

facts

Lecture 3: Experimental probes and signatures

Experimental probes and signatures

Lecture 4: Experimental aspects

Accelerator
Detectors

Signatures of the Quark-Gluon Plasma

Bose-Einstein interferometry

J/ψ suppression

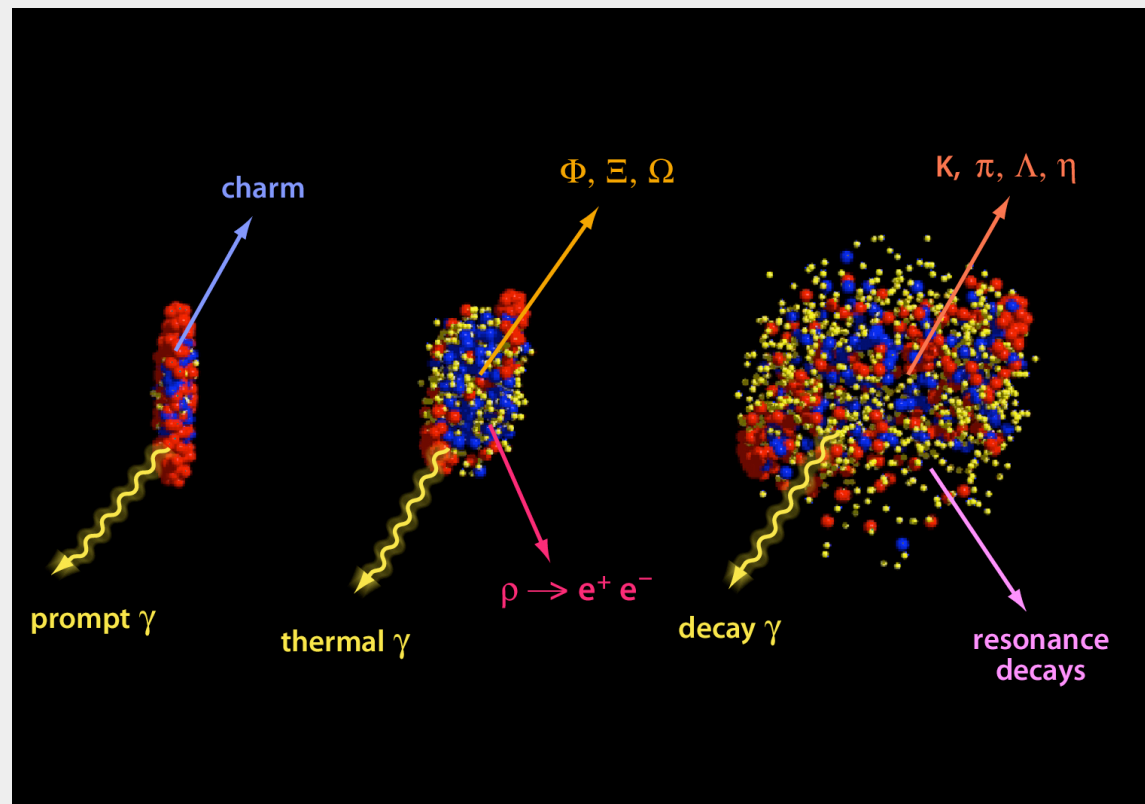
Jet quenching

Collective flow

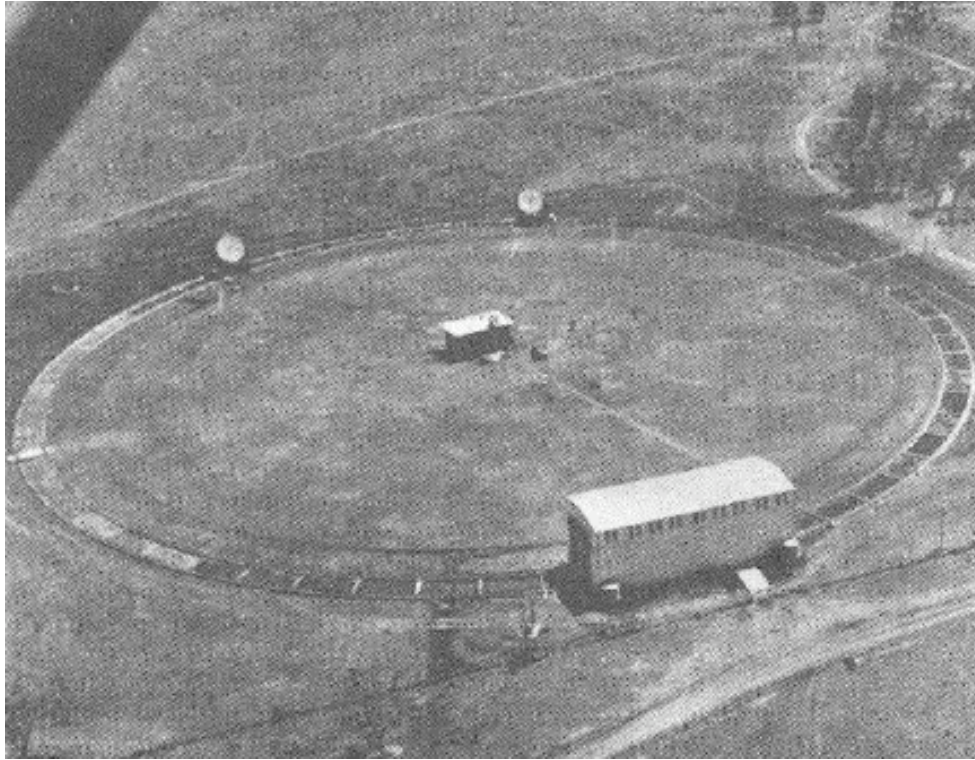
Direct photons

Dilepton production

Strange particle production



Bose-Einstein Correlations

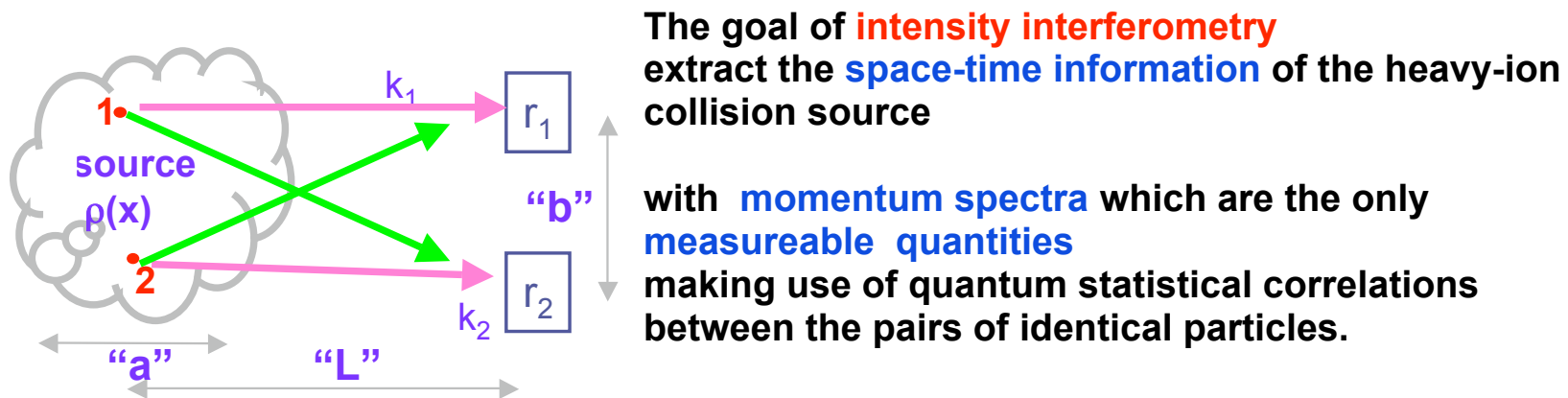


Robert Hanbury Brown and Richard Twiss

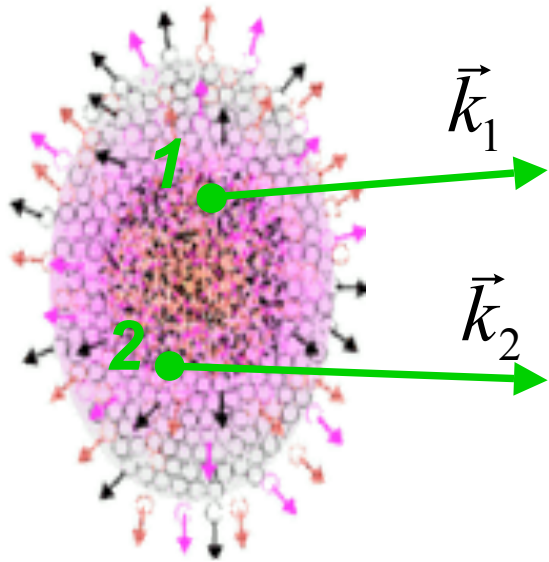
HBT Intensity Interferometry

Interference is a phenomenon associated with the superposition of two or more waves.

The two-particle correlations arise from the interference of particle wave-functions
... depend on whether the particles are bosons or fermions



Handbury, Brown and Twiss measured the diameter of Sirius
Nature 178(1956)1046



emission amplitud of two pions

$$A_{\pi\pi} = e^{i\vec{k}_1\vec{r}_1 + i\vec{k}_2\vec{r}_2} + e^{i\vec{k}_1\vec{r}_2 + i\vec{k}_2\vec{r}_1}$$

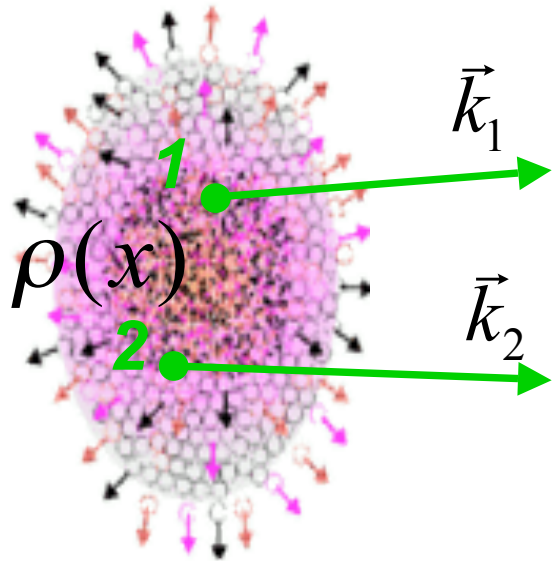
$$|A_{\pi\pi}|^2 = \left| e^{i\vec{k}_1\vec{r}_1 + i\vec{k}_2\vec{r}_2} + e^{i\vec{k}_1\vec{r}_2 + i\vec{k}_2\vec{r}_1} \right|^2$$

$$|A_{\pi\pi}|^2 = 1 + 1 + e^{i(\vec{k}_1 - \vec{k}_2)\vec{r}_1} + e^{i(\vec{k}_2 - \vec{k}_1)\vec{r}_2}$$

$$|A_{\pi\pi}|^2 = 2 + 2\cos(qr) \quad \text{where} \quad \begin{aligned} q &= |\vec{k}_1 - \vec{k}_2| \\ r &= |\vec{r}_1 - \vec{r}_2| \end{aligned}$$

$$\frac{|A_{\pi\pi}|^2}{|A_\pi||A_\pi|} = 1 + \cos(qr)$$

emission amplitud of two pions



$$A_{\pi\pi} = \rho(x)e^{i\vec{k}_1\vec{r}_1+i\vec{k}_2\vec{r}_2} + \rho(x)e^{i\vec{k}_1\vec{r}_2+i\vec{k}_2\vec{r}_1}$$

$$\int \rho(x)dx = 1$$

$$|A_{\pi\pi}|^2 = \rho^* \rho + \rho^* \rho e^{i(\vec{k}_1-\vec{k}_2)\vec{r}_1} + \rho^* \rho e^{i(\vec{k}_2-\vec{k}_1)\vec{r}_2}$$

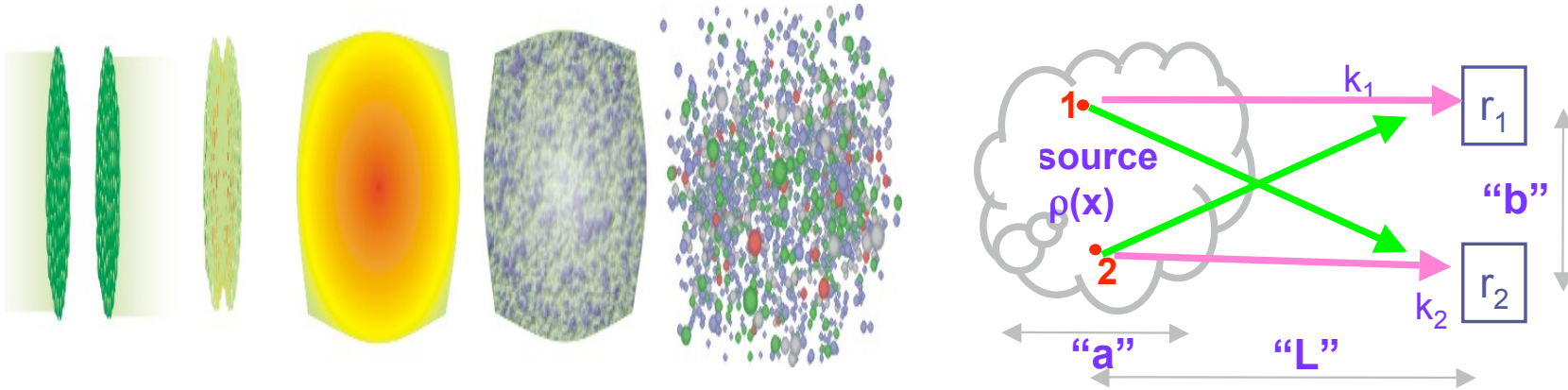
where

$$q = |\vec{k}_1 - \vec{k}_2|$$

$$r = |\vec{r}_1 - \vec{r}_2|$$

$$\frac{|A_{\pi\pi}|^2}{|A_\pi||A_\pi|} = 1 + |F(\rho)|^2$$

Probing source geometry through interferometry



The correlation function is defined as the ratio of the probability for the coincidence of k_1 and k_2 relative to the probability of observing k_1 and k_2 separately :

$$C(k_1, k_2) = \frac{P(k_1, k_2)}{P(k_1)P(k_2)} = 1 + |F(\rho)|^2$$

↑
↑
Measurable
Fourier Transform of pion source

Correlation function constructed experimentally, $C_2(q) = A(q) / B(q)$
 $A(q) \rightarrow$ is the pair distribution in momentum difference $q = p_2 - p_1$ for pairs of particles from the **same event**. $B(q) \rightarrow$ is the corresponding distribution for pairs of particles from **different events**.

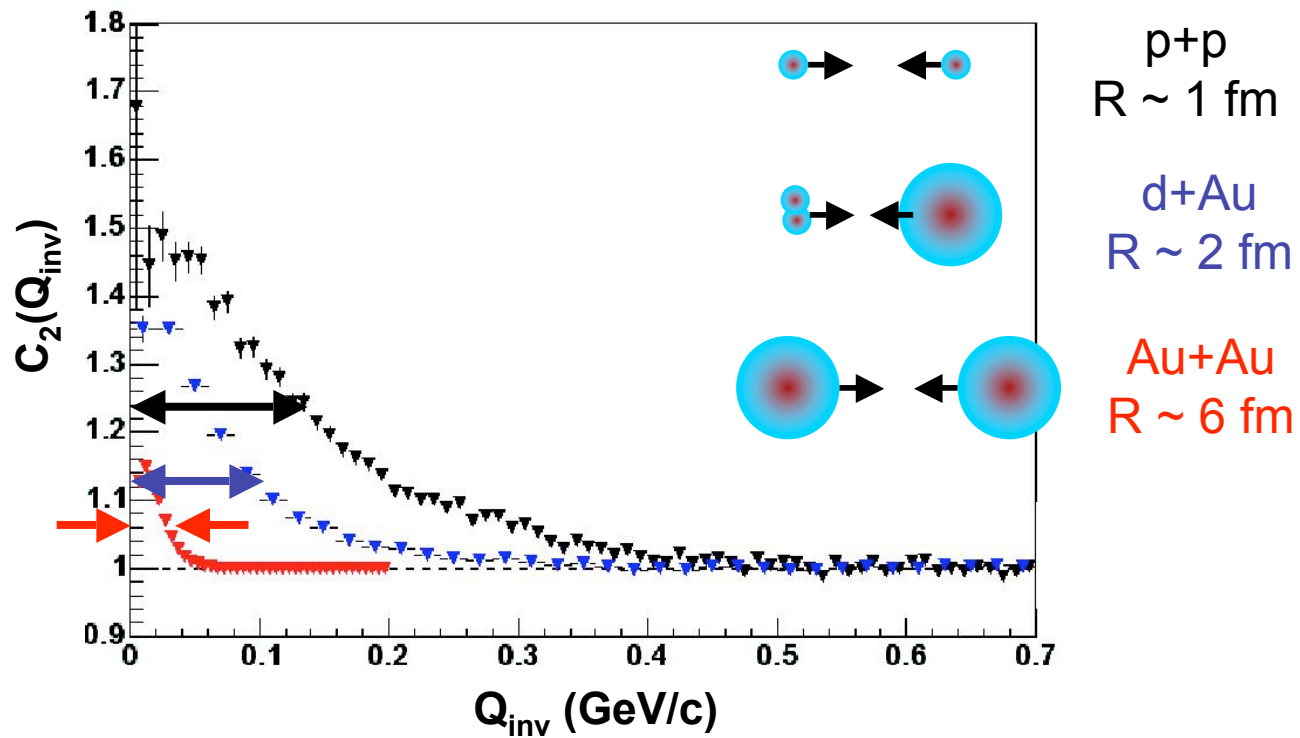
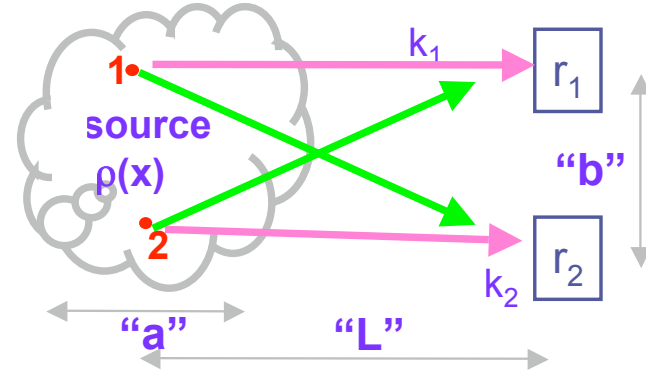
e.g.

