

Unified  
DM/DE

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Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

# A Unified Model of Dark Energy, Dark Matter, and Baryogenesis

Robert Brandenberger  
Physics Department, McGill University, Canada

CERN Colloquium, June 24 2020  
Work in collaboration with J. Fröhlich  
`arXiv:2004.10025`

# Outline

Unified  
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Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

- 1 Introduction
- 2 Standard Paradigm in Trouble
  - Dark Energy cannot be  $\Lambda$
  - Challenges for the WIMP Paradigm
- 3 Unified Model of DE, DM and Baryogenesis
- 4 Baryogenesis
- 5 Discussion and Conclusions

# Plan

Unified  
DM/DE

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berger

Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

- 1 Introduction
- 2 Standard Paradigm in Trouble
  - Dark Energy cannot be  $\Lambda$
  - Challenges for the WIMP Paradigm
- 3 Unified Model of DE, DM and Baryogenesis
- 4 Baryogenesis
- 5 Discussion and Conclusions

# Evidence for Dark Matter: Galaxy Velocity Rotation Curves

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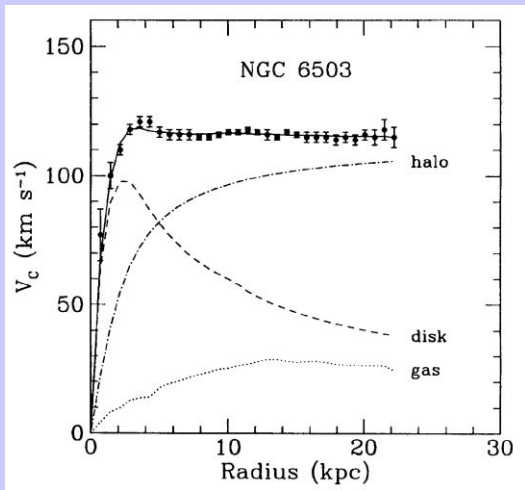
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Paradigm

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No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions



# Evidence for Dark Matter: Bullet Cluster

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## Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions



# Evidence for Dark Matter: Cosmological Structure Formation

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Introduction

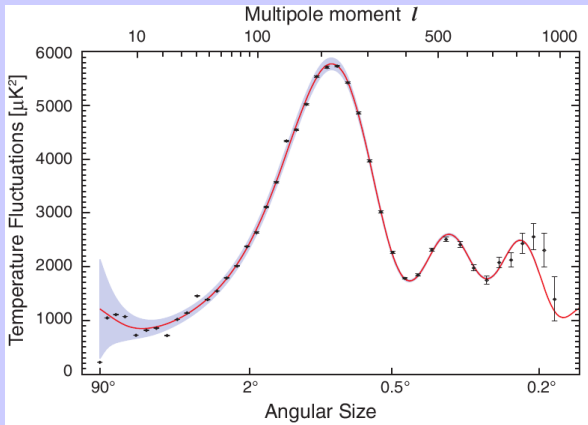
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DE is not  $\Lambda$   
No WIMP

Unified Dark Sector

Baryogenesis

Conclusions



Credit: NASA/WMAP Science Team

# Evidence for Dark Matter: Cosmological Structure Formation

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Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

- Without DM: matter fluctuations on galaxy scales couple to radiation for  $t_{eq} < t < t_{rec} \rightarrow$  decay as  $a(t)^{-1/2}$ .
- If DM dominates over baryonic matter: matter fluctuations grow  $\propto a(t)$  for  $t_{eq} < t < t_{rec}$ .
- Fix the matter power spectrum by observations.
- $\rightarrow$  without DM the matter fluctuations at  $t_{eq}$  were larger than in the case of dominant DM by a factor of  $(a(t_{rec})/a(t_{eq}))^{3/2}$
- $\rightarrow$  amplitude of the angular power spectrum of CMB anisotropies predicted to be  $\sim 10^{-3}$  (R. Sachs and A. Wolfe, Ap. J. 147, 73 (1967)).

