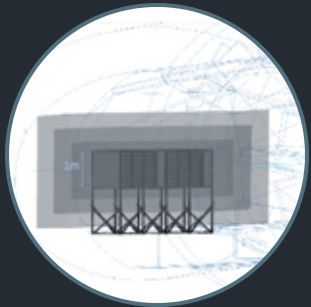


MoEDAL-MAPP Weekly Meets

July 31 2020

Michael Staelens
staelens@ualberta.ca

Overview



MAPP Detector Construction

Polishing and wrapping scintillator bars for MAPP-mCP (inner core of MAPP)

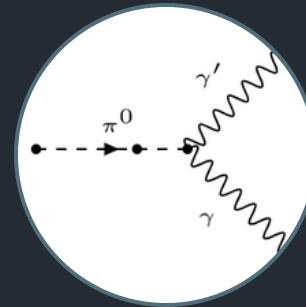
Update this week from Alejandro



MAPP - Fiducial Efficiency Calculations

Large # of simulations needed suggests using an alternative approach. (for LLPs)

Can use a mixture of simulations and numerical integration.. but we are still working this out.



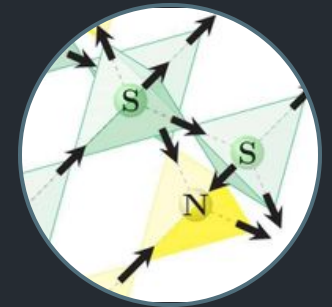
The Physics Case of the MAPP Detector

Updates on LLP Models:

Dark Photons

ALPs

mCP EPJ Paper Draft Shared



Emergent Monopoles

Modeling Emergent Monopole Excitations in spin ice.

Calculation Part 1: Second Quantized Hamiltonian

MAPP - Fiducial Efficiency Calculations

Fiducial Efficiency Calculations

The geometric acceptance of the CODEX-b box is $\sim 1\%$ (normalized to 4π). The LLP reach is attenuated further by the distribution of the LLP production and interplay between the LLP lifetime τ and the box depth. The number of LLP decay vertices expected in the box

$$N_{\text{box}} = \mathcal{L}_{\text{LHCb}} \times \sigma_{pp \rightarrow \varphi X} \times \int_{\text{vol}} \frac{d\varepsilon(r, \eta)}{dV} dV, \quad (1)$$

where the location of the box is specified by an azimuthal angle, the distance from the IP, r , and the pseudorapidity, η . In these coordinates, the differential fiducial efficiency is

$$\frac{d\varepsilon(r, \eta)}{dV} = \frac{1}{2\pi r^2 c\tau} \int d\beta w(\beta, \eta) \times \frac{e^{-r/(c\tau\beta\gamma)}}{\beta\gamma}. \quad (2)$$

Codex-b Method for Integrating over Box

c.f. Arxiv:1708.09395

Calculation of (3) from Arxiv:1708.09395 July 20, 2020

From (2), $\frac{d\varepsilon(r, \eta)}{dV} = \frac{1}{2\pi r^2 c\tau} \int d\beta w(\beta, \eta) \frac{e^{-r/(c\tau\beta\gamma)}}{\beta\gamma}$

Assume $w = \frac{S(\beta - \beta_0)}{2\Delta\beta_0}$ in the domain $\chi \in [0, 2, 0, 6]$
 $\Delta\beta_0 = 3$ w/ $\chi_0 = 5$

$\frac{d\varepsilon}{dV} = \frac{1}{2\pi r^2 c\tau} \int d\beta \frac{S(\beta - \beta_0)}{2\Delta\beta_0} \frac{e^{-r/(c\tau\beta\gamma)}}{\beta\gamma}$

$\frac{d\varepsilon}{dV} = \frac{1}{2\pi r^2 c\tau} \frac{1}{2\Delta\beta_0} \frac{e^{-r/(c\tau\beta_0\gamma)}}{\beta_0\gamma_0}$ let $r_0 = c\tau\beta_0\gamma_0$

$\frac{d\varepsilon}{dV} = \frac{1}{2\pi r^2 c\tau} \frac{1}{2\Delta\beta_0} e^{-r/r_0}$

Then, $\varepsilon_{\text{box}} = \int_{\text{vol}} \frac{1}{2\pi r^2 c\tau} \frac{1}{2\Delta\beta_0} e^{-r/r_0} dV(r, \eta)$
 $\varepsilon_{\text{box}} = \int_{r_0}^{r_2} \int_{\eta_0}^{\eta_2} \frac{1}{2\pi r^2 c\tau} \frac{1}{2\Delta\beta_0} e^{-r/r_0} r^2 dr d\eta$ (I'm ignoring the $\phi(\eta)$ bit in this calculation)