$$\rho(x, y, z) = \frac{N}{4\pi^2 R_x R_y R_z \sigma_t} e^{-\left(\frac{x^2}{2R_x^2} + \frac{y^2}{2R_y^2} + \frac{z^2}{2R_z^2} + \frac{t^2}{2\sigma_t^2}\right)}$$

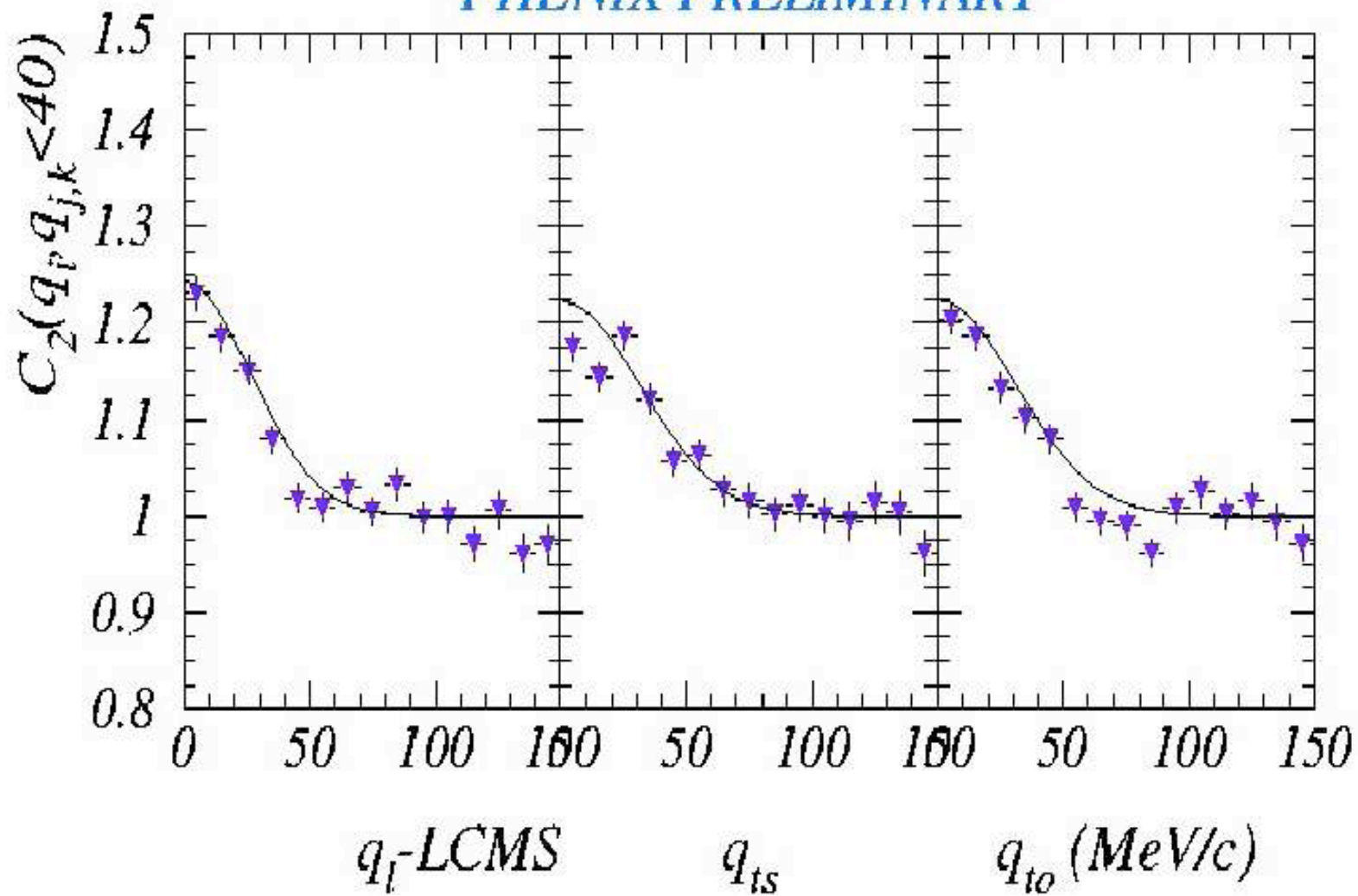
$$F[\rho](q) = N e^{-\left(\frac{R_x^2 q_x^2}{2} + \frac{R_y^2 q_y^2}{2} + \frac{R_z^2 q_z^2}{2} + \frac{\sigma_t^2 q_t^2}{2}\right)}$$

$$C(k_1, k_2) = 1 + |F(\rho)|^2$$

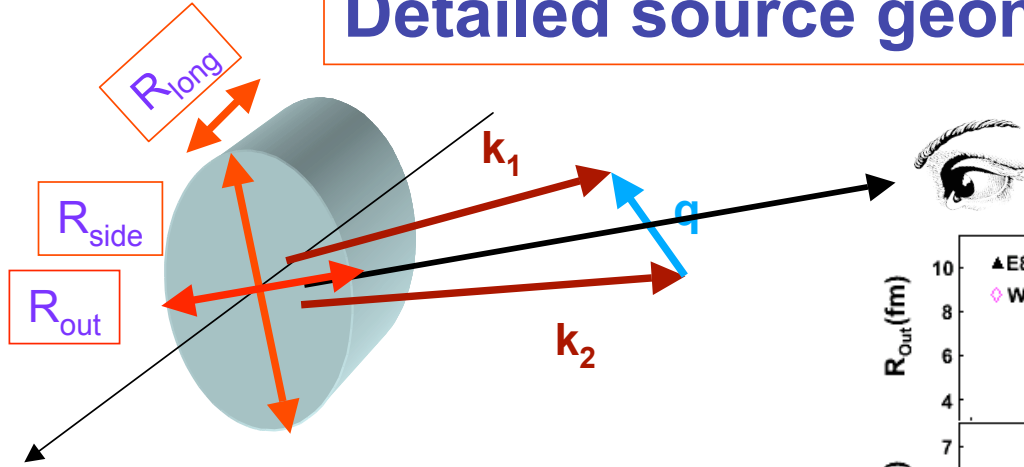
Source geometry



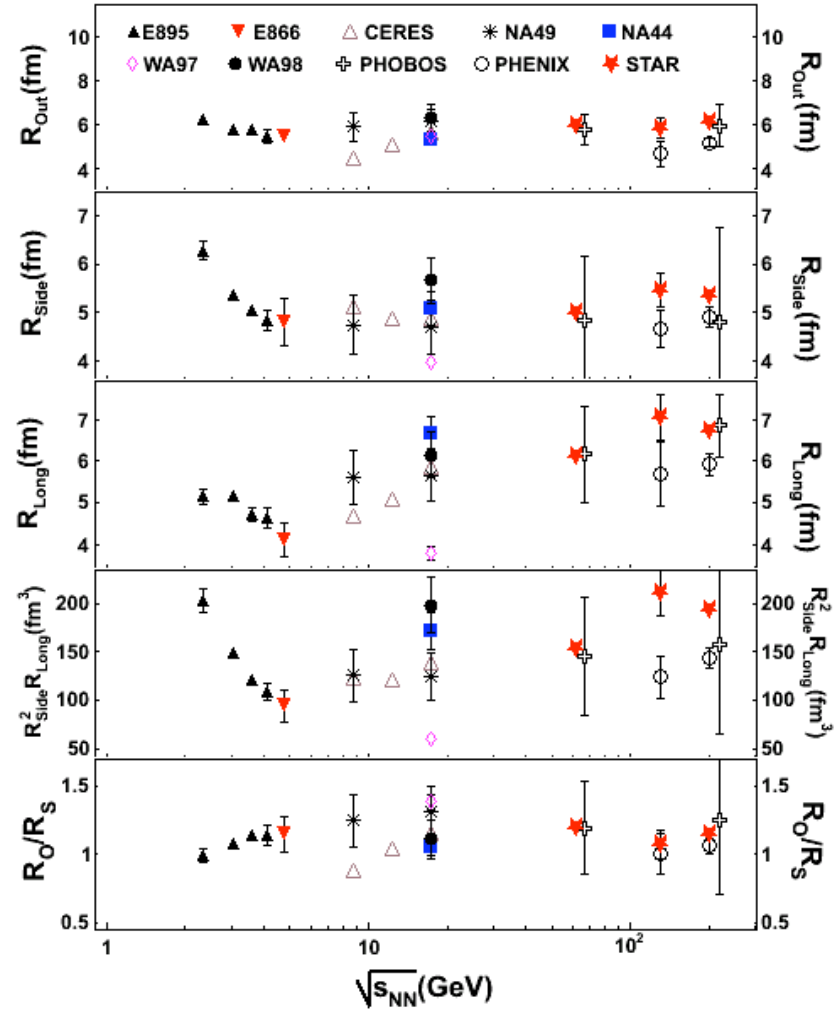
$Au+Au - 2\pi^+$
PHENIX PRELIMINARY



Detailed source geometry

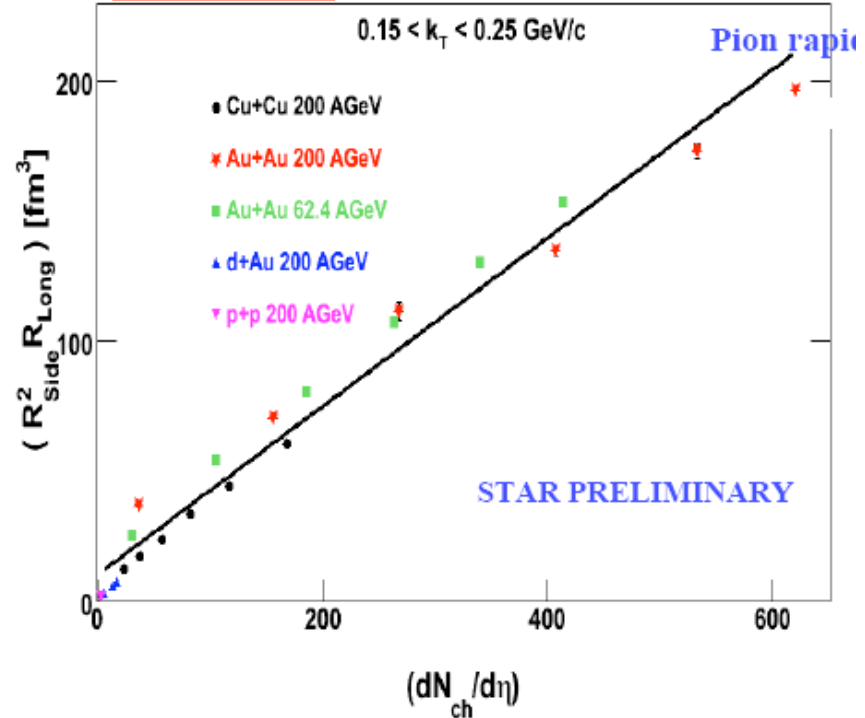
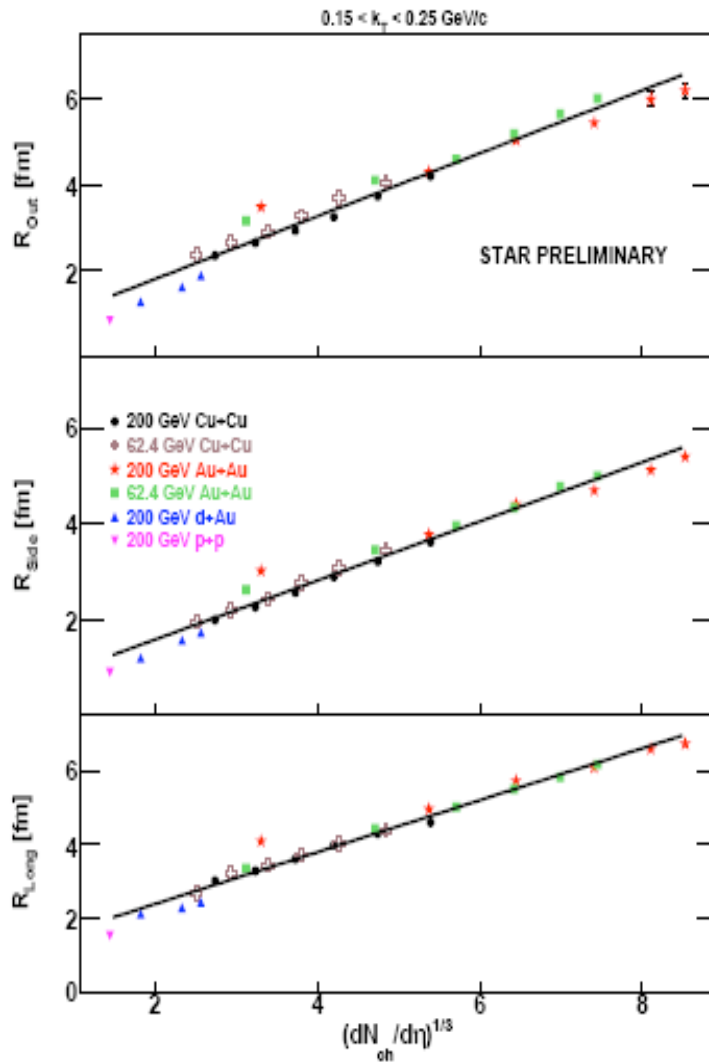


R_{long} – along beam direction
 R_{out} – along “line of sight”
 R_{side} – \perp “line of sight”



Beam energy dependence of pion HBT

STAR



Pion rapidity density is proportional to the freezeout volume => Constant Freezeout Volume (freezeout at a constant density).

Bose Einstein Correlations as a signature

If a Quark Gluon Plasma appears, pions will populate the central rapidity region

$$\text{Entropy of the QGP} = S_{QGP}$$

The Quark Gluon Plasma will hadronize

$$\text{Entropy of the hadronization phase} = S_{had}$$

total entropy

$$volume_{QGP} \times S_{QGP} \leq volume_{had} \times S_{had}$$

$$\text{since } S_{had} < S_{QGP} \quad \text{then } volume_{had} > volume_{QGP}$$

TABLE I. $\pi\pi$ -interference result [4] from the collision of an O beam (200 GeV/nucleon) on a stationary Au target.

	Rapidity Interval	$R_T(fm)$	Gaussian $R_L(fm)$	Λ
Central region (mid-rapidity)	$1 < y < 2$	4.3 ± 0.6	2.6 ± 0.6	$0.34^{+0.09}_{-0.06}$
		$R_T^{side} = 4.0 \pm 1.0 fm$	2.6 ± 0.6	$0.34^{+0.09}_{-0.06}$
	$2 < Y < 3$	$R_T^{out} = 4.4 \pm 1.0 fm$	2.6 ± 0.6	$0.34^{+0.09}_{-0.06}$
		8.1 ± 1.6	$5.6^{+1.2}_{-0.8}$	0.77 ± 0.19
		$R_T^{side} = 6.6 \pm 1.8 fm$	$5.6^{+1.2}_{-0.8}$	0.77 ± 0.19
		$R_T^{out} = 11.2 \pm 2.3 fm$	$5.6^{+1.2}_{-0.8}$	0.77 ± 0.19

$$R(\text{Oxygen}) \cong 3 fm$$

RHIC:

There is only a weak dependence on the energy

Results more or less what was observed at CERN SPS

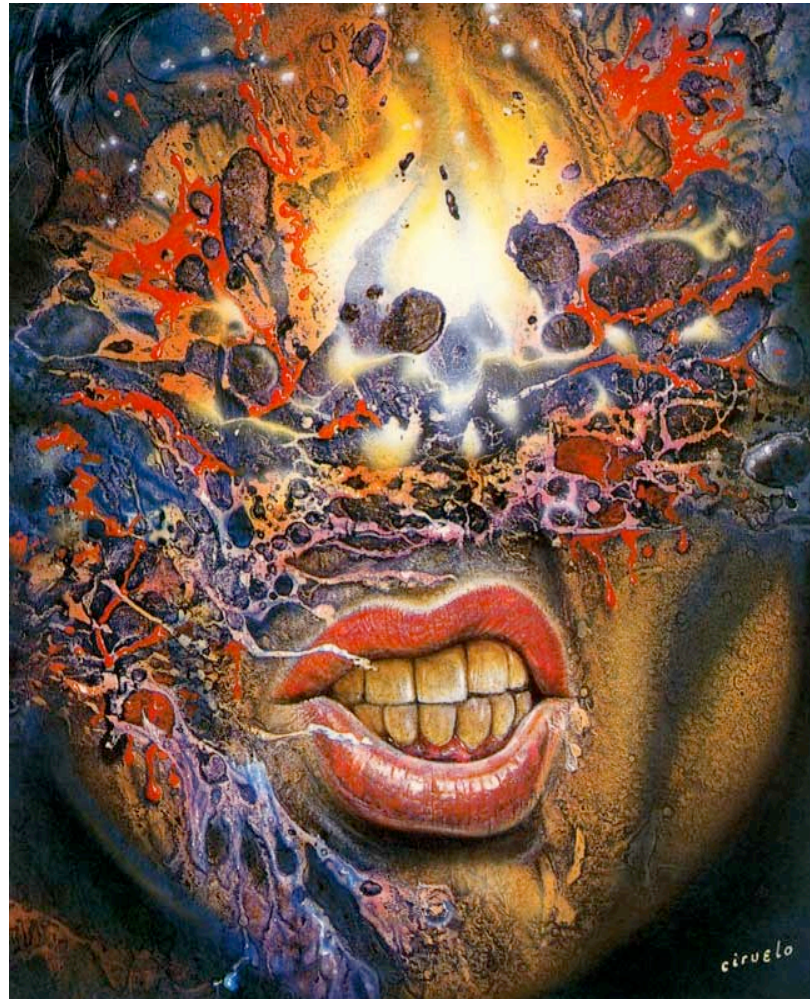
$$R_{out} / R_{side} \approx 1 \quad \text{when most models expect} \quad R_{out} / R_{side} \approx 2$$

no consistent description of data

something is missing in our picture of space time

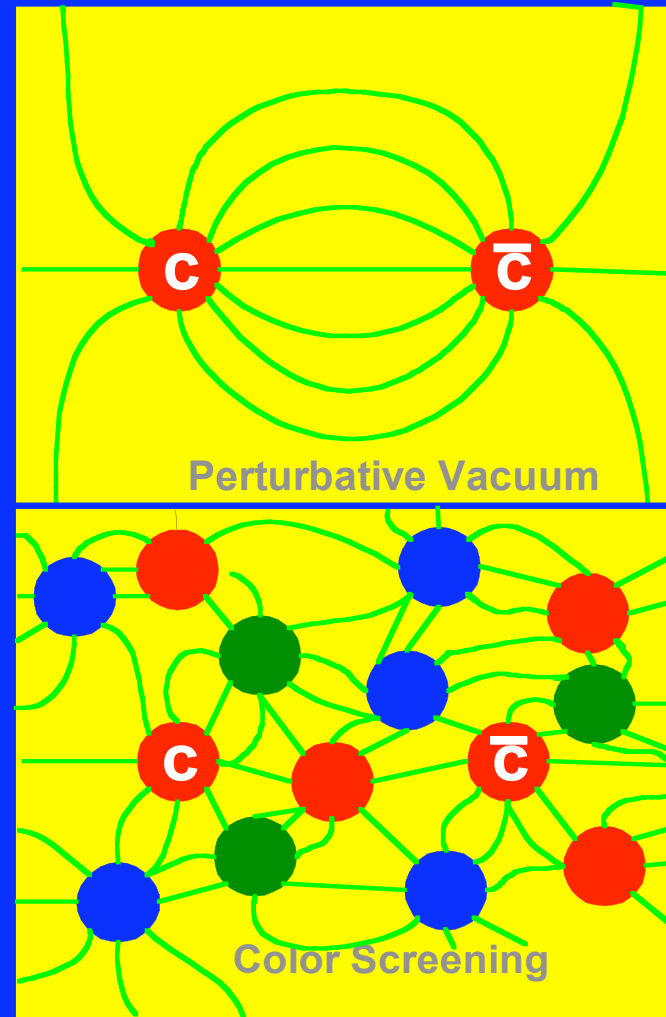
Charm suppression

Strageness enhacement



Suppression of the J/ψ production as a Signature of QGP

- J/ψ suppression because colour screening hinders the quarks from binding
- J/ψ interacts inelastically with some dense hadronic matter created in the collision



