



The CRESST-III Dark Matter Search

Status and Outlook

Christian Strandhagen | Universität Tübingen | IDM 2022, Vienna



The CRESST Collaboration



COMENIUS
UNIVERSITY
BRATISLAVA



Laboratori Nazionali del Gran Sasso



Technische Universität München



UNIVERSITY OF
OXFORD



MAX-PLANCK-INSTITUT
FÜR PHYSIK



EBERHARD KARLS
UNIVERSITÄT
TÜBINGEN



INSTITUTE OF HIGH ENERGY PHYSICS



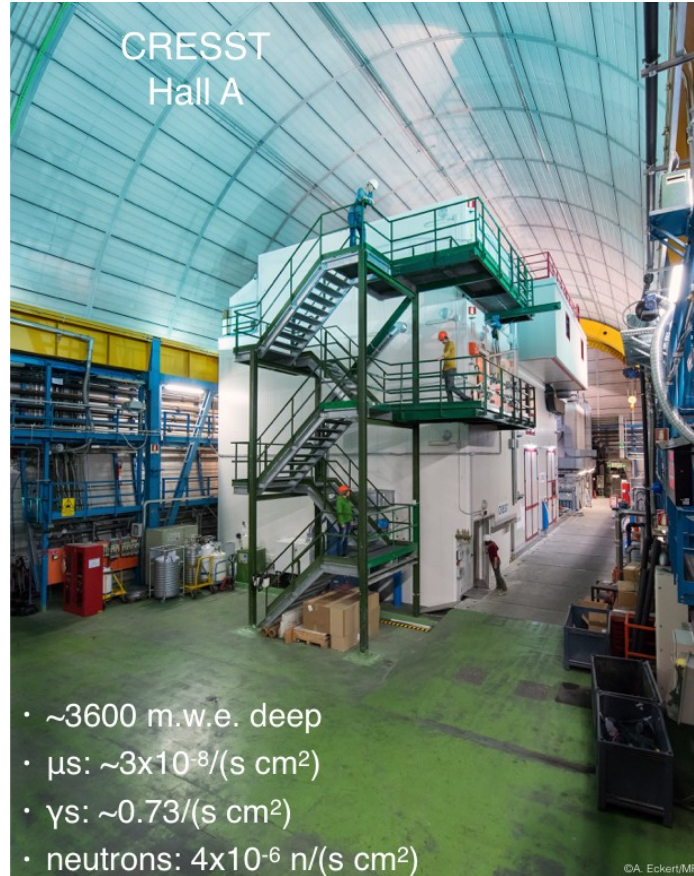
The CRESST Experiment



The CRESST experiment is located in the LNGS underground lab in Italy.



The CRESST Experiment

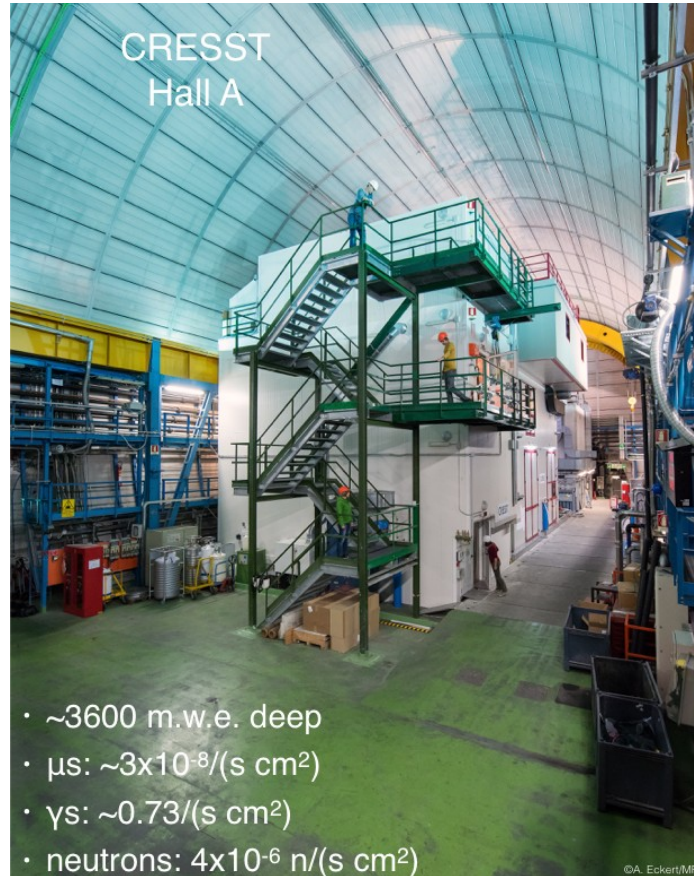
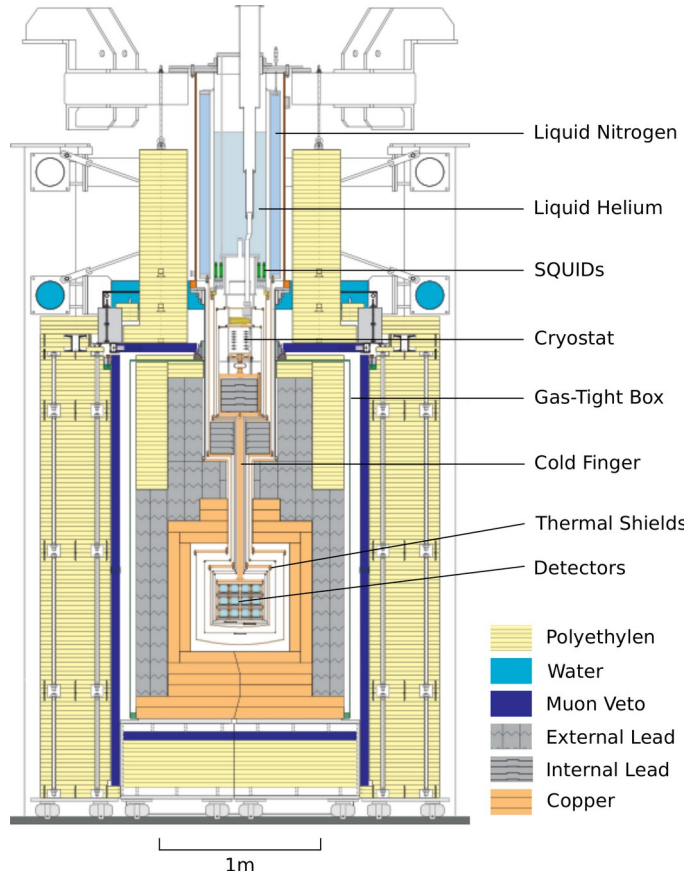


The CRESST experiment is located in the LNGS underground lab in Italy.

At least 1400 m of rock of the Gran Sasso d'Italia (2912 m) shield against cosmic muons.

- ~ 3600 m.w.e. deep
- μ s: $\sim 3 \times 10^{-8} / (\text{s cm}^2)$
- γ s: $\sim 0.73 / (\text{s cm}^2)$
- neutrons: $4 \times 10^{-6} \text{ n} / (\text{s cm}^2)$

The CRESST Experiment

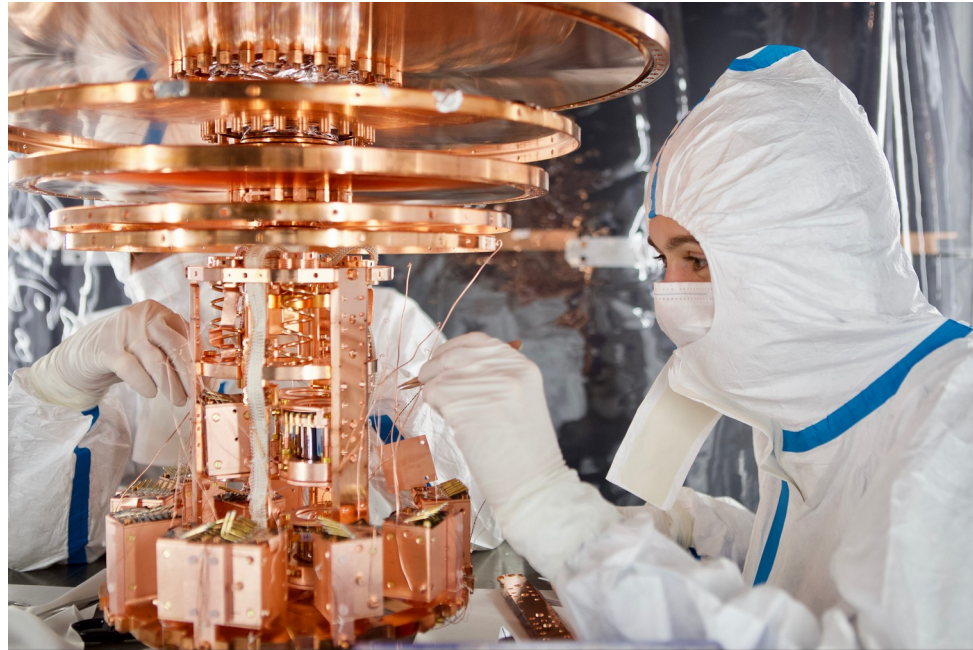
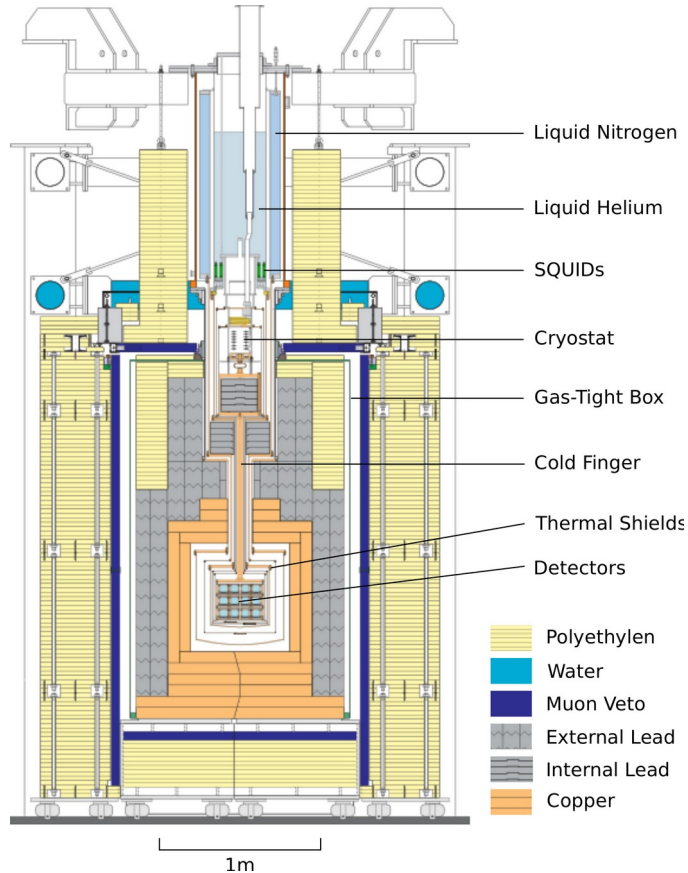


