

Statistics for High Energy Physics

1. Profile Likelihood, Asimov, CLs

2. Exclusion, Discovery

3. Look Elsewhere Effect

Eilam Gross,
Weizmann Institute of Science



Basically ITS ALL ABOUT NUMBERS

$$n_{expected} = L \cdot \sigma \cdot eff$$

$$[L] = events/cm^2$$

$$[\sigma] = cm^2$$



So Far



A counting experiment

- The Higgs hypothesis is that of signal $s(m_H)$

$$s(m_H) = L\sigma_{SM} \cdot \epsilon$$

For simplicity unless otherwise noted $s(m_H) = L\sigma_{SM}$

- In a counting experiment $n = \mu s(m_H) + b$

$$\mu = \frac{L\sigma_{obs}(m_H)}{L\sigma_{SM}(m_H)} = \frac{\sigma_{obs}(m_H)}{\sigma_{SM}(m_H)}$$

- μ is the strength of the signal
(with respect to the expected Standard Model one)
- The hypotheses are therefore denoted by H_μ
- H_1 is the SM with a Higgs, H_0 is the background only model



A Tale of Two Hypotheses

NULL

ALTERNATE

H_0 - SM w/o Higgs

H_1 - SM with Higgs

Rejecting H_0 in favour of $H_1(m_H)$
→ Discovery of a Higgs with a mass m_H

We quantify rejection by p-value (later)



Swapping Hypotheses \rightarrow exclusion

NULL

ALTERNATE

H_0 - SM w/o Higgs

H_1 - SM with Higgs

- Reject H_1 in favor of H_0

Excluding $H_1(m_H)$ \rightarrow Excluding the Higgs
with a mass m_H

We quantify rejection by p-value (later)



Likelihood

- Likelihood is the compatibility of the Hypothesis with a given data set.

But it depends on the data

$$L(H) = P(x | H)$$

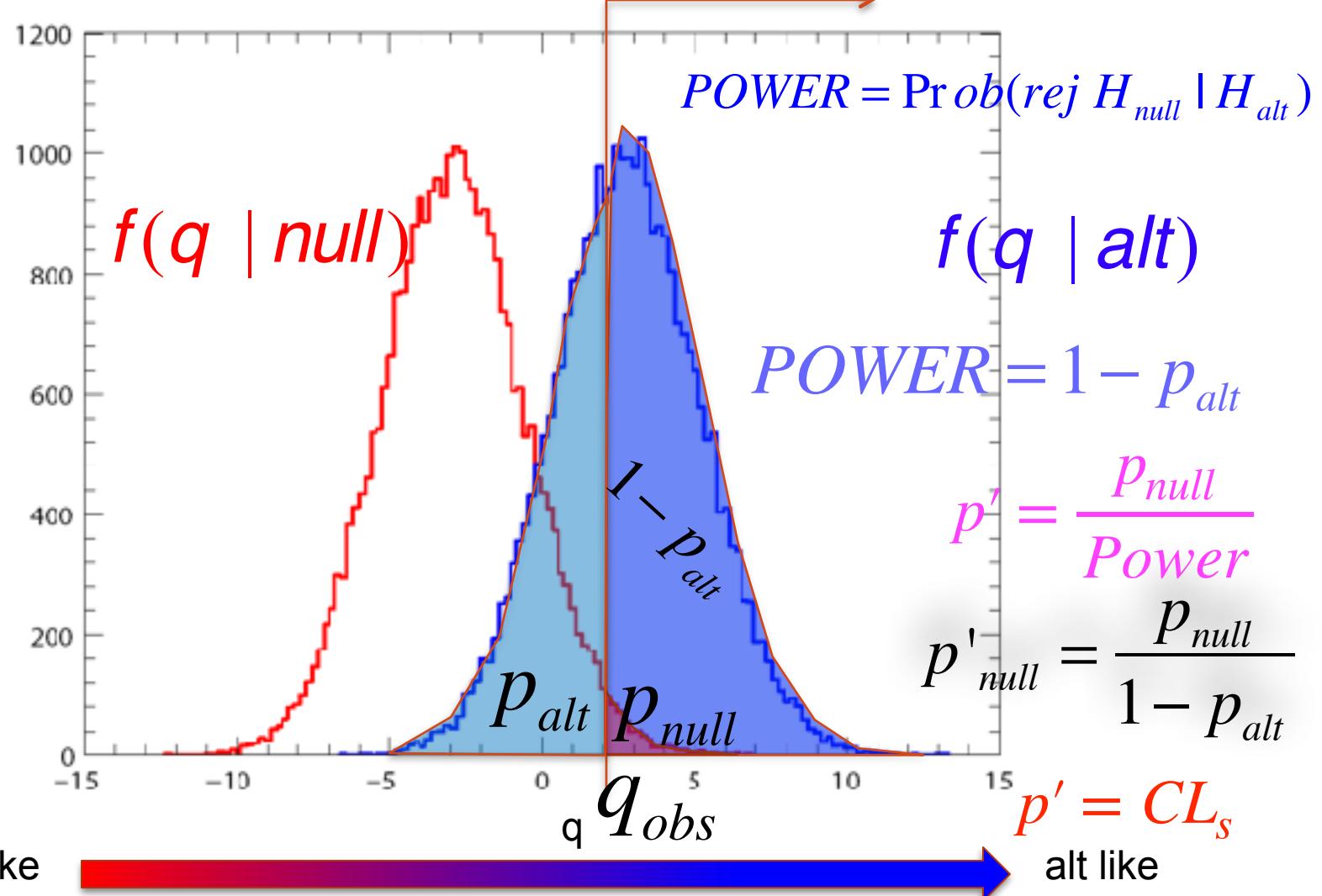


Test statistic and p-value



CLs

If $p \leq \alpha$ reject null



Test Spin 0 parity

$$H_0 = 0^+$$

$$H_1 = 0^-$$

$p_{H_1}(\text{exp} | H_0) = 0.37\%,$

$p_{H_1}(\text{obs}) = 1.5\%$

$p_{H_0}(\text{obs}) = 31\%$

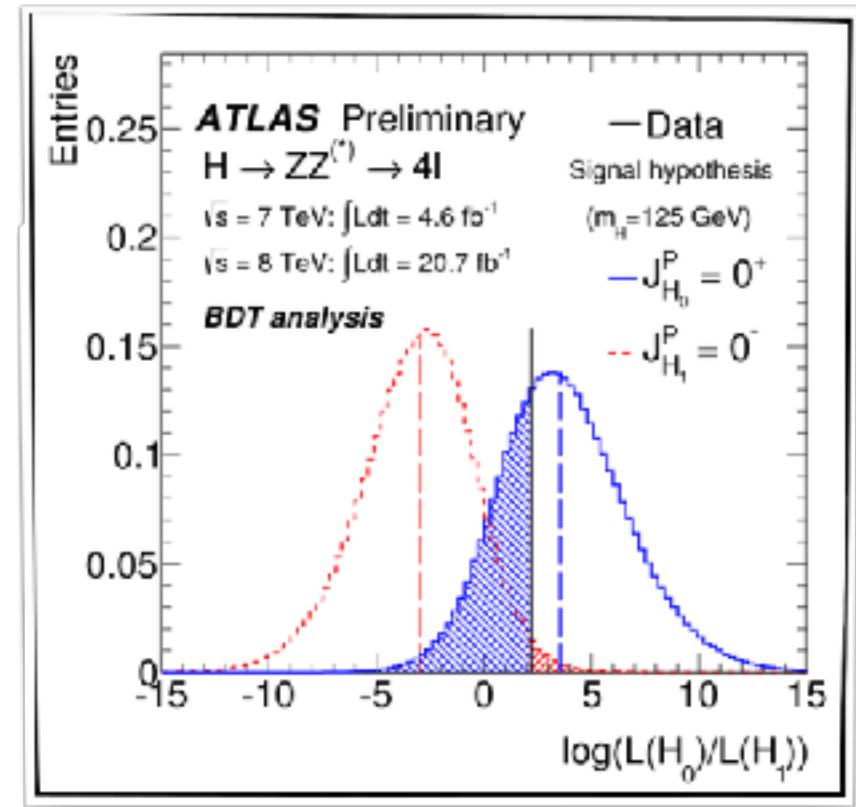
$p_{H_1}^{CL_s}(\text{obs}) = 2.2\%$

$$p_{H_1}^{CL_s} = \frac{p_{H_1}}{1 - p_{H_0}} = \frac{1.5\%}{1 - 0.31} = 2.2\%$$

Which means
 $J^p=0^-$ is excluded at the
 97.8% CL in favour of $J^p=0^+$

H_1 like

H_0 like



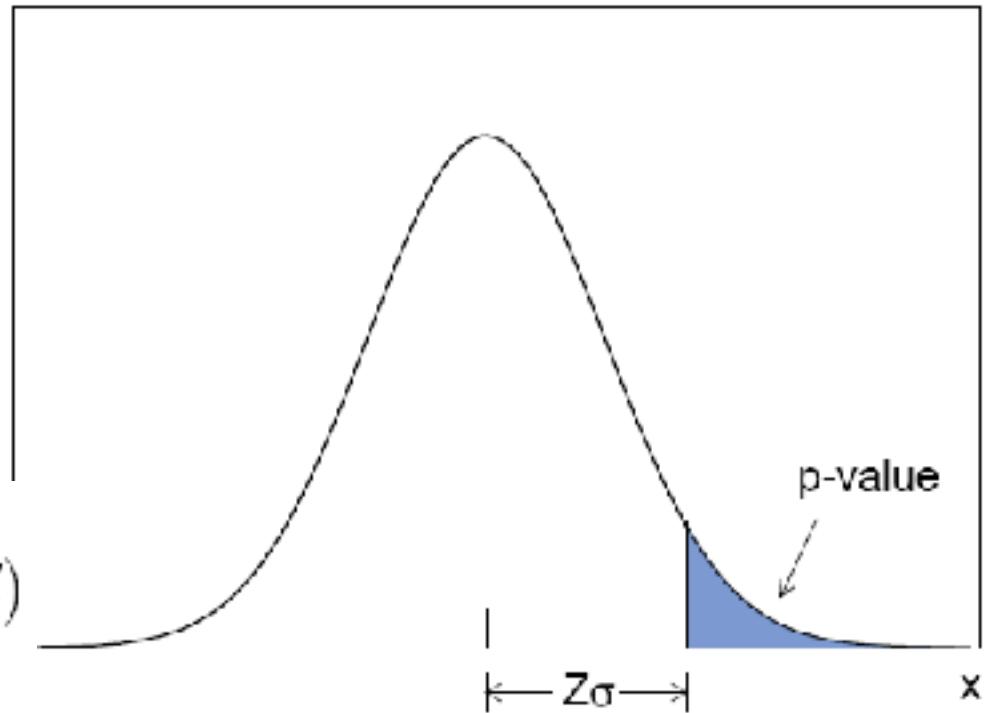
From p-values to Gaussian significance

It is a custom to express the p-value as the significance associated to it, had the pdf were Gaussians

$$p = \int_Z^\infty \frac{1}{\sqrt{2\pi}} e^{-x^2/2} dx = 1 - \Phi(Z)$$

$$Z = \Phi^{-1}(1 - p)$$

A significance of $Z = 5$ corresponds to $p = 2.87 \times 10^{-7}$



Beware of 1 vs 2-sided definitions!



p-value - testing the null hypothesis

When testing the b hypothesis ($\text{null}=b$), it is custom to set

$$\alpha = 2.9 \cdot 10^{-7}$$

→ if $p_b < 2.9 \cdot 10^{-7}$ the b hypothesis is rejected

→ Discovery

When testing the s+b hypothesis ($\text{null}=s+b$), set $\alpha = 5\%$

if $p_{s+b} < 5\%$ the signal hypothesis is rejected at the 95%

Confidence Level (CL)

→ Exclusion



Nuisance Parameters or Systematics



Nuisance Parameters (Systematics)

- There are two kinds of parameters:
 - Parameters of interest (signal strength... cross section... μ)
 - Nuisance parameters (background (b), signal efficiency, resolution, energy scale,...)
- The nuisance parameters carry systematic uncertainties



Pulls and Ranking of NPs

The pull of θ_i is given by $\frac{\hat{\theta}_i - \theta_{0,i}}{\sigma_0}$

without constraint $\sigma\left(\frac{\hat{\theta}_i - \theta_{0,i}}{\sigma_0}\right) = 1 \quad \left\langle \frac{\hat{\theta}_i - \theta_{0,i}}{\sigma_0} \right\rangle = 0$

It's a good habit to look at the pulls of the NPs and make sure that Nothing irregular is seen

In particular one would like to guarantee that the fits do not over constrain a NP in a non sensible way



Implementation of Nuisance Parameters

- Implement by marginalizing (Bayesian) or profiling (Frequentist)
- Marginalization (Integrating)
 - Integrate the Likelihood, L , over possible values of nuisance parameters (weighted by their prior belief functions -- Gaussian, gamma, others...)
 - $$L(\mu) = \int L(\mu, \theta) \pi(\theta) d\theta$$



