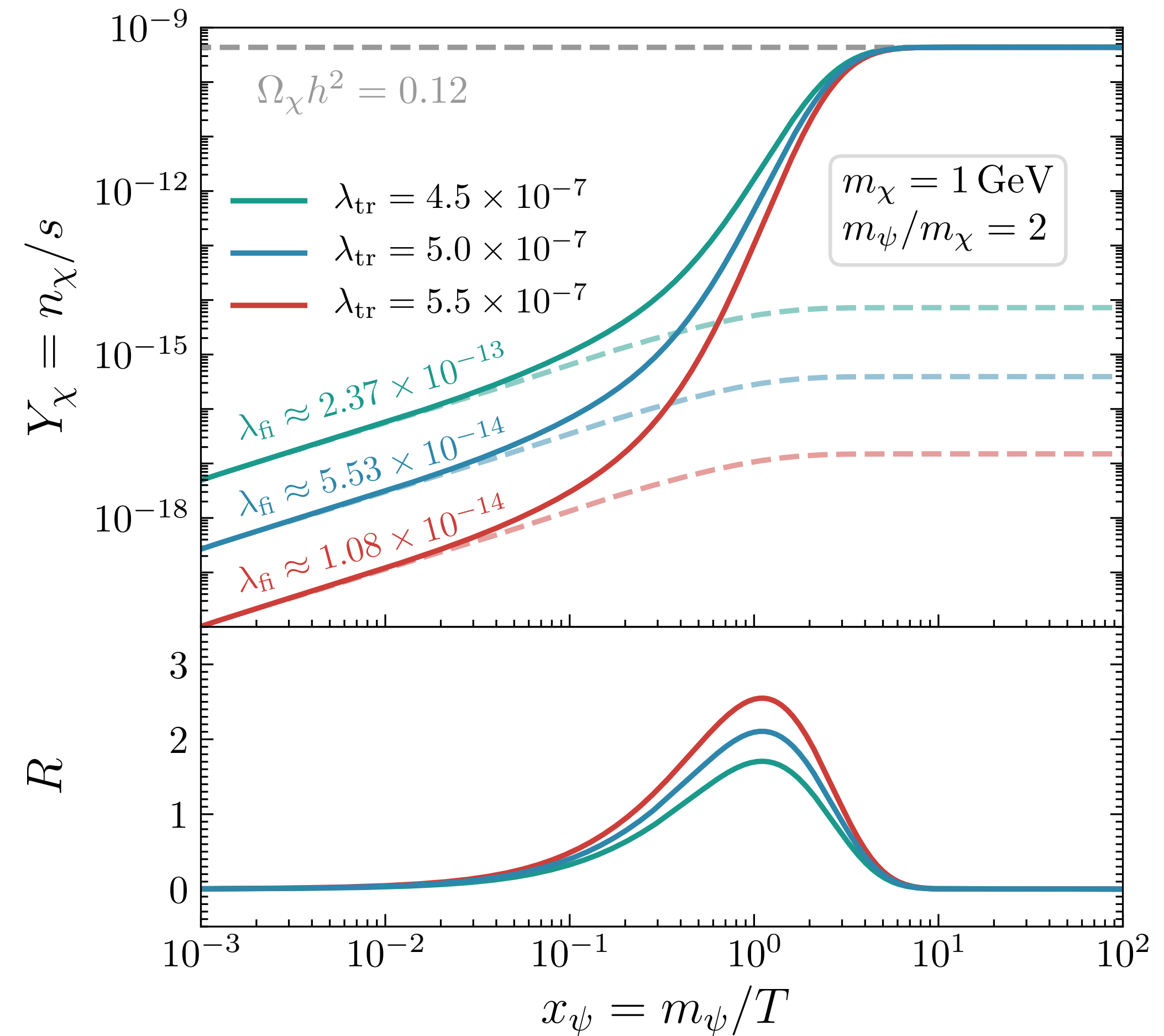


Evolution of DM abundance

Initial abundance from freeze-in



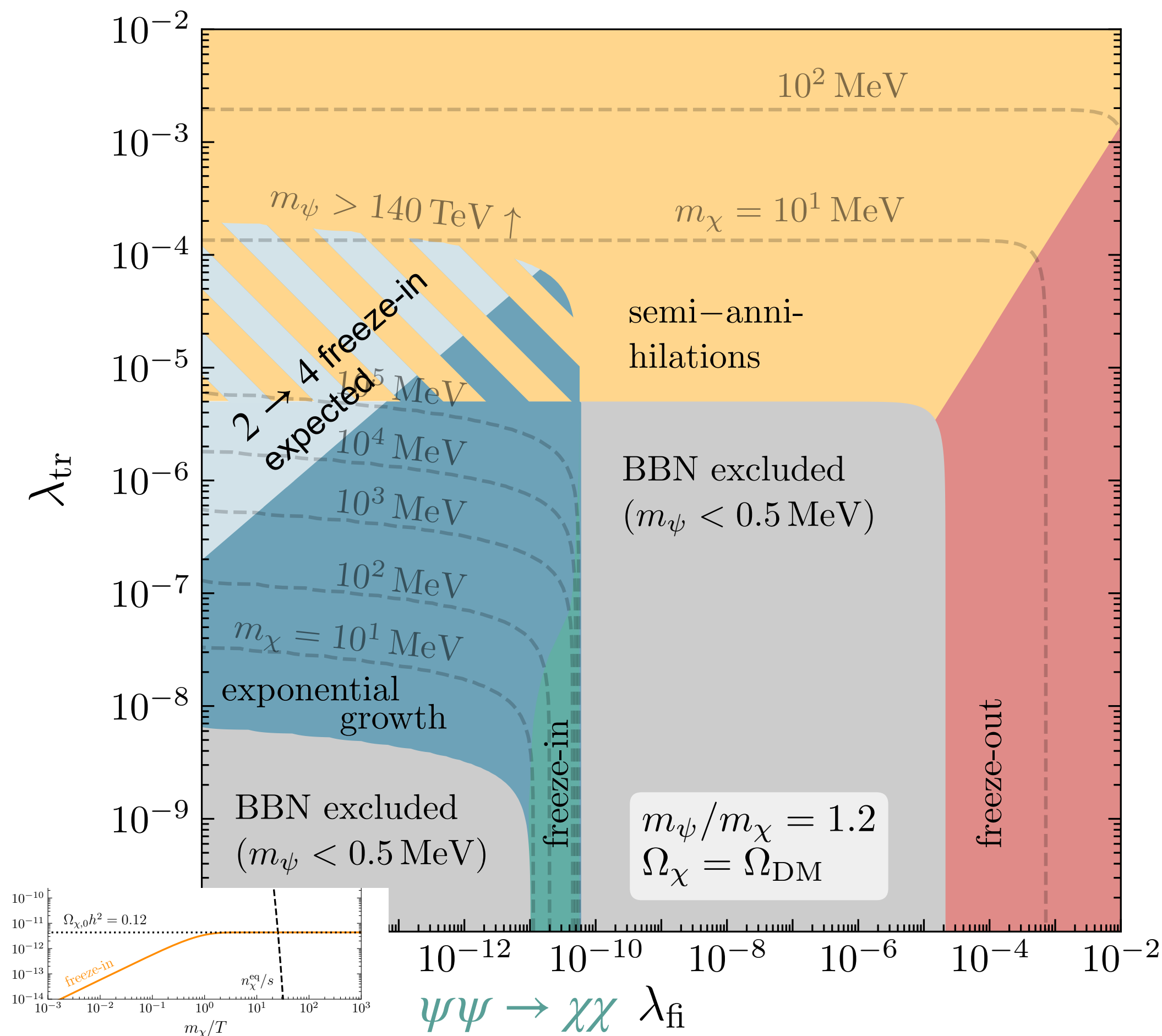
$$\dot{n}_\chi + 3Hn_\chi = \langle \sigma v \rangle_{\text{tr}} n_\psi^{\text{eq}} n_\chi + \langle \sigma v \rangle_{\text{fi}} (n_\psi^{\text{eq}})^2$$