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Unified  
DM/DE

R. Branden-  
berger

Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

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Unified  
DM/DE

R. Branden-  
berger

Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

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# Evidence for Dark Energy: Cosmological Probes

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## Introduction

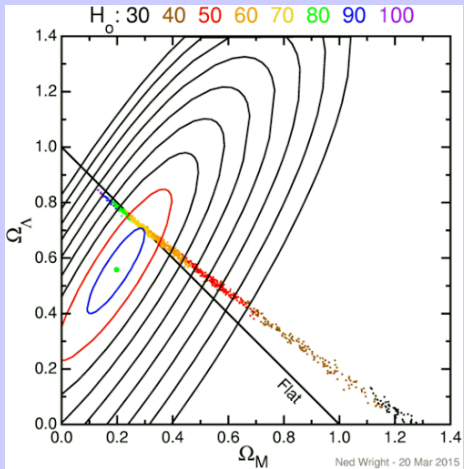
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No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions



# Evidence for Dark Energy: Supernovae

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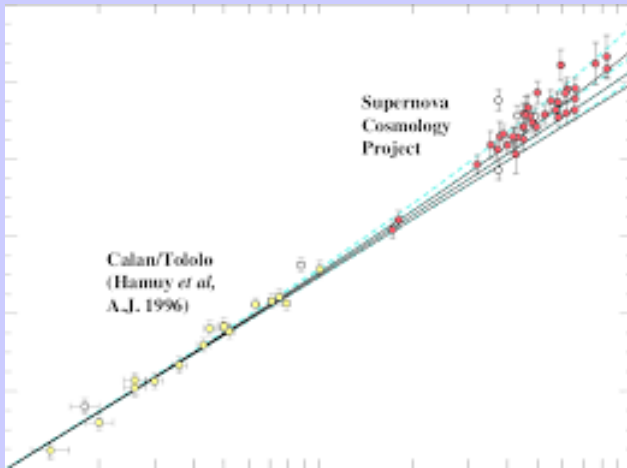
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Unified Dark  
Sector

Baryogenesis

Conclusions



# Evidence for Baryogenesis: Cosmic Rays

A. Cohen, A. De Rujula and S. Glashow., *Ap. J.* 495, 539 (1998)

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## Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

- Assume a patchwork of matter and antimatter regions.
- $\rightarrow$  after recombination annihilations at the regional boundaries.
- $\rightarrow$  cosmic diffuse  $\gamma$ -ray background.
- Excludes  $B = 0$  universe with domains smaller than Hubble volume.

# Evidence for Baryogenesis: Nucleosynthesis

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## Introduction

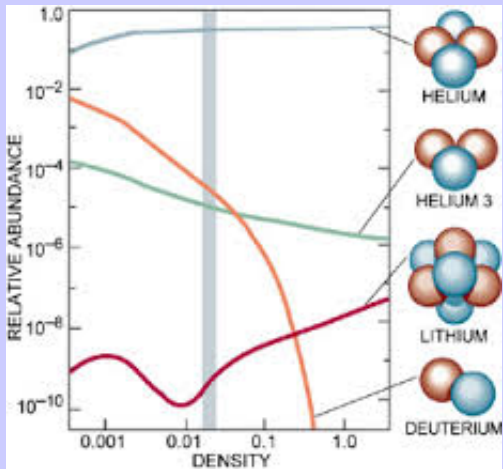
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Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions



# Standard Paradigm

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## Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

- Dark Energy is the Cosmological Constant  $\Lambda$ .
- Dark Matter is a WIMP (weakly interacting massive particle).

# Standard Paradigm for Dark Energy: LCDM

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R. Branden-  
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Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

- Dark Energy is a **cosmological constant**.
- No new parameters -  $\Lambda$  must arise in the low energy EFT of gravity.
- No conflict with data.
- **Cosmological constant problem**: why is  $\Lambda$  so small?
- **Coincidence problem**: why is  $\rho_\Lambda \sim \rho_m$  today?

