Current Indirect Detection Searches

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Les Houches Summer School on Dark Matter Lecture 4

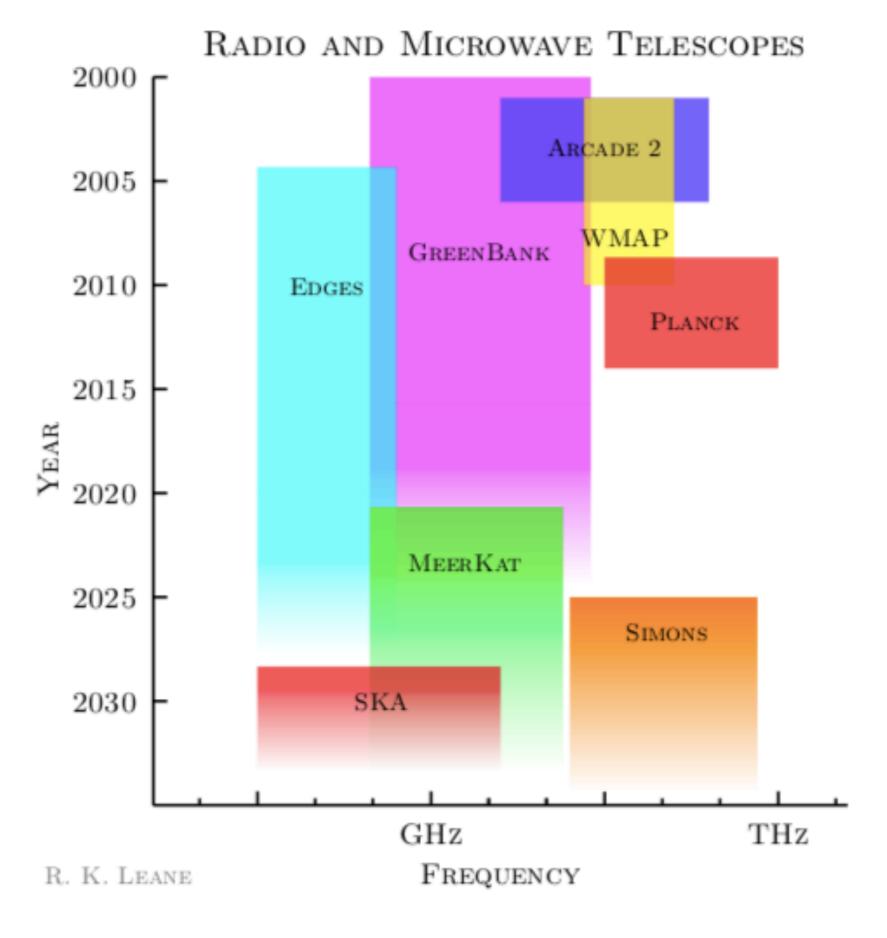


Goals for this section

- Know the general landscape of upcoming telescopes relevant for indirect detection (and where to look for more information)
- Discuss the strengths and weaknesses of different targets and channels for indirect detection
- Summarize leading indirect-detection constraints for both annihilation and decay, across a broad range of DM masses and final states

X-Ray and Gamma-Ray Telescopes 1990 1995 Comptel EGRET 2000 XMM Newton Integral Chandra 2005Suzaku Magic HESS 2010 NuStar FERMI VERITAS 2015HAWC 2020 Micro-X LHAASO X-Prism 2025EXTP CTA HERD 2030 AMEGO E-ASTROGAM ATHENA 2035 LYNX 2040 KEV PeVMeV GeV TeV R. K. Leane Energy

