Martiana 12 August 2020
Lagrongian

$$\begin{aligned} \mathcal{J} &= -\frac{n^{\alpha}n^{A}}{2n^{2}} g^{\mu\nu} (F_{\alpha\rho}^{A} F_{\mu\nu}^{A} + F_{\alpha\rho}^{B} F_{\mu\nu}^{B}) + \\ &+ \frac{n^{\alpha}n^{A}}{4n^{2}} e^{\mu\nu\nu\sigma} (F_{\alpha\nu}^{B} F_{\sigma\sigma}^{A} - F_{\alpha\nu}^{A} F_{\sigma\sigma}^{B}) - J_{\mu}A^{\mu} - K_{\mu}B^{\mu} \\ &- \frac{n^{\alpha}n^{A}}{4n^{2}} g^{\mu\nu} (F_{\alpha\rho}^{Ab} F_{\mu\nu}^{Ab} + F_{\alpha\rho}^{Bb} F_{\mu\nu}^{Bb}) \\ &+ \frac{eee_{D}}{n^{2}} g^{\mu\nu} (F_{\alpha\rho}^{Ab} F_{\mu\nu}^{Ab} - F_{\alpha\rho}^{Ab} F_{\mu\nu}^{Bb}) \\ &+ \frac{eee_{D}}{n^{2}} g^{\mu\nu} (F_{\alpha\rho}^{Ab} F_{\alpha\nu}^{Ab} - F_{\alpha\rho}^{Ab} F_{\mu\nu}^{Bb}) \\ &+ \frac{n^{\alpha}n^{A}}{4n^{2}} e^{\mu\nu\sigma} (F_{\alpha\nu}^{Ab} F_{\sigma\sigma}^{Ab} - F_{\alpha\nu}^{Ab} F_{\sigma\sigma}^{Bb}) - J_{\mu}^{D}A_{\mu}^{A} - K_{\mu}^{Ab} F_{\mu\nu}^{Ab} \\ &+ \frac{n^{\alpha}n^{A}}{4n^{2}} e^{\mu\nu\sigma} (F_{\alpha\nu}^{Ab} F_{\sigma\sigma}^{Ab} - F_{\alpha\nu}^{Ab} F_{\sigma\sigma}^{Bb}) - J_{\mu}^{D}A_{\mu}^{A} - K_{\mu}^{Ab} F_{\mu\nu}^{Ab} \\ &+ \frac{n^{\alpha}n^{A}}{4n^{2}} e^{\mu\nu\sigma} (F_{\alpha\nu}^{Ab} F_{\sigma\sigma}^{Ab} - F_{\alpha\nu}^{Ab} F_{\sigma\sigma}^{Bb}) - J_{\mu}^{D}A_{\mu}^{A} - K_{\mu}^{Ab} F_{\mu\nu}^{Ab} \\ &+ \frac{n^{\alpha}n^{A}}{4n^{2}} e^{\mu\nu\sigma} (F_{\alpha\nu}^{Ab} F_{\sigma\sigma}^{Ab} - F_{\alpha\nu}^{Ab} F_{\sigma\sigma}^{Bb}) - J_{\mu}^{D}A_{\mu}^{A} - K_{\mu}^{Ab} F_{\mu\nu}^{Ab} \\ &- K_{\mu}^{A} = e \overline{\psi} e^{\sigma} \Psi_{e} \\ &K_{\mu}^{A} = g \overline{\psi} g^{\sigma} \mathcal{H}_{\mu}^{Ab} \\ &J_{\mu}^{A} = e \overline{\psi} e^{\sigma} \mathcal{H}_{\mu}^{Ab} \\ &K_{\mu}^{D} = g_{D} \overline{\psi} g_{\sigma} \mathcal{H}_{\mu}^{Ab} \\ &K_{\mu}^{D} = g_{D} \overline{\psi} g_{\sigma} \mathcal{H}_{\mu}^{Ab} \\ &Note: normalization of hagrangian is chosen for diagonalization of kinetic terms \\ &g, g_{\mu} must match with Terming s on Verhaaren s \\ &\vdots H will amount to redifining F_{\mu\nu}^{A} to indude couplings. \\ \end{array}$$

