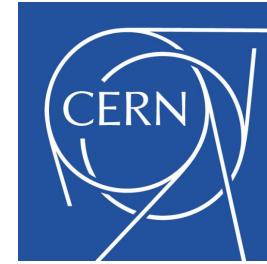


What can off-shell Higgs tell us?

Ennio Salvioni, CERN

Christophe Grojean, DESY/Humboldt

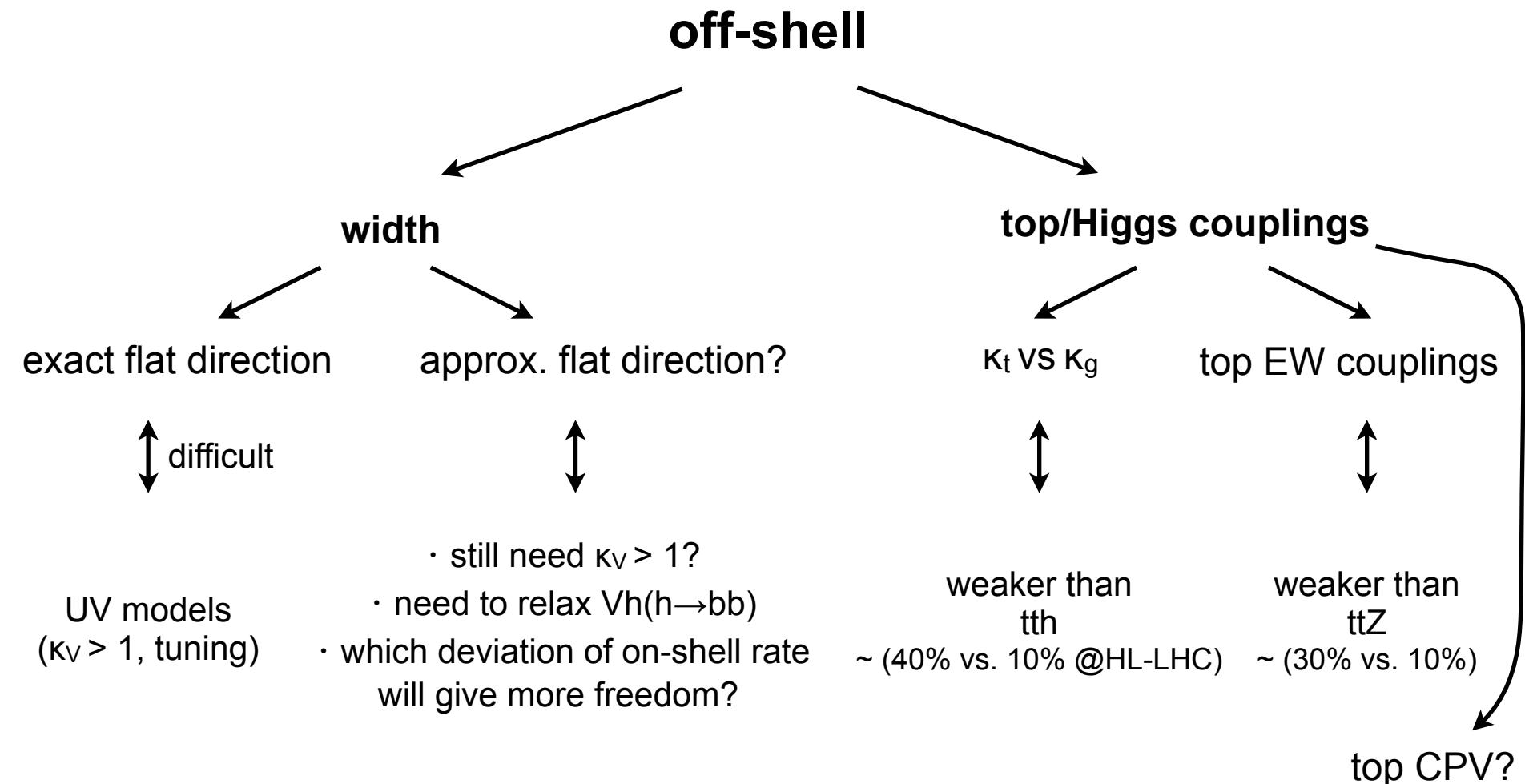
Alex Azatov, SISSA



SISSA
■ ■ ■

Off-shell interpretations task force
3rd meeting, July 8, 2020

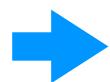
What can off-shell Higgs production tell us?



