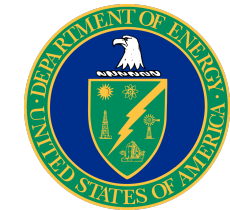


# Searching for Cosmological Concordance with New Physics in the Dark Sector: Hints and Challenges

Colin Hill

Columbia University

Copernicus Webinar  
6 June 2023



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2109.04451 w/ ACT Collaboration + 2112.10754 w/ La Posta, Louis, Garrido  
2210.14339 w/ F. McCarthy  
2112.09128 + **2212.08098** w/ M.-X. Lin, E. McDonough, W. Hu  
**2303.00746** w/ S. Goldstein, V. Irsic, B. Sherwin  
**2304.03750** w/ B. Bolliet



# “Tensions”

My personal view: observational situation remains unclear

Regardless, the situation has motivated us to think about many types of new physics in the cosmos that we otherwise (likely) would not have



How can we increase  $H_0$  inferred from the CMB and large-scale structure?  
...without worsening the  $S_8$  problem (if there is one!)



# “Tensions”

My personal view: observational situation remains unclear

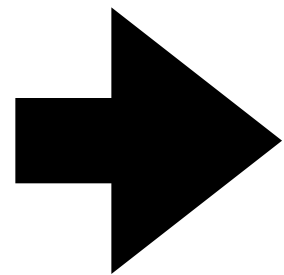
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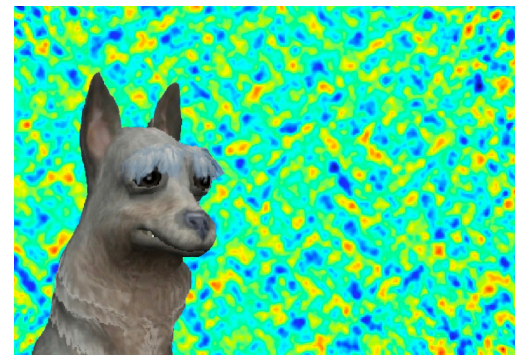
How can we increase  $H_0$  inferred from the CMB and large-scale structure?  
...without worsening the  $S_8$  problem (if there is one!)

Late-time ( $z < \text{few}$ ) theoretical modifications are highly constrained by (relative) expansion history data, e.g., BAO distances and SNIa distances

Such models often also conflict with integrated Sachs-Wolfe effect and CMB lensing data  
(e.g., [McCarthy & JCH \(2022\)](#): [2210.14339](#))



Viable models modify dynamics at high redshift



# Outline

- Classes of Viable Models
- Early Dark Energy
  - Hints? ACT DR4 (+SPT-3G)
  - Challenges —> Early Dark Sector
  - Severe Challenge: Lyman- $\alpha$  Forest
- Post-Recombination Reheating
- Outlook: ACT DR6 + Simons Observatory

# Classes of Models

## Viable paths to increase CMB-inferred $H_0$

- Pre-recombination energy injection (e.g., early dark energy and its variants)

Poulin+ (2019); Agrawal+ (2019); Lin+ (2019); Smith+ (2020); Knox & Millea (2020); JCH+ (2020); McDonough+ (2021); Lin+ (2022); ...

- Modified recombination (e.g., primordial magnetic fields; increased  $m_e$ ; or decreased  $\rho_\gamma$ )

Jedamzik & Pogosian (2018); Sekiguchi & Takahashi (2020); Hart & Chluba (2020); Chiang & Slosar (2018); Lee+ (2022); Ivanov+ (2020); JCH & Bolliet (2023)

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- Additional dark radiation species (beyond usual three neutrinos) with non-trivial dynamics/interactions

Buen-Abad+ (2015,2017,2023); Aloni+ (2021,2022); ...

- Strong neutrino interactions (delay  $\nu$  free-streaming)

Cyr-Racine & Sigurdson (2014); Lancaster+ (2017); Kreisch+ (2019); Escudero & Witte (2019)

If one of these models is actually realized in nature at a level that resolves the Hubble tension, we should soon see unambiguous evidence in CMB+LSS data



# $H_0$ and Searches for New Physics with the Atacama Cosmology Telescope

Aiola, ..., JCH, et al. (2020)

Choi, ..., JCH, et al. (2020)

JCH et al. (2021)

Thiele, Guan, JCH, Kosowsky,  
Spergel (2021)

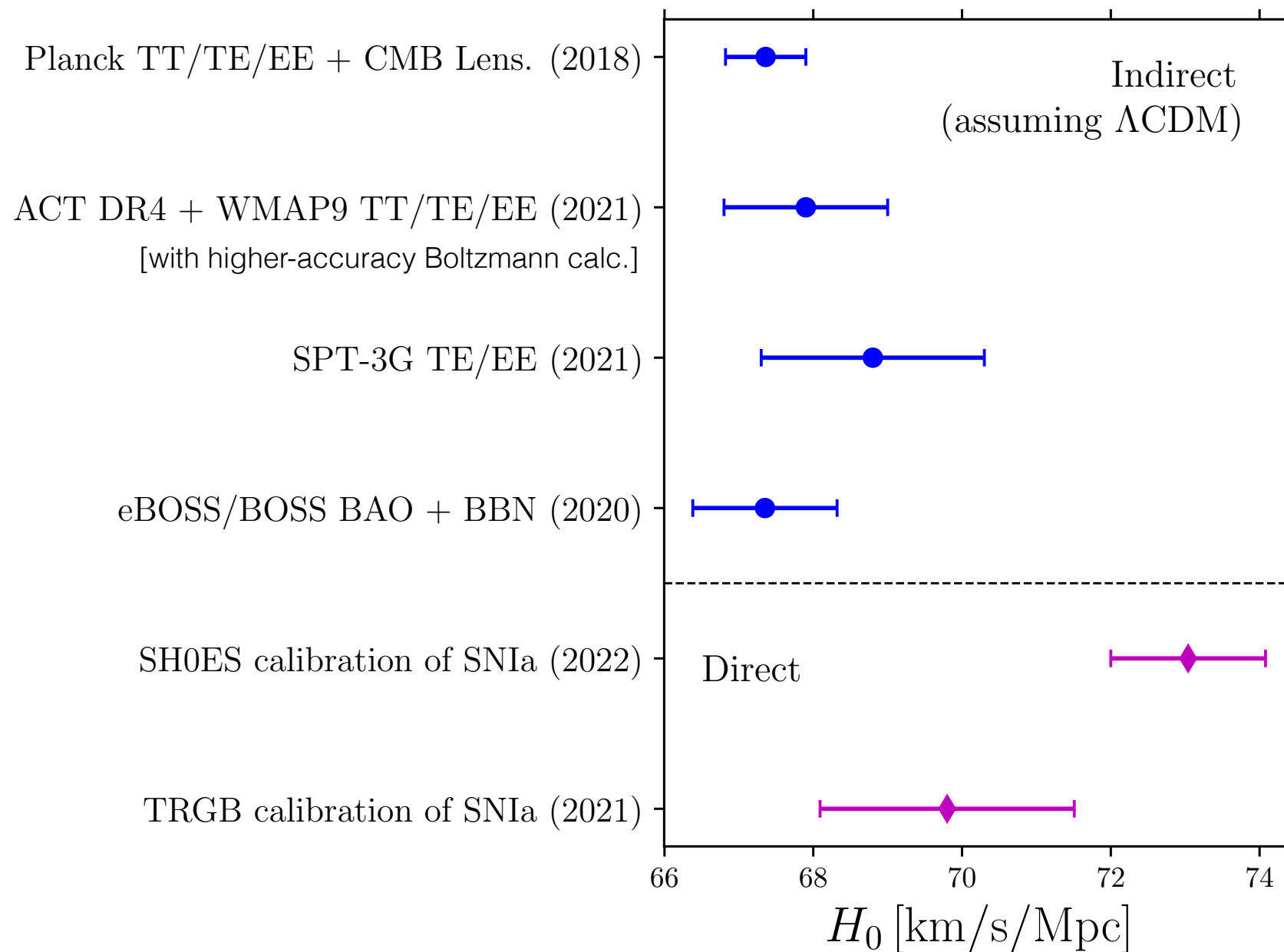
La Posta, Louis, Garrido, JCH (2022)

Kreisch, ..., JCH, et al. (2022)





# ACT DR4 Cosmology



~3.4 $\sigma$  difference between ACT+WMAP (high-acc.,  $\Lambda$ CDM) and Cepheid-calibrated SNIa (SH0ES 2022)

Agreement within ~1 $\sigma$  between ACT+WMAP and TRGB-calibrated SNIa

[Aiola et al. \(2020\)](#); [JCH et al. \(2021\)](#); see also [McCarthy, JCH, Madhavacheril \(2021\)](#)

# Classes of Models

- Modified recombination (e.g., varying  $m_e$ )
  - >some models mildly preferred by Planck
  - >most plausible model (baryon clumping due to primordial magnetic fields) is disfavored by ACT DR4

Thiele, Guan, JCH, Kosowsky, Spergel (2021)
- Additional dark radiation species with non-trivial dynamics/interactions
  - >some models mildly preferred by Planck (strong interactions and/or non-trivial time evolution is crucial)
  - >most such models weakly disfavored by ACT DR4

Schöneberg & Abellan (2022)

# Classes of Models

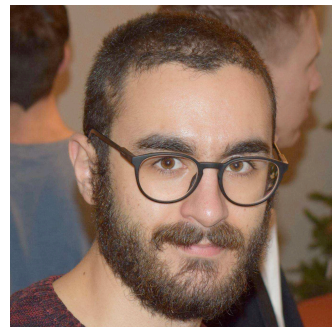
- Pre-recombination energy injection (e.g., early dark energy and its variants)
  - > not preferred by Planck or LSS data
  - > mild preference in ACT DR4

JCH et al. (2020); Ivanov, McDonough, JCH, et al. (2020); JCH et al. (2021)
- Strong neutrino interactions (delay  $\nu$  free-streaming)
  - > disfavored by Planck
  - > mild preference in ACT DR4

Kreisch, ..., JCH, et al. (2022)



# Constraints on Early Dark Energy



**JCH**, McDonough, Toomey, Alexander (2020, PRD Editors' Suggestion)  
Ivanov, McDonough, **JCH**, Simonovic, Toomey, Alexander, Zaldarriaga (2020)  
**JCH**, Calabrese, et al. [ACT Collaboration] (2021)  
La Posta, Louis, Garrido, **JCH** (2022)

# Early Dark Energy

Motivation: increase CMB-inferred  $H_0$

How does this work?

By decreasing the physical size of the sound horizon imprinted in the CMB

$$r_s^\star = \int_0^{t_\star} \frac{dt}{a(t)} c_s(t) = \int_{z_\star}^{\infty} \frac{dz}{H(z)} c_s(z)$$

scale factor      sound speed

# Early Dark Energy

Motivation: increase CMB-inferred  $H_0$

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$$r_s^* = \int_0^{t_*} \frac{dt}{a(t)} c_s(t) = \int_{z_*}^{\infty} \frac{dz}{H(z)} c_s(z)$$

Relevant ingredients in  $\Lambda$ CDM:  $\omega_b$ ,  $\omega_{\text{cdm}}$ ,  $\omega_v$ ,  $\omega_\gamma$  physical densities of  
baryons, CDM,  
neutrinos, photons

Angular sound horizon is (approx.) related to peak spacing:

$$\text{measured} \rightarrow \theta_s^* = \pi / \Delta\ell \longrightarrow D_A^* = r_s^* / \theta_s^* \longrightarrow H_0$$

$$D_A \sim 1/H_0$$

# Early Dark Energy

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$$r_s^* = \int_0^{t_*} \frac{dt}{a(t)} c_s(t) = \int_{z_*}^{\infty} \frac{dz}{H(z)} c_s(z)$$

Relevant ingredients in **EDE**:  $\omega_b$ ,  $\omega_m$ ,  $\omega_v$ ,  $\omega_\gamma$  **+ EDE parameters**

Angular sound horizon is (approx.) related to peak spacing:

$$\theta_s^* = \pi / \Delta\ell \longrightarrow D_A^* = r_s^* / \theta_s^* \longrightarrow H_0$$



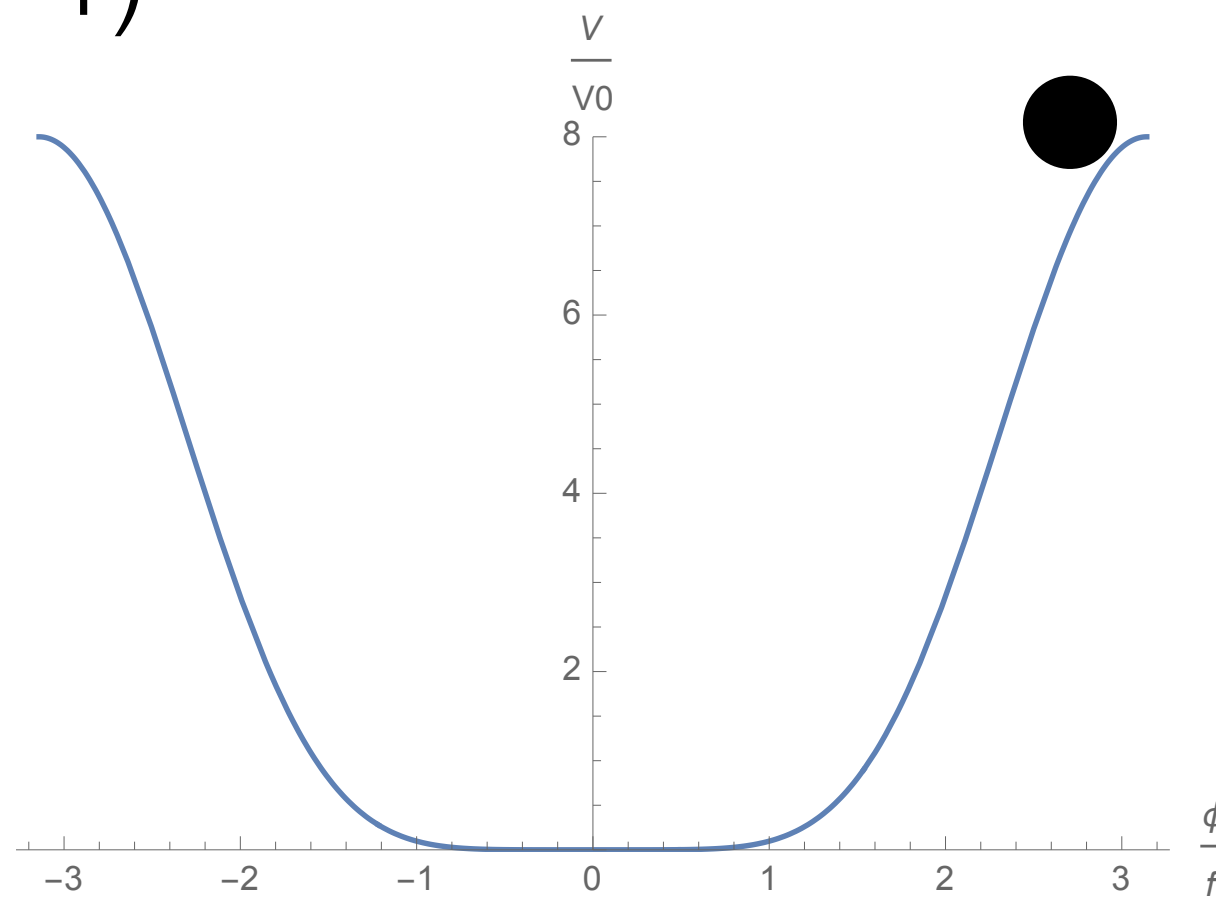
# Early Dark Energy

New component: (pseudo)-scalar field  $\phi$

# Early Dark Energy

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Idea: field initially frozen on its potential due to Hubble friction — acts as dark energy (equation of state  $P/\rho=w=-1$ )



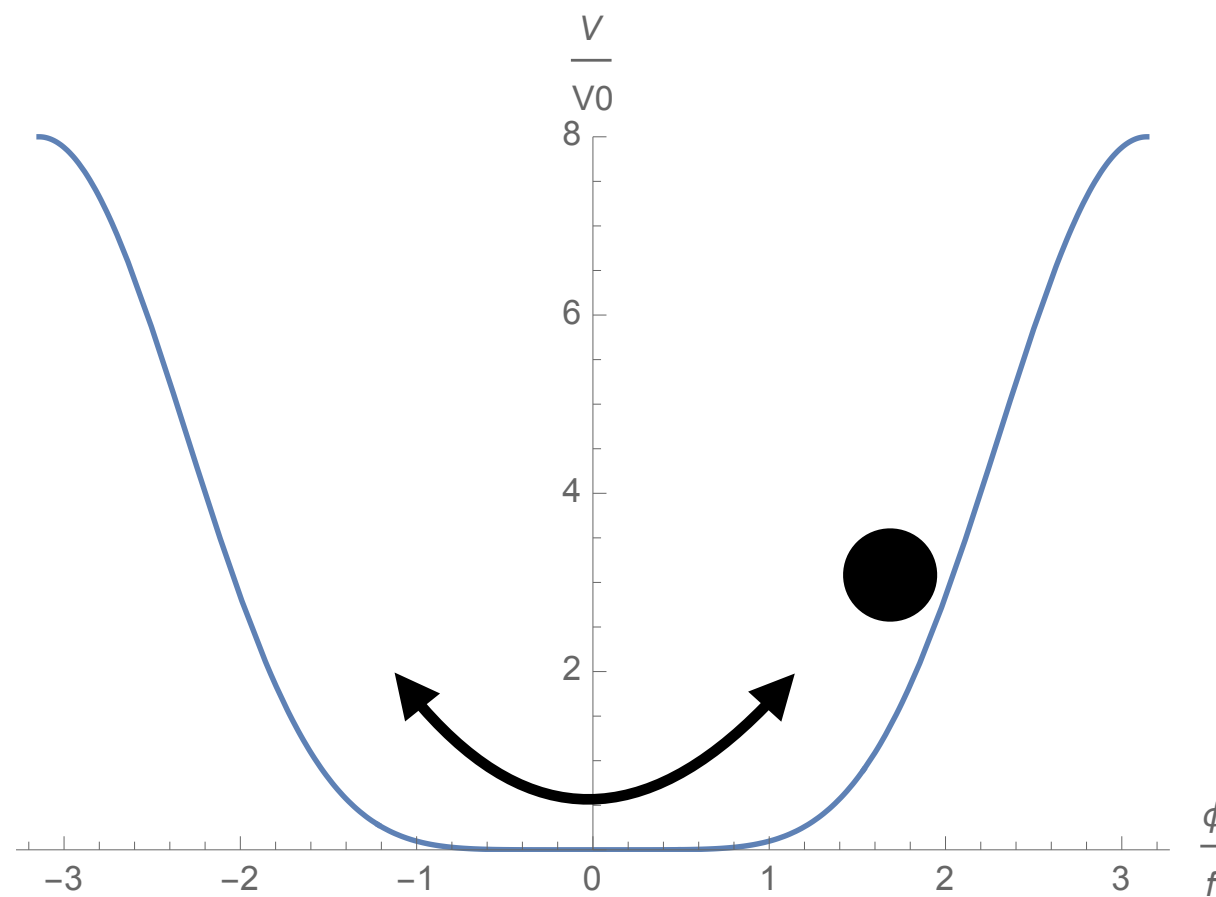
$$\ddot{\phi} + 3H\dot{\phi} + V'(\phi) = 0$$

$H \gg m$   
initially

# Early Dark Energy

New component: (pseudo)-scalar field  $\phi$

When  $H \sim m$  (field mass), it rolls down its potential and oscillates: effective EoS will depend on potential



For EDE, this must  
occur near  $\sim Z_{\text{CMB}}$



$$m \sim 10^{-27} \text{ eV}$$

e.g.,  $\phi(t) = \phi_i a^{-3/2} \cos(mt)$  if  $V(\phi) = m^2 \phi^2 / 2$

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Important: need late-time  $w>0$  so that EDE energy density decays faster than matter



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Canonical EDE  
Potential:

$$V(\phi) = m^2 f^2 (1 - \cos(\phi/f))^n$$

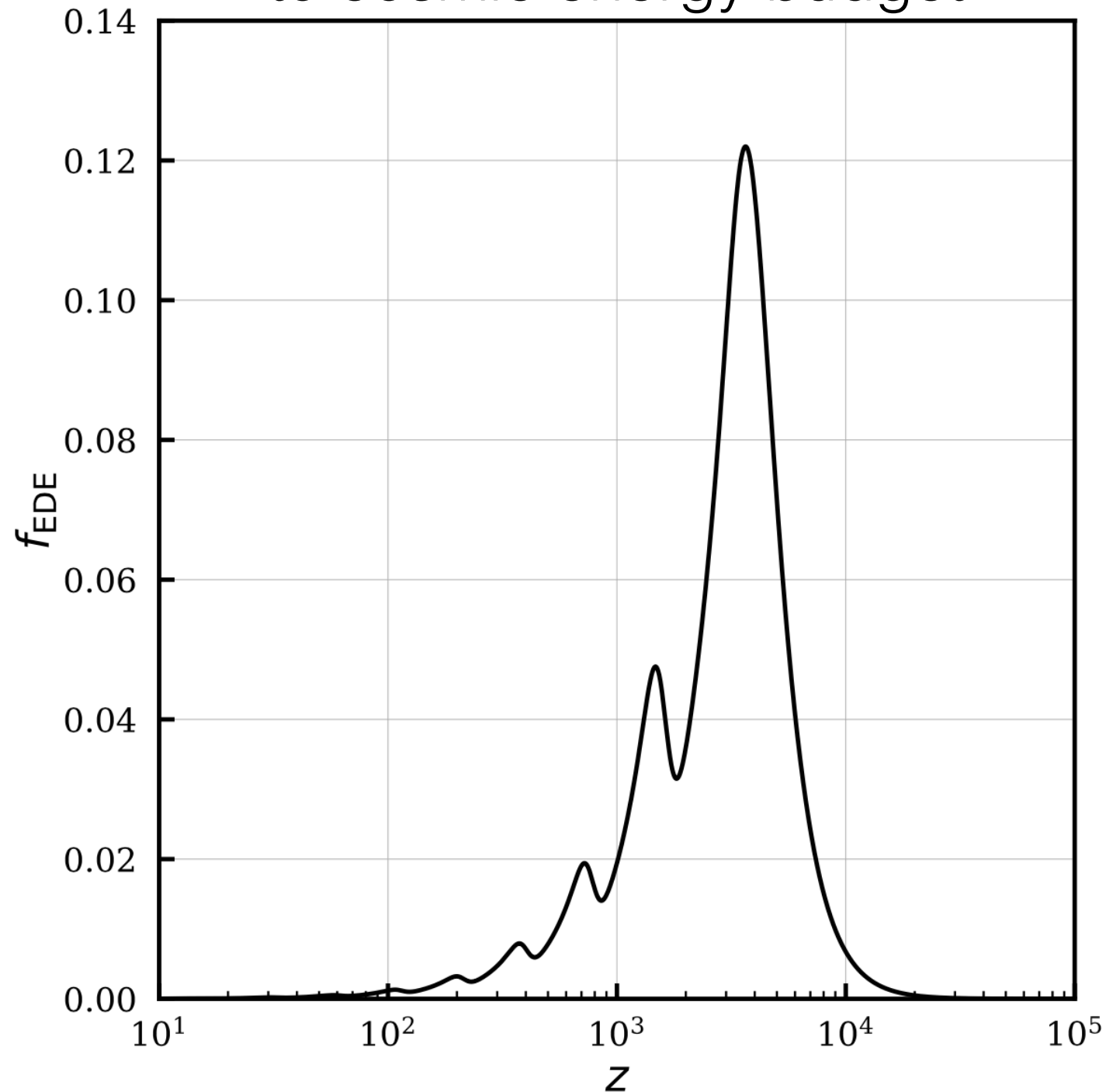
Near minimum,  $V \sim \phi^{2n} \longrightarrow w_\phi = \frac{n-1}{n+1}$

$m \sim 10^{-27}$  eV  
 $f \sim 10^{26-27}$  eV  
 $n \geq 2$  (we fix  
to 3 throughout)

# Early Dark Energy

## Parameterization

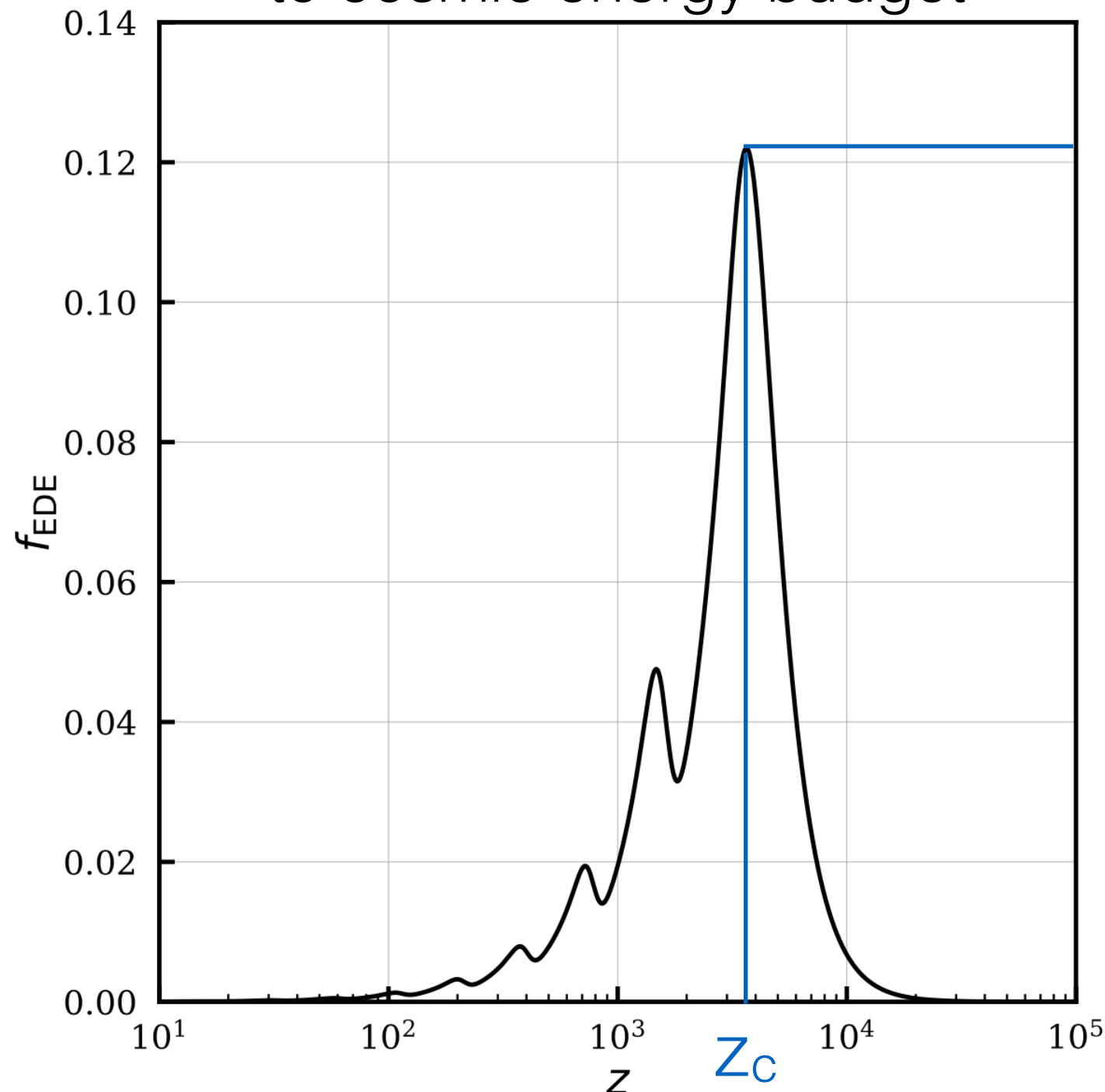
Fractional contribution of EDE  
to cosmic energy budget



# Early Dark Energy

## Parameterization

Fractional contribution of EDE  
to cosmic energy budget



Maximal contribution:

$$f_{\text{EDE}}(z_c) \equiv (\rho_{\text{EDE}}/3M_{pl}^2 H^2)|_{z_c}$$

which occurs at redshift  $z_c$

Final parameter:  $\theta_i = \phi_i/f$   
(initial field displacement)

➡  $\{f_{\text{EDE}}, z_c, \theta_i\}$

N.B.: highly non-linear  
relation to physical scalar  
field parameters

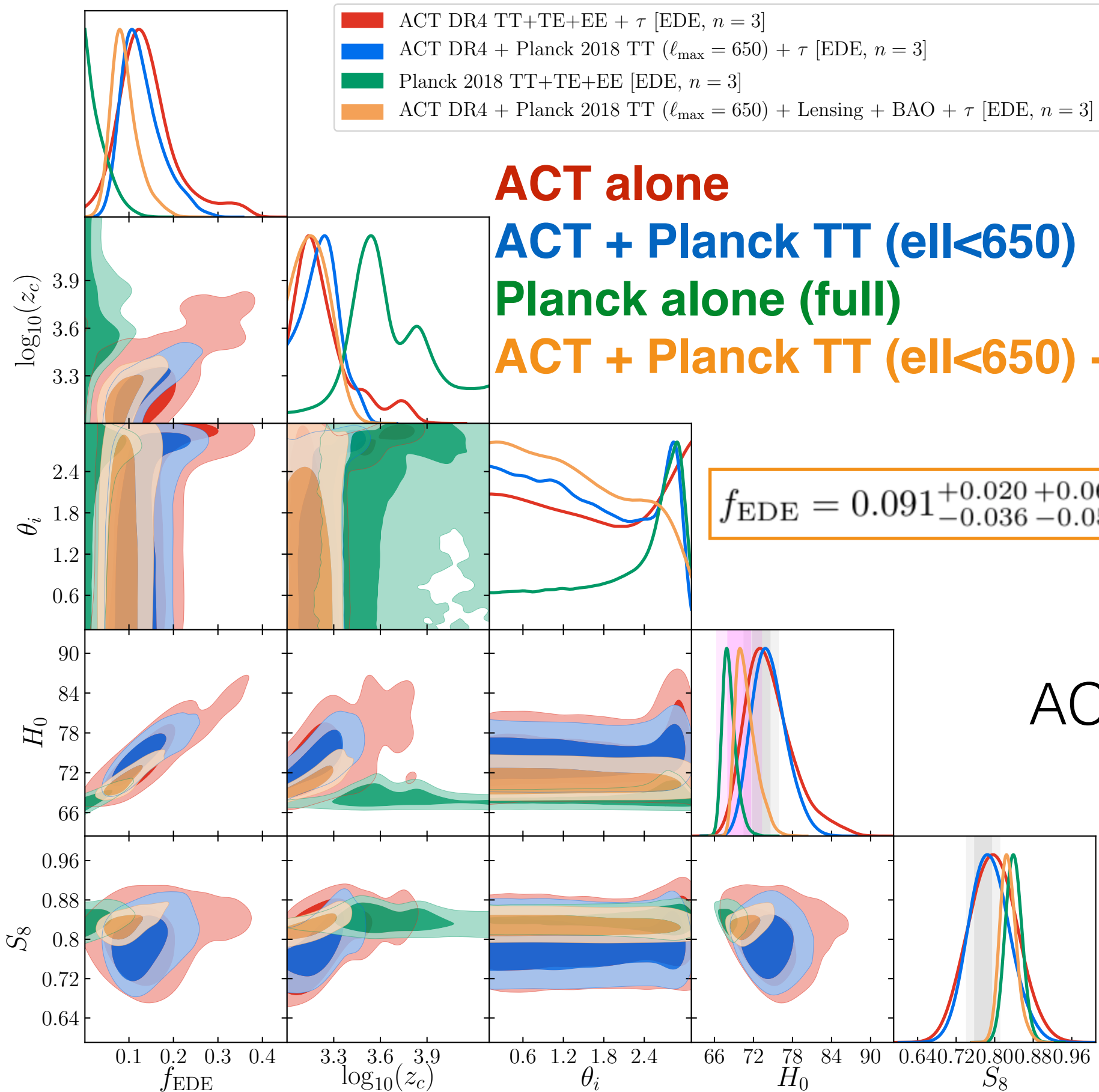
# ACT DR4 EDE Analysis

Colin Hill  
Columbia

The Atacama Cosmology Telescope: Constraints on Pre-Recombination Early Dark Energy

# ACT DR4 EDE Results

Colin Hill  
Columbia



**ACT alone**

**ACT + Planck TT ( $\ell < 650$ )**

**Planck alone (full)**

**ACT + Planck TT ( $\ell < 650$ ) + CMB Lensing + BAO**

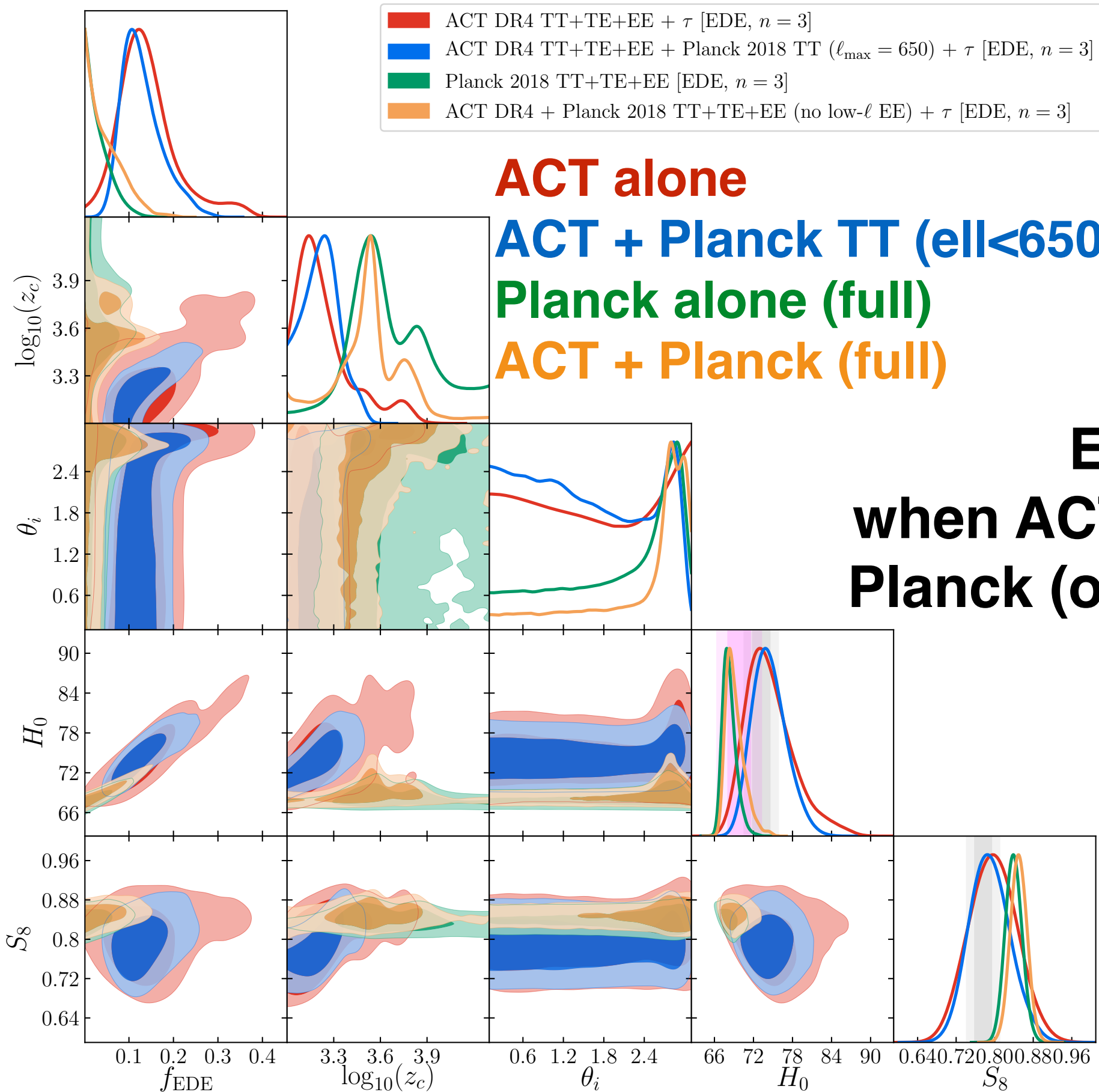
$$f_{\text{EDE}} = 0.091^{+0.020}_{-0.036} {}^{+0.069}_{-0.056} {}^{+0.11}_{-0.063} \quad (68\%/95\%/99.7\% \text{ CL})$$

$$H_0 = 70.9^{+1.0}_{-2.0} \text{ km/s/Mpc}$$

ACT drives preference  
for non-zero  $f_{\text{EDE}}$   
( $>99.7\%$  CL in joint  
fits)