Equations of motion (1) $\frac{n_{\alpha}}{h^{2}} \left(n \frac{\partial}{\partial y} F_{A}^{\alpha \gamma} - n \frac{\partial}{\partial y} F_{A}^{\alpha \mu} - \varepsilon \frac{\mu \gamma \kappa}{\beta} n_{y} \frac{\partial}{\partial y} F_{B}^{\beta \beta} \right)$ $+ \underbrace{eee}_{h^2} h_{x} \partial_{y} \left(h F_{Ap}^{xy} - n F_{Ap}^{xy} \right) = J^{k}$ (2) $\frac{n_{\alpha}}{n^{2}} \left(n^{\lambda} \partial_{y} F_{B}^{\alpha \nu} - n^{\nu} \partial_{y} F_{B}^{\alpha \mu} + \varepsilon^{\mu \nu \alpha} B n_{y} \partial_{y} F_{A}^{\beta \mu} \right)$ $-\underbrace{eee}_{n^2} n_x \partial_y \left(n^{\mu} F_{B_p}^{\alpha \nu} - n^{\nu} F_{B_p}^{\alpha \mu} \right) = K^{\mu}$ DARK $(3) \frac{n_{\alpha}}{n^{2}} \left(n^{2} \partial_{y} F_{Ap}^{\alpha \nu} - n^{2} \partial_{v} F_{Ap}^{\alpha \mu} - \varepsilon^{\mu \nu \alpha} p n_{y} \partial_{v} F_{Bp}^{\beta \mu} \right)$ $+ \underbrace{eee}_{n^2} h_{\alpha} \partial_{\nu} (n^{\mu} F_{A}^{\alpha \nu} - n^{\nu} F_{A}^{\alpha \mu}) = J_{D}^{\mu} + M_{AD}^2 A_{D}^{\mu} A_{D \mu}$ (4) $\frac{n_{\lambda}}{n^{2}}\left(n^{\lambda}\partial_{y}F_{B_{D}}^{\alpha\gamma}-n^{\gamma}\partial_{y}F_{B_{D}}^{\alpha\mu}+\varepsilon \frac{\mu_{\lambda}}{\beta}n_{\gamma}\partial_{y}F_{A_{D}}^{\gamma\beta}\right)$ $-\frac{ee_{B}}{n^{2}}\left(n_{\alpha}\partial_{y}n^{\mu}F_{B}^{\alpha y}-n_{y}^{\gamma}F_{B}^{\alpha \mu}\right)=K_{D}^{\mu}+M_{B_{p}}^{2}B_{D}^{\mu}B_{D\mu}$

2.

3,

Rotations $A_{\mu} = (\cos\phi - \epsilon ee_D \sin\phi) \overline{A}_{\mu} + (-\sin\phi - ee_B \cos\phi) \overline{A}_{\mu}^{D}$ $A_{D\mu} = \sin\phi \overline{A}_{\mu} + \cos\phi \overline{A}_{\mu}$ $B_{\mu} = \cos\phi \overline{B}_{\mu} - \sin\phi \overline{B}_{D\mu}$ $B_{D\mu} = (\sin\phi + \epsilon ee_B \cos\phi) \overline{B}_{\mu} + (\cos\phi - \epsilon ee_B \sin\phi) \overline{B}_{\mu}^{D}$