The Hybrid Cousins-Highland Marginalization

Cousins & Highland

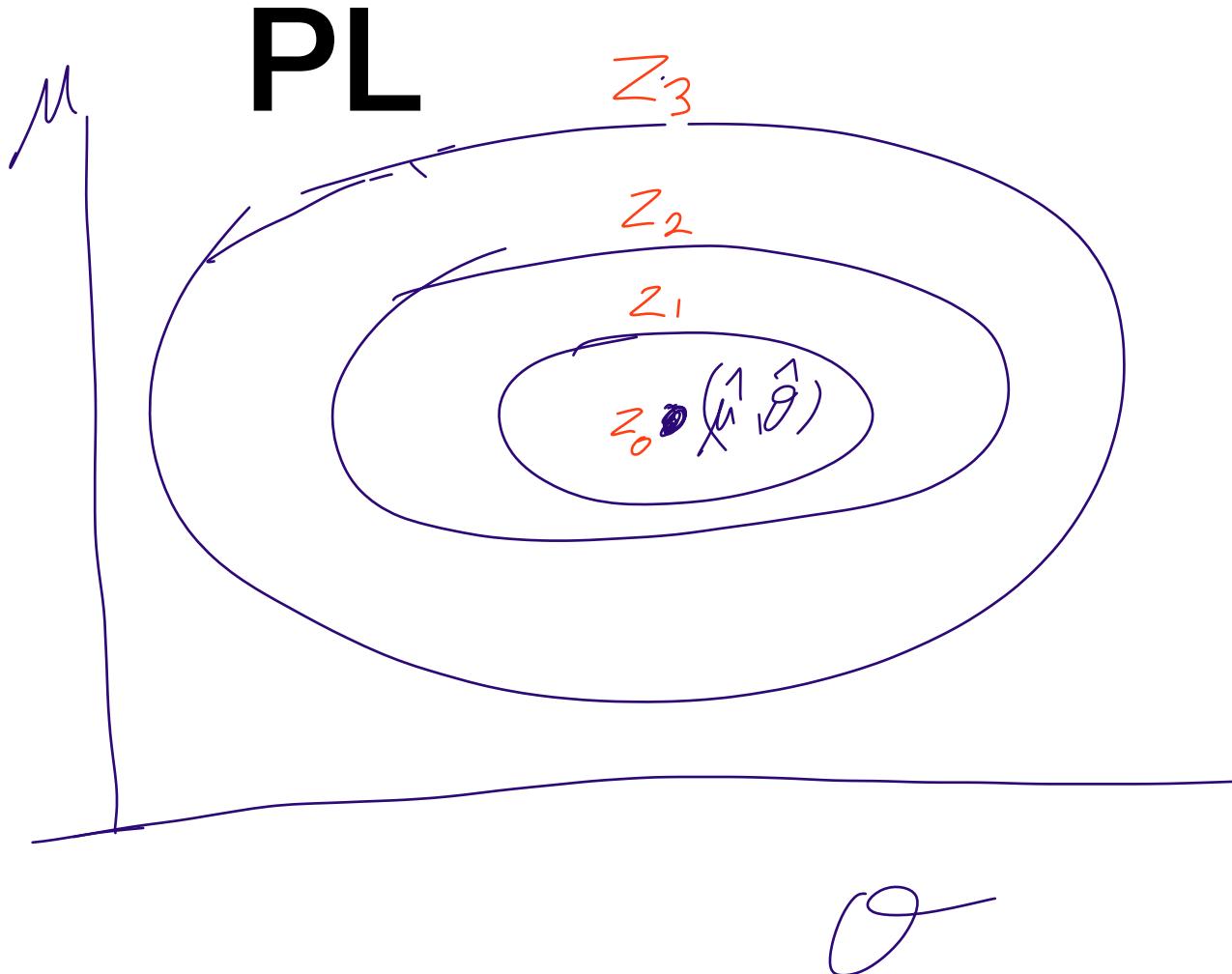
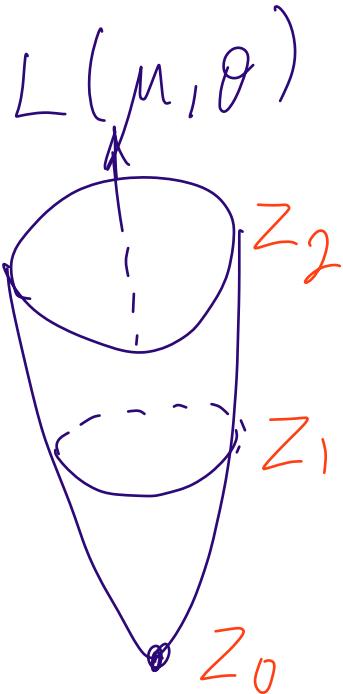
$$q = \frac{L(s + b(\theta))}{L(b(\theta))} \Rightarrow \frac{\int L(s + b(\theta)) \pi(\theta) d\theta}{\int L(b(\theta)) \pi(\theta) d\theta}$$

Profiling the NPs

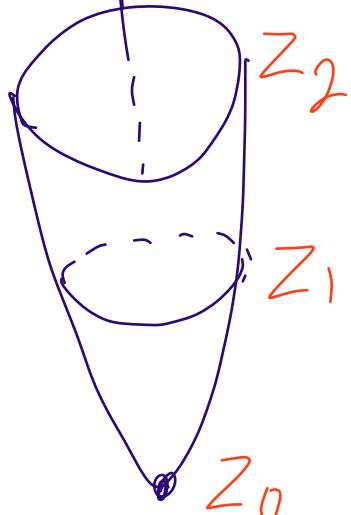
$$q = \frac{L(s + b(\theta))}{L(b(\theta))} \Rightarrow \frac{L(s + b(\hat{\hat{\theta}}_s))}{L(b(\hat{\hat{\theta}}_b))}$$

