# Determination of proton parton distribution functions using ATLAS data

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#### **Motivations**



Proton parton distribution functions (PDFs) are key ingredients to provide theoretical predictions at the LHC:

- dominant uncertainty for high precision measurements (W mass, weak-mixing angle)
- important limiting factor for Higgs production cross-sections
- discovery of new physics relies on precise knowledge of PDFs (i.e. high mass searches)
- → Frontier high-energy physics at colliders requires percent-level accuracy for PDFs

How LHC experiments may contribute to this effort?

- LHC measurements are used in global PDF fits providing important complementary information to datasets from DIS and fixed target experiments in corners of the phase space previously unexplored
- PDFs fits from LHC collaborations allow a first direct test of the PDF-sensitivity of the measurements and represent a precious "vademecum" for global PDF fitting collaborations

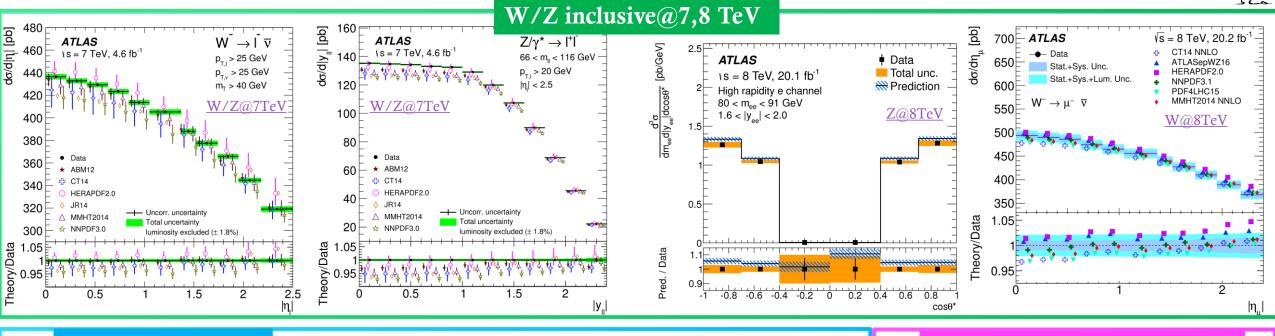
## **ATLAS PDFs fits**

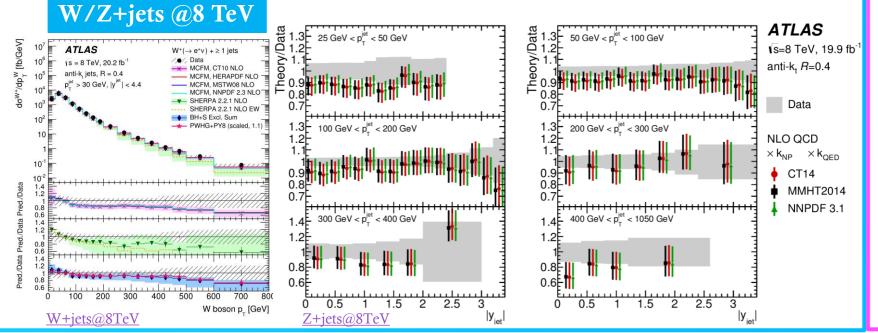


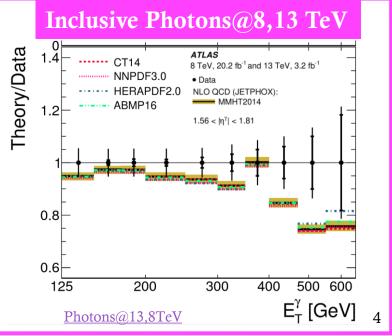
ATLAS PDF set (Publication)	Data	Sensitivity of ATLAS data
ATLASepWZ16 Eur.Phys.J.C 77 (2017) 367	HERA data + ATLAS inclusive W/Z cross-sections@ 7 TeV	Quark flavour separation
ATLASepWZVjet20 (JHEP 07 (2021) 223)	ATLASepWZ16 dataset + ATLAS W/Z+jets cross-sections@8 TeV	Quark flavour separation at higher <i>x</i>
ATLASepWZtop18 (ATL-PHYS-PUB-2018-017)	ATLASepWZ16 dataset + ATLAS $t\bar{t}$ cross-sections@ 8 TeV	Medium- and high-x gluon
ATLASpdf21 (arXiv:2112.11266 [hep-ex])	Topic of my talk today!	NEW
		NEW