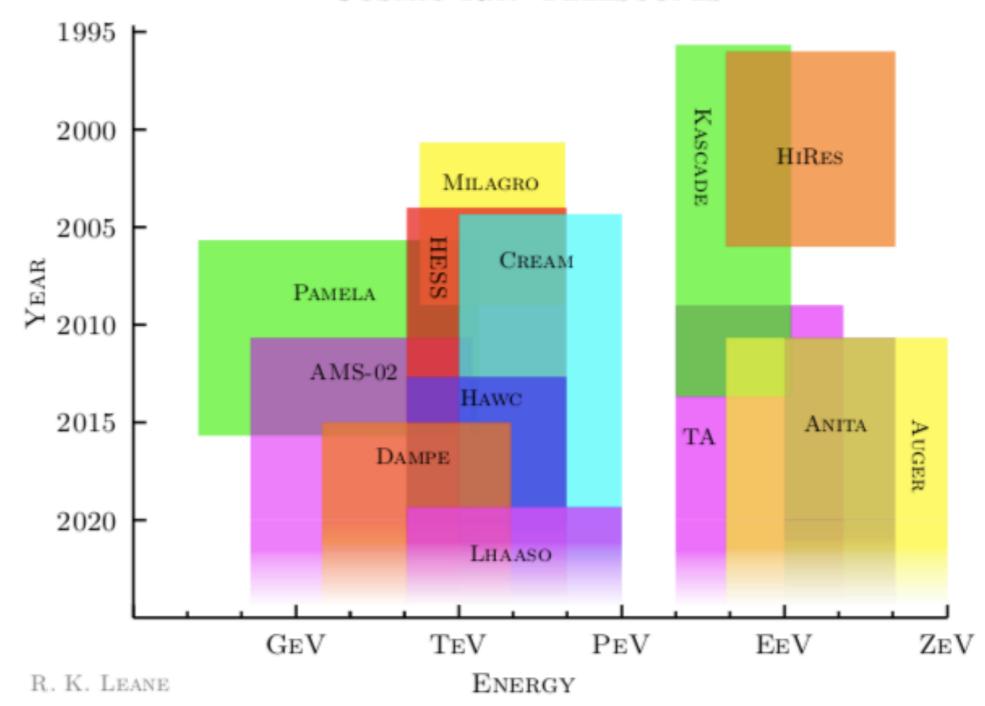
- At present, the Fermi Gamma-Ray Space Telescope has the highest sensitivity for O(0.1-100) GeV gamma rays
- At higher energies, groundbased telescopes such as HESS, VERITAS, HAWC, MAGIC take over due to larger area
- At lower energies, there are a number of sensitive X-ray experiments



- At even lower energies, radio and microwave telescopes measure the CMB
- These telescopes are also sensitive to:
 - synchrotron signals from electrons/ positrons
 - primordial 21cm radiation
 - pulsars, which serve as major backgrounds for many DM signals

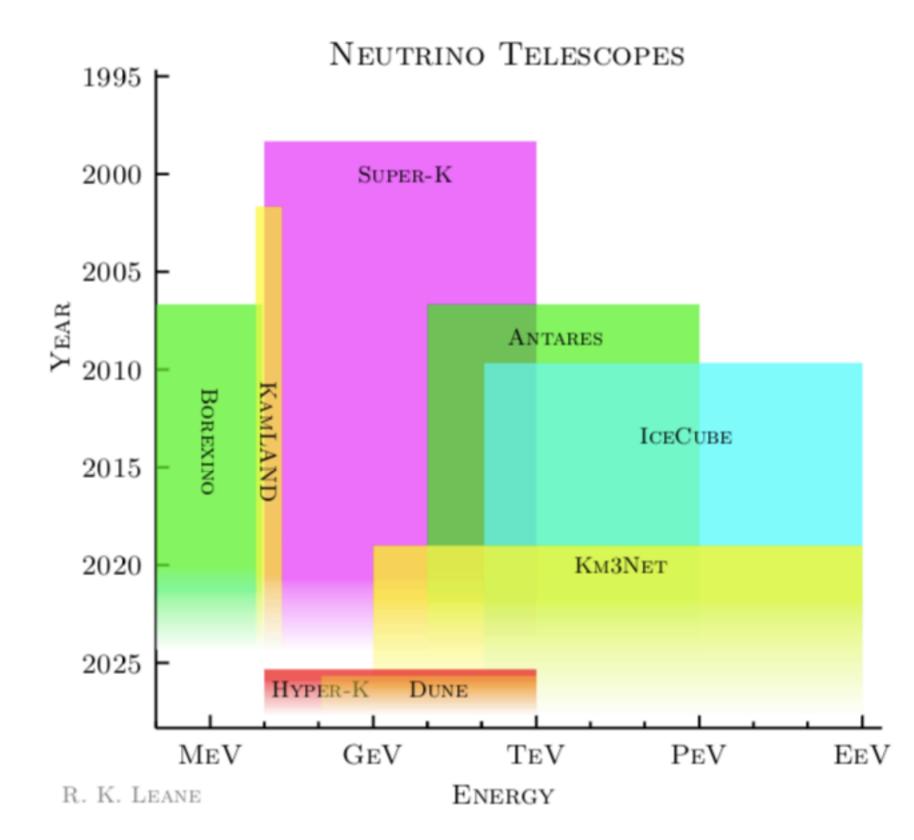
- AMS-02 reaches relatively low energies + has a magnet, allowing for charge discrimination
- AMS-02 typically sets most sensitive limits for classic WIMP DM
- Not on this plot:
 Voyager!
 Sensitive to
 even lower energy cosmic
 rays relevant
 for sub-GeV DM





