

On the Use of the Lossy Transmission Line Theory for the SRF Characterization

Pablo Vidal García



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Outline

The Lossy Transmission Line Theory

- Introduction
- Basis
- Analysis

RF and SRF cavities: a different approach

- Introduction
- The high-losses case
- The Lossy TLT connection

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Limitations. Proposals. Conclusions.

The Lossy Transmission Line Theory :

Introduction



Transmission Line Theory (TLT):

Objective: Complete characterization of **propagative waves** in **PEC** surrounded **waveguides** by means of **TL parameters**.



Mathematically: characterization of the **closed-form homogeneous solutions** to \square^2 when “**ideal**” **BCs** (Dirichlet & Neumann) and x_i - **invariance** are imposed.

Lossy TLT (LTLT):

Generalization: The lossless case is only a particular case, among others.



Mathematically: nonhomogeneous solutions to \square^2 with general BCs.

The Lossy Transmission Line Theory : *Introduction*

LTLT:

- TEM modes, i.e. plane waves \leftrightarrow analytic solutions in multiple connected planar (e.g. z -invariant) regions.

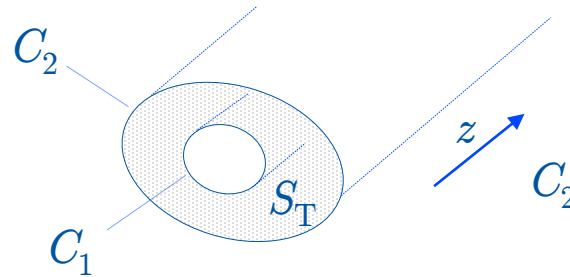


Fig. 1. Two-conductor waveguide that supports the TEM mode represented in some cylindrical coordinate system.

- Easily described in the frequency domain \leftrightarrow time (harmonic) \mathcal{FT} (coordinates on $e^{j\omega t}$ basis).
- Complex functions and parameters \rightarrow Complex Analysis.

The Lossy Transmission Line Theory : *Basis*

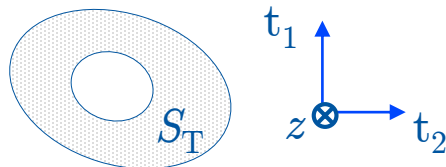


From the analyticity of TEM solutions in S_T , the integration of the *Maxwell equations* unequivocally leads to the *Telegrapher's equations* (written in the frequency domain):

Line parameters

$$\begin{cases} \frac{dV(z)}{dz} = - (R + j\omega L) \cdot I(z) ; & R = \omega\mu'' \cdot H_{12}/\ell_2 , L = \mu' \cdot H_{12}/\ell_2 \in \mathbb{R}^+ \\ \frac{dI(z)}{dz} = - (G + j\omega C) \cdot V(z) ; & G = \omega\varepsilon'' \cdot \ell_2/H_{12} , C = \varepsilon' \cdot \ell_2/H_{12} \in \mathbb{R}^+ \end{cases} ;$$

where, using generalized orthogonal coordinates on S_T : (t_1, t_2) ;
 $H_{12} = \int_{\langle t_1 \rangle} h_1/h_2 h dt_1$, $\ell_2 = \int_{\langle t_2 \rangle} dt_2$, (e.g. polar coord.: $t_1=r$, $t_2=\varphi$).



◀ Fig. 2. General orthogonal coordinate system.

The Lossy Transmission Line Theory : *Basis*



Solutions:

$$V(z) = V_0^+ e^{-kz} + V_0^- e^{+kz} = V_0^+ e^{-kz} (1 + \Gamma_0 e^{-2kz}),$$
$$I(z) = I_0^+ e^{-kz} - I_0^- e^{+kz} = (V_0^+ / Z_0) e^{-kz} (1 - \Gamma_0 e^{-2kz});$$

Basic parameters

$$Z_0 = V_0^+ / I_0^+ = V_0^- / I_0^- = \sqrt{\frac{R + j\omega L}{G + j\omega C}}, \quad k = \sqrt{(R + j\omega L)(G + j\omega C)} = \alpha + j\beta \in \mathbb{C}$$

Wave parameters

$$V(z) / I(z) = Z(z) = 1 / Y(z) = Z_0 \frac{1 + \Gamma(z)}{1 - \Gamma(z)}; \quad \Gamma(z) = \Gamma_0 e^{2kz}$$

The Lossy Transmission Line Theory : Analysis

Basic parameters

(fixed frequency analysis, ω_0)

