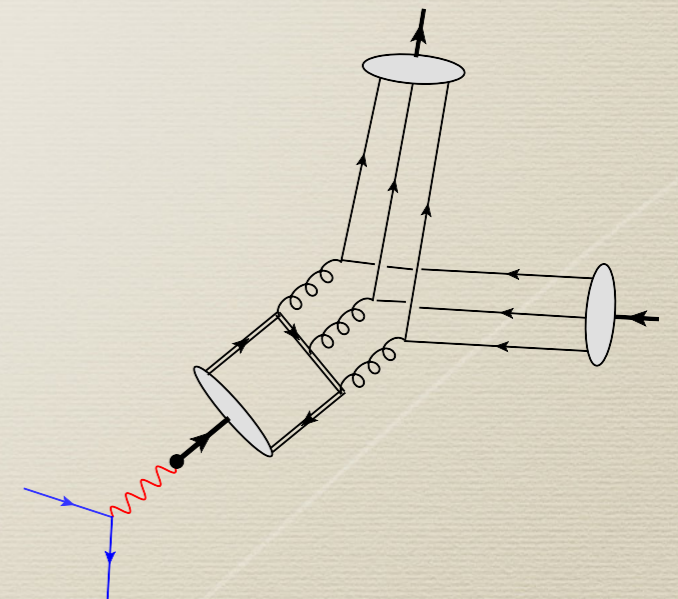


A study of J/ψ baryon decays beyond the leading twist accuracy

Nikolay Kivel



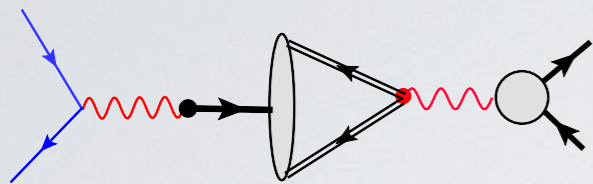
Eur.Phys.J. A 56(2020),
arXiv:2109.05847, to appear in Eur.Phys.J. A,
and new results (in preparation)



Light Cone 2021: Physics of hadrons on the light front,
November 29–December 4, 2021

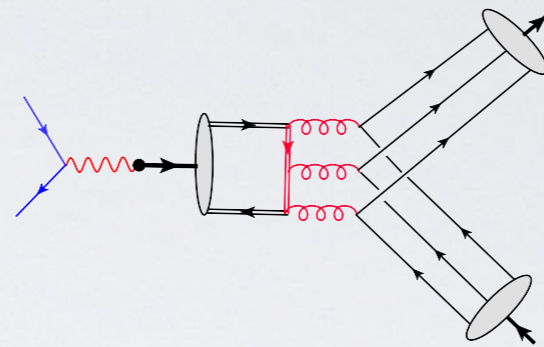
$J/\psi \rightarrow B\bar{B}$ decays

e.m. FFs



G_M, G_E

3-gluon annihilation

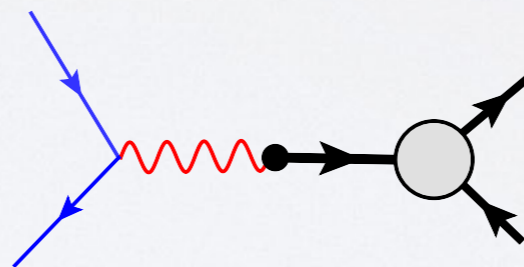


G_M^g, G_E^g

$$\text{Br}[J/\psi \rightarrow B\bar{B}] = \frac{M_\psi \beta}{12\pi\Gamma_{tot}} |G_M^g + G_M|^2 \left(1 + \frac{2m_B^2}{M_\psi^2} \gamma_B^2 \right)$$

$$\gamma_B = \left| \frac{G_E^g + G_E}{G_M^g + G_M} \right|$$

$$\frac{dN_B}{d\cos\theta} = \mathcal{N}(1 + \alpha_B \cos^2\theta)$$



$$\alpha_B = \frac{1 - 4m_B^2 \gamma_B^2 / M_\psi^2}{1 + 4m_B^2 \gamma_B^2 / M_\psi^2}$$

