

# Canonical Approach for Extreme QCD

## Seung-il Nam

Department of Physics, Pukyong National University (PKNU),  
Center for Extreme Nuclear Matters (CENuM), Korea University



## In collaboration with

Dr. Masayuki Wakayama (Kokushikan Univ., Japan)

Prof. A Hosaka (RCNP, Osaka Univ., Japan)

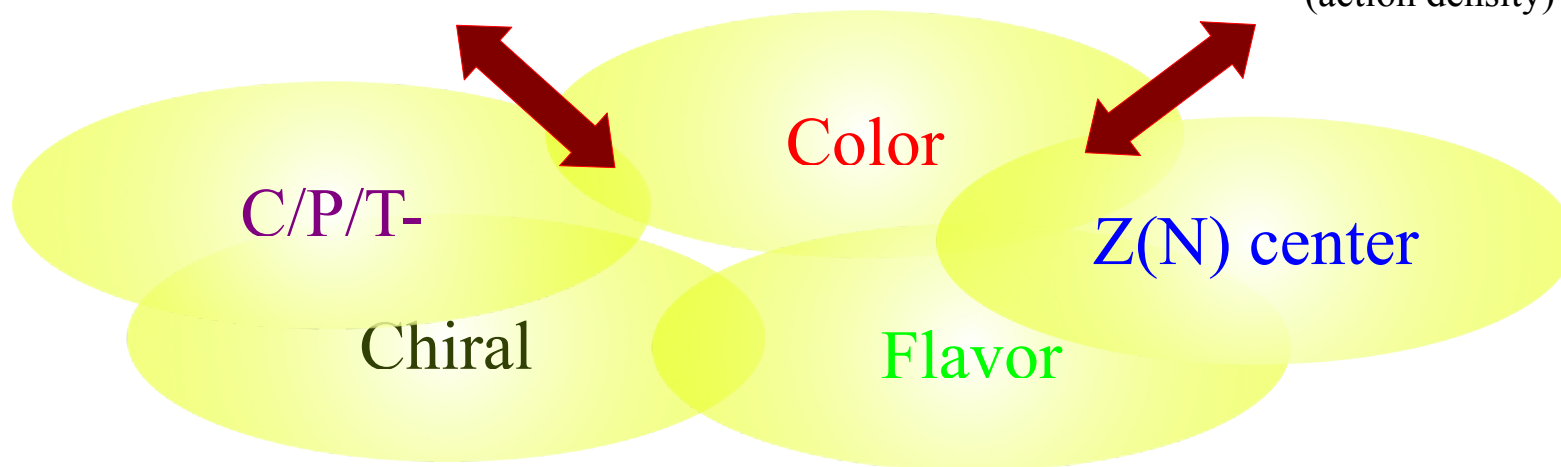
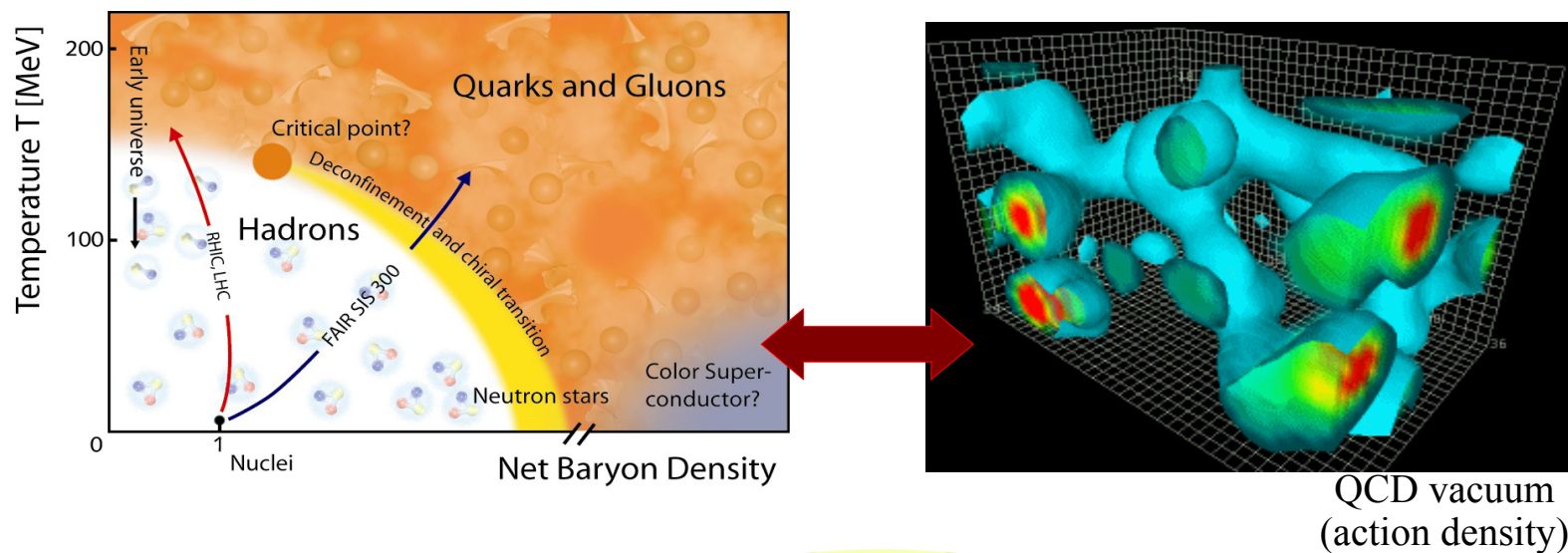


## Contents based on

Physical Review D 102, 034035 (2020)

# Introduction

- QCD has complicated phase structure: Not fully understood



## Introduction

- QCD has complicated phase structure: Not fully understood



C/P/T-

Chiral

Color

Flavor

Z(N) center

## Introduction

- LQCD suffers from **sign problem** at finite quark density
- Compute the following in LQCD with **probability**  $(\det D(U)) e^{-S_G[U]}$

$$\lim_{N \rightarrow \infty} \frac{1}{N} \sum_{i=1}^N \mathcal{O}'(U_i)$$

- If quark density is finite, the quark part of the weight becomes

$$D(\mu_q) = \not{D} + m + \mu_q \gamma_0$$

$$D^\dagger(\mu_q) = -\not{D} + m + \mu_q^* \gamma_0 = \gamma_5 D^\dagger(-\mu_q^*) \gamma_5$$

$$\{\det[D(\mu_q)]\}^* = \det[D^\dagger(\mu_q)] = \det[\gamma_5 D(-\mu_q^*) \gamma_5] = \det[D(-\mu_q^*)]$$