 At very high energies, non-detection of CRs sets limits on ultraheavy decaying DM

- **IceCube** (instrumented ice - South Pole) and **ANTARES** (instrumented water -Mediterranean) currently lead limits for decay/ annihilation to neutrinos
- Become
 competitive with
 gamma-ray
 searches for
 hadronic channels
 + sufficiently
 heavy DM

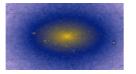


 Telescopes searching for lower-energy neutrinos can have sensitivity to DM capture in the Sun, "boosted DM" models where dark particles are produced relativistically

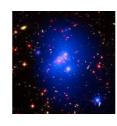
Dwarf galaxies



Galactic halo



Other galaxies and clusters

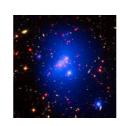


Dark matter subhalos



Dwarf galaxies
 low background, nearby

- Galactic center
- Galactic halo
- Other galaxies and clusters

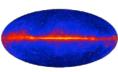


Dark matter subhalos



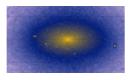
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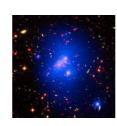


high signal, high background, sensitive to presence of density cusp/core

Galactic halo



Other galaxies and clusters



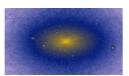
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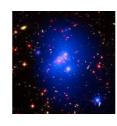


Galactic halo



large area, nearby, complex backgrounds

Other galaxies and clusters



Dark matter subhalos

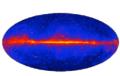


Dwarf galaxies



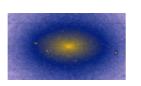
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Galactic center



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Galactic halo



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Other galaxies and clusters



large dark matter content, (potentially) hold redshift information, sensitive to amount of substructure

Dark matter subhalos

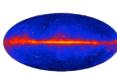


Dwarf galaxies



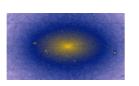
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Dark matter subhalos



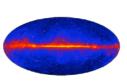
potentially numerous, probe small-scale structure

Dwarf galaxies



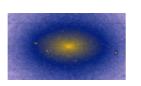
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Dark matter subhalos



potentially numerous, probe small-scale structure

Extragalactic background radiation

holds redshift information, probes halos at all scales

Return to cosmological limits

- We can work out the detailed version of the cosmological energyinjection bounds described in Lecture 1 (e.g. from CMB, Lyman-alpha)
- As for other signals, this requires us to work out:
 - the spectrum of stable SM particles produced by annihilation/decay
 - how that injected spectrum translates into an observable
- The first part is done as for other indirect searches
- The second part requires modeling how high-energy particles cool and deposit their energy in the early universe

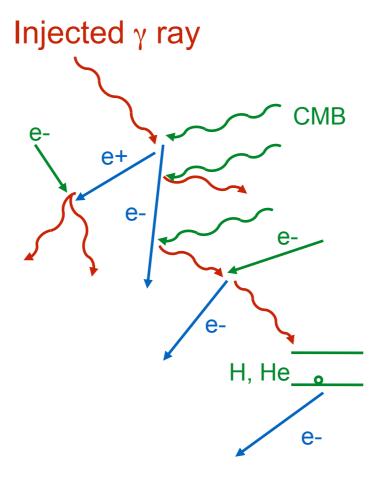
The energy deposition calculation (prelude)

- Simplifying approximations:
 - consider only electrons, positrons, photons (neutrinos are ~non-interacting, (anti)protons tend to give only a small contribution to energy deposition)
 - assume linearity: injected particles are rare enough they will not interact with each <u>other</u>, only with the background (although over time the extra heating/ionization may modify the background)
 - these two assumptions reduce the problem to considering the behavior of individual electrons/positrons/photons injected at some energy+redshift
- Note: this is purely SM physics all the DM physics goes into setting the spectrum and redshift-dependence of the particle injection

The energy deposition calculation

ELECTRONS

- Inverse Compton scattering on the CMB.
- Excitation, ionization, heating of electron/H/He gas.
- Positronium capture and annihilation.
- All processes fast relative to Hubble time: bulk of energy goes into photons via ICS.