J/Ψ in the plasma: the beginning

Volume 178, number 4

PHYSICS LETTERS B

9 October 1986

J/ψ SUPPRESSION BY QUARK–GLUON PLASMA FORMATION ☆

T. MATSUI

*Center for Theoretical Physics, Laboratory for Nuclear Science, Massachusetts Institute of Technology,
Cambridge, MA 02139, USA*

and

H. SATZ

*Fakultät für Physik, Universität Bielefeld, D-4800 Bielefeld, Fed. Rep. Germany
and Physics Department, Brookhaven National Laboratory, Upton, NY 11973, USA*

Received 17 July 1986

If high energy heavy ion collisions lead to the formation of a hot quark–gluon plasma, then colour screening prevents $c\bar{c}$ binding in the deconfined interior of the interaction region. To study this effect, the temperature dependence of the screening radius, as obtained from lattice QCD, is compared with the J/ψ radius calculated in charmonium models. The feasibility to detect this effect clearly in the dilepton mass spectrum is examined. It is concluded that J/ψ suppression in nuclear collisions should provide an unambiguous signature of quark–gluon plasma formation.

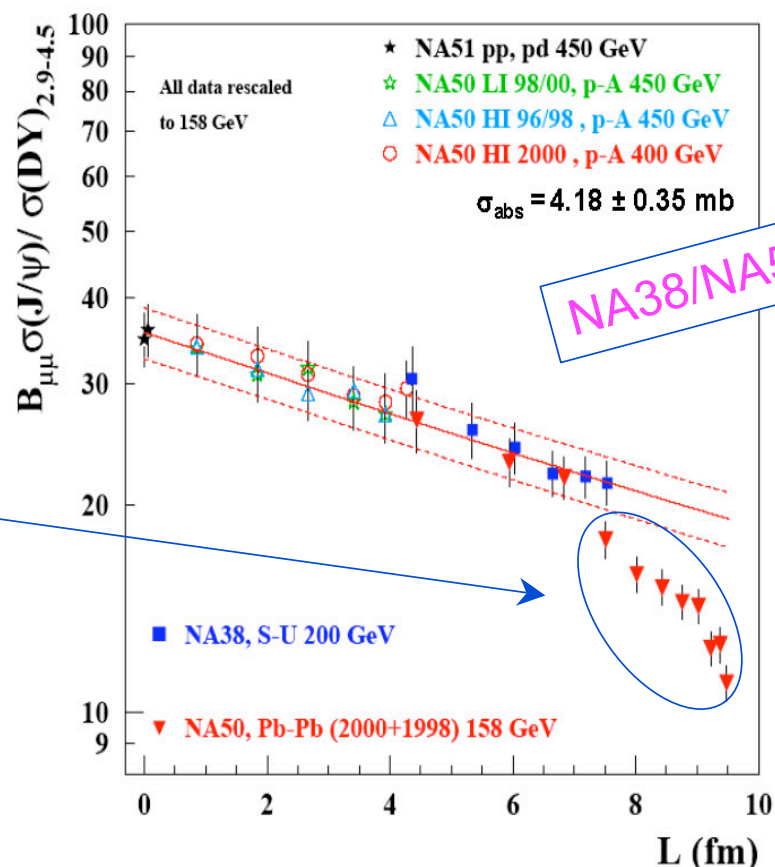
Anomalous suppression

In peripheral Pb-Pb collisions the $(J/\psi)/DY$ ratio is consistent with the normal suppression pattern

In central Pb-Pb collisions ($b < 8-8.5$ fm) a much stronger suppression is observed:

Anomalous suppression

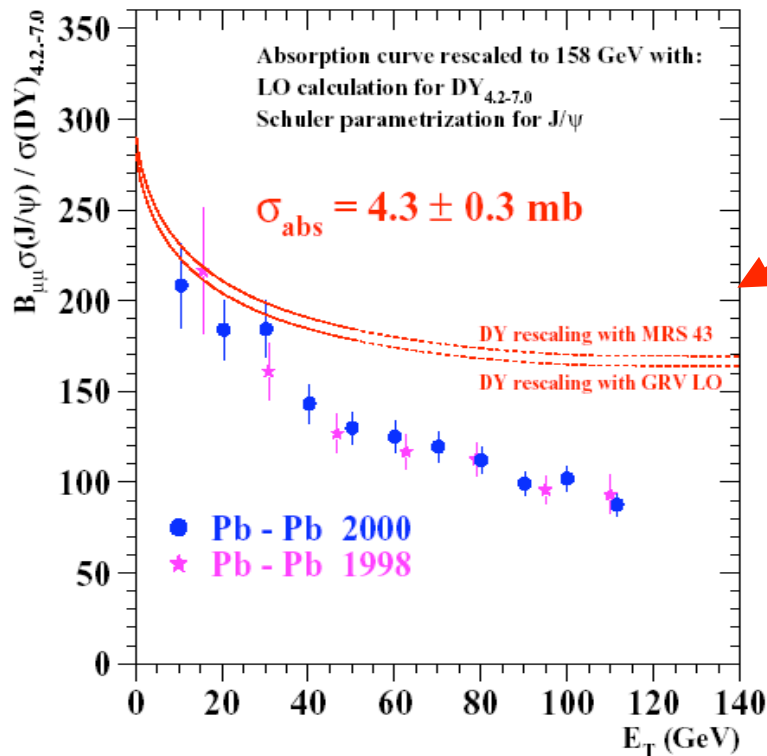
New In-In data follow the same pattern ! (NA60)



Anomalous J/Ψ suppression in A+A

Heavy flavor sector reflects the actual dynamics since heavy hadrons can only be formed in the very early phase of heavy-ion collisions !

■ Anomalous J/Ψ suppression in A+A (NA38/NA50/NA60)

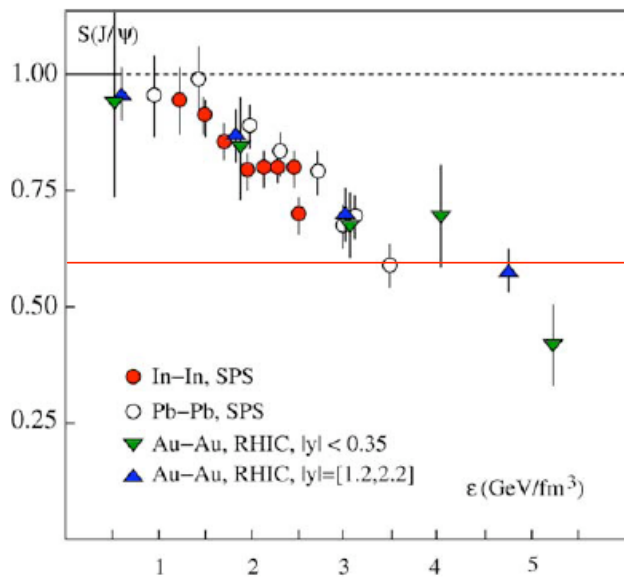


J/Ψ , normal' absorption
by nucleons
(Glauber model)

Experimental observation:
extra suppression in A+A collisions; increasing with centrality

Is that due to Debye screening of the confining potential in a QGP ?

or may be



Multiple rescattering ?

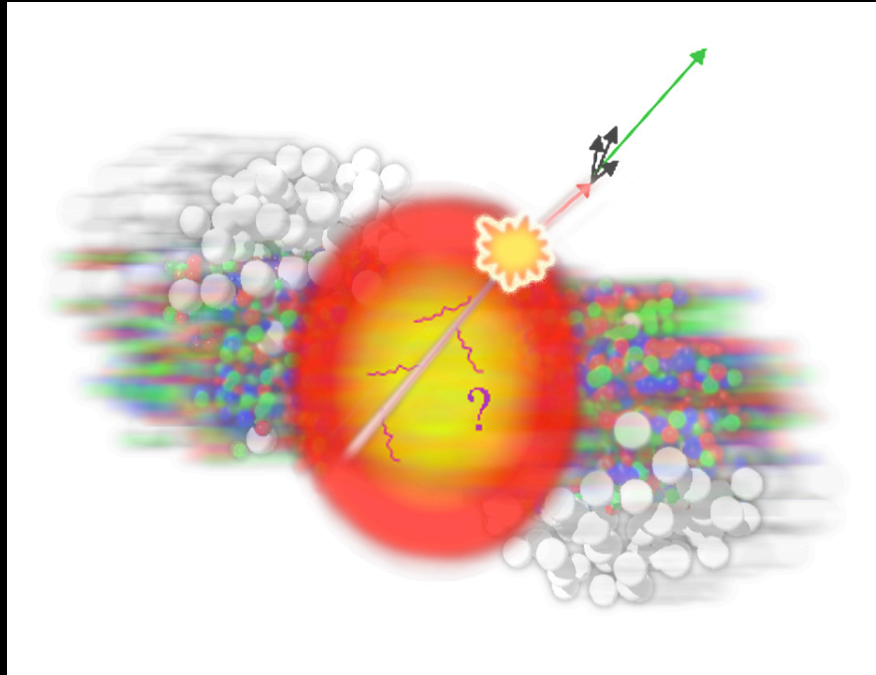
Changes in the gluon distribution function ?

Affected by other decays (excited state decays) ?

Affected by heavy quark/gluon energy loss ?

What is the production mechanism at RHIC energies (color singlet mode, color octet mode, color evaporation model) contribution from gluon fusion, heavy quark fragmentation, gluon fragmentation ?

Jet quenching



“Jet Quenching” – J.D.Bjorken 1982

Energy Loss of Energetic Partons in Quark-Gluon Plasma:
Possible Extinction of High p_T Jets in Hadron-Hadron Collisions.

J. D. BJORKEN
Fermi National Accelerator Laboratory
P.O. Box 500, Batavia, Illinois 60510

Abstract

High energy quarks and gluons propagating through quark-gluon plasma suffer differential energy loss via elastic scattering from quanta in the plasma.

...

An interesting signature may be events in which the hard collision occurs near the edge of the overlap region, with one jet escaping without absorption and the other fully absorbed.

FERMILAB-Pub-82/59-THY
August, 1982

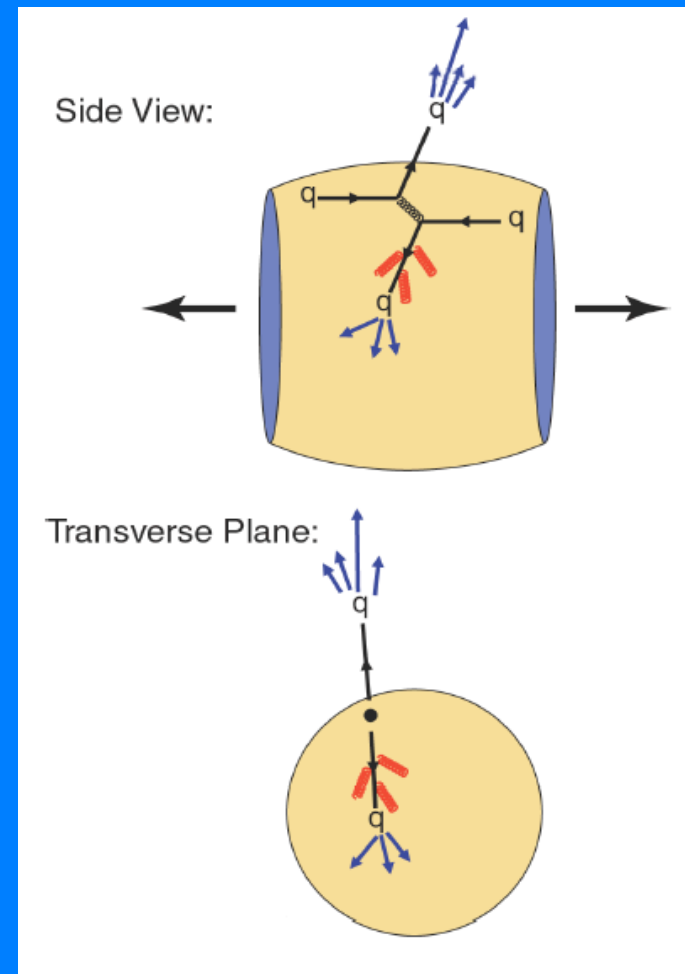
Jet Quenching as a Signature of QGP

leading
particle

hadrons

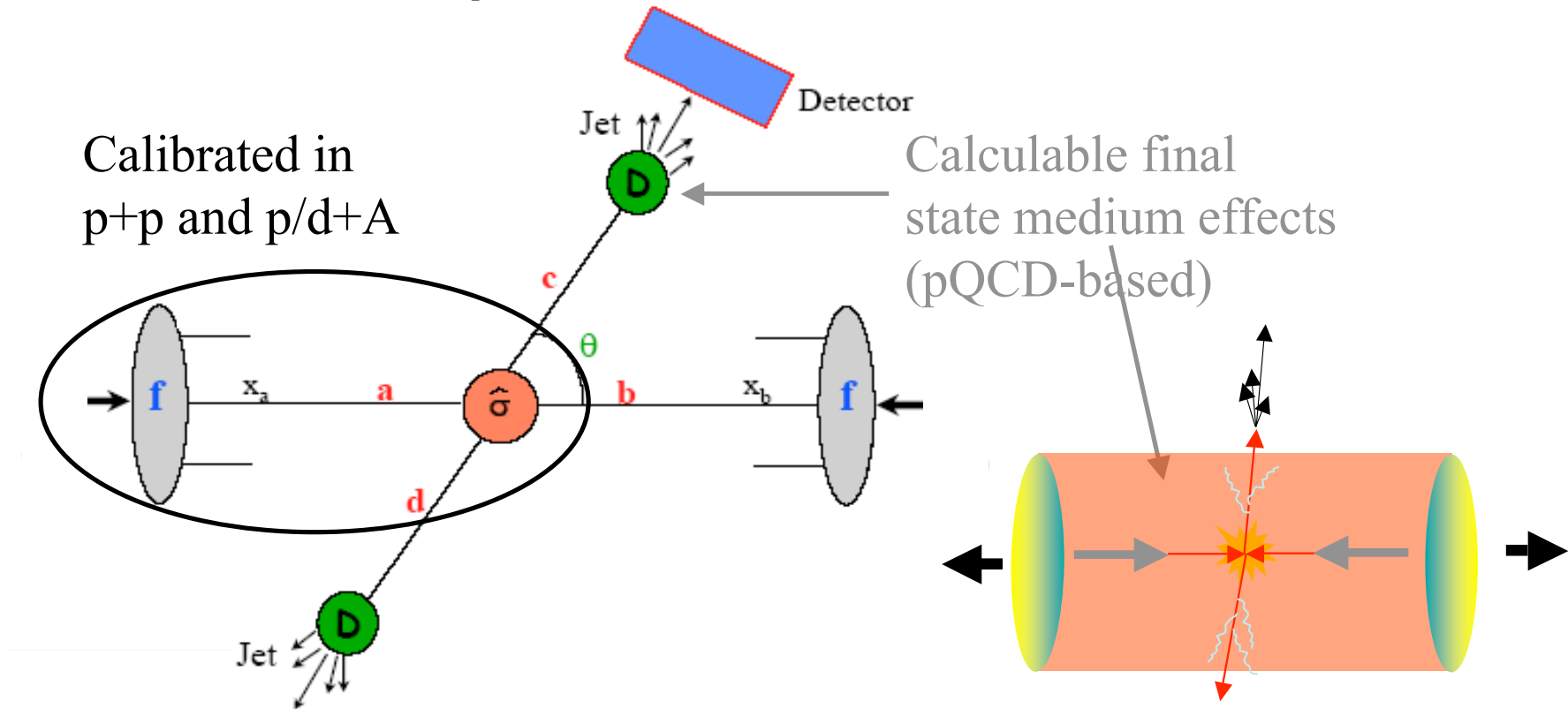
- Hard scatterings (HS) in nucleon collisions produce jets of particles
- In a colour deconfined medium the **partons** (quarks+gluons) strongly interact and lose energy ($\sim \text{GeV}/\text{fm}$) by gluon radiation
- HS near the surface can give a jet in one direction, while the other side is quenched

leading
particle



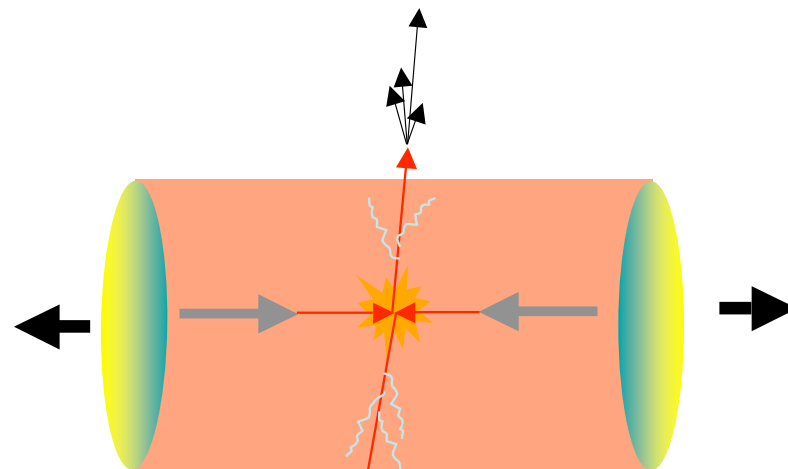
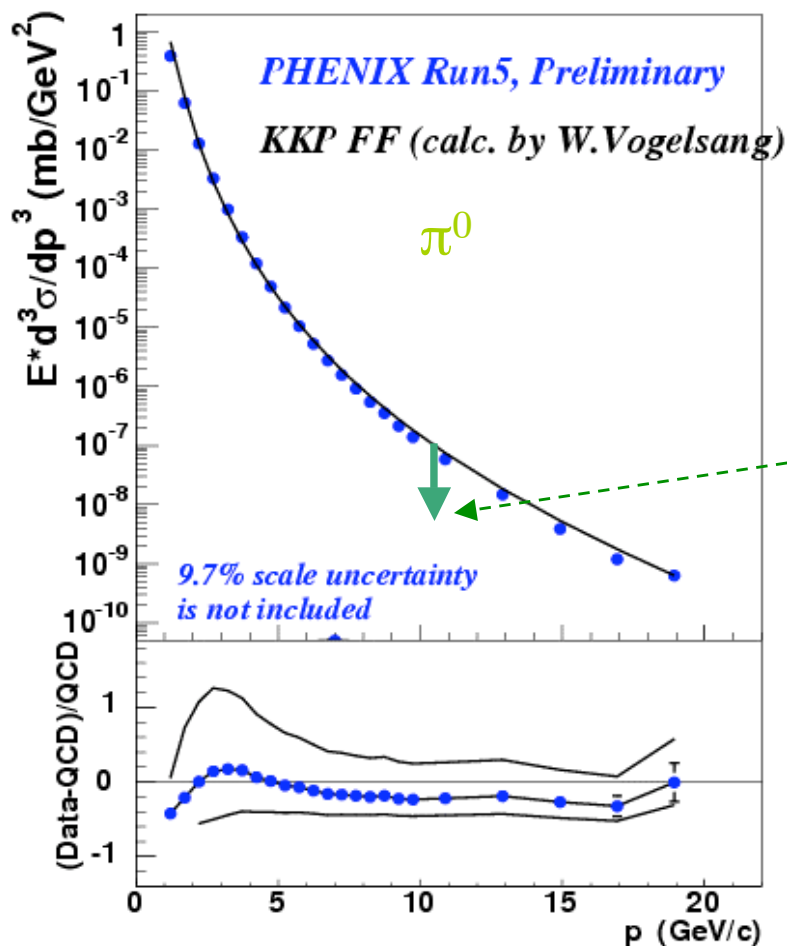
hadrons