# ACT DR4 EDE Results

Colin Hill  
Columbia



**ACT alone**

**ACT + Planck TT ( $\ell < 650$ )**

**Planck alone (full)**

**ACT + Planck (full)**

**EDE hint goes away  
when ACT is combined with  
Planck (overall constraining  
power still Planck-  
dominated)**



# Origin of ACT EDE Hint

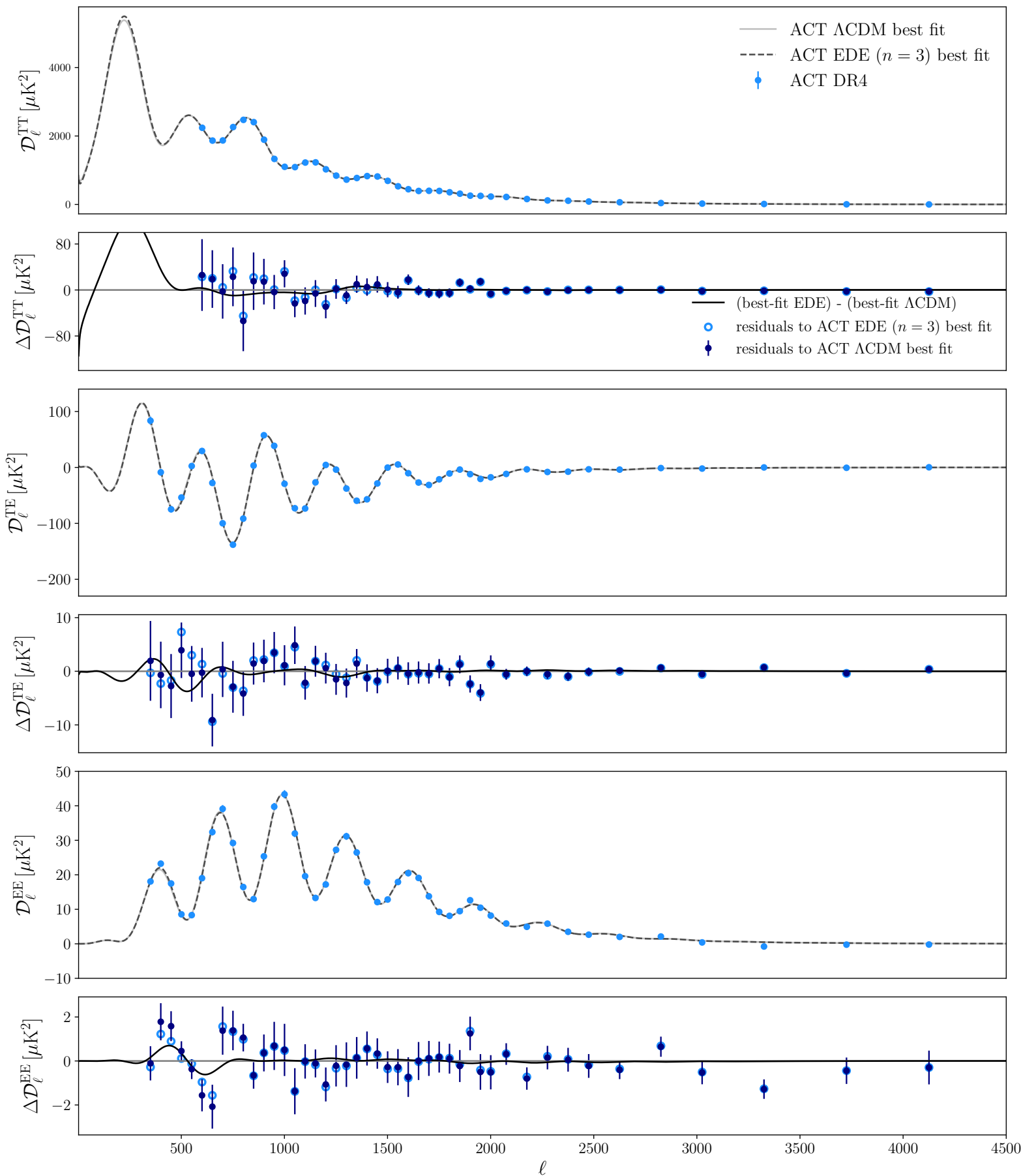
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# Origin of ACT EDE Hint

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LCDM  
residuals

EDE  
residuals



TT

TE

EE

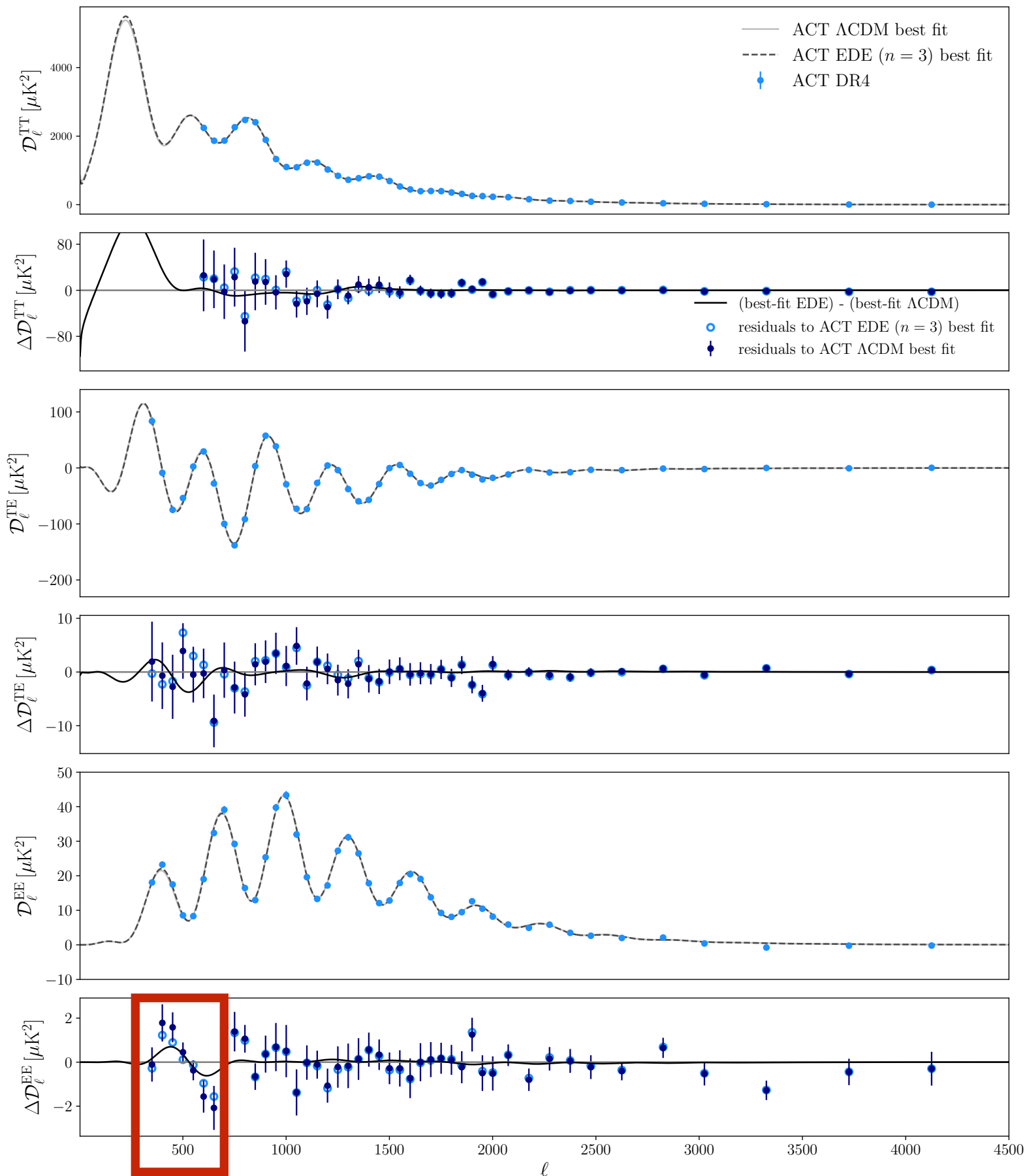
# Origin of ACT EDE Hint

LCDM  
residuals

EDE  
residuals

lowest  $\ell$   
bins in EE  
drive the  
preference

JCH et al. (2021)



TT

TE

EE

overall  
preference  
 $\sim 2.1\sigma$   
( $\Delta\chi^2 = -8.7$ )

# Origin of ACT EDE Hint

LCDM  
residuals

EDE  
residuals

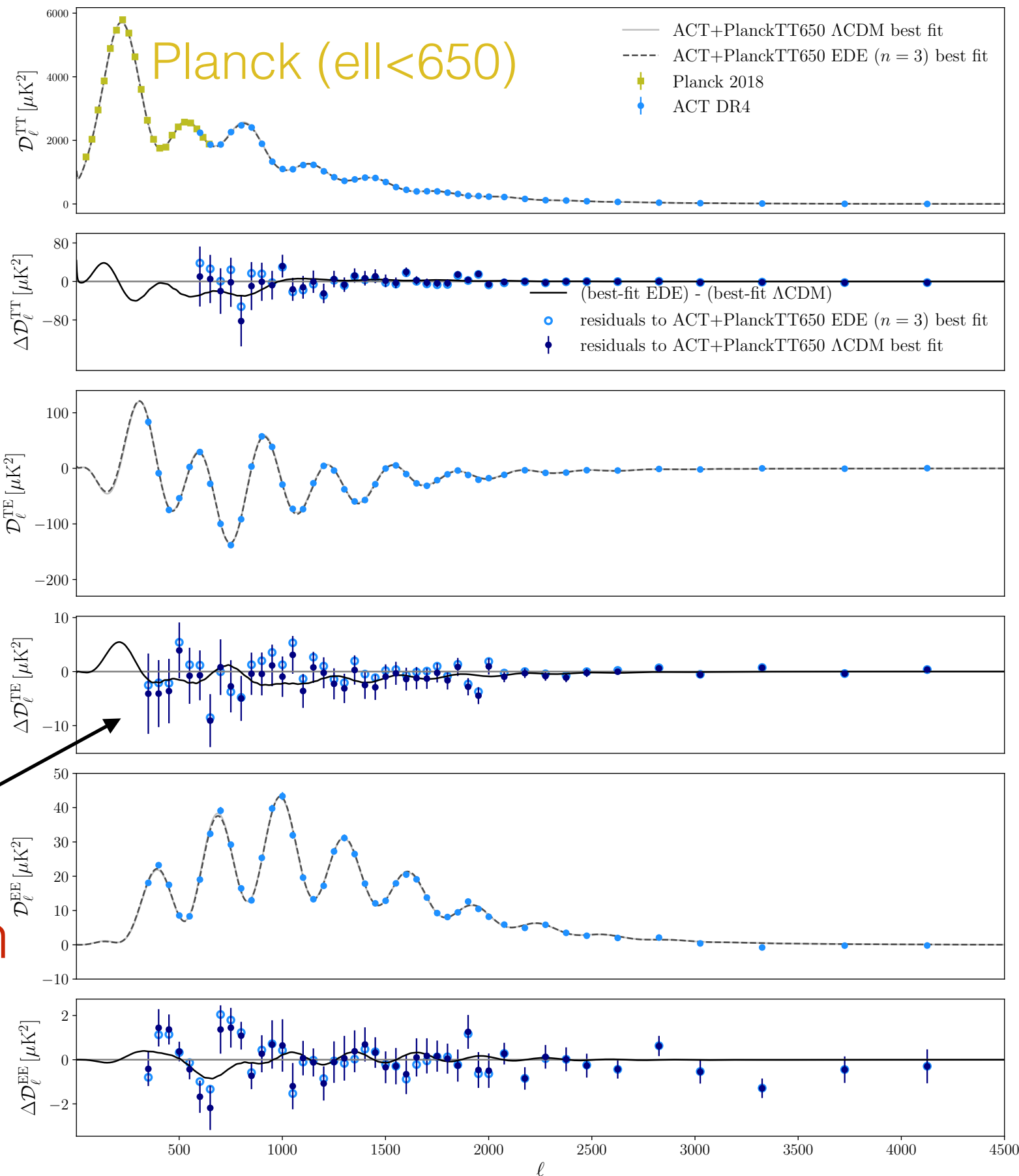
TE plays an  
important  
role in  
driving EDE  
preference in  
joint fits

TT

TE

EE

overall  
preference  
 $\sim 3.2\sigma$   
( $\Delta\chi^2 = -15.4$ )

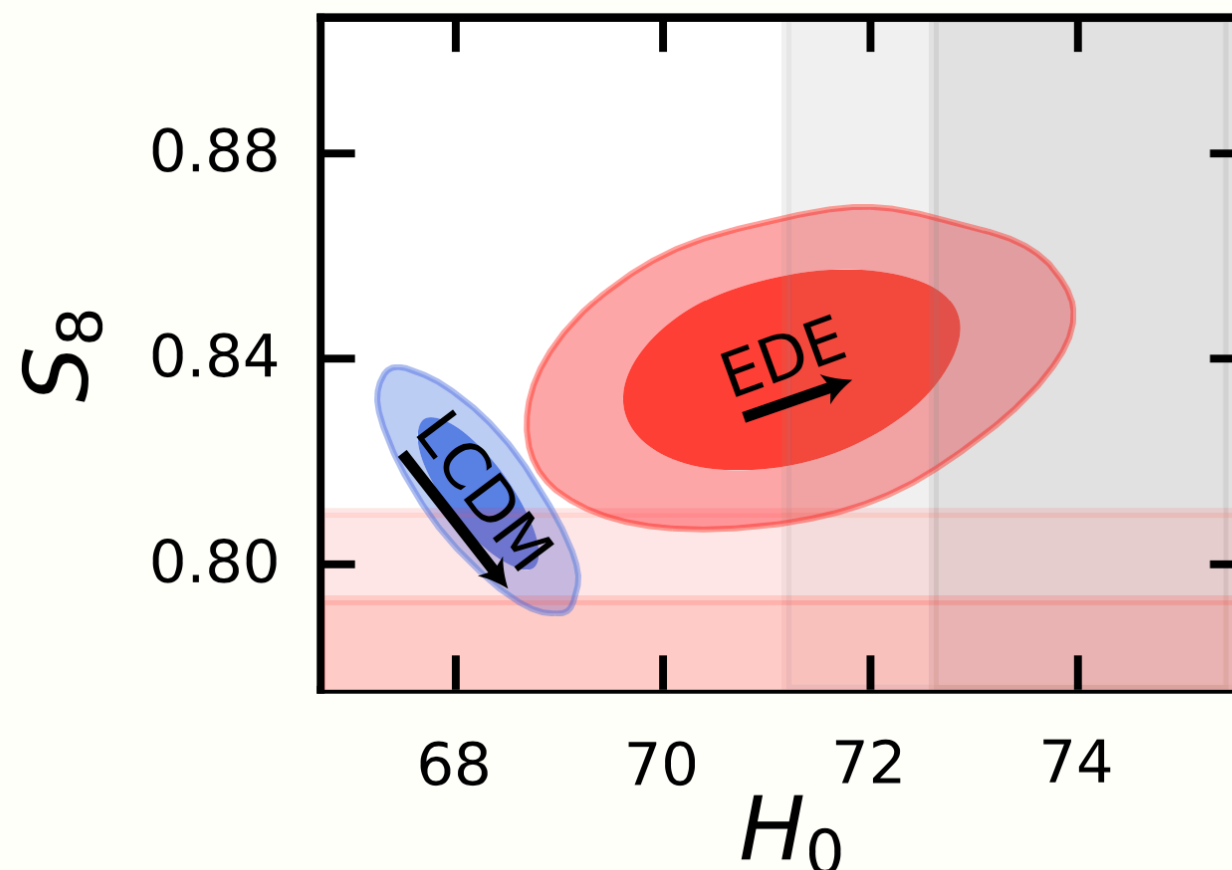


# EDE Puzzles & Problems



# EDE Puzzles & Problems

- Coincidence problem: why should these new dynamics appear near  $z_{\text{eq}}$ ? [ $\rightarrow V(\phi), V'(\phi)$ ]
- Initial conditions: axion-like field must start near top of cosine to fit Planck (e.g., Lin, Benevento, Hu, Raveri (2019)) [ $\rightarrow V''(\phi)$ ]
- “Tension-trading”:  $H_0$  increases in the CMB fit at the cost of adding significantly more dark matter and increasing  $n_s$ , hence raising  $S_8$



# EDE Puzzles & Problems

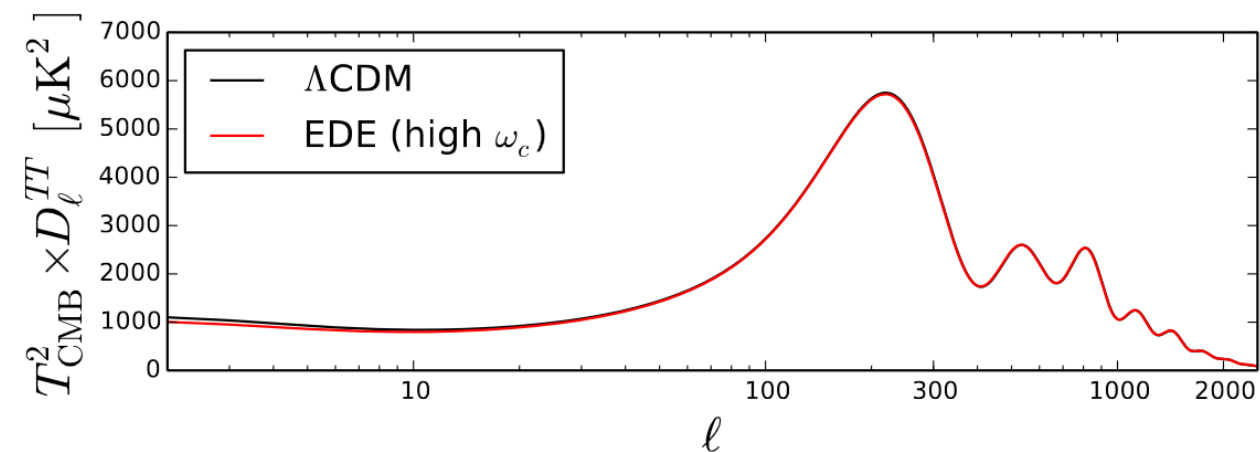
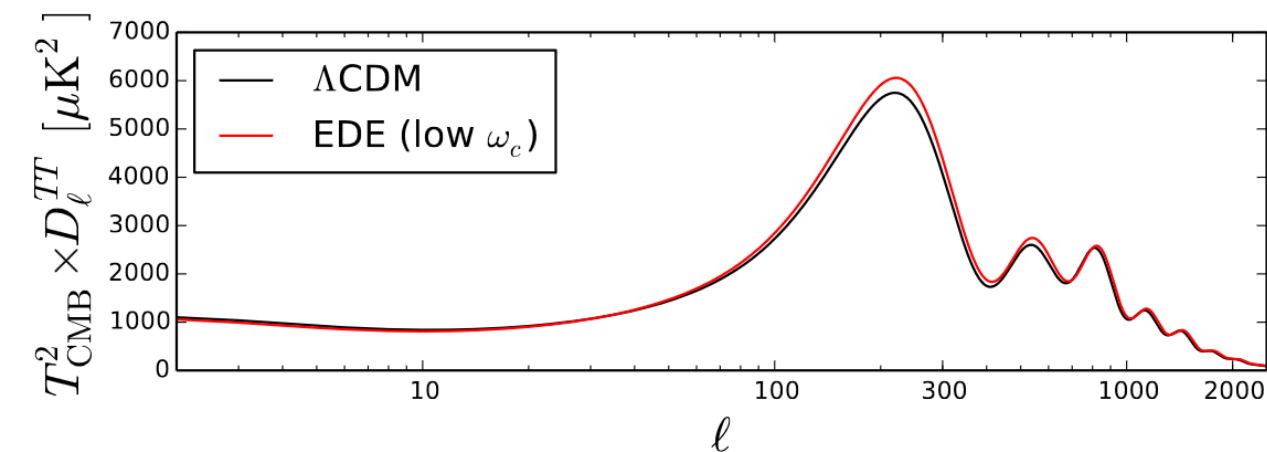
Why do  $\omega_c$  and  $n_s$  increase when fitting EDE to CMB data?

- Recall the integrated Sachs-Wolfe (ISW) effect: grav. potentials decay in a non-matter-dominated universe
- Early ISW arises because radiation is still important at  $z^*$   
—>Enhanced in an EDE cosmology (because the EDE is not matter)

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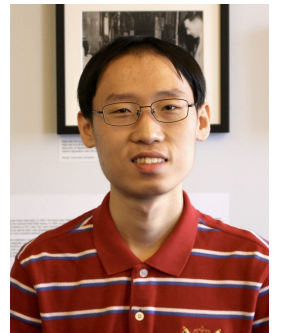
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—>Enhanced in an EDE cosmology (because the EDE is not matter)



primarily compensated by increasing the CDM density ( $\omega_c$ ), but also by increasing the slope of the power spectrum ( $n_s$ )

# Early Dark Sector

**A Dark Matter Trigger for Early Dark Energy Coincidence**



2112.09128 w/ Evan McDonough, Meng-Xiang Lin, Wayne Hu, Shengjia Zhou  
2212.08098 w/ Lin, McDonough, Hu

# Early Dark Sector

Goal: explain why EDE dynamics at  $z_c$  seem to be coincident with  $z_{eq}$  by coupling the EDE scalar  $\phi$  to the dark matter, such that DM triggers EDE evolution rather than the bare potential  $V(\phi)$

- Field dependent dark matter mass:  $m_{dm}(\phi)$
- Effective potential:  $V_{eff} = V_0 + m_{dm}(\phi)n_{dm}$



# Early Dark Sector

Colin Hill  
Columbia

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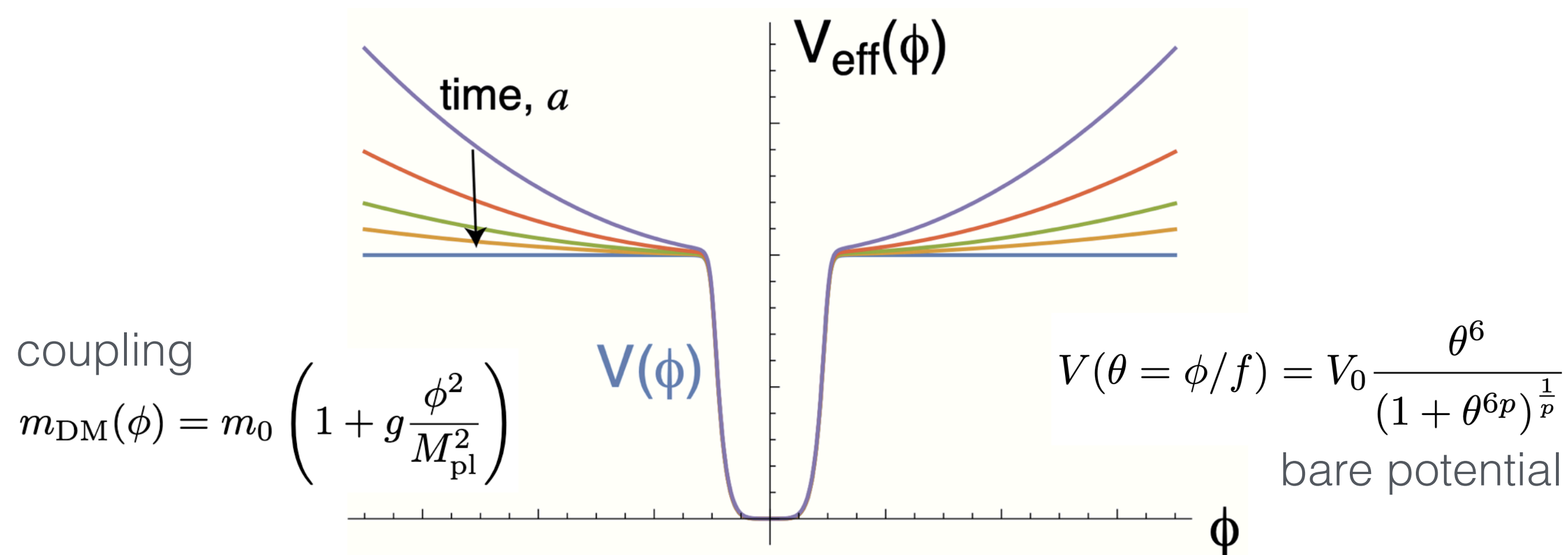
- Field dependent dark matter mass:  $m_{dm}(\phi)$
- Effective potential:  $V_{eff} = V_0 + m_{dm}(\phi)n_{dm}$

- Generically this produces evolution in the DM mass
- Problem for acceptable  $\Delta m_{DM}/m_{DM}$  and generic initial conditions: slope of bare potential in axion-like EDE is too steep to “trigger” off the EDE-DM coupling
- Solution: flatten  $V(\phi)$  into a plateau and choose EDE-DM coupling  $m(\phi)$  such that  $V_{eff}(\phi) \sim \rho_{DM}$  and  $\phi$  is released from Hubble friction near  $z_{eq}$

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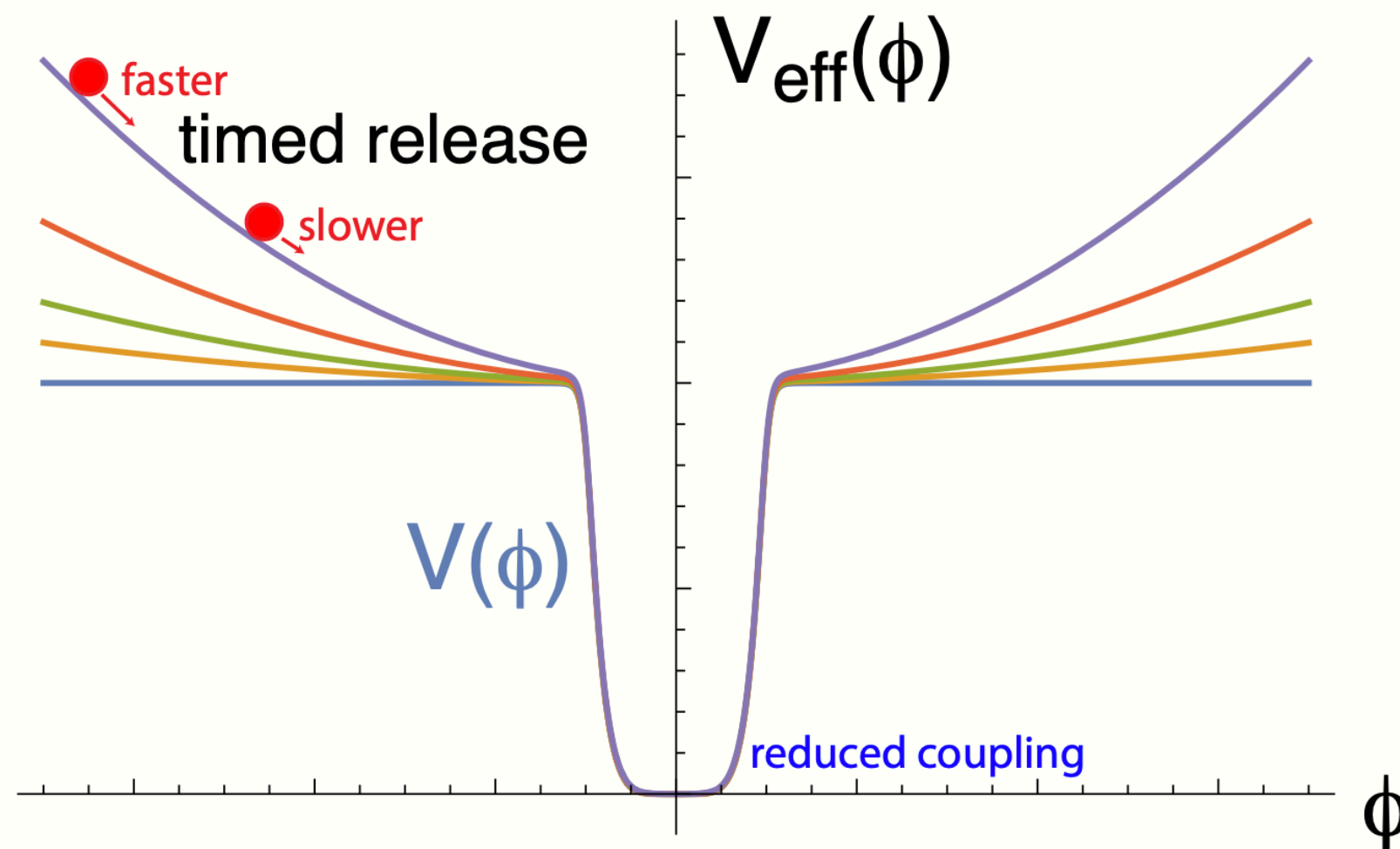
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# Early Dark Sector

## Solution

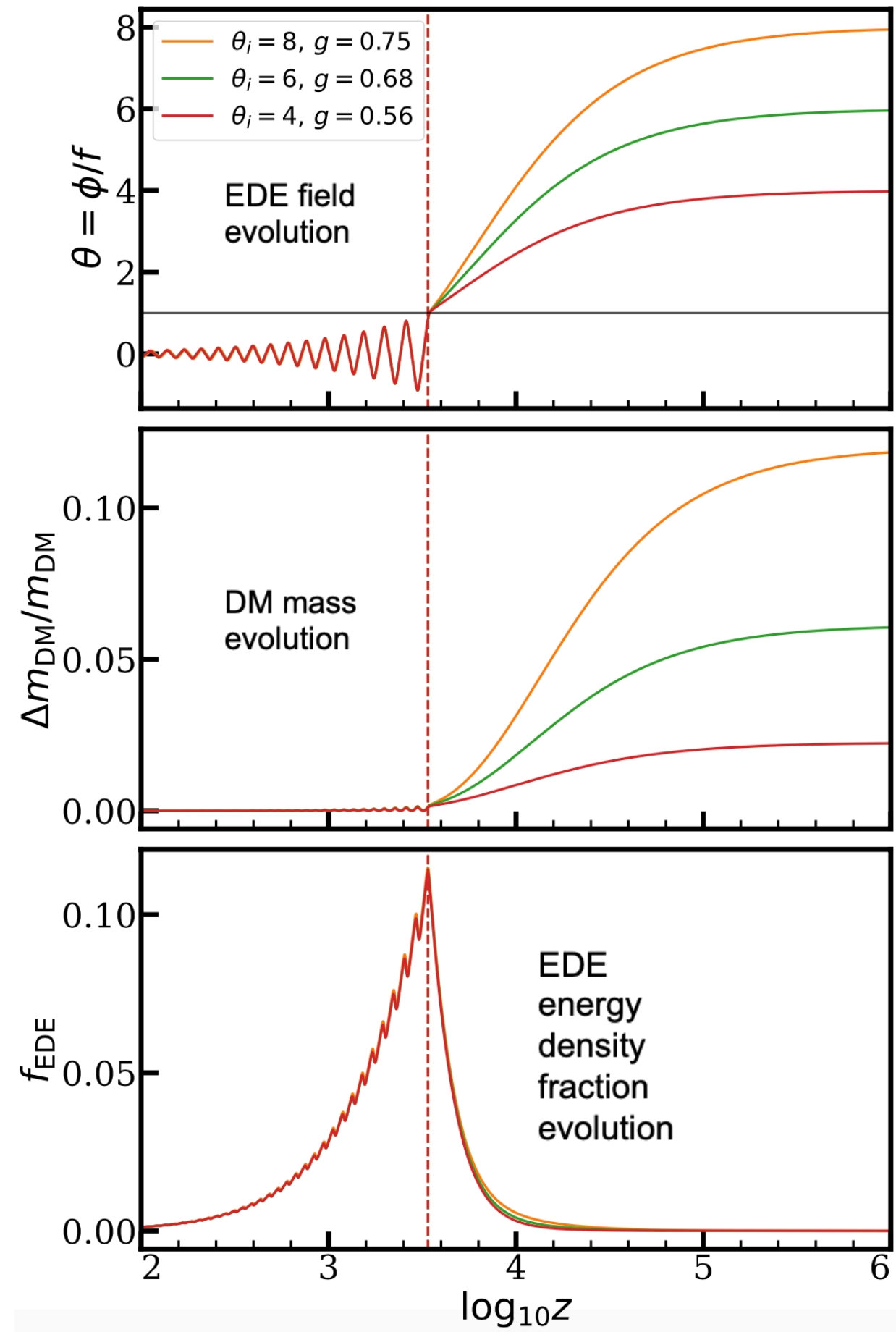
- Explains coincidence:  $\phi$  starts to roll because of equality
- No initial tuning needed:  $\phi$  rolls to edge of plateau from wide range of initial field values
- No late-time growth problems: for  $m(\phi) \sim 1 + g\phi^2$ , there is no fifth force since  $\phi \rightarrow 0$  at late times



# Early Dark Sector

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Basic validation: can successfully  
lower  $r_s$ , raise  $H_0 \sim 71.2$  km/s/Mpc



# Early Dark Sector

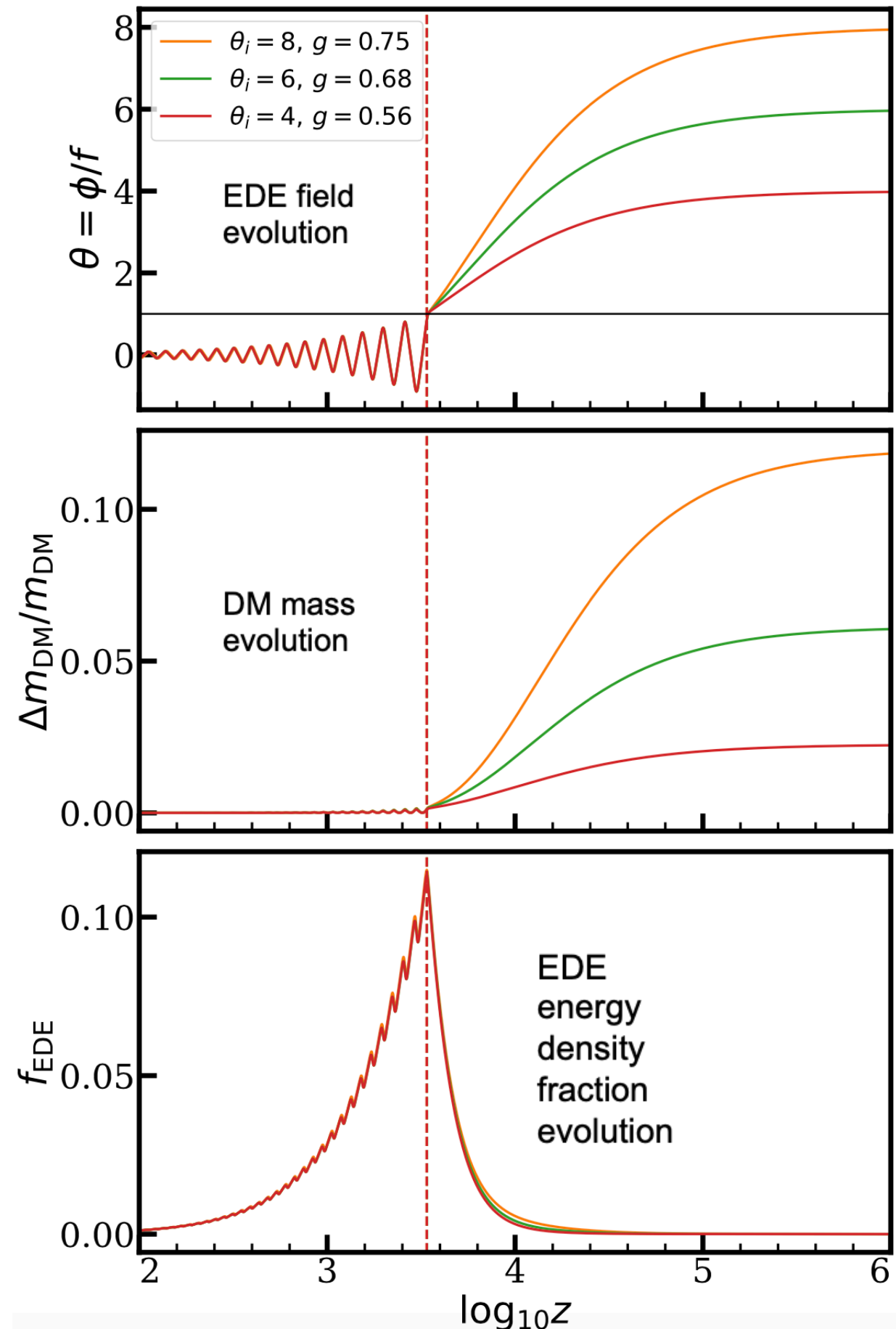
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Columbia

Basic validation: can successfully  
lower  $r_s$ , raise  $H_0 \sim 71.2$  km/s/Mpc

Best-fit parameters to  
Planck+BAO+SNIa+SH0ES  
+DES-Y3:

Model	EDE	tEDS( $p=8$ )
$f_{\text{EDE}}$	0.108	0.112
$\log_{10} z_c$	3.56	3.83
$H_0$	71.96	71.21
$S_8$	0.8236	0.8200
$n_s$	0.9894	0.9843

- Goodness-of-fit nearly identical to EDE
- Coincidence problem resolved
- Fine-tuning of initial conditions resolved
- $S_8$  problem partially ameliorated