Phase diagram

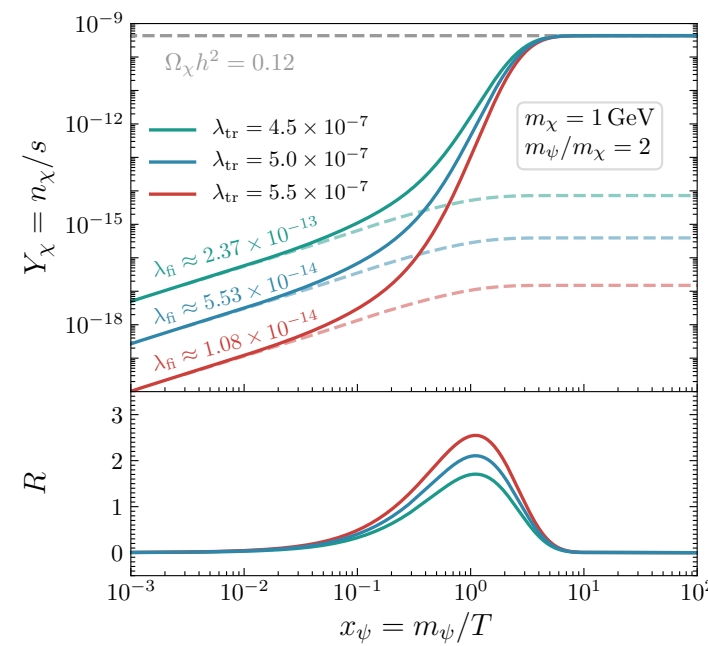
$$\dot{n}_\chi + 3Hn_\chi =$$

$$\langle \sigma v \rangle_{\text{fi}} [(n_\psi^{\text{eq}})^2]$$

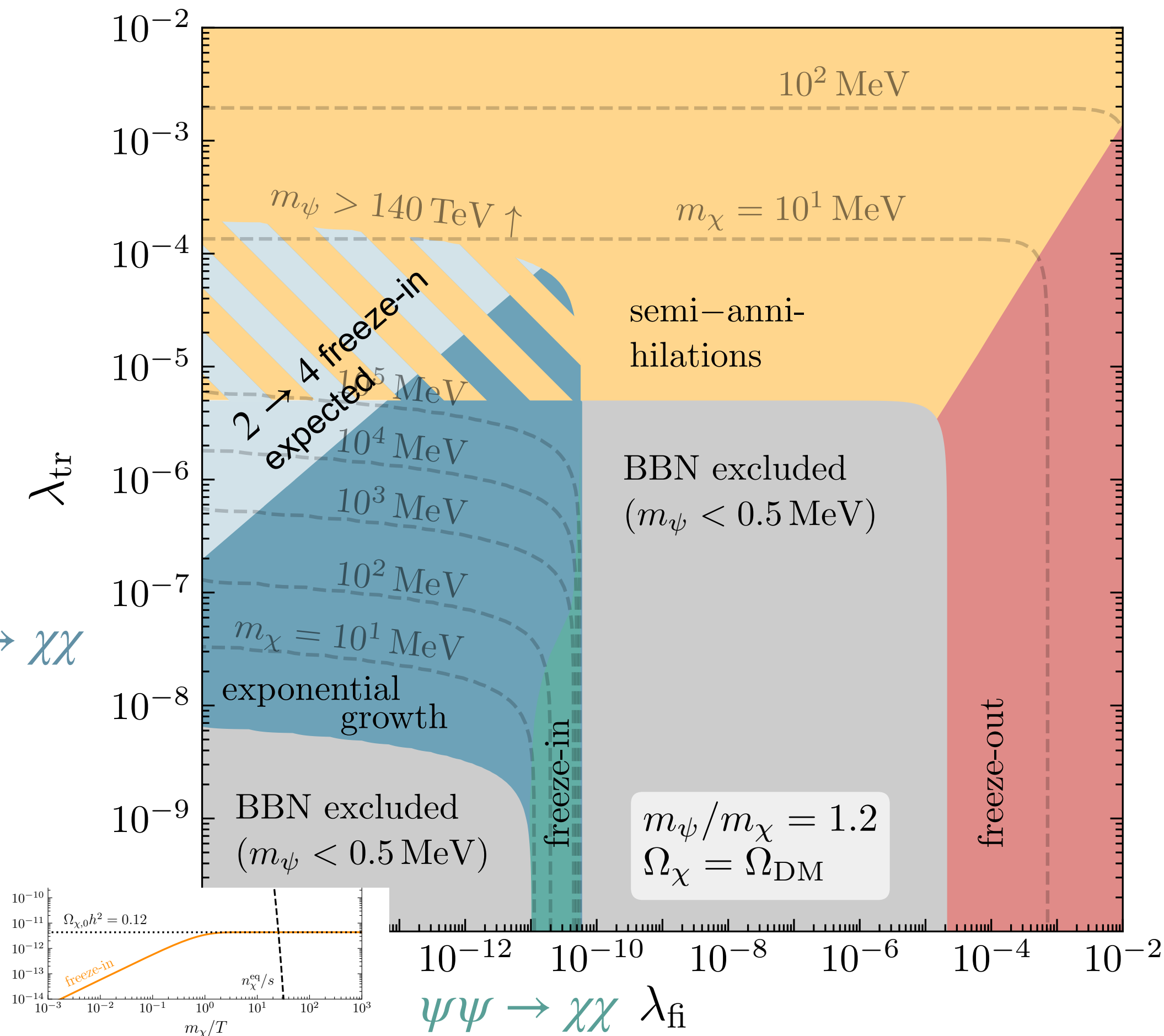


Phase diagram

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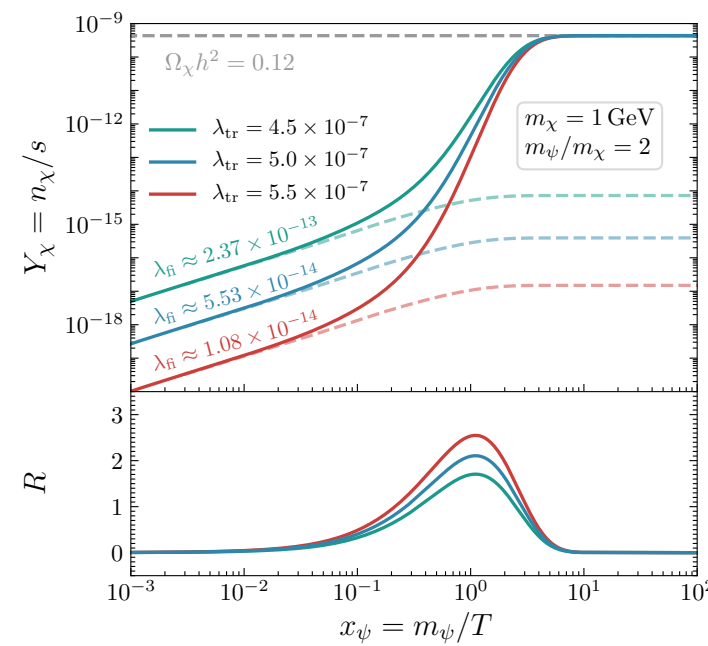
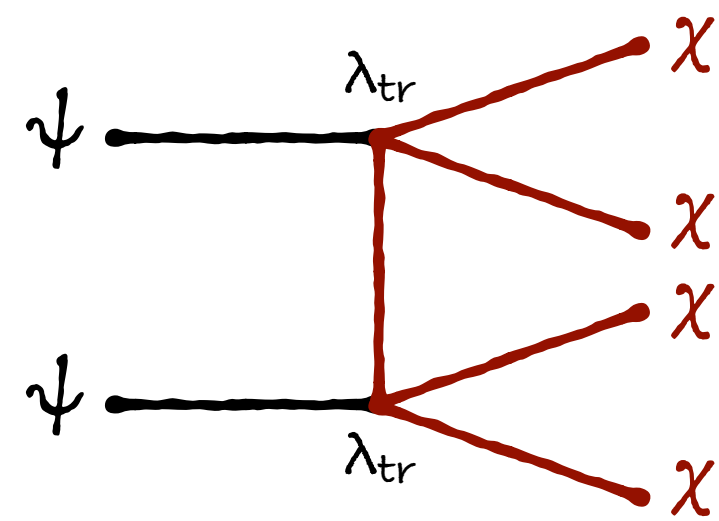