If  $\mu$  is real,  $\det[D]$  is not real but complex



## 2. Application Lattice **QCD**

- In addition, if it is a complex, then we have

$$\int dU O'(U)(R + iI)e^{-S_G} \sim \int dU O'(U)e^{-S_G + i\phi}$$

- It's oscillation to cancel out the integral: **Sign problem**
- Notorious problem in strongly interacting **fermion** systems even in condensed matter, QFT, and nuclear physics as well.
- How to solve the sign problem???
- So far, there have been no cures (NP-hard problem)
- Many indirect and approximated methods developed:
  - Small  $\mu$  expansion, complex-Langevin method, etc.

## Introduction

- **Canonical method** to overcome this problem:  $\mu_q \rightarrow i\mu_{qI}, \mu_{qI} \in \mathbb{R}$

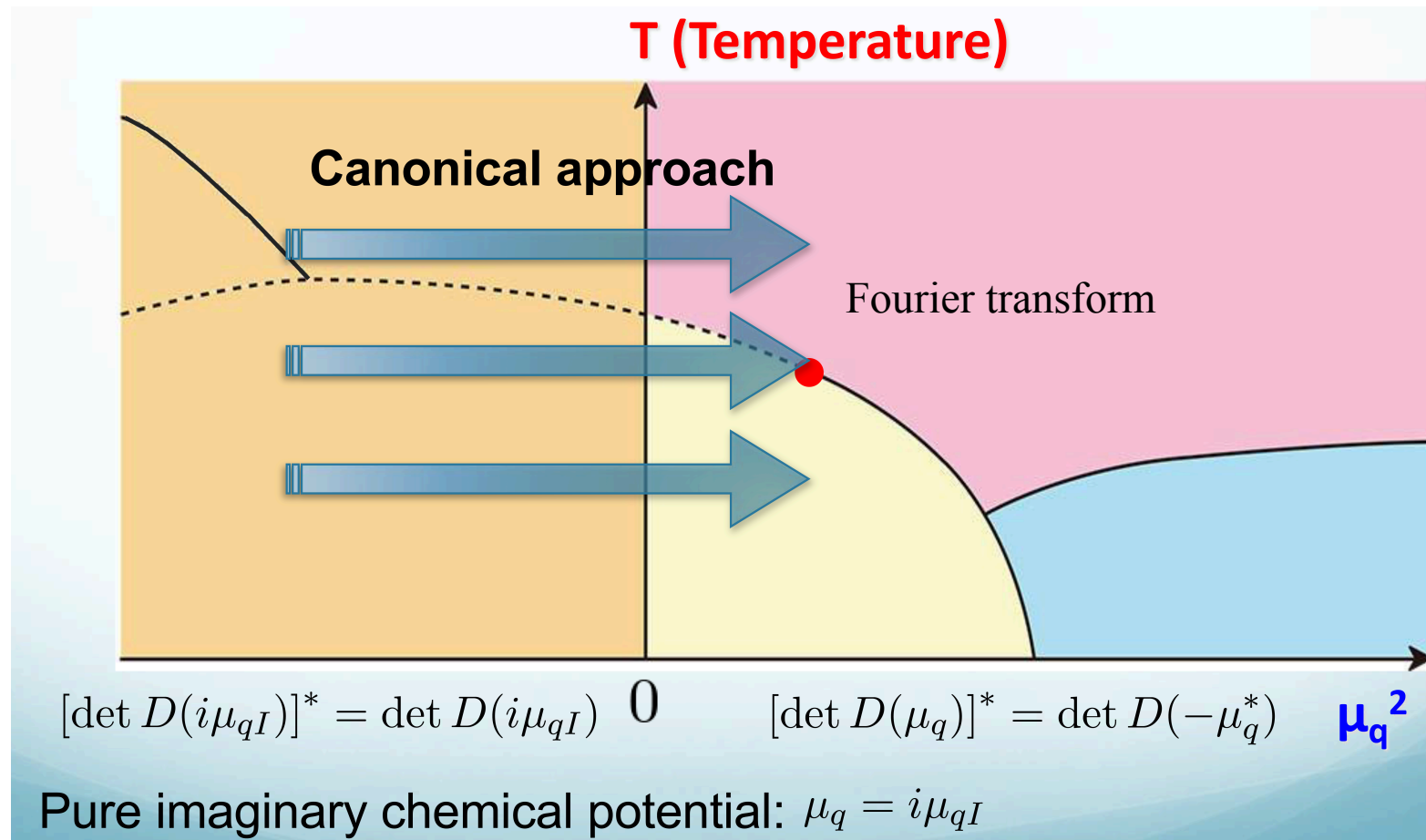


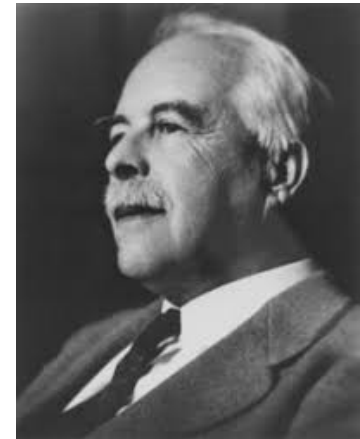
Figure by Dr. Wakayama

## 2. Application Lattice **Q**CD

- Fugacity expansion of grand canonical partition function

$$Z_{\text{GC}}(\mu, T, V) = \sum_{n=-\infty}^{\infty} Z_C(n, T, V) \xi^n$$

$\xi (\equiv e^{\mu/T})$   
Fugacity



Gilbert Newton Lewis

- Obtain canonical partition function by Fourier transform

$$Z_C(n, T, V) = \int_{-\pi}^{\pi} \frac{d\theta}{2\pi} e^{-in\theta} Z_{\text{GC}}(\mu = i\mu_I, T, V) \quad \theta = \mu_I/T.$$

