

The Little Neutral One

Brookhaven National Laboratory

Neutrinos: /

Finding $oldsymbol{
u}$ Cosmic rays and $oldsymbol{
u}$

Disappearing

u Mixing

CP Violaito

A Anne

Conclusions

The Little Neutral One A brief history of neutrinos

Mary Bishai Brookhaven National Laboratory

June 30th, 2020



Neutrinos in Popular Culture

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1964



From Symmetry Magazine, Feb 2013

Cosmic Gall

by John Updike

Neutrinos, they are very small.

They have no charge and have no mass And do not interact at all.

The earth is just a silly ball

To them, through which they simply pass,

Like dustmaids down a drafty hall

Or photons through a sheet of glass.

They snub the most exquisite gas,

Ignore the most substantial wall.

Cold-shoulder steel and sounding brass,

Insult the stallion in his stall,

And, scorning barriers of class,

Infiltrate you and me! Like tall

And painless guillotines, they fall

Down through our heads into the grass.

At night, they enter at Nepal

And pierce the lover and his lass

From underneath the bed-you call

It wonderful; I call it crass.



Neutrinos in Popular Culture

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Jan 7, 2011

NASA's silly sci-fi film list – 2012 the most flawed by Lin Edwards , PhysOrg.com (PhysOrg.com) – At a conference held at the Jet Propulsion Laboratory in California, NASA experts have voted 2012 the most scientifically flawed and absurd science fiction film ever made. The 2009 disaster film named 2012 was directed by Roland Emmerich and written by Emmerich and Harald Kloser and grossed almost \$ 800 million.....

NASAs Donald Yeomans, who headed the Near Earth Asteroid Rendezvous mission, called the film an "exceptional and extraordinary" example of bad science in Hollywood movies. He pointed out that neutrino particles cannot interact with physical substances, and there is no possible way neutrinos carried to Earth by solar flares as depicted could cook the planets core and cause hurricanes or earthquakes or produce tsunamis big enough to overwhelm Mount Everest as shown in the film.



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CONCEPTION OF THE NEUTRINO



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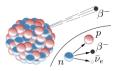
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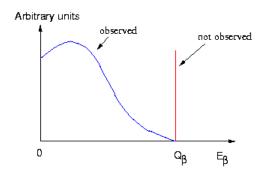
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Before 1930's: beta decay spectrum continuous - is this energy non-conservation?





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<u>Dec 1930:</u> Wolfgang Pauli's letter to physicists at a workshop in Tubingen:



Dear Radioactive Ladies and Gentlemen.

Wolfgang Pauli

......., I have hit upon a desparate remedy to save the "exchange theorem" of statistics and the law of conservation of energy. Namely, the possibility that there could exist in the nuclei electrically neutral particles, that I wish to call neutrons.... The mass of the neutrons should be of the same order of magnitude as the electron mass and in any event not larger than 0.01 proton masses. The continuous beta spectrum would then become understandable by the assumption that in beta decay a neutron is emitted in addition to the electron such that the sum of the energies of the neutron and the electron is constant........

Unfortunately, I cannot appear in Tubingen personally since I am indispensable here in Zurich because of a ball on the night of 6/7 December. With my best regards to you, and also to Mr Back. Your humble servant

. W. Pauli



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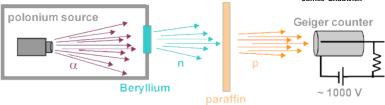
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1932: James Chadwick discovers the neutron, $\text{mass}_{\text{neutron}} = 1.0014 \times \text{mass}_{\text{proton}} \text{ - its too heavy - cant be Pauli's particle}$



James Chadwick





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Solvay Conference, Bruxelles 1933: Enrico Fermi proposes to name Pauli's particle the "neutrino".



Enrico Fermi



Particle physics units and symbols

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Symbols used for some common particles:

Symbol	Particle	
$\bar{\nu}, \bar{\nu}$	Neutrino and anti-neutrino	
γ	Photon	
\mathbf{e}^{-}	Electron	
$\mathbf{e^+}$	Anti-electron (positron)	
р	proton	
n	neutron	
N	nucleon - proton or neutron	



Mass is just a form of energy!

Particle physicists express masses in terms of energy, E = mc² Mass of proton = 1.67×10^{-24} g ≈ 1 billion (Giga) electron-volts (GeV) 1 thousand GeV = energy of a flying mosquito



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 $n \rightarrow p + e \rightarrow \bar{\nu}$

> 1933: Fermi builds his theory of weak interactions and beta decay **Neutral current Charged current interactions** interactions n or p interacts with Neutrino interacts neutrino or antineutrino Decay of neutron with neutron $vor \overline{v}$ n or p $vor \overline{v}$ n or p



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$$n \rightarrow p^+ + e^- + \bar{\nu}$$



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$$\mathbf{n} \rightarrow \mathbf{p}^+ + \mathbf{e}^- + \bar{\nu}$$

 $\mathbf{n} + \nu \rightarrow \mathbf{p}^+ + \mathbf{e}^-$



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Finding Neutrinos.... 1st attempt

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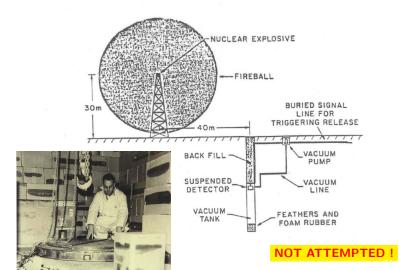
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1950's: Fredrick Reines, protege of Richard Feynman proposes to find neutrinos





Finding Neutrinos.... 2nd attempt

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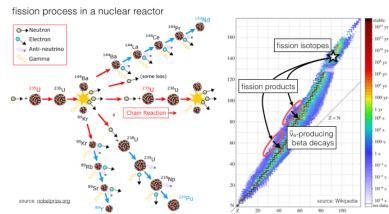
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Conclusion:

1950's: Fred Reines at Los Alamos and Clyde Cowan propose to use the Hanford nuclear reactor (1953) and the new Savannah River nuclear reactor (1955) to find neutrinos.