$\hat{\hat{\theta}}_s$ is the MLE of θ fixing s

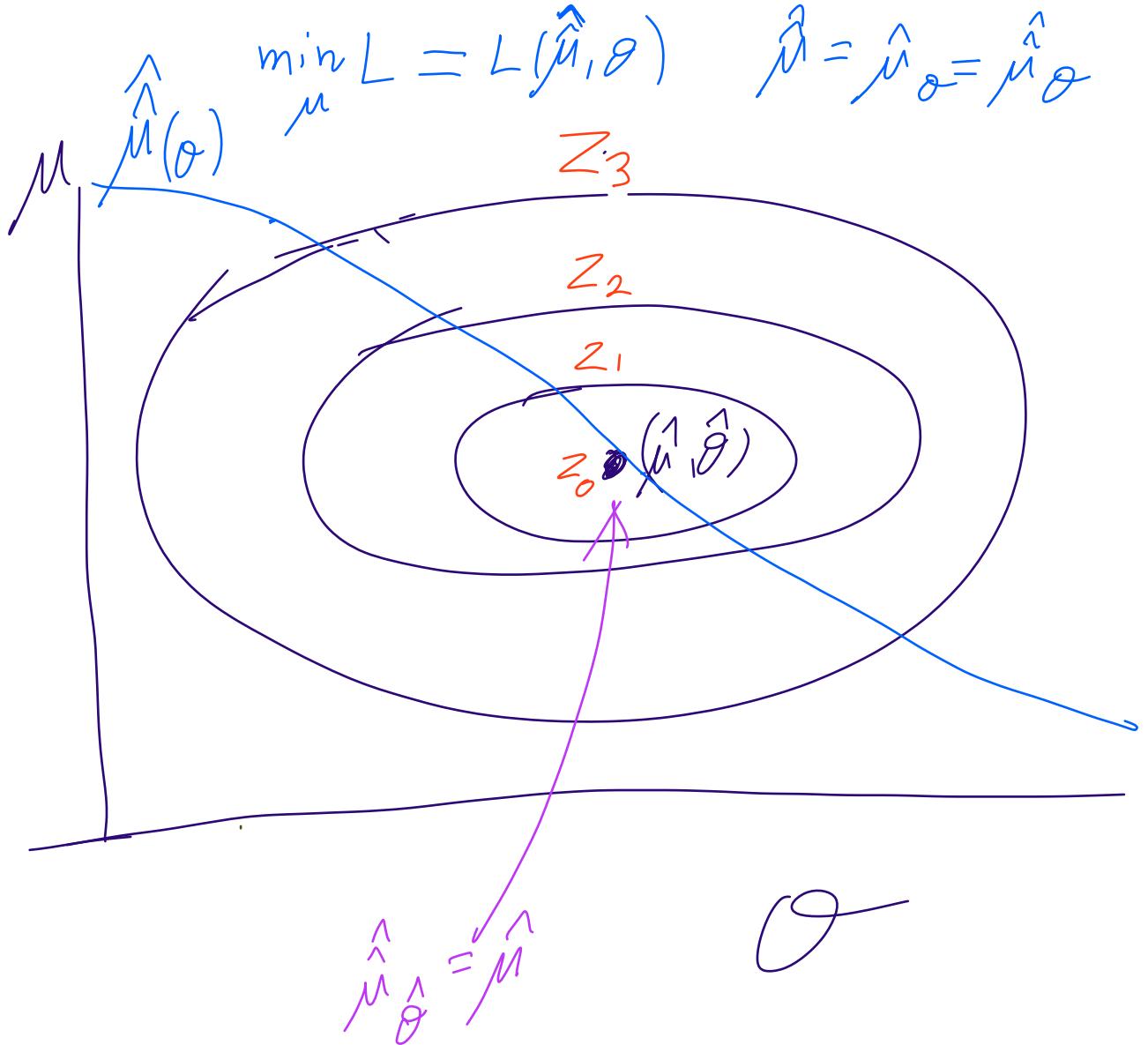




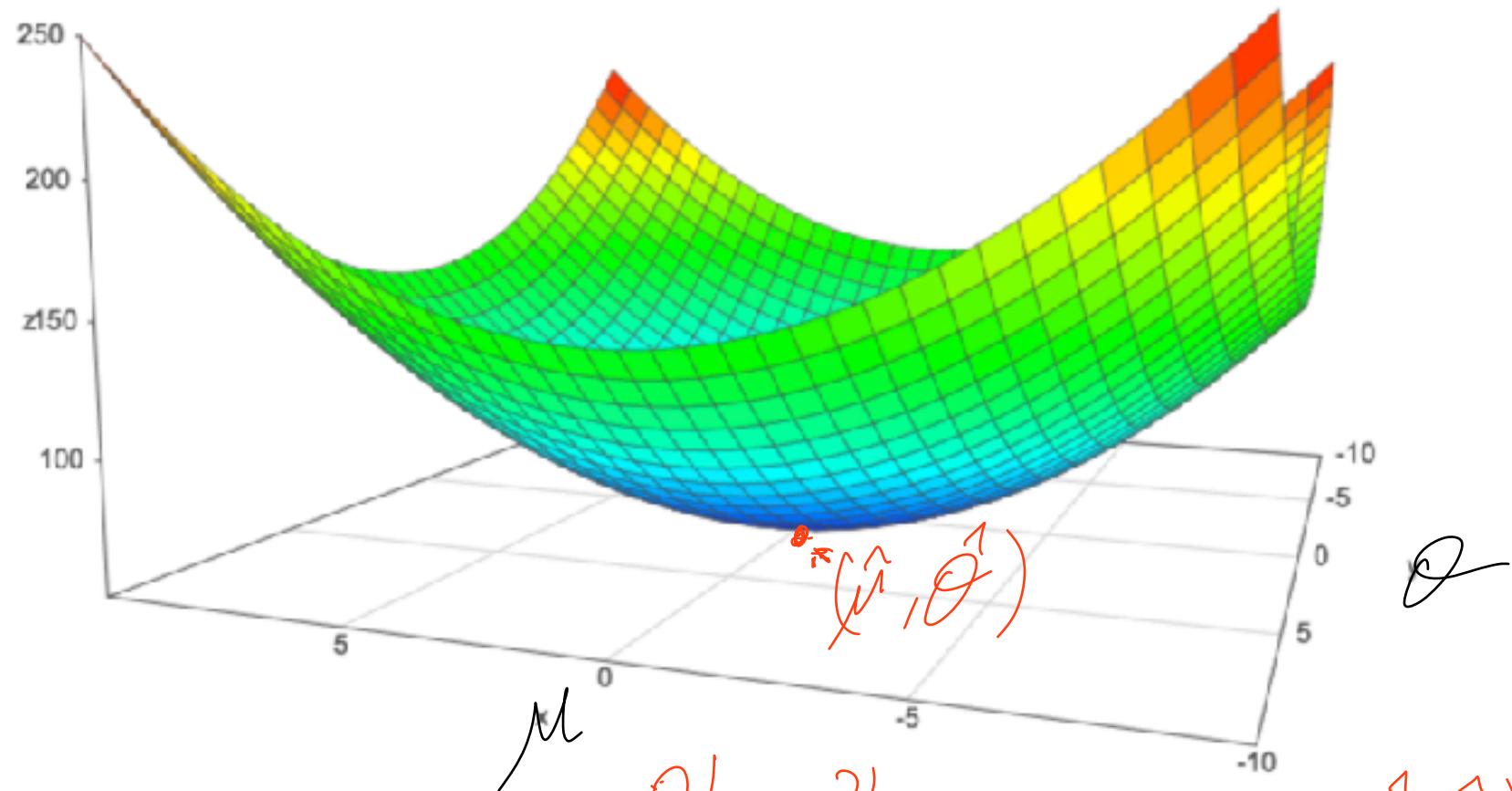
$$L(\mu, \theta)$$



$$\min_{\mu, \theta} L = L(\hat{\mu}, \hat{\theta})$$



$$L = L(\mu, \theta)$$

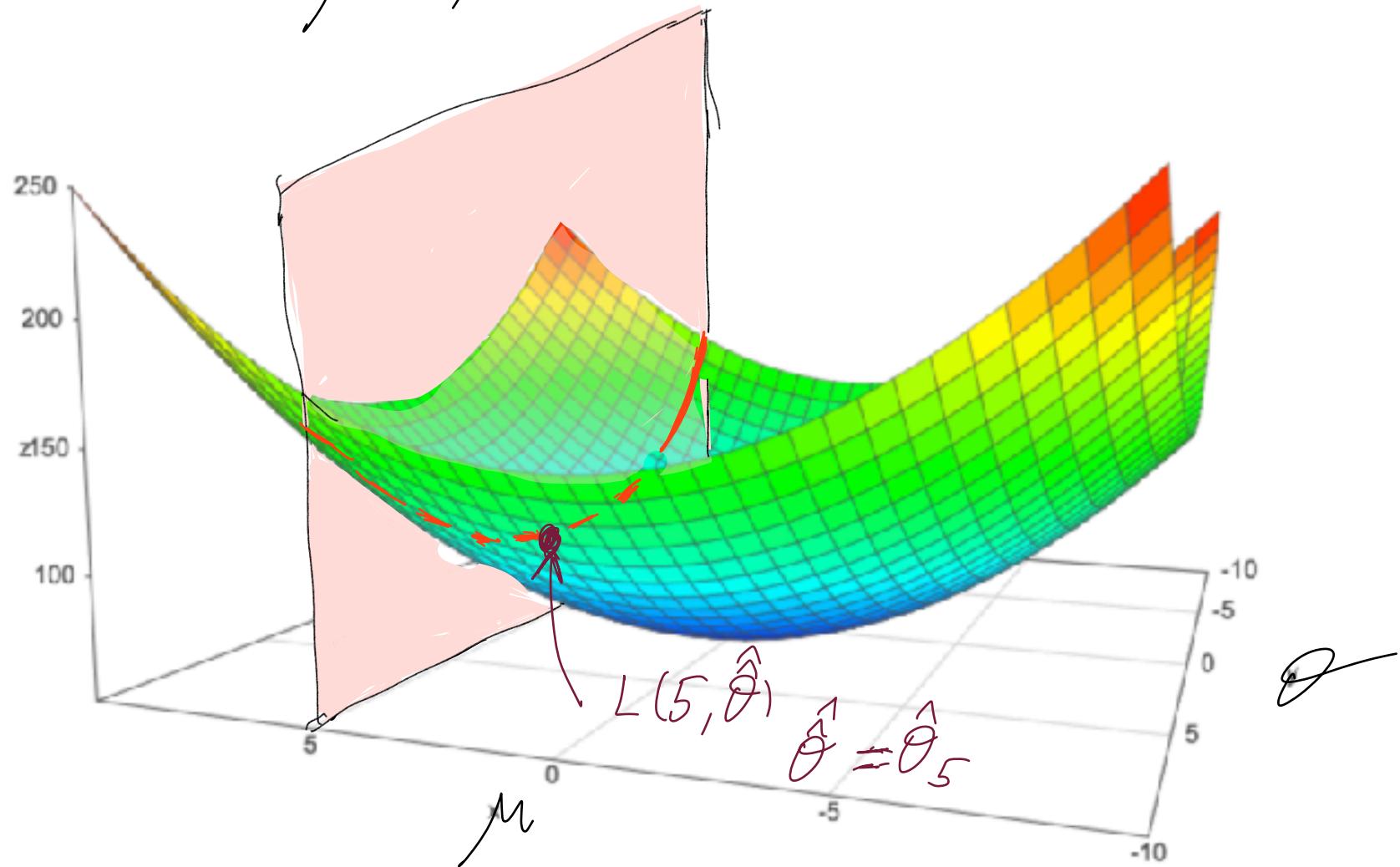


$$\frac{\partial L}{\partial \mu} = \frac{\partial L}{\partial \theta} = 0 \Rightarrow \min L = \hat{L}(\hat{\mu}, \hat{\theta})$$

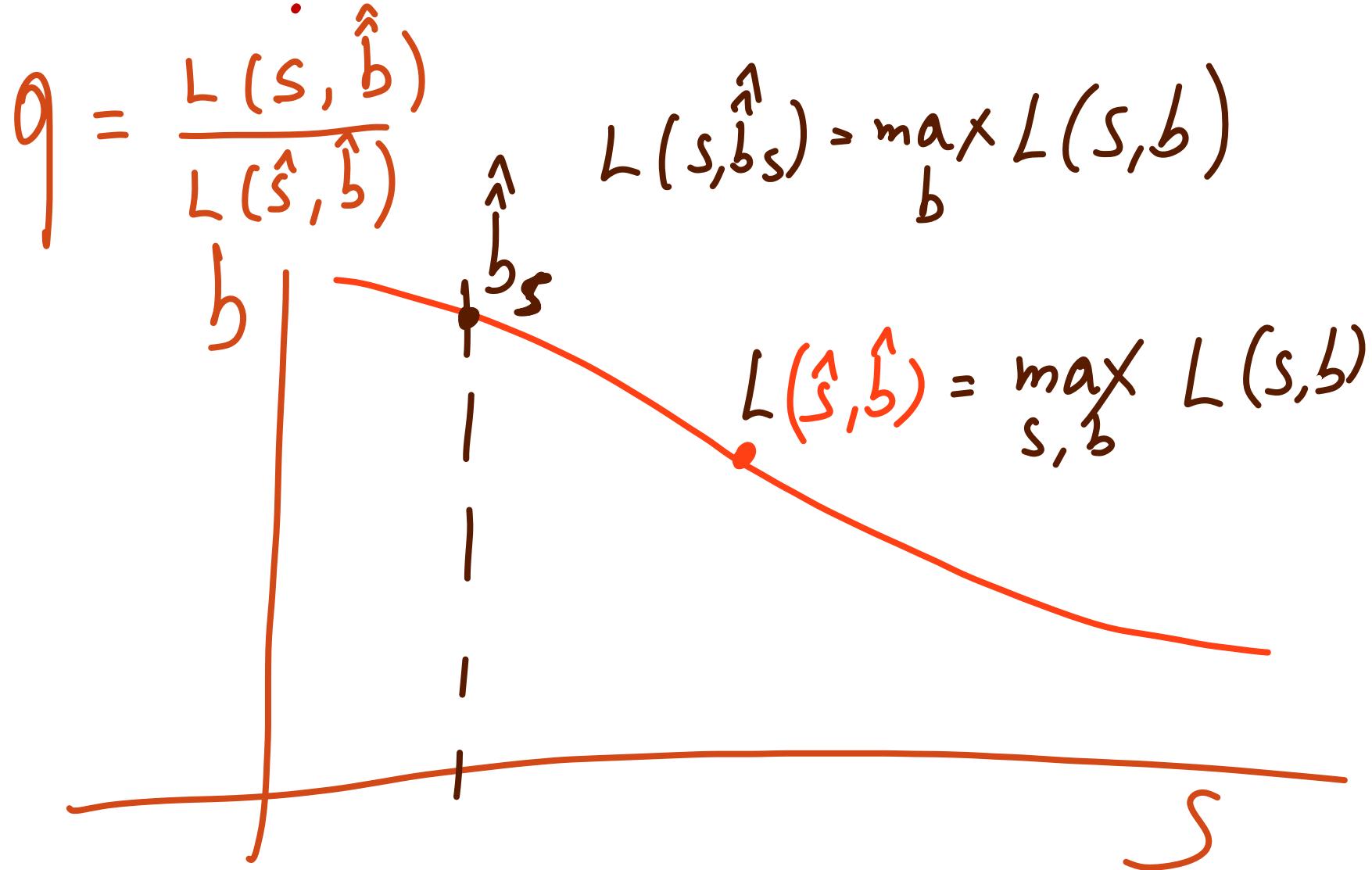


$$L = L(\mu, \theta)$$

fix $\mu=5$



Profile Likelihood



Wilks theorem in the presence of NPs

- Given n parameters of interest and any number of NPs, then

$$\lambda = \frac{L(\mu, \hat{\theta})}{L(\hat{\mu}, \hat{\theta})}$$

$$q(\mu, \theta) = -2 \ln \lambda(\mu) \sim \chi^2$$

$$q(\mu_i, \theta) = -2 \ln \lambda(\mu_i) \sim \chi_n^2 \quad \textit{Wilks Theorem}$$

Profile Likelihood with Nuisance Parameters

$$q_{\mu} = -2 \ln \frac{L(\mu s + \hat{b}_{\mu})}{L(\hat{\mu} s + \hat{b})}$$

$$q_{\mu} = -2 \ln \frac{\max_b L(\mu s + b)}{\max_{\mu, b} L(\mu s + b)}$$

$$q_{\mu} = q_{\mu}(\hat{\mu}) = -2 \ln \frac{L(\mu s + \hat{b}_{\mu})}{L(\hat{\mu} s + \hat{b})}$$

$\hat{\mu}$ MLE of μ

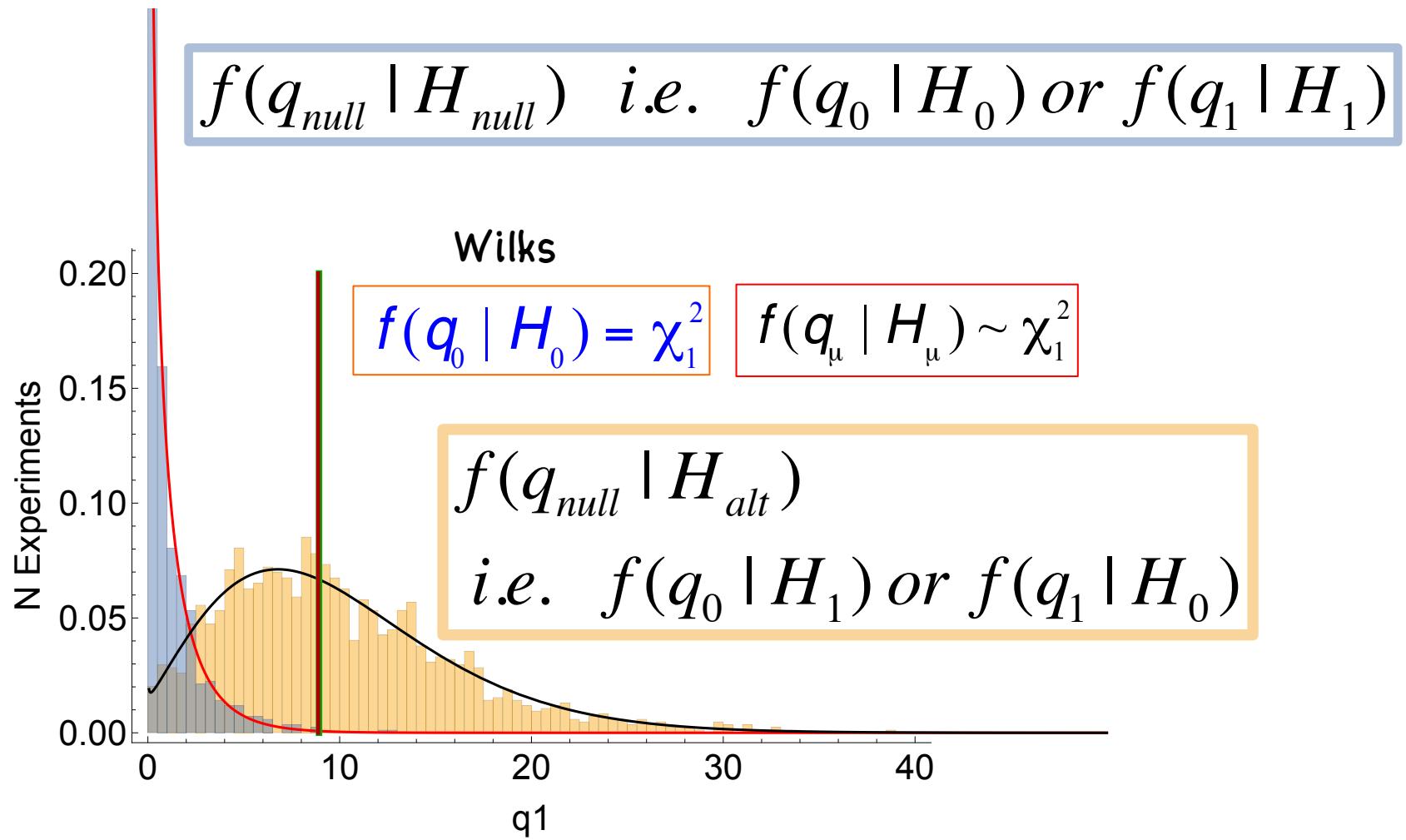
\hat{b} MLE of b

\hat{b}_{μ} MLE of b fixing μ

$\hat{\theta}_{\mu}$ MLE of θ fixing μ



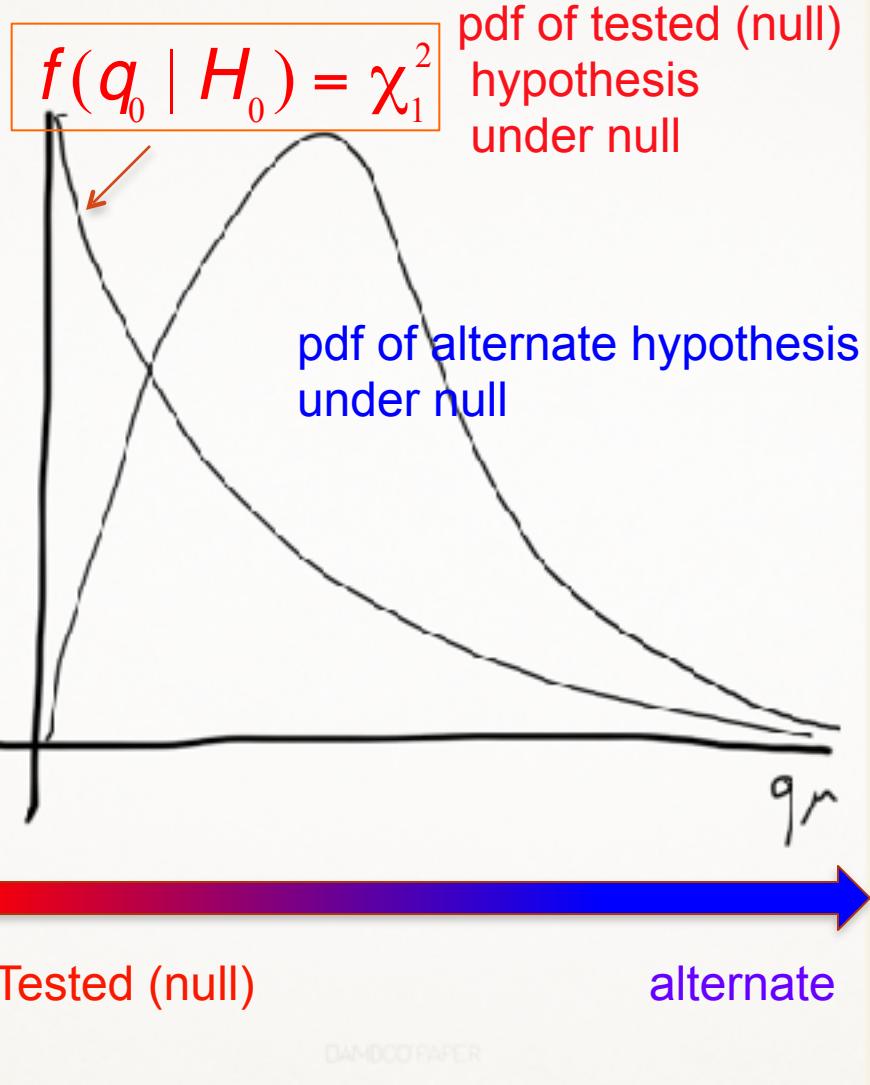
This Lecture's Questions



Wilks Theorem

S.S. Wilks, *The large-sample distribution of the ratio of the variance-covariance matrices of two populations when the sample sizes are unequal*, Ann. Math. Statist. 9 (1938) 60-2.