## Data sets



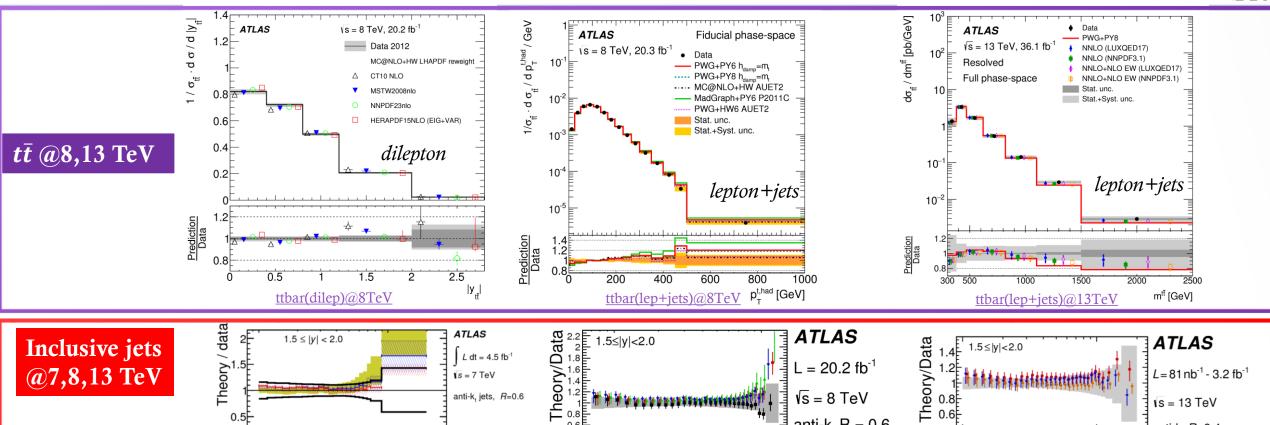




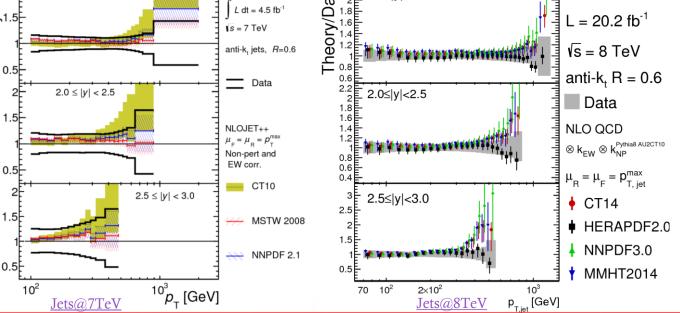


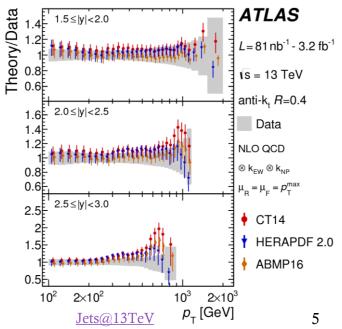
#### Data sets





- Data at 8 TeV in the nominal fit
- Data at 7 and 13 TeV 0.5 used for cross-check 2





## Data sets and theory



Summary of ATLAS data sets

		1		
Data set	$\sqrt{s}$ [TeV]	Luminosity [fb <sup>-1</sup> ]	Decay channel	Observables entering the fit
Inclusive $W, Z/\gamma^*$	7	4.6	$e, \mu$ combined	$\eta_{\ell}(W), y_{Z}(Z)$
Inclusive $Z/\gamma^*$	8	20.2	$e, \mu$ combined	$\cos \theta^*$ in bins of $y_{\ell\ell}$ , $m_{\ell\ell}$
Inclusive W	8	20.2	$\mu$	$\eta_{\mu}$
$W^{\pm}$ + jets	8	20.2	e	$p_{\mathrm{T}}^{W}$
Z + jets	8	20.2	e	$p_{\rm T}^{\rm jet}$ in bins of $ y^{\rm jet} $
$tar{t}$	8	20.2	lepton + jets, dilepton	$m_{t\bar{t}},p_{\mathrm{T}}^{t},y_{t\bar{t}}$
$tar{t}$	13	36	lepton + jets	$m_{t\bar{t}}, p_{\mathrm{T}}^t, y_t, y_{t\bar{t}}^{\mathrm{b}}$
Inclusive isolated $\gamma$	8, 13	20.2, 3.2	-	$E_{\rm T}^{\gamma}$ in bins of $\eta^{\gamma}$
Inclusive jets	7, 8, 13	4.5, 20.2, 3.2	-	$p_{\mathrm{T}}^{\mathrm{jet}}$ in bins of $ y^{\mathrm{jet}} $

<u>W,Z@7TeV</u>; <u>W@8TeV</u>; <u>Z@8TeV</u>; <u>W+jets@8TeV</u>; <u>Z+jets@8TeV</u>; <u>ttbar(dilep)@8TeV</u>; <u>ttbar(lep+jets)@8TeV</u>; <u>ttbar(lep+jets)@13TeV</u>; <u>Photons@13,8TeV</u>; <u>Jets@13TeV</u>.

**HERA data** - backbone of the fit – provide good constraints from  $\sim 10^{-4} < x < 0.1$ -0.4 (depending on which PDF) but they cannot provide:

- Flavour separation in the sea
- High-*x* gluon

these mainly come from ATLAS data

Fit to NNLO in QCD and NLO in EW achieved either by direct NNLO grids or by k-factor corrections

- current 'state of the art' – using <u>xFitter framework</u>

#### Fit Model



•Parametrisation of  $xu_v, xd_v, x\bar{u}, x\bar{d}, x\bar{s}$  and xg

$$xf_i(x) = A_i x^{B_i} (1-x)^{C_i} P_i(x)$$
 with  $P_i(x) = (1 + D_i x + E_i x^2 + F_i x^3)$ 

(with an extra negative term for the gluon  $-A'_g x^{B'g} (1-x)^{C'_g}$ )

- Normalisation  $(A_i)$  and slope parameters  $(B_i)$  of  $x\bar{u}, x\bar{d}, x\bar{s}$  all independent of each other (strategy only recently adopted by some global PDFs fitters)
- Assumed  $xs = x\bar{s}$
- $A_i$  parameters for the valence quarks and the gluon set by the sum-rules
- $D_{\nu}E_{\nu}F_{i}$  parameters are added until there is no further significant improvement in  $\chi^{2} \rightarrow 21$  free parameters
- $C'_g$  set to 25 to avoid negative contribution at high x