Finding Neutrinos.... 2nd attempt

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Conclusions

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THE UNIVERSITY OF CHICAGO
CHICAGO D' ILLINOIS
INSTITUTE FOR NUCLEAR STUDIES
October 8, 1952

Dr. Fred Reines Los Alamos Scientific Laboratory P.O. Box 1663 Los Alamos, New Mexico

Dear Fred

Thank you for your latter of October this by Clyde Gromn and yourself. I was very much interested in your mey plan for the detection of the neutrino. Certainly your new rathod should be much singler to carry out and have the great advantage that the measurement can be repeated any number of times. I shall be very interested in seeing hew your 10 cubic floot scintillation counter is going to work, but I do not know of any reason may it should not

Good luck.



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Finding Neutrinos.... 2nd attempt

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1950's: Fred Reines at Los Alamos and Clyde Cowan propose to use the Hanford nuclear reactor (1953) and the new Savannah River nuclear reactor (1955) to find neutrinos.

A detector filled with water with CdCl₂ in solution was located 11 meters from the reactor center and 12 meters underground.

The detection sequence was as follows:

- $1 \bar{\nu_e} + p \rightarrow n + e^+$
- 3 $n + ^{108} Cd \rightarrow ^{109} Cd* \rightarrow ^{109} Cd + \gamma$ $(\tau = 5\mu s).$



Neutrinos first detected using a nuclear reactor!

Reines shared 1995 Nobel for work on neutrino physics.





ν : A Truly Elusive Particle!

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Conclusion

Reines and Cowan were the first to estimate the interaction strength of neutrinos. The cross-section is $\sigma \sim 10^{-43} {\rm cm}^2$ per nucleon (N = n or p).

$$\nu$$
 mean free path = $\frac{1}{\sigma \times \text{number of nucleons per cm}^3}$

 ν Exercise: What is the mean free path of a neutrino in lead? (use Table of atomic and nuclear properties)

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$$= \frac{1}{10^{-43} \text{cm}^2 \times 11.4 \text{g/cm}^3 \times 6.02 \times 10^{23} \text{nucleons/g}}$$

$$\approx 1.5 \times 10^{16} \text{m}$$

How many light years is that? How does it compare to the distance from the sun to the moon?



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$$\approx 1.5 \times 10^{16} \text{m}$$

How many light years is that? How does it compare to the distance from the sun to the moon?

= 100,000 distance earth to sun

A proton has a mean free path of 10cm in lead



Reactor power and neutrinos

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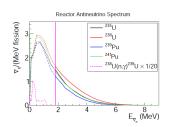
Conclusions

ν Exercise:

The following table shows the breakdown of energy released per fission from ²³⁵U:

Fission fragment	Energy (MeV)
Fission products	175
(2.44) neutrons	5
γ from fission	7
γ s and β s from beta decay	13
(6) neutrinos	10
Total	210

5% of a reactor's power is in neutrinos!



How many neutrinos are emitted per second from a 1 Gigawatt (thermal) reactor?

Reactor power and neutrinos

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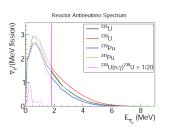
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How many neutrinos are emitted per second from a 1 Gigawatt (thermal) reactor?

$$1 \times 10^{9} \text{ Joules/sec} = 6.242 \times 10^{18} \text{ GeV/sec}$$

= $3 \times 10^{19} \text{ fissions/sec}$
 $\sim 2 \times 10^{20} \nu/\text{sec}$
= $1.6 \times 10^{13}/\text{m}^{2}/\text{sec}$ at 1 km



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ν Exercise:

Using the rate of neutrinos emitted from a reactor (= $2\times 10^{20}/\text{sec/GW})$ and the average cross-section of the inverse beta decay process ($\bar{\nu}_e+p\to e^++n$) is $\sigma=10^{-43}\mathrm{cm^2/proton}$, what is the rate of neutrino interactions per day in a detector containing 100 tons of scintillator (CH₂) located 1km from a 1GW reactor? Note that the IBD process only happens on free protons (H)



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interactions/day = flux $(\nu/\text{cm}^2/\text{day}) \times \sigma (\text{cm}^2/\text{p}) \times \text{protons/Nucleons} \times \text{Nucleons/gram} \times 10^8 \text{ g/100tons}$



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interactions/day = flux (
$$\nu/\text{cm}^2/\text{day}$$
) \times σ (cm²/p) \times protons/Nucleons \times Nucleons/gram \times 10⁸ g/100tons

$$\#$$
 interactions/day = 118

Precision ν expt: need 1 GW nuclear reactor (\$1B) + 100's tons



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FINDING NEUTRINOS in NATURE



Discovery of the Muon (μ)

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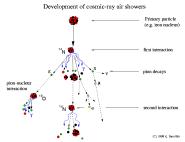
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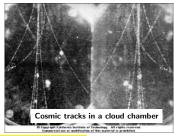
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Conclusions

1936: Carl Andersen, Seth Neddermeyer observed an unknown charged particle in cosmic rays with mass between that of the electron and the proton - called it the μ meson (now muons).







I. I Rabbi (founder of BNL): Who ordered THAT?



Discovery of the Pion: 1947

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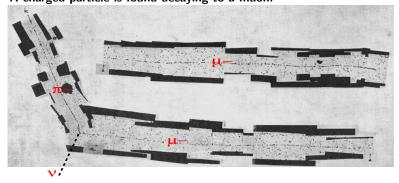
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Conclusion:

Cecil Powell takes emulsion photos aboard high altitude RAF flights. A charged particle is found decaying to a muon:



 $\mathsf{mass}_{\pi^-} = 0.1396~\mathsf{GeV/c^2}$, $\tau = 26$ nano-second (ns).

Pions are composite particles from the "hadron" family which includes protons and neutrons.