- Wilks' theorem says that the pdf of the statistic under the null hypothesis approaches a chi-square PDF for one degree of freedom



$$f(q_0 | H_0) = \chi^2_1$$

$$f(q_\mu | H_\mu) \sim \chi^2_1$$



Classification of Test Statistics

Test Stat.	Purpose	Expression	LR
q_0	discovery of positive signal	$q_0 = \begin{cases} -2 \ln \lambda(0) & \hat{\mu} \geq 0 \\ 0 & \hat{\mu} < 0 \end{cases}$	$\lambda(0) = \frac{L(0, \hat{\theta})}{L(\hat{\mu}, \hat{\theta})}$
t_μ	2-sided measurement	$t_\mu = -2 \ln \lambda(\mu)$	$\lambda(\mu) = \frac{L(\mu, \hat{\theta})}{L(\hat{\mu}, \hat{\theta})}$
\tilde{t}_μ	avoid negative signal (Feldman-Cousins)	$\tilde{t}_\mu = -2 \ln \tilde{\lambda}(\mu)$	$\tilde{\lambda}(\mu) = \begin{cases} \frac{L(\mu, \hat{\theta}'(\mu))}{L(\hat{\mu}, \hat{\theta})} & \hat{\mu} \geq 0 \\ \frac{L(\mu, \hat{\theta}'(\mu))}{L(0, \hat{\theta}(0))} & \hat{\mu} < 0 \end{cases}$
q_μ	exclusion	$q_\mu = \begin{cases} -2 \ln \lambda(\mu) & \hat{\mu} \leq \mu \\ 0 & \hat{\mu} > \mu \end{cases}$	
\bar{q}_μ	exclusion of positive signal	$\bar{q}_\mu = \begin{cases} -2 \ln \frac{L(\mu, \hat{\theta}'(\mu))}{L(0, \hat{\theta}(0))} & \hat{\mu} < 0, \\ -2 \ln \frac{L(\mu, \hat{\theta}'(\mu))}{L(\hat{\mu}, \hat{\theta})} & 0 \leq \hat{\mu} \leq \mu \\ 0 & \hat{\mu} > \mu \end{cases}$	



Discovery vs Exclusion

Test H_0 with q_0 , Reject $H_0 \Rightarrow$ Discovery

$$q_0 = \begin{cases} -2 \ln \lambda(0) & \hat{\mu} \geq 0 \\ 0 & \hat{\mu} < 0 \end{cases} \quad \lambda(0) = \frac{L(0, \hat{\theta}_0)}{L(\hat{\mu}, \hat{\theta})}$$

Test $H_\mu(m_H)$ with q_μ Reject $H_\mu(m_H) \Rightarrow$

Exclusion of a Higgs with $m_H \Rightarrow \mu_{up}(m_H)$

$$q_\mu = \begin{cases} -2 \ln \lambda(\mu) & \hat{\mu} \leq \mu \\ 0 & \hat{\mu} > \mu \end{cases} \quad \lambda(\mu) = \frac{L(\mu, \hat{\theta}_\mu)}{L(\hat{\mu}, \hat{\theta})}$$



Asymptotic Approximation

Asymptotic formulae for likelihood-based
tests of new physics

Glen Cowan (Royal Holloway, U. of London), Kyle
Cranmer (New York U.), Eilam Gross, Ofer Vitells
(Weizmann Inst.), Jul 10, 2010, 25 pp.
Published in Eur Phys J C71 (2011) 1554, Erratum:
Eur.Phys.J. C73 (2013) 2501

CCGV



Test Statistic $t_\mu = -2\ln\lambda(\mu)$

$$t_\mu = -2\ln\lambda(\mu) \quad \lambda(\mu) = \frac{L(\mu, \hat{\theta}_\mu)}{L(\hat{\mu}, \hat{\theta})}$$

Higher values of t_μ correspond to increasing incompatibility between the data and μ



Wald Theorem

$$\lambda(\mu) = \frac{L(\mu, \hat{\theta}_\mu)}{L(\hat{\mu}, \hat{\theta})} \quad t_\mu = -2 \ln \lambda(\mu) \quad \text{Wilks} \Rightarrow f(t_\mu \mid \mu) \sim \chi^2_1$$

How does t_μ distributes under $H_{\mu'} (\mu' \neq \mu)$

A. Wald, *Tests of Statistical Hypotheses Concerning Several Parameters When the Number of Observations is Large*, Transactions of the American Mathematical Society, Vol. 54, No. 3 (Nov., 1943), pp. 426-482.

$$t_\mu = -2 \ln \lambda(\mu) = \frac{(\mu - \hat{\mu})^2}{\sigma_{\hat{\mu}}^2} + O\left(1 / \sqrt{N}\right)$$

(Use the Asimov Dataset to estimate σ)

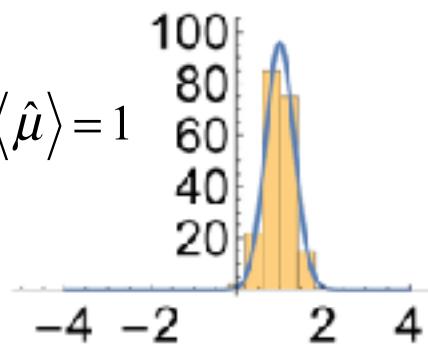
$f(t_\mu \mid \mu')$ follows a noncentral Chi squared distribution

with non-centrality parameter $\Lambda = \frac{(\mu - \mu')^2}{\sigma^2}$ with 1 d.o.f

where $\hat{\mu} \sim G(\mu', \sigma)$

N is the sample size

$$\mu' = 1 \Rightarrow \langle \hat{\mu} \rangle = 1$$



Wald Theorem

$$t_\mu = -2 \ln \lambda(\mu) = \frac{(\mu - \hat{\mu})^2}{\sigma_{\hat{\mu}}^2} + O\left(1/\sqrt{N}\right)$$

$$\hat{\mu} \sim G(\mu', \sigma)$$

N is the sample size

$f(t_\mu | \mu')$ follows a noncentral Chi squared distribution

with non-centrality parameter $\Lambda = \frac{(\mu - \mu')^2}{\sigma^2}$ with 1 d.o.f

$$f(t_\mu; \Lambda) = \frac{1}{2\sqrt{t_\mu}} \frac{1}{\sqrt{2\pi}} \left[\exp\left(-\frac{1}{2} (\sqrt{t_\mu} + \sqrt{\Lambda})^2\right) + \exp\left(-\frac{1}{2} (\sqrt{t_\mu} - \sqrt{\Lambda})^2\right) \right]$$

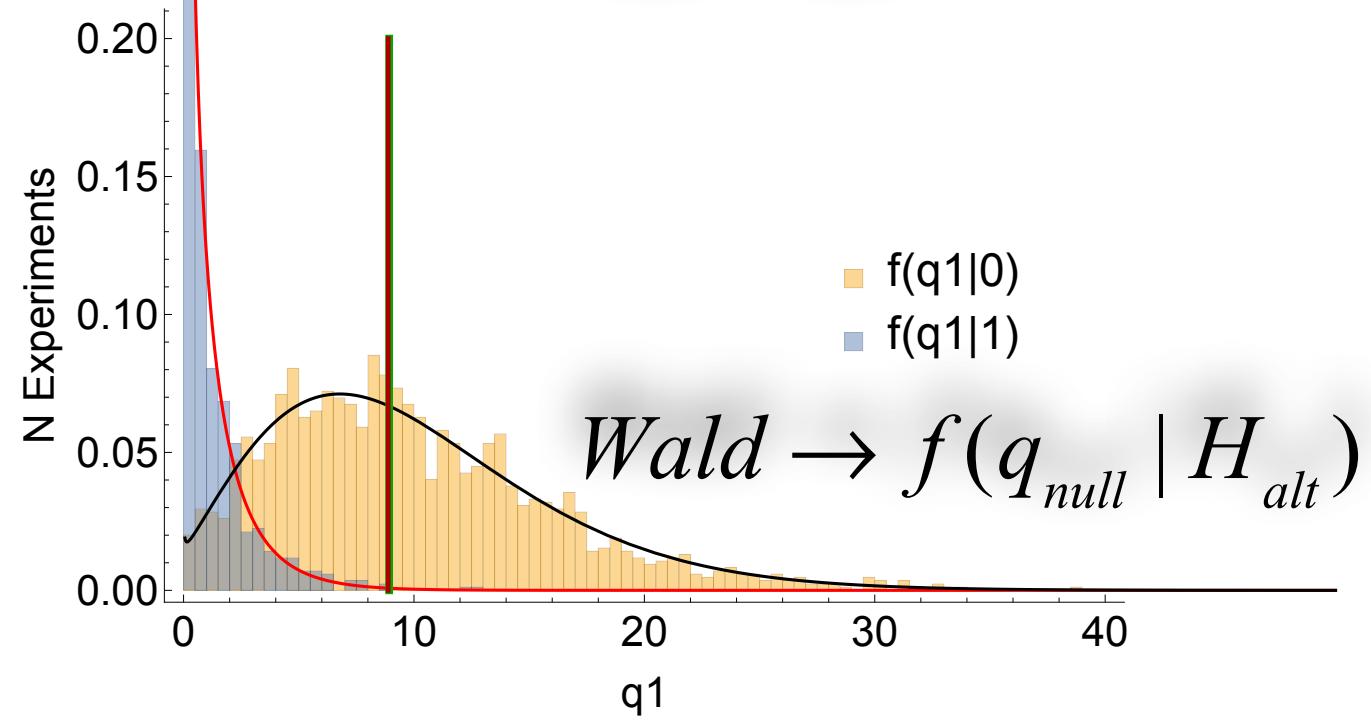
for $\mu' = \mu$ we retrieve Wilks' theorem

$$f(t_\mu) = \frac{1}{\sqrt{2\pi t_\mu}} e^{-\frac{1}{2}t_\mu} = \chi^2$$

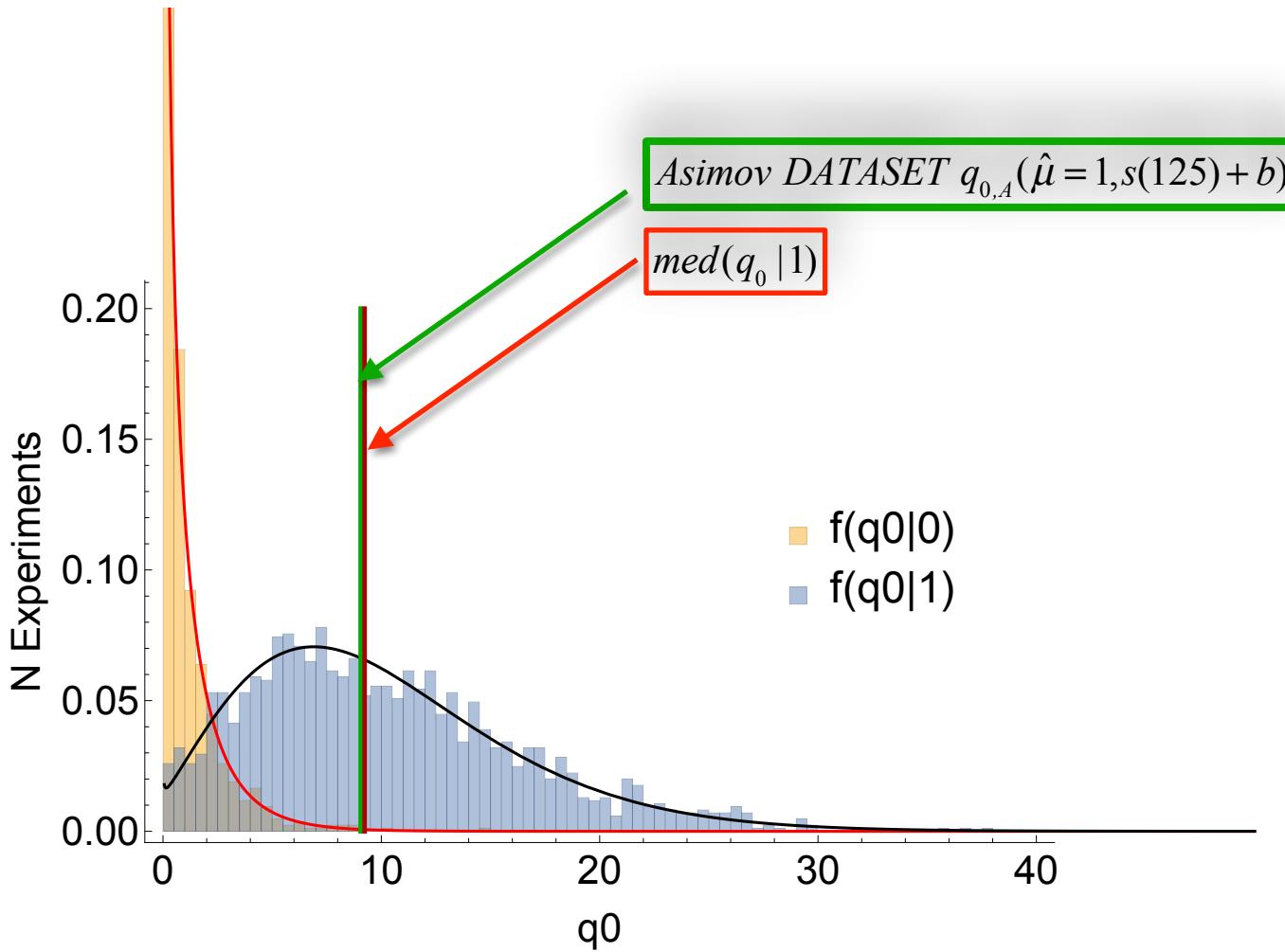


Asymptotics

Wilks $\rightarrow f(q_{null} \mid H_{null}) \sim \chi^2$



The Magic of Asimov



q_μ for exclusion

CCGV

$$q_\mu = \begin{cases} -2 \log \lambda(\mu) & \hat{\mu} \leq \mu \\ 0 & \hat{\mu} > \mu \end{cases}$$