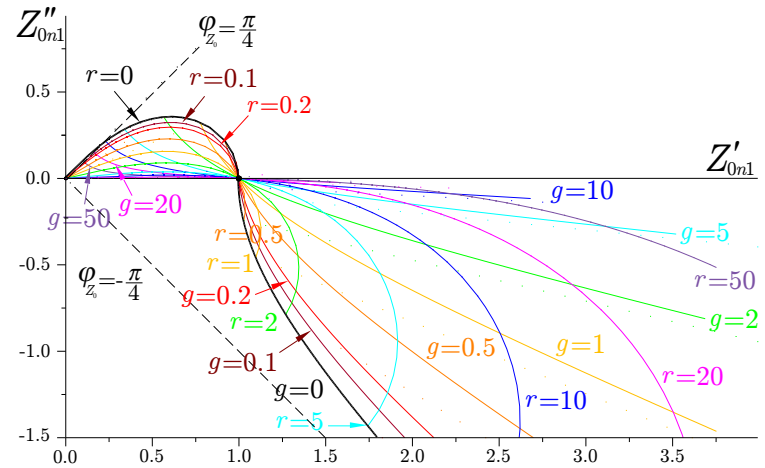
Parameterizations:

$$\begin{cases} r = R/(\omega_0 L) \\ g = G/(\omega_0 C) \end{cases}$$

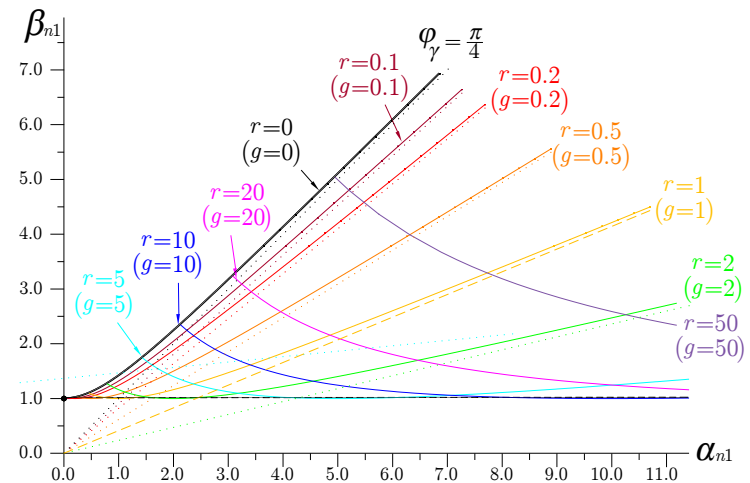
Normalizations :

$$Z_{0n1} = Z_0/Z_{0,lossless} = \sqrt{\frac{1-jr}{1-jg}}$$

$$k_{n1} = k/\beta_{lossless} = j\sqrt{(1-jr)(1-jg)}$$



(a)



(b)

Fig. 3. Normalized (a) characteristic impedance and (b) propagation constant.

The Lossy Transmission Line Theory : Analysis

Wave parameters

Normalizations :

$$Z_{n0} = Z/|Z_0|; Y_{n0} = Y/|Z_0|; \Gamma$$

$$Z_{0n} = Z_0/|Z_0| = e^{j\varphi_{Z_0}}$$

Parameterizations:

$$\left\{ \begin{array}{l} \varphi_{Z_0} \\ (Z'_{n0}, Z''_{n0}) \text{ or } (|Z_{n0}|, \varphi_Z) \end{array} \right. ;$$

$$\left\{ \begin{array}{l} \varphi_{Z_0} \\ (Y'_{n0}, Y''_{n0}) \text{ or } (|Y_{n0}|, \varphi_Y) \end{array} \right. ;$$

$$\left\{ \begin{array}{l} \varphi_{Z_0} \\ (\Gamma', \Gamma'') \text{ or } (|\Gamma|, \varphi_\Gamma) \end{array} \right. .$$

Usual Smith Chart: $\left\{ \begin{array}{l} \varphi_{Z_0} = 0 \\ (Z'_{n0}, Z''_{n0}) \end{array} \right.$