#### Fit Model



o Same  $\chi^2$  form as standard in HERA and ATLAS PDF fits

$$\chi^{2} = \sum_{ik} \left( D_{i} - T_{i} (1 - \sum_{j} \gamma_{ij} b_{j}) \right) C_{\text{stat,uncor},ik}^{-1}(D_{i}, D_{k}) \left( D_{k} - T_{k} (1 - \sum_{j} \gamma_{kj} b_{j}) \right) + \sum_{i} \log \frac{\delta_{i,\text{uncor}}^{2} T_{i}^{2} + \delta_{i,\text{stat}}^{2} D_{i} T_{i}}{\delta_{i,\text{uncor}}^{2} D_{i}^{2} + \delta_{i,\text{stat}}^{2} D_{i}^{2}} + \sum_{j} b_{j}^{2}$$

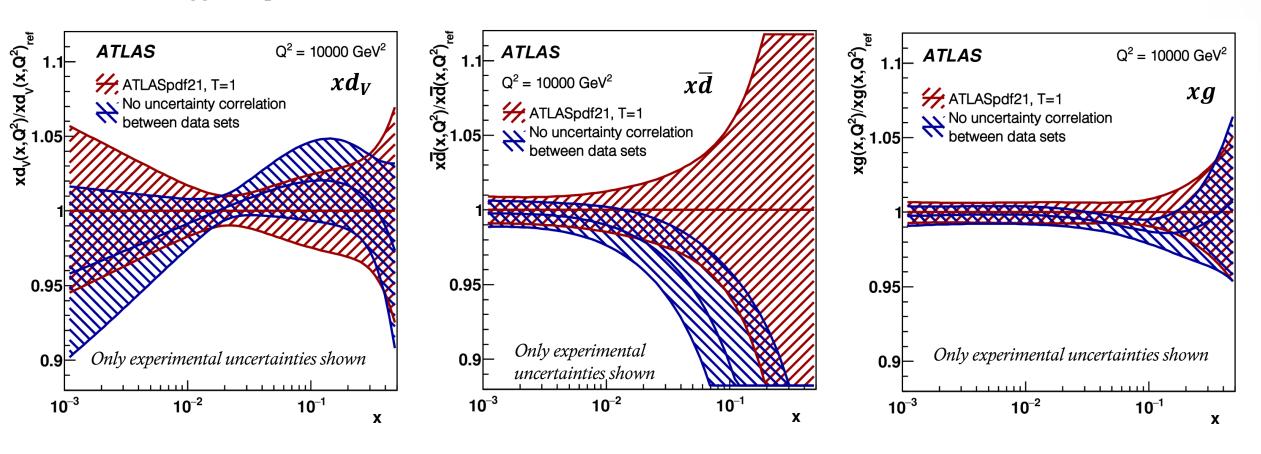
Considered correlations of statistical and systematic uncertainties within and between data: Jet Energy Scale correlations
play the dominant role

Measurements					
V+jets@8TeV	Z+jets@8TeV	tī@8TeV	<i>t</i> <del>t</del> <u>@</u> 13TeV	jets@8TeV	
JetScaleFlav2	Flavor Response	flavres-jes	JET29NP JET Flavour Response	syst JES Flavour Response*	
etScaleFlav1Known	Flavor Comp	flavcomp-jes	JET29NP JET Flavour Composition	syst JES Flavour Comp	
JetScalepunchT	Punch Through	punch-jes	-	syst JES PunchThrough MC15	
JetScalePileup2	PU OffsetMu	pileoffmu-jes	-	syst JES Pileup MuOffset	
-	PU Rho	pileoffrho-jes	JET29NP JET Pileup RhoTopology	syst JES Pileup Rho topology*	
JetScalePileup1	PU OffsetNPV	pileoffnpv-jes	JET29NP JET Pileup OffsetNPV	syst JES Pileup NPVOffset	
-	PU PtTerm	pileoffpt-jes	JET29NP JET Pileup PtTerm	syst JES Pileup Pt term	
JetJVFcut	JVF	jetvxfrac	-	syst JES Zjets JVF	
-	btag-jes J	JET29NP JET BJES Response	-	-	
-	jeten-res	JET JER SINGLE NP	-	-	
-	-	mup-scale		<u>-</u>	
-	-	muonms-res	MUON MS	<u>-</u>	
	-				
-	-	dibos-xsec			
	-	zjet-xsec	Zjets xsec	<del>-</del>	
	-	singletop-xsec	st xsec	<del>-</del>	
	JetScaleFlav2 etScaleFlav1Known JetScalepunchT JetScalePileup2  - JetScalePileup1  - JetJVFcut	JetScaleFlav2 Flavor Response etScaleFlav1Known JetScalepunchT Punch Through JetScalePileup2 PU OffsetMu - PU Rho JetScalePileup1 PU OffsetNPV - PU PtTerm JetJVFcut JVF - btag-jes - jeten-res	JetScaleFlav2 Flavor Response flavres-jes etScaleFlav1Known Flavor Comp flavcomp-jes JetScalepunchT Punch Through punch-jes JetScalePileup2 PU OffsetMu pileoffmu-jes  - PU Rho pileoffrho-jes  JetScalePileup1 PU OffsetNPV pileoffnpv-jes  - PU PtTerm pileoffpt-jes  JetJVFcut JVF jetvxfrac  - btag-jes JET29NP JET BJES Response  - jeten-res JET JER SINGLE NP  - mup-scale  - muonms-res  - muid-res  - dibos-xsec	JetScaleFlav2 Flavor Response flavres-jes JET29NP JET Flavour Response etScaleFlav1Known Flavor Comp flavcomp-jes JET29NP JET Flavour Composition JetScalepunchT Punch Through punch-jes -  JetScalePileup2 PU OffsetMu pileoffmu-jes -  PÜ Rho pileoffrho-jes JET29NP JET Pileup RhoTopology JetScalePileup1 PU OffsetNPV pileoffnpv-jes JET29NP JET Pileup OffsetNPV -  PU PtTerm pileoffpt-jes JET29NP JET Pileup PtTerm JetJVFcut JVF jetvxfrac -  btag-jes JET29NP JET BJES Response -  jeten-res JET JER SINGLE NP -  mup-scale MUON SCALE  muonms-res MUON MS  muid-res MUON ID  muid-res Diboson xsec  Zjets xsec	