blue/red = e+e-/photons carrying injected energy, green=background

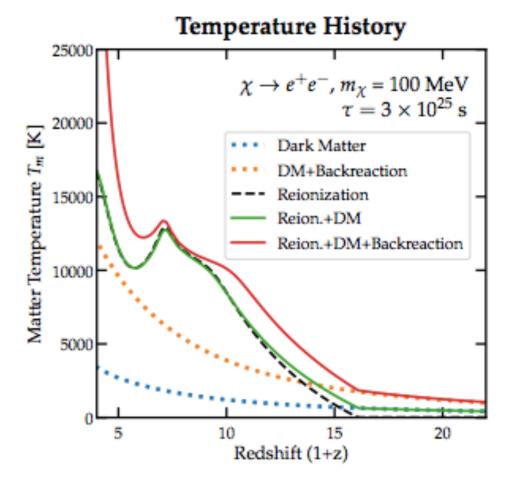
PHOTONS

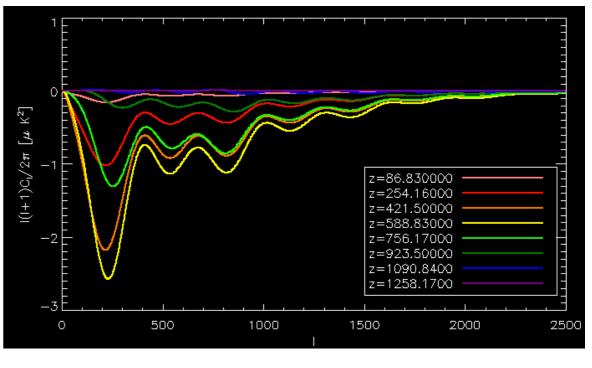
- Pair production on the CMB.
- Photon-photon scattering.
- Pair production on the H/ He gas.
- Compton scattering.
- Photoionization.
- Redshifting is important, energy can be deposited long after it was injected.

From energy deposition to

a signal

- Modeling this cascade tells us the power going into ionization, excitation, heating, distortions to CMB energy spectrum
- Results are tabulated (assuming standard background) in TRS 1506.03812; DarkHistory package [Liu, TRS et al '19] allows for variable ionization history + backreaction effects
- Feed these inputs into evolution equations for ionization/temperature (in public recombination codes RECFAST, HyREC, CosmoRec, DarkHistory)
- Compare temperature forecast to measurements, project CMB perturbations due to modified ionization history (via public codes CAMB, CLASS / ExoCLASS)

