



3rd MC-PAD Network Training Event,
Jožef Stefan Institute, Ljubljana, Slovenia
- 29 September 2010 -



Radiation Hardness of Semiconductor Detectors

- Radiation Effects and Detector Operation -

*... including an introduction to ongoing
radiation tolerant sensors developments*

Michael Moll (CERN/PH)



*Michael Moll – MC-PAD Network Training,
Ljubljana, 27.9.2010*





Outline



- **Motivation to develop radiation harder detectors**

- LHC upgrade and expected radiation levels
- Radiation induced degradation of detector performance

- **Radiation Damage in Silicon Detectors**

- Microscopic damage (crystal damage), NIEL
- Macroscopic damage (changes in detector properties)

- **Approaches to obtain radiation tolerant sensors**

- **Material Engineering**
 - New silicon materials – FZ, MCZ, DOFZ, EPI
 - Other semiconductors: Diamond, SiC, GaN,
- **Device Engineering**
 - p-in-n, n-in-n and n-in-p sensors, 3D sensors and thin devices

- **Silicon Sensors for the LHC upgrade**

..... *open questions and ongoing developments*

- **Summary**

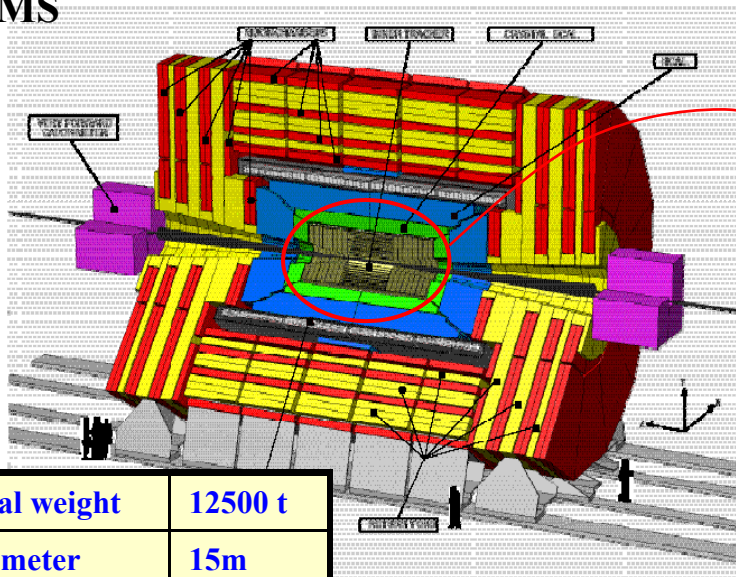
For details on
LHC detectors
see lecture of
Phil Allport

For details on
radiation induced defects
see lecture of
Ioana Pintilie

For **surface effects**
see lecture of
G.-F. Dalla Betta

For details on
sensor production
see lecture of
Manolo Lozano

■ CMS

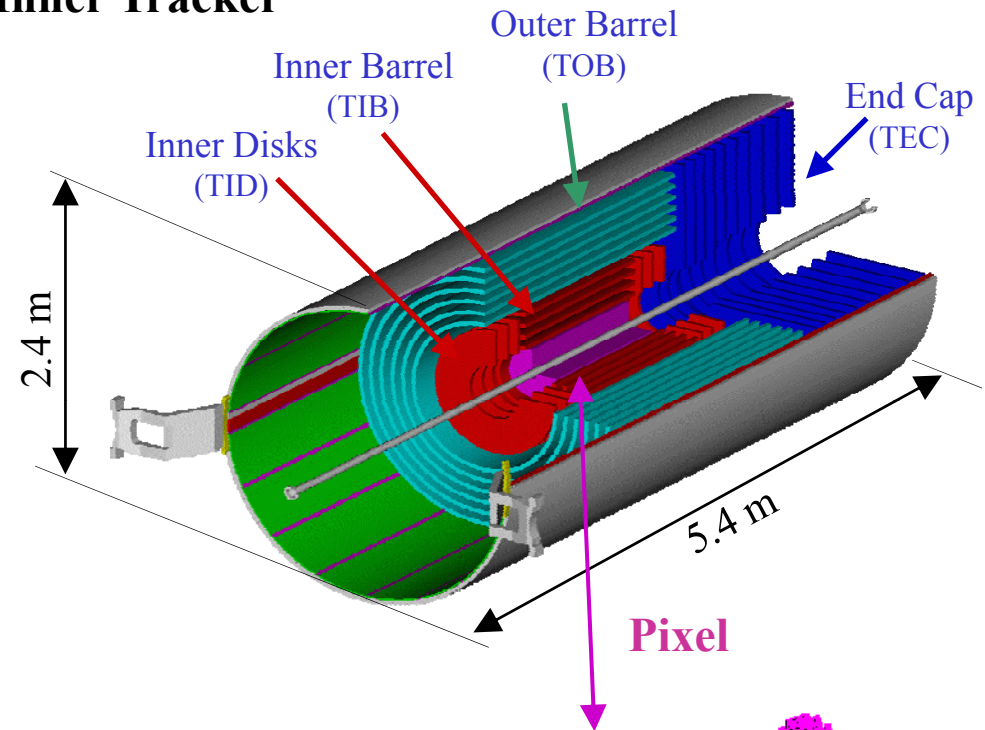


Total weight	12500 t
Diameter	15m
Length	21.6m
Magnetic field	4 T

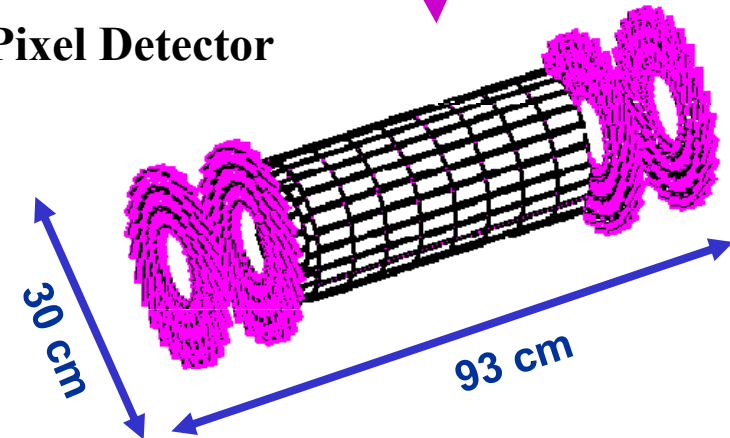
■ CMS – “Currently the Most Silicon”

- **Micro Strip:**
- ~ 214 m² of silicon strip sensors, 11.4 million strips
- **Pixel:**
- Inner 3 layers: silicon pixels (~ 1m²)
- 66 million pixels (100x150μm)
- Precision: $\sigma(r\phi) \sim \sigma(z) \sim 15\mu\text{m}$
- Most challenging operating environments (LHC)

■ Inner Tracker

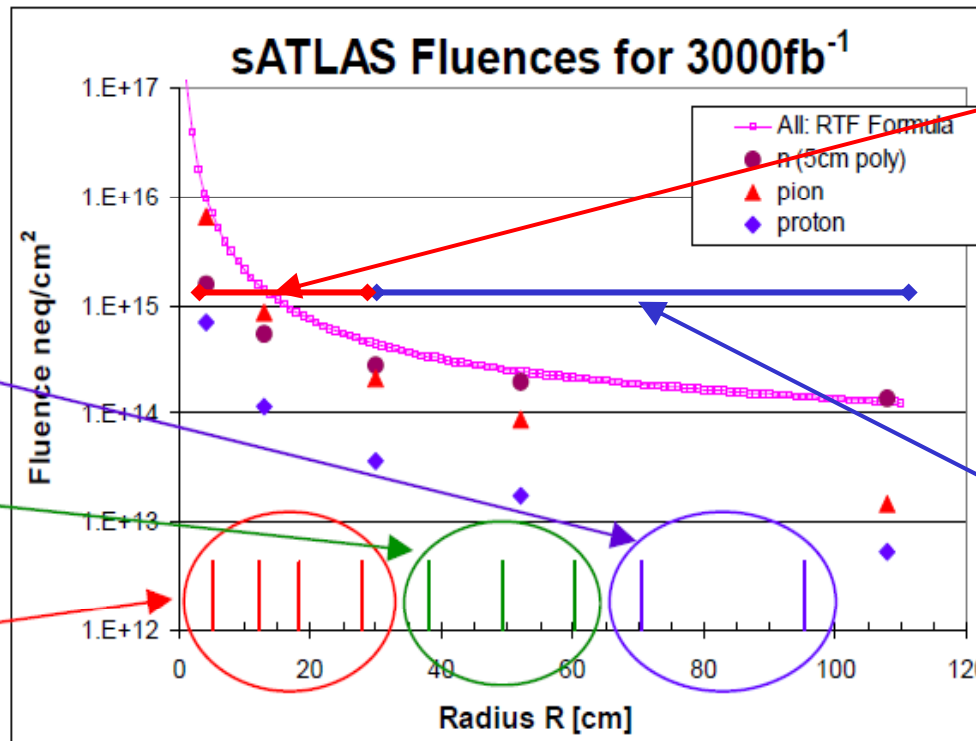


■ Pixel Detector



Radial distribution of sensors determined by Occupancy

- Long Strips (up to $4 \times 10^{14} \text{cm}^{-2}$)
- Short Strips (up to 10^{15}cm^{-2})
- Pixels (up to 10^{16}cm^{-2})



Dominated by pion damage

Dominated by neutron damage

ATLAS Radiation Taskforce [ATL-GEN-2005-01] & H.Sadrozinski [IEEE NSS 2007]

- **Radiation hardness requirements (including safety factor of 2)**

- $2 \times 10^{16} \text{ n}_{\text{eq}}/\text{cm}^2$ for the innermost pixel layers
- $7 \times 10^{14} \text{ n}_{\text{eq}}/\text{cm}^2$ for the innermost strip layers

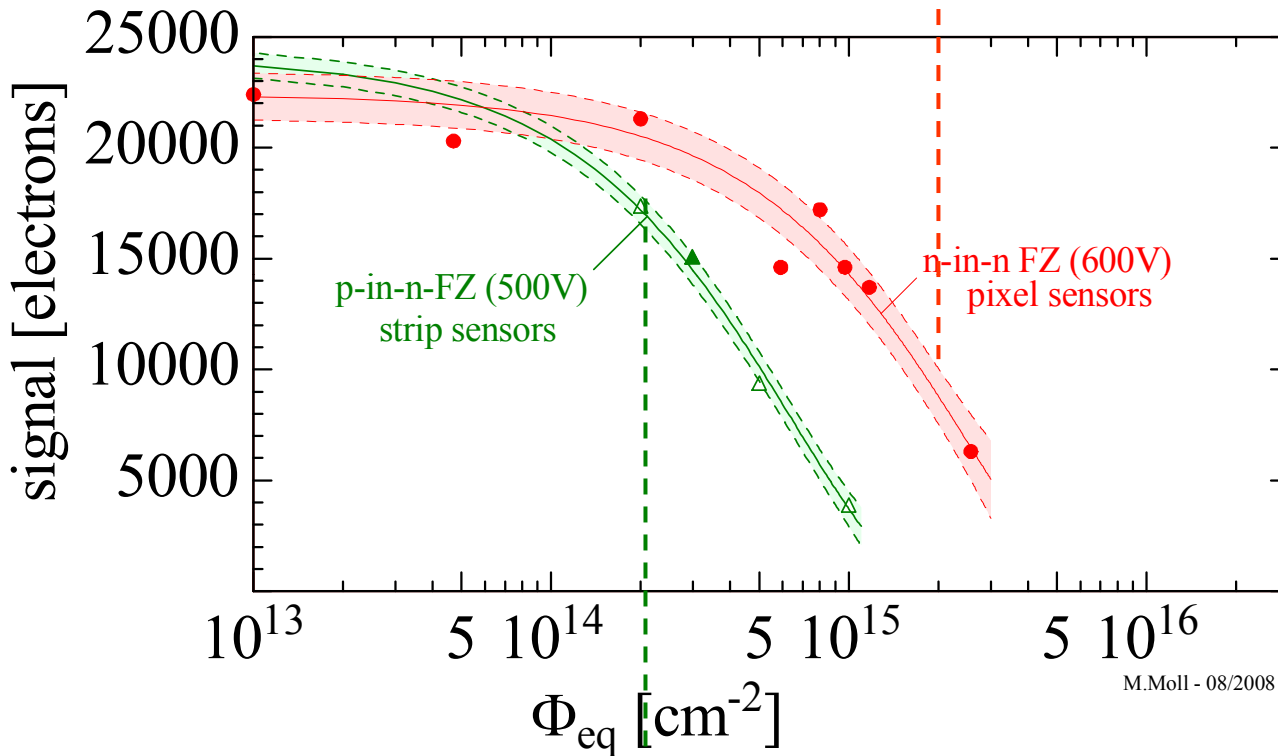
B-layer (R=3.7 cm):	$2.5 \times 10^{16} \text{ n}_{\text{eq}}/\text{cm}^2 = 1140 \text{ Mrad}$
2 nd Inner Pixel Layer (R=7 cm):	$7.8 \times 10^{15} \text{ n}_{\text{eq}}/\text{cm}^2 = 420 \text{ Mrad}$
1 st Outer Pixel Layer (R=11 cm):	$3.6 \times 10^{15} \text{ n}_{\text{eq}}/\text{cm}^2 = 207 \text{ Mrad}$
Short strips (R=38 cm):	$6.8 \times 10^{14} \text{ n}_{\text{eq}}/\text{cm}^2 = 30 \text{ Mrad}$
Long strips (R=85 cm):	$3.2 \times 10^{14} \text{ n}_{\text{eq}}/\text{cm}^2 = 8.4 \text{ Mrad}$



Signal degradation for LHC Silicon Sensors



Pixel sensors:
max. cumulated fluence for **LHC**



Note: Measured partly under different conditions!
Lines to guide the eye (no modeling)!

FZ Silicon
Strip and Pixel Sensors

- n-in-n (FZ), 285 μ m, 600V, 23 GeV p
- ▲ p-in-n (FZ), 300 μ m, 500V, 23GeV p
- △ p-in-n (FZ), 300 μ m, 500V, neutrons

References:

- [1] p/n-FZ, 300 μ m, (-30°C, 25ns), strip [Casse 2008]
- [2] n/n-FZ, 285 μ m, (-10°C, 40ns), pixel [Rohe et al. 2005]

M.Moll - 08/2008

Strip sensors:
max. cumulated fluence for **LHC**

Situation in 2005



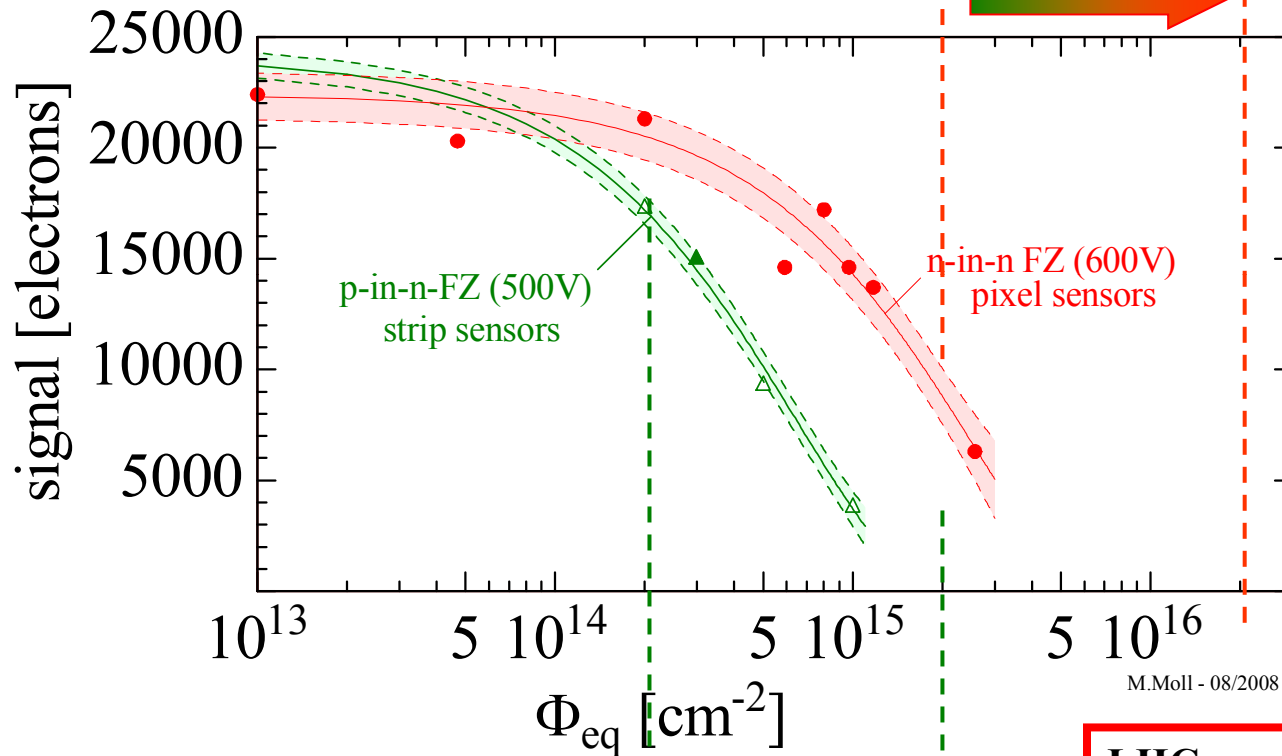
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Pixel sensors:

max. cumulated fluence for **LHC** and **LHC upgrade**

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FZ Silicon Strip and Pixel Sensors

- n-in-n (FZ), 285 μm , 600V, 23 GeV p
- ▲ p-in-n (FZ), 300 μm , 500V, 23 GeV p
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- [1] p/n-FZ, 300 μm , (-30°C, 25ns), strip [Casse 2008]
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Strip sensors:

max. cumulated fluence for **LHC** and **LHC upgrade**

LHC upgrade will need more radiation tolerant tracking detector concepts!