## **Correlations**



Study the effect of the correlations at a scale relevant for LHC physics and focus on the middling *x* range where W,Z and Higgs are produced



Differences in PDFs not large (mostly in the *d*-type sector) and PDFs uncertainties are similar (larger on the high- $x \bar{d}$  and *gluon* PDFs if correlations are not accounted for), but correlation can be important when the aim is percent level accuracy on PDFs

## Impact of different datasets

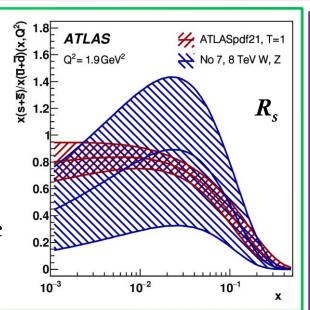


#### W/Z

Without W/Z data, ratio of strange to light sea quarks very poorly determined

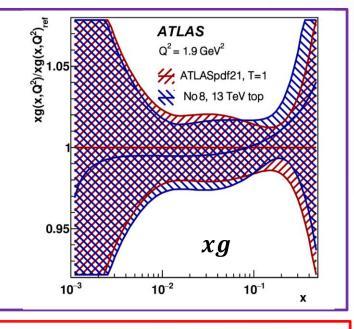
$$R_s = x(s+\bar{s})/x(\bar{u}+\bar{d})$$

Confirmed unsuppressed strange at low-*x* seen in previous ATLAS fits





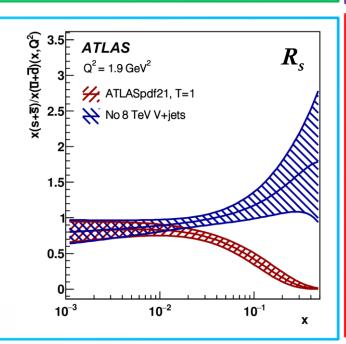
High-x gluon mildly softened when  $t\bar{t}$  data added and uncertainties of high-x gluon reduced



#### W/Z+jets

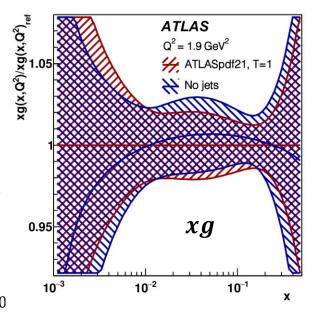
Without V+jets data little information of  $R_s$  at high x

V+jets data resolve a double minimum in the rest of data which marginally prefer  $x\bar{d}\sim x\bar{s}$  (blue).V+jets data increase  $\bar{d}$  and decrease  $\bar{s}$  at high-x (red)



#### Jets

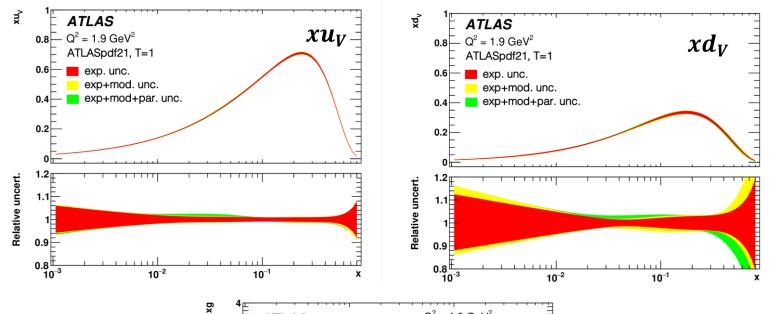
Impact of inclusive jet production (as direct photon) mostly on the high-x gluon



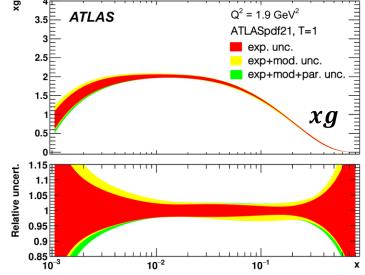
#### Uncertainties



Experimental uncertainties – dominant – also considered parametrisation uncertainties (extra  $F_{uv}$  and  $D_{dbar}$  parameters) and model ones (variation of  $Q_{min}^2 = 10 \text{ GeV}^2$  of HERA data, starting scale  $Q_0^2 = 1.9 \text{ GeV}^2$ , heavy quark masses including top)



Experimental uncertainties better determined in the *u*-type sector than for *d*-type one



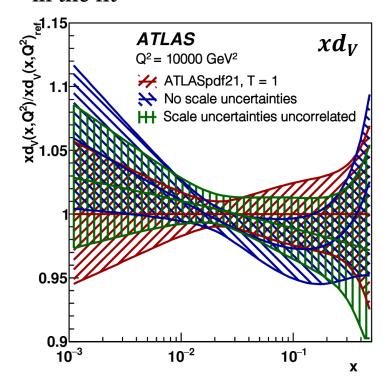
Gluon PDF well determined from 0.01 < x < 0.3 but not at low or high-x.