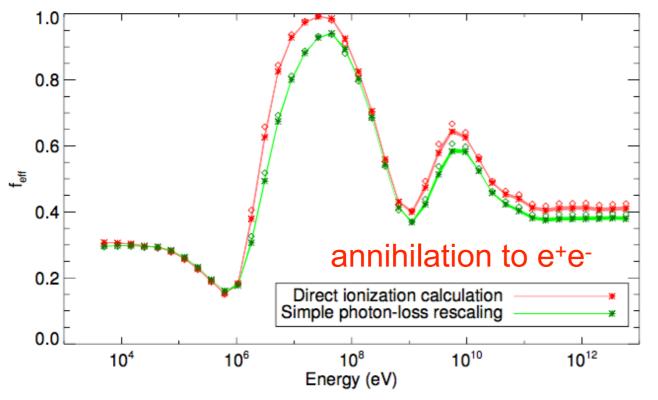


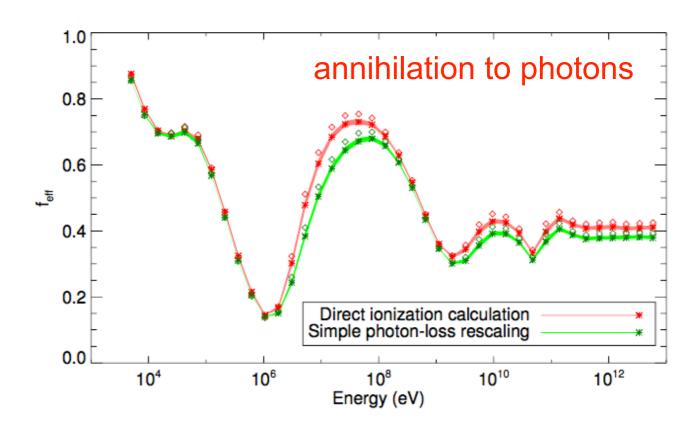


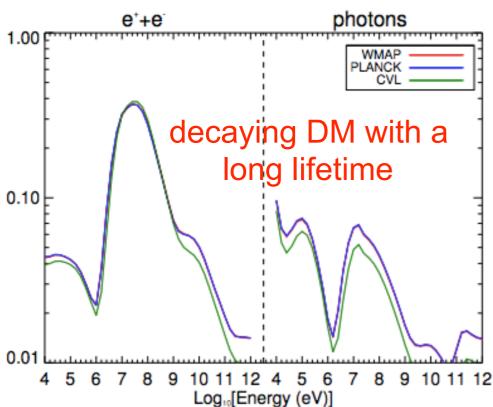
Signals of annihilation/decay in the CMB

- We can take linear combinations of results for single photons/e+/e- to work out the CMB perturbation for a given particle spectrum
- When we do this, we find that different annihilating DM models (with different spectra) give ~identical perturbations to the CMB anisotropy spectrum, up to a rescaling factor (f_{eff} = effective deposition efficiency)
- Decaying models are also very similar to each other
- Intuition: shape of CMB perturbation is mostly fixed by redshift-dependence of energy deposition - mostly controlled by (1+z)⁶ or (1+z)³ density-based scaling
- From projected perturbation to CMB, read off the f_{eff}(E) for photons/electrons/ positrons produced by DM
- CMB experimentalists can check for the distinctive imprint of an energy injection with rate scaling as (1+z)⁶ or (1+z)³ in the CMB, constrain its coefficient experimentally

Effective deposition efficiencies

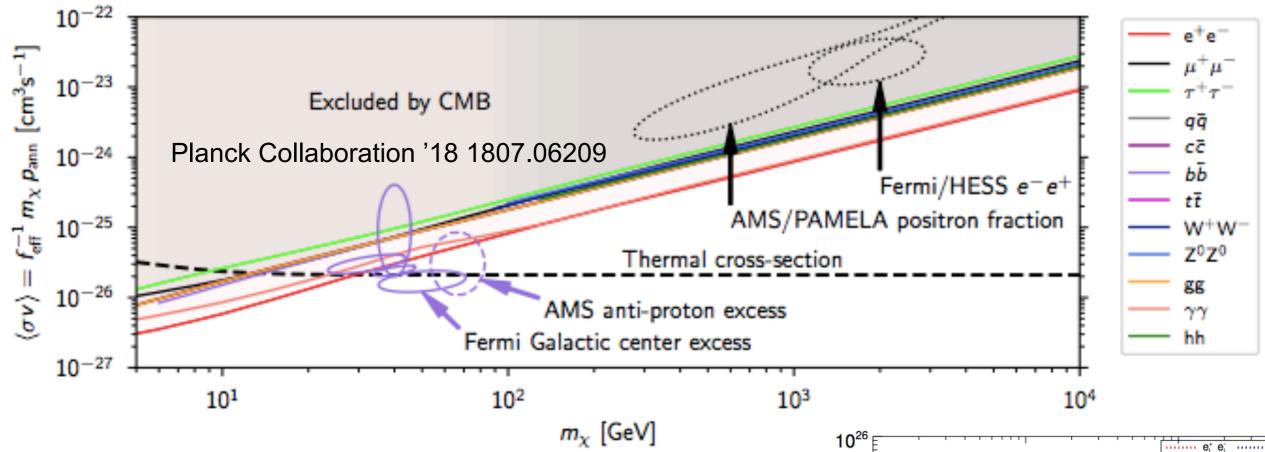






- We can get f_{eff} for a specific DM model by weighting these curves by the photon/e+e- spectra
- Planck '18 sets an upper bound for annihilation of $f_{\rm eff} \langle \sigma v \rangle / m_{\rm DM} < 3.2 \times 10^{-28} {\rm cm}^3/{\rm s/GeV}$
- In 1610.06933 we set a limit using 2015 Planck data of $\tau \gtrsim \frac{f_{\rm eff}}{f_{\rm eff}(30{\rm MeV}e^+e^-)} \times 2.6 \times 10^{25}s$

Limits on annihilation and decay



- At lower masses the available final states become limited, but down to keV masses:
 - the limit on xsec continues to scale as ~m_{DM} for annihilation,
 - the limit on lifetime stays ~constant for decay
- Temperature limits from Lyman-alpha are comparable to the CMB bounds for decay - weaker for annihilation