*Boundary conditions & other challenges:
Granularity, Powering, Cooling, Connectivity,
Triggering, Low mass, Low cost!*



Outline

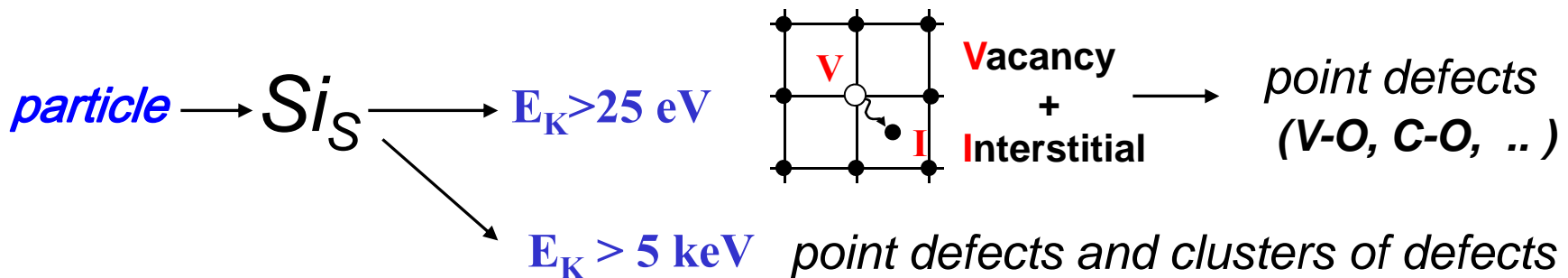
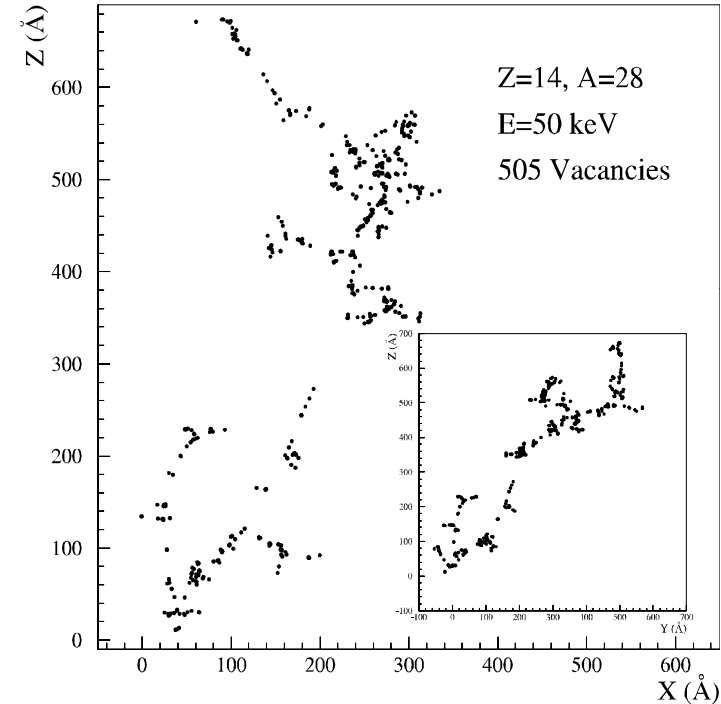
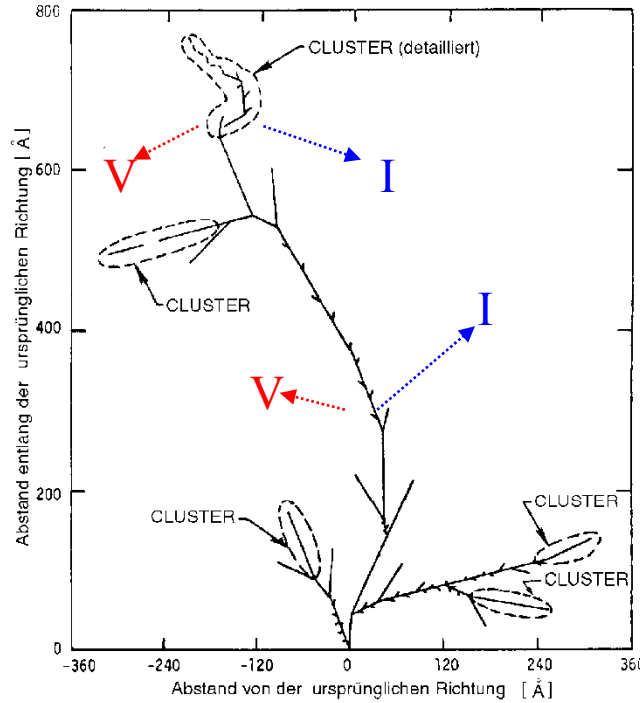


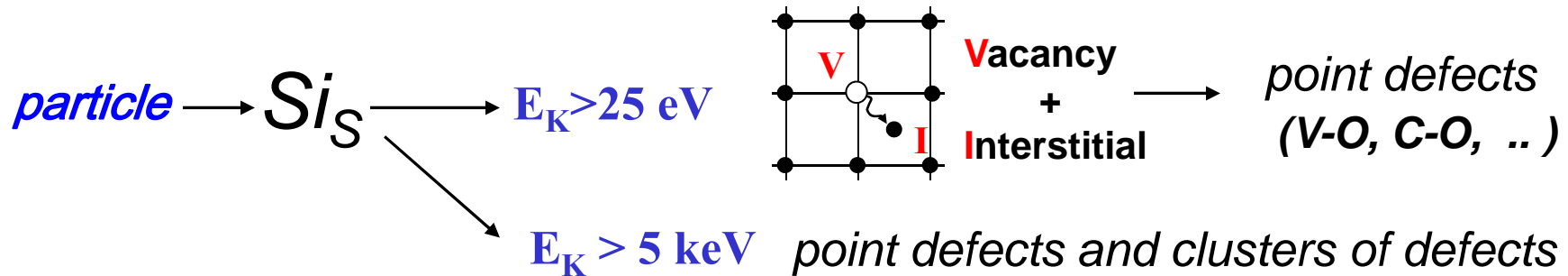
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 - **Microscopic damage (crystal damage), NIEL**
 - **Macroscopic damage (changes in detector properties)**
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 - **Material Engineering**
 - New silicon materials – FZ, MCZ, DOFZ, EPI
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..... open questions and ongoing developments
- **Summary**

- ◆ Spatial distribution of vacancies created by a 50 keV Si-ion in silicon.
(typical recoil energy for 1 MeV neutrons)

M.Huhtinen 2001

van Lint 1980





• **^{60}Co -gammas**

- Compton Electrons with max. $E_\gamma \approx 1 \text{ MeV}$ (no cluster production)

• **Electrons**

- $E_e > 255 \text{ keV}$ for displacement
- $E_e > 8 \text{ MeV}$ for cluster

• **Neutrons (elastic scattering)**

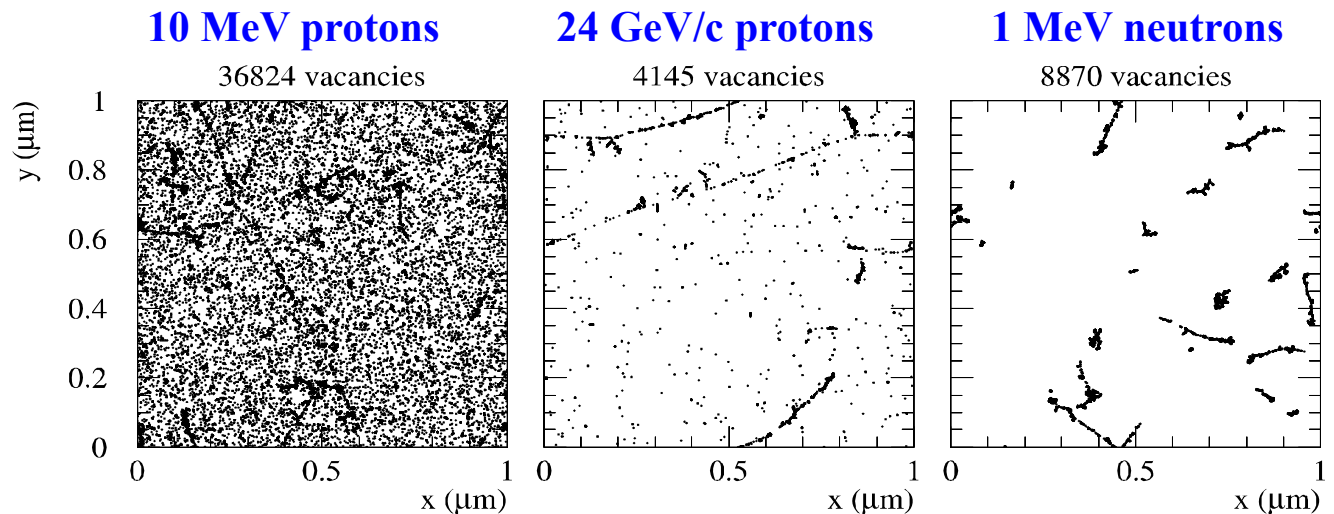
- $E_n > 185 \text{ eV}$ for displacement
- $E_n > 35 \text{ keV}$ for cluster

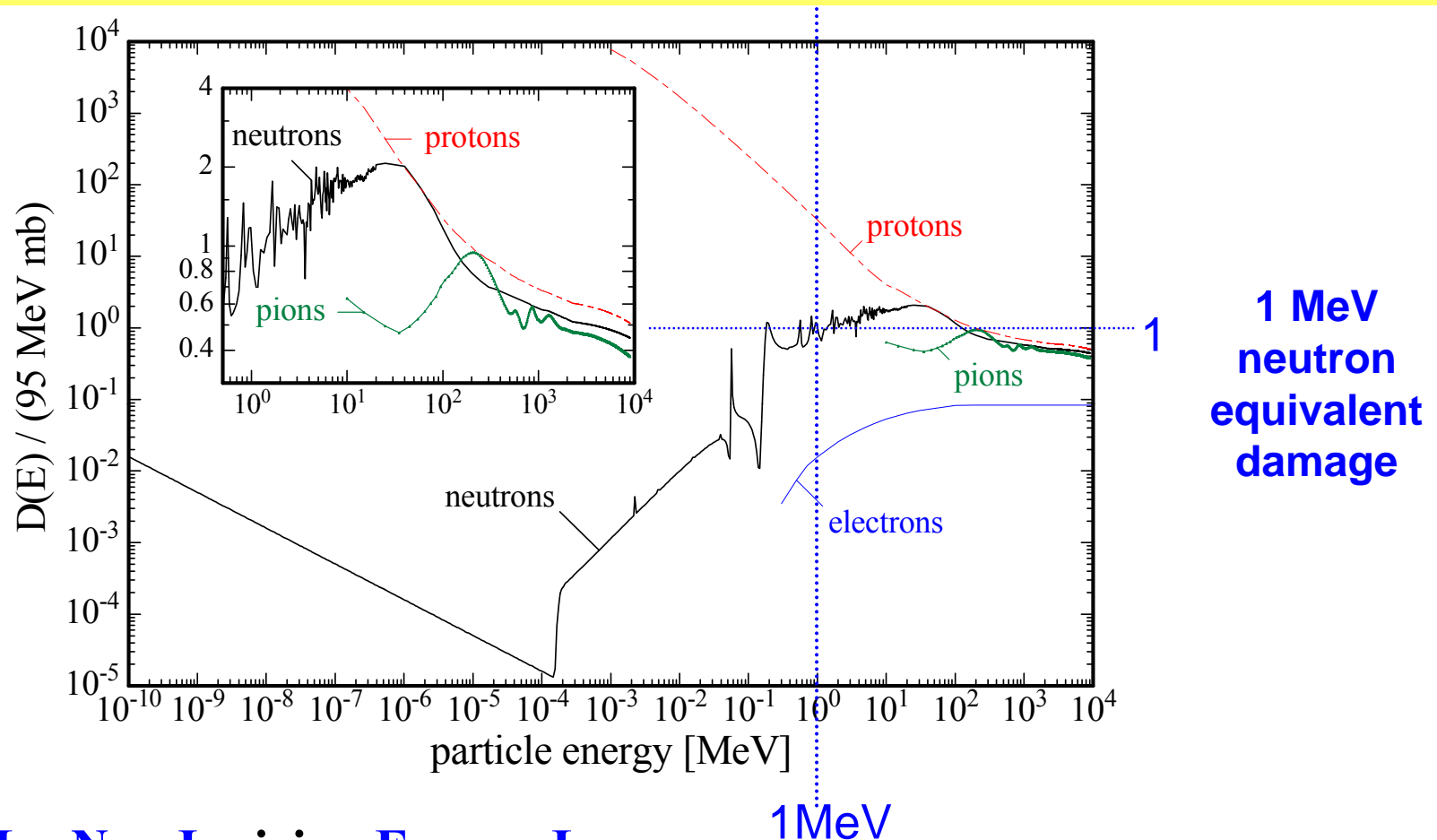
Only point defects \longleftrightarrow **point defects & clusters** \longleftrightarrow **Mainly clusters**

Simulation:

Initial distribution of vacancies in $(1\mu\text{m})^3$ after $10^{14} \text{ particles/cm}^2$

[Mika Huhtinen NIMA 491(2002) 194]





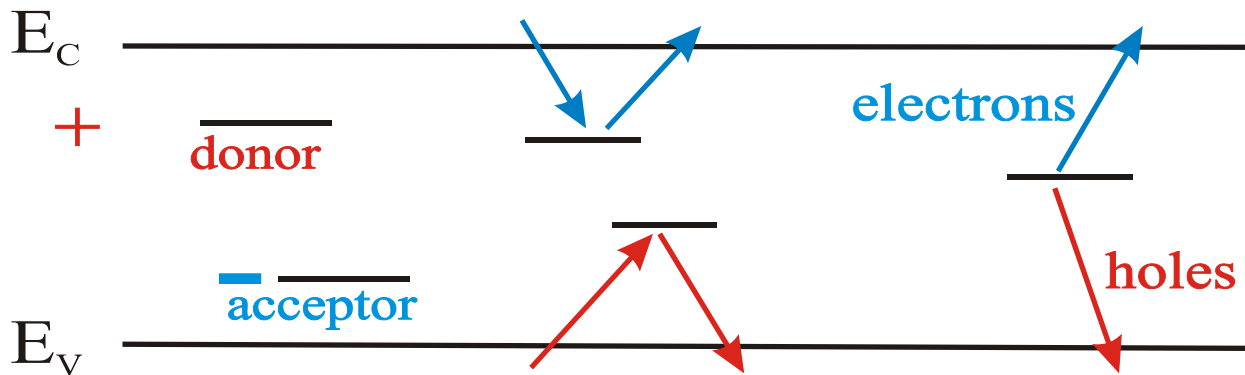
- **NIEL - Non Ionizing Energy Loss**
- **NIEL - Hypothesis: Damage parameters scale with the NIEL**
 - *Be careful, does not hold for all particles & damage parameters (see later)*



Impact of Defects on Detector Properties



Shockley-Read-Hall statistics



Details in next lecture by Ioana Pintilie on defects.

charged defects
 $\Rightarrow N_{\text{eff}}, V_{\text{dep}}$
 e.g. donors in upper
 and acceptors in
 lower half of band
 gap

Trapping (e and h)
 $\Rightarrow \text{CCE}$
 shallow defects do not
 contribute at room
 temperature due to fast
 detrapping

generation
 \Rightarrow **leakage current**
 Levels close to
 midgap
 most effective

Impact on detector properties can be calculated if all defect parameters are known:

$\sigma_{n,p}$: cross sections

ΔE : ionization energy

N_t : concentration



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- Detectors are basically **p⁺-n diodes made on high resistivity silicon**

- **Standard detector grade silicon:**

- **Float Zone silicon (FZ)**

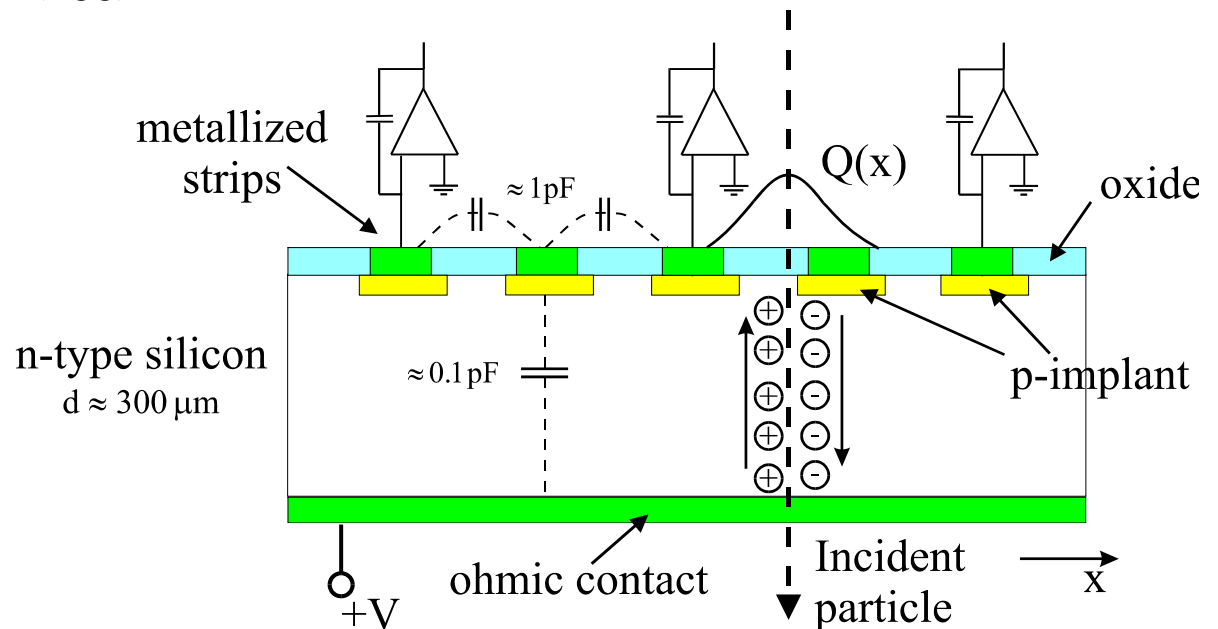
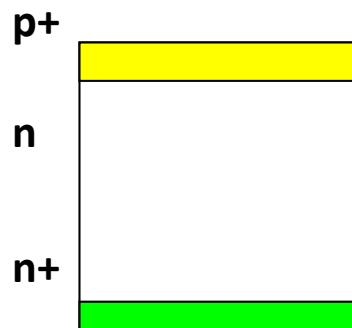
- **n-type: 2...20 KΩcm**

