



Northeastern  
University

STXS MEETING

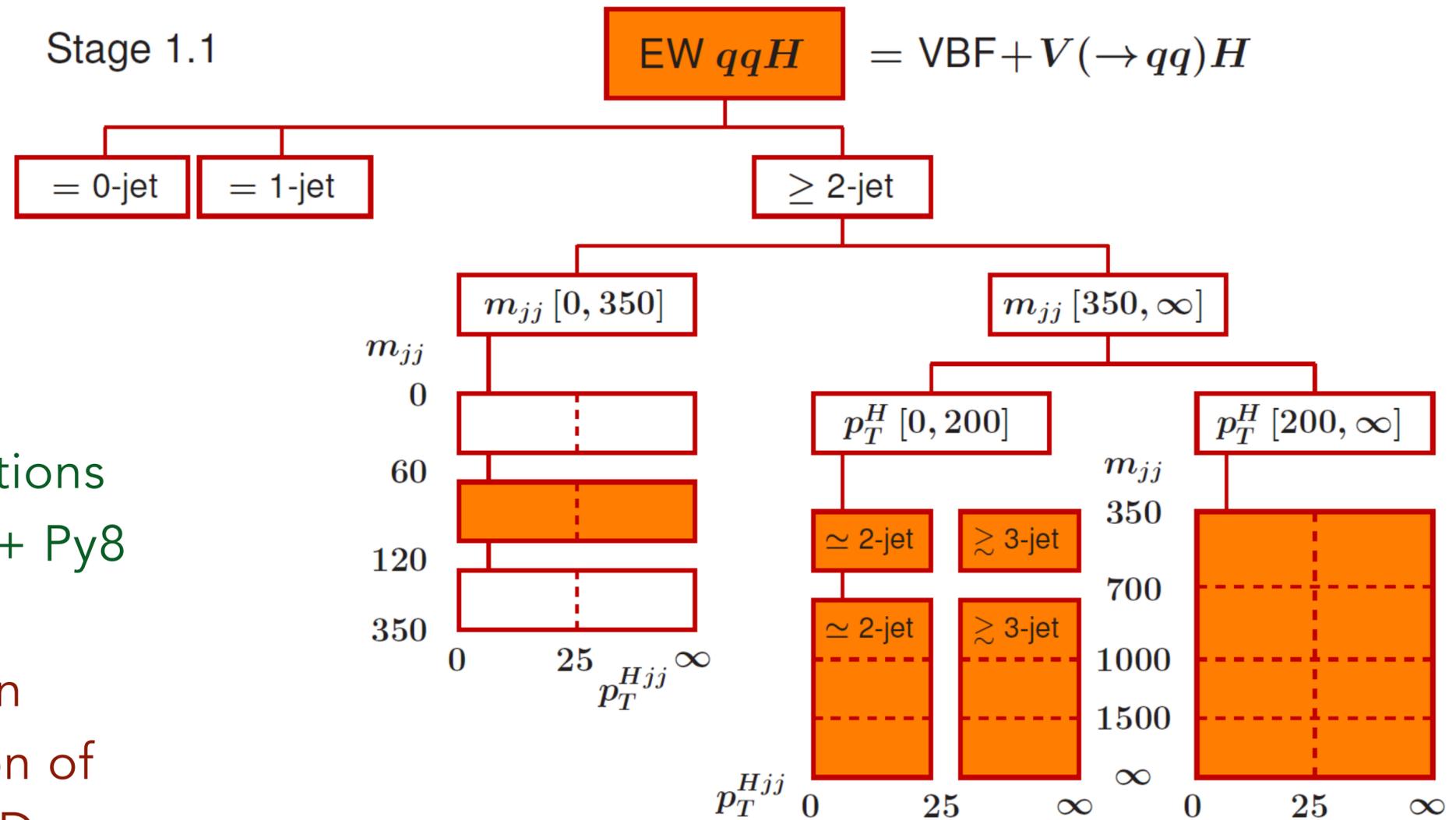
# EWK $H+2j$ STXS UNCERTAINTIES

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# INTRODUCTION

- **10 QCD-nuisances accounted:**
  - **1 yields uncertainty ("overall" NP) on the inclusive cross-section,  $\Delta_{\text{tot}}$**
  - **9 migration uncertainties**
    - $\Delta_{2\text{jet}}, \Delta_{200}$  and  $\Delta_{25}$
    - 6 NPs to describe  $M_{jj}$  spectrum
      - Estimated using ST method
- **So far we had:**
  - Uncertainties extracted using FO calculations
  - Acceptances estimated using POWHEG + Py8
- **Inclusion of VH hadronic** NEW
  - Checking validity of the VBF approximation
  - Updating uncertainties using full calculation of H+2jet and H+3jet production at NLO QCD
- **Electroweak Corrections** NEW
  - **EWK @NLO** correction applied for every STXS bin



Previous update: <https://indico.cern.ch/event/826136/contributions/3560473/attachments/1927391/3191007/STXS-uncertainties-VBF.pdf>

# LIMITS OF VBF APPROXIMATION

- Most the VBF generators in the market run with the VBF approximate (only t- and u-channel).
- More accurate EW Higgs + 2 jets requires the inclusion of the s-channel component
- Studies already initiated

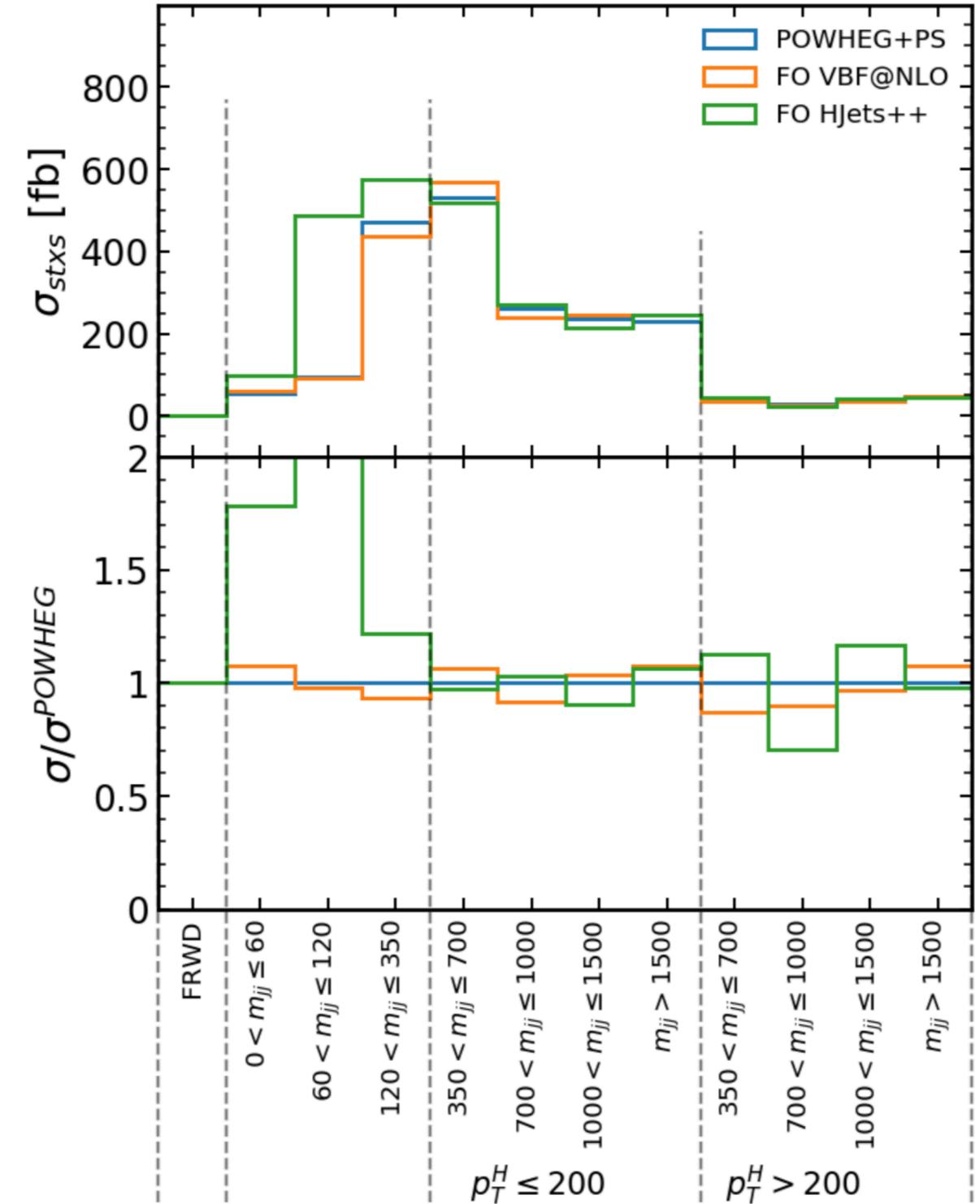
[Campanario, Figy, Plätzer, Sjö Dahl – PRL 111 (2013) 211802]

[Campanario, Figy, Plätzer, Rauch, Schichtel, Sjö Dahl – PRD 98 (2018) 033]

- HJets++ provides full EW H+2jet and H+3jet calculation at NLO QCD (VBF+VHhad)

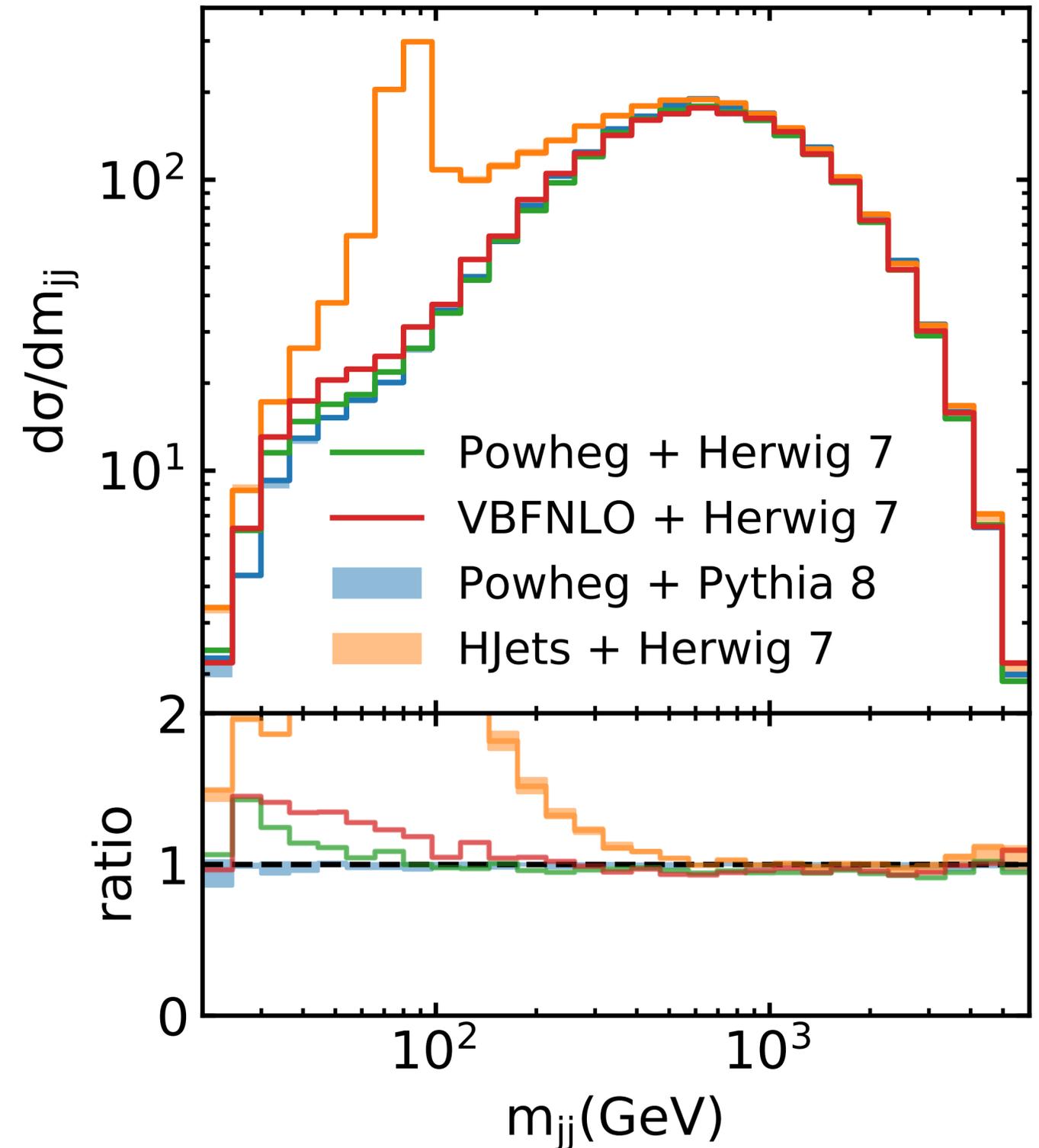
[Campanario, Figy, Plätzer, Sjö Dahl – PRL 111 (2013) 211802]

