Different Forms of the Generalised Crewther Relation in QCD : Concrete Consequences of Analytical Multiloop Calculations

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Why analytical perturbative calculations (APC) important:

1. Precise check of predictions of gauge theories realised in nature (QCD, QED, EW);
a) Comparison with the data of $e^{+} e^{-}$colliders (Novosibirsk, Beijing, Frascati); data of LEP still exist; b) Lepton-hadron DIS (CEBAF-JLAB working; lots of HERA, NOMAD (CERN) data still not studied, COMPASS (CERN); $\nu N$ DIS Minerva (Fermilab) works;
c) Tevatron (existing) and new LHC data are appearing
2. New Theoretical Effects/Relations found in APC

Among them Generalized quark-parton Crewther Relation (72) in $\bar{M} S$-scheme Broadhurst, Kataev (93) in NNLO;
Crewther (97) all orders proof; Baikov, Chetyrkin, Kuhn (10) in $\mathrm{N}^{3} \mathrm{LO}$; Kataev, Mikhailov (10) $\mathrm{N}^{3} \mathrm{LO} M \bar{M}$ resummed representation.
Physics which follows from this Generalized Relation between Characteristics of $e^{+} e^{-}$and DIS and Axial-Vector-Vector-Vector Triangle Diagram
a) Discovery of the Consequences of the Conformal Symmetryrelations between concrete parts of APC ;
b) Manifestation of Conformal Symmetry Breaking in QCD, related to exp. detectable Effect of Asymptotic Freedom in QCD ;
Responsible for Definition of Energy Momentum Tensor in High Orders PC- Important for Self-Consistent Formulation of Gauge Theories ; Both are encoded in Factor $\beta\left(a_{s}\right) / a_{s} a_{s}=\alpha_{s} / \pi$.

## Notion of Conformal Symmetry and its PT Breaking

Conformal Symmetry (CS) is the generalisation of Poincaré symmetry ; namely the symmetry under following transformations

- Space-time translations $x^{\prime}{ }^{\mu}=x^{\mu}+\alpha^{\mu}$;
- Lorentz transformations $x^{\prime \mu}=\Lambda_{\nu}^{\mu} x^{\nu}$;
- Special conformal transformations $x^{\prime}{ }^{\mu}=\frac{x^{\mu}+\beta^{\mu} x^{2}}{1+2 \beta x+\beta^{2} x^{2}}$
- Scale transformations $x^{\prime \mu}=\rho x^{\mu}$

In PT this symmetry is broken by the procedure of renormalizations and and non-zero $\mathrm{RG} \beta$-function.

Where in higher orders of APC one can discover the support of importance of both CS and CSB effects ???

Example of guess using notion on structure of CSB effects
Consider $D_{A}^{V}\left(a_{s}\right)=Q^{2} \int_{0}^{\infty} \frac{R^{e+e^{-}}(s)}{\left(s+Q^{2}\right)^{2}} d s=D_{A}^{N S}\left(a_{s}\right)+D_{A}^{S I}\left(a_{s}\right)$
First part evaluated analytically up to $a_{s}^{4}$ in $S U\left(N_{c}\right)$ (BChK (10))
$D_{A}^{N S}\left(a_{s}\right)=d_{R}\left(\sum_{F} Q_{F}^{2}\right) C_{A}^{N S}\left(a_{s}\right)=d_{R}\left(\sum_{F} Q_{F}^{2}\right)\left(1+\sum_{n=1}^{4} d_{n} a_{s}^{n}+\ldots\right)$
5 -loop diagrams for Green function of VV currents ;
$D_{A}^{S I}\left(a_{s}\right)=\left(\sum_{F} Q_{F}\right)^{2} C_{A}^{S I}\left(a_{s}\right)=\left(\sum_{F} Q_{F}\right)\left(\sum_{n=3}^{4} d_{n}^{S I} a_{s}^{n}\right)$
SI-contributions to Green function of VV currents, $d_{3}^{S I}$ known;
Gorishny,Kataev,Larin (91); Surguladze, Samuel (91)
!! Example of guess on structure of CSB effects in generalized
Crewther relation !! Theoretical prediction Kataev (11)
$d_{4}^{S I}=d^{a b c} d^{a b c}\left(C_{F}\left(-\frac{13}{64}-\frac{\zeta_{3}}{4}+\frac{5 \zeta_{5}}{8}\right)+C_{A}\left(\frac{481}{1152}-\frac{971}{1152} \zeta_{3}+\frac{295}{576} \zeta_{5}-\frac{11}{32} \zeta_{3}^{2}\right)\right.$
$\left.+\left(T_{F} N_{F}\right)\left(-\frac{119}{1152}+\frac{67}{288} \zeta_{3}-\frac{35}{144} \zeta_{5}+\frac{1}{8} \zeta_{3}^{2}\right)\right)$ from guess on CSB effects and Baikov, Chetyrkin, Kuhn (10) evaluations of $a_{s}^{4}$ corrections to the Gross-Llewellyn Smith sum rule for $\nu N$ DIS Waiting for the direct APC by BChK- work in progress !!