$$\begin{aligned} \Rightarrow & \text{in}_{\mu} \left(1^{\prime} \right) F_{A}^{\mu\nu} + \epsilon e e_{b} F_{A}^{\mu\nu} \\ &= \left(\cos \phi - \epsilon e e_{b} \sin \phi \right) F_{A}^{\mu\nu\nu} + \epsilon e e_{b} \left(\sin \phi F_{A}^{\mu\nu\nu} + \cos \phi F_{A}^{\mu\nu\nu} \right) \\ \left(- \sin \phi - \epsilon e e_{g} \left(\cos \phi \right) F_{B}^{\mu\nu\nu} \right) &= \cos \phi F_{A}^{\mu\nu\nu} + \sin \phi F_{A}^{\mu\nu\nu} \\ \left(1^{\prime\prime} \right) \frac{n_{d}}{n^{2}} \left[n^{\mu} \frac{\partial_{\gamma} \left(\cos \phi F_{A}^{\mu\nu\nu} + \sin \phi F_{B}^{\mu\nu\nu} \right) - n^{\nu} \frac{\partial_{\gamma} \left(\cos \phi F_{A}^{\mu\nu\nu} + \sin \phi F_{B}^{\mu\nu\nu} \right) \right] \\ - n_{\gamma} e^{\mu\nu\nu}}{p^{\mu}} \frac{\partial_{\gamma} \left(\cos \phi F_{B}^{\mu\nu} - \sin \phi F_{B}^{\mu\nu} \right) - n^{\nu} \frac{\partial_{\gamma} \left(\cos \phi F_{A}^{\mu\nu\nu} + \sin \phi F_{B}^{\mu\nu\nu} \right) \right] \\ \Rightarrow J_{\mu} = \cos \phi \overline{J_{\mu}} + \sin \phi \overline{J_{\mu}}^{\mu\nu} \\ - n_{\gamma} e^{\mu\nu\nu}}{p^{\mu}} \frac{\partial_{\gamma} \left(\cos \phi F_{B}^{\mu\nu} - \sin \phi F_{B}^{\mu\nu} \right) - \epsilon e e_{b} \left(\sin \phi + \epsilon e e_{b} \cos \phi \right) F_{B}^{\mu\nu\nu} \\ - e e e_{b} \left(\cos \phi - e e e_{b} \sin \phi \right) \overline{F_{B}}^{\mu\nu} - \epsilon e e_{b} \left(\sin \phi + e e e_{b} \cos \phi \right) \overline{F_{B}}^{\mu\nu} \right] \\ \left(2^{\mu} \right) \frac{n_{\mu}}{n^{2}} \left[n^{\mu} \frac{\partial_{\gamma} \left[(\cos \phi - e e e_{b} \sin \phi) \overline{F_{B}}^{\mu\nu} - (\sin \phi + e e e_{b} \cos \phi) \overline{F_{B}}^{\mu\nu} \right] \\ - n^{\nu} \frac{\partial_{\gamma} \left[(\cos \phi - e e e_{b} \sin \phi) \overline{F_{B}}^{\mu\nu} - (\sin \phi + e e e_{b} \cos \phi) \overline{F_{B}}^{\mu\nu} \right] \\ + n_{\gamma} e^{\mu\nu\mu}} \frac{\partial_{\gamma} \left[(\cos \phi - e e e_{b} \sin \phi) \overline{F_{B}}^{\mu\mu} - (\sin \phi + e e e_{b} \cos \phi) \overline{F_{B}}^{\mu\nu} \right] \\ = K^{\mu} \\ \Rightarrow K_{\mu} = \left(\cos \phi - e e e_{b} \sin \phi \right) \overline{K_{\mu}} - \left(\sin \phi + e e e_{b} \cos \phi \right) \overline{K_{\mu}} \right] \end{aligned}$$