- Little impact for VBF selection, more significant changes  $m_{jj} < 350$  GeV

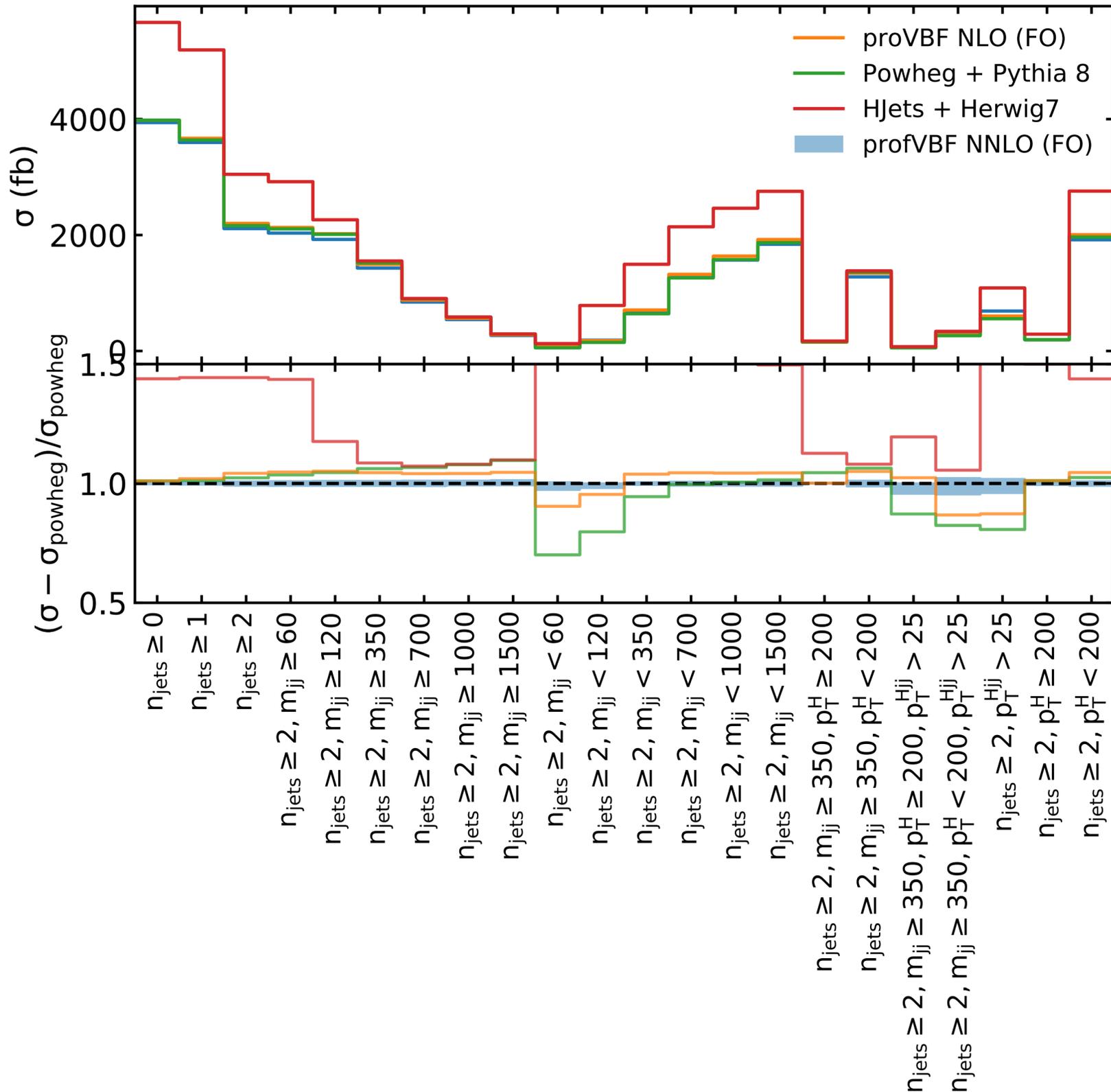


# H+2J EWK: ADDING THE S-CHANNEL

- Compared various combinations of ME+PS:
  - POWHEG: NLO-QCD (3rd jet LO from PS)
    - VBF approximated: only t/u-channels
    - Interfaced with Herwig7 and Pythia8 (with `dipoleRecoil=on` )
  - VBFNLO : NLO-QCD (3rd jet LO from PS)
    - VBF approximated: only t/u-channels
    - Interfaced with Herwig7
  - HJets++ : NLO-QCD (3rd jet LO from PS)
    - Full EWK H+2 jets calculation
    - Interfaced with Herwig7
- **Stage 1.1 acceptances need to be updated to account for the s-channel contribution using HJet++**

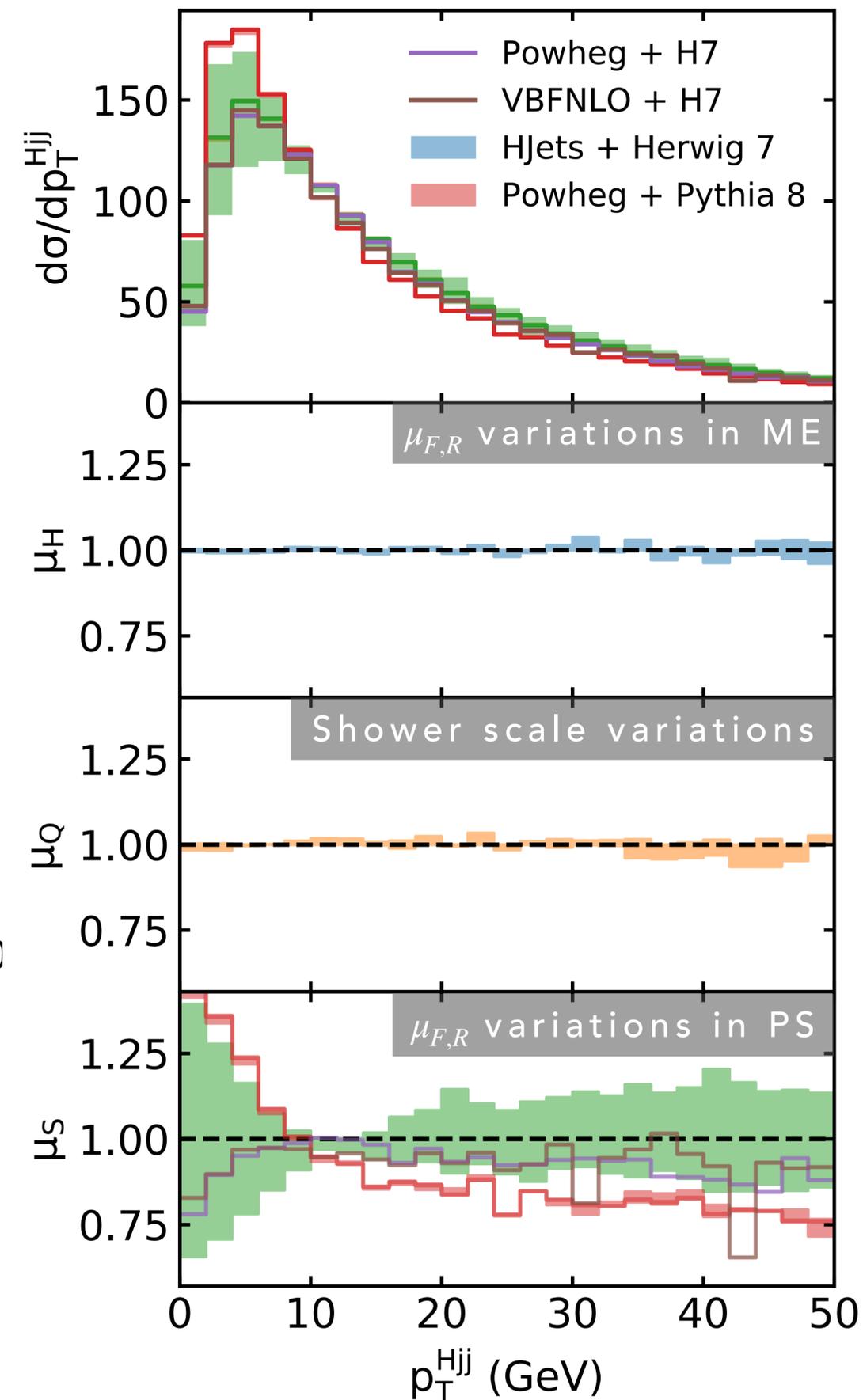
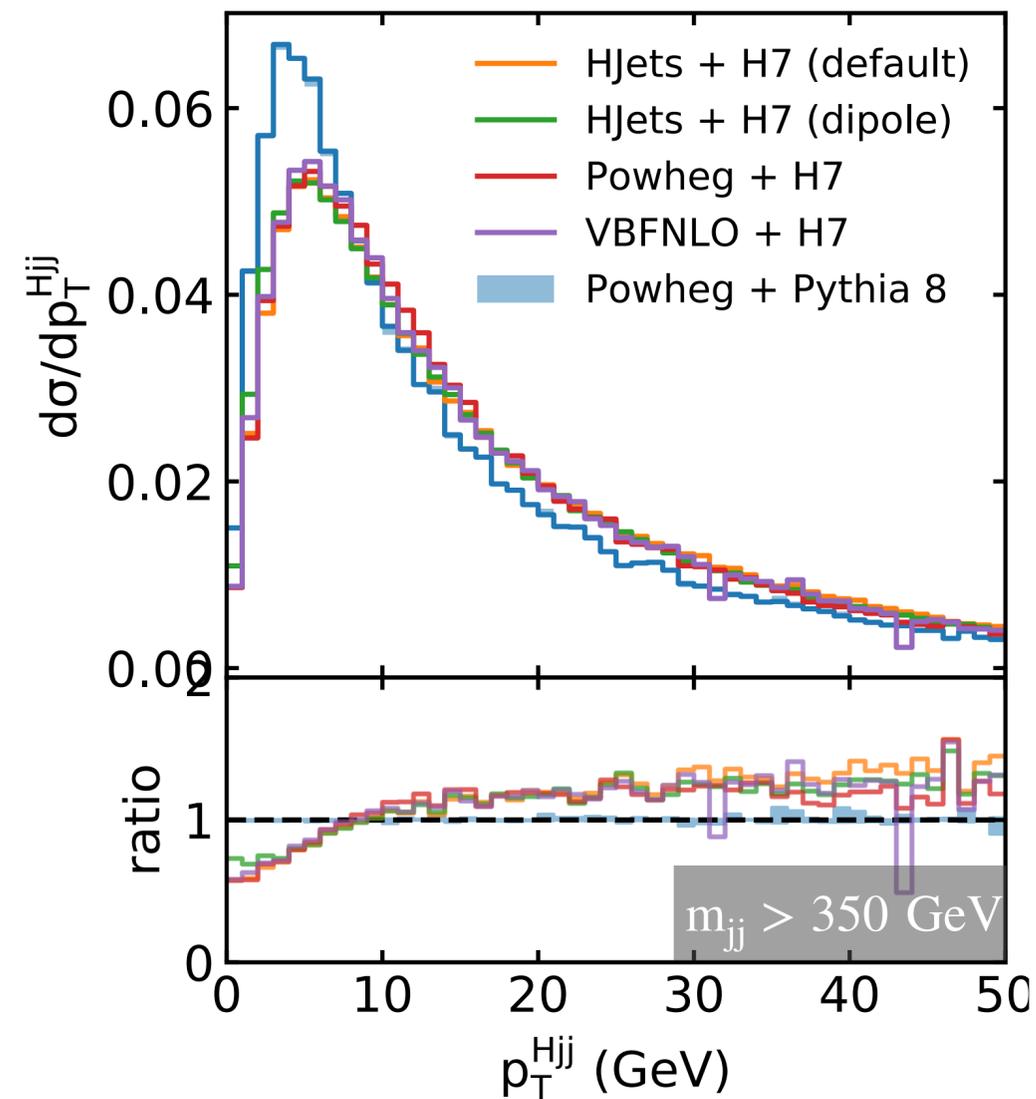
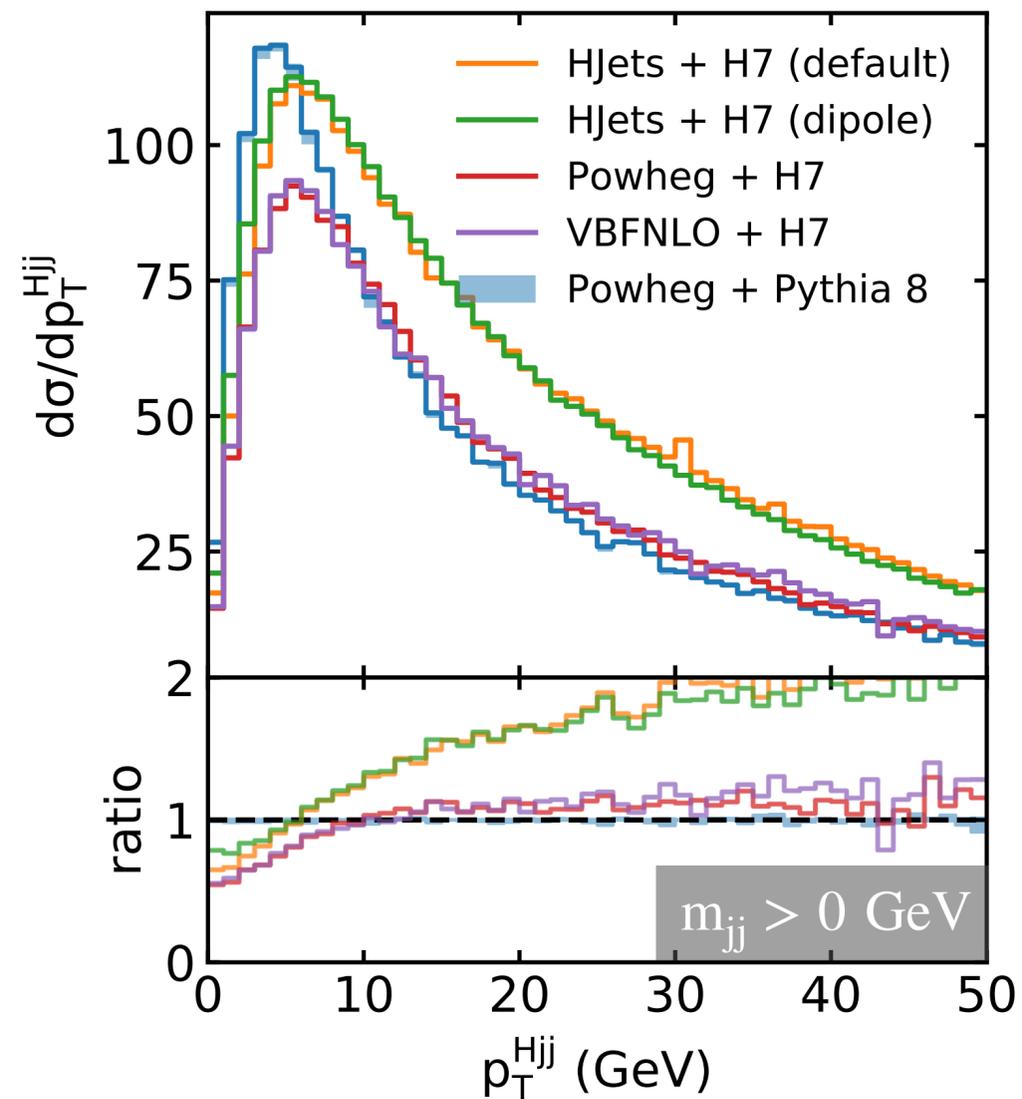