$$m_Q \rightarrow \infty \quad \alpha = 1$$

Brodsky, Lepage 1981

$$\frac{4m_N^2}{M_\psi^2} \simeq 0.37$$

$$\frac{4m_\Upsilon^2}{M_\psi^2} \simeq 0.04$$

J/ψ baryonic decays

Data: BESIII/2008/2012/2016/2017/2019/2020/ & PDG

| B | $\text{Br}[J/\psi \rightarrow B\bar{B}] \times 10^3$ | α_B | $\gamma_B = \frac{ G_E^g + G_E }{ G_M^g + G_M }$ |
|------------|--|------------|--|
| p | 2.12(3) | 0.59(1) | 0.83 |
| n | 2.1(2) | 0.50(4) | 0.95 |
| Λ | 1.89(9) | 0.47(3) | 0.83 |
| Σ^0 | 1.17(3) | -0.45(2) | 2.11 ! |
| Σ^+ | 1.5(3) | -0.51(2) | 2.27 ! |
| Ξ^+ | 0.97(8) | 0.58(4) | 0.61 |
| Ξ^0 | 1.17(3) | 0.66(3) | 0.53 |

Theoretical estimates

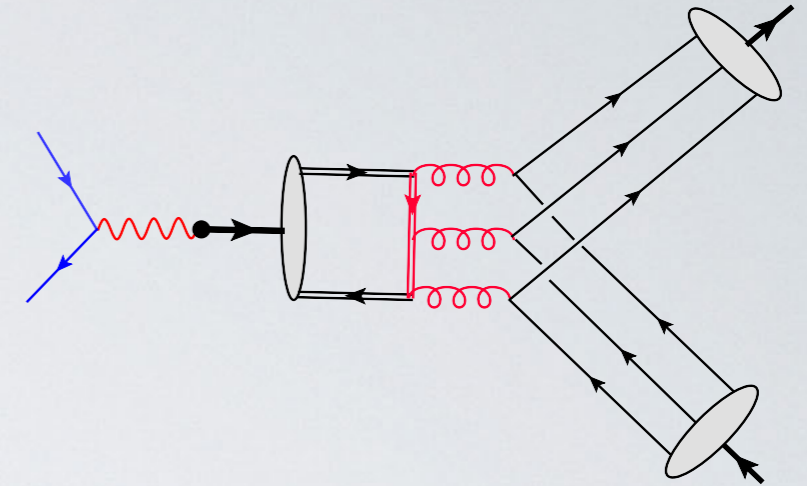
Largest amount of papers are dedicated to Br's and constrains of baryon DAs

Carimalo, 1987 $\alpha_N = 0.70$ constituent quarks, non-relativistic nucleon WF

Murgia, Melis, 1995 $\alpha_N = 0.561 - 0.963$ constituent quark mass $m_i = x_i m_N$
& the LT factorisation formula

J/ψ baryonic decays

Theory: NRQCD + QCD collinear factorisation



$$G_M^g \simeq R_{10}(0) \left\{ \underbrace{\phi^{\text{tw}3} * T_{33} * \phi^{\text{tw}3}}_{\text{LP}} + \underbrace{\phi^{\text{tw}3} * T_{35} * \phi^{\text{tw}5} + \phi^{\text{tw}4} * T_{44} * \phi^{\text{tw}4}}_{\text{NLP} \sim \Lambda^2/m_c^2} \right\}$$

$$G_E^g \simeq R_{10}(0) \phi^{\text{tw}3} * T_{34} * \phi^{\text{tw}4}$$

$G_M^g(\text{LP})$ Brodsky, Lepage 1981/ Chernyak et al, 1984

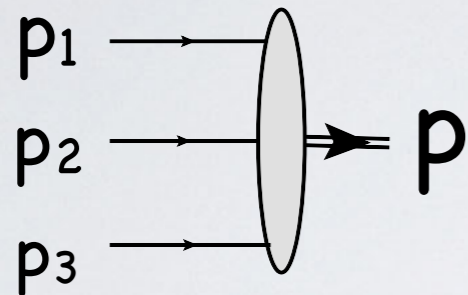
$G_M^g(\text{NLP}), G_E^g$ NK, 2020/2021 & 2021 (in preparation)

$$\gamma_B^g = \left| \frac{G_E^g}{G_M^g} \right| \quad \text{dominates \& sensitive to a baryon structure !}$$

