

Discussion group report **Multi-leg NLO**

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Contributions

Convenor: G. Passarino

Scalar one-loop 4-point functions with complex masses A. Denner

NLO tensor reduction, where we are J. Fleischer

Automatic calculations, HELAC NLO status report C. Papadopoulos

Interference between Higgs signal and background L. Dixon



A. Denner

Complex mass scheme

Motivation

- Almost all interesting elementary particles are unstable
- description of resonance requires resummation of propagator corrections
- mixing of perturbative order violates gauge invariance in general formally of higher orders but can be drastically enhanced
- many finite-width schemes have been proposed all have some drawbacks
- practical and consistent scheme: complex-mass scheme



The complex-mass scheme

Denner, Dittmaier, Roth, Wackeroth '99, Denner, Dittmaier, Roth, Wieders '05

Basic idea: (renormalized mass)² = location of propagator pole in complex p^2 plane

replace
$$M_{\mathrm{W}}^2 \to \mu_{\mathrm{W}}^2 = M_{\mathrm{W}}^2 - \mathrm{i} M_{\mathrm{W}} \Gamma_{\mathrm{W}}, \qquad M_{\mathrm{Z}}^2 \to \mu_{\mathrm{Z}}^2 = M_{\mathrm{Z}}^2 - \mathrm{i} M_{\mathrm{Z}} \Gamma_{\mathrm{Z}}$$
 and define (complex) weak mixing angle via $\cos^2 \theta_{\mathrm{w}} \equiv c_{\mathrm{w}}^2 = 1 - s_{\mathrm{w}}^2 = \frac{\mu_{\mathrm{W}}^2}{\mu_{\mathrm{Z}}^2}$ \hookrightarrow complex mass renormalization: $M_{\mathrm{W},0}^2 = \mu_{\mathrm{W}}^2 + \delta \mu_{\mathrm{W}}^2, \ldots$

virtues

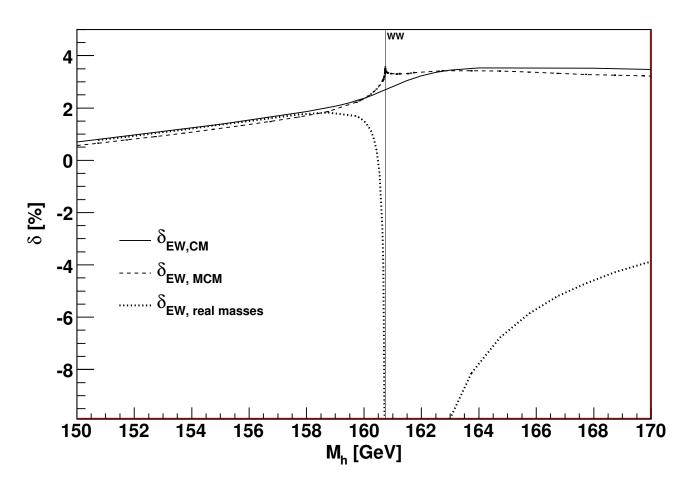
- gauge invariance (Ward identities), gauge-parameter independence hold exactly
- perturbative calculations as usual (with complex masses and counterterms)
- no double counting (bare Lagrangian unchanged!)

drawbacks

- spurious terms of $\mathcal{O}(\Gamma/M)$, of higher order $[\mathcal{O}(\alpha^2)]$ in $\mathcal{O}(\alpha)$ calculation]
- loop integrals with complex masses, now available!



Example: electroweak two-loop corrections to $H \to \gamma \gamma$



Actis, Passarino, Sturm, Uccirati '08

CM = complex-mass scheme

MCM = "minimal" complex-mass scheme

⇒ consistent use of complex masses mandatory!



Scalar one-loop integrals for complex masses

All NLO amplitudes reducible to scalar 1-, 2-, 3-, and 4-point integrals

- regular 1-, 2-, 3-point integrals 't Hooft, Veltman '79:
- regular 4-point integral
 - Dao, Le '09: result in terms of 108 dilogarithms following 't Hooft, Veltman '79
 - Denner, Dittmaier '10: result in terms of 72 dilogarithms based on improved approach of 't Hooft, Veltman '79
 - Denner, Dittmaier '10: result in terms of 32 dilogarithms generalizing approach of Denner, Nierste, Scharf '91
- soft and/or collinear singular 4-point integrals in dimensional, mass and mixed regularizations Denner, Dittmaier '10