Flat direction in on-shell Higgs measurements

$$\sigma(XX \rightarrow h) \text{BR}(h \rightarrow YY) \propto g_{hXX}^2 \frac{g_{hYY}^2}{\Gamma_h}$$

$$g_{hii} = \kappa_{\text{univ}} g_{hii}^{\text{SM}}$$



all on-shell rates unchanged

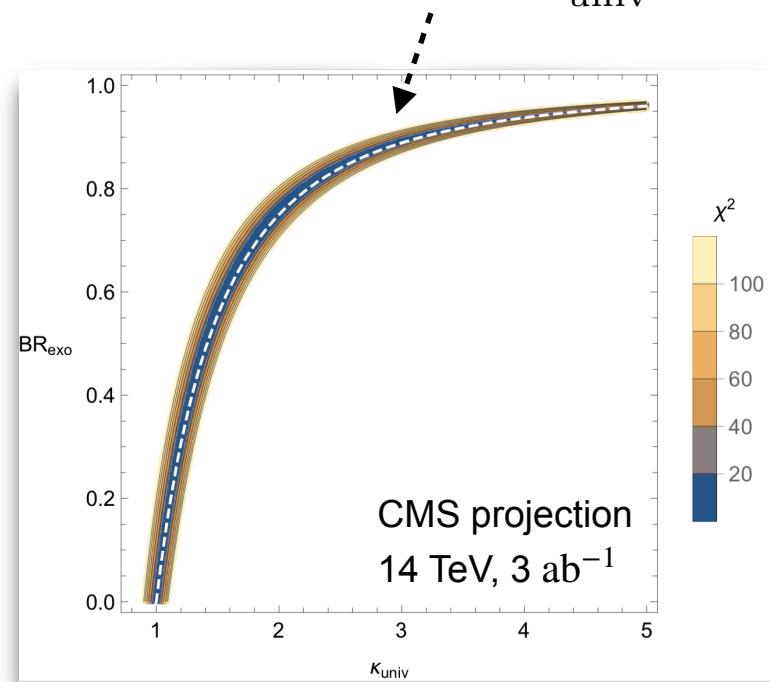
$$\Gamma_h = \kappa_{\text{univ}}^4 \Gamma_h^{\text{SM}}$$

$$\Gamma_h = \kappa_{\text{univ}}^2 \Gamma_h^{\text{SM}} + \Gamma_h^{\text{exo}} \quad \rightarrow \quad \text{BR}_{\text{exo}} = \frac{\kappa_{\text{univ}}^2 - 1}{\kappa_{\text{univ}}^2}$$

untagged BR gives most leeway,
invisible BR already strongly constrained

$\text{BR}(\text{Higgs} \rightarrow \text{invisible}) < 13\% \text{ (95\% CL)}$

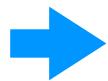
[ATLAS-CONF-2020-008]



Flat direction vs off-shell Higgs

$$\sigma(XX \rightarrow h) \text{BR}(h \rightarrow YY) \propto g_{hXX}^2 \frac{g_{hYY}^2}{\Gamma_h}$$

$$g_{hii} = \kappa_{\text{univ}} g_{hii}^{\text{SM}}$$



all on-shell rates unchanged

$$\Gamma_h = \kappa_{\text{univ}}^4 \Gamma_h^{\text{SM}}$$

$$\Gamma_h = \kappa_{\text{univ}}^2 \Gamma_h^{\text{SM}} + \Gamma_h^{\text{exo}} \quad \rightarrow \quad \text{BR}_{\text{exo}} = \frac{\kappa_{\text{univ}}^2 - 1}{\kappa_{\text{univ}}^2}$$

Going to off-shell region:

$$\frac{d\sigma_{gg \rightarrow h^* \rightarrow ZZ}}{ds} \propto \frac{g_{hgg}^2 g_{hZZ}^2}{(s - m_h^2)^2 + \Gamma_h^2 m_h^2} \sim \frac{g_{hgg}^2 g_{hZZ}^2}{s^2} \quad \text{which grows as } \kappa_{\text{univ}}^4$$



upper bound on off-shell cross section translates to upper bound on width

[Caola and Melnikov 1307.4935]