$\varepsilon_{\text{box}} = \frac{|\eta_2 - \eta_0|}{2\Delta\beta_0} \frac{|\theta_2 - \theta_0|}{2\pi} \left[e^{-r_1/r_0} - e^{-r_2/r_0} \right]$
 On this box domain $\chi \in [0, 2, 0, 6]$, we get

$\varepsilon_{\text{box}} = \frac{0.4}{2\Delta\beta_0} \frac{|\theta_2 - \theta_0|}{2\pi} \left[e^{-r_1/r_0} - e^{-r_2/r_0} \right] \quad (3)$

Calculation of (3)

Jacobian in χ - ϕ space for change from $V(x, y, z) \rightarrow V(r, \eta, \phi)$

χ = pseudorapidity from $-\infty$ to ∞ $\eta = \frac{1}{2} \ln \left[\frac{E + p_z}{E - p_z} \right]$
 ϕ = azimuthal angle from 0 to 2π
 r = the usual from 0 to ∞

$x = r \cos \phi \sin(2 \arctan(e^\eta))$
 $y = r \sin \phi \sin(2 \arctan(e^\eta))$
 $z = r \cos(2 \arctan(e^\eta))$

Def $x_\mu = \frac{\partial x}{\partial \mu}$

The Jacobian is (r, η, ϕ) , $J = \begin{vmatrix} x_r & x_\eta & x_\phi \\ y_r & y_\eta & y_\phi \\ z_r & z_\eta & z_\phi \end{vmatrix}$

$x_\eta = r \cos \phi \left[-\frac{2e^\eta \cos(2 \arctan(e^\eta))}{e^{2\eta} + 1} \right]$
 $y_\eta = r \sin \phi \left[-\frac{2e^\eta \cos(2 \arctan(e^\eta))}{e^{2\eta} + 1} \right]$
 $z_\eta = \frac{2e^\eta \sin(2 \arctan(e^\eta))}{e^{2\eta} + 1} r$

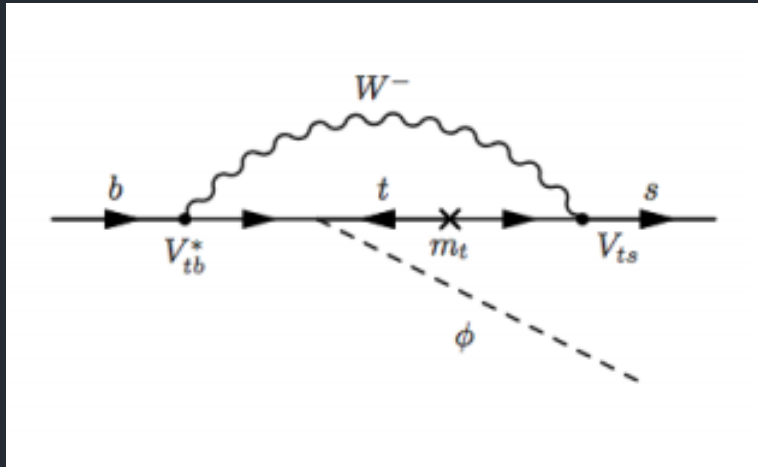
so, $J = \begin{vmatrix} \cos \phi \sin(2 \arctan(e^\eta)) & r \cos \phi \left[-\frac{2e^\eta \cos(2 \arctan(e^\eta))}{e^{2\eta} + 1} \right] & -r \sin \phi \sin(2 \arctan(e^\eta)) \\ \sin \phi \sin(2 \arctan(e^\eta)) & r \sin \phi \left[-\frac{2e^\eta \cos(2 \arctan(e^\eta))}{e^{2\eta} + 1} \right] & r \cos \phi \sin(2 \arctan(e^\eta)) \\ \cos(2 \arctan(e^\eta)) & r \left[\frac{2e^\eta \sin(2 \arctan(e^\eta))}{e^{2\eta} + 1} \right] & 0 \end{vmatrix}$

Using Method 1: $J = -r^2 \sin(2 \arctan(e^\eta)) \text{sech}(\eta) [\sin^2(2 \arctan(e^\eta)) + \cos(2 \arctan(e^\eta)) \tanh(\eta)]$

Jacobian for change to eta-phi space

Physics Performance of the MoEDAL-MAPP Detector

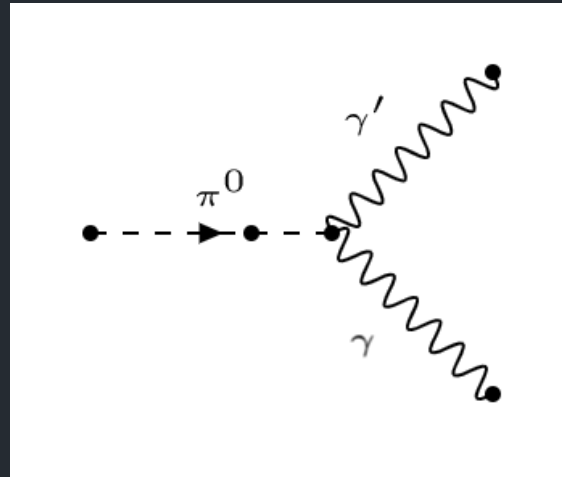
MAPP-LLP Physics Performance Cases



Light Scalar Portal (Dark Higgs)