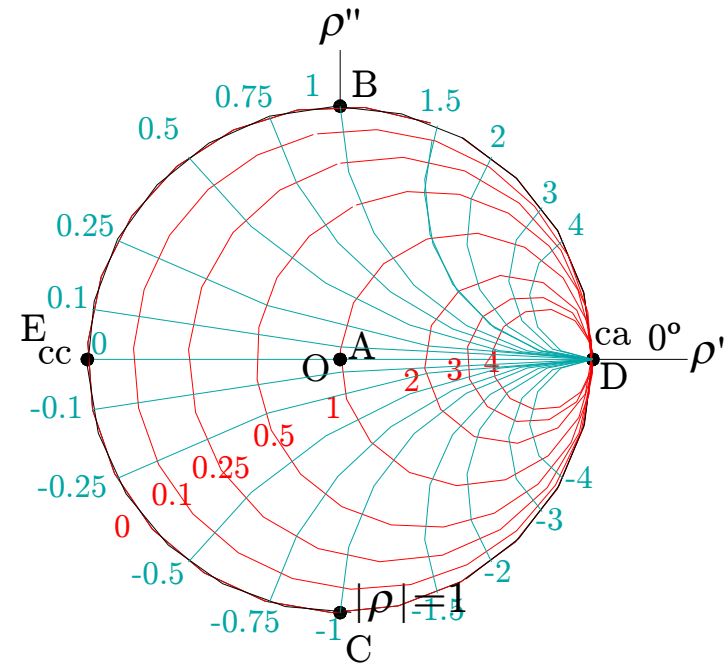


Fig. 4. Transformations of the real and imaginary parts of the Z_{n0} -plane into the Γ -plane.

The Lossy Transmission Line Theory : *Analysis*

Generalized Smith Chart: $\left\{ \begin{array}{l} \varphi_{Z_0} = -25^\circ \\ (Z'_{n0}, Z''_{n0}) \end{array} \right.$

Wave parameters

Normalizations :

$$Z_{n0} = Z/|Z_0|; Y_{n0} = Y/|Z_0|; \Gamma$$

$$Z_{0n} = Z_0/|Z_0| = e^{j\varphi_{Z_0}}$$

Parameterizations:

$$\left\{ \begin{array}{l} \varphi_{Z_0} \\ (Z'_{n0}, Z''_{n0}) \text{ or } (|Z_{n0}|, \varphi_Z) \end{array} \right. ;$$

$$\left\{ \begin{array}{l} \varphi_{Z_0} \\ (Y'_{n0}, Y''_{n0}) \text{ or } (|Y_{n0}|, \varphi_Y) \end{array} \right. ;$$

$$\left\{ \begin{array}{l} \varphi_{Z_0} \\ (\Gamma', \Gamma'') \text{ or } (|\Gamma|, \varphi_\Gamma) \end{array} \right. .$$

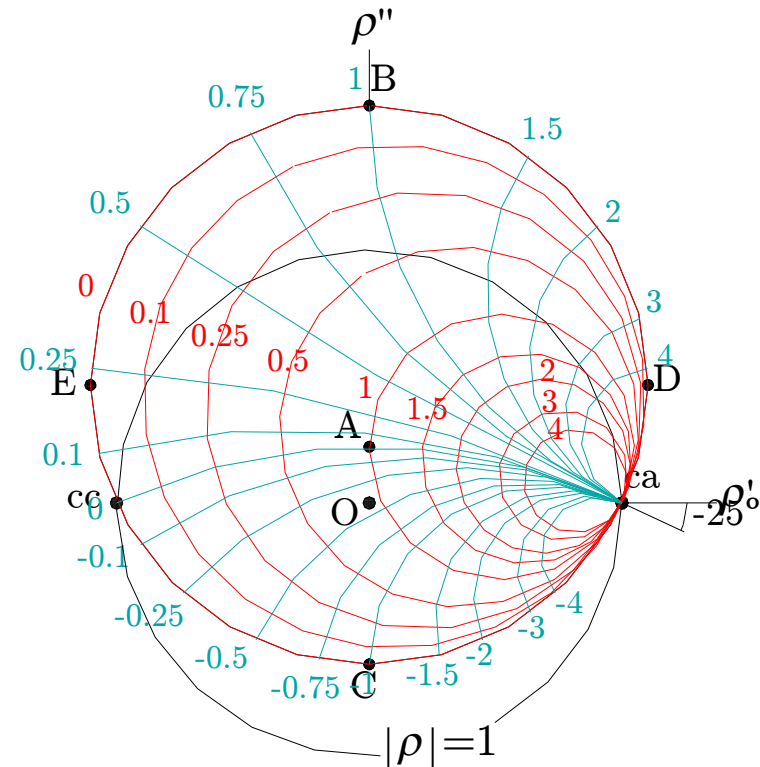


Fig. 4. Transformations of the real and imaginary parts of the Z_{n0} -plane into the Γ -plane.