- **[P] = 20...2×10¹¹ cm⁻³ (very low concentration !! below 1ppba = 5×10¹³cm⁻³)**

- **[O] ≈ several 10¹⁵cm⁻³**

- **[C] ≈ some 10¹⁵cm⁻³, usually [C] < [O]**

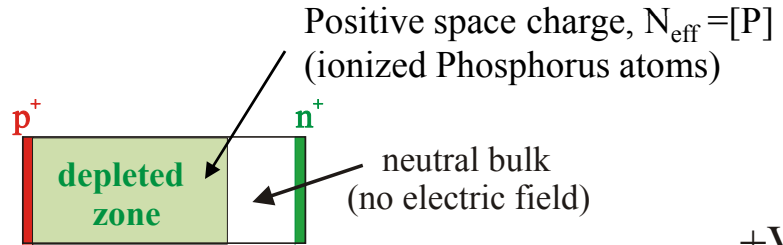
- **crystal orientation: <111> or <100>**



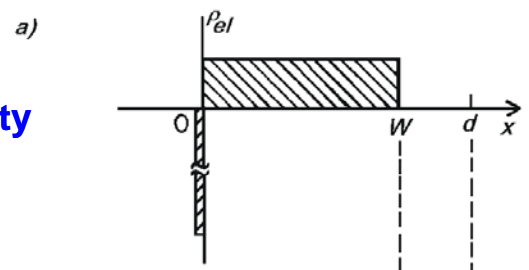
Reminder: Reverse biased abrupt p⁺-n junction

Poisson's equation

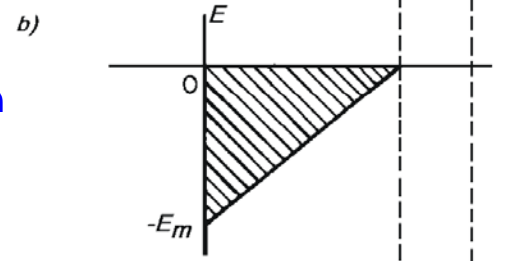
$$-\frac{d^2}{dx^2} \phi(x) = \frac{q_0}{\epsilon\epsilon_0} \cdot N_{eff}$$



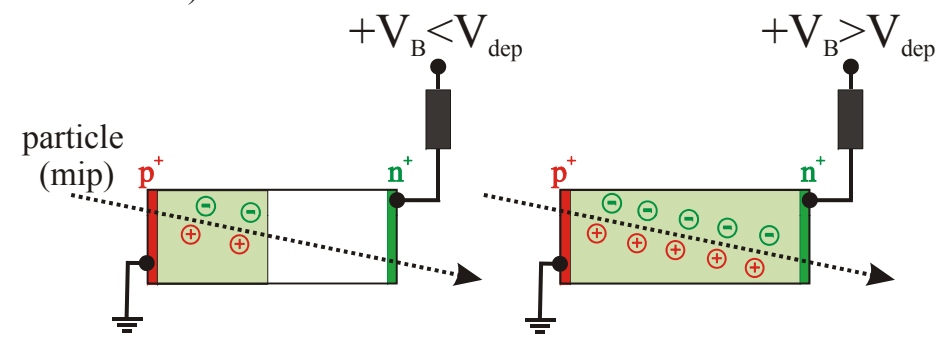
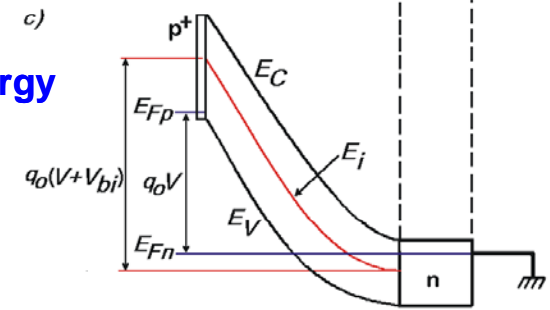
Electrical charge density



Electrical field strength



Electron potential energy



Full charge collection only for $V_B > V_{dep}$!

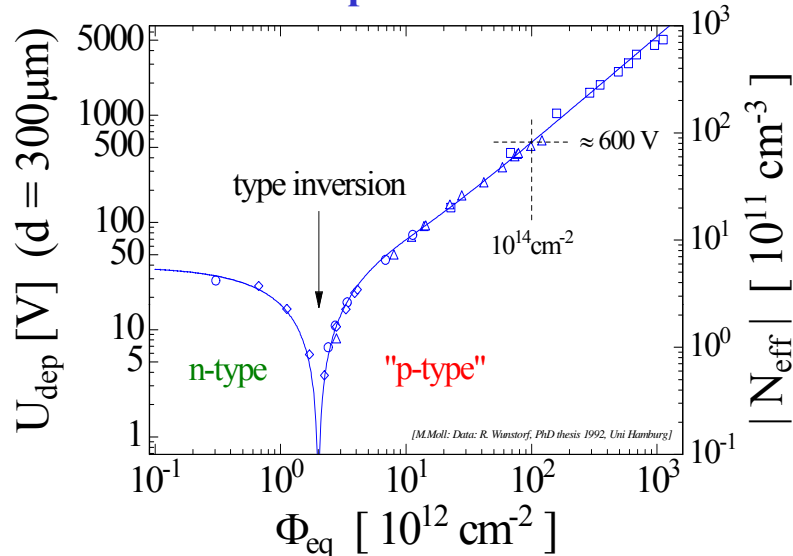
depletion voltage

$$V_{dep} = \frac{q_0}{\epsilon\epsilon_0} \cdot |N_{eff}| \cdot d^2$$

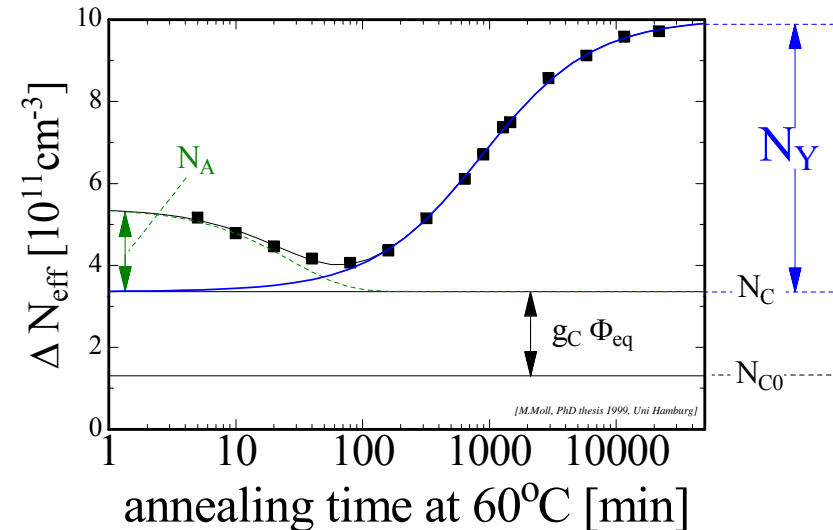
effective space charge density

Change of Depletion Voltage V_{dep} (N_{eff})

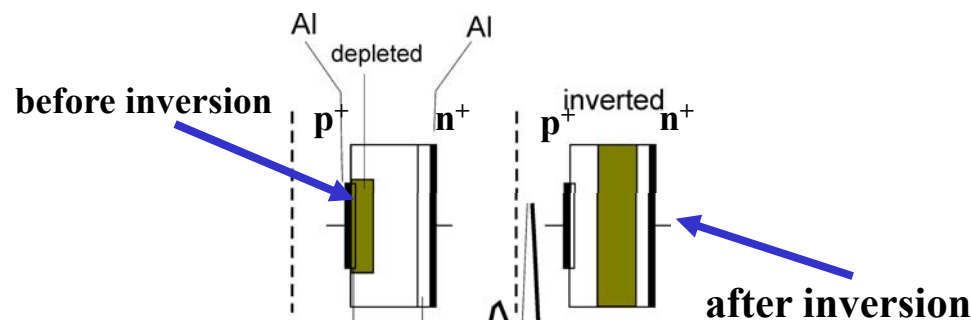
.... with particle fluence:



.... with time (annealing):



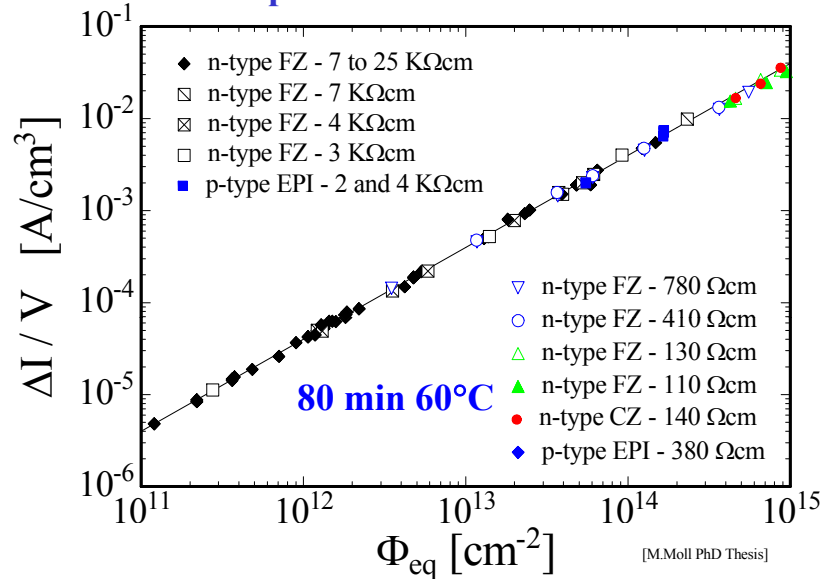
- “**Type inversion**”: N_{eff} changes from positive to negative (Space Charge Sign Inversion)



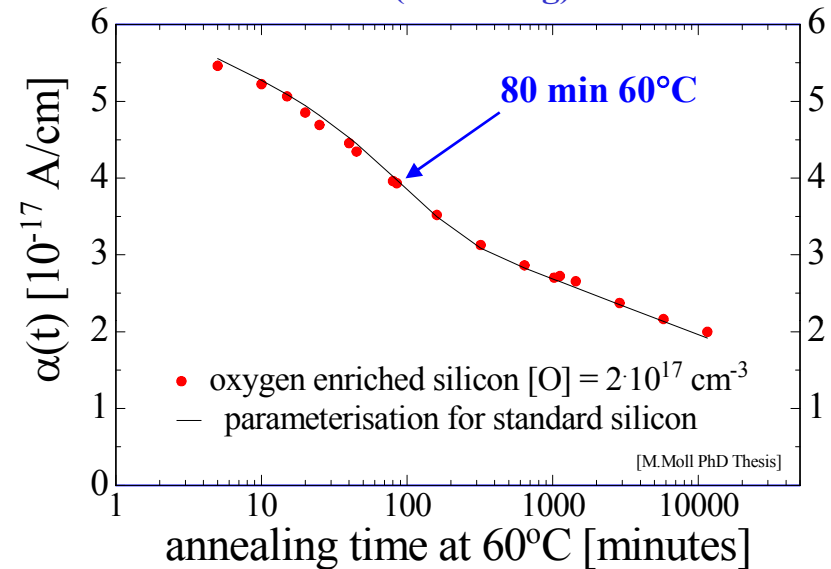
- Short term: “**Beneficial annealing**”
- Long term: “**Reverse annealing**”
- time constant depends on temperature:
 - ~ 500 years (-10°C)
 - ~ 500 days (20°C)
 - ~ 21 hours (60°C)
- Consequence: **Detectors must be cooled even when the experiment is not running!**

Change of Leakage Current (after hadron irradiation)

.... with particle fluence:



.... with time (annealing):



- Damage parameter α (slope in figure)

$$\alpha = \frac{\Delta I}{V \cdot \Phi_{eq}}$$

Leakage current
per unit volume
and particle fluence

- α is constant over several orders of fluence and independent of impurity concentration in Si
⇒ can be used for fluence measurement

- Leakage current decreasing in time (depending on temperature)
- Strong temperature dependence

$$I \propto \exp\left(-\frac{E_g}{2k_B T}\right)$$

Consequence:

Cool detectors during operation!
Example: $I(-10^\circ\text{C}) \sim 1/16 I(20^\circ\text{C})$



Radiation Damage – III. CCE (Trapping)

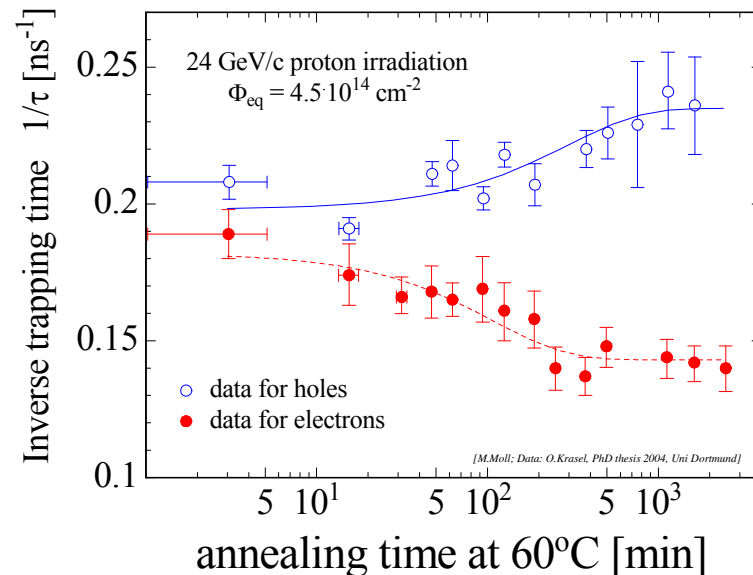
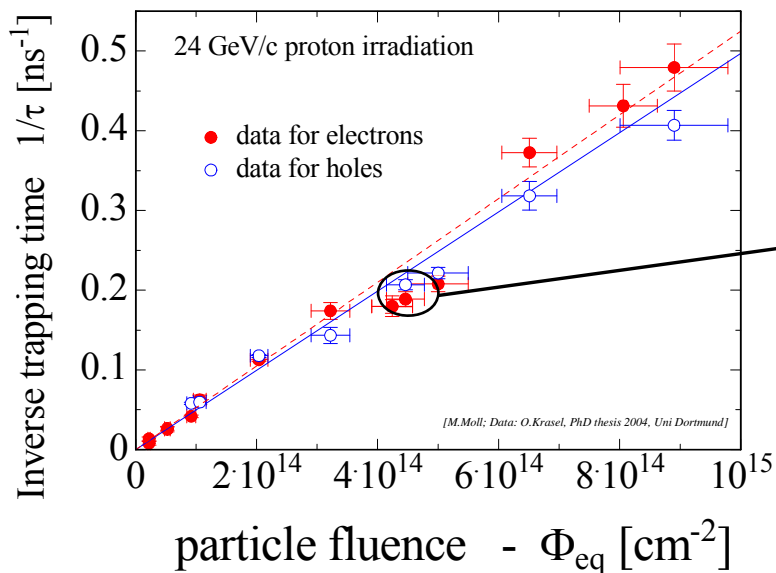


■ Deterioration of Charge Collection Efficiency (CCE) by trapping

Trapping is characterized by an effective trapping time τ_{eff} for electrons and holes:

$$Q_{e,h}(t) = Q_{0,e,h} \exp\left(-\frac{1}{\tau_{\text{eff } e,h}} \cdot t\right) \quad \text{where} \quad \frac{1}{\tau_{\text{eff } e,h}} \propto N_{\text{defects}}$$

Increase of inverse trapping time ($1/\tau$) with fluence and change with time (annealing):





Summary: Radiation Damage in Silicon Sensors



■ Two general types of radiation damage to the detector materials:

● Bulk (Crystal) damage due to Non Ionizing Energy Loss (NIEL)

- displacement damage, built up of crystal defects –

Influenced
by impurities
in Si – Defect
Engineering
is possible!

I. Change of **effective doping concentration** (higher depletion voltage, under- depletion)

II. Increase of **leakage current** (increase of shot noise, thermal runaway)

III. Increase of **charge carrier trapping** (loss of charge)

Same for
all tested
Silicon

materials!

● Surface damage due to Ionizing Energy Loss (IEL)

- accumulation of positive in the oxide (SiO_2) and the Si/ SiO_2 interface –
affects: interstrip capacitance (noise factor), breakdown behavior, ...

■ Impact on detector performance and Charge Collection Efficiency (depending on detector type and geometry and readout electronics!)