# FIXED ORDER CALCULATION VS ME+PS



- Compared ME+PS to FO calculation from proVBF-H (NNLO-QCD)
- FO-NNLO-QCD cross-section estimate is consistent with POWHEG
  - Discrepancy at low  $m_{jj}$  is due to soft emissions present in the FO NNLO calculation
  - Good agreement of at large  $m_{jj}$  values

# THE $p_T^{H_{jj}}$ OBSERVABLE

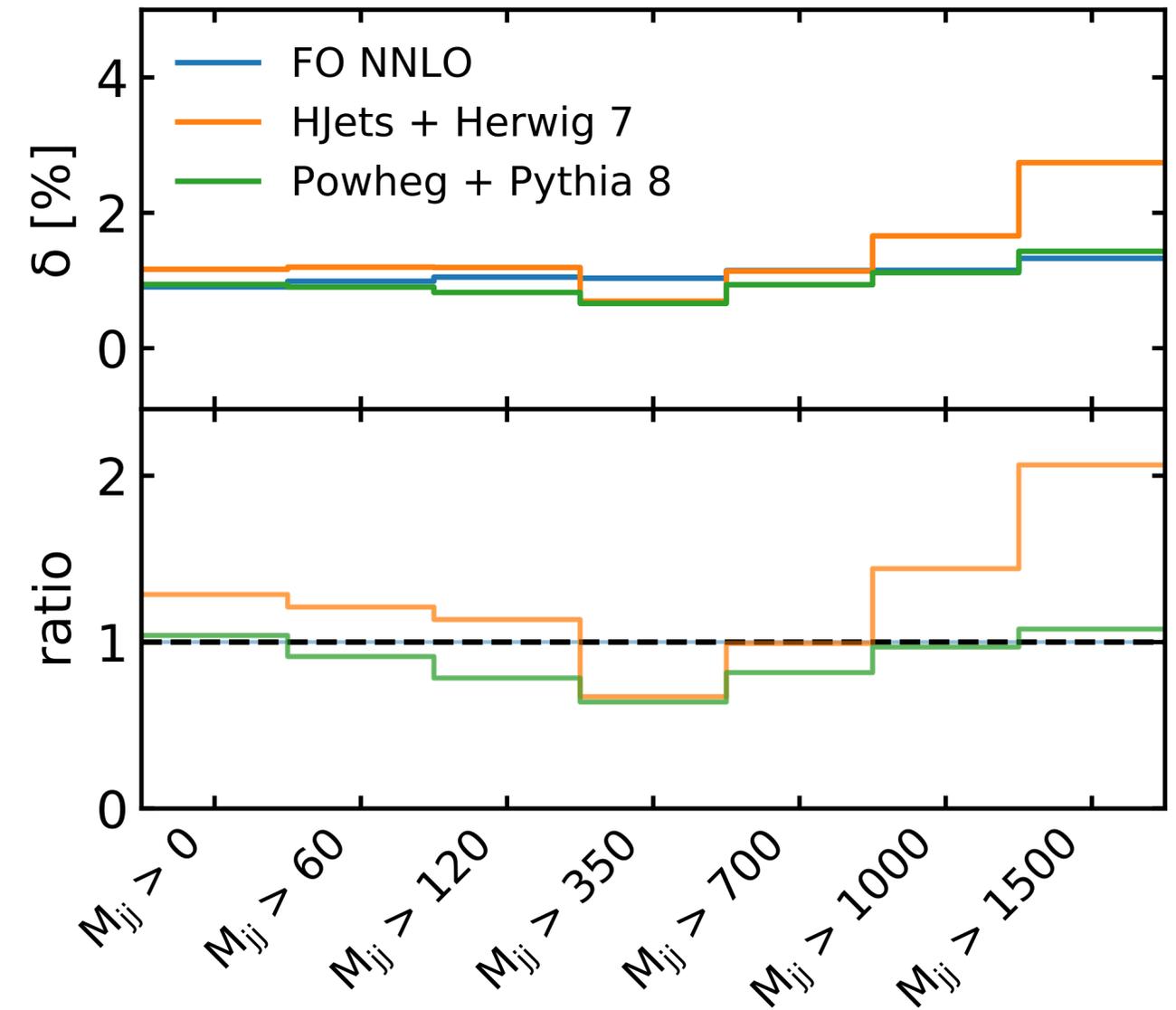
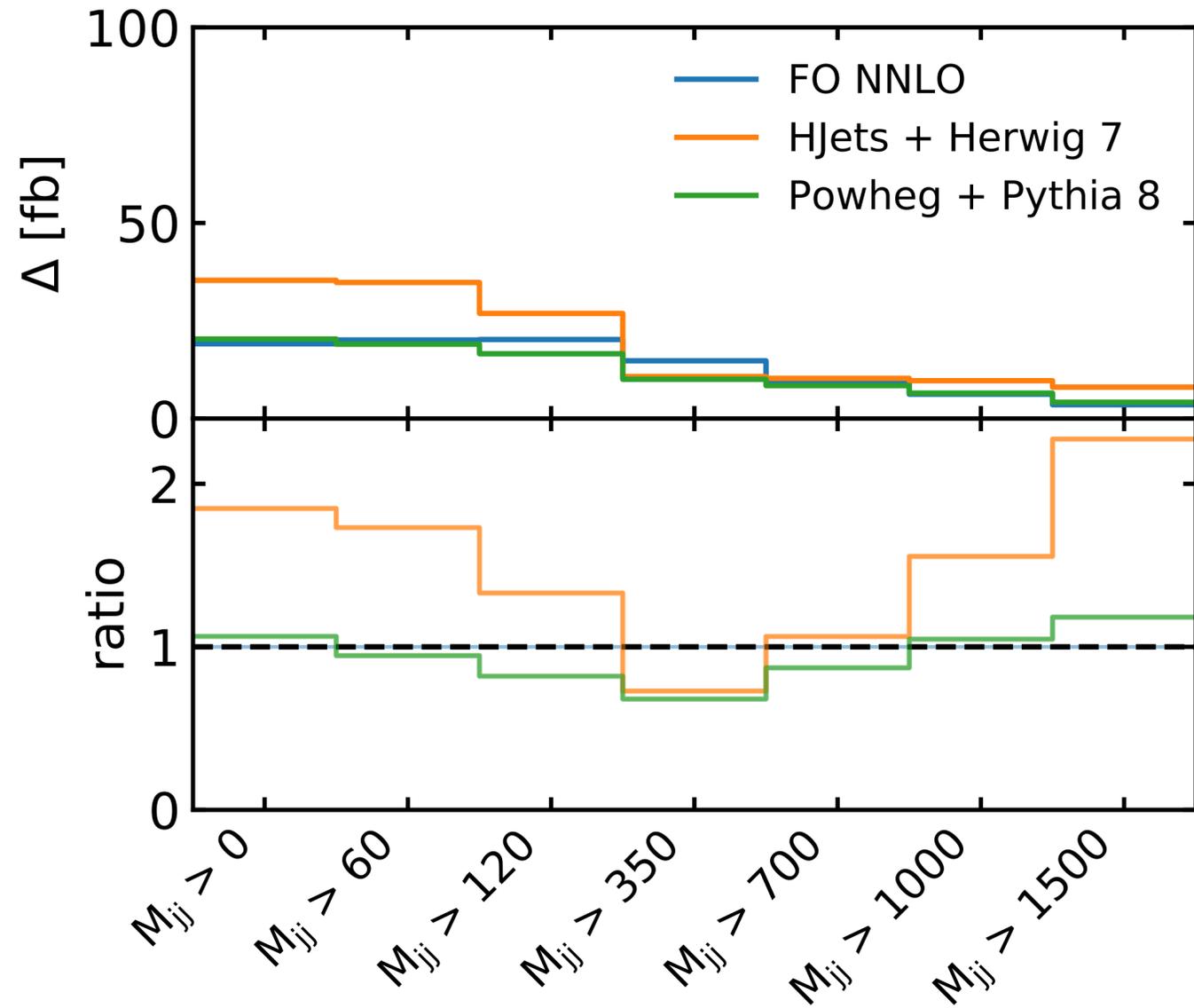


- **Over 40% discrepancy between H7 and Pythia8**
  - Region dominated by soft emissions
  - Expect large PS  $\mu_{F,R}$  scale variation
  - Still investigating these differences with PS authors



# COMPARING UNCERTAINTIES VS $m_{jj}$

- Hard process scale variations for every  $m_{jj}$  cut



# UNCERTAINTIES PROPAGATION SCHEME

- **9 migration uncertainties**  $\Delta_{2\text{jets}}$ ,  $\Delta_{200}$  and  $\Delta_{25}$ 
  - $\Delta_{2\text{jets}}$ ,  $\Delta_{200}$  and  $\Delta_{25}$
  - 6 NPs to describe  $M_{jj}$  spectrum
- **1 yield uncertainty on the inclusive cross-section,  $\Delta_{\text{tot}}$**
- **Uncertainties computed by varying the QCD scales**
  - Extracted using ME or FO NNLO
- **Bins acceptance computed with ME+PS**
  - Moved from the fraction of the  $\Delta$  distributed across STSX bins to  $\sigma$

$$\Delta_{\text{tot}} = \sigma_{\text{tot}} \times \delta_{\text{tot}}$$

$$\Delta_{2j} = \sigma_{2j} \times (\delta_{2j}^2 - \delta_{\text{tot}}^2)^{1/2}$$

$$\Delta_{60} = \sigma_{m_{jj}>60} \times (\delta_{m_{jj}>60}^2 - \delta_{2j}^2)^{1/2}$$

$$\dots = \dots$$

$$\Delta_{350} = \sigma_{m_{jj}>350} \times (\delta_{m_{jj}>350}^2 - \delta_{m_{jj}>120}^2)^{1/2}$$

$$\dots = \dots$$

- If undefined uncertainty (ex:  $\delta_{[350,\infty]} < \delta_{[120,\infty]}$ )  
→ Replace with:

$$\Delta_{350} = \sigma_{m_{jj}>350} \times \rho \cdot \delta_{m_{jj}>350}$$

- $\rho$  value of 1/2 is assumed for the remaining talk



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- $\Delta_{2\text{jets}}$ ,  $\Delta_{200}$  and  $\Delta_{25}$
- 6 NPs to describe  $M_{jj}$  spectrum

- **1 yield uncertainty on the inclusive cross-section,  $\Delta_{\text{tot}}$**

$$\Delta_{\text{tot}} = \sigma_{\text{tot}} \times \delta_{\text{tot}}$$

$$\Delta_{2j} = \sigma_{2j} \times (\delta_{2j}^2 - \delta_{\text{tot}}^2)^{1/2}$$

$$\Delta_{60} = \sigma_{m_{jj}>60} \times (\delta_{m_{jj}>60}^2 - \delta_{2j}^2)^{1/2}$$

$$\dots = \dots$$

$$\Delta_{350} = \sigma_{m_{jj}>350} \times (\delta_{m_{jj}>350}^2 - \delta_{m_{jj}>120}^2)^{1/2}$$

$$\dots = \dots$$

- Total yield uncertainty taken from YR4 the  $\delta_{\text{tot}} \sim 0.38\%$  (need to be updated)
- From  $\delta_{2\text{jet}}$  remove contribution from  $\delta_{\text{tot}}$
- The effect of each migration  $\Delta$  is anti-correlated for bins above/below