The Lossy Transmission Line Theory : *Analysis*

Wave parameters

Normalizations :

$$Z_{n0} = Z/|Z_0|; Y_{n0} = Y|Z_0|; \Gamma$$

$$Z_{0n} = Z_0/|Z_0| = e^{j\varphi_{Z_0}}$$

Parameterizations:

$$\left\{ \begin{array}{l} \varphi_{Z_0} \\ (Z'_{n0}, Z''_{n0}) \text{ or } (|Z_{n0}|, \varphi_Z) \end{array} \right. ;$$

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$$\left\{ \begin{array}{l} \varphi_{Z_0} \\ (\Gamma', \Gamma'') \text{ or } (|\Gamma|, \varphi_\Gamma) \end{array} \right. .$$

$$\left\{ \begin{array}{l} \varphi_{Z_0} = -25^\circ \\ (|Z_{n0}|, \varphi_Z) \end{array} \right.$$

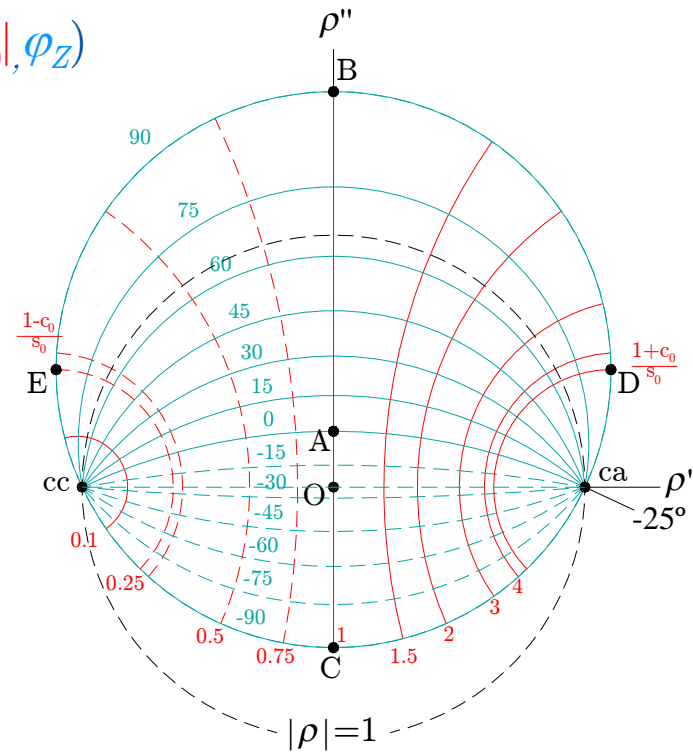


Fig. 5. Transformations of the modulus and phase of the Z_{n0} -plane into the Γ -plane.

The Lossy Transmission Line Theory : *Analysis*

Example of analysis: Wave parameters along the TL

$$\Gamma = \Gamma_L e^{-2kl} \quad (\text{logarithmic spiral in the } \Gamma\text{-plane})$$

$$\Gamma_L = 0.756 \angle -4.118^\circ, \quad \varphi_{Z_0} = 25.526^\circ$$

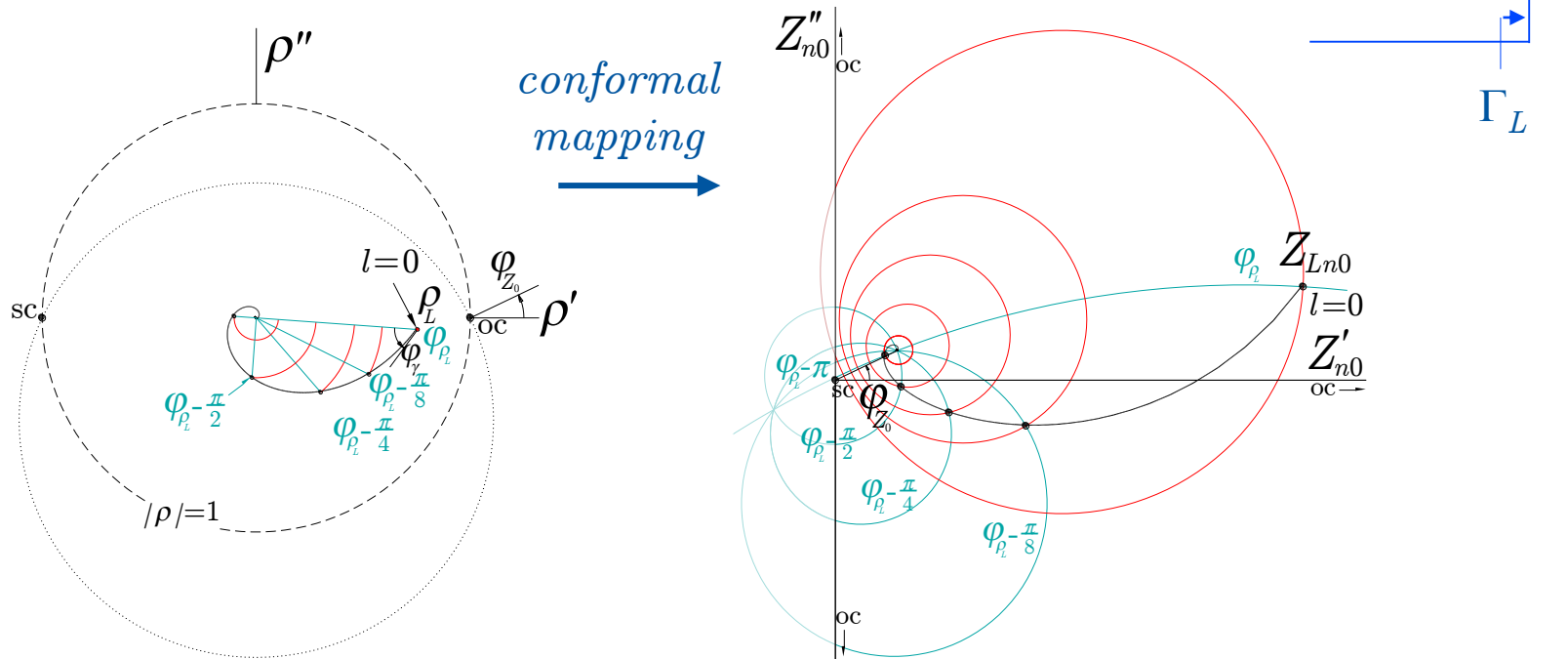
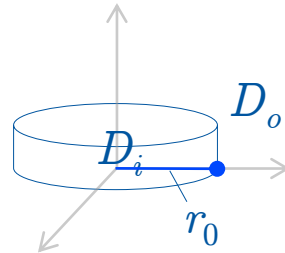