Baryon Light-Cone Distribution Amplitudes

Long distance physics associated with the nucleon WFs

Light-cone Distribution Amplitude $\phi(x_i)$ describes how the long. momentum is shared between the constituents



$$\phi^{\text{tw}3} \equiv \phi_3(x_i, \mu) \sim \int_{|p_{i\perp}| \lesssim \mu} dp_{i\perp} \Psi(x_i, p_{i\perp})$$

$p_i = x_i p + p_{\perp i}$ $p_{\perp i} \sim \Lambda$ Defined as a light-cone matrix element

$$\langle 0 | \varepsilon^{abc} q^a(z_1) q^b(z_2) q^c(z_3) | B(p) \rangle_{z_i^2=0} \sim \text{FT}[\phi_3(x_i)]$$

$$\phi_3(x_i, \mu) = f_B(\mu) x_1 x_2 x_3 (1 + \phi_{10}^{(3)}(\mu) P_{10}(x_i) + \phi_{11}^{(3)}(\mu) P_{11}(x_i) + \dots)$$

$P_{kn}(x_i)$ homogeneous orthogonal polynomials of degree k

$f_B, \phi_{ij}^{(3)}$ non-perturbative moments, m.e.'s of local operators

Baryon Light-Cone DAs: 3-quark operators only

| $B = N, \Lambda, \Sigma, \Xi$ | Γ chiral even | Γ chiral odd | $\langle 0 [q(z_1) C \Gamma q(z_2)] q(z_3) B(p) \rangle_{z_i^2=0}$ | |
|-------------------------------|----------------------------|---------------------------|--|--------------------------|
| twist 3 | $\Phi_{3\pm}^B(x_i)$ | $\Pi_3^B(x_i)$ | | |
| twist 4 | $\Phi_{4\pm}^B(x_i)$ | $\Xi_{4\pm}^B(x_i)$ | $\Pi_4^B(x_i)$ | $\Upsilon_4^B(x_i)$ |
| twist 5 | $\Phi_{5\pm}^B(x_i)$ | $\Xi_{5\pm}^B(x_i)$ | $\Pi_{5\pm}^B(x_i)$ | $\Upsilon_{5\pm}^B(x_i)$ |

Definitions/ SU(3) relations/ Conformal expansions/
Evolution/ Wandzura-Wilzeck decompositions are known

Braun et al, 2000, 2008, 2013

Manashov, Anikin 2013

Shäfer, Wein 2015

Anikin 2015

Baryon Light-Cone DAs

Models: first few term of the conformal expansion:

Twist-3

$$\phi_{3+}^B(x_i) = f_B x_1 x_2 x_3 (1 + \phi_{11}^B 21(x_1 - 2x_2 + x_3))$$

$$\phi_{3-}^B(x_i) = f_B x_1 x_2 x_3 \phi_{10}^B 21(x_1 - x_3)$$

$$\phi^{\text{tw}4}(x_i, \mu) = \mathcal{P}[x_i; \phi_{ij}^{(3)}] + \bar{\phi}^{\text{tw}4}(x_i, \mu)$$

Twist-4

$$\bar{\Phi}_{4+}^B(x_i) = \lambda_1^B 24x_1 x_2 (- \eta_{11}^B (2x_1 - x_2 - 2x_3))$$

$$\bar{\Phi}_{4-}^B(x_i) = \lambda_1^B 24x_1 x_2 (1 + \eta_{10}^B (4 - 10x_2))$$

quark-gluon moments contribute starting from

η_{2i}^B and therefore can be neglected

Twist-5

$$\phi^{\text{tw}5}(x_i, \mu) = \mathcal{P}[x_i; \phi_{ij}^{(3)}, \phi_{ij}^{(4)}] + \bar{\phi}^{\text{tw}5}(x_i, \mu)$$

Baryon Light-Cone DAs

What do we know about the moments?