# Standard Paradigm for Dark Energy: LCDM

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berger

## Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

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berger

## Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

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# Plan

Unified  
DM/DE

R. Branden-  
berger

Introduction

**Standard  
Paradigm**

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

- 1 Introduction
- 2 **Standard Paradigm in Trouble**
  - Dark Energy cannot be  $\Lambda$
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- 3 Unified Model of DE, DM and Baryogenesis
- 4 Baryogenesis
- 5 Discussion and Conclusions

# Trans-Planckian Problem

J. Martin and R.B., *Phys. Rev. D*63, 123501 (2002)

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Introduction

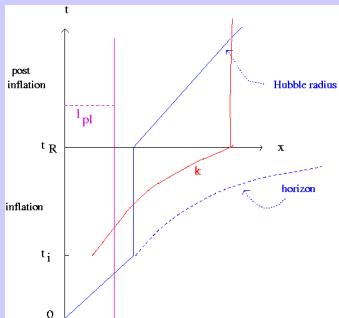
Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions



- **Success of inflation:** At early times scales are inside the Hubble radius  $\rightarrow$  causal generation mechanism is possible.
- **Problem:** If time period of inflation is more than  $70H^{-1}$ , then  $\lambda_p(t) < l_{pl}$  at the beginning of inflation.
- $\rightarrow$  new physics **MUST** enter into the calculation of the fluctuations.

# Trans-Planckian Censorship Conjecture (TCC)

A. Bedroya and C. Vafa., arXiv:1909.11063

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Introduction

Standard  
Paradigm

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No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

No trans-Planckian modes ever exit the Hubble horizon.

$$ds^2 = dt^2 - a(t)^2 dx^2$$

$$H(t) \equiv \frac{\dot{a}}{a}(t)$$

$$\frac{a(t_R)}{a(t_i)} \Big|_{pl} < H(t_R)^{-1}$$

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Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

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# Implications for Dark Energy

Unified  
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R. Branden-  
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Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

## Dark Energy cannot be a Cosmological Constant

- If DE is  $\Lambda$ , then  $H^{-1}$  constant to the future.
- $\rightarrow$  Planck scale today eventually becomes larger than  $H^{-1}$ .

# Implications for Dark Energy

Unified  
DM/DE

R. Branden-  
berger

Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

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# Sub-Extremal Black Hole

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Introduction

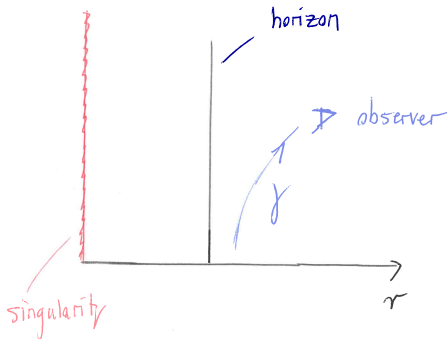
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Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions





# Super-Extremal Black Hole

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Introduction

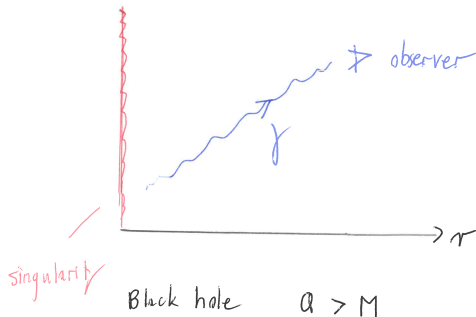
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Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions



# Penrose's Cosmic Censorship Hypothesis

R. Penrose, *Riv. Nuovo Cim.* 1, 252 (1969)

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Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

**Time-like singularities are hidden by horizons.**

- Singularity hidden from external observer.
- Non-unitarity and acausality hidden from external observer.

Note: **Effective field theory** of GR allows for naked singularities, but **UV complete theory** does not.

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Unified  
DM/DE

R. Branden-  
berger

Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

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# Non-Unitarity of Effective Field Theory (EFT) in an Expanding Background

N. Weiss, *Phys. Rev. D.* 32, 3228 (1985)

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Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
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Unified Dark  
Sector

Baryogenesis

Conclusions

- Hilbert space of EFT:  $\mathcal{H} = \prod_k \mathcal{H}_k$
- where  $\mathcal{H}_k$  is the Hilbert space of the  $k$ 'th **comoving** mode.
- **UV cutoff** required in **physical coordinates**
- Corresponding comoving UV cutoff is time-dependent
- $k_{UV}(t) \sim a(t)$
- $\rightarrow$  **non-unitarity**.