Fig. 6. Parameterization of $\Gamma(l)$ and its transformation to the Z_{n0} -plane.

RF and SRF cavities: a different approach

Introduction

Example: TM_{010} -like mode on pill-box cavity.



◀ Fig. 7. Geometry and domains of the canonical pill-box problem.

Solve : $\nabla^2 E_z + k^2 E_z = 0$, which particularizes to :

$$r^2 \partial_r^2 [\mathbf{R}(r)] + r \partial_r [\mathbf{R}(r)] + r^2 k^2 \mathbf{R}(r) = 0, \mathbf{R}(r) \in \mathcal{C}^2$$

where $k_i = \omega/c \in \mathbb{R}^+$ on $D_i = [0, r_0) \times [-\pi, \pi) \times [0, d]$,
 and $k_o = \omega \sqrt{\epsilon_c \mu_0} \in \mathbb{C}$ on $D_o = (r_0, \infty) \times [-\pi, \pi) \times [0, d]$;
 subject to $|\mathbf{R}(0)| < \infty$ (*finiteness*)
 $(|\mathbf{R}(\infty)| < \infty$ automatically verifies since $k_o \in \mathbb{C}$).

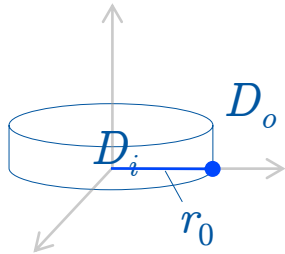
General Sol :
$$\begin{cases} A J_0(ar) \text{ on } D_i \\ B_1 H_0^{(1)}(br) + B_2 H_0^{(2)}(br) \text{ on } D_o \end{cases};$$

RF and SRF cavities: a different approach

The high-losses case

In D_o : $|\sigma|/(\omega\epsilon_0) \gg 1 \Rightarrow$ asymptotic analysis of solutions:

$$\begin{aligned}
 B_1 \mathbf{H}_0^{(1)}(br) + B_2 \mathbf{H}_0^{(2)}(br) &\sim (B_1/\sqrt{r})e^{-k_0 r} + (B_2/\sqrt{r})e^{+k_0 r} \sim \\
 &\equiv B_1(e^{-\alpha_0 r}/\sqrt{r})e^{-j\beta_0 r} + B_2(e^{+\alpha_0 r}/\sqrt{r})e^{+j\beta_0 r} \sim \\
 &\sim \underbrace{B_1 e^{-k_0 r} + B_2 e^{+k_0 r}}
 \end{aligned}$$

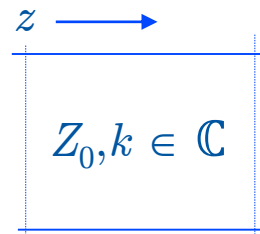


TL-like propagative waves!

◀ Fig. 7. Geometry and domains of the canonical pill-box problem.

Proposition 1: Asymptotically, the behavior of any kind of waves in high-lossy media may be studied by means of LTLT.