Hard probes of QCD matter



Calculable interactions of energetic partons with the medium
⇒ calibrated, penetrating *tomographic* probes

Jet quenching: hadron suppression



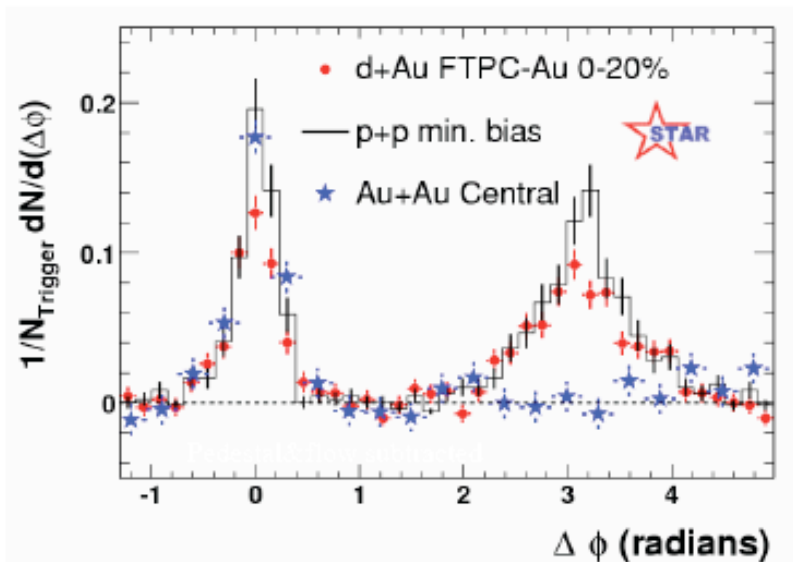
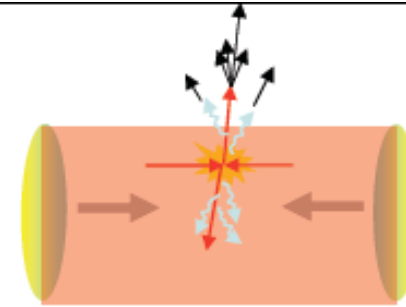
Energy loss \Rightarrow softening of fragmentation \Rightarrow suppression of leading hadron yield

$$R_{AA}(p_T) = \frac{d^2 N^{AA} / dp_T d\eta}{T_{AA} d^2 \sigma^{NN} / dp_T d\eta}$$

Binary collision scaling

p+p reference

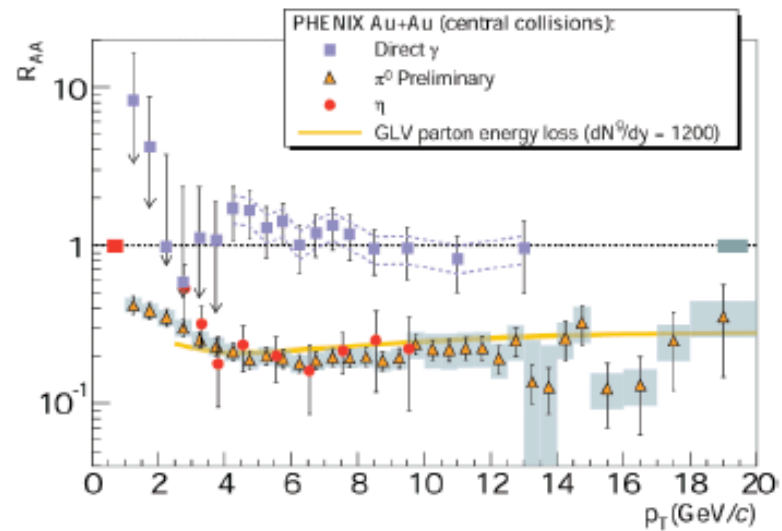
Jets Modified by the Medium



Away side suppression

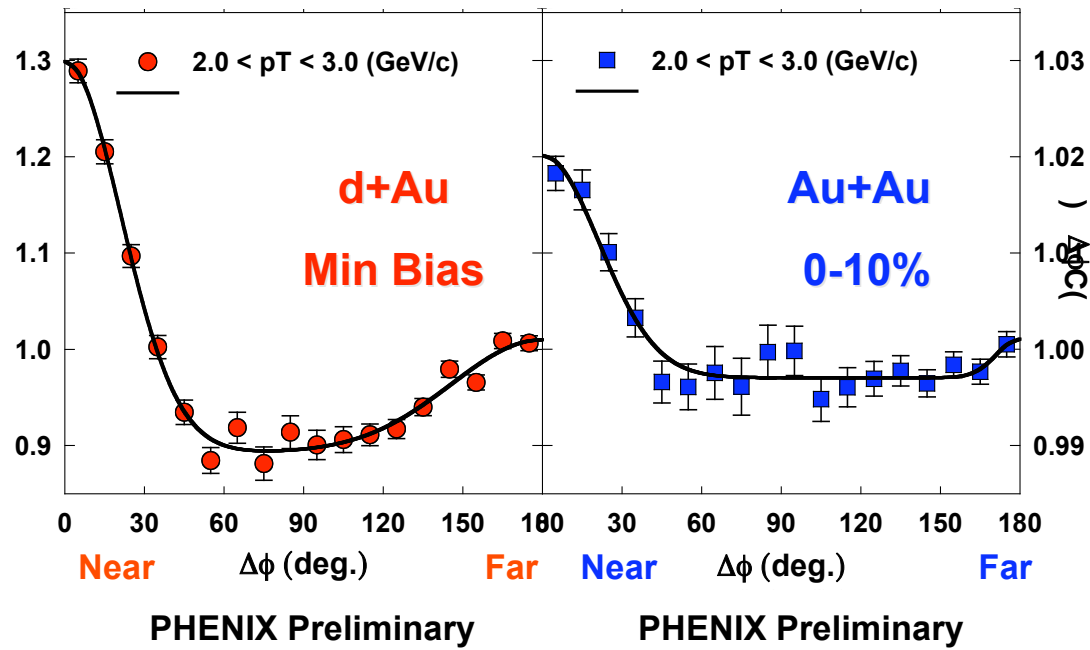
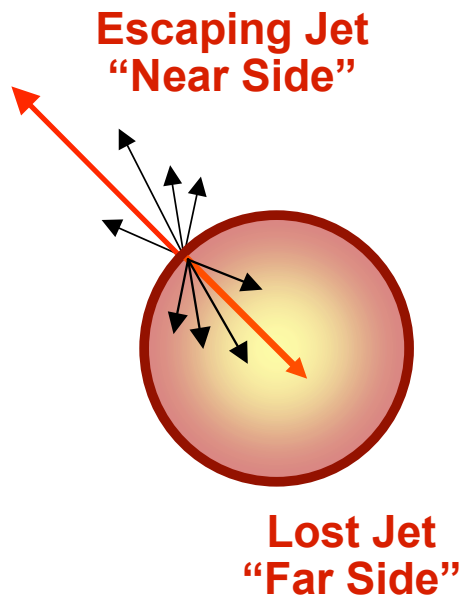
$4 < p_T(\text{trig}) < 6 \text{ GeV}/c$

$p_T(\text{assoc}) > 2 \text{ GeV}/c$



PHENIX

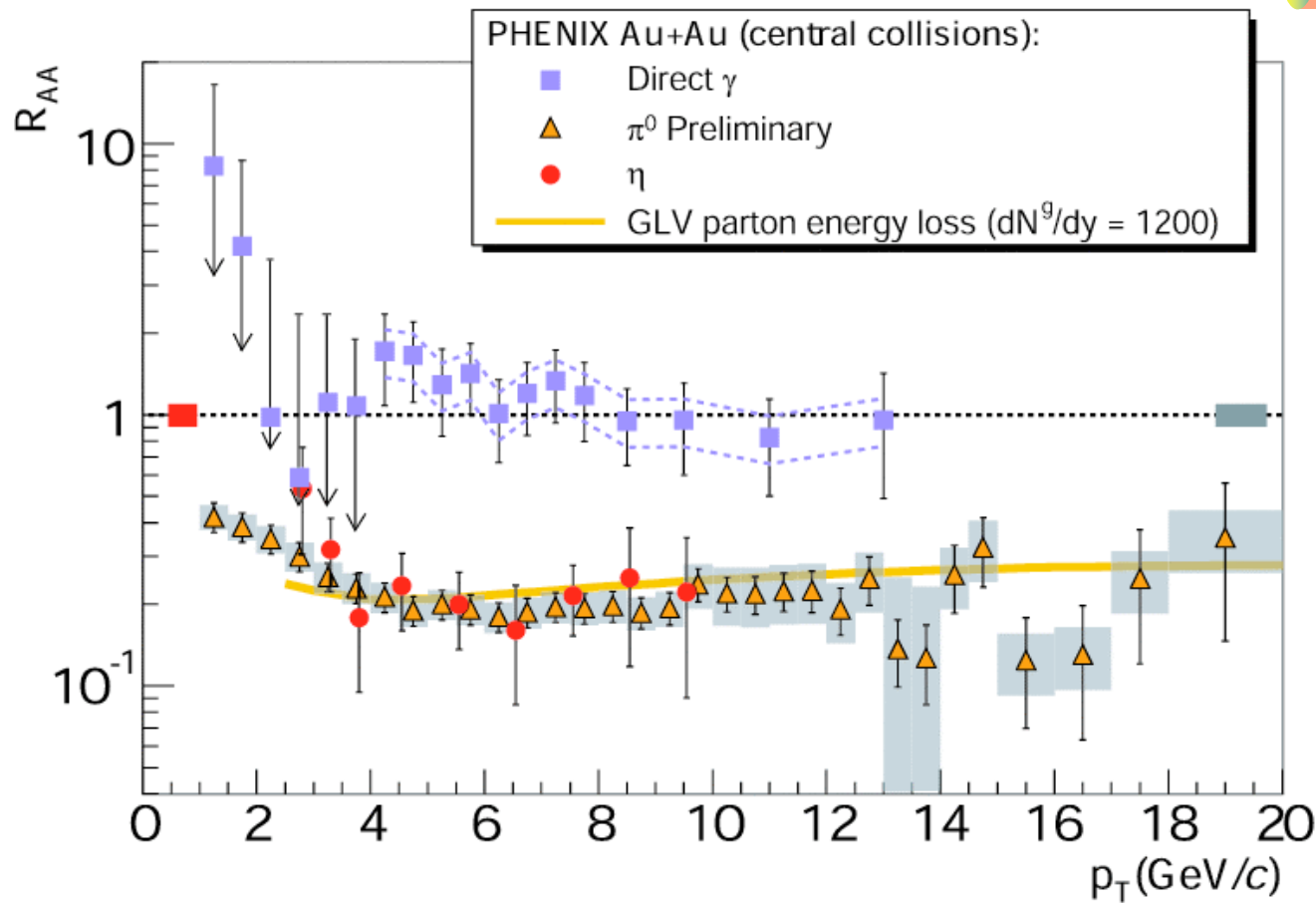
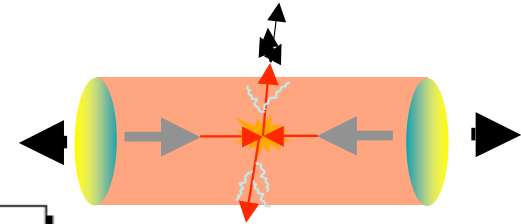
The “Away-Side” Jet



- Jets produced on the periphery of the collision zone coming out should survive.
- However, their partner jet will necessarily be pointed into the collision zone and be absorbed.
- Peripheral Au+Au similar to d+Au
- Central Au+Au shows distinct reduction in far side correlation.
- Away-side Jet is missing in Au+Au

Jet quenching I: hadrons are suppressed, photons are not

$$R_{AA}(p_T) = \frac{d^2 N^{AA} / dp_T d\eta}{T_{AA} d^2 \sigma^{NN} / dp_T d\eta}$$



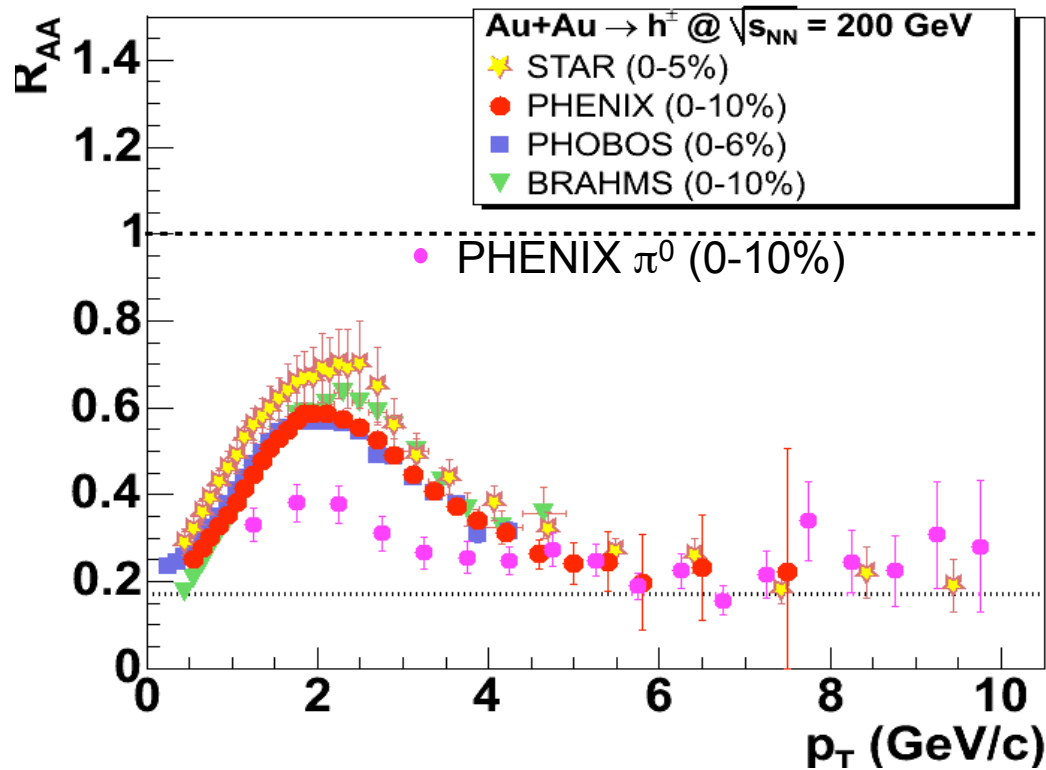
Photons (color-neutral)

Jets (color-charged)

Discoveries at RHIC: high- p_t suppression

- Nuclear modification factor of p_t distributions:

$$R_{AA}(p_t) = \frac{1}{\langle N_{coll} \rangle_C} \times \frac{d^2 N_{AA}^C / dp_t d\eta}{d^2 N_{pp} / dp_t d\eta}$$

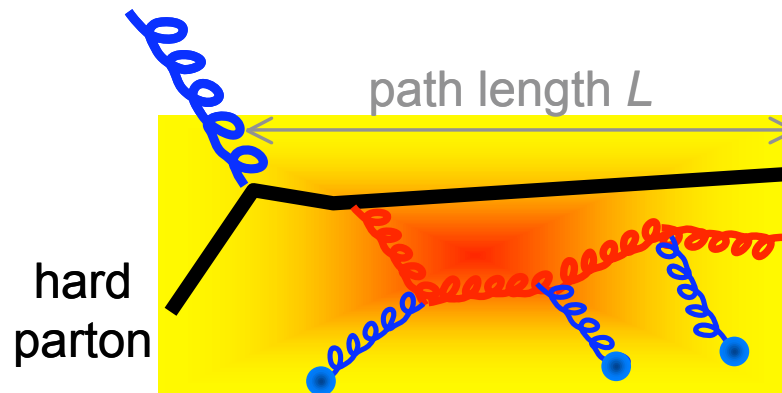


factor 5
suppression!