Without sign problem

- For imaginary chemical potential, there is no SIGN problem

One can do MCMC or Metropolis-Hastings MC

Then, we obtain  $Z_{\text{GC}}$  on LQCD

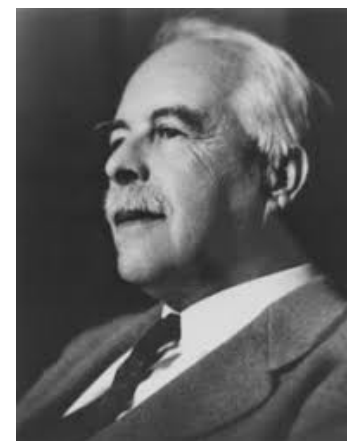
## 2. Application Lattice **Q**CD

- Fugacity expansion of grand canonical partition function

$$Z_{\text{GC}}(\mu, T, V) = \sum_{n=-\infty}^{\infty} Z_C(n, T, V) \xi^n$$

$$\xi (\equiv e^{\mu/T})$$

Fugacity



Gilbert Newton Lewis

- Obtain canonical partition function by Fourier transform

$$Z_C(n, T, V) = \int_{-\pi}^{\pi} \frac{d\theta}{2\pi} e^{-in\theta} Z_{\text{GC}}(\mu = i\mu_I, T, V) \quad \theta = \mu_I/T.$$

Without sign problem

- For imaginary chemical potential, there is no SIGN problem

One can do MCMC or Metropolis-Hastings MC

Then, we obtain  $Z_{\text{GC}}$  on LQCD

## 2. Application Lattice **QCD**

Canonical approach developed

### **Lattice QCD**

$$n_q(\mu_q = i\mu_{qI}, T, V)$$

Number density formulation  
V. Bornyakov et al., PRD95(2017)

$$\frac{n_q}{T^3} = \frac{1}{VT^2} \frac{\partial}{\partial \mu_q} \ln Z_{GC}$$

$$Z_{GC}(\mu_q = i\mu_{qI}, T, V)$$

Fourier transform

$$Z(n, T, V)$$

$$Z_{GC}(\mu_q, T, V) = \sum_{n=-\infty}^{\infty} Z(n, T, V) \xi^n \quad \xi = e^{\mu_q/T}$$

If we get  $Z_n$  for all  $n$ , we can search at **ANY** density!

## 2. Application Lattice **QCD**

Canonical approach developed

### **Lattice QCD**

$$n_q(\mu_q = i\mu_{qI}, T, V)$$

Number density formulation  
V. Bornyakov et al., PRD95(2017)

$$\frac{n_q}{T^3} = \frac{1}{VT^2} \frac{\partial}{\partial \mu_q} \ln Z_{GC}$$

$$Z_{GC}(\mu_q = i\mu_{qI}, T, V)$$

Fourier transform

$$Z(n, T, V)$$

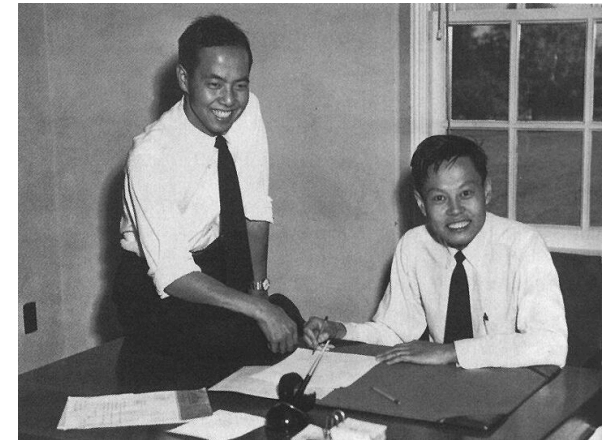
$$Z_{GC}(\mu_q, T, V) = \sum_{n=-N_{\max}}^{N_{\max}} Z(n, T, V) \xi^n \quad \xi = e^{\mu_q/T}$$

In numerical calculations,  $n$  is **finite**.

## 2. Application Lattice **QCD**

- Application of canonical method: **Lee-Yang zeros**
- Zeros of  $Z_{GC}$  so-called Lee-Yang Zeros (LYZ) contain a valuable information on the phase transitions of a system.

