

The status of the SANC project for polarized electron-positron beams

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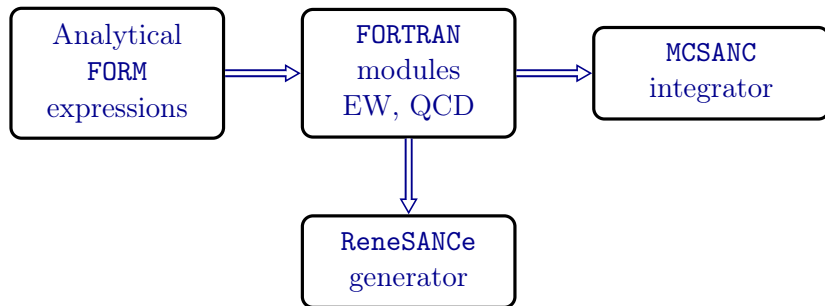
4rd FCC Physics and Experiments Workshop
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Outline

- Review of SANC products for e^+e^-
- Results for implemented processes at NLO EW
- Higher order improvements

The SANC/ARIEl framework and products family



Publications:

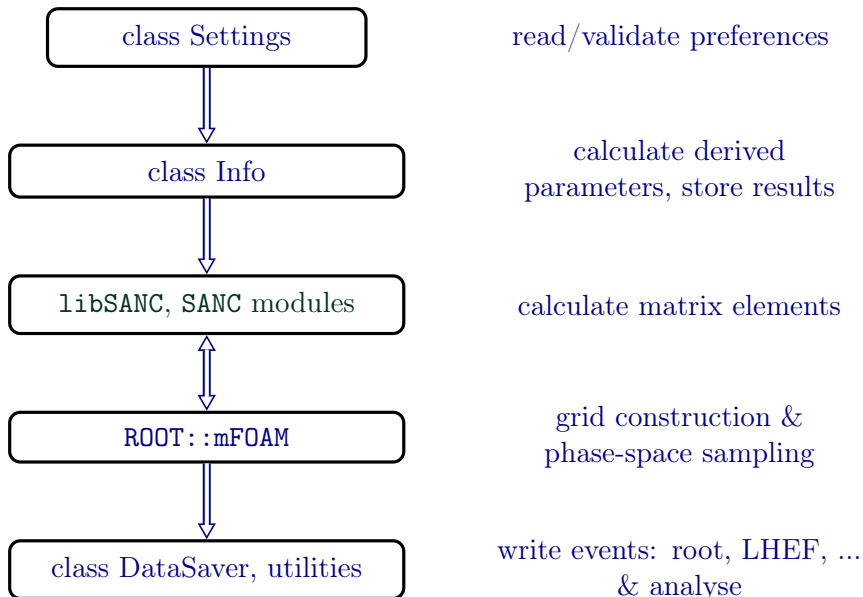
SANC – Comput.Phys.Commun. 174 (2006), 481-517.

MCSANC (pp-mode) – Comput.Phys.Commun. 184 (2013), 2343-2350;
JETP Letters 103 (2016), 131-136.

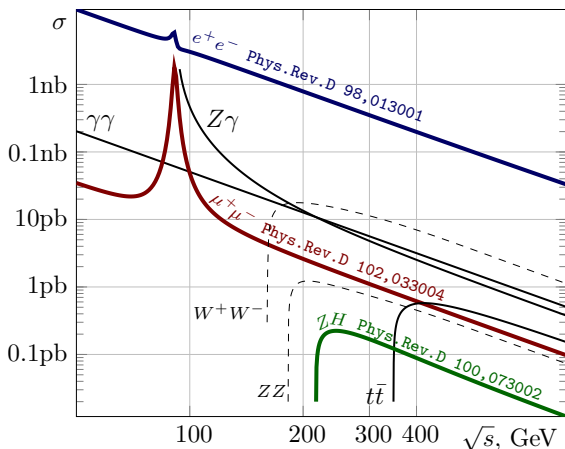
ReneSANCe – Comput.Phys.Commun. 256 (2020), 107445.

SANC products are available at <http://sanc.jinr.ru/download.php>.

Scheme of ReneSANCe event generator structure



Basic processes of SM for e^+e^- annihilation



The cross sections are given for polar angles between $10^\circ < \theta < 170^\circ$ in the final state.