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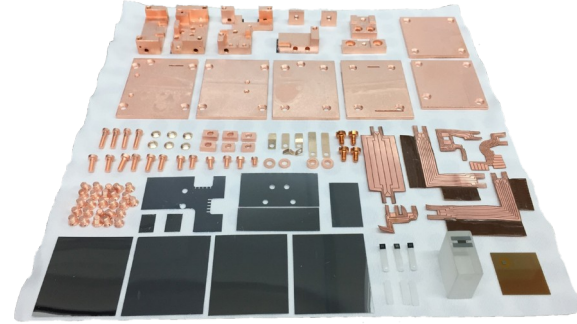
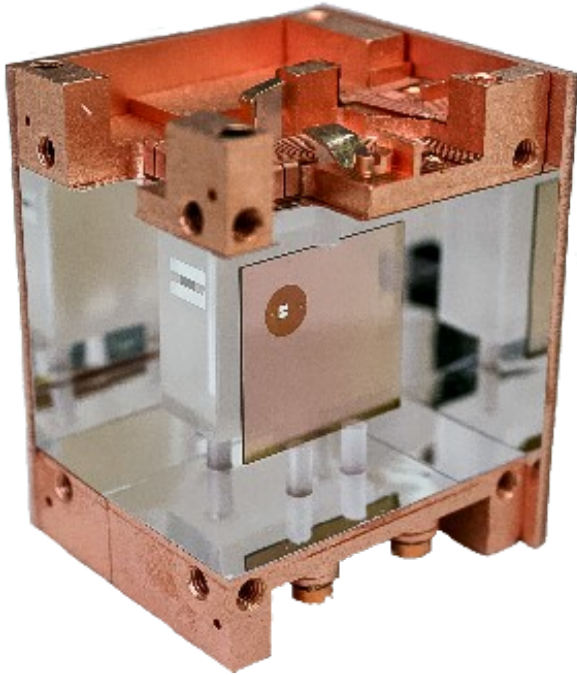
At least 1400 m of rock of the Gran Sasso d'Italia (2912 m) shield against cosmic muons.

The cryostat is enclosed by a layered shielding (PE, Pb, Cu) and an active muon veto.

The CRESST Experiment



The detector modules are mounted in the so-called “carousel” and cooled to ~ 10 mK.

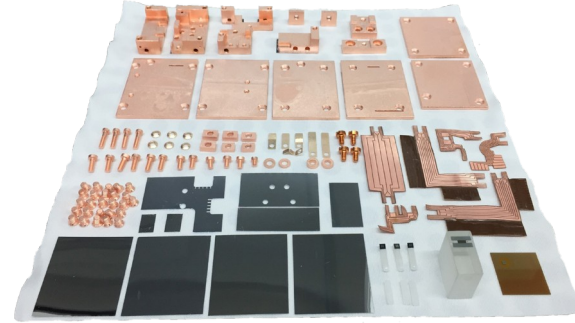
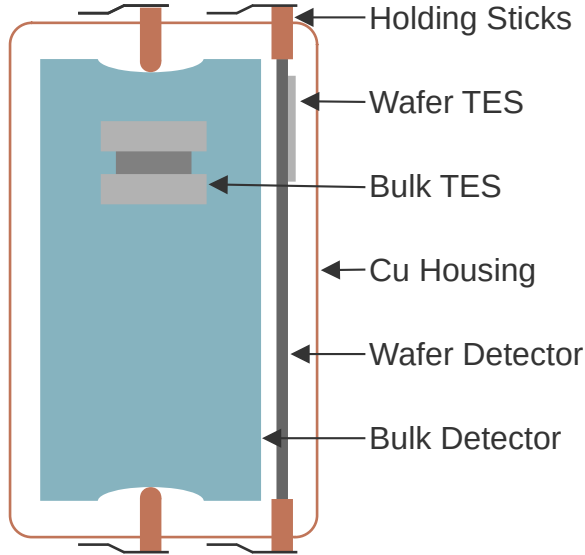


Main absorber: $(2 \times 2 \times 1)$ cm³, broad choice of materials
e.g. CaWO₄ (24 g), Al₂O₃ - sapphire (16 g), LiAlO₂ (10 g), Si (9 g)

Thin wafer detector: $(2 \times 2 \times 0.04)$ cm³, Si or silicon-on-sapphire (SOS)
serves as light detector for scintillating absorbers

Holding structure: light-tight copper housing, scintillating reflector foil
detectors held by sticks from CaWO₄ or copper

Sensors: W-TES directly evaporated on the crystals



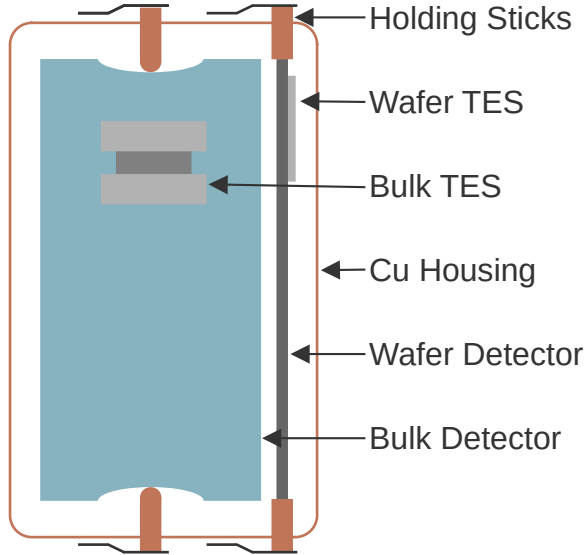
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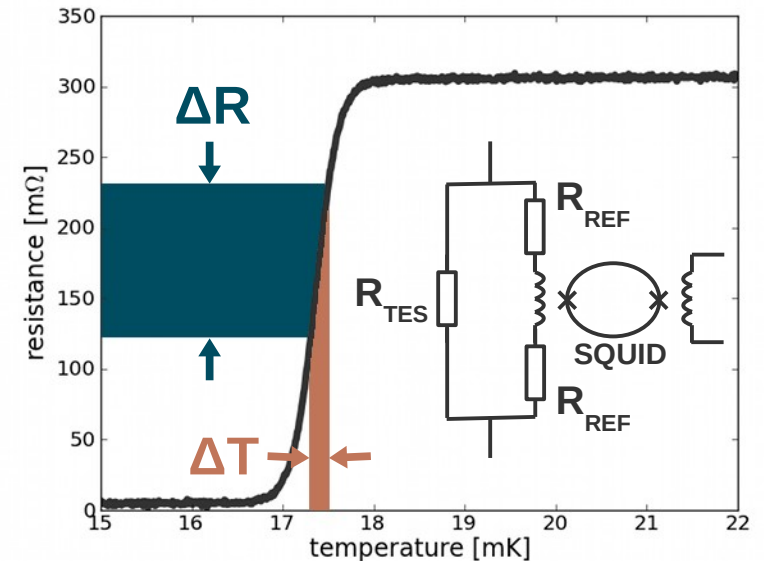
Transition Edge Sensor (TES)



small temperature rise
→ large increase in resistance

Read out by SQUID-based electronics.

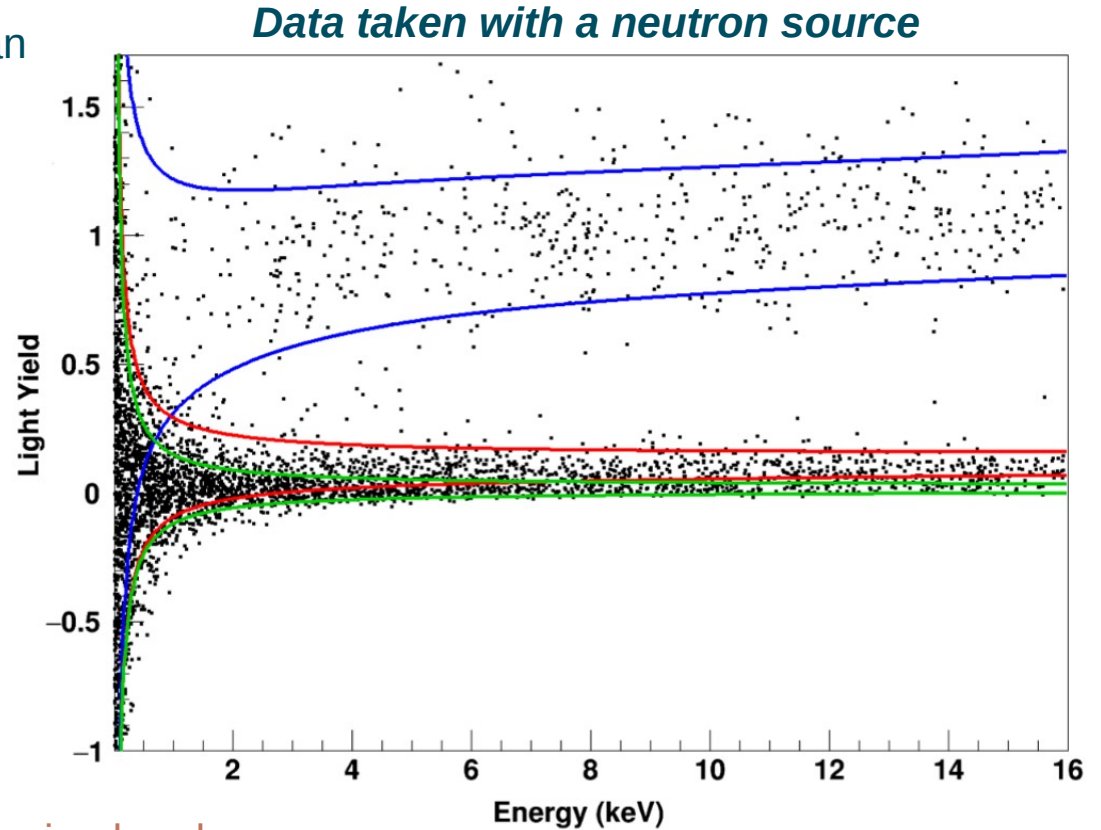
Pulses injected through a separate heater are used to stabilize the operating point and probe the detector response over time.