Generally: $\lambda_{\text{fi}} \ll \lambda_{\text{tr}}$

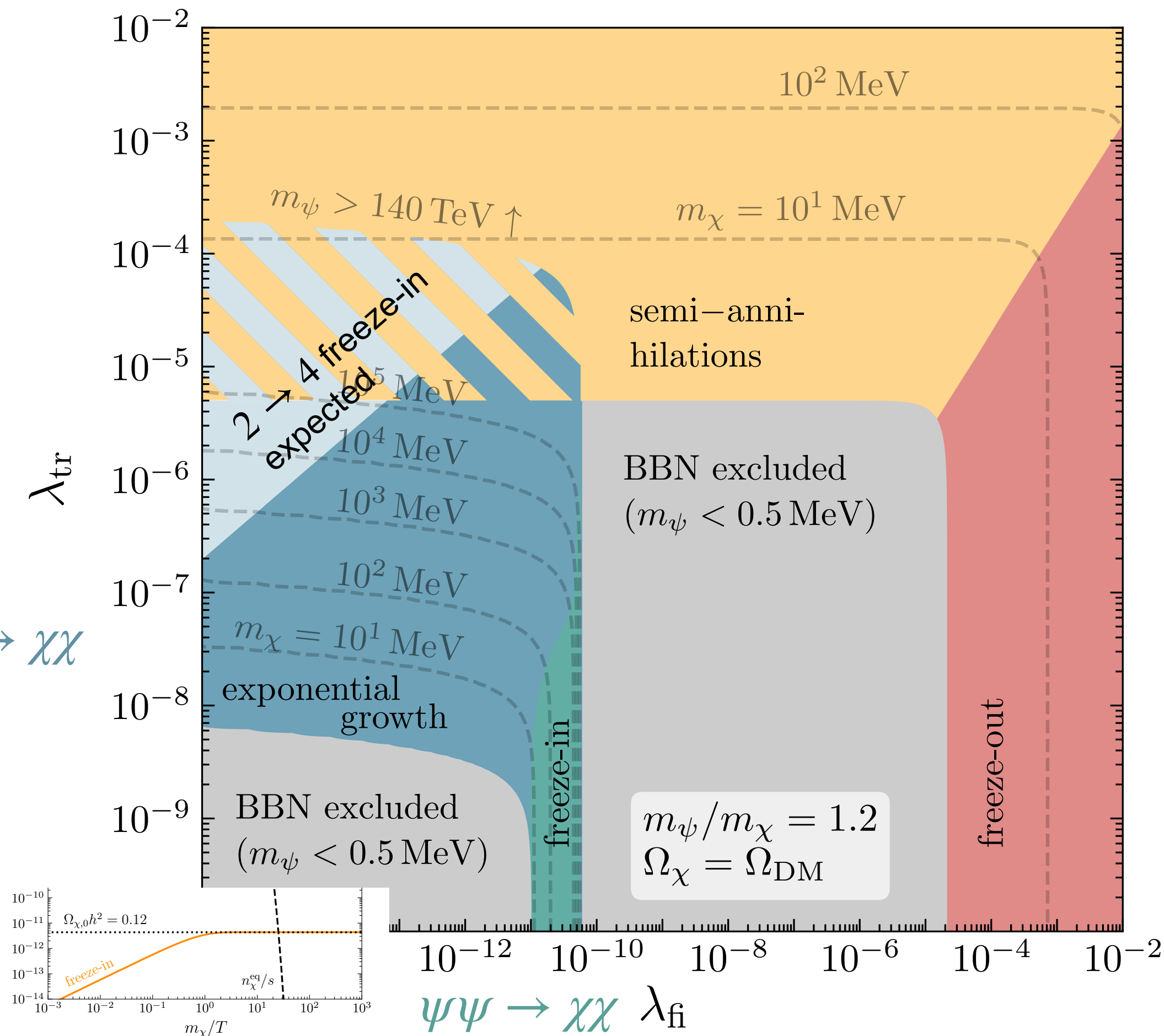


Phase diagram

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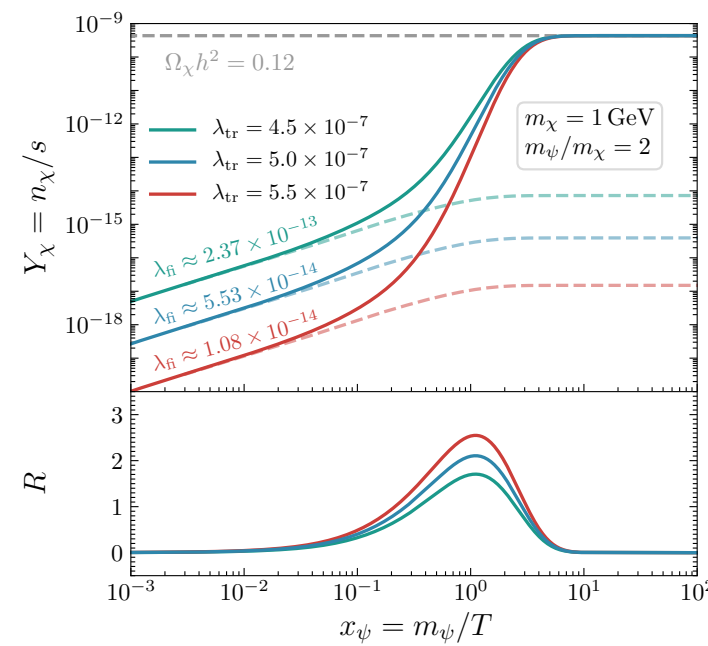
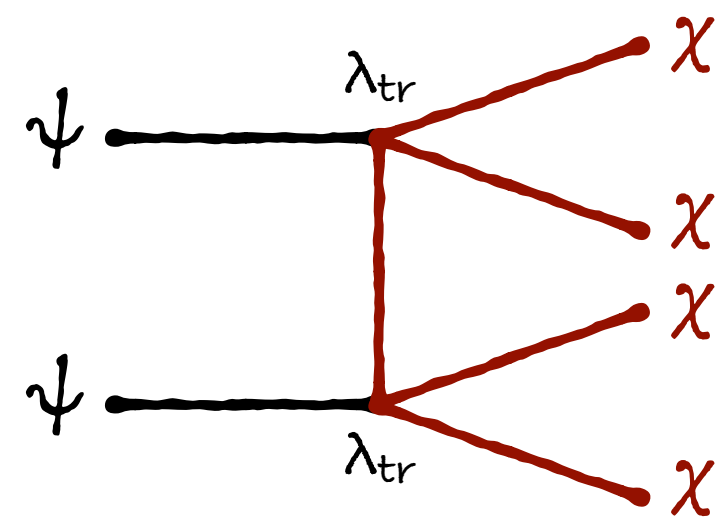
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Phase diagram