J. Fleischer

NLO tensor reduction

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Motivation

- Conventional techniques for NLO calculations rely on reduction of tensor integrals
- simple Passarino—Veltman reduction breaks down when Gram determinants become small or zero
- improved methods are required to circumvent this problem
- various methods are available in the literature
- publically available packages are not complete

goal: provide opensource package for reduction of tensor integrals



Method of Fleischer/Riemann

Fleischer/Riemann arXiv:1009.4436 [hep-ph]

- Tensors are expressed by scalar integrals in higher dimensions following Davydychev '91
- notations for determinants according to Melrose '65
- use recurrence relations for higher-dimensional scalar integrals Tarasov '96, Fleischer '99
- reduce rank r pentagons to rank r-1 pentagons and boxes without inverse Gram determinants
- for small Gram determinants: expansion in Gram determinant improve convergence with Padé approximants

remark AD: used methods are similar to those of other groups

Binoth et al, Denner/Dittmaier, R.K. Ellis et al

differences in details

Introduction

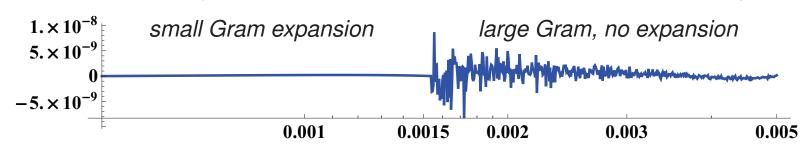
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Numerical implementation of described algorithms: C++ package **PJFry** by **V. Yundin** [in preparation]

- Reduction of 5-point 1-loop tensor integrals up to rank 5
- No limitations on internal/external masses combinations
- Small Gram determinants treatment by expansion
- Interfaces for C, C++, FORTRAN and MATHEMATICA

Example:

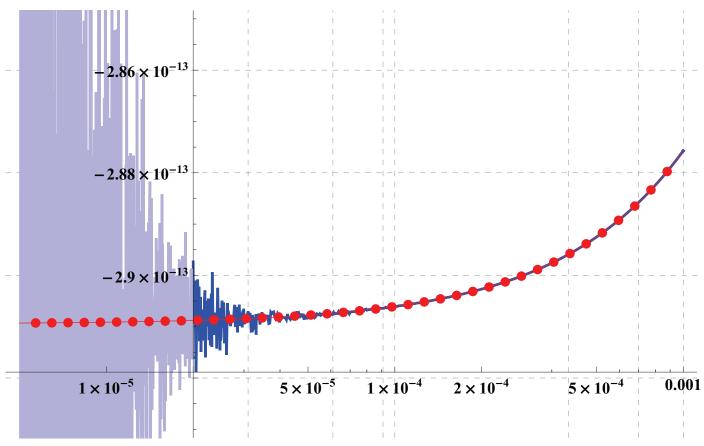
Relative accuracy of E_{3333} coef. around small Gram4 region



PJFry — small Gram region example

Example: E_{3333} coefficient in small Gram region (x = 0)

Comparison of Regular and Expansion formulae:



x=0: $E_{3333}(0,0,-6\times10^4,0,0,10^4,-3.5\times10^4,2\times10^4,-4\times10^4,1.5\times10^4,0,6550,0,0,8315)$

Introduction

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C. Papadopoulos

HELAC-NLO

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HELAC-NLO

Authors: Bevilacqua, Czakon, Garzelli, van Hameren, Malamos, Papadopoulos, Pittau, Worek

General, automatized NLO calculator

ingredients:

- Dyson-Schwinger recursive equations
- Feynman rules in colour-flow representation
- OPP method for virtual corrections (CutTools)
- Rational terms (R2) via "counterterms"
- OneLOop library for scalar one-loop integrals in dimensional regularization
- HELAC-Dipoles: automatic production of Catani-Seymour dipoles
 - for arbitrary QCD and EW processes
 - massive and massless external states
 - ightharpoonup restrictions on phase space (via α parameter)

HOW HELAC-NLO WORKS-VIRTUAL

Generate w=1 events (Les Houches format) using HELAC at tree order. Information included: LH + color assignment, helicity. Optimization!

Calculate using HELAC-1L virtual part for each w=1 event. Produce a new LH file including virtual corrections. Includes UV renormlization

The final LH file can now be used to produce any kinematical distribution!