# Generalizing Penrose's Cosmic Censorship Hypothesis

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Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

- Singularity  $\rightarrow$  trans-Planckian region.
- Black hole horizon  $\rightarrow$  Hubble horizon.

**Demand:** The external observer is shielded from non-unitarity associated with the trans-Planckian region by the Hubble horizon.

Note:

- EFT of cosmology allows for violations of TCC.
- Conjecture: **UV physics prohibits violations of TCC.**

# Generalizing Penrose's Cosmic Censorship Hypothesis

Unified  
DM/DE

R. Branden-  
berger

Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

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Unified  
DM/DE

R. Branden-  
berger

Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

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# Application to Inflation

A. Bedroia, R.B., M. Loverde and C. Vafa., arXiv:1909.11106

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Introduction

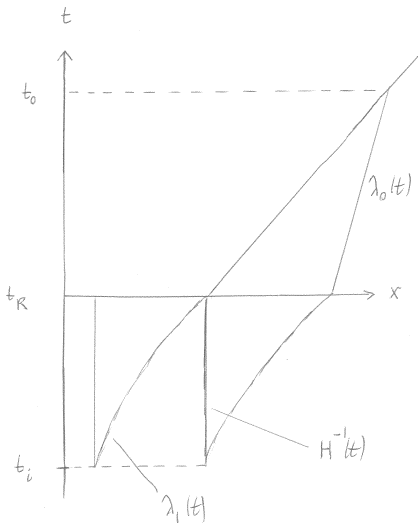
Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions





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A. Bedroya, R.B., M. Loverde and C. Vafa., arXiv:1909.11106

Unified  
DM/DE

R. Branden-  
berger

Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

TCC implies:

$$\frac{a(t_R)}{a(t_*)} |_{pl} < H(t_R)^{-1}$$

Demanding that inflation yields a causal mechanism for generating CMB anisotropies implies:

$$H_0^{-1} \frac{a(t_0)}{a(t_R)} \frac{a(t_R)}{a(t_*)} < H^{-1}(t_*)$$

# Implications for Inflation

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Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

**Upper bound** on the **energy scale of inflation**:

$$V^{1/4} < 3 \times 10^9 \text{ GeV}$$

→ **upper bound** on the **primary tensor to scalar ratio**  $r$ :

$$r < 10^{-30}$$

Note: Secondary gravitational waves are generated from scalar fluctuations.

# Implications for Inflation

Unified  
DM/DE

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Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

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# Swampland

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Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

- There is a vast **landscape** of **effective field theories**.
- Any space-time dimension, and number of fields, any shape of the potential, any field range, and value of  $\Lambda$ .
- **Superstring theory** is very **constraining**.
- Only a **small subset** of all EFTs is consistent with string theory.
- The rest lie in the **swampland**.

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Introduction

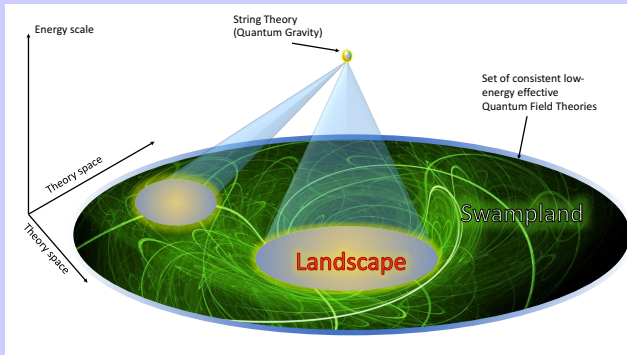
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Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions



# Swampland Conjectures

H. Ooguri and C. Vafa, hep-th/0605264; G. Obied, H. Ooguri, L. Spodyneiko and C. Vafa, arXiv:1806.08362; H. Ooguri, E. Palti, G. Shiu and C. Vafa, arXiv:1810.05506.

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Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

What are conditions for habitable islands sticking out from the swamp?

- Bare  $\Lambda \leq 0$ .
- The effective field theory is only valid for  $\Delta\varphi < dm_{pl}$  (field range condition).
- The potential of  $\varphi$  obeys (de Sitter conjecture)

$$\left| \frac{V'}{V} \right| m_{pl} \geq c_1 \text{ or}$$
$$\frac{V''}{V} m_{pl}^2 \leq -c_2$$

Note:  $d, c_1, c_2$  constants of order 1.