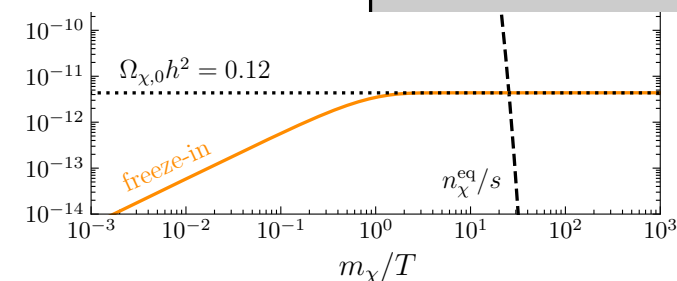
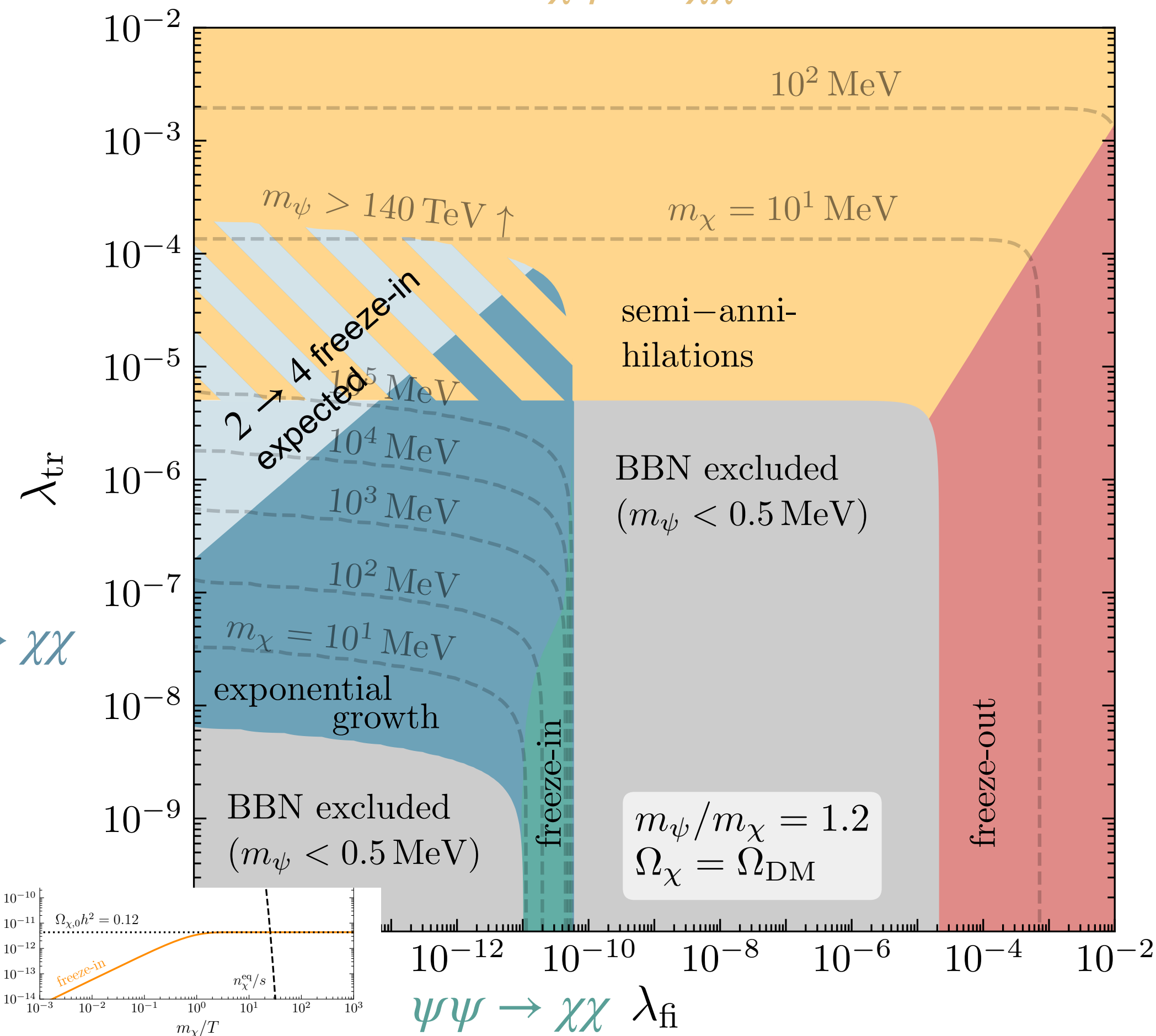
$$\dot{n}_\chi + 3Hn_\chi = \langle\sigma v\rangle_{\text{tr}} [n_\psi^{\text{eq}}n_\chi - n_\chi^2 n_\psi^{\text{eq}}/n_\chi^{\text{eq}}] + \langle\sigma v\rangle_{\text{fi}} [(n_\psi^{\text{eq}})^2 - (n_\chi n_\psi^{\text{eq}}/n_\chi^{\text{eq}})^2]$$