Proposal to find Atmospheric Neutrinos

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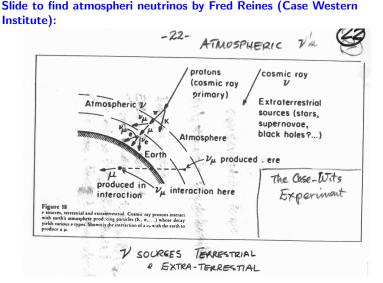
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The CWI-SAND Experiment

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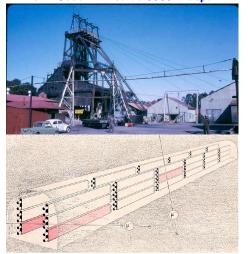
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1964: The Case Western Institute-South Africa Neutrino Detector (CWI-SAND) and a search for atmospheric ν_μ at the East Rand gold mine in South Africa at 3585m depth







The CWI-SAND Experiment

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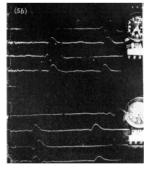
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Conclusions

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Downward-going Muon (background)

Horizontal Muon (neutrino signal)

Detection of the first neutrino in nature!



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FLAVORS OF NEUTRINOS



Producing Neutrinos from an Accelerator

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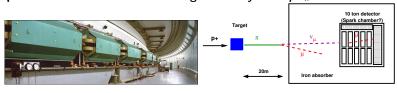
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Conclusions



1962: Leon Lederman, Melvin Schwartz and Jack Steinberger use a proton beam from BNL's Alternating Gradient Synchrotron (AGS) to produce a beam of neutrinos using the decay $\pi \to \mu \nu_x$



Making ν 's



The Two-Neutrino Experiment

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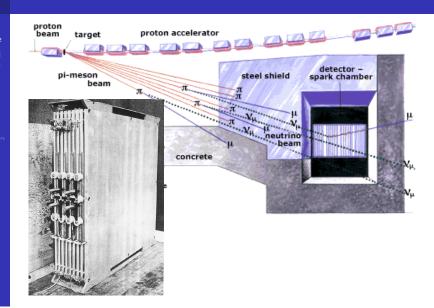
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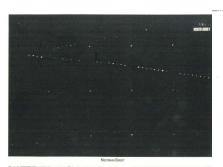
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Conclusion:







Classification of "Events"

Single Traces							
$p_{_{\rm H}}$ < 300 MeV/c ^h	49						
p _µ > 300	34						
> 400	19						
> 500	8						
> 600	3						
> 700	2						
Total "single Muor Events"	34						

Vertex Brents

Visible Energy Released < 1 SeV 15

Visible Energy Released > 1 SeV __7

Total vertex events 22

"Shower" Events

aergy	of	"electron"	-	300	±	100	MeV	3
				220				1
				240				3
				280				1
Tot	al	"stower ev	en:	ts"b				6

- These are not included in the "event" count.
- The two shower events which are so located that their potential energy release in the chamber corresponds to smoots of less than 300 MeV/c are not included here.

The first event!



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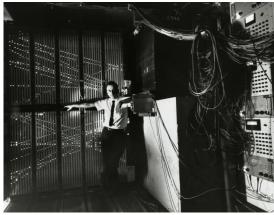
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μ Δnns

Conclusions



Result: 40 neutrino interactions recorded in the detector, 6 of the resultant particles where identified as background and 34 identified as

$$\mu \Rightarrow \nu_{x} = \nu_{\mu}$$

The first successful accelerator neutrino experiment was at Brookhaven Lab.

X is other particles

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Conclusion

To produce neutrinos from accelerators

 ${\sf p}^+ + {\sf A} \to \pi^\pm + {\sf X}, \quad \pi^\pm \to \mu^\pm + \nu_\mu/\bar{\nu}_\mu$ where A = Carbon (Graphite), Berillyium, Tungsten,

u Exercise: The Main Injector accelerator at Fermilab delivers $4.86 \times 10^{13}~120~\text{GeV}$ protons in a 10 microsecond pulse every 1.33 seconds to the NuMI neutrino beamline target. What is the average power of the proton beam delivered in megawatts?

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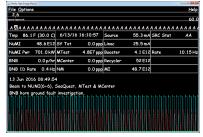
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To produce neutrinos from accelerators

$$p^+ + A \rightarrow \pi^{\pm} + X$$
, $\pi^{\pm} \rightarrow \mu^{\pm} + \nu_{\mu}/\bar{\nu}_{\mu}$ where $A =$ Carbon (Graphite), Berillyium, Tungsten, X is other particles

u Exercise: The Main Injector accelerator at Fermilab delivers 4.86×10^{13} 120 GeV protons in a 10 microsecond pulse every 1.33 seconds to the NuMI neutrino beamline target. What is the average power of the proton beam delivered in megawatts?

Power = 120 GeV \times 4.86 10^{13} protons \times 1.6 10^{-10} Joules/GeV \times 1/1.33s = 702 kW





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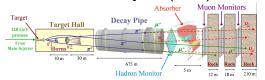
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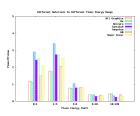
Ci Violation

Conclusions

Neutrinos at the Main Injector (NuMI) Beamline



The result of a FLUKA simulation of pion production from 120 GeV protons is shown on the right:



(work by Yi LU, a HS student intern)

u Exercise: What fraction of 6 GeV pions from the NuMI target will decay to neutrinos before reaching the end of the evacuated NuMI decay pipe of 675m long? The π^+ rest mass and lifetime are 140 MeV and 26 ns

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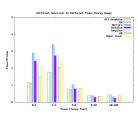
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6 GeV
$$\pi^+$$
 lifetime: $\tau=\gamma\tau_0=\frac{\rm E}{\rm m_0c^2}\times 26\rm ns=1.1\mu s,\ c\tau=334\ m$

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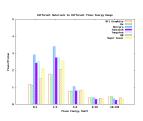
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Conclusions

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6 GeV
$$\pi^+$$
 lifetime: $\tau = \gamma \tau_0 = \frac{E}{m_0 c^2} \times 26 \text{ns} = 1.1 \mu \text{s}, \ c\tau = 334 \text{ m}$

$$\mathsf{F}_{\mathsf{decays}} = (1 - \exp^{-1/c\tau}) \mathcal{F}(\pi \to \mu \nu_\mu) = (0.87)(0.99) = 0.86$$



The Lepton Family and Flavors

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Conclusion:

The muon and the electron are different "flavors" of the same family of elementary particles called leptons.

