

# Pion off-shell electromagnetic form factors with the light-front approach models

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# Contents at a glance

- Motivations
- Off Shell Electromagnetic Form Factors
- Extraction of pion form factors through experimental cross section
- Results
- Conclusions

- **Hadronic form factors:** important source information for hadrons structure
- **Off-shell effects:** electroproduction  $\implies {}^1H(e, e'\pi^+)n$  (**Fig. 1**)
- **Cross section:**  $\sigma_L, \sigma_T, \sigma_{LT}$  e  $\sigma_{TT}$
- **Extraction of the pion's electromagnetic form factors:**  $F_1(Q^2, t)$  e  $F_2(Q^2, t)$

## Ref:

Ho-Meoyng Choi et al. "Pion off-shell electromagnetic form factors: Data extraction and model analysis."  
In: *Physical Review D* 100.11 (2019), p. 116020

# Motivations

## Electroproduction experiment

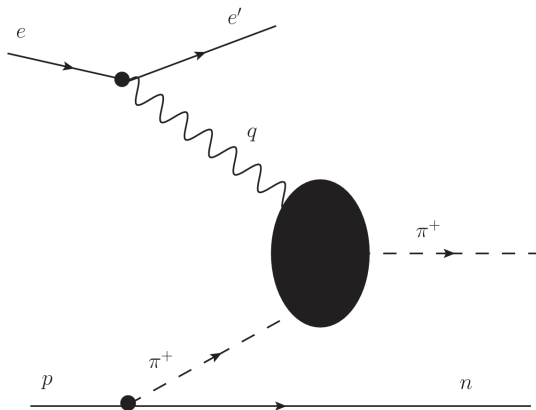


Figure 1:  $e p \rightarrow e' \pi^+ n$  scattering

# Off Shell Electromagnetic Form Factors

- Off-shell case  $\implies$  **Two Electromagnetic Form Factors:**  $F_1$  and  $F_2$
- Most simple structure:

$$\langle p' | \mathcal{O} | p \rangle = (p' + p)^\mu F_1(Q^2, t) + (p' - p)^\mu F_2(Q^2, t)$$

$\implies$   $\mathcal{O}$  is the electromagnetic current operator

- **Vertex:**

$$\Gamma_\mu = [(p + p')_\mu G_1 + (p' - p)_\mu G_2]$$

and,

$$q^\mu \Gamma_\mu = [\Delta^{-1}(p') - \Delta(p)^{-1}]$$

where  $\Delta(p) = \frac{1}{p^2 - m_\pi^2 - \Pi(p) + i\epsilon}$

# Off Shell Electromagnetic Form Factors

From the Ward-Takahashi identity, we have to:

$$(p'^2 - p^2)G_1(q^2, p^2, p'^2) + q^2 G_2(q^2, p^2, p'^2) = \Delta^{-1}(p') - \Delta^{-1}(p)$$

that, with the normalization  $G_1(q^2 = 0) = 1$ :

$$G_2(q^2, p^2, p'^2) = \frac{(p'^2 - p^2)[G_1(0, p^2, p'^2) - G_1(q^2, p^2, p'^2)]}{q^2}$$

In the case, where  $p^2 = t$  and  $p'^2 = m_\pi^2$ :

$$F_2(Q^2, t) = \frac{t - m_\pi^2}{Q^2} [F_1(0, t) - F_1(Q^2, t)]$$

when  $F_i(Q^2, t) \equiv G_i(q^2, t, m_\pi^2)$  ( $i = 1, 2$ ) and  $Q^2 = -q^2$

# Off Shell Electromagnetic Form Factors

Solving, we obtain:

$$\Gamma_\mu = (p' + p)_\mu F_1(Q^2, t) + q_\mu \frac{(t - m_\pi^2)}{Q^2} [F_1(0, t) - F_1(Q^2, t)]$$

And defining


$$g(Q^2, t) \equiv \frac{F_2(Q^2, t)}{t - m_\pi^2}$$

we obtain:

$$F_1(Q^2, t) - F_1(0, t) + Q^2 g(Q^2, t) = 0$$

This Equation allows us to extract the form factors from the off-shell pion. The form factors  $F_2(Q^2, t)$  and  $g(Q^2, t)$  disappear in the on-shell limit  $t = m_\pi^2$ .