*T.D. Lee & C.N. Yang, Phys. Rev. 87, 404&410 (1952)*



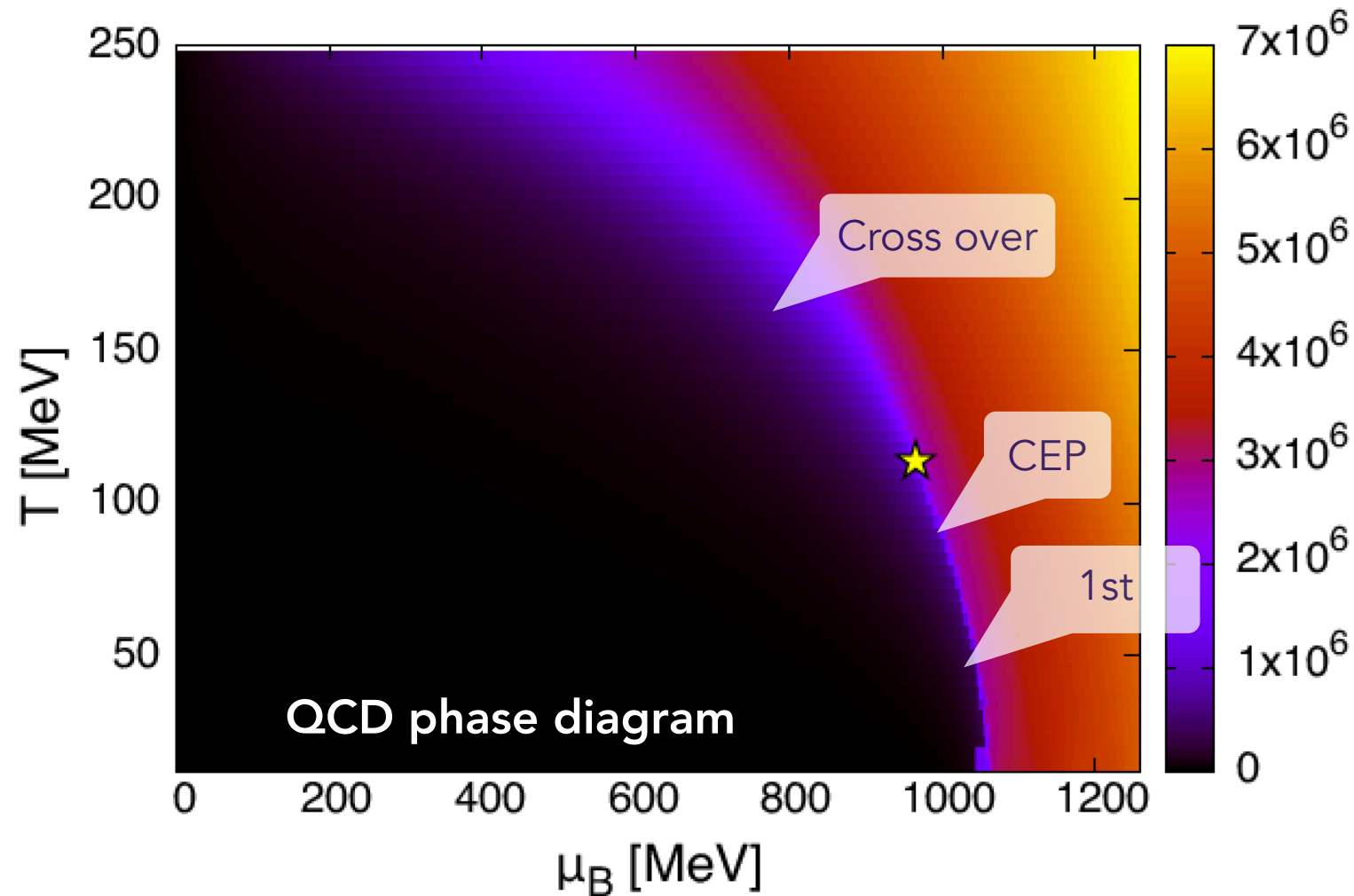
$$Z_{GC}(\mu_q, T, V) = \sum_{n=-N_{\max}}^{N_{\max}} Z_c(n, T, V) \xi^n = 0$$

Physically, at LYZ, critical-end point (CEP) appears!!



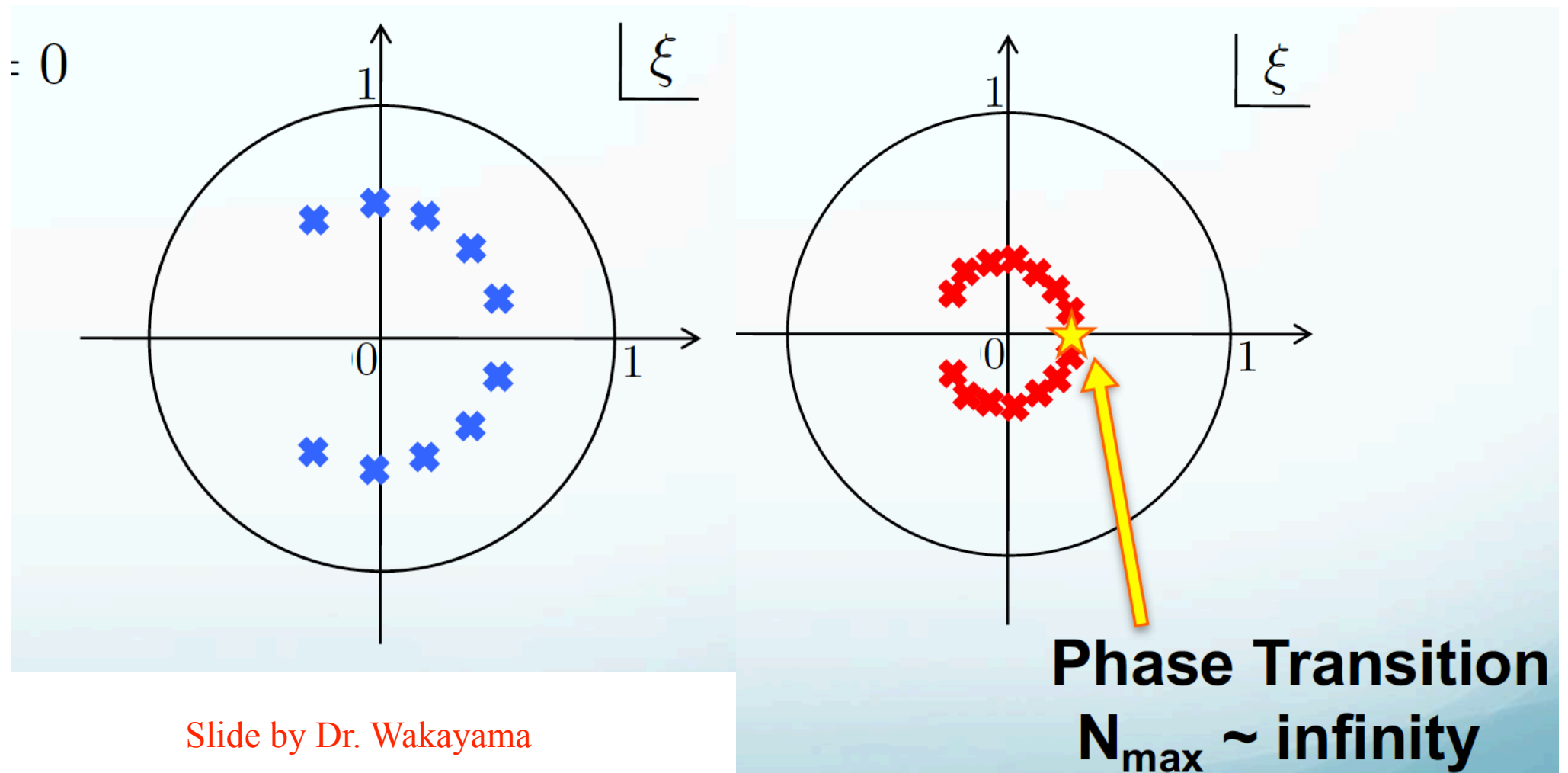
## 2. Application Lattice QCD

- What is critical-end point (CEP)??