Generation	l l	II	111
Lepton	e ⁻	$oldsymbol{\mu}$	$oldsymbol{ au}$
Mass (GeV)	0.000511	0.1057	1.78
Lifetime (sec)	stable	2.2×10^{-6}	2.9×10^{-13}

Neutrinos are neutral leptons.



Number of Neutrino Flavors: Particle Colliders

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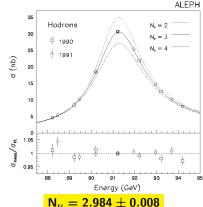
Conclusion

<u>1980's - 90's:</u> The number of neutrino types is precisely determined from studies of Z^0 boson properties produced in e^+e^- colliders.

The LEP e⁺e⁻ collider at CERN, Switzerland



The 27km LEP ring was reused to build the Large Hadron Collider





The Particle Zoo

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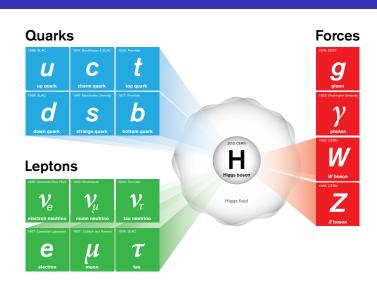
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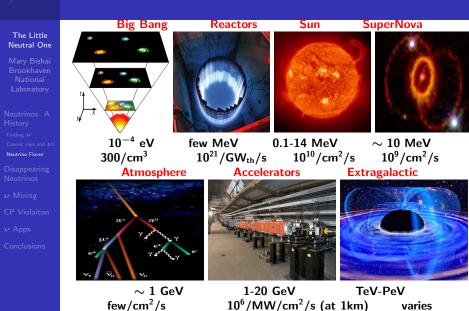
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Conclusion:





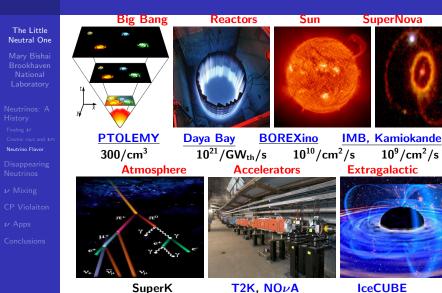
Neutrino Sources





Neutrino Experiments

few/cm²/s



 $\frac{\text{T2K}}{10^6/\text{MW/cm}^2/\text{s}}$ (at 1km) varies



Neutrinos and Todays Universe

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NEUTRINO MIXING AND OSCILLATIONS



Solar Neutrinos

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u}$ Neutrino Flavor

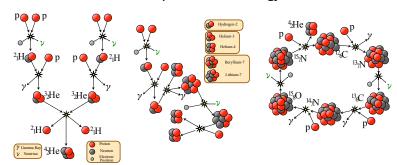
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Fusion of nuclei in the Sun produces solar energy and neutrinos



The Homestake Experiment

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1967: Ray Davis from BNL installs a large detector, containing 615 tons of tetrachloroethylene (cleaning fluid), 1.6km underground in Homestake mine, SD.

1
$$\nu_{\rm e}^{\rm sun} + ^{37}{\rm CL} \rightarrow {\rm e}^- + ^{37}{\rm Ar}, \ \tau(^{37}{\rm Ar}) = 35 {\rm \ days}.$$

2 Number of Ar atoms \approx number of $\nu_{\rm e}^{\rm sun}$ interactions.



Ray Davis



Results: 1969 - 1993 Measured 2.5 \pm 0.2 SNU (1 SNU = 1 neutrino interaction per second for 10^{36} target atoms) while theory predicts 8 SNU. This is a ν_e^{sun} deficit of 69%.

Where did the suns ν_e 's go?



2002 Nobel Prize

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Ray Davis Brookhaven Lab, USA (Homestake experiment)



Masatoshi Koshiba University of Tokyo, USA (Kamiokande experiment)

The Nobel Prize in Physics 2002 was awarded 1/4 to Ray Davis and 1/4 Masatoshi Koshiba "for pioneering contributions to astrophysics, in particular for the detection of cosmic neutrinos."



The Super-Kamiokande Experiment. Kamioka Mine, Japan

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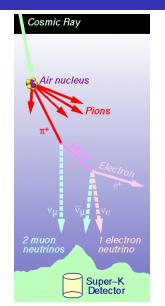
Disappearing Neutrinos

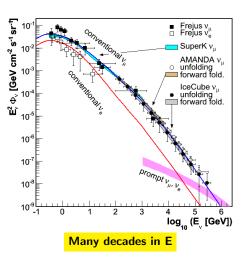
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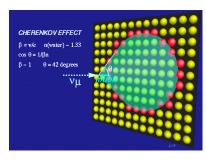
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Conclusior



50kT double layered tank of ultra pure water surrounded by 11,146 20" diameter photomultiplier tubes.

Neutrinos are identified by using CC interaction $\nu_{\mu,\mathrm{e}} \to \mathrm{e}^{\pm}, \mu^{\pm} \mathrm{X}.$ The lepton produces Cherenkov light as it goes through the detector:





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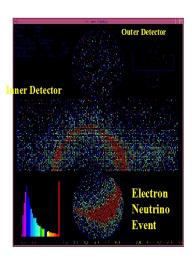
Disappearing Neutrinos

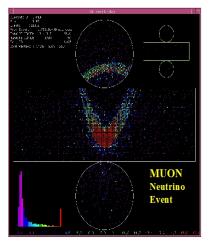
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More Disappearing Neutrinos!!

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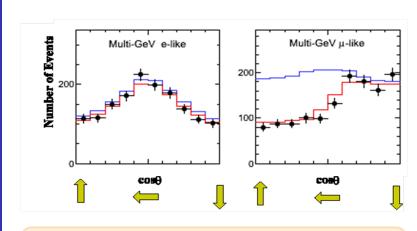
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Conclusions



All the $\nu_{\rm e}$ are there! But what happened to the ν_{μ} ??