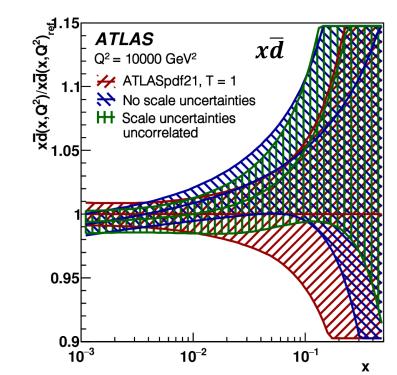
#### Uncertainties

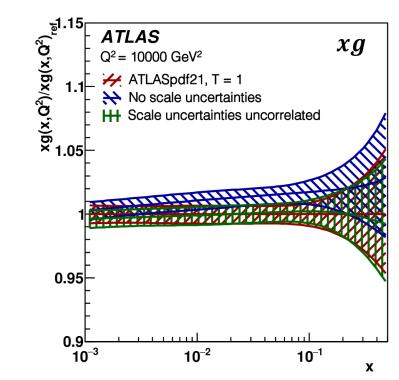


#### Also considered impact of scale uncertainties:

Scale uncertainties compatible with experimental ones in inclusive W/Z production ( $\sim 0.5\%$ ) so they have been included in the fit







Effect of scale uncertainties not large but can be important when the aim is percent level accuracy on PDFs Impact of scale uncertainties on other data sets also considered and found negligible

# $\chi^2$ tolerance

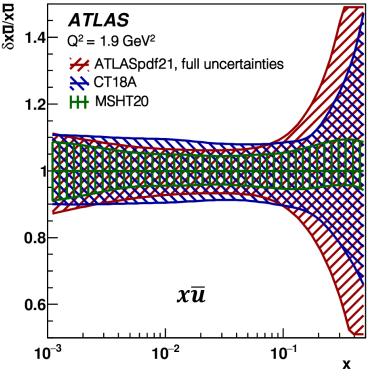


Experimental uncertainties in previous ATLAS+HERA fits evaluated with  $\chi^2$  tolerance  $T = \sqrt{\Delta \chi^2} = 1$ , but now more disparate ATLAS data

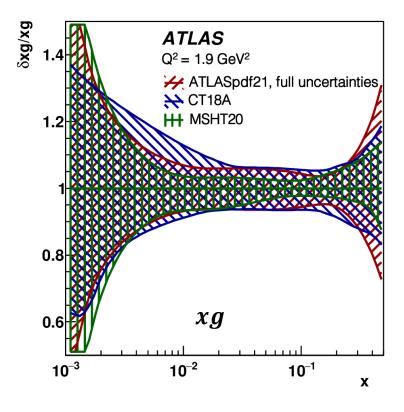
Some global PDF fitting groups (CT and MSHT) use enhanced  $\chi^2$  tolerances that are dynamic. Other groups use different methods but obtain similar uncertainties

Performed a study on an appropriate tolerance applying the procedure of the MSHT group and adopted T = 3

Similar uncertainties with some understood difference



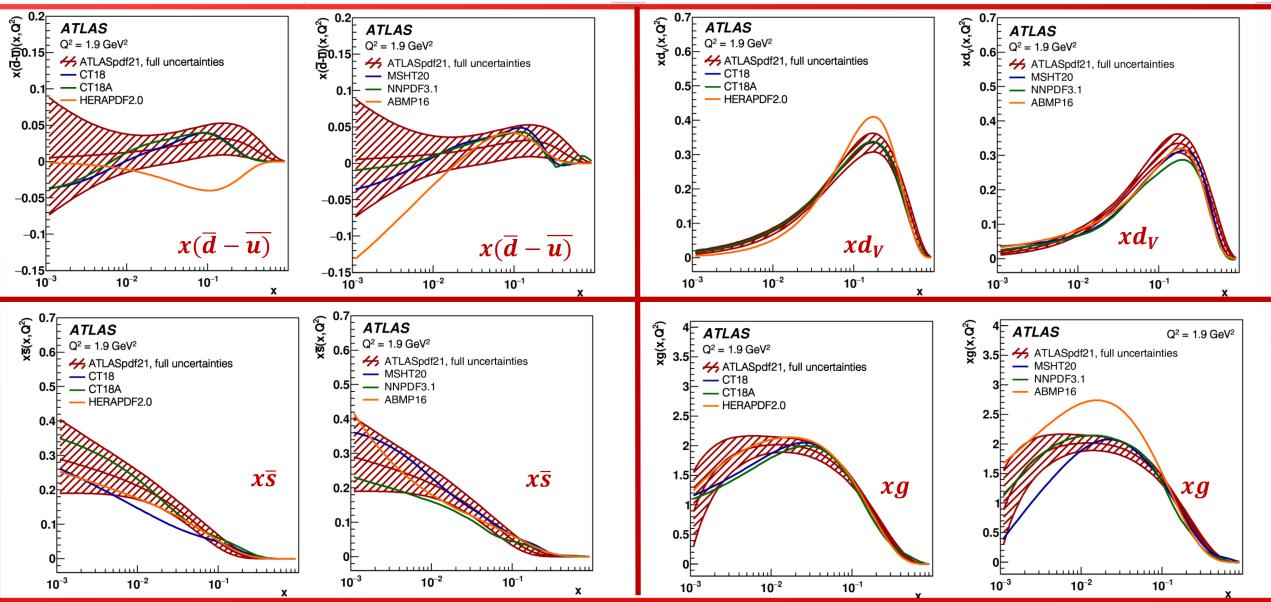
No fixed-target Drell-Yan data in ATLASpdf21