QCD Sum Rules: twist-3,4 moments

Chernyak et al, 1989

Braun et al, 2000

QCD Light-Cone SR: twist-3,4 moments
(data fit, ABOI model for N)

Anikin et al, 2013

Lattice: twist-3,4 moments

Bali et al, 2019

relations in the $SU(3)_f$ limit + data fit

The moments used in calculations

$$\mu^2 = 1.5 \text{ GeV}^2$$

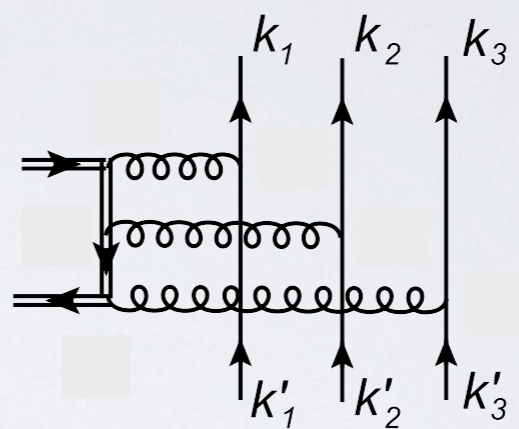
| B | $f_B, \text{ GeV}^2$ | ϕ_{10} | ϕ_{11} | ϕ_{20} | ϕ_{21} | ϕ_{22} | $f_{\perp}^B, \text{ GeV}^2$ | π_{10}^B | π_{11}^B |
|-----------|-----------------------|-------------|-------------|-------------|-------------|-------------|------------------------------|--------------|--------------|
| N | 4.94×10^{-3} | 0.051 | 0.051 | 0.078 | -0.028 | 0.179 | — | — | — |
| Λ | 5.60×10^{-3} | 0.120 | 0.048 | 0 | 0 | 0 | — | 0.042 | — |
| Σ | 4.71×10^{-3} | 0.021 | 0.047 | 0 | 0 | 0 | 4.61×10^{-3} | — | -0.022 |
| Ξ | 4.94×10^{-3} | 0.078 | -0.003 | 0 | 0 | 0 | 4.83×10^{-3} | — | 0.093 |

| B | $\lambda_1^B, \text{ GeV}^2$ | η_{10}^B | η_{11}^B | $\lambda_{\perp}^B, \text{ GeV}^2$ | ζ_{10}^B | ζ_{11}^B | $\lambda_2, \text{ GeV}^2$ | ξ_{10} |
|-----------|------------------------------|---------------|---------------|------------------------------------|----------------|----------------|----------------------------|------------|
| N | -28×10^{-3} | -0.040 | 0.163 | — | — | 0.127 | 53×10^{-3} | -0.27 |
| Λ | -39×10^{-3} | -0.040 | 0.147 | -49×10^{-3} | -0.040 | — | 92×10^{-3} | — |
| Σ | -43×10^{-3} | -0.051 | 0.262 | — | — | 0.262 | 93×10^{-3} | -0.169 |
| Ξ | -46×10^{-3} | -0.040 | 0.147 | — | — | 0.147 | 92.6×10^{-3} | -0.27 |

Hard kernels

Power corrections:

contributions with derivatives with respect to \perp components of quark momenta:

$$\sum_{i,j=1,2} \left\{ A(x_i) + B(x_i) \frac{\partial}{\partial k_{\perp i}} + C(x_i) \frac{\partial}{\partial k_{\perp i}} \frac{\partial}{\partial k_{\perp j}} + \dots \right\}$$