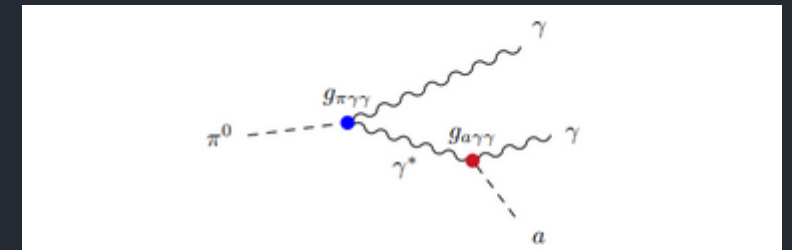
We are working on solutions to the large # of events required to generate MAPP-2, 3 ab-1, limit curves.

c.f. Arxiv:1708.09395



Vector Portal (Dark Photon)

Limit Plots are being updated (we can now show the full limits curves from 10-300 MeV or so)



Axion-Like Particle (ALPs)

Branching Ratio Formula Doesn't Match Plot

Calculated it w/ Mathematica to clarify.

c.f. Arxiv:1806.02348 (Published version still has same issue)

LLPs @ MAPP via Rare Meson Decays: A Summary

B, π and η decays are being studied.

Dark Higgs

$$N_{\phi,events} = \mathcal{L}_{LHCb} \times \sigma_{pp \rightarrow \phi X} \times BR(B \rightarrow X_s \phi) \quad (1)$$

$$N_{\phi,CodeX-b} = \mathcal{L}_{LHCb} \times \sigma_{b\bar{b}} \times BR(B \rightarrow X_s \phi) \times \epsilon_{fid,CodeX-b} \quad (2)$$

$$N_{\phi,MAPP} = \mathcal{L}_{LHCb} \times \sigma_{b\bar{b}} \times BR(B \rightarrow X_s \phi) \times \epsilon_{fid,MAPP} \quad (3)$$

where $\sigma_{b\bar{b}} \sim 500 \mu\text{b}$ for $\sqrt{s}=14 \text{ TeV}$, and $\mathcal{L}_{LHCb}=300 \text{ fb}^{-1}$. Here ϵ_{fid} should take into account $BR(\phi \rightarrow \mu^+ \mu^-)$. (Any cuts as well. e.g. tracking requirements.) In any case, ϵ_{fid} is obtained by simulations (Pythia8).

Dark Photons

$$N_{\gamma',events} = \mathcal{L}_{LHCb} \times \sigma_{pp,inel} \times N_{M/pp} \times BR(M \rightarrow \gamma' \gamma) \quad (4)$$

$$N_{\gamma',MAPP} = N_{M,tot} \times BR(M \rightarrow \gamma' \gamma) \times \epsilon_{fid,MAPP} \quad (5)$$

where $N_{M,tot}$ is the total number of expected mesons M given $\mathcal{L}_{LHCb}=300 \text{ fb}^{-1}$ and $\sigma_{pp,inel} \sim 75 \text{ mb}$ at $\sqrt{s}=13 \text{ TeV}$. Here ϵ_{fid} should take into account $BR(\gamma' \rightarrow l^+ l^-)$ for $(l = e, \mu)$.

More precisely $\sigma_{pp,inel} = 75.4 \pm 3.0 \pm 4.5 \text{ mb}$.
(Taken from <https://arxiv.org/pdf/1803.10974.pdf>)

N.B. I calculated $N_{\pi^0,tot} \sim 4.7 * 10^{17}$ (which matches the multiplicity of pions found by both Pythia simulations and the paper from ACTA PHYSICA POLONICA Vol. B4 1973). Eta multiplicities are in alignment as well. $N_{\eta,tot} \sim 5.1 * 10^{16}$.

Axion-like Particles

$$N_{a,events} = \mathcal{L}_{LHCb} \times \sigma_{pp,inel} \times N_{M/pp} \times BR(M \rightarrow a \gamma \gamma) \quad (6)$$

$$N_{a,MAPP} = N_{M,tot} \times BR(M \rightarrow a \gamma \gamma) \times \epsilon_{fid,MAPP} \quad (7)$$

Light Scalar Portal (Dark Higgs)

Limit plots enforce 2 requirements:

1. 4 Decays $\phi \rightarrow \mu^+ \mu^-$ inside MAPP-LLP-1
2. Both μ hit at least 2 planes of the detector (~80% eff.?)