No Tevatron jet data in ATLASpdf21

# Comparison with global PDFs





**ATLASpdf21** fit agrees with the global fits as well as they agree with each other and ATLAS data bring the PDFs much closer to the global PDFs than to HERAPDF2.0

#### **Conclusions**



ATLASpdf21 fits to all PDF-sensitive ATLAS data sets that have information on correlated systematic uncertainties and NNLO QCD +NLO EW predictions:

- W/Z data allow to remove constraints on  $x\bar{u}$ ,  $x\bar{d}$ ,  $x\bar{s}$  and confirm unsuppressed strangeness at low x
- V+jets data constraint light-sea quarks and gluon at higher x
- jets,  $t\bar{t}$  and photon data constrain gluon PDF

An enhanced tolerance is used for realistic PDF uncertainty estimation and results with both choices (T=1 and T=3) made public

There is reasonable agreement with modern PDF sets

Correlations between data sets are included and information made public for global PDF fitters Scale uncertainties considered and included when comparable to experimental uncertainties:  $\rightarrow$ Both aspects relatively small but important for future high precision measurements as  $m_W$  and  $\sin^2\theta_W$ 

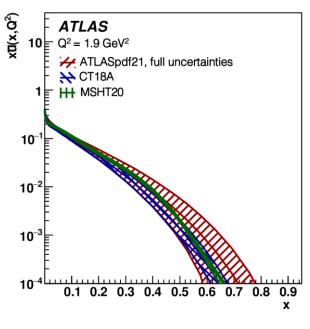
ATLASpdf21 (arXiv:2112.11266 [hep-ex]) represents a precious "vademecum" on the treatment of ATLAS data for global PDF fitting collaborations

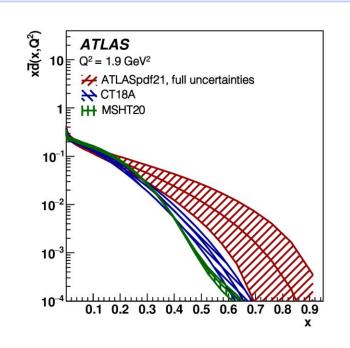


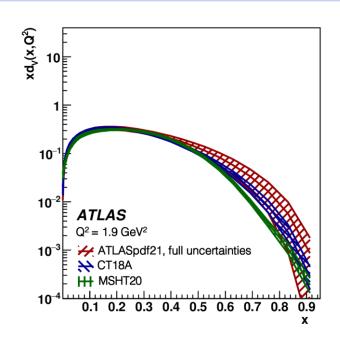
## **BACKUP**

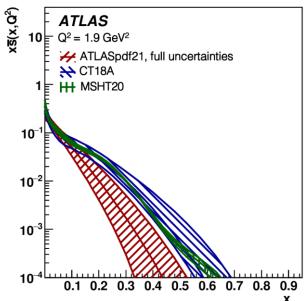
## Comparison with global PDFs at high x

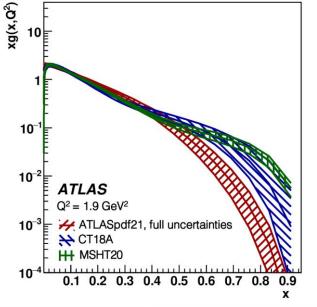












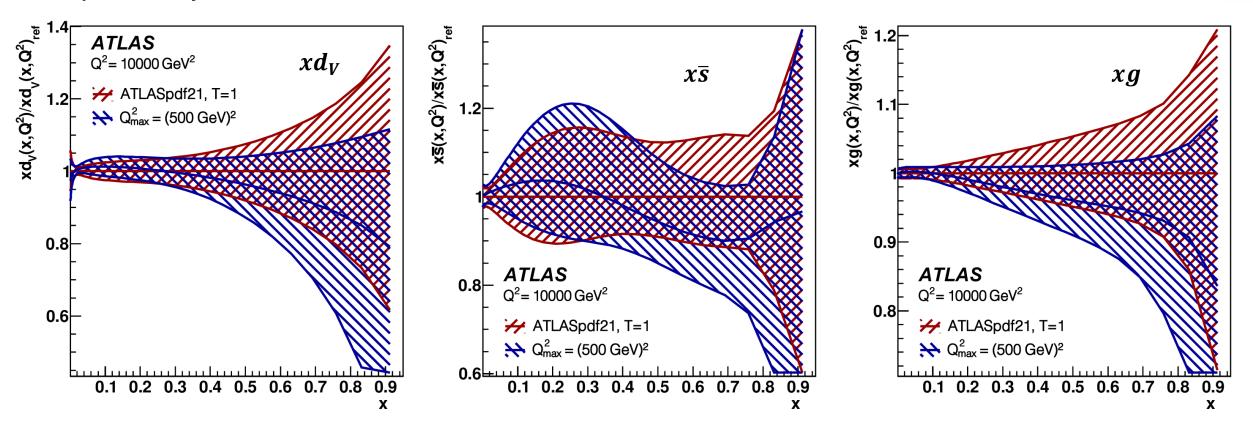
Some discrepancies for  $x \ge 0.7$ , where no data available to determine any of the PDFs. Nevertheless agreement within  $\sim 2\sigma$ .