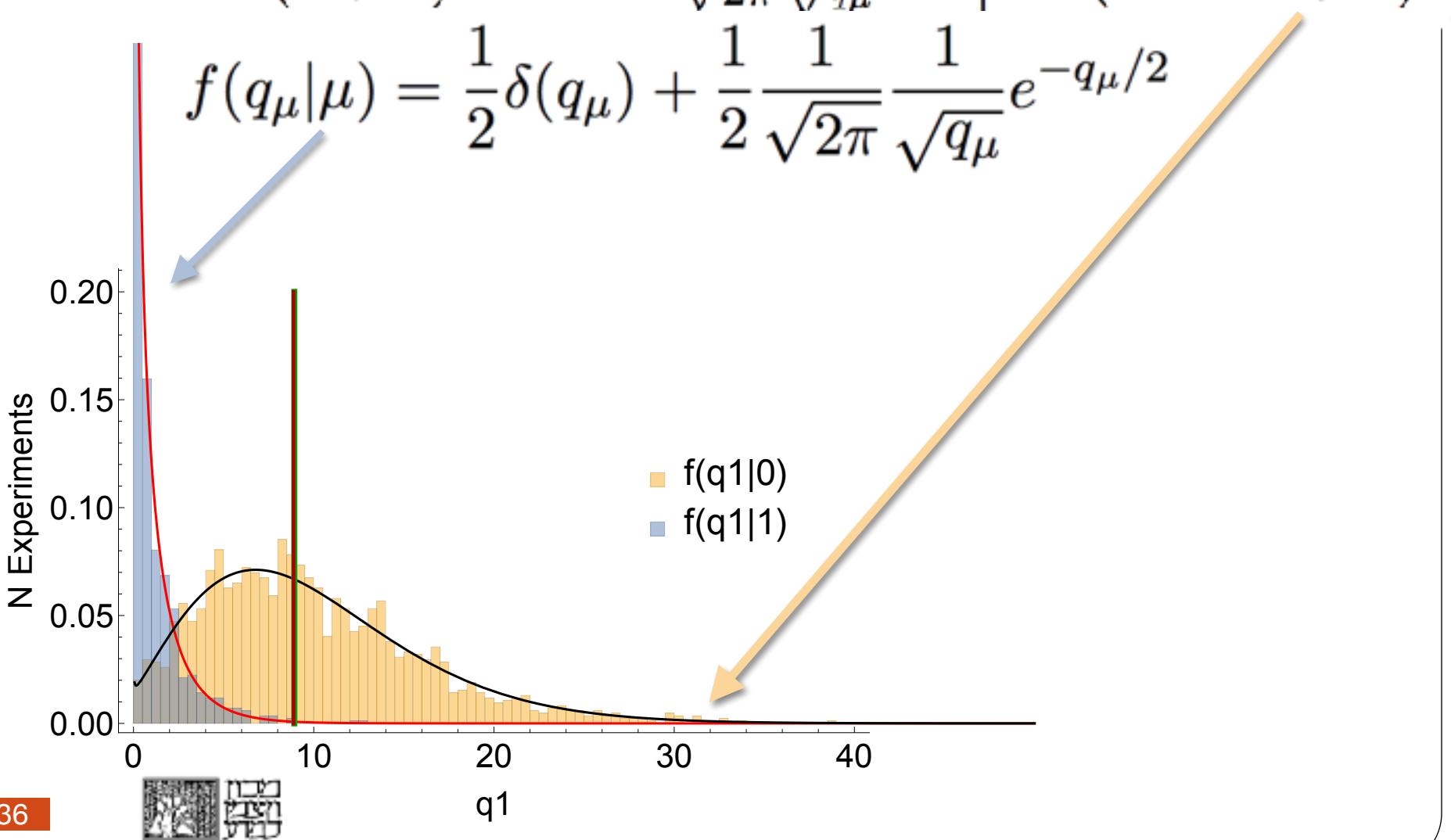
Upward fluctuations of the signal
do not serve as an evidence against the signal



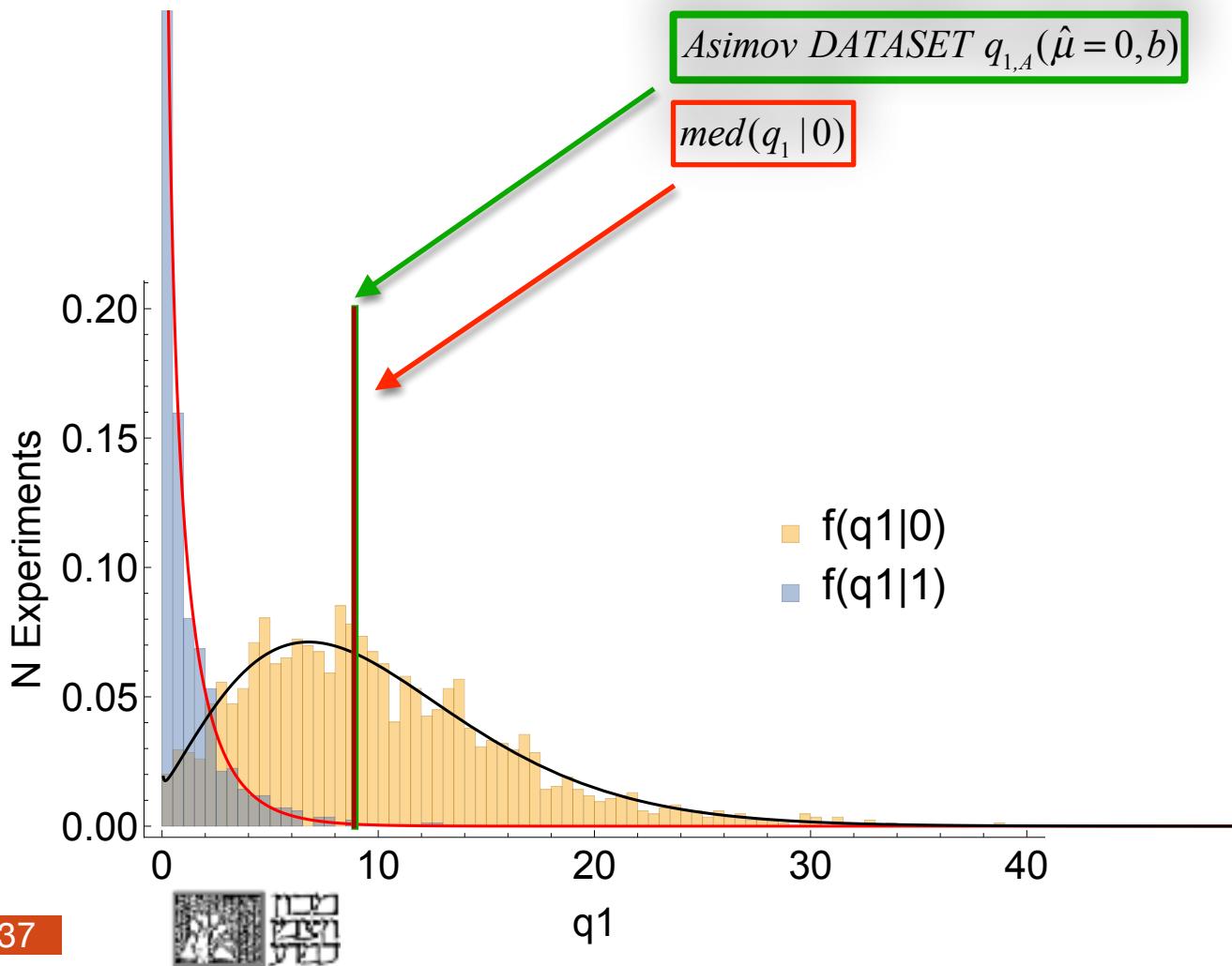
PDF of $(q_1|1)$ and $(q_1|0)$

$$f(q_\mu|\mu') = \Phi\left(\frac{\mu' - \mu}{\sigma}\right) \delta(q_\mu) + \frac{1}{2} \frac{1}{\sqrt{2\pi}} \frac{1}{\sqrt{q_\mu}} \exp\left[-\frac{1}{2} \left(\sqrt{q_\mu} - \frac{\mu - \mu'}{\sigma}\right)^2\right]$$

$$f(q_\mu|\mu) = \frac{1}{2} \delta(q_\mu) + \frac{1}{2} \frac{1}{\sqrt{2\pi}} \frac{1}{\sqrt{q_\mu}} e^{-q_\mu/2}$$

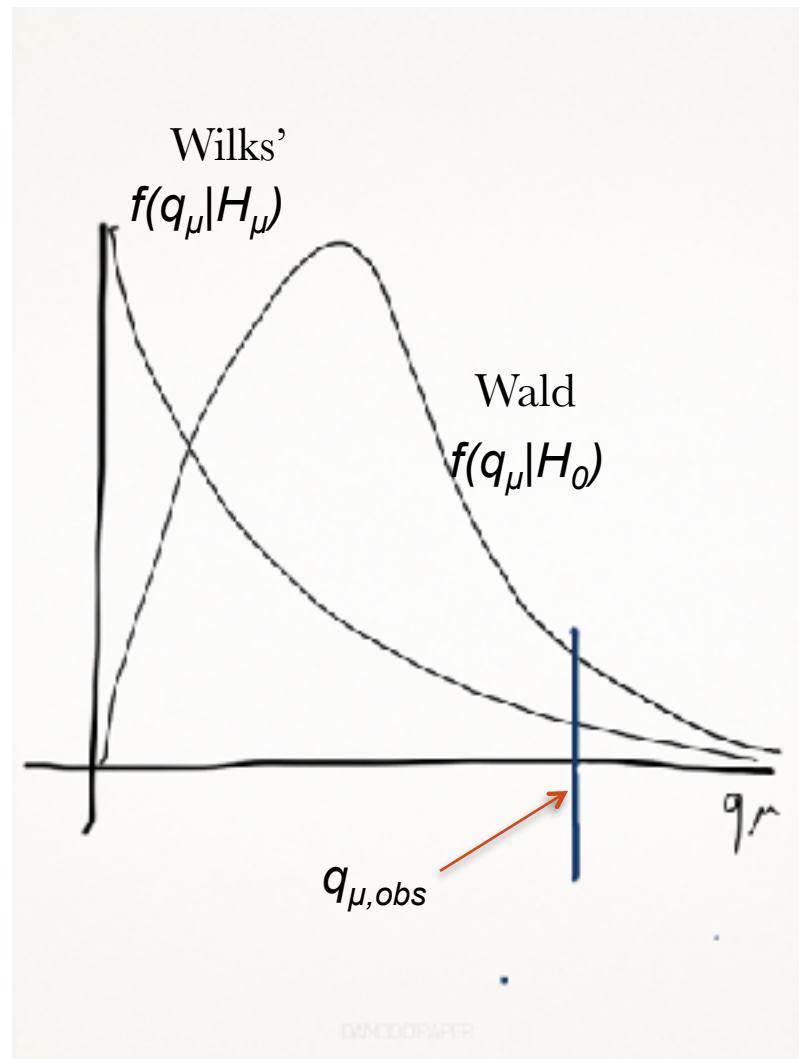


PDF of $(q_1|1)$ and $(q_1|0)$



Exclusion at 95% CL

- We test hypothesis H_μ
- We calculate the PL (profile likelihood) ratio with the one observed data
- $q_{\mu,obs}$

