

Nucleon quasi parton distributions in the large N_c limit

with

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Introduction

Parton distribution functions (PDFs)

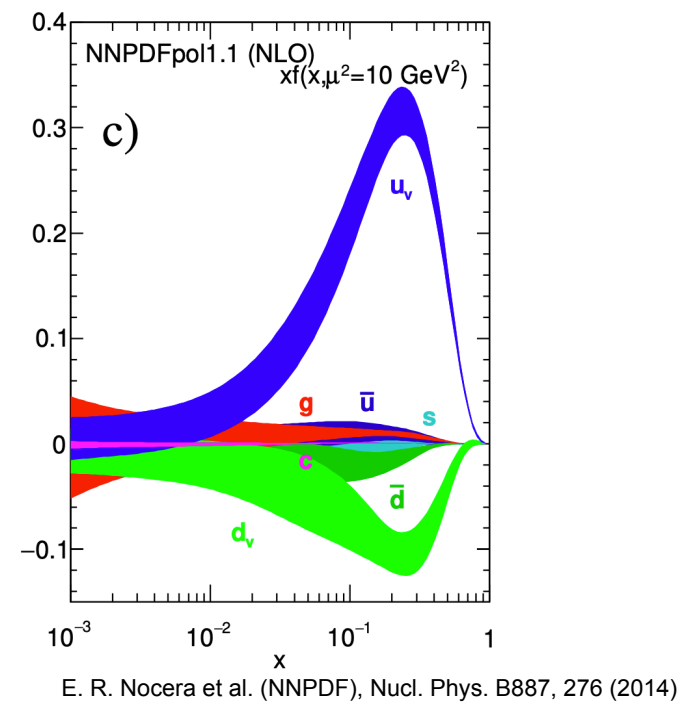
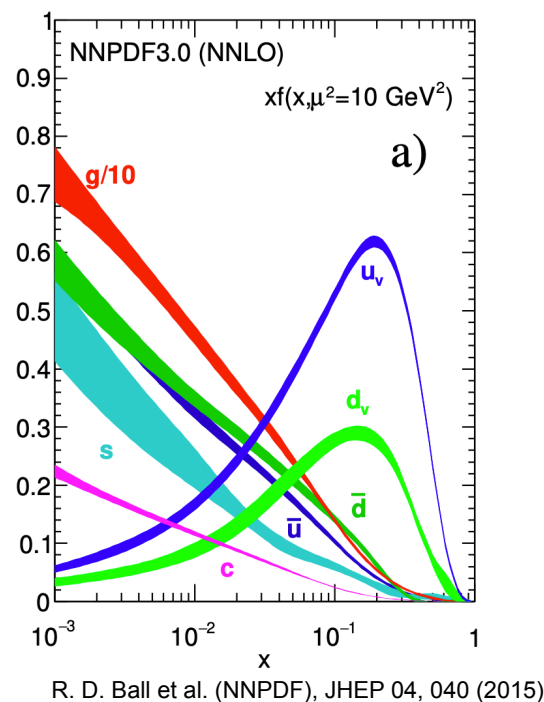
How partons (quarks and gluons) are distributed inside a hadron

Probability density on the light-cone

Non-perturbative QCD matrix elements

Factorization & universality

Plots from PDG 2021



Parton distribution functions (PDFs)

How partons (quarks and gluons) are distributed inside a hadron

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Factorization & universality

Theoretical understanding of PDFs

Effective models at low renormalization scale

- providing initial conditions of the QCD evolution
- to understand the properties of PDFs (qualitative and quantitative predictions)
- Chiral quark-soliton model (polarized antiquark asymmetry)

[D. Diakonov, V. Petrov, and P. Pobylitsa, Nucl. Phys. B 306, 809 (1988)]

[D. Diakonov, V. Y. Petrov, P. V. Pobylitsa, M. Polyakov, and C. Weiss, Nuclear Physics B 480, 341 (1996)]

Parton distribution functions (PDFs)

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Factorization & universality

Theoretical understanding of PDFs

Lattice QCD

- no direct computation is possible due to Euclidean nature
- studies on the Mellin moments of the PDFs, mostly

Quasi parton distribution function

Xiangdong Ji, Phys. Rev. Lett. 110, 262002 (2013)

$$q(x, \mu, P^z) = \int \frac{dz}{4\pi} e^{-ixP^z z} \langle P | \bar{\psi}(0) \gamma^z \exp \left[-ig \int_0^z dz' A^z(z') \right] \psi(z) | P \rangle + \mathcal{O}\left(\frac{\Lambda_{\text{QCD}}^2}{(P^z)^2}, \frac{M_N^2}{(P^z)^2}\right)$$

$$x \in (-\infty, +\infty)$$

μ : renormalization scale

P_z : nucleon momentum

Large Momentum Effective Theory

Jiunn-Wei Chen's talk on Monday for an overview

Spacelike equal-time matrix element \rightarrow can be calculated on the Lattice

No unique definition $\rightarrow \Gamma = \gamma^3$ or $\Gamma = \gamma^0$

Approaches to PDFs in the limit $P_z \rightarrow \infty$, or $v \rightarrow 1$.