4,

 $\rightarrow \ln(3')$ $F_{Ab}^{\alpha\gamma} + eeg F_{A}^{\alpha\gamma} =$ $\sin\phi \overline{F}_{A}^{xv} + \cos\phi \overline{F}_{A}^{xv} + \epsilon e e_{\delta} (\cos\phi - \epsilon e e_{\delta} \sin\phi) \overline{F}_{A}^{xv}$ $-\epsilon ee_{\phi}(\sin\phi + \epsilon ee_{\phi}\cos\phi) F_{A}^{\alpha\nu}$ $\cong (\sin\phi + \epsilon e e_{b} \cos\phi) \overline{F}_{A}^{*\nu} + (\cos\phi - \epsilon e e_{b} \sin\phi) \overline{F}_{A}^{*\nu}$ $(3'') \left| \frac{n_{x}}{n^{2}} n^{\mu} \partial_{\mu} \left[(\sin \phi + eee_{\mu} \cos \phi) F_{A}^{(\alpha)} + (\cos \phi - eee_{\mu} \sin \phi) F_{A}^{(\alpha)} \right] \right]$ $-n^{\nu}\partial_{\gamma} [(\sin\phi + eeep \cos\phi)F_{A} + (\cos\phi - eeep \sin\phi)F_{A} + f_{A} + (\cos\phi - eeep \sin\phi)F_{A} + f_{A} + f_{A}$ $-n_{y} \in \mathcal{A}_{p} = \partial_{y} \left[(\sin \phi + \epsilon e_{p} \cos \phi) F_{B}^{\gamma \beta} + (\cos \phi - \epsilon e_{p} \sin \phi) F_{Ap}^{\gamma \beta} \right]$ = Jm + MA (sin \$ Au + cos \$ Auo + 2sin\$ cos\$ Au Auo) $J_{p} = (\sin \phi + eeg \cos \phi) J_{p} + (\cos \phi - eeg \sin \phi) J_{p}$ in(4') $F_{BD}^{(1)} - \epsilon ee_{D}F_{B}^{(2)} =$ $(\sin\phi + eee_b\cos\phi) \overline{F}_{B}^{AV} + (\cos\phi - eee_b\sin\phi) \overline{F}_{B_b}^{AV}$ - Eeg cos \$ FR + Eeg sin \$ FBD = sin \$ FB + cos \$ FBD $(4'') \frac{n_{\alpha}}{n^{2}} \left[h^{\mu} \partial_{\nu} \left(\cos \phi F_{B_{D}}^{\alpha \nu} + \sin \phi F_{B}^{\alpha \nu} \right) - n^{\nu} \partial_{\nu} \left(\cos \phi F_{B_{D}}^{\alpha \mu} + \sin \phi F_{B}^{\alpha \mu} \right) \right]$ + NYE B DU (cos \$ FAD + sin\$ FAB)] = = $K_D^{A} + M_{B_D}^2 \left(\sin \phi B_{\mu}^2 + \cos \phi B_{D\mu}^2 + 2\cos \phi \sin \phi B_{D\mu} B^{A} \right)$ Kon= cos & Kon + sin & Kn

5,

1 Normalization of Lagrangian terms with visible and dark sectors - 12 August 2020

Our current Lagrangian. Until now, we have worked with Lagrangian terms collected from Eqs. (2.4), (4.1) and (4.2) (with $\theta = 0$) in Ref. [1] (or Eq. (2.1) in Ref. [2])

$$\mathcal{L} = - \frac{n^{\alpha}n^{\mu}}{8\pi n^{2}}g^{\beta\nu}g\left(F_{\alpha\beta}^{A}F_{\mu\nu}^{A} + F_{\alpha\beta}^{B}F_{\mu\nu}^{B}\right) + \frac{n^{\alpha}n_{\mu}}{16\pi n^{2}}\epsilon^{\mu\nu\gamma\delta}g\left(F_{\alpha\nu}^{B}F_{\gamma\delta}^{A} - F_{\alpha\nu}^{A}F_{\gamma\delta}^{B}\right) - J_{\mu}A^{\mu} - gK_{\mu}B^{\mu} - \frac{n^{\alpha}n^{\mu}}{8\pi n^{2}}g^{\beta\nu}g_{D}\left(F_{\alpha\beta}^{A}F_{\mu\nu}^{A} + F_{\alpha\beta}^{B}F_{\mu\nu}^{B}\right) + \frac{n^{\alpha}n_{\mu}}{16\pi n^{2}}\epsilon^{\mu\nu\gamma\delta}g_{D}\left(F_{\alpha\nu}^{B}F_{\gamma\delta}^{A} - F_{\alpha\nu}^{A}F_{\gamma\delta}^{B}\right) - e_{D}J_{\mu}A_{D}^{\mu} - e_{D}g_{D}K_{\mu}B_{D}^{\mu} - \frac{1}{2}M_{A_{D}}^{2}A_{D}^{2} - \frac{1}{2}M_{B_{D}}^{2}B_{D}^{2} + \epsilon e e_{D}\frac{n^{\alpha}n^{\mu}}{n^{2}}g^{\beta\nu}\left(F_{D\alpha\beta}^{A}F_{\mu\nu}^{A} - F_{D\alpha\beta}^{B}F_{\mu\nu}^{B}\right)$$
(1)

where $g = \frac{4\pi}{e^2}$, $g_D = \frac{4\pi}{e_D^2}$ and $F_{\mu\nu}^X = \partial_\mu X_\nu - \partial_\nu X_\mu$ (X = A, B, A_D, B_D). Note that I added an extra e_D here in the last term of line 2: $-e_D g_D K_\mu B_D^\mu$, to agree with Eq. (4.1) of Ref. [1]. The mass terms are from Eqs. (5.1) and (5.3) in Ref. [1] and Ref. [3].