HELAC-NLO in a Nutshell

- **□** HELAC-PHEGAS
 - > Event generator for all parton level processes @ LO
- **□** HELAC-1LOOP
 - > Evaluation of virtual one-loop amplitudes, based on **HELAC**
- □ CUTTOOLS
 - Reduction of tensor integrals and determination of coefficients via OPP reduction method
- □ ONELOOP
 - > Evaluation of scalar integrals (divergent and finite scalar integrals)
- **□** HELAC-DIPOLES
 - ➤ Catani-Seymour dipole subtraction for massless and massive cases
 - Phase space integration of subtracted real radiation and integrated dipoles
 - Arbitrary polarizations & phase space restriction on dipoles contribution

http://helac-phegas.web.cern.ch/helac-phegas/

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HELAC-NLO

HELAC-NLO has been applied to

- pp $\rightarrow VVV$
- $pp \rightarrow t\bar{t}b\bar{b}$
- $pp \rightarrow t\bar{t}jj$

goal: provide NLO calculator for $2 \rightarrow n$ processes with n = 6, 7 particles attached to the loop

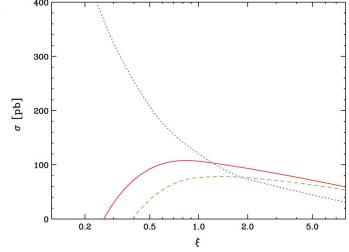
complete software package for NLO-QCD at LHC including virtual and real corrections

pp -> ttjj @ LHC

☐ Scale dependence & integrated cross sections

Bevilacqua, Czakon, Papadopoulos, Worek 10

Process	$\sigma^{ m LO}$ [pb]	Contribution
$pp \to t\bar{t}jj$	120.17(8)	100%
$qg \rightarrow t\bar{t}qg$	56.59(5)	47.1%
$gg \rightarrow t\bar{t}gg$	52.70(6)	43.8%
$qq' \to t\bar{t}qq', q\bar{q} \to t\bar{t}q'\bar{q}'$	7.475(8)	6.2%
$gg \to t\bar{t}q\bar{q}$	1.981(3)	1.6%
$q\bar{q} \to t\bar{t}gg$	1.429(1)	1.2%



 $\sigma_{
m LO} = (120.17 \pm 0.08) \;
m pb$

 $\sigma_{
m NLO} = (106.94 \pm 0.17) \;
m pb$

 $\sigma_{\rm NLO}^{\rm veto} = (76.58 \pm 0.17)~\rm pb$

Scale dependence reduced:

72% @ LO down to 13% @ NLO 54% @ NLO with jet veto of 50 GeV

K factor of K = 0.89 (K = 0.64) Negative shift of 11% (36%)

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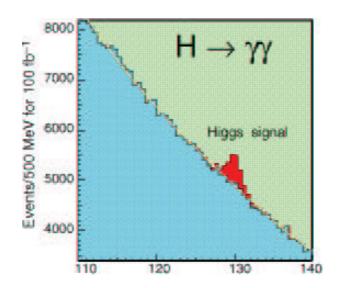


L. Dixon

Higgs signal-background interference

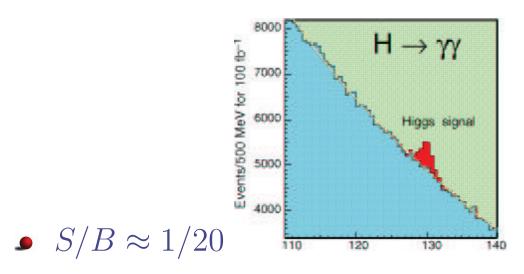
Motivation

- For $M_{
 m H} < 140\,{
 m GeV}$ best decay mode is ${
 m H} o \gamma \gamma$
- QCD background is huge
 - can it contaminate the signal through interference?



- how big is resonance-continuum interference for $gg \to H \to \gamma \gamma$? Dixon, Siu
- general issue when extracting couplings from experimental signal for various production/decay channels.

Back-of-envelope calculation



CMS

- But expt'l resolution $\approx 1 \text{ GeV} \approx 1000 \times \Gamma_H$
- Also, only 1/3 or so of B is from $gg \to \gamma \gamma$
- So intrinsic $S/B \approx 1/20 \times 1000 \times 3 \approx 150$
- Interference effect is $\approx 2\sqrt{B/S} \approx 15\%$