Implemented processes

In the current version of the generator and the integrator for polarized e^+e^- physics implemented:

- Bhabha ($e^+e^- \rightarrow e^-e^+$), *Phys. Rev. D* 98 (2018), 013001.
- ZH ($e^+e^- \rightarrow ZH$), *Phys. Rev. D* 100 (2019), 073002.
- s-channel ($e^+e^- \rightarrow \mu^-\mu^+$, $e^+e^- \rightarrow \tau^-\tau^+$), *Phys. Rev. D* 102 (2020), 033004.

Based on the SANC (Support for Analytic and Numeric Calculations for experiments at colliders) modules the SANC tools take into account complete one-loop and some higher-order electroweak radiative corrections (RC), all the particle masses and polarizations.

Cross section structure

The cross section of processes at one-loop can be divided into:

$$\sigma^{1\text{-loop}} = \sigma^{\text{Born}} + \sigma^{\text{virt}}(\lambda) + \sigma^{\text{soft}}(\lambda, \omega) + \sigma^{\text{hard}}(\omega),$$

Contributions due to:

σ^{Born} — Born level cross section,

σ^{virt} — virtual(loop) corrections,

σ^{soft} — soft photon Bremsstrahlung,

σ^{hard} — hard photon Bremsstrahlung (with energy $E_\gamma > \omega$).

Auxiliary parameters λ ("photon mass") and ω cancel out after summation.

$$\delta = (\sigma^{1\text{-loop}} / \sigma^{\text{Born}} - 1) \cdot 100\%.$$

We use the helicity approach for all contributions.

Results for $e^+e^- \rightarrow ZH(\gamma)$

P_{e^-}, P_{e^+}	0, 0	-0.8, 0	-0.8, -0.6	-0.8, 0.6
$\sqrt{s} = 250 \text{ GeV}$				
$\sigma_{e^+e^-}^{\text{Born}}, \text{ pb}$	225.59(1)	266.05(1)	127.42(1)	404.69(5)
$\sigma_{e^+e^-}^{1\text{-loop}}, \text{ pb}$	206.77(1)	223.33(2)	111.67(2)	334.99(1)
$\delta, \%$	-8.3(1)%	-16.1(1)%	-12.4(1)%	-17.2(1)%
$\sqrt{s} = 500 \text{ GeV}$				
$\sigma_{e^+e^-}^{\text{Born}}, \text{ pb}$	53.74(1)	63.38(1)	30.35(1)	96.40(1)
$\sigma_{e^+e^-}^{1\text{-loop}}, \text{ pb}$	62.42(1)	68.31(1)	34.04(1)	102.58(1)
$\delta, \%$	16.7(1)%	7.8(1)%	12.1(1)%	6.4(1)%
$\sqrt{s} = 1000 \text{ GeV}$				
$\sigma_{e^+e^-}^{\text{Born}}, \text{ pb}$	12.05(1)	14.217(1)	6.809(1)	21.62(1)
$\sigma_{e^+e^-}^{1\text{-loop}}, \text{ pb}$	14.56(1)	15.80(1)	7.95(1)	23.66(1)
$\delta, \%$	20.8(1)%	11.1(1)%	16.7(1)%	9.4(1)%

Planned improvements for Higgs-strahlung

- Implementation of the final particle decay.
- Calculation of complete NLO EW matrix elements for $e^+e^- \rightarrow H f \bar{f}$ process.

Results for $e^+e^- \rightarrow e^+e^-(\gamma)$ (Bhabha)

P_{e^-}, P_{e^+}	0, 0	-0.8, 0	-0.8, -0.6	-0.8, 0.6
$\sqrt{s} = 250$ GeV				
$\sigma_{e^+e^-}^{\text{Born}}$, pb	56.6763(1)	57.7738(1)	56.2725(4)	59.2753(5)
$\sigma_{e^+e^-}^{1\text{-loop}}$, pb	61.731(6)	62.587(6)	61.878(6)	63.287(7)
δ , %	8.92(1)%	8.33(1)%	9.96(1)%	6.77(1)%
$\sqrt{s} = 500$ GeV				
$\sigma_{e^+e^-}^{\text{Born}}$, pb	14.3789(1)	15.0305(1)	12.7061(1)	17.3550(2)
$\sigma_{e^+e^-}^{1\text{-loop}}$, pb	15.465(2)	15.870(2)	13.861(1)	17.884(2)
δ , %	7.56(1)%	5.59(1)%	9.09(1)%	3.05(1)%
$\sqrt{s} = 1000$ GeV				
$\sigma_{e^+e^-}^{\text{Born}}$, pb	3.67921(1)	3.90568(1)	3.03577(3)	4.77562(5)
$\sigma_{e^+e^-}^{1\text{-loop}}$, pb	3.8637(4)	3.9445(4)	3.2332(3)	4.6542(7)
δ , %	5.02(1)%	0.99(1)%	6.50(1)%	-2.54(1)%

cuts are: $|\cos\theta_{e^-}| < 0.9$, $|\cos\theta_{e^+}| < 0.9$.