- If undefined uncertainty (ex:  $\delta_{[350,\infty]} < \delta_{[120,\infty]}$ )  
→ Replace with:

$$\Delta_{350} = \sigma_{m_{jj}>350} \times \rho \cdot \delta_{m_{jj}>350}$$

- $\rho$  value of 1/2 is assumed for the remaining talk

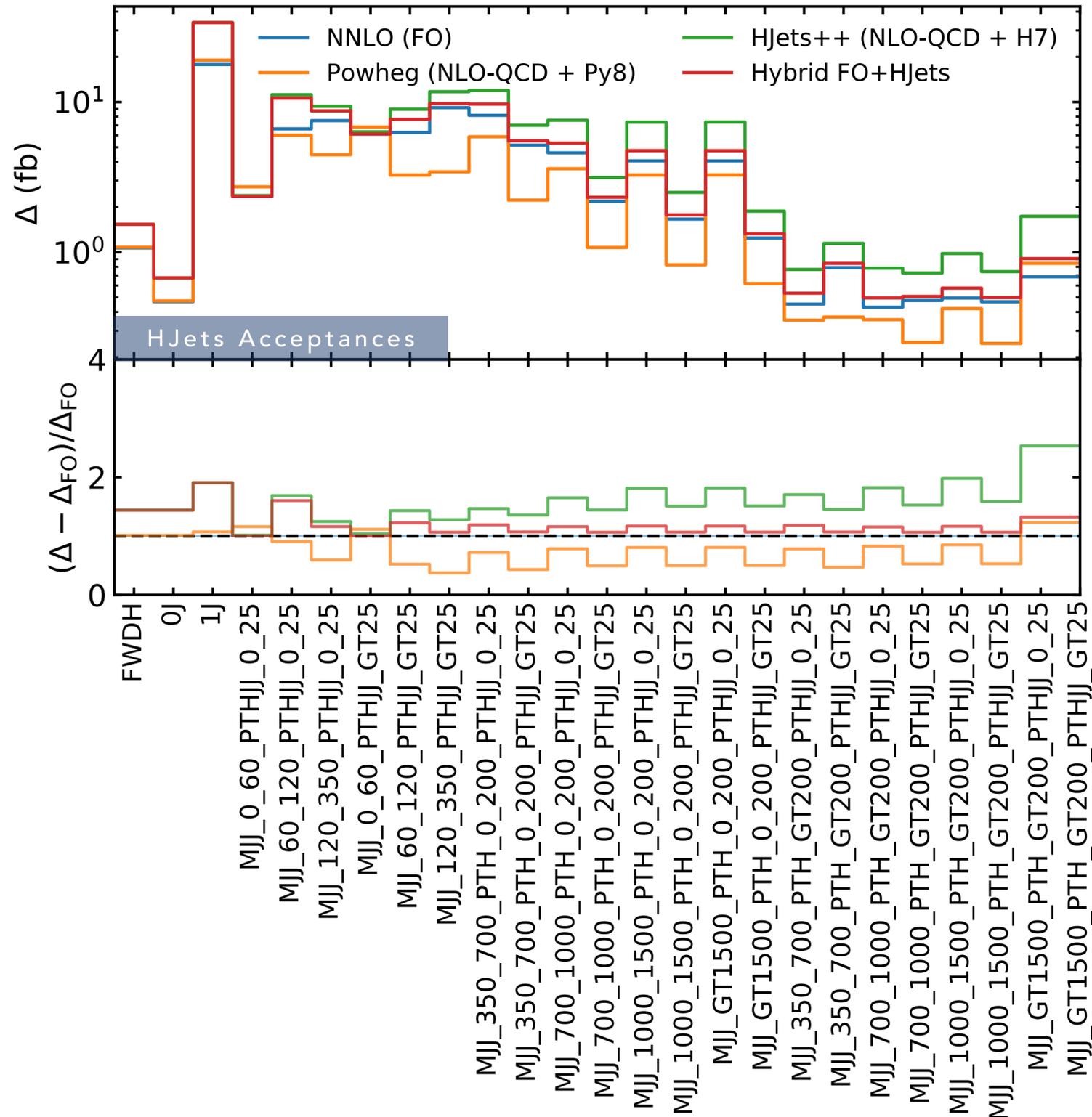
# DEFINITION OF ACCEPTANCES

- Basic definition: bin cross-section divided by cross-section in the NP phase space
  - Exception:  $\Delta_{XX}$  for bin with a cut :  $m_{jj} < X$
- The effect of each migration  $\Delta$  is anti-correlated for bins above/below

	$\Delta_{tot}$	$\Delta_{200}$	$\Delta_{60}$	$\Delta_{120}$	$\Delta_{350}$	$\Delta_{700}$	$\Delta_{1000}$	$\Delta_{1500}$	$\Delta_{25}$	$\Delta_{2jets}$
JET01	$\frac{\sigma_{N_{jets}<2}}{\sigma_{tot}}$	0	0	0	0	0	0	0	0	-1
MJJ_0_60_JET3	$\frac{\sigma_{0<m_{jj}<60&p_T^{Hjj}>25}}{\sigma_{tot}}$	0	$\frac{\sigma_{0<m_{jj}<60&p_T^{Hjj}>25}}{\sigma_{0<m_{jj}<60}}$	<div style="background-color: #e0f7fa; padding: 10px; border: 1px solid #00796b; display: flex; align-items: center; justify-content: center;"> <math>\sigma_{0&lt;m_{jj}&lt;60 \ \&amp; \ p_T^{Hjj}&lt;25}</math>  <math>\sigma_{0&lt;m_{jj}&lt;60}</math> </div>						$\frac{\sigma_{0<m_{jj}<60&p_T^{Hjj}>25}}{\sigma_{N_{jets}>2}}$
MJJ_0_60_JET3VETO	$\frac{\sigma_{0<m_{jj}<60&p_T^{Hjj}<25}}{\sigma_{tot}}$	0	$\frac{\sigma_{0<m_{jj}<60&p_T^{Hjj}<25}}{\sigma_{0<m_{jj}<60}}$							$\frac{\sigma_{0<m_{jj}<60&p_T^{Hjj}>25}}{\sigma_{p_T^{Hjj}>25}}$
MJJ_60_120_JET3	$\frac{\sigma_{60<m_{jj}<120&p_T^{Hjj}>25}}{\sigma_{tot}}$	0	$\frac{\sigma_{60<m_{jj}<120&p_T^{Hjj}>25}}{\sigma_{m_{jj}>60}}$	<div style="background-color: #fff9c4; padding: 10px; border: 1px solid #f96; display: flex; align-items: center; justify-content: center;"> <math>\sigma_{60&lt;m_{jj}&lt;120 \ \&amp; \ p_T^{Hjj}&lt;25}</math>  <math>\sigma_{m_{jj}&gt;60}</math> </div>						$\frac{\sigma_{60<m_{jj}<120&p_T^{Hjj}>25}}{\sigma_{N_{jets}>2}}$
MJJ_60_120_JET3VETO	$\frac{\sigma_{60<m_{jj}<120&p_T^{Hjj}<25}}{\sigma_{tot}}$	0	$\frac{\sigma_{60<m_{jj}<120&p_T^{Hjj}<25}}{\sigma_{m_{jj}>60}}$							$\frac{\sigma_{60<m_{jj}<120&p_T^{Hjj}>25}}{\sigma_{p_T^{Hjj}>25}}$
MJJ_120_350_JET3	$\frac{\sigma_{120<m_{jj}<350&p_T^{Hjj}>25}}{\sigma_{tot}}$	0	$\frac{\sigma_{120<m_{jj}<350&p_T^{Hjj}>25}}{\sigma_{m_{jj}>60}}$	...	...	...	0	...	...	...
MJJ_120_350_JET3VETO	$\frac{\sigma_{120<m_{jj}<350&p_T^{Hjj}<25}}{\sigma_{tot}}$	0	$\frac{\sigma_{120<m_{jj}<350&p_T^{Hjj}<25}}{\sigma_{m_{jj}>60}}$	...	...	...	0	...	...	...
MJJ_350_700_pTH_0_200_JET3	$\frac{\sigma_{350<m_{jj}<700&p_T^H<200&p_T^{Hjj}>25}}{\sigma_{tot}}$	$\frac{\sigma_{350<m_{jj}<700&p_T^H<200&p_T^{Hjj}>25}}{\sigma_{p_T^H<200}}$	$\frac{\sigma_{350<m_{jj}<700&p_T^H<200&p_T^{Hjj}>25}}{\sigma_{m_{jj}>60}}$	...	...	...	...	...	...	$\frac{\sigma_{350<m_{jj}<700&p_T^H<200&p_T^{Hjj}>25}}{\sigma_{N_{jets}>2}}$
MJJ_350_700_pTH_0_200_JET3VETO	$\frac{\sigma_{350<m_{jj}<700&p_T^H<200&p_T^{Hjj}<25}}{\sigma_{tot}}$	$\frac{\sigma_{350<m_{jj}<700&p_T^H<200&p_T^{Hjj}<25}}{\sigma_{p_T^H<200}}$	$\frac{\sigma_{350<m_{jj}<700&p_T^H<200&p_T^{Hjj}<25}}{\sigma_{m_{jj}>60}}$	...	...	...	...	...	...	$\frac{\sigma_{350<m_{jj}<700&p_T^H<200&p_T^{Hjj}<25}}{\sigma_{p_T^{Hjj}>25}}$
...	...	...	...	...	...	...	...	...	...	...
MJJ_700_1000_pTH_0_200_JET3	$\frac{\sigma_{700<m_{jj}<1000&p_T^H<200&p_T^{Hjj}>25}}{\sigma_{tot}}$	$\frac{\sigma_{700<m_{jj}<1000&p_T^H<200&p_T^{Hjj}>25}}{\sigma_{p_T^H<200}}$	$\frac{\sigma_{700<m_{jj}<1000&p_T^H<200&p_T^{Hjj}>25}}{\sigma_{m_{jj}>60}}$	...	...	...	$\frac{\sigma_{700<m_{jj}<1000&p_T^H<200&p_T^{Hjj}>25}}{\sigma_{m_{jj}<1000}}$	...	...	...
...	...	...	...	...	...	...	...	...	...	...
MJJ_1500_pTH_0_200_JET3	$\frac{\sigma_{m_{jj}>1500&p_T^H<200&p_T^{Hjj}>25}}{\sigma_{tot}}$	$\frac{\sigma_{m_{jj}>1500&p_T^H<200&p_T^{Hjj}>25}}{\sigma_{p_T^H<200}}$	$\frac{\sigma_{m_{jj}>1500&p_T^H<200&p_T^{Hjj}>25}}{\sigma_{m_{jj}>60}}$	...	...	...	$\frac{\sigma_{m_{jj}>1500&p_T^H<200&p_T^{Hjj}>25}}{\sigma_{m_{jj}>1000}}$	...	...	...
...	...	...	...	...	...	...	...	...	...	...
MJJ_350_700_pTH_gt200_JET3	$\frac{\sigma_{350<m_{jj}<700&p_T^H>200&p_T^{Hjj}>25}}{\sigma_{tot}}$	$\frac{\sigma_{350<m_{jj}<700&p_T^H>200&p_T^{Hjj}>25}}{\sigma_{p_T^H>200}}$	$\frac{\sigma_{350<m_{jj}<700&p_T^H>200&p_T^{Hjj}>25}}{\sigma_{m_{jj}>60}}$	...	...	...	0	...	...	...
MJJ_350_700_pTH_gt200_JET3VETO	$\frac{\sigma_{350<m_{jj}<700&p_T^H>200&p_T^{Hjj}<25}}{\sigma_{tot}}$	$\frac{\sigma_{350<m_{jj}<700&p_T^H>200&p_T^{Hjj}<25}}{\sigma_{p_T^H>200}}$	$\frac{\sigma_{350<m_{jj}<700&p_T^H>200&p_T^{Hjj}<25}}{\sigma_{m_{jj}>60}}$	...	...	...	0	...	...	...
...	...	...	...	...	...	...	...	...	...	...
MJJ_700_1000_pTH_gt200_JET3	$\frac{\sigma_{700<m_{jj}<1000&p_T^H>200&p_T^{Hjj}>25}}{\sigma_{tot}}$	$\frac{\sigma_{700<m_{jj}<1000&p_T^H>200&p_T^{Hjj}>25}}{\sigma_{p_T^H>200}}$	$\frac{\sigma_{700<m_{jj}<1000&p_T^H>200&p_T^{Hjj}>25}}{\sigma_{m_{jj}>60}}$	...	...	...	$\frac{\sigma_{700<m_{jj}<1000&p_T^H>200&p_T^{Hjj}>25}}{\sigma_{m_{jj}<1000}}$	...	...	...
...	...	...	...	...	...	...	...	...	...	...
MJJ_1500_pTH_gt200_JET3	$\frac{\sigma_{m_{jj}>1500&p_T^H>200&p_T^{Hjj}>25}}{\sigma_{tot}}$	$\frac{\sigma_{m_{jj}>1500&p_T^H>200&p_T^{Hjj}>25}}{\sigma_{p_T^H>200}}$	$\frac{\sigma_{m_{jj}>1500&p_T^H>200&p_T^{Hjj}>25}}{\sigma_{m_{jj}>60}}$	...	...	...	$\frac{\sigma_{m_{jj}>1500&p_T^H>200&p_T^{Hjj}>25}}{\sigma_{m_{jj}>1000}}$	...	...	$\frac{\sigma_{m_{jj}>1500&p_T^H>200&p_T^{Hjj}>25}}{\sigma_{p_T^{Hjj}>25}}$
...	...	...	...	...	...	...	...	...	...	...