Fig. 8. Equivalent TL representing the high-lossy media. ▶



$$|\sigma|/(\omega\epsilon_0) \gg 1 \gg \mu_0$$

$$|Z_0| \rightarrow 0$$

$$|k| \rightarrow ? \text{ (in general: } |k| \gg 1 \text{)}$$

RF and SRF cavities: a different approach

The high-losses case

$R=0 \Rightarrow r=0$; $g = \sigma / (-\sigma')$ ($<0!$) ($\sigma \in \mathbb{C}$ from the two-fluid model)

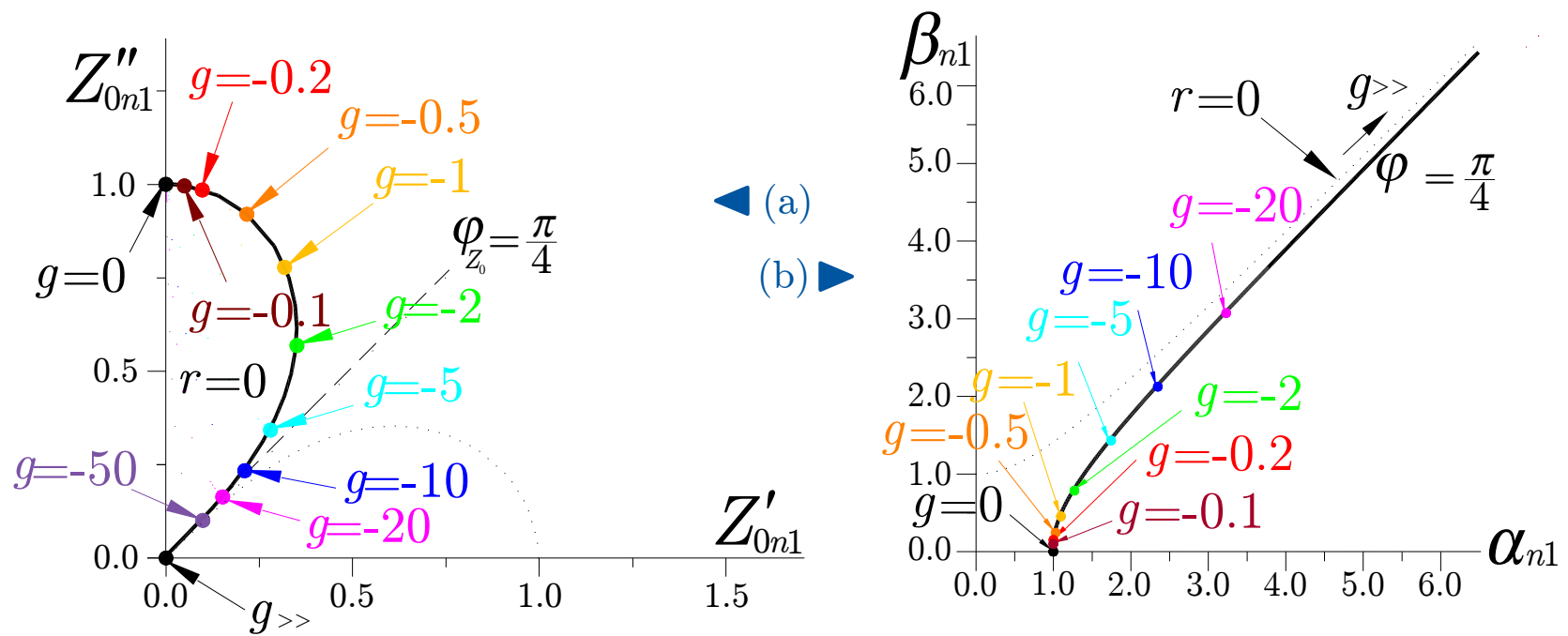


Fig. 9. Normalized (a) characteristic impedance and (b) propagation constant for good conductors.

Particular cases :

Normal conductor ($\varphi_{Z_0} \rightarrow \pi/4, \varphi_k \rightarrow \pi/4$): $Z_0 \rightarrow R_s(1+j), k \rightarrow (1/\delta)(1+j)$

Superconductor ($\varphi_{Z_0} \rightarrow \pi/2, \varphi_k \rightarrow 0$) $\rightarrow Z_0 \rightarrow R_o e^{j\pi/2}, k \rightarrow 1/\lambda_L$

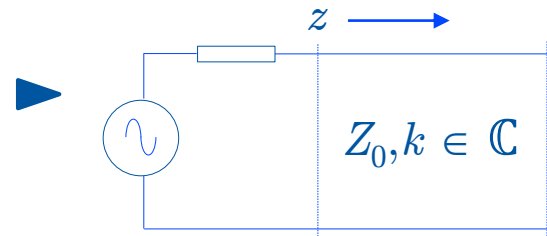
RF and SRF cavities: a different approach

The Lossy TLT connection



- (i) The surface impedance of the cavity walls is the characteristic impedance of the asymptotic approximation of the waves which propagate along lossy media.
- (ii) The cavity itself behaves as a generator for these waves (as long as $\beta_e \gg 1$, regarding the “antenna” which keeps the fields inside static).
- (iii) The impedance of the generator would play the role of the impedance of that “antenna”. Then, the maximum transfer of power (*Jacobi's law*) would represent the matching between the antenna and the cavity.