Vector Portal (Dark Photon)

Limit plots here require only 3 decays of $\gamma' \rightarrow l^+ l^-$ ($l=e, \mu$) inside MAPP-LLP-1.

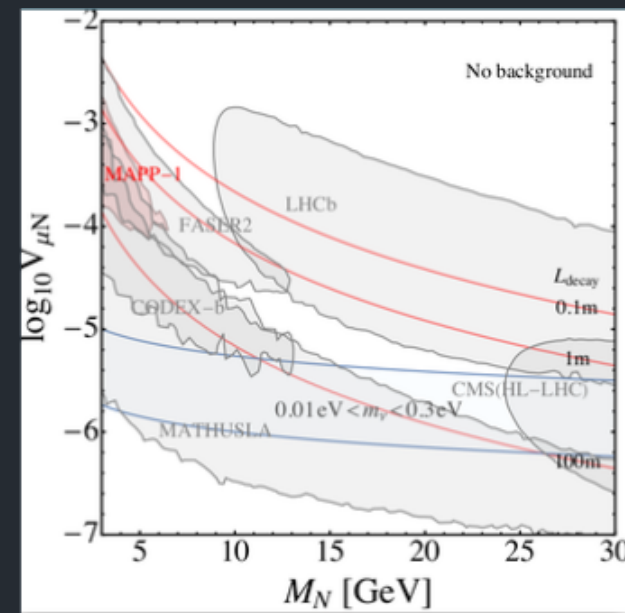
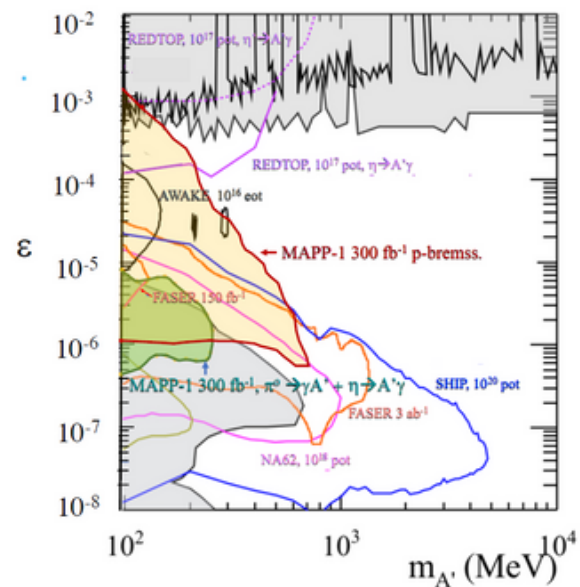
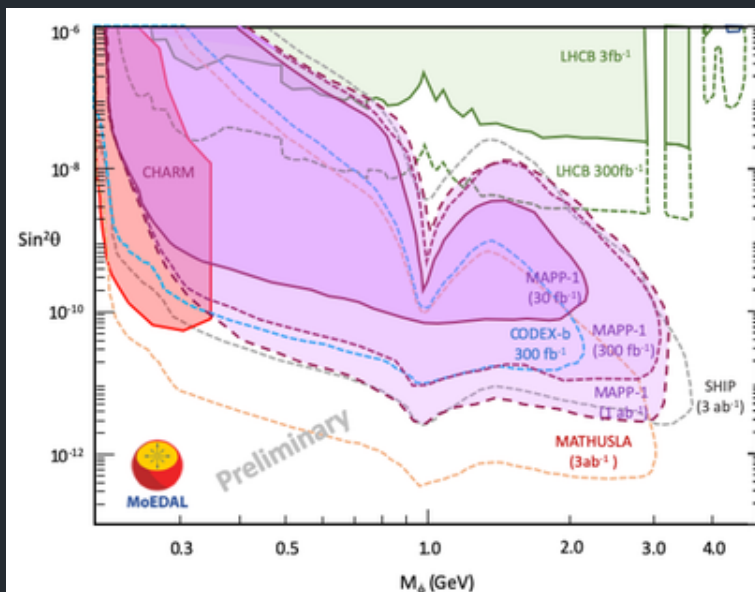
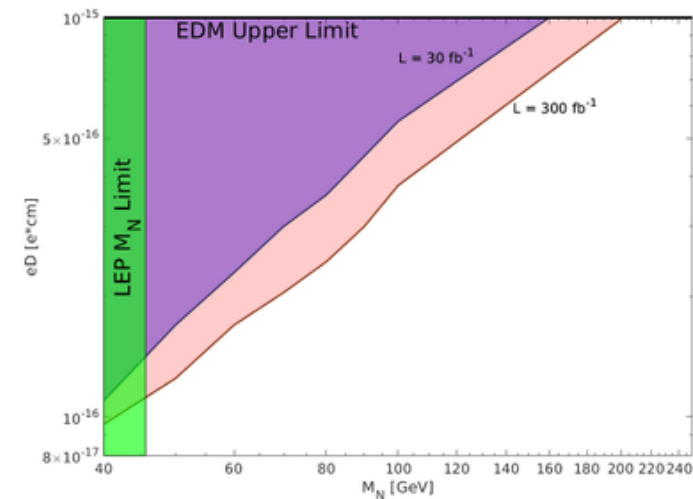
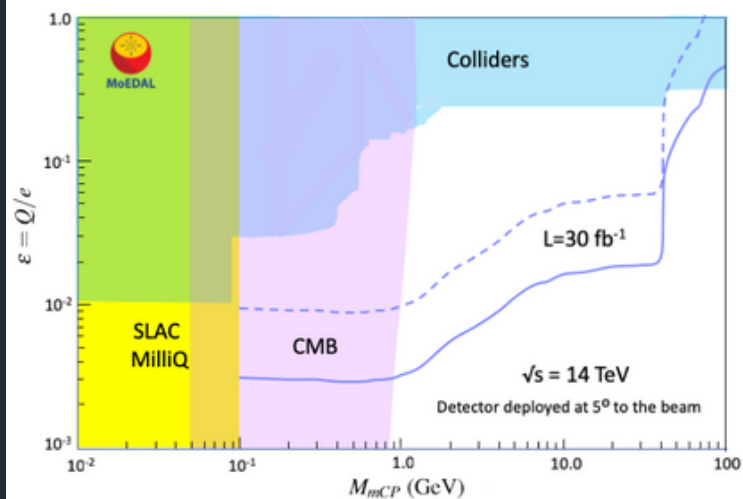
No cuts.