Sensors: W-TES directly evaporated on the crystals

Phonon-Light-Technique

For scintillating absorbers, the **light yield*** of an event depends on the interaction type:



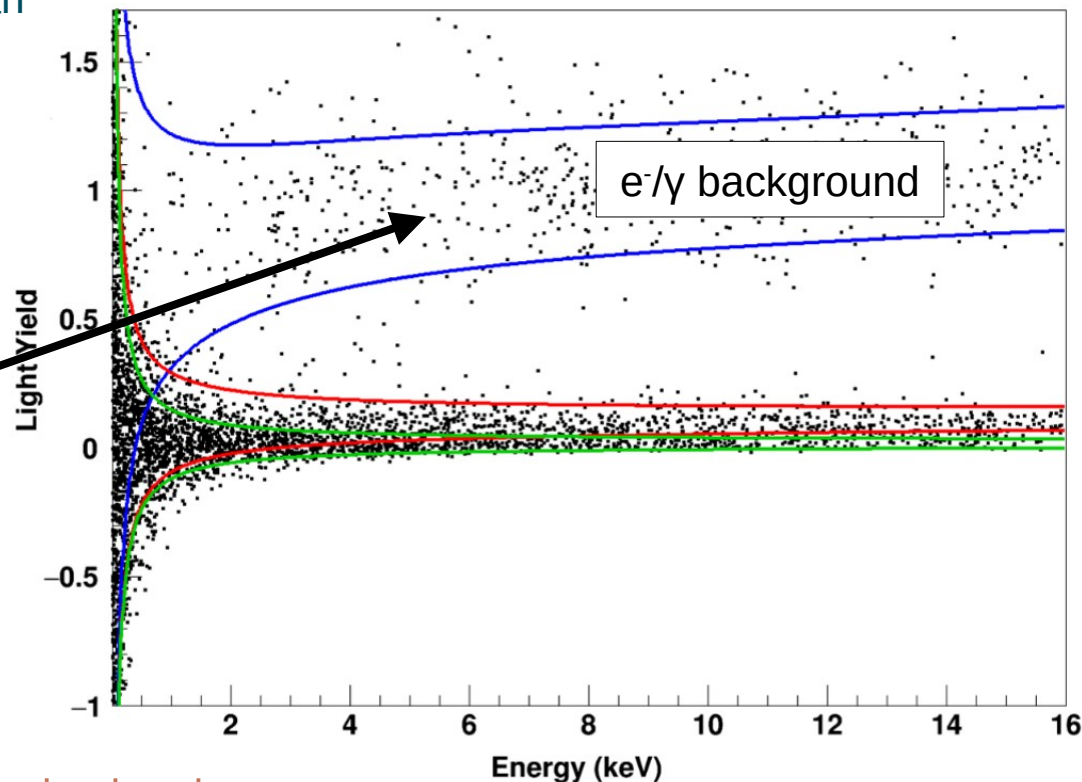
*light yield = detected light energy / energy in main absorber

Phonon-Light-Technique

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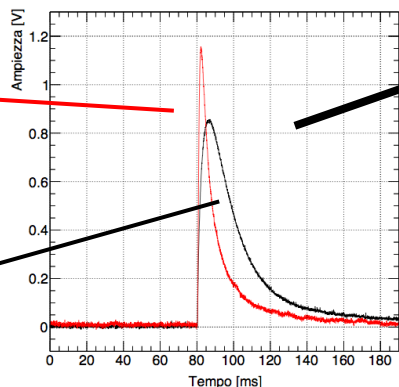
- **e⁻/γ background** has a higher light yield

Data taken with a neutron source



light energy

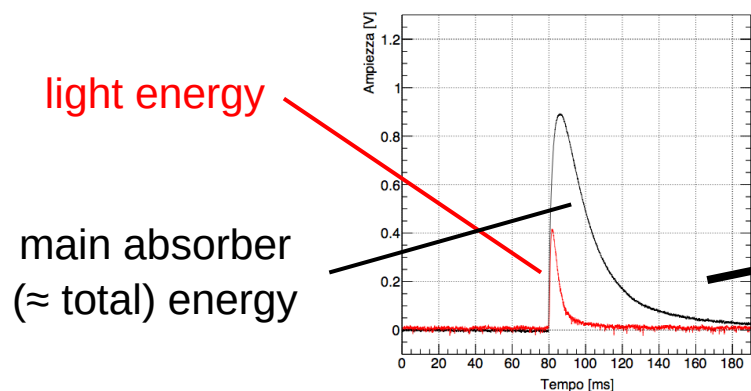
main absorber
(\approx total) energy



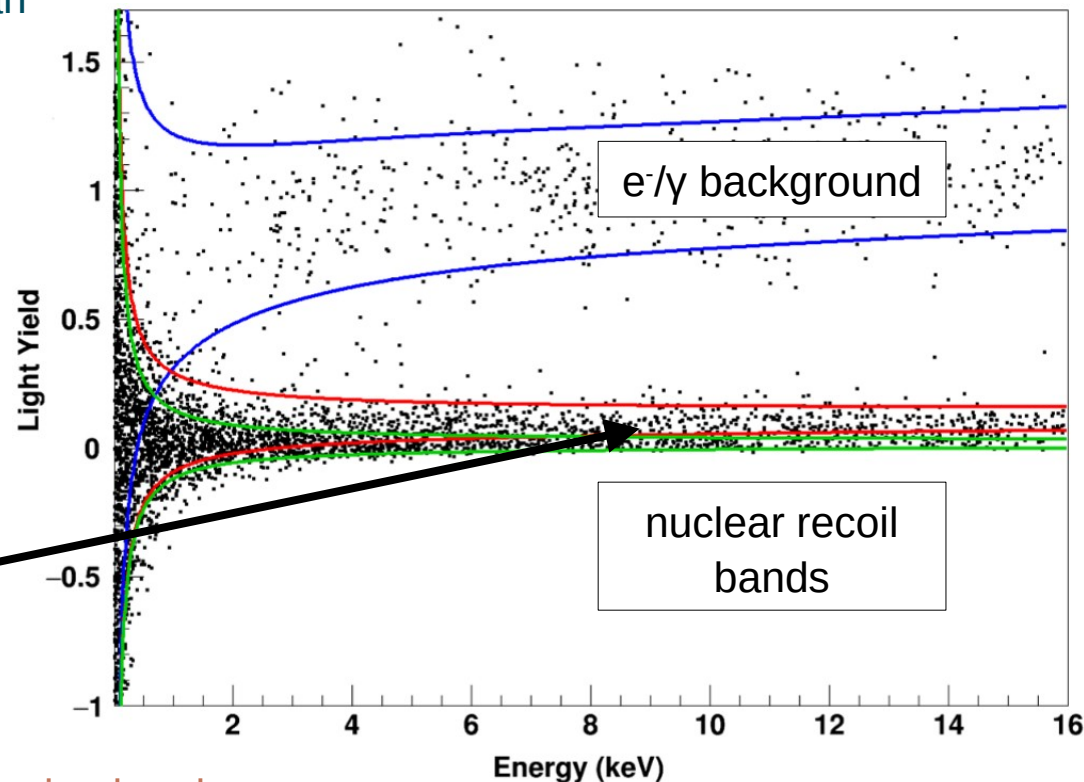
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For scintillating absorbers, the **light yield*** of an event depends on the interaction type:

- **e⁻/γ background** has a higher light yield
- **nuclear recoils** have lower light yield



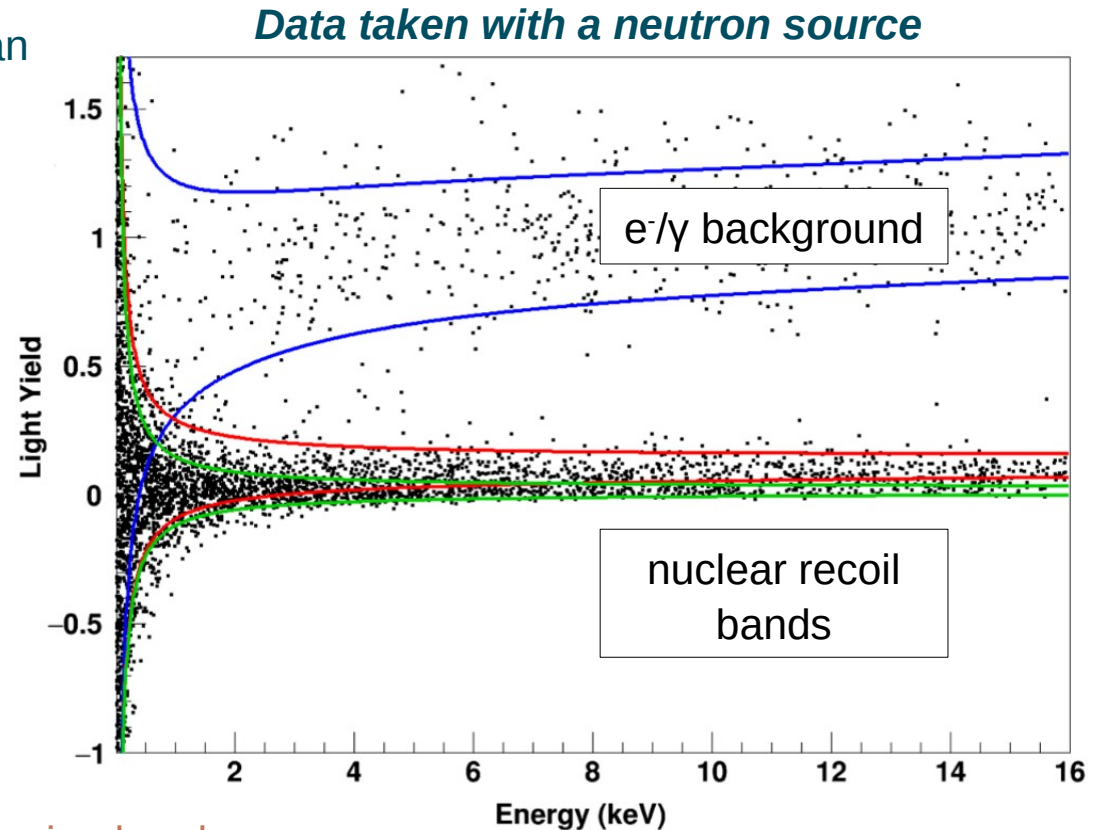
Data taken with a neutron source



*light yield = detected light energy / energy in main absorber

For scintillating absorbers, the **light yield*** of an event depends on the interaction type:

- **e⁻/γ background** has a higher light yield
- **nuclear recoils** have lower light yield
- events cluster in bands around the mean light yield; width of bands given by light detector resolution
- bands are less separated at lower energies
- position and width of the bands can be obtained from neutron calibration data