Fig. 10. Proposed equivalent circuit for the cavity characterization.



Example of application

Optimization of thin films in terms of R_s

Problem: To analyze the R_s at the input (cavity-film interface) in terms of film thickness (d) and choose the minimum (optimum).

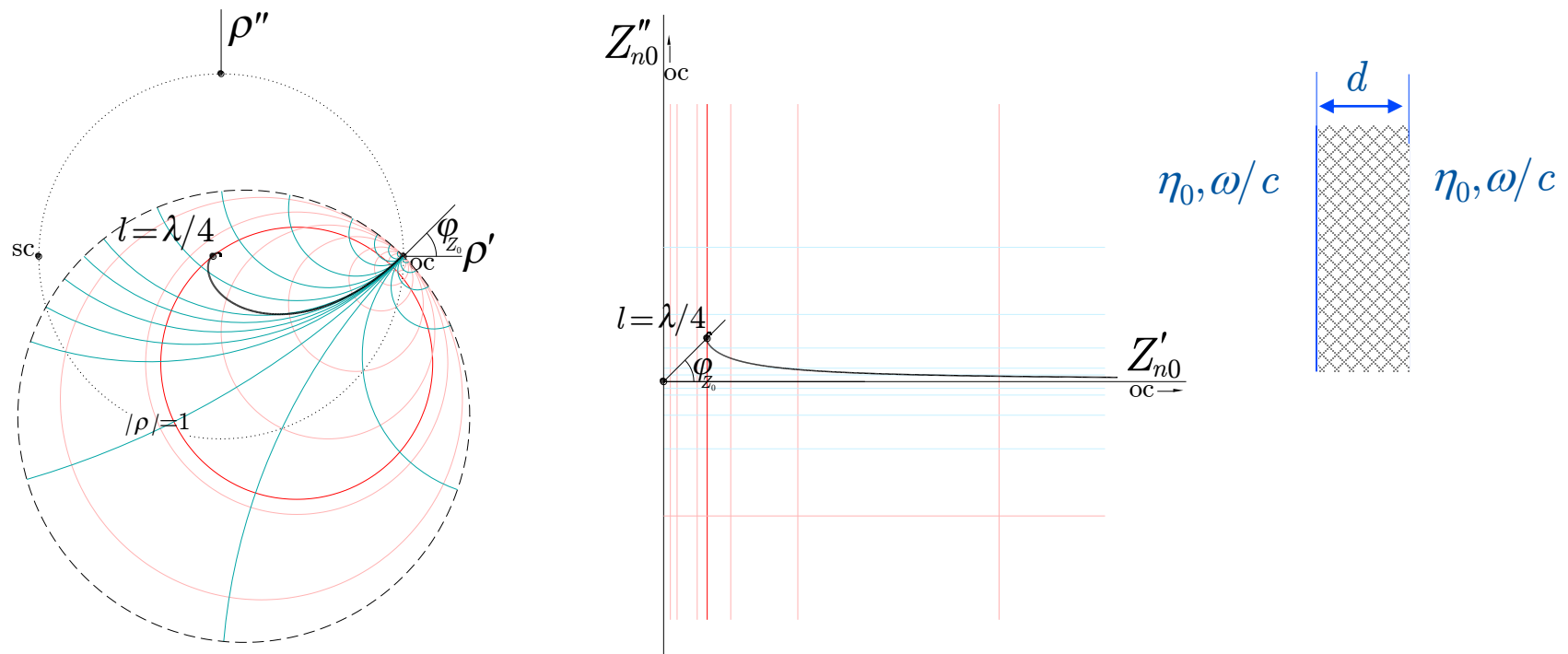


Fig. 11. Wave parameter characterization of the vacuum-thin film-vacuum structure.