## 2. Application Lattice QCD

- There are  $2N_{\max}$  LYZs in complex fugacity plane



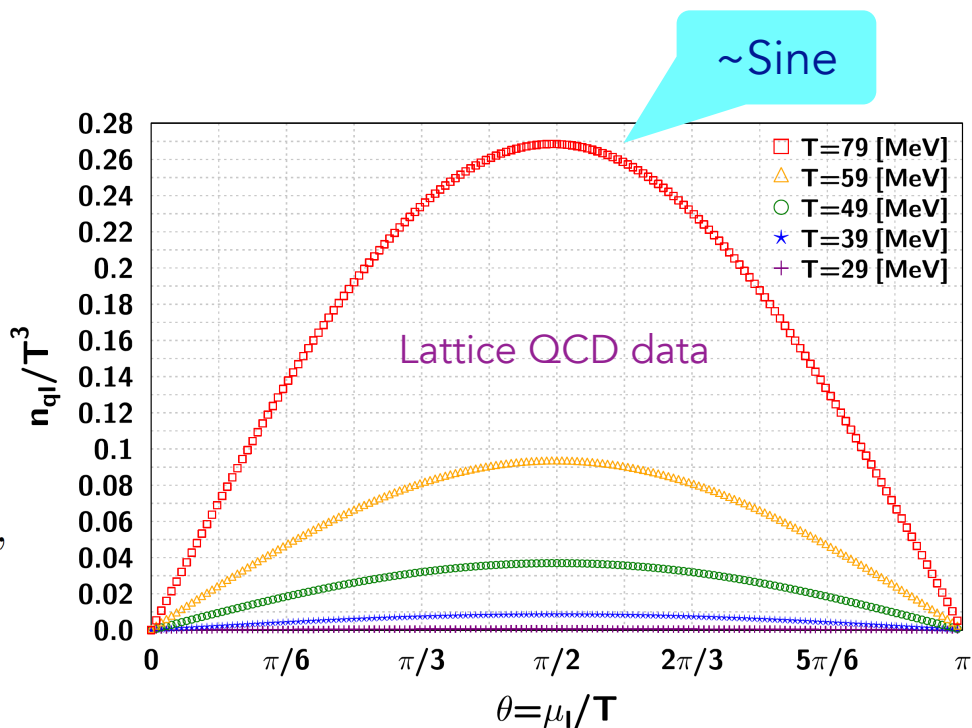
Slide by Dr. Wakayama

## 2. Application Lattice **QCD**

- First, we parameterize number density with sine function for more reliable numerical treatment in lattice QCD

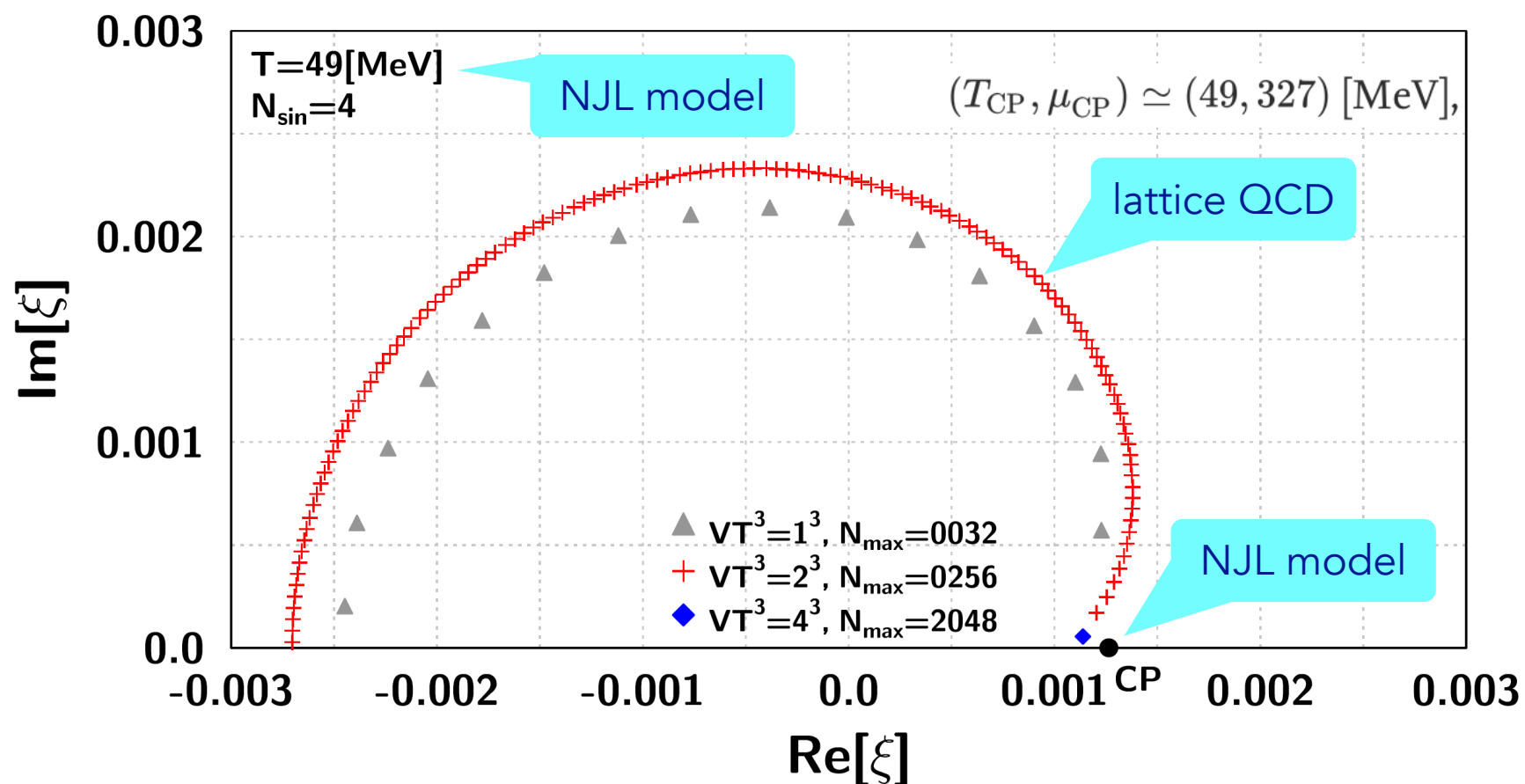
$$\frac{n_{qI}}{T^3}(\theta) = \sum_{k=1}^{N_{\text{sin}}} f_k \sin(k\theta)$$

$$\begin{aligned} Z_{\text{GC}}(i\mu_I, T, V) &= C \exp \left\{ -V \int_0^\theta d\theta' n_{qI}(\theta') \right\} \\ &= C \exp \left\{ VT^3 \sum_{k=1}^{N_{\text{sin}}} \frac{f_k}{k} \cos(k\theta) \right\}, \end{aligned}$$