Quasi parton distribution function

Xiangdong Ji, Phys. Rev. Lett. 110, 262002 (2013)

$$q(x, \mu_R, P^z) = \int_{-1}^1 \frac{dy}{|y|} \underbrace{C\left(\frac{x}{y}, \frac{\mu_R}{\mu}, \frac{\mu}{p^z}\right)}_{\text{Perturbative matching coefficients}} q(y, \mu) + \mathcal{O}\left(\frac{\Lambda_{\text{QCD}}^2}{(P^z)^2}, \frac{M_N^2}{(P^z)^2}\right)$$

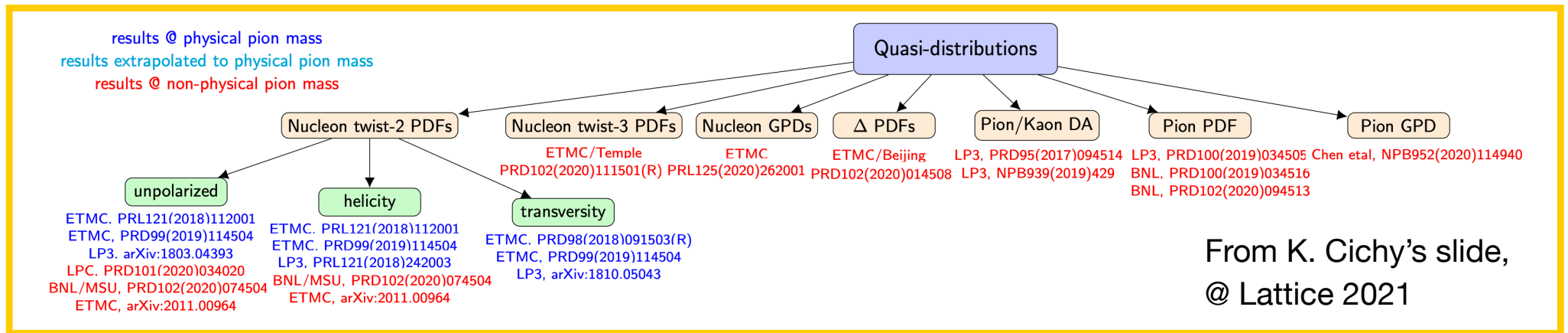
Perturbative matching coefficients

Divergences, Renormalization, matching coefficients...

Market results $P_z \sim 2\text{-}3$ GeV

Giunn-Wei Chen's talk on Monday for an overview

N, π , K / PDFs, DAs, GPDs ...



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Divergences, Renormalization, matching coefficients...

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N, π , K / PDFs, DAs, GPDs ...

Enough accuracy and uncertainty for actual application?

Reliable model computations on quasi-PDFs is needed

Review: K. Cichy and M. Constantinou, Adv.High Energy Phys. 2019 (2019) 3036904

Community report: M. Constantinou et al, Prog.Part.Nucl.Phys. 121 (2021) 103908
and many more..

(Quasi-)PDFs in the chiral quark-soliton model

Nucleon PDFs at low renormalization scale: a successful description

[D. Diakonov, V. Y. Petrov, P. V. Pobylitsa, M. Polyakov, and C. Weiss, Nuclear Physics B 480, 341 (1996)]

Nucleon matrix element in Euclidean separation

Lorentz boost \rightarrow PDFs \sim quasi-PDF

Properties of qPDFs for quarks and antiquarks in the nucleon:

Sum-rules, positivity, evolution in P_z

Gravitational form factor \bar{c}^q is related to the momentum sum-rule:

$\bar{c}^q \sim$ non-convergence of the separate quark EMT operator

Mass decomposition of the nucleon

Interaction strength between the quark- and gluon- subsystems

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[Maxim Polyakov and HDS, JHEP 09 (2018) 156]

In this talk,

Description of the chiral quark-soliton model

quasi-PDFs within the chiral quark-soliton model

Properties: **sum-rules** for the quasi-PDFs as their Mellin moments

Numerical results for the **unpolarized** and **longitudinally polarized** quasi-quark distributions

Comparison of the Dirac structures defining the quasi-PDFs: γ_0 vs γ_3

Summary

quasi-PDFs in the χ QSM

Quark distribution functions: large components

In general, in the large N_c limit:

Isosinglet unpolarized	$u(x) + d(x)$	$\sim N_c^2 \rho(N_c x)$
Isovector polarized	$\Delta u(x) - \Delta d(x)$	

Isovector unpolarized	$u(x) - d(x)$	$\sim N_c \rho(N_c x)$
Isosinglet polarized	$\Delta u(x) + \Delta d(x)$	

quasi-PDFs acquire overall factor of v
→ follow the same N_c ordering

Chiral quark-soliton model

[D. Diakonov, V. Petrov, and P. Pobylitsa, Nucl. Phys. B 306, 809 (1988)]

$$Z = \int \mathcal{D}\pi^a d\psi^\dagger d\psi \exp \int d^4x \psi^\dagger(x) (i\not{\partial} + iMU\gamma^5) \psi(x)$$

$$U^{\gamma^5}(x) = U(x) \frac{1 + \gamma^5}{2} + U^\dagger(x) \frac{1 - \gamma^5}{2} \quad U(x) = \exp \left[\frac{i}{F_\pi} \pi^a(x) \tau^a \right]$$

From QCD to the low energy effective theory via the instantons

Parameters: $\bar{\rho} \sim 1/3$ fm & $\bar{R} \sim 1$ fm

Intrinsic renormalization scale $\Lambda \sim 1/\bar{\rho} \approx 600$ MeV

Dynamically generated quark mass $M = 350$ MeV

Nucleon as the soliton made of N_c valence quarks in the self-consistent mean-field at large N_c

Interplays the quark-model and (topological) soliton picture of the baryons

Fully field theoretic: successively describes a wide class of baryon properties

Quasi-PDFs in the χ QSM

Isoscalar unpolarized $x \in (-\infty, \infty)$

$$\sum_f q_f(x, v) = N_c M_N v \sum_{n, occ} \int \frac{d^3 k}{(2\pi)^3} \delta(k^3 + vE_n - vM_N x) \left[\Phi_n^\dagger(\vec{k}) (1 + v\gamma^0 \gamma^3) \gamma_0 \Gamma \Phi_n(\vec{k}) \right]$$

Isvector polarized (helicity)

$$\Delta u(x, v) - \Delta d(x, v) = -\frac{1}{3} (2T^3) N_c M_N v \sum_{n, occ} \int \frac{d^3 k}{(2\pi)^3} \delta(k^3 + vE_n - vM_N x) \left[\Phi_n^\dagger(\vec{k}) (1 + v\gamma^0 \gamma^3) \gamma_0 \Gamma \tau^3 \gamma^5 \Phi_n(\vec{k}) \right]$$