Universal flat direction

How to realize a **universal** flat direction? Toy setup:

$$\mathcal{L}_{\text{BSM}} = \frac{c_H}{f^2} (\partial_\mu |H|^2)^2 + \lambda_{H\phi} |H|^2 \phi^2$$

$$\frac{g_{hii}}{g_{hii}^{\text{SM}}} = 1 - c_H \frac{v^2}{f^2}$$

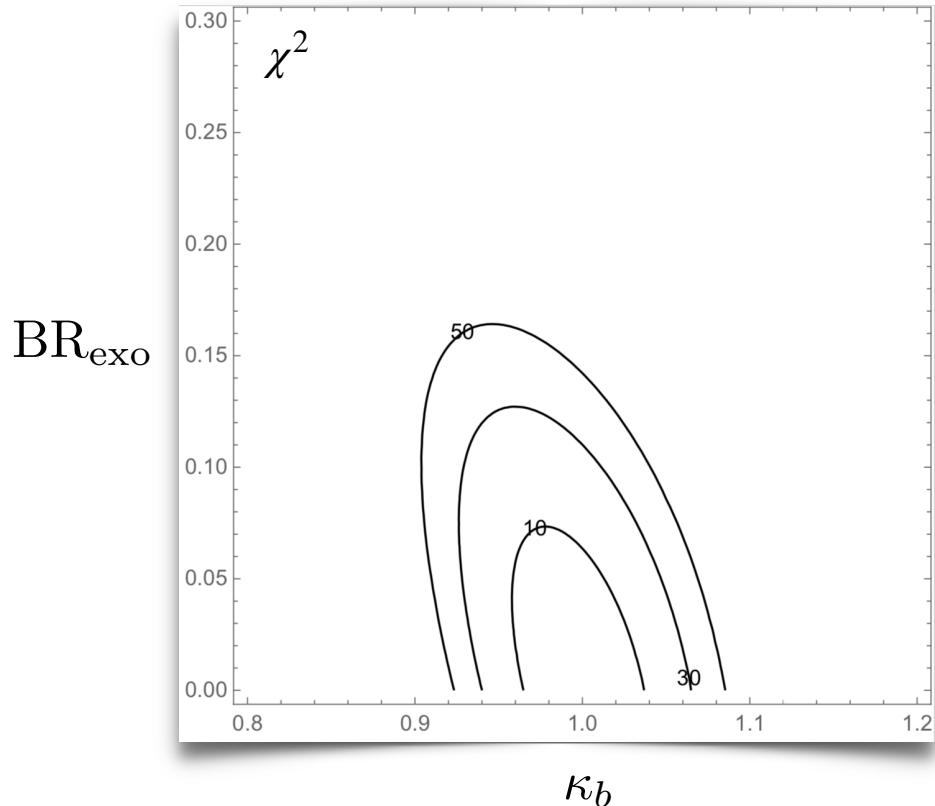
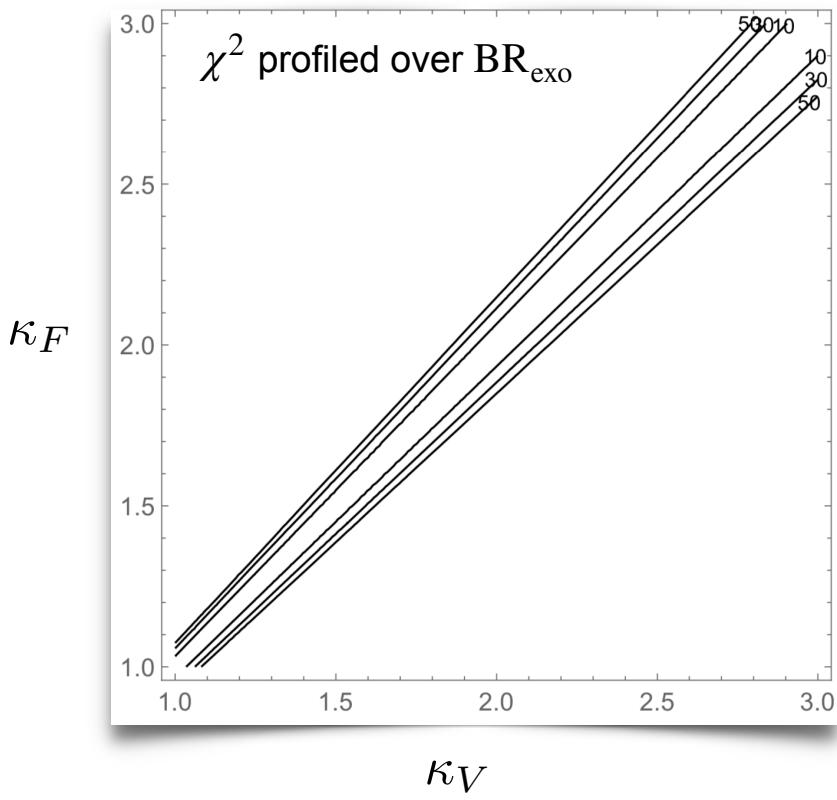
$$\Gamma(h \rightarrow \phi\phi) \sim \frac{\lambda_{H\phi}^2 v^2}{8\pi m_h}$$