Signal/noise ratio is the quantity to watch

⇒ Sensors can fail from radiation damage !

Can be
optimized!

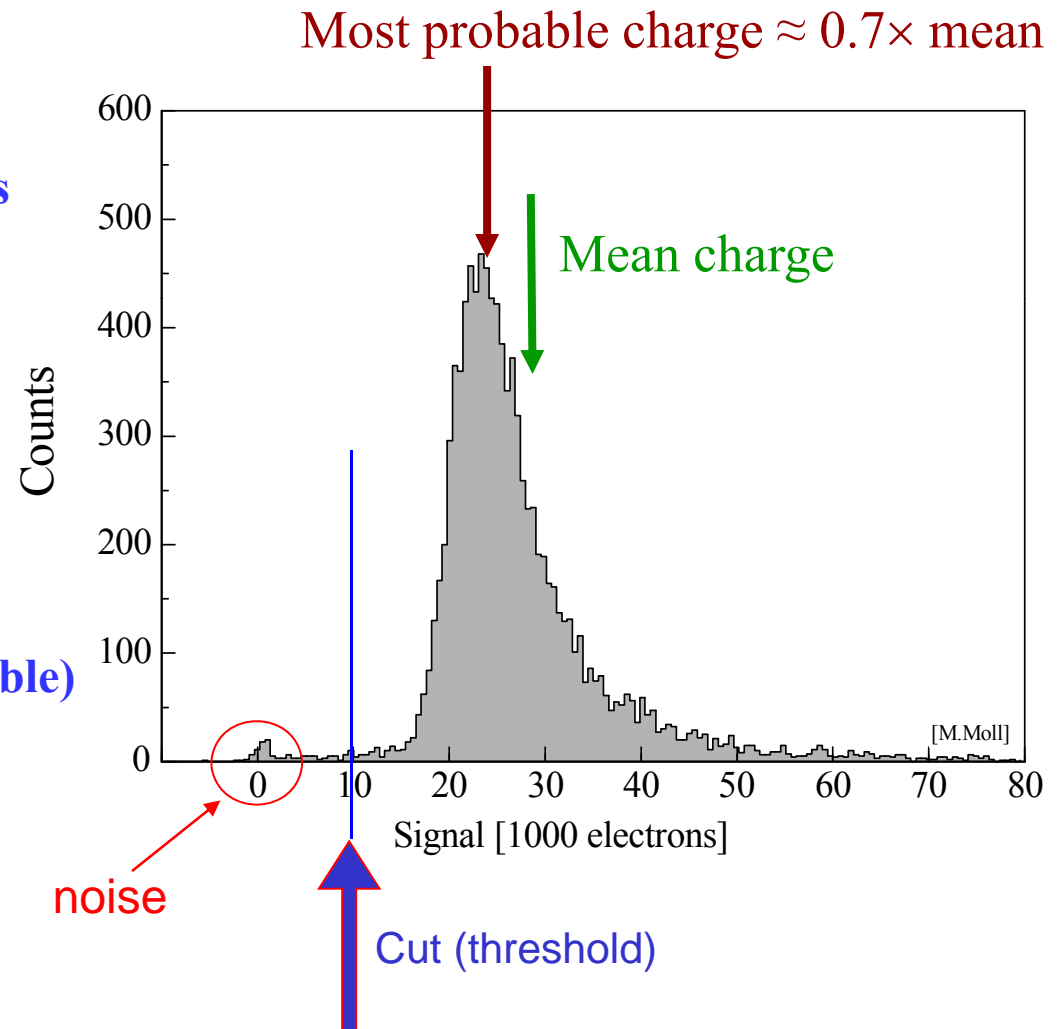


The charge signal



■ Collected Charge for a Minimum Ionizing Particle (MIP)

- Mean energy loss
 dE/dx (Si) = 3.88 MeV/cm
 \Rightarrow 116 keV for 300 μ m thickness
- Most probable energy loss
 $\approx 0.7 \times$ mean
 \Rightarrow 81 keV
- 3.6 eV to create an e-h pair
 \Rightarrow 72 e-h / μ m (mean)
 \Rightarrow 108 e-h / μ m (most probable)
- Most probable charge (300 μ m)
 ≈ 22500 e ≈ 3.6 fC





Signal to Noise ratio



- Landau distribution has a low energy tail
 - becomes even lower by noise broadening
- Noise sources: (ENC = Equivalent Noise Charge)

- Capacitance $ENC \propto C_d$

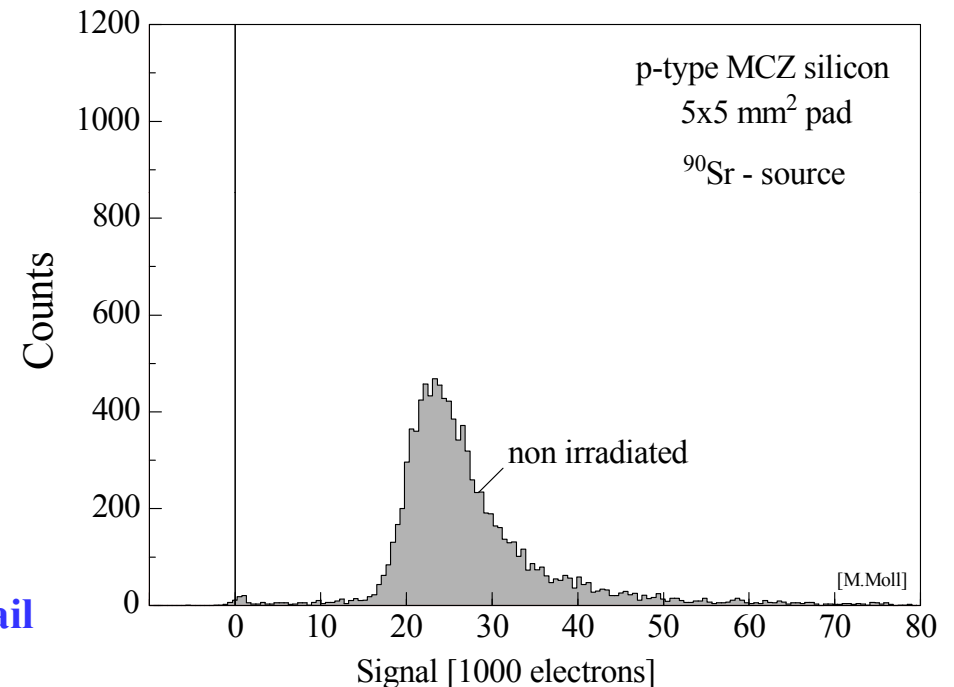
- Leakage Current $ENC \propto \sqrt{I}$

- Thermal Noise (bias resistor) $ENC \propto \sqrt{k_B T / R}$

- Good hits selected by requiring NADC > noise tail
 - If cut too high \Rightarrow efficiency loss
 - If cut too low \Rightarrow noise occupancy

■ Figure of Merit: Signal-to-Noise Ratio S/N

- Typical values >10-15, people get nervous below 10.
Radiation damage severely degrades the S/N.





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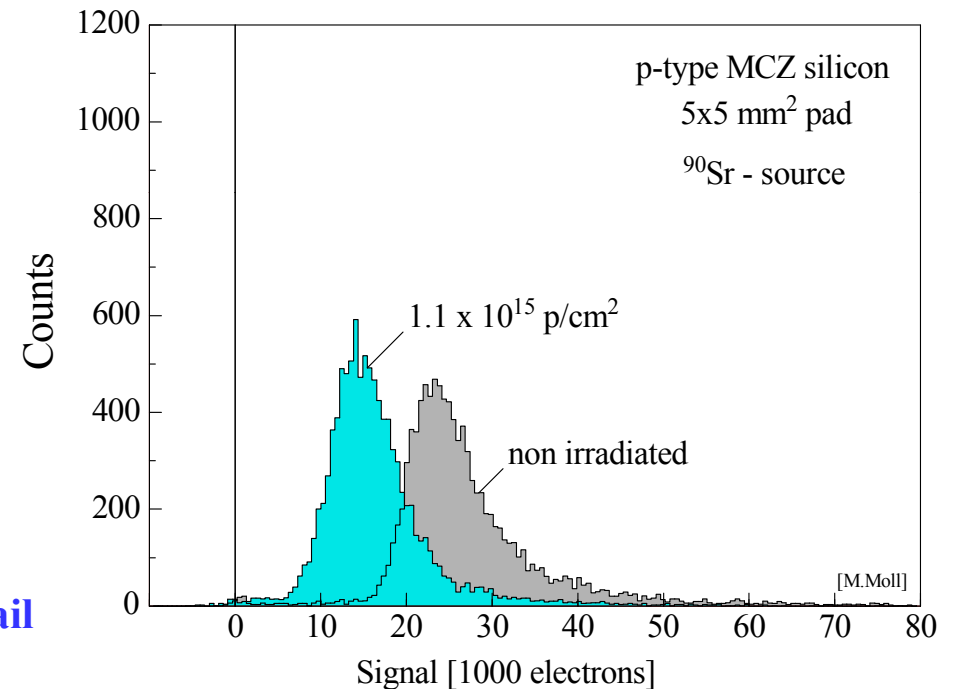
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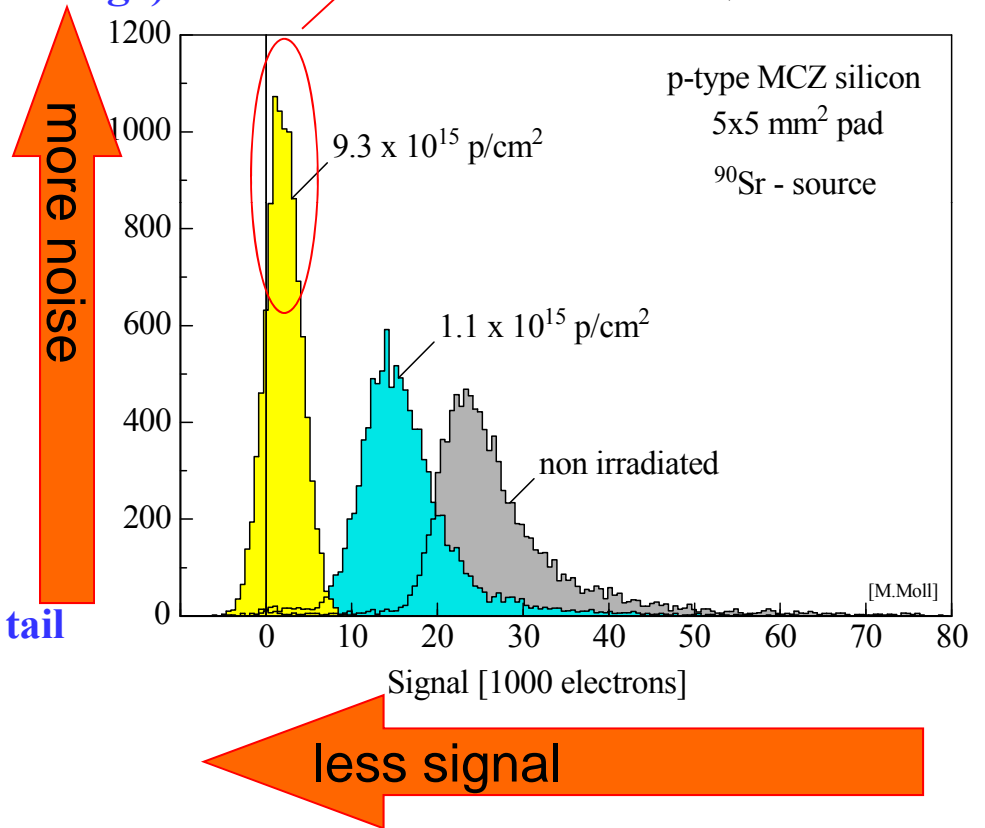
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What is signal and what is noise?





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Approaches to develop radiation harder solid state tracking detectors



Scientific strategies:

- I. Material engineering
- II. Device engineering
- III. Change of detector operational conditions

CERN-RD39
“Cryogenic Tracking Detectors”
operation at 100-200K
to reduce charge loss

- Defect Engineering of Silicon

Deliberate incorporation of impurities or defects into the silicon bulk to improve radiation tolerance of detectors

- **Needs:** Profound understanding of radiation damage
 - microscopic defects, macroscopic parameters
 - dependence on particle type and energy
 - defect formation kinetics and annealing
- **Examples:**
 - Oxygen rich Silicon (DOFZ, Cz, MCZ, EPI)
 - Oxygen dimer & hydrogen enriched Si
 - Pre-irradiated Si
 - Influence of processing technology

- New Materials

- Silicon Carbide (SiC), Gallium Nitride (GaN)
- Diamond (CERN RD42 Collaboration)
- Amorphous silicon

- Device Engineering (New Detector Designs)

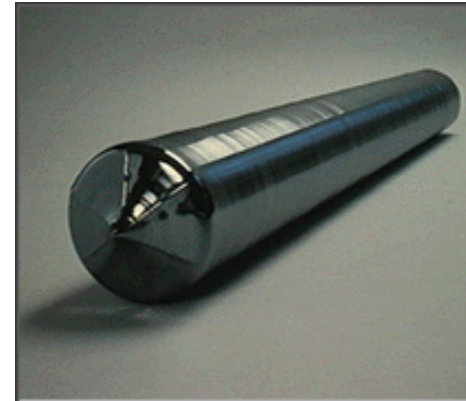
- p-type silicon detectors (n-in-p)
- thin detectors, epitaxial detectors
- 3D detectors and Semi 3D detectors, Stripixels
- Cost effective detectors
- Monolithic devices

■ Float Zone process

- Using a single Si crystal seed, melt the vertically oriented rod onto the seed using RF power and “pull” the **monocrystalline ingot**



■ Mono-crystalline Ingot



■ Wafer production

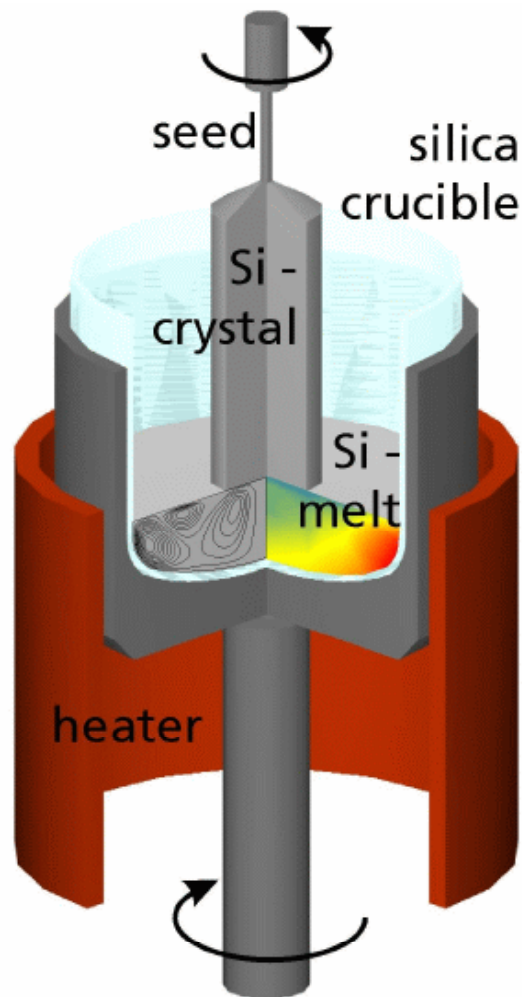
- Slicing, lapping, etching, polishing



■ Oxygen enrichment (DOFZ)

- Oxidation of wafer at high temperatures

■ Czochralski silicon



- Pull Si-crystal from a Si-melt contained in a silica crucible while rotating.
- Silica crucible is dissolving oxygen into the melt ⇒ **high concentration of O in CZ**
- Material used by IC industry (cheap)
- Recent developments (~5 years) made CZ available in sufficiently high purity (resistivity) to allow for use as particle detector.