Axion-like Particle (ALPs)

First limit plots here will be for MAPP-mCP-1 and MAPP-LLP-1 (Phase 1 section), requiring 3 fiducial decays of $a \rightarrow \gamma \gamma$.

Slide will be updated to include LLP decay length/width formulae as well.

MAPP-1.. Results So Far



Progress with ALPs @ MAPP-1 (via rare meson decays)

The differential decay width is, then,

$$\frac{d\Gamma(\pi^0 \rightarrow a\gamma\gamma)}{dE_1 dE_2} = \frac{1}{2} \frac{1}{(4\pi)^3 M} \sum_{\text{pol}} |M|^2 = \frac{(g_{\pi^0 \gamma\gamma} g_{a\gamma\gamma})^2}{2(4\pi)^3 M} f(E_1, E_2), \quad (\text{B4})$$

where

$$f(E_1, E_2) = \frac{E_1^2 E_2^2 [M^2 (1 + \cos^2 \theta_{12}) + 2E_1(E_1 - M)(1 - \cos \theta_{12})^2]}{(M - 2E_1)^2} + 2 \frac{E_1^2 E_2^2 [M^2 (1 + \cos^2 \theta_{12}) + [2E_1 E_2 - M(E_1 + E_2)](1 - \cos \theta_{12})^2]}{(M - 2E_1)(M - 2E_2)} + \frac{E_1^2 E_2^2 [M^2 (1 + \cos^2 \theta_{12}) + 2E_2(E_2 - M)(1 - \cos \theta_{12})^2]}{(M - 2E_2)^2}. \quad (\text{B5})$$