SNO Experiment: Solar ν Measurments $1 \leftrightarrow 2 \text{ mix ing}$

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Conclusions

2001-02: Sudbury Neutrino Observatory. Water Čerenkov detector with 1 kT heavy water (0.5 B\$ worth on loan from Atomic Energy of Canada Ltd.) located 2Km below ground in INCO's Creighton nickel mine near Sudbury, Ontario. Can detect the following $\nu^{\rm sun}$ interactions:

- 1) $\nu_{\rm e} + {\rm d} \rightarrow {\rm e}^- + {\rm p} + {\rm p}$ (CC).
- 2) $\nu_{\rm e,x}+{\rm e}^-\to{\rm e}^-+\nu_{\rm x},~\sigma^{ES}_{\rm nu_e}:\sigma^{ES}_{\nu_{\rm x}}=6:1$ (ES).
- 3) $\nu_{\rm x}+{\rm d} \rightarrow {\rm p}+{\rm n}+\nu_{\rm x}, \ {\rm x}={\rm e},\mu,\tau$ (NC).



SNO measured:

$$\overline{\phi_{\text{SNO}}^{\text{CC}}(\nu_{\text{e}}) = 1.75 \pm 0.07 (\text{stat})_{-0.11}^{+0.12} (\text{sys.}) \pm 0.05 (\text{theor}) \times 10^6 \text{cm}^{-2} \text{s}^{-1}}$$

$$\phi_{\mathsf{SNO}}^{\mathsf{ES}}(
u_{\mathsf{x}}) = 2.39 \pm 0.34 (\mathrm{stat})_{-0.14}^{+0.16} (\mathrm{sys.}) \pm \times 10^6 \mathrm{cm}^{-2} \mathrm{s}^{-1}$$

$$\phi_{\mathsf{SNO}}^{\mathsf{NC}}(\nu_{\mathsf{x}}) = 5.09 \pm 0.44 (\mathrm{stat})_{-0.43}^{+0.46} (\mathrm{sys.}) \pm \times 10^6 \mathrm{cm}^{-2} \mathrm{s}^{-1}$$

All the solar ν 's are there but ν_e appears as ν_x !



Some Quantum Mechanics

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Conclusion

1924: Louis-Victor-Pierre-Raymond, 7th duc de Broglie proposes in his doctoral thesis that all matter has wave-like and particle-like properties.

For highly relativistic particles : energy \approx momentum



De Broglie

Wavelength (nm)
$$\approx \frac{1.24 \times 10^{-6} \text{ GeV.nm}}{\text{Energy (GeV)}}$$



Neutrino Mixing

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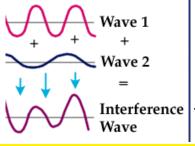
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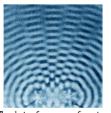
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ν App

Conclusions

1957,1967: B. Pontecorvo proposes that neutrinos of a particular flavor are a mix of quantum states with different masses that propagate with different phases:





The interference of water waves coming from two sources.

The inteference pattern depends on the difference in masses

Neutrino Mixing ⇒ Oscillations

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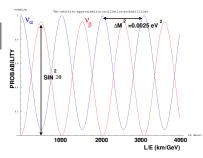
$$\left(\begin{array}{c} \mathbf{\nu_a} \\ \mathbf{\nu_b} \end{array}\right) = \left(\begin{array}{cc} \cos(\theta) & \sin(\theta) \\ -\sin(\theta) & \cos(\theta) \end{array}\right) \left(\begin{array}{c} \nu_1 \\ \nu_2 \end{array}\right)$$

$$\begin{split} \nu_a(t) &= \cos(\theta)\nu_1(t) + \sin(\theta)\nu_2(t) \\ P(\nu_a \to \nu_b) &= |<\nu_b|\nu_a(t)>|^2 \\ &= \sin^2(\theta)\cos^2(\theta)|e^{-iE_2t} - e^{-iE_1t}|^2 \end{split}$$

$$\mathsf{P}(\textcolor{red}{\nu_{\mathsf{a}}} \rightarrow \textcolor{red}{\nu_{\mathsf{b}}}) = \sin^2 2\theta \sin^2 \frac{1.27\Delta \mathsf{m}_{21}^2 \mathsf{L}}{\mathsf{E}}$$

where $\Delta m_{21}^2 = (m_2^2 - m_1^2)$ in eV², L (km) and E (GeV).

Observation of oscillations implies non-zero mass eigenstates



Two Different Mass Scales!



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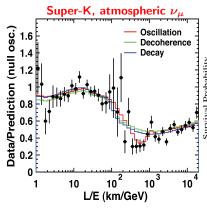
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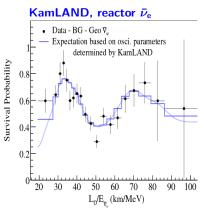
Conclusions



Global fit 2013:

$$\Delta m^2_{
m atm} = 2.43^{+0.06}_{-0.10} imes 10^{-3} \
m eV^2 \
m sin^2 \ heta_{
m atm} = 0.386^{+0.24}_{-0.21}$$

Atmospheric L/E ~ 500 km/GeV



Global fit 2013:

$$\Delta m_{
m solar}^2 = 7.54_{-0.22}^{+0.26} \times 10^{-5} \text{ eV}^2
\sin^2 \theta_{
m solar} = 0.307_{-0.16}^{+0.18}$$

Solar L/E \sim 15,000 km/GeV



2015 Nobel Prize

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Takaaki Kajita University of Tokyo, Japan (SuperKamiokande)



Arthur B. MacDonald Queens University, Canada (SNO)

The Nobel Prize in Physics 2015 was awarded jointly to Takaaki Kajita and Arthur B. McDonald "for the discovery of neutrino oscillations, which shows that neutrinos have mass"



