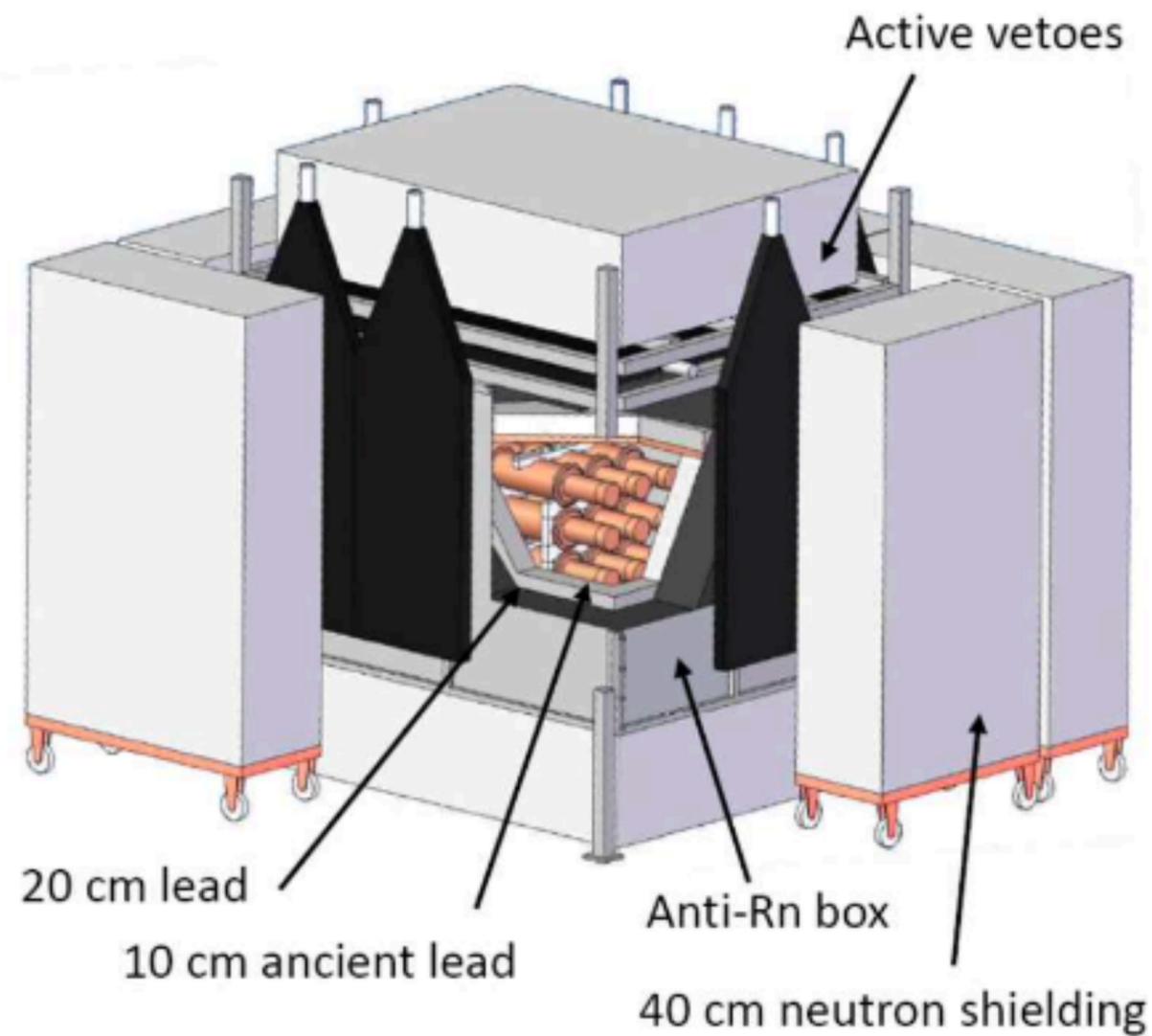


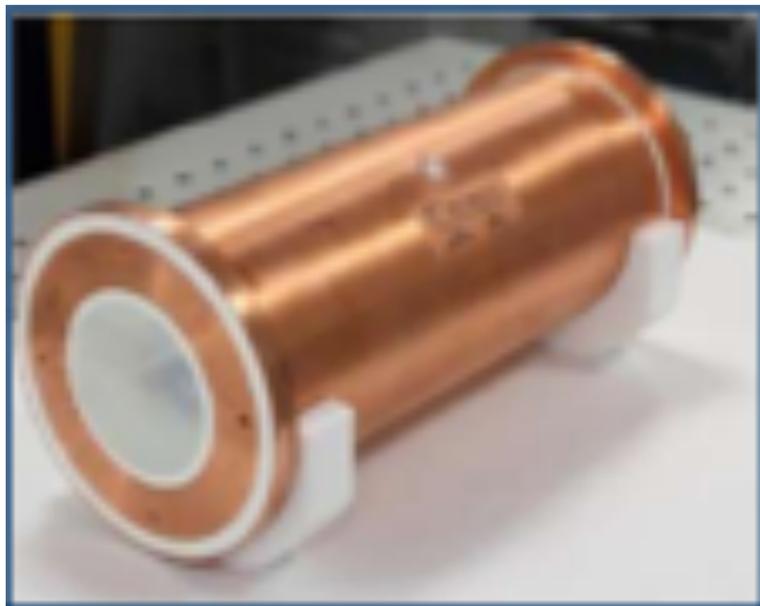
DARK MATTER DIRECT DETECTION OF CLASSICAL WIMPS – LECTURE 3

Jodi Cooley
SMU

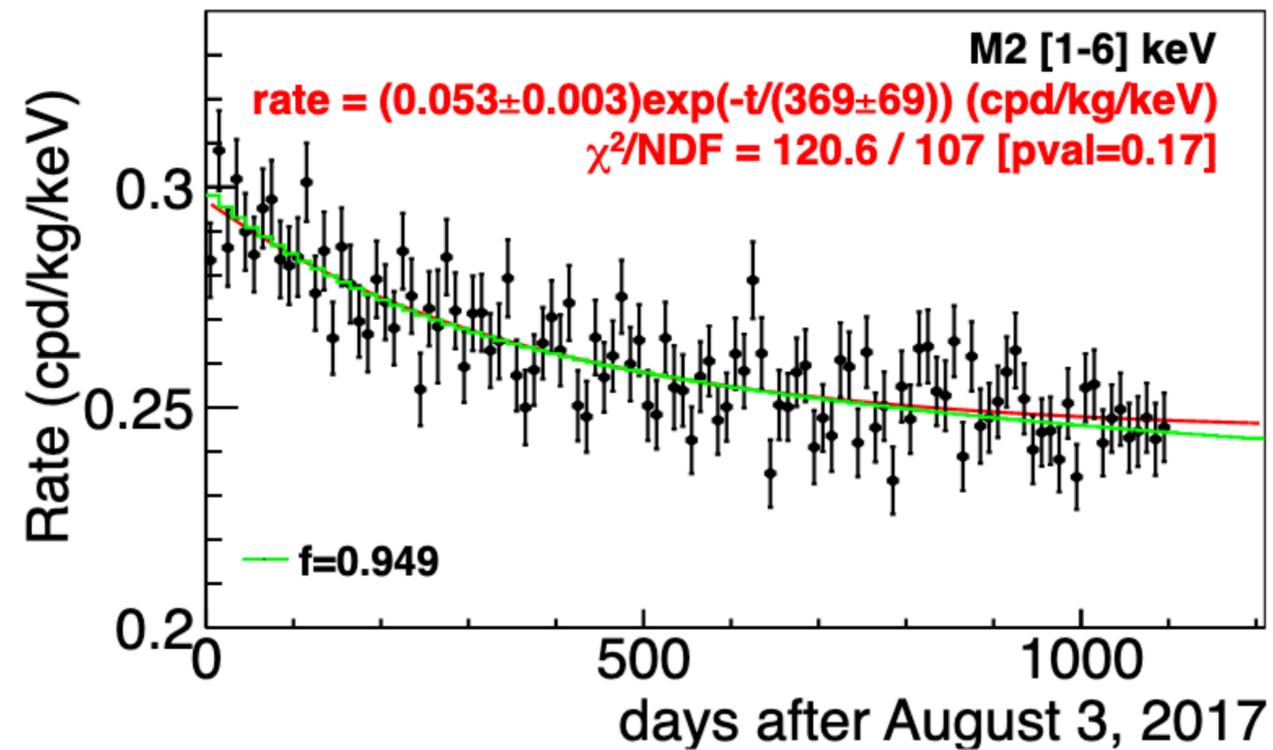
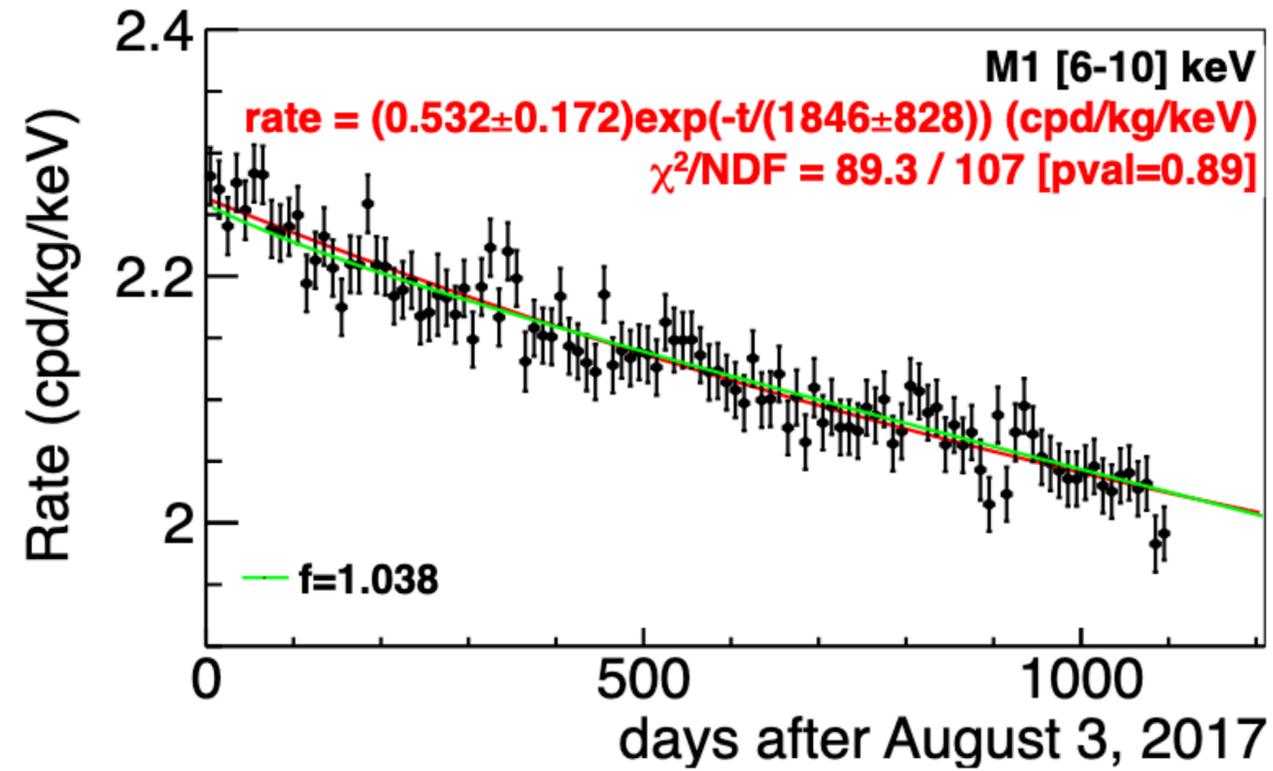
ANAIS 112



- Located in Hall B at the Canfranc Laboratory (2450 mwe).
- NaI(Tl) crystals (12.5 kg each) grown from ultra pure NaI powder and housed in OFE copper.
- 112.5 kg of NaI(Tl), distributed in a 3×3 array of modules.
- Mylar window for low energy calibration
- Two Hamamatsu R12669SEL2 photomultipliers
- Low background, high quantum efficiency.



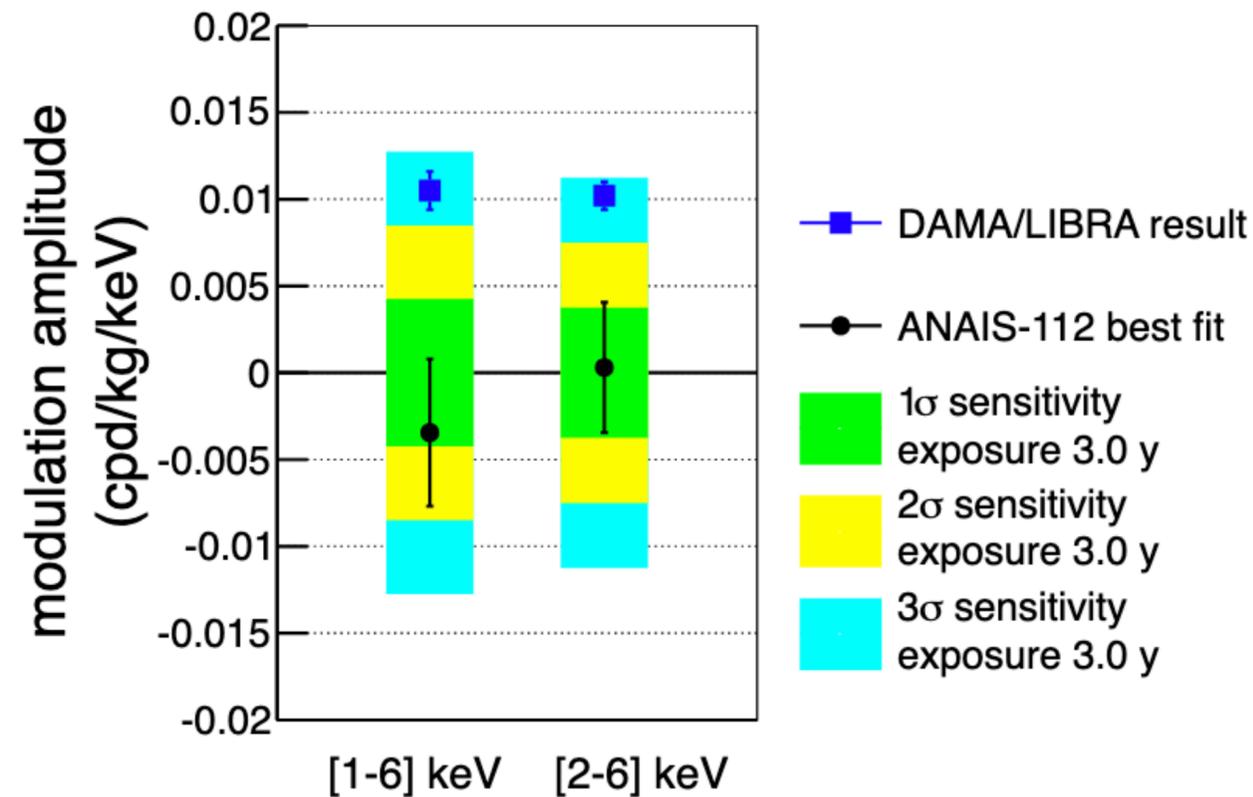
ANAIS112: 3-YEAR BACKGROUND MODELS:



- Three independent background modeling procedures:
 - Exponentially decaying background
 - Probability distribution function derived from background mode
 - Probability distribution function for every detector to account for possible systematic effects related with the different backgrounds and efficiencies of the different modules.

ANAIS 112: 3 YEAR RESULTS

Energy region	Model	χ^2 /NDF null hyp	nuisance params	S_m cpd/kg/keV	p-value mod	p-value null
[1-6] keV	1	132 / 107	3	-0.0045 ± 0.0044	0.051	0.051
	2	143.1 / 108	2	-0.0036 ± 0.0044	0.012	0.013
	3	1076 / 972	18	-0.0034 ± 0.0042	0.011	0.011
[2-6] keV	1	115.7 / 107	3	-0.0008 ± 0.0039	0.25	0.27
	2	120.8 / 108	2	0.0004 ± 0.0039	0.17	0.19
	3	1018 / 972	18	0.0003 ± 0.0037	0.14	0.15



- Data support the absence of modulation in both energy region and three background models.
- Best fits are incompatible with DAMA/LIBRA at 3.3σ in the [1-6] keV region and 2.6σ in the [2-6]keV region

LIQUID NOBLE EXPERIMENTS

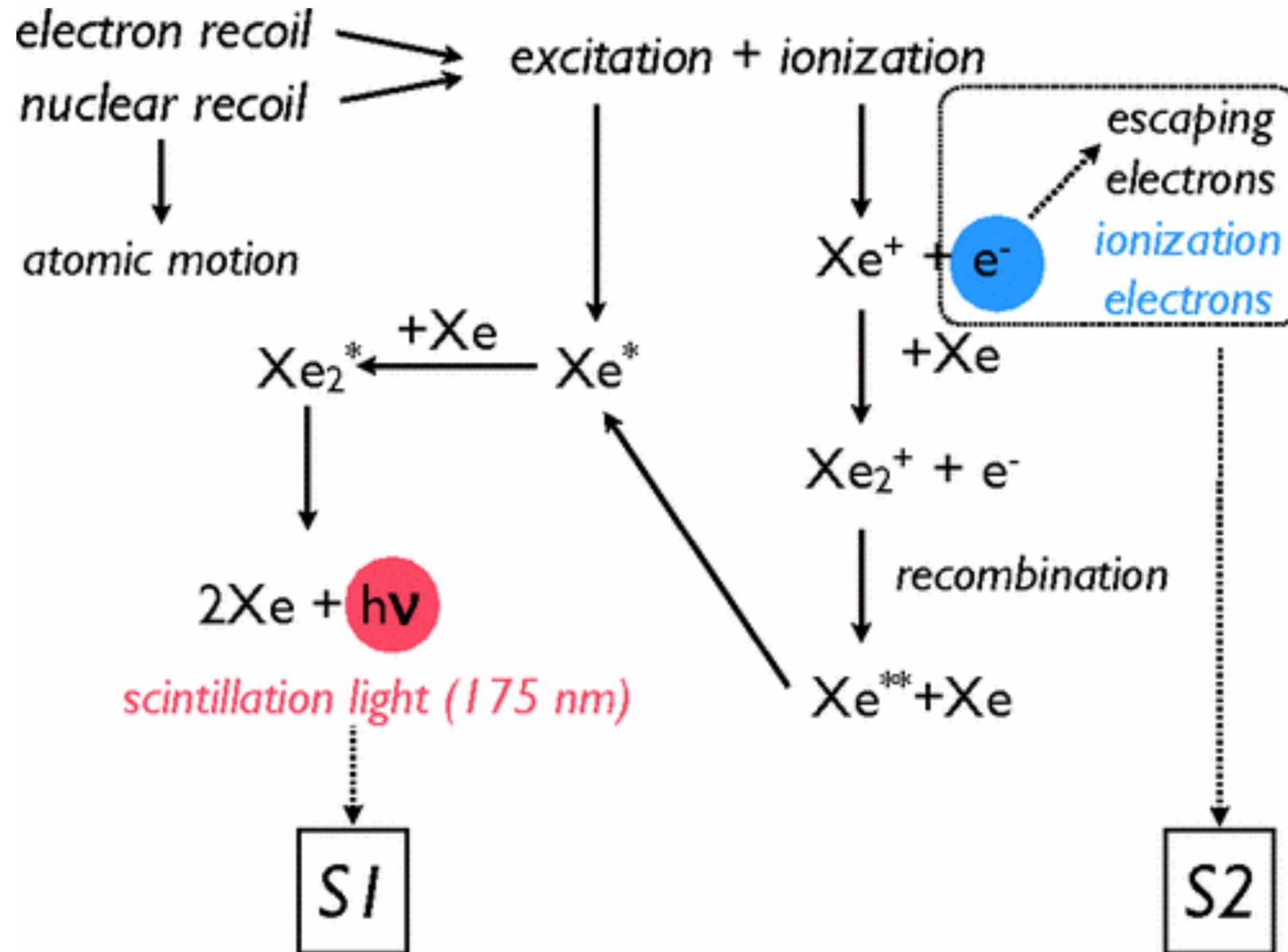
LIQUID NOBLE PROPERTIES

Property (unit)	Xe	Ar	Ne
Atomic Number	54	18	10
Mean relative atomic mass	131.3	40.0	20.2
Boiling Point T_b (K)	165.0	87.3	27.1
Melting Point T_m (K)	161.4	83.8	24.6
Liquid density at T_b (g cm^{-3})	2.94	1.40	1.21
Volume fraction in Earth's atmosphere (ppm)	0.09	9340	18.2
Scintillation light wavelength (nm)	175	128	78
Triplet lifetime (ns)	27	1600	15000
Singlet lifetime (ns)	3	7	<18
Electron mobility ($\text{cm}^2 \text{V}^{-1} \text{s}^{-1}$)	2200	400	low
Scintillation yield (photons/keV)	42	40	30

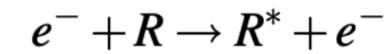
Material	Ar	Kr	Xe
Gas			
Ionization potential I (eV)	15.75	14.00	12.13
W values (eV)	26.4 ^a	24.2 ^a	22.0 ^a
Liquid			
Gap energy (eV)	14.3	11.7	9.28
W value (eV)	23.6±0.3 ^b	18.4±0.3 ^c	15.6±0.3 ^d

- ▶ Three different noble liquids have been considered for dark matter detection over the past few decades.
- ▶ Properties of the noble liquids determine many practical aspects of the detectors. For example, Xe has a high density and a large target mass (favorable) but it is not very abundant in the atmosphere (more expensive).
- ▶ The energy loss of an incident particle in noble liquids is shared between excitation, ionization and sub-excitation electrons liberated in the ionization process
- ▶ The average energy loss in ionization is slightly larger than the ionization potential or the gap energy, because it includes multiple ionization processes
- ▶ As a result, the ration of the W -value (= average energy required to produce an electron-ion pair) to the ionization potential or gap energy = 1.6 - 1.7

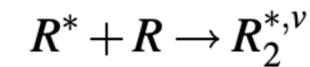
LIQUID NOBLE SIGNAL PRODUCTION



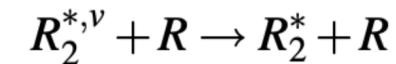
- Energy is transferred to a particle by excitation, ionization or heat (atomic motion).