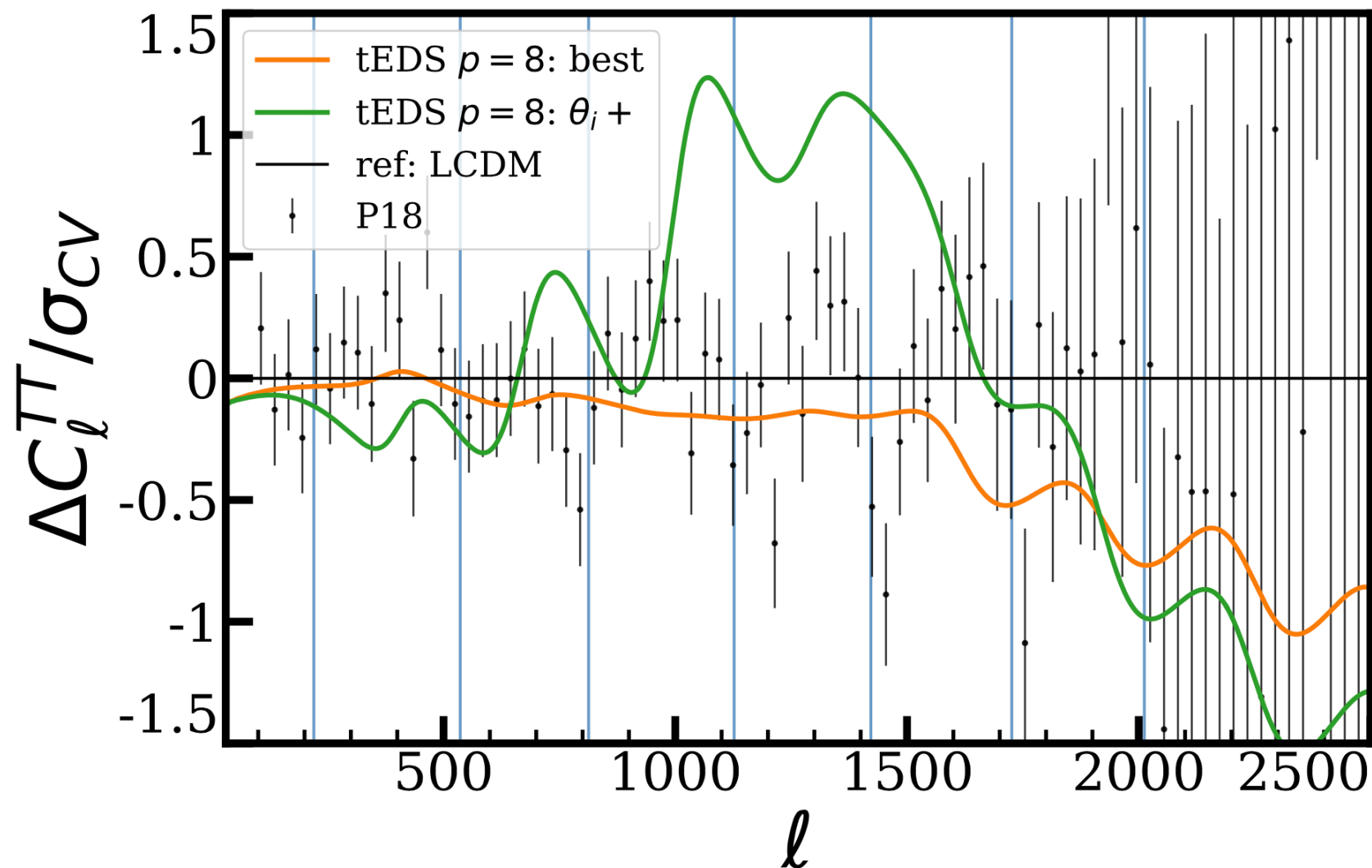




# Early Dark Sector

However: excess field fluctuations induced by rolling in  $V_{\text{eff}}(\phi)$

Consider increase in initial field position ( $\theta_i$ ), hold  $z_c$  and  $V(\phi)$  fixed



Result: data pick out specific  $\theta_i$  to achieve dynamical balance

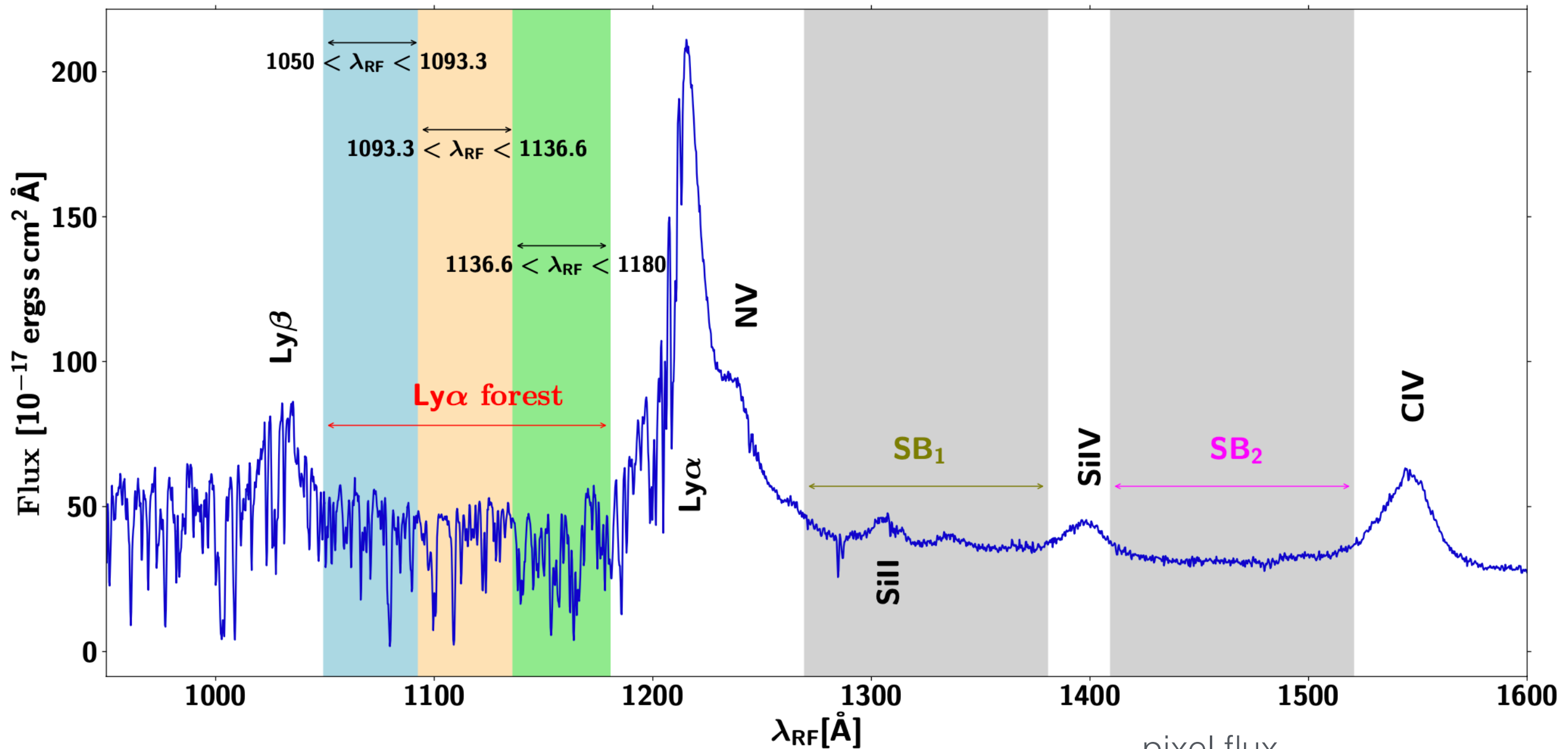
Next: MCMC/further model improvements

# Challenge: Lyman- $\alpha$ Forest



# Ly $\alpha$ Forest

Absorption lines due to HI clouds along the LOS to a distant quasar



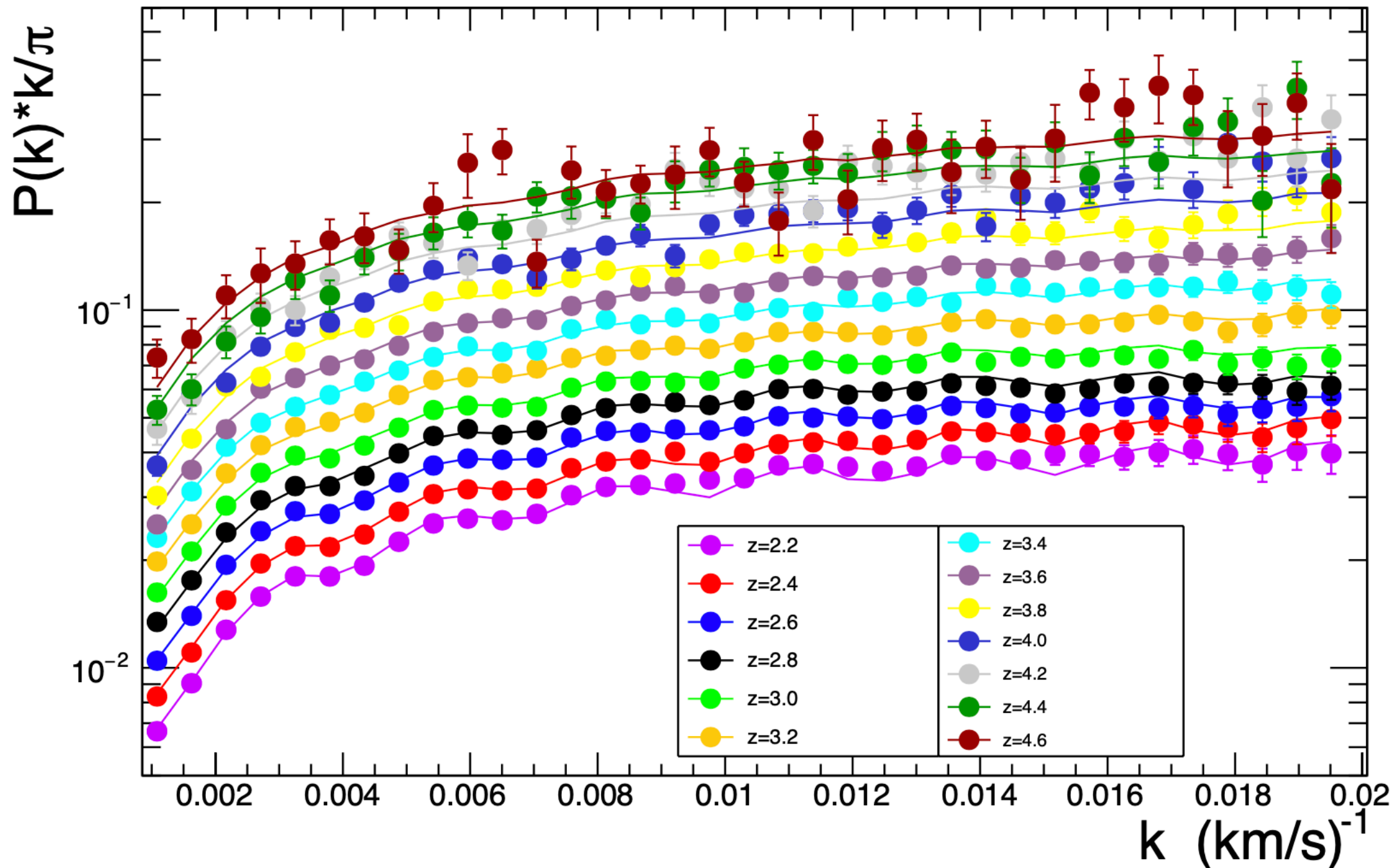
$$\delta(\lambda) = \frac{\text{pixel flux } f(\lambda)}{\text{quasar continuum } C_q(\lambda) \overline{F}(z_{\text{Ly}\alpha})} - 1$$

mean trans. flux fraction

# Lya Forest

Observable: 1D flux power spectrum

BOSS/eBOSS: ~44,000 quasar spectra (moderate S/N)

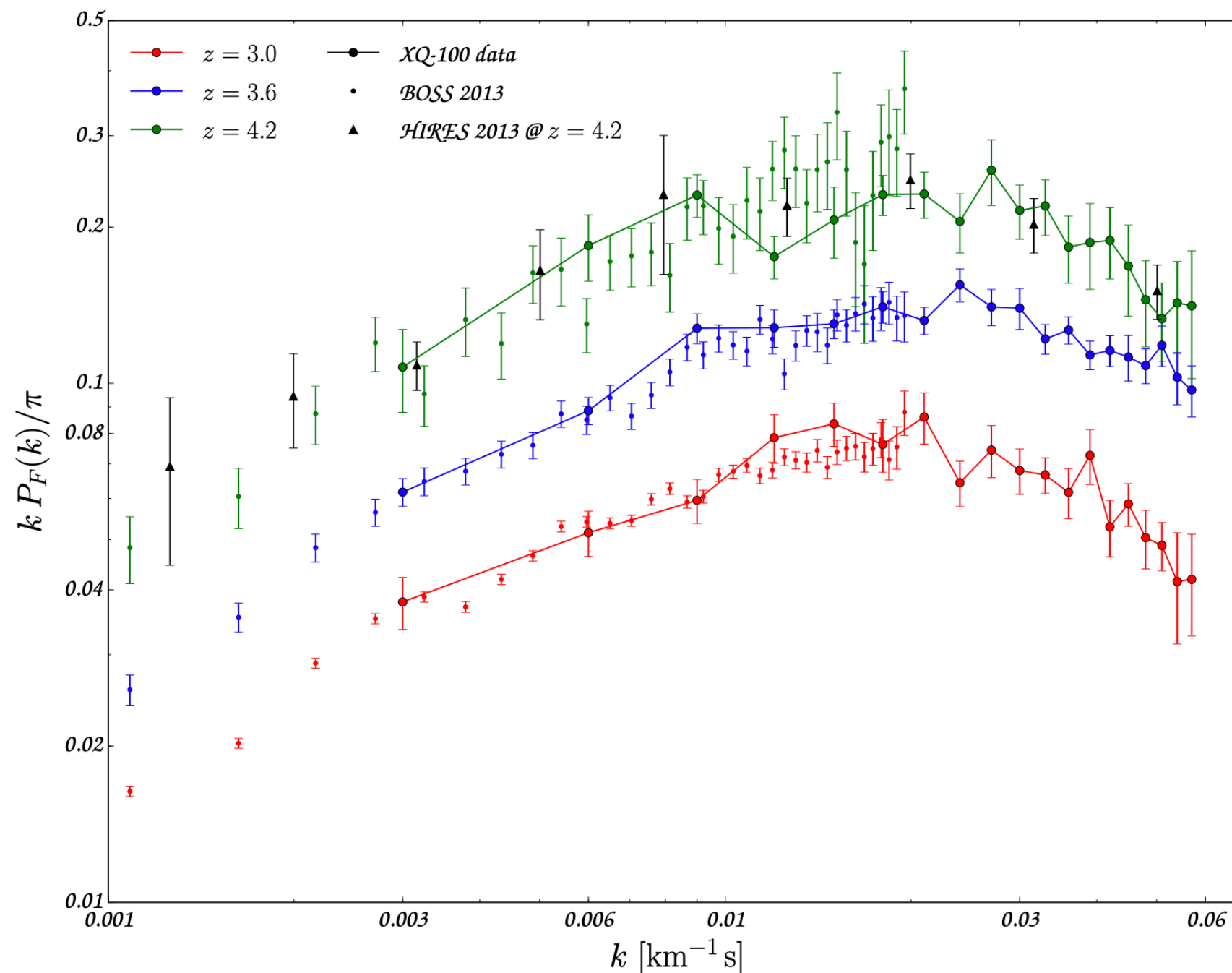


$k \sim 0.01 \text{ s/km} \longleftrightarrow k \sim 1 \text{ h/Mpc at } z=3$

# Lya Forest

Observable: 1D flux power spectrum

XQ100: 100 quasar spectra (high-S/N)

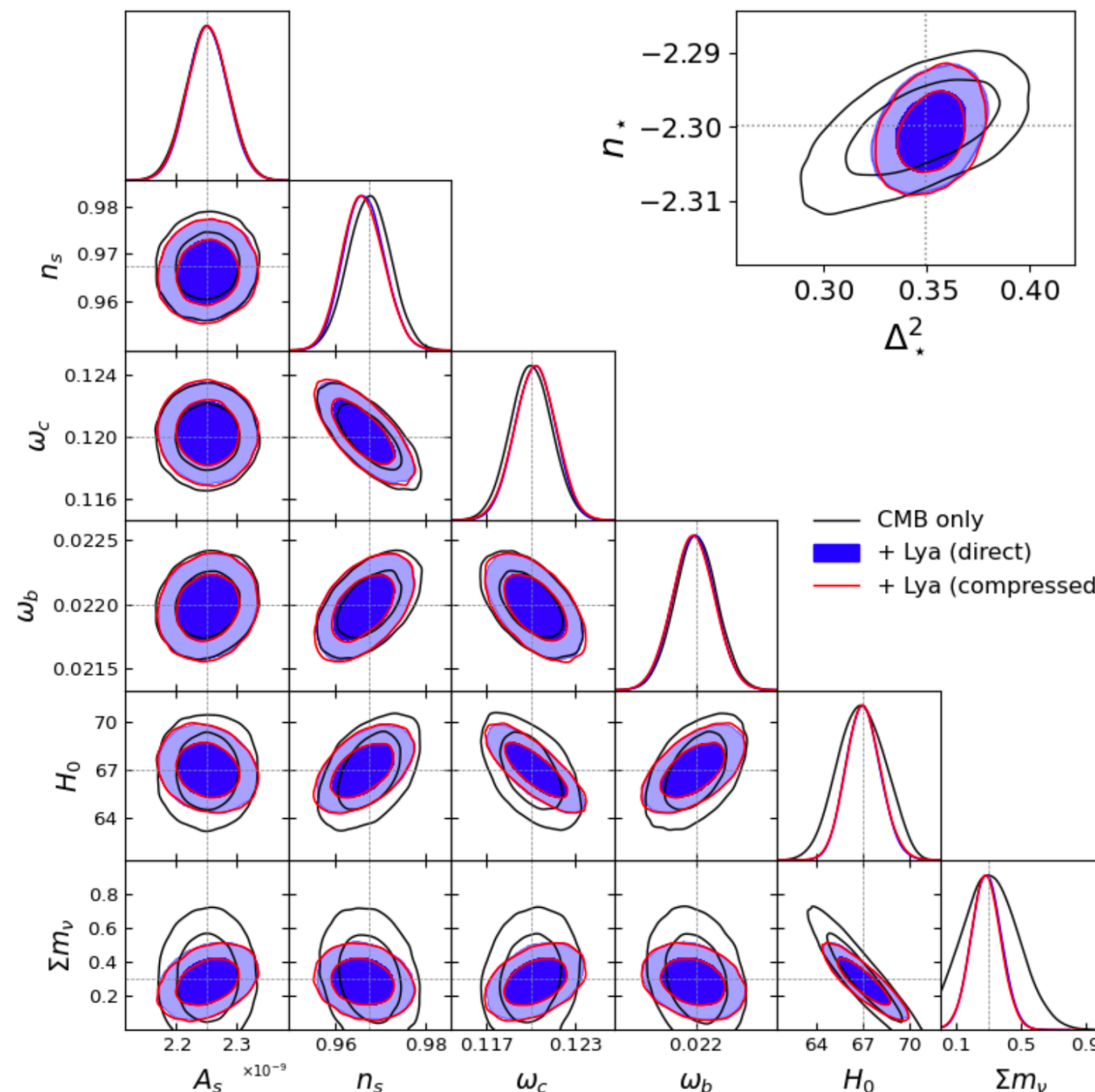


+ MIKE/  
HIRES  
spectra at  
 $z=4.2 - 5.4$

# Lya Forest

Information content fully contained in compressed 2D likelihood

amplitude  $\Delta_L^2 \equiv k^3 P_{\text{lin}}(k_p, z_p)/(2\pi^2)$  at  $k_p = 0.009 \text{ s/km}$   
 slope  $n_L \equiv d \ln P_{\text{lin}}(k_p, z_p)/d \ln k$  and  $z_p = 3$

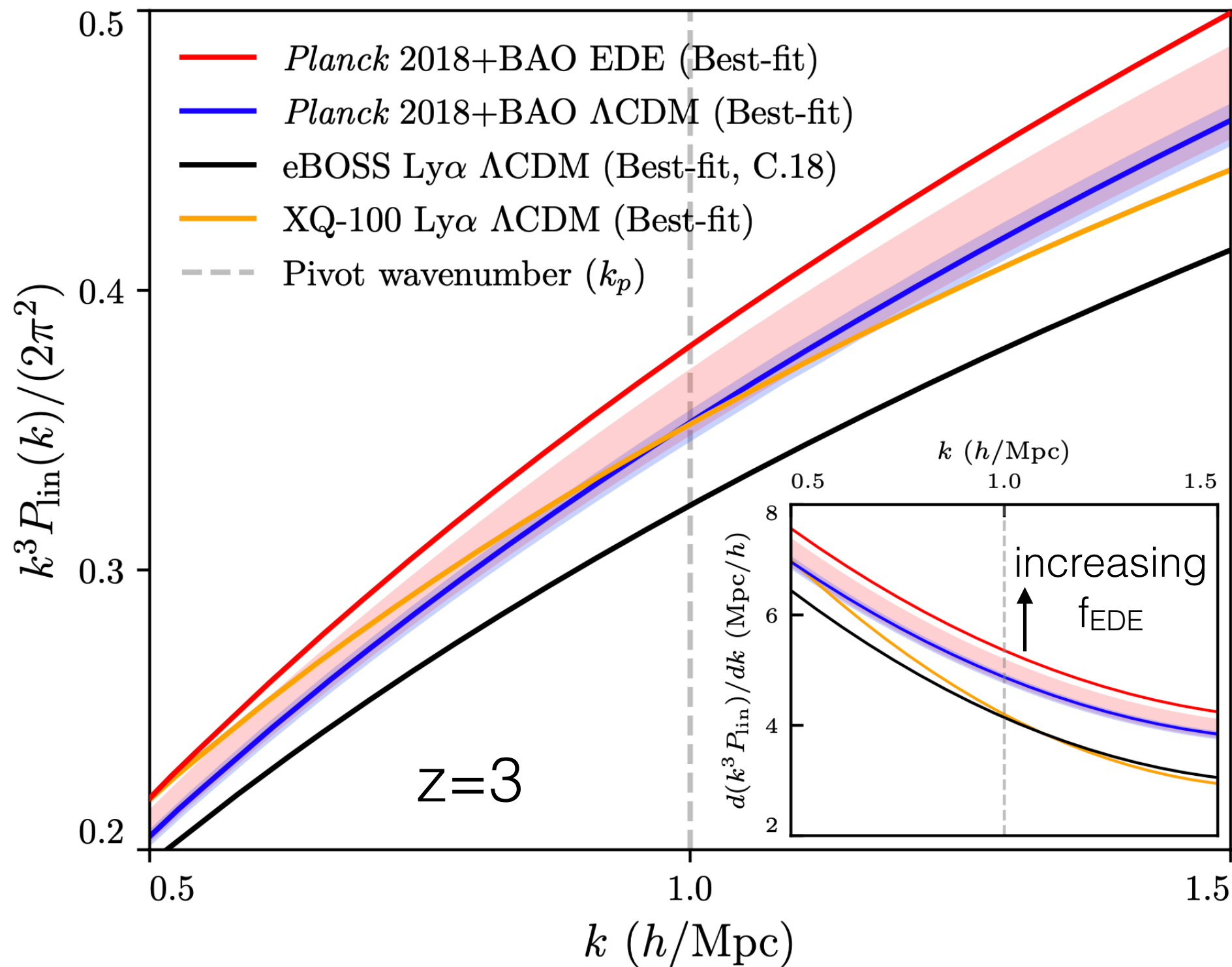


(validation  
using sims  
here)



# Ly $\alpha$ Forest

Inconsistent with prediction of EDE model fit to Planck CMB + BAO



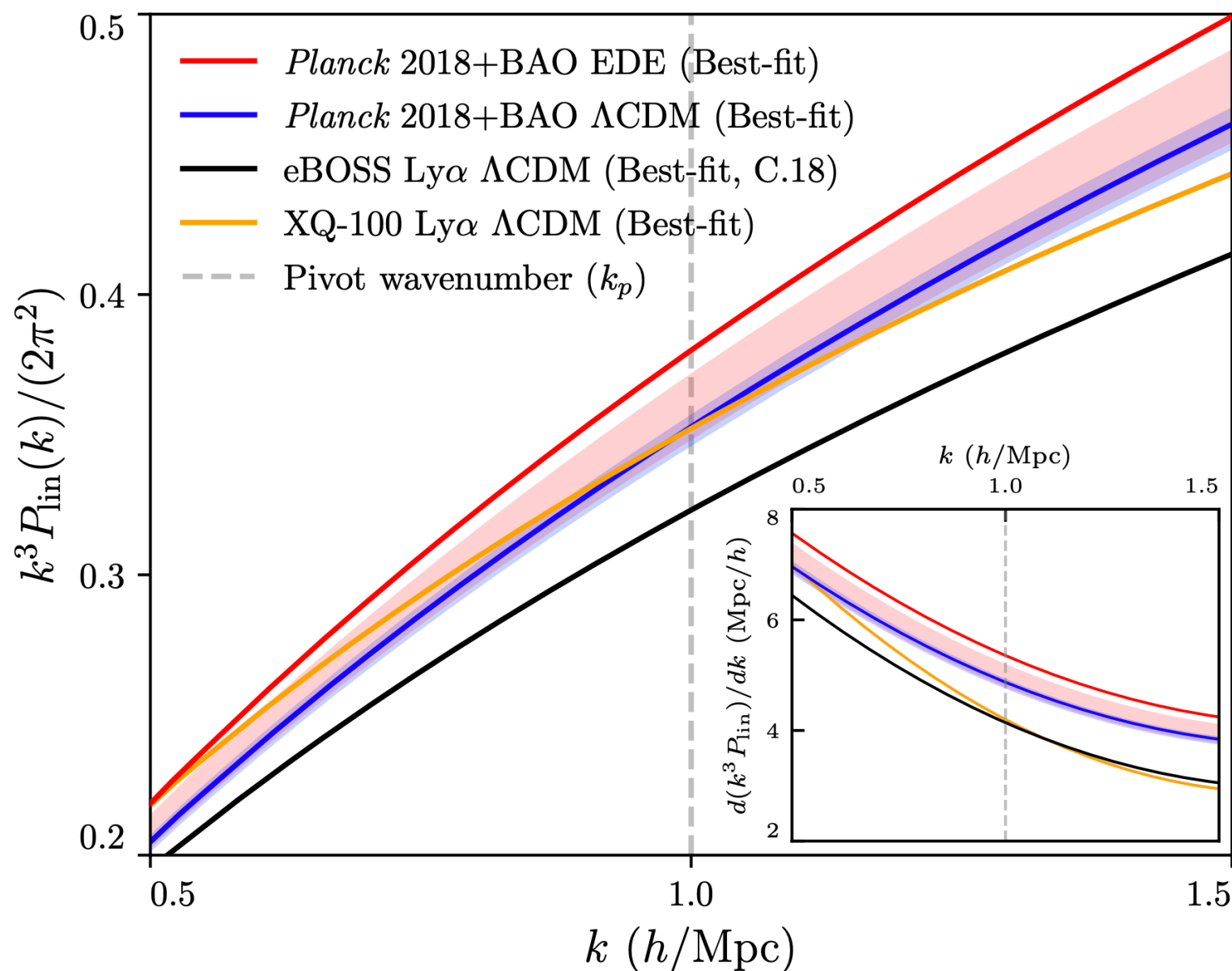
# Origin of EDE Changes to $P(k)$

Colin Hill  
Columbia

Why? Parameter shifts necessary to compensate  
enhanced early ISW effect in EDE cosmologies

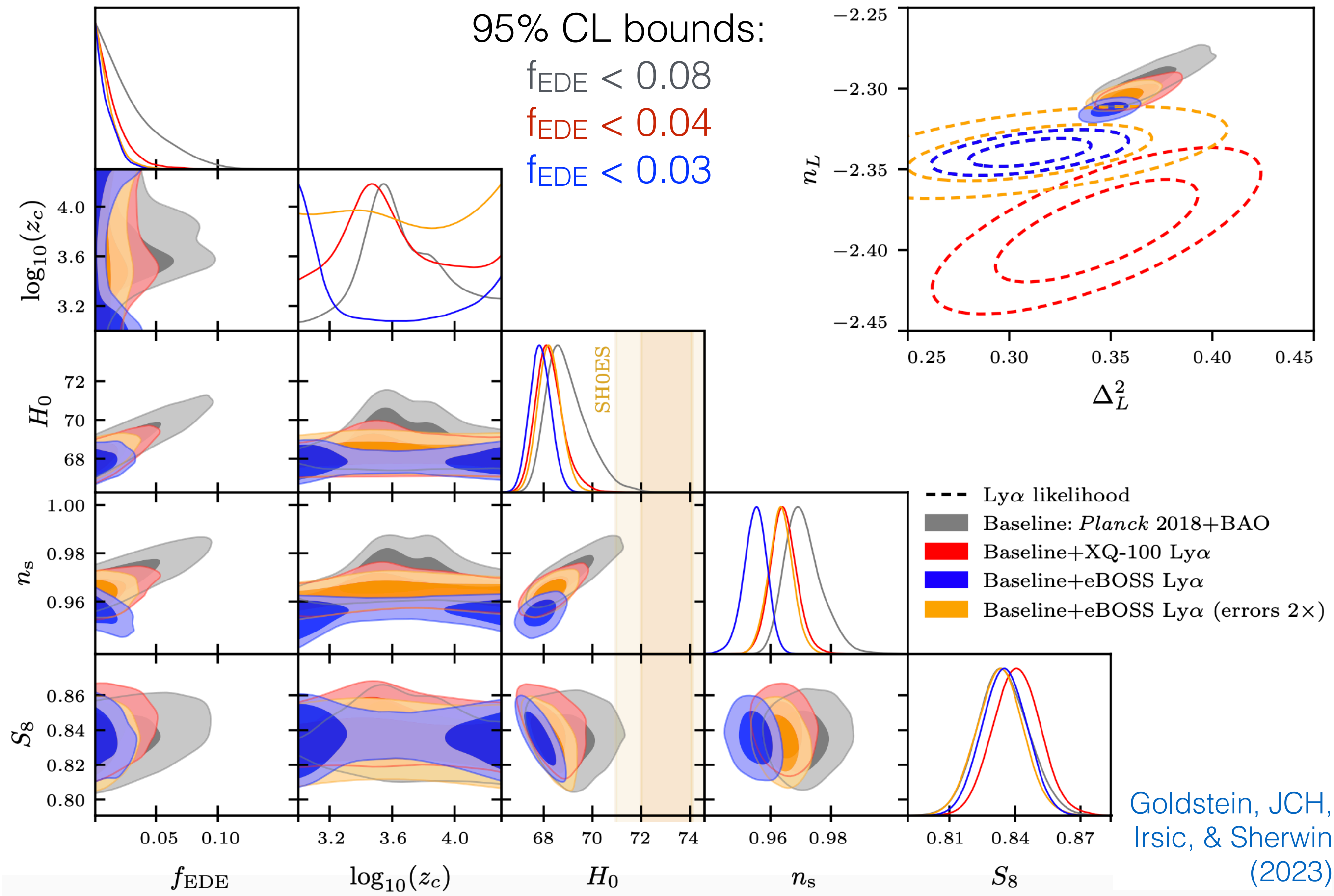
large increase in  $\omega_{\text{cdm}}$  and increase in  $n_s$

which act to increase  $P(k)$  and its slope on Ly $\alpha$  scales



# Ly $\alpha$ Forest: EDE Constraints

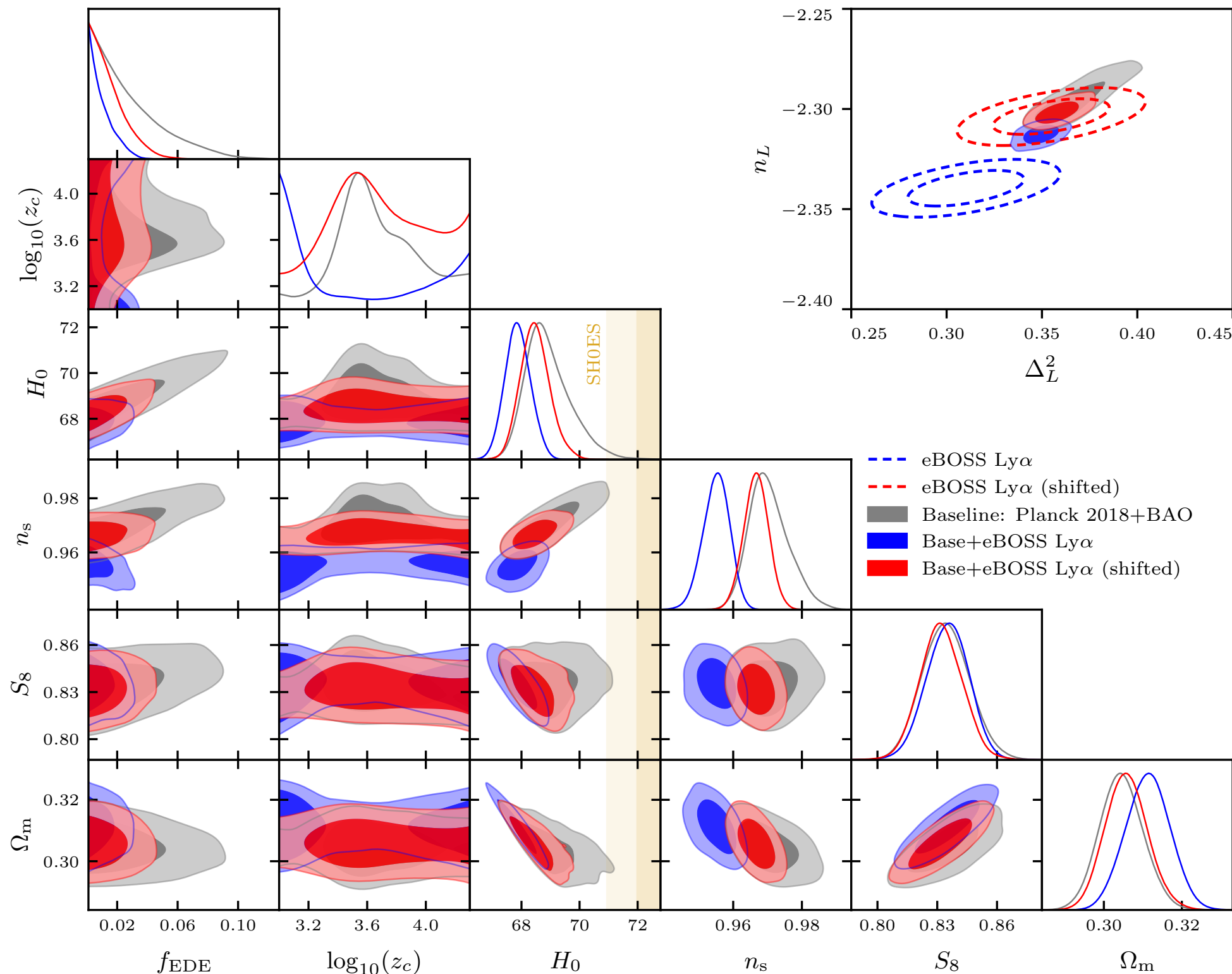
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# Ly $\alpha$ Forest: EDE Constraints

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Columbia

Constraints are still very tight even if Ly $\alpha$  data are artificially shifted



# Ly $\alpha$ Forest: EDE Constraints

Taken at face value, Ly $\alpha$  forest excludes EDE resolution of  $H_0$  tension

- Baseline (CMB+BAO):  $H_0 = 69.0^{+0.6}_{-1.0}$  (best-fit = 70.1) km/s/Mpc
- Baseline + eBOSS:  $H_0 = 67.9 \pm 0.4$  (best-fit = 67.9) km/s/Mpc
- Baseline + XQ-100:  $H_0 = 68.2^{+0.5}_{-0.6}$  (best-fit = 68.2) km/s/Mpc

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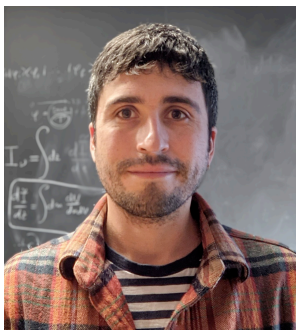
Even direct inclusion of SH0ES ( $H_0 = 73.04 \pm 1.04$  km/s/Mpc) hardly moves the Ly $\alpha$  EDE posteriors

Note that the hydro simulations used to construct Ly $\alpha$  likelihoods *do* include  $P(k)$  that well-represent EDE models

Are the BOSS/eBOSS/XQ100/MIKE/HIRES Ly $\alpha$  data fully secure? There is already some tension w.r.t. Planck even in  $\Lambda$ CDM. Our results motivate close scrutiny!



# Post-Recombination Reheating (PRR)



# $H_0$ and $T_{\text{CMB},0}$

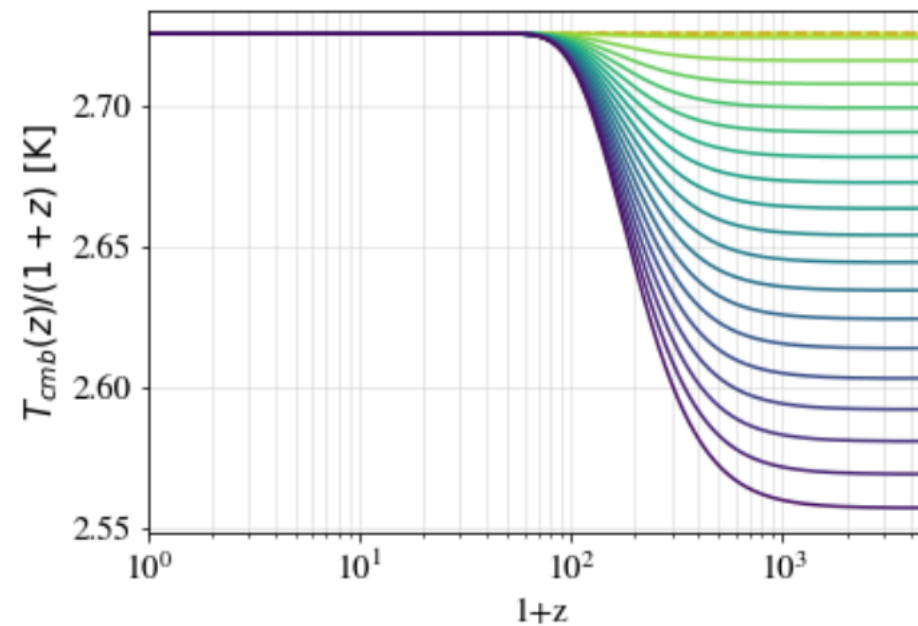
$H_0$  tension or  $T_0$  tension?