Example of application

Optimization of thin films in terms of R_s

Example (good conductor): $\sigma=5.8e7$ [S·m⁻¹]

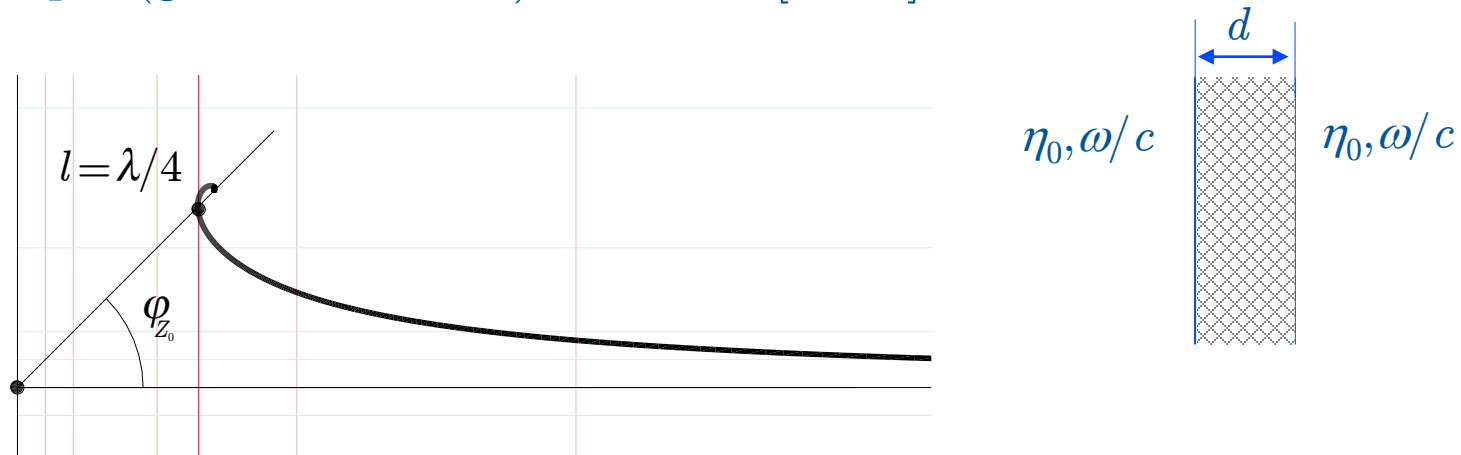


Fig. 11. Wave parameter characterization of the vacuum-thin film-vacuum structure.

A thin film of $d=\lambda/4$ ($\sim 4.83\mu\text{m}$ at 400MHz) leads to an input resistance of $0.917\cdot R_s$!!

Proposition 2 : The structure vacuum-(super)conductor-vacuum optimizes the R_s when a thin film of $d=\lambda/4$ is used in the middle.

Limitations. Proposals. Conclusions

Limitations

- (i) Depending on the SC state (given by $|g| = \sigma' / \sigma''$), it may be difficult to synthesize a thin-film $\sim \lambda/4$. (from μm to nm).
- (ii) However, the more SC the thin-film is, the more accurate is the asymptotic approximation.
- (iii) The actual equivalent R_s is not the R_s at the surface ($E(0)/H(0)$) anymore but it is given by the definite integral of the ratio of the fields in the range $[0, d]$, that is, the wave impedance along the TL.

Proposals

- (i) To use this analysis to study multilayer structures that lead to minimize the integrated R_s .
- (ii) To study the viability of using more complex TL structures (e.g. parallel stubs) and use the asymptotic analysis to study more designs (e.g. (and especially) periodic structures).

Limitations. Proposals. Conclusions

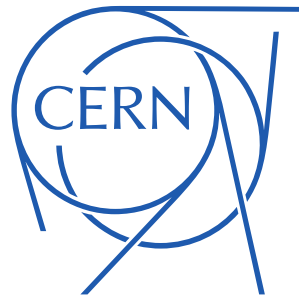
Conclusions

- (i) The LTLT is connected to the characterization of propagative waves in high lossy media where the cavity -inevitably- “radiates” the inner fields. The link is by means of the asymptotic analysis of the solutions, which is possible thanks to the high losses. In this sense, a complete equivalence between the cavity and the TL has been given.
- (ii) The usual definition of R_s at the cavity-wall interface is only valid as limit when $d \gg \lambda$. As explained and exemplified, thin-films behave differently.
- (iii) The geometrical analysis of the curves represented in the planes associated to each TL parameters becomes very helpful.
- (iv) A simple but illustrative example with practical application on thin-films deposition has been presented. From the example, and using the angle conservation property of conformal mappings, any multilayer structure may be analyzed.

References



- [1] P. Vidal-Garcia, *Generalized Study of the Complex Analysis of the Transmission Line Theory and its Application to Real Electromagnetic Systems*, PhD. Thesis at University of Oviedo, (Spain), 2019.
- [2] E. Gago-Ribas, *Complex Transmission Line Analysis Handbook*, Vol. GW-I, “Electromagnetics & Signal Theory Notebooks” series, GREditores, 2001.



Thank you for your attention!

Contact me at pablo.vidal.garcia@cern.ch