# Lesson

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Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

Dark Energy cannot be  $\Lambda$ .

# Reasoning

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R. Branden-  
berger

Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

- $\varphi$  is a string theory modulus field.
- **Field range condition:** move  $\delta\varphi \sim m_{pl} \rightarrow$  tower of string states becomes massless and must be included in low energy EFT.
- **De Sitter conjecture:** moduli stabilization  $\rightarrow$  steep potential.
- Example: S. Laliberte and R.B., arXiv:1911.00199:



# Reasoning

Unified  
DM/DE

R. Branden-  
berger

Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

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# Current Limits on WIMPs

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Introduction

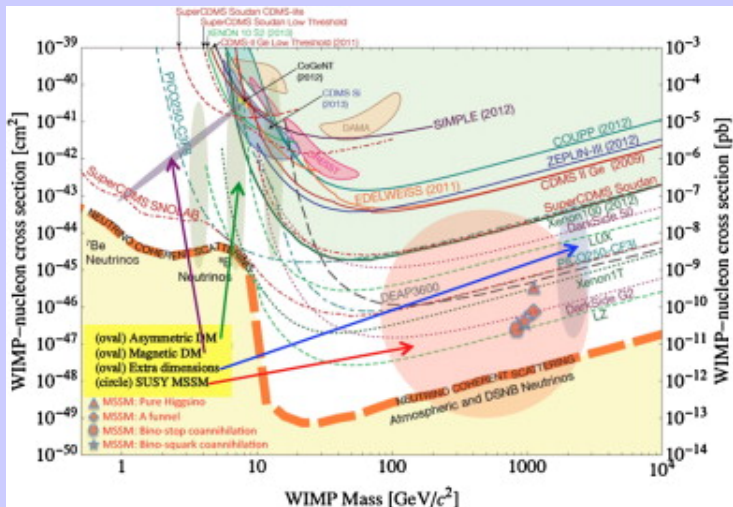
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Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions



# Alternatives to WIMP DM Exist

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Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

## Alternatives to WIMP DM exist:

- **Axions** Well motivated, arising in the solution of the Strong CP Problem.
- **Superfluid Dark Matter** (see e.g. E. Ferreira arXiv:2005.03254 for a review)
- **Ultralight Dark Matter**
- .....

# Plan

Unified  
DM/DE

R. Branden-  
berger

Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

- 1 Introduction
- 2 Standard Paradigm in Trouble
  - Dark Energy cannot be  $\Lambda$
  - Challenges for the WIMP Paradigm
- 3 Unified Model of DE, DM and Baryogenesis
- 4 Baryogenesis
- 5 Discussion and Conclusions

# Motivation

R.B. and J. Fröhlich, arXiv:2004.10025

Unified  
DM/DE

R. Branden-  
berger

Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

- Can a single (complex) scalar field describe both dark matter and dark energy?
- Can one obtain baryogenesis as an added bonus?

See also R.B., J.F. and R. Namba, *Unified Dark Matter, Dark Energy and baryogenesis via a cosmological wetting transition* JCAP **1909**, 069 (2019) [arXiv:1907.06353] .

See also S. Alexander, R.B. and J.F. *Tracking Dark Energy from Axion-Gauge Field Couplings*, arXiv:1601.00057 [hep-th].

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R.B. and J. Fröhlich, arXiv:2004.10025

Unified  
DM/DE

R. Branden-  
berger

Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

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# Model

R.B. and J. Fröhlich, arXiv:2004.10025

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Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

Consider a **complex scalar field**  $\zeta$

$$\zeta \equiv e^{-(\varphi+i\theta)/f}$$

whose angular variable  $\theta$  is an **axion**.