Geometric degeneracy: background and linear perturbation evolution depend (almost) only on parameter combination  $H_0 T_{\text{CMB},0}^{1.2}$

- Cosmology would be very different without COBE-FIRAS data!
- Ignoring FIRAS, Planck+SH0ES can be fit with  $T_{\text{CMB},0} = 2.56 \pm 0.05$  K
- BAO breaks degeneracies: Planck + BAO yield  $T_{\text{CMB},0} = 2.71 \pm 0.02$  K
- FIRAS result:  $T_{\text{CMB},0} = 2.72548 \pm 0.00057$  K
- Can we build on this idea while maintaining agreement with FIRAS?

# Did the Universe (Slightly) Reheat after Recombination?

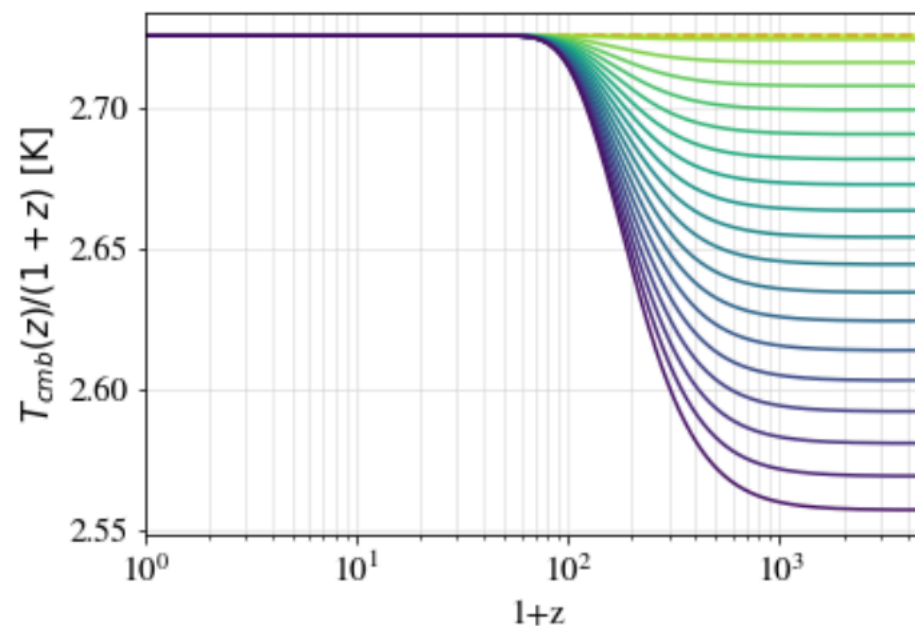
- Suppose  $T_{\text{CMB}}(z^*) < T_{\text{CMB}}(z^*)^{\Lambda\text{CDM}}$ , but a process injects energy at  $z < z^*$



Direct constraints  
on  $T_{\text{CMB}}(z)$  only  
exist at  $z < \sim 3$

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Direct constraints  
on  $T_{\text{CMB}}(z)$  only  
exist at  $z < \sim 3$

- To keep  $k_{\text{eq}}$  fixed,  $\Omega_m$  must decrease, and hence  $S_8$  decreases
- A conspiracy of integrals leads to higher  $H_0$ : the sound horizon is not decreased (in fact it increases due to lower  $H(z)$  at early times)

$$r_s = \int_{z_*}^{\infty} \frac{dz}{H(z)} c_s(z)$$

- To keep  $\theta_s^*$  fixed,  $D_A^*$  must increase, but since  $\Omega_m$  decreases, one must increase  $H_0$  to compensate otherwise increased value of the integral:

$$D_A^* = \int_0^{z_*} c dz / H(z)$$

# Did the Universe (Slightly) Reheat after Recombination?

Concrete model: sub-component of CDM decays into photons after  $z^*$

- Background evolution similar to usual decaying DM- $\rightarrow$ DR:

$$\begin{aligned}\rho'_{\text{DCDM}} &= -3aH\rho_{\text{DCDM}} - a\Gamma\rho_{\text{DCDM}} \\ \rho'_\gamma &= -4aH\rho_\gamma + a\Gamma\rho_{\text{DCDM}},\end{aligned}$$

- New parameters:

$\Gamma$  = decay rate

$\omega_{\text{DCDM,ini}}$  = initial decaying CDM density

( $T_{\text{CMB,ini}}$  = initial CMB monopole temperature) — not really new

- Perturbation evolution equations for photons acquire new terms not present in  $\Lambda$ CDM or usual DCDM- $\rightarrow$ DR model

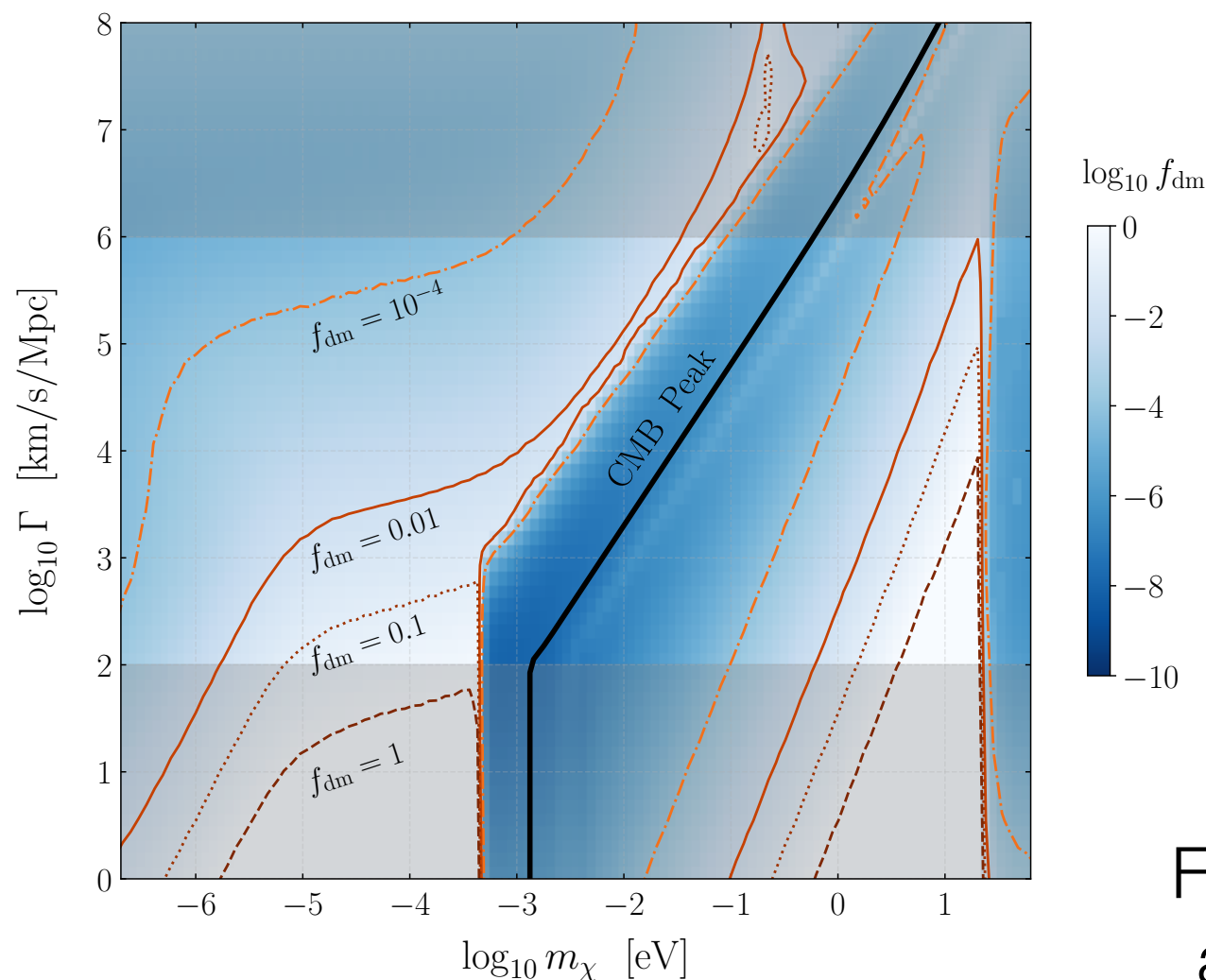
# Did the Universe (Slightly) Reheat after Recombination?

- Key feature: only a tiny amount of decaying CDM is needed to increase  $T_{\text{CMB}}$  by the necessary magnitude (e.g., for decay at  $z=22$ , only 0.02% of CDM decaying will increase  $T_{\text{CMB}}$  by 1%)



# Did the Universe (Slightly) Reheat after Recombination?

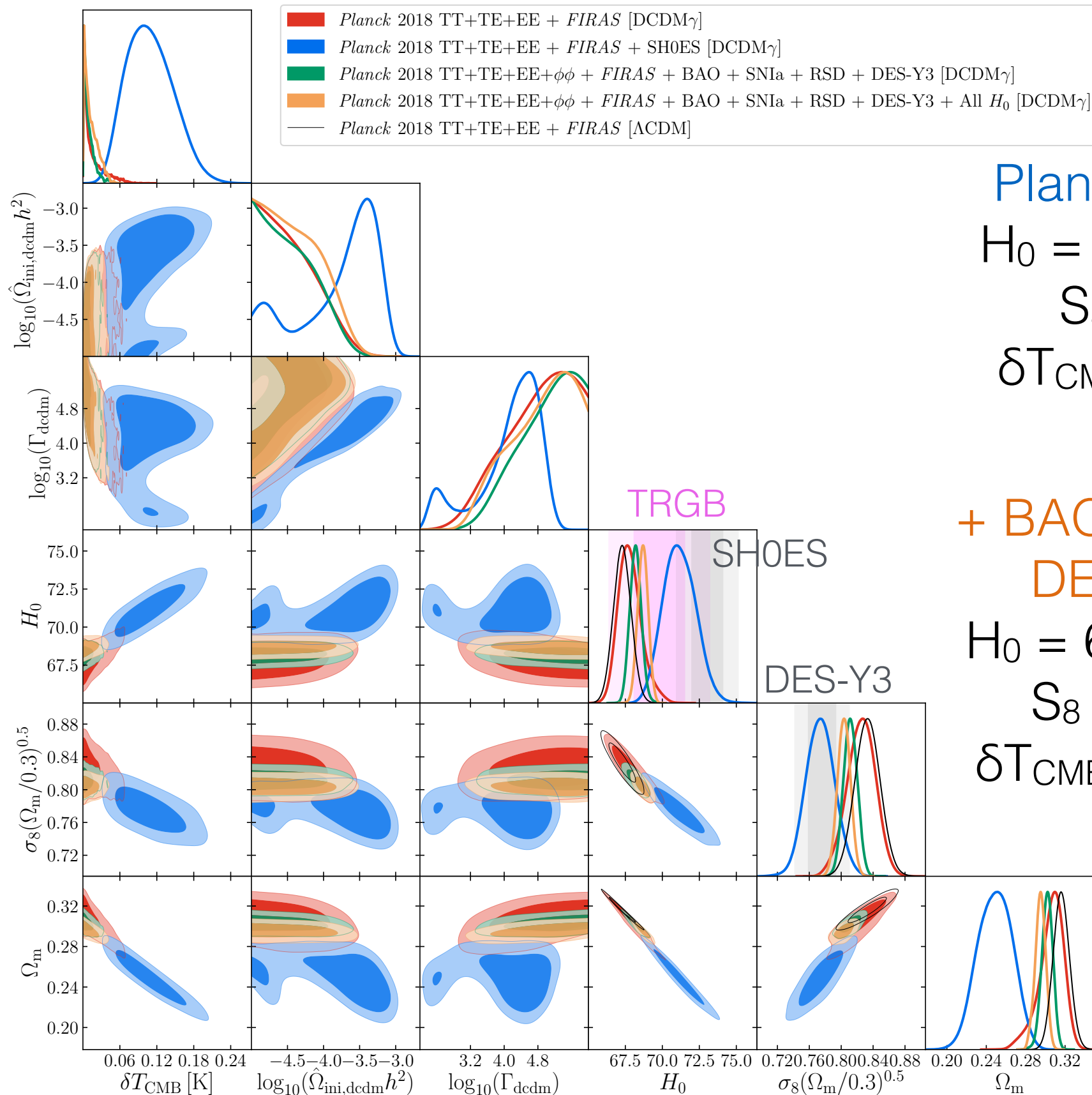
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- What about spectral distortion constraints?



$f_{\text{DM}}$  = fraction  
of DM that  
decays

For simplicity here, we  
assume that injected  
photons are thermal

# PRR Analysis



Planck + FIRAS + SH0ES:

$$H_0 = 71.2 \pm 1.1 \text{ km/s/Mpc}$$

$$S_8 = 0.774 \pm 0.018$$

$$\delta T_{\text{CMB}} = 0.109^{+0.033}_{-0.044} \text{ K}$$

+ BAO + SNIa +  $\phi\phi$  + RSD +  
DES-Y3 + MCP + SBF:

$$H_0 = 68.7 \pm 0.35 \text{ km/s/Mpc}$$

$$S_8 = 0.8035 \pm 0.0081$$

$$\delta T_{\text{CMB}} < 0.0342 \text{ K [95\% CL]}$$

$\chi^2_{\text{Planck}}$  is as  
good as that of  
 $\Lambda$ CDM fit to  
Planck alone

# PRR Takeaways

- Main obstruction to success of the model: conflict with  $\Omega_m$  constraint from BAO and SNIa
- One tweak: allow  $N_{\text{eff}}$  to 'restore' early-universe radiation density back to its normal value, and mitigate decrease in  $\Omega_m$  (thus giving up  $S_8$  fix and lessening increase in  $H_0$ ) — such a model fits data better than  $\Lambda\text{CDM}+N_{\text{eff}}$ , but only slightly
- Most plausible route to avoid SD constraints: suppress the spectrum at frequencies  $< 60$  GHz in early universe (non-trivial)
- Key points:
  - There are large swathes of cosmic history where (semi-)dramatic changes to the model could still lurk
  - Seemingly small changes ( $\delta T_{\text{CMB}} \sim 10\text{-}100$  mK) can have big effects
  - We should measure CMB spectral distortions much better than FIRAS!Strong motivation for PIXIE, FOSSIL, BISOU, etc.

# Next: ACT DR6

(target: later this year)

# ACT DR6 Forecasts

***ACT TT + TE + EE : precision cosmology beyond Planck***

	ACT DR4	ACT DR4 + WMAP	Planck	Planck + ACT DR6
$\sigma(H_0)$	1.5	1.1	0.5	0.4
$\sigma(n_s)$	0.015	0.006	0.004	0.003
$\sigma(N_{\text{eff}})$	0.4	0.3	0.2	0.1

Large improvements in beyond- $\Lambda$ CDM parameters:  
~2x increase in sensitivity to new light relic particles



***PRELIMINARY FORECAST***

**Upcoming ACT DR6 precision cosmology constraints will surpass those from Planck ( $H_0$ ,  $N_{\text{eff}}$ ,  $\Sigma m_\nu$ ,  $\sigma_8$ , + beyond- $\Lambda$ CDM models) — stay tuned!**

# Discovering EDE in the CMB?

Colin Hill  
Columbia



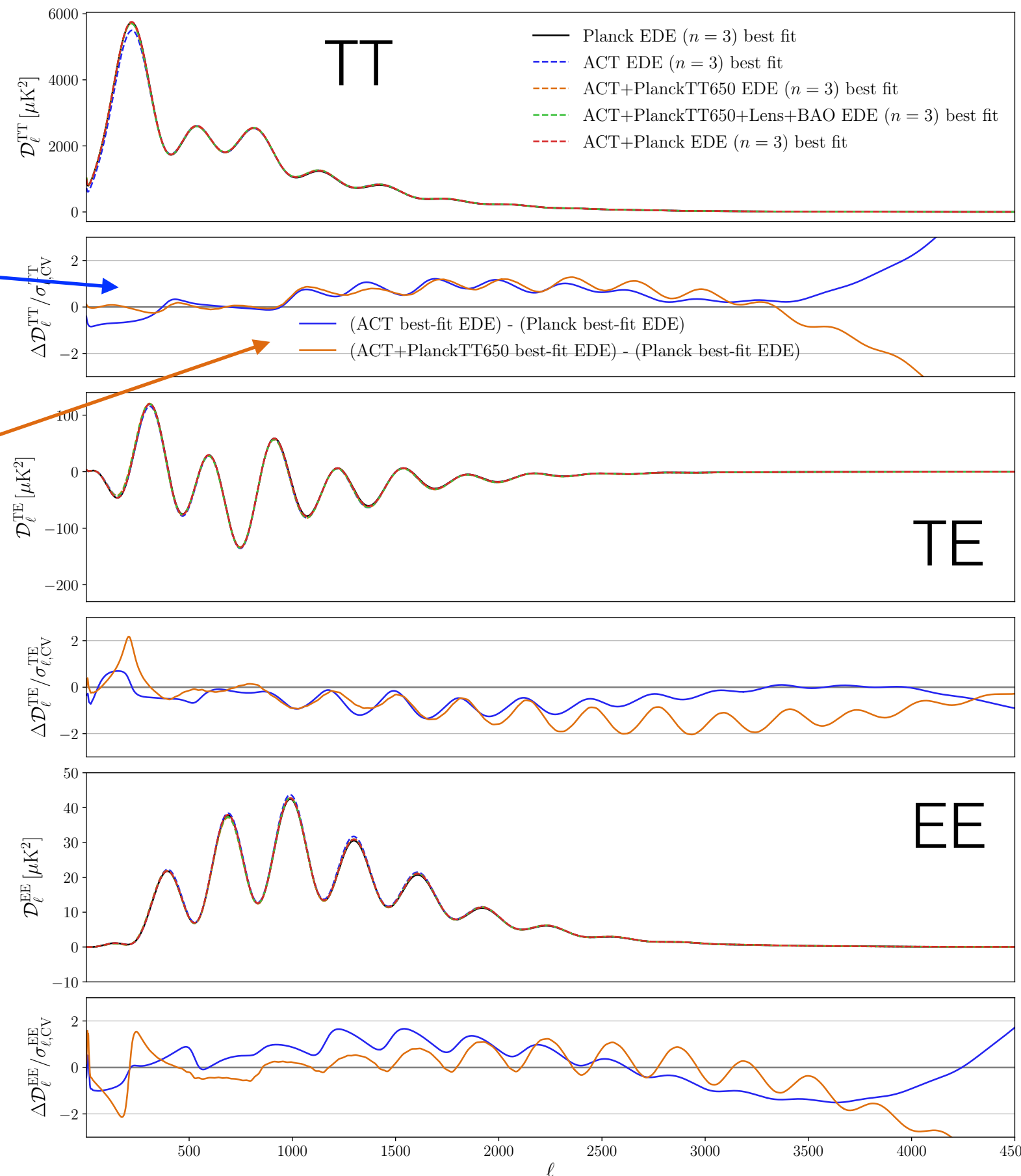
# Discovering EDE in the CMB?

Colin Hill  
Columbia

ACT best-fit EDE -  
Planck EDE

ACT+P18TT650 EDE -  
Planck EDE

Imminent potential  
discovery with upcoming  
ACT DR6 (~2023): the  
models shown  
here can be  
distinguished at  $\sim 15\text{-}20\sigma$

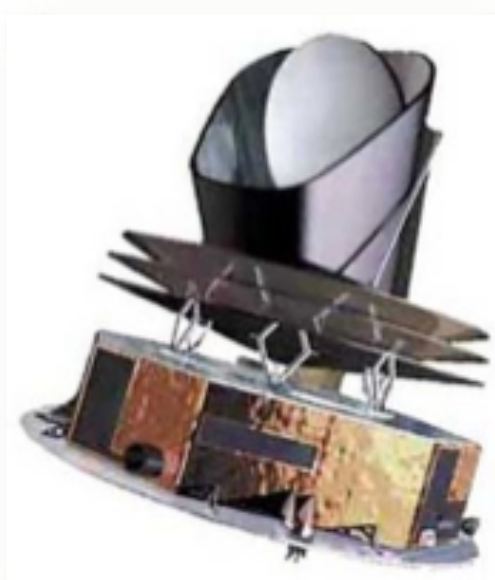


		SO-Pre		SO-Nominal Operations					
2020	2021	2022	2023	2024	2025	2026	2027	2028	

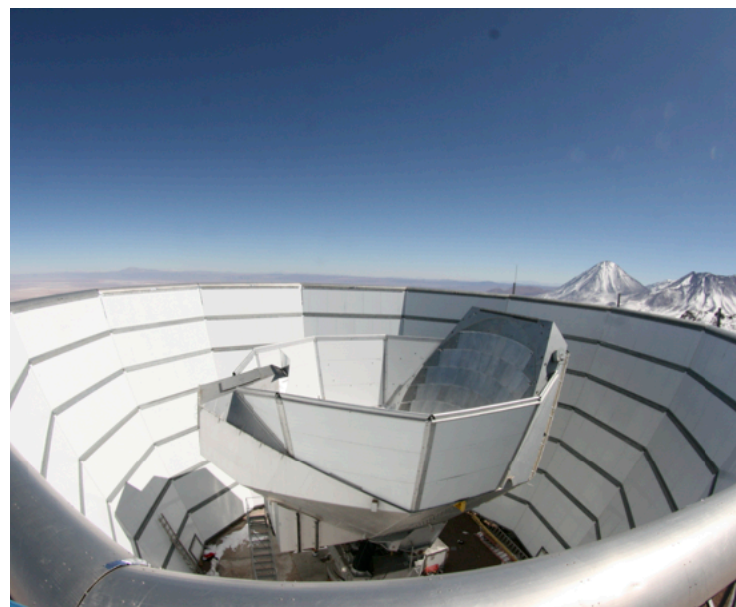
Planck



ACT

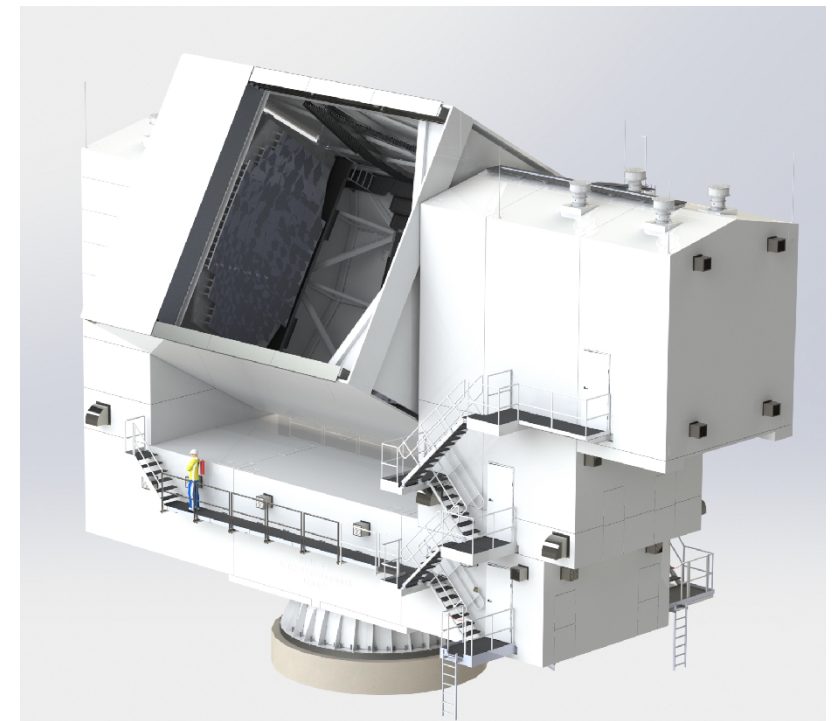
**SO** Large Aperture Telescope

Final data 2018  
100% sky  
0.35 — 10 mm (9 bands)  
5 — 33' resolution



Observations until 2022  
40% sky  
Noise ~3 times < Planck  
1.4 — 10 mm (5 bands)  
1 — 7' resolution

*[South Pole Telescope - same  
timeframe]*



Observations 2024 - ~29  
40% sky  
Noise ~3 times < ACT  
1 — 10 mm (6 bands)  
1 — 7' resolution

*[CMB-S4 would start observing after  
this, with multiple telescopes]*

# **Detecting faint traces of universe's explosive birth is aim of NSF-supported Advanced Simons Observatory**

Funded by NSF in  
May 2023

- \$52.7M NSF investment to double LAT detector count; build robust data pipeline (including transient alerts); build solar array at site to generate 70% of energy for the observatory
- SO observations: 2024-2029
- ASO observations: 2028-2033
- Significant gains in scientific capabilities
- PI: Mark Devlin (Penn)
- Co-PIs: Jo Dunkley (Princeton), Jeff McMahon (Chicago), Suzanne Staggs (Princeton)
- Co-Project Scientists: JCH (Columbia), Susan Clark (Stanford)



# Take-Home Messages

- 1) Cosmic discordance? Observational situation unclear, but viable models will be confirmed or refuted imminently
- 2) ACT and Planck prefer somewhat different EDE model parameters, with ACT yielding higher  $f_{\text{EDE}}$  and  $H_0$
- 3) Challenge: Ly $\alpha$  forest severely constrains canonical EDE
- 4) Early-universe  $H_0$  /  $S_8$  resolutions generically predict clear deviations from  $\Lambda$ CDM in the CMB — imminently testable with ACT DR6, SPT-3G, Simons Observatory



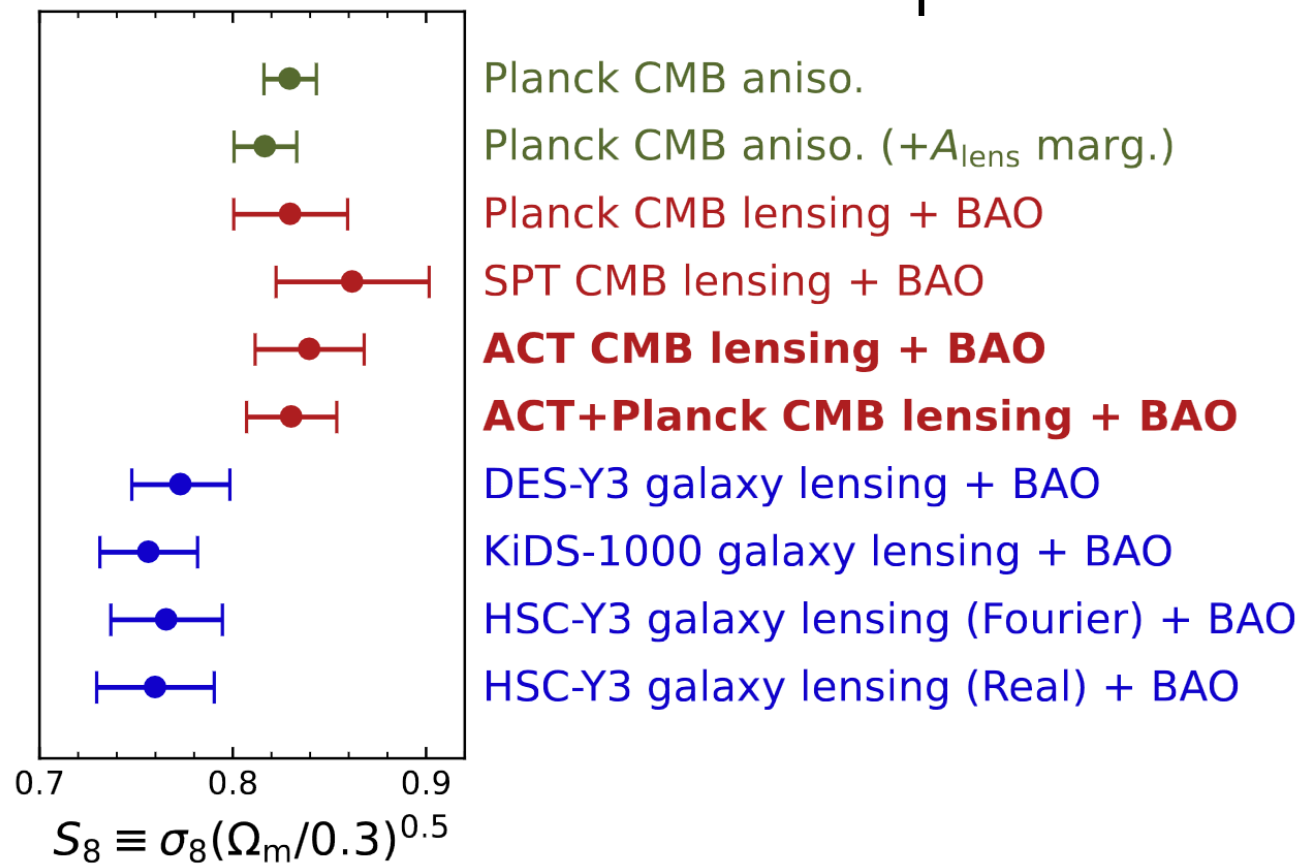
**Thanks!**



# Bonus

# ACT DR6 CMB Lensing

## Implications for EDE?

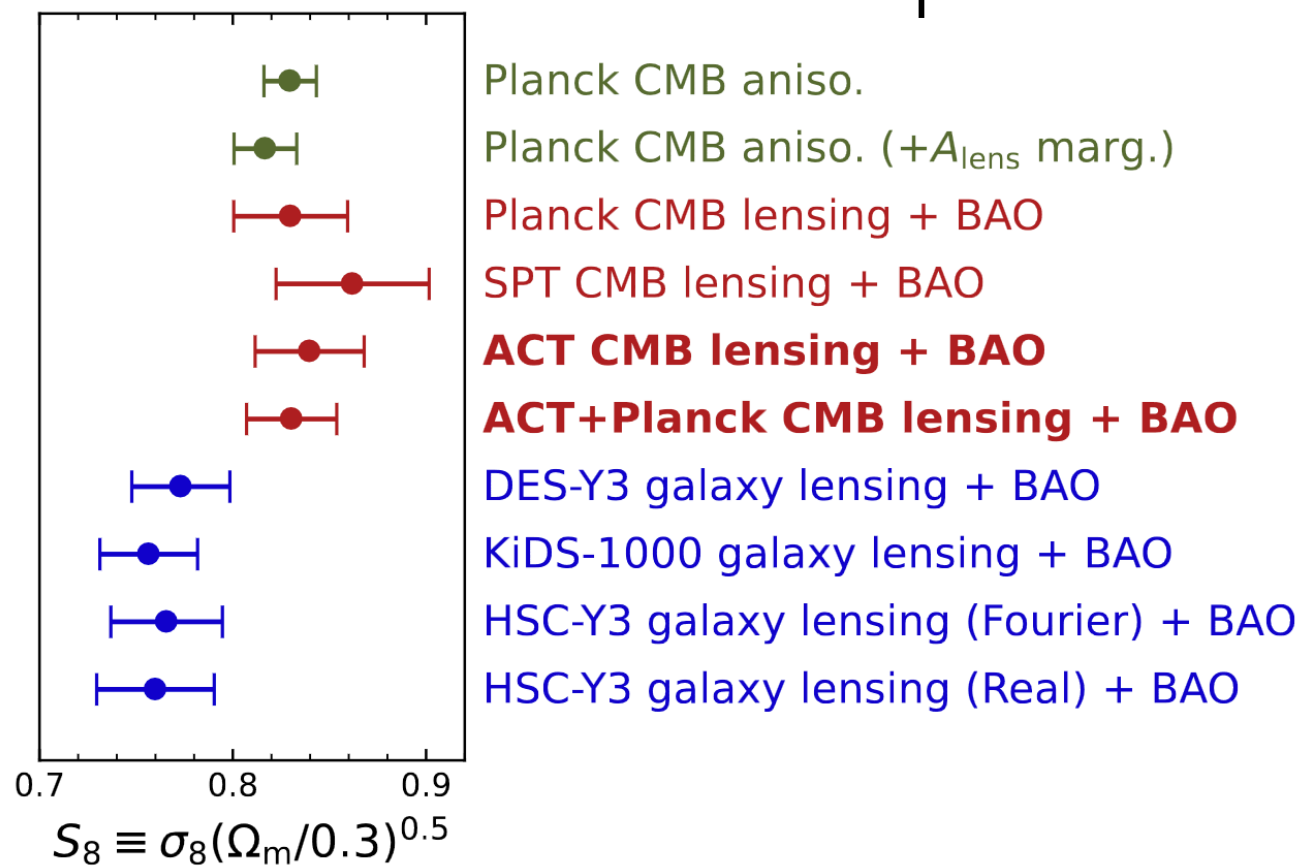


$S_8$  is not low — reduces  
discordance with EDE  
prediction from  
CMB+SH0ES fit



# ACT DR6 CMB Lensing

## Implications for EDE?

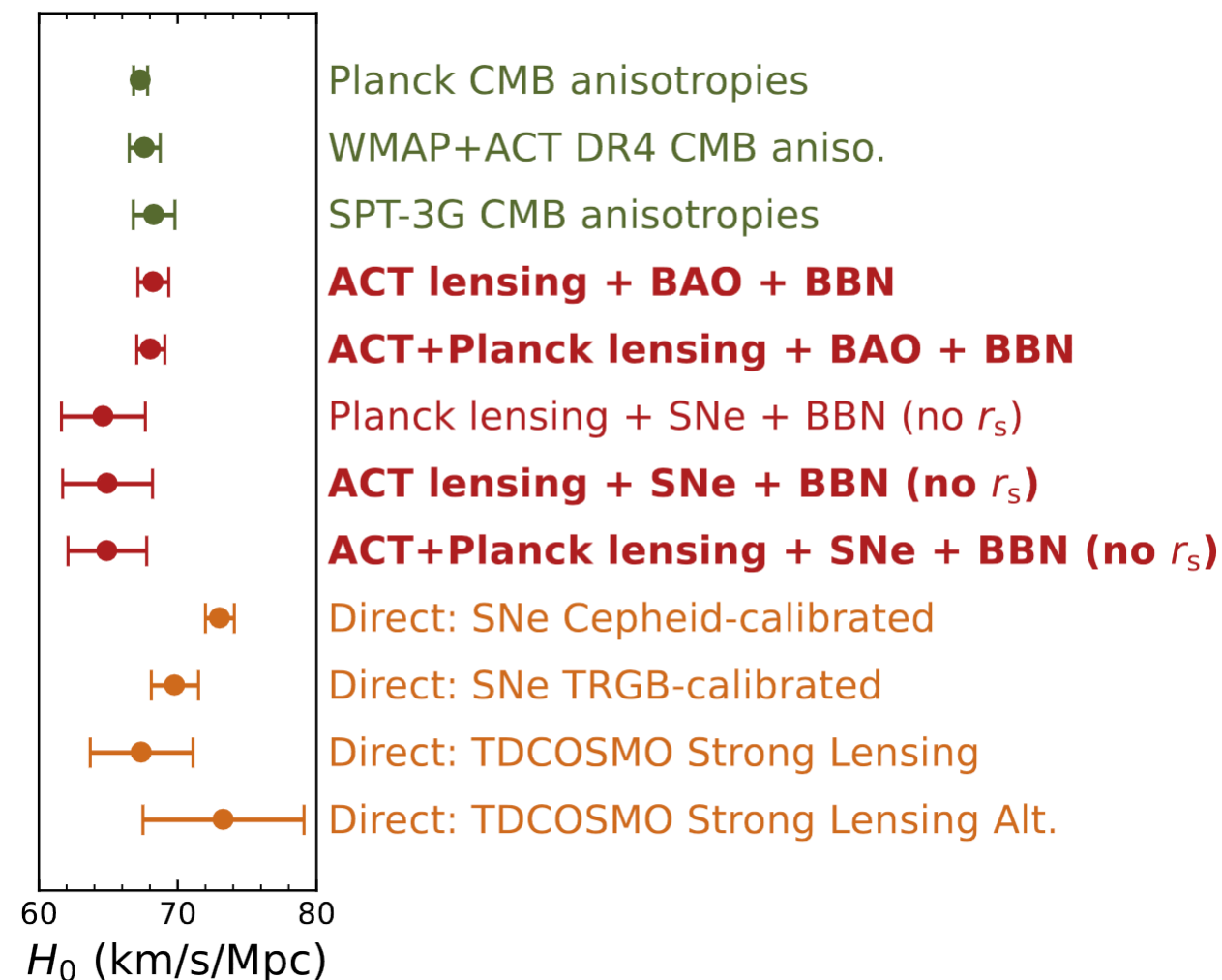


$S_8$  is not low — reduces discordance with EDE prediction from CMB+SH0ES fit

... But sound-horizon-free  $H_0$  constraint is  $3\sigma$  lower than SH0ES

$$H_0 = 64.9 \pm 2.8 \text{ km s}^{-1} \text{ Mpc}^{-1}$$

Success for  $\Lambda$ CDM at  $z > z^*$



# PRR: Particle Physics Model(s)

Colin Hill  
Columbia

Constraints on light particles coupled to photons from, e.g., white dwarfs can be evaded via models with an excited dark matter state