Integrating over phase space results in the total decay width

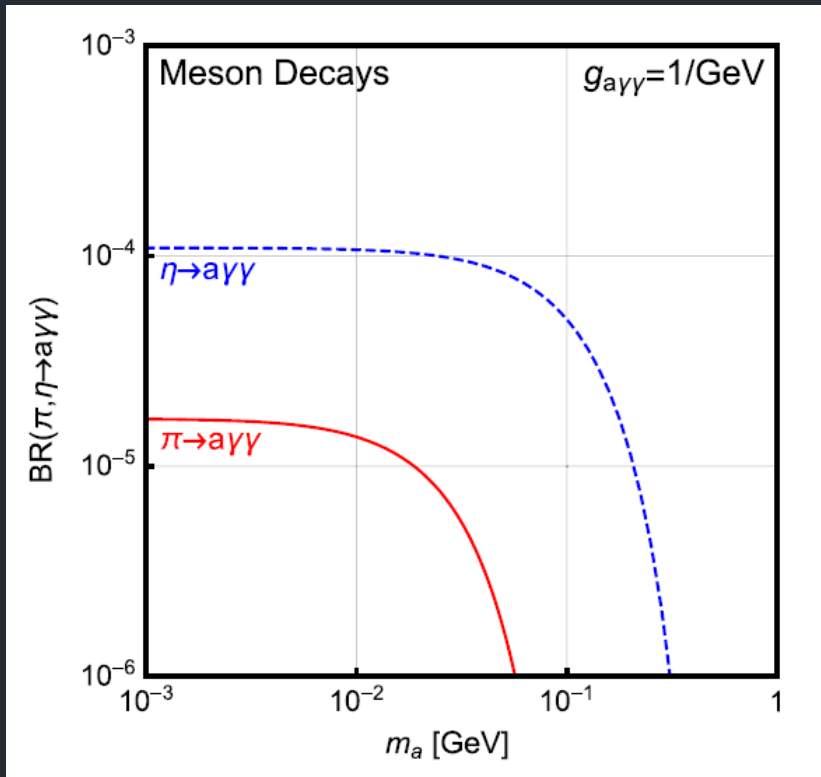
$$\Gamma(\pi^0 \rightarrow a\gamma\gamma) = \int_0^{\frac{M^2 - m^2}{2M}} dE_1 \int_{\frac{M^2 - m^2}{2M} - E_1}^{\frac{M^2 + m^2}{2M}} dE_2 \frac{d\Gamma(\pi^0 \rightarrow a\gamma\gamma)}{dE_1 dE_2} = \frac{(g_{\pi^0 \gamma\gamma} g_{a\gamma\gamma})^2}{768(4\pi)^3 M^3} F(M, m), \quad (\text{B6})$$

where

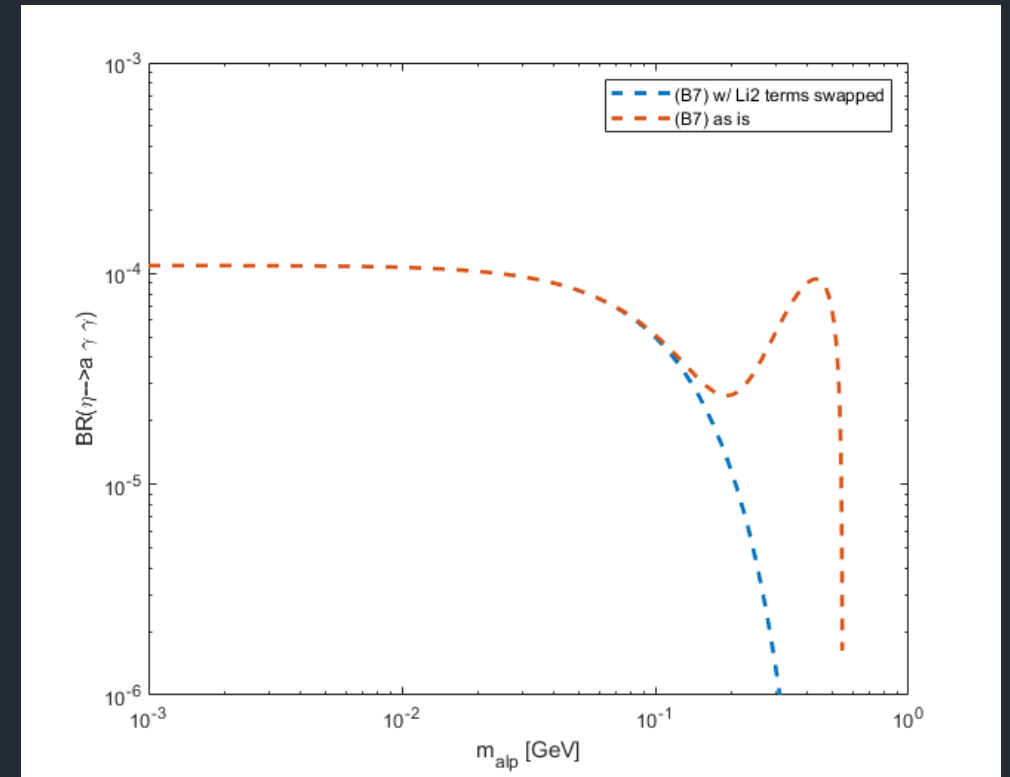
$$F(M, m) = 24 \log\left(\frac{m}{M}\right) \left[6m^2 M^2 (M^4 + m^4) + 15m^4 M^4 + 2m^4 M^4 \log\left(\frac{mM}{m^2 + M^2}\right) \right] + 7(M^8 - m^8) + 148M^2 m^2 (M^4 - m^4) + 24m^4 M^4 \left[\text{Li}_2\left(\frac{m^2}{m^2 + M^2}\right) - \text{Li}_2\left(\frac{M^2}{m^2 + M^2}\right) \right]. \quad (\text{B7})$$

- B6 is plotted & calculated in [PhysRevD.98.055021](#)
- In-house script to plot B6 as validation of the calculation does not work (unless one switches the order of the Li2 terms)
- Plots on the next slide show this..
- B6 & B7 will be calculated to check this.