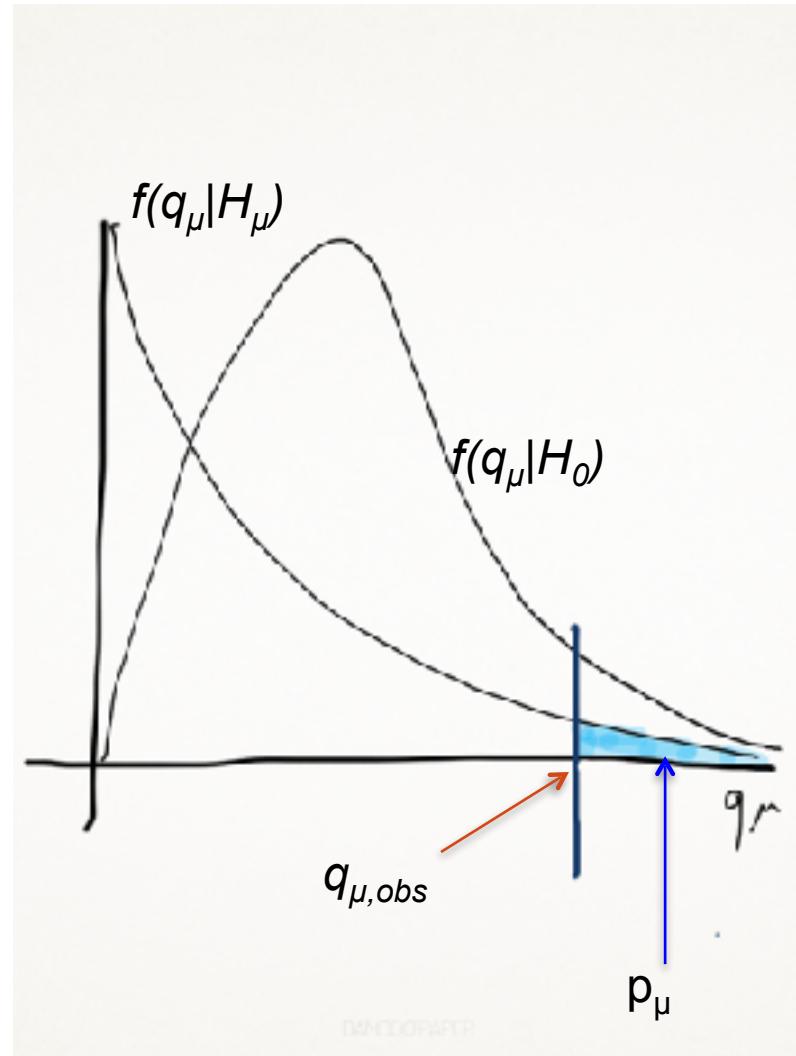


Exclusion at the 95% CL

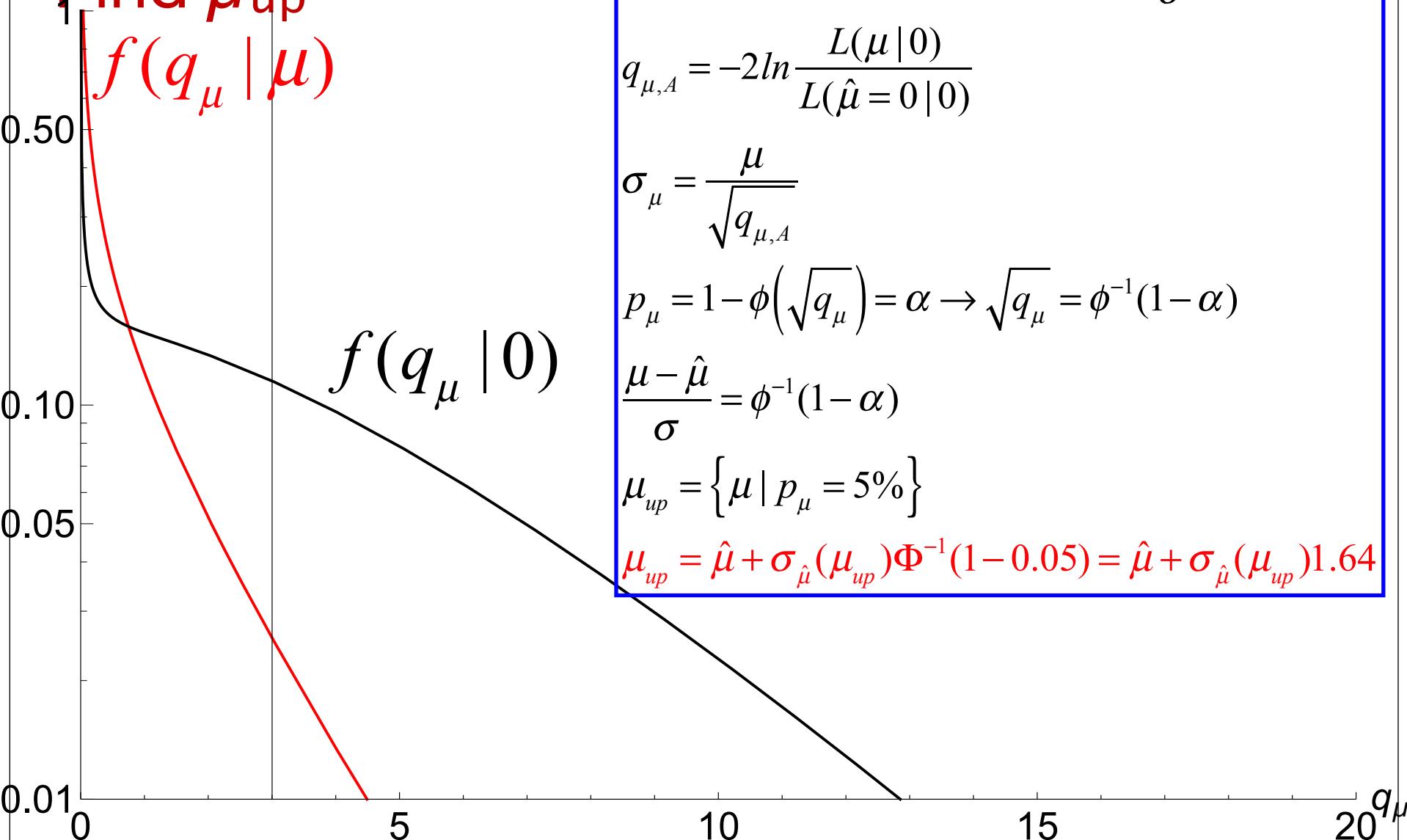
- Find the p-value of the signal hypothesis H_μ

$$p_\mu = \int_{q_{\mu,obs}}^{\infty} f(q_\mu | H_\mu) dq_\mu$$

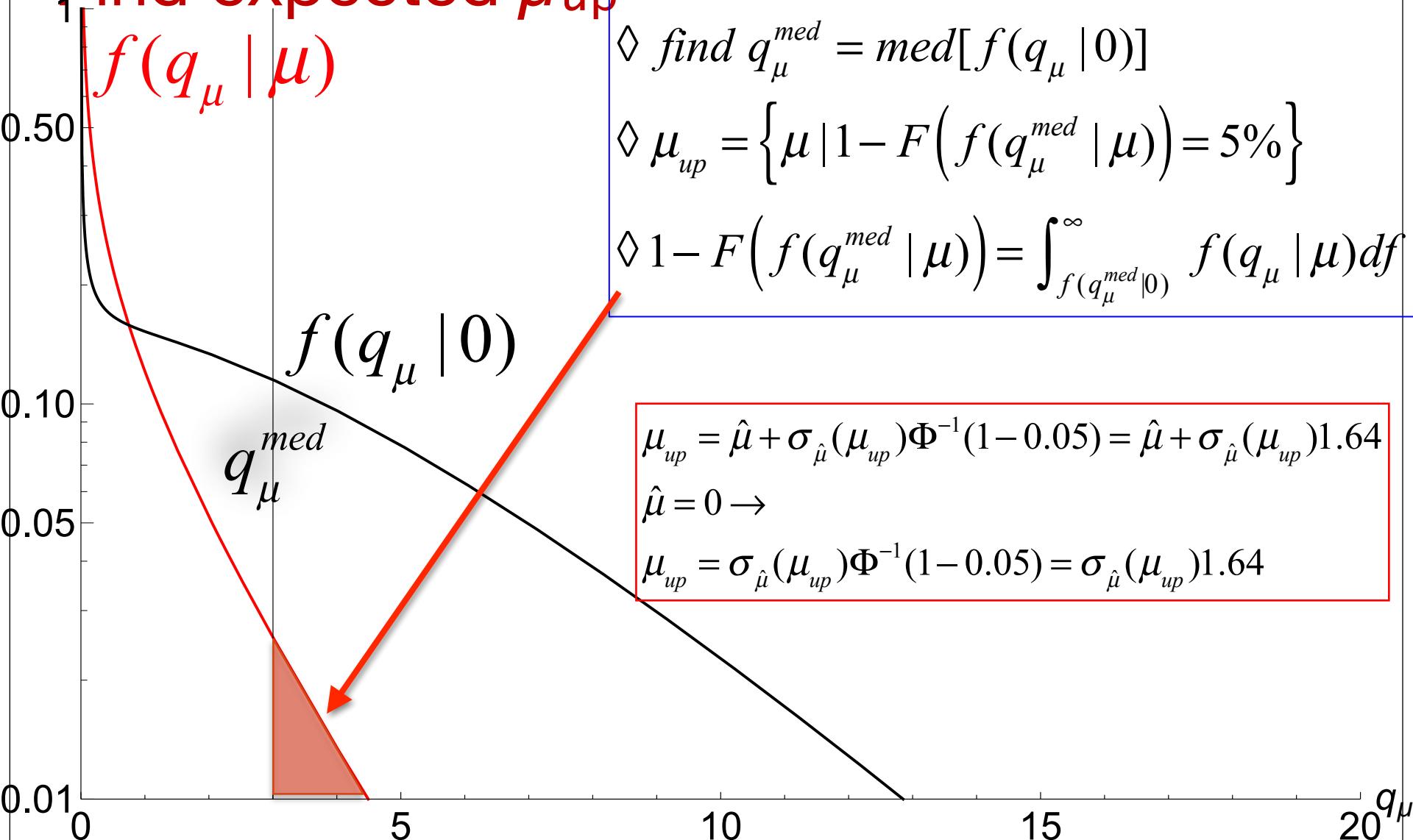
- In principle if $p_\mu < 5\%$, H_μ hypothesis is excluded at the 95% CL
- Note that H_μ is for a given Higgs mass m_H