Planned 2-loop improvements

Implementation of the 2-loop (NNLO) QED results on RC to Bhabha and annihilation cross sections and comparison with existing works:

- J. Blümlein et al., “The $O(\alpha^2)$ initial state QED corrections to $e^+e^- \rightarrow \gamma^*/Z_0^*$ ”, Nucl. Phys. B 956 (2020), 115055
- F. Campanario et al., “Complete QED NLO contributions to the reaction $e^+e^- \rightarrow \mu^+\mu^-\gamma \dots$ ”, JHEP 02 (2014), 114
- C.M. Carloni Calame et al., “NNLO massive corrections to Bhabha scattering and theoretical precision of BabaYaga@NLO”, Nucl. Phys. B Proc. Suppl. 225-227 (2012), 293
- C. Carloni Calame et al., “NNLO leptonic and hadronic corrections to Bhabha scattering ...”, JHEP 07 (2011), 126
- S. Actis et al., “Virtual hadronic and leptonic contributions to Bhabha scattering”, Phys. Rev. Lett. 100 (2008), 131602
- A. Penin and G. Ryan, “Two-loop electroweak corrections to high energy large-angle Bhabha scattering”, JHEP 11 (2011), 081

$e^+e^- \rightarrow l^+l^-(\gamma)$: HA for Born and Virtual parts

$$\mathcal{H}_{-++-} = -e^2 c_+ \left(Q_e Q_l \mathcal{F}_\gamma(s) + \chi_Z^s \delta_e \left[\beta^- I_l^{(3)} \mathcal{F}_{QL} + \delta_l \mathcal{F}_{QQ} \right] \right),$$

$$\mathcal{H}_{-+ \pm \pm} = e^2 \frac{2m_l}{\sqrt{s}} \sin \theta \left(Q_e Q_l \mathcal{F}_\gamma(s) + \chi_Z^s \delta_e \left[I_l^{(3)} \mathcal{F}_{QL} + \delta_l \mathcal{F}_{QQ} + \frac{s}{2} \beta_l^2 I_l^{(3)} \mathcal{F}_{\mathbf{QD}} \right] \right),$$

$$\mathcal{H}_{+- \pm \pm} = -e^2 \frac{2m_l}{\sqrt{s}} \sin \theta \left(Q_e Q_l \mathcal{F}_\gamma(s) + \chi_Z^s \left[2I_e^{(3)} I_l^{(3)} \mathcal{F}_{LL} + 2I_e^{(3)} \delta_l \mathcal{F}_{LQ} + \delta_e I_l^{(3)} \mathcal{F}_{QL} + \delta_e \delta_l \mathcal{F}_{QQ} + \frac{s}{2} \beta_l^2 I_l^{(3)} \left(2I_e^{(3)} \mathcal{F}_{LD} + \delta_e \mathcal{F}_{\mathbf{QD}} \right) \right] \right),$$

$$\mathcal{H}_{+---} = -e^2 c_+ \left(Q_e Q_l \mathcal{F}_\gamma(s) + \chi_Z^s \left[\beta^+ I_l^{(3)} \left(2I_e^{(3)} \mathcal{F}_{LL} + \delta_e \mathcal{F}_{QL} \right) + \delta_l \left(2I_e^{(3)} \mathcal{F}_{LQ} + \delta_e \mathcal{F}_{QQ} \right) \right] \right),$$

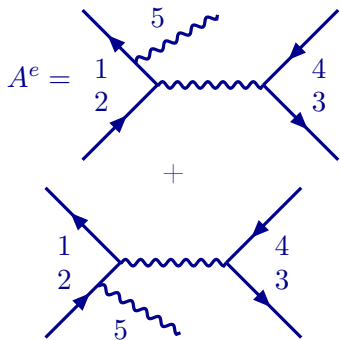
$$\mathcal{H}_{-+--+} = \mathcal{H}_{-++-}(c_+ \rightarrow c_-, \beta^- \rightarrow \beta^+),$$

$$\mathcal{H}_{+-+--} = \mathcal{H}_{+---}(c_+ \rightarrow c_-, \beta^- \rightarrow \beta^+)$$