Requires:

- $c_H < 0$, possible but somewhat exotic. for example [Georgi and Machacek 1985]
E.g., pNGB of non-compact coset, or models with larger scalar reps.
- ϕ decays to hadrons, for example color octet or singlet $\rightarrow gg$
- A “conspiracy” relating c_H and $\lambda_{H\phi}$

Relaxing coupling universality?



Flat direction is lifted **very rapidly** when we depart
from universal rescaling of Higgs couplings



Preliminary - in progress

Beyond the Higgs width: (SM)EFT

- $gg \rightarrow ZZ$ [Higgs-mediated + continuum] is sensitive to Higgs/top operators.

In Higgs basis:

9 CP-even coefficients

$$\{c_{gg}, (\delta y_u)_{33}, \delta c_z, c_{zz}, c_{z\square}, (\delta g_L^{Zu})_{33}, (\delta g_R^{Zu})_{33}, \text{Re}(d_{Zu})_{33}, \text{Re}(d_{Gu})_{33}\}$$

5 CP-odd coefficients

$$\{\tilde{c}_{gg}, (\phi_u)_{33}, \tilde{c}_{zz}, \text{Im}(d_{Zu})_{33}, \text{Im}(d_{Gu})_{33}\}$$

$$\begin{aligned} \mathcal{L} = & \frac{h}{v} \left[\tilde{c}_{gg} \frac{g_s^2}{4} G_{\mu\nu}^a \tilde{G}_{\mu\nu}^a - m_t (\delta y_u)_{33} e^{i(\phi_u)_{33}} \bar{t}_R t_L + \text{h.c.} + \delta c_z \frac{g_Z^2 v^2}{4} Z_\mu Z^\mu + c_{zz} \frac{g_Z^2}{4} Z_{\mu\nu} Z^{\mu\nu} + c_{z\square} g^2 Z_\mu \partial_\nu Z^{\mu\nu} + \tilde{c}_{gg} \frac{g_s^2}{4} G_{\mu\nu}^a \tilde{G}_{\mu\nu}^a + \tilde{c}_{zz} \frac{g_Z^2}{4} Z_{\mu\nu} \tilde{Z}^{\mu\nu} \right] \\ & + g_Z (\delta g_L^{Zu})_{33} Z_\mu \bar{t}_L \gamma^\mu t_L + g_Z (\delta g_R^{Zu})_{33} Z_\mu \bar{t}_R \gamma^\mu t_R - \frac{m_t}{4v^2} \left[\underline{g_s \bar{t}_L \sigma^{\mu\nu} T^a} \underline{(d_{Gu})_{33} t_R G^{\mu\nu a}} + g_Z \bar{t}_L \sigma^{\mu\nu} T^a \underline{(d_{Zu})_{33} t_R Z^{\mu\nu}} \right] + \text{h.c.} \end{aligned}$$