# High Q<sup>2</sup> data



Subtle effects of BSM physics may be present in high-scale data. If these are included in a PDF fit, then the estimates of background in searches could be distorted.

Check if the PDFs differ when we cut out possible hidden new physics in the high scale data  $(Q_{MAX}^2 = (500 \text{ GeV})^2)$ Mainly removed jet data



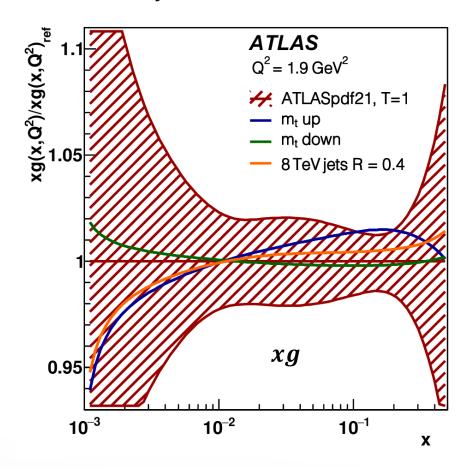
PDF shapes and uncertainties not strongly affected  $\rightarrow$ maximum  $Q^2$  cut not considered in the fit

#### Jet radius



Model uncertainty also includes Jet radius choice:

W/Z+jets and  $t\bar{t}$  measurements performed with a jet radius of R=0.4 instead inclusive jets measurement performed with R=0.6. Applied a jet radius choice uncertainty considering the results of the fit obtained with R=0.4 instead of R=0.6 for inclusive jets

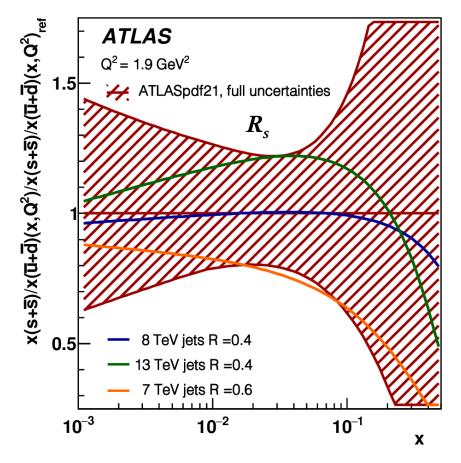


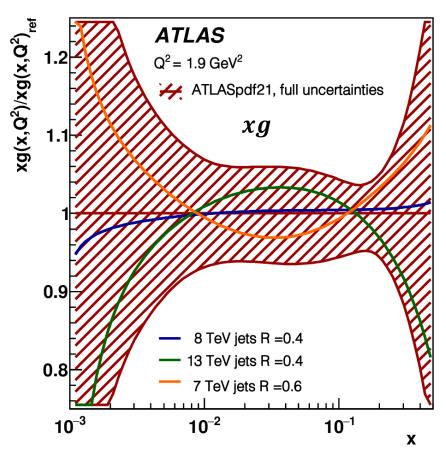
Only significant change is in the gluon PDF

## Jet data



Inclusive jet data@7,8 and 13 TeV not included simultaneously in the fit since full experimental systematic uncertainty correlations between these data sets are not known





Differences in PDFs dominated by the change in centre-of-mass energy.

PDFs using 8 TeV jet data lie between those of the 7 TeV jet data and the 13 TeV jet data at 13 TeV. Differences are not significant compared to the full PDFs uncertainties





$$\chi^{2} = \sum_{ik} \left( D_{i} - T_{i} (1 - \sum_{j} \gamma_{ij} b_{j}) \right) C_{\text{stat,uncor},ik}^{-1}(D_{i}, D_{k}) \left( D_{k} - T_{k} (1 - \sum_{j} \gamma_{kj} b_{j}) \right) + \sum_{i} \log \frac{\delta_{i,\text{uncor}}^{2} T_{i}^{2} + \delta_{i,\text{stat}}^{2} D_{i} T_{i}}{\delta_{i,\text{uncor}}^{2} D_{i}^{2} + \delta_{i,\text{stat}}^{2} D_{i}^{2}} + \sum_{j} b_{j}^{2}$$

Partial  $\chi^2$ : for each data set 1<sup>st</sup> contribution being the main, 2<sup>nd</sup> being a small bias correction Correlated term between data sets

 $D_i$  = measured data;  $T_i$ = theory prediction;  $\gamma_{ij}$  correlated systematics accounted for using nuisance parameters  $b_j$   $C_{stat,uncor,ik}$  = statistical (plus uncorrelated) covariance matrix;