# COMPARING UNCERTAINTIES

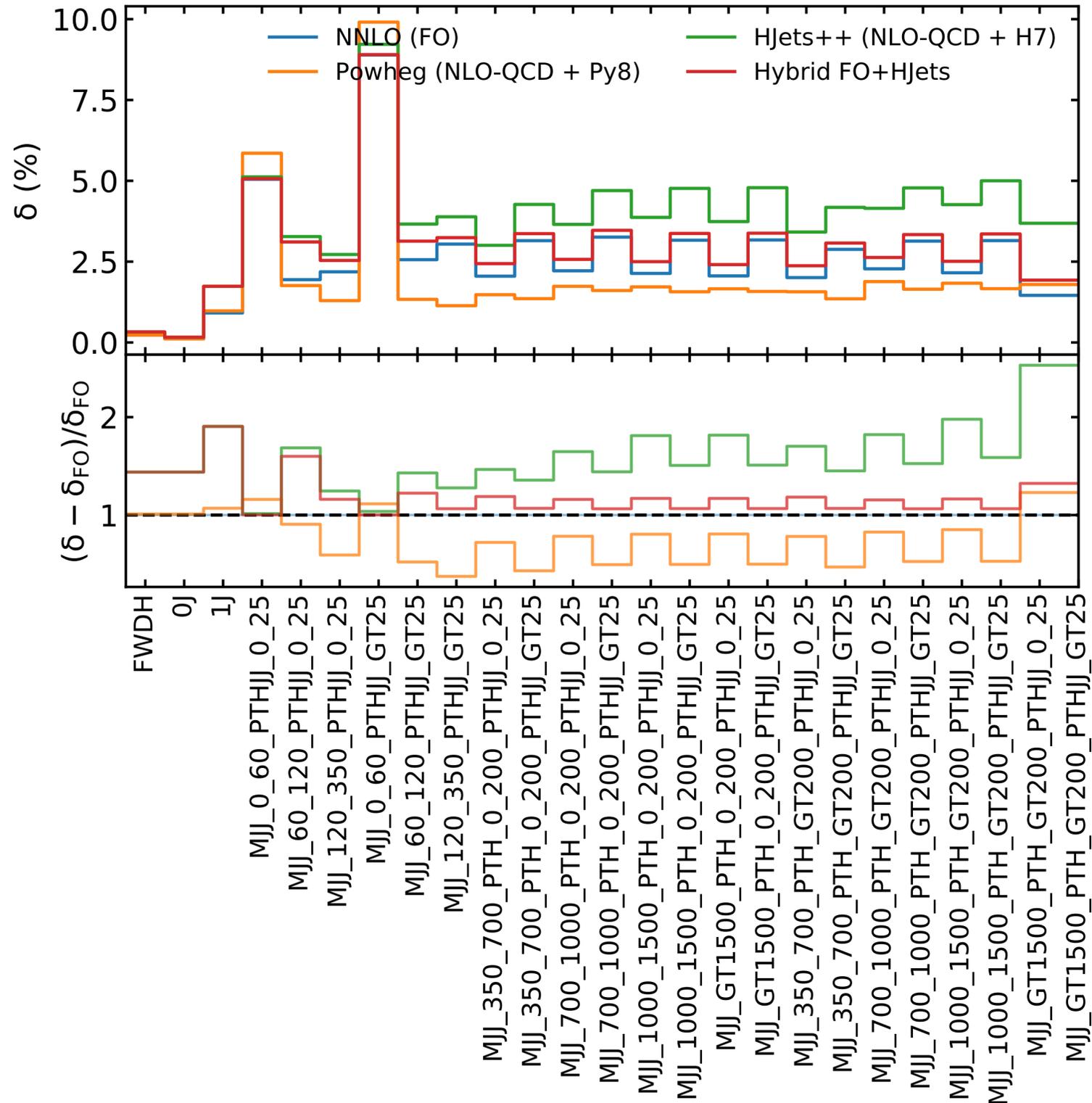


- **Acceptances updated using HJets + Herwig 7**
- **What to use for the uncertainty sources (the bing  $\Delta$ )?**
  - The 3rd jet is generated in HJets/Powheg at LO and from PS, Hence the HJets/Powheg QCD scale uncertainties in the 2j/3j bins are not reliable. **FO estimation should be used to estimate the uncertainties  $\Delta_{25}$**
- **Hybrid sources solution:**
  - S-channel contributes only in the low Mjj region  $< 350$  GeV, FO can be used for  $\Delta_{350-1500}$ , HJets for  $\Delta_{60-120}$  and  $\Delta_{2j}$
  - $\Delta_{25}$  and  $\Delta_{200}$  from FO

source [fb]	FO NNLO	POWHEG NLO	HJets NLO	MIXED
Delta_tot	14.972	15.131	21.539	21.539
Delta_200	0.622	1.081	2.989	0.622
Delta_Mjj60	8.057	9.511	8.003	8.003
Delta_Mjj120	6.84	8.286	13.446	13.446
Delta_Mjj350	7.389	5.025	5.385	7.389
Delta_Mjj700	4.201	5.973	8.158	4.201
Delta_Mjj1000	3.115	3.545	7.045	3.115
Delta_Mjj1500	1.764	2.614	6.404	1.764
Delta_25	27.387	2.674	35.46	27.387
Delta_2jet	17.355	18.617	33.412	33.412



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Delta_25	27.387	2.674	35.46	27.387
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# ELECTROWEAK CORRECTIONS IN STXS BINS

- **The state of the art calculation from HAWK 2.0**

[Denner, Dittmaier, SK, Muck [arXiv:1412.5390]]

- Provides complete NLO QCD and EWK corrections and includes s-channel and interferences
- provides predictions for partonic channels with incoming photons as part of NLO EW corrections
- EW corrections of 5-10% in VBF production
- Enhanced electroweak corrections at high energies: driven by Sudakov  $\log \alpha \rightarrow \alpha \log(Q/M_w)$  at Higgs  $p_T$  tail

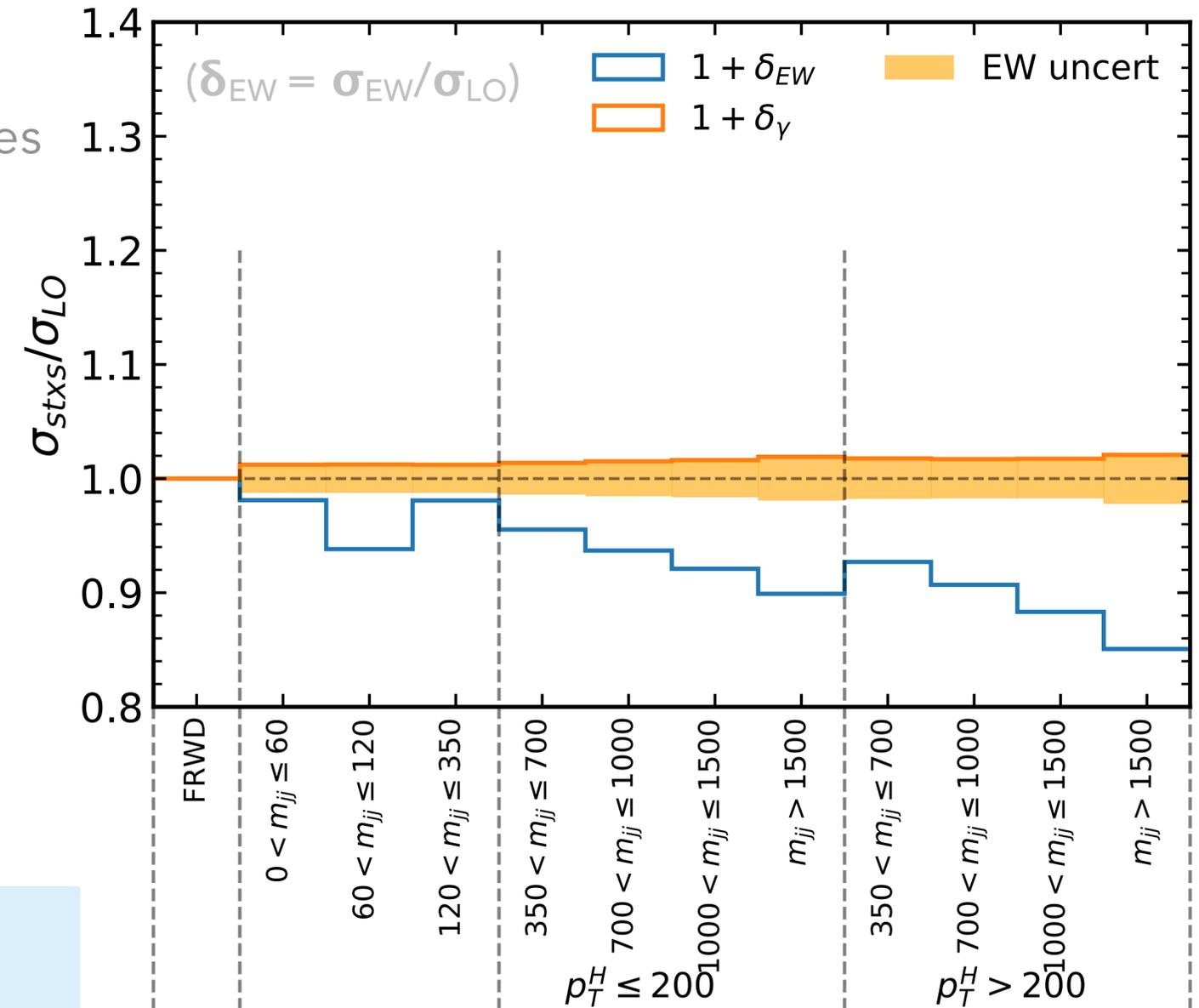
[Ciccolini, Denner, Dittmaier [arXiv:0710.4749]]

- Uncertainty estimated following the same prescription as in the Yellow Report 4

$$\Delta_{EW} = \max\{0.5\%, \delta_{EW}^2, \sigma_\gamma/\sigma_{VBF}\}$$

- **Proposition:**

- Since EW correction is driven by Sudakov log we can consider  $\delta_{EW}^2$  as the pure Sudakov nuisance:  $\Delta_{sud}$
- $\delta_\gamma$  can be considered as a separate nuisance for non-Sudakov nuisance:  $\Delta_\gamma$



# HOW TO INCLUDE THE EW CORRECTIONS ?

	STXS bin	$\sigma_{LO}(\text{fb})$	$(1 + \delta_{EW})$	$\sigma_{\gamma}(\text{fb})$	$\Delta_{EW}$
	$0 < m_{jj} \leq 60$	6.67	0.981	0.081	0.012
	$60 < m_{jj} \leq 120$	601.78	0.938	7.440	0.012
	$120 < m_{jj} \leq 350$	540.59	0.981	6.567	0.012
$p_T^H \leq 200$	$350 < m_{jj} \leq 700$	659.75	0.955	9.056	0.014
	$700 < m_{jj} \leq 1000$	318.83	0.937	4.820	0.015
	$1000 < m_{jj} \leq 1500$	275.94	0.921	4.481	0.016
	$m_{jj} > 1500$	251.33	0.899	4.798	0.019
$p_T^H > 200$	$350 < m_{jj} \leq 700$	45.72	0.927	0.807	0.018
	$700 < m_{jj} \leq 1000$	37.91	0.907	0.647	0.017
	$1000 < m_{jj} \leq 1500$	44.03	0.883	0.765	0.017
	$m_{jj} > 1500$	55.99	0.851	1.165	0.022

- Start with best QCD prediction for VBF and assume approximate factorisation of corrections:

$$\sigma_{\text{VBF}} = \sigma_{\text{best}}(1 + \delta_{EW}) + \sigma_{\text{gamma}} \text{ with } \sigma_{EW} = \sigma_{EW}/\sigma_{LO}$$

- These correction are now implanted in the VBF-uncertainty tool



# CONCLUSION

- Update uncertainties and acceptances are now available VBF uncertainty standalone tool [[here](#)]
- The tool apply uncertainties as event weights (same strategy as ggH)
- Acceptance and uncertainties updated with full EW H+2j calculation
- Hybrid uncertainties using FO NNLO + HJets is set as default
  - Other configuration with only FO/POWHEG/HJets++ also available
- **Large differences observed in  $p_T^{Hjj}$  observable:**
  - **Values might change in the coming months depending on the PS authors inputs**
- Electroweak correction have been estimated in STXS bins
  - Corrections of 5-10% in VBF production
  - Available in the VBF uncertainty tool
  - **We still need input from EW expert on uncertainties**