**Ref:**

Choi et al., "Pion off-shell electromagnetic form factors: Data extraction and model analysis" 

# Extraction of pion form factors through experimental cross section

The electromagnetic form factor of the pion is given as a function of the cross section:

$$F_\pi^2 = \frac{N}{4\hbar c} \frac{1}{(eG_{\pi NN}(t))^2} \frac{(t - m_\pi^2)^2}{-Q^2 t} \frac{d\sigma_L}{dt}$$

The cross section is given as:

$$\boxed{\frac{d\sigma_L}{dt} = \frac{16\pi}{137N} \frac{-tQ^2}{(t - m_\pi^2)} G_{\pi NN}^2(t) F_\pi^2(Q^2, t) \left(\frac{197.3}{10}\right)^2}$$

We also have to:

$$\frac{d\sigma_L}{dt} = \frac{16\pi}{137N} \frac{-tQ^2}{(t - m_\pi^2)} [H]^2 \left(\frac{197.3}{10}\right)^2$$

where  $H \equiv F_\pi \cdot G_{\pi NN}$



# Extraction of pion form factors through experimental cross section

$H$  can be extracted from the cross section:

$$H^2 = \frac{137N}{16\pi} \frac{(t - m_\pi^2)^2}{-tQ^2} \frac{d\sigma_L}{dt} \left( \frac{10}{197.3} \right)^2$$

Where  $G_{\pi NN}(t)$  is given by:

$$G_{\pi NN}(t) = G_{\pi NN}(m_\pi^2) \left( \frac{\Lambda_\pi^2 - m_\pi^2}{\Lambda_\pi^2 - t} \right)^n$$

with  $G_{\pi NN}(m_\pi^2) = 13.4$ ,  $\Lambda_\pi = 0.80$  GeV and  $n = 1$ .

Therefore:

$$F_\pi^2 = \frac{137N}{16\pi} \frac{1}{G_{\pi NN}(t)} \frac{(t - m_\pi^2)^2}{-Q^2 t} \frac{d\sigma_L}{dt} \left( \frac{10}{197.3} \right)^2$$

# Constituent Quark Models in Light Front

The elements of the electromagnetic current matrix:

$$J^\mu = -i2e \frac{m^2}{f_\pi^2} N_c \int \frac{d^4k}{(2\pi)^4} \text{Tr} [S(k)\gamma^5 S(k-p')\gamma^\mu S(k-p)\gamma^5] \Gamma(k, p')\Gamma(k, p)$$

where  $S(p) = \frac{1}{\not{p} - m + i\epsilon}$  and  $N_c = 3$ .

For the non-symmetric vertex model:

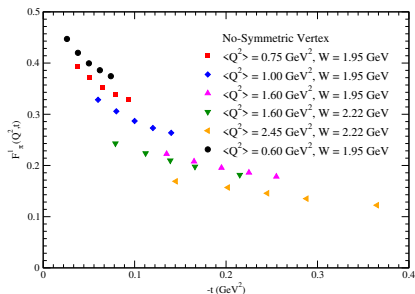
$$\Gamma^{NSV}(k, p) = \frac{N}{(p-k)^2 - m_R^2 + i\epsilon}$$

For the symmetric vertex model:

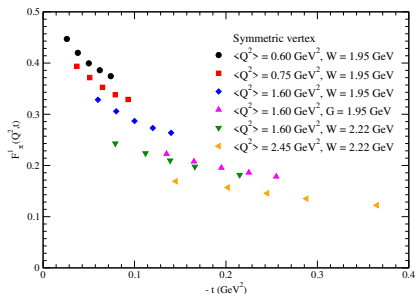
$$\Gamma^{SV}(k, p) = \frac{C}{k^2 - m_R^2 + i\epsilon} + \frac{C}{(p-k)^2 - m_R^2 + i\epsilon}$$