The diagram shows a hard kernel with three quark lines (represented by straight lines) and a gluon loop (represented by a wavy line). The quark lines are labeled with momenta k_1, k_2, k_3 at the top and k'_1, k'_2, k'_3 at the bottom. The gluon loop is connected to the quark lines. The diagram is part of a larger expression involving derivatives with respect to the perpendicular components of the quark momenta.

$$k_i = x_i k + k_{\perp i}$$
$$k'_i = y_i k' + k'_{\perp i}$$

this usually produces the singular terms $\sim 1/x_i^n$ and often yields IR-divergencies, that usually indicates about a violation of the collinear factorization

however a soft gluon emission is suppressed by power of the small heavy quark velocity v and therefore such region is subleading, hence factorization holds and for the power corrections associated with higher twist baryon DAs

a separation of the 3-quark and 3-quark+gluon contributions for the twist-5 matrix elements, is not simple and, probably, not unique more details in the paper

NK, 2021 (in preparation)

Amplitudes analytical expressions are not complicated but somewhat lengthy,

Nucleon **amplitude**:

$$G_E^g = \frac{f_\psi}{m_c^6} (\pi\alpha_s)^3 \frac{20}{81} \int Dy_i \frac{1}{y_1 y_2 y_3} \int Dx_i \frac{1}{x_1 x_2 x_3} \frac{1}{D_1 D_2 D_3} \{ \dots \}$$

$$D_i = x_i(1 - y_i) + (1 - x_i)y_i$$

tw.4

tw.3

$$\begin{aligned} \{ \dots \} = & (\mathcal{A}_1 - \mathcal{V}_1) (x_{123}) \phi_3(x_{213}) x_1 (x_2 (y_2 - y_3) - \bar{y}_1 y_2) \\ & + (\mathcal{A}_1 + \mathcal{V}_1) (x_{123}) \phi_3(y_{123}) x_2 (x_2 - y_2) (y_1 - y_3) \\ & + (\mathcal{T}_{21} - \mathcal{T}_{41}) (x_{123}) (\phi_3(y_{132}) + \phi_3(y_{231})) x_3 (x_2 (y_1 - y_2) + y_2 \bar{y}_3) \end{aligned}$$

$\mathcal{V}_1, \mathcal{A}_1, \mathcal{T}_{21} - \mathcal{T}_{41}$ **convenient auxiliary tw.4 DAs** $\mathcal{V}_1(x_i) \equiv \mathcal{V}_1[\Phi_{4\pm}(x_i)], \dots$

$$\phi_3(y_i) \sim (y_1 y_2 y_3) \times \text{polynomials in } y_i$$

$$\{ \mathcal{V}_1, \mathcal{A}_1, \mathcal{T}_{21} - \mathcal{T}_{41} \} \sim (x_1 x_2 x_3) \times \text{polynomials in } x_i$$

$$\Rightarrow \int Dy_i \frac{1}{y_1 y_2 y_3} \int Dx_i \frac{1}{x_1 x_2 x_3} \frac{1}{D_1 D_2 D_3} \{ \dots \} \quad \text{well defined}$$

Phenomenological analysis

Power corrections

$$G_M^g \simeq R_{10}(0) \left\{ \underbrace{\phi^{\text{tw}3} * T_{33} * \phi^{\text{tw}3}}_{\text{LP}} + \underbrace{\phi^{\text{tw}3} * T_{35} * \phi^{\text{tw}5} + \phi^{\text{tw}4} * T_{44} * \phi^{\text{tw}4}}_{\text{NLP} \sim \Lambda^2/m_c^2} \right\}$$

| | N | Λ | Σ | Ξ |
|---|------|-----------|----------|-------|
| $\frac{G_M^g(\text{NLP})}{G_M^g(\text{LP})} \times 100\%$ | 8% | 4.5% | 13.3% | 35.6% |
| $\frac{m_B^2}{M_\psi^2}$ | 9.2% | 13% | 14.8% | 17.5% |

Not very large!