Action:

$$S = \int d^4x \sqrt{-g} \left\{ \frac{1}{2} \partial_\mu \varphi \partial^\mu \varphi + \frac{1}{2} \partial_\mu \theta \partial^\mu \theta - \Lambda e^{-2\varphi/f} - V(\varphi, \theta) + \dots \right\},$$

Potential energy function:

$$V(\varphi, \theta) = \frac{1}{2} \mu^4 \sin^2(\theta/f) e^{-2\varphi/f}$$

# Model

R.B. and J. Fröhlich, arXiv:2004.10025

Unified  
DM/DE

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Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

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berger

Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

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# Equations of Motion

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Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

Variational equations in an expanding cosmological background:

$$\ddot{\varphi} + 3H\dot{\varphi} = \left[ \frac{2}{f}\Lambda + \frac{\mu^4}{f}\sin^2\frac{\theta}{f} \right] e^{-2\varphi/f},$$
$$\ddot{\theta} + 3H\dot{\theta} = -\frac{\mu^4}{f}\sin\left(\frac{\theta}{f}\right)\cos\left(\frac{\theta}{f}\right)e^{-2\varphi/f},$$

Roles of the fields:

- **Dark matter:**  $\theta$  oscillates about  $\theta = 0$ .
- **Dark energy:**  $\varphi$  rolls at late times towards  $\infty$ .

# Equations of Motion

Unified  
DM/DE

R. Branden-  
berger

Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

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# Dark Energy Phase

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Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

After sufficient relaxation of  $\theta$ :

$$\frac{1}{2}\mu^4 \sin^2(\theta/f) \ll \Lambda$$

→  $\varphi$  evolves as a standard quintessence field:

$$\varphi(t) = f \ln(\beta t),$$

$$\beta^2 \simeq \frac{4}{3} \frac{\Lambda}{f^2} \left(\frac{m_{pl}}{f}\right)^2,$$

Leads to an equation of state  $w \equiv p/\rho$  with

$$w \simeq -1 + \frac{4}{3} \left(\frac{m_{pl}}{f}\right)^2.$$

→  $\varphi$  behaves as DE for  $f > m_{pl}$ .

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Unified  
DM/DE

R. Branden-  
berger

Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

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Unified  
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Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

## While

$$\frac{1}{2}\mu^4 \sin^2(\theta/f) > \Lambda$$

- The axion dominates the energy-momentum tensor
- The axion undergoes damped oscillations with amplitude  $\mathcal{A}(t)$

$$\mathcal{A}(t) \sim a(t)^{-3/2} \sim T(t)^{3/2}.$$

- The axion dominates the source term in the equation of motion for  $\varphi$
- Resulting evolution of  $\varphi$ :

$$\varphi(t) = \alpha \ln\left(\frac{t}{t_{eq}}\right) + \varphi(t_{eq}).$$

# Dark Matter Phase

Unified  
DM/DE

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Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

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# Early Phase

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Introduction

Standard  
Paradigm

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Unified Dark  
Sector

Baryogenesis

Conclusions

Assume an early confining force which traps  $\theta$  at  $\theta = -\pi f/2$ .

→ **early slow-rolling phase** of  $\theta$ :

$$\theta = -\frac{\pi f}{2} + \Delta\theta.$$

$$(\Delta\theta)'' + \frac{3}{2t}(\Delta\theta)' \simeq \frac{\mu^4}{f^2} e^{-2\varphi/f} \Delta\theta.$$

$$\Delta\theta(t) \sim \frac{t^{1/2}}{t_c^{1/2}} \Delta\theta(t_c),$$



# Early Phase II

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Sector**

Baryogenesis

Conclusions

Slow rolling of  $\varphi$  during the early phase:

$$e^{-2\varphi/f} = \frac{f^2}{2\mu^4} \frac{1}{t^2}.$$

# Numerical Evolution: $\theta$ field

Thanks to Z. Wang for the numerical work

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Introduction

Standard  
Paradigm

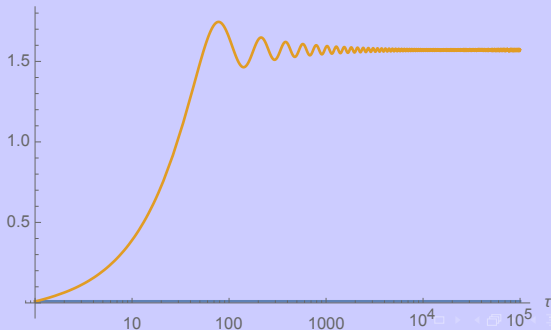
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**Unified Dark  
Sector**

