

1 From Marc - 2 Sept 2020

P. 2 : I rewrote Mariana's Lagrangian with some minor corrections (e.g. indices) and the connection with Terning's paper (with some constants that we might have to consider).

P. 3: I checked my results sent previously, and fixed my 1/2 factor! Done for visible A only. The red arrows indicate differences with Mariana's equation (1) (from her 12 August notes, here in P. 4)

P. 4 From Mariana's notes sent on 12 August.

PP. 5-6: Pages sent previously, applied to the equations of motion of P. 3. If this is correct, it suggest that Terning's diagonalization basis in Eqs. (4.4) - (4.7) is correct.

Next steps: check equations for B , A_D and B_D and compare with Mariana's results on P. 7 (in particular the unwanted term $\propto B_D^\mu$ in Eq. (2d)). Work out the Feynman rules. MadGraph.

Marc 2 SEPTEMBER 2020

LAGRANGIAN (From MARIANA, 12 August, with minor corrections)

EoS in TERNING (2018)

$$\mathcal{L} = -\frac{n^A n^B}{2n^2} g^{\mu\nu} (F_{AB}^A F_{\mu\nu}^A + F_{AB}^B F_{\mu\nu}^B)$$

(2.4) $\times e^2$

$$+ \frac{n^A n^B}{4n^2} \epsilon^{\mu\nu\delta\sigma} (F_{\alpha\delta}^B F_{\gamma\sigma}^A - F_{\alpha\delta}^A F_{\gamma\sigma}^B)$$

$$- J_\mu^A A^\mu - K_\mu^B B^\mu$$

$$(2.4); K_\mu \times \frac{e^2}{4\pi}$$

$$- \frac{n^A n^B}{2n^2} g^{\mu\nu} (F_{AB}^{AD} F_{\mu\nu}^{AD} + F_{AB}^{BD} F_{\mu\nu}^{BD})$$

(4.1)

$$+ \frac{n^A n^B}{4n^2} \epsilon^{\mu\nu\delta\sigma} (F_{\alpha\delta}^{BD} F_{\gamma\sigma}^{AD} - F_{\alpha\delta}^{AD} F_{\gamma\sigma}^{BD})$$

$$- J_\mu^D A_D^\mu - K_\mu^D B_D^\mu$$

$$J^D \leftarrow \frac{(4.1)}{e_D} \quad K^D \leftarrow \frac{(4.1) \times e_D}{4\pi}$$

$$+ e e_D \frac{n^A n^B}{n^2} g^{\mu\nu} (F_{AB}^{AD} F_{\mu\nu}^A - F_{AB}^{BD} F_{\mu\nu}^B)$$

(4.2)

$$- \frac{1}{2} m_A^2 A_{\mu\nu} A^\mu - \frac{1}{2} m_B^2 B_{\mu\nu} B^\mu$$

(5.1) (5.3)

$$F_{\mu\nu}^X \equiv \partial_\mu X_\nu - \partial_\nu X_\mu$$

$$J_{(D)}^\mu \equiv e_{(D)} \overline{\psi}_{e_{(D)}} \gamma^\mu \psi_{e_{(D)}} \quad K_{(D)}^\mu \equiv g_{(D)} \overline{\psi}_{g_{(D)}} \gamma^\mu \psi_{g_{(D)}}$$

$$\text{is } g_{(D)} = \frac{4\pi}{e_{(D)}^2} \text{ or } \frac{4\pi}{e_D} ?$$

2.