JPBC de Melo et al. "Pair term in the electromagnetic current within the Front-Form dynamics: spin-0 case." In: *Nuclear Physics A* 707.3-4 (2002), pp. 399-424

JPBC de Melo. "Covariância na frente de luz." PhD thesis. 1998



(a)



(b)

Figure 2: (a) Form factor  $F_{\pi}^1(Q^2, t)$  for the non-symmetric model; (b) Form factor  $F_{\pi}^1(Q^2, t)$  for the symmetric model

## Ref:

HP Blok et al. "Charged pion form factor between  $Q^2 = 0.60$  and  $2.45 \text{ GeV}^2$ . I. Measurements of the cross section for the  $1 \text{ H} (e, e^{\pi+}) n$  reaction." In: *Physical Review C* 78.4 (2008), p. 045202

GM Huber et al. "Charged pion form factor between  $Q^2 = 0.60$  and  $2.45 \text{ GeV}^2$ . II. Determination of, and results for, the pion form factor." In: *Physical Review C* 78.4 (2008), p. 045203

Choi et al., "Pion off-shell electromagnetic form factors: Data extraction and model analysis" 11/19

# First results

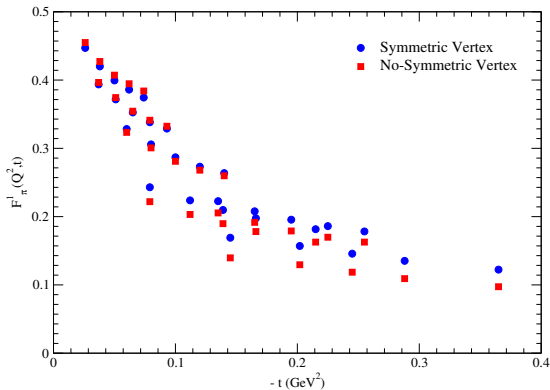
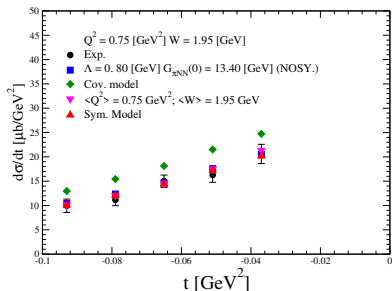
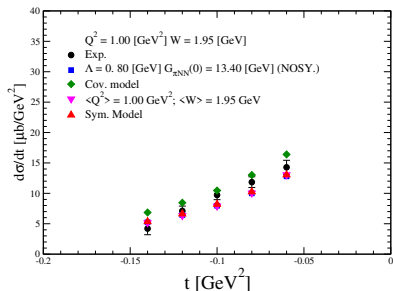


Figure 3: Comparison between the form factors for the two models **symmetric** and **non-symmetric**

# First results



(a)



(b)

Figure 4: Cross section

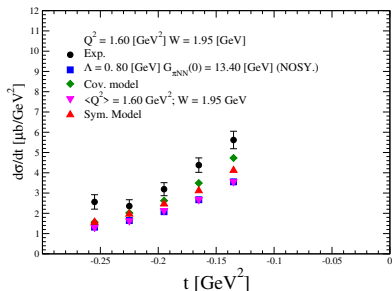
## Ref:

Blok et al., “Charged pion form factor between  $Q^2 = 0.60$  and  $2.45$  GeV<sup>2</sup>. I. Measurements of the cross section for the  $1\text{H}(e, e'\pi^+)n$  reaction”

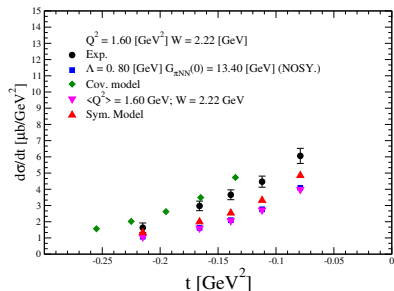
Huber et al., “Charged pion form factor between  $Q^2 = 0.60$  and  $2.45$  GeV<sup>2</sup>. II. Determination of, and results for, the pion form factor”

Choi et al., “Pion off-shell electromagnetic form factors: Data extraction and model analysis”