DIS characteristics: Bjorken polarized sum rule
$\int_{0}^{1} g_{1}^{l p-l n}\left(x, Q^{2}\right) d x=\frac{g_{a}}{6} C_{B j p}^{N S}\left(a_{s}\right)=\frac{g_{a}}{6}\left(1+\sum_{n}^{4} a_{s}^{n} c_{n}+\ldots\right)$
Measured at CEBAF at intermediate and low $Q^{2}$; Like the coefficients $d_{1}-d_{4}$ of $e^{+} e^{-}$-annihilation characteristic $D_{A}^{N S}\left(a_{s}\right)$
$c_{1}-c_{4}$ known analytically, depend from powers of $S U\left(N_{c}\right)$ structures
$C_{F}, C_{A}, T_{F} N_{F}, C_{F}^{2}, C_{F} C_{A}, C_{F} T_{F} N_{F}$, and $\frac{d_{F}^{a b c d} d_{A}^{a b c d}}{d_{R}}, \frac{N_{F} d_{F}^{a b c d} d_{F}^{a b c d}}{d_{R}}$
(in $c_{4}$ and $d_{4}$ )
GLS sum rule : $\frac{1}{2} \int_{0}^{1} F_{3}^{\nu p+\bar{\nu} p}\left(x, Q^{2}\right) d x=3 C_{G L S}\left(a_{s}\right)$
$C_{G L S}\left(a_{s}\right)=C_{G L S}^{N S}\left(a_{s}\right)+C_{G L S}^{S I}\left(a_{s}\right)=C_{B j p}^{N S}\left(a_{s}\right)+C_{G L S}^{S I}\left(a_{s}\right)$,
$C_{G L S}^{S I}=\sum_{n=3}^{4} c_{n}^{S I} a_{s}^{n}+\ldots($ BChK (10))
Ellis-Jaffe polarized sum rule
$\int_{0}^{1} g_{1}^{l p}\left(x, Q^{2}\right) d x=C_{B j p}^{N S}\left(a_{s}\right)\left(\frac{1}{12} a_{3}+\frac{1}{36} a_{8}\right)+C_{E J p}^{S I}\left(a_{s}\right) \frac{1}{9} \Delta \Sigma\left(Q^{2}\right)$
where $a_{3}, a_{8}$ and $\Delta \Sigma$ in the $\bar{M} S$ are defined through polarized parton distributions $\Delta u, \Delta d, \Delta s, C_{E J P}^{S I}\left(a_{s}\right)=1+\sum_{n=1}^{3} c_{n}^{E J, S I} a_{s}^{3}$ Larin, van Ritbergen and Vermaseren (97) - expansion of $\int_{0}^{a_{s}}\left(\gamma_{S I}(a) / \beta(a)\right) d a$ taken into account in defining $c_{n}^{E J,}{ }^{S I}$.