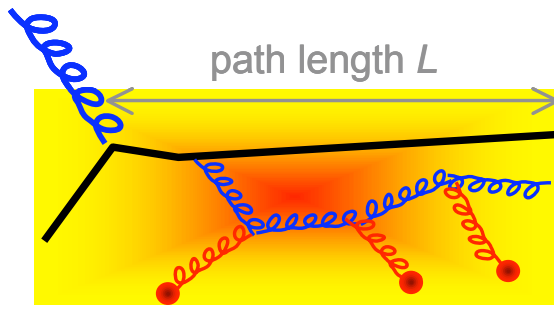
Parton Energy Loss

- Partons travel ~ 4 fm in the high colour-density medium
- Bjorken ('82): energy loss due to elastic scattering
- Successive calculations ('92 \rightarrow): a QCD mechanism dominates, **medium-induced gluon radiation**
- Coherent **wave-function gluon** accumulates k_T due to multiple inelastic scatterings in the medium; it decoheres and is radiated



Bjorken, Gyulassy, Plumer, Wang, Baier, Dokshitzer, Mueller, Pagne, Schiff, Levai, Vitev, Zhakarov, Wang, Salgado, Wiedemann,...

Calculating Parton Energy Loss



$$\langle \Delta E \rangle \approx \int d\omega \omega \frac{dI}{d\omega} \propto \alpha_s C_R \hat{q} L^2 \quad (\text{BDMPS})$$

Casimir coupling factor:
 4/3 for quarks
 3 for gluons

Medium transport coefficient
 \propto gluon density and momenta

Probe the medium

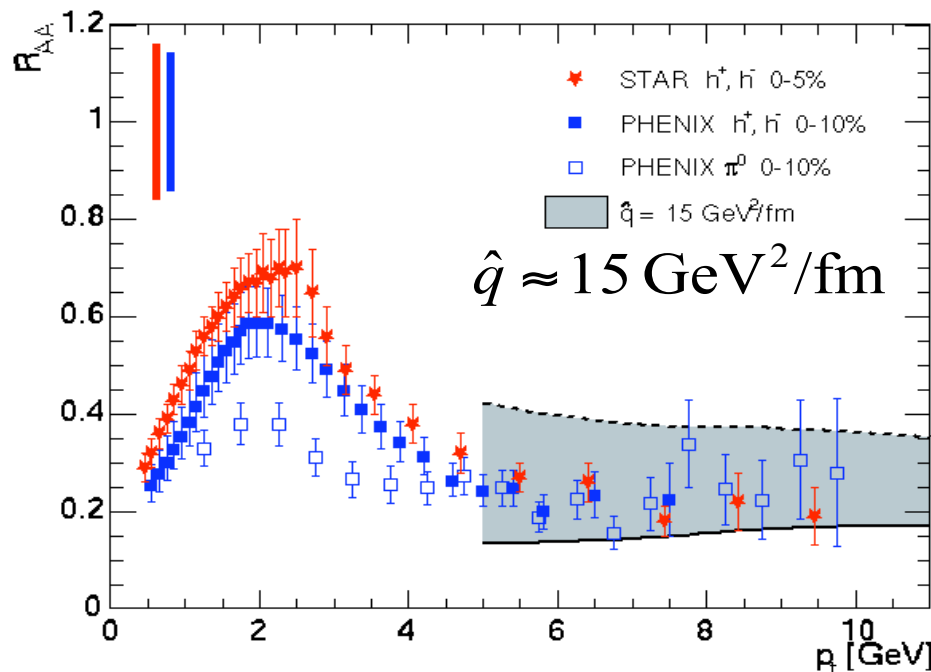
- To analyse RHIC data, need more than mean ΔE
- Quenching weights: energy loss probability distributions

$$P(\Delta E; C_R, \hat{q}, L)$$

- Main limitation: calculated in $E \rightarrow \infty$ approx. \rightarrow uncertainties

Model vs RHIC data

- Density (\hat{q}) “tuned” to match R_{AA} in central Au-Au at 200 GeV
- Extrapolations
 - in centrality: according to Glauber-model collision geometry
 - in \sqrt{s} : assuming $\hat{q} \propto N^{\text{gluons}}/\text{volume} \propto (\sqrt{s})^{0.6}$ (saturation model)



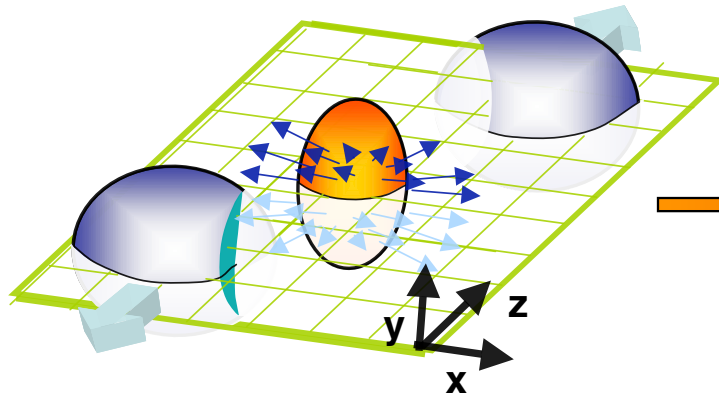
Caveat:

- No initial-state effects and in-medium hadronization: results given for $p_t > 5$ GeV
- Band represents theoretical uncertainty

➔ naturally matches p_t -independence of suppression at high p_t

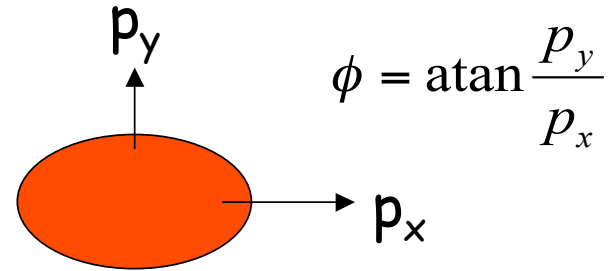
Saturation model: K.J.Eskola, et al, Nucl. Phys. **B570** (2000) 379 [hep-ph/9909456].

Collective Flow of QCD Matter



$$\varepsilon = \frac{\langle y^2 \rangle - \langle x^2 \rangle}{\langle y^2 \rangle + \langle x^2 \rangle}$$

Initial spatial anisotropy



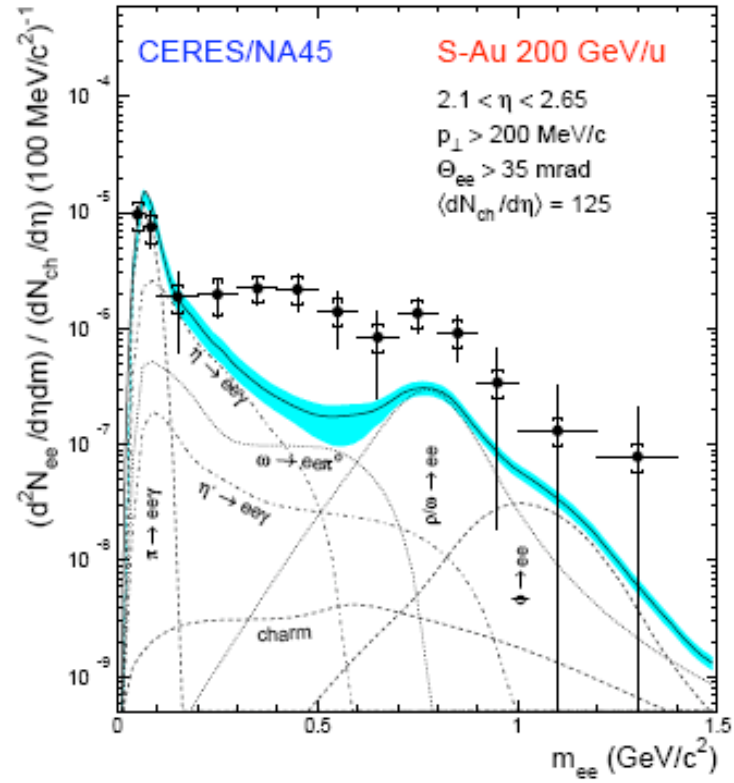
$$v_2 = \frac{\langle p_x^2 \rangle - \langle p_y^2 \rangle}{\langle p_x^2 \rangle + \langle p_y^2 \rangle}$$

Final momentum anisotropy

Reaction plane defined by
“soft” (low p_T) particles

Elliptic flow $\frac{dN}{d\phi} \propto 1 + 2v_2 \cos[2(\phi - \Psi_R)]$

Dilepton spectra



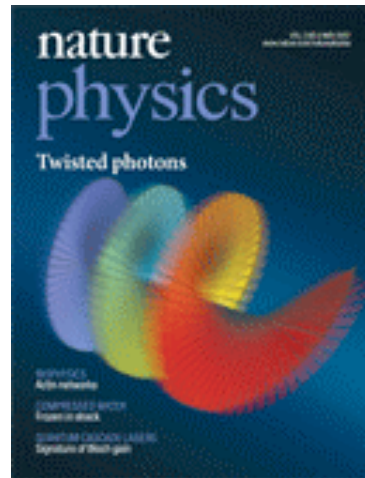
News and Views

***Nature Physics* 3, 297 - 298**

(2007)

Heavy-ion collisions: Where string theory meets reality?

Makoto Natsuume



Heavy ion collisions and black holes in Anti-de Sitter space.

[John R. Ellis](#) ([CERN](#)) . CERN-TH-98-391, Oct 1998. 18pp.

Invited talk at RHIC Physics and Beyond: Kay Kay Gee Day, Upton, NY, 23 Oct 1998.

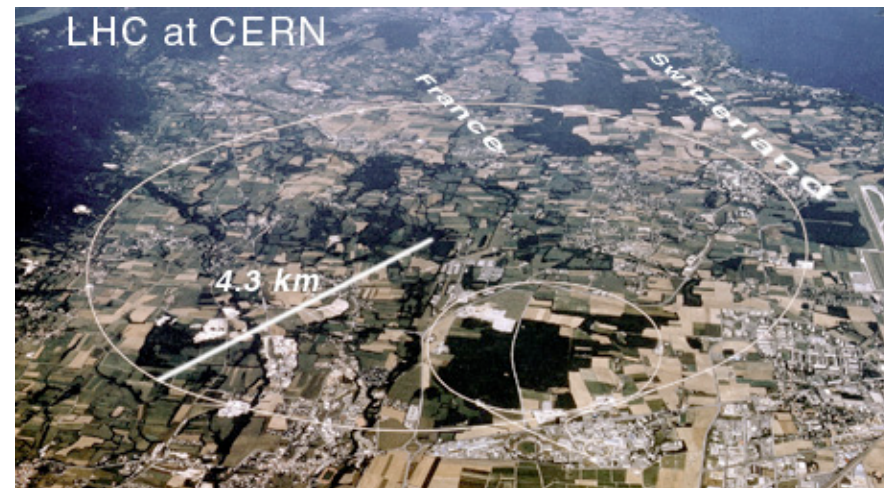
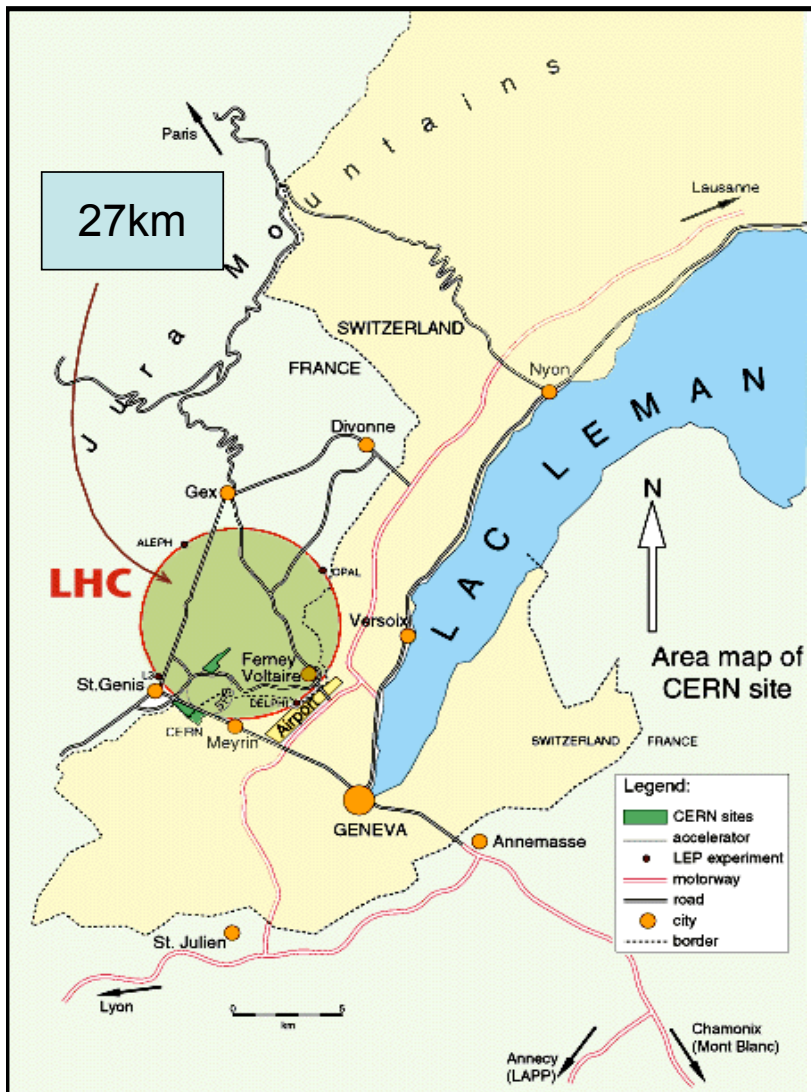
In *Brookhaven 1998, RHIC physics and beyond* 35-52.

e-Print: **nucl-th/9902061**

Experimental aspects

Accelerator



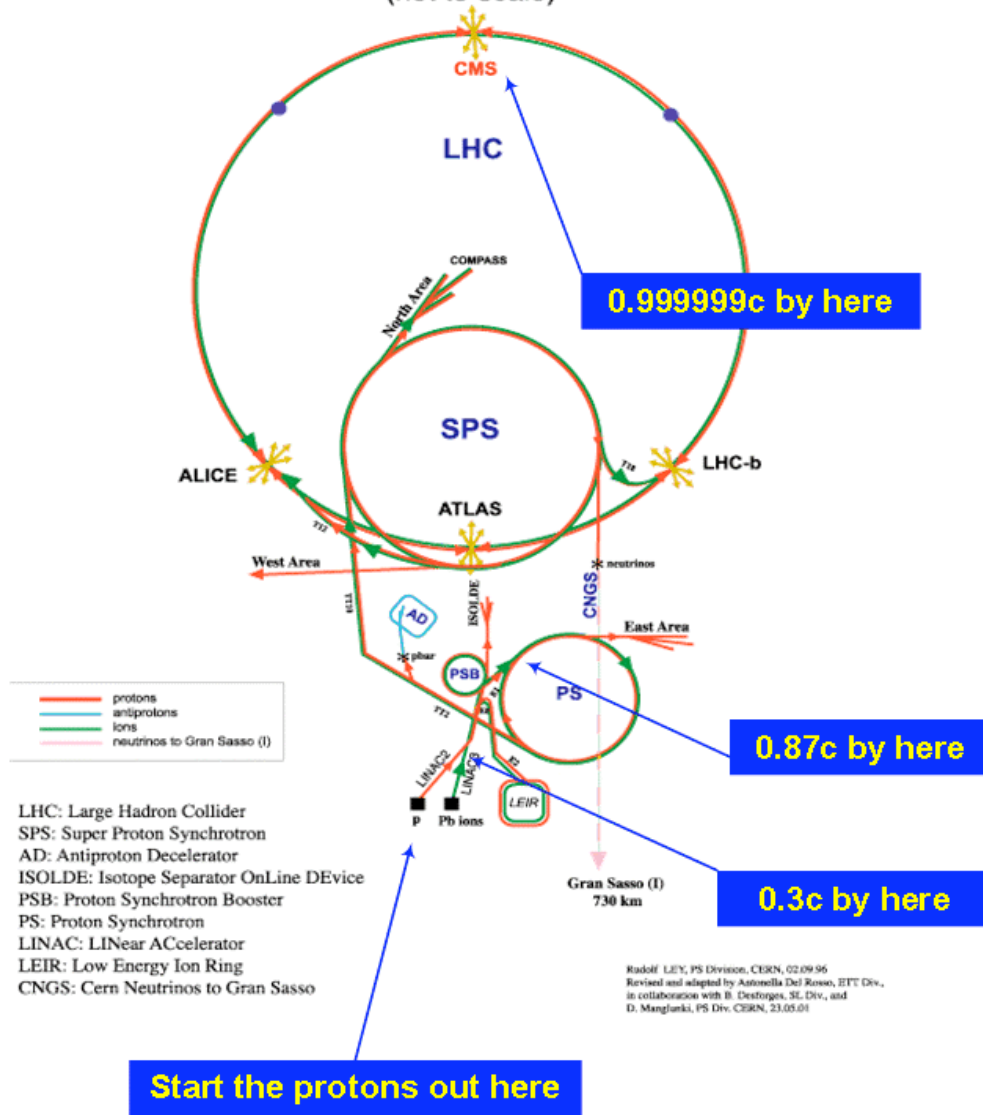


Note : \sqrt{s} limited by needed bending power.