Beyond the Higgs width: (SM)EFT

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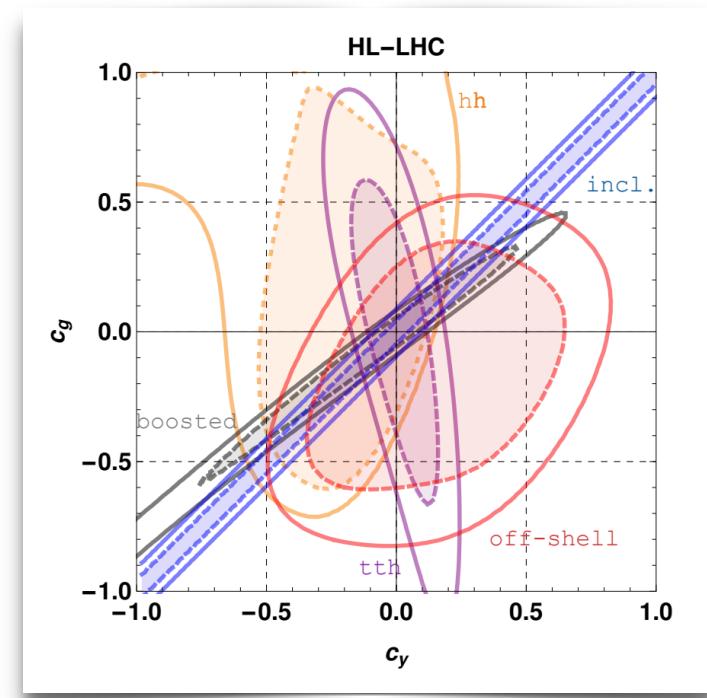
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$$\frac{\sigma_{gg \rightarrow h}}{\sigma_{gg \rightarrow h}^{\text{SM}}} \sim [1 + c_{gg} - (\delta y_u)_{33}]^2$$

blind direction in on-shell gluon fusion production

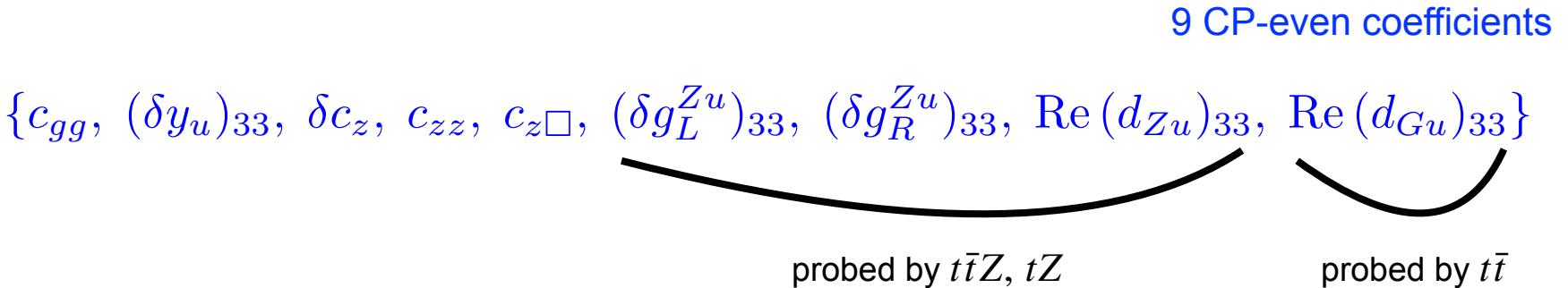


[Azatov, Grojean, Paul, Salvioni 1608.00977]

Beyond the Higgs width: (SM)EFT

- $gg \rightarrow ZZ$ [Higgs-mediated + continuum] is sensitive to Higgs/top operators.

In Higgs basis:



68% CL bounds on $\frac{\Lambda}{\sqrt{C_i}} [\text{TeV}]$

$C_{\phi Q}^- : [-0.50, 0.37]$
$C_{\phi t} : [-0.24, 0.55]$
$C_{tZ} : [-0.66, 0.66]$
$C_{tG} : [1.2, 1.8]$

current bound on $C_{tG} \leftrightarrow \text{Re}(d_{Gu})_{33}$ is significantly stronger than for others

[Brivio, Bruggisser, Maltoni, Moutafis, Plehn, Vryonidou, Westhoff, Zhang 1910.03606]

Beyond the Higgs width: (SM)EFT

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In Higgs basis:

9 CP-even coefficients

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- $gg \rightarrow ZZ$ is mostly sensitive to **axial vector** coupling $\propto (\delta g_L^{Zu})_{33} - (\delta g_R^{Zu})_{33}$

[Azatov, Grojean, Paul, Salvioni 1608.00977]

- Angular distributions can be exploited to isolate **LL** pol. mode, which is enhanced for BSM couplings

[Cao, Yan, Yuan, Zhang 2004.02031]

$$\frac{g}{2c_w} \bar{t} (v_t - a_t \gamma_5) t Z^\mu$$

$$a_t^{\text{SM}} = 1/2$$