Neutrino Mixing: 3 flavors, 3 amplitudes, 2 mass scales

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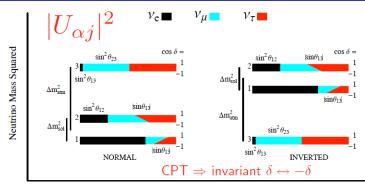
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Fractional Flavor Content varying $\cos \delta$

The "mixing angles" $(\theta_{13},\theta_{12},\theta_{23})$ represent the fraction of ν_e,ν_μ in the 3 mass states. They determine the probability of oscillation from one flavor to the other

 $\sin^2 \theta_{12} pprox \sin^2 \theta_{
m solar}$, $\sin^2 \theta_{23} pprox \sin^2 \theta_{
m atmospheric}$

3 quantum states interfering \Rightarrow phase δ



CP Violation in PMNS (leptons) and CKM (quarks)

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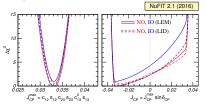
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In 3-flavor mixing the degree of CP violation is determined by the Jarlskog invariant:

$$\mathsf{J}_{\mathsf{CP}}^{\mathsf{PMNS}} \equiv rac{1}{8} \sin 2 heta_{12} \sin 2 heta_{13} \sin 2 heta_{23} \cos heta_{13} \sin \delta_{\mathsf{CP}}.$$



(JHEP 11 (2014) 052, arXiv:1409.5439)

Given the current best-fit values of the ν mixing angles :

$$J_{\text{CP}}^{\mathrm{PMNS}} pprox 3 imes 10^{-2} \sin \delta_{\mathrm{CP}}.$$

For CKM (mixing among the 3 quark generations):

$$J_{CP}^{\rm CKM} \approx 3 \times 10^{-5},$$

despite the large value of $\delta_{CP}^{CKM} \approx 70^{\circ}$.

$u_{\mu} ightarrow u_{ m e}$ Oscillations in the 3-flavor u SM

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Conclusions

In the ν 3-flavor model matter/anti-matter asymmetries in neutrinos are best probed using $\nu_{\mu}/\bar{\nu}_{\mu} \to \nu_{\rm e}/\bar{\nu}_{\rm e}$ oscillations (or vice versa). With terms up to second order in $\alpha \equiv \Delta m_{21}^2/\Delta m_{31}^2 = 0.03$ and $\sin^2\theta_{13} = 0.02$, (M. Freund. Phys. Rev. D 64, 053003):

$$P(\nu_{\mu} \rightarrow \nu_{e}) \cong P(\nu_{e} \rightarrow \nu_{\mu}) \cong \underbrace{P_{0}}_{\theta_{13}} + \underbrace{P_{\sin\delta}}_{\text{Violating}} + \underbrace{P_{\cos\delta}}_{\text{CP conserving}} + \underbrace{P_{3}}_{\text{solar oscillation}}$$

where for oscillations in vacuum:

$$P_0 = \sin^2 \theta_{23} \sin^2 2\theta_{13} \sin^2(\Delta),$$

$$P_{\sin \delta} = \alpha 8 J_{cp} \sin^3(\Delta),$$

$$P_{\cos \delta} = \alpha 8 J_{cp} \cot \delta_{CP} \cos \Delta \sin^2(\Delta),$$

$$P_3 = \alpha^2 \cos^2 \theta_{23} \sin^2 2\theta_{12} \sin^2(\Delta),$$

where $\Delta = 1.27 \Delta m_{31}^2 (eV^2) L(km) / E(GeV)$

For
$$ar{
u}_{\mu}
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u}_{ ext{e}}, \underbrace{\mathsf{P}_{\sin\delta}
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$$P(\nu_{\mu} \to \nu_{e}) \cong P(\nu_{e} \to \nu_{\mu}) \cong \underbrace{P_{0}}_{\theta_{13}} + \underbrace{P_{\sin\delta}}_{\text{CP violating}} + \underbrace{P_{\cos\delta}}_{\text{conserving solar oscillation}} + \underbrace{P_{3}}_{\text{Solar oscillation}}$$

where for oscillations in matter with constant density:

$$\begin{split} \mathsf{P}_0 &= \sin^2\theta_{23} \frac{\sin^22\theta_{13}}{(\mathsf{A}-1)^2} \sin^2[(\mathsf{A}-1)\Delta], \\ \mathsf{P}_{\sin\delta} &= \alpha \frac{8\mathsf{J}_{\mathrm{cp}}}{\mathsf{A}(1-\mathsf{A})} \sin\Delta\sin(\mathsf{A}\Delta) \sin[(1-\mathsf{A})\Delta], \\ \mathsf{P}_{\cos\delta} &= \alpha \frac{8\mathsf{J}_{\mathrm{cp}} \cot\delta_{\mathrm{CP}}}{\mathsf{A}(1-\mathsf{A})} \cos\Delta\sin(\mathsf{A}\Delta) \sin[(1-\mathsf{A})\Delta], \\ \mathsf{P}_3 &= \alpha^2 \cos^2\theta_{23} \frac{\sin^22\theta_{12}}{\mathsf{A}^2} \sin^2(\mathsf{A}\Delta), \end{split}$$

where $\Delta=1.27\Delta m_{31}^2 (eV^2) L(km)/E(GeV)~$ and $A=\sqrt{2}G_FN_e2E/\Delta m_{31}^2.$

For
$$ar{
u}_{\mu}
ightarrow ar{
u}_{e}$$
, $egin{array}{c} P_{\sin\delta}
ightarrow - P_{\sin\delta} \\ \hline ext{CP asymmetry} \end{array}$ matter asymmetry

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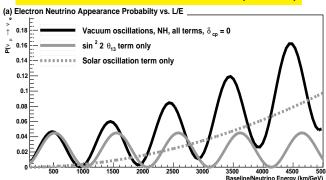
Conclusion:

 ν Exercise: Use your favorite plotting package and reproduce the plots shown below

The $\nu_{\mu}
ightarrow
u_{\rm e}$ oscillation probability maxima occur at

$$\frac{L~(\mathrm{km})}{E_n(\mathrm{GeV})} = \left(\frac{\pi}{2}\right) \frac{(2n-1)}{1.27 \times \Delta m_{31}^2 (\mathrm{eV}^2)} \approx (2n-1) \times \frac{515~\mathrm{km}}{\mathrm{GeV}}$$