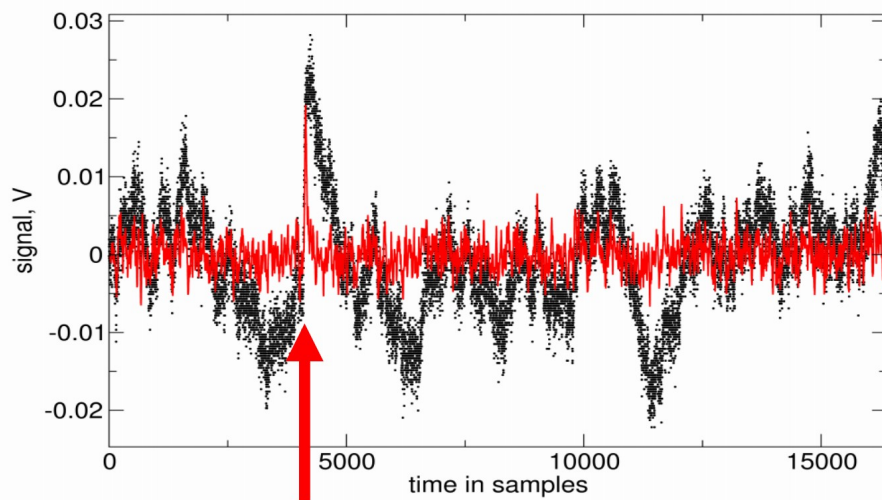
*light yield = detected light energy / energy in main absorber

We record the **full data stream** on disk

→ apply **optimum (matched) filter** to the stream

Optimum filter output is used for:

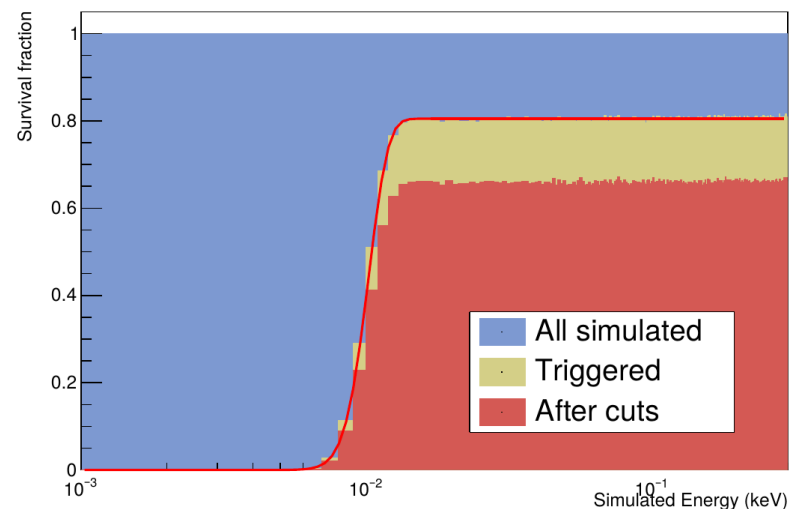
- offline triggering
- energy reconstruction



All analysis steps (calibration, data selection, ...) are defined on a small training set

→ **DM search** is performed on a **blind data set**

Survival probability of signal events is estimated with **artificial pulses** injected on the data stream and passed through the analysis pipeline

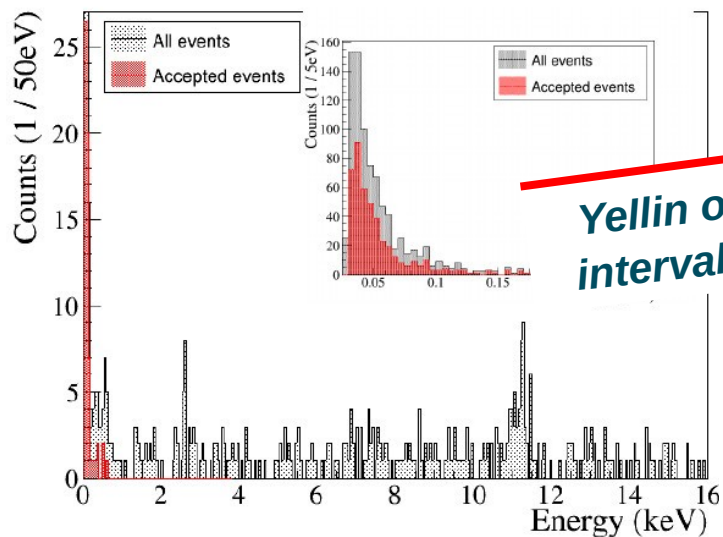


First Results from CRESST-III (2019)

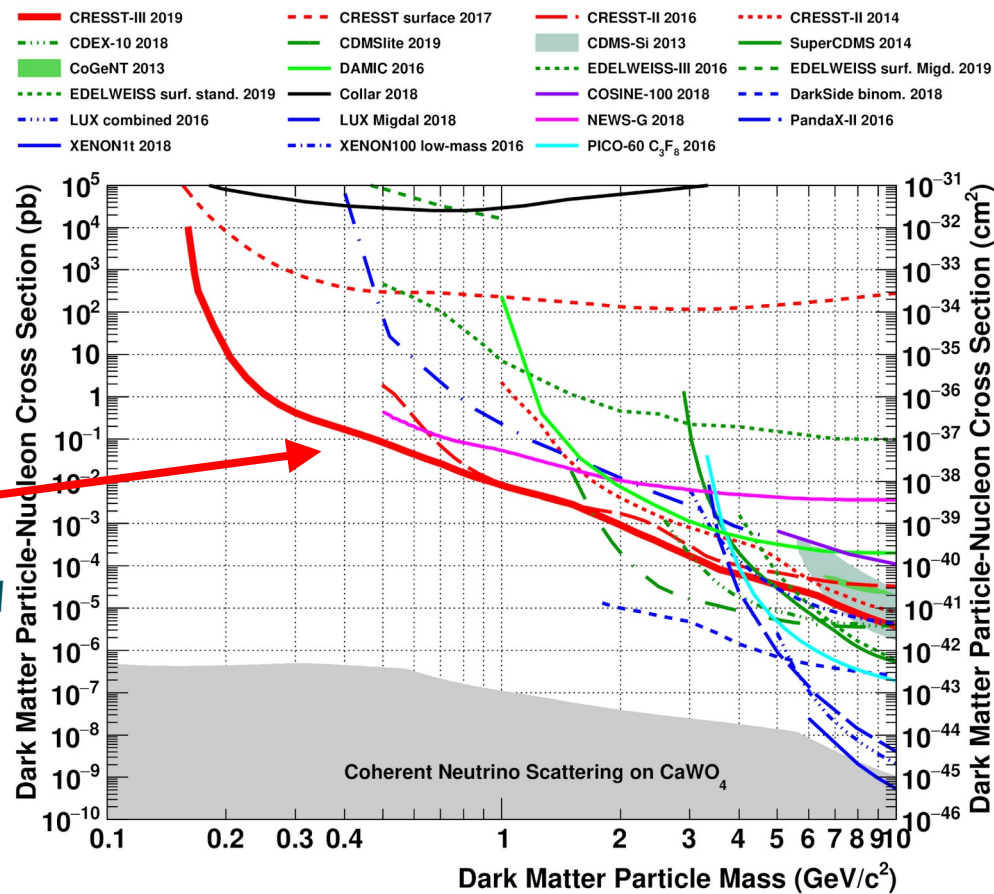


Detector A – 23.6 g CaWO₄

data taking period Oct 2016 – Jan 2018
 exposure 5.698 kg · days
 baseline resolution 4.6 eV
 nuclear recoil threshold 30.1 eV



Yellin optimum interval method



Phys. Rev. D 100, 102002 / arXiv:1904.00498

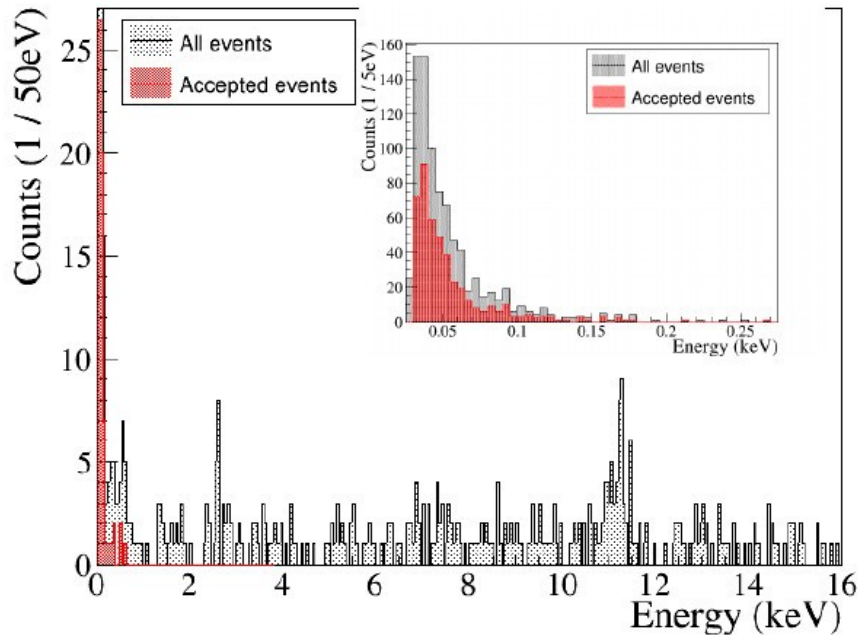
The Low Energy Excess (LEE)