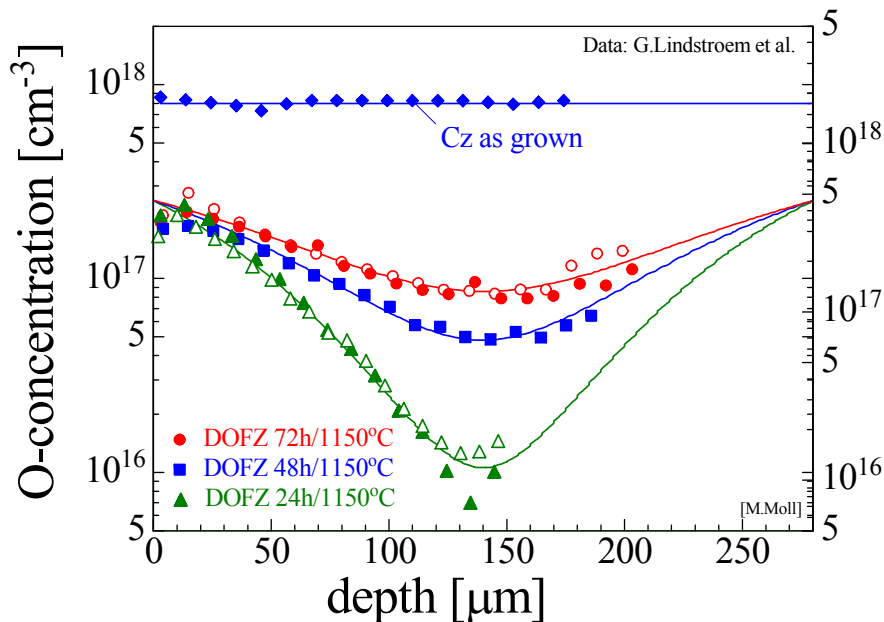


■ Epitaxial silicon

- Chemical-Vapor Deposition (CVD) of Silicon
- CZ silicon substrate used ⇒ **in-diffusion of oxygen**
- growth rate about 1 μm/min
- excellent homogeneity of resistivity
- up to 150 μm thick layers produced (thicker is possible)
- price depending on thickness of epi-layer but not extending ~ 3 x price of FZ wafer

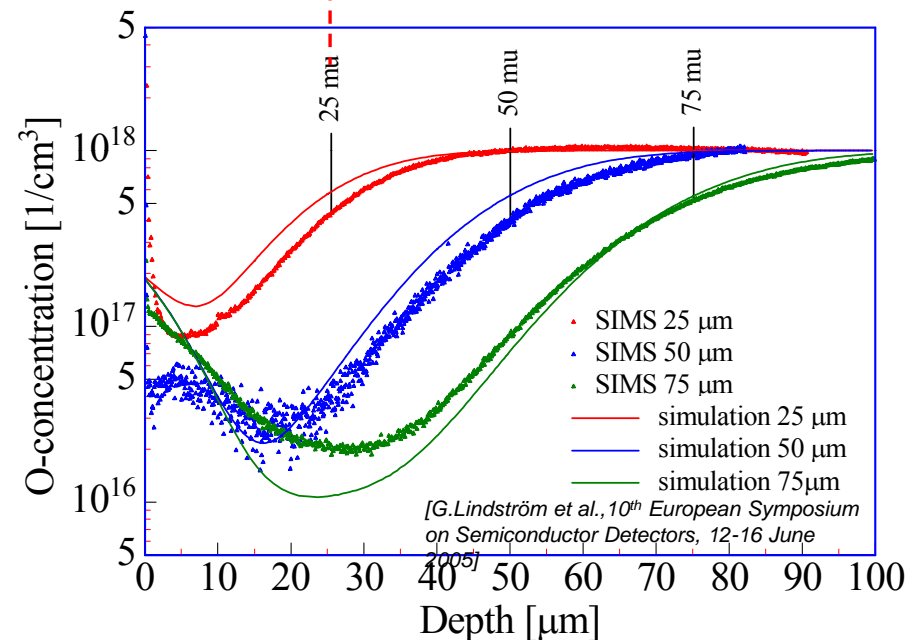
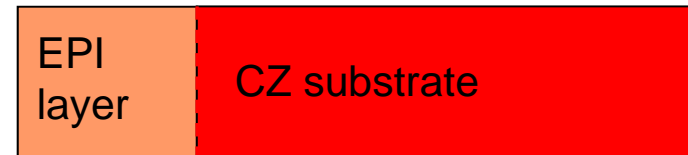
DOFZ and CZ silicon

- DOFZ: inhomogeneous oxygen distribution
- DOFZ: oxygen content increasing with time at high temperature



- CZ: high O_i (oxygen) and O_{2i} (oxygen dimer) concentration (homogeneous)
- CZ: formation of Thermal Donors possible !

Epitaxial silicon



- EPI: O_i and O_{2i} (?) diffusion from substrate into epi-layer during production
- EPI: in-homogeneous oxygen distribution



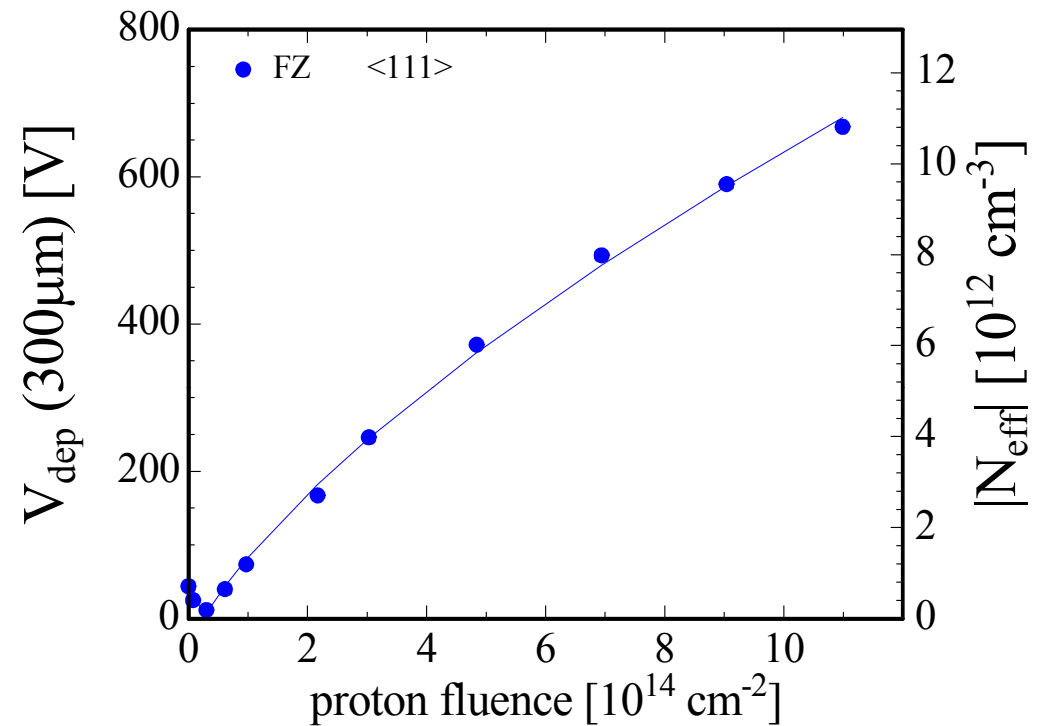
Standard FZ, DOFZ, MCz and Cz silicon



24 GeV/c proton irradiation

• Standard FZ silicon

- type inversion at $\sim 2 \times 10^{13}$ p/cm²
- strong N_{eff} increase at high fluence





Standard FZ, DOFZ, MCz and Cz silicon



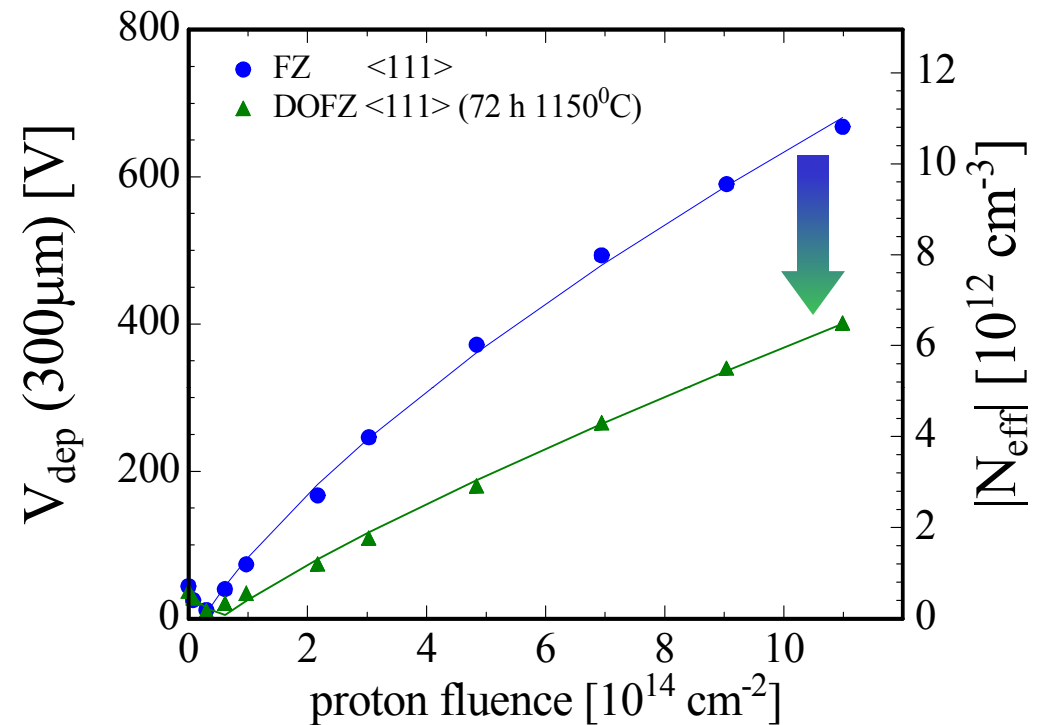
24 GeV/c proton irradiation

• Standard FZ silicon

- type inversion at $\sim 2 \times 10^{13}$ p/cm²
- strong N_{eff} increase at high fluence

• Oxygenated FZ (DOFZ)

- type inversion at $\sim 2 \times 10^{13}$ p/cm²
- reduced N_{eff} increase at high fluence



24 GeV/c proton irradiation

• Standard FZ silicon

- type inversion at $\sim 2 \times 10^{13}$ p/cm²
- strong N_{eff} increase at high fluence

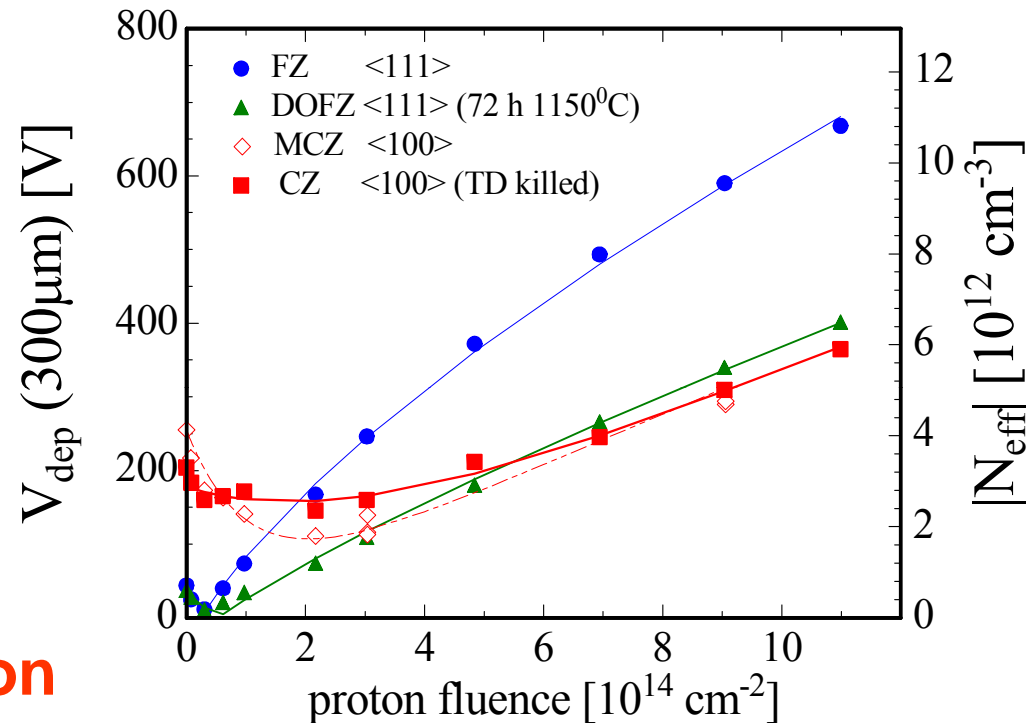
• Oxygenated FZ (DOFZ)

- type inversion at $\sim 2 \times 10^{13}$ p/cm²
- reduced N_{eff} increase at high fluence

• CZ silicon and MCZ silicon

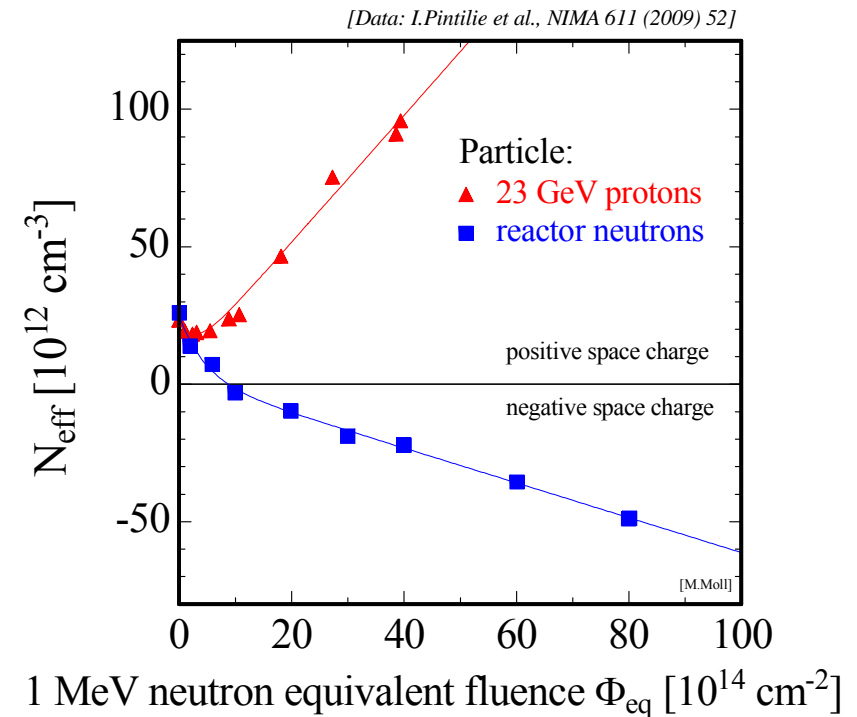
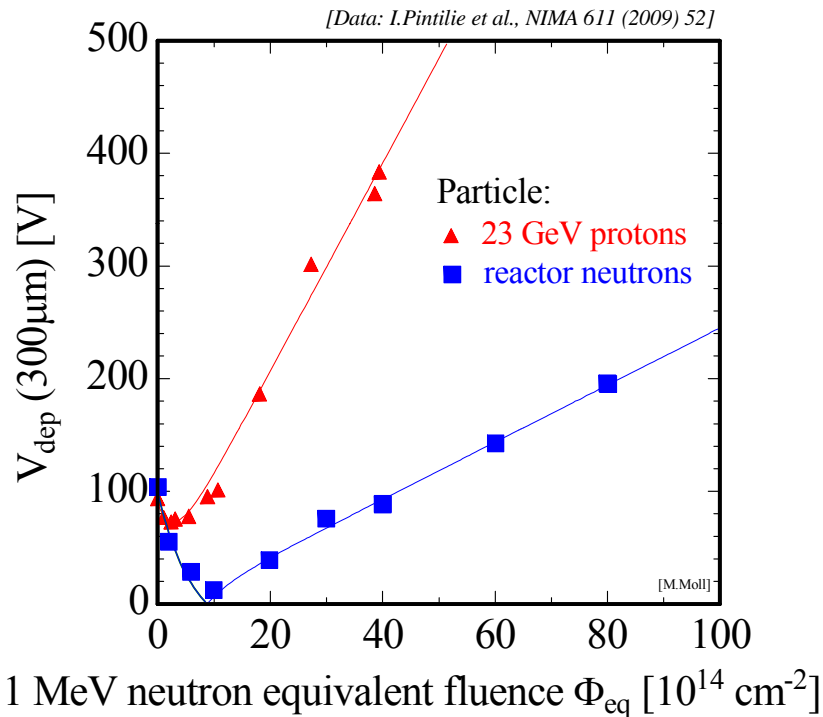
- “no type inversion” in the overall fluence range

(for experts: there is no “real” type inversion, a more clear understanding of the observed effects is obtained by investigating directly the internal electric field; look for: TCT, MCZ, double junction)



- **Common to all materials** *(after hadron irradiation, not after γ irradiation)*:
 - reverse current increase
 - increase of trapping (electrons and holes) within $\sim 20\%$

- **Epitaxial silicon** (*EPI-DO*, 72 μm , 170 Ωcm , diodes)
irradiated with **23 GeV protons** or **reactor neutrons**

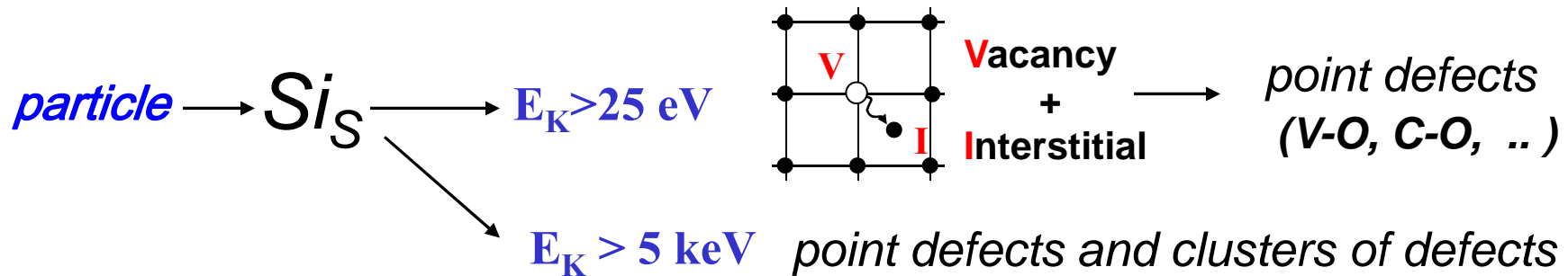


depletion voltage \rightarrow

$$V_{\text{dep}} = \frac{q_0}{\epsilon\epsilon_0} \cdot |N_{\text{eff}}| \cdot d^2$$

\leftarrow absolute effective space charge density

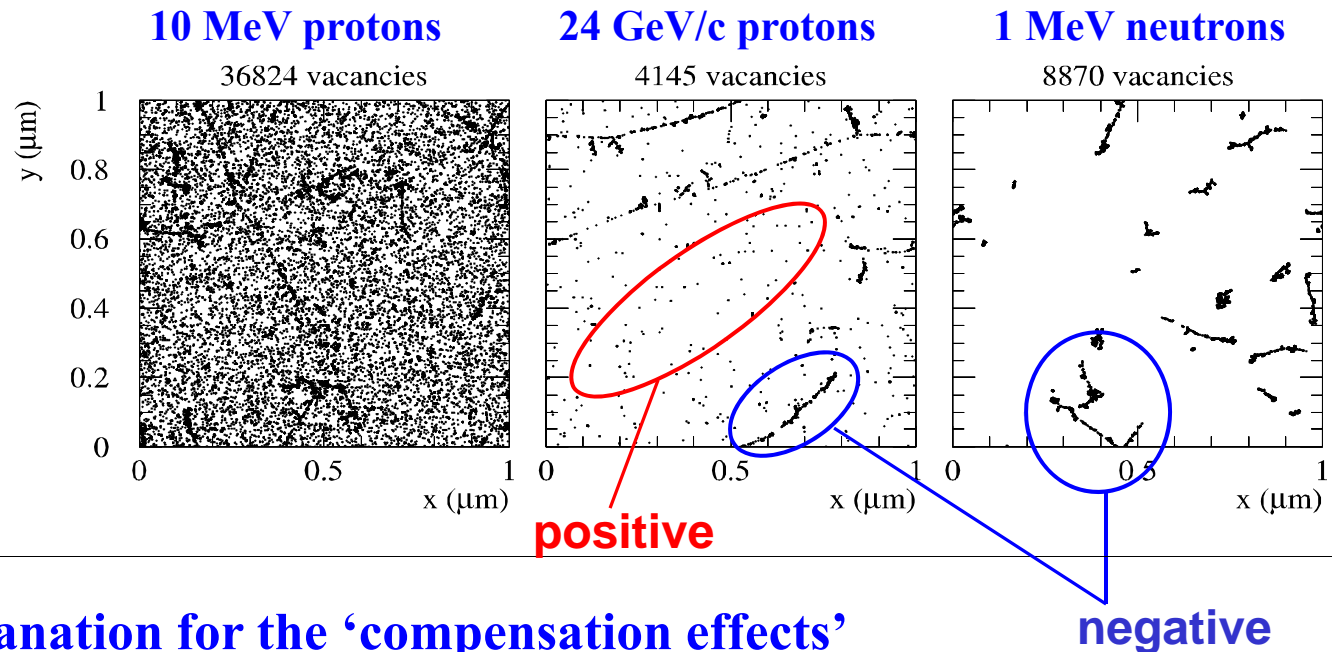
Why is proton and neutron damage different?