BACKUP

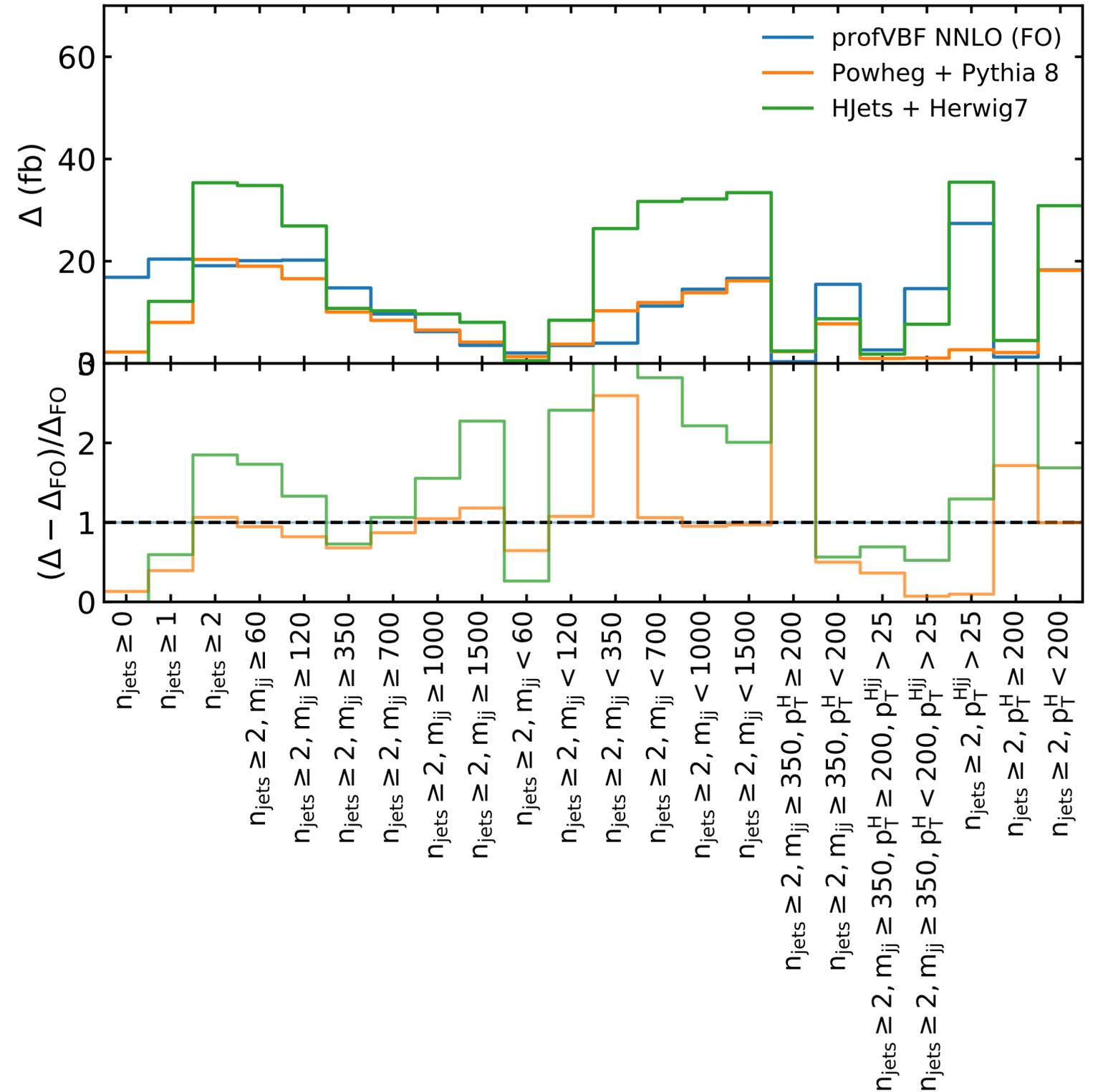
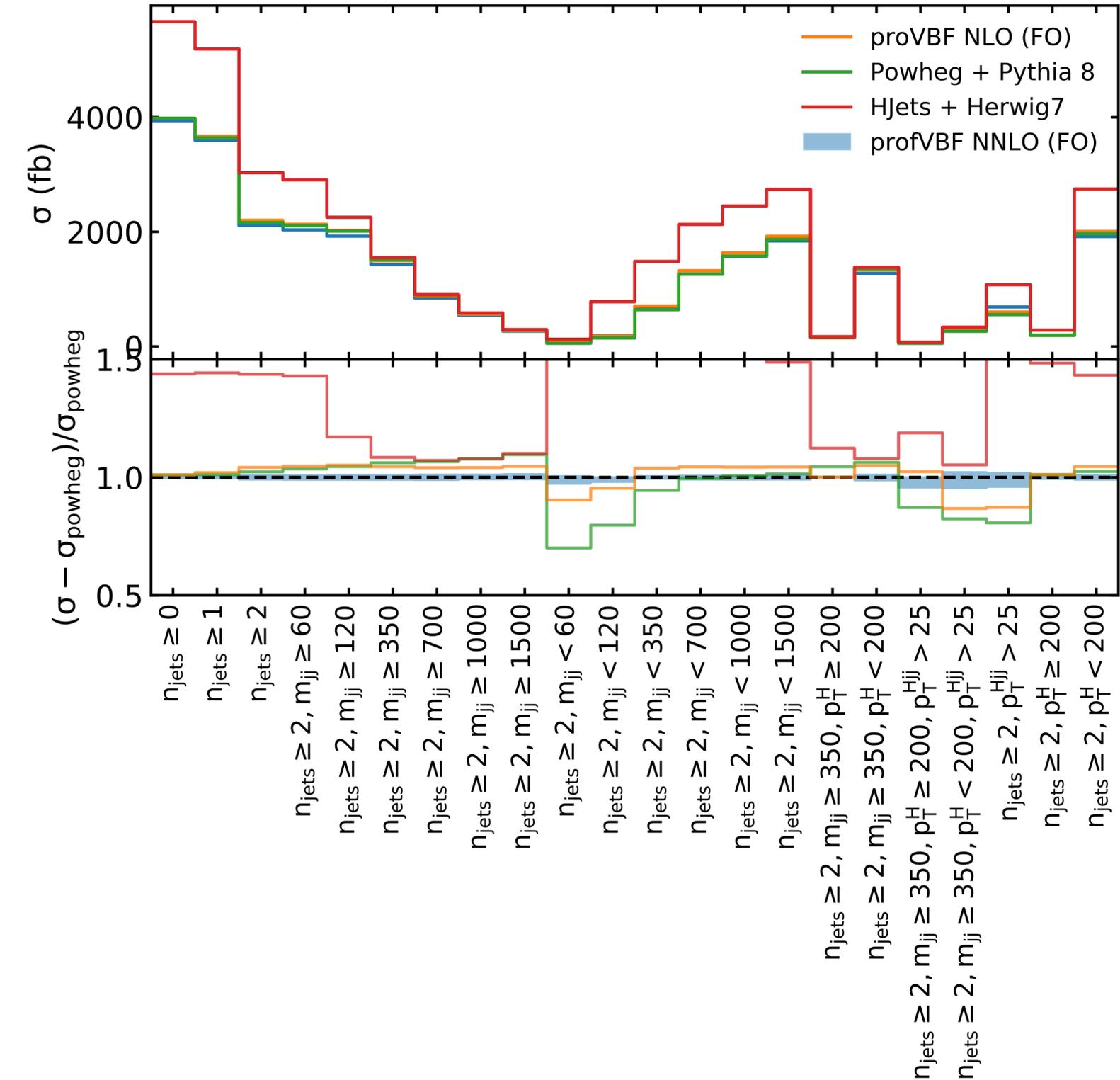
# SETUP AND DEFINITIONS

- Extracted using QCD variations of the renormalisation and factorisation scales  $\mu_r$  ,  $\mu_f$  from POWHEG + PYTHIA 8
- Keeping only variations with  $1/2 \leq \mu_r$  ,  $\mu_f \leq 2$  ,  $1/2 \leq \mu_r/\mu_f \leq 2$
- Take uncertainty envelope
- Uncertainty propagation based on Stewart-Tackmann method [1] :

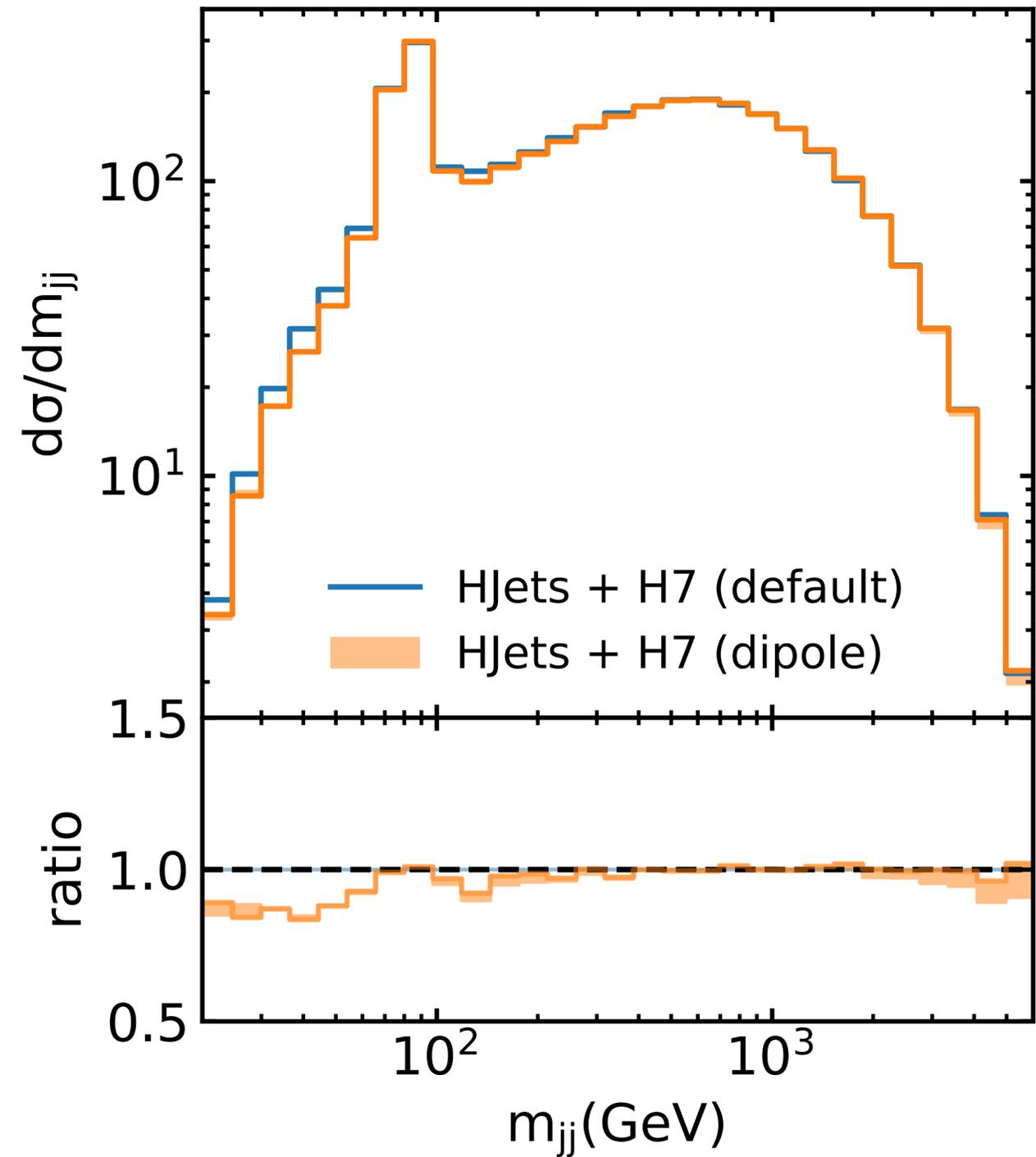
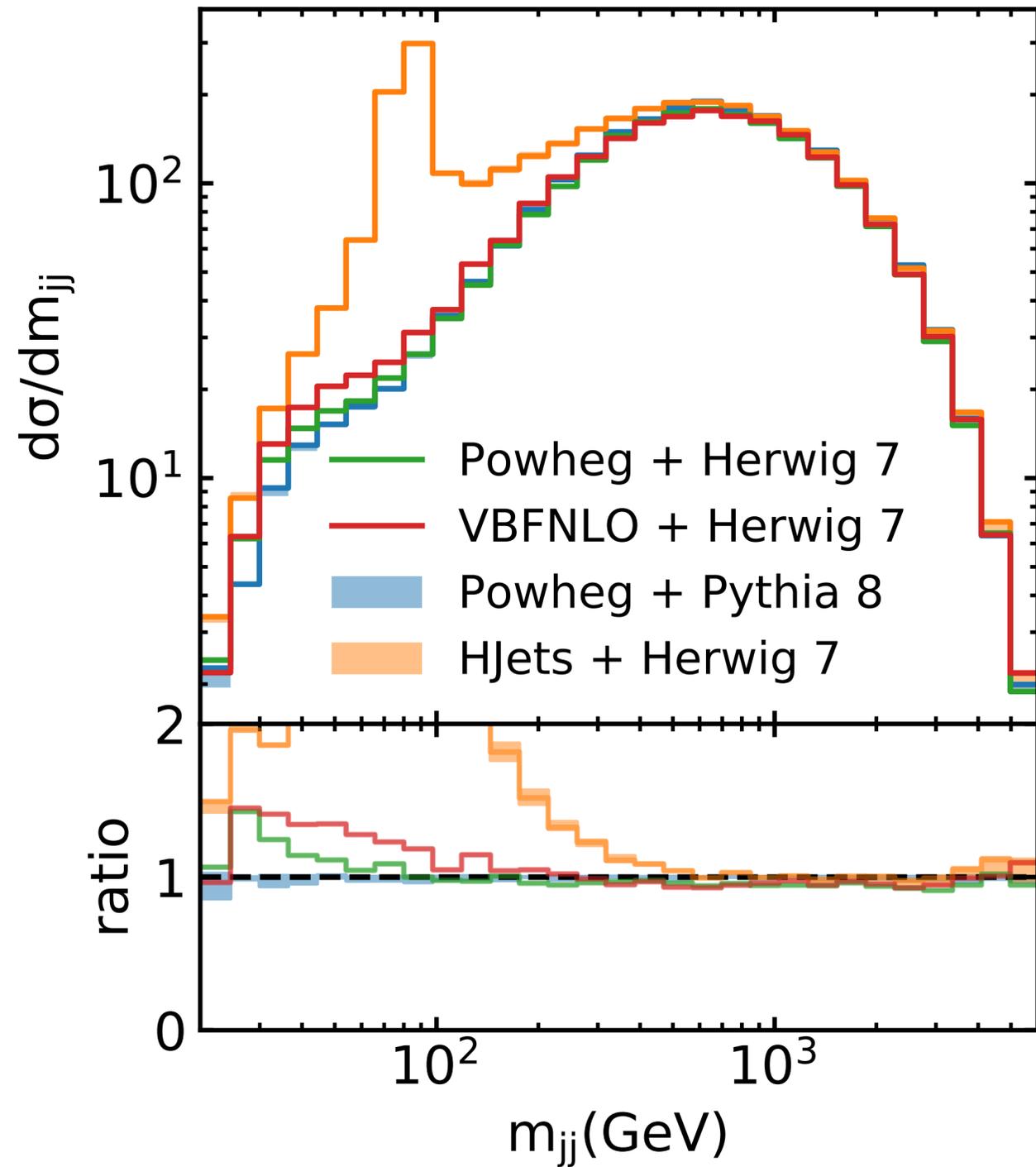
$$C(\{\sigma_0, \sigma_{\geq 1}\}) = \begin{pmatrix} (\Delta_0^y)^2 & \Delta_0^y \Delta_{\geq 1}^y \\ \Delta_0^y \Delta_{\geq 1}^y & (\Delta_{\geq 1}^y)^2 \end{pmatrix} + \begin{pmatrix} \Delta_{cut}^2 & -\Delta_{cut}^2 \\ -\Delta_{cut}^2 & \Delta_{cut}^2 \end{pmatrix}$$

- Jet definition :
  - Higgs decay products are ignored
  - Jets built using anti- $k_T$   $R = 0.4$  from all stable particles
  - Only jet with  $p_T > 30$  GeV and  $|\eta| < 4.7$