First observations: 2016 – 2018

Unexplained event population at low energies

- high count rate
- steep rise in energy below ~ 200 eV

Detector A

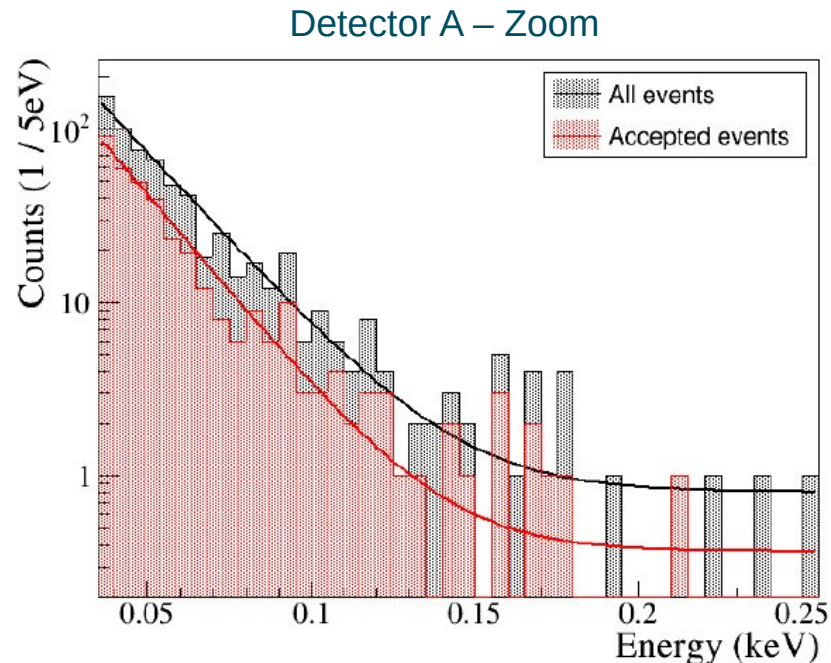


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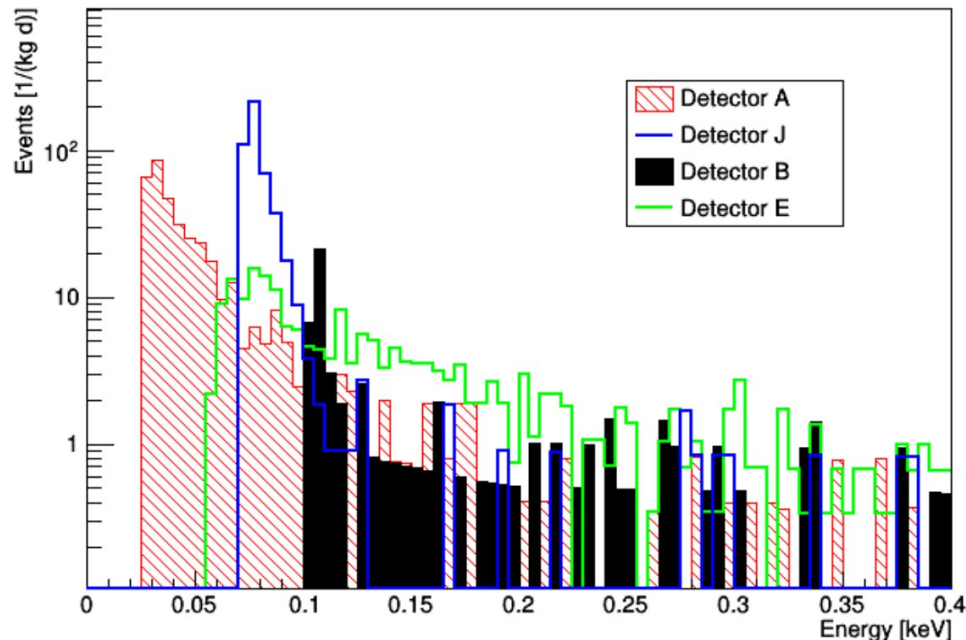


The Low Energy Excess (LEE)

First observations: 2016 – 2018

Unexplained event population at low energies

Overview of CaWO₄ Detectors

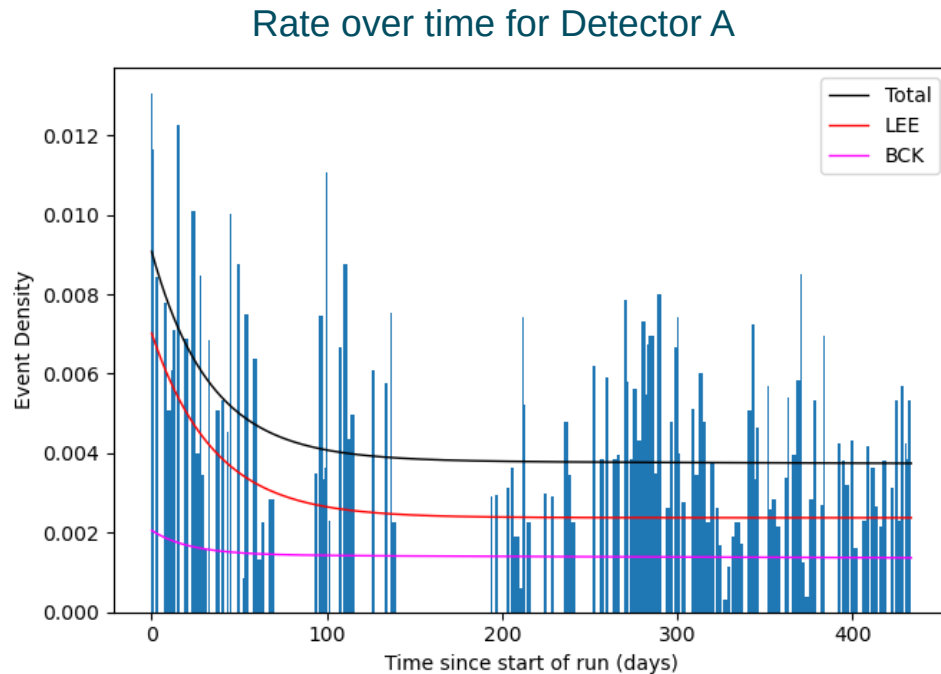


- high count rate
- steep rise in energy below ~200 eV
- different shape in different detectors

Detector	Threshold
Det-A	30.1 eV
Det-B	120 eV
Det-E	64.8 eV
Det-J	83.4 eV

The Low Energy Excess (LEE)

First observations: 2016 – 2018



Unexplained event population at low energies

- high count rate
- steep rise in energy below ~ 200 eV
- different shape in different detectors
- count rate decreases over time



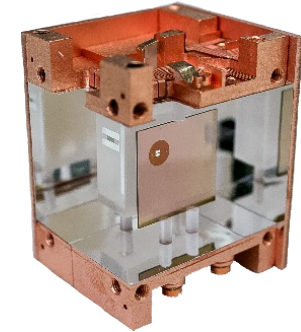
Shown at the first EXCESS
Workshop in June 2021

Current Measurement Campaign

Managed to cool-down the detectors in summer 2020 – despite the pandemic

Took data for DM search from Nov 2020 until Aug 2021

Followed by neutron calibration and measurements of the LEE



Dedicated modifications to probe LEE:

- different target materials
- change how crystals are held
- remove scintillating components

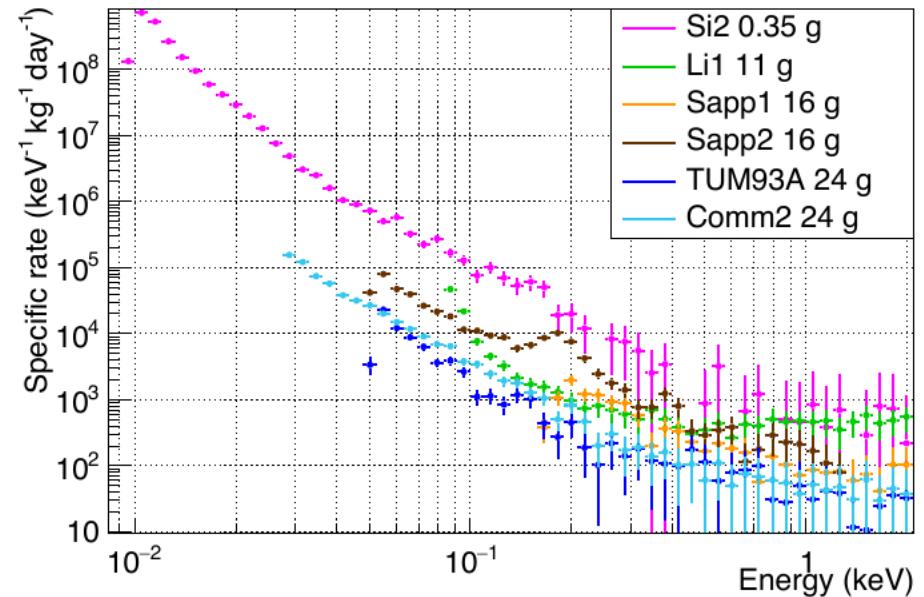
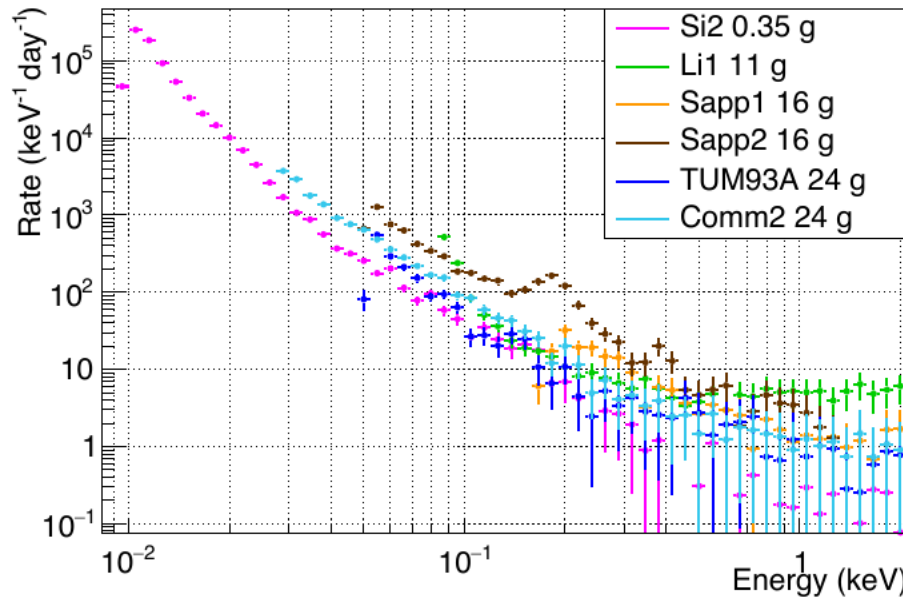
Routinely achieve thresholds < 100 eV

Name	Material	Holding	Foil	Mass	Threshold
Comm2	CaWO ₄	bronze clamps	no	24.5 g	29 eV
TUM93A	CaWO ₄	2 Cu + 1 CaWO ₄	yes	24.5 g	54 eV
Sapp1	Al ₂ O ₃	Cu sticks	no	15.9 g	157 eV
Sapp2	Al ₂ O ₃	Cu sticks	yes	15.9 g	52 eV
Li1	LiAlO ₂	Cu sticks	yes	11.2 g	84 eV
Si2	Si	Cu sticks	no	0.35 g	10 eV