## 2. Application Lattice QCD

- We observe LYZs cross the  $\text{Im}[\xi]=0$  line: CEP



## 2. Application Lattice **QCD**

- Application of canonical method: **QCD phase structure**
- This method is not full lattice QCD but mimics it closely
- Then, can we describe QCD phase diagram???: Yes!!!
- Before doing lattice QCD with canonical method,  
we test it in effective models: NJL and PNJL

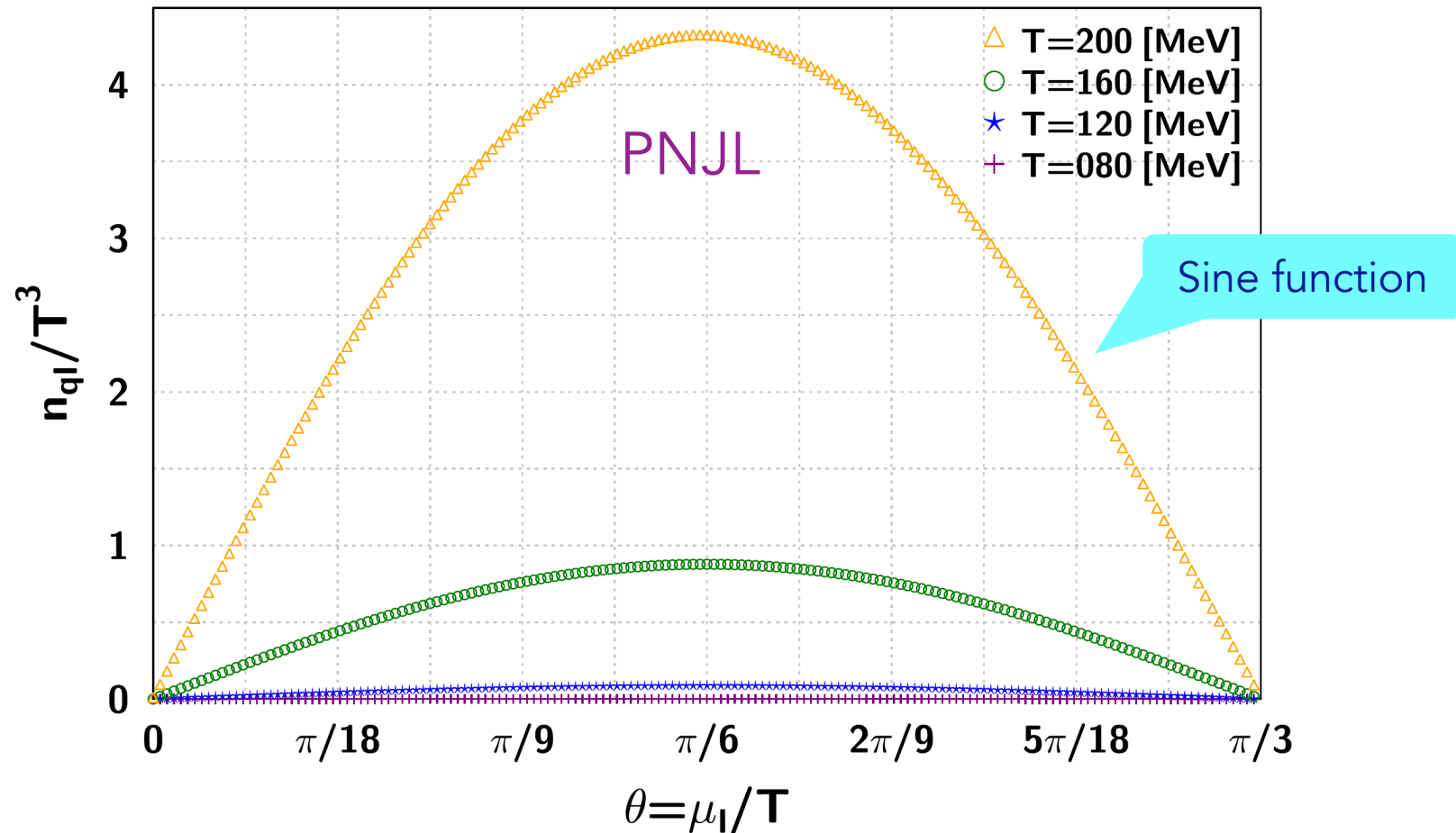
Thermodynamic potential of PNJL

$$\omega = \frac{1}{2G} (M - m_q)^2 - 2N_c N_f \int \frac{d^3p}{(2\pi)^3} E_p - 2N_f T \int \frac{d^3p}{(2\pi)^3} \left\{ \text{Tr}_c \ln \left[ 1 + L e^{-\frac{E_p - \mu}{T}} \right] \right. \\ \left. + \text{Tr}_c \ln \left[ 1 + L^\dagger e^{-\frac{E_p + \mu}{T}} \right] \right\} + T^4 \left[ -\frac{b_2(T)}{2} \ell \bar{\ell} - \frac{b_3}{6} (\ell^3 + \bar{\ell}^3) + \frac{b_4}{4} (\ell \bar{\ell})^2 \right]$$

Quark (pointing to  $2N_c N_f$ )  
 Quark-Gluon (pointing to  $\text{Tr}_c \ln$ )  
 Mass gap (pointing to  $(M - m_q)^2$ )  
 Quark-Gluon (pointing to  $\ell \bar{\ell}$ )  
 Gluon  $\sim Z(N_c)$  (pointing to  $b_3$ )

## Numerical results

- QCD phase structure from Polyakov-loop NJL model
- Quark number density from PNJL at  $i\mu$