# H+2J EWK: ADDING THE S-CHANNEL

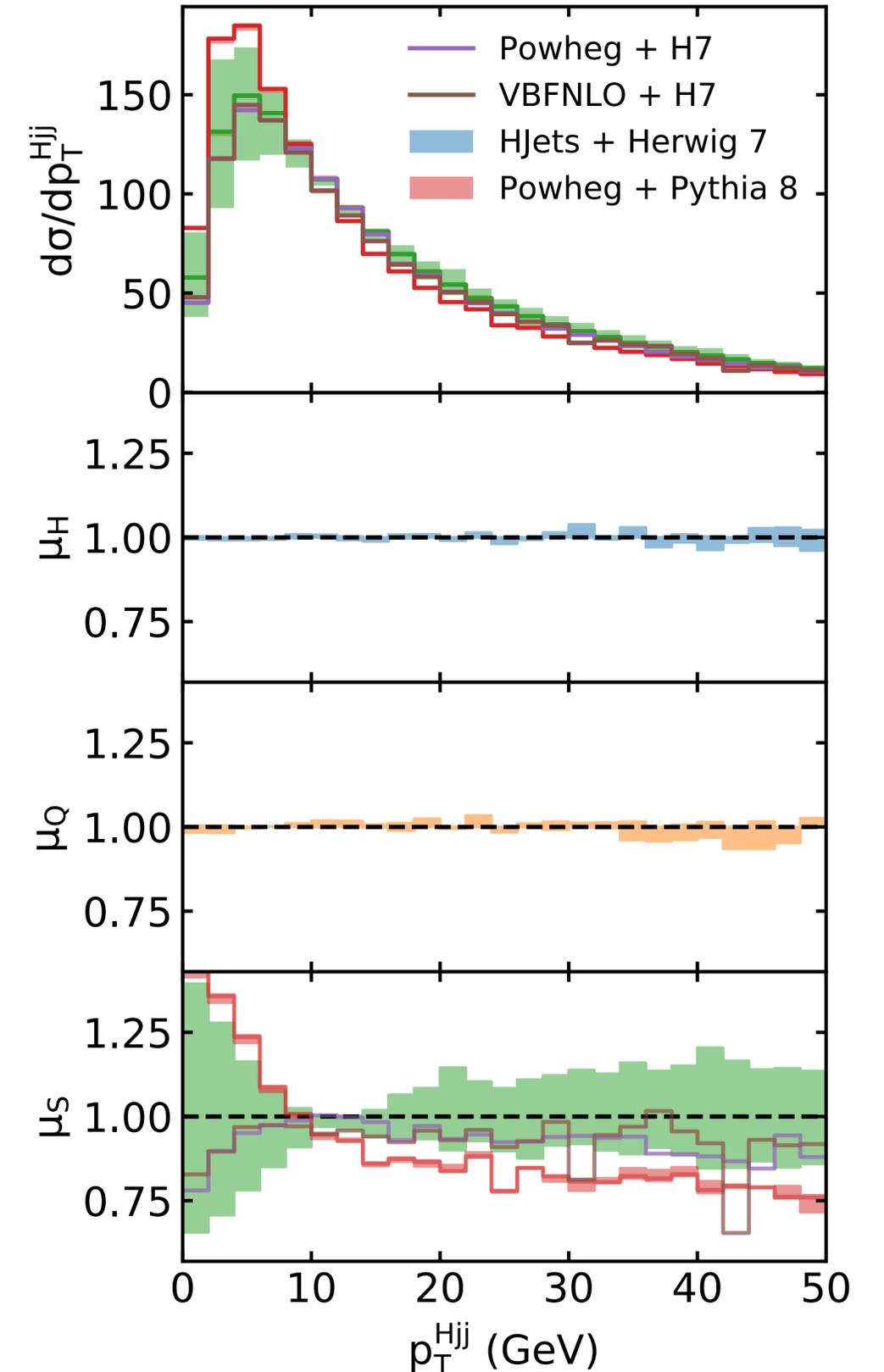
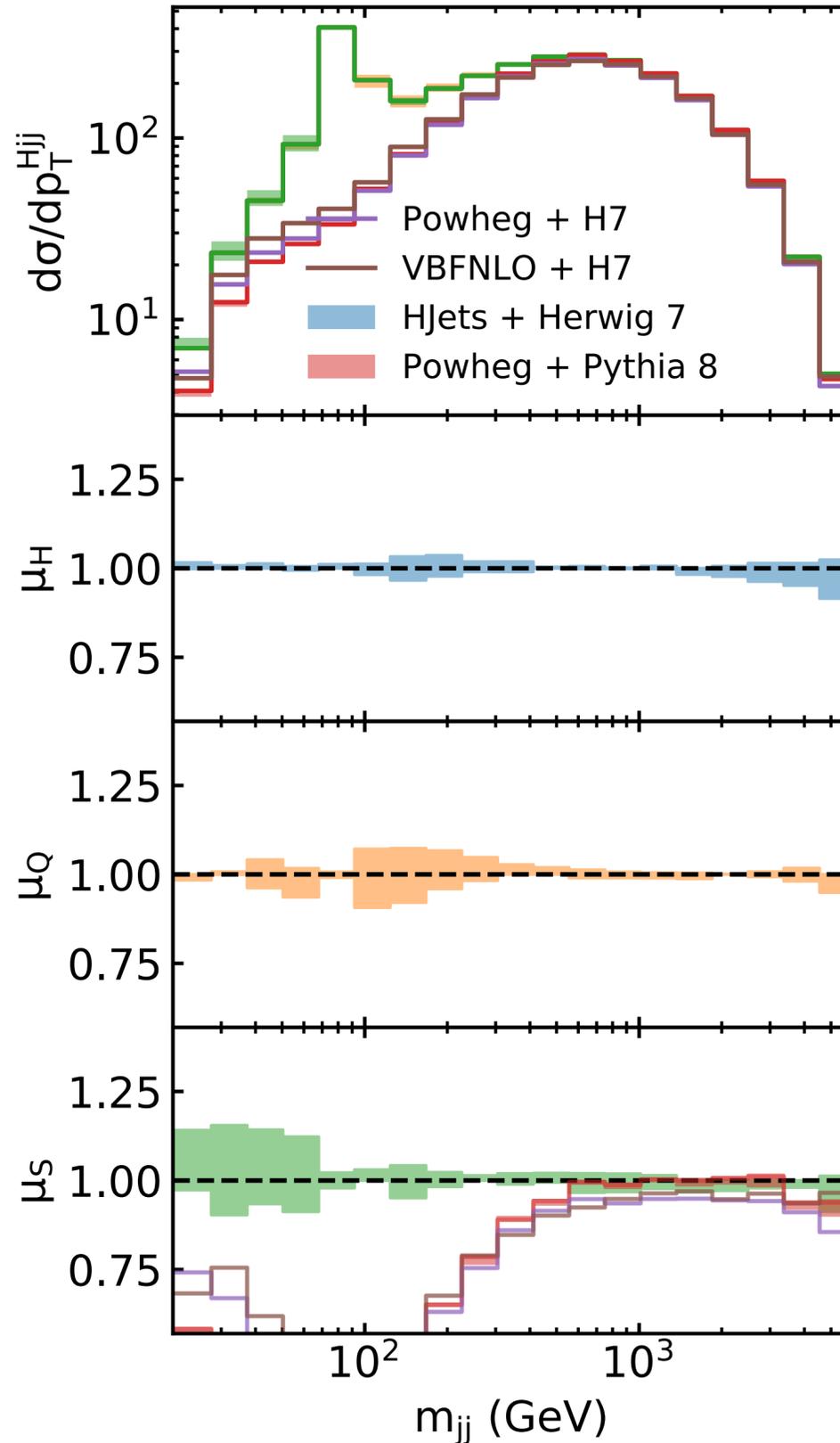


# COMPARISON WITH DIPOLE SHOWER



# PARTON SHOWER UNCERTAINTIES

- $\mu_H$ : variation of the renormalisation and factorisation scales in the hard process (Matrix Element)
- $\mu_S$ : Shower scale, it is the argument of  $\alpha_s$  and PDFs in the PS.
  - The band also includes variations in the hard process as well as the variation in renormalisation and factorisation scales for the NLO Matching and Merging
- $\mu_Q$ : veto scale: the boundary of the hardness of emissions in the PS
- low  $p_T^{\text{Hjj}}$  the PS uncertainties are huge, but expected, reaching  $\sim 40\%$ . This is also visible on the  $m_{jj}$  distribution where the PS variation gets larger in the low values where also soft emissions dominate. However, the discrepancy between Pythia8 and Herwig is till larger than the uncertainty.



# ELECTROWEAK UNCERTAINTIES

The scale uncertainty,  $\Delta_{\text{scale}}$ , results from a variation of the factorization and renormalization scales (I.5.3) by a factor of 2 keeping  $\mu_F = \mu_R$ , as indicated above, and the combined PDF $\oplus\alpha_s$  uncertainty  $\Delta_{\text{PDF}\oplus\alpha_s}$  is obtained following the PDF4LHC recipe [35]. Both  $\Delta_{\text{scale}}$  and  $\Delta_{\text{PDF}\oplus\alpha_s}$  are actually obtained from  $\sigma_{\text{NNLOQCD}}^{\text{DIS}}$ , but this QCD-driven uncertainties can be taken over as uncertainty estimates for  $\sigma^{\text{VBF}}$  as well. The theoretical uncertainties of integrated cross sections originating from unknown higher-order EW effects can be estimated by

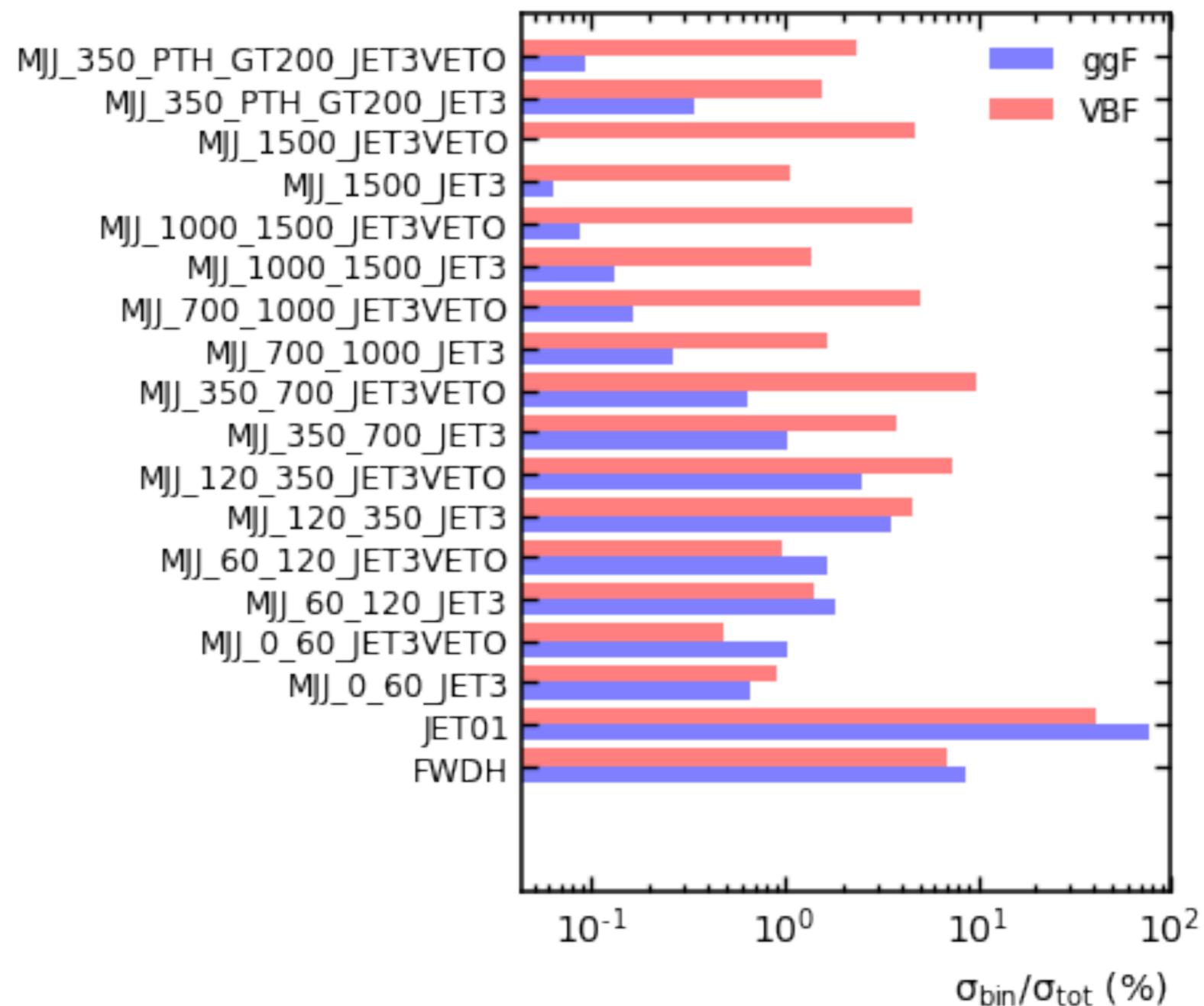
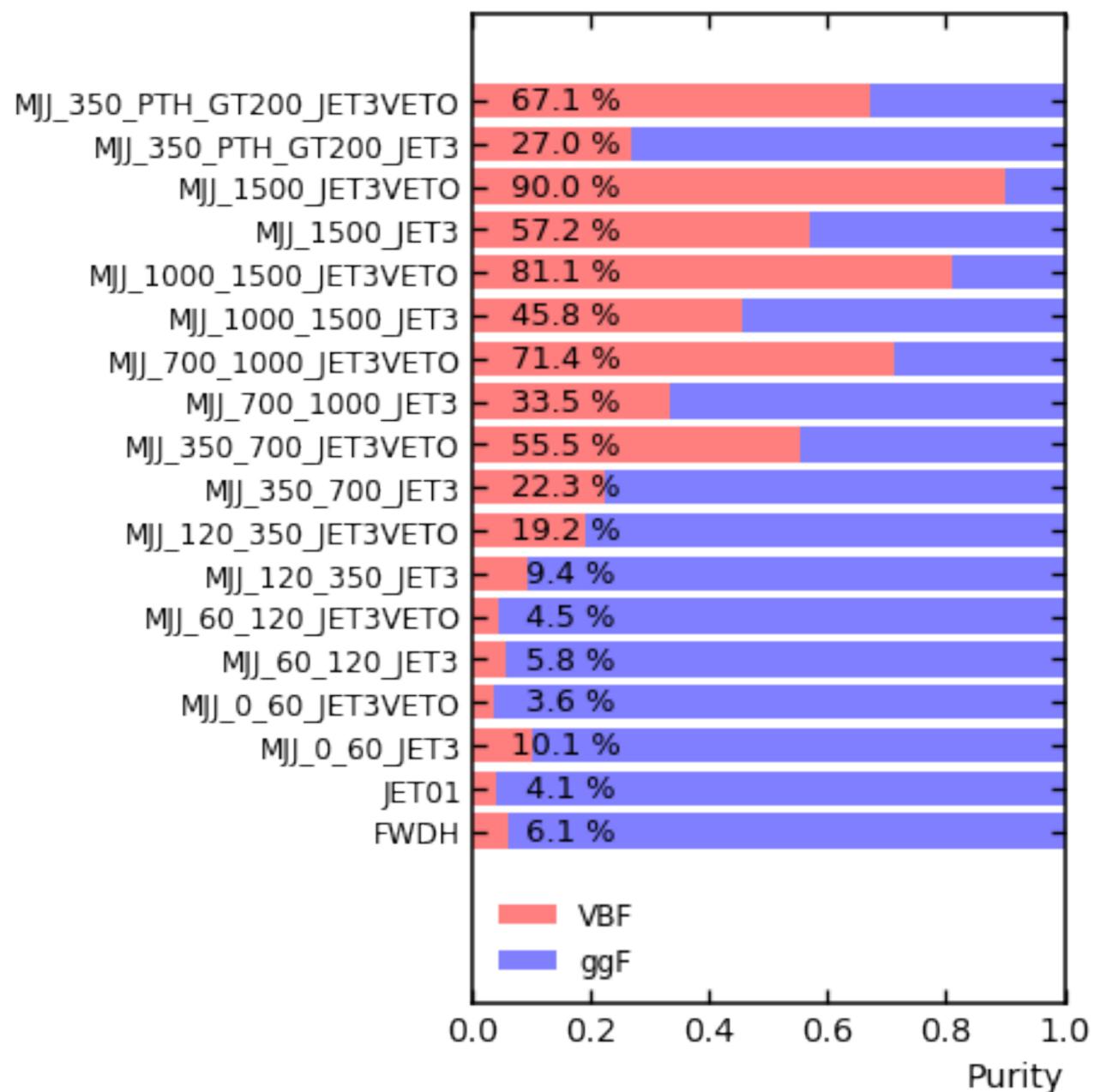
$$\Delta_{\text{EW}} = \max\{0.5\%, \delta_{\text{EW}}^2, \sigma_\gamma/\sigma^{\text{VBF}}\}. \quad (\text{I.5.7})$$