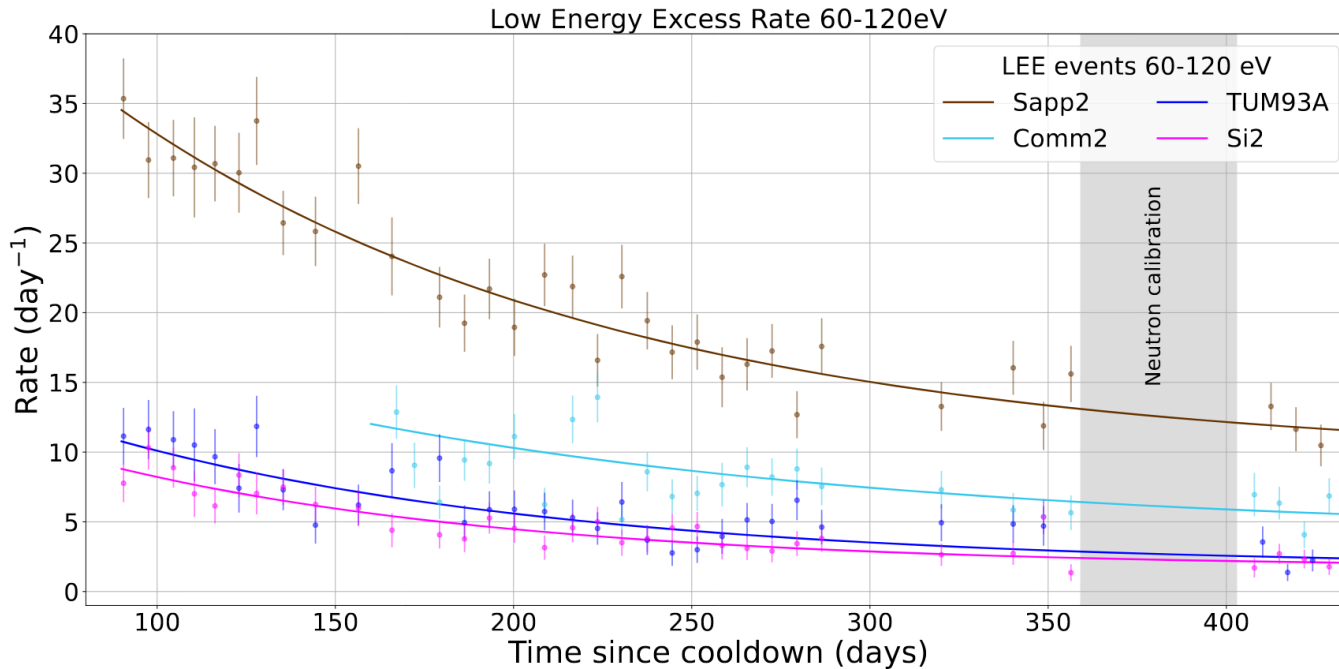
Observations on LEE

The LEE is observed in all these detectors (different materials and geometries).

Scaling the count rate by the absorber mass (right plot) does not improve the agreement between the count rates in different detectors.

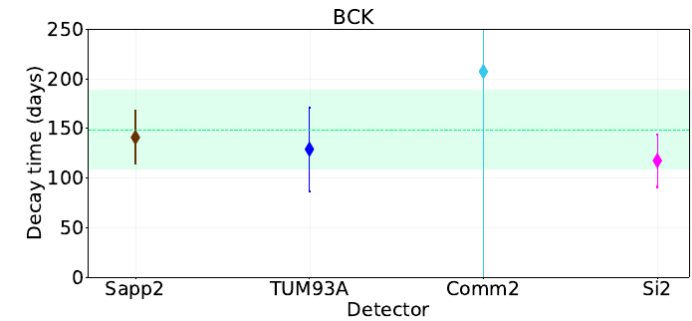


Time-evolution of LEE

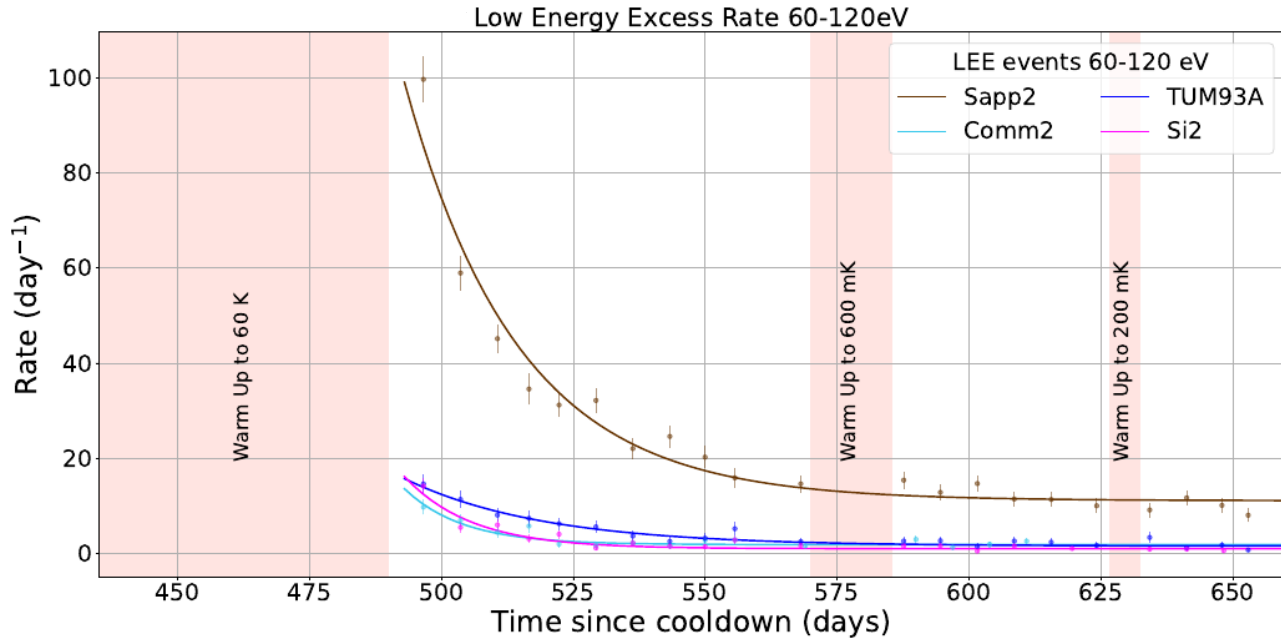


LEE count rate decreases over time in the DM data set ($\tau \approx 150$ d)

Subsequent neutron calibration has no effect on the LEE rate

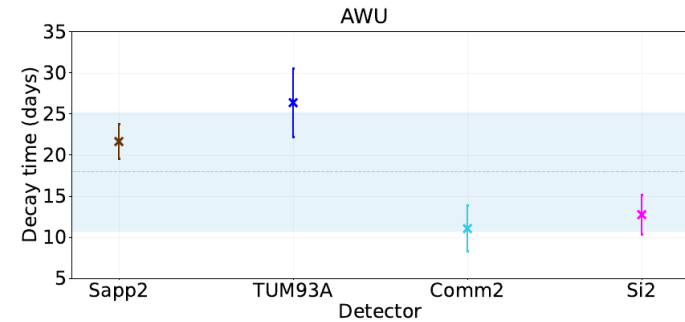


Time-evolution of LEE

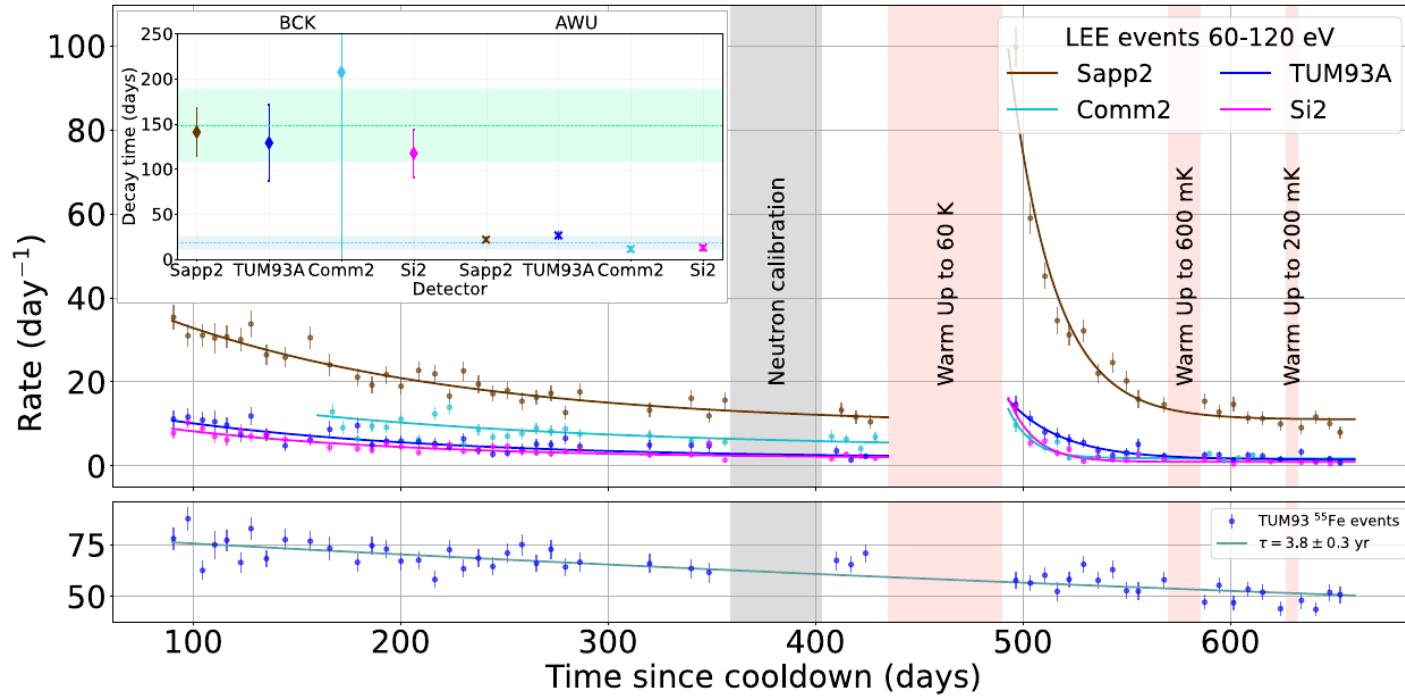


Warm-up to ~ 60 K leads to a sudden increase of the LEE rate, which then decreases again relatively fast ($\tau \approx 18$ d)

Warm-ups to ~ 200 mK and ~ 600 mK have no effect on the LEE rate



Time-evolution of LEE



LEE count rate decreases over time in the DM data set ($\tau \approx 150$ d)

Subsequent neutron calibration has no effect on the LEE rate

Warm-up to ~ 60 K leads to a sudden increase of the LEE rate, which then decreases again relatively fast ($\tau \approx 15$ d)