Generalized $\bar{M} S$-scheme Crewther relation(s)
$C_{A}^{N S}\left(a_{s}\left(Q^{2}\right)\right) \times C_{B j p}^{N S}\left(a_{s}\left(Q^{2}\right)\right)=1+\Delta_{C S B}\left(a_{s}\left(Q^{2}\right)\right)$
$\Delta_{\mathrm{CSB}}\left(a_{s}\right)=\left(\frac{\beta\left(a_{s}\right)}{a_{s}}\right) \mathcal{P}\left(a_{s}\right)$ where $\mathcal{P}\left(a_{s}\right)=\sum_{m \geq 1} K_{m} a_{s}^{m}$ (Note
$\left(\beta\left(a_{s}\right) / a_{s}\right)$ enters into conformal anomaly- anomalous dimension of
Trace of Energy Momentum Tensor!) $K_{1}=K_{1}[1,0,0] \mathrm{C}_{\mathrm{F}}$,
$K_{2}=K_{2}[2,0,0] \mathrm{C}_{\mathrm{F}}^{2}+K_{2}[1,1,0] \mathrm{C}_{\mathrm{F}} \mathrm{C}_{\mathrm{A}}+K_{2}[1,0,1] \mathrm{C}_{\mathrm{F}} \mathrm{T}_{\mathrm{F}} \mathrm{N}_{\mathrm{F}}$ notice
$\mathrm{T}_{\mathrm{F}} \mathrm{N}_{\mathrm{F}}$-dependence (!) Broadhurst, Kataev (93)
$K_{3}=K_{3}[3,0,0] \mathrm{C}_{\mathrm{F}}^{3}+K_{3}[2,1,0] \mathrm{C}_{\mathrm{F}}^{2} \mathrm{C}_{\mathrm{A}}+K_{3}[1,2,0] \mathrm{C}_{\mathrm{F}} \mathrm{C}_{\mathrm{A}}^{2}+$
$K_{3}[2,0,1] \mathrm{C}_{\mathrm{F}}^{2} \mathrm{~T}_{\mathrm{F}} \mathrm{N}_{\mathrm{F}}+K_{3}[1,1,1] \mathrm{C}_{\mathrm{F}} \mathrm{C}_{\mathrm{A}} \mathrm{T}_{\mathrm{F}} \mathrm{N}_{\mathrm{F}}+K_{3}[1,0,2] \mathrm{C}_{\mathrm{F}} \mathrm{T}_{\mathrm{F}}^{2} \mathrm{~N}_{\mathrm{F}}^{2}$
BChK (10), contain $K_{3}[1,0,2]$-term evaluated in BK (93)
$K_{1}$ contain rational numbers, $\zeta_{3} ; K_{2}$ - rational, $\zeta_{3}, \zeta_{5}$,
$K_{3}$ - rational, $\zeta_{3}, \zeta_{5}, \zeta_{7}$.
In this approximation $\beta\left(a_{s}\right)$ is 3 -loop QCD $\beta$-function, evaluated by Tarasov, Vladimirov, Zharkov (1980).

How to use for checks: 1) Strong check of validity of complicated calculations of $a_{s}^{4}$ coefficients in $C_{A}^{N S}\left(a_{s}\left(Q^{2}\right)\right)$ and $C_{B J P}^{N S}\left(a_{s}\left(Q^{2}\right)\right)$ by BChK (10))
a) cancellation of $\frac{d_{F}^{a b c d} d_{A}^{a b c d}}{d_{R}}, \frac{N_{F} d_{F}^{a b c d} d_{F}^{a b c d}}{d_{R}}$ and $C_{F}^{4}$ and of some other contributions to $d_{4}$ and $c_{4}$ - consequence of Crewther relation, valid in the Conformal Invariant Limit $\left(\beta\left(a_{s}\right)=0\right)$
b) Possibility to compare with resumming $\mathbf{T}_{\mathbf{F}} \mathbf{N}_{\mathbf{F}}$-contributions to $\mathcal{P}\left(a_{s}\right)$ polynomial representation of generalized Crewther
$\Delta_{\mathrm{CSB}}=\sum_{n \geq 1} \sum_{r \geq 1}\left(\frac{\beta\left(a_{s}\right)}{a_{s}}\right)^{n} P_{n}^{(r)}[\mathrm{k}, \mathrm{m}] C_{F}^{\mathrm{k}} C_{A}^{\mathrm{m}} a_{s}^{r}, r=\mathrm{k}+\mathrm{m}$
Kataev, Mikhailov (10) valid at $a_{s}^{4}$-level In high-orders may (or may not) get additional $d^{a b c d} d^{a b c d}$ contributions. c): Consequence: fixation of $\beta\left(N_{F}=n_{0}\right)=0$, namely choice $\mathrm{T}_{\mathrm{F}} \mathrm{N}_{\mathrm{F}}=(11 / 4) \mathrm{C}_{\mathrm{A}}$ leads to $\mathrm{d}_{4}\left(n_{0}\right)+\mathrm{k}_{4}\left(n_{0}\right)=-\frac{333}{1024} \mathrm{C}_{\mathrm{F}}^{4}+\mathrm{C}_{\mathrm{A}} \mathrm{C}_{\mathrm{F}}^{3}\left(-\frac{1661}{3072}+\frac{1309}{128} \zeta_{3}-\frac{165}{16} \zeta_{5}\right)+$ $+\mathrm{C}_{\mathrm{A}}^{2} \mathrm{C}_{\mathrm{F}}^{2}\left(-\frac{3337}{1536}+\frac{7}{2} \zeta_{3}-\frac{105}{16} \zeta_{5}\right)+\mathrm{C}_{\mathrm{A}}^{3} \mathrm{C}_{\mathrm{F}}\left(-\frac{28931}{12288}+\frac{1351}{512} \zeta_{3}\right)$
This expression agrees with result, which follows from BChK (10)