Oscillations in vacuum - different terms ($\delta_{\rm CP}=0$)



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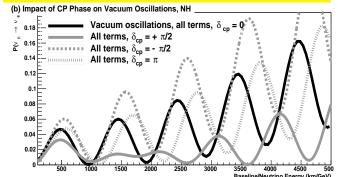
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Impact of δ_{CP} on oscillations in vacuum, $\Delta m_{31}^2 > 0$ (NH)



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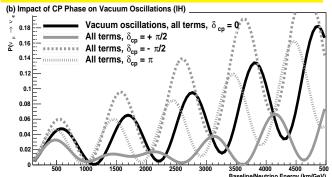
Conclusions

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The $u_{\mu}
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$$\frac{L~(\mathrm{km})}{E_n(\mathrm{GeV})} = \left(\frac{\pi}{2}\right) \frac{(2n-1)}{1.27 \times \Delta m_{31}^2 (\mathrm{eV}^2)} \approx (2n-1) \times \frac{515~\mathrm{km}}{\mathrm{GeV}}$$

Impact of δ_{CP} on oscillations in vacuum, $\Delta m_{31}^2 < 0$ (IH)



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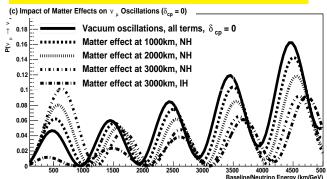
Conclusions

 ν Exercise: Use your favorite plotting package and reproduce the plots shown below

The $u_{\mu} \rightarrow
u_{e}$ oscillation probability maxima occur at

$$\frac{L~(\mathrm{km})}{E_n(\mathrm{GeV})} = \left(\frac{\pi}{2}\right) \frac{(2n-1)}{1.27 \times \Delta m_{21}^2 (\mathrm{eV}^2)} \approx (2n-1) \times \frac{515~\mathrm{km}}{\mathrm{GeV}}$$

Impact of matter effect on ν_{μ} oscillations ($\delta_{\rm CP}=0$)



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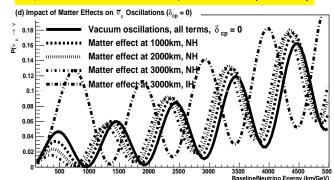
Conclusions

u Exercise: Use your favorite plotting package and reproduce the plots shown below

The $u_{\mu} \rightarrow
u_{\rm e}$ oscillation probability maxima occur at

$$\frac{L~(\mathrm{km})}{E_n(\mathrm{GeV})} = \left(\frac{\pi}{2}\right) \frac{(2n-1)}{1.27 \times \Delta m_{31}^2 (\mathrm{eV}^2)} \approx (2n-1) \times \frac{515~\mathrm{km}}{\mathrm{GeV}}$$

Impact of matter effect on $\bar{\nu}_{\mu}$ oscillations ($\delta_{\rm CP}=0$)



Expected Appearance Signal Event Rates

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 ν Exercise: The total number of electron neutrino appearance events expected for a given exposure from a muon neutrino source as a function of baseline is given as

$$\mathsf{N}_{\nu_{\mathrm{e}}}^{\mathrm{appear}}(\mathsf{L}) = \int \Phi^{\nu_{\mu}}(\mathsf{E}_{\nu},\mathsf{L}) \times \mathsf{P}^{\nu_{\mu} \to \nu_{\mathrm{e}}}(\mathsf{E}_{\nu},\mathsf{L}) \times \sigma^{\nu_{\mathrm{e}}}(\mathsf{E}_{\nu}) \mathsf{d}\mathsf{E}_{\nu}$$

Assume the neutrino source produces a flux that is constant in energy and using only the dominant term in the probability(no matter effect)

$$\begin{array}{lcl} \Phi^{\nu_{\mu}}(\mathsf{E}_{\nu},\mathsf{L}) & \approx & \dfrac{\mathsf{C}}{\mathsf{L}^2}, \quad \mathsf{C} = \mathrm{number\ of}\ \nu_{\mu}/\mathrm{m}^2/\mathrm{GeV/sec\ at}\ 1\ \mathrm{km} \\ \\ \mathsf{P}^{\nu_{\mu}\to\nu_{e}}(\mathsf{E}_{\nu},\mathsf{L}) & \approx & \underbrace{\sin^2\theta_{23}\sin^22\theta_{13}\sin^2(1.27\Delta m_{31}^2\mathsf{L}/\mathsf{E}_{\nu})}_{\mathsf{P}_{0}} \\ \\ \sigma^{\nu_{e}}(\mathsf{E}_{\nu}) & = & 0.7\times10^{-42}(\mathrm{m}^2/\mathrm{GeV/N})\times\mathsf{E}_{\nu}, \quad \mathsf{E}_{\nu} > 1\ \mathrm{GeV} \end{array}$$

Prove that the rate of ν_e appearing integrated over a constant range of L/E is independent of baseline!



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$$\begin{split} N_{\nu_e}^{\rm appear}(L) & \propto {\rm constant~term} \times \int \frac{\sin^2(ax)}{x^3} dx, \\ x & \equiv L/E_{\nu},~a \equiv 1.27 \Delta m_{31}^2~{\rm GeV/(eV^2.km)} \end{split}$$

 ν Exercise:

 $\text{C}\approx 1\times 10^{17}~\nu_{\mu}/\mathrm{m}^2/\mathrm{GeV/yr}$ at 1 km (from 1MW accelerator) $\sin^22\theta_{13}=0.084,~\sin^2\theta_{23}=0.5, \Delta m_{31}^2=2.4\times 10^{-3}\mathrm{eV}^2$

Calculate the rate of ν_e events observed per kton of detector integrating over the region x = 100 km/GeV to 2000 km/GeV. Use ROOT to do the integral!