Warm-ups to ~ 200 mK and ~ 600 mK have no effect on the LEE rate

Time-evolution of LEE



LEE count rate decreases over time in the DM data set ($\tau \approx 150 \text{ d}$)

Subsequent neutron calibration has no effect on the LEE rate

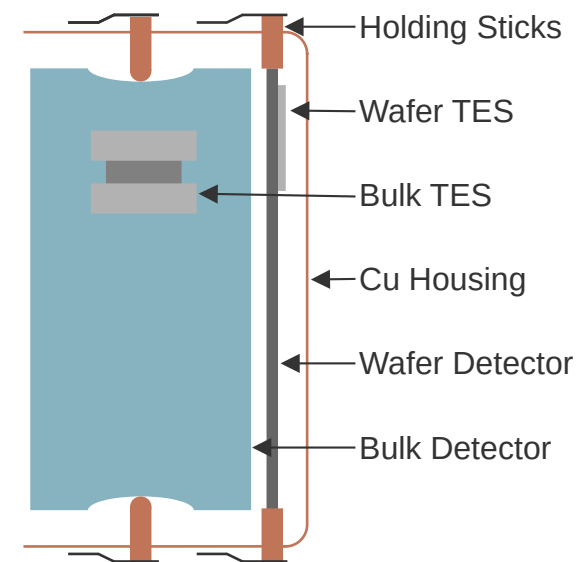
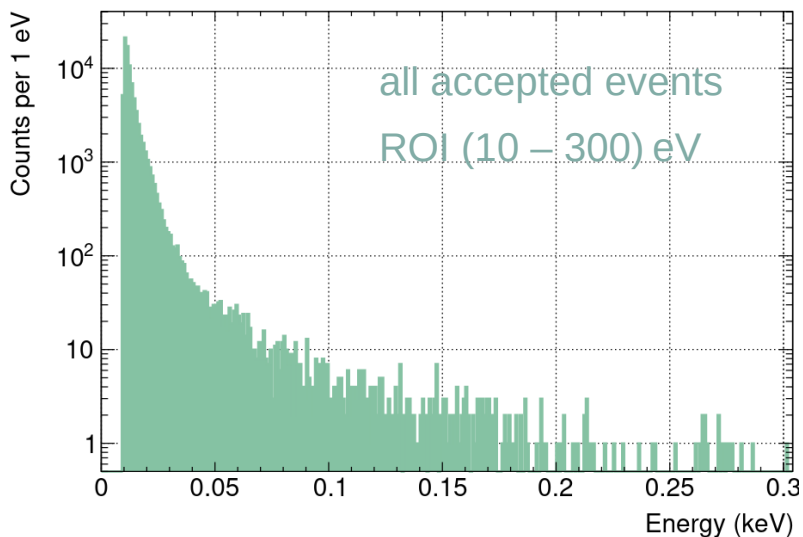
Warm-up to $\sim 60 \text{ K}$ leads to a sudden increase of the LEE rate, which then decreases again relatively fast ($\tau \approx 15 \text{ d}$)

Warm-ups to $\sim 200 \text{ mK}$ and $\sim 600 \text{ mK}$ have no effect on the LEE rate

Results from a Si Wafer Detector

Si2 wafer detector – 0.35 g Si

data taking period Nov 2020 – Aug 2021
exposure 55.06 g days
baseline resolution 1.36 eV
nuclear recoil threshold 10.0 eV

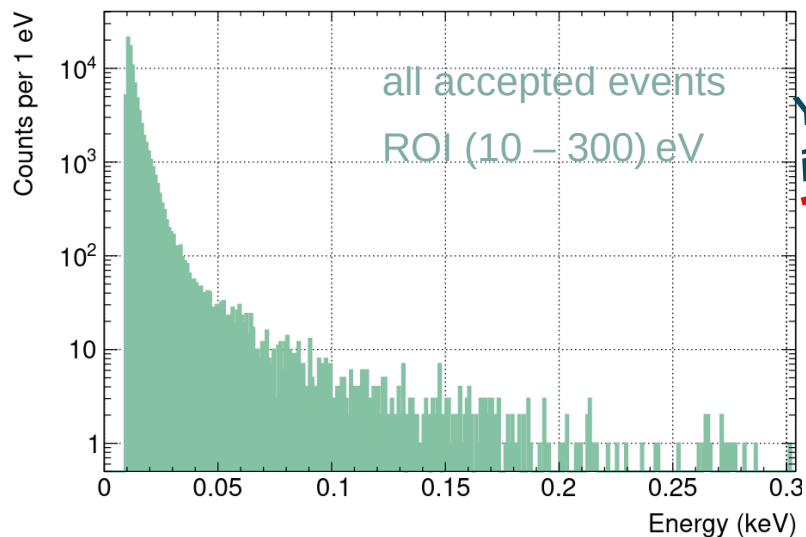


Using thin wafer detector as target and bulky detector as veto detector.

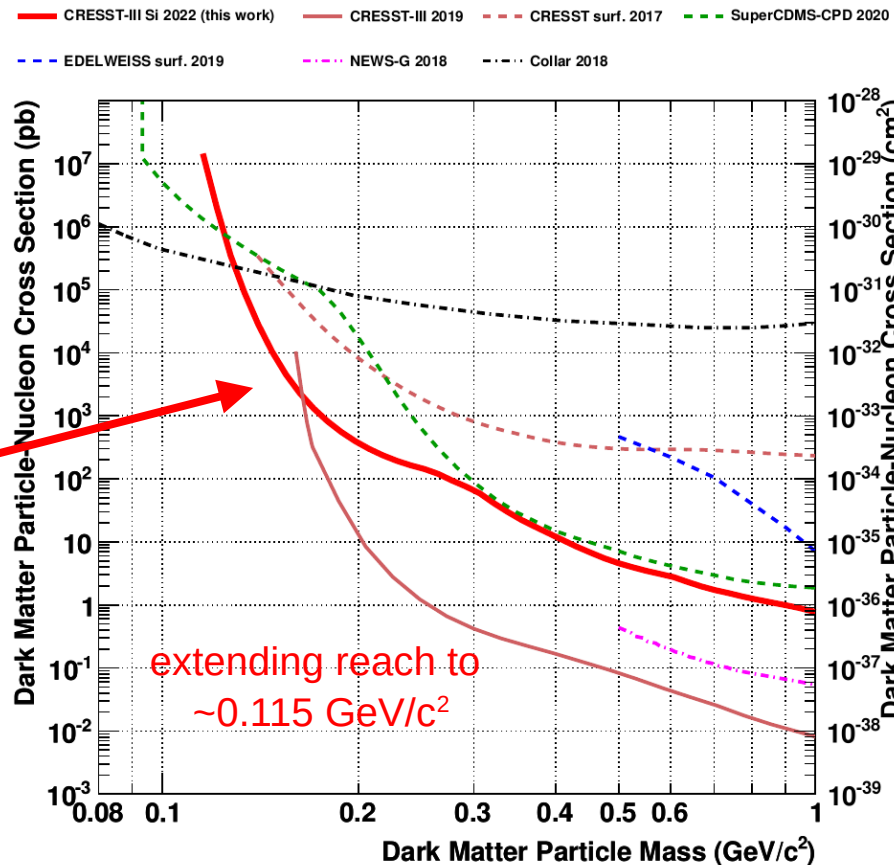
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Yellin optimum
interval method



publication in preparation

SD Limits with LiAlO₂ Detectors

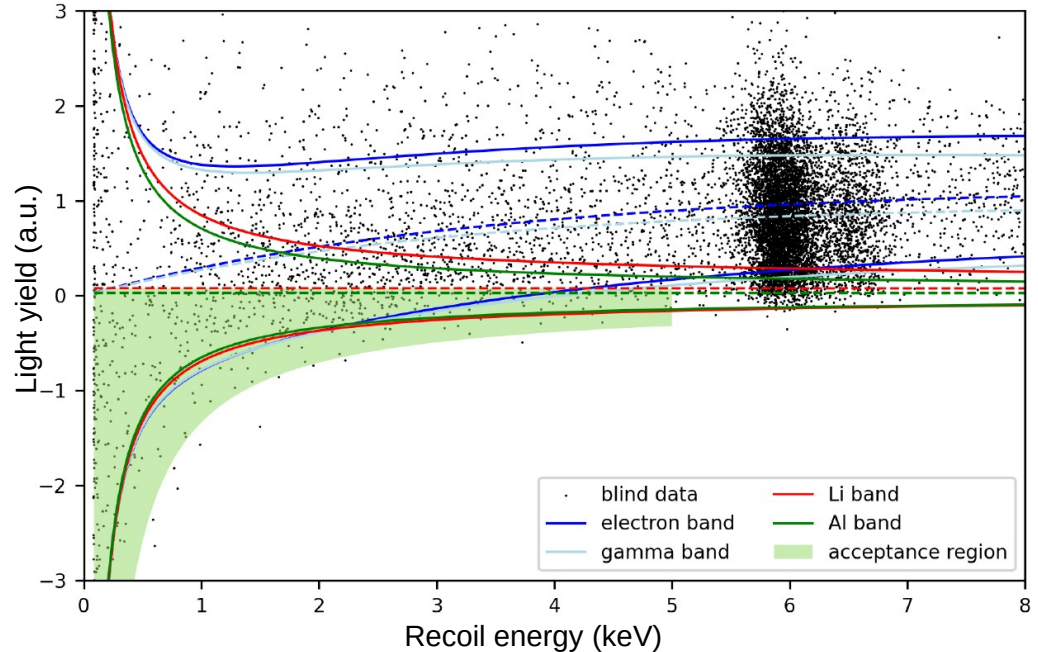
Li1 detector – 11.2 g LiAlO₂

data taking period Nov 2020 – Aug 2021
 exposure 1.161 kg days
 baseline resolution 12.8 eV
 nuclear recoil threshold 83.6 eV

$$\sigma_0^{SD} \propto \mu_N^2 \cdot \frac{J_N + 1}{J_N} \cdot [a_p \cdot \langle S^p \rangle + a_n \cdot \langle S^n \rangle]^2$$

Isotopes sensitive to SD interactions:

Isotope	$\langle S_p \rangle$	$\langle S_n \rangle$
⁶ Li	0.472	0.472
⁷ Li	0.497	---
²⁷ Al	0.343	0.0296