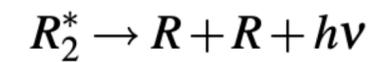
impact excitation



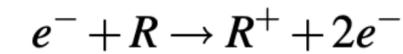
excimer formation



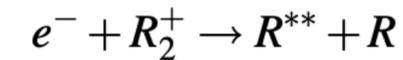
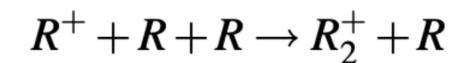
relaxation



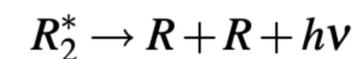
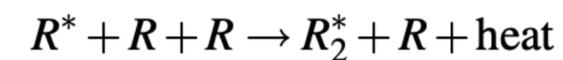
VUV emission



ionization

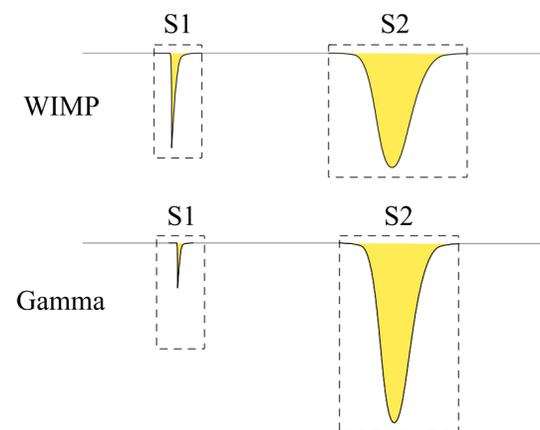
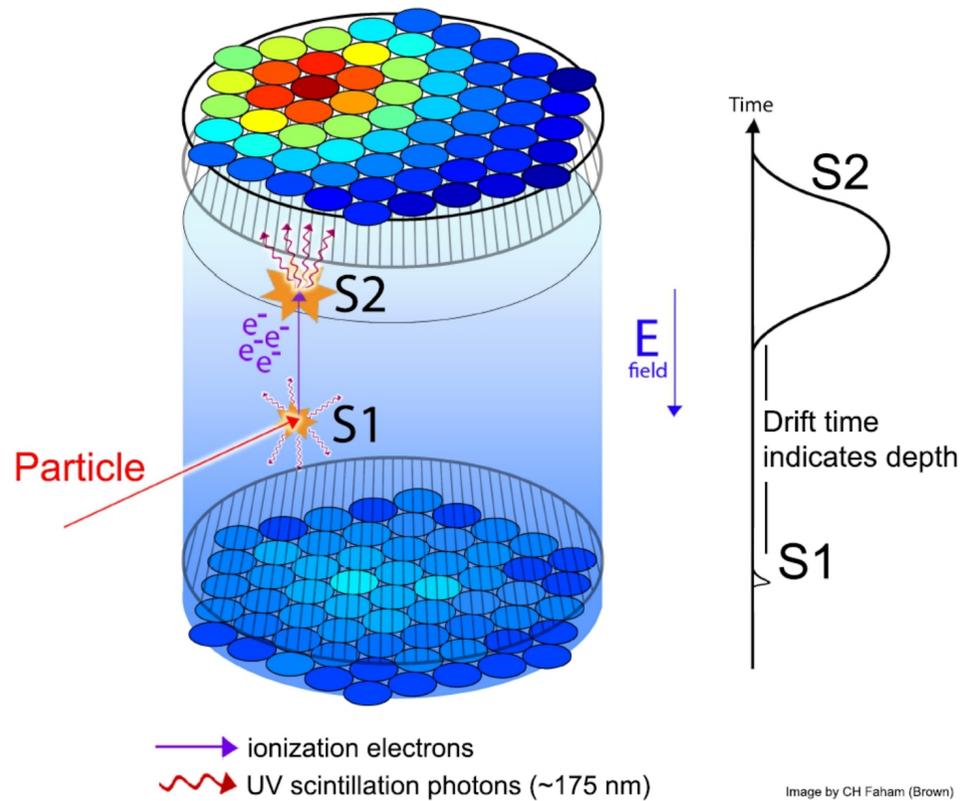


recombination



VUV emission

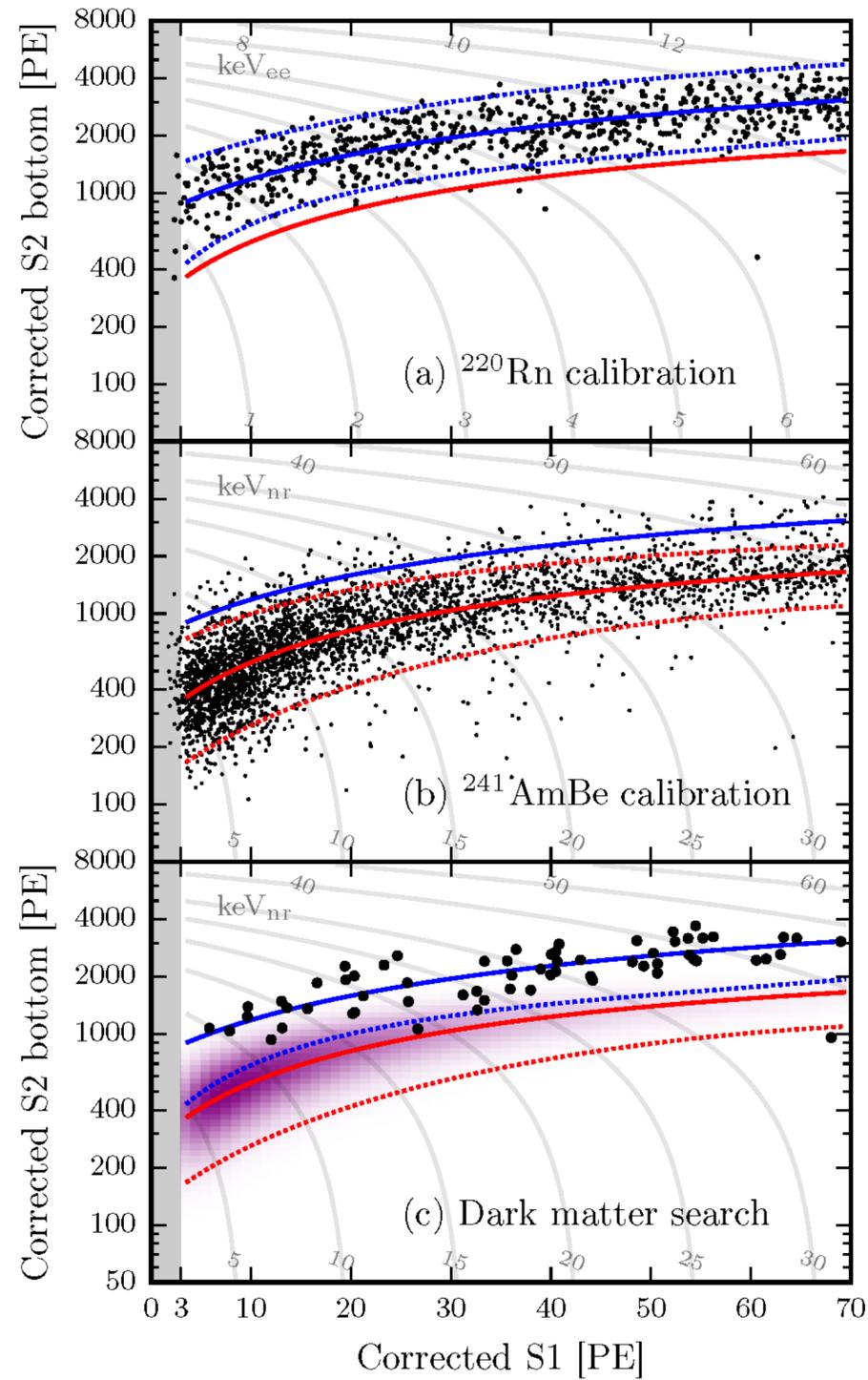
LIQUID NOBLE DETECTORS



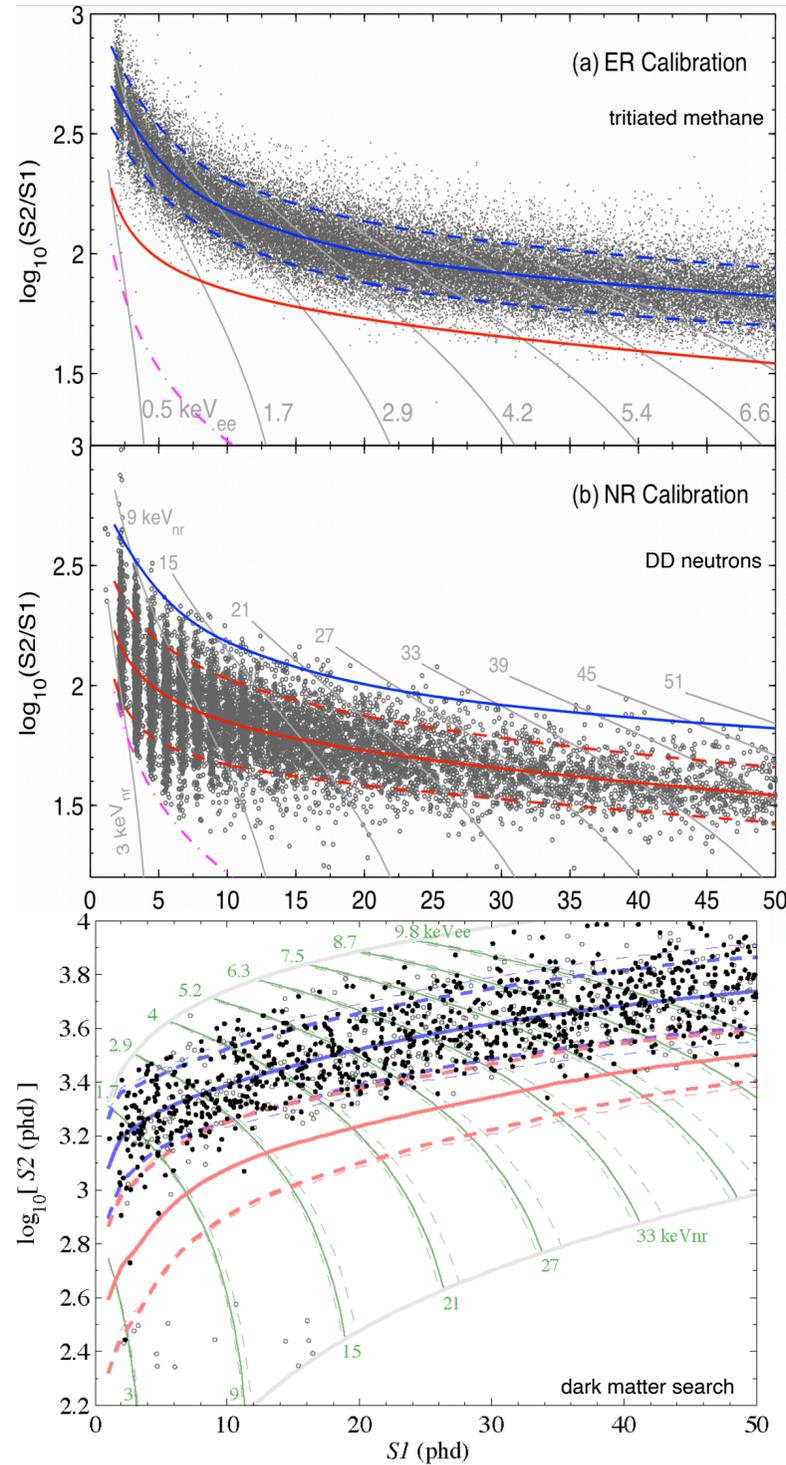
Dual Phase TPCs (XENON, LUX/LZ, Darkside PandaX, etc)

- Interactions in the liquid produce excitation and ionization.
- Excitation leads to scintillation light emission
- Ionization electrons are drifted with an applied electric field into the gas phase (S1).
- In the gas phase, electrons are further accelerated producing proportional scintillation (S2).
- PMTs on the bottom and top of the chamber record scintillation signals.
- Distribution of S2 give xy coordinates, drift time gives z coordinates
- Ratio of S2/S1 discriminates electron and nuclear recoils

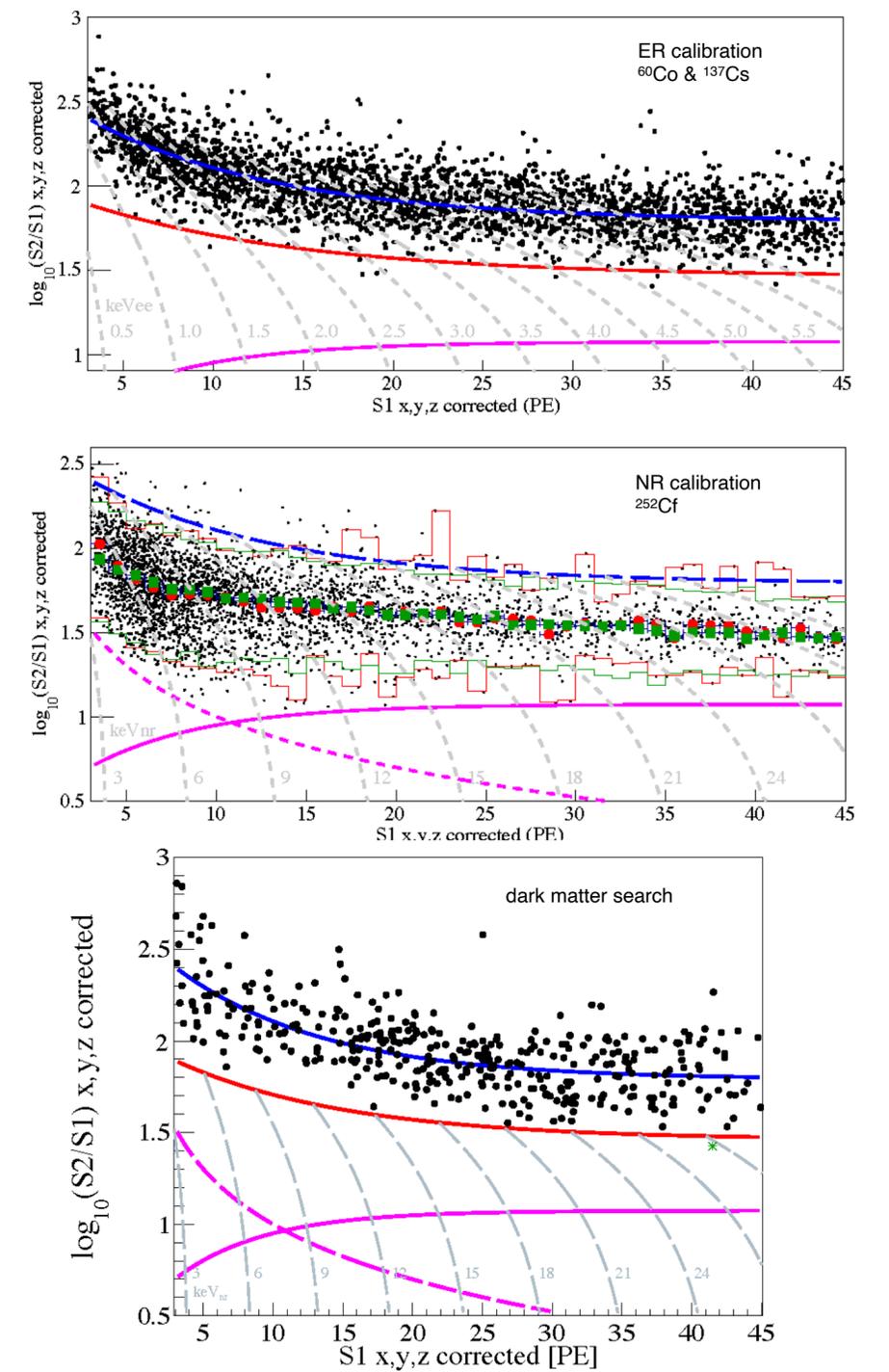
Xenon1T



LUX



PandaX-II



ENERGY

Nuclear recoils are measured through a combination of scintillation light and ionization. The nuclear recoil energy is related to S1 by

$$E_{nr} = \frac{S1}{L_y L_{eff}} \times \frac{S_e}{S_r}$$

[keV_{nr}] → E_{nr} observed scintillation [PE] → S1
 light yield [PE/keV_{ee}] → L_y scintillation efficiency of NR in LXe → L_{eff}
 suppression of scintillation signal from electric field for ER and NR events → $\frac{S_e}{S_r}$

L_{eff} accounts for the quenching of the scintillation signal for a nuclear recoil.

$$L_{eff} \equiv \frac{S1(E_{nr})/E_{nr}}{S1(122keV_{ee})/122keV_{ee}}$$

122 γ line from ⁵⁷Co source

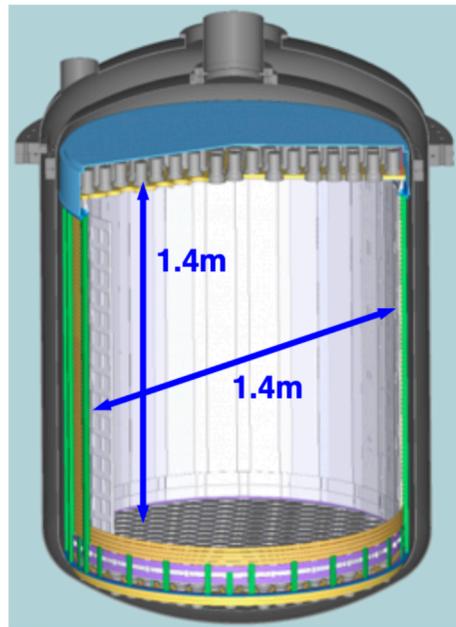
The nuclear recoil energy is related to S2 by

$$E = \frac{S2}{Y} \frac{1}{Q_y(E)}$$

[keV_{nr}] → E observed scintillation [PE] → S2
 secondary amplification factor [pe/e] → Y number of free electrons per unit energy → $Q_y(E)$

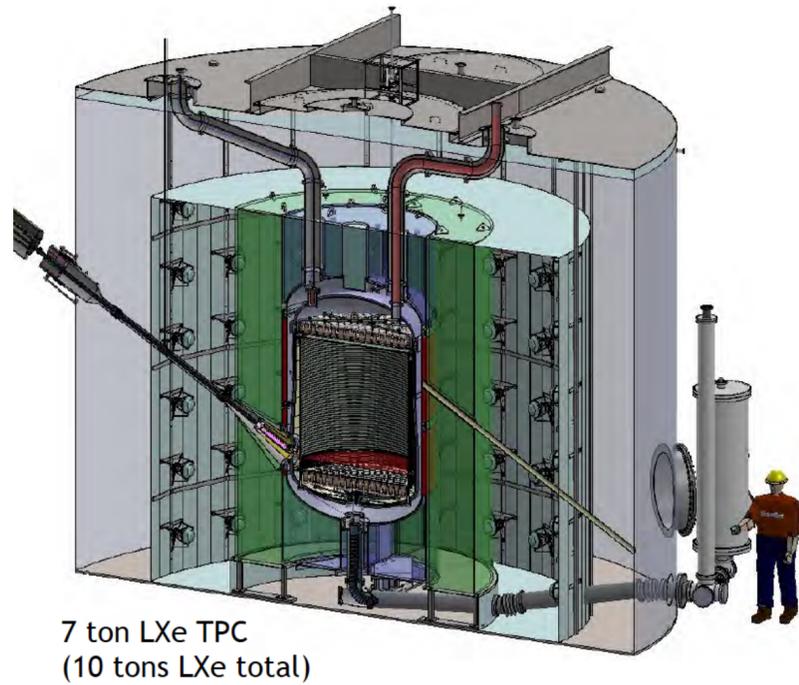
STATUS OF CURRENT TPC DARK MATTER EXPERIMENTS

XENONnT



Data Taking
2019-2025
8t LXe

LZ



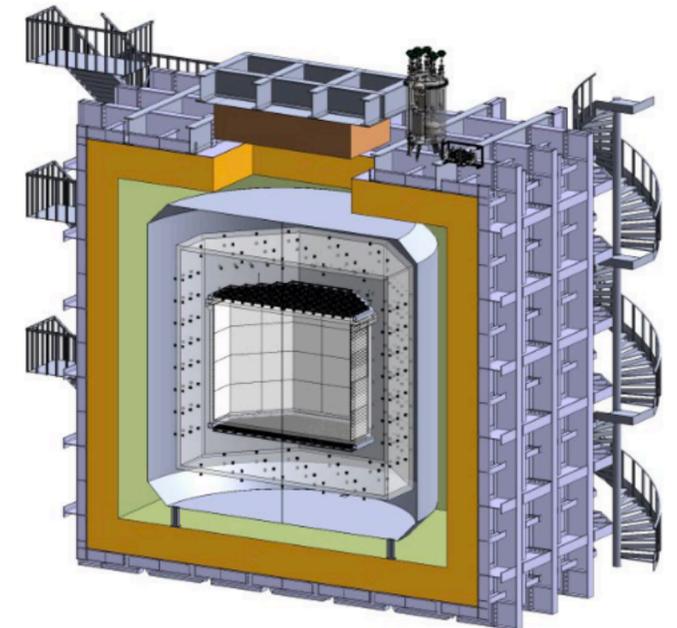
Commissioning
2019-2025 ?
7t LXe

PandaX-4T



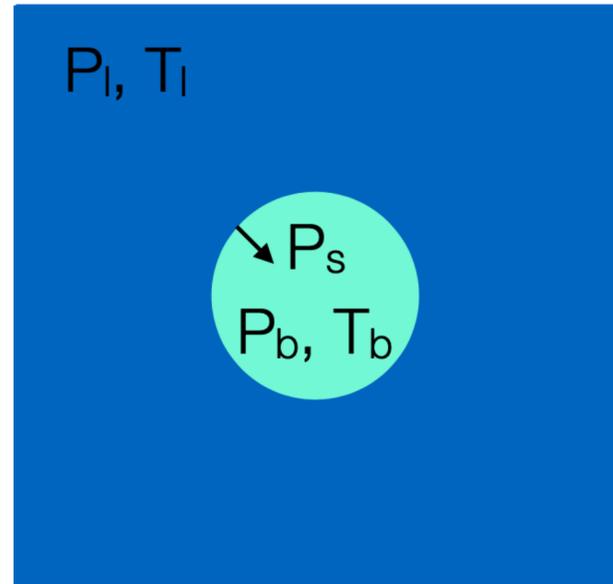
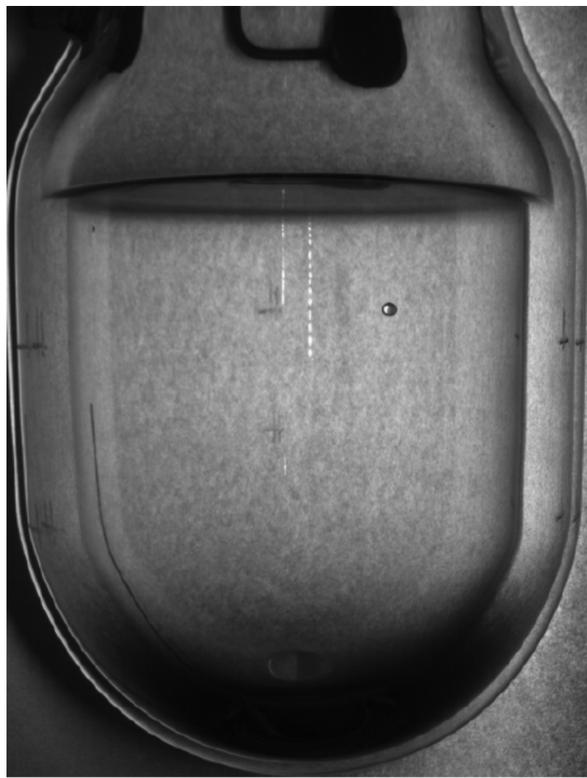
Commissioning
2019 - ?
4t LXe

DarkSide-20K



Under Design
50t LAr

BUBBLE CHAMBERS



HOW DO BUBBLE CHAMBERS WORK?