FASER Decay width plot vs. mine



FASER



Me

(Eta process)

Calculation of B6 & B7 in Mathematica

The differential decay width is, then,

$$\frac{d\Gamma(\pi^0 \rightarrow a\gamma\gamma)}{dE_1 dE_2} = \frac{1}{2(4\pi)^3 M} \sum_{\text{pol}} |M|^2 = \frac{(g_{\pi^0\gamma\gamma} g_{a\gamma\gamma})^2}{2(4\pi)^3 M} f(E_1, E_2), \quad (\text{B4})$$

where

$$f(E_1, E_2) = \frac{E_1^2 E_2^2 [M^2 (1 + \cos^2 \theta_{12}) + 2E_1(E_1 - M)(1 - \cos \theta_{12})^2]}{(M - 2E_1)^2} + 2 \frac{E_1^2 E_2^2 [M^2 (1 + \cos^2 \theta_{12}) + [2E_1 E_2 - M(E_1 + E_2)](1 - \cos \theta_{12})^2]}{(M - 2E_1)(M - 2E_2)} + \frac{E_1^2 E_2^2 [M^2 (1 + \cos^2 \theta_{12}) + 2E_2(E_2 - M)(1 - \cos \theta_{12})^2]}{(M - 2E_2)^2}. \quad (\text{B5})$$

Integrating over phase space results in the total decay width

$$\Gamma(\pi^0 \rightarrow a\gamma\gamma) = \int_0^{\frac{M^2 - m^2}{2M}} dE_1 \int_{\frac{m^2 - M^2}{2E_1}}^{\frac{M^2 - m^2}{2M}} dE_2 \frac{d\Gamma(\pi^0 \rightarrow a\gamma\gamma)}{dE_1 dE_2} = \frac{(g_{\pi^0\gamma\gamma} g_{a\gamma\gamma})^2}{768(4\pi)^3 M^3} F(M, m), \quad (\text{B6})$$

where

$$F(M, m) = 24 \log\left(\frac{m}{M}\right) \left[6m^2 M^2 (M^4 + m^4) + 15m^4 M^4 + 2m^4 M^4 \log\left(\frac{mM}{m^2 + M^2}\right) \right] + 7(M^8 - m^8) + 148M^2 m^2 (M^4 - m^4) + 24m^4 M^4 \left[\text{Li}_2\left(\frac{m^2}{m^2 + M^2}\right) - \text{Li}_2\left(\frac{M^2}{m^2 + M^2}\right) \right]. \quad (\text{B7})$$

- Feeding the integral to Mathematica
- Need to go term by term..
- ~12 integrals to do.. will take a bit more time.
- Also, still need the Full Decay width for pi0..

Is Phys. Rev. Lett. 33, 1400 reliable? (8.02±0.42 eV???)

```
(* Try term by term... *)
(* Row 1 - Term 1 *)
Integrate[ $\frac{1}{(M - 2 * E1)^2} E1^2 E2^2 * (M^2 * (1 + (\frac{1}{2 * E1 * E2} * (M^2 - m^2 - 2 * M * (E1 + E2) + 2 * E1 * E2))^2))$ , {E2,  $\frac{M^2 - m^2}{2 * M} - E1, \frac{M}{2} + \frac{m^2}{4 * E1 - M}$ }, {E1,  $\frac{m^2 + (2 * E1 - M) * M}{3 * (2 * E1 - M)^5 * M}$ }]

(* Row 1 - Term 2 *)
Integrate[ $\frac{1}{(M - 2 * E1)^2} E1^2 E2^2 * (2 * E1 * (E1 - M) * (1 - \frac{1}{2 * E1 * E2} * (M^2 - m^2 - 2 * M * (E1 + E2) + 2 * E1 * E2))^2)$ , {E2,  $\frac{M^2 - m^2}{2 * M} - E1, \frac{M}{2} + \frac{m^2}{4 * E1 - M}$ }, {E1,  $\frac{2 * E1^4 * (E1 - M) * (m^2 + (2 * E1 - M) * M)^3}{3 * M * (-2 * E1 + M)^5}$ }]

FullSimplify[ $\frac{E1^3 * (m^2 + (2 * E1 - M) * M)^3 * (2 * E1^2 - 2 * E1 * M + M^2)}{3 * (2 * E1 - M)^5 * M} - \frac{2 * E1^4 * (E1 - M) * (m^2 + (2 * E1 - M) * M)^3}{3 * M * (-2 * E1 + M)^5}$ ]

 $\frac{E1^3 * (m^2 + (2 * E1 - M) * M)^3}{3 * (2 * E1 - M)^5 * M}$ 
```