Dirac eq. $H\Phi_n(\vec{x}) = E_n\Phi_n(\vec{x})$

$\Gamma = \gamma^0$ or $\Gamma = \gamma^3$

Sum-rules

Baryon number

$$\int_{-\infty}^{\infty} dx q(x, v) = \begin{cases} N_c B, & \Gamma = \gamma^0 \\ v N_c B, & \Gamma = \gamma^3 \end{cases}$$

Momentum

$$\int_{-\infty}^{\infty} dx x q(x, v) = \begin{cases} 1, & \Gamma = \gamma^0 \\ v, & \Gamma = \gamma^3 \end{cases}$$

Bjorken

$$\int_{-\infty}^{\infty} dx (\Delta u(x, v) - \Delta d(x, v)) = \begin{cases} v g_A^{(3)}, & \Gamma = \gamma^0 \\ g_A^{(3)}, & \Gamma = \gamma^3 \end{cases}$$

→ **better Dirac structure for P_N convergence ?**

→ **Interpretation of the QCD symmetry currents**

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U(1): charge density (γ^0) vs flux (γ^3)

Sum-rules

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Momentum sum-rule is satisfied only by quarks

Energy-momentum tensor: momentum flux ($T^{30} \sim \partial_3 \gamma^0$) vs pressure ($T^{33} \sim \partial_3 \gamma^3$)

In general, $M_2^q(\Gamma = \gamma^3) = v \left(A^q(0) - \frac{1 - v^2}{v^2} \bar{c}^q(0) \right)$

S. Bhattacharya et al, Phys.Rev.D 102 (2020) 5, 054021

Maxim Polyakov and HDS, JHEP 09 (2018) 156

Sum-rules

Baryon number

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Bjorken

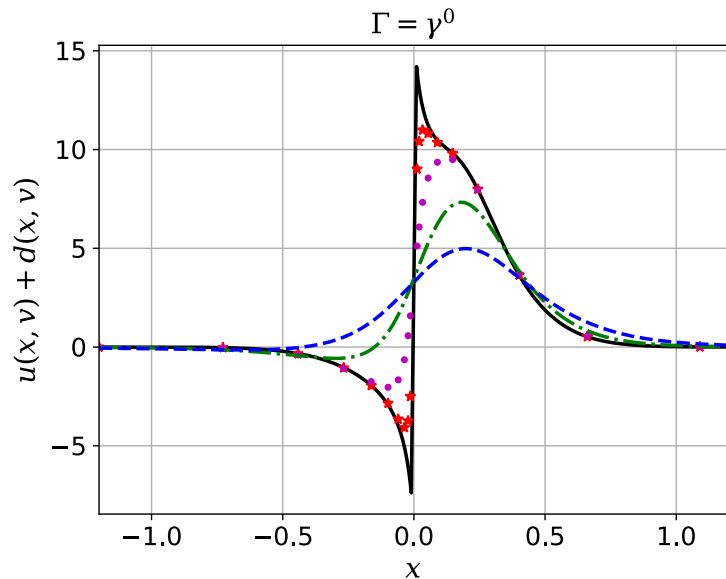
$$\int_{-\infty}^{\infty} dx (\Delta u(x, v) - \Delta d(x, v)) = \begin{cases} v g_A^{(3)}, & \Gamma = \gamma^0 \\ g_A^{(3)}, & \Gamma = \gamma^3 \end{cases}$$

Axial current: $\gamma^3 \sim S^3 g_A^{(3)}$ vs $\gamma^0 \sim \vec{S} \cdot \vec{v} g_A^{(3)}$

Numerical results

$u + d$

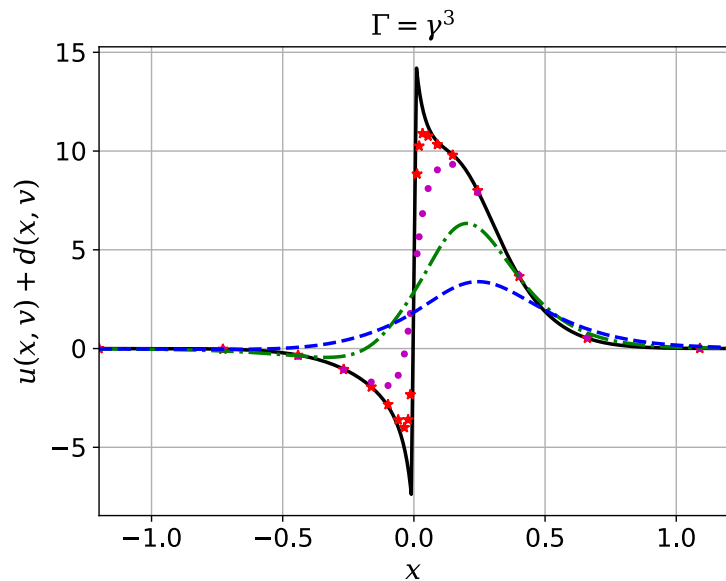
—	$v=1$	***	$v=0.999$	•••	$v=0.99$	-.-.-	$v=0.9$	- - -	$v=0.7$
$P_N/M_N=\infty$		22.3		7.0		2.1		1.0	



$$\bar{q}(x) = -q(-x) \text{ (LC PDF)}$$

Strong v dependence at small x : due to smearing of the quark and antiquark parts

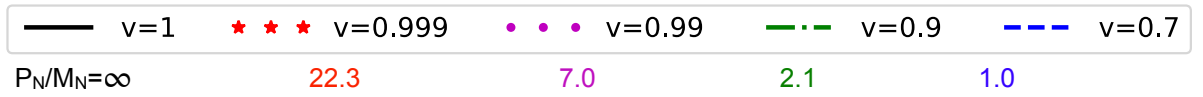
Antiquark part (negative x) breaks the positivity