In search of a phase

• Total $gg \to \gamma \gamma$ amplitude

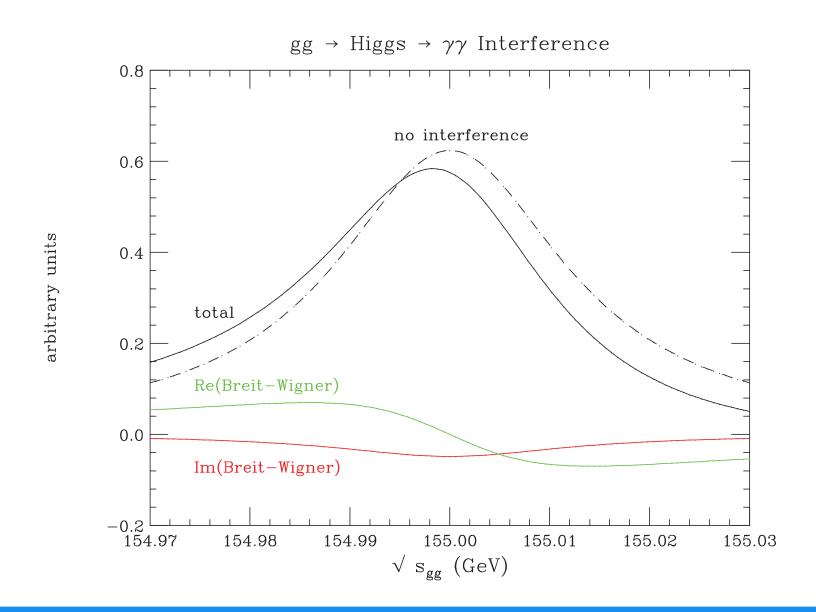
$$\mathcal{A}_{gg\to\gamma\gamma} = \frac{-\mathcal{A}_{gg\to H}\mathcal{A}_{H\to\gamma\gamma}}{\hat{s} - m_H^2 + im_H\Gamma_H} + \mathcal{A}_{cont}$$

Interference term has 2 pieces

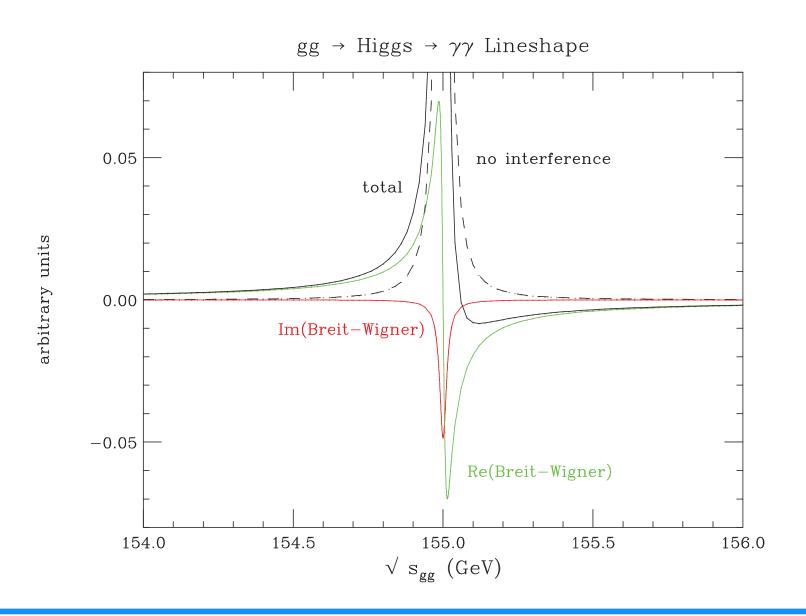
$$\delta \hat{\sigma}_{gg \to H \to \gamma \gamma} = -2(\hat{s} - m_H^2) \frac{\text{Re} (\mathcal{A}_{gg \to H} \mathcal{A}_{H \to \gamma \gamma} \mathcal{A}_{\text{cont}}^*)}{(\hat{s} - m_H^2)^2 + m_H^2 \Gamma_H^2}$$
$$-2m_H \Gamma_H \frac{\text{Im} (\mathcal{A}_{gg \to H} \mathcal{A}_{H \to \gamma \gamma} \mathcal{A}_{\text{cont}}^*)}{(\hat{s} - m_H^2)^2 + m_H^2 \Gamma_H^2}$$

- "Re" term vanishes upon integration over \hat{s} , provided that $\mathcal{A}_{gg\to H}, \, \mathcal{A}_{H\to\gamma\gamma}, \, \mathcal{A}_{\mathrm{cont}}$ do not vary too quickly. Dicus, Willenbrock
- "Im" term needs relative phase, resonance vs. continuum.

Close-up of Higgs resonance



Not-so-close-up of Higgs resonance



Conclusions

- Interference of $gg \to H \to \gamma \gamma$ signal with continuum background at LHC relies on an interplay between phases in QCD and electroweak amplitudes.
- In the SM, 2-loop QCD phases are more important.
- Effect is destructive in the SM, though -(2-6)% is below presently anticipated th. & expt. uncertainties.
- Effect can get larger in MSSM, for example, in regions where the $\gamma\gamma$ signal is still visible.
- Further study of effects in MSSM, other Higgs models, and in selected other channels in the (MS)SM is warranted.