- Start with a bubble in a liquid in thermal and chemical equilibrium

$$T_l = T_b$$

- If $P_b > P_l$ the bubble will expand (assuming no surface tension).

- Include surface tension, $P_s = 2\sigma/r$, bubble grows when

$$P_b > P_l + P_s$$

and $r > r_c = \frac{2\sigma}{P_b - P_l}$

- Bubbles that do not meet this criteria collapse

- The threshold for bubble nucleation is given by

$$E_T = r\pi r_c^2 \left(\sigma - T \left[\frac{d\sigma}{dT} \right]_{\mu} \right) + \frac{4\pi}{3} r_c^3 \rho_b (h_b - h_l) - \frac{4\pi}{3} r_c^3 (P_b - P_l)$$

surface energy

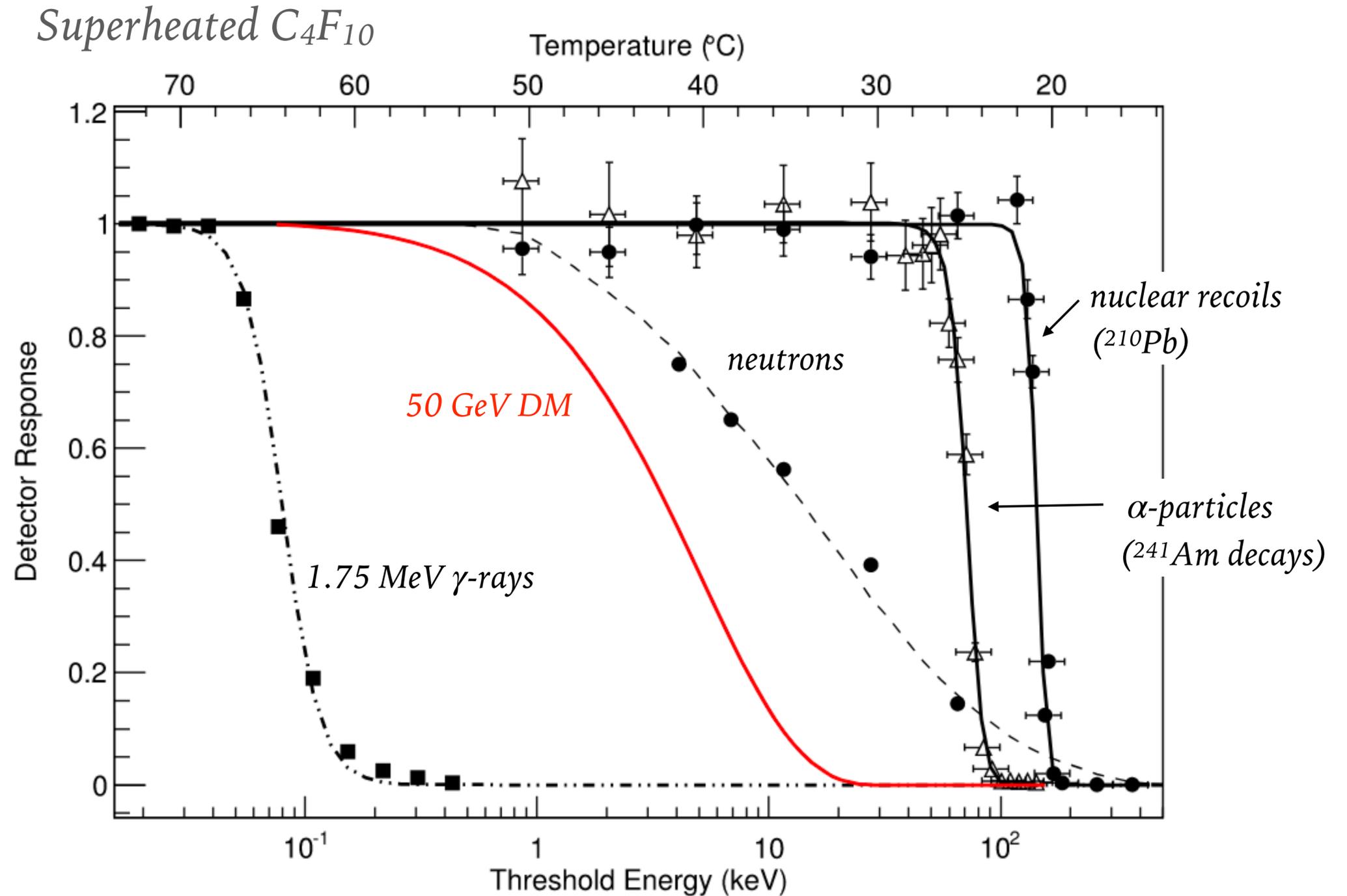
bulk energy

reversible work

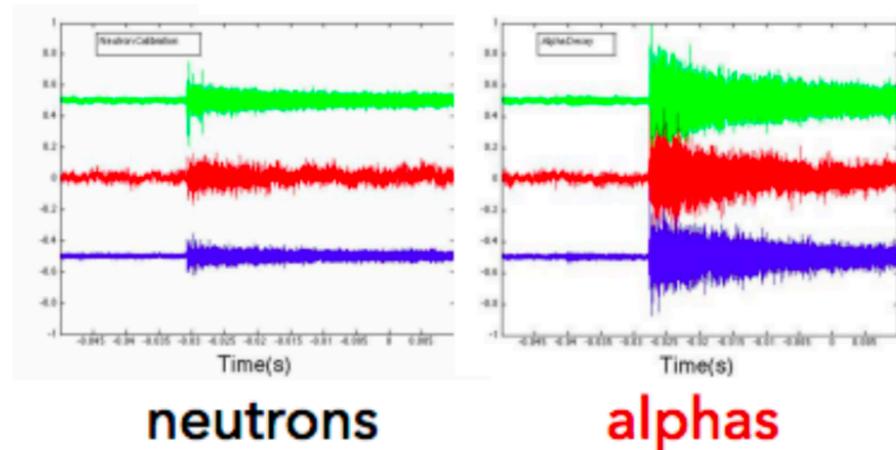
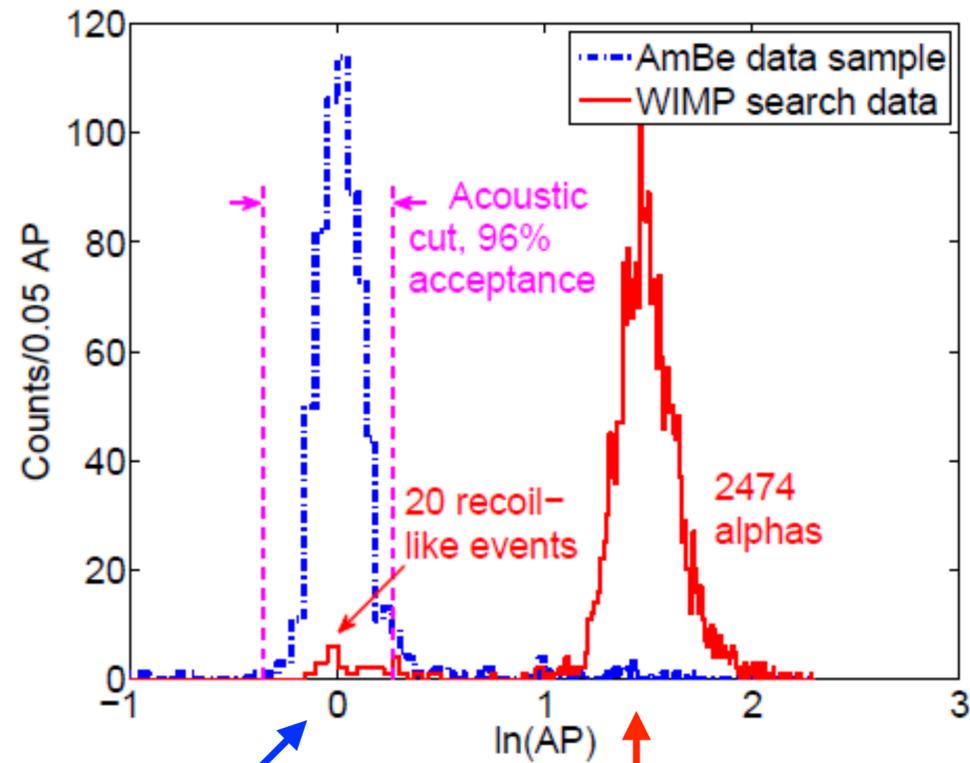
$\rho = \text{density}$ and $h = \text{specific heat}$

DETECTOR RESPONSE

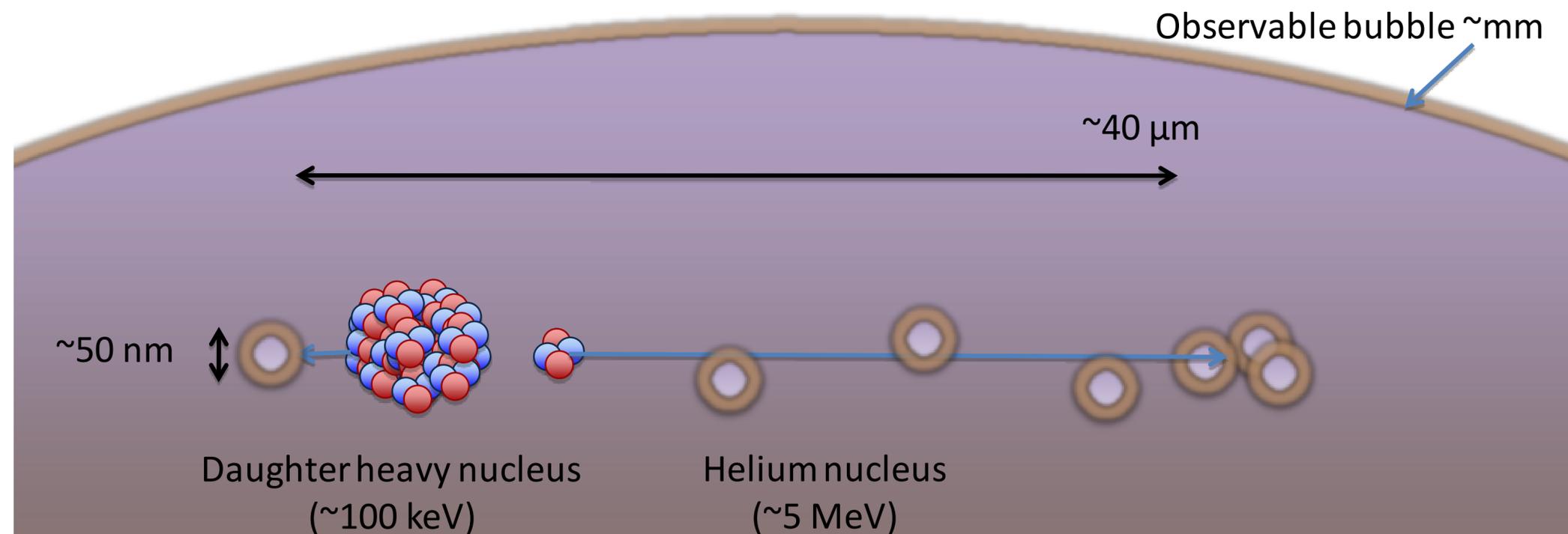
- Heavier particles have higher thresholds
- Tune the chamber to be unresponsive to most backgrounds (ER).
- Underground location and shielding to mitigate neutrons.
- But what about alphas?



ACOUSTIC DISCRIMINATION

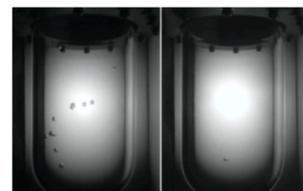
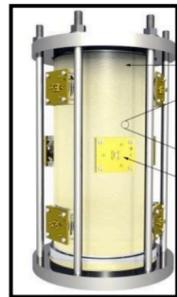


- Alphas deposit their energy over 10s of microns
- Nuclear recoils deposit their energy over 10s of nanometers
- Alpha particles are ~ 4 times louder than NR. This can be measured by piezoelectric sensors



PICO PROGRAM

PICASSO

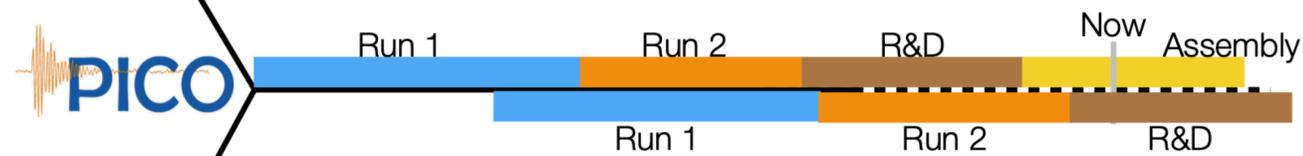


COUPP

PICO-2L



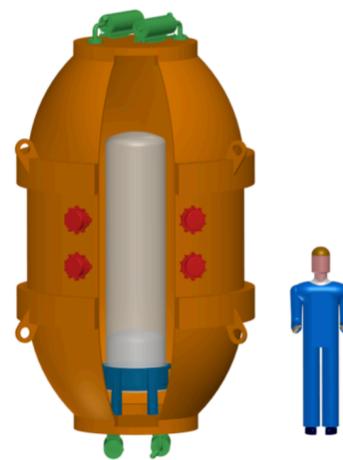
PICO-40L



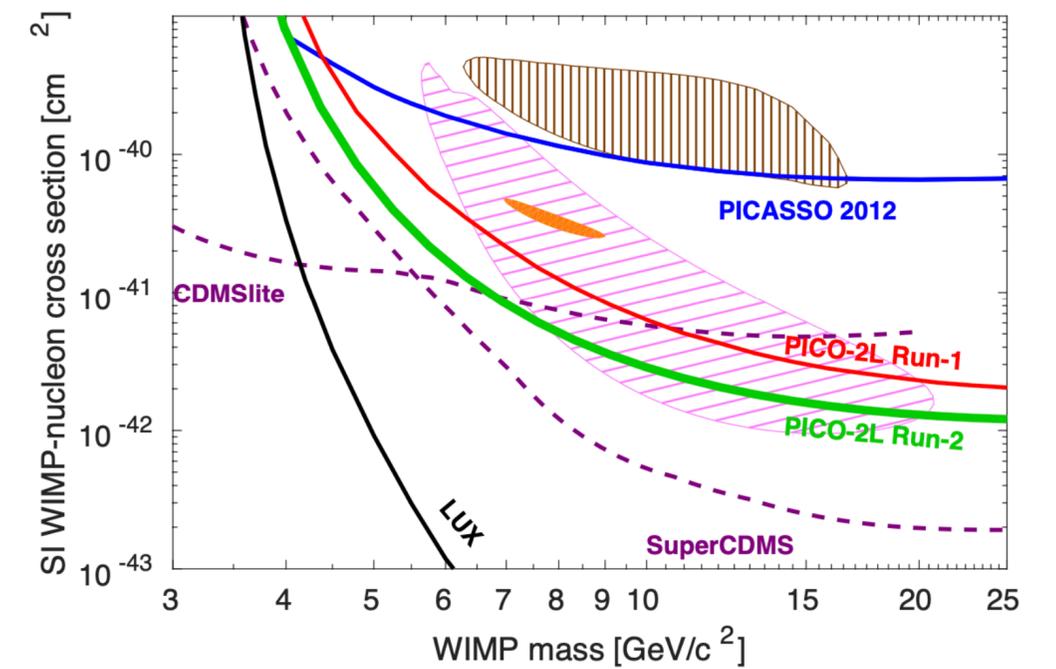
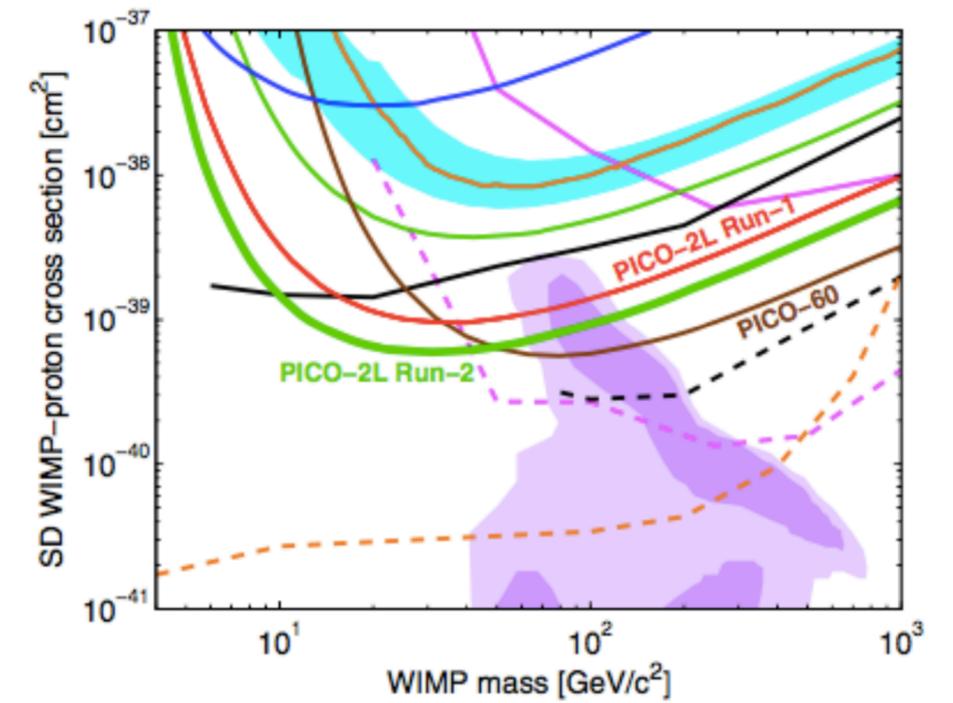
PICO-60