We did not use them yet but it is not clear how the currents would be defined; we tentatively suggested

$$J_{\mu} = e \overline{\Psi}_{e} \gamma_{\mu} \Psi_{e}, \qquad K_{\mu} = g \overline{\Psi}_{g} \gamma_{\mu} \Psi_{g},$$

$$J_{D\mu} = e_{D} \overline{\Psi}_{De} \gamma_{\mu} \Psi_{De}, \qquad K_{D\mu} = g_{D} \overline{\Psi}_{Dg} \gamma_{\mu} \Psi_{Dg}, \qquad (2)$$

where Ψ_e and Ψ_g are some sort of projections on the electric and magnetic components.

Would this be a better normalization? The lack of symmetry with the charges in the interaction terms of Eq. (1) seems to have been remedied in p.6 of Ref. [4]. Also, in p. 4 (as well, see Eq. (3.9) or Ref. [1], they use the magnetic coupling $b = \frac{4\pi}{e}$ and likewise for b_D , rather than g. But I did not check the kinetic terms of Verhaaren's p. 6. It looks like Eq. (1) would be replaced with

$$\mathcal{L} = - \frac{n^{\alpha}n^{\mu}}{8\pi n^{2}}g^{\beta\nu}b\left(F_{\alpha\beta}^{A}F_{\mu\nu}^{A} + F_{\alpha\beta}^{B}F_{\mu\nu}^{B}\right) + \frac{n^{\alpha}n_{\mu}}{16\pi n^{2}}\epsilon^{\mu\nu\gamma\delta}b\left(F_{\alpha\nu}^{B}F_{\gamma\delta}^{A} - F_{\alpha\nu}^{A}F_{\gamma\delta}^{B}\right) - eJ_{\mu}A^{\mu} - bK_{\mu}B^{\mu} \\ - \frac{n^{\alpha}n^{\mu}}{8\pi n^{2}}g^{\beta\nu}b_{D}\left(F_{\alpha\beta}^{A}F_{\mu\nu}^{A} + F_{\alpha\beta}^{B}F_{\mu\nu}^{B}\right) + \frac{n^{\alpha}n_{\mu}}{16\pi n^{2}}\epsilon^{\mu\nu\gamma\delta}b_{D}\left(F_{\alpha\nu}^{B}F_{\gamma\delta}^{A} - F_{\alpha\nu}^{A}F_{\gamma\delta}^{B}\right) - e_{D}J_{\mu}A_{D}^{\mu} - b_{D}K_{\mu}B_{D}^{\mu} \\ - \frac{1}{2}M_{A_{D}}^{2}A_{D}^{2} - \frac{1}{2}M_{B_{D}}^{2}B_{D}^{2} + \epsilon e e_{D}\frac{n^{\alpha}n^{\mu}}{n^{2}}g^{\beta\nu}\left(F_{D\alpha\beta}^{A}F_{\mu\nu}^{A} - F_{D\alpha\beta}^{B}F_{\mu\nu}^{B}\right)$$
(3)

where $b = \frac{4\pi}{e}$, $b_D = \frac{4\pi}{e_D}$. It seems some terms in Eq. (3) are obtained by multiplying terms in Eq. (1) by e (but not all?!?) Now with the charges included explicitly in the interaction terms, should we use We did not use them yet but it is not clear how the currents would be defined; we tentatively suggested

$$J_{\mu} = \overline{\Psi}_{e} \gamma_{\mu} \Psi_{e}, \qquad K_{\mu} = \overline{\Psi}_{g} \gamma_{\mu} \Psi_{g},$$

$$J_{D\mu} = \overline{\Psi}_{De} \gamma_{\mu} \Psi_{De}, \qquad K_{D\mu} = \overline{\Psi}_{Dg} \gamma_{\mu} \Psi_{Dg}, \qquad (4)$$

rather than Eq. (2)?