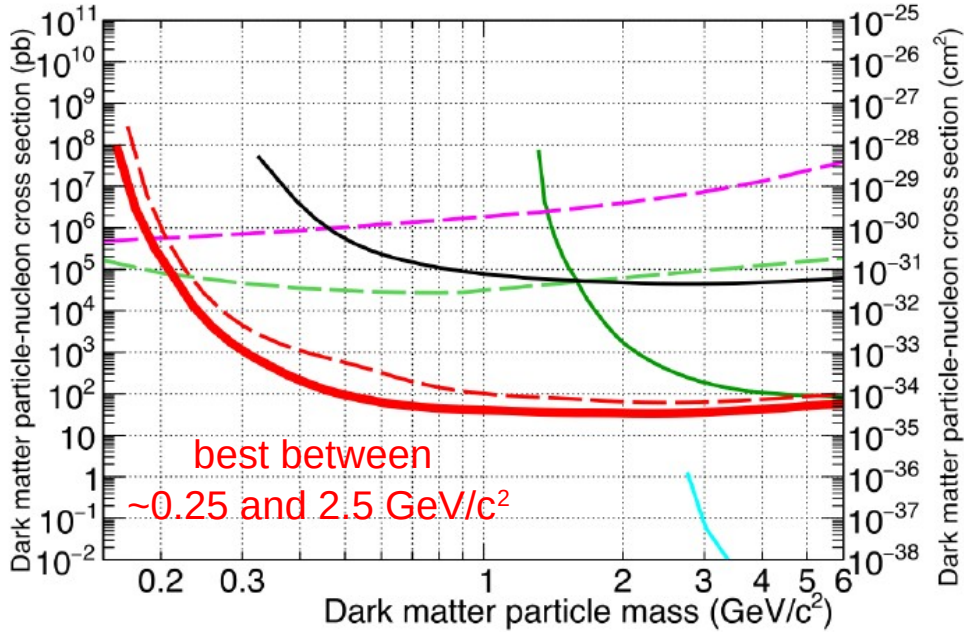
arXiv:2207.07640

SD Limits with LiAlO₂ Detectors



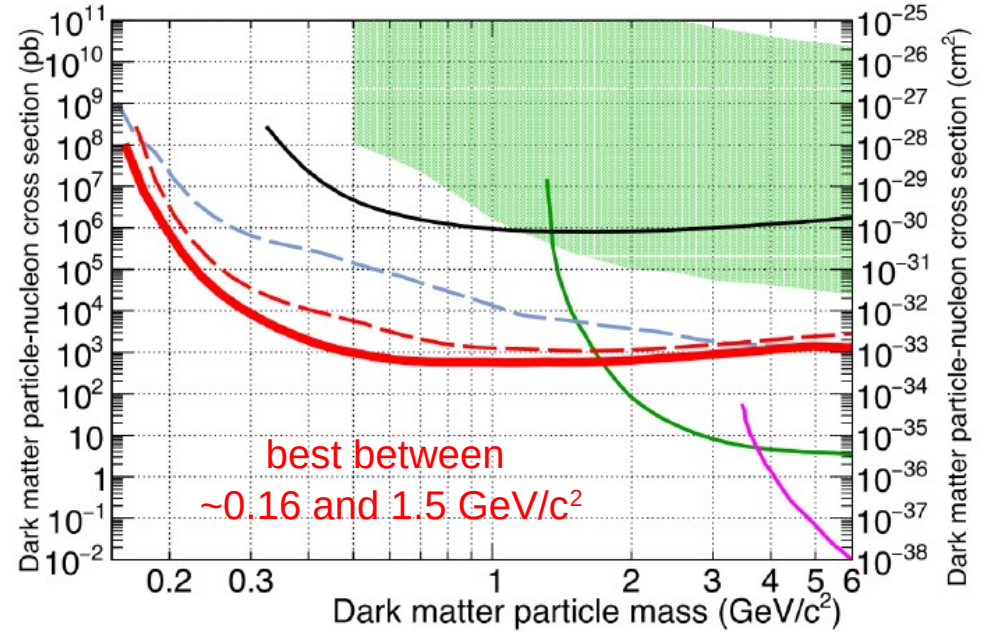
Proton-only

- CDMSlite 2017
- Collar 2018
- CREST LiAlO₂, 2020 above ground
- CREST-III Li¹ 2022
- Borexino 2019
- PICO-60 2019
- CREST-III Li₂ 2022



Neutron-only

- EDELWEISS-Surf 2019
- LUX 2016
- CREST LiAlO₂, 2020 above ground
- CREST-III Li¹ 2022
- CDMSlite 2017
- CREST-III 2019
- CREST-III Li₂ 2022



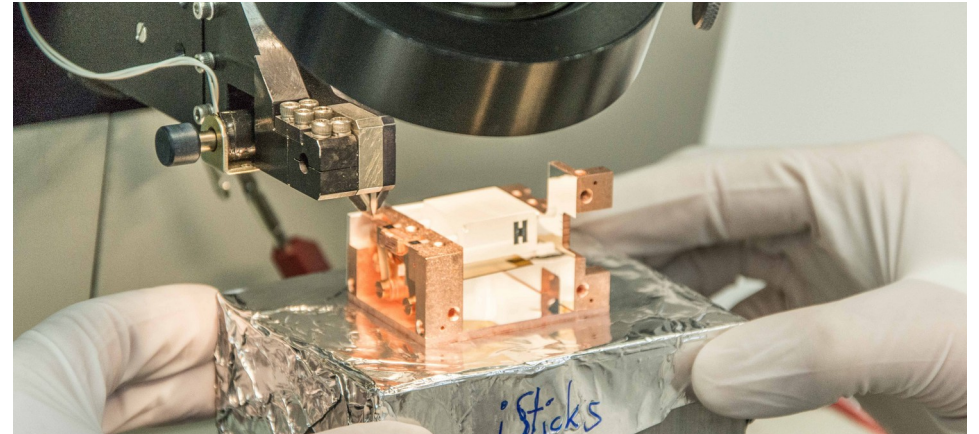
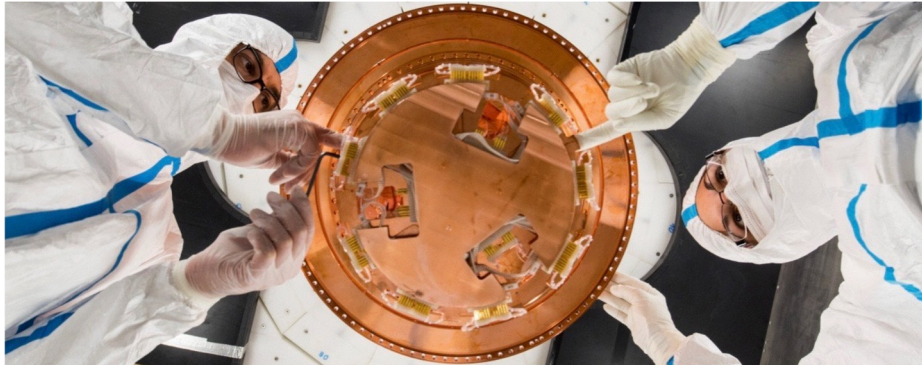
Great improvement over previous above-ground results! Extending reach to 0.16 GeV / c²

What happens next?

Perform **additional warm-up tests** to collect additional information about the LEE.

Detector R&D is ongoing:

- even lower thresholds
- complementary materials
- stress-free holding
- mass production and reproducibility



Upgrade of the **read-out system to 288 channels** is being prepared:

- SQUIDs and wiring is already procured
- new DAQ and bias electronics is designed
- goal: installation at LNGS in 2023

Exciting **new dark matter results**:

- extending reach to lower masses with **10 eV threshold** Si detector
- unique ability to use a **wide-range of target materials** enables SD results with LiAlO_2 absorbers

Broad R&D program to **identify the origin of the LEE** and mitigate its effects

Ongoing efforts to **improve our detectors** (lower thresholds, increase radiopurity of crystals, ...)

Preparing for a **major infrastructure upgrade** of the read-out electronics to accommodate up to 100 detector modules



Thanks for your Attention!

Please also check out the other CRESST contributions @ IDM:

Talks

“Probing Li targets in CRESST-III”

S. Gupta

up next!

“Characterization of a low background CaWO₄ crystal for CRESST-III”

A. Kinast

Thu 14:40

Posters

“Testing spin-dependent dark matter with lithium targets in CRESST-III” – F.Wagner

“Neutron simulation studies and their implications for CRESST” – A. Fuss

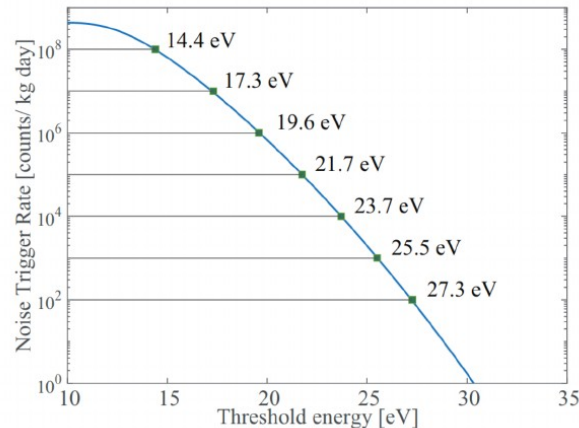
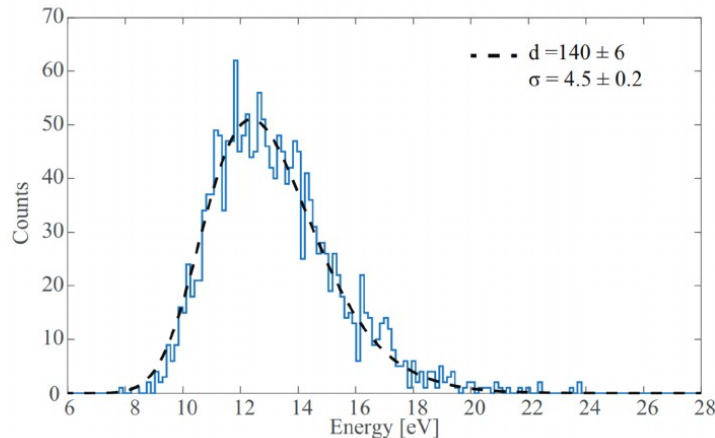
“Using a likelihood fit to identify radioactive background components in the CRESST experiment” – J. Burkhart

Threshold Determination

Analytical description of **amplitude distribution** of “empty” noise traces

→ see NIM A **940**, 492 (2019) / arXiv:1711.11459

$$NTR(x_{thr}) = \frac{1}{t_{win} \cdot m_{det}} \cdot \int_{x_{thr}}^{\infty} P_d(x_{max})$$



Allows to define the threshold as a function of allowed noise trigger rate

Default choice in CRESST:
1 noise trigger / (kg day)

LEE Pulse Shape

Average of all pulses just above threshold (“LEE template”) agrees with expected signal pulse shape from particle interactions.