PICO-500



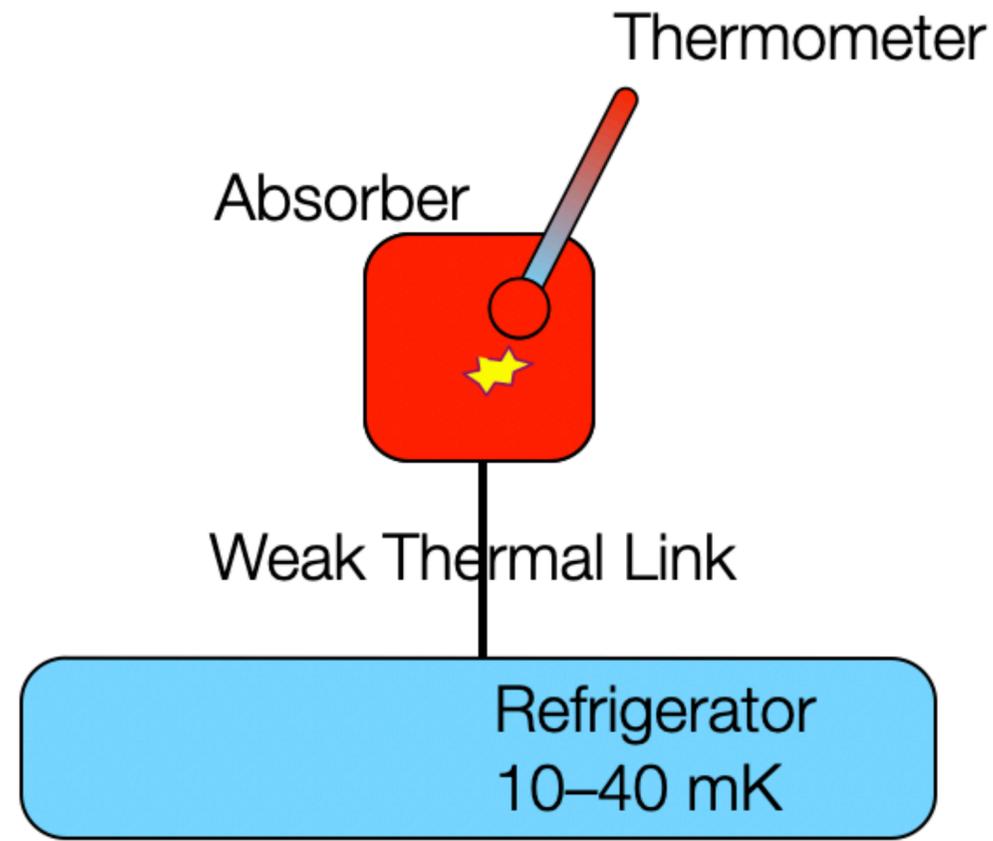
PICO-2L Results



Ken Clark

CRYOGENIC SOLID STATE DETECTORS

CRYOGENIC DETECTORS: PHONON AND HEAT SIGNALS

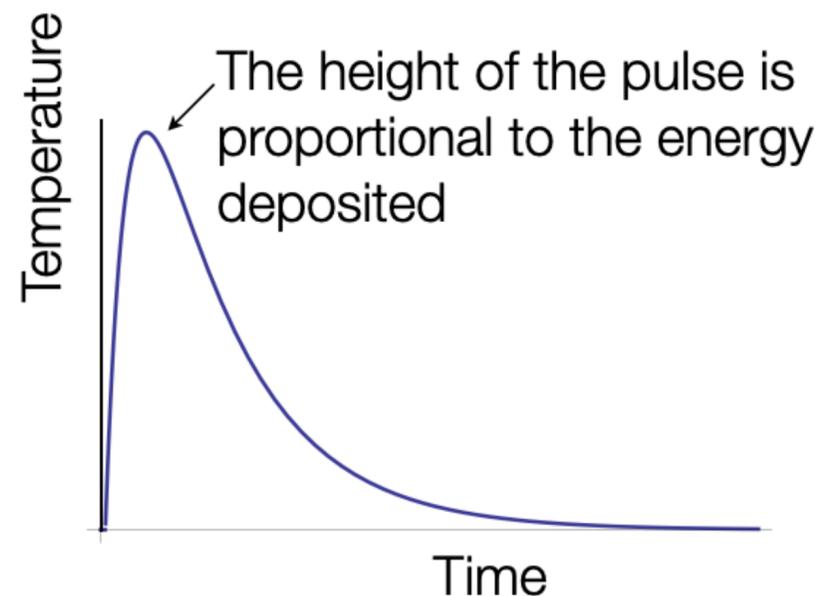


$$\Delta T = \frac{E}{C(T)} e^{-\frac{t}{\tau}}$$

$$\tau = \frac{C(T)}{G(T)}$$

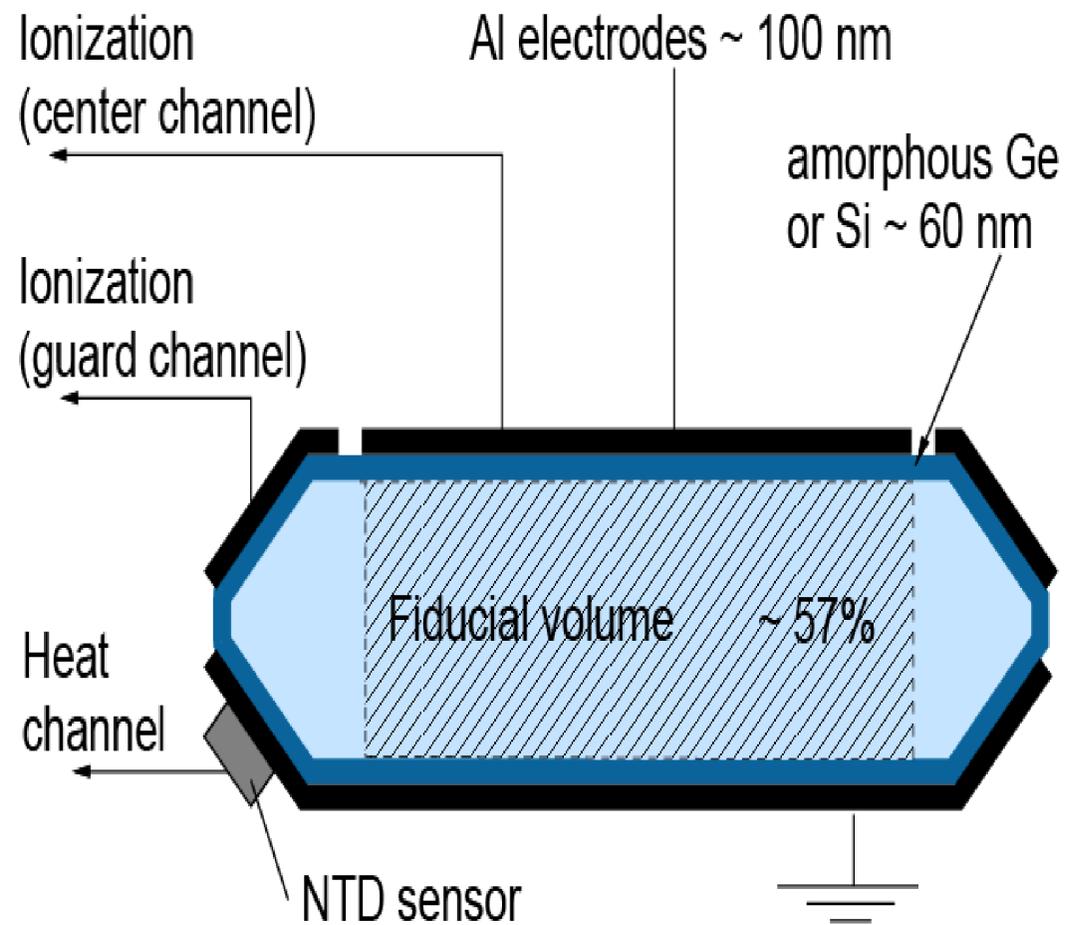
$C(T)$ = heat capacity
of absorber

$G(T)$ = thermal
conductance of the
link between absorber
and reservoir at $T_{\{0\}}$



- Two families of sensors for phonon signal: thermal and athermal
- Thermal sensors - wait for the full thermalization of the phonons within the bulk of the detector and the sensor itself
- Athermal sensors detect fast, non-equilibrium phonons
- Temperature increase is equal to the deposited energy over the heat capacity of the system.
- Two most widely used technologies to measure these signals are neutron doped germanium sensors (NTD) and transition edge sensors (TES)

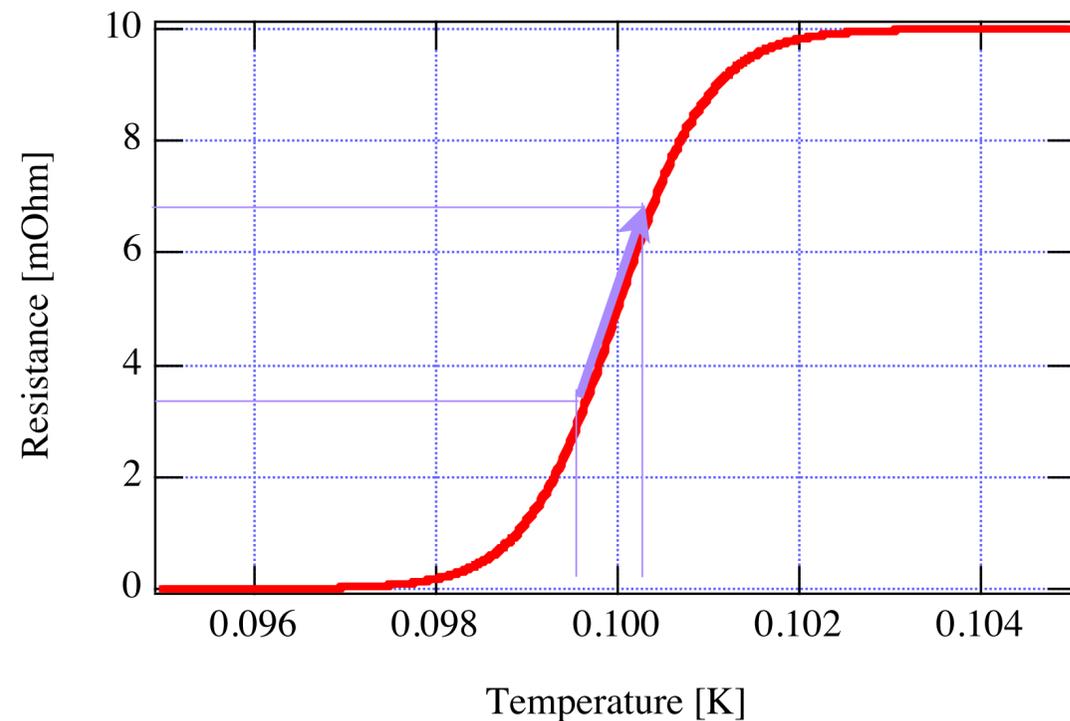
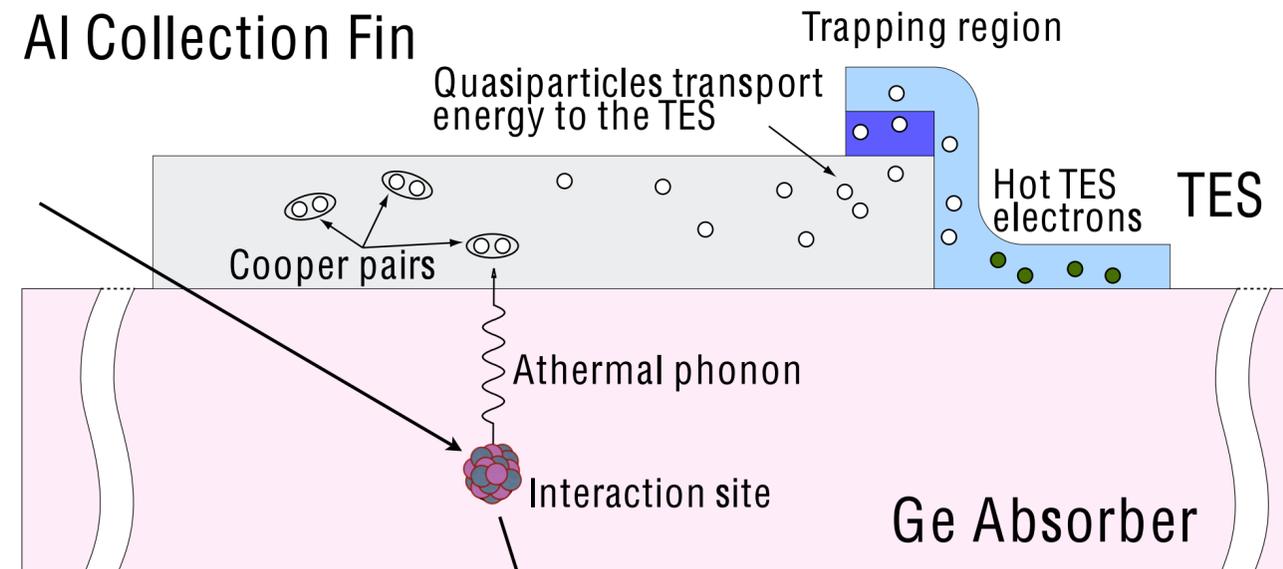
NTDS



Schematic "Ge-NTD"
EDELWEISS detector

- NTDs are small Ge semiconductor crystals that have been exposed to a neutron flux to make a large, controlled density of impurity.
- NTD measures small temperature variations relative to T_0 , which is set to be on the transition from superconducting and resistance regime with dependence of the resistance with temperature T
- Resistance is continuously measured by flowing current through it and measuring the resulting voltage.
- Sensors are glued onto detector.

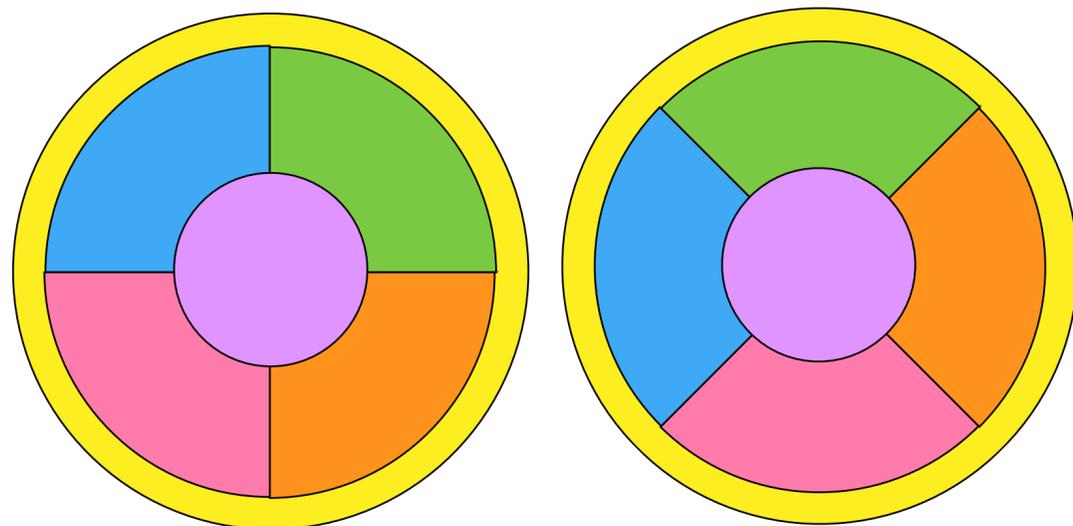
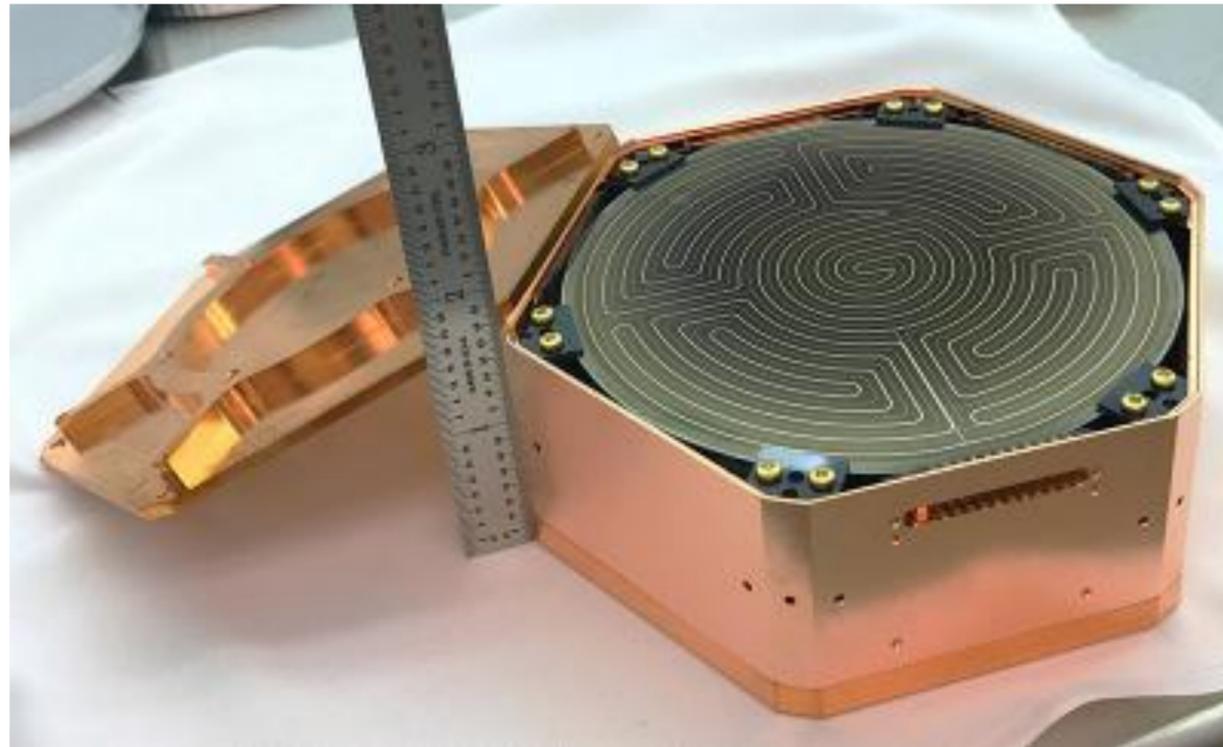
TRANSITION EDGE SENSORS



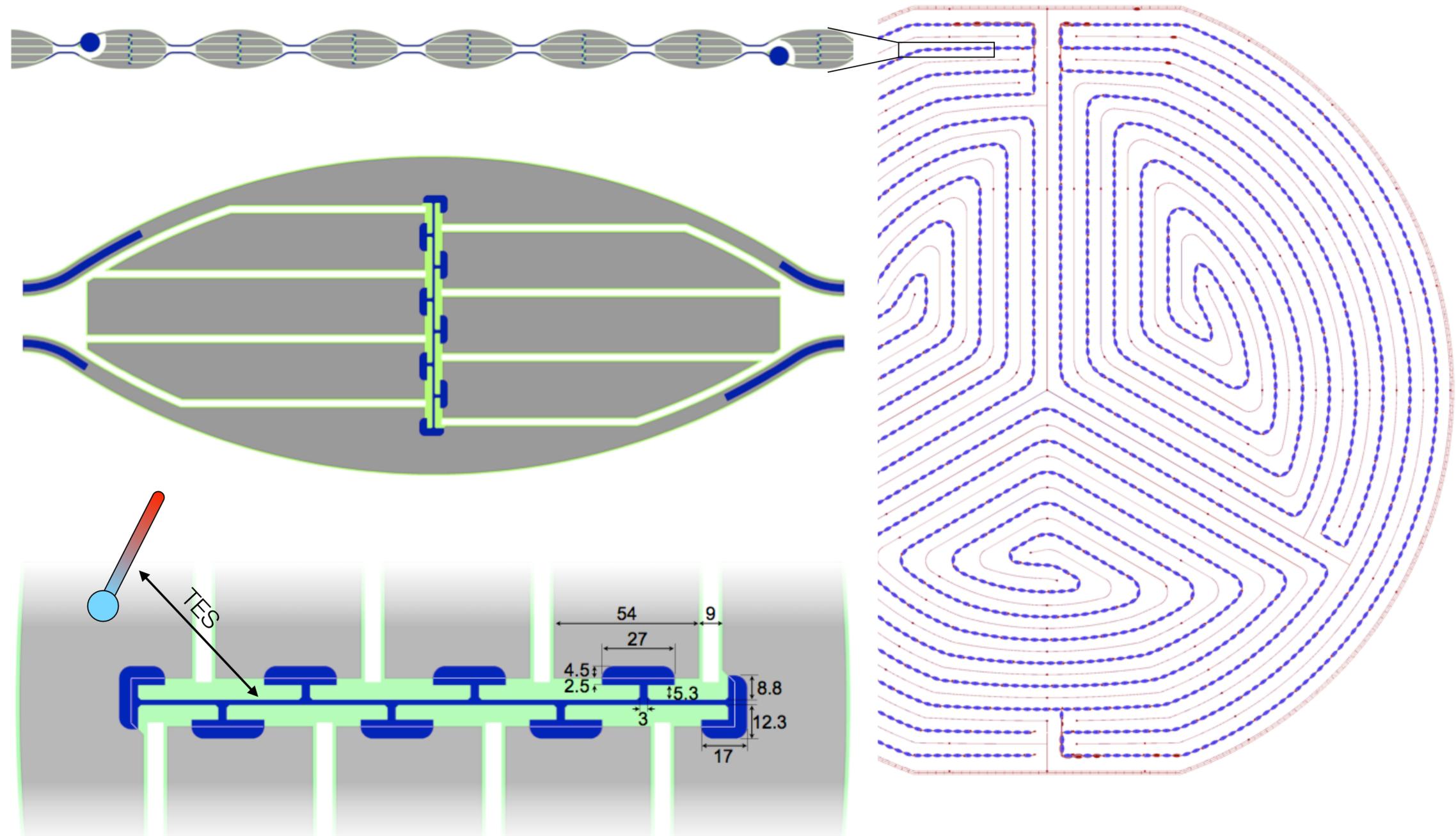
- TES is a thin superconducting film operated near its T_c .
- Refrigerator temperature needs to be close to absolute zero.
- A heater with an electrothermal feedback system maintains temperature at superconducting edge.
- Temperature changes are detected by a change in the feedback current, collected by a SQUID.