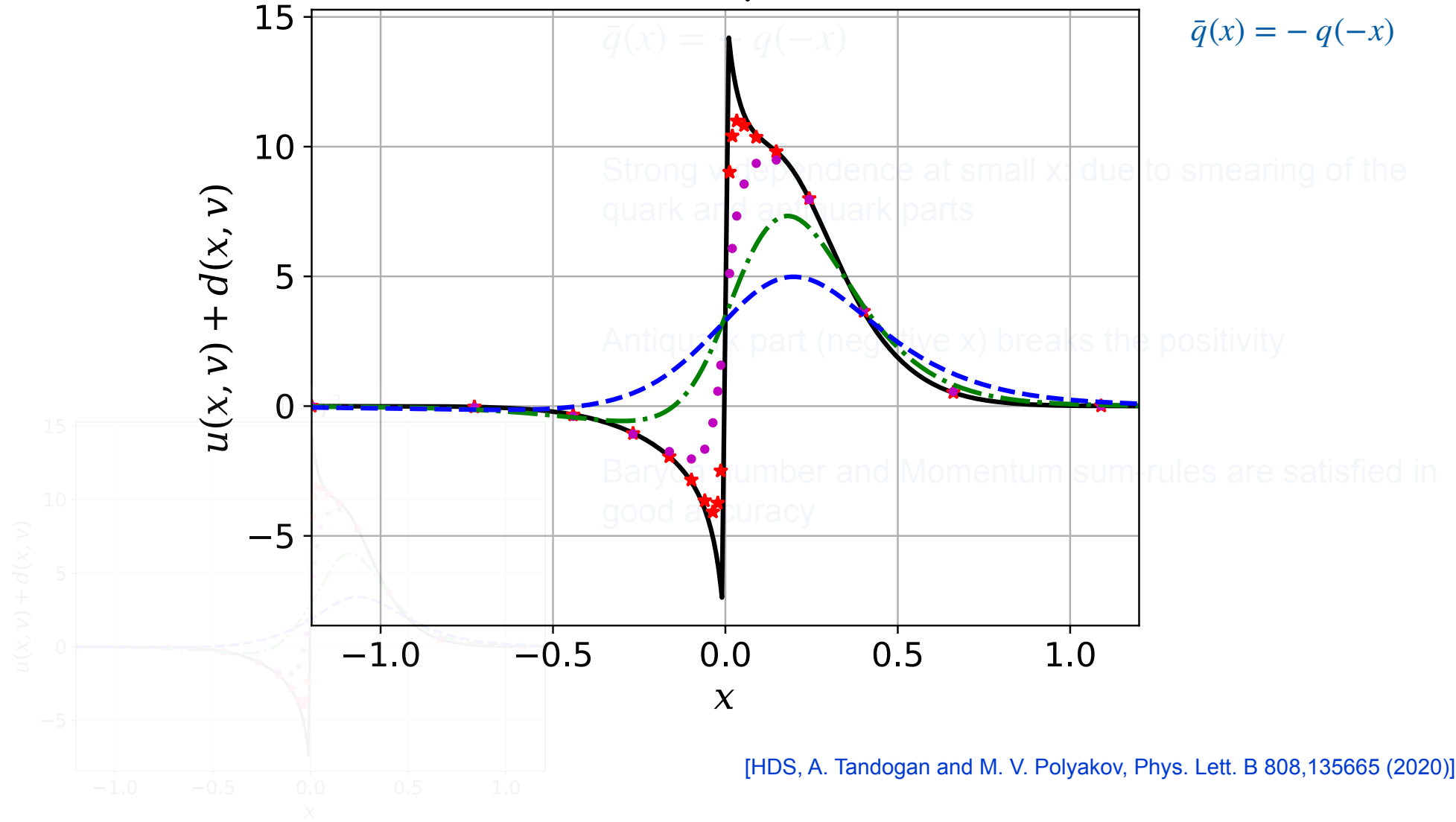
Baryon number and Momentum sum-rules are satisfied in good accuracy

[HDS, A. Tandogan and M. V. Polyakov, Phys. Lett. B 808,135665 (2020)]