LHC : 1232 superconducting dipoles with $B = 8.4 \text{ T}$ working at 1.9 Kelvin (the largest cryogenic system in the world)

Accelerators and LHC experiments at CERN

CERN Accelerators
(not to scale)



Energies:

Linac 50 MeV

PSB 1.4 GeV

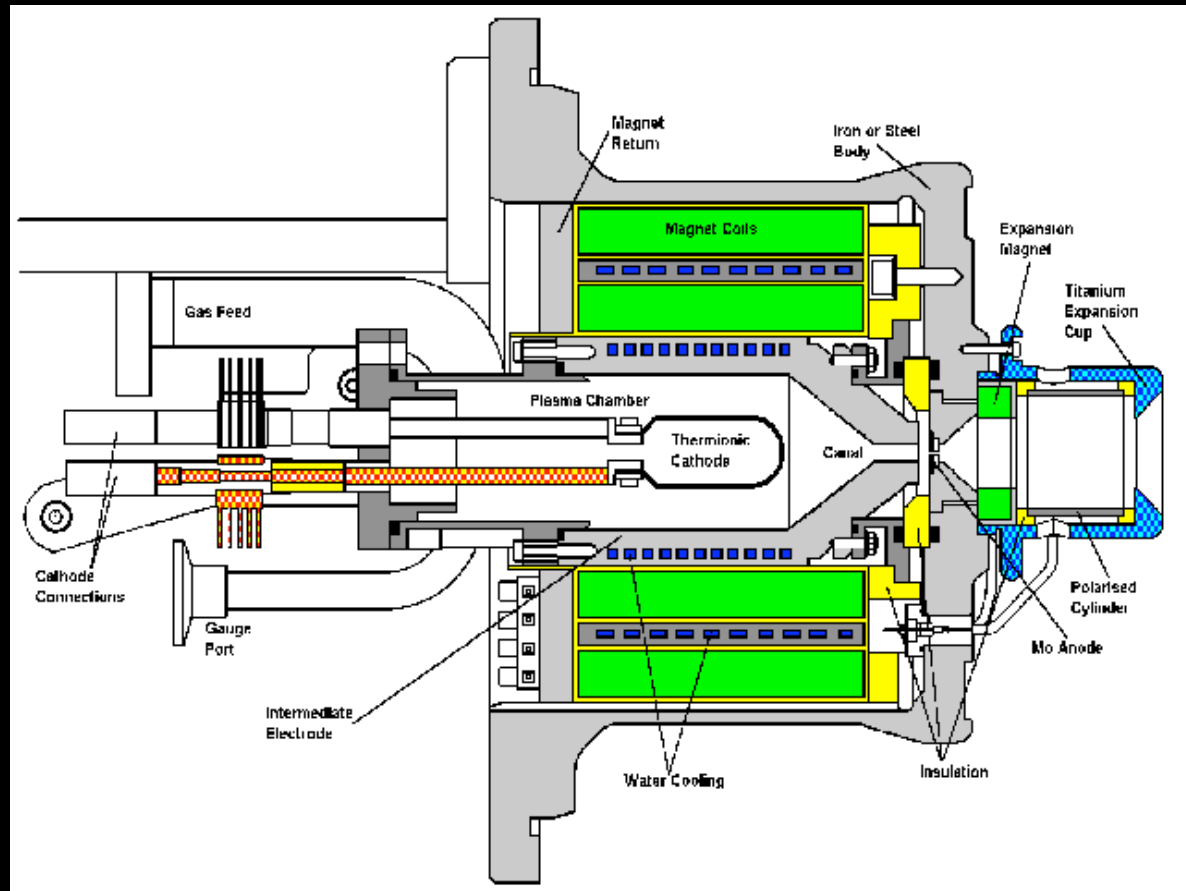
PS 28 GeV

SPS 450 GeV

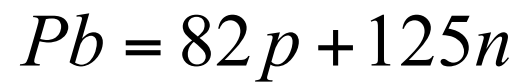
LHC 7 TeV

Radolf LEY, PS Division, CERN, 02.09.96
 Revised and adapted by Antonella Del Rio, IFTT Div.,
 in collaboration with B. Desforges, SE Div., and
 D. Manglani, PS Div. CERN, 21.05.01

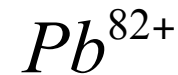
Proton source duoplasmatron



2.76 TeV/nucleon



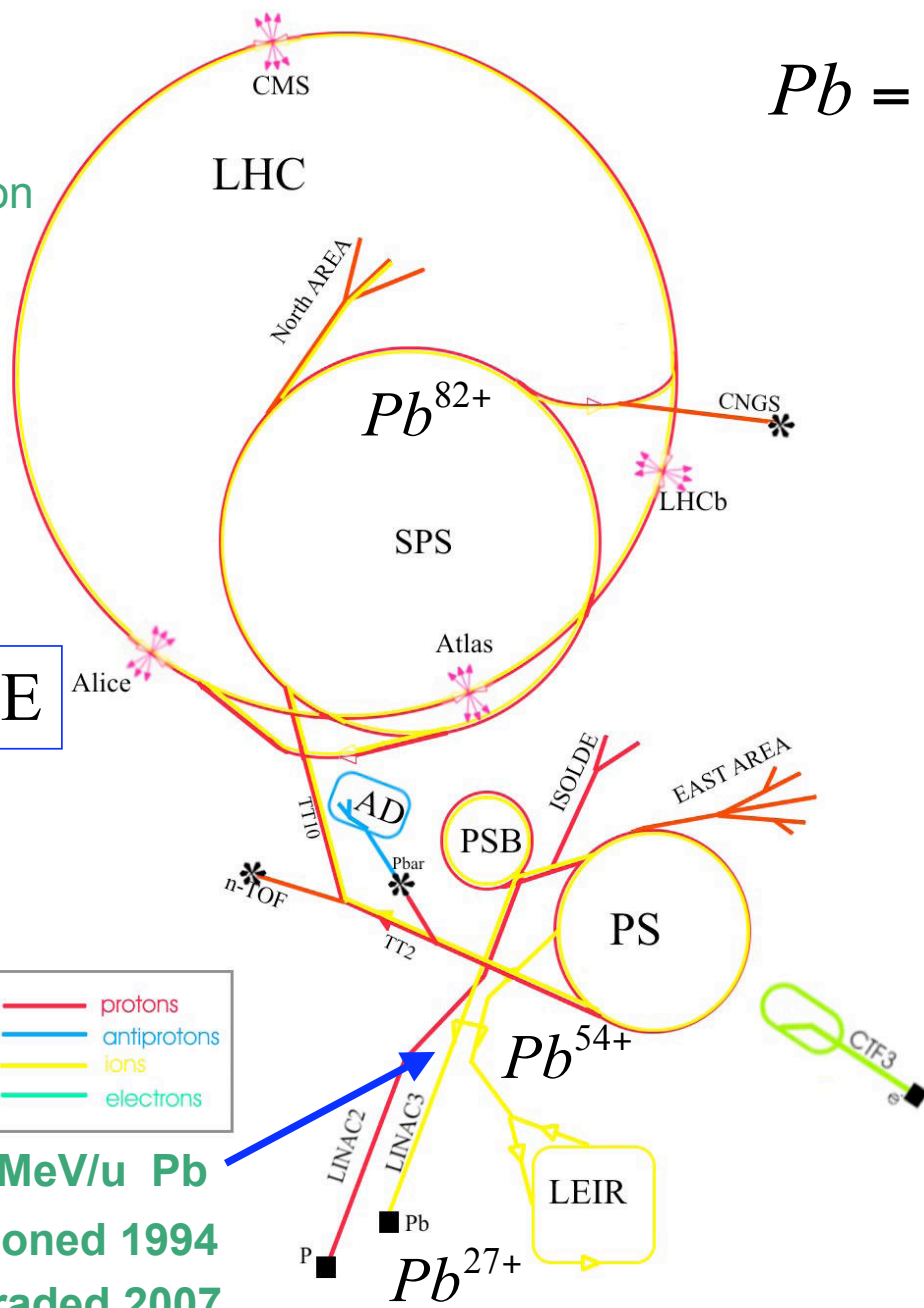
fully stripped
Pb ions



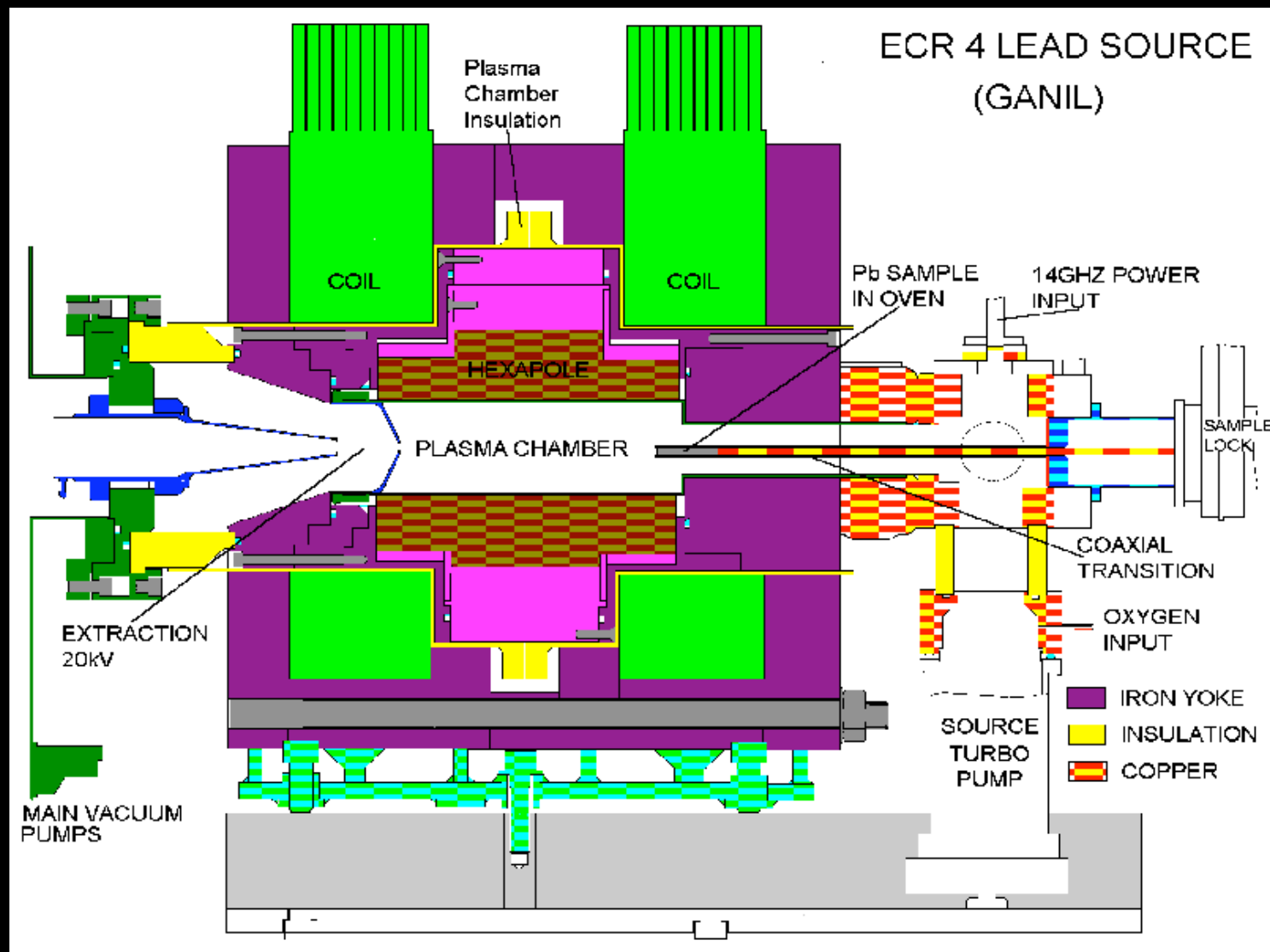
ALICE

- protons
- antiprotons
- ions
- electrons

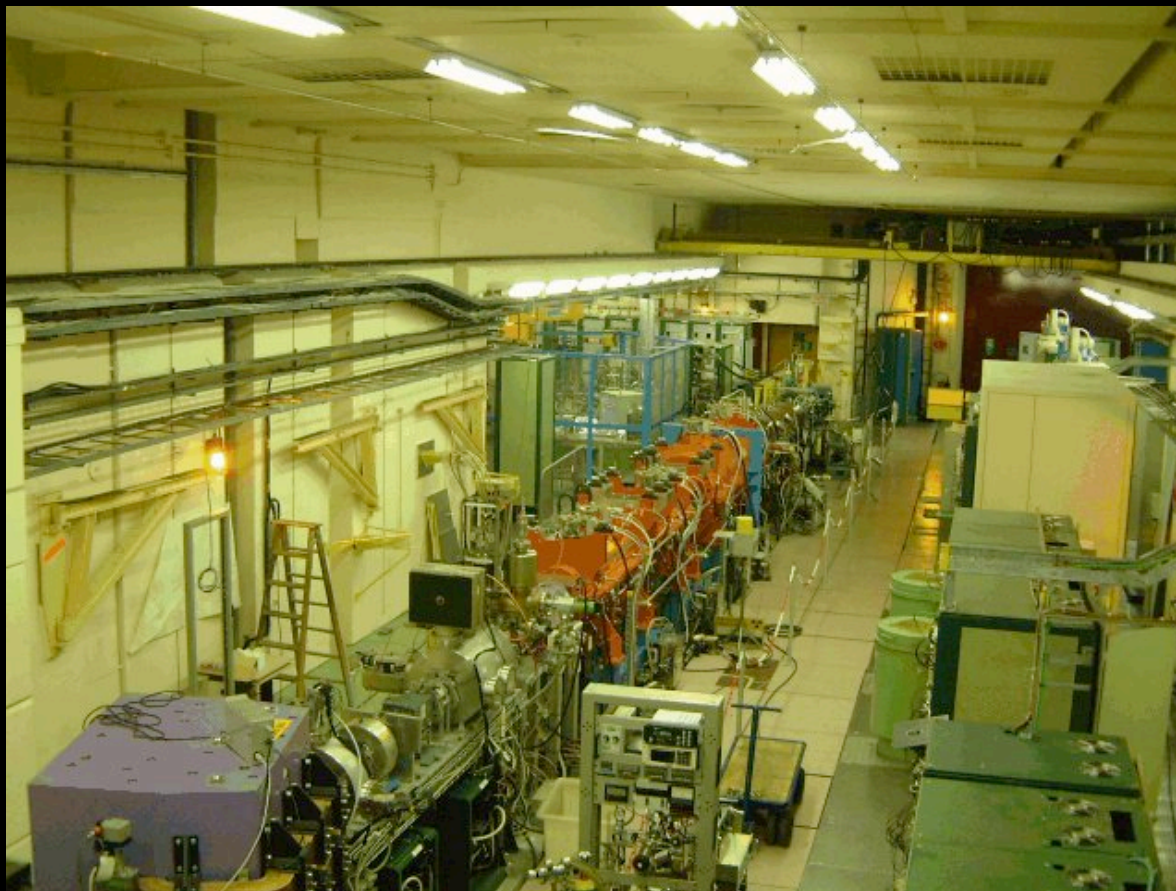
4.2 MeV/u Pb
commissioned 1994
upgraded 2007



Electron Cyclotron Resonance ECR



Heavy ion linac



LHC approved by the CERN Council on december 1994

... as a 2-stage project:

First Stage	10 TeV	2004
Second Stage	14 TeV	2008



Conceptual Design Report "Yellow Book" october 1995

first machine built at CERN with substantial contribution from non-Member states



Canada, India, Japan, Russia, USA

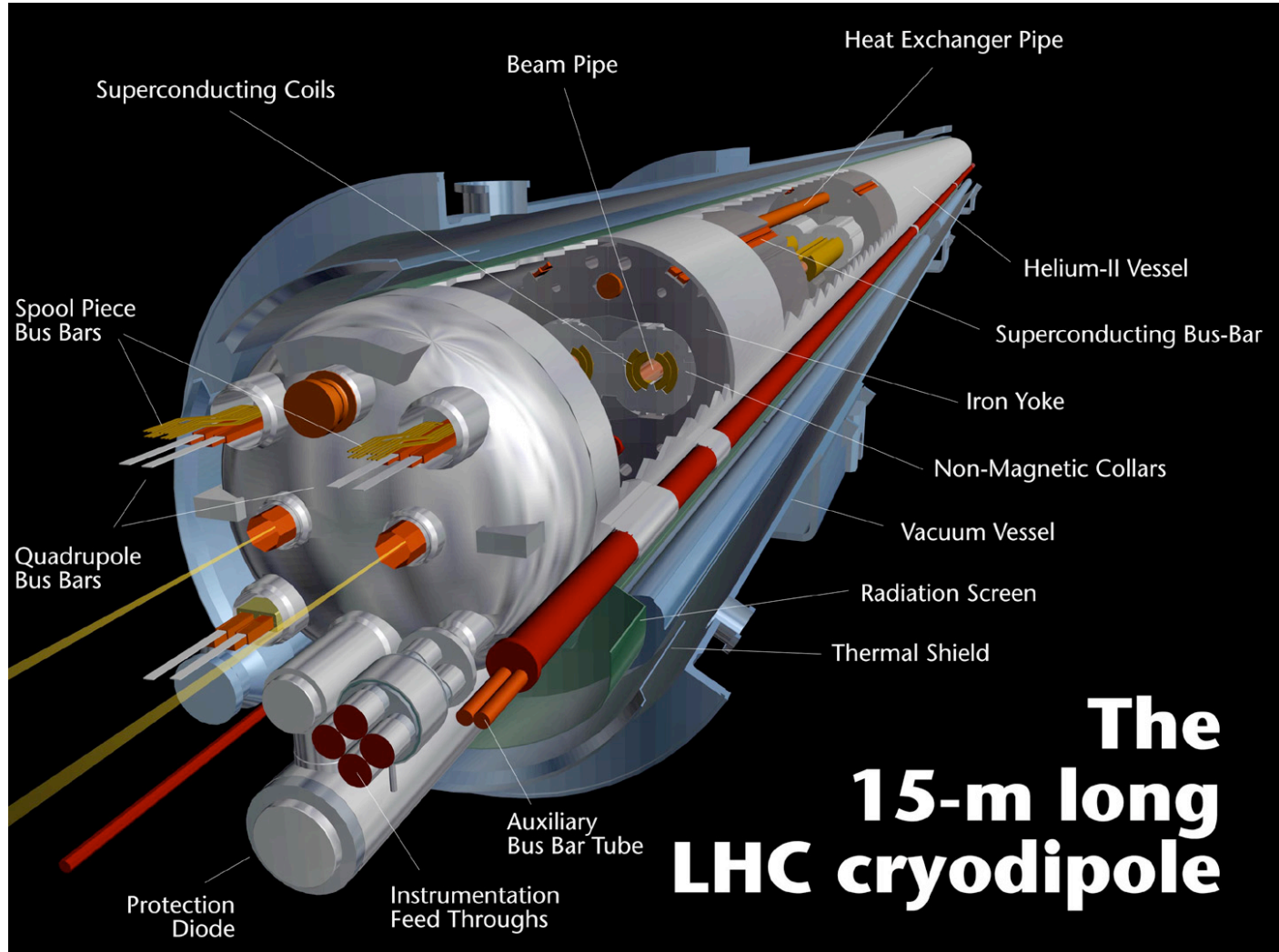
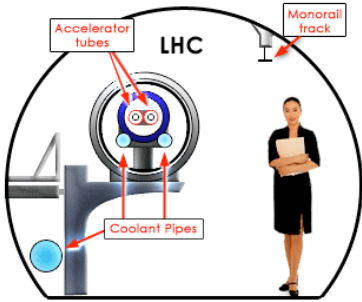
0.5 million US dollars each