$\chi\psi \leftrightarrow \chi\chi$

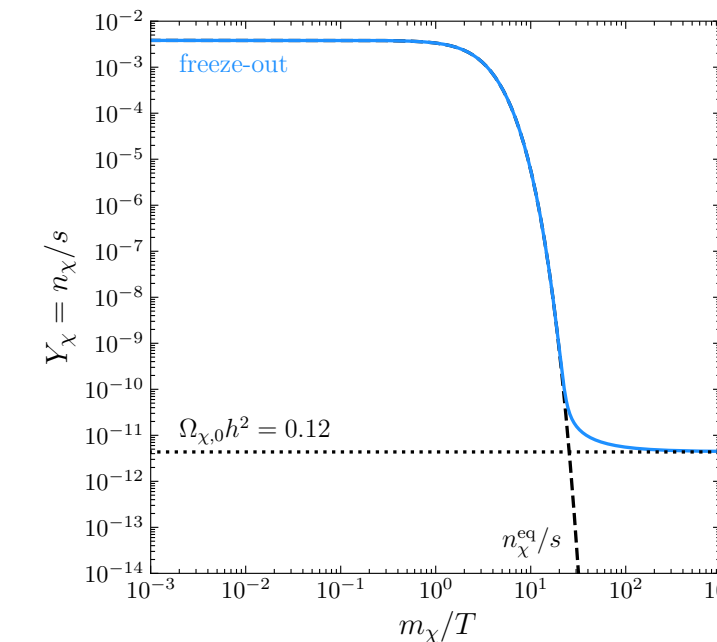


$\chi\psi \rightarrow \chi\chi$

Generally: $\lambda_{\text{fi}} \ll \lambda_{\text{tr}}$



$\psi\psi \rightarrow \chi\chi$ λ_{fi}



$\psi\psi \leftrightarrow \chi\chi$



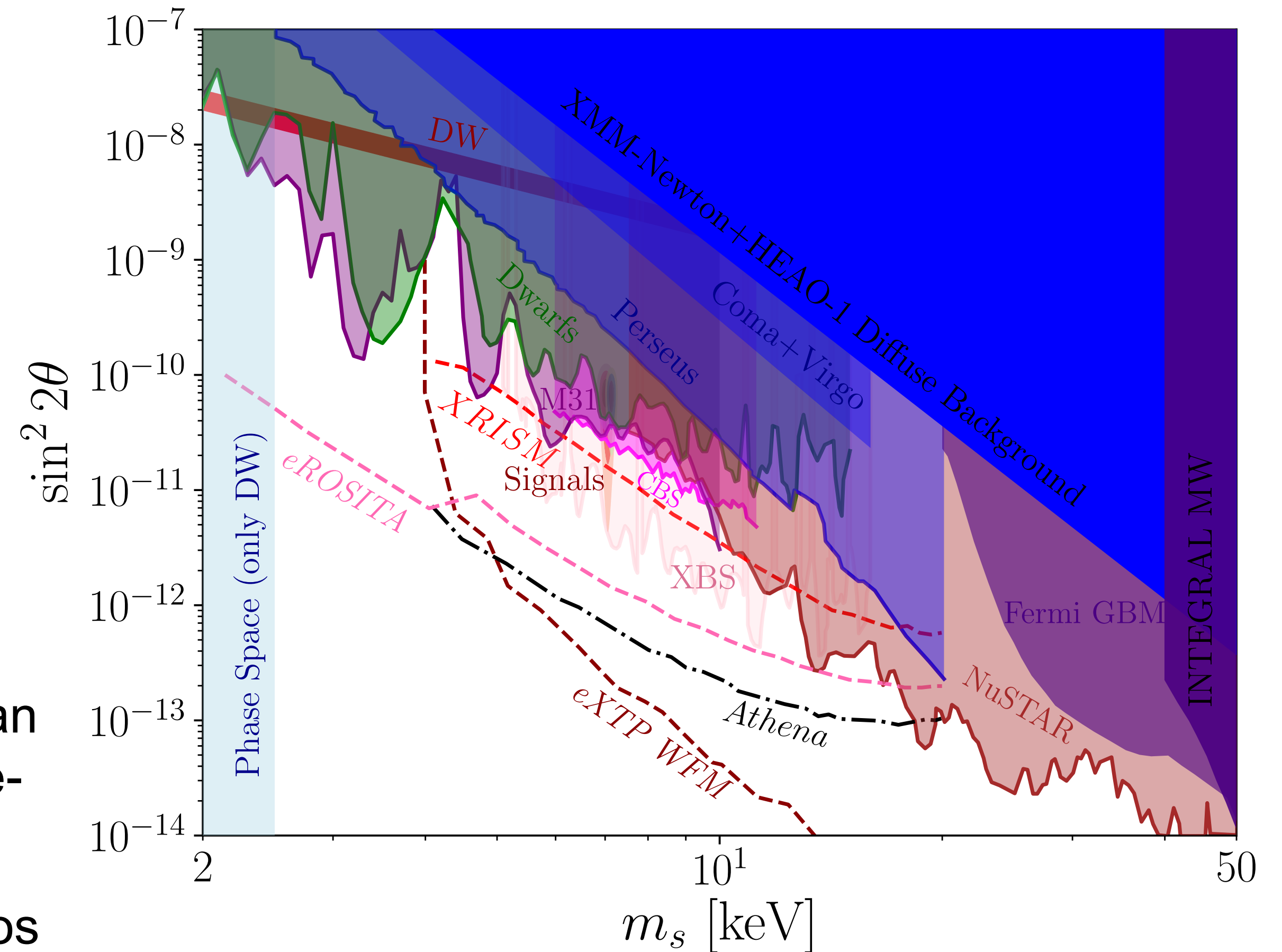
Coupling hierarchy through mass mixing

- Convenient way to realize $\lambda_{\text{fi}} \ll \lambda_{\text{tr}}$: Two fermions with tiny mass mixing angle θ , only one (mostly χ) interacts with some mediator ϕ via Yukawa coupling
- After mass diagonalization:
 - $\bar{\chi}\chi$ vertices $\propto \cos^2 \theta \sim 1$
 - $\bar{\psi}\chi$ vertices $\propto \cos \theta \sin \theta \sim \theta$
 - $\bar{\psi}\psi$ vertices $\propto \sin^2 \theta \sim \theta^2$
- Transformation amplitude $\propto \theta$
- Freeze-in amplitude $\propto \theta^2$



A new life for sterile neutrino DM after exp. growth

- What if ψ further is in the SM?
- Sterile neutrino ($\chi = \nu_s$), mass-mixing with active neutrino ($\psi = \nu_\alpha$)
- Yukawa coupling between mediator ϕ and ν_s in flavor-space:
 - $\mathcal{L}_{\text{int}} \supset \frac{y}{2} \phi \bar{\nu}_s^c \nu_s + \text{h.c.}$
 - $\rightarrow \frac{y}{2} \phi (\cos^2 \theta \bar{\nu}_s^c \nu_s - \sin(2\theta) \bar{\nu}_\alpha^c \nu_s + \sin^2 \theta \bar{\nu}_\alpha^c \nu_\alpha) + \text{h.c.}$
- Allows for sterile neutrino DM at θ much smaller than Dodelson-Widrow mechanism (production by active-sterile oscillations) in a very simple model
- Generally present for self-interacting sterile neutrinos