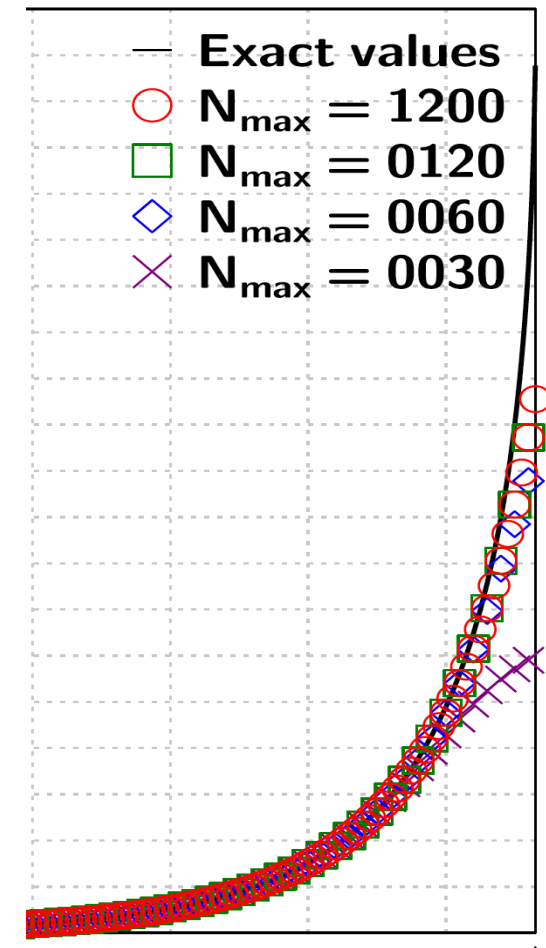


## 2. Application Lattice QCD

Wakayama, Nam, and Hosaka, PRD (2020)

- Application of canonical method: **QCD phase structure**
- As  $N_{\max}$  increases, results from canonical method reaches to exact value
- Nonetheless, canonical method does not coincide with exact one: limitation of the present method.
- Then, how do we quantify phase transition in this method?: Taking tolerance

$$\frac{n_B^{\text{PNJL}}}{n_B^{\text{Canonical}}} < 10\%$$

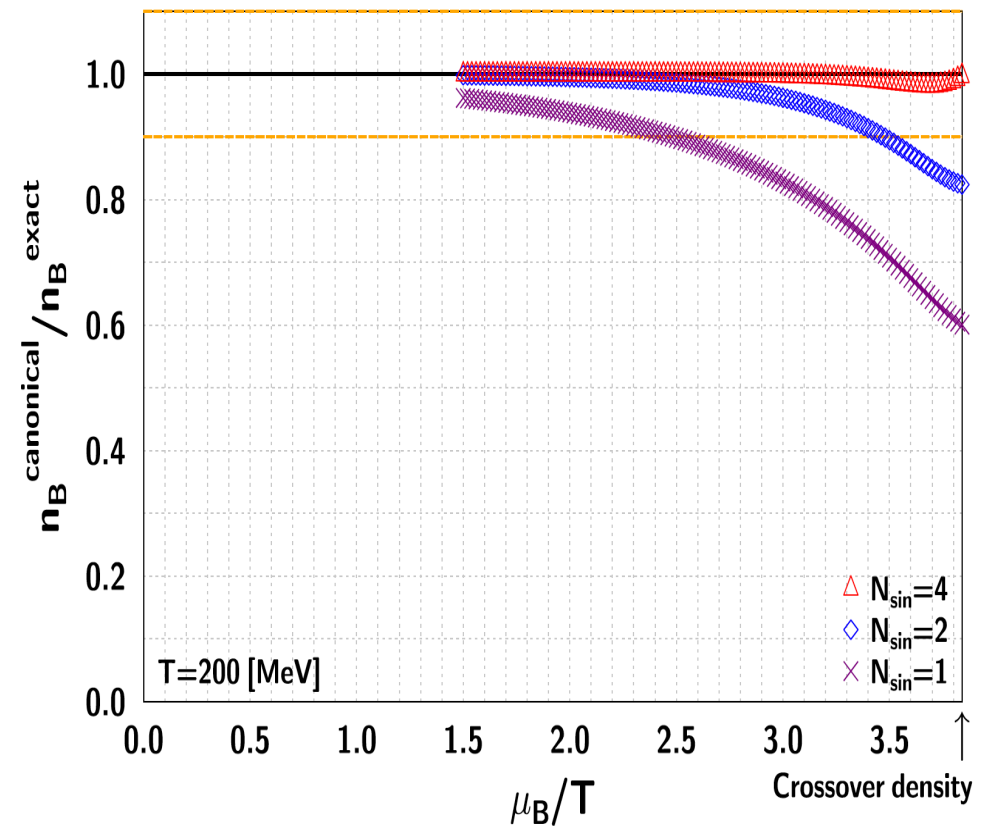
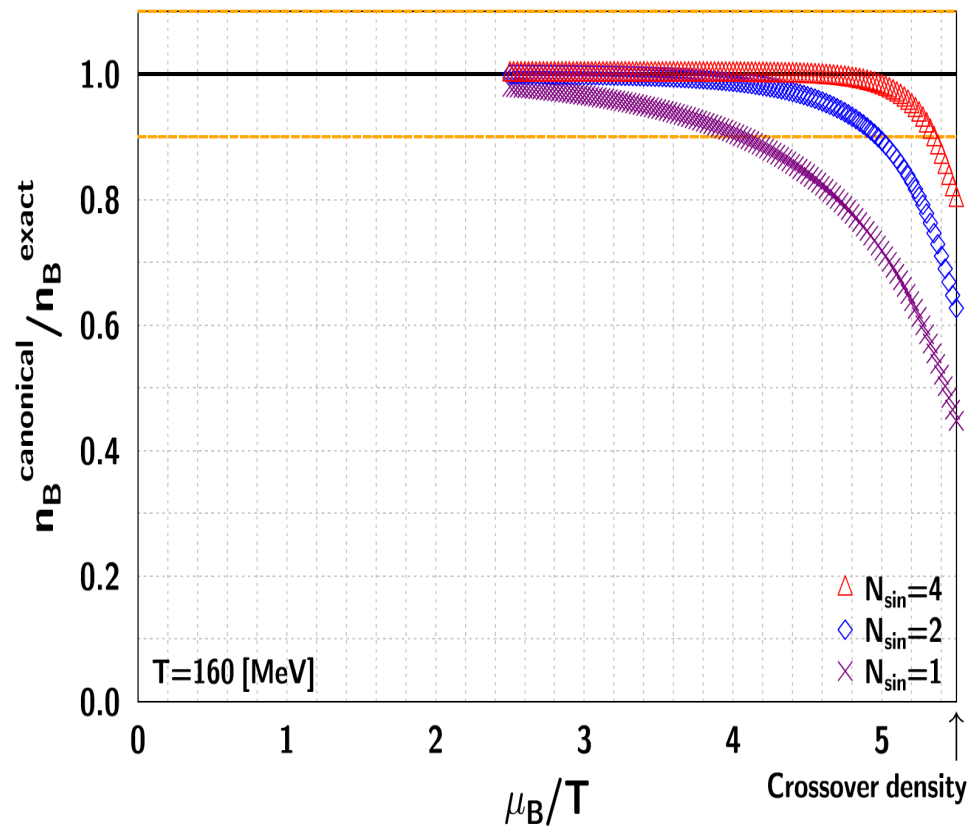