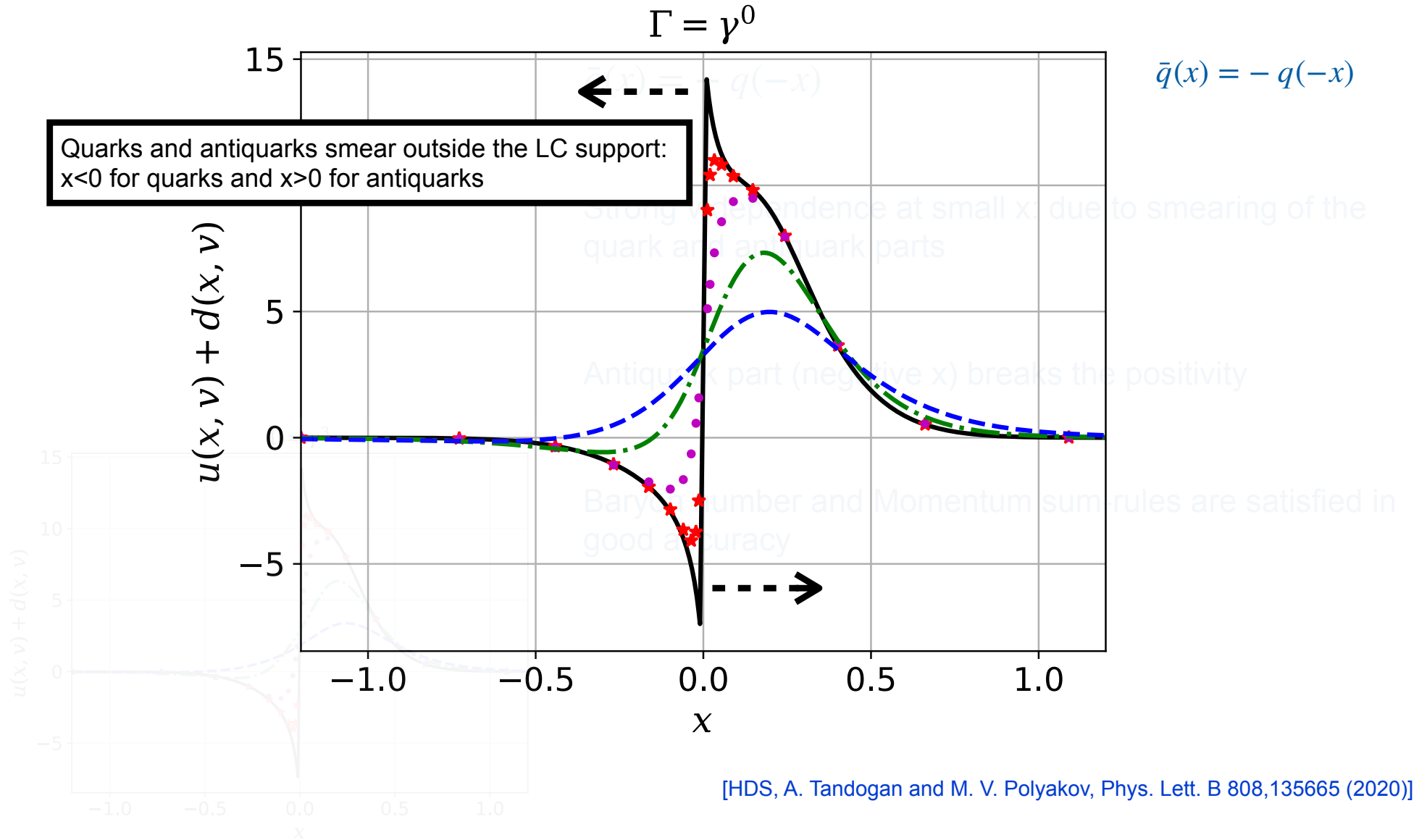
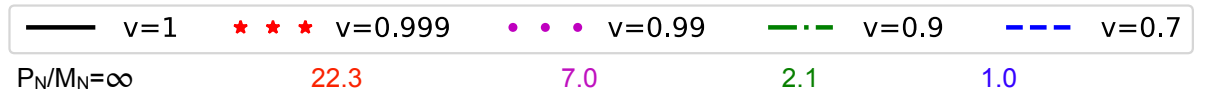
$u + d$