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Calculate the rate of ν_e events observed per kton of detector integrating over the region x = 100 km/GeV to 2000 km/GeV. Use ROOT to do the integral!

$$N_{\nu_{\rm e}}^{\rm appear}(L) \approx (2\times 10^6 {\rm events/kton/yr}) \cdot ({\rm km/GeV})^2 \int_{x_0}^{x_1} \frac{\sin^2(ax)}{x^3} dx,$$

 $N_{
u_e}^{
m appear}(L) \sim \mathcal{O}(20-30) \ {
m events/kton/yr}$



The Deep Underground Neutrino Experiment (DUNE) Coming online ~ 2028-2030

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Neutrinos: History

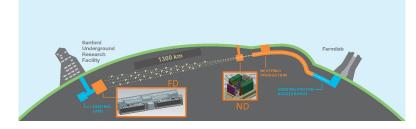
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- A very long baseline experiment: 1300km from Fermilab in Batavia, IL to the Sanford Underground Research Facility (former Homestake Mine) in Lead, SD.
- A highly capable near detector at Fermilab.
- A very deep (1 mile underground) far detector: massive 40-kton Liquid Argon Time-Projection-Chamber with state-of-the-art instrumentation.
- High intensity tunable wide-band neutrino beam from LBNF produced from upgraded MW-class proton accelerator at Fermilab.



The DUNE Scientific Collaboration

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As of Jan 2018:

60 % non-US

1061 collaborators from 175 institutions in 31 nations

Armenia, Brazil, Bulgaria, Canada, CERN, Chile, China, Colombia, Czech Republic, Finland, France, Greece, India, Iran, Italy, Japan, Madagascar, Mexico, Netherlands, Paraguay, Peru, Poland, Romania, Russia, South Korea, Spain, Sweden, Switzerland, Turkey, UK, Ukraine, USA





Scientific Objectives of DUNE

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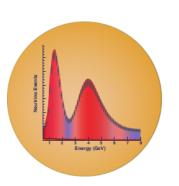
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- I precision measurements of the parameters that govern $\nu_{\mu} \rightarrow \nu_{e}$ oscillations; this includes precision measurement of the third mixing angle θ_{13} , measurement of the charge-parity (CP) violating phase δ_{CP} , and determination of the neutrino mass ordering (the sign of $\Delta m_{31}^2 = m_3^2 m_1^2$), the so-called mass hierarchy
- 2 precision measurements of the mixing angle θ_{23} , including the determination of the octant in which this angle lies, and the value of the mass difference, $-\Delta m_{32}^2$, in $\nu_{\mu} \rightarrow \nu_{e,\mu}$ oscillations



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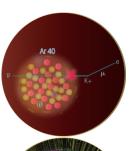
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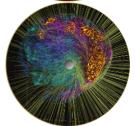
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- 3 search for proton decay, yielding significant improvement in the current limits on the partial lifetime of the proton (τ/BR) in one or more important candidate decay modes, e.g., $p \to K^+ \overline{\nu}$
- 4 detection and measurement of the neutrino flux from a core-collapse supernova within our galaxy, should one occur during the lifetime of DUNE



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PRACTICAL APPLICATIONS of u



Practical Applications of Technologies for u Experiments

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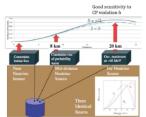
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Synergies and Applications - Examples

Cyclotrons for neutrino physics (and industrial applications)





Neutrino detectors for reactor monitoring and non-proliferation





remote discovery of undeclared nuclear reactors with large detectors at km scale



US Short-Baseline Experiment

reactor antineutrino studies at short baselines



Multi-MW Accelerators Driving Thorium Reactors

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First proposed by Carlo Rubbia in 1995 (1984 Nobel Prize winner)



Global energy resources in ZetaJoules

Resource	Туре	Yearly consumption (1999) ZJ	Resources ZI	Consumed until
Oil	Conventional	0.13	12.08	4.85
	Unconventional	0.01	20.35	0.29
	Total oil	0.14	32.42	5.14
Natural gas	Conventional	0.08	16.56	2.35
	Unconventional	0.00	33.23	0.03
	Total gas	0.08	49.79	2.38
Coal	Total coal	0.09	199.67	5.99
Total Fossils		0.31	281.88	13.51
Uranium	Thermal reactors	0.04	5.41 (2'000, sw)	
	Breeder	0	324 (120'000, sw)	
Thorium			1'300'000	

sw: including sea water
1 ZI (Zetaloule)= 103 EI(Exaloule)= 1021 I(Joule)

Requires proton accelerators with powers of 10 MW. Currently neutrino and neutron experiments are driving the technology of high power MW class proton beams.

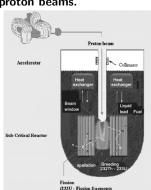


Figure 1. Schematic representation of Energy Amplifier proposed



Neutrinos and Earth's Geology

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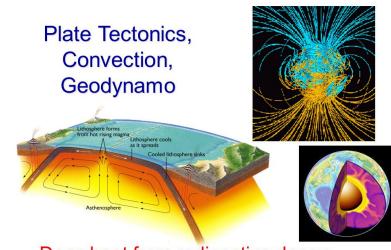
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Does heat from radioactive decay drive the Earth's engine?



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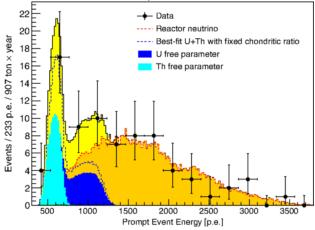
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Signal of $\bar{\nu_e}$ from radioactive decays of U/TH in the earth observed in the BOREXINO solar neutrino experiment:





Summary and Further Study

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- Neutrinos have been at the forefront of fundamental discoveries in particle physics for decades.
- Discoveries of neutrino properties like the very small mass, large almost maximal mixing, are the ONLY direct evidence for physics beyond the Standard Model of particle physics, and new hidden symmetries.
- The future T2HK and LBNF/DUNE projects are ambitious multi-billion dollar, multi-national neutrino experiments designed to probe matter/anti-matter asymmetries, neutrino oscillations and cosmological neutrinos with unprecedented precision.
- Studying neutrinos is advancing new technologies in accelerators, non-proliferation, geology...etc



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