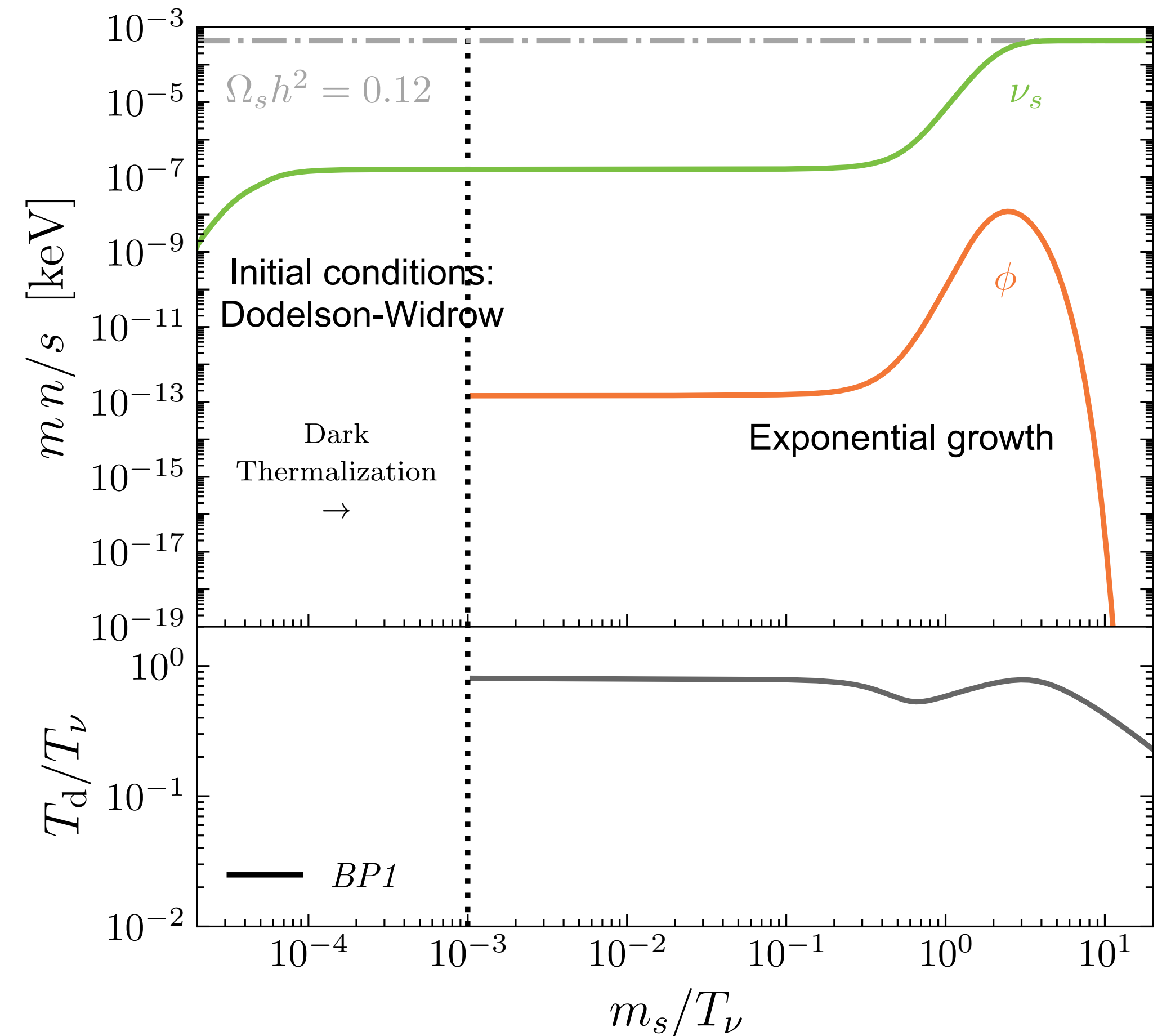
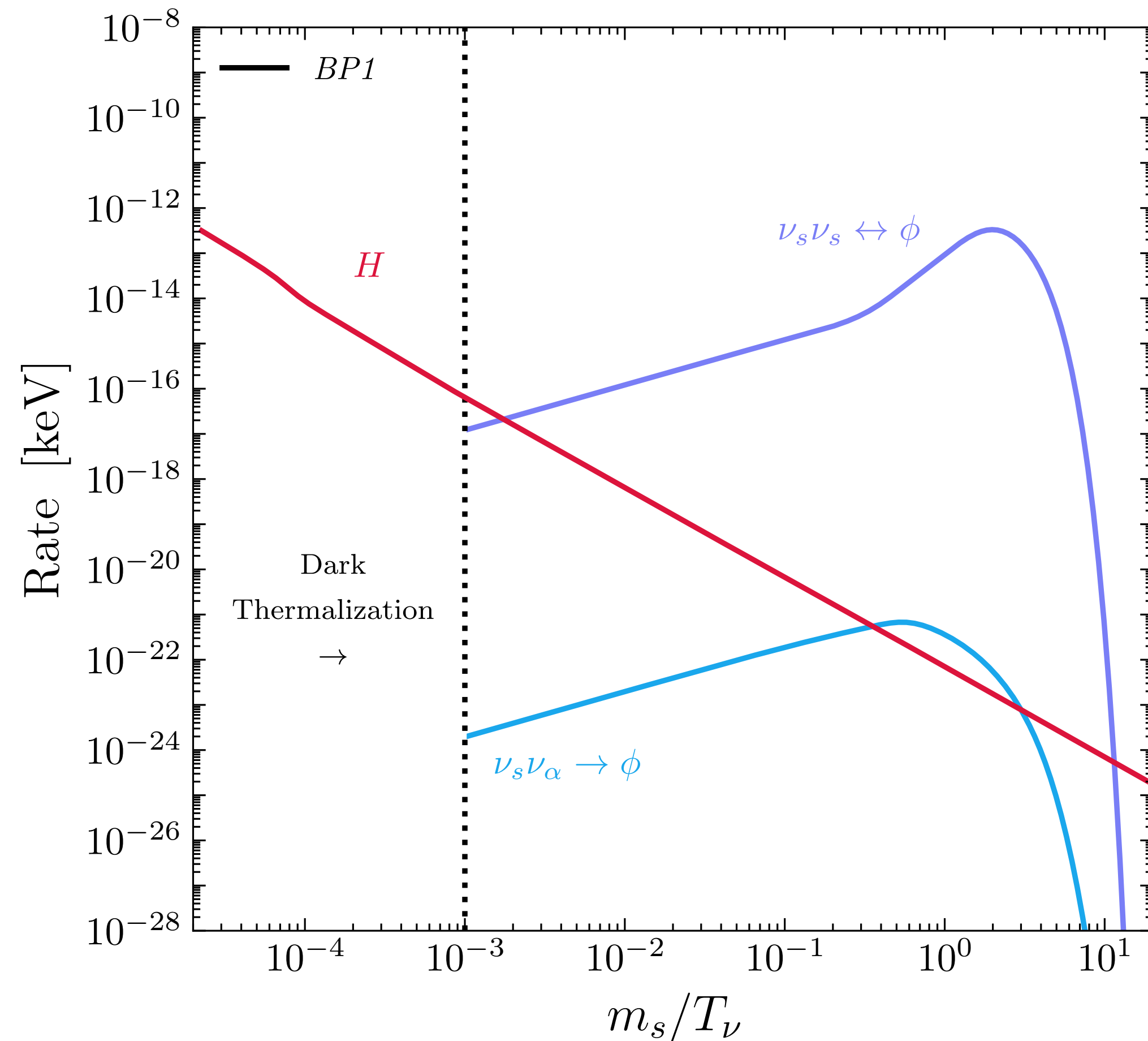