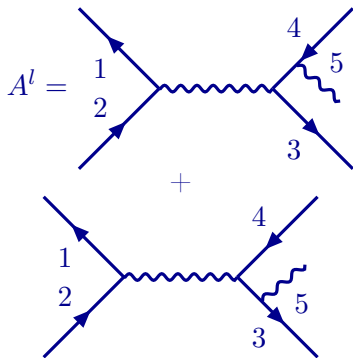
Gauge-invariant diagram sets

$$A = A^e + A^l$$

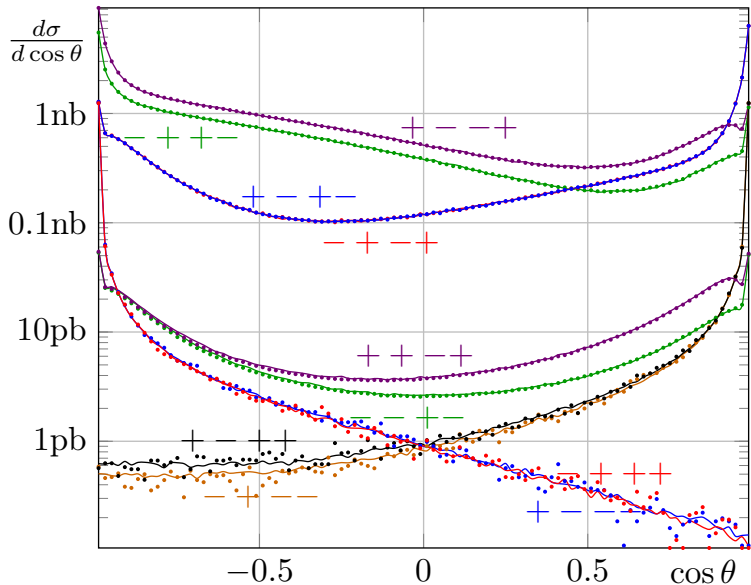
ISR

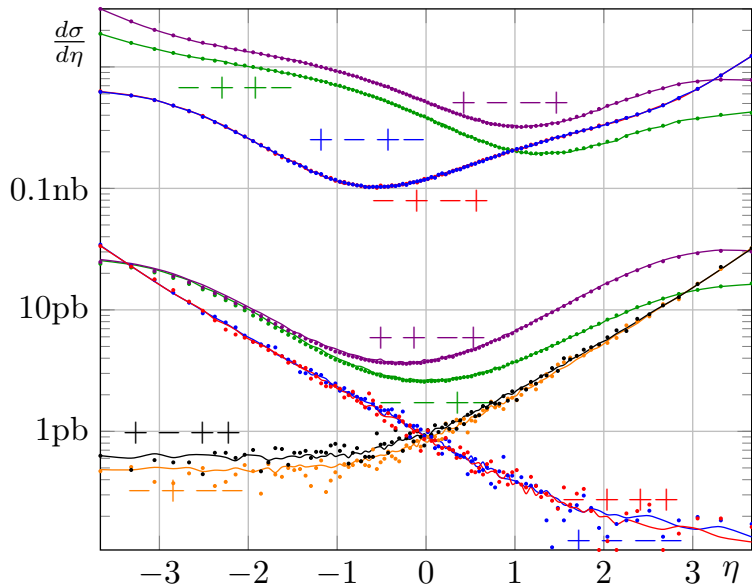


FSR



$$A^l_{\chi_1 \chi_2 \chi_3 \chi_4 \chi_5}(p_1, p_2, p_3, p_4, p_5) = A^e_{\chi_4 \chi_3 \chi_2 \chi_1 \chi_5}(p_4, p_3, p_2, p_1, p_5) |_{m_e \leftrightarrow m_l}$$