The first entry represents the generic size of NNLO EW corrections, while the second accounts for potential enhancement effects. Note that the whole photon-induced cross-section contribution  $\sigma_\gamma$  is treated as uncertainty here, because the PDF uncertainty of  $\sigma_\gamma$  is estimated to be 100% with the NNPDF2.3QED PDF set. At present, this source, which is about 1.5%, dominates the EW uncertainty of the integrated VBF cross section



# PURITY OF EW qqH BINS

- Only **ggF** and **VBF** production considered so far
- Higher VBF purity obtained thanks to the high M<sub>jj</sub> split



# ACCEPTANCES

```
// acceptances for VBH+VHHad: extracted from NJets (NLO) + H7
// it includes the full H+2Jets EWK calculation
static std::map<int, std::vector<double> > stxs_acc =
{ //stxs  tot      ptH200  mjj60   mjj120  mjj350  mjj700  mjj1000  mjj1500  ptHjj25  jet2
  { 200, {0.083 , 0.0   , 0.0   , 0.0   , 0.0   , 0.0   , 0.0   , 0.0   , 0.0   , 0.0   , 0.0   }}, // FWD
  { 201, {0.0735, 0.0   , 0.0   , 0.0   , 0.0   , 0.0   , 0.0   , 0.0   , 0.0   , -0.1762 }}, // Jet0
  { 202, {0.3438, 0.0   , 0.0   , 0.0   , 0.0   , 0.0   , 0.0   , 0.0   , 0.0   , -0.8238 }}, // Jet1
  { 203, {0.0082, 0.0   , -0.4038, 0.0   , 0.0   , 0.0   , 0.0   , 0.0   , -0.0256, 0.0164 }}, // Mjj 0-60,      PTHjj 0-25
  { 204, {0.0603, 0.0   , 0.1258, -0.5825, 0.0   , 0.0   , 0.0   , 0.0   , -0.1876, 0.1206 }}, // Mjj 60-120,   PTHjj 0-25
  { 205, {0.0608, 0.0   , 0.1268, 0.1617, -0.5332, 0.0   , 0.0   , 0.0   , -0.1891, 0.1216 }}, // Mjj 120-350, PTHjj 0-25
  { 206, {0.0121, 0.0   , -0.5962, 0.0   , 0.0   , 0.0   , 0.0   , 0.0   , 0.0681, 0.0243 }}, // Mjj 350-700, PTHjj 0-25   , pTH 0-200
  { 207, {0.0432, 0.0   , 0.0901, -0.4175, 0.0   , 0.0   , 0.0   , 0.0   , 0.2423, 0.0865 }}, // Mjj 700-1000, PTHjj 0-25   , pTH 0-200
  { 208, {0.0532, 0.0   , 0.111  , 0.1416, -0.4668, 0.0   , 0.0   , 0.0   , 0.2986, 0.1065 }}, // Mjj 1000-1500, PTHjj 0-25   , pTH 0-200
  { 209, {0.0702, -0.3026, 0.1465, 0.1868, 0.2682, -0.6504, 0.0   , 0.0   , -0.2185, 0.1405 }}, // Mjj 1500-inf , PTHjj 0-25   , pTH 0-200
  { 210, {0.0289, -0.1247, 0.0604, 0.077  , 0.1105, -0.2681, 0.0   , 0.0   , 0.1624, 0.0579 }}, // Mjj 350-700, PTHjj 0-25   , pTH 200-inf
  { 211, {0.0366, -0.1576, 0.0763, 0.0973, 0.1397, 0.2377, -0.6724, 0.0   , -0.1138, 0.0732 }}, // Mjj 700-1000, PTHjj 0-25   , pTH 200-inf
  { 212, {0.0118, -0.0509, 0.0246, 0.0314, 0.0451, 0.0767, -0.217  , 0.0   , 0.0662, 0.0236 }}, // Mjj 1000-1500, PTHjj 0-25   , pTH 200-inf
  { 213, {0.0335, -0.1445, 0.07   , 0.0892, 0.1281, 0.218  , 0.3371, -0.6777, -0.1043, 0.0671 }}, // Mjj 1500-inf , PTHjj 0-25   , pTH 200-inf
  { 214, {0.0093, -0.04   , 0.0193, 0.0247, 0.0354, 0.0603, 0.0932, -0.1874, 0.052  , 0.0186 }}, // Mjj 0-60,     PTHjj 25-inf
  { 215, {0.0348, -0.1498, 0.0725, 0.0925, 0.1328, 0.226  , 0.3495, 0.6955, -0.1082, 0.0696 }}, // Mjj 60-120,   PTHjj 25-inf
  { 216, {0.0069, -0.0298, 0.0144, 0.0184, 0.0264, 0.045  , 0.0695, 0.1384, 0.0388, 0.0138 }}, // Mjj 120-350, PTHjj 25-inf
  { 217, {0.004  , 0.1332, 0.0083, 0.0106, 0.0152, -0.0368, 0.0   , 0.0   , -0.0123, 0.0079 }}, // Mjj 350-700, PTHjj 25-inf , pTH 0-200
  { 218, {0.0048, 0.1623, 0.0101, 0.0129, 0.0185, -0.0448, 0.0   , 0.0   , 0.0271, 0.0097 }}, // Mjj 700-1000, PTHjj 25-inf , pTH 0-200
  { 219, {0.0033, 0.1118, 0.0069, 0.0089, 0.0127, 0.0216, -0.0612, 0.0   , -0.0104, 0.0067 }}, // Mjj 1000-1500, PTHjj 25-inf , pTH 0-200
  { 220, {0.0027, 0.0901, 0.0056, 0.0071, 0.0103, 0.0175, -0.0494, 0.0   , 0.0151, 0.0054 }}, // Mjj 1500-inf , PTHjj 25-inf , pTH 0-200
  { 221, {0.0041, 0.1361, 0.0085, 0.0108, 0.0155, 0.0264, 0.0408, -0.082  , -0.0126, 0.0081 }}, // Mjj 350-700, PTHjj 25-inf , pTH 200-inf
  { 222, {0.0026, 0.0879, 0.0055, 0.007  , 0.01  , 0.017  , 0.0263, -0.0529, 0.0147, 0.0052 }}, // Mjj 700-1000, PTHjj 25-inf , pTH 200-inf
  { 223, {0.0057, 0.19   , 0.0118, 0.0151, 0.0216, 0.0368, 0.0569, 0.1133, -0.0176, 0.0113 }}, // Mjj 1000-1500, PTHjj 25-inf , pTH 200-inf
  { 224, {0.0026, 0.0886, 0.0055, 0.007  , 0.0101, 0.0172, 0.0265, 0.0528, 0.0148, 0.0053 }}, // Mjj 1500-inf , PTHjj 25-inf , pTH 200-inf
};
```



# ACCEPTANCES

```
// bin acceptances extracted from POWHEG VBFH (NLO)
static std::map<int, std::vector<double> > stxs_acc_powheg = {
  //stxs    total    ptH_200  mjj_60   mjj_120  mjj_350  mjj_700  mjj_1000  mjj_1500  ptHjj_25  njets_30_2
  { 200 , {0.0668, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000}},
  { 201 , {0.0765, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, -0.1821}},
  { 202 , {0.3435, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, -0.8179}},
  { 203 , {0.0048, 0.0000, -0.3761, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, -0.0126, 0.0093}},
  { 204 , {0.0096, 0.0000, 0.0192, -0.4400, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, -0.0253, 0.0187}},
  { 205 , {0.0782, 0.0000, 0.1564, 0.1635, -0.6859, 0.0000, 0.0000, 0.0000, 0.0000, -0.2056, 0.1525}},
  { 206 , {0.0079, 0.0000, -0.6239, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0599, 0.0155}},
  { 207 , {0.0122, 0.0000, 0.0245, -0.5600, 0.0000, 0.0000, 0.0000, 0.0000, 0.0000, 0.0923, 0.0239}},
  { 208 , {0.0358, 0.0000, 0.0716, 0.0749, -0.3141, 0.0000, 0.0000, 0.0000, 0.0000, 0.2701, 0.0698}},
  { 209 , {0.1061, -0.3265, 0.2121, 0.2218, 0.2912, -0.7233, 0.0000, 0.0000, 0.0000, -0.2789, 0.2068}},
  { 210 , {0.0306, -0.0940, 0.0611, 0.0639, 0.0838, -0.2083, 0.0000, 0.0000, 0.0000, 0.2304, 0.0595}},
  { 211 , {0.0545, -0.1678, 0.1090, 0.1140, 0.1497, 0.2505, -0.7179, 0.0000, 0.0000, -0.1434, 0.1063}},
  { 212 , {0.0136, -0.0417, 0.0271, 0.0283, 0.0372, 0.0623, -0.1784, 0.0000, 0.0000, 0.1022, 0.0264}},
  { 213 , {0.0504, -0.1550, 0.1007, 0.1052, 0.1382, 0.2313, 0.3553, -0.7105, -0.1324, 0.0982}},
  { 214 , {0.0111, -0.0341, 0.0222, 0.0232, 0.0305, 0.0510, 0.0783, -0.1566, 0.0837, 0.0216}},
  { 215 , {0.0507, -0.1560, 0.1013, 0.1060, 0.1391, 0.2328, 0.3576, 0.7154, -0.1333, 0.0988}},
  { 216 , {0.0081, -0.0248, 0.0161, 0.0168, 0.0221, 0.0370, 0.0568, 0.1136, 0.0607, 0.0157}},
  { 217 , {0.0058, 0.1466, 0.0116, 0.0121, 0.0159, -0.0394, 0.0000, 0.0000, -0.0152, 0.0113}},
  { 218 , {0.0042, 0.1076, 0.0085, 0.0089, 0.0117, -0.0290, 0.0000, 0.0000, 0.0320, 0.0083}},
  { 219 , {0.0050, 0.1273, 0.0100, 0.0105, 0.0138, 0.0231, -0.0661, 0.0000, -0.0132, 0.0098}},
  { 220 , {0.0029, 0.0724, 0.0057, 0.0060, 0.0078, 0.0131, -0.0376, 0.0000, 0.0215, 0.0056}},
  { 221 , {0.0064, 0.1628, 0.0128, 0.0134, 0.0176, 0.0295, 0.0453, -0.0906, -0.0169, 0.0125}},
  { 222 , {0.0030, 0.0763, 0.0060, 0.0063, 0.0083, 0.0138, 0.0212, -0.0424, 0.0227, 0.0059}},
  { 223 , {0.0089, 0.2249, 0.0177, 0.0185, 0.0243, 0.0408, 0.0626, 0.1252, -0.0233, 0.0173}},
  { 224 , {0.0032, 0.0821, 0.0065, 0.0068, 0.0089, 0.0149, 0.0229, 0.0457, 0.0244, 0.0063}}
};
```