Consider DM with non-zero dipole moment, coupled to SM sector through a kinetically mixed massive dark photon (DP)

The DP allows for transitions between ground and excited states of the DM

If the energy splitting of the states is  $\sim 0.1$  eV, the excited state is metastable for  $\sim 1$ -10 Myr as we need in the PRR scenario

If  $m_{\text{DP}} > \text{MeV}$ , it is not produced in stars or supernovae

In general our results suggest that cosmological implications of such scenarios should be considered

Most plausible route to avoid SD constraints: suppress the spectrum at frequencies  $< 60$  GHz in early universe (non-trivial)

(e.g.) Baryakhtar+ (2020)

# Consistent with SPT-3G?

Colin Hill  
Columbia

Analysis using public SPT-3G TE/EE data

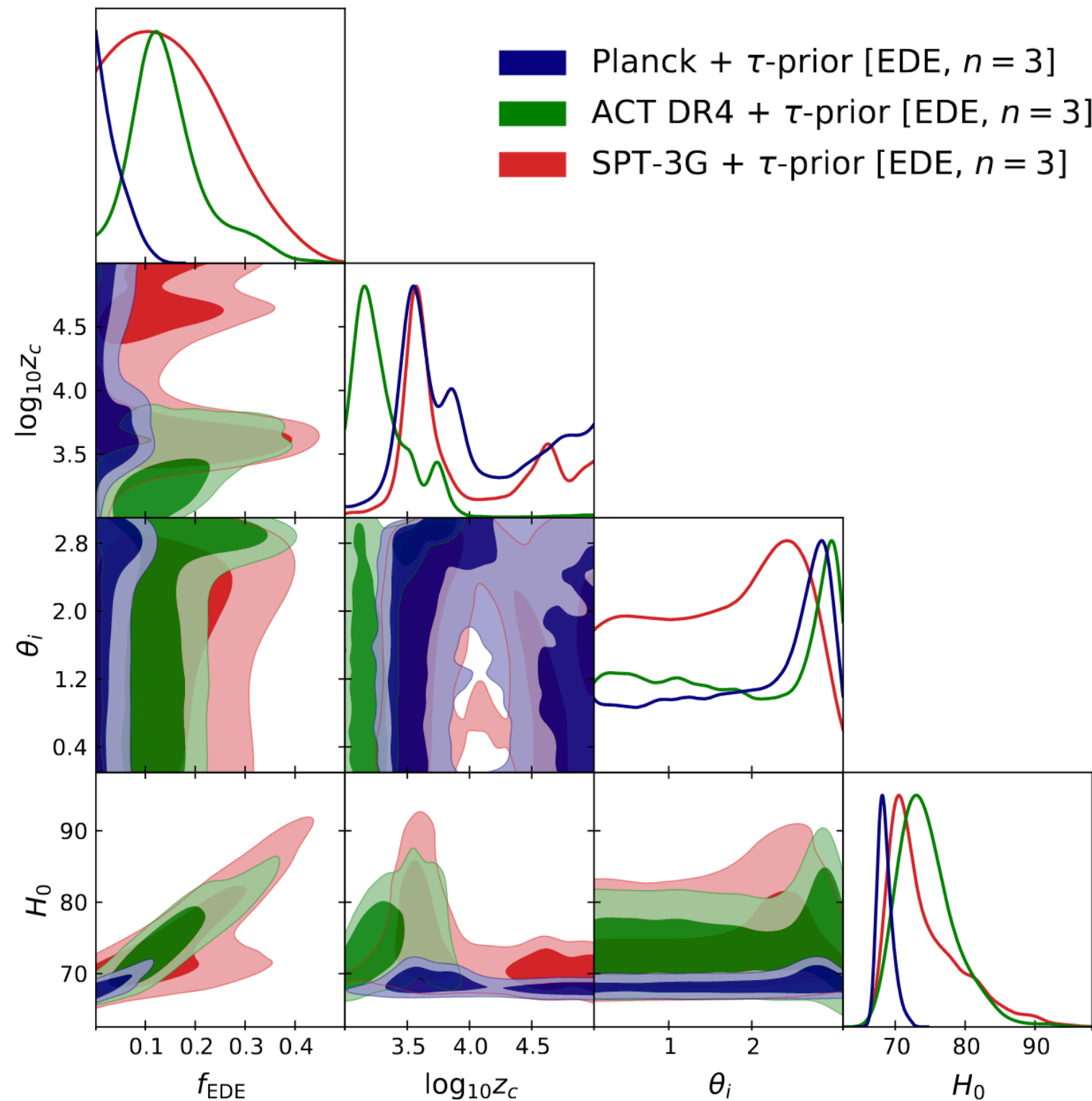
<https://arxiv.org/abs/2112.10754>; Phys. Rev. D 105, 083519

La Posta, Louis, Garrido, **JCH** (2022) ; SPT-3G data from Dutcher et al. (2021)

# Consistent with SPT-3G

Colin Hill  
Columbia

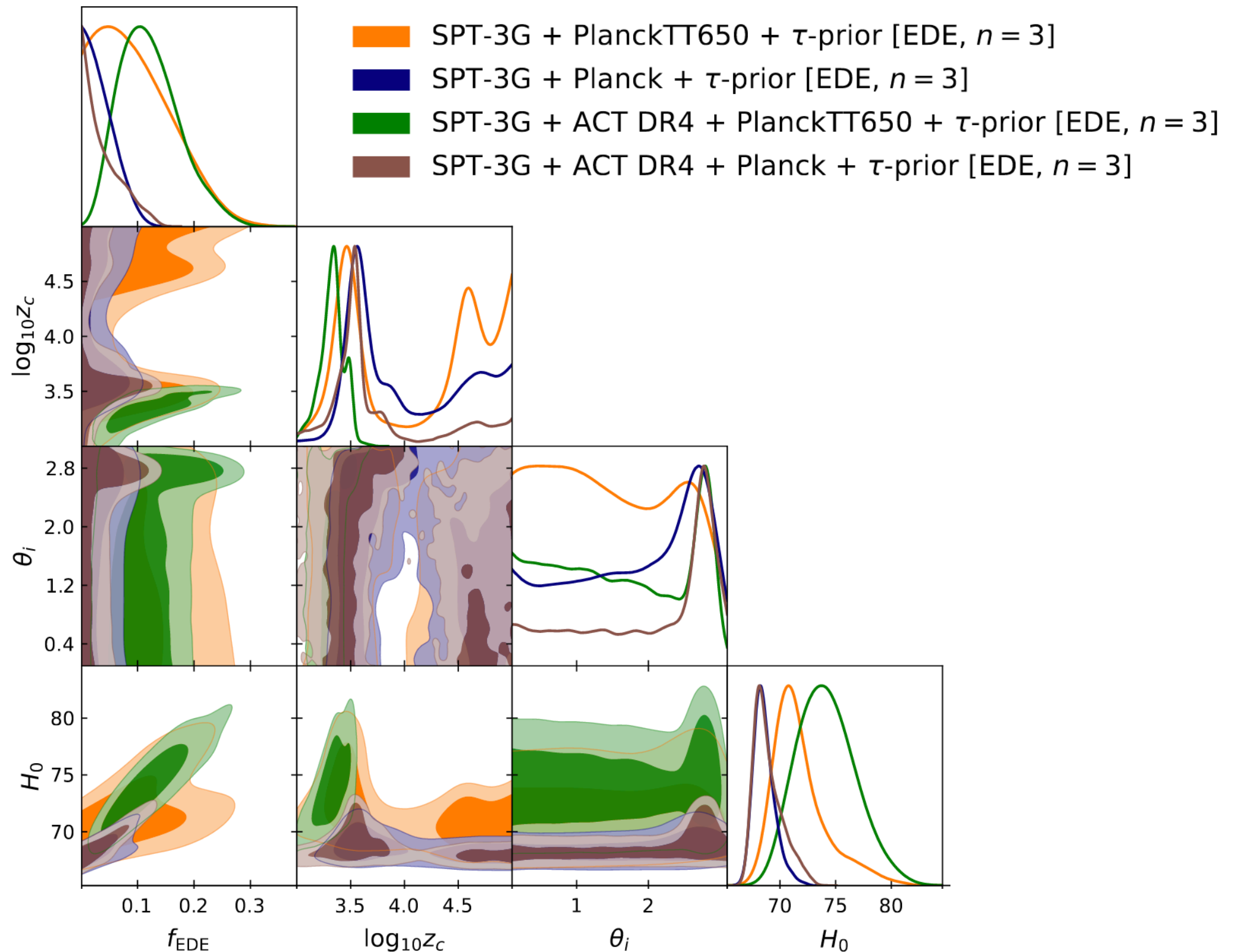
Analysis using public SPT-3G TE/EE data



# Consistent with SPT-3G

Colin Hill  
Columbia

Analysis using public SPT-3G TE/EE data



# Consistent with SPT-3G

Analysis using public SPT-3G TE/EE data

Inclusion of full Planck TT data still dominates overall constraining power and removes preference for non-zero EDE

Parameters	SPT-3G + Planck	SPT-3G + PlanckTT650 + ACT DR4	SPT-3G + Planck + ACT DR4
$f_{\text{EDE}}$	$< 0.088$	$0.121^{+0.040}_{-0.064}$	$< 0.107$
$H_0$ [km/s/Mpc]	$68.6^{+0.7}_{-1.1}$	$74.2^{+2.3}_{-3.0}$	$68.9^{+0.7}_{-1.6}$

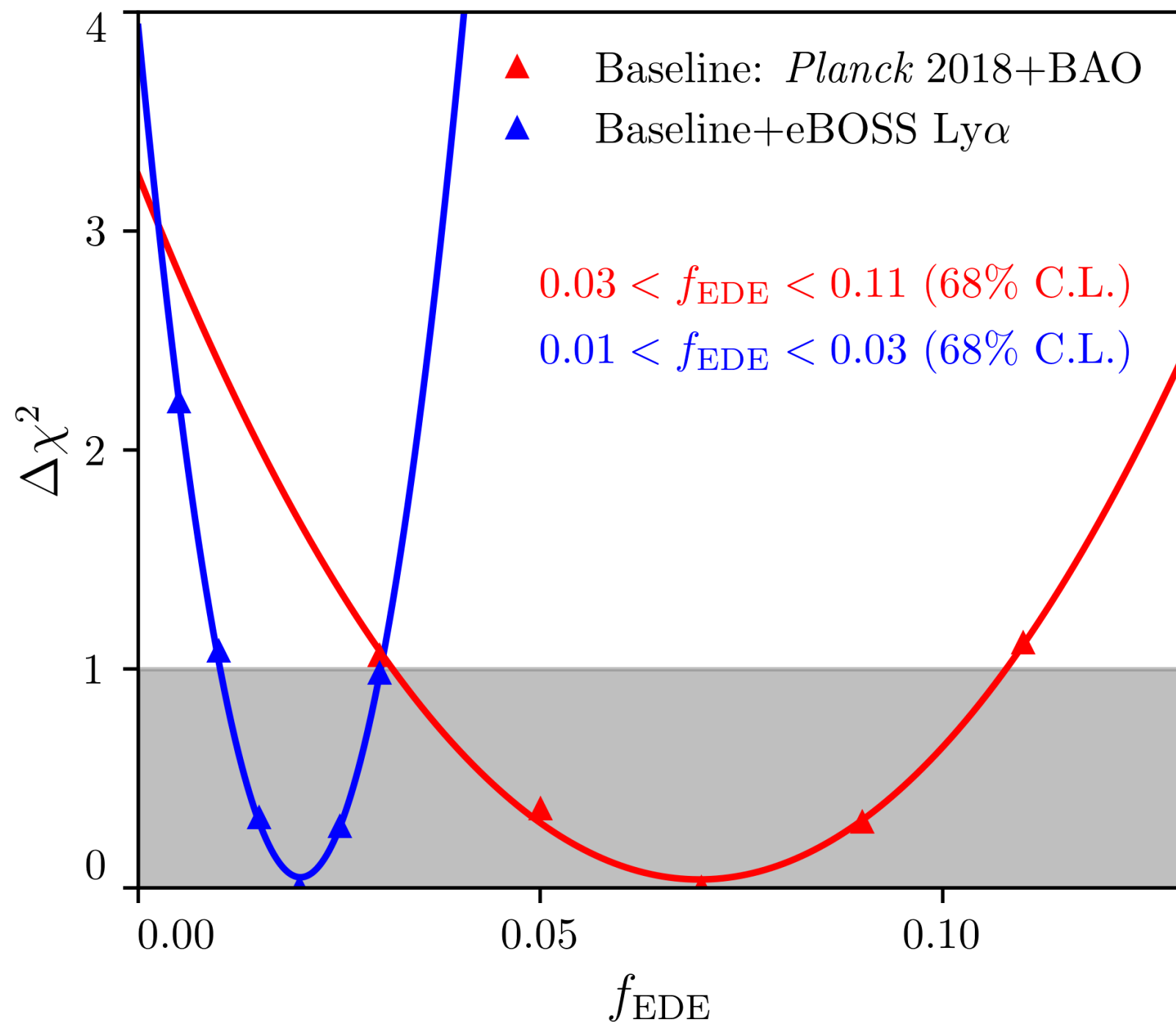
↓  
2.6 $\sigma$  hint

Upcoming data from ACT + SPT will be very interesting!



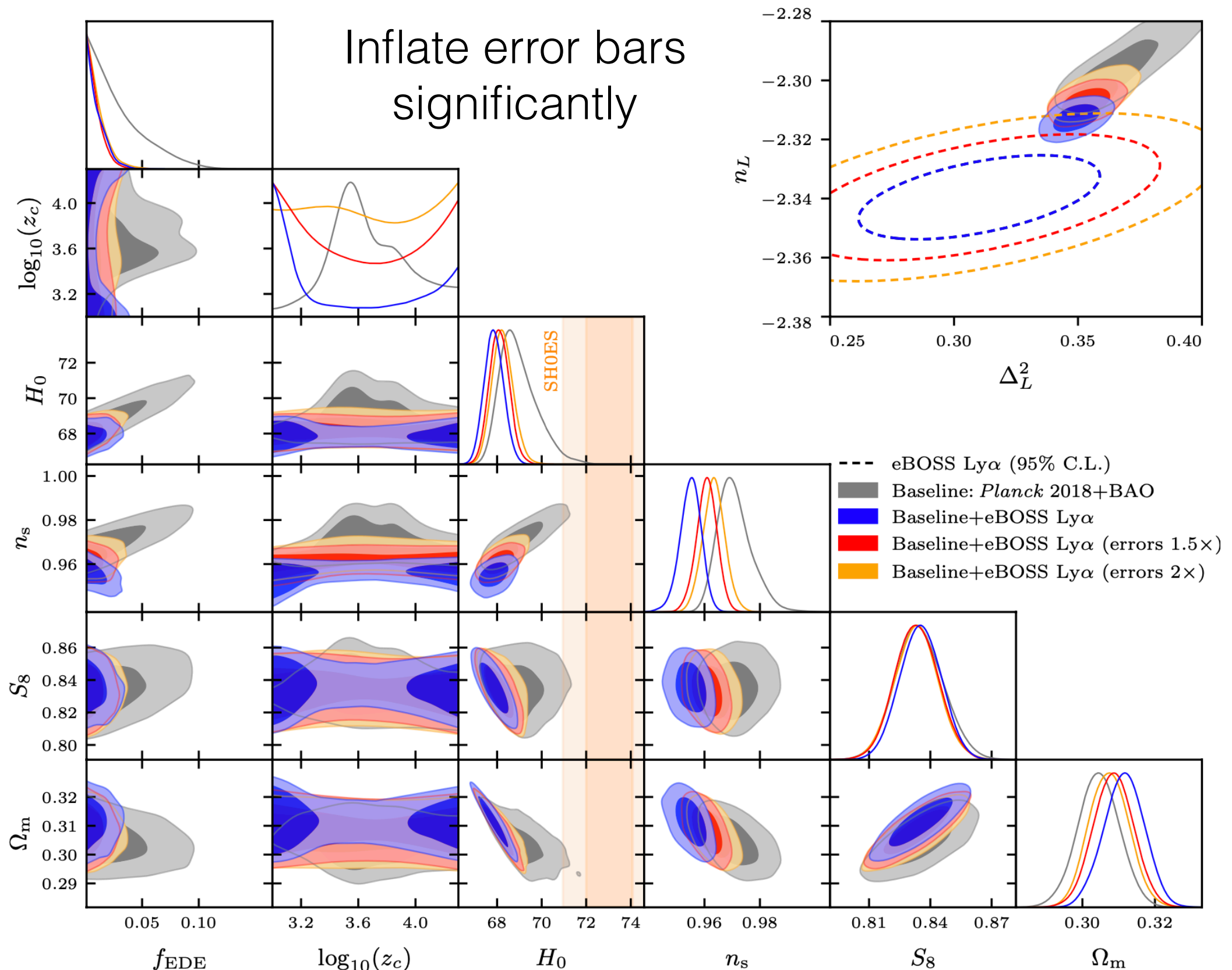
# Ly $\alpha$ Forest: EDE Constraints

No biases due to “prior volume effects”: profile likelihood analysis yields similarly tight constraints



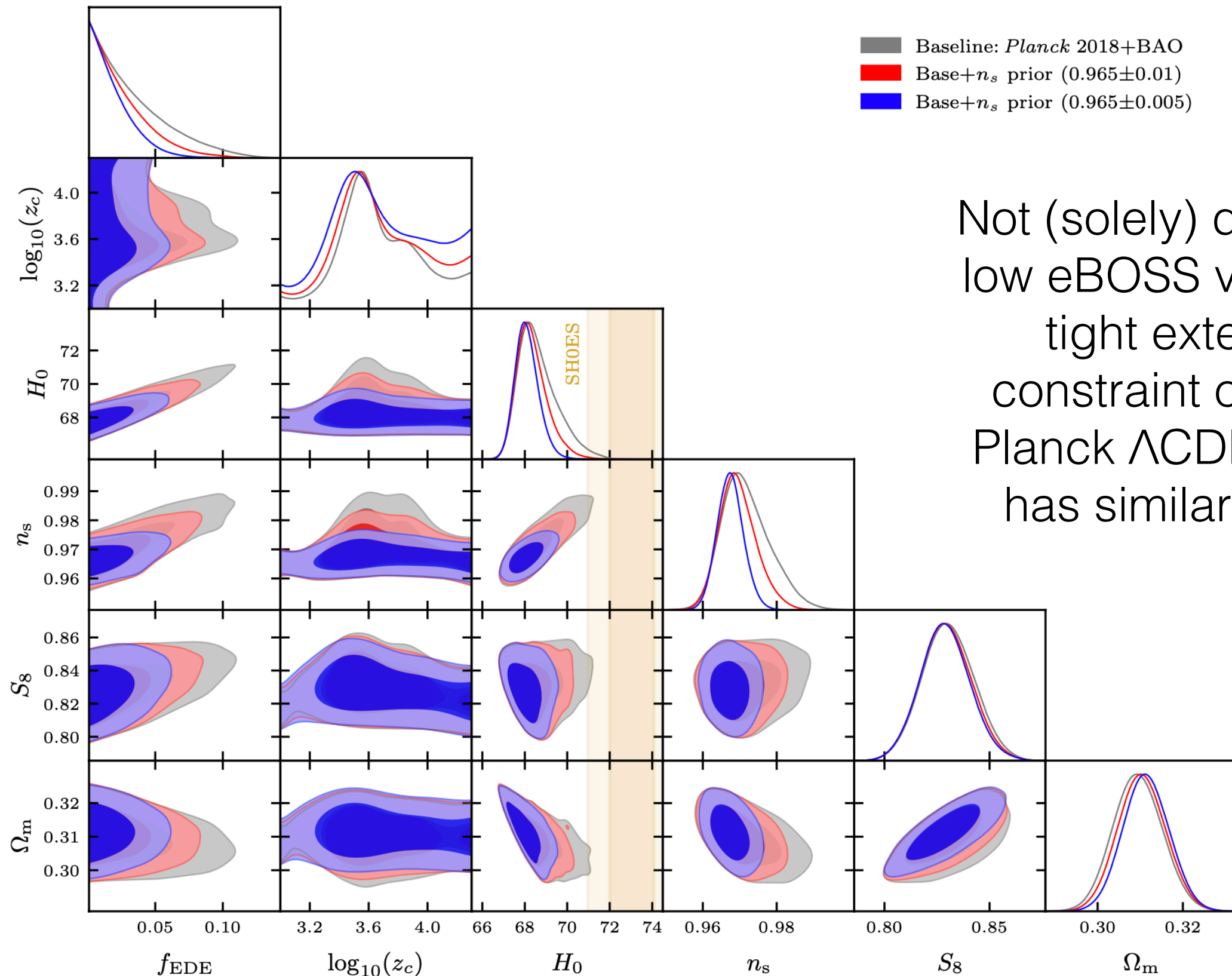
# Ly $\alpha$ EDE Validation

Colin Hill  
Columbia



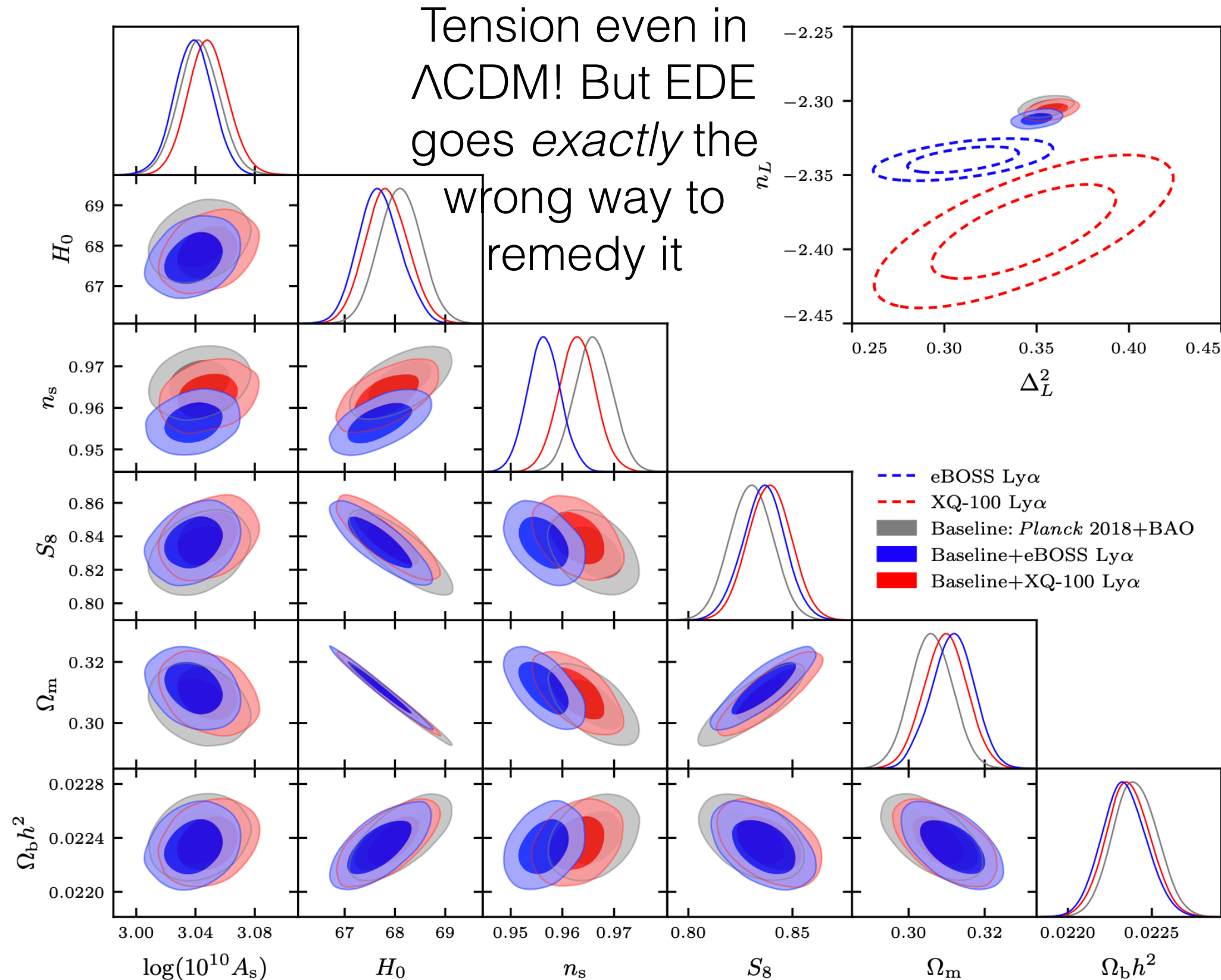
# Ly $\alpha$ EDE Validation

Colin Hill  
Columbia



Not (solely) driven by  
low eBOSS values: a  
tight external  
constraint on  $n_s$  at  
Planck  $\Lambda$ CDM value  
has similar effect

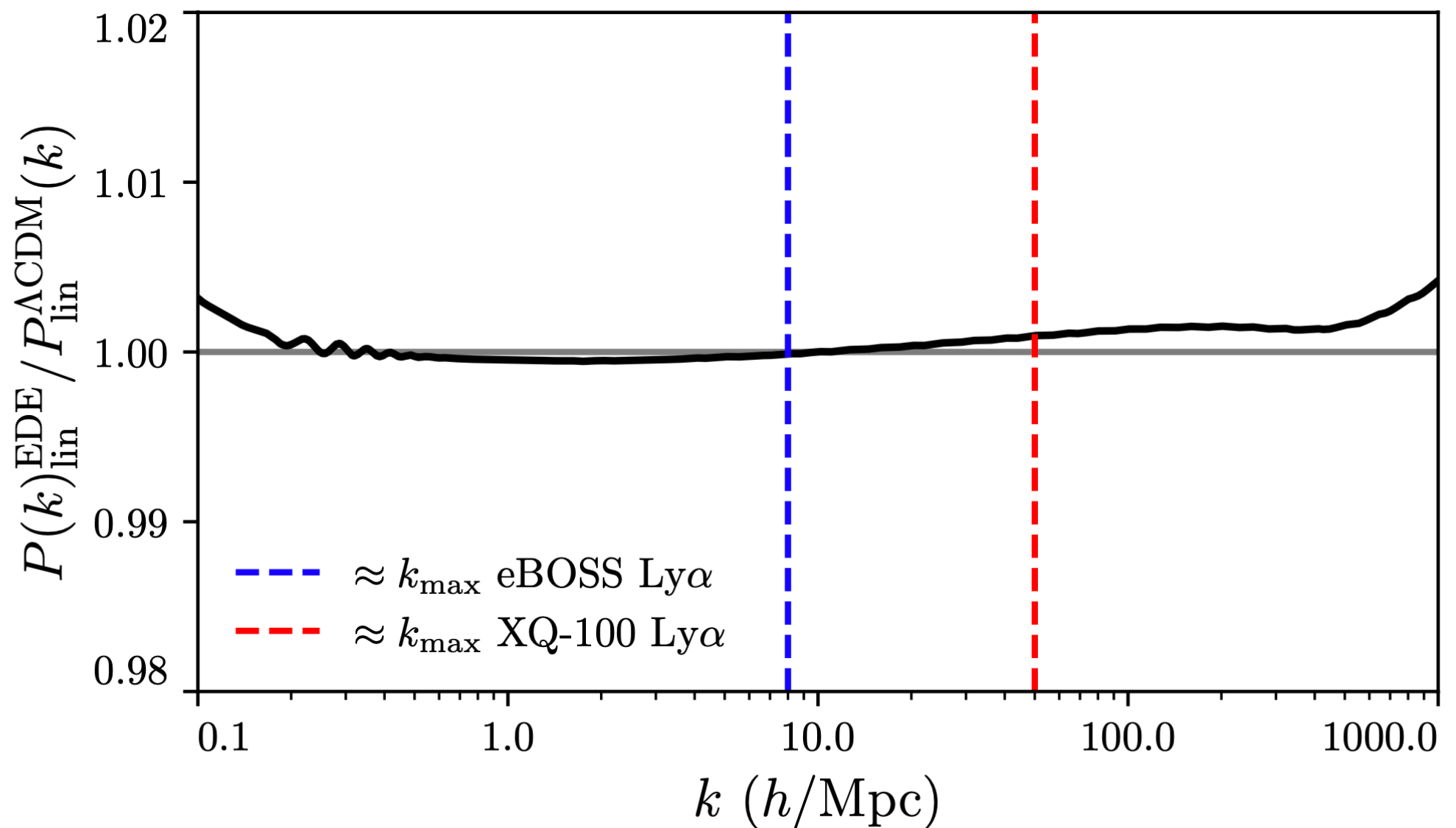
# Ly $\alpha$ EDE Validation



# Ly $\alpha$ EDE Validation

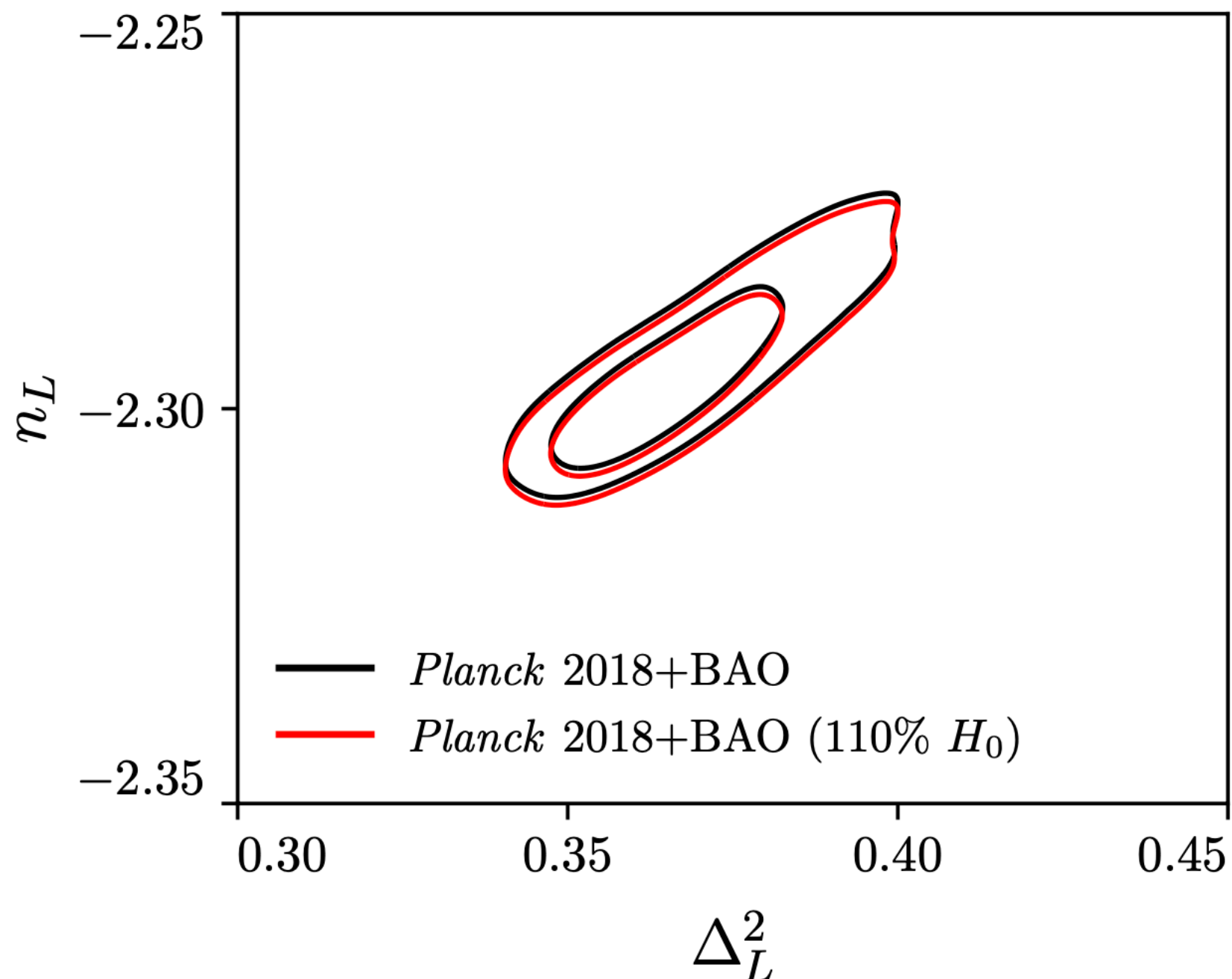
Do the hydro sim grids used in the Ly $\alpha$  likelihood construction cover relevant  $P_{\text{lin}}(k)$  for EDE analysis? Yes

Best-fit baseline EDE  $P(k)$  at  $z=3$  can be very accurately mimicked by  $\Lambda$ CDM  $P(k)$  with slightly tweaked parameters



# Ly $\alpha$ EDE Validation

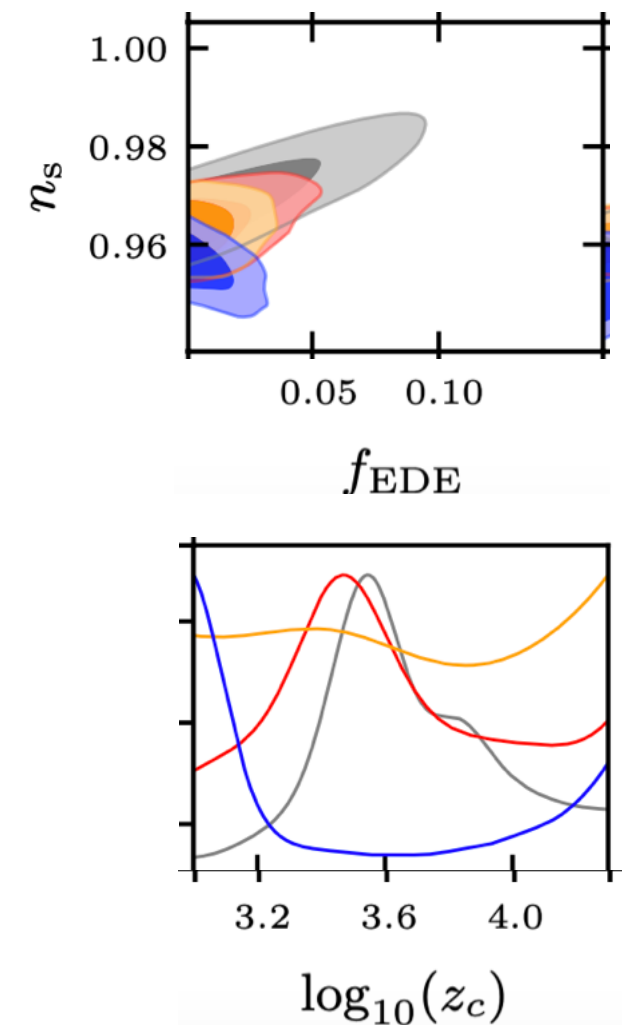
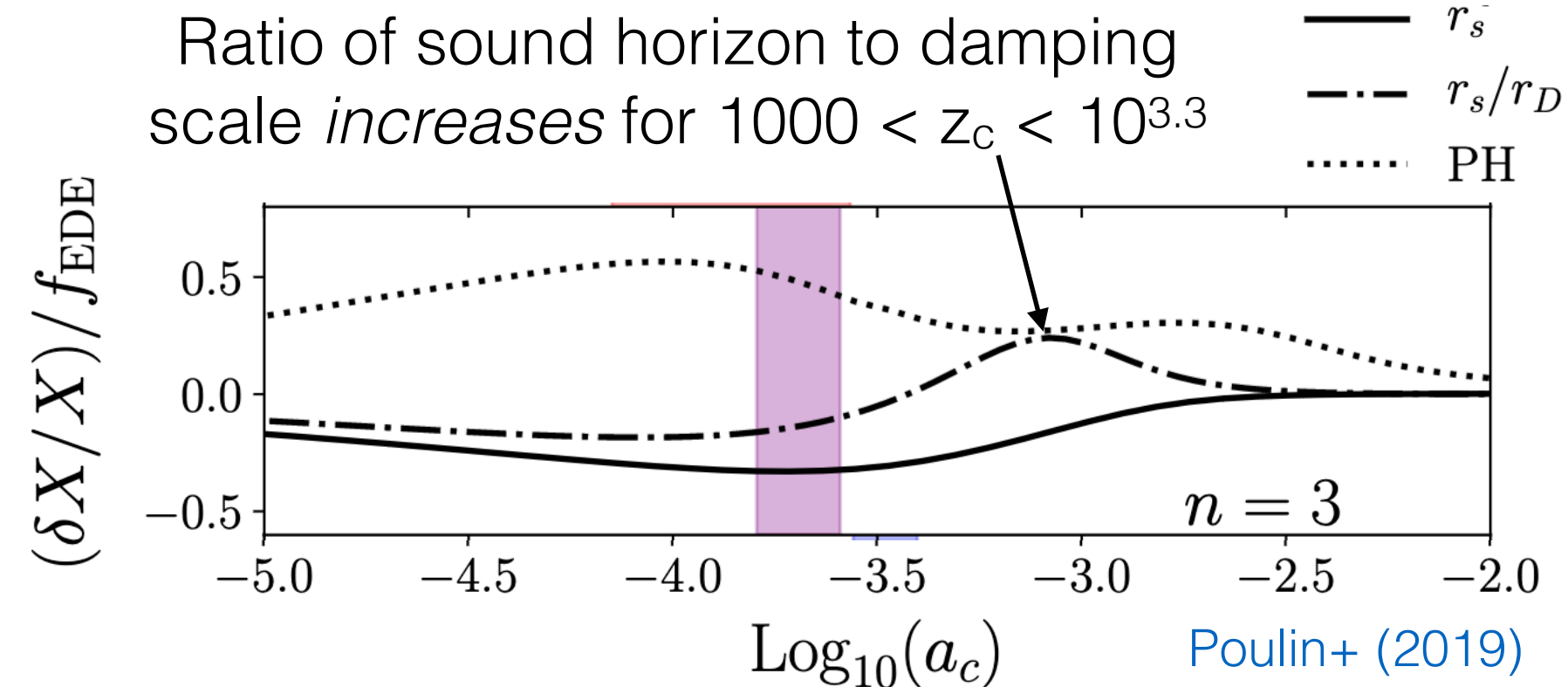
Do the priors used in the Ly $\alpha$  likelihood construction have any impact on the compressed parameter likelihoods used in EDE analysis? No





# Ly $\alpha$ EDE Validation

Origin of the  $n_s$  -  $f_{\text{EDE}}$  anti-correlation for the [baseline+eBOSS analysis](#)



Thus  $\theta_s / \theta_d$  increases; but  $\theta_s$  is fixed by observations, so  $\theta_d$  *decreases*, i.e.,  $\ell_d$  increases. Hence less damping at a given  $\ell$ , so  $n_s$  decreases to compensate.