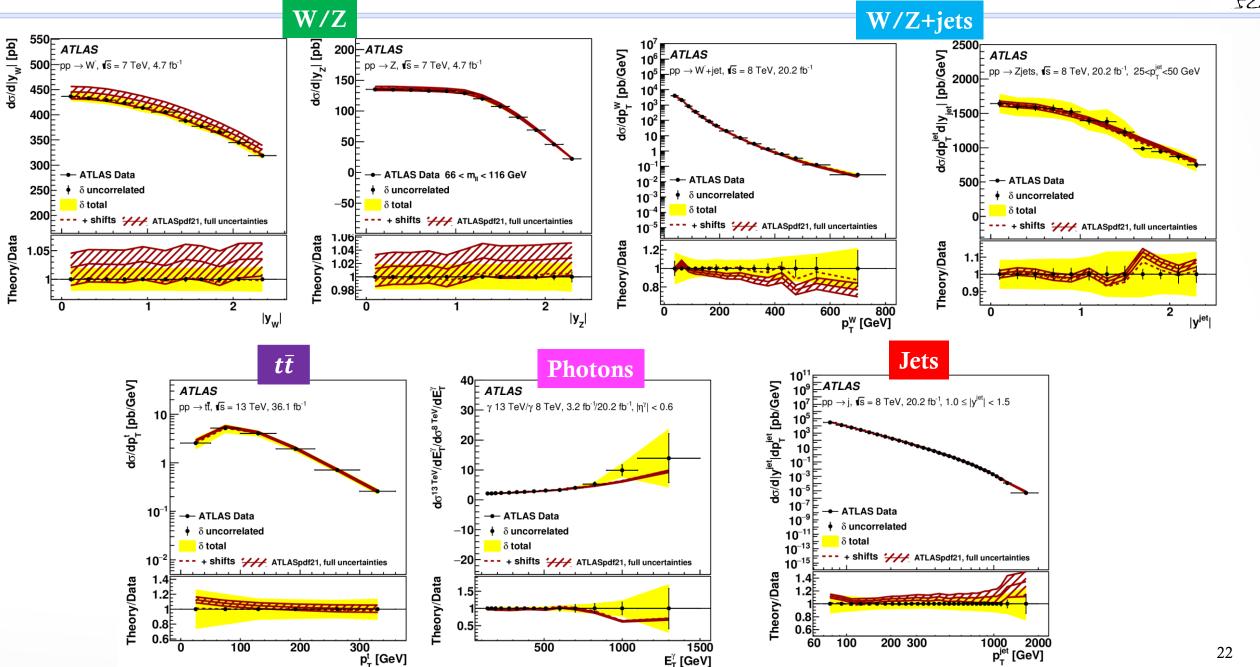
 $\delta_{i,uncor}$ =uncorrelated systematics uncertainties;

 $\delta_{i,stat}$ = statistical uncertainties

Total $\chi^2/NDF$	2010/1620	
HERA $\chi^2/NDP$	1112/1016	
HERA correlated term	50	
ATLAS $W$ , $Z$ 7 TeV $\chi^2$ /NDP	68/55	
ATLAS $Z/\gamma^*$ 8 TeV $\chi^2/\text{NDP}$	208/184	
ATLAS W 8 TeV $\chi^2$ /NDP	31/22	
ATLAS W and $Z/\gamma^*$ 7 and 8 TeV		
correlated term	71 = (38 + 33)	
ATLAS direct $\gamma$ 13/8 TeV $\chi^2$ /NDP	27/47	
ATLAS direct $\gamma$ 13/8 TeV		
correlated term	6	
ATLAS V+ jets 8 TeV $\chi^2$ /NDP	105/93	
ATLAS $t\bar{t}$ 8 TeV $\chi^2$ /NDP	13/20	
ATLAS $t\bar{t}$ 13 TeV $\chi^2/\text{NDP}$	25/29	
ATLAS inclusive jets 8 TeV $\chi^2$ /NDF	207/171	
ATLAS V+jets 8 TeV and		
$t\bar{t}$ + jets 8,13 TeV and		
R = 0.6 inclusive jets 8 TeV correlated term	87 = (16 + 9 + 21 + 41)	

## Data sets

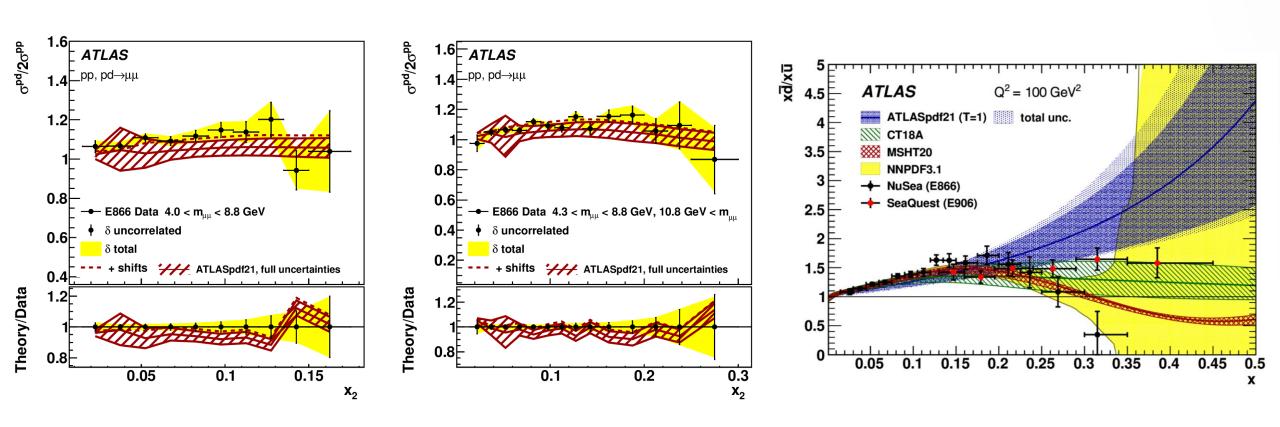




## Fixed-target Drell-Yan data



Fixed-target Drell-Yan data from E866 and E906, uniquely able to constrain the difference  $x(\bar{d}-\bar{u})$  at high x



Quite good description of the E866 data, even better agreement with the new data from E906