Simulation:

Initial distribution of vacancies in $(1\mu\text{m})^3$ after 10^{14} particles/cm²

[Mika Huhtinen NIMA 491(2002) 194]



- A ‘simplified’ explanation for the ‘compensation effects’
 - Defect clusters produce predominantly **negative space charge**
 - Point defects produce predominantly **positive space charge** (in ‘oxygen rich’ silicon)

For the experts: Note the NIEL violation

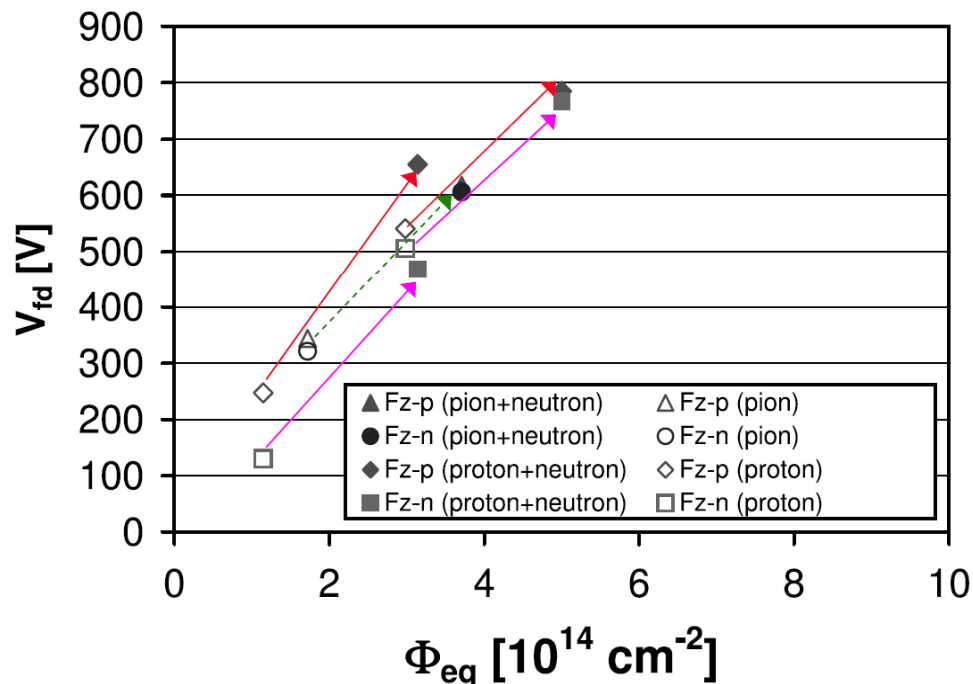


Mixed irradiations – Change of N_{eff}

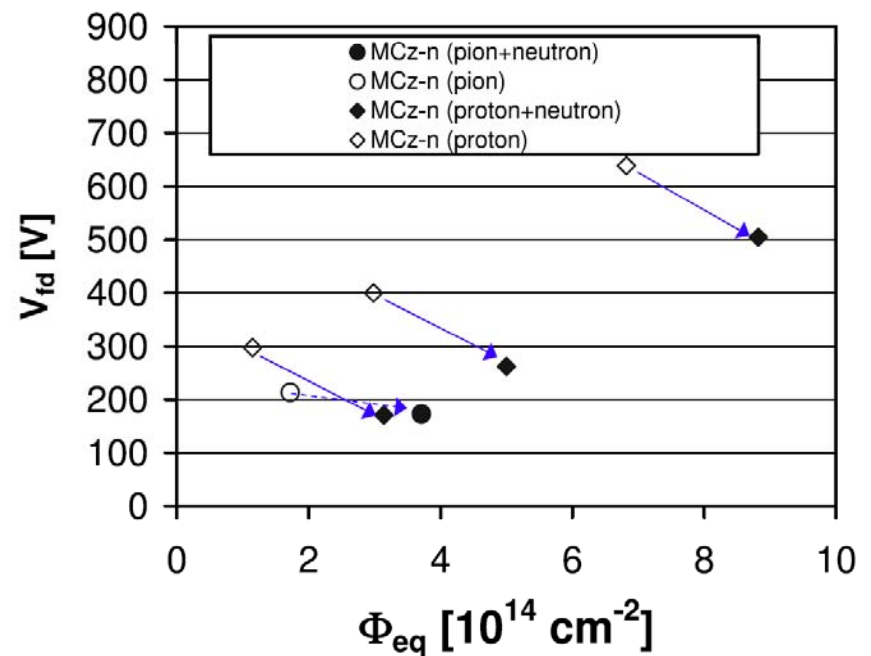


- Exposure of FZ & MCZ silicon sensors to ‘mixed’ irradiations
 - First step: Irradiation with protons or pions
 - Second step: Irradiation with neutrons

FZ: Accumulation of damage



MCZ: Compensation of damage



[G.Kramberger et al., “Performance of silicon pad detectors after mixed irradiations with neutrons and fast charged hadrons”, NIMA 609 (2009) 142-148]

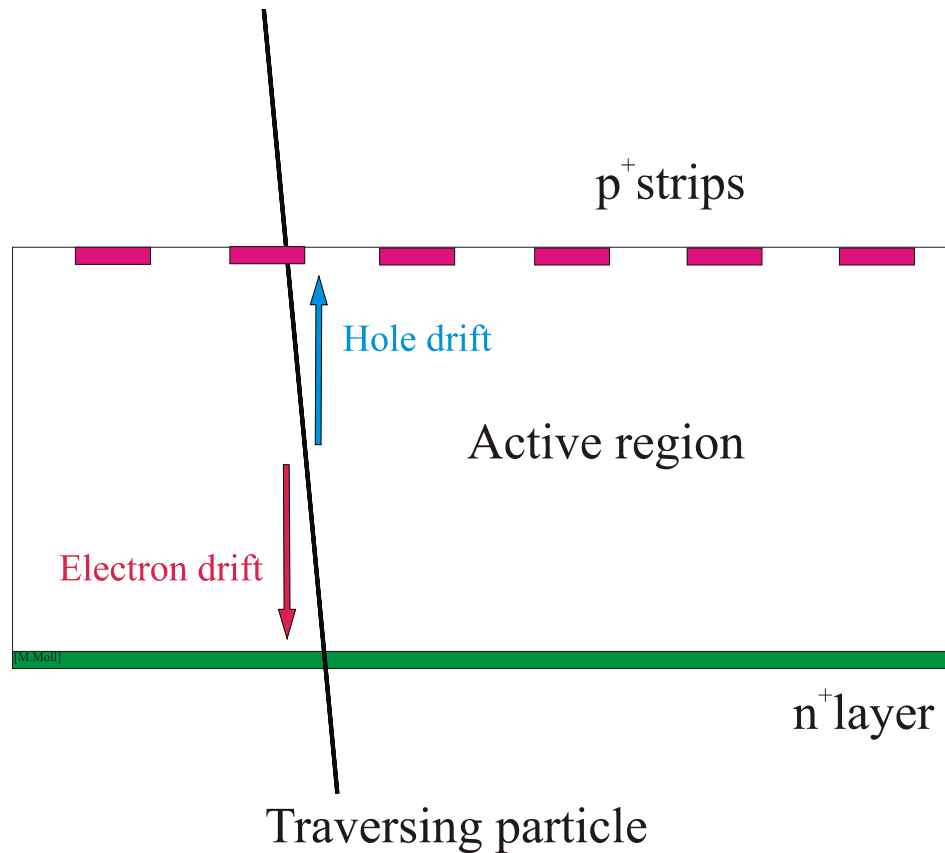


Advantage of non-inverting material

p-in-n detectors (schematic figures!)



Fully depleted detector
(non – irradiated):





Advantage of non-inverting material

p-in-n detectors (schematic figures!)



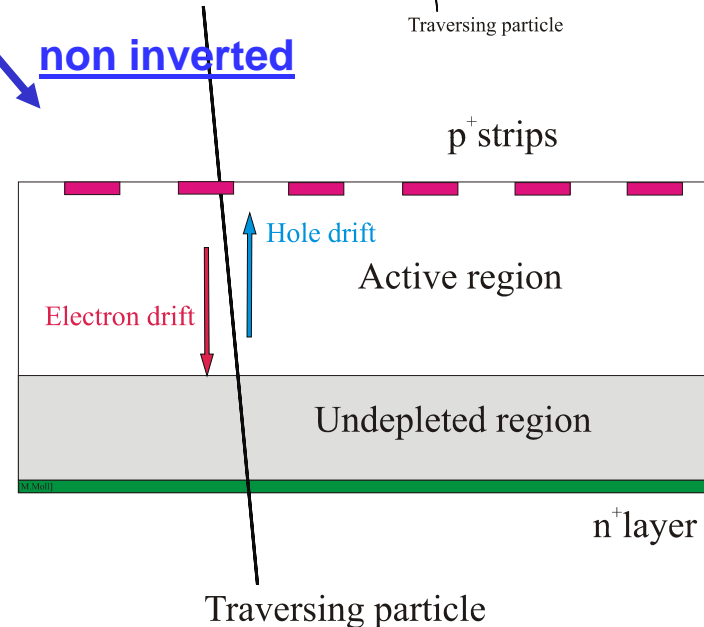
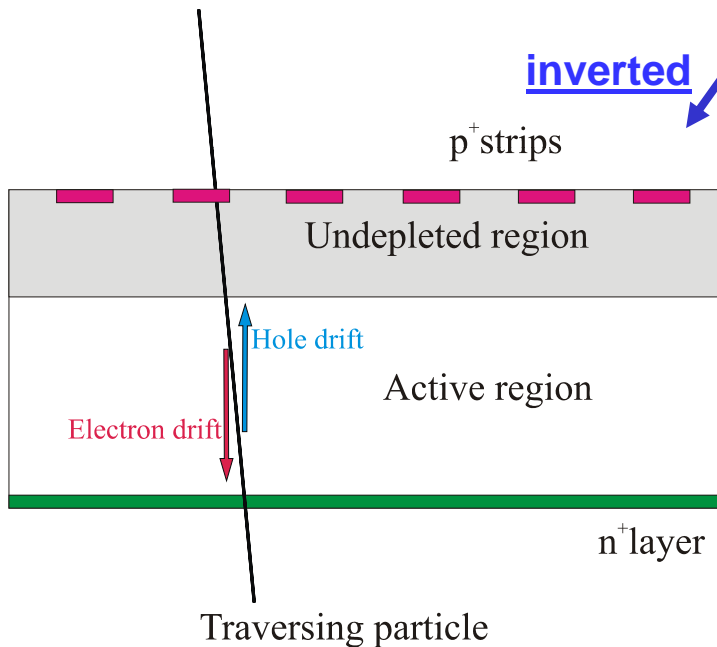
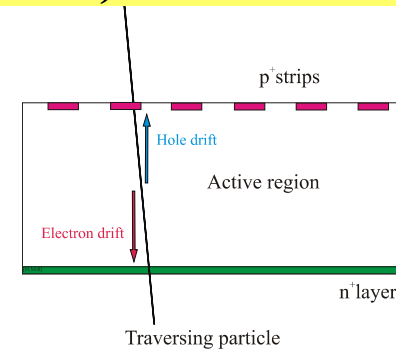
Be careful, this is a very schematic explanation, reality is more complex !

Fully depleted detector (non – irradiated):

heavy irradiation

inverted

non inverted



inverted to “p-type”, under-depleted:

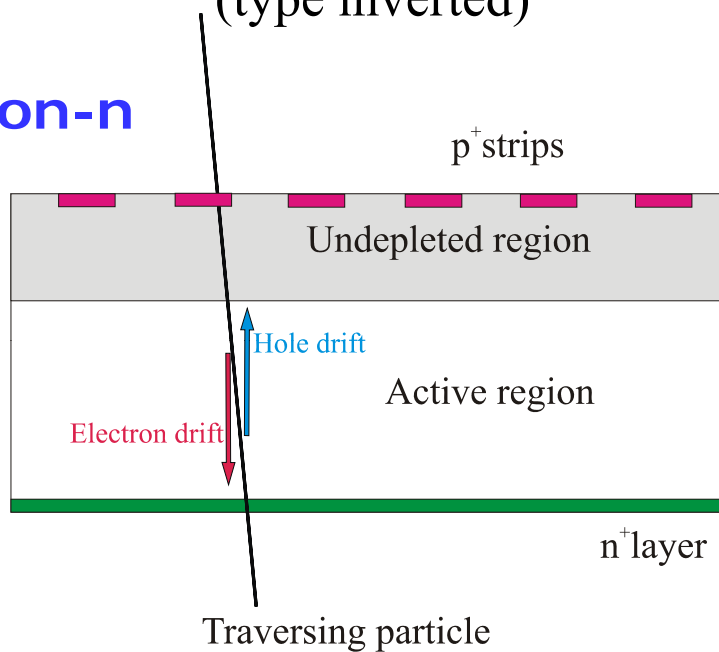
- Charge spread – degraded resolution
- Charge loss – reduced CCE

non-inverted, under-depleted:

- Limited loss in CCE
- Less degradation with under-depletion

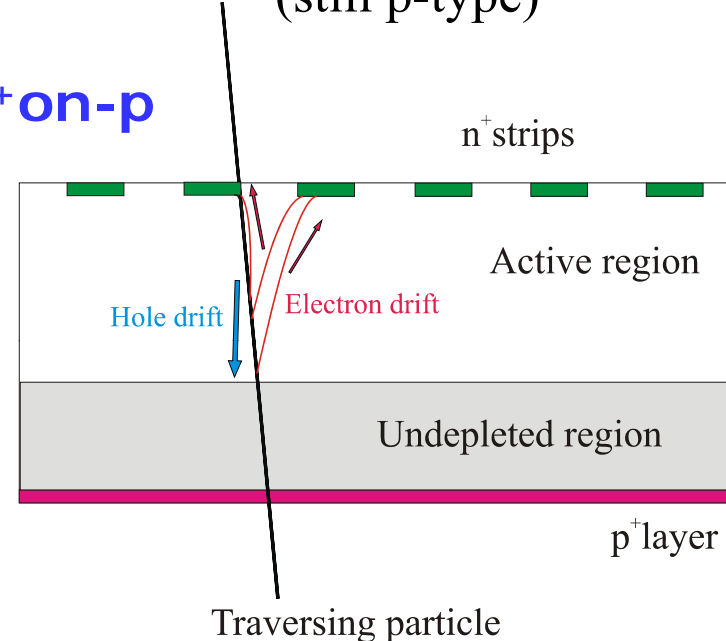
n-type silicon after high fluences:
(type inverted)

p⁺on-n



p-type silicon after high fluences:
(still p-type)

n⁺on-p



p-on-n silicon, under-depleted:

- Charge spread – degraded resolution
- Charge loss – reduced CCE

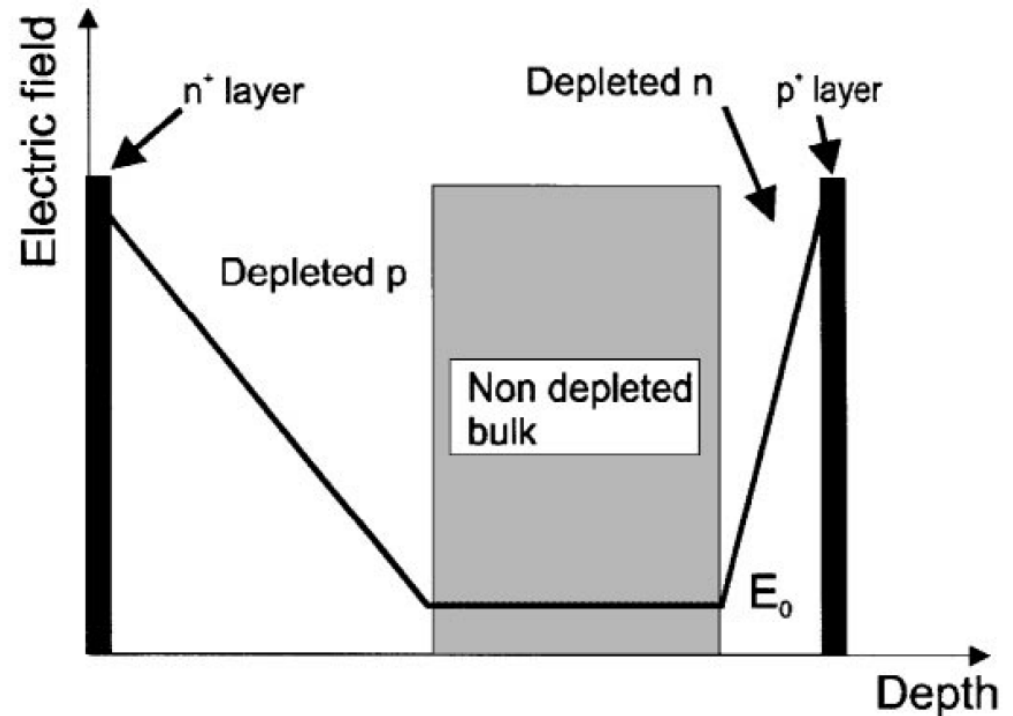
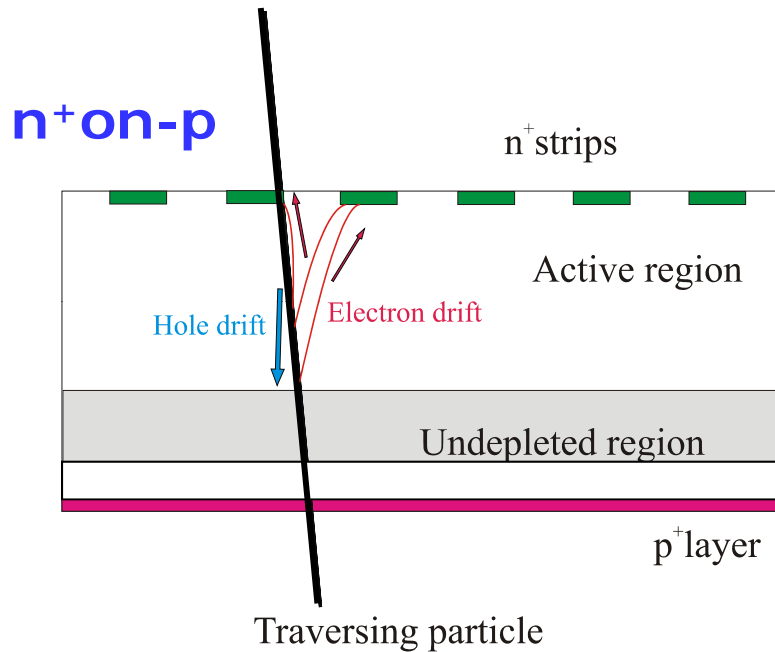
n-on-p silicon, under-depleted:

- Limited loss in CCE
- Less degradation with under-depletion
- Collect electrons (3 x faster than holes)

Comments:

- Instead of n-on-p also n-on-n devices could be used

p-type silicon after high fluences:
(still “p-type”)



- **Dominant junction close to n⁺ readout strip for FZ n-in-p**
- **For MCZ p-in-n even more complex fields have been reported:**
 - no “type inversion”(SCSI) = dominant field remains at p implant
 - “equal double junctions” with almost symmetrical fields on both sides

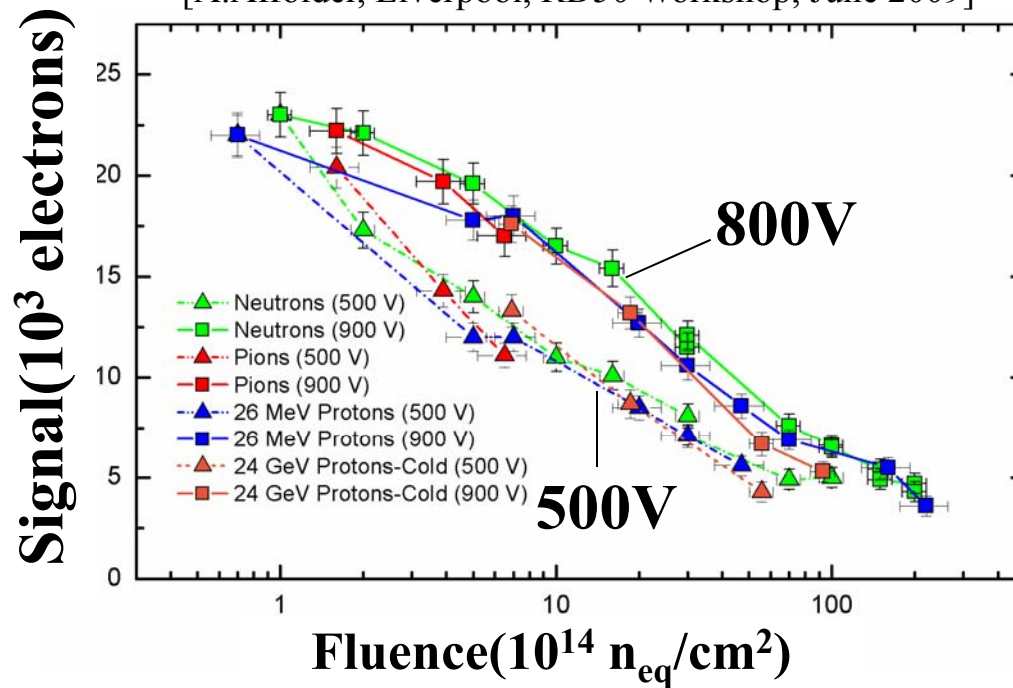


FZ n-in-p microstrip detectors (n, p, p - irradi)



- **n-in-p microstrip p-type FZ detectors** (Micron, 280 or 300 μm thick, 80 μm pitch, 18 μm implant)
- **Detectors read-out with 40MHz** (SCT 128A)

[A.Affolder, Liverpool, RD50 Workshop, June 2009]



- **CCE: $\sim 7300e$ ($\sim 30\%$)**
after $\sim 1 \times 10^{16} \text{cm}^{-2}$ 800V
- **n-in-p sensors are strongly considered**
for ATLAS upgrade (previously p-in-n used)

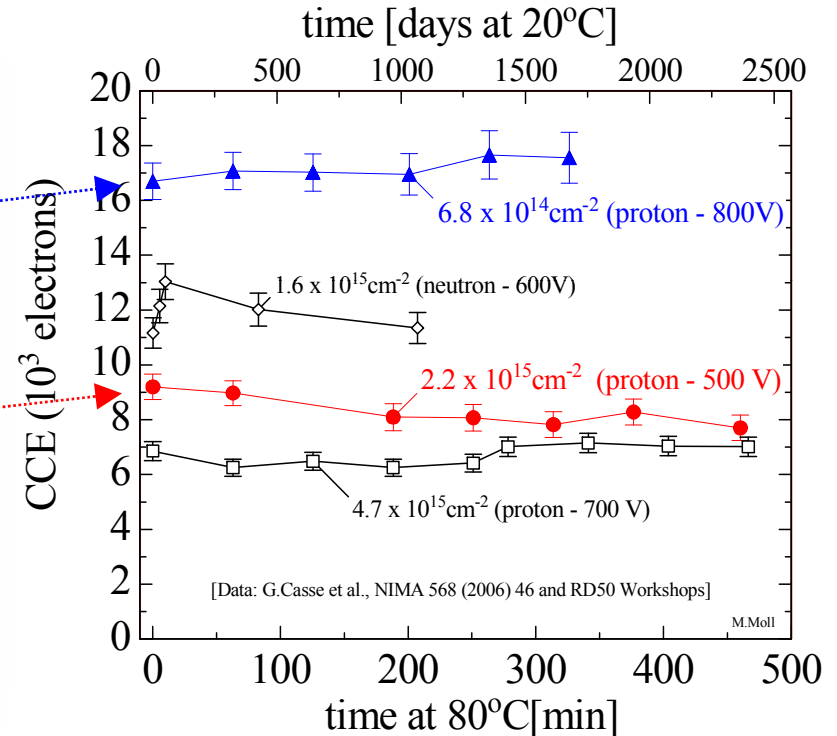
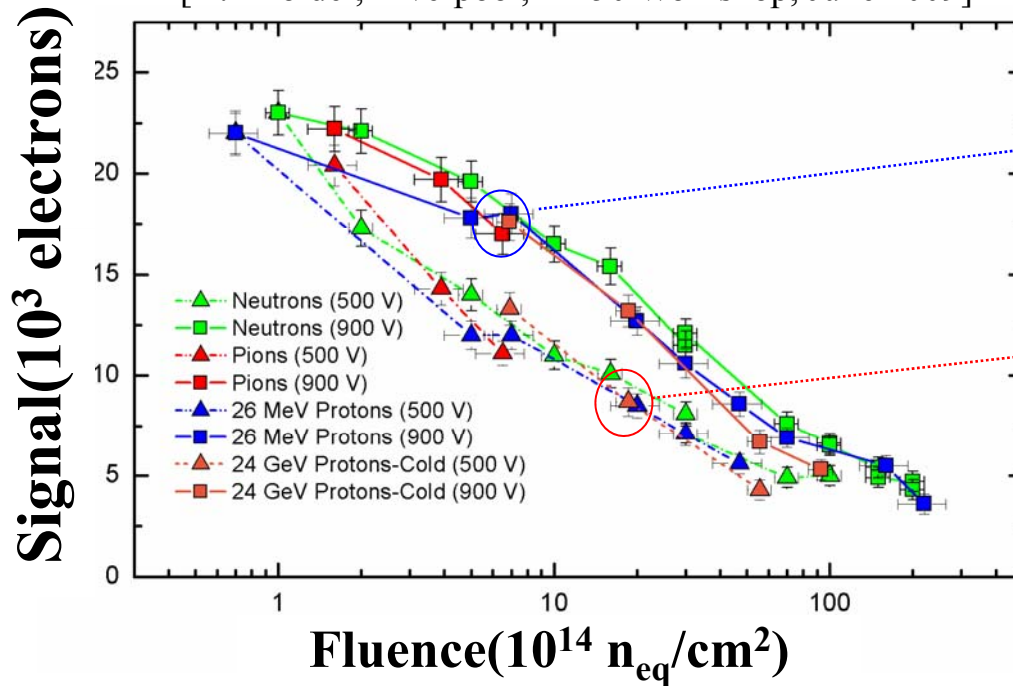


FZ n-in-p microstrip detectors (n, p, π – irradiad)



- **n-in-p microstrip p-type FZ detectors** (Micron, 280 or 300 μ m thick, 80 μ m pitch, 18 μ m implant)
- **Detectors read-out with 40MHz** (SCT 128A)

[A.Affolder, Liverpool, RD50 Workshop, June 2009]



- **CCE: $\sim 7300e$ ($\sim 30\%$)**
after $\sim 1 \times 10^{16} cm^{-2}$ 800V
- **n-in-p sensors are strongly considered for ATLAS upgrade** (previously p-in-n used)

- **no reverse annealing in CCE measurements for neutron and proton irradiated detectors**

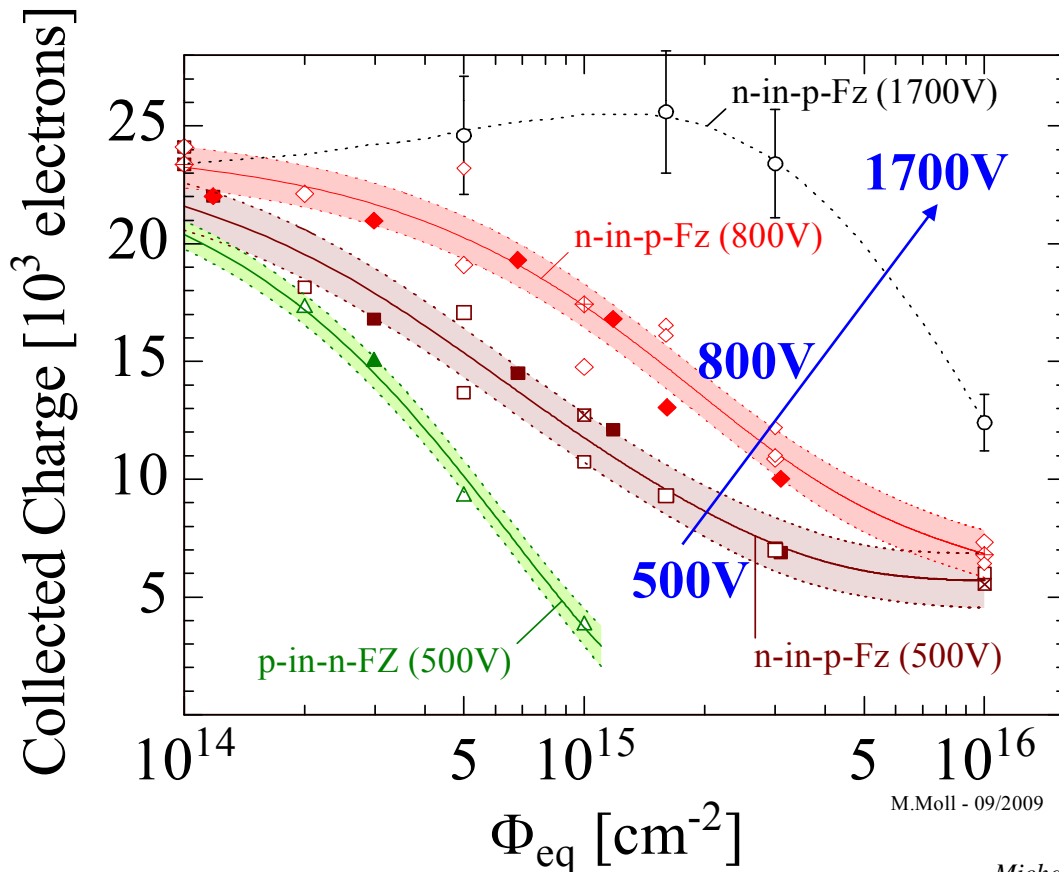


Good performance of planar sensors at high fluence



• Why do planar silicon sensors with n-strip readout give such high signals after high levels ($>10^{15} \text{ cm}^{-2} \text{ p/cm}^2$) of irradiation?

- Extrapolation of charge trapping parameters obtained at lower fluences would predict much lower signal!
- Assumption: ‘Charge multiplication effects’ as even $\text{CCE} > 1$ was observed



FZ Silicon Strip Sensors

- n-in-p (FZ), 300 μm , 500V, 23GeV p [1]
- n-in-p (FZ), 300 μm , 500V, neutrons [1,2]
- ⊠ n-in-p (FZ), 300 μm , 500V, 26MeV p [1]
- ◆ n-in-p (FZ), 300 μm , 800V, 23GeV p [1]
- ◇ n-in-p (FZ), 300 μm , 800V, neutrons [1,2]
- ◊ n-in-p (FZ), 300 μm , 800V, 26MeV p [1]
- n-in-p (FZ), 300 μm , 1700V, neutrons [2]
- ▲ p-in-n (FZ), 300 μm , 500V, 23GeV p [1]
- △ p-in-n (FZ), 300 μm , 500V, neutrons [1]

References:

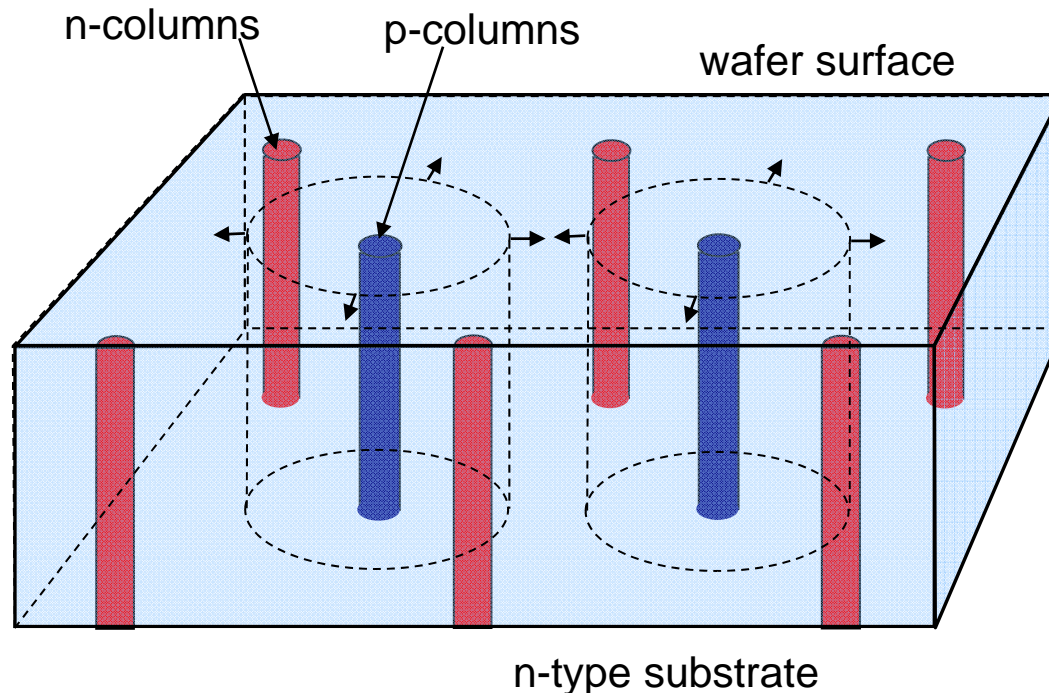
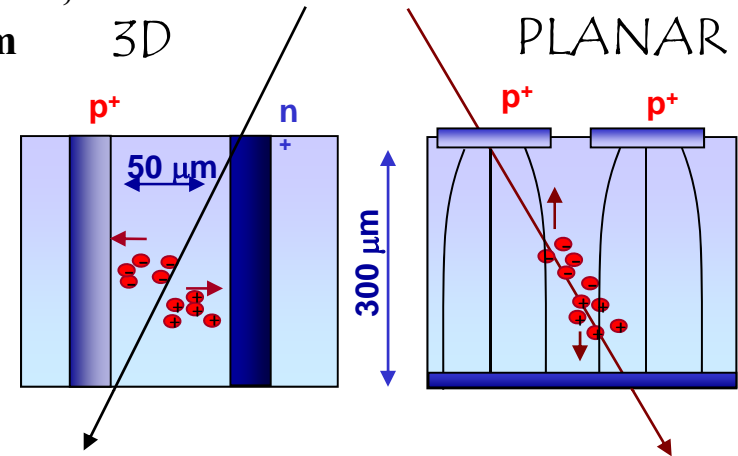
[1] G.Casse, VERTEX 2008
(p/n-FZ, 300 μm , (-30°C, 25ns)

[2] I.Mandic et al., NIMA 603 (2009) 263
(p-FZ, 300 μm , -20°C to -40°C, 25ns)

• Which voltage can be applied?