Find μ_{up}



Find expected μ_{up}

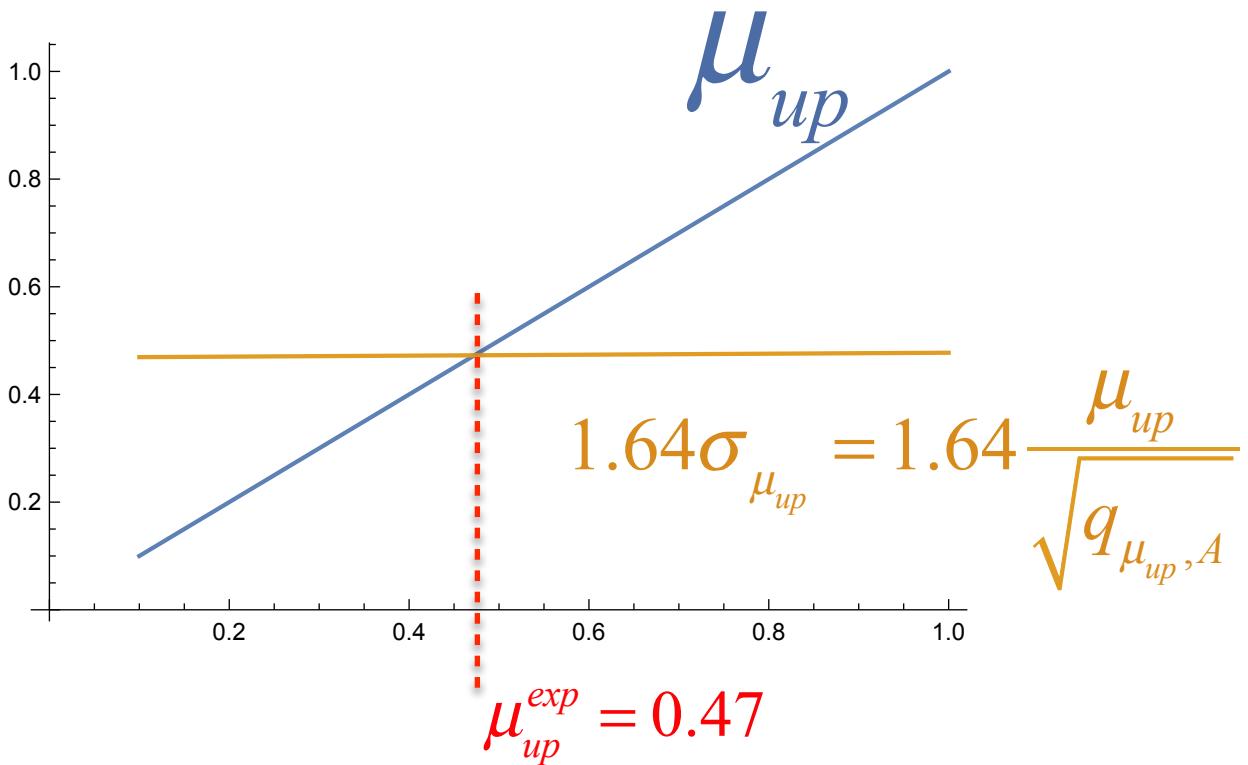


Find expected μ_{up}

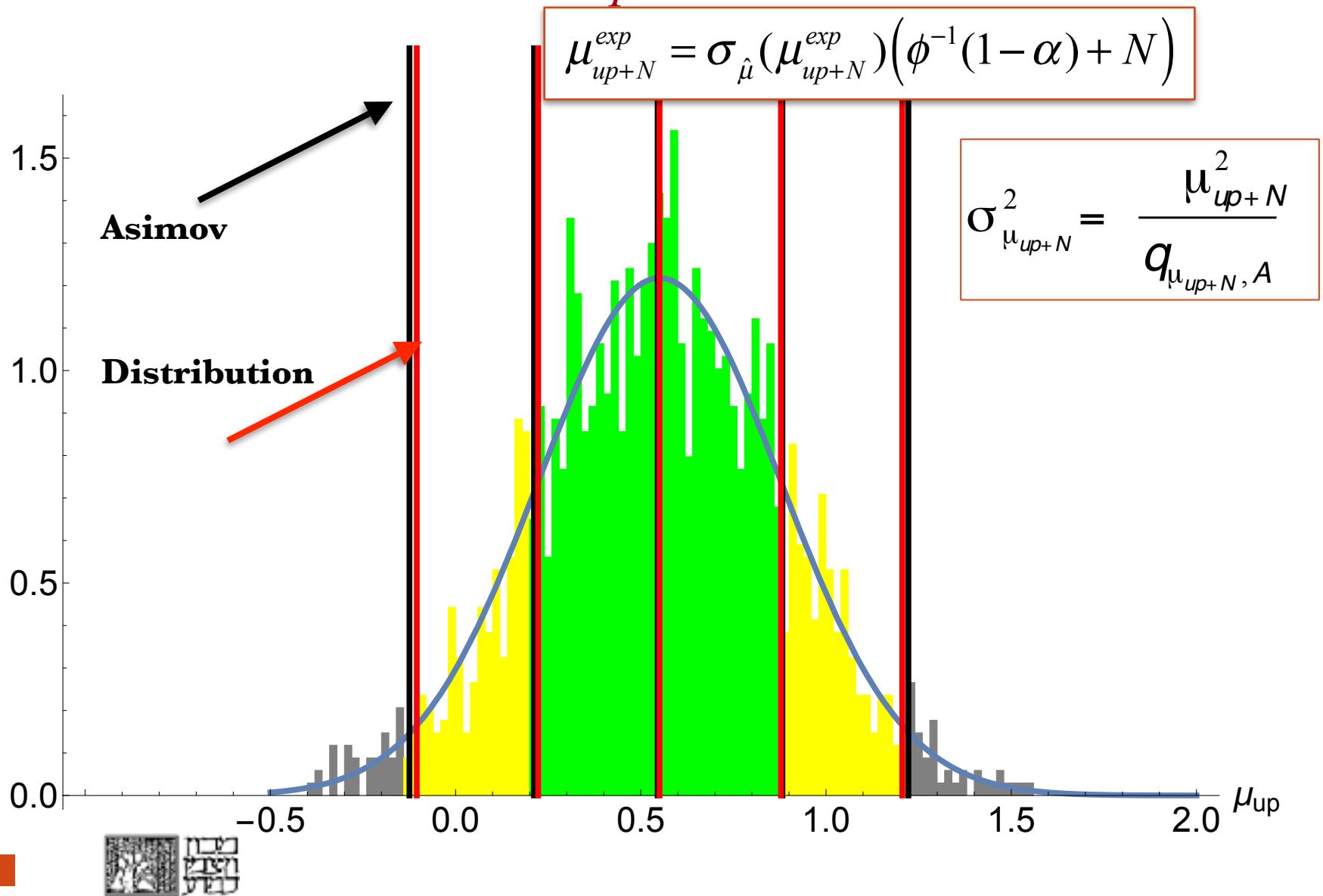
$$\mu_{up} = \hat{\mu} + \sigma_{\hat{\mu}}(\mu_{up})\Phi^{-1}(1 - 0.05) = \hat{\mu} + \sigma_{\hat{\mu}}(\mu_{up})1.64$$

$$\hat{\mu}_A = 0 \rightarrow$$

$$\mu_{up}^{exp} = \sigma_{\hat{\mu}}(\mu_{up})\Phi^{-1}(1 - 0.05) = \sigma_{\hat{\mu}}(\mu_{up}^{exp})1.64$$



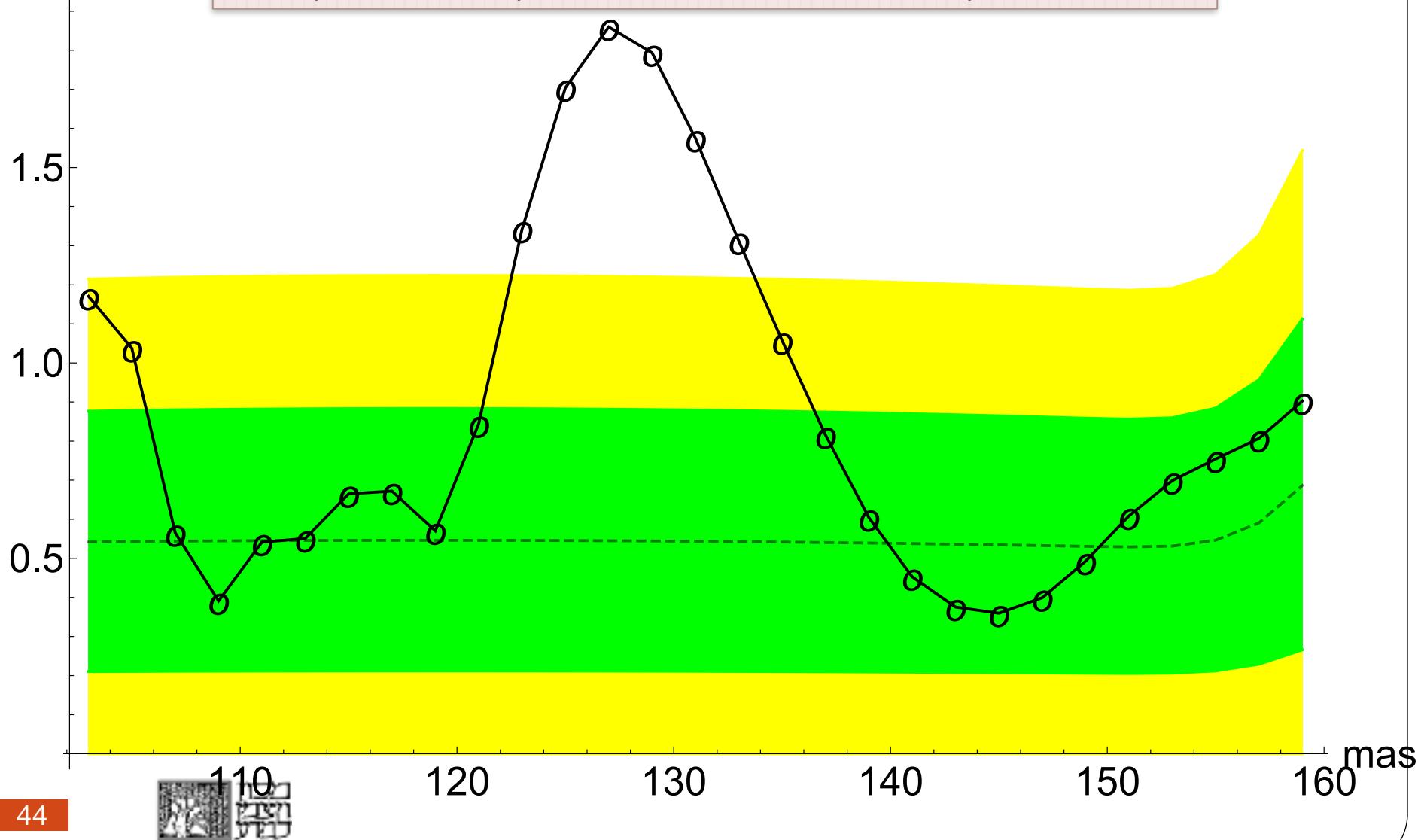
Expected μ_{up} Bands at $m=125$



Brazil Plot

μ_{up}

Every Discovery starts with the inability to exclude



$$q_0 \equiv \tilde{t}_0 = \begin{cases} -2 \ln \lambda(0) & \hat{\mu} \geq 0 \\ 0 & \hat{\mu} < 0 \end{cases} \quad \lambda(0) = \frac{L(\mu=0)}{L(\hat{\mu})} = \frac{L(\hat{b}_{\mu=0})}{L(\hat{\mu}s + \hat{b})} = \frac{L(\hat{b}_{\mu=0})}{L(\hat{s} + \hat{b})}$$

q_0 for discovery

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$$q_0 \equiv \tilde{t}_0 = \begin{cases} -2 \ln \lambda(0) & \hat{\mu} \geq 0 \\ 0 & \hat{\mu} < 0 \end{cases}$$

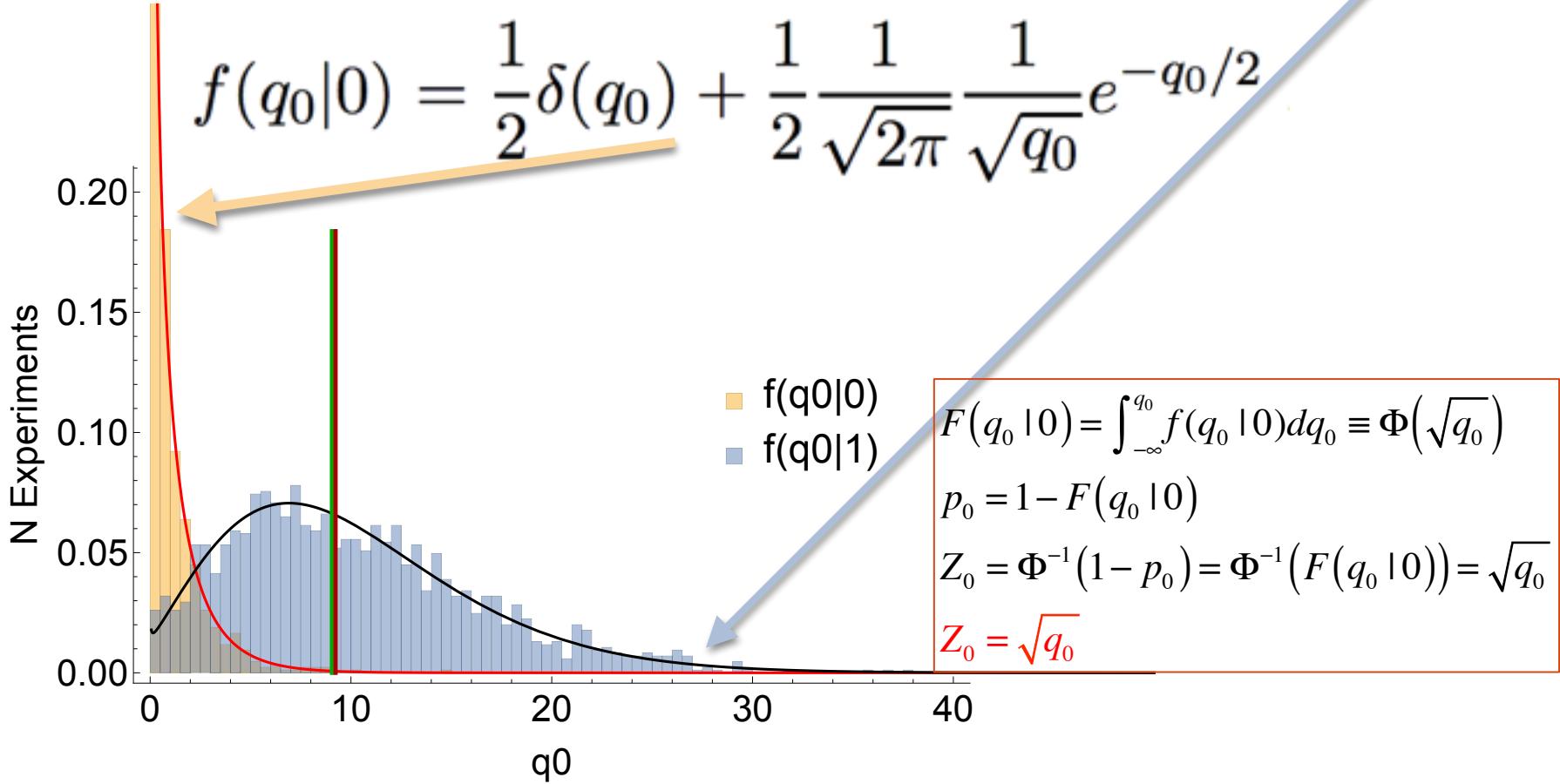
Downward fluctuations of the background
do not serve as an evidence against the background



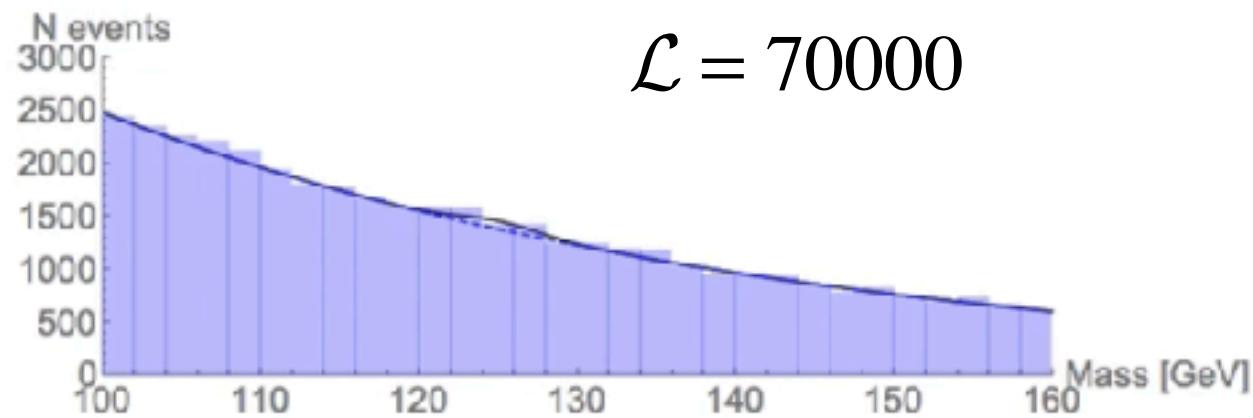
PDF of $(q_0|0)$ and $(q_0|1)$

$$f(q_0|\mu') = \left(1 - \Phi\left(\frac{\mu'}{\sigma}\right)\right) \delta(q_0) + \frac{1}{2} \frac{1}{\sqrt{2\pi}} \frac{1}{\sqrt{q_0}} \exp\left[-\frac{1}{2} \left(\sqrt{q_0} - \frac{\mu'}{\sigma}\right)^2\right]$$