- additional confirmation of the result of complicated calculations

How to use for predictions: 1) The comparison of generalized Crewther relations
$C_{A}^{N S}\left(a_{s}\left(Q^{2}\right)\right) \times C_{B j p}^{N S}\left(a_{s}\left(Q^{2}\right)\right)=1+\Delta_{c s b}^{N S}\left(a_{s}\left(Q^{2}\right)\right)$
$C_{A}^{V}\left(a_{s}\left(Q^{2}\right)\right) \times C_{G L S}\left(Q^{2}\right)=1+\Delta_{c s b}^{V, G L S}\left(a_{s}\left(Q^{2}\right)\right)$
and the guess of "universality"
$C_{A}^{N S}\left(a_{s}\left(Q^{2}\right)\right) \times C_{B j p}^{N S}\left(a_{s}\left(Q^{2}\right)\right)=C_{A}^{V}\left(a_{s}\left(Q^{2}\right)\right) \times C_{G L S}\left(Q^{2}\right)$
which seems to follow from application of OPE method to AVV triangle amplitude allows to get already described theoretical prediction for $a_{s}^{4}$ correction to SI contribution to $D_{A}^{V}\left(a_{s}\right)=Q^{2} \int_{0}^{\infty} \frac{R^{e^{+} e^{-}}(s)}{\left(s+Q^{2}\right)^{2}} d s=D_{A}^{N S}\left(a_{s}\right)+D_{A}^{S I}\left(a_{s}\right)$ Kataev
2) Application of the generalized Crewther relation in the SI
channel gives possibility to prove the following identity in theCI-limit $C_{A}^{S I}\left(a_{s}\right) \times C_{E J p}^{S I}\left(a_{s}\right)=C_{A}^{N S}\left(a_{s}\right) \times C_{B j p}^{N S}\left(a_{s}\right)$ which gives $\left.C_{E J p}^{S I}\left(a_{s}\right)\right|_{C I-l i m i t}=\left.C_{B j p}^{N S}\left(a_{s}\right)\right|_{C I-l i m i t}$ Kataev (10) Planned calculations by Baikov, Chetyrkin, Kuhn of $a_{s}^{4}$ terms to Ellis-Jaffe sum rule have prediction for control of future APC.

## Conclusion

There are a lot of interesting problems for future APC (analytical perturbative calculations):

1) Study of the consequences of Conformal Symmetry and

Conformal Symmetry Breaking
2) Closely related to study of the possibility of the evaluation of high order corrections to Triangle Diagrams with two large external momentum (at present some studies already exist, but at 2-loop level- (Braguta and Onishchenko (04); F. Jegerlehner and O. V. Tarasov (06)), etc.- CI-limit.
3) Studies of generalized Crewther relation may give additional insight on some special features of schemes and scheme-dependence problem (Garkusha and Kataev (11))