# Primordial Magnetic Fields and Baryon Clumping

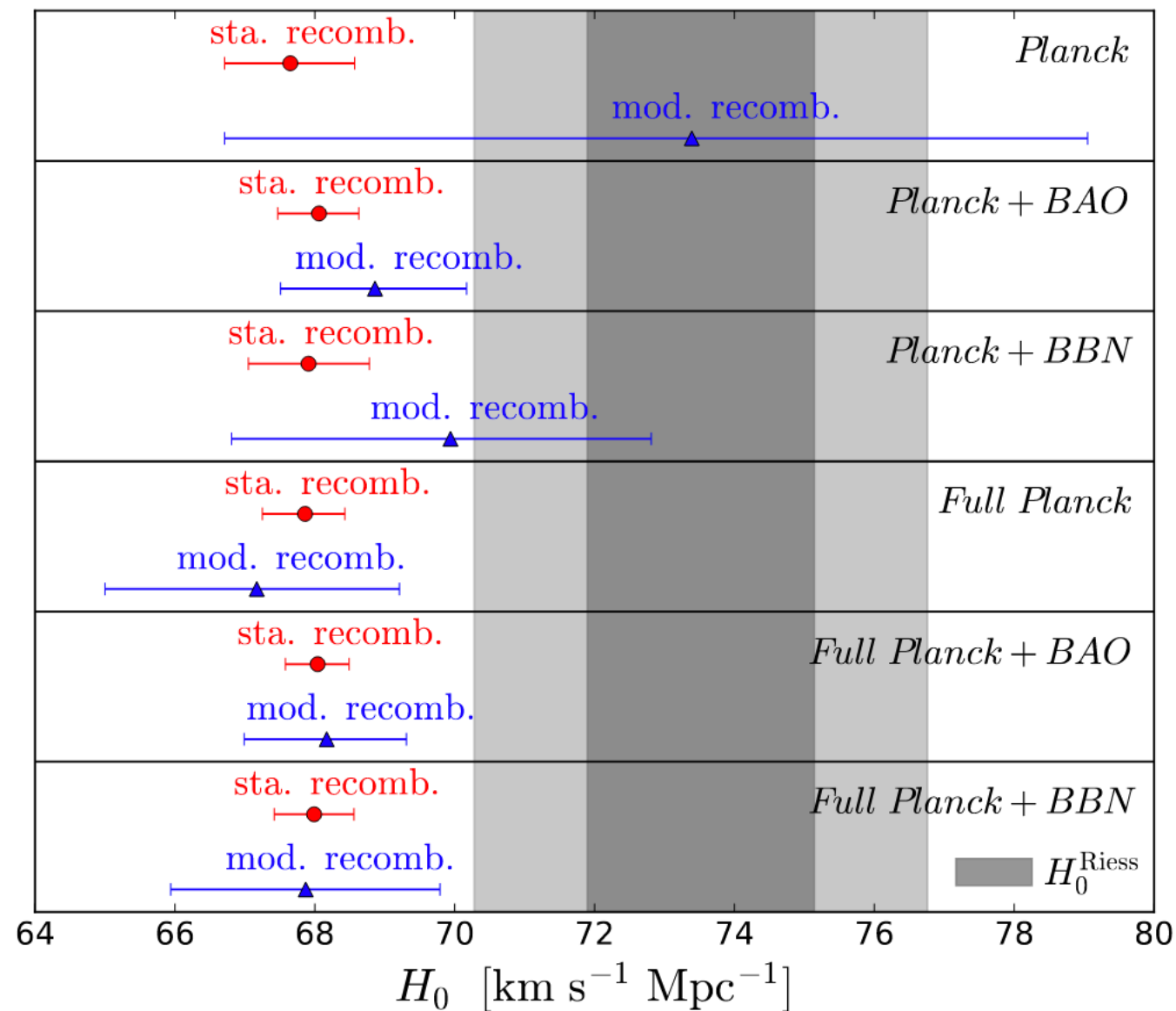


**2105.03003 w/ L. Thiele, Y. Guan, A. Kosowsky, D. Spergel**

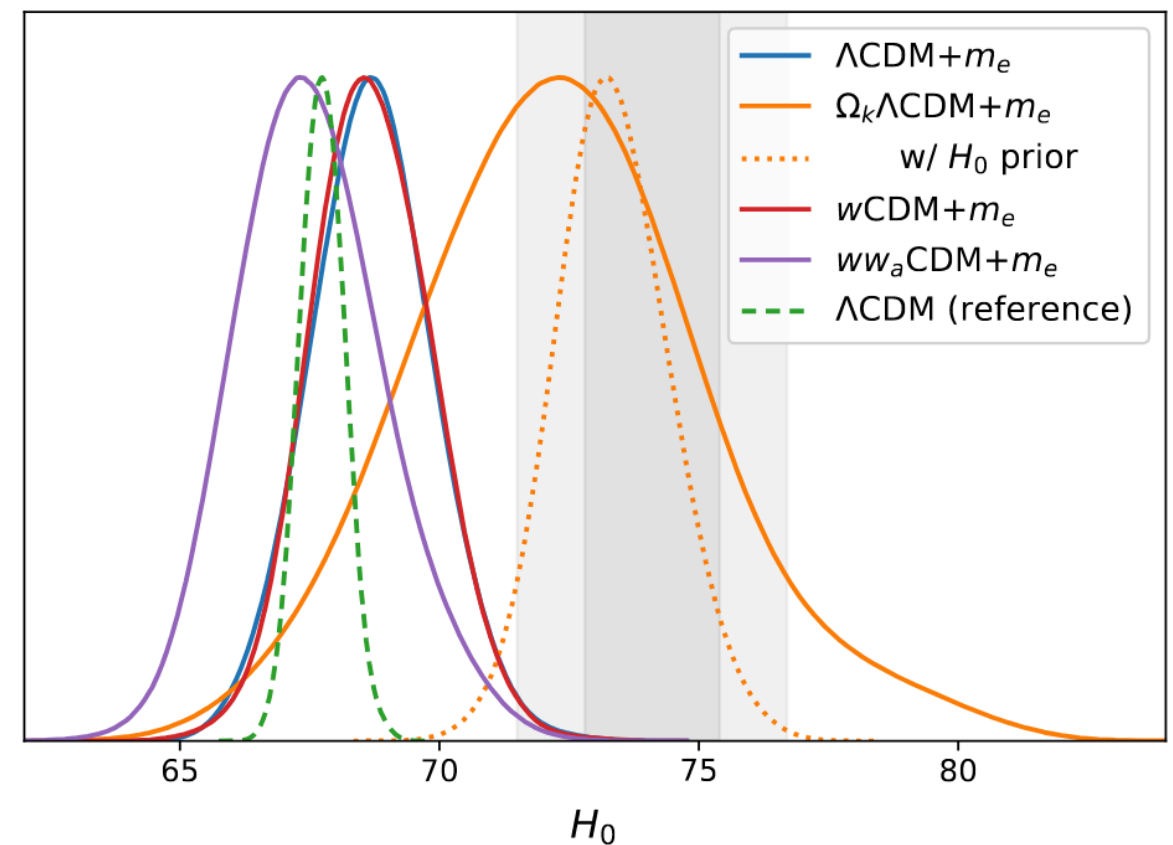
# Early Recombination

Success depends on the details of the physical implementation

## Phenomenological model



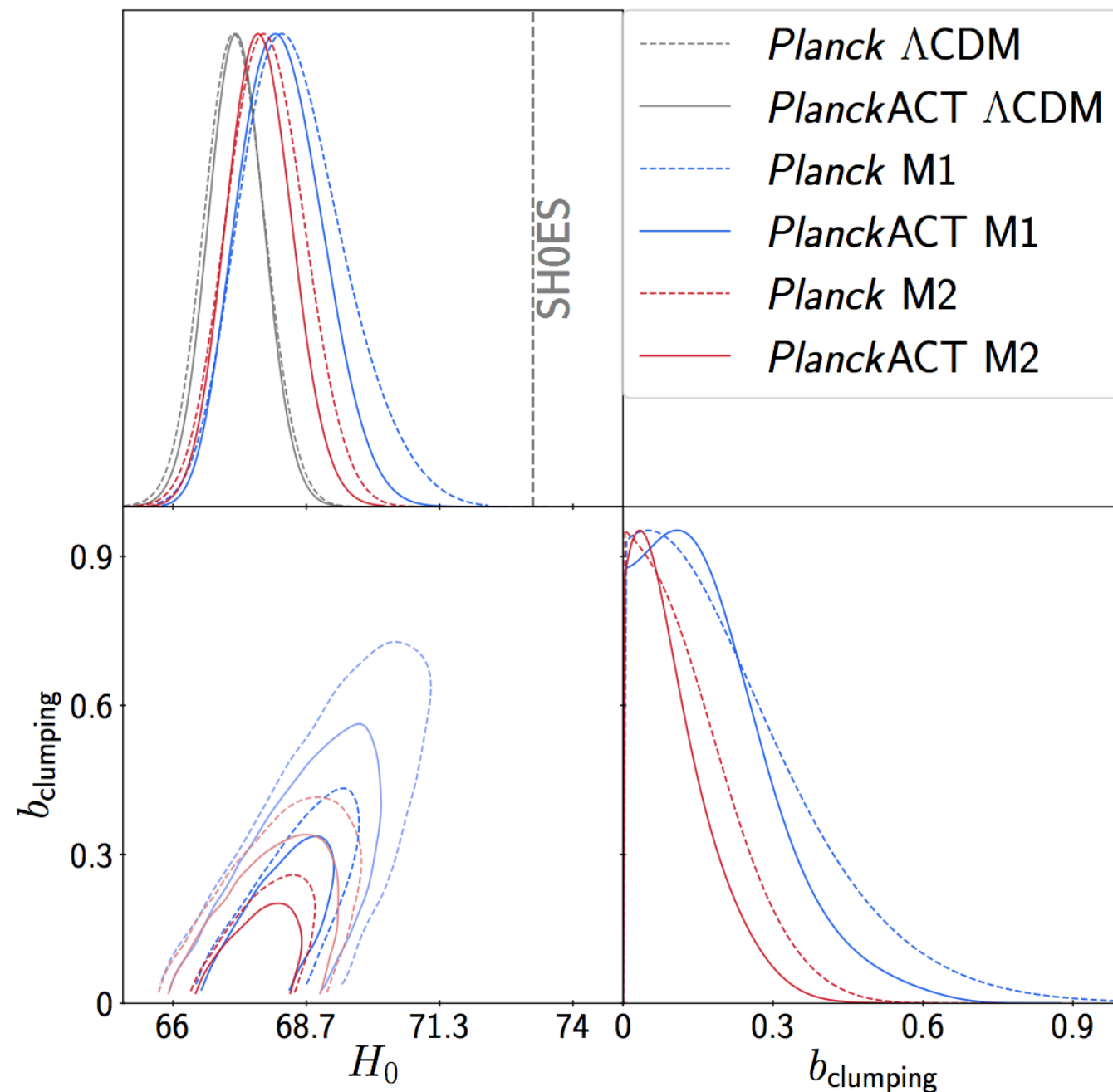
## “Physical” model (varying $m_e$ )



# Early Recombination

Success depends on the details of the physical implementation

**“Physical” model (primordial magnetic fields) — effectively excluded by ACT DR4**



$\Lambda\text{CDM}$	$67.26 \pm 0.60$
M1	$68.18 \pm 0.87$
M2	$67.74 \pm 0.71$

# Baryon Clumping

Colin Hill  
Columbia/CCA

- Basic idea: push recombination to earlier time, thus yielding a decreased sound horizon (and higher  $H_0$  from CMB data)
- Recombination rate:  $\dot{n}_e \propto -n_e^2$
- Thus, small-scale clumping s.t.  $\langle n_e^2 \rangle > \langle n_e \rangle^2$  yields faster recombination and smaller  $r_s$

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Colin Hill  
Columbia/CCA

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- Recombination rate:  $\dot{n}_e \propto -n_e^2$
- Thus, small-scale clumping s.t.  $\langle n_e^2 \rangle > \langle n_e \rangle^2$  yields faster recombination and smaller  $r_s$
- This is naturally achieved with primordial magnetic fields, which can dynamically sustain large inhomogeneities on small scales (otherwise suppressed by diffusion/Silk damping) — but we can agnostically consider “baryon clumping” on small scales in general

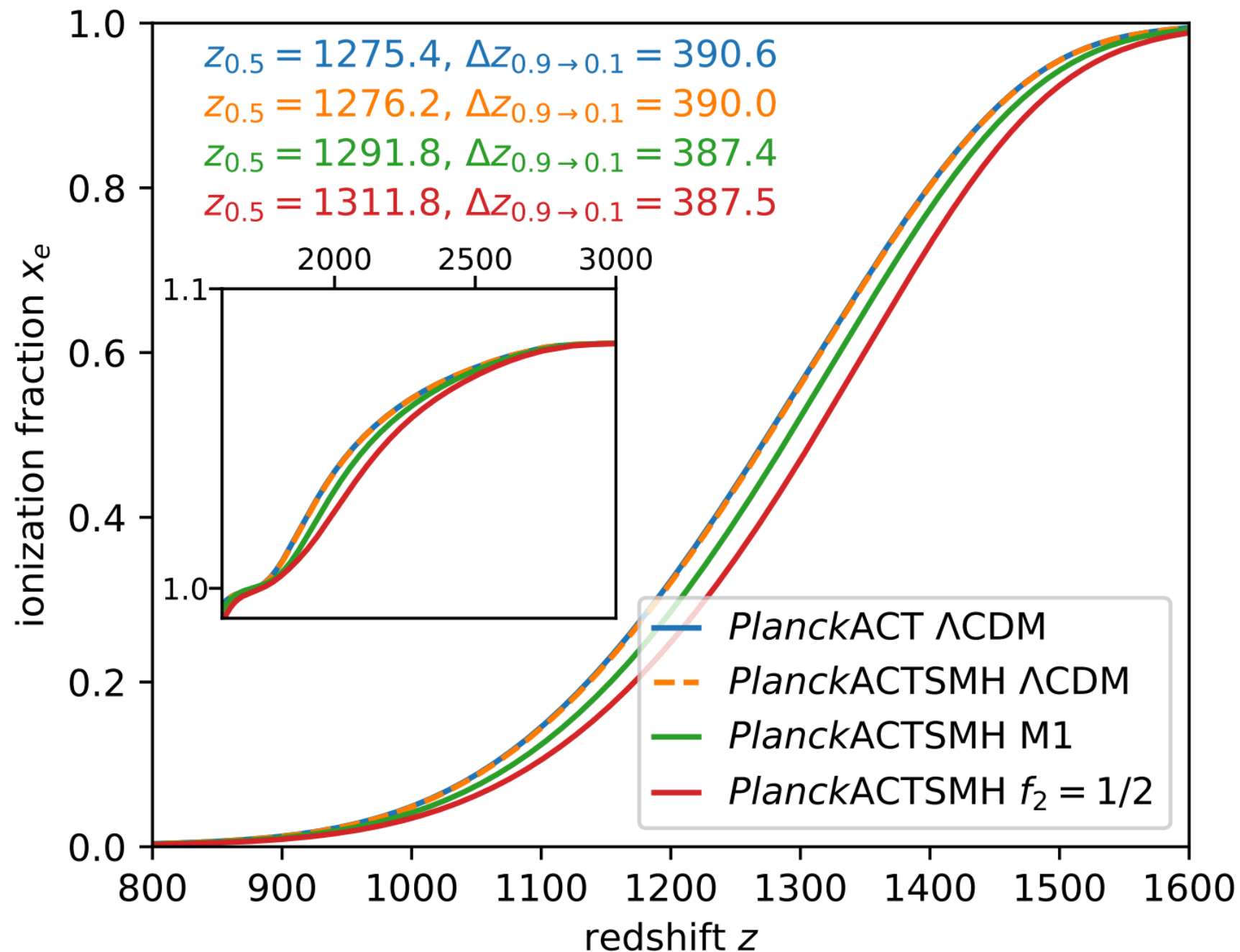


# Baryon Clumping

Colin Hill  
Columbia/CCA

Illustration of the mechanism

Ionization history of the universe





# Baryon Clumping

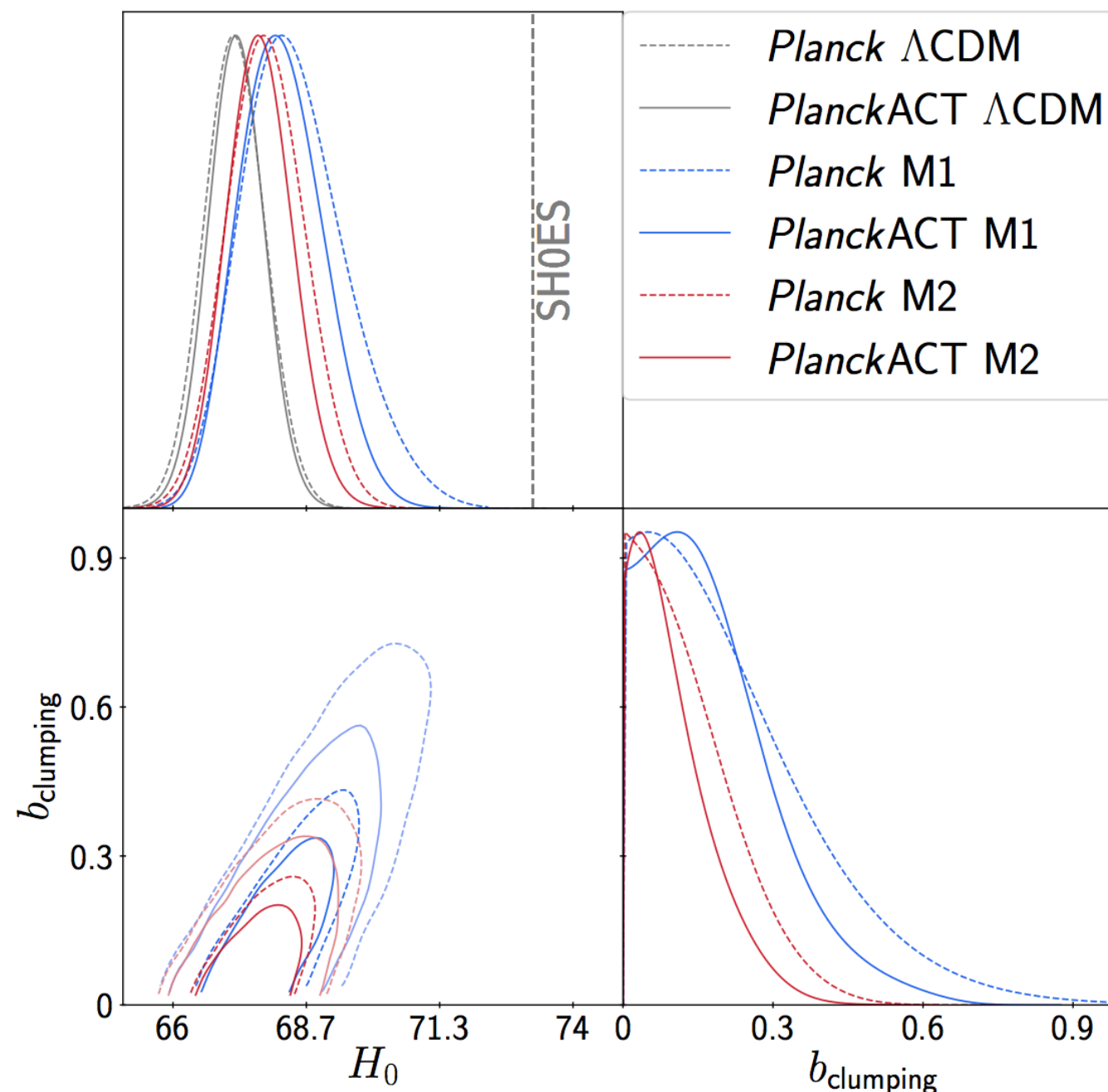
Colin Hill  
Columbia/CCA

ACT: independent test of this scenario with high- $\ell$  CMB data,  
probing scales that Planck does not (and higher S/N in TE + EE)

# Baryon Clumping

Colin Hill  
Columbia/CCA

ACT: independent test of this scenario with high- $\ell$  CMB data, probing scales that Planck does not (and higher S/N in TE + EE)



*Planck*+ACT

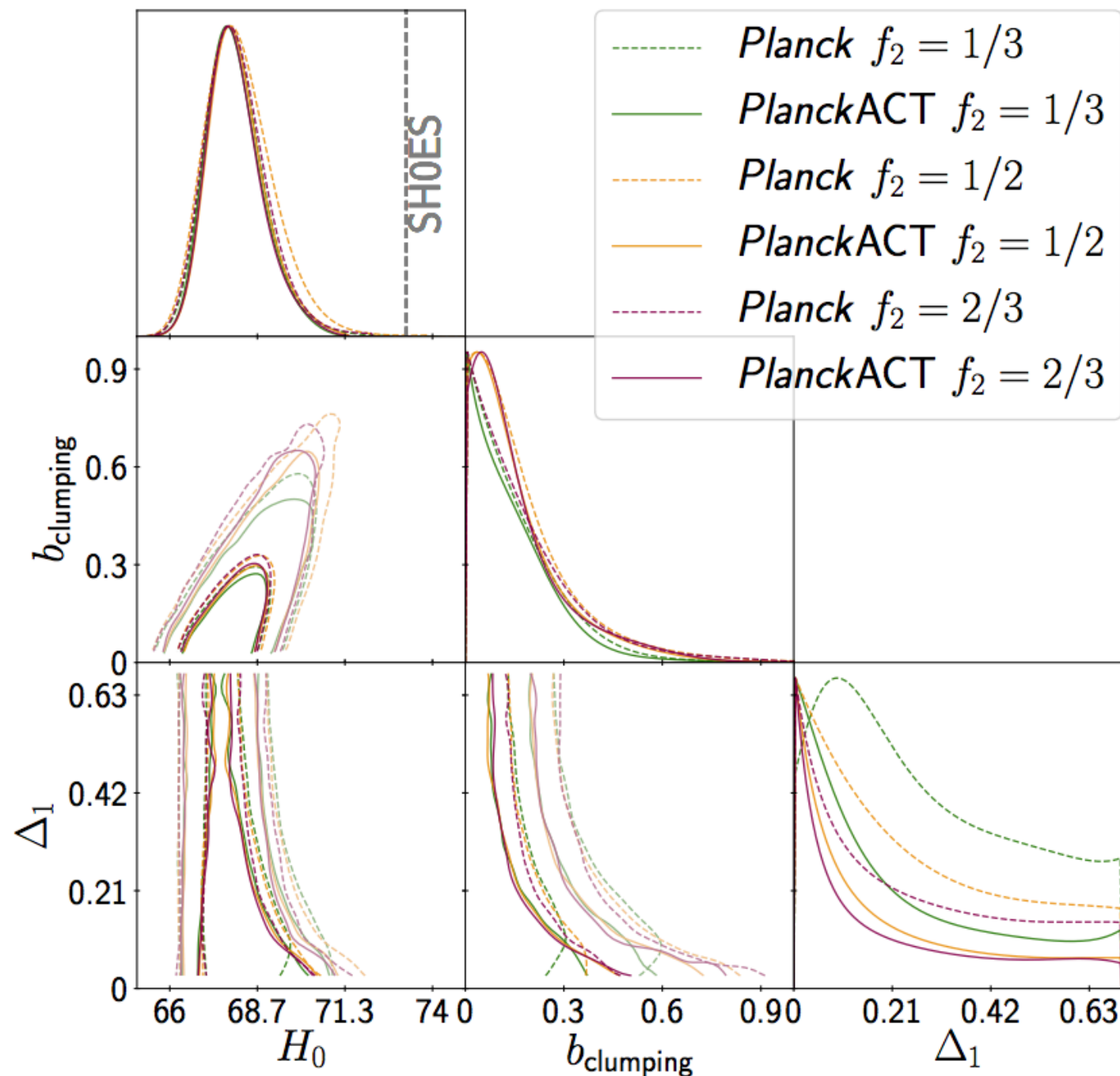
model	$H_0$ [km/s/Mpc]	$b$
$\Lambda$ CDM	$67.26 \pm 0.60$	—
M1	$68.18 \pm 0.87$	$<0.42$
M2	$67.74 \pm 0.71$	$<0.26$

no relief of tension  
with SH0ES — one  
should thus be  
wary of combining

# Baryon Clumping

Colin Hill  
Columbia/CCA

Results persist even in broadened parameter space (more freedom in the small-scale baryon PDF)



# Baryon Clumping

Colin Hill  
Columbia/CCA

## Takeaways and Outlook

- Baryon clumping models not preferred by ACT + Planck (or ACT + Planck + BAO)
- CMB  $H_0$  still in tension with SH0ES even in the context of this model, so one should be wary of combining likelihoods — if one tries to do so, the CMB  $\chi^2$  worsens by  $\sim 6-7$  relative to  $\Lambda$ CDM — thus the model does not appear to restore concordance
- Perhaps further model extensions could help — e.g., time-dependence in the small-scale baryon PDF (although see Lee & Ali-Haïmoud 2021)

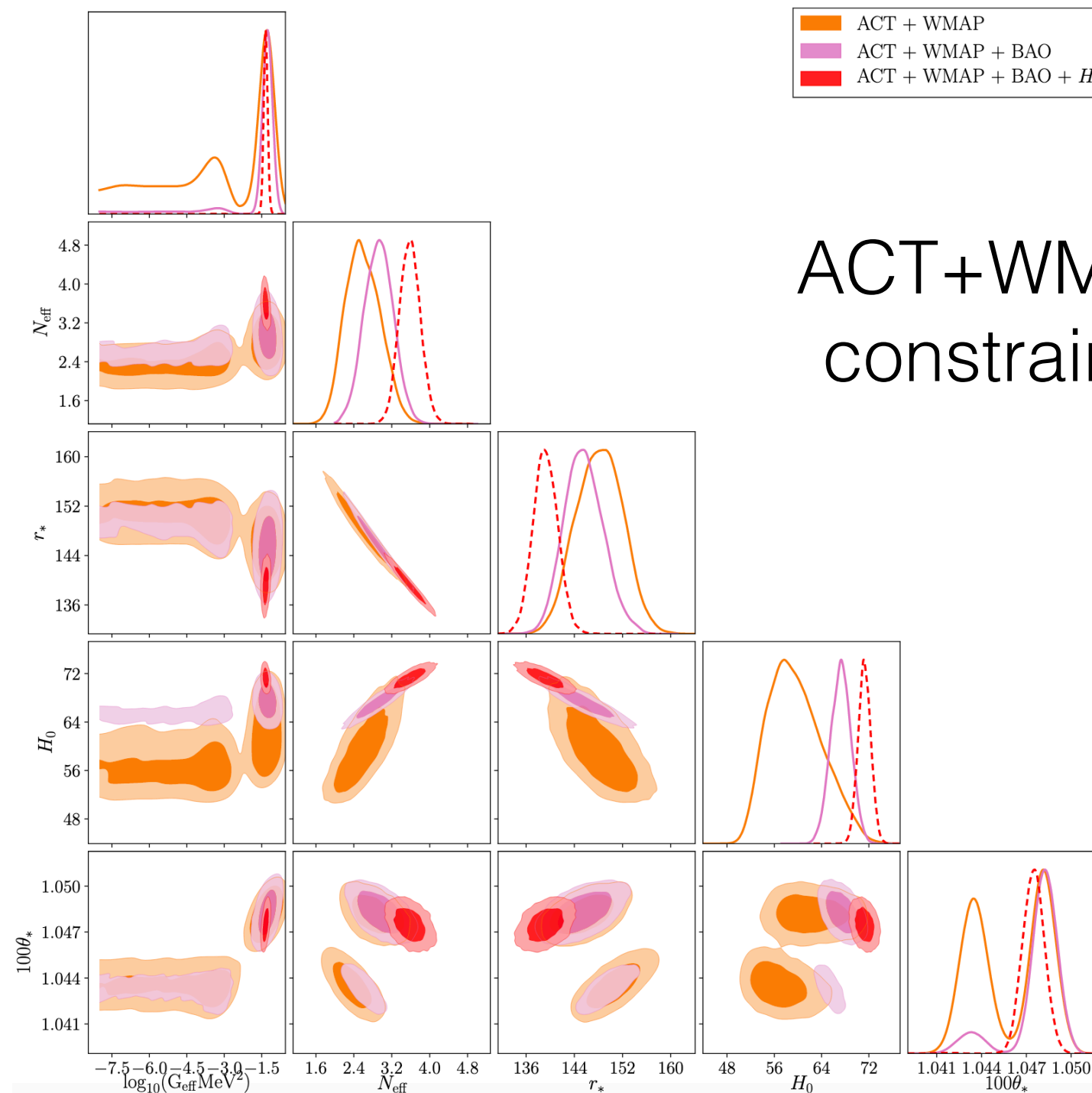
# SI Neutrinos

Colin Hill  
Columbia/CCA

Idea: delay onset of neutrino free-streaming until  $z \sim z^*$

$$\dot{\tau}_\nu = -aG_{\text{eff}}^2 T_\nu^5$$

$$g_\nu(\tau) \equiv -\dot{\tau}_\nu e^{-\tau_\nu}$$



ACT+WMAP  
constraints

Kreisch+ (2022); see  
also earlier work from  
Kreisch, Cyr-Racine,  
Dore, Sigurdson, ++

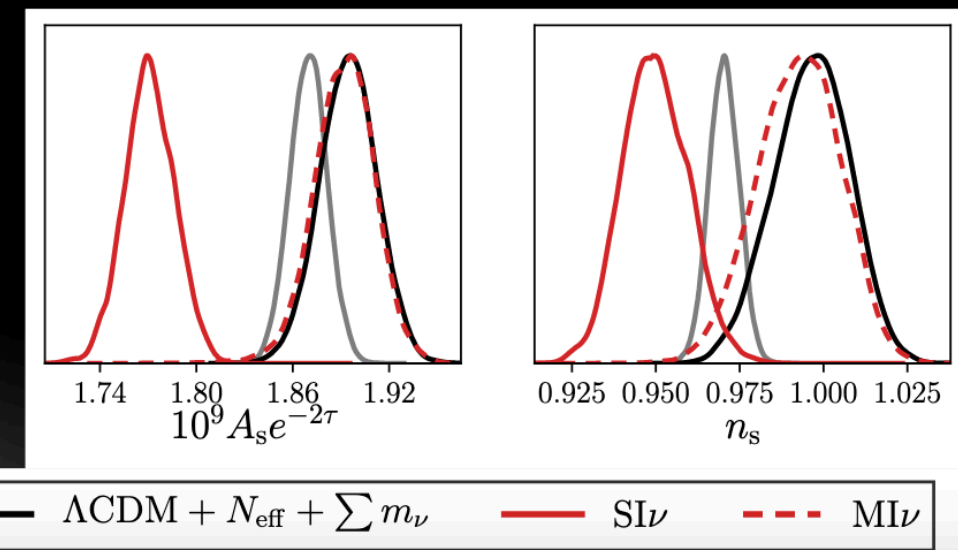
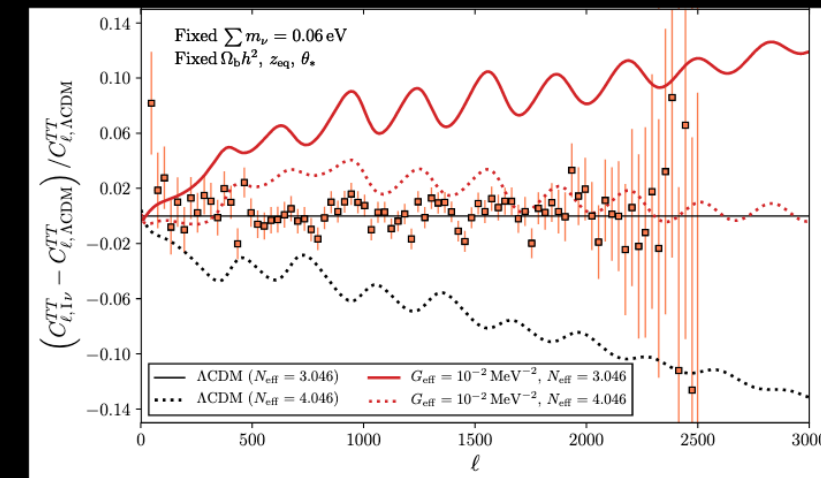
# SI Neutrinos

Colin Hill  
Columbia/CCA

Why does this model work?