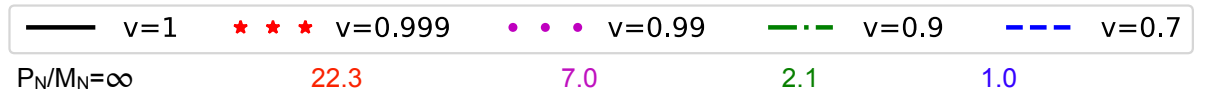
$\Gamma = \gamma^0$



$u + d$



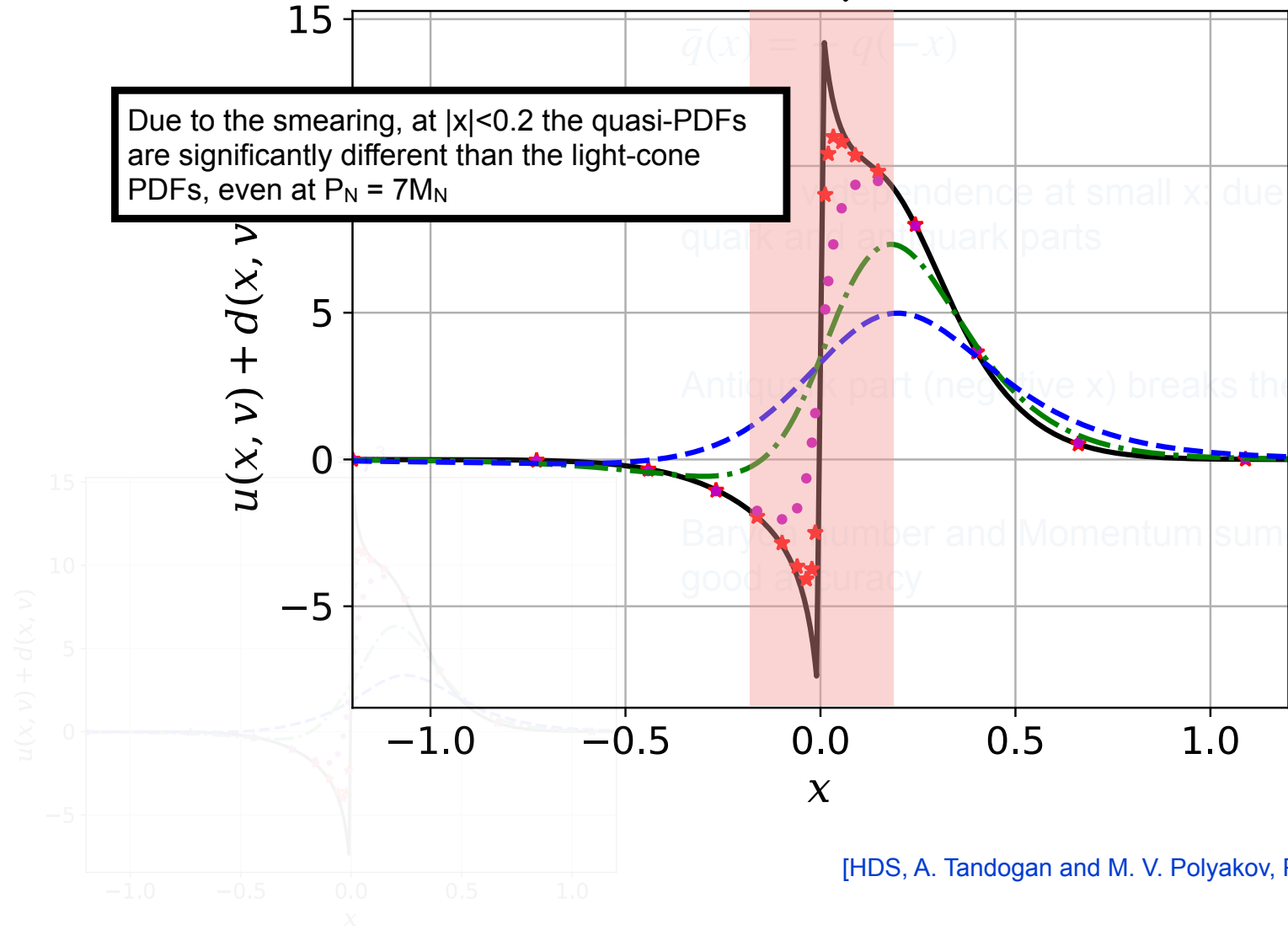
$u + d$



$\Gamma = \gamma^0$

$\bar{q}(x) = -q(-x)$

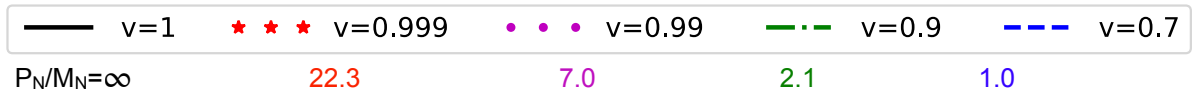
Due to the smearing, at $|x| < 0.2$ the quasi-PDFs are significantly different than the light-cone PDFs, even at $P_N = 7M_N$



[HDS, A. Tandogan and M. V. Polyakov, Phys. Lett. B 808,135665 (2020)]