SUPERCDMS SNOLAB DETECTORS

- Initial payload 4 towers, each w/6 detectors (1.39 kg Ge crystals, 0.61 kg Is crystals) each 100 mm diameter, 33.3 mm thick:
 - 2 HV (4 Ge + 2 Si)
 - 2 iZIP (6 Ge & 4 Ge + 2 Si)
- iZIP detectors
 - 8 phonon channels + 2 charge sensors each side
- HV detectors
 - 6 phonon channels on each side

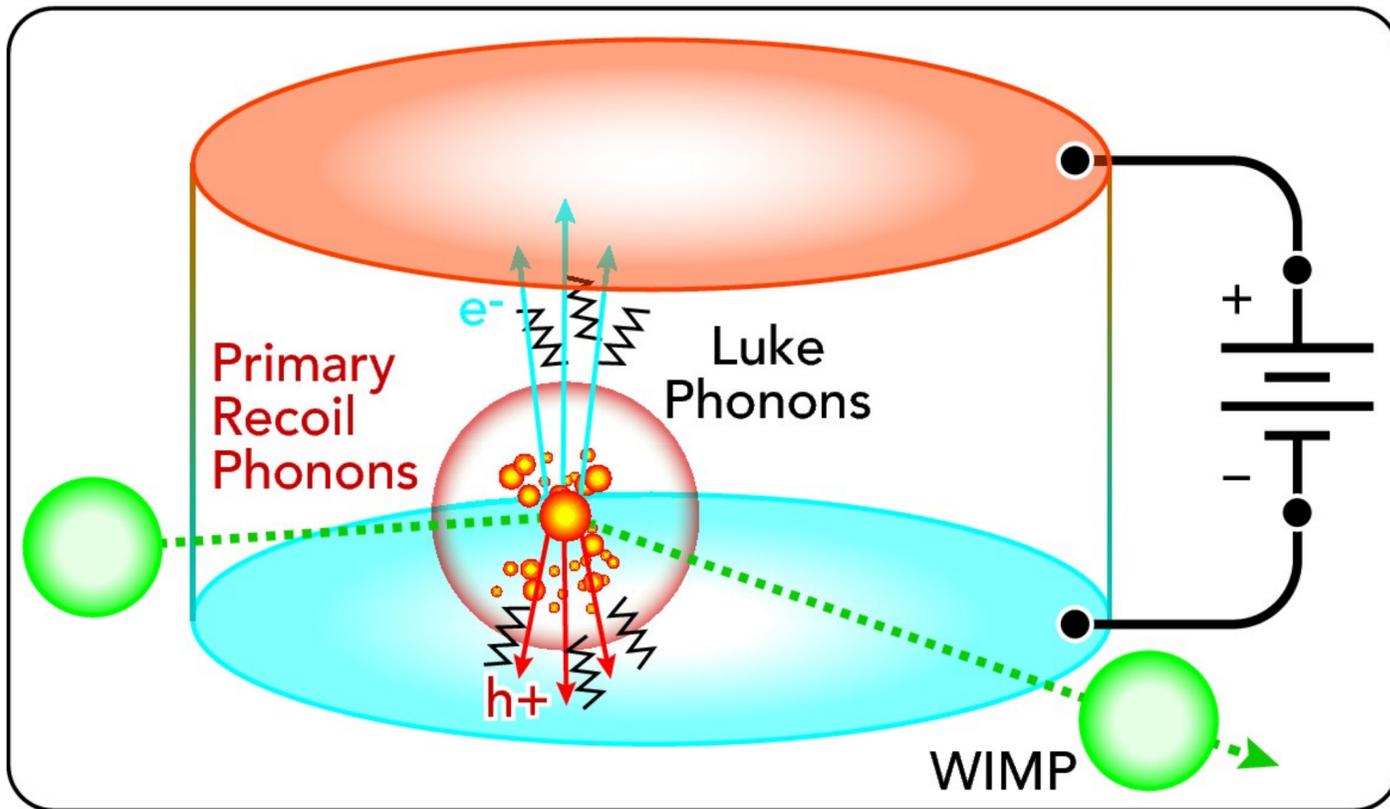


IZIP DETECTOR

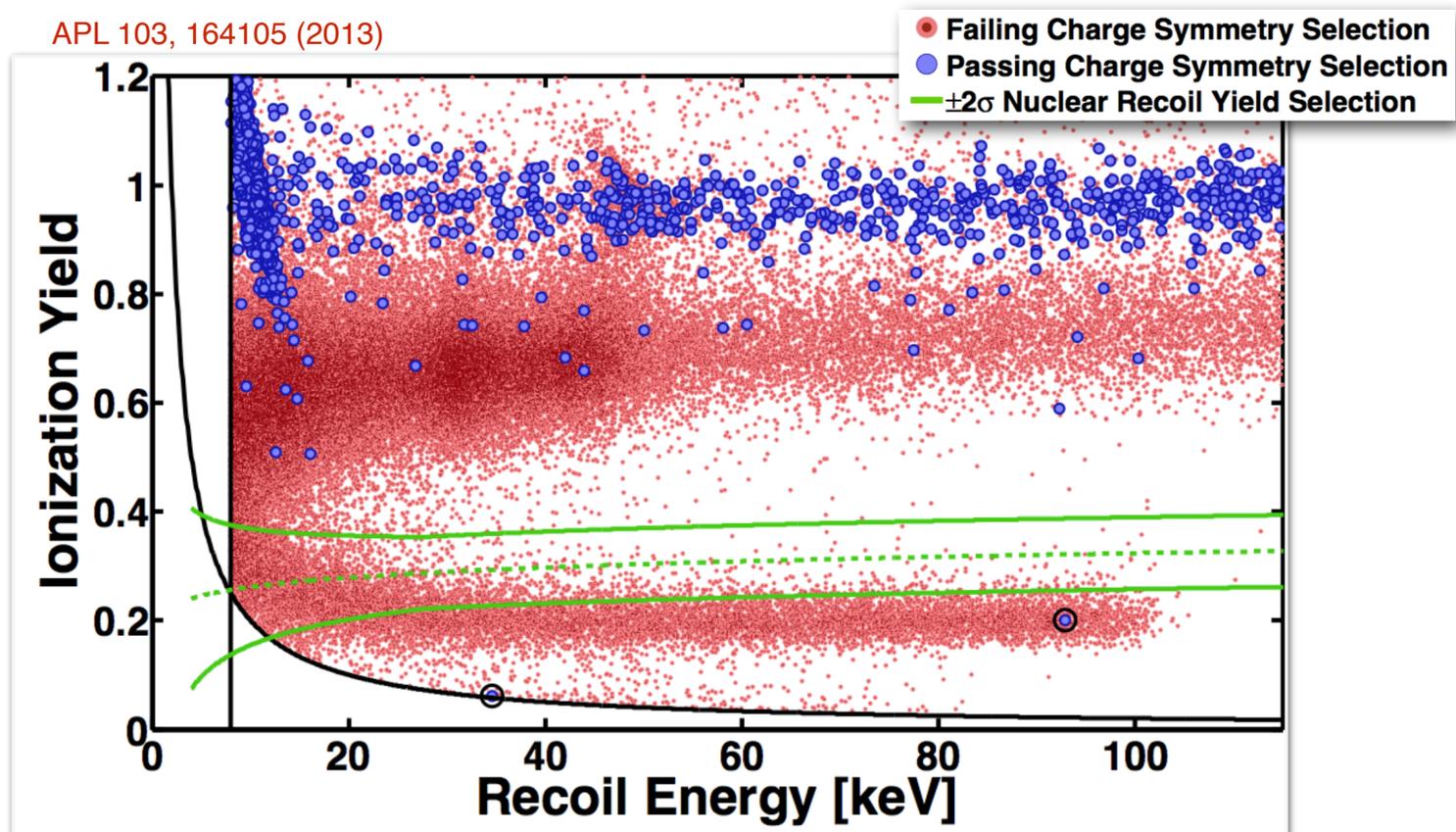


SUPERCDMS - STANDARD IZIP MODE

- ▶ Primary (prompt) phonon and ionization signals allow for discrimination between NR and ER events
- ▶ High resolution phonon and charge readout
- ▶ All surface and ER backgrounds above a few keV can be easily removed with selection criteria.

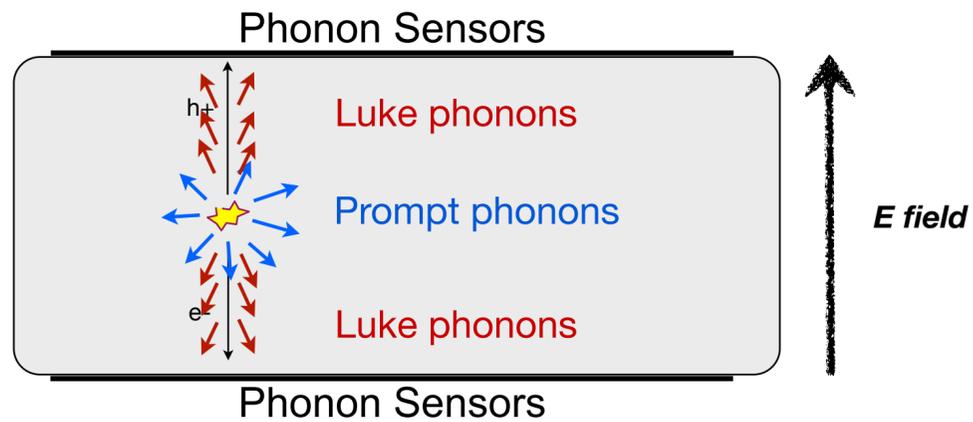
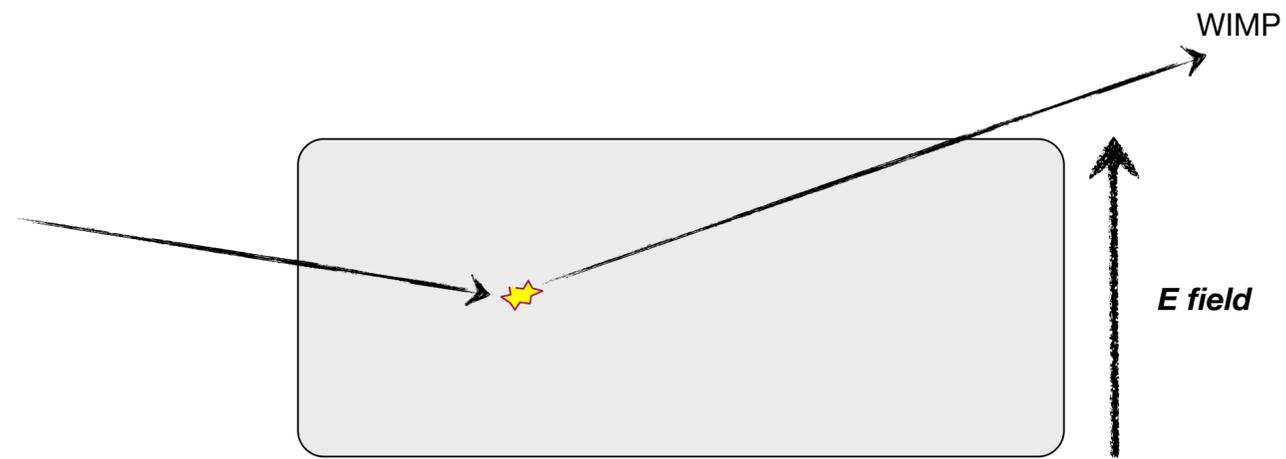


APL 103, 164105 (2013)



SUPERCDMS HIGH-VOLTAGE DETECTORS

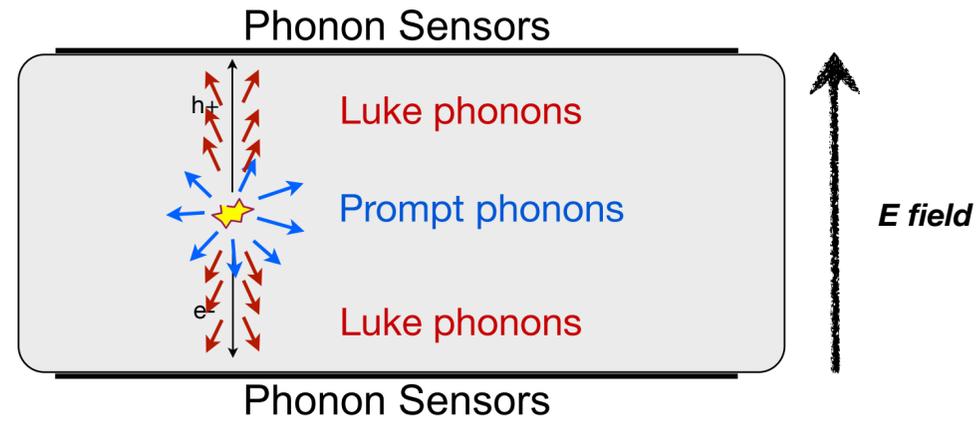
- ▶ Drifting electrons across a potential (V) generates a large number of phonons (NLT phonons)



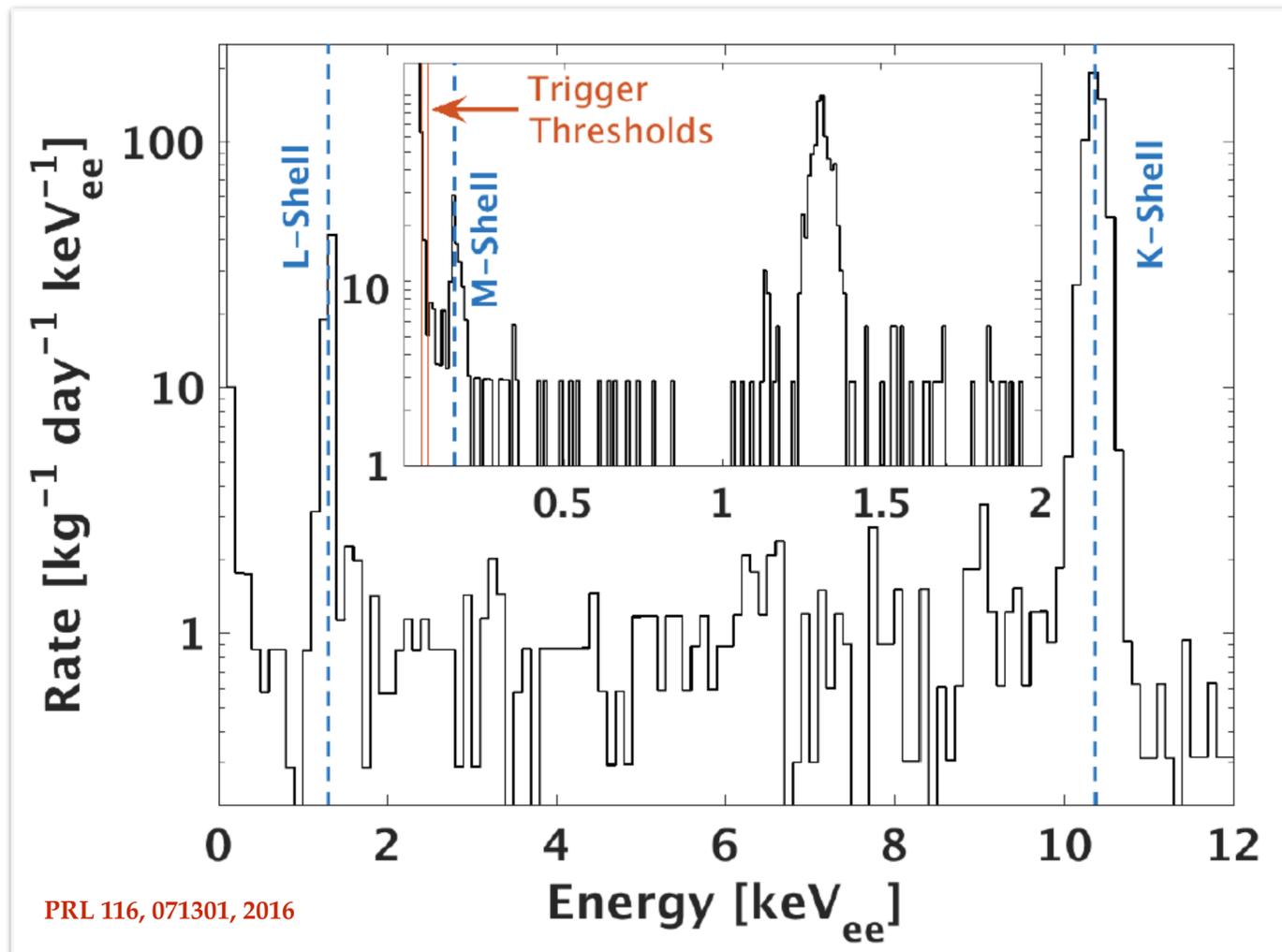
$$E_t = E_r + N_{eh}eV_b$$

E_t (total phonon energy) = E_r (primary recoil energy) + $N_{eh}eV_b$ (Luke phonon energy)

SUPERCDMS HIGH-VOLTAGE DETECTORS



- Drifting electrons across a potential (V) generates a large number of phonons (N phonons)



$$E_t = E_r + N_{eh}eV_b$$

E_t → total phonon energy
 E_r → primary recoil energy
 $N_{eh}eV_b$ → Luke phonon energy

- Ultra high resolution indirect charge measurement
- Thresholds 75 eV_{ee} and 56 eV_{ee}
- No yield or detector face discrimination

ON UNITS

We know that NTL phonon energy is given by

$$E_{NTL} = N_{eh} V_b$$

The number of electron hole pairs generated in an interaction is given by

$$N_{eh} = \frac{E_r}{\epsilon} \quad \epsilon G_e = 3.0 \text{ eV}$$

The total energy (phonon) is given by

$$E_t = E_r + eV_b N_{eh}$$

NR produce eh-pairs less efficiently than ER. Take this into account, define $Y \equiv 1$ for ER.

$$N_{eh} = Y(E_r) \frac{E_r}{\epsilon}$$

The total energy can then be written

$$E_{tot} = E_r \left(1 + Y(E_r) \frac{eV_b}{\epsilon} \right)$$

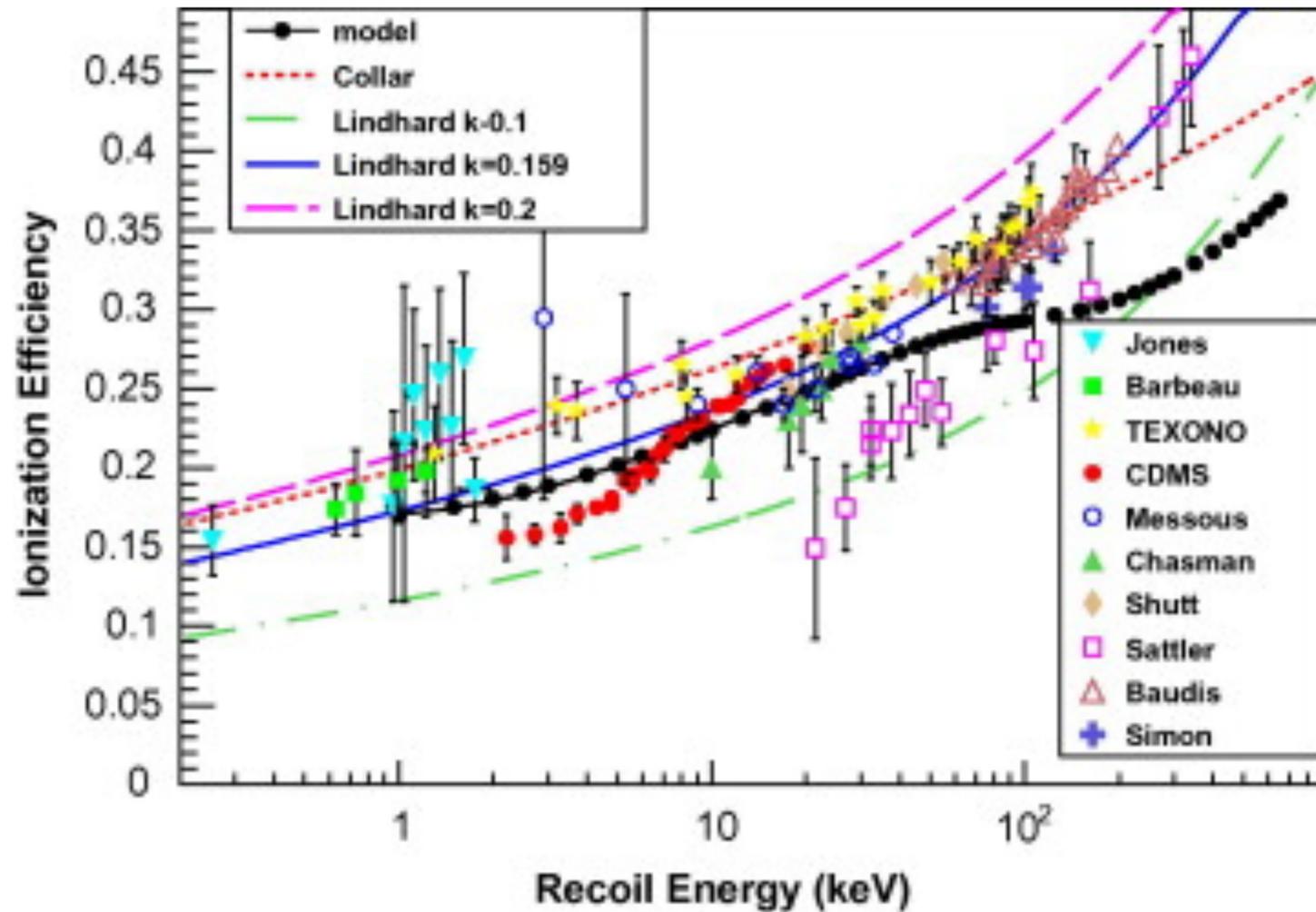
If we calibrate detectors using ER, the resulting energy scale is keV_{ee} to convert to keV_{nr} equate for NR and ER.

$$E_{nr} \left(1 + Y(E_{nr}) \frac{eV_b}{\epsilon} \right) = E_{ee} \left(1 + \underbrace{Y(E_{ee})}_{\text{recall } Y = 1 \text{ for ER}} \frac{eV_b}{\epsilon} \right)$$

recall $Y = 1$ for ER

$$E_{nr} = E_{ee} \left(\frac{1 + eV_b/\epsilon}{1 + Y(E_{nr})eV_b/\epsilon} \right)$$

HOW TO DETERMINE Y?