SANC vs. WHIZARD (dots): all-polarized $e^+e^- \rightarrow \tau^+\tau^-\gamma$ 

SANC vs. WHIZARD (dots): all-polarized $e^+e^- \rightarrow \tau^+\tau^-\gamma$ 

$e^+e^- \rightarrow \mu^+\mu^-(\gamma)$: NLO EW results, no cuts

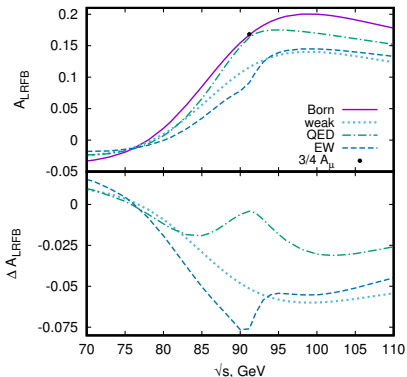
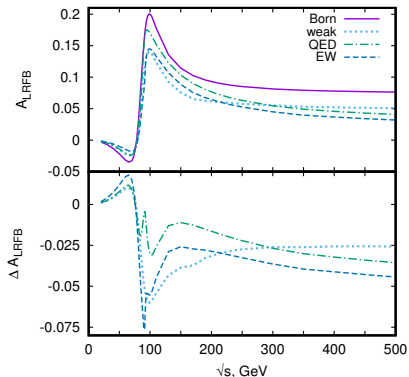
\sqrt{s}	P_{e^-}, P_{e^+}	0, 0	-0.8, 0	-0.8, 0.6	-0.8, -0.6
250 GeV	$\sigma_{\mu^+\mu^-}^{\text{born}}$, fb	1653.7(1)	1804.0(1)	2710.5(1)	897.5(1)
	$\sigma_{\mu^+\mu^-}^{\text{ew}}$, fb	4526.3(2)	4915.2(2)	7298.3(4)	2532.0(1)
	$\delta, \%$	173.7(1)%	172.4(1)%	169.3(1)%	182.1(1)%
500 GeV	$\sigma_{\mu^+\mu^-}^{\text{born}}$, fb	400.85(1)	433.51(1)	650.41(1)	216.61(1)
	$\sigma_{\mu^+\mu^-}^{\text{ew}}$, fb	1138.9(1)	1227.7(1)	1818.2(1)	637.2(1)
	$\delta, \%$	184.1(1)%	183.2(1)%	179.5(1)%	194.2(1)%
1000 GeV	$\sigma_{\mu^+\mu^-}^{\text{born}}$, fb	99.57(1)	107.47(1)	161.20(1)	53.75(1)
	$\sigma_{\mu^+\mu^-}^{\text{ew}}$, fb	296.70(2)	318.74(3)	471.61(4)	165.87(1)
	$\delta, \%$	198.0(1)%	196.6(1)%	192.6(1)%	208.6(1)%

$e^+e^- \rightarrow \mu^+\mu^-(\gamma)$: NLO EW results with cuts

\sqrt{s}	P_{e^-}, P_{e^+}	0, 0	-0.8, 0	-0.8, 0.6	-0.8, -0.6
250 GeV	$\sigma_{\mu^+\mu^-}^{\text{born}}$, fb	1417.6(1)	1546.5(1)	2323.5(1)	769.37(2)
	$\sigma_{\mu^+\mu^-}^{\text{ew}}$, fb	2399(1)	2614(1)	3909(1)	1318(1)
	$\delta, \%$	69.2(1)%	69.0(1)%	68.2(1) %	71.3(1)%
500 GeV	$\sigma_{\mu^+\mu^-}^{\text{born}}$, fb	343.63(1)	371.62(1)	557.56(1)	185.69(1)
	$\sigma_{\mu^+\mu^-}^{\text{ew}}$, fb	469.8(4)	495.4(5)	739.3(7)	251.5(2)
	$\delta, \%$	36.7(1)%	33.3(1) %	32.6(1) %	35.4(1)%
1000 GeV	$\sigma_{\mu^+\mu^-}^{\text{born}}$, fb	85.355(3)	92.131(5)	138.18(1)	46.079(2)
	$\sigma_{\mu^+\mu^-}^{\text{ew}}$, fb	116.2(1)	121.1(1)	180.3(1)	61.83(2)
	$\delta, \%$	36.2(1) %	31.4(1) %	30.5(1)%	34.2(1)%

cuts are: $|\cos\theta_{\mu^-}| < 0.9$, $|\cos\theta_{\mu^+}| < 0.9$.