1232 needed



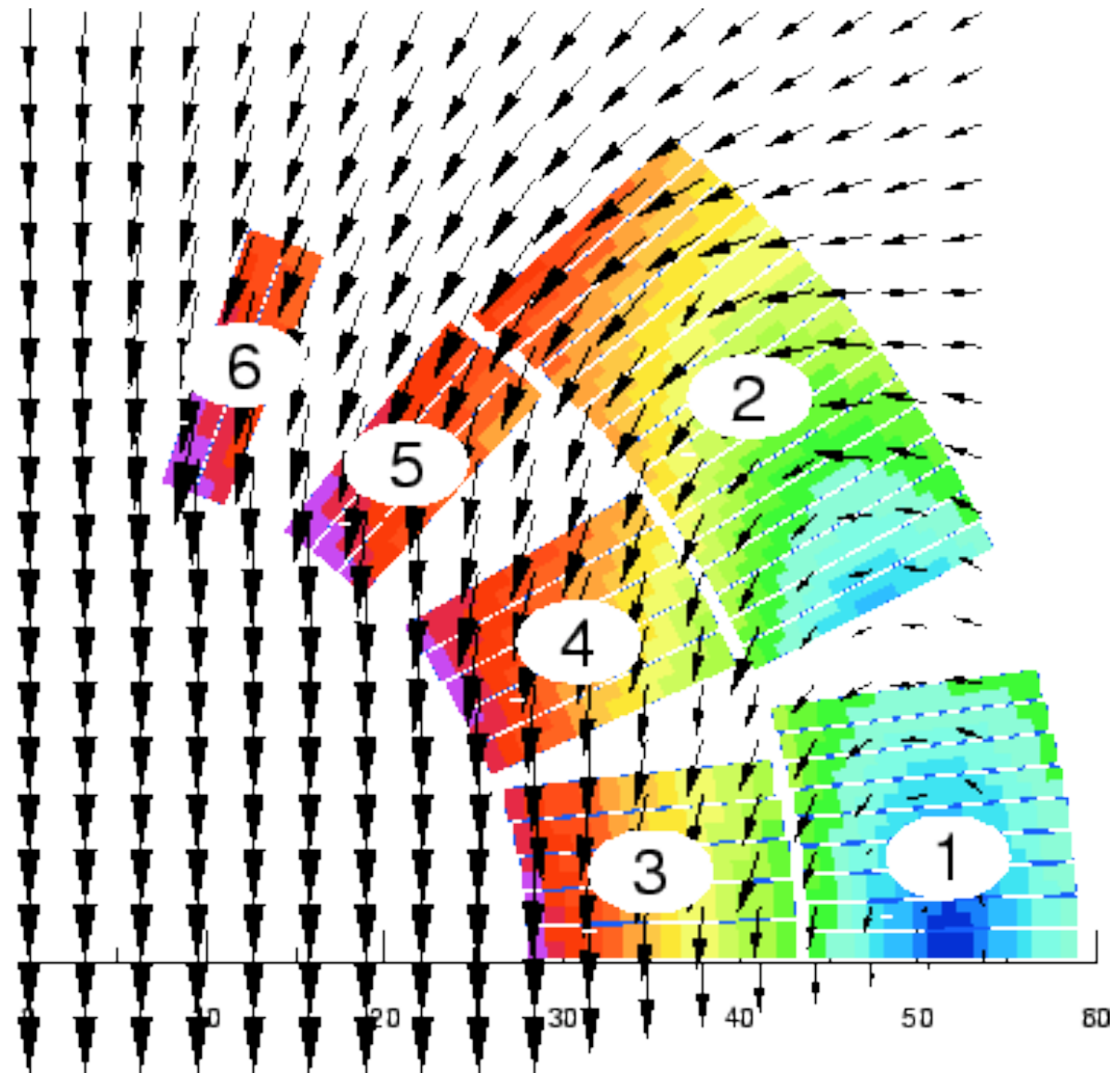
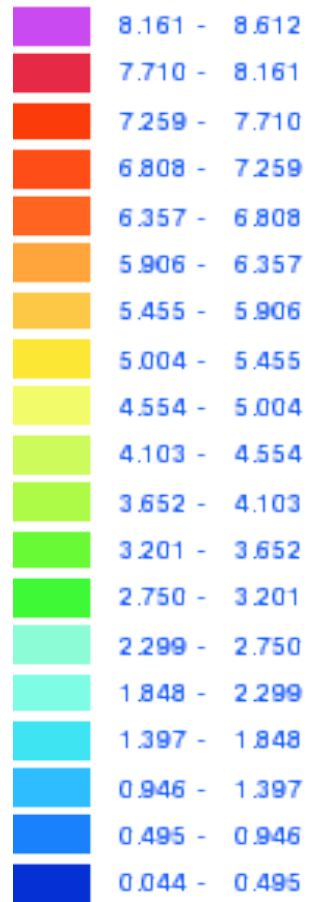
Dipole Magnet





Magnetic Field

$|B|$ (T)

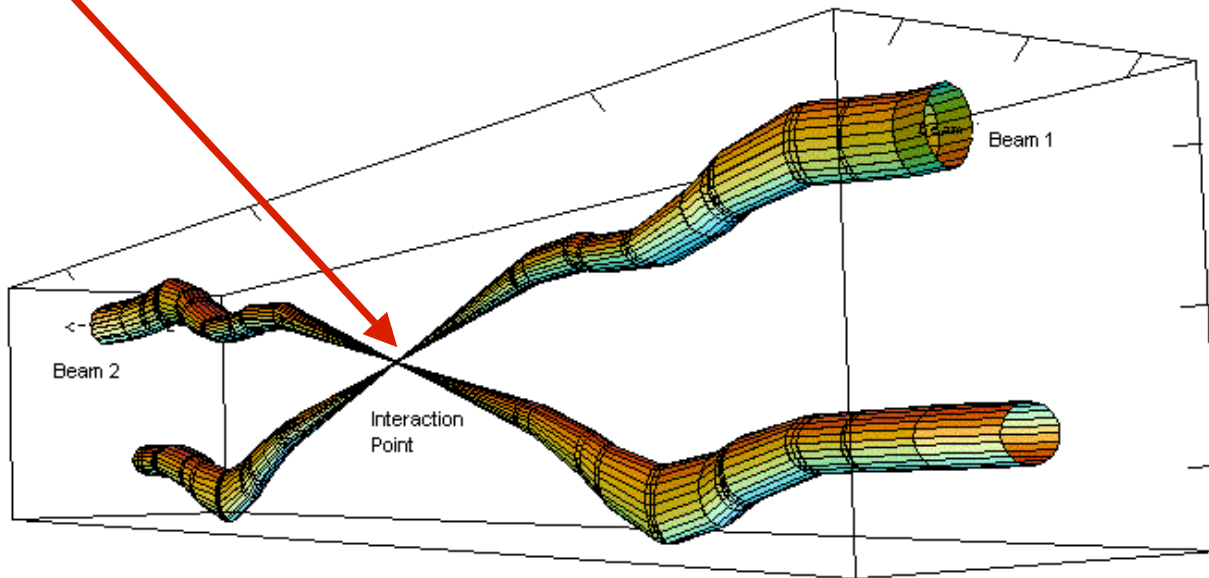


total proton - proton cross section at 7 TeV \approx 110 mbarn

$$\sigma_{inel} = 60mb$$

$64\mu m$ (width of a human hair)

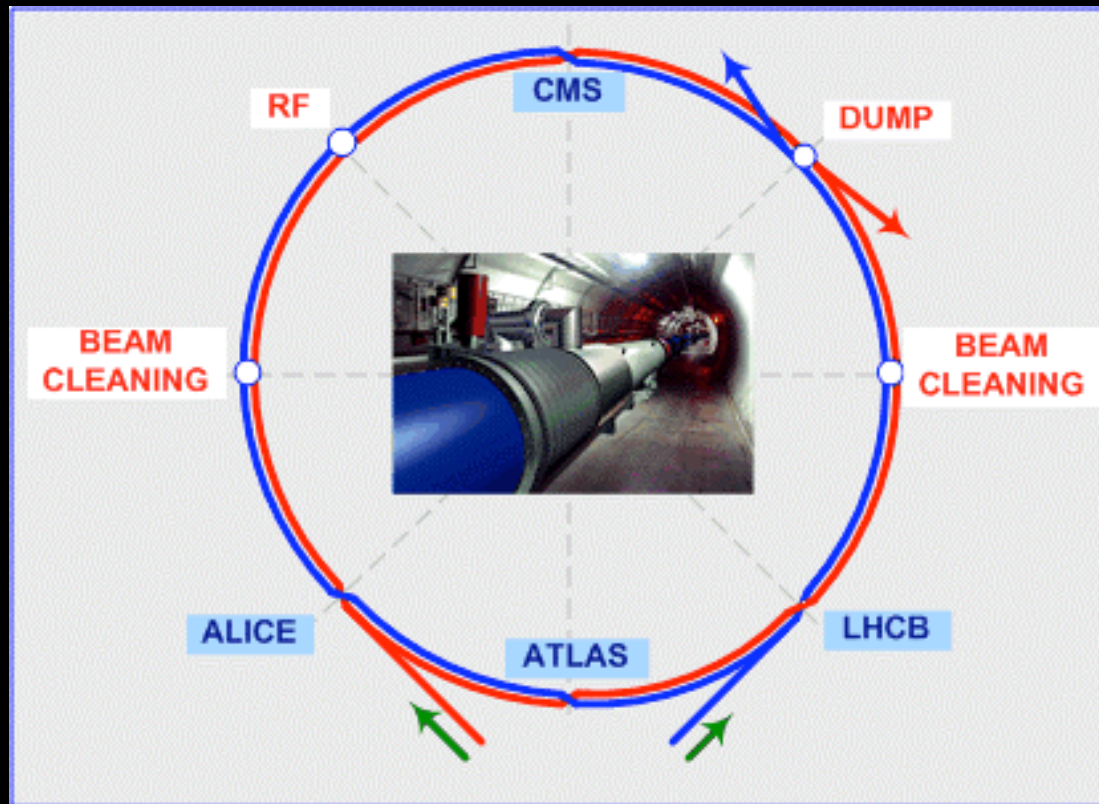
1 barn = 10^{-24} cm



Relative beam sizes around IP1 (Atlas) in collision

event rate = luminosity x cross section

$$\text{event rate} = 10^{34} \cdot 60 \cdot 10^{-3} \cdot 10^{-24} = 600 \text{ millions/sec}$$



injected at 450 GeV, accelerated to 7 TeV

For our fully-stripped ${}_{208}\text{Pb}^{82+}$ nuclei ("ions")
in the same magnetic field as 7 TeV protons:

$$\begin{aligned} E_{\text{pb}} &= Z E_p = 82 \times 7 \text{ TeV} = 7Z \text{ TeV} = 574 \text{ TeV} \\ &= A E_n = 208 \times 2.76 \text{ TeV} = 2.76A \text{ TeV} = 574 \text{ TeV} \\ &= 82 \times (\text{kinetic energy of mosquito at 1m/s}) \\ &= (\text{kinetic energy of 1 mm diameter grain of sand at 40 km/h}) \end{aligned}$$

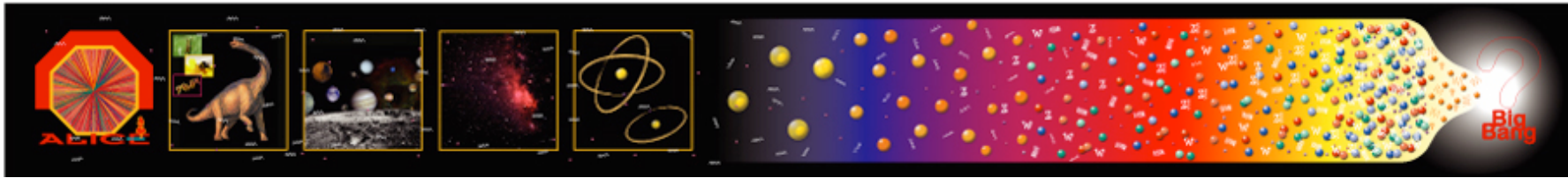
Key Parameters of “Early” Pb Ion Beam (from LHC Design Report)

Parameter	Units	Early Beam	Nominal
Energy per nucleon	TeV	2.76	2.76
Initial ion-ion Luminosity L_0	$\text{cm}^{-2} \text{s}^{-1}$	$\sim 5 \times 10^{25}$	1×10^{27}
No. bunches, k_b		62	592
Minimum bunch spacing	ns	1350	99.8
β^*	m	1.0	0.5 / 0.55
Number of Pb ions/bunch		7×10^7	7×10^7
Transv. norm. RMS emittance	μm	1.5	1.5
Longitudinal emittance	eV s/charge	2.5	2.5
Luminosity half-life (1,2,3 expts.)	h	14, 7.5, 5.5	8, 4.5, 3
<p>At full energy, luminosity lifetime is determined mainly by collisions (“burn-off” from ultraperipheral electromagnetic interactions) $\sigma \approx 520 \text{ barn}$</p>		Only possibility for 2009 or early 2010	Goal for 2-3 years (?) beyond

Experimental aspects

Detector

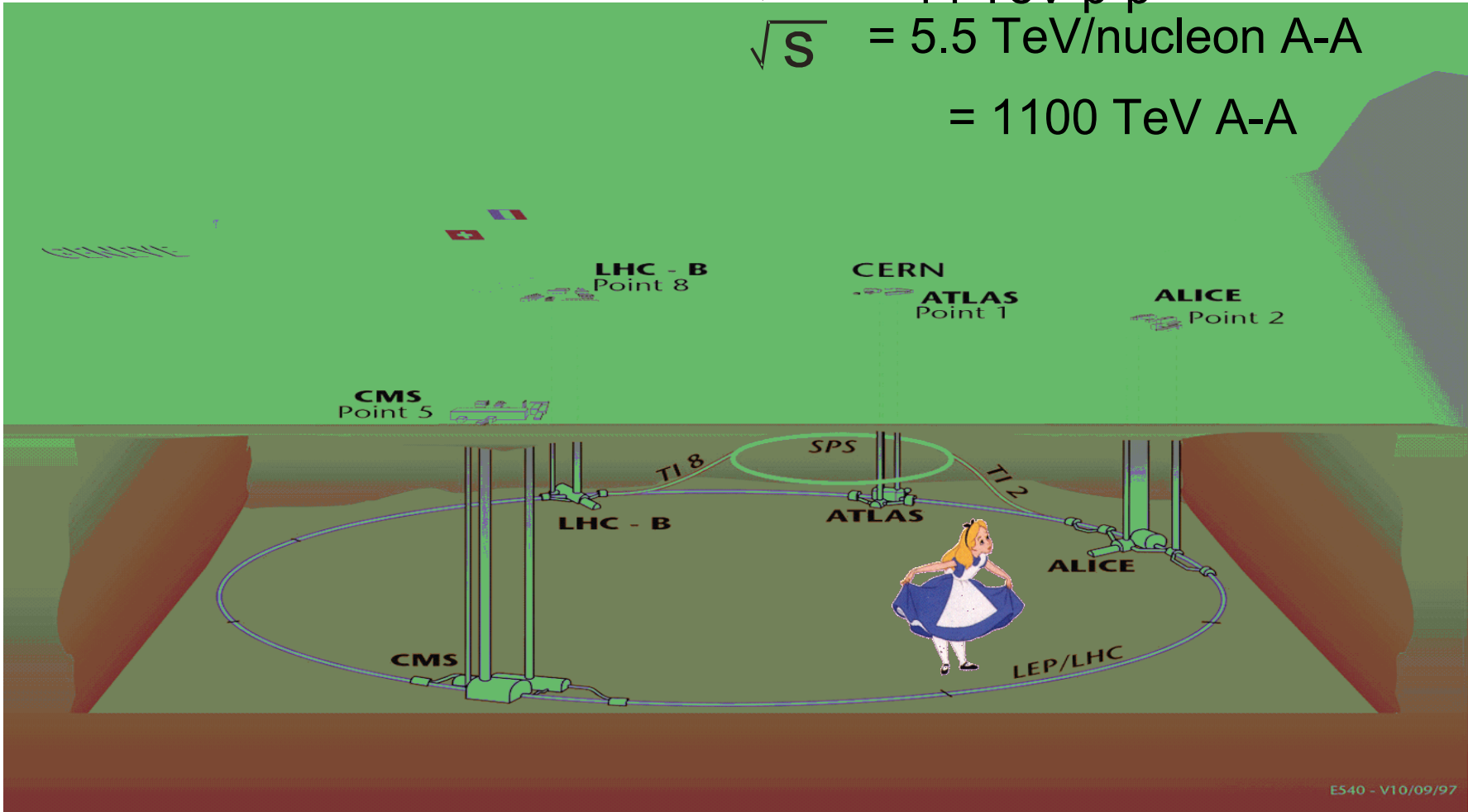




$$\sqrt{S} = 14 \text{ TeV p-p}$$

$$\sqrt{S} = 5.5 \text{ TeV/nucleon A-A}$$

$$= 1100 \text{ TeV A-A}$$



E540 - V10/09/97

$$5.5 \text{ TeV} \times 100 \text{ nucleones} \times 2 \text{ (cms)} = 1100 \text{ TeV}$$

ALICE Physics goals all aspects in the same experiment!