- Either you need to measure it directly or model it.
- The most utilized model is from Lindhard.

$$Y(E_{nr}) = \frac{k \cdot g(\epsilon)}{1 + k \cdot g(\epsilon)}$$

where

$$g(\epsilon) = 3\epsilon^{0.15+0.7\epsilon^{0.6}+\epsilon}$$

$$\epsilon = 11.5E_{nr}(\text{keV})Z^{-7/3}$$

$Z = \text{atomic number}$

ASIDE: ENERGY

The total energy (phonon) is given by

$$E_{tot} = E_r + \frac{eV_b N_Q}{\epsilon}$$

recoil energy
[keV_{nr}]
total phonon
energy Neganov-Luke
Phonons

➤ Assuming that an event is an ER and that the detector bias voltage is 3V, the recoil energy in [keV_{ee}] can be expressed as

$$\begin{aligned} E_r &= E_{tot} - eV_b N_Q \\ &= E_{tot} - eV_b \frac{E_Q}{\epsilon} \\ &= E_{tot} - E_Q \end{aligned}$$

ϵ_{Ge} = 3.0 eV

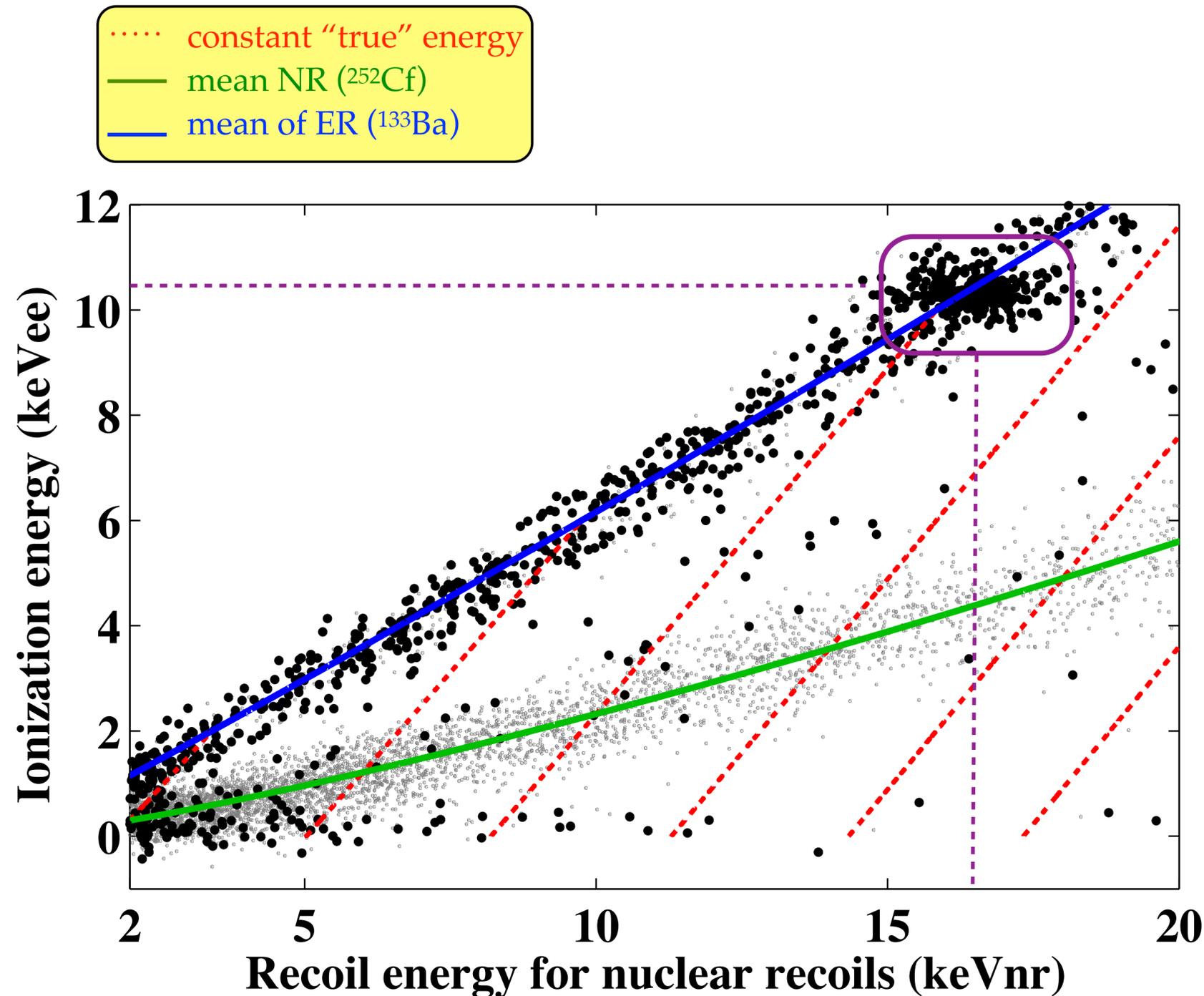
➤ Assuming that an event is a NR, a smaller correction for the Luke phonons is applied. The mean ionization energy for nuclear recoils ($\mu_{Q,nr}(p_t)$) is determined using calibration data from a 252Cf source.

$$E_r(p_t) = p_t - \mu_{Q,NR}(p_t)$$

[keV_{nr}] total phonon
energy - Luke
energy

where $\mu_{Q,NR} = A E_r^B$

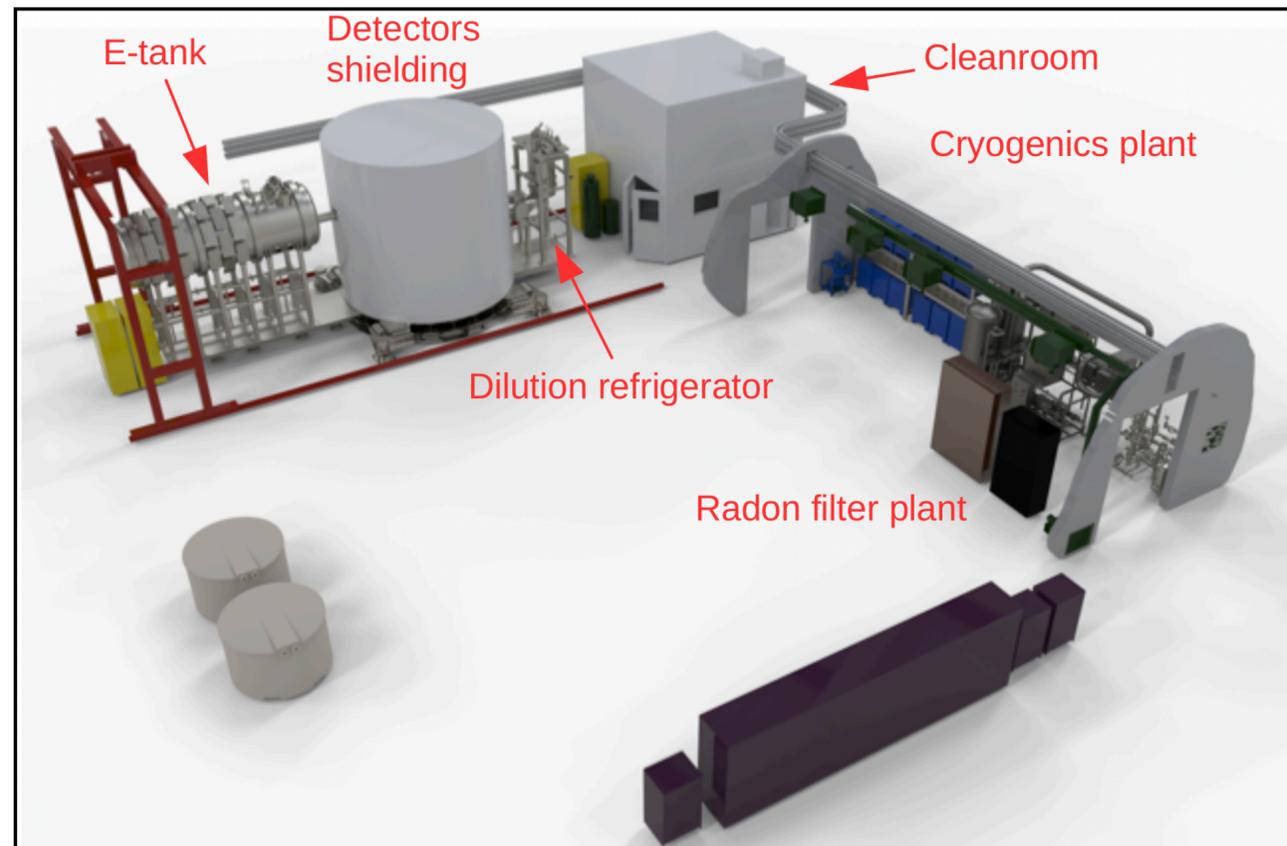
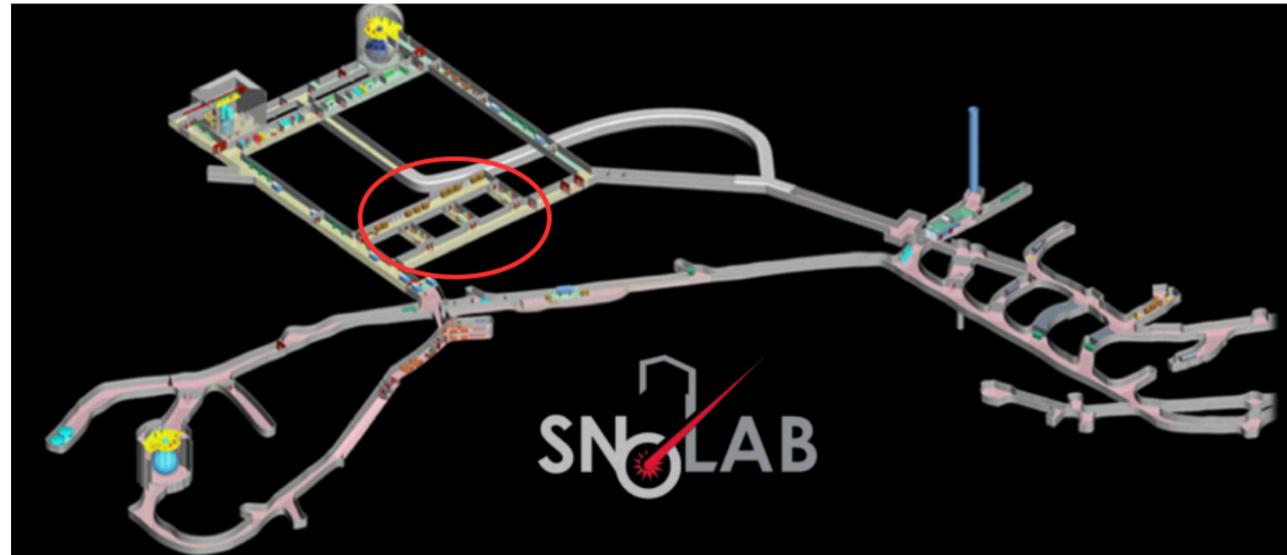
KEV_{EE} VS KEV_{NR}



- ▶ Ionization energy vs recoil energy assuming NR scale consistent with Luke phonon contributions for NR.
- ▶ ER recoils are pushed to higher energies using the NR scale.
- ▶ Example - 10.4 keV_{ee} ER line appears at ~16 keV_{nr}

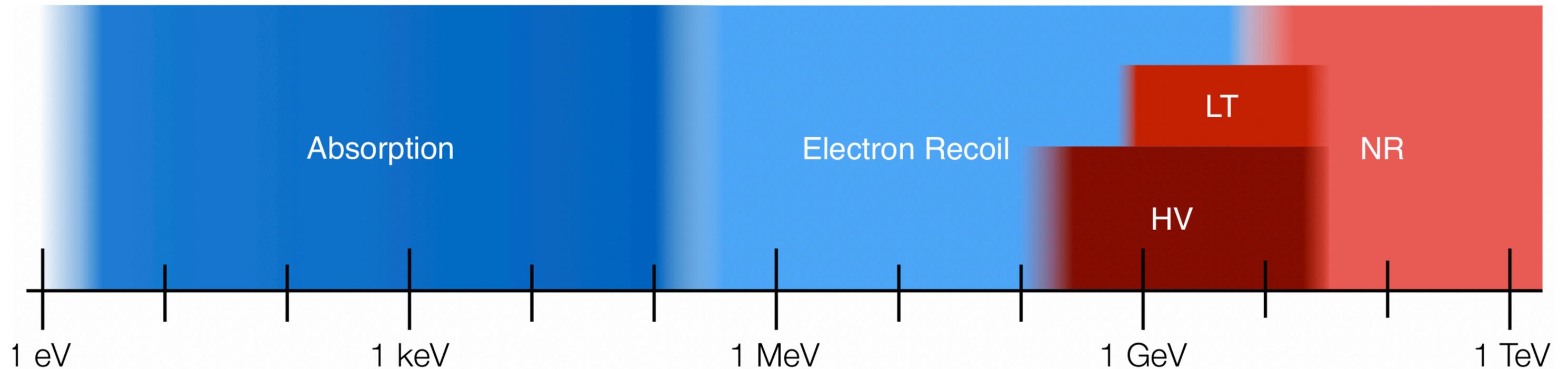
*A good reference is David Moore's thesis, Chapters 3 and 4
<http://thesis.library.caltech.edu/7043/>

SUPERCDMS SNOLAB



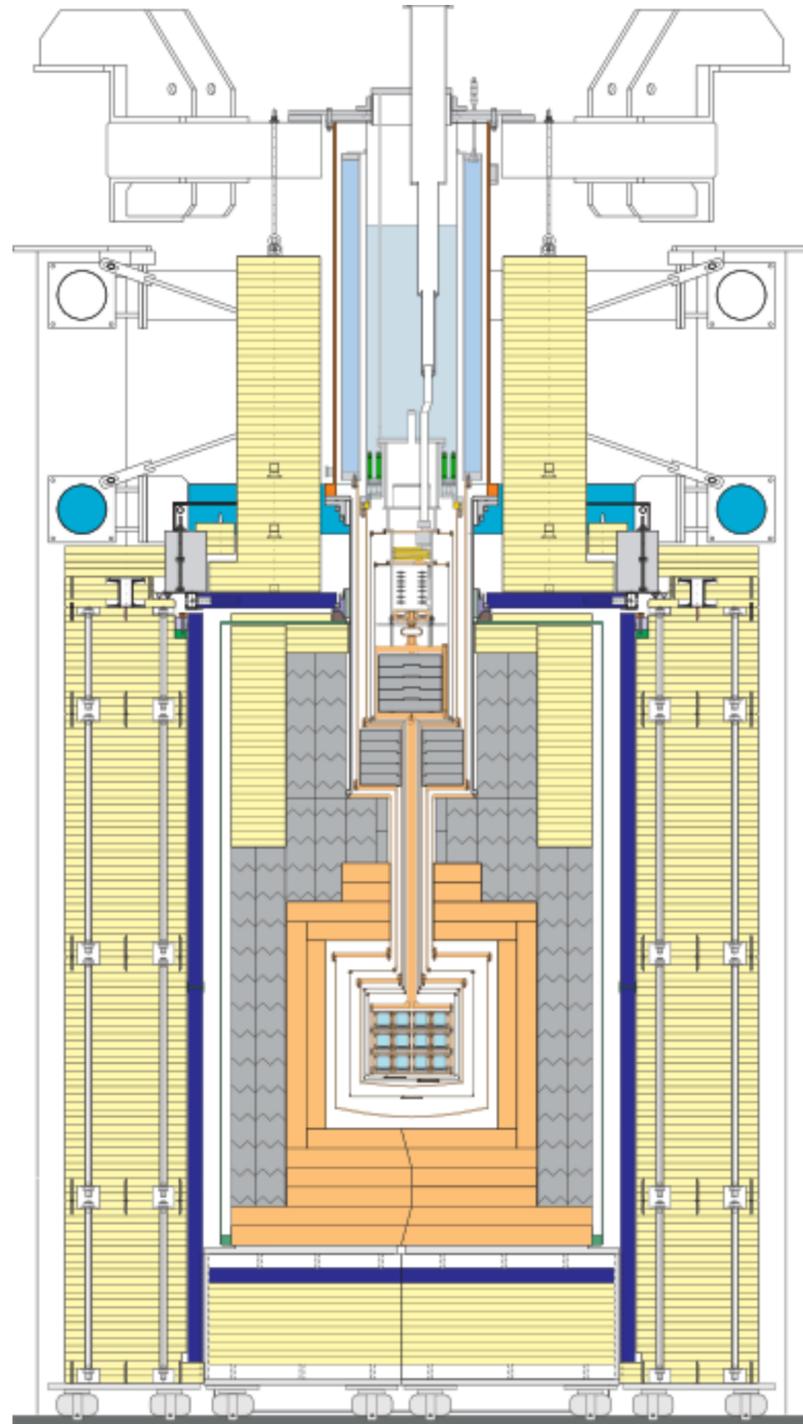
- Generation-2 dark matter experiment under construction at SNOLAB
- Infrastructure:
 - depth ~ 6900 mwe (results in a factor 100 reduction in muon flux from cosmic rays as compared to Soudan)
 - class 2000 or better cleanroom
 - Cryostat will be able to accommodate up to 7 towers
 - (0.1) dru gamma background
 - 15 mK base temperature
 - vibration isolation
- Initial payload: ~ 30 kg total, 4 towers with 6 detectors per tower (12 iZIP, 12 HV)

SUPERCDCMS DARK MATTER SENSITIVITY

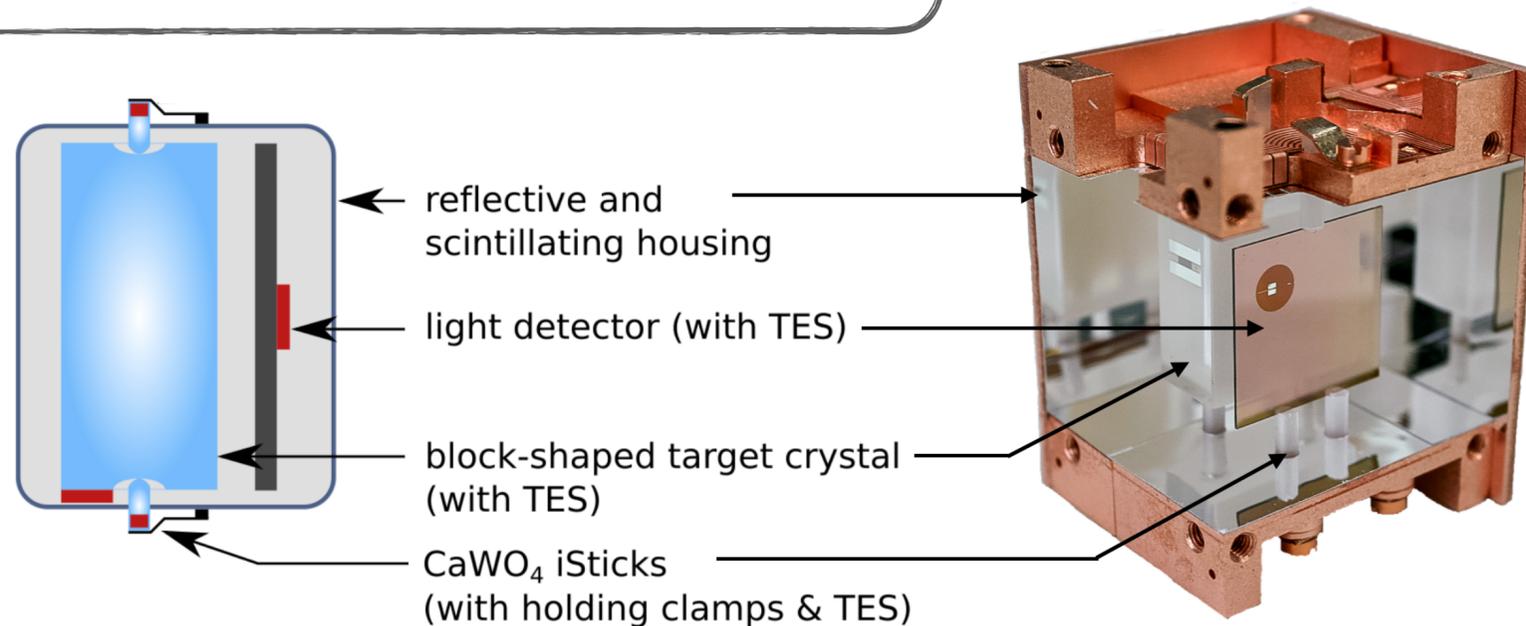
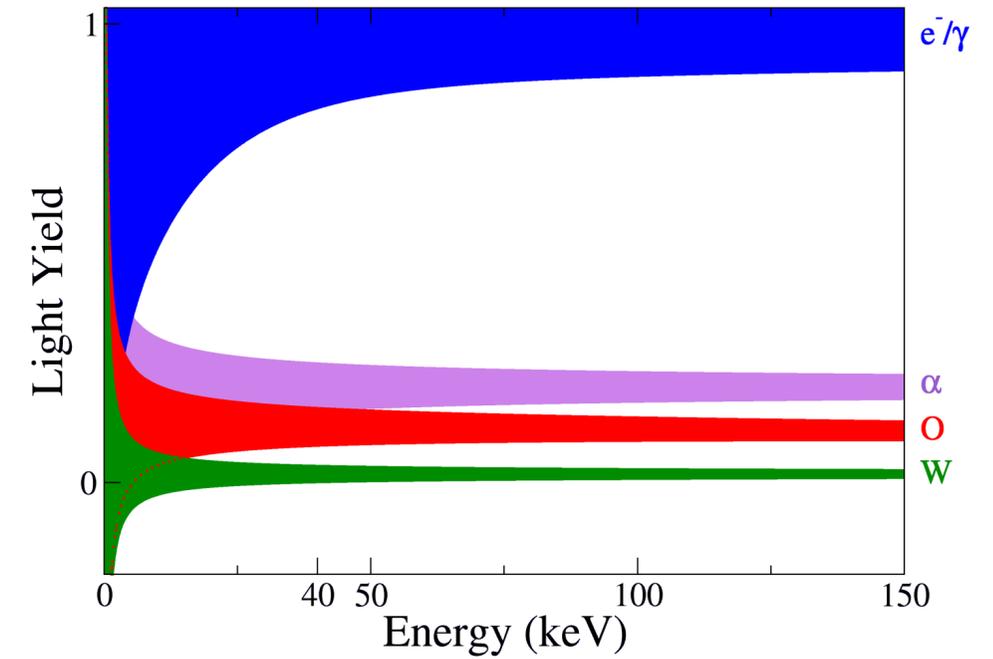


Traditional NR:	iZIP, Background free	>5 GeV
Low Threshold NR:	iZIP, limited discrimination	>1 GeV
HV Mode:	HV, no discrimination	~0.3 - 10 GeV
Electron Recoil:	HV, no discrimination	~0.5 MeV - 10 GeV
Absorption (Dark Photons, ALPs)	HV, no discrimination	~1 eV - 500 keV (peak search)