Baryogenesis

Conclusions

Rescaled Field Value



# Numerical Evolution - Field Values

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Introduction

Standard  
Paradigm

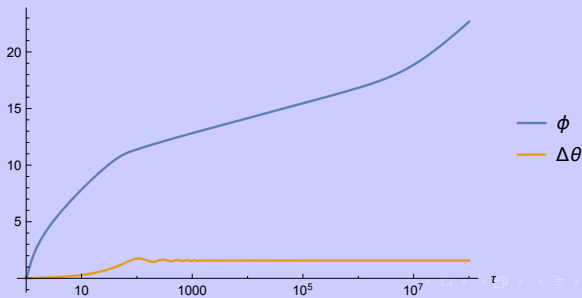
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Unified Dark  
Sector

Baryogenesis

Conclusions

Rescaled Field Value



# Numerical Evolution - Equation of State

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Introduction

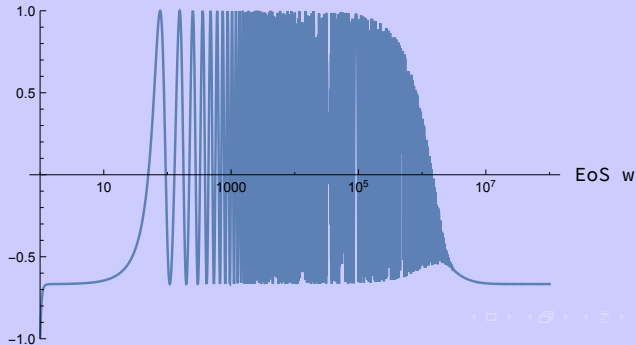
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Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions



# Parameter Requirements

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Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

Parameters:  $\Lambda, \mu, T_C, f$

- $f$  determines the equation of state in the DE phase:  
 $f \sim m_{pl}$ .
- $\Lambda$  has the right value to yield DE today.
- Note: no solution of the coincidence problem!
- $\mu$  and  $T_C$  chosen such that we get the observed matter energy density today.
- Note: dark matter mass set by  $\mu$

$$m_{DM} \sim \frac{\mu^2}{f}$$

# Parameter Requirements

Unified  
DM/DE

R. Branden-  
berger

Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

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Unified  
DM/DE

R. Branden-  
berger

Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

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# Parameter Choices

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Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

**Unified Dark  
Sector**

Baryogenesis

Conclusions

$$m_{DM} \equiv m_a 1\text{eV},$$

$$\begin{aligned}\mu &\sim m_a^{1/2} 10^5 \text{GeV} \\ T_c &\sim m_a 10^{14} \text{GeV}.\end{aligned}$$



# Plan

Unified  
DM/DE

R. Branden-  
berger

Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

**Baryogenesis**

Conclusions

- 1 Introduction
- 2 Standard Paradigm in Trouble
  - Dark Energy cannot be  $\Lambda$
  - Challenges for the WIMP Paradigm
- 3 Unified Model of DE, DM and Baryogenesis
- 4 **Baryogenesis**
- 5 Discussion and Conclusions

# Coupling to the Baryon Current

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Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

Assume a coupling between  $\zeta$  and the baryon current:

$$\delta\mathcal{L} = \tilde{\alpha}\partial_\mu\mathfrak{S}\zeta j_B^\mu,$$

Can be generated by the chiral anomaly for the baryon current:

$$\partial_\mu j_B^\mu \sim \frac{g^2}{16\pi^2} F \wedge F,$$

# Spontaneous Baryogenesis

A. Cohen and D. Kaplan, Phys. Lett. 199B, 251 (1987)

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Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
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Unified Dark  
Sector

Baryogenesis

Conclusions

During the early phase when both  $\varphi$  and  $\theta$  are uniformly rolling a chemical potential for baryon number is induced:

$$\mu_B = \tilde{\alpha}(\mathfrak{S}\zeta) \cdot = \frac{\tilde{\alpha}}{f} \left[ \dot{\theta} - \frac{\dot{\varphi}}{f} \theta \right] e^{-\varphi/f}.$$