2 Diagonalization of the kinetic term

The other equations that are not clear to me are the transformation that alledgedly lead to the diagonal basis \overline{V} (with $V = A, B, A_D, B_D$) of the kinetic term:

$$A_{\mu} = (\cos \phi + \epsilon e e_D \sin \phi) \overline{A}_{\mu} + (-\sin \phi + \epsilon e e_D \cos \phi) \overline{A}_{D\mu},$$

$$A_{D\mu} = \sin \phi \overline{A}_{\mu} + \cos \phi \overline{A}_{D\mu},$$

$$B_{\mu} = \cos \phi \overline{B}_{\mu} - \sin \phi \overline{B}_{D\mu},$$

$$B_{D\mu} = (\sin \phi - \epsilon e e_D \cos \phi) \overline{B}_{\mu} + (\cos \phi + \epsilon e e_D \sin \phi) \overline{B}_{D\mu},$$
(5)

along with the diagonal currents

$$e\overline{J}_{\mu} = (\cos\phi + \epsilon ee_{D}\sin\phi) eJ_{\mu} + \sin\phi e_{D}J_{D\mu},$$

$$e_{D}\overline{J}_{D\mu} = (-\sin\phi + \epsilon ee_{D}\cos\phi) eJ_{\mu} + \cos\phi e_{D}J_{D\mu},$$

$$\frac{1}{e}\overline{K}_{\mu} = \cos\phi \frac{1}{e}K_{\mu} + (\sin\phi - \epsilon ee_{D}\cos\phi) \frac{1}{e_{D}}K_{D\mu},$$

$$\frac{1}{e_{D}}\overline{K}_{D\mu} = -\sin\phi \frac{1}{e}K_{\mu} + (\cos\phi + \epsilon ee_{D}\sin\phi) \frac{1}{e_{D}}\overline{K}_{D\mu}.$$
(6)

By solving for the currents in the non-diagonal basis, I obtained (to be verified)

$$eJ_{\mu} = \cos\phi \ eJ_{\mu} - \sin\phi \ e_{D}J_{D\mu},$$

$$e_{D}J_{D\mu} = (\sin\phi - \epsilon ee_{D}\cos\phi) \ e\overline{J}_{\mu} + (\cos\phi + \epsilon ee_{D}\sin\phi) \ e_{D}\overline{J}_{D\mu},$$

$$\frac{1}{e}K_{\mu} = (\cos\phi + \epsilon ee_{D}\sin\phi) \ \frac{1}{e}\overline{K}_{\mu} + (-\sin\phi + \epsilon ee_{D}\cos\phi) \ \frac{1}{e_{D}}\overline{K}_{D\mu},$$

$$\frac{1}{e_{D}}K_{D\mu} = \sin\phi \ \frac{1}{e}\overline{K}_{\mu} + \cos\phi \ \frac{1}{e_{D}}\overline{K}_{D\mu}.$$
(7)

Now should Eqs. (5) and (6) diagonalize the mixing of the kinetic term in Eq. (1)? I tried –probably the wrong way– and it did not work. And how does one obtain Eqs. (5) and (6) anyways?

References

- J. Terning and C. B. Verhaaren, JHEP 12, 123 (2018) doi:10.1007/JHEP12(2018)123 [arXiv:1808.09459 [hep-th]].
- [2] J. Terning and C. B. Verhaaren, JHEP 12, 152 (2019) doi: 10.1007/JHEP12(2019)152 [arXiv: arXiv:1906.00014v2 [hep-ph]].
- [3] A. Hook and J. Huang, Phys. Rev. D 96, no.5, 055010 (2017) doi:10.1103/PhysRevD.96.055010
 [arXiv:1705.01107 [hep-ph]].
- [4] C. Verhaaren, Magnetic Monopole Dark Matter, Talk given at the International Conference on Neutrinos and Dark Matter (NDM-2020), Hurghada, Egypt, 13 January 2020.