Eqs of Motion w/r A_μ

$$\begin{aligned}
 -J^\sigma &= -\frac{n_\mu}{m^2} (n^\delta \partial_S F_A^{\mu\delta} + n^\sigma \partial_S F_A^{\sigma\mu}) \\
 &\quad + \epsilon_{\nu\sigma} \frac{n_\lambda}{m^2} (n^\delta \partial_S F_{AB}^{\lambda\delta} - n^\sigma \partial_S F_{AB}^{\lambda\sigma}) \\
 &\quad + \frac{n_\mu}{4m^2} \left[2n^\alpha \epsilon^{\mu\nu\delta\sigma} \partial_S F_{AB}^B - (n^\delta \epsilon^{\mu\nu\delta\sigma} - \underbrace{n^\sigma \epsilon^{\mu\nu\delta\delta}}_{\text{BY CONTRACTION}}) \partial_S F_{BB}^B \right]
 \end{aligned}$$

LINE 3 is $\frac{n_\mu n^\delta}{2m^2} \overbrace{\epsilon^{\mu\nu\delta\sigma} \partial_S F_{AB}^B}^{\text{RENAMING + PERMUTING indices...}} - \frac{n_\mu n^\delta}{4m^2} \epsilon^{\mu\nu\delta\delta} \partial_S F_{BB}^B$

$$\Rightarrow \frac{n_\mu n^\delta}{2m^2} \epsilon^{\mu\nu\delta\sigma} \partial_S \partial_\sigma B_\nu + \frac{n_\mu n^\delta}{4m^2} \epsilon^{\mu\nu\delta\sigma} \partial_\sigma (\partial_S B_\nu - \partial_\nu B_S) \quad \textcircled{*}$$

IN TERM 1 of $\textcircled{*}$ with $\mu=0, \sigma=1$: $\epsilon^{0\nu 1\delta} \partial_S B_\nu = \partial_3 B_2 - \partial_2 B_3 = F_{32}$

IN TERM 2: $\epsilon^{0\nu 1\delta} (\partial_S B_\nu - \partial_\nu B_S) = \partial_3 B_2 - \partial_2 B_3 - (\partial_2 B_3 - \partial_3 B_2) = 2F_{32}$

ETC. SO THE TWO TERMS IN $\textcircled{*}$ ARE EQUAL SO THAT $\textcircled{*} = \text{LINE 3}$ IS

SIMPLY $\frac{n_\mu n^\delta}{m^2} \epsilon^{\mu\nu\delta\sigma} \partial_S F_{AB}^B$. Eqs of motion BECOMES

$$\begin{aligned}
 \frac{n_\mu}{m^2} (-n^\mu \partial_\nu F_A^{\nu\lambda} + n^\nu \partial_\nu F_A^{\mu\lambda}) + \frac{\epsilon_{\nu\sigma} n_\lambda}{m^2} (n^\mu \partial_\nu F_{AB}^{\lambda\mu} - n^\nu \partial_\nu F_{AB}^{\lambda\mu}) \\
 - \frac{n_\lambda}{m^2} n^\gamma \epsilon^{\lambda\beta\delta\mu} \partial_\nu F_{BB}^B = J^\mu
 \end{aligned}$$

$$\frac{n_\mu}{m^2} \left(\cancel{-n^\mu \partial_\nu F_A^{\nu\lambda}} + \cancel{n^\nu \partial_\nu F_A^{\mu\lambda}} + n^\delta \epsilon^{\mu\nu\delta\lambda} \partial_\nu F_{BB}^B \right)$$

$$+ \epsilon_{\nu\sigma} \frac{n_\lambda}{m^2} \partial_\nu (n^\mu F_{AB}^{\lambda\mu} - n^\nu F_{AB}^{\lambda\mu}) = J^\mu$$