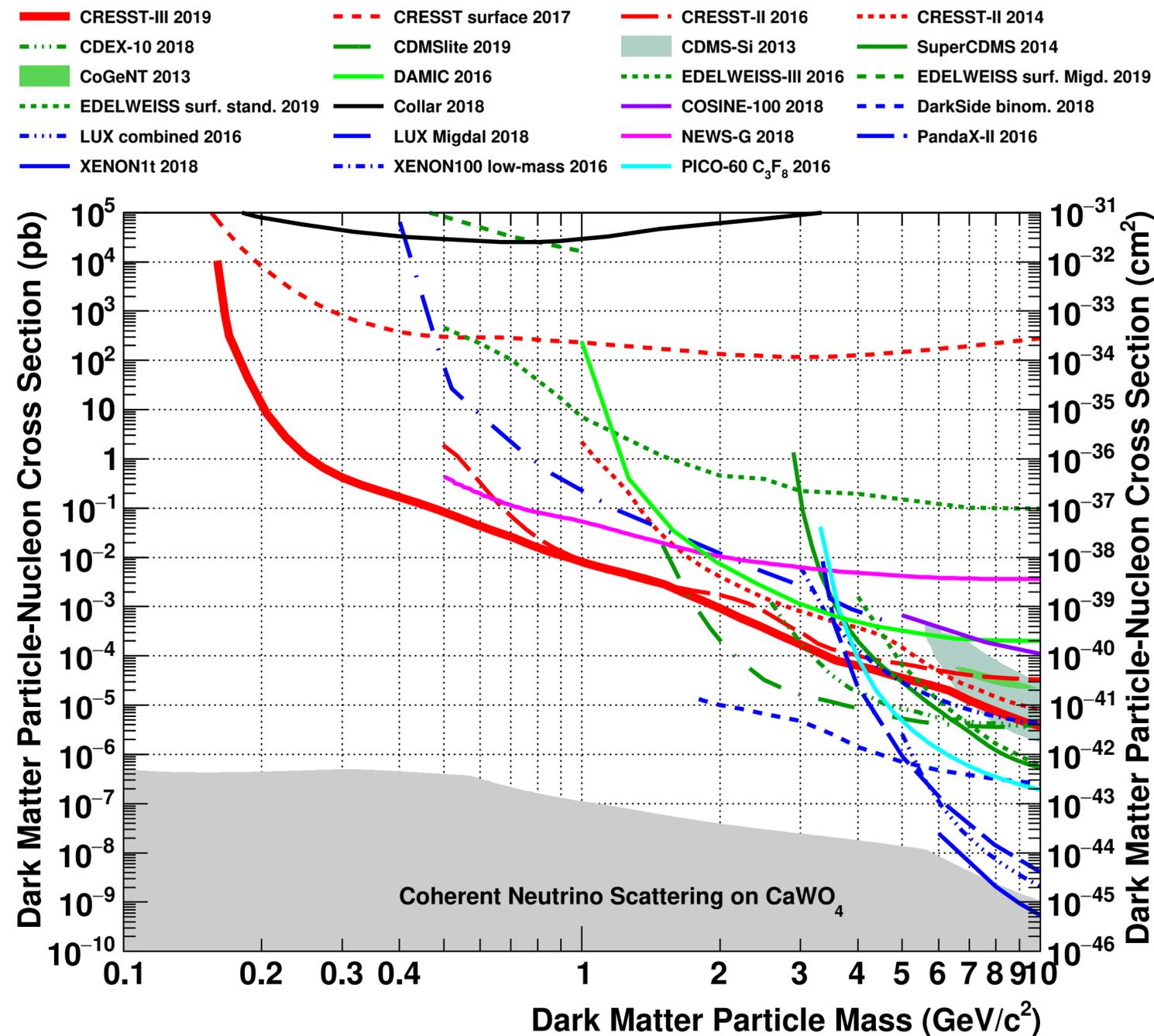
CRESST EXPERIMENT OPERATION PRINCIPLES



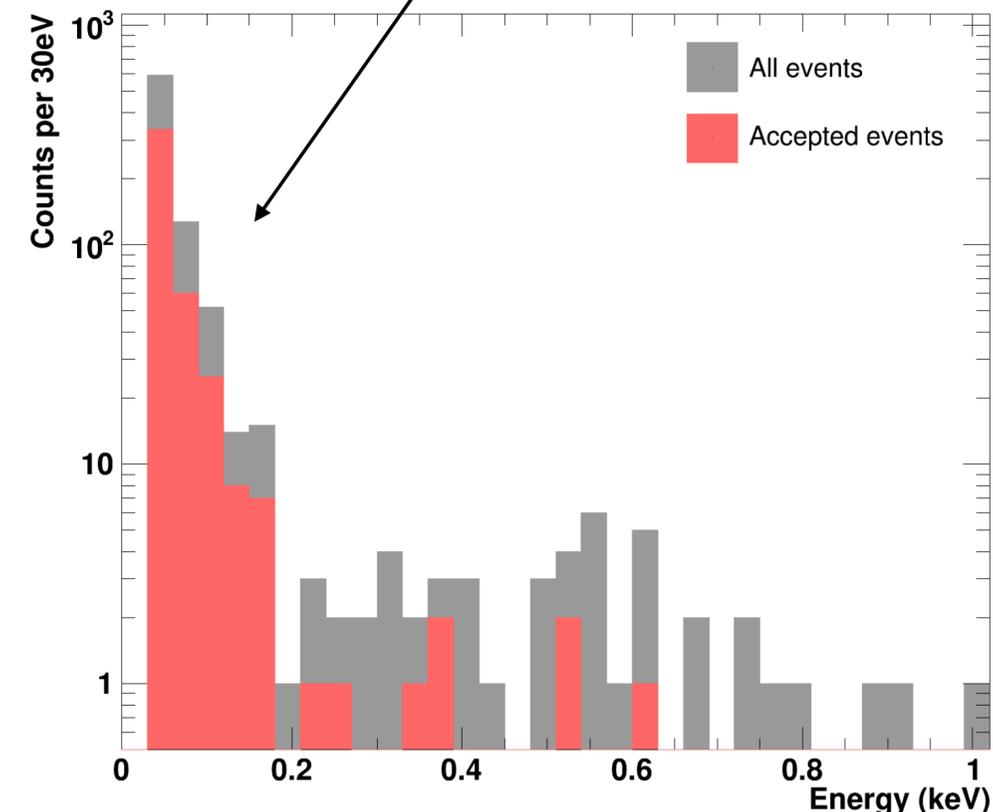
- Search of light DM direct interactions with CaWO_4 cryogenic detectors
- Operating temperature ~ 15 mK
- Second cryogenic detector to collect emitted scintillation light: particle identification
- Single detector mass ~ 24 g
- Energy Threshold: 30 eV

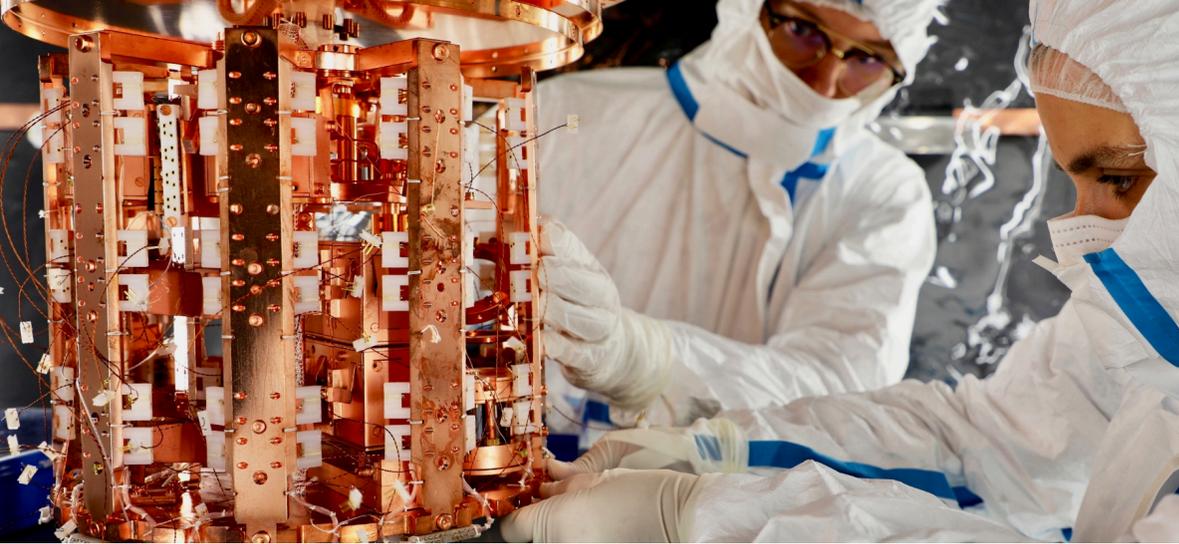


LIMITATIONS: CRESST-III RECENT RESULTS



- More than one order of magnitude improvement at 0.5 GeV/c²
- Extended reach from 0.5 GeV/c² to 0.16 GeV/c²
- Sensitivity limited by unknown background below 200 eV





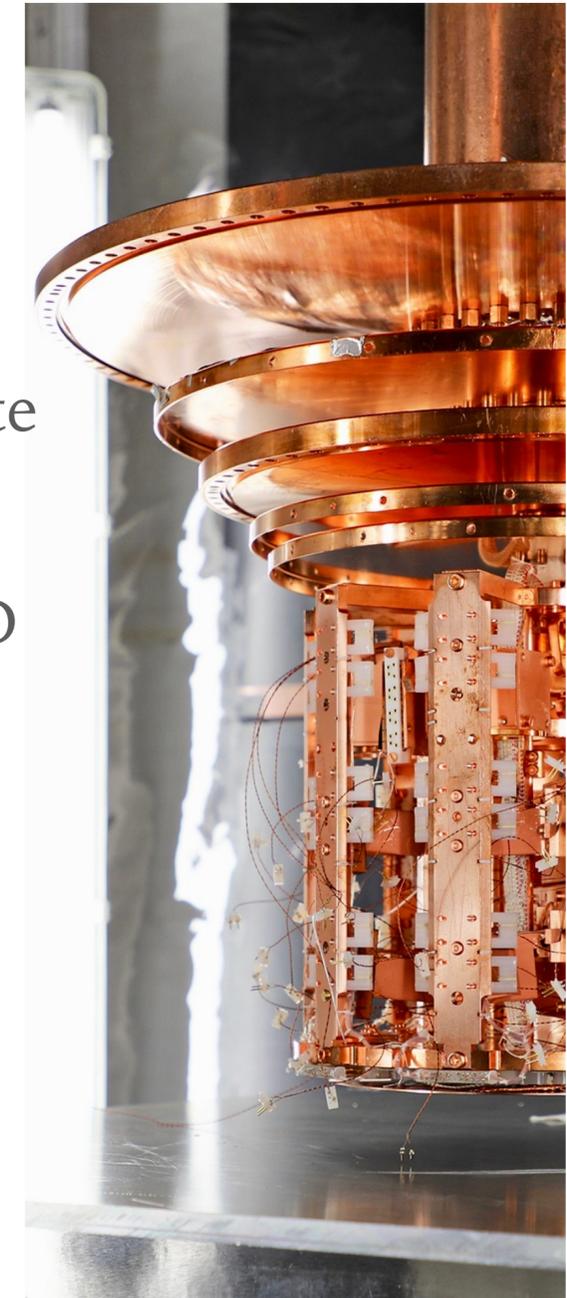
CRESST-III FUTURE PLANS

Run3 2020 - 2021

- 2nd round with additional modifications
- Successful cool-down in 03/2020, but stopped due to Corona virus pandemic
- Cool-down started July 20, 2020
- Detector commissioning Aug - Oct 2020
- November 5, 2020 science data!
 - Preliminary analysis of calibration data is very promising!

CRESST Upgrade 2020 - 2023

- Upgrade to 288 readout channels to accommodate 100 modules for O(2 kg) target mass
- Final planning, prototyping and testing of SQUID read-out electronics, biasing system and DAQ
- Sensor development to further push detector threshold (10 eV)
- Complementary detector materials (LiAlO_2 , Al_2O_3) which also yield sensitivity for spin-dependent interactions
- 2023 science data taking starts!



Single e^-/h^+ Pair Sensitivity

SENSEI

Sub-Electron-Noise SkipperCCD
Experimental Instrument

Pi. Javier Tiffenberg
javiert@fnal.gov

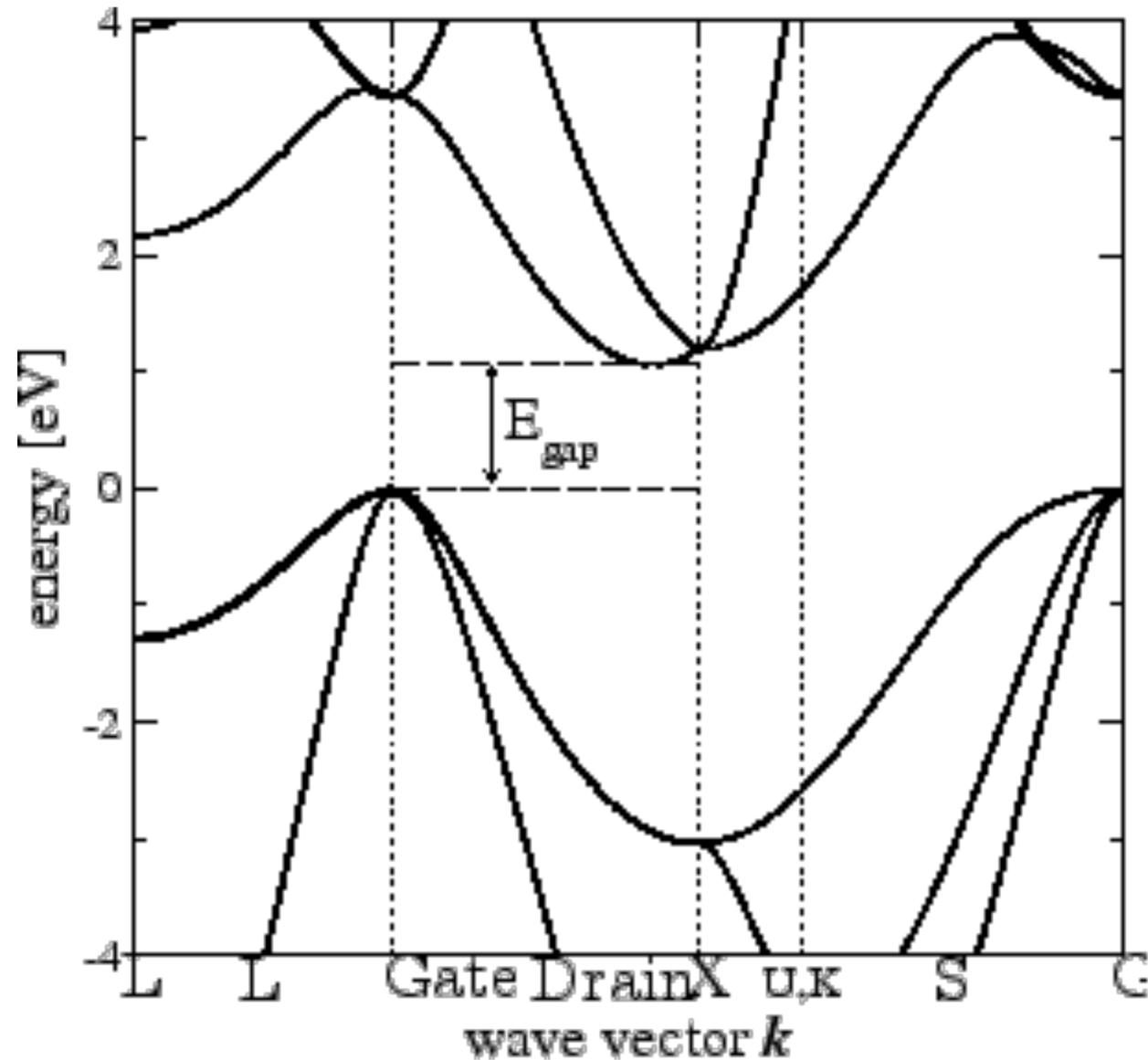
Fermilab

4



ENERGY SCALE IN SEMICONDUCTORS

Band Diagram for Si



- ▶ e- excitation momentum and energy scales in semiconductors can be exploited to search for light mass dark matter
- ▶ Si $E_{\text{gap}} \sim 1.2$ eV
 - ▶ Indirect band gap requires phonon for transition to happen.
 - ▶ Temperature dependent
- ▶ $\epsilon_{\text{Si}} \sim 3.6$ eV
 - ▶ Average energy to produce e/h pair
 - ▶ Temperature dependent
- ▶ Sensitive to energy deposits of $\mathcal{O}(\text{eV})$ (electron scattering) to $\mathcal{O}(10 \text{ eV})$ (nuclear scattering)

REALM OF SOLID STATE PHYSICS

Solid state physics

$E < 30 \text{ eV}$

Multi-body system

Allowed energies/momenta given
by dispersion relation

Particles may have effective
masses

Particle physics

$E > \text{keV}$

Free particles

$E = p^2/2m$

Particle masses well defined

CHALLENGES

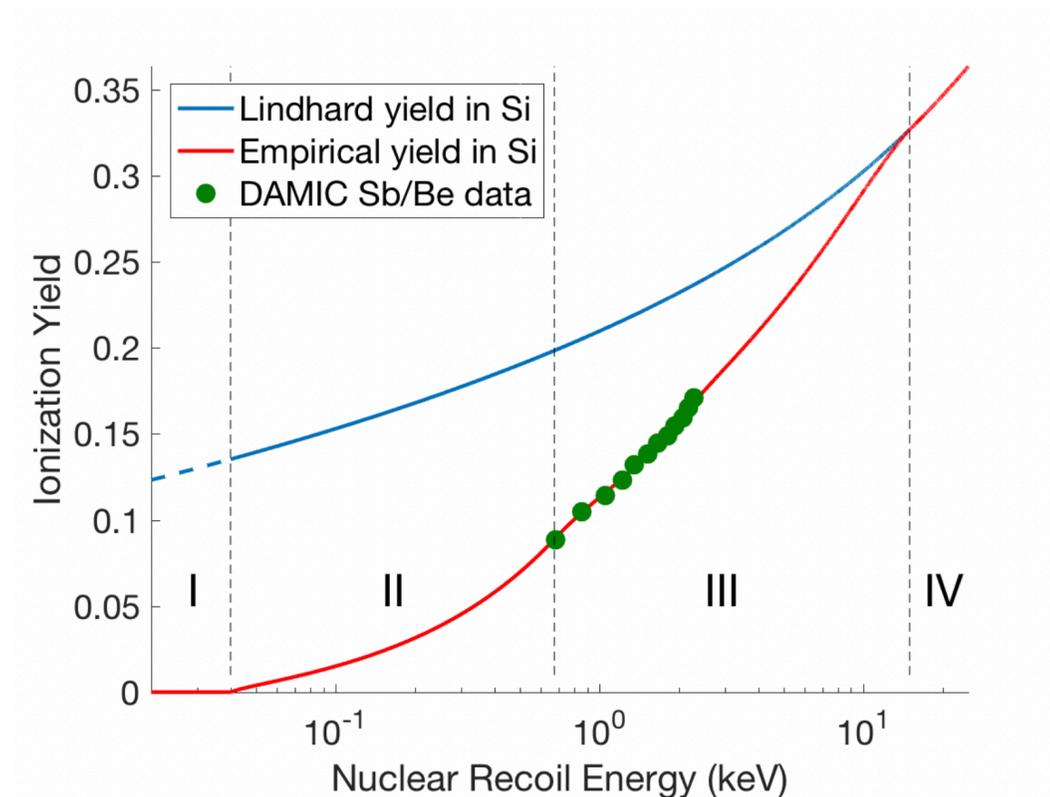
- **Detector Response**

- Details of the band structure become increasingly important

CHALLENGES

► Detector Response

- Details of the band structure become increasingly important
- PDF to get the numbers of e/h pairs given an energy deposition required, $P(n_{eh} | E_{dep})$
 - Fano statistics (dispersion probabilities)
 - For NR: quenching (ionization yield < 1)



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 - Fano statistics (dispersion probabilities)
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- Crystal impurities can lead to partial energy deposits \rightarrow gives events between quantization peaks

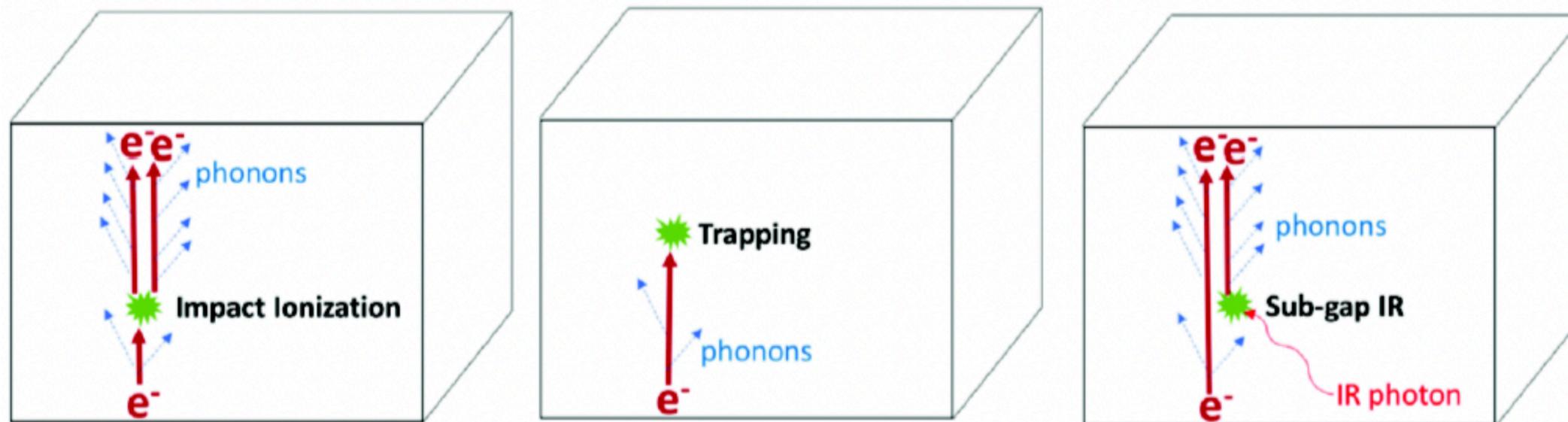


Figure: R.K. Romani

CHALLENGES

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- Spectral information about radioactive decays at eV scale required.