Evaluation:

$$\mu_B \sim \tilde{\alpha} \frac{T^4}{\mu^2 m_{pl}},$$

Induced baryon number density:

$$n_B \sim \mu_B T^2,$$

Evaluation:

$$\frac{n_B}{s} \sim \tilde{\alpha} \left( \frac{T}{\mu} \right)^2 \frac{T}{m_{pl}},$$

# Baryogenesis from Hypermagnetic Helicity

Unified  
DM/DE

R. Branden-  
berger

Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

**Baryogenesis**

Conclusions

Baryon current anomaly  $\rightarrow$

$$\Delta N_B = C_y \frac{\alpha_y}{8\pi} \Delta \mathcal{H},$$

Helicity density:

$$\dot{h} = -2 \langle \mathbf{E} \cdot \mathbf{B} \rangle,$$

Magnetohydrodynamics  $\rightarrow$

$$\mathbf{E} \cdot \mathbf{B} = \frac{1}{\sigma} \mathbf{B} \cdot (\nabla \wedge \mathbf{B}),$$

$\sigma \sim T$ : conductivity

$$\langle \mathbf{B} \cdot (\nabla \wedge \mathbf{B}) \rangle = \int_k \frac{d^3 k}{(2\pi)^3} |k|^3 (|A_{k,+}|^2 - |A_{k,-}|^2),$$

# Baryogenesis from Hypermagnetic Helicity

Unified  
DM/DE

R. Branden-  
berger

Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

**Baryogenesis**

Conclusions

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# Growth of Helicity

Unified  
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Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

Assume the coupling

$$\delta\mathcal{L}_2 = \alpha \frac{\mathfrak{S}\zeta}{4} \tilde{Y}_{\mu\nu} Y^{\mu\nu},$$

Equation of motion for the gauge field becomes:

$$\ddot{A}_{k,\pm} + (k^2 \pm \alpha k (\mathfrak{S}\zeta)') A_{k,\pm} = 0,$$

$$\ddot{A}_k + (k^2 \pm \alpha k e^{-\varphi/f} \frac{\dot{\theta}}{f}) A_k = 0.$$

Exponential growth for  $k < k_c$ :

$$k_c = \alpha \tau^{-1} \left( \frac{T_c}{\mu} \right)^2.$$

# Resulting Baryon to Entropy Ratio

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Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

Duration of the instability is set by **back-reaction**: density in the produced gauge particles must remain smaller than the pre-existing radiation density.

Resulting **baryon to entropy ratio**:

$$\frac{\Delta n_B}{s} \sim C_y \frac{\alpha_y}{8\pi} \alpha^{-1/2}.$$

# Plan

Unified  
DM/DE

R. Branden-  
berger

Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

- 1 Introduction
- 2 Standard Paradigm in Trouble
  - Dark Energy cannot be  $\Lambda$
  - Challenges for the WIMP Paradigm
- 3 Unified Model of DE, DM and Baryogenesis
- 4 Baryogenesis
- 5 Discussion and Conclusions



# Embedding of the Model in a UV Theory

Unified  
DM/DE

R. Branden-  
berger

Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

- Example:  $\zeta$  is the **axion-dilaton field** of string theory:

$$\zeta = e^{-\Phi} + iC_0$$

- $\Phi$ : dilaton;  $C_0$ : Ramond-Ramond zero form axion.
- Classical level: potential flat in angular direction.
- Quantum level: instantons break the symmetry and generate a potential.
- Example:  $\zeta$  is a complex scalar field whose real part describes the size of the compact space.
- Negative curvature of the internal manifold  $\rightarrow$  positive exponential potential for the real part.

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Unified  
DM/DE

R. Branden-  
berger

Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

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# Conclusions

Unified  
DM/DE

R. Branden-  
berger

Introduction

Standard  
Paradigm

DE is not  $\Lambda$   
No WIMP

Unified Dark  
Sector

Baryogenesis

Conclusions

- **Unified model of dark energy and dark matter**
- **Complex scalar field  $\zeta$**
- **Radial component** yields **dark energy**.
- **Angular component** yields an **axion** acting as **dark matter**.
- **Early evolution: uniform motion of  $\zeta \rightarrow$  spontaneous breaking of CPT symmetry  $\rightarrow$  **spontaneous baryogenesis**.**