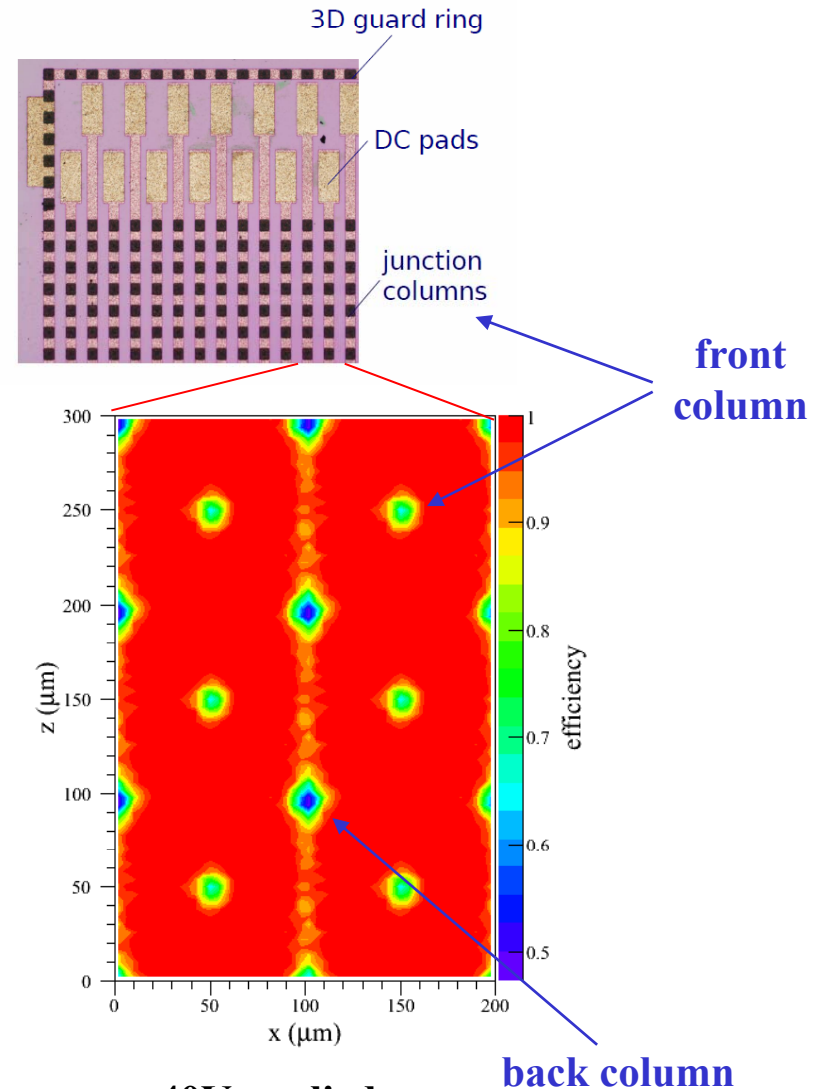
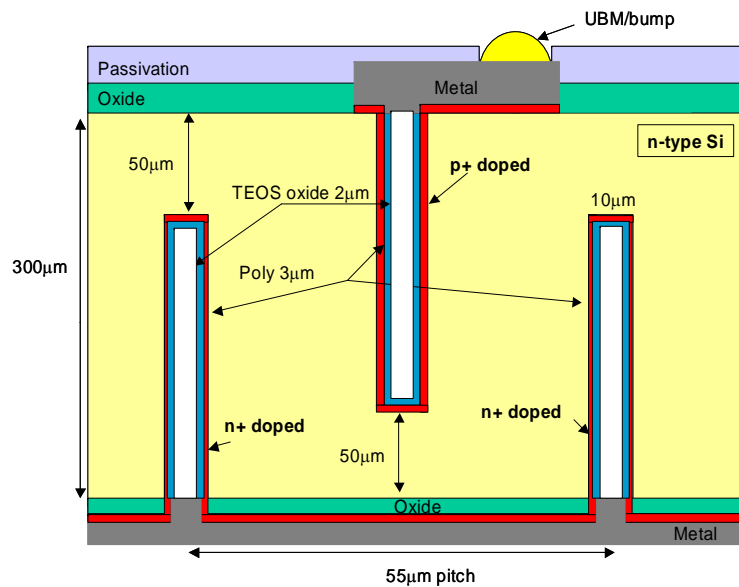
3D detector - concept

- **“3D” electrodes:**
 - narrow columns along detector thickness,
 - diameter: $10\mu\text{m}$, distance: $50 - 100\mu\text{m}$
- **Lateral depletion:**
 - lower depletion voltage needed
 - thicker detectors possible
 - fast signal
 - radiation hard



- **DDTC – Double sided double type column**

[G.Fleta, RD50 Workshop, June 2007]



40V applied
~98% efficiency

- **Testbeam data – Example: efficiency map**
[M.Koehler, Freiburg Uni, RD50 Workshop June 09]

- **Processing of 3D sensors is challenging, but many good devices with reasonable production yield produced.**

see lecture by
Manuel Lozano

- **Competing e.g. for ATLAS IBL pixel sensors**

Property	Diamond	GaN	4H SiC	Si
E_g [eV]	5.5	3.39	3.3	1.12
$E_{\text{breakdown}}$ [V/cm]	10^7	$4 \cdot 10^6$	$2.2 \cdot 10^6$	$3 \cdot 10^5$
μ_e [cm^2/Vs]	1800	1000	800	1450
μ_h [cm^2/Vs]	1200	30	115	450
v_{sat} [cm/s]	$2.2 \cdot 10^7$	-	$2 \cdot 10^7$	$0.8 \cdot 10^7$
e-h energy [eV]	13	8.9	7.6-8.4	3.6
e-h pairs/ X_0	4.4	~2-3	4.5	10.1

- **Diamond: wider bandgap**
 \Rightarrow lower leakage current
 \Rightarrow less cooling needed
 \Rightarrow less noise

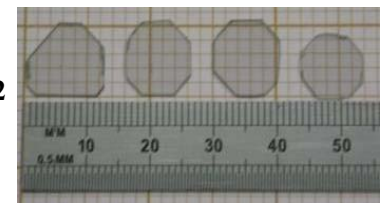
- **Signal produced by m.i.p:**
Diamond 36 e/ μm
Si 89 e/ μm
 \Rightarrow Si gives more charge than diamond

- **GaAs, SiC and GaN** \Rightarrow strong radiation damage observed
 \Rightarrow no potential material for LHC upgrade detectors
(judging on the investigated material)
- **Diamond (RD42)** \Rightarrow good radiation tolerance *(see later)*
 \Rightarrow already used in LHC beam condition monitoring systems
 \Rightarrow considered as potential detector material for sLHC pixel sensors

poly-CVD Diamond
 -16 chip ATLAS
 pixel module



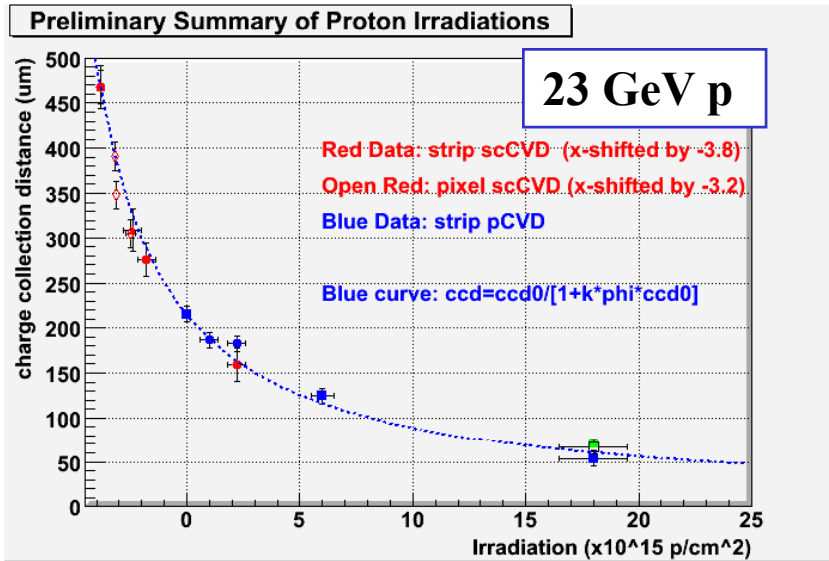
single crystal CVD
 Diamond of few cm^2



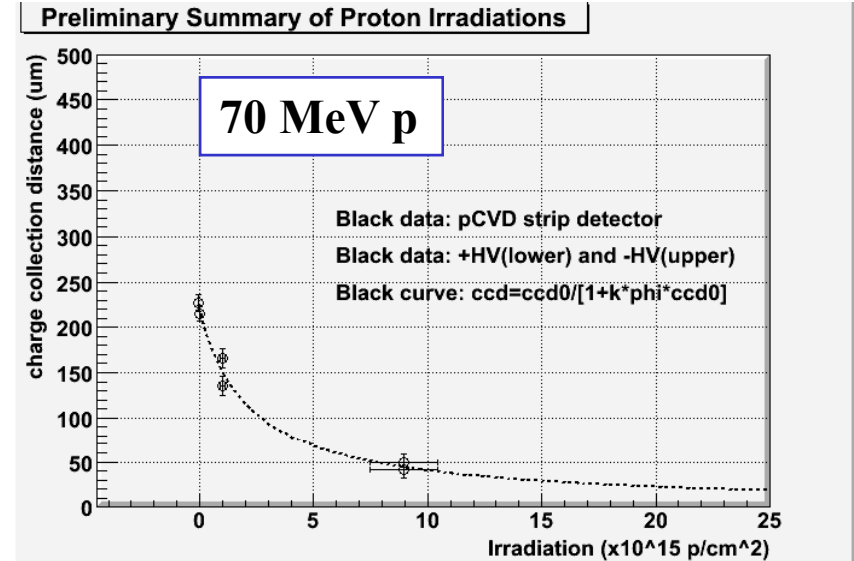
**Diamond sensors are heavily used in
 LHC Experiments for Beam Monitoring**



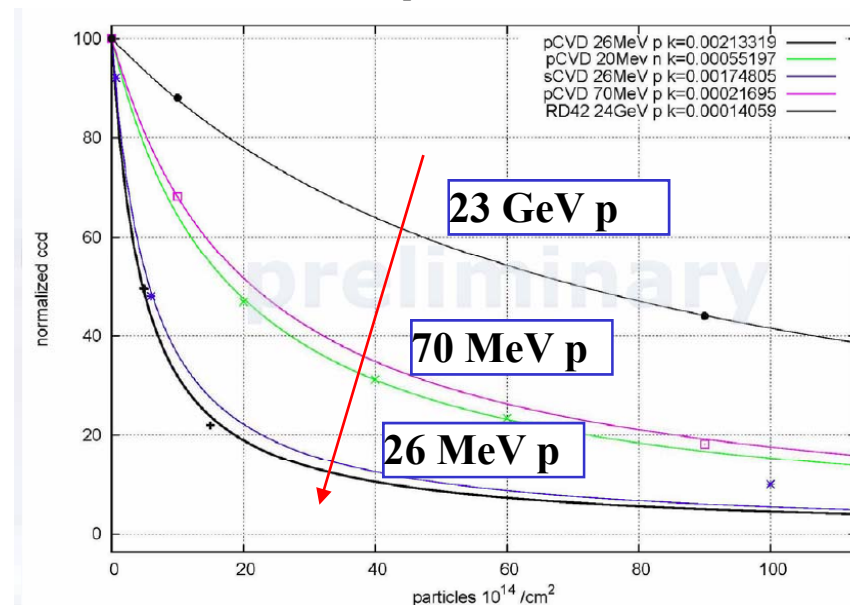
Are diamond sensors radiation hard?



[RD42, LHCC Status Report, Feb. 2010]



[RD42, LHCC Status Report, Feb. 2010]



[V.Rylov, CERN ESE Seminar 9.11.2009]

- Most published results on 23 GeV protons
- 70 MeV protons 3 times more damaging than 23 GeV protons
- 25 MeV protons seem to be even more damaging (Preliminary: RD42 about to cross check the data shown to the left)
- In line with NIEL calc. for Diamond [W. de Boer et al. Phys.Status Solidi 204:3009,2007]



Summary – Radiation Damage



- **Radiation Damage in Silicon Detectors**

- Change of **Depletion Voltage** (internal electric field modifications, “type inversion”, reverse annealing, loss of active volume, ...) (can be influenced by defect engineering!)
- Increase of **Leakage Current** (same for all silicon materials)
- Increase of **Charge Trapping** (same for all silicon materials)

Signal to Noise ratio is quantity to watch (material + geometry + electronics)

- **Microscopic defects & Damage scaling factors**

- Microscopic crystal defects are the origin to detector degradation.
- NIEL – Hypothesis used to scale damage of different particles with different energy
- Different particles produce different types of defects! (NIEL – violation!)
- There has been an enormous progress in the last 5 years in understanding defects.

**Details in next lecture by
Ioana Pintilie on defects.**

- **Approaches to obtain radiation tolerant devices:**

- **Material Engineering:**
 - explore and develop new silicon materials (oxygenated Si)
 - use of other semiconductors (Diamond)
- **Device Engineering:**
 - look for other sensor geometries
 - 3D, thin sensors, n-in-p, n-in-n, ...



Detectors for the LHC upgrade



- **At fluences up to 10^{15}cm^{-2} (outer layers – ministrip sensors):**

The change of the depletion voltage and the large area to be covered by detectors are major problems.

- **n-MCZ silicon detectors** show good performance in mixed fields due to compensation of charged hadron damage and neutron damage (N_{eff} compensation) (more work needed)

- **p-type silicon microstrip detectors show very encouraging results**

$\text{CCE} \approx 6500 \text{ e}$; $\Phi_{\text{eq}} = 4 \times 10^{15} \text{ cm}^{-2}$, $300 \mu\text{m}$, immunity against reverse annealing!

This is presently the “most considered option” for the ATLAS SCT upgrade

- **At fluences $> 10^{15}\text{cm}^{-2}$ (Innermost tracking layers – pixel sensors)**

The active thickness of any silicon material is significantly reduced due to trapping.

Collection of electrons at electrodes essential: Use n-in-p or n-in-n detectors!

- **Recent results show that planar silicon sensors might still give sufficient signal,** (still some interest in epitaxial silicon and thin sensor options)

- **3D detectors** : looks promising, drawback: technology has to be optimized!

Many collaborations and sensor producers working on this.

- **Diamond** has become an interesting option (*Higher damage due to low energy protons?*)

- **Questions to be answered:**

- **a) Can we profit from avalanche effects and control them?**

- **b) Can we profit from compensation effects in mixed fields?**

- **c) Can we understand detector performance on the basis of simulations using defect parameters as input?**



Details in lecture by
Phil Allport



Acknowledgements & References



- **Many thanks to the MC-PAD Network for the invitation to give this lecture**
- **Most references to particular works given on the slides**
- **Some additional material taken from the following presentations:**
 - *RD50 presentations: <http://www.cern.ch/rd50/>*
 - *Anthony Affolder: Presentations on the RD50 Workshop in June 2009 (sATLAS fluence levels)*
 - *Frank Hartmann: Presentation at the VCI conference in February 2010 (Diamond results)*
- **Books containing chapters about radiation damage in silicon sensors**
 - *Helmuth Spieler, “Semiconductor Detector Systems”, Oxford University Press 2005*
 - *Frank Hartmann, “Evolution of silicon sensor technology in particle physics”, Springer 2009*
 - *L.Rossi, P.Fischer, T.Rohe, N.Wermes “Pixel Detectors”, Springer, 2006*
 - *Gerhard Lutz, “Semiconductor radiation detectors”, Springer 1999*
- **Research collaborations and web sites**
 - *The RD50 collaboration (<http://www.cern.ch/rd50>) - Radiation Tolerant Silicon Sensors*
 - *The RD39 collaboration – Cryogenic operation of Silicon Sensors*
 - *The RD42 collaboration – Diamond detectors*
 - *ATLAS IBL, ATLAS and CMS upgrade groups*