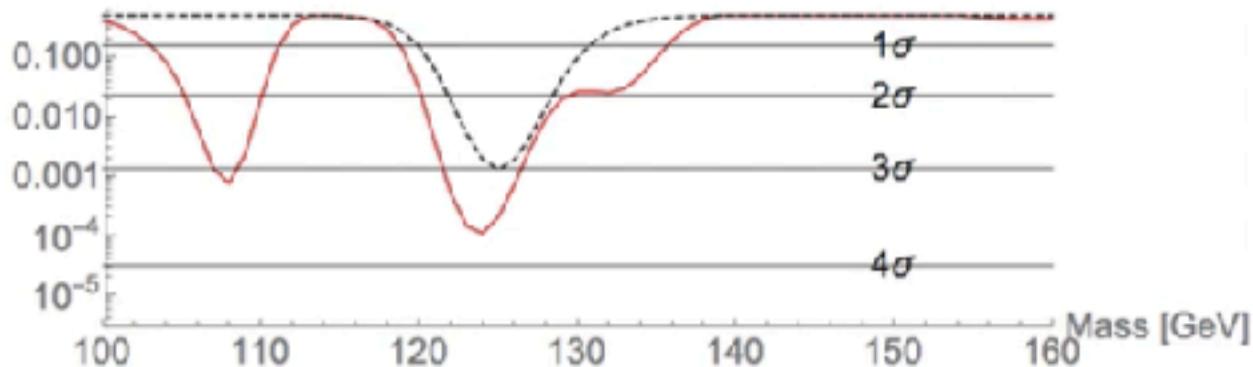
$$f(q_0|0) = \frac{1}{2} \delta(q_0) + \frac{1}{2} \frac{1}{\sqrt{2\pi}} \frac{1}{\sqrt{q_0}} e^{-q_0/2}$$



p-value

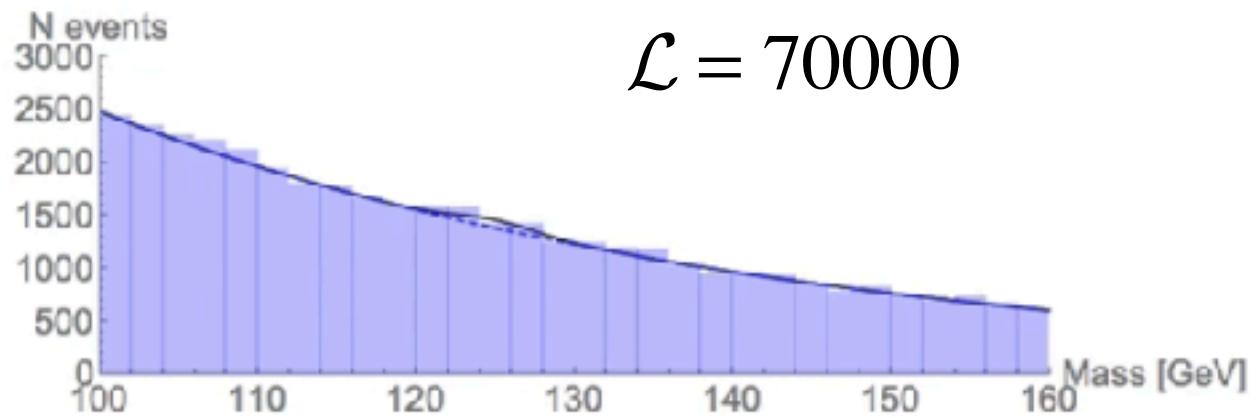


$$p = \text{prob}(q_0 \geq q_{0,obs})$$

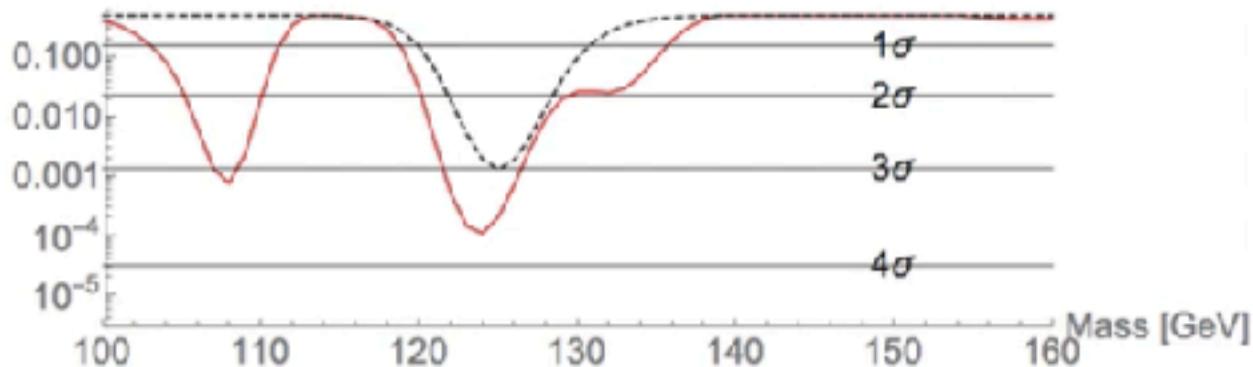


p-value

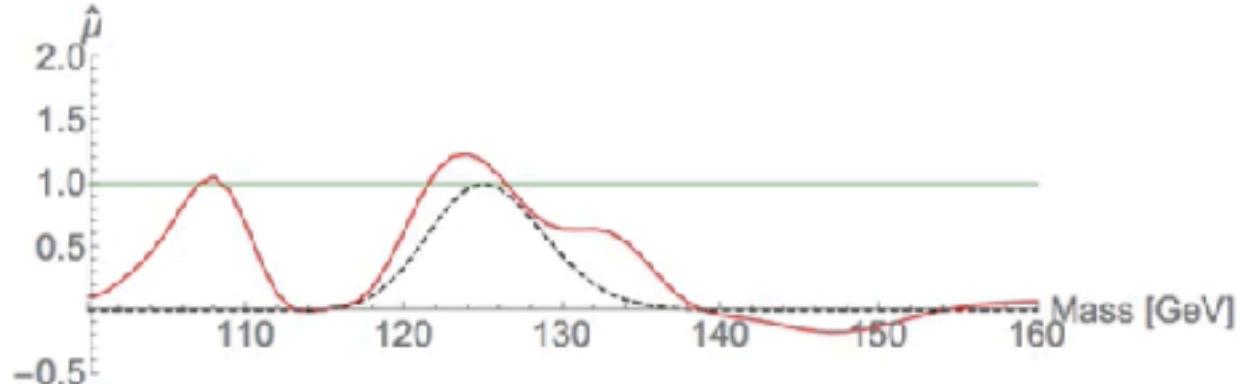
$$\mathcal{L} = 70000$$



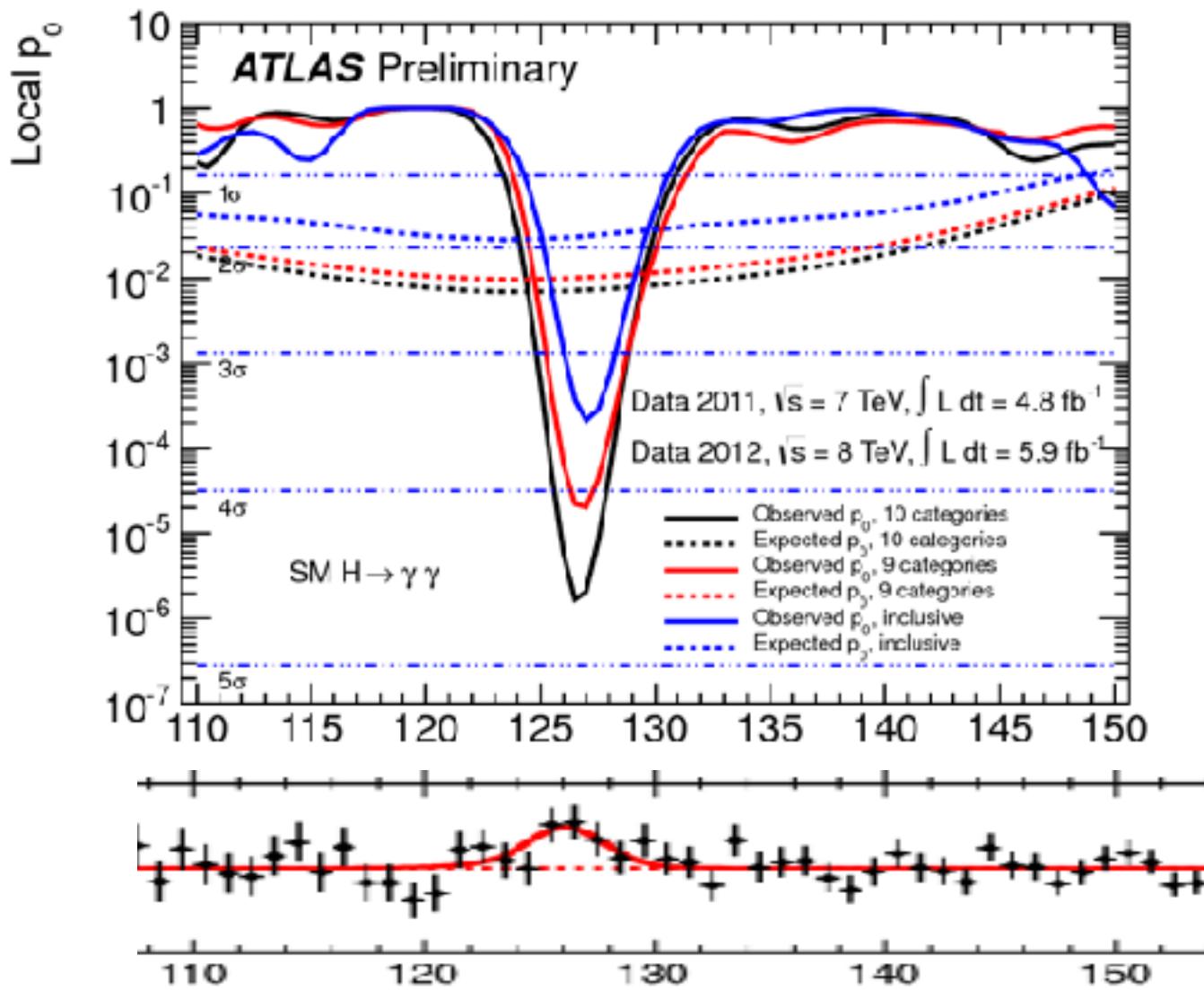
$$p = \text{prob}(q_0 \geq q_{0,obs})$$



$$\hat{\mu} = \frac{\sigma_{obs}}{\sigma_{exp}}$$



$H \rightarrow \gamma\gamma$



More Magic



The New s/\sqrt{b}

The new s/\sqrt{b}

$$Z_A = \sqrt{q_{0,A}}$$

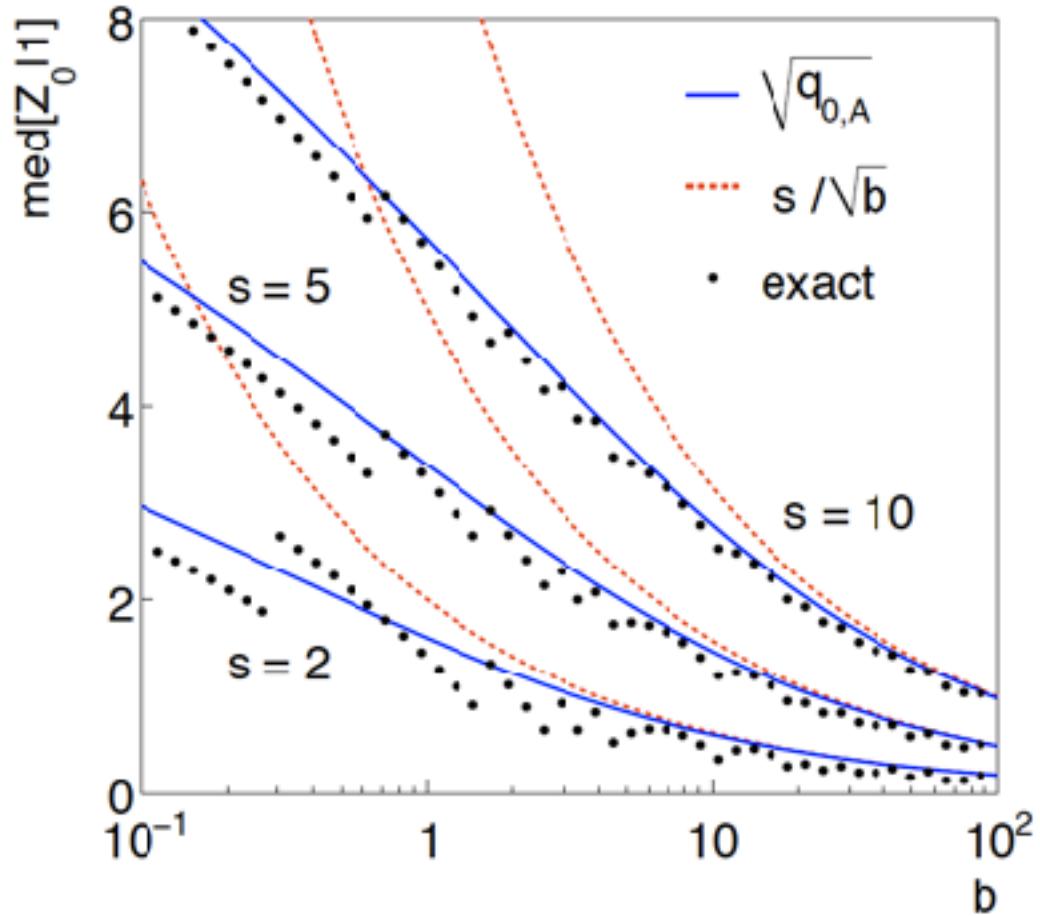
$$\text{med}[Z_0|1] = \sqrt{q_{0,A}} = \sqrt{2((s+b)\ln(1+s/b) - s)}$$

$$Z_A = \sqrt{q_{0,A}} \xrightarrow{s/b \ll 1} \frac{s}{\sqrt{b}} + O(s/b)$$



The New s/\sqrt{b}

s/\sqrt{b} ?



The new s/\sqrt{b}

$$\text{med}[Z_0|1] = \sqrt{q_{0,A}} = \sqrt{2((s+b)\ln(1+s/b) - s)}$$



Next: Look Elsewhere Effect

Next: Look Elsewhere Effect



E.G., O. Vitells “Trial factors for the look elsewhere effect in high energy physics”,

Eur. Phys. J. C 70 (2010) 525

O. Vitells and E. G., Estimating the significance of a signal in a multi-dimensional search,

1669 Astropart. Phys. 35 (2011) 230, arXiv:1105.4355