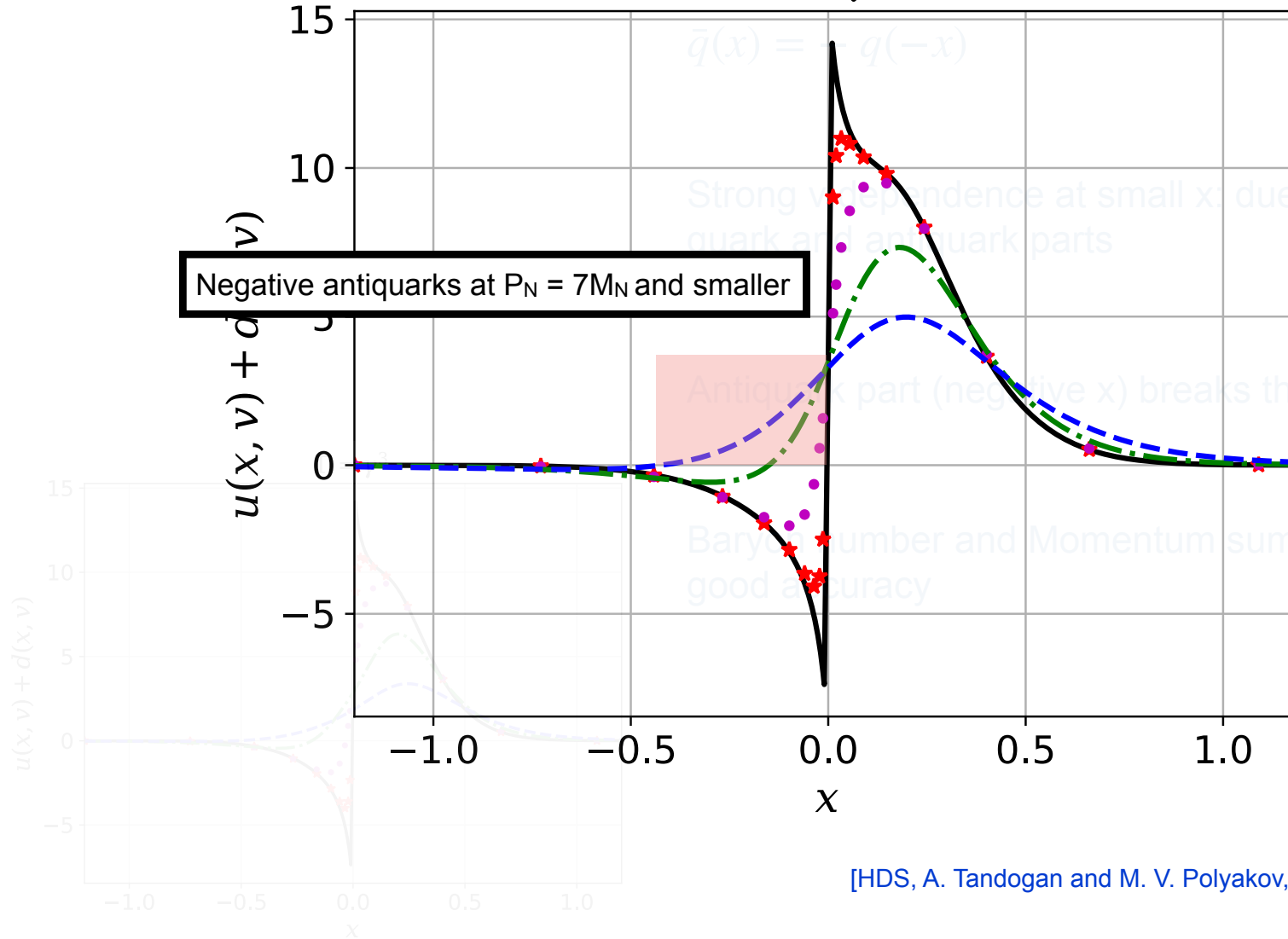
$u + d$



$\Gamma = \gamma^0$

$\bar{q}(x) = -q(-x)$

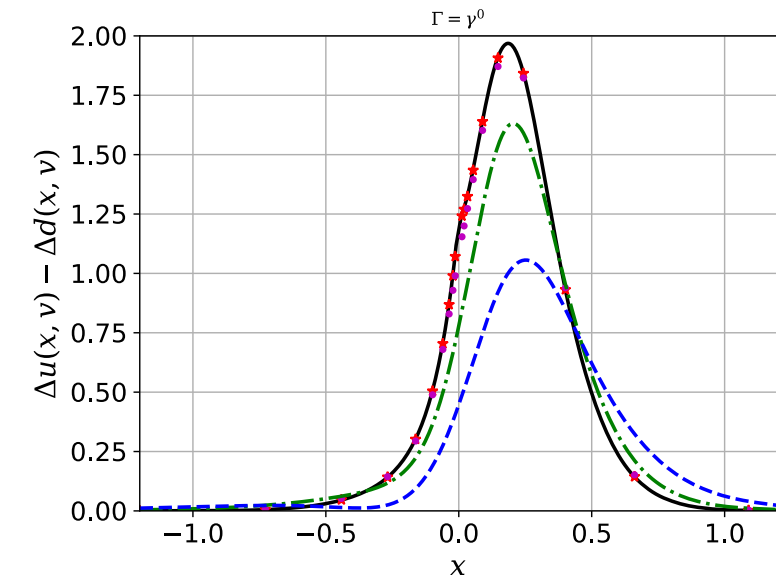
Negative antiquarks at $P_N = 7M_N$ and smaller



[HDS, A. Tandogan and M. V. Polyakov, Phys. Lett. B 808,135665 (2020)]

$\Delta u - \Delta d$

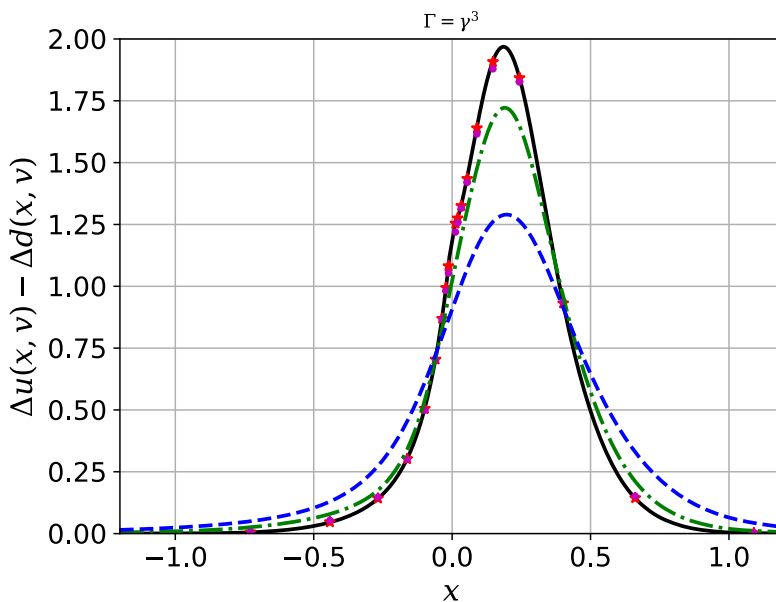
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$P_N/M_N=\infty$		22.3		7.0		2.1		1.0	