Equations of motion

VISIBLE

$$(1) \quad \frac{n_\alpha}{n^2} \left(n^\mu \partial_\nu F_A^{\alpha\nu} - n^\nu \partial_\nu F_A^{\alpha\mu} - \epsilon^{\mu\nu\alpha}_\beta n_\gamma \partial_\nu F_B^{\gamma\beta} \right)$$

$$+ \frac{e e e_0}{n^2} n_\alpha \partial_\nu \left(n^\mu F_{A_D}^{\alpha\nu} - n^\nu F_{A_D}^{\alpha\mu} \right) = J^\mu$$

$$(2) \quad \frac{n_\alpha}{n^2} \left(n^\mu \partial_\nu F_B^{\alpha\nu} - n^\nu \partial_\nu F_B^{\alpha\mu} + \epsilon^{\mu\nu\alpha}_\beta n_\gamma \partial_\nu F_A^{\gamma\beta} \right)$$

$$- \frac{e e e_0}{n^2} n_\alpha \partial_\nu \left(n^\mu F_{B_D}^{\alpha\nu} - n^\nu F_{B_D}^{\alpha\mu} \right) = K^\mu$$

DARK

$$(3) \quad \frac{n_\alpha}{n^2} \left(n^\nu \partial_\nu F_{A_D}^{\alpha\nu} - n^\nu \partial_\nu F_{A_D}^{\alpha\mu} - \epsilon^{\mu\nu\alpha}_\beta n_\gamma \partial_\nu F_{B_D}^{\gamma\beta} \right)$$

$$+ \frac{e e e_0}{n^2} n_\alpha \partial_\nu \left(n^\mu F_A^{\alpha\nu} - n^\nu F_A^{\alpha\mu} \right) = J_D^\mu + M_{A_D}^2 A_D^\mu A_{D\mu}$$

$$(4) \quad \frac{n_\alpha}{n^2} \left(n^\mu \partial_\nu F_{B_D}^{\alpha\nu} - n^\nu \partial_\nu F_{B_D}^{\alpha\mu} + \epsilon^{\mu\nu\alpha}_\beta n_\gamma \partial_\nu F_{A_D}^{\gamma\beta} \right)$$

$$- \frac{e e e_0}{n^2} \left(n_\alpha \partial_\nu n^\mu F_B^{\alpha\nu} - n_\nu F_B^{\alpha\mu} \right) = K_D^\mu + M_{B_D}^2 B_D^\mu B_{D\mu}$$

Proceeding as in previous p.3 with my signs but without η_2 :

$$(1) \quad \frac{m_A}{m^2} \left(m^2 F_A^{AB} - m^2 F_A^{AB} + e \cos \theta_m F_B^{AB} \right)$$

$$+ \frac{e \cos \theta_m}{m^2} \left(m^2 F_{AB} - m^2 F_{AB} \right) = J^A$$

$$J^A = \frac{m_A}{m^2} \left[m^2 F_A^{AB} - m^2 F_A^{AB} + e \cos m^2 F_{AB} - e \cos m^2 F_{AB} \right]$$

$$= \frac{m_A}{m^2} \left[m^2 \left(F_A^{AB} - e \cos F_A^{AB} \right) - m^2 \left(F_A^{AB} - e \cos F_A^{AB} \right) + m^2 e \cos \theta_m F_B^{AB} \right]$$

From turning Eqs (4.4,5) $F_A^{AB} = (\cos \phi + e \cos \sin \phi) F_A^{AB} + (-\sin \phi + e \cos \cos \phi) F_A^{AB}$

$$F_A^{AB} = \sin \phi F_A^{AB} + \cos \phi F_A^{AB}$$

$$F_B^{AB} = \cos \phi F_B^{AB} - \sin \phi F_B^{AB}$$

$$\begin{aligned} F_A^{AB} - e \cos \theta_m F_A^{AB} &= \cos \phi F_A^{AB} + e \cos \sin \phi F_A^{AB} - \sin \phi F_A^{AB} + e \cos \cos \phi F_A^{AB} \\ &\quad - e \cos \sin \phi F_A^{AB} - e \cos \cos \phi F_A^{AB} \\ &= \cos \phi F_A^{AB} - \sin \phi F_A^{AB} \end{aligned}$$

$$J^A = \frac{m_A}{m^2} \left[m^2 \left(\cos \phi F_A^{AB} - \sin \phi F_A^{AB} \right) - m^2 \left(\cos \phi F_A^{AB} - \sin \phi F_A^{AB} \right) \right. \\ \left. + m^2 e \cos \theta_m F_B^{AB} - m^2 e \cos \theta_m F_B^{AB} \right]$$