• Global observables:

Multiplicities, η distributions

• Degrees of freedom as a function of Temperature

hadron ratios and spectra, dilepton continuum, direct photons

• Early state manifestation of collective effects: elliptic flow

• Energy loss of partons in quark gluon plasma:
jet quenching, high pt spectra, open charm and open beauty

• Deconfinement:

charmonium and bottomonium spectroscopy

• Chiral symmetry restoration:
neutral to charged ratios, res. decays

• Fluctuation phenomena - critical behavior:

event-by-event particle comp. and spectra

• Geometry of the emitting source:
HBT, impact parameter via zero-degree energy flow

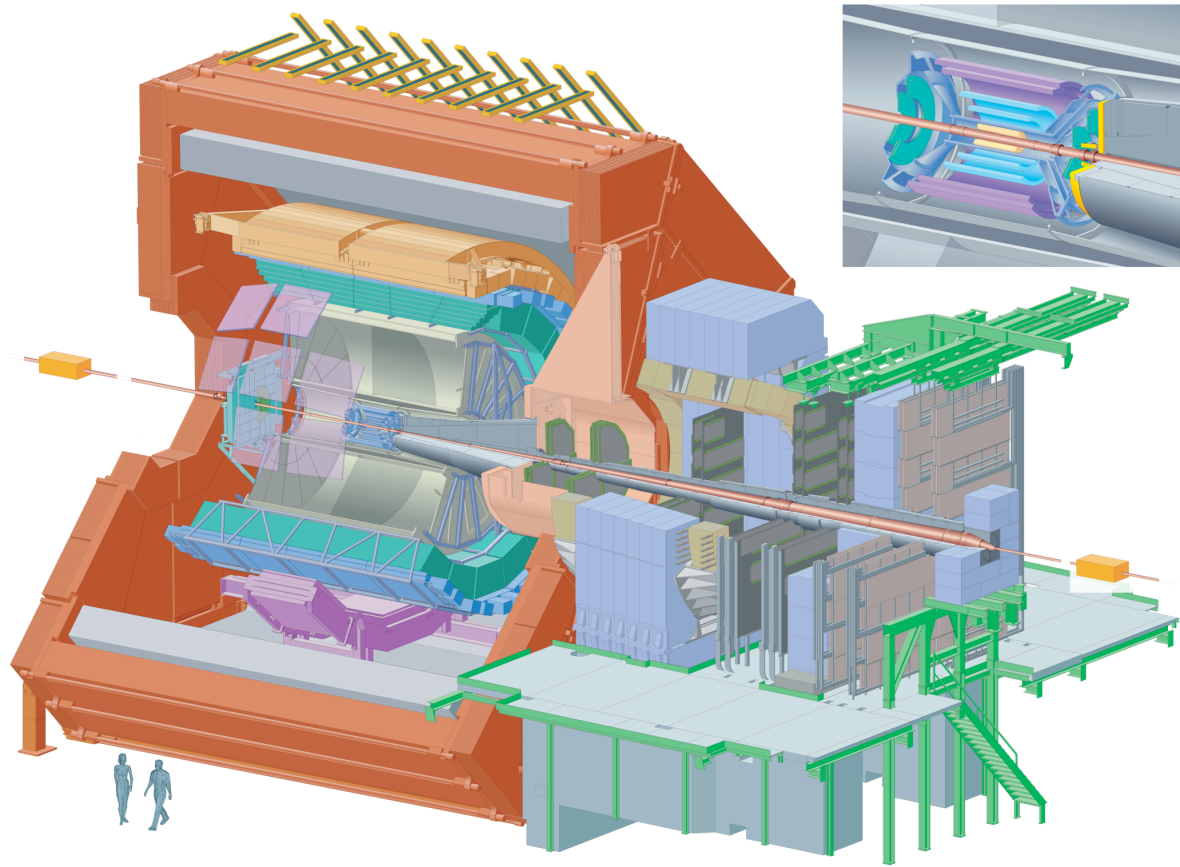
• pp collisions in a new energy domain

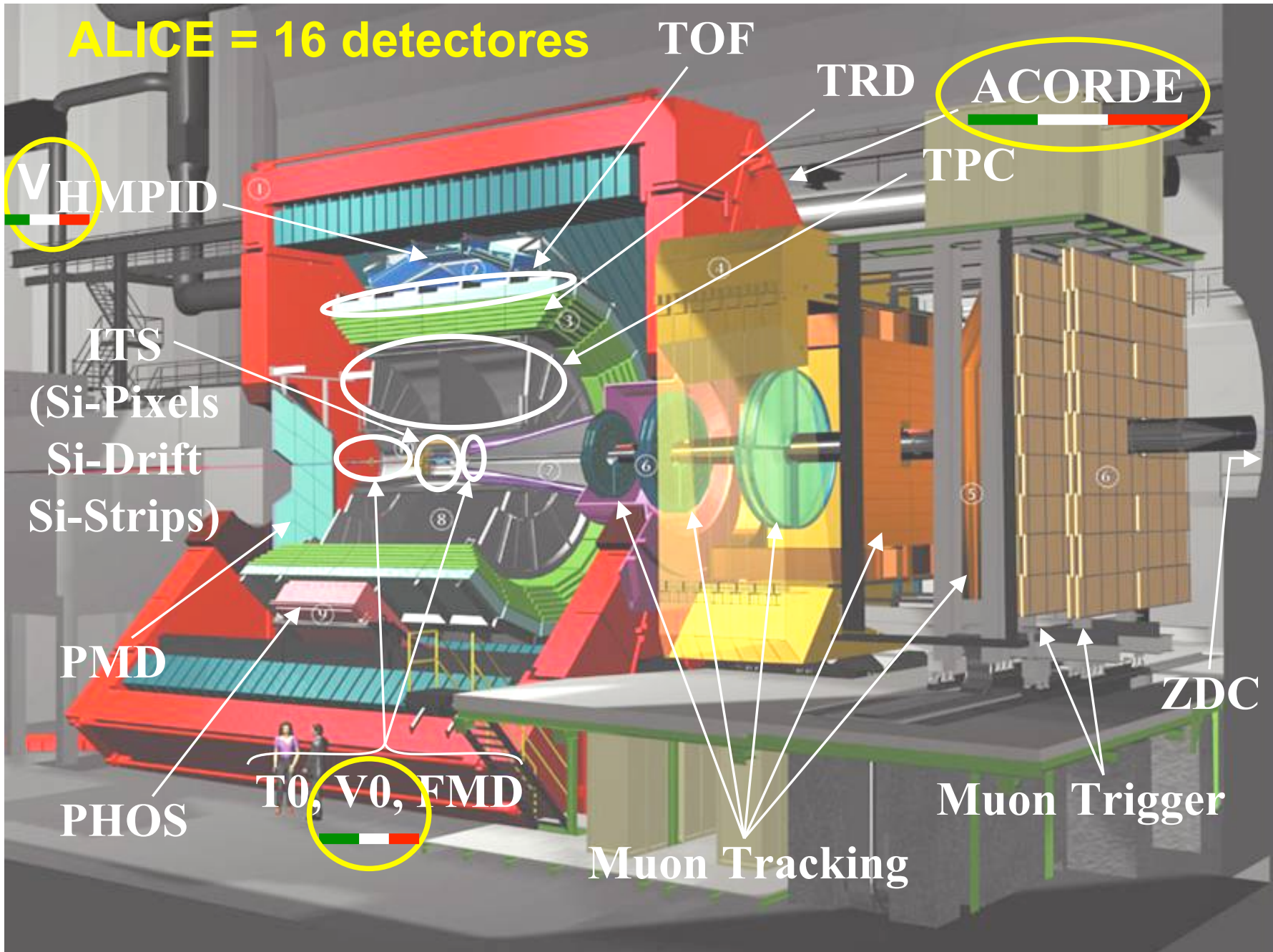
- Large acceptance
- Good tracking capabilities
- Selective triggering
- Excellent granularity

- Wide momentum coverage
- P.I.D. of hadrons and leptons
- Good sec. vertex reconstr.
- Photon Detection



use a variety of experimental techniques!

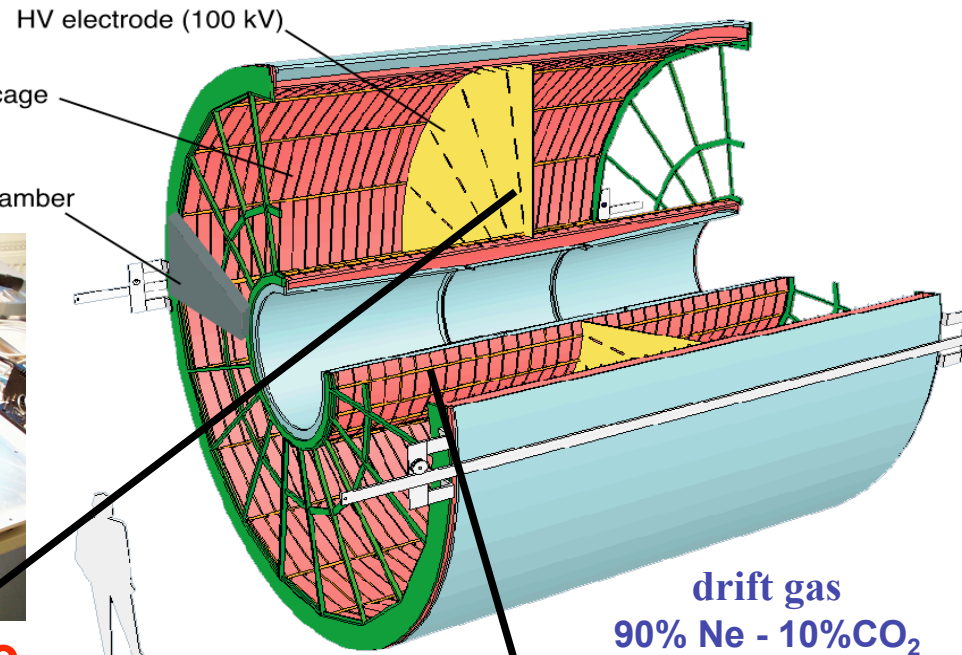
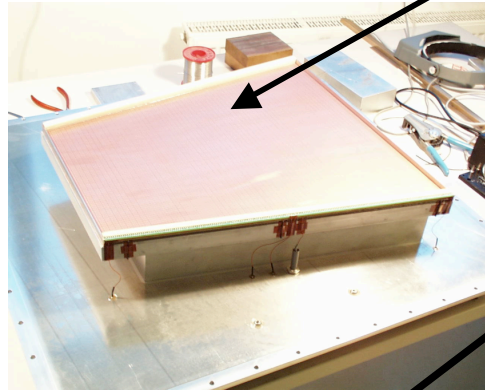




Time Projection Chamber

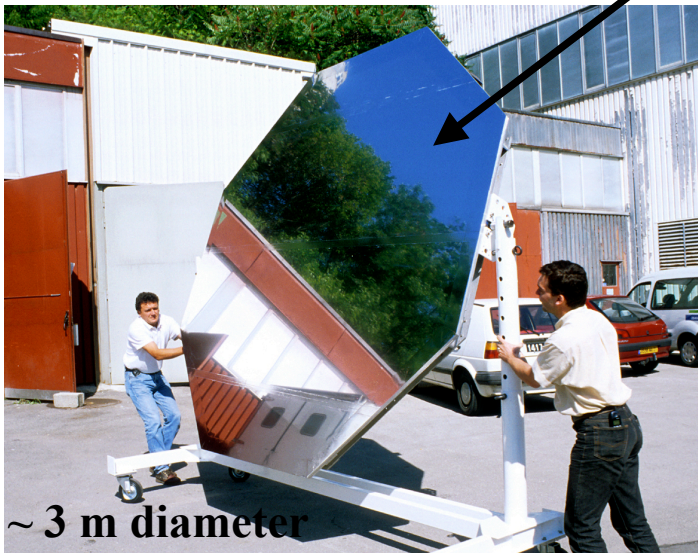
largest ever: 88 m³,
570 k channels

for tracking
and PID via
dE/dx
- 0.9 < η < 0.9

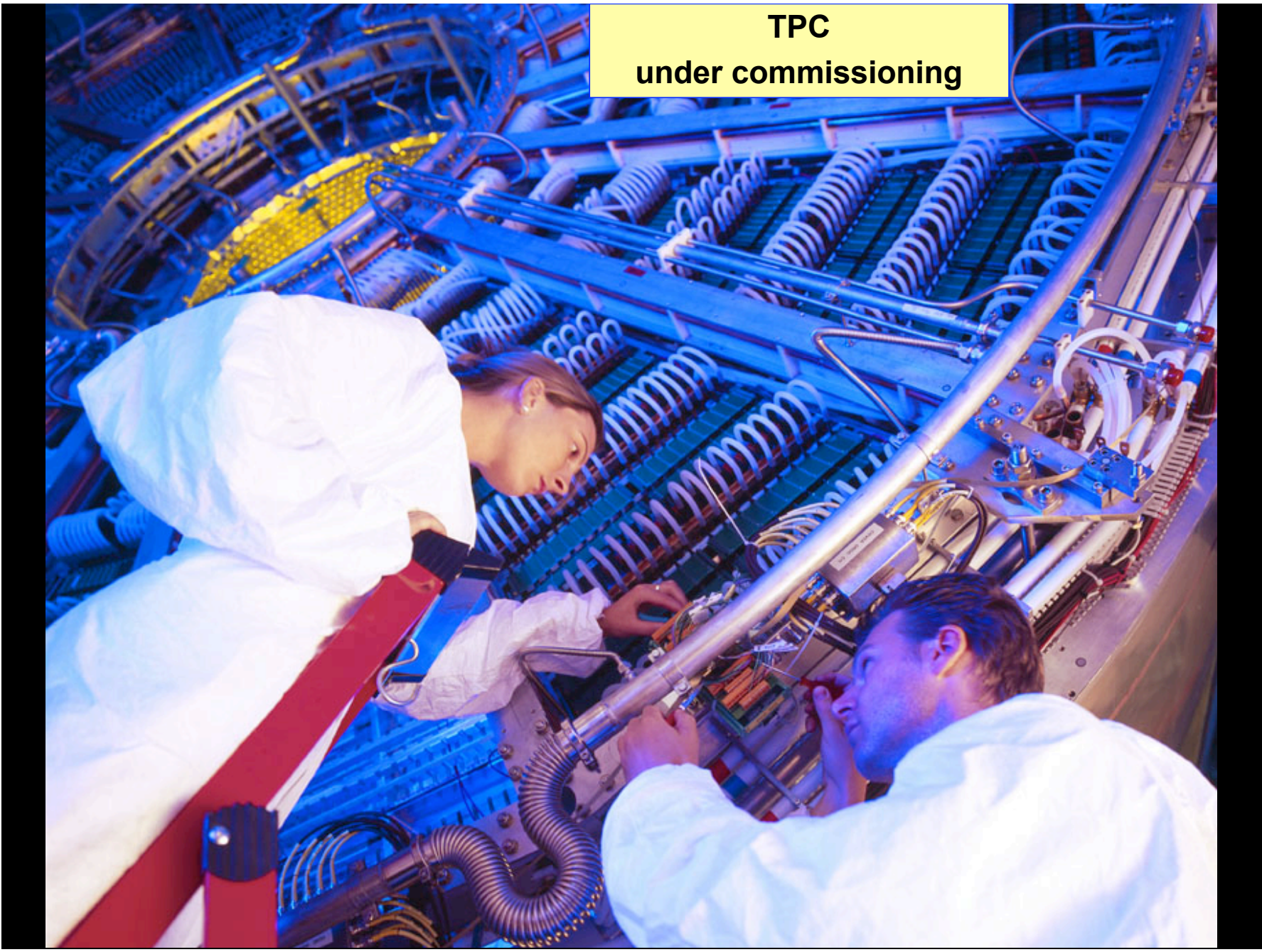


Central Electrode Prototype

25 μ m aluminized Mylar on Al frame

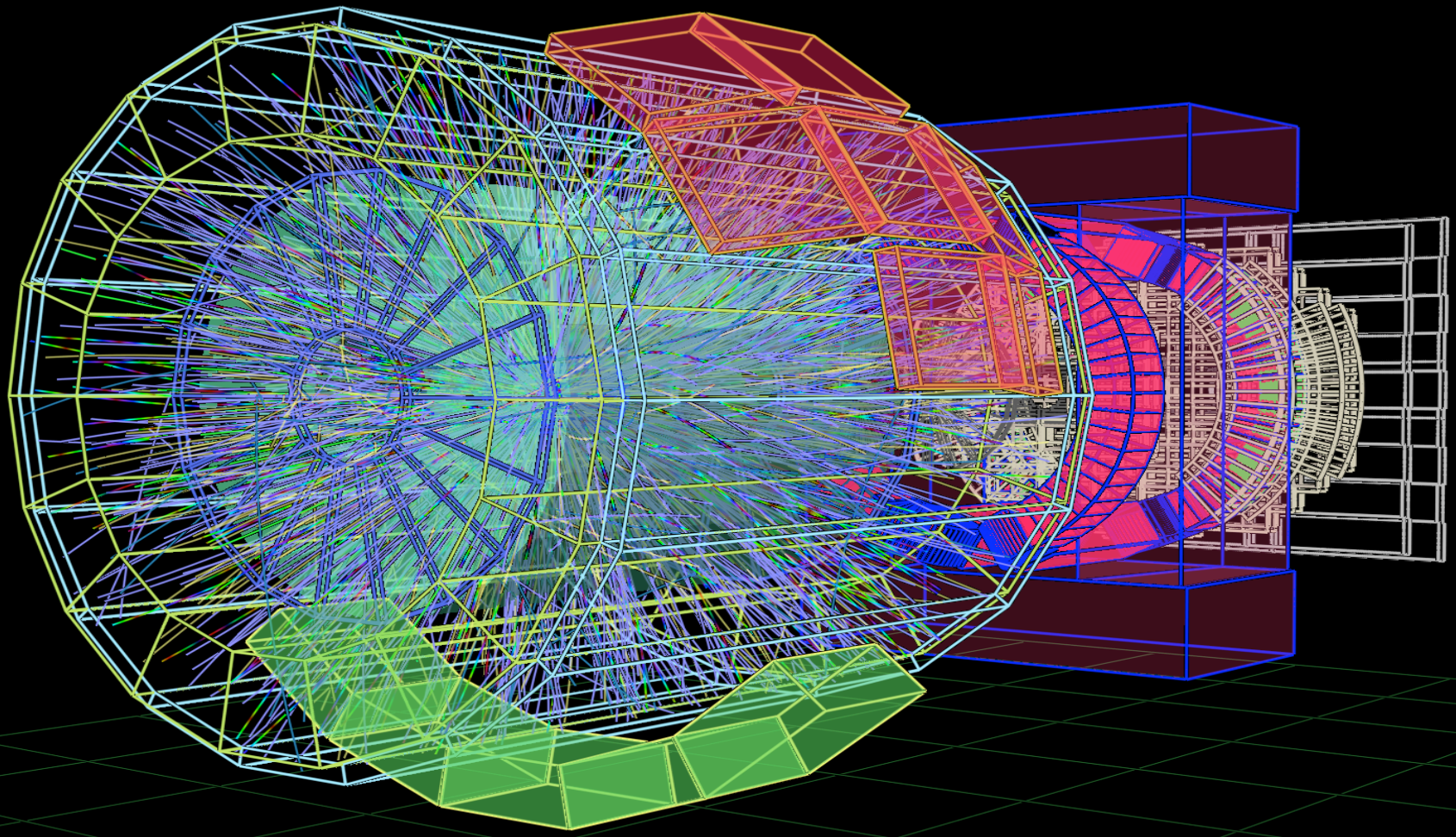


TPC
under commissioning



Summary

- LHC is the ultimate machine for Heavy Ion Collisions
 - very **significant step** beyond RHIC
 - excellent conditions for **experiment & theory** (QCD)
- ALICE is a powerful next generation detector
 - first truly **general purpose** HI experiment
 - addresses most relevant observables: from **super-soft to ultra-hard**
 - many **advances** in technology
- Both the experiment and the machine are getting ready
... Physics at last is coming



thanks