Then integrating over E1....

```
(* If you scale it by the constants out front... (first line of eq. B5 only) *)
```

$$\frac{g_{\pi^0\gamma\gamma}^2 g_{a\gamma\gamma}^2}{384 (4\pi)^3 M^3} * \left[-m^8 - 28 m^6 M^2 + 28 m^2 M^6 + M^8 + 12 m^2 M^2 (m^4 + 3 m^2 M^2 + M^4) \left(\text{Log}\left[-\frac{m^2}{M}\right] - \text{Log}[-M] \right) \right]$$

Plan Ahead

- 1. Continuing to investigate the integration formula for the fiducial efficiency with Ameir.
- 2. Run ALP simulations on office PC to get 95% CL for MAPP-mCP-1 & MAPP-LLP-1
- 3. Run Dark Photon simulations on laptop, to complete 95% C.L. for MAPP-LLP-1
- 4. Will also give an update next week on CM monopoles.
- 5. Some words on the mIP paper for EPJ...

GEANT4 Simulations to come soon, we chat w/ Matti this morning.

Questions?

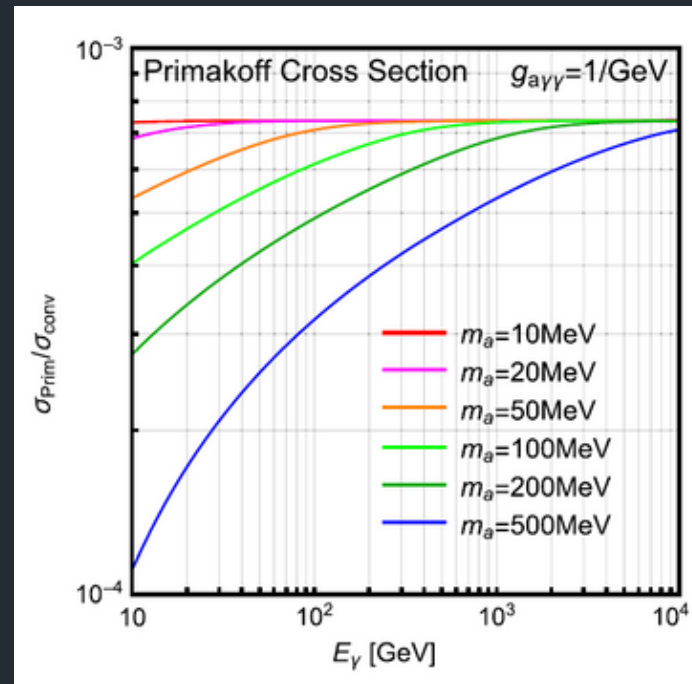
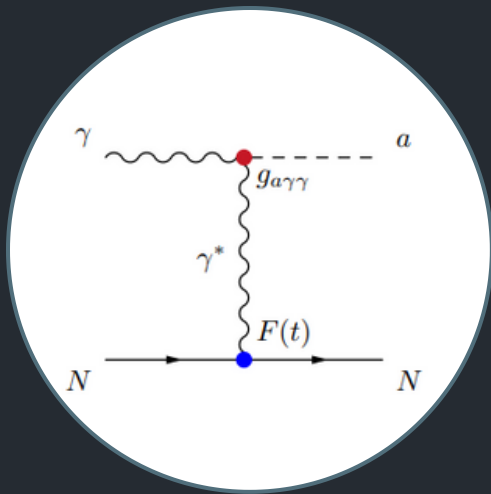
Modeling Emergent Monopoles

Using the Method of Coherent Structures (MCS)

Second Quantized Hamiltonian next week.

Future work.. ALPs - Production via Primakoff Process

For the Primakoff Process; resonant production of neutral pseudoscalar mesons by high-energy photons interacting with an atomic nucleus. ($\gamma N \rightarrow a N$)



Ratio between the Primakoff and pair-production cross section (in iron)

The relevant pair-production cross section in iron for the photon energies of interest is of the order of $\sigma_{\text{conv}} = 5$ barn [Faser ALP paper]

MAPP-mCP Physics Performance Cases



Case 1: Minicharged (mCP) Particles in Dark QED

Kinetic mixing leads to fractionally charged particles

Mass range of 0.1-100 GeV

Charges could be as low as $\sim e/500$



Case 2: “Heavy Neutrinos” With a Large EDM

EDM of a new heavy neutral Dirac fermion could be as large as $\sim 10^{-15} e \cdot \text{cm}$ (predicted by 5-D EFT).

Mass range of ~ 45.6 -200 GeV

EDMs as low as 10^{-16}



Case 3: (Dark) Mini-Dyons via “Magnetic Mixing”

WIP.. Still in the model building & computational implementation phase.