Abazajian et al. 2203.07377



Evolution

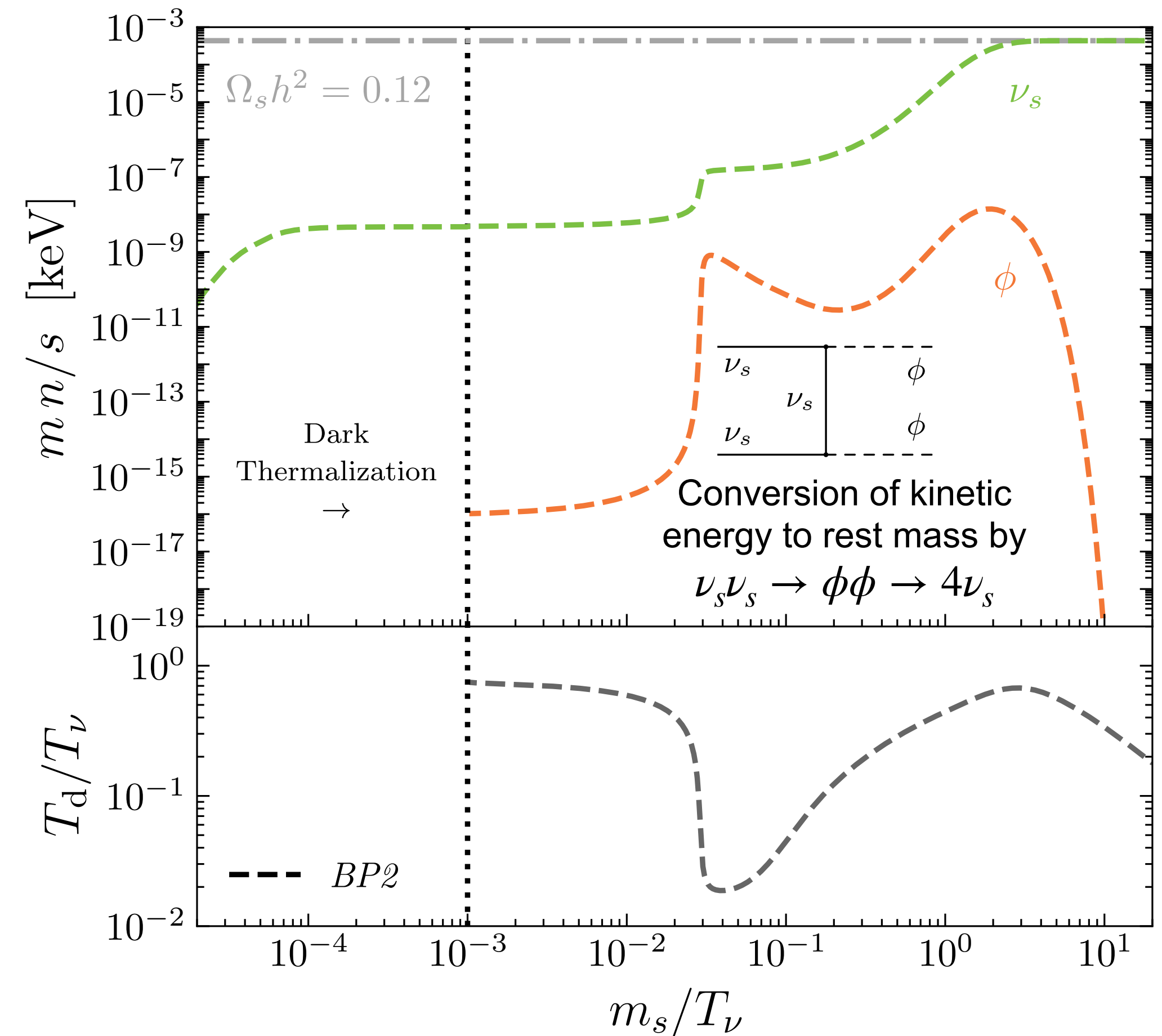
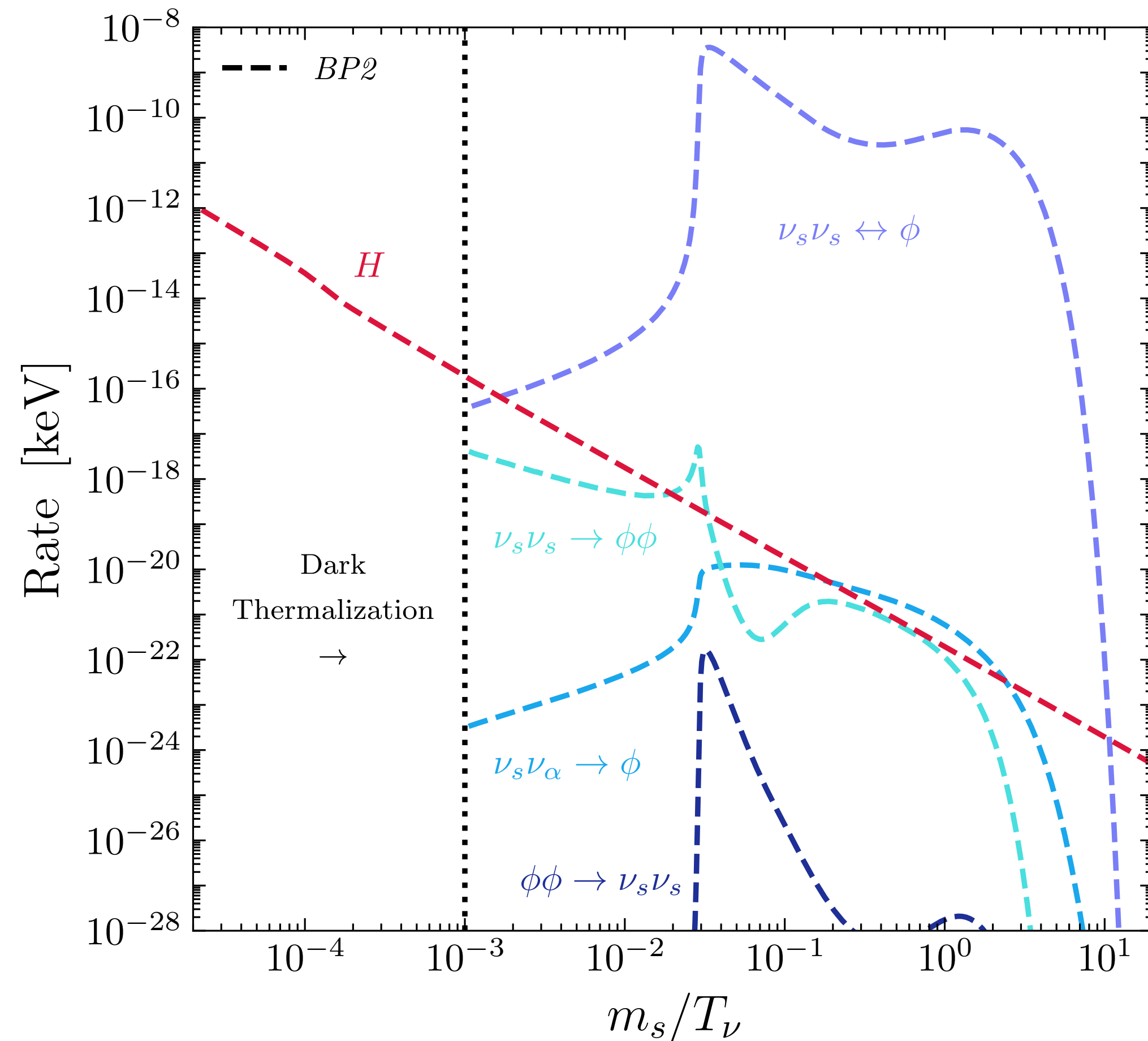
	m_s	m_ϕ	$\sin^2(2\theta)$	y
<i>BP1</i>	12 keV	36 keV	2.5×10^{-13}	1.905×10^{-4}



Evolution

	m_s	m_ϕ	$\sin^2(2\theta)$	y
<i>BP1</i>	12 keV	36 keV	2.5×10^{-13}	1.905×10^{-4}
<i>BP2</i>	20 keV	60 keV	3.0×10^{-15}	1.602×10^{-3}