# First results



(a)



(b)

Figure 5: Cross section

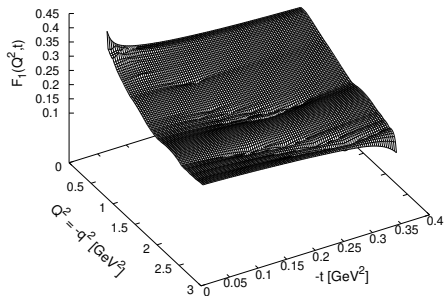
## Ref:

Blok et al., “Charged pion form factor between  $Q^2 = 0.60$  and  $2.45 \text{ GeV}^2$ . I. Measurements of the cross section for the  $1 \text{ H} (e, e' \pi^+) n$  reaction”

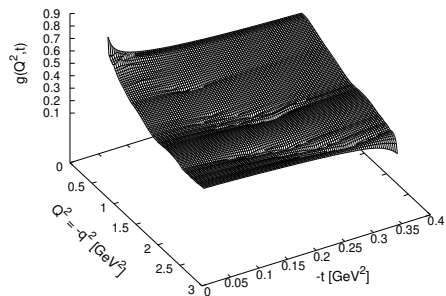
Huber et al., “Charged pion form factor between  $Q^2 = 0.60$  and  $2.45 \text{ GeV}^2$ . II. Determination of, and results for, the pion form factor”

Choi et al., “Pion off-shell electromagnetic form factors: Data extraction and model analysis”

# First results

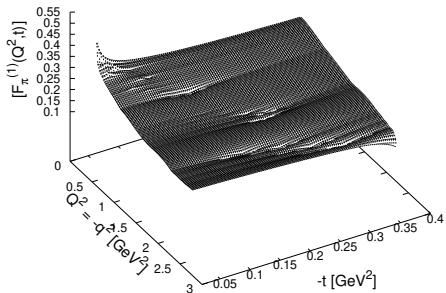


(a)

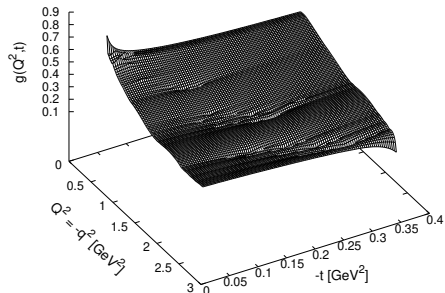


(b)

Figure 6: (a) Form factor  $F^1(Q^2, t)$  and (b)  $g(Q^2, t)$  for the **non-symmetric model**



(a)



(b)

Figure 7: (a) Form factor  $F^1(Q^2, t)$  and (b)  $g(Q^2, t)$  for the [symmetric model](#)



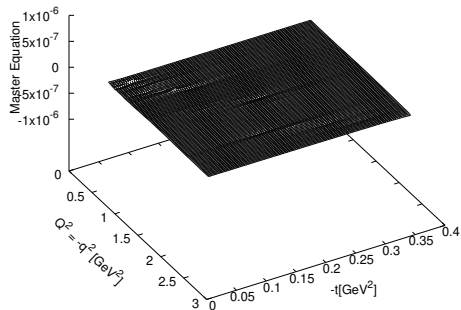


Figure 8: The master equation - [symmetric model](#)

# Conclusions

- We reviewed some of the calculations for (Choi et. al., 2019);
- We verified that our calculations indicate that they do not depend on models;
- The graphs show that our result agrees with the references;
- We will continue our study for other models and parameters;
- The present approach will also be applied to other pseudoscalar mesons, for example: the kaon.

Choi et al., “Pion off-shell electromagnetic form factors: Data extraction and model analysis”

Blok et al., “Charged pion form factor between  $Q^2 = 0.60$  and  $2.45 \text{ GeV}^2$ . I. Measurements of the cross section for the  $1 \text{ H} (e, e' \pi^+) n$  reaction”

Huber et al., “Charged pion form factor between  $Q^2 = 0.60$  and  $2.45 \text{ GeV}^2$ . II. Determination of, and results for, the pion form factor”

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