- $N_{\text{eff}}$  increases Hubble at early times, hence reducing the sound horizon.
- The tightly-coupled neutrinos do not over damp or phase shift the photon-baryon fluctuations.
- Changes in the primordial spectrum of fluctuations ( $n_s$ ,  $A_s$ ) absorbs the remainder of the changes.

$$r_s = \int_0^{a_d} da \frac{c_s(a)}{a^2 H(a)}$$



—  $\Lambda\text{CDM}$     —  $\Lambda\text{CDM} + N_{\text{eff}} + \sum m_\nu$     —  $\text{SI}\nu$     - - -  $\text{MI}\nu$

# Generalized Dark Matter $\longrightarrow$ Dark Radiation Conversion

See

<https://arxiv.org/abs/2210.14339>





# DM $\longrightarrow$ DR

Analytic ansatz for evolution of ‘decaying’ dark matter (DDM) density:

$$\rho_\chi(a) = \frac{\rho_\chi^0}{a^3} \left[ 1 + \zeta \frac{1 - a^\kappa}{1 + (a/a_t)^\kappa} \right]$$

- Assumption: DDM  $\longrightarrow$  massless dark radiation

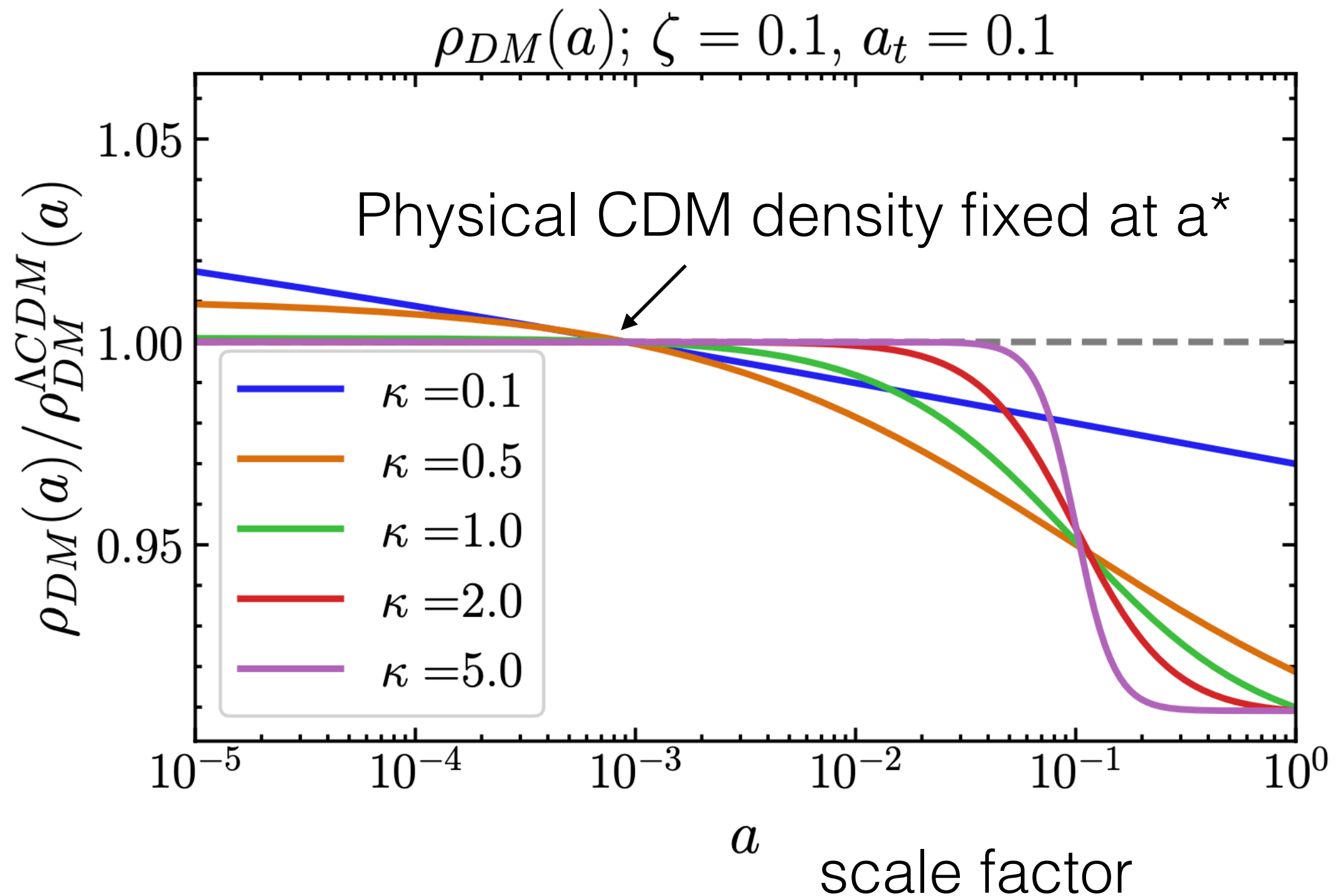
$$\begin{aligned} \rho_\phi(a) = & \zeta \frac{\rho_\chi^0}{a^3} \frac{(1 + a_t^\kappa)}{(a^\kappa + a_t^\kappa)} \\ & \times \left( (a^\kappa + a_t^\kappa) {}_2F_1 \left[ 1, \frac{1}{\kappa}; 1 + \frac{1}{\kappa}; - \left( \frac{a}{a_t} \right)^\kappa \right] - a_t^\kappa \right) \end{aligned}$$

- Comoving DM density decreases by a factor of  $(1+\zeta)$
- Transition is centered at  $a = a_t$ 
  - e.g., for  $a_t \ll 10^{-3}$ , effects of this model are similar to  $N_{\text{eff}}$
- $\kappa$  determines the rate of the transition
  - e.g.,  $\kappa=2 \longrightarrow$  standard exponential decay with  $\Gamma \sim H(a_t)$
  - e.g.,  $\kappa=1 \longrightarrow$  Sommerfeld-enhanced annihilation

# DM $\longrightarrow$ DR

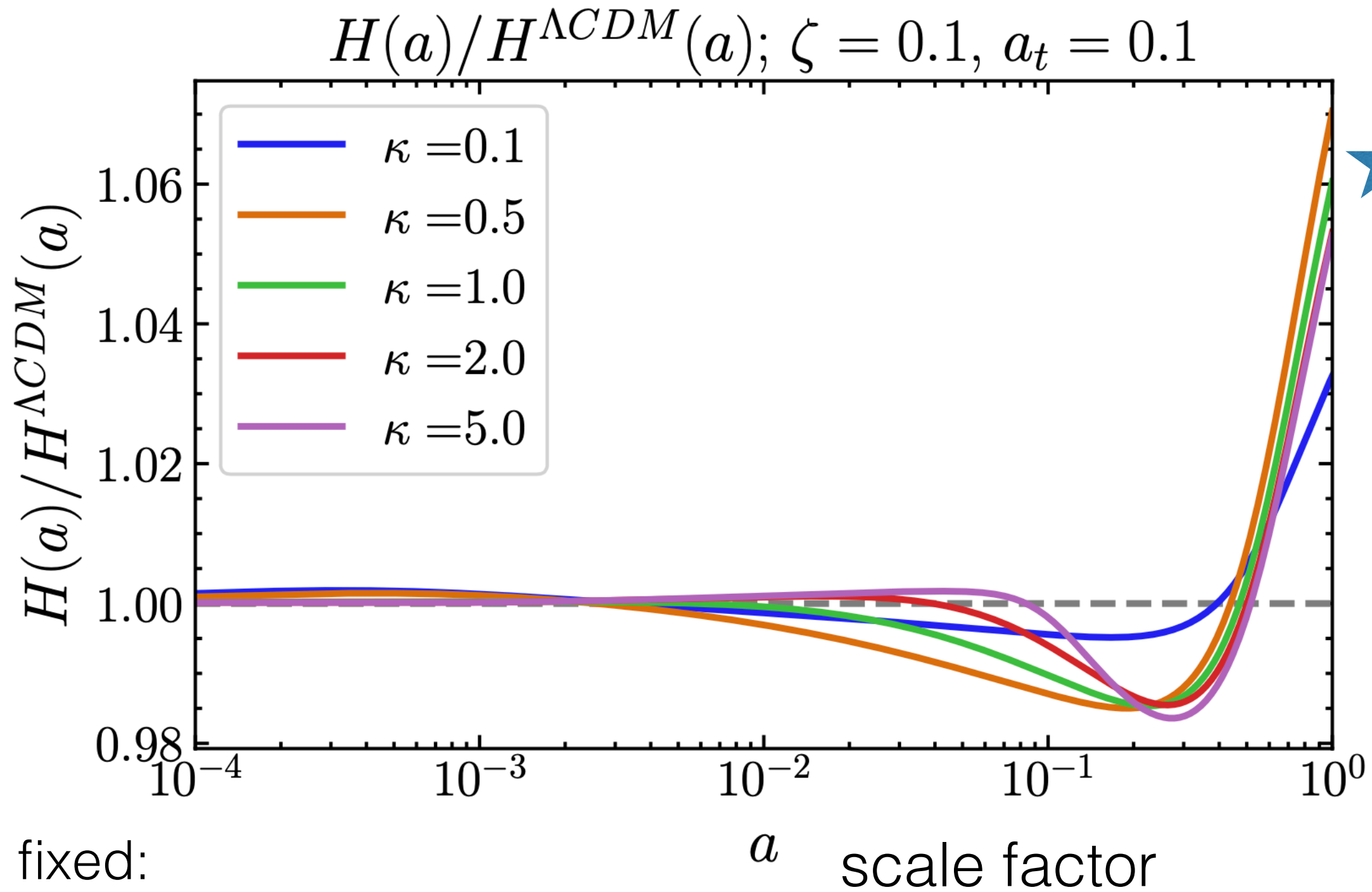
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Analytic ansatz for DDM density



# DM $\longrightarrow$ DR

Evolution of Hubble parameter



Held fixed:

$\omega_b, \theta^*, A_s, n_s, \omega_c(z^*)$

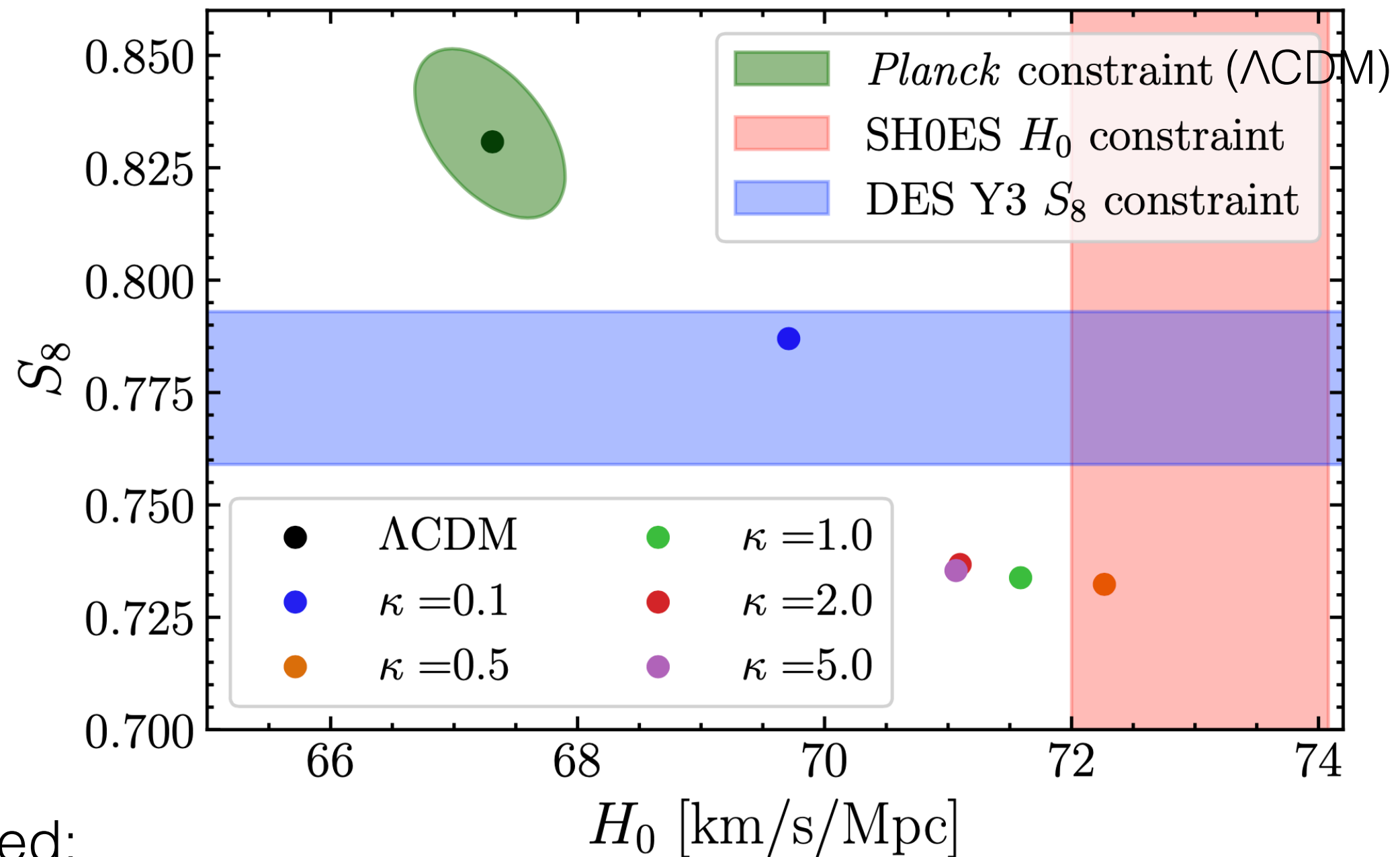
Bringmann+ (2018); McCarthy & JCH (2022)

how?  $z_\Lambda$  increases

# DM $\longrightarrow$ DR

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Motivation: such scenarios (can) naturally move both  $H_0$  and  $S_8$  in the “right direction”



Held fixed:  
 $\omega_b, \theta^*, A_s, n_s, \omega_c(z^*)$

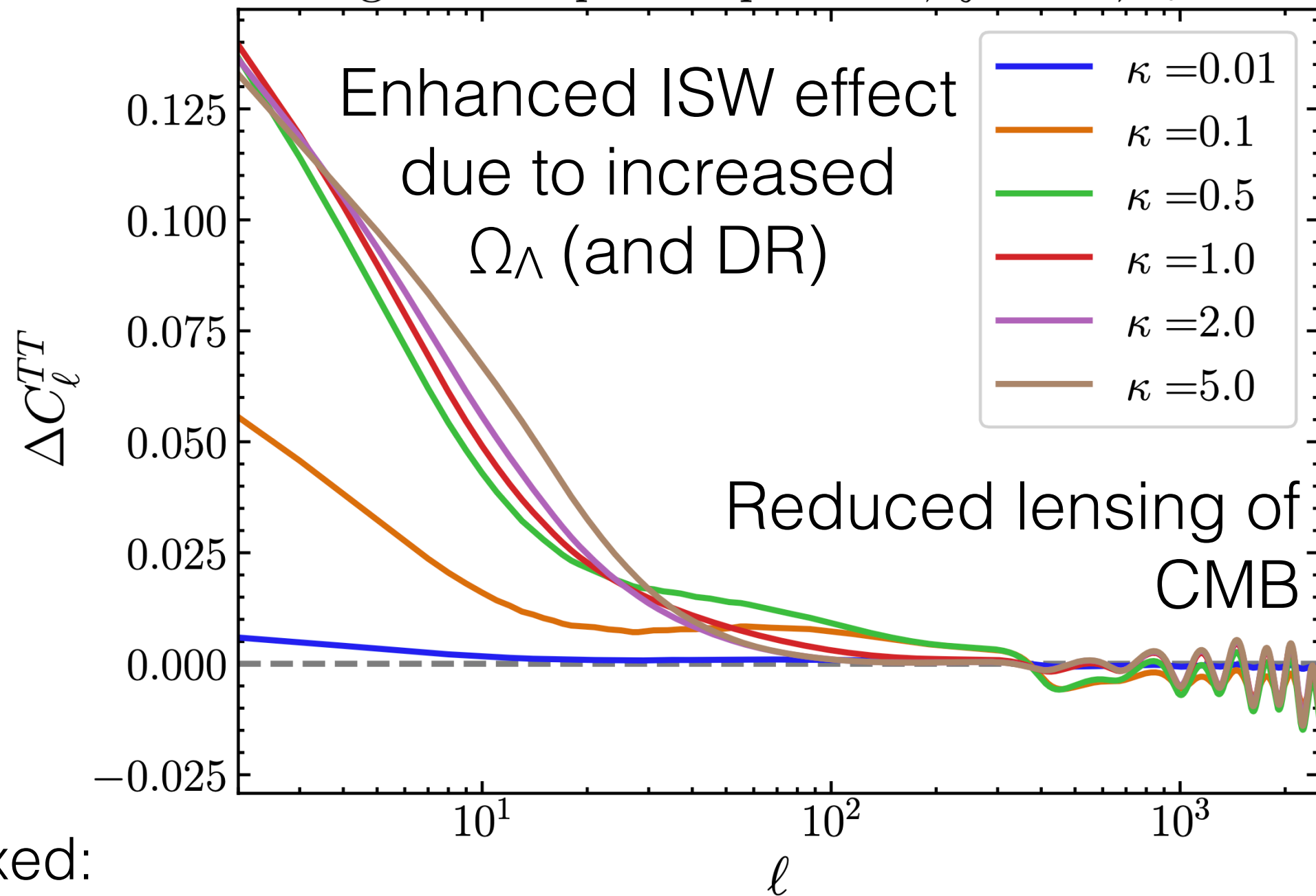
Bringmann+ (2018); McCarthy & JCH (2022); see also Chen et al. (2021) [DES Collaboration]

# DM $\longrightarrow$ DR

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Observables:  $TT$  power spectrum

Change in  $TT$  power spectrum;  $\zeta = 0.1, a_t = 0.1$



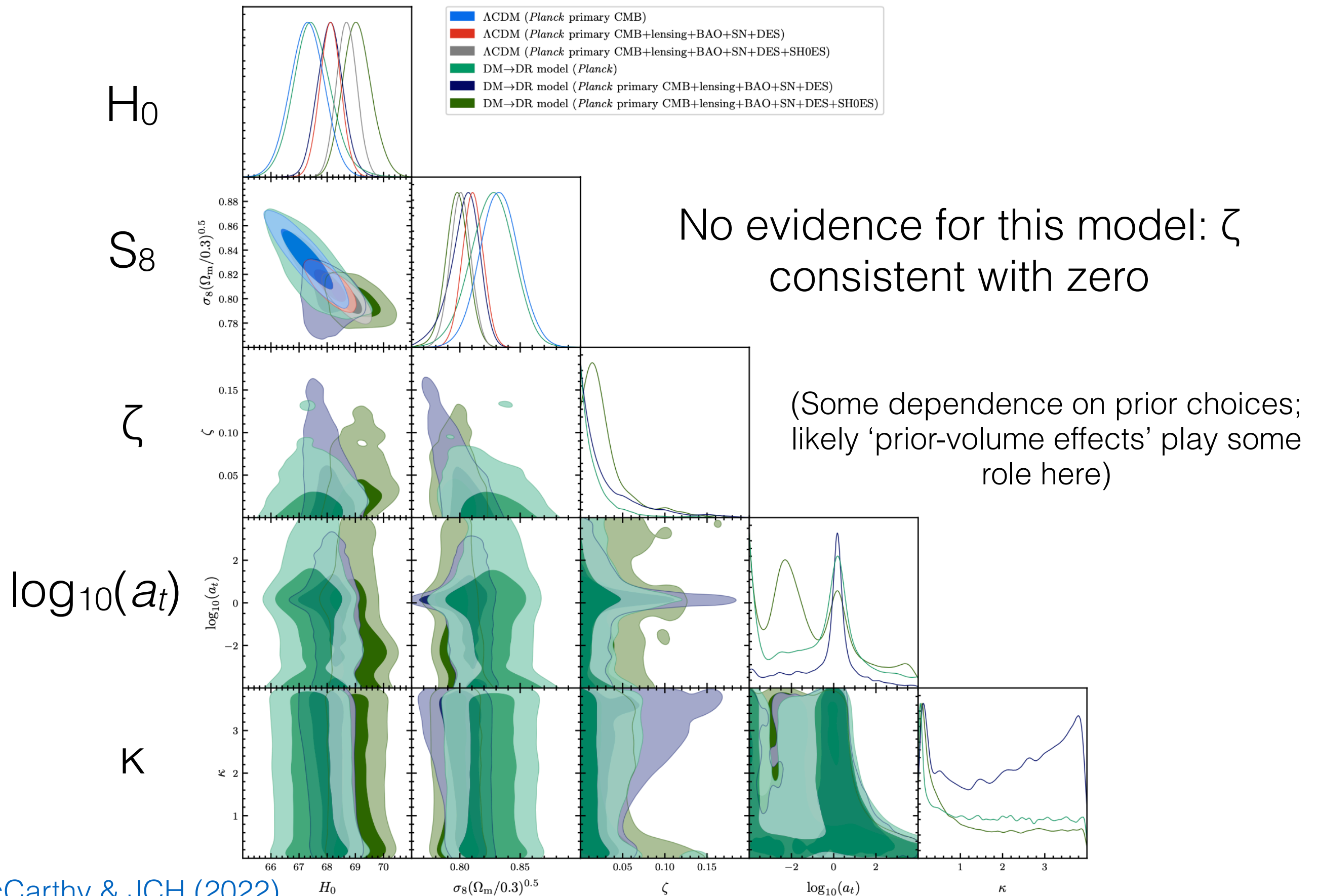
Held fixed:

$\omega_b, \theta^*, A_s, n_s, \omega_c(z^*)$

Bringmann+ (2018); McCarthy & JCH (2022)

# DM $\rightarrow$ DR: Results

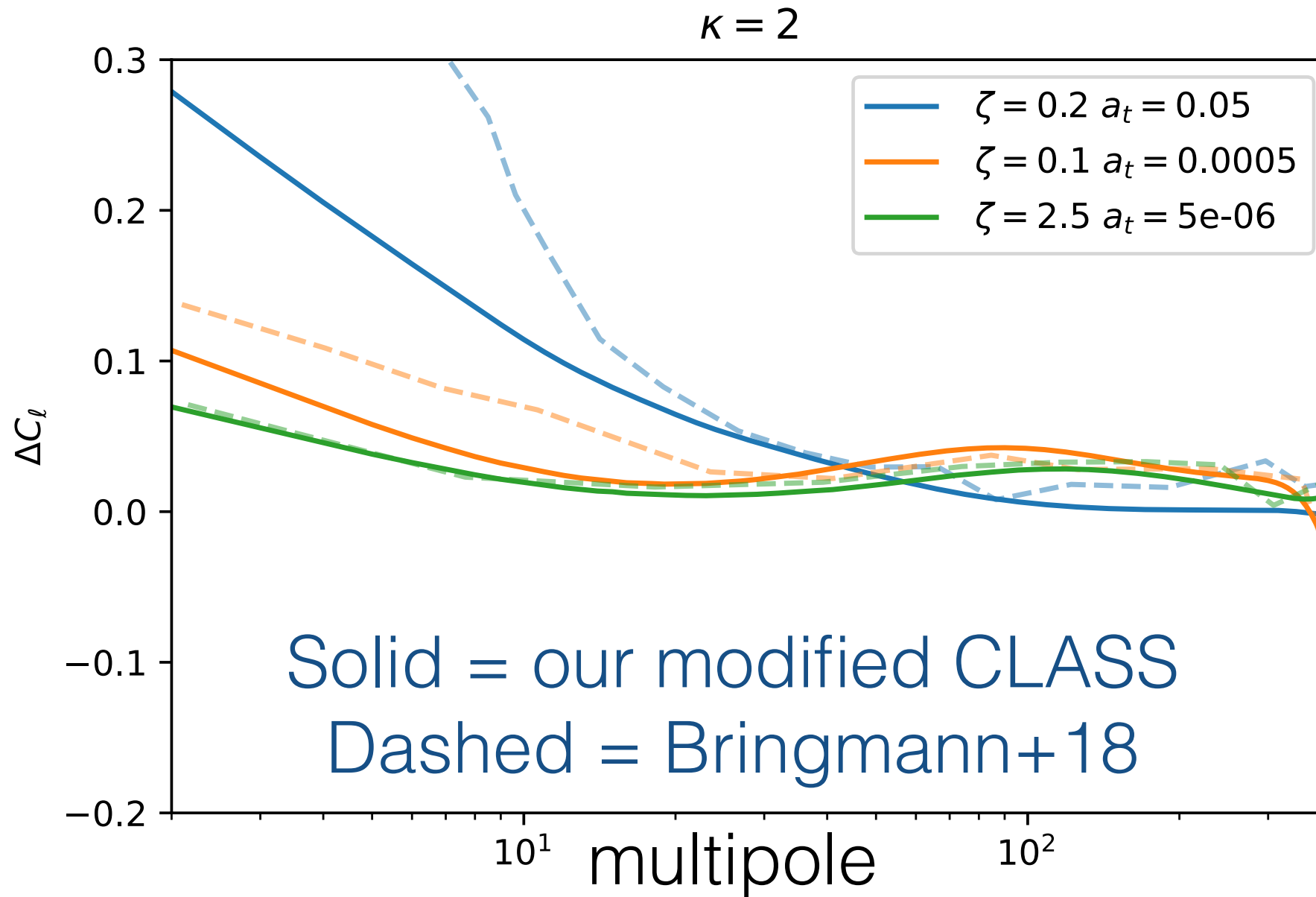
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# DM $\rightarrow$ DR

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Earlier work (Bringmann+18) found  $\sim 2\sigma$  hint for this model in Planck15: what changed?  
Diffs. in theory calc. (Boltzmann hierarchy for the DR decay products) + datasets/priors



A couple of bugs also found in code used in Chen+21 via cross-check;  
we have now iterated with A. Chen and D. Huterer and have two independent  
codes (modified CLASS and modified CAMB) that agree very well

Bringmann+ (2018); McCarthy & JCH (2022); see also Chen et al. (2021) [DES Collaboration]



# DM $\longrightarrow$ DR

## Takeaways

- Fit to Planck TTTEEE $\kappa\kappa$  + BAO + Pantheon + SH0ES + DES-Y3:  
 $H_0 = 69.1 \pm 0.5$  km/s/Mpc  
 $S_8 = 0.799 \pm 0.009$   
*however*  
Planck CMB  $\chi^2$  is significantly worse than a fit of  $\Lambda$ CDM to Planck ( $\Delta\chi^2 \sim +15$ )
- Why? Excess ISW signal in low- $\ell$  TT and reduced lensing
- Even this very general DM $\longrightarrow$ DR model lacks sufficient flexibility to accommodate late-time  $H_0$  and  $S_8$  while maintaining good fit to Planck

# Did the Universe (Slightly) Reheat after Recombination?

## Linear perturbation theory

$$\delta'_{\text{DCDM}} = -\theta_{\text{DCDM}} - m_{\text{cont}} - a\Gamma m_{\psi} \quad (3)$$

$$\theta'_{\text{DCDM}} = -\frac{a'}{a}\theta_{\text{DCDM}} + k^2 m_{\psi} \quad (4)$$

$$\delta'_{\gamma} = -\frac{4}{3}\theta_{\gamma} - \frac{4}{3}m_{\text{cont}} + a\Gamma \frac{\rho_{\text{DCDM}}}{\rho_{\gamma}} (\delta_{\text{DCDM}} - \delta_{\gamma} + m_{\psi}) \quad (5)$$

$$\theta'_{\gamma} = k^2 \left( \frac{1}{4}\delta_{\gamma} - \sigma_{\gamma} \right) + k^2 m_{\psi} + an_e \sigma_{\text{T}} (\theta_b - \theta_{\gamma}) - \frac{3}{4}a\Gamma \frac{\rho_{\text{DCDM}}}{\rho_{\gamma}} \left( \frac{4}{3}\theta_{\gamma} - \theta_{\text{DCDM}} \right) \quad (6)$$

$$F'_{\gamma,2} = 2\sigma'_{\gamma} = \frac{8}{15}\theta_{\gamma} - \frac{3k}{5}F_{\gamma,3} + \frac{8}{15}m_{\text{shear}} - \frac{9}{5}an_e \sigma_{\text{T}} \sigma_{\gamma} + \frac{1}{10}an_e \sigma_{\text{T}} (G_{\gamma,0} + G_{\gamma,2}) - 2\sigma_{\gamma}a\Gamma \frac{\rho_{\text{DCDM}}}{\rho_{\gamma}} \quad (7)$$

$$F'_{\gamma,\ell} = \frac{k}{2\ell+1} [\ell F_{\gamma,\ell-1} - (\ell+1) F_{\gamma,\ell+1}] - a\Gamma F_{\gamma,\ell} \frac{\rho_{\text{DCDM}}}{\rho_{\gamma}} \quad (8)$$

All terms  
containing  $\Gamma$  in  
photon  
perturbation  
equations are  
new

Gauge	Synchronous	Newtonian
$m_{\text{cont}}$	$h'/2$	$-3\phi'$
$m_{\psi}$	0	$\psi$
$m_{\text{shear}}$	$(h' + 6\eta')/2$	0

# Did the Universe (Slightly) Reheat after Recombination?

“Hat” variables: all cosmological quantities at a given  $T_{\text{CMB}}$  depend on quantities proportional to baryon-to-photon and dark matter-to-photon number ratios

$$\begin{aligned}\hat{\omega}_{\text{b}} &\equiv \omega_{\text{b}} \left( \frac{T_{\text{CMB,ini}}}{T_{\text{FIRAS}}} \right)^{-3} \\ \hat{\omega}_{\text{c}} &\equiv \omega_{\text{c}} \left( \frac{T_{\text{CMB,ini}}}{T_{\text{FIRAS}}} \right)^{-3} \\ \hat{\omega}_{\text{DCDM,ini}} &\equiv \omega_{\text{DCDM,ini}} \left( \frac{T_{\text{CMB,ini}}}{T_{\text{FIRAS}}} \right)^{-3} \\ \hat{A}_{\text{s}} &\equiv A_{\text{s}} \left( \frac{T_{\text{CMB,ini}}}{T_{\text{FIRAS}}} \right)^{n_{\text{s}}-1}.\end{aligned}$$

BBN abundances depend only on  $\hat{\omega}_{\text{b}}$

Consistency of this parameter with Planck thus maintains (approximate) consistency with BBN [modulo varying  $N_{\text{eff}}$ , etc.]

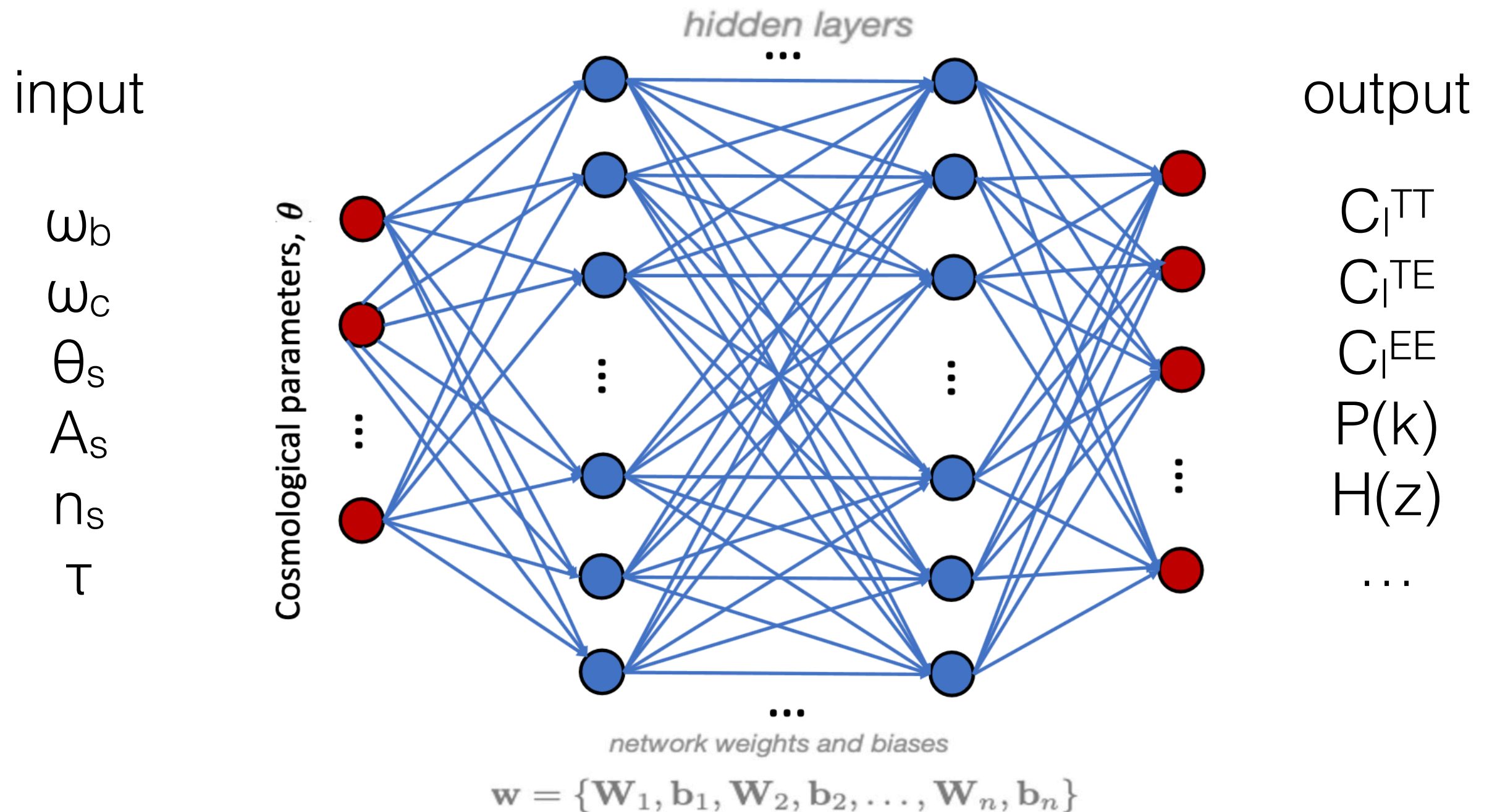
# Einstein-Boltzmann Emulators



Bolliet, Spurio Mancini, JCH, Madhavacheril, et al. (2023)

# CosmoPower

Cosmological observables are smooth functions of the input parameters:  
easy to emulate at high accuracy with modern neural networks



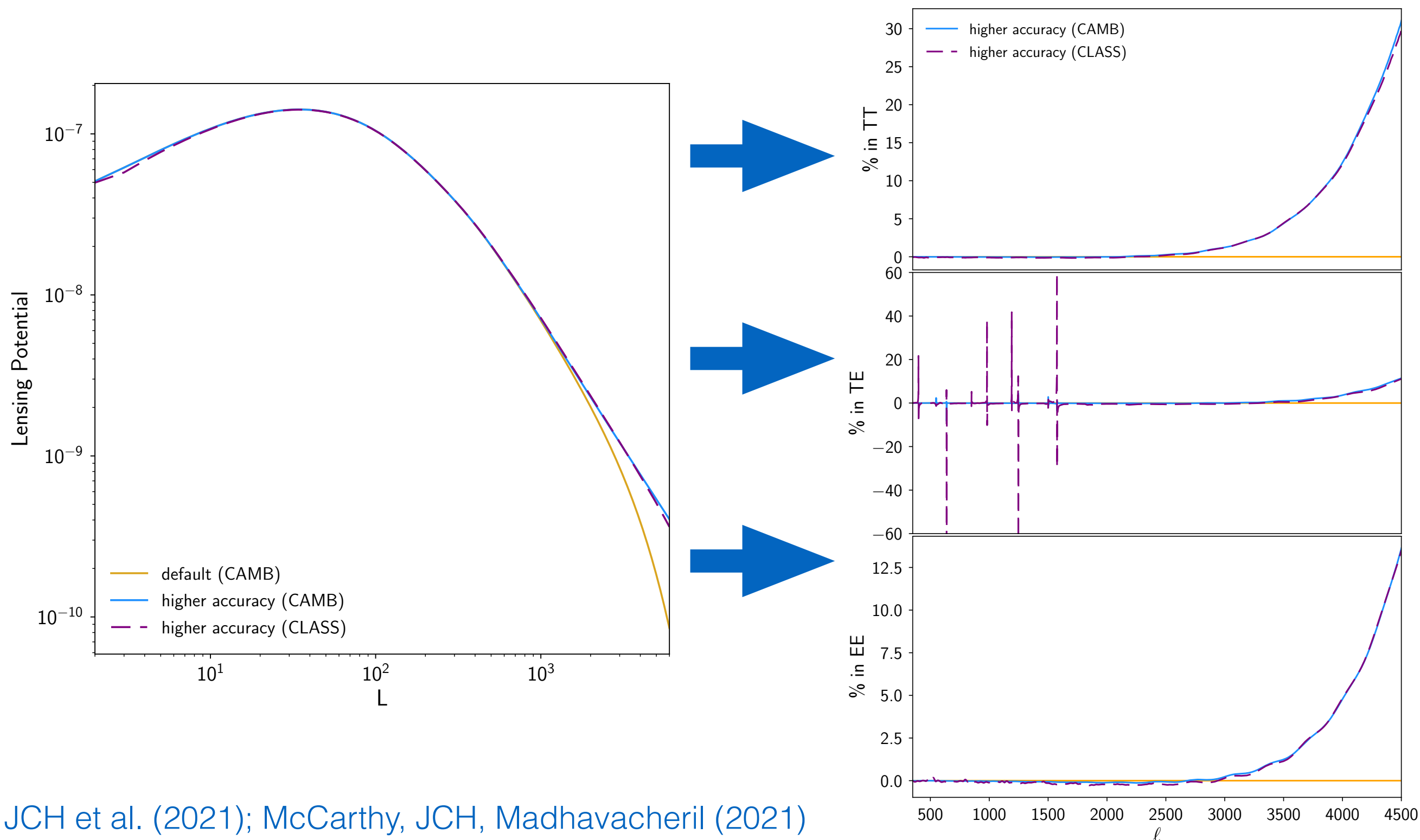
# Theoretical Accuracy

Are the default accuracy settings in CAMB/CLASS OK for ACT/SO?

Almost, but not quite! Higher accuracy needed in lensing calc.

# Theoretical Accuracy

Are the default accuracy settings in CAMB/CLASS OK for ACT/SO?  
Almost, but not quite! Higher accuracy needed in lensing calc.