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CHALLENGES

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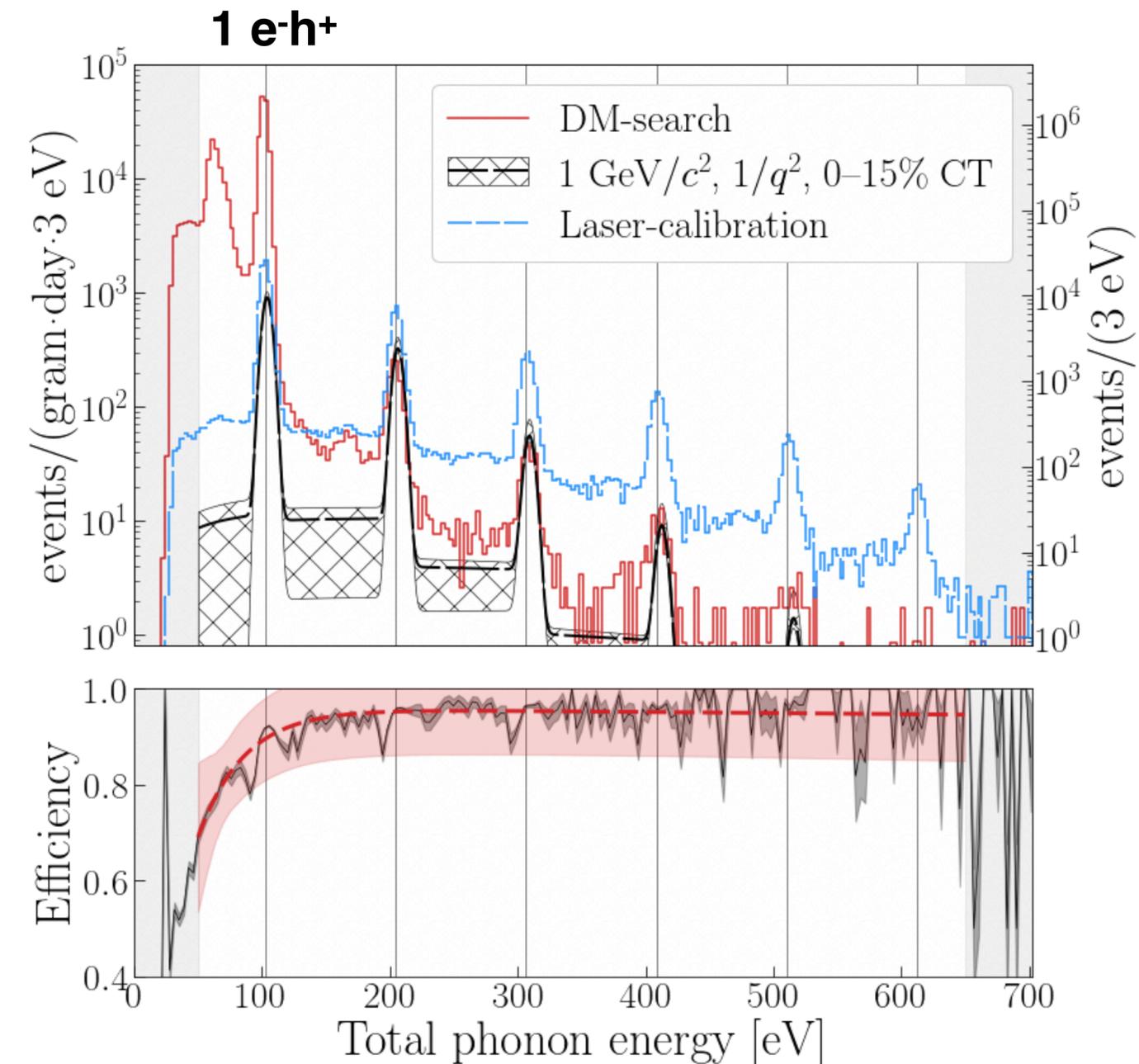
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➤ **Backgrounds**

- Spectral information about radioactive decays at eV scale required.
 - Relevance is exposure dependent
- IR and optical photons become significant backgrounds at lowest energies.
- Dark/leakage current can be significant, dominant background at lowest energies.

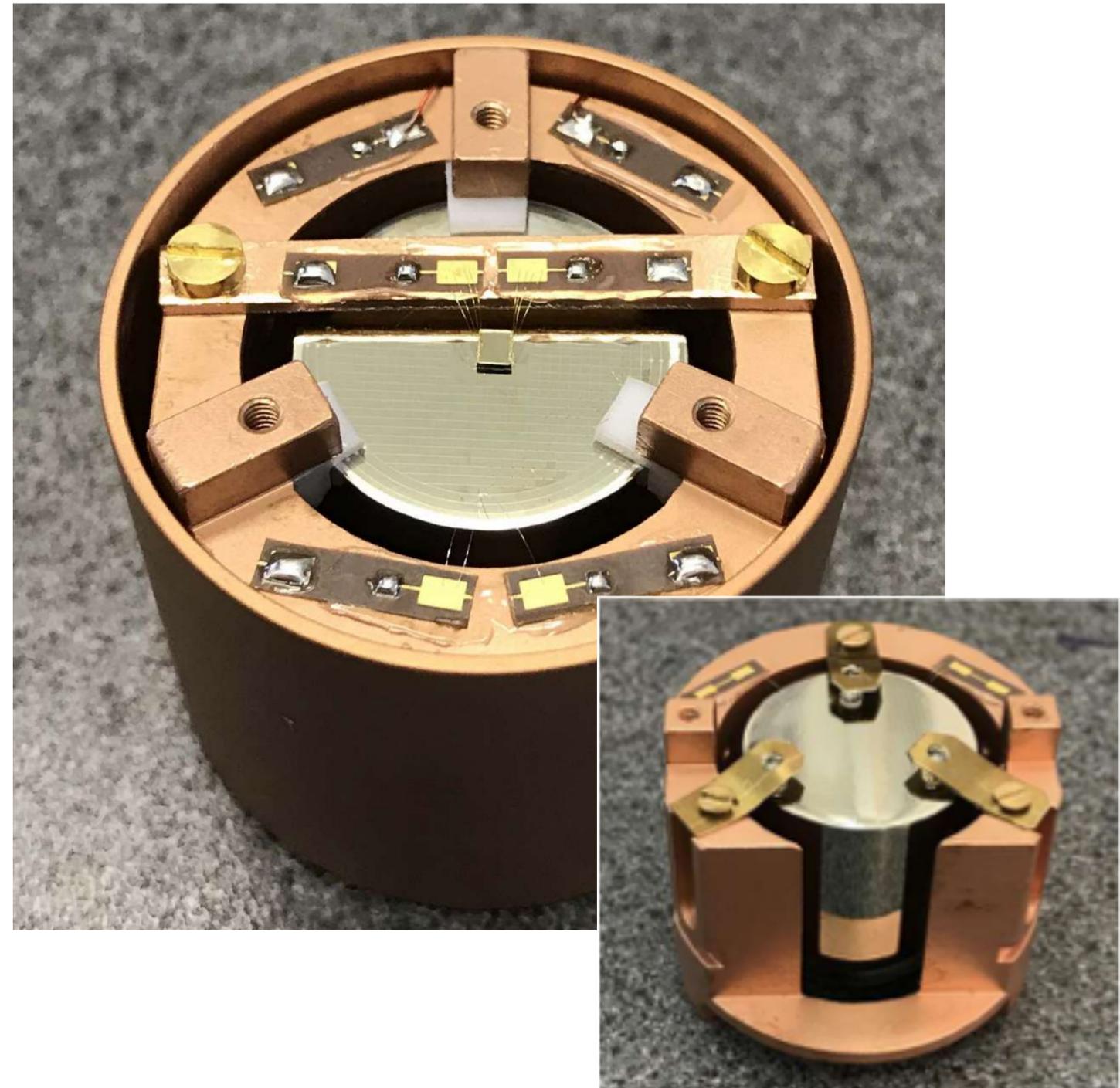
HVEV DETECTORS

- Single-hole e/h-pair resolution devices will have sensitivity to a variety of sub-GeV DM models with g^*d exposures
- 0.93 g Si crystal ($1 \times 1 \times 0.4 \text{ cm}^3$) operated at 50-52 mK at a surface test facility.
- Exposure: 3.0 gram-days (collected over 3 days)
 - operation voltage: 100 V
 - energy resolution: $\sigma_{\text{ph}} = 3 \text{ eV}$
 - charge resolution: $\sigma_{\text{eh}} = 0.03 \text{ e-h}^+$
- Calibrations with in-run monochromatic 635 nm laser fiber-coupled to room temperature.
- Data selection criteria were applied to remove leakage and surface events.



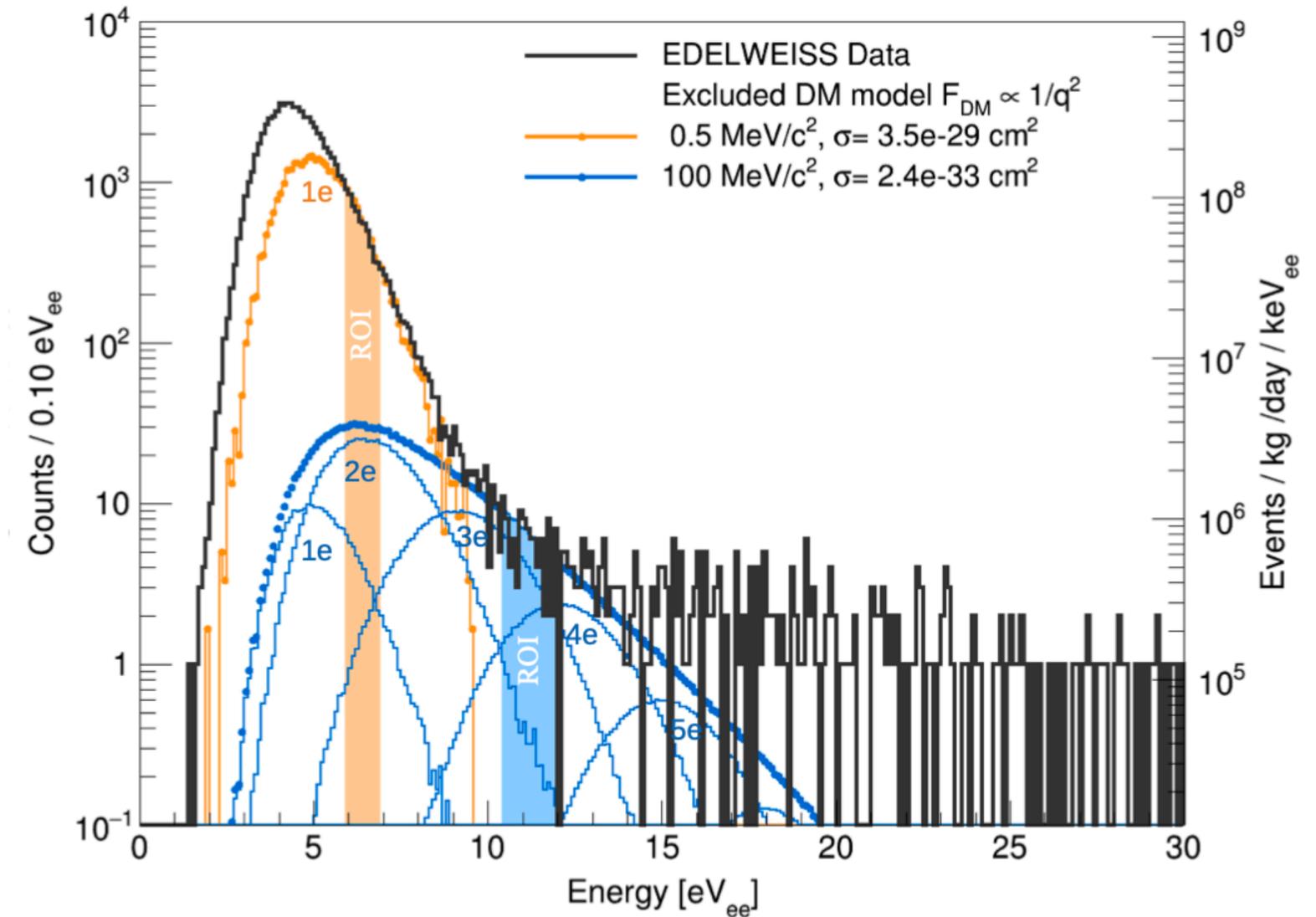
EDELWEISS RED 30 DETECTOR: HV OPERATION

- 33.4 g (20 x 20 mm) Ge bolometer with NTD sensor and electrodes operated in LSM ($5 \mu\text{m}^2/\text{d}$)
- Exposure: 2.44 days
 - operation voltage: 78 V
 - energy resolution: $\sigma_{\text{ph}} = 44 \text{ eV}$ (1.6 eVee)
 - charge resolution: $\sigma_{\text{eh}} = 0.53 \text{ e-h}^+$
- Calibrations using ^{71}Ge KLM (0.16, 1,30 and 10.37 keV) activation lines from AmBe neutron source.
- Data selection criteria were applied to remove events occurring in the NTD (instead of the crystal).



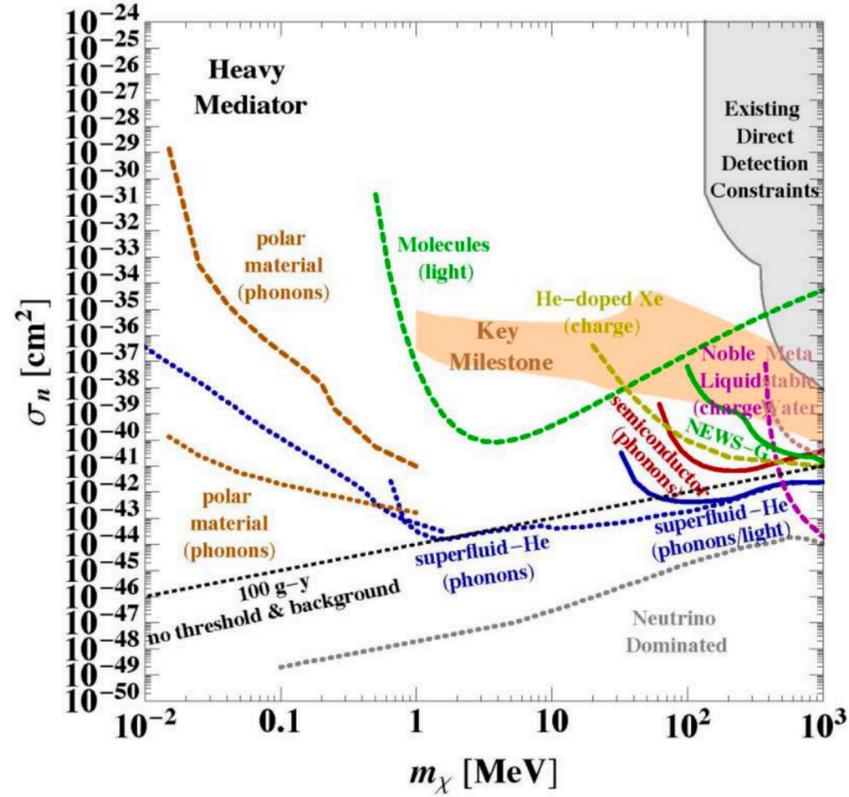
EDELWEISS RED 30 DETECTOR: HV OPERATION

- Heat only events (those not affected by NTL amplification) are the main source of backgrounds.
 - 10^6 DRU @ 10 eV_{ee}
 - 1.5×10^5 DRU @ 25 eV_{ee}
- Dominant limitation for >3 e- signals
- May hypothesis have been studied as to the origin. No single contributor has been found.
 - These events are probably multiple sources.

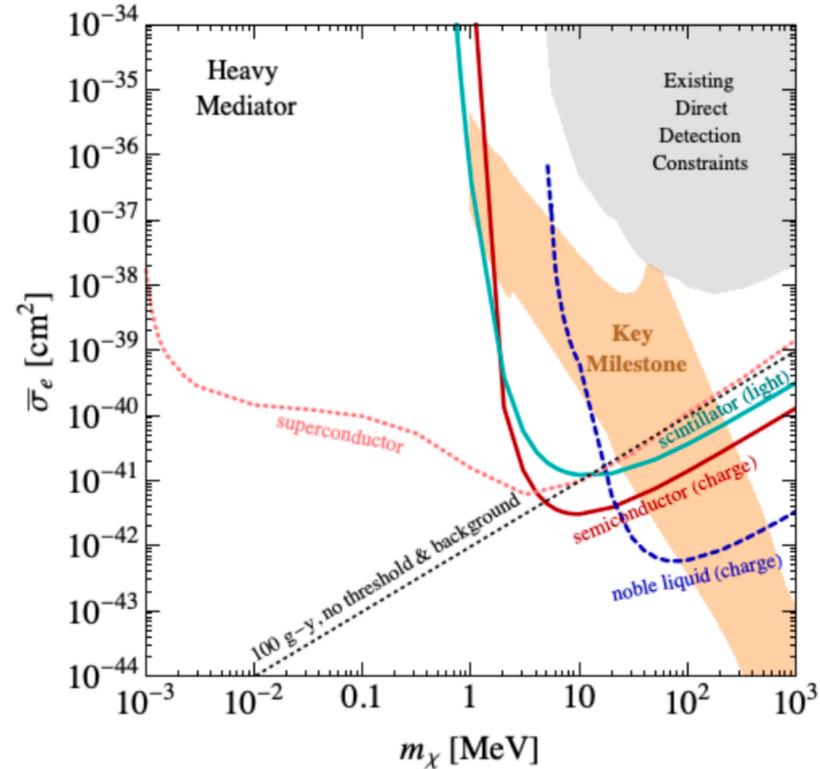


(dru = event/day/keV/kg)

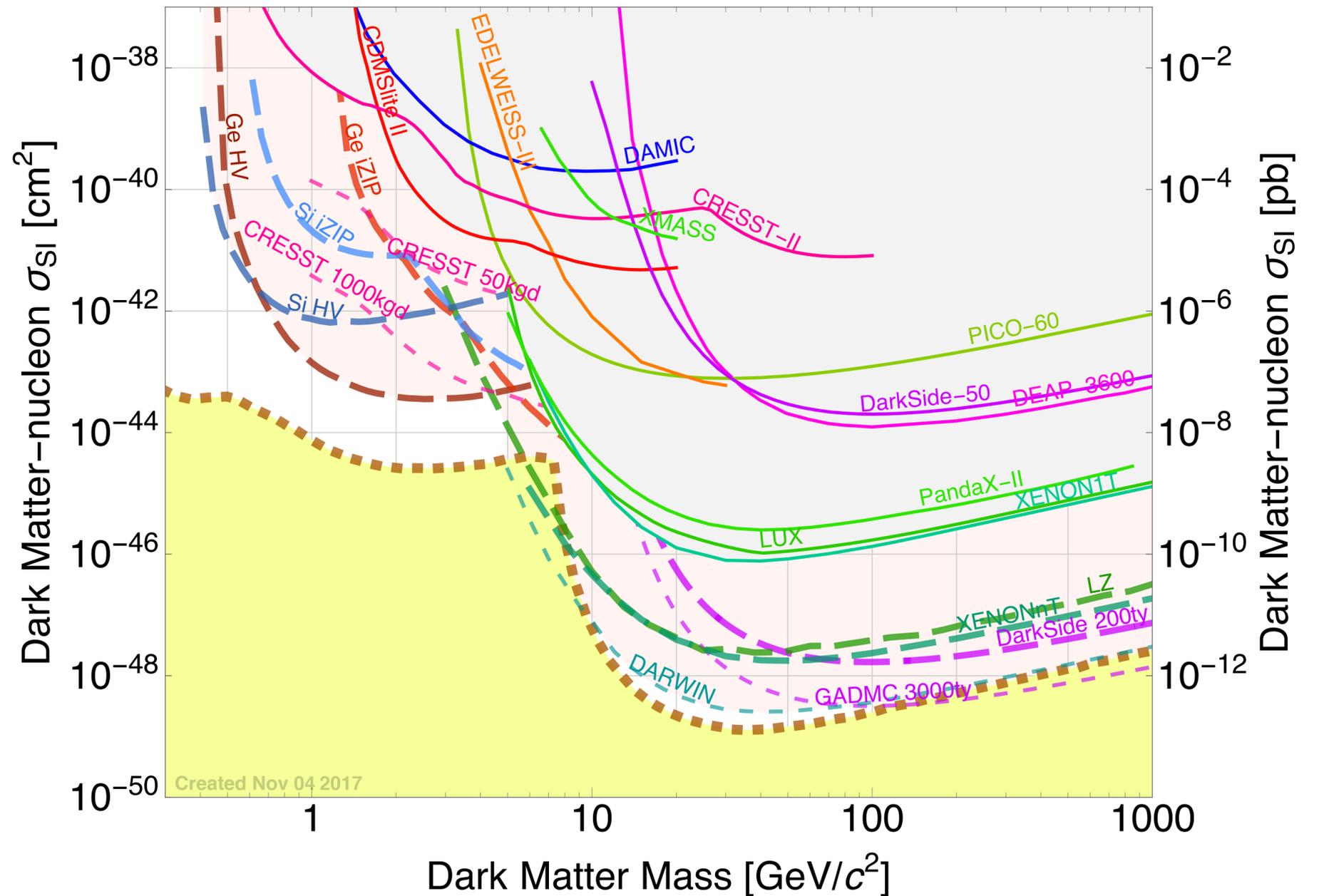
Scattering off nuclei



Scattering off electrons



WHERE ARE WE GOING?



CONCLUSIONS

- The next decade will be very exciting for dark matter direct detection. Various G2 Experiments will come online, covering a lot of new parameter space.
- Although WIMPs remain a very interesting dark matter candidate, other scenarios are gaining traction in the theoretical community, while new ideas for direct searches have been proposed and are gaining momentum.
- Given the wealth of theoretical possibilities, a diversity of experimental designs and targets will be needed to constrain the theory and couplings of any discovered signal.