$$\begin{aligned}
 J^\mu &= \frac{m}{m^2} \left[\cos(\pi^1) \partial_\nu F_{\bar{A}}^{\nu\mu} - m^2 \partial_\nu F_{\bar{A}}^{\mu\nu} + m^2 e^{\mu\nu\alpha\beta} \partial_\nu F_{\bar{B}}^{\nu\beta} \right] \\
 &\quad - \sin(\pi^1) \left(m^2 F_{\bar{A}\nu}^{\mu\nu} - m^2 F_{\bar{A}\nu}^{\nu\mu} + m^2 e^{\mu\nu\alpha\beta} F_{\bar{B}\nu}^{\nu\beta} \right) \\
 &= \cos \bar{J}^\mu - \sin \frac{\pi^1}{e} \bar{J}_D^\mu \quad \text{from } \text{renorm } E_9 \text{ eq}
 \end{aligned}$$

with success

$$\begin{aligned}
 \bar{J}^\mu &= \frac{m}{m^2} \left(m^2 \partial_\nu F_{\bar{A}}^{\nu\mu} - m^2 \partial_\nu F_{\bar{A}}^{\mu\nu} + m^2 e^{\mu\nu\alpha\beta} \partial_\nu F_{\bar{B}}^{\nu\beta} \right) \\
 e_D \bar{J}_D^\mu &= \frac{e m}{m^2} \left(m^2 \partial_\nu F_{\bar{A}\nu}^{\mu\nu} - m^2 \partial_\nu F_{\bar{A}\nu}^{\nu\mu} + m^2 e^{\mu\nu\alpha\beta} \partial_\nu F_{\bar{B}\nu}^{\nu\beta} \right)
 \end{aligned}$$

So the (almost) decoupled equations are

$$(1d) \quad \frac{n_\alpha}{n^2} \partial_\nu \left[n^\mu \bar{F}_A^{\alpha\nu} - n^\nu \bar{F}_A^{\alpha\mu} - n_\gamma \epsilon^{\mu\nu\lambda} {}_\beta \bar{F}_B^{\gamma\lambda} \right] = \bar{J}_\mu^\mu$$

$$(2d) \quad \frac{n_\alpha}{n^2} \partial_\nu \left[n^\mu \bar{F}_B^{\alpha\nu} - n^\nu \bar{F}_B^{\alpha\mu} + n_\gamma \epsilon^{\mu\nu\lambda} {}_\beta \bar{F}_A^{\gamma\lambda} \right] = \bar{K}_\mu^\mu + eee_D M_{B_D}^2 B_D^{2\mu}$$

$$(3d) \quad \frac{n_\alpha}{n^2} \partial_\nu \left[n^\mu \bar{F}_{A_D}^{\alpha\nu} - n^\nu \bar{F}_{A_D}^{\alpha\mu} - n_\gamma \epsilon^{\mu\nu\lambda} {}_\beta \bar{F}_{B_D}^{\gamma\lambda} \right] = \bar{J}_\mu^{D\mu} + M_{A_D}^2 A_D^{2\mu}$$

$$(4d) \quad \frac{n_\alpha}{n^2} \partial_\nu \left[n^\mu \bar{F}_{B_D}^{\alpha\nu} - n^\nu \bar{F}_{B_D}^{\alpha\mu} + n_\gamma \epsilon^{\mu\nu\lambda} {}_\beta \bar{F}_{A_D}^{\gamma\lambda} \right] = \bar{K}_\mu^{D\mu} + M_{B_D}^2 B_D^{2\mu}$$

Perhaps we can think of $eee_D M_{B_D}^2 B_D^{2\mu}$ as some residual mass term from the dark sector into the visible one.
 (unfortunately it depends on B_D^μ :))