$$\Delta \bar{q}(x) = \Delta q(-x)$$

Antiquark asymmetry $\Delta \bar{u} - \Delta \bar{d} \neq 0$

At $v=0.9$ ($P \sim 2$ GeV), qPDF \sim PDF

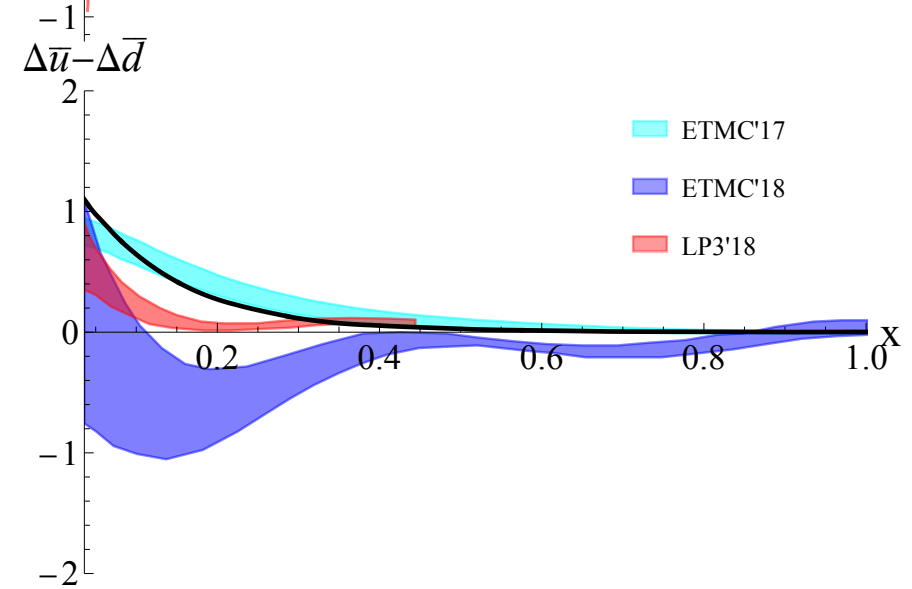
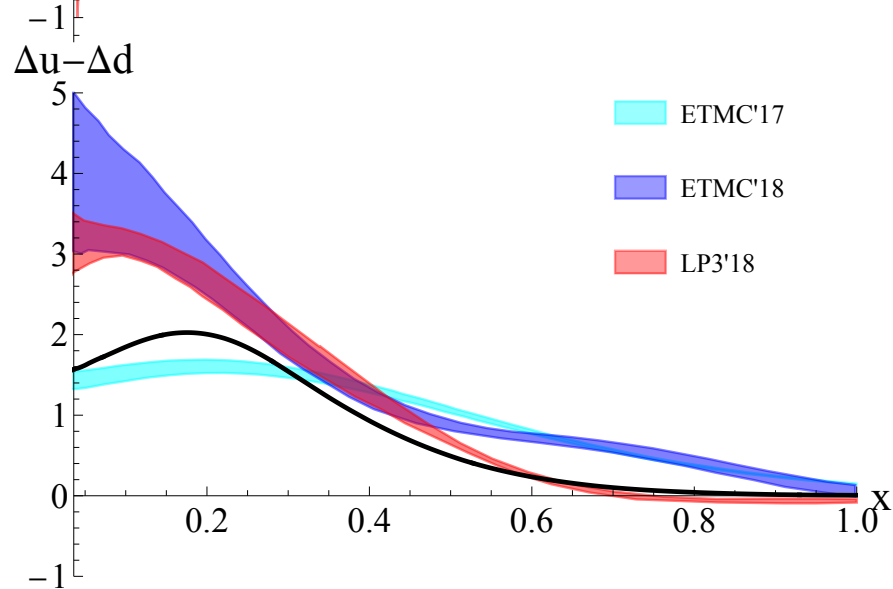
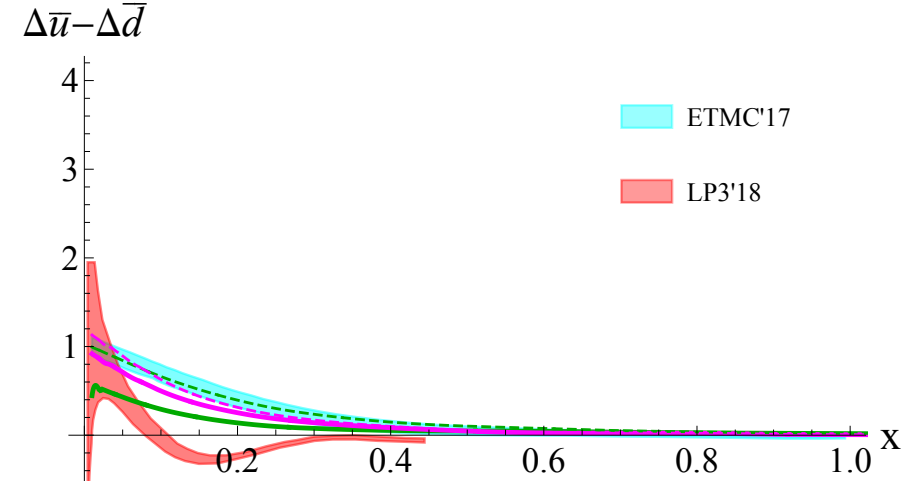
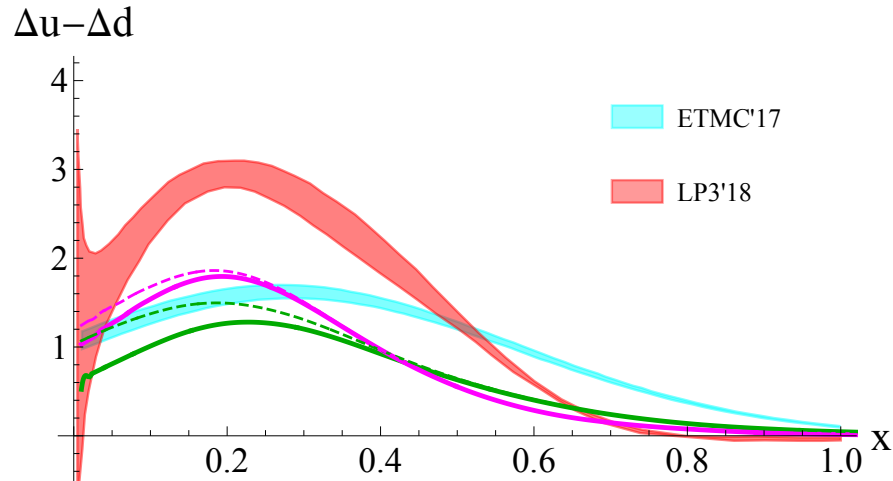


$\Gamma = \gamma^3$ qPDF approaches faster to the lightcone PDF

$$\int_{-\infty}^{\infty} dx (\Delta u(x, v) - \Delta d(x, v)) = \begin{cases} v g_A^{(3)}, & \Gamma = \gamma^0 \\ g_A^{(3)}. & \Gamma = \gamma^3 \end{cases}$$

vs. Lattice results

— $v = 1$ — $[v = 0.93, \Gamma = \gamma^0]$ - - - $[v = 0.93, \Gamma = \gamma^3]$ — $[v = 0.77, \Gamma = \gamma^0]$ - - - $[v = 0.77, \Gamma = \gamma^3]$
 $P_N/M_N = \infty$ 3.0 GeV 1.4 GeV



$(m_\pi, P_z, \mu) = (0.37, 1.4, 2.0)$ [ETMC'17 Alexandrou et al. Phys. Rev. D, vol. 96, no. 1, p. 014513, 2017]

$(0.13, 1.4, 2.0)$ [ETMC'18 Alexandrou et al. Phys.Rev.Lett. 121 (2018) 11, 112001, 2018]

$(0.135, 3.0, 3.0)$ [LP3'18 Lin et al. Phys. Rev. Lett., vol. 121, no. 24, p. 242003, 2018]

Closing remarks

Summary

χ QSM: a working framework for understanding the (quasi-)PDFs

Sum-rules: \bar{c}^q , 'better' Γ for the convergence to the PDFs

Good convergence of the isovector polarized

vs. poor $P_z \rightarrow \infty$ convergence for the isoscalar unpolarized quasi-PDF

Antiquark flavor asymmetry is predicted

Future tasks

Transversity PDF and other small components in the large N_c

Quantitative comparison vs Lattice: scale evolution, large x , $m_\pi \neq 0$

gluon PDFs (the EIC physics)

Thank you very much!