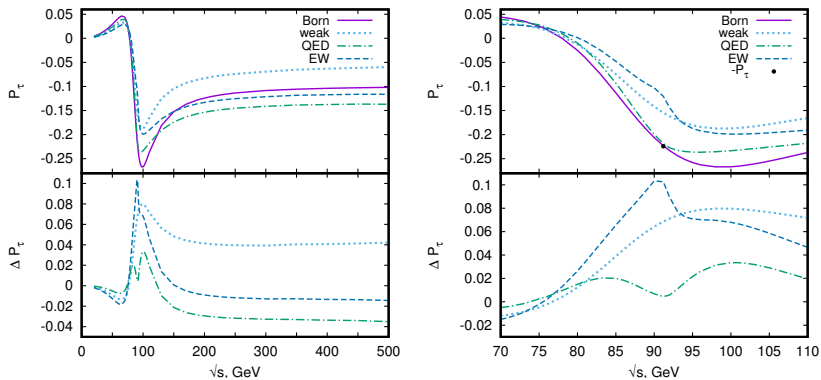
Left–Right Forward–Backward Asymmetry



(**Left**) The A_{LRFB} asymmetry in the Born and 1-loop (weak, QED, EW) approximations and ΔA_{LRFB} for c.m.s. energy range; (**Right**) the same for the Z peak region.

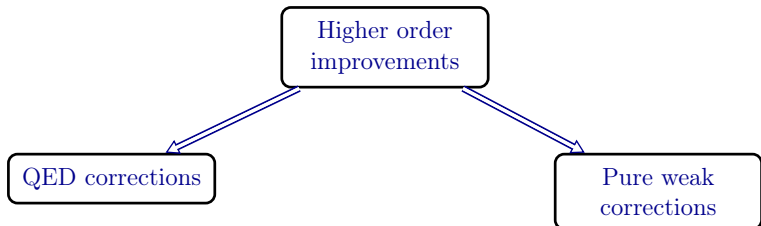
$$A_{LRFB} = \frac{(\sigma_{L_e} - \sigma_{R_e})_F - (\sigma_{L_e} - \sigma_{R_e})_B}{(\sigma_{L_e} + \sigma_{R_e})_F + (\sigma_{L_e} + \sigma_{R_e})_B},$$

Final-State Fermion Polarization



(Left) The P_τ polarization in the Born and 1-loop (weak, pure QED, and EW) approximations and ΔP_τ vs. c.m.s. energy in a wide range; **(Right)** the same for the Z peak region. The black dot indicates the Born value P_τ at the Z resonance.

Higher order improvements



- Leading logarithmic (LL) approximation.
- Corrections to $\Delta\alpha$.
- Shower with matching.
- Corrections to $\Delta\rho$.
- Leading Sudakov logarithms.

Higher order improvements, QED

The leading log in the annihilation channel is $L = \ln \frac{s}{m_l^2}$.

$O(1)$		1		
$O(\alpha)$		αL	α	
$O(\alpha^2)$		$\frac{1}{2}\alpha^2 L^2$	$\frac{1}{2}\alpha^2 L$	$\frac{1}{2}\alpha^2$
$O(\alpha^3)$		$\frac{1}{6}\alpha^3 L^3$	$\frac{1}{6}\alpha^3 L^2$...

In the LL approximation we can separate pure photonic (marked “ γ ”) and the rest corrections which include pure pair and mixed photon-pair effects (marked as “pair”).

ISR corrections in LL approx. for $e^+e^- \rightarrow \mu^+\mu^-(\gamma)$ **PRELIMINARY**

$\sqrt{s} = 250$ GeV, Born cross section $\sigma_0 = 1417.6(1)$ fb.

$\Sigma \equiv \sum_{n=2}^4 \mathcal{O}(\alpha^n L^n)$, $\delta = \delta_{\text{ISR LLA}} \equiv \delta\sigma_{\text{ISR LLA}}/\sigma_0$, $\bar{\delta}$ calculated at scale $= 4s$ and $\underline{\delta}$ at scale $= s/4$.