Phenomenological analysis

$$\text{Br}[J/\psi \rightarrow B\bar{B}] = \frac{M_\psi \beta}{12\pi\Gamma_{tot}} |G_M^g + G_M|^2 \left(1 + \frac{2m_B^2}{M_\psi^2} \gamma_B^2 \right)$$

$$\gamma_B = \left| \frac{G_E^g + G_E}{G_M^g + G_M} \right|$$

e.m. FFs estimate

$$|G_M| \approx |G_{eff}|$$

$$e^+e^- \rightarrow B\bar{B}$$

BABAR 2007

BESIII 2017, 2019, 2020, 2021

$$G_M = |G_M| e^{i\phi_M}$$

$\cos[\phi_M]$ is estimated from Br data

$$G_E = |G_E| e^{i\phi_E}$$

$0 < |G_E^B| \cos \varphi_E \leq 1.5 G_{\text{eff}}^B$ assumption

Phenomenological analysis

normalization $\mu^2 = 1.5 \text{ GeV}^2$ $\alpha_s(\mu^2) = 0.35$

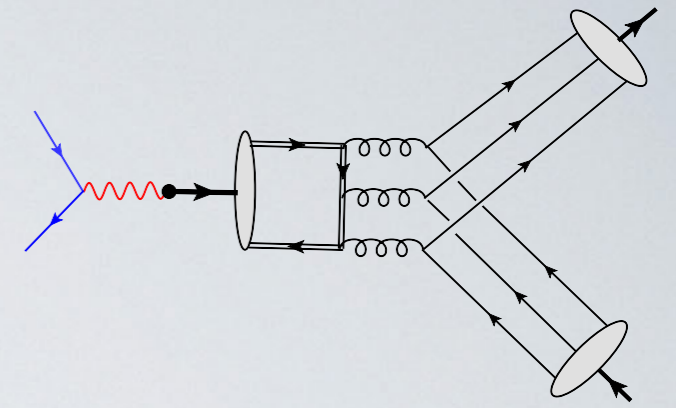
charmonium WF $|R_{10}(0)|^2 \simeq 0.81 \text{ GeV}^3$ $m_c = 1.48 \text{ GeV}$ Eichten Quigg, 1995

| | $ G_M^B $ | $\cos \phi_M$ | $Br[J/\psi \rightarrow B\bar{B}]$ | $Br[\text{data}]$ | γ_B | $\gamma_B[\text{data}]$ |
|--------------------------|-----------|---------------|-----------------------------------|---------------------|-------------------------|-------------------------|
| p | 3.47 | 0.45 | 2.13 | 2.12(3) | $0.67^{+0.16}_{-0.08}$ | 0.83(2) |
| n | 2.10 | 0.65 | 2.10 | 2.09(2) | $0.67^{+0.13}_{-0.06}$ | 0.95(6) |
| Λ | 2.29 | 0.50 | 1.71 | 1.89(9) | $0.67^{+0.17}_{-0.06}$ | 0.83(4) |
| Σ^0 Σ^+ | 2.00 | 0.50 | 1.23 | 1.17(3) 1.50(3) | $1.75^{+0.10}_{-0.19}$ | 2.11(5) 2.27(5) |
| Ξ^+ Ξ^0 | 1.60 | 0.50 | 1.14 | 0.97(8) 1.16(44) | $0.47^{+0.11}_{-0.003}$ | 0.61(5) 0.66(3) |

The numerical estimates are in agreement with the exp. data within 10-30% accuracy

Conclusions

All decay amplitudes associated with the $3g$ annihilation, are computed within the QCD EFT framework.



The considered models for baryon DAs provide reliable description of Br's for the relatively low norm. scale only $\mu^2 = 1.5 \text{ GeV}^2$

The interference of hadronic amplitudes and e.m. baryon FF's is important.

The numerical ratios of hadronic amplitudes are in agreement with the exp. data within 10-30% accuracy

A description of the Σ -channels requires a strong SU(3) violation for parameters of twist-4 DAs

The power corrections of order Λ^2/m_c^2 are also computed (higher twist 3q DAs).

They are about 4-15% for N/ Λ / Σ and 35% for Ξ , that is not very large.

Thanks!