## 2. Application Lattice QCD

Wakayama, Nam, and Hosaka, PRD (2020)

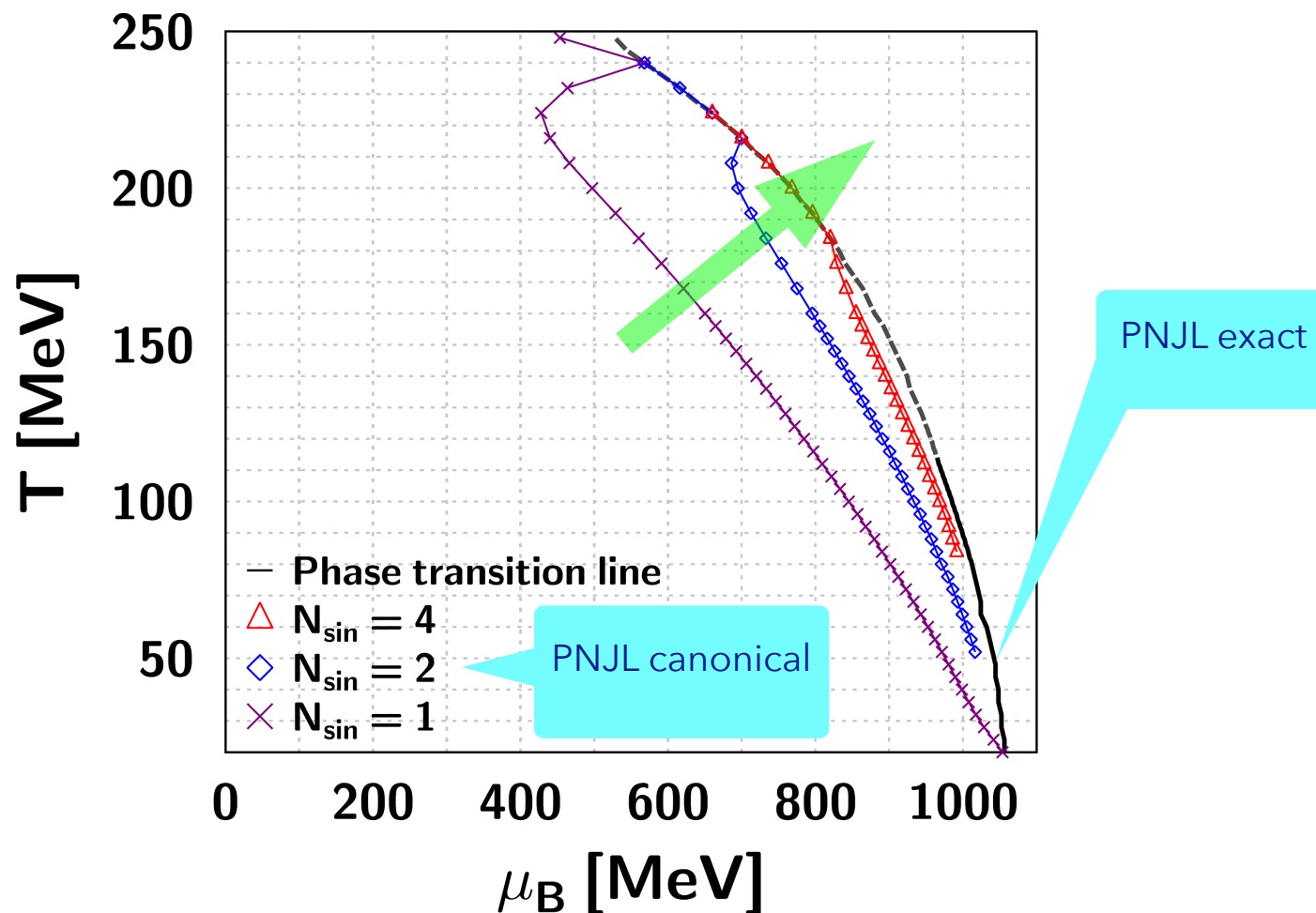
### Application of canonical method: QCD phase structure



$$\frac{n_B^{\text{PNJL}}}{n_B^{\text{Canonical}}} < 10\%$$

## Numerical results

- QCD phase structure from Polyakov-loop NJL model
- Chiral boundary is described very well!!!



## Summary

- QCD phase diagram investigated via **canonical method**
- To verify the method, we compare  $\text{PNJL}_{\text{exact}}$  and  $\text{PNJL}_{\text{canonical}}$
- Taking 10% tolerance between them
- Sufficiently small parameters can reproduce QCD phases: 4!
- At finite  $\mu$ , not on phase transition line, LQCD can be used!
- **QCD phase diagram via LQCD** (in progress)
- **LQCD data for nuclear matter** (in progress)
- **Interesting new phase found at  $i\mu$**  (in progress)



# Thank you for your attention!!

Supported by the National Research Foundation of Korea (NRF) grants:

**No.2018R1A5A1025563 and No.2019R1A2C1005697**