	$\mathcal{O}(\alpha^2 L^2)$		$\mathcal{O}(\alpha^3 L^3)$		$\mathcal{O}(\alpha^4 L^4)$	Σ
	$[\gamma]$	[pair]	$[\gamma]$	[pair]	$[\gamma]$	
Cuts: $ \cos\theta_{\mu^\pm} < 0.9$, $M_{\mu^+\mu^-} > 10$ GeV.						
$\delta\sigma$, fb	108.2(1)	53.70(1)	-0.49(3)	3.47(1)	-0.23(1)	164.7(1)
δ , %	7.63(1)%	3.79(1)%	-0.035(2)%	0.245(1)%	-0.017(1)%	11.62(1)%
$\bar{\delta}$, %						13.91(1)%
$\underline{\delta}$, %						11.15(1)%
Cuts: $ \cos\theta_{\mu^\pm} < 0.9$, $M_{\mu^+\mu^-} > 100$ GeV.						
$\delta\sigma$, fb	4.8(1)	5.9(1)	-0.76(2)	0.00(1)	0.00(1)	9.9(1)
δ , %	0.34(1)%	0.42(1)%	-0.053(1)%	0.00(1)%	0.00(1)%	0.70(1)%
$\bar{\delta}$, %						0.78(1)%
$\underline{\delta}$, %						0.63(1)%

FSR corrections in LL approx. for $e^+e^- \rightarrow \mu^+\mu^-(\gamma)$

PRELIMINARY

$\sqrt{s} = 250$ GeV, Born cross section $\sigma_0 = 1417.6(1)$ fb.

$\Sigma \equiv \sum_{n=2}^4 \mathcal{O}(\alpha^n L^n)$, $\delta = \delta_{\text{FSR LLA}} \equiv \delta\sigma_{\text{FSR LLA}}/\sigma_0$, $\bar{\delta}$ calculated at scale $= 4s$ and $\underline{\delta}$ at scale $= s/4$.

	$\mathcal{O}(\alpha^2 L^2)$		$\mathcal{O}(\alpha^3 L^3)$		$\mathcal{O}(\alpha^4 L^4)$	Σ
	$[\gamma]$	[pair]	$[\gamma]$	[pair]	$[\gamma]$	
Cuts: $ \cos\theta_{\mu^\pm} < 0.9$, $M_{\mu^+\mu^-} > 10$ GeV.						
$\delta\sigma$, fb	0.00(1)	7.64(1)	0.00(1)	-0.129(1)	0.00(1)	7.50(1)
δ , %	0.00(1)%	0.539(1)%	0.00(1)%	-0.0091(1)%	0.00(1)%	0.529(1)%
$\bar{\delta}$, %						0.613(1)%
$\underline{\delta}$, %						0.450(1)%
Cuts: $ \cos\theta_{\mu^\pm} < 0.9$, $M_{\mu^+\mu^-} > 100$ GeV.						
$\delta\sigma$, fb	-0.54(1)	0.87(1)	0.00(1)	-0.069(1)	0.00(1)	0.26(1)
δ , %	-0.038(1)%	0.061(1)%	0.00(1)%	-0.005(1)%	0.00(1)%	0.018(1)%
$\bar{\delta}$, %						0.019(1)%
$\underline{\delta}$, %						0.017(1)%

Higher order improvements, weak

Higher order improvements added throw $\Delta\rho$ parameter:

$$s_W^2 \rightarrow \bar{s}_W^2 \equiv s_W^2 + \Delta\rho c_W^2.$$

At the two-loop level, the quantity $\Delta\rho$ contains two contributions:

$$\Delta\rho = N_c X_t \left[1 + \rho^{(2)} \left(M_H^2/m_t^2 \right) X_t \right] \left[1 - \frac{2\alpha_s(M_Z^2)}{9\pi} (\pi^2 + 3) \right],$$

where $X_t = \frac{\sqrt{2}G_F m_t^2}{16\pi^2}$.