# Theoretical Accuracy

For ACT DR4, this correction shifts some parameters by  $\sim 0.2-0.3\sigma$

primary parameters affected are

$\Omega_c h^2$  and  $n_s$

but this propagates to  $H_0$  and  $\sigma_8$

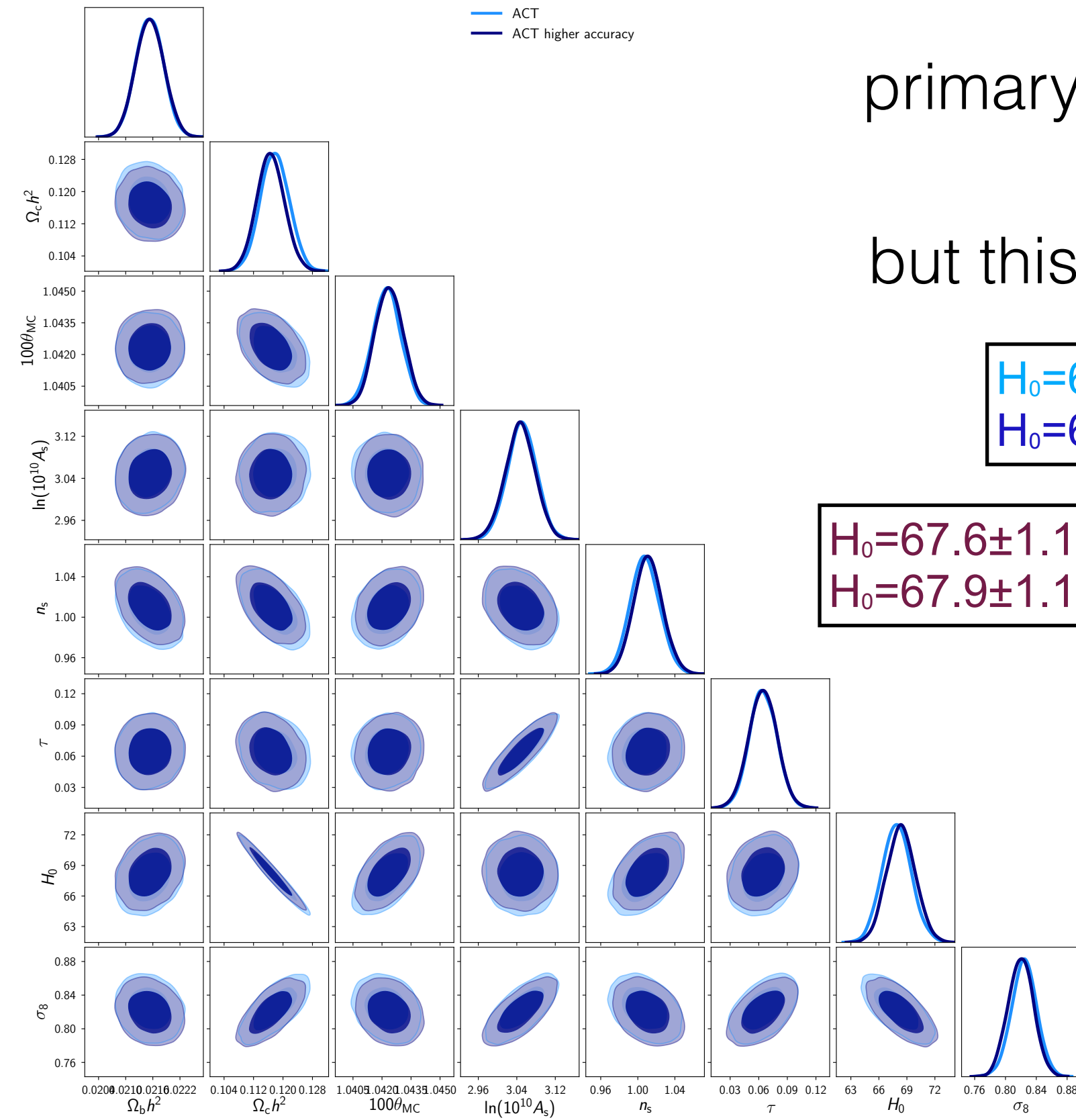
$H_0 = 67.9 \pm 1.5$  km/s/Mpc *ACT (original)*  
 $H_0 = 68.4 \pm 1.5$  km/s/Mpc *ACT (high-acc.)*

$H_0 = 67.6 \pm 1.1$  km/s/Mpc *ACT+WMAP (original)*  
 $H_0 = 67.9 \pm 1.1$  km/s/Mpc *ACT+WMAP (high-acc.)*

JCH et al. (2021)

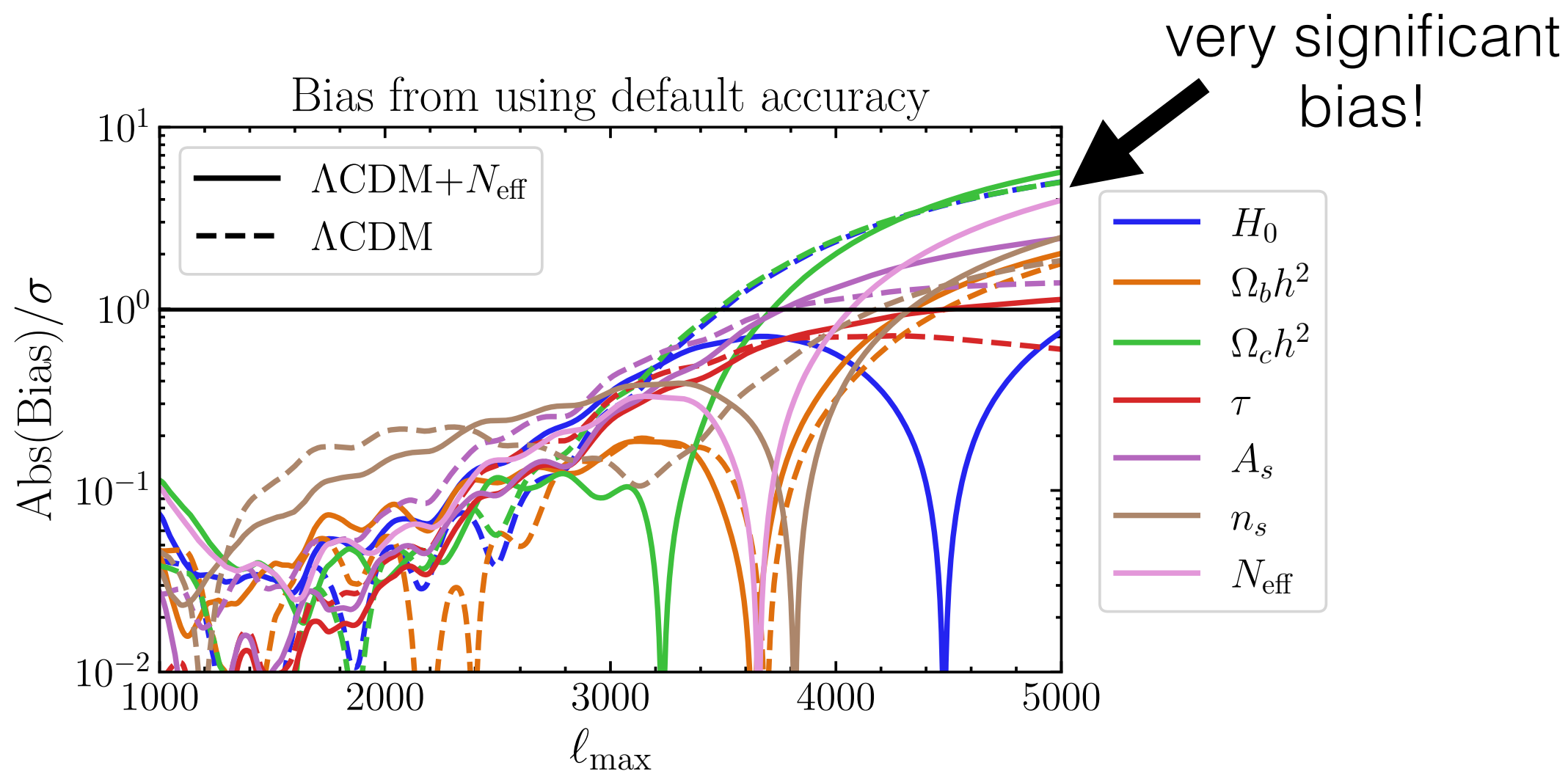
Can be multi- $\sigma$  bias for  
upcoming experiments

McCarthy, JCH,  
Madhavacheril (2021)



# Boltzmann Accuracy

For upcoming experiments, the bias can be very large! Easy to fix: just run CAMB/CLASS with high-accuracy settings (and test them)



Related: baryonic feedback corrections to CMB lensing also become important in this context

McCarthy, JCH, Madhavacheril (2021): [arXiv:2103.05582](https://arxiv.org/abs/2103.05582)



# CosmoPower++

Goal: build emulators using very high-precision CLASS calculations  
— these require 1 minute per evaluation (much slower than default!)

- CMB TT/TE/EE power spectra accurate to  $< 0.5\%$  at all multipoles  $< 10^4$
- Linear  $P(k)$  accurate to  $< 0.5\%$  at all  $k < 50 \text{ h/Mpc}$
- Distance-redshift relation;  $H(z)$
- BAO observables
- Derived parameters ( $\sigma_8$ ,  $\theta_s$ , etc.)
- Factor of 100-1000x speedup per Boltzmann call in MCMC
- NNs are fully differentiable (can be used in gradient-based inference)
- Can be run on GPUs for further acceleration

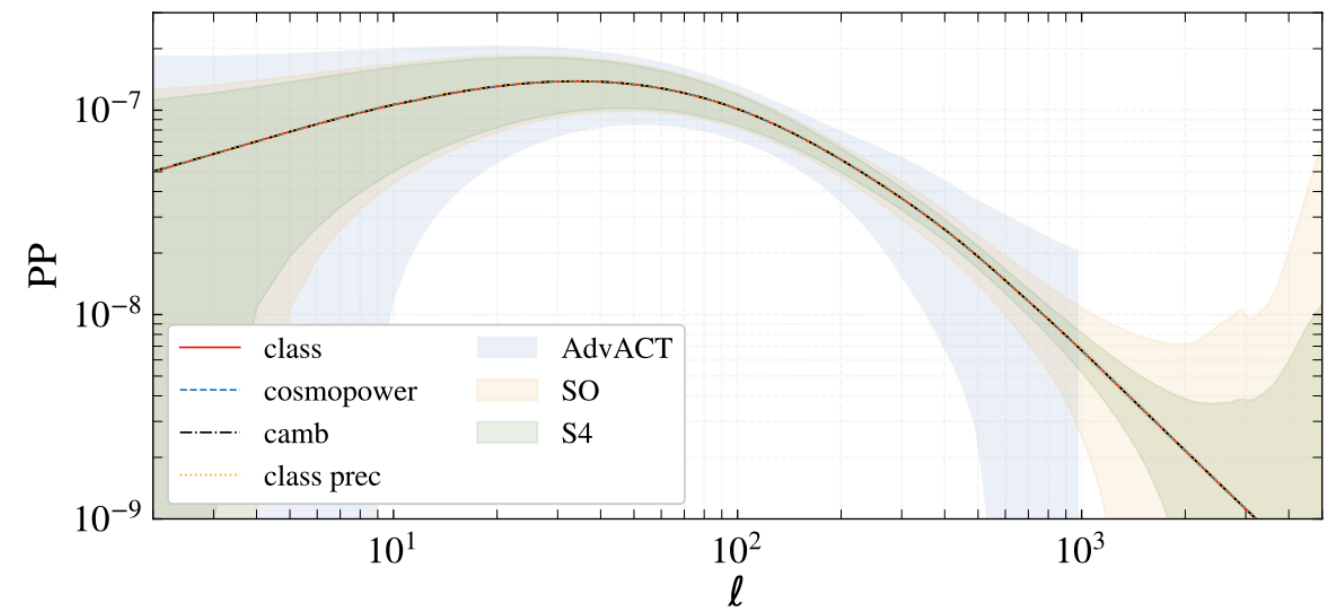
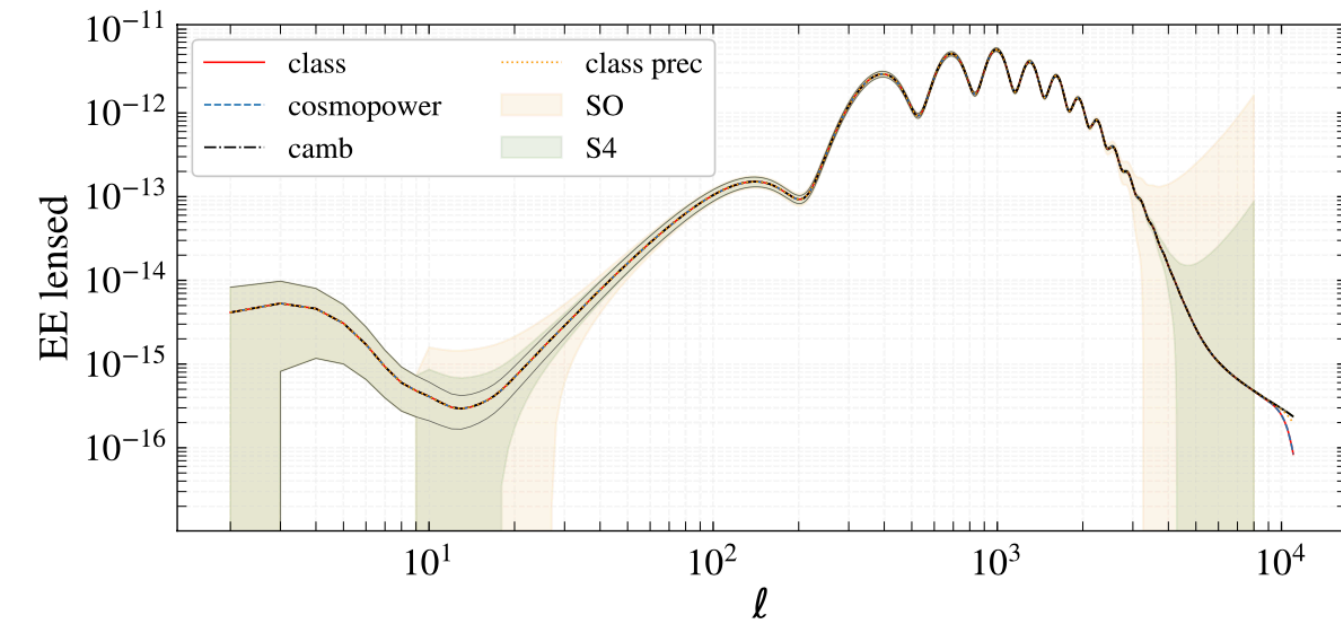
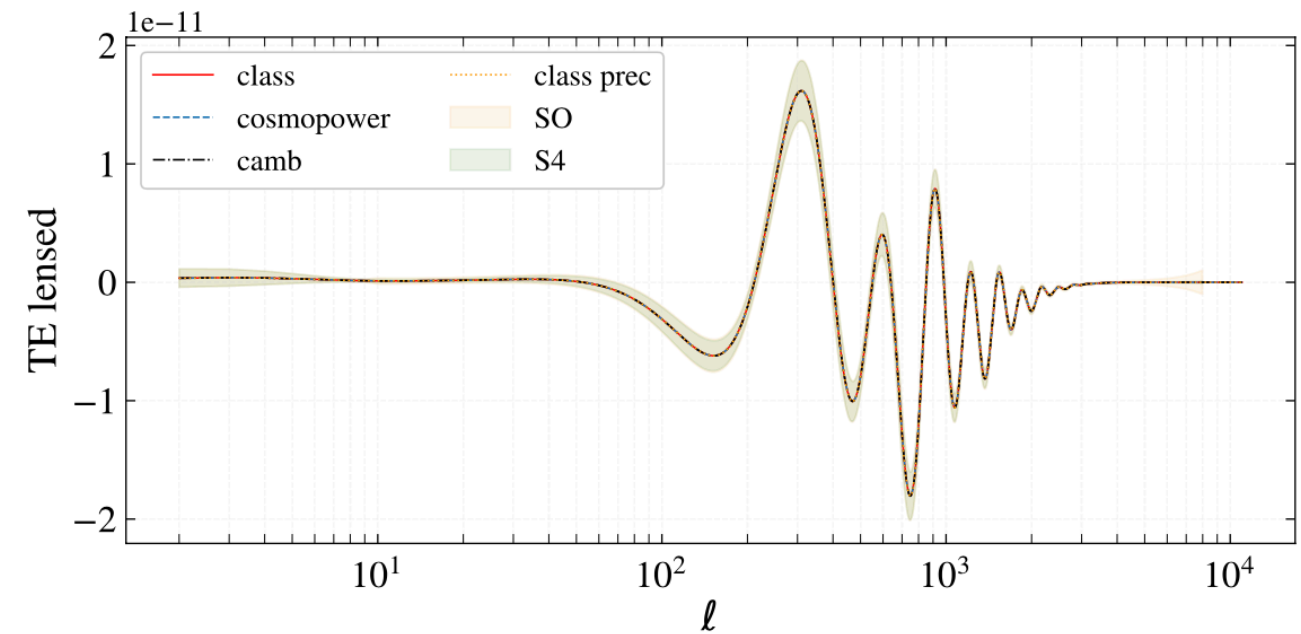
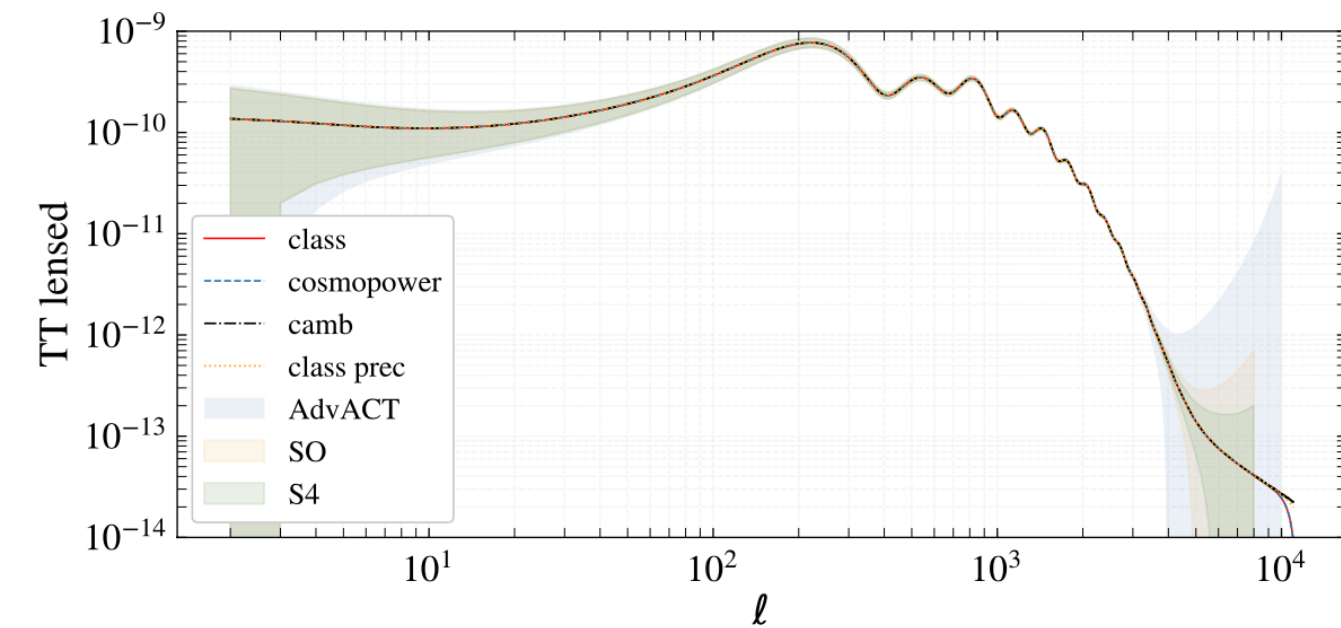
Models run thus far (128,000 parameter sets each):

$\Lambda\text{CDM}$ ,  $+N_{\text{eff}}$ ,  $+M_v$ ,  $+w$

# CosmoPower++

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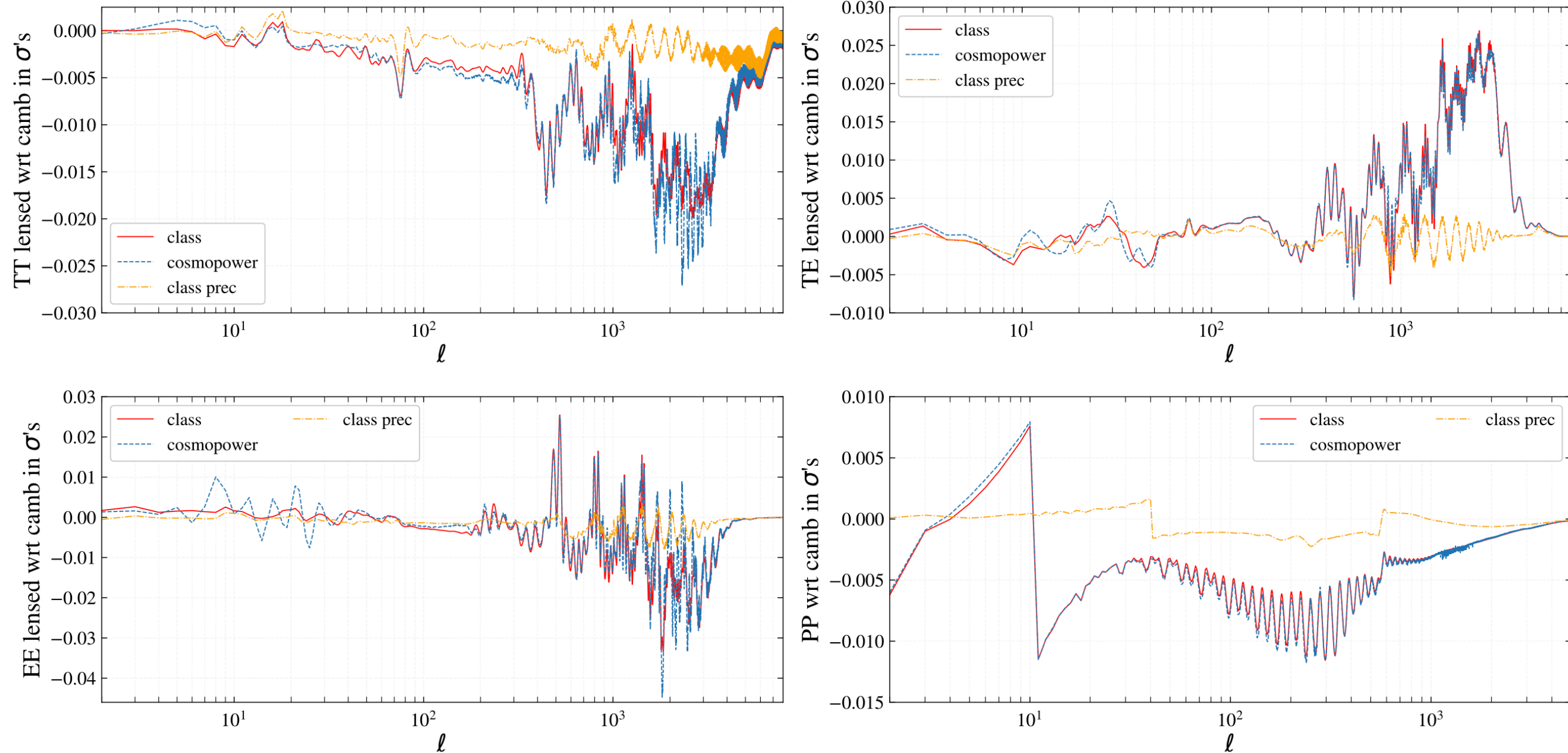
It works :-)





# CosmoPower++

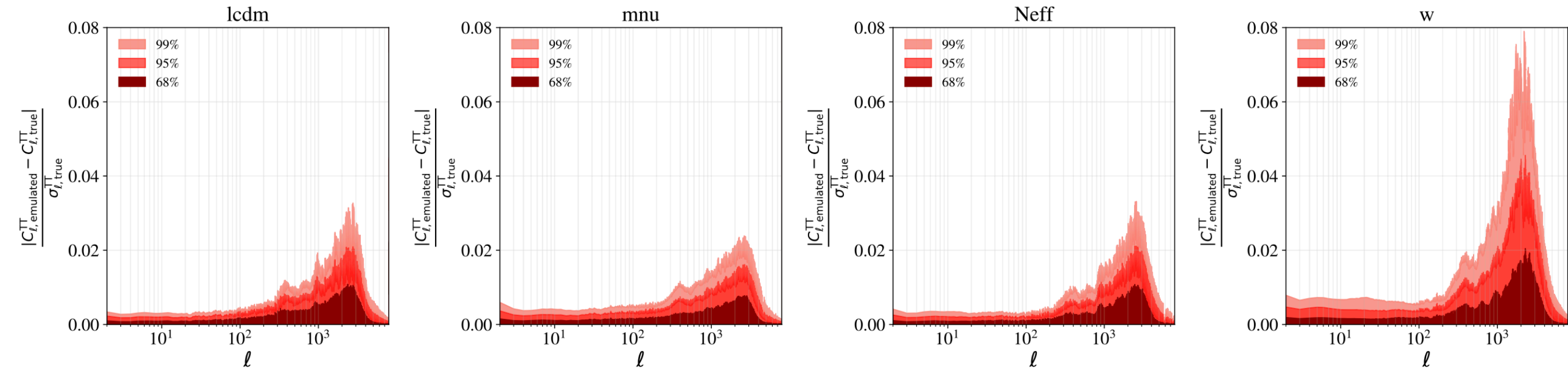
Assess accuracy in terms of forecast CMB-S4 error bars:  $< 0.05\sigma$ !



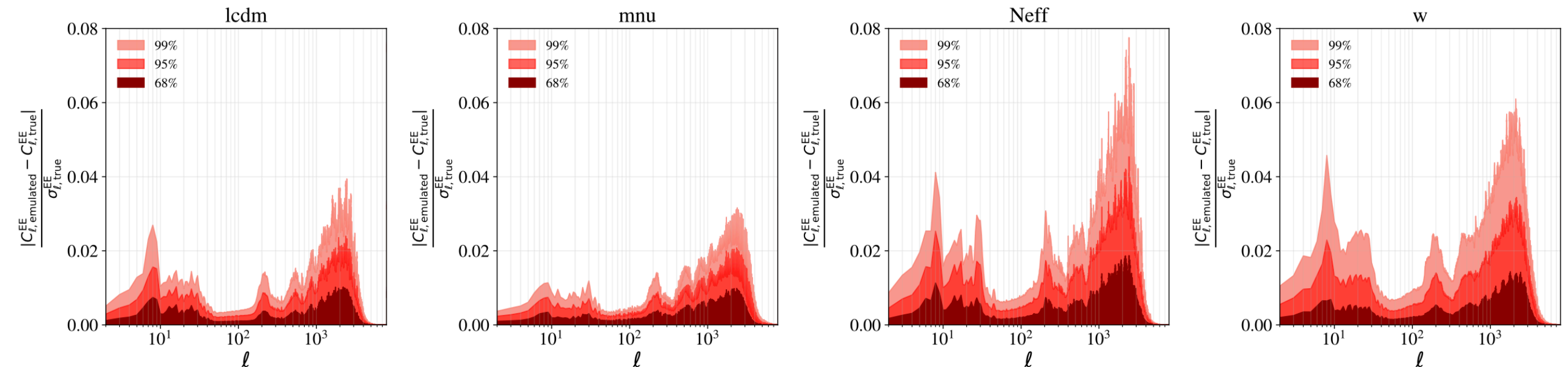
# Validation on Test Set

Assess accuracy in terms of forecast CMB-S4 error bars:  $< 0.07\sigma$ !

## TT power spectrum



## EE power spectrum



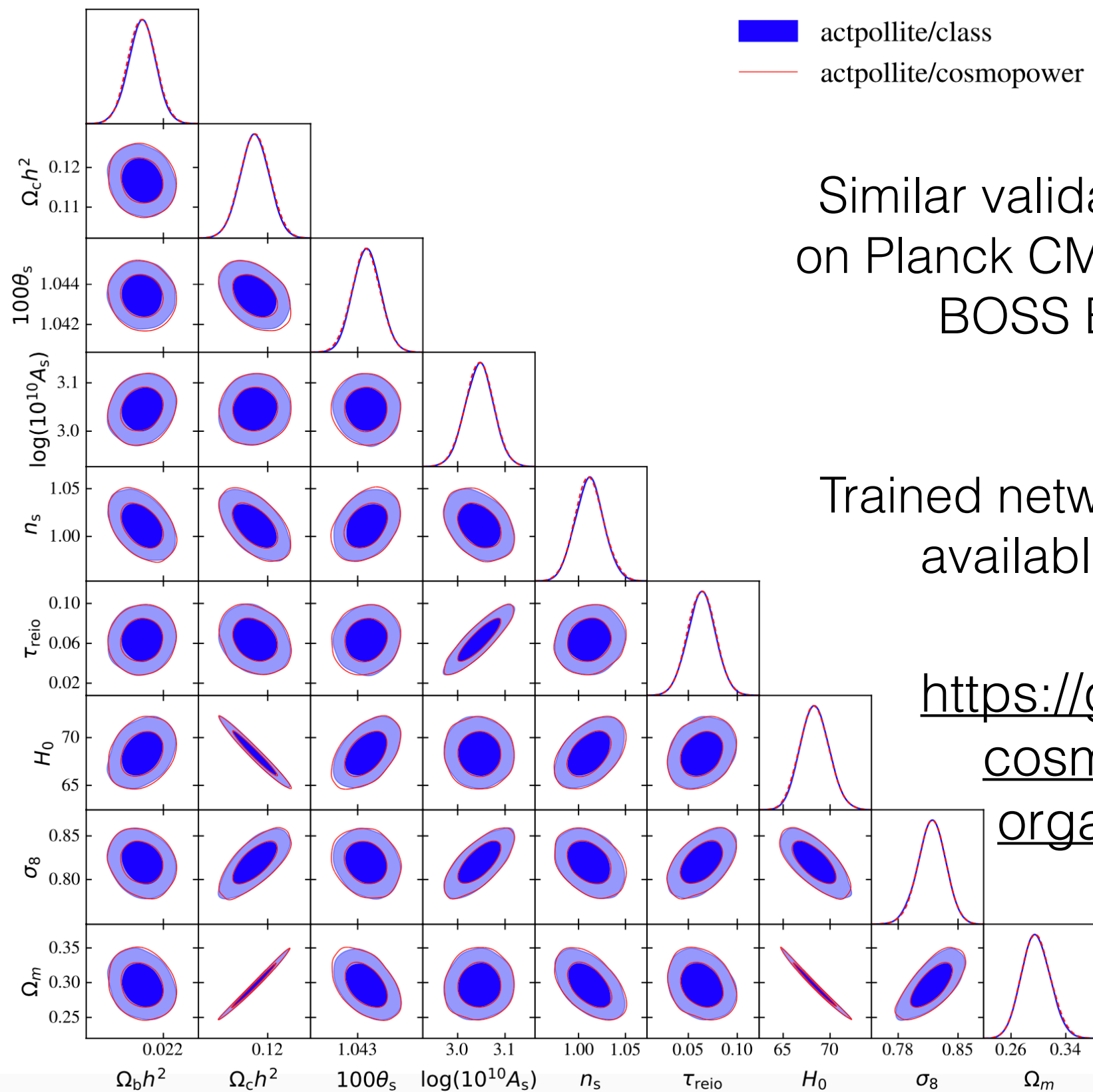


# ACT DR4 Reproduction

~few minutes on laptop vs. ~few days on CCA cluster (!)

# ACT DR4 Reproduction

~few minutes on laptop vs. ~few days on CCA cluster (!)



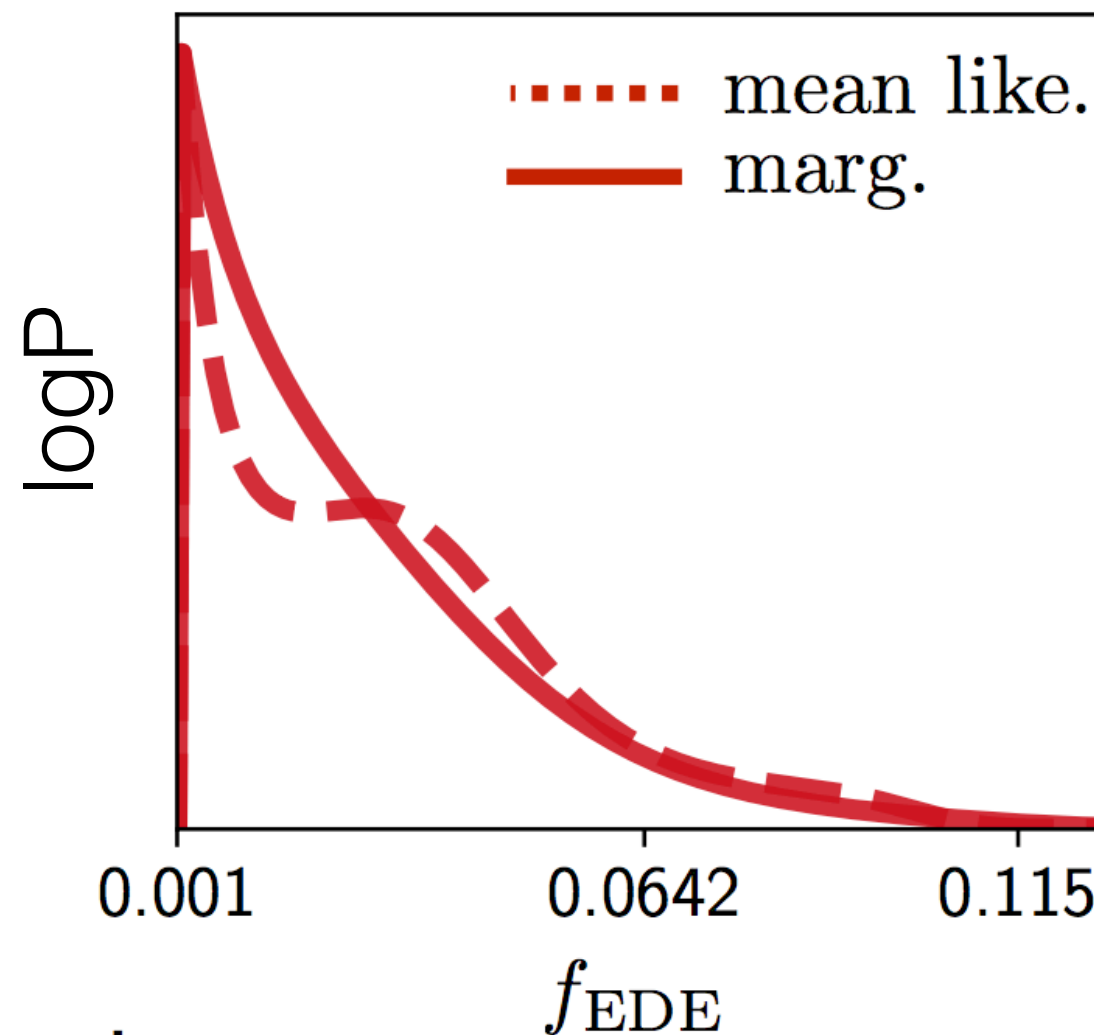
Similar validation performed  
on Planck CMB, CMB lensing,  
BOSS BAO+RSD

Trained networks are publicly  
available via GitHub

[https://github.com/  
cosmopower-  
organization](https://github.com/cosmopower-organization)

# Prior Volume Effects?

non-preference for  $H_0$ -resolving EDE in Planck is robust in either frequentist or Bayesian methodology

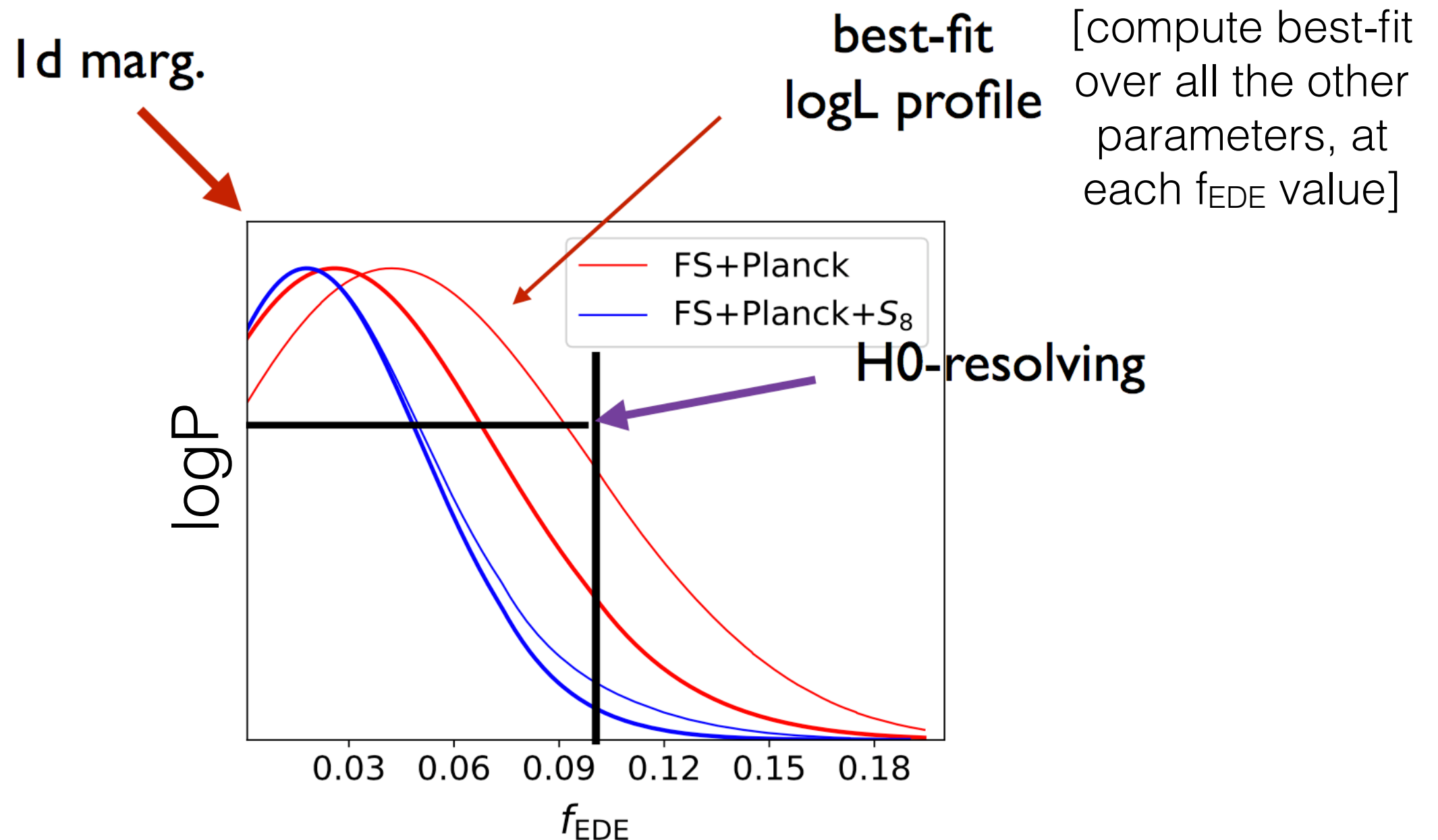


mean like. =  
[compute average  
log(likelihood),  
averaging over  
all the other  
parameters, at  
each  $f_{\text{EDE}}$  value]

**FS+Planck only**

# Prior Volume Effects?

non-preference for  $H_0$ -resolving EDE in Planck is robust in either frequentist or Bayesian methodology



# Prior Volume Effects?

non-preference for H0-resolving EDE in Planck is robust in either frequentist or Bayesian methodology