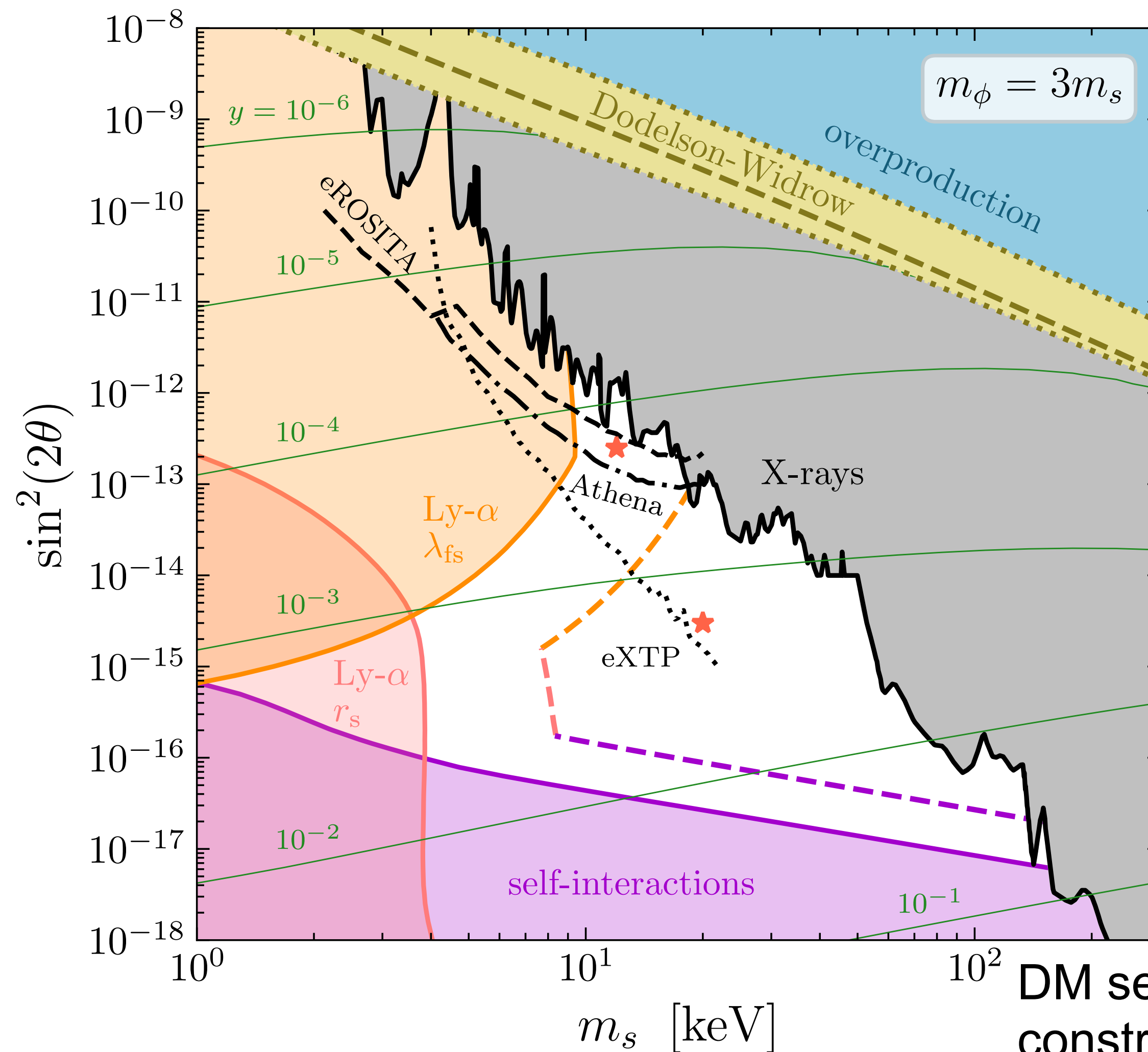
Smaller $\theta \Rightarrow$ larger $y \Rightarrow$ additional processes



Parameter space

Lyman- α forest constraints from suppression of small-scale structure:

- DM self-scatterings before kinetic decoupling
→ structures below sound horizon r_s suppressed
- DM free-streaming after kinetic decoupling
→ structures below free-streaming length λ_{fs} suppressed



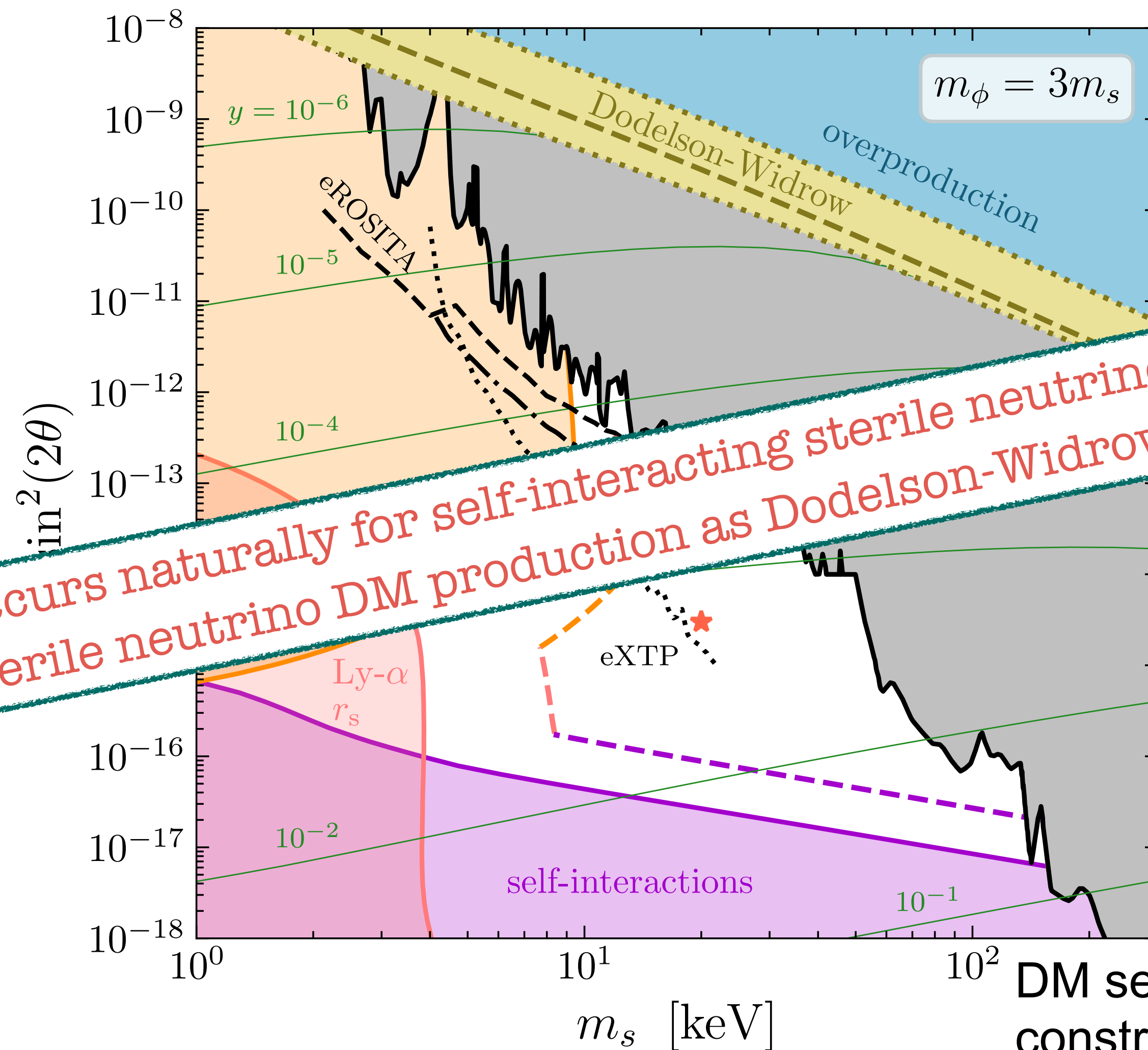
X-ray constraints from DM decays

DM self-interactions constrained by astrophysical observations at late times

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X-ray constraints from DM decays

DM self-interactions constrained by astrophysical observations at late times

Simplest model for sterile neutrino DM production as Dodelson-Widrow scenario is excluded!

Occurs naturally for self-interacting sterile neutrinos!

Conclusions

- New non-thermal DM production mechanism involving exponential growth
- Complements freeze-in and freeze-out scenarios
- Interesting phenomenological consequences
- A new life for sterile neutrino DM after exponential growth
 - Exponential growth regime occurs naturally for self-interacting sterile neutrinos
 - Allows for mixing angle much smaller than in Dodelson-Widrow scenario
 - Simplest allowed model for sterile neutrino DM production as Dodelson-Widrow is excluded
 - Much of parameter space is testable in the foreseeable future



Thank you!



Phenomenological consequences

Higgs portal $\lambda_{h\psi} |H|^2 \psi^2 / 2$