Higher order improvements, weak for $e^+e^- \rightarrow \mu^+\mu^-$

P_{e^+}, P_{e^-}	0, 0	0,-0.8	0.3,-0.8	0,0.8	-0.3,0.8
$\sigma_{\alpha(0)}^{\text{Born}}, \text{pb}$	1.41763(1)	1.54645(1)	1.93499(1)	1.28880(1)	1.58073(1)
$\sigma_{G_\mu}^{\text{Born}}, \text{pb}$	1.50971(1)	1.64690(1)	2.06068(1)	1.37252(1)	1.68341(1)
$\sigma_{\alpha(M_Z^2)}^{\text{Born}}, \text{pb}$	1.59923(1)	1.74456(1)	2.18287(1)	1.45391(1)	1.78323(1)
$\delta\sigma_{\alpha(0)}^{\text{weak}}, \text{pb}$	0.15525(1)	0.11883(1)	0.14243(1)	0.19167(1)	0.242587(1)
$\delta\sigma_{G_\mu}^{\text{weak}}, \text{pb}$	0.07911(1)	0.03249(1)	0.03400(1)	0.12574(1)	0.162206(1)
$\delta\sigma_{\alpha(M_Z^2)}^{\text{weak}}, \text{pb}$	-0.01194(1)	-0.07003(1)	-0.09468(1)	0.46147(1)	0.06506(1)
$\delta\sigma_{\alpha(0)}^{\text{ho}}, \text{pb}$	0.02122(1)	0.02304(1)	0.02882(1)	0.01940(1)	0.02380(1)
$\delta\sigma_{G_\mu}^{\text{ho}}, \text{pb}$	-0.00555(1)	-0.00351(1)	-0.00407(1)	-0.00759(1)	-0.00969(1)
$\delta\sigma_{\alpha(M_Z^2)}^{\text{ho}}, \text{pb}$	0.00387(1)	0.00898(1)	0.01183(1)	-0.00124(1)	-0.00222(1)

cuts are: $|\cos\theta_{\mu^-}| < 0.9, \quad |\cos\theta_{\mu^+}| < 0.9.$

Impact of all corrections for $e^+e^- \rightarrow \mu^+\mu^-(\gamma)$

	Born	+QED (1-loop)	+WEAK (1-loop)	+WEAK ($\Delta\rho$)	+QED (LL)
σ , pb	1.50971(1)	cut1: + 0.829(1) cut2: + 0.197(1)	+0.07911(1)	-0.00555(1)	cut1: + 0.1837(1) cut2: + 0.0108(1)
δ , %	100%	cut1: + 54.9(1)% cut2: + 13.1(1)%	+5.24(1)%	-0.37(1)%	cut1: + 12.17(1)% cut2: + 0.72(1)%

Calculated in G_μ EW scheme, $\sqrt{s} = 250$ GeV.

Cuts are: $|\cos\theta_{\mu^-}| < 0.9$, $|\cos\theta_{\mu^+}| < 0.9$,

cut1: $M_{ll} > 10$ GeV,

cut2: $M_{ll} > 100$ GeV.

RESUME: SANC

- Monte Carlo integrator MCSANcEE and event generator ReneSANCe are being improved
 - Own library of complete one-loop EW corrections
 - Initial & final state polarization
 - LL-accuracy QED improvements to cross section
 - Higher order WEAK improvements throw $\Delta\rho$
 - PROCESSES, **DONE**: e^+e^- , $\mu^+\mu^-$, $\tau^+\tau^-$, ZH
 - PROCESSES, **UNDERWAY**: $Z\gamma$, $\gamma\gamma$, $t\bar{t}$.
Next version of SANC.
 - **UNDERWAY**: resonance approximation, showers, Sudakov logarithms.

Numerical results: Setup for tuned comparison

We performed a tuned comparison of our results for polarized Born and hard Bremsstrahlung with the results `WHIZARD` and `CalcHEP` programs.

Initial parameters

$$\begin{aligned}
 \alpha^{-1}(0) &= 137.03599976, & M_W &= 80.451495 \text{ GeV}, & \Gamma_W &= 2.0836 \text{ GeV}, \\
 M_H &= 125.0 \text{ GeV}, & M_Z &= 91.1867 \text{ GeV}, & \Gamma_Z &= 2.49977 \text{ GeV}, \\
 m_e &= 0.5109990 \text{ MeV}, & m_\mu &= 0.105658 \text{ GeV}, & m_\tau &= 1.77705 \text{ GeV}, \\
 m_d &= 0.083 \text{ GeV}, & m_s &= 0.215 \text{ GeV}, & m_b &= 4.7 \text{ GeV}, \\
 m_u &= 0.062 \text{ GeV}, & m_c &= 1.5 \text{ GeV}, & m_t &= 173.8 \text{ GeV}.
 \end{aligned}$$

with cuts $|\cos\theta| < 0.9$, $E_\gamma > 1 \text{ GeV}$

`WHIZARD` and `CalcHEP`

- W. Kilian, T. Ohl, J. Reuter, *Eur.Phys.J.C*71 (2011) 1742,
- A.Belyaev, N.Christensen,A.Pukhov, *Comp. Phys. Comm.* 184 (2013), pp. 1729-1769